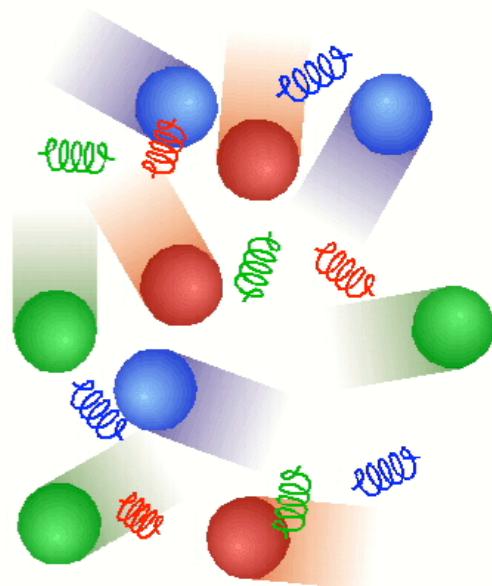


# What I have learned at RHIC

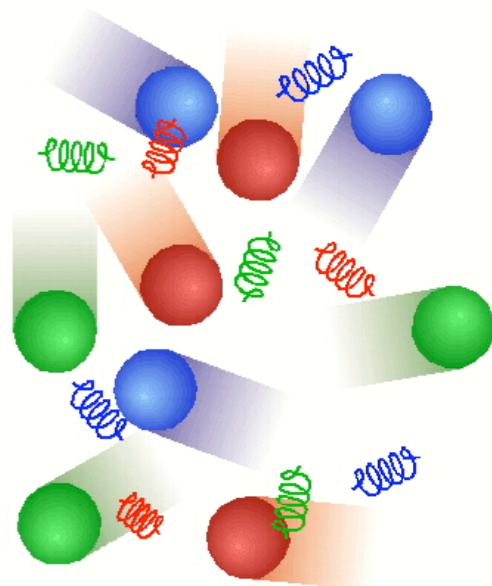
**Yasuo MIAKE**  
**Univ. of Tsukuba**



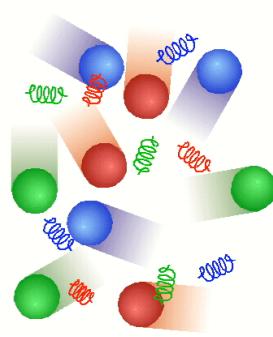
I think

# What I have learned at RHIC

**Yasuo MIAKE**  
**Univ. of Tsukuba**



# Text book for graduate students



CAMBRIDGE | Catalogue

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## Quark-Gluon Plasma

Series: [Cambridge Monographs on Particle Physics, Nuclear Physics and Cosmology](#)

Kohsuke Yagi

*Urawa University, Japan*

Tetsuo Hatsuda

*University of Tokyo*

Yasuo Miake

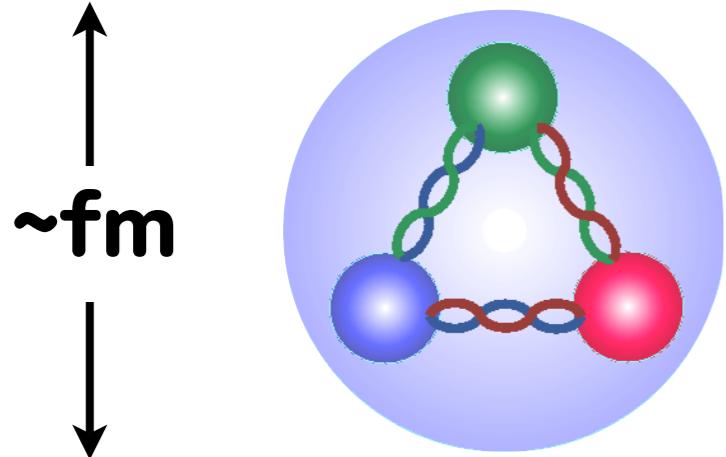
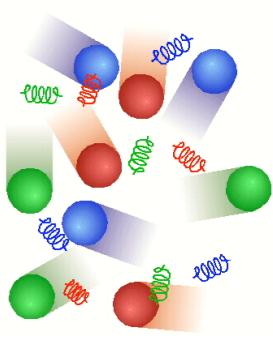
*University of Tsukuba, Japan*

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**Hardback** (ISBN-10: 0521561086 | ISBN-13: 9780521561088)

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# Quark Gluon Plasma

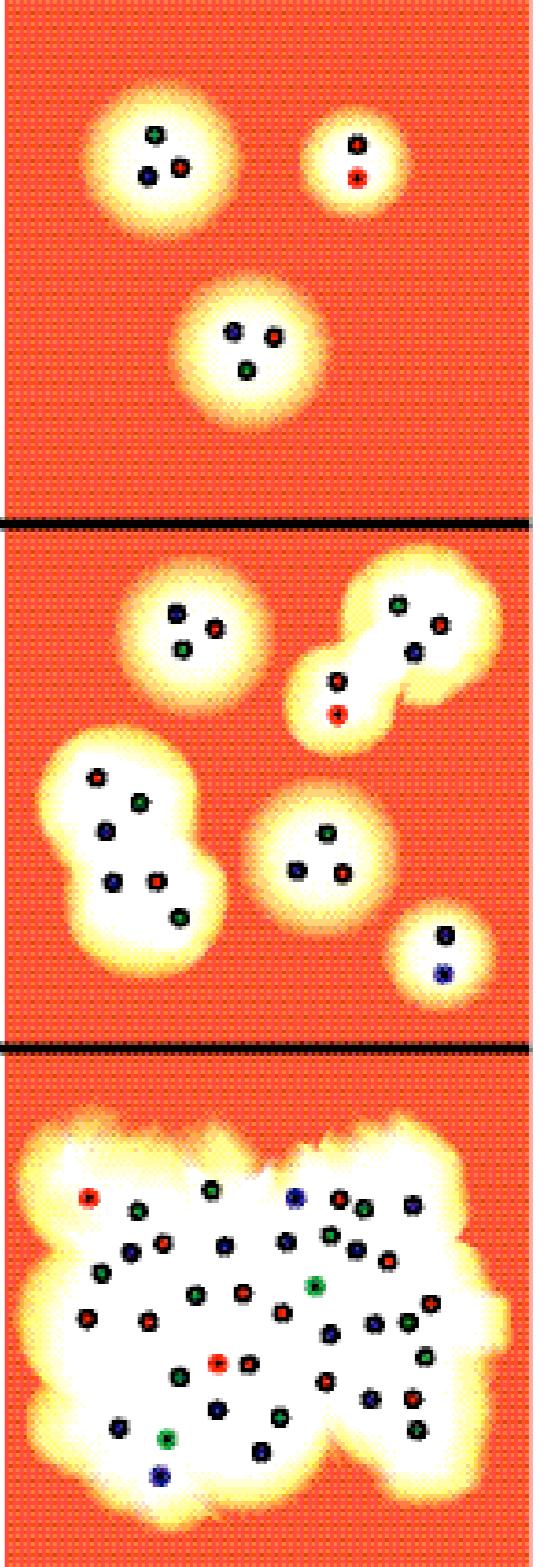


- ✓ **Size of hadrons  $\sim$  fm**
- ◆ **At high  $\rho$  or high  $T$ ,  
hadron gas to quark-gluon plasma.**
- **Phase transition**

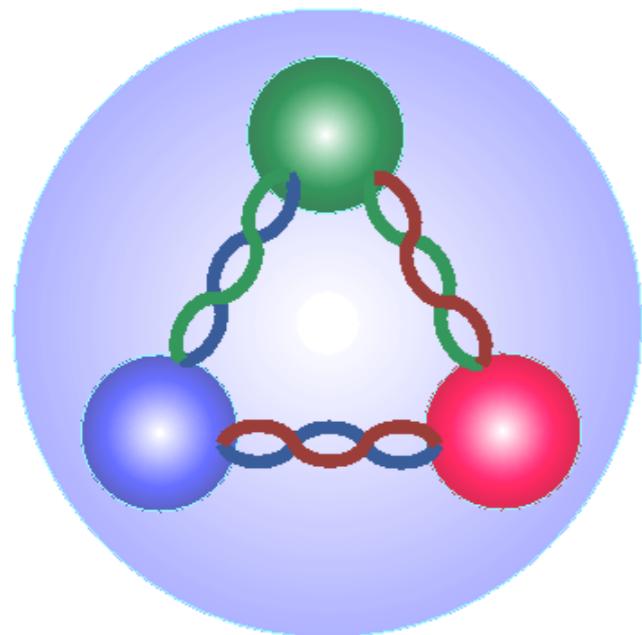
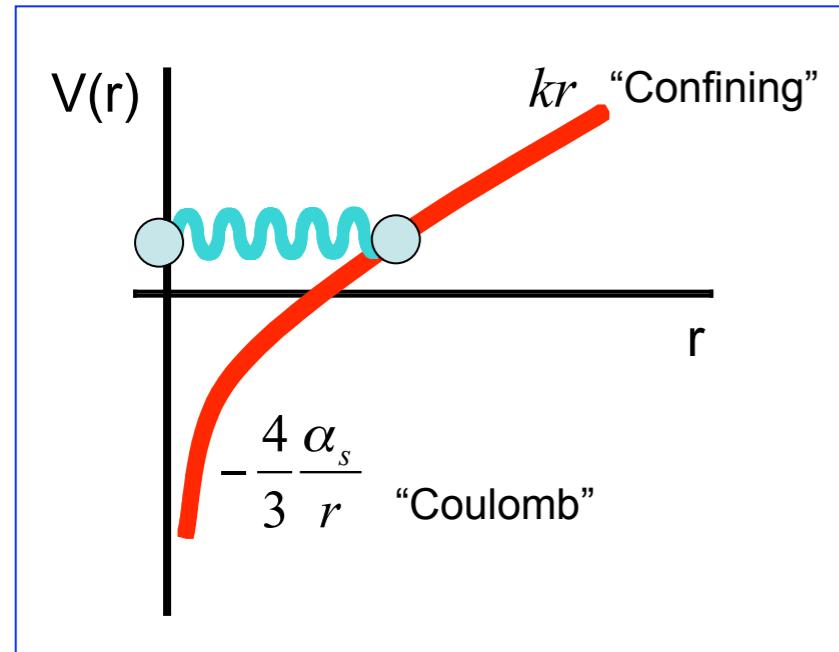
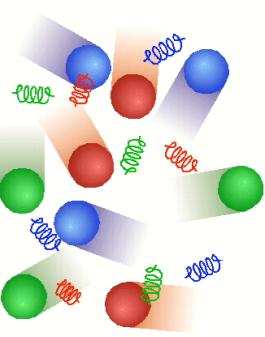
hadron gas  
 $T, \rho$  low

phase transition  
 $T, \rho$  critical

quark-gluon-plasma  
 $T, \rho$  high

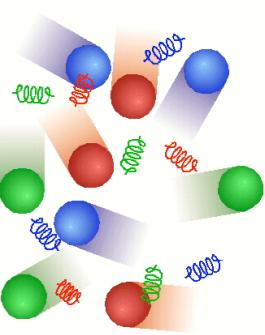


# Is it QGP inside a proton?



- ✓ Asymptotic freedom means a QGP inside a proton?  
→ No!
- ✓ Free moving quarks & gluons for a large volume
- ✓ Applicability of **Statistical Physics** is essential!

# What we expect: Statistical Nature



$$\epsilon_{\text{QGP}} \sim 2 \text{ [GeV/fm}^3\text{]} \xleftarrow{\hspace{1cm}} \text{Ex. Lattice QCD}$$

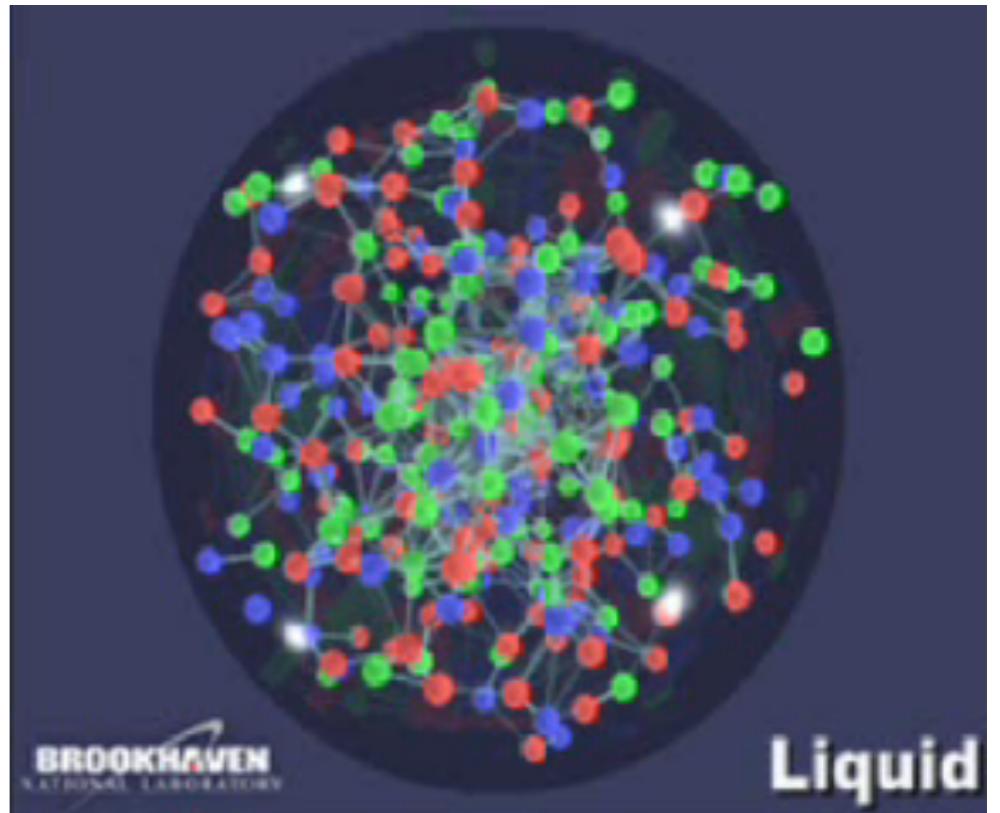
$$\langle n_{q,\bar{q}} \rangle \sim \frac{\epsilon_{\text{QGP}}}{\langle m_T \rangle} \sim \frac{2\text{GeV}}{0.4\text{GeV}} \sim 5$$

$$\lambda_q = \frac{1}{n\sigma_{qq}}$$

$$\sim \frac{1}{5 \times 0.4} = 0.5 \text{ [fm]}$$

$$\lambda_q \ll R_{\text{system}}$$

$$\therefore \sigma_{qq} \sim \frac{\sigma_{NN}}{3 \times 3} \sim \frac{4[\text{fm}^2]}{9} \sim 0.4$$

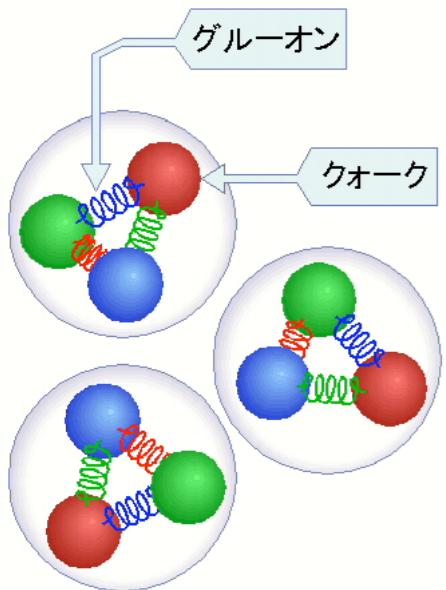


Animation by Jeffery Mitchell (Brookhaven National Laboratory). Simulation by the UrQMD Collaboration

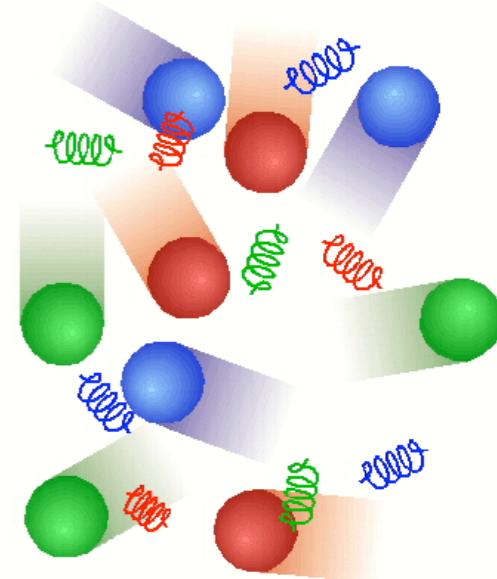
✓ What we expect:

◆ Statistical physics at quark level

◆ Hydrodynamical behavior at quark level

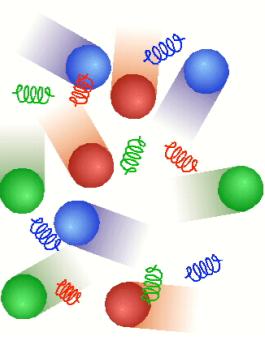


# Statistical nature observed

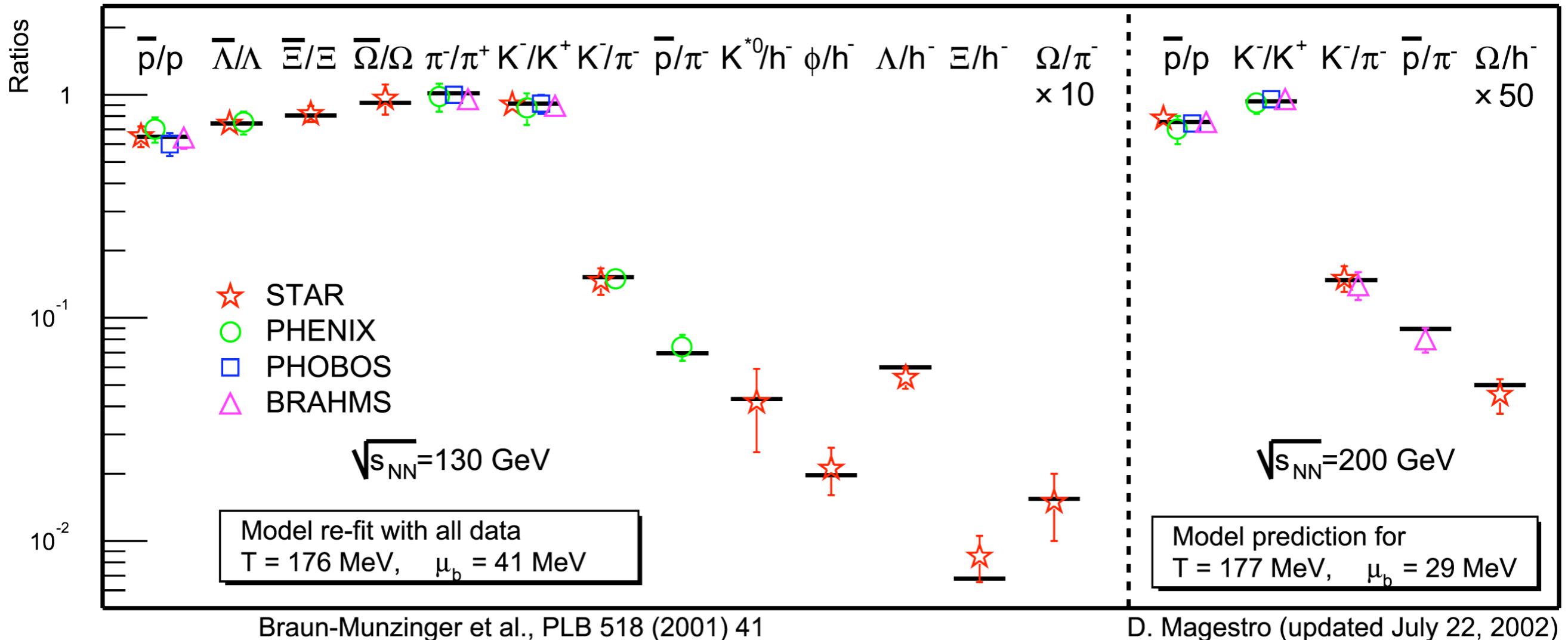


- ★ Chemical equilibrium from particle yield ratio
- ★ Kinematical equilibrium from transverse distr.

# Chemical Eq. from particle yield ratios



$$n_i = \frac{g_i}{2\pi^2} \int_0^\infty \frac{p^2 dp}{e^{(E_i - \mu_i)/T} \pm 1}$$

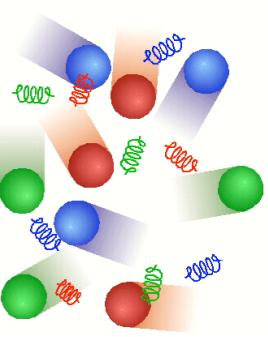


Braun-Munzinger et al., PLB 518 (2001) 41

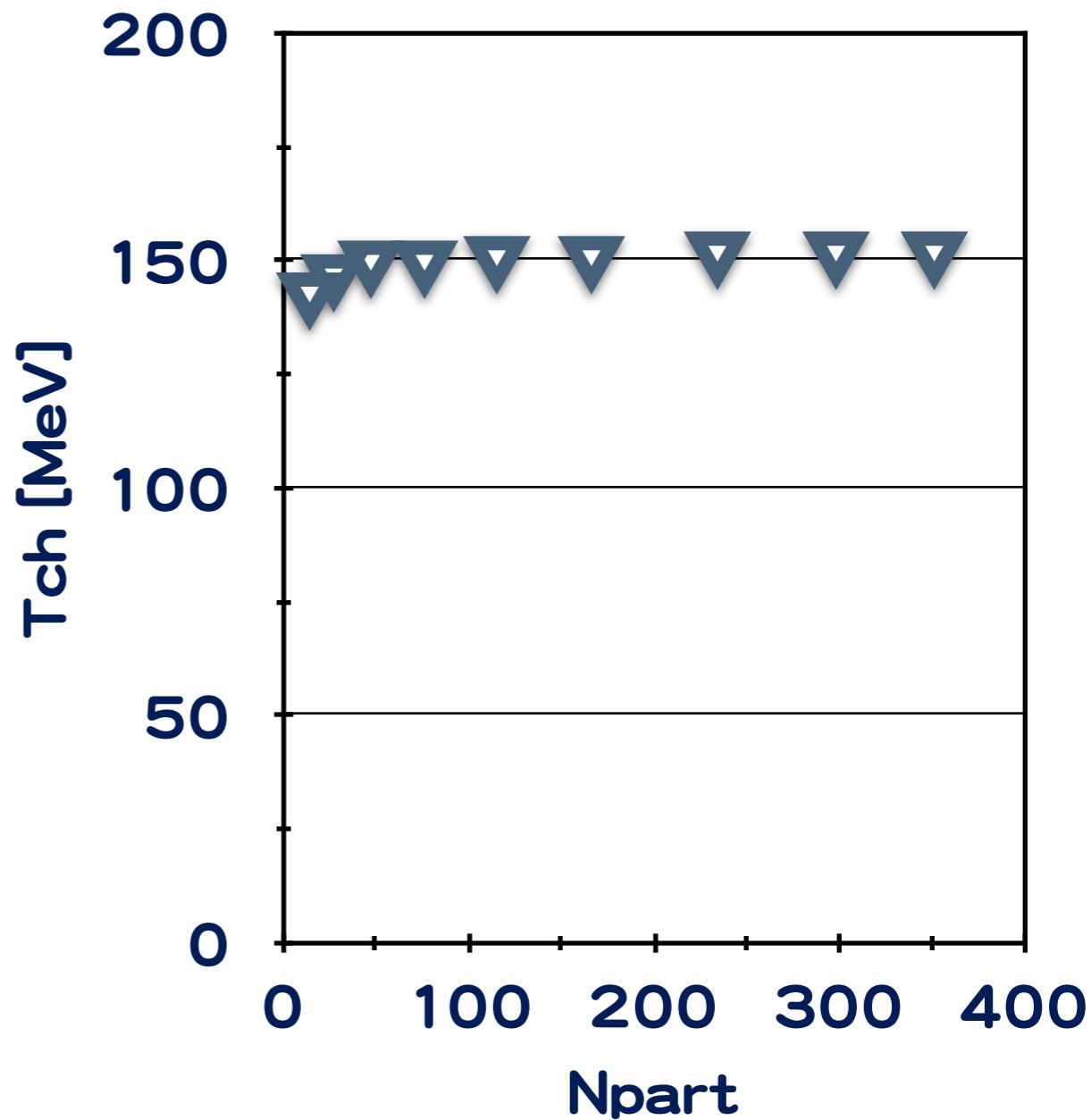
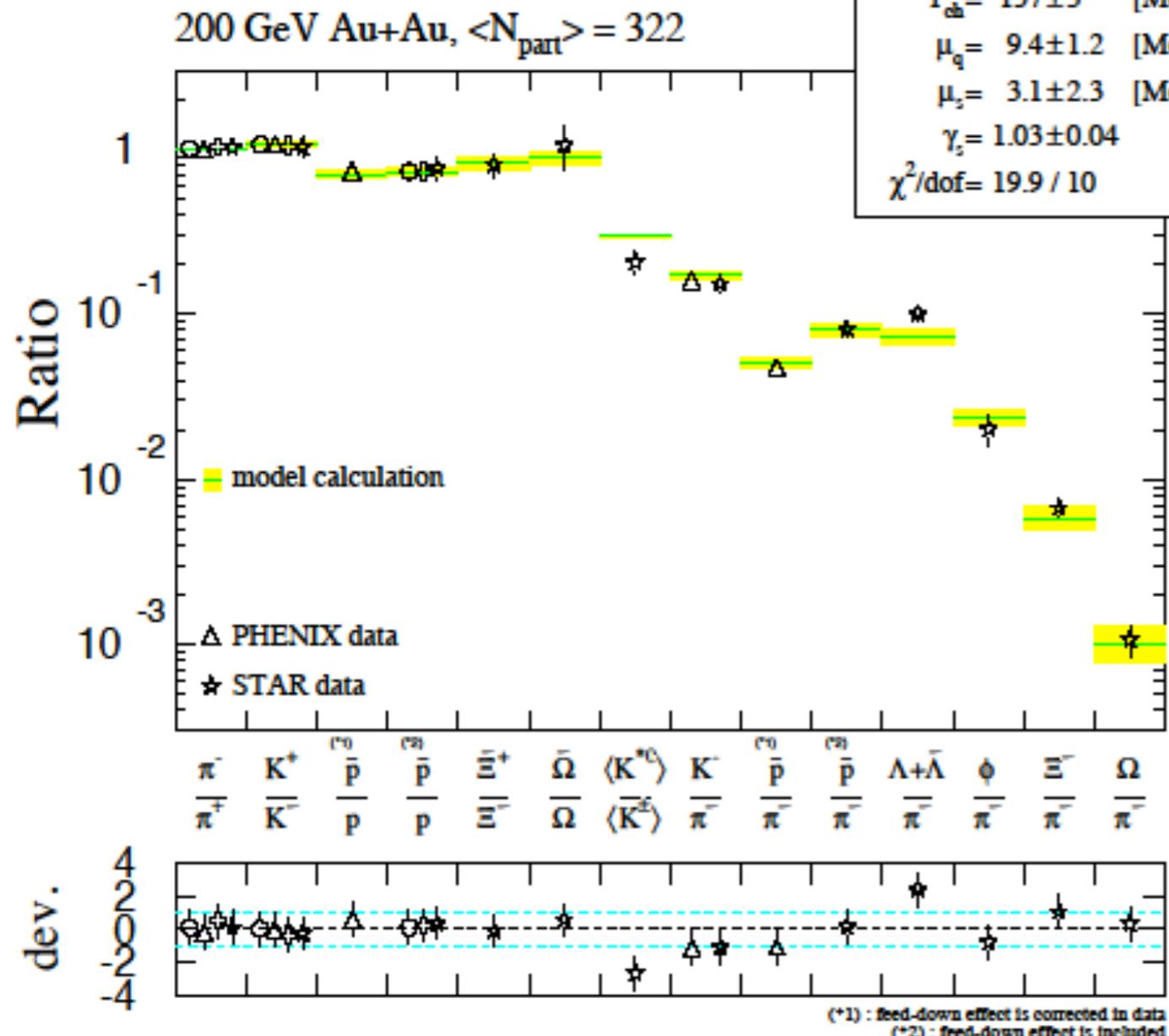
D. Magestro (updated July 22, 2002)

✓ Only two parameters fit every ratio very well !

# Centrality dep. of $T_{ch}$

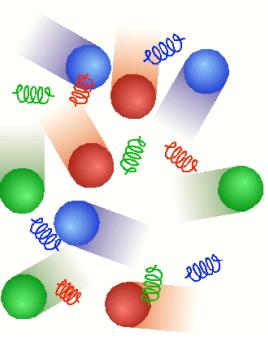


M.Kaneta, N.Xu, nucl-th/0405068



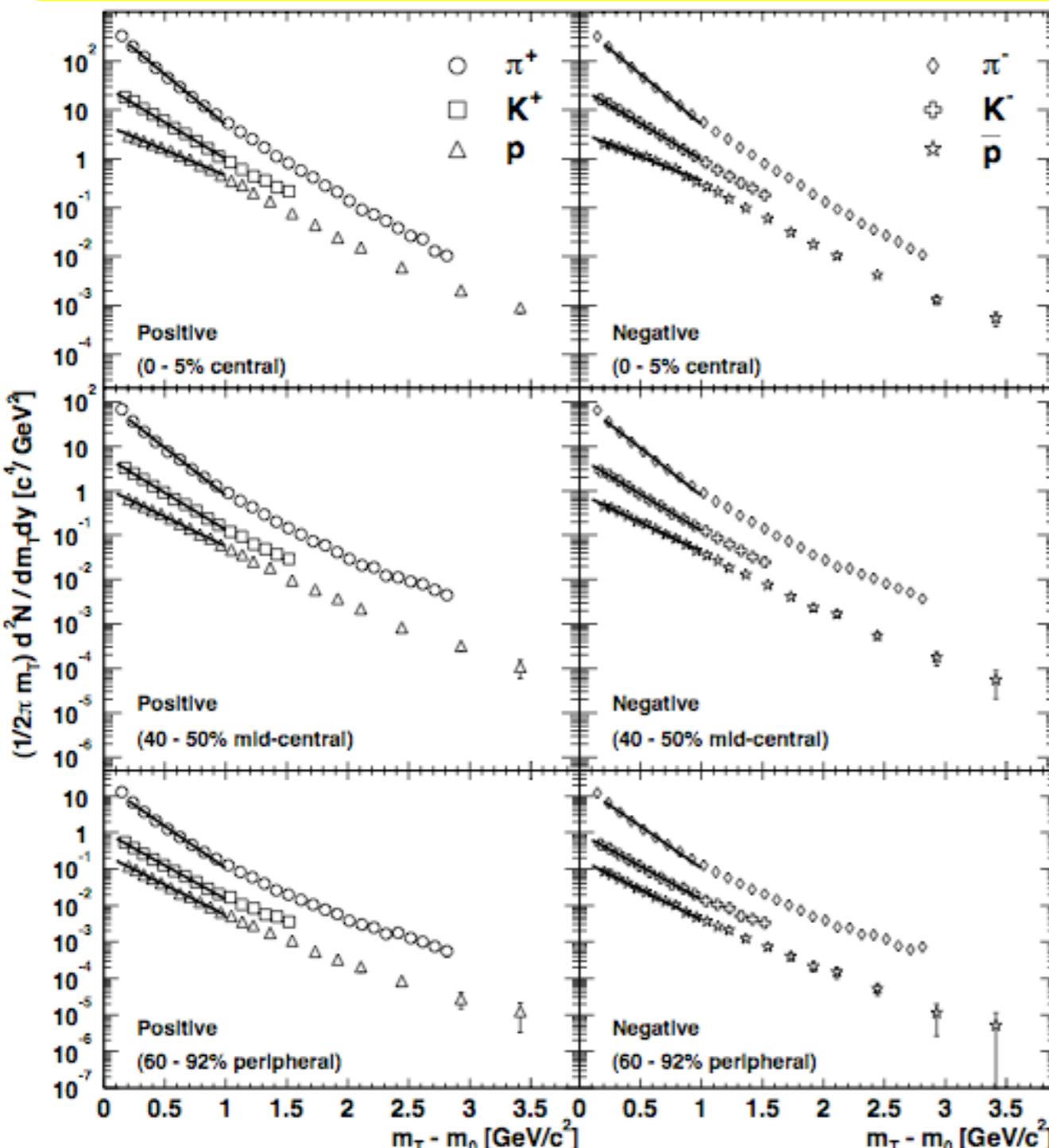
✓  $T_{ch}$  stays amazingly constant from peripheral to central collisions

# Transverse mass distr.

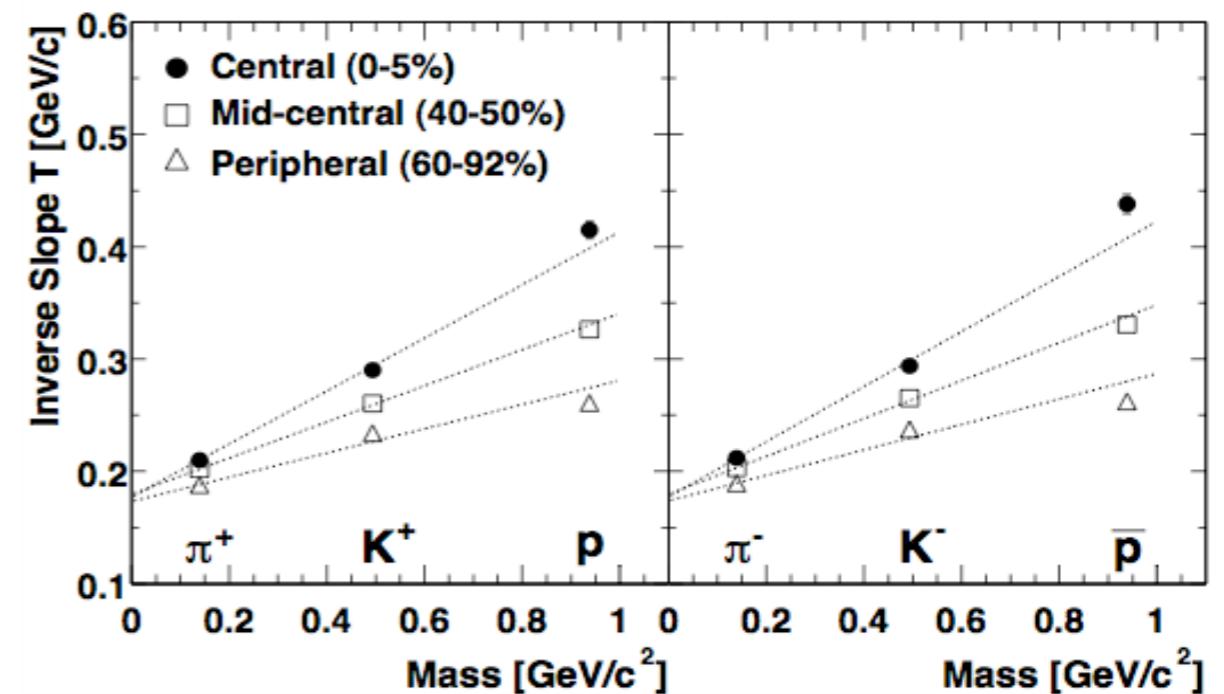


$$m_t = \sqrt{p_t^2 + m^2}$$

$$\frac{d^2N}{2\pi m_T dm_T dy} = \frac{1}{2\pi T(T+m_0)} A \exp\left(-\frac{m_T - m_0}{T}\right)$$



PHENIX, PRC69,034909(2004)

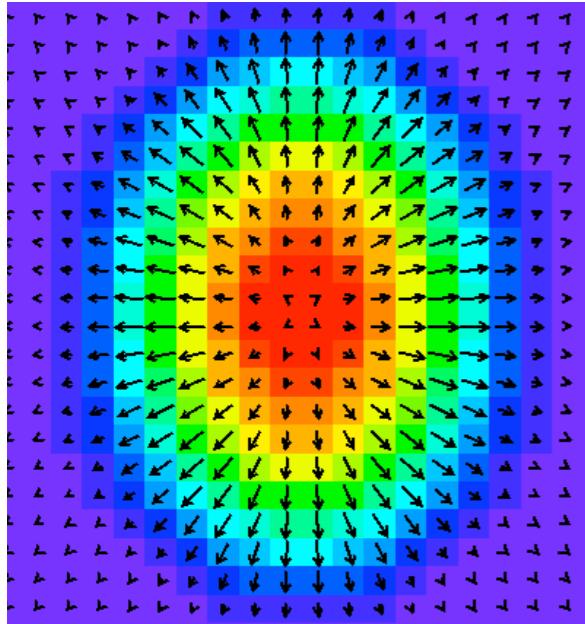
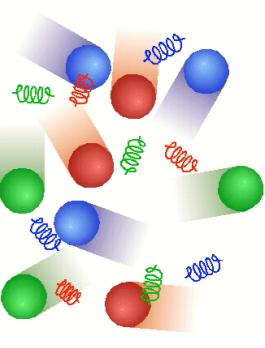


- ✓ Exponential in  $m_T$
- ◆ Known as  $m_T$  scaling
- ◆ Thermal distr.
- ✓ Flatter  $m_T$  distr for heavier particle mass
- ◆ Mass Ordering of Slope param.
- ◆ Effect of Collective Flow

$$T \approx T_0 + \frac{1}{2} m \langle v_r \rangle^2$$

Collective Flow

# Blast Wave Model



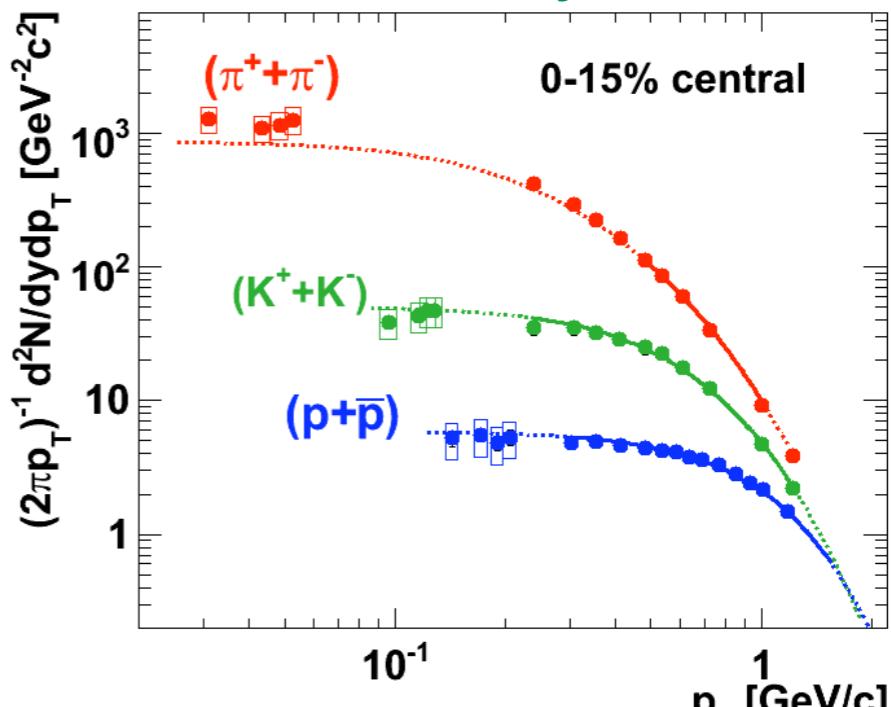
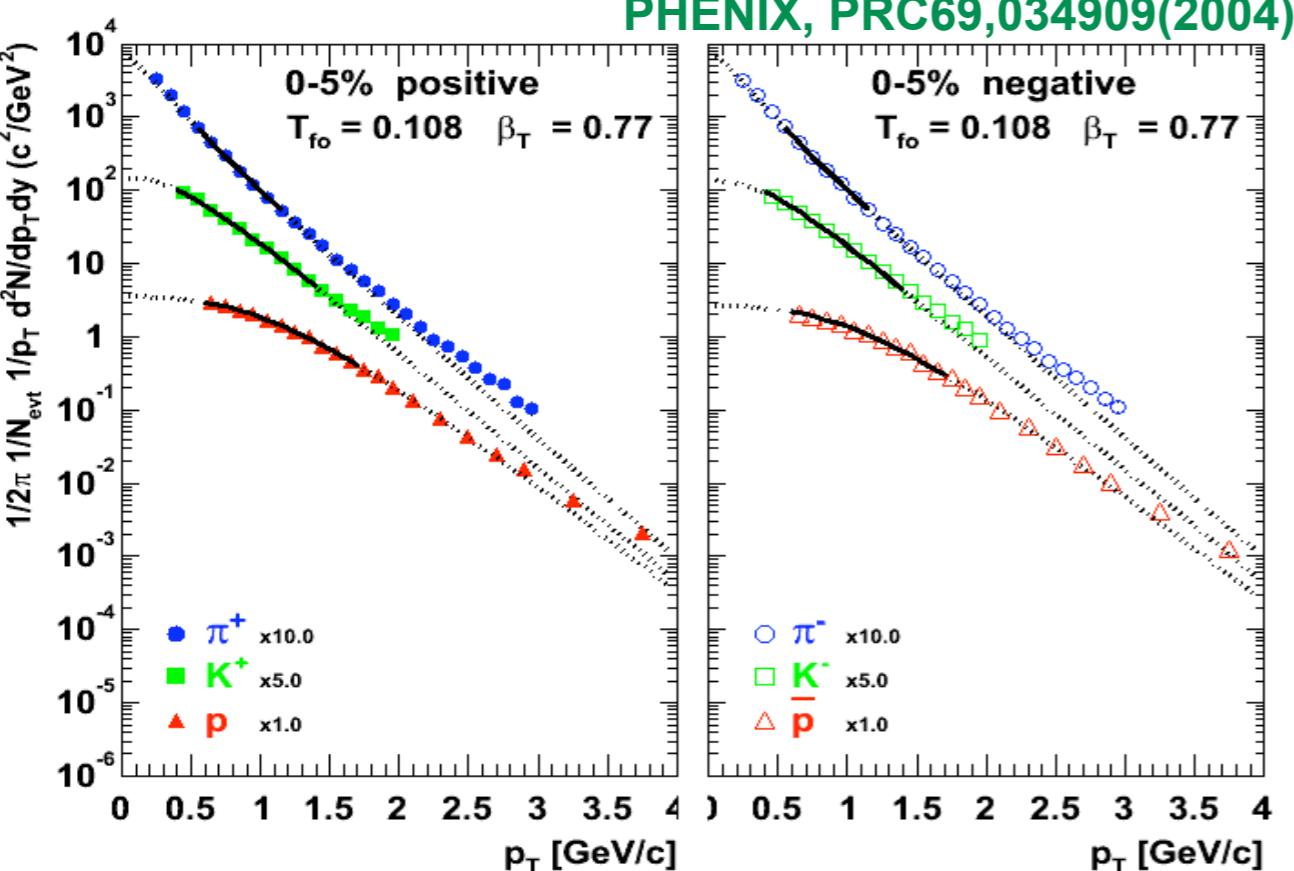
$$\frac{1}{m_T} \frac{dN}{dm_T} = A \int_0^R f(r) r dr m_T I_0\left(\frac{p_T \sinh \rho}{T_{fo}}\right) K_1\left(\frac{m_T \cosh \rho}{T_{fo}}\right)$$

**PRC48(1993)2462.**

$$\rho(r) = \tanh^{-1}(\beta_T) \cdot r/R$$

$I_0$ ,  $K_1$ : modified Bessel function

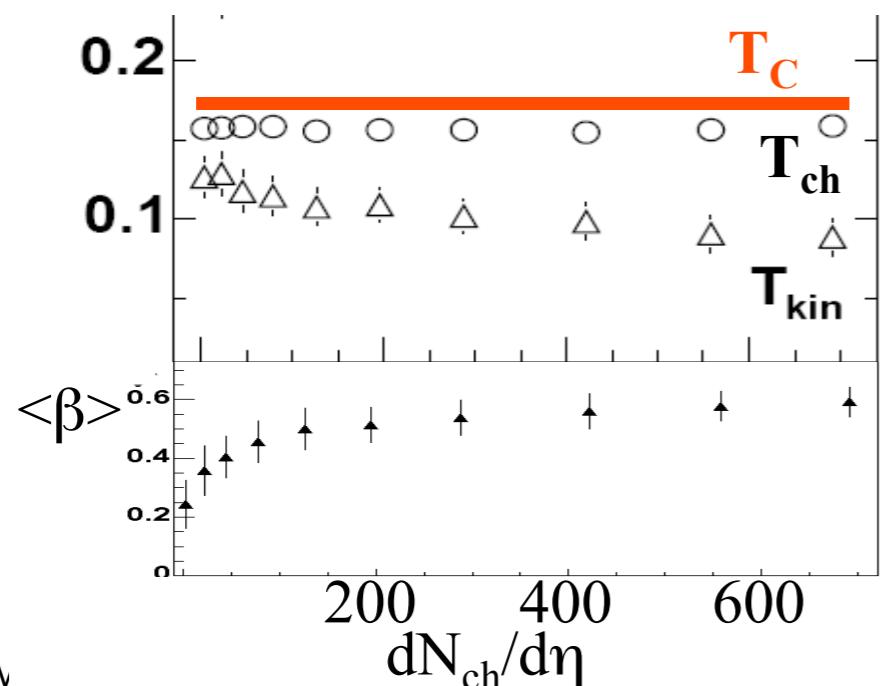
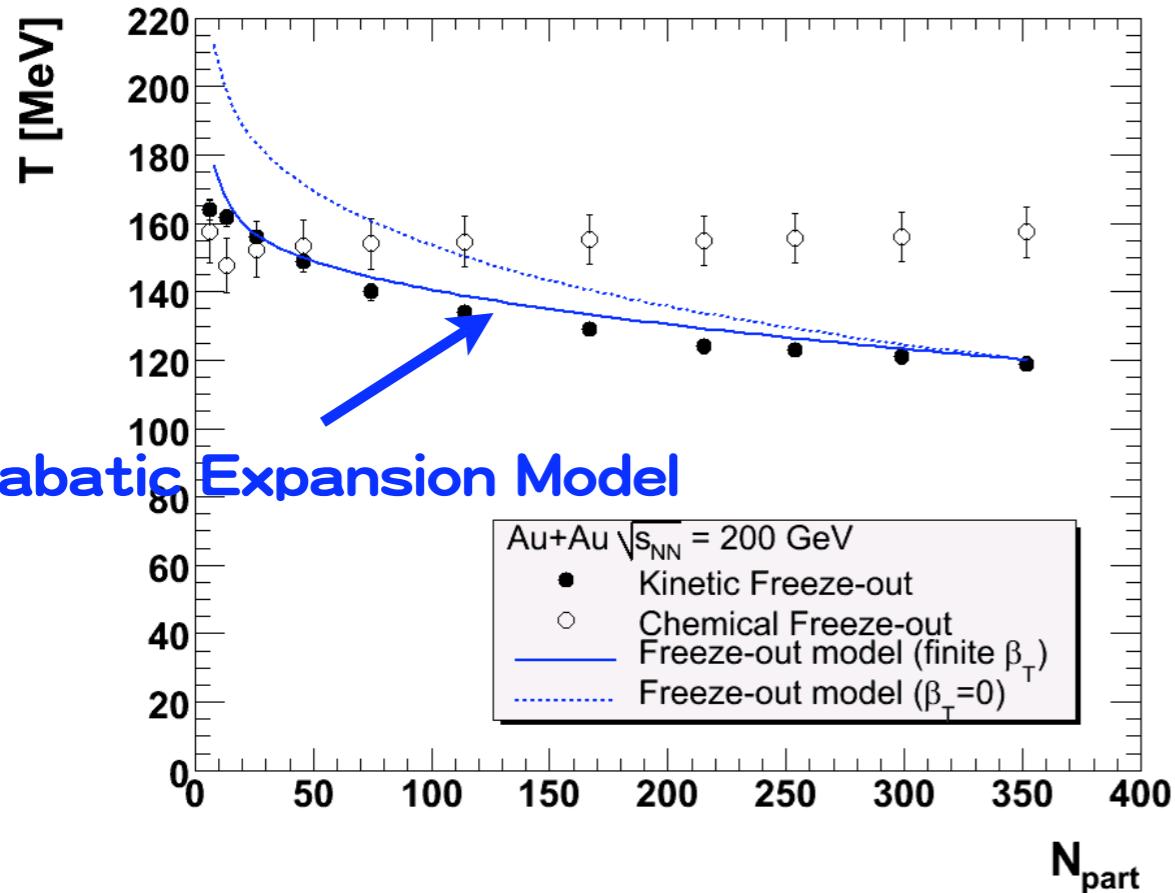
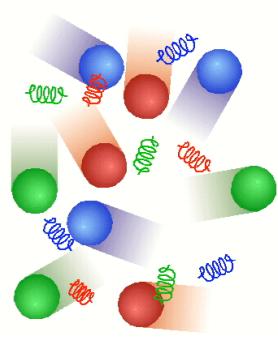
**Phobos, J.Phys.G34,S1103-7(2007)**



✓ Good tool to separate thermal and collective

✓ Well describe < 2 GeV/c

# Freeze-out Conditions



✓ Kinematical & Chemical freeze-out show difference in centrality dependence!

◆ Kinematical :

$$\rightarrow T_{ch}^{\text{cent.}} < T_{ch}^{\text{per.}} < T_{ch}$$

Freeze-out with  $\lambda \sim R$

◆ Chemical :

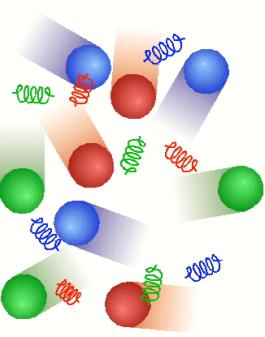
$$\rightarrow T_{ch}^{\text{cent.}} \sim T_{ch}^{\text{per.}} \sim 170 \text{ MeV}$$

Freeze-out with  $\varepsilon \sim \varepsilon_{\text{crit}}$

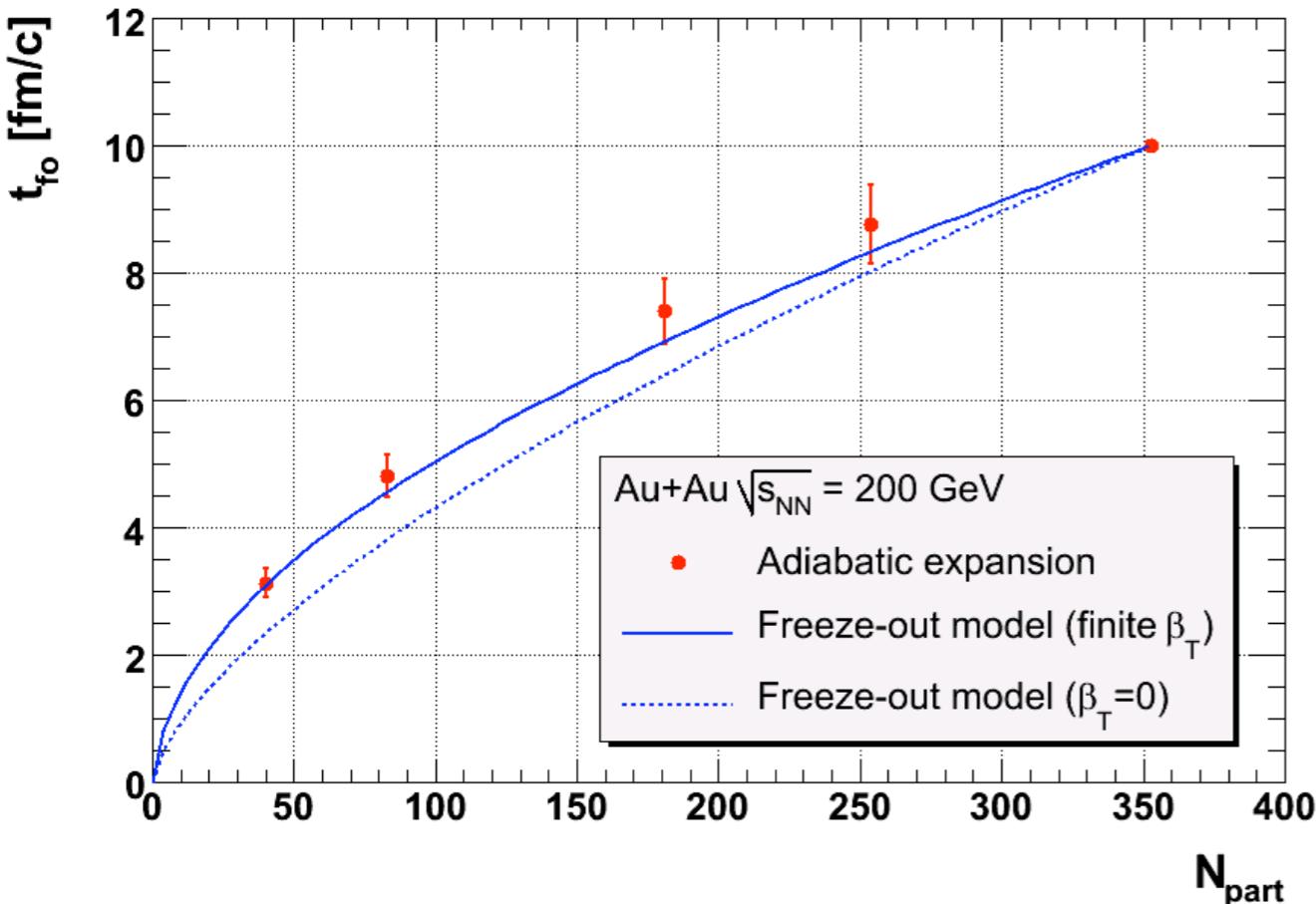
✓ Nature of Freeze-out

◆ Kinematical freeze-out is collisional, while chemical is not.

# Adiabatic Expansion Model



## ✓ Assumptions;

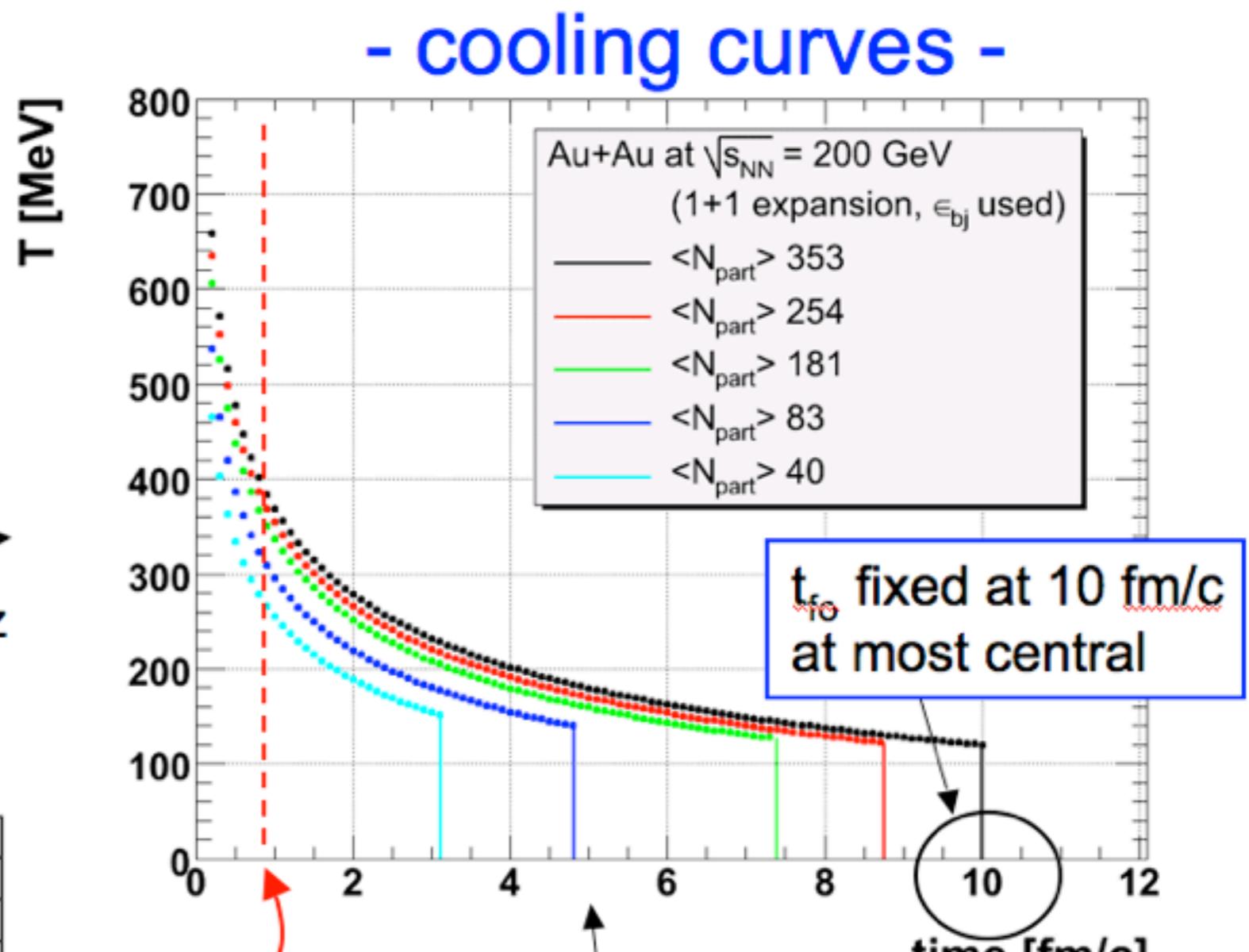
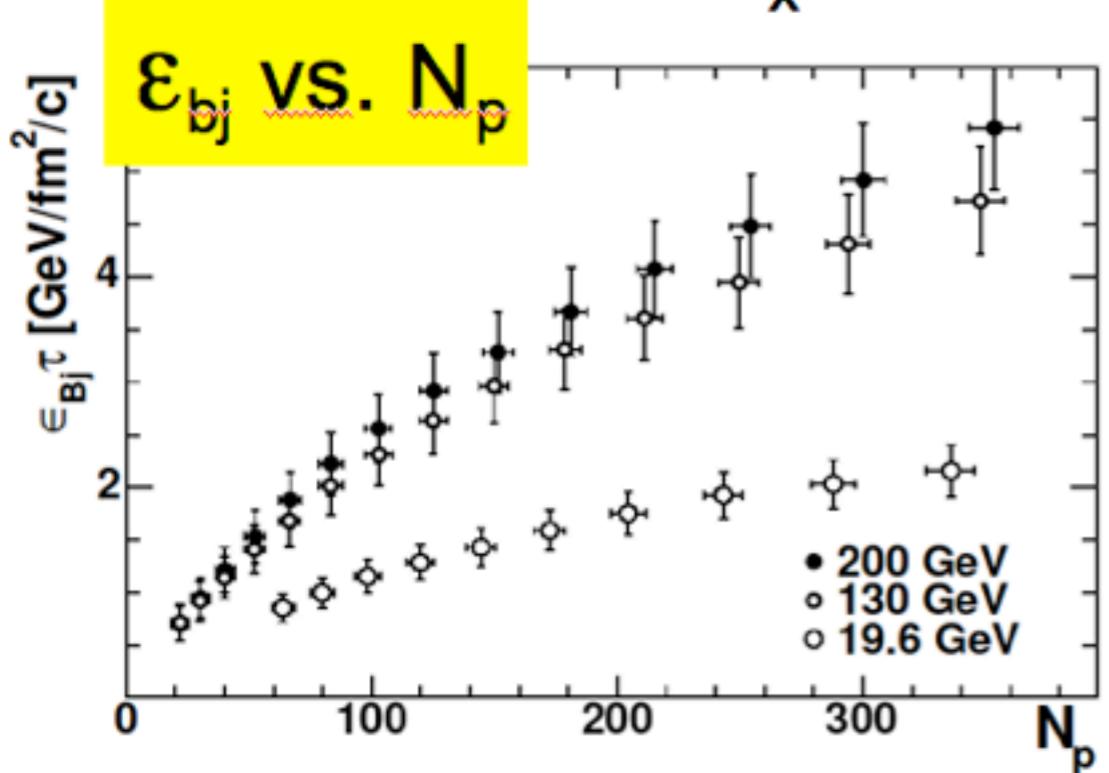
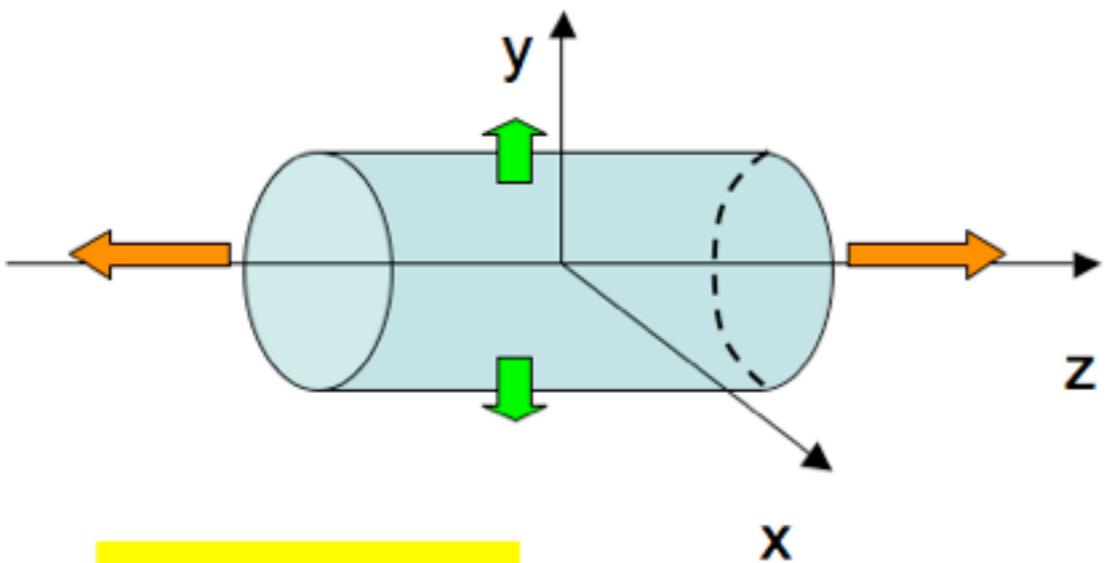


- ◆ Perfect fluid/Ideal gas
- ◆ Entropy conservation
- ◆ Bjorken Formula for  $\epsilon$
- ◆ Transverse & longitudinal expansion
- ◆  $\lambda = R$  as freezeout condition

## ✓ Central collisions freeze-out later than

# 1+1D Adiabatic Expansion

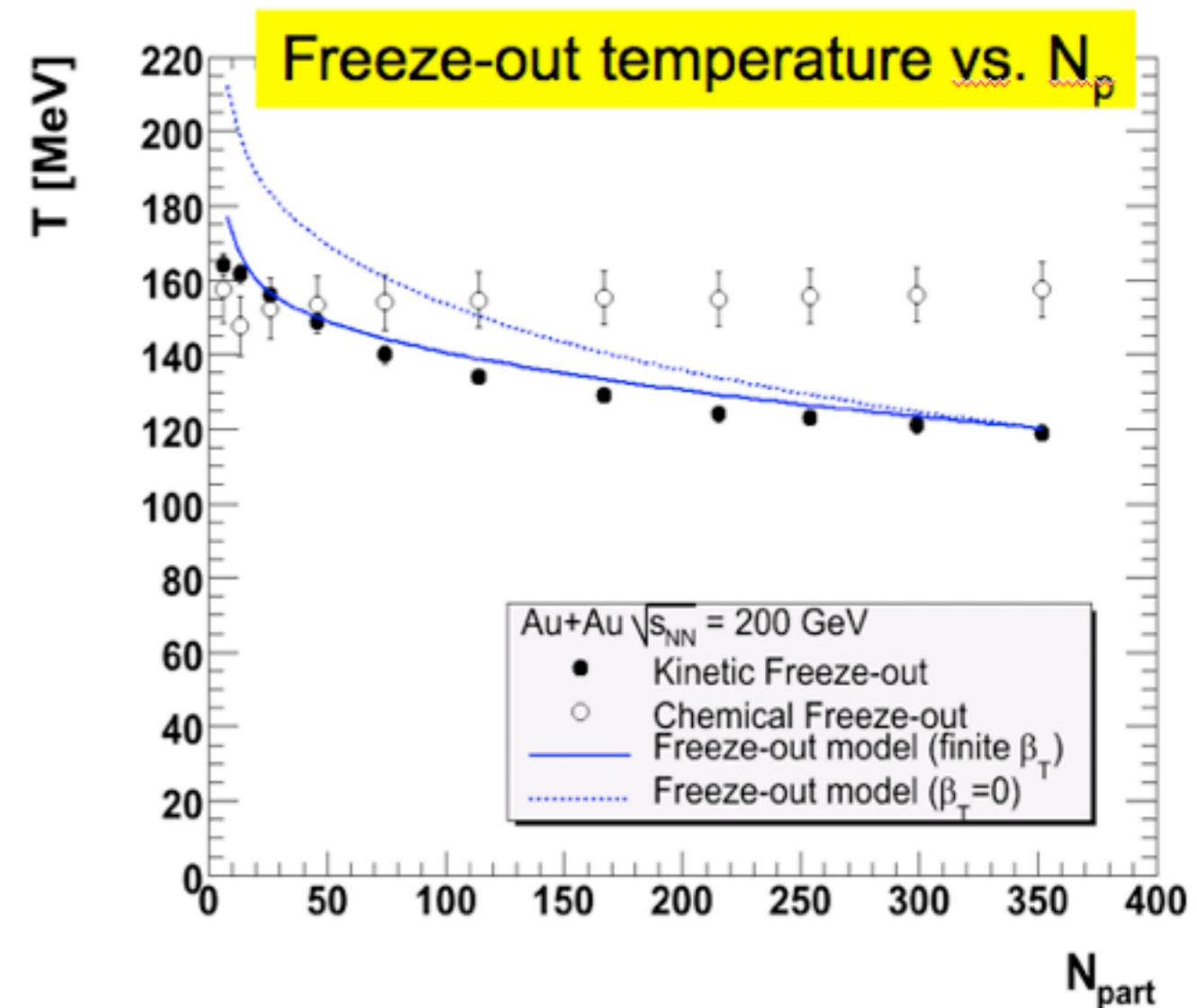
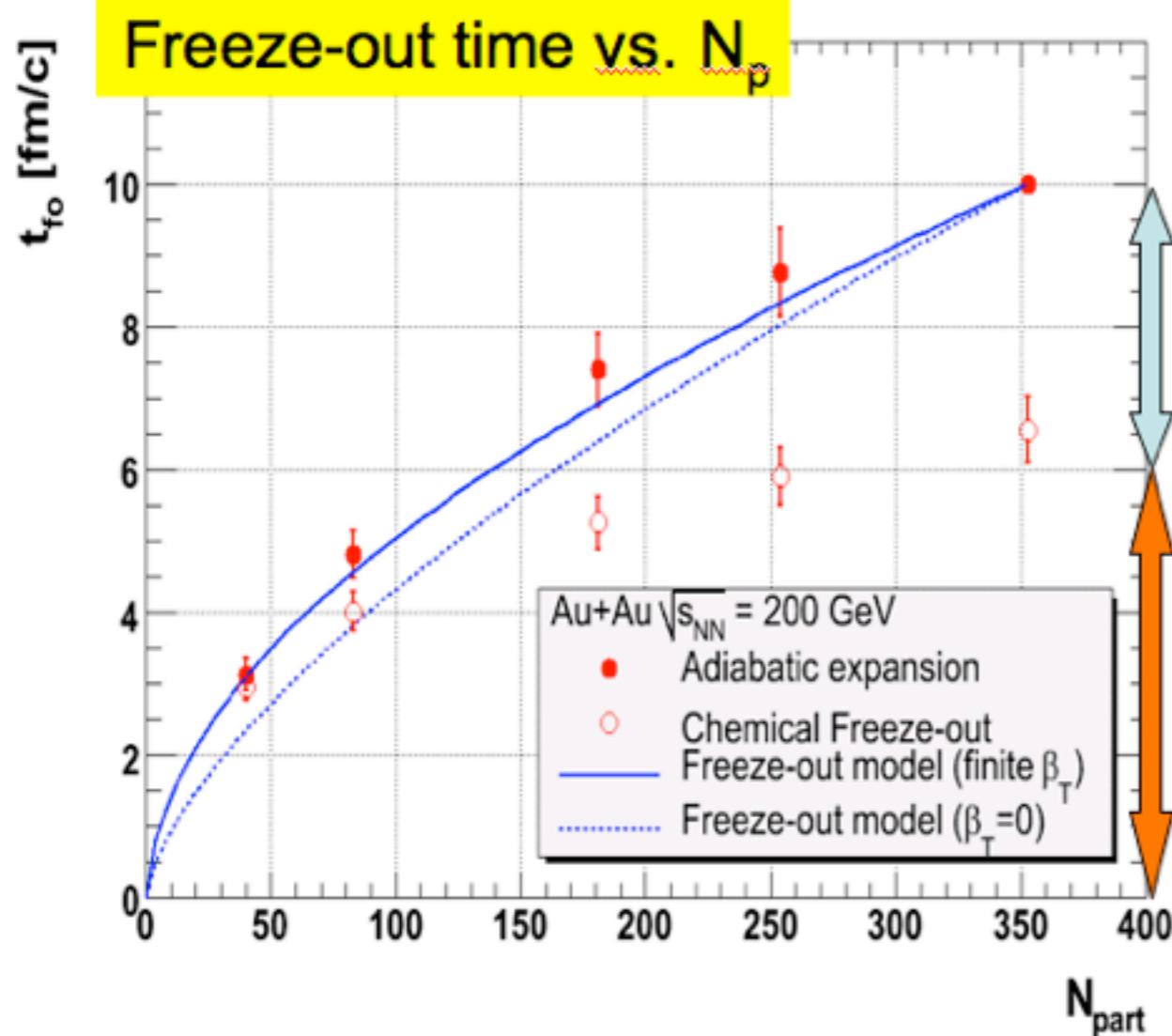
- Ideal gas:  $P=(1/3)\epsilon$
- Entropy conservation
- Longitudinal expansion & Transverse expansion



T scaled with  $(\epsilon_{bj})^{1/4}$  at  $t = 1 \text{ fm/c}$

Cooling stopped at  $T_{fo}$

# Freeze-out Time & Temperature

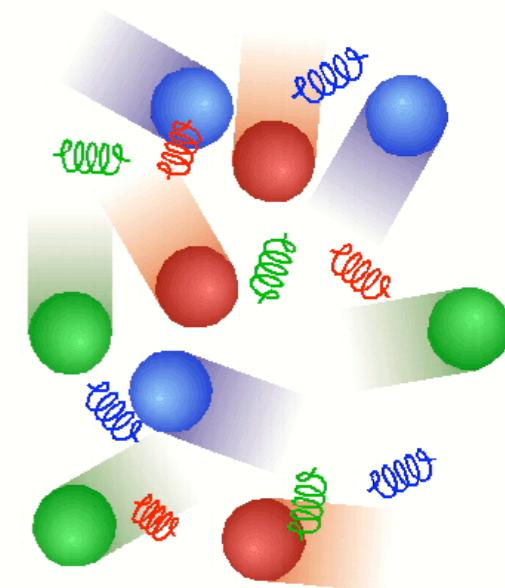
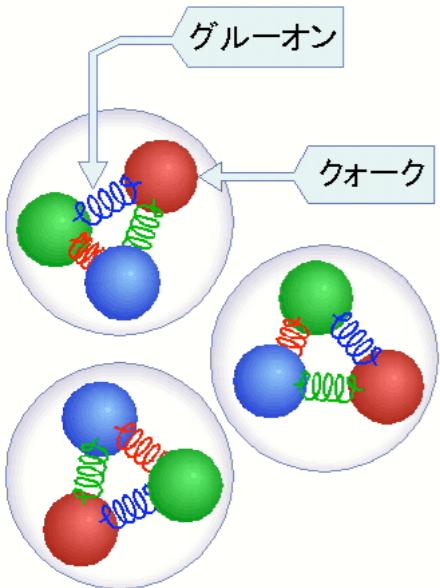


- More central collisions freeze out later at lower temperature.
- Consistent with freeze-out condition:  $\lambda(t)=R(t)$
- Even if quark phase is created before hadronization, **hadronic scattering** should be taken into account.

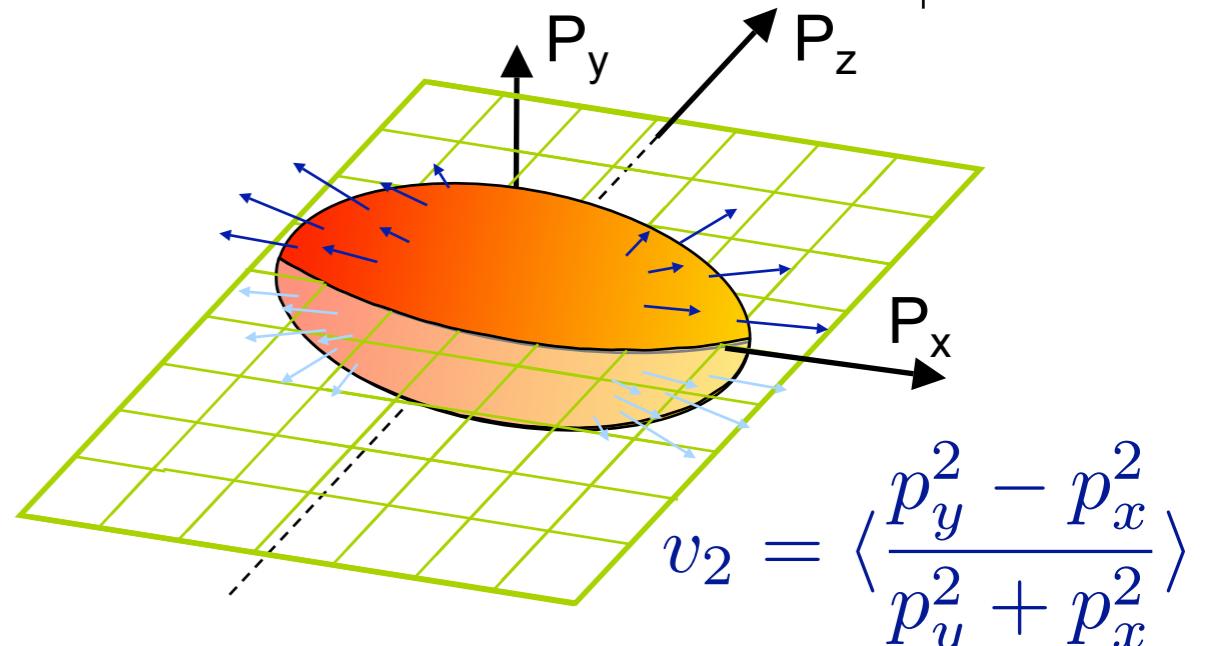
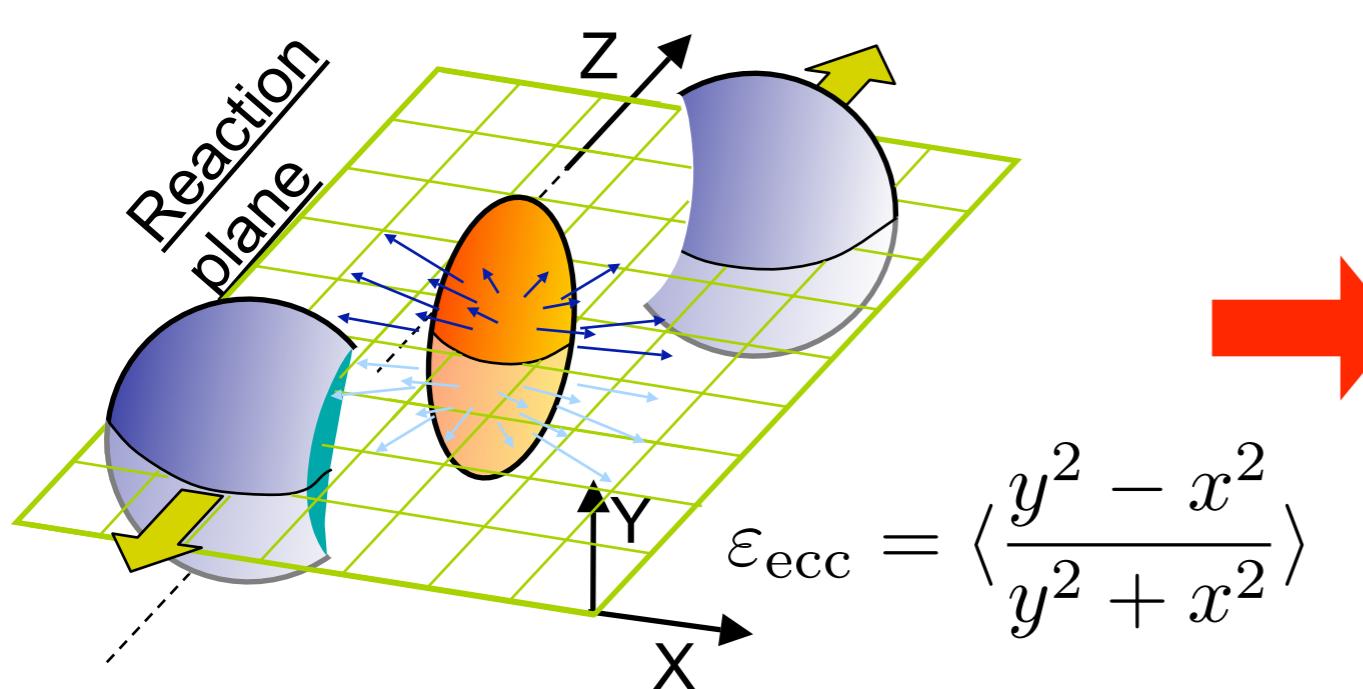
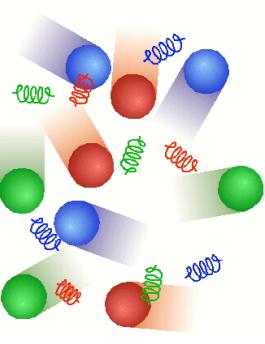
- As expected,  $T_{\text{fo}}$  is lower than  $T_{\text{ch}}$ . Different centrality dependence.
- $T_{\text{fo}}$  dropping is consistent with 1+1D adiabatic expansion.
- $T_{\text{c}} \sim T_{\text{ch}} \Rightarrow$  the observed chemical eq. not via **hadronic scatterings**.

# Two Major Discoveries at RHIC

- 1) Large Elliptic Flow
- 2) Jet Modification



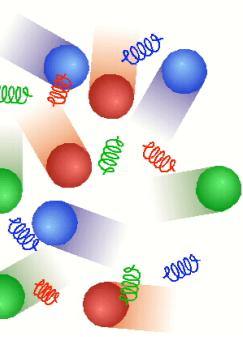
# Elliptic Flow, v<sub>2</sub> (Azimuthal Anisotropy)



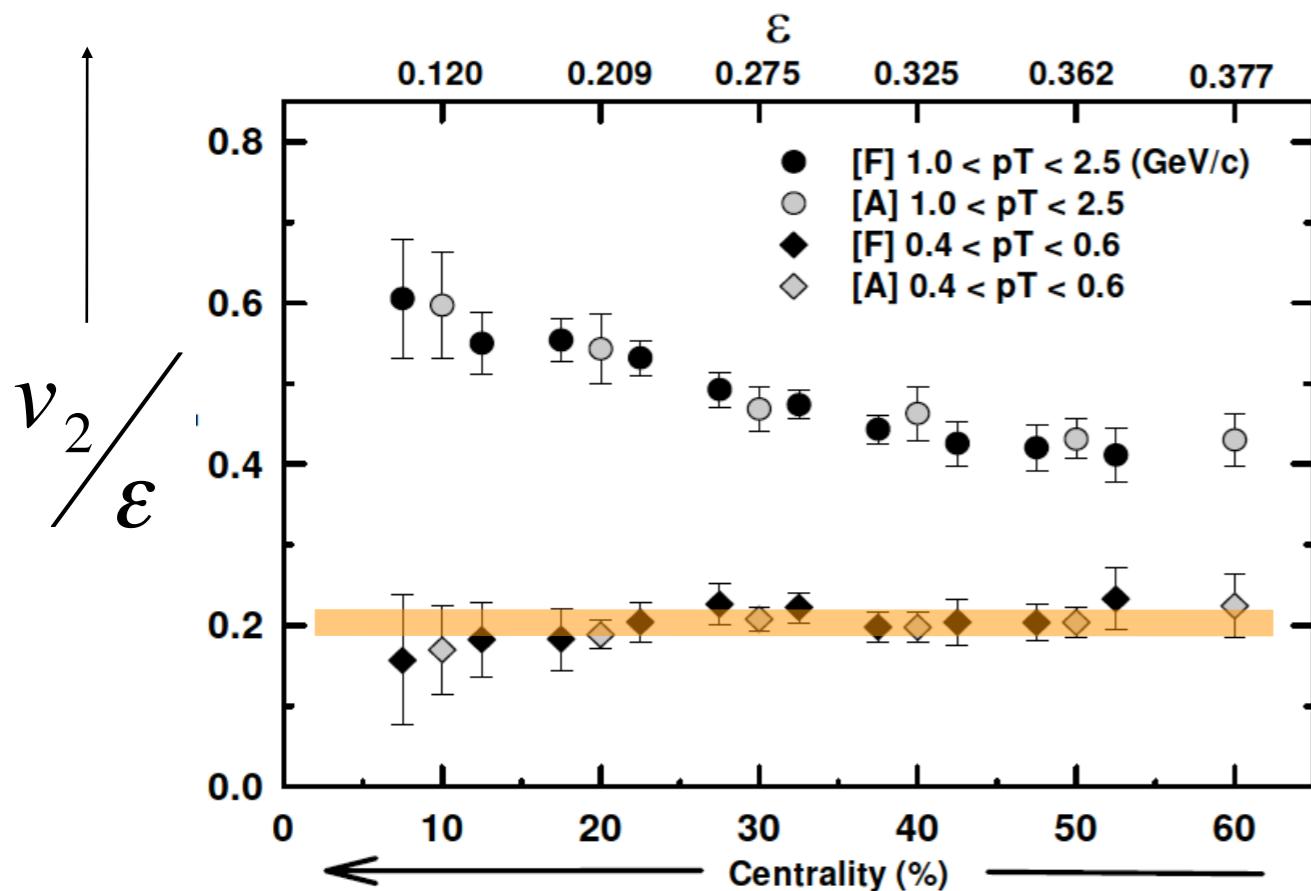
- ✓ In non-central collisions, participant region has almond shape.  
→ azimuthal anisotropy in coordinate space
- ✓ If  $\lambda \ll R$ , azimuthal anisotropy of the coordinate space is converted to that of the momentum space.  
→  $v_2$  ; second Fourier harmonics of azimuthal distribution
- ✓ Goodies :
  - ◆ Clear origin of the signal
  - ◆ Collision geometry can be determined experimentally

$$N(\phi) = N_0 \left\{ 1 + 2v_1 \cos(\phi - \Psi_0) + 2v_2 \cos(2(\phi - \Psi_0)) \right\}$$

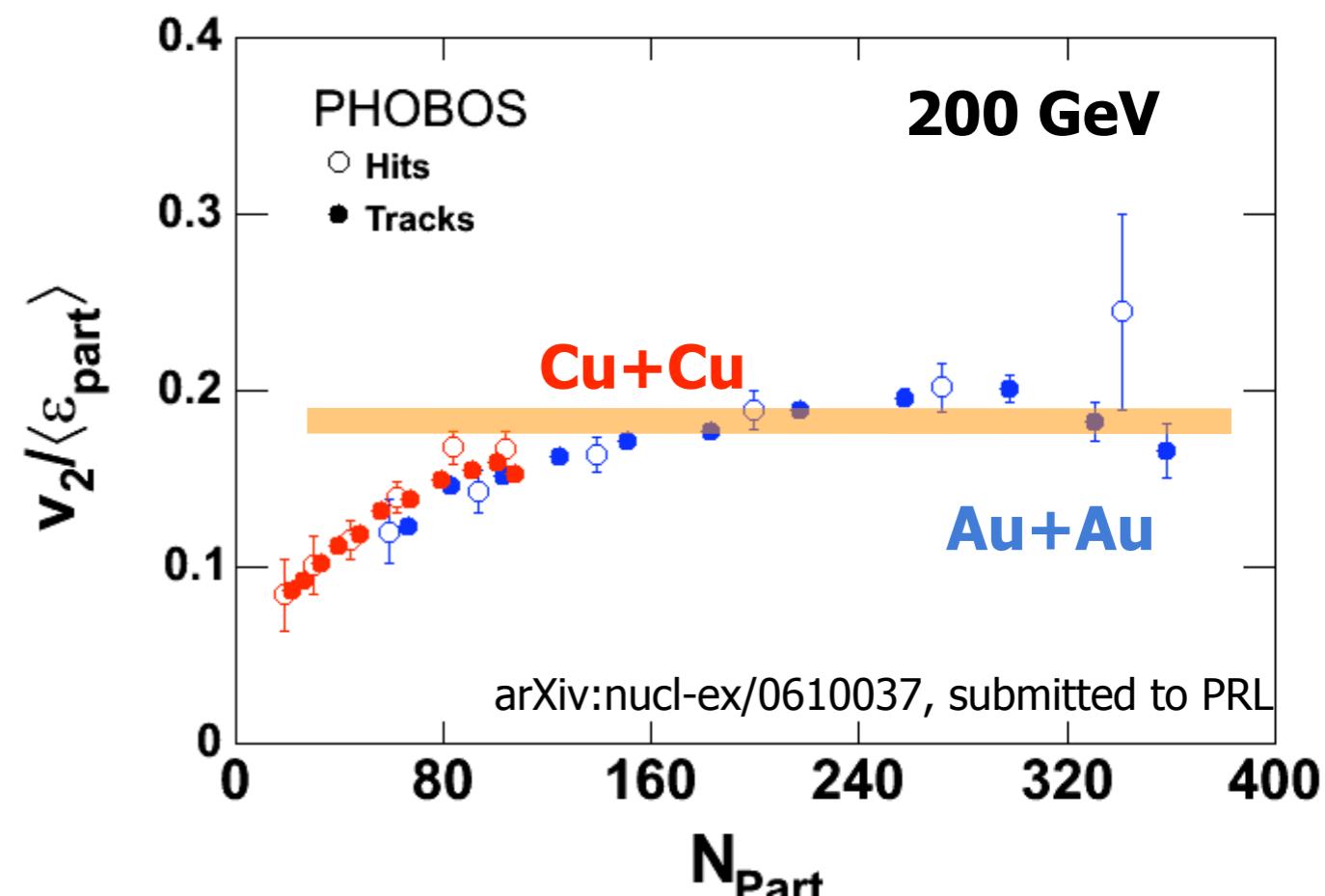
# $v_2$ vs. Eccentricity



Phenix; PRL 89(2002)212301



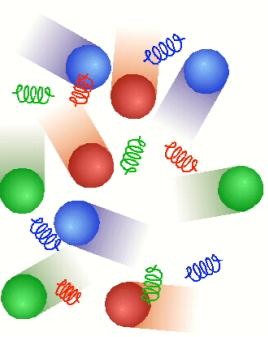
Phobos; nucl-ex/0610037



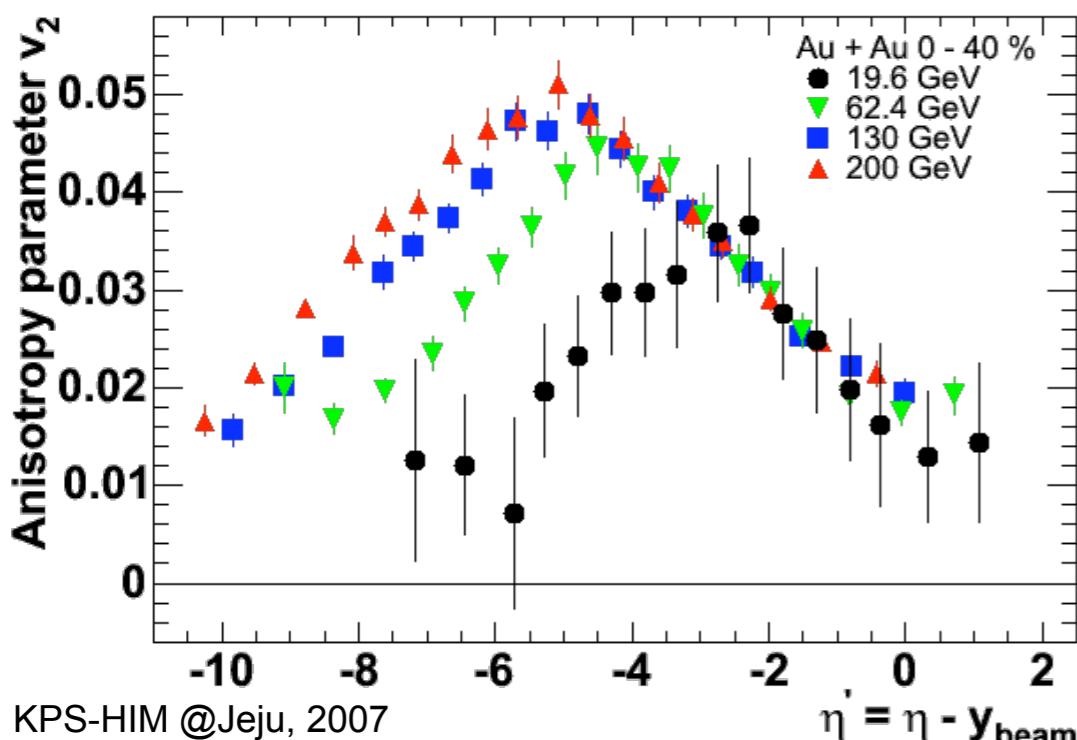
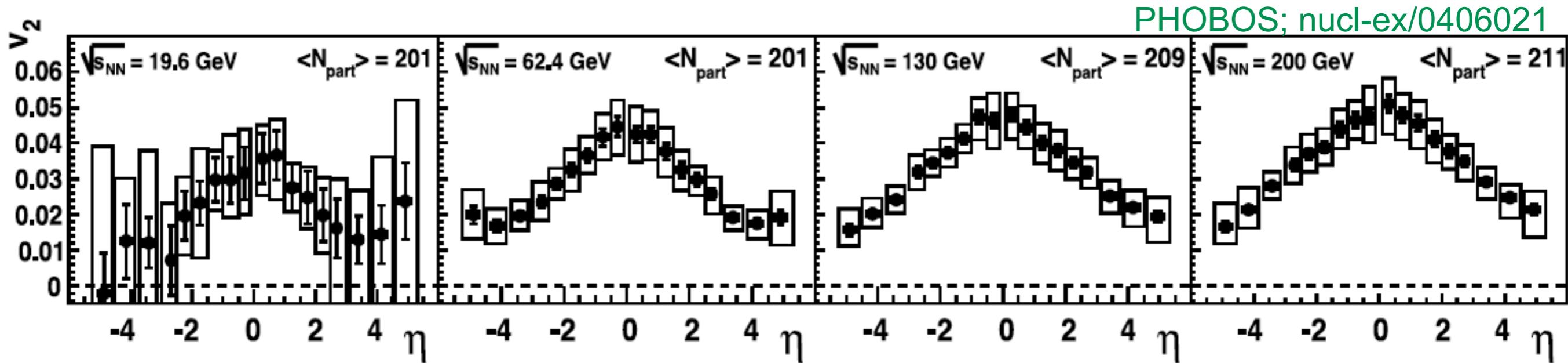
$$\therefore \langle v_2 \rangle \propto \epsilon$$

- ✓ Eccentricity is evaluated from centrality of collisions
- ✓ Ratio stays ~constant
- ◆ Eccentricity scaling observed in comparison of Au+Au, Cu+Cu
  - | → Scaling with eccentricity shows  $v_2$  builds up at early stage

# Large azimuthal anisotropy



$$N(\phi) = N_0 \left\{ 1 + 2v_1 \cos(\phi - \Psi_0) + 2v_2 \cos(2(\phi - \Psi_0)) \right\}$$

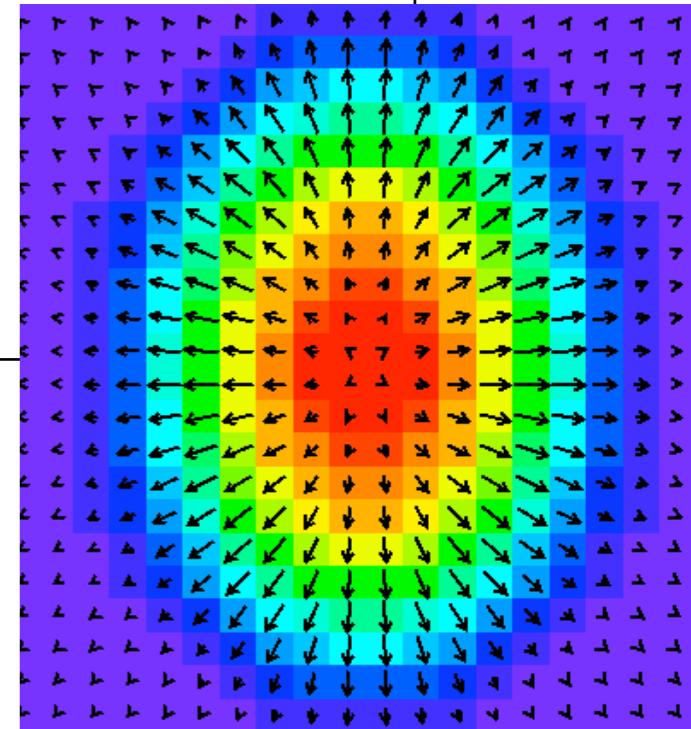
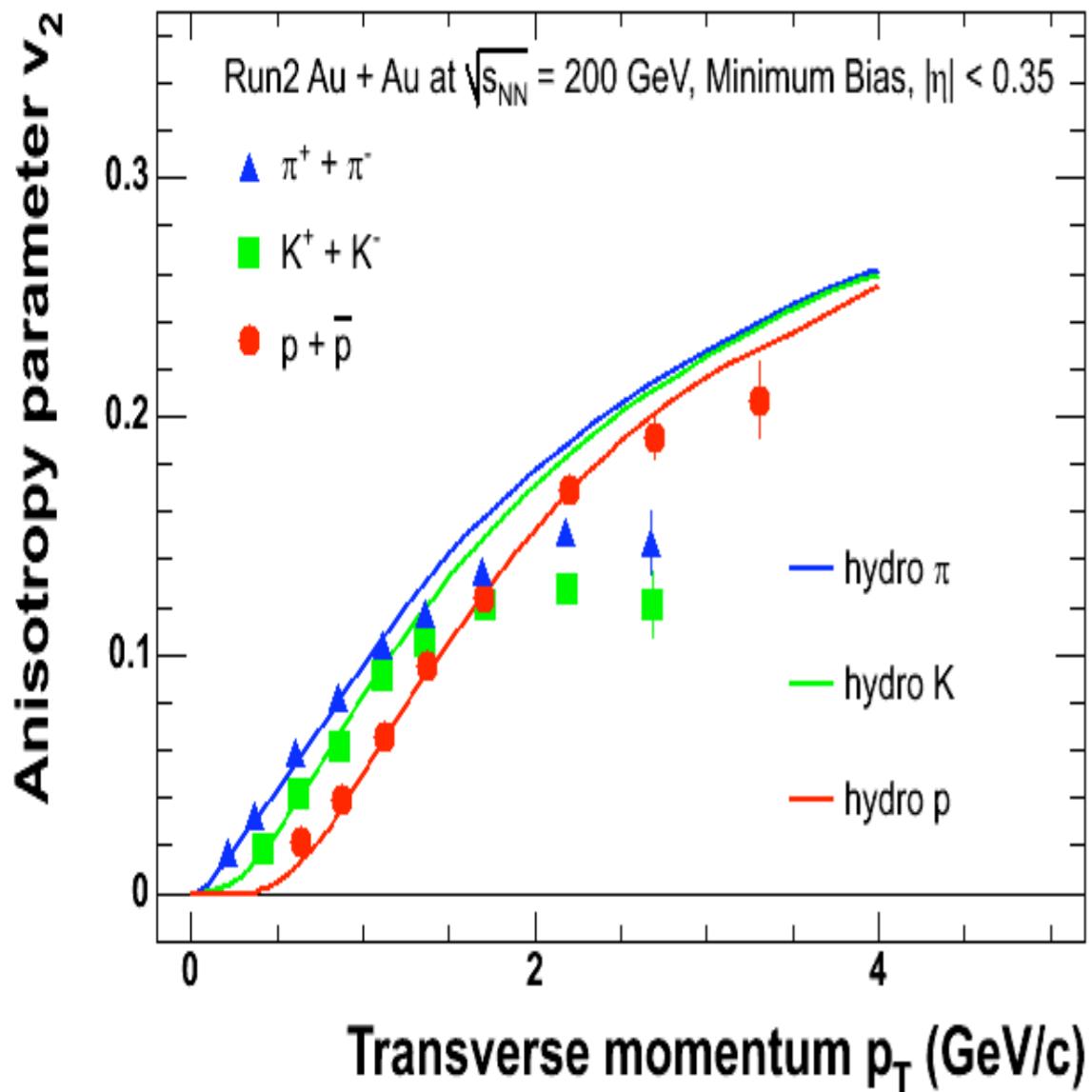


✓ Getting larger & larger in higher energies.

◆ Scaling w.  $\eta - y_{beam}$  !?

# Collective Flow

PHENIX : P.R.L. 91, 182301 (2003)



✓ Mass Ordering of  $v_2$  at low  $p_T$  region:

- Existence of collective flow
- ♦ Good agreement with hydrodynamics of perfect fluid
  - Early thermalization ( $\sim 0.6$  fm/c)
  - High energy density ( $\sim 20$  GeV/fm $^3$ )
  - Low viscosity

✓ Departure at high  $p_T$  region ( $> 1.5$  GeV/c);

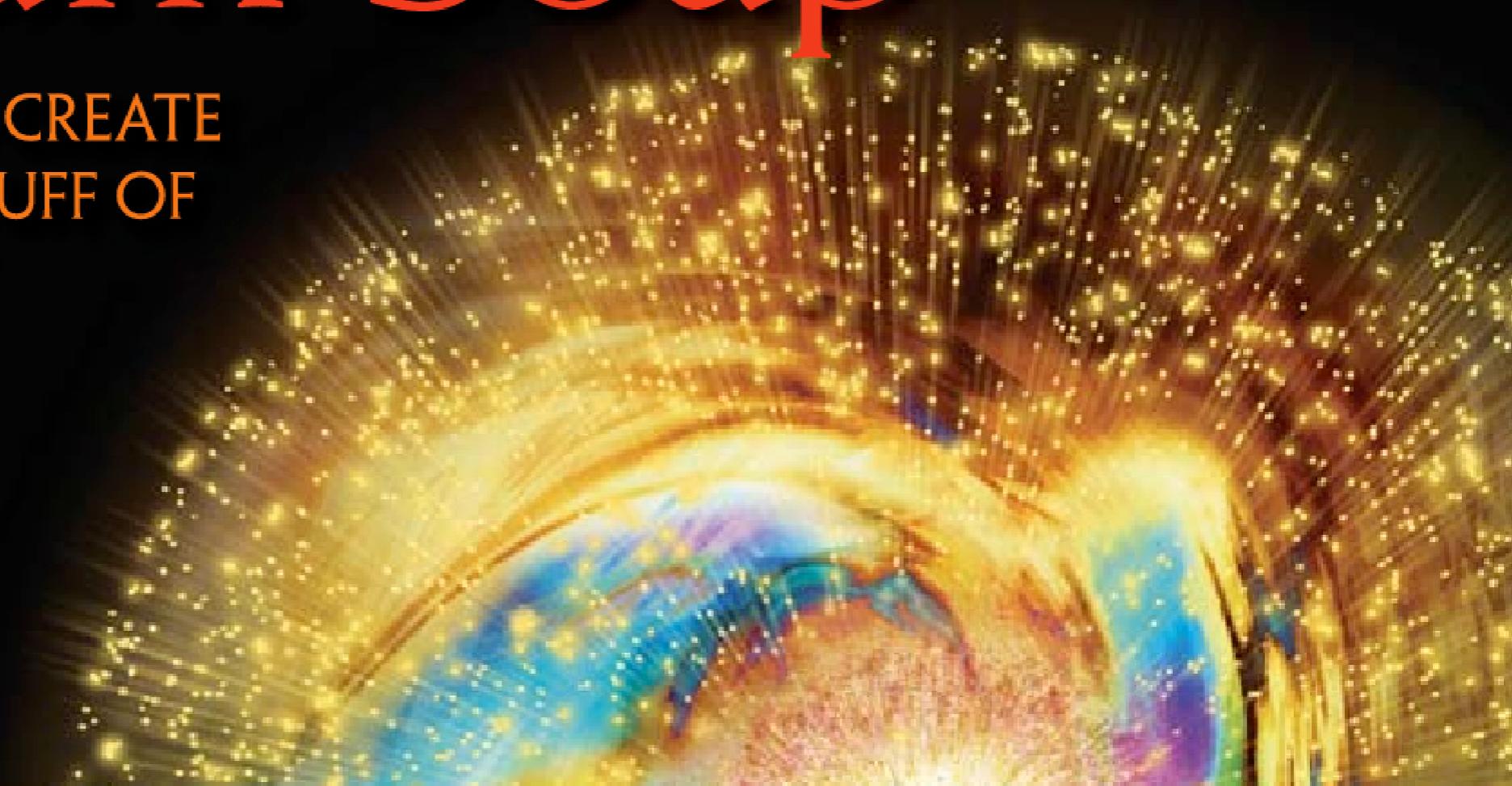
→ Other mechanism?

# SCIENTIFIC AMERICAN

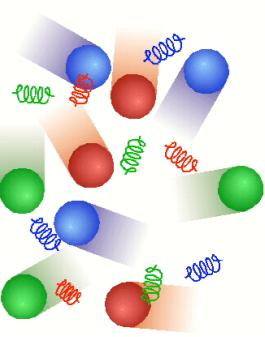
MAY 2006  
[WWW.SCIAM.COM](http://WWW.SCIAM.COM)

# Quark Soup

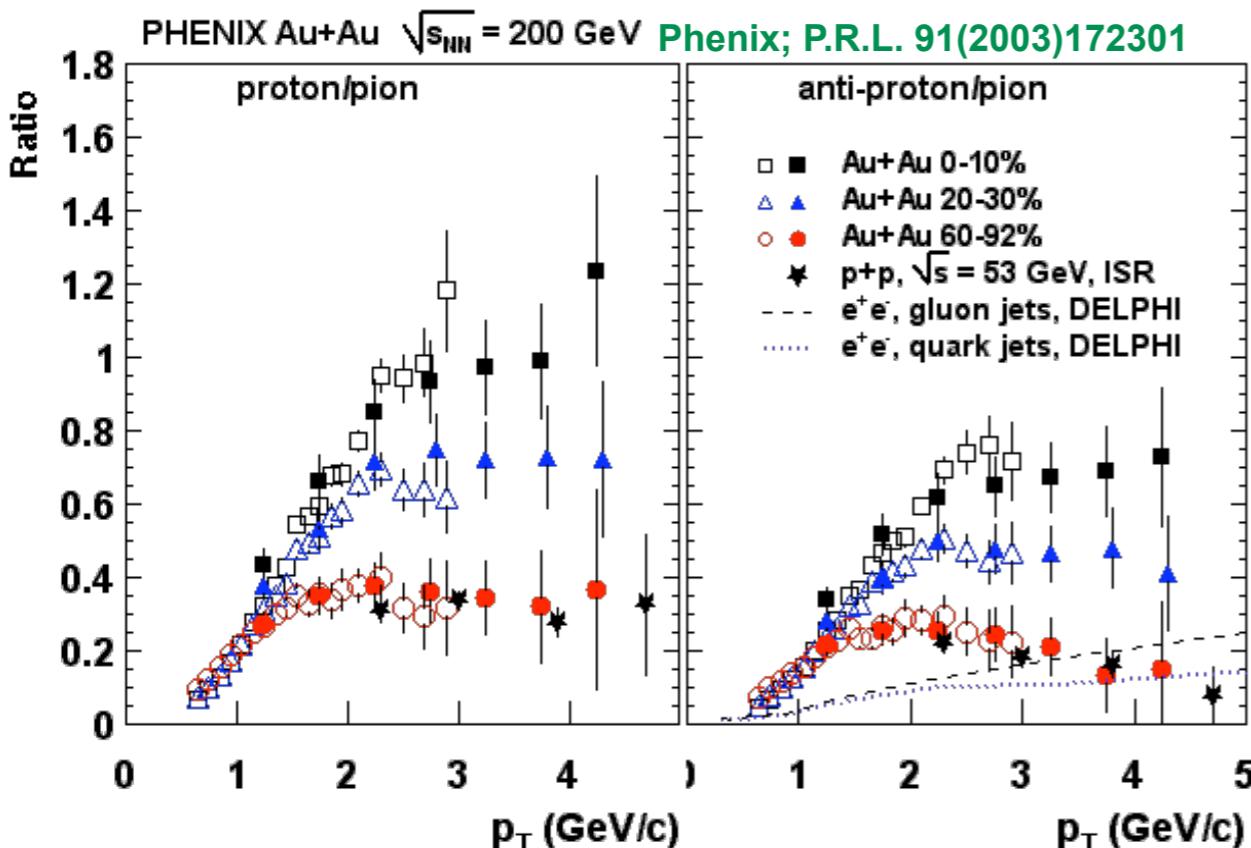
PHYSICISTS RE-CREATE  
THE LIQUID STUFF OF  
**THE EARLIEST  
UNIVERSE**



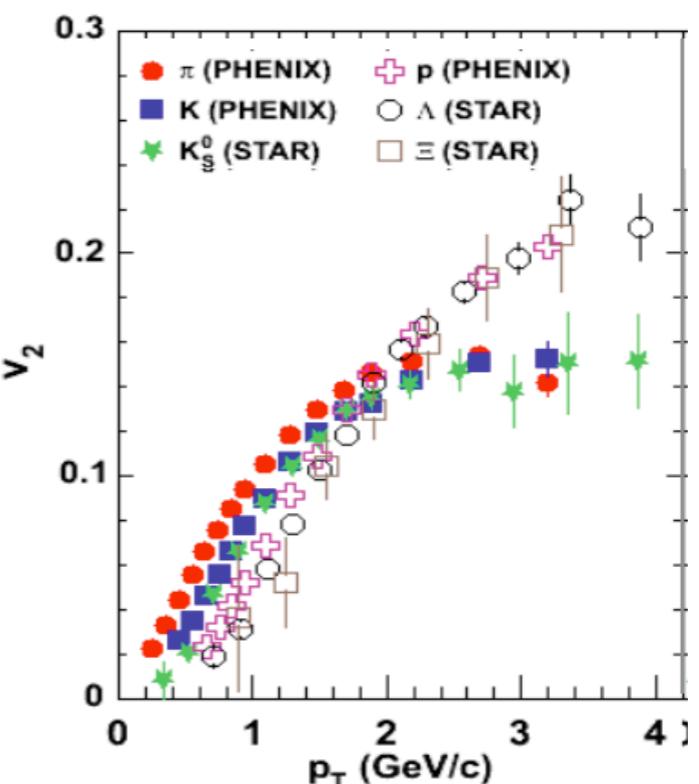
# Puzzles in the mid-pt region



$p/\pi$  enhances above 1.5 GeV/c



$v_2$  deviates from the mass ordering above 2 GeV/c



✓ In central col.,  $p/\pi$  ratio is very large, while in peripheral,  $p/\pi$  ratio similar to those in ee/pp suggesting fragmentation process.

| Fragmentation process should show  $n_p < n_\pi$  as seen in ee/pp.

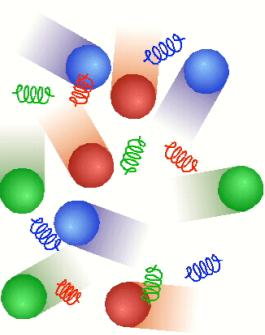
✓ While mass ordering of  $v_2$  seen at low pt region, clear departure observed.

✓ Suggesting other production mechanism.

↓

**Quark Recombination Model  
(Quark Coalescence Model)**

# Quark recombination model (RECO)

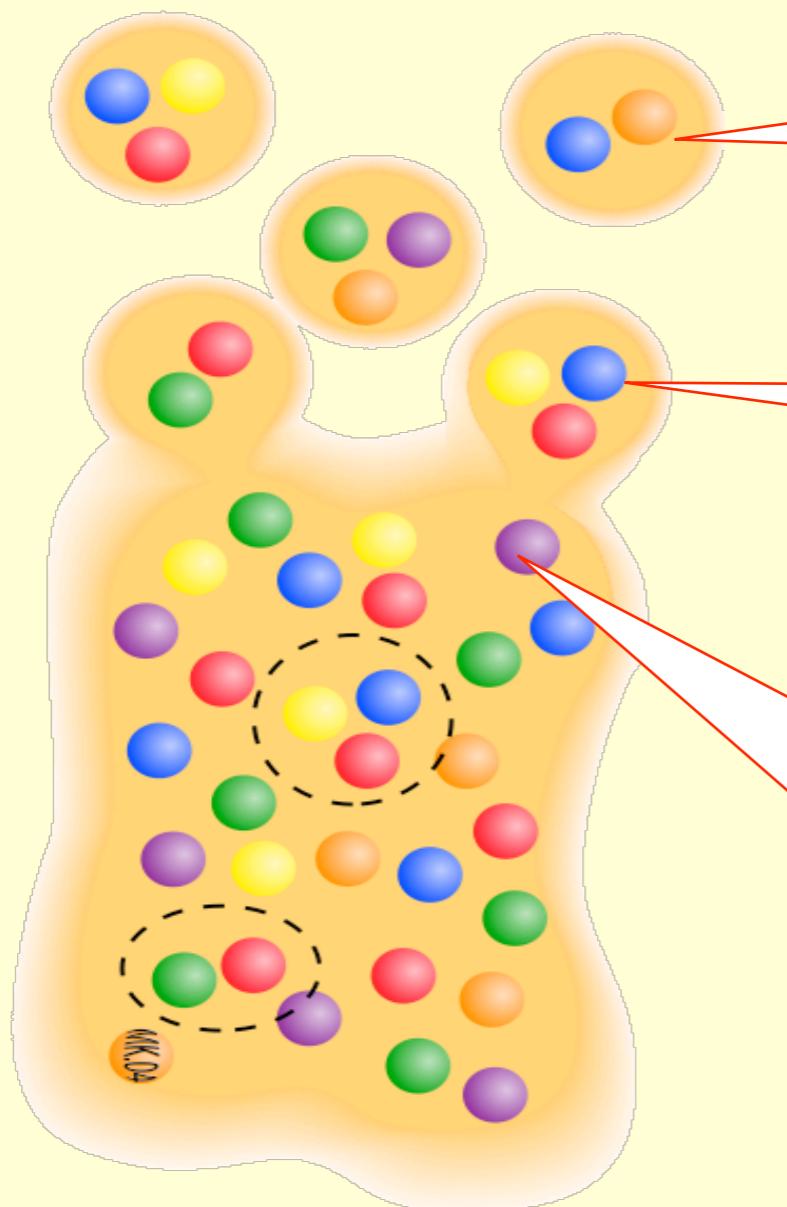


Hadron



QGP

Because of the steep distr. of  $w(pt)$ , RECO wins at high  $pt$  even w. small  $C_x$ .



✓ Quarks, anti-quarks combine to form mesons and baryons from universal quark distribution,  $w(pt)$ .

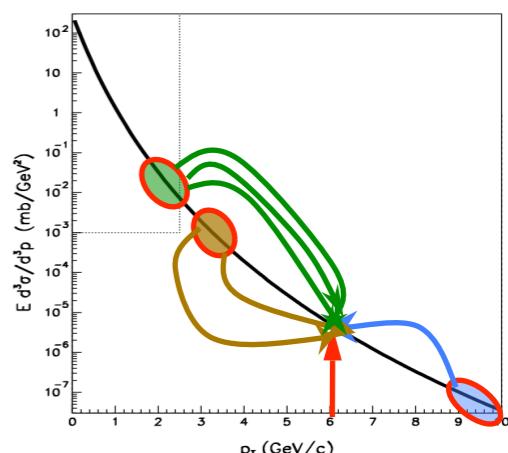
Mom. distr. of meson (2q);

$$W_M(p_t) \approx C_M \cdot w^2(p_t/2)$$

Mom. distr. of baryon (3q);

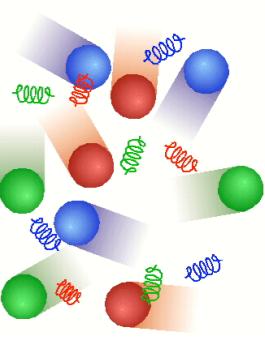
$$W_B(p_t) \approx C_B \cdot w^3(p_t/3)$$

$w(pt)$ ;  
Universal mom.  
distr. of quarks  
*[steep in pt]*



Characteristic scaling features expected.  
→ Quark number scaling

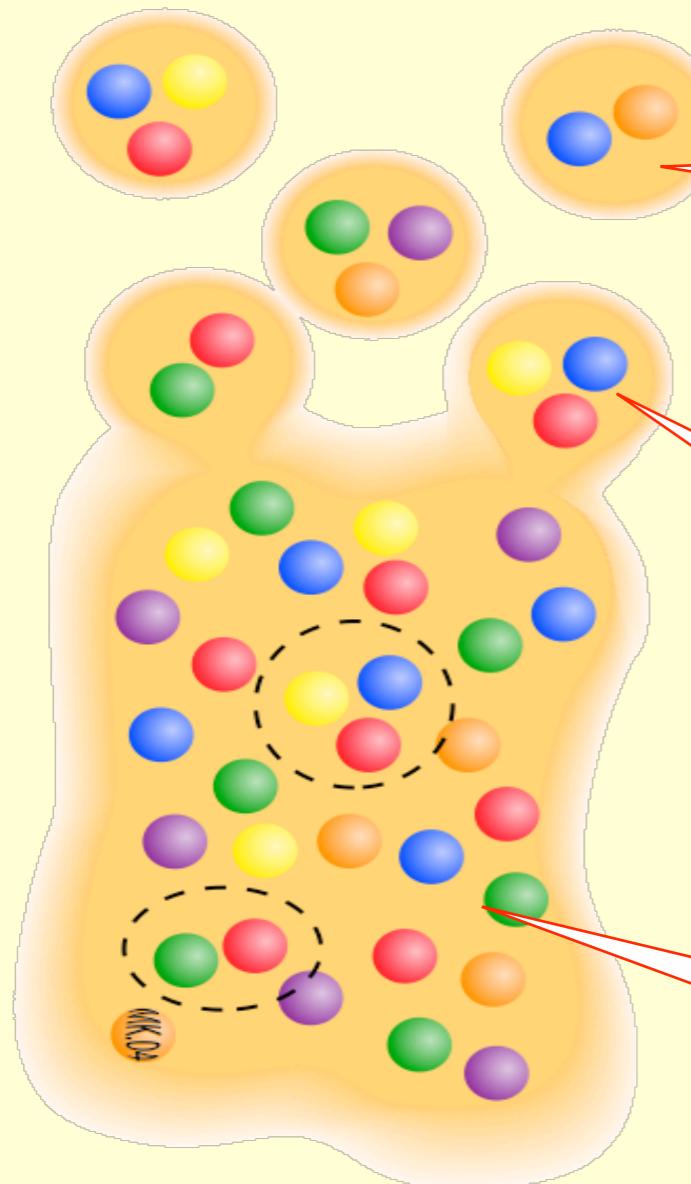
# $V_2$ from RECO



Hadron



QGP



✓ Characteristic scaling behavior

→ Quark Number Scaling

Azimuthal distr. of meson (2q);

$$\frac{dN_M}{d\phi} \propto w^2 = (1 + 2v_{2,q} \cos 2\phi)^2$$

$$\approx (1 + 4v_{2,q} \cos 2\phi)$$

Azimuthal distr. of baryon (3q);

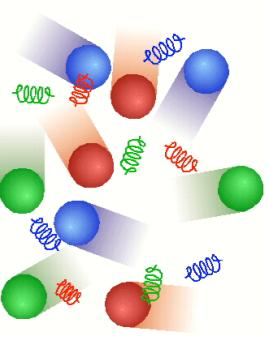
$$\frac{dN_B}{d\phi} \propto w^3 = (1 + 2v_{2,q} \cos 2\phi)^3$$

$$\approx (1 + 6v_{2,q} \cos 2\phi)$$

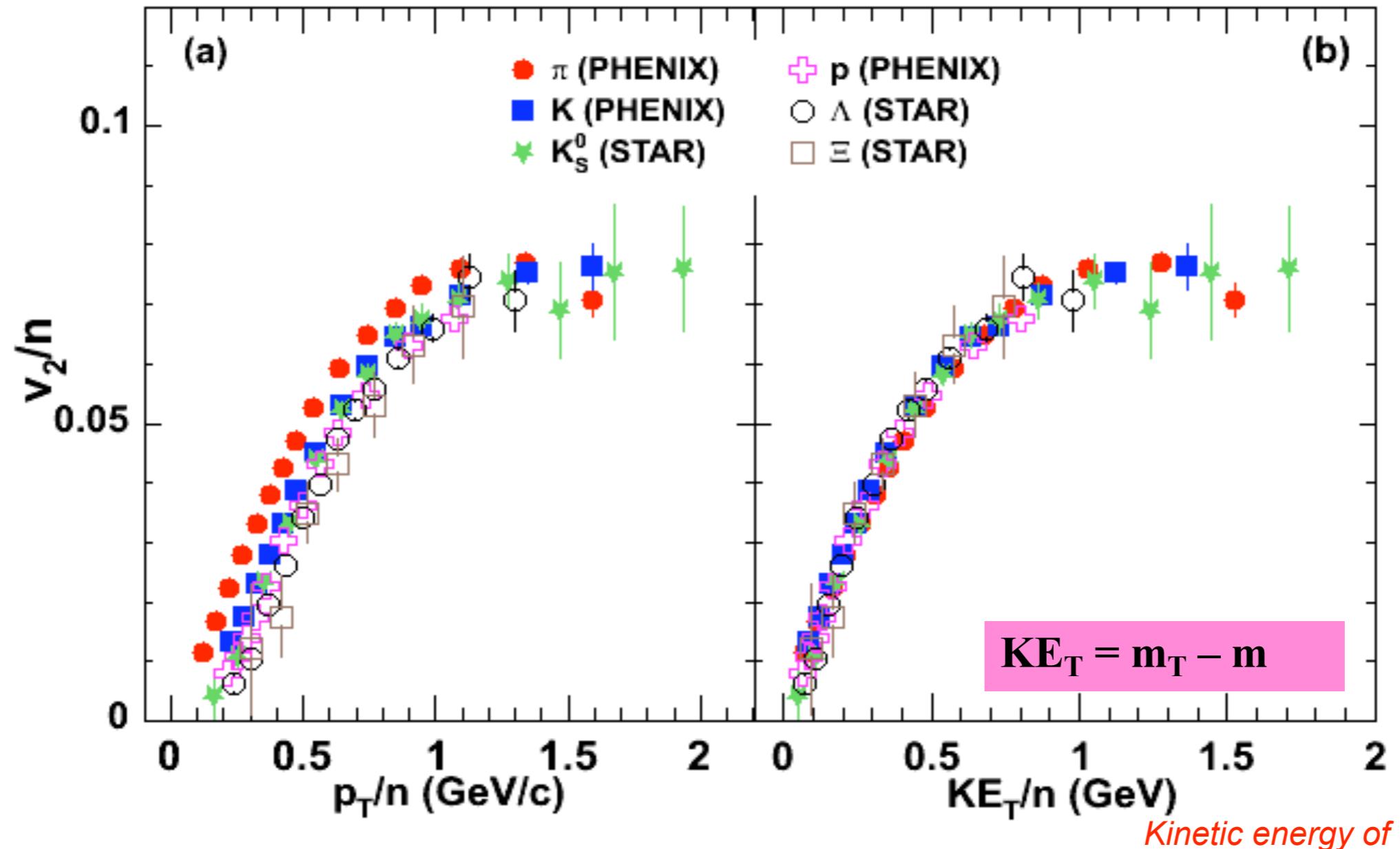
Azimuthal distr. of quark;  $w$

$$w \propto (1 + 2v_{2,q} \cos 2\phi)$$

# Quark Number Scaling !!

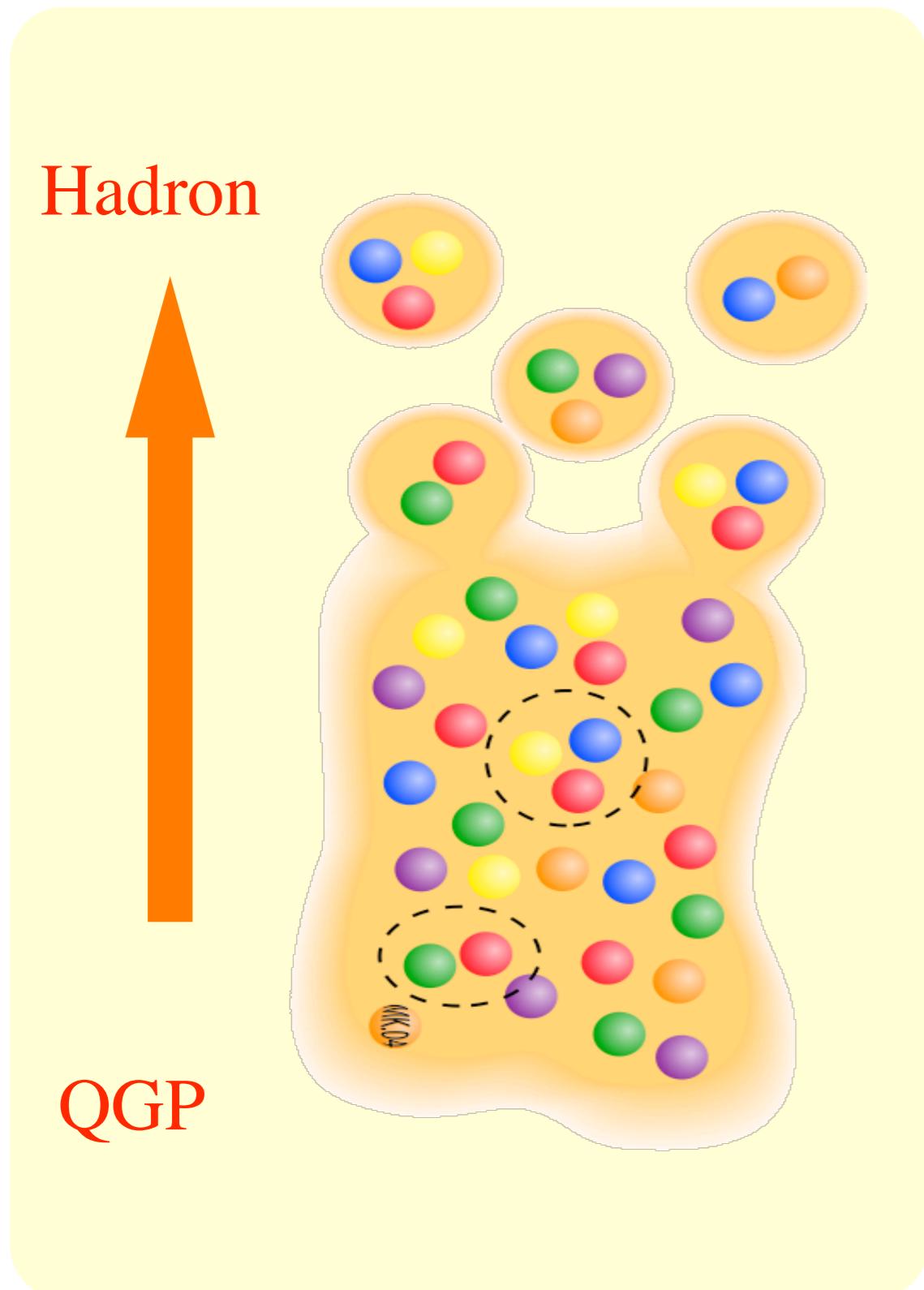


M. Issah, A. Taranenko, nucl-ex/0604011



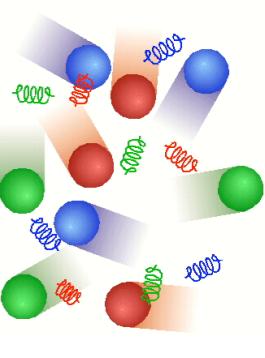
- ✓ Mesons and baryons made of light quarks seem to be consistent with the recombination model and there seem to be universal quark distribution,  $w(pt, \Phi)$ .

# Somewhat weird feeling of $KE_T$ scaling



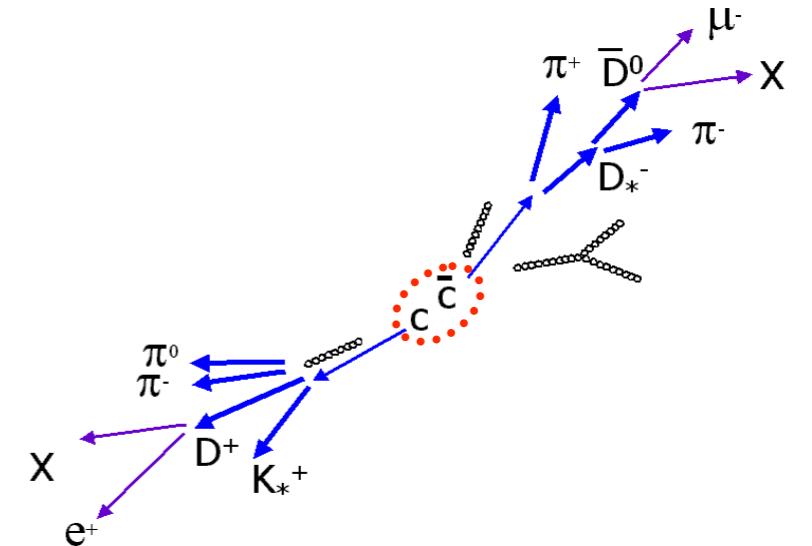
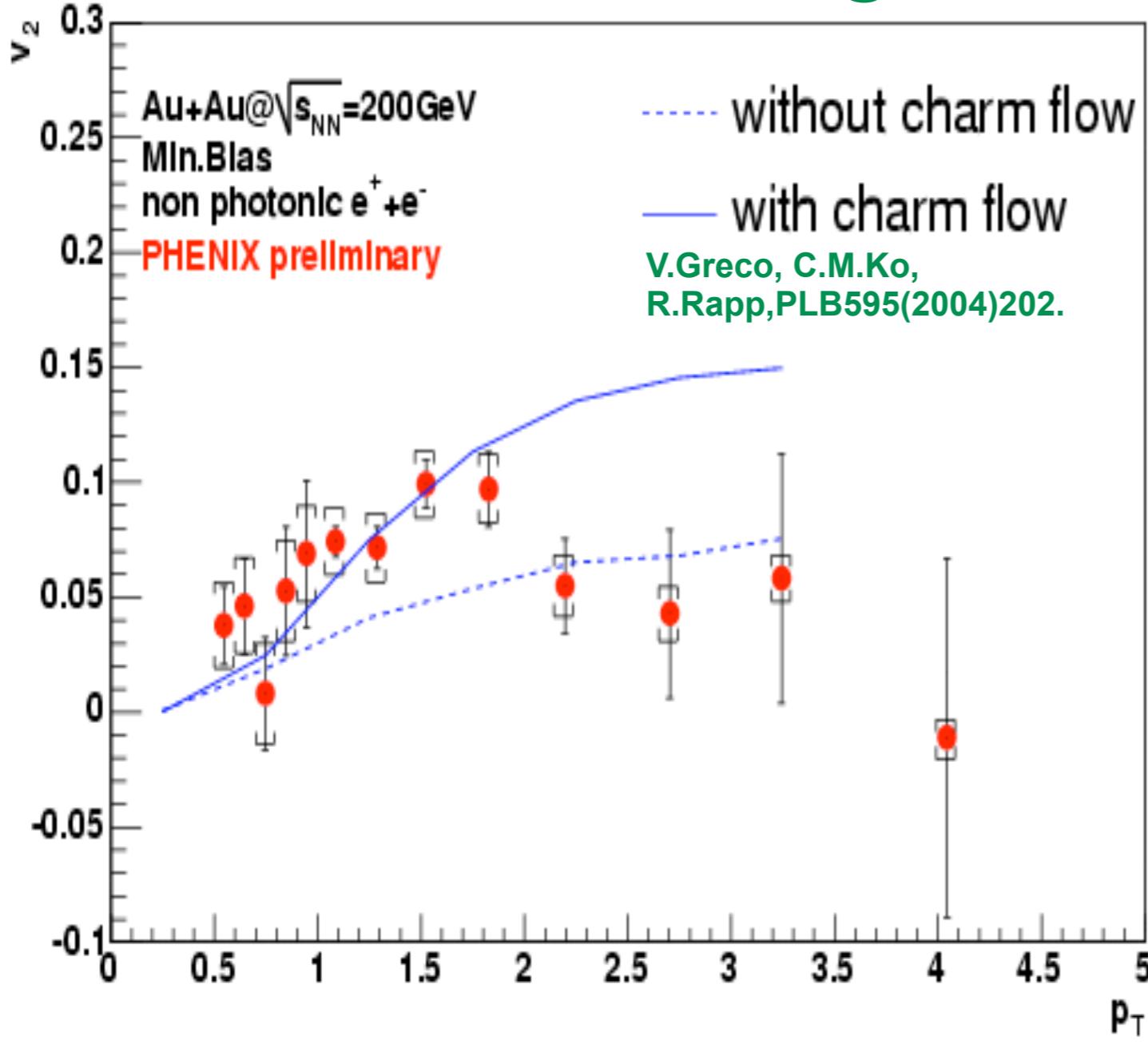
- Where or what is the role of gluons in Reco?
- Universal behavior with kinetic energy of effective mass of quarks ?
  - Quarks wear the gluons first, knowing which hadron they will be.(Effective mass)
  - Then, coalesced and boosted by the pressure behind.

# Even charm flows!



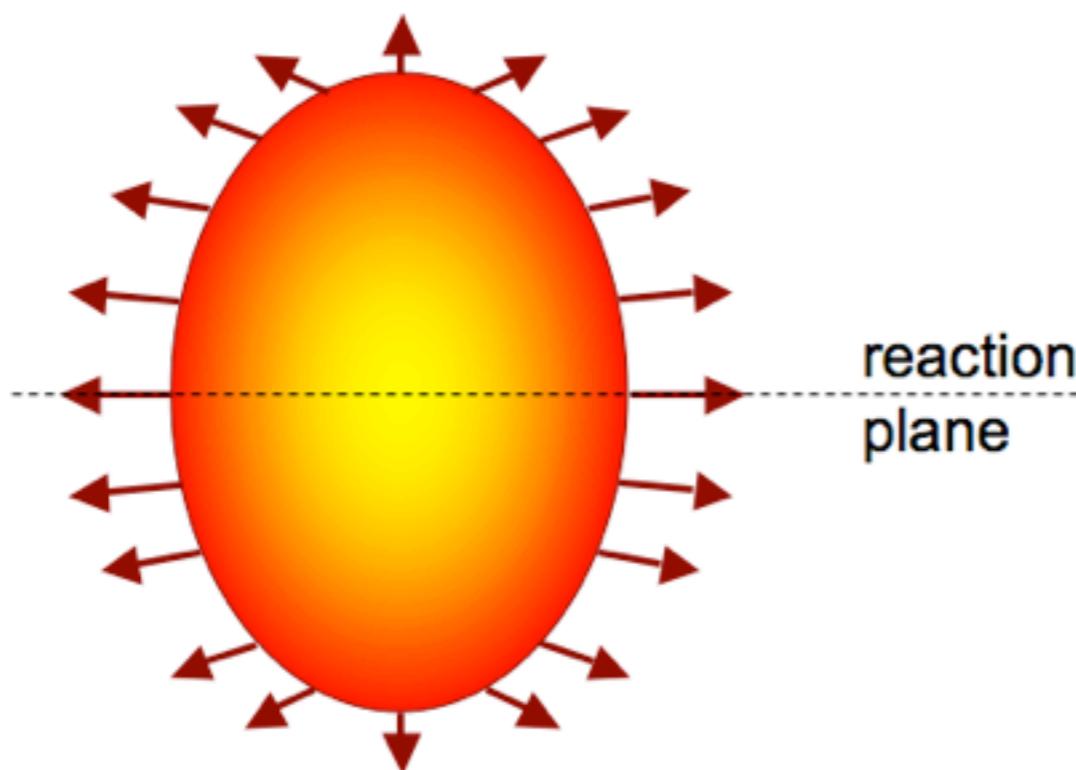
v2 of single electrons

S.Esumi & S. Sakai@SQM2006



- ✓ Charm decay produces high energy electrons.
- ✓  $v_2$  of single electrons are measured.
- ✓ Observed data favors flow of charm, suggesting thermalization of heavy quarks.
- ✓ This supports quark-coalescence & formation of QGP.
- ✓ Best to measure  $v_2$  of  $J/\psi$

# Extended Blast-wave model



$$E \frac{d^3N}{dp^3} = \int_{\Sigma_f} f(x, p) p^\mu d\Sigma_\mu$$

$$f(x, p) = \frac{1}{(2\pi)^3} \frac{1}{e^{[(p_\mu u^\mu(x) - \mu(x))/T(x)]} \mp 1}$$

$$v_2 = \frac{\int_0^R r dr \int_0^{2\pi} d\phi \cos(2\phi) I_2(\alpha) K_1(\beta)}{\int_0^R r dr \int_0^{2\pi} d\phi I_0(\alpha) K_1(\beta)}$$

$$\alpha = \left(\frac{p_T}{T}\right) \sinh \rho, \quad \beta = \left(\frac{m_T}{T}\right) \cosh \rho$$

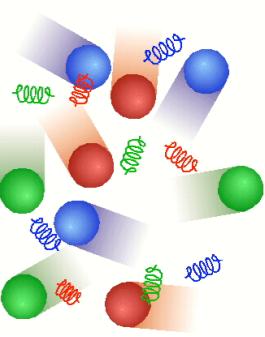
$$\beta_T(r) = \tanh \rho$$

- Extend standard Blast-wave parameterization in order to fit  $v_2(p_T)$  in non-central collisions
- Assumptions
  - Use density distributions from initial geometry overlap instead of uniform density
  - Velocity profile
    - Use Density gradient distributions
  - Eccentricity ( $\epsilon$ ) is fixed by initial overlap density
  - Velocity anisotropy ( $\beta_2$ ) is fixed by the velocity profile

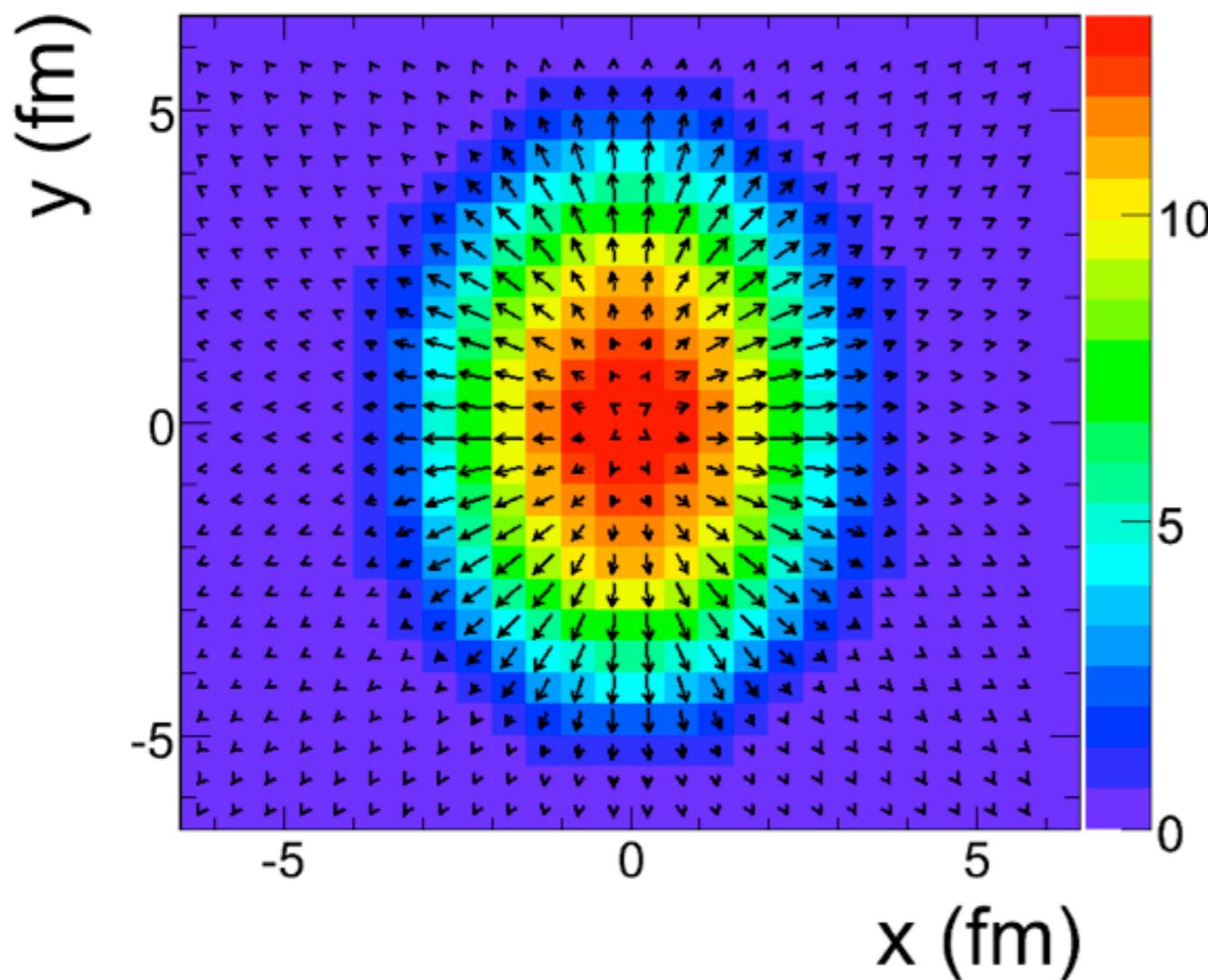
## 2 free parameters

- Temperature: T
- Magnitude of transverse boost velocity:  $\beta_T$

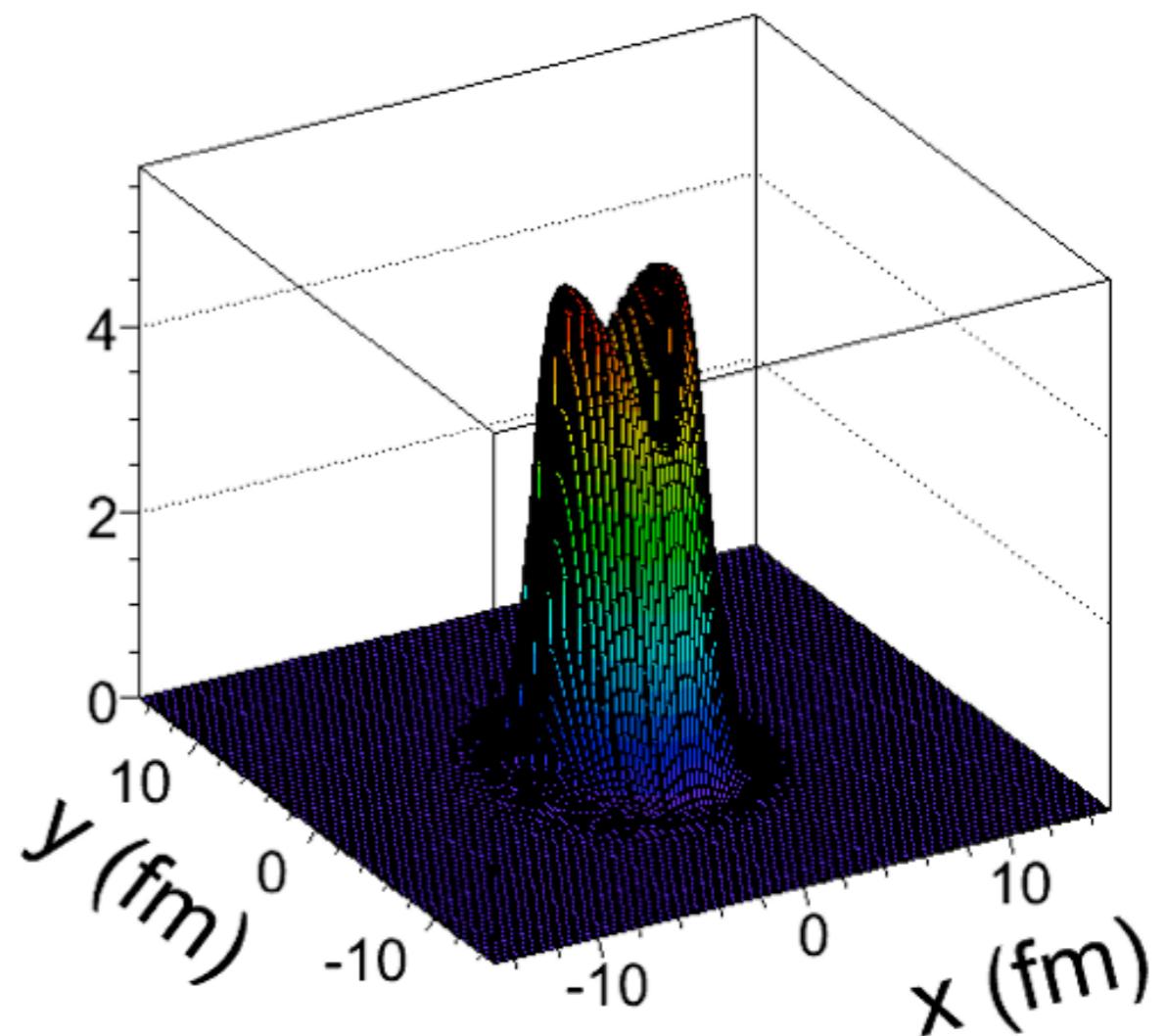
# N<sub>part</sub> profile & N<sub>coll</sub> profile



N<sub>coll</sub>, 20 - 30 % centrality

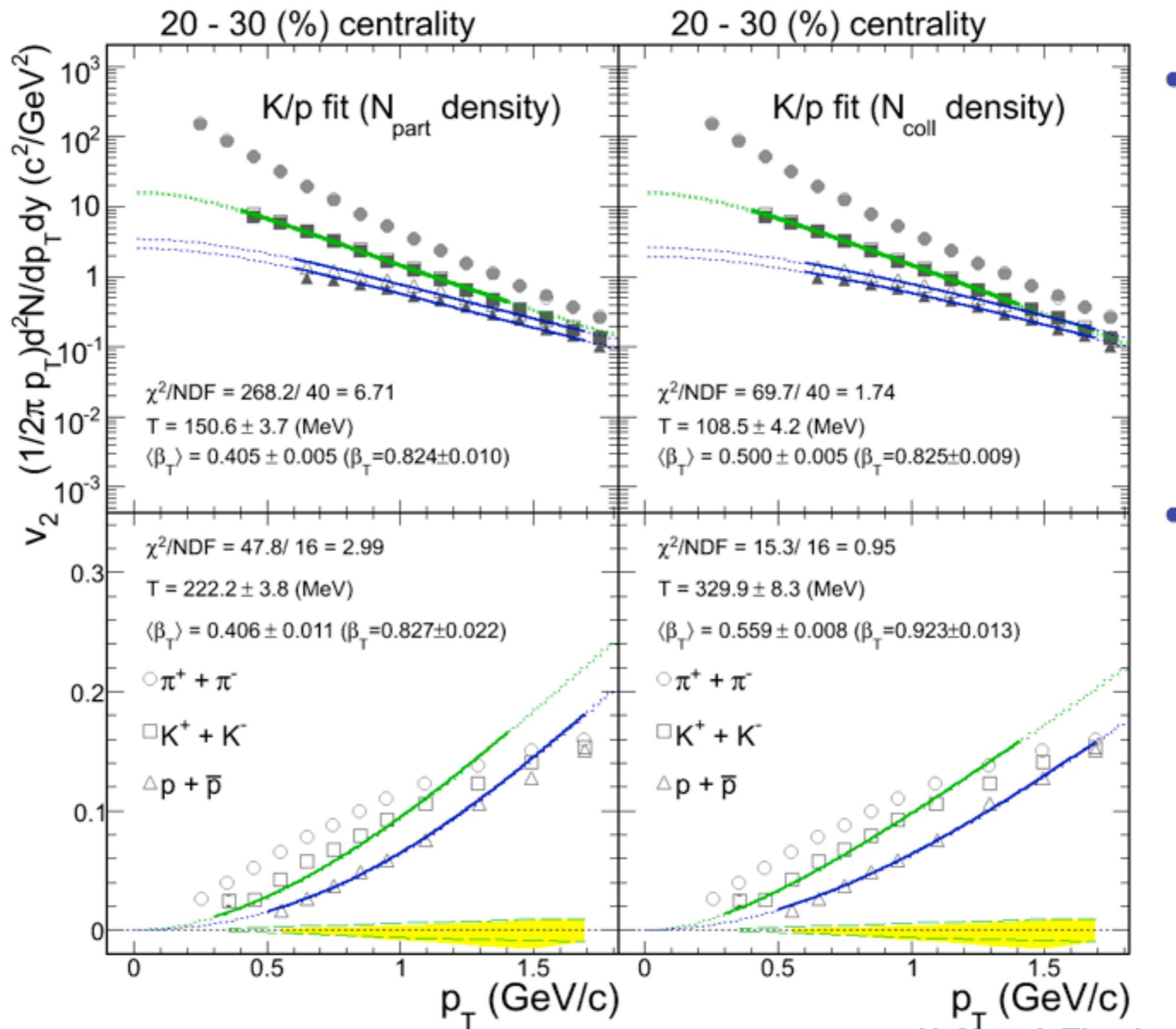
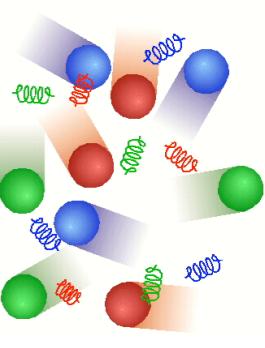


N<sub>coll</sub> gradient distribution, 20 - 30 % centrality



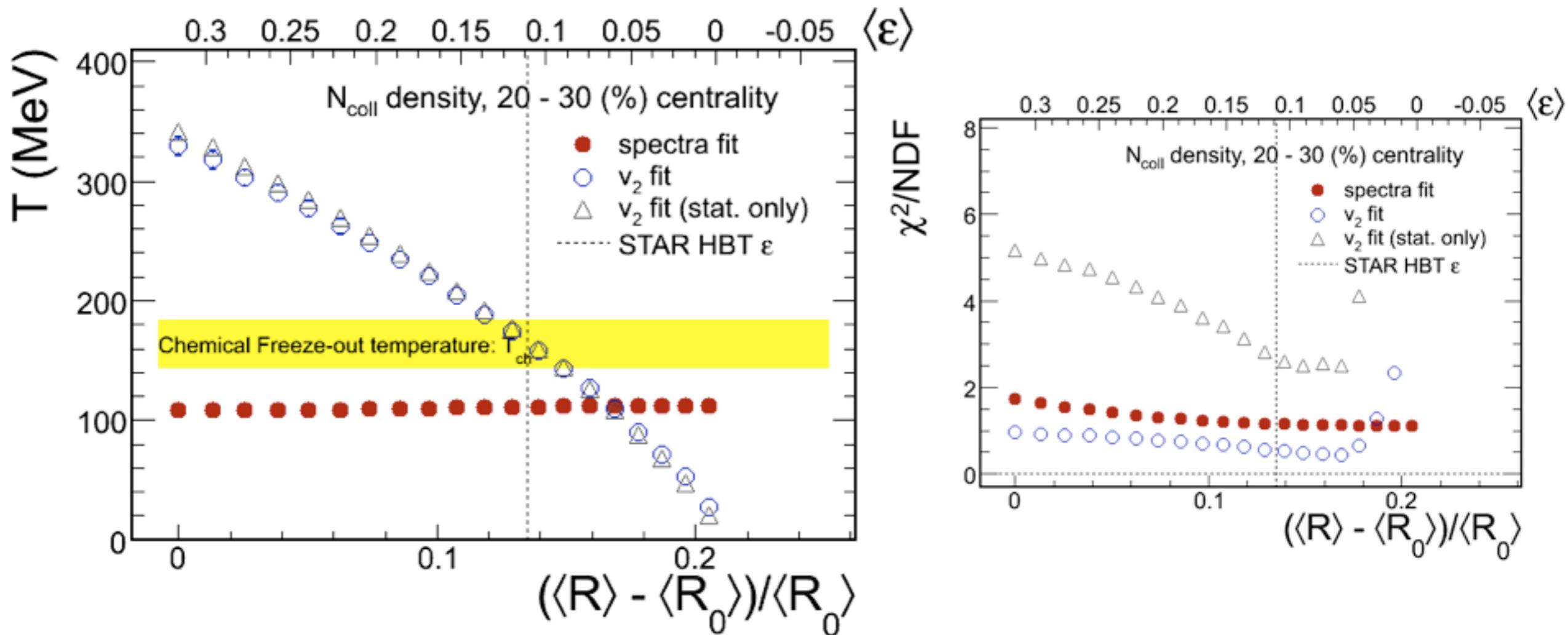
- Calculate N<sub>part</sub>(x, y) and N<sub>coll</sub>(x, y) from Woods-saxon density profile
  - Direction of density gradient  $\Rightarrow$  direction of boost
    - Length = magnitude of boost

# N<sub>part</sub> profile vs N<sub>coll</sub> profile

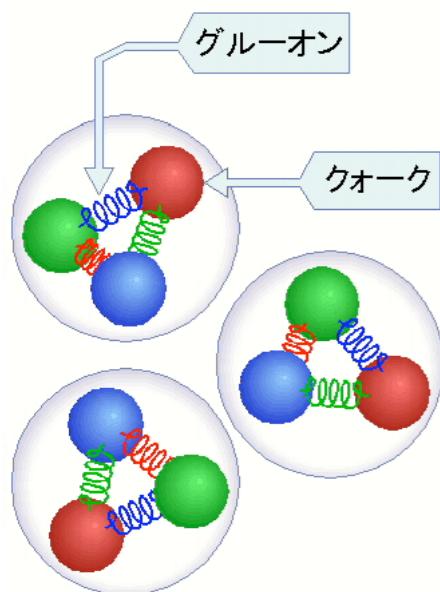


# Results

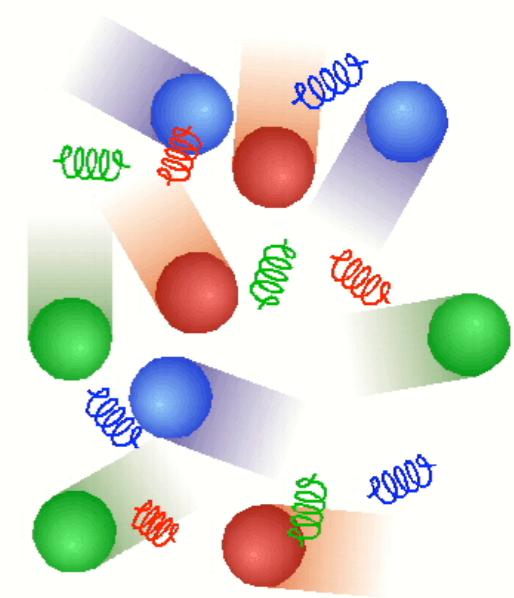
36



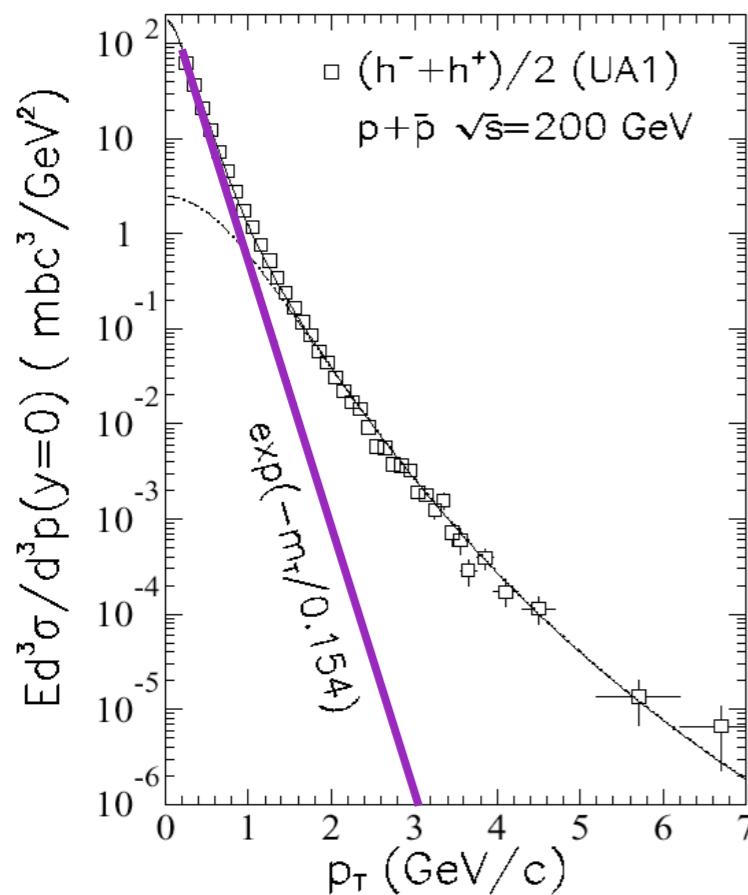
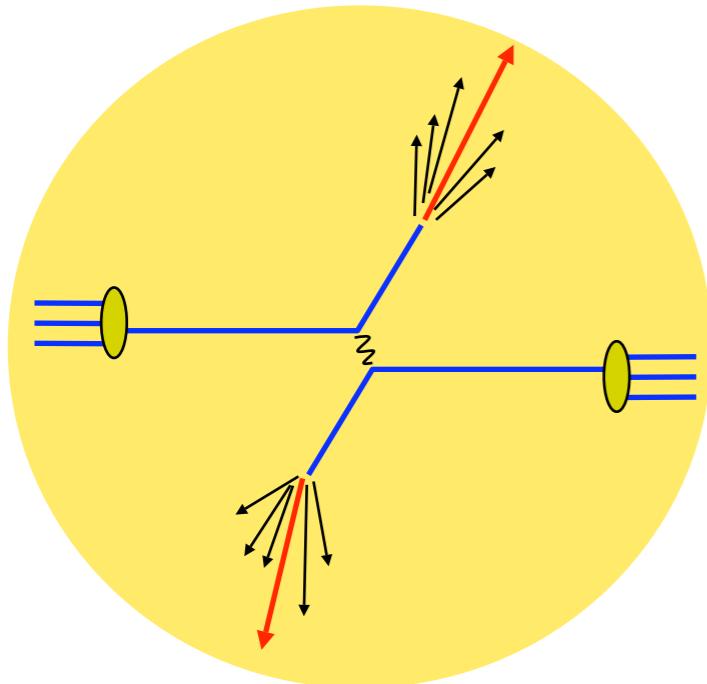
- Radial velocity and  $T$  from spectra are not changed with  $\langle \epsilon \rangle$
- Temperature from  $v_2$  fit strongly decrease with  $\langle \epsilon \rangle$ 
  - $\chi^2/\text{NDF} < 1$  up to  $\epsilon \sim 0.05$  with systematic error
- $T(v_2) \sim T_{ch}$  (150 MeV) >  $T(\text{spectra})$  at the eccentricity extracted by HBT analysis  $\Rightarrow$  early freeze-out of  $v_2$  ?



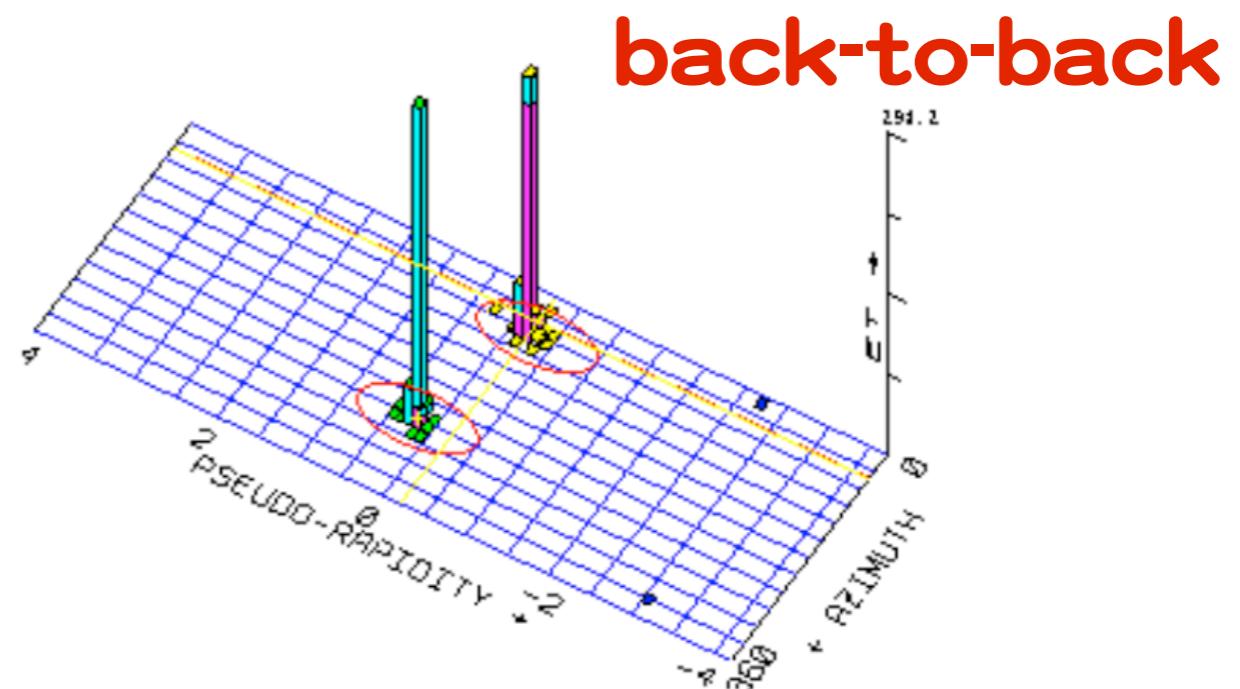
# Jet Modification



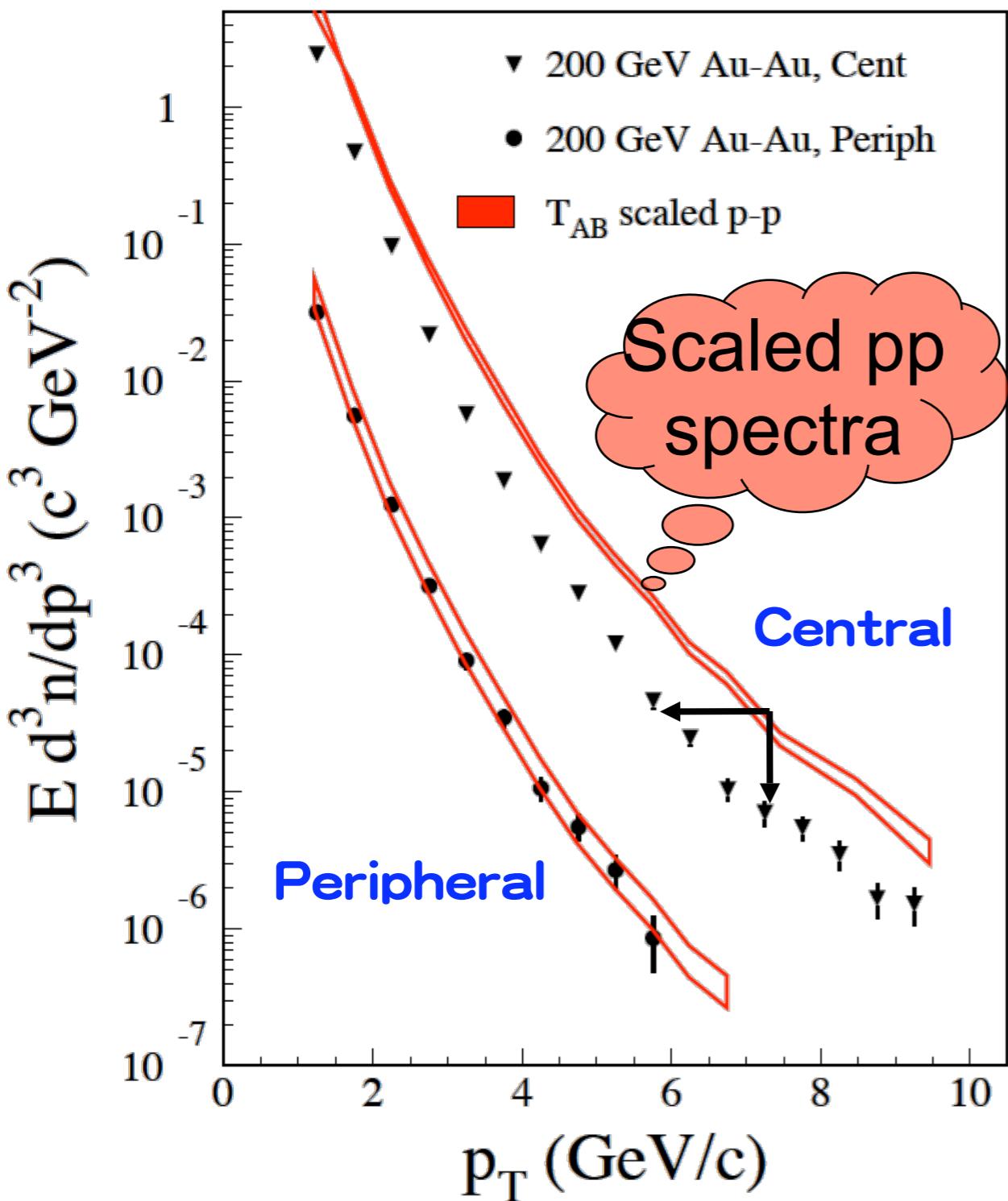
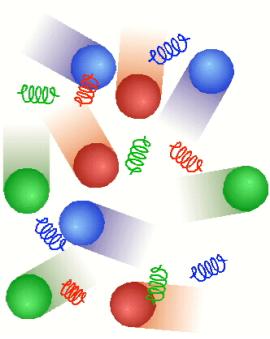
# What is Jet ?



- ✓ At ISR in 1972, deviation from the  **$m_T$  scaling** at high  $p_T$  region is observed as a first time.
- ✓ Binary parton scattering followed by fragmentation produces **back-to-back** jet.
- ✓ Main source of high  $p_T$  particles.

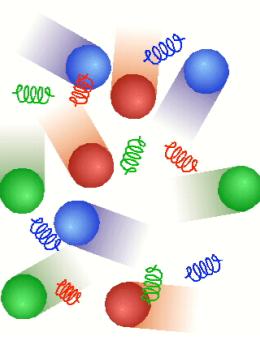


# Comparison of Au+Au and pp



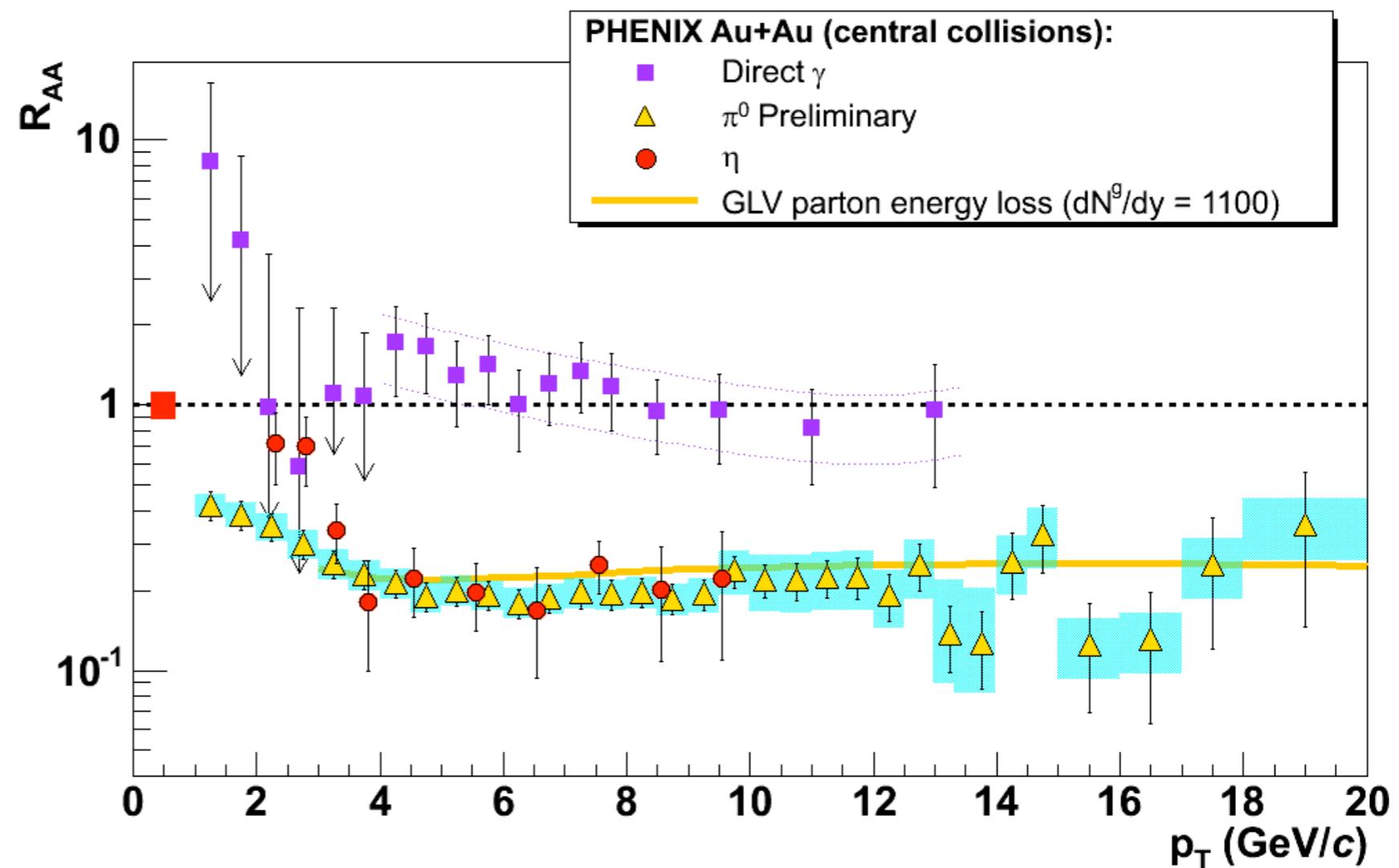
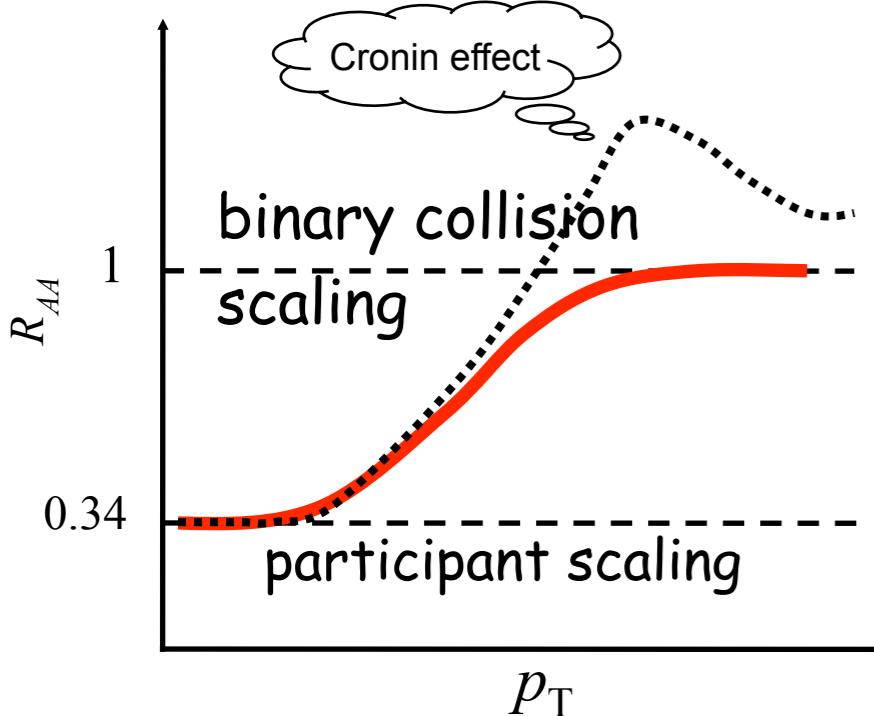
- ✓ For comparison, Au +Au & pp spectra scaled by  $N_{\text{binary}}$ •
- ✓ In peripheral collisions,  
 $\text{Au+Au} \sim \text{pp}$
- ✓ In central collisions,  
 $\text{Au+Au} < \text{pp}$ 
  - █ Suppression of yield ?
  - █ Loss of  $p_T$  ?
  - █ Jet Quench?

# Suppression of high $p_T$ particles



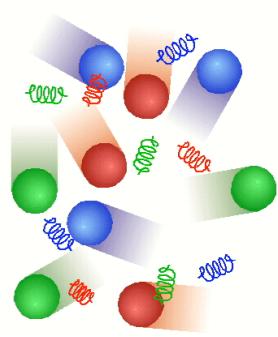
Nuclear  
Modification  
Factor

$$R_{Au+Au} = \frac{dn_{Au+Au}/dp_T dy}{\langle N_{binary} \rangle \cdot dn_{pp}/dp_T dy}$$

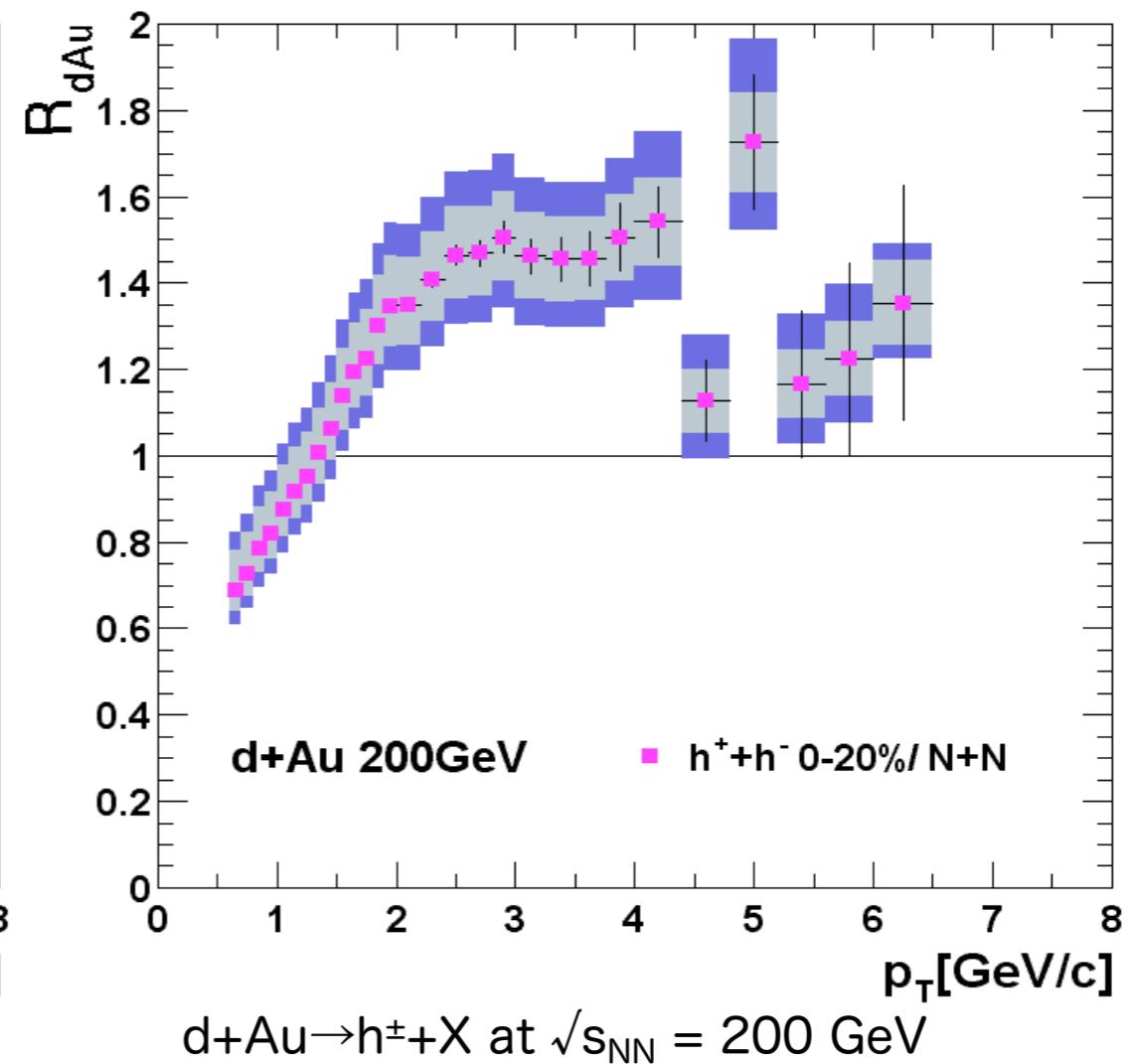
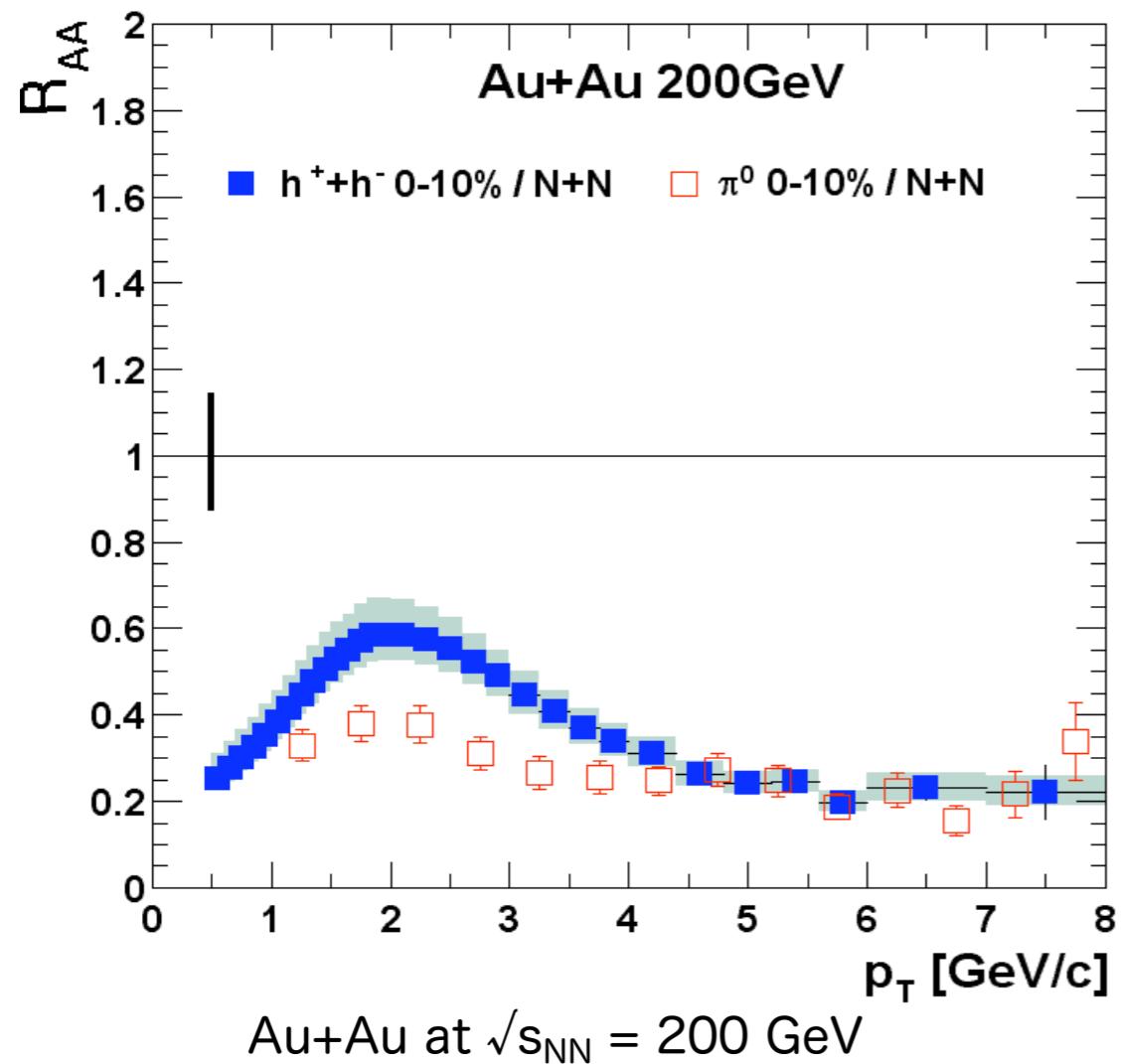


✓ Pions are suppressed, direct photons are not

# Au+Au vs d+Au

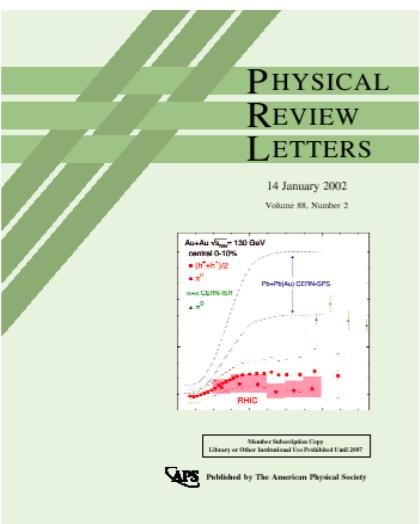


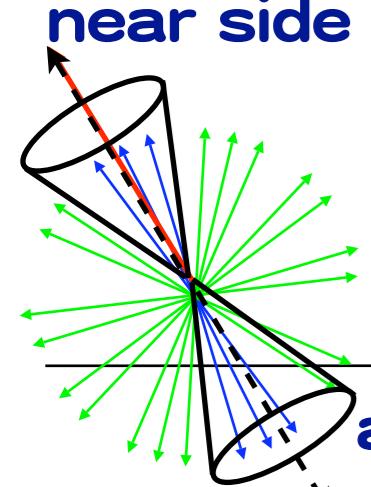
Phenix; P.R.L. 91, 072303 (2003)



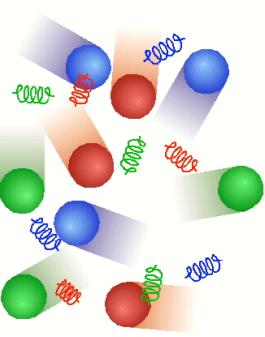
✓ High  $p_T$  suppression in Au+Au, while not observed in d+Au.

→ Effect is not due to initial state, but final state.

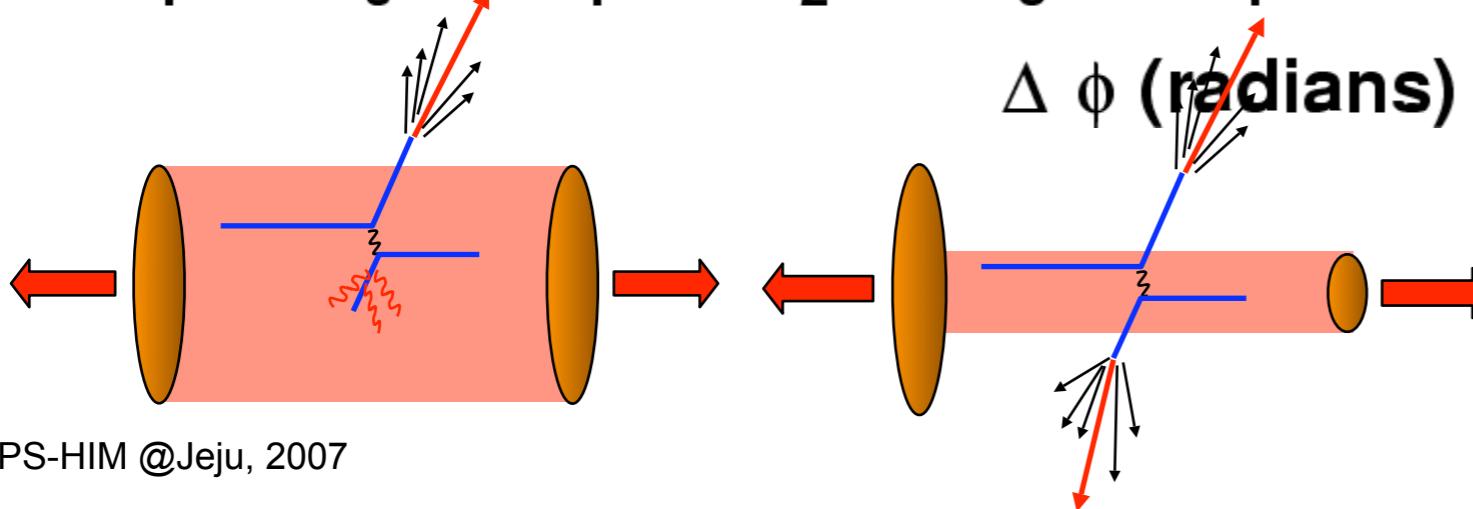
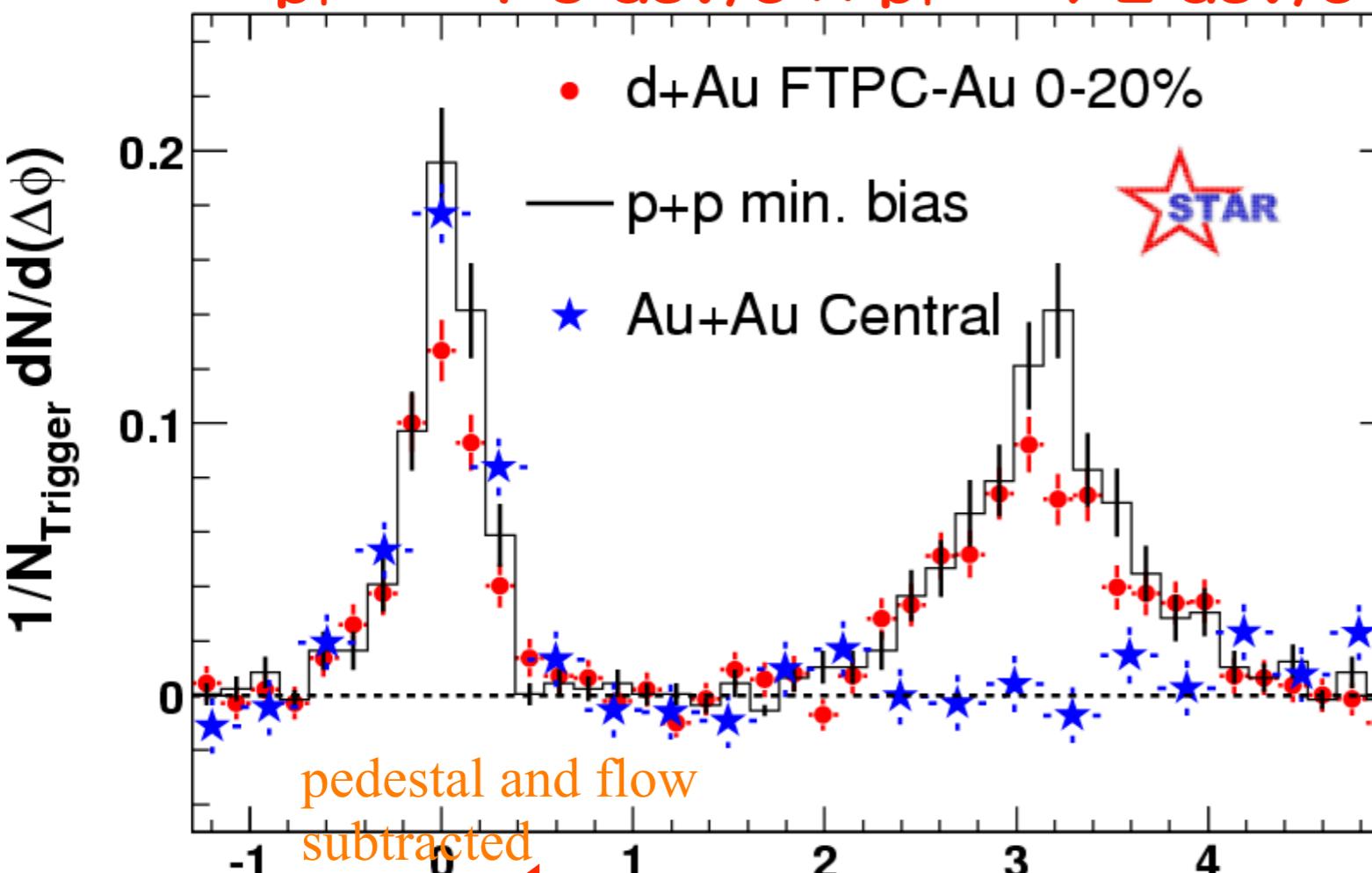




# Disappearance of back-to-back corr.



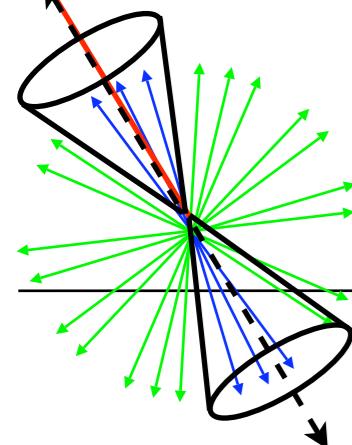
**away side**      Star; P.R.L. 91, 72304 (2003)  
 $p_T^{\text{trig}} = 4\text{--}6 \text{ GeV}/c \times p_T^{\text{assoc}} > 2 \text{ GeV}/c$



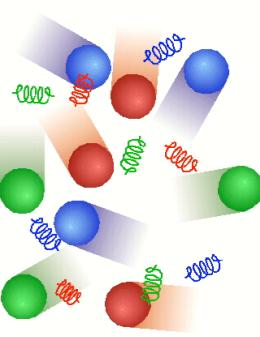
KPS-HIM @Jeju, 2007

- ✓ Direct evidence of loss of ‘jet’
- ✓ Azimuthal correlation w.r.t. high pt leading particle (trigger).
  - | pp ; clean di-jet
  - | dAu; similar to pp
  - | Au+Au; Similar on the same side (suggesting jet-like mechanism), but b-to-b disappeared
    - | Effect is not in initial but in final stage
    - | Energy loss of partons in dense matter created in Au+Au

near side

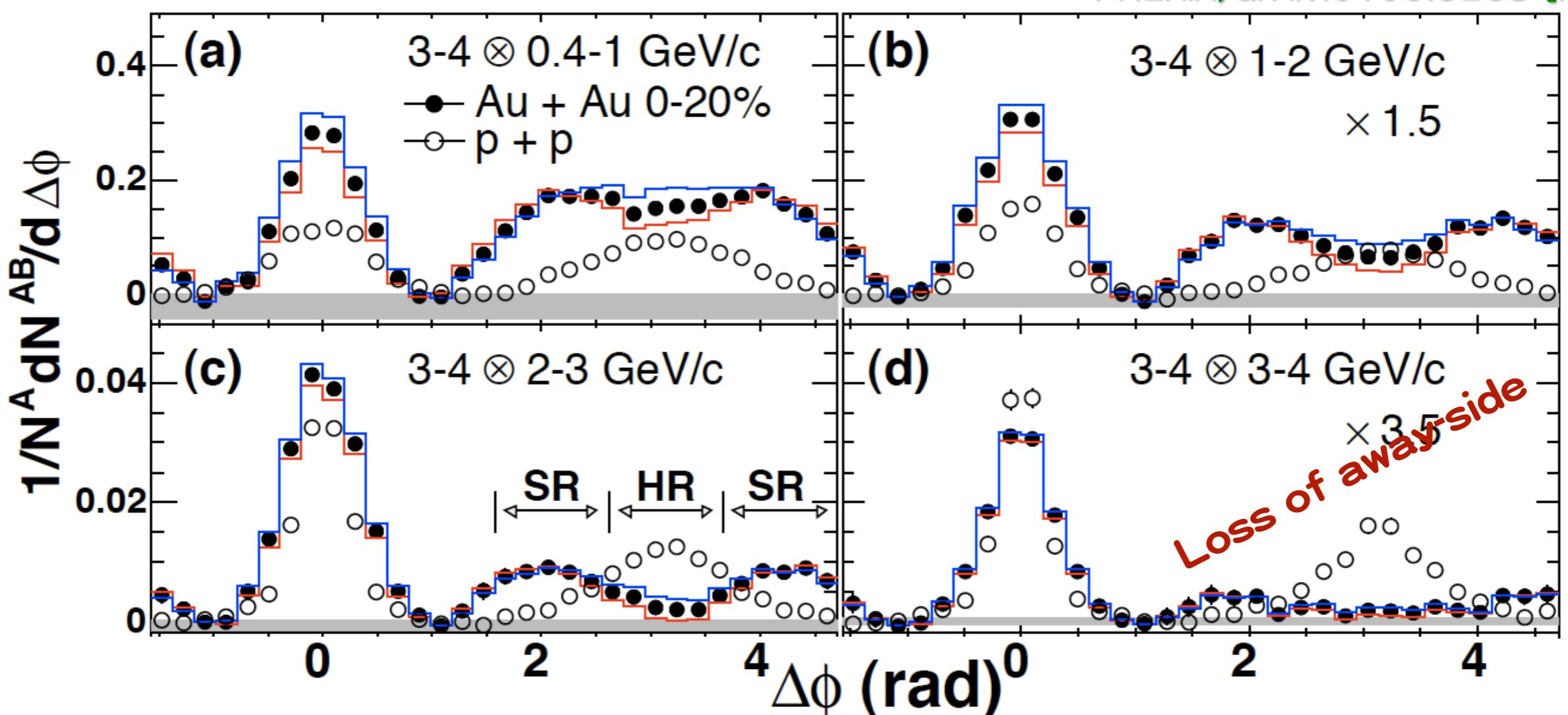


# Shape change of away-side



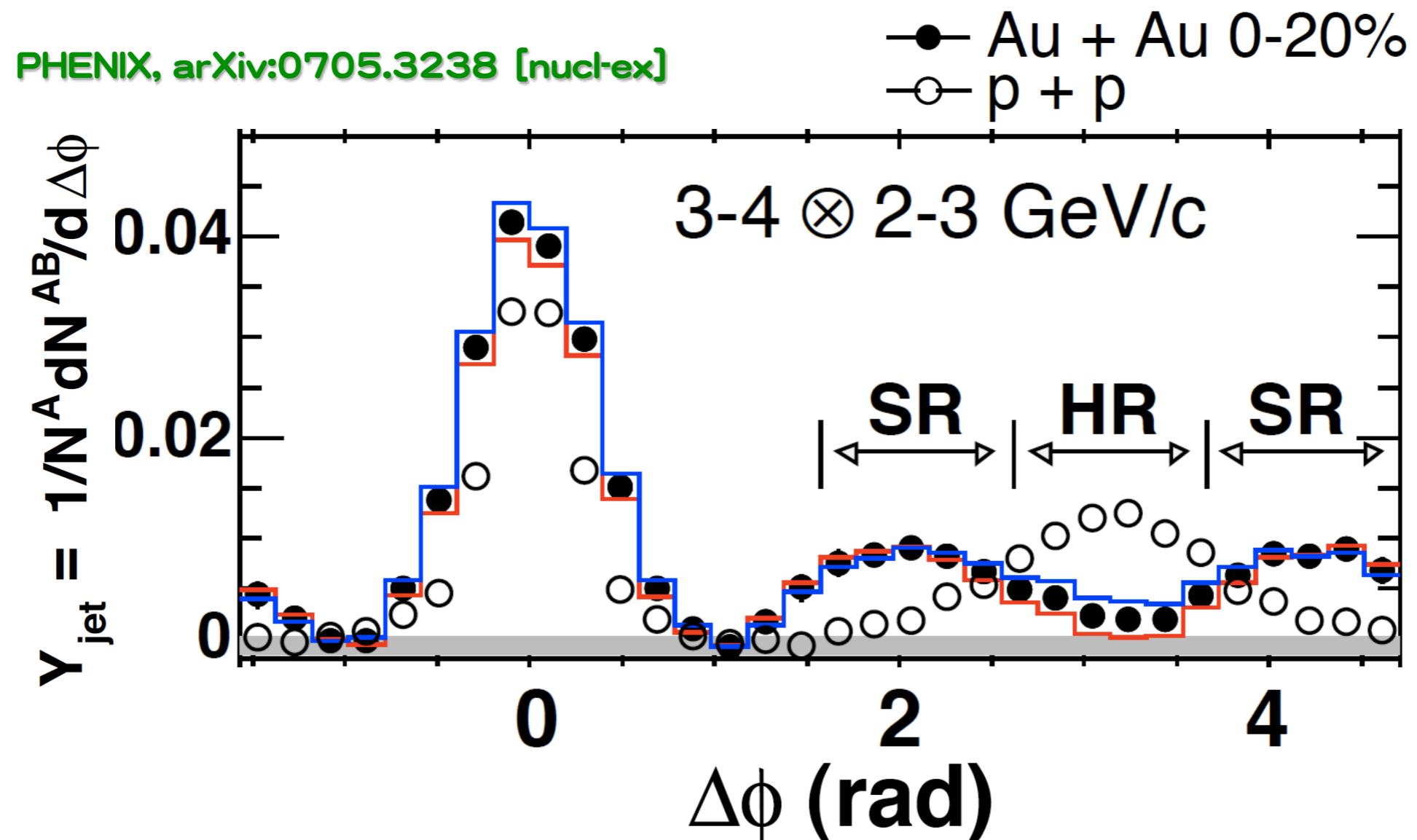
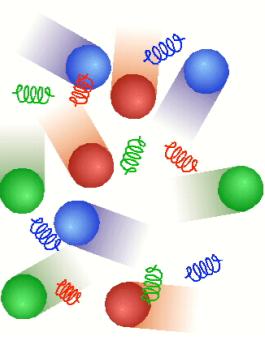
$$p_T^{\text{trig}} = 3\text{--}4 \text{ GeV}/c \times p_T^{\text{assoc}}$$

PHENIX, arXiv:0705.3238 [nucl-ex]



✓ From broad/none to distinct two shoulders  
at  $\Delta\Phi = \pi \pm 1.1$  with decreasing momentum.

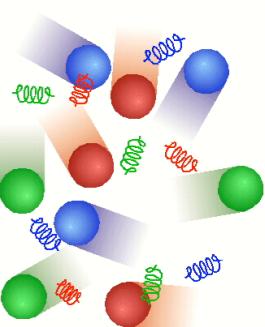
# Shoulders at $\Delta\phi = \pi \pm 1.1$



✓ Location &  $\langle pt \rangle$  of shoulder seem to be independent of centrality and  $pt$ .

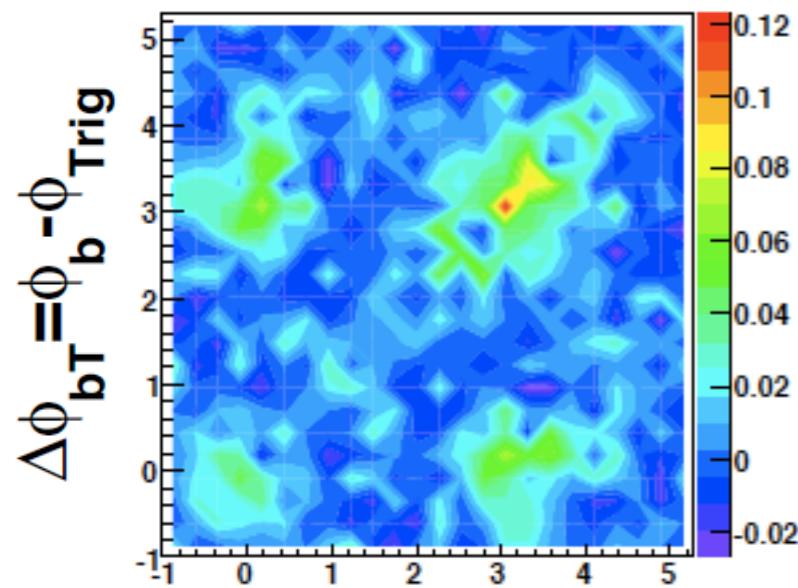
◆ Not cherenkov, not deflection, but, Shock wave !?

# Sophisticated 3-particle correlation analysis

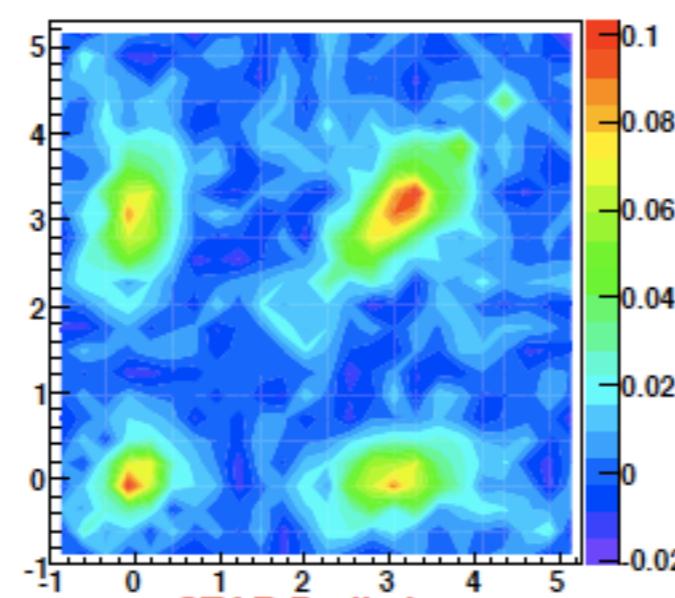


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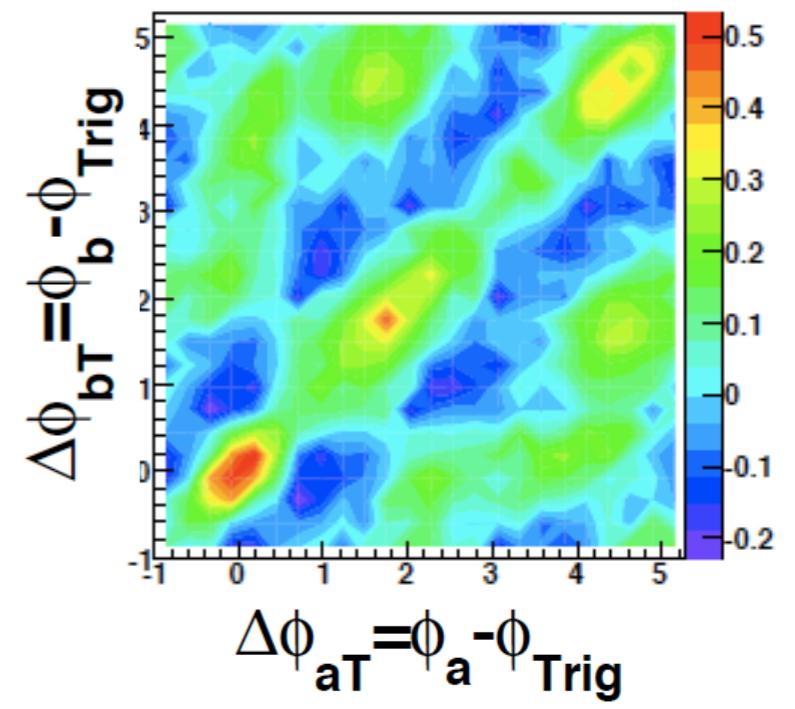
pp



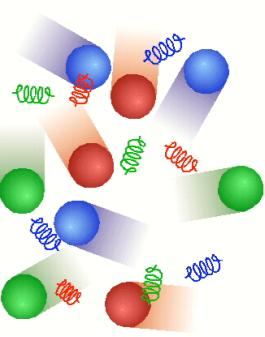
dAu



AuAu(cent)

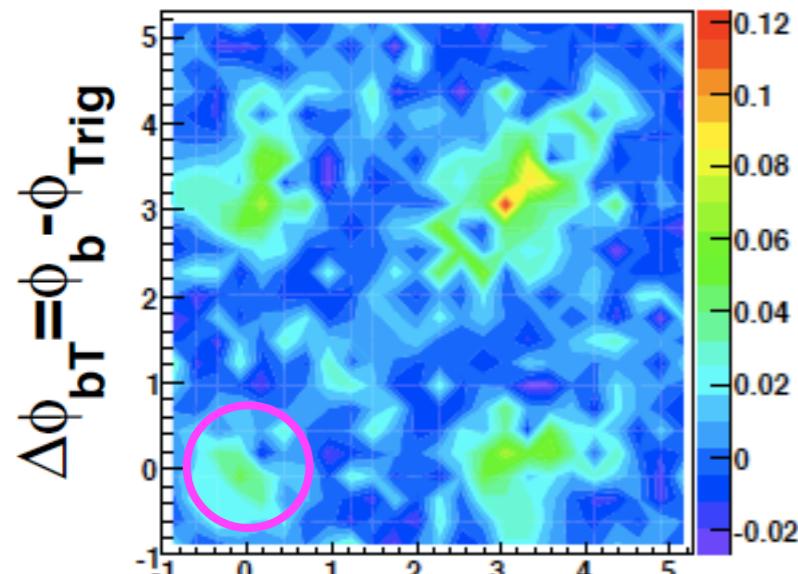


# Sophisticated 3-particle correlation analysis

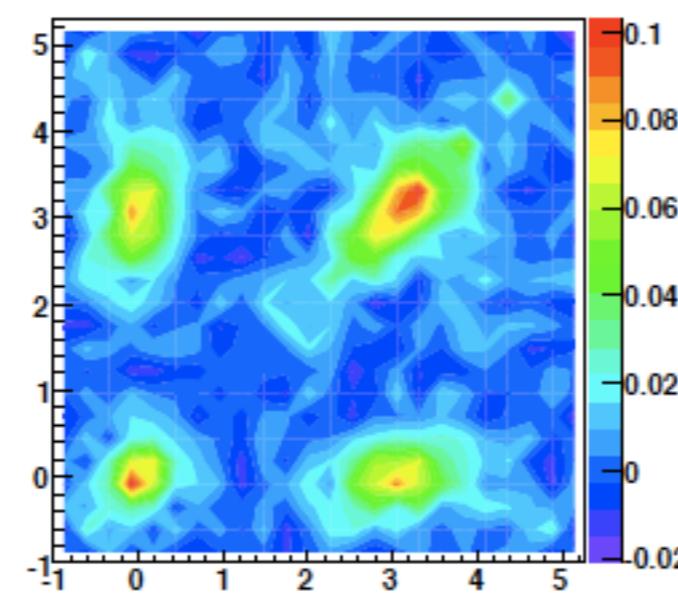


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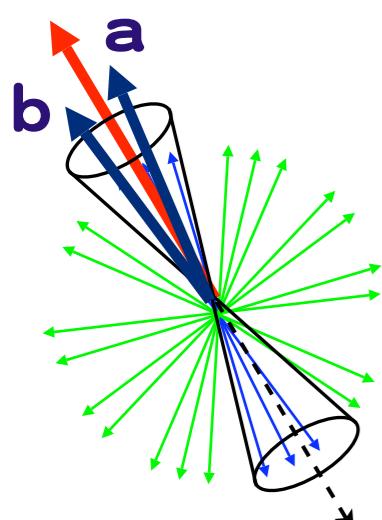
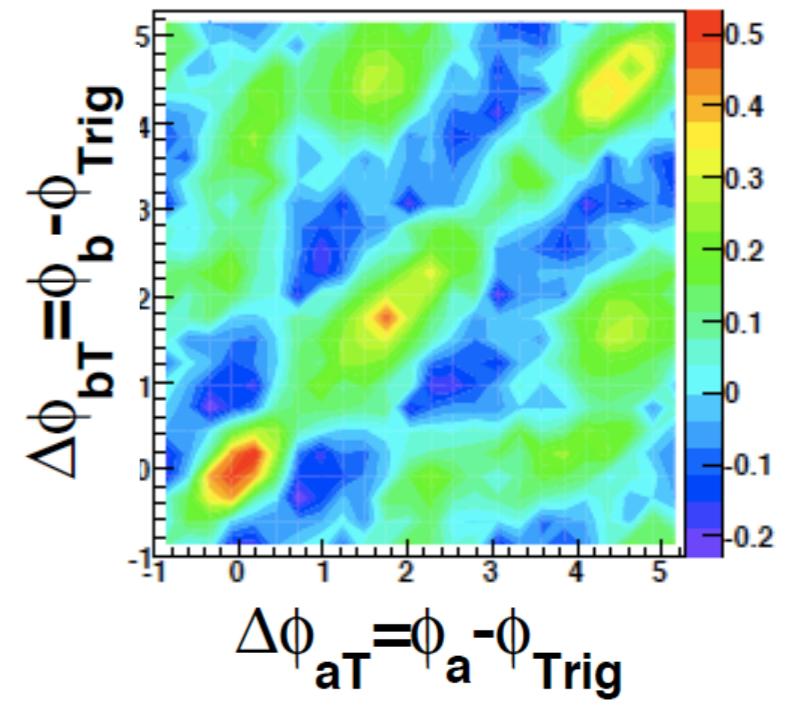
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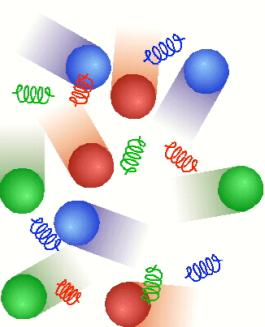
dAu



AuAu(cent)

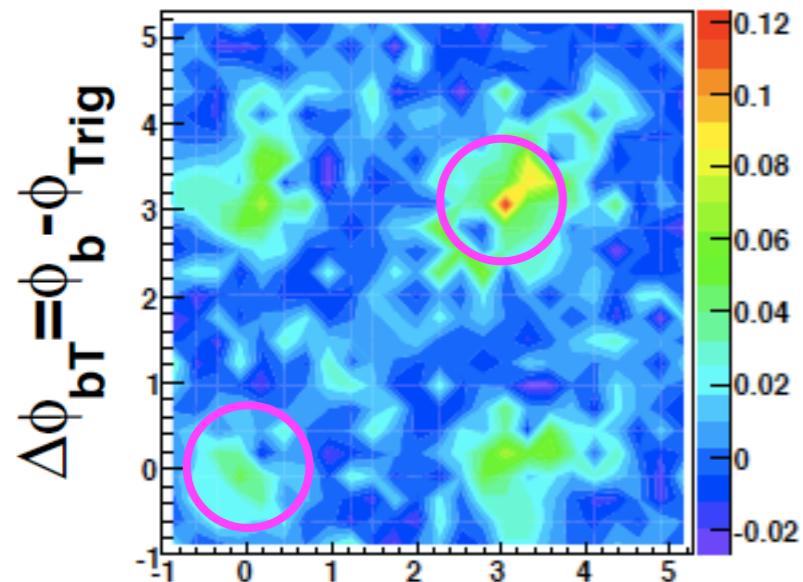


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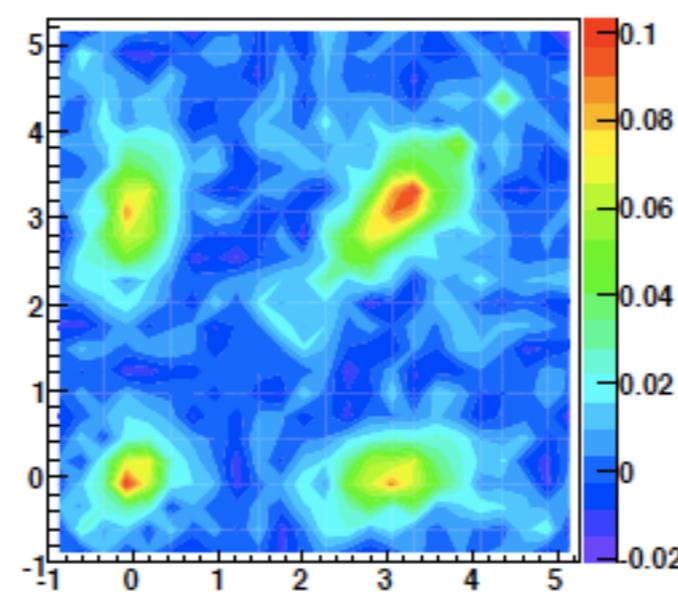


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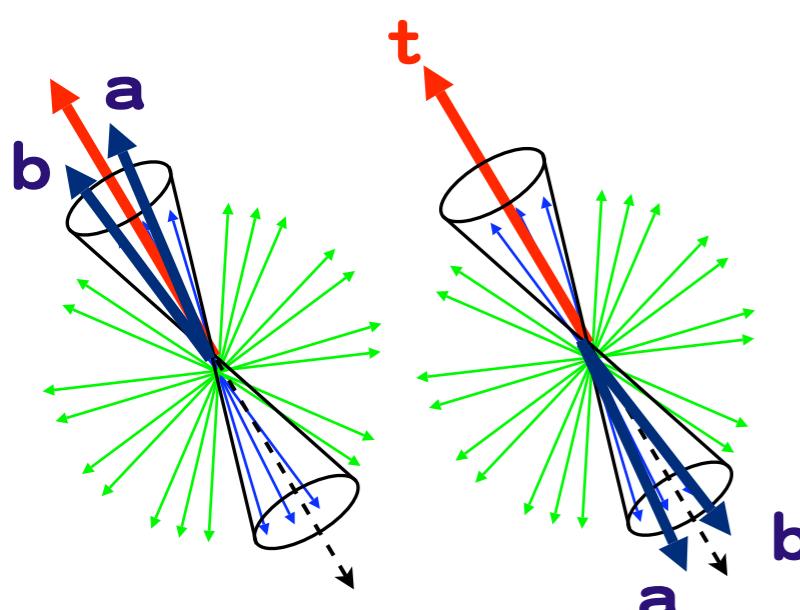
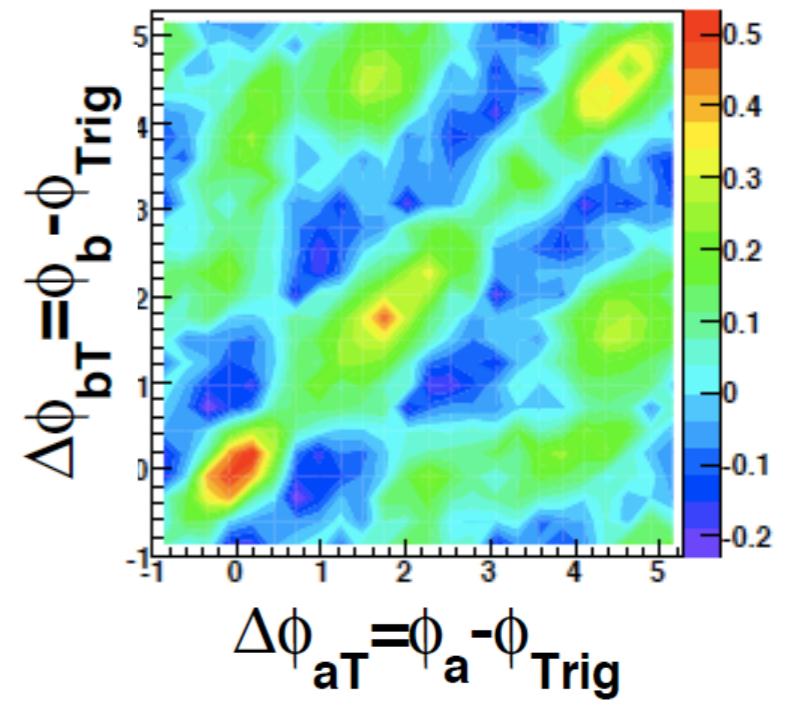
pp



dAu

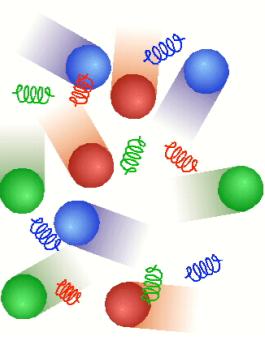


AuAu(cent)



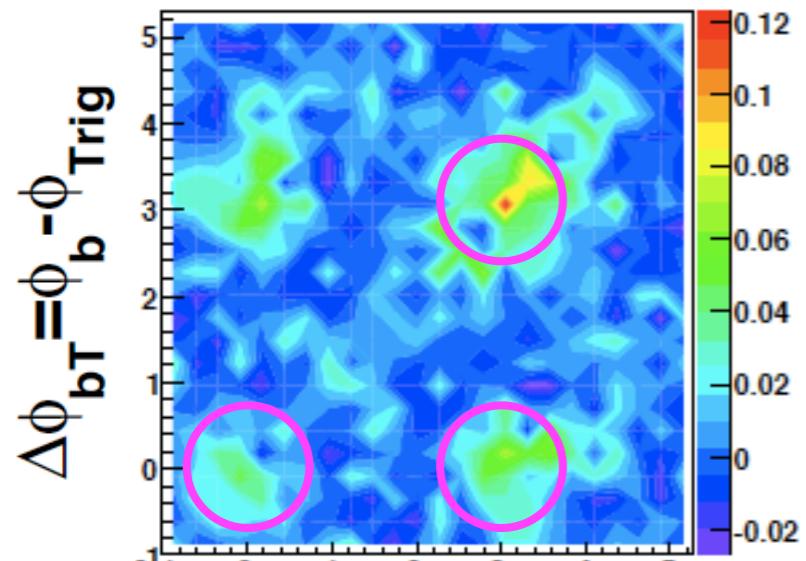
KPS-HIM @Jeju, 2007

# Sophisticated 3-particle correlation analysis

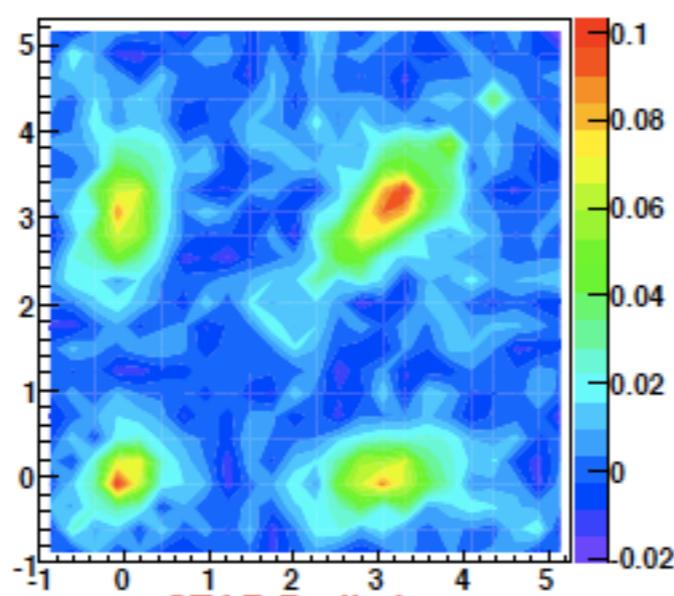


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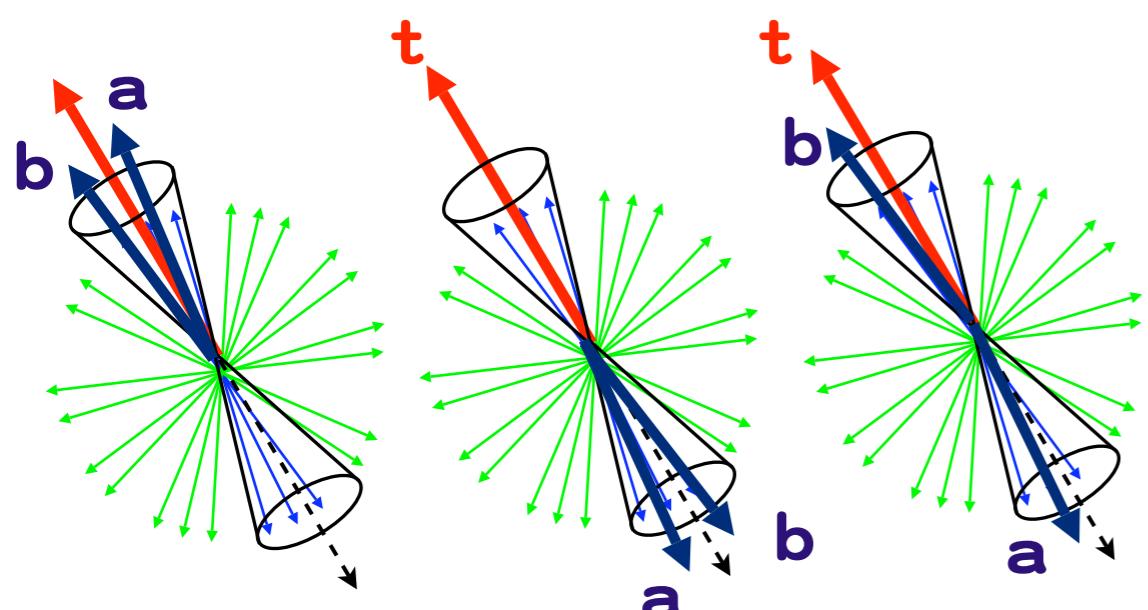
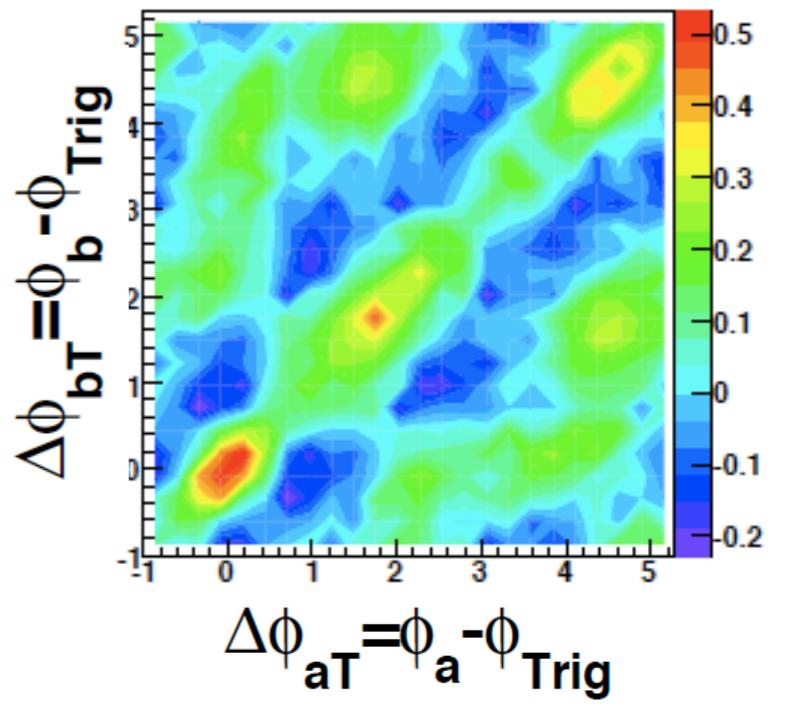
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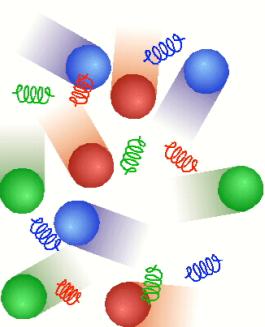
dAu



AuAu(cent)

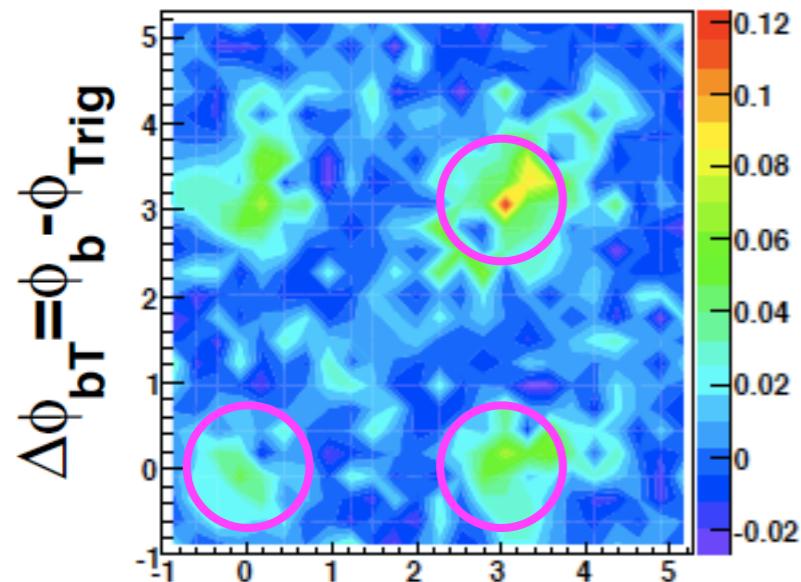


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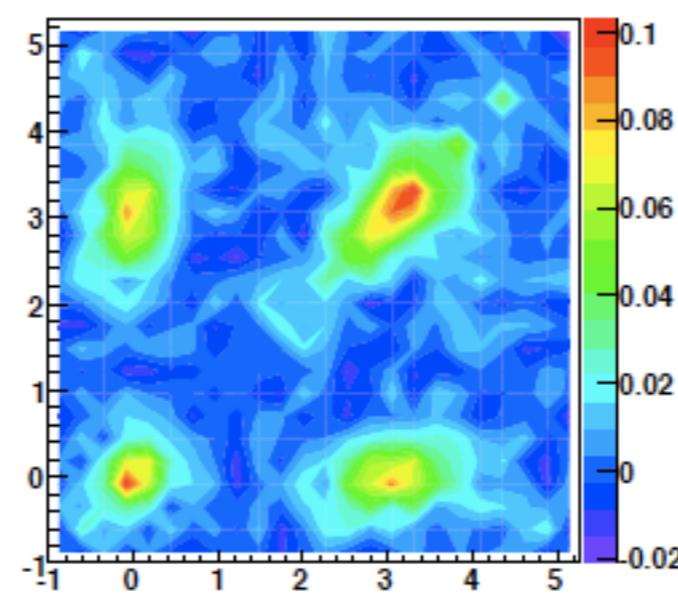


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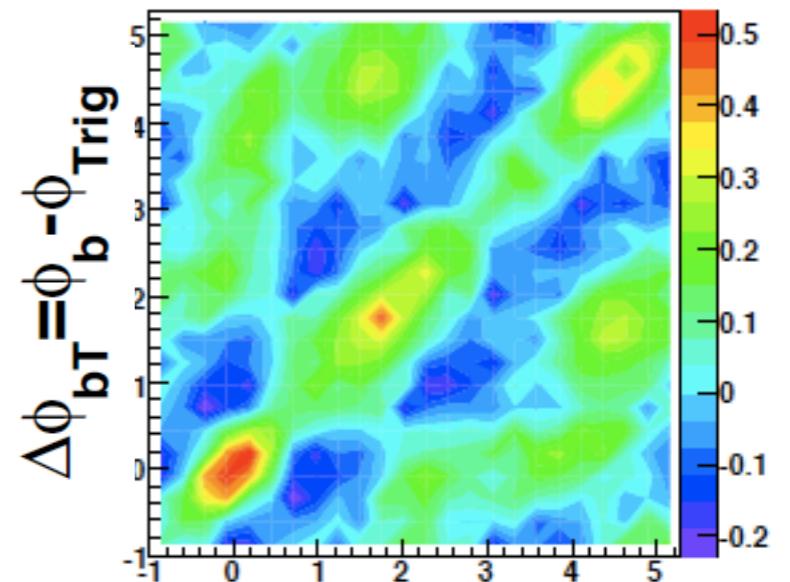
**pp**



**dAu**

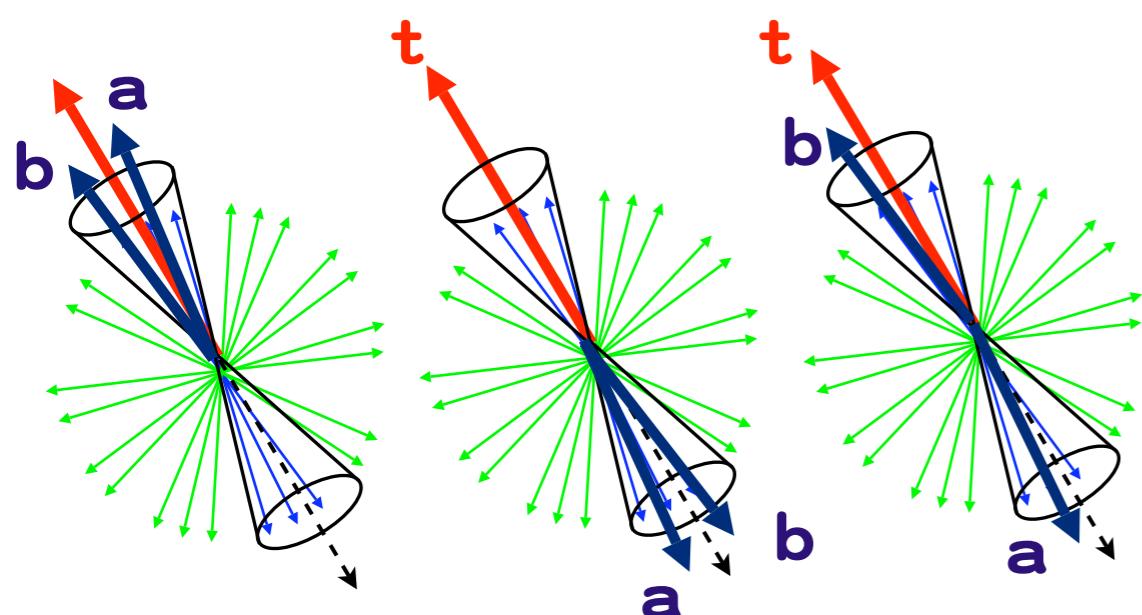


**AuAu(cent)**

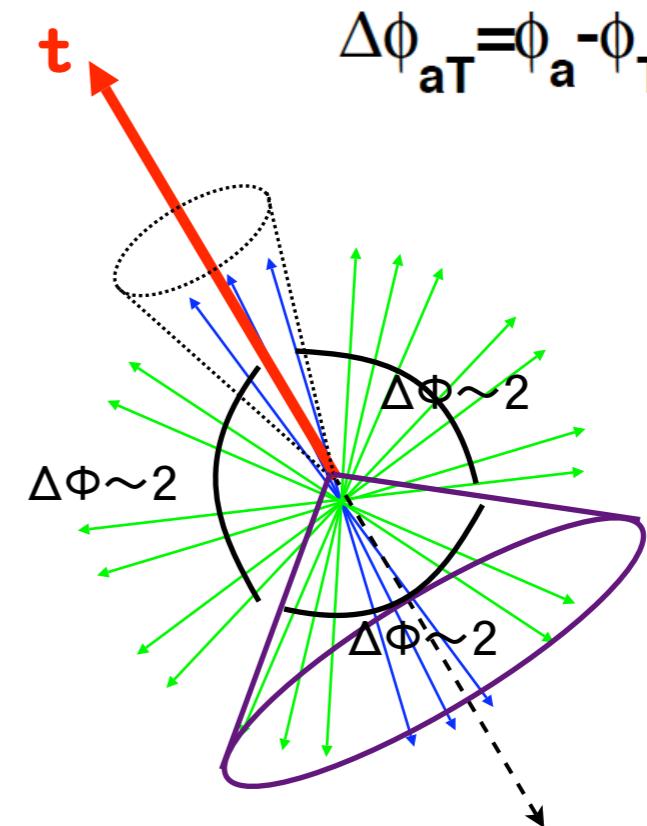


$\Delta\phi_{aT} = \phi_a - \phi_{\text{Trig}}$

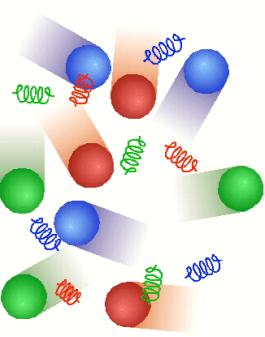
$\Delta\phi_{aT} = \phi_a - \phi_{\text{Trig}}$



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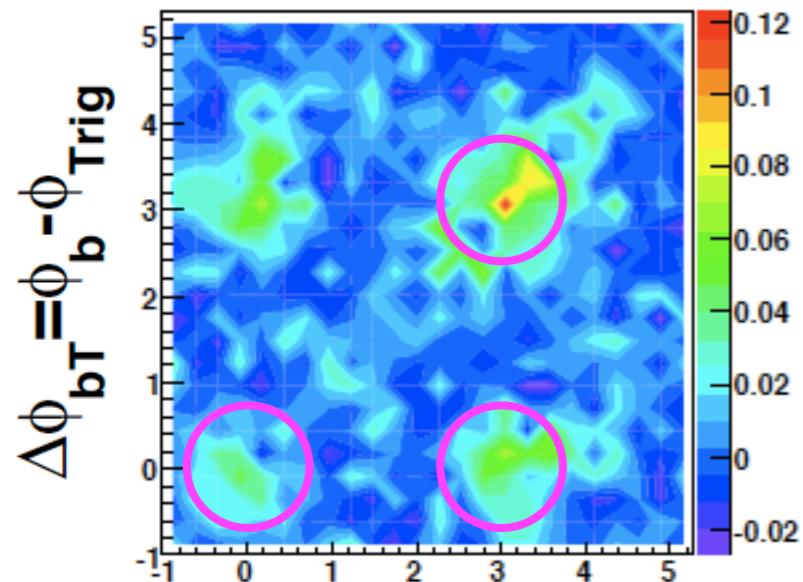


# Sophisticated 3-particle correlation analysis

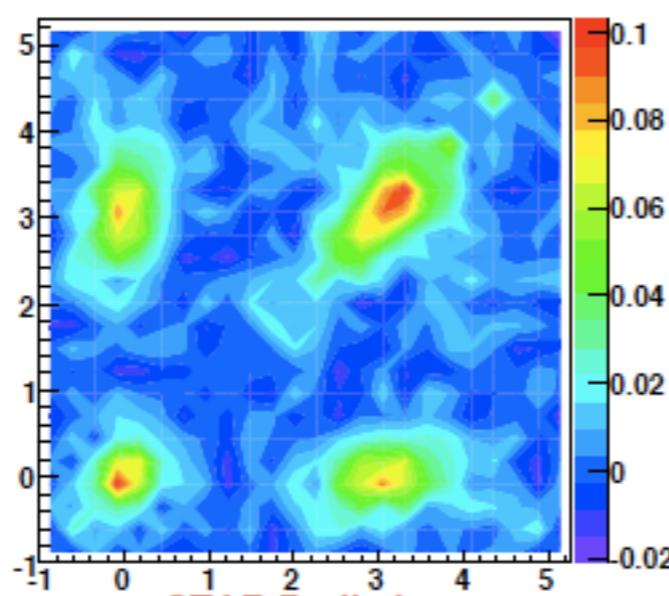


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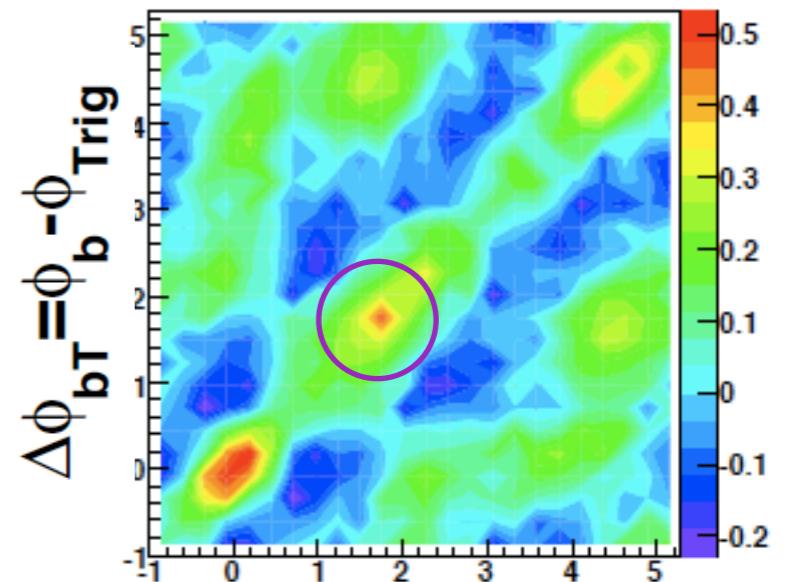
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**dAu**

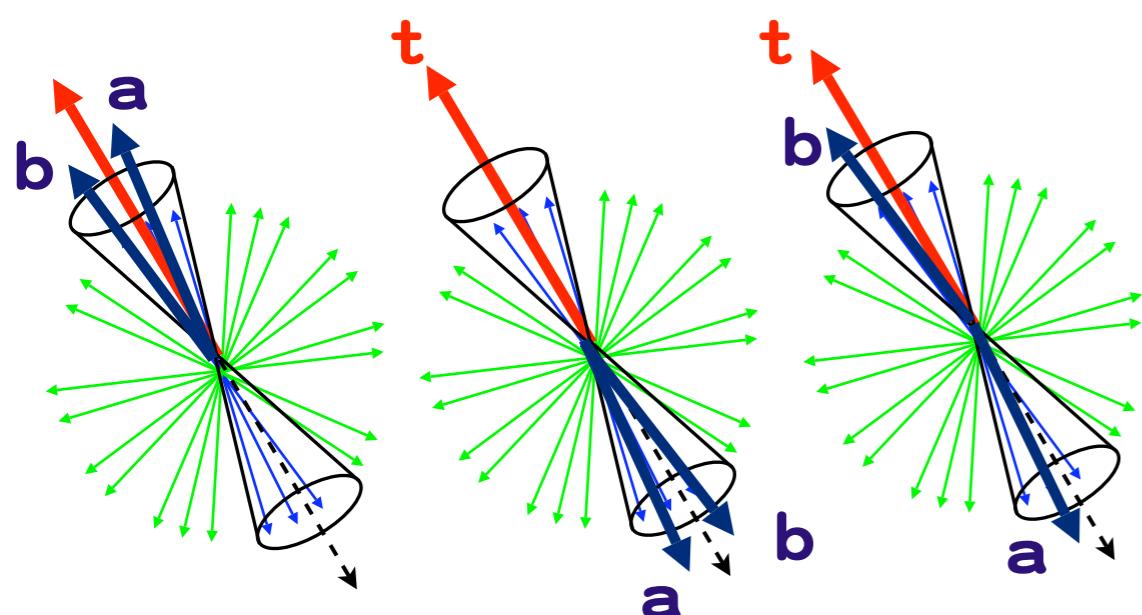


**AuAu(cent)**

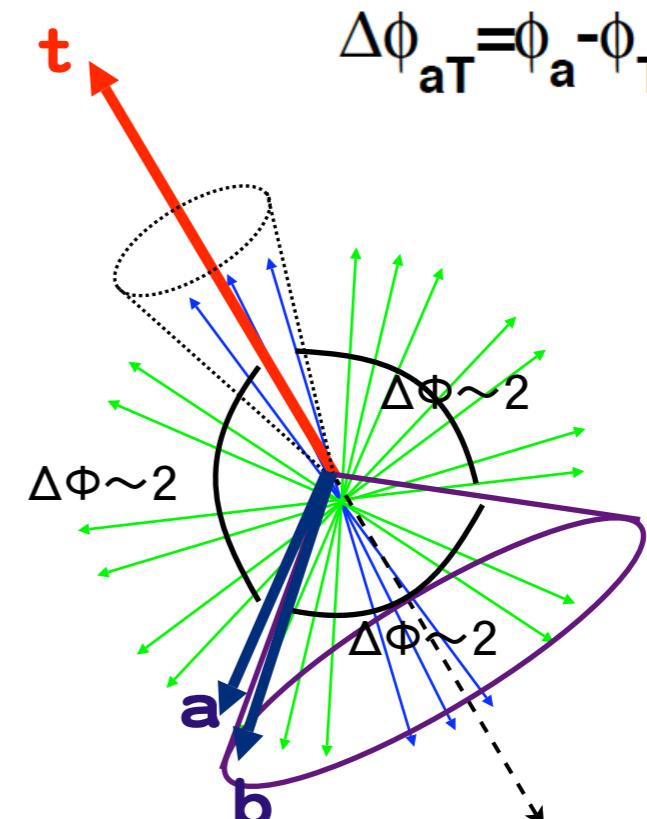


$\Delta\phi_{aT} = \phi_a - \phi_{\text{Trig}}$

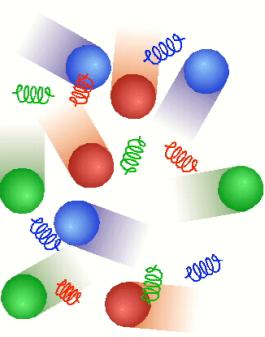
$\Delta\phi_{aT} = \phi_a - \phi_{\text{Trig}}$



KPS-HIM @Jeju, 2007

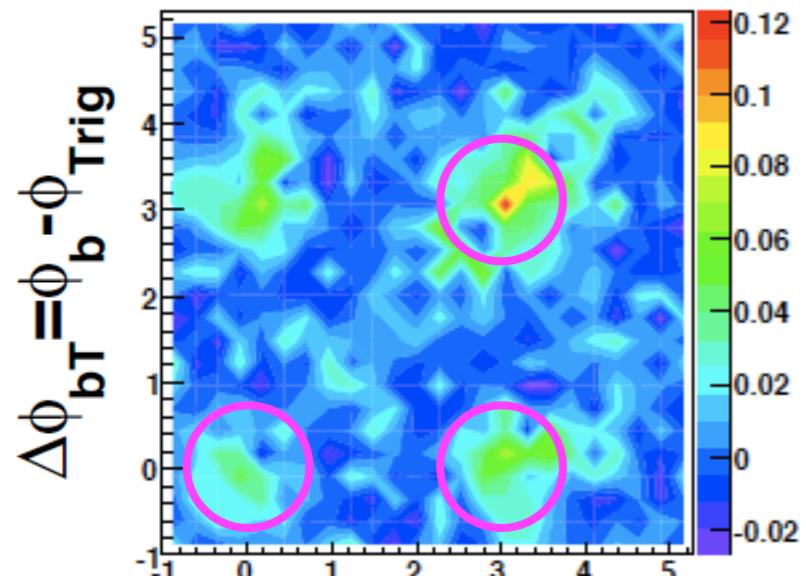


# Sophisticated 3-particle correlation analysis

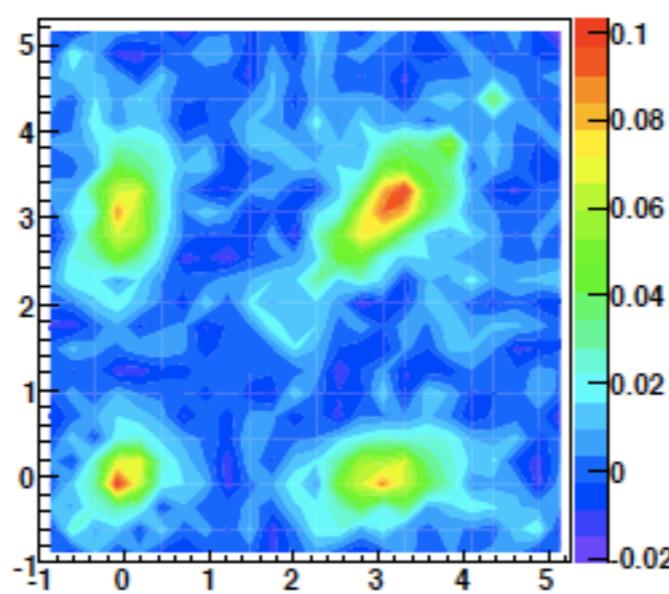


STAR@QM06, nucl-ex/7030 10

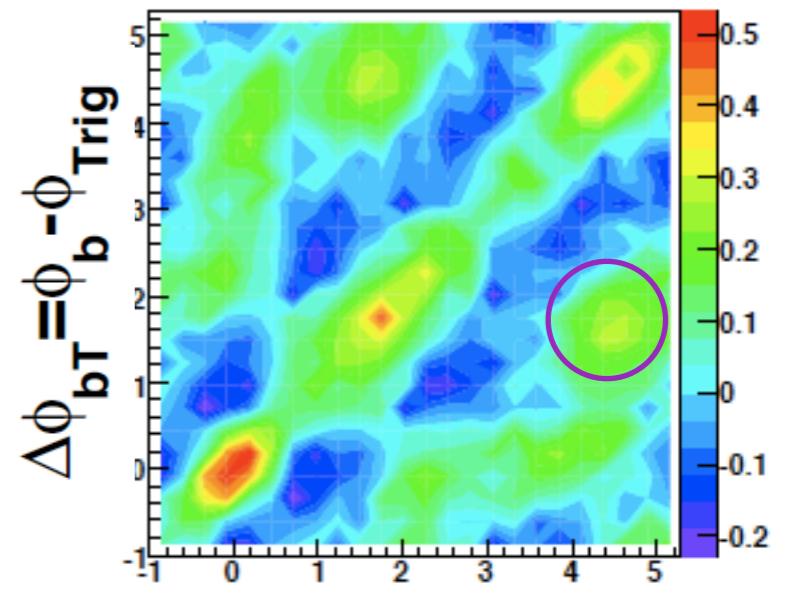
**pp**



**dAu**

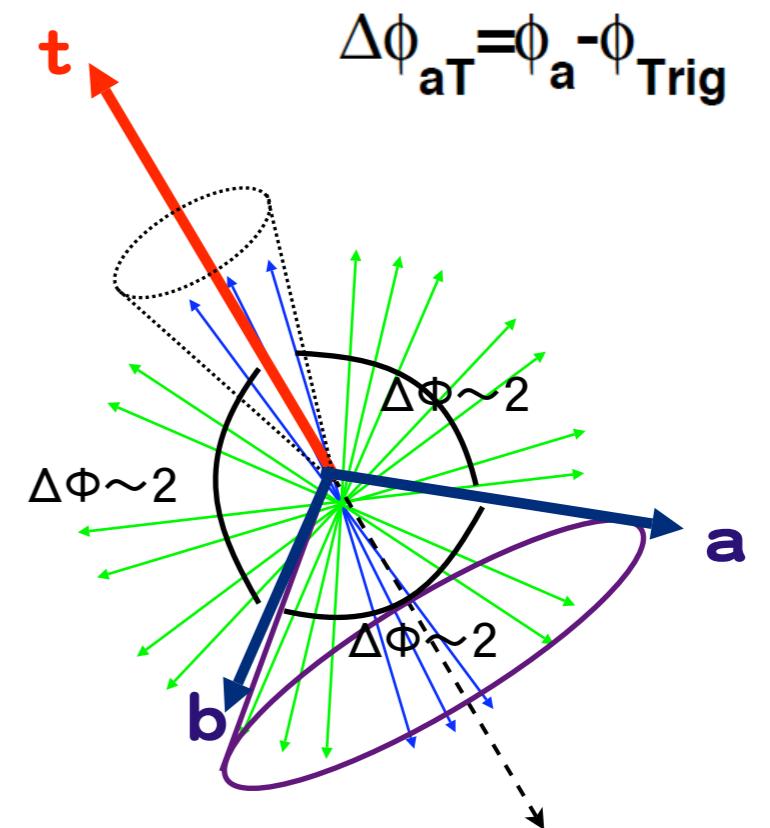
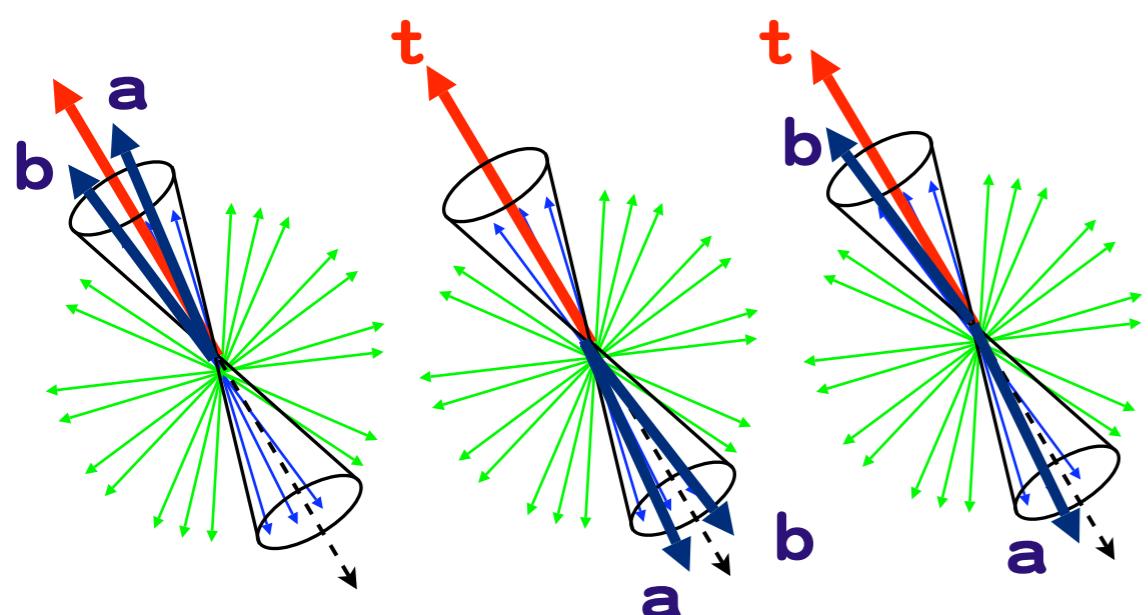


**AuAu(cent)**

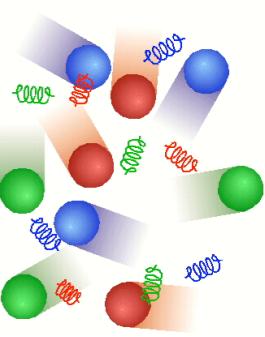


$\Delta\phi_{aT} = \phi_a - \phi_{\text{Trig}}$

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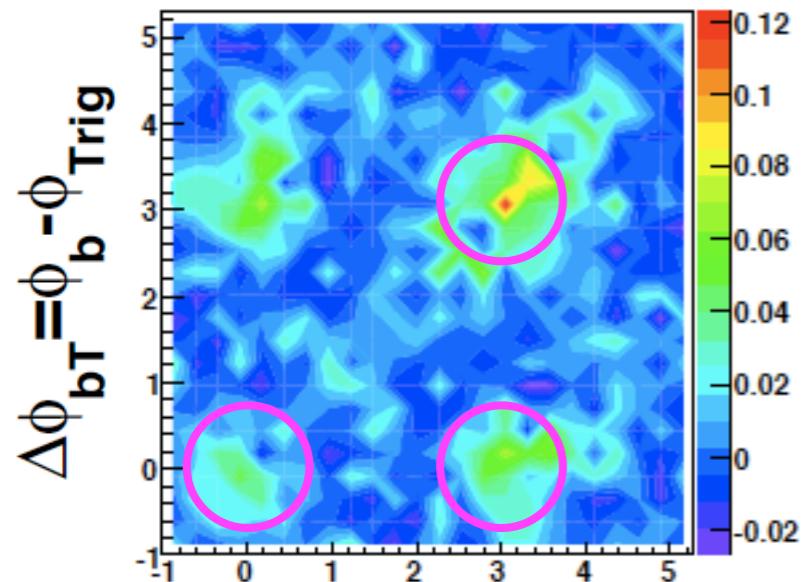


# Sophisticated 3-particle correlation analysis

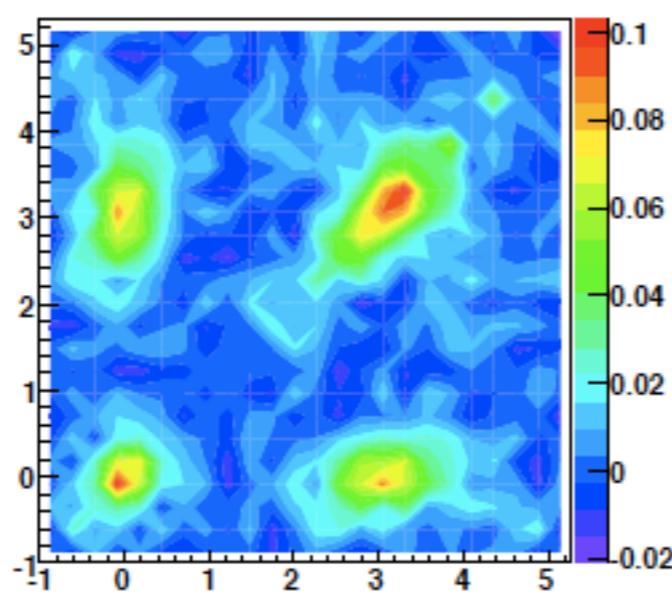


STAR@QM06, nucl-ex/7030 10

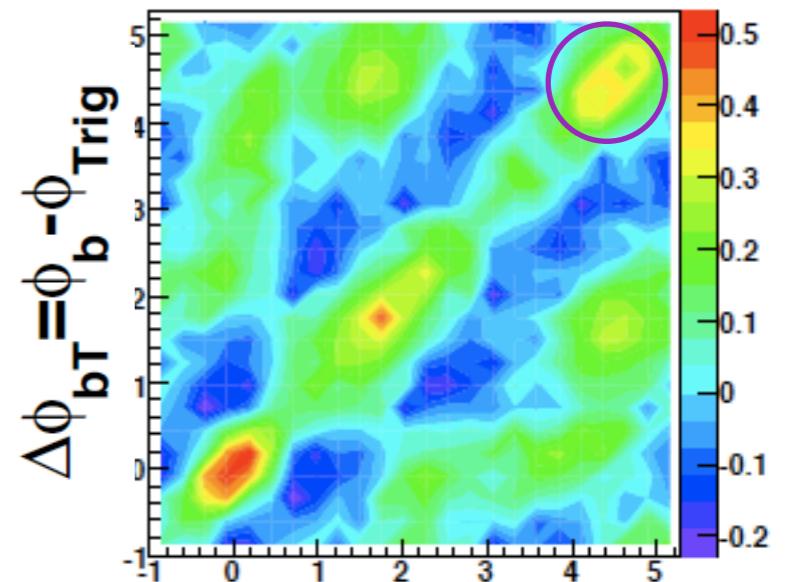
**pp**



**dAu**

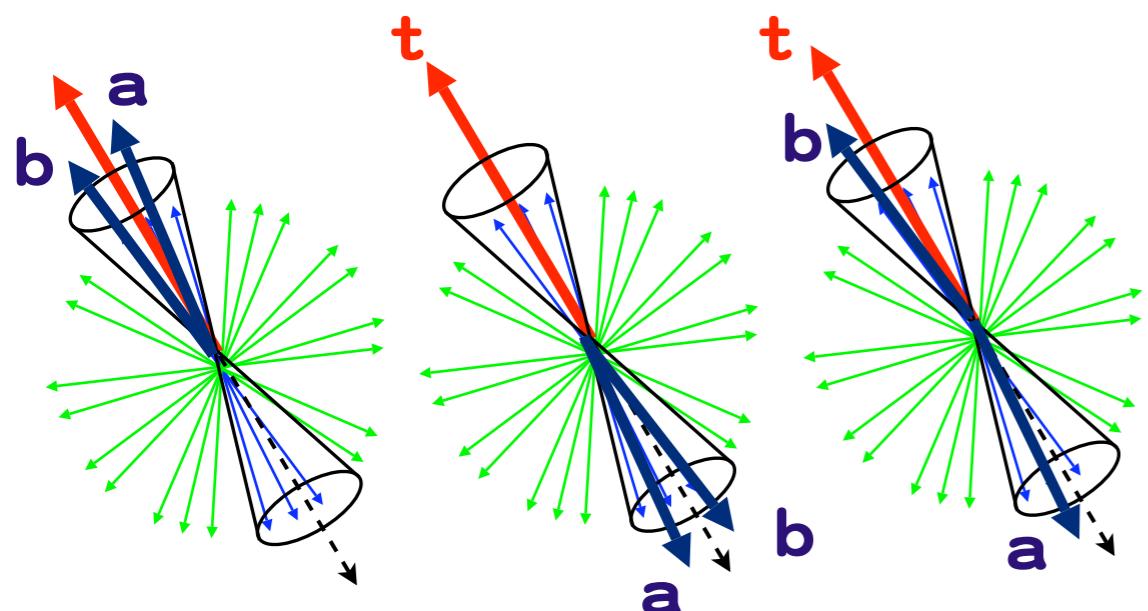


**AuAu(cent)**

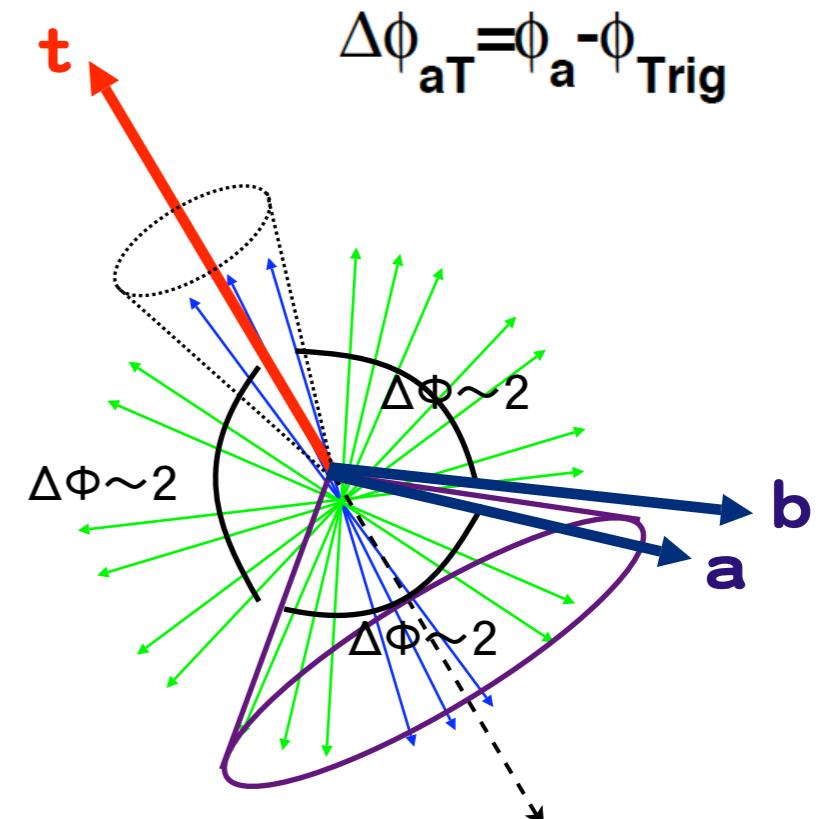


$\Delta\phi_{aT} = \phi_a - \phi_{\text{Trig}}$

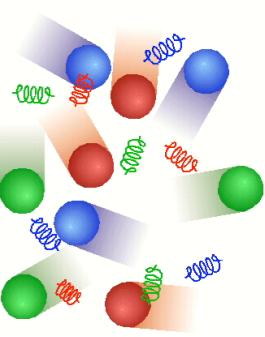
$\Delta\phi_{aT} = \phi_a - \phi_{\text{Trig}}$



KPS-HIM @Jeju, 2007

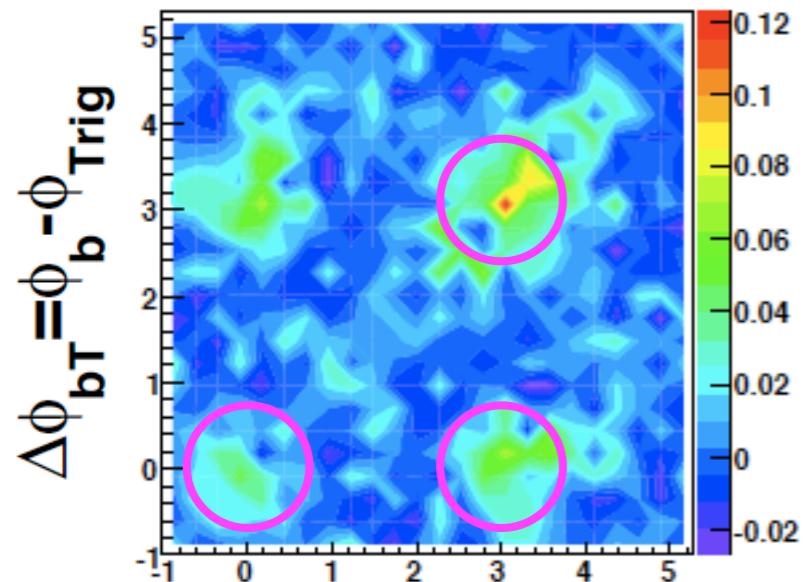


# Sophisticated 3-particle correlation analysis

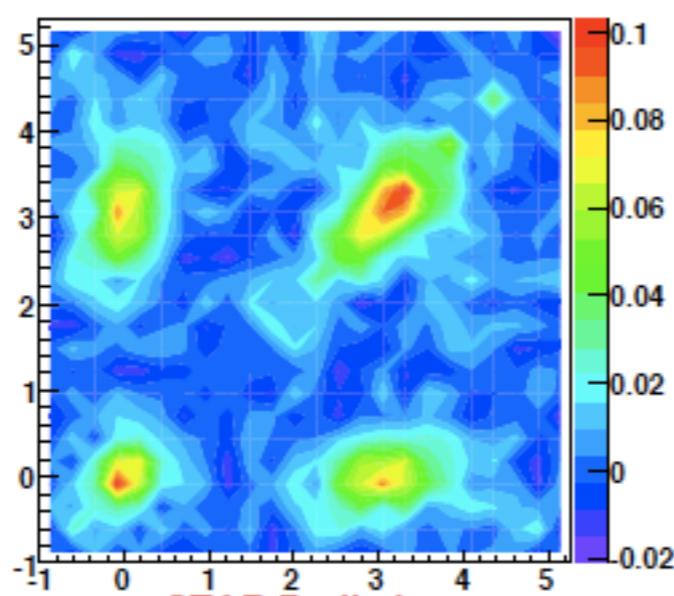


STAR@QM06, nucl-ex/7030 10

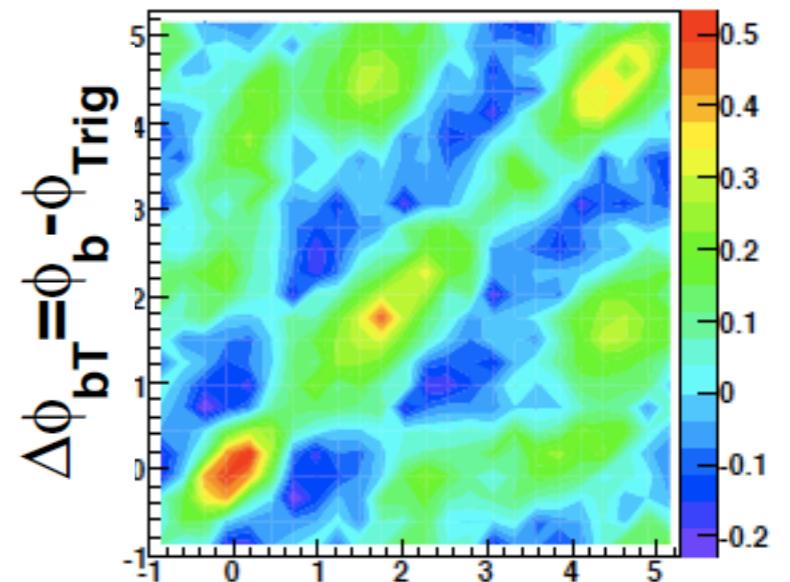
**pp**



**dAu**

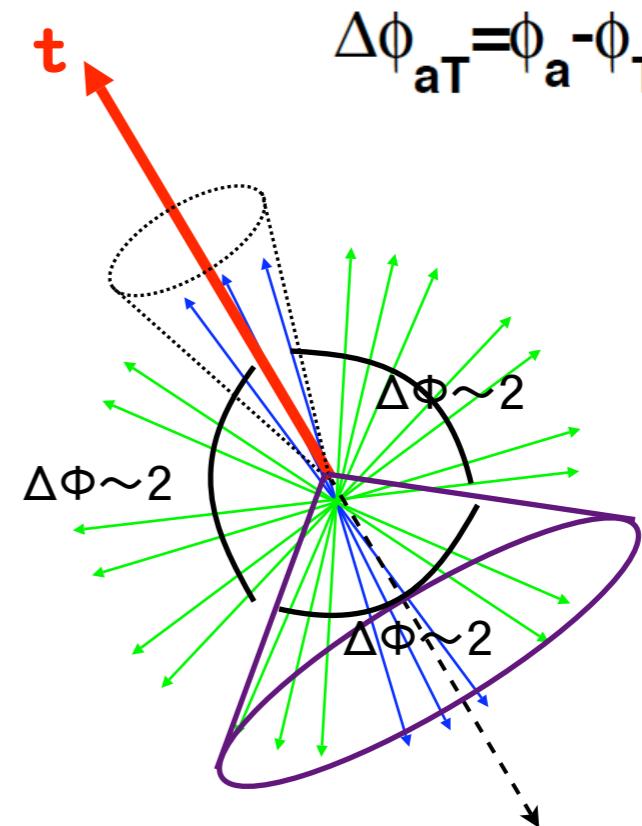
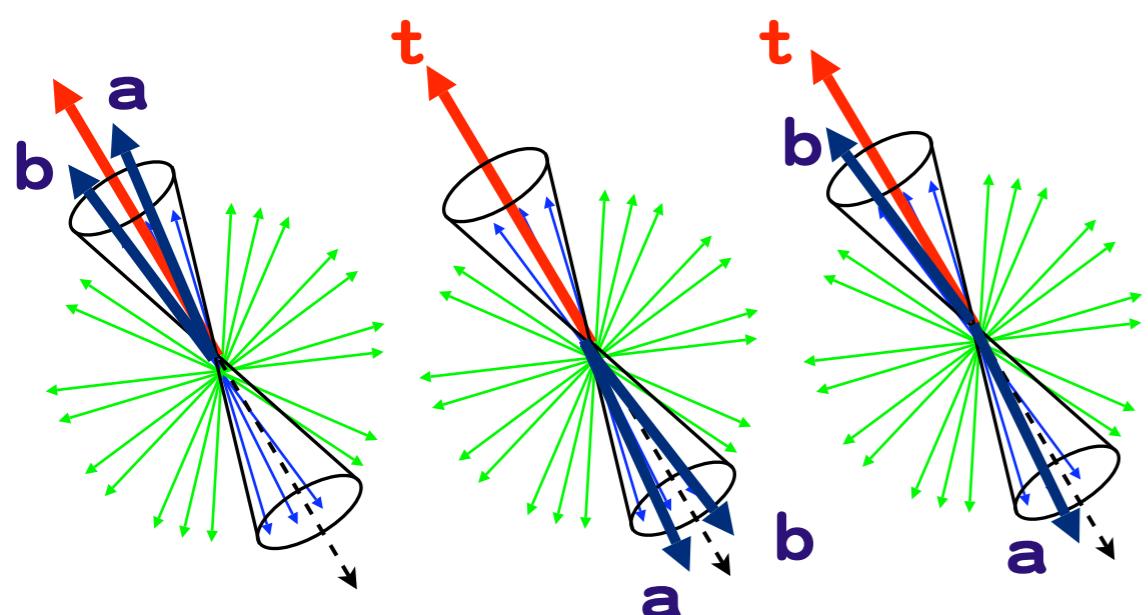


**AuAu(cent)**

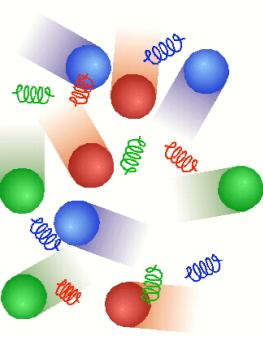


$\Delta\phi_{aT} = \phi_a - \phi_{\text{Trig}}$

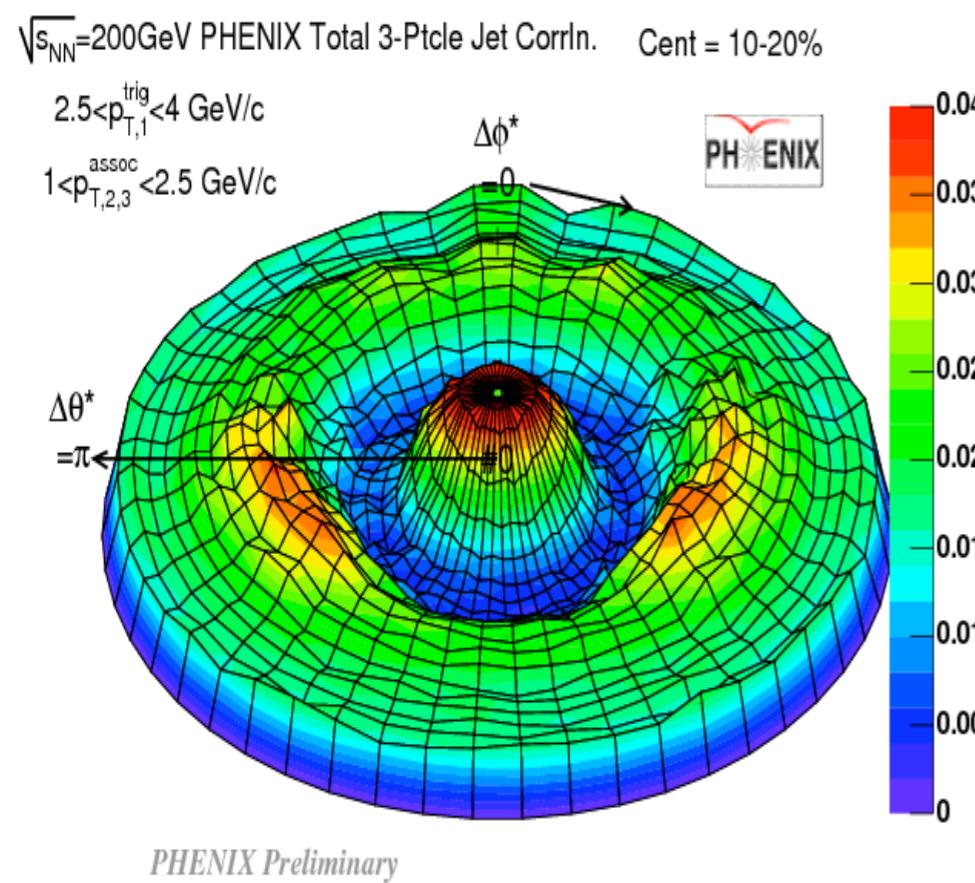
$\Delta\phi_{aT} = \phi_a - \phi_{\text{Trig}}$



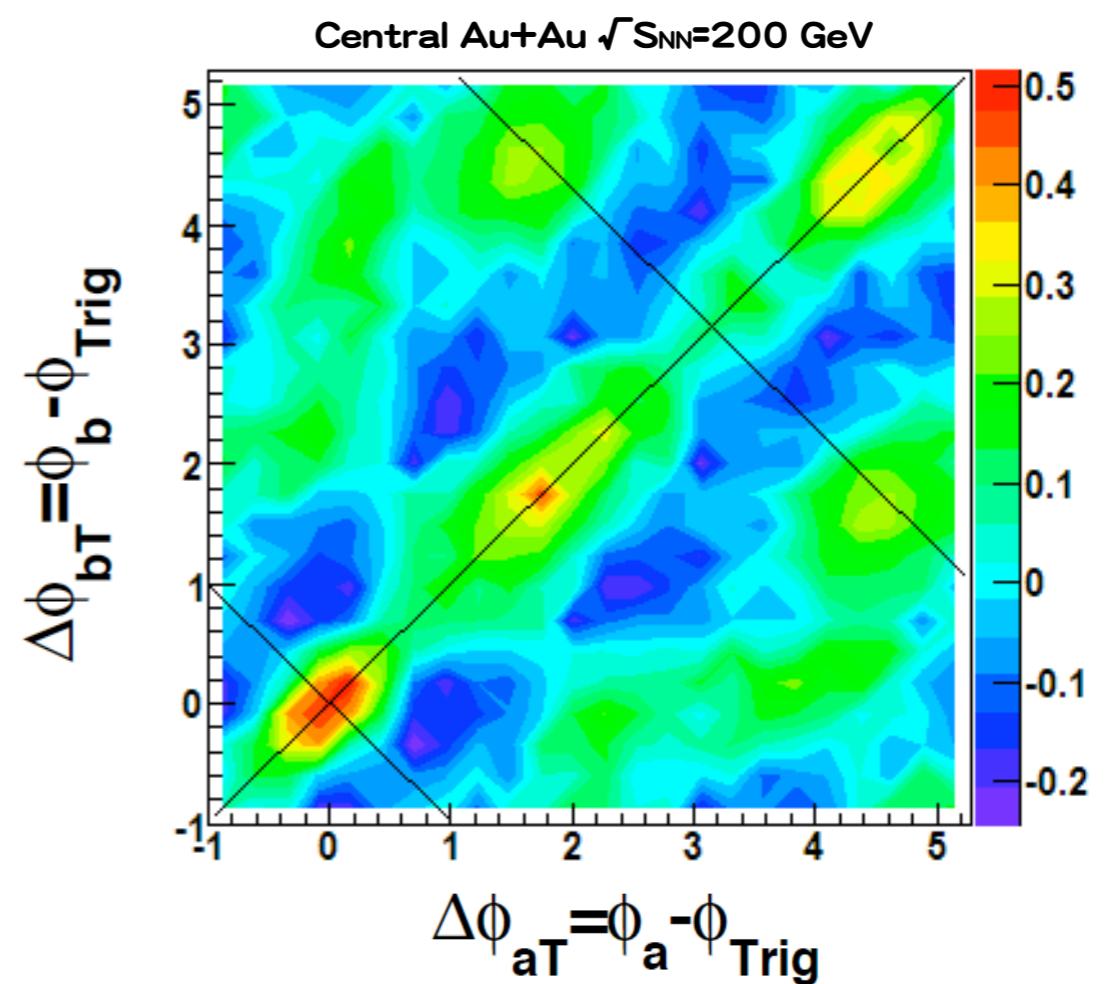
# Difficulty of the analysis



Phenix@QM06, J.Phys.G 34(2007)S191



STAR@QM06, nucl-ex/703010

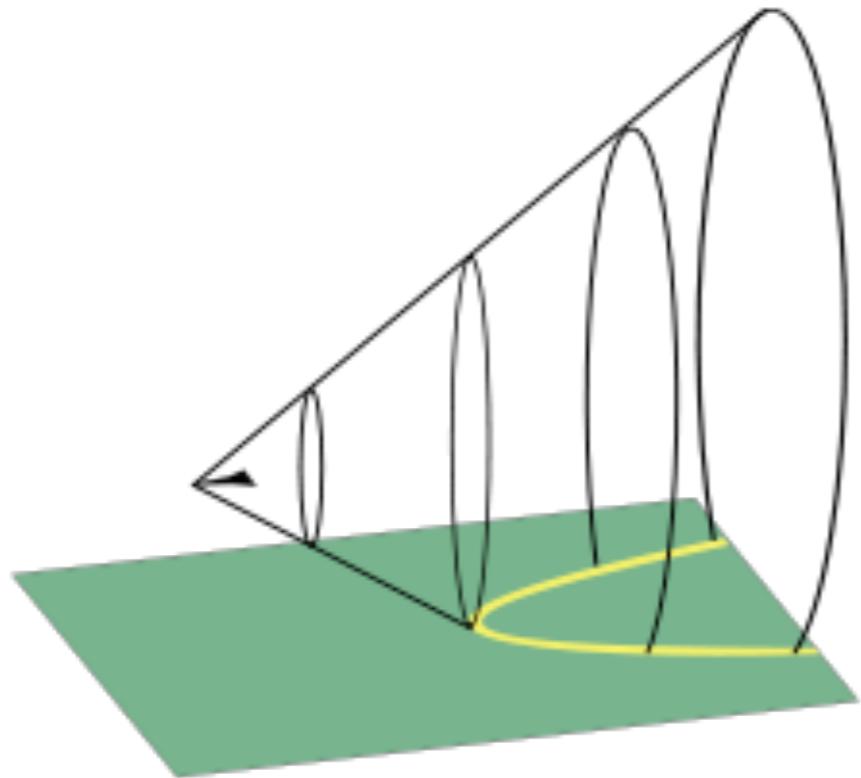
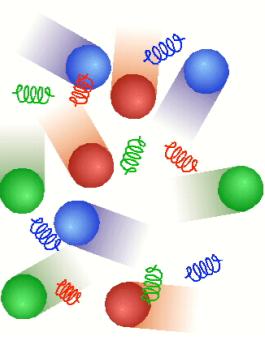


✓ Large azimuthal anisotropy is the major background.

- “Discovery of yesterday is background of today and calibration of tomorrow”
- More precise determination of the elliptic flow is important

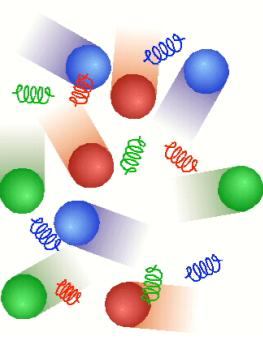
✓ Both group say it is consistent with Mach Cone.

# Shock Wave !?



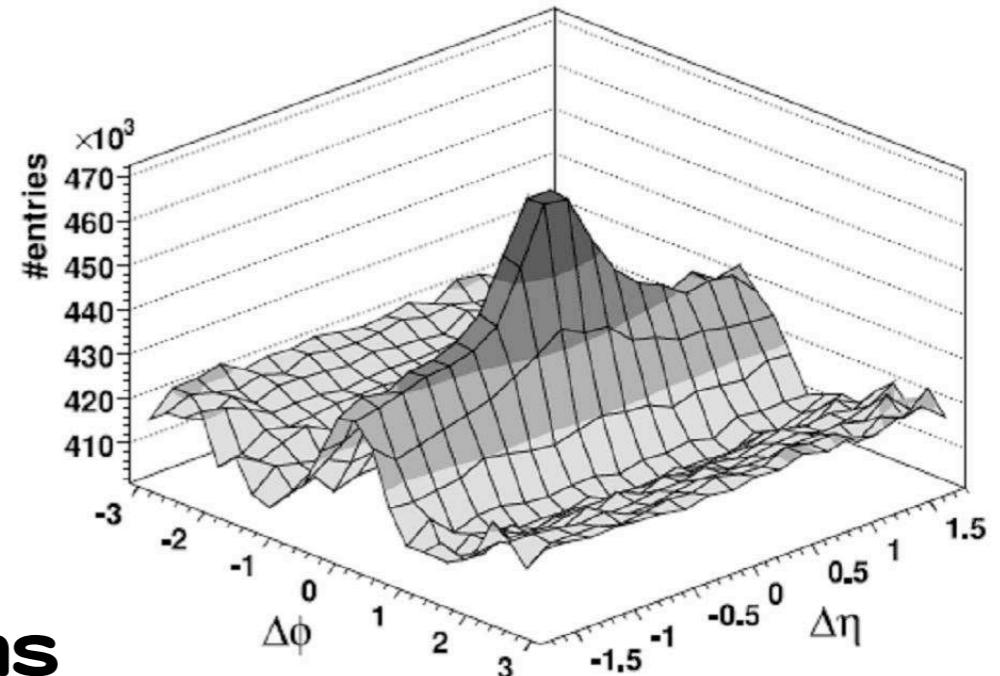
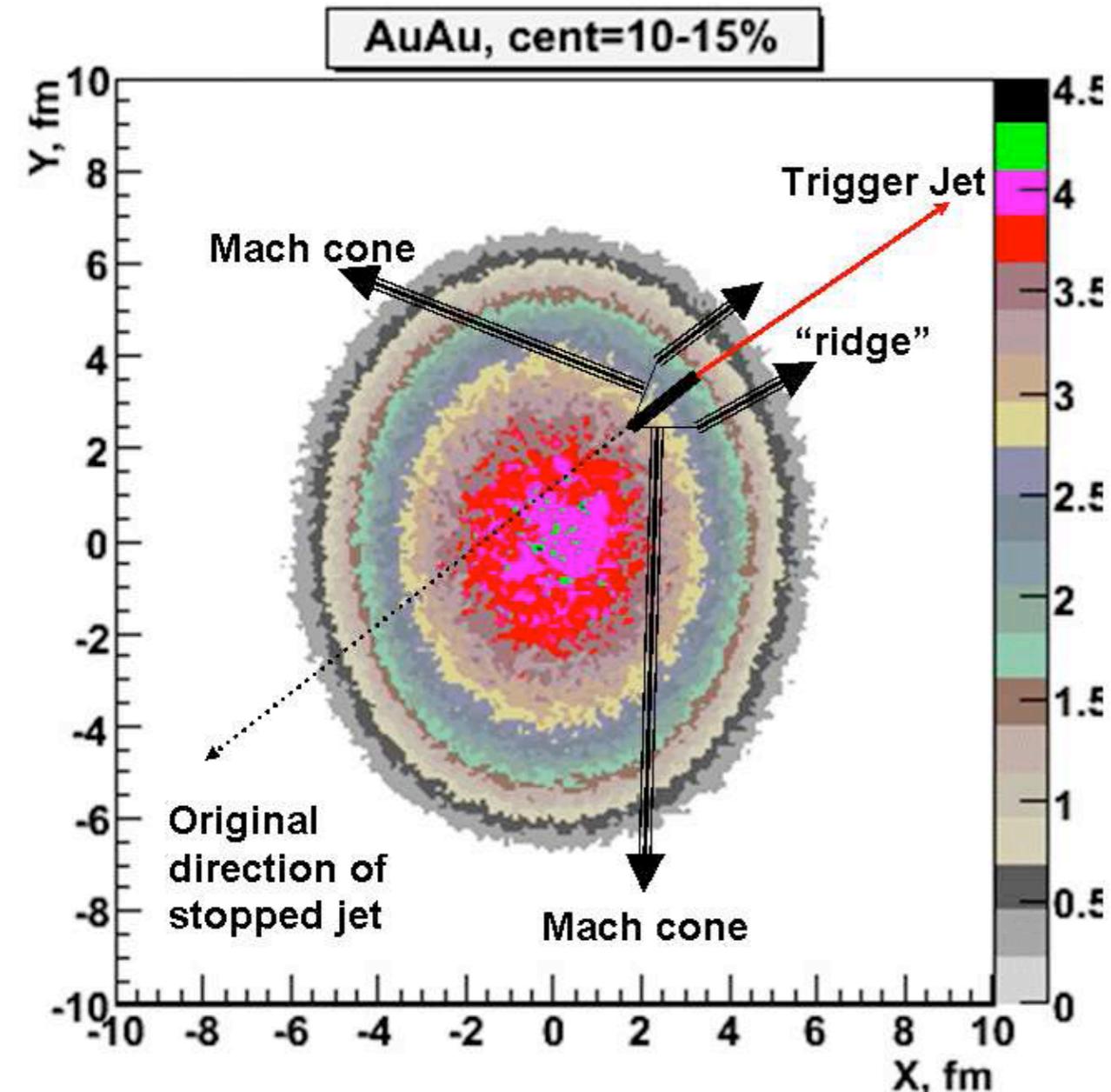
✓ If confirmed, it is breakthrough from the era of QGP discovery to the study of property, such as sound velocity in the plasma

# Near side ridge as splash



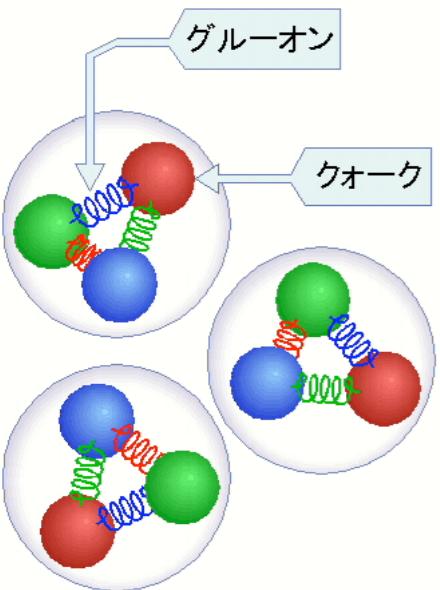
V.S. Pantuev, arXiv:hep-ph/0701.1882v1

STAR, arXiv:nucl-ex/0701074

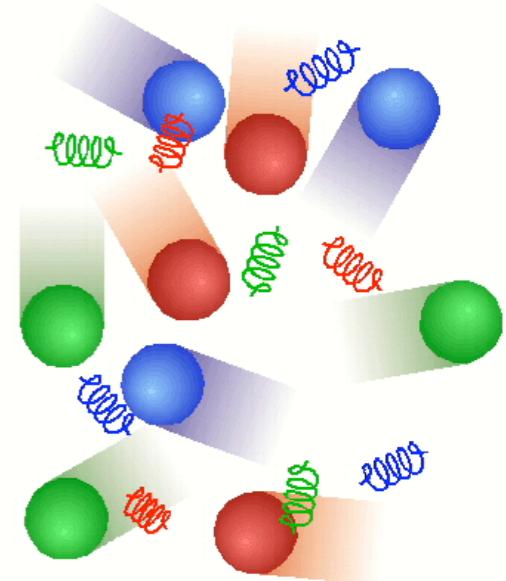


## ✓ Claims

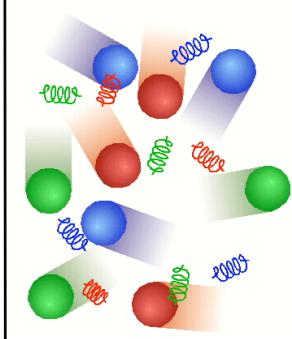
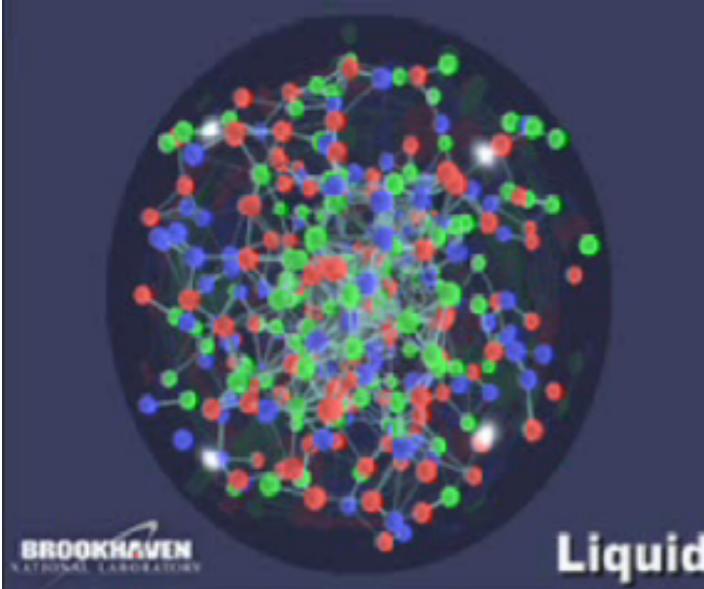
- ◆ Stopped parton is the source of ridge as well as Mach cone.
- ◆ This happens only at the surface
- ◆ There should be velocity boost at the ridge region ( $\pi, K, p$ )
- ◆  $v_3$  components independent of R.P.



If it is “Liquid”,  
why not a back  
splash?



# Summary



- ✓ We have seen partonic matter, ie, a QGP!
- ✓ Successful description of the system in terms of statistical thermo-dynamics;
  - ◆ Particle ratios in  $T_{ch}$ ,  $\mu$ , Kinematical distr. in  $T_{th}$  and  $\beta$
- ✓ Partonic
  - ◆ Large azimuthal anisotropy cannot be created with hadronic process.
  - ◆ High  $p_t$  suppression and disappearance of back-to-back is at parton level.
  - ◆ Successful description of quark recombination;
    - Phenomenological, but universal quark distribution function!
- ✓ We are in the state of studying property of plasma, like  $c_s$ .