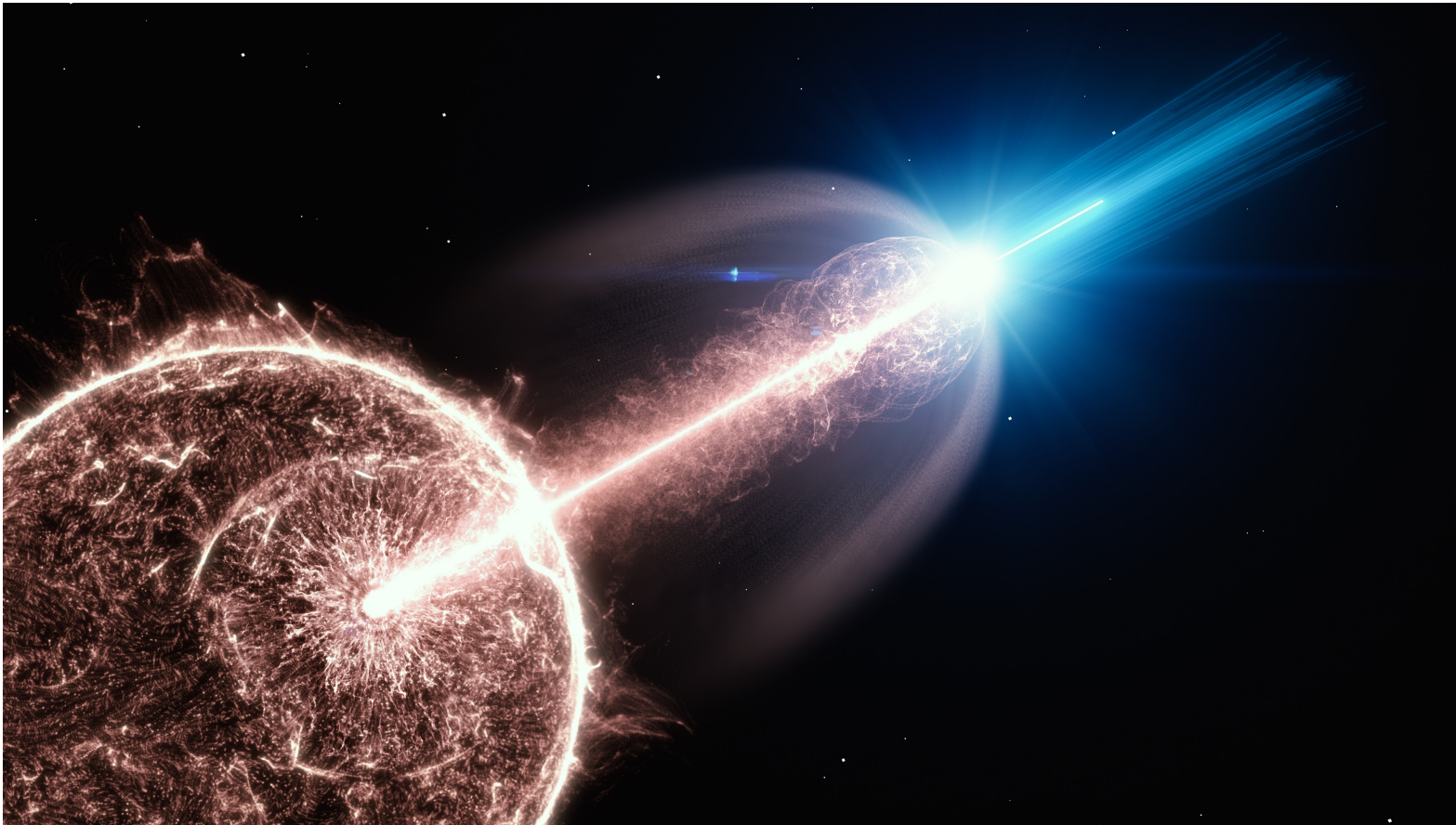
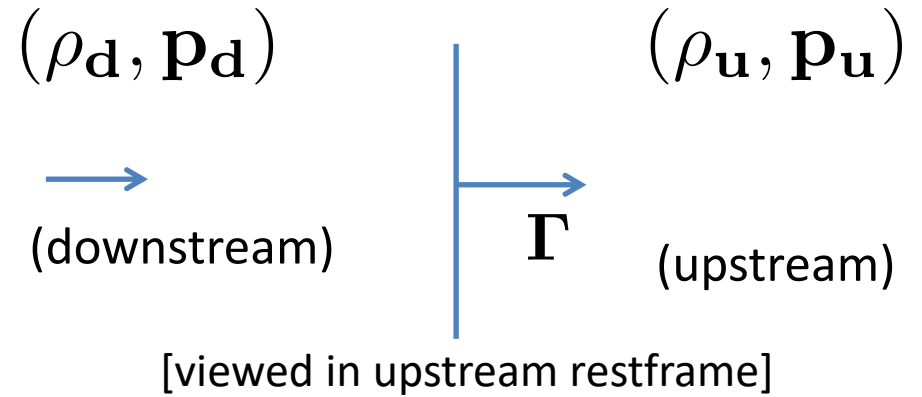
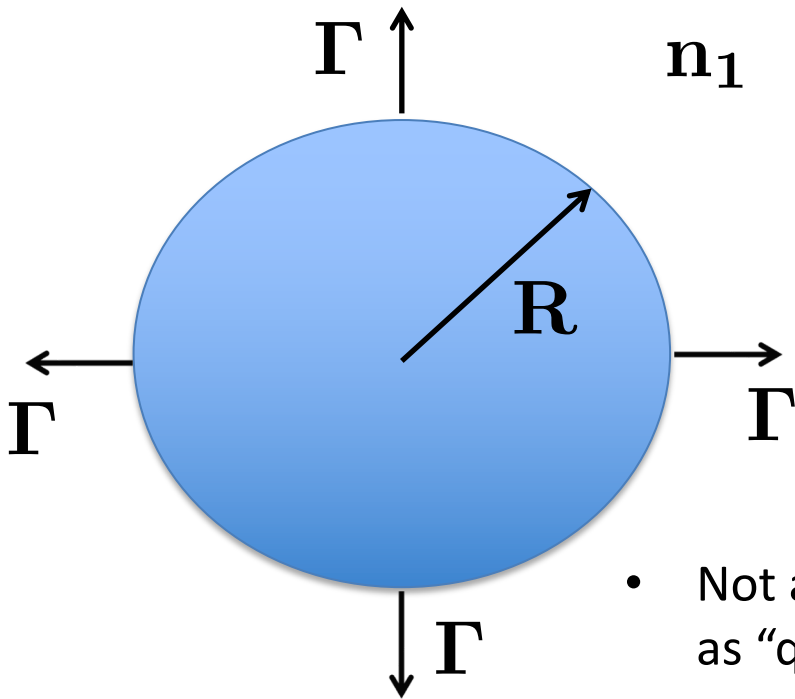


# Gamma-Ray Emission from GRBs Outflows- An Overview



**Caveat!** The Gamma-Ray emission I will be focusing on here is the radiation observed during the “afterglow” phase of the GRB

# What Are GRBs?



- Not actually isotropic outflows, but can be considered as “quasi-isotropic” since  $\theta_{\text{jet}} > 1/\Gamma$
- Isotropic equivalent energy in gamma-rays,  $E_{\text{iso}}$ , around  $10^{54}$  erg, is close to Gravitational binding energy limit
- Extremely efficient emitters in terms of converting kinetic energy flux to radiation

# Evolutionary Phases of Blastwave

Assuming shock is radiative (ie. incoming KE flux radiated away)

[R. Blandford + McKee 1976]

$$\frac{dE_k}{dt} = -4\pi R^2 \beta (\Gamma^2 \rho - \Gamma \rho)$$

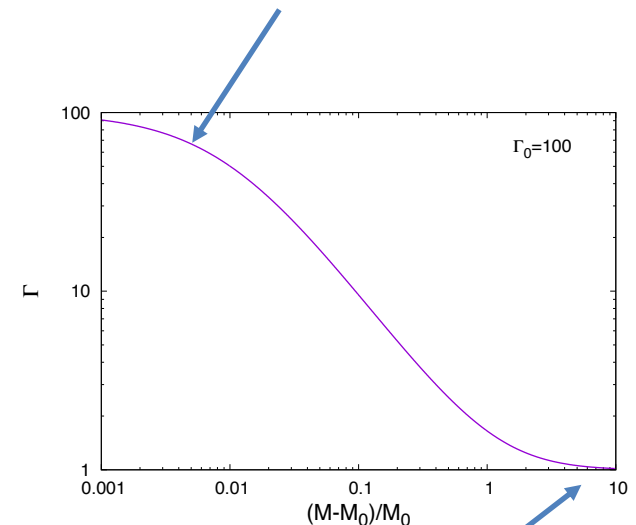
$$E_k = \frac{\rho 4\pi R^3}{3} (\Gamma - 1)$$

$$\frac{d\Gamma}{dM} = -\frac{(\Gamma^2 - 1)}{M}$$

This has the solution

$$\Gamma - 1 = 2 \left( \frac{M^2 (\Gamma_0 + 1)}{M_0^2 (\Gamma_0 - 1)} - 1 \right)^{-1}$$

Critical mass where free expansion changes to deceleration phase

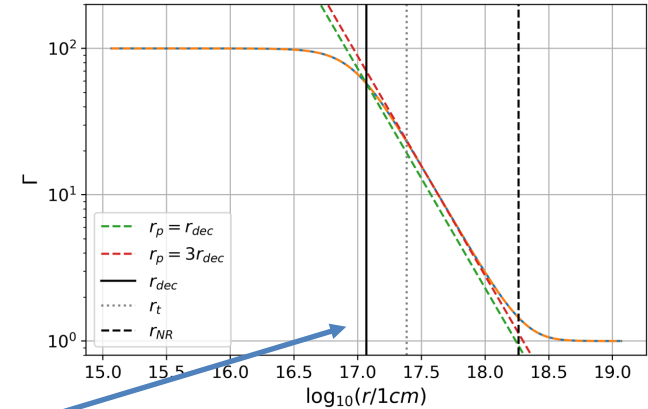


Blast wave becomes non-relativistic

# Temporal Compression of Observed Signal

For a constant density medium, during the deceleration phase,

$$\Gamma \propto r^{-3/2}$$



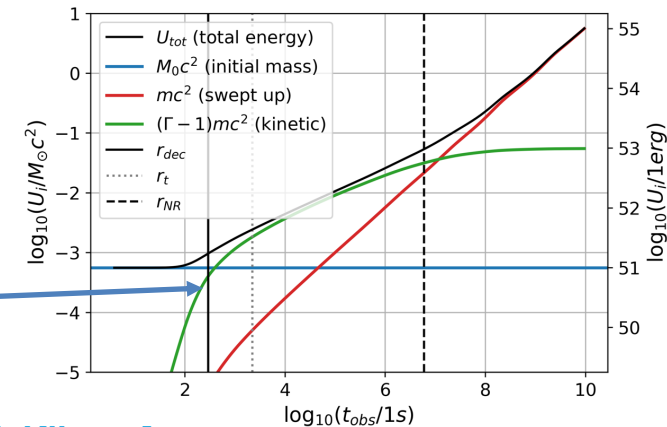
$$R_{dec} \approx 10^{17} \left( \frac{E_{iso}}{10^{53} \text{ erg}} \right)^{1/3} \left( \frac{n}{1 \text{ cm}^{-3}} \right)^{-1/3} \left( \frac{\Gamma}{100} \right)^{-2/3} \text{ cm}$$

Since moving emitter is observed along the beam direction

$$cdt_{obs} = (1 - \beta)dr \approx dr/2\Gamma^2$$

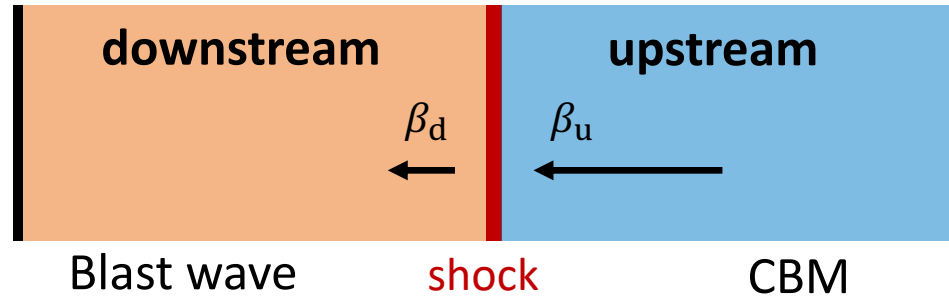
$$\Gamma \propto t_{obs}^{-3/8}$$

$$t_{dec}^{obs} \approx \frac{R}{c\Gamma^2} = 300 \left( \frac{R}{10^{17} \text{ cm}} \right) \left( \frac{\Gamma}{100} \right)^{-2} \text{ s}$$



# Relativistic Hydro Shocks

What's the compression ratio for relativistic shocks?



Mass Flux:

$$\rho_u \beta_u \Gamma_u = \rho_d \beta_d \Gamma_d$$

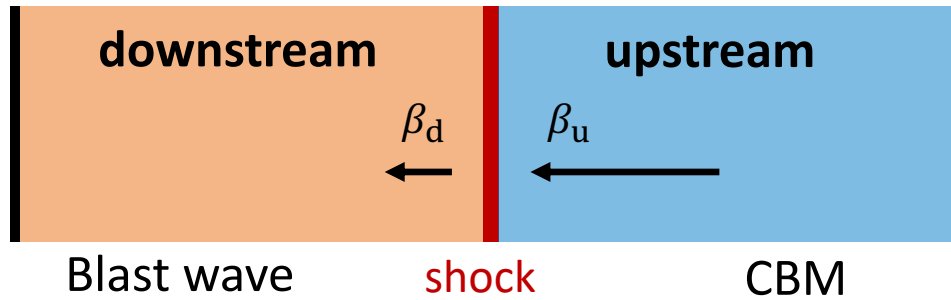
Momentum Flux:  $\mathbf{p}_u + \mathbf{w}_u \beta_u^2 \Gamma_u^2 = \mathbf{p}_d + \mathbf{w}_d \beta_d^2 \Gamma_d^2$

Energy Flux:

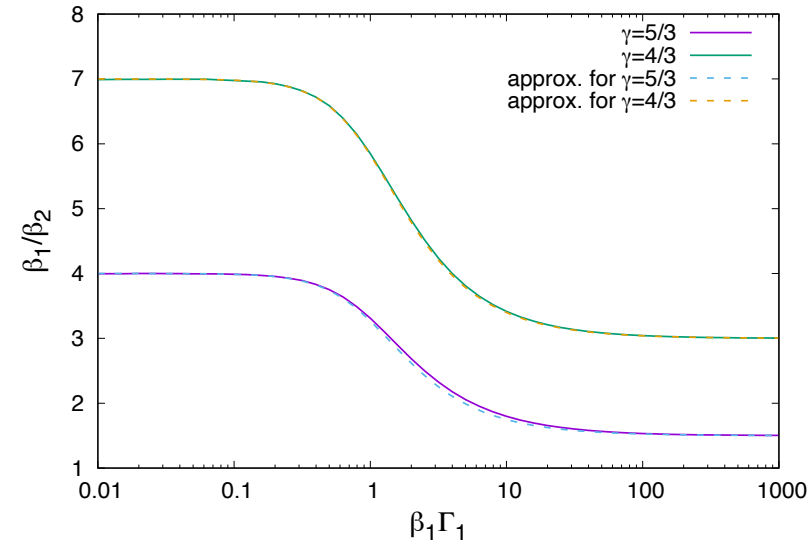
$$\mathbf{w}_u \beta_u \Gamma_u^2 = \mathbf{w}_d \beta_d \Gamma_d^2$$

$$\mathbf{w}_{\text{rel.}} = \frac{\gamma}{\gamma - 1} \mathbf{p} + \rho$$

# Rel. Hydro Shock- Downstream Partition of the Upstream Ram Pressure



[viewed in shock restframe]

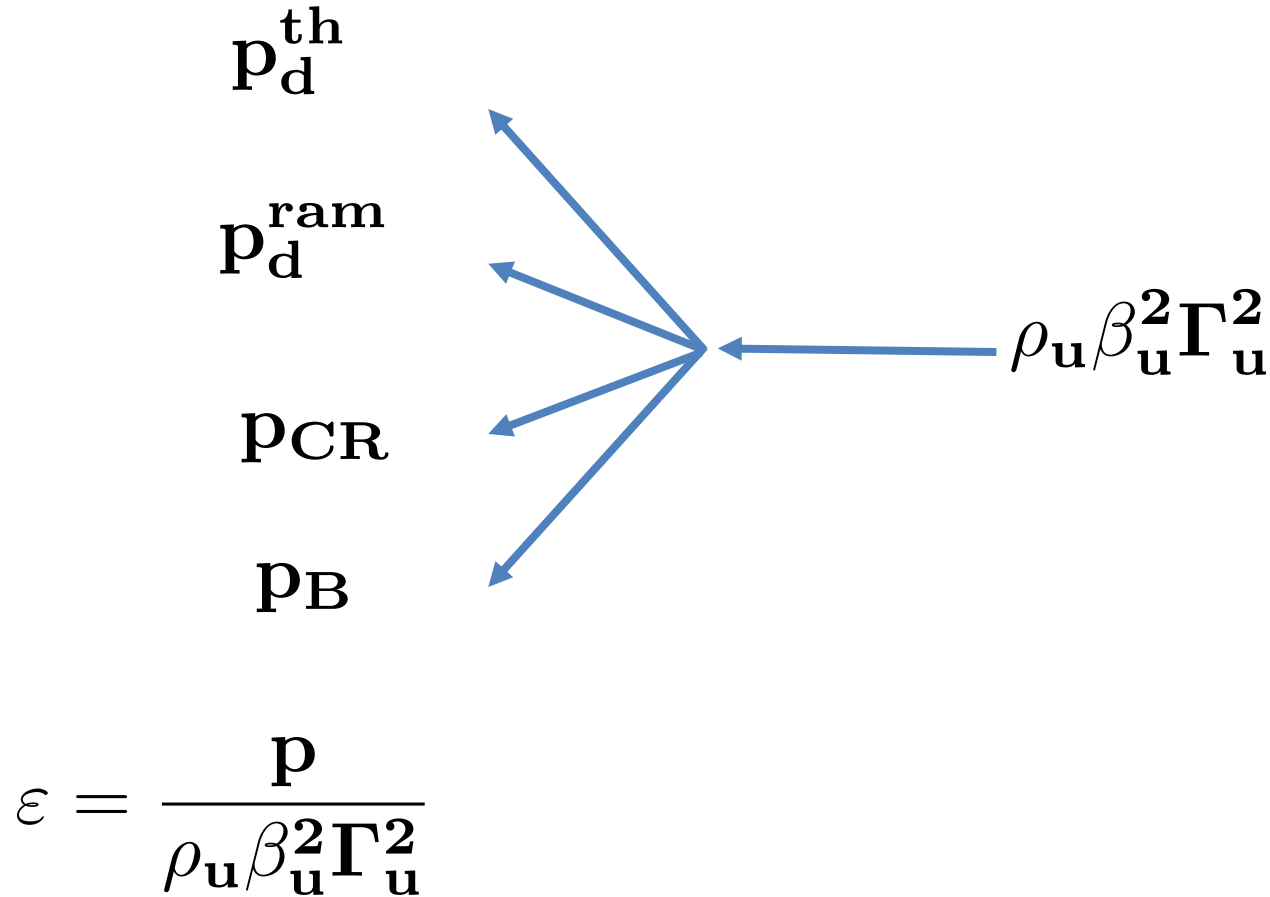


$$P_d = \frac{2}{3} \rho_u \beta_u^2 \Gamma_u^2$$

$$\rho_d \beta_d^2 \Gamma_d^2 = \frac{1}{3} \rho_u \beta_u^2 \Gamma_u^2$$



# Rel. MHD Shock- Downstream Partition of the Upstream Ram Pressure



# Relativistic MHD Shocks

Downstream magnetic field partition of upstream ram pressure:

$$\varepsilon_B = \frac{U_B}{\rho_u \beta_u^2 \Gamma_u^2}$$

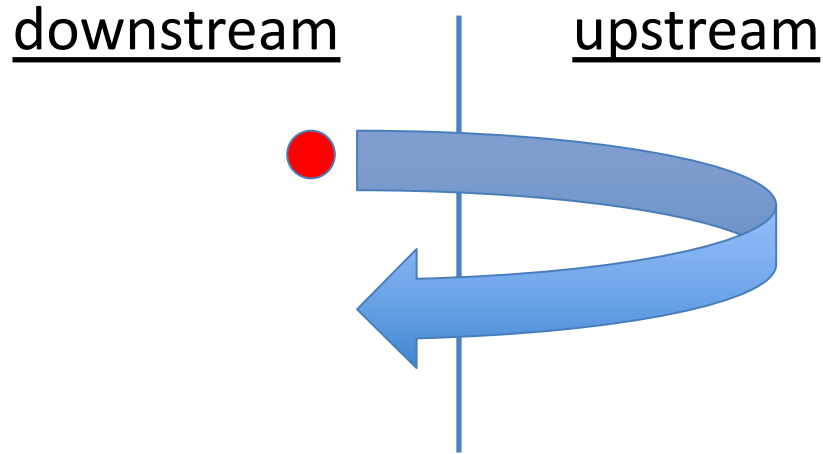
For

$$\varepsilon_B = 0.1 \quad n_u = 1 \text{ cm}^{-3} \quad \beta_u \Gamma_u = 10$$
$$B \approx 0.6 \text{ G}$$

$$\varepsilon_B = 10^{-5} \quad n_u = 1 \text{ cm}^{-3} \quad \beta_u \Gamma_u = 10$$
$$B \approx 6 \text{ mG}$$



# Particle Acceleration and Magnetic Turbulence



$$\begin{aligned} t_{\text{acc.}} &= \Delta t_{\text{cyc}} (\mathbf{E} / \Delta \mathbf{E}_{\text{cyc}}) \\ &= t_{\text{scat}} / \beta^2 \end{aligned}$$

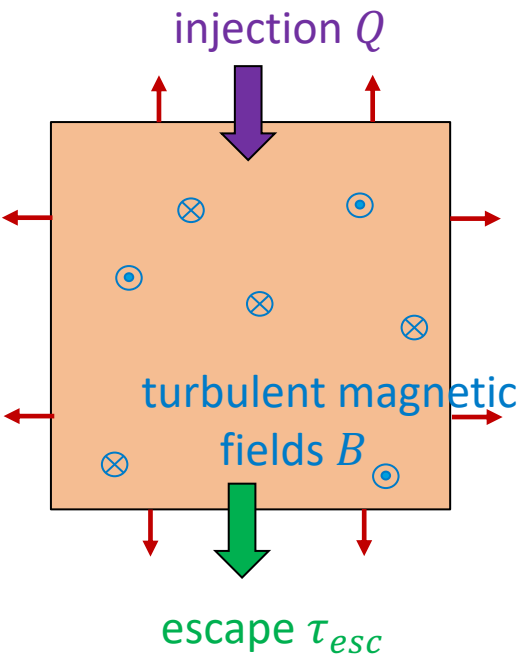
- Isotropisation is caused by magnetic turbulence, its rate is described by the scattering time, which in Larmor time units is  $\eta$

$$t_{\text{scat}} = \eta \frac{R_{\text{lar}}}{c}$$

- Scattering agent velocity  $\beta$  dictates energy gain each crossing cycle

# One Zone Model

$$\frac{\partial n_{\mathbf{p}}}{\partial t} = -\nabla_{\mathbf{p}} \cdot \left[ \frac{\mathbf{p}}{\tau_{\text{acc}}(\mathbf{p})} n_{\mathbf{p}} - \frac{\mathbf{p}}{\tau_{\text{loss}}(\mathbf{p})} n_{\mathbf{p}} \right] - \frac{n_{\mathbf{p}}}{\tau_{\text{esc}}(\mathbf{p})} + \mathbf{Q}$$



Note the absence of spatial information in the transport equation

[Diagram Courtesy of M. Klinger]

# Hadronic Particle Acceleration in Sources

$$\frac{\partial n_{\mathbf{p}}}{\partial t} = -\nabla_{\mathbf{p}} \cdot \left[ \frac{\mathbf{p}}{\tau_{\text{acc}}(\mathbf{p})} n_{\mathbf{p}} - \frac{\mathbf{p}}{\tau_{\text{loss}}(\mathbf{p})} n_{\mathbf{p}} \right] - \frac{n_{\mathbf{p}}}{\tau_{\text{esc}}(\mathbf{p})} + Q$$

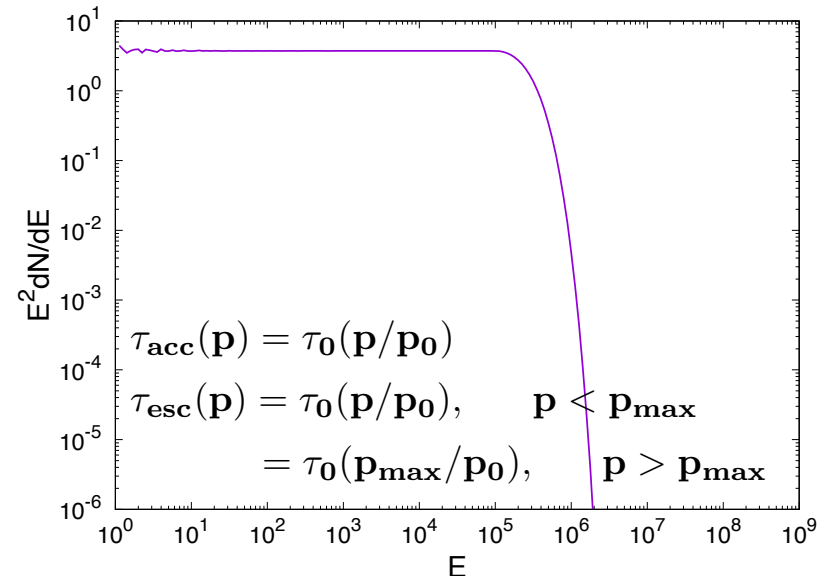
Steady state

No losses

Delta injection

$$n_{\mathbf{p}} = Q \left( \frac{\mathbf{p}}{\mathbf{p}_0} \right)^{-\left(1 + \frac{\tau_{\text{acc}}}{\tau_{\text{esc}}}\right)}$$

Note- shock acceleration is not the only acceleration process on the block

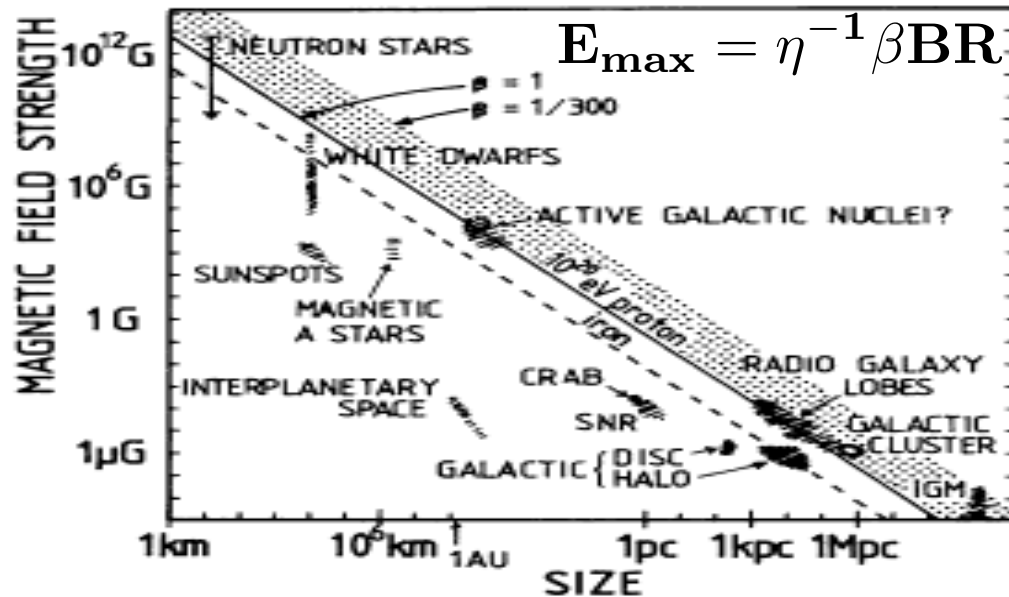


# Cosmic Ray Source Requirements

$$t_{\text{acc}} = \eta \frac{R_{\text{lar}}}{c\beta^2}$$

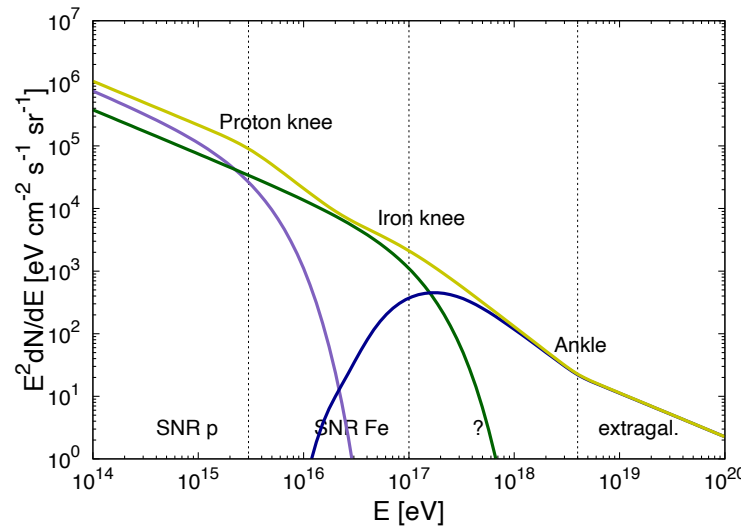
$$t_{\text{esc.}} = \frac{R}{c\beta}$$

[AM Hillas (1984)]

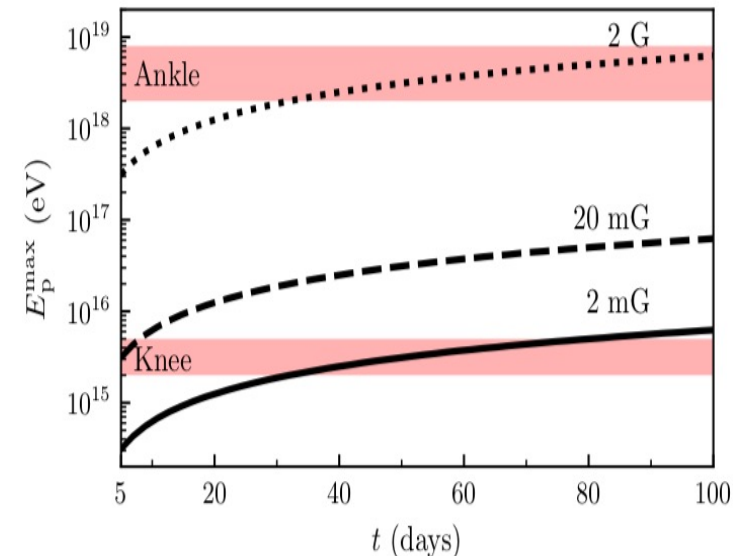


Not many objects appear capable of accelerating cosmic rays up to EeV energies. Blackhole related phenomena seem most promising- **AGN** and **GRB**

# GRB Outflows as a Cosmic Ray Sources



- As the source expands, **CRs** can be accelerated to energies between the **knee and the ankle**
- If the  $B$ -field is as large as  $\sim G$   $\rightarrow$  possibility of **UHECRs**



[X. Rodrigues, A. Taylor, et al., ApJ 2019]

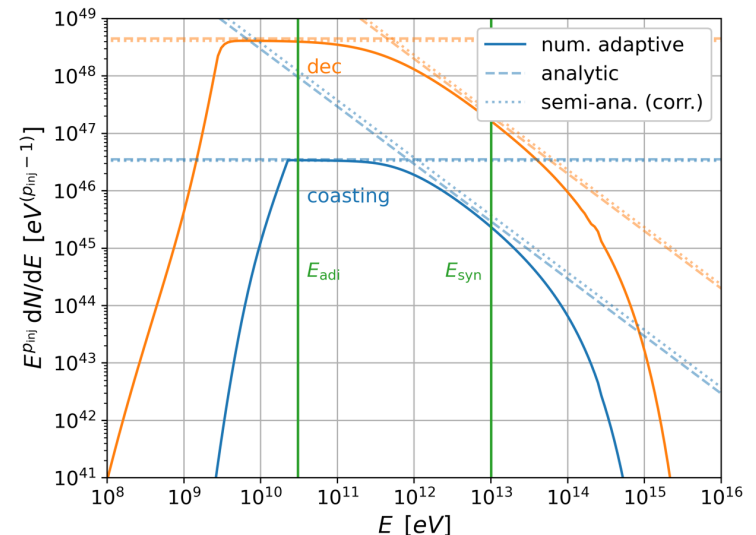
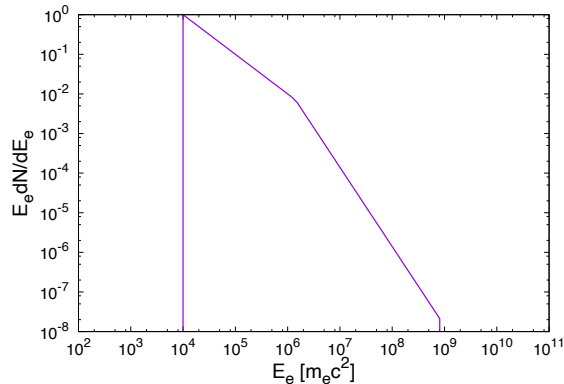
# Electron Spectrum Produced in Sources

$$\frac{\partial n_{\mathbf{p}}}{\partial t} = -\nabla_{\mathbf{p}} \cdot \left( -\frac{\mathbf{p}}{\tau_{\text{loss}}(\mathbf{p})} n_{\mathbf{p}} \right) - \frac{n_{\mathbf{p}}}{\tau_{\text{esc}}(\mathbf{p})} + \mathbf{Q}$$

Steady state

$$\tau_{\text{eff}} = \left( \tau_{\text{loss}}^{-1} + \tau_{\text{esc}}^{-1} \right)^{-1}$$

$$n_{\mathbf{p}} \approx \mathbf{Q} \tau_{\text{eff}}$$



# Electron Acceleration with Cooling

$$t_{\text{acc}} = \eta \frac{R_{\text{lar}}}{c\beta^2}$$

$$t_{\text{cool}} = \frac{9}{8\pi\alpha} \left( \frac{U_{\text{Bcrit}}}{U_{\text{B}}} \right) \left( \frac{h}{E_e} \right)$$

$$B_{\text{crit}} = 4 \times 10^{13} \text{ G}$$

$$E_e^{\text{max}} = \left( \frac{\eta^{-1/2}}{\alpha^{1/2} (B/B_{\text{crit}})^{1/2}} \right) m_e c^2$$

Maximum synchrotron energy tells us how efficient accelerator is!

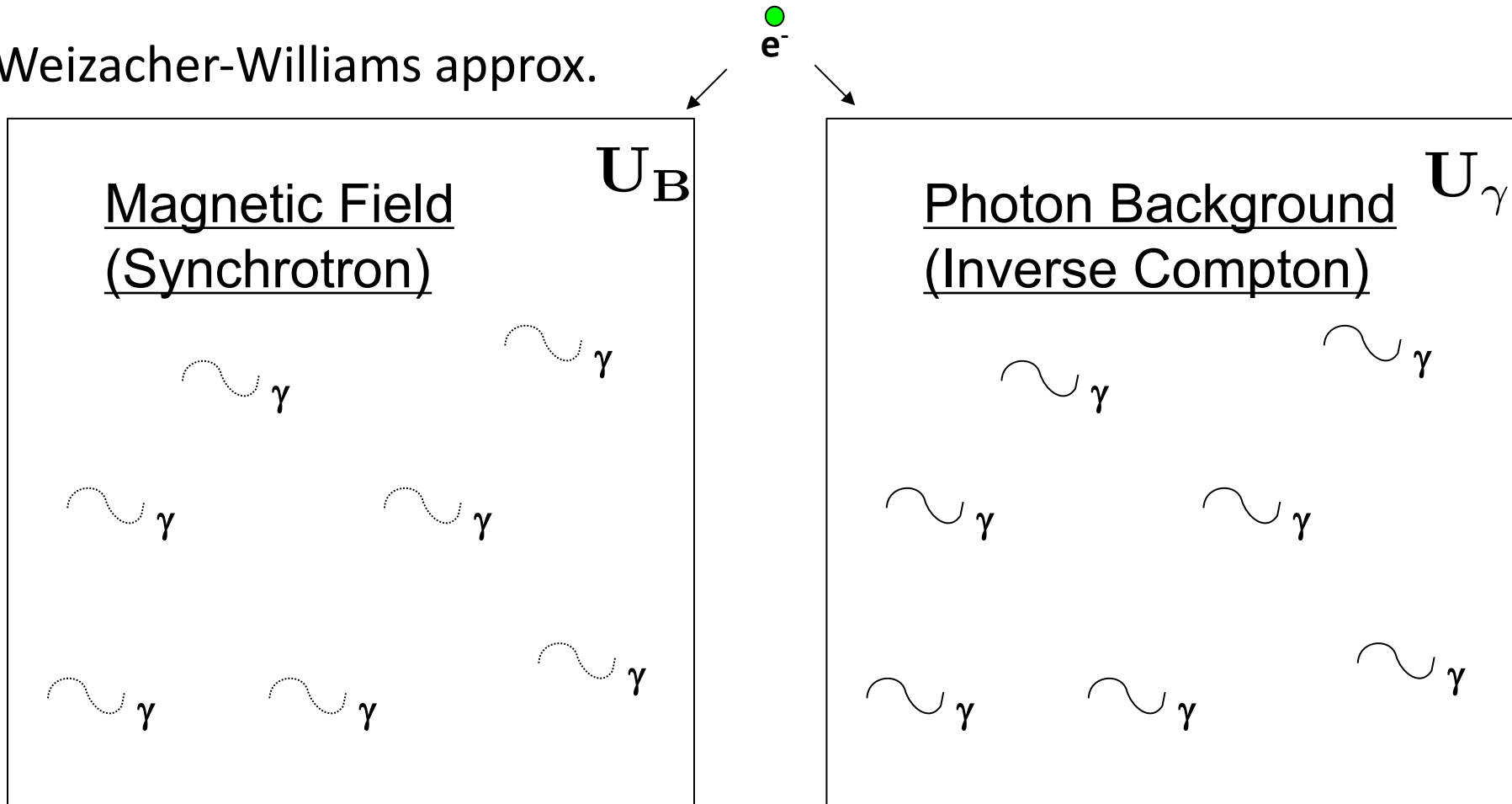
$$E_{\gamma}^{\text{sync}} \approx \frac{9}{4} \eta^{-1} \beta^2 \frac{m_e}{\alpha}$$



Where do synchrotron cutoffs for **AGN** and **GRB** sit in energy?

# Possible VHE Emission Processes

Weizacher-Williams approx.



Virtual Photons

Real Photons

$$E_\gamma^{\text{target}} = \left( \frac{B}{B_{\text{crit}}} \right) m_e c^2$$

Andrew Taylor

$$E_\gamma^{\text{target}}$$



# Efficiency Transfer Efficiency for Inverse Compton Emission

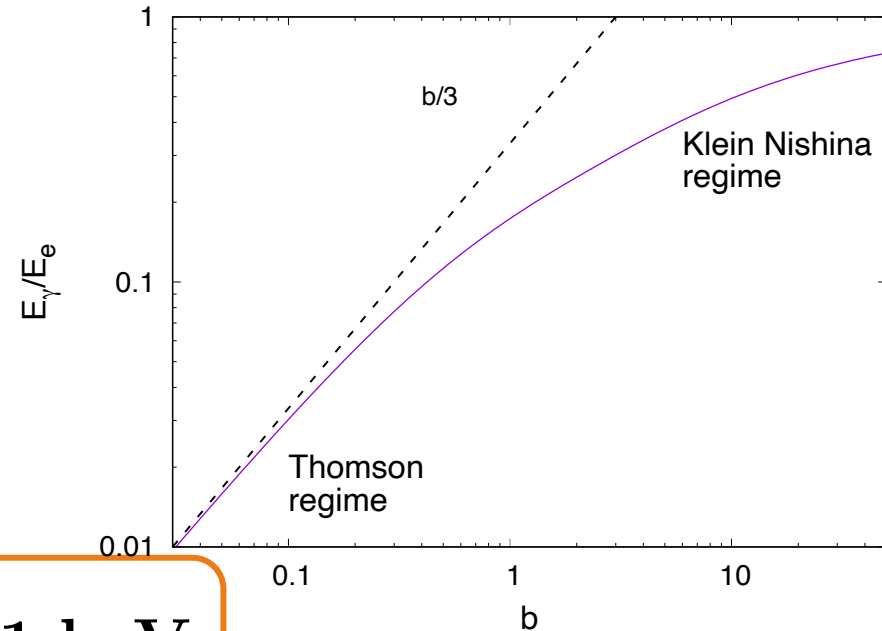
$$E_{\gamma}^{\text{IC}} \approx \left( \frac{b}{1+b} \right) E_e$$

$$b = \frac{4E_e E_{\gamma}^{\text{target}}}{(m_e c^2)^2}$$

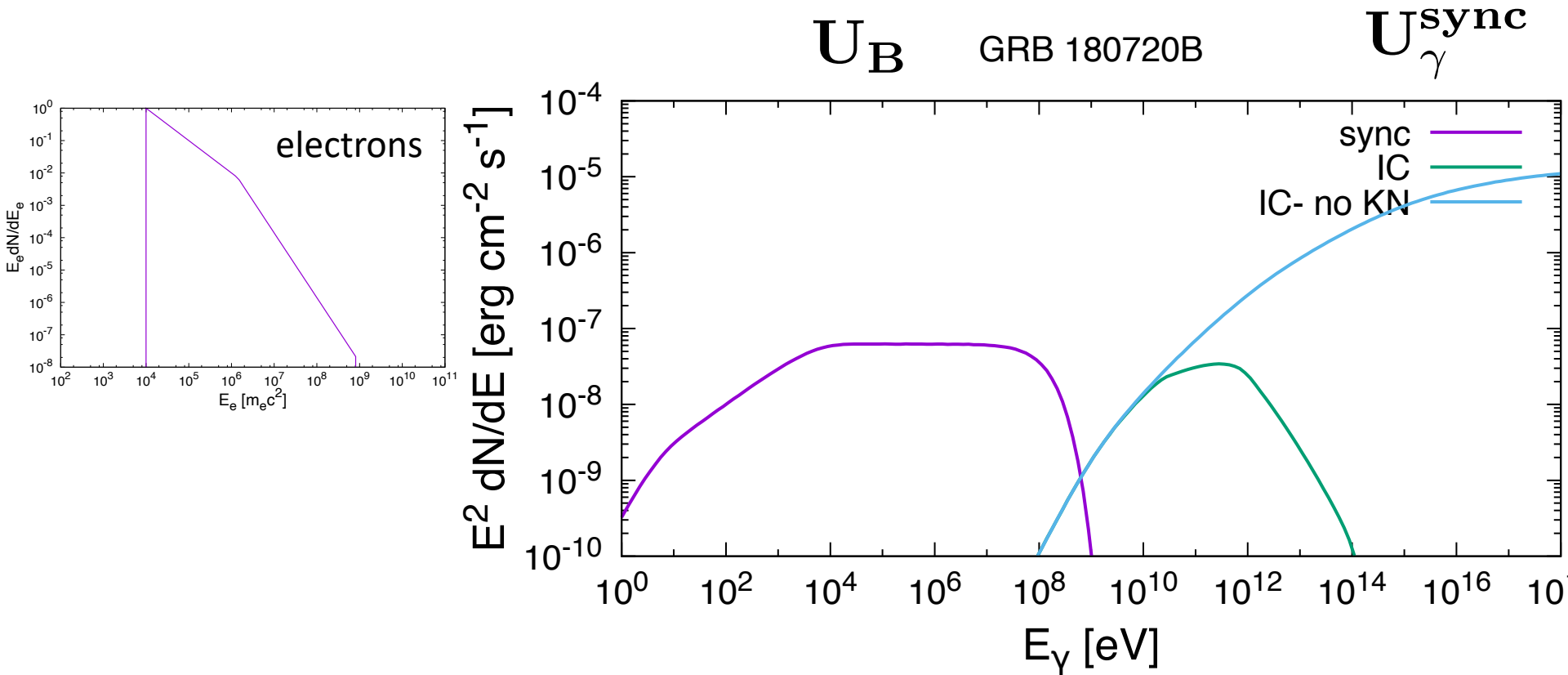
$$E_e = 1 \text{ TeV}$$

$$E_{\gamma}^{\text{target}} = 1 \text{ keV}$$

$$\left( \frac{b}{1+b} \right) \approx 1$$



# Afterglow GRB SED- Expected from SSC Model



Without KN effects, the ratio of the heights of the IC to Synchrotron bumps would scale with  $U_e/U_B$  (ie.  $\epsilon_e/\epsilon_B$ )

An SSC origin of the VHE emission has been adopted by others to describe early time VHE emission

[Nature 575, 459-463 (2019)]

# GRB Energy Flux Histogram

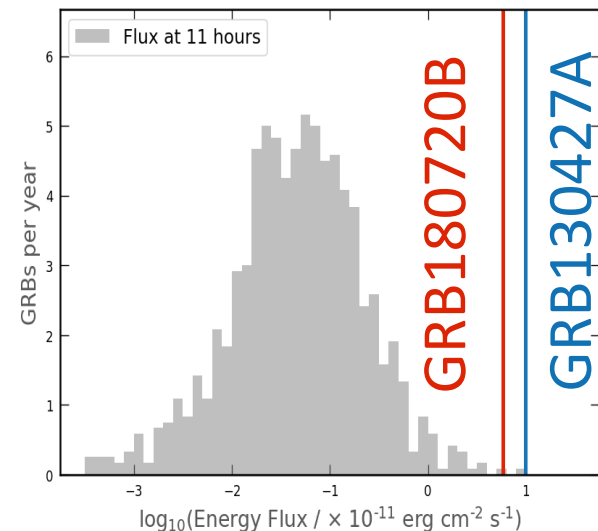
- GRBs at HE and VHE:  
~12 GRBs per year Fermi-LAT
- However, most science learnt from brightest event- GRB130427A: 94 GeV max energy photon.

VHE emission has been a decades-long mystery

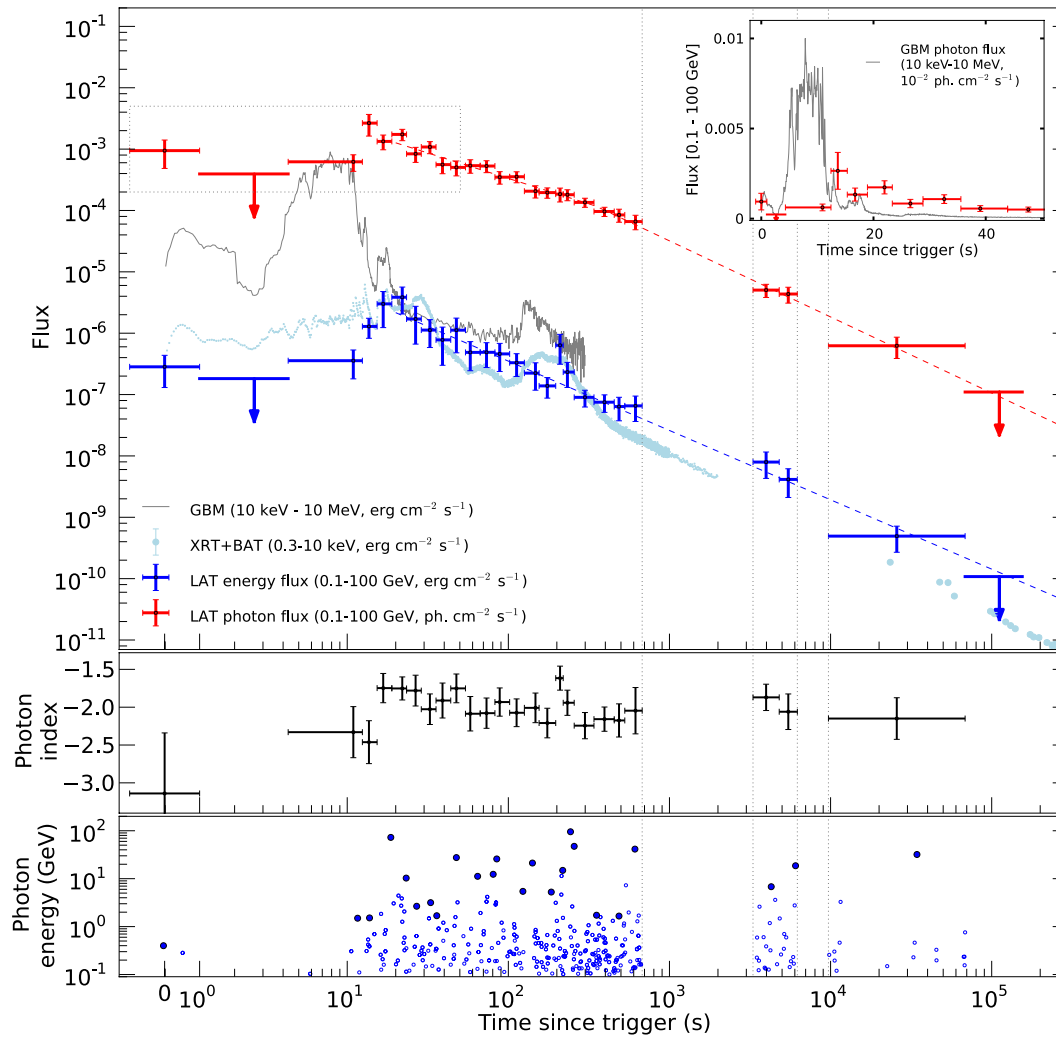
$t_{90}^{\text{GBM}} \sim 138 \text{ s}$ ,  $t_{90}^{\text{BAT}} \sim 163 \text{ s}$   
 $z = 0.34$

- Fermi-LAT detection from  $T_0$  to  $T_0+10000 \text{ s}$  (max. energy photon  $>90 \text{ GeV}$ ).
- Extremely bright burst:
  - 2nd brightest afterglow measured by Swift-XRT.

Swift-XRT GRBs  
energy flux distribution at 11 hours

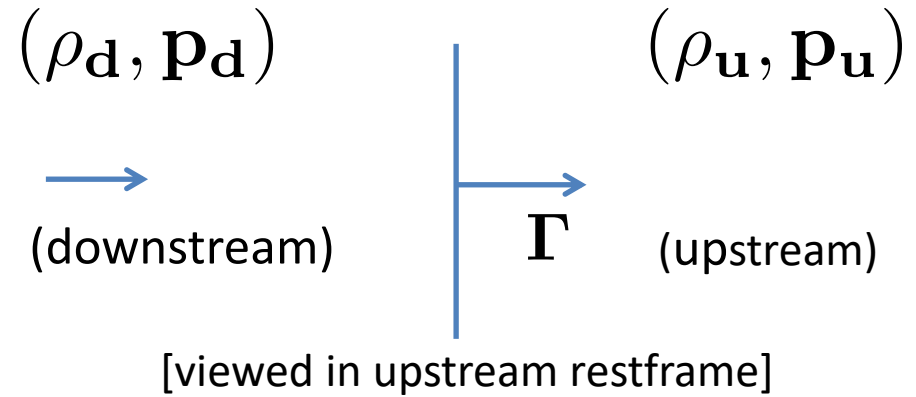
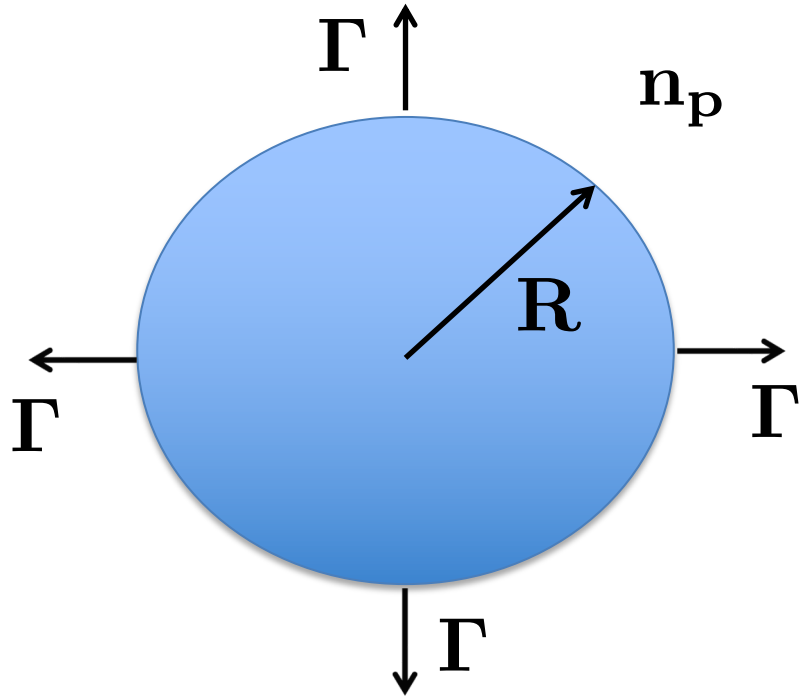


# GRB 130427A Lightcurve



$$L_{\text{XRT}} \propto t^{-\alpha}$$
$$\alpha = 1.17 \pm 0.06$$

# Origin of Temporal Decay Structure



Assuming  $\eta_\gamma$  is constant in time.....

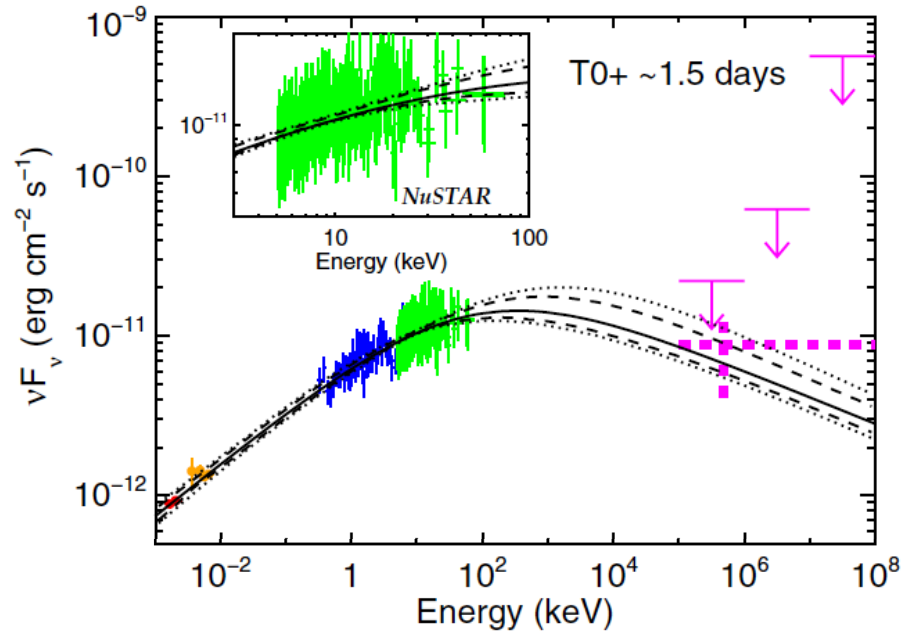
$$\frac{L_{\text{sync}}^{\text{iso}}}{4\pi\Gamma^2 R^2 c} = \epsilon_{\text{rad}} \Gamma^2 n_p m_p c^2$$

$$\Gamma \propto t^{-3/8} \quad R \propto t^{1/4}$$

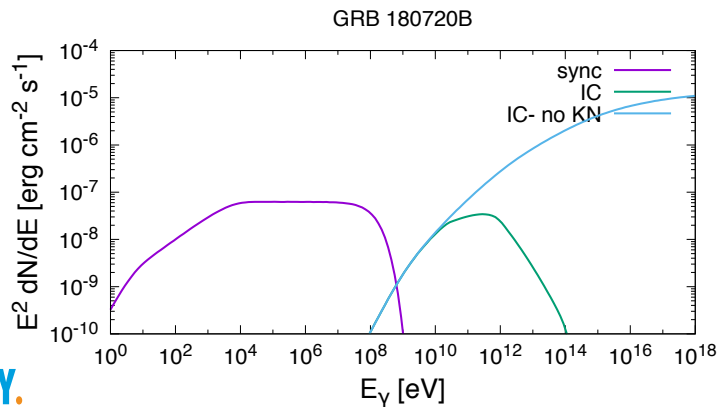
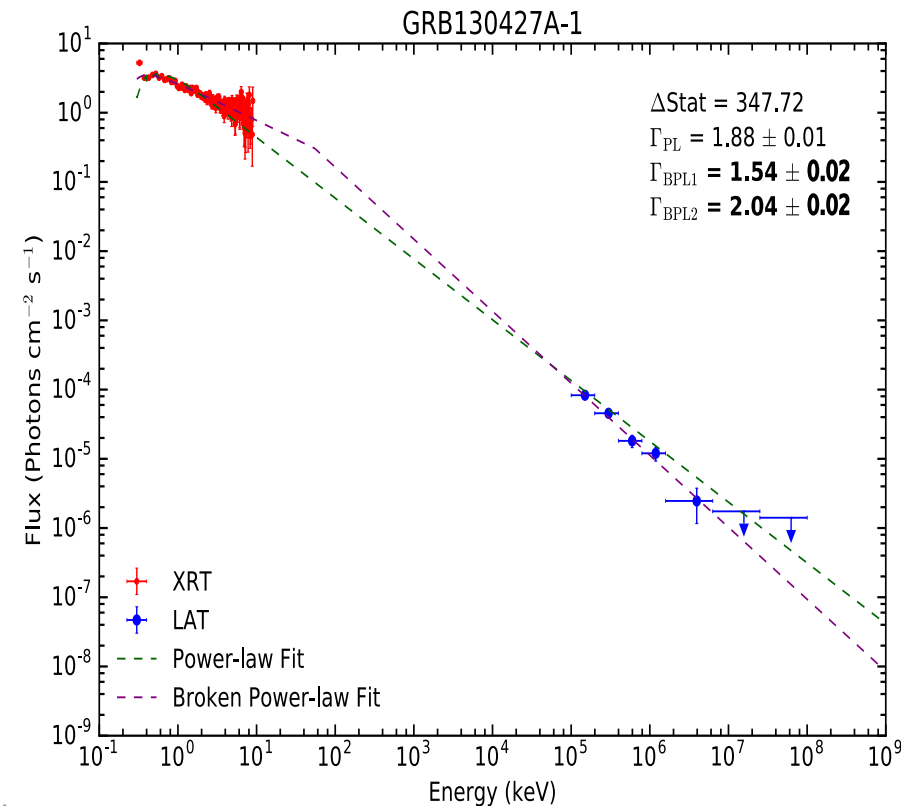
$$L_{\text{sync}}^{\text{iso}} \propto t^{-1}$$

# No Synchrotron Cutoff of GRB 130427A Seen in X-rays and Gamma-Rays

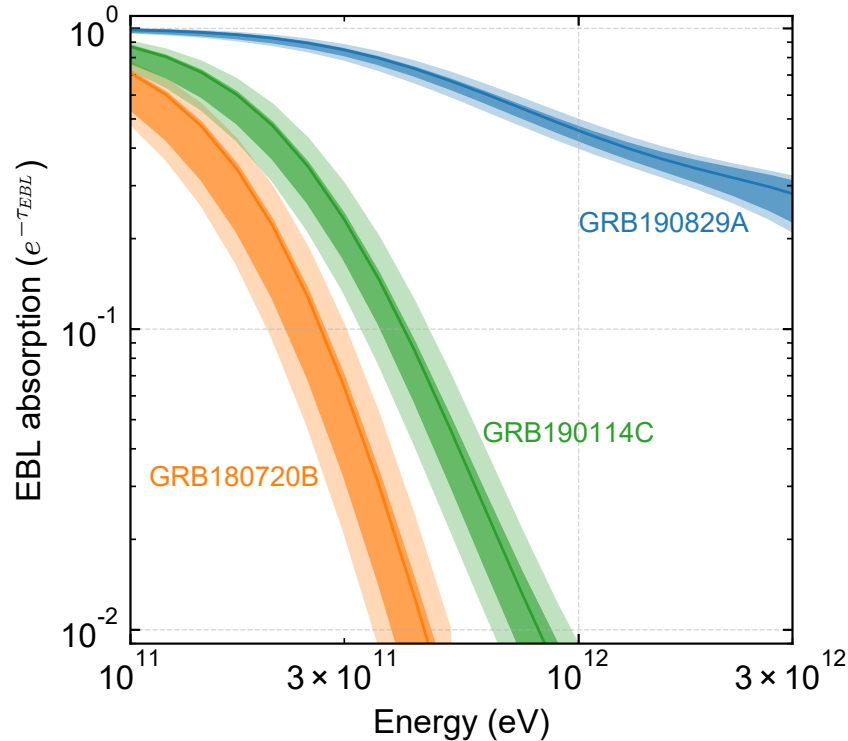
[Kouveliotou et al., ApJL 779 (2013)]



[Ajello et al., ApJ 863 138 (2018)]



# Energy Spectrum Information



The effect of the EBL on the (optically thin) attenuation for a nearby ( $z=0.08$ ) source for  $E_\gamma < 6$  TeV is a softening of the spectrum by around  $\Delta\Gamma \approx 0.5$ , starting around 250 GeV.

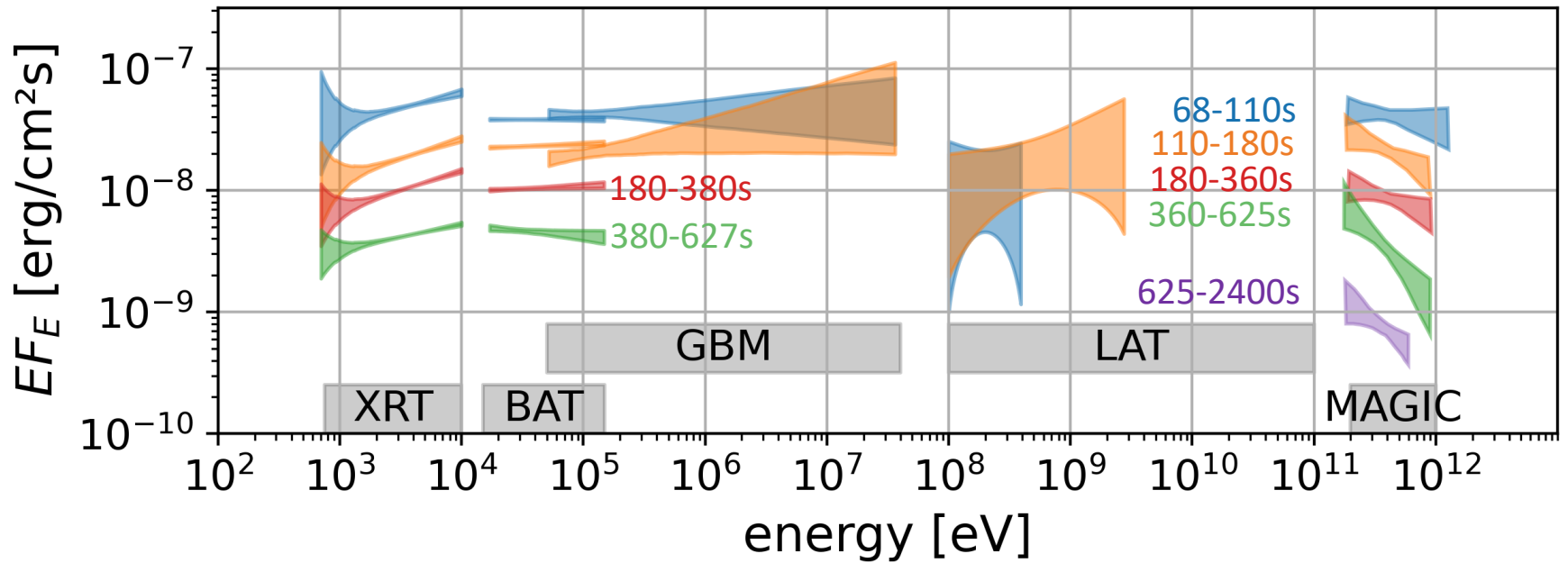
[HESS- A. Taylor, et al., Science 2021]

# Conclusions

- ◆ Fast shocks from massive energy release events are the most viable sources of extragalactic cosmic rays
- ◆ Synchrotron emission from long GRB tell us directly how efficient these sources operate as cosmic ray accelerators
- ◆ We are finally starting to probe the very high energy (TeV) gamma-ray emission from GRB, allowing us to start probing the magnetic fields in the source
- ◆ Whether a new component in the GRB spectrum is present remains unclear- the VHE GRB detections appear compatible with a continuation of the synchrotron emission beyond the expected supposed theoretical limit



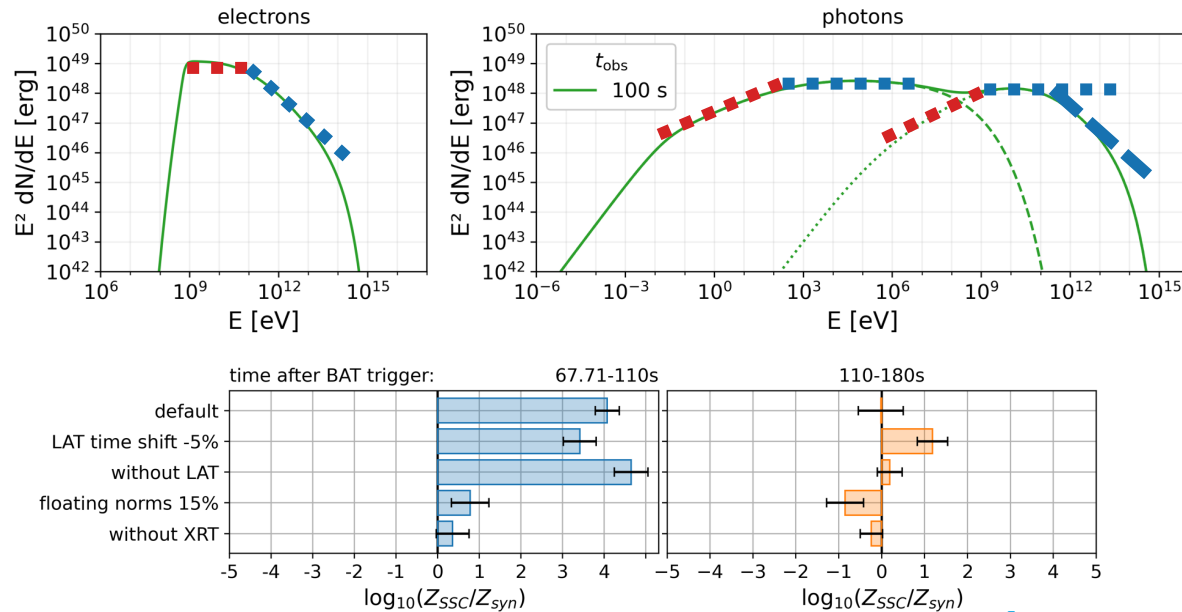
# GRB 190114C (Detected by MAGIC)



[Nature 575, 459-463 (2019)]

- remarkably flat over 9 orders of magnitude in energy!

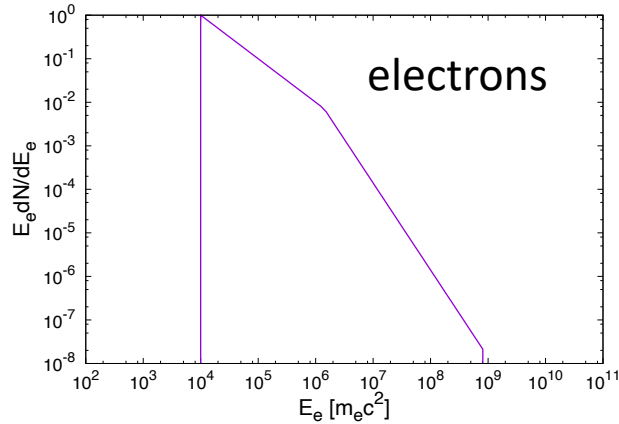
# Evidence for a New Component?



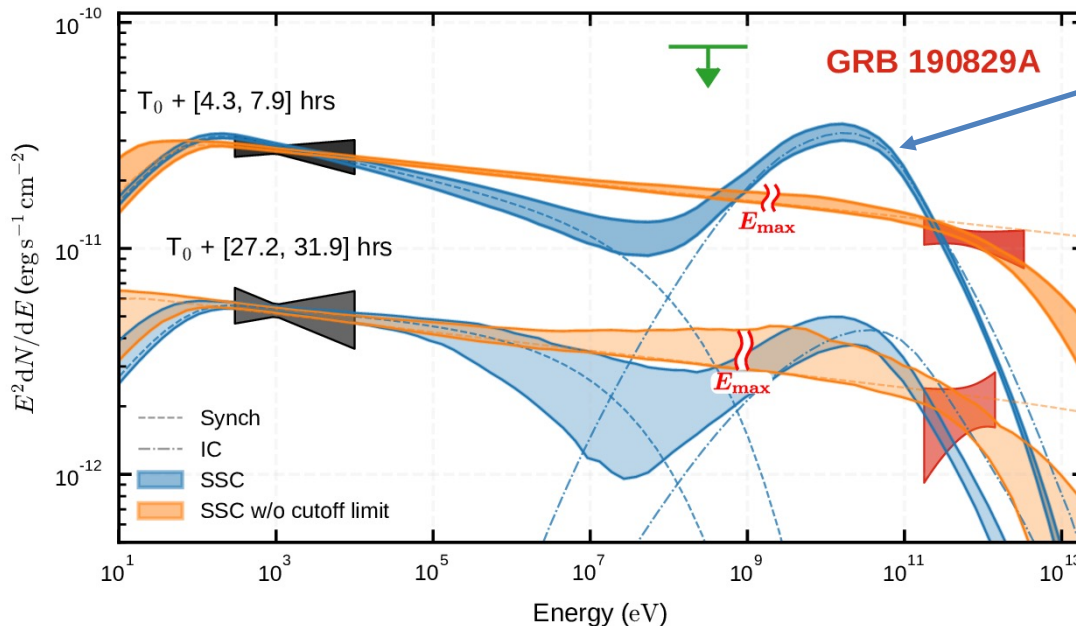
[M. Kinger et al., MNRAS 501 2023]

- SSC spectra are mirroring a smoothly BPL electron distribution
- We need more **bright, nearby** GRBs (without moonlight!)
- GRB 190114C shows no clear evidence for the onset of a new component

# GRB 190829A- Testing the “Standard” and Non-Standard VHE Emission Scenarios



MCMC fits to Night 1

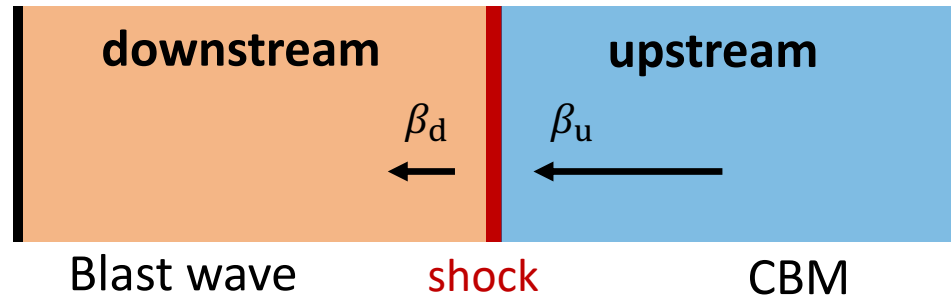


Synchrotron + SSC  
 $E_e > 400 \text{ GeV}$

Synchrotron Only  
 $E_e > 1 \text{ PeV}$

[HESS- A. Taylor, et al., Science 2021]

# Non-Rel. Hydro Shock- Downstream Partition of the Upstream Ram Pressure



[viewed in shock restframe]

$$p_d = \frac{3}{4} \rho_u \beta_u^2$$
$$\rho_d \beta_d^2 = \frac{1}{4} \rho_u \beta_u^2$$

A diagram illustrating the momentum flux vectors in the shock restframe. A horizontal blue arrow on the right points left towards the shock, labeled  $\rho_u \beta_u^2$ . Two blue arrows branch out from the shock point towards the left, representing the downstream momentum flux.

# Energy Transfer Efficiency for Synchrotron Emission

$$E_{\gamma}^{\text{sync}} \approx \frac{b}{3} E_e$$

$$b = \frac{4E_e E_{\gamma}^{\text{target}}}{(m_e c^2)^2}$$

$$E_{\gamma}^{\text{target}} = \left( \frac{B}{B_{\text{crit}}} \right) m_e c^2$$

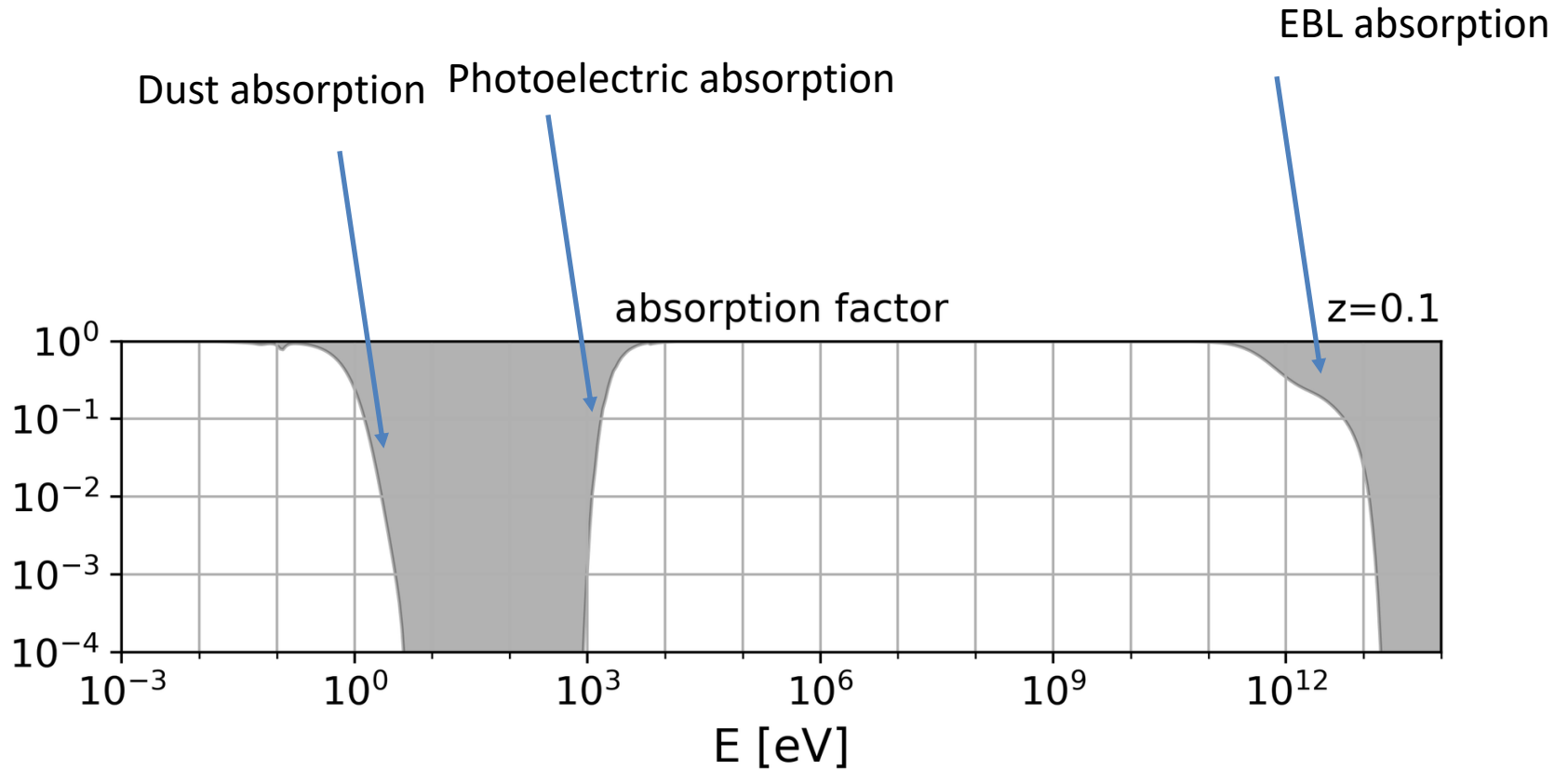
( $B_{\text{crit}} = 4 \times 10^{13}$  G)

$$E_e = 1 \text{ TeV}$$

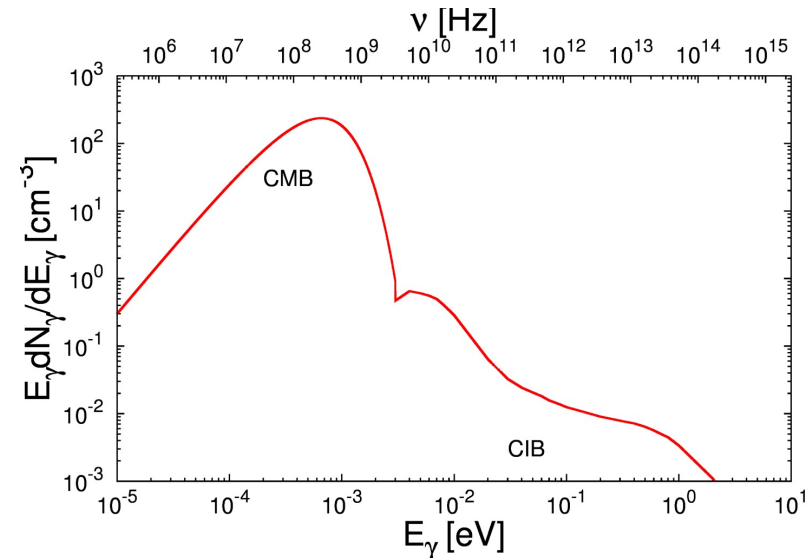
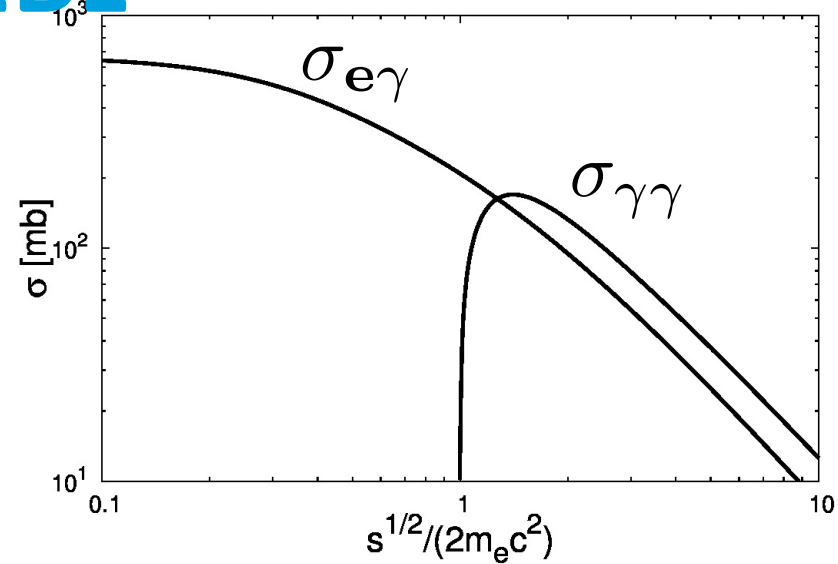
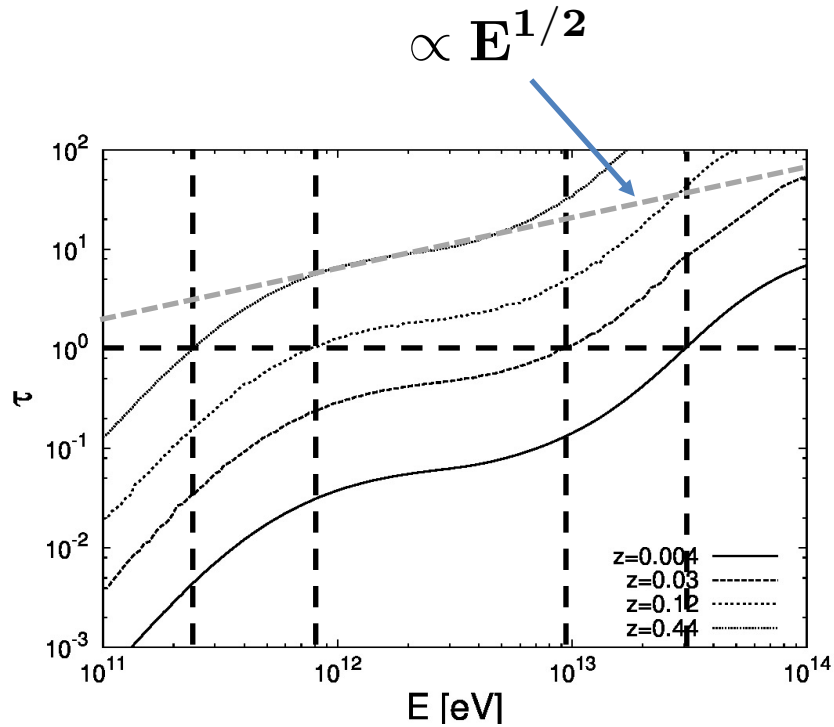
$$B = 1 \text{ G}$$

$$b \approx 10^{-7}$$

# The Observational Challenges for GRBs Absorption!

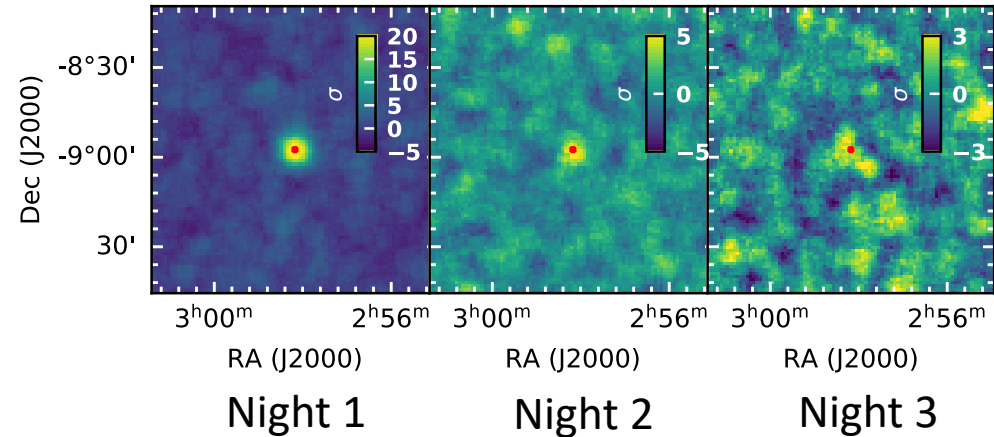
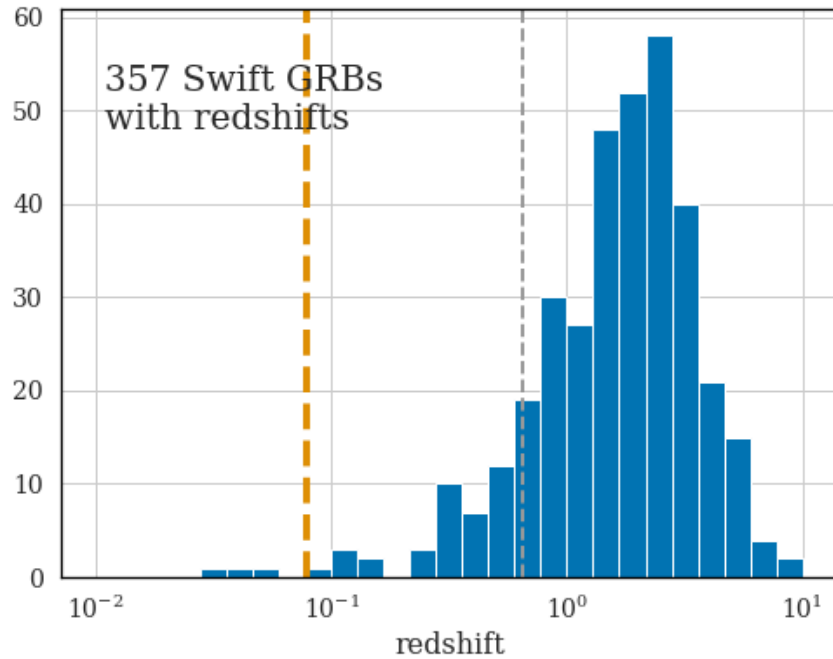


# Attenuation through Pair Production on the EBL



# HESS Detection of GRB 190829A

First detection of a GRB in VHE band for multiple nights

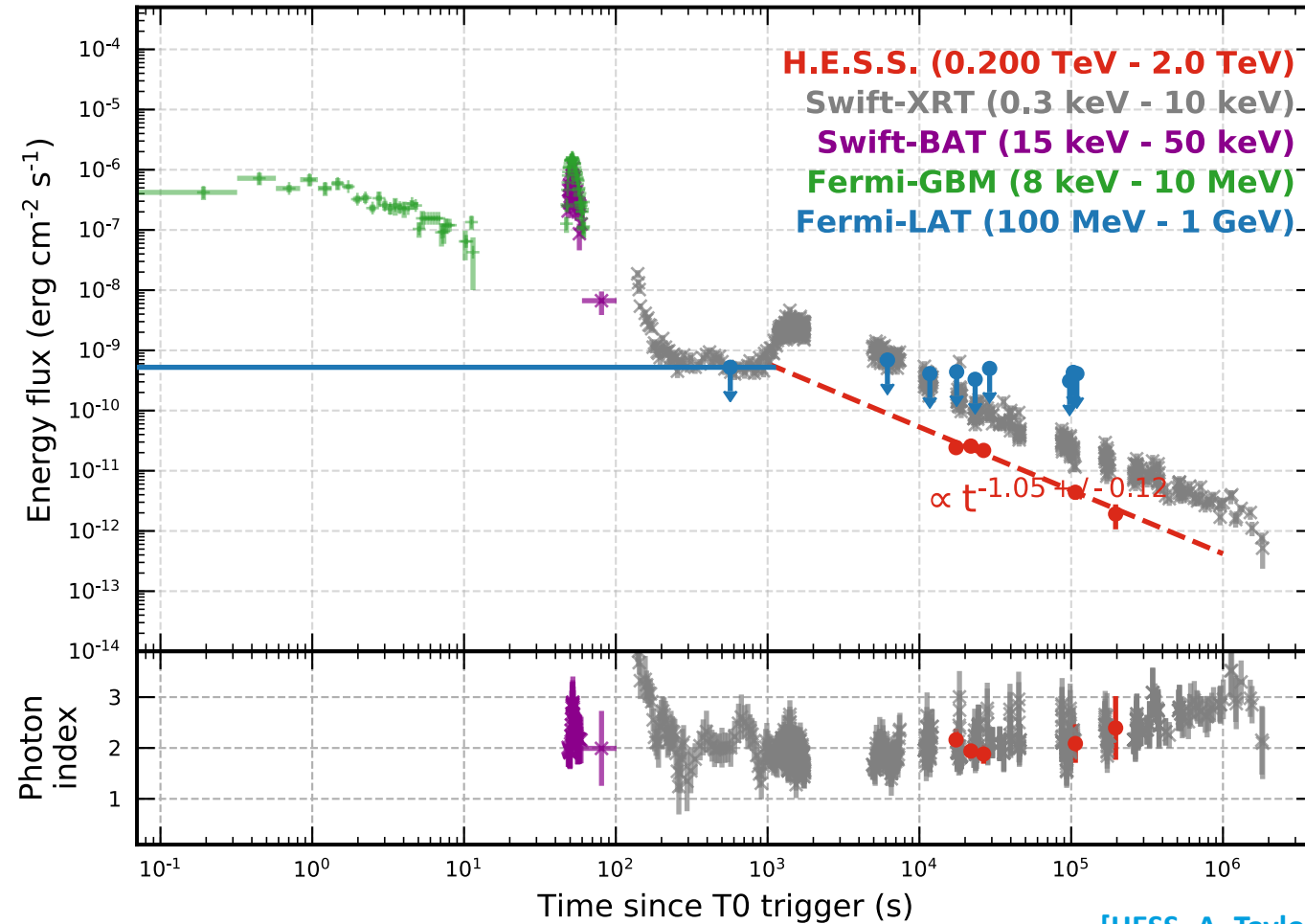


[HESS- A. Taylor, et al., Science 2021]

$t_{90}^{\text{GBM}} \sim 60 \text{ s}$ ,  $t_{90}^{\text{BAT}} \sim 60 \text{ s}$   
 $z = 0.078$



# MWL Energy Flux Lightcurve



GRB was not detected by Fermi-LAT

[HESS- A. Taylor, et al., Science 2021]

X-ray and Gamma-ray energy fluxes decay in a remarkably similar way-

$$\alpha_{\text{XRT}} = 1.09 \pm 0.04$$

$$\alpha_{\text{HESS}} = 1.05 \pm 0.12$$

DESY.

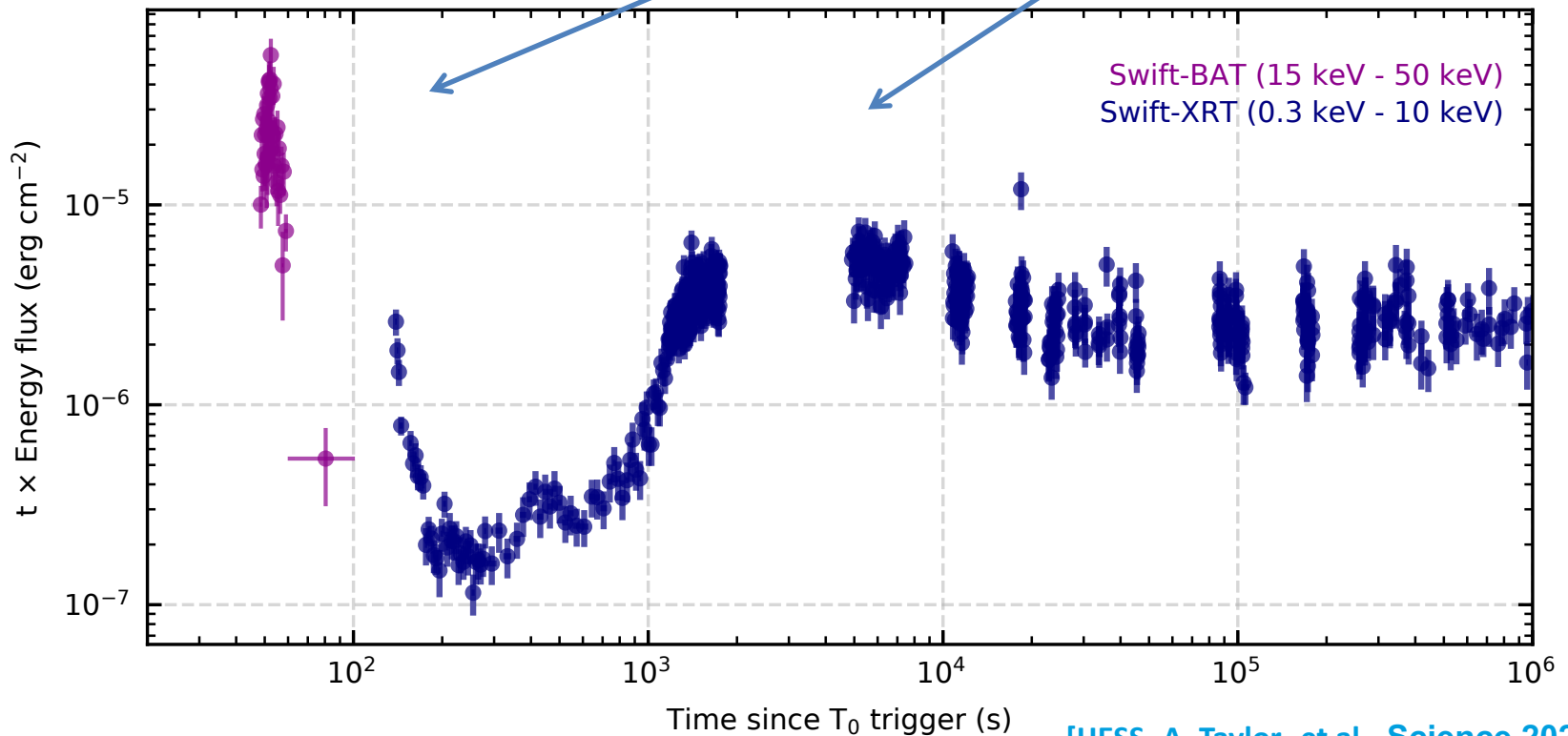
$$\mathbf{F}(t) \propto t^{-\alpha}$$

# When Does the Afterglow Fluence Saturate?

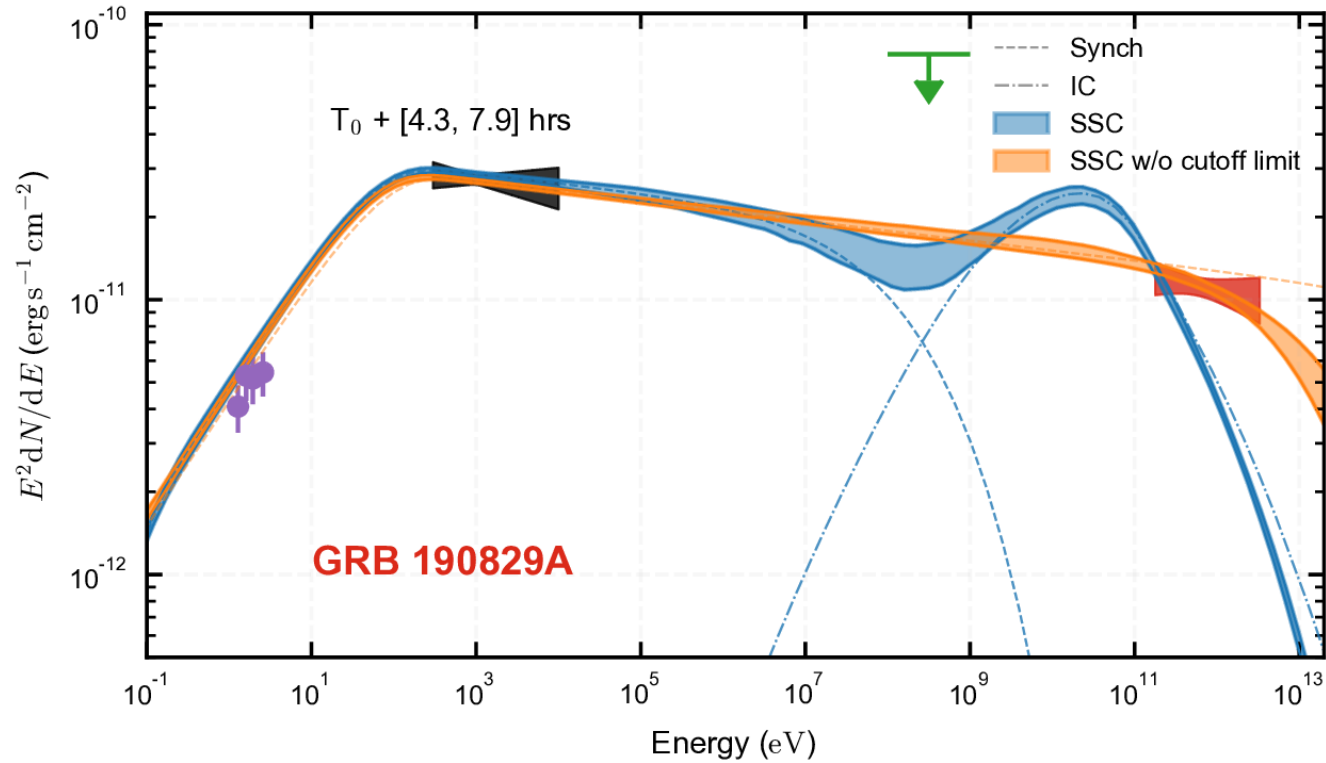
$$\text{fluence} = \int_{t_{\min}}^{t_{\max}} F(t) dt$$

$$E_{\text{GBM}}^{\text{iso}} = 2 \times 10^{50} \text{ erg}$$

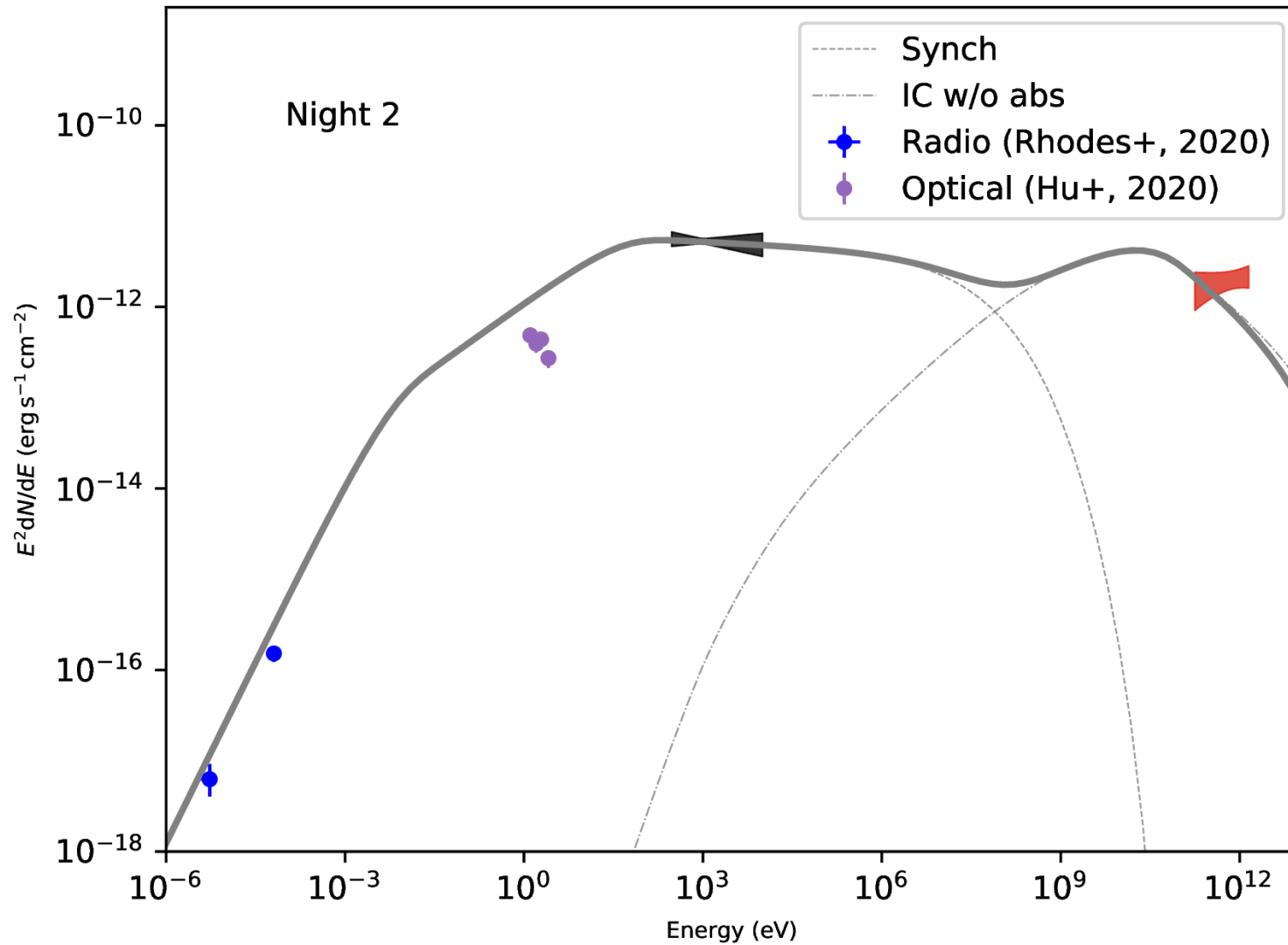
$$E_{\text{XRT}}^{\text{iso}} = 6 \times 10^{50} \text{ erg}$$



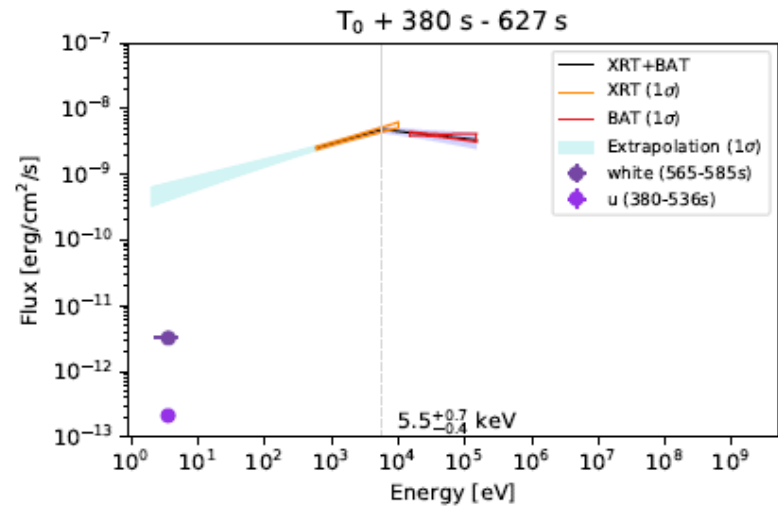
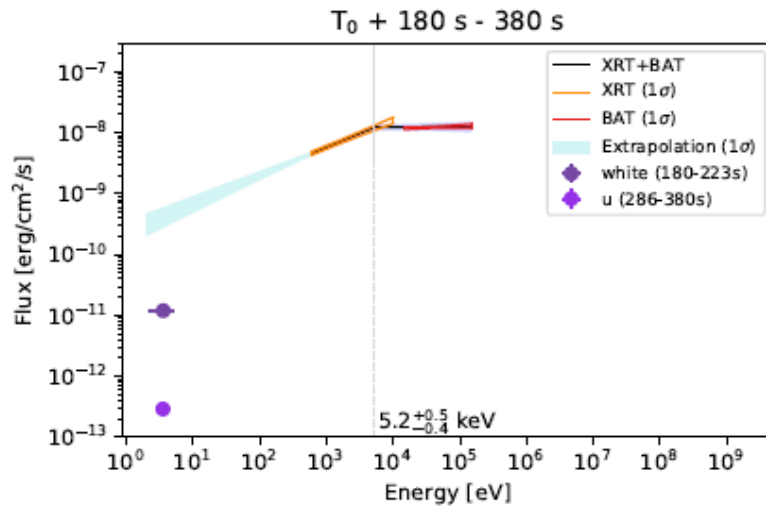
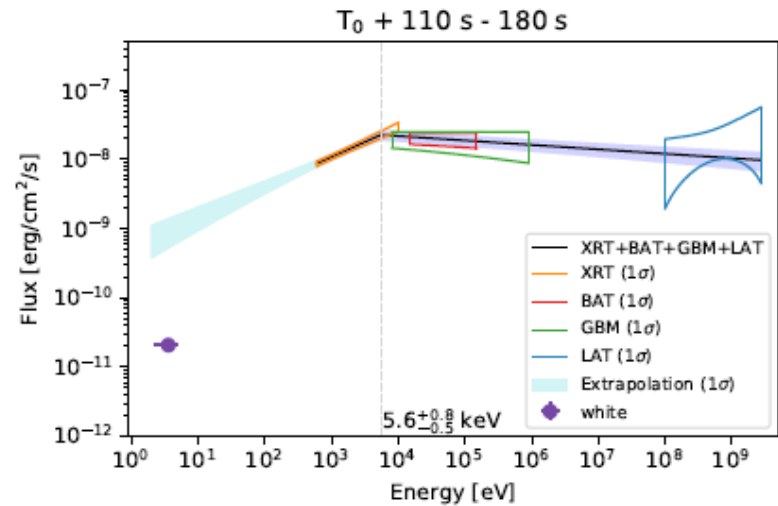
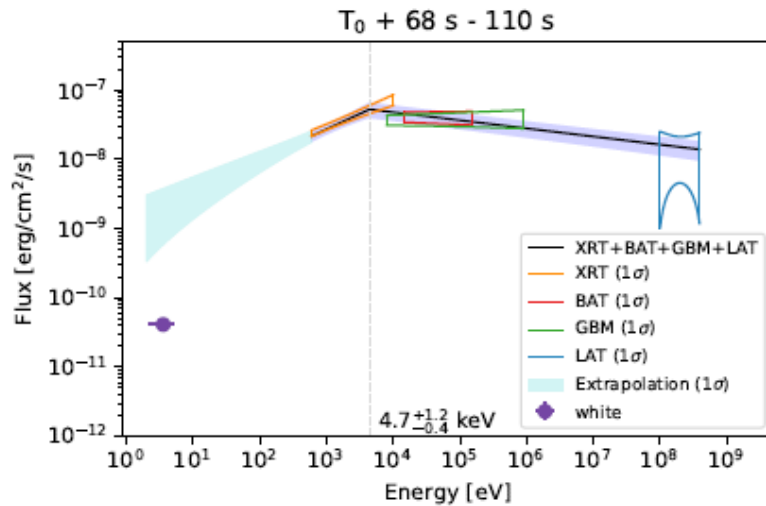
# GRB 190829A- Optical Data



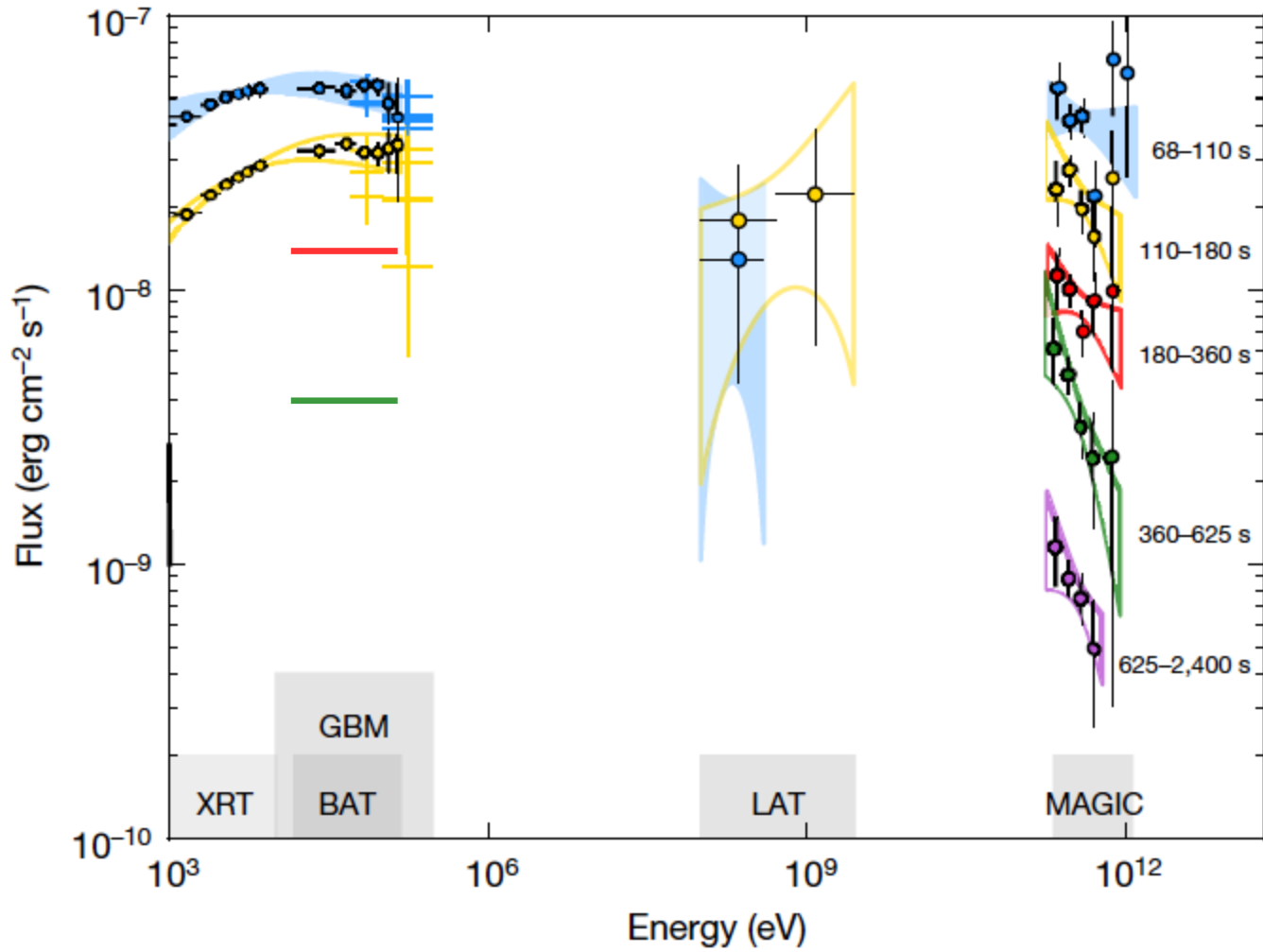
# GRB 190829A- Radio Data



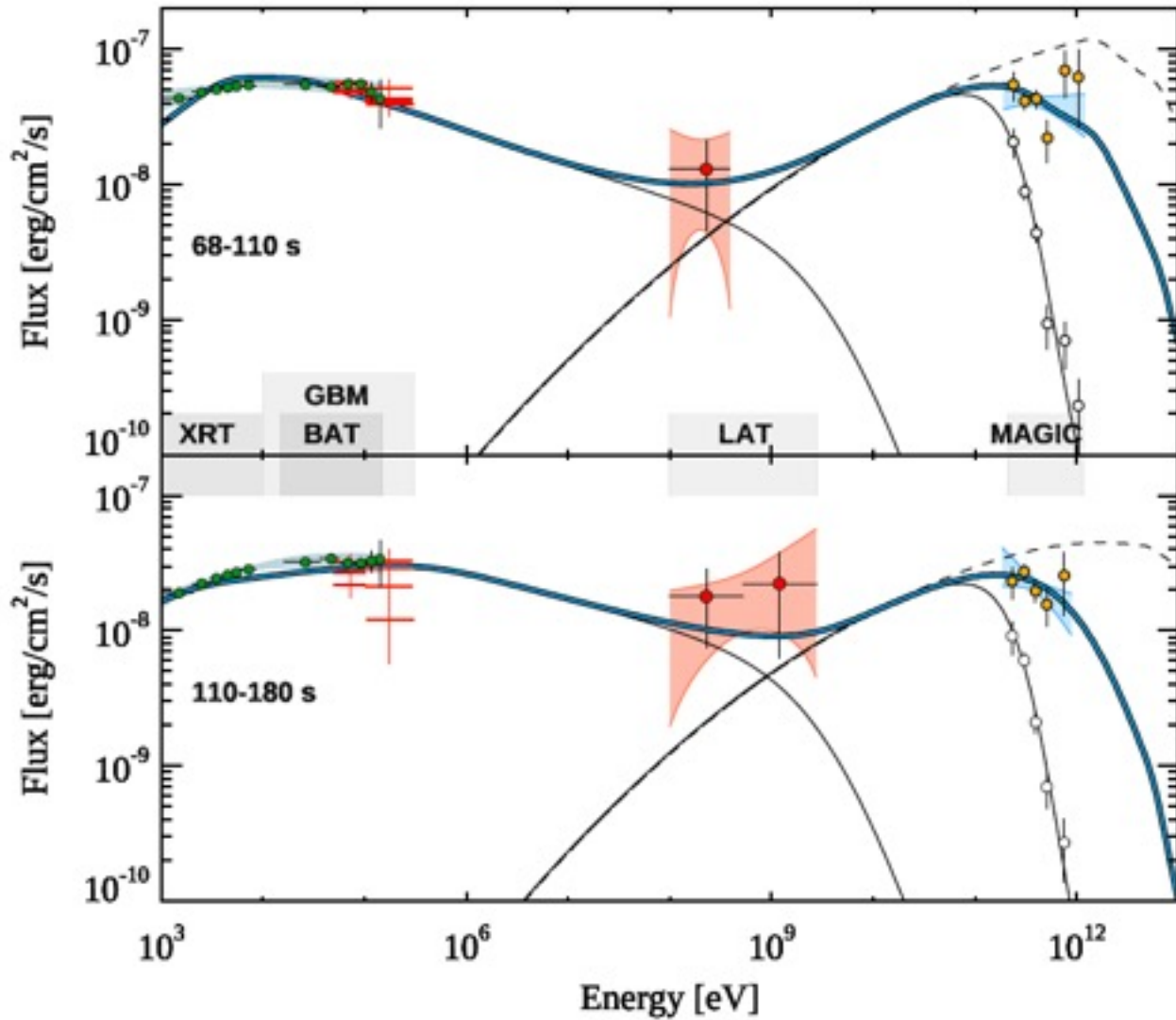
# GRB 190114C



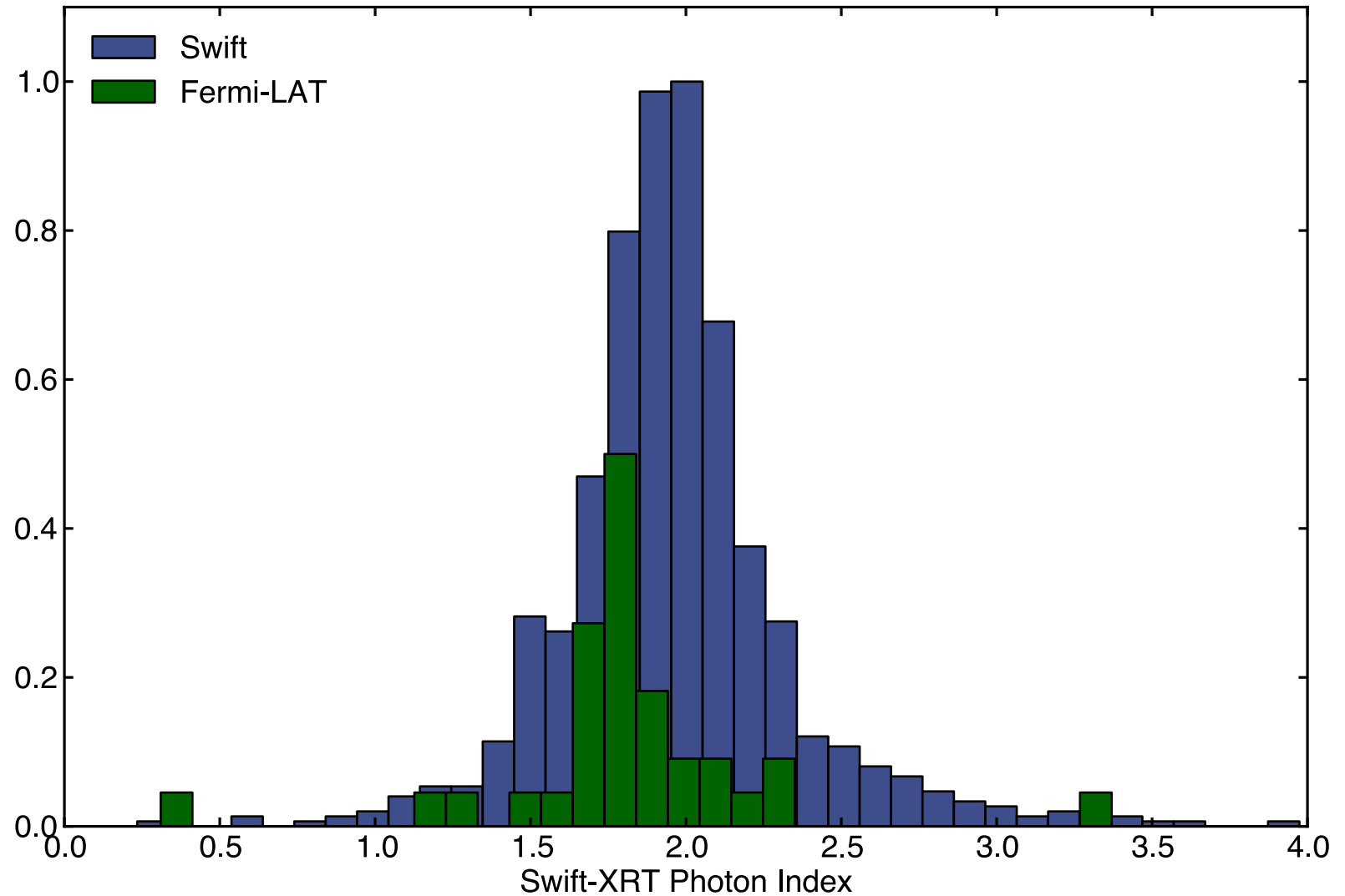
# GRB 190114C



# GRB 190114C



# Swift XRT Photon Index Distribution

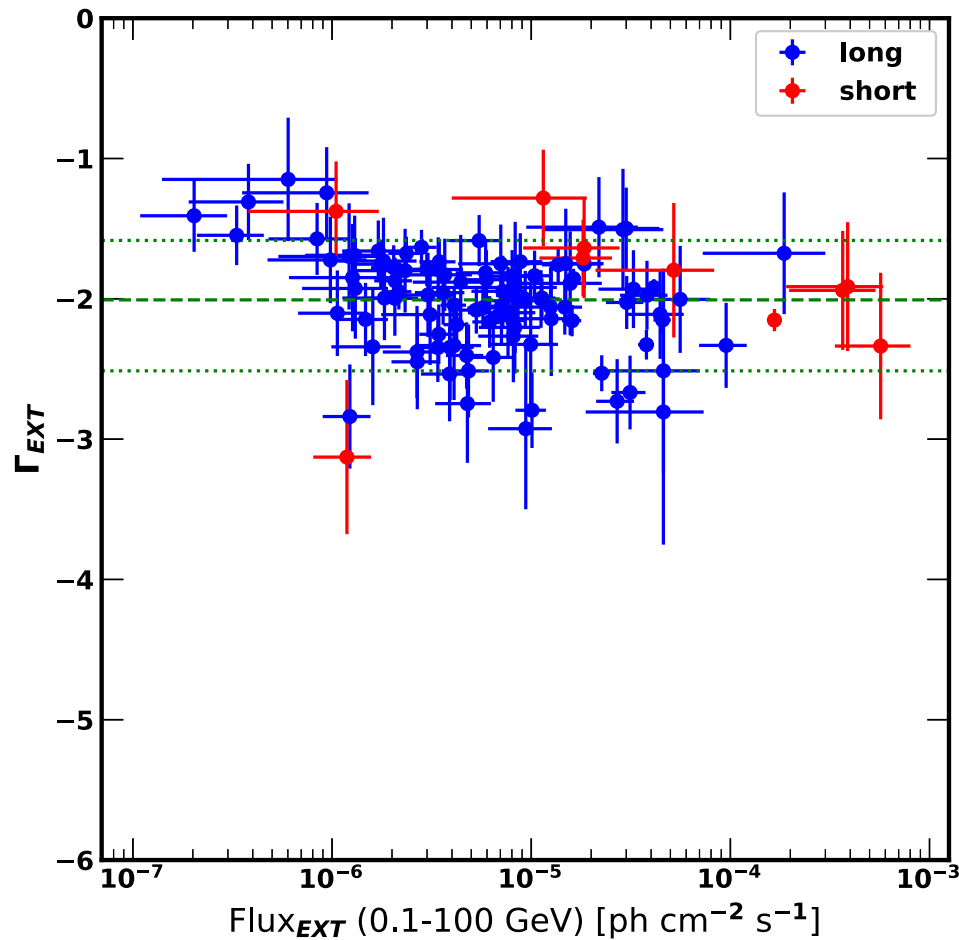


[Ajello et al., Ap. J., 863 138, 2018]

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# Fermi-LAT Photon Index Distribution



[Ajello et al., Ap. J., 878:52, 2019]

# Relativistic Shocks

Momentum Flux:

$$\mathbf{p}_1 + \left( \frac{\gamma}{\gamma - 1} \mathbf{p}_1 + \rho_1 \right) \beta_1^2 \Gamma_1^2 = \mathbf{p}_2 + \left( \frac{\gamma}{\gamma - 1} \mathbf{p}_2 + \rho_2 \right) \beta_2^2 \Gamma_2^2$$

Energy Flux:

$$\left( \frac{\gamma}{\gamma - 1} \mathbf{p}_1 + \rho_1 \right) \beta_1 \Gamma_1^2 = \left( \frac{\gamma}{\gamma - 1} \mathbf{p}_2 + \rho_2 \right) \beta_2 \Gamma_2^2$$

# Cold Relativistic Shocks

Momentum Flux:

$$\rho_1 \beta_1^2 \Gamma_1^2 = \mathbf{p}_2 + \left( \frac{\gamma}{\gamma - 1} \mathbf{p}_2 + \rho_2 \right) \beta_2^2 \Gamma_2^2$$

$$\rho_1 \beta_1^2 \Gamma_1^2 - \rho_2 \beta_2^2 \Gamma_2^2 = \mathbf{p}_2 \left[ \mathbf{1} + \left( \frac{\gamma}{\gamma - 1} \right) \beta_2^2 \Gamma_2^2 \right]$$

Energy Flux:

$$\rho_1 \beta_1 \Gamma_1^2 = \left( \frac{\gamma}{\gamma - 1} \mathbf{p}_2 + \rho_2 \right) \beta_2 \Gamma_2^2$$

$$\rho_1 \beta_1 \Gamma_1 (\Gamma_1 - 1) = \frac{\gamma}{\gamma - 1} \mathbf{p}_2 \beta_2 \Gamma_2^2 + \rho_2 \beta_2 \Gamma_2 (\Gamma_2 - 1)$$



# Relativistic Shocks

Momentum Flux:

$$\frac{\mathbf{p}_2}{\Gamma_1^2 \beta_1^2 \rho_1} \left[ \mathbf{1} + \Gamma_2^2 \beta_2^2 \left( \frac{\gamma}{\gamma - 1} \right) \right] = \left( \mathbf{1} - \frac{\Gamma_2 \beta_2}{\Gamma_1 \beta_1} \right)$$

Energy Flux:

$$\left( \frac{\gamma}{\gamma - 1} \right) \frac{\Gamma_2^2 \mathbf{p}_2 \beta_2}{\Gamma_1^2 \rho_1 \beta_1} = \left( \mathbf{1} - \frac{(\Gamma_2 - 1)}{(\Gamma_1 - 1)} \right)$$

# Relativistic Shocks

$$\frac{1 - \frac{\Gamma_2 \beta_2}{\Gamma_1 \beta_1}}{1 + \Gamma_2^2 \beta_2^2 \frac{\gamma}{\gamma - 1}} = \frac{1 - \frac{\Gamma_2 - 1}{\Gamma_1 - 1}}{\Gamma_2^2 \beta_2^2 \frac{\gamma}{\gamma - 1}}$$

$$1 + \Gamma_2^2 \beta_2^2 \left( \frac{\gamma}{\gamma - 1} \right) = \Gamma_2^2 \beta_2^2 \left( \frac{\gamma}{\gamma - 1} \right)$$

$$(\beta_2 - 1)(\beta_2 - (\gamma - 1)) = 0$$

Eg:  $\gamma = \frac{4}{3} \quad \rightarrow \quad \frac{\beta_2}{\beta_1} = \frac{1}{3}$