

NEUTRINOLESS DOUBLE BETA DECAY : PRESENT AND FUTURE

Ben Jones

University of Texas at Arlington

Rencontres De Blois 2023



UNIVERSITY OF
TEXAS
ARLINGTON

SM as a low energy effective theory:

$$L = L_{SM} + \frac{1}{E_{new}} L_1 + \frac{1}{E_{new}^2} L_2 + \dots$$

SM is a low energy effective theory:

$$L = L_{SM} + \frac{1}{E_{new}} L_1 + \frac{1}{E_{new}^2} L_2 + \dots$$

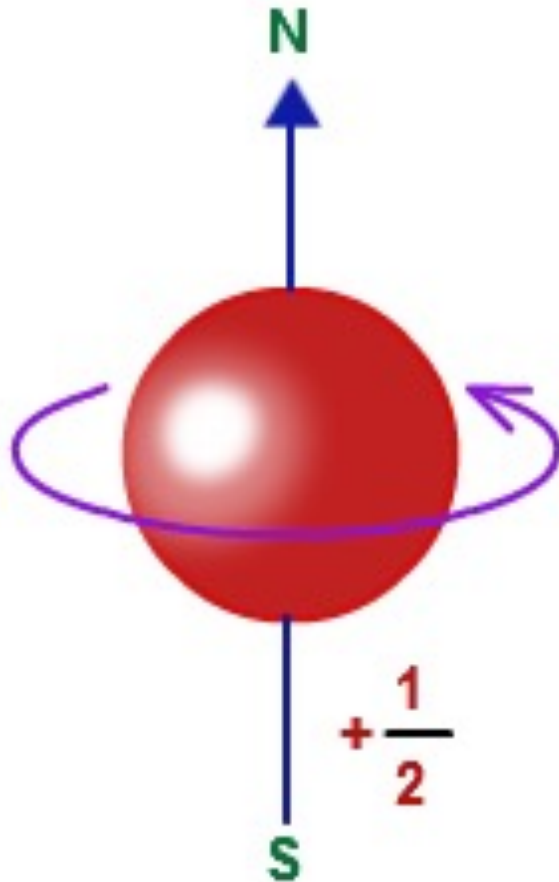
- The only dimension-5 operator one can add obeying SM gauge symmetry:

$$\frac{L_1}{E_{new}} = y_{ij} \frac{\nu^i H \nu^j H}{E_{new}}$$

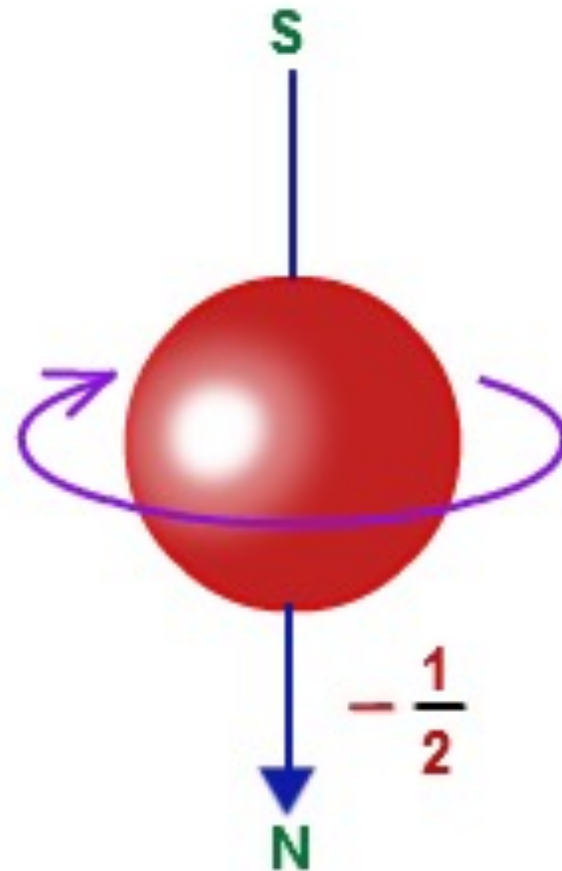
Weinberg 1979.

- This term does ~one observable thing - it makes neutrinos Majorana particles, with mass suppressed by new physics scale.
- And it makes the theory non-renormalizable – implying there must be something else at high scale too (see also: Seesaw)

Majorana Neutrinos in a Nutshell



Behaves like matter



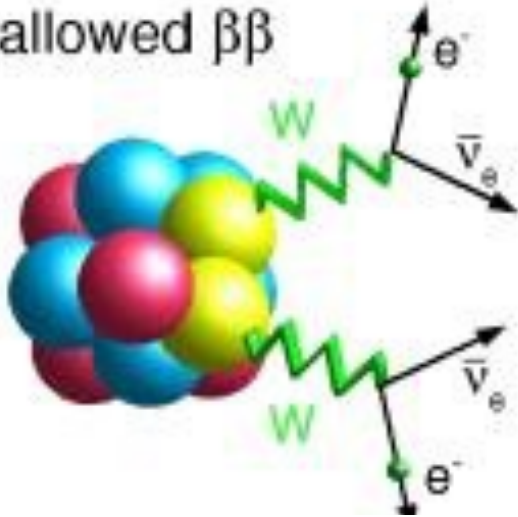
Behaves like antimatter

Robust observation of Majorana neutrinos would tell us 5 things about nature before breakfast:

- 1) Lepton number conservation is violated.**
- 2) Massive fermions exist that are neither matter or antimatter but something else (Majorana fermions)**
- 3) The SM with the Majorana term is non-renormalizable \rightarrow SM is definitely a low energy effective theory.**
- 4) There are other mass generating mechanisms in nature beyond the Higgs mechanism (though the Higgs may be involved, nonetheless).**
- 5) Majorana neutrinos are a prediction of the theory of Leptogenesis that may generate observed matter/anti-matter asymmetry of the Universe (given leptonic CPV – see also: HyperK / DUNE)**

Double beta decay

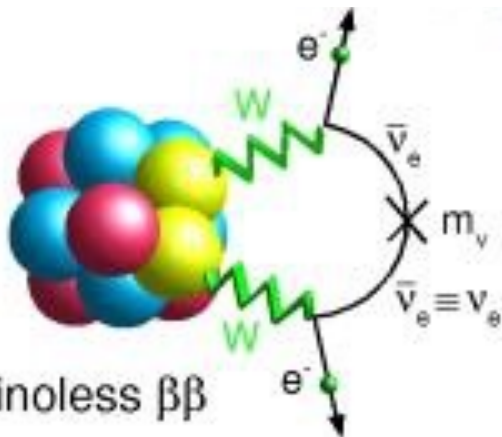
allowed $\beta\beta$



**A known standard model process
and an important calibration tool**

$$T_{\frac{1}{2}} = 10^{19} - 10^{21} \text{ yr}$$

Final state: $e^- e^- \bar{\nu}_e \bar{\nu}_e$



neutrinoless $\beta\beta$

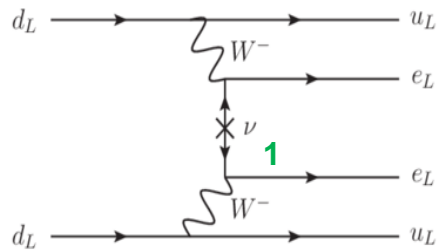
**Observation would prove that the
neutrino is a Majorana fermion**

$$T_{\frac{1}{2}} = [G \times |M|^2 \times m_{\beta\beta}^2]^{-1}$$

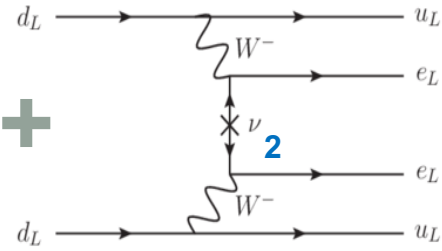
Final state: $e^- e^-$

The "minimal" mechanism:

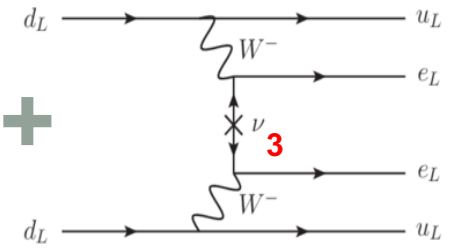
Γ α



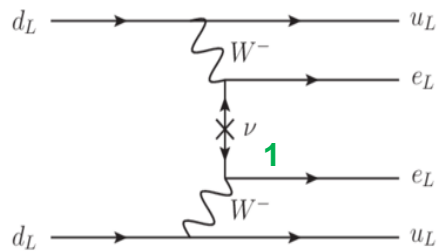
+



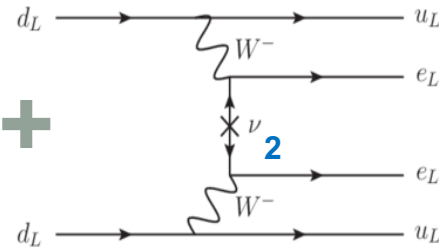
+



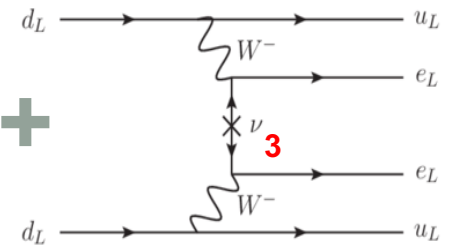
2

$\Gamma \alpha$


+



+



2

=

 $(U_{e1})^2 m_1$

+

 $(U_{e2})^2 m_2$

+

 $(U_{e3})^2 m_3$

2

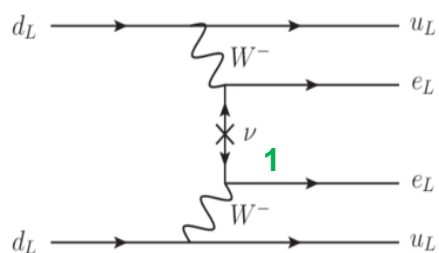
$$\Gamma_\alpha = \left| \begin{array}{c} d_L \longrightarrow \text{---} u_L \\ \text{---} W^- \text{---} \\ \text{---} e_L \\ \text{---} \nu_1 \\ \text{---} e_L \\ \text{---} W^- \text{---} \\ d_L \longrightarrow \text{---} u_L \end{array} \right. + \left| \begin{array}{c} d_L \longrightarrow \text{---} u_L \\ \text{---} W^- \text{---} \\ \text{---} e_L \\ \text{---} \nu_2 \\ \text{---} e_L \\ \text{---} W^- \text{---} \\ d_L \longrightarrow \text{---} u_L \end{array} \right. + \left| \begin{array}{c} d_L \longrightarrow \text{---} u_L \\ \text{---} W^- \text{---} \\ \text{---} e_L \\ \text{---} \nu_3 \\ \text{---} e_L \\ \text{---} W^- \text{---} \\ d_L \longrightarrow \text{---} u_L \end{array} \right. \Bigg|^2$$

$$= \left| (U_{e1})^2 m_1 + (U_{e2})^2 m_2 + (U_{e3})^2 m_3 \right|^2$$

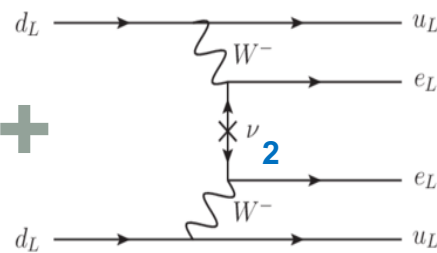
$$= \left| c_{12}^2 c_{13}^2 e^{2i\alpha} m_1 + c_{12}^2 s_{12}^2 e^{2i\beta} m_2 + s_{13}^2 m_3 \right|^2$$

$$\begin{aligned}
 \Gamma_\alpha &= \left| \begin{array}{c} \text{Diagram 1} \\ \text{Diagram 2} \\ \text{Diagram 3} \end{array} \right|^2 \\
 &= \left| (\mathbf{U}_{e1})^2 m_1 + (\mathbf{U}_{e2})^2 m_2 + (\mathbf{U}_{e3})^2 m_3 \right|^2 \\
 &= \left| c_{12}^2 c_{13}^2 e^{2i\alpha} m_1 + c_{12}^2 s_{12}^2 e^{2i\beta} m_2 + s_{13}^2 m_3 \right|^2 \\
 &= \left| m_{\beta\beta} \right|^2 \quad T_{\frac{1}{2}} = [G \times |M|^2 \times m_{\beta\beta}^2]^{-1}
 \end{aligned}$$

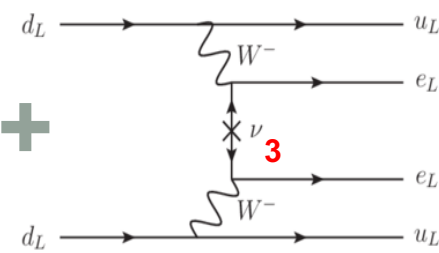
The diagrams show the exchange of a \$W^-\$ boson between two \$d_L\$ quarks and two \$e_L\$ neutrinos. The neutrinos are labeled \$\nu_1\$, \$\nu_2\$, and \$\nu_3\$ in green, blue, and red respectively. The diagrams are summed and then squared to give the final expression for \$\Gamma_\alpha\$.

$\Gamma \alpha$


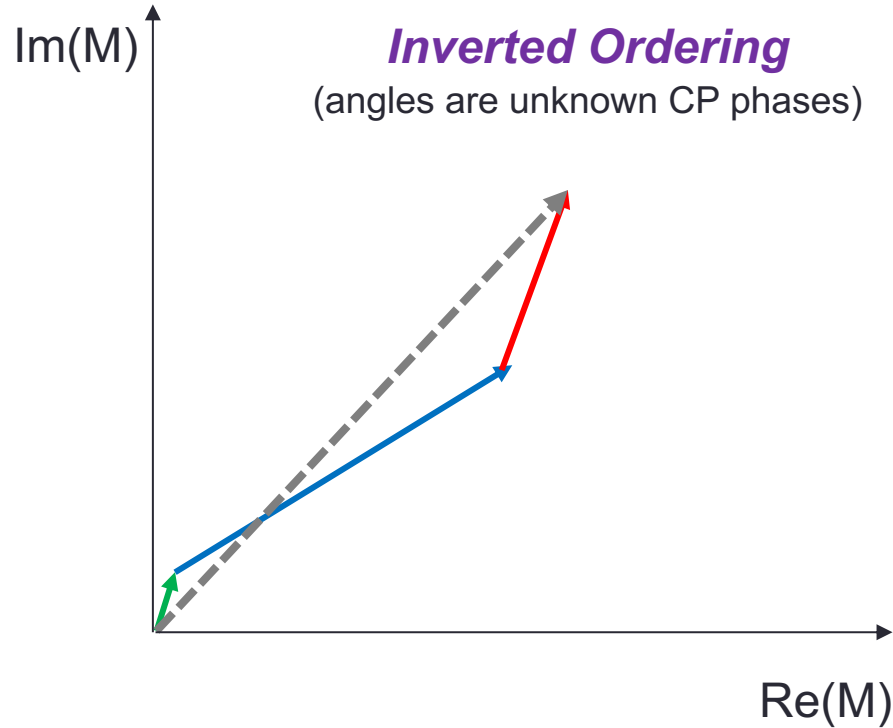
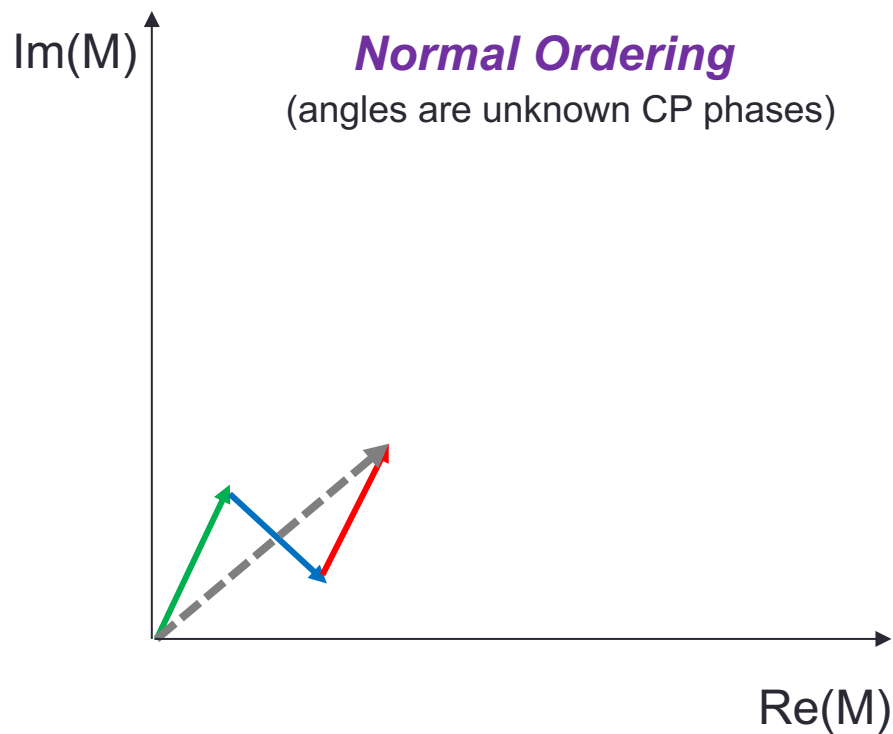
+

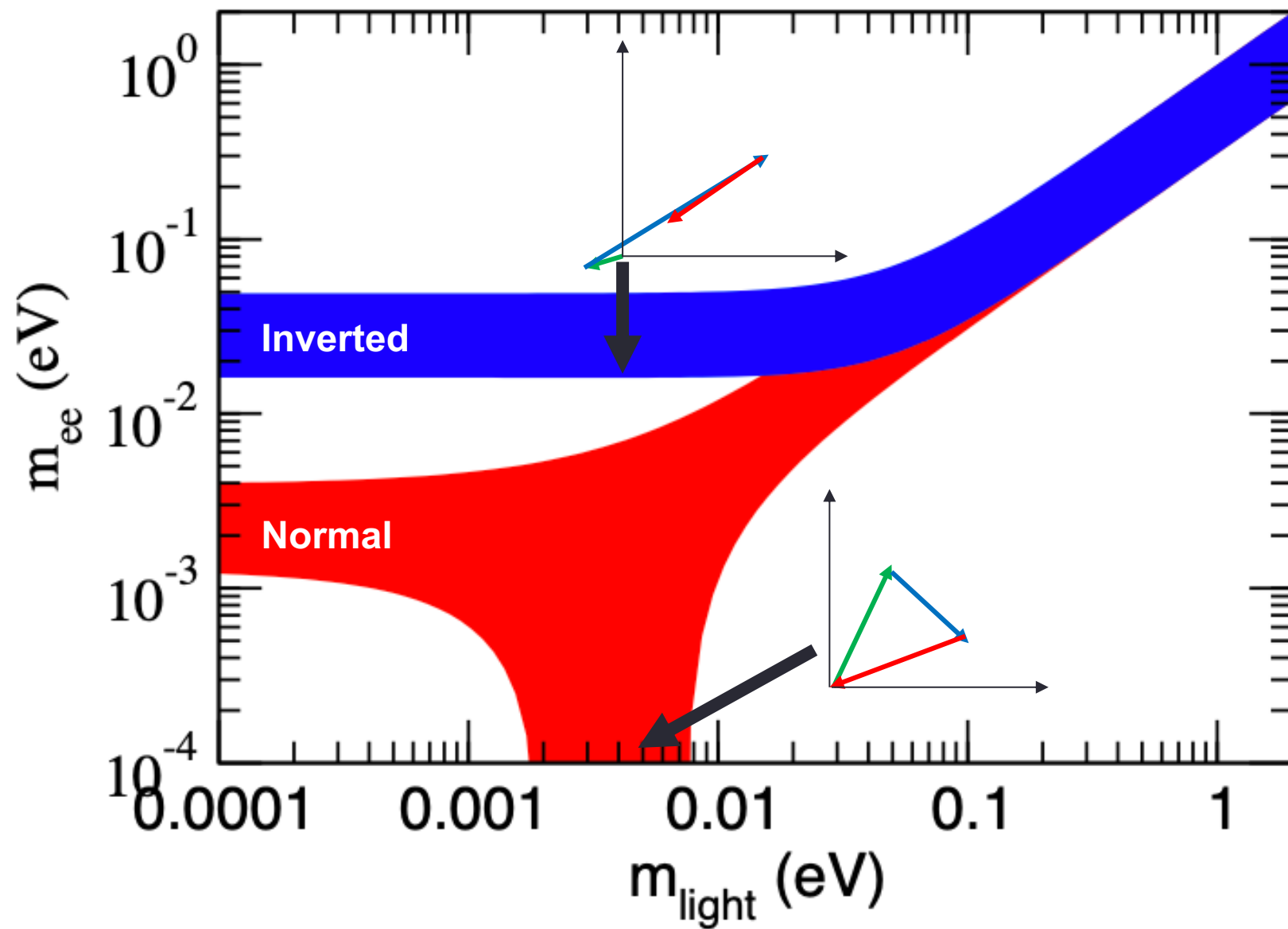


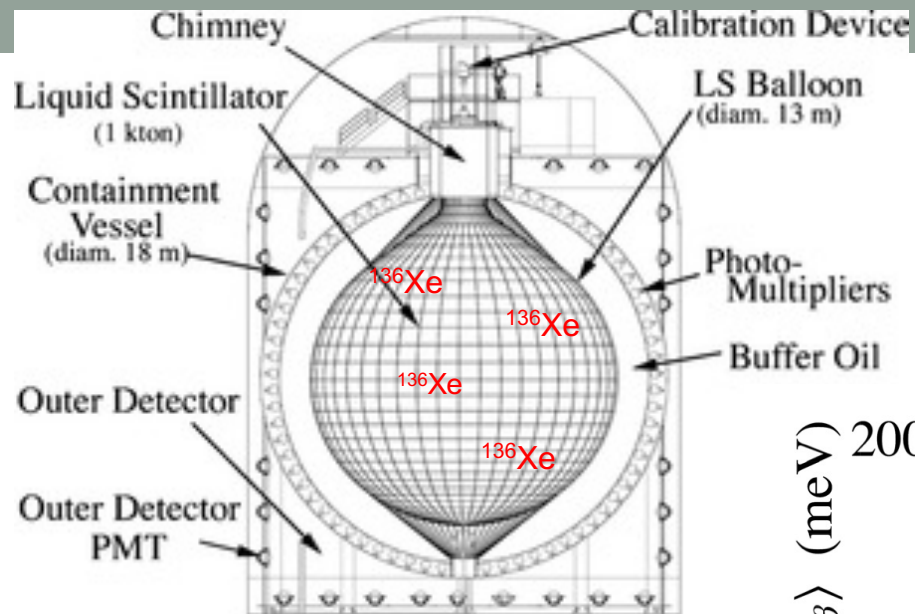
+



2





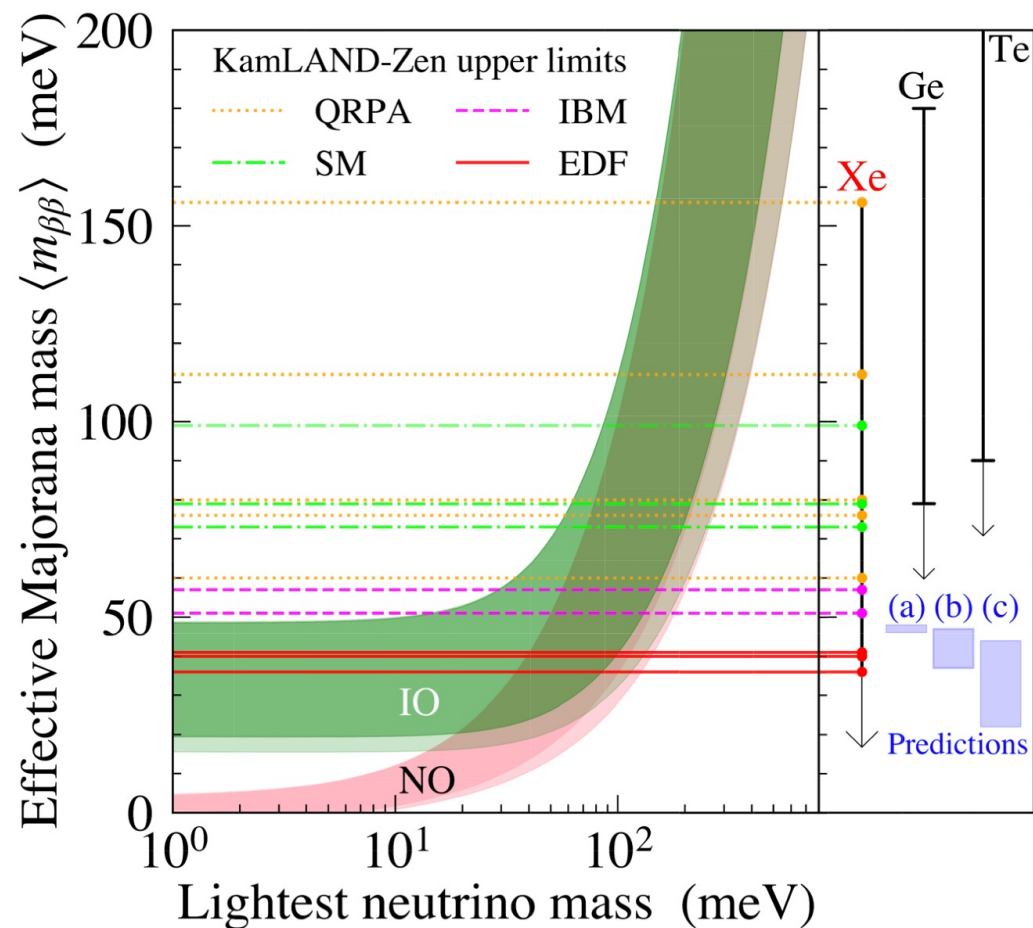


Current limits

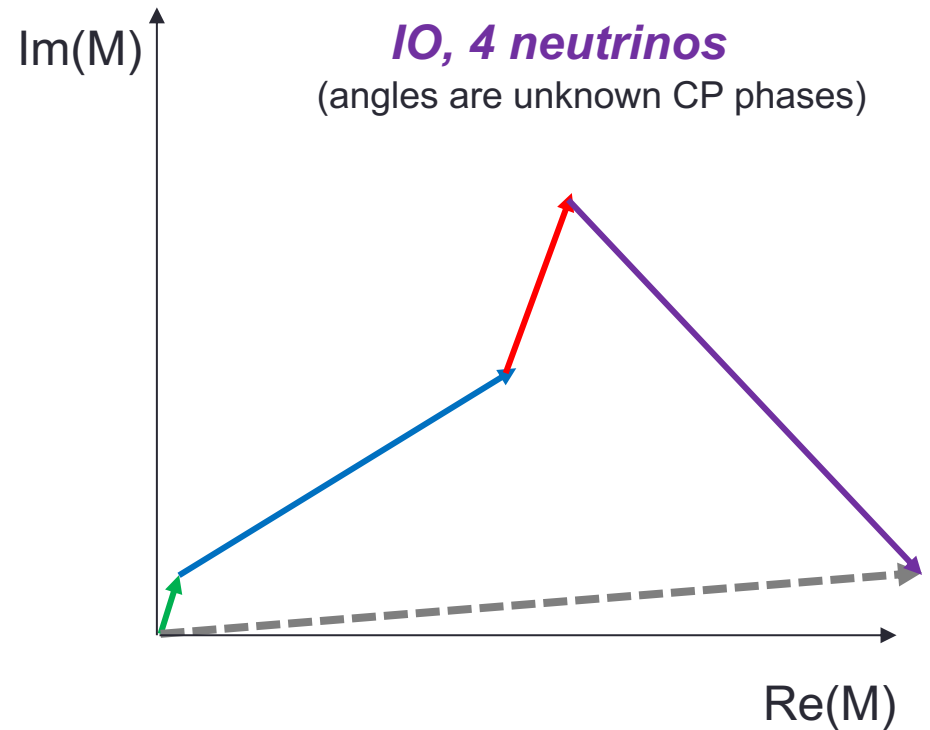
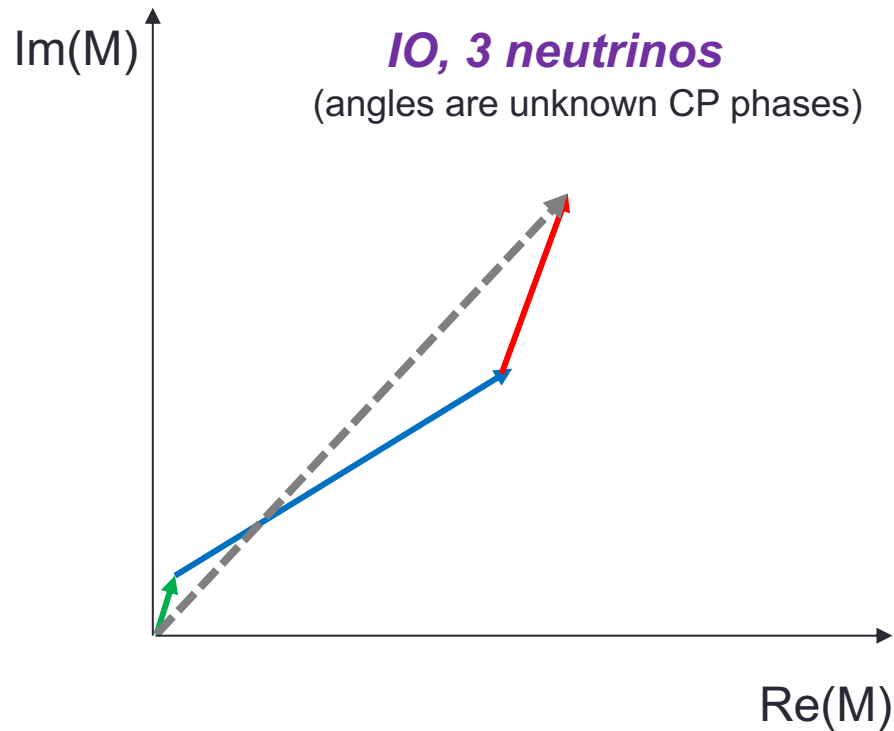
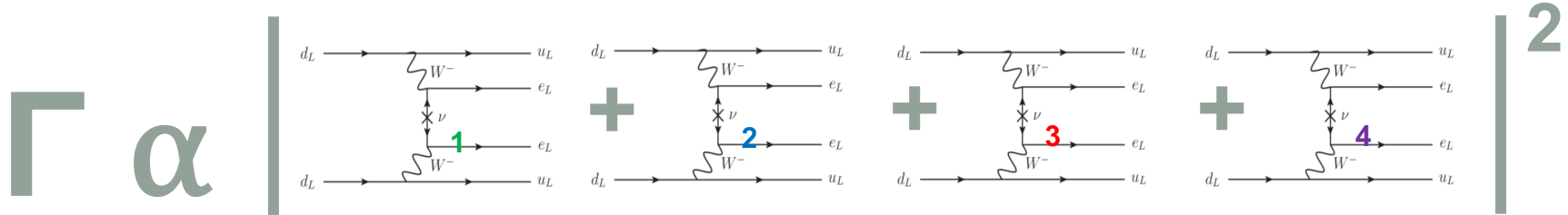
KZ leads with:
 $T_{1/2} > 2.3 \times 10^{26} \text{ yr}$

Leading limits in each isotope:

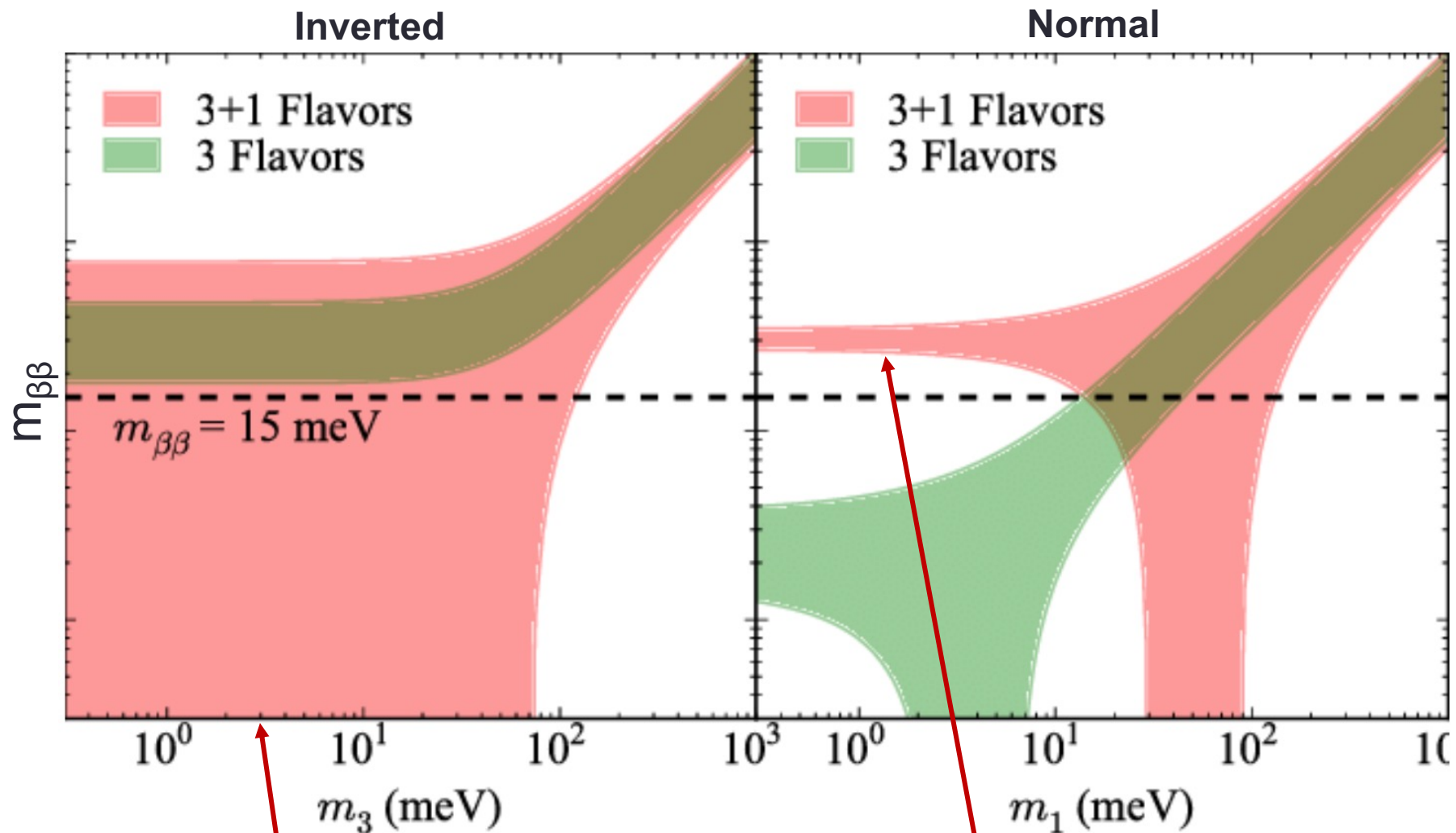
Experiment	Isotope	Exposure [kg yr]	$T_{1/2}^{0\nu}$ [10^{25} yr]	$m_{\beta\beta}$ [meV]
Gerda	^{76}Ge	127.2	18	79-180
Majorana	^{76}Ge	26	2.7	200-433
KamLAND-Zen	^{136}Xe	970	23	36-156
EXO-200	^{136}Xe	234.1	3.5	93-286
CUORE	^{130}Te	1038.4	2.2	90-305



Non-minimal mechanisms:



With a sterile neutrino:



Now you can close the loop here after all

But this gets less difficult!

Higher order operators

$$L = L_{SM} - \frac{1}{E_{new}} L_1 + \frac{1}{E_{new}^2} L_2 + \dots$$

NDBD can also come from up here

Weinberg operator ("minimal" model)

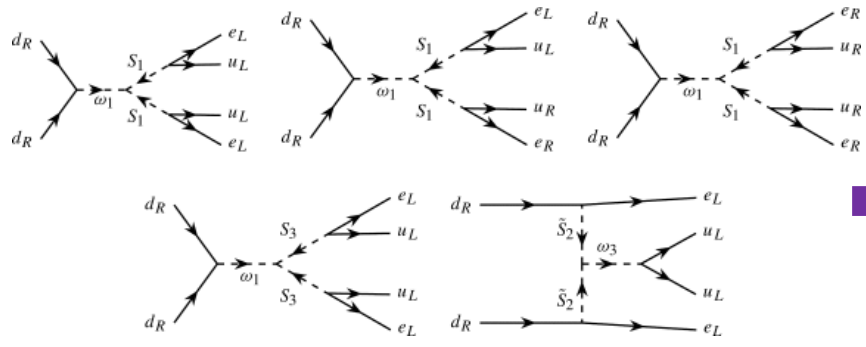
Class 1	$\psi^2 H^4$	Class 5	$\psi^4 D$
\mathcal{O}_{LH}	$\epsilon_{ij} \epsilon_{mn} (L_i^T C L_m) H_j H_n (H^\dagger H)$	$\mathcal{O}_{LL\bar{d}uD}^{(1)}$	$\epsilon_{ij} (\bar{d}\gamma_\mu u) (L_i^T C (D^\mu L)_j)$
Class 2	$\psi^2 H^2 D^2$	Class 6	$\psi^4 H$
$\mathcal{O}_{LHD}^{(1)}$	$\epsilon_{ij} \epsilon_{mn} (L_i^T C (D_\mu L)_j) H_m (D^\mu H)_n$	$\mathcal{O}_{LL\bar{e}H}$	$\epsilon_{ij} \epsilon_{mn} (\bar{e} L_i) (L_j^T C L_m) H_n$
$\mathcal{O}_{LHD}^{(2)}$	$\epsilon_{im} \epsilon_{jn} (L_i^T C (D_\mu L)_j) H_m (D^\mu H)_n$	$\mathcal{O}_{LLQ\bar{d}H}^{(1)}$	$\epsilon_{ij} \epsilon_{mn} (\bar{d} L_i) (Q_j^T C L_m) H_n$
Class 3	$\psi^2 H^3 D$	$\mathcal{O}_{LLQ\bar{d}H}^{(2)}$	$\epsilon_{im} \epsilon_{jn} (\bar{d} L_i) (Q_j^T C L_m) H_n$
\mathcal{O}_{LHDe}	$\epsilon_{ij} \epsilon_{mn} (L_i^T C \gamma_\mu e) H_j H_m (D^\mu H)_n$	$\mathcal{O}_{LL\bar{Q}uH}$	$\epsilon_{ij} (\bar{Q}_m u) (L_m^T C L_i) H_j$
Class 4	$\psi^2 H^2 X$	$\mathcal{O}_{Leu\bar{d}H}$	$\epsilon_{ij} (L_i^T C \gamma_\mu e) (\bar{d}\gamma^\mu u) H_j$
\mathcal{O}_{LHB}	$\epsilon_{ij} \epsilon_{mn} g' (L_i^T C \sigma^{\mu\nu} L_m) H_j H_n B_{\mu\nu}$		
\mathcal{O}_{LHW}	$\epsilon_{ij} (\epsilon\tau^I)_{mn} g (L_i^T C \sigma^{\mu\nu} L_m) H_j H_n W_{\mu\nu}^I$		

<https://arxiv.org/pdf/1708.09390.pdf>

Higher order operators

$$L = L_{SM} + \frac{1}{E_{new}} L_1 + \frac{1}{E_{new}^2} L_2 + \dots$$

Your favorite LNV theory



Class 1	$\psi^2 H^4$	Class 5	$\psi^4 D$
\mathcal{O}_{LH}	$\epsilon_{ij} \epsilon_{mn} (L_i^T C L_m) H_j H_n (H^\dagger H)$	$\mathcal{O}_{LL\bar{d}uD}^{(1)}$	$\epsilon_{ij} (\bar{d}\gamma_\mu u) (L_i^T C (D^\mu L)_j)$
Class 2	$\psi^2 H^2 D^2$	Class 6	$\psi^4 H$
$\mathcal{O}_{LHD}^{(1)}$	$\epsilon_{ij} \epsilon_{mn} (L_i^T C (D_\mu L)_j) H_m (D^\mu H)_n$	$\mathcal{O}_{LL\bar{e}H}$	$\epsilon_{ij} \epsilon_{mn} (\bar{e} L_i) (L_j^T C L_m) H_n$
$\mathcal{O}_{LHD}^{(2)}$	$\epsilon_{im} \epsilon_{jn} (L_i^T C (D_\mu L)_j) H_m (D^\mu H)_n$	$\mathcal{O}_{LLQ\bar{d}H}^{(1)}$	$\epsilon_{ij} \epsilon_{mn} (\bar{d} L_i) (Q_j^T C L_m) H_n$
Class 3	$\psi^2 H^3 D$	$\mathcal{O}_{LLQ\bar{d}H}^{(2)}$	$\epsilon_{im} \epsilon_{jn} (\bar{d} L_i) (Q_j^T C L_m) H_n$
\mathcal{O}_{LHDe}	$\epsilon_{ij} \epsilon_{mn} (L_i^T C \gamma_\mu e) H_j H_m (D^\mu H)_n$	$\mathcal{O}_{LLQ\bar{u}H}$	$\epsilon_{ij} (\bar{Q}_m u) (L_m^T C L_i) H_j$
Class 4	$\psi^2 H^2 X$	$\mathcal{O}_{Leu\bar{d}H}$	$\epsilon_{ij} (L_i^T C \gamma_\mu e) (\bar{d}\gamma^\mu u) H_j$
\mathcal{O}_{LHB}	$\epsilon_{ij} \epsilon_{mn} g' (L_i^T C \sigma^{\mu\nu} L_m) H_j H_n B_{\mu\nu}$		
\mathcal{O}_{LHW}	$\epsilon_{ij} (\epsilon\tau^I)_{mn} g (L_i^T C \sigma^{\mu\nu} L_m) H_j H_n W_{\mu\nu}^I$		

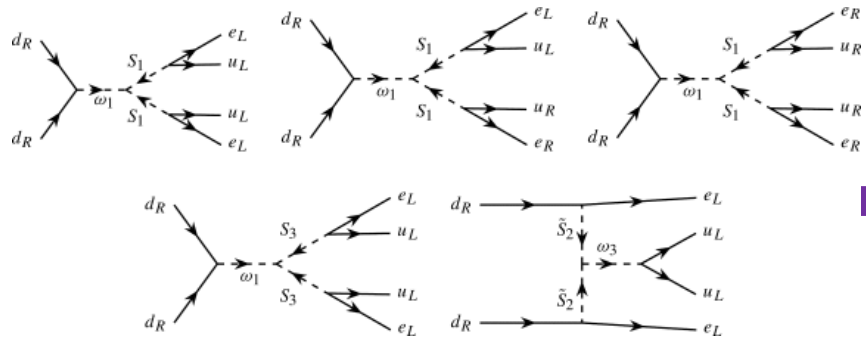
Leptoquarks, SUSY, Kaluza Klein, LR symmetric couplings, etc etc

<https://arxiv.org/pdf/1708.09390.pdf>

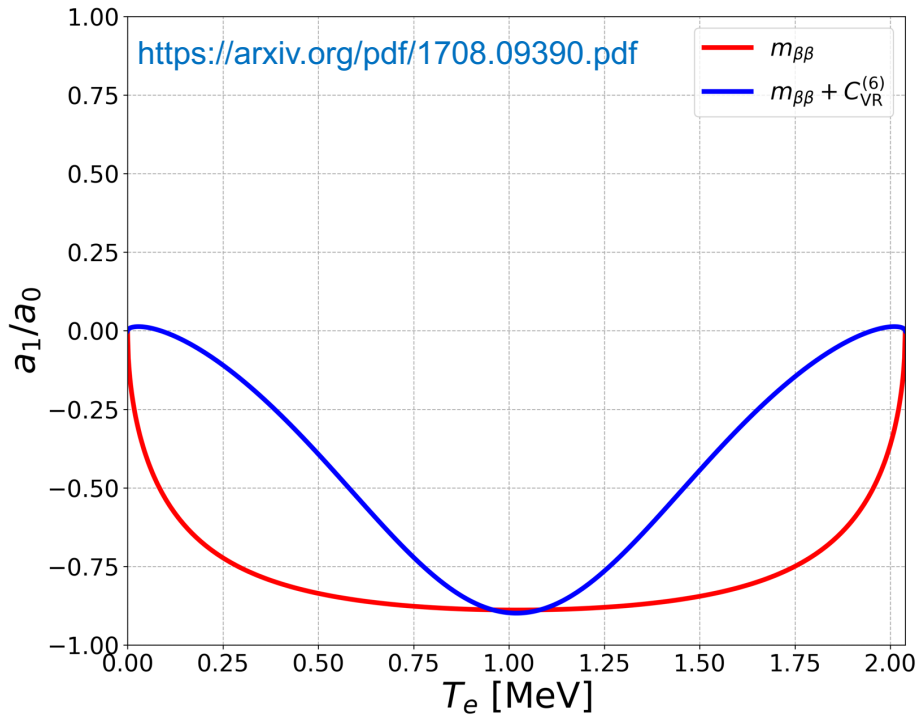
Higher order operators

$$L = L_{SM} + \frac{1}{E_{new}} L_1 + \frac{1}{E_{new}^2} L_2 + \dots$$

Your favorite LNV theory

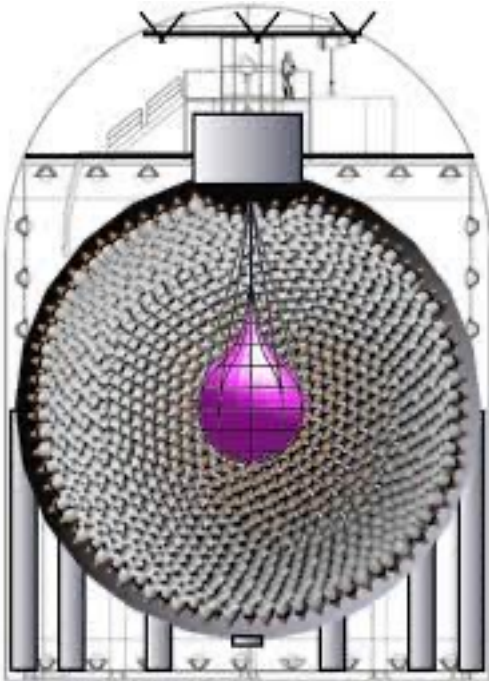


Leptoquarks, SUSY, Kaluza Klein, LR symmetric couplings, etc etc



If your detector can resolve the electron tracks, you can tell the difference...!

But we need to see it first...



**Kamland-Zen
backgrounds →**

TABLE I: Summary of the estimated and best-fit background contributions for the frequentist and Bayesian analyses in the energy region $2.35 < E < 2.70$ MeV within the 1.57-m-radius spherical volume. In total, 24 events were observed.

Background	Estimated	Best-fit	
		Frequentist	Bayesian
$^{136}\text{Xe } 2\nu\beta\beta$	-	11.98	11.95
Residual radioactivity in Xe-LS			
^{238}U series	0.14 ± 0.04	0.14	0.09
^{232}Th series	-	0.84	0.87
External (Radioactivity in IB)			
^{238}U series	-	3.05	3.46
^{232}Th series	-	0.01	0.01
Neutrino interactions			
^8B solar νe^- ES	1.65 ± 0.04	1.65	1.65
Spallation products			
Long-lived	$7.75 \pm 0.57^\dagger$	12.52	11.80
^{10}C	0.00 ± 0.05	0.00	0.00
^6He	0.20 ± 0.13	0.22	0.21
^{137}Xe	0.33 ± 0.28	0.34	0.34

[†] Estimation based on the spallation MC study. This event rate constraint is not applied to the spectrum fit.

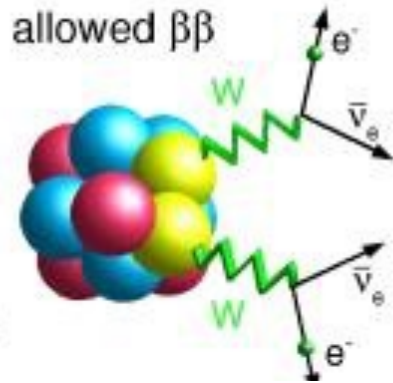
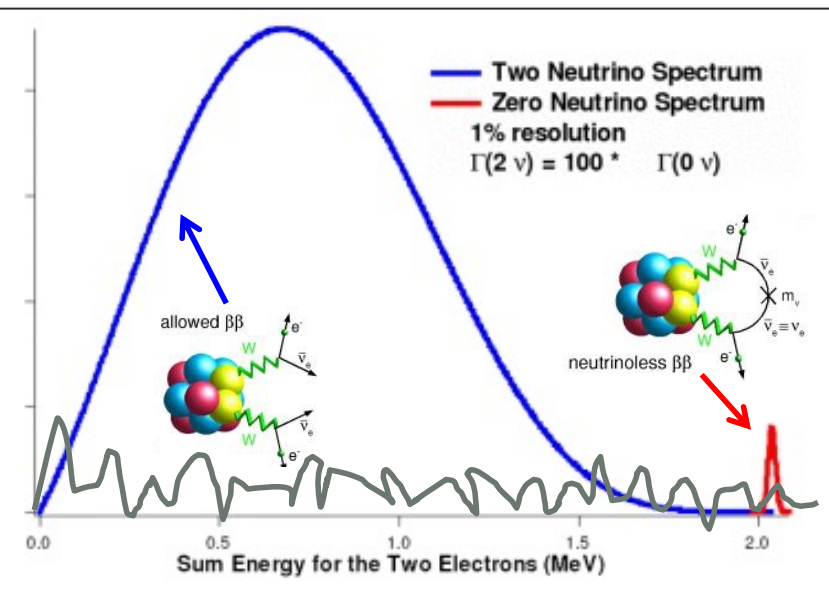
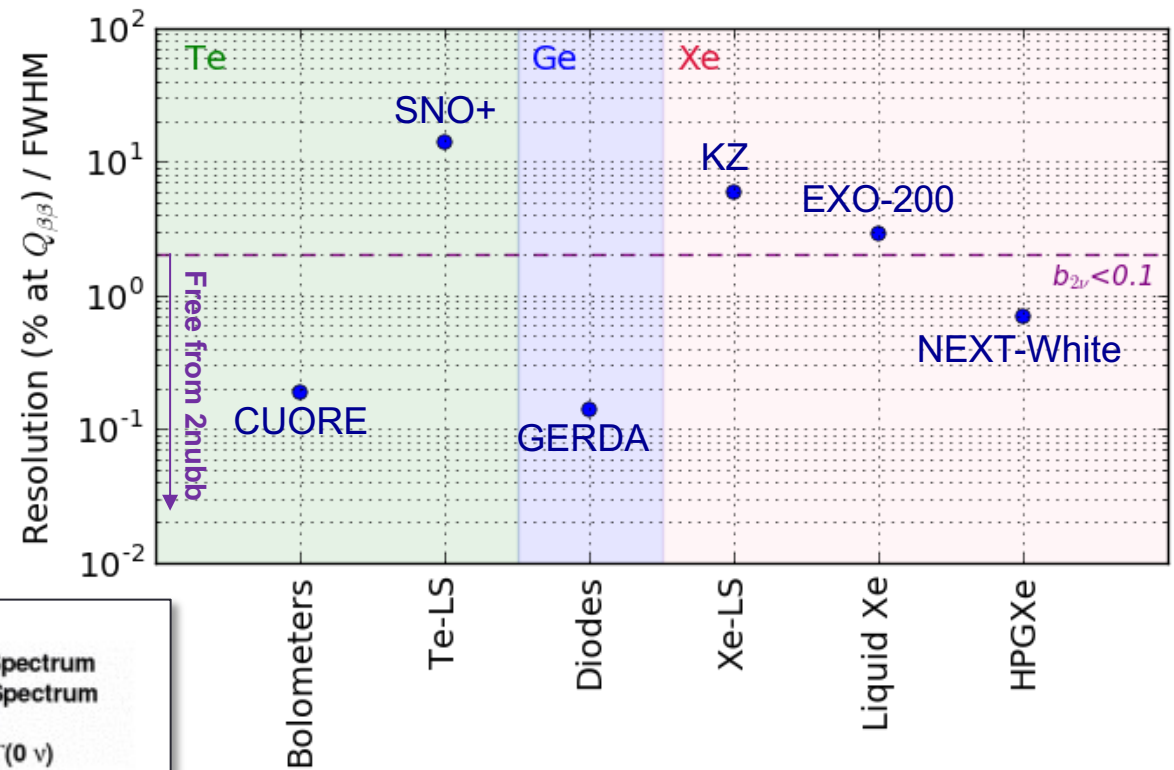


TABLE I: Summary of the estimated and best-fit background contributions for the frequentist and Bayesian analyses in the energy region $2.35 < E < 2.70$ MeV within the 1.57-m-radius spherical volume. In total, 24 events were observed.

Background	Estimated	Best-fit	
		Frequentist	Bayesian
$^{136}\text{Xe } 2\nu\beta\beta$	-	11.98	11.95
Residual radioactivity in Xe-LS			
^{238}U series	0.14 ± 0.04	0.14	0.09
^{232}Th series	-	0.84	0.87
External (Radioactivity in IB)			
^{238}U series	-	3.05	3.46
^{232}Th series	-	0.01	0.01
Neutrino interactions			
^8B solar νe^- ES	1.65 ± 0.04	1.65	1.65
Spallation products			
Long-lived	$7.75 \pm 0.57^\dagger$	12.52	11.80
^{10}C	0.00 ± 0.05	0.00	0.00
^6He	0.20 ± 0.13	0.22	0.21
^{137}Xe	0.33 ± 0.28	0.34	0.34

† Estimation based on the spallation MC study. This event rate constraint is not applied to the spectrum fit.

Measuring Energy



2nubb can be efficiently rejected with sufficient energy resolution alone.

Where are the backgrounds from?

Radiogenics

Solar neutrinos(!)

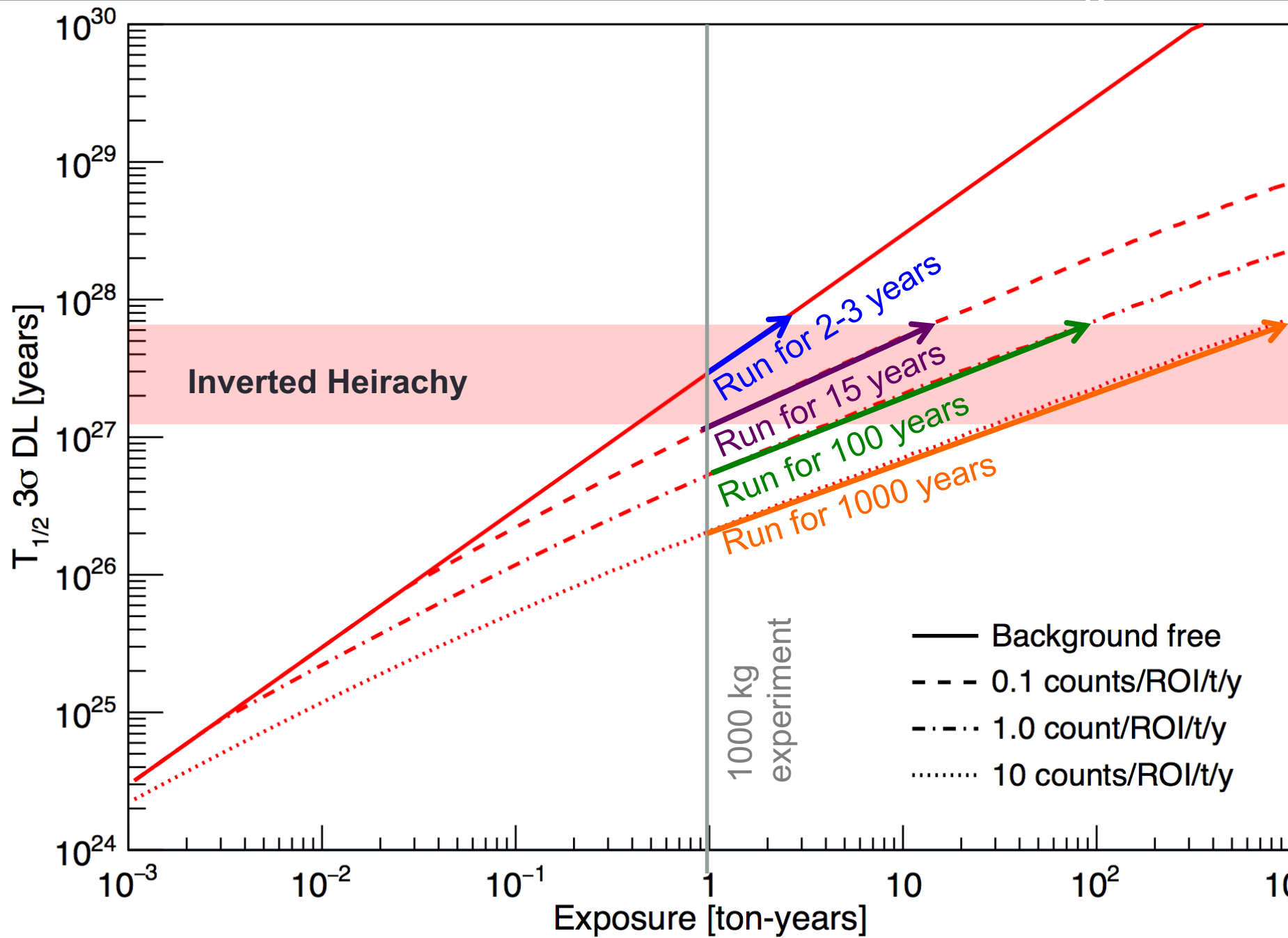
Cosmogenics

These will limit all future experiments.

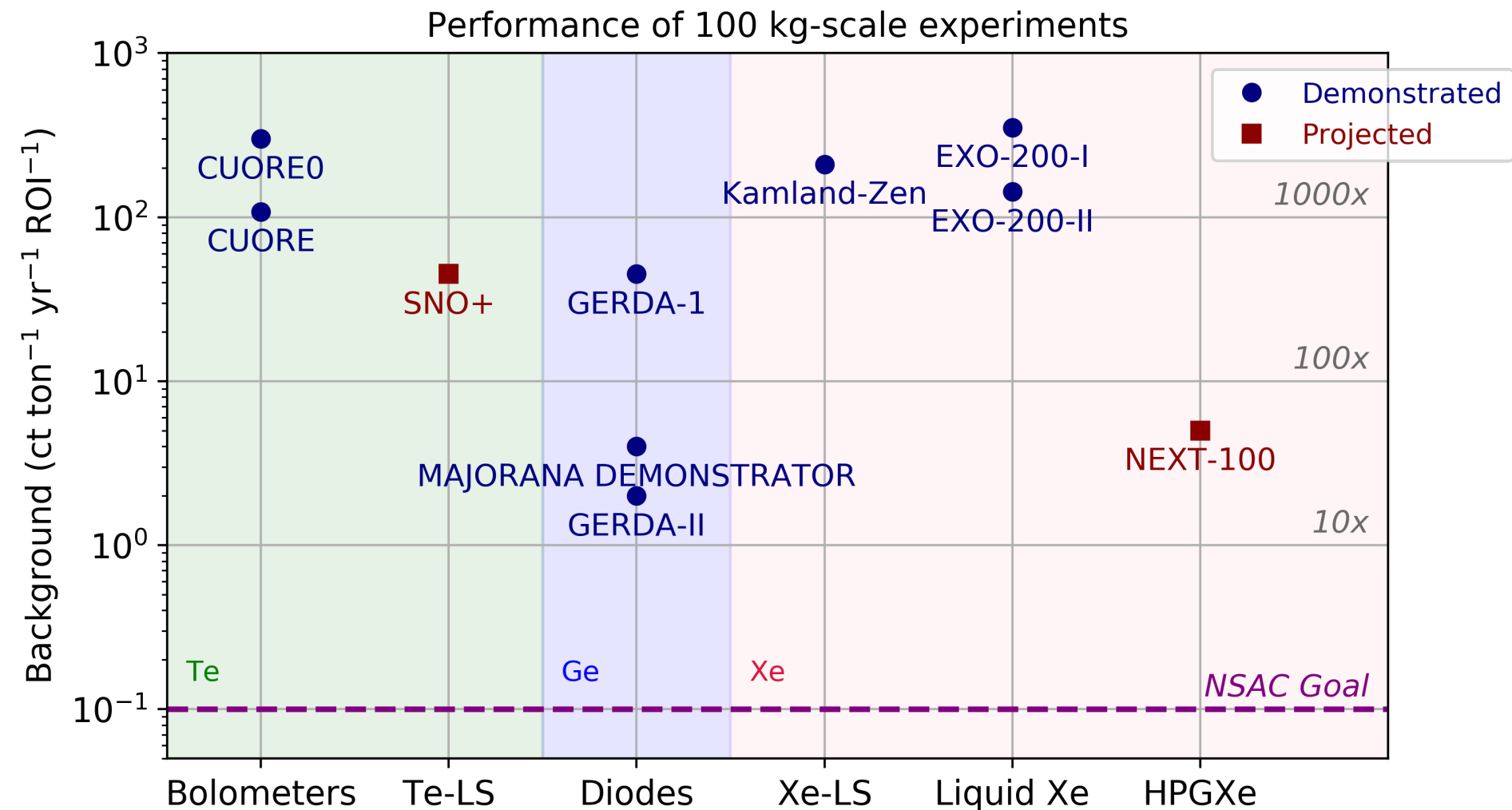
TABLE I: Summary of the estimated and best-fit background contributions for the frequentist and Bayesian analyses in the energy region $2.35 < E < 2.70$ MeV within the 1.57-m-radius spherical volume. In total, 24 events were observed.

Background	Estimated	Best-fit	
		Frequentist	Bayesian
$^{136}\text{Xe } 2\nu\beta\beta$	-	11.98	11.95
Residual radioactivity in Xe-LS			
^{238}U series	0.14 ± 0.04	0.14	0.09
^{232}Th series	-	0.84	0.87
External (Radioactivity in IB)			
^{238}U series	-	3.05	3.46
^{232}Th series	-	0.01	0.01
Neutrino interactions			
^8B solar νe^- ES	1.65 ± 0.04	1.65	1.65
Spallation products			
Long-lived	$7.75 \pm 0.57^\dagger$	12.52	11.80
^{10}C	0.00 ± 0.05	0.00	0.00
^6He	0.20 ± 0.13	0.22	0.21
^{137}Xe	0.33 ± 0.28	0.34	0.34

[†] Estimation based on the spallation MC study. This event rate constraint is not applied to the spectrum fit.

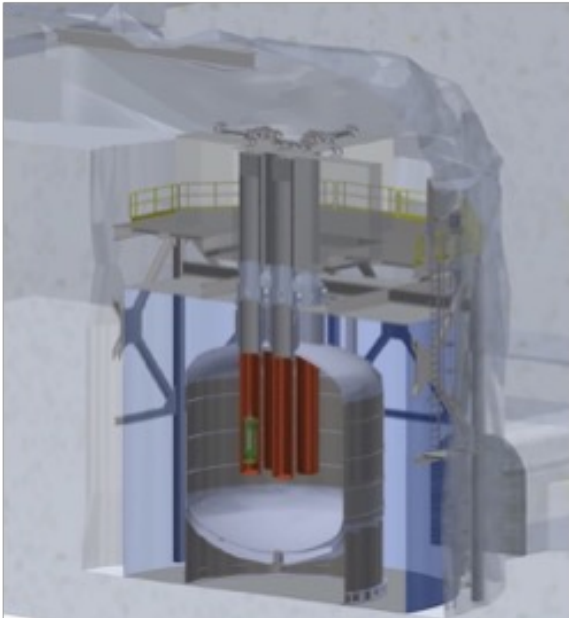


“100kg-class” experiments:



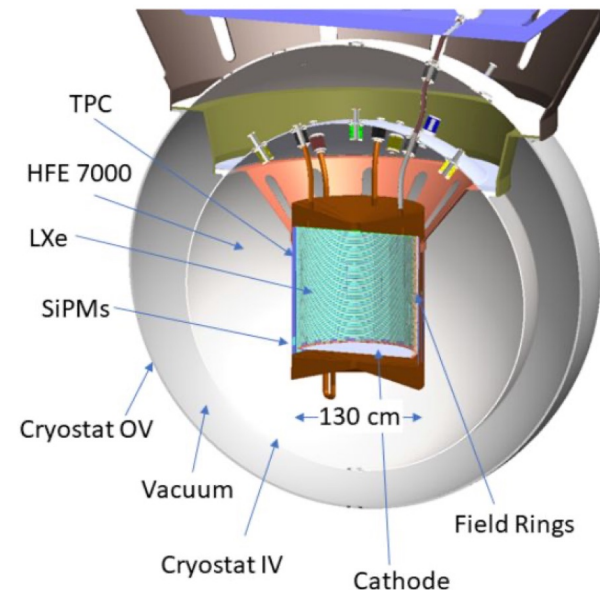
LEGEND

- Germanium diodes immersed in low background liquid argon.
- Lowest demonstrated background from any technology with GERDA and Majorana.
- 200 kg phase in operation, 1 Ton phase proposed.
- Superb energy resolution, $\sim 0.2\%$ FWHM
- Deployable in stages as isotope becomes available.



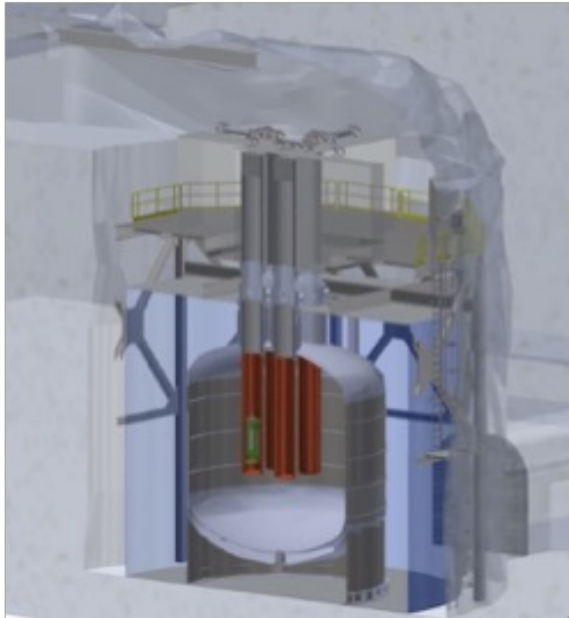
nEXO

- 5 Ton liquid xenon time projection chamber using enriched ^{136}Xe .
- Several signal-sensitive variables combined in a multivariate manner for final signal metric.
- Demonstrated energy res around 2-3% FWHM, aiming to improve through ongoing R&D.
- Outer tons of ^{136}Xe self-shield the middle ~ 2 tons to give a clean, very sensitive inner region.



LEGEND

- Germanium diodes immersed in low background liquid argon.
- Lowest demonstrated background from any technology with GERDA and Majorana.
- 200 kg phase in operation, 1 Ton phase proposed
- Superb energy resolution, $\sim 0.2\%$ FWHM
- Deployable in stages as isotope becomes available

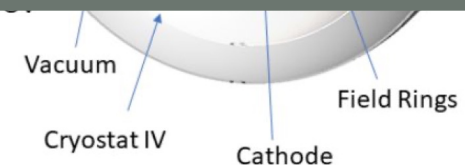
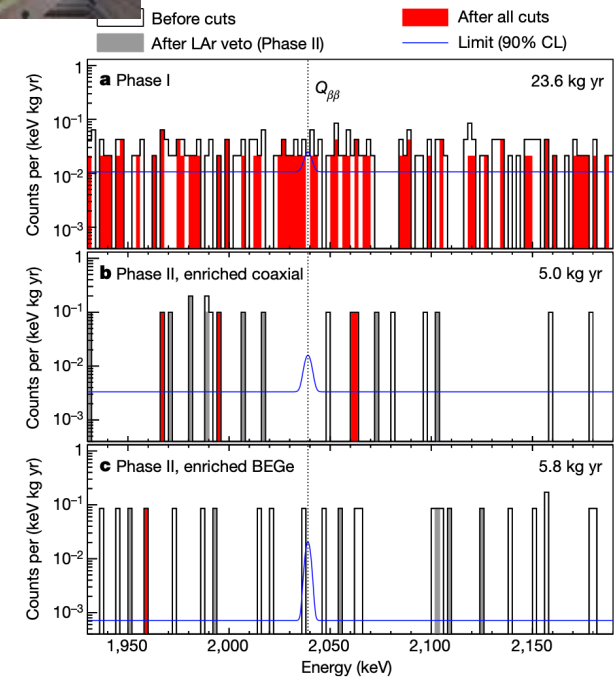


nEXO



Proven essentially background free using energy reconstruction alone by MJD / GERDA

Scaling is mostly a matter of repetition.



LEGEND

- Germanium diodes immersed in low background

Background suppression from self shielding emerges at large scale, and power comes from a multivariate analysis of many somewhat-signal-sensitive variables.

nEXO

- 5 Ton liquid xenon time projection chamber using enriched ^{136}Xe .

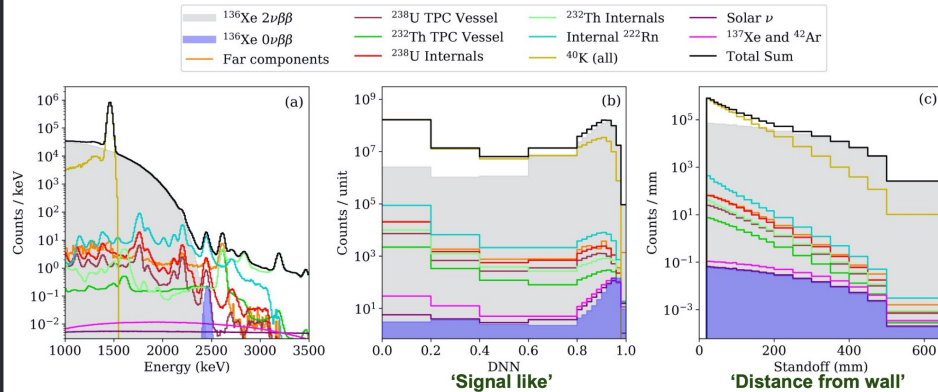
Several signal-sensitive variables combined in a multivariate manner for final signal metric.

Demonstrated energy res around 2-3% FWHM, aiming to improve through ongoing R&D.

Outer tons of ^{136}Xe self-shield the middle ~2 tons to give a clean, very sensitive inner region.

Expected full event distributions

Signal: $0.7 \times 10^{28} \text{y}$

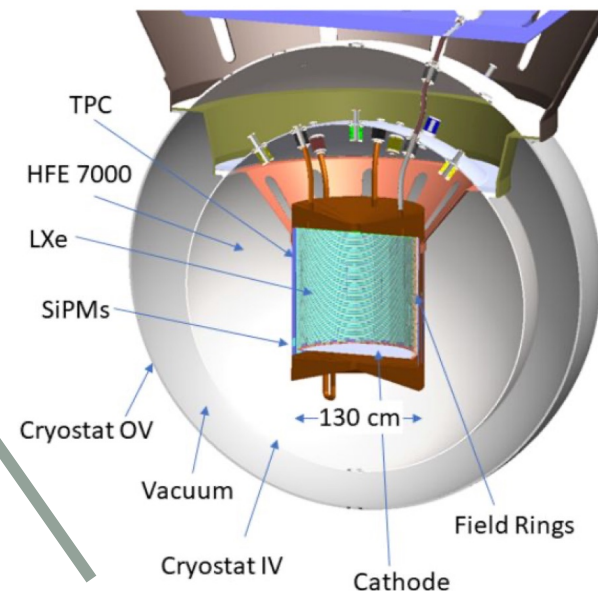


P.A. (Sander) Breur

NDM 2022

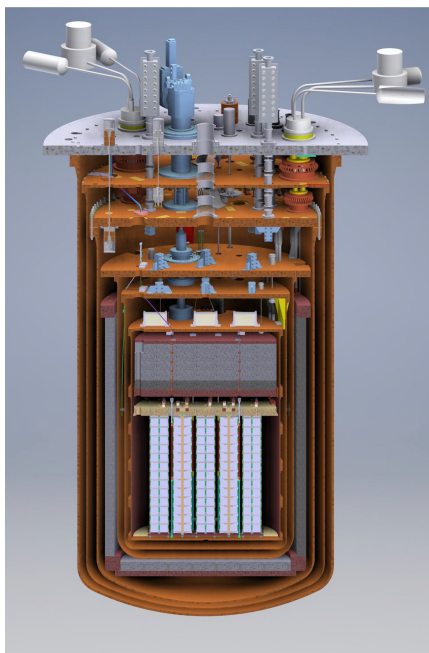
May 18, 2022

11



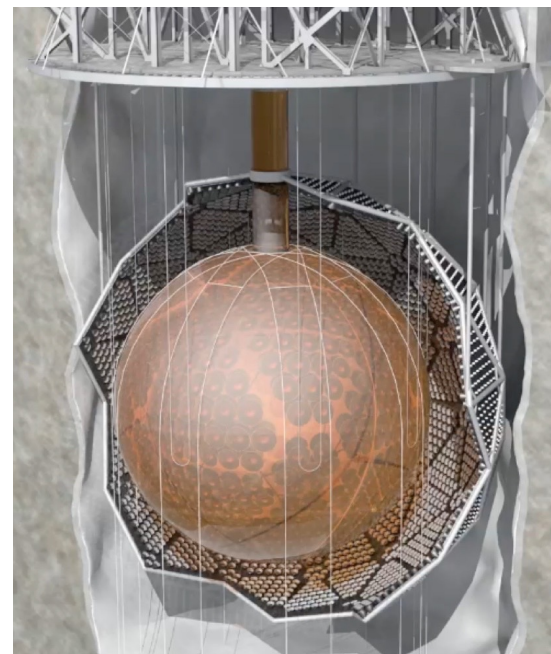
CUPID

- ^{100}Mo enriched scintillating Li_2MoO_4 bolometers read out with transition edge sensors. 280kg of isotope.
- Excellent demonstrated energy resolution of 0.2% FWHM.
- Scintillation allows separation of surface backgrounds from bulk, which ultimately limited CUORE sensitivity.
- New crystals have been operated in CUPID-Mo demonstrator, showing strong performance.



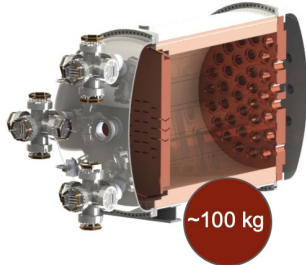
SNO+

- 3.9 tons of tellurium loaded in 780 tons of liquid scintillator in an upgraded SNO detector
- Water-filled phase, scintillator-phase in progress.
- Isotope to be added in stages until maximal loading without detriment to scintillator optical properties.
- Modest energy resolution characteristic of scintillator detectors.
- Large isotope mass enabled by high natural abundance of ^{134}Te - no enrichment needed.



NEXT-100

2022-2026

Neutrinoless double beta decay
search in ^{136}Xe **NEXT-HD**

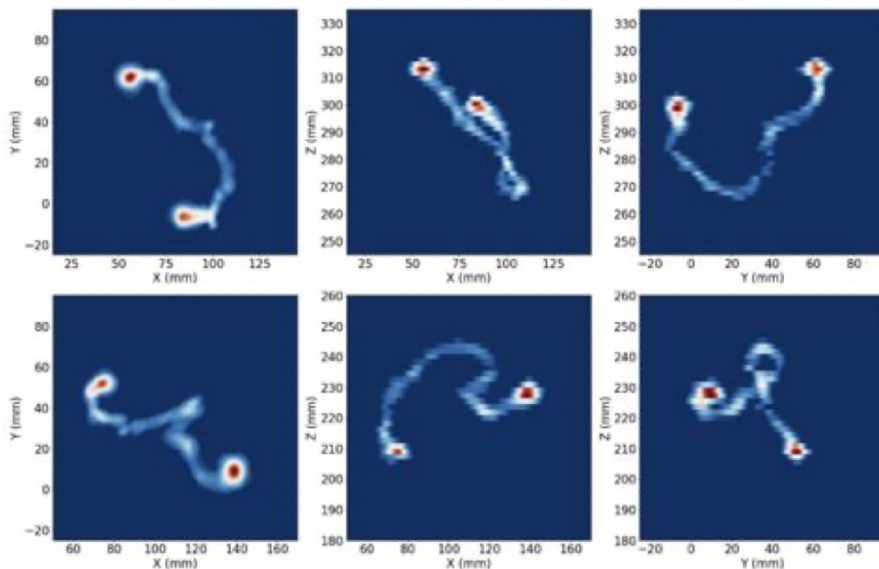
2027?

Neutrinoless double beta decay
search through inverted neutrino
mass ordering**NEXT-BOLD**Barium tagging for background-free
experiment

15

NEXT

- Xenon gas (10-15 bar) time projection chambers with enriched ^{136}Xe .
- Recombination-less ionization readout in gas gives leading energy resolution in xenon, $\sim 0.9\%$ FWHM.
- Tracks can be imaged topologically to separate signals (2 blobs) from backgrounds (1 blob)
- Full detector volume is active isotope (no self shielding)
- NEXT-100 will run in 2023, followed by ton-scale phases.

 ^{208}Tl double escape Events: RL deconvolved

Please indulge me while I talk about my own stuff for 2 slides... 🧐

Barium Tagging

In the decay:



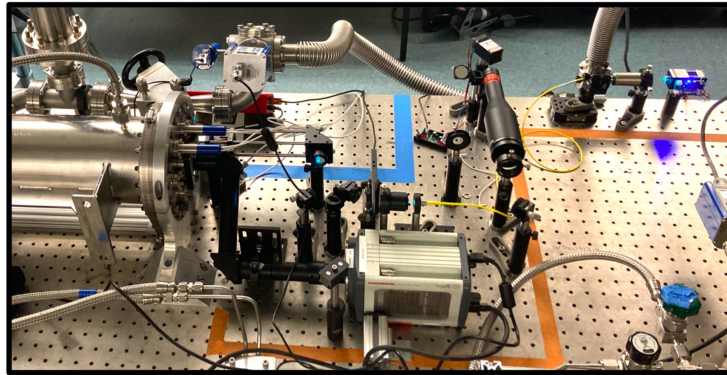
Barium does not accompany any of the major cosmogenic or radiogenic backgrounds.

Single ion imaging may enable background-free multi-ton scale techniques.

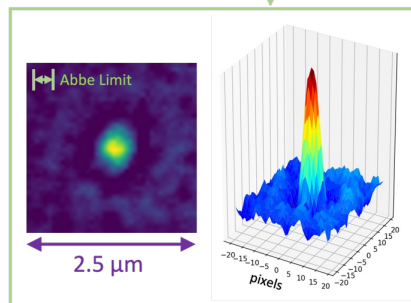
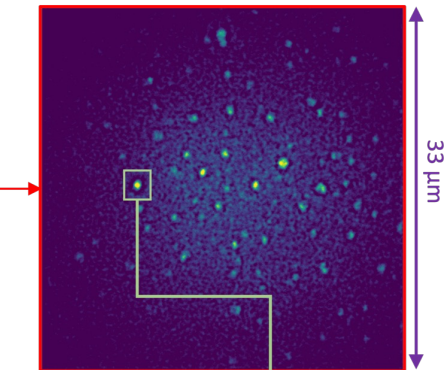
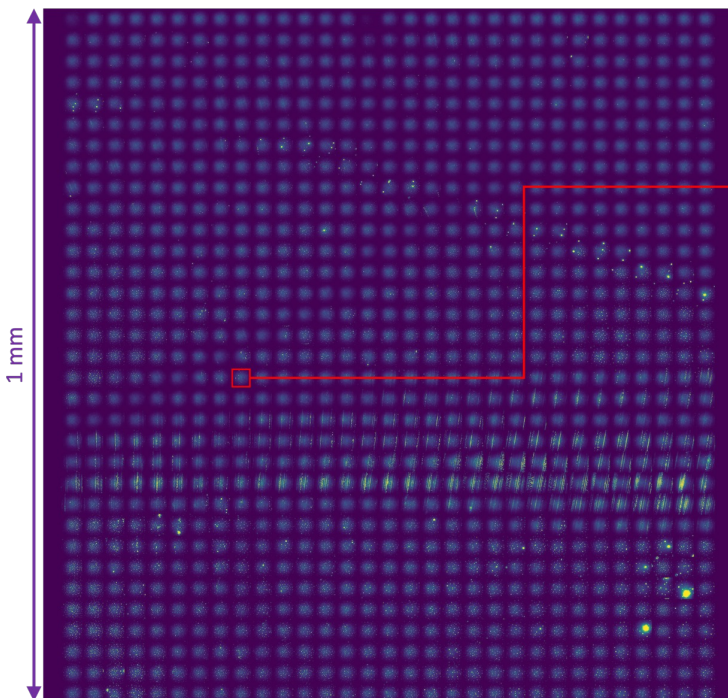
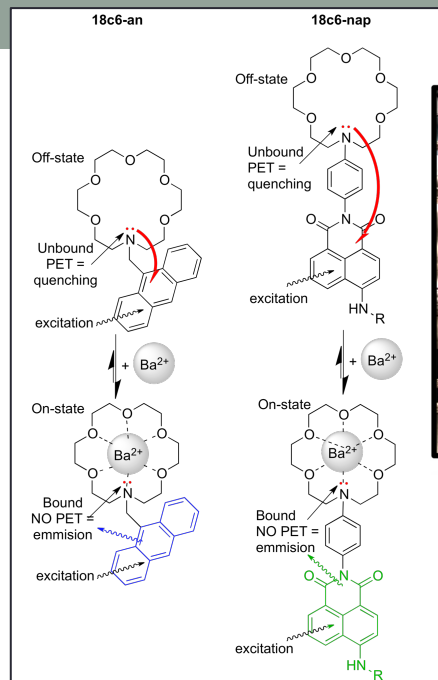
Efficient identification of one ion in a ton of material is an extreme technological challenge.

Progress on imaging is being made rapidly in both liquid and gaseous xenon.

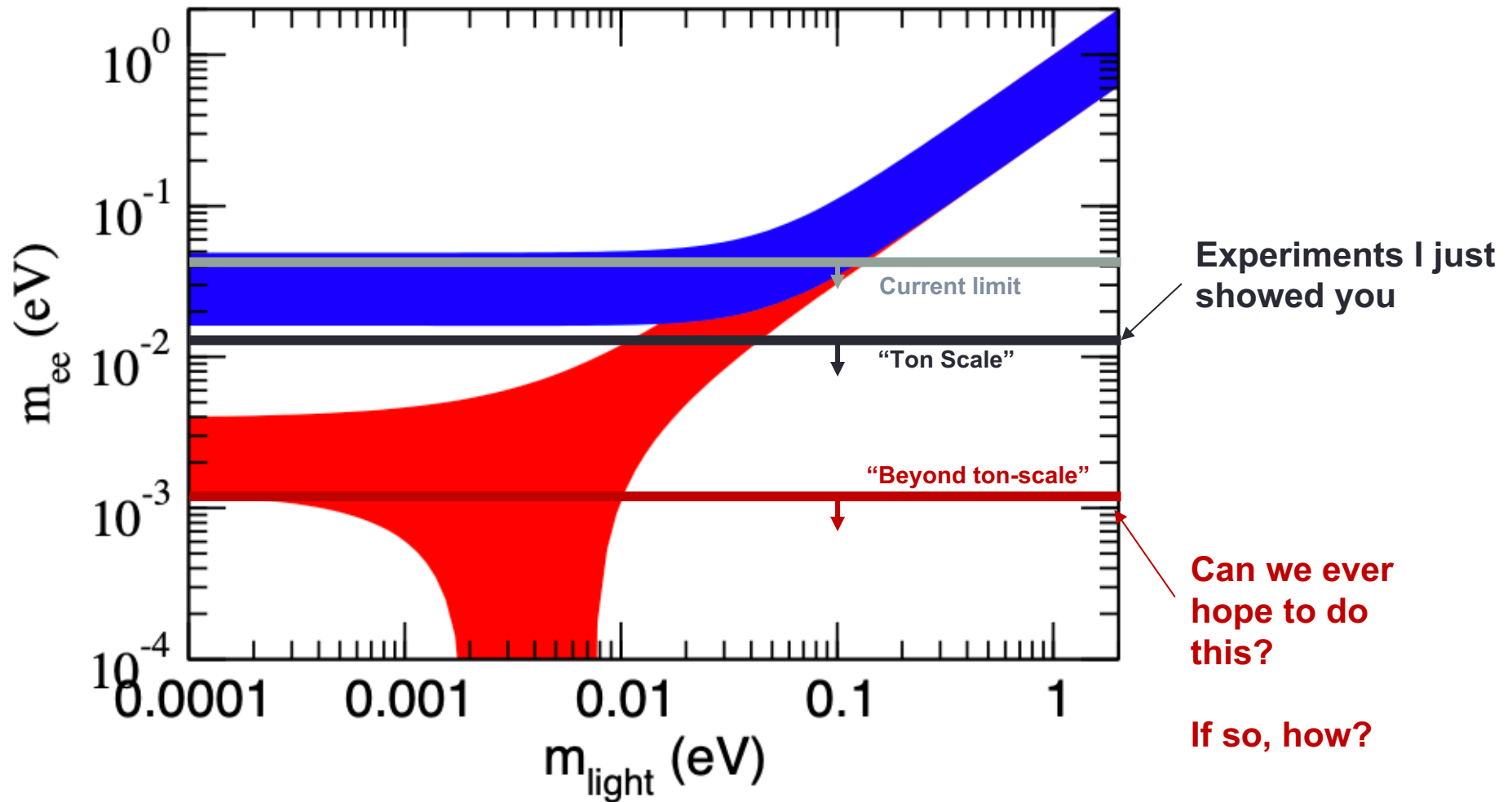
Collection from detector volume is the next major R&D challenge.



Eg: Single Ba²⁺ ions imaged over mm² surfaces in high pressure xenon gas for NEXT.



Toward Normal Ordering



The path to normal ordering is far from clear. Ultra-low background, very large scale detectors, with hundreds of tons of isotope, are needed.

We would need to do it if:

- Neutrino masses are normal ordered
- The same mechanism generates neutrino mass and drives θ_{MNSP}
- The absolute mass of the lightest neutrino is <0.05 eV

These may very well all be true statements...

What then?

Isotopes

- Already a challenging problem for ton-scale experiments, but would need hundreds of tons of isotope to go to normal ordering.
- Difficulties are associated with both **acquisition** of raw material and **enrichment**.
- Some notable factoids:
 - **Tellurium:**
Comes naturally enriched to 34%. So natural tellurium is the most viable for an unenriched experiment.
 - **Molybdenum:**
New capacity for enrichment of Molybdenum in 100Mo for nuclear medicine is needed. Onubb may be parasitic?
 - **Germanium:**
Semiconductor industry enriches germanium already; 76Ge can in principle be extracted as byproduct?
 - **Xenon:**
Atmospheric carbon capture technology based on metal organic frameworks has plausible extendibility to capture atmospheric Xe. Free from steel industry capacity limit?

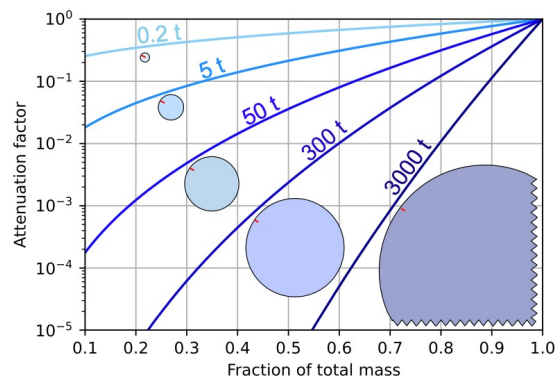
R&D and new facilities would be needed to produce isotope at the scale needed for a normal-ordering experiment.



Giant TPCs

Kiloton-scale xenon detectors for neutrinoless double beta decay and other new physics searches

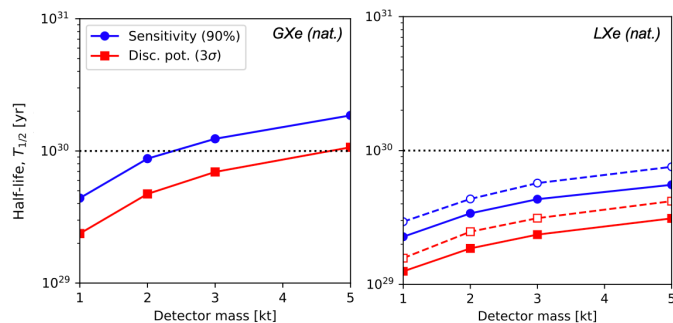
A. Avasthi,¹ T.W. Bowyer,² C. Bray,³ T. Brunner,^{4,5} N. Catarineu,⁶ E. Church,² R. Guenette,⁷ S.J. Haselschwardt,⁸ J.C. Hayes,² M. Heffner,^{6,*} S.A. Hertel,⁹ P.H. Humble,² A. Jamil,¹⁰ S.H. Kim,⁶ R.F. Lang,¹¹ K.G. Leach,³ B.G. Lenardo,¹² W.H. Lippincott,¹³ A. Marino,³ D.N. McKinsey,^{14,8} E.H. Miller,^{15,16} D.C. Moore,^{10,†} B. Mong,¹⁵ B. Monreal,¹ M.E. Monzani,^{15,16} I. Olcina,^{8,14} J.L. Orrell,² S. Pang,⁶ A. Pocar,⁹ P.C. Rowson,¹⁵ R. Saldanha,² S. Sangiorgio,⁶ C. Stanford,⁷ and A. Visser⁶



← Radiogenics become totally irrelevant due to self shielding

Cosmogenics and solar neutrinos become a serious concern!

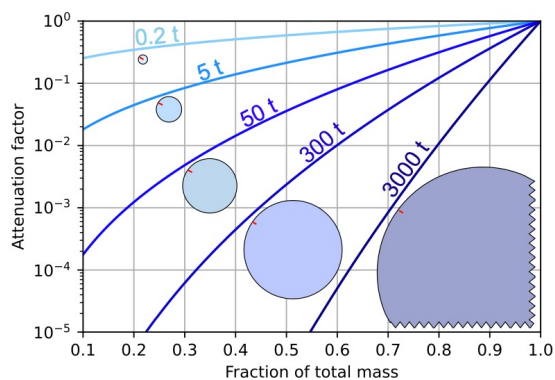
If energy resolution achievable at scale, with kiloton masses, normal ordering parameter space is accessible.



Giant TPCs ...in salt caverns?

Kiloton-scale xenon detectors for neutrinoless double beta decay and other new physics searches

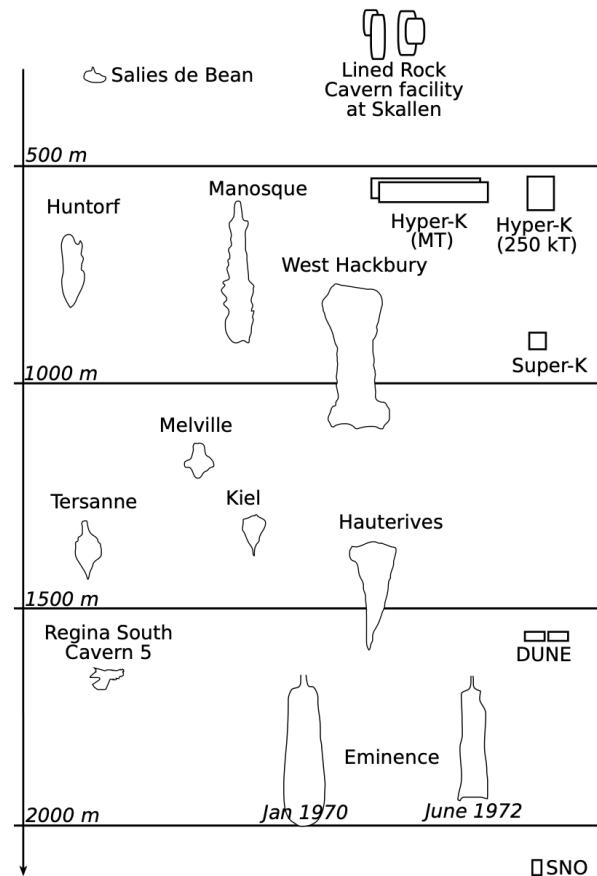
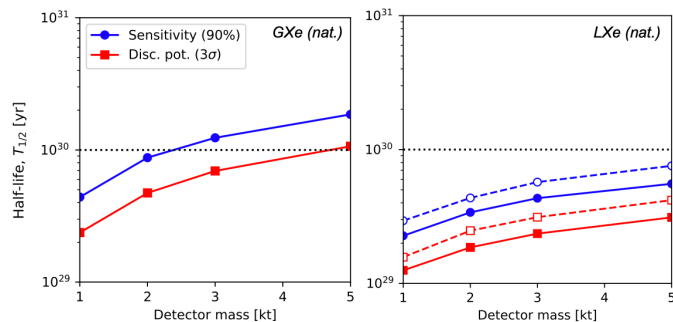
A. Avasthi,¹ T.W. Bowyer,² C. Bray,³ T. Brunner,^{4,5} N. Catarineu,⁶ E. Church,² R. Guenette,⁷ S.J. Haselschwardt,⁸ J.C. Hayes,² M. Heffner,^{6,*} S.A. Hertel,⁹ P.H. Humble,² A. Jamil,¹⁰ S.H. Kim,⁶ R.F. Lang,¹¹ K.G. Leach,³ B.G. Lenardo,¹² W.H. Lippincott,¹³ A. Marino,³ D.N. McKinsey,^{14,8} E.H. Miller,^{15,16} D.C. Moore,^{10,†} B. Mong,¹⁵ B. Monreal,¹ M.E. Monzani,^{15,16} I. Olcina,^{8,14} J.L. Orrell,² S. Pang,⁶ A. Pocar,⁹ P.C. Rowson,¹⁵ R. Saldanha,² S. Sangiorgio,⁶ C. Stanford,⁷ and A. Visser⁶



← Radiogenics become totally irrelevant due to self shielding

Cosmogenics and solar neutrinos become a serious concern!

If energy resolution achievable at scale, with kiloton masses, normal ordering parameter space is accessible.



High-pressure TPCs in pressurized caverns: opportunities in dark matter and neutrino physics

Author:

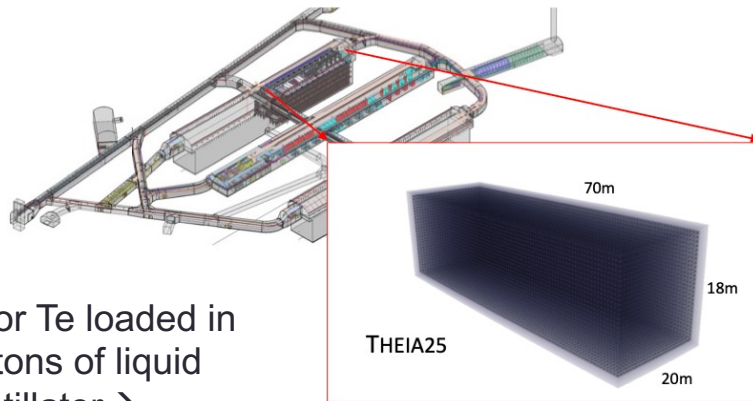
Benjamin Monreal (Case Western Reserve U.) [benjamin.monreal@case.edu]

Isotope in kTons of liquid scintillator

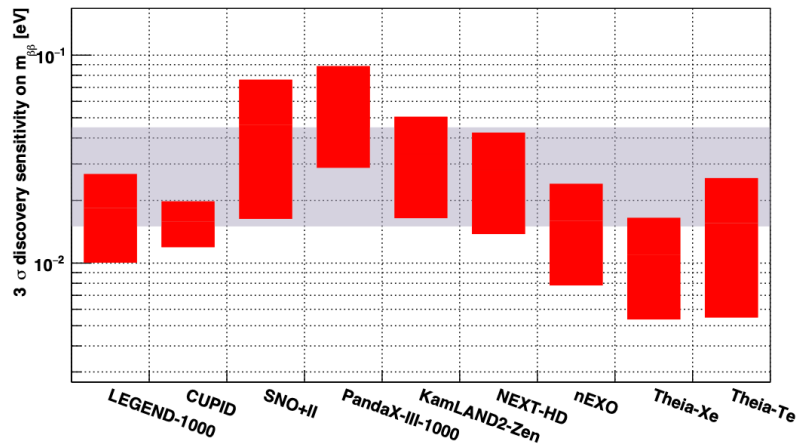
THEIA: Summary of physics program

Snowmass White Paper Submission

M. Askins,^{1,2} Z. Bagdasarian,^{1,2} N. Barros,^{3,4,5} E.W. Beier,³ A. Bernstein,⁶ M. Böhles,⁷ E. Blucher,⁸



Xe or Te loaded in
kilotons of liquid
scintillator →

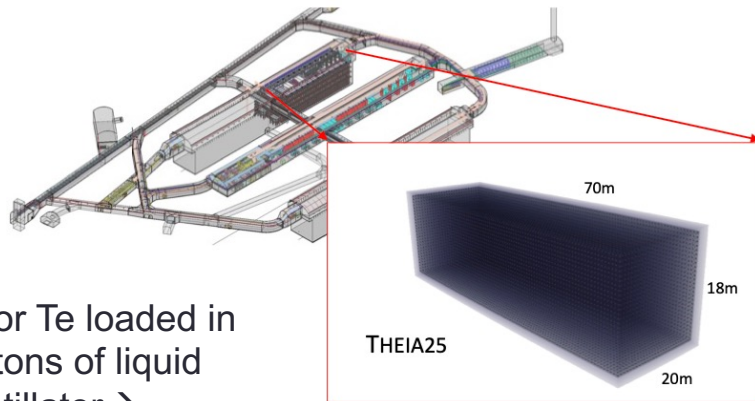


Isotope in kTons of liquid scintillator

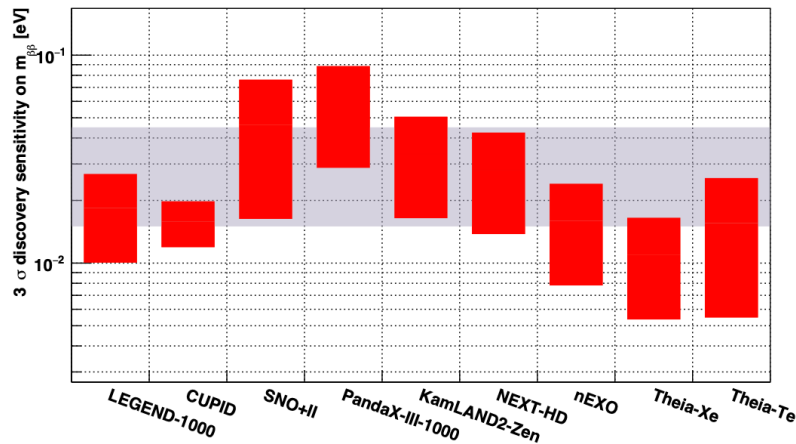
THEIA: Summary of physics program

Snowmass White Paper Submission

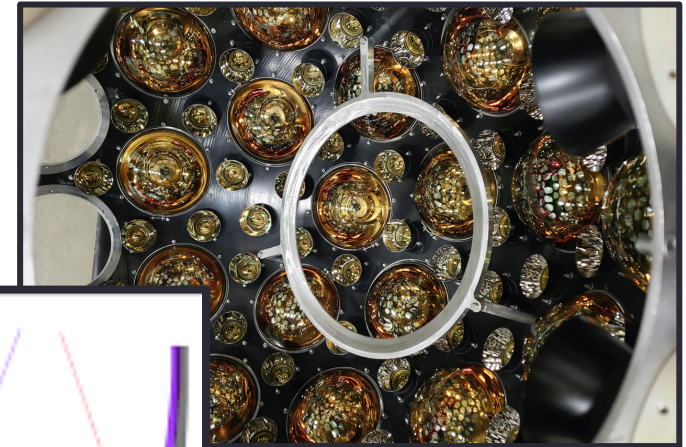
M. Askins,^{1,2} Z. Bagdasarian,^{1,2} N. Barros,^{3,4,5} E.W. Beier,³ A. Bernstein,⁶ M. Böhles,⁷ E. Blucher,⁸



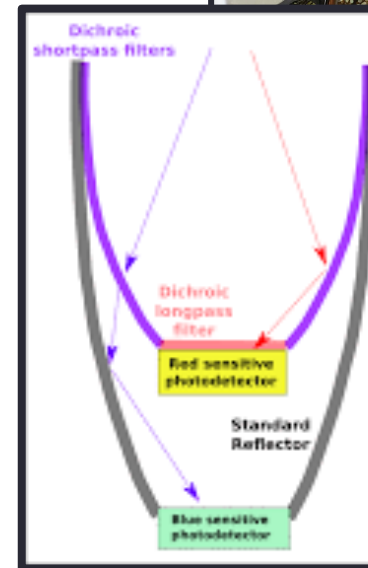
Xe or Te loaded in kilotons of liquid scintillator →



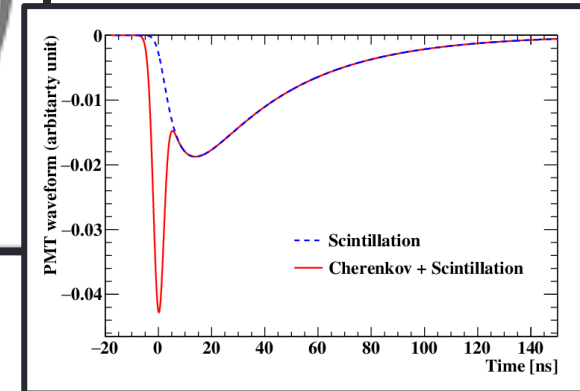
Much ongoing R&D...



Quantum dot doped liquid scintillator @NuDot



Dichroicon



Slow scintillator for Cherenkov separation

Conclusions

- NDBD is the only sensitive known way to probe the Majorana nature of the neutrino.
- Experiments at the 100kg scale have demonstrated background indices in the range 2-200 ct/ton/ky/yr
- Ton-scale experiments plan to reduce backgrounds by 1.5-3 orders of magnitude relative to 100kg phases and probe the inverted mass ordering range of parameter space.
- Beyond-ton-scale will require huge, ultra-low background detectors that we don't yet know how to build, but need to figure out!

Thank you for your attention

