

J/ ψ production in pPb and PbPb collisions by CMS



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for the CMS collaboration



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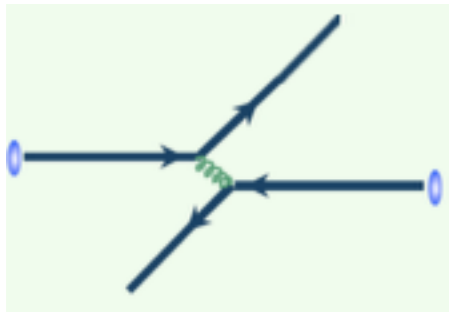
- ④ **Introduction – collision systems**

- ④ **J/ψ in PbPb collisions**
 - Final state effects
 - CMS detector & analysis method
 - Nuclear modification factor R_{AA}

- ④ **J/ψ in pPb(dAu) collisions**
 - Initial state effects
 - Nuclear modification factor R_{pPb} and R_{dAu}
 - Status of J/psi analysis in pPb with CMS

- ④ **System size study – Cu+Au collisions**

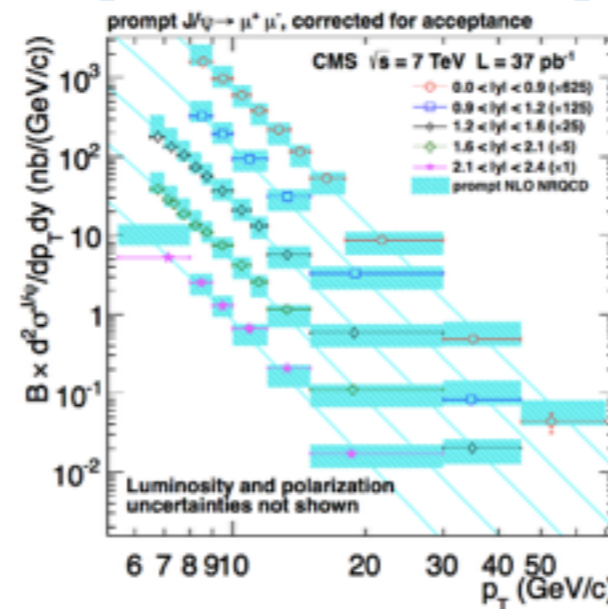
- ④ **Summary**



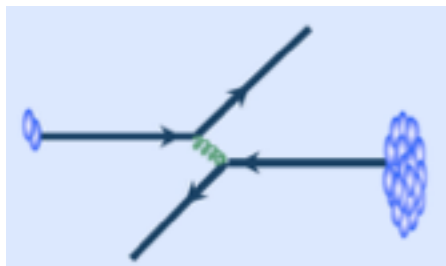
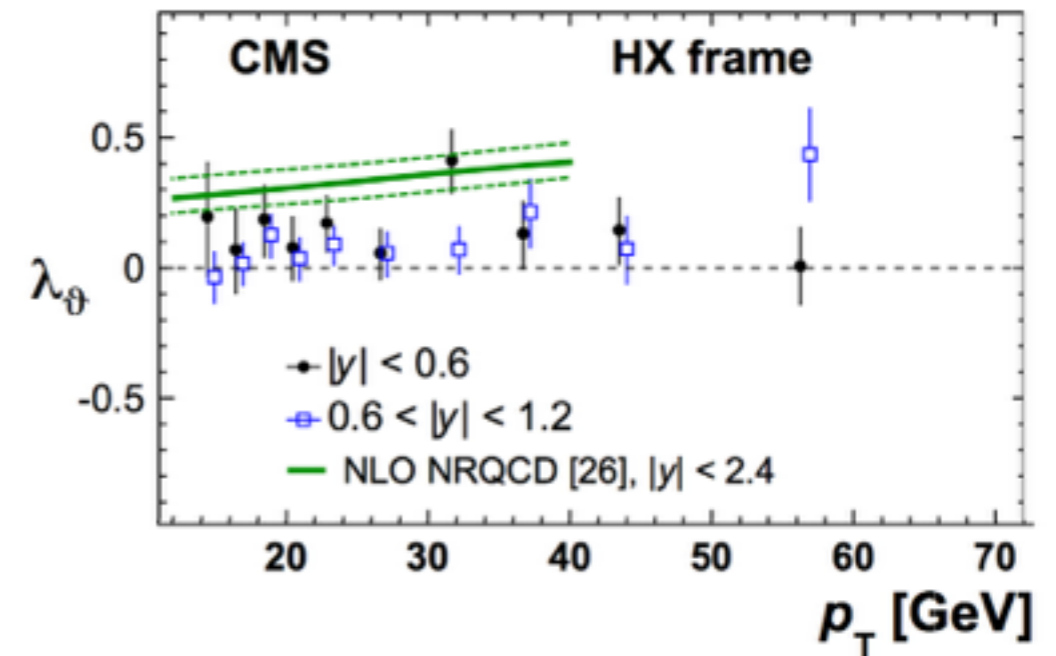
⊗ pp collisions

- No single hadro-production model describes all quarkonia data.

[cross section]

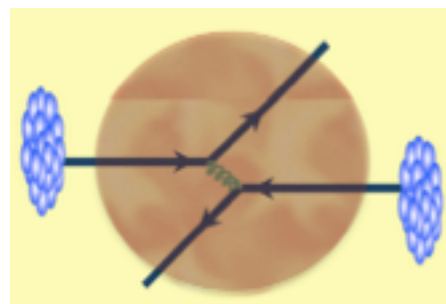


[polarization]



⊗ pPb collisions

- Cold Nuclear Matter effects
- p_T broadening, Initial energy loss, absorption, nPDF, etc.



⊗ PbPb collisions

- Quark-Gluon plasma formation in central collisions (deconfined state of strongly interacting matter)
- Sequential suppression, recombination, etc.

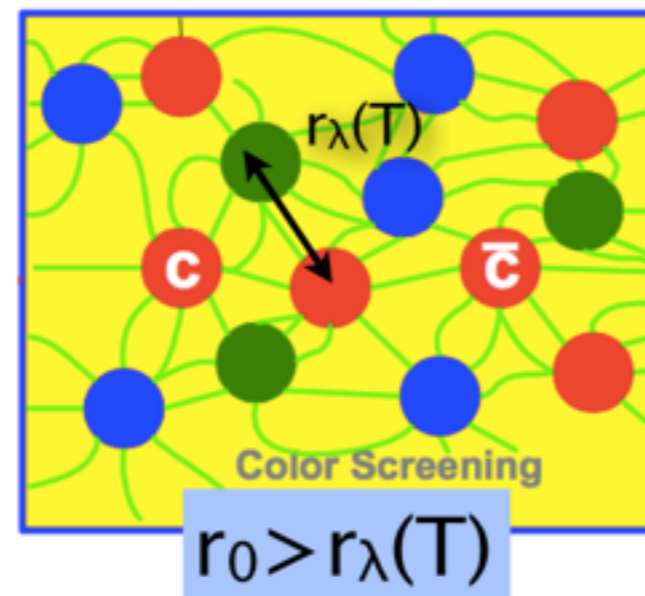
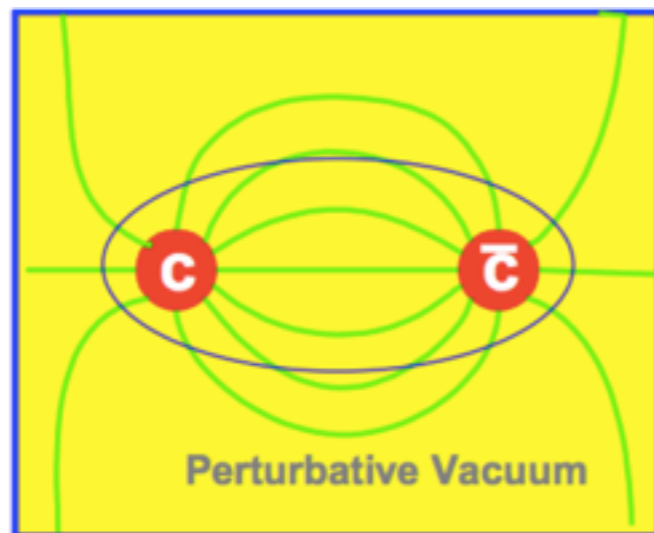
⊕ Quarkonia



- Bound states of heavy quark and antiquark
- Large mass requires a large momentum transfer during the **early stage**.
⇒ Powerful tool to probe QGP

⊕ Debye screening (suppression)

- Loosely bound states (with smaller binding energies) melt at lower temperature.
- Sequential melting of the quarkonia ⇒ Thermometer of QGP

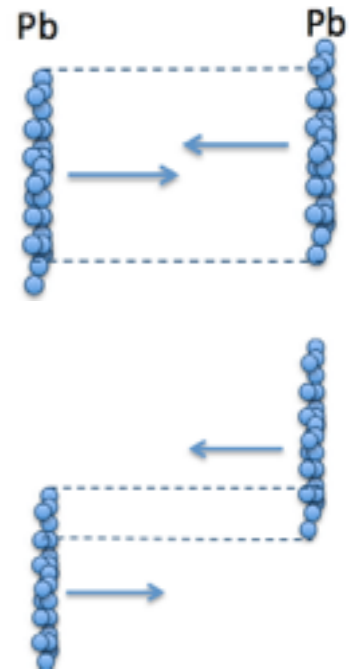


⊕ Recombination (enhancement)

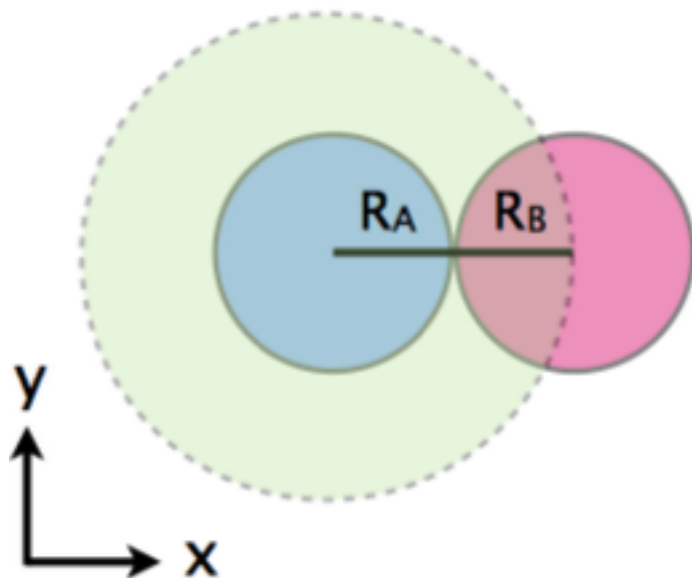
- combination of quarks and antiquarks which are initially produced in “different” nucleon–nucleon collisions.

Centrality

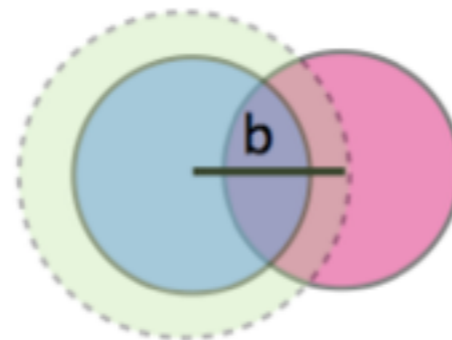
- related to the overlap fraction of the geometrical cross sections
- 0% is the most central, 100% is the most peripheral collisions.
- In CMS experiment – defined by transverse energy deposit in HF



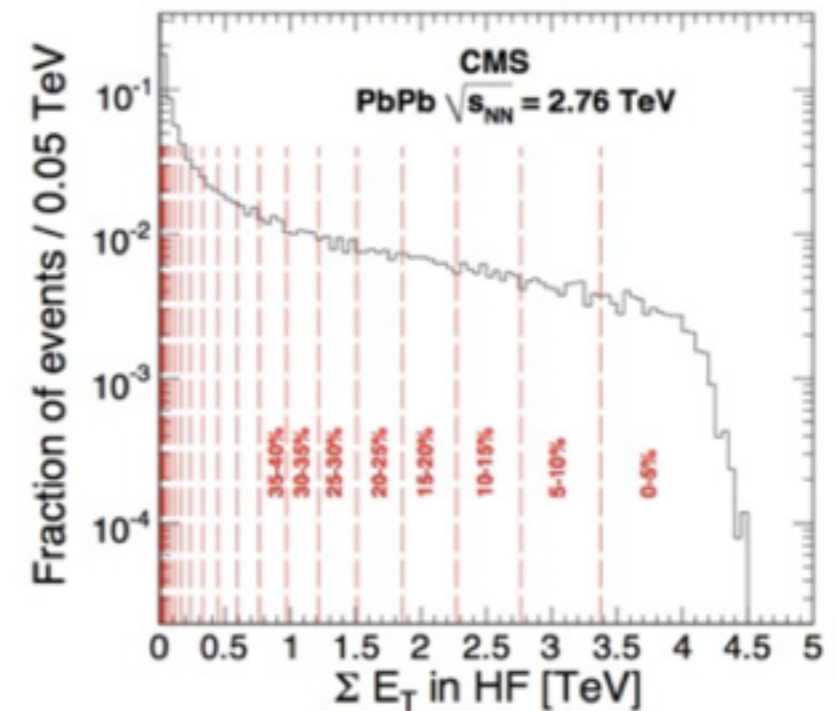
$$[\sigma_{geom} \sim \pi(R_A + R_B)^2]$$

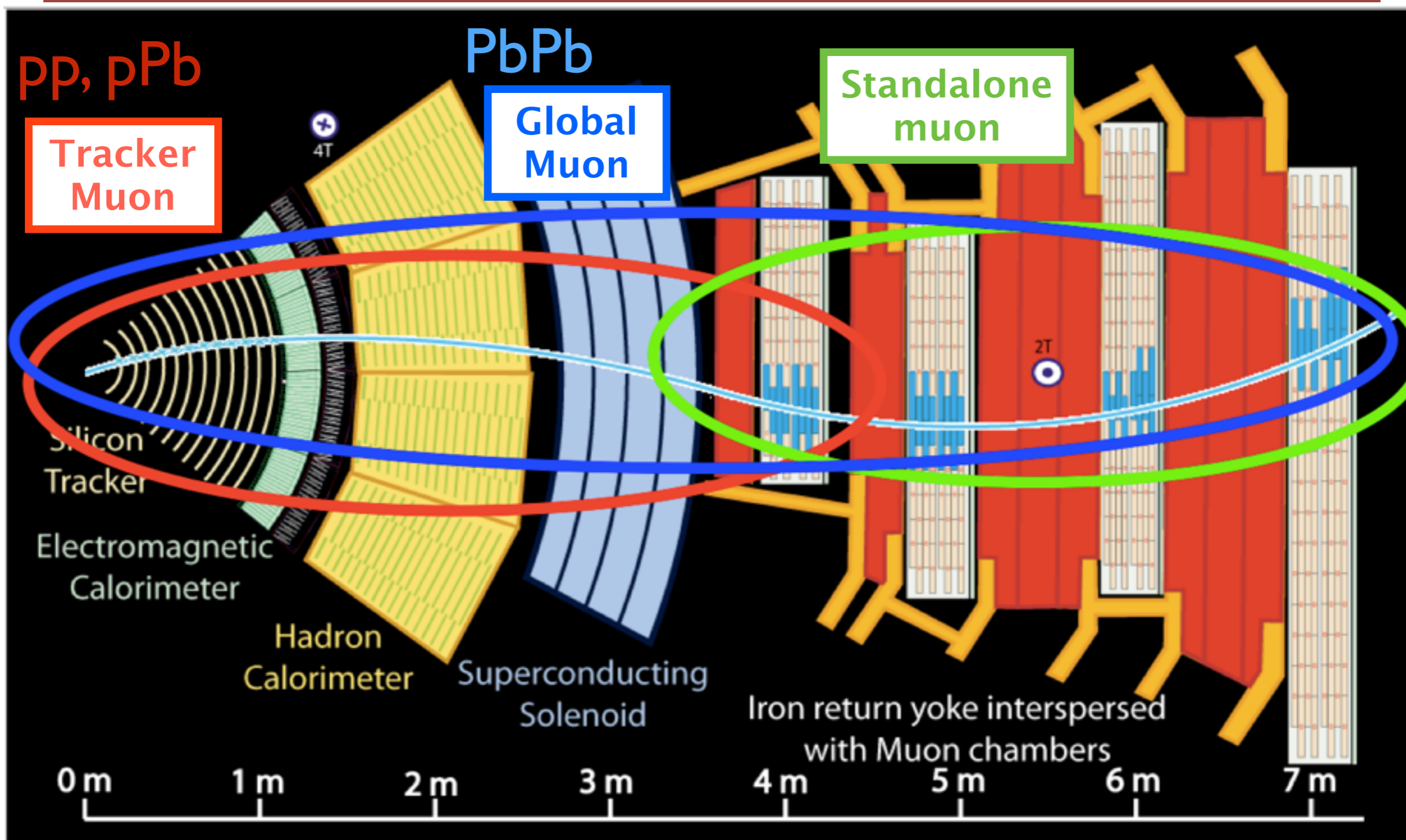


$$[\sigma_{reac} \sim \pi b^2]$$



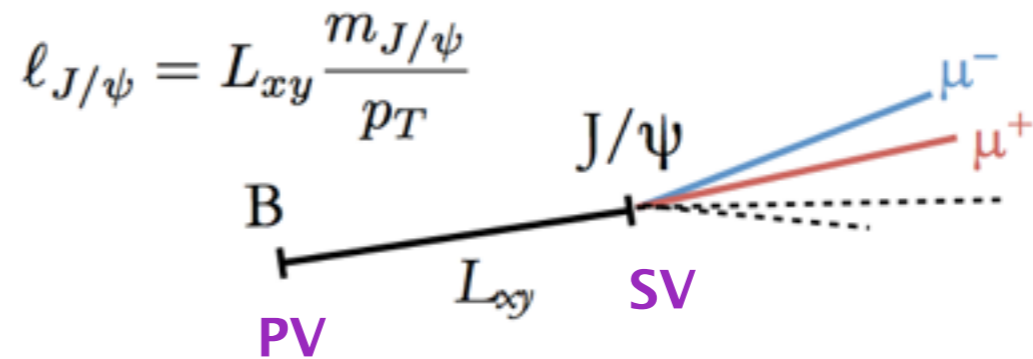
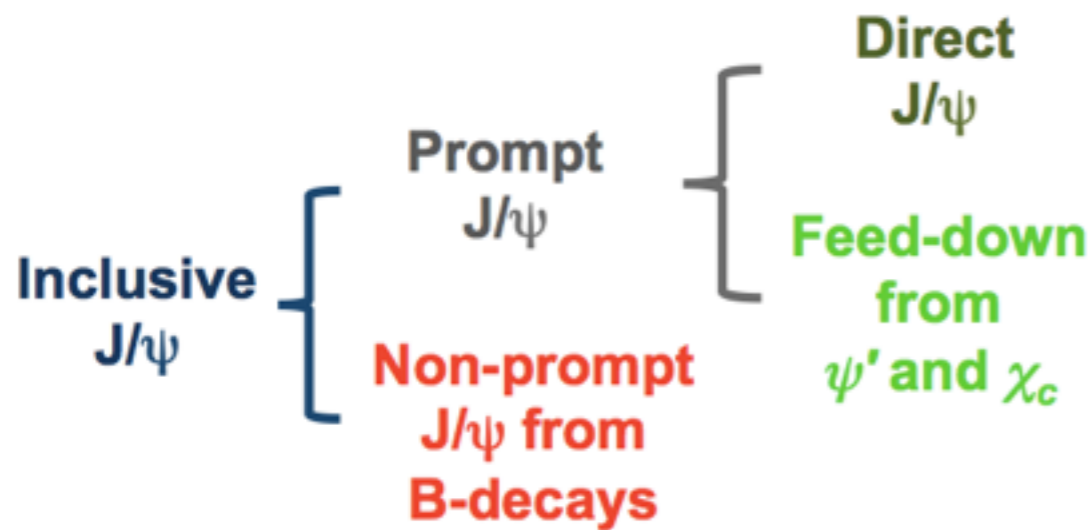
$$centrality(\%) = \frac{\sigma_{reac}}{\sigma_{geom}} \times 100$$



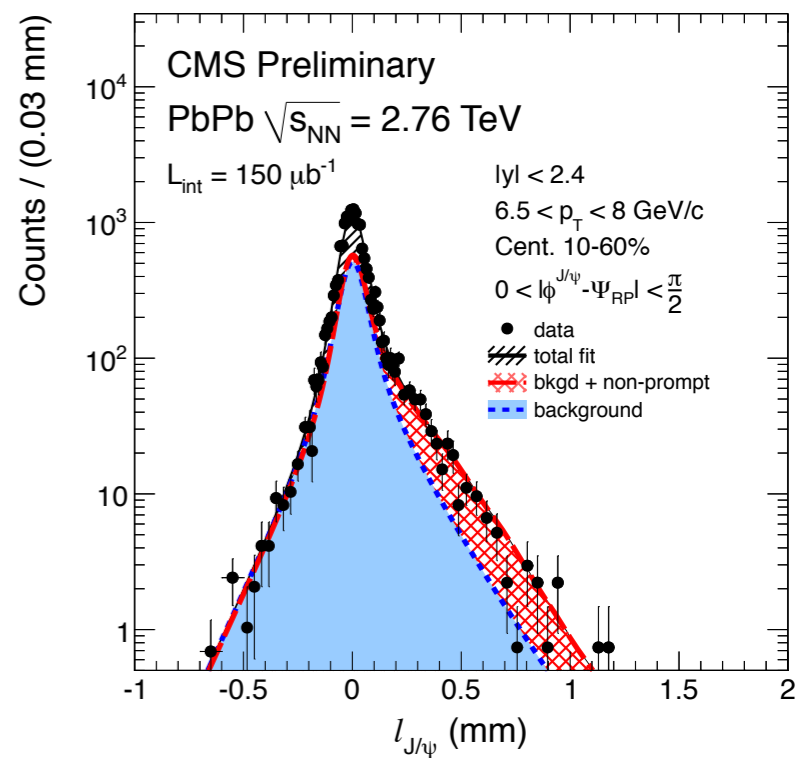
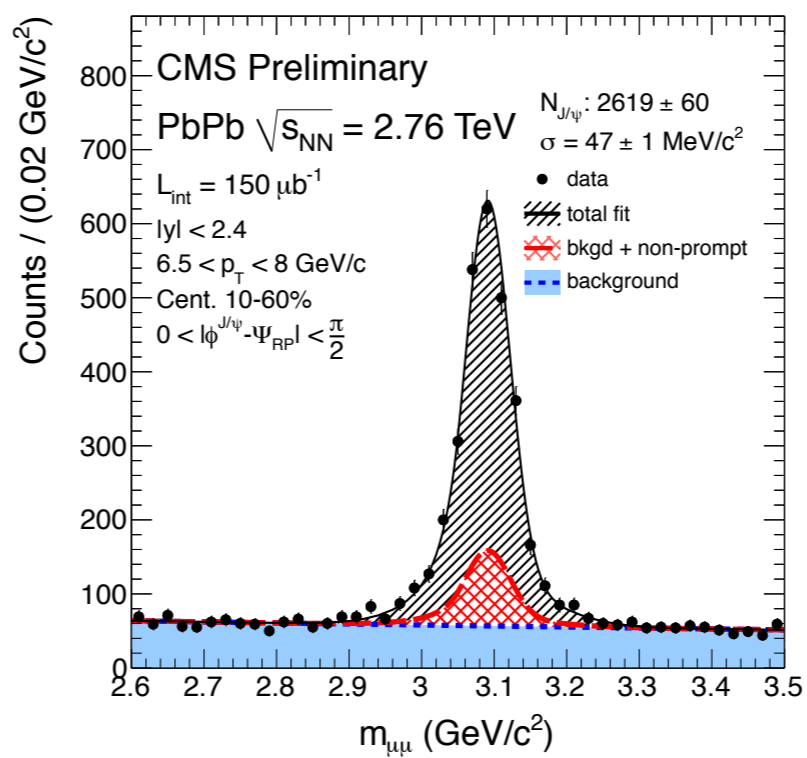
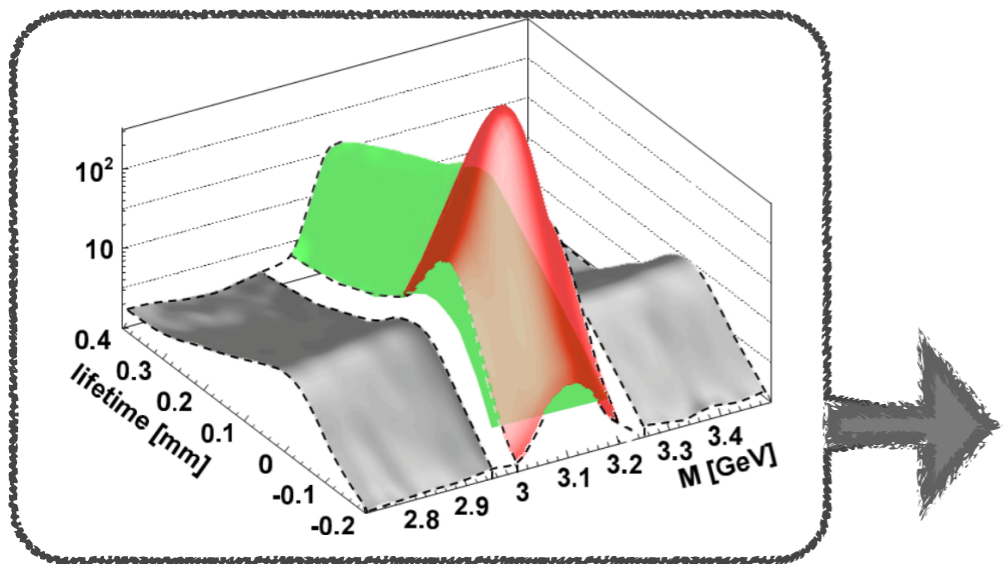


- ⊕ Excellent muon Identification and triggering in the muon system
- ⊕ High momentum and vertex resolution of the tracking system

- ⊕ **Separation of prompt J/ψ and non-prompt J/ψ**
 - 2-Dimensional simultaneous fit for $m_{\mu\mu}$ & $l_{J/\psi}$



lifetime of the b hadrons ($\mathcal{O}(500) \mu\text{m}/c$)



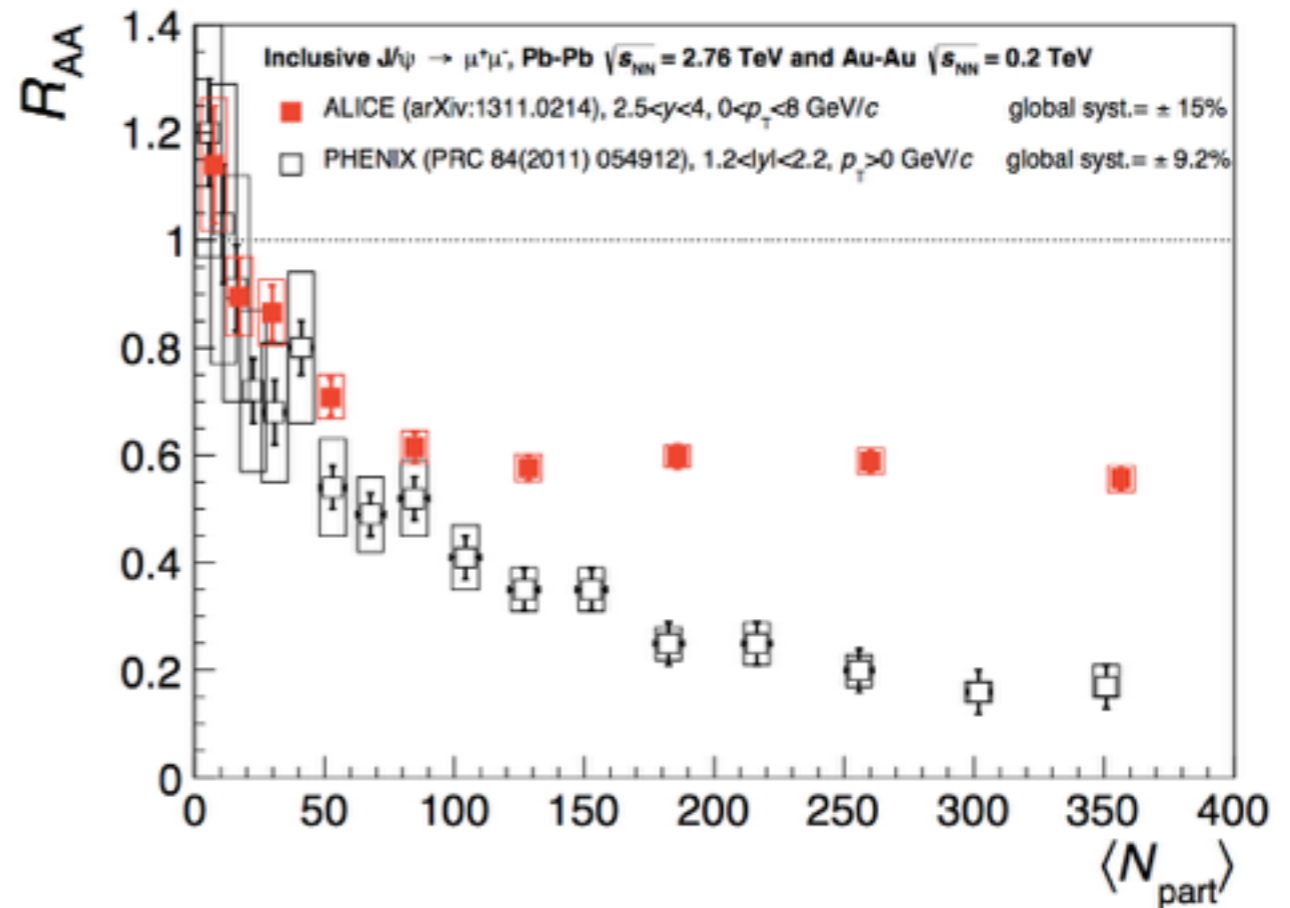
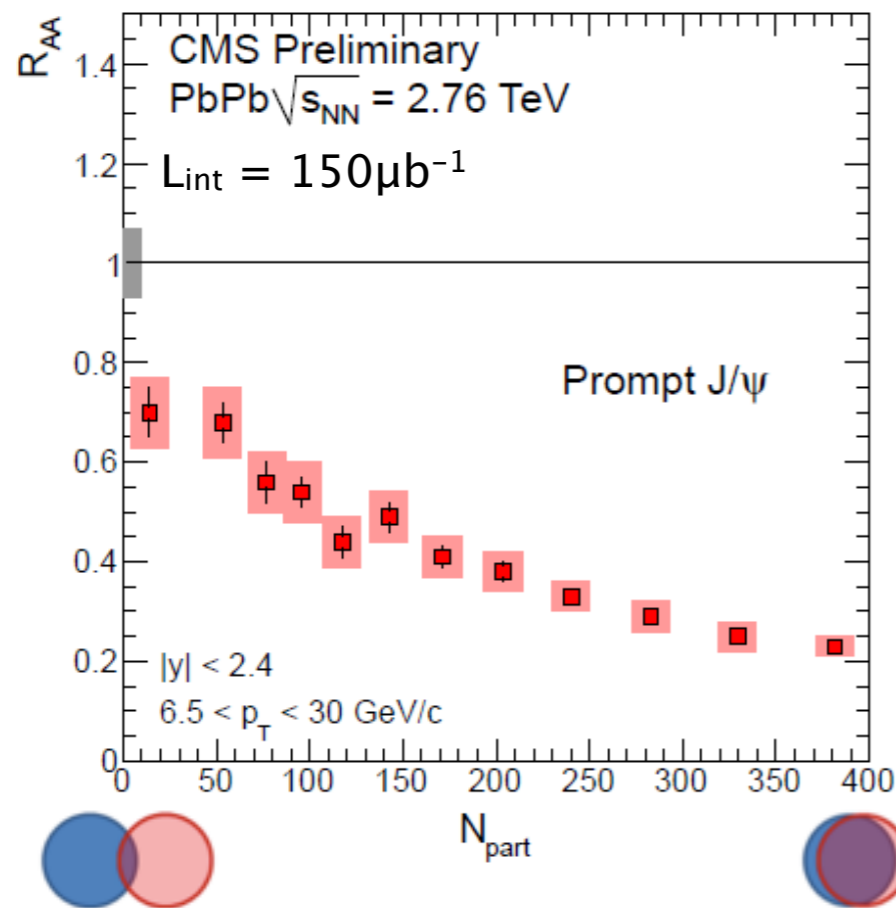
CMS-PAS HIN-12-014
 JHEP 05 (2012) 063

⊕ Nuclear modification factor

$$R_{AA} = \frac{1}{\langle N_{coll} \rangle} \frac{N_{AA}}{N_{pp}}$$

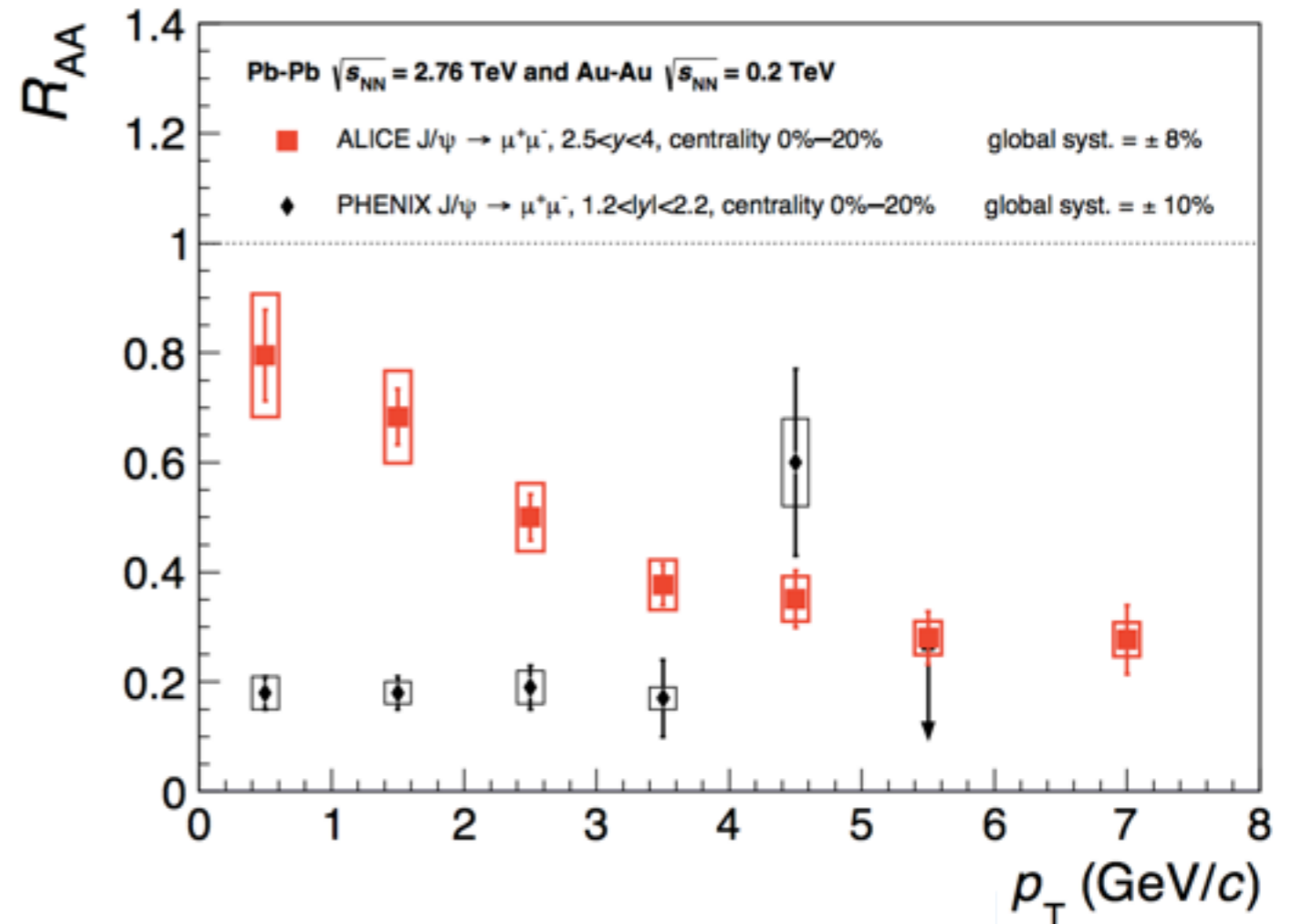
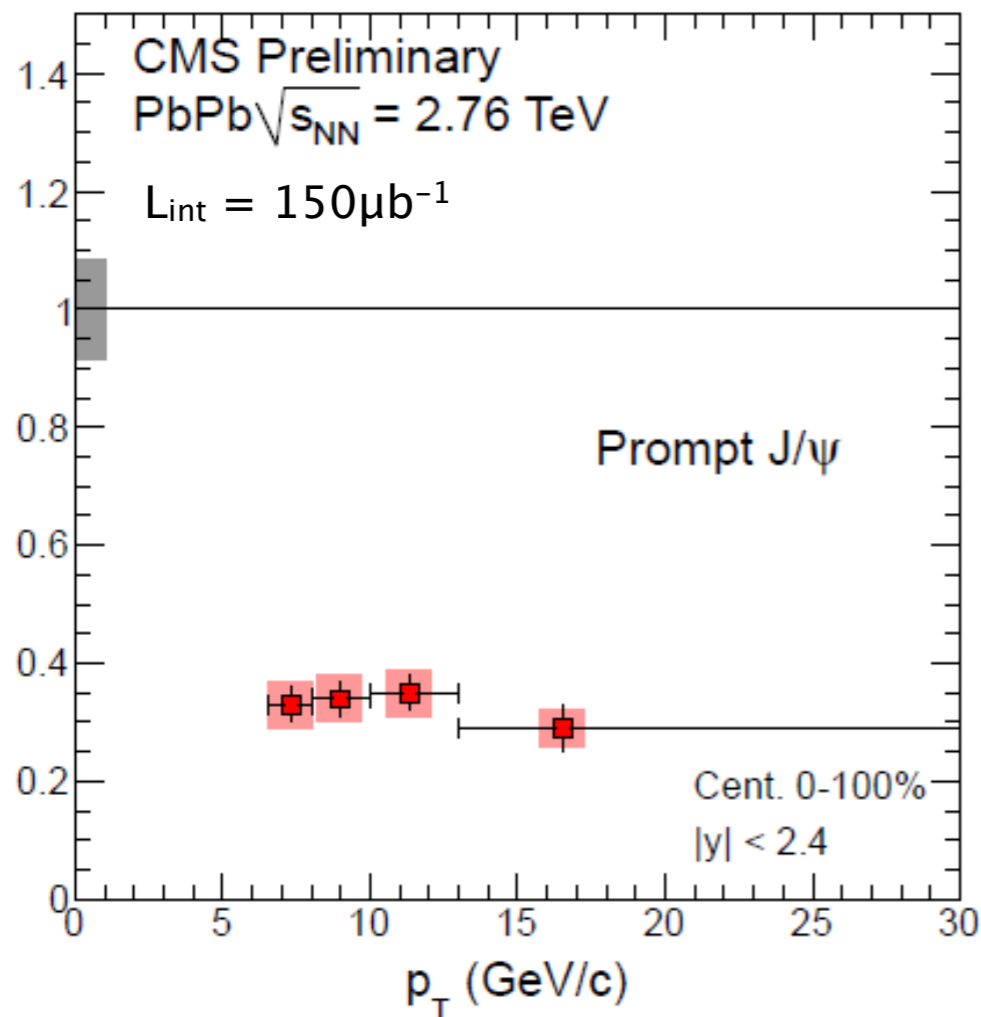
: $R_{AA} = 1$ No modification compared to pp collisions

⊕ Centrality dependence



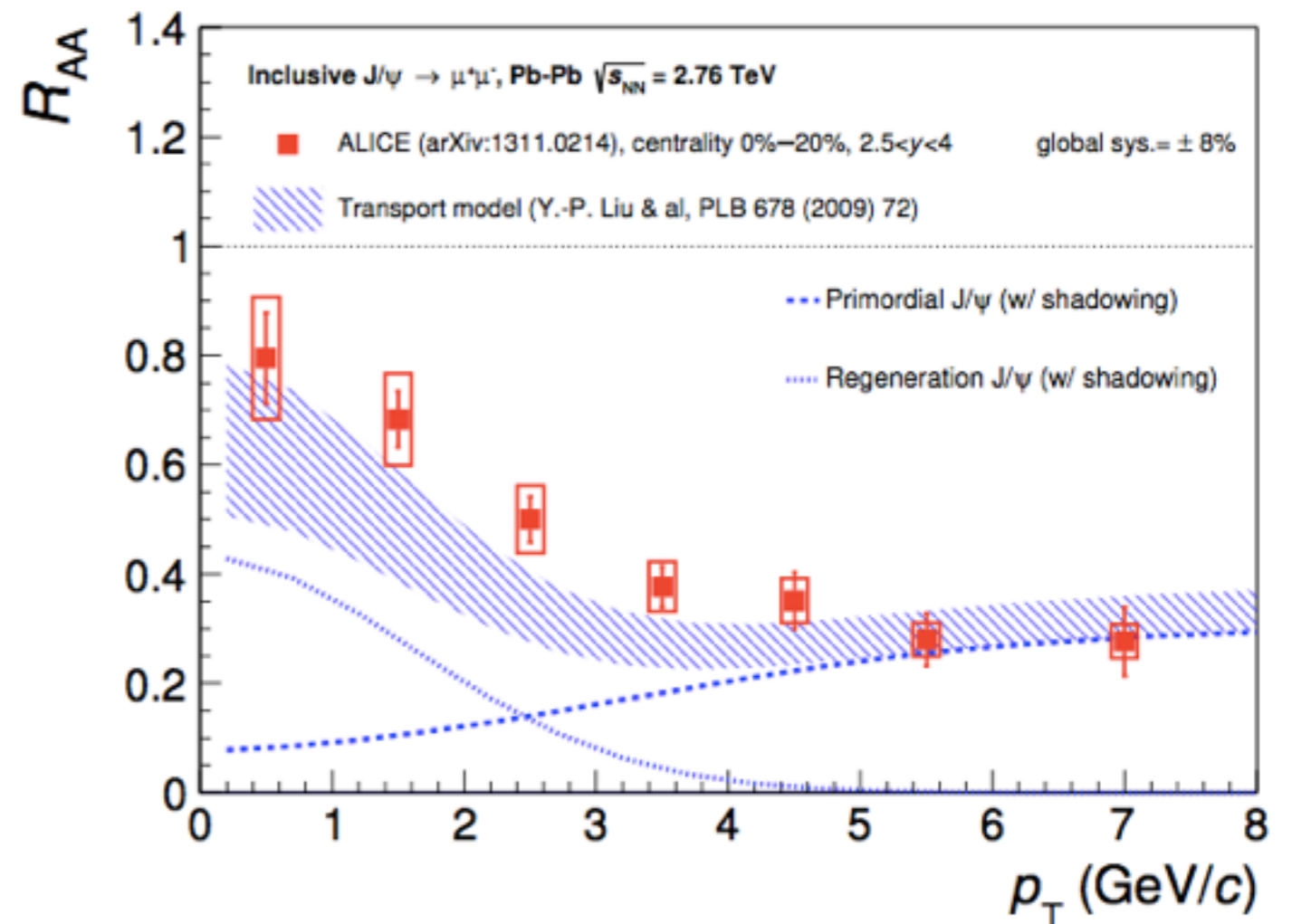
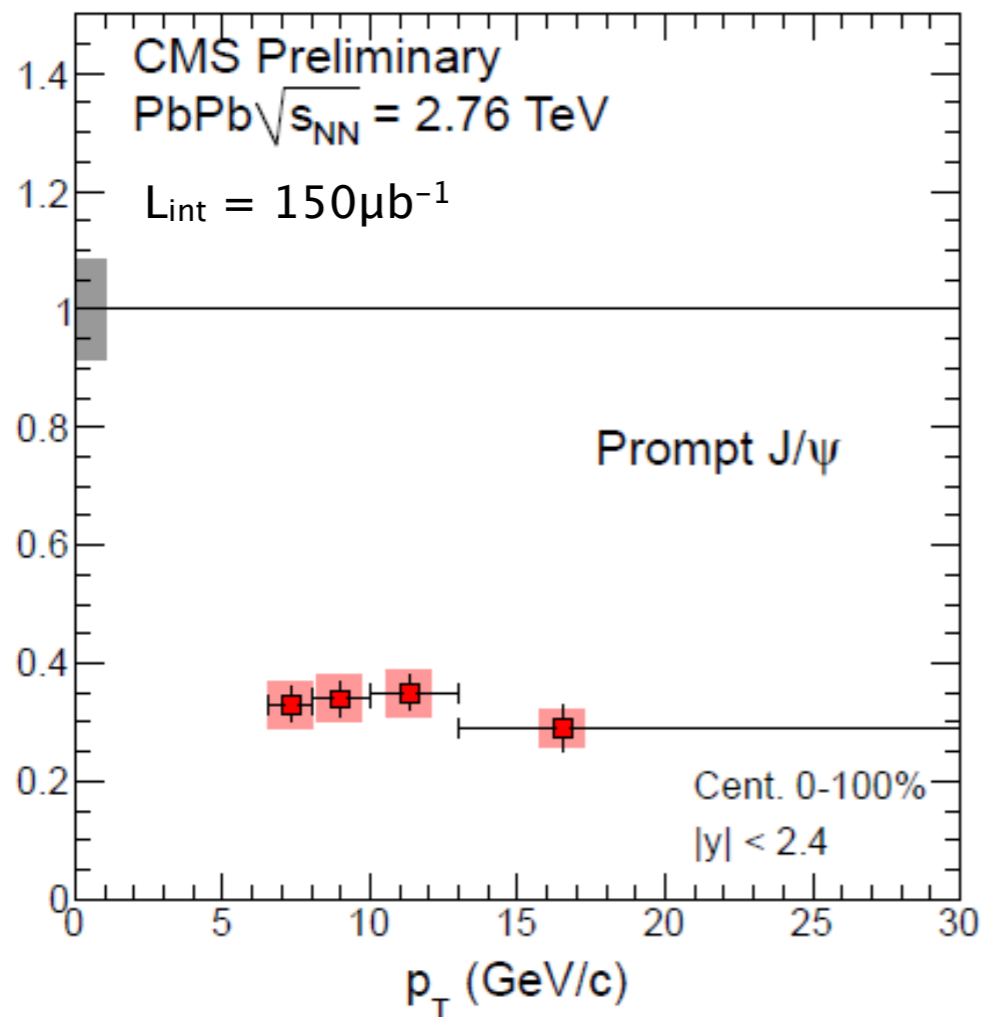
- Different kinematic ranges, different collision energies, prompt vs inclusive

⊕ p_T dependence



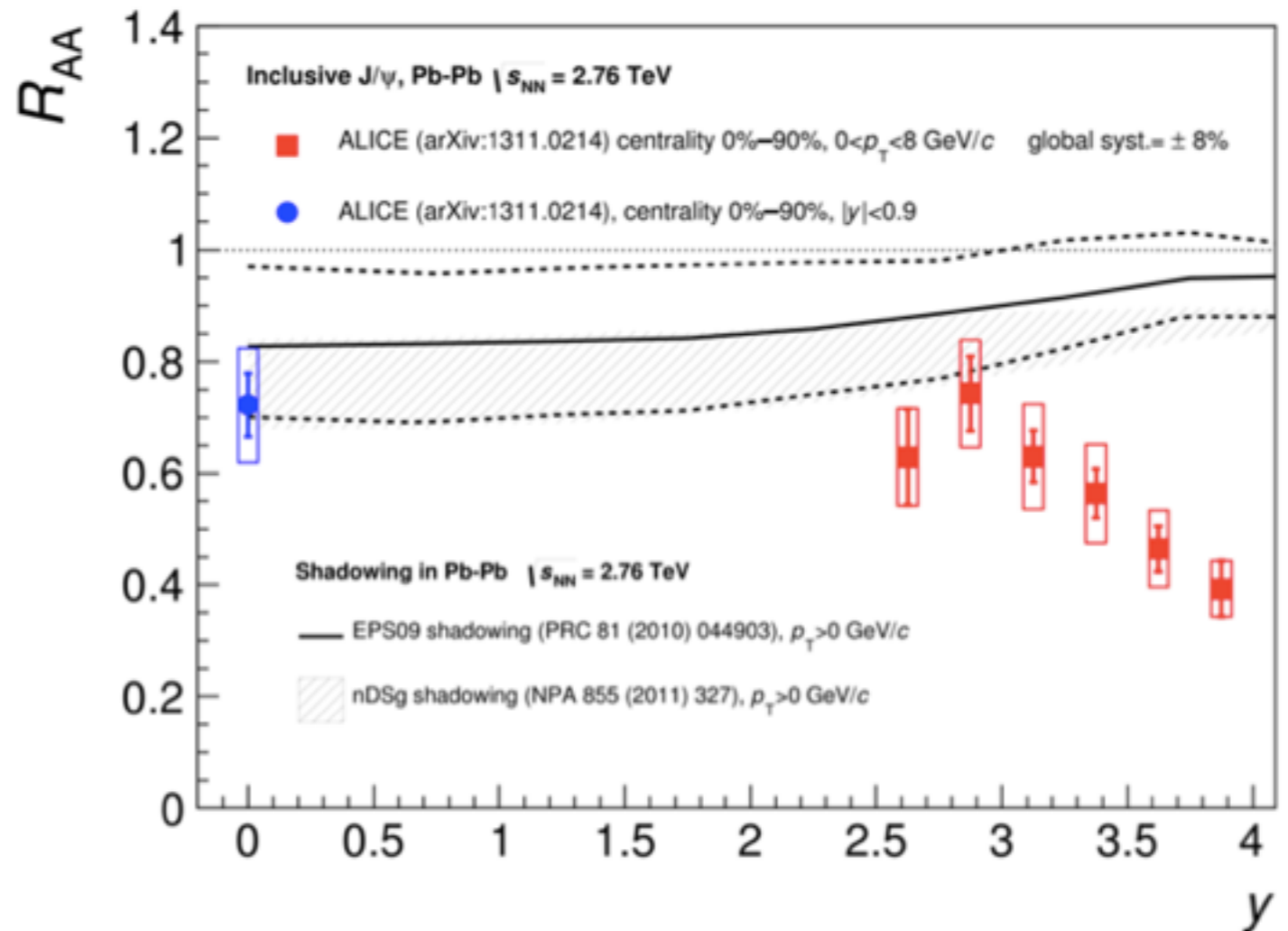
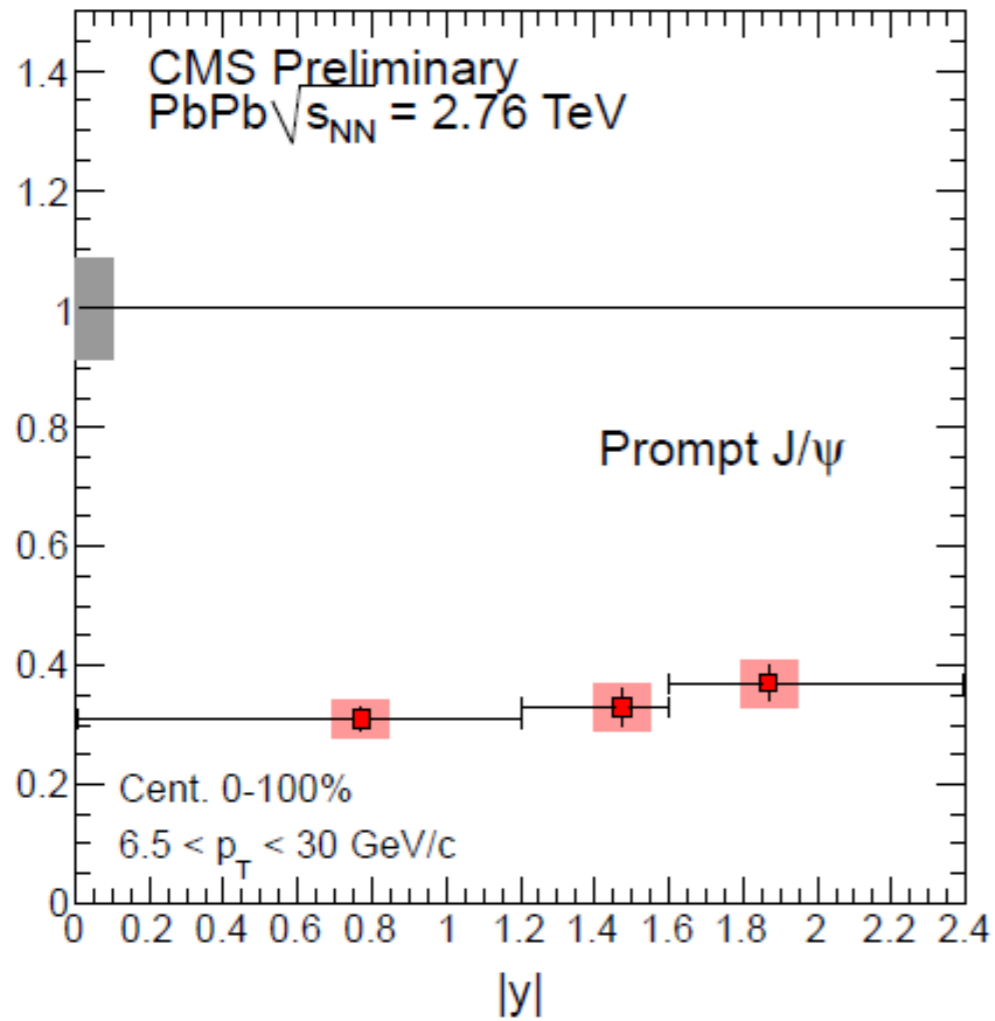
- ⊕ CMS and ALICE show similar suppression in the overlapping p_T range.
- ⊕ ALICE shows less suppression at low p_T compared to PHENIX.

⊕ p_T dependence



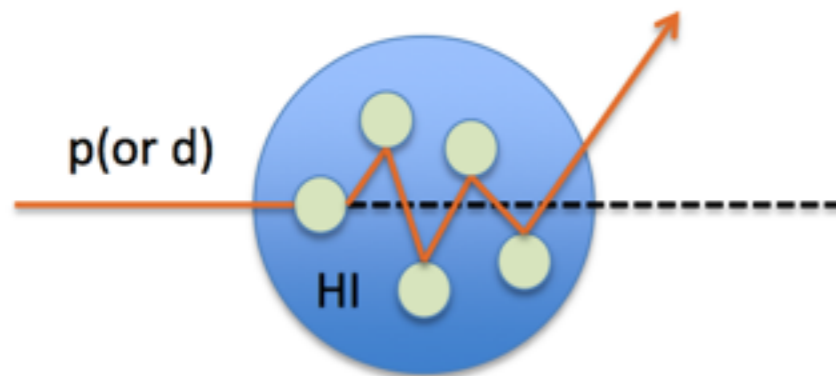
- ⊕ CMS and ALICE show similar suppression in the overlapping p_T range.
- ⊕ ALICE shows less suppression at low p_T compared to PHENIX.
→ Recombination contribution important at low p_T ?

⊕ Rapidity dependence



- ⊕ CMS – No strong rapidity dependence
- ⊕ ALICE – More suppression at forward rapidity
– CNM models fail to describe data.

① p_T broadening

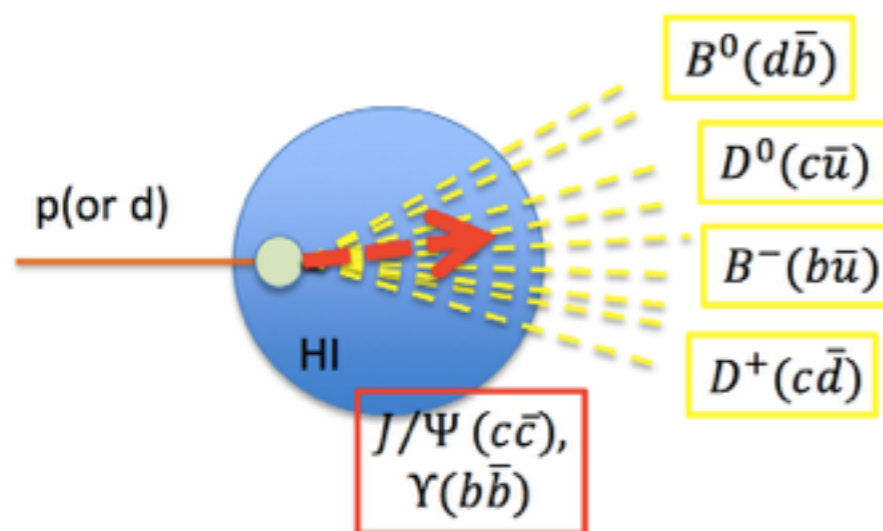


- Before the hard scattering, multiple elastic scatterings can occur.
 - Projectile can get p_T from soft scattering.
 - Yields decrease at low p_T , while increase at intermediate p_T .

② Initial State Energy Loss

- Energy loss by soft scattering reduces the initial energy of hard scattering.
- Yields decrease at forward rapidity, while increase at mid-rapidity regions.

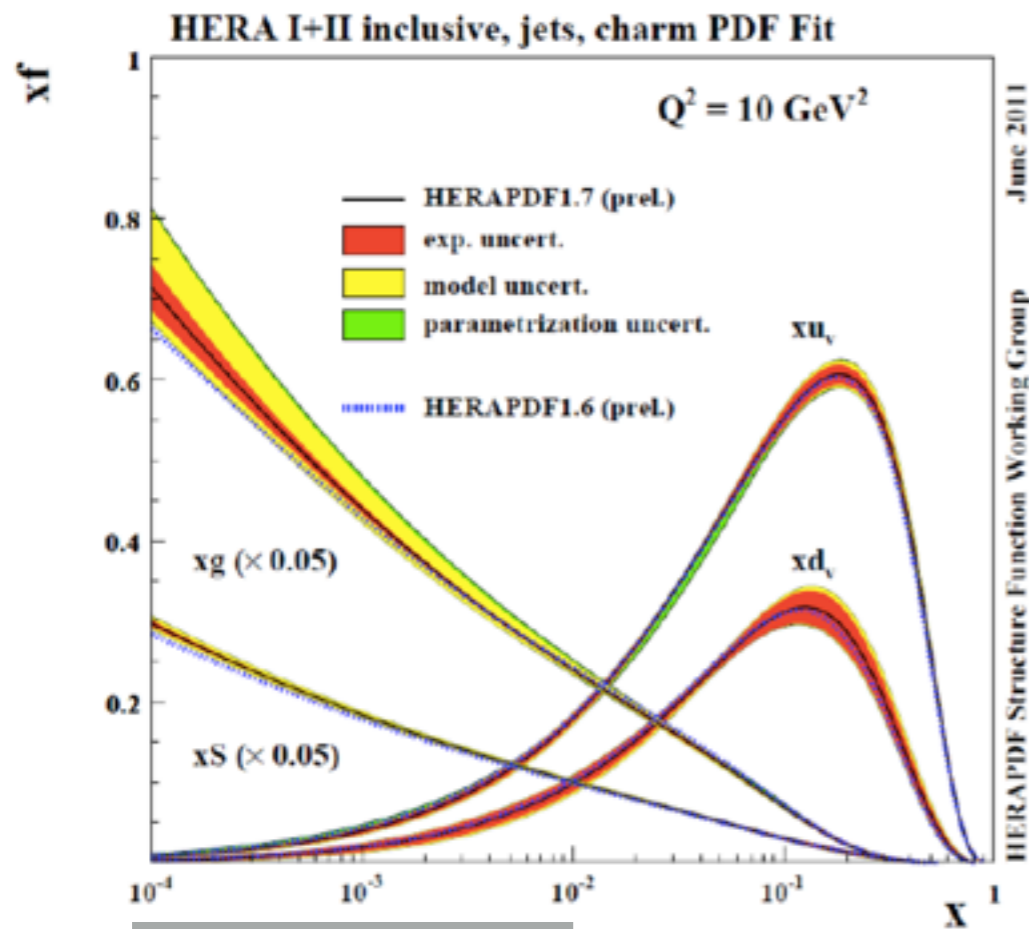
③ Nuclear absorption



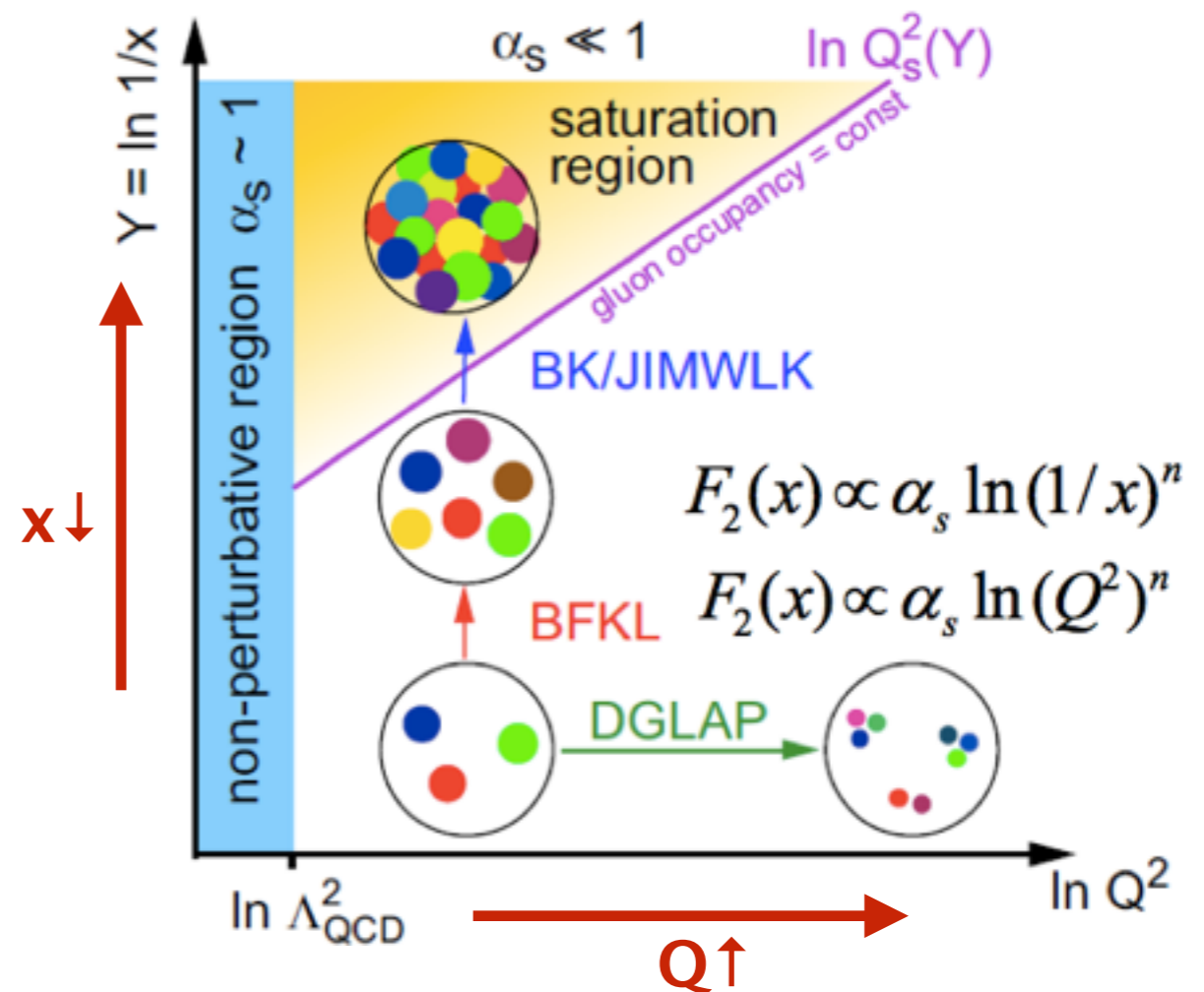
- Right after $c\bar{c}$ pairs are created, they interact with surrounding nucleons and break up.
- negligible at LHC energy (crossing time shorter than quarkonia formation time)

Parton Distribution Function in nucleon

- Q^2 : momentum transfer
- x : fraction of nucleon momentum carried by parton



arXiv.1111.5432v4

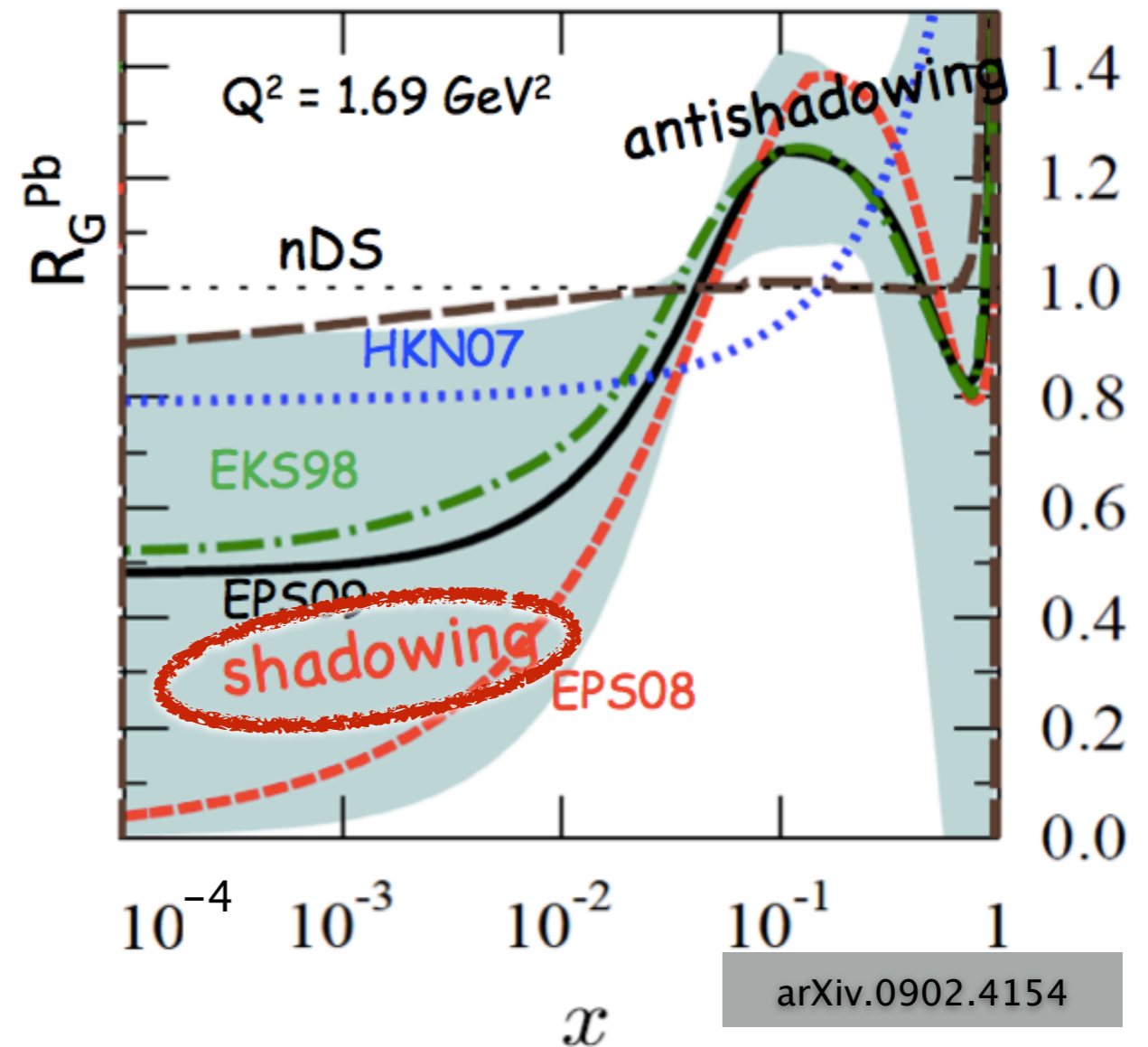


- Distributions of partons are related to the momentum fraction, **Bjorken x**.
- At low x , gluons are saturated and occupancy becomes constant.

⊕ nuclear Parton Distribution Function

$$R_G^{Pb}(x, Q^2) = \frac{xG_A(x, Q^2)}{AxG_p(x, Q^2)}$$

- A : mass number of nucleus
- x : Bjorken x
- G_p : gluon structure function in proton
- G_A : gluon structure function in nucleus



- $R_i^{pPb} \neq 1$
: Parton distributions in nucleus are different from those in free nucleons.
- nPDF is different for various model predictions (especially at low x).

⊕ J/ψ in pPb collisions

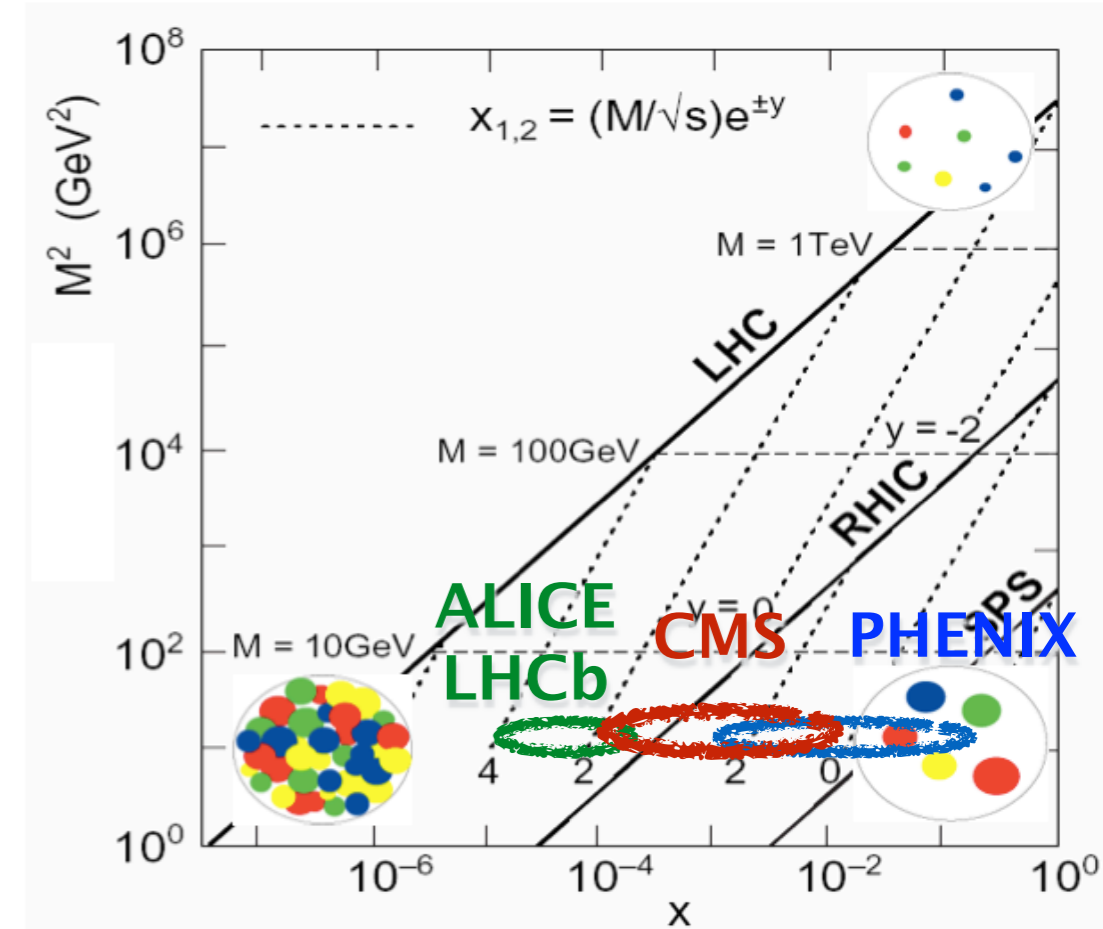
- J/ψ production measurements help in constraining nPDF of gluons, which dominate the production process.
- In case of 2 → 1 process, x_2 (of Pb nucleus) can be approximately :

$$x_2 = \frac{Q}{\sqrt{S_{NN}}} e^{-y}$$

⊕ J/ψ with CMS at LHC

- $M = 3.09 \text{ GeV}/c^2$
- $\sqrt{S_{NN}} = 5.02 \text{ TeV}$
- $-2.87 < y_{CM} < 1.93$
- $0 < p_T < 30 \text{ GeV}/c$

➔ $10^{-4} < x_2 < 10^{-2}$



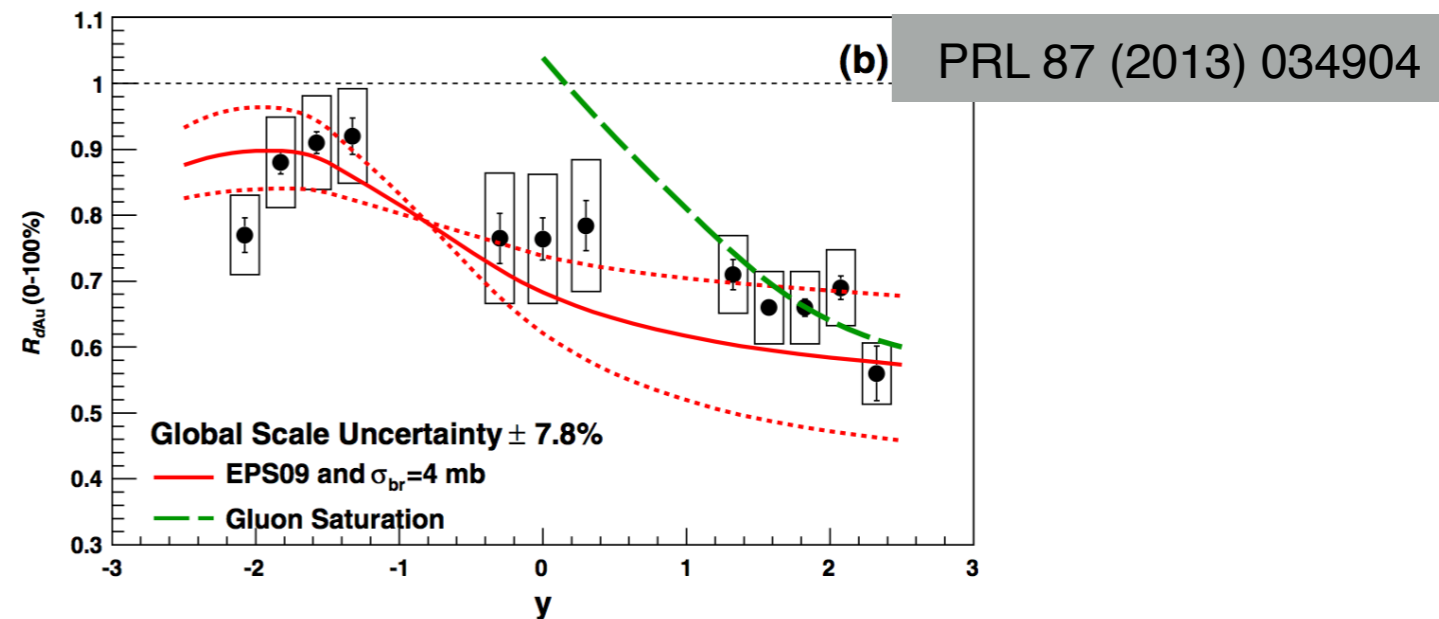
- We expect to investigate **smaller x** regions than RHIC
- Different kinematics range (**mid-rapidity, higher p_T**) with ALICE, LHCb
→ Better understanding & more constraints on nPDF models

⊕ Nuclear modification factor

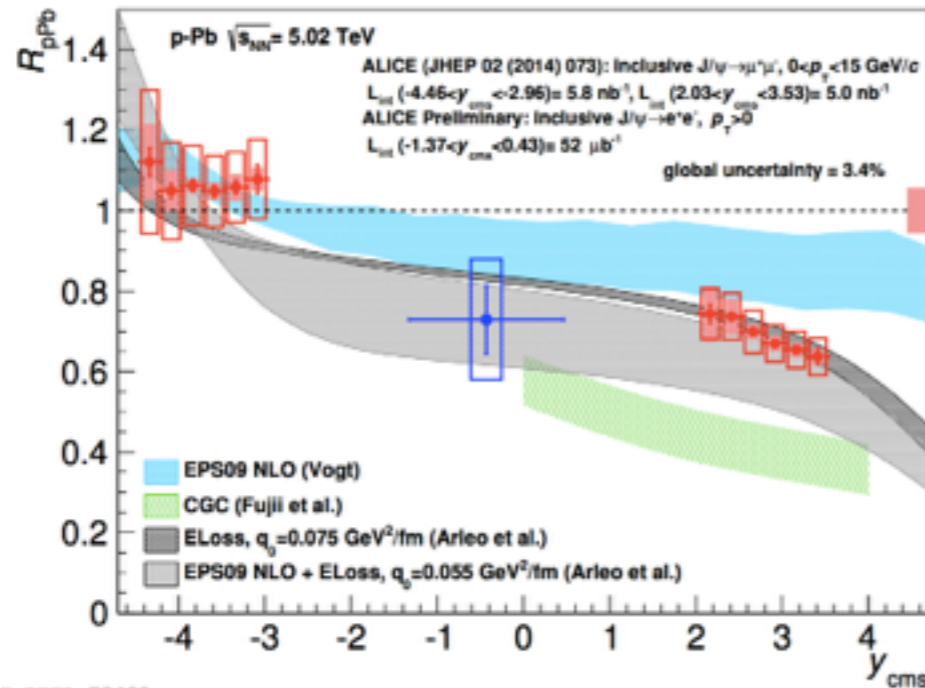
$$R_{pA} = \frac{1}{A} \cdot \frac{\sigma_{pA}}{\sigma_{pp}}$$

JHEP 02 (2013) 072

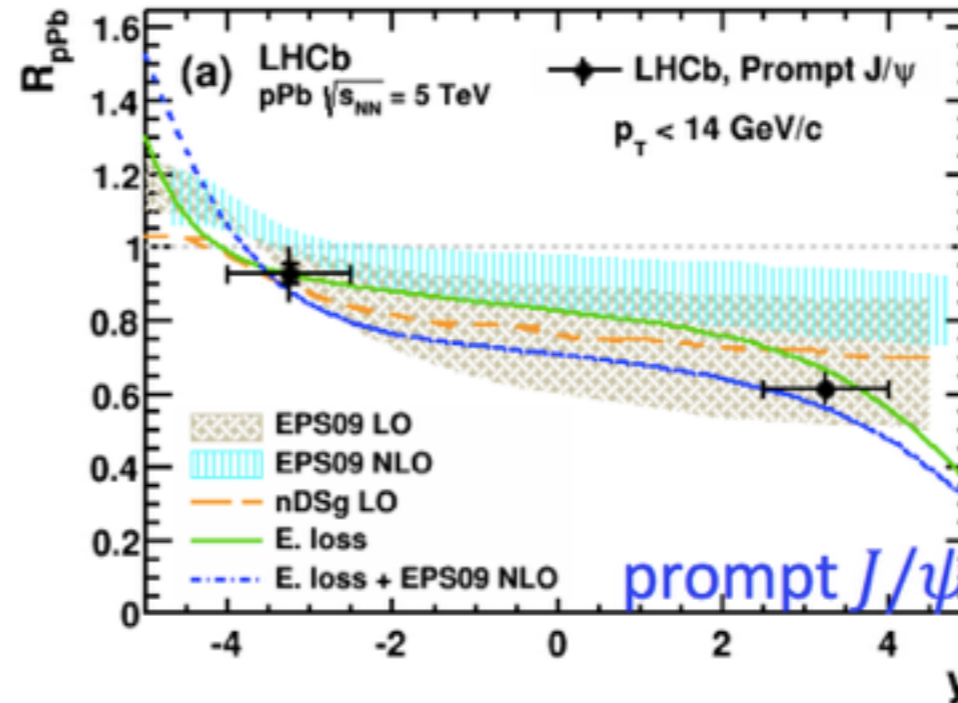
JHEP 02 (2014) 073



⊕ [Shadowing + Breakup] reproduces PHENIX data.

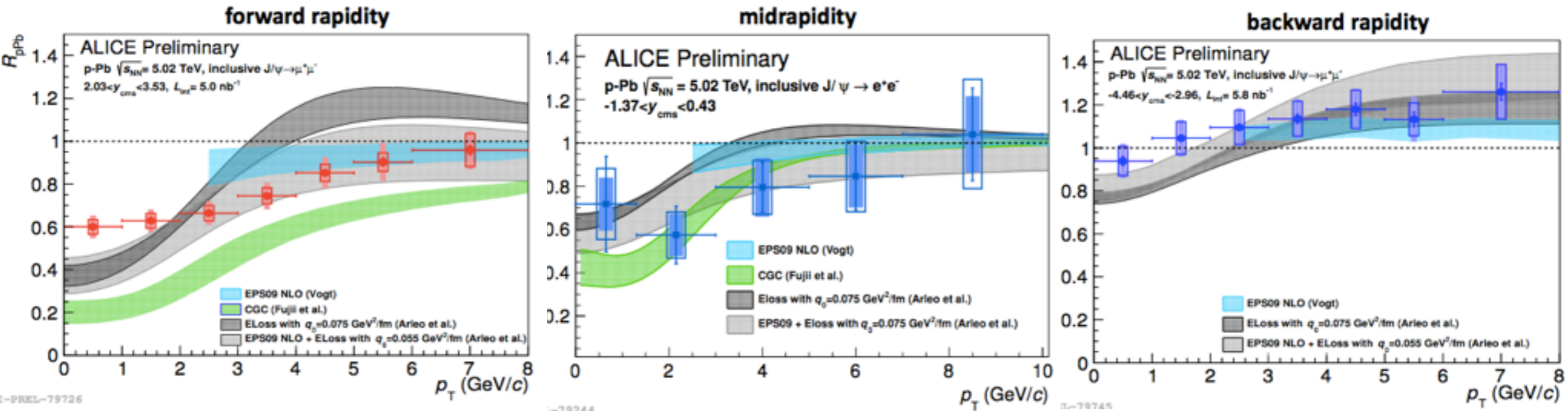


1-PREL-73492



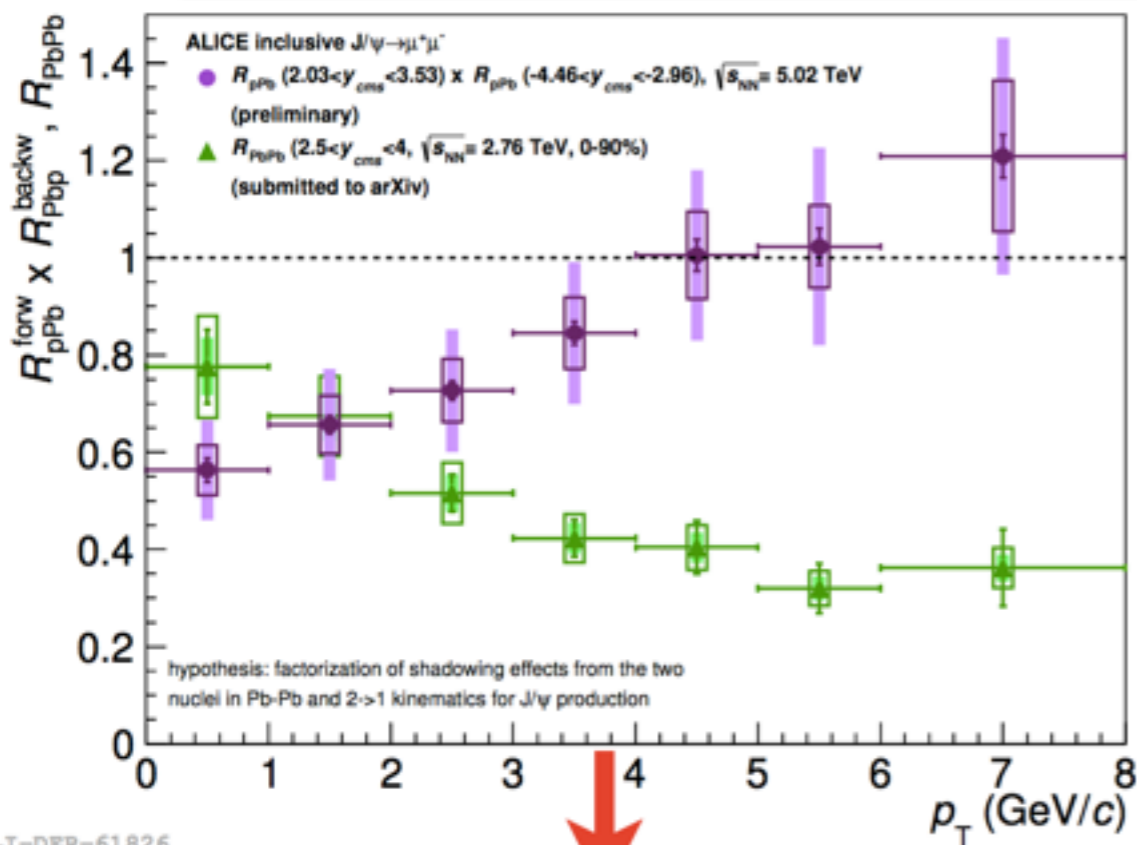
⊕ [Shadowing + Energy loss] describes ALICE and LHCb well.

⊕ [CGC inspired model] in poorer agreement with data



- ⊗ Forward : R_{pPb} increases with p_T , goes to unity for $p_T > 5 \text{ GeV}$
- ⊗ Mid-rapidity & backward : small p_T dependence, compatible with unity

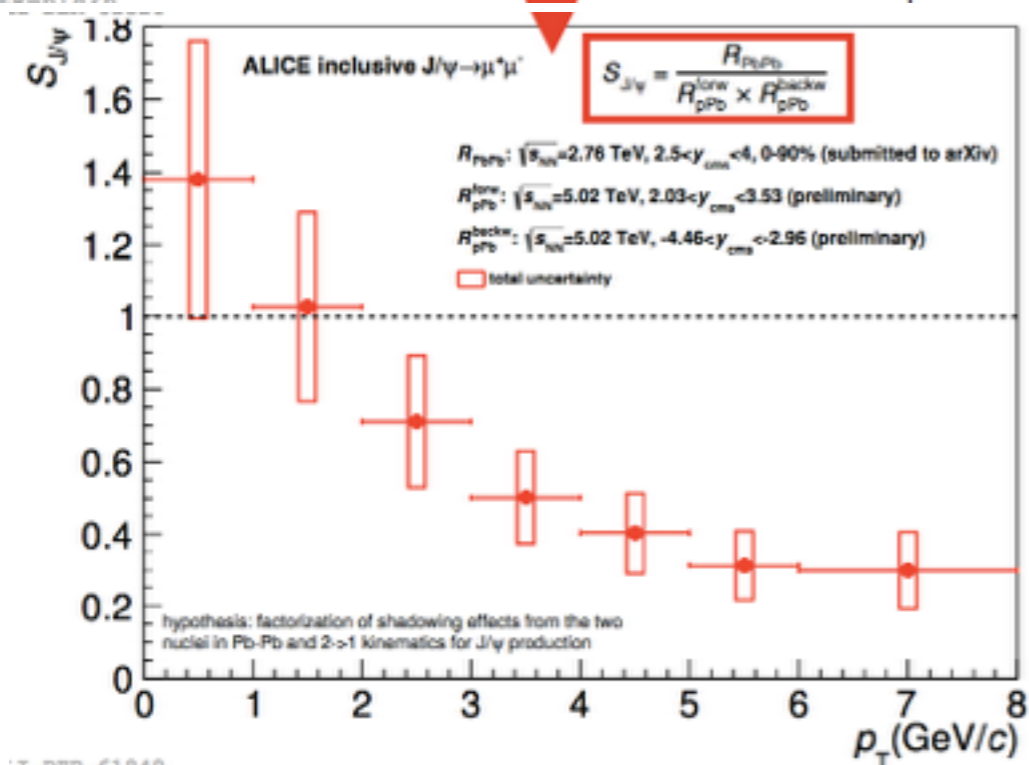
- EPS09 calculation is consistent with data ($p_T > 2.5 \text{ GeV}$).
- Energy loss shows disagreement at forward rapidity and low p_T .
- CGC overestimates suppression at forward rapidity.



⊙ $R_{pPb}^{forward} \times R_{pPb}^{backward} = R_{PbPb} ?$

■ Assumption

- Production mechanism $g+g \rightarrow J/\psi$
- CNM effects factorize in p-nucleus and are dominated by shadowing



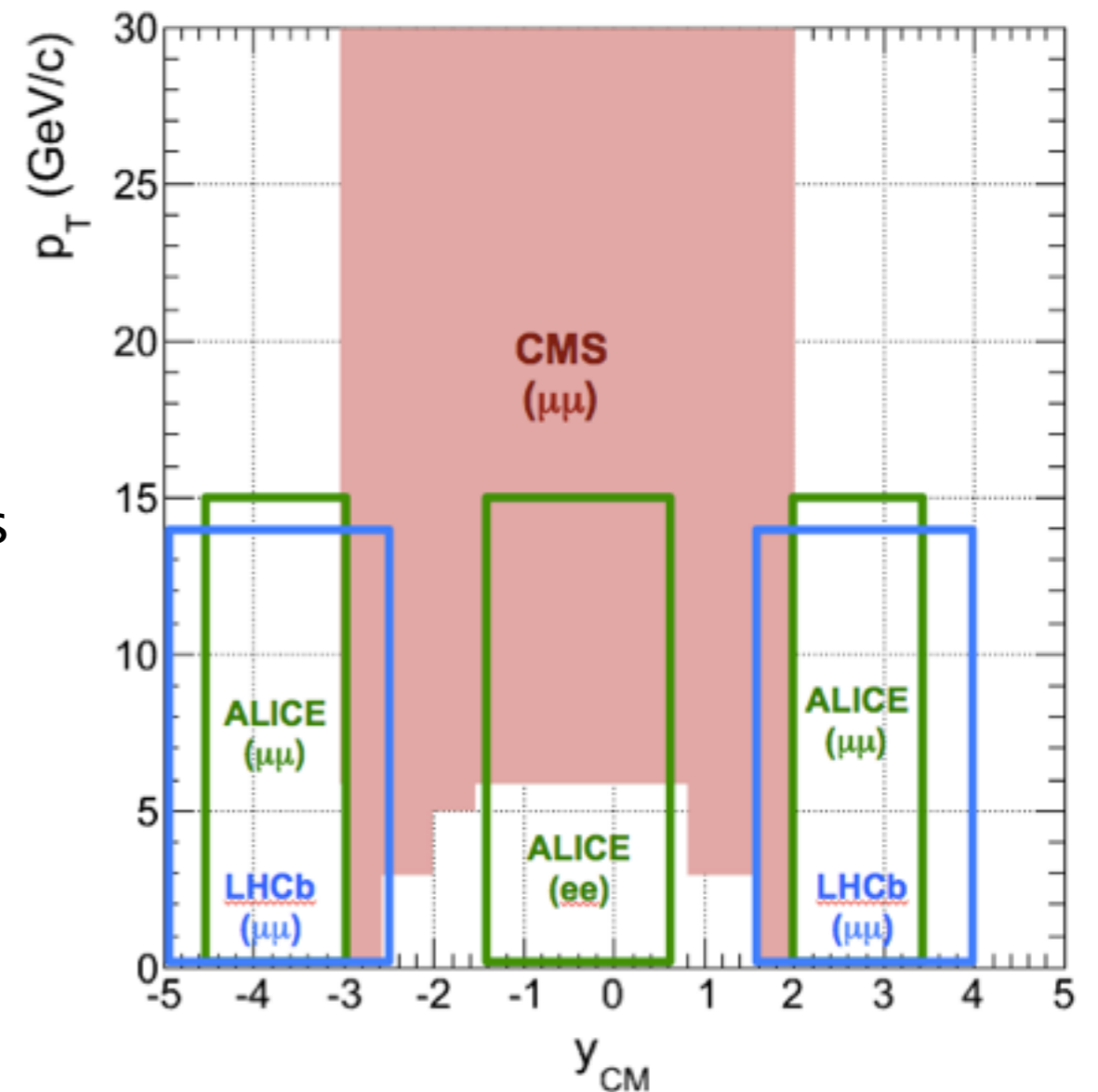
- ⊙ At low p_T , enhancement (recombination scenario?)
- ⊙ At high p_T , strong suppression

⊕ Observables

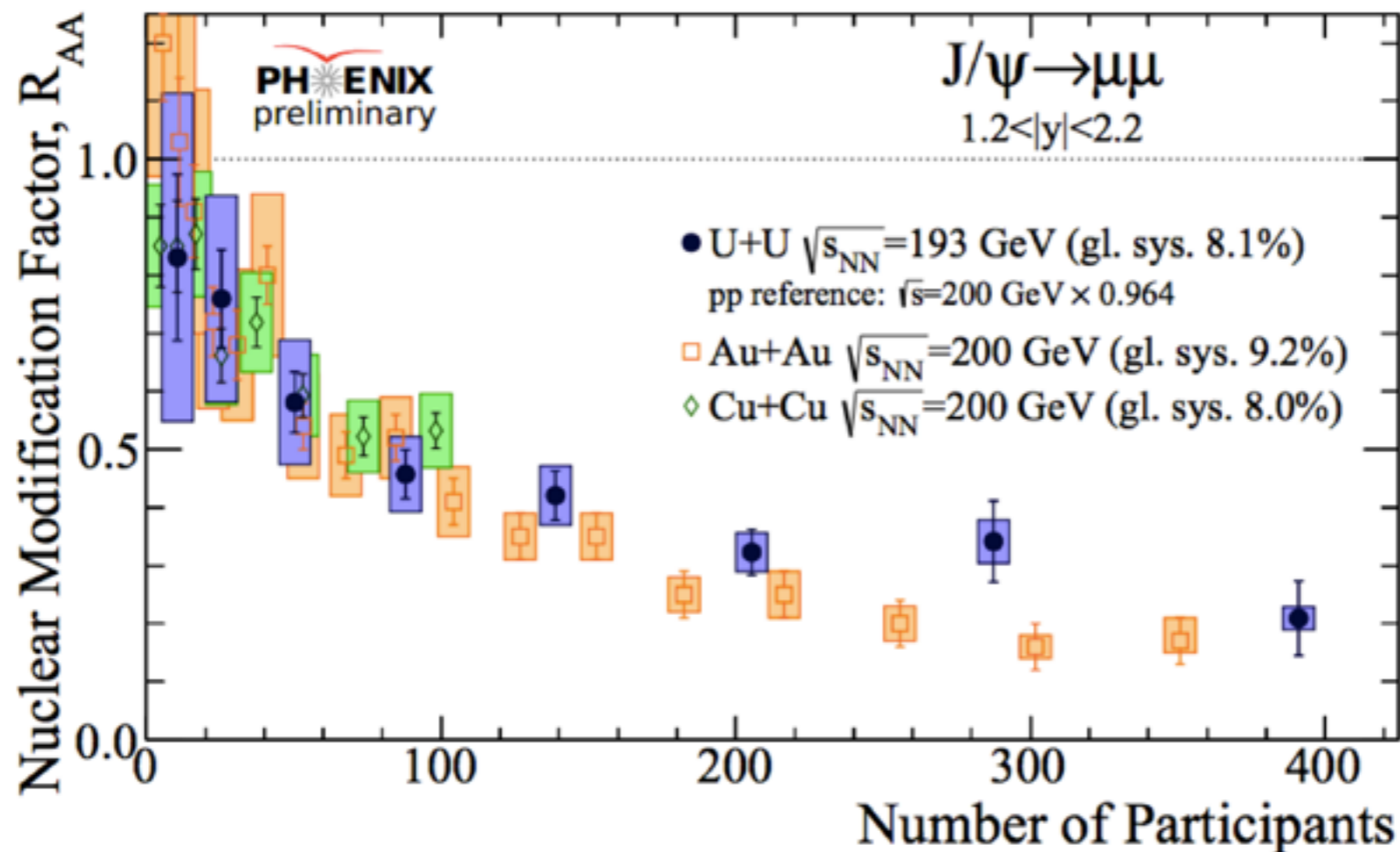
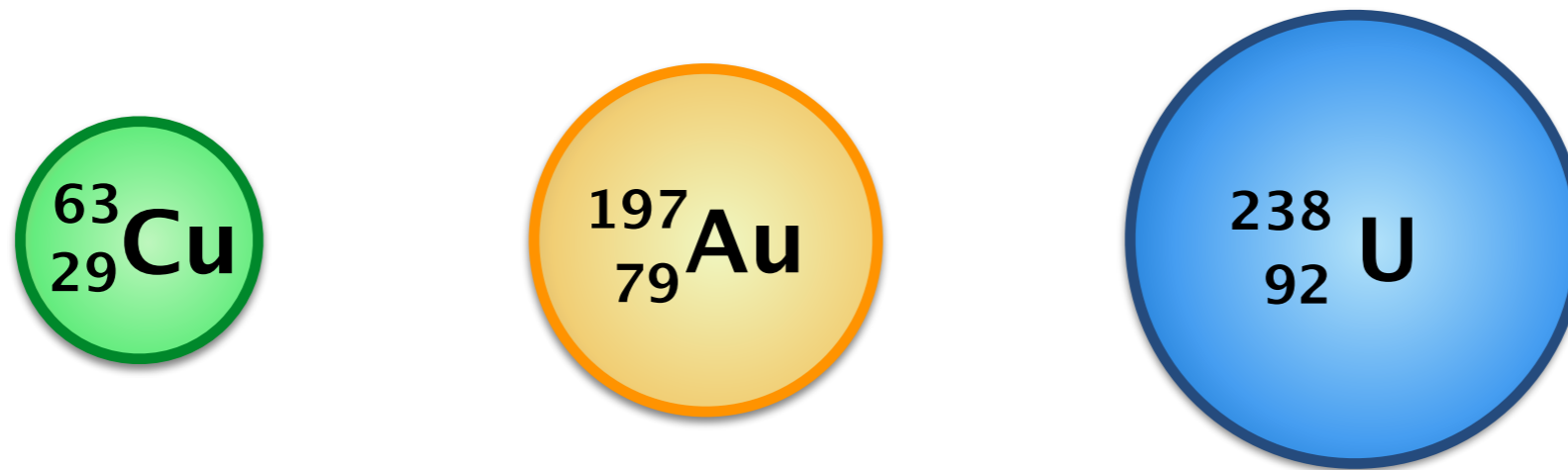
- Production cross-sections
- Forward to backward ratios R_{FB}
- Event-activity dependence

⊕ Down to $p_T \sim 0$ GeV/c

- Less background compared to PbPb collisions
 - Tracker-tracker muon pairs can be used
 - Basic technique is similar with PbPb analysis
- **J/ ψ Acceptance**
 - at mid-rapidity : $p_T > 6.5$ GeV/c
 - at forward rapidity : $p_T > 0$ GeV/c



- Both initial and final state effects depend on system size.



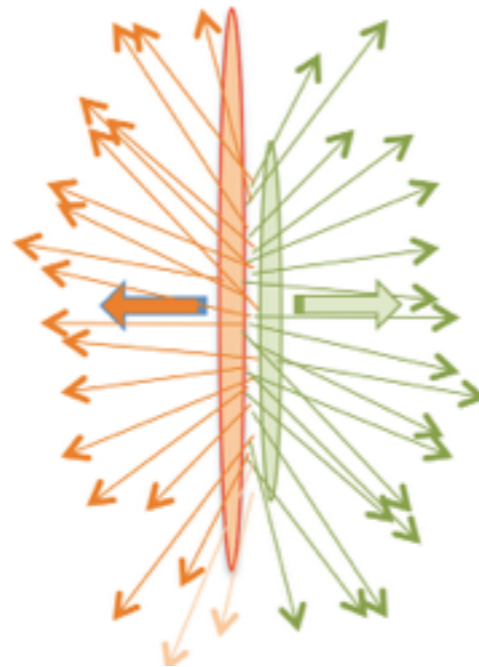
- R_{AA} consistent with each other

Backward (Au-going)

Final state effects

- Higher energy density
 - more suppression
 - (or more enhancement?)

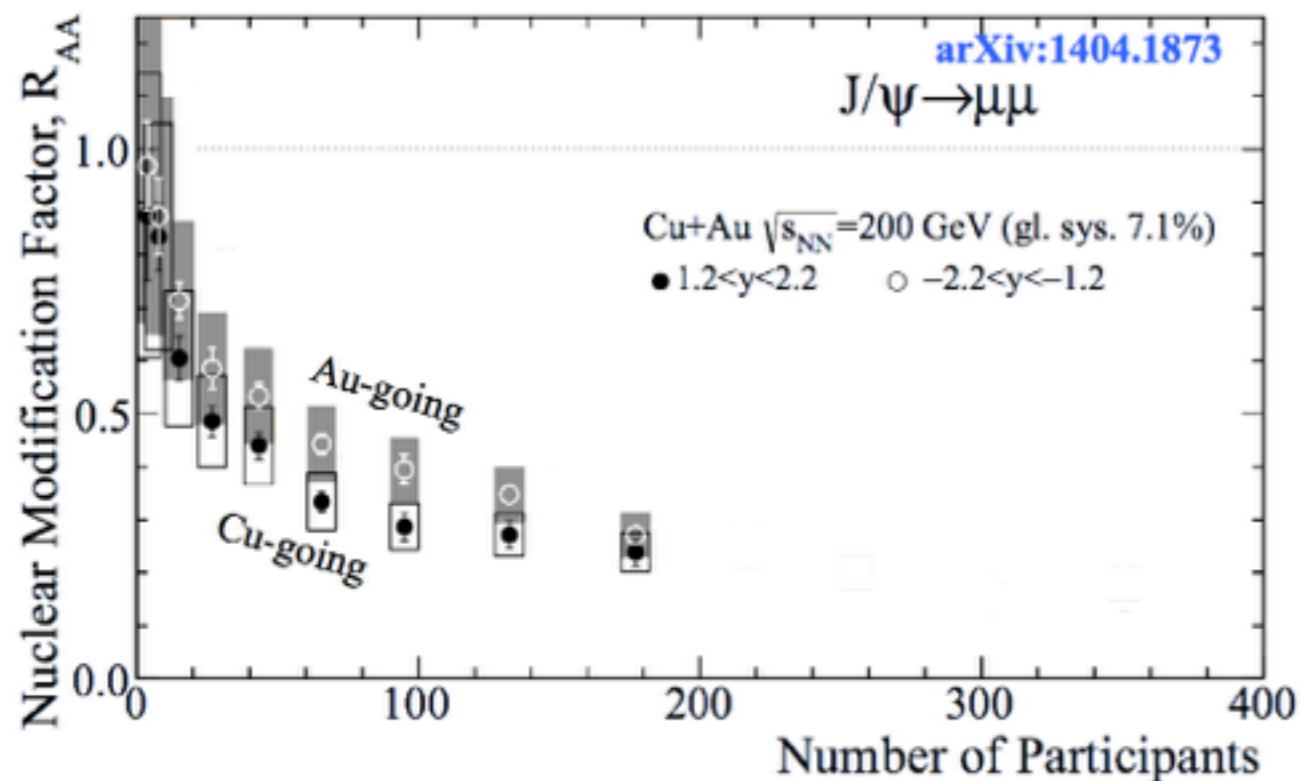
Cu+Au



Forward (Cu-going)

Initial state effects

- probes small x in Au
 - PDF more strongly modified in heavier Au compared to Cu.
- Longer crossing time for Cu than Au
 - energy loss depends on Au
 - break-up depends on Cu



arXiv.1404.1873

R_{AA} is also asymmetric?

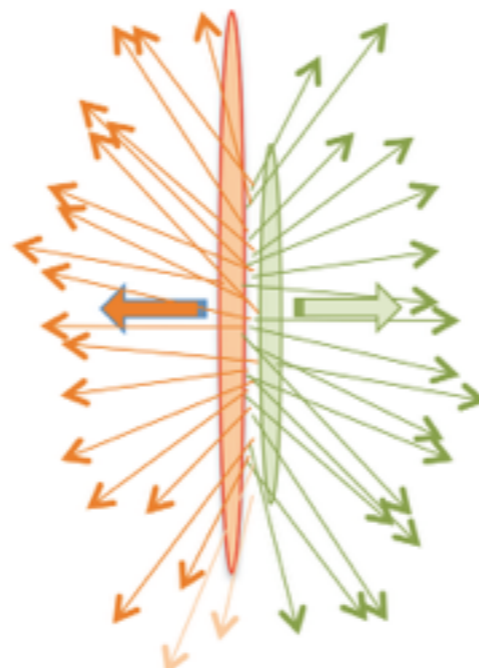
- Their ratio R (forward/backward) in the next slide

Backward (Au-going)

Final state effects

- Higher energy density
 - more suppression
 - (or more enhancement?)

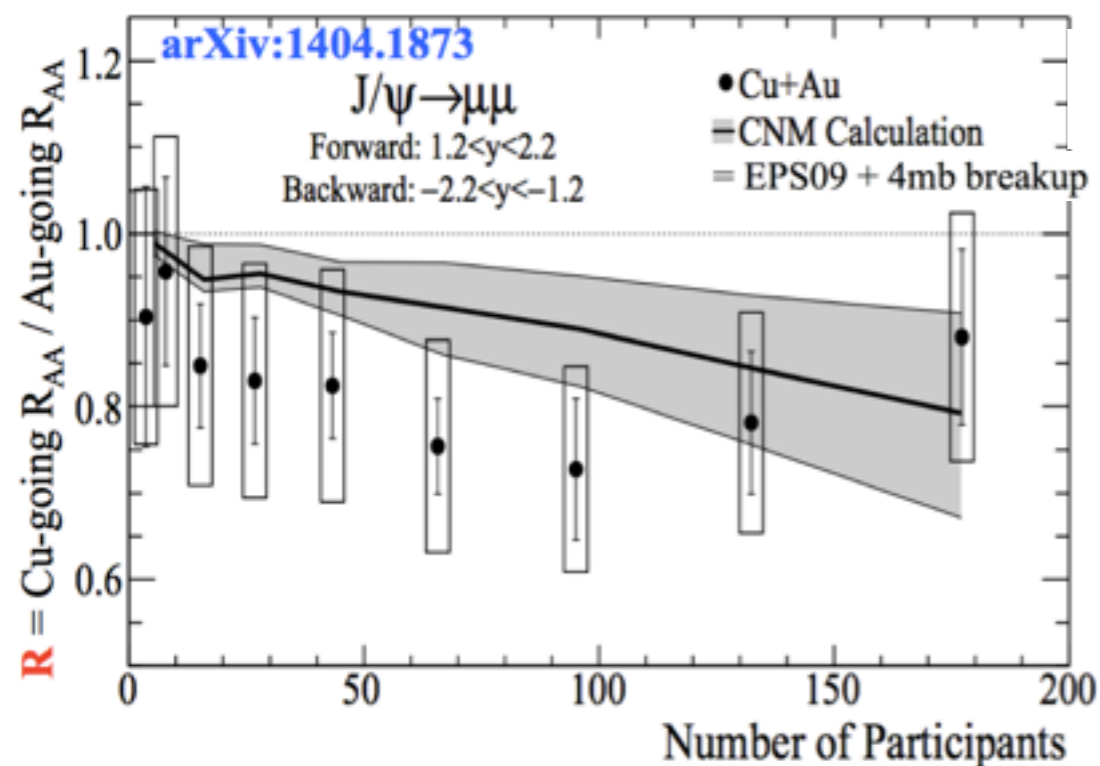
Cu+Au



Forward (Cu-going)

Initial state effects

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 - PDF more strongly modified in heavier Au compared to Cu.
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arXiv.1404.1873

- Initial state effects cause R decrease with centrality
- Final state effects would increase R with centrality.

⊕ This study may provide insight on the balance of hot and cold nuclear effects.

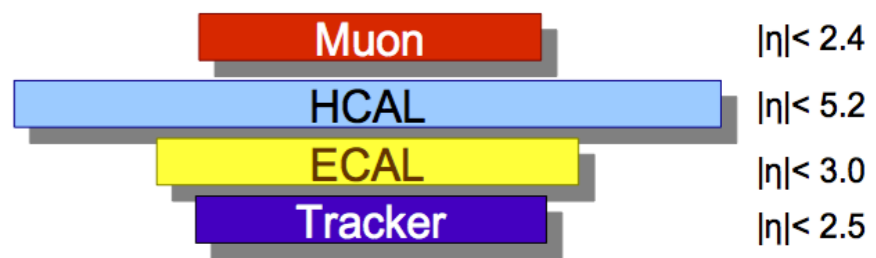
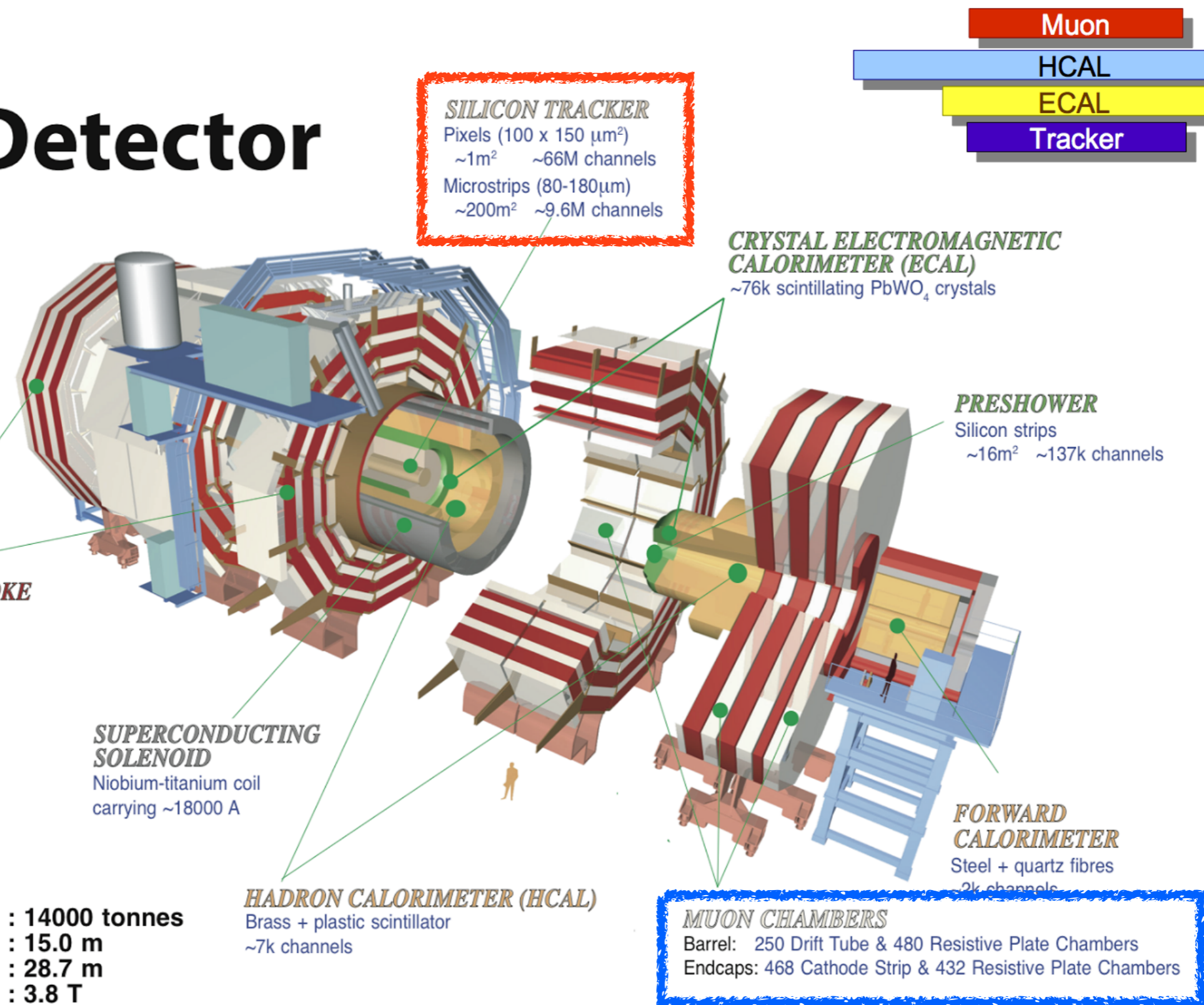
- ⊕ **J/ ψ measurements in pp, pPb and PbPb collisions play a key role to understand and factorize the initial and final state effects.**
- ⊕ **In PbPb, suppression has been observed and its dependence on centrality, p_T and rapidity in different kinematic ranges and beam energies need to be interpreted further.**
- ⊕ **In pPb, data can be described by several CNM models, especially nuclear shadowing, but no single models can describe all results.**
- ⊕ **System size can help to investigate the different sources of J/ ψ modification in heavy ion collisions.**



BACK-UP

CMS Detector

Pixels
 Tracker
 ECAL
 HCAL
 Solenoid
 Steel Yoke
 Muons



Total weight : 14000 tonnes
Overall diameter : 15.0 m
Overall length : 28.7 m
Magnetic field : 3.8 T

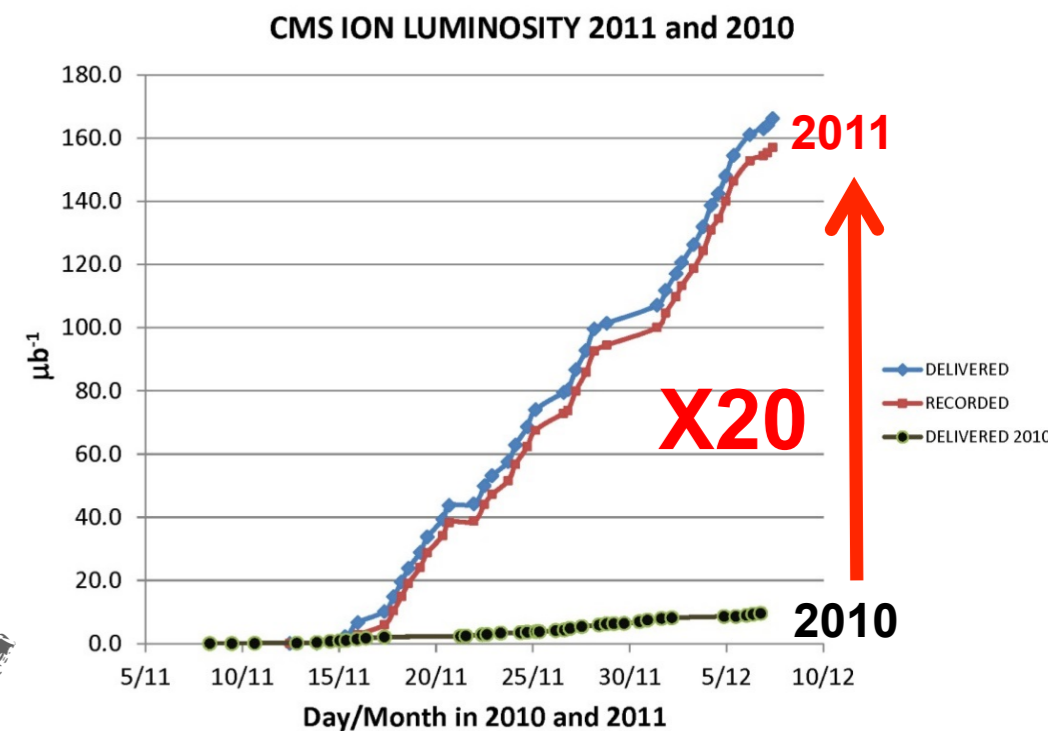
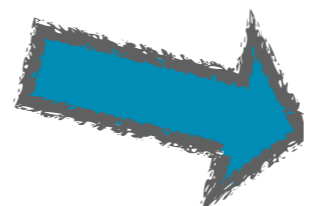
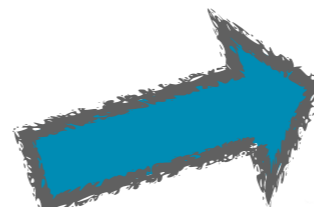
- ① **1st PbPb run @ $\sqrt{s_{NN}} = 2.76$ TeV**
 - Nov. – Dec. 2010
 - Recorded luminosity by CMS : $7.28 \mu\text{b}^{-1}$

- ① **1st pp run @ $\sqrt{s_{NN}} = 2.76$ TeV**
 - March 2011
 - Recorded luminosity by CMS : 225 nb^{-1}

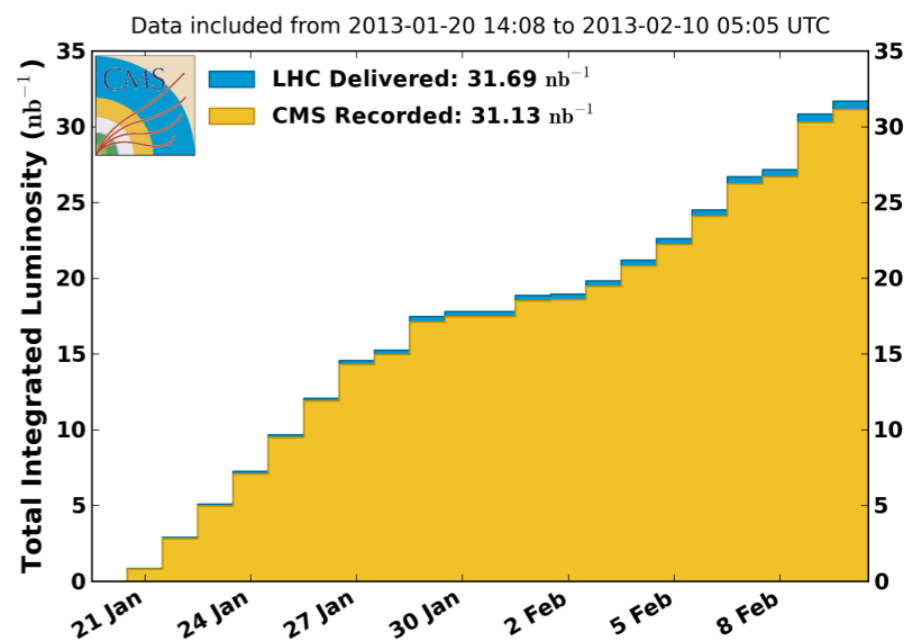
- ① **2nd PbPb run @ $\sqrt{s_{NN}} = 2.76$ TeV**
 - Nov. – Dec. 2011
 - Recorded luminosity by CMS : $150 \mu\text{b}^{-1}$

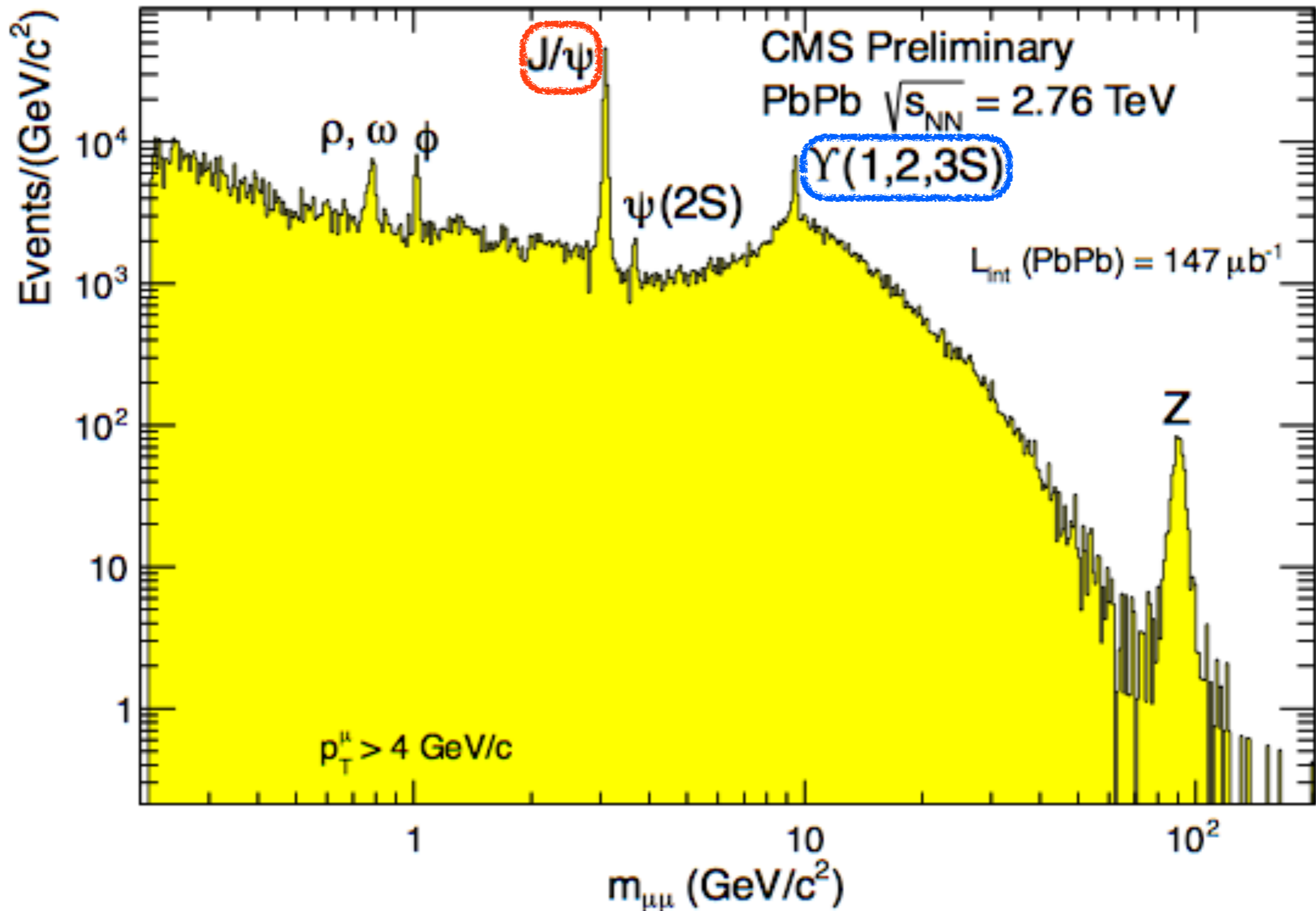
- ① **pPb run @ $\sqrt{s_{NN}} = 5.02$ TeV**
 - Jan. – Feb. 2013
 - Recorded luminosity by CMS : 35 nb^{-1}

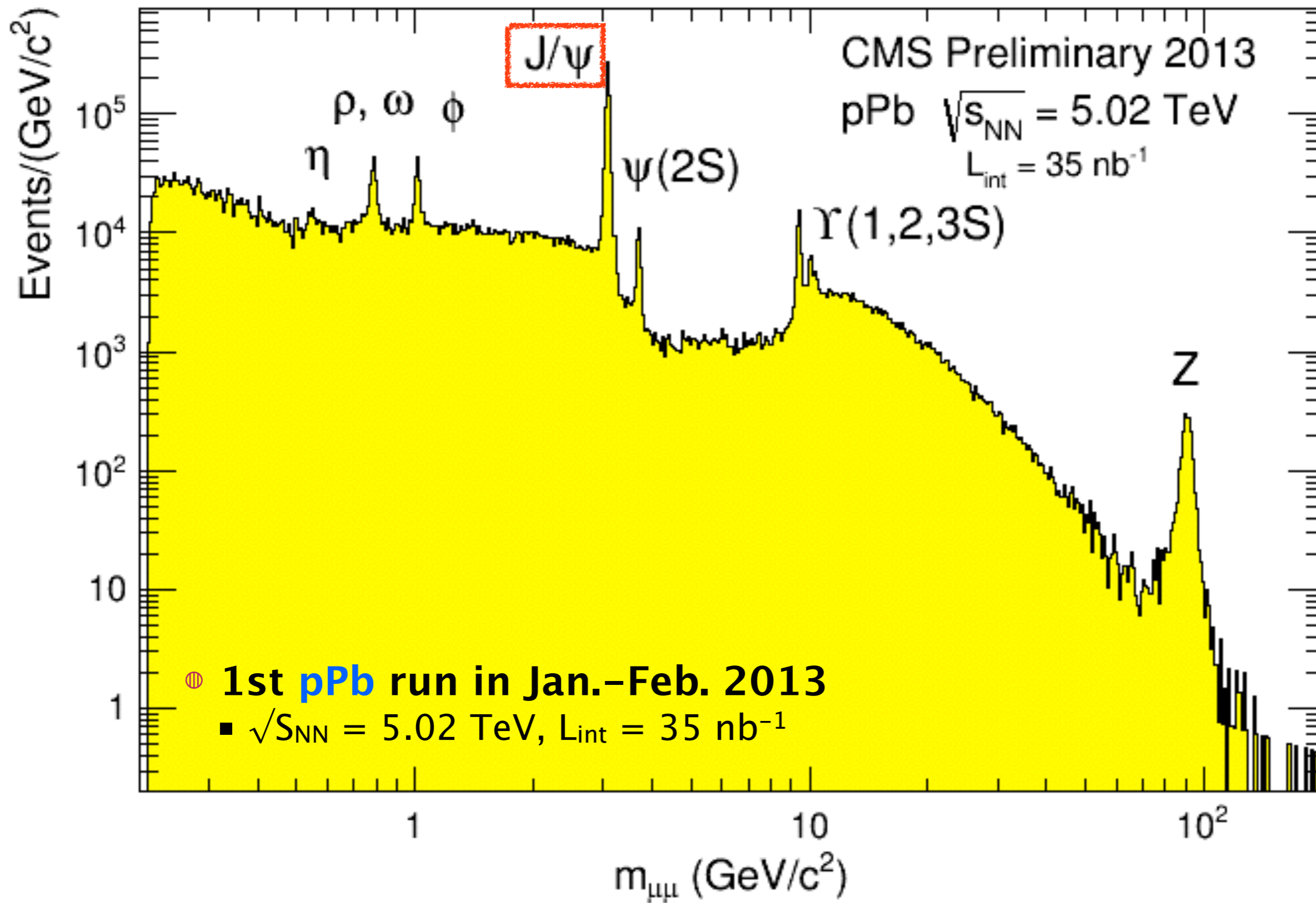
- ① **2nd pp run @ $\sqrt{s_{NN}} = 2.76$ TeV**
 - Feb. 2013 (3 days)
 - Recorded luminosity by CMS : 5.41 pb^{-1}



CMS Integrated Luminosity, pPb, 2013, $\sqrt{s} = 5.02$ TeV/nucleon





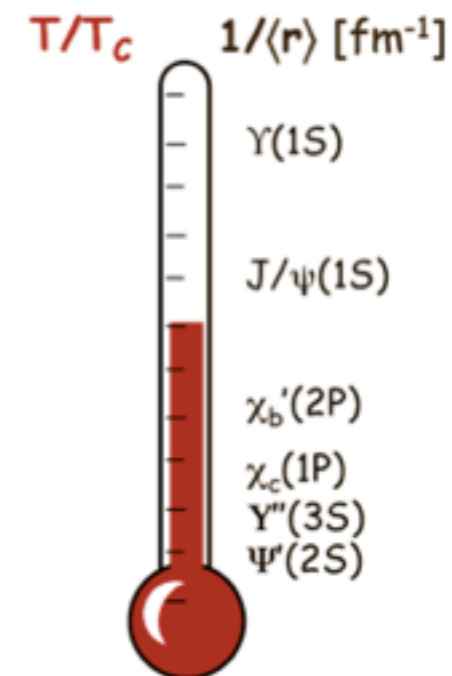
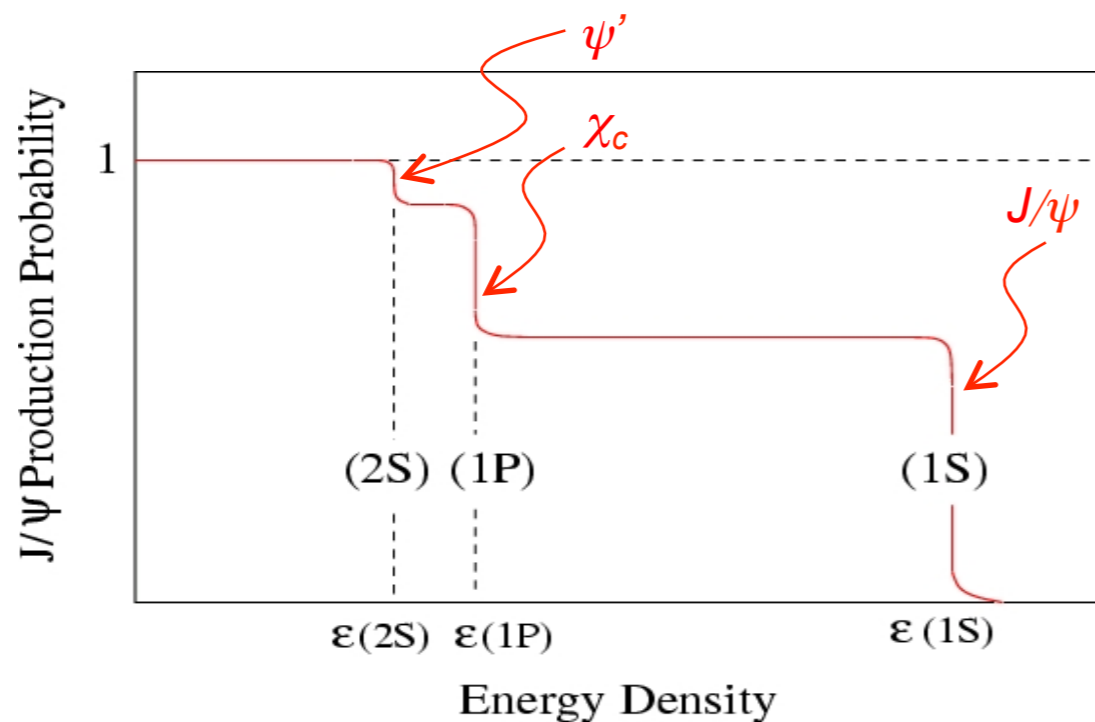


Debye screening

Debye screening

- The larger the binding energy, the higher the dissociation temperature T_d .
- As temperature goes up, Debye length $r_\lambda(T)$ decreases.

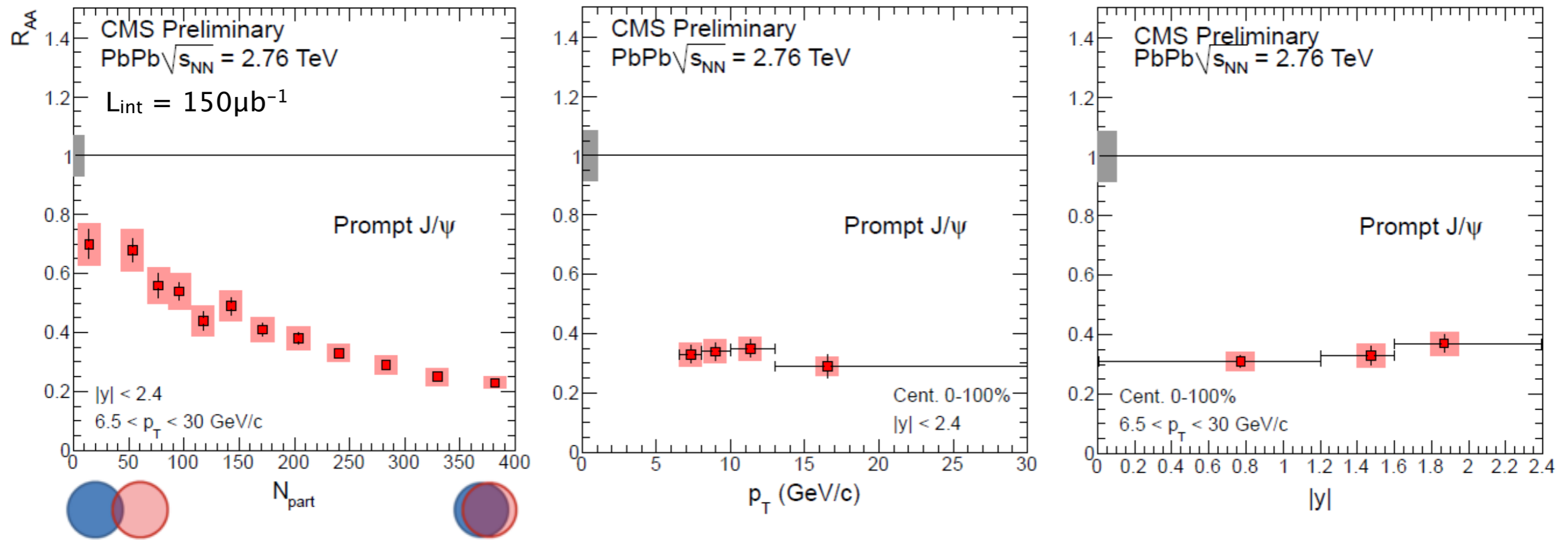
| Resonance | J/ψ | ψ' | Υ(1S) | Υ(2S) | Υ(3S) |
|-------------|------|------|-------|-------|-------|
| Mass [GeV] | 3.10 | 3.68 | 9.46 | 10.02 | 10.36 |
| ΔE [GeV] | 0.64 | 0.05 | 1.10 | 0.54 | 0.20 |
| Radius [fm] | 0.25 | 0.45 | 0.14 | 0.28 | 0.39 |



Mocsy, EPJC 61 (2009) 705

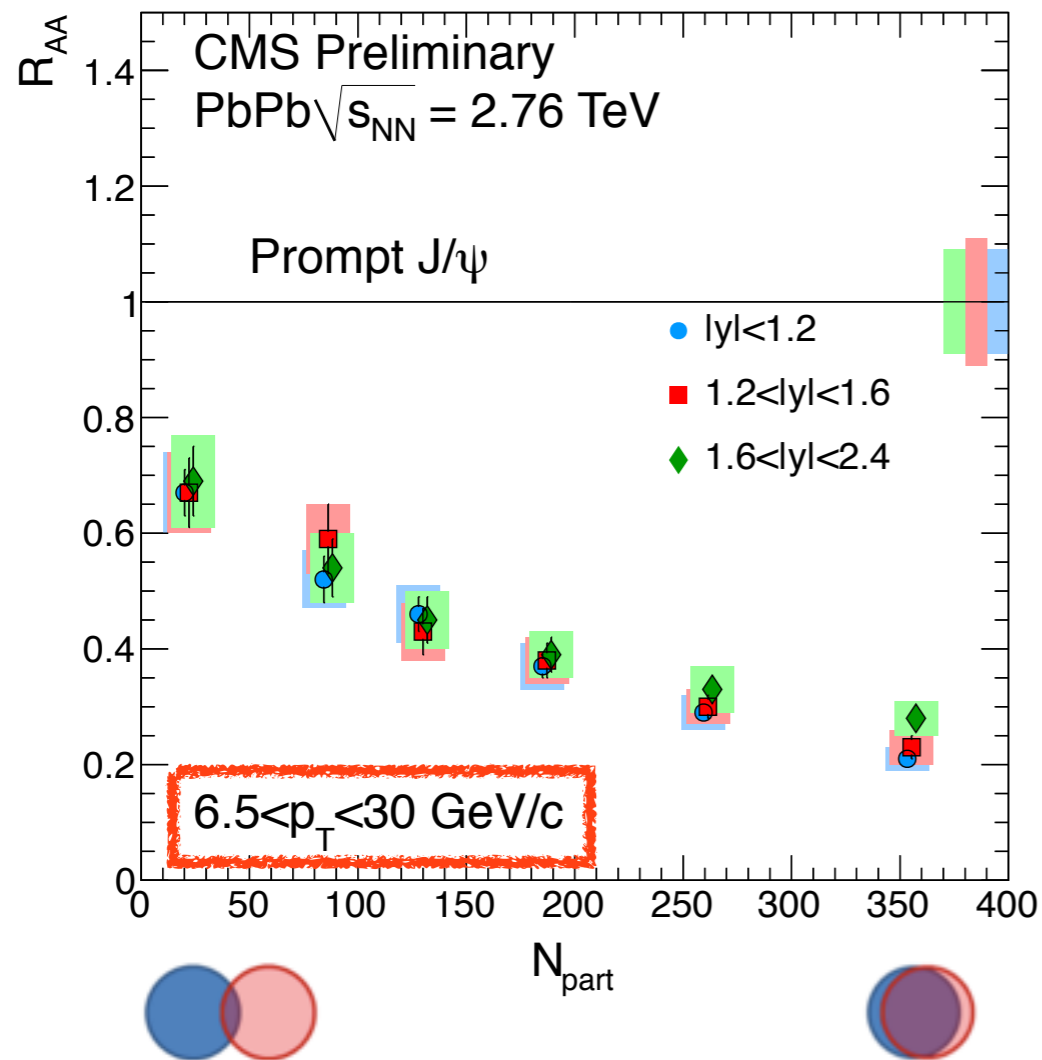
⊕ Nuclear modification factor

$$R_{AA} = \frac{\mathcal{L}_{pp}}{T_{AA} N_{MB}} \frac{N_{PbPb}}{N_{pp}} \cdot \frac{\epsilon_{pp}}{\epsilon_{PbPb}} \quad : R_{AA} = 1 \text{ No modification compared to pp collisions}$$

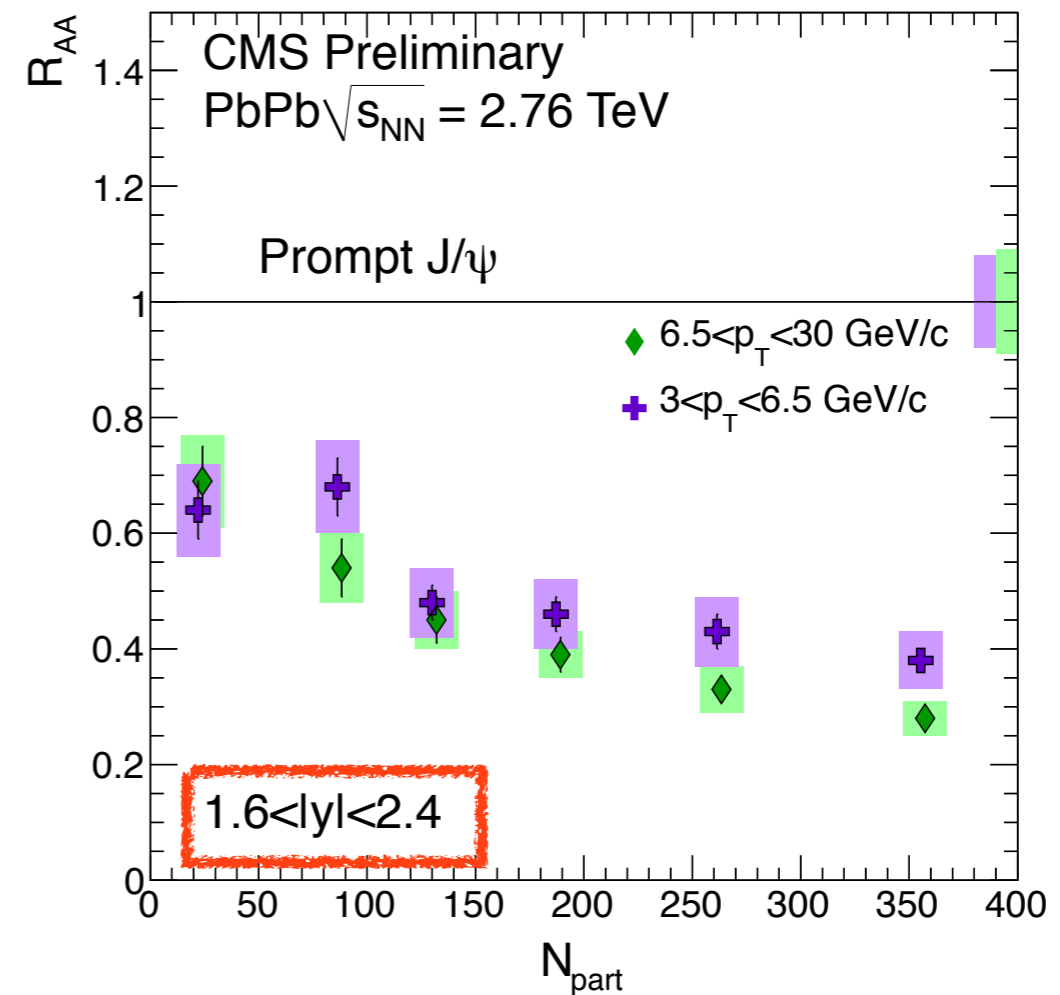


- ⊕ Suppressed by factor ~ 5 in the most central bin
- ⊕ No p_T and y dependent suppression is observed.

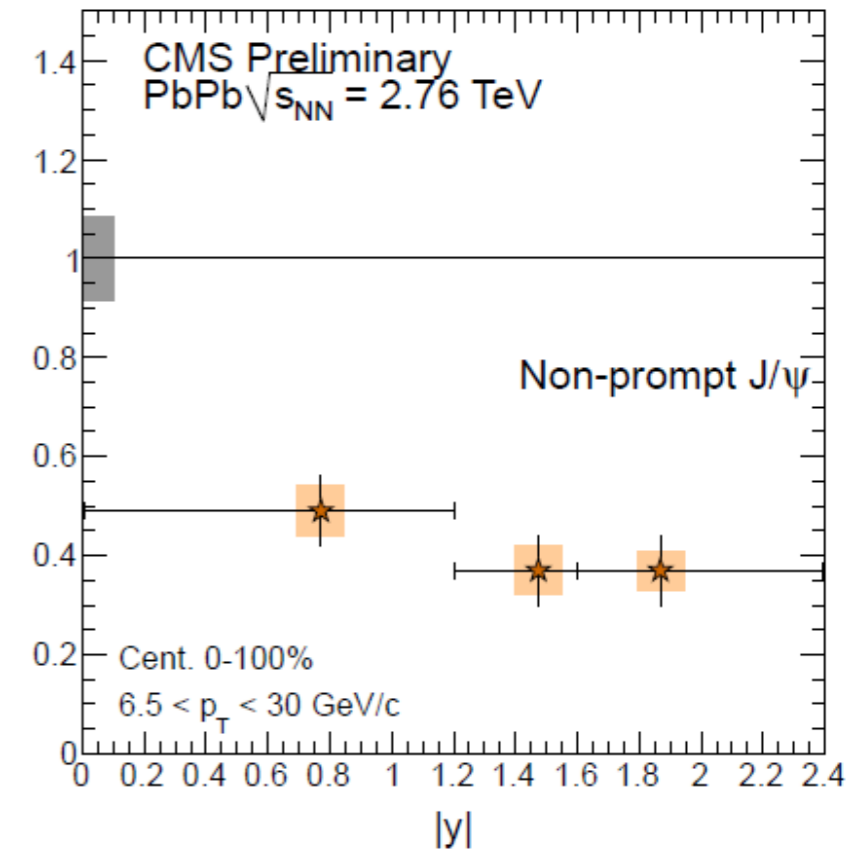
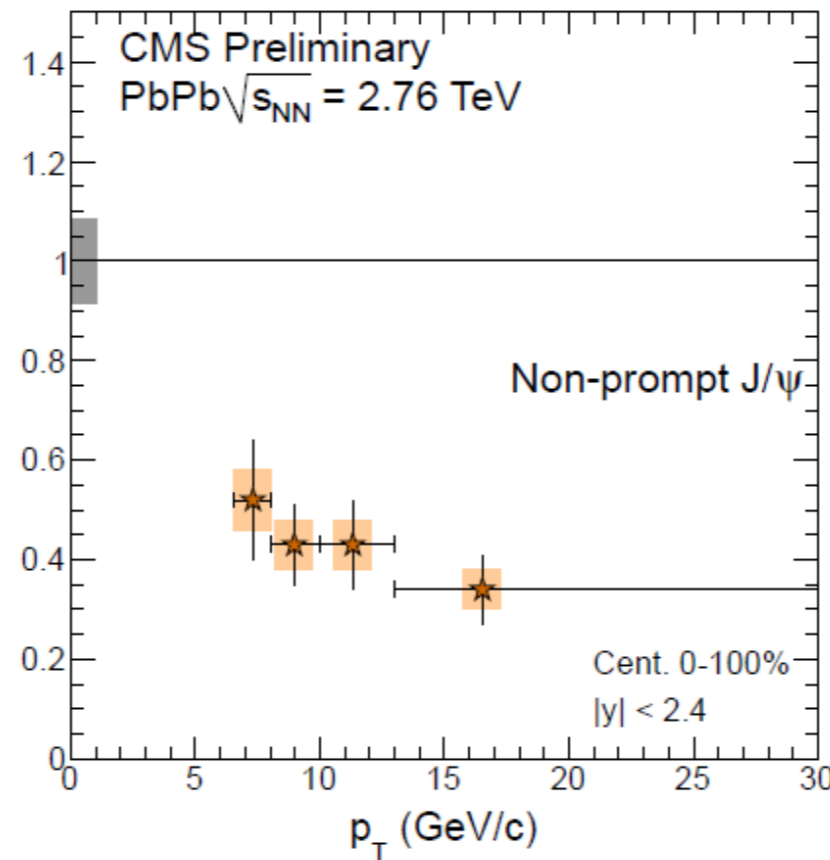
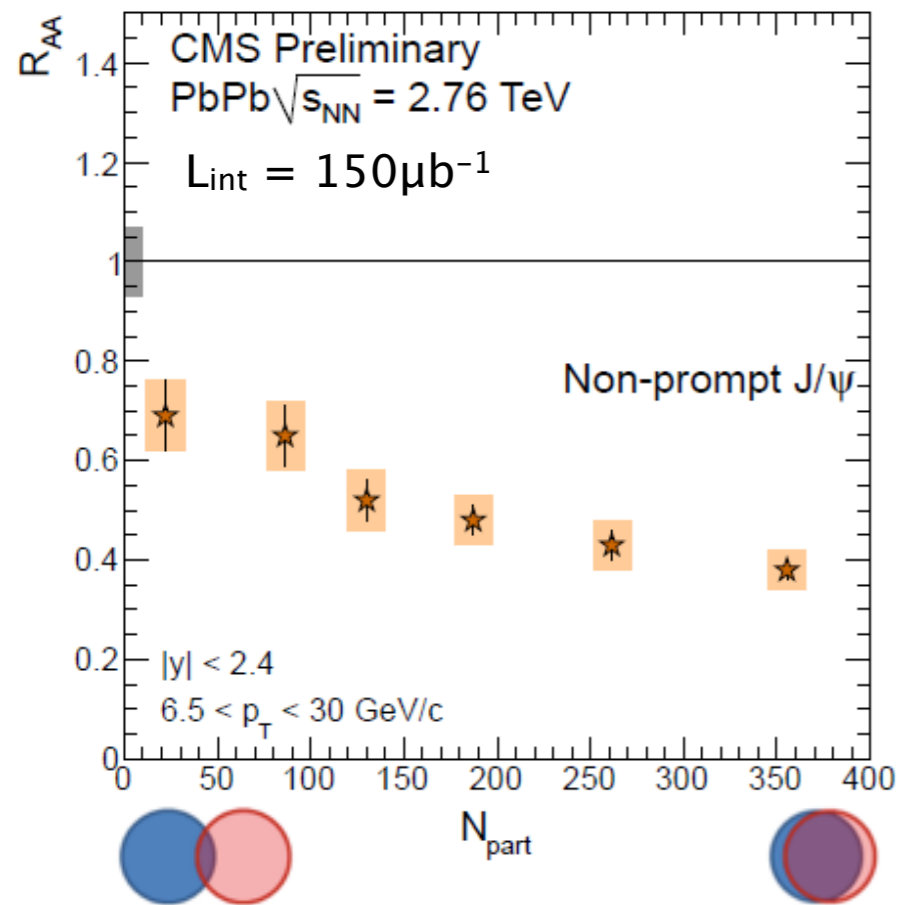
⊗ Rapidity dependence



⊗ p_T dependence



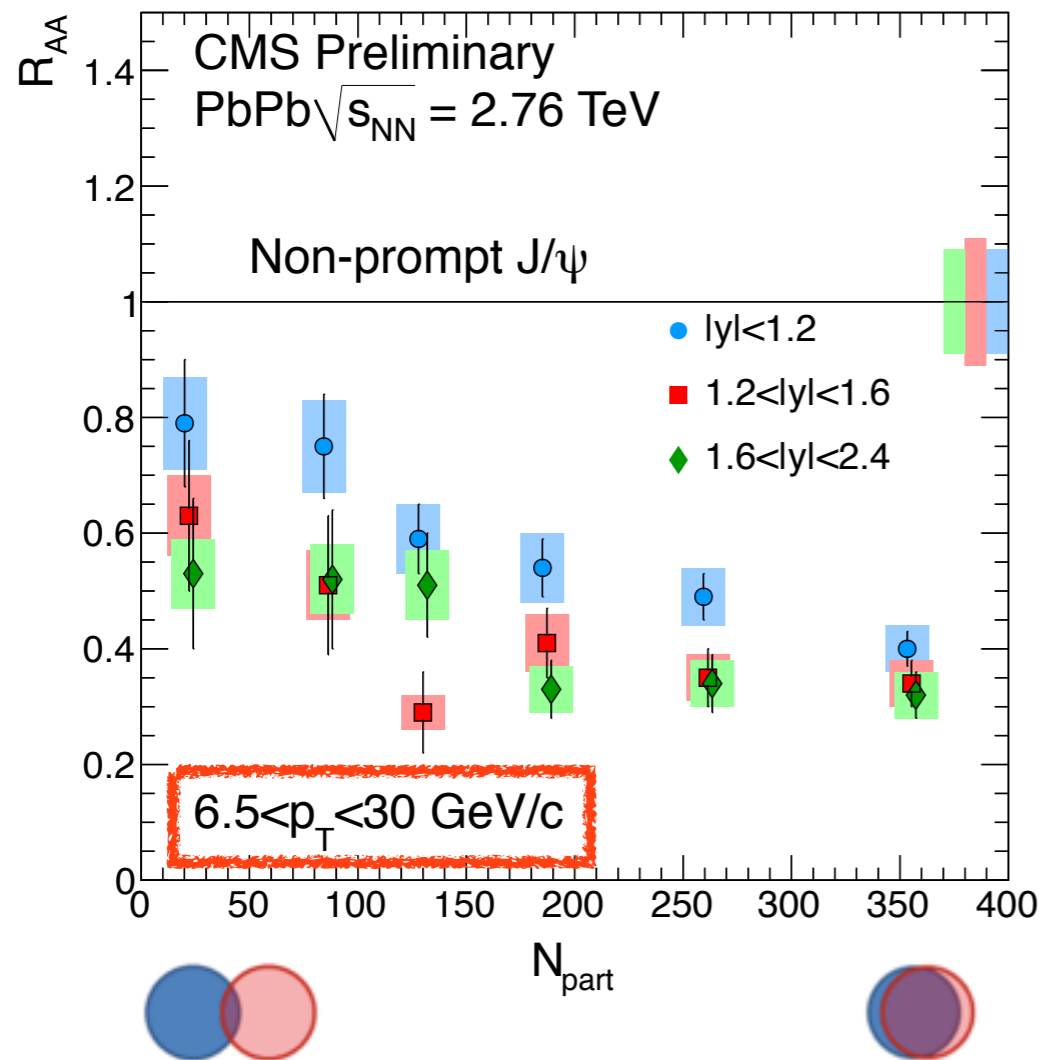
- ⊗ Left : No strong dependence on rapidity at high p_T region
- ⊗ Right : At forward rapidity region, lower p_T J/ψ is slightly less suppressed in the most central bins.



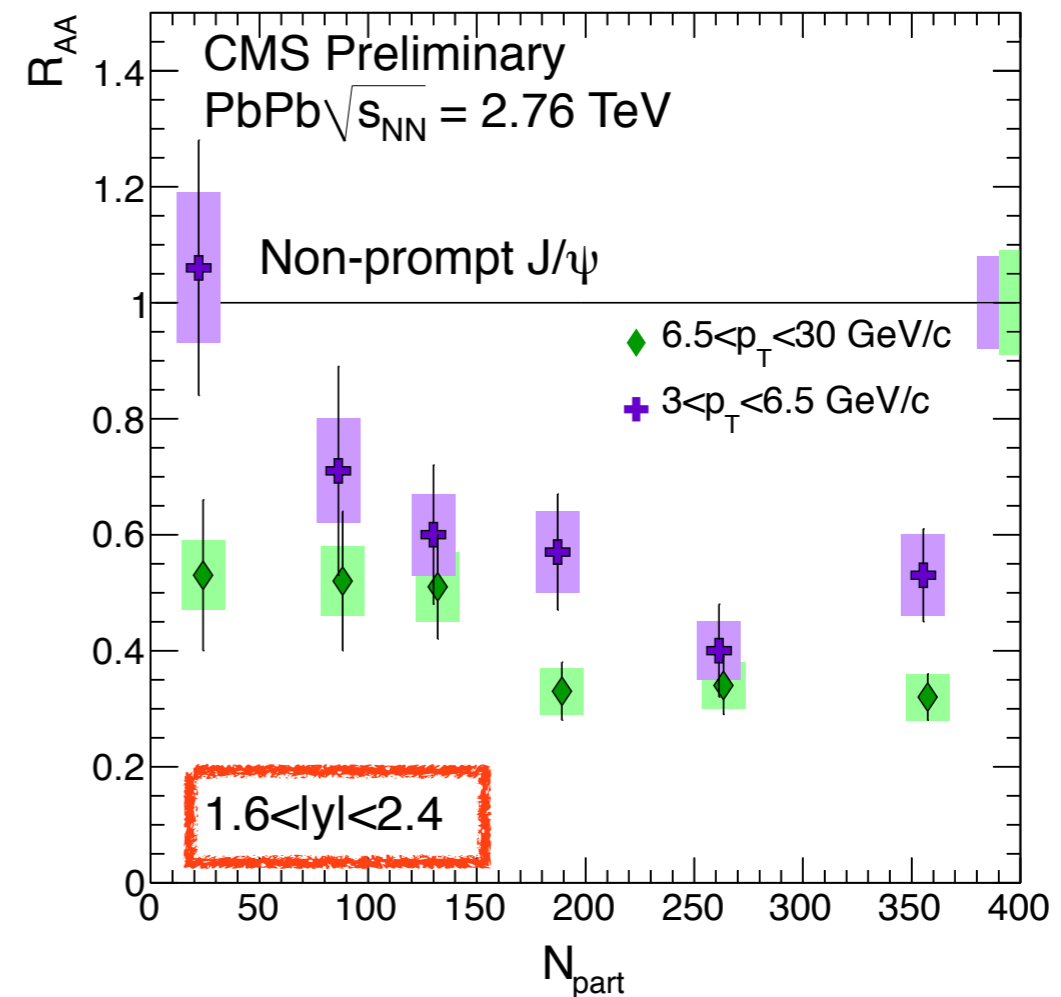
- ⊗ Suppressed by factor ~3 in the most central bin
- ⊗ Hints of smaller suppression at lower p_T region, mid-rapidity region

Information on the **b-quark energy loss** in medium

⊗ Rapidity dependence



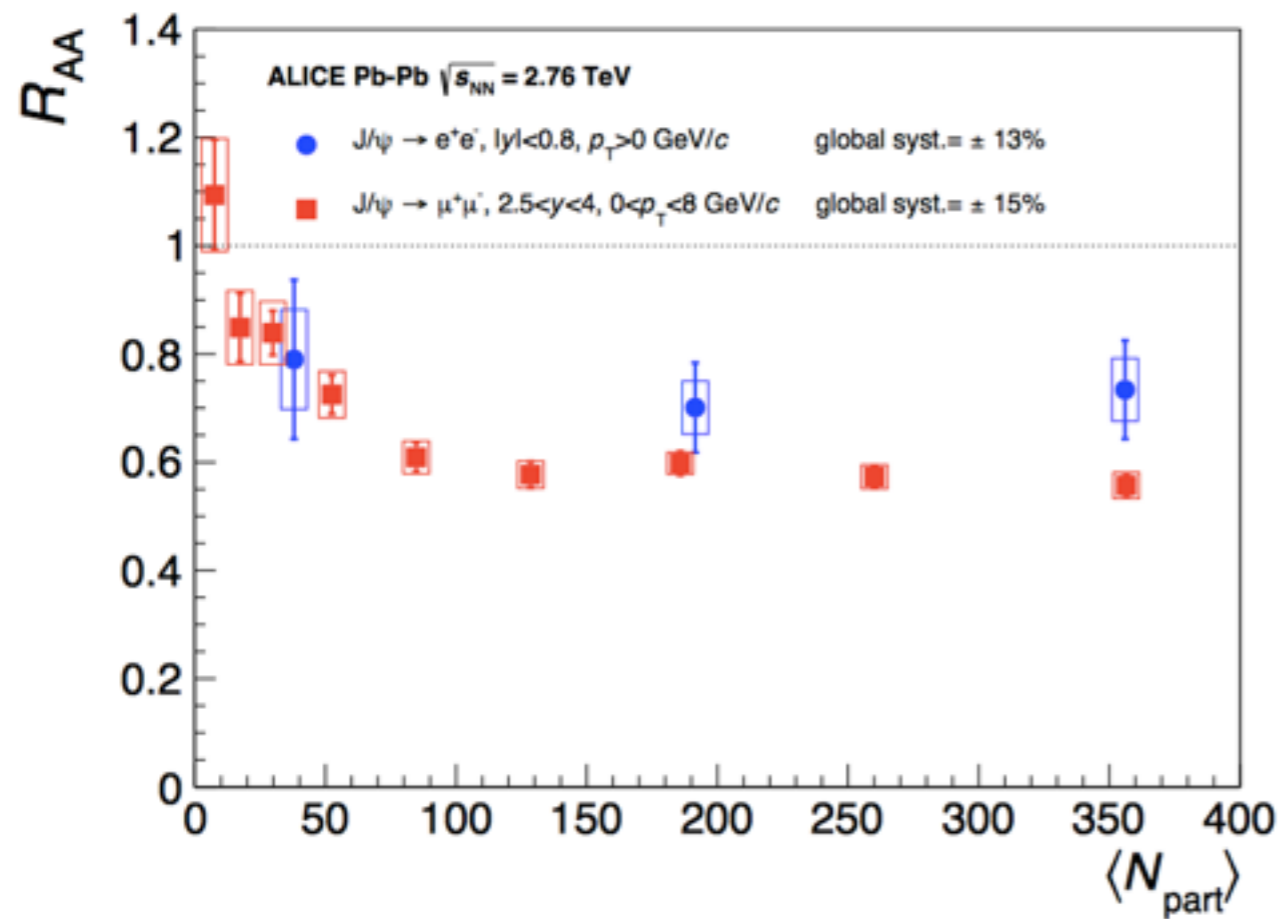
⊗ p_T dependence



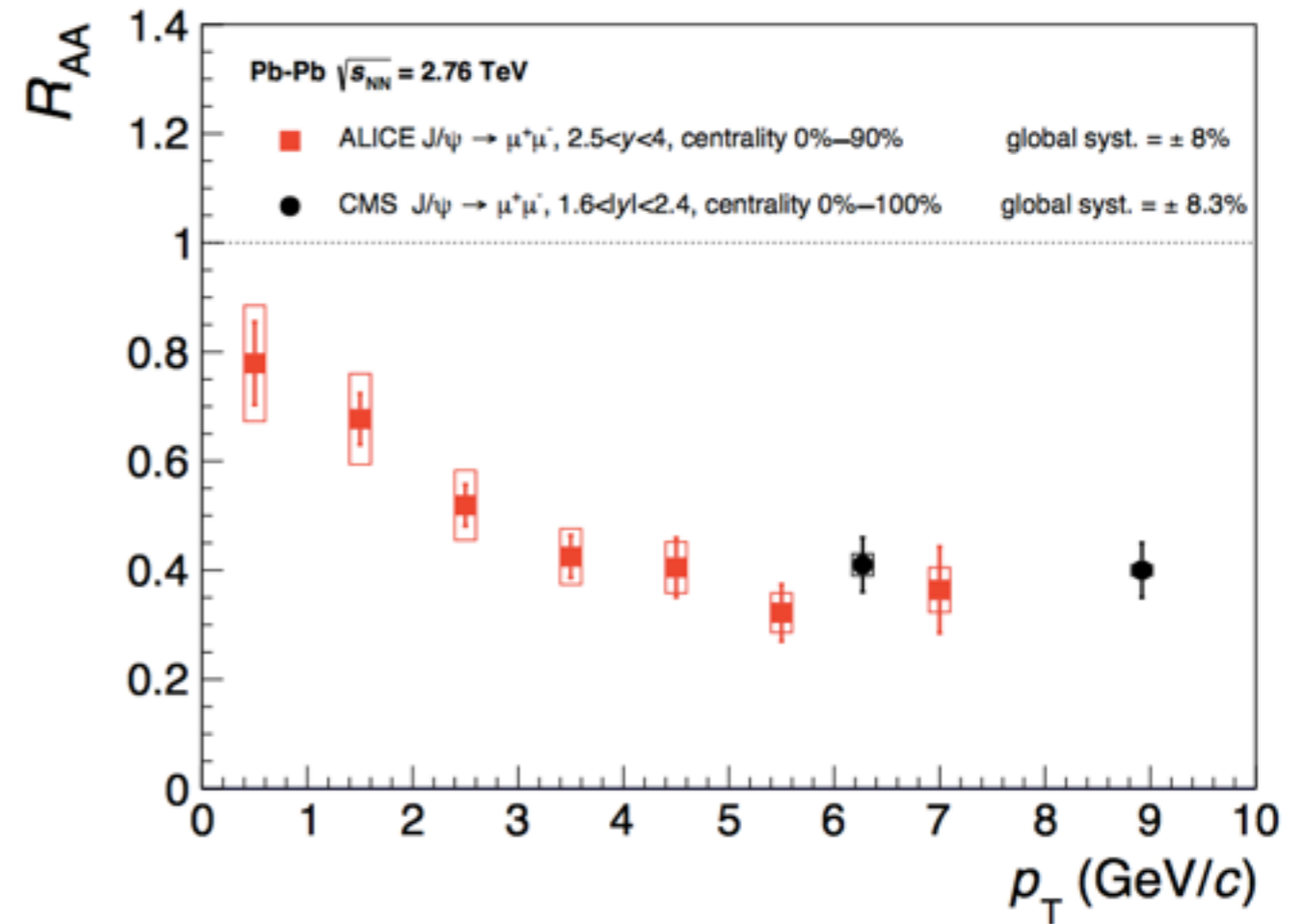
⊗ **Left** : In all rapidity bins at high p_T region, centrality dependent suppression is shown.

⊗ **Right** : In the forward region, lower p_T J/ψ has strong centrality dependence and less suppressed than high p_T J/ψ.

ALICE : $\mu\mu$ & ee channel

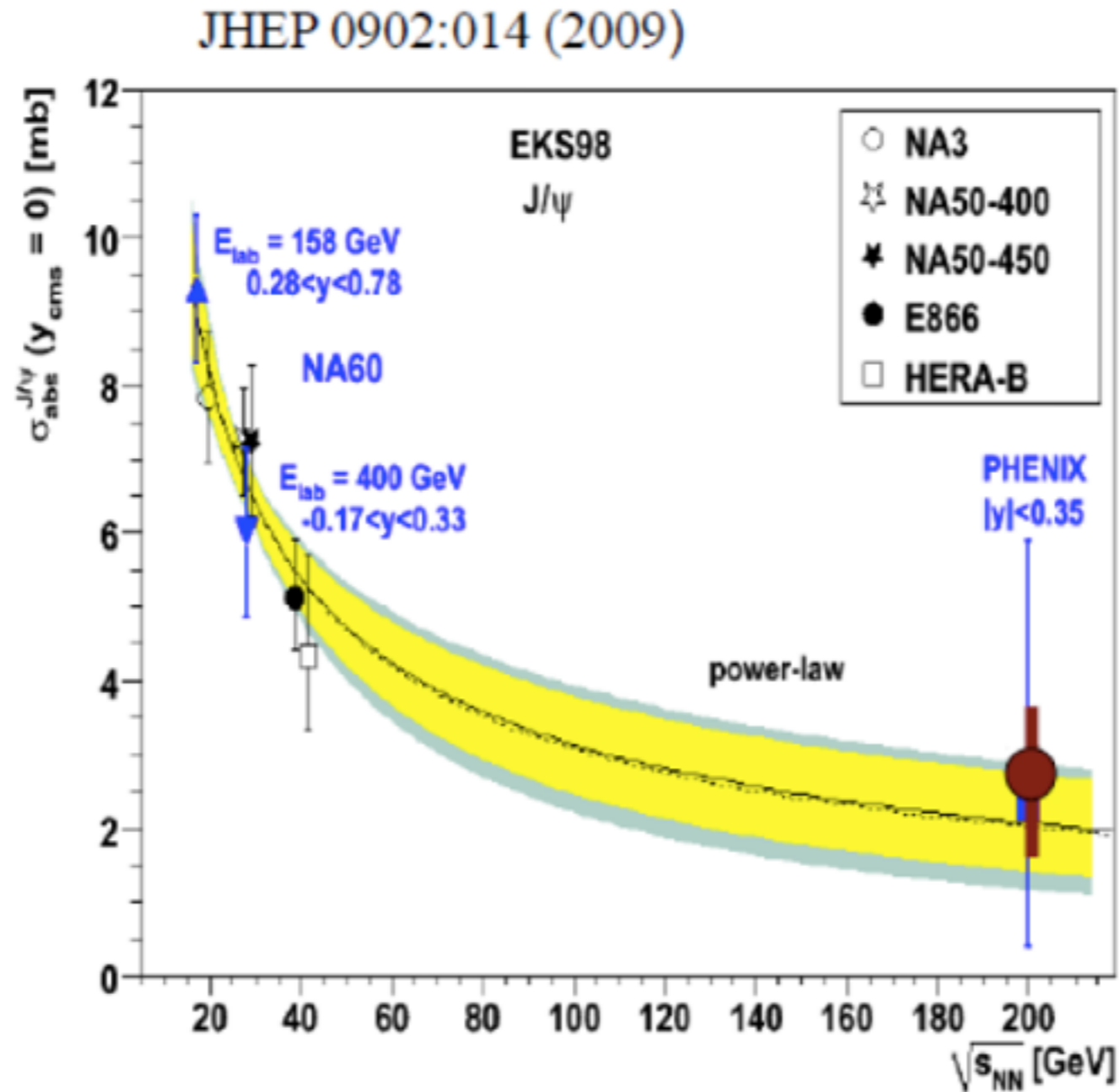


ALICE vs CMS



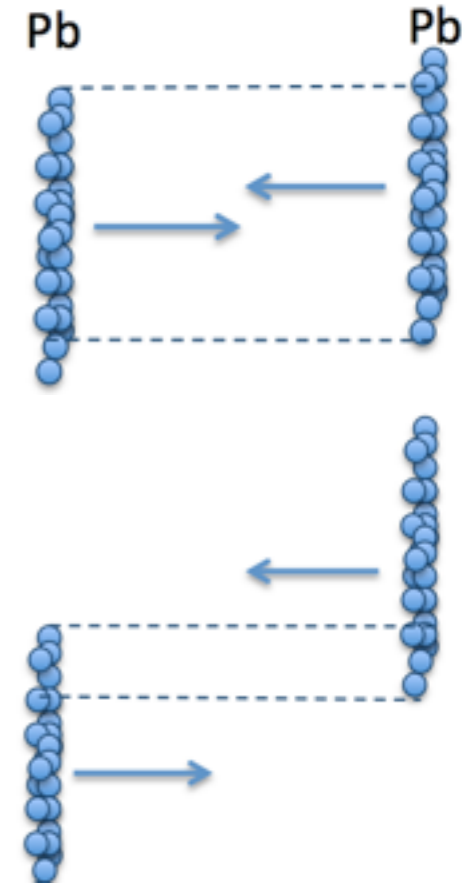
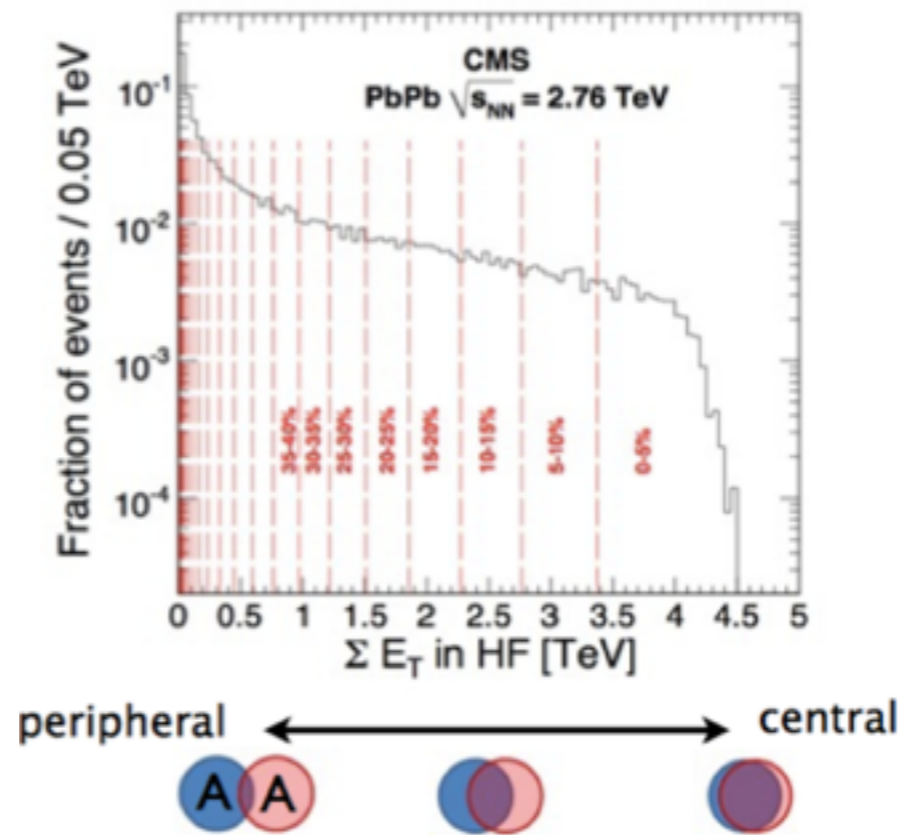
ALICE : centrality 0–90%

⊕ Nuclear absorption (break-up)



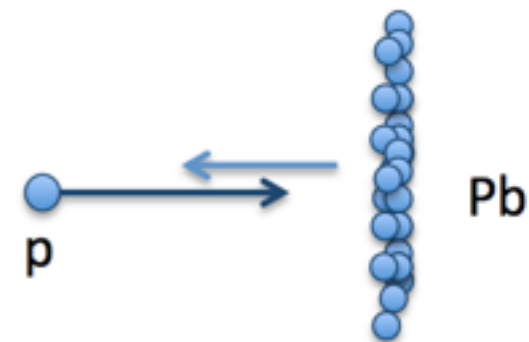
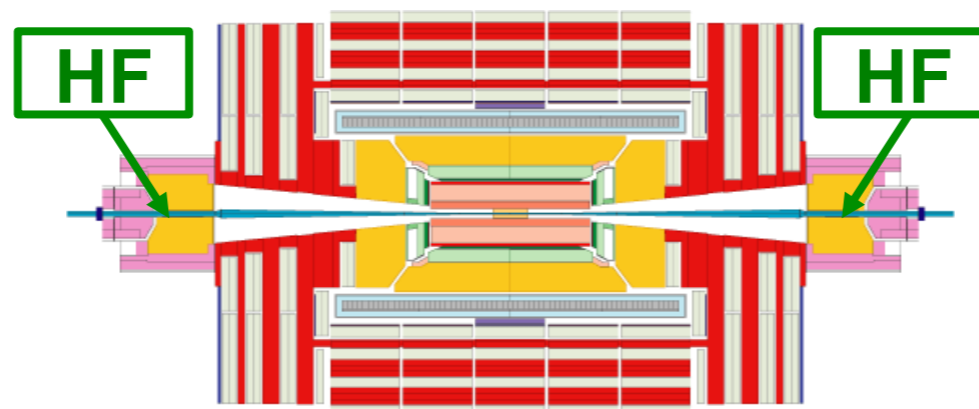
Centrality in PbPb

- related to the overlap fraction of the geometrical cross sections



Event-activity variables in pPb

- E_T^{HF} : raw transverse energy deposited in forward region HF ($4 < |\eta| < 5.2$)





Mandelstam invariant

$$S = (E_1 + E_2)^2 - (\vec{p}_1 + \vec{p}_2)^2$$

$$m_1 \sim m_2 \sim 1 \text{ GeV}/c^2$$

$$m_{1,2} \ll E_{1,2}$$

$$\Rightarrow S = (m_1^2 + \vec{p}_1^2) + (m_2^2 + \vec{p}_2^2) + 2E_1E_2 - (\vec{p}_1^2 + \vec{p}_2^2 + 2\vec{p}_1 \cdot \vec{p}_2)$$

$$\Rightarrow S = m_1^2 + m_2^2 + 2E_1E_2 - 2\vec{p}_1 \cdot \vec{p}_2$$

$$= -2|\vec{p}_1||\vec{p}_2|\cos(180^\circ)$$

$$\sim 2E_1E_2$$

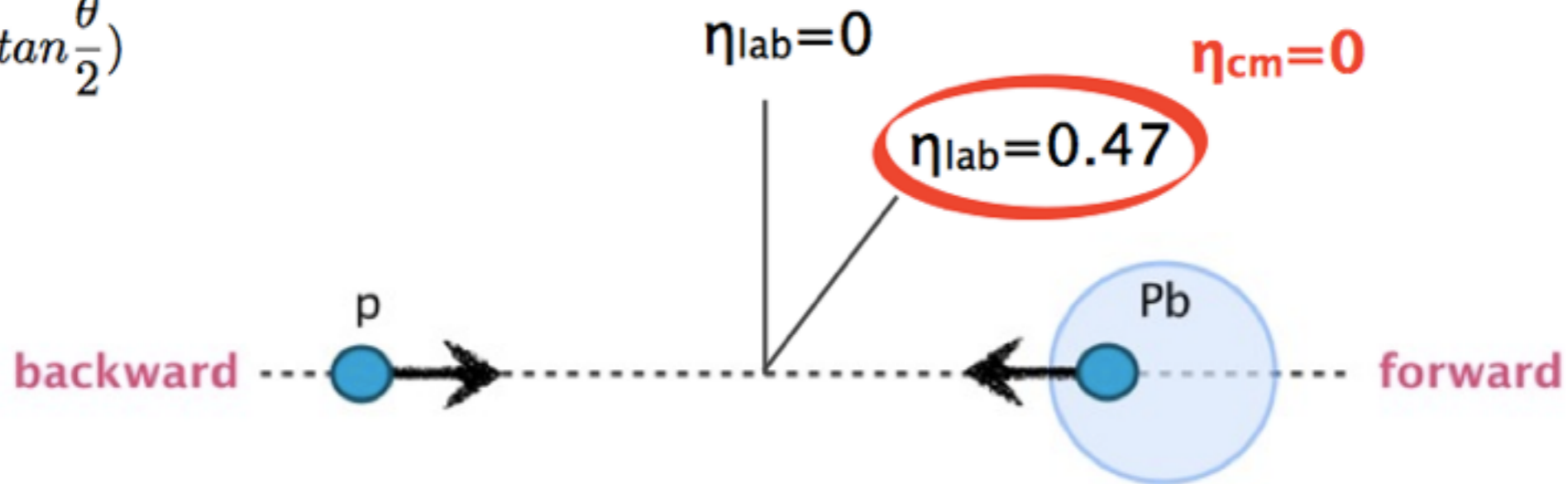
$$\Rightarrow S \simeq 4E_1E_2 \quad \Rightarrow \sqrt{S} = \sqrt{4 \times 4 \times 1.58} \simeq 5.02 \text{ TeV}$$

rapidity

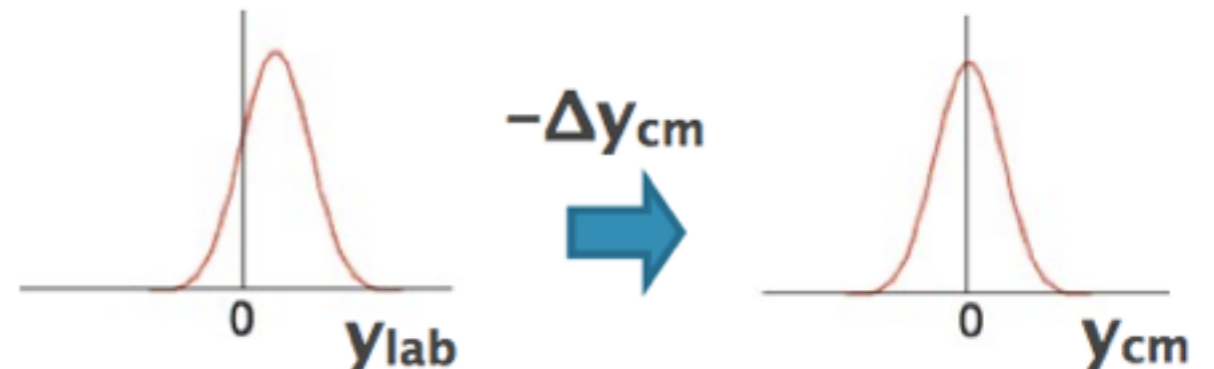
$$y = \frac{1}{2} \ln \left(\frac{E + p_z}{E - p_z} \right) = \frac{1}{2} \ln \left(\frac{1 + \beta}{1 - \beta} \right)$$

$$\Rightarrow \gamma = \frac{E}{mc^2} \Rightarrow \beta = \sqrt{1 - \frac{1}{\gamma^2}} \Rightarrow \begin{matrix} y_p = 9.05 \\ y_{Pb} = 8.12 \end{matrix} \Rightarrow \Delta y_{cm} = \frac{y_p - y_{Pb}}{2} \sim 0.47$$

$$\eta = -\ln \left(\tan \frac{\theta}{2} \right)$$



● rapidity distribution is invariant!



PRL 87 (2013) 034904

arXiv.0902.4154

- **Kopeliovich et. al.**

- : Shadowing – nDSg nPDFs
- : Cronin effect (parameterized from low energy data)
- : $\sigma_{\text{break-up}}$ (color dipole σ_{cc} from HERA)

- **Landberg et. al.**

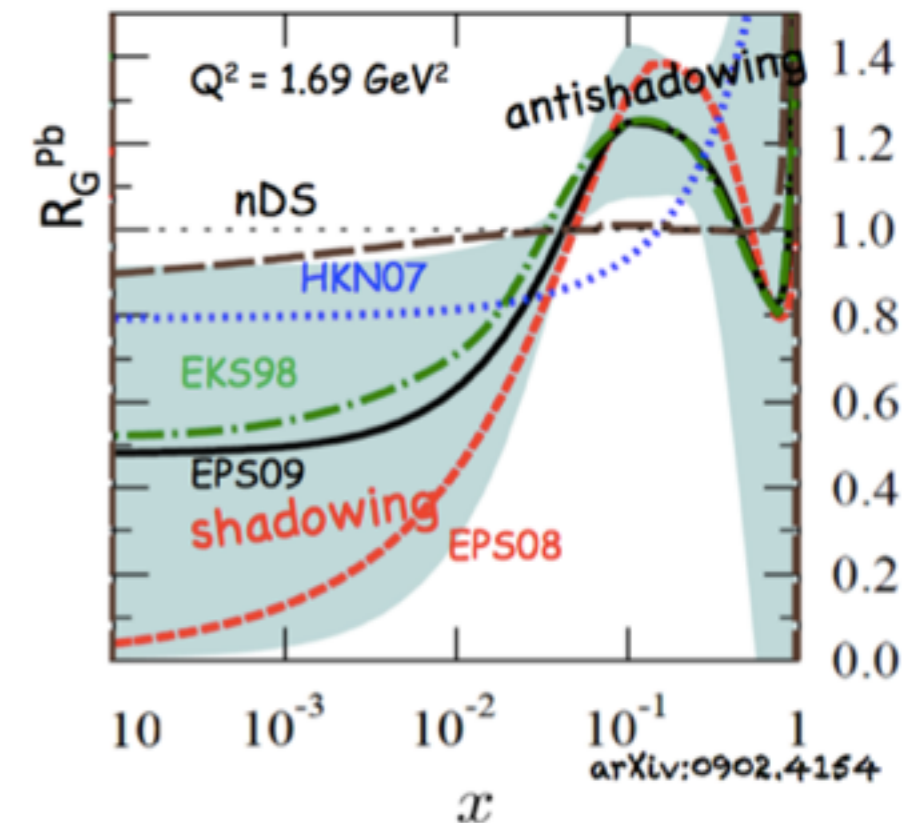
- : Shadowing – EKS98, nDSg, or EPS08 nPDFs
- : effective $\sigma_{\text{break-up}}$ (0, 2.6, 4.2, 6 mb)
- : No cronin effect
- : No initial state energy loss

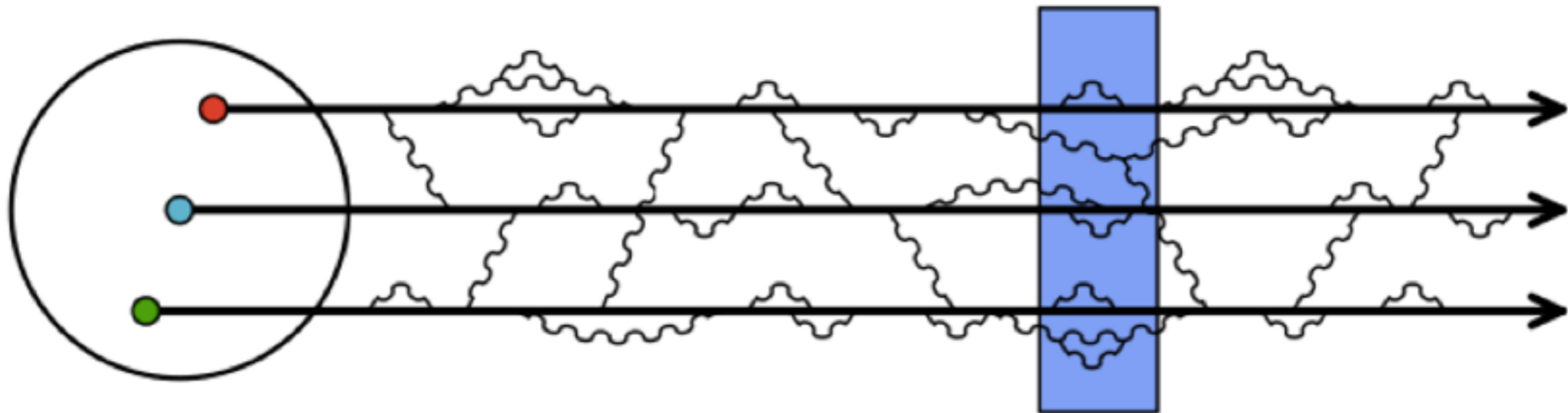
- **Nagle et. al.**

- : Shadowing – EPS09 nPDFs
- : $\sigma_{\text{break-up}}$ (0–20 mb)
- : Tried initial state energy loss

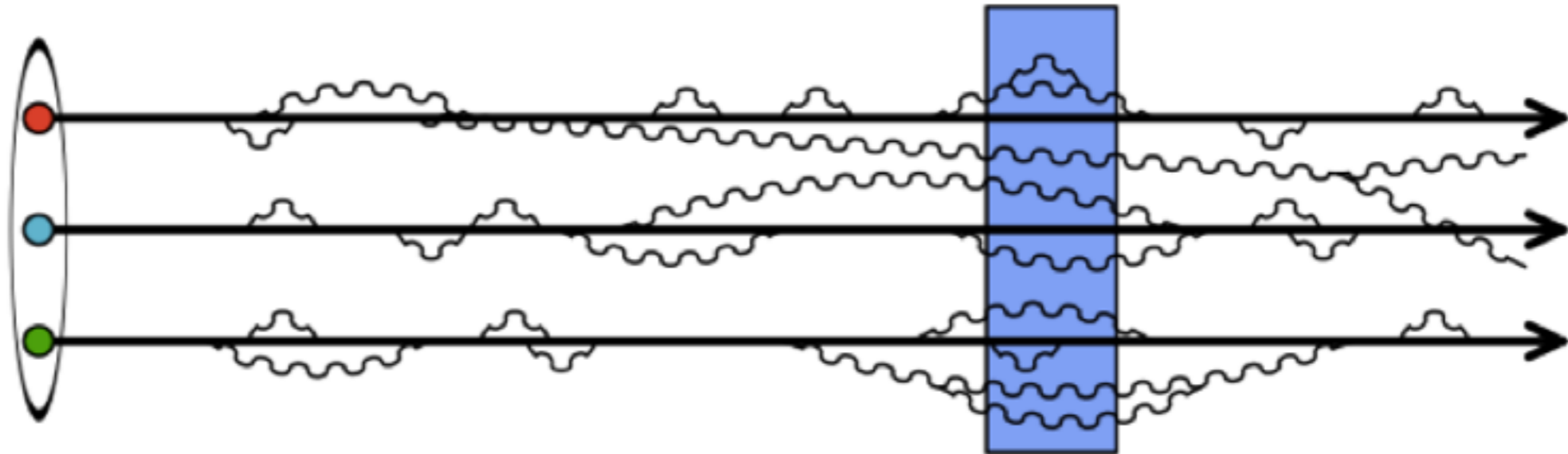
- **Kharzeev et. al.**

- : Shadowing – Coherent scattering / Color Glass Condensate model



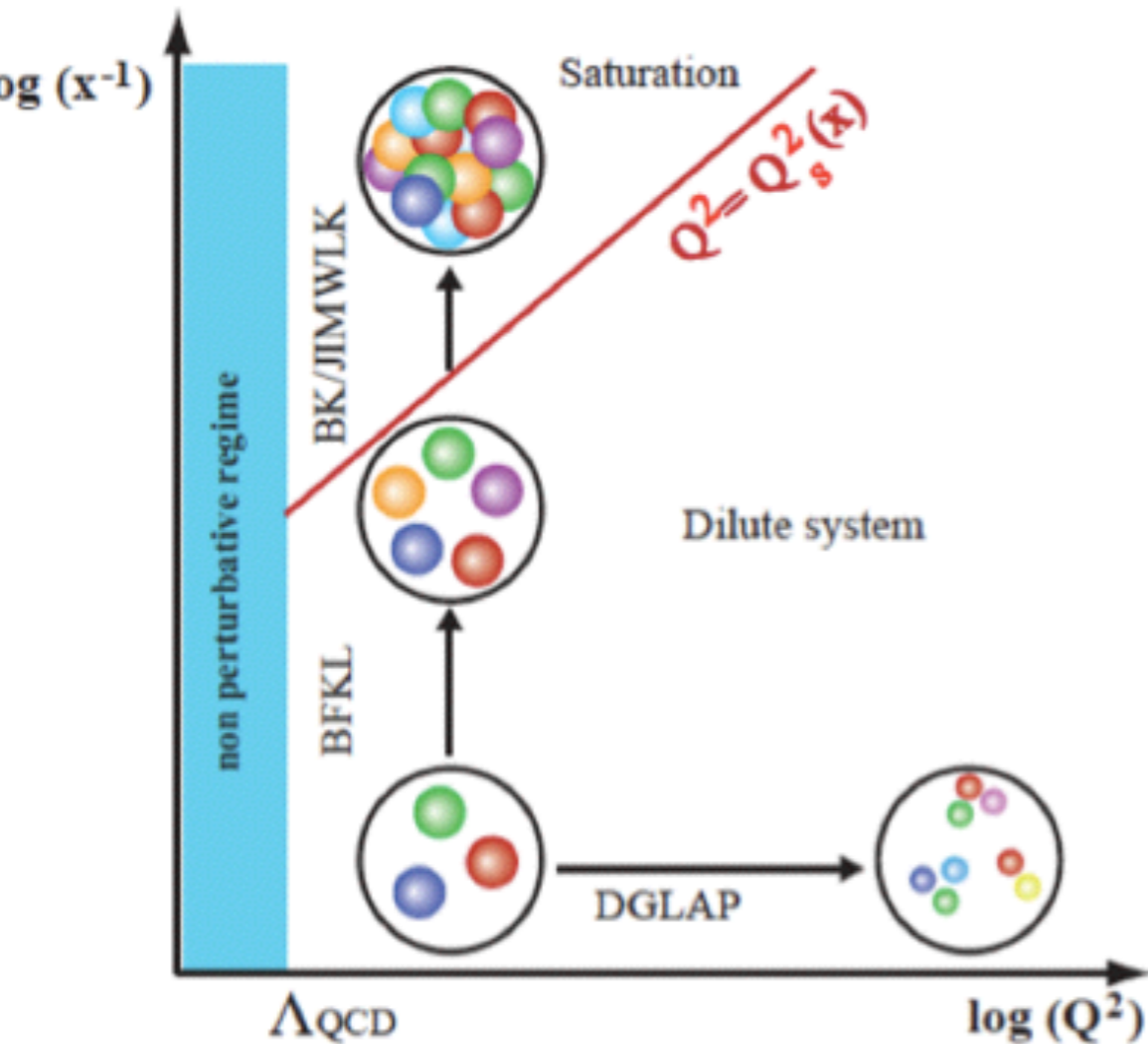


- Very complicated **non-perturbative** object...
- Contains **fluctuations at all space-time scales** smaller than its own size
- Only the fluctuations that are longer lived than the external probe participate in the interaction process
- The only role of short lived fluctuations is to renormalize the masses and couplings
- Interactions are very complicated if the constituents of the nucleon have a non trivial dynamics over time-scales comparable to those of the probe



- **Dilation** of all internal time-scales of the nucleon
- Interactions among constituents now take place over time-scales that are longer than the characteristic time-scale of the probe
 - ▷ **the constituents behave as if they were free**
- Many fluctuations live long enough to be seen by the probe. The nucleon appears **denser at high energy** (it contains more gluons)
- Pre-existing fluctuations are totally frozen over the time-scale of the probe, and act as static sources of new partons

Nucleon structure: x and Q^2 evolution



Saturation regime at low x and Q^2

Gribov, Levin, Ryskin, *Phys.Rept.*100,1983
 Mueller, Qiu, *Nucl.Phys.*B268,1986
 Blaizot, Mueller, *Nucl.Phys.*B289,1987



x -evolution

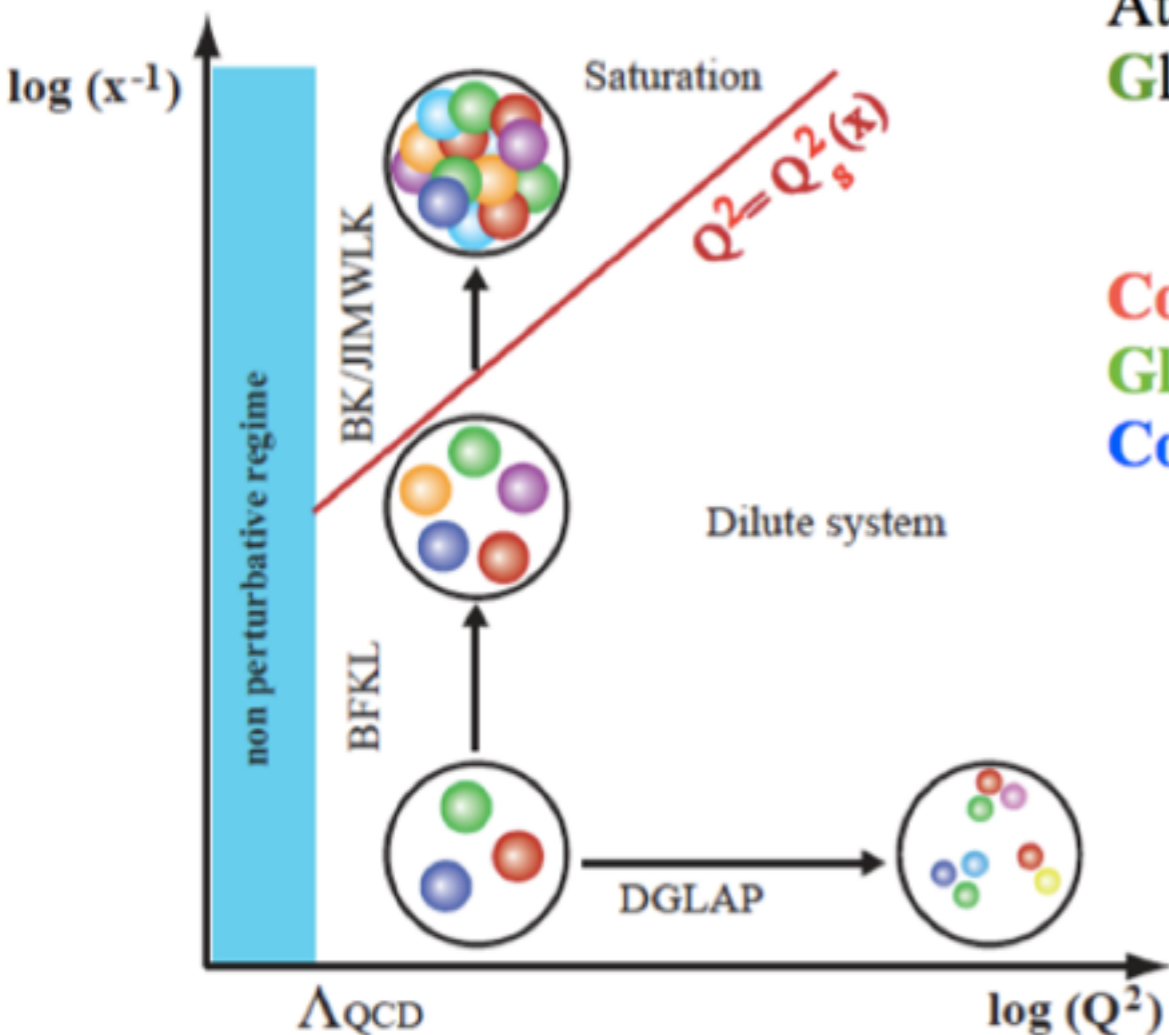
Large- x , linear \rightarrow BFKL

Balitsky, Fadin, Kuraev, Lipatov

Small- x , non linear \rightarrow JIMWLK/BK

*Jalilian-Marian, Iancu, McLerran, Weigert, Leonidov, Kovner
 Balitsky, Kovchegov*

$Q_s^2(x)$ defines the scale below which the gluon density saturates



At small- x , hadronic matter can be described by the **C**olour **G**lass **C**ondensate (CGC) model

McLerran, Venugopalan, Iancu, Leonidov, Mueller, ...

Colour gluons are coloured
Glass fast partons as “frozen” sources of soft gluons
Condensate high gluon occupancy

If **large** Q_s^2 : weakly coupled QCD-based theory but not necessarily perturbative \rightarrow CGC is a classical effective theory

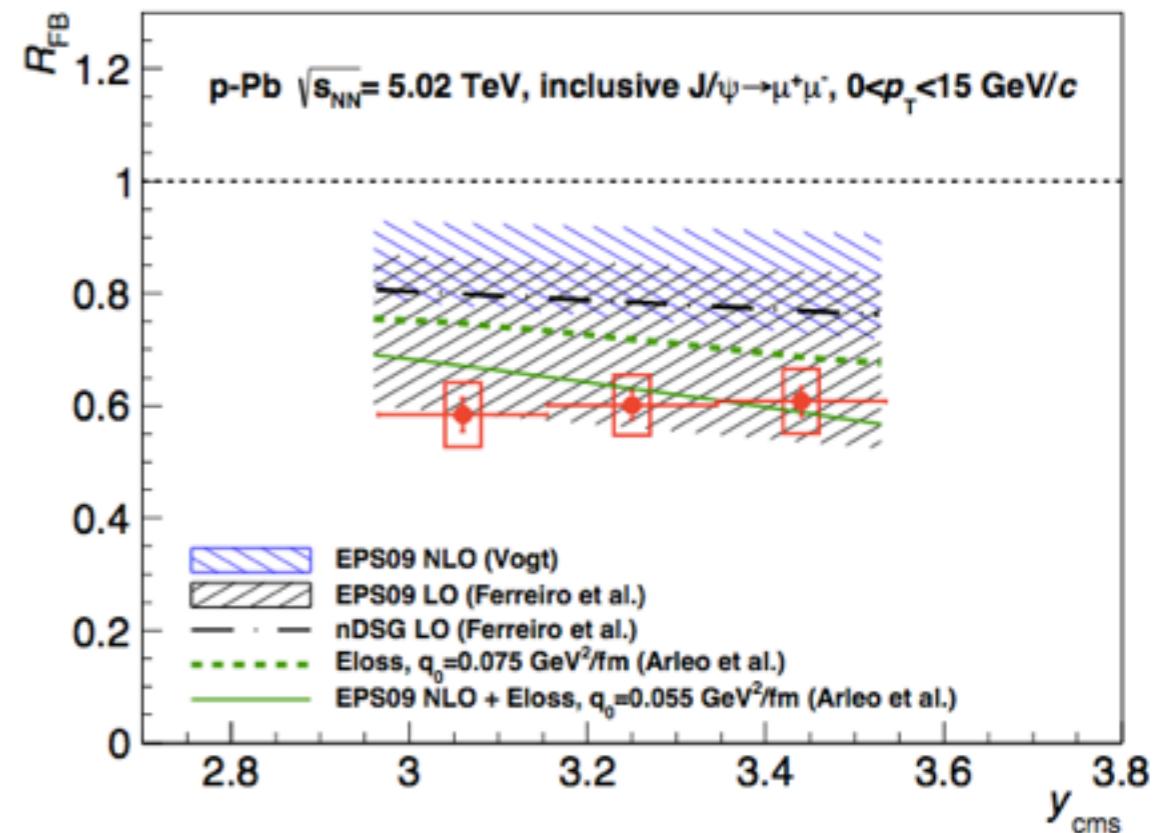
$$Q_s^2 \propto x^{-0.3} A^{1/3} \rightarrow \text{saturation for low } x \text{ (high } \sqrt{s}=1/x), \text{ large } A$$

$$\rightarrow Q_{s,\text{LHC}}^2 = 3 Q_{s,\text{RHIC}}^2$$

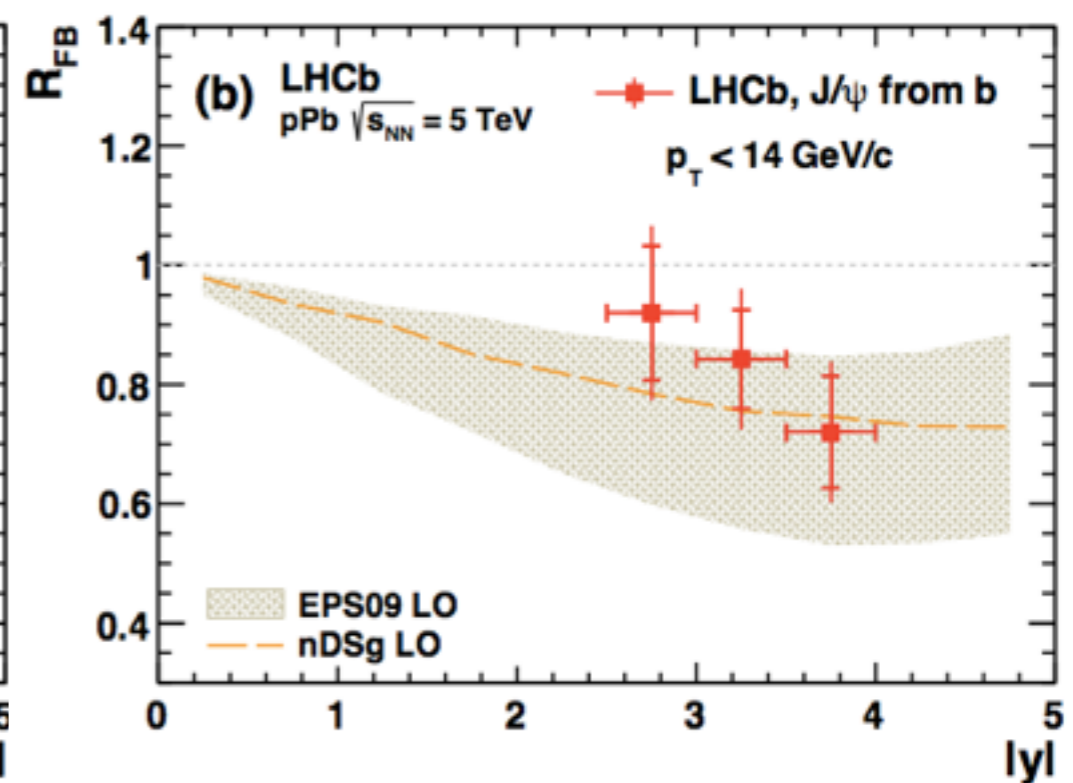
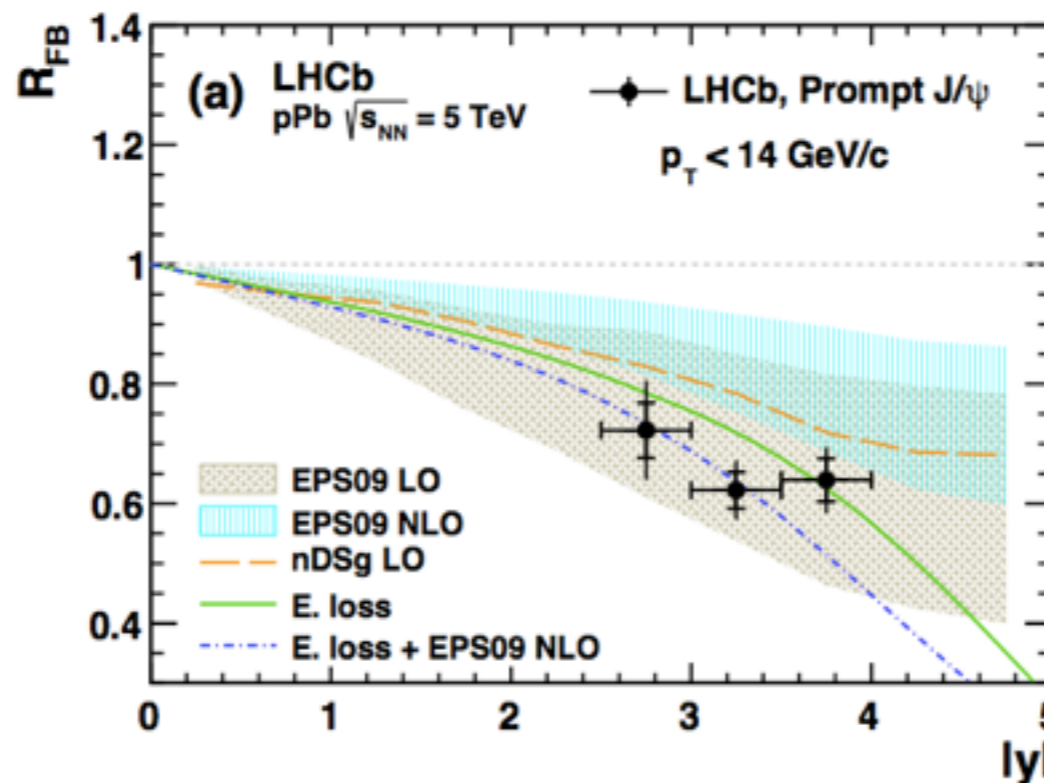
$$\rightarrow Q_{s,\text{Pb}}^2 = 6 Q_{s,\text{p}}^2$$

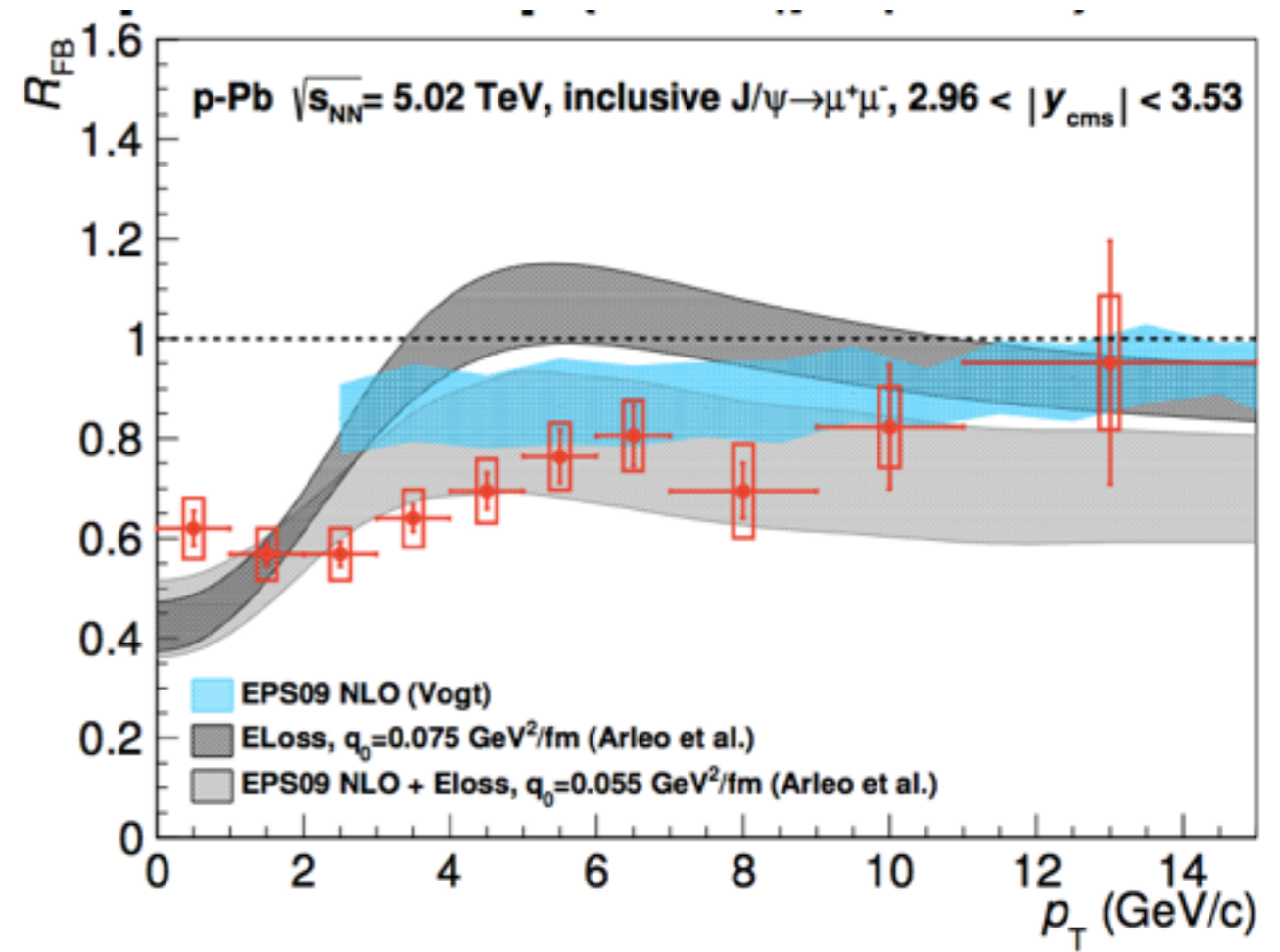
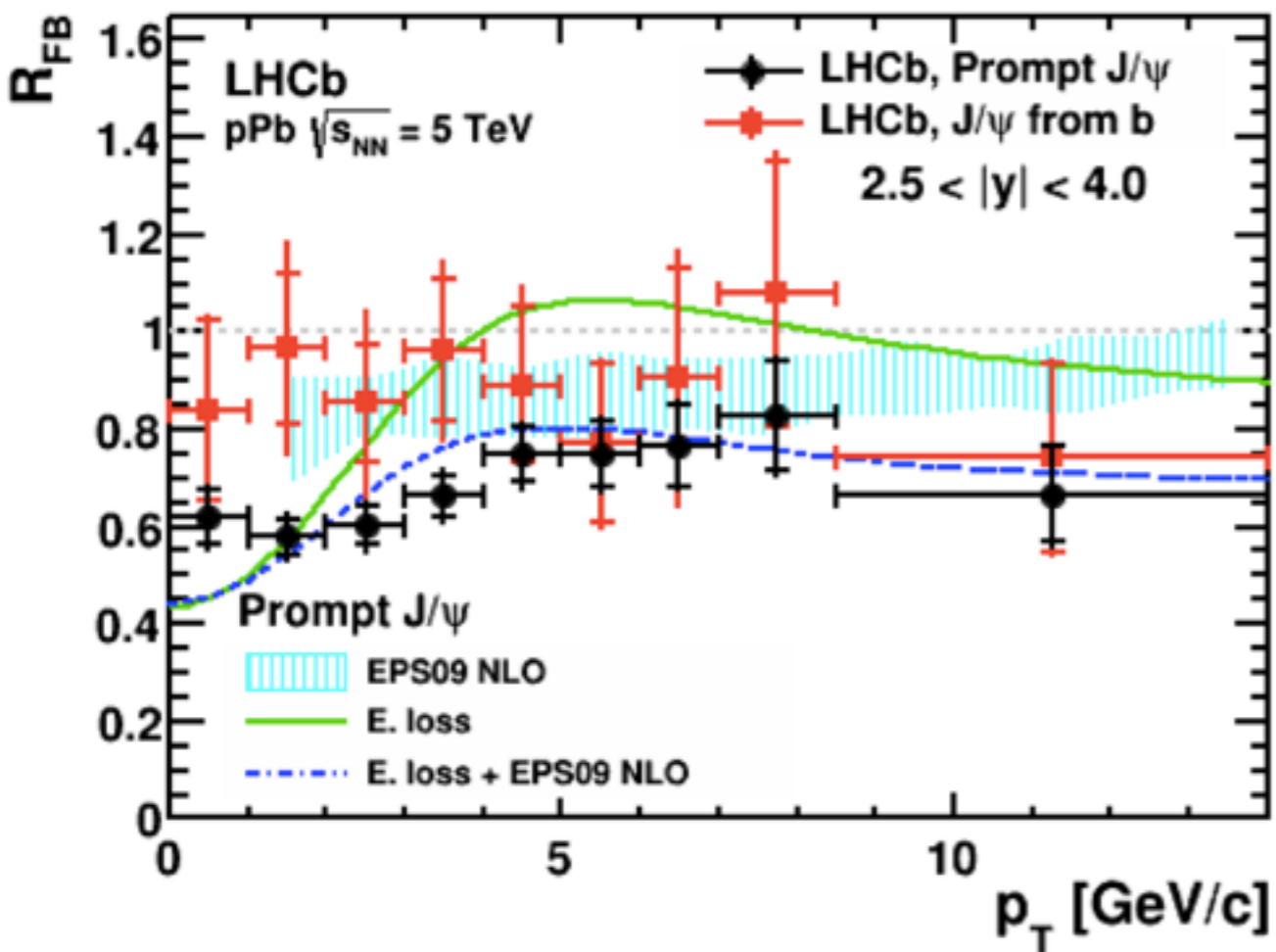
$$R_{FB} = \frac{N^{Forward}}{N^{Backward}}$$

- **ALICE**
 - $p_T < 15 \text{ GeV}/c$
 - $2.96 < |y_{CM}| < 3.53$



- **LHCb**
 - $p_T < 14 \text{ GeV}/c$
 - $2 < |y_{CM}| < 5$





-DER-73144

- ⊕ Shadowing + energy loss in good agreement with data
- ⊕ p_T dependence (suppression stronger at low p_T ?)

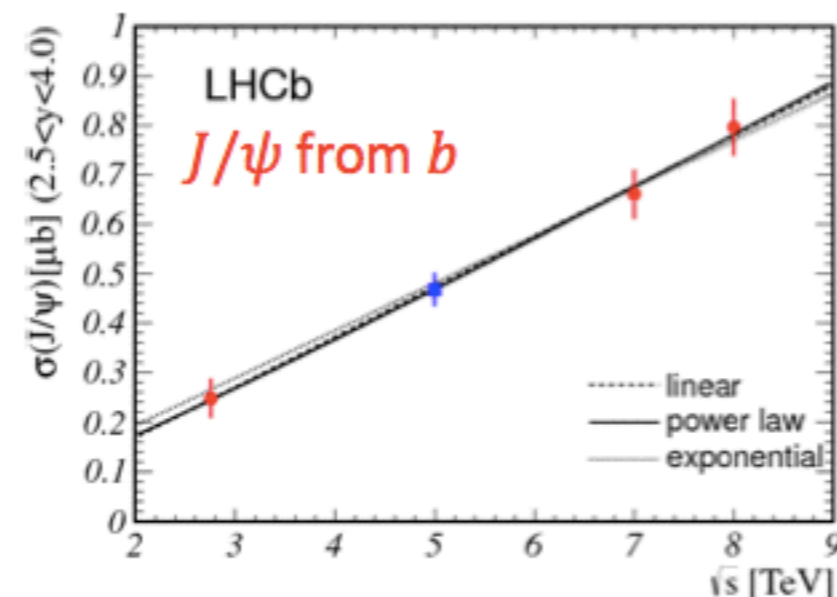
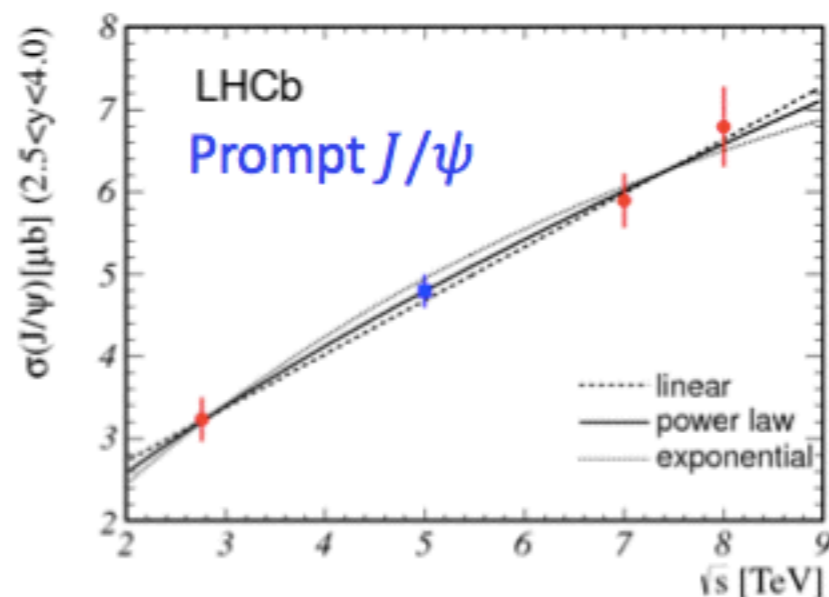
Reference cross-sections in pp at $\sqrt{s} = 5$ TeV

LHCb-CONF-2013-013

- Input to the determination of the nuclear modification factor R_{pPb}
- Interpolated from measurements at 2.76 TeV, 7 TeV and 8 TeV
- Three different fit functions used to interpolate

$$\begin{aligned}
 & (\sqrt{s}/p_0)^{p_1} \quad \longrightarrow \quad \text{adopted as nominal} \\
 & p_0 + p_1\sqrt{s} \\
 & p_0(1 - e^{p_1\sqrt{s}})
 \end{aligned}$$

- Discrepancy between the three interpolated values taken as systematics
- Checked against functions from LO-CEM and FONLL



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- **pp reference cross-section:** no measurement at required energy, need to interpolate. Two step procedure in both analyses.
 - Dimuon analysis [1]:
 - 1) \sqrt{s} interpolation: performed bin-per-bin in rapidity (or p_T) using available ALICE pp results at 2.76 and 7 TeV
 - 2) Further rapidity extrapolation: due to rapidity shift, p-Pb y_{cms} -range lies slightly outside the pp y_{cms} -range
 - Dielectron analysis:
 - 1) \sqrt{s} interpolation: performed using available PHENIX, CDF and ALICE results at $y \approx 0$
 - 2) p_T dependence: phenomenological scaling inspired by reference [2]