

A Pathway to a Transport Model for RAON

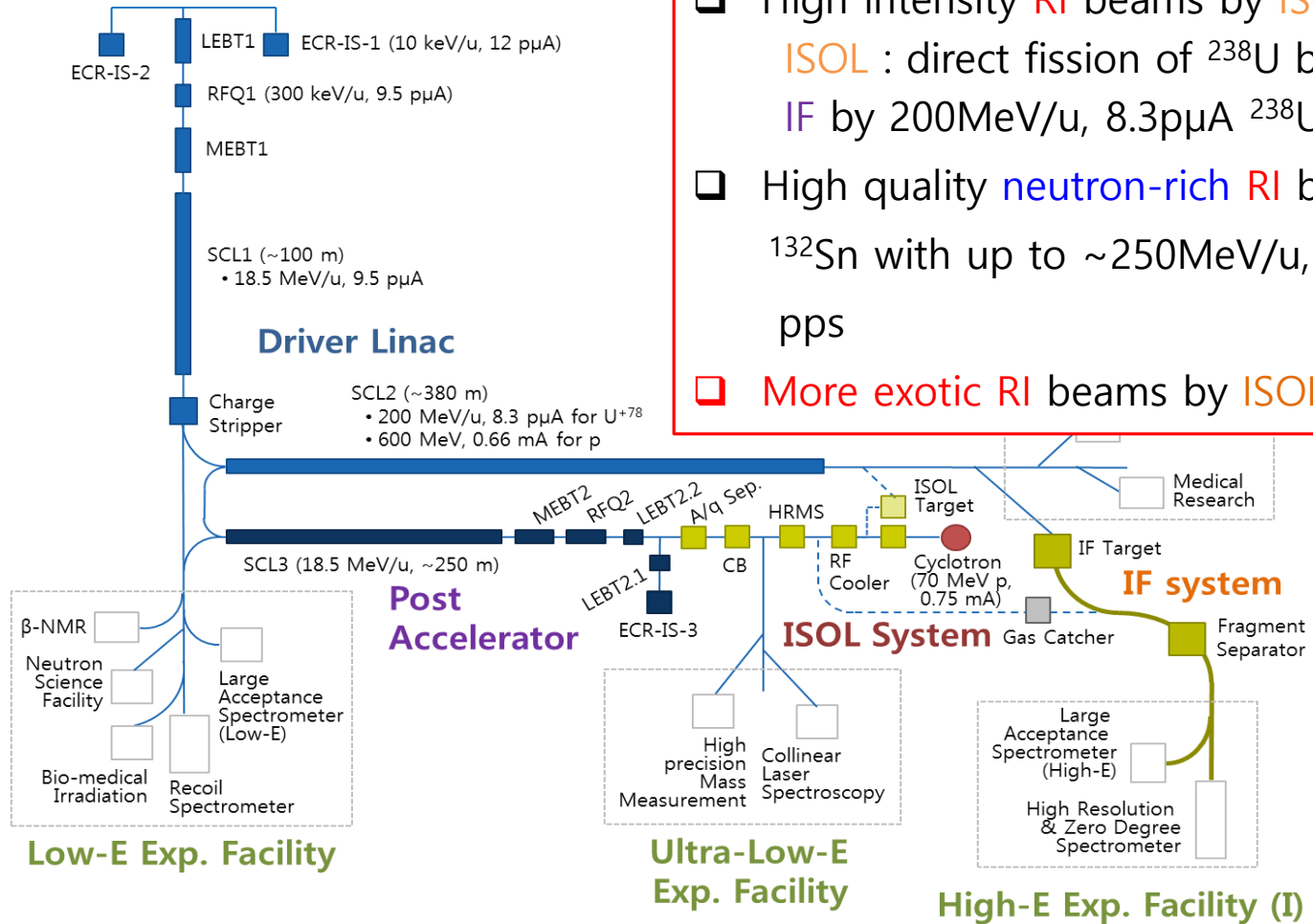
Kyungil Kim

Dec. 7th ,2013 ,HIM

Contents

- I. RAON and Nuclear Reactions
- II. Introduction to Transport Model
- III. Equation of State and Symmetry Energy
- IV. Transport Model for RAON
- V. Summary and Outlook

RAON and Nuclear Reactions



- ❑ High intensity RI beams by ISOL & IF
 - ISOL : direct fission of ^{238}U by p 70MeV
 - IF by 200MeV/u, 8.3pμA ^{238}U
- ❑ High quality neutron-rich RI beams
 - ^{132}Sn with up to ~250MeV/u, up to 10^7 pps
- ❑ More exotic RI beams by ISOL+IF

RAON and Nuclear Reactions

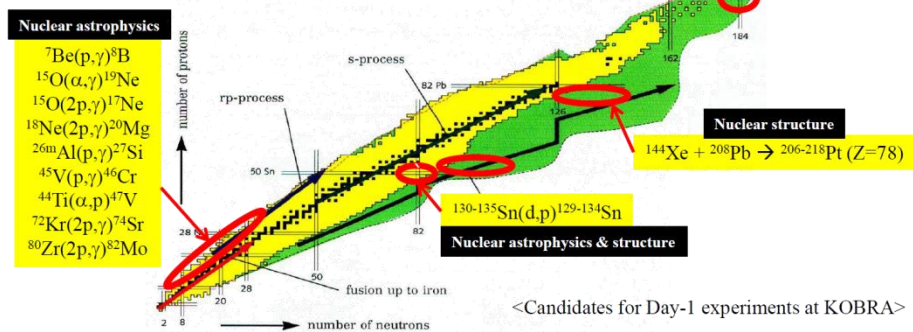
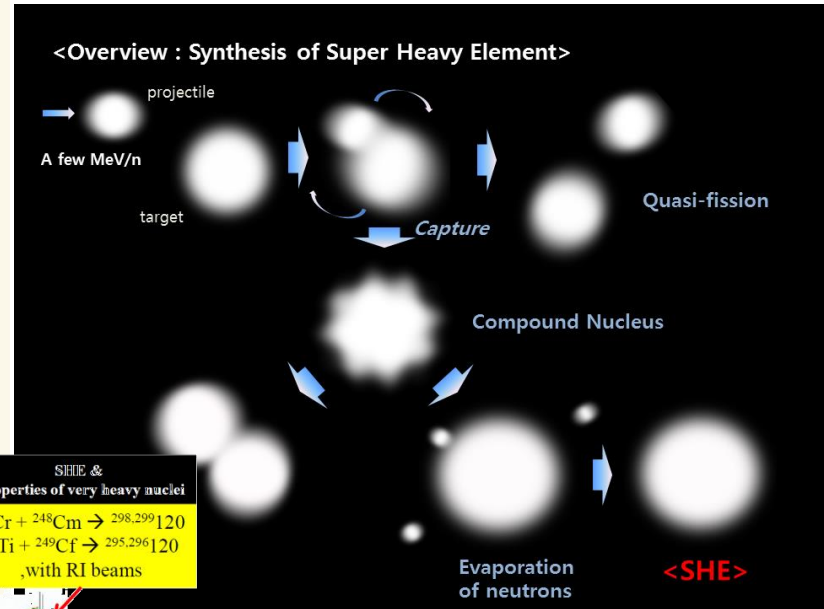
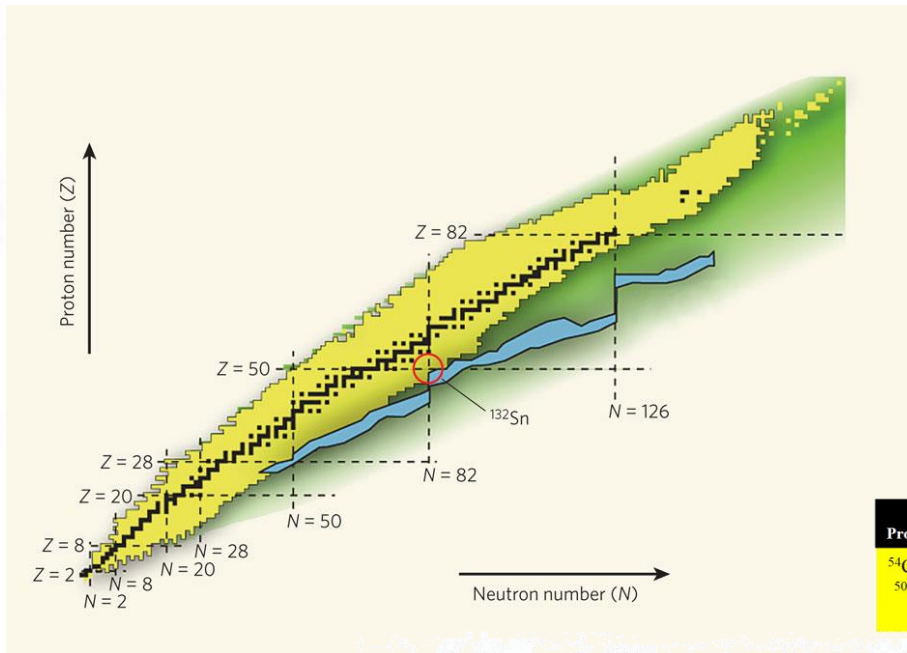
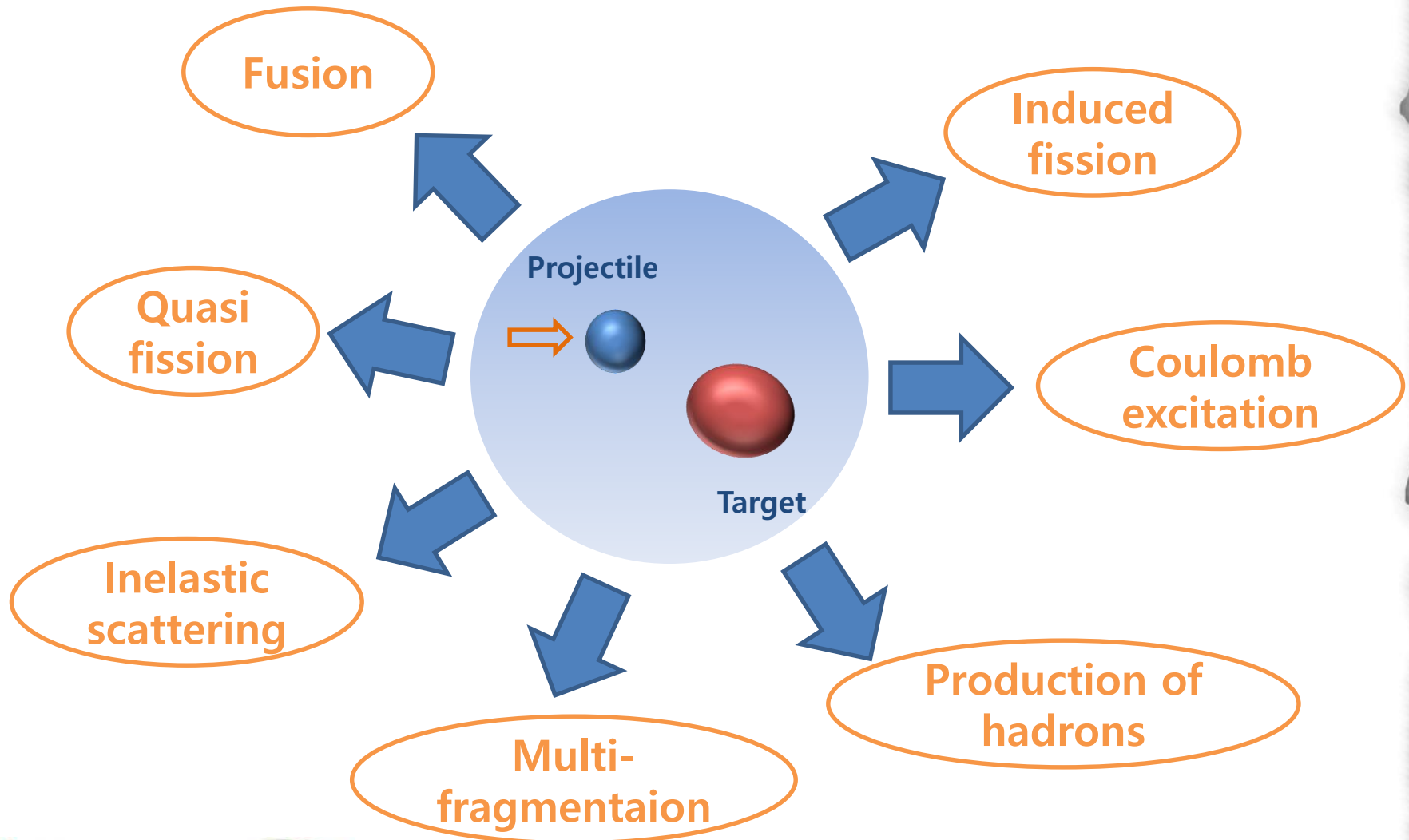


figure from Y.K.Kim's presentation

RAON and Nuclear Reactions



RAON and Nuclear Reactions

Direct Reactions

Compound Nuclear Reactions

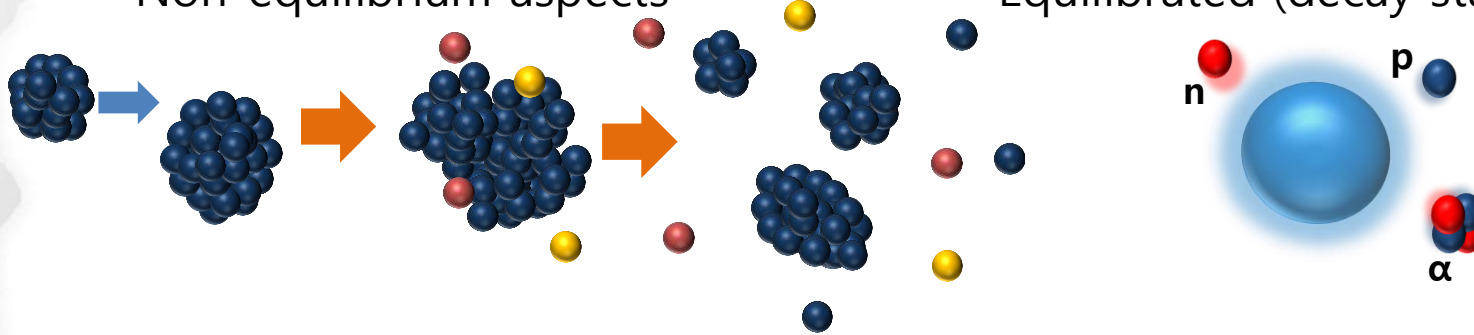
$\sim 10^{-20}$ sec

Time

$\sim 10^{-16}$ sec

Non-equilibrium aspects

Equilibrated (decay statistically)



Low and intermediate energy regime : $\sim 100 \text{ MeV/n}$

- *fusion, quasi-fission (multinucleon transfer), fragmentation*

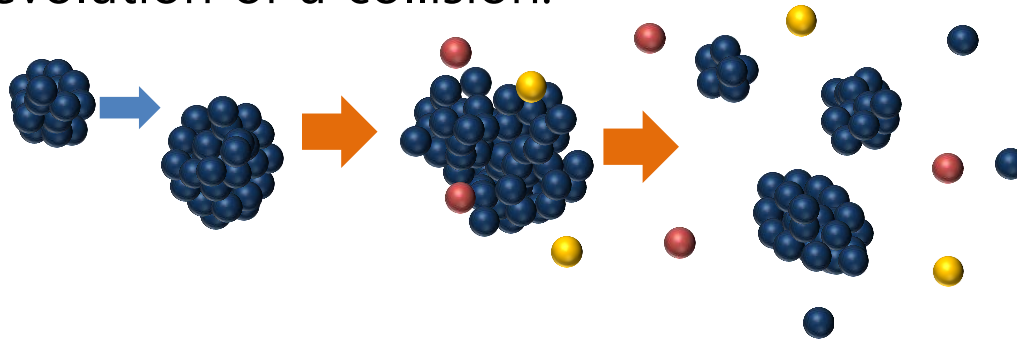
Relativistic regime : $100 \text{ MeV/n} \sim \text{a few GeV/n}$

- *production of hadrons, collective flow phenomena*

Kinetic Energy	$\lambda = h/p$
1 MeV	$3 \times 10^{-14} \text{ m}$
10 MeV	$9 \times 10^{-15} \text{ m}$
100 MeV	$3 \times 10^{-15} \text{ m}$
1 GeV	$7 \times 10^{-16} \text{ m}$

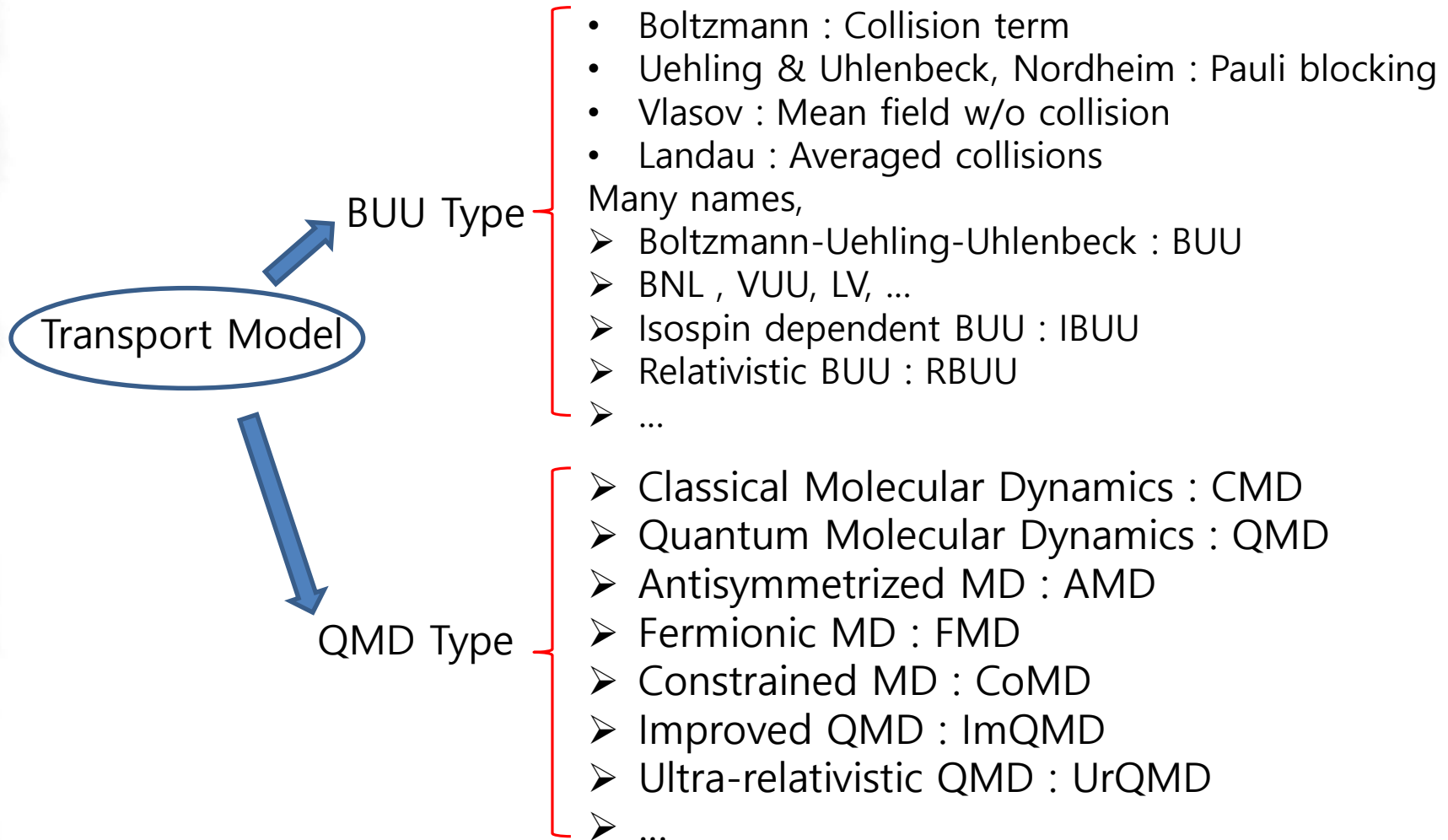
Transport Model?

Transport model: Model to treat non-equilibrium aspects of the temporal evolution of a collision.



- ✓ Many-body problem with nucleons
- ✓ Numerical simulation
- ✓ Direct reaction regime (compound nuclear reaction -> statistical model)
- ✓ Different methods with different energies

Various Codes

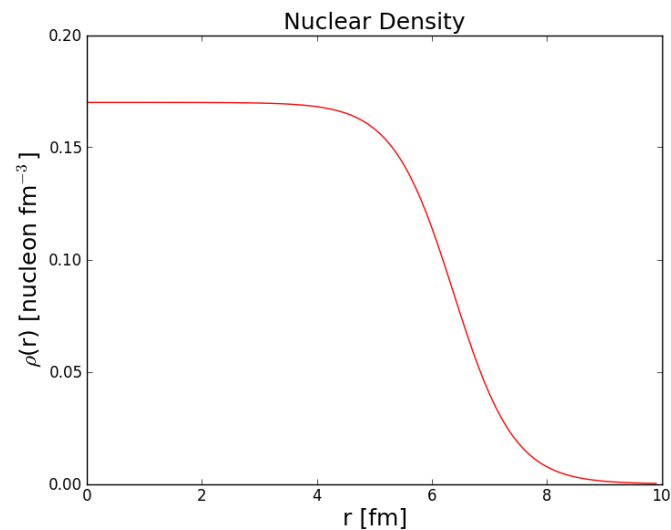


Introduction to Transport Model

1. Initialization

$$\rho = \rho_0 \left[1 + \exp \left\{ \frac{r - R}{a} \right\} \right]^{-1}$$

$$R = 1.12A^{1/3} \quad a = 0.53 \text{ [fm]}$$



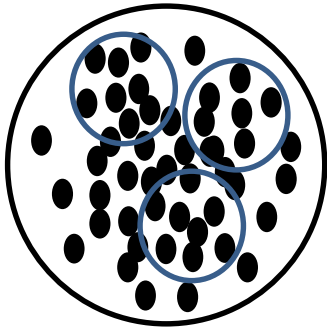
<Phase Space Density and Wigner Transformation>

- Both a position and a momentum of nucleons are needed.
 - **Phase space density!!**
- **Wigner transformation**

$$f(r, p; t) = \int d^4\zeta \exp(ip_\mu \zeta^\mu) \tilde{f} \left(r + \frac{\zeta}{2}, r - \frac{\zeta}{2} \right), \quad \zeta = r_1 - r_2, \quad r = \frac{r_1 + r_2}{2}$$

Introduction to Transport Model

➤ BUU Model



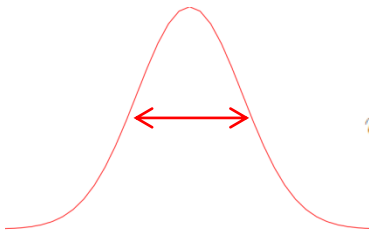
- Semi-classical approach : point particles
- Test particle method : 1~500 test particles per nucleon

$$f(\mathbf{r}, \mathbf{p}; t) = \frac{1}{N_{TP}} \sum_{i=1}^{AN_{TP}} \delta(\mathbf{r} - \mathbf{r}_i(t)) \delta(\mathbf{p} - \mathbf{p}_i(t))$$



Randomly distributed !

➤ QMD Model



$L = 1.08 \text{ fm}^2$
 $\rightarrow r_N = 1.8 \text{ fm}$

- Gaussian wave packets

$$\psi_i(\mathbf{r}, \mathbf{p}_{i0}, \mathbf{r}_{i0}, t) = \frac{\exp\{i[\mathbf{p}_{i0} \cdot (\mathbf{r} - \mathbf{r}_{i0}) - p_{i0}^2 t / 2m]\}}{[\sqrt{\pi/2L} 2L(t)]^{2/3}} \exp\{-[\mathbf{r} - \mathbf{r}_{i0} - \mathbf{p}_{i0} t / m]^2 / 4L(t)\}$$

$$f(\mathbf{r}, \mathbf{p}, t) = \frac{1}{(2\pi)^3} \int e^{-i\mathbf{p} \cdot \mathbf{r}_{12}} \psi_i(\mathbf{r} + \mathbf{r}_{12}/2, t) \psi_i^*(\mathbf{r} - \mathbf{r}_{12}/2, t) d^3 r_{12}$$

$$= \frac{1}{\pi^3} \exp[-(\mathbf{r} - \mathbf{r}_{i0} - \mathbf{p}_{i0} t / m)^2 / sL - (\mathbf{p} - \mathbf{p}_{i0})^2 \cdot 2L]$$

Introduction to Transport Model

2. Propagation

Boltzmann-Uehling-Uhlenbeck equation

$$\frac{df}{dt} = \underbrace{\frac{\partial f}{\partial t} + \frac{\vec{p}}{m} \vec{\nabla}^{(r)} f}_{\text{drift term}} - \underbrace{\vec{\nabla}^{(r)} U \vec{\nabla}^{(p)} f}_{\text{accelerating term}} = \underbrace{I_{coll}}_{\text{collision term}}$$

cf) RBUU eq. : $[p_{\mu}^* \partial_x^{\mu} + (p_{\nu}^* F^{\mu\nu} + m^* (\partial_x^{\mu} m^*)) \partial_{\mu}^{p^*}] f(x, p^*) = \mathcal{I}_{coll}$

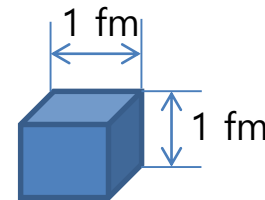
➤ BUU Model

Equation of motion

$$\frac{\partial r_i}{\partial t} = \frac{p_i}{m}, \quad \frac{\partial p_i}{\partial t} = -\nabla U|_{r_i}$$

Density dependent mean field

$$U(\rho) = \alpha(\rho/\rho_0) + \beta(\rho/\rho_0)^{\gamma}$$



(Unit box for density estimation)

Introduction to Transport Model

➤ QMD Model

Equation of motion

$$\frac{d}{dt}\mathbf{r}_i = \{\mathbf{r}_i, \mathcal{H}\}, \quad \frac{d}{dt}\mathbf{p}_i = \{\mathbf{p}_i, \mathcal{H}\} \quad \mathcal{H}\{\mathbf{r}_n, \mathbf{p}_n\} = \sum_{i=1}^A \frac{\mathbf{p}_i^2}{2m_i} + \sum_{i<j} V(|\mathbf{r}_i - \mathbf{r}_j|)$$

Nucleon-Nucleon interaction

: Skyrme , Volkov , Gogny , ...

ex) Skyrme force

$$V^{\text{loc}} = t_1 \delta(\mathbf{r}_1 - \mathbf{r}_2) + t_2 \delta(\mathbf{r}_1 - \mathbf{r}_2) \delta(\mathbf{r}_1 - \mathbf{r}_3)$$

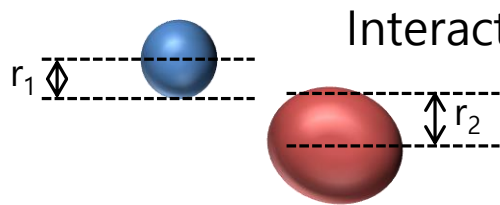
➡ $U^{\text{loc}} = \alpha(\rho/\rho_0) + \beta(\rho/\rho_0)^\gamma$

ex) Gogny force

$$v_{ij} = \sum_{k=1}^2 v_{0k} (W_k + B_k P_\sigma - H_k P_\tau - M_k P_\sigma P_\tau) \exp[-(\mathbf{r}_i - \mathbf{r}_j)^2/a_k^2] \\ + \frac{t_\rho}{6} (W_\rho + B_\rho P_\sigma - H_\rho P_\tau - M_\rho P_\sigma P_\tau) \rho(\mathbf{r}_i)^\sigma \delta(\mathbf{r}_i - \mathbf{r}_j),$$

Introduction to Transport Model

3. Collision



Interaction radius = $\pi(r_1 + r_2)^2$

$$b < \frac{1}{\pi} \sqrt{\sigma^{tot}(\sqrt{s})} \Rightarrow \text{Scatter with probability 1 !}$$

Pauli blocking factor

$$I_{coll} = \int d\mathbf{v}_2 d\mathbf{v}_{1'} d\mathbf{v}_{2'} |\mathbf{v}_2 - \mathbf{v}_1| \sigma(\Omega) (2\pi)^3 \delta(\mathbf{p}_1 + \mathbf{p}_2 - \mathbf{p}_{1'} - \mathbf{p}_{2'})$$

$$\underline{[f_{1'} f_{2'} (1 - f_1)(1 - f_2) - f_1 f_2 (1 - f_{1'}) (1 - f_{2'})]}$$

In-medium cross-section

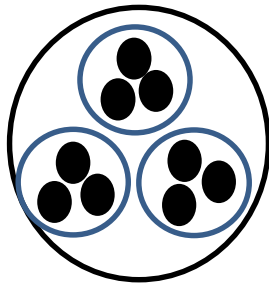
<Elastic and In-elastic scattering>

NN \rightarrow NN , NN \rightarrow NΛK , NN \rightarrow NΔ ,
 NΔ \rightarrow NΔ , Δ \rightarrow Nπ , ...

Introduction to Transport Model

➤ BUU Model

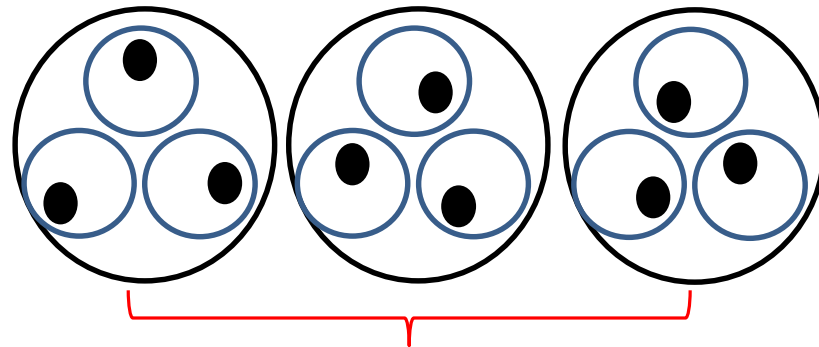
<Full ensemble method>



$$\sigma_{TP} = \sigma_{NN} / N_{TP}$$

$$\sim (AN_{TP})^2$$

<Parallel ensemble method>

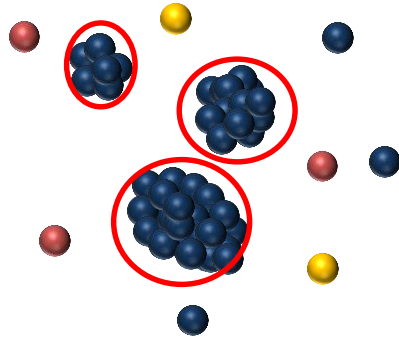


N_{TP} ensembles

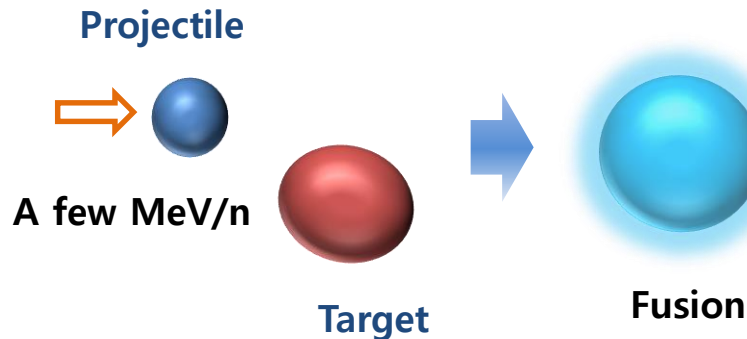
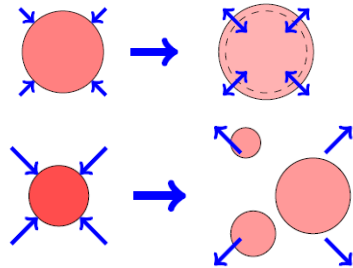
$$\sim N_{TP} A^2$$

Introduction to Transport Model

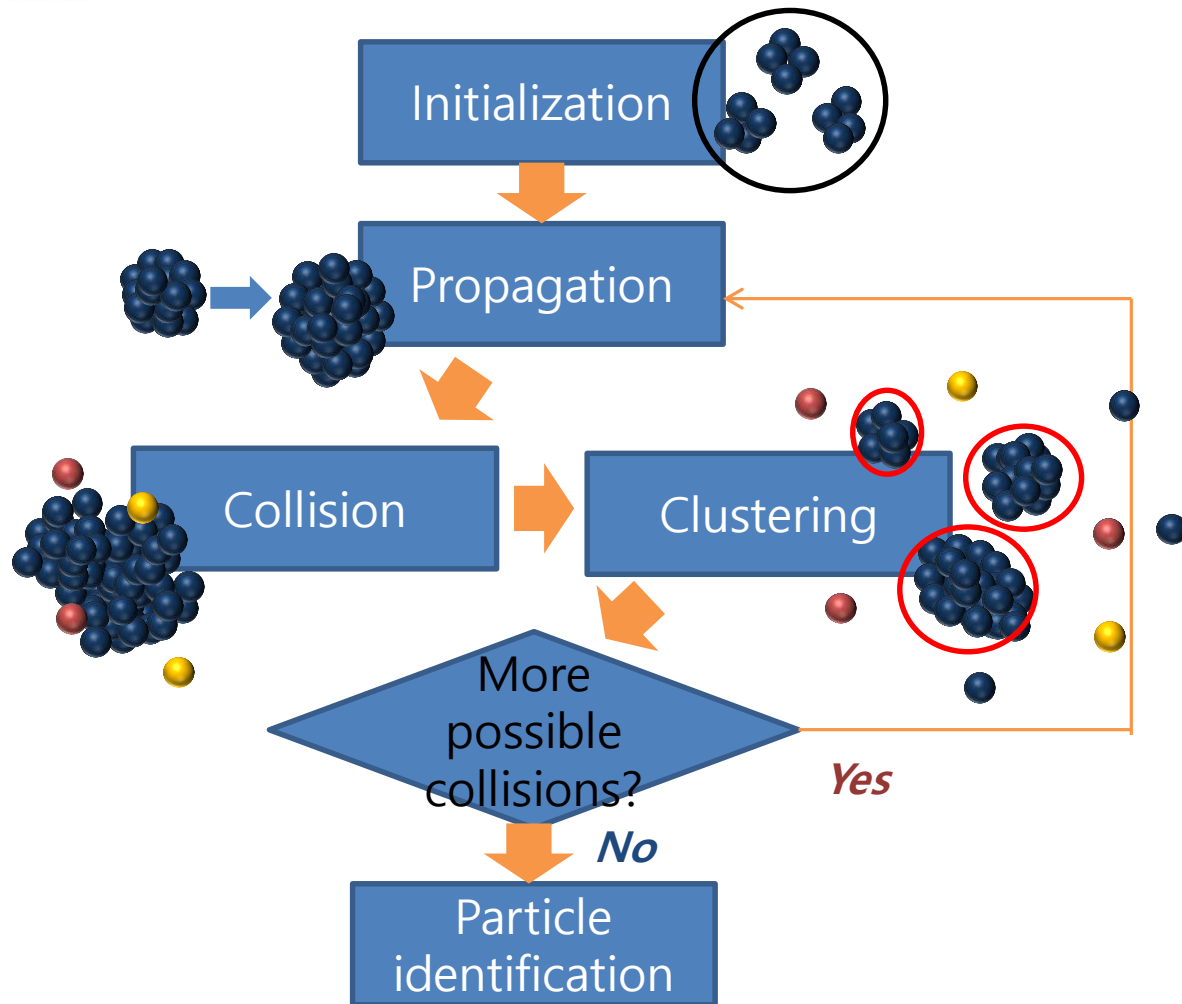
4. Clustering



Identifications of collision fragments are performed by clustering nearby nucleons.



Introduction to Transport Model



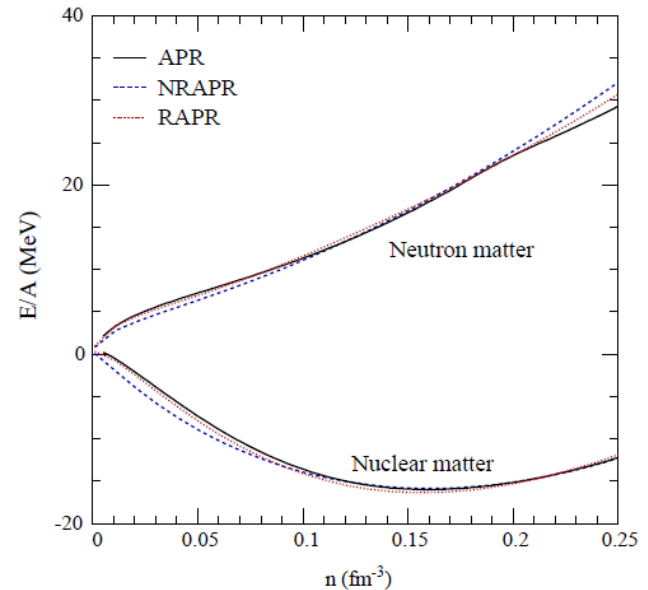
EOS and Symmetry Energy

Equation of State

$$E(\rho, \delta) = E(\rho, \delta = 0) + E_{sym}(\rho)\delta^2 + \mathcal{O}(\delta^4)$$

$$\text{where, } \delta = \frac{\rho_n - \rho_p}{\rho_n + \rho_p}$$

- Astrophysics (super novae, neutron star)
- Giant monopole resonance
- Heavy-ion collisions



Ref.) A.Steiner et al. P.Rept 411 (2005) 325

Incompressibility of symmetry nuclear matter at its saturation density $\rho \approx 0.16 \text{ fm}^{-3}$

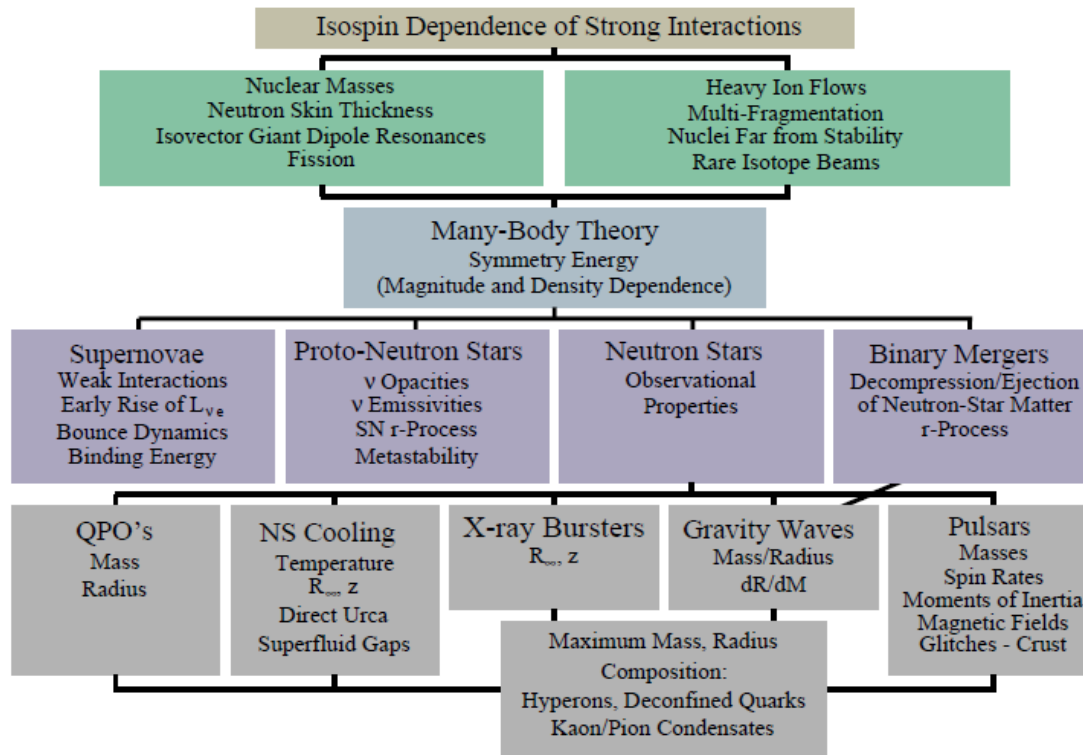
$$K_0 = 231 \pm 5 \text{ [MeV]}$$



from Giant Monopole Resonance

EOS and Symmetry Energy

<The multifaceted influence of the nuclear symmetry energy>



To explore the EOS of isospin asymmetric matter from heavy-ion reactions induced by neutron-rich beams, we need appropriate theoretical tools.

-> Transport Model !!

Ref.) A.Steiner et al. P.Rept 411 (2005) 325

CoMD and ImQMD

Typical QMD + constraints -> less CPU time

Constrained Molecular Dynamics (CoMD)

$$\bar{f}_i \leq 1 \quad (\text{for all } i),$$

$$\bar{f}_i \equiv \sum_j \delta_{\tau_i, \tau_j} \delta_{s_i, s_j} \int_{h^3} f_j(\mathbf{r}, \mathbf{p}) d^3r d^3p.$$



*Restriction
on phase space density*

Improved Quantum Molecular Dynamics (ImQMD)

$$U_{loc} = \frac{\alpha}{2} \sum_i \left\langle \frac{\rho}{\rho_0} \right\rangle_i + \frac{\beta}{3} \sum_i \left\langle \frac{\rho^2}{\rho_0^2} \right\rangle_i + \frac{C_s}{2} \int \frac{(\rho_p - \rho_n)^2}{\rho_0} d^3\mathbf{r} + \int \frac{g_1}{2} (\nabla \rho)^2 d^3\mathbf{r}.$$

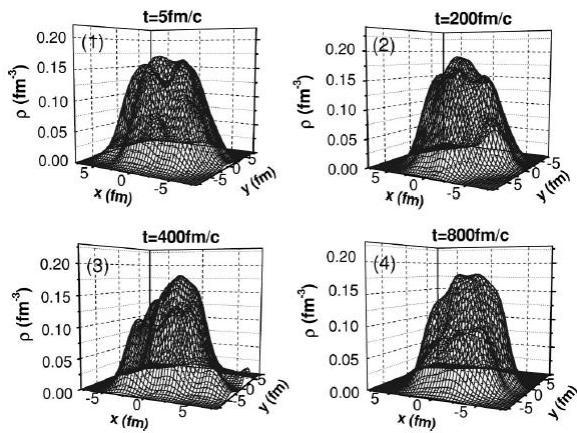
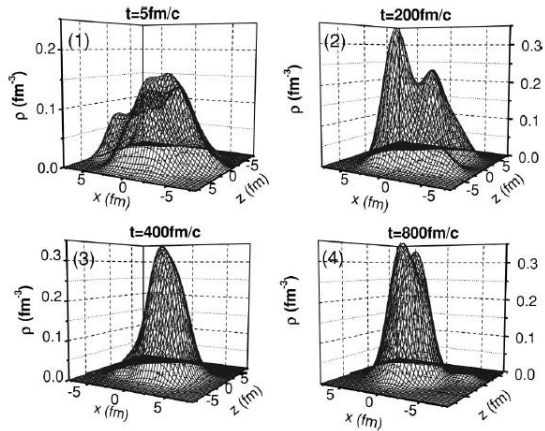
$$\sum_i \left\langle \frac{\rho^2}{\rho_0^2} \right\rangle_i \approx \sum_i \left\langle \frac{\rho}{\rho_0} \right\rangle_i^2 + \int \frac{g_2}{2} (\nabla \rho)^2 d^3\mathbf{r},$$



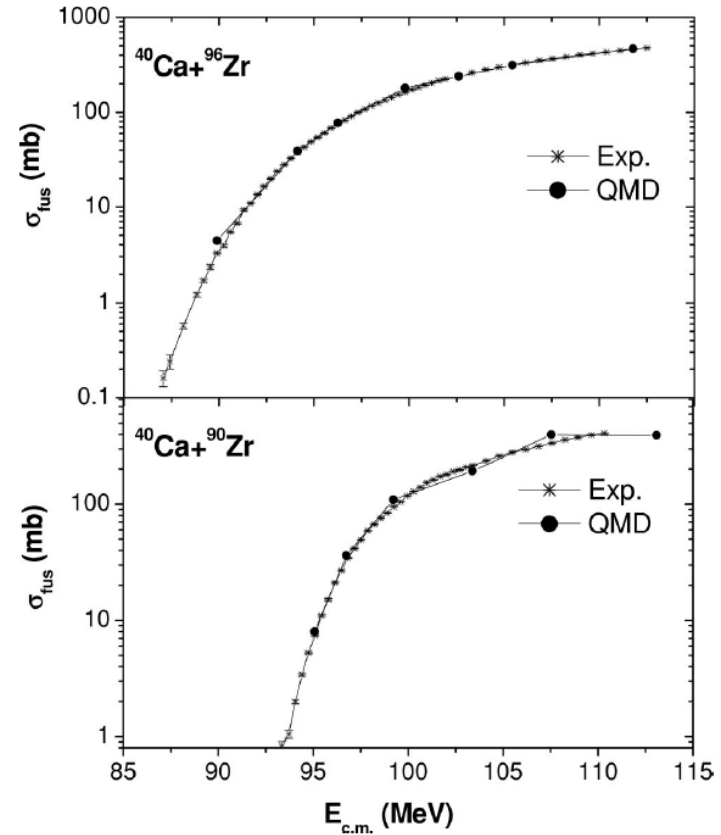
Consideration of surface energy

CoMD and ImQMD

Testing stability of nuclei



Fusion cross-sections



N.Wang – Phys.Rev.C 65 (2002) 064608

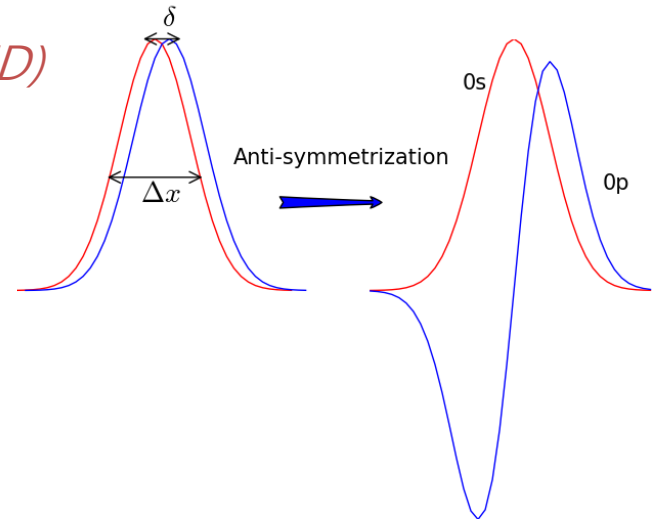
FMD and AMD

Fermionic Molecular Dynamics (FMD)
Antisymmetrized Molecular Dynamics (AMD)

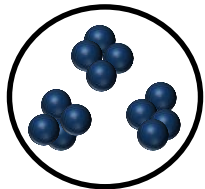
AMD wave function

$$|\Phi(Z)\rangle = \det_{ij} \left[\exp \left\{ -\nu \left(\mathbf{r} - \frac{\mathbf{Z}}{\sqrt{\nu}} \right)^2 \right\} \chi_{\alpha_i}(j) \right]$$

Slater Determinant



Ex.) ^{12}C in AMD



3α structure

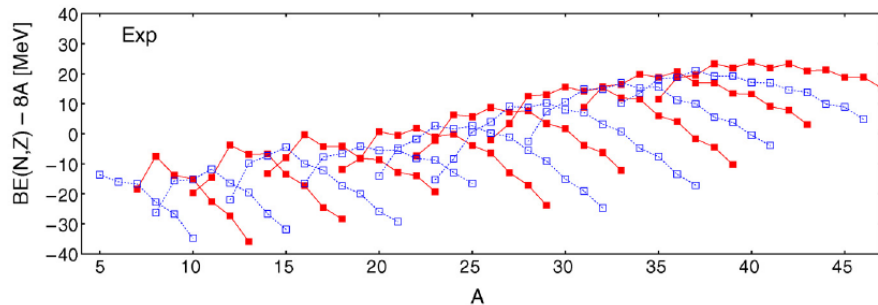
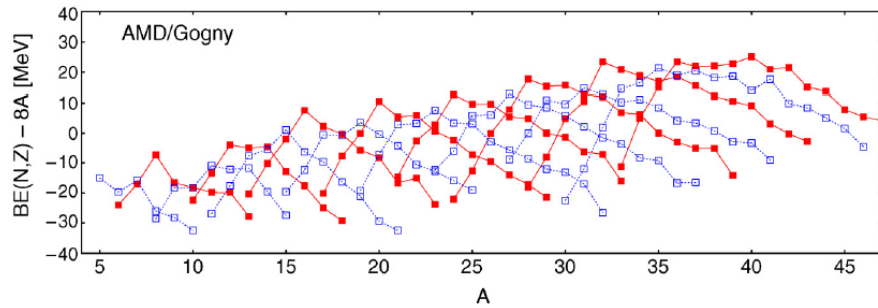
Equation of motion for the wave packet centroids Z

$$\frac{d}{dt} Z_i = \{Z_i, \mathcal{H}\}_{PB} \quad \text{and} \quad i\hbar \sum_{j=1}^A \sum_{\tau=x,y,z} C_{i\sigma,j\tau} \frac{dZ_{j\tau}}{dt} = \frac{\partial \mathcal{H}}{\partial Z_{i\sigma}}$$

(c.f. $C_{i\sigma,j\tau} = \delta_{ij}\delta_{\sigma\tau}$ in QMD)

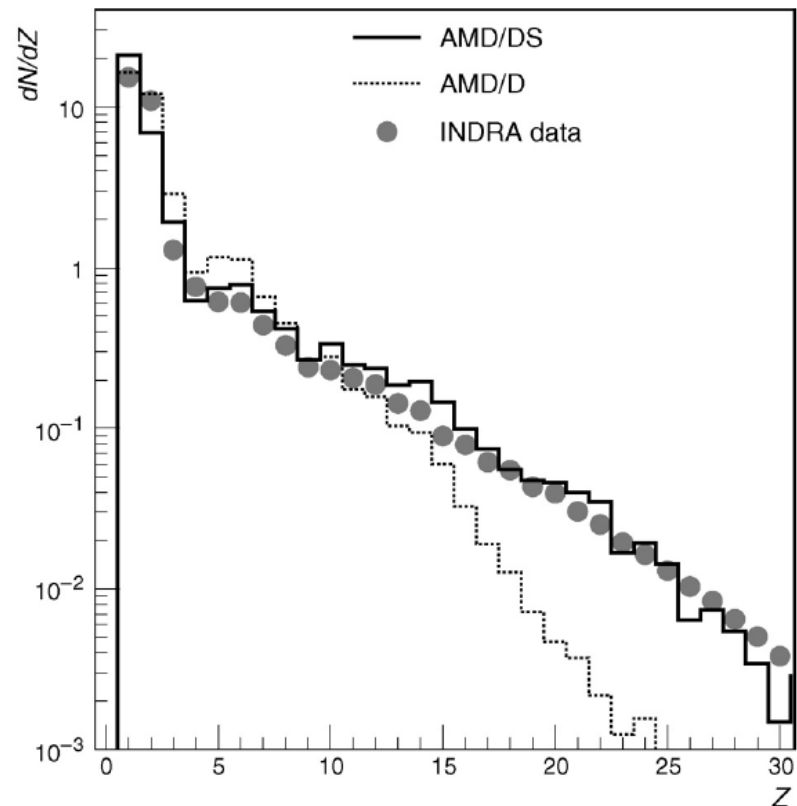
Better fermionic nature -> more CPU time!

FMD and AMD



<Binding energies of nuclei>

The charge distribution of the produced clusters in $^{129}\text{Xe}+\text{Sn}$ at 50 MeV/n



A.Ono - Prog.Part.Nucl.Phys 53 (2004) 501

Transport Model for RAON

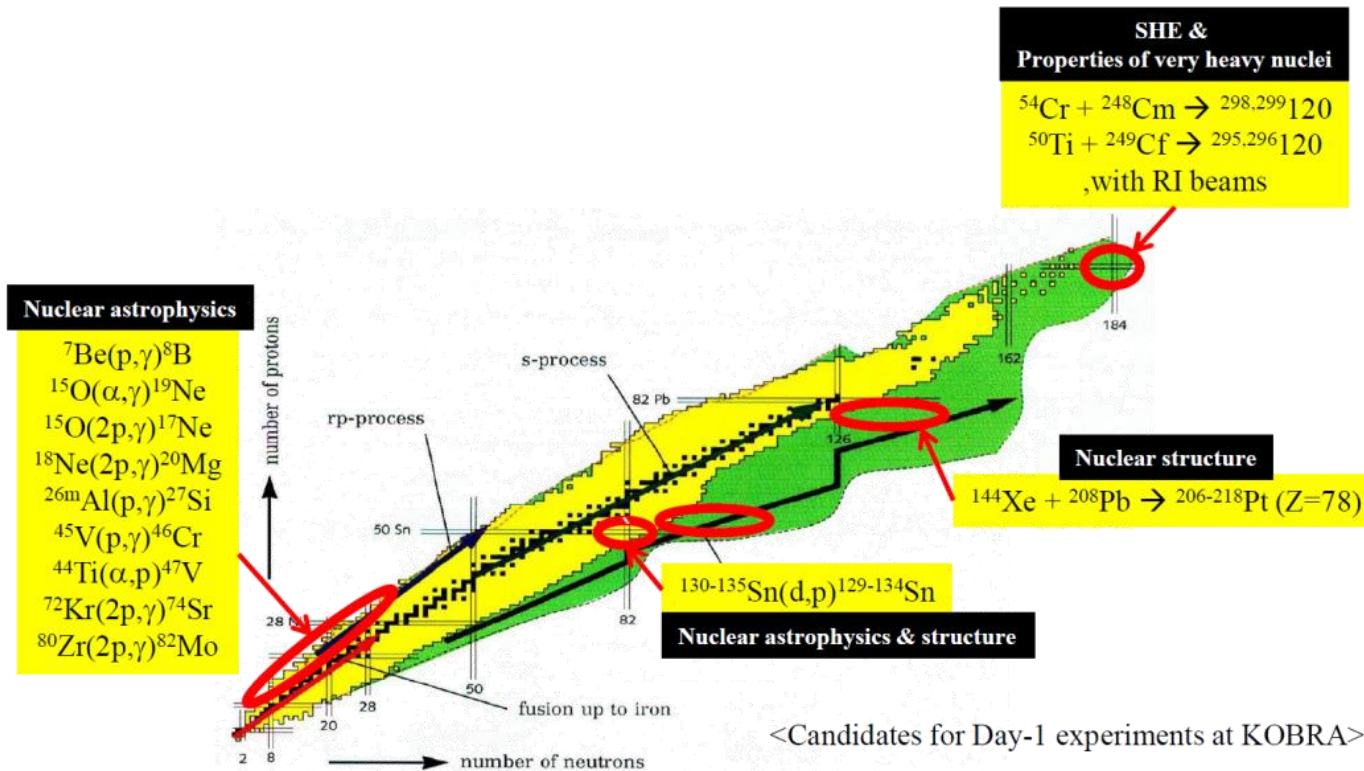
Unstable nuclei



$$r \neq 1.12A^{1/3}$$

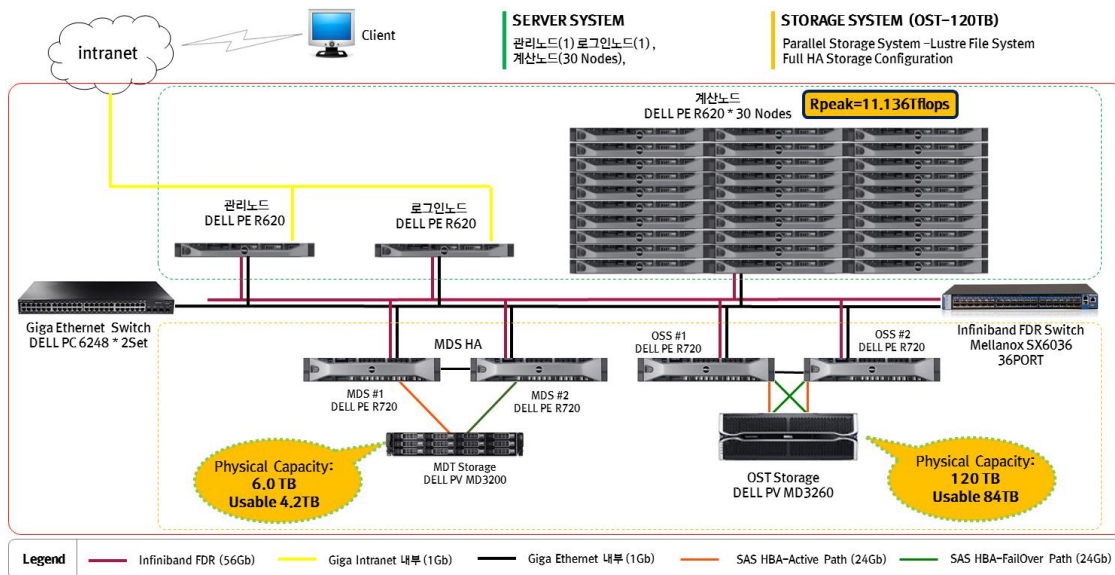
e.g. ^{11}Li is bigger than ^{208}Pb

Nuclear Structure !!



Transport Codes and Super Computer

Cluster at RISP : 480 CPU cores



Tachyon II at KISTI : 25408 CPU cores



Summary and Outlook

- ❖ RAON facility will provide opportunities to study isospin asymmetric matters and exotic nuclei by heavy-ion collisions induced by neutron-rich beams.
- ❖ As a reliable theoretical tools, we need a transport model.
- ❖ Transport model is a model to treat non-equilibrium aspects of the temporal evolution of a collision.
- ❖ Many transport model codes are developed for different energy regions and observables.
- ❖ To simulate HIC at RAON, we need a transport model which describes well low energy reactions and nuclear structures.

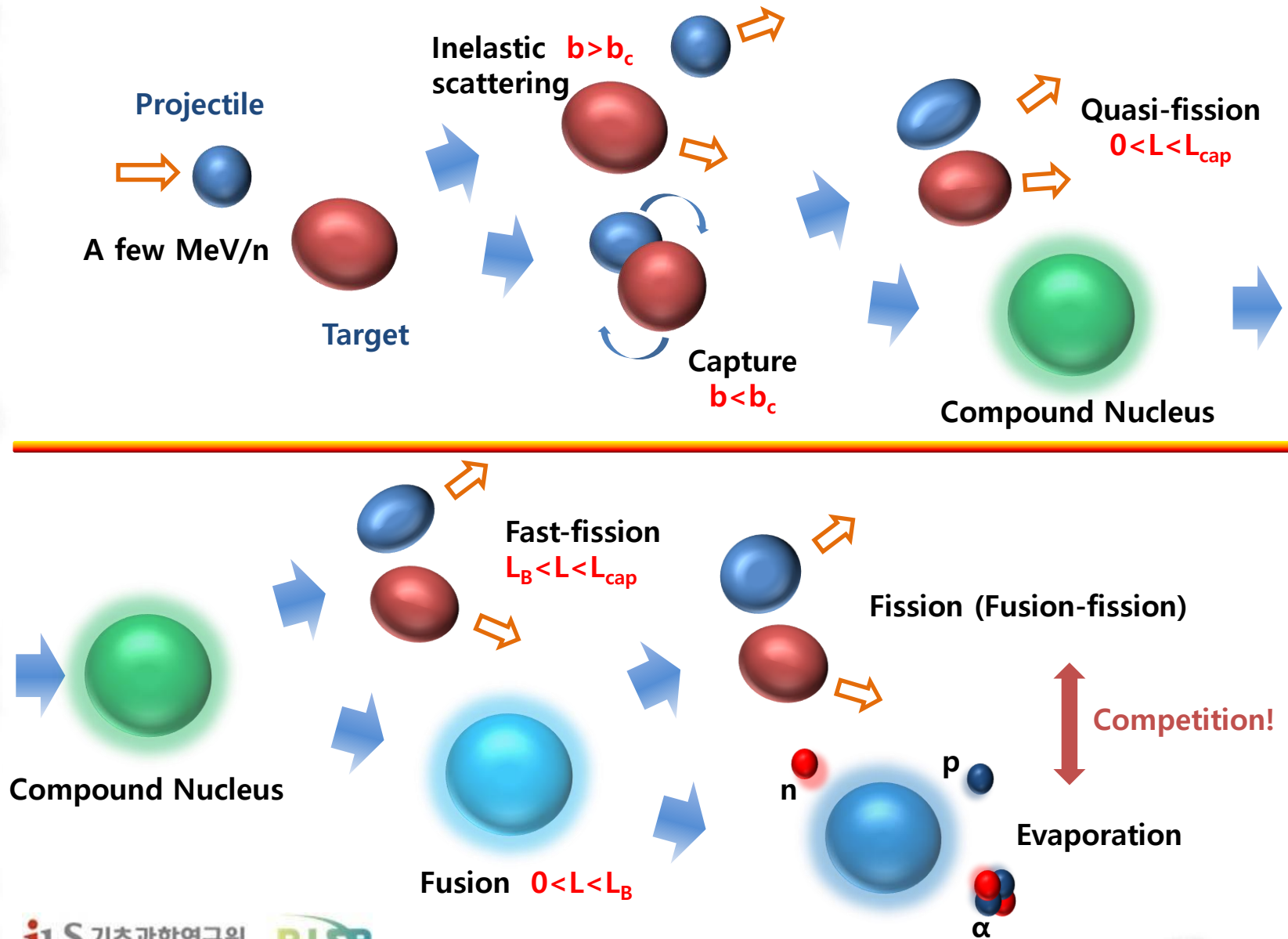
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- ✓ ImQMD
 - N.Wang – Phys.Rev.C 65 (2002) 064608
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Thank you for your attention!!

Backup Slides

<Schematic View of the Fusion Process>



Super Heavy Elements at RAON
