

# The Nelson-Barr Axion Dark Matter

(converting the LHC+B-factories into quantum sensors)

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# Outline

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- Intro. the strong CP problem + ultralight dark-matter
- The Nelson-Barr solution to the strong CP problem
- The strong-CP axion
- New pheno' new signatures for the LHC+B-factories
- Challenges of the model

# How serious is the strong CP problem ?

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● 3 levels of formulating the strong CP problem:

(i)  $\bar{\theta} = \theta - \arg \left[ \det (Y_u Y_d) \right] \lesssim 10^{-10}$ , is it a problem?

(who knows?)

(ii)  $\bar{\theta} = \lesssim 10^{-10} \ll \theta_{\text{KM}} = \arg \left\{ \det \left[ Y_u Y_u^\dagger, Y_d Y_d^\dagger \right] \right\} = \mathcal{O}(1)$ , is it a problem?

(not if these are natural/protected and sequestered)

(iii)  $\bar{\theta} = \lesssim 10^{-10} \ll \theta_{\text{KM}}$ , but  $\bar{\theta} = \bar{\theta}_{\text{bare}} + \epsilon \theta_{\text{KM}} \ln (\Lambda_{\text{UV}}/M_W)$ , is it a problem?

( $\epsilon$  appears in 7 loops and contains several other suppression factor)

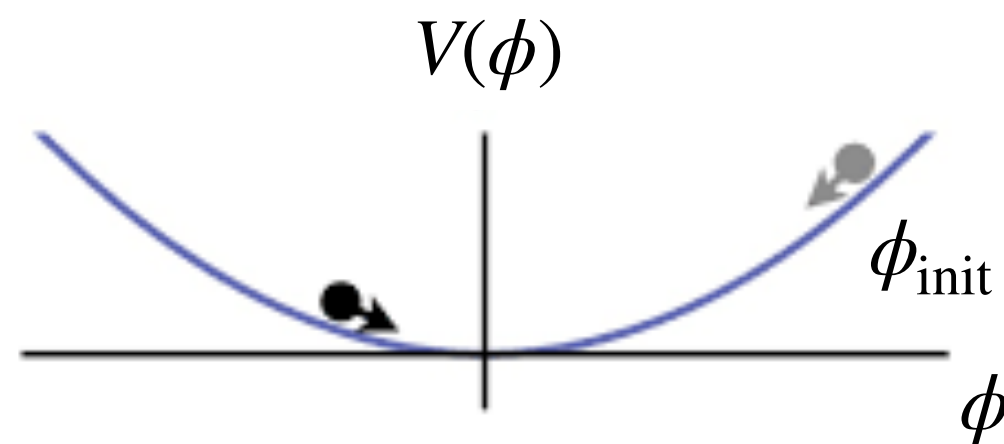
● Should we be cautious [at least till we reach  $\mathcal{O}(10^{-16})$  precision]

# We nevertheless focus on axions

- Begin with ultralight dark matter (UDM), minimal model would be just a free massive scalar:

$$\mathcal{L} \in m_\phi^2 \phi^2, \rho_{\text{Eq}}^{\text{DM}} \sim \text{eV}^4 \sim m_\phi^2 \phi_{\text{Eq}}^2 = m_\phi^2 \phi_{\text{init}}^2 (\text{eV}/T_{\text{osc}})^3 \quad \left[ T_{\text{os}} \sim \sqrt{M_{\text{Pl}} m_\phi} \right]$$

minimal misalignment mechanism



- Assuming (“best case”) MeV reheating:  $\phi_{\text{init}} (f_{\text{min}}) = \begin{cases} 10^{17} \text{ GeV} \left( \frac{10^{-27} \text{ eV}}{m_\phi} \right)^{\frac{1}{4}} & m_\phi \lesssim 10^{-15} \text{ eV} \\ 10^{15} \text{ GeV} \left( \frac{10^{-15} \text{ eV}}{m_\phi} \right) & m_\phi \gtrsim 10^{-15} \text{ eV} \end{cases}$

# Solving the QCD problem *not* with QCD axion

- There's a class of models where CP is UV-sym' and at tree level we find:

$$\bar{\theta} = \theta - \arg \left[ \det (Y_u Y_d) \right] = 0 \quad \& \quad \theta_{\text{KM}} = \arg \left\{ \det \left[ Y_u Y_u^\dagger, Y_d Y_d^\dagger \right] \right\} = \mathcal{O}(1)$$

- This is realized if:

- Yukawas are Hermitian (left-right models or wave function renorm')

Hiller & Schmaltz (01); Harnik, GP, Schwartz & Shirman (04); Cheung, Fitzpatrick & Randall (08)

- Structure/sym.  $\Rightarrow$  det(0), concretely, Nelson-Barr (NB)

Nelson; Barr (84)

- We focus on NB, which are easy to control & of higher quality

# Nelson-Barr (crash course)

- $\mathcal{L}_{\text{NB}} = \mu q^c q + (g_i \Phi + \tilde{g}_i \Phi^*) d_i^c q + Y_d H Q d^c + Y_u \tilde{H} Q u^c$  (with  $q, q^c, \Phi \subset Z_2$  – odd)

- Assume that theory is real and only  $\Phi = \frac{f + \rho}{\sqrt{2}} \exp\left(\frac{ia}{f}\right)$ ;  $\langle a \rangle \neq 0$  breaks CP, then:

1.  $\mathcal{M}_d = \begin{pmatrix} \mu & B_i \\ 0 & m_d \end{pmatrix}$ ;  $m_d \equiv Y_d v$ ;  $B_i \equiv (g_i \Phi + \tilde{g}_i \Phi^*) \Rightarrow \det [\mathcal{M}_d] \in \text{Real}$

2. At low energy ( $v \ll \mu, g_i f$ ), effective  $m_d$  satisfies  $m_d^{\text{eff}} m_d^{\text{eff}\dagger} = m_d \left( \mathbf{1}_3 + \frac{B_i^* B_j}{\mu^2 + B_f B_f^\dagger} \right) m_d^\dagger$ ,

which if  $g_i$  isn't parallel to  $\tilde{g}_i$  and  $\mu \lesssim B_i$  lead to  $\theta_{\text{KM}} = \mathcal{O}(1)$



# Nelson-Barr axion-like pheno for the CP breaking

Discussion with: M. Dine, Y. Nir, W. Ratzinger, I. Savoray

- $\mathcal{L}_{\text{NB}} = \mu q^c q + (g_i \Phi + \tilde{g}_i \Phi^*) d_i^c q + Y_d H Q d^c + Y_u \tilde{H} Q u^c$  (with  $q, q^c, \Phi \in Z_2$  – odd)
- Assume approx' flavor sym' such that  $g_i \propto (1,0,0)$  &  $\tilde{g}_i \propto (0,0,1)$
- Then  $a$  is a pseudo-Nambu-Goldstone-boson, with suppressed potential, but with  $\langle a \rangle = 0$
- Furthermore, one can show that  $\theta_{\text{KM}} = \frac{a}{f}$  Involved the 3rd generation

$$\left\{ m_d^{\text{eff}} m_d^{\text{eff}\dagger} \sim m_d \left[ \mathbf{1}_3 + r \begin{pmatrix} 1 & 0 & e^{\frac{2ia}{f}} \\ 0 & 0 & 0 \\ e^{-\frac{2ia}{f}} & 0 & 1 \end{pmatrix} \right] m_d^T \right\}$$
- Also, mixing angles develop quadratic dependence on  $a$  (but not masses)

# Nelson-Barr ultralight-DM pheno

With: M. Dine, Y. Nir, W. Ratzinger, I. Savoray (also discussion with Surjeet)

- In case another sector breaks the shift sym' (say Planck suppress or other) then the minimum of potential generically would lead to  $\langle a \rangle \neq 0$  and spontaneous breaking of CP  $\Rightarrow \bar{\theta} = 0$  &  $\theta_{\text{KM}} = \mathcal{O}(1)$ 
  - Relaxion: Graham, Kaplan & Rajendran (15)
  - NB-relaxion - Davidi, Gupta, GP, Redigolo, & Shalit (17)
- Now if we tip the NB-axion from it's minimum it'd behave as a new type of ultralight DM



New type of pheno: *time dependent CKM angles*

While the strong CP is always zero



# NB-UDM signature & parameter space

- What is the size of the effect?  $\delta\theta_{\text{KM}} \sim \frac{\sqrt{\rho_{\text{DM}}}}{m_{\text{NB}}f} \cos(m_{\text{NB}}t) \sim 10^{-3} \times \frac{10^{10} \text{ GeV}}{f} \times \frac{10^{-19} \text{ eV}}{m_{\text{NB}}} \times \cos(m_{\text{NB}}t)$
- Currently (PDG):  $\theta_{\text{KM}} = 1.14 \pm 0.03$
- How to search such signal? Need time dependence CP violation, perfect for  $B$ -asym
- Bound from EP:  $\frac{\Delta m_u}{m_u} \approx \frac{3}{32\pi^2} y_b^2 |V_{ub}^{\text{SM}}|^2 \frac{a}{f} \Rightarrow f \gtrsim 10^{10} \text{ GeV}$
- Minimal misalignment DM bound, can't be satisfied:  $f \gtrsim 10^{15} \text{ GeV} \left( \frac{10^{-19} \text{ eV}}{m_\phi} \right)^{\frac{1}{4}}$
- Naive naturalness  $\Rightarrow$  sub-MeV cutoff,  $\Delta m_a \approx \frac{y_b |V_{ub}| m_u \Lambda_{\text{UV}}}{16\pi^2 f} \Rightarrow$  current B-factories probe finely tuned region

# Conclusions

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- The strong CP problem is marginally a problem but it is interesting to try and solve it
- Nelson-Barr models account for the smallness of the strong CP phase & the fact the KM phase is order one, which requires spontaneous CP violation
- Spontaneous breaking lead to the presence of a light axion-like field
- If this field consist of ultralight dark matter it'd lead to new type of pheno', with time-dependent CKM angles (involving the 3rd gen')
- Maybe be probed by the LHC/B-factories