

# Rotating Neutron Stars

Universal Relations for the Increase in the Mass and Radius of a Rotating Neutron Star

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ALBERTA**

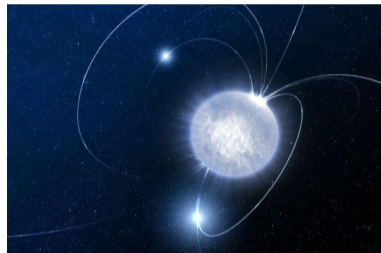
- **Part I** : Basics of Neutron Stars  
(Structure and EOSs)
- 

- **Part II** : Rotating neutron stars and the  
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- **Part III** : Results & Applications
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- **Part IV** : Why Universality?
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- **Part V** : Conclusions



**Figure:** Neutron star. Credits:  
ESO/LUÍS CALÇADA

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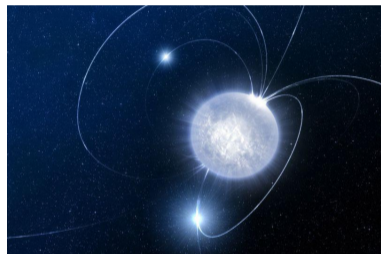


Figure: Neutron star. Credits:  
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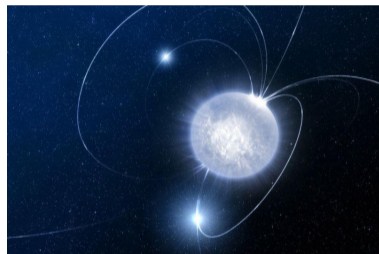


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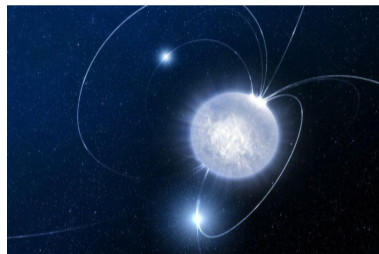


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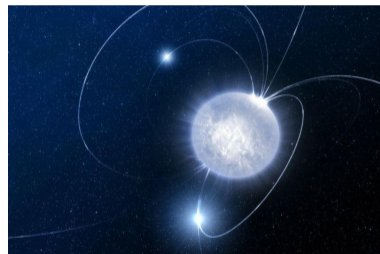


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# Part I : Basics of Neutron Stars

- **Atmosphere**

- slim plasma envelope

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- **Outer crust** ( $\sim 10^{11} \text{ g cm}^{-3}$ )

- Z ions and free e

- **Inner crust** ( $\sim 10^{14} \text{ g cm}^{-3}$ )

- free n and e

- neutron-rich atomic nuclei

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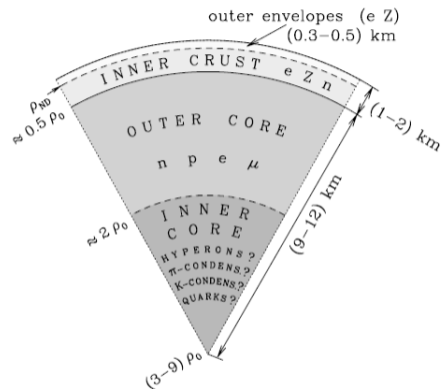
- **Outer core** ( $\sim 10^{14} \text{ g cm}^{-3}$ )

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- **Inner core** (larger than  $10^{15} \text{ g cm}^{-3}$ )

- hyperons, free quarks, exotic matter (???)



**Figure:** Internal structure of a neutron star  
( $\rho_0 = 2.28 \cdot 10^{14} \text{ g cm}^{-3}$ )  
[Haensel et al., 2007]



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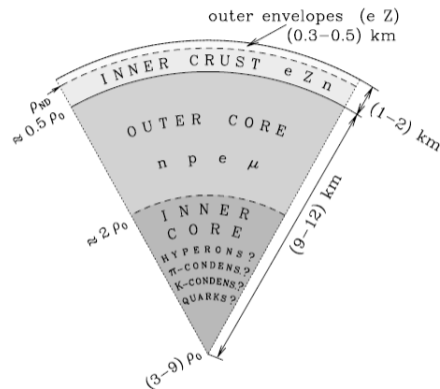
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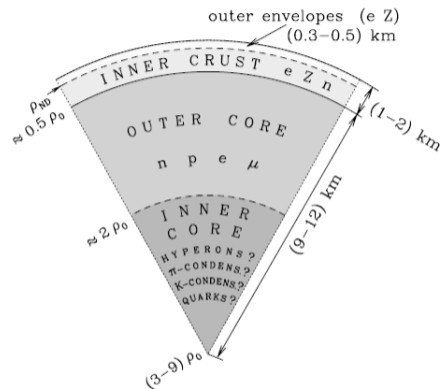
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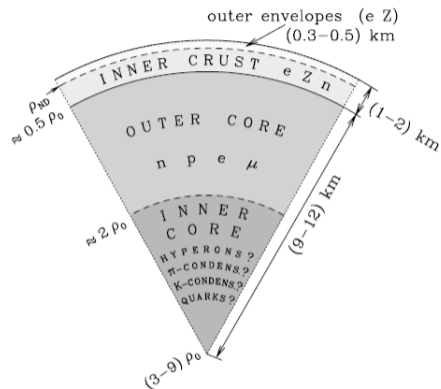
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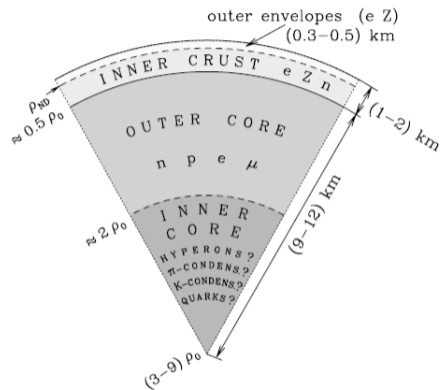
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**Figure:** Internal structure of a neutron star  
( $\rho_0 = 2.28 \cdot 10^{14} \text{ g cm}^{-3}$ )  
[Haensel et al., 2007]

- For a spherical symmetric (non-rotating) NS we use the TOV equations

$$\frac{dP(r)}{dr} = - \frac{G(\epsilon(r)c^2 + P(r))(m(r) + 4r^3 P(r)/c^2)}{rc^2[r - 2Gm(r)/c^2]} \quad (1)$$

$$\frac{dm(r)}{dr} = 4\pi r^2 \epsilon(r) \quad (2)$$

$$EOS : P(\epsilon) \quad (3)$$

- where P: pressure,  $\epsilon$ : energy density, m: mass coordinate and r: the radial distance

- The blue area in the plots is our expectations based on **observations** and **experiments**
- Phenomenological EOS:
  - Piecewise polytropes
  - Speed of sound model
- EOS models:
  - Nuclear EOS : APR, BBB, HLPS
  - Hybrid EOS: QHC21, ABPR
  - Hyperons: H0
  - $\pi^0$  condensate: L

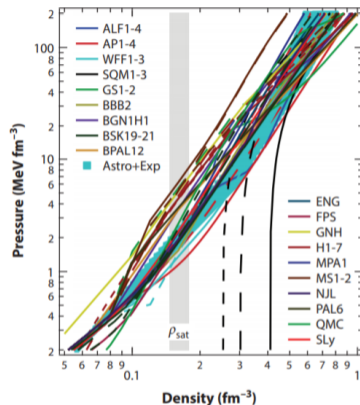


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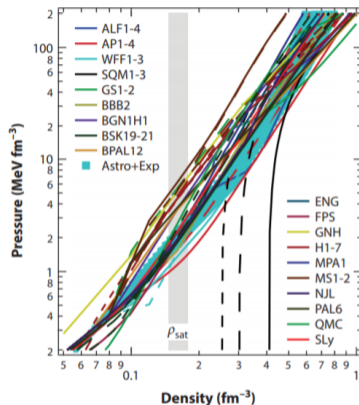


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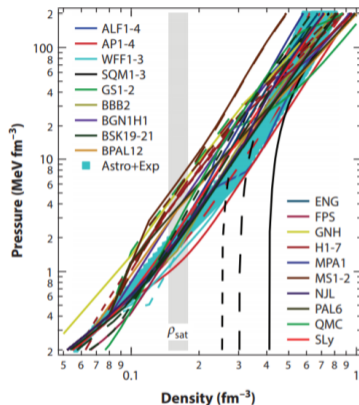
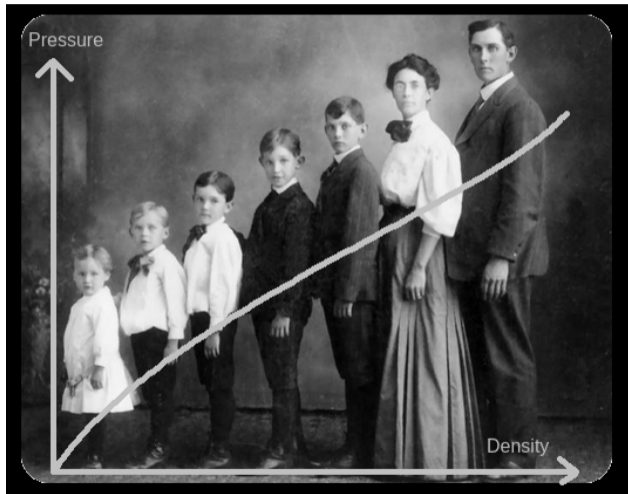


Figure: EOS models  
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# The Equation of State Families



- [Read et al., 2009],

$$P(\rho) = \begin{cases} P_{crust} & \rho \leq \rho_{crust} \\ K_1 \rho^{\Gamma_1} & \rho_{crust} \leq \rho \leq \rho_1 \\ K_2 \rho^{\Gamma_2} & \rho_1 \leq \rho \leq \rho_2 \\ K_3 \rho^{\Gamma_3} & \rho_2 \leq \rho \end{cases} \quad (4)$$

- [Hebeler et al., 2013] a) Causality  
b) Supports the heaviest observed NS

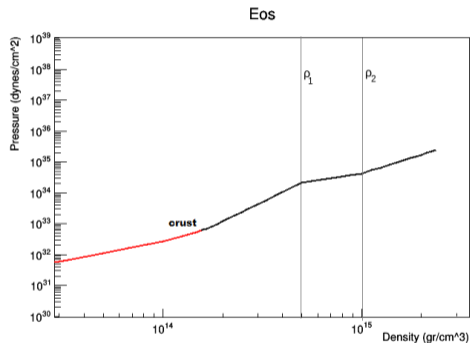


Figure: Piecewise Polytrope

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- [Hebeler et al., 2013]

a) **Causality**

b) Supports the **heaviest** observed NS

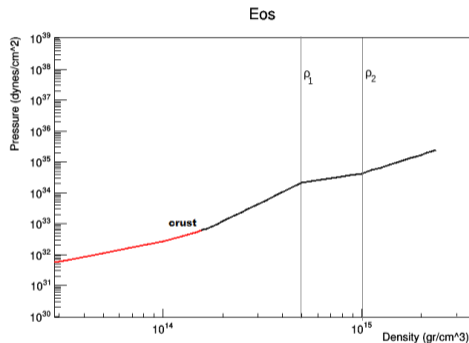
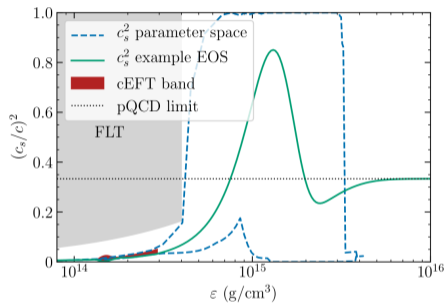


Figure: Piecewise Polytrope

- [Greif et al., 2019],

$$P(\epsilon) = \begin{cases} P_{crust} & n \leq 0.5n_0 \\ P_{cEFT} & 0.5n_0 \leq n \leq 1.1n_0 \\ \int_0^\epsilon (c_s(\epsilon')/c)^2 d\epsilon' & 1.1n_0 \leq n \end{cases} \quad (5)$$

- Causality  $0 \leq c_s \leq c$
- Supports the **heaviest** observed NS
- perturbative quantum chromodynamics (**pQCD**):  $\frac{c_s}{c}$  approaches  $\frac{1}{\sqrt{3}}$  from below when  $\sim 10^{16}$  g/cm<sup>3</sup>)



**Figure:** Speed of sound profile from [Greif et al., 2019]

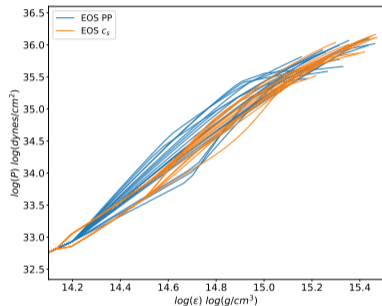


Figure: EOS

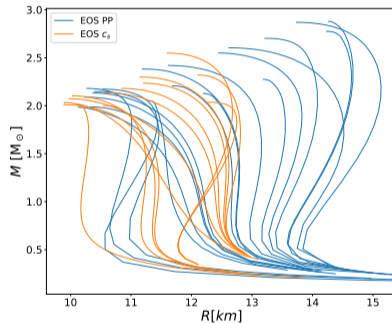


Figure: Mass - Radius curves

# Part II : Rotating Neutron Stars And Universal Relations



- **Cold (Temperatures less than the neutrons Fermi energy)**

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- Rigidly rotating

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- Low magnetic field

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- Described by a stationary metric

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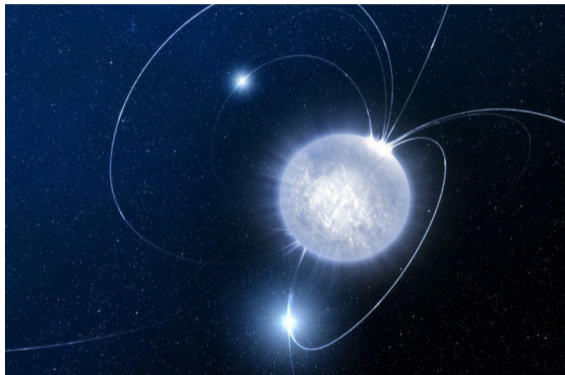


Figure: Neutron star. Credits: ESO/LUÍS CALÇADA

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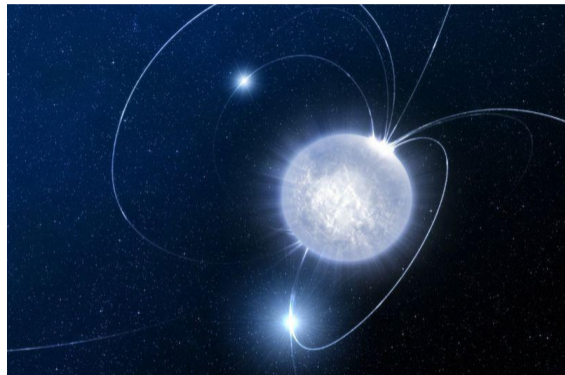


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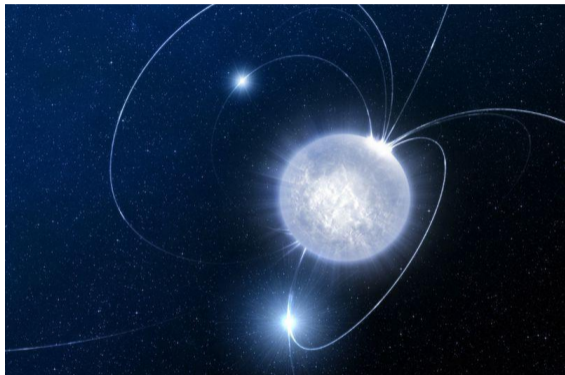


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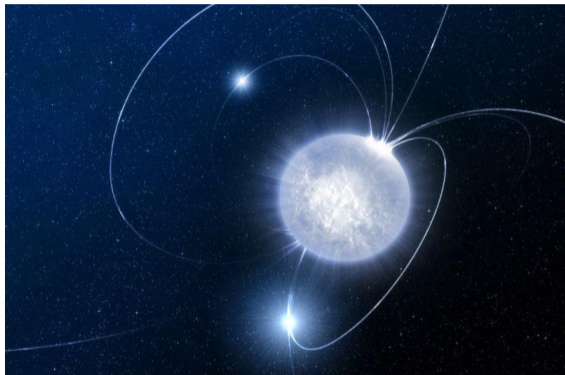


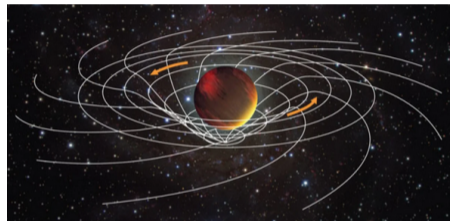
Figure: Neutron star. Credits: ESO/LUÍS CALÇADA

- **Axis symmetric metric**

$$ds^2 = -e^{\gamma+\lambda} dt^2 + e^{2\alpha} (dr^2 + r^2 d\theta^2) + e^{\gamma-\lambda} r^2 \sin^2(\theta) (d\phi - \omega dt)^2$$

$\lambda, \gamma, \alpha$  and  $\omega$  are called metric potentials and depend on  $r$  and  $\theta$ . [Cook et al., 1994]

- Use **RNS** code.  
[Stergioulas and Friedman, 1995]  
Numerically solves the GR equations for one rotating NS



**Figure:** Curved space-time for rotating systems

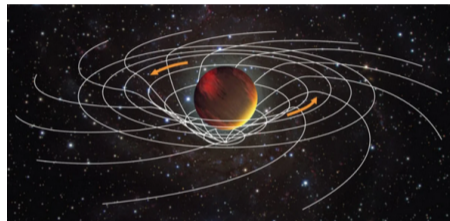
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- Create sequences of rotating NSs with constant  $\epsilon_c$  in order to see the relation among the rotation and the EOS
- $M$ : Total mass (in  $M_\odot$ )
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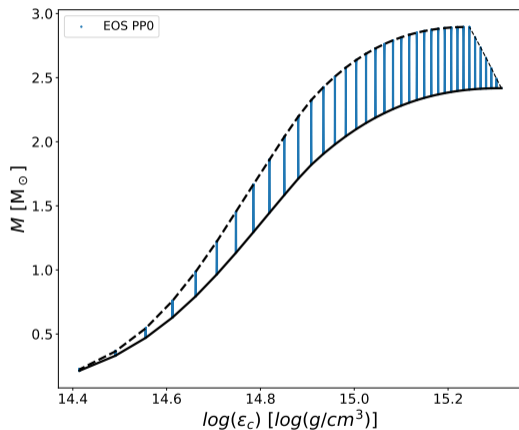


Figure: Total Mass vs central energy density

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- Are there any universal relations?
- If yes,

Are they useful for modelling the rotating neutron stars?

Are they useful for our observations?

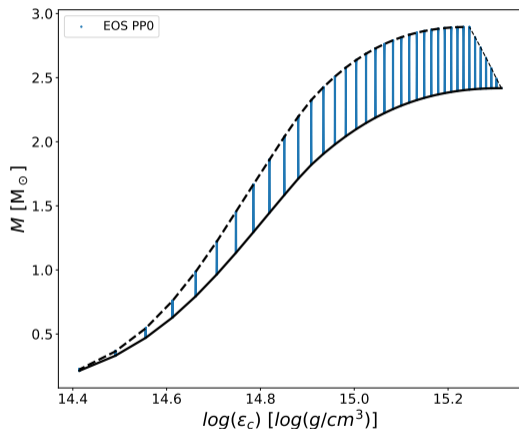


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- **Non-Universal Relations:**  
Equations that are strongly affected by the EOS structure.
- i.e. Mass-Radius relations.

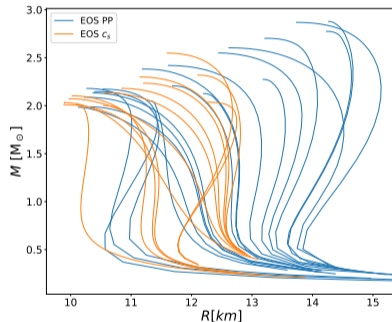


Figure: Mass - Radius curves

- **Universal Relations:**  
Equations that are **NOT** strongly affected by the EOS structure.

- **Some Examples:**  
BE vs  $M/R$  [Lattimer & Prakash (2001)],

Oblateness [Morsink et al. (2007), Silva et al. (2021)];

$I/M^3$  vs  $M/R$   
[Breu and Rezzolla, 2016],

I-Love-Q [Yagi and Yunes, 2013].

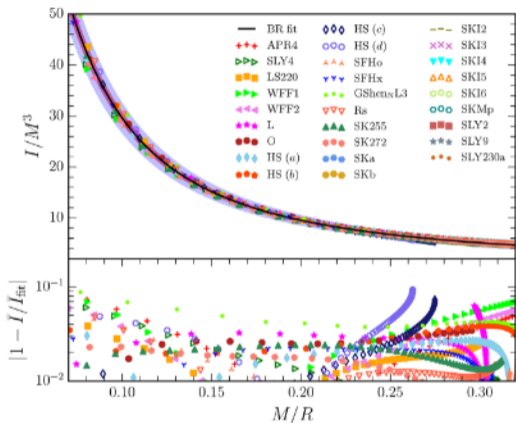


Figure:  $I/M^3$  vs  $M/R$  [Breu and Rezzolla, 2016]

# Part III : Results & Applications



- $\Omega_n = \frac{\Omega}{\Omega_K(\text{empirical})}$

- Empirical approximation:

Previous work: Haensel & Zdunik (1989), Friedman et al. (1989), Haensel et al. (2009), Haskell et al. (2018), and Koliogiannis & Moustakidis (2020)

Our work:

$$\Omega_K(\text{empirical}) = N_\Omega(C_*) \sqrt{\frac{GM_*}{R_*^3}}$$

- $N_\Omega(C_*) = (a_1 C_*^4 + a_2 C_*^3 + a_3 C_*^2 + a_4 C_* + a_5)$

$$\text{Dev}(\Omega_K)_{\text{max}} = 1.60\%$$

- $\text{Dev}(Z) = |Z_{\text{data}} - Z_{\text{best fit}}| / Z_{\text{data}}$

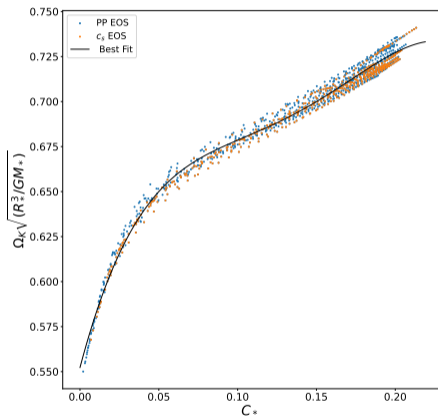


Figure: Kepler frequency

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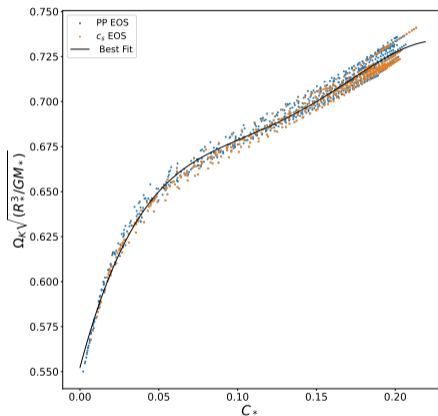


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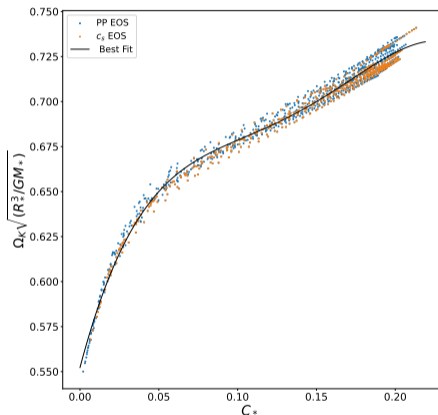


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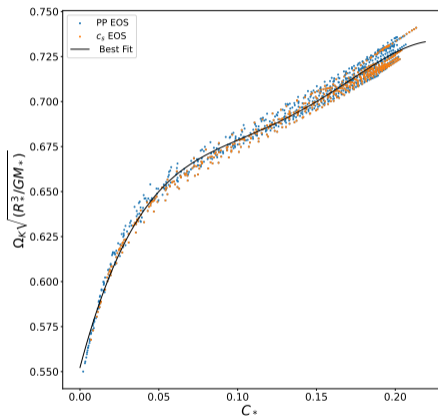


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- $C_e$ : Equatorial compactness  $\frac{M}{R_e} \frac{km}{M_\odot}$
- Compactness is strongly related to the central energy density
- $C_*$ : Initial compactness  $\frac{M_*}{R_*} \frac{km}{M_\odot}$

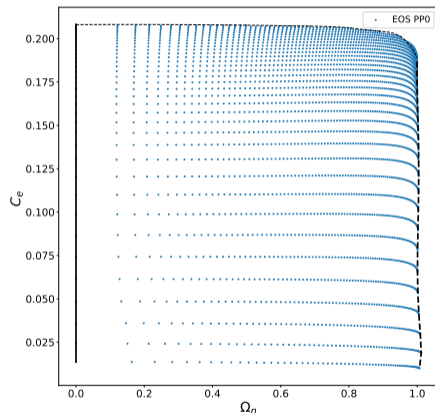


Figure:  $C_e$  vs  $\Omega_n$  for PP EOS 0



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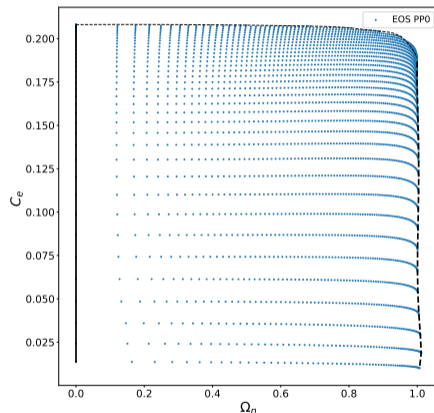


Figure:  $C_e$  vs  $\Omega_n$  for PP EOS 0

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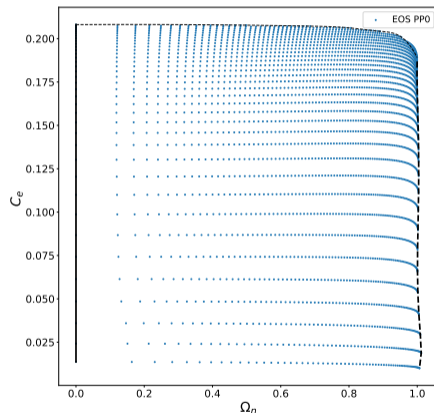


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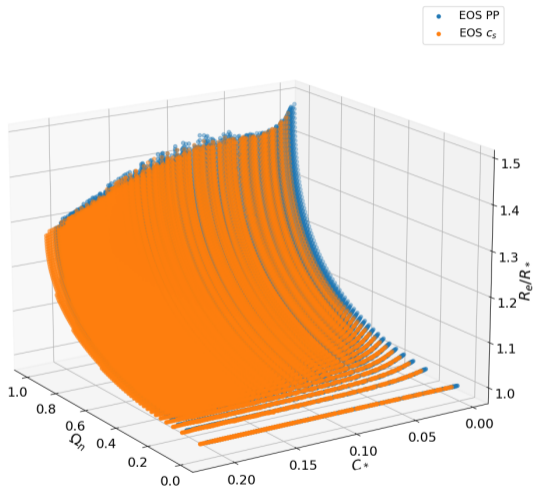
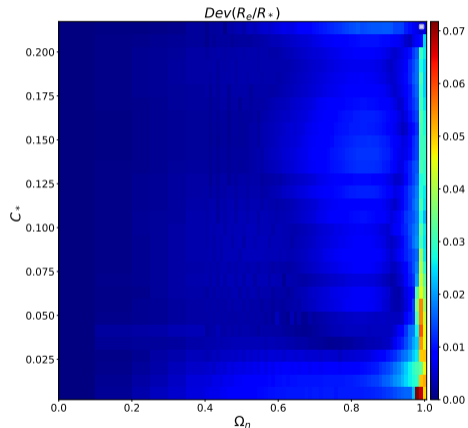


Figure:  $R_e/R_*$  vs  $C_*$  vs  $\Omega_n$



$$\frac{R-R_*}{R_*} = (e^{\alpha r} \Omega_n^2 - 1 - \beta_r \ln(1 - (\frac{\Omega_n}{1.1})^4))^2 \times (1 + \sum_{i=1}^5 r_i C_*^i)$$

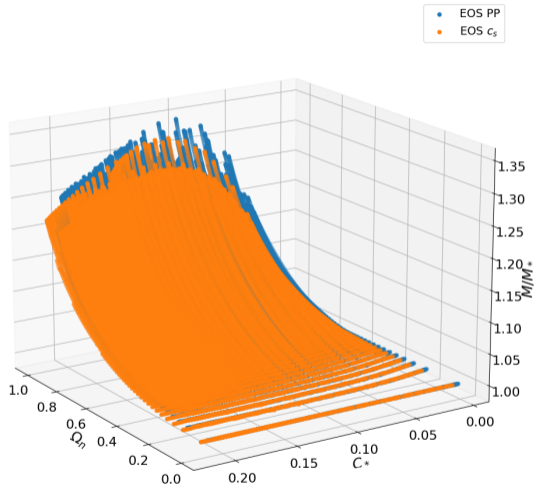
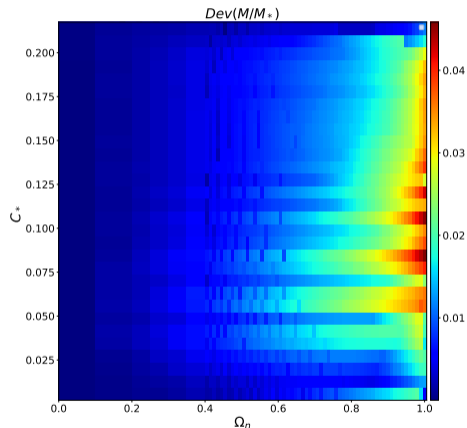


Figure:  $M/M_*$  vs  $C_*$  vs  $\Omega_n$



$$\frac{M - M_*}{M_*} = (e^{\alpha_m \Omega_n^2} - 1) \times \left( \sum_{i=0}^4 m_i C_*^i \right)$$

- Solve the TOV Equations for a specific EOS (Black solid line)  
This means that we know  $M_*$ ,  $R_*$  and  $C_*$
- Use the **best fit surfaces** for  $M/M_*$  and  $R_e/R_*$  for  $\Omega_n = 0.95$  (Orange solid line)

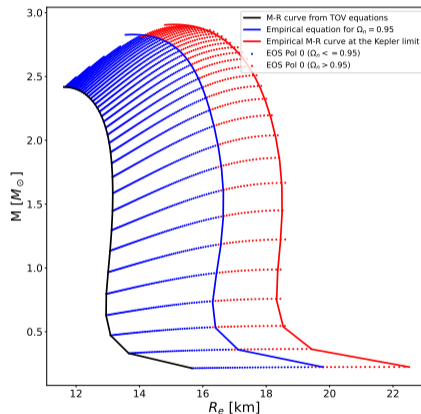


Figure: Mass - Radius curves for rotating NSs

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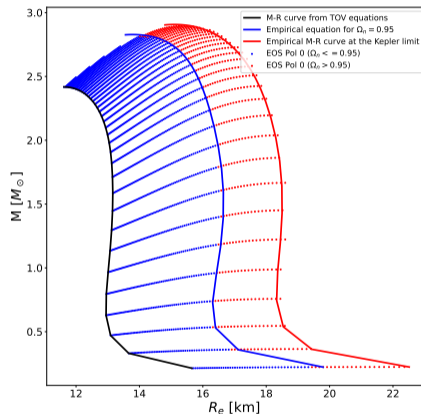


Figure: Mass - Radius curves for rotating NSs

- $\Omega_n(\Omega, R_*, M_*) \rightarrow \Omega_n(\Omega, R_e, M)$  (Max Dev  $\approx 1.5\%$ )

$$\frac{M}{M_*}(\Omega, R_*, M_*) \rightarrow \frac{M}{M_*}(\Omega, R_e, M) \text{ (Max Dev } \approx 5\%)$$

$$\frac{R_e}{R_*}(\Omega, R_*, M_*) \rightarrow \frac{R_e}{R_*}(\Omega, R_e, M) \text{ (Max Dev } \approx 1.8\%)$$

- We can find the non rotating NS with the same central energy density (Orange points)

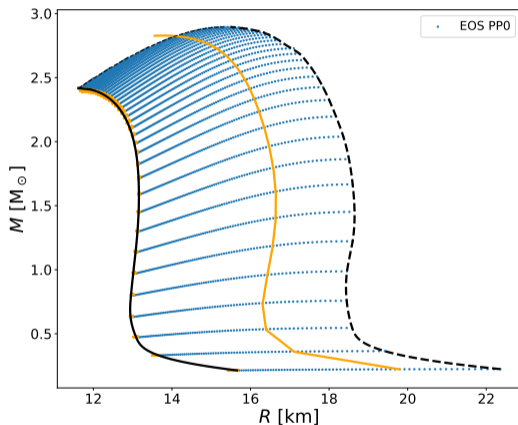


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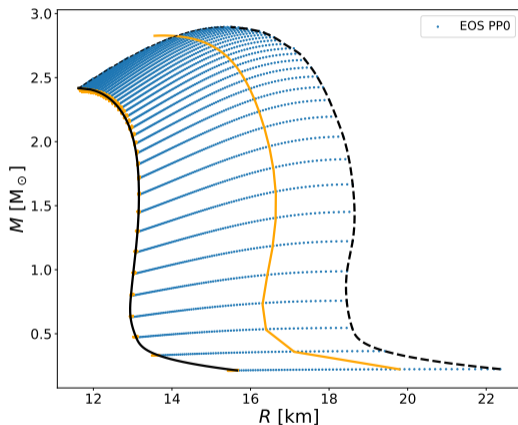
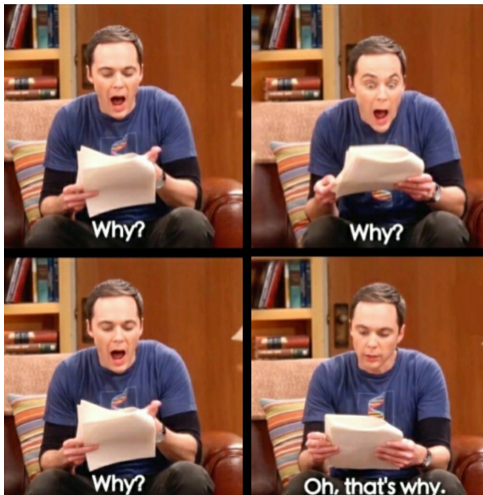
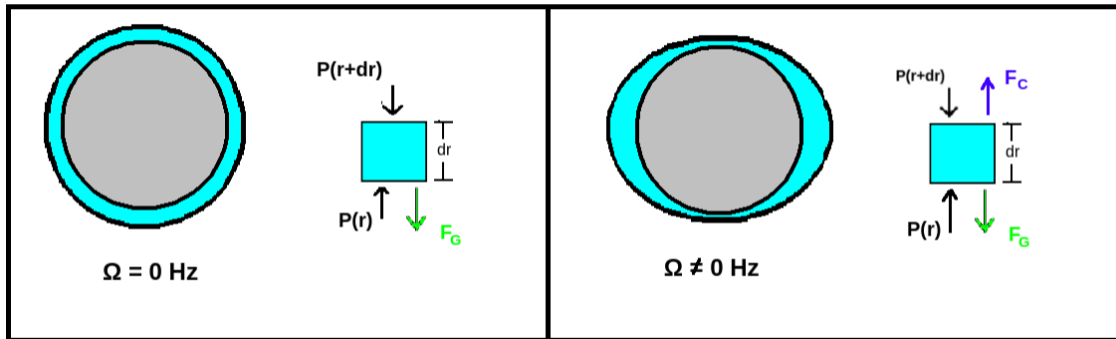


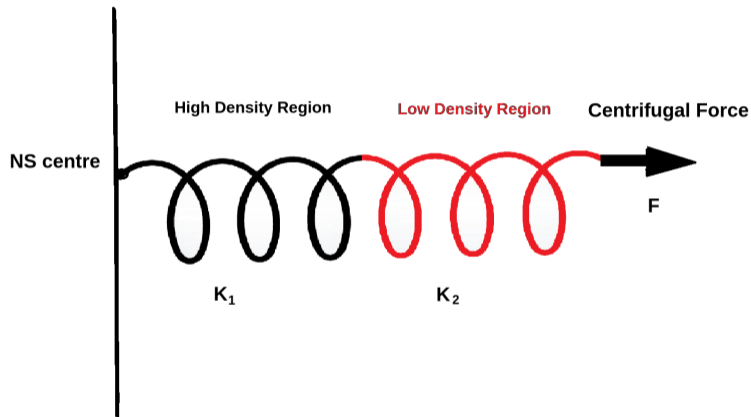
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# Part IV : Why Universality?







- $K_{eff} = \frac{K_1 K_2}{K_1 + K_2}$

- $K_{eff} \rightarrow K_2$ , when  $K_2 \ll K_1$

- A star with a uniform constant density  $\rho_0$

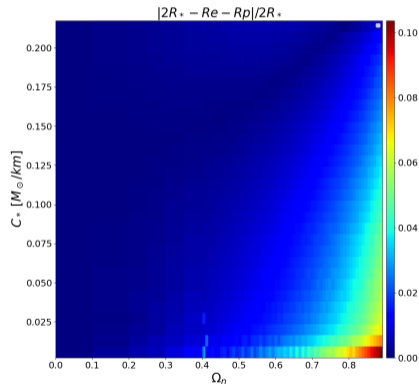
- $M = \frac{4\pi\rho_0}{3}R_e^2R_p$

- $R_e = R_* + \delta R_e(\Omega)$   
 $R_p = R_* - \delta R_p(\Omega)$

- $V_i \sim V_f$ , and  $\frac{\delta R_e(\Omega)}{R_*} \ll 1$   
 $\delta R_e(\Omega) \sim 2 \times \delta R_p(\Omega)$

- For  $\frac{\delta R_e(\Omega)}{R_*} \ll 1$ ,  
 $R_p \times R_e = R_*(1 + \alpha_e(\Omega)) \times$   
 $R_*(1 - \alpha_p(\Omega)) \approx R_*^2$

- $M \approx \frac{4\pi\rho_0}{3}R_*^2R_e \Rightarrow \frac{M}{M_*} \approx \frac{R_e}{R_*}$



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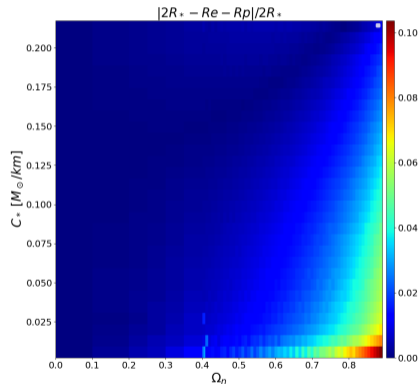
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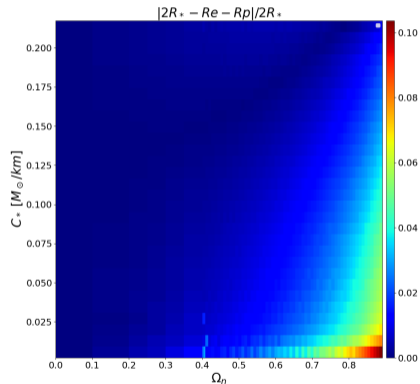
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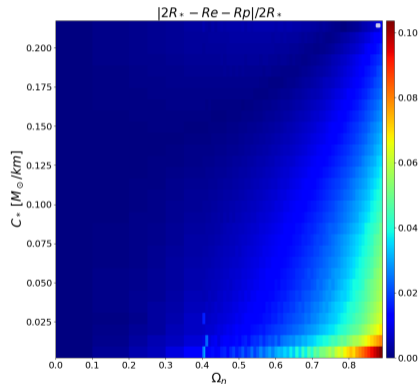
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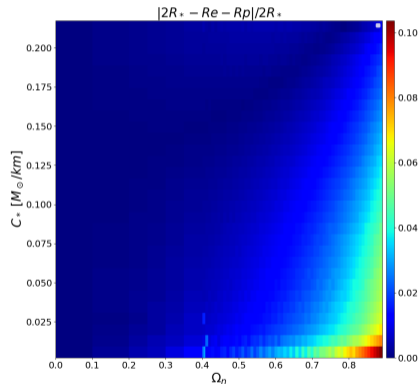
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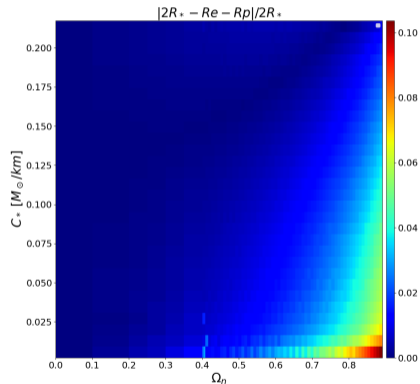


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- $\rho = \rho_{center} [1 - G(\sqrt{x^2/R_e^2 + y^2/R_e^2 + z^2/R_p^2})]$
- $M = \rho_{center} \int_{-R_e}^{R_e} \int_{-R_e \sqrt{1 - \frac{x^2}{R_e^2}}}^{R_e \sqrt{1 - \frac{x^2}{R_e^2}}} \int_{-R_p \sqrt{1 - \frac{x^2}{R_e^2} - \frac{y^2}{R_e^2}}}^{R_p \sqrt{1 - \frac{x^2}{R_e^2} - \frac{y^2}{R_e^2}}} [1 - G(\sqrt{x^2/R_e^2 + y^2/R_e^2 + z^2/R_p^2})] dx dy dz$
- We can set  $\frac{x}{R_e} = u$ ,  $\frac{y}{R_e} = v$  and  $\frac{z}{R_p} = w$
- $M = \rho_{center} R_e^2 R_p [\frac{4\pi}{3} - I]$ , where  $I = \int_{-1}^1 \int_{-\sqrt{1-u^2}}^{\sqrt{1-u^2}} \int_{-\sqrt{1-u^2-v^2}}^{\sqrt{1-u^2-v^2}} [G(\sqrt{u^2 + v^2 + w^2})] du dv dw$
- $M \approx \rho_{center} R_*^2 [\frac{4\pi}{3} - I] R_e$
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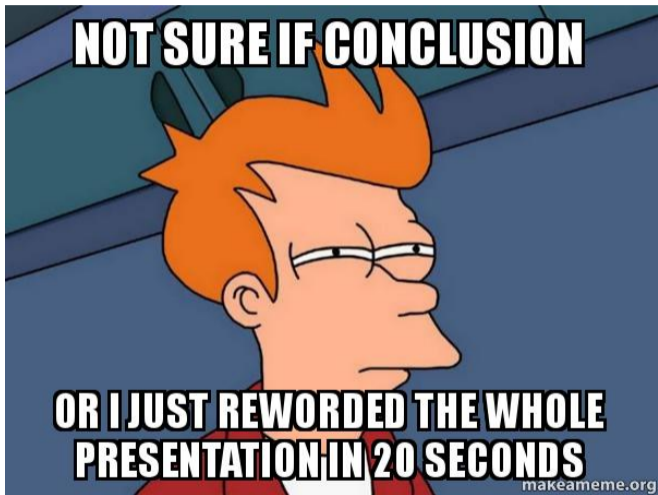
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## Part V : Conclusions





- We used rns code to create sequences of rapidly rotating NSs with constant  $\epsilon_c$ .
- We provide a new empirical equation for the Kepler frequency.
- $M/M_*$  and  $R_e/R_*$  are universal.
- We provide best fit equations for them, so we can recreate the Mass-Radius curves for the rotating NSs.
- We can use our equations, in order to find the non rotating neutron star with the same central energy density.

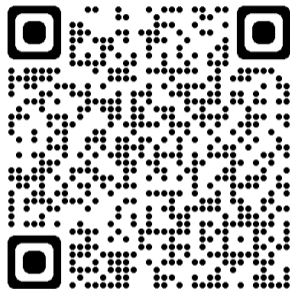


Figure: Universal Relations for the Increase in the Mass and Radius of a Rotating Neutron Star (Konstantinou Morsink 2022)

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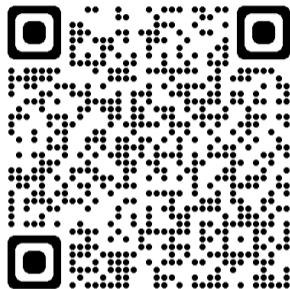


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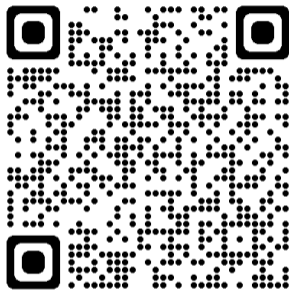


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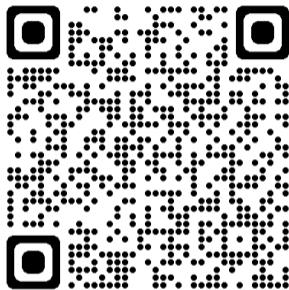


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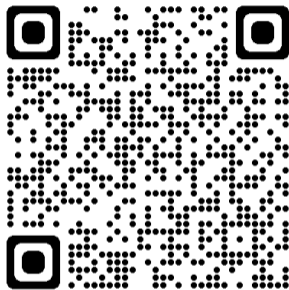









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Thank you for your attention!

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Maximum mass, moment of inertia and compactness of relativistic stars.  
*Physical Review D*, 459(1):646–656.
-  Cook, G. B., Shapiro, S. L., and Teukolsky, S. A. (1994).  
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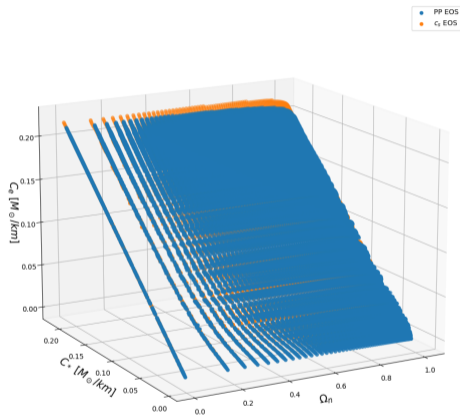
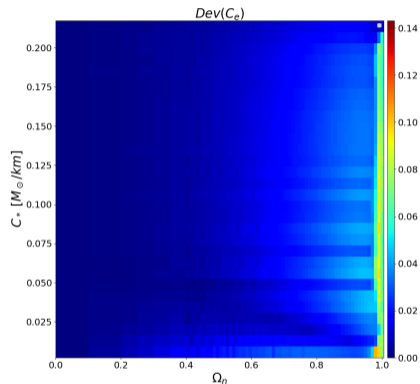


Figure:  $C_e$  vs  $C_*$  vs  $\Omega_n$



$$C_e \approx C_* + D_0 \times \ln\left(1 - \left(\frac{\Omega_n}{1.1}\right)^3\right) \times \left(1 + D_1 \times C_* + D_2 \times C_*^2 + D_3 \times C_*^4 + D_5 \times C_*^6\right)$$



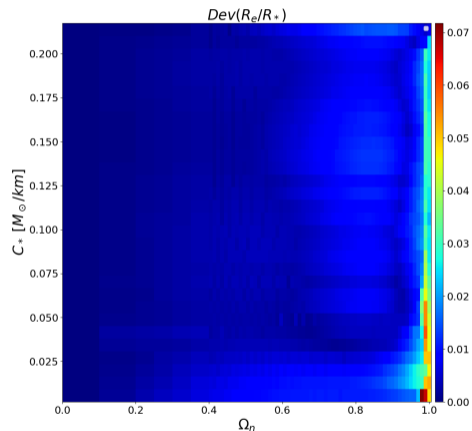
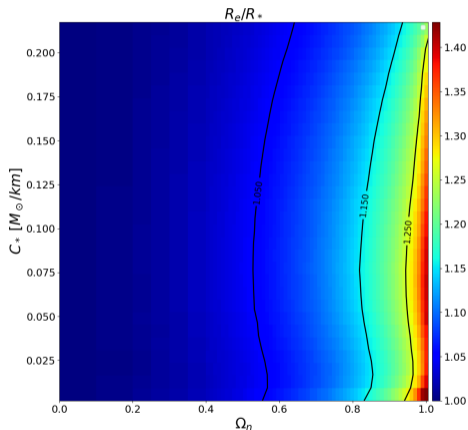


Figure:  $R_e/R_*$  vs  $C_*$  vs  $\Omega_n$

$$\frac{R-R_*}{R_*} = (e^{\alpha r} \Omega_n^2 - 1 - \beta_r \ln(1 - (\frac{\Omega_n}{1.1})^4))^2 \times (1 + \sum_{i=1}^5 r_i C_*^i)$$

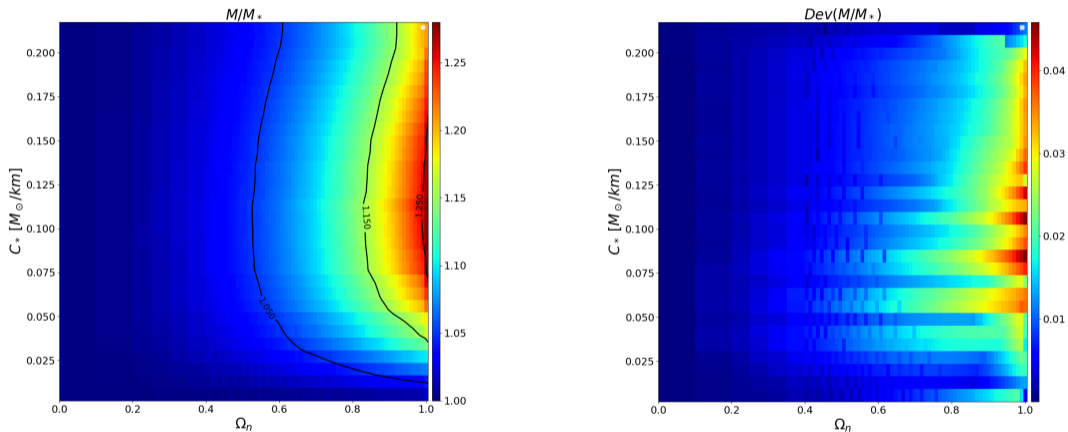


Figure:  $M/M_*$  vs  $C_*$  vs  $\Omega_n$

$$\frac{M-M_*}{M_*} = (e^{\alpha_m \Omega_n^2} - 1) \times \left( \sum_{i=0}^4 m_i C_*^i \right)$$