



Hadronic Freedom

**approaching from first
principles**

Mannque Rho, Saclay/Hanyang

Chiral Nuclear Dynamics II

From Quarks to Nuclei
to Compact Stars

Mannque Rho





Weinberg 'folk theorem'

('F-theorem')

“What is quantum field theory, and what did we think it is?”
hep-th/9702027.

“When you use quantum field theory to study low-energy phenomena, then according to the **folk theorem**, you're not really making any assumption that could be wrong, unless of course Lorentz invariance or quantum mechanics or cluster decomposition is wrong, provided you don't say specifically what the Lagrangian is.



'F-theorem' continued

As long as you let it be the most general possible Lagrangian consistent with the symmetries of the theory, you're simply writing down the most general theory you could possibly write down. ... “

“F-proof”: It's hard to see how it can go wrong



Objective of Fundamental Principles in Nuclear Physics

- Recover and sharpen the standard nuclear physics approach, put it in the framework of the Standard Model.
- Make precise predictions that play a key ingredient in other areas of science, e.g., solar evolution and neutrino mass.
- Quest for new states of matter created under extreme conditions



QCD is the First Principle

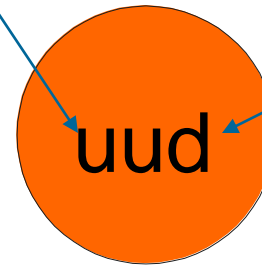
$$\mathcal{L}_{QCD} = \bar{\psi}(i\gamma_{\mu}D^{\mu} - M)\psi - \frac{1}{2}\text{Tr}G_{\mu\nu}G^{\mu\nu}$$

QCD Nucleon

MIT Bag (1970's)

Proton

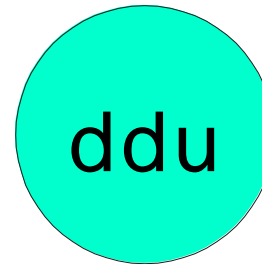
“Up” quark



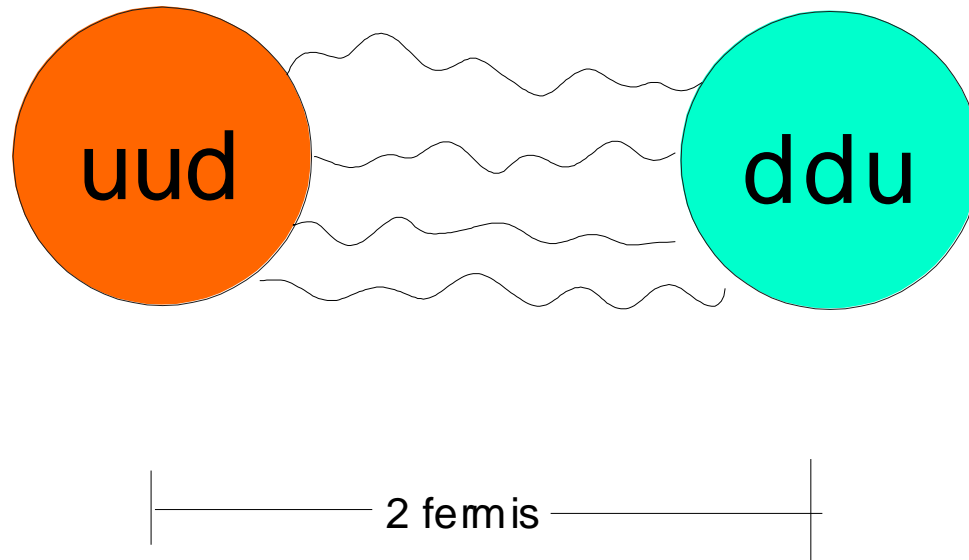
“Down” quark

$R \sim 1 \text{ fm}$

Neutron



DEUTERON



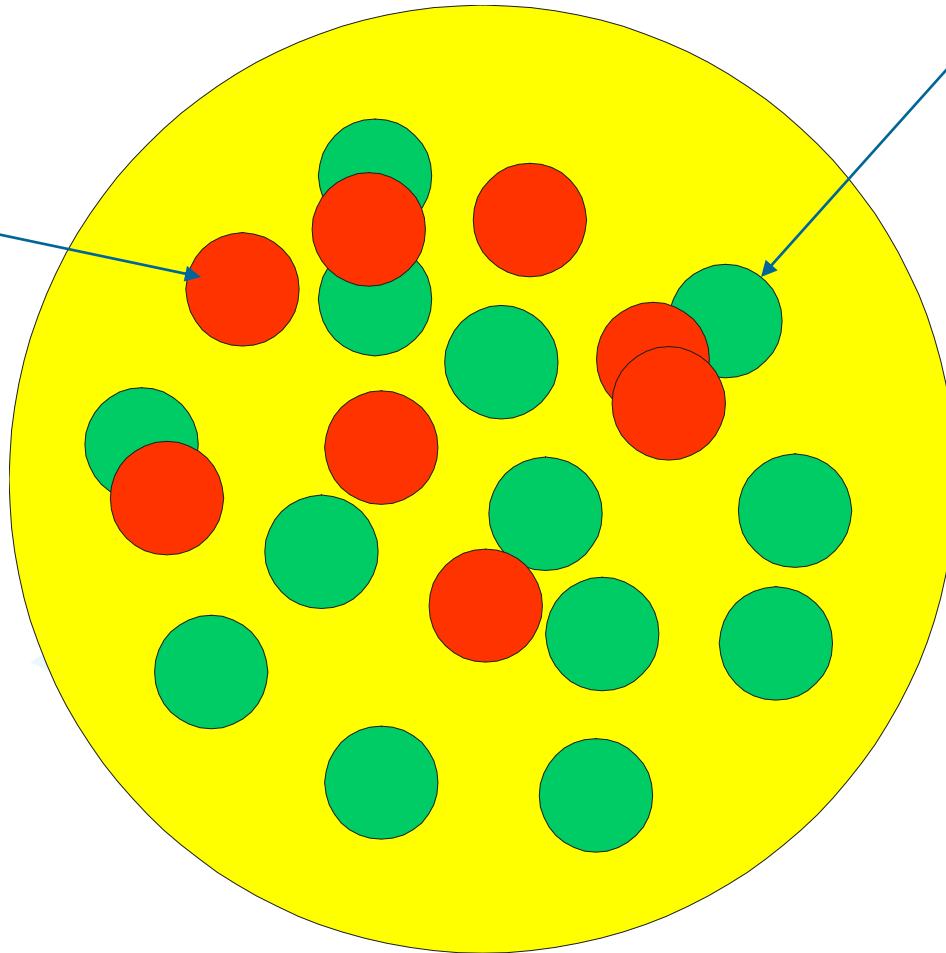
Do the bags of $R \sim 1$ fm overlap?

Heavy Nucleus

Grapefruits in the salad bowl !!!???

PROTON

NEUTRON

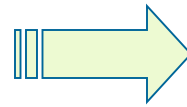


***SIZE
CRISIS?***

Size Problem



MIT bags

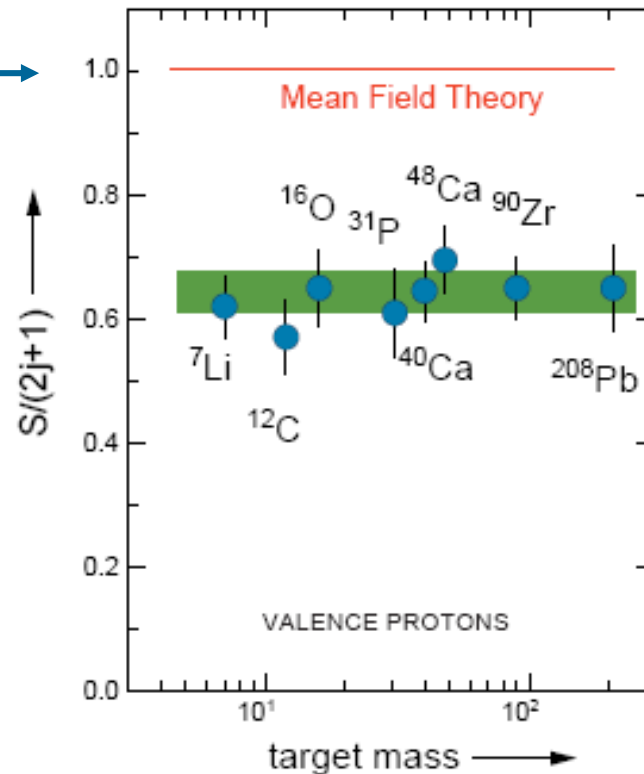


pea soup in ^{208}Pb ?

But
shell model

Spectroscopic
Factor \sim single
particleness

Something
amiss





A Way out

Cheshire cat



“Origin” of the proton mass

Cheshire Cat

*Alice in the
wonderland*

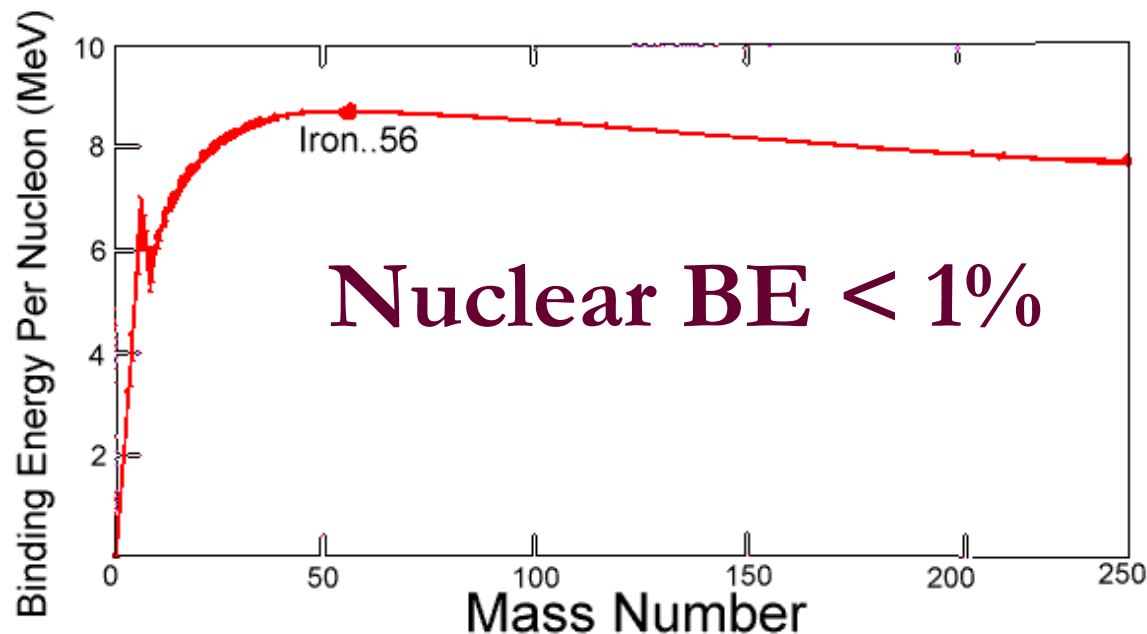


Where does the mass come from?

For Molecules, Atoms, Nuclei

Constituents: protons, neutrons, electrons

Masses = sum of masses of constituents
+ tiny binding energy





A 'Mass' Problem

- Proton/Neutron Mass = 938/940 MeV

Constituents: Quarks and gluons

- Proton = uud ; Neutron = udd

Sum of "current-quark" masses \approx 10 MeV

Where do \sim 99% of the mass come from?

QCD Answer

- QCD on lattice explains the proton mass within $\sim 10\%$.

F. Wilczek

“Energy stored in the motion of the (nearly) massless quarks and energy in massless gluons that connect them”

Proton mass ≈ 1 GeV

“Mass without mass”

- Technically, “chiral symmetry spontaneously broken (χ SB)”

à la Nambu/Goldstone

Order Parameter

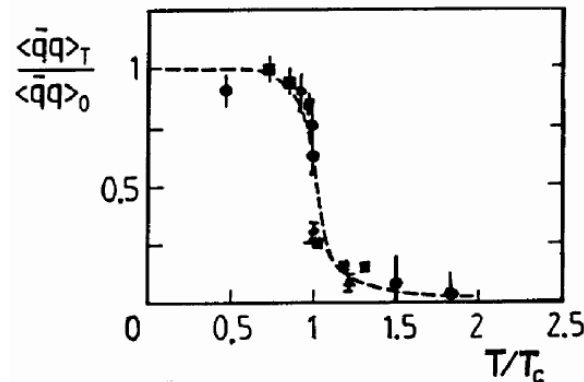
Quark condensate: $\langle \bar{q}q \rangle$

$\neq 0$ χS broken

$= 0$ χS restored

- $\langle \bar{q}q \rangle \approx - (0.23 \pm 0.03 \text{ GeV})^3 \rightarrow$ Proton mass $\approx 1 \text{ GeV}$
- Mass disappears when $\langle \bar{q}q \rangle \rightarrow 0$?

**Lattice
QCD**



Stony Brook "Little Bag"

G.E. Brown and MR 1979

Shrink the bag to $\sim 1/3$ fm from ~ 1 fm

● How?

$\chi_{SB} \rightarrow$ pions as (pseudo)Goldstone bosons



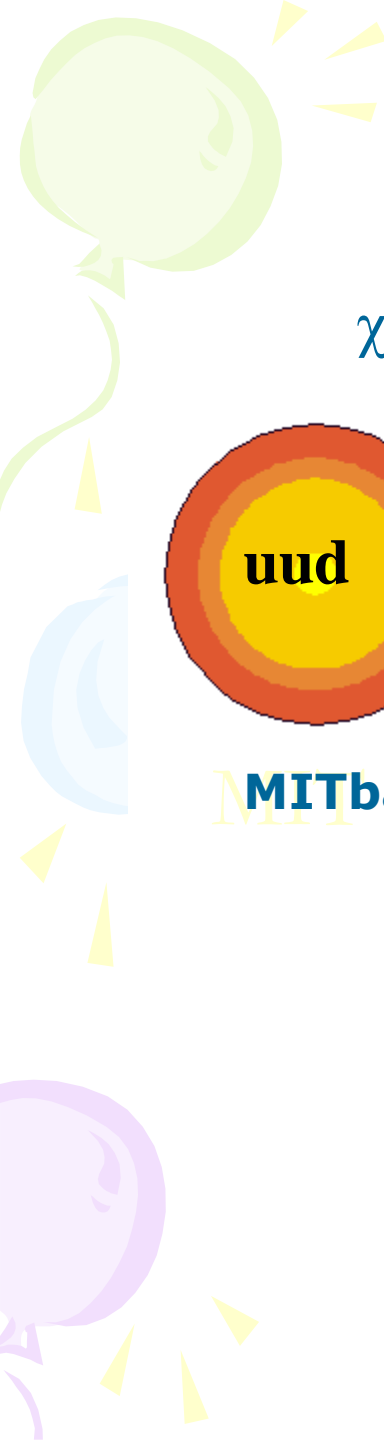
● This reasoning was *not* quite correct!

Enter Cheshire Cat in *Infinite Hotel*

Nadkarni, Nielsen and Zahed 1985

- Bag radius (confinement radius) is a *gauge* (“redundant”) degree of freedom
- ∴ Low-energy physics should not depend upon the bag or confinement size
- R can be shrunk to zero → skyrmion

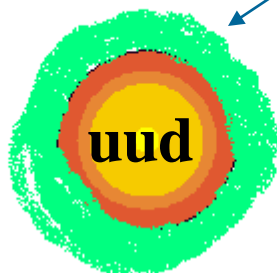
Quarks/gluons \longleftrightarrow **“Smile of the Cheshire Cat”**



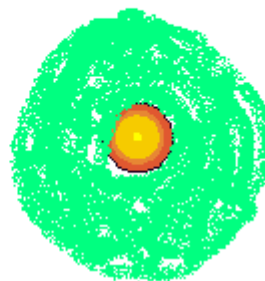
χ SB & anomaly



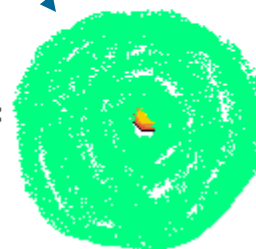
MITbag



"cloudy" bag

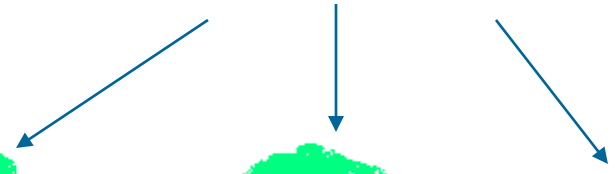


SB little bag

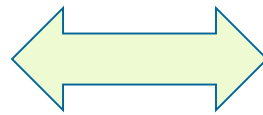


skyrmion

Nambu/Goldstone
(Pion) Cloud



*Equivalent description
of the proton*

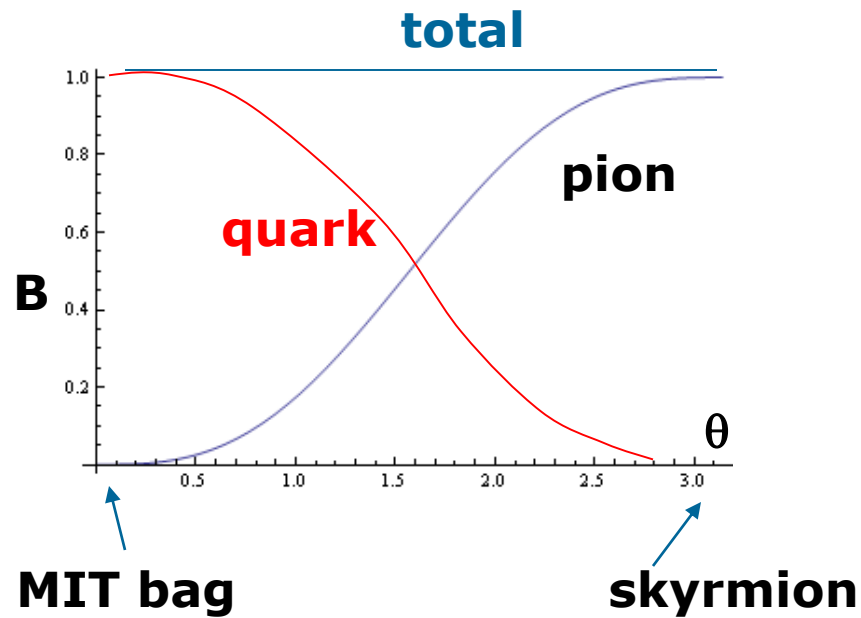


MIT

Stony Brook

Baryon Number

❖ Topological invariant



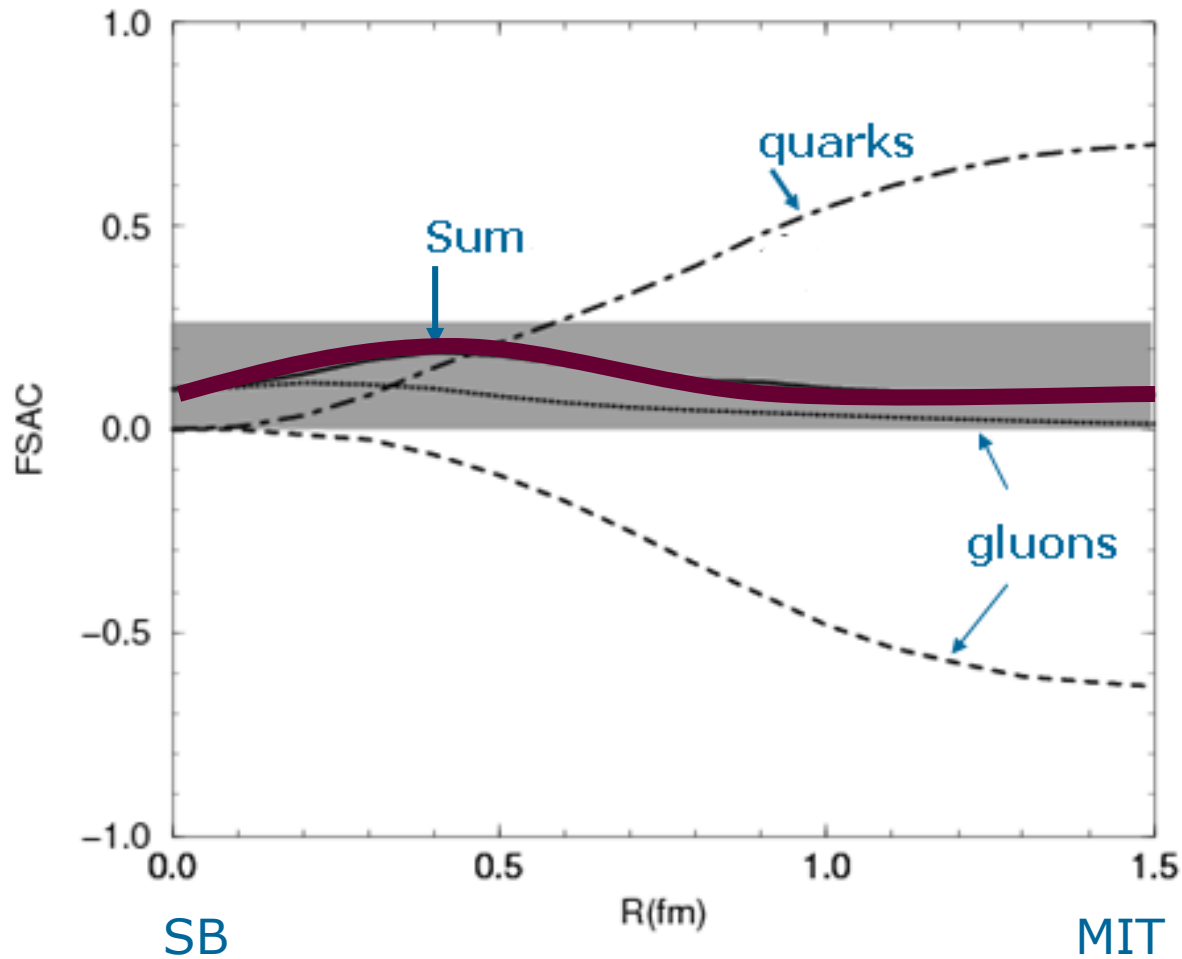
$$L_{\text{QCD}} = \bar{\psi}(i\gamma_{\mu}D^{\mu} - m)\psi - \frac{1}{2}\text{Tr}G_{\mu\nu}G^{\mu\nu}$$

$$L_{\text{EFT}} = \frac{f_{\pi}^2}{4}\text{Tr}(\partial_{\mu}U\partial^{\mu}U^{\dagger}) + \dots$$

$$U = \exp(i\tau \cdot \pi / f_{\pi})$$

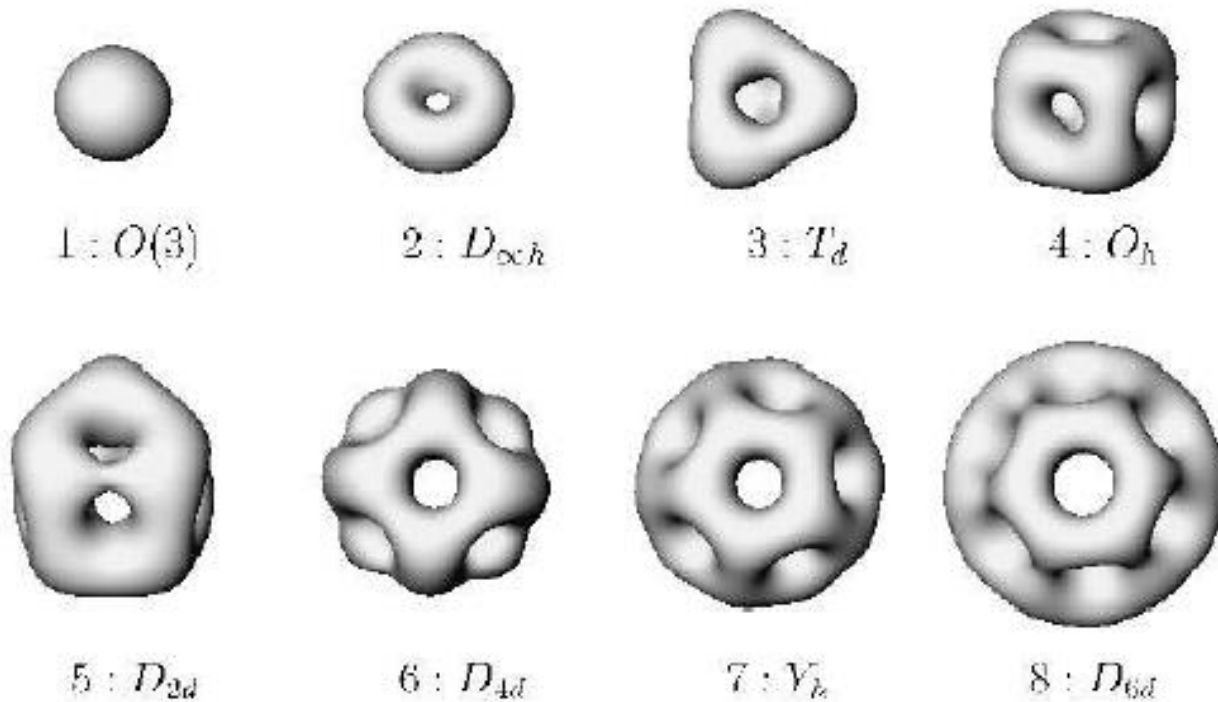
$g_A^0 \sim$ "Proton spin"

Non-topological \sim dynamical



Nuclei as skyrmions

Manton, Sutcliffe et al 2008



Classical, need to be quantized (in progress)

'F-theorem' applied to nuclei

Relevant degrees of freedom: Low-mass hadrons

$$\pi (140), \rho (770), \omega (780), \dots, N (940)$$

• For $E \gg m_\pi (140) \gg m_N (940)$

$$\mathcal{L}_N = N^\dagger (i\partial_t + \nabla^2/2M) N + c(N^\dagger N)^2 + \dots \quad \text{"Pionless Lagrangian"}$$

← Local field
galilean invariance etc.

• For $E \sim m_\pi \sim m_N$

$$\mathcal{L} = \mathcal{L}_N + \mathcal{L}_\pi + \mathcal{L}_{\pi N}$$


$$\mathcal{L}_\pi = (f_\pi^2/4) \text{Tr}(\partial_\mu U \partial^\mu U^\dagger) + \dots$$

$$U = \exp(2i\pi/f_\pi)$$

Chiral invariance, Lorentz invariance ..

Strategy Chiral Lagrangian

- ❖ Pions play a crucial role à la Weinberg
- ❖ Applicable for $E < m_\rho = 770 \text{ MeV}$
- ❖ Match to highly sophisticated 'standard nuclear physics approach' refined since decades:



Weinberg F-corollary “ ... *it allows one to show in a fairly convincing way that what they've been doing all along is the correct first step in a consistent approximation scheme*”

1990 – 2000 : QCD to EFT of nuclei

How does it fare with Nature?

- *Parameter free calculations*

accurate to better than 97%

- ❖ Thermal $n+p \rightarrow d+\gamma$:

$$\sigma^{\text{th}} = 334 \pm 2 \text{ mb} \quad (\text{exp: } 334.2 \pm 0.5 \text{ mb})$$

- ❖ $\mu^- + {}^3\text{He} \rightarrow \nu_\mu + {}^3\text{H}$

$$\Gamma^{\text{th}} = 1499 \pm 16 \text{ Hz} \quad (\text{exp: } 1496 \pm 4 \text{ Hz})$$

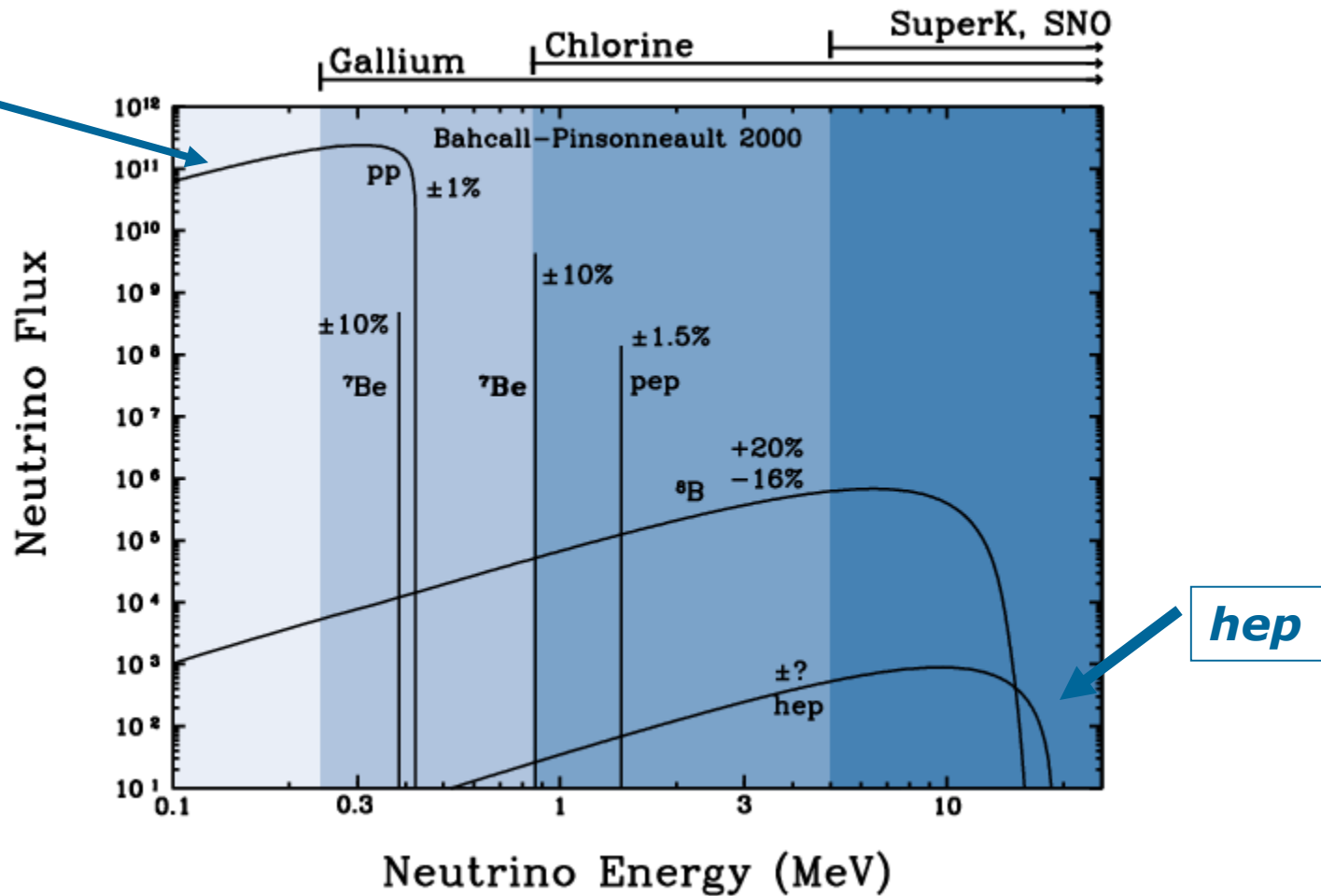
- ❖ $\mu^{\text{th}}({}^3\text{H}) = 3.035 \pm 0.013 \quad (\text{exp: } 2.979 \pm \dots)$

$$\mu^{\text{th}}({}^3\text{He}) = -2.198 \pm 0.013 \quad (\text{exp: } -2.128 \pm \dots)$$

- *Predictions: solar neutrinos*

Solar Neutrino Spectrum

pp



Tortuous History of *hep* Theory

1950-2001

S-factor in 10^{-20} MeV-b unit

| | | |
|-------------------------|------------|---|
| '52 (Salpeter) | 630 | Single particle model |
| '67 (Werntz) | 3.7 | Symmetry group consideration |
| '73 (Werntz) | 8.1 | Better wave functions (P-wave) |
| '83 (Tegner) | 4~25 | D-state & MEC |
| '89 (Wolfs) | 15.3±4.7 | Analogy to $^3\text{He}+n$ |
| '91 (Wervelman) | 57 | $^3\text{He}+n$ with shell-model |
| '91 (Carlson et al.) | 1.3 | VMC with Av14 |
| '92 (Schiavilla et al.) | 1.4-3.1 | VMC with Av28 (N+ Δ) |
| '01 (Marcucci et al.) | 9.64 | CHH with Av18 (N+ Δ) + <i>p</i> -wave |

Serious wave “function overlap” problem



Bahcall's challenge to nuclear physics

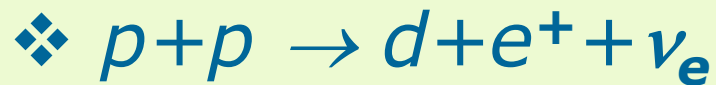
J. Bahcall, hep-ex/0002018

“The most important unsolved problem in theoretical nuclear physics related to solar neutrinos is the range of values allowed by **fundamental physics** for the *hep* production cross section”

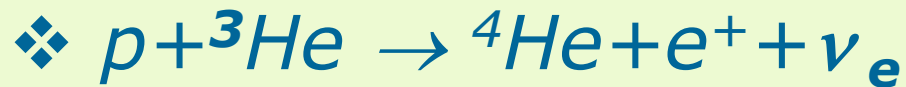
Predictions

T.S. Park et al, 2001

Solar neutrino processes



$$S_{pp} = 3.94 \times (1 \pm 0.0025) \times 10^{-25} \text{ MeV-b}$$



$$S_{hep} = (8.6 \pm 1.3) \times 10^{-20} \text{ keV-b}$$

Awaits experiment!

The background features several large, overlapping, colorful swirls in shades of purple, green, and blue. Scattered throughout are numerous small, yellow, triangular shapes that resemble sparks or particles. The overall aesthetic is dynamic and scientific.

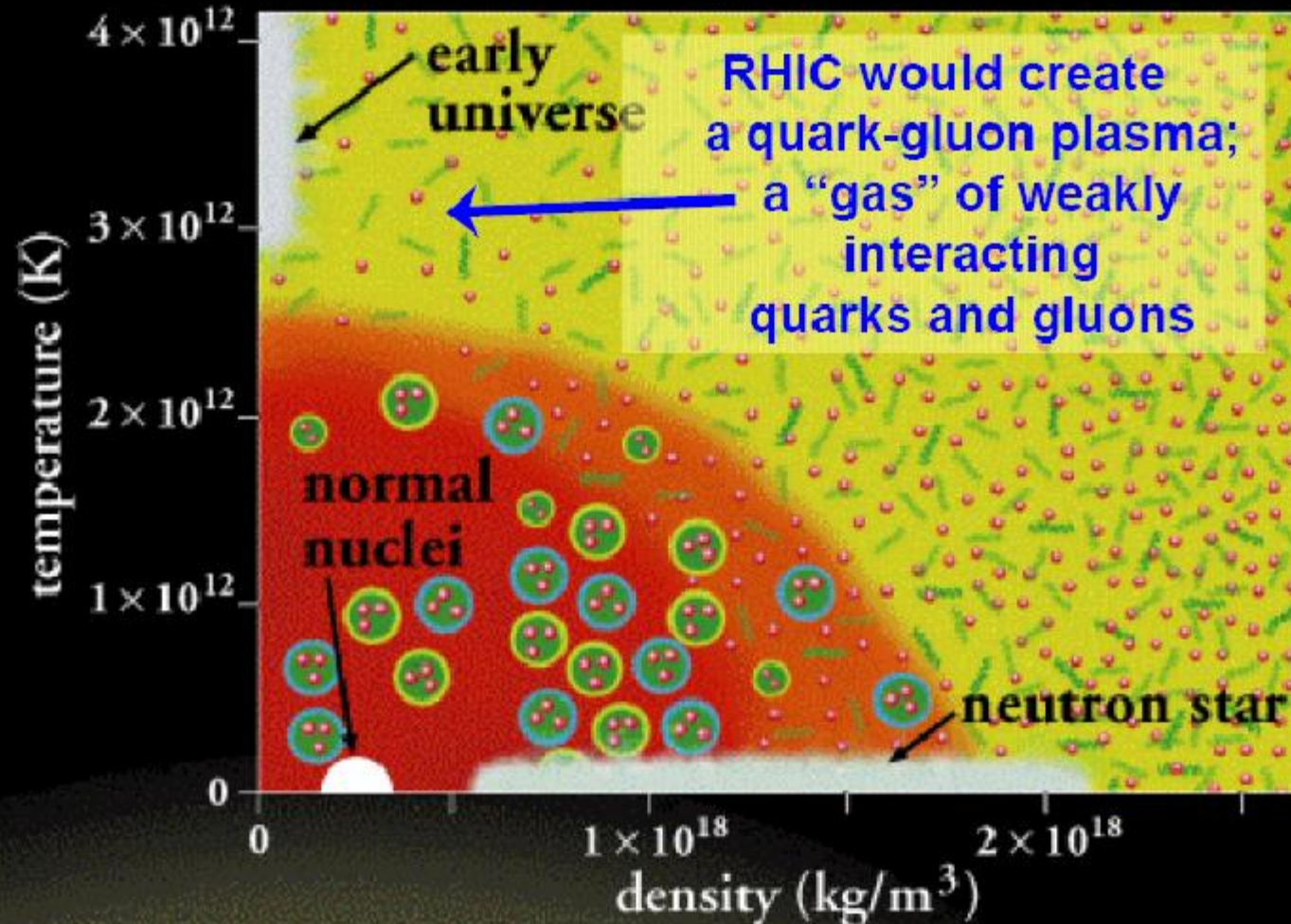
Matter under extreme conditions


**Quest for new states of
matter – New physics**

'Phase diagram'



Expectations circa 2000





What happens as
 $\langle \bar{q}q \rangle \rightarrow 0$?

One possibility is that other
light degrees of freedom than
the pions start figuring

Hidden/emergent gauge symmetries

- ❖ At very low energies, only pions figure

$$\mathcal{L} = (f_\pi^2/4) \text{Tr}[\partial^\mu U \partial_\mu U^\dagger] + \dots \quad \text{“Current algebra”}$$

$$U = \exp(2i\pi/f_\pi) \in SU(N)_L \times SU(N)_R / SU(N)_{V=L+R}$$

Nucleons emerge as skyrmions

- ❖ As energy increases, exploit “gauge symmetry”

Vector mesons $\rho, \rho', \dots, \omega, \omega', \dots$ figure with dropping masses à la Brown-Rho

Nucleons emerge as instantons or skyrions

Gauge symmetry is a redundancy

Famous case: **charge-spin separation** of electron

$e(x) \equiv$ electron, $f(x) \equiv$ "new electron," $b(x) \equiv$ "boson"

$$e(x) = b(x)f^+(x)$$

❖ Invariance: $b(x) \rightarrow e^{ih(x)}b(x)$, $f(x) \rightarrow e^{ih(x)}f(x)$

❖ Endow with a gauge field: $a_\mu \rightarrow a_\mu + \partial_\mu h(x)$
"emergent" gauge field

What we are concerned with

Emerging ρ (770) (and ω)

$$U(x) = e^{2i\pi(x)/f_\pi} = \xi_L^+ \xi_R, \quad \xi_{L/R} = e^{i\sigma(x)/f_\sigma} e^{\mp i\pi(x)/f_\pi}$$

❖ Invariance under $\xi_{L/R} \rightarrow h(x)\xi_{L/R}$ $h(x) \in SU(N)_{L+R}$

❖ "Emergent" SU(N) gauge fields $\rho_\mu \rightarrow h(x)(\rho_\mu + i\partial_\mu)h^+(x)$

Excitation energy $\rightarrow m_\rho \sim 800$ MeV

Bando et al 1986

Harada & Yamawaki 2003



- Emerging “infinite tower” of vectors

$$\rho, \rho', \dots, \omega, \omega', \dots, a_1 \dots$$

$$U(x) = e^{2i\pi/f_\pi} = \Sigma_0 \Sigma_1 \Sigma_2 \bullet \bullet \bullet \Sigma_\infty$$

- ❖ 5-Dimensionally deconstructed QCD (?)(Son & Stephanov 04)

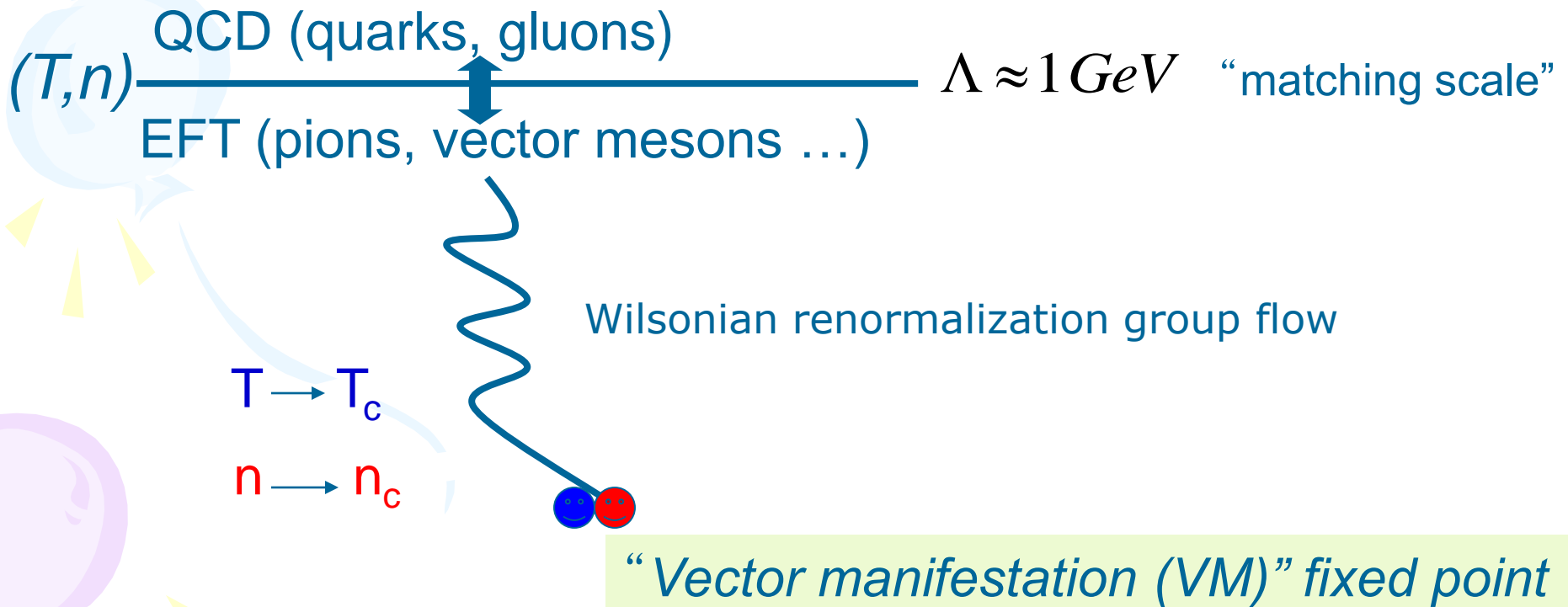
$$S = \int d^4x dz \frac{-1}{2g(z)^2} \sqrt{g} \text{Tr}(F_{AB} F^{AB}) + \bullet \bullet \bullet$$

$$A, B = 0, 1, 2, 3, z$$

- This form descends ALSO from string theory!
- Harada-Yamawaki theory is a truncated HLS theory at the lowest vector mesons ρ, ω .

Matching HLS to QCD

Masayasu Harada &
Koichi Yamawaki
Phys. Rep. 381 (2003) 1-233



Vector Manifestation

In the chiral limit

$$As (T, n) \rightarrow (T_c, n_c)$$

$$m_\rho \sim g \sim m_{\text{const quark}} \propto \langle \bar{q}q \rangle \rightarrow 0$$

$$f_\pi = g = m_\rho = m_\pi = 0 \quad \text{“VM fixed point”}$$

$$a = 1$$

All light-quark hadrons lose mass at the VM point

“VM (or BR) scaling”



VM scaling in nuclei?

Dropping mass tagged to $\langle \bar{q}q \rangle$
Precursor in nuclear structure



- ❖ Warburton ratio
- ❖ carbon-14 dating
- ❖ others

Warburton Ratio ϵ_{MEC}

E. Warburton 91

Warburton defined/measured in nuclei

$$\epsilon_{MEC} \equiv \langle f | A_0 | i \rangle_{\text{exp}} / \langle f | A_0 | i \rangle_{\text{impulseapprox}}$$

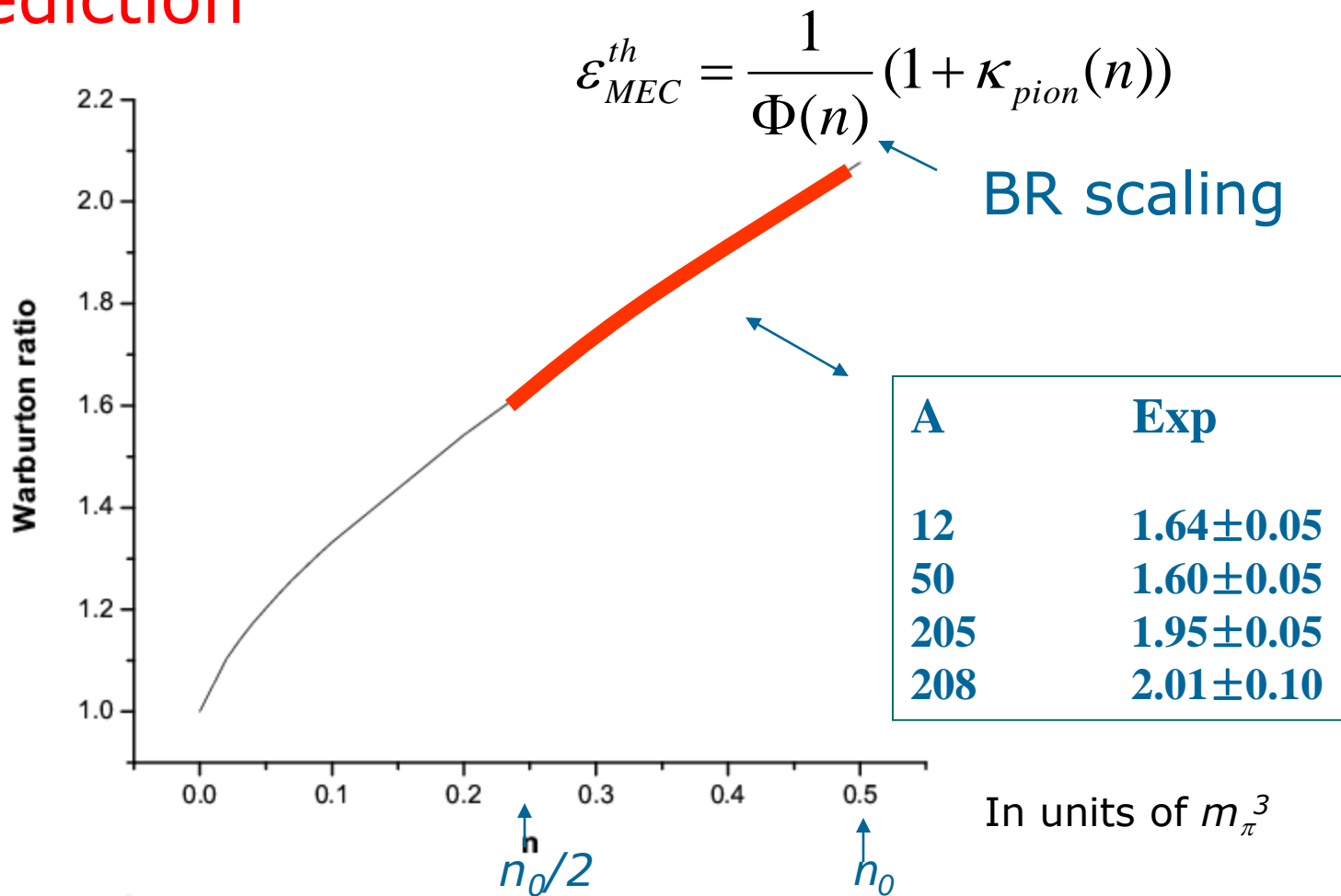
for the weak axial-charge transition

$$A(J^{+/-}) \rightarrow A(J^{-/+}) e \nu \quad \Delta T = 1$$

Found large enhancement in heavy nuclei

$$\epsilon_{MEC} = 1.9 \leftrightarrow 2.1$$

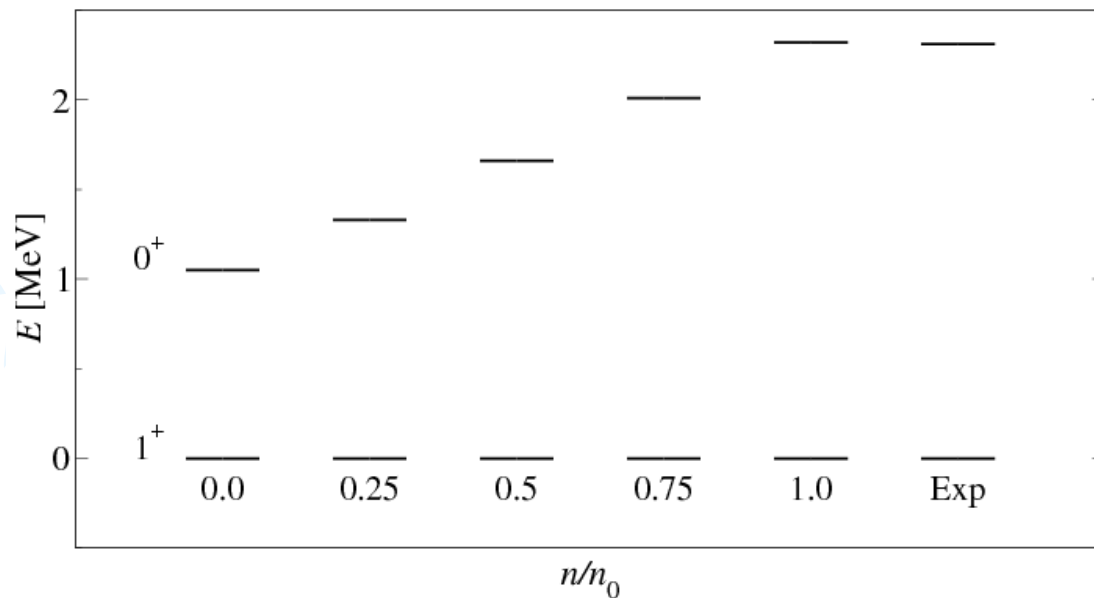
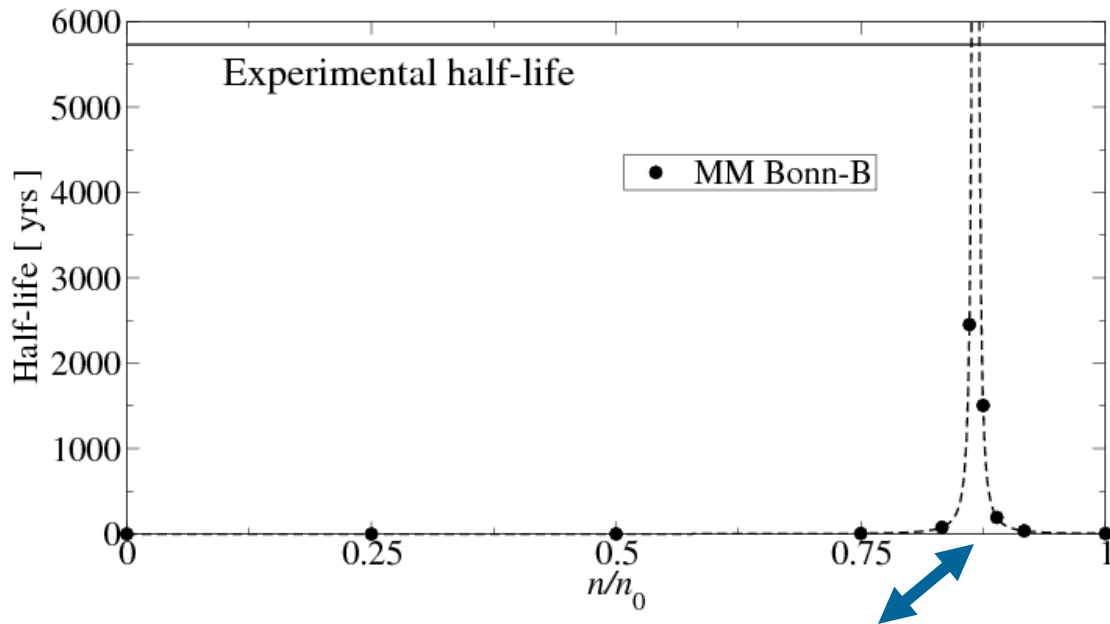
Prediction



Carbon-14 dating

Tensor force
fine-tuned by
BR scaling!

Holt et al 2008



A decorative graphic on the left side of the slide features three balloons: a green one at the top, a light blue one in the middle, and a purple one at the bottom. Each balloon is attached to a streamer and has several small yellow triangular shapes around it, resembling confetti or streamer ends.

**Hadronic matter at high
temperature and/or density**



Large efforts in heavy-ion collisions
at CERN and RHIC and in the space

No smoking gun signal yet

But there are two neat predictions from
VM!! “Hadronic Freedom”



“Hadronic Freedom”

VM implies that near the phase transition (PT) approaching from below, hadronic interactions become very weak

Assume between the PT point and the “flash point” (at which hadrons become strongly interacting), hadrons flow “freely” with little interaction



Brown Rule (after Bethe): ***Set equal to zero!***

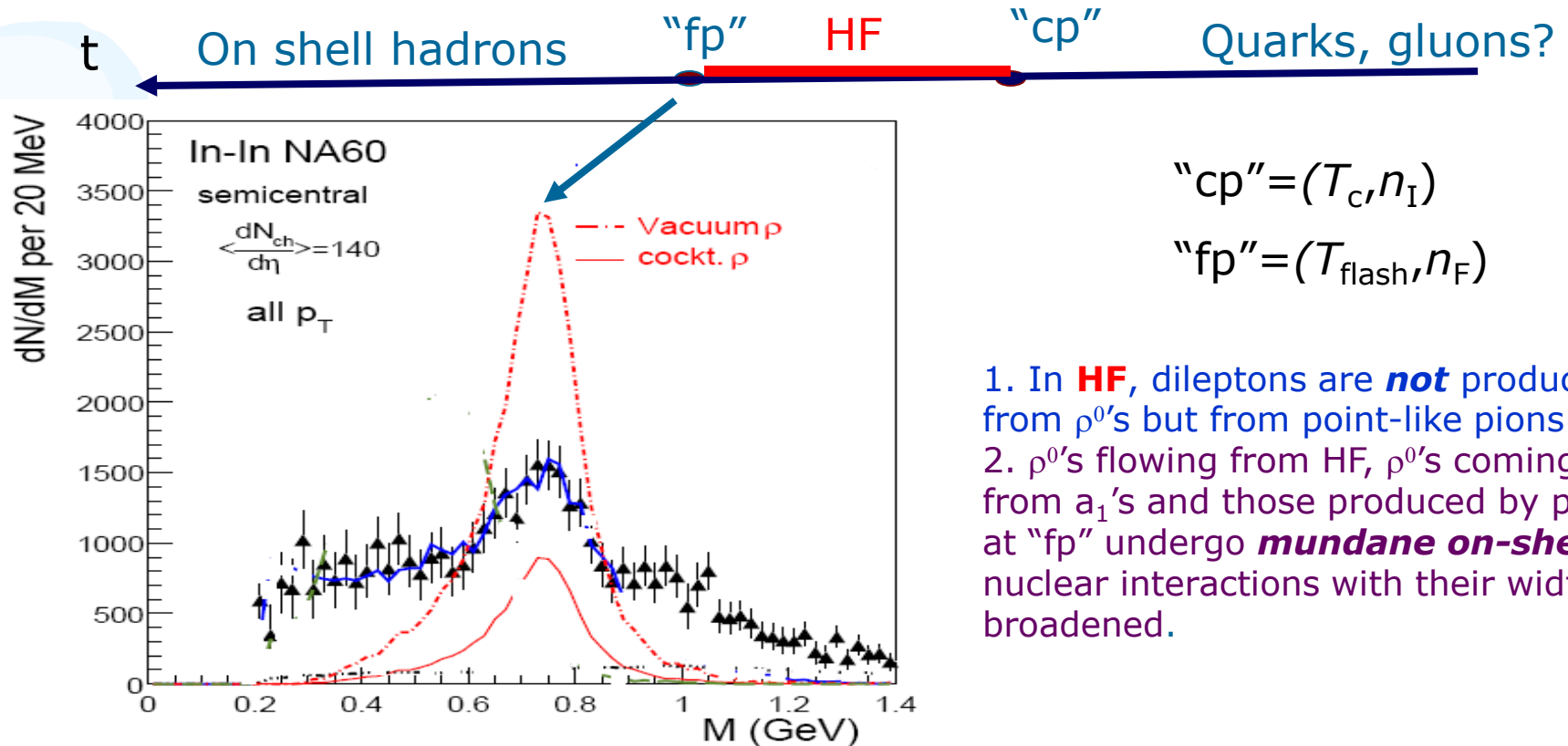
Predictions

1. Gives simple explanation of dilepton productions from heavy-ion collisions: “triviality”
2. Sets maximum stable neutron star mass M_{max}^{BB}

Dileptons

Dileptons are blind to the vector mesons in HF

Brown, Holt, Harada, Rho and Sasaki, arXiv:0901.1513



$$\text{"cp"} = (T_c, n_I)$$

$$\text{"fp"} = (T_{\text{flash}}, n_F)$$

1. In **HF**, dileptons are **not** produced from ρ^0 's but from point-like pions
2. ρ^0 's flowing from HF, ρ^0 's coming from a_1 's and those produced by pi-pi at "fp" undergo **mundane on-shell** nuclear interactions with their widths broadened.



How does one see VM (or BR scaling or precursor to chiral restoration)?

Mesure direct $\pi^+\pi^-\rightarrow$ dileptons

Subtract all the cocktails that include the on-shell broadened ρ^0 's \rightarrow flat distribution coming from HF!!

Will check HF and VM/BR



High Density Regime

Compact stars and Black Holes

Questions:

- ❖ What happens as density increases to that of **compact stars**?
- ❖ Does **hadronic physics** matter for the collapse of stars?
- ❖ Are the plethora of **high density matter** observable?

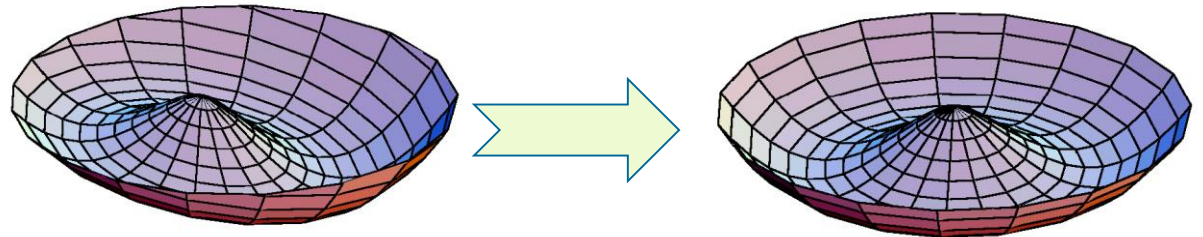
Assertion:

- ❖ The first – and possibly last (?) – phase change is that **kaons condense** at relatively **low density near the “flash density”**

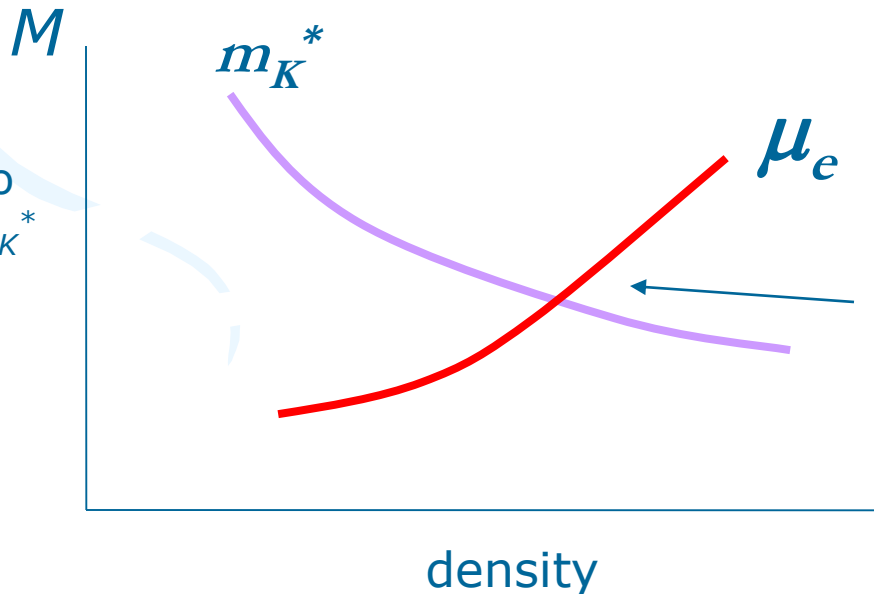
Kaons condense in compact stars

$$m_\pi \sim 0, m_K \sim 1/2 \text{ GeV}$$

Dropping mass
"restores" $SU(3)$
symmetry



HF allows to
calculate m_K^*



$$n_c^K \approx 3n_0 < n_{\langle \bar{q}q \rangle} \rightarrow 0$$

Kaons condense

Consequences

A scenario proposed

PRL 101, 091101 (2008)

PHYSICAL REVIEW LETTERS

week ending
29 AUGUST 2008

Kaon Condensation, Black Holes, and Cosmological Natural Selection

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²*Department of Physics, Pusan National University, Busan 609-735, Korea*

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(Received 29 February 2008; published 28 August 2008)

It is argued that a well-measured double neutron-star binary in which the two neutron stars are more than 4% different from each other in mass or a massive neutron star with mass $M \lesssim 2M_{\odot}$ would put in serious doubt or simply falsify the following chain of predictions: (1) a nearly vanishing vector meson mass at chiral restoration, (2) kaon condensation at a density $n \sim 3n_0$, (3) the Brown-Bethe maximum neutron-star mass $M_{\max} \approx 1.5M_{\odot}$, and (4) Smolin's "cosmological natural selection" hypothesis.

DOI: [10.1103/PhysRevLett.101.091101](https://doi.org/10.1103/PhysRevLett.101.091101)

PACS numbers: 97.60.Jd, 97.60.Lf, 98.80.Bp, 98.80.Qc

- i. A lot of light-mass black holes in the Universe
- ii. "BH-Nothingness" after kaon condensation

Bethe-Brown Mass

“Stars more massive than $M_{max}^{BB} \approx 1.6 M_{\odot}$ collapse into black holes”

Why? Because such massive stars have condensed kaons which soften the EOS and trigger instability.

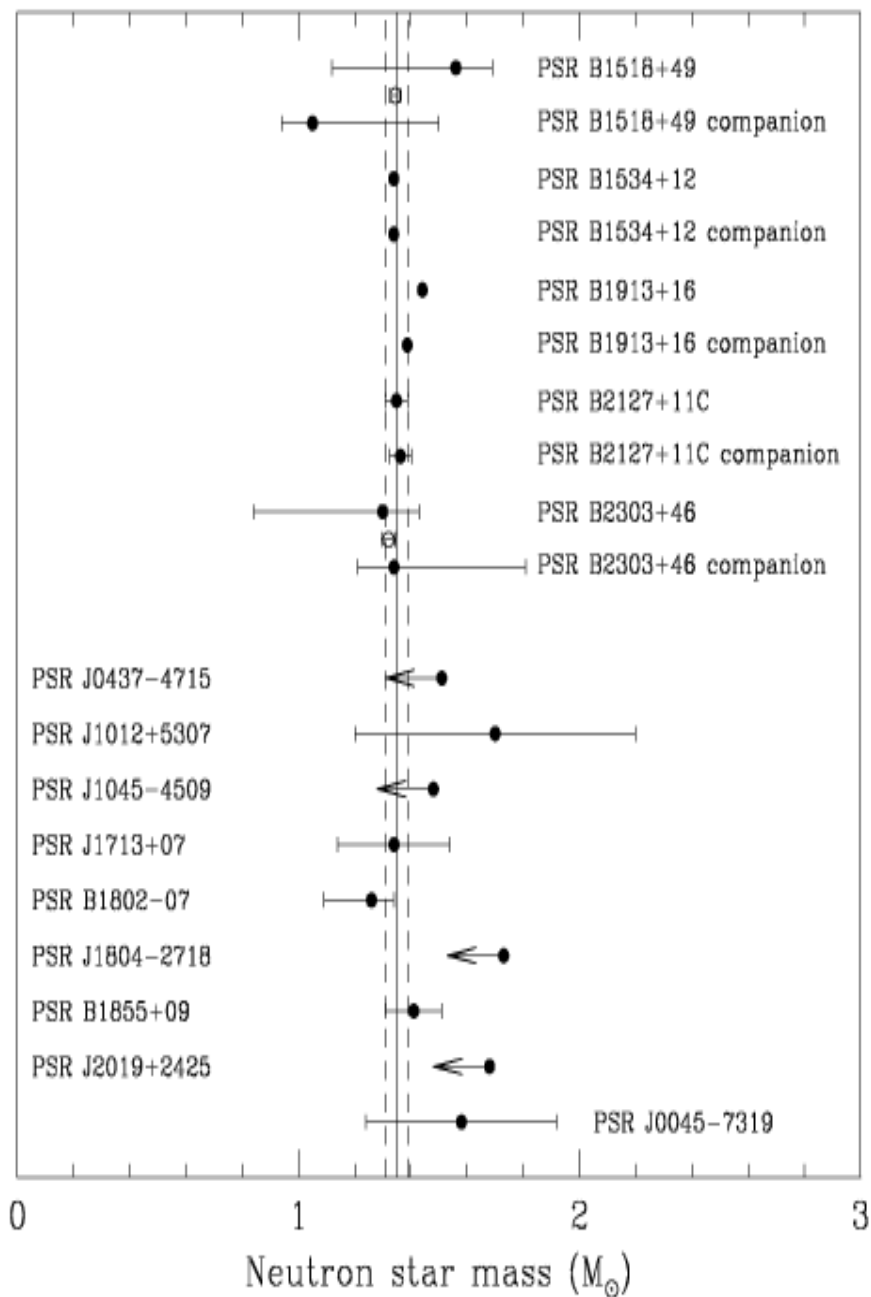
“No proof. It’s a conjecture to be checked by nature .”

What to do?

- a) “Find a compact star with mass $M > M_{max}^{BB}$ ”
- b) “Find binary pulsars with mass difference $> 4\%$ ”

If found, the following will be invalidated

- a) Maximization of black holes in the Universe
- b) Mechanism for “Cosmological Natural Selection”
- c) Kaon condensation, VM, “hadronic freedom”



X-ray Binaries

| | | | |
|-----------|--|---------------|------------------------|
| 4U1700-37 | $2.44^{+0.27}_{-0.27}$ | Vela X-1 | $1.86^{+0.16}_{-0.16}$ |
| Cyg X-1 | $1.78^{+0.23}_{-0.23}$ | 4U1538-52 | $0.96^{+0.19}_{-0.16}$ |
| SMC X-1 | $1.17^{+0.16}_{-0.16}$, 1.05 ± 0.09 | XTE J2123-058 | $1.53^{+0.30}_{-0.42}$ |
| LMC X-4 | $1.47^{+0.22}_{-0.19}$, 1.31 ± 0.14 | Her X-1 | $1.47^{+0.12}_{-0.18}$ |
| Cen X-3 | $1.09^{+0.30}_{-0.26}$, 1.24 ± 0.24 | 2A 1822-371 | > 0.73 |

Neutron Star - Neutron Star Binaries

| | | | |
|-------------|------------------------------|----------------------|------------------------------|
| 1518+49 | $1.56^{+0.13}_{-0.44}$ | 1518+49 companion | $1.05^{+0.45}_{-0.11}$ |
| 1534+12 | $1.3332^{+0.0010}_{-0.0010}$ | 1534+12 companion | $1.3452^{+0.0010}_{-0.0010}$ |
| 1913+16 | $1.4408^{+0.0003}_{-0.0003}$ | 1913+16 companion | $1.3873^{+0.0003}_{-0.0003}$ |
| 2127+11C | $1.349^{+0.040}_{-0.040}$ | 2127+11C companion | $1.363^{+0.040}_{-0.040}$ |
| J0737-3039A | $1.337^{+0.005}_{-0.005}$ | J0737-3039B | $1.250^{+0.005}_{-0.005}$ |
| J1756-2251 | $1.40^{+0.02}_{-0.03}$ | J1756-2251 companion | $1.18^{+0.03}_{-0.02}$ |

Neutron Star - White Dwarf Binaries

| | | | |
|------------|--------------------------|------------|------------------------|
| B2303+46 | $1.38^{+0.06}_{-0.10}$ | J1012+5307 | $1.68^{+0.22}_{-0.22}$ |
| J1713+0747 | $1.54^{+0.007}_{-0.008}$ | B1802-07 | $1.26^{+0.08}_{-0.17}$ |
| B1855+09 | $1.57^{+0.12}_{-0.11}$ | J0621+1002 | $1.70^{+0.32}_{-0.29}$ |
| J0751+1807 | $2.20^{+0.20}_{-0.20}$ | J0437-4715 | $1.58^{+0.18}_{-0.18}$ |
| J1141-6545 | $1.30^{+0.02}_{-0.02}$ | J1045-4509 | < 1.48 |
| J1804-2718 | < 1.70 | J2019+2425 | < 1.51 |

Neutron Star - Main Sequence Binaries

| | |
|------------|------------------------|
| J0045-7319 | $1.58^{+0.34}_{-0.34}$ |
|------------|------------------------|

J0751+1807

Nice et al 2005

Observation in neutron star–white dwarf binary of $2.2 \pm 0.2 m_{\oplus}$ led to *pitched activities*

- ❖ strong repulsive N-nucleon forces (with $N \geq 3$)
- ❖ crystalline color-superconducting stars
- ❖ etc etc producing \sim one paper a week

This would unambiguously “kill” the BB conjecture and also VM

But (!) new analysis in 2007 corrects the 2005 value to $1.26 + 0.14 / - 0.12$!!

BB still OK!

Summary

- We went to **skyrmions** from quarks
- We went to **nuclei** via skyrmions via F-theorem
- We went to **HF** to **compact stars** via nuclear matter via hidden local symmetry
- Enter **string theory**:

Sakai and Sugimoto showed (2005) that hadrons at low energy $E < M_{KK}$ could be described by the 5D action **top-down** from AdS/CFT:

$$S = -\int d^4x dz \frac{1}{4e^2(z)} \text{Tr}[F_{AB}F^{AB}] + \dots + S_{CS}$$

Arises also **bottom-up** from current algebra by “deconstruction”



Thanks for the attention!