

Linear Colliders Lecture 3 Subsystems II



Frank Tecker – CERN

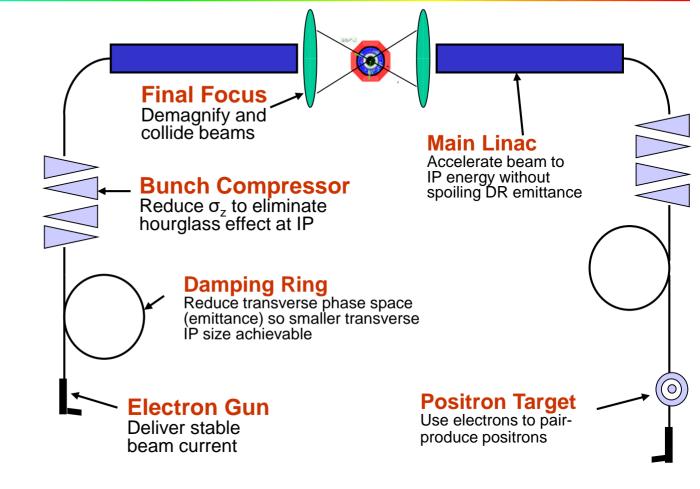
- Main Linac (cont.)
 - RF system and technology
 - Accelerating gradient
- Beam / bunch structure
- Beam Delivery System
- Alignment and Stabilization



Previous Lecture



- Particle production
- Damping rings with wiggler magnets
- Bunch compressor
 with magnetic chicane
- ⇒ small, short bunches
 to be accelerated
 w/o emittance blowup
- Main linac:
 - longitudinal wakefields => energy spread, chromatic effects
 - Transverse wakefields, minimized by structure design
- Now: Acceleration in the linac

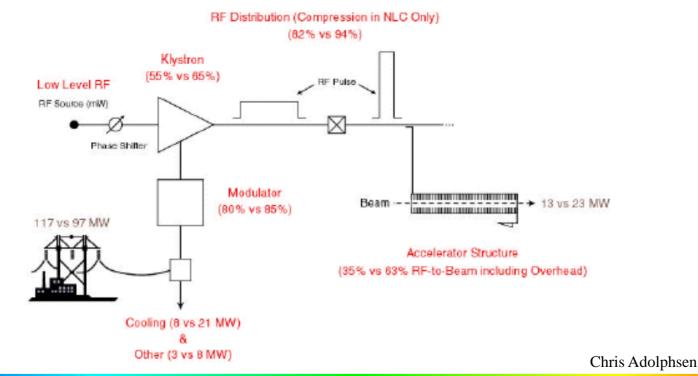




RF systems



- Need efficient acceleration in main linac
- 4 primary components:
 - Modulators: convert line $AC \rightarrow pulsed DC$ for klystrons
 - ◆ Klystrons: convert DC → RF at given frequency
 - ◆ RF distribution: transport RF power → accelerating structures evtl. RF pulse compression
 - ◆ Accelerating structures: transfer RF power → beam



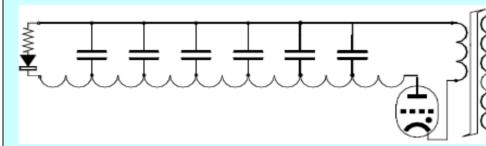


RF systems



Modulator

Energy storage in capacitors charged up to 20-50 kV (between pulses)



High voltage switching and voltage transformer rise time > 300 ns

Or solid state device

Klystron

U 150 -500 kV *I* 100 -500 A *f* 0.2 -20 GHz

 $P_{ave} < 1.5 \text{ MW}$ $P_{peak} < 150 \text{ MW}$

efficiency 40-70%

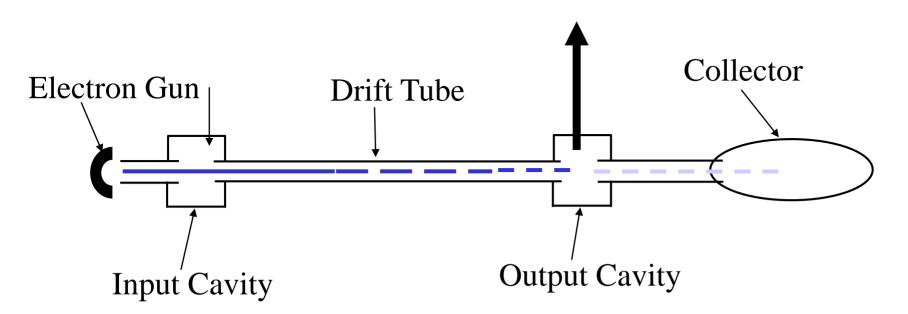
=> for power efficient operation pulse length $t_P >> 300$ ns favourable



Klystrons



- narrow-band vacuum-tube amplifier at microwave frequencies (an electron-beam device).
- low-power signal at the design frequency excites input cavity
- Velocity modulation becomes time modulation in the drift tube
- Bunched beam excites output cavity





RF efficiency: cavities



- Fields established after cavity fill time
- Only then the beam pulse can start
- Steady state: power to beam, cavity losses, and (for TW) output coupler

$$hotal P_{RF o beam} = \frac{P_{beam}}{P_{beam} + P_{loss} + P_{out}} \frac{T_{beam}}{T_{fill} + T_{beam}}$$

 ≈ 1 for SC SW cavities

• NC TW cavities have smaller fill time T_{fill}



SC Technology



- In the past, SC gradient typically 5 MV/m and expensive cryogenic equipment
- TESLA development: new material specs, new cleaning and fabrication techniques, new processing techniques
- Significant cost reduction
- Gradient substantially increased
- Electropolishing technique has reached ~35 MV/m in 9-cell cavities
- 31.5 MV/m ILC baseline
- limited by critical magnetic field H_{crit} above which no superconductivity exists



Chemical polish



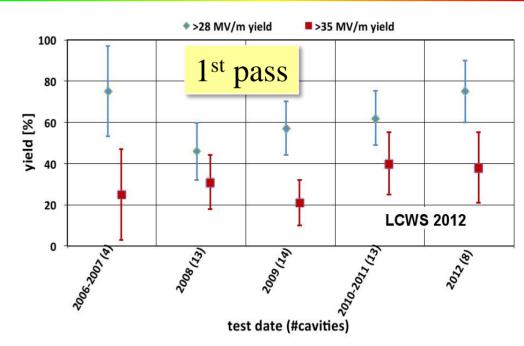
Electropolishing

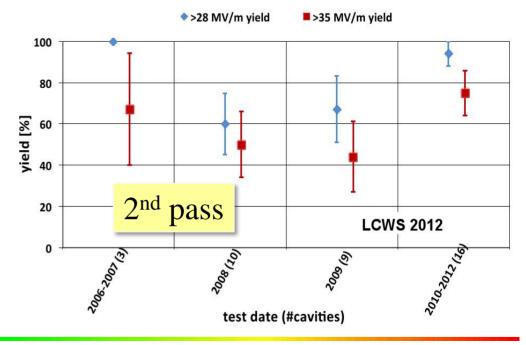


Achieved SC accelerating gradients



- Large progress by R&D program to systematically understand and set procedures for the production process
- reached goal for a 50% yield at 35 MV/m by the end of 2010
- 90% yield at 28 MV/m exceeded in 2012
- Tests for higher gradient ongoing
- limited certainly below
 MV/m (H_{crit})
- X-FEL running with 23.6 MV/m







Limitations of Gradient E_{acc}



Surface magnetic field

- SC structures become normal conducting above H_{crit}
- NC: Pulsed surface heating => material fatigue => cracks
- Field emission due to surface electric field
 - Vacuum arcs RF break downs
 - Break down rate => Operation efficiency
 - Local plasma triggered by field emission => Erosion of surface
 - Dark current capture
 - => Efficiency reduction, activation, detector backgrounds

RF power flow

• RF power flow and/or iris aperture apparently have a strong impact on achievable E_{acc} and on surface erosion. Mechanism not fully understood

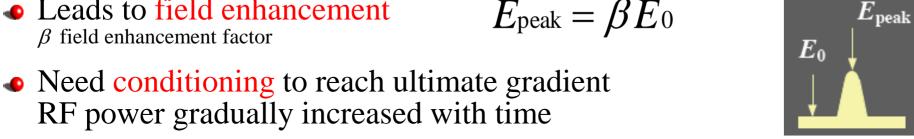


NC Structure conditioning

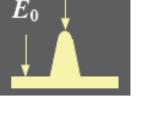


- Material surface has some intrinsic roughness (from machining)
- Leads to field enhancement β field enhancement factor

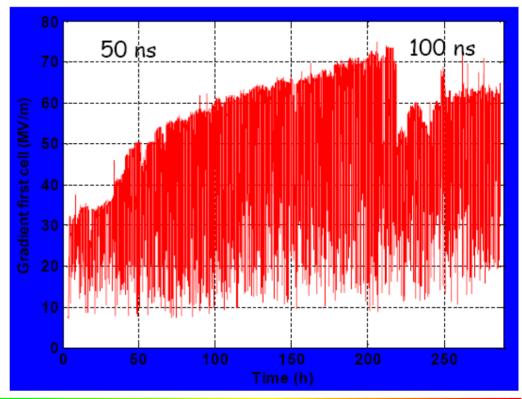
$$E_{\text{peak}} = \beta E_0$$



- RF processing can melt field emission points
 - Surface becomes smoother
 - field enhancement reduced
 - => higher fields less breakdowns
- More energy: Molten surface splatters and generates new field emission points!
- Excessive fields can also damage the structures



from S.Doebert





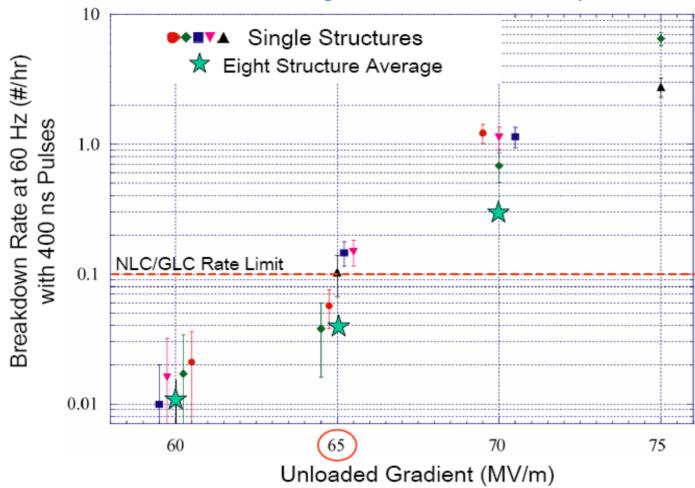
Breakdown-rate vs gradient



Strong increase of breakdown rate for higher gradient

High Gradient Performance

5 Structures after ~ 500 hr of Operation and 8 Structure Average after > 1500 hr of Operation



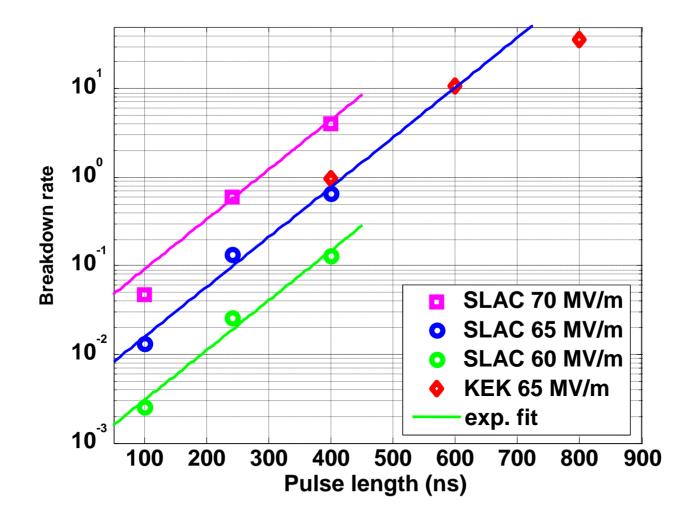
C. Adolphsen /SLAC



Breakdown-rate vs pulse length



• Higher breakdown rate for longer RF pulses



• Summary: breakdown rate limits pulse length and gradient

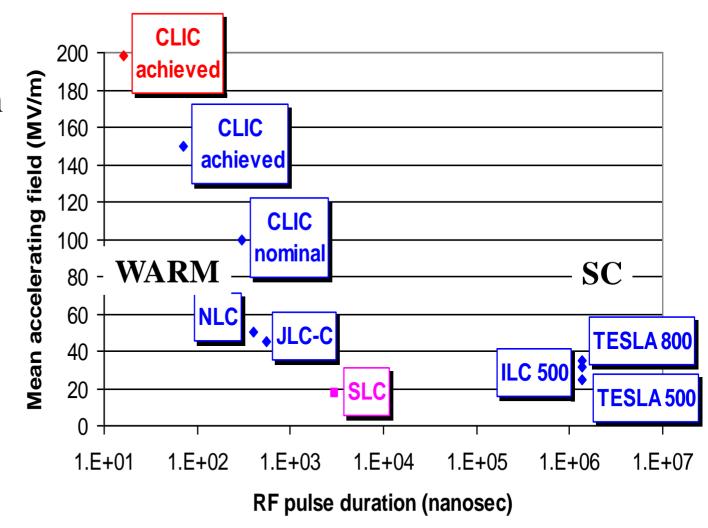


Accelerating gradient



- Normal conducting cavities have higher gradient with shorter RF pulse length
- Superconducting cavities have lower gradient (fundamental limit) with long RF pulse

Accelerating fields in Linear Colliders

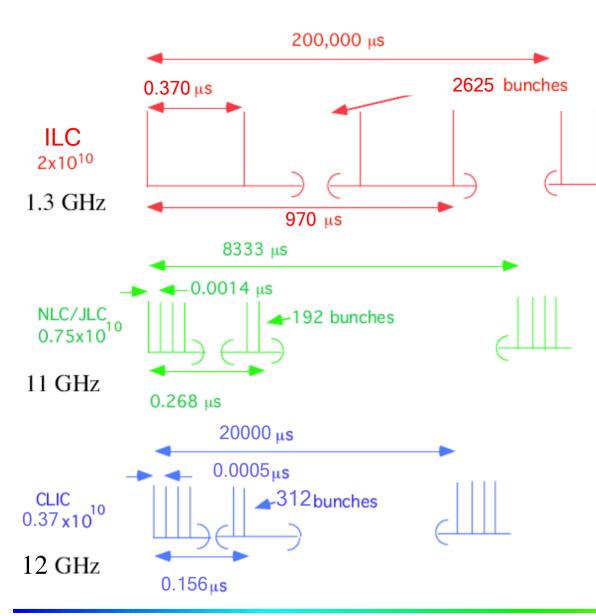




Bunch structure



• SC allows long pulse, NC needs short pulse with smaller bunch charge

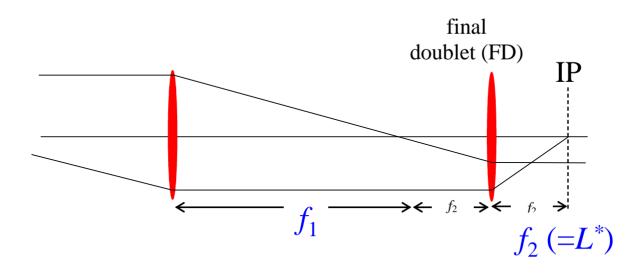


The different RF technologies used by ILC, NLC/JLC and CLIC require different packaging for the beam power



Beam Delivery: Final Focus





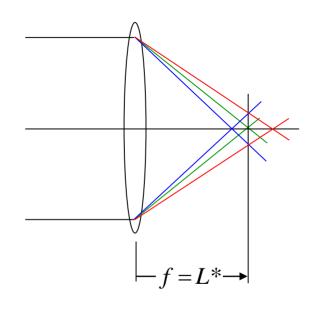
- Need large demagnification of the (mainly vertical) beam size $M = \sqrt{\beta_{linac}/\beta_y^*} = f_1/f_2$ typical value ≈ 300
- β_y^* of the order of the bunch length σ_z (hour-glass effect)
- Need free space around the IP for physics detector
- Assume $f_2 = 2 \text{ m} => f_1 \approx 600 \text{ m}$
- Can make shorter design but this roughly sets the length scale



Final Focus: chromaticity



- Need strong quadrupole magnets for the final doublet
- Typically hundreds of Tesla/m
- Get strong chromatic aberations



for a *thin-lens* of length *l*:

$$\frac{1}{f} \approx k_1 l$$

change in deflection:

$$Dy'_{quad} \approx -k_1 l y_{quad} \frac{d}{1+d} \approx -k_1 l y_{quad} d$$

change in IP position:

$$Dy_{IP} \approx f Dy'_{quad} = y_{quad} d$$

RMS spot size:

$$\langle Dy_{IP}^2 \rangle = \langle y_{quad}^2 \rangle \langle O^2 \rangle = D_{quad} e_y O_{rms}^2$$



Final focus: Chromaticity



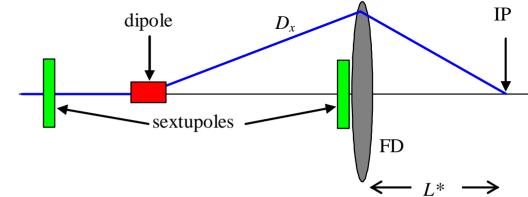
- Small $\beta^* => \beta_{FD}$ very large (~ 100 km)
- for $\frac{7M}{\text{rms}}$ ~ 0.3%

$$\sqrt{\left\langle \Delta y_{IP}^2 \right\rangle} \approx 20 - 40 \text{ nm}$$

- Definitely much too large
- We need to correct chromatic effects
- => introduce sextupole magnets

$$B_x = s x y$$

$$B_y = \frac{1}{2} s \left(x^2 + y^2 \right)$$



• Use dispersion *D*:

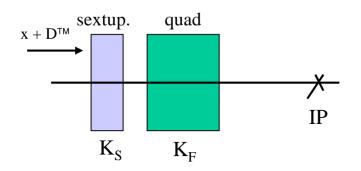
$$x = x_o + D\delta$$



Chromaticity correction



• Combine quadrupole with sextupole and dispersion



Second order

y plane straightforward x plane more tricky

$$Dx' = \frac{K_F}{(1+O)}(x+DO) \Rightarrow K_F(-Ox-DO^2)$$

$$Dx' = \frac{K_s}{2}(x + Dd)^2 \implies K_s D(dx + \frac{Dd^2}{2})$$

Quad:
$$Dx' = \frac{K_F}{(1+d)}(x+Dd) \Rightarrow K_F(-dx-Dd^2)$$

$$\text{chromaticity} \qquad \downarrow \qquad \downarrow$$

$$\text{chromaticity} \qquad \downarrow \qquad \downarrow$$

$$\text{Sextupole:} \qquad Dx' = \frac{K_S}{2}(x+Dd)^2 \Rightarrow K_SD(dx+\frac{Dd^2}{2})$$

$$Dx' = \frac{K_F}{(1+d)}(x+Dd) + \frac{K_{b\text{-match}}}{(1+d)}x \Rightarrow 2K_F(-dx-\frac{Dd^2}{2})$$

Could require $K_S = K_F/D$

=> ½ of second order dispersion left

Create as much chromaticity as FD upstream

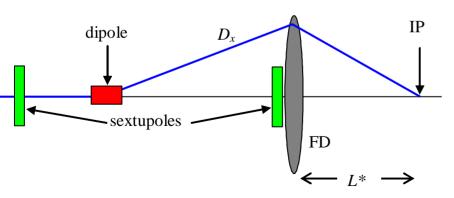
$$K_{b\text{-match}} = K_F$$
 $K_S = \frac{2K_F}{D}$

=> second order dispersion corrected

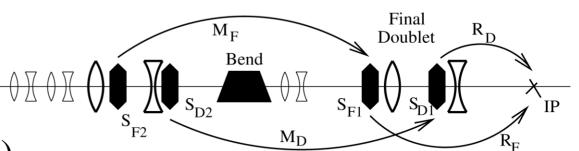


Final Focus: Chromatic Correction

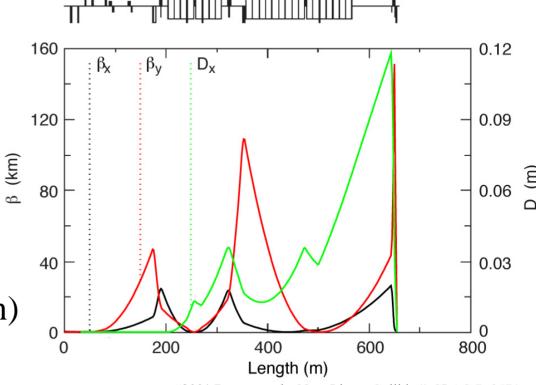




Correction in both planes



- Relatively short (few 100 m)
- Local chromaticity correction
- High bandwidth (energy acceptance)
- FF tested at ATF2 (KEK Japan)
 - 44 nm achieved (37 nm design)
 - scales to 6 nm at ILC (5 nm)



"2001 Report on the Next Linear Collider", SLAC-R-0571



Final focus: fundamental limits



- From the hour-glass effect: $b_y \circ S_z$
- For highest energies, additional fundamental limit: synchrotron radiation in the final focusing quadrupoles => beamsize growth at the IP
- so-called *Oide Effect*:

minimum beam size:
$$\sigma \approx 1.83 \left(\frac{r_e \lambda_e}{2\pi} F\right)^{1/7} \varepsilon_n^{5/7}$$

• for

$$\beta \approx 2.39 \left(\frac{r_e \lambda_e}{2\pi} F\right)^{2/7} \varepsilon_n^{3/7}$$

 λ_e is the Compton wavelength of the electron

F is a function of the focusing optics: typically $F \sim 7$ (minimum value ~ 0.1)

• $\sigma_{\text{oide}} = 0.85 \text{ nm for } 3 \text{ TeV CLIC}$



Stability and Alignment



- Tiny emittance beams, nm vertical beam size at collision
- Any quadrupole misalignment and jitter will cause orbit oscillations and displacement at the IP

$$\Delta y^* = \sum_{i}^{Quads} k_{Q,i} \Delta y_{Q,i} \sqrt{\frac{\gamma_i}{\gamma^*}} \sqrt{\beta_i \beta^*} \sin(\Delta \phi_i)$$

- Tight component tolerances
 - Field quality
 - Alignment
- Vibration and Ground Motion issues
- Active stabilisation
- Feedback systems

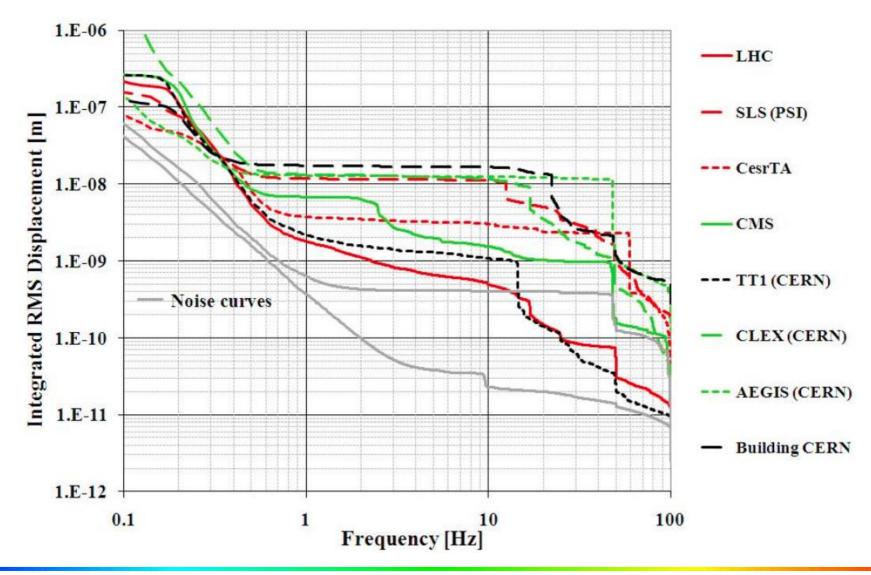
- Some numbers (CLIC):
 - Cavity alignment (RMS) 17 μm
 - Main Beam quad alignment: 14 μm
 - vert. MB quad stability:1.5 nm @>1 Hz
 - hor. MB quad stability: 5 nm @>1 Hz
 - Final quadrupole: 0.15 nm @>4 Hz !!!



Ground Motion



 Site dependent ground motion with decreasing amplitude for higher frequencies





Ground motion: ATL law



- Need to consider short and long term stability of the collider
- Ground motion model: ATL law

$$\langle \Delta y^2 \rangle = ATL$$

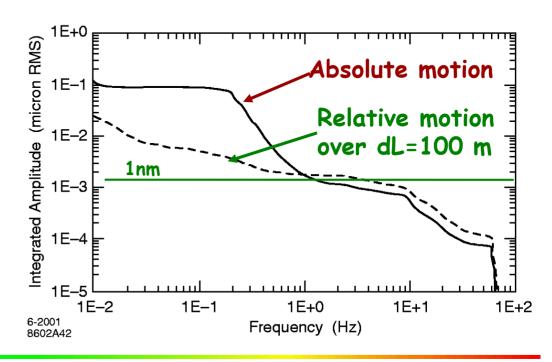
A site dependent constant

T time

L distance

A range 10^{-5} to $10^{-7} \, m\text{m}^2/\text{m/s}$

- This allows you to simulate ground motion effects
- Relative motion smaller
- Long range motion less disturbing

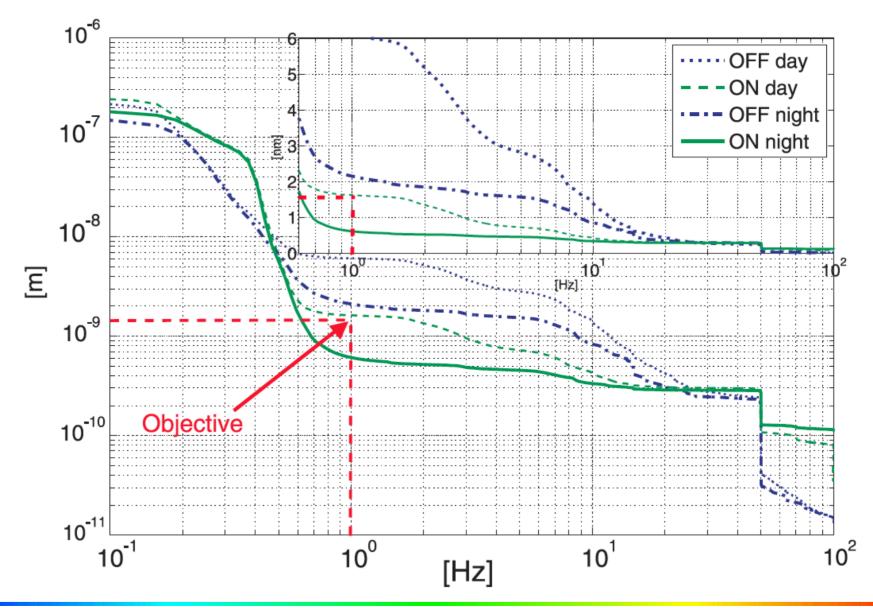




Active stabilization



• Test bench reaches required stability of CLIC MB quadrupole

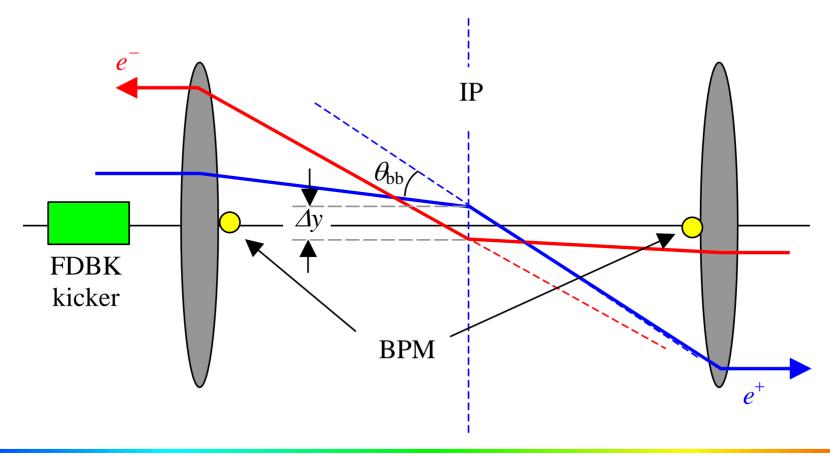




Beam-Beam feedback



- Use the strong beam-beam deflection kick for keeping beams in collision
- Sub-nm offsets at IP cause well detectable offsets (micron scale)
 a few meters downstream





Other IP issues



Collimation:

- Beam halo will create background in detector
- Collimation section to eliminate off-energy and off-orbit particle
- Material and wakefield issues

• Crossing angle:

- NC small bunch spacing requires crossing angle at IP to avoid parasitic beam-beam deflections
- Luminosity loss ($\approx 10\%$ when $(= f_x/f_z)$

Crab cavities

Introduce additional time dependent transverse kick to improve collision

Spent beam

- Large energy spread after collision
- Design for spent beam line not easy



Post-Collision Line (CLIC)



R.B. Appleby, A. Ferrari, M.D. Salt and V. Ziemann, Phys. Rev. ST Accel. Beams 12 (2009) 021001.

Baseline: vertical chicane with 2x4 dipoles

- Separation by dipole magnets of the disrupted beam, beamstrahlung photons and particles with opposite sign from coherent pairs, from low energy tails
 - → Short line to prevent the transverse beam size from growing too much
 - → Intermediate dumps and collimator systems
- Back-bending region with dipoles to direct the beam onto the final dump
 - → Long line allowing non-colliding beam to grow to acceptable size

