

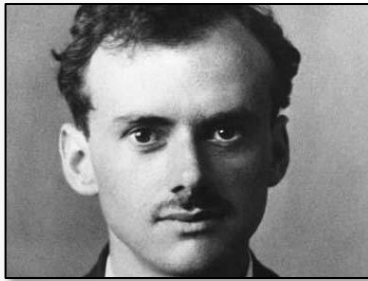
# Sterile Neutrinos in Neutrinoless Double Beta Decay

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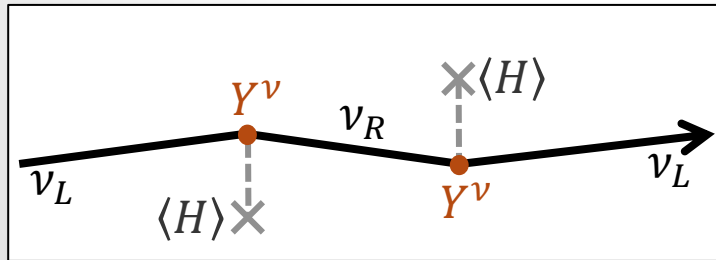
University College London

# Dirac versus Majorana

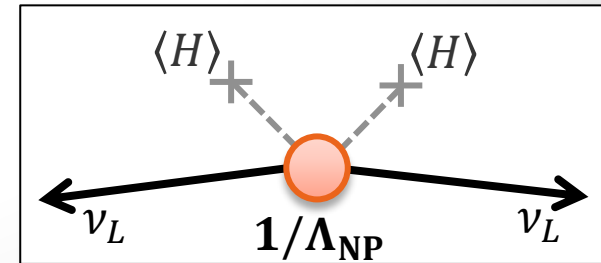
- ▶ Origin of neutrino masses beyond the Standard Model
- ▶ Two possibilities to define neutrino mass



Dirac mass analogous to other fermions but with  $m_\nu / \Lambda_{EW} \approx 10^{-12}$  couplings to Higgs



Majorana mass, using only a left-handed neutrino  
→ Lepton Number Violation

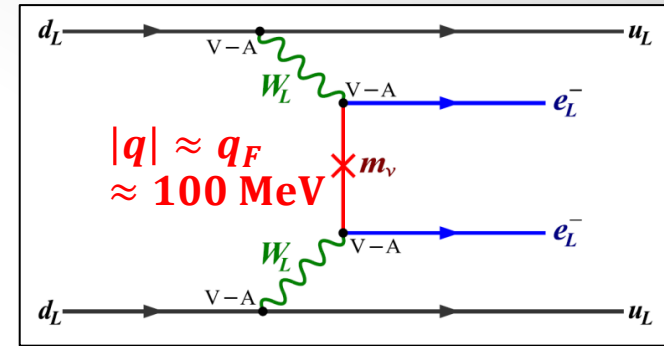


# Neutrinoless Double $\beta$ Decay

▶ Half-life

$$T_{1/2}^{-1} = |m_{\beta\beta}|^2 G^{0\nu} |M^{0\nu}|^2$$

▶ Particle Physics



$$\mathcal{A}_{\mu\nu}^{lep} = \frac{G_F^2}{4} \sum_{i=1}^3 U_{ei}^2 \gamma_\mu (1 + \gamma_5) \frac{\not{q} + m_{\nu_i}}{q^2 - m_{\nu_i}^2} \gamma_\nu (1 - \gamma_5) \approx G_F^2 \frac{\gamma_\mu (1 + \gamma_5) \gamma_\nu}{4q^2} \sum_{i=1}^3 U_{ei}^2 m_{\nu_i} \rightarrow m_{\beta\beta}$$

▶ Atomic Physics

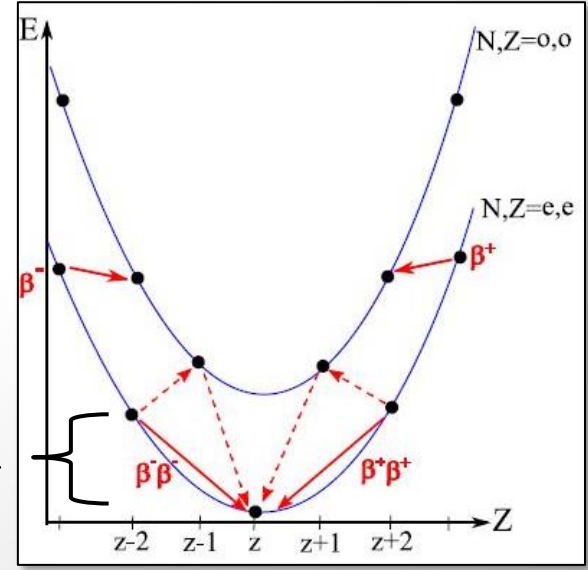
- Leptonic phase space  $G^{0\nu} \propto Q^5$

$$\Gamma_{0\nu\beta\beta} \sim m_\nu^2 G_F^4 q_F^2 Q^5 \sim (m_\nu / 0.1 \text{ eV})^2 (10^{26} \text{ yr})^{-1}$$

▶ Nuclear Physics

- Nuclear matrix element  $M^{0\nu} \approx 1$  but large uncertainties

$Q + 2m_e \approx 3-5 \text{ MeV}$

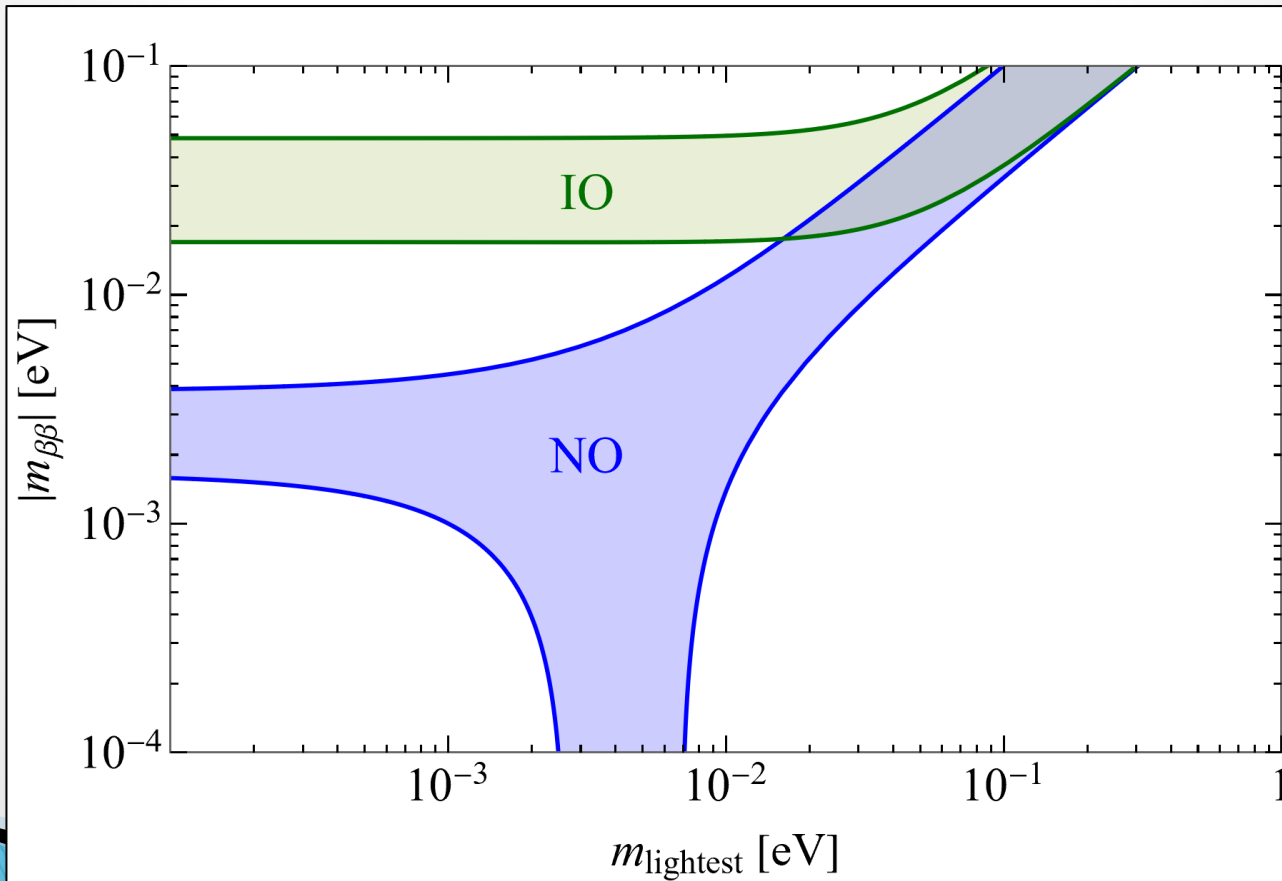


# Three Active Neutrinos

## ▶ Effective $0\nu\beta\beta$ Mass

degenerate &  $\theta_{13} \approx 0$

$$|m_{\beta\beta}| = |c_{12}^2 c_{13}^2 m_{\nu_1} + s_{12}^2 c_{13}^2 m_{\nu_2} e^{i\phi_{12}} + s_{13}^2 m_{\nu_3} e^{i\phi_{13}}| \approx m_{\nu} \sqrt{1 - \sin^2(2\theta_{12}) \sin^2(\phi_{12}/2)}$$

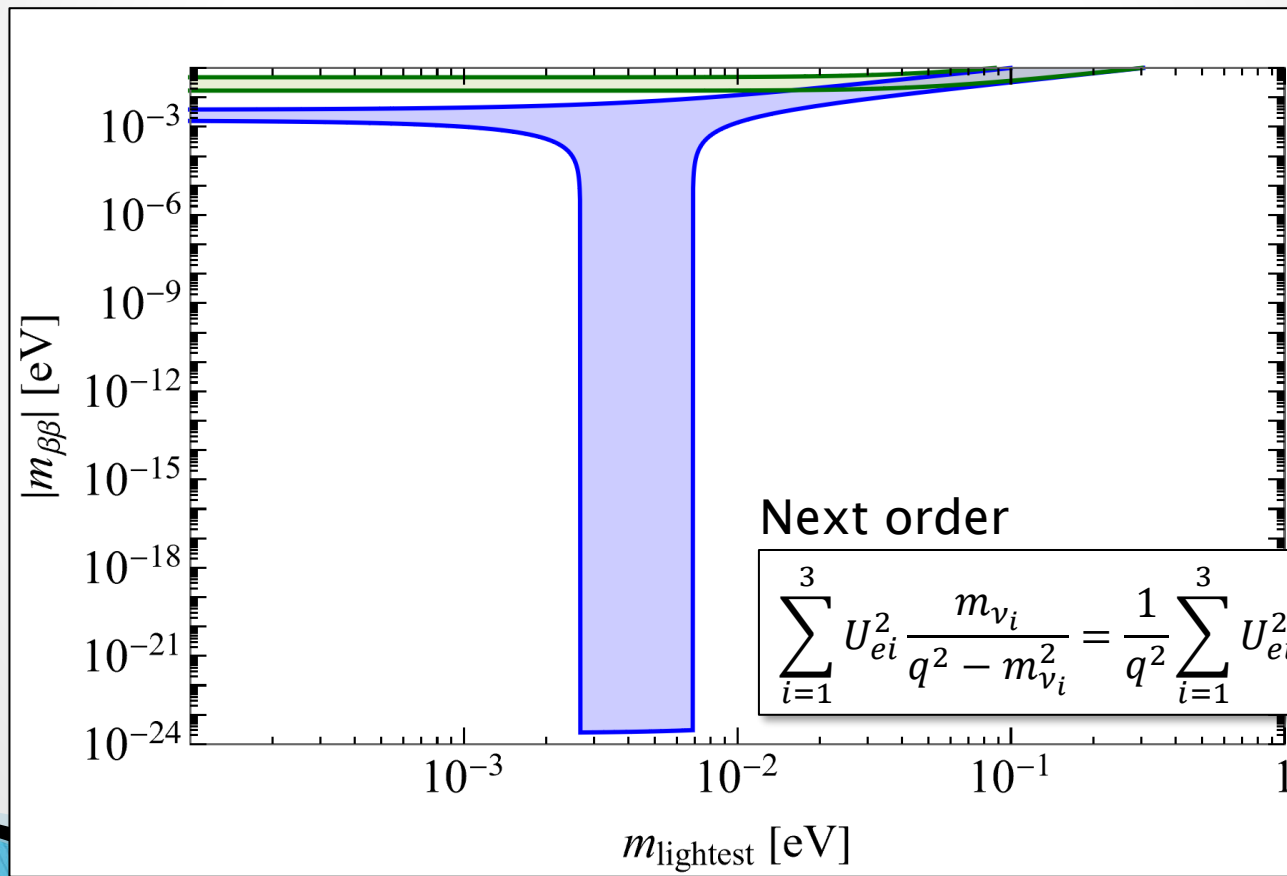


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# Sterile Neutrinos

## ▶ SM + Sterile Neutrinos

$$\mathcal{L} = \mathcal{L}_{\text{SM}} + i\bar{N}_{iR}\not{\partial}N_{iR} - (Y_\nu)_{\alpha i}\bar{L}_\alpha\tilde{H}N_{iR} - \frac{1}{2}(\mathcal{M}_S)_{ij}\bar{N}_{iR}^cN_{jR} + \text{h.c.}$$

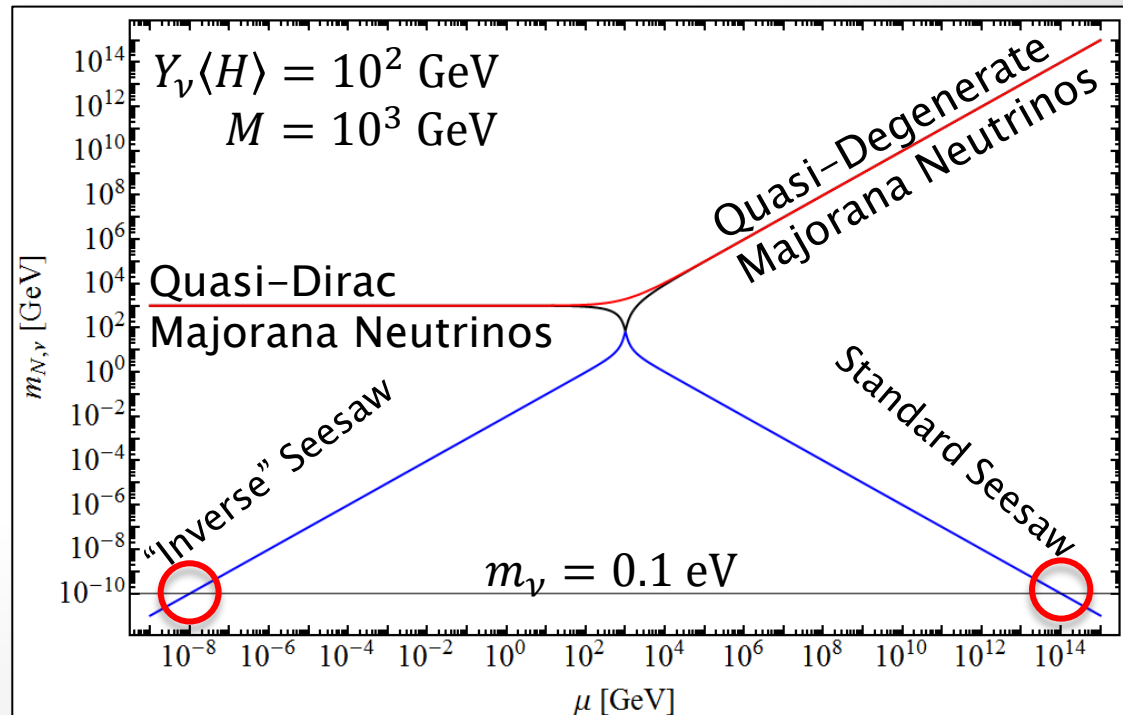
- Seesaw Mechanism with TeV scale heavy neutrinos
  - Standard Seesaw with small Yukawa couplings

$$V^{\nu N} \approx Y_\nu \approx 10^{-6} \sqrt{M_N/\text{TeV}}$$

- “Bent” Seesaw mechanisms
  - Decouple  $\Lambda_{\text{LNV}}$  from heavy neutrino mass

$\nu$	$N_1$	$N_2$
$0$	$Y_\nu\langle H \rangle$	$0$
$Y_\nu\langle H \rangle$	$\mu$	$M$
$0$	$M$	$\mu$

$$\mathcal{M} = \begin{pmatrix} 0 & Y_\nu\langle H \rangle & 0 \\ Y_\nu\langle H \rangle & \mu & M \\ 0 & M & \mu \end{pmatrix}$$

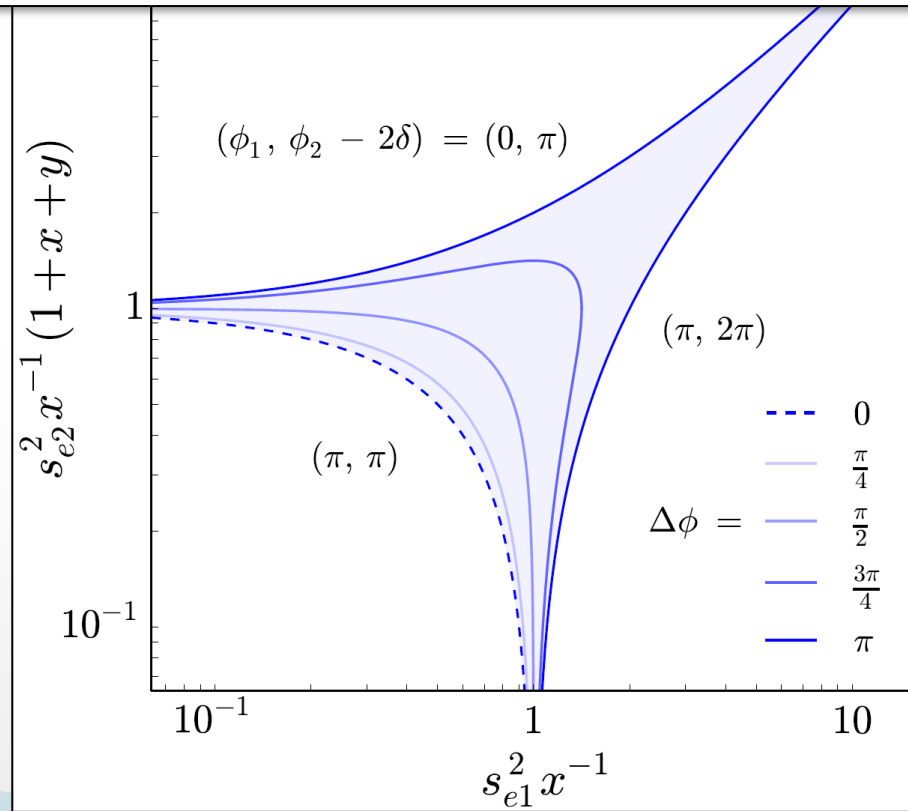




# 1ν + 2N: Pheno Parametrization

- ▶ Relation between active–sterile mixing angles (tree level)

$$(\mathcal{M}_\nu)_{\alpha\beta} = 0 \quad \Rightarrow \quad 0 = \sum_{i=1}^{\mathcal{N}_A} m_i U_{\alpha i} U_{\beta i} + \sum_{k=1}^{\mathcal{N}_S} m_{N_k} U_{\alpha N_k} U_{\beta N_k}$$



# Sterile Neutrinos in $0\nu\beta\beta$

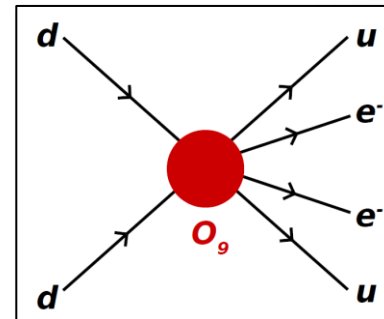
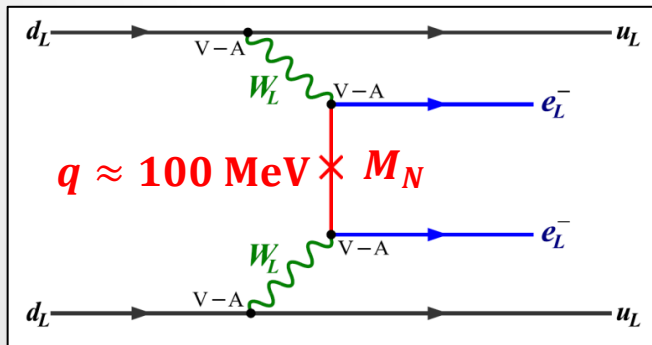
- ▶ Masses lighter than  $\approx 100$  MeV

$$|m_{\beta\beta}| = |c_{12}^2 c_{13}^2 m_{\nu_1} + s_{12}^2 c_{13}^2 m_{\nu_2} e^{i\phi_{12}} + s_{13}^2 m_{\nu_3} e^{i\phi_{13}} + s_{14}^2 m_{\nu_4} e^{i\phi_{14}} + \dots|$$

- ▶ Masses heavier than  $\approx 100$  MeV

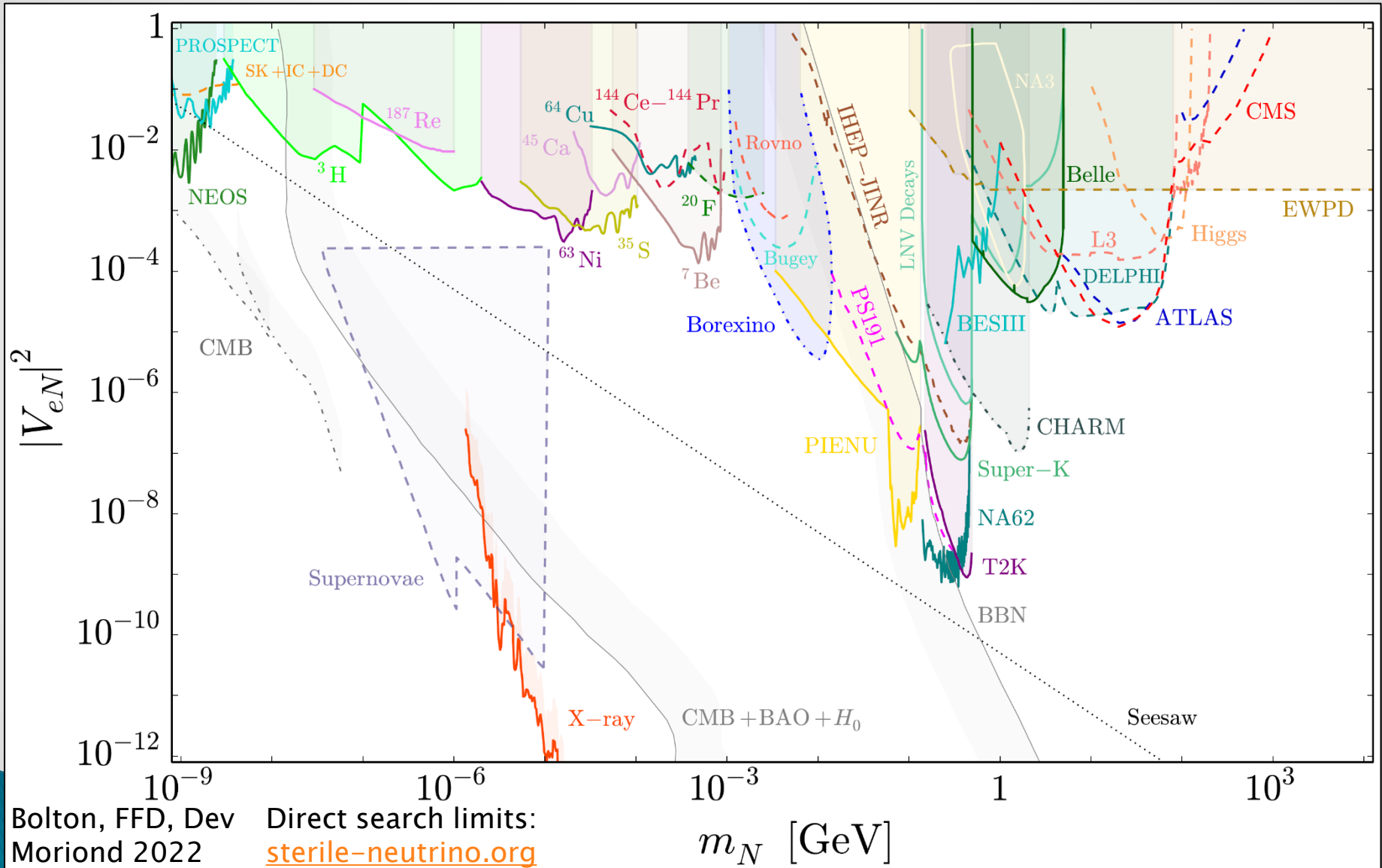
$$\mathcal{A}_{\mu\nu}^{lep} = \frac{1}{4} \sum_{i=1}^3 V_{ei}^2 \gamma_\mu (1 + \gamma_5) \frac{\not{q} + M_{N_i}}{q^2 - M_{N_i}^2} \gamma_\nu (1 - \gamma_5) \approx \frac{-\gamma_\mu (1 + \gamma_5) \gamma_\nu}{4} \sum_{i=1}^3 \frac{V_{ei}^2}{M_{N_i}} \rightarrow \left\langle \frac{1}{M_N} \right\rangle_{\beta\beta}$$

- ▶ Short-distance on nuclear scale

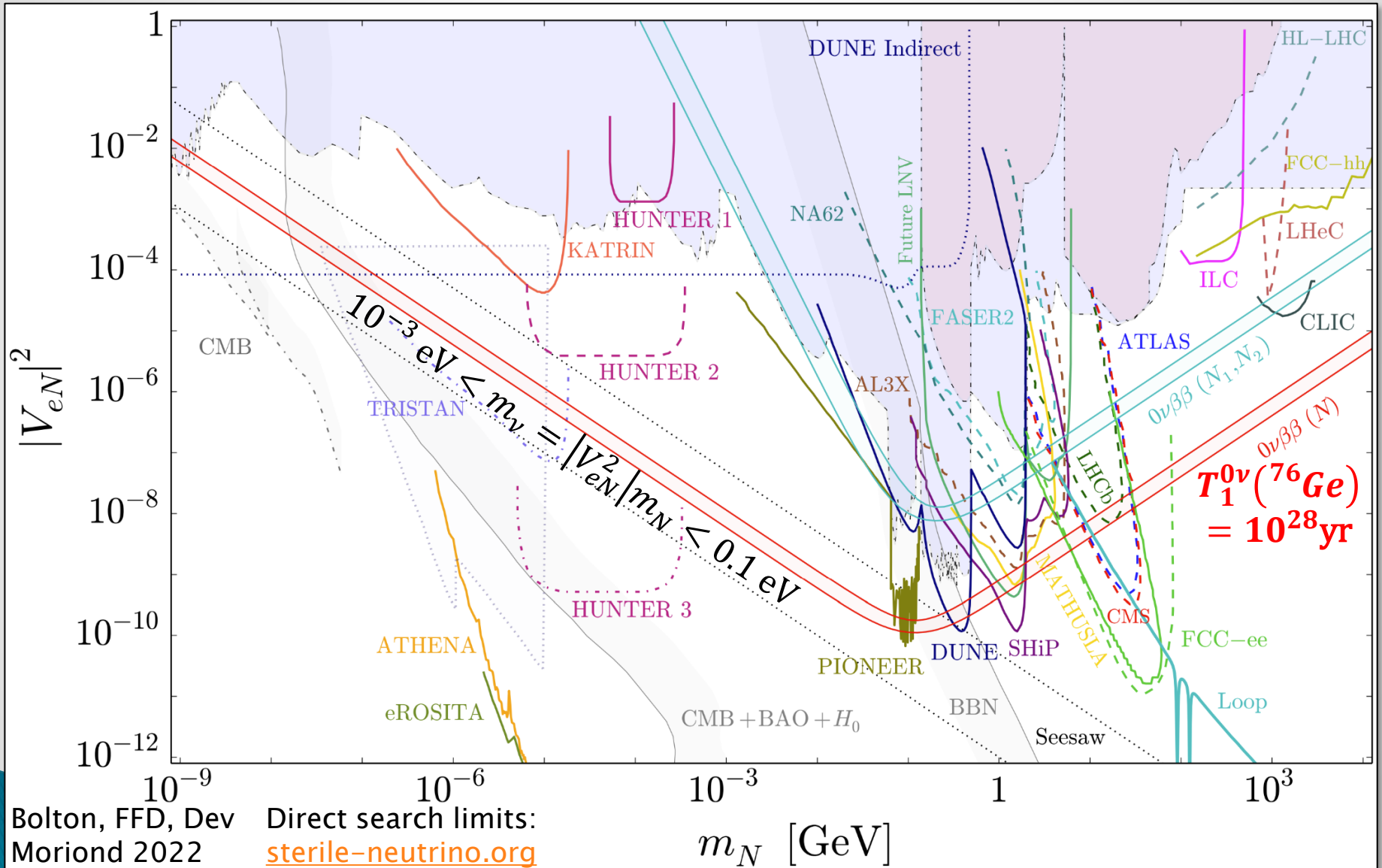




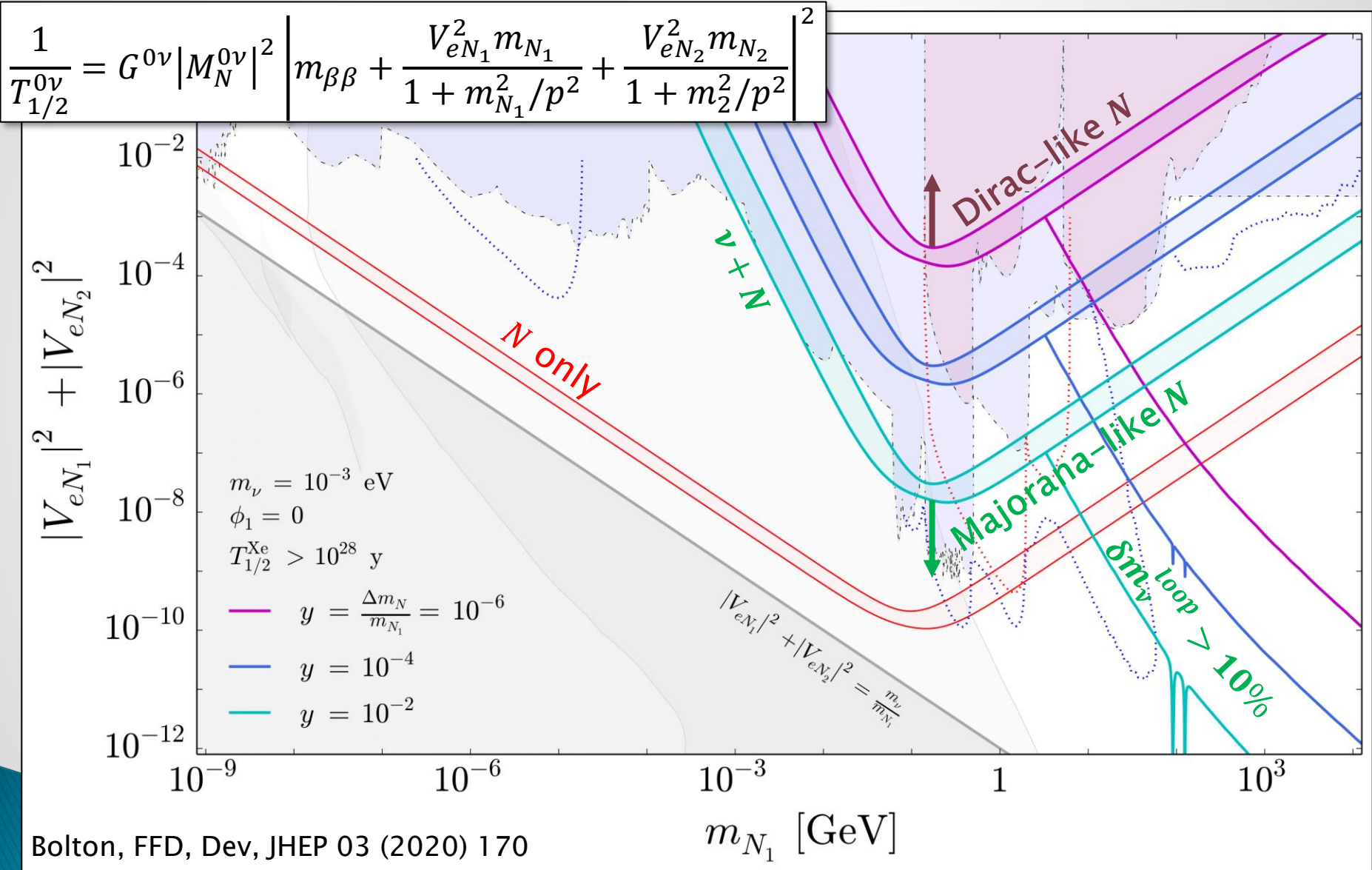
# HNL - Current Limits



# HNL – Future Sensitivities

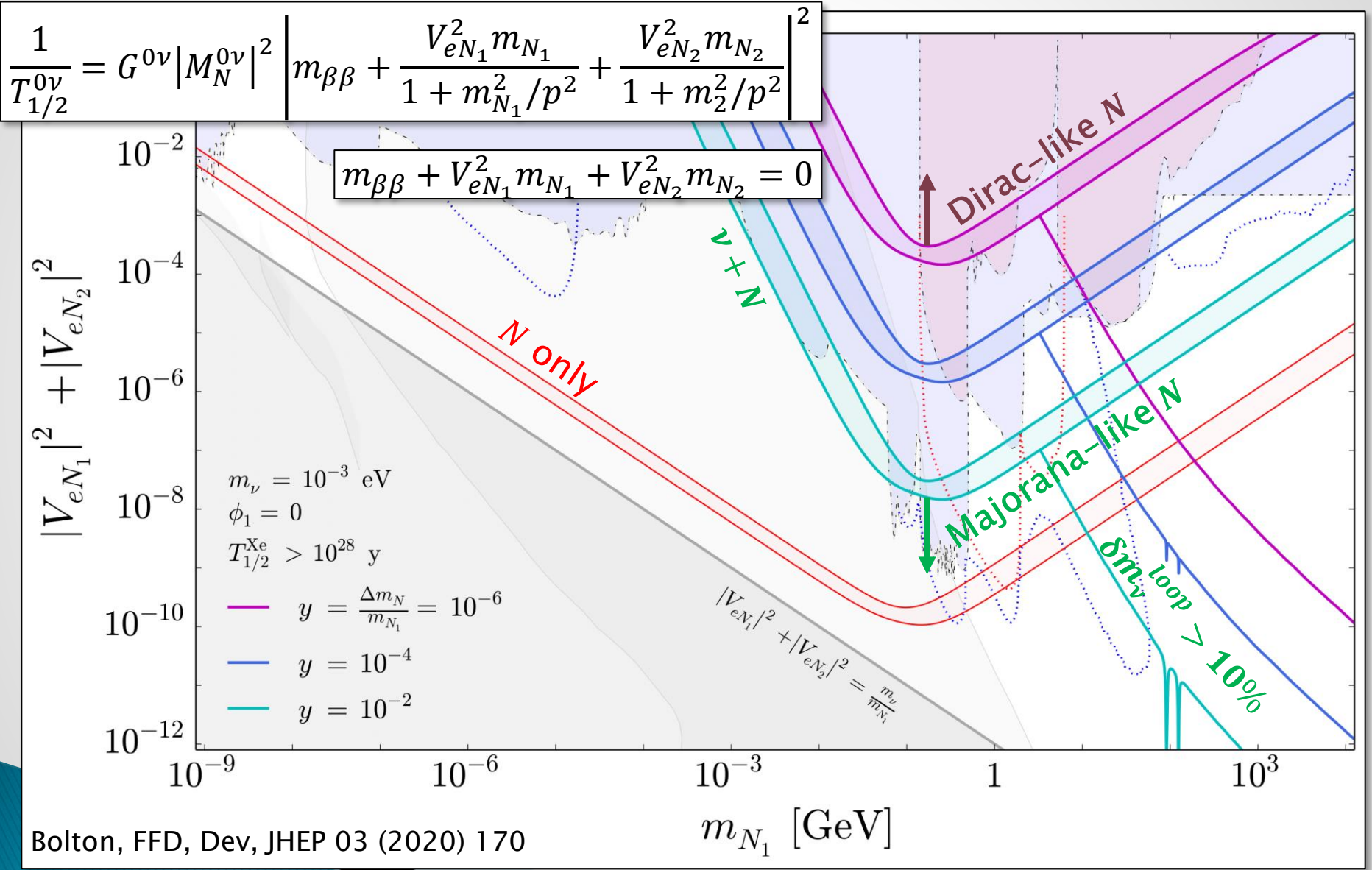


# HNL - Comparison with $0\nu\beta\beta$



Bolton, FFD, Dev, JHEP 03 (2020) 170

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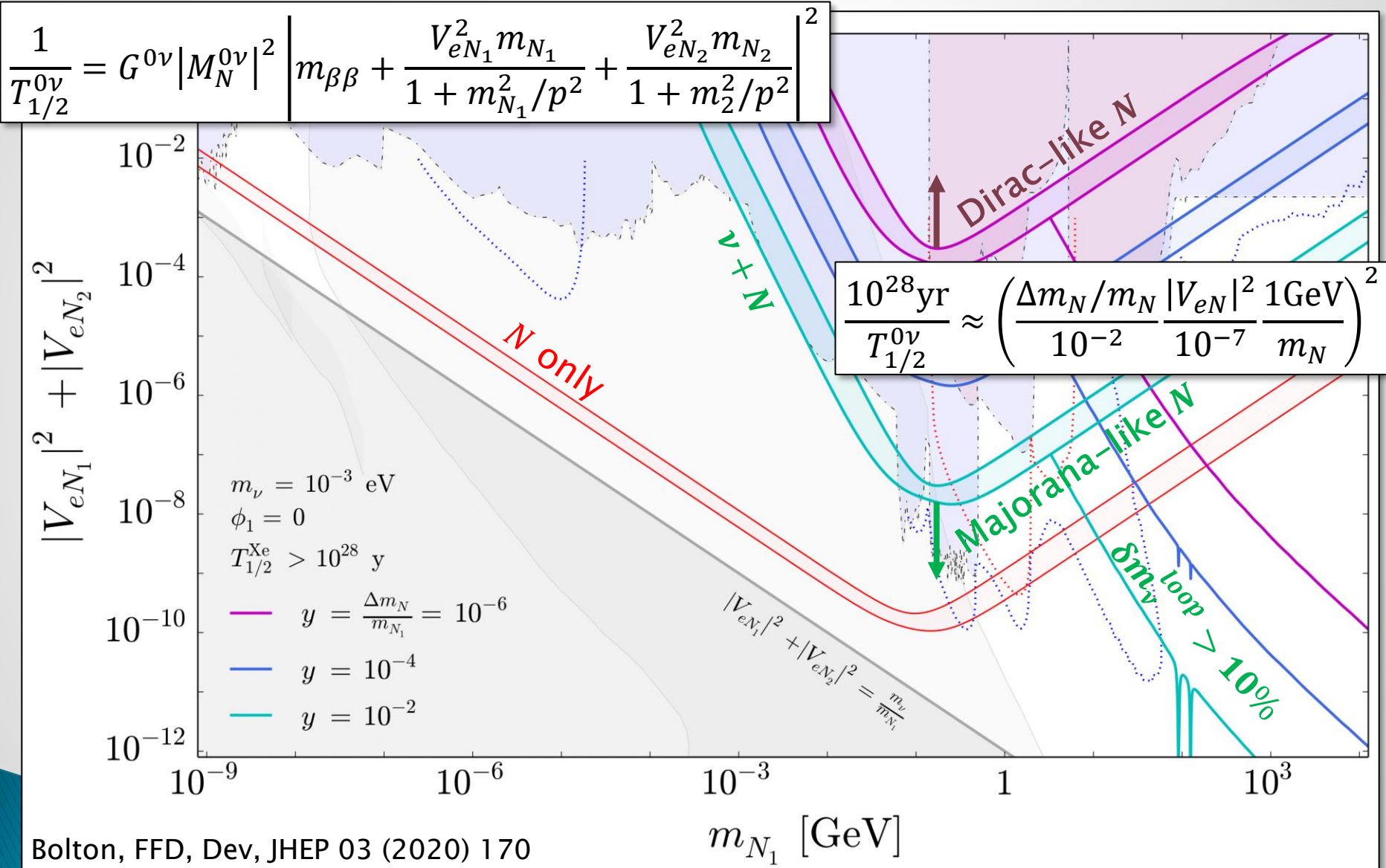


Bolton, FFD, Dev, JHEP 03 (2020) 170

$m_{N_1}$  [GeV]



# HNL - Comparison with $0\nu\beta\beta$



Bolton, FFD, Dev, JHEP 03 (2020) 170

# Sterile Neutrinos in $\chi$ EFT

Dekens, de Vries, Mereghetti, Menéndez, Soriano, Zhou, Phys.Rev.C 108 (2023) 4, 045501

- ▶ Improved calculation based on chiral EFT
  - Enhanced contributions from soft and hard neutrino exchange
  - Inclusion of excited intermediate nucleus states
  - Smaller suppression of Seesaw cancellation term

$$\frac{m_N^2}{k_F^2}$$

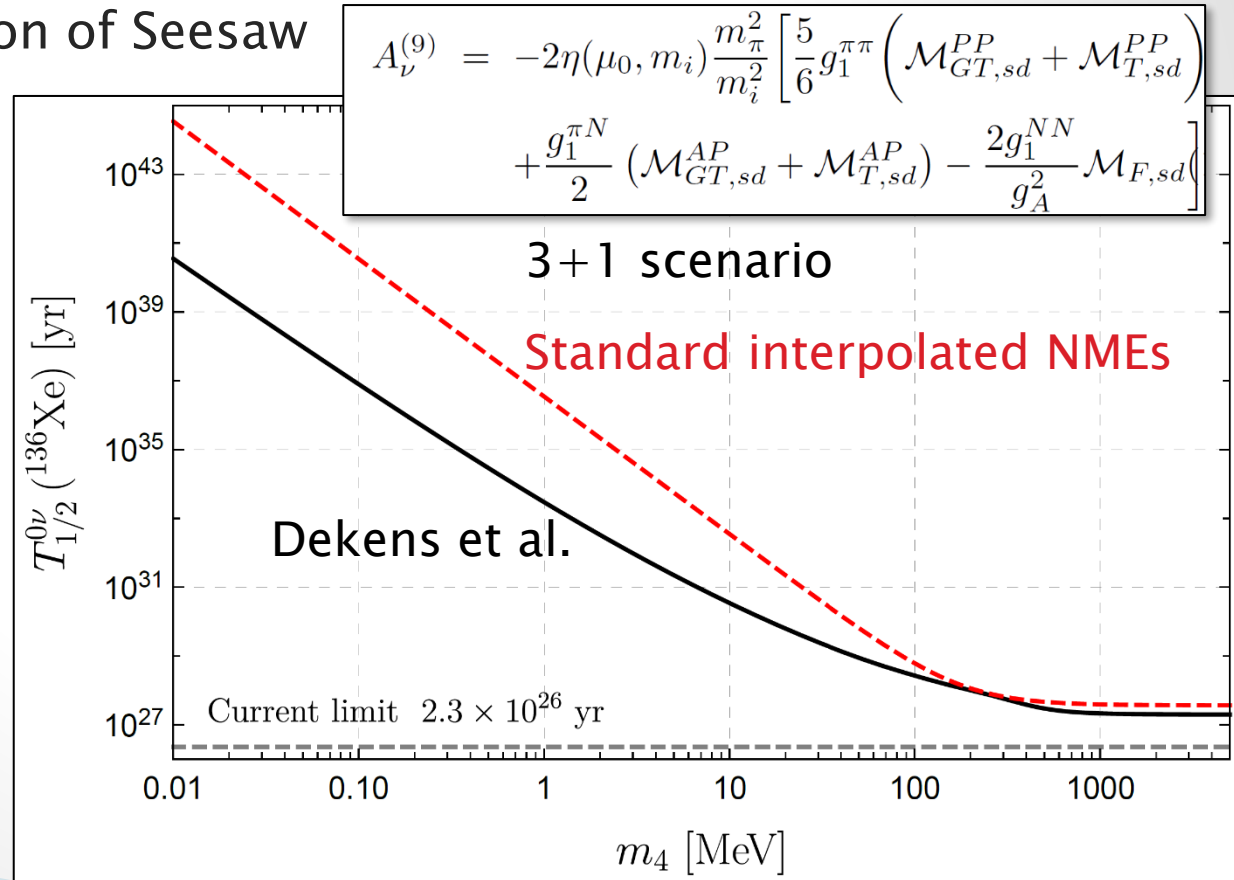
instead of

$$\frac{m_N^4}{k_F^4}$$

for  $m_N \ll k_F$

- Does not apply to bare sterile neutrino contrib.

$$A_\nu^{(9)} = -2\eta(\mu_0, m_i) \frac{m_\pi^2}{m_i^2} \left[ \frac{5}{6} g_1^{\pi\pi} \left( \mathcal{M}_{GT, sd}^{PP} + \mathcal{M}_{T, sd}^{PP} \right) + \frac{g_1^{\pi N}}{2} \left( \mathcal{M}_{GT, sd}^{AP} + \mathcal{M}_{T, sd}^{AP} \right) - \frac{2g_1^{NN}}{g_A^2} \mathcal{M}_{F, sd} \right]$$





# Sterile Neutrinos in $\chi$ EFT

Dekens, de Vries, Mereghetti, Menéndez, Soriano, Zhou, Phys.Rev.C 108 (2023) 4, 045501

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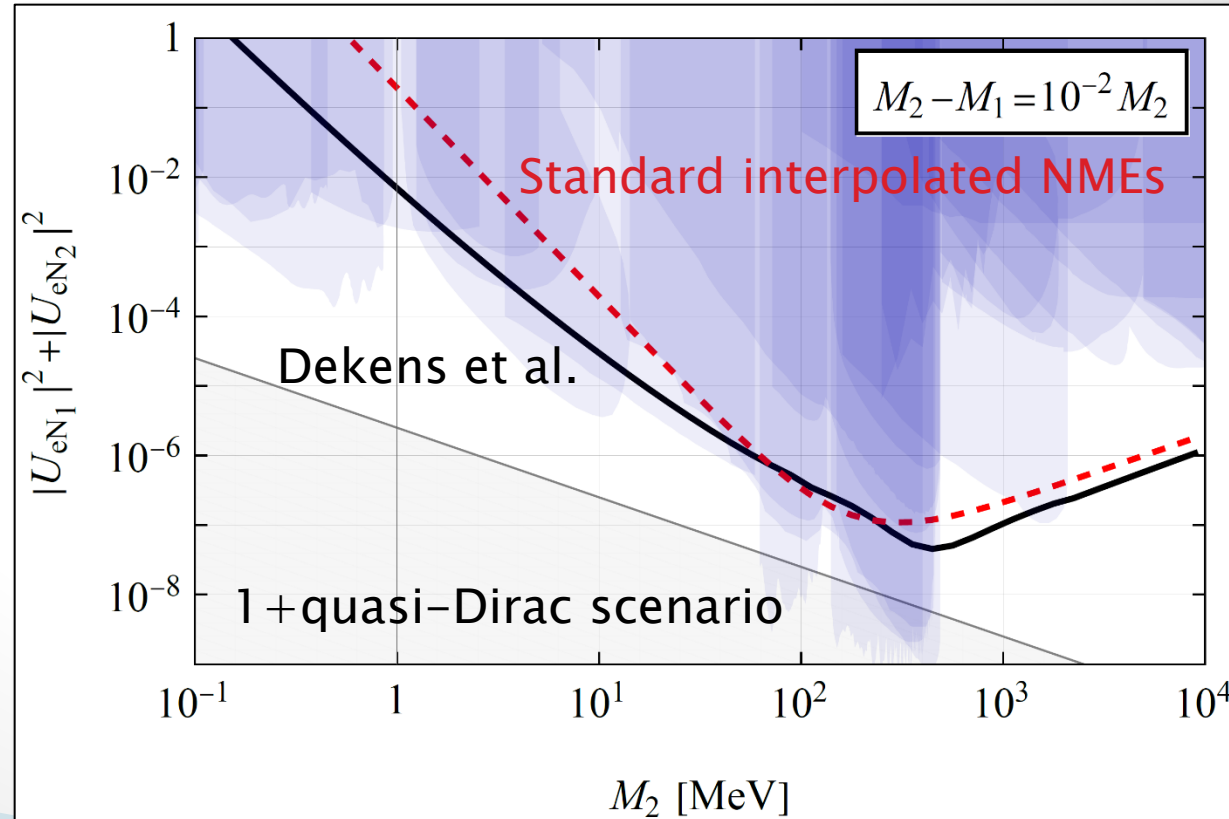
$$\frac{m_N^2}{k_F^2}$$

instead of

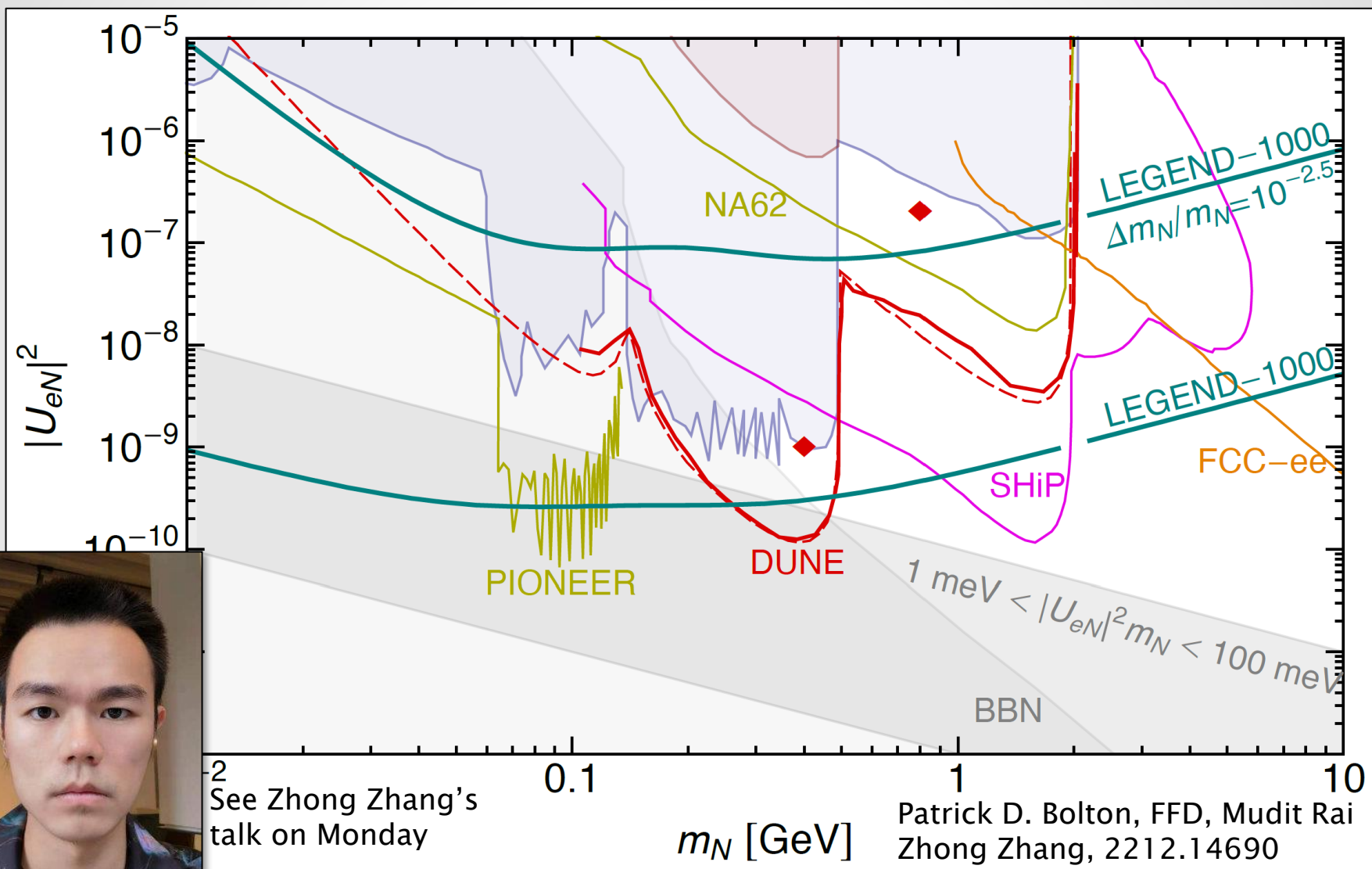
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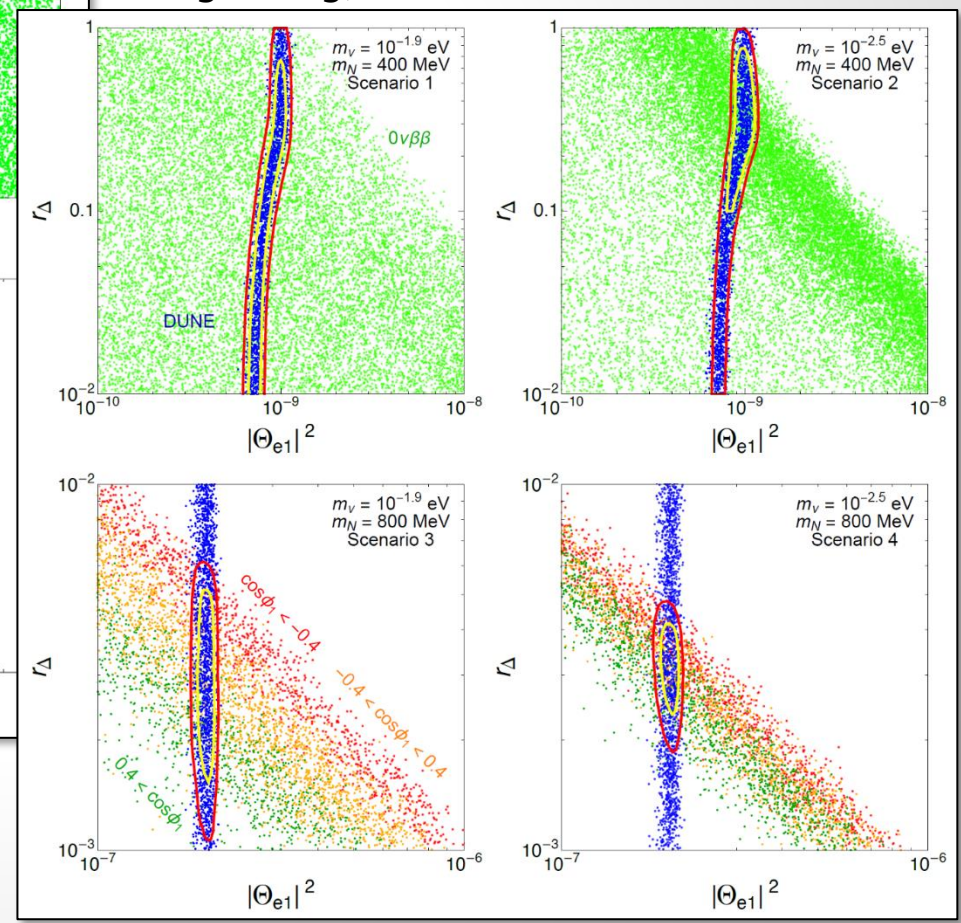
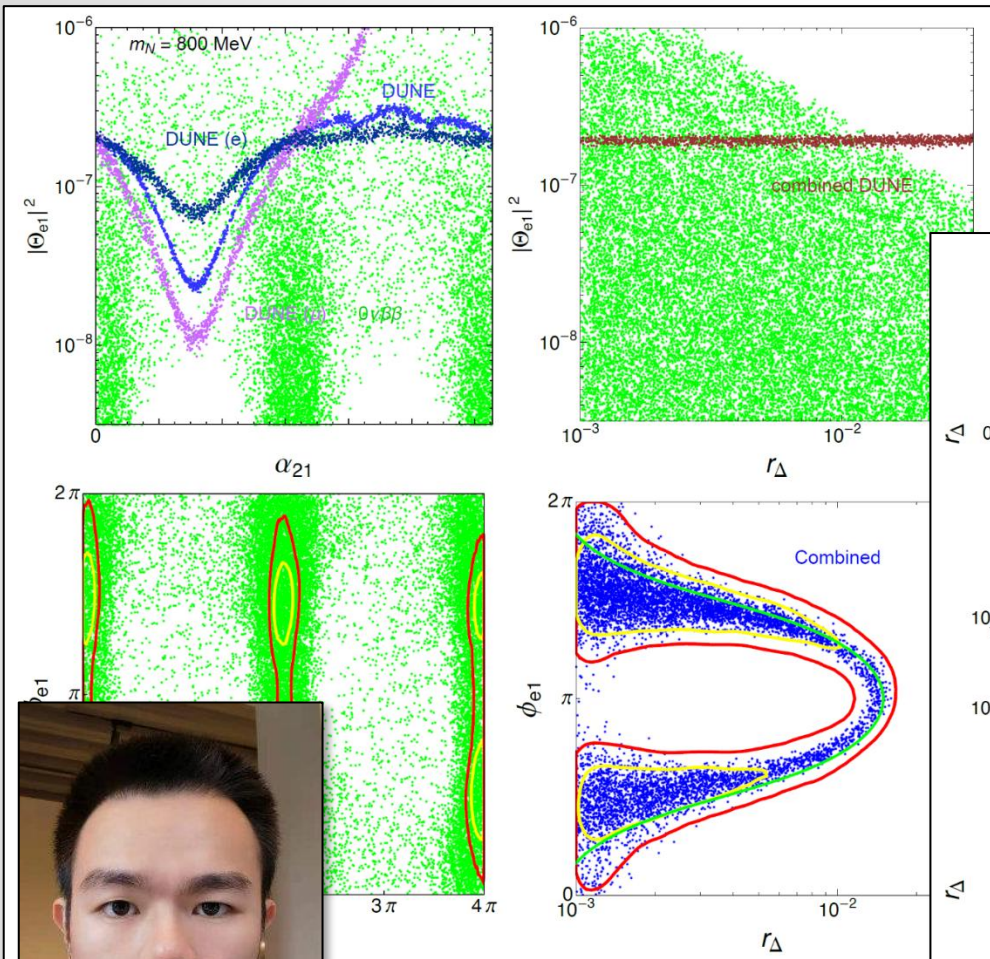


# Complementarity with Direct Searches



# Complementarity with Direct Searches

Patrick D. Bolton, FFD, Mudit Rai  
Zhong Zhang, 2212.14690



See Zhong Zhang's talk on Monday



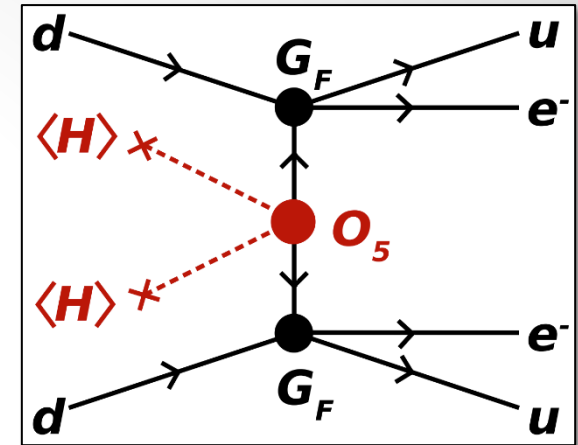
# Conclusion

▶  **$0\nu\beta\beta$  is crucial probe for BSM physics**

- Effective  $0\nu\beta\beta$  mass  $m_{\beta\beta} \approx 10 - 100$  meV probes New Physics near GUT scale
- Sensitive to Majorana HNL for  $0.1 \text{ eV} < m_N < 10 \text{ TeV}$
- Sensitive to nature of HNL: Dirac vs Quasi-Dirac vs Majorana

▶ **HNLs aka sterile neutrinos are natural extension to SM**

- Sensitivity in  $0\nu\beta\beta$  and direct searches such as DUNE
- Complementarity in probing properties of HNLs
- Precise prediction of HNL contribution and interference with  $m_{\beta\beta}$
- Strong enhancement for “light” Seesaw with  $m_N \ll 100 \text{ MeV}$
- Lifted  $0\nu\beta\beta$  “floor”?



$$\frac{T_{1/2}^{0\nu\beta\beta}}{10^{28} \text{ y}} \approx \left( \frac{\Lambda_{\text{NP}}}{10^{15} \text{ GeV}} \right)^2$$

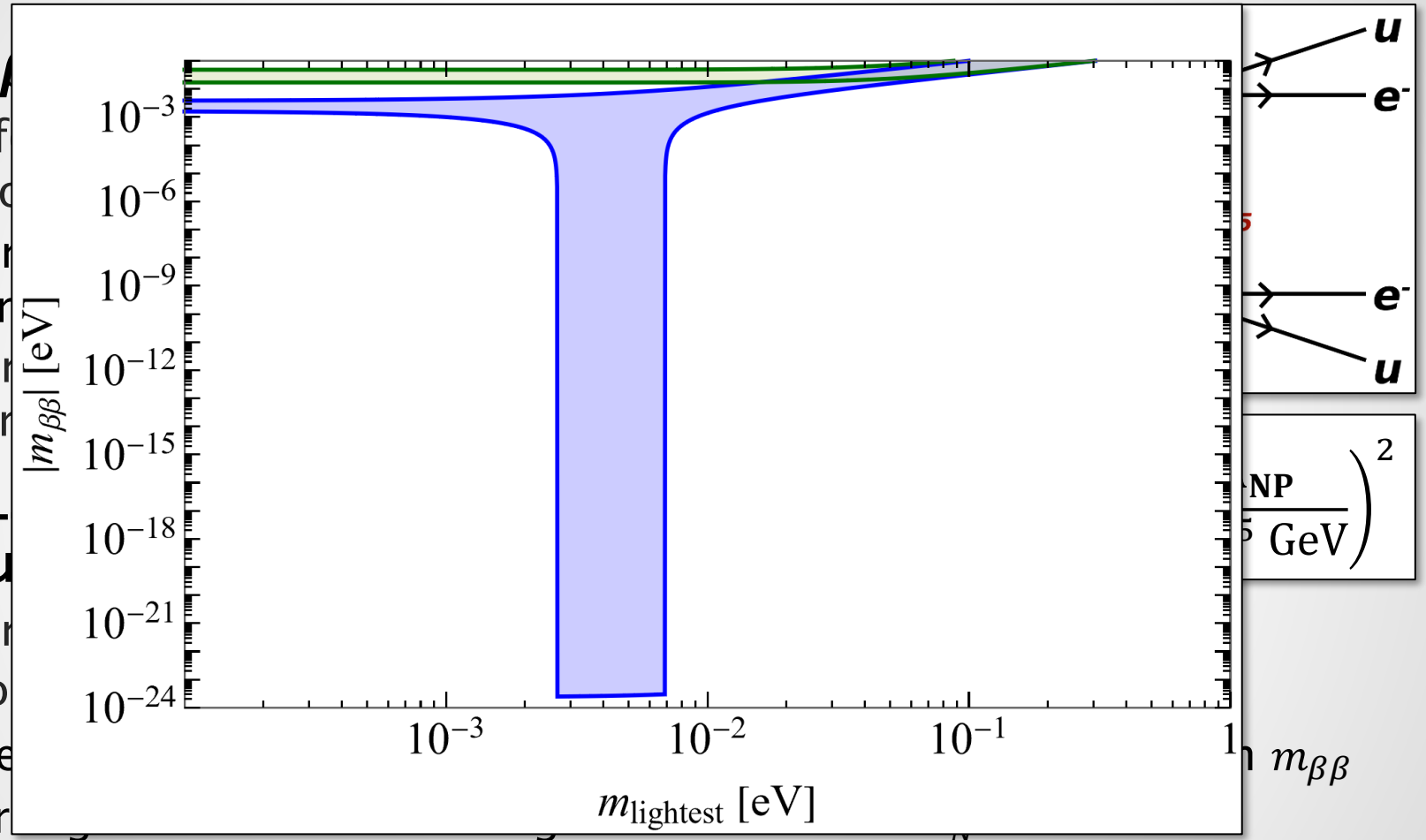
# Conclusion

▶  $0\nu\beta\beta$

- Eff
- pro
- Ser
- for
- Ser
- Dir

▶ HNL  
natu

- Ser
- Co
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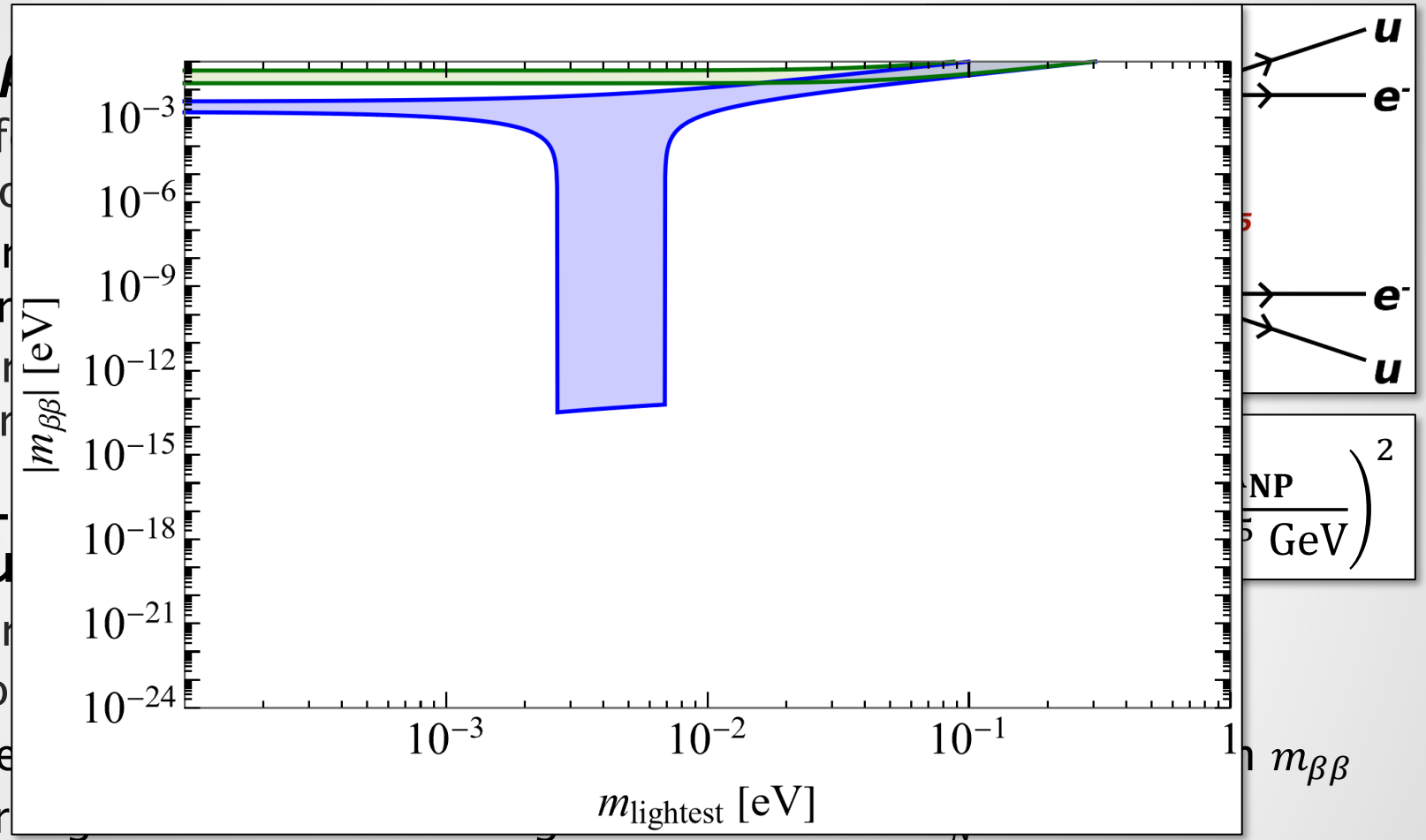
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# Conclusion

“Searching for sterile neutrinos at Christmas”  
according to Bing AI  
Image Generator

