String Theoretic Evidence for the Weak Gravity Conjecture

Seung-Joo Lee

b Institute for Basic Science

CAU BSM Workshop 2023, CAU

21-02-2023

Disclaimer

What this talk is (not) for

- I was told to talk about the current status of the Weak Gravity
 Conjecture (WGC) a prominent example of swampland conjectures.
- There have been so many developments on the WGC that a single talk cannot cover them all.
- I will focus on the string theoretic aspects of the WGC (mostly on the recent developments I myself was part of through [S.-J.L., Lerche, Weigand]'s).
 I will also address a few other swampland conjectures.
- For other aspects, see e.g. the comprehensive reviews: [Palti '19] (on swampland conjectures); [Harlow, Heidenreich, Reece, Rudelius '22] (on the WGC). In particular, for the particle physics aspects, see [talk by Reece] just before mine.

Motivation

Quantum Gravity and String Theory

• Swampland Conjectures

- Which effective field theories (EFTs) couple to Quantum Gravity?
 - Swampland v.s. Landscape
- EFTs in the Landscape subject to universal consistency constraints
- Swampland Conjectures
 -> general, useful, but not fully understood
- Stringy Realization
 - Quantitative verification of the explicit conjectures
 - Manifestations in string geometry
 - Refinement

picture from [Palti '19]

Swampland

Quantum Gravity

Landscape

Motivation

Quantum Gravity and String Theory

• Swampland Conjectures

- Which effective field theories (EFTs) couple to Quantum Gravity?
 - Swampland v.s. Landscape
- EFTs in the Landscape subject to universal consistency constraints
 - Swampland Conjectures
 -> general, useful, but **not** fully understood
- Stringy Realization
 - Quantitative verification of the explicit conjectures
 - Manifestations in string geometry
 - Refinement

Quantum Gravity

Landscape

Swampland

picture from [Palti '19]

Motivation

Quantum Gravity and String Theory

• Swampland Conjectures

- Which effective field theories (EFTs) couple to Quantum Gravity?
 - Swampland v.s. Landscape
- EFTs in the Landscape subject to universal consistency constraints
 - Swampland Conjectures
 -> general, useful, but **not** fully understood

Stringy Realization

- Quantitative verification of the explicit conjectures
- Manifestations in string geometry
- Refinement

picture from [Palti '19]

Swampland

Quantum Gravity / String Theory

Landscape











The Claim

• The WGC, minimal ver. [Arkani-Hamed, Motl, Nicolis, Vafa '06]

$$S = \int d^4x \sqrt{-G} \left(M_{\rm Pl}^2 R - \frac{1}{4g^2} F^2 + \cdots \right)$$

There must exist a particle with mass "smaller than" charge:
$$\left. \frac{g^2 q^2}{m^2} \right|_{\text{ExtBH}} = \left. \frac{\mu_{\text{ext}}}{M_{\text{Pl}}^2} \right|_{\text{ExtBH}}$$

The Claim

• The WGC, minimal ver. [Arkani-Hamed, Motl, Nicolis, Vafa '06]

$$S = \int d^4x \sqrt{-G} \left(M_{\rm Pl}^2 R - \frac{1}{4g^2} F^2 + \cdots \right) \qquad \text{Weak-Gravity (WG) Particle}$$

There must exist a particle with mass "smaller than" charge:
$$\frac{g^2 q^2}{m^2} \ge \left. \frac{g^2 Q^2}{M^2} \right|_{\rm ExtBH} = \left. \frac{\mu_{\rm ext}}{M_{\rm Pl}^2} \right|_{\rm ExtBH}$$

The Claim

• The WGC, minimal ver. [Arkani-Hamed, Motl, Nicolis, Vafa '06]



• The WGC, tower/sublattice ver.

There must exist a <u>tower</u>/<u>sublattice</u> of such particles with mass scale $m_0 \sim g M_{\rm Pl}$



The Claim

• The WGC, minimal ver. [Arkani-Hamed, Motl, Nicolis, Vafa '06]



• The WGC, tower/sublattice ver.

There must exist a <u>tower</u>/<u>sublattice</u> of such particles with mass scale $m_0 \sim g M_{\rm Pl}$





First Course

Swampland Conjectures & String Theory The Weak Gravity Conjecture(s)

> Second Course Preliminary Evidence (16 Qs) - Heterotic String on a Torus -

> > Stronger Evidence (8 Qs) - 6d F-theory Vacua -

Strongest Evidence (4 Qs) - 4d F-theory Vacua -

..... Dessert

Evidence for Other Conjectures





The Menu

First Course Swampland Conjectures & String Theory The Weak Gravity Conjecture(s)

> Second Course Preliminary Evidence (16 Qs) - Heterotic String on a Torus -

> > Main Course Stronger Evidence (8 Qs) - 6d F-theory Vacua -

> > Strongest Evidence (4 Qs) - 4d F-theory Vacua -

> > > Dessert

Evidence for Other Conjectures



Preliminary Evidence - Heterotic String Theory -

Preliminary Evidence: Upshot

Heterotic String on a Torus

- Heterotic string theory is a **closed-string** theory and reduces at very low energies to 10d SUGRA coupled to 10d SYM w/ G=SO(32) (or $E_8 \times E_8$).
- Such a 10d EFT alone does not serve as a testing ground for the WGC.

• The heterotic string also produces a tower of massive excitations; we can quantize the heterotic string on a flat space, such as $\mathbb{R}^{1,9}$ or $\mathbb{R}^{1,3} \times T^6$.

• The idea is to check if some of those excitations serve as "WG particles", i.e., if they obey $\frac{g^2 q^2}{m^2} \ge \frac{g^2 Q^2}{M^2} \bigg|_{\text{ExtBH}} = \frac{\mu_{\text{ext}}}{M_{\text{Pl}}^2}$; the WGC will be tested this way.

Preliminary Evidence: Upshot

Heterotic String on a Torus

- Heterotic string theory is a closed-string theory and reduces at very low energies to 10d SUGRA coupled to 10d SYM w/ G=SO(32) (or E8 x E8).
- Such a 10d EFT alone does not serve as a testing ground for the WGC.
- The heterotic string also produces a tower of massive excitations; we can quantize the heterotic string on a flat space, such as $\mathbb{R}^{1,9}$ or $\mathbb{R}^{1,3} \times T^6$.
- The idea is to check if some of those excitations serve as "WG particles", i.e., if they obey $\frac{g^2 q^2}{m^2} \ge \left. \frac{g^2 Q^2}{M^2} \right|_{\text{ExtBH}} = \left. \frac{\mu_{\text{ext}}}{M_{\text{Pl}}^2} \right|_{\text{ExtBH}} = \left. \frac{\mu_{\text{ext}}}{M_{\text{Pl}}^2} \right|_{\text{ExtBH}}$

10d Supergravity

- Low-Energy EFT
 - 10d N=1 SUGRA coupled to 10d SYM w/ G=SO(32) (\supset U(1)¹⁶)

$$S = \int_{\mathbb{R}^{1,9}} \frac{M_{\rm Pl}^8}{2} \left(\sqrt{-GR} - \frac{1}{2} \mathrm{d}\phi \wedge *\mathrm{d}\phi\right) - \frac{1}{2g_0^2} e^{-\phi/2} \operatorname{Tr} F \wedge *F$$

The extremal charge-to-mass ratio of dilatonic RN BHs

$$\frac{g^2 Q^2}{M^2} \bigg|_{\text{ExtBH}} = \frac{\mu_{\text{ext}}}{M_{\text{Pl}}^8} \quad \text{where} \quad \mu_{\text{ext}} = \frac{d-3}{d-2} + \frac{\alpha^2}{2} = \frac{7}{8} + \frac{1}{8} = 1$$

10d Supergravity

- Low-Energy EFT
 - 10d N=1 SUGRA coupled to 10d SYM w/ G=SO(32) (\supset U(1)¹⁶)

$$S = \int_{\mathbb{R}^{1,9}} \frac{M_{\rm Pl}^8}{2} \left(\sqrt{-GR} - \frac{1}{2} \mathrm{d}\phi \wedge *\mathrm{d}\phi\right) - \frac{1}{2g_0^2} e^{-\phi/2} \operatorname{Tr} F \wedge *F$$

The extremal charge-to-mass ratio of dilatonic RN BHs

 $\frac{g^2 Q^2}{M^2} \bigg|_{\text{ExtBH}} = \frac{\mu_{\text{ext}}}{M_{\text{Pl}}^8} \quad \text{where} \quad \mu_{\text{ext}} = \frac{d-3}{d-2} + \frac{\alpha^2}{2} = \frac{7}{8} + \frac{1}{8} = 1$

10d Supergravity

- Low-Energy EFT
 - 10d N=1 SUGRA coupled to 10d SYM w/ G=SO(32) (\supset U(1)¹⁶)

$$S = \int_{\mathbb{R}^{1,9}} \frac{M_{\rm Pl}^8}{2} (\sqrt{-GR} - \frac{1}{2} \mathrm{d}\phi \wedge *\mathrm{d}\phi) - \frac{1}{2g_0^2} e^{-\frac{\phi/2}{2}} \mathrm{Tr} \, F \wedge *F$$

• The extremal charge-to-mass ratio of dilatonic RN BHs e.g. [Heidenreich, Reece, Rudelius '15]

$$\frac{g^2 Q^2}{M^2} \Big|_{\text{ExtBH}} = \frac{\mu_{\text{ext}}}{M_{\text{Pl}}^8} \quad \text{where} \quad \mu_{\text{ext}} = \frac{d-3}{d-2} + \frac{\alpha^2}{2} = \frac{7}{8} + \frac{1}{8} = 1$$

10d Supergravity

• Low-Energy EFT

• 10d N=1 SUGRA coupled to 10d SYM w/ G=SO(32) (\supset U(1)¹⁶)

$$S = \int_{\mathbb{R}^{1,9}} \frac{M_{\rm Pl}^8}{2} (\sqrt{-GR} - \frac{1}{2} \mathrm{d}\phi \wedge *\mathrm{d}\phi) - \frac{1}{2g_0^2} e^{-\frac{\phi/2}{2}} \mathrm{Tr} \, F \wedge *F$$

• The extremal charge-to-mass ratio of dilatonic RN BHs e.g. [Heidenreich, Reece, Rudelius '15]

$$\frac{g^2 Q^2}{M^2} \bigg|_{\text{ExtBH}} = \frac{\mu_{\text{ext}}}{M_{\text{Pl}}^8} \quad \text{where} \quad \mu_{\text{ext}} = \frac{d-3}{d-2} + \frac{\alpha^2}{2} = \frac{7}{8} + \frac{1}{8} = 1$$

String Excitations

Mass spectrum:

$$m^2 = 8\pi (n - 1 + \frac{1}{2} |\vec{q}|^2) M_{\text{str}}^2 , \ n \in \mathbb{Z}_{\geq 0}$$

Charge spectrum:

$$\vec{q} = (q_1 + \frac{c}{2}, \dots, q_{16} + \frac{c}{2})$$

with $c = 0$ or $1, q_i \in \mathbb{Z}, \sum_i q_i \in 2\mathbb{Z}$

10d Supergravity

• Low-Energy EFT

• 10d N=1 SUGRA coupled to 10d SYM w/ G=SO(32) (\supset U(1)¹⁶)

$$S = \int_{\mathbb{R}^{1,9}} \frac{M_{\rm Pl}^8}{2} (\sqrt{-G}R - \frac{1}{2} \mathrm{d}\phi \wedge *\mathrm{d}\phi) - \frac{1}{2g_0^2} e^{-\frac{\phi/2}{2}} \mathrm{Tr} \, F \wedge *F$$

The extremal charge-to-mass ratio of dilatonic RN BHs e.g. [Heidenreich, Reece, Rudelius '15]

$$\frac{g^2 Q^2}{M^2} \Big|_{\text{ExtBH}} = \frac{\mu_{\text{ext}}}{M_{\text{Pl}}^8} \quad \text{where} \quad \mu_{\text{ext}} = \frac{d-3}{d-2} + \frac{\alpha^2}{2} = \frac{7}{8} + \frac{1}{8} = 1$$

String Excitations

Mass spectrum:

$$m^2 = 8\pi (n - 1 + \frac{1}{2} |\vec{q}|^2) M_{\text{str}}^2 , \ n \in \mathbb{Z}_{\geq 0}$$

Charge spectrum:

$$\vec{q} = (q_1 + \frac{c}{2}, \dots, q_{16} + \frac{c}{2})$$

with $c = 0$ or $1, q_i \in \mathbb{Z}, \sum_i q_i \in 2\mathbb{Z}$

WG Particles?

Find states in the spectrum with

$$\frac{g^2 q^2}{m^2} \ge \frac{\mu_{\text{ext}}}{M_{\text{Pl}}^8} = \frac{1}{M_{\text{Pl}}^8}$$

• Charge-to-mass ratio (with n=0):

$$\frac{g^2q^2}{m^2} = \frac{1}{M_{\rm Pl}^8} \cdot \frac{|\vec{q}\,|^2}{|\vec{q}\,|^2 - 2} > \frac{1}{M_{\rm Pl}^8}\,, \ \, \forall \vec{q}$$

4d N=4 Supergravity (16 Q's)

• Low-Energy EFT

• 10d N=1 SUGRA coupled to 10d SYM w/ G=SO(32)*

$$S = \int_{\mathbb{R}^{1,9}} \frac{M_{\rm Pl}^8}{2} \left(\sqrt{-G}R - \frac{1}{2} \mathrm{d}\phi \wedge *\mathrm{d}\phi\right) - \frac{1}{2g_0^2} e^{-\phi/2} \operatorname{Tr} F \wedge *F$$

- Reduces to 4d N=4 SUGRA coupled to $16 \text{ U}(1)\text{s}^*$ (w/Wilson lines)
- The extremal charge-to-mass ratio is stable under the reduction [Heidenreich, Reece, Rudelius '15] $\frac{g^2 Q^2}{M^2}\Big|_{\text{ExtBH}} = \frac{\mu_{\text{ext}}}{M_{\text{Pl}}^{d-2}}, \quad \mu_{\text{ext}} = \frac{d-3}{d-2} + \frac{\alpha^2}{2}$
- $M_{
 m Pl}$ and g acquire a common volume factor
- Changes in d and α cancel in μ_{ext} !

WG Particles

• The 10d WG particles in the (sub)lattice w.r.t. the Cartan lead to 4d WG particles

*The same results follow for $E_8 \times E_8$ heterotic string.

4d N=4 Supergravity (16 Q's)

• Low-Energy EFT

• 10d N=1 SUGRA coupled to 10d SYM w/ G=SO(32)*

$$S = \int_{\mathbb{R}^{1,9}} \frac{M_{\rm Pl}^8}{2} (\sqrt{-G}R - \frac{1}{2} \mathrm{d}\phi \wedge * \mathrm{d}\phi) - \frac{1}{2g_0^2} e^{-\phi/2} \operatorname{Tr} F \wedge *F$$

- Reduces to 4d N=4 SUGRA coupled to $16 \text{ U}(1)\text{s}^*$ (w/Wilson lines)
- The extremal charge-to-mass ratio is stable under the reduction [Heidenreich, Reece, Rudelius '15]

$$\frac{g^2 Q^2}{M^2} \Big|_{\text{ExtBH}} = \frac{\mu_{\text{ext}}}{M_{\text{Pl}}^{d-2}}, \quad \mu_{\text{ext}} = \frac{d-3}{d-2} + \frac{\alpha^2}{2}$$

- $M_{
 m Pl}$ and g acquire a common volume factor
- Changes in d and $\,lpha$ cancel in $\,\mu_{
 m ext}!$

WG Particles

• The 10d WG particles in the (sub)lattice w.r.t. the Cartan lead to 4d WG particles

*The same results follow for $E_8 \times E_8$ heterotic string.

4d N=4 Supergravity (16 Q's)

• Low-Energy EFT

• 10d N=1 SUGRA coupled to 10d SYM w/ G=SO(32)*

$$S = \int_{\mathbb{R}^{1,9}} \frac{M_{\rm Pl}^8}{2} (\sqrt{-G}R - \frac{1}{2} \mathrm{d}\phi \wedge * \mathrm{d}\phi) - \frac{1}{2g_0^2} e^{-\phi/2} \operatorname{Tr} F \wedge *F$$

- Reduces to 4d N=4 SUGRA coupled to $16 U(1)s^*$ (w/Wilson lines)
- The extremal charge-to-mass ratio is stable under the reduction [Heidenreich, Reece, Rudelius '15]

$$\frac{g^2 Q^2}{M^2} \bigg|_{\text{ExtBH}} = \frac{\mu_{\text{ext}}}{M_{\text{Pl}}^{d-2}}, \quad \mu_{\text{ext}} = \frac{d-3}{d-2} + \frac{\alpha^2}{2}$$

- $M_{\rm Pl}$ and g acquire a common volume factor
- Changes in d and α cancel in $\mu_{\mathrm{ext}}!$

• WG Particles

• The 10d WG particles in the (sub)lattice w.r.t. the Cartan lead to 4d WG particles

*The same results follow for $E_8 \times E_8$ heterotic string.

4d N=4 Supergravity (16 Q's)

• Low-Energy EFT

• 10d N=1 SUGRA coupled to 10d SYM w/ G=SO(32)*

$$S = \int_{\mathbb{R}^{1,9}} \frac{M_{\rm Pl}^8}{2} (\sqrt{-G}R - \frac{1}{2} \mathrm{d}\phi \wedge * \mathrm{d}\phi) - \frac{1}{2g_0^2} e^{-\phi/2} \operatorname{Tr} F \wedge *F$$

- Reduces to 4d N=4 SUGRA coupled to $16 U(1)s^*$ (w/Wilson lines)
- The extremal charge-to-mass ratio is stable under the reduction [Heidenreich, Reece, Rudelius '15]

$$\frac{g^2 Q^2}{M^2} \bigg|_{\text{ExtBH}} = \frac{\mu_{\text{ext}}}{M_{\text{Pl}}^{d-2}}, \quad \mu_{\text{ext}} = \frac{d-3}{d-2} + \frac{\alpha^2}{2}$$

- $M_{\rm Pl}$ and g acquire a common volume factor
- Changes in d and $\,\alpha$ cancel in $\,\mu_{\rm ext}!$

WG Particles

• The 10d WG particles in the (sub)lattice w.r.t. the Cartan lead to 4d WG particles

*The same results follow for $E_8 \times E_8$ heterotic string.



The Menu

First Course Swampland Conjectures & String Theory The Weak Gravity Conjecture(s) Second Course Preliminary Evidence (16 Qs) - Heterotic String on a Torus -

> Main Course Stronger Evidence (8 Qs) - 6d F-theory Vacua -

Strongest Evidence (4 Qs) - 4d F-theory Vacua -

Dessert

Evidence for Other Conjectures



Stronger Evidence - 6d F-theory -

6d F-theory

 F-theory is another name for Type IIB string theory with 7-branes and hence, with a non-trivial dilation profile along the internal compact directions.

• E.g. 6d F-theory is the 10d IIB string theory put on a compact (curved) 4manifold; it preserves 8 Q's (to be contrasted with the heterotic on a torus preserving 16 Q's).

The 7-branes carry gauge algebras and charged particles arise from the open string ending on them with a fixed "charge factor". It may look impossible to fulfill the sublattice WGC.

6d F-theory

 F-theory is another name for Type IIB string theory with 7-branes and hence, with a non-trivial dilation profile along the internal compact directions.

• E.g. 6d F-theory is the 10d IIB string theory put on a compact (curved) 4manifold; it preserves 8 Q's (to be contrasted with the heterotic on a torus preserving 16 Q's).

 The 7-branes carry gauge algebras and charged particles arise from the open string ending on them with a fixed "charge factor". It may look impossible to fulfill the sublattice WGC.

6d F-theory

 The day is saved by another ingredient of F-theory — the 3-brane, which can wrap an internal 2-cycle to produce an effective string in the EFT.

• A <u>special</u> 2-cycle can be identified so that the resulting string is the (dual) heterotic string; its excitations are captured by the (index-ver.) partition function, which we can control though the internal space is <u>not</u> flat!

 At weak gauge coupling, the heterotic string is also weakly coupled and its mass spectrum can be reliably computed. One can thus explicitly check if the excitations are WG particles.

6d F-theory

The day is saved by another ingredient of F-theory — the 3-brane, which can wrap an internal 2-cycle to produce an effective string in the EFT.

 A <u>special</u> 2-cycle can be identified so that the resulting string is the (dual) heterotic string; its excitations are captured by the (index-ver.) partition function, which we can control though the internal space is <u>not</u> flat!

 At weak gauge coupling, the heterotic string is also weakly coupled and its mass spectrum can be reliably computed. One can thus explicitly check if the excitations are WG particles.

6d F-theory

 The day is saved by another ingredient of F-theory — the 3-brane, which can wrap an internal 2-cycle to produce an effective string in the EFT.

 A <u>special</u> 2-cycle can be identified so that the resulting string is the (dual) heterotic string; its excitations are captured by the (index-ver.) partition function, which we can control though the internal space is <u>not</u> flat!

 At weak gauge coupling, the heterotic string is also weakly coupled and its mass spectrum can be reliably computed. One can thus explicitly check if the excitations are WG particles.

F-theory in 6 Dimensions

Couplings via Kahler Moduli

• 6d F-theory

IIB string theory on a compact 4-manifold X4

with 7-branes on an internal 2-dim'l surface S_2

- external 6-dim'l gauge fields
- a non-trivial dilaton profile
- Physics via Geometry
 - We will look into the "Kahler moduli" of X_4
 - Govern the cycle volumes and in turn, the couplings
 - gravity: $\left(M_{\rm Pl}/M_{\rm IIB}
 ight)^4 = 4\pi\mathcal{V}_{X\!4}$
 - gauge: $(1/g^2)/M_{\rm IIB}^2 = (2\pi)^{-1}\mathcal{V}_{\rm S2}$
 - Weak gauge coupling (w/ gravity fixed): $V_{S_2} \rightarrow \infty (V_{X_4} \sim 1)$


Couplings via Kahler Moduli

- 6d F-theory
 - IIB string theory on a compact 4-manifold X_4

with 7-branes on an internal 2-dim'l surface S_2

- external 6-dim'l gauge fields
- a non-trivial dilaton profile (described by an elliptic fibration : $Y_6 \longrightarrow X_4$)
- Physics via Geometry
 - We will look into the "Kahler moduli" of X_4
 - Govern the cycle volumes and in turn, the couplings
 - gravity: $\left(M_{\rm Pl}/M_{\rm IIB}
 ight)^4 = 4\pi \mathcal{V}_{X4}$
 - gauge: $(1/g^2)/M_{\rm IIB}^2 = (2\pi)^{-1}\mathcal{V}_{\rm S2}$
 - Weak gauge coupling (w/ gravity fixed): $V_{S_2} \rightarrow \infty (V_{X_4} \sim 1)$



Couplings via Kahler Moduli

• 6d F-theory

IIB string theory on a compact 4-manifold X4

with 7-branes on an internal 2-dim'l surface S_2

- external 6-dim'l gauge fields
- a non-trivial dilaton profile (described by an elliptic fibration: $Y_6 \longrightarrow X_4$)
- Physics via Geometry
 - We will look into the "Kahler moduli" of X_4
 - Govern the cycle volumes and in turn, the couplings
 - gravity: $\left(M_{\rm Pl}/M_{\rm IIB}
 ight)^4 = 4\pi \mathcal{V}_{X\!4}$
 - gauge: $(1/g^2)/M_{\rm IIB}^2 = (2\pi)^{-1} \mathcal{V}_{\rm S2}$
 - Weak gauge coupling (w/ gravity fixed): $V_{S_2} \rightarrow \infty (V_{X_4} \sim 1)$



Couplings via Kahler Moduli

• 6d F-theory

IIB string theory on a compact 4-manifold X4

with 7-branes on an internal 2-dim'l surface S_2

- external 6-dim'l gauge fields
- a non-trivial dilaton profile (described by an elliptic fibration : $Y_6 \longrightarrow X_4$)

- We will look into the "Kahler moduli" of X_4
- Govern the cycle volumes and in turn, the couplings
- gravity: $\left(M_{\rm Pl}/M_{\rm IIB}
 ight)^4 = 4\pi \mathcal{V}_{X\!4}$
- gauge: $(1/g^2)/M_{\rm IIB}^2 = (2\pi)^{-1} \mathcal{V}_{\rm S2}$
- Weak gauge coupling (w/ gravity fixed): $\mathcal{V}_{S_2} \rightarrow \infty \quad (\mathcal{V}_{X_4} \sim 1)$



Couplings via Kahler Moduli

• 6d F-theory

IIB string theory on a compact 4-manifold X4

with 7-branes on an internal 2-dim'l surface S_2

- external 6-dim'l gauge fields
- a non-trivial dilaton profile (described by an elliptic fibration : $Y_6 \longrightarrow X_4$)

- We will look into the "Kahler moduli" of X_4
- Govern the cycle volumes and in turn, the couplings
- gravity: $\left(M_{\rm Pl}/M_{\rm IIB}
 ight)^4 = 4\pi \mathcal{V}_{X\!4}$
- gauge: $(1/g^2)/M_{\rm IIB}^2 = (2\pi)^{-1} \mathcal{V}_{\rm S2}$
- Weak gauge coupling (w/ gravity fixed): $\mathcal{V}_{S_2} \rightarrow \infty \ (\mathcal{V}_{X_4} \sim 1)$



Couplings via Kahler Moduli

• 6d F-theory

IIB string theory on a compact 4-manifold X4

with 7-branes on an internal 2-dim'l surface S_2

- external 6-dim'l gauge fields
- a non-trivial dilaton profile (described by an elliptic fibration : $Y_6 \longrightarrow X_4$)

- We will look into the "Kahler moduli" of X_4
- Govern the cycle volumes and in turn, the couplings
- gravity: $\left(M_{\rm Pl}/M_{\rm IIB}
 ight)^4 = 4\pi \mathcal{V}_{X\!4}$
- gauge: $(1/g^2)/M_{\rm IIB}^2 = (2\pi)^{-1}\mathcal{V}_{\rm S2}$
- Weak gauge coupling (w/ gravity fixed): $\mathcal{V}_{S_2} \rightarrow \infty \ (\mathcal{V}_{X_4} \sim 1) \implies$ exists C_2 with $\mathcal{V}_{C_2} \rightarrow 0$



Couplings via Kahler Moduli

• 6d F-theory

IIB string theory on a compact 4-manifold X4

with 7-branes on an internal 2-dim'l surface S_2

- external 6-dim'l gauge fields
- a non-trivial dilaton profile (described by an elliptic fibration : $Y_6 \longrightarrow X_4$)

- We will look into the "Kahler moduli" of X_4
- Govern the cycle volumes and in turn, the couplings
- gravity: $\left(M_{\rm Pl}/M_{\rm IIB}
 ight)^4 = 4\pi\mathcal{V}_{X\!4}$
- gauge: $(1/g^2)/M_{\rm IIB}^2 = (2\pi)^{-1} \mathcal{V}_{\rm S2}$
- Weak gauge coupling (w/ gravity fixed): $\mathcal{V}_{S_2} \rightarrow \infty \ (\mathcal{V}_{X_4} \sim 1) \implies$ exists C_2 with $\mathcal{V}_{C_2} \rightarrow 0$





D3 Brane on Shrinking 2-cycle [S.-J.L., Lerche, Weigand, '18]

• Geometry

• Dialing $g \rightarrow 0$ deforms X_4 such that a certain 2-cycle C_2 therein shrinks, where C_2 is

(a) the unique shrinking cycle as a class*

(b) a 2-sphere

(C) a fiber

* To be precise, other 2-cycles may also shrink, but only at finite distance.



• Physics

• D3-brane on C2 leads to heterotic string with a vanishing tension

$$\frac{M_{\rm str}^2}{M_{\rm Pl}^2} \sim \frac{\mathcal{V}_{C_2}}{\mathcal{V}_{X_4}^{1/2}} \to 0$$

D3 Brane on Shrinking 2-cycle [S.-J.L., Lerche, Weigand, '18]

• Geometry

• Dialing $g \rightarrow 0$ deforms X_4 such that a certain 2-cycle C_2 therein shrinks, where C_2 is

(a) the unique shrinking cycle as a class*

(b) a 2-sphere

(C) a fiber

* To be precise, other 2-cycles may also shrink, but only at finite distance.



• D3-brane on C_2 leads to heterotic string with a vanishing tension

$$\frac{M_{\rm str}^2}{M_{\rm Pl}^2} \sim \frac{\mathcal{V}_{C_2}}{\mathcal{V}_{X_4}^{1/2}} \to 0$$



D3 Brane on Shrinking 2-cycle [S.-J.L., Lerche, Weigand, '18]

• Geometry

• Dialing $g \rightarrow 0$ deforms X_4 such that a certain 2-cycle C_2 therein shrinks, where C_2 is

(a) the unique shrinking cycle as a class*

(b) a 2-sphere

(C) a fiber

* To be precise, other 2-cycles may also shrink, but only at finite distance.



• Physics

• D3-brane on C_2 leads to heterotic string with a vanishing tension

$$\frac{M_{\rm str}^2}{M_{\rm Pl}^2} \sim \frac{\mathcal{V}_{C_2}}{\mathcal{V}_{X_4}^{1/2}} \to 0$$

D3 Brane on Shrinking 2-cycle [S.-J.L., Lerche, Weigand, '18]

• Geometry

• Dialing $g \rightarrow 0$ deforms X_4 such that a certain 2-cycle C_2 therein shrinks, where C_2 is

(a) the unique shrinking cycle as a class*

(b) a 2-sphere

(C) a fiber

* To be precise, other 2-cycles may also shrink, but only at finite distance.



Physics

• D3-brane on C_2 leads to heterotic string with a vanishing tension



D3 Brane on Shrinking 2-cycle [S.-J.L., Lerche, Weigand, '18]

• Geometry

• Dialing $g \rightarrow 0$ deforms X_4 such that a certain 2-cycle C_2 therein shrinks, where C_2 is

(a) the unique shrinking cycle as a class*

(b) a 2-sphere

(C) a fiber

* To be precise, other 2-cycles may also shrink, but only at finite distance.



• Physics

• D3-brane on C_2 leads to heterotic string with a vanishing tension

$$\frac{M_{\rm str}^2}{M_{\rm Pl}^2} \sim \frac{\mathcal{V}_{C_2}}{\mathcal{V}_{X_4}^{1/2}} \to 0$$

D3 Brane on Shrinking 2-cycle [S.-J.L., Lerche, Weigand, '18]

Geometry

• Dialing $g \to 0$ deforms X_4 such that a certain 2-cycle C_2 therein shrinks, where C_2 is

(a) the unique shrinking cycle as a class*

(b) a 2-sphere

(C) a fiber

* To be precise, other 2-cycles may also shrink, but only at finite distance.



Physics

• D3-brane on C_2 leads to heterotic string with a vanishing tension

$$rac{M_{
m str}^2}{M_{
m Pl}^2}\sim rac{\mathcal{V}_{C_2}}{\mathcal{V}_{X_4}^{1/2}}
ightarrow 0$$
 Light tower of strin

g excitations

D3 Brane on Shrinking 2-cycle [S.-J.L., Lerche, Weigand, '18]

• Geometry

• Dialing $g \rightarrow 0$ deforms X_4 such that a certain 2-cycle C_2 therein shrinks, where C_2 is

(a) the unique shrinking cycle as a class*

(b) a 2-sphere

(C) a fiber

* To be precise, other 2-cycles may also shrink, but only at finite distance.



• Physics

• D3-brane on C_2 leads to heterotic string with a vanishing tension

$$rac{M_{
m str}^2}{M_{
m Pl}^2} \sim rac{\mathcal{V}_{C_2}}{\mathcal{V}_{X_4}^{1/2}} \to 0$$
 Light tower of string excitations

D3 Brane on Shrinking 2-cycle [S.-J.L., Lerche, Weigand, '18]

• Geometry

• Dialing $g \rightarrow 0$ deforms X_4 such that a certain 2-cycle C_2 therein shrinks, where C_2 is

(a) the unique shrinking cycle as a class*

(b) a 2-sphere

(C) a fiber

* To be precise, other 2-cycles may also shrink, but only at finite distance.



• Physics

• D3-brane on C_2 leads to heterotic string with a vanishing tension



The WGC at Weak Gauge Coupling

Main Ideas

- The 6d Effective Action [S.-J.L., Lerche, Weigand, '18]
 - Asymptotically turns into a dilatonic Einstein-Maxwell theory w/ dil. coupling $\alpha = 1$

$$S = \int_{\mathbb{R}^{1,5}} \frac{M_{\rm Pl}^4}{2} (\sqrt{-GR} - \mathrm{d}x \wedge *\mathrm{d}x) - \frac{1}{2g_0^2} e^{\alpha x} F \wedge *F + \cdots$$

The extremal charge-to-mass ratio of dilatonic RN BHs

$$\frac{g^2 Q^2}{M^2} \Big|_{\text{ExtBH}} = \frac{\mu_{\text{ext}}}{M_{\text{Pl}}^4} \quad \text{where} \quad \mu_{\text{ext}} = \frac{d-3}{d-2} + \frac{\alpha^2}{4} = \frac{3}{4} + \frac{1}{4} = 1$$

The WGC at Weak Gauge Coupling

Main Ideas

- The 6d Effective Action [S.-J.L., Lerche, Weigand, '18]
 - Asymptotically turns into a dilatonic Einstein-Maxwell theory w/ dil. coupling $\alpha = 1$

$$S = \int_{\mathbb{R}^{1,5}} \frac{M_{\rm Pl}^4}{2} (\sqrt{-GR} - \mathrm{d}x \wedge *\mathrm{d}x) - \frac{1}{2g_0^2} e^{\alpha x} F \wedge *F + \cdots$$

The extremal charge-to-mass ratio of dilatonic RN BHs

 $\frac{g^2 Q^2}{M^2} \Big|_{\text{ExtBH}} = \frac{\mu_{\text{ext}}}{M_{\text{Pl}}^4} \quad \text{where} \quad \mu_{\text{ext}} = \frac{d-3}{d-2} + \frac{\alpha^2}{4} = \frac{3}{4} + \frac{1}{4} = 1$

String Excitations

• Mass spectrum of the light tower:

 $m^2 = 8\pi(n-1)M_{\rm str}^2$

• Charge spectrum at each level *n*:

 $q \in [0, q^{\max}(n)]$

The WGC at Weak Gauge Coupling

Main Ideas

- The 6d Effective Action [S.-J.L., Lerche, Weigand, '18]
 - Asymptotically turns into a dilatonic Einstein-Maxwell theory w/ dil. coupling $\alpha=1$

$$S = \int_{\mathbb{R}^{1,5}} \frac{M_{\rm Pl}^4}{2} (\sqrt{-GR} - \mathrm{d}x \wedge *\mathrm{d}x) - \frac{1}{2g_0^2} e^{\alpha x} F \wedge *F + \cdots$$

The extremal charge-to-mass ratio of dilatonic RN BHs

$$\frac{g^2 Q^2}{M^2} \Big|_{\text{ExtBH}} = \frac{\mu_{\text{ext}}}{M_{\text{Pl}}^4} \quad \text{where} \quad \mu_{\text{ext}} = \frac{d-3}{d-2} + \frac{\alpha^2}{4} = \frac{3}{4} + \frac{1}{4} = 1$$

String Excitations

• Mass spectrum of the light tower:

 $m^2 = 8\pi(n-1)M_{\rm str}^2$

Charge spectrum at each level n:
 q ∈ [0, q^{max}(n)]

WG Particles?

- Find $(n, q^{\max}(n))$ in the spectrum w/

$$\frac{g^2 q^{\max}(n)^2}{m(n)^2} \ge \frac{\mu_{\text{ext}}}{M_{\text{Pl}}^4} = \frac{1}{M_{\text{Pl}}^4}$$

- Test if such $q^{\max}(n)$ fill a sublattice

Universal Feature of the Excitation Spectrum

• The Spectrum

- Info needed of the maximal charge q^{\max} at a given mass level n
- Captured by the "partition function" (or its index version) of the 6d heterotic string: $Z(\nu, \eta) = -\nu^{-1} \sum N(n, q) \nu^n \eta^q$

where N(n, q) = "Degeneracy at excitation mass level n and charge q"

- A Distinguished Sublattice [S.-J.L., Lerche, Weigand '18]
 - A general property of $\mathcal{Z}(\nu,\eta)$: "quasi-modularity"

 $\dots N(n,q)$ non-zero for charges up to $q_k^{\max} = 2rk \ (n_k = rk^2)$ for $k \in \mathbb{Z}$

- This charge sublattice (of index 2r) is filled by physical states
- The states lying in this sublattice turn out to be WG particles!

Universal Feature of the Excitation Spectrum

• The Spectrum

- Info needed of the maximal charge q^{\max} at a given mass level n
- Captured by the "partition function" (or its index version) of the 6d heterotic string: $Z(\nu, \eta) = -\nu^{-1} \sum N(n, q) \nu^n \eta^q$

where N(n, q) = "Degeneracy at excitation mass level *n* and charge *q*"

• A Distinguished Sublattice [S.-J.L., Lerche, Weigand '18]

- A general property of $\mathcal{Z}(\boldsymbol{\nu},\eta)$: "quasi-modularity"
 - $\dots \rightarrow N(n,q)$ non-zero for charges up to $q_k^{\max} = 2rk \ (n_k = rk^2)$ for $k \in \mathbb{Z}$
- This charge sublattice (of index 2r) is filled by physical states
- The states lying in this sublattice turn out to be WG particles!

Universal Feature of the Excitation Spectrum

• The Spectrum

- Info needed of the maximal charge q^{\max} at a given mass level n
- Captured by the "partition function" (or its index version) of the 6d heterotic string: $Z(\nu, \eta) = -\nu^{-1} \sum N(n, q) \nu^n \eta^q$

where N(n, q) = "Degeneracy at excitation mass level *n* and charge *q*"

• A Distinguished Sublattice [S.-J.L., Lerche, Weigand '18] $r = \frac{1}{2}C_2 \cdot \mathbf{S}_2$

- A general property of $\mathcal{Z}(\mathbf{\nu}, \eta)$: "quasi-modularity"
 - N(n,q) non-zero for charges up to $(q_k^{\max} = 2rk)(n_k = rk^2)$ for $k \in \mathbb{Z}$
- This charge sublattice (of index 2r) is filled by physical states
- The states lying in this sublattice turn out to be WG particles!

Universal Feature of the Excitation Spectrum

• The Spectrum

- Info needed of the maximal charge q^{\max} at a given mass level n
- Captured by the "partition function" (or its index version) of the 6d heterotic string: $Z(\nu, \eta) = -\nu^{-1} \sum N(n, q) \nu^n \eta^q$

where N(n, q) = "Degeneracy at excitation mass level n and charge q"

• A Distinguished Sublattice [S.-J.L., Lerche, Weigand '18]

$$r = \frac{1}{2}C_2 \cdot \mathbf{S}_2$$

- A general property of $\mathcal{Z}(\nu, \eta)$: "quasi-modularity"
 - N(n,q) non-zero for charges up to $q_k^{\max} = 2rk(n_k = rk^2)$ for $k \in \mathbb{Z}$
- This charge sublattice (of index 2r) is filled by physical states
- The states lying in this sublattice turn out to be WG particles!

Universal Feature of the Excitation Spectrum

• The Spectrum

- Info needed of the maximal charge q^{\max} at a given mass level n
- Captured by the "partition function" (or its index version) of the 6d heterotic string: $Z(\nu, \eta) = -\nu^{-1} \sum N(n, q) \nu^n \eta^q$

where N(n, q) = "Degeneracy at excitation mass level *n* and charge *q*"

• A Distinguished Sublattice [S.-J.L., Lerche, Weigand '18]

$$r = \frac{1}{2}C_2 \cdot \mathbf{S}_2$$

• A general property of $\mathcal{Z}(\boldsymbol{\nu},\eta)$: "quasi-modularity"

N(n,q) non-zero for charges up to $q_k^{\max} = 2rk(n_k = rk^2)$ for $k \in \mathbb{Z}$

- This charge sublattice (of index 2r) is filled by physical states
- The states lying in this sublattice turn out to be WG particles!

Verification for 6d F-theory [S.-J.L., Lerche, Weigand '18]

• The Distinguished Sublattice

Quasi-Jacobi nature of the heterotic worldsheet partition function guarantees that

 $q_k^{\max} = 2rk \ (n_k = rk^2) \text{ for } k \in \mathbb{Z}$ are filled by physical states, with the index $2r = C_2 \cdot \mathbf{S}_2$



$$q_k^{\max} = \sqrt{4 r n_k} \ge \sqrt{4 r (n_k - 1)}$$

(an example w/ sublattice index 2r=4)

Verification for 6d F-theory [S.-J.L., Lerche, Weigand '18]

• The Distinguished Sublattice

Quasi-Jacobi nature of the heterotic worldsheet partition function guarantees that



Verification for 6d F-theory [S.-J.L., Lerche, Weigand '18]

• The Distinguished Sublattice

Quasi-Jacobi nature of the heterotic worldsheet partition function guarantees that



Verification for 6d F-theory [S.-J.L., Lerche, Weigand '18]

• The Distinguished Sublattice

Quasi-Jacobi nature of the heterotic worldsheet partition function guarantees that



Verification for 6d F-theory [S.-J.L., Lerche, Weigand '18]

• The Distinguished Sublattice

Quasi-Jacobi nature of the heterotic worldsheet partition function guarantees that



Verification for 6d F-theory [S.-J.L., Lerche, Weigand '18]

• The Distinguished Sublattice

Quasi-Jacobi nature of the heterotic worldsheet partition function guarantees that



Verification for 6d F-theory [S.-J.L., Lerche, Weigand '18]

• The Distinguished Sublattice

Quasi-Jacobi nature of the heterotic worldsheet partition function guarantees that



Strongest Evidence - 4d F-theory -

4d F-theory

 4d F-theory is the 10d IIB string theory put on a compact (curved) 6manifold, it preserves 4 Q's (to be contrasted with the 6d F-theory, which preserves 8 Q's).

 Once again, a special internal 2-cycle can be identified so that the 3-brane wrapping it produces an effective heterotic string in the 4d EFT.

• With the SUSY reduced, its (index-ver.) partition function is much less controlled than in the 6d F-theory case, but we still have some control.

4d F-theory

 4d F-theory is the 10d IIB string theory put on a compact (curved) 6manifold, it preserves 4 Q's (to be contrasted with the 6d F-theory, which preserves 8 Q's).

 Once again, a special internal 2-cycle can be identified so that the 3-brane wrapping it produces an effective heterotic string in the 4d EFT.

• With the SUSY reduced, its (index-ver.) partition function is much less controlled than in the 6d F-theory case, but we still have some control.

4d F-theory

 4d F-theory is the 10d IIB string theory put on a compact (curved) 6manifold, it preserves 4 Q's (to be contrasted with the 6d F-theory, which preserves 8 Q's).

 Once again, a special internal 2-cycle can be identified so that the 3-brane wrapping it produces an effective heterotic string in the 4d EFT.

• With the SUSY reduced, its (index-ver.) partition function is much less controlled than in the 6d F-theory case, but we still have some control.

4d F-theory

 Technically, the heterotic partition functions in 6d and 4d F-theory are, respectively, *quasi-modular* and *quasi-Jacobi*; the latter is much less
 constraining the partition function (charges for each mass level).

But the excitation mass is again reliably computed at weak gauge coupling, as the emergent heterotic string in 4d is weakly coupled as in 6d.

 Having seen in 6d that the quasi-modular property distinguishes a charge sublattice for the sublattice WGC, we can test if even the reduced control in 4d from the quasi-Jacobi nature suffices.

4d F-theory

 Technically, the heterotic partition functions in 6d and 4d F-theory are, respectively, *quasi-modular* and *quasi-Jacobi*; the latter is much less
 constraining the partition function (charges for each mass level).

 But the excitation mass is again reliably computed at weak gauge coupling, as the emergent heterotic string in 4d is weakly coupled as in 6d.

Having seen in 6d that the *quasi-modular* property distinguishes a charge sublattice for the sublattice WGC, we can test if even the reduced control in 4d from the *quasi-Jacobi* nature suffices.

Couplings via Kahler Moduli

• 6d F-theory

• IIB string theory on a compact 4-manifold X_4

with 7-branes on an internal 2-dim'l cycle S_2

- external 6-dim'l gauge fields
- a non-trivial dilaton profile (described by an elliptic fibration : $Y_6 \longrightarrow X_4$)

- We will look into the "Kahler moduli" of X_4
- Govern the cycle volumes and in turn, the couplings
- gravity: $\left(M_{\rm Pl}/M_{\rm IIB}
 ight)^4 = 4\pi \mathcal{V}_{X\!4}$
- gauge: $(1/g^2)/M_{\rm IIB}^2 = (2\pi)^{-1} \mathcal{V}_{\rm S2}$
- Weak gauge coupling (w/ gravity fixed): $\mathcal{V}_{S_2} \rightarrow \infty \ (\mathcal{V}_{X_4} \sim 1) \implies$ exists C_2 with $\mathcal{V}_{C_2} \rightarrow 0$




F-theory in 4 Dimensions

Couplings via Kahler Moduli

• 4d F-theory

• IIB string theory on a compact 6-manifold X_6

with 7-branes on an internal 4-dim'l cycle S_4

- external 4-dim'l gauge fields
- a non-trivial dilaton profile (described by an elliptic fibration : $\gamma_8 \longrightarrow \chi_6$)

• Physics via Geometry

- We will look into the "Kahler moduli" of X_6
- Govern the cycle volumes and in turn, the couplings
- gravity: $\left(M_{\rm Pl}/M_{\rm IIB}\right)^2 = 4\pi \mathcal{V}_{X_6}$
- gauge: $(1/g^2) = (2\pi)^{-1} \mathcal{V}_{S4}$
- Weak gauge coupling (w/ gravity fixed): $\mathcal{V}_{S_4} \rightarrow \infty \ (\mathcal{V}_{X_6} \sim 1) \implies$ exists C_2 with $\mathcal{V}_{C_2} \rightarrow 0$





The WGC at Weak Gauge Coupling

Main Ideas: Same as in 6 Dimensions

- The 4d Effective Action [S.-J.L., Lerche, Weigand '19]
 - Asymptotically turns into a dilatonic Einstein-Maxwell theory w/ dil. coupling $lpha=\sqrt{2}$

$$S = \int_{\mathbb{R}^{1,3}} \frac{M_{\mathrm{Pl}}^2}{2} (\sqrt{-GR} - \mathrm{d}x \wedge *\mathrm{d}x) - \frac{1}{2g_0^2} e^{\alpha x} F \wedge *F + \cdots$$

The extremal charge-to-mass ratio of dilatonic RN BHs

$$\frac{g^2 Q^2}{M^2} \Big|_{\text{ExtBH}} = \frac{\mu_{\text{ext}}}{M_{\text{Pl}}^2} \quad \text{where} \quad \mu_{\text{ext}} = \frac{d-3}{d-2} + \frac{\alpha^2}{4} = \frac{1}{2} + \frac{1}{2} = 1$$

String Excitations

• Mass spectrum of the light tower:

 $m^2 = 8\pi(n-1)M_{\rm str}^2$

Charge spectrum at each level n:
 q ∈ [0, q^{max}(n)]

WG Particles?

- Find $(n, q^{\max}(n))$ in the spectrum w/

$$\frac{g^2 q^{\max}(n)^2}{m(n)^2} \ge \frac{\mu_{\text{ext}}}{M_{\text{Pl}}^2} = \frac{1}{M_{\text{Pl}}^2}$$

- Test if such $q^{\max}(n)$ fill a sublattice

The Sublattice WGC

Verification for 4d F-theory [S.-J.L., Lerche, Weigand '19] [Klawer, S.-J.L., Weigand, Wiesner '20]

• The Distinguished Sublattice

Quasi-Jacobi nature of the heterotic worldsheet partition function guarantees that

 $q_k^{\max} = 2rk \ (n_k = rk^2)$ for $k \in \mathbb{Z}$ are filled by physical states, with the index $2r = C_2 \cdot \mathbf{S}_4$



Less Protection with Less SUSY [Klawer, S.-J.L., Weigand, Wiesner '20]

• Origin of Corrections

- Shrinking cycles in principle lead to corrections to the classical geometric calculations
- 4d F-theory w/ 4 Q's does suffer from quantum corrections (unlike in 6d w/ 8 Q's)

• Fate of the "WG Inequality"

- The sublattice is filled by WG particles if $\mu_{str} := 2r \frac{\mathcal{V}_{X_6}}{\mathcal{V}_{S_*}\mathcal{V}_{C_5}} \stackrel{\textbf{?}}{=} 1$ where $2r = C_2 \cdot \mathbf{S}_4$
- Classical $\mu_{\text{str}}^{\text{cl}} = 1$ in other words: $\mathcal{V}_{\mathbf{S}_4}^{\text{cl}} \mathcal{V}_{C_2}^{\text{cl}} \stackrel{\mathbf{l}}{=} (C_2 \cdot \mathbf{S}_4) \mathcal{V}_{X_6}^{\text{cl}}$
- Quantum $\mu_{
 m str} = 1 \Delta$

leading corrections to the volumes cancel in the ratio, i.e., $\,\Delta o 0$ at strict weak gauge coupling

$$\frac{q^2 q^2}{m^2} \stackrel{?}{\geq} \frac{1-\Delta}{M_{\rm Pl}^2} \quad \begin{array}{l} \text{Mass renormalization ansatz: } \frac{m^2}{M_{\rm Pl}^2} = 8\pi(n-1)\frac{M_{\rm str}^2}{M_{\rm Pl}^2}(1+\rho) \\ \text{If the sublattice still provides WG particles: } \rho = -\frac{1}{2}\Delta \end{array} \quad \begin{array}{l} \underbrace{\frac{g^2 q^2}{m^2} \geq \frac{1-\frac{1}{2}\Delta}{M_{\rm Pl}^2}}{m^2} \\ \end{array}$$

Less Protection with Less SUSY [Klawer, S.-J.L., Weigand, Wiesner '20]

• Origin of Corrections

- Shrinking cycles in principle lead to corrections to the classical geometric calculations
- 4d F-theory w/ 4 Q's does suffer from quantum corrections (unlike in 6d w/ 8 Q's)

• Fate of the "WG Inequality"

- The sublattice is filled by WG particles if $\mu_{str} := 2r \frac{V_{X_6}}{V_{S} V_{C_2}} = 1$ where $2r = C_2 \cdot S_4$
- Classical $\mu_{\text{str}}^{\text{cl}} = 1$ in other words: $\mathcal{V}_{\mathbf{S}_4}^{\text{cl}} \mathcal{V}_{C_2}^{\text{cl}} \stackrel{!}{=} (C_2 \cdot \mathbf{S}_4) \mathcal{V}_{X_6}^{\text{cl}}$
- Quantum $\mu_{\rm str} = 1 \Delta$

leading corrections to the volumes cancel in the ratio, i.e., $\Delta o 0$ at strict weak gauge coupling

Less Protection with Less SUSY [Klawer, S.-J.L., Weigand, Wiesner '20]

• Origin of Corrections

- Shrinking cycles in principle lead to corrections to the classical geometric calculations
- 4d F-theory w/ 4 Q's does suffer from quantum corrections (unlike in 6d w/ 8 Q's)

• Fate of the "WG Inequality"

- The sublattice is filled by WG particles if $\mu_{\text{str}} := 2r \frac{\mathcal{V}_{X_6}}{\mathcal{V}_{\mathbf{S}_4} \mathcal{V}_{C_2}} = 1$ where $2r = C_2 \cdot \mathbf{S}_4$
- Classical $\mu_{\text{str}}^{\text{cl}} = 1$ in other words: $\mathcal{V}_{\mathbf{S}_4}^{\text{cl}} \mathcal{V}_{C_2}^{\text{cl}} \stackrel{!}{=} (C_2 \cdot \mathbf{S}_4) \mathcal{V}_{X_6}^{\text{cl}}$
- Quantum $\mu_{\rm str} = 1 \Delta$

leading corrections to the volumes cancel in the ratio, i.e., $\Delta
ightarrow 0$ at strict weak gauge coupling

Less Protection with Less SUSY [Klawer, S.-J.L., Weigand, Wiesner '20]

• Origin of Corrections

- Shrinking cycles in principle lead to corrections to the classical geometric calculations
- 4d F-theory w/ 4 Q's does suffer from quantum corrections (unlike in 6d w/ 8 Q's)

• Fate of the "WG Inequality"

- The sublattice is filled by WG particles if $\mu_{str} := 2r \frac{\mathcal{V}_{X_6}}{\mathcal{V}_{S_4} \mathcal{V}_{C_5}} = 1$ where $2r = C_2 \cdot S_4$
- Classical $\mu_{\text{str}}^{\text{cl}} = 1$ in other words: $\mathcal{V}_{\mathbf{S}_4}^{\text{cl}} \mathcal{V}_{C_2}^{\text{cl}} \stackrel{!}{=} (C_2 \cdot \mathbf{S}_4) \mathcal{V}_{X_6}^{\text{cl}}$
- Quantum $\mu_{\rm str} = 1 \Delta$

leading corrections to the volumes cancel in the ratio, i.e., $\Delta \to 0$ at strict weak gauge coupling



The Menu

First Course Swampland Conjectures & String Theory The Weak Gravity Conjecture(s) Second Course Preliminary Evidence (16 Qs) - Heterotic String on a Torus -

> Main Course Stronger Evidence (8 Qs) - 6d F-theory Vacua -

Strongest Evidence (4 Qs) - 4d F-theory Vacua -

..... Dessert

Evidence for Other Conjectures



The Distance Conjecture [Ooguri, Vafa, '06]

- Our F-theoretic test of the WGC concerns weak gauge coupling, which is achievable at infinite distance in the EFT parameter space.
- An independent conjecture (the Distance Conjecture) had claimed that at infinite distance in moduli space an infinite tower of states must become light.
- Later on this was refined to the Emergent String Conjecture [S.-J.L., Lerche, Weigand, '19]:

At infinite distance in moduli space a quantum gravity theory <u>either</u> decompactifies or reduces to a weakly coupled string theory

The Distance Conjecture [Ooguri, Vafa, '06]

- Our F-theoretic test of the WGC concerns weak gauge coupling, which is achievable at infinite distance in the EFT parameter space.
- An independent conjecture (the Distance Conjecture) had claimed that at infinite distance in moduli space an infinite tower of states must become light.
- Later on this was refined to the Emergent String Conjecture [S.-J.L., Lerche, Weigand, '19]:

At infinite distance in moduli space a quantum gravity theory <u>either</u> decompactifies <u>or</u> reduces to a <u>weakly coupled string theory</u>

The Distance Conjecture [Ooguri, Vafa, '06]

- Our F-theoretic test of the WGC concerns weak gauge coupling, which is achievable at infinite distance in the EFT parameter space.
- An independent conjecture (the Distance Conjecture) had claimed that at infinite distance in moduli space an infinite tower of states must become light.
- Later on this was refined to the Emergent String Conjecture [S.-J.L., Lerche, Weigand, '19]:



The Emergent String Conjecture [S.-J.L., Lerche, Weigand, '19]

At infinite distance in moduli space a quantum gravity theory <u>either</u> decompactifies

or reduces to a weakly coupled string theory

 Many non-trivial checks have recently been made on infinite distance limits of string EFTs, supporting the string emergence proposal.

Kahler Moduli (SIZE)

"emergence of unique critical tensionless string!"
F/M/IIA theory in 6/5/4d [S.-J.L., Lerche, Weigand '18-'20]
IIA/IIB hyper moduli in 4d [(Baume,) Marchesano, Wiesner '19]
M-theory in 4d [Xu '20]
F-theory in 4d, classical & quantum
[S.-J.L., Lerche, Weigand '19], [Klawer, S.-J.L., Weigand, Wiesner '20]

Complex Structure Moduli (SHAPE)

"decompactification via (dual) KK-like tower!"

Type II theory in 4d (closed string sector) [Grimm, Palti, Valenzuela '18], [Grimm, Li, Palti '18], [Klemm, Joshi '19], [Grimm, Li, Valenzuela '19], ... F-theory in 8d/6d (open string sector) [S.-J.L., (Lerche,) Weigand '21] / [Alvarez-Garcia, S.-J.L, Weigand '23]

The Emergent String Conjecture [S.-J.L., Lerche, Weigand, '19]

At infinite distance in moduli space a quantum gravity theory <u>either</u> decompactifies

or reduces to a weakly coupled string theory

 Many non-trivial checks have recently been made on infinite distance limits of string EFTs, supporting the string emergence proposal.

Kahler Moduli (SIZE)

"emergence of unique critical tensionless string!"
F/M/IIA theory in 6/5/4d [S.-J.L., Lerche, Weigand '18-'20]
IIA/IIB hyper moduli in 4d [(Baume,) Marchesano, Wiesner '19]
M-theory in 4d [Xu '20]
F-theory in 4d, classical & quantum
[S.-J.L., Lerche, Weigand '19], [Klawer, S.-J.L., Weigand, Wiesner '20]

Complex Structure Moduli (SHAPE)

"decompactification via (dual) KK-like tower!"

Type II theory in 4d (closed string sector) [Grimm, Palti, Valenzuela '18], [Grimm, Li, Palti '18], [Klemm, Joshi '19], [Grimm, Li, Valenzuela '19], ... F-theory in 8d/6d (open string sector) [S.-J.L., (Lerche,) Weigand '21] / [Alvarez-Garcia, S.-J.L, Weigand '23]

Completeness [Polchinski, '03]

 The excitations of our emergent heterotic string provide massive particles with an arbitrary charge.



The number theoretic (quasi-modular and quasi-Jacobi, resp., in 6d and 4d) property of the heterotic partition function allows no gaps in the charge spectrum.*

Completeness [Polchinski, '03]

 The excitations of our emergent heterotic string provide massive particles with an arbitrary charge.



The number theoretic (quasi-modular and quasi-Jacobi, resp., in 6d and 4d) property of the heterotic partition function allows no gaps in the charge spectrum.*

Completeness [Polchinski, '03]

 The excitations of our emergent heterotic string provide massive particles with an arbitrary charge.



• The number theoretic (quasi-modular and quasi-Jacobi, resp., in 6d and 4d) property of the heterotic partition function allows no gaps in the charge spectrum.*

Completeness [Polchinski, '03]

 The excitations of our emergent heterotic string provide massive particles with an arbitrary charge.



• The number theoretic (quasi-modular and quasi-Jacobi, resp., in 6d and 4d) property of the heterotic partition function allows no gaps in the charge spectrum.*

No Global Symmetry [Banks, Dixon, '88]

 At weak gauge coupling, we have observed an infinite tower of particles (the heterotic excitations) become light.

 This indicates that the effective description of the physics breaks down and in particular that the strict weak gauge coupling limit is not part of the moduli space.

In other words, the U(1) gauge vector <u>cannot</u> lead to a strict global U(1) symmetry; this supports the No Global Symmetry Conjecture.

SUMMARY

The WGC claims weakness of gravity as a general feature of quantum gravity.

 The minimal WGC predicts a WG particle (w/ mass smaller than charge), while the sublattice version conjectures a charge sublattice to be filled by WG particles.

 We have discussed strong evidence for the WGC from string theory.
 Microscopically, the sublattice of WG particles arise from the excitations of an (emergent) critical string.

Preliminary evidence is found for the heterotic EFTs (w/ 16 Q's).

The WGC claims weakness of gravity as a general feature of quantum gravity.

 The minimal WGC predicts a WG particle (w/ mass smaller than charge), while the sublattice version conjectures a charge sublattice to be filled by WG particles.

 We have discussed strong evidence for the WGC from string theory.
 Microscopically, the sublattice of WG particles arise from the excitations of an (emergent) critical string.

Preliminary evidence is found for the heterotic EFTs (w/ 16 Q's).

- The WGC claims weakness of gravity as a general feature of quantum gravity.
- The minimal WGC predicts a WG particle (w/ mass smaller than charge), while the sublattice version conjectures a charge sublattice to be filled by WG particles.

 We have discussed strong evidence for the WGC from string theory.
 Microscopically, the sublattice of WG particles arise from the excitations of an (emergent) critical string.

• Preliminary evidence is found for the heterotic EFTs (w/ 16 Q's).

 Much stronger evidence is found in the weak gauge coupling limits of the F-theoretic EFTs (w/ 8 or 4 Q's), where a light critical string necessarily emerges to realize the sublattice WGC!

Specifically, the realization has been conspired by the <u>universal</u> features of

- Asymptotic Effective Action;
- Internal Geometry;
- String Partition Function.

 Our F-theoretic analysis has also served as evidence for some other conjectures (String Emergence, No Global Symmetry and Completeness).

 Much stronger evidence is found in the weak gauge coupling limits of the F-theoretic EFTs (w/ 8 or 4 Q's), where a light critical string necessarily
 emerges to realize the sublattice WGC!

- Specifically, the realization has been conspired by the <u>universal</u> features of
 - Asymptotic Effective Action;
 - Internal Geometry;
 - String Partition Function.

 Our F-theoretic analysis has also served as evidence for some other conjectures (String Emergence, No Global Symmetry and Completeness).

 Much stronger evidence is found in the weak gauge coupling limits of the F-theoretic EFTs (w/ 8 or 4 Q's), where a light critical string necessarily
 emerges to realize the sublattice WGC!

- Specifically, the realization has been conspired by the <u>universal</u> features of
 - Asymptotic Effective Action;
 - Internal Geometry;
 - String Partition Function.
- Our F-theoretic analysis has also served as evidence for some other conjectures (String Emergence, No Global Symmetry and Completeness).

 Much stronger evidence is found in the weak gauge coupling limits of the F-theoretic EFTs (w/ 8 or 4 Q's), where a light critical string necessarily
 emerges to realize the sublattice WGC!

- Specifically, the realization has been conspired by the <u>universal</u> features of
 - Asymptotic Effective Action;
 - Internal Geometry;
 - String Partition Function.
- Our F-theoretic analysis has also served as evidence for some other conjectures (String Emergence, No Global Symmetry and Completeness).

Ghank You