Production of strange particles in jets in p–p, p–Pb and Pb–Pb collisions measured with ALICE

Vít Kučera¹

¹Inha University

25 May 2018

Heavy Ion Meeting 2018-05 Kangwon Nat'l University, Chuncheon

Motivation for PID in jets

- Baryon-to-meson ratio is enhanced in A–A and p–A collisions (RHIC, LHC).
- This phenomenon cannot be explained by fragmentation in vacuum.
- What is the effect of QGP on hadronization mechanism(s) in jets?
- What are the mechanisms (parton recombination)?



Comparison of data with models



Phys. Rev. Lett. 111 (2013) 222301

Motivation for PID in jets

We aim to understand the origin(s) of the Λ/K_S^0 enhancement by separating hadrons produced in hard processes (jets) from hadrons produced in soft processes (underlying event).

Is the baryon-to-meson ratio enhanced due to the collective effects in the plasma (parton recombination, radial flow,...) or is it (also) due to a modification of the jet fragmentation in the medium?

- jet fragmentation
 A high-p_T parton from hard scattering fragments into hadrons.
- parton recombination Multiple partons cluster together to form a hadron.



ALICE

- ▶ collisions studied: p–p at $\sqrt{s} = 7$ TeV, p–Pb at $\sqrt{s_{\rm NN}} = 5.02$ TeV, Pb–Pb at $\sqrt{s_{\rm NN}} = 2.76$ TeV
- tracking of charged particles by ITS & TPC in magnetic field of 0.5 T
- centrality estimated from the multiplicity of charged particles in the detectors at forward and backward pseudorapidities



Analysis of charged jets

- track selection
 - charged primary particles
 - *p*_T^{track} > 150 MeV/c
 - uniform in $\phi imes \eta$, $|\eta_{
 m track}| < 0.9$
- raw-jet reconstruction
 - anti-k_t algorithm
 - resolution parameter R = 0.2, (0.3, 0.4)
- subtraction of average soft background
 - average background density ρ estimated from the median $k_{\rm t}$ cluster
- signal-jet selection (good candidates for hard scattering)
 - $p_{T}^{\text{leading track}} > 5 \text{ GeV/c} \text{ (only Pb-Pb)}$
 - $A_{\rm jet,ch} > 0.6\pi R^2$
- further $p_{\rm T}^{\rm jet,ch}$ corrections
 - background anisotropy (intra-event p_T fluctuations)
 - detector response

Jet spectra



Larger $R \Rightarrow$ harder spectrum (but softer jets at a given p_{T}^{jet}).

Analysis of neutral strange particles

Strange neutral particles decaying into two charged daughter particles

- $\blacktriangleright \text{ meson } \mathsf{K}^0_S \rightarrow \pi^+ + \pi^- \text{ (BR 69\%)}$
- baryon $\Lambda \rightarrow p + \pi^-$ (BR 64%)

Mother V^0 particle reconstructed using topology of its V-shaped decay.

Combinatorial background suppressed by cuts on decay parameters. Signal yield extracted from the invariant-mass distribution.



Strange particles in jets

Analysis steps

- V⁰ candidate selection
- candidate-jet matching (V⁰s in jet cones)

$$\sqrt{(\phi_{\mathsf{V}^0}-\phi_{\mathsf{jet},\mathsf{ch}})^2+(\eta_{\mathsf{V}^0}-\eta_{\mathsf{jet},\mathsf{ch}})^2} < R,$$

 $|\eta_{\rm jet,ch}|^{\rm max} < |\eta_{\rm V^0}|^{\rm max} - R$

- candidate-UE matching (V⁰s in events without selected jets with p_T^{jet,ch} > 5 GeV/c)
- signal extraction (invariant-mass distribution)
- efficiency correction (in jet cones, in UE)
- subtraction of V⁰s in UE
- subtraction of V⁰s coming from decays of jet constituents (Ξ → Λ), i.e. "feed-down" correction



Estimation of V^0 s in the underlying event

- ▶ no-jet events: V⁰s in events with no selected jets
- outside cones: V⁰s outside jet cones
- ▶ random cones: V⁰s in a randomly oriented cone
- median-cluster cones: V^0 s in the cone of the median k_t -cluster
- ▶ perpendicular cones: V⁰s in cones perpendicular to the jet in azimuth

Methods differ in regions, events, statistics, efficiency.



 K_S^0



Reconstruction efficiency of V^0 particles

- Reconstruction efficiency depends strongly on $p_{\rm T}^{\rm V^0}$ and $\eta_{\rm V^0}$.
- Shape of the measured $\eta_{\mathcal{N}^0}$ distribution depends on the selection criteria.
- ▶ Not enough statistics to apply efficiency correction in 2D ($p_T^{V^0} \times \eta_{V^0}$).

 \Rightarrow Efficiency of inclusive V^0s is scaled (in 2D) to get efficiency in jet cones and UE (in 1D).



Feed-down in jets

Feed-down fraction of $\boldsymbol{\Lambda}$ in jets estimated from:

- inclusive Λ (Pb–Pb-like),
- ▶ jets generated by PYTHIA 8 (p-p-like).



Estimation of systematic uncertainties

The systematic uncertainties are studied for the following sources:

- ▶ reconstruction efficiency of V^0 s (selection cuts applied on V^0 candidates),
- signal extraction (fitting parameters),
- subtraction of spectra of V⁰s in UE (5 methods),
- subtraction of feed-down in jets (inclusive vs PYTHIA),
- material budget (detector model),
- fluctuations of UE (jet embedding).

Open issue: $\Lambda - \overline{\Lambda}$ asymmetry

Discrepancy between inclusive spectra of Λ and $\overline{\Lambda}$.



Strong dependence on the polarity of magnetic field and the sign of η . Additional 6% (symmetric) considered as systematics.

Open issue: discrepancy in MC between runs 2010, 2011

Differences in spectra traced back to the reconstruction efficiencies.



Effects of $\lesssim 10~\%$ for $K_S^0, \lesssim 20~\%$ for $\Lambda,$ partially cancel out in $\Lambda/K_S^0.$ Observed also in other analyses (charged-particle spectra, correlations). Additional 10% (symmetric) considered as systematics.

Systematics: combined

$$R=0.2,~p_{
m T}^{
m jet,ch}>10\,{
m GeV/c},~(\Lambda+\overline{\Lambda})/2{
m K}_{
m S}^0$$



Uncertainties from different sources (except feed-down) are combined in squares and considered symmetric.

V⁰ vertexer problem

 There may be two candidates for the point of closest approach ("cowboy/sailor" configuration).



- Old vertexer: sailor misidentified as cowboy ightarrow CPA pprox -1
 ightarrow rejected.
- New vertexer: Select the point with the smallest DCA calculated in 3D.



- ► ⇒ Better MC-real data matching.
- Cause of loosing sailors with the old vertexer still unclear.

$\Lambda/\mathrm{K}^0_\mathrm{S}$ ratio in jets in p–p at $\sqrt{s}=$ 7 TeV and 13 TeV



- The ratio in UE is consistent with the inclusive ratio.
- The ratio in jets is clearly different from the inclusive ratio at low and intermediate p_T^{V⁰}.
- ► A slight increase of the ratio in jets with increasing *R*.

Λ/K_S^0 ratio in jets in p–Pb at $\sqrt{s_{NN}}=5.02\,\text{TeV}$ (high-multiplicity collisions, 0–10%)



The ratio in jets

- ▶ is clearly different from the inclusive ratio at low and intermediate $p_T^{V^0}$,
- is different from the inclusive ratio in PYTHIA (black line),
- is similar to the ratios in PYTHIA jets (red dashed lines),
- shows no significant dependence on $p_{\rm T}^{\rm jet,ch}$ and a slight dependence on R.

 Λ/K_S^0 ratio in jets in Pb–Pb at $\sqrt{s_{NN}} = 2.76 \text{ TeV}$ (7.4 × 10⁶ central collisions, 0–10%)



The ratio in jets

- ▶ is clearly different from the inclusive ratio at low and intermediate $p_{\rm T}^{V^0}$,
- shows no significant dependence on p_T^{jet,ch}
- ▶ is consistent with the ratio in jets in p–Pb and p–p at $p_T^{V^0} > 4 \text{ GeV}/c$.



- Same slopes of spectra from measurement and from PYTHIA.
- Enhancement for Λ at $p_{\rm T}^{\rm V^0} < 4 \,{\rm GeV/c}$.

Summary and outlook

ALICE has performed the first measurement of the Λ/K_S^0 ratio in charged jets in p–p, p–Pb and Pb–Pb collisions at the LHC.

Results

- ▶ In every collision system, the Λ/K_S^0 ratio in jets is significantly smaller than the inclusive ratio (and the UE).
- ► The Λ/K_S^0 ratios in jets are consistent within uncertainties in all collision systems for $p_T^{V^0} > 4 \text{ GeV/c}$.
- The dominant source of the enhancement are soft processes associated with collective behaviour.
- ► A potential modification of jet fragmentation seems to be restricted to the region $p_T^{V^0} < 4 \text{ GeV/c}$ and manifest by an enhancement of the Λ yields.

Outlook

- ▶ Solve the 2011/2010 issue.
- Comparison with more models (JEWEL).

Thank you for your attention.

Backup

 p/π ratio in Au–Au at $\sqrt{s_{NN}}=200\,\text{GeV}$



Phys. Rev. Lett. 97 (2006) 152301

 Ξ/Λ ratio in jets in p-p at $\sqrt{s_{\rm NN}} = 13 \,{\rm TeV}$



Pengyao Cui for the ALICE Collaboration, Quark Matter 2018

Details on the $\Lambda/\overline{\Lambda}$ discrepancy



V^0 candidate selection

| Cut variable | Value |
|--|--|
| Daughter tracks | |
| TPC refit | true |
| type of production vertex | not kKink |
| DCA to the primary vertex | $\geq 0.1{ m cm}$ |
| DCA between daughters | $\leq 1\sigma_{TPC}$ |
| $ \eta $ | _≤ 0.8 |
| V ⁰ candidate | |
| reconstruction method | offline |
| cosine of the pointing angle (CPA) | \geq 0.998 |
| radius of the decay vertex | 5–100 cm |
| $ \eta $ | ≤ 0.7 |
| transverse proper lifetime | $\leq 5	au$ |
| Armenteros–Podolanski cut (K ⁰ _S) | $p_{\mathrm{T}}^{\mathrm{Arm.}} \geq 0.2 \alpha^{\mathrm{Arm.}} $ |

Jet algorithms

A sequential recombination jet finder is defined according to this general scheme:

1. $\forall i, j$: calculate distances d_{ij} and d_{iB} (NB $k_t \equiv p_T$):

$$d_{ij} = \min\left(k_{t,i}^{2p}, k_{t,j}^{2p}
ight) rac{\Delta_{ij}^2}{R^2}, \quad \Delta_{ij}^2 = (y_i - y_j)^2 + (\phi_i - \phi_j)^2, \quad d_{iB} = k_{t,i}^{2p}$$

2. Find d_{min}:

$$d_{\min} = \min\left(d_{ij}, d_{iB}\right).$$

- If ∃ i, j : d_{min} = d_{ij}, merge particles i and j into a single particle and combine their momenta.
- If $\exists i : d_{\min} = d_{iB}$, declare particle *i* to be a final jet and remove it from the list.

These steps are repeated until no particles are left.

$$p = \begin{cases} 1 & k_{\rm t} \text{ (background estimation} \\ 0 & \text{Cambridge/Aachen} \\ -1 & \text{anti-}k_{\rm t} \text{ (signal jets)} \end{cases}$$



Matteo Cacciari et al. JHEP 0804 (2008) 063

Background in Pb-Pb

Production of soft particles by underlying-event processes.

average background density ρ :

► *k*_t jets w/o 2 hardest

each event: $\rho = \text{median} \left\{ p_{\text{T}}^{\text{jet}} / A_{\text{jet}} \right\}$ each jet: $p_{\text{T,jet}}^{\text{corrected}} = p_{\text{T,jet}}^{\text{raw}} - \rho A_{\text{jet}}$



 ρ anisotropy in events (fluctuations):

$$\blacktriangleright \ \delta p_{\mathsf{T}} = p_{\mathsf{T},\mathsf{probe}}^{\mathsf{raw}} - \rho A_{\mathsf{probe}}$$

• response matrix \rightarrow deconvolution



ibid.

ALICE, JHEP 1203 (2012) 053

Scaling of the reconstruction efficiency

- $\blacktriangleright~\epsilon$ reconstruction efficiency of inclusive particles
- ▶ ϵ_s reconstruction efficiency of particles of interest (scaled ϵ)
- ▶ *a_s* yield of associated particles of interest
- ▶ g_s yield of generated particles of interest
- m uncorrected yield of measured particles (candidates) of interest
- t yield of true (corrected) particles of interest
- P signal purity

Signal extraction in JC, UE (assume that $P_{\text{inclusive}}(p_T^{V^0}, \eta_{V^0})$ is the same as for V⁰s of interest):

$$m(p_{\mathsf{T}}^{\mathsf{V}^0},\eta_{\mathsf{V}^0})=m_{\mathsf{raw}}(p_{\mathsf{T}}^{\mathsf{V}^0},\eta_{\mathsf{V}^0})|_{\mathsf{peak region}}\cdot P_{\mathsf{inclusive}}(p_{\mathsf{T}}^{\mathsf{V}^0},\eta_{\mathsf{V}^0})|_{\mathsf{peak region}}$$

Efficiency calculation:

$$\begin{aligned} a_{s} \equiv m, \quad \sigma_{a_{s}} \equiv 0, \qquad g_{s} = a_{s}/\epsilon \\ \frac{1}{\epsilon_{s}(p_{T}^{V^{0}})} = \frac{\sum_{\eta_{\vee 0_{j}}} g_{s}(\eta_{\vee 0_{j}}, p_{T}^{V^{0}})}{\sum_{\eta_{\vee 0_{j}}} a_{s}(\eta_{\vee 0_{j}}, p_{T}^{V^{0}})} = \sum_{\eta_{\vee 0_{j}}} \frac{a_{s}(\eta_{\vee 0_{j}}, p_{T}^{V^{0}})}{\sum_{\eta_{\vee 0_{j}}} a_{s}(\eta_{\vee 0_{j}}, p_{T}^{V^{0}})} \frac{1}{\epsilon(\eta_{\vee 0_{j}}, p_{T}^{V^{0}})} \end{aligned}$$

Spectra correction:

$$t = m/\epsilon_s$$