Production of strange particles in jets in p–p, p–Pb and Pb–Pb collisions measured with ALICE

Vít Kučera<sup>1</sup>

<sup>1</sup>Inha University

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# Motivation for PID in jets

- <sup>I</sup> Baryon-to-meson ratio is enhanced in A–A and p–A collisions (RHIC, LHC).
- This phenomenon cannot be explained by fragmentation in vacuum.
- What is the effect of QGP on hadronization mechanism(s) in jets?
- $\triangleright$  What are the mechanisms (parton recombination)?



#### Comparison of data with models



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# Motivation for PID in jets

We aim to understand the origin(s) of the  $\Lambda/\mathsf{K}^0_\mathsf{S}$  enhancement by separating hadrons produced in hard processes (jets) from hadrons produced in soft processes (underlying event).

Is the baryon-to-meson ratio enhanced due to the collective effects in the plasma (parton recombination, radial flow,. . . ) or is it (also) due to a modification of the jet fragmentation in the medium?

- $\blacktriangleright$  jet fragmentation A high- $p_T$  parton from hard scattering fragments into hadrons.
- $\blacktriangleright$  parton recombination Multiple partons cluster together to form a hadron.



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# ALICE

- ► collisions studied: p–p at  $\sqrt{s} = 7$  TeV, p–Pb at  $\sqrt{s_{\sf NN}} = 5.02$  TeV, Pb–Pb at  $\sqrt{s_{NN}} = 2.76$  TeV
- ▶ tracking of charged particles by ITS & TPC in magnetic field of 0.5 T
- centrality estimated from the multiplicity of charged particles in the detectors at forward and backward pseudorapidities



# Analysis of charged jets

- $\blacktriangleright$  track selection
	- $\blacktriangleright$  charged primary particles
	- $\blacktriangleright$   $p_{\text{T}}^{\text{track}} > 150$  MeV/c
	- $\blacktriangleright$  uniform in  $\phi \times η$ ,  $|\eta_{\text{track}}|$  < 0.9
- $\blacktriangleright$  raw-jet reconstruction
	- anti- $k_t$  algorithm
	- resolution parameter  $R = 0.2$ , (0.3, 0.4)
- $\triangleright$  subtraction of average soft background
	- **Ex** average background density  $\rho$  estimated from the median  $k_t$  cluster
	- $\blacktriangleright$   $p_{\text{T}}^{\text{jet,ch,corr}} = p_{\text{T}}^{\text{jet,ch,raw}}$ (where  $A_{\text{jet,ch}}$  is jet area)
- $\triangleright$  signal-jet selection (good candidates for hard scattering)
	- $\blacktriangleright$   $p_T$ <sup>leading track</sup>  $>$  5 GeV/c (only Pb–Pb)
	- $\blacktriangleright$   $A_\mathsf{jet,ch}>0.6$ π $R^2$
- In further  $p_T^{\text{jet,ch}}$  corrections
	- **background anisotropy (intra-event**  $p<sub>T</sub>$  **fluctuations)**
	- $\blacktriangleright$  detector response

#### Jet spectra



Larger  $R \Rightarrow$  harder spectrum (but softer jets at a given  $p^{\text{jet}}_{\text{T}}$ ).

# Analysis of neutral strange particles

Strange neutral particles decaying into two charged daughter particles

- $\blacktriangleright$  meson K<sub>S</sub><sup>0</sup> → π<sup>+</sup> + π<sup>-</sup> (BR 69%)
- $\blacktriangleright$  baryon  $\Lambda \to p + \pi^-$  (BR 64%)

Mother  $V^0$  particle reconstructed using topology of its V-shaped decay.

Combinatorial background suppressed by cuts on decay parameters. Signal yield extracted from the invariant-mass distribution.



# Strange particles in jets

Analysis steps

- $\blacktriangleright \; \mathsf{V}^0$  candidate selection
- **D** candidate-jet matching ( $V^0$ s in jet cones)

$$
\sqrt{(\phi_{\mathsf{V}^0} - \phi_{\mathsf{jet},\mathsf{ch}})^2 + (\eta_{\mathsf{V}^0} - \eta_{\mathsf{jet},\mathsf{ch}})^2} < R,
$$

$$
|\eta_{\text{jet,ch}}|^{\text{max}} < |\eta_{\text{V}^0}|^{\text{max}} - R
$$

- **Exerc**idate–UE matching ( $V^0$ s in events without selected jets with  $p_{\text{T}}^{\text{jet,ch}} > 5 \,\text{GeV/c}$ )
- $\triangleright$  signal extraction (invariant-mass distribution)
- efficiency correction (in jet cones, in UE)
- $\blacktriangleright$  subtraction of  $V^0$ s in UE
- ightharpoonup subtraction of  $V^0$ s coming from decays of jet constituents ( $\Xi \rightarrow \Lambda$ ), i.e. "feed-down" correction



# Estimation of  $V^0$ s in the underlying event

- $\blacktriangleright$  no-jet events:  $V^0$ s in events with no selected jets
- outside cones:  $V^0$ s outside jet cones
- **P** random cones:  $V^0$ s in a randomly oriented cone
- **IF** median-cluster cones:  $V^0$ s in the cone of the median  $k_t$ -cluster
- $\blacktriangleright$  perpendicular cones:  $V^0$ s in cones perpendicular to the jet in azimuth

Methods differ in regions, events, statistics, efficiency.



# Reconstruction efficiency of  $V^0$  particles

- Reconstruction efficiency depends strongly on  $p_T^{\vee}$  and  $\eta_{V^0}$ .
- **In Shape of the measured**  $\eta_{\mathcal{N}^0}$  distribution depends on the selection criteria.
- ► Not enough statistics to apply efficiency correction in 2D  $(p_T^{V^0} \times \eta_{V^0})$ .

 $\Rightarrow$  Efficiency of inclusive  $\mathsf{V}^0$ s is scaled (in 2D) to get efficiency in jet cones and UE (in 1D).



#### Feed-down in jets

Feed-down fraction of Λ in jets estimated from:

- $\blacktriangleright$  inclusive  $\Lambda$  (Pb–Pb-like),
- iets generated by PYTHIA 8 (p-p-like).



### Estimation of systematic uncertainties

The systematic uncertainties are studied for the following sources:

- reconstruction efficiency of  $V^0$ s (selection cuts applied on  $V^0$  candidates),
- $\triangleright$  signal extraction (fitting parameters),
- ightharpoonup subtraction of spectra of  $V^0$ s in UE (5 methods),
- $\triangleright$  subtraction of feed-down in jets (inclusive vs PYTHIA),
- $\triangleright$  material budget (detector model),
- $\blacktriangleright$  fluctuations of UE (jet embedding).

# Open issue: Λ–Λ asymmetry

Discrepancy between inclusive spectra of  $\Lambda$  and  $\overline{\Lambda}$ .



Strong dependence on the polarity of magnetic field and the sign of *η*. Additional 6 % (symmetric) considered as systematics.

#### Open issue: discrepancy in MC between runs 2010, 2011

Differences in spectra traced back to the reconstruction efficiencies.



Effects of  $\lesssim$  10 % for K $^0_\mathsf{S}, \lesssim$  20 % for Λ, partially cancel out in Λ/K $^0_\mathsf{S}.$ Observed also in other analyses (charged-particle spectra, correlations). Additional 10 % (symmetric) considered as systematics.

#### Systematics: combined

$$
R = 0.2, p_T^{\text{jet,ch}} > 10 \,\text{GeV/c}, \, (\Lambda + \overline{\Lambda})/2 \text{K}_\text{S}^0
$$



Uncertainties from different sources (except feed-down) are combined in squares and considered symmetric.

#### V 0 vertexer problem

 $\blacktriangleright$  There may be two candidates for the point of closest approach ("cowboy/sailor" configuration).



- Old vertexer: sailor misidentified as cowboy  $\rightarrow$  CPA  $\approx$  -1  $\rightarrow$  rejected.
- New vertexer: Select the point with the smallest DCA calculated in 3D.



- $\Rightarrow$  Better MC–real data matching.
- Cause of loosing sailors with the old vertexer still unclear.

#### $\mathsf{\Lambda}/\mathsf{K}^0_\mathsf{S}$  $\frac{0}{\mathsf{S}}$  ratio in jets in p–p at  $\sqrt{\mathsf{s}} = 7$  TeV and 13 TeV



- The ratio in UE is consistent with the inclusive ratio.
- $\triangleright$  The ratio in jets is clearly different from the inclusive ratio at low and intermediate  $p_T^{\vee^0}$ .
- A slight increase of the ratio in jets with increasing  $R$ .

#### $\mathsf{\Lambda}/\mathsf{K}^0_\mathsf{S}$  $\frac{0}{\mathsf{S}}$  ratio in jets in p–Pb at  $\sqrt{s_\mathsf{NN}} = 5.02\,\mathsf{TeV}$ (high-multiplicity collisions, 0–10 %)



The ratio in jets

- is clearly different from the inclusive ratio at low and intermediate  $p_T^{\vee}$ ,
- $\triangleright$  is different from the inclusive ratio in PYTHIA (black line),
- $\triangleright$  is similar to the ratios in PYTHIA jets (red dashed lines),
- In shows no significant dependence on  $p_T^{\text{jet,ch}}$  and a slight dependence on R.

 $\mathsf{\Lambda}/\mathsf{K}^0_\mathsf{S}$  $\frac{0}{\mathsf{S}}$  ratio in jets in Pb–Pb at  $\sqrt{s_\mathsf{NN}} = 2.76$  TeV  $(7.4 \times 10^6$  central collisions, 0–10 %)



The ratio in jets

- is clearly different from the inclusive ratio at low and intermediate  $p_T^{\vee}$ ,
- ightharpoonup shows no significant dependence on  $p_T^{\text{jet}, \text{ch}}$ ,
- is consistent with the ratio in jets in p-Pb and p-p at  $p_T^{V^0} > 4$  GeV/c.

#### $\mathsf{K}^0_\mathsf{S}$  $<sup>0</sup>$ , Λ spectra in jets in Pb–Pb</sup> comparison to PYTHIA smeared with  $p_{\text{T}}^{\text{jet},\text{ch}}$  fluctuations  $p_{\textsf{T}}^{\textsf{jet},\textsf{ch}}>10\,\textsf{GeV/c}$  $p_{-}$  (GeV/c) 2 3 4 5 6 7 8 9 10 ۱/(N<sub>jets</sub>πR<sup>e</sup>) dWdp<sub>r</sub> (c/GeV)<br> <sup>−</sup><sup>3</sup> 10  $10^{-2}$ <sup>−</sup><sup>1</sup> 10 1 F  $K_s^0$ , stat. unc., (x 1.5)  $(\overline{\Lambda}+\overline{\Lambda})/2$ , stat. unc. syst. unc. full markers  $\mathsf{K}_{\mathrm{s}}^0$ ,  $(\mathrm{x}\,1.5)$ open markers (Λ+Λ)/2 PYTHIA 8 - tune Monash PYTHIA 6 - tune Perugia 2011 PYTHIA 6 - tune Perugia NoCR  $\frac{1}{\tau}$ smeared with true  $\sigma(\delta p)$ T  $\rho_+^{\text{\tiny{per}}}$  :  $p_{\tau}^{\rm jet, ch}$  > 10 GeV/ $c$ T ALICE Preliminary  $\frac{\text{track}}{\tau} > 150 \text{ MeV}/c$  $\rho_{_{\,{\sf T}}}^{\rm ^{wave}}$  $> 5$  GeV/ $c$ leading track  $p_{T}^{max}$  $| < 0.5$ jet,ch  $|\eta|$ anti- $k_t$ ,  $R = 0.2$  $|\eta_{\gamma\rho}| < 0.7$ Pb−Pb, s NN = 2.76 TeV, 0−10 %  $p_{\text{T}}^{\text{jet,ch}} > 20 \text{ GeV/c}$  $p_{n}$  (GeV/ $c$ 2 3 4 5 6 7 8 9 10 ۱/(N<sub>|ets</sub>πR<sup>e</sup>) dWdp<sub>T</sub> (c/GeV)<br> 10 1  $10^{-2}$ <sup>−</sup><sup>1</sup> 10 1F stat. unc.,  $(x 1.5)$ K<sup>α</sup>, <u>st</u>at. unc., (x 1.5<br>(Λ+Λ)/2, stat. unc. syst. unc. tull markers K;, (x 1.5)<br>open markers (Λ+Λ)/2 full markers  $K_s^0$ , (: PYTHIA 8 - tune Monash PYTHIA 6 - tune Perugia 2011 PYTHIA 6 - tune Perugia NoCR  $\frac{1}{\tau}$ smeared with true  $\sigma$ (δρ $'$ T  $\rho^{\text{\tiny jet}}_{\scriptscriptstyle +}$  : ALICE Preliminary  $p_{\perp}^\mathrm{jet,ch}$  > 20 GeV/c T  $\frac{\text{track}}{\text{r}} > 150 \text{ MeV}/c$  $\rho_{_{\rm T}}^{\rm{max}}$  > 5 GeV/c leading track  $p_T^{\text{max}}$  $| < 0$ jet,ch  $|\eta|$ anti- $k_t$ ,  $R = 0.2$  $|\eta_{\rm v0}| < 0.7$ Pb−Pb, s NN = 2.76 TeV, 0−10 %

**ALI−PREL−112798**

**ALI−PREL−112802**

- ▶ Same slopes of spectra from measurement and from PYTHIA.
- **E**nhancement for  $\Lambda$  at  $p_T^{\vee^0} < 4$  GeV/c.

## Summary and outlook

ALICE has performed the first measurement of the  $\Lambda/\mathrm{K^0_S}$  ratio in charged jets in p–p, p–Pb and Pb–Pb collisions at the LHC.

Results

- $\blacktriangleright$  In every collision system, the  $\Lambda/K_S^0$  ratio in jets is significantly smaller than the inclusive ratio (and the UE).
- $\blacktriangleright$  The  $\Lambda/K_S^0$  ratios in jets are consistent within uncertainties in all collision systems for  $p_{\text{T}}^{\text{V}^0} > 4 \text{ GeV/c}.$
- $\blacktriangleright$  The dominant source of the enhancement are soft processes associated with collective behaviour.
- $\triangleright$  A potential modification of jet fragmentation seems to be restricted to the region  $p_T^{\vee}$ <sup>0</sup> < 4 GeV/c and manifest by an enhancement of the  $\Lambda$  yields.

Outlook

- $\triangleright$  Solve the 2011/2010 issue.
- $\triangleright$  Comparison with more models (JEWEL).

# Thank you for your attention.

# Backup

p/π ratio in Au–Au at  $\sqrt{s_\text{NN}} = 200$  GeV



Phys. Rev. Lett. **97** (2006) 152301

 $\overline{\Xi}/\overline{\Lambda}$  ratio in jets in p–p at  $\sqrt{s_\text{NN}}=13\,\text{TeV}$ 



Pengyao Cui for the ALICE Collaboration, Quark Matter 2018

### Details on the  $\Lambda/\overline{\Lambda}$  discrepancy



# $V^0$  candidate selection



#### Jet algorithms

A sequential recombination jet finder is defined according to this general scheme:

1. ∀ *i*, *j* : calculate distances  $d_{ij}$  and  $d_{iB}$  (NB  $k_t \equiv p_T$ ):

$$
d_{ij} = \min \left( k_{\text{t},i}^{2p}, k_{\text{t},j}^{2p} \right) \frac{\Delta_{ij}^2}{R^2}, \quad \Delta_{ij}^2 = (y_i - y_j)^2 + \left( \phi_i - \phi_j \right)^2, \quad d_{iB} = k_{\text{t},i}^{2p}
$$

2. Find  $d_{\text{min}}$ :

$$
d_{\min}=\min\left(d_{ij},d_{iB}\right).
$$

- ► If  $\exists i, j : d_{\text{min}} = d_{ij}$ , merge particles *i* and *j* into a single particle and combine their momenta.
- If  $\exists i$  :  $d_{\text{min}} = d_{iB}$ , declare particle *i* to be a final jet and remove it from the list.

These steps are repeated until no particles are left.

$$
p = \left\{ \begin{array}{cl} 1 & k_t \text{ (background estimation)} \\ 0 & \text{Cambridge/Aachen} \\ -1 & \text{anti-}k_t \text{ (signal jets)} \end{array} \right.
$$



Matteo Cacciari et al. JHEP **0804** (2008) 063

# Background in Pb–Pb

Production of soft particles by underlying-event processes.

average background density *ρ*:

 $\blacktriangleright$   $k_t$  jets w/o 2 hardest

each event:  $\rho =$  median  $\left\{ \rho_T^{\rm jet}/A_{\rm jet} \right\}$ each jet:  $\rho_{\mathsf{T},\mathsf{jet}}^{\mathsf{corrected}} = \rho_{\mathsf{T},\mathsf{jet}}^{\mathsf{raw}} - \rho \mathcal{A}_{\mathsf{jet}}$ 



*ρ* anisotropy in events (fluctuations):

$$
\blacktriangleright \delta p_{\mathsf{T}} = p_{\mathsf{T},\mathsf{probe}}^{\mathsf{raw}} - \rho A_{\mathsf{probe}}
$$

response matrix  $\rightarrow$  deconvolution



ibid.

ALICE, JHEP **1203** (2012) 053

### Scaling of the reconstruction efficiency

- $\bullet$   $\epsilon$  reconstruction efficiency of inclusive particles
- $\bullet \epsilon_{s}$  reconstruction efficiency of particles of interest (scaled  $\epsilon$ )
- $\triangleright$  a<sub>s</sub> yield of associated particles of interest
- $g_s$  yield of generated particles of interest
- $\blacktriangleright$  m uncorrected yield of measured particles (candidates) of interest
- $\blacktriangleright$  t yield of true (corrected) particles of interest
- $\blacktriangleright$  P signal purity

Signal extraction in JC, UE (assume that  $P_{\text{inclusive}}(p_{\text{T}}^{\text{V}^0},\eta_{\text{V}^0})$  is the same as for  $\text{V}^0\text{s}$ of interest):

$$
m(\rho_T^{V^0},\eta_{V^0})=m_{\sf raw}(\rho_T^{V^0},\eta_{V^0})|_{\sf peak\ region}\cdot P_{\sf inclusive}(\rho_T^{V^0},\eta_{V^0})|_{\sf peak\ region}
$$

Efficiency calculation:

$$
a_s \equiv m, \quad \sigma_{a_s} \equiv 0, \qquad g_s = a_s/\epsilon
$$

$$
\frac{1}{\epsilon_s(p_T^{\vee 0})} = \frac{\sum_{\eta_{\sqrt{0}}}}^{} g_s(\eta_{\sqrt{0}}^{}, p_T^{\sqrt{0}}^{})}{\sum_{\eta_{\sqrt{0}}^{}, g_s(\eta_{\sqrt{0}}^{}, p_T^{\sqrt{0}}^{})}} = \sum_{\eta_{\sqrt{0}}^{}, \frac{a_s(\eta_{\sqrt{0}}^{}, p_T^{\sqrt{0}}^{})}{\sum_{\eta_{\sqrt{0}}^{}, g_s(\eta_{\sqrt{0}}^{}, p_T^{\sqrt{0}}^{})}} = \frac{1}{(\eta_{\sqrt{0}}^{}, p_T^{\sqrt{0}}^{})}
$$

Spectra correction:

$$
t=m/\epsilon_s
$$