Atom Interferometry for dark matter and gravitational waves

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in collaboration with Leonardo Badurina, Ankit Beniwal, Diego Blas, John Carlton, Val Gibson, Jeremiah Mitchell, and others in AION



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Setting the scene

Light pulse atom interferometry (physical-space)

- Launch ultra-cold cloud of atoms into an atomic fountain
- Sequence of optical pulses manipulate the atoms
- Quantum superposition over macroscopic distances (>50cm achieved)
- Interfere using a final optical pulse when they spatially overlap
- Image the two interferometer output ports







Light pulse atom interferometry (space-time)



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Two-level system separated by optical frequency difference ω_a

Initial pulse: 'beamsplitter' Middle pulse: 'mirror' Final pulse: 'beamsplitter (interfere)'

Atom evolves extra clock phase:

$$\frac{1}{\sqrt{2}}|1\rangle + \frac{1}{\sqrt{2}}|2\rangle e^{-i\omega_a T}$$

Phase sensitive to changes in timings, atomic structure, and local accelerations



New atom interferometers across the world coming online



MAGIS-100, arXiv:2104.02835; MIGA, arXiv:1703.02490; AION, arXiv:1911.11755; VLBAI, arXiv:2003.04875; ZAIGA, arXiv:1903.09288



AION: Atom Interferometer Observatory and Network



7 institutes in the UK



AION, arXiv:1911.11755

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Collaboration ~50 people Cold atom: fundamental physics ratio is ~50:50





Badurina, CM, et al (AION), arXiv:1911.11755 Image from Abe et al (MAGIS-100) arXiv:2104.02835

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Operate in gradiometer configuration: run two atom interferometers simultaneously with the same laser

Pushing state-of-the-art single photon strontium atom interferometry

Partnering with MAGIS-100 in the US

Most sensitive to 'mid-band' (0.1 - 10 Hz)



AION: envisaged as a multi-stage project



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Stage 1: AION-10

~10m tower in the Beecroft building in Oxford (new, low-vibration building)

Now: 5 new Sr labs and design '24-'26: construction '26-'27: commissioning 2028+: science







AION: envisaged as a multi-stage project

AION-10 2020s ~10m instrument in Oxford AION-100 2030s ~100m instrument at Boulby/CERN/...?





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CERN study: arXiv:2304.00614 ; AEDGE, arXiv:1908.00802; Roadmap, arXiv:2201.07789









Near-term aim: probe dark matter

Badurina, Blas, **CM**, PRD, arXiv:2109.10965 Badurina, ..., **CM**, et al, Phil.Trans.Roy.Soc.Lond., arXiv:2108.02468





DM landscape: classifying by mass



Ultra-light dark matter

DM lighter than ~few eV behaves as a classical wave



Angular frequency set by the ULDM mass: $\omega \simeq m_{arphi} \left(1 + \mathcal{O}(v^2)\right)$

e.g., Foster et al arXiv: 1711.10489 Derevianko arXiv:1605.09717





Classifying atom interferometer signals

Difficulty: very high Careful analysis of systematic effects needed, which may be hard to quantify

Focus initially on time-dependent signals

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ULDM-induced signal Static vs Time-dependent

Difficulty: medium Characteristic DM signal allows for greater signal discrimination



An oscillating ULDM field can induce several signals testable with Als:

- 1. Changes in fundamental constants (scalar ULDM)
- 2. Accelerations on test masses (vector ULDM)
- 3. Precession of spins (pseudoscalar ULDM)



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Time-dependent signals





Changes in fundamental constants (Scalar)

$$\mathcal{L} \supset \sqrt{4\pi G_N} \phi igg[d_{m_e} m_e ar{e} e - rac{d_e}{4} F_{\mu
u} F^{\mu
u} igg]
ightarrow m_e(t,\mathbf{x}) = m_e igg[1 + d_{m_e} \sqrt{4\pi G_N} \phi(t,\mathbf{x}) igg]
ightarrow lpha(t,\mathbf{x}) = lpha igg[1 + d_e \sqrt{4\pi G_N} \phi(t,\mathbf{x}) igg]$$

Oscillations in the field lead to oscillations in optical transitions:



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$$\int \Delta \omega_A(t) \sim \left[d_{m_e} + \xi_A d_e \right] \cos(m_\phi t + \xi_A d_e)$$

See e.g., Geraci et al, arXiv:1605.04048 and Arvanitaki et al, arXiv:1606.04541







Phase is accumulated by the excited state relative to the ground state along all paths:

$$\Phi_{t_1}^{t_2}(\mathbf{r}) = \int_{t_1}^{t_2} \Delta \omega_a(t, \mathbf{r}) dt$$

 $\Delta \omega_A(t) \sim \left[d_{m_e} + \xi_A d_e \right] \cos(m_\phi t + \theta)$

 t_1, t_2 = time in excited state

Scalar ULDM signal







Many parameters to tune to reach sensitivity

$$d_{m_e}^{\text{best}} \sim \left(\frac{1}{T}\right)^{5/4} \frac{1}{C n \Delta r} \left(\frac{\Delta t}{N_a}\right)^{1/2} \left(\frac{1}{T_{\text{int}}}\right)^1$$

Handles to optimise (in order of priority):

 $T \sim Is$ (interrogation time) $C \sim 0.1 - I$ (constrast) $n \sim 1000$ (LMT) $\Delta r \sim AI$ separation $\Delta t \sim sampling$ time $N_a \sim atoms$ in cloud $T_{int} \sim 10^7 s$ (integration time)







Sensitivity	\mathbf{L}	T_{int}	$\delta \phi_{ m noise}$	\mathbf{LMT}
Scenario	[m]	[sec]	$[1/\sqrt{\text{Hz}}]$	$[\mathrm{number}\ n]$
AION-10 (initial)	10	1.4	10^{-3}	100
AION-10 (goal)	10	1.4	10^{-4}	1000
AION-100 (initial)	100	1.4	10^{-4}	1000
AION-100 (goal)	100	1.4	10^{-5}	40000
AION-km	2000	5	$0.3 imes10^{-5}$	40000

Badurina, CM, et al, arXiv:1911.11755, 2108.02468

Near- and long-term prospects (Scalar)



Long-term aim: Gravitational wave searches

Badurina, Buchmueller, Ellis, Lewicki, CM, Vaskonen Phil.Trans.Roy.Soc.Lond., arXiv:2108.02468

Gravitational wave detection



Dimopoulos et al, PRD arXiv:0802.4098, PRD arXiv:0806.2125 Graham et al, PRL arXiv:1206.0818, PRD arXiv:1606.01860

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$$\Phi \propto hL \sin^2\left(rac{\omega T}{2}
ight)$$

Sensitive to GW frequencies ~ 1/T ~ Hz



GW soundscape today





Conventional GW soundscape ~2040



CERN 08/21



GW soundscape (~2040s) with atom interferometers



Badurina, Buchmueller, Ellis, Lewicki, CM, Vaskonen Phil.Trans.Roy.Soc.Lond., arXiv:2108.02468



Example: sensitivity to binary mergers (equal masses)



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Badurina, Buchmueller, Ellis, Lewicki, CM, Vaskonen Phil.Trans.Roy.Soc.Lond., arXiv:2108.02468



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Ongoing work: mitigating backgrounds

Badurina, Gibson, CM, Mitchell, PRD, arXiv:2211.01854 Carlton, CM, to appear



Short-termer challenge: operating in a university building



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Many moving 'test masses'

ULDM searches run for many months

Could the busy environment hide a fundamental physics signal?



Mitigation through data analysis



Preliminary: Carlton, CM, to appear

Simple strategy works: mask noisy periods in analysis

Loss in sensitivity small since:

$$d_{m_e}^{\text{best}} \sim \left(\frac{1}{T_{\text{int}}}\right)^{1/4}$$

Recover shot-noise limited sensitivity

Exploring methods to keep all data



Seismic activity induces Gravity Gradient Noise (GGN)

Expectation: will limit low-frequency searches

Rayleigh waves give the largest density variations so considered the most dangerous

Harms, Living Rev.Rel.18 (2015) 3, arXiv:1507.05850; Baker et al, arXiv:1201.5656; Vetrano et al, arXiv:1304.1702; Harms et al, arXiv:1308.2074; Chaibi et al, arXiv:1601.00417; Junca et al (MIGA), arXiv:1902.05337

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Longer-term challenge: seismic noise













(Partially) mitigated with multi-gradiometer configuration

GGN signal decays exponentially from the surface $\Phi_{\text{Rayleigh}} = \left(\widetilde{A}e^{-qkz_0} + \widetilde{B}e^{-kz_0} \right)$

ULDM (or GW) signals scale linearly with AI separation

$$\Phi_{\rm ULDM} \sim \frac{\Delta z}{L}$$

Cross-correlation methods to search for the linear signal

Badurina, CM, et al, PRD, arXiv:2211.01854

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Multi-gradiometer: probe depth-scaling of signal and background



Badurina, CM, et al, PRD, arXiv:2211.01854

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ULDM Projections for km-baseline

ASN = best-case sensitivity

Blue: New High Noise Model with **two** interferometers

Other curves: New High Noise Model with *five* interferometers

Increased sensitivity for ~0.1 to 1 Hz



Historically, new observational techniques have led to new discoveries

Ultralight dark matter probe

- Mass $< 10^{-12} \, \text{eV}$

- Scalar-, vector- and pseudoscalar-coupled DM candidates

Mid-band gravitational wave detection

- LIGO sources before they reach LIGO band
- Early-Universe cosmological sources

And more...

- Tests of quantum mechanics at macroscopic scales
- Probe of seismic activity...

Summary

- Time-varying energy shifts, EP-violating new forces, spin-coupled effects







Science and Technology Facilities Council



A wide landscape of DM candidates



US Cosmic Visions





Speed distribution in our galaxy

Many models also predict some substructure in the distribution, see e.g., O'Hare et al arXiv:1807.09004, 1810.11468, 1909.04684



Coherence of the field



Impact of the speed distribution apparent over long time-scales: field amplitude evolves with a 'coherence time' $\tau \sim (m_{\rm DM} \sigma_v^2)^{-1}$

Al signals depend on the field amplitude \Rightarrow will also vary with a coherence time



B – L coupled vector appears in many extensions of the Standard Model

As ULDM, this generates background 'dark electric field':

In a *dual-species interferometer*, isotopes experience a different forces (accelerations):

Graham et al arXiv:1512.06165

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 $E_{B-L} \sim \cos(m_{\rm DM}t + \theta)$

$$\Delta F_{B-L} \sim g_{B-L} \left(\frac{Z_1}{A_1} - \frac{Z_2}{A_2}\right) E_{B-L}$$





Other ULDM signals (1): Near- and long-term prospects (Vector)



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log₁₀[*m*/eV]

Abe et al arXiv:2104.02835



Other ULDM signals (2): Precession of spins (Pseudoscalar)

Light pseudoscalar (axions) are ubiquitous in extensions of the Standard Model

In a dual-species interferometer, pseudoscalars couple to the different spin of the isotopes:

Phase
$$\sim (m_{S,1} - m_{S,2}) \cos(m_a t + \theta)$$

Challenging: km-baseline, highrepetition rate (10 Hz), long interrogation time, good control of magnetic fields $\delta B \sim 10^{-15} \mathrm{T}$



Graham et al arXiv:1709.07852

Beyond oscillating field signals...



See also Riedel et al arXiv:1212.3061, 1609.04145

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Some work in the direction of the DM abundance

e.g. 'Thermal misalignment of scalar dark matter' Batell & Ghalsaia, arXiv:2109.04476







Complementarity with atomic clocks





Multi-gradiometer configuration





Lots of space for multiple atom interferometers on km-baseline!



Effect of Rayleigh waves

Model wave travelling across the surface as:

$$\vec{\xi}(\varrho,\theta,z,t) = \left(\xi_H(z)\hat{k} - \xi_V(z)\vec{e}_z\right)e^{i(k\varrho\cos(\theta-\theta')-\theta')}$$

Horizontal displacement Vertical displacement

Induces density fluctuations below the surface:

 $\frac{\delta\rho(z>0)}{\rho_0} = \left[\xi_V\delta(z) + \mathcal{R}(z)\right]e^{i(k\rho\cos(\theta-\theta')-\omega t)}$

$$\mathcal{R}(z) = k\xi_V \frac{(q^2 - 1)}{q} \left(\frac{1 + s^2}{1 - s^2}\right) e^{-qkz} \quad \text{wher}$$



re $q, s \sim \mathcal{O}(1)$



Effect of Rayleigh waves

Density fluctuations imply a time dependent gravitational potential:

$$\left\langle \delta\phi\left(z_{0},t\right)\right\rangle = -2\pi G\rho_{0}\,\xi_{V}\,e^{-i\omega t}\,\frac{1}{qk}\,\left($$

Vertical displacement

Induces a phase in the interferometers:

$$\Phi_{\text{Rayleigh}} = \left(\widetilde{A}e^{-qkz_0} + \widetilde{A} \right)$$

Amplitude decays exponentially

 $\left(\frac{1+s^2}{1-s^2}\right)\left((1+\sqrt{q/s})e^{-kz_0}-2e^{-qkz_0}\right)$

Amplitude decays exponentially with depth

 $\widetilde{B}e^{-kz_0}$ $\xi_V \cos(\omega T + \Theta)$

Vertical displacement





Badurina et al, arXiv:2211.01854

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Projections for km-long baseline

ASN = target sensitivity

Orange: Peterson's New High Noise Model

Blue: Peterson's New Low Noise Model

с_н parameterises decay length of Rayleigh wave density variation:

$$\lambda_{\rm GGN} = \frac{c_H}{\omega_a} \simeq 100 \,\mathrm{m} \,\left(\frac{250 \,\mathrm{m\,s}^{-1}}{c_H}\right)^{-1} \,\left(\frac{2.5 \,\mathrm{Hz}}{\omega_a}\right)^{-1} \,$$

 10^{1}







