

THE SEARCH FOR AXIONS CIRCA 2023

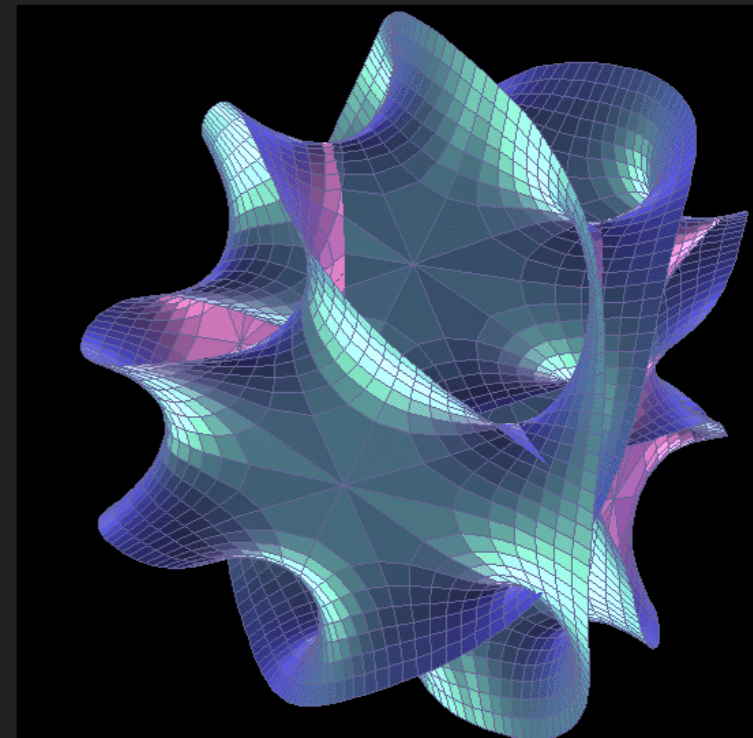
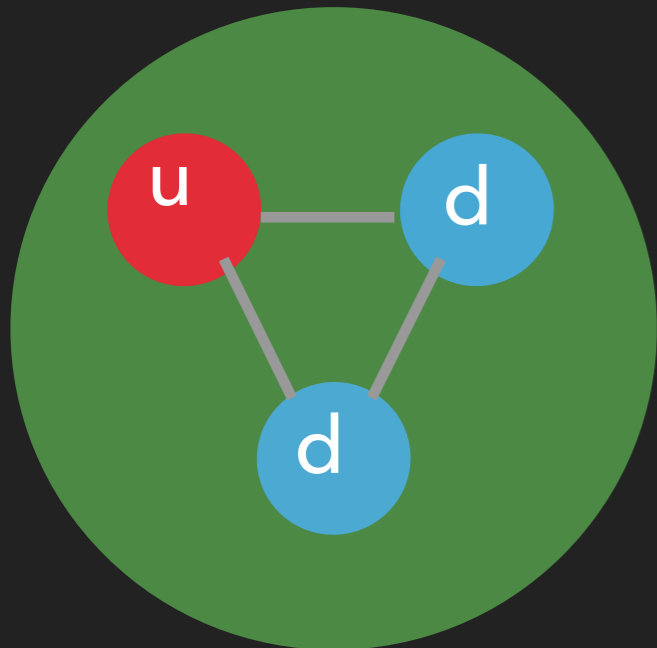
BEN SAFDI

**BERKELEY CENTER FOR THEORETICAL PHYSICS
UNIVERSITY OF CALIFORNIA, BERKELEY**

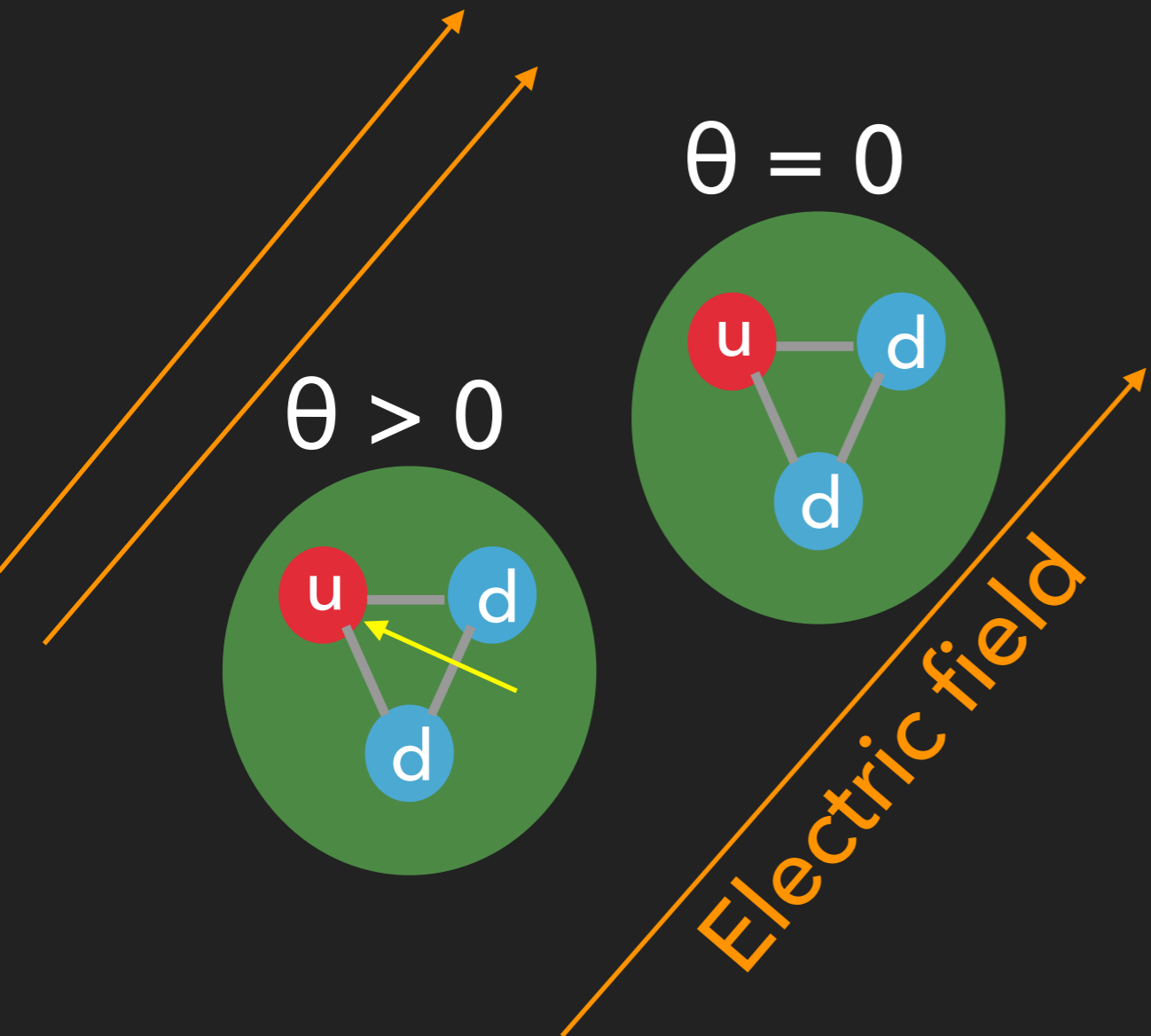
CHUNG-ANG UNIVERSITY BEYOND THE STANDARD MODEL WORKSHOP

Reasons to think about axions

1. They solve the Strong-CP Problem
2. They can explain the observed dark matter
3. They arise generically in string theory
4. They may be connected to deep aspects of quantum gravity (see Matt Reece's talk)

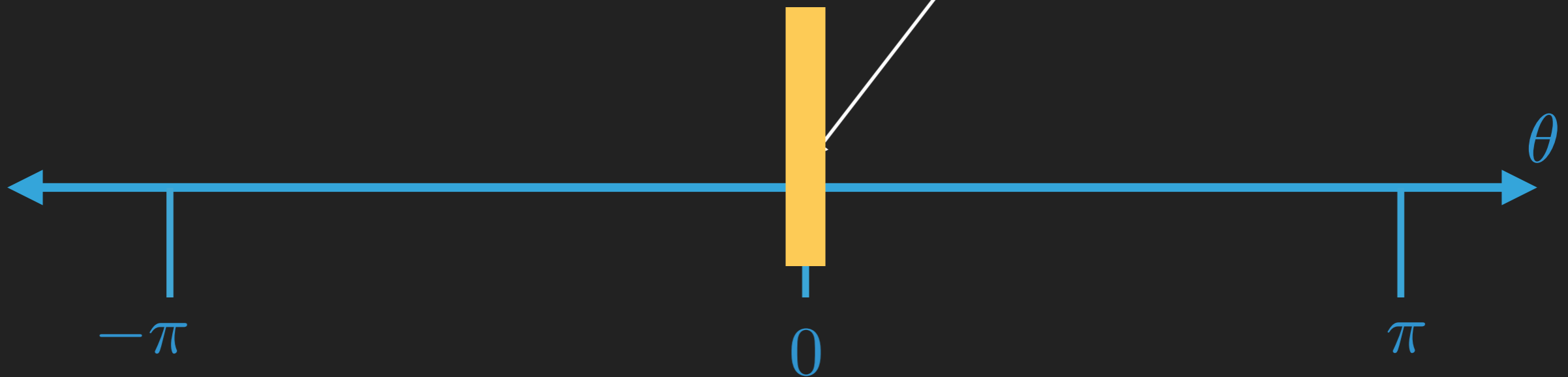


The Strong CP Problem: no neutron EDM?



why do we live here?

1/1000000000



The neutron electric dipole moment puzzle

Roberto Peccei



1977

Helen Quinn



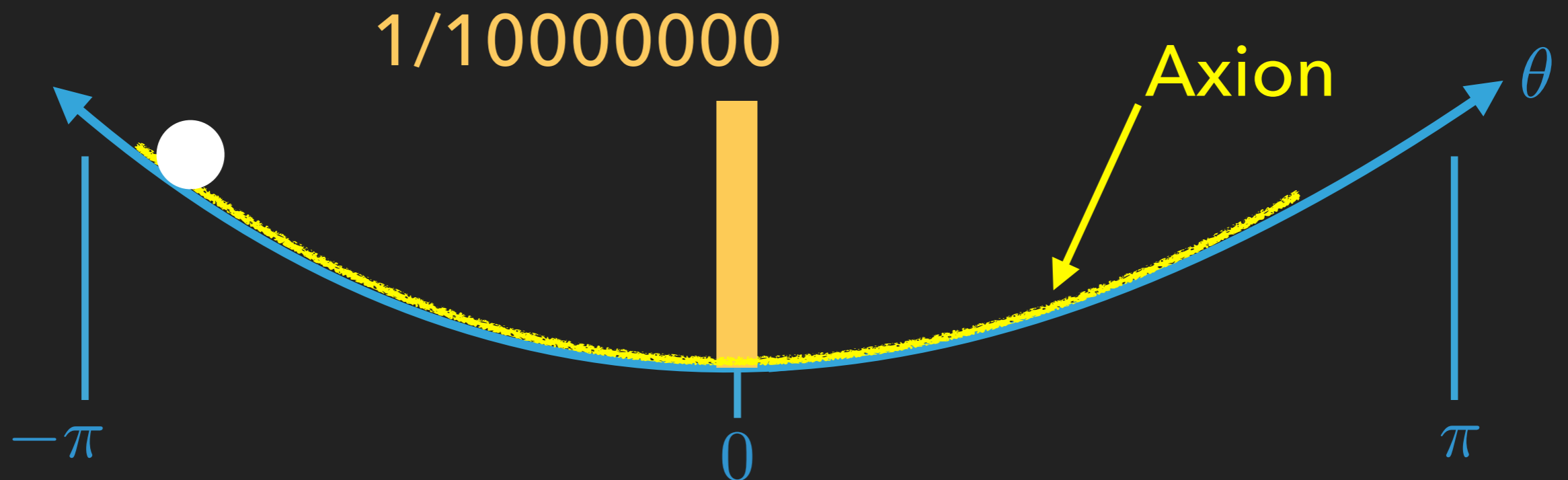
Steven Weinberg



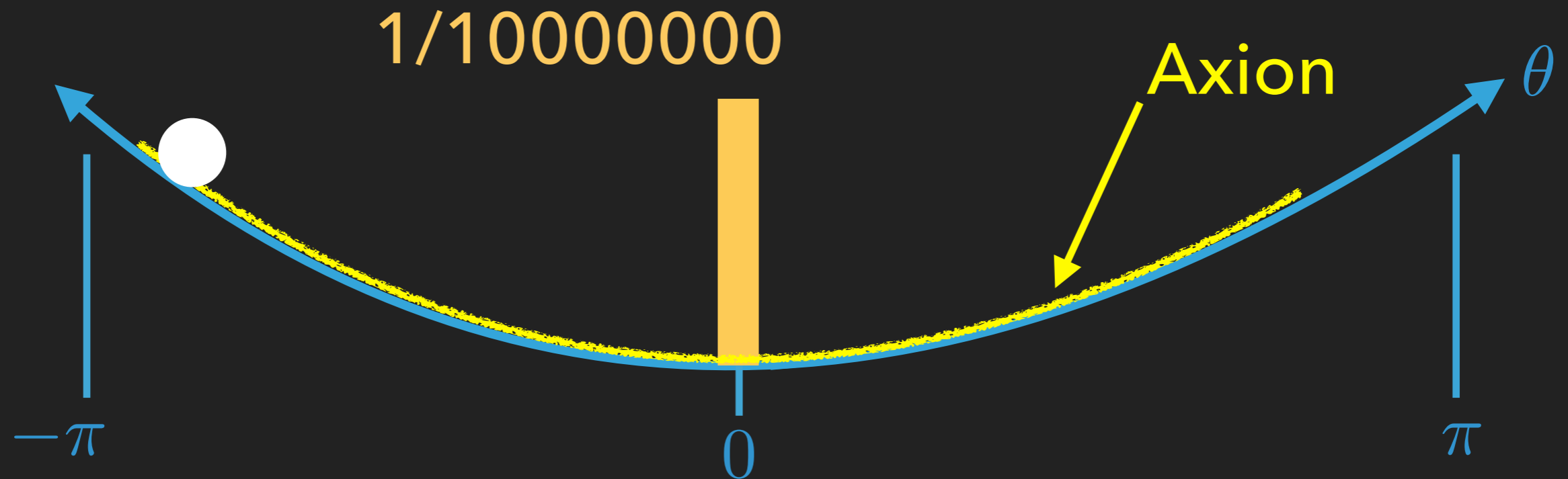
Frank Wilczek



1978



Axion Solution to Strong CP (more precisely)



$$\mathcal{L} = -\frac{g^2}{32\pi^2} \left(\bar{\theta} + \frac{a}{f_a} \right) G_{\mu\nu} \tilde{G}^{\mu\nu} \quad f_a \gtrsim 10^9 \text{ GeV}$$

$$d_n \propto \left(\bar{\theta} + \frac{a}{f_a} \right)$$

$$V(a) \approx \frac{1}{2} \Lambda_{\text{QCD}}^4 \left(\bar{\theta} + \frac{a}{f_a} \right)^2$$

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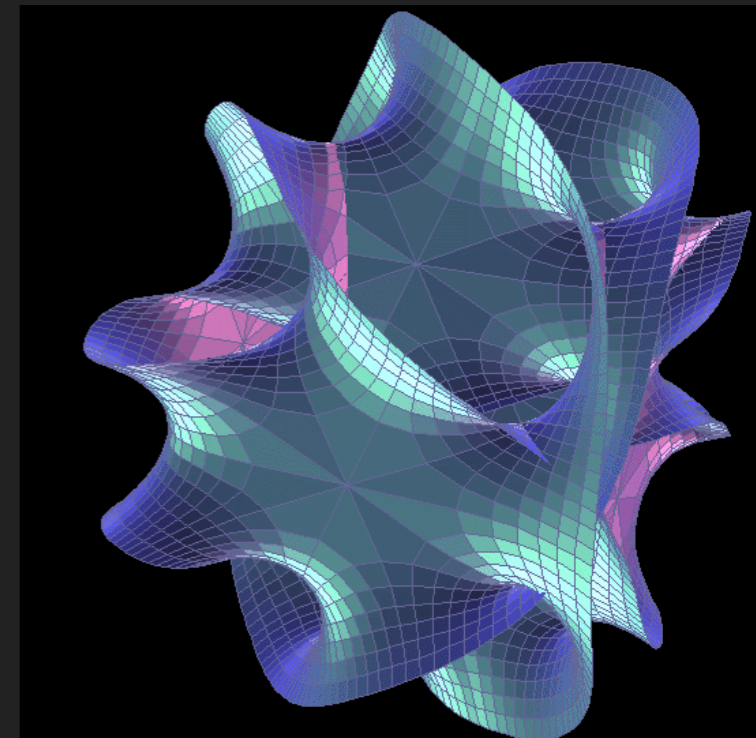
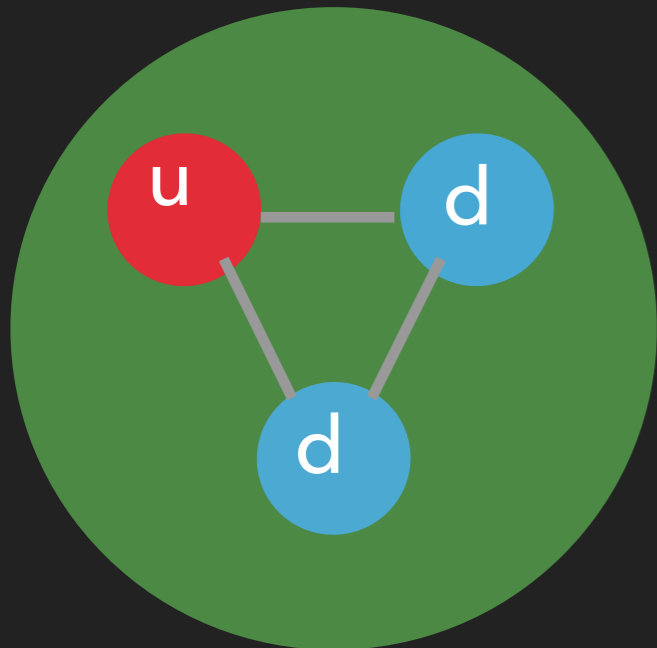
Axion mass: $m_a \approx \frac{\Lambda_{\text{QCD}}^2}{f_a} \approx 10^{-5} \text{ eV} \left(\frac{10^{12} \text{ GeV}}{f_a} \right)$

Axions also couple to EM: $\mathcal{L} = -g_{a\gamma\gamma} \frac{a F \tilde{F}}{4} = g_{a\gamma\gamma} a \mathbf{E} \cdot \mathbf{B}$

$$g_{a\gamma\gamma} = \frac{C_\gamma \alpha_{\text{EM}}}{2\pi f_a}, \quad C_\gamma \sim \mathcal{O}(1)$$

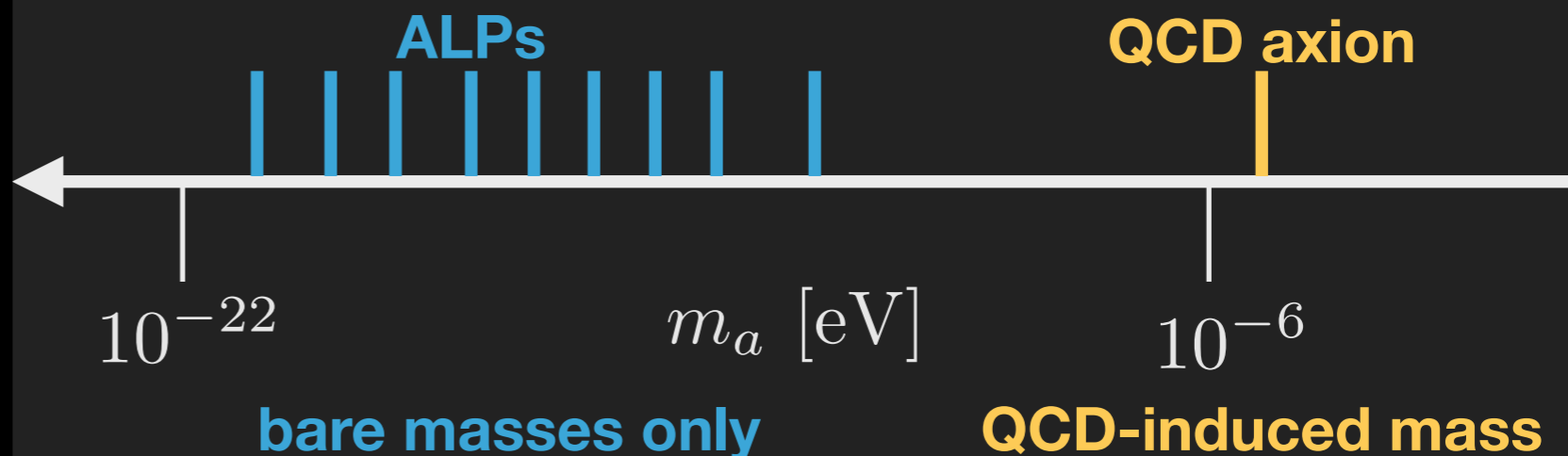
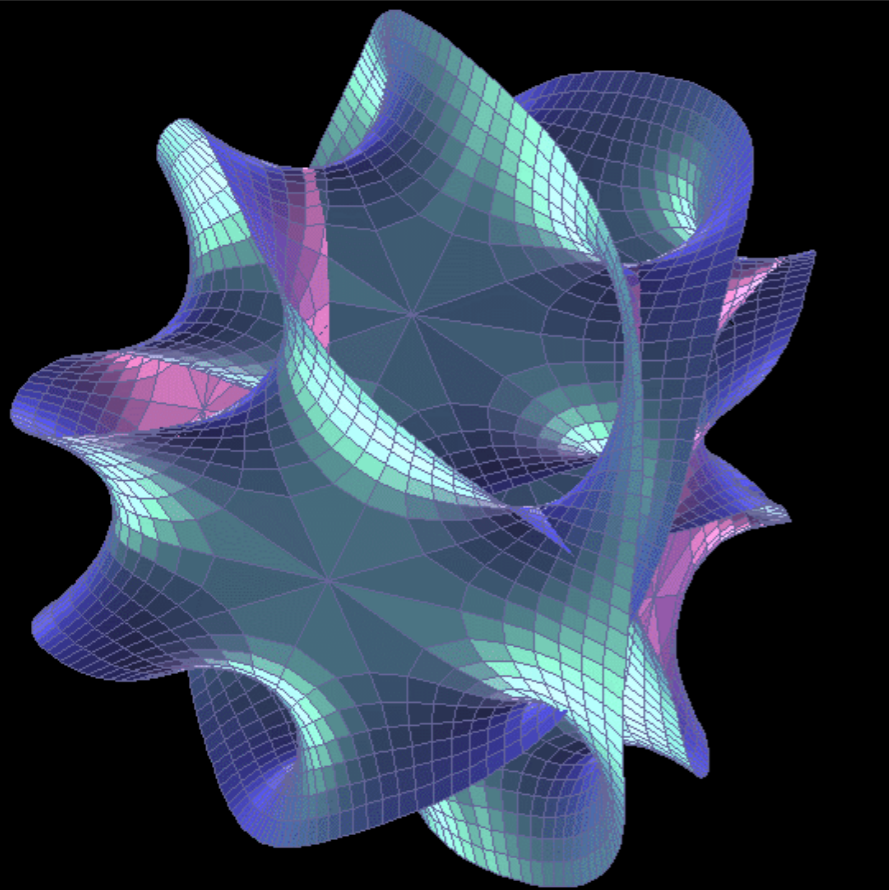
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Axions Appear in String Theory

- ▶ Axions arise from compactified 2-form fields
- ▶ Perturbative masses protected by 10D gauge symmetry
- ▶ “Ultralight” non-perturbative masses (e.g. string instantons)



Axion-like particles versus QCD Axion

$$\mathcal{L} = \frac{-g^2}{32\pi^2} \frac{a}{f_a} G_{\mu\nu} \tilde{G}^{\mu\nu} - \frac{C_\gamma \alpha_{\text{EM}}}{8\pi f_a} a F_{\mu\nu} \tilde{F}^{\mu\nu} + \sum_f \frac{C_f}{2f_a} \partial_\mu a \bar{f} \gamma^\mu \gamma_5 f - \frac{1}{2} m_{\text{bare}}^2 a^2$$

~~QCD axion only!~~
 $g_{a\gamma\gamma}$
 g_{aff}
 $m_a^{\text{QCD}} \gg m_{\text{bare}}$
 small



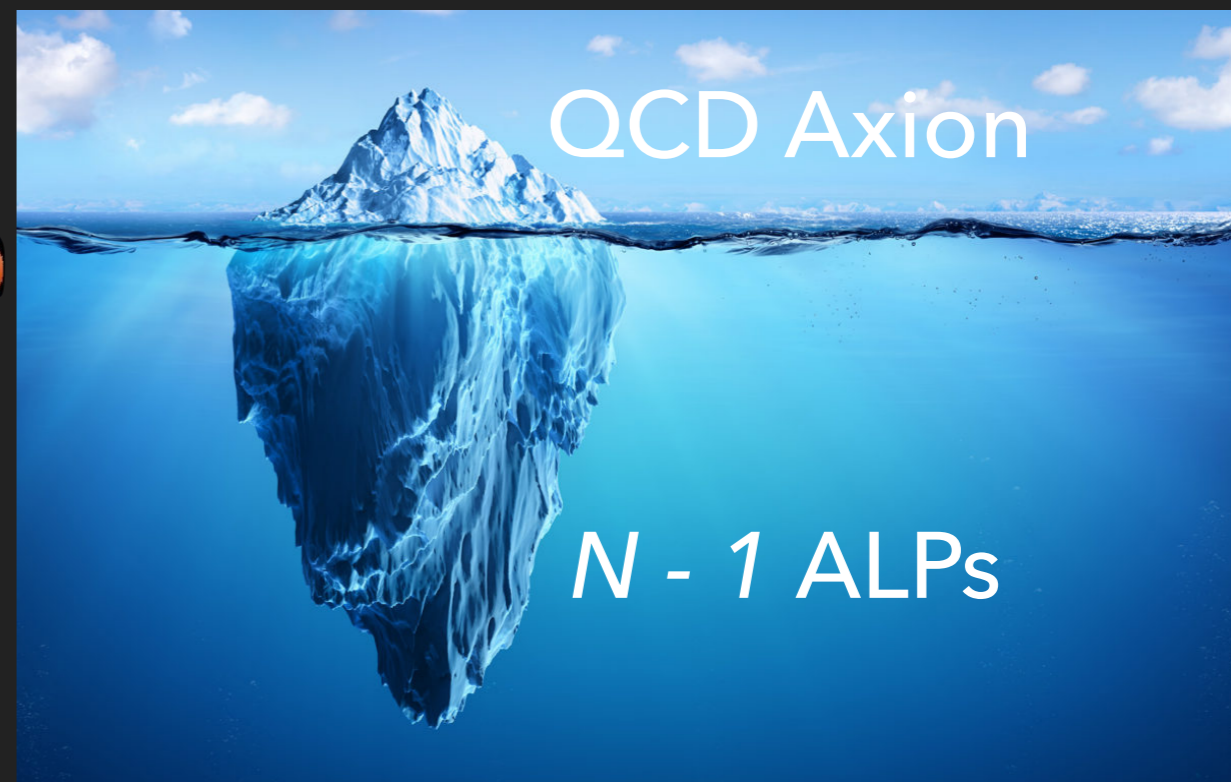
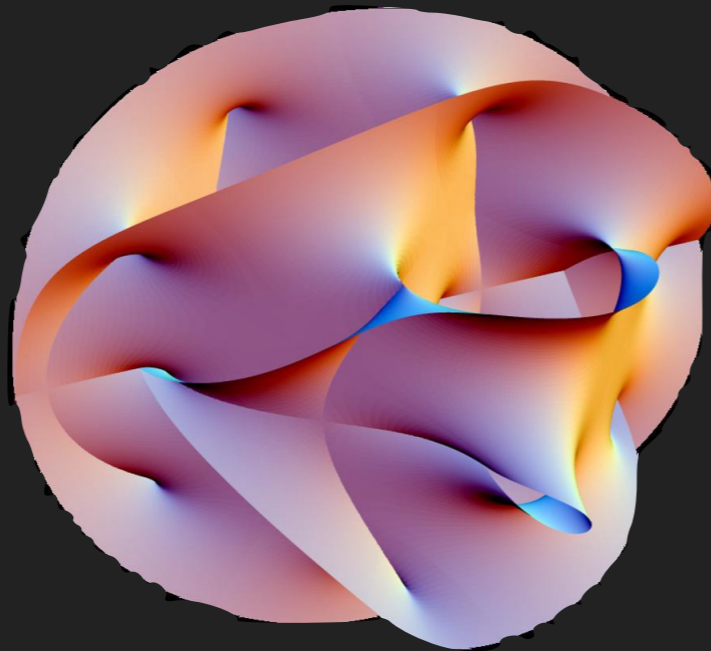
String Axiverse: N pseudo-scalars $\rightarrow N-1$ ALPs + 1 QCD axion

String theory provides answer to: why axions?



Axiverse: N pseudo-scalars $\rightarrow N-1$ ALPs + 1 QCD axion

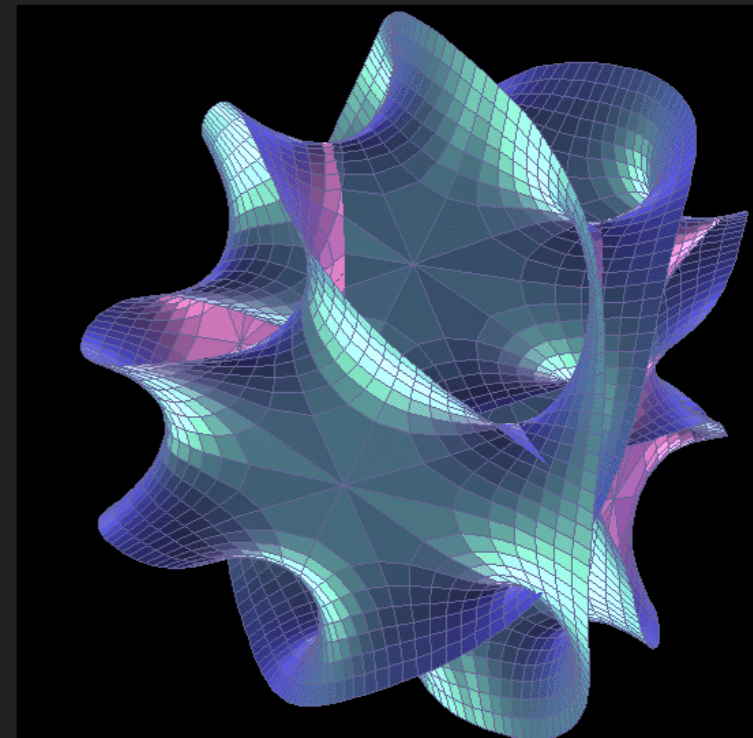
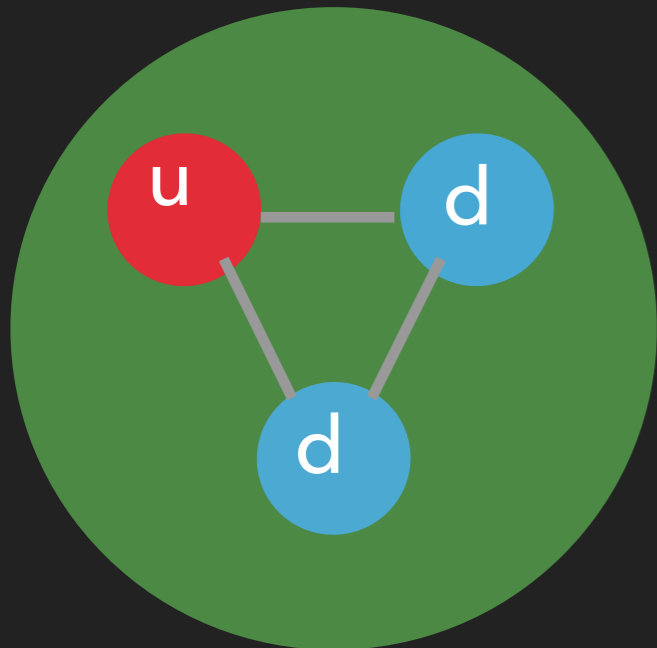
Zero modes of 2-form fields generate **10's - 100's of ALPs** from compactification on internal space



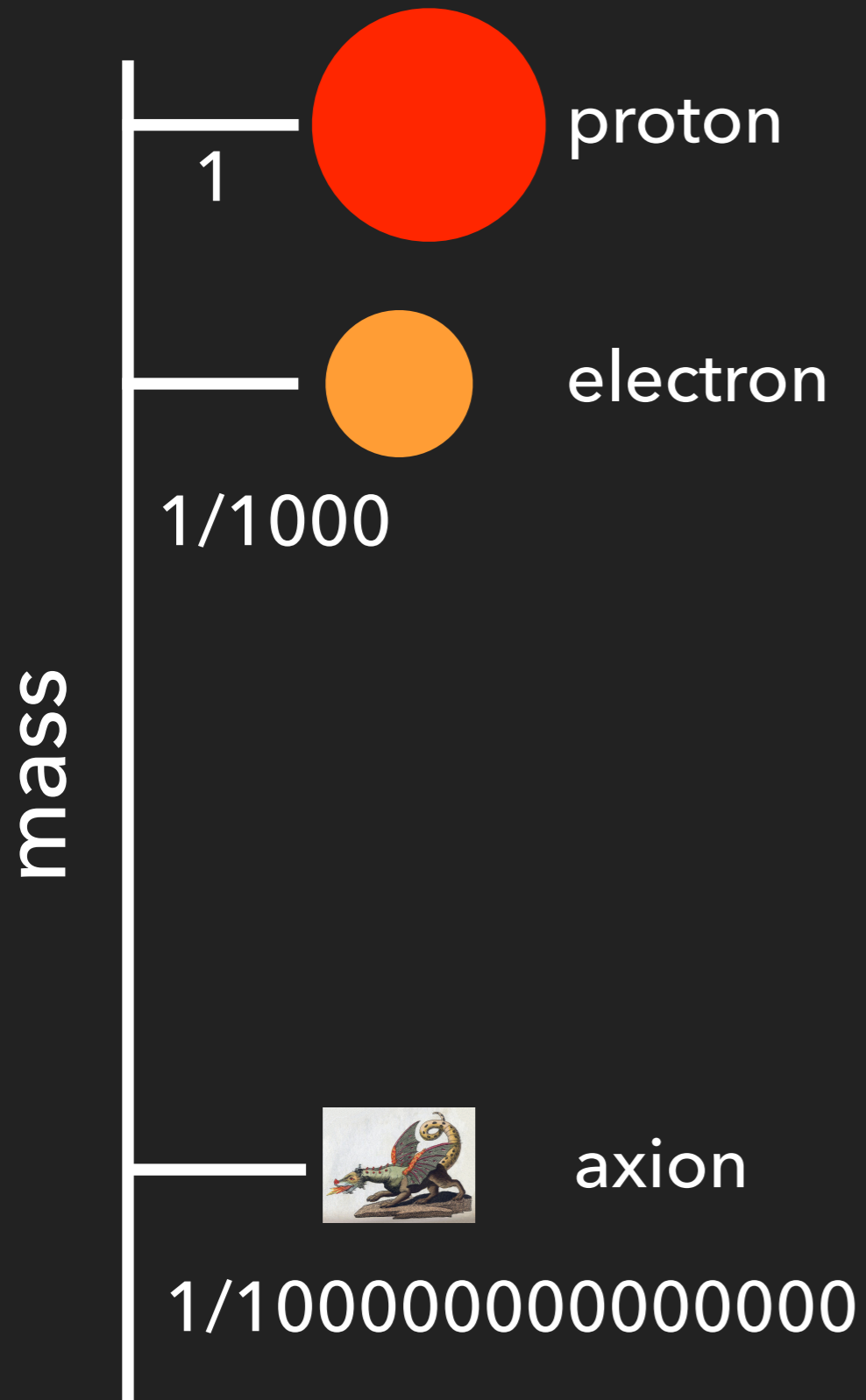
Axions could be only way to experimentally probe string theory compactifications!

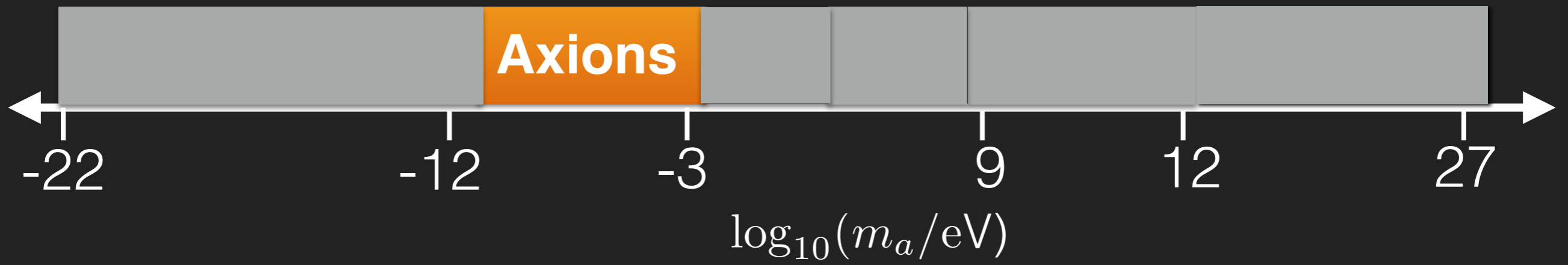
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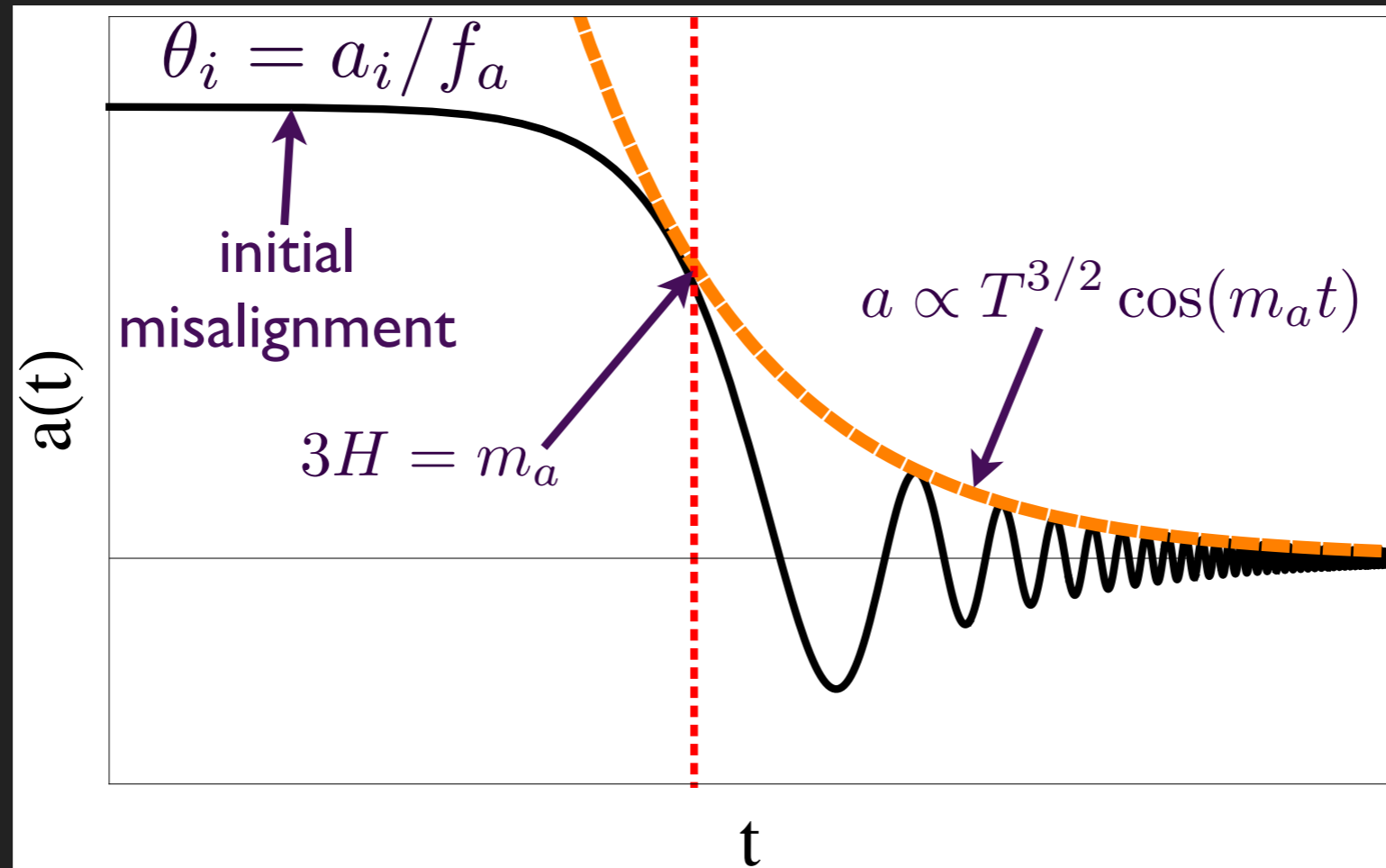


axion dark matter



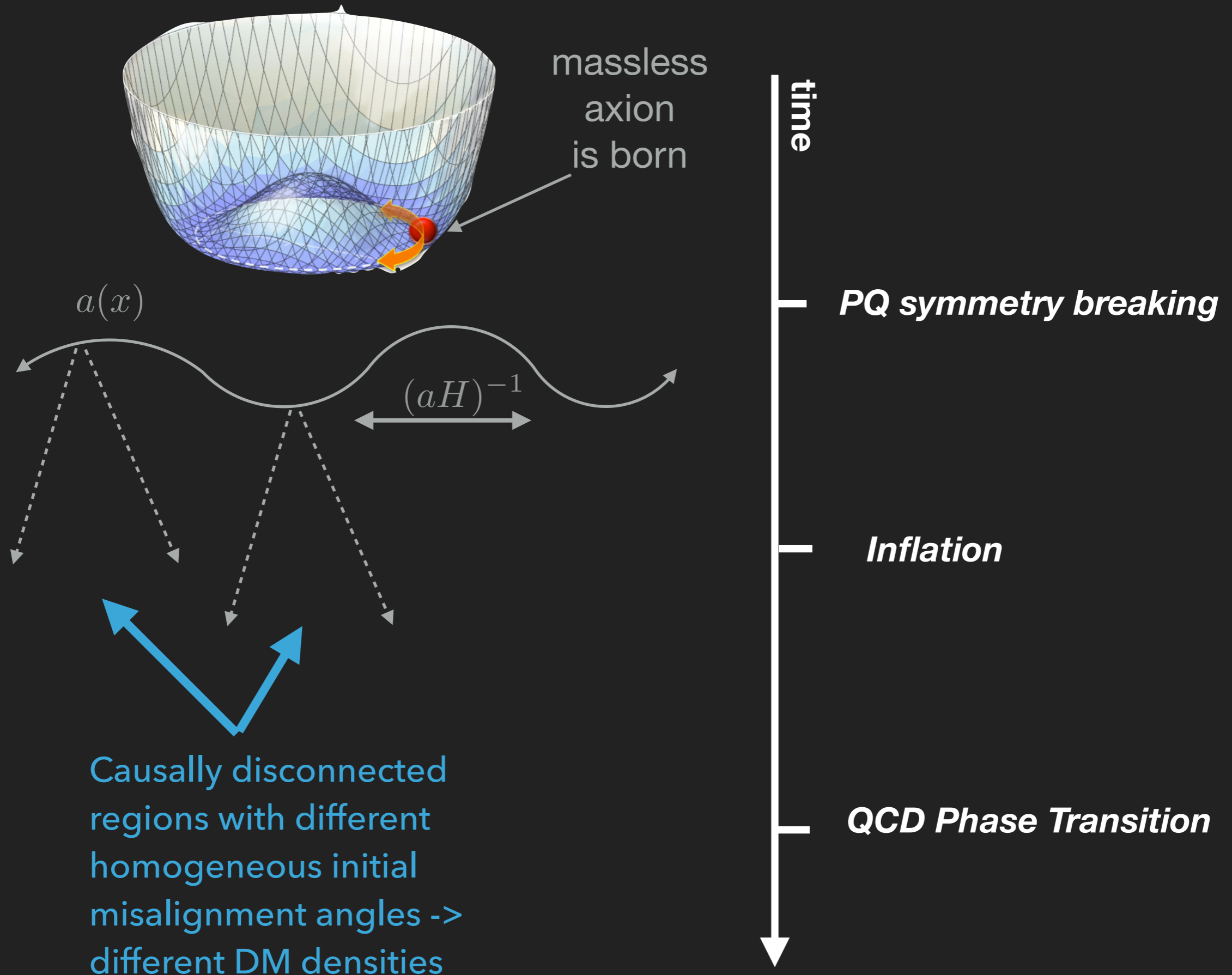


$$\sin a/f_a \sim a/f_a : \quad \ddot{a} + 3H\dot{a} + m_a^2 a = 0$$

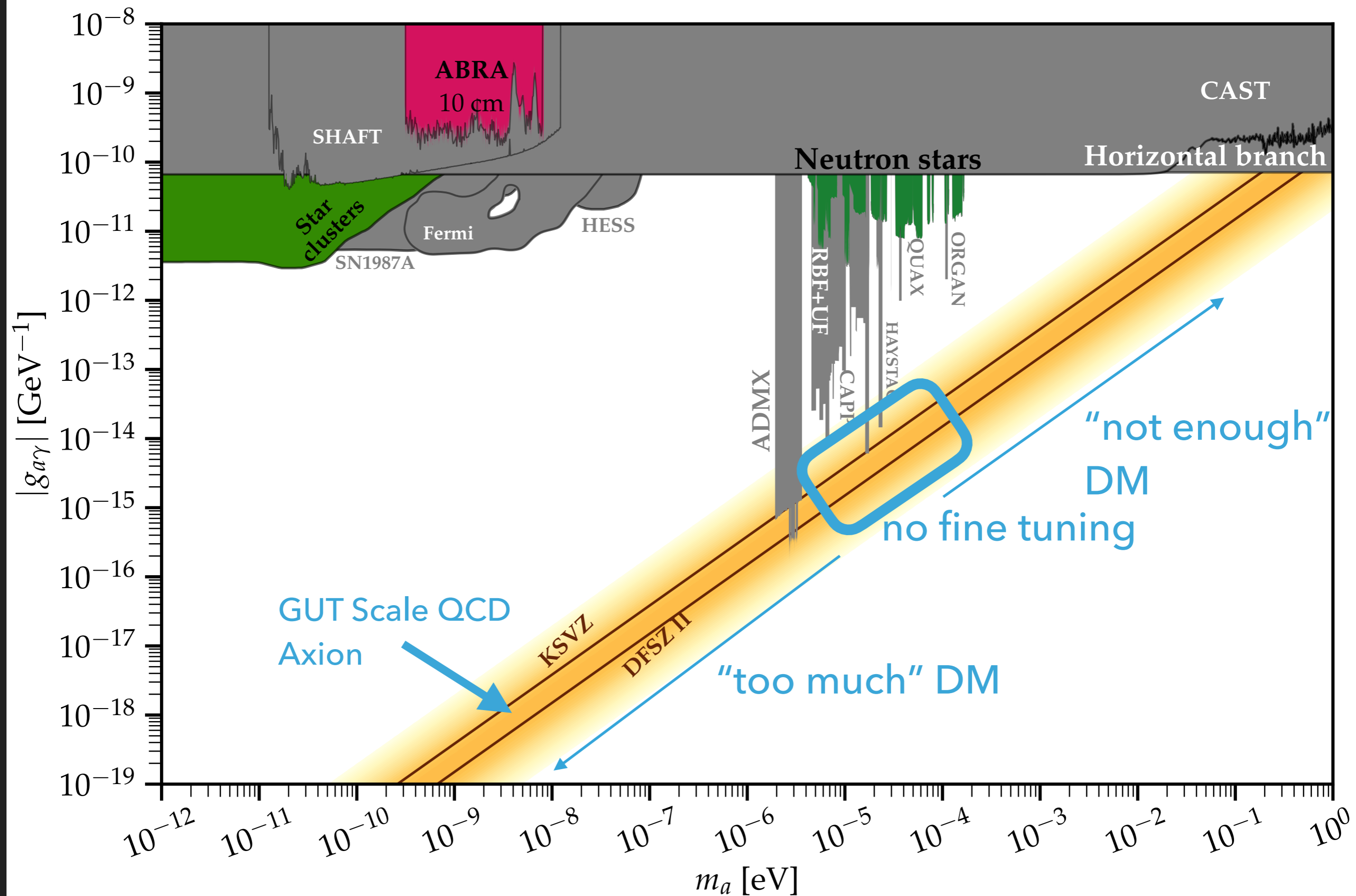


$$\Omega_a h^2 \sim 0.1 \left(\frac{f_a}{10^{12} \text{ GeV}} \right)^{7/6} \theta_i^2$$

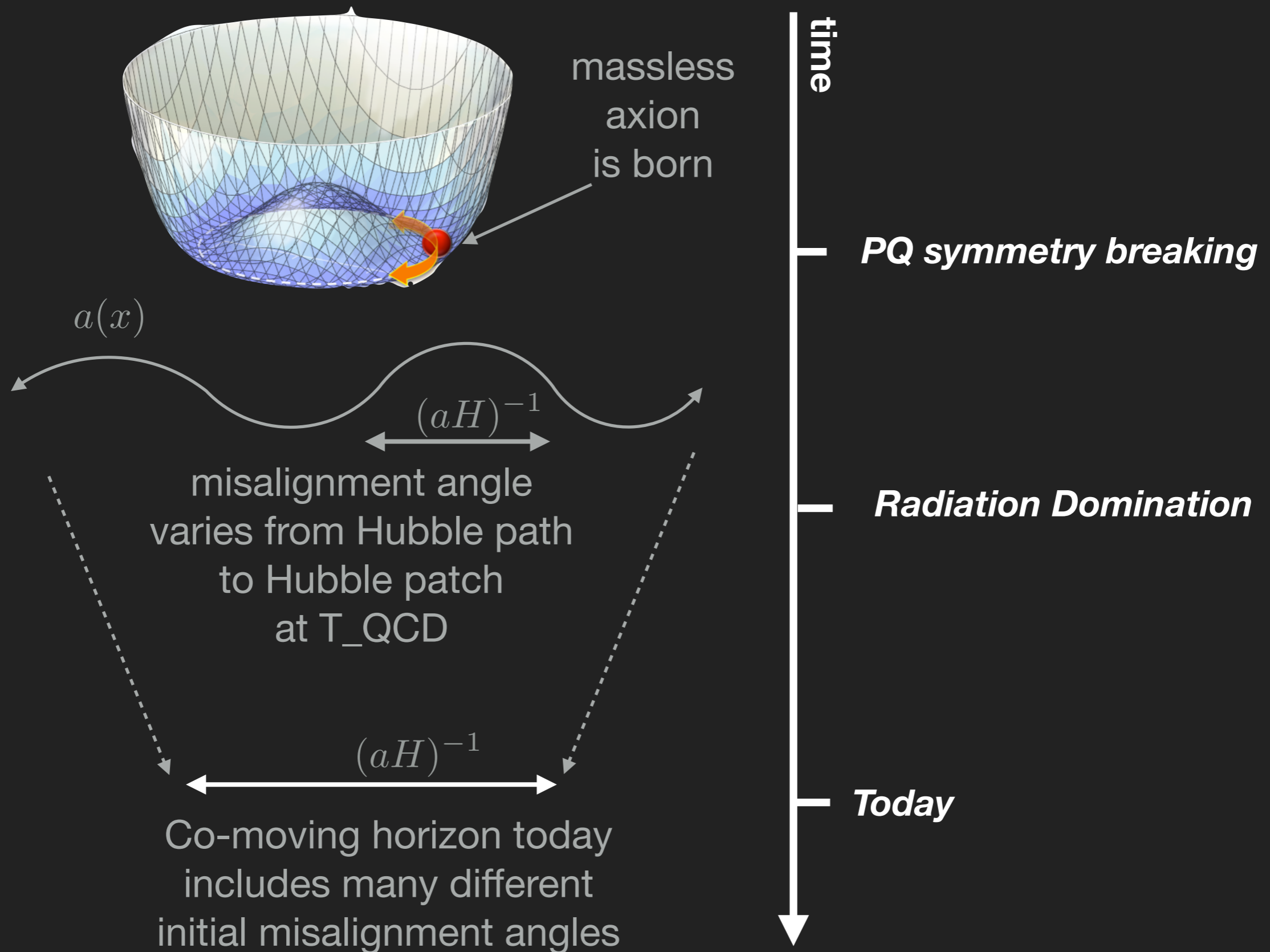
Axion generated before inflation



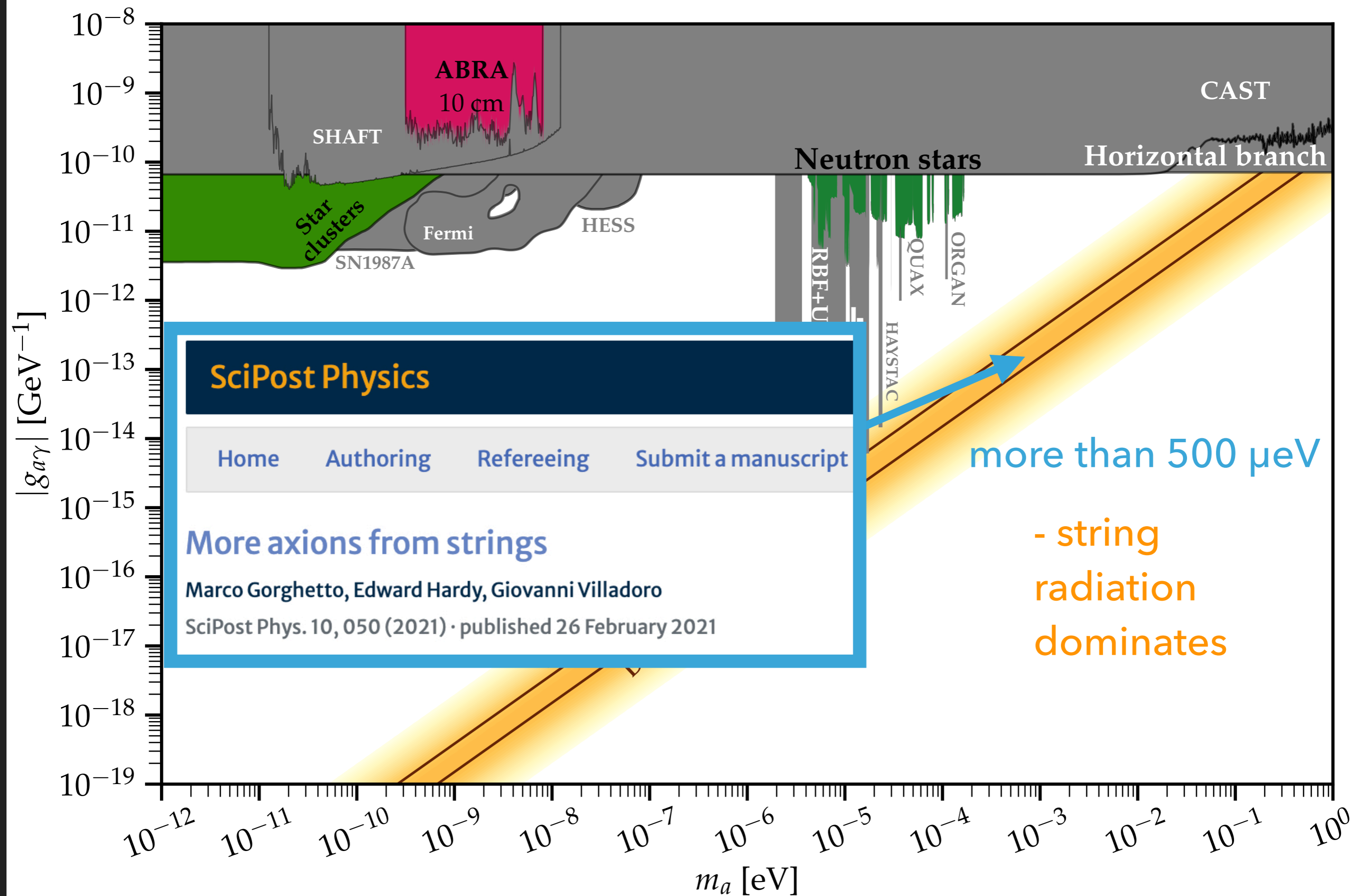
PQ Broken Before Inflation



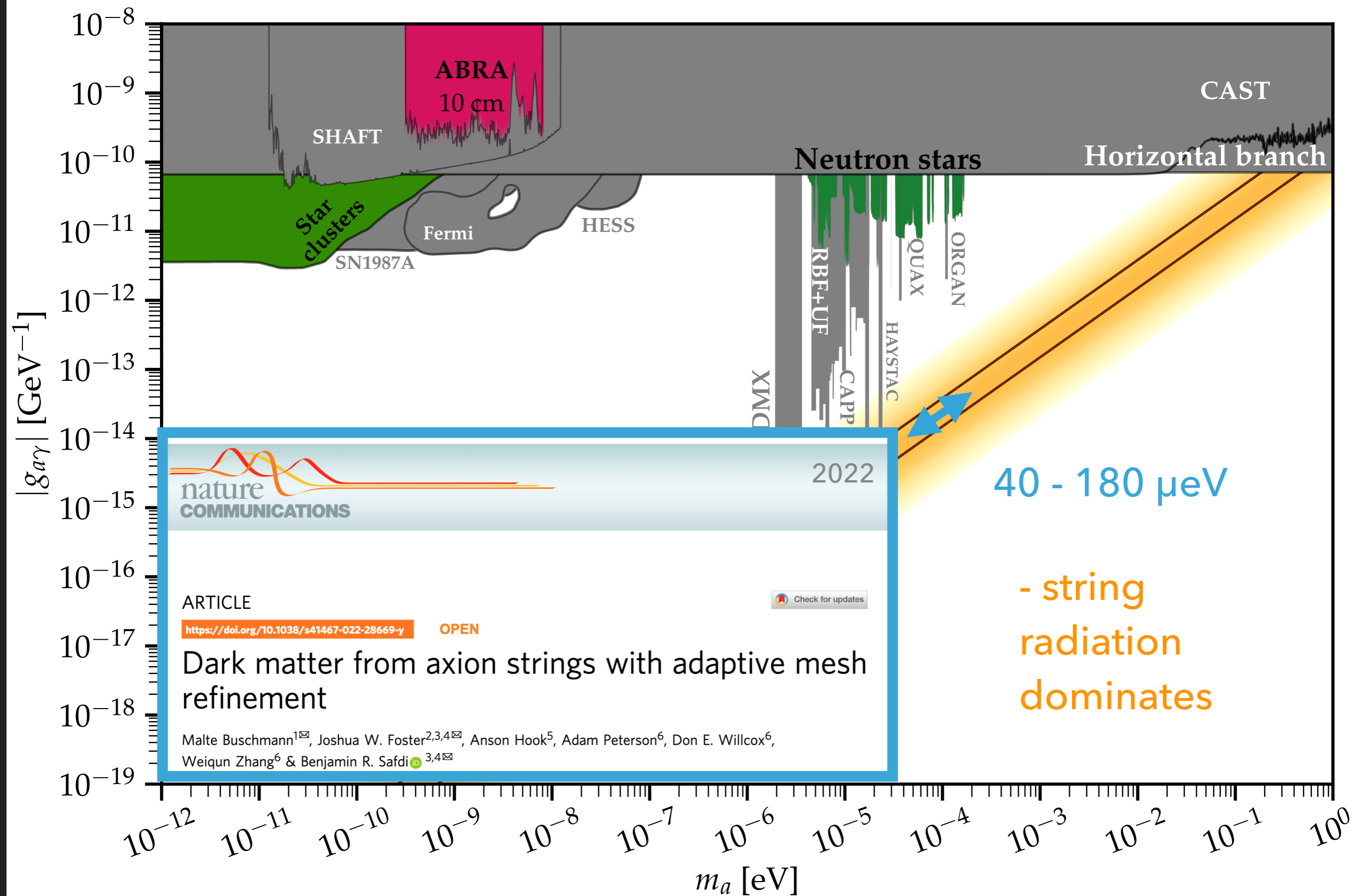
Axion generated after inflation



What does the literature say?



What does the literature say?



Axion generated after inflation

M. Buschmann, J. Foster, **B.S.** PRL 2020

Simulate on static grid with $\sim 10^{10}$ sites

Simulate from PQ phase transition to matter-radiation equality

M. Buschmann, J. Foster, **B.S.**, AMReX Collaboration, Nat. Comm. 2022

Simulate on adaptive grid equiv. to static grid with $\sim 10^{15}$ sites



National Energy Research
Scientific Computing Center



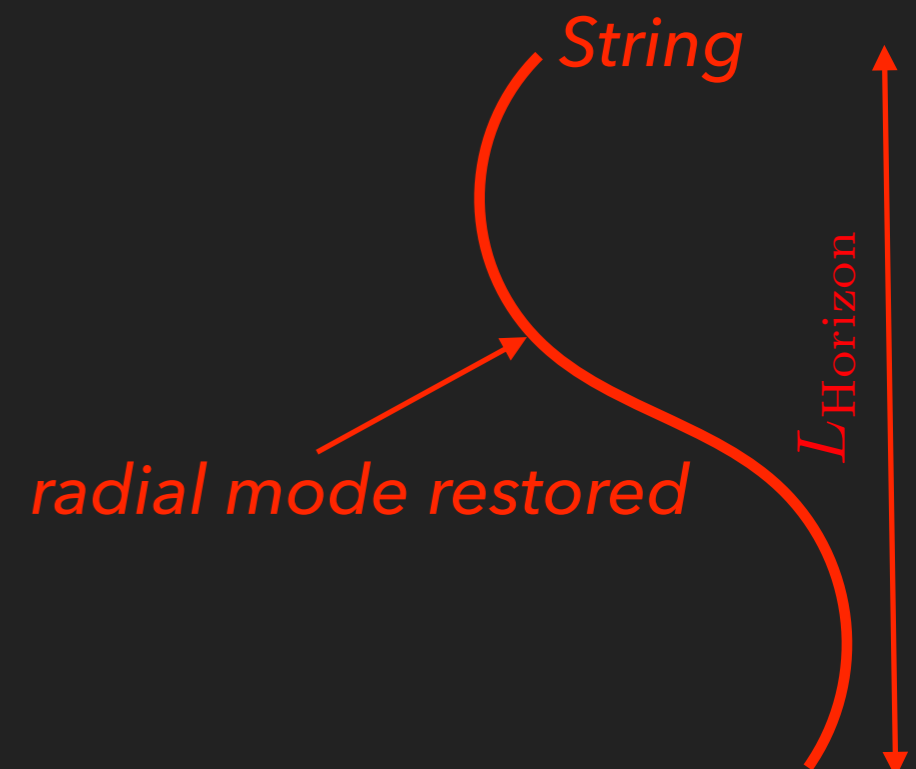
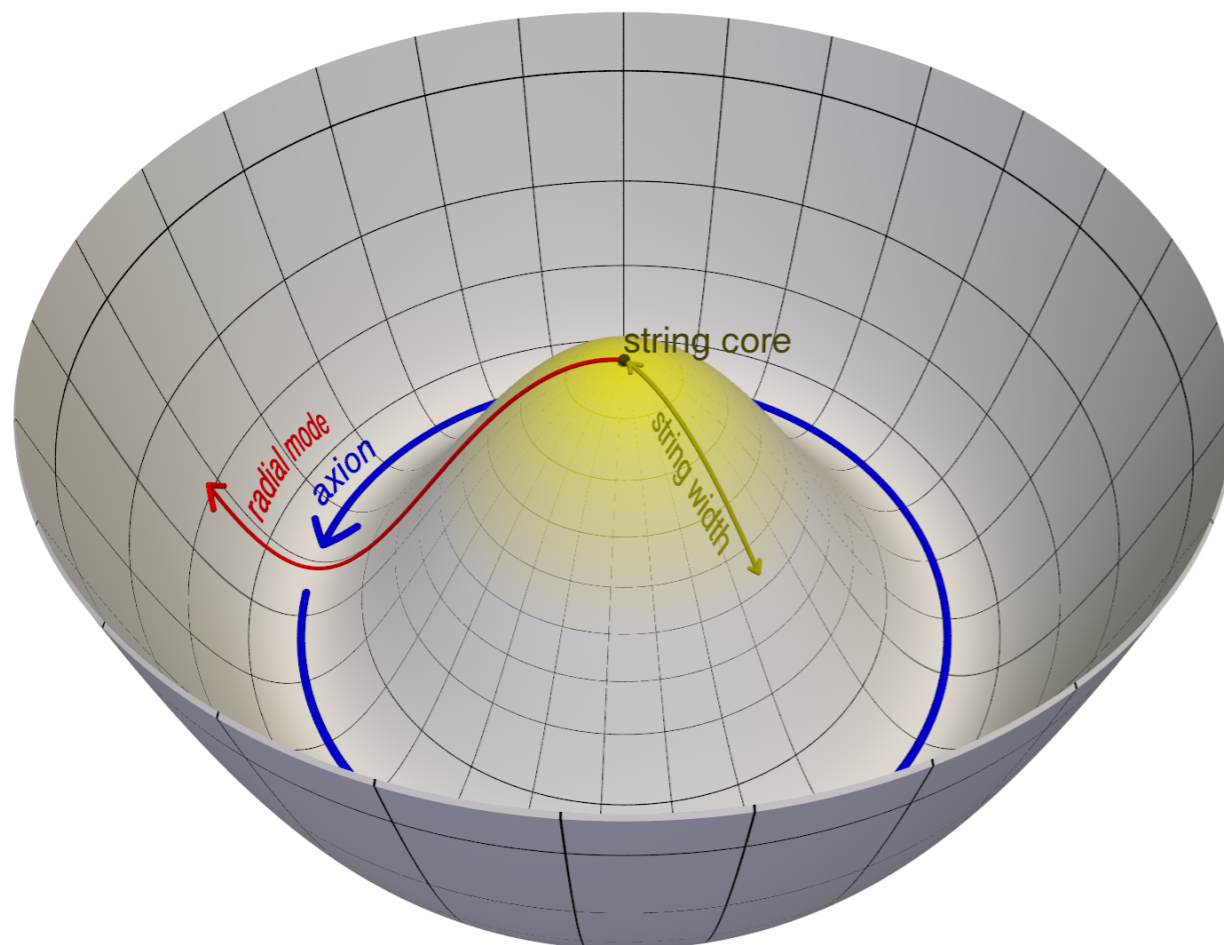
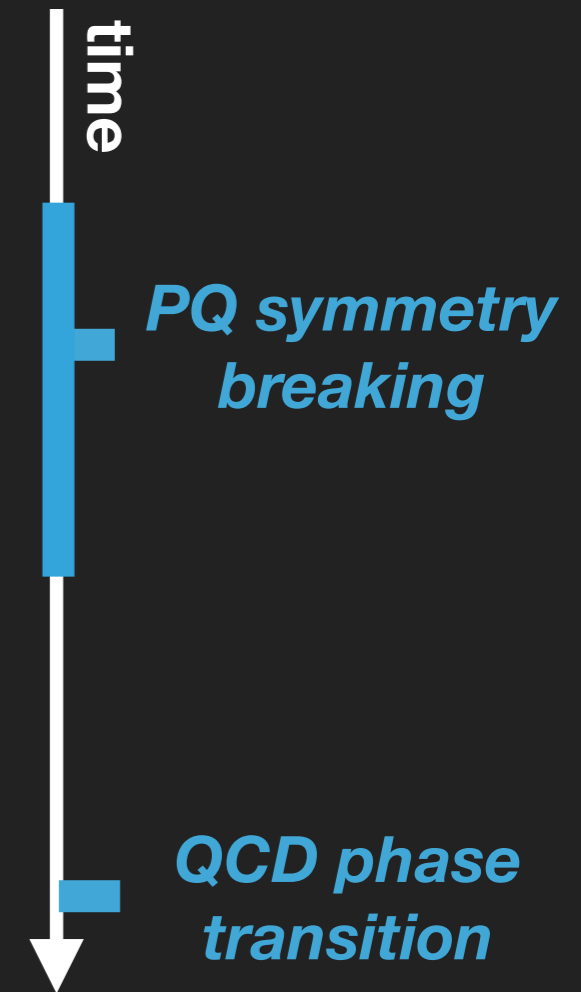
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ENERGY

Office of
Science

Axion dark matter primarily produced from strings

$$\mathcal{L}_{\text{PQ}} = \frac{1}{2} |\partial\Phi|^2 - \frac{\lambda}{4} (|\Phi|^2 - f_a^2)^2 - \frac{\lambda T^2}{6} |\Phi|^2$$

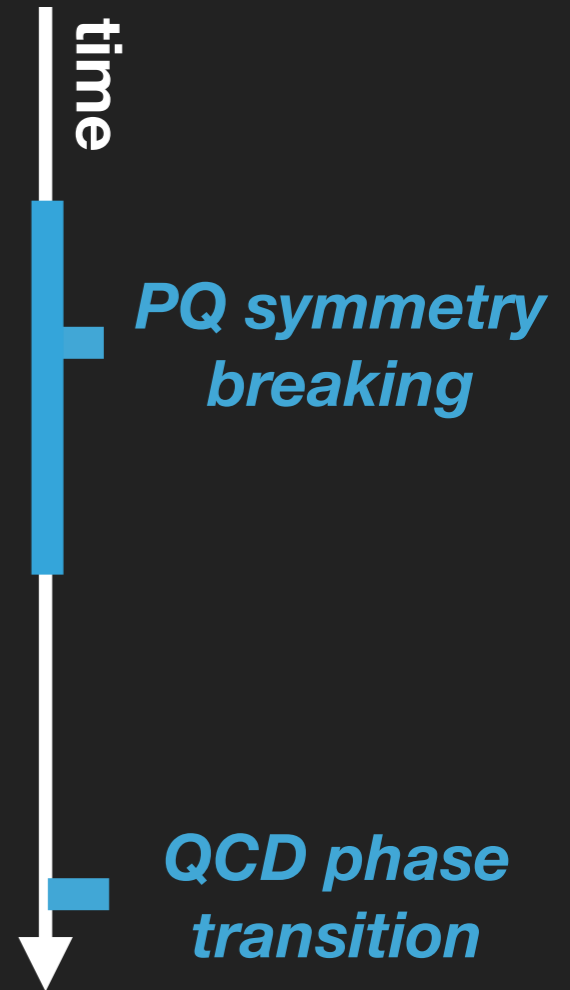
- ▶ symmetry broken for $T \lesssim f_a$
- ▶ write $\Phi(x) = \frac{r(x) + f_a}{\sqrt{2}} e^{ia(x)/f_a}$
- ▶ radial mode acquires mass: $m_r \sim \sqrt{\lambda} f_a$ (take $\lambda \sim 1$)



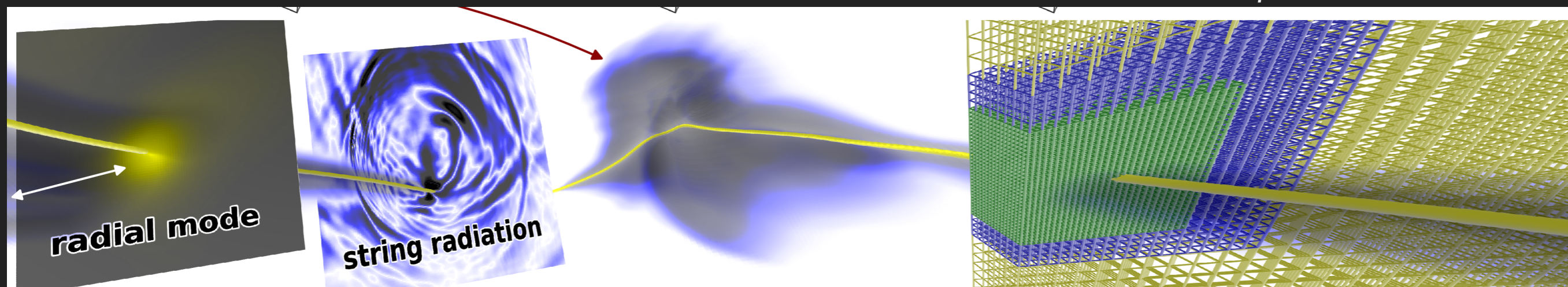
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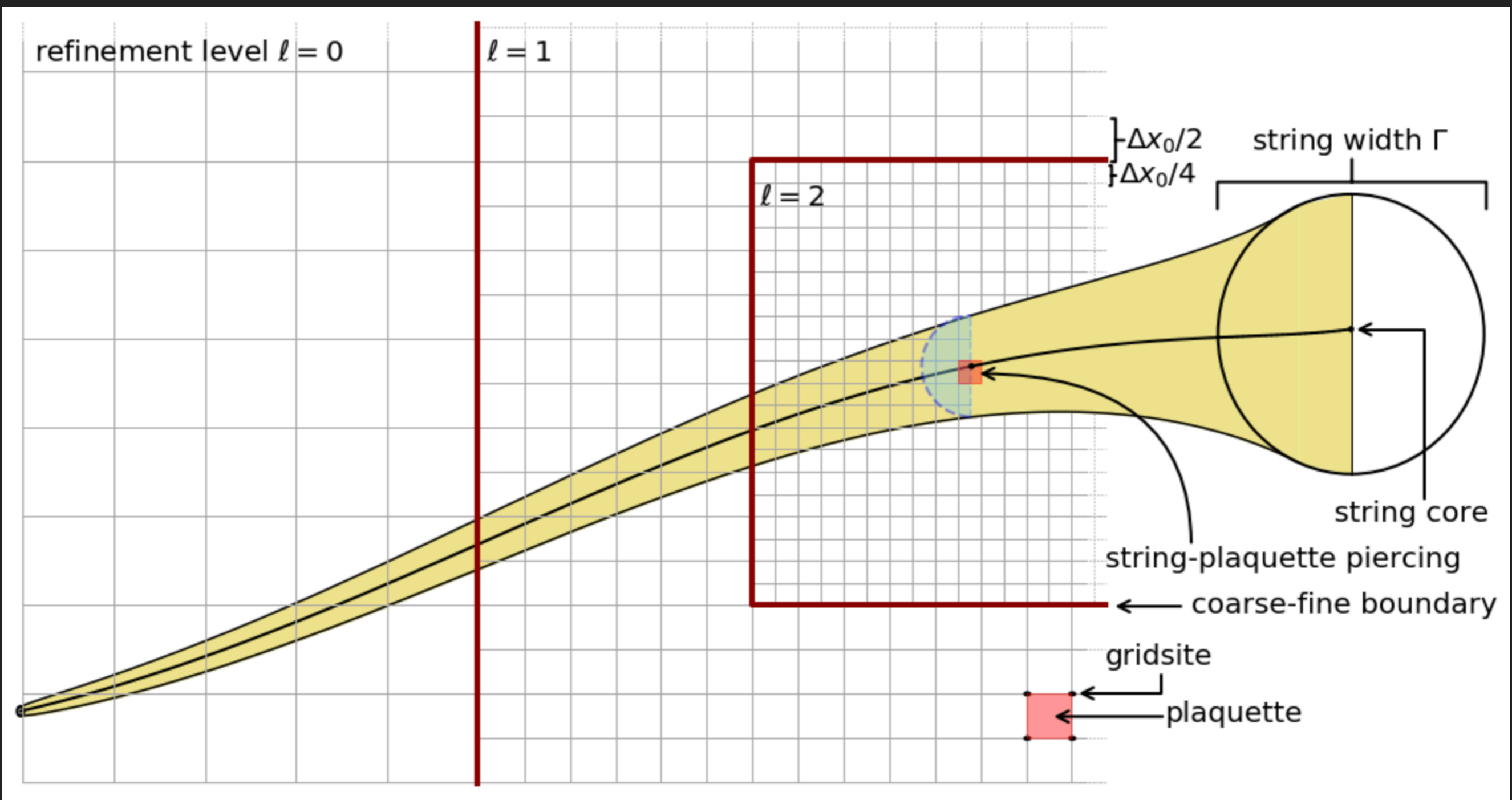
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- ▶ Adaptive mesh to maintain high resolution around strings
- ▶ 1e7 CPU-hours over ~100,000 CPUs and ~100 TB RAM
- ▶ Goal: measure axion radiation from strings to compute DM abundance



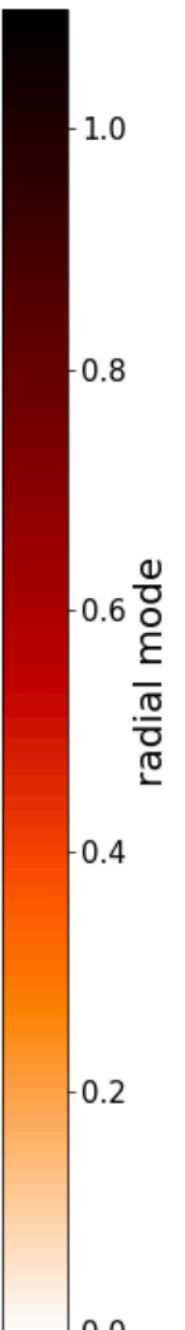
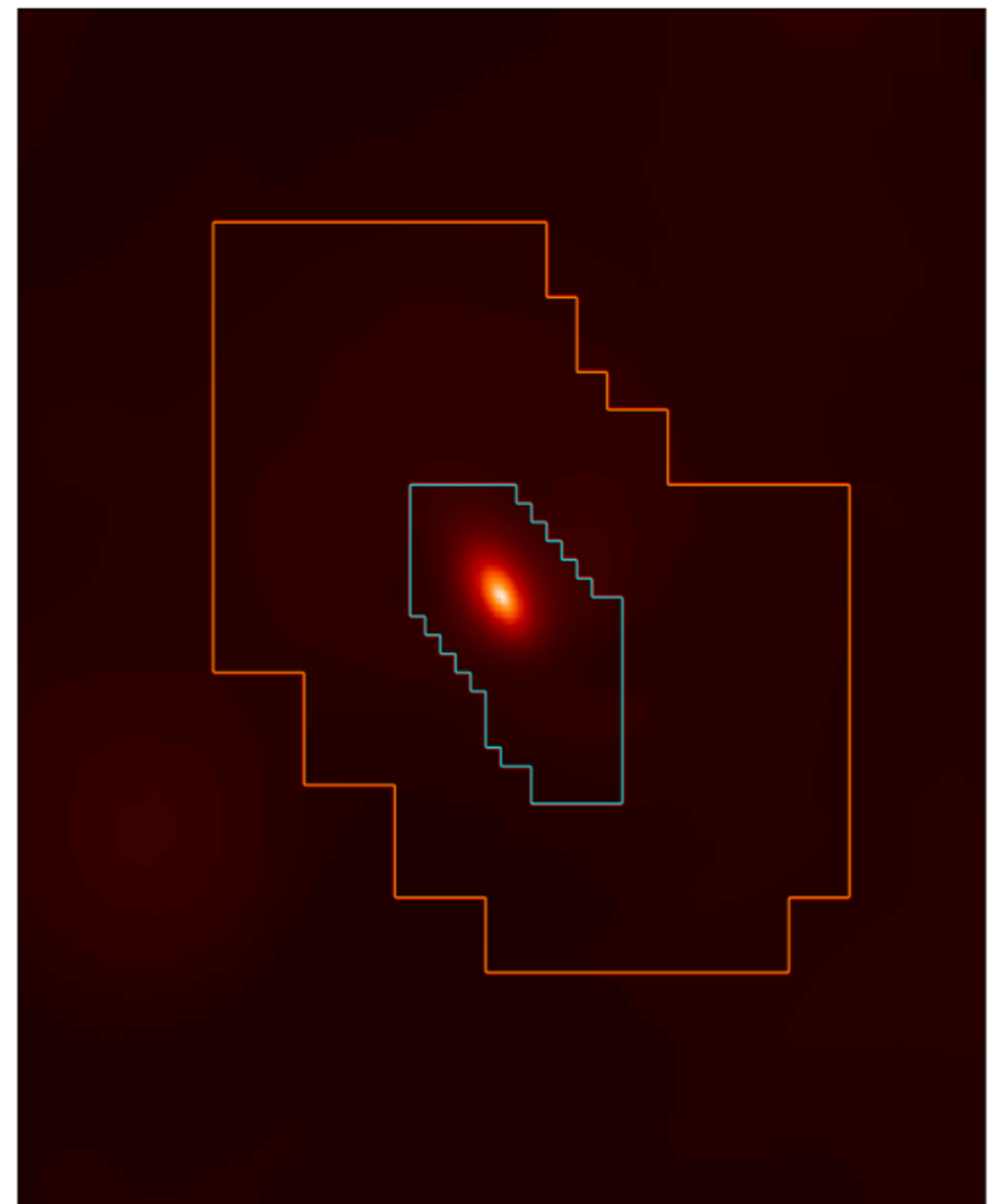
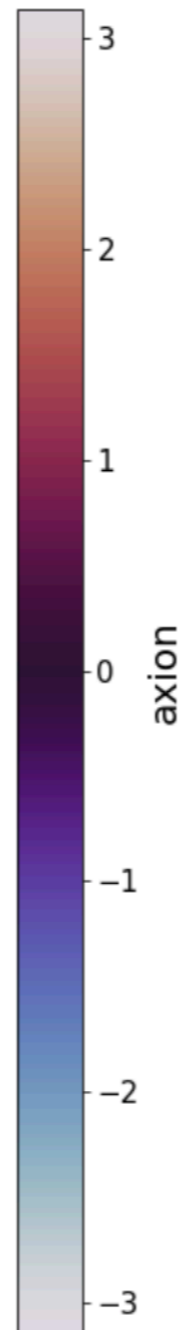
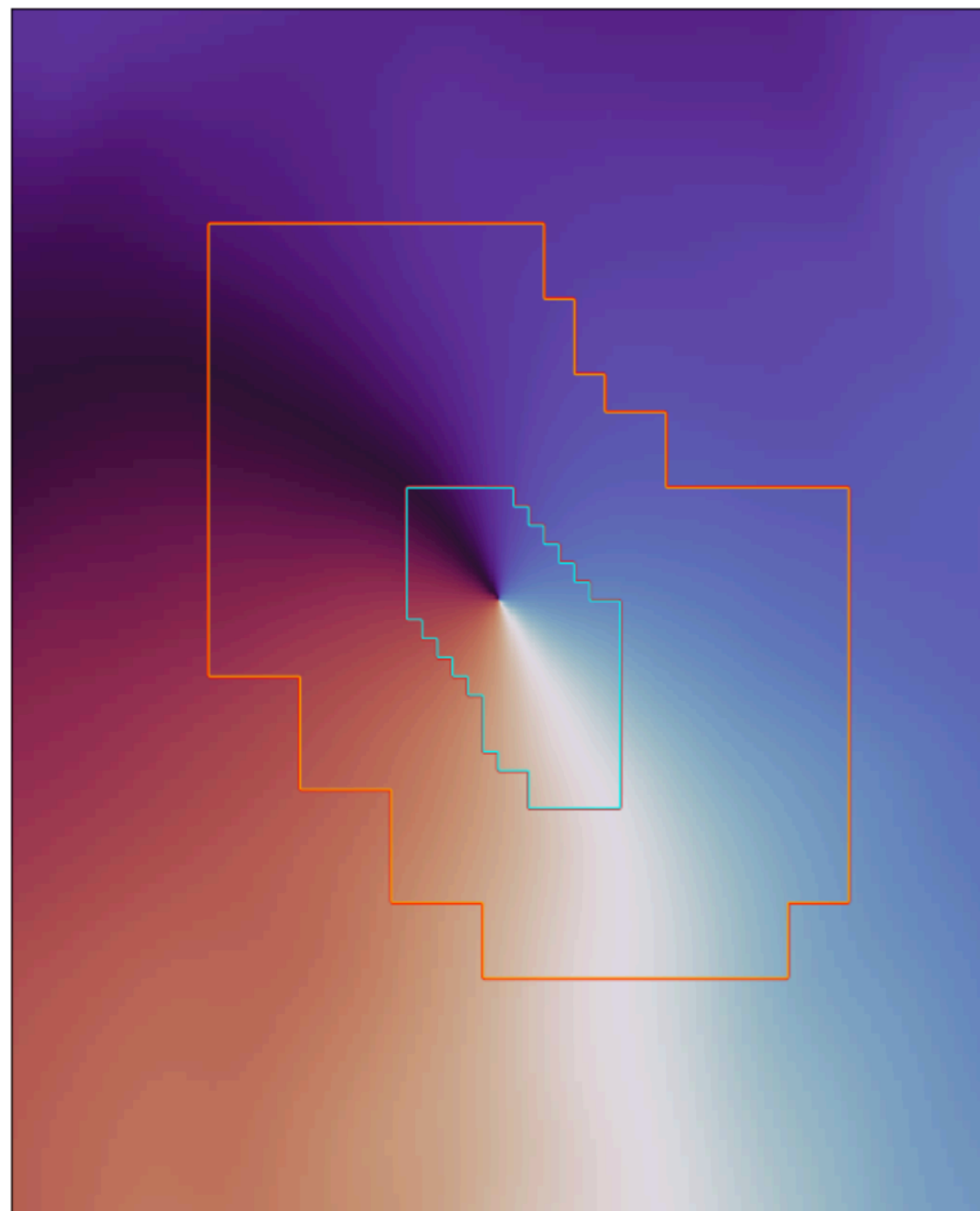
Adaptive mesh



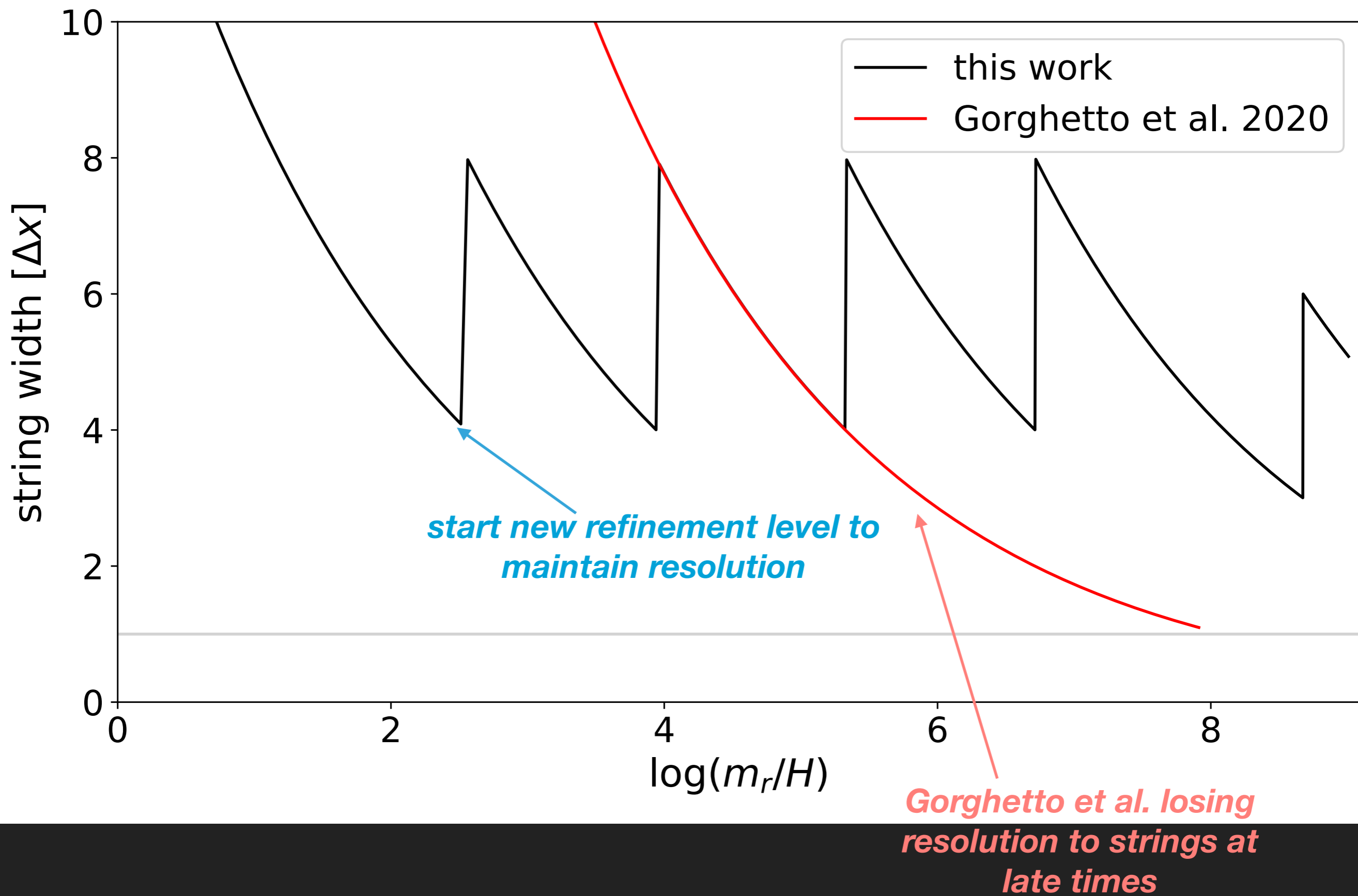
Answer with Adaptive Mesh Refinement Simulations (AMReX)

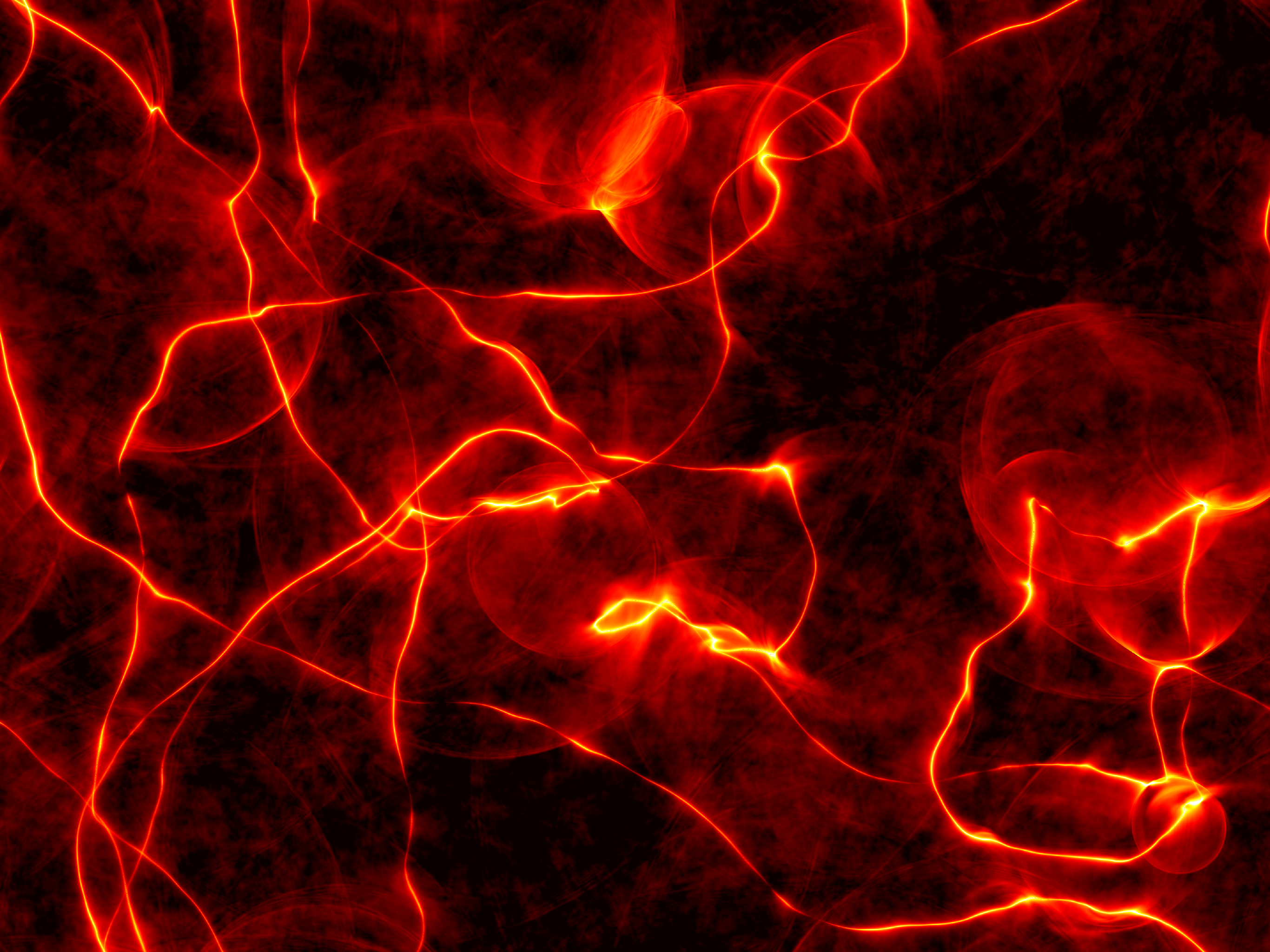


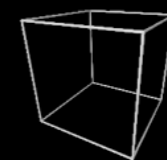
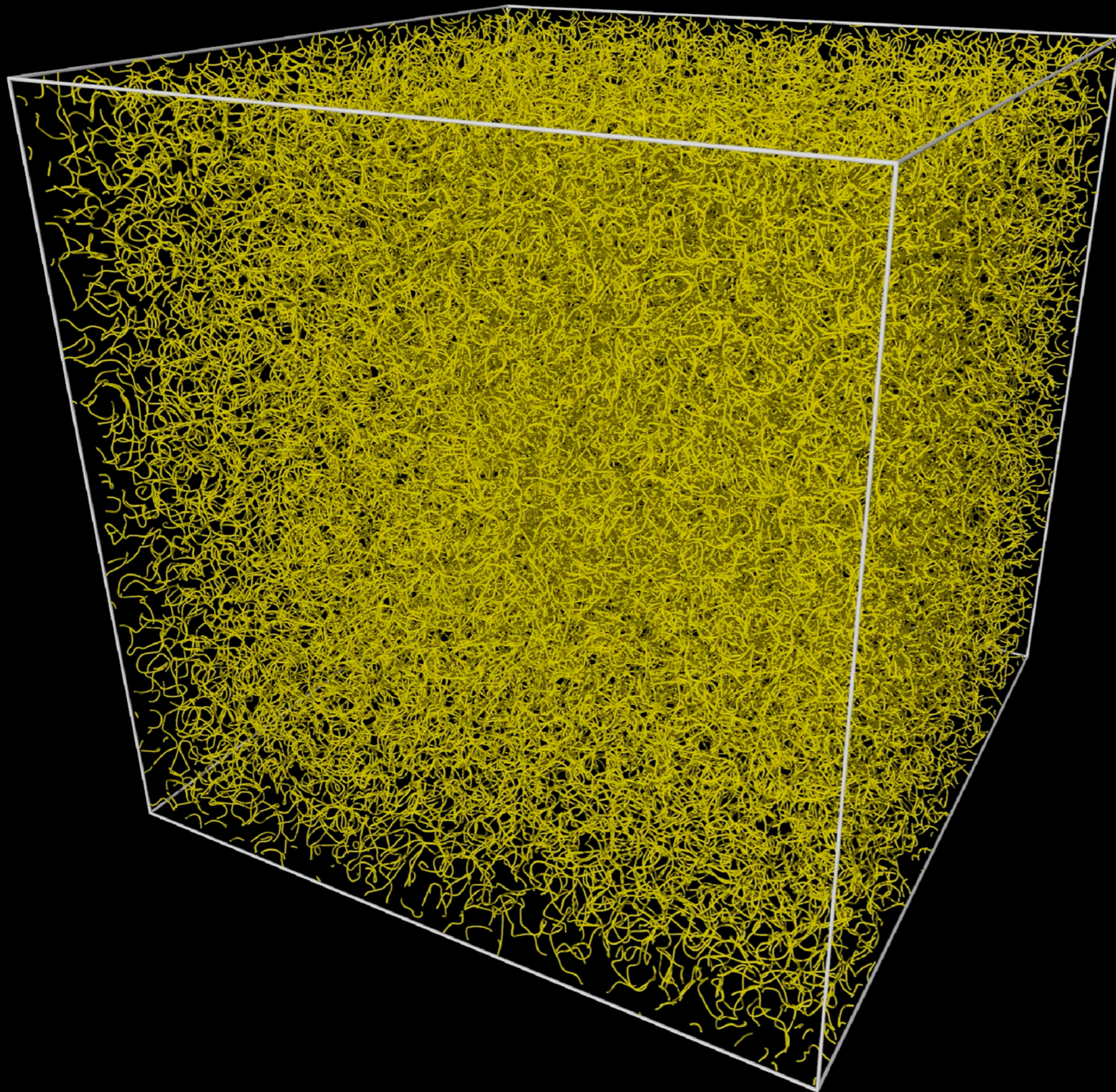
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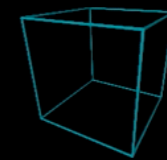
Axions with AMReX







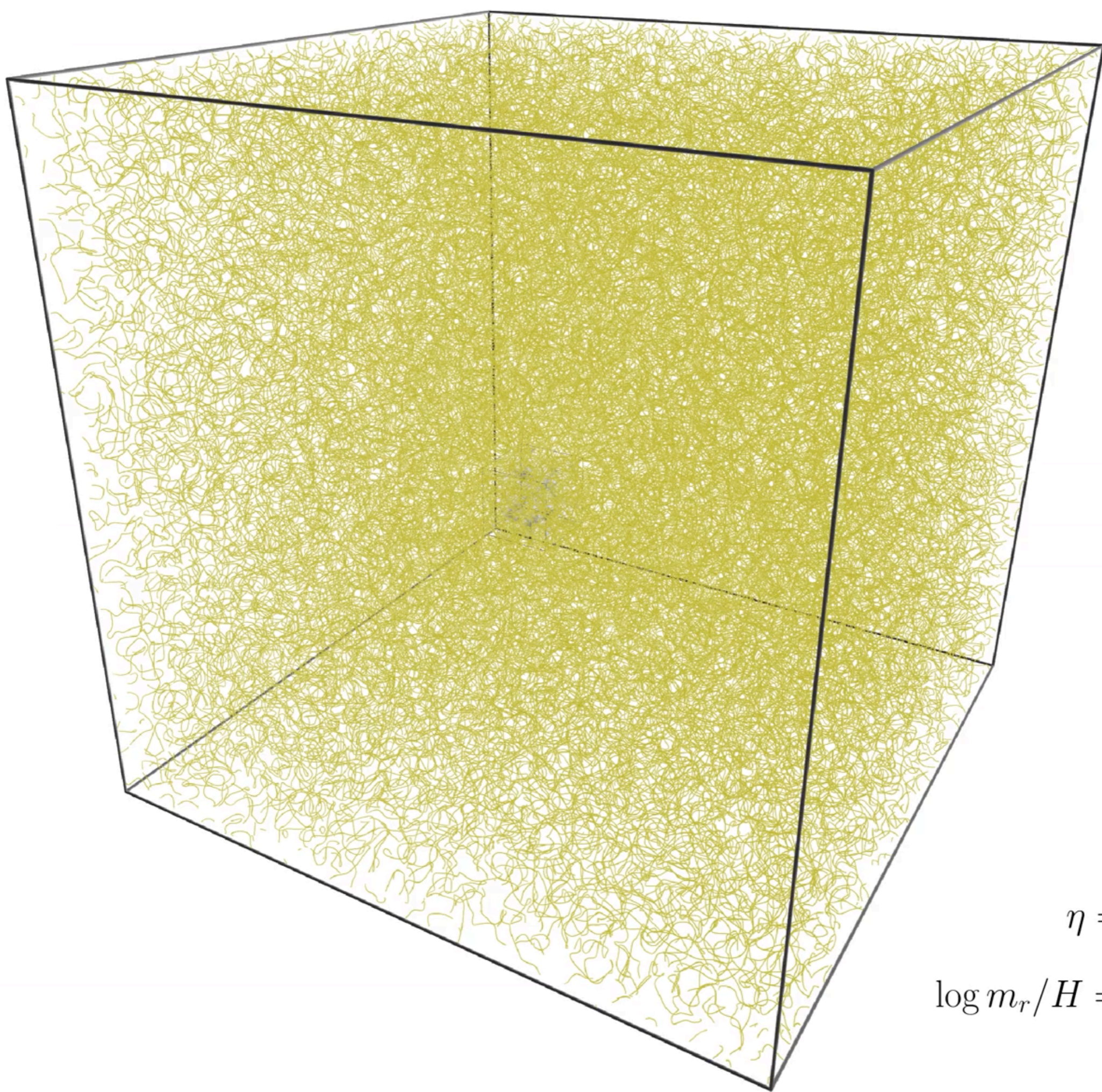
Simulation volume



Hubble volume

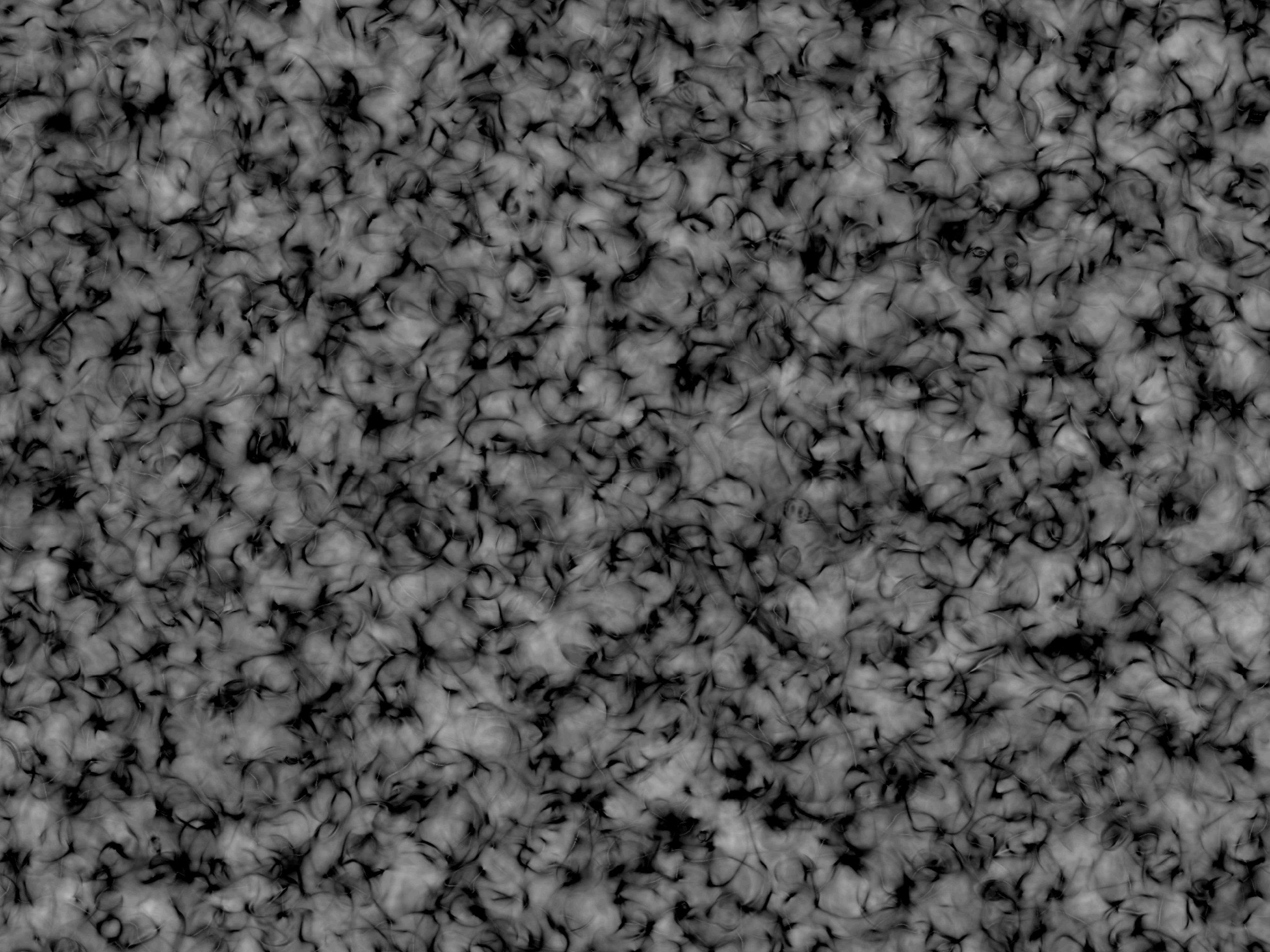
$$\eta = 0.1$$

$$\log m_r/H = -3.58$$



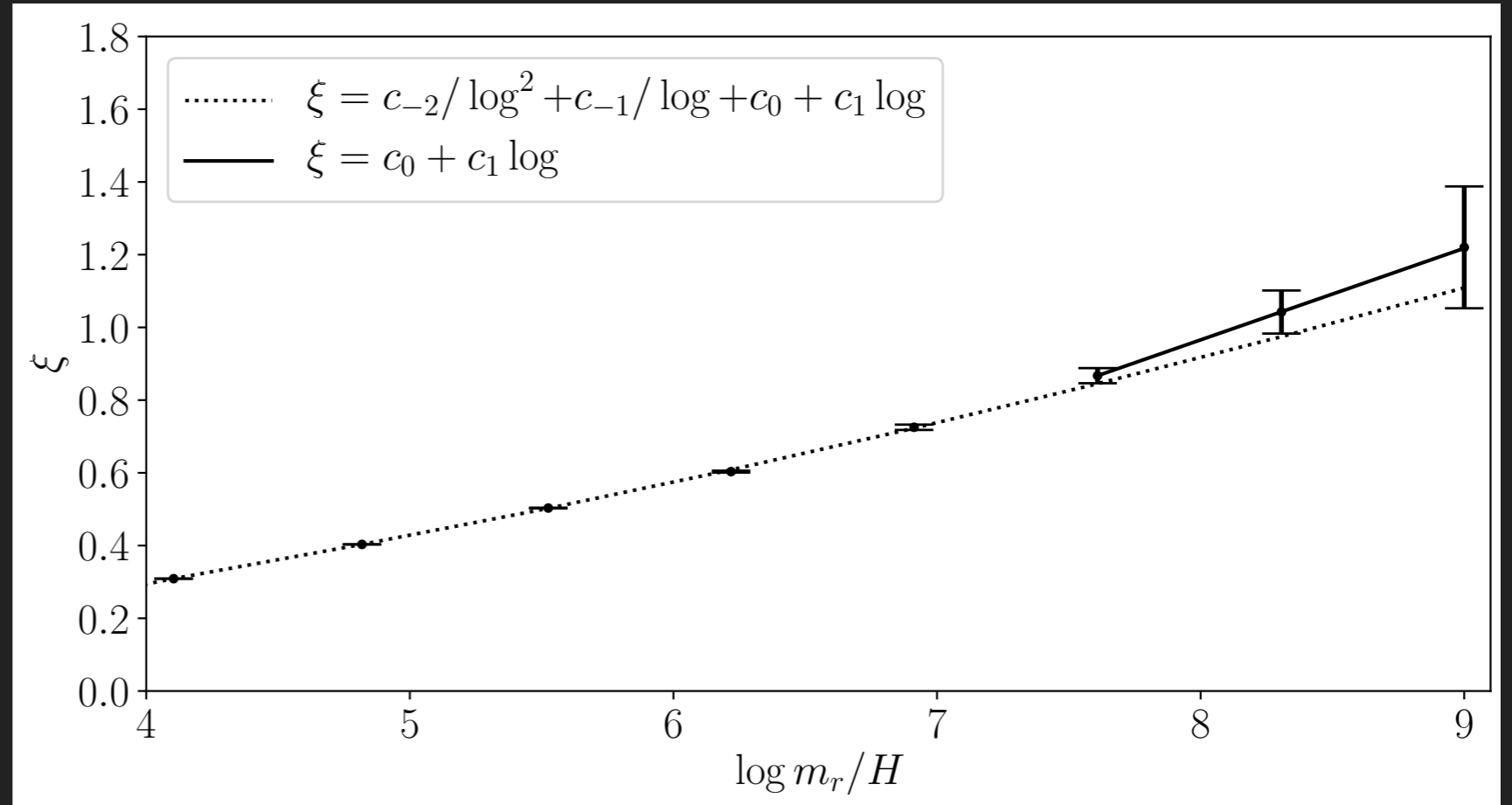
$$\eta = 0.1$$

$$\log m_r/H = -3.58$$



What do we find?

1. strings per Hubble increases logarithmically with time (and understand this analytically)



SciPost Physics

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More axions from strings

Marco Gorghetto, Edward Hardy, Giovanni Villadoro

SciPost Phys. 10, 050 (2021) · published 26 February 2021

Pointed out here, and we confirm this!

What do we find?

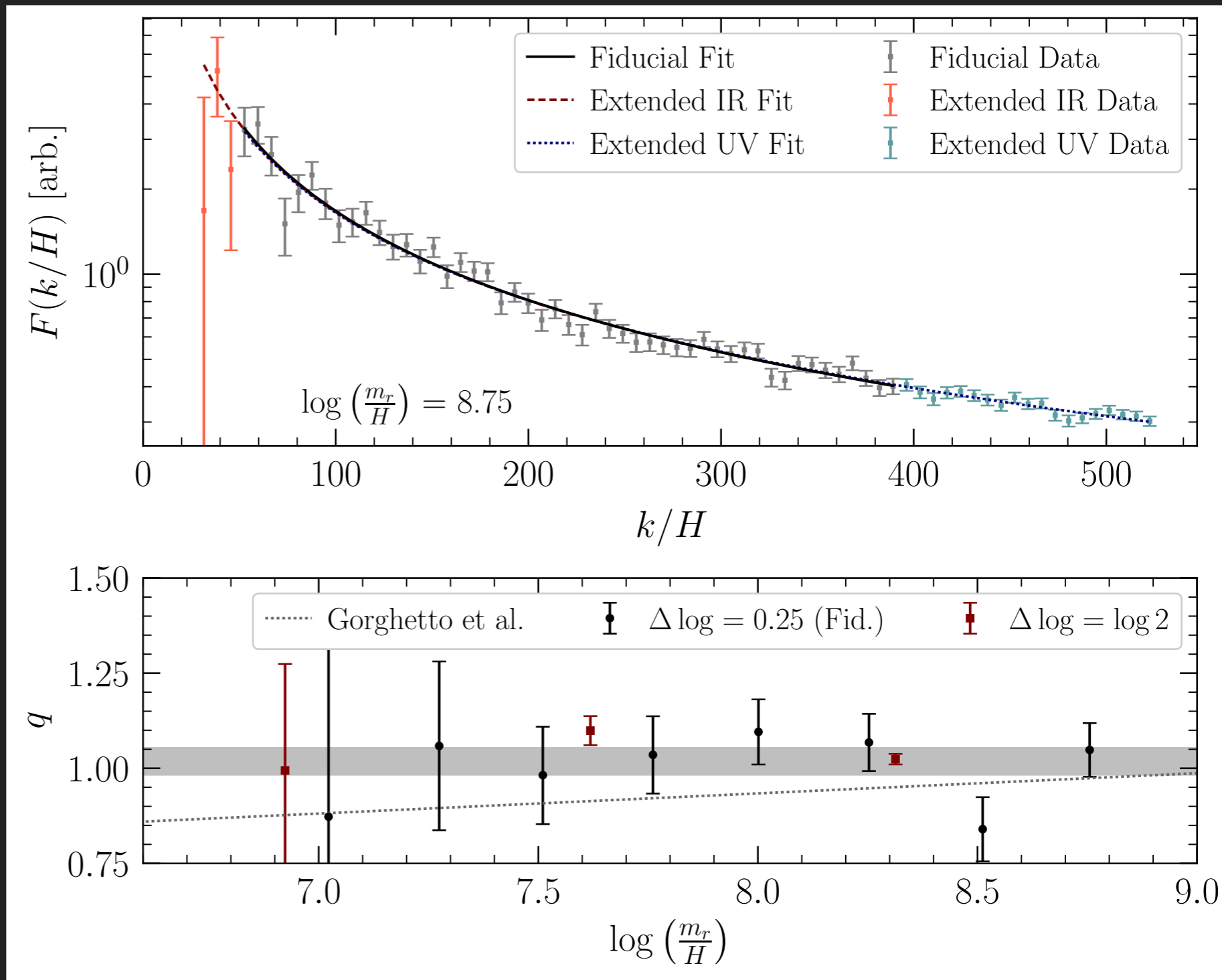
2. radiation does not become IR dominated – conformal to within few %

$$\frac{\partial \rho}{\partial k} \sim \frac{1}{k^q}$$

$$q \in (0.98, 1.04)$$



$$m_a \in (40, 180) \mu\text{eV}$$



$$m_a = 65 \pm 6 \mu\text{eV} \longleftarrow q = 1$$

What are we doing now?

Perlmutter + GPU acceleration



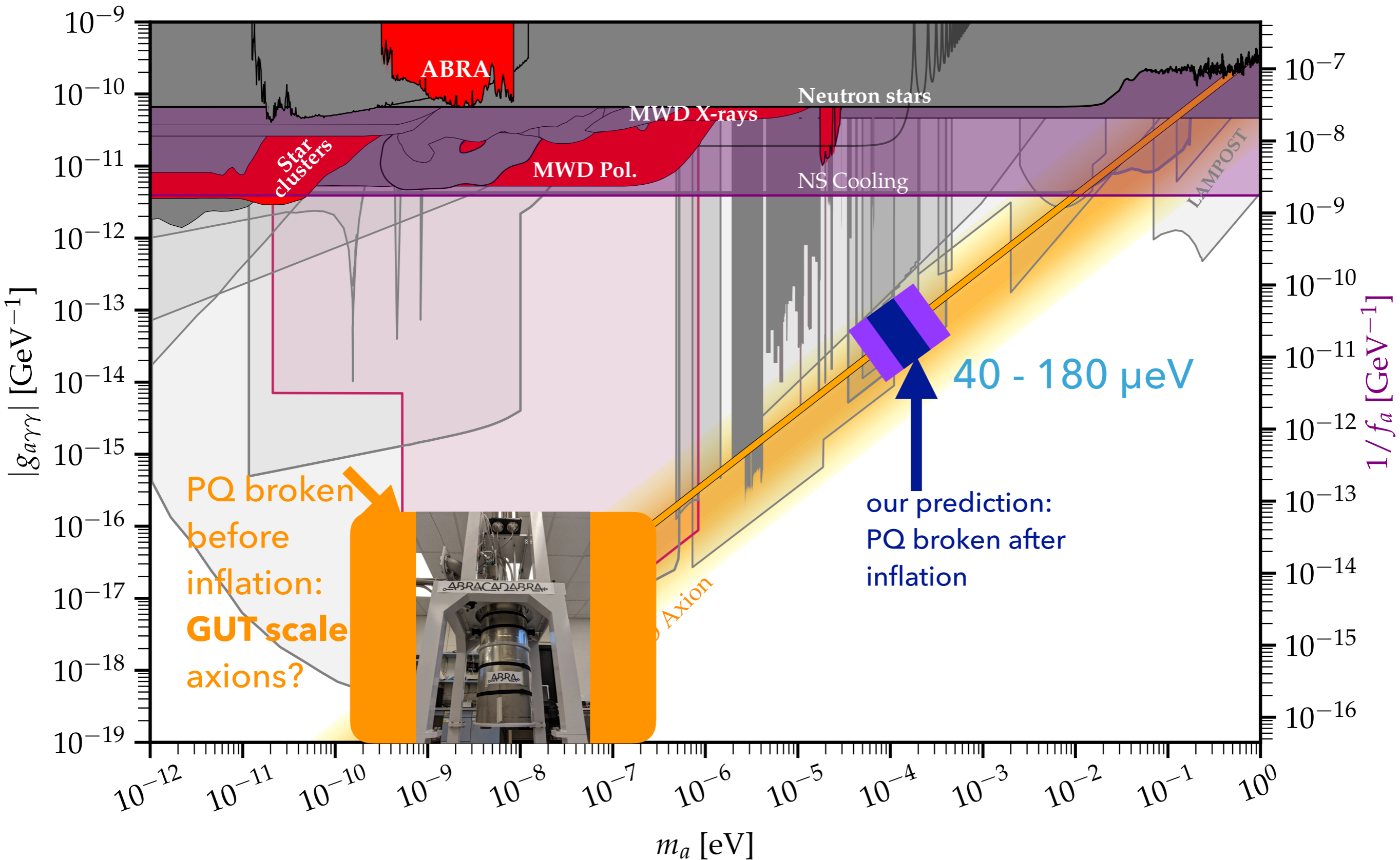
1. GPU cluster being commissioned now. Already 5th most powerful supercomputer in world

2. our plan: ~10x increase in dynamic range

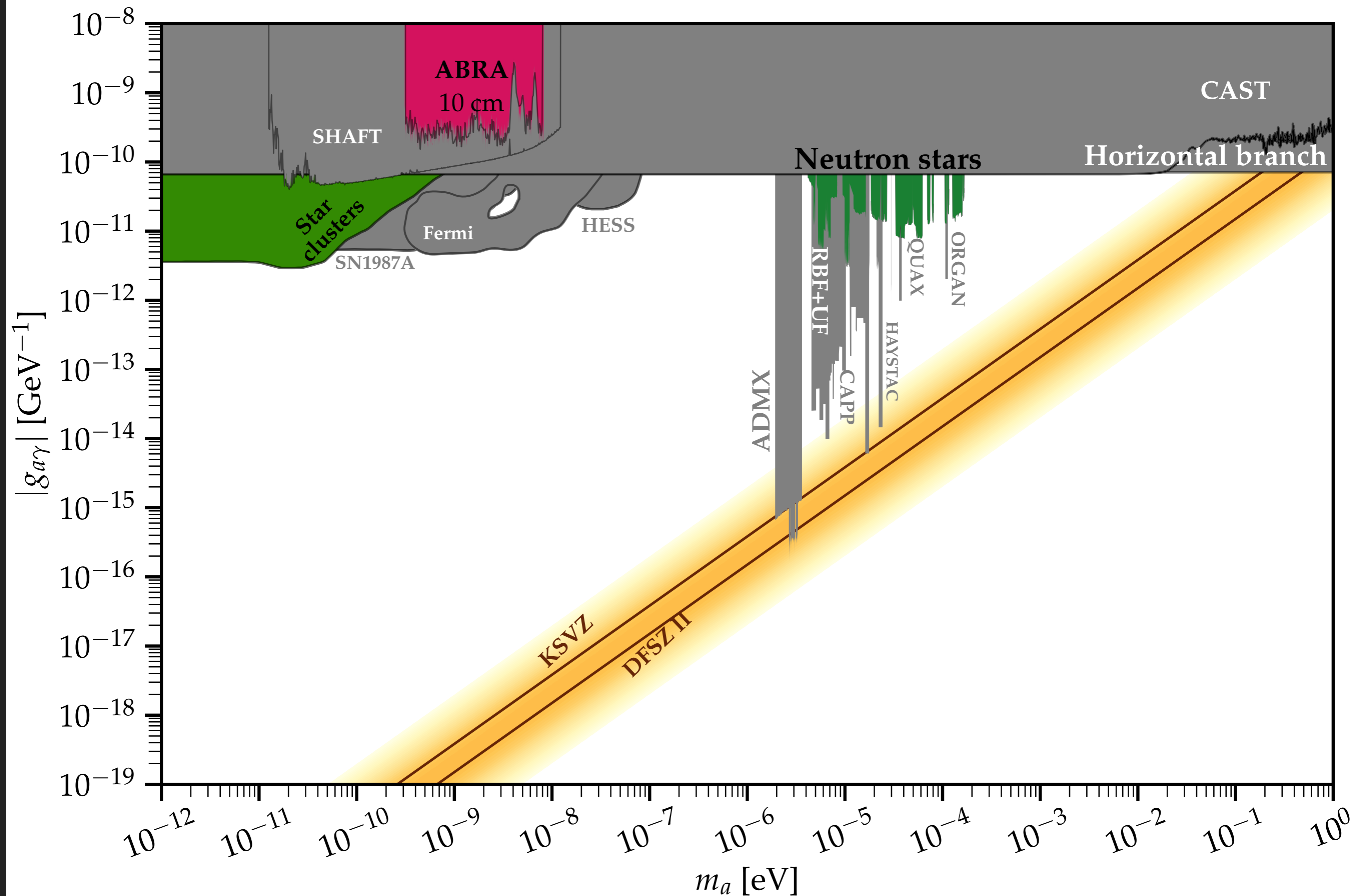
Why does this matter?

Motivated axion dark matter mass ranges

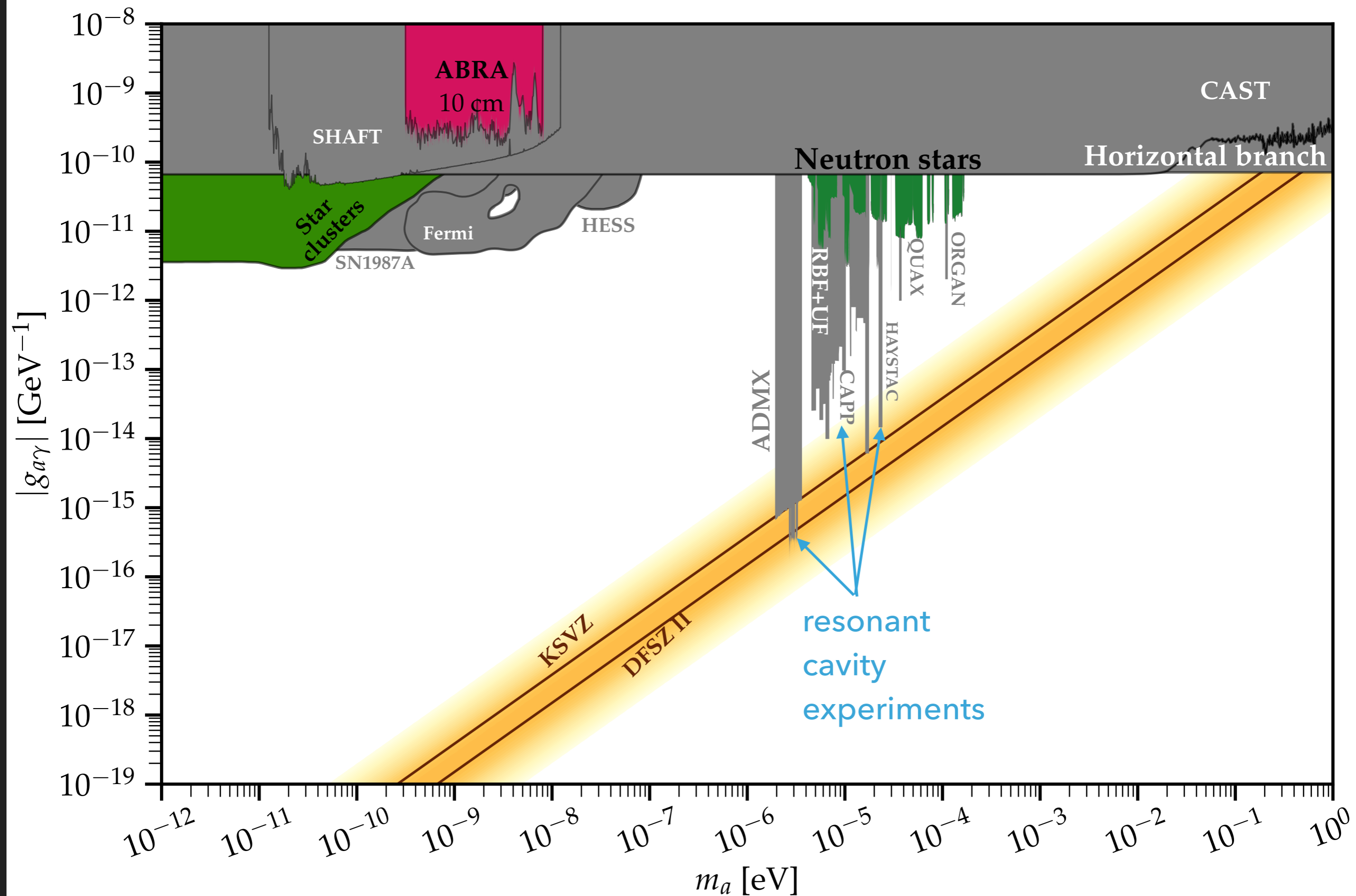
$$\mathcal{L} = g_{a\gamma\gamma} a \mathbf{E} \cdot \mathbf{B}$$



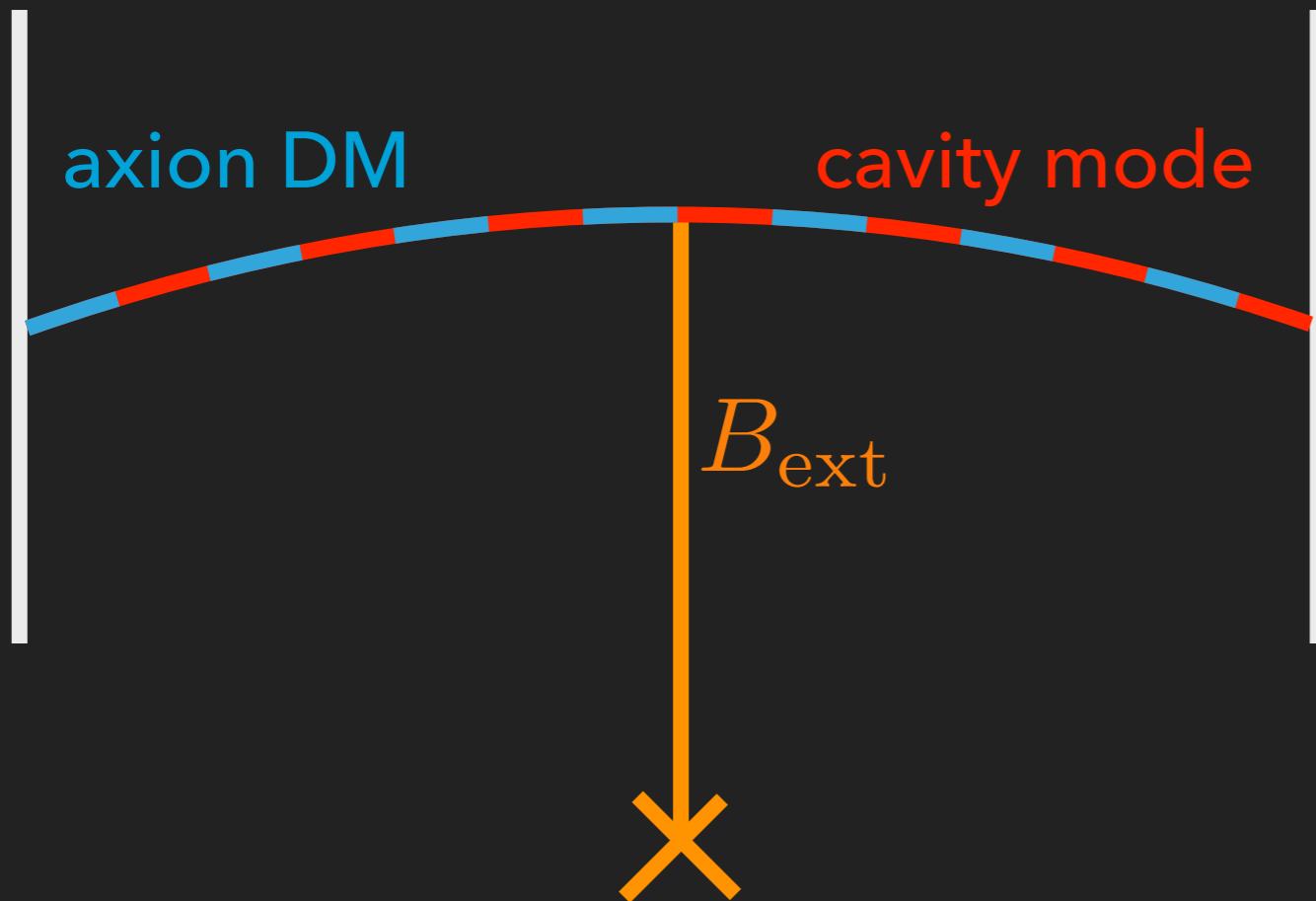
Existing Constraints: $\mathcal{L} = -g_{a\gamma\gamma} \frac{aF\tilde{F}}{4}$



Existing Constraints: $\mathcal{L} = -g_{a\gamma\gamma} \frac{aF\tilde{F}}{4}$



$$\mathcal{L} = -g_{a\gamma\gamma} \frac{aF\tilde{F}}{4} = g_{a\gamma\gamma} a \mathbf{E} \cdot \mathbf{B}$$



1. Axion coherence time limited by

$$\delta E/E \sim \left(\frac{v_{\text{DM}}}{c}\right)^2 \sim 10^{-6}$$

2. $\rightarrow Q \sim 10^6$ cavities ideal

3. $m_a = 25 \mu\text{eV} \rightarrow 5 \text{ cm cavity}$

conversion probability – $p_{a \rightarrow \gamma} = g_{a\gamma\gamma}^2 B_{\text{ext}}^2 L^2$

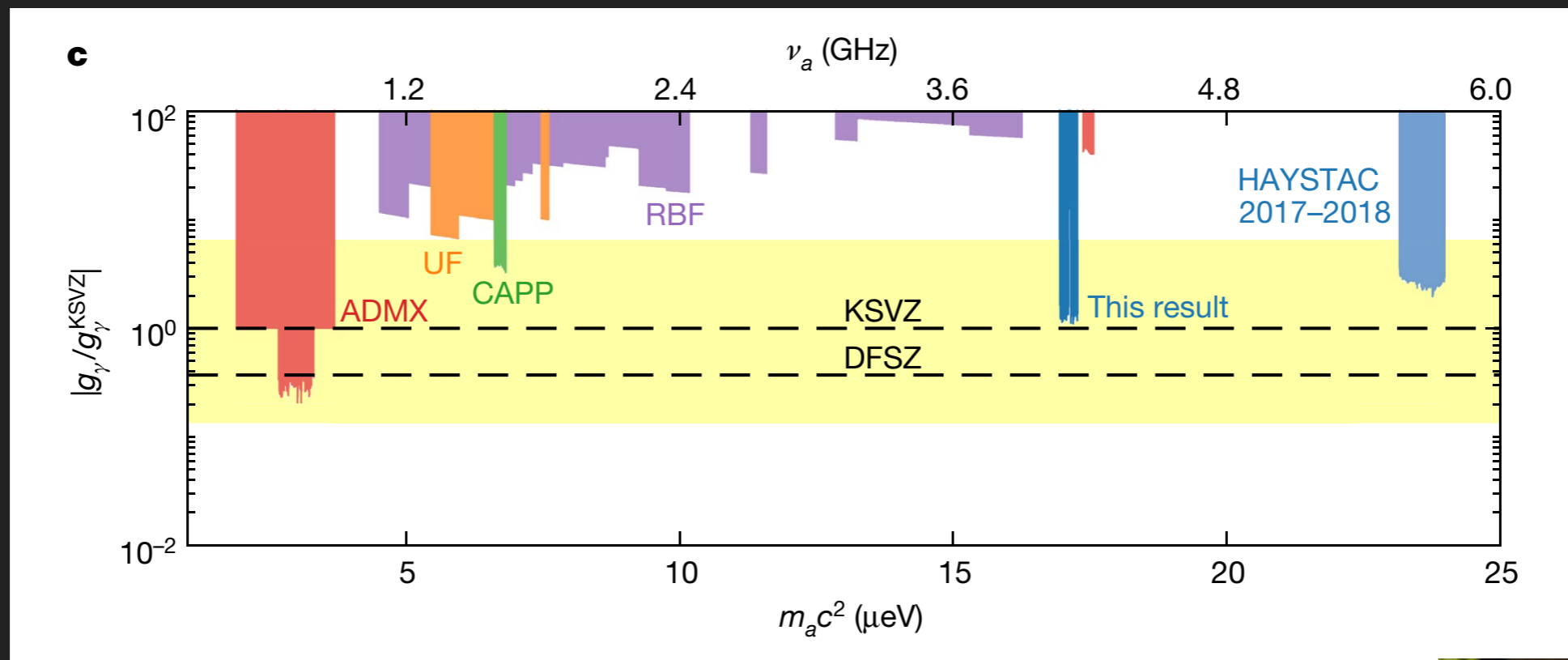
cavity: converted power – $P_{a \rightarrow \gamma} \propto \frac{\rho_{\text{DM}}}{m_a} g_{a\gamma\gamma}^2 B_{\text{ext}}^2 V Q$

Example: HAYSTAC Experiment

Article

Nature 2021

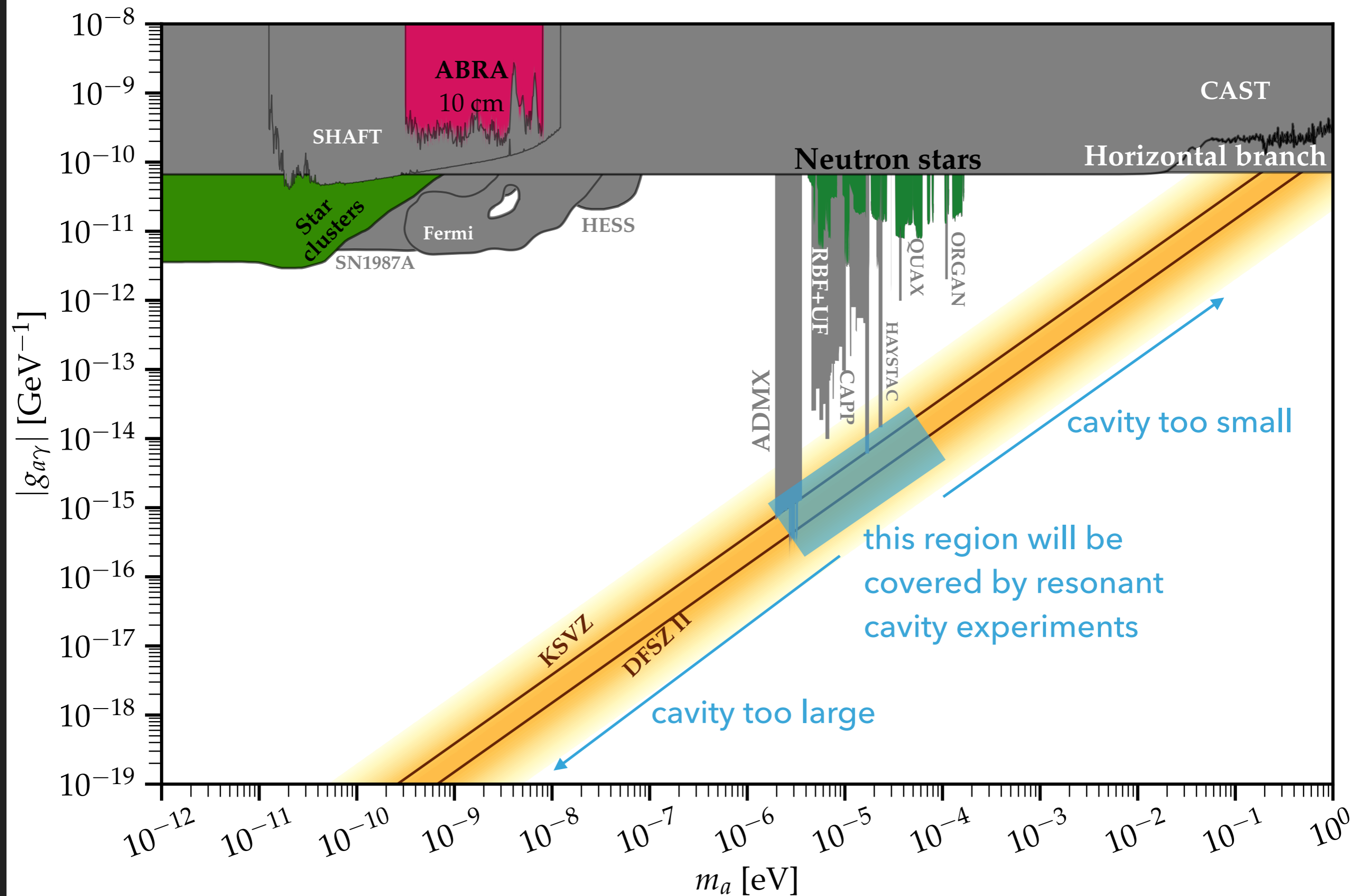
A quantum enhanced search for dark matter axions



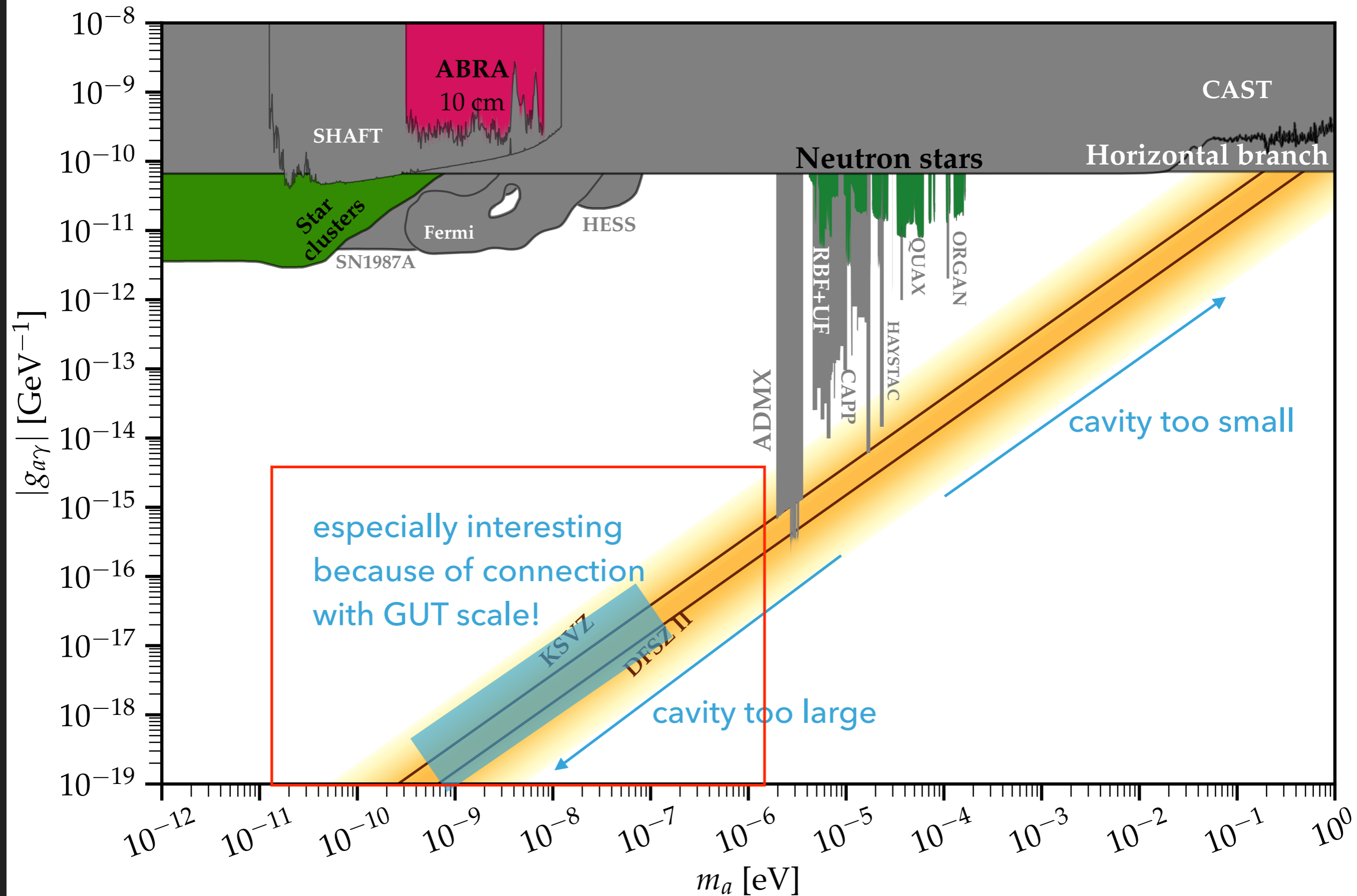
Used quantum squeezed states in cavity to be more sensitive to axion signals



Going beyond the resonant cavity



Going beyond the resonant cavity



A Broadband / Resonant Approach to Cosmic Axion Detection with an Amplifying B-field Ring Apparatus

Y. Kahn, B.S., J. Thaler (PRL 2016)



Axion wavelength much longer than size of detector



Detecting axion dark matter beyond the magnetoquasistatic approximation
 Joshua N. Benabou,^{1,2} Joshua W. Foster,³ Yonatan Kahn,⁴ Benjamin R. Safdi,^{1,2} and Chiara P. Salemi^{5,6,7}

$a(t)$

$(m_a t)$

$$\frac{1}{2} m_a^2 a_0^2 = \rho_{\text{DM}}$$

$$\nabla \times \mathbf{B} = \frac{\partial \mathbf{E}}{\partial t} - g_{a\gamma\gamma} \left(\mathbf{E} \times \nabla a - \mathbf{B} \frac{\partial a}{\partial t} \right)$$

$$\nabla \cdot \mathbf{E} = -g_{a\gamma\gamma} \nabla a \cdot \mathbf{B}$$

Suppressed by DM velocity
 $v/c \sim 0.001$

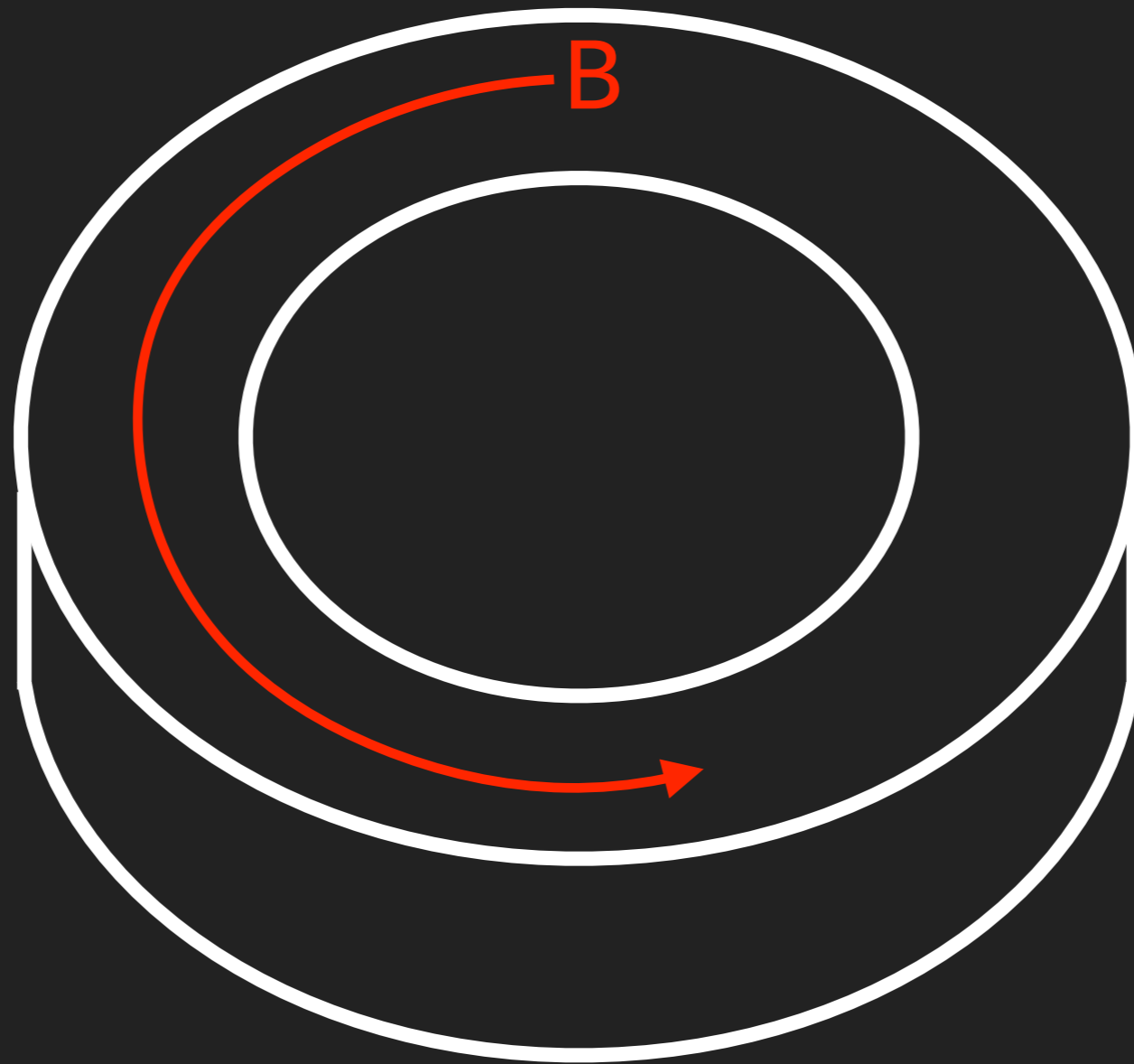
But see
B.S. et al.
 2022!

$$\nabla \times \mathbf{E} = -\frac{\partial \mathbf{B}}{\partial t}$$

Drop these terms in MQS
 approximation – long axion
 Compton wavelength

$$\nabla \cdot \mathbf{B} = 0$$

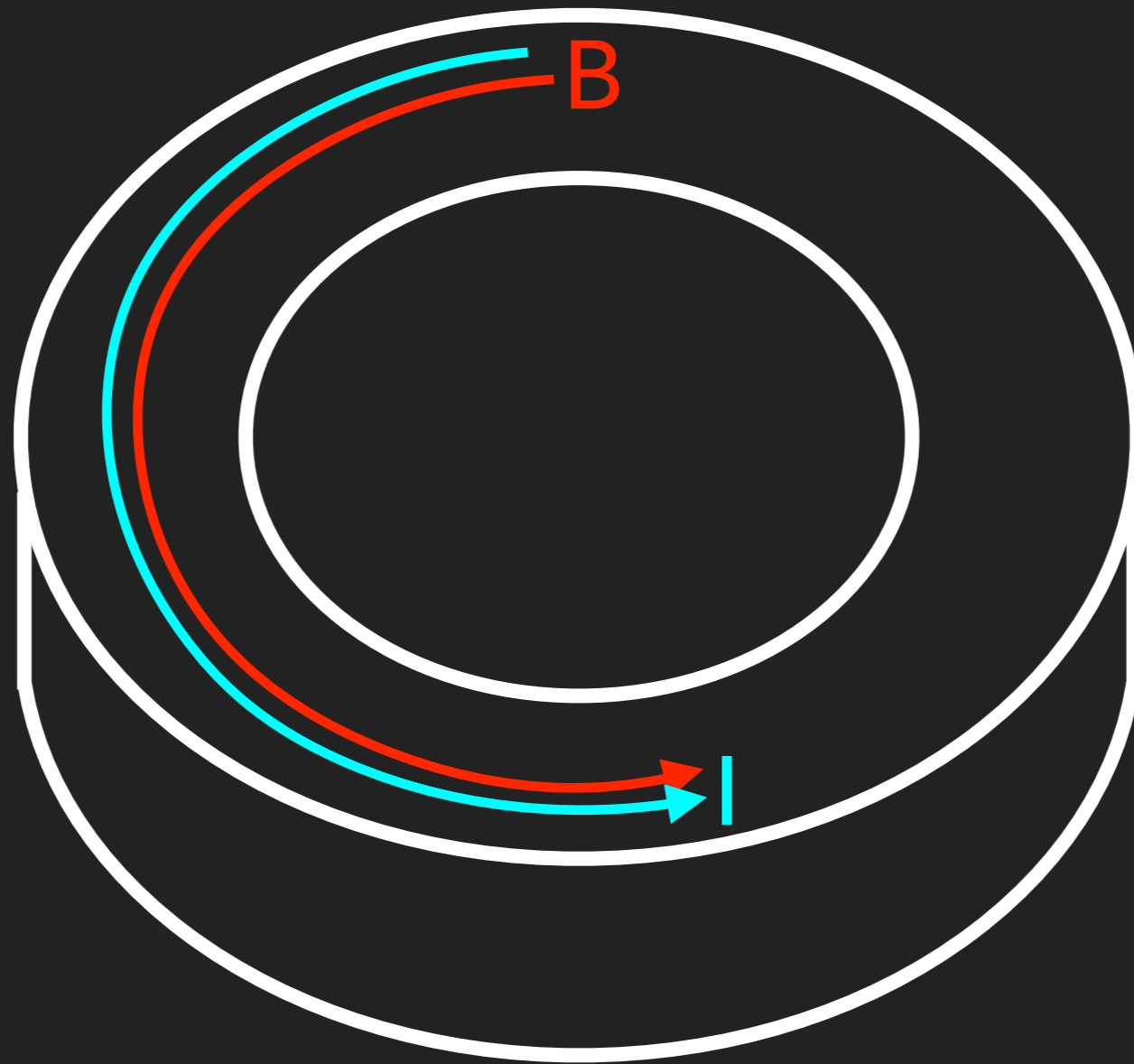
Toroidal Magnetic Field: \mathbf{B}



$$\mathcal{L} = g_{a\gamma\gamma} a \mathbf{E} \cdot \mathbf{B}$$

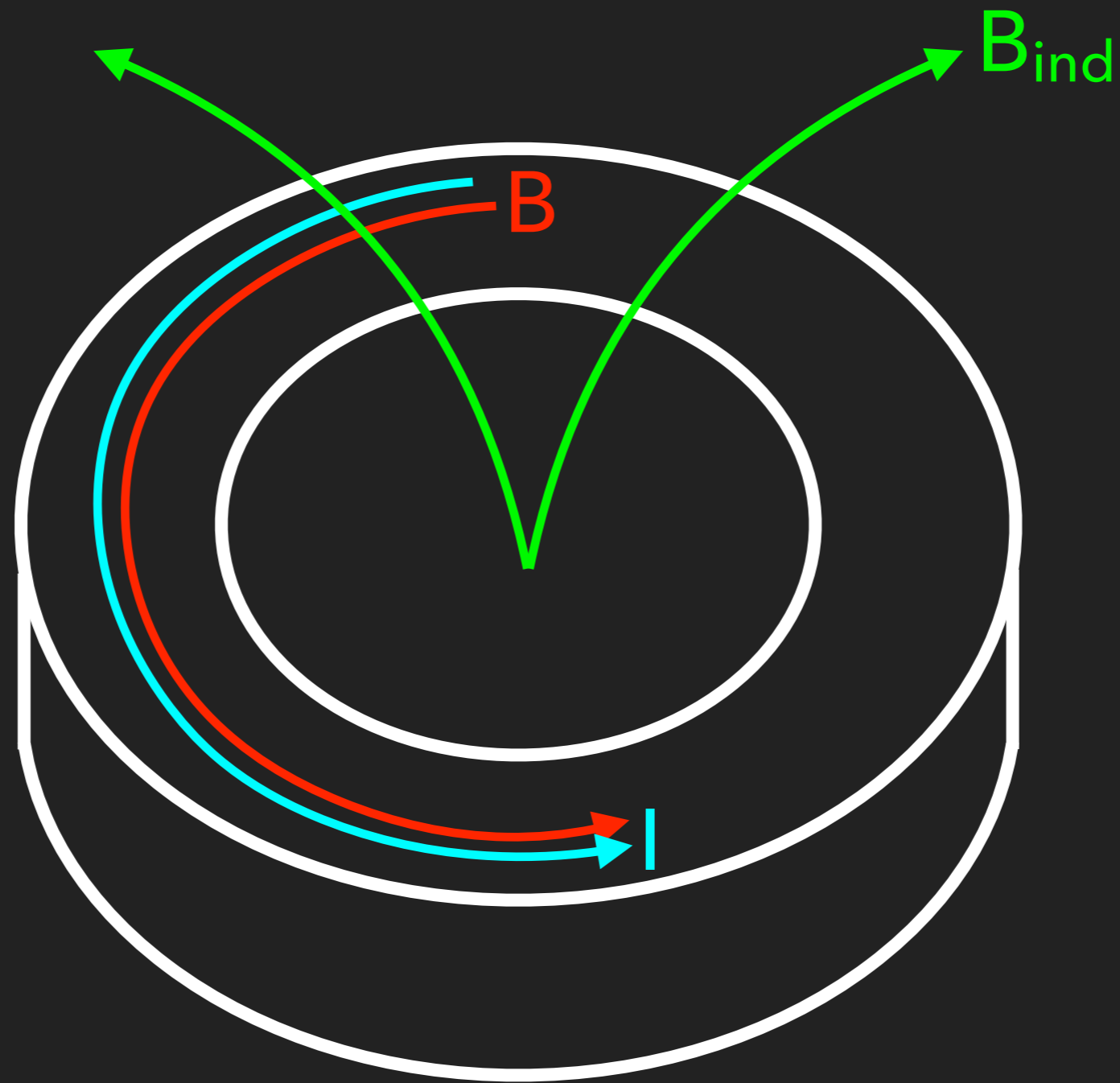
$m_a^{-1} \gg$ size of experiment: $\nabla \times \mathbf{B} = g_{a\gamma\gamma} \mathbf{B} \frac{\partial a}{\partial t}$

Axion Effective Electric Current: I



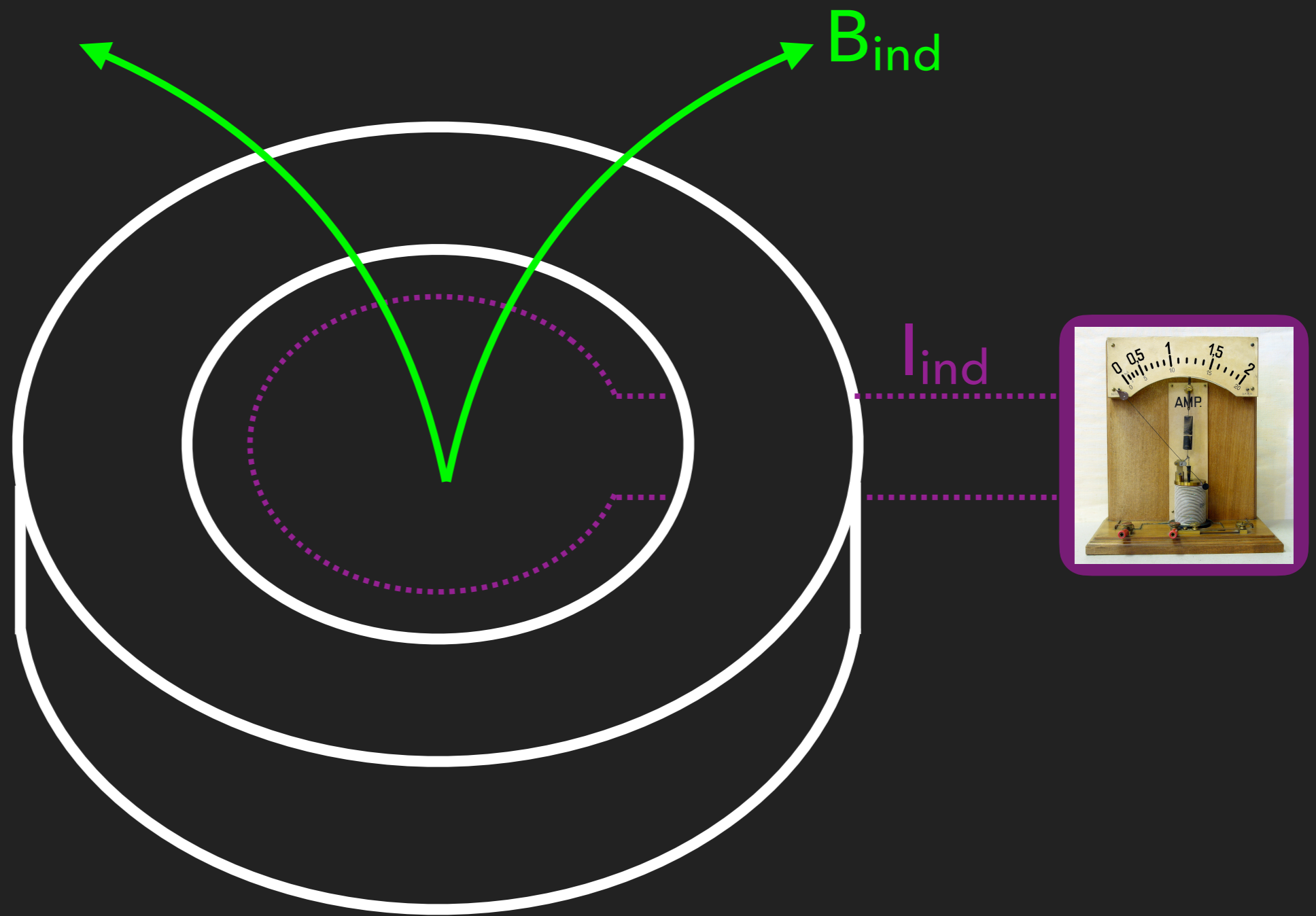
$$\mathcal{L} = g_{a\gamma\gamma} a \mathbf{E} \cdot \mathbf{B}$$

$m_a^{-1} \gg \text{size of experiment: } \nabla \times \mathbf{B} = g_{a\gamma\gamma} \mathbf{B} \frac{\partial a}{\partial t}$



Secondary axion-induced B-field: B_{ind}

$m_a^{-1} \gg \text{size of experiment: } \nabla \times \mathbf{B} = g_{a\gamma\gamma} \mathbf{B} \frac{\partial a}{\partial t}$



Pickup-loop current: I_{ind}

$m_a^{-1} \gg \text{size of experiment: } \nabla \times \mathbf{B} = g_{a\gamma\gamma} \mathbf{B} \frac{\partial a}{\partial t}$

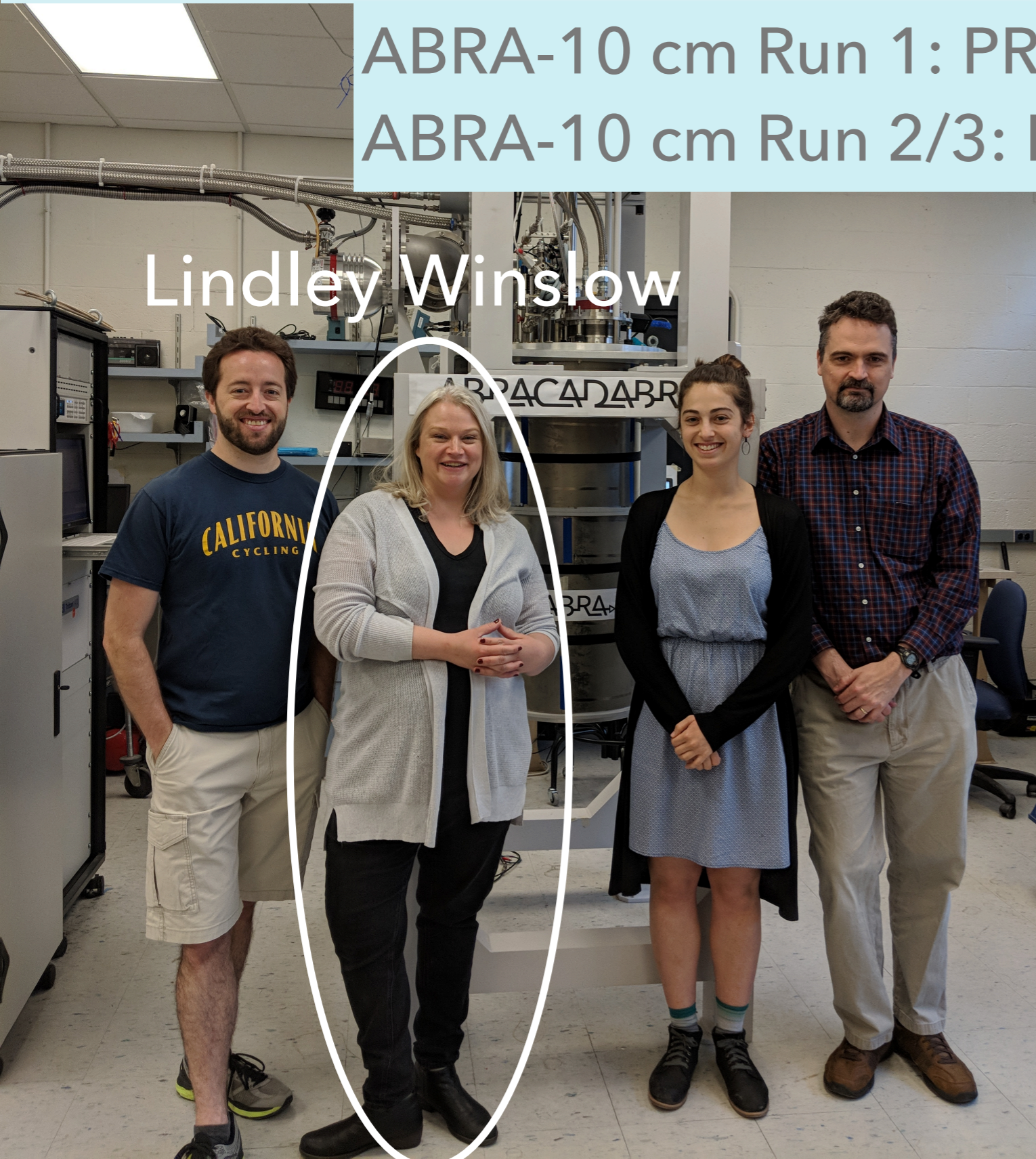


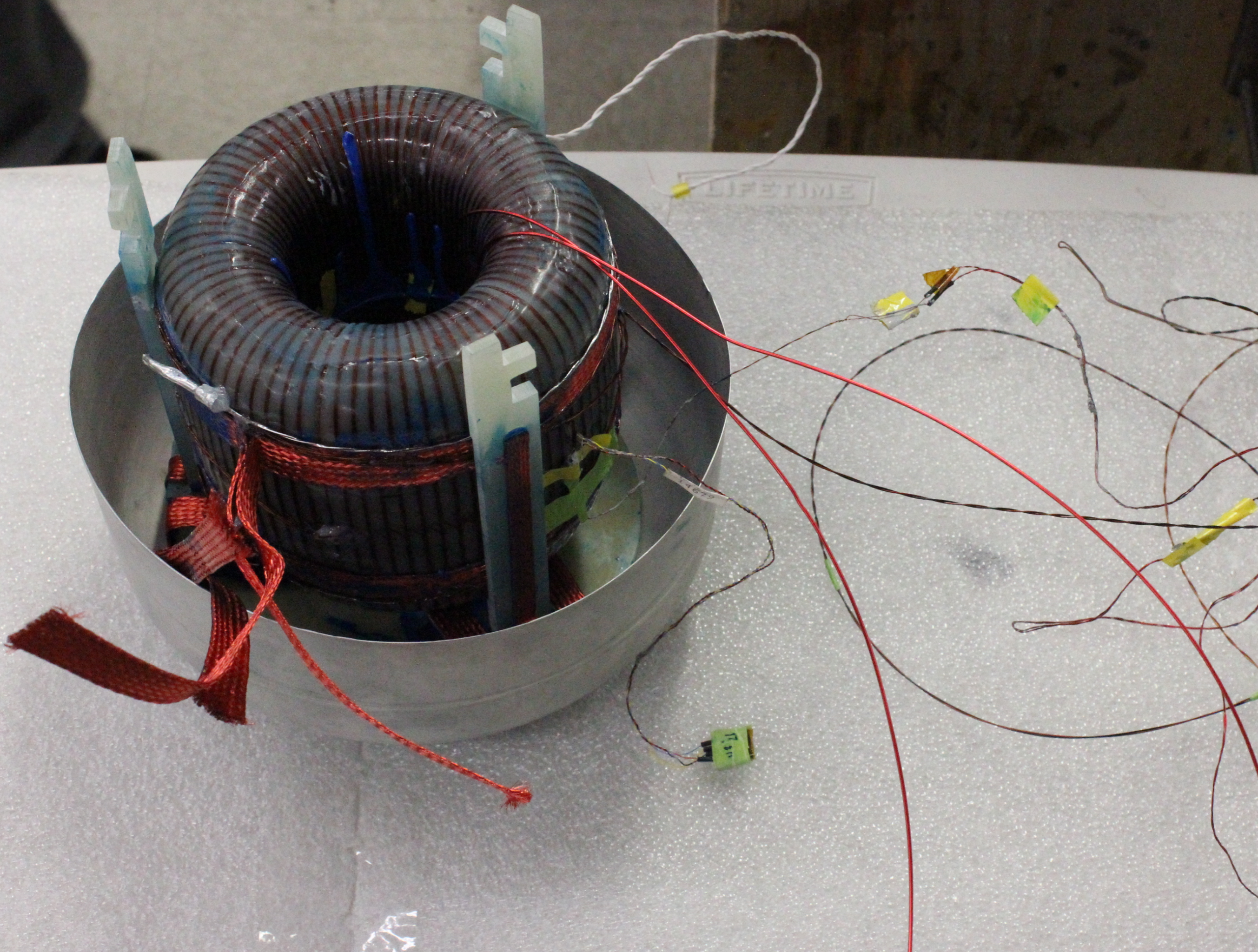
THE UNIVERSITY
of NORTH CAROLINA
at CHAPEL HILL



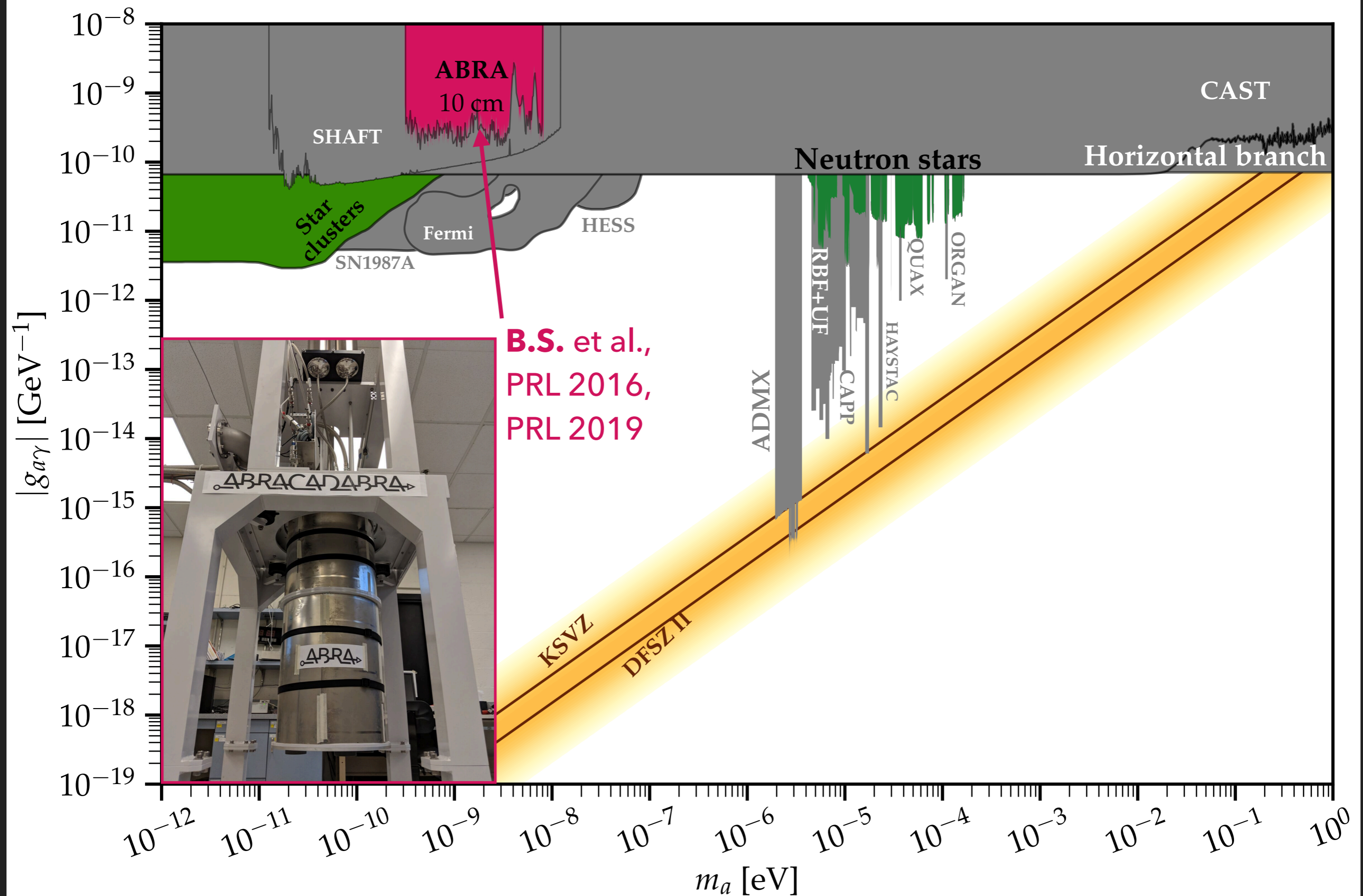
ABRA-10 cm Run 1: PRL 2018, PRD 2019
ABRA-10 cm Run 2/3: PRL 2021

Lindley Winslow

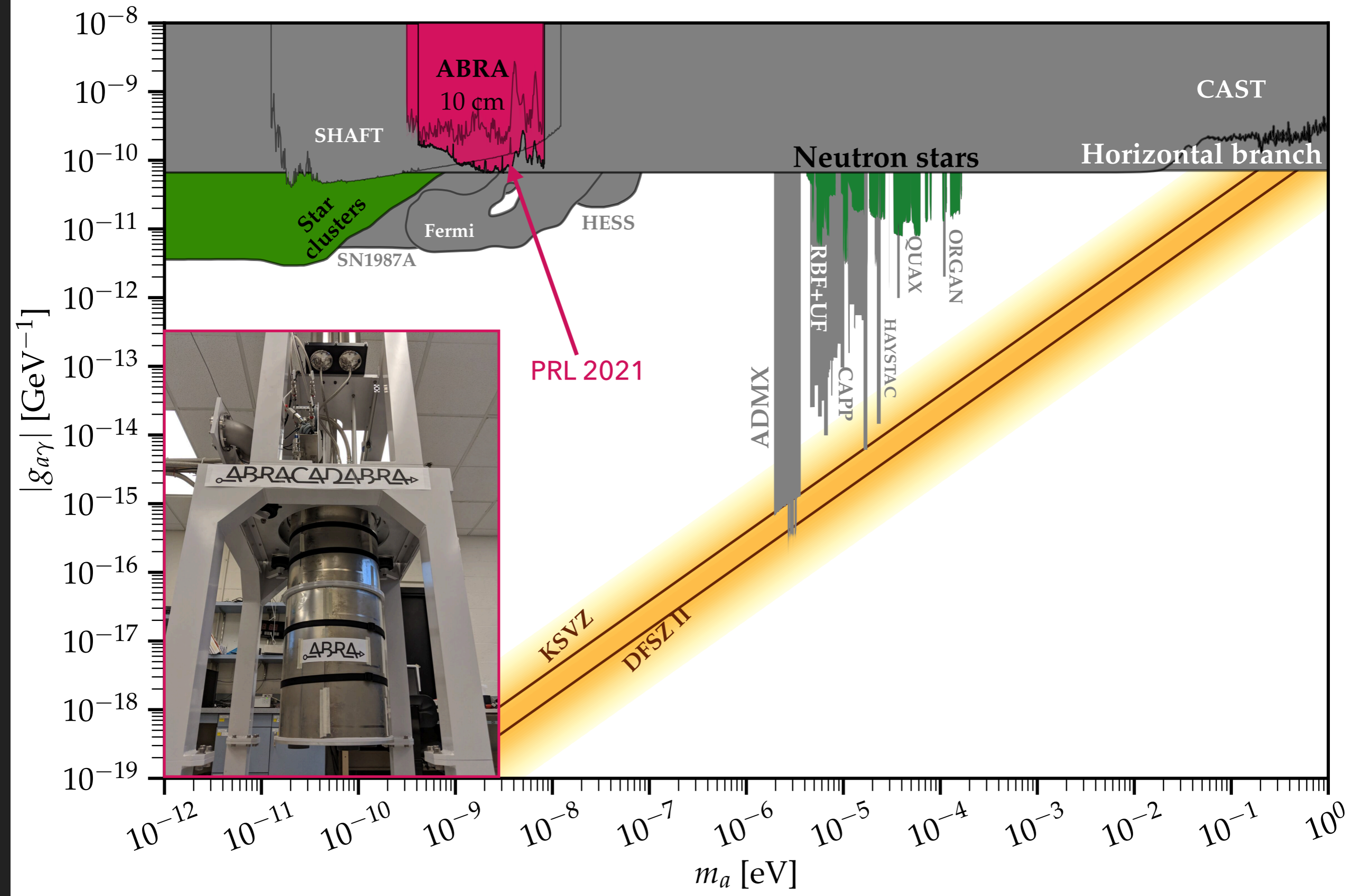




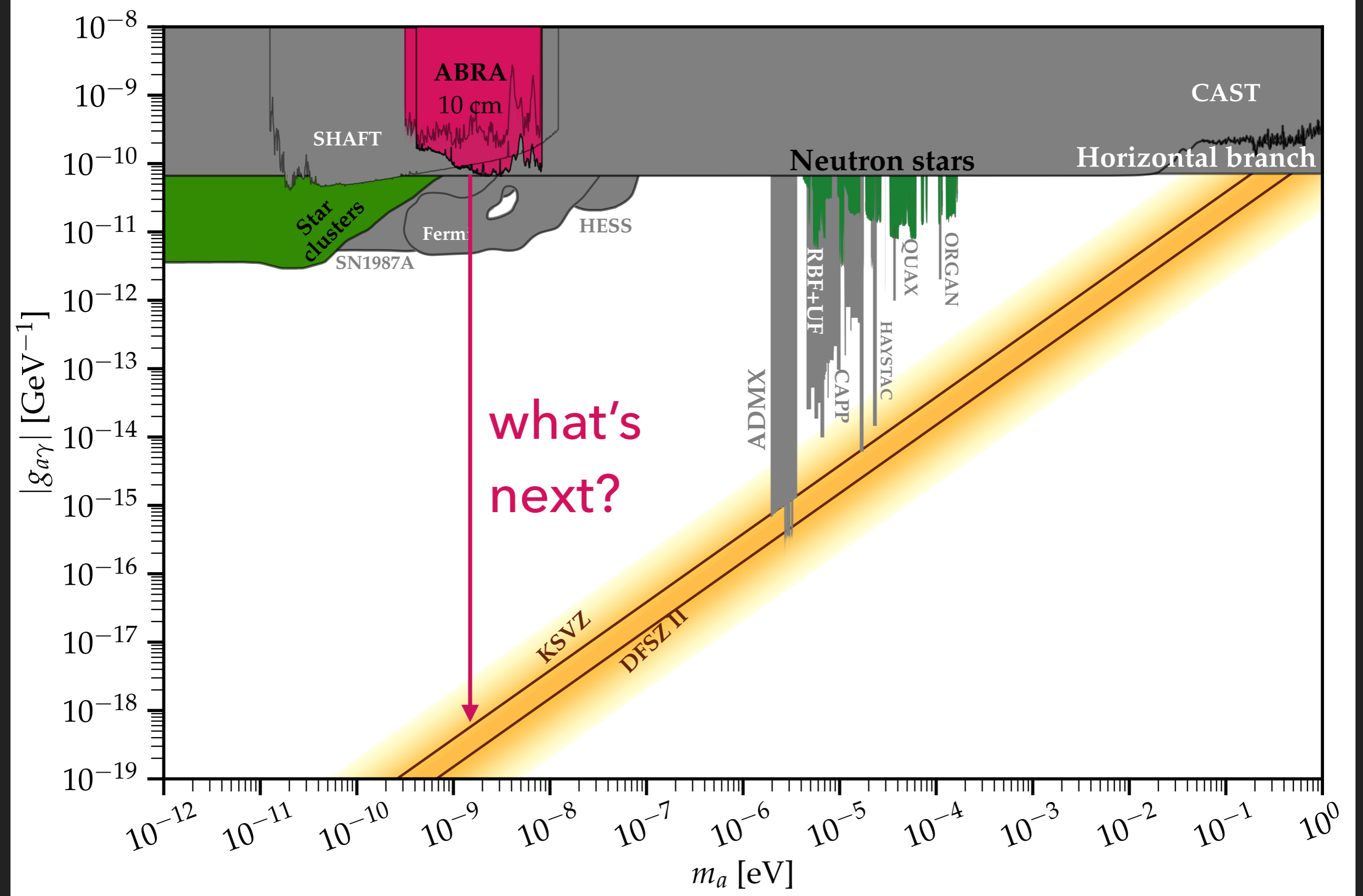
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$$\mathcal{L} = -g_{a\gamma\gamma} \frac{aF\tilde{F}}{4}$$



DMRadio Collaboration

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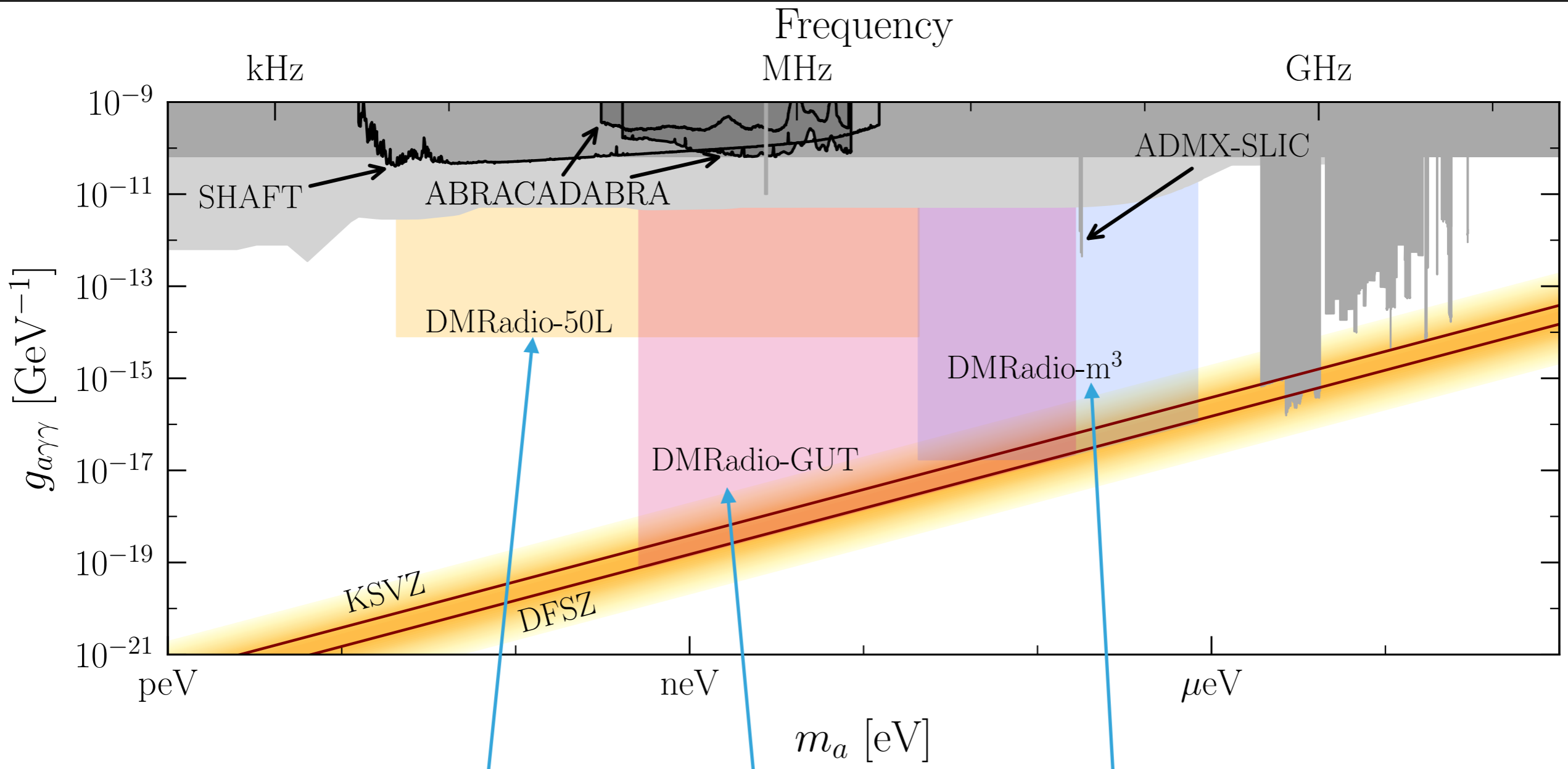
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DMRadio Science Plan

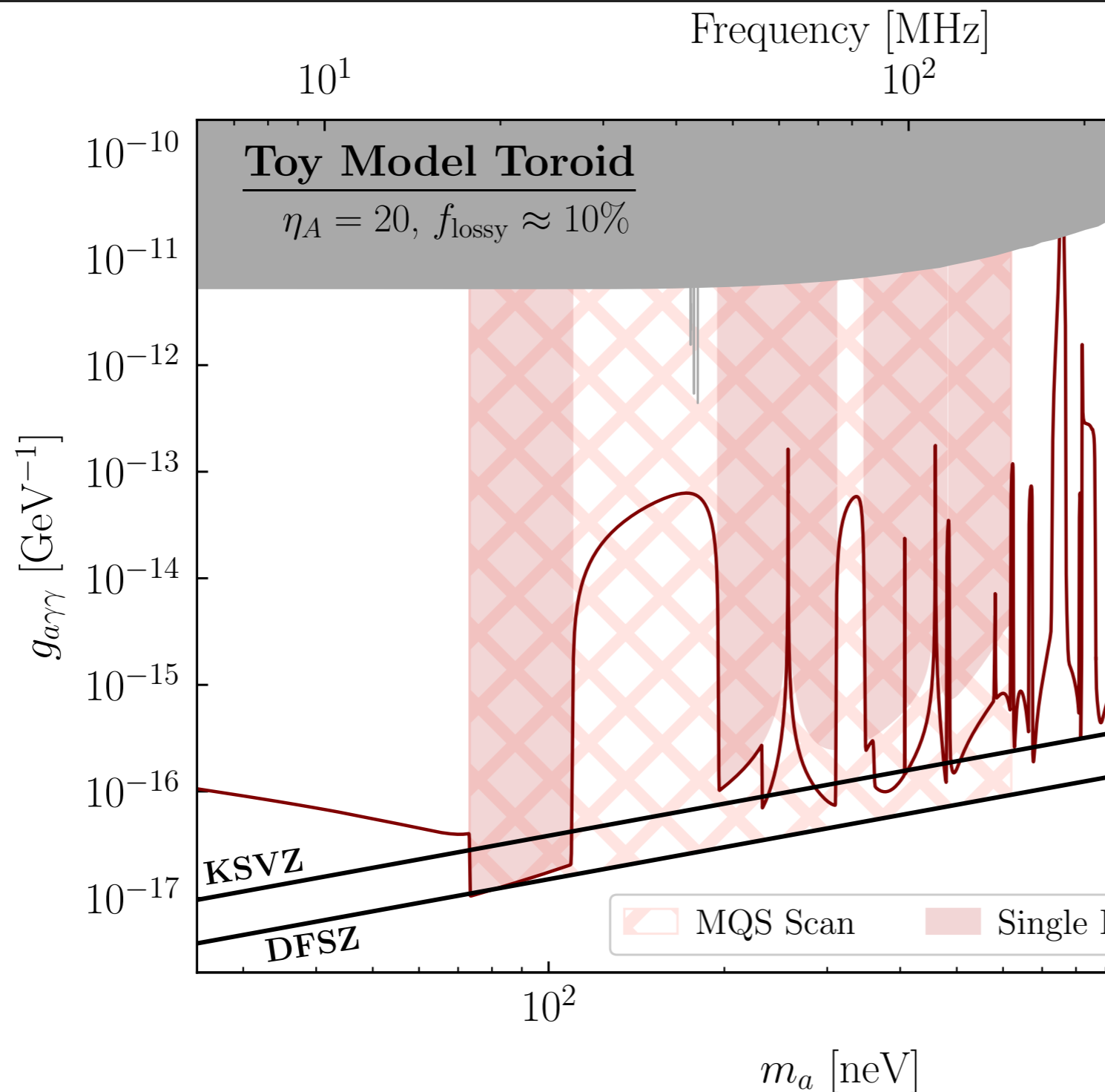


happening now

future
aspirations

designing now

Caveat: Our simulations show beyond MQS Approximation, challenges emerge



DM Radio is actively engineering solutions to mitigate these previously unaccounted for effects

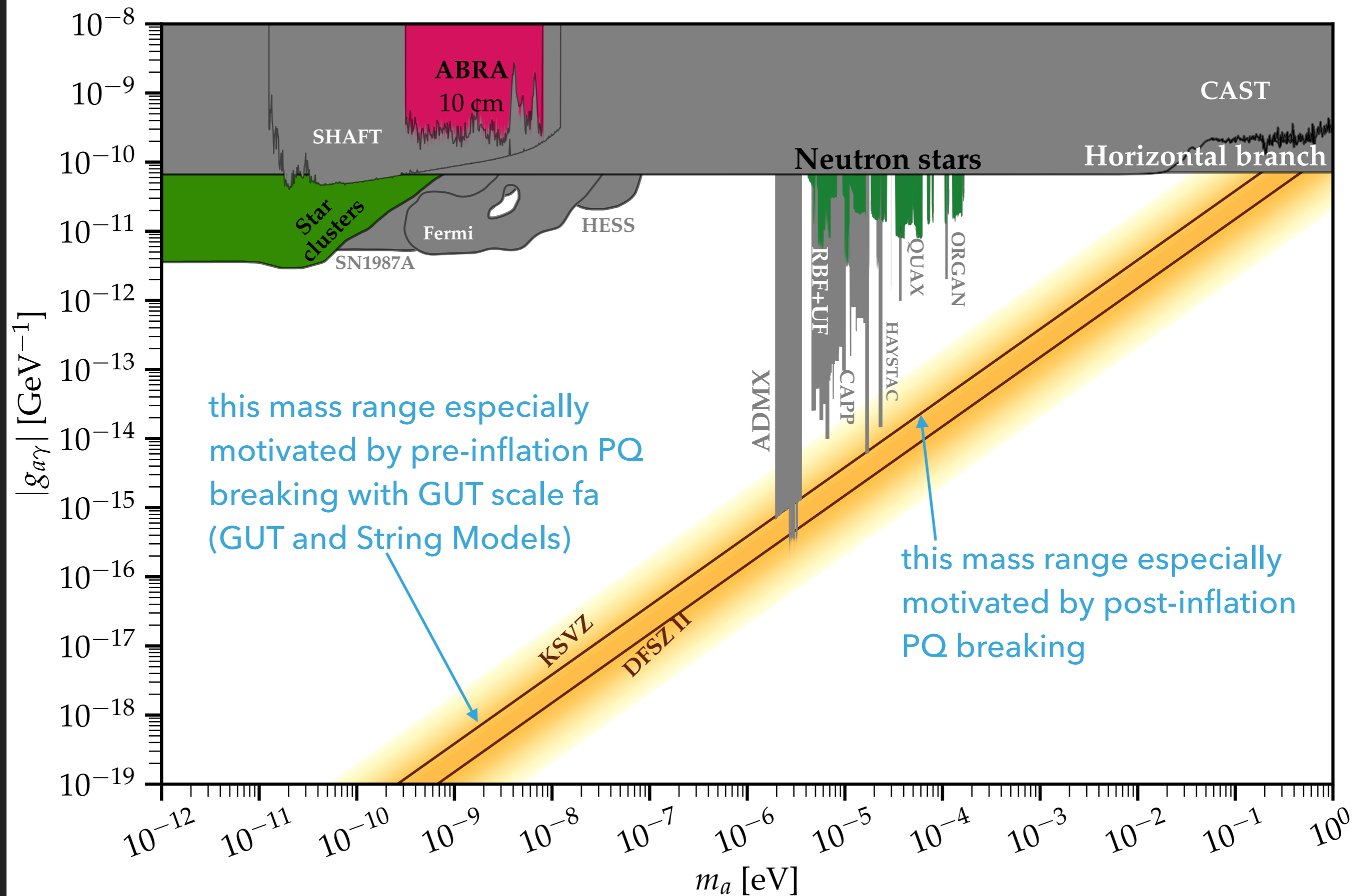
1. Radiation Losses
2. Toroid not ideal inductor

1. At some frequencies toroid is capacitive!

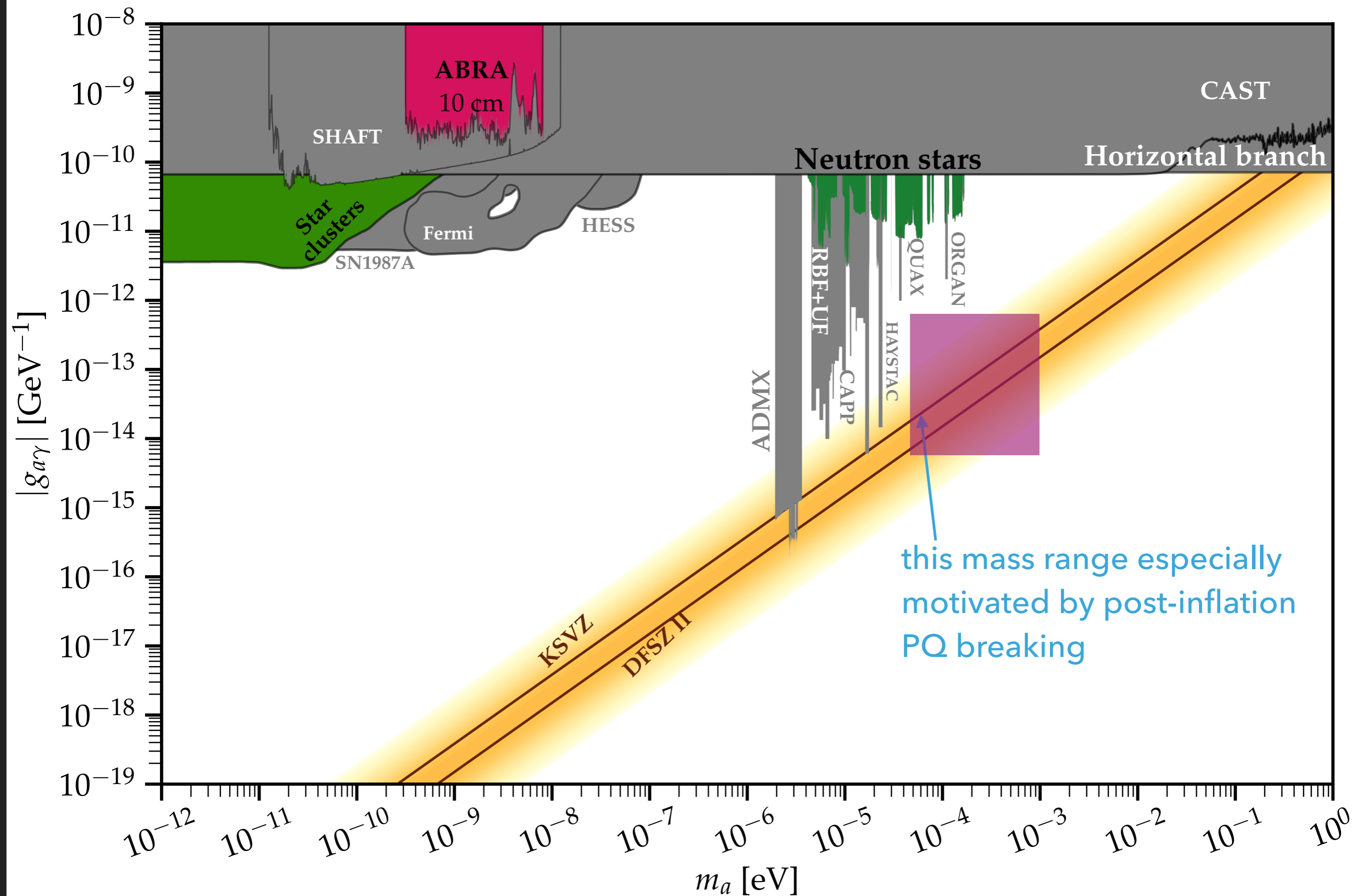
Detecting axion dark matter beyond the magnetoquasistatic approximation

Joshua N. Benabou,^{1,2} Joshua W. Foster,³ Yonatan Kahn,⁴ Benjamin R. Safdi,^{1,2} and Chiara P. Salemi^{5,6,7}

Motivated mass ranges



High-mass example: MADMAX



MADMAX is happening! DESY, Hamburg

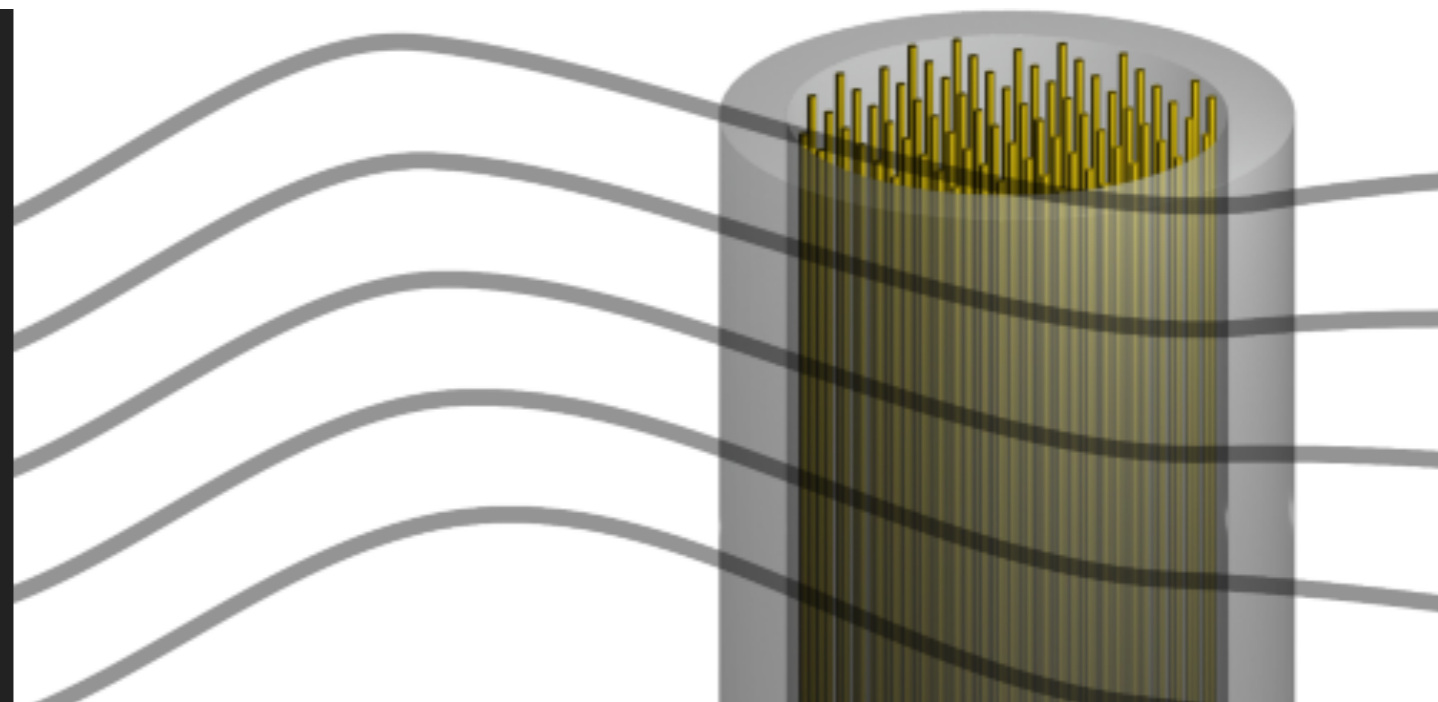


Other experiments in this range also:

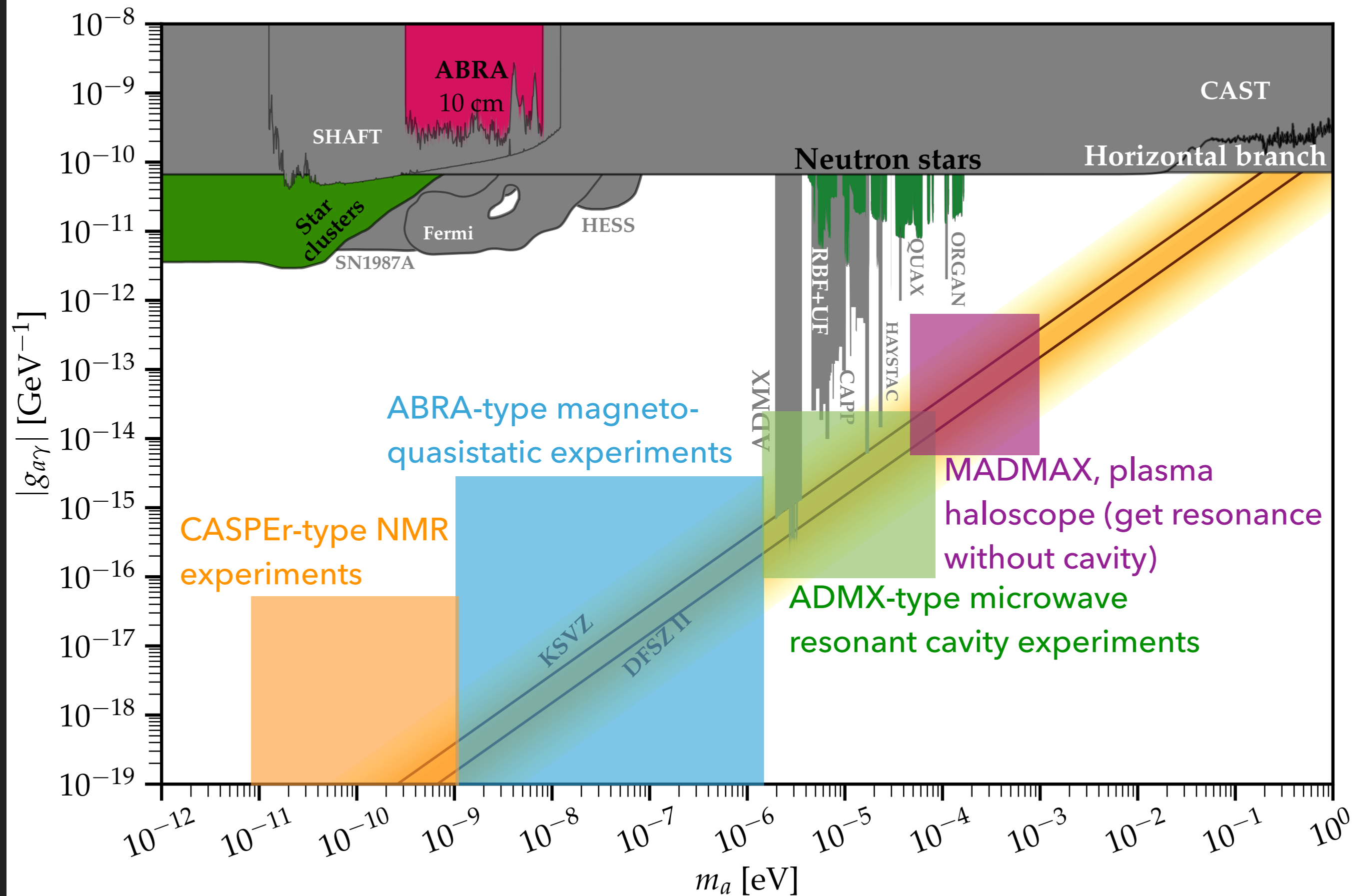
Array of wires instead of dielectric layers (morally accomplishes same result)

Tunable axion plasma haloscopes

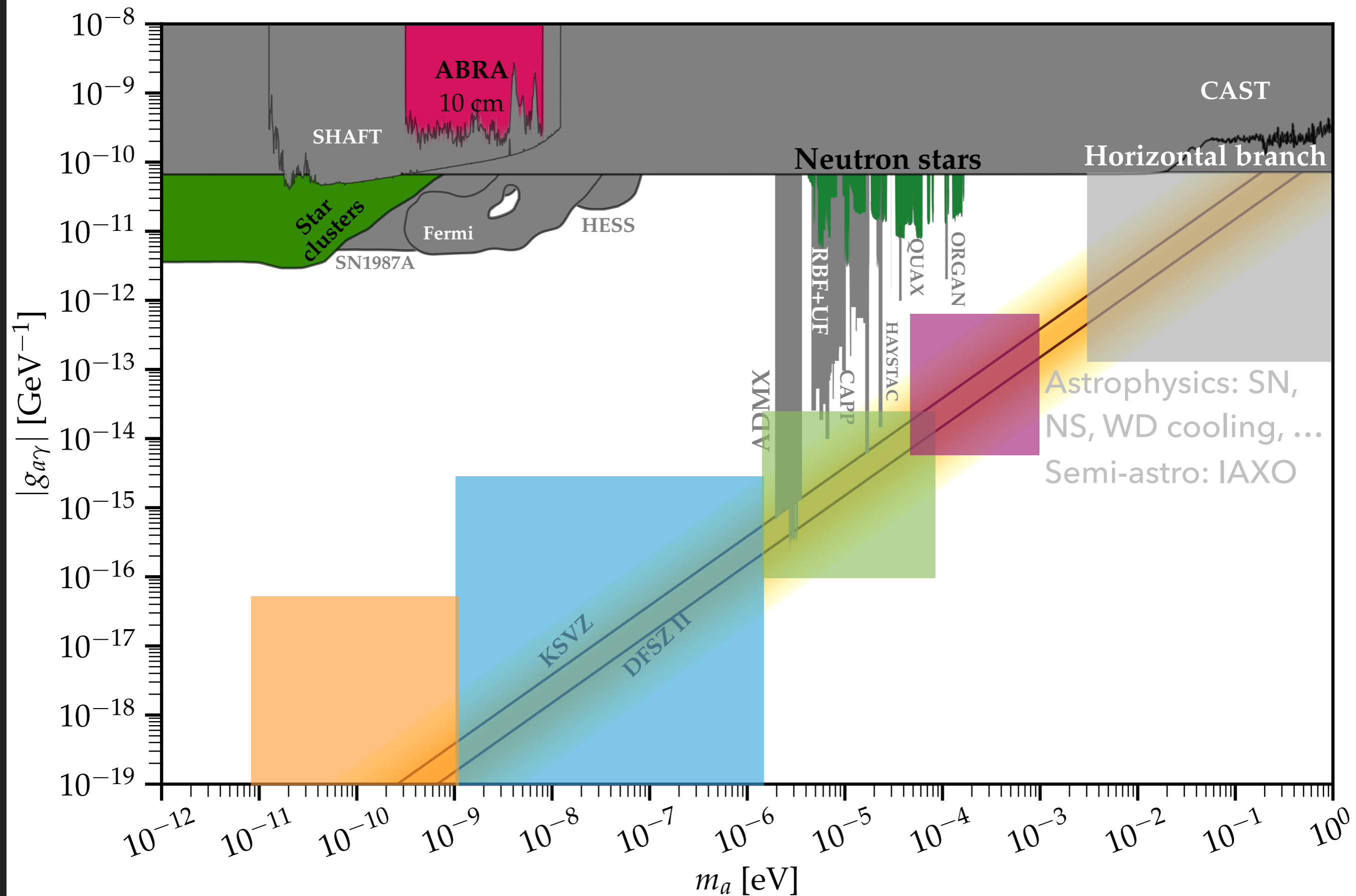
Matthew Lawson, Alexander J. Millar, Matteo Pancaldi, Edoardo Vitagliano, Frank Wilczek



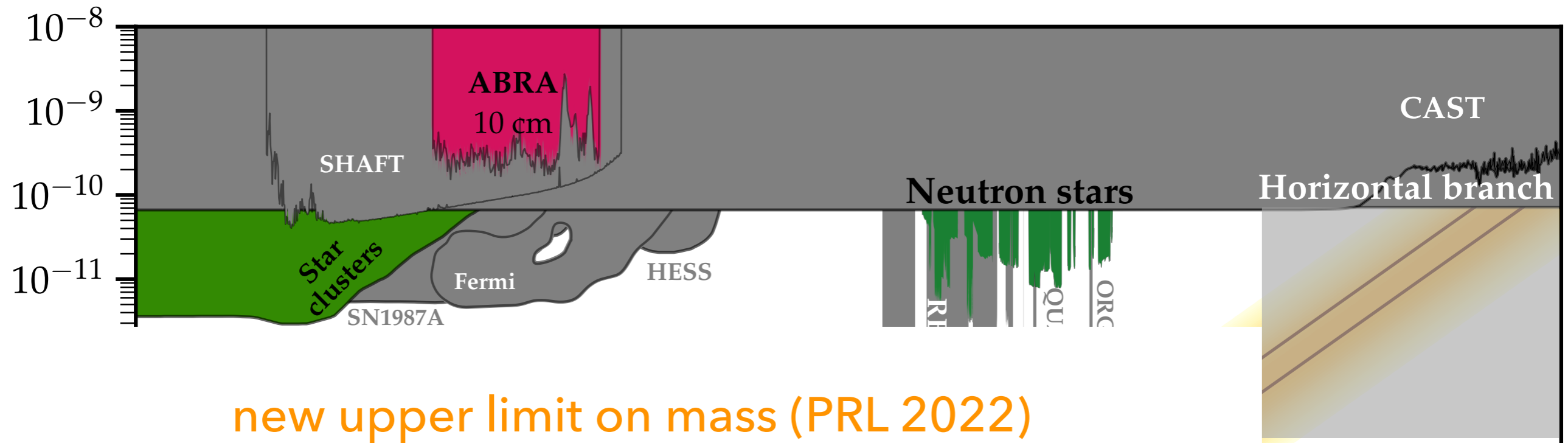
QCD Axion Direct Detection Summary



Astrophysical probes most relevant at high masses

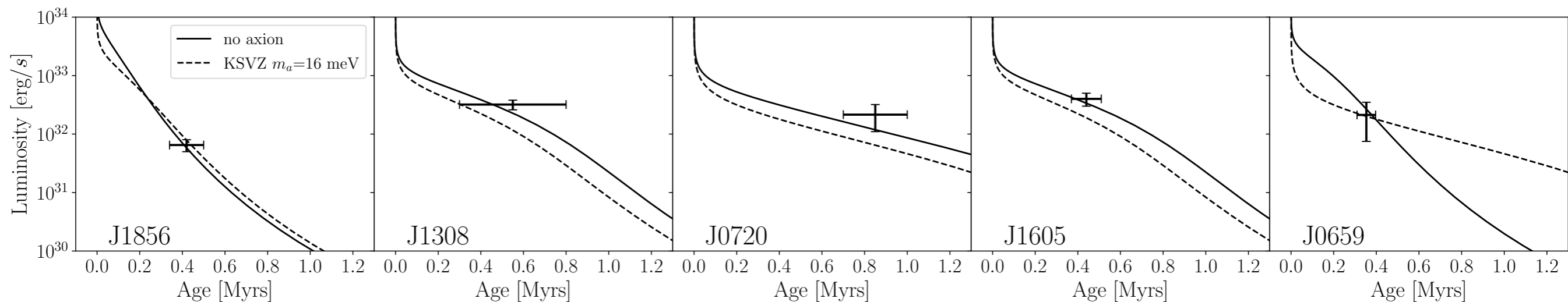


Astrophysical probes most relevant at high masses



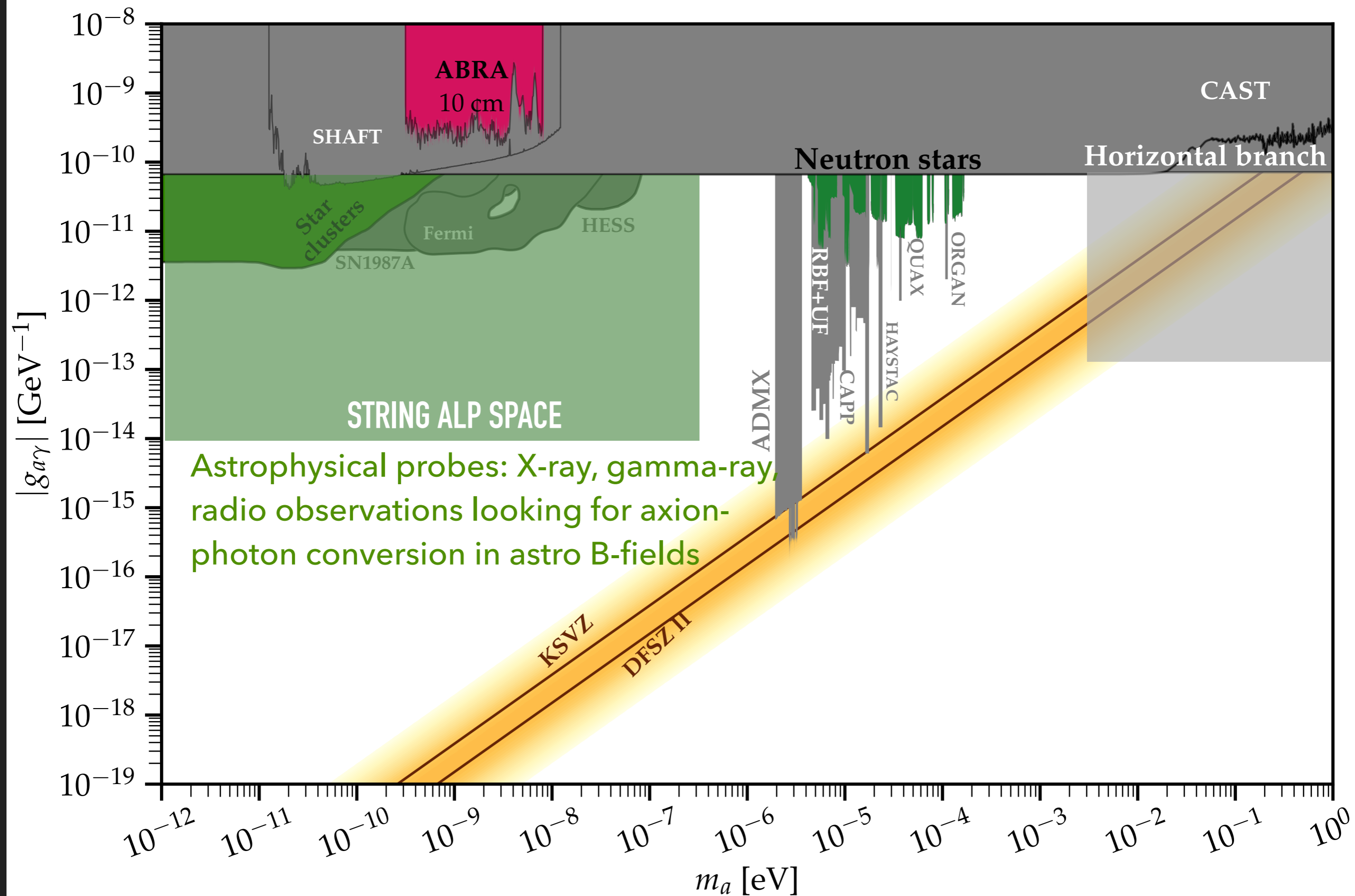
Upper Limit on the QCD Axion Mass from Isolated Neutron Star Cooling

Malte Buschmann,¹ Christopher Dessert,^{2,3,4} Joshua W. Foster,⁵ Andrew J. Long,⁶ and Benjamin R. Safdi^{3,4}



$m_a < (10 - 30) \text{ meV}$ (depending of KSVZ/DFSZ)

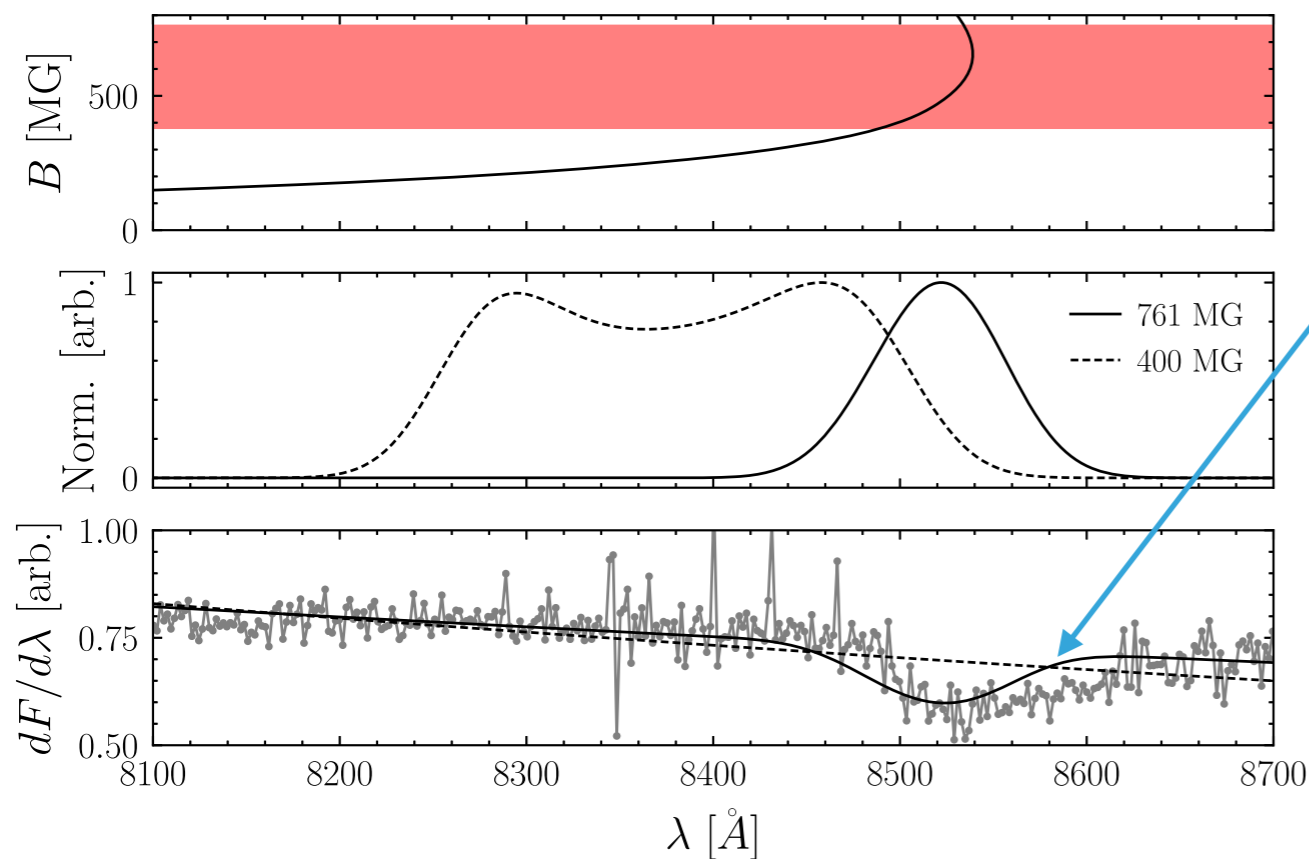
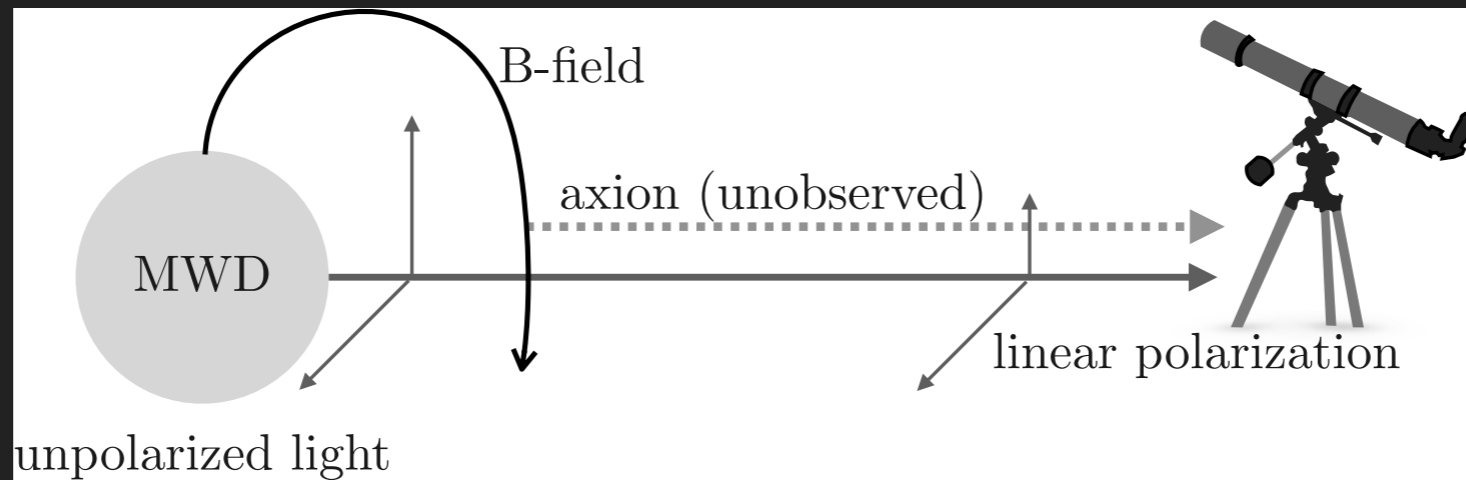
Low-mass ALPs motivated by string theory



Example: Magnetic White Dwarf Polarization

C. Dessert, D. Dunsky, **B.S.**, 2022

Robust probes of ultra-light ALPs



Magnetic fields measured by Zeeman effect for absorption lines in MWD atmosphere

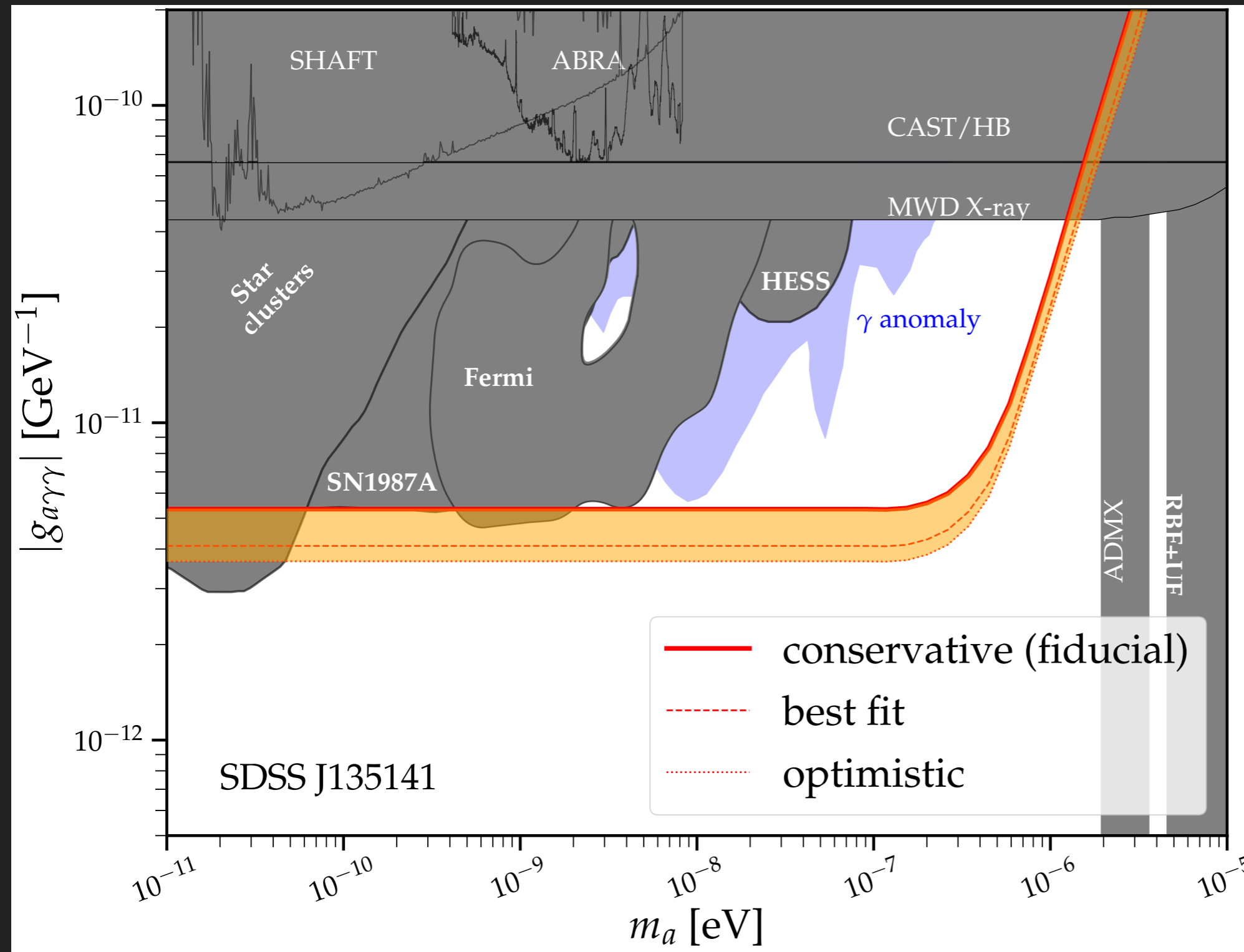
Astro contribution to linear polarization is negligible

1. current data is instrument-limited upper limit on polarization

Example: Magnetic White Dwarf Polarization

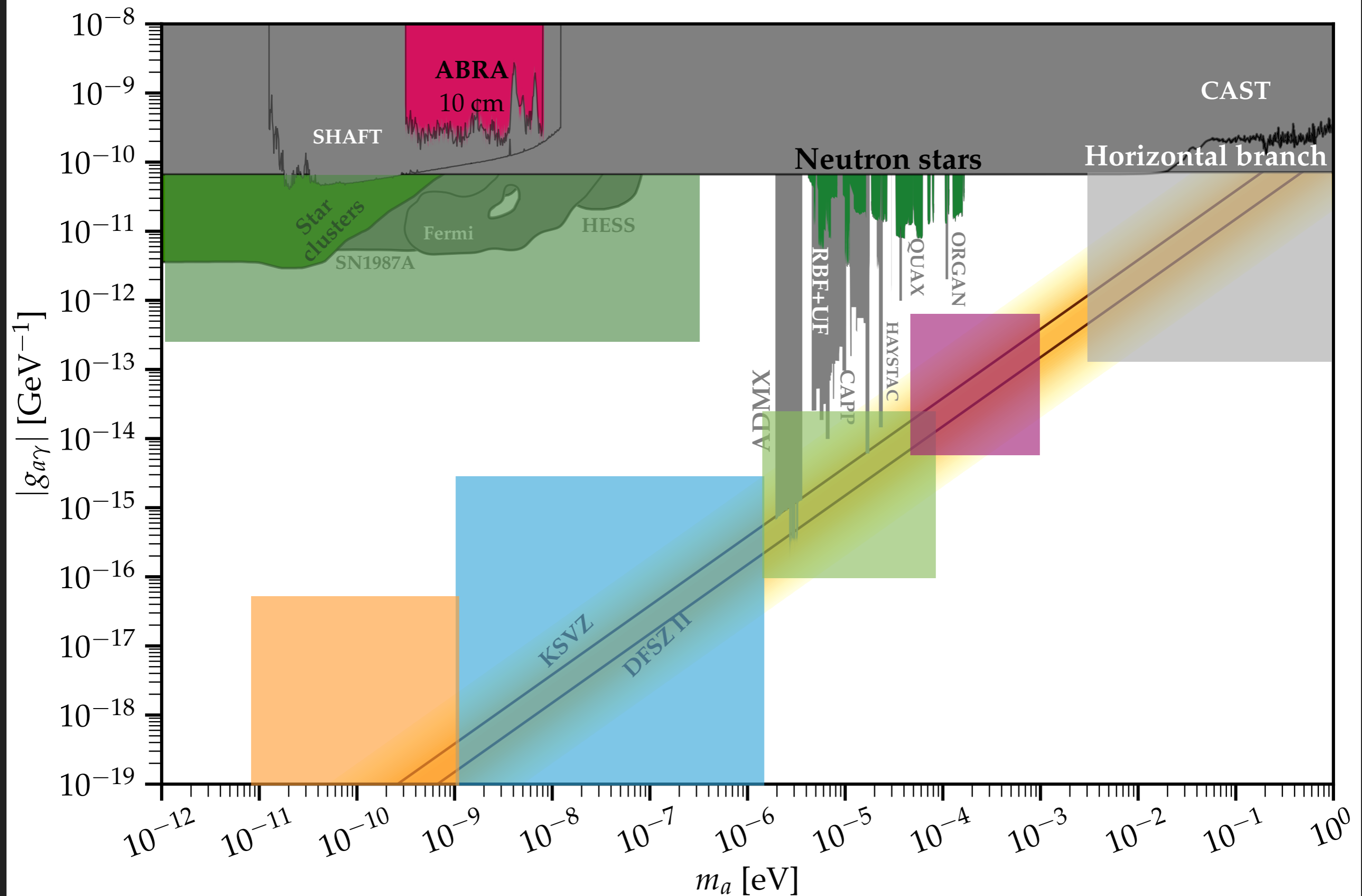
C. Dessert, D. Dunsky, **B.S.**, 2022

Robust probes of ultra-light ALPs



1. Order of magnitude improvement possible with dedicated data
2. How do we move to parametrically smaller couplings?

Summary



IF IT'S OUT THERE, WE WILL FIND IT!

QUESTIONS?