$\Lambda_{c} \text{ Enhancement from} \\ \textbf{Strongly Coupled QGP} \\$

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- Where are diquarks?
- Diquark observation in QGP
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Diquarks

1. Construction of flavor multiplets of baryons Gell-Mann, Ida-Kobayashi, Lichtenberg *et al*.

Diquarks

- 1. Construction of flavor multiplets of baryons
- 2. Exotic structure of hadrons

 σ meson as tetraquark Jaffe

Diquarks

- 1. Construction of flavor multiplets of baryons
- 2. Exotic structure of hadrons
- 3. Dynamical description of diquark in baryons

Ebert, Feldmann, Friedlich, Reinhardt Nagata, Hosaka, Abu-Raddad,

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- 1. Construction of flavor multiplets of baryons
- 2. Exotic structure of hadrons
- 3. Dynamical description of diquark in baryons
- 4. Lattice QCD

Alexandrou, de Forcrand, Lucini

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- 4. Lattice QCD
- 5. Dense matter QCD
 - Color Superconductivity Bose-Einstein condensate

Diquarks

- 1. Construction of flavor multiplets of baryons
- 2. Exotic structure of hadrons
- 3. Dynamical description of diquark in baryons
- 4. Lattice QCD
- 5. Dense matter QCD
- 6. Strongly coupled QGP (sQGP)
 - Color non-singlet bound states Shuryak, Zahed

Diquarks in sQGP

 Perfect fluid behavior
 ☑ Exp. Small viscosity Collective flow
 ☑ Theor. Analysis from AdS/CFT

2. Strong correlations
☑ c̄c bound state at T>> T_c
☑ Variety of bound states
qq, gg, QQ
☑ color non-singlet bound states ??
qq, qqq







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Assumptions a. Bound states of diquarks in sQGP b. Diquark picture in heavy baryons *Qqq*

Heavy baryons (Qqq)

heavy quark (c or b): color spectator



In $m_{c,b} \rightarrow \infty$, short distance forces \square one-gluon exchange \square instanton-induced interaction are suppressed.

> H.J. Lipkin, PL70B, 113 (1977) R.L.Jaffe, PRD72, 074508 (2005)

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Diquarks decouple from heavy quark.

What is an experimental observable for diquarks in sQGP?

Assumptions a. Bound state of diquarks in sQGP b. Diquark picture in heavy baryons *Qqq*

1. Stability of diquark

color-spin interaction

Diquark observation in QGP Diquarks in QGP

- 1. Stability of diquark
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S.Yasui, Chiral07, Osaka

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Diquark (ud) enhances heavy baryon (Λ_c) yield.

Numerical estimate

1. fire-cylinder model p π te das restrictes de 1-01-73 http:// QGP T_c-170 MeV Localdi Perete

fire-cylinder model (thermal distribution T_c~170 MeV)
 quark/diqaurk mass at T_c
 ☑ quark m_u = m_d~ 300 MeV
 ☑ diquark m_{ud} = ???

1. fire-cylinder model (thermal distribution T_c ~170 MeV) 2. quark/diqaurk mass at T_c \square quark $m_u = m_d \sim 300$ MeV \square diquark $m_{ud} = m_u + m_d - \Delta E_{color-spin}$ a. Threshold $\Delta E_{color-spin} \sim 0$ MeV b. Most deeply $\Delta E_{color-spin} \sim 150$ MeV

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fire-cylinder model (thermal distribution T_c~170 MeV)
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 quark m_u = m_d ~ 300 MeV
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 the coalescence model
 hadron yield ∝ (quark thermal distribution in QGP) ×(Wigner function of hadron)

Greco Ko Levai PRL90 (2003)

P)



Numerical estimate

- $\Lambda_{\rm c}$ yield from the coalescence model
 - $N_{\Lambda_{c}(\text{cud})}^{\text{coal}} = g_{\Lambda_{c}(\text{cud})} \int_{\sigma_{C}} \prod_{i=1}^{n=3} \frac{p_{i} \cdot d\sigma_{i} d^{3} \mathbf{p}_{i}}{(2\pi)^{3} E_{i}} f_{q}(x_{i}, p_{i})$ $\times f_{\Lambda_{c}}^{W}(x_{1}..x_{n}; p_{1}..p_{n}), \quad \begin{array}{l} \text{quark thermal distribution} \\ \text{Wigner function of hadron} \end{array}$
- 3-body collision (without diquark)

$$f_{\Lambda_c(\mathrm{cud})}^{\mathrm{W}}(x;p) = 8^2 \exp\left(-\sum_{i=1}^2 \frac{\mathbf{y}_i^2}{\sigma_i^2} - \sum_{i=1}^2 \mathbf{k}_i^2 \sigma_i^2\right)$$

2-body collision (with diquark)

$$f_{\Lambda_{c}(c[ud])}^{W}(x;p) = 8 \exp\left(-\frac{\mathbf{y}^{2}}{\sigma_{c[ud]}^{2}} - \mathbf{k}^{2}\sigma_{c[ud]}^{2}\right)$$

Numerical estimate

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- 3-body collision (without diquark) $N_{\Lambda_{c}(\text{cud})}^{\text{coal}} \simeq g_{\Lambda_{c}(\text{cud})} N_{c} N_{u} N_{d} \prod_{i=1}^{2} \frac{(4\pi\sigma_{i}^{2})^{3/2}}{V_{c}(1+2\mu_{i}T_{C}\sigma_{i}^{2})}$

2-body collision (with diquark)

$$N_{\Lambda_c(c[ud])}^{\text{coal}} \simeq g_{\Lambda_c(c[ud])} N_c N_{[ud]} \frac{(4\pi\sigma_{c[ud]}^2)^{3/2}}{V_c(1+2\mu_{c[ud]}T_C\sigma_{c[ud]}^2)}$$

1. fire-cylinder model (thermal distribution T_c -170 MeV) 2. quark/diqaurk mass at T_c \square quark $m_u = m_d - 300 \text{ MeV}$ \square diquark $m_{ud} = m_u + m_d - \Delta E_{color-spin} - 450-600 \text{ MeV}$ 3. the coalescence model \square hadron yield \propto (quark thermal distribution in QGP) ×(Wigner function of hadron) 4. normalization by Λ_c/D^0 $\square D^0$ is not affected by diquark correlation in QGP.





without diquark









Comparison with another phenomena

<u>Without</u> diquark correlation 1. statistical model : $\Lambda_c/D^0 = 0.09$

 $\exp(-(m_{\Lambda_c} - m_{D^0})/T_{\rm C}) \simeq 0.09$

Comparison with another phenomena

Without diquark correlation

1. statistical model : $\Lambda_{\rm c}/D^0 = 0.09$

2. *pp* collisions : $\Lambda_c/D^0 = 0.159$ (SELEX)



Comparison with another phenomena

<u>Without</u> diquark correlation 1. statistical model : $\Lambda_c/D^0 = 0.09$ 2. *pp* collisions : $\Lambda_c/D^0 = 0.159$ (SELEX) 3. *B* decay : $\Lambda_c/D^0 = 0.03$



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Experiments

1. much abundance of *c* quarks in heavy ion collisions 1.1. estimate by initial hard scattering of nucleons $N_c = 3$ by Au+Au collisions at $s_{\rm NN}^{1/2} = 200$ GeV $N_c = 20$ by Pb+Pb collisions at $s_{\rm NN}^{1/2} = 5.5$ TeV

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1. much abundance of c quarks in heavy ion collisions 1.1. estimate by initial hard scattering of nucleons $N_c = 3$ by Au+Au collisions at $s_{NN}^{1/2} = 200 \text{ GeV}$ $N_c = 20$ by Pb+Pb collisions at $s_{\rm NN}^{1/2} = 5.5$ TeV 1.2. free c quarks in QGP $\overline{c}c$ pairs are resolved as J/ψ suppression 2. enhanced tracking system for charmed hadrons sensitivity to short lifetime ($c\tau$ -60mm) ALICE at LHC STAR and PHENIX at RHIC 48 S.Yasui, Chiral07, Osaka

Summary

- Diquarks in QGP enhance $\Lambda_{c,b}$ yield from heavy ion collisions.
- We propose to measure Λ_c/D^0 in LHC and RHIC.
 - \square New way to observe the existence of QGP.
 - ☑ Experimental approach to study diquark correlation.
 - ☑ Diquark structure in heavy baryons with single heavy quark.