

Hadron production in heavy ion collisions

Francesca Bellini, CERN (Switzerland) down to Heavy Ion Meeting – Seoul, July 4 ... and systems

QCD and light flavors



B. Muller, NPA 750 (2005) 84-97 1000000 □ QCD mass 100000 Higgs mass 10000 1000 100 10 1 b d t u s С

Confining property of QCD \Rightarrow colorless hadrons Asymptotic freedom \Rightarrow deconfinement, QGP Confinement and chiral transitions both around the same T_c

Lattice QCD locates the phase transition at $T_c \sim 145 - 164 \text{ MeV}$ [A. Bazavov et al., PRD 90 (2014) 094503]

Light flavor (LF) \equiv composed by **u**, **d** and **s**

$$\begin{array}{c} m_{u} \approx 2.2 \text{ MeV} \\ m_{d} \approx 4.7 \text{ MeV} \\ m_{s} \approx 96 \text{ MeV} \end{array} \right\} \quad < \Lambda_{QCD} << m_{c} \approx 1.3 \text{ GeV}$$

u, d, s can be **thermally produced in QGP**, as m_{u,d,s} < T_c

(Some of) The fundamental questions

- I. Can we produce QGP in AA collisions? And in pp, pA? What is the smallest system where we can produce QGP?
- 2. Can we measure the **chemical freeze-out** temperature T_{ch} ? Do all (light) flavours freeze-out at the same temperature T_{ch} ?
- 4. How does the hadronic phase affect the measured observables?
- 5. How are **composite "fragile" objects** such as anti-nuclei formed in AA?

Questions addressed by a comprehensive set of measurements of **identified particle production in all collision systems**: $\Pi^{\pm}, K^{\pm}, K^{0}{}_{S}, p, \Lambda, \Xi, \wedge, \rho, K^{*0}, \phi, \Sigma^{*\pm}, \Lambda^{*}, \Xi^{*0},$ (anti)d, (anti)³H, (anti)³He, (anti)⁴He, (anti)³_{\Lambda}H

It's all about production mechanisms

Thermal production

Measure p_{T} integrated yields and particle ratios

Hydrodynamics (flow)

Fragmentation + recombination

Measure p_T spectra, mean p_T, p_T -differential ratios, flow harmonics, multi-particle correlations

Composite objects (Anti-)(hyper-)nuclei

Strangeness

System size ⇔ charged particle multiplicity

Event multiplicity/centrality **classes** defined "slicing" on the signal amplitude measured in the forward region / detectors (In ALICE, in V0 scintillators, placed at $2.8 < \eta < 5.1$ (V0A) and $-3.7 < \eta < -1.7$ (V0C))

 $\langle dN_{ch}/d\eta \rangle$, N_{ch} measured in **at mid-rapidity** \rightarrow avoid "auto-biases" in multiplicity determination, especially in small systems



New data in Xe-Xe: test system dependence

A short (~6h) run with ¹²⁹Xe at $\sqrt{s_{NN}} = 5.44 \text{ TeV}$ in Nov. 2017 ¹²⁹Xe is smaller nucleus, different geometry wrt ²⁰⁸Pb



Charged particle production in AA collisions



In Pb-Pb $\langle dN_{ch}/d\eta \rangle / \langle N_{part} \rangle$ increases with \sqrt{s} following a steeper power law than pp

About 20% increase from 2.76 TeV to 5.02 TeV (similar $\langle N_{part} \rangle$)

Similar $\langle dN_{ch}/d\eta \rangle / N_{part}$ in most central Xe-Xe and Pb-Pb (little difference in energy)

Xe-Xe vs Pb-Pb



Consistency between **Pb-Pb** at different energies

Scaling violations:

- N_{part} scaling violated
- Central collisions of medium-size nuclei (Xe) produce more particles per N_{part} than mid-central collisions of large nuclei (Pb) at the same N_{part} \rightarrow to be understood

π, K, p spectra in AA collisions

πKp constitute the bulk of particle production \rightarrow fundamental input!

- *p*_T spectra:
 - Compare to hydrodynamic models
 - Quantify radial flow
 - Discuss flow effect vs recombination at intermediate p_{T}
- p_{T} -integrated particle ratios
 - Study hadro-chemistry and extract chemical freeze-out temperature
- Nuclear modification factors (R_{AA}, R_{pA})
- Reference for resonance and nuclei production studies



"First-order" look at spectral shapes: $\langle p_T \rangle$

First-moment of the distribution provides first-order characterization of spectral shapes

→ "day-I" benchmark for models

Hydrodynamics predicts

[G. Giacalone et al., PRC 97, 034904 (2018)]

- weak dependence on centrality in AA collisions
- small difference between Xe-Xe and Pb-Pb, O(3%)
- mass scaling of $\langle p_{\rm T} \rangle$
- → New data in Xe-Xe for charged hadrons show consistency with Pb-Pb and scaling of charged hadron production with charged particle multiplicity



$\langle p_{\rm T} \rangle$ of identified hadrons – Pb-Pb, Xe-Xe

 $\langle p_T \rangle$ of **identified hadrons** consistent between Xe-Xe and Pb-Pb at similar multiplicities

- Increasing of mean p_T with multiplicity
- Mass ordering observed, as featured by a system experiencing radial flow



$\langle p_{\rm T} \rangle$ of identified hadrons - p-Pb

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Similar trend with mult. in pp, pA (even steeper!)



ALI-PUB-103929

$\langle p_{\rm T} \rangle$ of identified hadrons - pp

$\langle p_{\rm T} \rangle$ of identified hadrons

consistent between Xe-Xe and Pb-Pb at similar multiplicities

- Increasing of mean p_T with multiplicity
- Mass ordering observed, as featured by a system experiencing radial flow

Similar trend with mult. in pp, pA (even steeper!)

BUT also models with initial stage effects (e.g. CGC + PYTHIA) can reproduce the mass ordering (also for v_2 !!!)



For details on the model, see backup

Mass ordering of $\langle p_T \rangle$ (and v_2) not be an exclusive product of hydrodynamical flow

 $\langle p_T \rangle$ only provides first-order characterization of spectral shapes, physics is in the full spectrum!

Blast-Wave model fit to spectra

Boltzmann-Gibbs Blast-Wave used to

quantify radial flow

[E. Schnedermann et al., Phys. Rev. C48 (1993) 2462]

A simplified hydrodynamic model with 3 free fit parameters,

- T_{kin} = kinetic freeze-out temperature
- $\langle \beta_T \rangle$: transverse radial flow velocity
- *n*: velocity profile

to describe particle production from a thermalized source + radial flow boost

Simultaneous fit to π , K, p spectra

- increase of $\langle \beta_T \rangle$ with centrality in AA
- Xe-Xe and Pb-Pb consistent



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- in pp and p-Pb, similar evolution with multiplicity
- at similar multiplicity, $\left<\beta_T\right>$ is larger for smaller systems

Hydro model comparison – Pb-Pb 0-5%





iEbyE + VISHNU + Trento/AMPT: viscous hydro with different initial conditions

different initial conditions [arXiv:1703.10792v1, PRC 92, (2015) 014903, PRC 92 (2015) 011901(R)]

McGill: MUSIC viscous hydro with IP-Glasma initial conditions [PRC 95 (2017) 064913]

"Full-hydro" models reproduce features of particle spectra and ratios below 2 GeV/c at 20-30% level

Hydro model comparison – Pb-Pb 0-5%





EPOS-LHC: core (hydro) + corona [*PRC* 92 (2015) 034906] does not reproduce spectra but **qualitatively** captures the spectral **ratios up to higher** p_T

For tested models, agreement worsens towards peripheral events / low multiplicity (see backup)

→ Improvements are needed to reproduce the evolution with multiplicity, also in view of extension of hydro description to small systems

Production mechanisms

Baryon-to-meson ratios are a powerful tool to test production mechanisms and their interplay

- **Low-***p*_T rise described by hydro
- Models where recombination involves only soft thermal radially flowing partons consistent with data
- High-p_T p/π is the same in pp and Pb-Pb collisions

 \rightarrow fragmentation dominates



Role of recombination at intermediate p_{T}

Behaviour in Xe-Xe confirms observations in Pb-Pb at 5.02 TeV (ratios compared at similar multiplicity)

Pivotal role of the ϕ -meson, that has similar mass as the proton

 the flatness of the p/φ ratio is consistent with hydro but can be accommodated by models with recombination





→ Still an **open point** on whether **recombination or flow** determine the spectral shape **at intermediate** p_T

System size evolution of dynamics



Particle ratios in Pb-Pb at 5.02 TeV and Xe-Xe at 5.44 TeV are **consistent** within uncertainties once **compared at the same multiplicity** (and not just centrality percentile)

Evolution with centrality in AA collisions and the shape of the p/π ratio are consistent with the presence of **radial flow** (mass dependence of spectral shapes)

...down to small systems



Across the three systems the baryon-to-meson ratios evolve with multiplicity in a qualitatively similar way

Baryon-to-meson – heavy-flavor sector



Heavy flavor baryon-to-meson ratio in small systems similar to light-flavors (Λ/K_{S}^{0})

• Not reproduced by models (PYTHIA, etc... see backup)



We observe a continuous evolution of spectral shapes and ratios with multiplicity across different systems

In HI hydrodynamics works at low $p_{T,}$ interplay with recombination at intermediate p_{T} is still an open point

In small systems, origin of collectivity is still to be understood

Statistical hadronisation model in a nutshell

Thermal fits map heavy-ion collisions to the QCD phase diagram and allow for comparison with lattice-QCD results.

Conventional picture: (ideal) hadron-resonance gas model in chemical equilibrium (based on Grand Canonical ensemble)

$$n_i = N_i/V = -\frac{T}{V}\frac{\partial \ln Z_i}{\partial \mu} = \frac{g_i}{2\pi^2} \int_0^\infty \frac{p^2 \mathrm{d}p}{\exp[(E_i - \mu_i)/T] \pm 1}$$

- Fit to yields: parameters $\mu_{\rm B}$, $T_{\rm ch}$, V
- Thermal model fit to yields: V cancels out
- Fits based on minimization of χ^2
- Deviations from (GC) equilibrium through empirical under(over)-saturation parameters for strange, charm or light quarks (γ_s , γ_c , γ_q)



Thermal model fit to Pb-Pb 2.76 TeV (0-10%)



Production of (most) lightflavour hadrons described $(\chi^2/ndf \sim 2)$ by thermal models with a **single chemical freeze-out** temperature, $T_{ch} \approx 156 \text{ MeV}$

Deviation for short-lived K^{*0} resonance Tensions between protons and multi-strange Nuclei included

Figure from ALICE, Nucl. Phys. A 971 (2018) 1-20 THERMUS: Wheaton et al, Comput.Phys.Commun, 180 84 GSI-Heidelberg: Andronic et al, Phys. Lett. B 673 (2011) 142 SHARE: Petran et al, arXiv:1310.5108

Thermal model fit to Pb-Pb 5.02 TeV (0-10%)



New preliminary ALICE data in 0-10% Pb-Pb at 5.02 TeV can be fitted with a slightly lower temperature, $T \approx 153 \text{ MeV}$ and higher $\chi^2/\text{ndf} \sim 4-6$

Tensions between protons and multi-strange are confirmed at the new energy

 Strange particles prefer a higher temperature.

Figure from ALICE, Nucl. Phys. A 971 (2018) 1-20 THERMUS: Wheaton et al, Comput.Phys.Commun, 180 84 GSI-Heidelberg: Andronic et al, Phys. Lett. B 673 (2011) 142 SHARE: Petran et al, arXiv:1310.5108

System size evolution of hadrochemistry



Chemistry is driven by charged particle multiplicity, i.e. the size of the system (regardless of its type and \sqrt{s})

 \rightarrow A further test for thermal model will be the fit to the Xe-Xe data, to check dependence on system size of Tch

Resonance suppression in central AA

Short-lived resonance ratios to long-lived particles are **suppressed as centrality increases** in AA collisions

ρ(770)/π (ρ lifetime = 1.3 fm/c)
 K(892)⁰/K (K* lifetime = 4.5 fm/c)
 Λ(1520)/Λ (Λ* lifetime = 12.5 fm/c)
 Ξ(1530)/Ξ (Ξ* lifetime = 22.5 fm/c)

- **Χ φ(1020)/K** (φ lifetime = 45 fm/c)
- Re-scattering effects expected to be stronger in central collisions, as the medium is denser and lasts longer
- Depending on the species, regeneration effects might be dominant → measure Sigma*!



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In thermal fits, differences between protons and strangeness sector observed at 2.76 TeV are confirmed at 5.02 TeV



Entering an era of precision tests for thermal model(s)

Important role of the hadronic phase for resonances. Is there more to the story (baryon annihilation for p)?

(Anti-)nuclei production

2 main models give predictions consistent in order of magnitude but based on different parameters:

Thermal production at chemical freeze-out/phase boundary

• Key parameters are mass and chemical freeze-out temperature:

 $dN/dy \sim \exp(-m/T_{ch})$

• Model provides yields but no p_{T} spectra (no dynamics)

Coalescence of nucleons at kinetic freeze-out

- Key parameters are nuclear wave functions, size of the (hyper)nucleus
- Production probability quantified by coalescence parameter B_A
- "Simple" coalescence model limited -- source size to be included in modeling
- Model provides spectra

$$E_i \frac{\mathrm{d}^3 N_i}{\mathrm{d} p_i^3} = B_A \left(E_\mathrm{p} \frac{\mathrm{d}^3 N_\mathrm{p}}{\mathrm{d} p_\mathrm{p}^3} \right)^A$$

Key measurements: multiplicity-dependence of nuclei with different sizes relative to the system size [F. Bellini, A. Kalweit, in preparation]

Nuclei are rare objects



pp/p-Pb collisions:

"penalty factor" for adding one nucleon ~600-10³

Anti-matter / matter $\sim I$ at the LHC

Pb-Pb collisions:

- "penalty factor" for adding one nucleon
- Consistency with thermal model expectations



Deuteron flows in AA

Deuterons develop v_2 , surviving in the fireball despite their low binding energy ($B_E \sim 2.2 \text{ MeV}$)

→ Proposed deuteron as "6-quarks bag" formed thermally (A. Andronic et al.)

Simple coalescence fails in reproducing the deuteron B_2 and v_2



Nuclei production across collision systems



Smooth evolution of d/p ratio across systems and no significant $\sqrt{s_{NN}}$ dependence at $\mu_B = 0$

 B_2 decrease with centrality in Pb-Pb is explained as an increase in the source volume (in "full" coalescence models, e.g. Scheibl-Heinz, Sato-Yazaki, Nagle, ...)



Many new results in the nuclei sector at the LHC contribute to the systematic study of production mechanisms

Key measurements: multiplicitydependence of nuclei with different sizes relative to the system size to probe coalescence vs thermal production
Strangeness production

Enhancement of strangeness from low to high multiplicity pp, p-Pb collisions, until saturation in Pb-Pb [ALICE, *Nat. Phys. 13, 535–539 (2017)*]

ightarrow confirmed with new data from LHC Run II

Ongoing efforts to explain behavior with models

- Lund string, color ropes (PYTHIA, DIPSY)
- core-corona (EPOS-LHC)
- thermal-statistical (canonical suppression) [V.Vislavicius, A. Kalweit, aXiv:1610.03001]
 - Conventional pp generators successful, with MPI + CR generating some collectivity, but now cracks.
 - Need new framework for baryon production.

T. Sjostrand, Quark Matter 2018



Strangeness canonical suppression

In equilibrium SHM models strangeness enhancement is a result of the **canonical suppression** of strangeness production **in small systems** due to the explicit **conservation** of the **strangeness** quantum number in a finite system

First comparisons to model calculations based on THERMUS code

→ agreement with data within uncertainties, except for φ meson (also "immune" to canonical suppression)



The special role of ϕ meson

As a s-sbar pair (S=0) with the same mass as the proton, the ϕ meson is "special" \rightarrow **Does** ϕ **behave like a S=0 or S=2 particle?**

- Indications of increase of ϕ/π ratio with multiplicity in small systems
- Flat Ξ/φ for multiplicities between ~6 and ~700? Or slightly increasing in pp, p-Pb vs multiplicity?

→ Need more precision from experiment!



System size evolution of hadrochemistry



Particle composition evolves smoothly across collision systems, depending on charged particle multiplicity.



Common origin in all systems?

For MC generators, work is still needed to reproduce evolution with system size in view of a unified description of all collision systems

Collectivity but no jet quenching?

Similarities are observed for flow observables between peripheral Pb-Pb and high multiplicity p-Pb collisions.

New and more precise measurements from ALL experiments on nuclear modification factors.

In (minimum bias) **p-Pb**, no suppression at high- p_T is observed, contrary to **peripheral Pb-Pb**.

 \rightarrow Do we understand this?



ALICE, arXiv:1802.09145

Nuclear modification in very peripheral collisions

New measurement in very fine centrality bins!

Strong change of behaviour of R_{AA} beyond 80% centrality





Nuclear modification in very peripheral collisions

New measurement in very fine centrality bins!

Strong change of behaviour of R_{AA} beyond 80% centrality

→ reproduced by HG-PYTHIA with **biases** in event selection and collision geometry, and no nuclear modification.

Considering this, the jet quenching signal is smaller than typical systematics above ~80% centrality \rightarrow consistent with R_{pPb}



Conclusions and outlook



Continuity in chemistry and dynamics across collision systems (dependence on charged particle **multiplicity**) is observed.

pp, p-nucleus collisions are much more than a "reference" for heavy-ion collisions



MC generators can generate collective-like behaviour but **fail in the details** of hadron (baryon) production as a function of multiplicity.

Absence of jet-quenching in small systems remains as the main challenge to the final-state effect interpretation.

Origin of collectivity in small systems is still to be understood.

STRANGENESS

Strangeness from the HI perspective

- ~300 MeV are enough to create an $s\overline{s}$ pair^[1] (even less if $m_s^{\text{QCD}} \rightarrow m_s^{\text{Higgs}}$ by restoration of chiral symmetry in the QGP phase)
- Strange quarks are dominantly produced by (thermalised^[2]) gluon fusion in QGP
- Strangeness enhacement wrt pp collisions historically proposed as signature for a deconfined QGP^[3]
- pp collisions as reference

[1] PDG group, Chin. Phys. C38 (2014) 090001
[2] E. Shuryak, Phys. Rev. Lett. 68 (1992) 3270
[3] J. Rafelski and B. Muller, PRL 48 (1982) 1066



Strangeness from the pp perspective

Producing strangeness is "expensive"
 → threshold problem



E.g. in a hadron gas $\pi + \pi \rightarrow \pi + \pi + \Lambda + \Lambda$ -bar, $E_{th} \sim 2200 \text{ MeV}$

- Measurements of strange hadron production used as input for tuning Monte Carlo generators
- Contribute to the understanding of the rich structure of the underlying event arising from MPI in pp, p-Pb collisions



Strangeness production in QGP

- ~300 MeV are enough to create an s-sbar pair (even less if $m_s^{QCD} \rightarrow m_s^{Higgs}$ by restoration of chiral symmetry)
- gluon fusion (a) is the dominant mechanism for strangeness production over quark annihilation (b)
 - Gluons quickly thermalise in t < 1 fm/c
 [E. Shuryak, Phys. Rev. Lett. 68 (1992) 3270]
- The backward reaction of (b) depends on the s quark density, thus on the QGP lifetime → saturation of strangeness abundance
- After hadronisation, the abundance of (multi)strange hadrons reflects that of strangeness in the partonic phase
 - If the hadronic phase is short enough to avoid re-diffusion
 - For small hadronic cross sections



FIG. 3. Time evolution of the relative strange-quark to baryon-number abundance in the plasma for various

Strangeness enhancement in pp



No increase for p/π is observed **Hierarchy** of the increase associated with the strangeness content

What's next?

Does strangeness keep increasing with multiplicity or saturate?

High multiplicity-triggered data sample in pp 13 TeV (2016, 2017) being analysed

Measure in p-Pb at 8.16 TeV, Xe-Xe at 5.44 TeV, more differential in peripheral Pb-Pb collisions (2018)





What's next?

2 Can we relate high multiplicity with soft- or hard-QCD dominated processes? Use event shapes as tools to select jetty/isotropic events in high multiplicity PP

2 Can the φ meson provide further insights on strangeness production vs multiplicity? Measure more differential (event shapes?), extract φ/π, improve precision

3 New observables...



COLLECTIVITY IN SMALL SYSTEMS

Collectivity (in short)

"Loose" definition: correlations of (more than 2) particles across rapidity due to a common source

Origin of collectivity:

- Initial state correlations → among hadrons in the final state arise from momentum correlations at partonic level [gluon saturation, CGC]
- Final state correlations → anisotropies and correlations in space converted into anisotropies in momentum space, e.g. via hydrodynamic flow [established in Pb-Pb collisions]

Flow in heavy-ion collisions



$$E\frac{\mathrm{d}^3N}{\mathrm{d}p^3} = \frac{1}{2\pi} \frac{\mathrm{d}^2N}{p_{\mathrm{T}}\mathrm{d}p_{\mathrm{T}}\mathrm{d}y} (1 + 2\sum_{n=1}^{\infty} v_n) \cos[n(\varphi - \Psi_n)]),$$

Characteristic features:

- Multiplicity dependence
- Higher harmonics azimuthal flow
- Mass scaling of v₂
- Correlations between harmonics

The hallmarks of flow in AA collisions

Increase in mean p_T with increasing centrality \rightarrow Push from radial flow affects low p_T part of spectra



Baryon-to-meson ratios (with Δm) \rightarrow sensitive to particle production mechanisms (radial flow at low p_T , recombination at mid- p_T)



The hallmarks of flow in AA collisions

Centrality / multiplicity dependence \rightarrow reflects the degree of "anisotropy" in the initial geometry of the collision Non-zero higher-order flow coefficients ("harmonics") \rightarrow sensitivity to fluctuations of initial geometry



The hallmarks of flow in AA collisions

Mass scaling of flow coefficients \rightarrow Expansion under a common velocity field

Correlations between harmonics \rightarrow Sensitivity to fluctuations in initial geometry (v_2 , v_3) and mediumtransport properties (v_2 , v_4)



Signs of collectivity in small systems



Signs of collectivity in **small systems** "discovered" at the LHC in terms of long-range ($2 < |\Delta \eta| < 4$) near-side ($\Delta \phi = 0$) "ridge" in 2-particle correlations, visible in **high multiplicity** pp, p-Pb, Pb-p collisions

Are these long-range correlations coming from (hydrodynamic) flow?

→ Investigated with new measurements with run 2 data, new analysis techniques

The challenge of removing "non-flow"

In **small systems** the contribution of **non-flow** cannot be neglected:

- Different contribution from jets
- Larger fluctuations in the number of particle sources

A word of **caution**:

- Sensitivity to the event class definitions used in analysis [ATLAS, EPJ C (2017) 77-428]
- Sensitivity to strategy for nonflow background subtraction





Non-flow subtraction / suppression is a delicate business in pp, p-Pb! Big effort ongoing in defining "smart" observables / new techniques

If collectivity, it involves more than 2 particles

Measure elliptic flow v_2 using correlations among k particles in a single event, subtracting correlations from smaller number of particles [A. Bilanzic et al., PRC 83 (2011) 044913]



If long-range, correlations stay across sub-events



True collectivity in small systems!





 $v_2{4} \approx v_2{6} \approx v_2{8} \rightarrow \text{true collectivity (even) in}$ smallest systems $v_2{2} \text{ larger } \rightarrow \text{ residual "non-flow"}$

Light-flavor particle v_2



Heavy-flavor particle v_2 in p-Pb



ADDITIONAL SLIDES

 $R_{\Lambda\Lambda}$ in Pb-Pb and Xe-Xe



R_{AA} in Xe-Xe and Pb-Pb consistent if compared at similar charged particle multiplicity

Identified particle production in |y| < 0.5



Probing the hadronic phase with resonances

Key measurements:

- Resonance yields and ratios to long-lived particles vs. centrality
 - Re-scattering effects expected to be stronger in central collisions, as the medium is denser and lasts longer
 - Depending on the species, regeneration effects might be dominant (e.g. Σ*)
- Spectra down to low *p*_T
 - Improve precision on the yields by minimising the extrapolated fraction
 - UrQMD predicts the largest effects for $p_T < 2 \text{ GeV}/c$



SHARE fit – Pb-Pb 5.02 TeV (0-10%)



SHARE: M. Petran et al.. PRC88 (2013) 3, 034907

Hydro model comparison – semi-central Pb-Pb



Hydro model comparison - peripheral Pb-Pb



Blast-Wave fits to particle spectra

Boltzmann-Gibbs Blast-Wave model [E. Schnedermann et al., Phys. Rev. C48 (1993) 2462] A simplified hydrodynamic model with 3 free fit parameters:

- β_T : transverse radial flow velocity
- T_{kin} = kinetic freeze-out temperature
- n: velocity profile


Blast-Wave fits - Pb-Pb

Caveats:

- Sensitivity to fit range: low p_T particle spectra affected by resonance decay in the hadronic phase

- Sensitivity to the set of particles included in the fit





Λ, Ξ, **Ω**, **φ** compared to prediction obtained from fit parameters resulting from π,K,p (gray bands indicate range in which fit works ok in Pb-Pb 2.76 TeV)

Blast-Wave fits - Xe-Xe

1.4 data / BW-fit Fit range 1.3 1.2 1.1 1 0.9 0.8 0.7 ALICE Preliminary, Xe-Xe $\sqrt{s_{NN}} = 5.44 \text{ TeV}$ 0.6 -0-5% -5-10% -10-20% $\pi^++\pi^-$ 0.5 70-90% 50-60% 0.4 0.4 0.6 0.8 $p_{_{\rm T}}^{^1} ({\rm GeV}/c)^{^{1.2}}$ 0.2 ALI-PREL-155927



v_2 of identified hadrons in Pb-Pb 5.02 TeV

arXiv:1805.04390



Light-flavor particle v_2



Clear mass ordering at low p_T in p-Pb from new results on v_2^{sub} for identified hadrons

 \rightarrow Consistent with hydrodynamics (and AA)

BUT it could also be due to other effects

- Initial stage effects (CGC + PYTHIA)
- Parton escape (AMPT)
- Hadronic rescattering (UrQMD)



Is mass ordering no longer an exclusive product of hydrodynamic flow?

CGC + Pythia

CGC meets Lund String fragmentation of PYTHIA

CGC + PYTHIA : A new approach to simulate p+p & p+A collisions

- 1) Output distribution of Gluons from CGC based **IP-Glasma** model
- 2) Sample gluons in momentum space
- 3) Connect the gluons close in phase space to color neutral strings
- 4) Input to **PYTHIA** and fragment into final particles



PYTHIA (realistic mechanism of hadronization) \rightarrow partonic correlations from initial state dynamics \rightarrow correlated production of final-state particles

P.Tribedy, Quark Matter 2018, Venice, Italy

Bayesian analysis, Pb-Pb and Xe-Xe



Comparison to Bayesian analysis



Mean p_T and integrated yields in Pb-Pb at 2.76 TeV are used as input to a multiparameter **Bayesian analysis** to extract predictions for 5.02 TeV [G. Bernhardt et al., NPA 967 (2017) 293–296]

Comparison with ALICE results at 5 TeV indicate that

- the K/ π ratio is reproduced in most central collisions, but description fails in peripheral
- p/π and mean transverse momenta overestimated, trend with multiplicity described qualitatively \rightarrow model improvements since 2017

Lc/D ratio in small systems vs MC



The ALICE detector at the LHC



Particle identification in ALICE



Identification of light flavour hadrons and light (anti-)nuclei via practically ' all known PID techniques in 0.1 GeV/c $< p_T < 30$ GeV/c