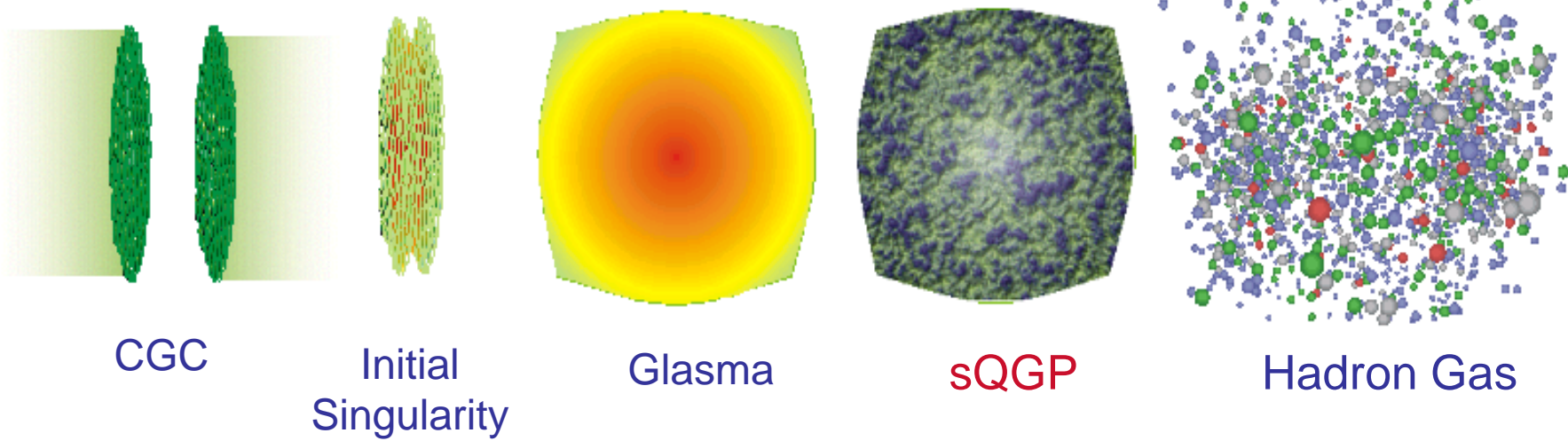


Theoretical Aspects of the Color Glass Condensate and Glasma

Art due to S. Bass



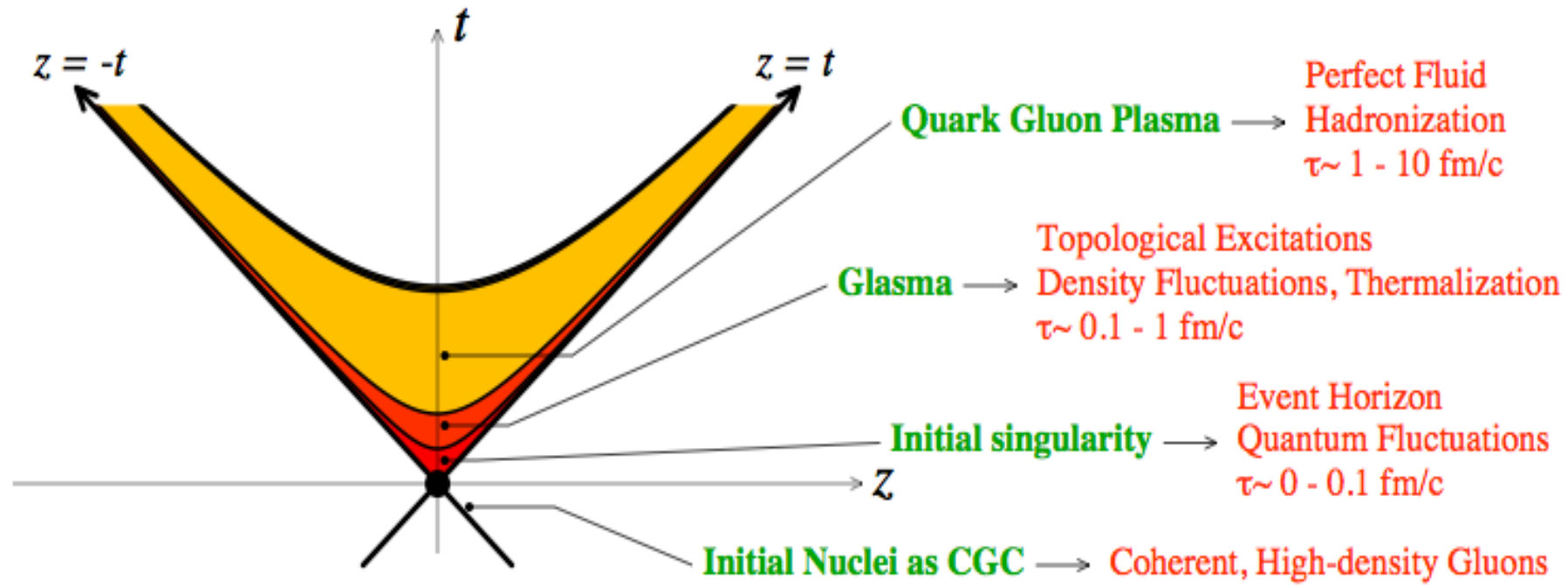
CGC

Initial Singularity

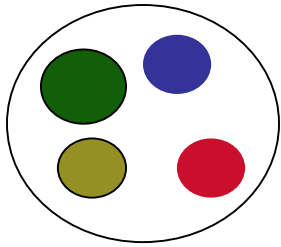
Glasma

sQGP

Hadron Gas



The Hadron Wavefunction at High Energy



Baryon:

3 quarks

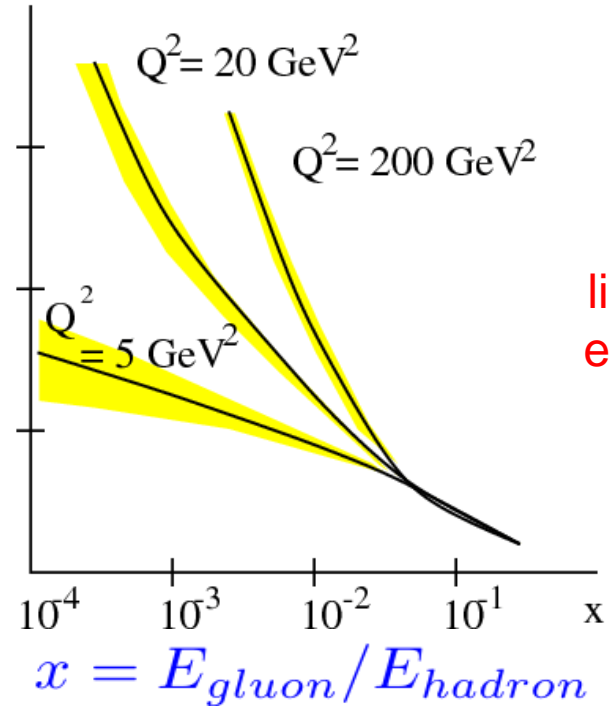
3 quarks 1
gluon

.....

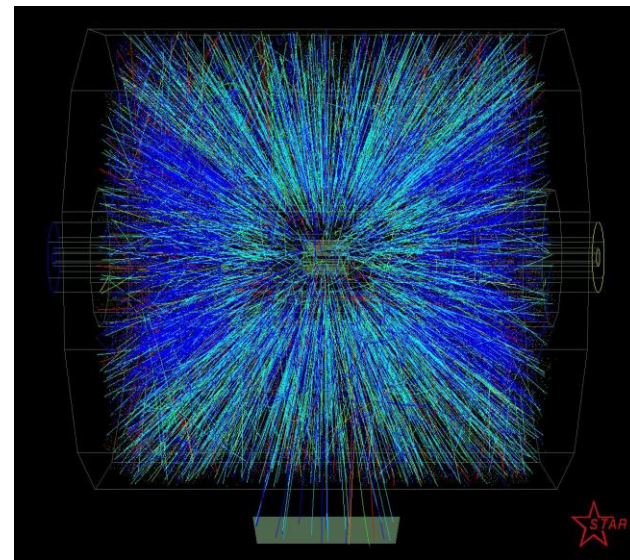
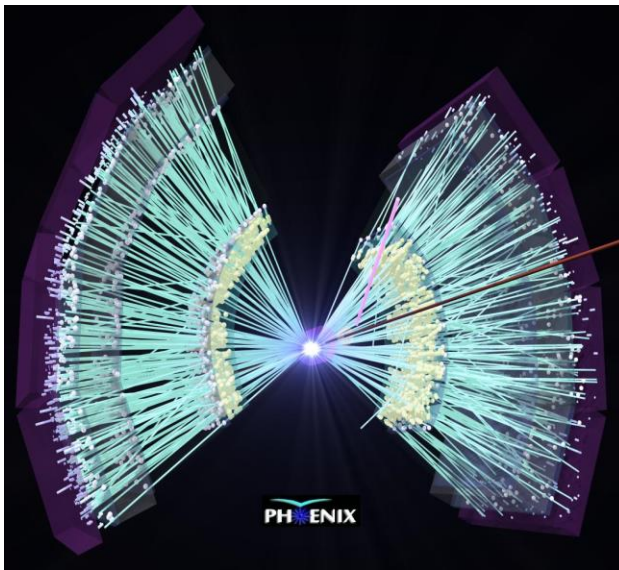
3 quarks and
lots of gluons



$xG(x, Q^2)$

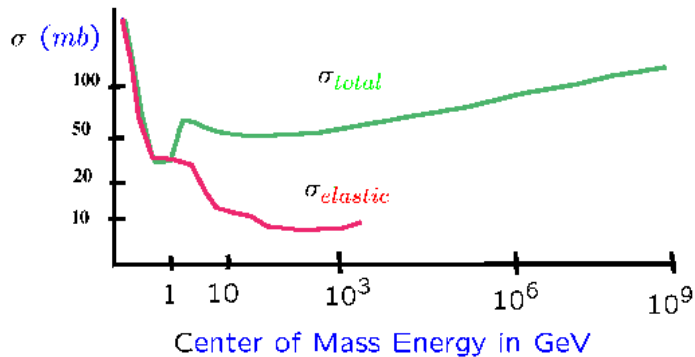


Small x
limit is high
energy limit



Where do all the gluons go?

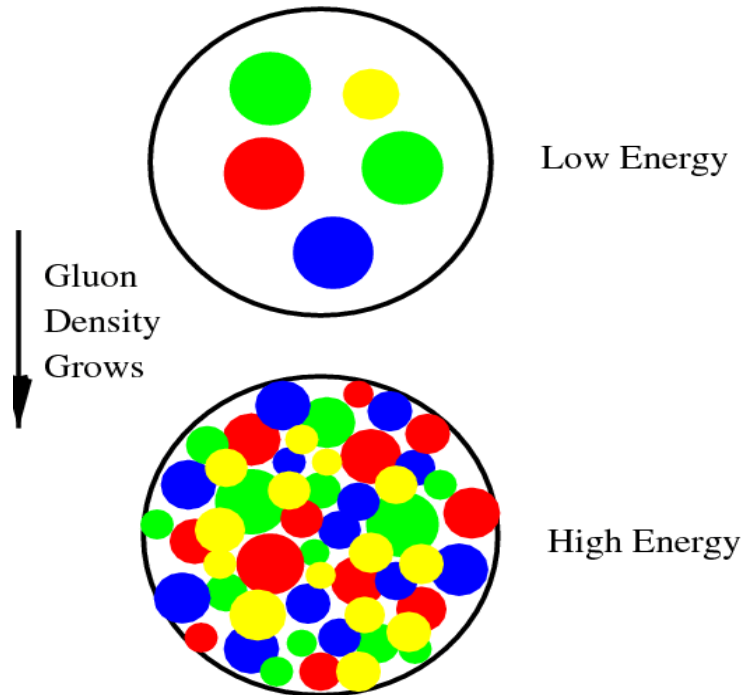
The total hadronic cross section:



Cross sections for hadrons rise very slowly with energy

$$\sigma_{tot} \sim \ln^2(E/\Lambda_{QCD})$$

$$\Lambda_{QCD} \sim 200 \text{ MeV}$$



But the gluon density rises much more rapidly!

The high energy limit is the high gluon density limit.

Surely the density must saturate for fixed sizes of gluons at high energy.

What is the Color Glass Condensate?

Glue at large x generates glue at small x

Glue at small x is classical field

Time dilation \rightarrow Classical field is glassy

High phase space density \rightarrow Condensate

Phase space density: $\frac{dN}{dyd^2p_Td^2x_T} = \rho$ $y = \ln(1/x)$

Attractive potential $V \sim -\rho$ Repulsive interactions $\sim \alpha_{strong}\rho^2$

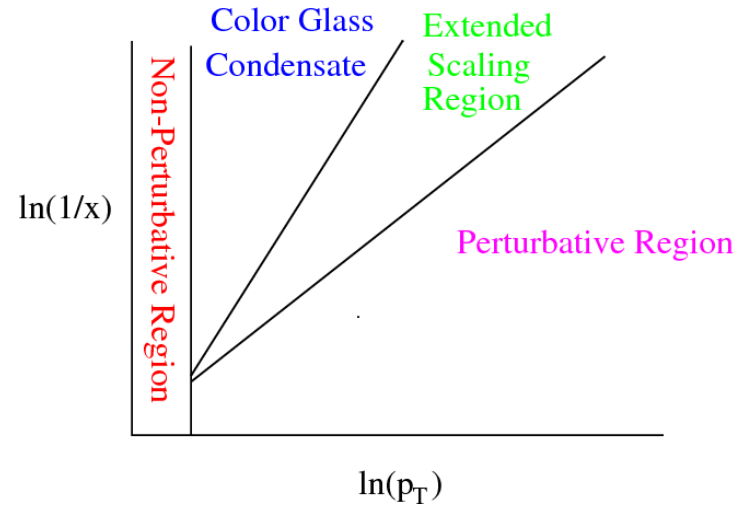
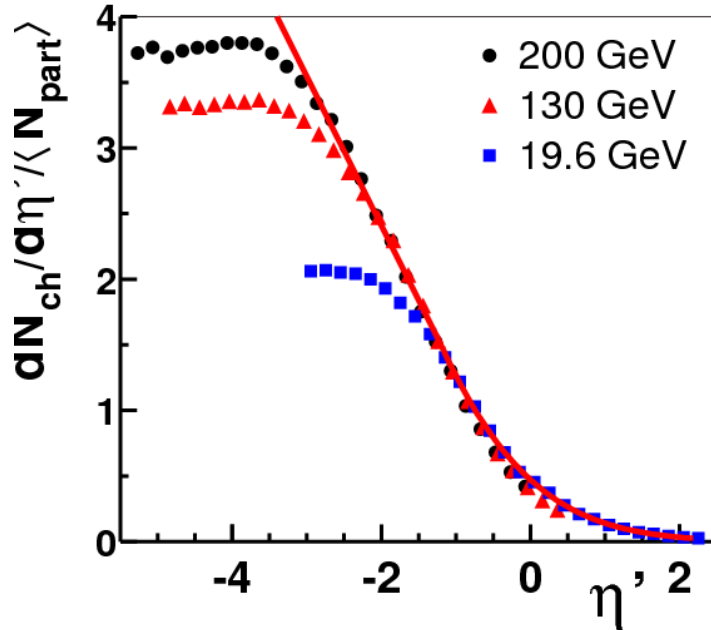
Density as high as it can be $\rho \sim 1/\alpha_{strong}$

Because the density is high α_{strong} is small

ρ is big

There must be a renormalization group

The x which separates high x sources from small x fields is arbitrary



Phobos multiplicity data

High energy QCD “phase” diagram

$$\frac{dN}{dyd^2r_T} \sim \int d^2p_T \frac{dN}{dyd^2p_T d^2r_T} \sim \frac{1}{\alpha_{strong}} Q_{sat}^2$$

The CGC Path Integral:

$$Z = \int_{\Lambda} [dA][d\rho] \exp\{iS[A, \rho] - F[\rho]\}$$

The current source:

$$J^{\mu} = \delta^{\mu+} \rho(x_T, y)$$

Rapidity:

$$y = \ln(x_0^- / x^-) \sim \ln(1/x) \sim \frac{1}{2} \ln(p^+ / p^-)$$

The separation scale is in rapidity or
longitudinal momentum

Λ

The Renormalization Group Equation:

$$Z_0 = e^{-F[\rho]}$$

$$\frac{d}{dy} Z_0 = -H[d/d\rho, \rho] Z_0$$

For strong and intermediate strength fields: H is second order in

$$d/d\rho$$

It has no potential, and a non-linear kinetic energy term

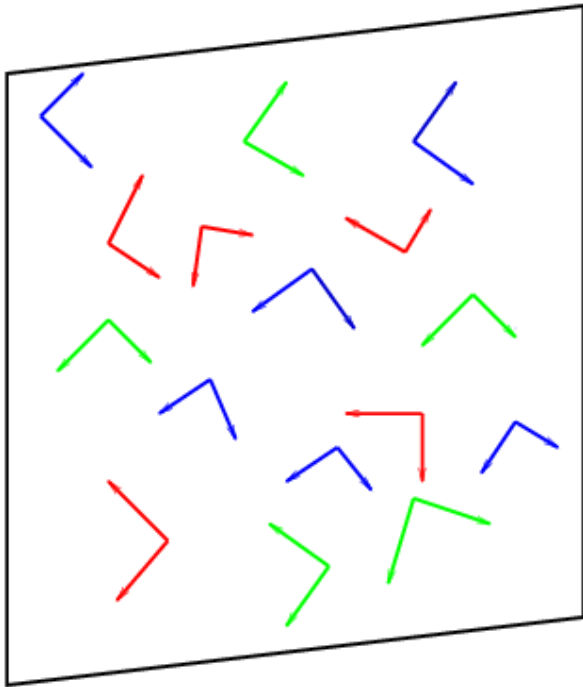
Like diffusion

$$d/dt \psi = -p^2/2 \psi$$

$$\psi \sim e^{-x^2/2t}$$

Wavefunction spreads for all time, and has universal limit:
Universality at high energy

What does a sheet of Colored Glass look like?



$$\vec{E} \perp \vec{B} \perp \vec{z}$$

On the sheet $x^- = t - z$ is small

Independent of $x^+ = t + z$

$$F^{i-} = E - B \quad \text{small}$$

$$F^{i+} = E + B \quad \text{big}$$

$$F^{ij}$$

Lienard-Wiechart potentials

Random Color

Density of gluons per unit area $\frac{1}{\pi R^2} \frac{dN}{dy} \sim \frac{1}{\alpha_{strong}} Q_{sat}^2$

Fields in longitudinal space:

$$F^{i+}$$

is a delta function on scales less than the
inverse longitudinal cutoff

The diagram shows two intersecting lines representing the longitudinal axes $x^- = 0$ and $x^+ = 0$. The $x^- = 0$ line is labeled at its top-left end. The $x^+ = 0$ line is labeled at its top-right end. Two gauge field expressions are written along these lines:

- Along the $x^+ = 0$ line: $A^j = \frac{1}{i} U_2 \nabla^j U_2^\dagger$
- Along the $x^- = 0$ line: $A^j = \frac{1}{i} U_1 \nabla^j U_1^\dagger$

Gluon distribution is at
scales larger than the
cutoff

$$G(k) \sim 1/p^+$$

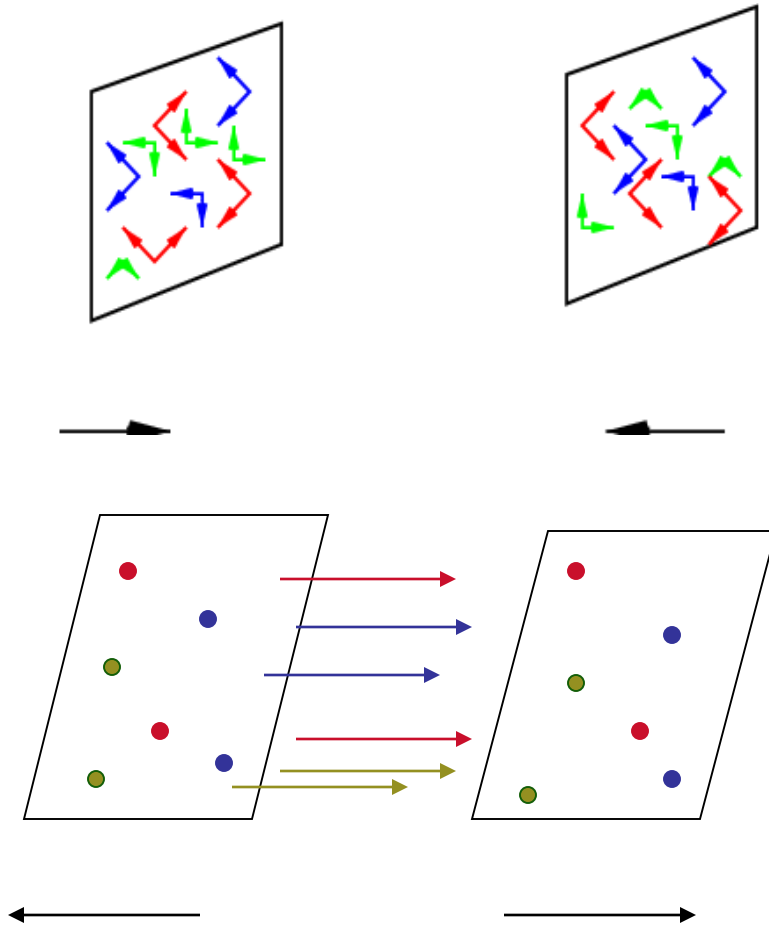
$$G(k) = \langle a^\dagger(k) a(k) \rangle \sim \langle A(k) A(-k) \rangle$$

CGC Gives Initial Conditions for QGP in Heavy Ion Collisions

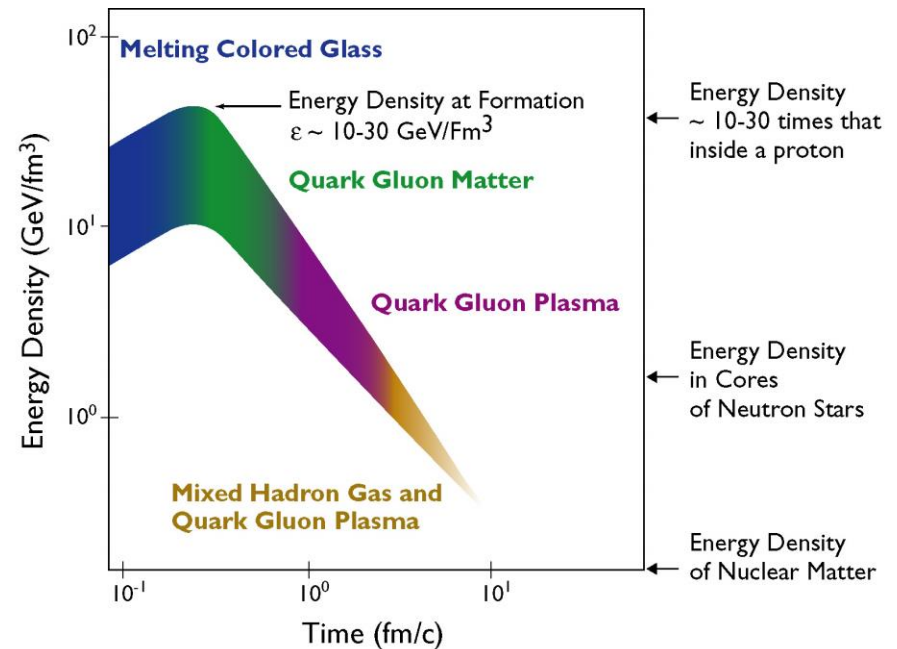
Two sheets of colored glass collide

Glass melts into gluons and thermalize

QGP is made which expands into a mixed phase of QGP and hadrons



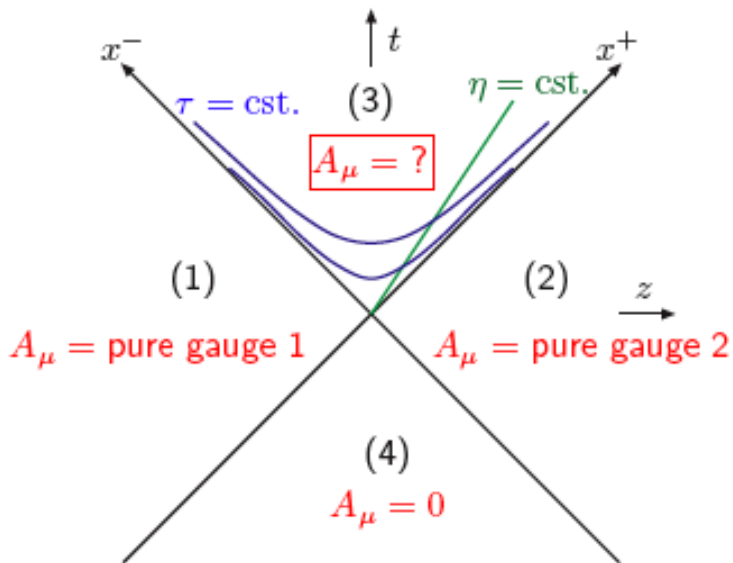
“Instantaneously” develop longitudinal color E and B fields



Before the collision only transverse E and B
CGC fields

Color electric and magnetic monopoles

Almost instantaneous phase change
to longitudinal E and B



In forward light cone

$$A_1^i + A_2^i$$

generates correct sources on
the light cone

$$\nabla \cdot E = A \cdot E$$

$$\nabla \cdot B = A \cdot B$$

$$A_1 \cdot E_2$$

$$A_1 \cdot B_2$$

Equal strength for magnetic and
electric charge on average!

$$\partial^\mu J_\mu^5 = \kappa E \cdot B + O(m_{quark})$$

Different signs

Generate different chiralities
and vorticity in the fluid.

Violates P and CP on an event
by event basis

Integral vanishes initially

Topological charge density is
maximal:

Anomalous mass generation

In cosmology:

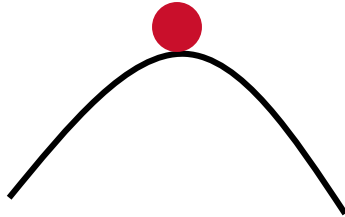
Anomalous Baryogenesis

Classical equation do
not generate net
topological charge.

Instabilities in these
solutions will generate
such charge, and can
thermalize the system



Before collision, stability



After collisions, unstable

Interactions of evaporated gluons
with classical field is $g \times 1/g \sim 1$ is
strong

Thermalization?

$$W[p, X] = \int dz e^{-ipz} \psi^*(X + z/2) \psi(X - z/2)$$

Quantum fluctuations can become as
big as the classical field

Quantum fluctuations analogous to
Hawking Radiation

Growth of instability generates
turbulence \Rightarrow Kolmogorov spectrum

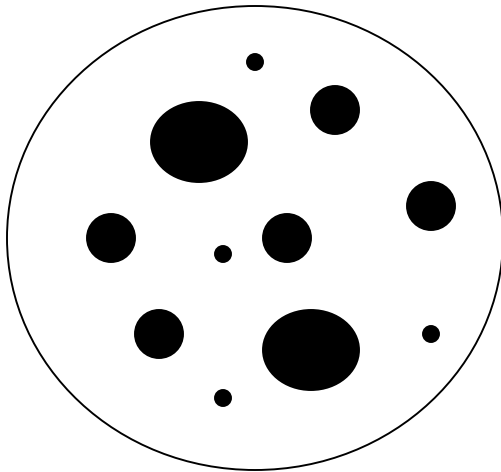
Analogous to Zeldovich spectrum of
density fluctuations in cosmology

Topological mass generation

Wigner distribution of initial
wavefunction gives seeds of
fluctuations.

These seeds grow when inserted
into classical equations

Fluctuation in saturated regions and 2-d Quantum Gravity?



How to compute
scattering?

What is the distribution of the
sizes of saturated spots of
gluonic matter with size less
than the saturation size
scale?

Such fluctuations can in
principle dominate the
scattering out to very large
momentum scales:
breakdown of factorization of
hadron-hadron relative to
lepton hadron

2-d field theory: Dimensionless scalar theory

$$\phi(x_T) = \ln(Q_{sat}^2(x_T)/Q_{ave}^2)$$

$$\nabla^2 \sim Q_{sat}^2(x_T) \sim M^2 e^{\phi(x_T)}$$

$$S = \frac{1}{g^2} \int d^2 x_T ((\nabla\phi)^2 + M^2 e^{\phi})$$

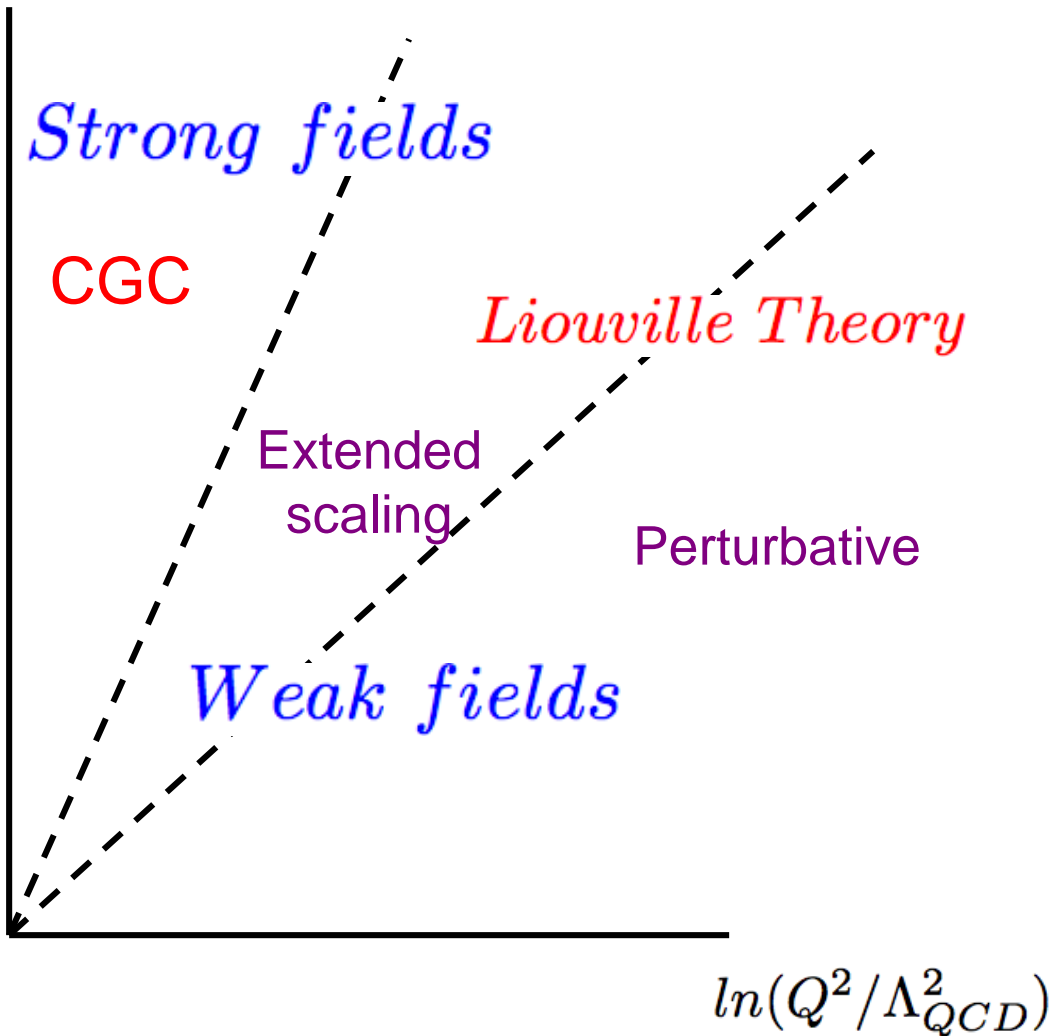
Liouville theory: 2-d conformal invariance:

$$Q_{ave}^2 = M^2 \langle e^{\phi} \rangle = 0$$

Insert a source term: $-M^2\phi$
 $\langle e^{\phi} \rangle = 1$

Theory appears to be UV finite (unlike Liouville which is renormalizable) and conformal on distance scales less than the saturation scale!

$\ln(1/x)$



Successes:

Geometric scaling in DIS

Diffractive DIS

Shadowing in dA

Multiplicity in AA

Limiting fragmentation

Long range correlations

Total cross section

Pomeron, reggeon, odderon