

Parton Production in nn and AA collision

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0. Introduction

- QM05:
 - We have a parton system which is not just a bunch of partons but has collective effects: flow, jet suppression and very dense.
 - Which means that the classical parton cascade may have only very limited application! → quantum Boltzmann or PC+Hydro
 - Before we move on, we need to understand very carefully the partonic system: ‘How many parton’, ‘How are they evolving’, ‘How are they breaking up’, ‘How are the particles formed’

1. Particle production

- Depends on the colliding particles:
 - ep
 - nn
 - AA
- Depends on the available energy:
 - KeV
 - MeV
 - GeV
 - TeV

CROSS SECTION:

1) Physical meaning:

$$a. \quad d\sigma = D(\theta, \phi) d\Omega \rightarrow D(\theta, \phi) = \frac{d\sigma}{d\Omega} = \left| \frac{b}{\sin \theta} \left(\frac{db}{d\theta} \right) \right|$$

$$b. \quad dN = L d\sigma \rightarrow \frac{d\sigma}{d\Omega} = \frac{1}{L} \frac{dN}{d\Omega}$$

2) Fermi Golden Rule:

$$1 + 2 \rightarrow 3 + 4 + \dots + n$$

$$d\sigma = |M|^2 \frac{\hbar^2 S}{4\sqrt{(p_1 + p_2)^2 - (m_1 m_2 c^2)^2}} \left[\left(\frac{cd^3 p_3}{(2\pi)^2 2E_3} \right) \right] \left[\left(\frac{cd^3 p_4}{(2\pi)^2 2E_4} \right) \right] \dots \left[\left(\frac{cd^3 p_n}{(2\pi)^2 2En_n} \right) \right] \\ (2\pi)^4 \delta^4(p_1 + p_2 - p_3 - \dots - p_n)$$

3. ep Collisions

Depending on energy:

- Movie at **KeV**:
- Movie at **MeV**:
- Movie at **GeV**:
- Movie at **TeV**:

- ep at KeV: see only electric field: Rutherford scattering, Mott formula

$$\left(\frac{d\sigma}{d\Omega}\right)_{\text{mott}} = \left(\frac{\alpha\hbar}{2\vec{p}^2 \sin^2(\theta/2)}\right)^2 \left[(mc)^2 + \vec{p}^2 \cos^2 \frac{\theta}{2} \right]$$

$$\left(\frac{d\sigma}{d\Omega}\right)_{\text{Rutherford}} = \left(\frac{e^2}{2mv^2 \sin^2(\theta/2)}\right)^2$$

- ep at MeV: Rosenbluth formula(Elastic scattering)

$$\frac{d\sigma}{d\Omega} = \left(\frac{\hbar E'}{8\pi McE}\right)^2 \langle |M|^2 \rangle$$

$$\langle |M|^2 \rangle = \frac{g_e^4}{q^4} L_{\text{electron}}^{\mu\nu} K_{\mu\nu, \text{proton}}$$

$$L_{\text{electron}}^{\mu\nu} = 2 \left\{ p_1^\mu p_3^\nu + p_3^\mu p_1^\nu + g^{\mu\nu} \left[(mc)^2 + (p_1 \cdot p_3) \right] \right\}$$

$$K_{\text{proton}}^{\mu\nu} = -K_1 g^{\mu\nu} + \frac{K_2}{(Mc)^2} p_2^\mu p_2^\nu + \frac{K_3}{(Mc)^2} q_2^\mu q_2^\nu + \frac{K_5}{(Mc)^2} (p_2^\mu q^\nu + p_2^\nu q^\mu)$$

Thus,

$$\frac{d\sigma}{d\Omega})_{\text{Rosenbluth}} = \left(\frac{\alpha\hbar}{24ME \sin^2(\theta/2)} \right)^2 \frac{E'}{E} \left[2K_1 \sin^2(\theta/2) + K_2 \cos^2(\theta/2) \right]$$

- ep at GeV: Inelastic collisions

$$\frac{d\sigma}{dE' d\Omega} = \left(\frac{\alpha \hbar}{cq^2} \right)^2 \frac{E'}{E} \langle |M|^2 \rangle$$

$$\langle |M|^2 \rangle = \frac{g_e^4}{q^4} L_{electron}^{\mu\nu} W_{\mu\nu, proton}$$

$$W_{proton}^{\mu\nu} = -W_1 g^{\mu\nu} + \frac{W_2}{(Mc)^2} p_2^\mu p_2^\nu + \frac{W_4}{(Mc)^2} q_2^\mu q_2^\nu + \frac{W_5}{(Mc)^2} (p_2^\mu q_2^\nu + p_2^\nu q_2^\mu)$$

$$\frac{d\sigma}{dE' d\Omega} = \left(\frac{\alpha \hbar}{2E \sin^2(\theta/2)} \right)^2 \left[2W_1(q^2, x) \sin^2(\theta/2) + W_2(q^2, x) \cos^2(\theta/2) \right]$$

$$\text{where...} x = -\frac{q^2}{2q \cdot p}$$

<Bjorken function>

$$MW_1(q^2, x) \rightarrow F_1(x),$$

$$\frac{-q^2}{2Mc^2x} W_2(q^2, x) \rightarrow F_2(x)$$

<Callan & Gross relation>

$$2xF_1(x) = F_2(x)$$

- ep at TeV:

Hard particles: More than 80%, coming from elastic collision between e-partons.

Soft particles: relatively small number.

→ the parton distribution of nucleon is very important.

At RHIC energy, about half of them are hard particles and the other half soft particles.

4. nn, AA scattering

- pp scattering at MeV:
 - Mostly elastic collision
- pp scattering at GeV:
 - Inelastic collision with lots of soft particles
- pp scattering at TeV:
 - Elastic collisions among partons with minijets

Models

- Schwinger mechanism:
- Inside-Outside cascade:

- Classical String Model:
- Dual Parton Model:

- Parton-Parton elastic scattering (factorization):
- CGC shattering(Yang-Mills equation):

Toy models

- Schwinger mechanism:
 - two infinite plates with charges(condensers):
 - constant electric field between them
 - vacuum persistence probability: P_V
 - particle production= $1-P_V$

Assume the constant color electric field(flux) between q and q_{bar} .

If field energy $>$ threshold energy, (qq_{bar}) pair produced. \rightarrow

Production rate: quark field equation in the potential \rightarrow Klein-Gordon equation \rightarrow Scheodinger equation to identify the effective potential \rightarrow sea quark tunnels the potential to leave antiquark \rightarrow WKB approximation gives the tunneling rate.

- String model(Lund model):

- Meson: q-q_bar connected with string in yo-yo state.
- Suppose a string connected to quark and antiquark separating each other

- See Figure

- Quark-antiquark born at V_i

- $V_i = (x,t) \rightarrow (t+x, t-x) = (\Gamma, y), \Gamma = \sqrt{\kappa\tau}, y = \textit{rapidity}$

- Vertices are on the proper time: At high energy, the produced particles are independent of rapidity

- The probability element: for example, 2 vertices production

- V_0 : $dP_0 = \rho(\Gamma_0, y_0) d\Gamma_0 dy_0$

- V_1 : $dP_1 = f_{01}(z_+, M_{01}^2) g(M_{01}^2) dz_+ dM_{01}^2$

Where

$$z_+ = (p_{0+} - p_{1+}) / p_{0+}, \quad M_{01}^2 = (p_{0+} - p_{1+})(p_{1-} - p_{0-})$$

figure 1

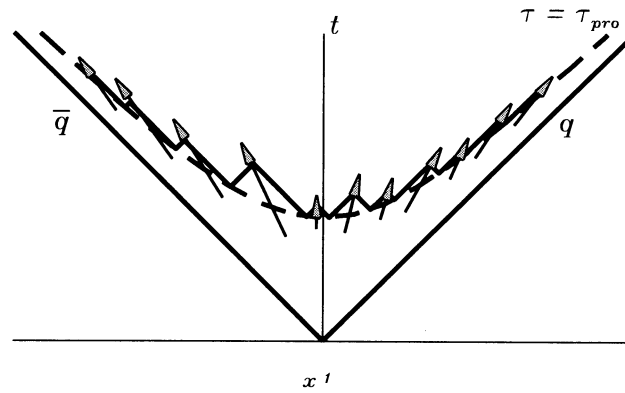


Fig. 6.4 A schematic picture of particle production in the *inside – outside cascade* picture.

– The solution:

$$g(M_{01}^2) = \sum_{\lambda} c_{\lambda} \delta(M_{01}^2 - m_{\lambda}^2)$$

$$f(z_+, m^2) = N \frac{(1 - z_+)^a}{z_+} e^{-bm^2/z_+}$$

$$\rho(\Gamma) = C\Gamma^a e^{-b\Gamma}$$

– FRITOF etc

- Dual Resonance Model:

- Consider H–H collision at high energy. The exchange of a particle of mass M , spin J : $A \propto s^J$

- Sum of J : Veneziano scattering amplitude

$$A(s, t) = g^2 \frac{\Gamma(-\alpha(s))\Gamma(-\alpha(t))}{\Gamma(-\alpha(s) - \alpha(t))}$$

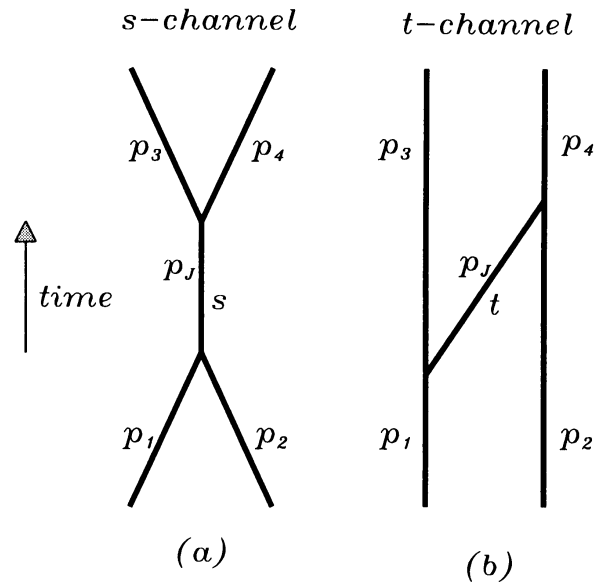
where... $\alpha(x) = \alpha_0 + \alpha' x$.

- For fixed t , expand in s . Each term corresponds to a t -channel Feynman diagram with exchange of a particle:

$$J = \alpha_0 + \alpha' M_J^2$$

- Regge trajectory \rightarrow Reggeon
- Same for s -channel:

- figure2



Symmetrize over possible channels:

$$A_4 = A(s, t) + A(t, u) + A(u, s)$$
$$\approx -g^2 \pi e^{-\alpha(t)} \frac{(1 + e^{-i\pi\alpha(t)})}{\Gamma(1 + \alpha(t))} \frac{(\alpha' s)^{\alpha(t)}}{\sin(\pi\alpha(t))}$$

- Basic properties of HH collisions at high energy:
 - Inelastic processes with no quantum number flow=dominant
 - Forward scattering amplitude \sim purely imaginary
 - Total cross section \sim constant.
- Out of Regge trajectory, $\alpha_0 = 1$ satisfies these.
 - \Rightarrow Pomeron trajectory: Pomeron

Elastic scattering between partons:

- Simple and Clear!
- Inclusive gluon production:

$$\frac{d\sigma}{dyd^2p_t} = \frac{16N_c}{N_c^2 - 1} \frac{1}{p_t^2} \int \alpha_s(\Omega^2) F_1(x_1, k_1^2) F_2(x_2, k_2^2) \delta(\vec{k}_1 + \vec{k}_2 - \vec{p}_t) d^2k_1 d^2k_2$$

where

$$x_{1/2} = \frac{p_t}{\sqrt{s}} \exp(\pm y)$$

$$\Omega^2 = \max(k_1^2, k_2^2, p_t^2)$$

Pros: Pretty good with minijets

Cons: no soft particles

- Minijets:

$$\frac{dN^{jet}}{dydp_t} = KT(b) \int dy_2 \frac{2\pi p_t}{\hat{s}} \sum_{ijkl} x_1 f_{i/A}(x_1, p_t^2) x_2 f_{j/B}(x_2, p_t^2) \sigma_{ij \rightarrow kl}(\hat{s}, \hat{t}, u)$$

- Unintegrated gluon distribution:

- BFKL: solution of BFKL equation

$$f(x, k_t^2) = \frac{C}{x^\lambda} \left(\frac{k_t^2}{q_0^2} \right)^{1/2} \frac{\phi_0}{\sqrt{2\pi\lambda'' \ln(1/x)}} \exp \left[-\frac{\ln^2(k_t^2 / \bar{q}^2)}{2r\lambda'' \ln(1/x)} \right],$$

where

$$\lambda = 4\bar{\alpha}_s \ln 2, \lambda'' = 28\bar{\alpha}_s \zeta(3), \bar{\alpha}_s = 3\alpha_s / \pi, \zeta(3) = 1.202,$$

$$C\bar{\phi}_0 = 1.19, r = 0.15$$

- Kimber–Martin–Ryskin: DGLAP evolution

$$F(x, k^2) = \frac{\partial}{\partial Q^2} \left[xg(x, Q^2) \right]_{Q^2=k^2},$$

$$f_a(x, k^2, \mu^2) = T_a(k^2, \mu^2) \frac{\alpha_s(k^2)}{2\rho} \sum_{a'} \int_x^{1-\delta} P_{aa'}(z) \frac{x}{z} a' \left(\frac{x}{z}, k^2 \right) dz$$

- Kharzeev–Levin: based on CGC

$$F(x, k^2) = \begin{cases} f_0 \dots \dots \dots \text{if } k^2 > Q_s^2 \\ f_0 \frac{Q_s^2}{k^2} \dots \text{if } k^2 > Q_s^2 \end{cases}$$

where

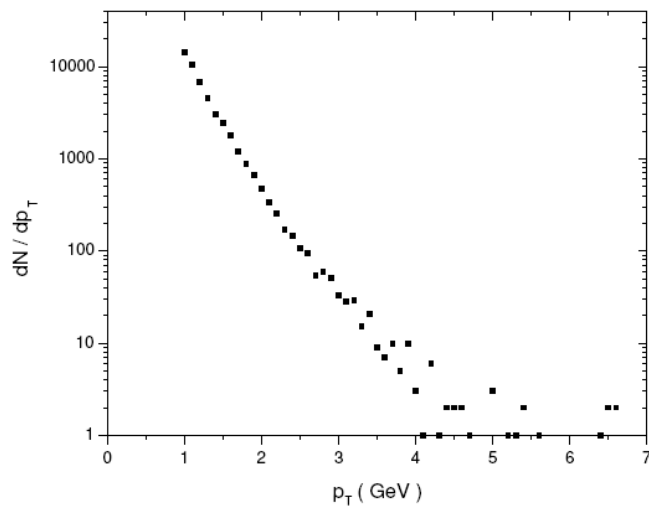
$$Q_s^2(x) = 1(\text{GeV})^2 \left(\frac{x_0}{x} \right)^2, f_0 = 170 \text{mb}$$

- GRV98: we have the code!

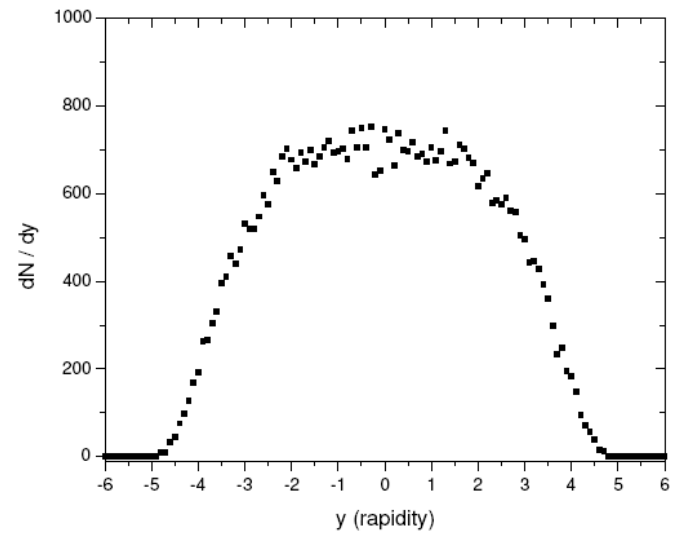
– CGC distribution: Triantafyllopoulos

$$\varphi(k_t, y) = \begin{cases} \frac{1}{\alpha_s} \ln \frac{Q_s^2}{k_t^2} \dots \dots \dots \text{if } k_t \ll Q_s \\ \frac{1}{\alpha_s} \left(\frac{Q_s^2}{k_t^2} \right) \left(\ln \frac{k_t^2}{Q_s^2} + \Delta \right) \dots \dots \dots \text{if } k_t \geq Q_s \\ \frac{Q_s^2}{k_t^2} I_0 \left(\sqrt{4\alpha_s y \ln \frac{k_t^2}{Q_s^2}} \right) \dots \dots \dots k_t \gg Q_s \end{cases}$$

- GRV98 & EKS98+minijets:



P_T distribution



rapidity distribution

- DGLAP evolution:(Dokshitzer-Gribov-Lipatov-Altarelli-Parisi)

$$\frac{\partial a(x, \mu^2)}{\partial \ln \mu^2} = \frac{\alpha_s(\mu^2)}{2\pi} \sum_{b=q,g} \int_x^1 dz P_{ab}(z, \alpha_s(\mu^2)) b\left(\frac{x}{z}, \mu^2\right),$$

where

$P_{ab} : b \rightarrow ac$, branching..kernel,

a & b : parton..distribution

- Modified DGLAP:

$$\frac{\partial a(x, \mu^2)}{\partial \ln \mu^2} = \frac{\alpha_s(\mu^2)}{2\pi} \sum_{b=q,g} \left[\int_x^1 dz P_{ab}(z, \alpha_s(\mu^2)) b\left(\frac{x}{z}, \mu^2\right) - a(x, \mu^2) \int_0^1 d\zeta \zeta P_{ba}(\zeta) \right]$$

- *BFKL equation:(Balitsky-Fadin-Kuraev-Lipatov)*

$$\frac{\partial f_g(x, k_t^2)}{\partial \ln(1/x)} = \frac{C_A \alpha_s}{\pi} k_t^2 \int_0^\infty \frac{dk_t'^2}{k_t'^2} \left[\frac{f_g(x, k_t'^2) - f_g(x, k_t^2)}{|k_t'^2 - k_t^2|} + \frac{f_g(x, k_t^2)}{(4k_t'^4 + k_t^4)^{1/2}} \right]$$

Where

$$a(x, \mu^2) = \int_0^{\mu^2} \frac{dk_t^2}{k_t^2} f_a(x, k_t^2)$$

Yang–Mills Equation:

- Yang–Mills equation of motion:
 - Krasnitz, Nara and Venugopalan; Lappi

$$\left[D_\mu, F^{\mu\nu} \right]^a = J^{\nu,a}$$

where

$$J^\nu = \rho_1 \delta(x^-) \delta^{\nu+} + \rho_2 \delta(x^+) \delta^{\nu-}$$

- Solve the equation and identify the number density:

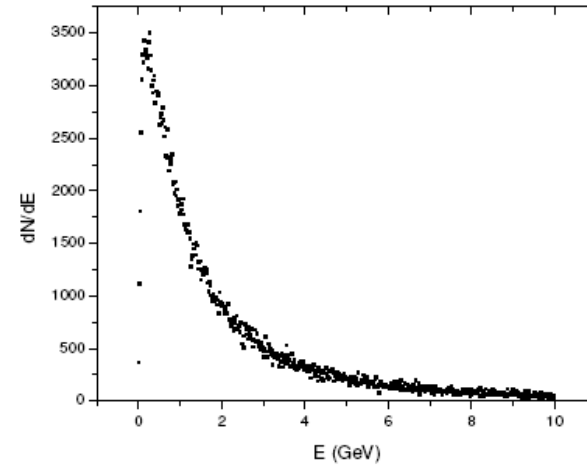
$$f^{glue} = \frac{1}{g^2} \begin{cases} a_1 \left[\exp(\sqrt{k_t^2 + m^2} / T_{eff}) - 1 \right]^{-1} \dots (k_t / \Lambda_s > 1.5) \\ a_2 \Lambda_s^4 \log(4\pi k_t / \Lambda_s) k_t^{-4} \dots \dots \dots (k_t / \Lambda_s > 1.5) \end{cases}$$

where

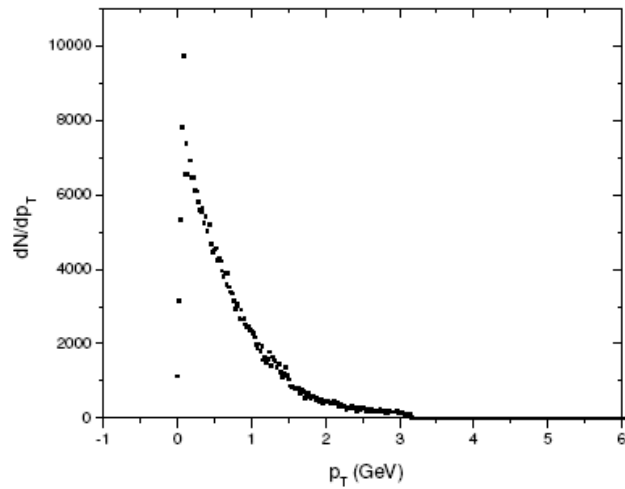
$$a_1 = 0.137, a_2 = 0.0087,$$

$$m = 0.0358 \Lambda_s, T_{eff} = 0.465 \Lambda_s$$

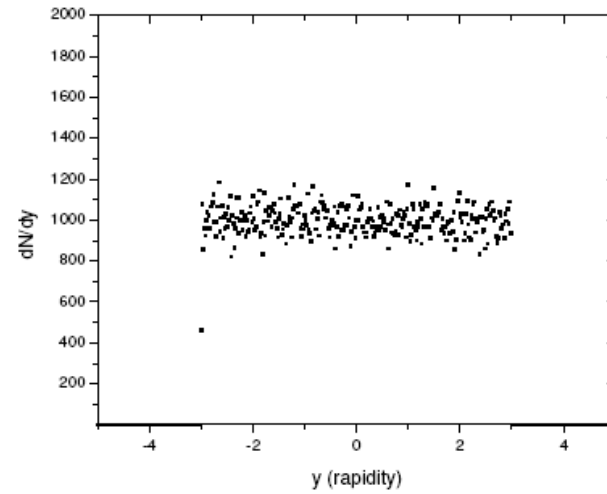
- KNV:



energy distribution



Number distribution



rapidity distribution

- Pros and cons:
 - Good: based on the first principle(QCD)
 - Bad: infinite flat plates, only gluons,
no information on space

5. Summary

- Needs lots more work on particle production.
- Of course the work is one of great problems to challenge.
- Need a new method to solve physics problems? I.e. we have a finite proton-proton cross section but we cannot calculate that in current formalism.