

Physics Opportunities with heavy quark system at FAIR

Su Houng Lee
Yonsei Univ., Korea

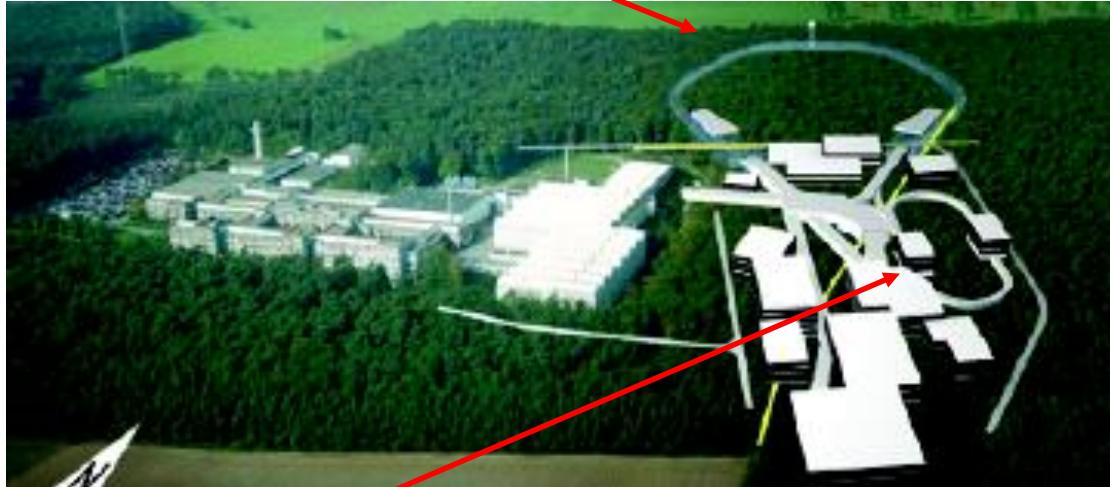
1. Introduction to FAIR
2. Heavy quark system in nuclear medium
3. Heavy exotics from heavy Ion collision
4. Summary

Introduction to FAIR

Facility for Antiproton and Ion Research
at GSI

FAIR

1. CBM 2-35 GeV Heavy Ion (Gold) projectile
QCD phase diagram at high baryon density



2. PANDA: anti proton project (1-15 GeV)
charmonium spectroscopy, origin of hadron mass

Heavy quark system in nuclear medium

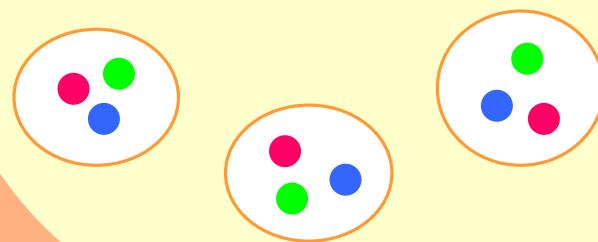
Physics with antiproton beam

QCD Vacuum

$$\langle \bar{q}q \rangle \approx -(250 \text{ MeV})^3$$

Nuclear matter

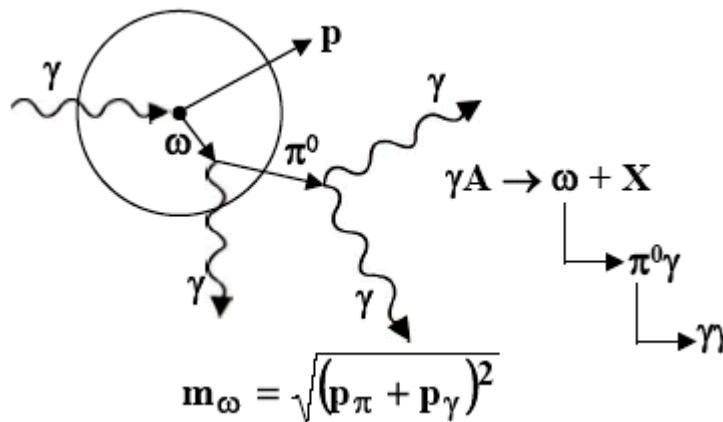
$$\langle \bar{q}q \rangle \approx -0.8 \times (250 \text{ MeV})^3$$



Partial Chiral symmetry restoration in nuclear matter

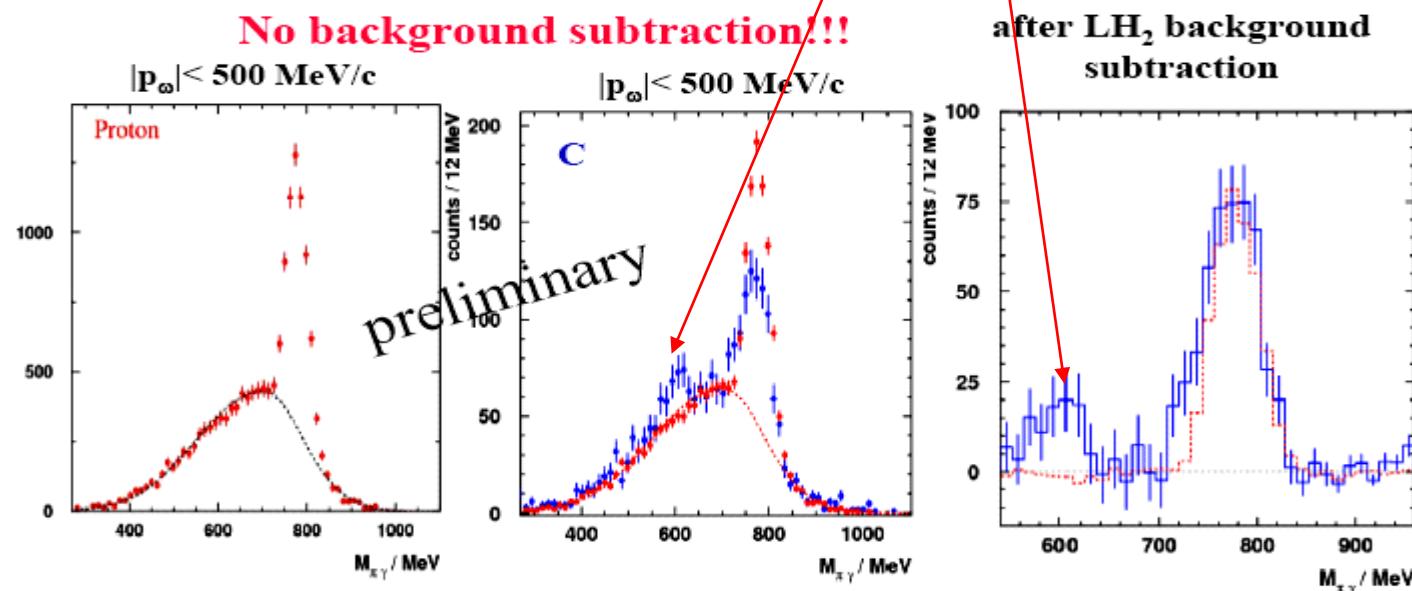
→ Could be probed by light vector mesons

From Volker Metag's recent talk



Shocking result !!

New structure appearing?



QCD Vacuum

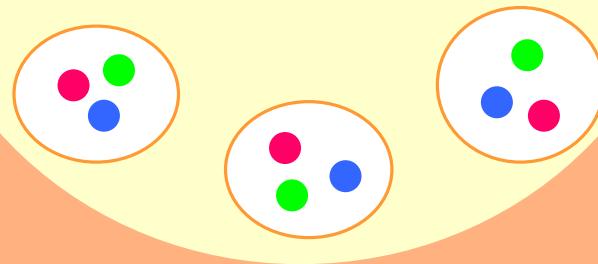
$$\langle \frac{\alpha_s}{\pi} B^2 \rangle \approx 380 \text{ MeV/fm}^3$$

$$\langle \frac{\alpha_s}{\pi} E^2 \rangle \approx -380 \text{ MeV/fm}^3$$

Nuclear matter

$$\langle \frac{\alpha_s}{\pi} B^2 \rangle \approx 365 \text{ MeV/fm}^3$$

$$\langle \frac{\alpha_s}{\pi} E^2 \rangle \approx -335 \text{ MeV/fm}^3$$



Change of gluonic background in nuclear matter

→ could be probed by heavy quark system

Why Heavy Quark system can probe this gluonic change

1. Heavy quark propagation

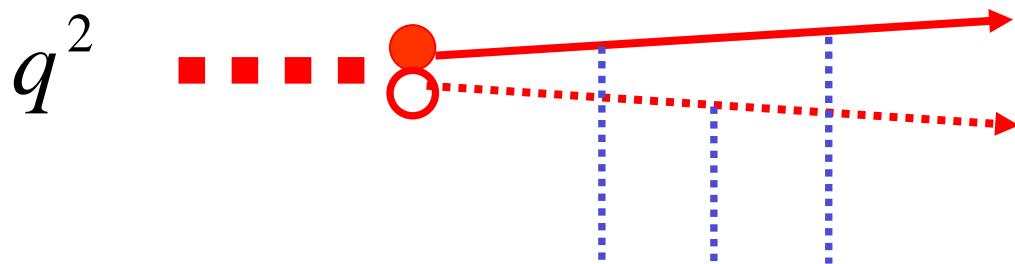


$$S_G(q) = S(q) + S(q)GS(q) + \dots \quad \text{where,} \quad S(q) = \frac{1}{q - m}$$

Perturbative treatment are possible

because $m - q \gg \Lambda_{QCD}$ even for $q \rightarrow 0$

2. System with two heavy quarks

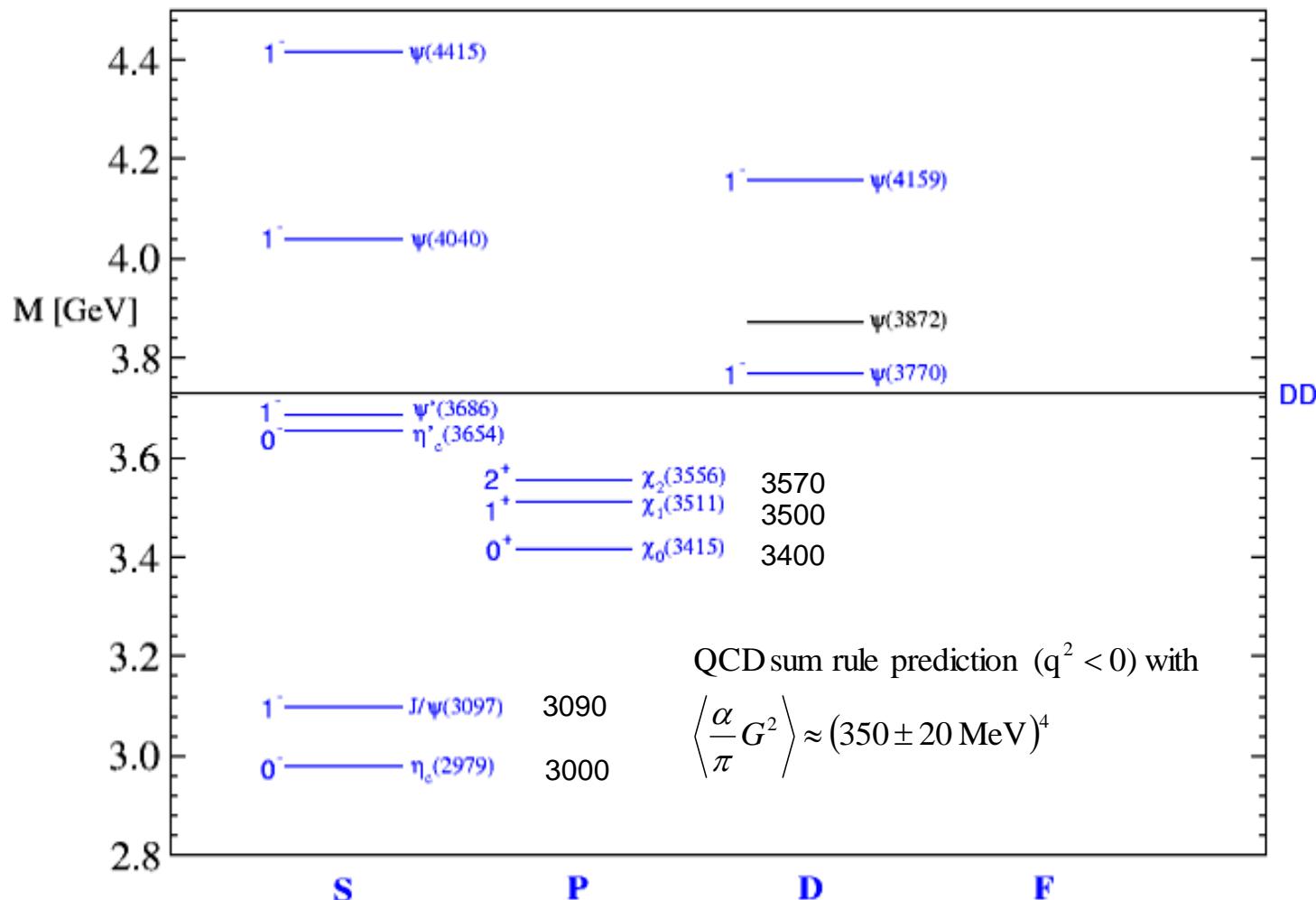


$$\Pi(q) = \dots + \int_0^1 dx \frac{F(q^2, x)}{(4m^2 - q^2 - (x - 1/2)^2 q^2)^n} \cdot \langle G^n \rangle + \dots$$

Perturbative treatment are possible when

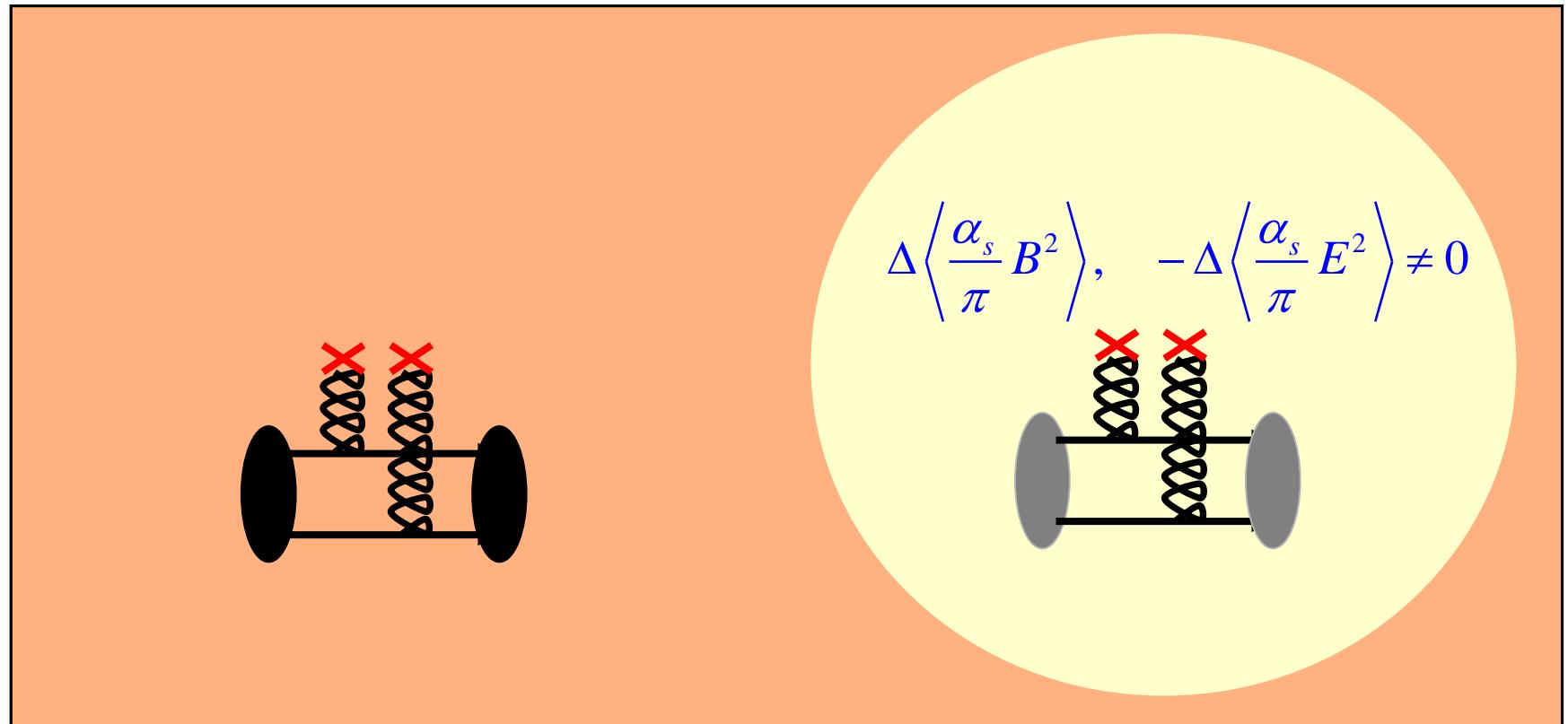
$$4m^2 - q^2 \gg \Lambda_{QCD}^2$$

$q^2 < 0$: QCD sum rules for charmonium (cc)



Heavy quark propagator is sensitive to only gluon fields

and therefore the mass of heavy quark system will change
in nuclear medium



Approaches for charmonium mass shift in nuclear matter

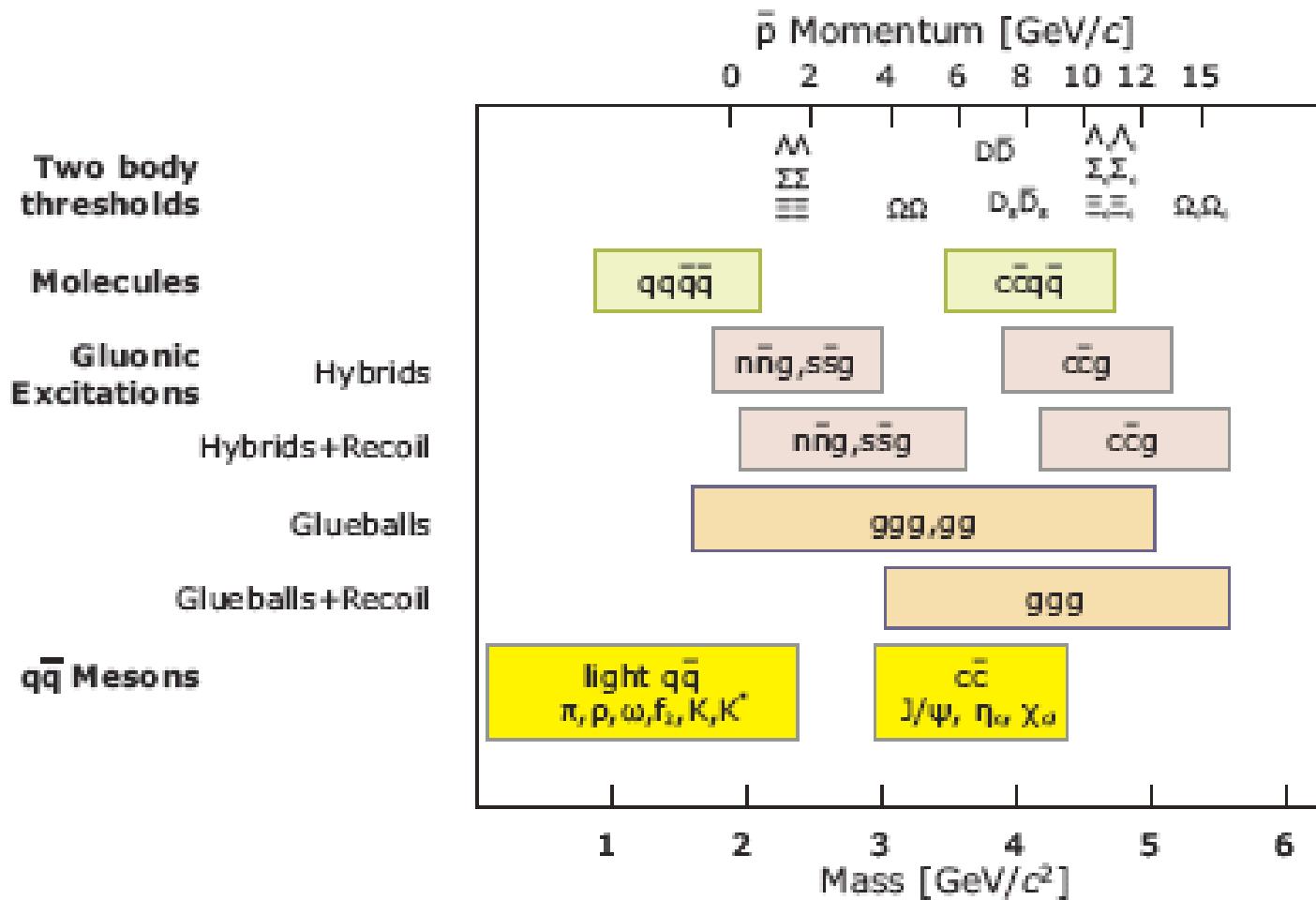
	Quantum numbers	QCD 2 nd Stark eff.	Potential model	QCD sum rules	Effects of DD loop
H_c	0^{-+}	-8 MeV		-5 MeV (Klingl, SHL, Weise)	No effect (SHL, Ko)
J/ψ	1^{--}	-8 MeV (Peskin, Luke)	-10 MeV (Brodsky et al.).	-7 MeV (Klingl, SHL, Weise)	<2 MeV (SHL, Ko)
$\chi_{0,1,2}$	$0, 1, 2^{++}$	-40 MeV (SHL)		-60 MeV (SHL)	No effect on χ_1
$\psi(3686)$	1^{--}	-100 MeV (SHL)			< 30 MeV (SHL, Ko)
$\psi(3770)$	1^{--}	-140 MeV (SHL)			< 30 MeV (SHL, Ko)

Such experiment can be done at

1. GSI future accelerator facility

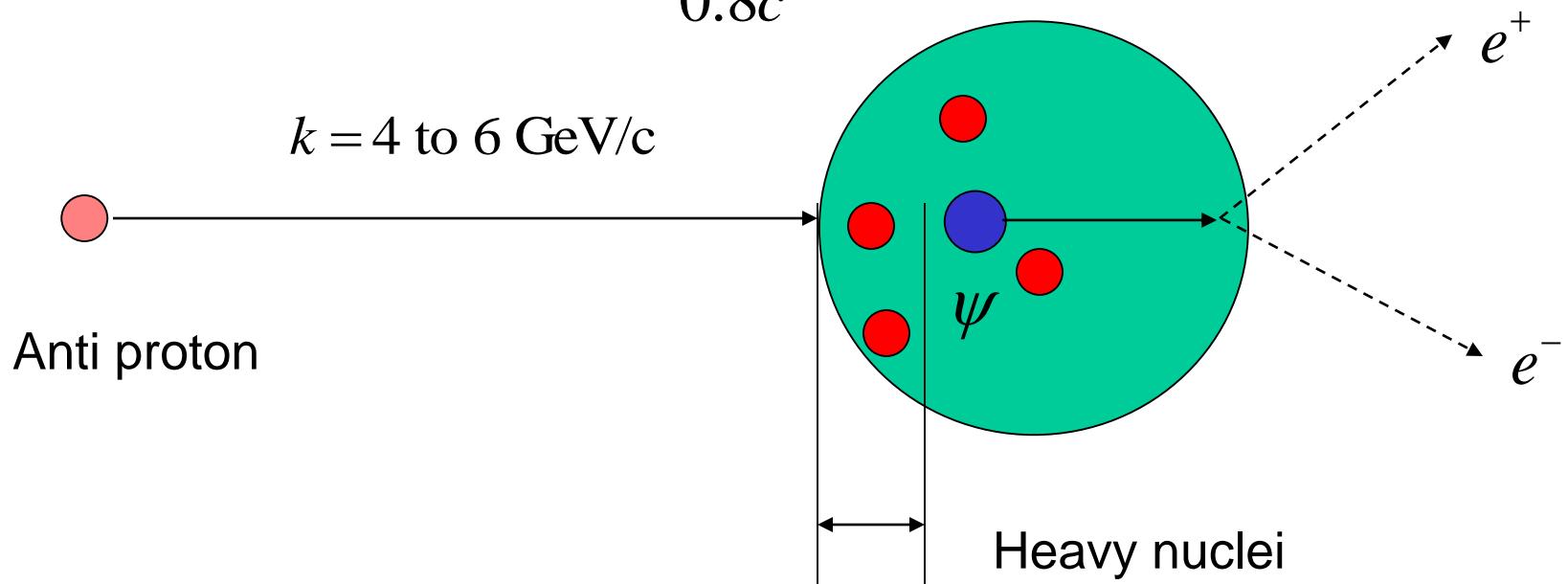
⇒ anti proton project (1-15 GeV)





Anti proton project at GSI

$$t = \frac{2 \times 1.3 A^{1/3} \gamma^{-1}}{0.8c} = 10 \text{ fm}/c \quad \text{for } A=125$$



$$\frac{1}{0.17 \text{ fm}^{-3} \times 5 \text{ fm}^2} = 1.2 \text{ fm}$$

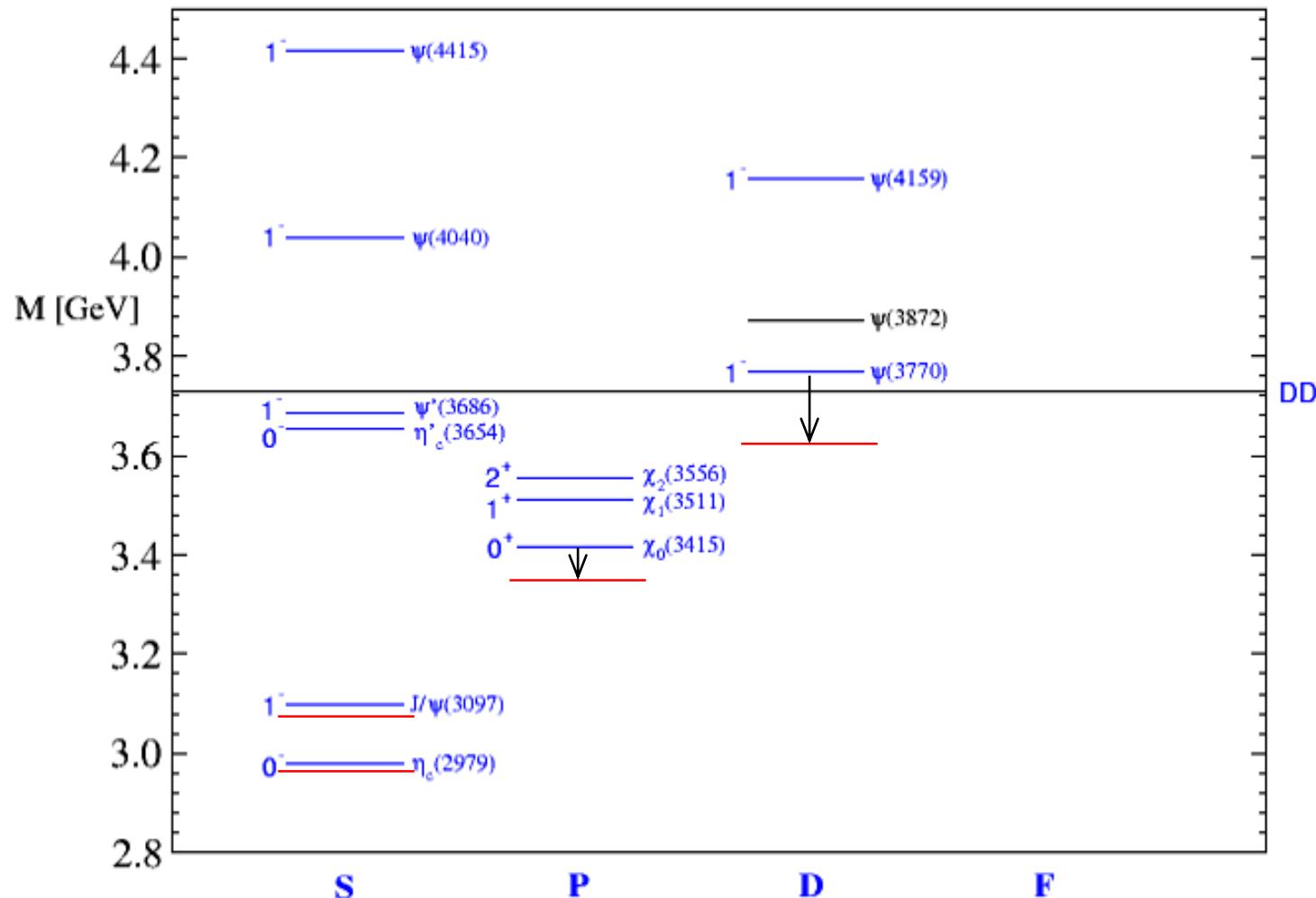
Anti proton will be absorbed at surface and
Charmonium will decay inside the heavy nuclei

Production rate: Luminosity = $2 \times 10^{32} \text{ cm}^{-2} \text{ s}^{-1}$

$$\sigma_{BW} = \frac{2J+1}{(2s_1+1)(2s_2+1)} \frac{\pi}{k^2} \frac{B_{in} B_{out} \Gamma_{Total}^2}{(E - E_R)^2 + \Gamma_{Total+medium}^2 / 4}$$

	J ^{PC}	Mass shift	Final state	σ to final state	Events per day
J/ ψ	1 --	-8 MeV	e ⁺ + e ⁻	6 pb	100
$\psi(3686)$	1 --	-100 MeV	e ⁺ + e ⁻	0.6 pb	10
$\psi(3770)$	1 ---	-100 MeV	e ⁺ + e ⁻	1 pb	17
χ_0	0 ++	-50 MeV	J/ ψ + γ	200 pb	3400
χ_1	1 ++	-50 MeV	J/ ψ + γ	80 pb	1360
χ_{12}	0 ++	-50 MeV	J/ ψ + γ	350 pb	5950

Charmonium spectrum



Similar to discoveries in E&M

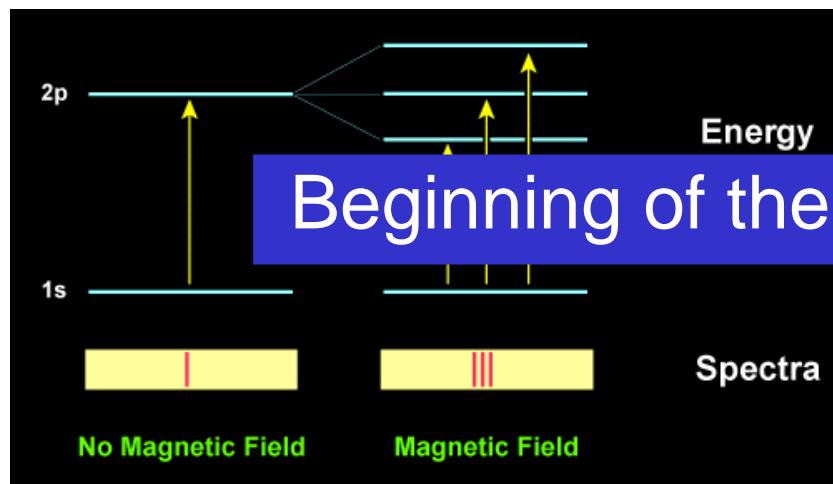
Pieter Zeeman (1865-1943)



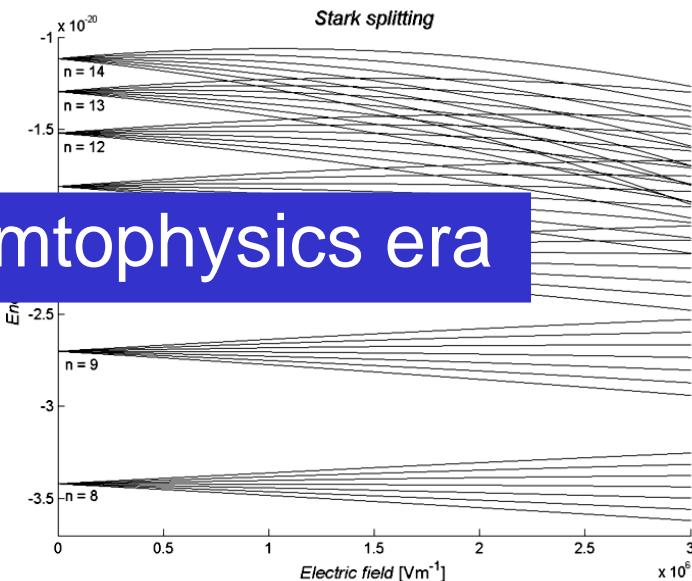
Johannes Stark (1874-1957)



Zeeman effect: Nobel prize 1902



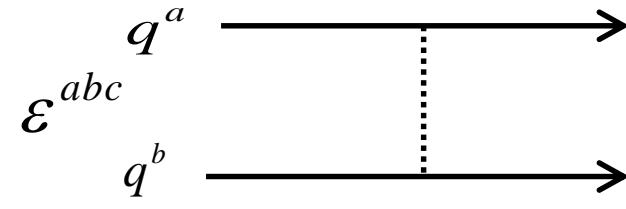
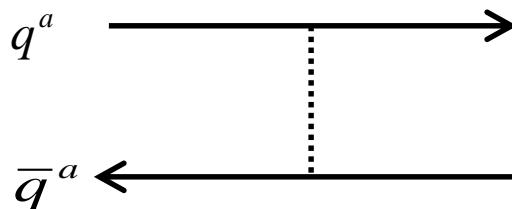
Stark effect: Nobel prize 1919



Heavy exotics from Heavy Ion collision

Physics with CBM

Attractions in quark-antiquark vs. diquark



$$-\frac{4\alpha}{3} \frac{1}{r}$$

$$E \propto -\frac{m_i m_j}{m_i + m_j} \times Z^2$$

$$-\boxed{\frac{1}{2}} \frac{4\alpha}{3} \frac{1}{r}$$

$$\sigma \times r$$

$$E \propto -\sqrt{\frac{m_i + m_j}{m_i m_j}} \times Z^{1/2}$$

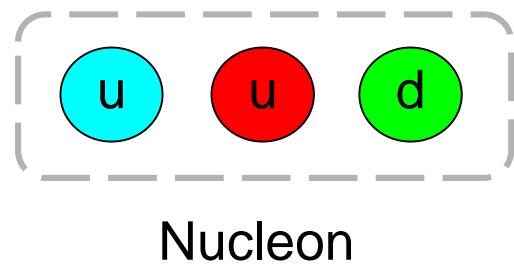
$$\boxed{\frac{1}{2}} \sigma \times r$$

$$\frac{C}{m_i m_k} \sum S_i \cdot S_k$$

Phenomenological fit to
hadron spectrum
color spin interaction

$$\boxed{\frac{1}{3}} \frac{C}{m_i m_k} \sum S_i \cdot S_k$$

Color spin interaction explains hadron spectrum



$$\begin{aligned} & \frac{C_B}{m_i m_k} [s_u \cdot s_u + s_u \cdot s_d + s_u \cdot s_d] \\ &= \frac{C_B}{m_q^2} \frac{1}{2} [(s_u + s_u + s_d)^2 - s_u^2 - s_u^2 - s_d^2] \end{aligned}$$

Baryon Mass difference

Mass Difference	$M_\Delta - M_N$	$M_\Sigma - M_\Lambda$	$M_{\Sigma_c} - M_{\Lambda_c}$
Formula	$\frac{3C_B}{2m_c^2}$	$\frac{C_B}{m_u^2}(1 - \frac{m_u}{m_s})$	$\frac{C_B}{m_u^2}(1 - \frac{m_u}{m_c})$
Fit	290 MeV	77 MeV	154 MeV
Experiment	290 MeV	75 MeV	170 MeV

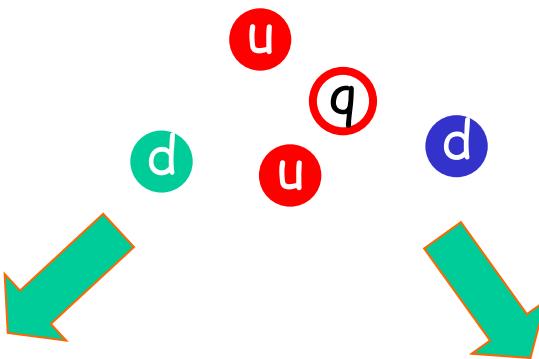
Meson Mass difference

Mass Difference	$M_\rho - M_\pi$	$M_{K^*} - M_K$	$M_{D^*} - M_D$
Formula	$\frac{C_M}{m_u^2}$	$\frac{C_M}{m_u m_s}$	$\frac{C_M}{m_u m_c}$
Fit	635 MeV	381 MeV	127 MeV
Experiment	635 MeV	397 MeV	137 MeV

Works very well with $3C_B = C_M = \text{constant}$

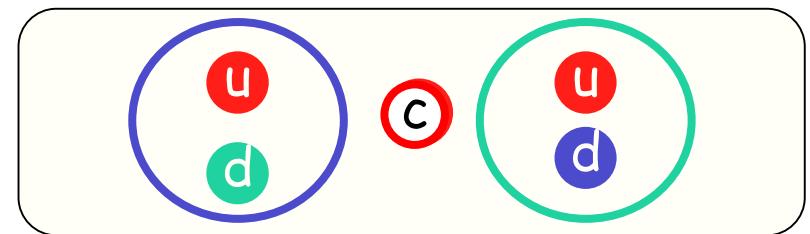
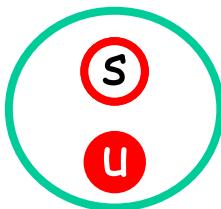
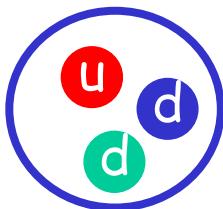
$$m_u = m_d = 300 \text{ MeV}, \quad m_s = 500 \text{ MeV}, \quad m_c = 1500 \text{ MeV}$$

When will a pentaquark form ?



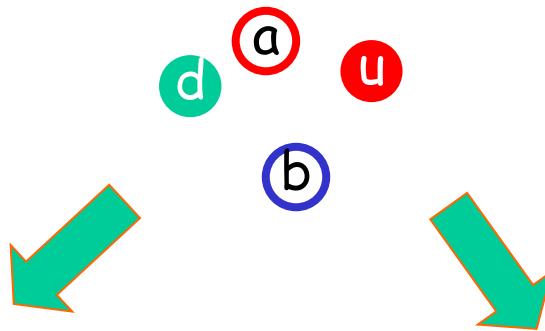
Strong quark-antiquark
attraction

Weak quark-antiquark
attraction



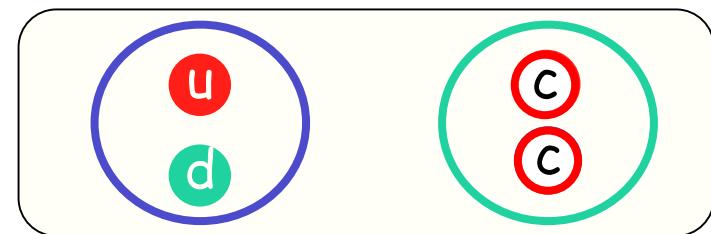
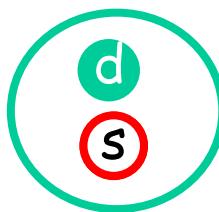
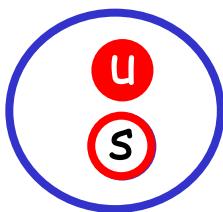
$$E_{\text{pentaquark}} - E_{\text{meson-Baryon}} = -50 \text{ MeV}$$

When will a tetraquark form ?



Strong quark-antiquark
attraction

Weak quark-antiquark
attraction



$$E_{\text{tetraquark}} - E_{\text{2-meson}} = +38 \text{ MeV}$$

Success of statistical model

P. Braun-Munzinger, J. Stachel (95 ...)

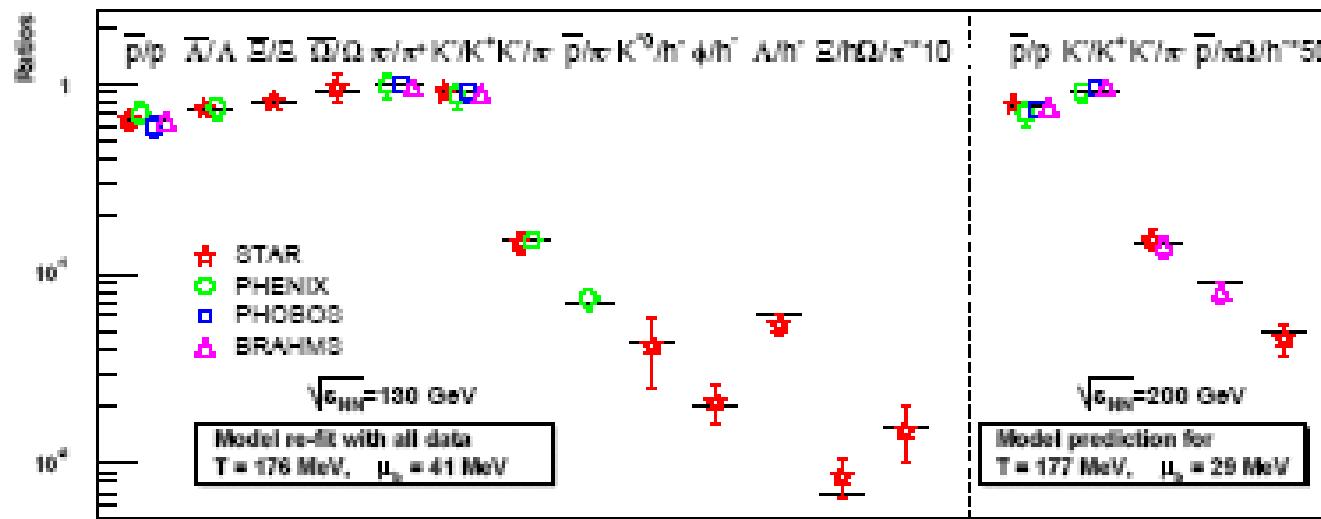
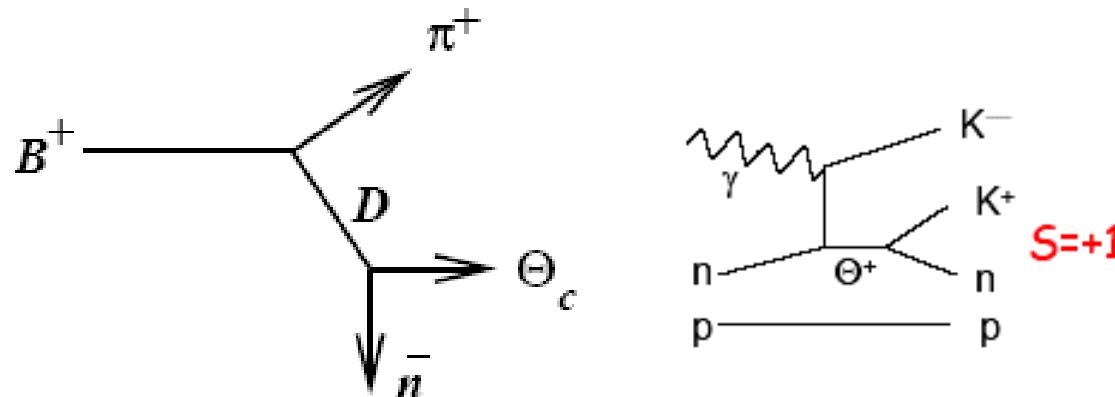


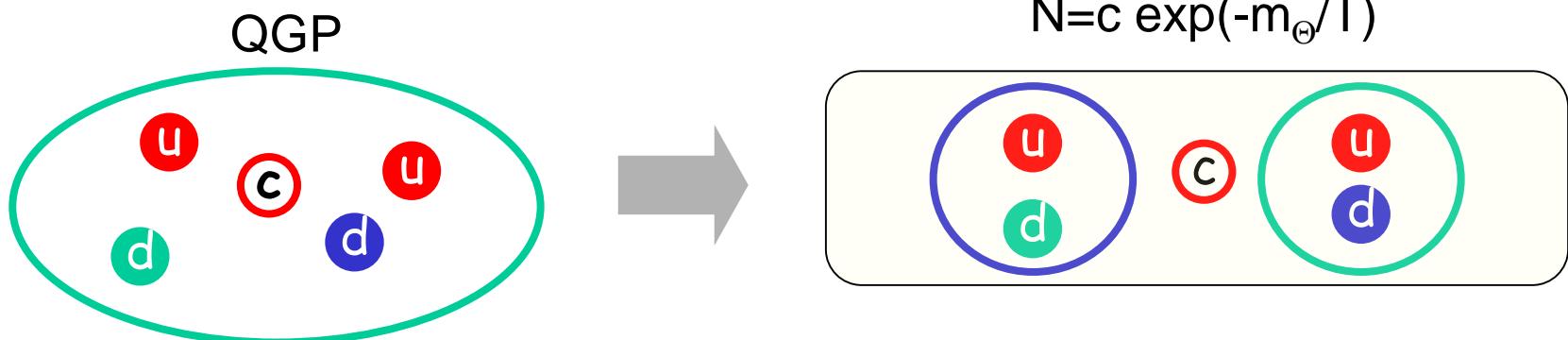
Figure 1. Fit of particle ratios for Au-Au collisions measured at RHIC energies. The measurements are the symbols, the thermal model values are the lines [6, 10]

Exotic particle production: elementary vs HIC

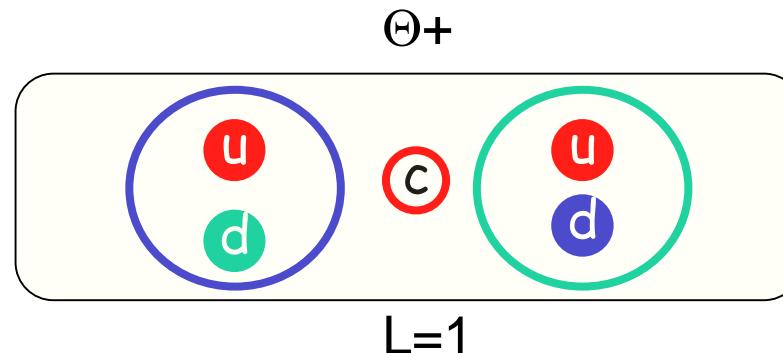
Exotic particle production from elementary processes



Exotic particle production from Heavy Ion collision

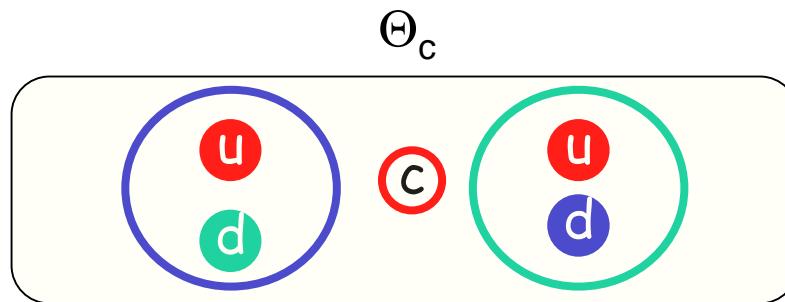


Possible Decay mode of charmed Pentaquark



exotic	Decay mode	Final states	Branching ratio
udus <u>c</u>	uds(Λ) + <u>c</u> (D^0)	$\Lambda + k^0 \pi^0$ ($k^+ \pi^-$) $\Lambda + k^0 \pi^- \pi^+$	2.3 % (3.8 %) 5.97 %
	uud(p) + <u>s</u> <u>c</u> (Ds-)	$P + k^0 \pi^-$ $P + k^+ \pi^- \pi^-$	2.82 % 9.2 %
udud <u>c</u>	udu (p)+ <u>d</u> <u>c</u> (D-)	$P + k^0 \pi^-$ $P + k^+ \pi^- \pi^-$	2.82 % 9.2 %
	udd (n)+ <u>c</u> (D^0)	$n + k^0 \pi^0$ ($k^+ \pi^-$) $n + k^0 \pi^- \pi^+$	2.3 % (3.8 %) 5.97 %

Rough estimate of events at RHIC



Final state	Decay mode	Final states	Branching ratio
udud <u>c</u>	udu (p)+d <u>c</u> (D-) p + k ⁰ π ⁻	P + k ⁺ π ⁻ π ⁻	2.82 %
	udd (n)+u <u>c</u> (<u>D</u> ⁰)	n + k ⁰ π ⁰ (k ⁺ π ⁻) n + k ⁰ π ⁻ π ⁺	9.2 % 2.3 % (3.8 %) 5.97 %

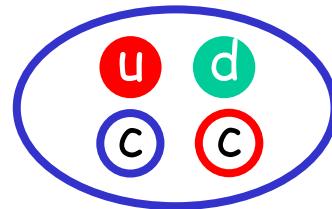
Central collision (0-10%)

$dN_{cc}/dy = 2.2$ (phenix)

$$dN_{\Theta_c} / dy = dN_D / dy \times \exp[-m_p / T] \times \text{Branching ratio}$$

$$= 2.2 \times \frac{1}{148} \times 0.03 = 0.00045$$

Possible Decay mode of Tetra-quark



exotic	Decay mode	Final states	Branching ratio
<u>uscc</u>	$u\underline{c} (\underline{D^0}) + s\underline{c} (D_s^-)$	$\underline{D^0} (k^0 \pi^0, k^+ \pi^-)$ $k^0 \pi^- \pi^+$ $D_s^- (k^0 k^-, k^- k^+ \pi^-)$	3.6 % (4.4 %)
<u>udcc</u>	$d\underline{c} (D^-) + u\underline{c} (\underline{D^0})$	$D^- (k^0 \pi^- k^+ \pi^- \pi^-)$ + $\underline{D^0} (k^0 \pi^0, k^+ \pi^-)$ $k^0 \pi^- \pi^+$	

If STAR can do it, CBM can do much better

STAR

$$D^0 \rightarrow k^- \pi^+$$

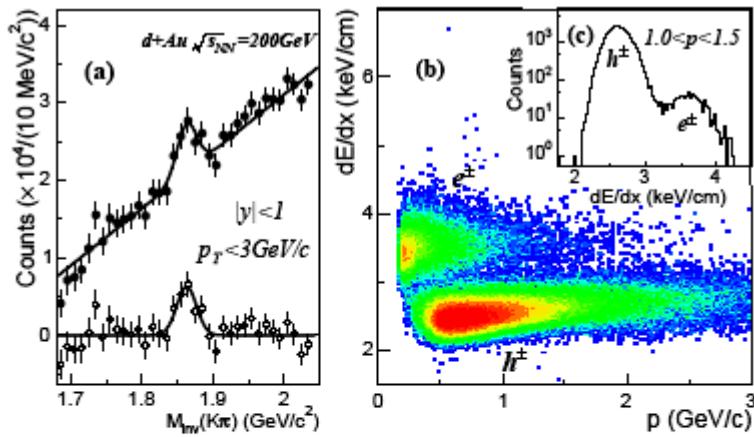


FIG. 1: (a) Invariant mass distributions of kaon-pion pairs from d+Au collisions. The solid circles depict the signal after mixed-event background subtraction, the open circles after subtraction of the residual background using a linear parametrization. (b) dE/dx in the TPC vs. particle momentum (p) with a TOF cut of $|1/\beta - 1| \leq 0.03$. Insert: projection on the dE/dx axis for particle momenta $1 < p < 1.5 \text{ GeV}/c$.

CBM

4 month Au+Au at 25 AGeV CBM

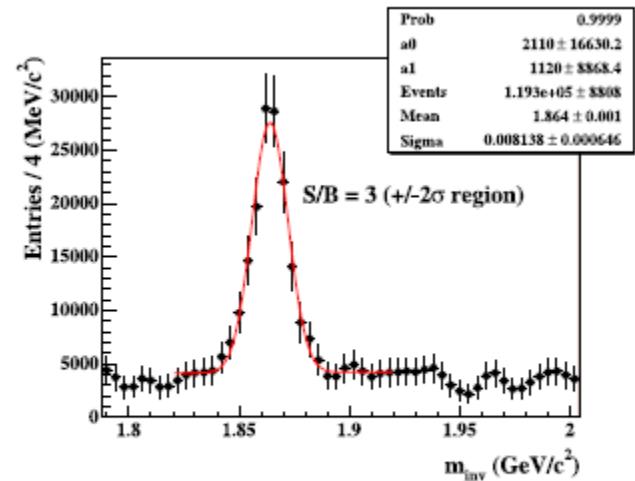
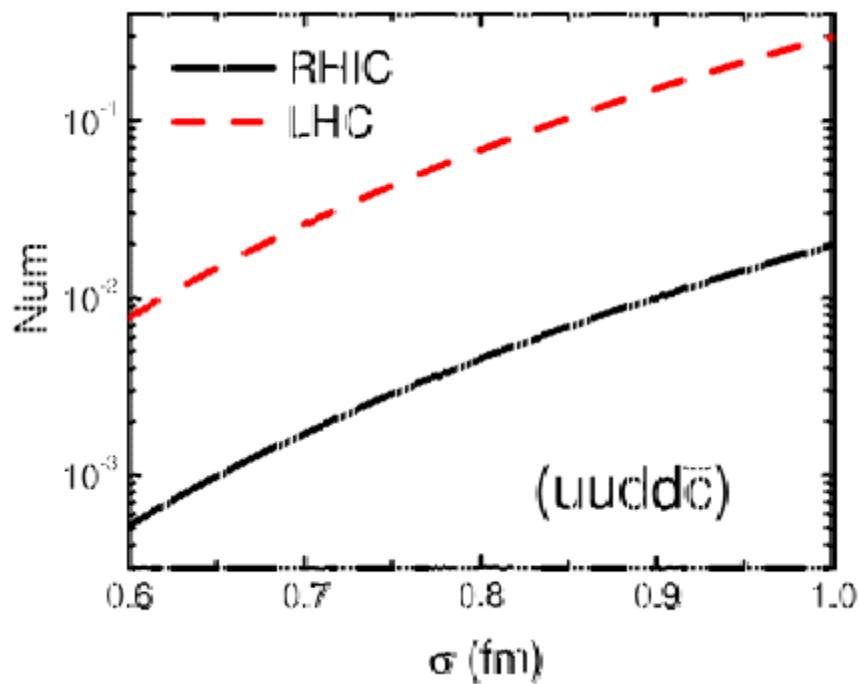


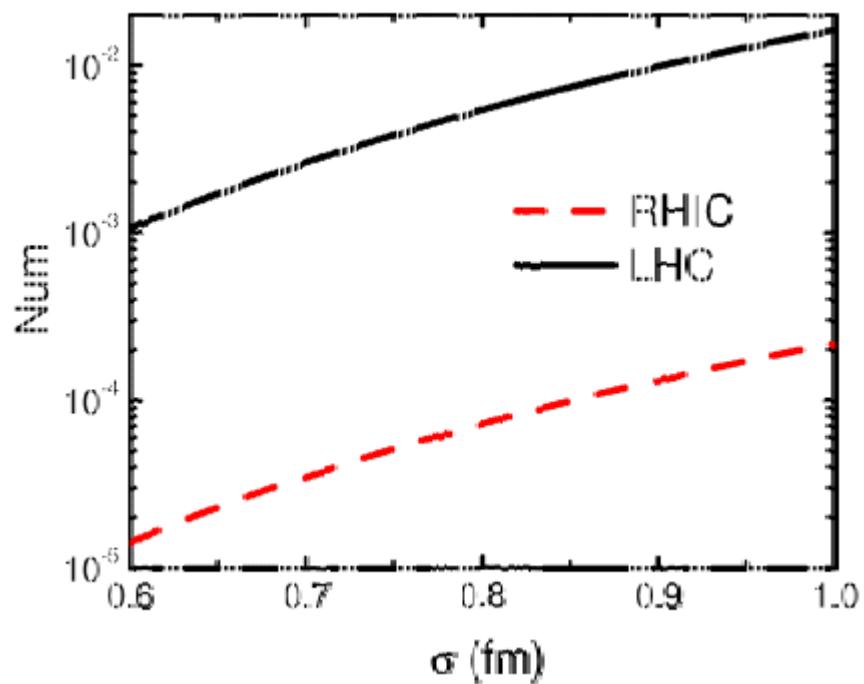
Figure 19.12: Invariant mass spectrum for the simulated signal plus background. The D^0 peak is clearly visible. Estimated data taking time is about 4 months at 0.1 MHz interaction rate.

And also at LHC (W.Liu, C.M.Ko, SHL 07)

udud c



ud cc



Summary

1. From PANDA (anti proton beam), can look at the charmonium mass shift in nuclear medium
→ Beginning of Femto Physics
Hint to confinement and QCD vacuum
3. From CBM (heavy Ion Physics), can look at heavy Exotics
→ New exotic hadron in QCD
If found the first real exotic ever, will tell us about QCD and dense matter → color superconductivity

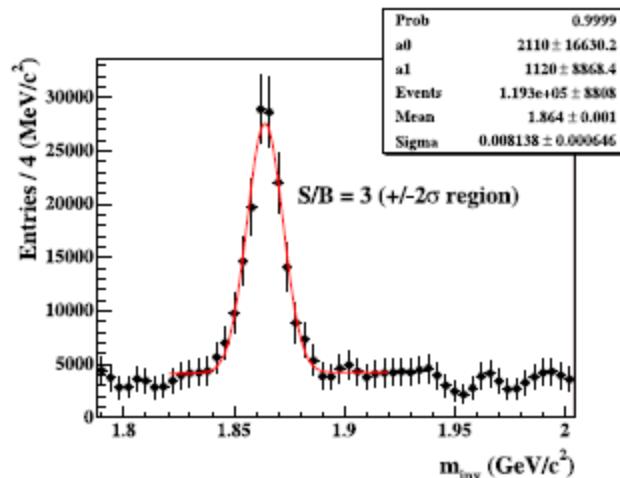
beam energy (AGeV)	J/ ψ mult. min. bias	J/ ψ yield detected per week	runtime (weeks)
10	$5 \cdot 10^{-8}$	$1.8 \cdot 10^3$	10
15	$6 \cdot 10^{-7}$	$2.2 \cdot 10^4$	10
20	$2 \cdot 10^{-6}$	$7 \cdot 10^4$	10
25	$5 \cdot 10^{-6}$	$1.8 \cdot 10^3$	5
30	$1.0 \cdot 10^{-5}$	$3.6 \cdot 10^5$	2
35	$1.5 \cdot 10^{-5}$	$5.4 \cdot 10^5$	1

Table 12.1: Expected statistics and run time for J/ ψ measurements in minimum bias Au+Au collisions

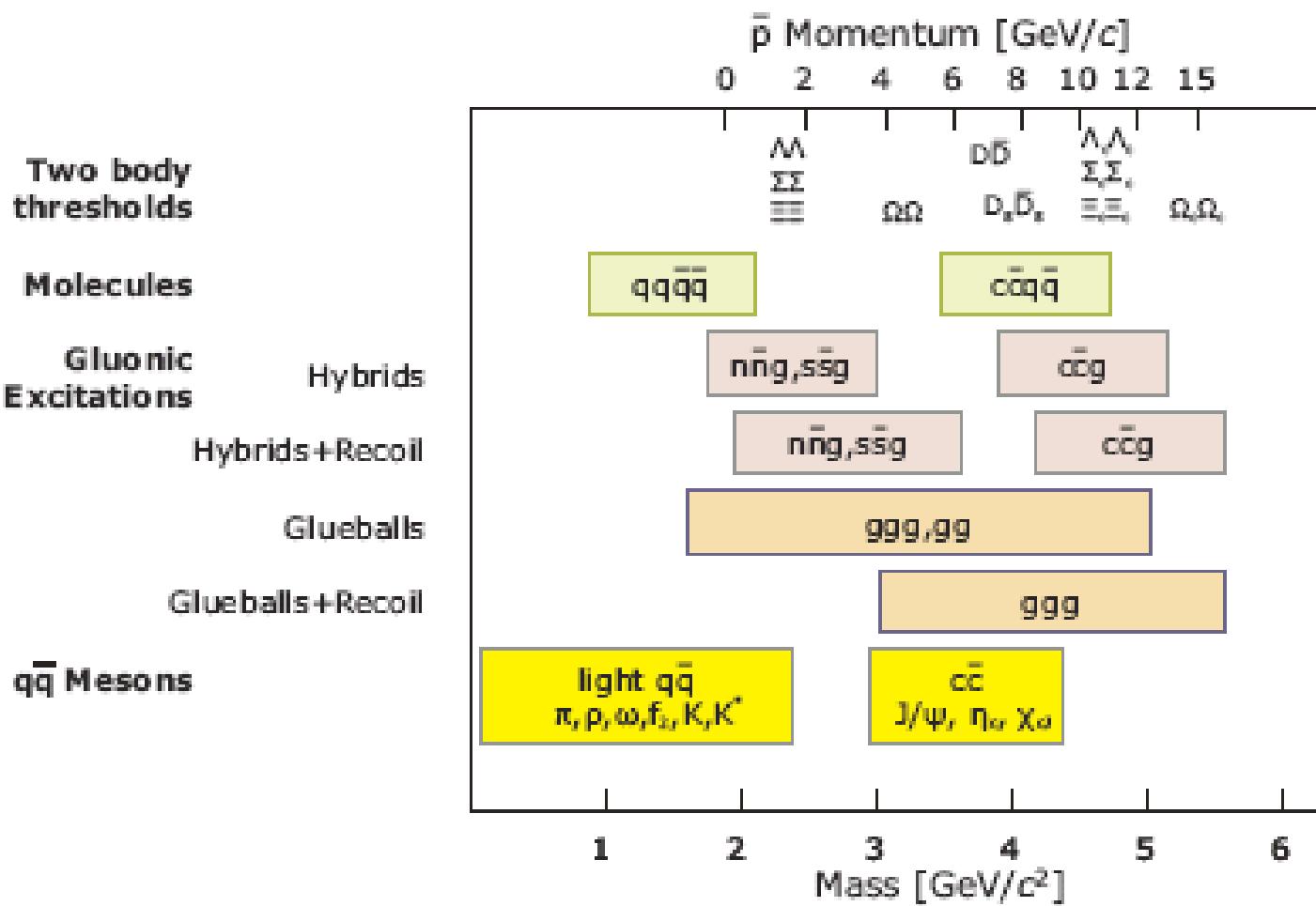
beam energy (AGeV)	D ⁰ mult. min. bias	D ⁰ yield detected per week	runtime (weeks)
15	$7 \cdot 10^{-7}$	$8 \cdot 10^3$	10
20	$7 \cdot 10^{-6}$	$8 \cdot 10^4$	10
25	$3 \cdot 10^{-5}$	$3.3 \cdot 10^5$	5
30	$7 \cdot 10^{-5}$	$8 \cdot 10^5$	2
35	$1.3 \cdot 10^{-4}$	$1.5 \cdot 10^6$	1

Table 12.2: Expected statistics and run time for D⁰ meson measurements in minimum bias Au+Au collisions

4 month Au+Au at 25 AGeV CBM



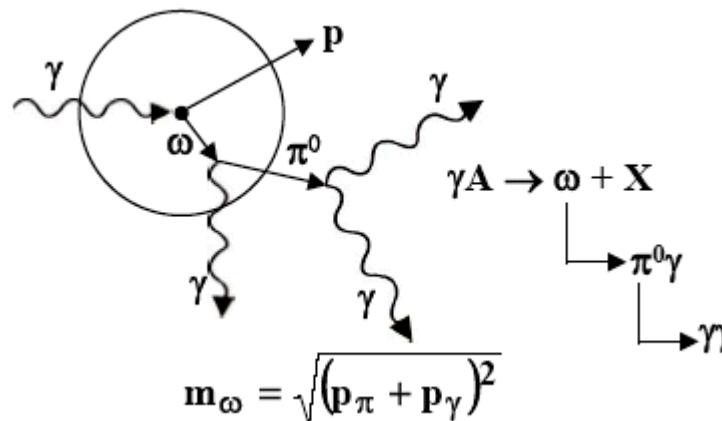
12: Invariant mass spectrum for the simulated signal plus background. The D⁰ peak is clearly visible. The taking time is about 4 months at 0.1 MHz interaction rate.



Perturbative treatment are possible when $\mu^2 = 4m^2 - q^2 \gg \Lambda_{QCD}^2$

q^2	process	expansion parameter
negative	QCD sum rules for heavy quarks	$\frac{\Lambda_{QCD}^2}{4m^2 + Q^2}$
0	Photo production of open charm	$\frac{\Lambda_{QCD}^2}{4m^2}$
$m_{J/\psi}^2 > 0$	Dissociation cross section of bound states	$\frac{\Lambda_{QCD}^2}{(2m + m_{J/\psi}) \cdot \epsilon_0}$

Example: from Volker Metag's recent talk



Shocking result !!

New structure appearing?

