

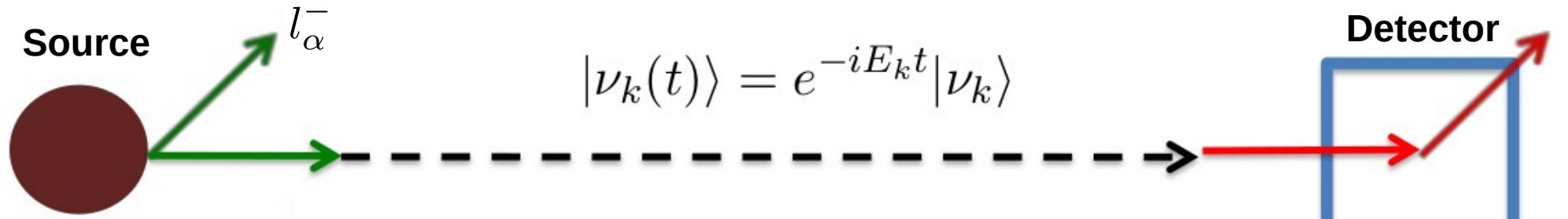
# Development in reactor neutrino models



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NuPhys 2023, London  
December 19<sup>th</sup> 2023

# Neutrino oscillations



$$|\nu_\alpha\rangle = \sum_k U_{\alpha k}^* |\nu_k\rangle$$

$$|\nu_k(t)\rangle = e^{-iE_k t} |\nu_k\rangle$$

$$\langle \nu_\beta | \nu_\alpha(t) \rangle = \sum_k U_{\alpha k}^* U_{\beta k} e^{-iE_k t}$$

$$P_{\nu_\alpha \rightarrow \nu_\beta}(t) = |A_{\nu_\alpha \rightarrow \nu_\beta}(t)|^2 = \sum_{k,j} U_{\alpha k}^* U_{\beta k} U_{\alpha j} U_{\beta j}^* e^{-i(E_k - E_j)t}$$

# Three-neutrino oscillations

Neutrino mixing matrix

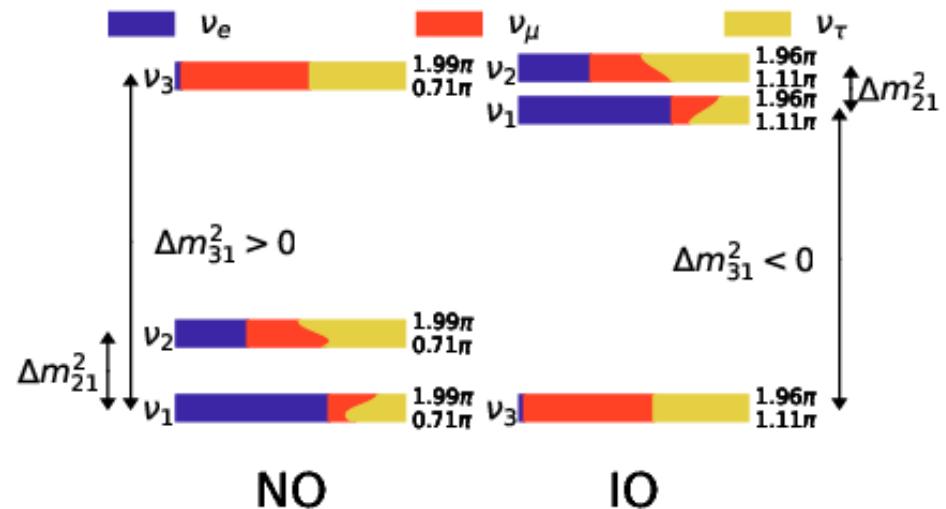
$$U = \begin{pmatrix} 1 & 0 & 0 \\ 0 & c_{23} & s_{23} \\ 0 & s_{23} & c_{23} \end{pmatrix} \begin{pmatrix} c_{13} & 0 & s_{13}e^{-i\delta} \\ 0 & 1 & 0 \\ -s_{13}e^{i\delta} & 0 & c_{13} \end{pmatrix} \begin{pmatrix} c_{12} & s_{12} & 0 \\ s_{12} & c_{12} & 0 \\ 0 & 0 & 1 \end{pmatrix}$$

Three mixing angles  $\theta_{12}, \theta_{13}, \theta_{23}$

1 Dirac + 2 Majorana CP-phases

Three masses  $m_1, m_2, m_3$  for which two orderings are possible

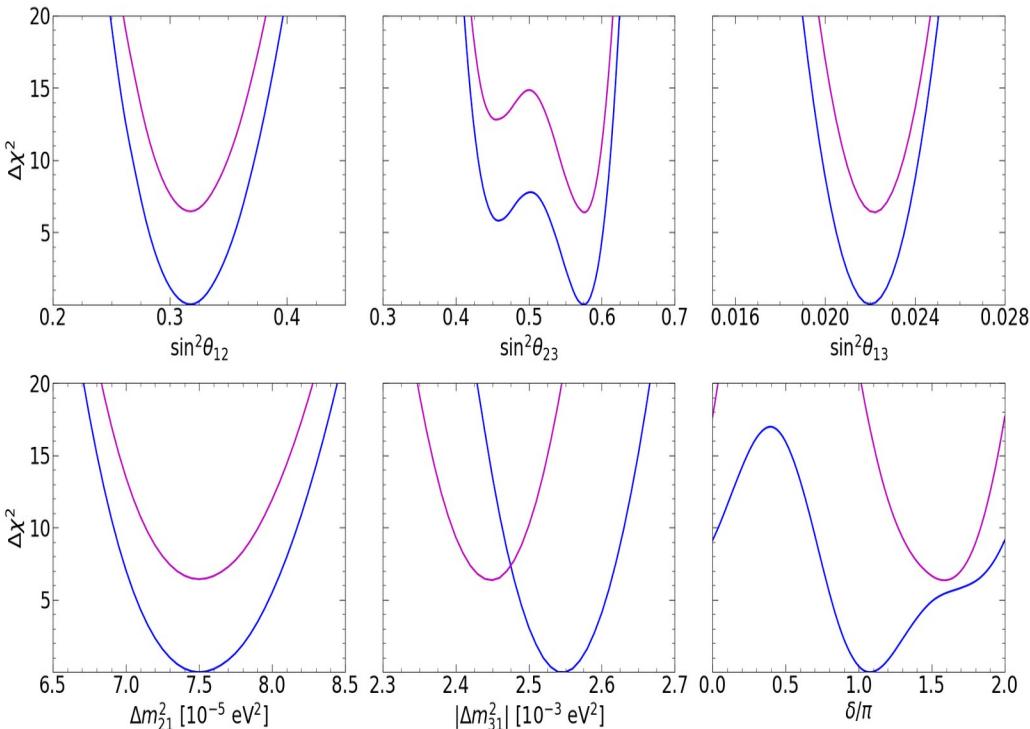
Oscillations are only sensitive to mass splittings



# Three-neutrino oscillations

Valencia - Global Fit, 2006.11237, JHEP 2021

Talk by Mariam Tortola!



parameter	best fit $\pm 1\sigma$	$2\sigma$ range	$3\sigma$ range
$\Delta m^2_{21} [10^{-5} \text{ eV}^2]$	$7.50^{+0.22}_{-0.20}$	7.12–7.93	6.94–8.14
$ \Delta m^2_{31}  [10^{-3} \text{ eV}^2]$ (NO)	$2.55^{+0.02}_{-0.03}$	2.49–2.60	2.47–2.63
$ \Delta m^2_{31}  [10^{-3} \text{ eV}^2]$ (IO)	$2.45^{+0.02}_{-0.03}$	2.39–2.50	2.37–2.53
$\sin^2 \theta_{12}/10^{-1}$	$3.18 \pm 0.16$	2.86–3.52	2.71–3.69
$\sin^2 \theta_{23}/10^{-1}$ (NO)	$5.74 \pm 0.14$	5.41–5.99	4.34–6.10
$\sin^2 \theta_{23}/10^{-1}$ (IO)	$5.78^{+0.10}_{-0.17}$	5.41–5.98	4.33–6.08
$\sin^2 \theta_{13}/10^{-2}$ (NO)	$2.200^{+0.069}_{-0.062}$	2.069–2.337	2.000–2.405
$\sin^2 \theta_{13}/10^{-2}$ (IO)	$2.225^{+0.064}_{-0.070}$	2.086–2.356	2.018–2.424
$\delta/\pi$ (NO)	$1.08^{+0.13}_{-0.12}$	0.84–1.42	0.71–1.99
$\delta/\pi$ (IO)	$1.58^{+0.15}_{-0.16}$	1.26–1.85	1.11–1.96

See also:

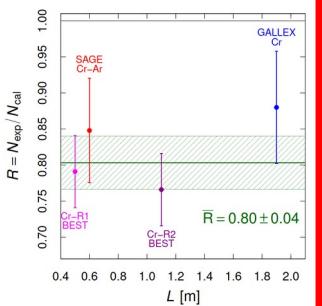
Bari – 2107.00532, PRD 2021

See also:

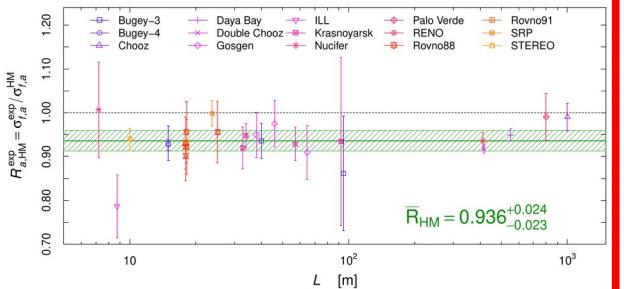
NuFit - 2111.03086 , Universe 2021

# Anomalies

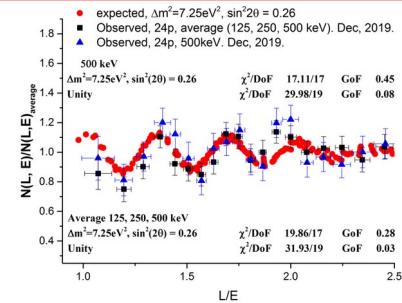
Gallium



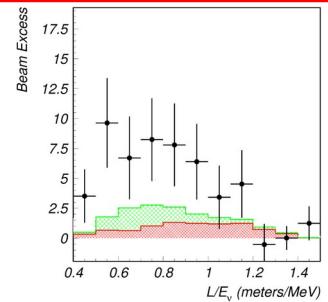
RAA



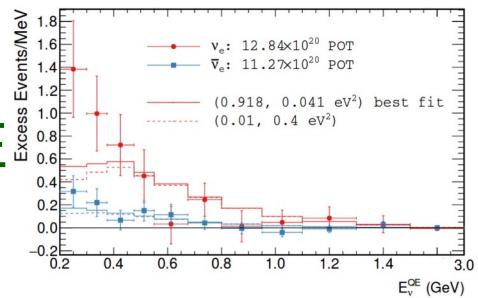
Neutrino-4



LSND



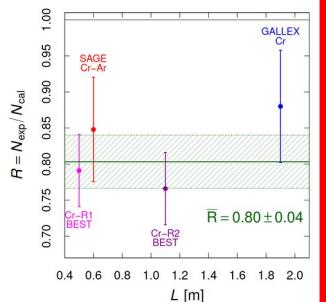
MiniBooNE



# Anomalies

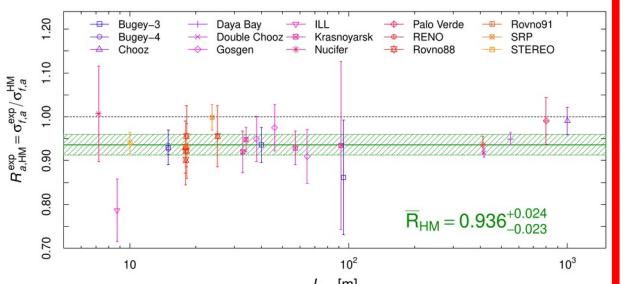
5-6 $\sigma$

Gallium



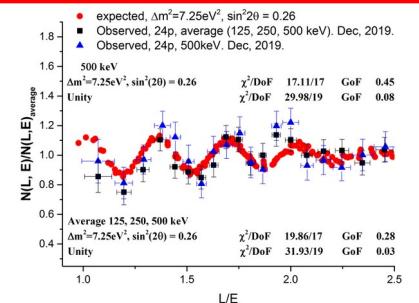
1-3 $\sigma$

RAA



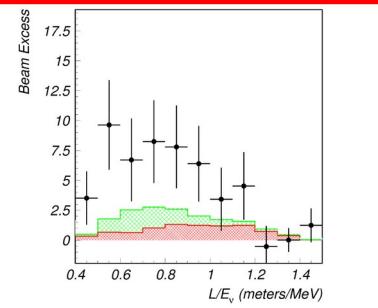
2-3 $\sigma$

Neutrino-4



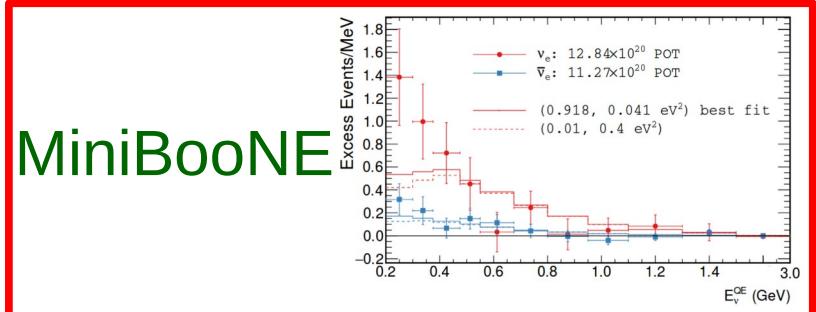
$\sim 4\sigma$

LSND



$\sim 5\sigma$

MiniBooNE



# Anomalies

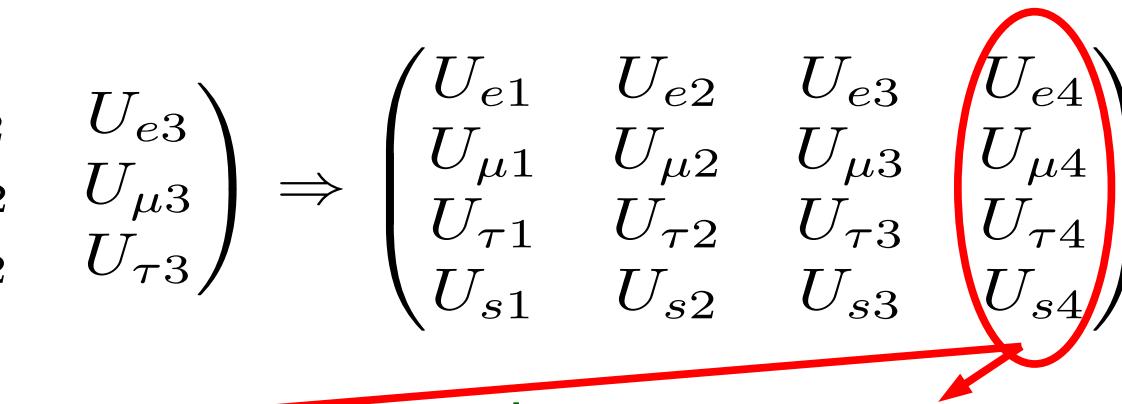
Three-neutrino oscillations can not account for short baseline anomalies

$$L_{kj}^{\text{osc}} = \frac{4\pi E}{\Delta m_{kj}^2}$$
$$L_{21}^{\text{osc}} \gtrsim 50 \text{ km} \frac{E}{\text{MeV}}$$
$$L_{31}^{\text{osc}} \gtrsim 1 \text{ km} \frac{E}{\text{MeV}}$$

Short baseline oscillations require:

$$\frac{L}{E} \lesssim 10 \text{ m/MeV} \implies \Delta m^2 \gtrsim 0.1 \text{ eV}^2$$

# 3+1 neutrino oscillations

$$\begin{pmatrix} U_{e1} & U_{e2} & U_{e3} \\ U_{\mu 1} & U_{\mu 2} & U_{\mu 3} \\ U_{\tau 1} & U_{\tau 2} & U_{\tau 3} \end{pmatrix} \Rightarrow \begin{pmatrix} U_{e1} & U_{e2} & U_{e3} & U_{e4} \\ U_{\mu 1} & U_{\mu 2} & U_{\mu 3} & U_{\mu 4} \\ U_{\tau 1} & U_{\tau 2} & U_{\tau 3} & U_{\tau 4} \\ U_{s1} & U_{s2} & U_{s3} & U_{s4} \end{pmatrix}$$


Appearance

$$P_{\alpha\beta}^{\text{SBL}} \approx \sin^2(2\theta_{\alpha\beta}) \sin^2 \left( \frac{\Delta m_{41}^2 L}{4E} \right)$$

$$\sin^2(2\theta_{\alpha\beta}) = 4|U_{\alpha 4}|^2 |U_{\beta 4}|^2$$

Disappearance

$$P_{\alpha\alpha}^{\text{SBL}} \approx 1 - \sin^2(2\theta_{\alpha\alpha}) \sin^2 \left( \frac{\Delta m_{41}^2 L}{4E} \right)$$

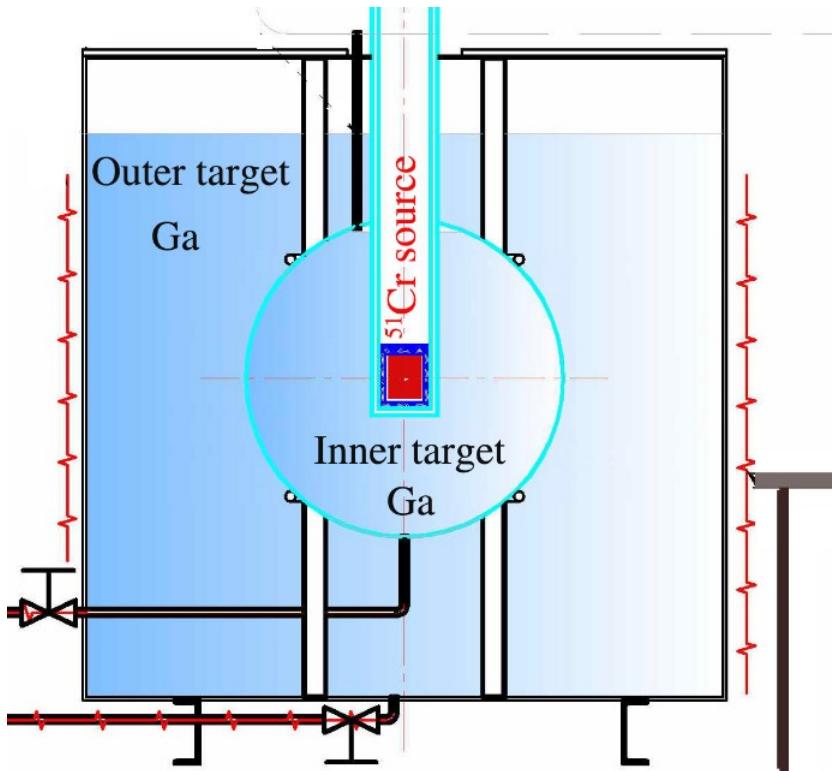
$$\sin^2(2\theta_{\alpha\alpha}) = 4|U_{\alpha 4}|^2 (1 - |U_{\alpha 4}|^2)$$

@LSND, Karmen, MiniBooNE,  
Opera

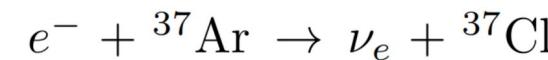
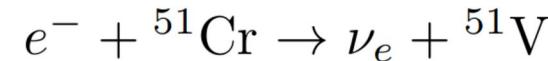
@Reactors and Gallium  
@atmospherics and accelerators

# The Gallium anomaly

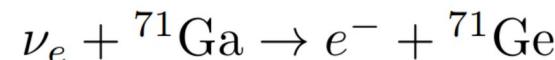
BEST, 2109.11482, PRL 2022



Intense sources of electron neutrinos are placed into the detector volume



The neutrino interact with the detector material



# The Gallium anomaly

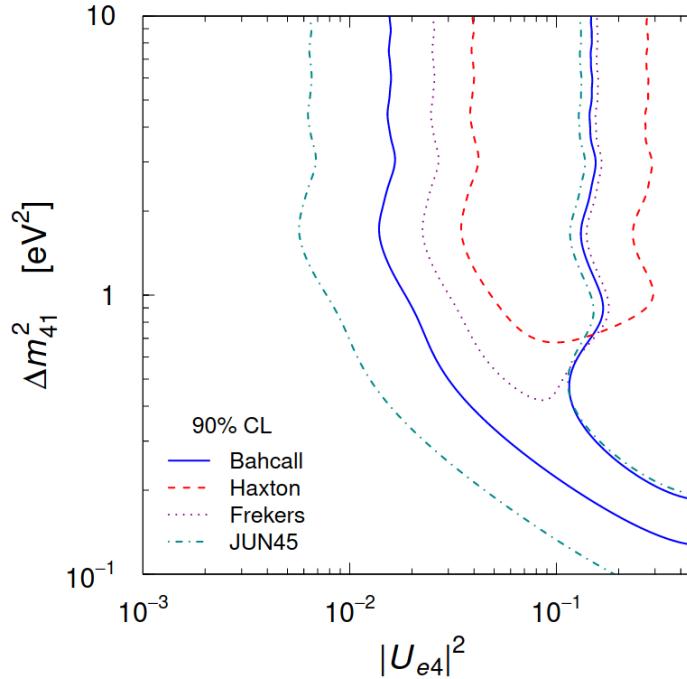
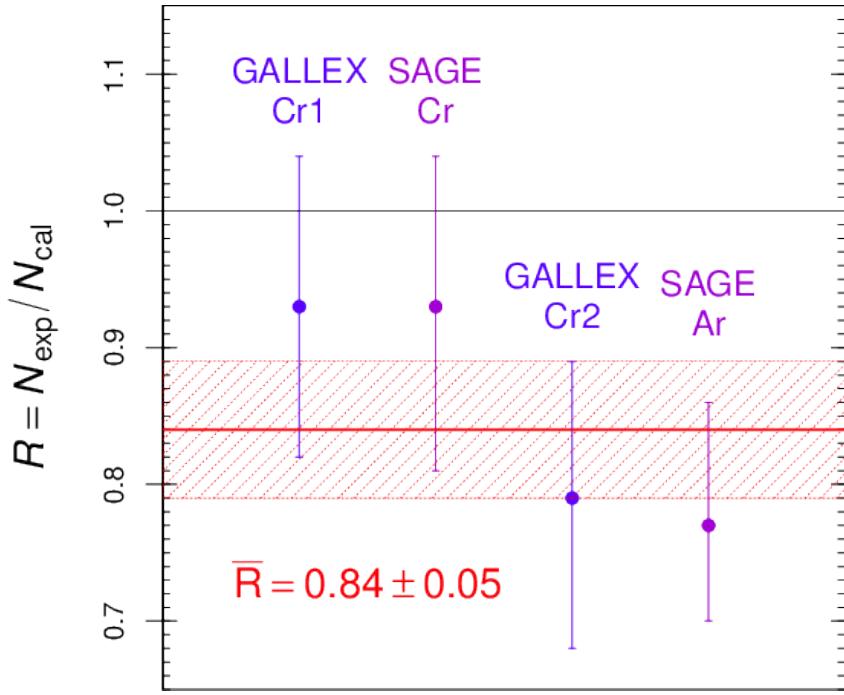
Model	Method	$^{51}\text{Cr}$		$^{37}\text{Ar}$	
		$\sigma_{\text{tot}}$	$\delta_{\text{exc}}$	$\sigma_{\text{tot}}$	$\delta_{\text{exc}}$
Ground State	$T_{1/2}(^{71}\text{Ge})$	$5.539 \pm 0.019$	–	$6.625 \pm 0.023$	–
Bahcall (1997)	$^{71}\text{Ga}(p, n)^{71}\text{Ge}$	$5.81 \pm 0.16$	4.7%	$7.00 \pm 0.21$	5.4%
Haxton (1998)	Shell Model	$6.39 \pm 0.65$	13.3%	$7.72 \pm 0.81$	14.2%
Frekers et al. (2015)	$^{71}\text{Ga}({}^3\text{He}, {}^3\text{H})^{71}\text{Ge}$	$5.92 \pm 0.11$	6.4%	$7.15 \pm 0.14$	7.3%
Kostensalo et al. (2019)	Shell Model	$5.67 \pm 0.06$	2.3%	$6.80 \pm 0.08$	2.6%
Semenov (2020)	$^{71}\text{Ga}({}^3\text{He}, {}^3\text{H})^{71}\text{Ge}$	$5.938 \pm 0.116$	6.7%	$7.169 \pm 0.147$	7.6%

Giunti, Li, Ternes, Tyagi, Zhao, 2209.00916, JHEP 2022

Slightly different values for the different cross section models

# The Gallium anomaly

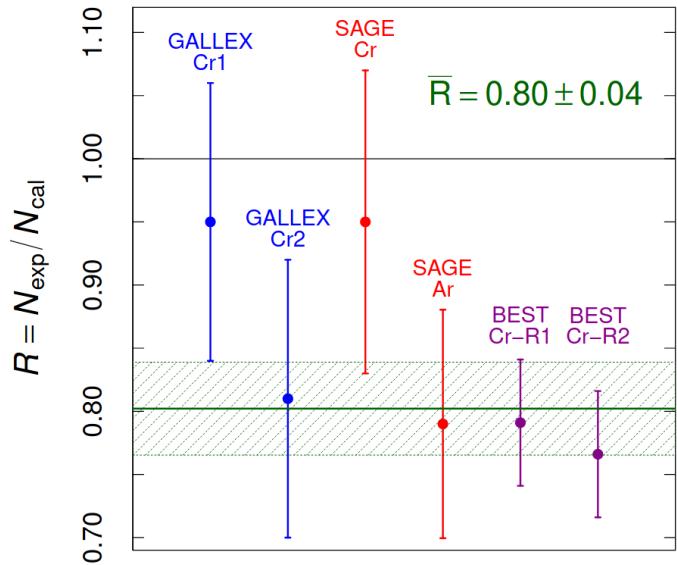
Kostensalo, Suhonen, Giunti, Srivastava, 1906.10980, PLB 2019



The significance of the “old” Gallium anomaly varied between 2.3 and 3.0 $\sigma$ , depending on the cross section model

# The Gallium anomaly

Giunti, Li, Ternes, Tyagi, Zhao,  
2209.00916, JHEP 2022

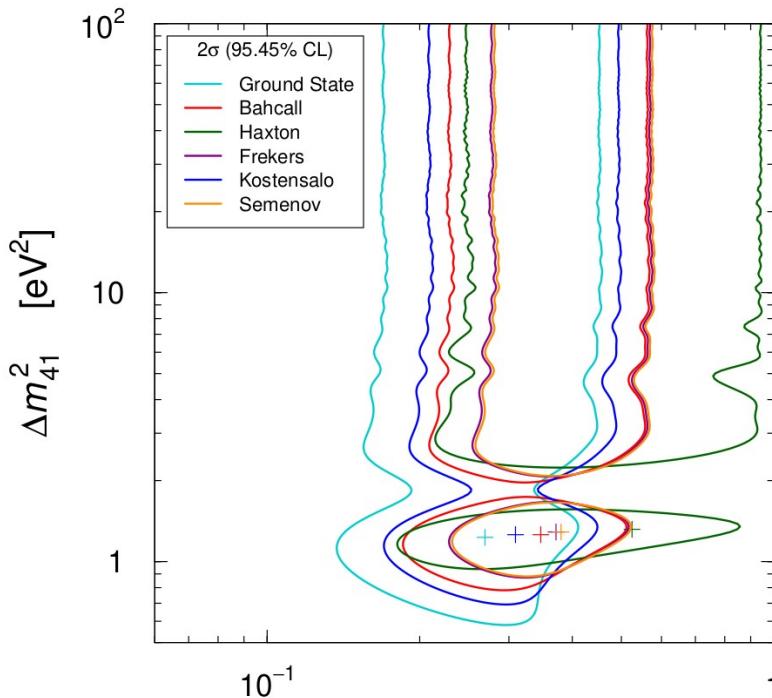


Strong indication for short baseline (SBL) oscillations?

See also:

Barinov, Gorbunov, 2109.14654, PRD2022

Berryman, Coloma, Huber, Schwetz, Zhao, 2111.12530, JHEP 2022



Model	Method	$\bar{R}$	GA
Ground State	$T_{1/2}(^{71}\text{Ge})$	$0.845^{+0.031}_{-0.031}$	5.0
Bahcall (1997)	$^{71}\text{Ga}(p, n)^{71}\text{Ge}$	$0.804^{+0.037}_{-0.036}$	5.2
Haxton (1998)	Shell Model	$0.731^{+0.088}_{-0.072}$	5.1
Frekers et al. (2015)	$^{71}\text{Ga}(^{3}\text{He}, ^{3}\text{H})^{71}\text{Ge}$	$0.789^{+0.033}_{-0.032}$	6.1
Kostensalo et al. (2019)	Shell Model	$0.825^{+0.031}_{-0.031}$	5.5
Semenov (2020)	$^{71}\text{Ga}(^{3}\text{He}, ^{3}\text{H})^{71}\text{Ge}$	$0.787^{+0.033}_{-0.032}$	6.1

# The reactor antineutrino anomaly: Flux calculations

The neutrino spectrum is produced from the beta decays of the fission products of  $^{235}\text{U}$ ,  $^{238}\text{U}$ ,  $^{239}\text{Pu}$ , and  $^{241}\text{Pu}$

## Summation method

$$S_{\text{tot}}(E, t) = \sum_k F_k(t) S_k(E)$$

fission fractions

$$S_k(E) = \sum_n Y_n^k \quad \sum_b BR_n^b \quad S_n^b(E) \leftarrow$$

cumulative  
fission  
yield

branching  
ratio

allowed or  
forbidden  
decay  
spectrum

There are more than 1000 beta spectra and branching ratios

Nuclear data bases might be incomplete or inaccurate

## Conversion method

Measure beta spectra of  $^{235}\text{U}$ ,  $^{238}\text{U}$ ,  $^{239}\text{Pu}$ ,  $^{241}\text{Pu}$  and use empirical method to get

$$S_k^e(E_e) \rightarrow S_k^\nu(E_\nu)$$

# 5 MeV bump

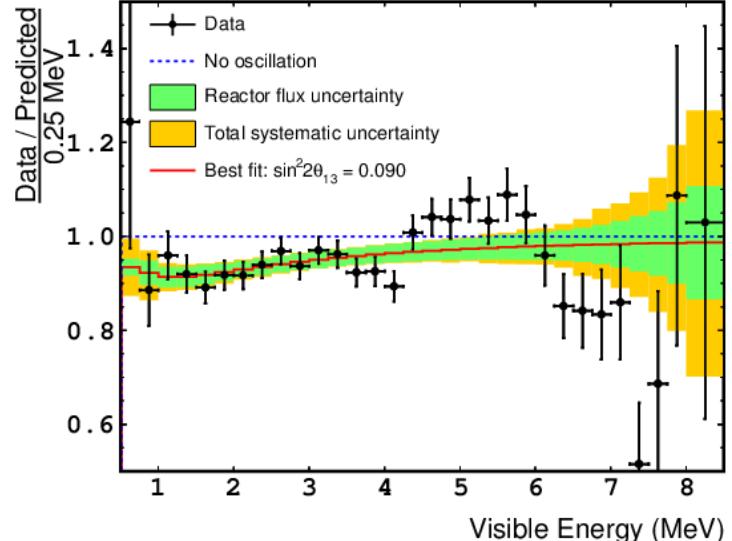
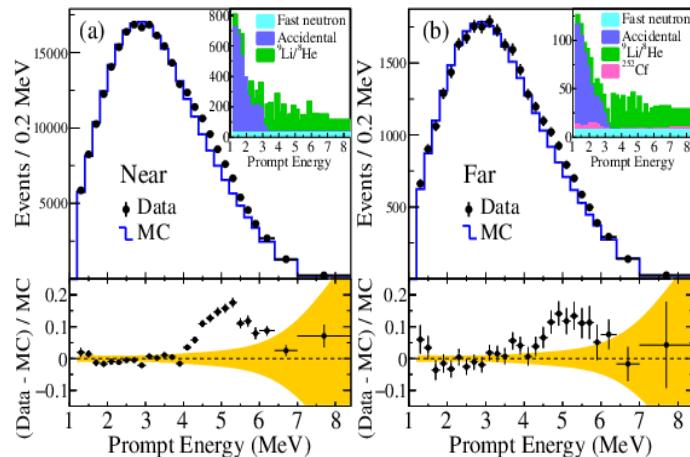
Double Chooz, 1406.7763, JHEP 2015

5 MeV bump discovered in 2014

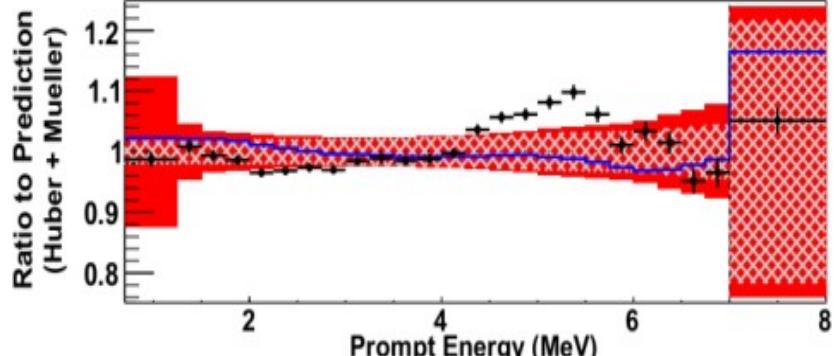
Can not been explained with SBL oscillations

Proof of our incomplete understanding of nuclear reactor fluxes

RENO, 1511.05849, PRL 2016



Daya Bay, 1508.04233, PRL 2016



# Rate calculation

Calculate inverse beta yields for each isotope

We use the Strumia-Vissani IBD cross section

Strumia, Vissani, astro-ph/0302055, PLB 2003

$$\sigma_i = \int_{E_\nu^{\text{thr}}}^{E_\nu^{\text{max}}} dE_\nu \Phi_i(E_\nu) \sigma_{\text{IBD}}(E_\nu)$$

Yields depend on  
the neutrino flux

Giunti, Li, Ternes, Xin, 2110.06820, PLB 2022

Model	$\sigma_{235}$	$\sigma_{238}$	$\sigma_{239}$	$\sigma_{241}$
HM	$6.74 \pm 0.17$	$10.19 \pm 0.83$	$4.40 \pm 0.13$	$6.10 \pm 0.16$
EF	$6.29 \pm 0.31$	$10.16 \pm 1.02$	$4.42 \pm 0.22$	$6.23 \pm 0.31$
HKSS	$6.82 \pm 0.18$	$10.28 \pm 0.84$	$4.45 \pm 0.13$	$6.17 \pm 0.16$
KI	$6.41 \pm 0.14$	$9.53 \pm 0.48$	$4.40 \pm 0.13$	$6.10 \pm 0.16$

$$\sigma_{f,a} = \sum_i f_i^a \sigma_i$$

Berryman, Huber, 2005.01756, JHEP 2021

Model	$\sigma_{235}$	$\sigma_{238}$	$\sigma_{239}$	$\sigma_{241}$
HM	$6.60 \pm 0.14$	$10.00 \pm 1.12$	$4.33 \pm 0.11$	$6.01 \pm 0.13$
EF	$6.17 \pm 0.13$	$9.94 \pm 1.09$	$4.32 \pm 0.11$	$6.10 \pm 0.13$
HKSS	$6.67 \pm 0.15$	$10.08 \pm 1.14$	$4.37 \pm 0.12$	$6.06 \pm 0.14$

# Compare against measurements

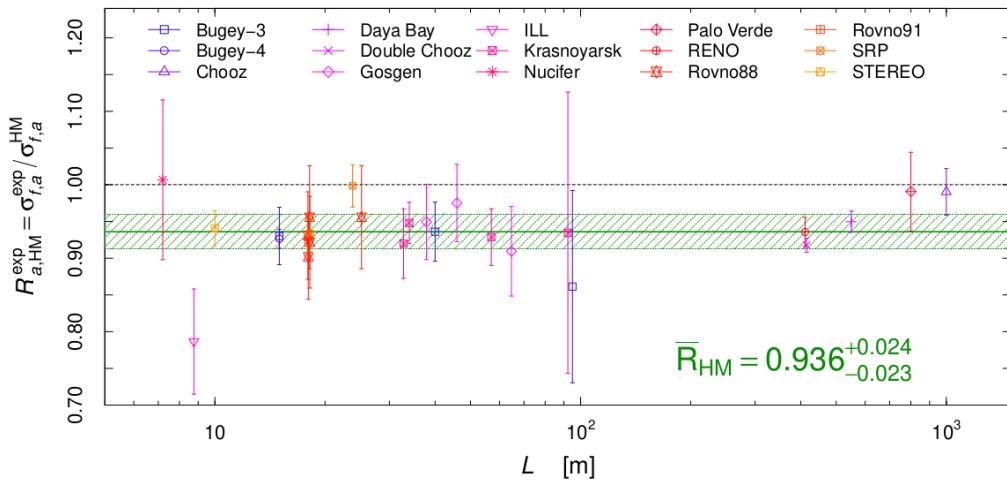
Giunti, Li, Ternes, Xin, 2110.06820

$a$	Experiment	$f_{235}^a$	$f_{238}^a$	$f_{239}^a$	$f_{241}^a$	$\sigma_{f,a}^{\text{exp}}$	$R_{a,\text{HM}}^{\text{exp}}$	$R_{a,\text{EF}}^{\text{exp}}$	$R_{a,\text{HKSS}}^{\text{exp}}$	$R_{a,\text{KI}}^{\text{exp}}$	$R_{a,\text{HKSS-KI}}^{\text{exp}}$	$\delta_a^{\text{exp}} [\%]$	$\delta_a^{\text{cor}} [\%]$	$L_a [\text{m}]$
1	Bugey-4	0.538	0.078	0.328	0.056	5.75	0.927	0.962	0.916	0.962	0.944	1.4	1.4	15
2	Rovno91	0.614	0.074	0.274	0.038	5.85	0.924	0.965	0.914	0.962	0.945	2.8		18
3	Rovno88-II	0.607	0.074	0.277	0.042	5.70	0.902	0.941	0.892	0.939	0.921	6.4	3.1	18
4	Rovno88-2I	0.603	0.076	0.276	0.045	5.89	0.931	0.971	0.920	0.969	0.951	6.4		17.96
5	Rovno88-1S	0.606	0.074	0.277	0.043	6.04	0.956	0.997	0.945	0.995	0.976	7.3	2.2	18.15
6	Rovno88-2S	0.557	0.076	0.313	0.054	5.96	0.956	0.994	0.945	0.993	0.974	7.3		25.17
7	Rovno88-3S	0.606	0.074	0.274	0.046	5.83	0.922	0.962	0.911	0.960	0.942	6.8	3.1	18.18
8	Bugey-3-15	0.538	0.078	0.328	0.056	5.77	0.930	0.966	0.920	0.966	0.947	4.2		15
9	Bugey-3-40	0.538	0.078	0.328	0.056	5.81	0.936	0.972	0.926	0.972	0.953	4.3	4.0	40
10	Bugey-3-95	0.538	0.078	0.328	0.056	5.35	0.861	0.895	0.852	0.894	0.877	15.2		95
11	Gosgen-38	0.619	0.067	0.272	0.042	5.99	0.949	0.992	0.939	0.988	0.971	5.4	2.0	37.9
12	Gosgen-46	0.584	0.068	0.298	0.050	6.09	0.975	1.016	0.964	1.014	0.995	5.4		45.9
13	Gosgen-65	0.543	0.070	0.329	0.058	5.62	0.909	0.945	0.899	0.944	0.927	6.7	3.8	64.7
14	ILL	1.000	0.000	0.000	0.000	5.30	0.787	0.843	0.777	0.827	0.818	9.1		8.76
15	Krasnoyarsk87-33	1	0	0	0	6.20	0.920	0.986	0.909	0.967	0.957	5.2	4.1	32.8
16	Krasnoyarsk87-92	1	0	0	0	6.30	0.935	1.002	0.924	0.983	0.972	20.5		92.3
17	Krasnoyarsk94-57	1	0	0	0	6.26	0.929	0.995	0.918	0.977	0.966	4.2	0	57
18	Krasnoyarsk99-34	1	0	0	0	6.39	0.948	1.016	0.937	0.997	0.986	3.0	0	34
19	SRP-18	1	0	0	0	6.29	0.934	1.000	0.923	0.982	0.971	2.8	0	18.2
20	SRP-24	1	0	0	0	6.73	0.998	1.070	0.987	1.050	1.038	2.9	0	23.8
21	Nucifer	0.926	0.008	0.061	0.005	6.67	1.007	1.074	0.995	1.056	1.044	10.8	0	7.2
22	Chooz	0.496	0.087	0.351	0.066	6.12	0.990	1.025	0.979	1.027	1.007	3.2	0	$\approx 1000$
23	Palo Verde	0.600	0.070	0.270	0.060	6.25	0.991	1.033	0.980	1.031	1.012	5.4	0	$\approx 800$
24	Daya Bay	0.564	0.076	0.304	0.056	5.94	0.950	0.988	0.939	0.987	0.968	1.5	0	$\approx 550$
25	RENO	0.571	0.073	0.300	0.056	5.85	0.936	0.974	0.925	0.973	0.954	2.1	0	$\approx 411$
26	Double Chooz	0.520	0.087	0.333	0.060	5.71	0.918	0.952	0.907	0.953	0.934	1.1	0	$\approx 415$
27	STEREO	1	0	0	0	6.34	0.941	1.008	0.930	0.989	0.978	2.5	0	9 – 11

# 2011 Huber Mueller fluxes

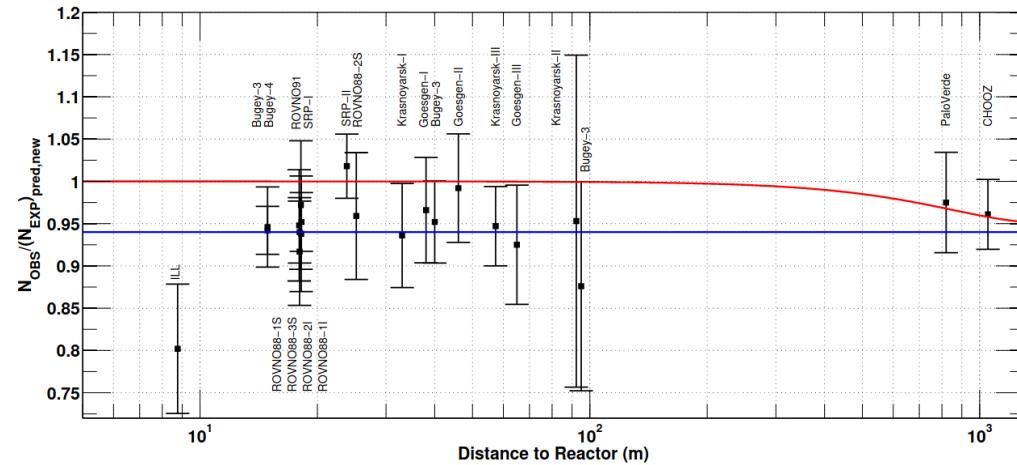
Huber, 1106.0687, PRC 2012

Mueller, Lhuillier, Fallot, Letourneau, Cormon, 1101.2663, PRC 2012



HM flux gives  $2.5\sigma$  anomaly

Giunti, Li, Ternes, Xin, 2110.06820, PLB 2022

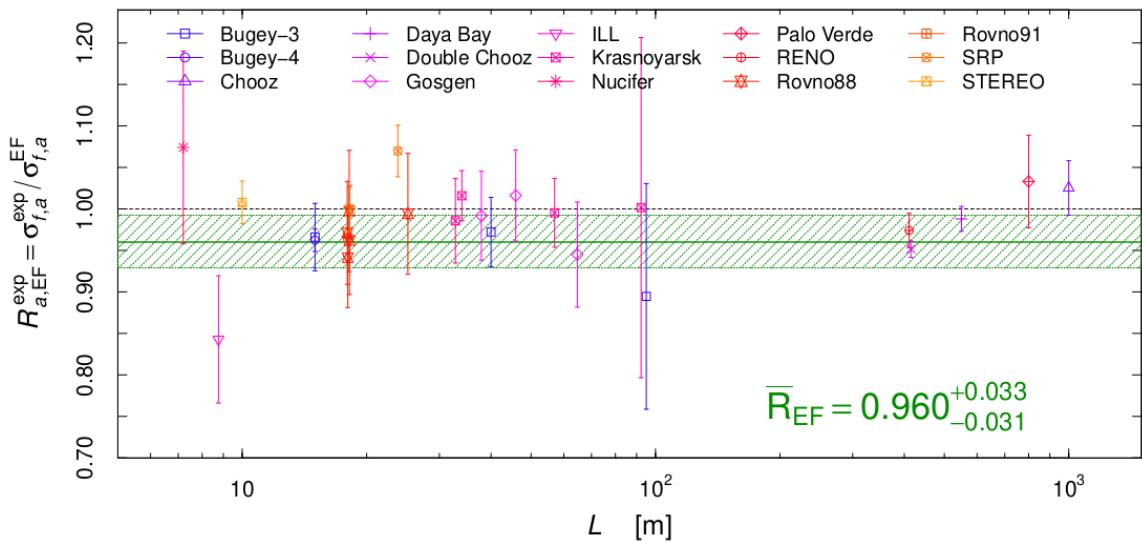


Original RAA was also  $2.5\sigma$

Mention, Fechner, Lasserre, Mueller, Lhuillier,  
1101.2755, PRD 2011

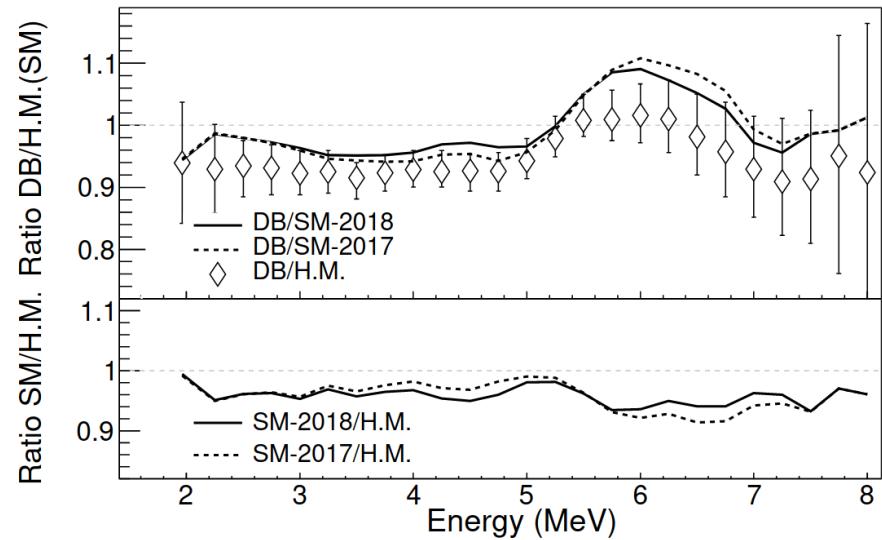
# 2019 summation method fluxes

Estienne, Fallot, et al, 1904.09358, PRL 2019



1.2 $\sigma$  deficit, no anomaly!

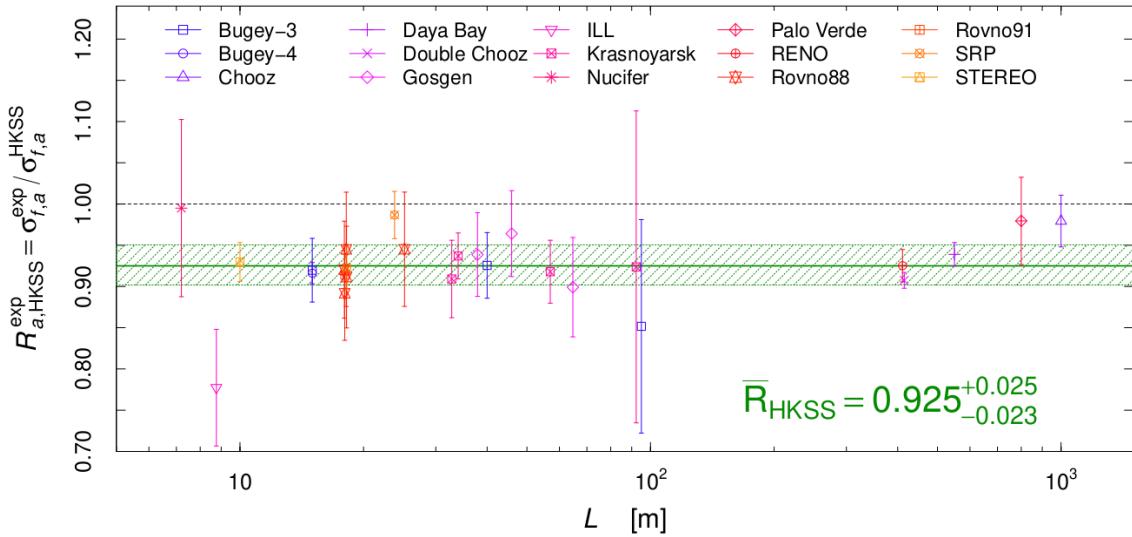
Giunti, Li, Ternes, Xin, 2110.06820, PLB 2022



Ratio reduced with respect  
to HM for all energies!

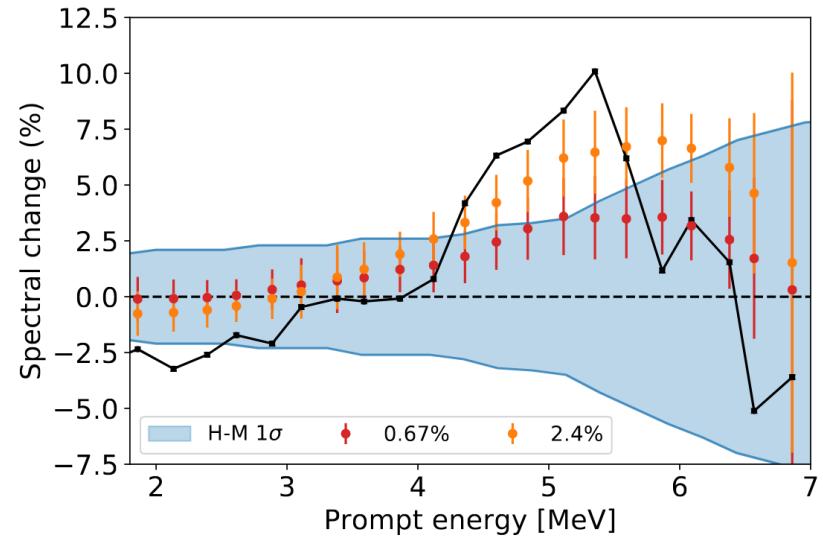
# 2019 new converted fluxes

Hayen, Kostensalo, Severijns, Suhonen, 1908.08302, PRC 2019



HKSS flux results in  $2.9\sigma$  anomaly!

Giunti, Li, Ternes, Xin, 2110.06820, PLB 2022

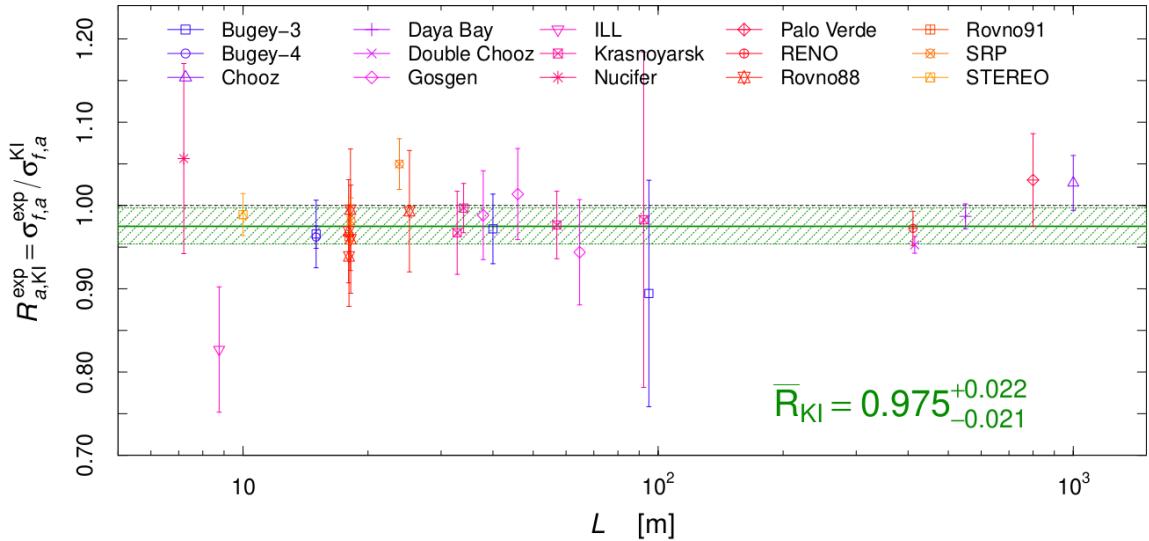


Better prediction for the energies of the bump!

# 2021 new converted fluxes

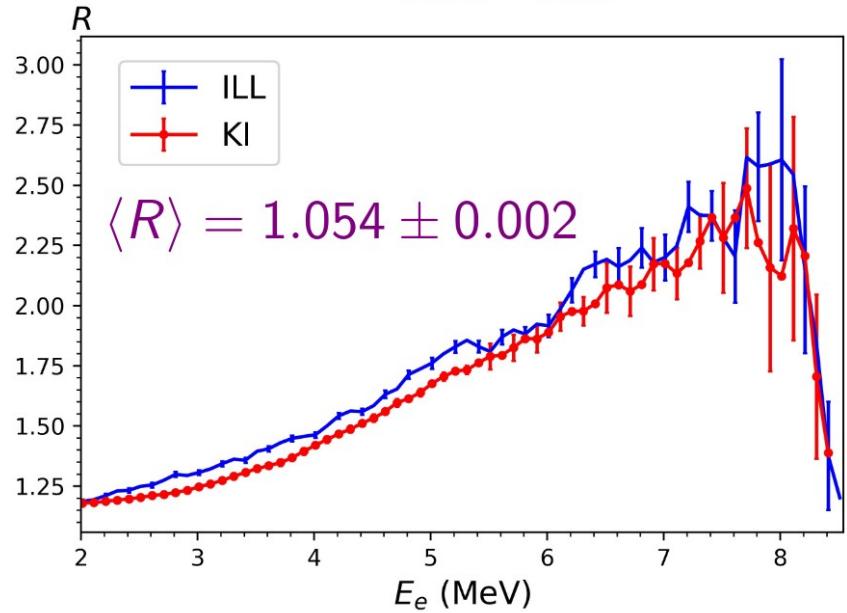
Kurchatov Institute: Kopeikin, Skorokhvatov, Titov, 2103.01684, PRD 2021

$$R = S_{235}^{(e)} / S_{239}^{(e)}$$



No anomaly ( $1.1\sigma$ ) with KI flux!

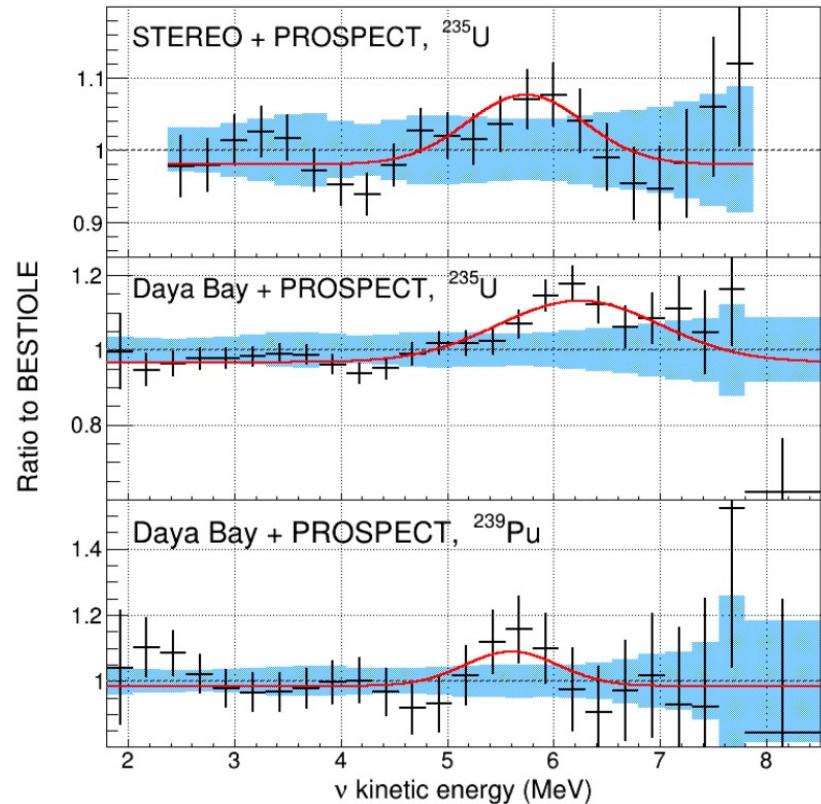
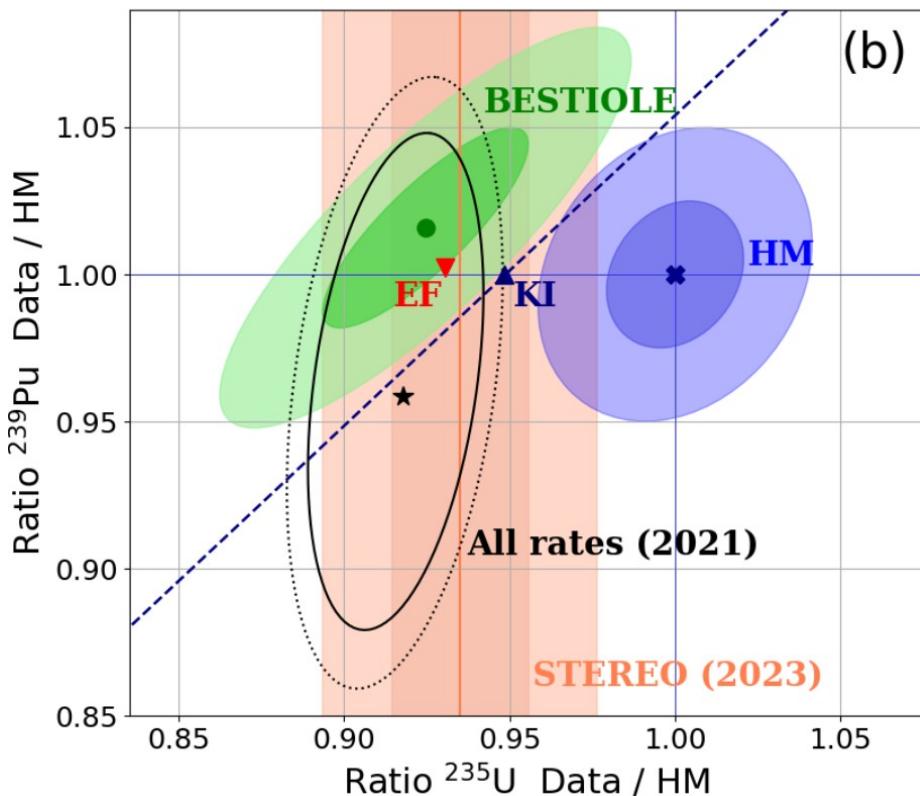
Giunti, Li, Ternes, Xin, 2110.06820, PLB 2022



New measurement  
suggests a reduction in  
the flux of  $^{235}\text{U}$

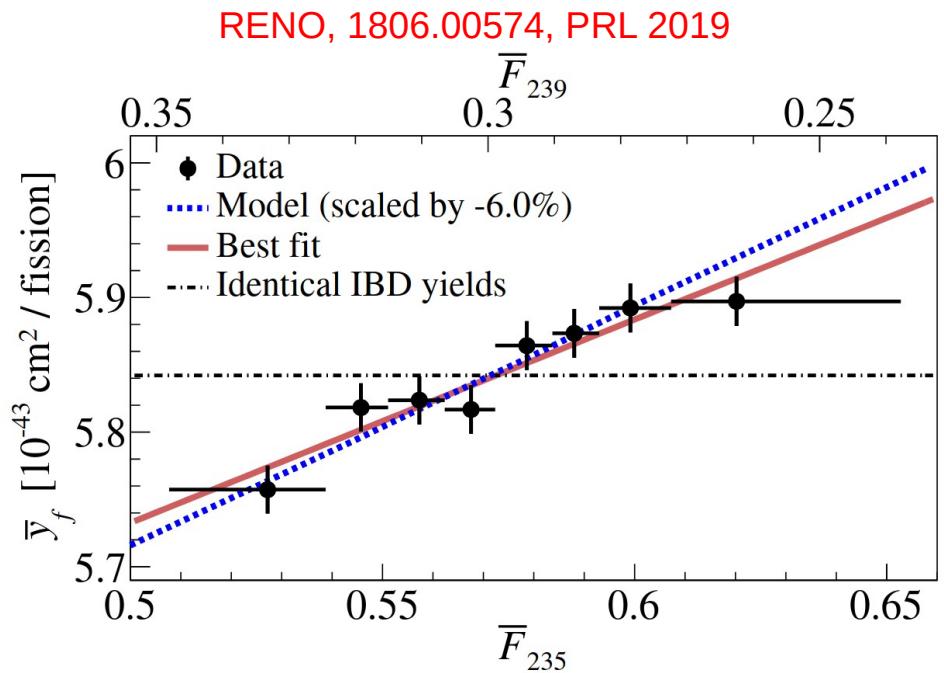
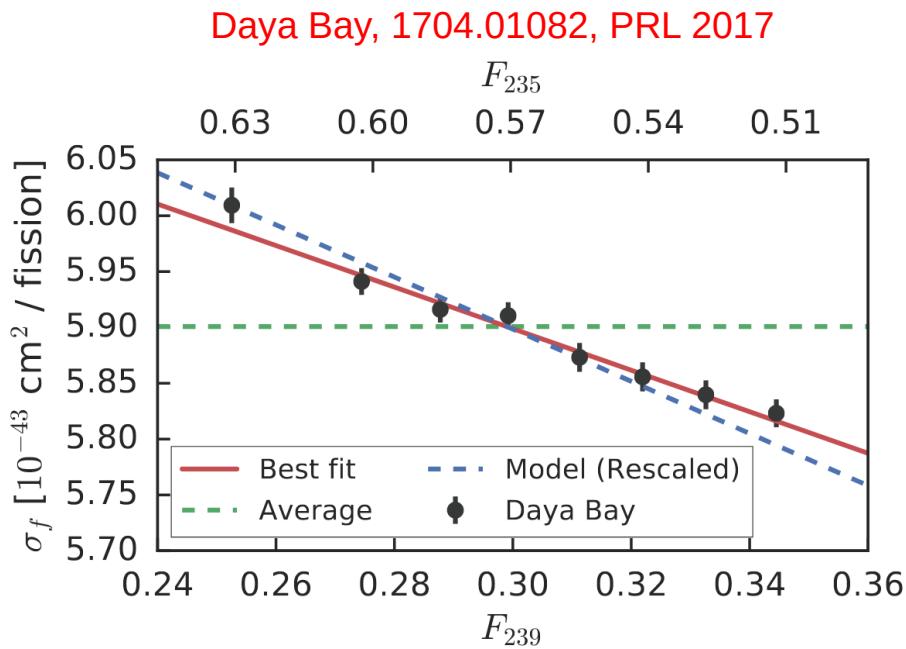
# BESTIOLE: 2023 new flux from summation method

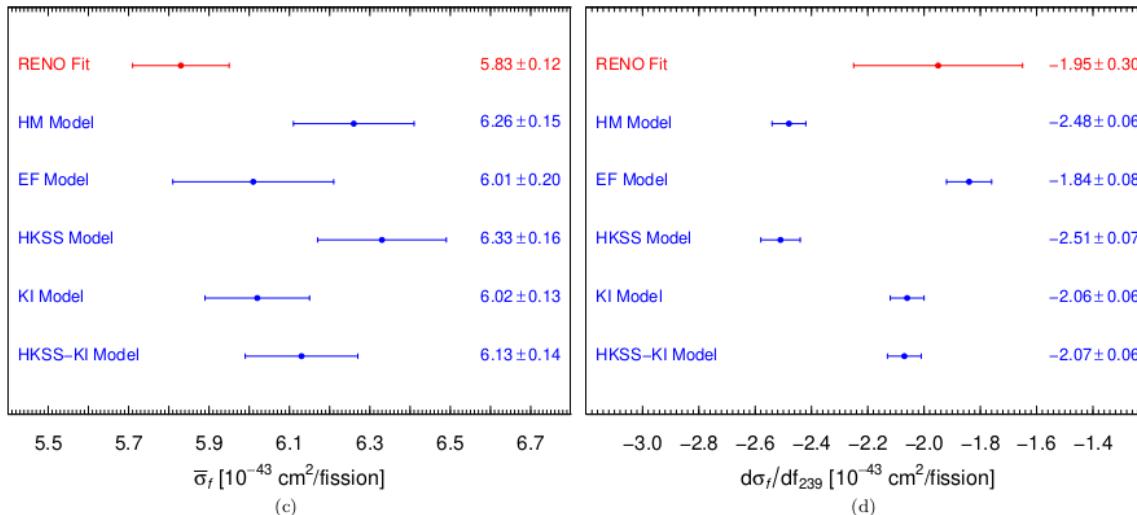
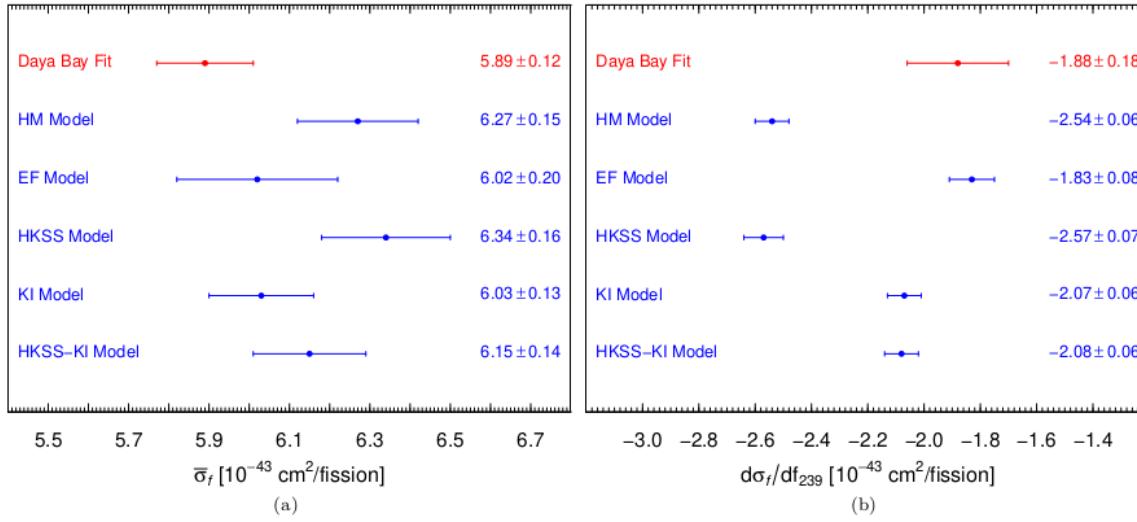
BESTIOLE: Périssé, el al, 2304.14992



# Evolution data

Measure rates at different stages of reactor cycle





Giunti, Li, Ternes, Xin, 2110.06820, PLB 2022

We get additional information from the measurement of the slope parameter

# Evolution data

## Effect of evolution data on RAA

Model	Rates		Evolution		Rates + Evolution	
	$\bar{R}_{\text{mod}}$	RAA	$\bar{R}_{\text{mod}}$	RAA	$\bar{R}_{\text{mod}}$	RAA
<b>HM</b>	$0.936^{+0.024}_{-0.023}$	$2.5\sigma$	$0.933^{+0.025}_{-0.024}$	$2.6\sigma$	$0.930^{+0.024}_{-0.023}$	$2.8\sigma$
<b>EF</b>	$0.960^{+0.033}_{-0.031}$	$1.2\sigma$	$0.975^{+0.032}_{-0.030}$	$0.8\sigma$	$0.975^{+0.032}_{-0.030}$	$0.8\sigma$
<b>HKSS</b>	$0.925^{+0.025}_{-0.023}$	$2.9\sigma$	$0.925^{+0.026}_{-0.024}$	$2.8\sigma$	$0.922^{+0.024}_{-0.023}$	$3.0\sigma$
<b>KI</b>	$0.975^{+0.022}_{-0.021}$	$1.1\sigma$	$0.973^{+0.023}_{-0.022}$	$1.2\sigma$	$0.970 \pm 0.021$	$1.4\sigma$
<b>HKSS-KI</b>	$0.964^{+0.023}_{-0.022}$	$1.5\sigma$	$0.955^{+0.024}_{-0.023}$	$1.9\sigma$	$0.960^{+0.022}_{-0.021}$	$1.8\sigma$

Giunti, Li, Ternes, Xin, 2110.06820, PLB 2022

# Best fit reactor flux model

We perform several statistical tests for the best fit flux model

We find that the recent KI model is the best among the conversion models

The EF model is equally good as the KI model

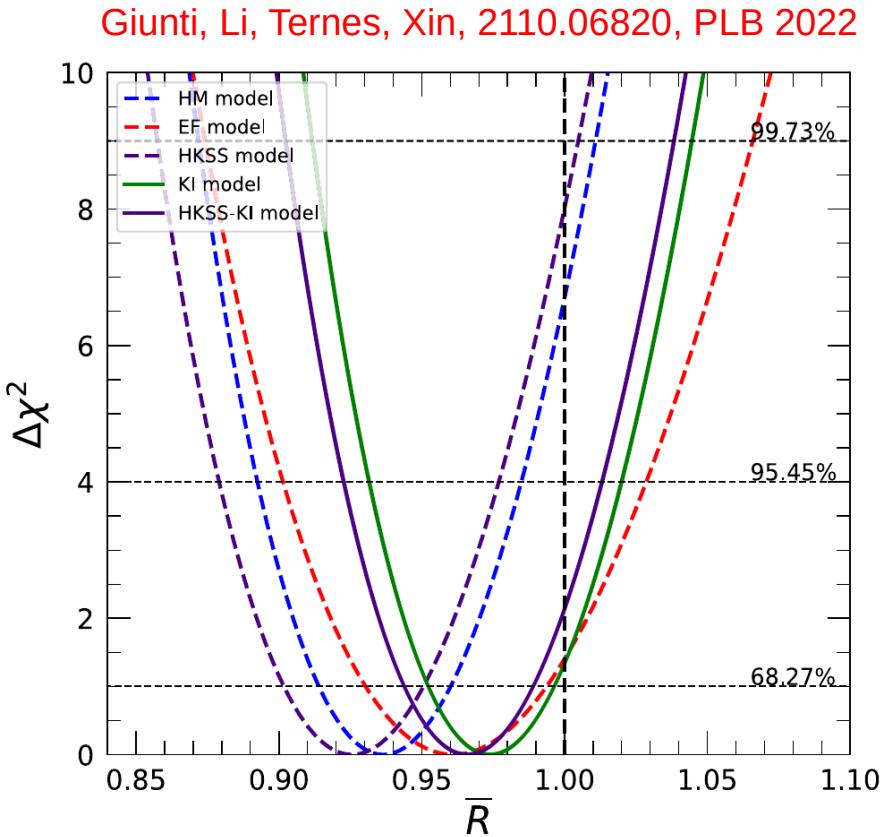
	Rates + Evolution				
$\chi^2$	0.13	0.22	0.08	0.68	0.44
SW	0.32	0.13	0.35	0.59	0.41
sign	0.03	0.38	0.006	0.38	0.11
KS	0.04	0.84	0.02	0.39	0.20
CVM	0.02	0.67	0.006	0.38	0.14
AD	0.02	0.57	0.006	0.40	0.13
$Z_K$	$< 10^{-3}$	0.05	$< 10^{-3}$	0.05	0.008
$Z_C$	0.02	0.11	0.005	0.55	0.15
$Z_A$	0.03	0.20	0.01	0.41	0.12

# The reactor rate anomaly

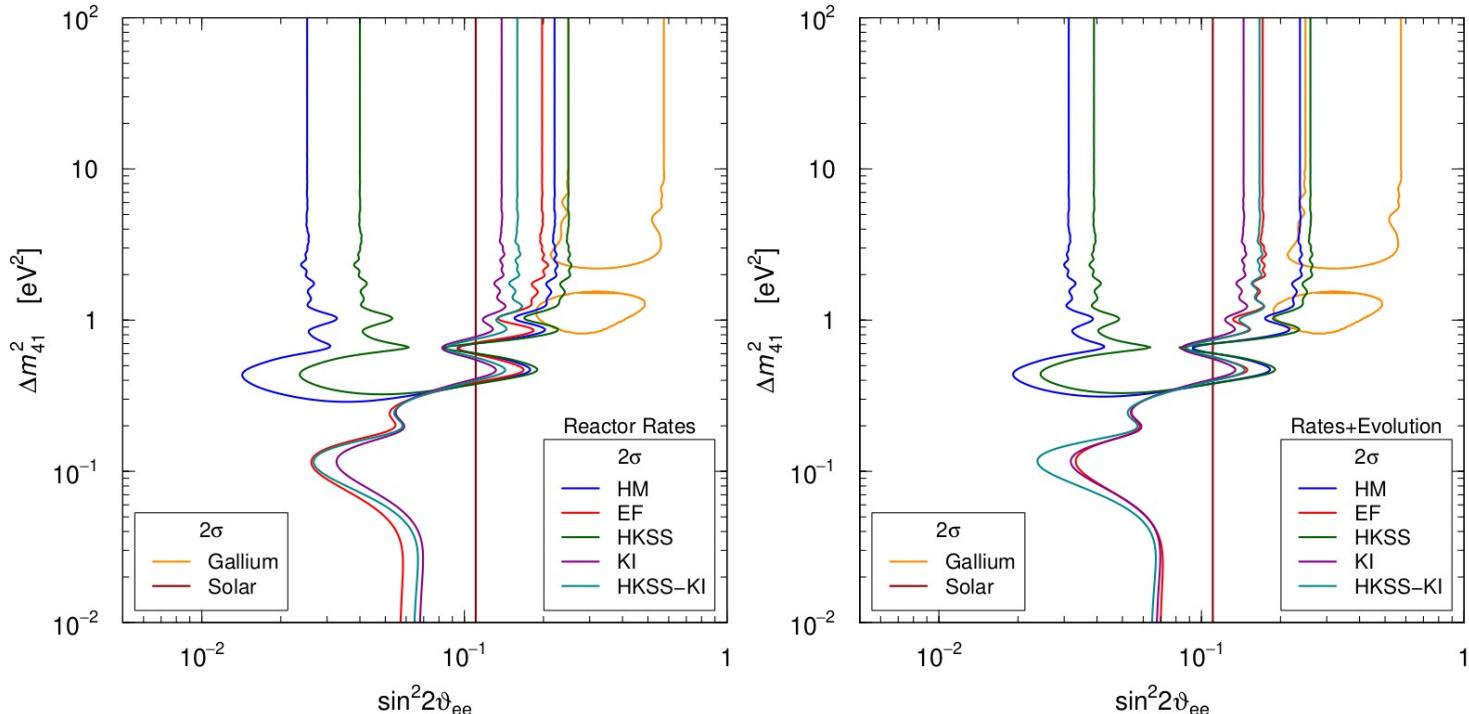
The significance of the RAA  
depends on the input flux model

The EF and KI models have no  
anomaly

Mention, Fechner, Lasserre, Mueller, Lhuillier,  
1101.2755, PRD 2011 Huber, 1106.0687, PRC 2012  
Mueller, Lhuillier, Fallot, Letourneau, Cormon, 1101.2663, PRC 2012  
Estienne, Fallot, et al, 1904.09358, PRL 2019  
Hayen, Kostensalo, Severijns, Suhonen, 1908.08302, PRC 2019  
Kurchatov Institute: Kopeikin, Skorokhvatov, Titov, 2103.01684, PRD 2021



# Impact on neutrino oscillations



Rate data are (for any flux model) in tension with Gallium data

BEST coll., 2109.11482

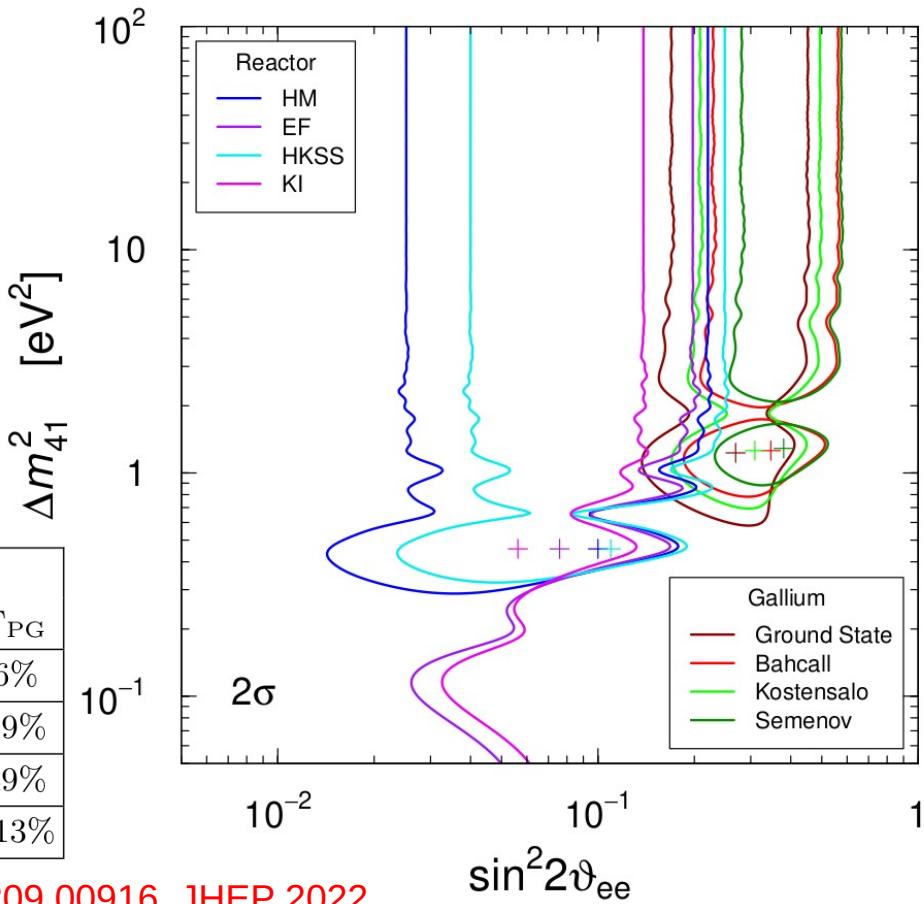
Rate data are in agreement with solar neutrino bounds

Goldhagen, et al., 2109.14898, EPJC 2022

# Tension between reactor rate and Gallium data

Severe tension between reactor  
rate and Gallium data!

	HM		HKSS		EF		KI	
	$\Delta\chi^2_{\text{PG}}$	GoF <sub>PG</sub>						
Ground State	7.2	2.8%	5.4	6.8%	9.1	1.1%	11.9	0.26%
Bahcall	10.9	0.42%	8.9	1.2%	12.9	0.16%	16.3	0.029%
Kostensalo	9.6	0.83%	7.5	2.4%	11.5	0.31%	15.3	0.049%
Semenov	15.1	0.052%	12.6	0.18%	17.0	0.02%	22.5	0.0013%

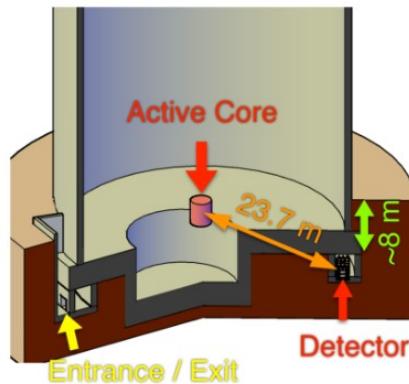


Giunti, Li, Ternes, Tyagi, Zhao, 2209.00916, JHEP 2022

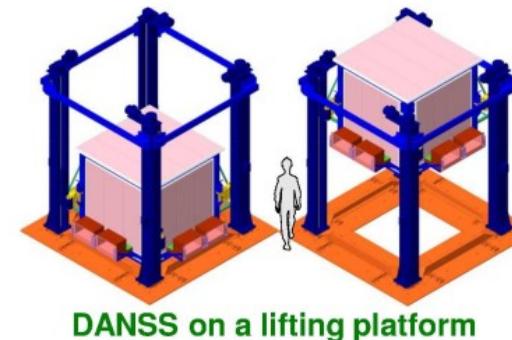
# Ratio analysis

Talk by Bryce Littlejohn!

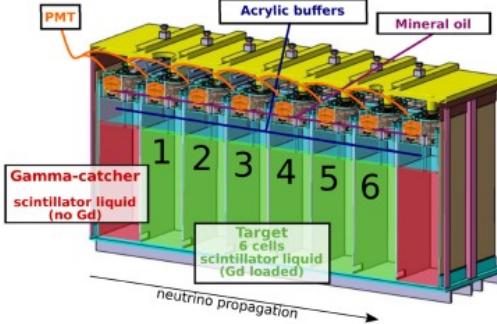
NEOS



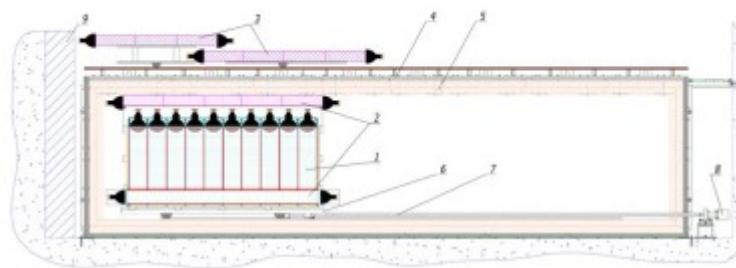
DANSS



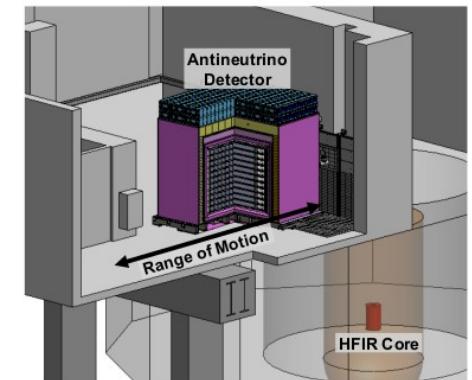
STEREO



Neutrino-4



PROSPECT



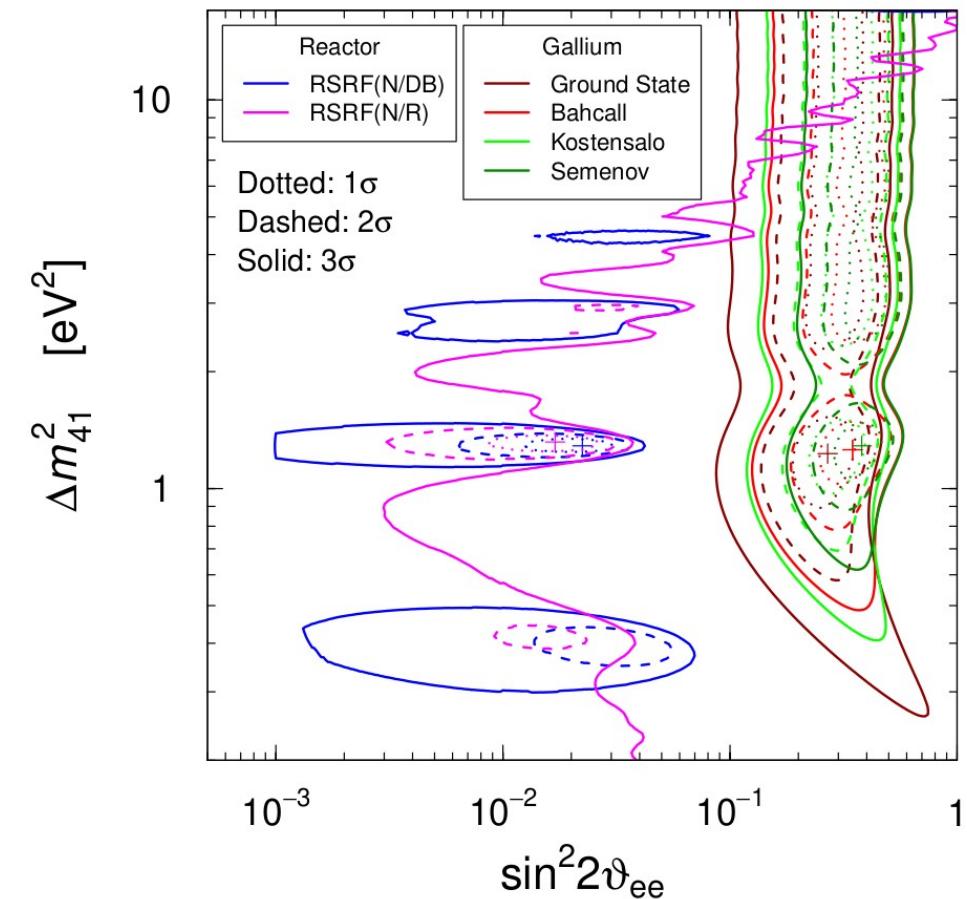
# Ratio analysis

Giunti, Li, Ternes, Tyagi, Zhao,  
2209.00916, JHEP 2022

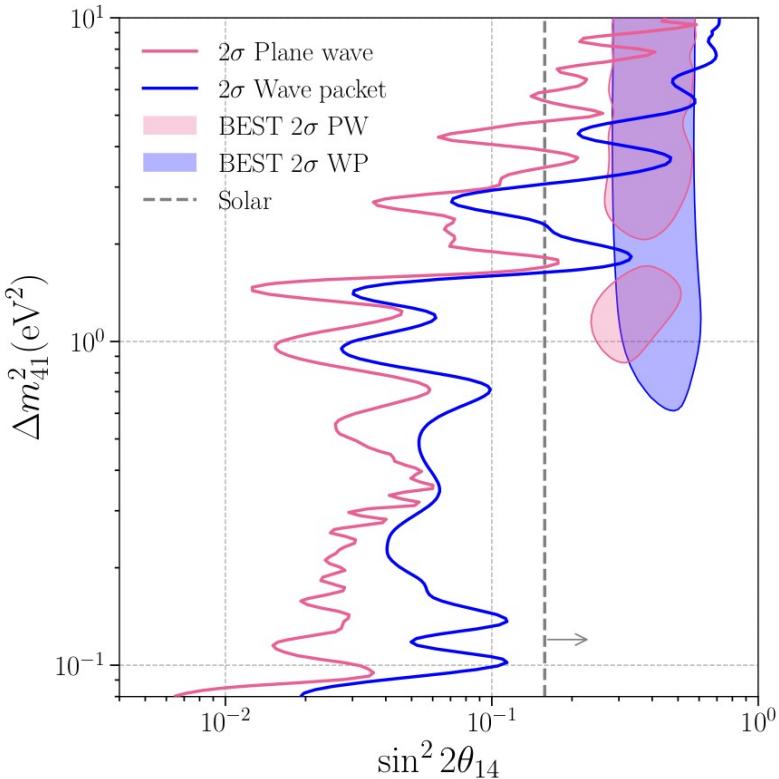
Severe tension between  
RSRF(N/DB) and Gallium  
data!

No good fit for RSRF(N/R)  
either.

	RSRF(N/DB)		RSRF(N/R)	
	$\Delta\chi^2_{\text{PG}}$	$\text{GoF}_{\text{PG}}$	$\Delta\chi^2_{\text{PG}}$	$\text{GoF}_{\text{PG}}$
Ground State	12.95	0.15%	8.91	1.2%
Bahcall	12.86	0.16%	8.74	1.3%
Kostensalo	12.91	0.16%	8.89	1.2%
Semenov	12.88	0.16%	8.70	1.3%



# Wavepackets



It was argued that if the size of the neutrino wavepacket is small the tension is reduced

$$P_{ee}^{\text{WP}} = 1 - \frac{1}{2} \sin^2 2\vartheta_{ee} \left[ 1 - \cos \left( \frac{\Delta m_{41}^2 L}{2E} \right) e^{-\left( \frac{L}{L_{\text{coh}}} \right)^2} \right]$$

$$L_{\text{coh}} = \frac{4\sqrt{2}E^2}{|\Delta m_{41}^2|} \sigma$$

But: The estimated size of the wavepacket is very large

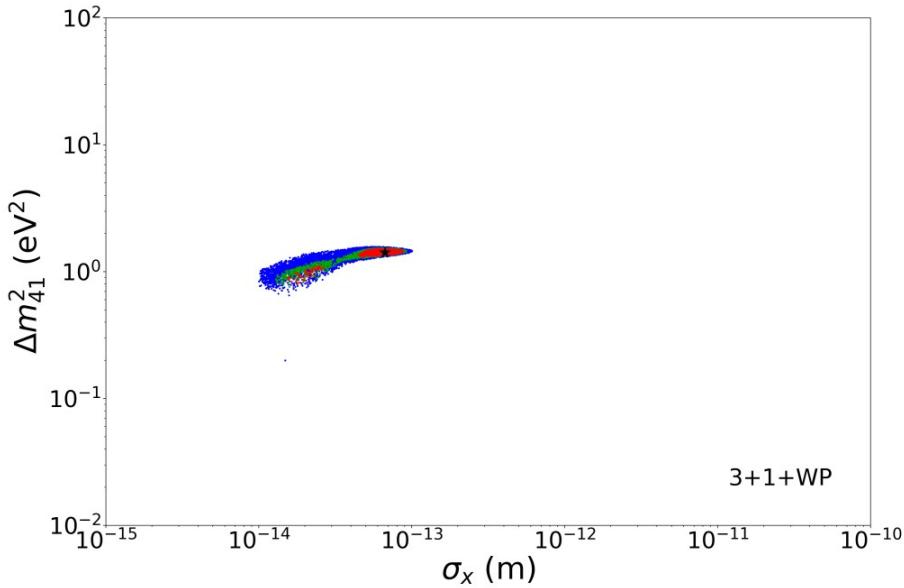
$$\sigma_x = (2 \times 10^{-5} - 1.4 \times 10^{-4}) \text{ cm}$$

Argüelles, Bertález-Martínez, Salvado,  
2201.05108, PRD 2023

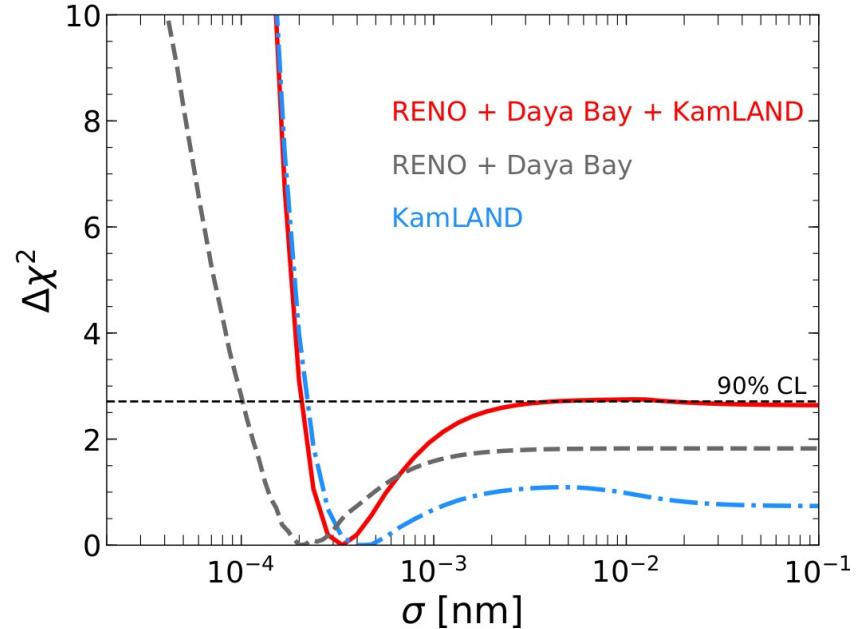
Akhmedov, Smirnov, 2208.03736, JHEP 2022

# Wavepackets

Hardin, et al, 2211.02610, JHEP 2023



de Gouvêa, De Romeri, Ternes, 2104.05806, JHEP 2021



The required size of the wavepacket is in tension with other bounds

# Neutrino decay or broad neutrinos?

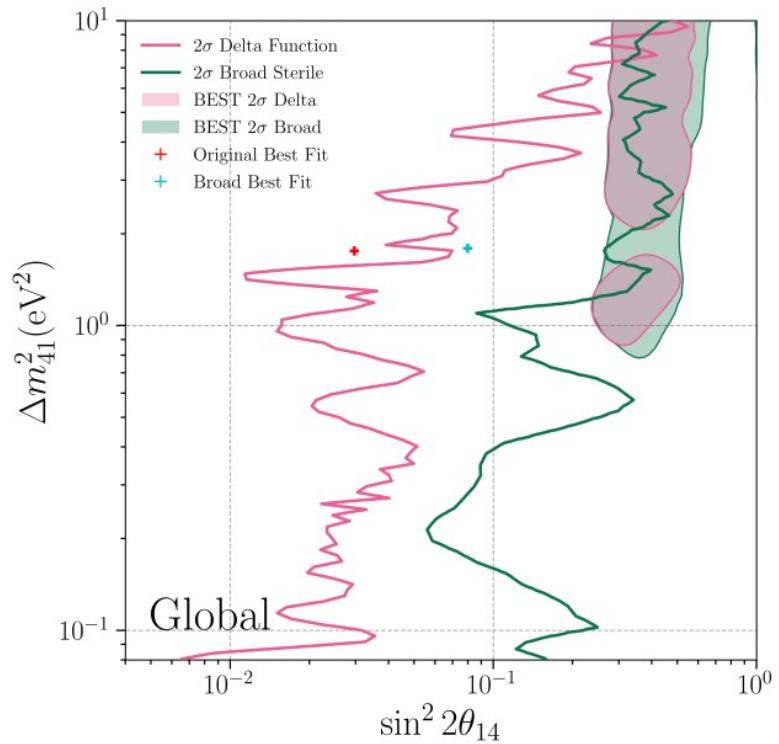
Hardin, et al, 2211.02610, JHEP 2023

Banks, Kelly, McCullough, Zhou, 2311.06352

$$P_{ee}^{\text{dec.}} = 2|U_{e4}|^2(1 - |U_{e4}|^2)e^{-\frac{\Gamma m_4 L}{2E}} \cos\left(\frac{\Delta m_{41}^2 L}{2E}\right) + |U_{e4}|^4 e^{-\frac{\Gamma m_4 L}{E}} + (1 - |U_{e4}|^2)^2$$

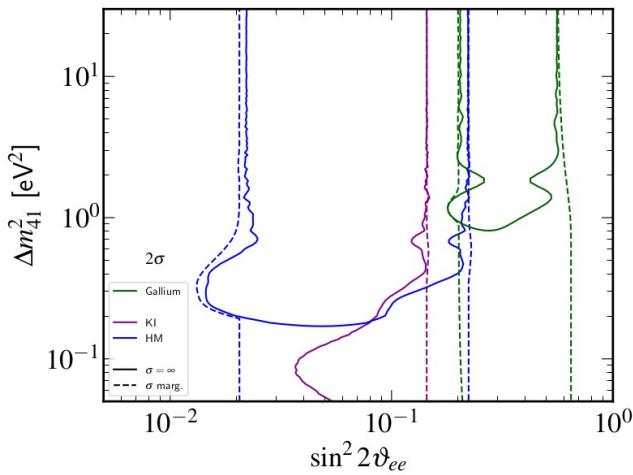
$$P_{ee}^b = \left(1 + \left(\text{sinc}\left(\frac{bL}{4E}\right) - 1\right)|U_{e4}|^2\right)^2 - 4|U_{e4}|^2(1 - |U_{e4}|^2)\sin^2\left(\frac{\Delta m_{41}^2 L}{4E}\right)\text{sinc}\left(\frac{bL}{4E}\right)$$

Alternatively, the neutrino decay or broad neutrino mass distributions have been suggested to alleviate the tension

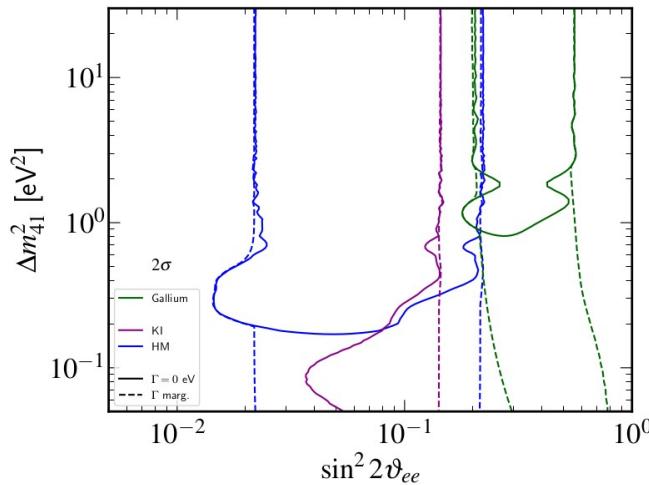


# Gallium-Reactor tension

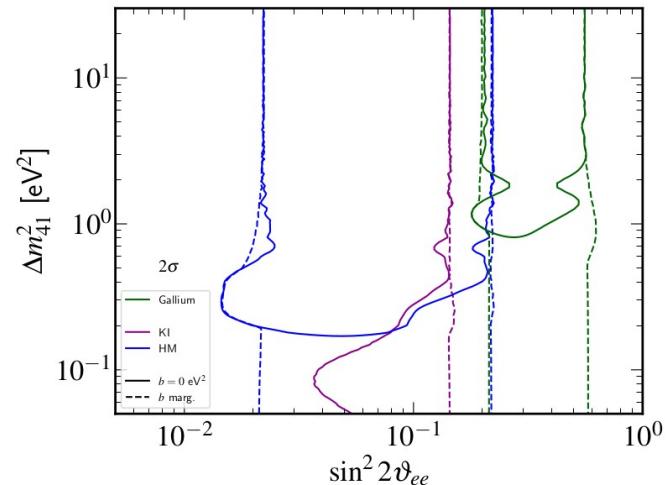
Reactor rate data was not considered in any of the previously discussed analyses



Wavepacket!



Decay!



Broad neutrinos!

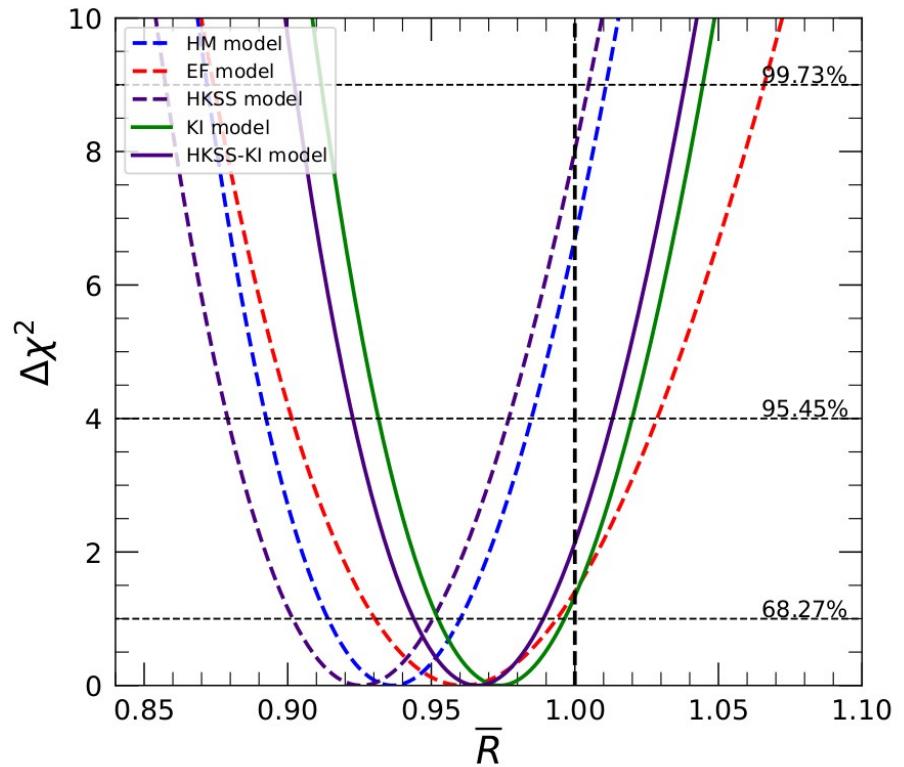
# Conclusions

RAA mostly resolved for some flux models

Good agreement between KI and EF predictions! The argument that reactor rate data is not reliable is outdated!

Reactor rate data excludes the sterile neutrino explanation of the Gallium anomaly

Reactor rate data eliminates several models proposed to alleviate tensions among data sets



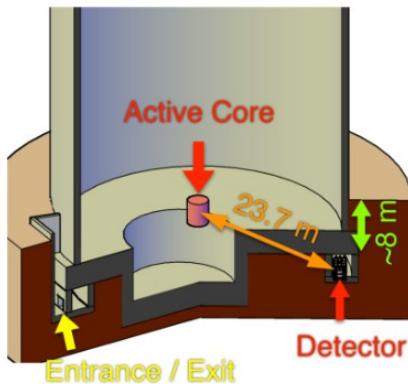


Thanks!

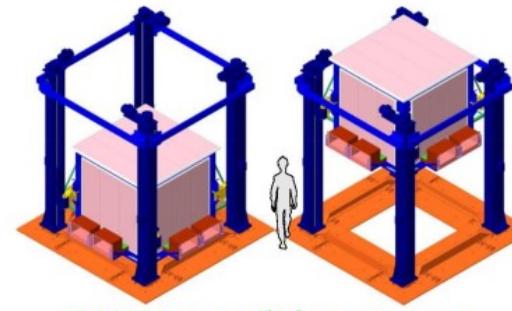


# Ratio analysis

NEOS

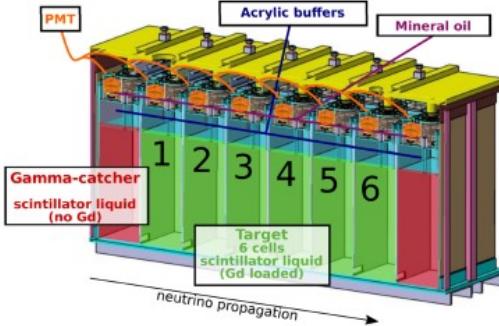


DANSS

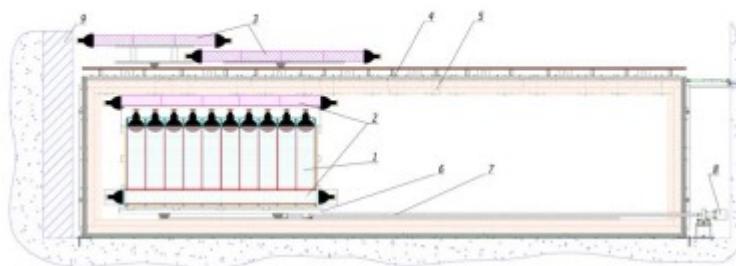


DANSS on a lifting platform

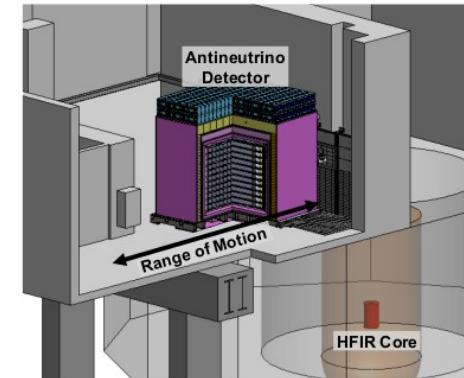
STEREO



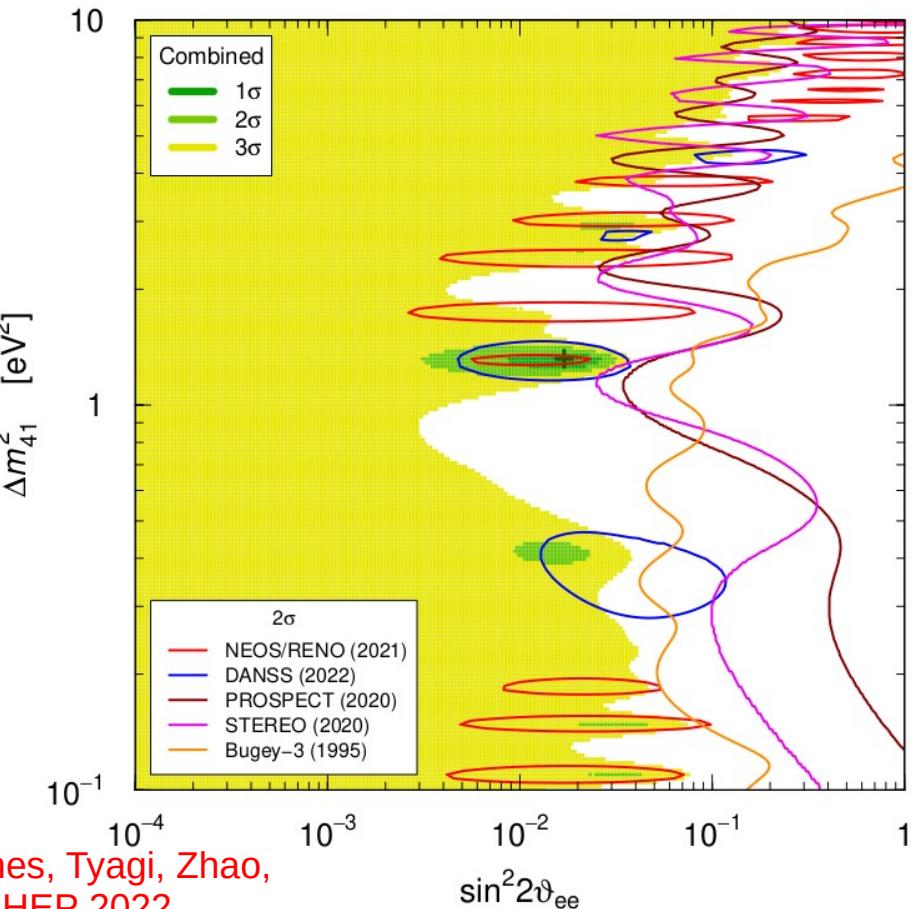
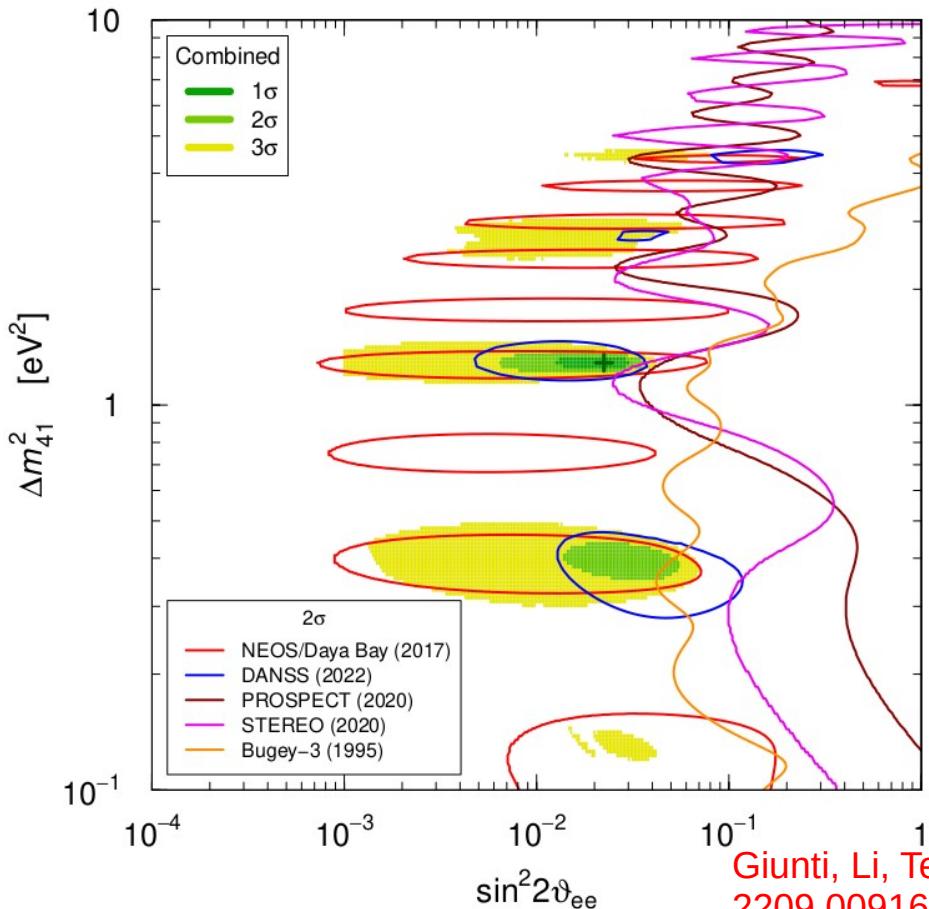
Neutrino-4



PROSPECT

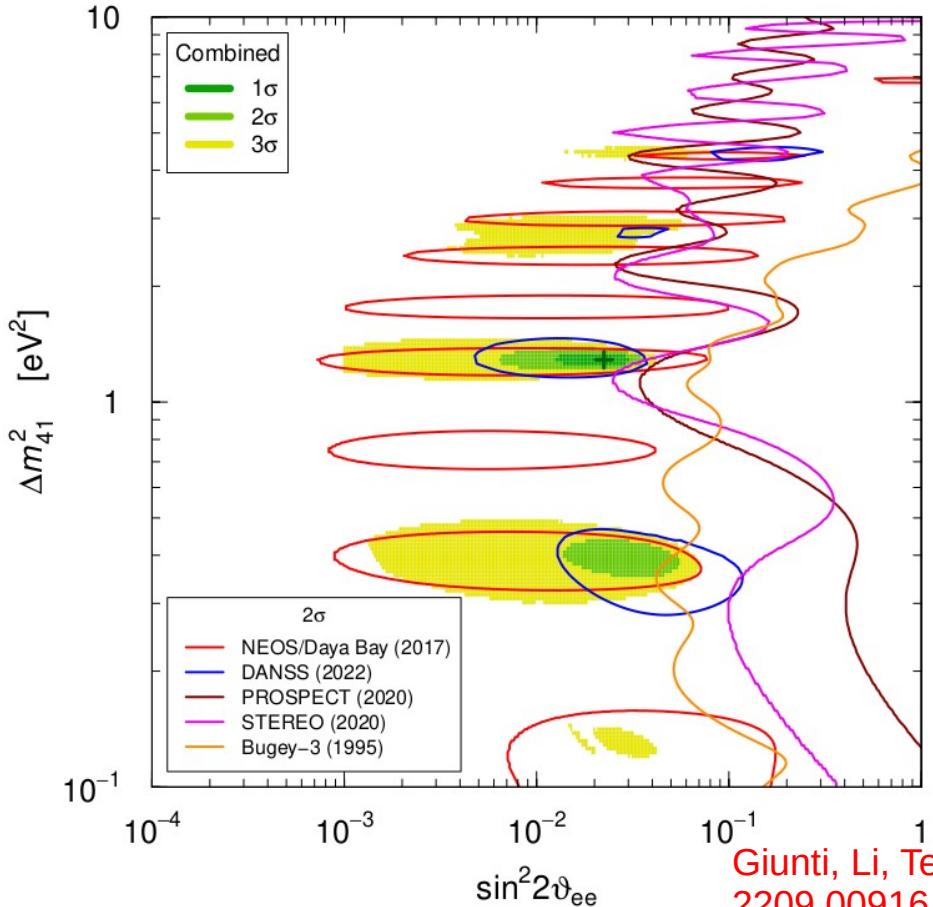


# Ratio analysis



Giunti, Li, Ternes, Tyagi, Zhao,  
2209.00916, JHEP 2022

# Ratio analysis



The NEOS collaboration performed an analysis using the Daya Bay spectrum as a reference spectrum

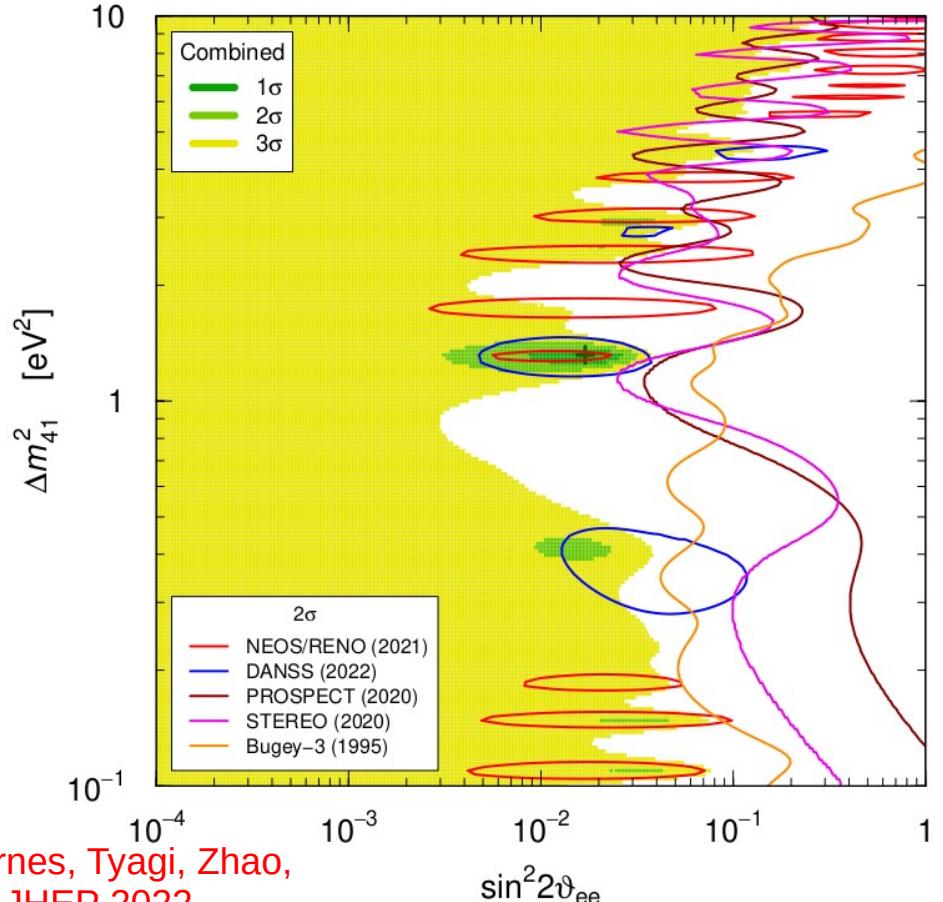
Many many events at Daya Bay!

Giunti, Li, Ternes, Tyagi, Zhao,  
2209.00916, JHEP 2022

# Ratio analysis

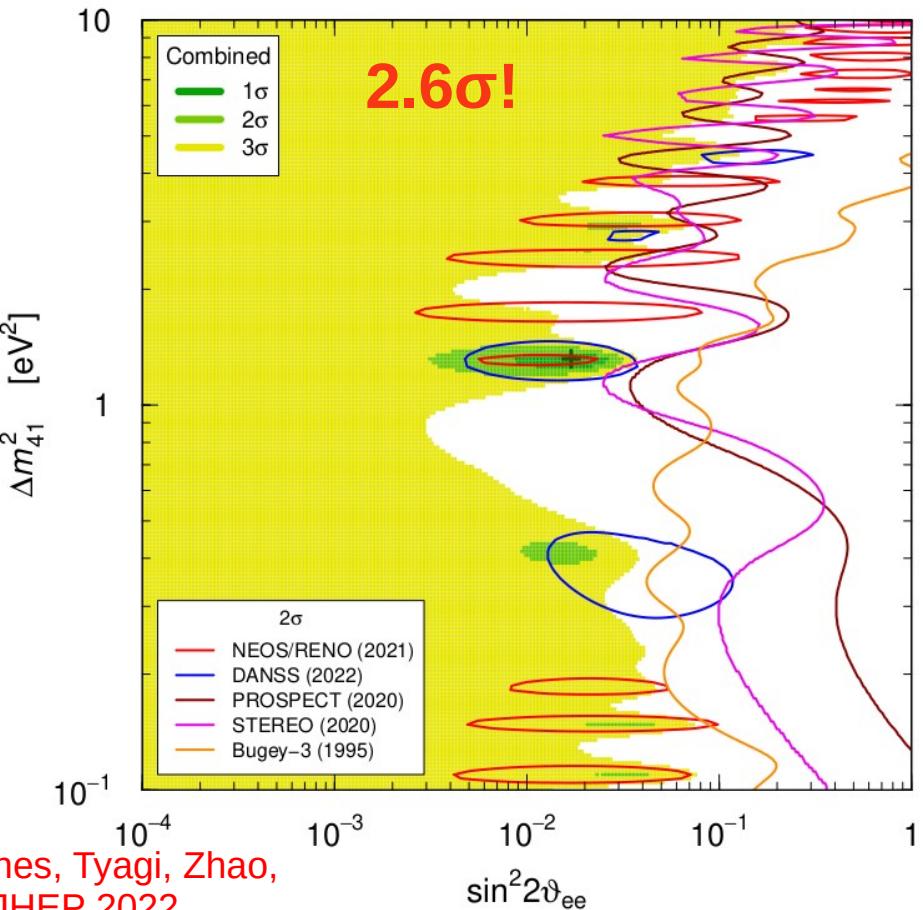
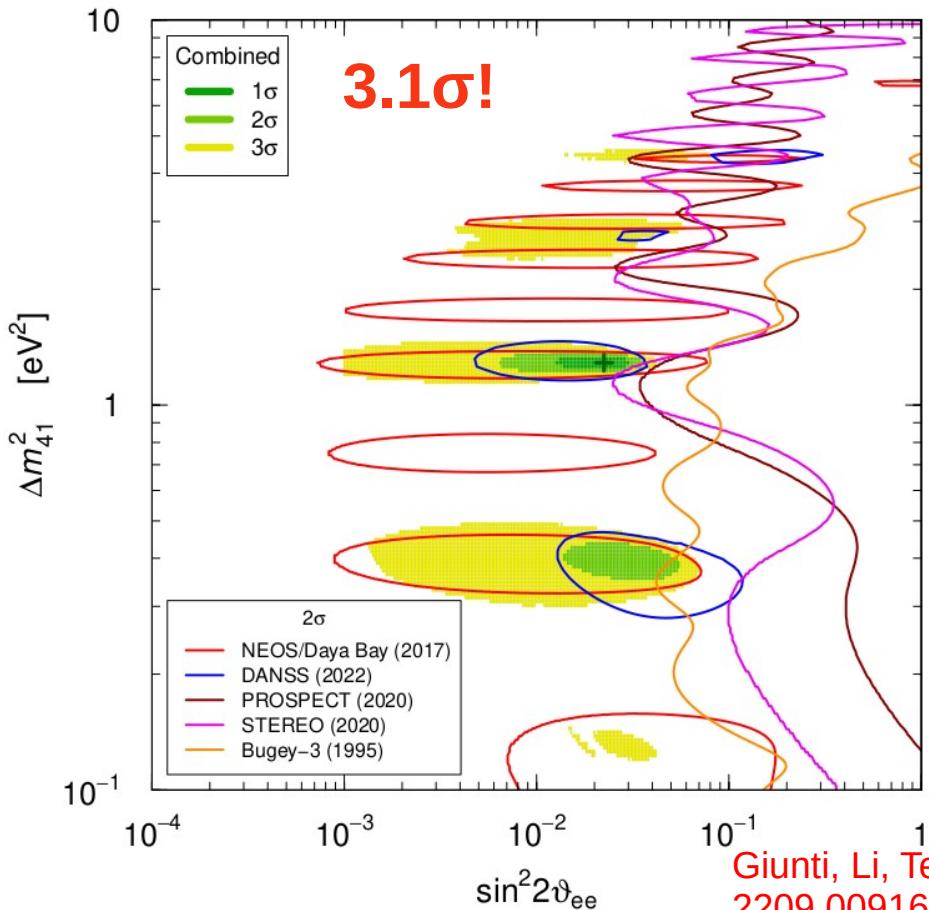
The NEOS collaboration also performed an analysis using the RENO spectrum as a reference spectrum

Same reactor complex,  
better control of systematic  
uncertainties!



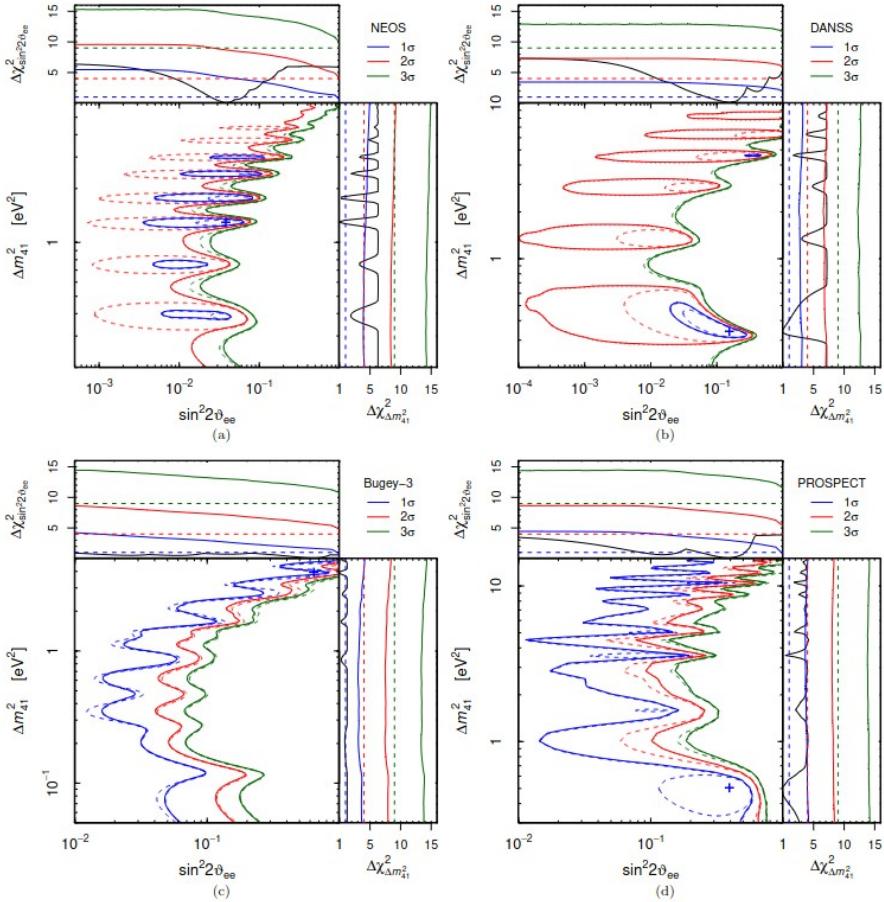
Giunti, Li, Ternes, Tyagi, Zhao,  
2209.00916, JHEP 2022

# Ratio analysis



Giunti, Li, Ternes, Tyagi, Zhao,  
2209.00916, JHEP 2022

# Caution!



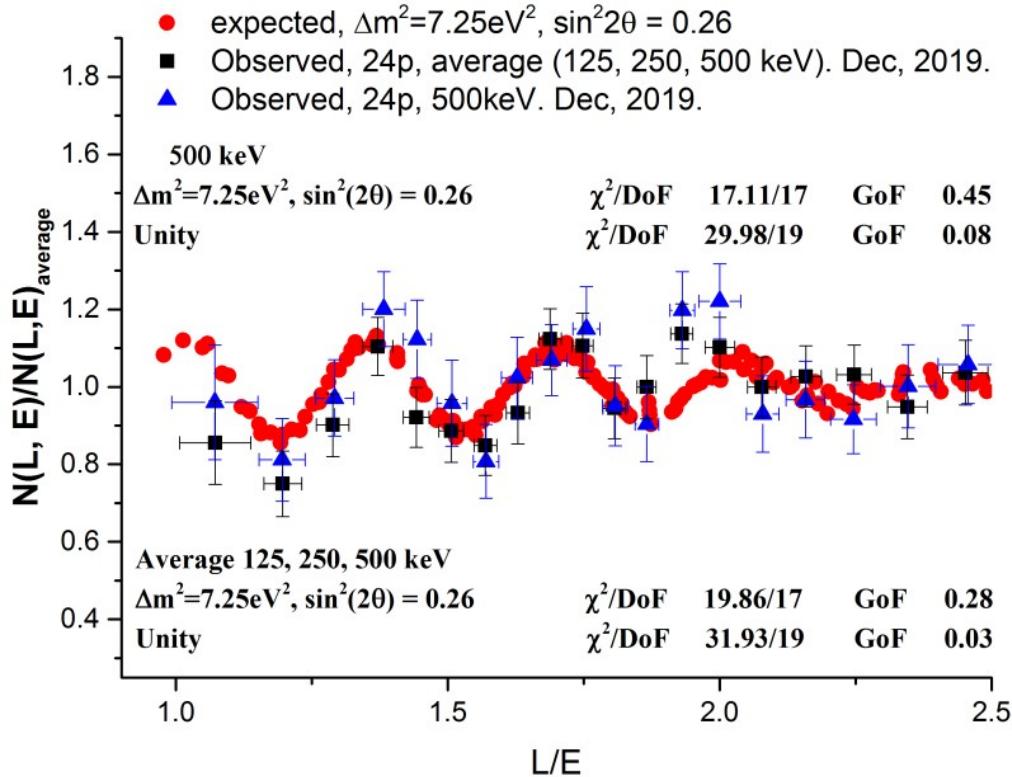
We performed simple  $\chi^2$  analyses

If one takes into account statistical fluctuation of the data the significance can be reduced

Giunti, 2004.07577, PRD 2020

# Neutrino-4

Neutrino-4, 2005.05301, PRD 2021



Neutrino-4 observes  
sterile oscillations  
at about  $3\sigma$

Very large mixing  
In tension with solar  
data

# Neutrino-4

Neutrino-4, 2005.05301, PRD 2021

- [v1] Sat, 9 May 2020 08:02:58 UTC (4,608 KB)
- [v2] Thu, 18 Jun 2020 19:22:37 UTC (4,634 KB)
- [v3] Fri, 31 Jul 2020 15:14:06 UTC (5,803 KB)
- [v4] Sun, 16 Aug 2020 19:05:32 UTC (5,849 KB)
- [v5] Sun, 14 Feb 2021 10:27:34 UTC (4,406 KB)
- [v6] Sun, 21 Feb 2021 07:51:12 UTC (4,405 KB)
- [v7]** Mon, 5 Apr 2021 15:21:56 UTC (5,488 KB)
- [v8] Tue, 25 May 2021 15:21:59 UTC (5,479 KB)

# Neutrino-4

Neutrino-4, 2005.05301, PRD 2021



- [v1] Sat, 9 May 2020 08:02:58 UTC (4,608 KB)  $2.8\sigma$
- [v2] Thu, 18 Jun 2020 19:22:37 UTC (4,634 KB)  $2.8\sigma$
- [v3] Fri, 31 Jul 2020 15:14:06 UTC (5,803 KB)  $4.6\sigma$  (added Gallium data)
- [v4] Sun, 16 Aug 2020 19:05:32 UTC (5,849 KB)  $4.6\sigma$
- [v5] Sun, 14 Feb 2021 10:27:34 UTC (4,406 KB)  $2.4\sigma$  (removed Gallium data)
- [v6] Sun, 21 Feb 2021 07:51:12 UTC (4,405 KB)  $3.2\sigma$  (?????)
- [v7] Mon, 5 Apr 2021 15:21:56 UTC (5,488 KB)  $2.9\sigma$
- [v8] Tue, 25 May 2021 15:21:59 UTC (5,479 KB)  $2.7\sigma$ - $2.9\sigma$

# Neutrino-4

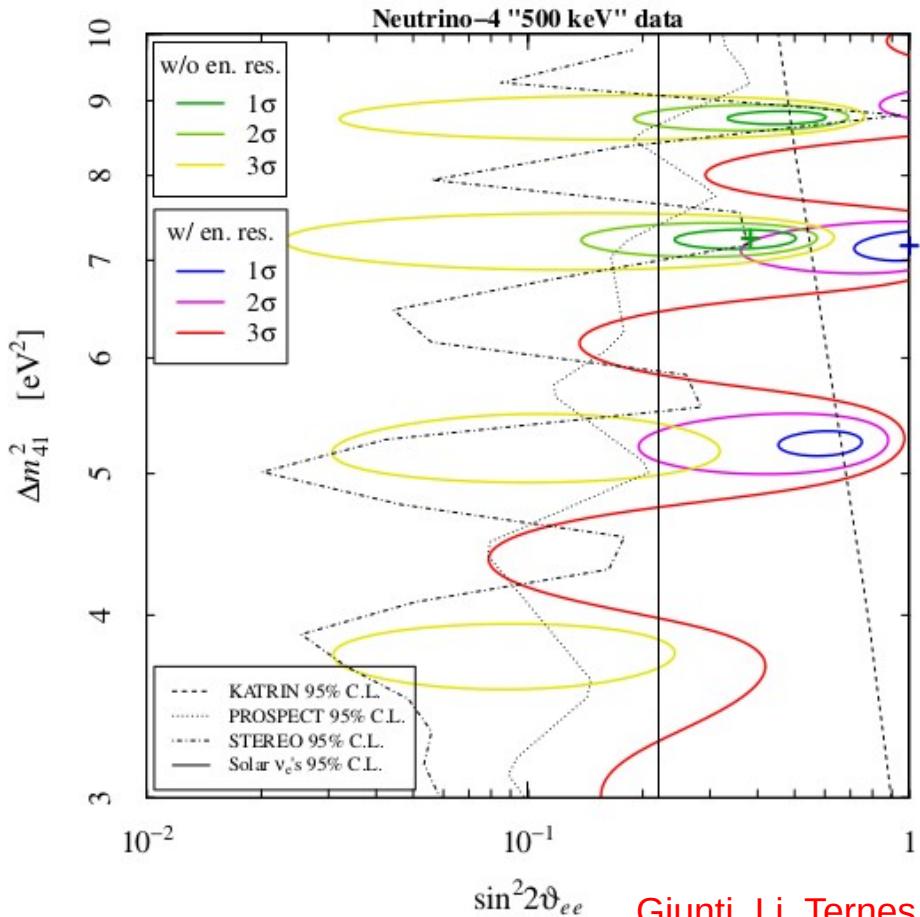
Averaging contains integration over flux, distance, detector resolution

$$\left\langle \sin^2 \left( \frac{\Delta m_{41}^2 L}{4E} \right) \right\rangle_{ik} = \frac{\int_{L_k^{\min}}^{L_k^{\max}} dL L^{-2} \int_{E_i^{\min}}^{E_i^{\max}} dE'_p \int dE_p R(E_p, E'_p) \sin^2 \left( \frac{\Delta m_{41}^2 L}{4E} \right) \phi_{\bar{\nu}_e}(E) \sigma_{\bar{\nu}_e p}(E)}{\int_{L_k^{\min}}^{L_k^{\max}} dL L^{-2} \int_{E_i^{\min}}^{E_i^{\max}} dE'_p \int dE_p R(E_p, E'_p) \phi_{\bar{\nu}_e}(E) \sigma_{\bar{\nu}_e p}(E)}$$

Using energy calibration information from 2005.05301 we extract the approximate energy resolution function

$$R(E_p, E'_p) = \frac{1}{\sqrt{2\pi}\sigma_{E_p}} \exp \left( -\frac{(E_p - E'_p)^2}{2\sigma_{E_p}^2} \right) \quad \sigma_{E_p} = 0.19 \sqrt{\frac{E_p}{\text{MeV}}} \text{ MeV.}$$

# Neutrino-4

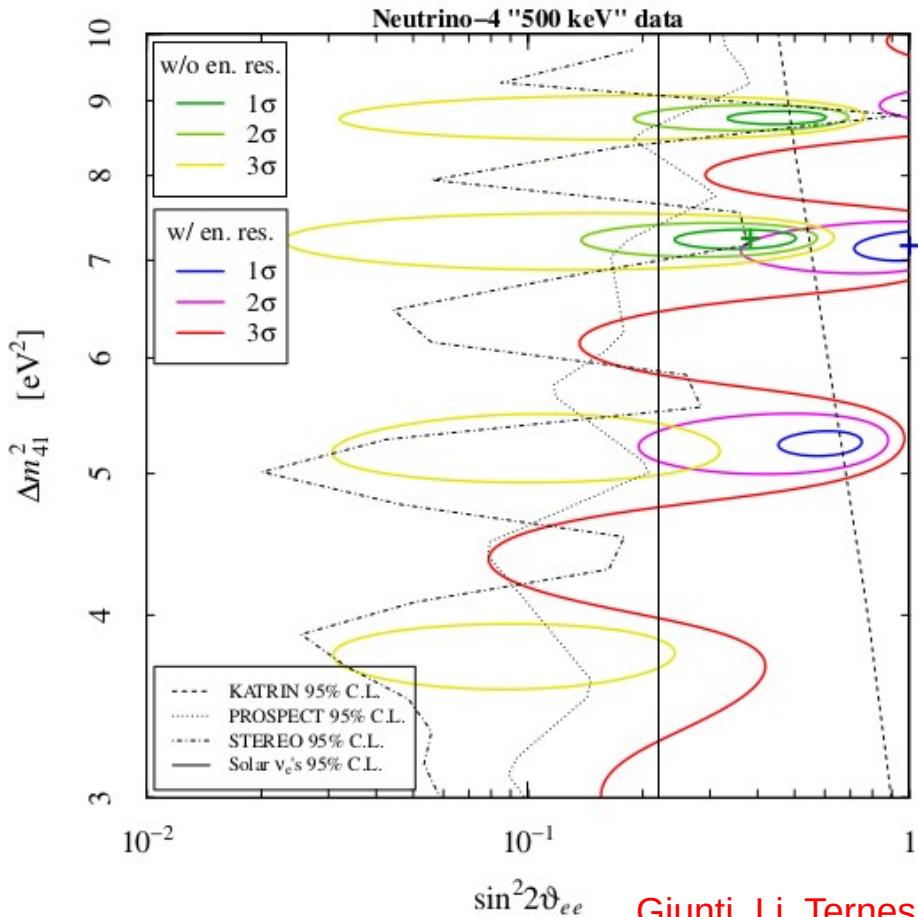


We can only reproduce Neutrino-4 confidence regions when not including energy resolution

Inclusion shifts the best fit to even larger values, but reduces the preference for sterile oscillations

Giunti, Li, Ternes, Zhang, 2101.06785, PLB 2021

# Neutrino-4



Giunti, Li, Ternes, Zhang, 2101.06785, PLB 2021

We can only reproduce Neutrino-4 confidence regions when not including energy resolution  
Inclusion shifts the best fit to even larger values, but reduces the preference for sterile oscillations

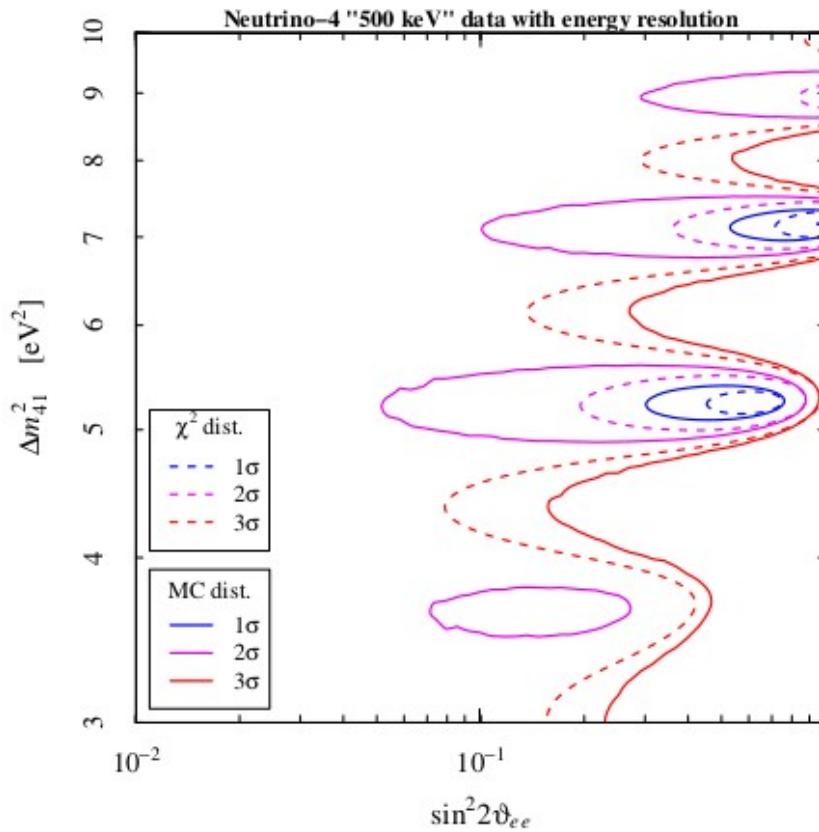
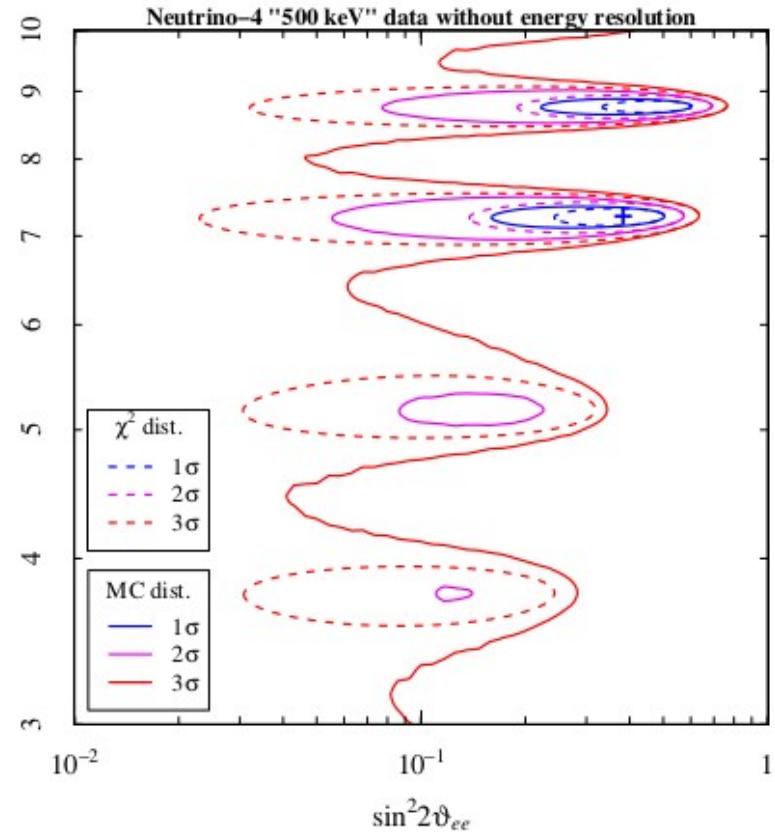
$$\chi^2 = \sum_{j=1}^{19} \left( \frac{R_j^{\text{the}} - R_j^{\text{exp}}}{\Delta R_j^{\text{exp}}} \right)^2$$

# Monte Carlo analysis

Giunti, Li, Ternes, Zhang, 2101.06785, PLB 2021

# Neutrino-4

See also: Coloma, Huber, Schwetz,  
2008.06083, EPJC 2021



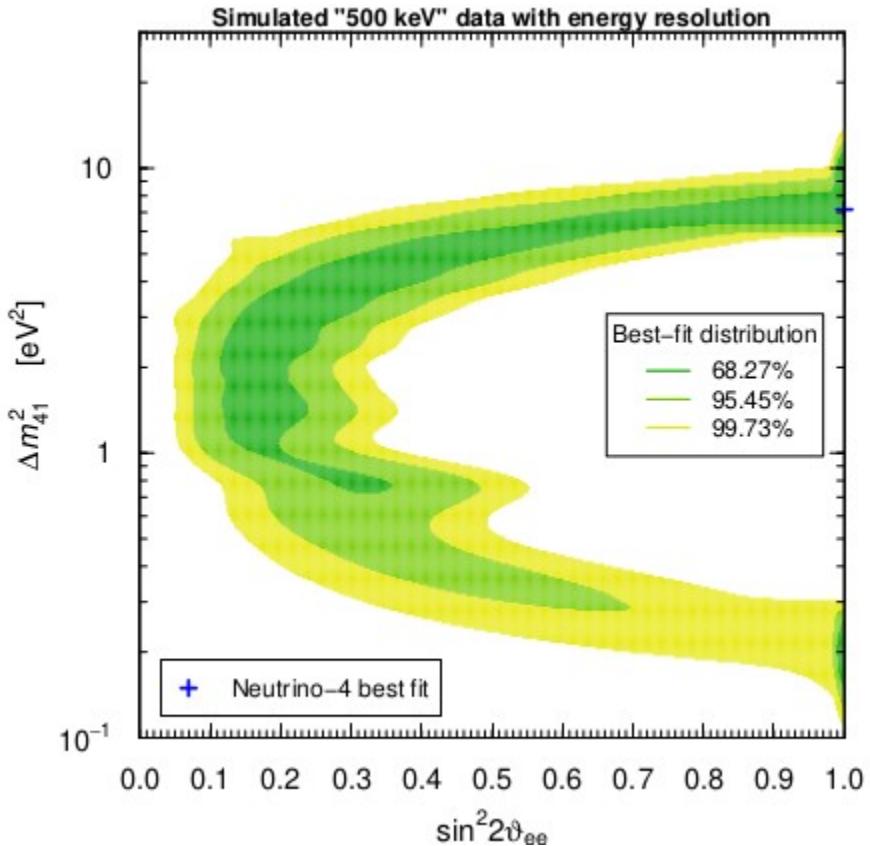
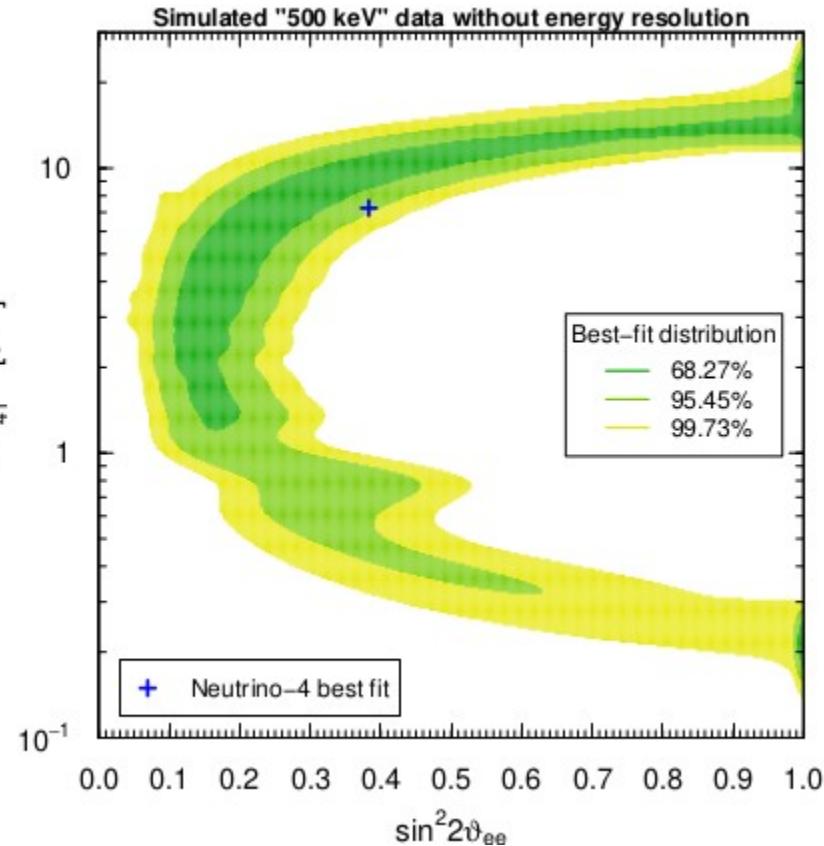
# Neutrino-4

## Summary

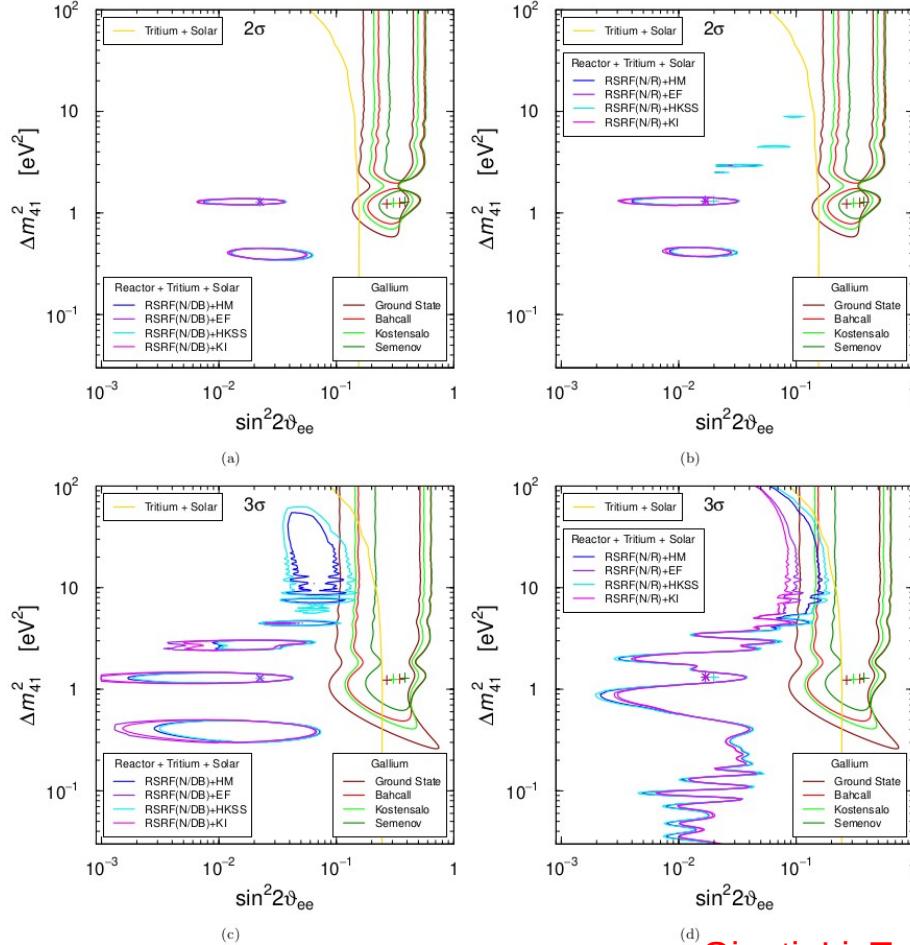
Neutrino-4	"500 keV" data		"125-250-500 keV" data	
	without en. res.	with en. res.	without en. res.	with en. res.
$\chi^2_{\text{min}}$	14.9	18.2	21.9	21.1
GoF	60%	37%	19%	22%
$(\sin^2 2\vartheta_{ee})_{\text{bf}}$	0.38	1.0	0.27	0.93
$(\Delta m^2_{41})_{\text{bf}}$	7.2	7.2	8.8	7.2
$\Delta\chi^2_{\text{NO}}$	13.1	9.8	9.9	10.7
$\chi^2$ distribution				
$p$ -value	0.0014	0.0075	0.0072	0.0048
$\sigma$ -value	3.2	2.7	2.7	2.8
Monte Carlo distribution				
$p$ -value	0.011	0.028	0.087	0.026
$\sigma$ -value	2.5	2.2	1.7	2.2

# Neutrino-4

## Distribution of best fit points without oscillations



# Combined reactor, Tritium, and solar data



Combination of all data!

Severe and unacceptable tension  
for any combination with Gallium  
data!

Global Fit: RSRF(N/DB) + Reactor Rates + Tritium + Solar								
	HM		HKSS		EF		KI	
	$\Delta\chi^2_{PG}$	GoF <sub>PG</sub>						
Ground State	21.54	0.0021%	19.51	0.0058%	21.92	0.0017%	21.90	0.0018%
Bahcall	25.99	0.00023%	23.88	0.00065%	26.13	0.00021%	26.11	0.00021%
Kostensalo	25.05	0.00036%	22.77	0.0011%	27.62	0.0001%	27.60	0.0001%
Semenov	32.52	0.0000087%	29.93	0.000032%	37.69	0.00000065%	38.81	0.00000037%
Global Fit: RSRF(N/R) + Reactor Rates + Tritium + Solar								
	HM		HKSS		EF		KI	
	$\Delta\chi^2_{PG}$	GoF <sub>PG</sub>						
Ground State	17.61	0.015%	15.53	0.042%	22.56	0.0013%	22.66	0.0012%
Bahcall	22.07	0.0016%	19.90	0.0048%	26.82	0.00015%	26.80	0.00015%
Kostensalo	21.11	0.0026%	18.77	0.0084%	26.27	0.0002%	28.45	0.000066%
Semenov	28.57	0.000062%	25.93	0.00023%	34.00	0.0000041%	38.24	0.0000005%

Giunti, Li, Ternes, Tyagi, Zhao, 2209.00916, JHEP 2022

Christoph Ternes

# Combined reactor, Tritium, and solar data

Global RSRF(N/DB) Fit				
	HM	HKSS	EF	KI
$\chi^2_{\text{min}}$	393.5	395.2	391.2	391.4
GoF	43%	40%	46%	46%
$(\sin^2 2\vartheta_{ee})_{\text{b.f.}}$	0.022	0.022	0.022	0.022
$(\Delta m^2_{41})_{\text{b.f.}}/\text{eV}^2$	1.29	1.29	1.29	1.29
$\Delta\chi^2_{4\nu-3\nu}$	13.8	14.1	12.6	12.9
$n\sigma_{4\nu-3\nu}$	3.3	3.3	3.1	3.2

Global RSRF(N/R) Fit				
	HM	HKSS	EF	KI
$\chi^2_{\text{min}}$	386.5	388.3	384.0	384.2
GoF	53%	50%	56%	56%
$(\sin^2 2\vartheta_{ee})_{\text{b.f.}}$	0.017	0.019	0.017	0.017
$(\Delta m^2_{41})_{\text{b.f.}}/\text{eV}^2$	1.32	1.32	1.32	1.32
$\Delta\chi^2_{4\nu-3\nu}$	10.1	10.3	9.1	9.3
$n\sigma_{4\nu-3\nu}$	2.7	2.8	2.6	2.6

Global fit (without Gallium data)  
has a preference between  $2.6\sigma$   
and  $3.3\sigma$  in favor of 3+1  
oscillations!

Due to new reactor ratio data

# More on the Gallium anomaly

Giunti, Li, Ternes, Tyagi, Zhao, 2209.00916, JHEP 2022

Model	Method	$^{51}\text{Cr}$		$^{37}\text{Ar}$	
		$\sigma_{\text{tot}}$	$\delta_{\text{exc}}$	$\sigma_{\text{tot}}$	$\delta_{\text{exc}}$
Ground State	$T_{1/2}(^{71}\text{Ge})$	$5.539 \pm 0.019$	–	$6.625 \pm 0.023$	–
Bahcall (1997)	$^{71}\text{Ga}(p, n)^{71}\text{Ge}$	$5.81 \pm 0.16$	4.7%	$7.00 \pm 0.21$	5.4%
Haxton (1998)	Shell Model	$6.39 \pm 0.65$	13.3%	$7.72 \pm 0.81$	14.2%
Frekers et al. (2015)	$^{71}\text{Ga}({}^3\text{He}, {}^3\text{H})^{71}\text{Ge}$	$5.92 \pm 0.11$	6.4%	$7.15 \pm 0.14$	7.3%
Kostensalo et al. (2019)	Shell Model	$5.67 \pm 0.06$	2.3%	$6.80 \pm 0.08$	2.6%
Semenov (2020)	$^{71}\text{Ga}({}^3\text{He}, {}^3\text{H})^{71}\text{Ge}$	$5.938 \pm 0.116$	6.7%	$7.169 \pm 0.147$	7.6%

# More on the Gallium anomaly

Giunti, Li, Ternes, Tyagi, Zhao, 2209.00916, JHEP 2022

Model	Method	$^{51}\text{Cr}$		$^{37}\text{Ar}$	
		$\sigma_{\text{tot}}$	$\delta_{\text{exc}}$	$\sigma_{\text{tot}}$	$\delta_{\text{exc}}$
Ground State	$T_{1/2}(^{71}\text{Ge})$	$5.539 \pm 0.019$		$6.625 \pm 0.023$	–
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Semenov (2020)	$^{71}\text{Ga}({}^3\text{He}, {}^3\text{H})^{71}\text{Ge}$	$5.938 \pm 0.116$	6.7%	$7.169 \pm 0.147$	7.6%

Small

# More on the Gallium anomaly

Giunti, Li, Ternes, Tyagi, Zhao, 2209.00916, JHEP 2022

Model	Method	$^{51}\text{Cr}$		$^{37}\text{Ar}$	
		$\sigma_{\text{tot}}$	$\delta_{\text{exc}}$	$\sigma_{\text{tot}}$	$\delta_{\text{exc}}$
Ground State	$T_{1/2}(^{71}\text{Ge})$	$5.539 \pm 0.019$	–	$6.625 \pm 0.023$	–
Bahcall (1997)	$^{71}\text{Ga}(p, n)^{71}\text{Ge}$	$5.81 \pm 0.16$	4.7%	$7.00 \pm 0.21$	5.4%
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Semenov (2020)	$^{71}\text{Ga}({}^3\text{He}, {}^3\text{H})^{71}\text{Ge}$	$5.938 \pm 0.116$	6.7%	$7.169 \pm 0.147$	7.6%

Large

# More on the Gallium anomaly

Giunti, Li, Ternes, Tyagi, Zhao, 2209.00916, JHEP 2022

Model	Method	$^{51}\text{Cr}$		$^{37}\text{Ar}$	
		$\sigma_{\text{tot}}$	$\delta_{\text{exc}}$	$\sigma_{\text{tot}}$	$\delta_{\text{exc}}$
Ground State	$T_{1/2}(^{71}\text{Ge})$	$5.539 \pm 0.019$	–	$6.625 \pm 0.023$	–
Bahcall (1997)	$^{71}\text{Ga}(p, n)^{71}\text{Ge}$	$5.81 \pm 0.16$	4.7%	$7.00 \pm 0.21$	5.4%
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Semenov (2020)	$^{71}\text{Ga}({}^3\text{He}, {}^3\text{H})^{71}\text{Ge}$	$5.938 \pm 0.116$	6.7%	$7.169 \pm 0.147$	7.6%

Main  
contribution

$$\sigma_{\text{tot}} = \sigma_{\text{gs}} \left( 1 + \xi_{5/2^-} \frac{\text{BGT}_{5/2^-}}{\text{BGT}_{\text{gs}}} + \xi_{3/2^-} \frac{\text{BGT}_{3/2^-}}{\text{BGT}_{\text{gs}}} + \xi_{5/2^+} \frac{\text{BGT}_{5/2^+}}{\text{BGT}_{\text{gs}}} \right)$$

Corrections

# More on the Gallium anomaly

The ground-state cross section is obtained from the half life measurement

$$\sigma_{\text{gs}} = \frac{G_F^2 \cos^2 \vartheta_C}{\pi} g_A^2 \text{BGT}_{\text{gs}} \langle p_e E_e F(Z_{\text{Ge}}, E_e) \rangle = \frac{\pi^2 \ln 2}{m_e^5 f t_{1/2}(^{71}\text{Ge})} \langle p_e E_e F(Z_{\text{Ge}}, E_e) \rangle$$

Different results obtained in the past

$$T_{1/2}^{\text{BGZZ}}(^{71}\text{Ge}) = 12.5 \pm 0.1 \text{ d} \quad (\text{Bisi, Germagnoli, Zappa, and Zimmer, 1955})$$

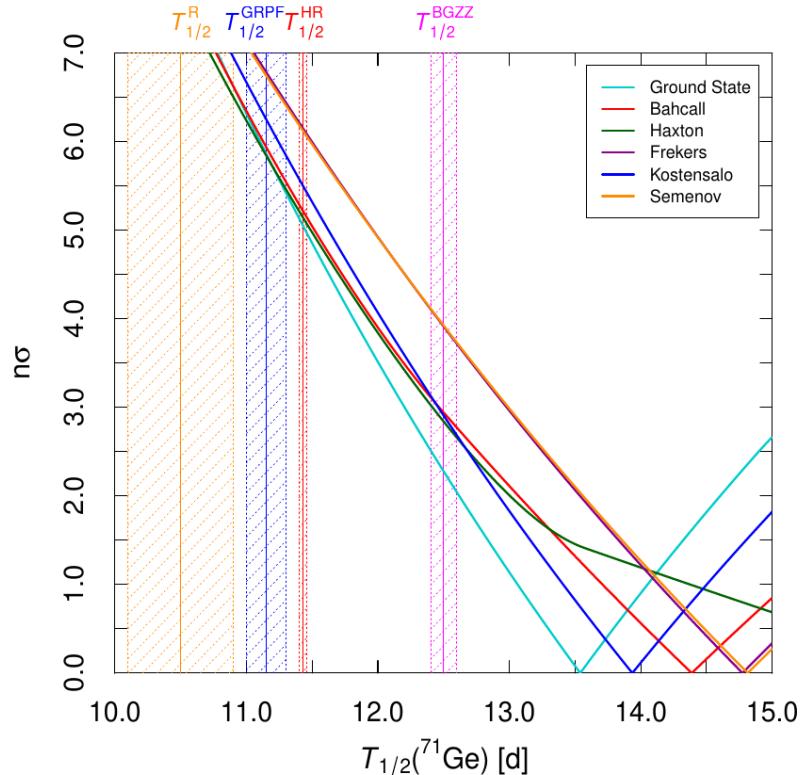
$$T_{1/2}^{\text{R}}(^{71}\text{Ge}) = 10.5 \pm 0.4 \text{ d} \quad (\text{Rudstam, 1956})$$

$$T_{1/2}^{\text{GRPF}}(^{71}\text{Ge}) = 11.15 \pm 0.15 \text{ d} \quad (\text{Genz, Renier, Pengra, and Fink, 1971})$$

$$T_{1/2}^{\text{HR}}(^{71}\text{Ge}) = 11.43 \pm 0.03 \text{ d} \quad (\text{Hampel and Remsberg, 1985})$$

# More on the Gallium anomaly

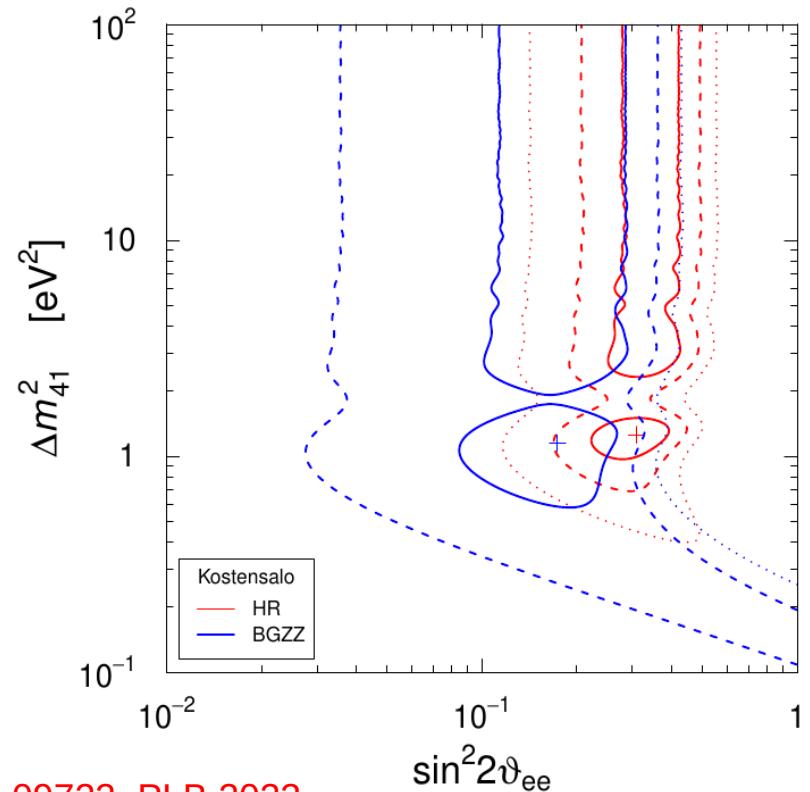
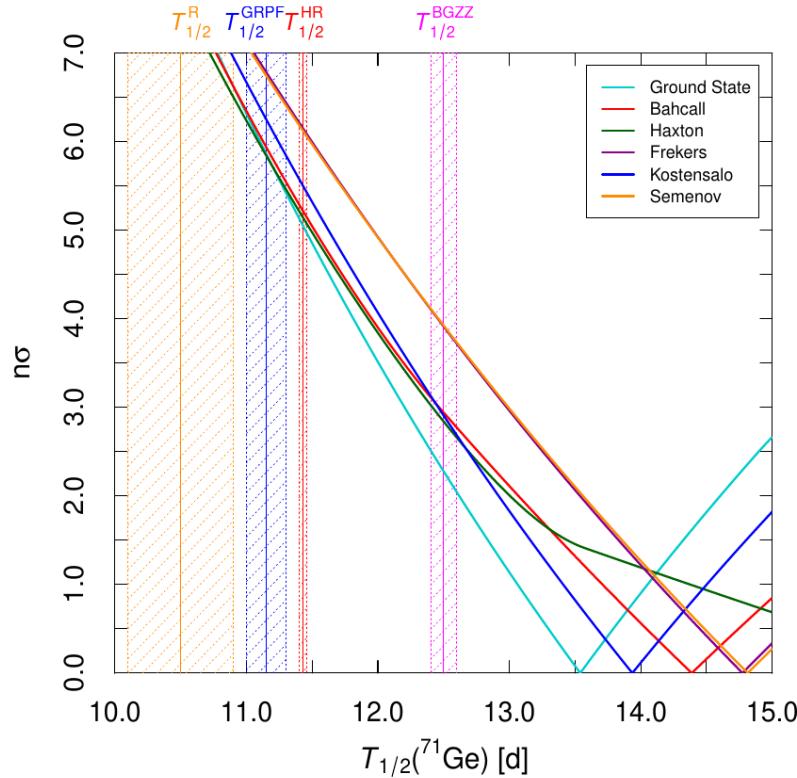
Fit the Germanium half life using data from Gallium experiments



Giunti, Li, Ternes, Zhao, 2212.09722, PLB 2023

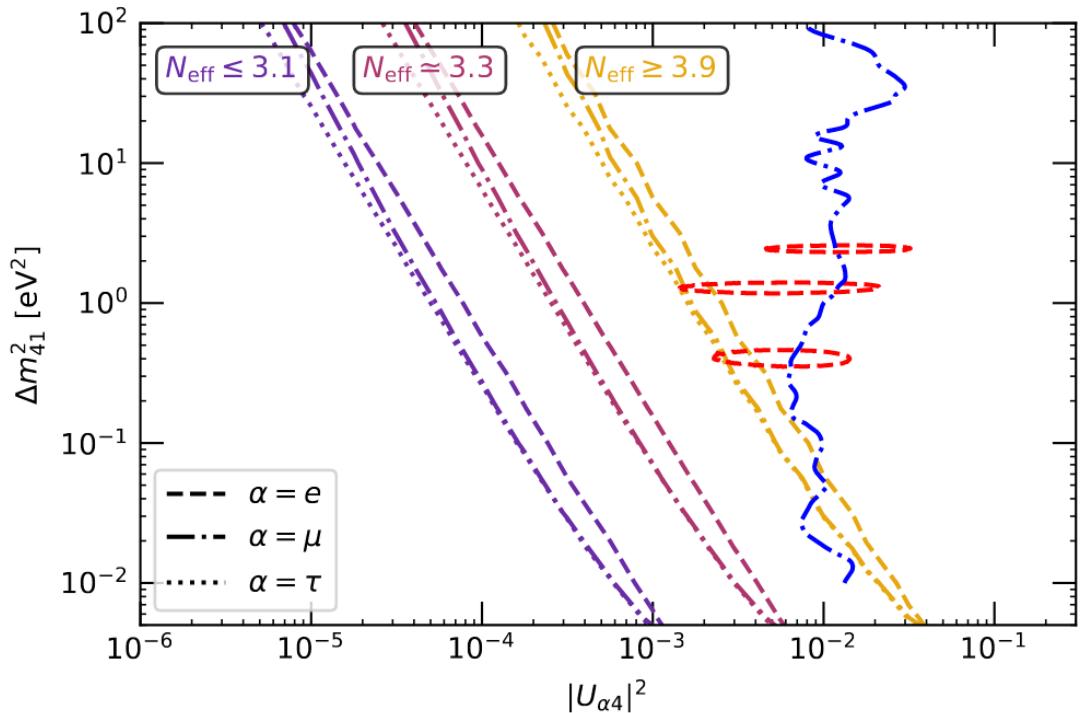
# More on the Gallium anomaly

Fit the Germanium half life using data from Gallium experiments



Giunti, Li, Ternes, Zhao, 2212.09722, PLB 2023

# Cosmology



Cosmology can set strong bounds on sterile parameter space