

# Dynamical Recombination model of QGP

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- Introduction – recombination model
- Dynamic recombination calculation :  
Hydrodynamic evolution of QGP +  
+ Hadronization via recombination  
+ Hadronic rescattering by URQMD
- Results
- Summary and Outlook

collaborators:

- S. Bass
- B. Müller
- C. Nonaka

# Recombination : long history

## High Energy Physics Phenomenology:

- K.P. Das & R.C. Hwa, Phys. Lett. B68, 459 (1977)

*Quark-Antiquark Recombination in the Fragmentation Region*

➤ description of leading particle effect

- T. Ochiai, Prog. Theo. Phys. 75, 1184 (1986)
- E. Braaten, Y. Jia & T. Mehen, Phys. Rev. Lett. 89, 122002 (2002)
- R. Rapp & E.V. Shuryak, Phys. Rev. D67, 074036 (2003)

## Heavy-Ion Phenomenology:

- T. S. Biro, P. Levai & J. Zimanyi, Phys. Lett. B347, 6 (1995)

*ALCOR: a dynamical model for hadronization*

➤ yields and ratios via counting of constituent quarks

- R.C. Hwa & C.B. Yang, Phys. Rev. C66, 025205 (2002)
- R. Fries, B. Mueller, C. Nonaka & S.A. Bass, Phys. Rev. Lett. 90
- R. Fries, B. Mueller, C. Nonaka & S.A. Bass, Phys. Rev. C68, 044902 (2003)
- V. Greco, C.M. Ko and P. Levai, Phys. Rev. Lett. 90 - **AMPT**

## Anisotropic flow:

- S. Voloshin, QM2002, nucl-ex/020014
- Z.W. Lin & C.M. Ko, Phys. Rev. Lett 89, 202302 (2002)

ksle D. Molnar & S. Voloshin, nucl-th/0302014

# Recombination+Fragmentation Model

R.J. Fries, C. Nonaka, B. Mueller & S.A. Bass, PRL 90 202303 (2003) ; PRC68, 044902 (2003)

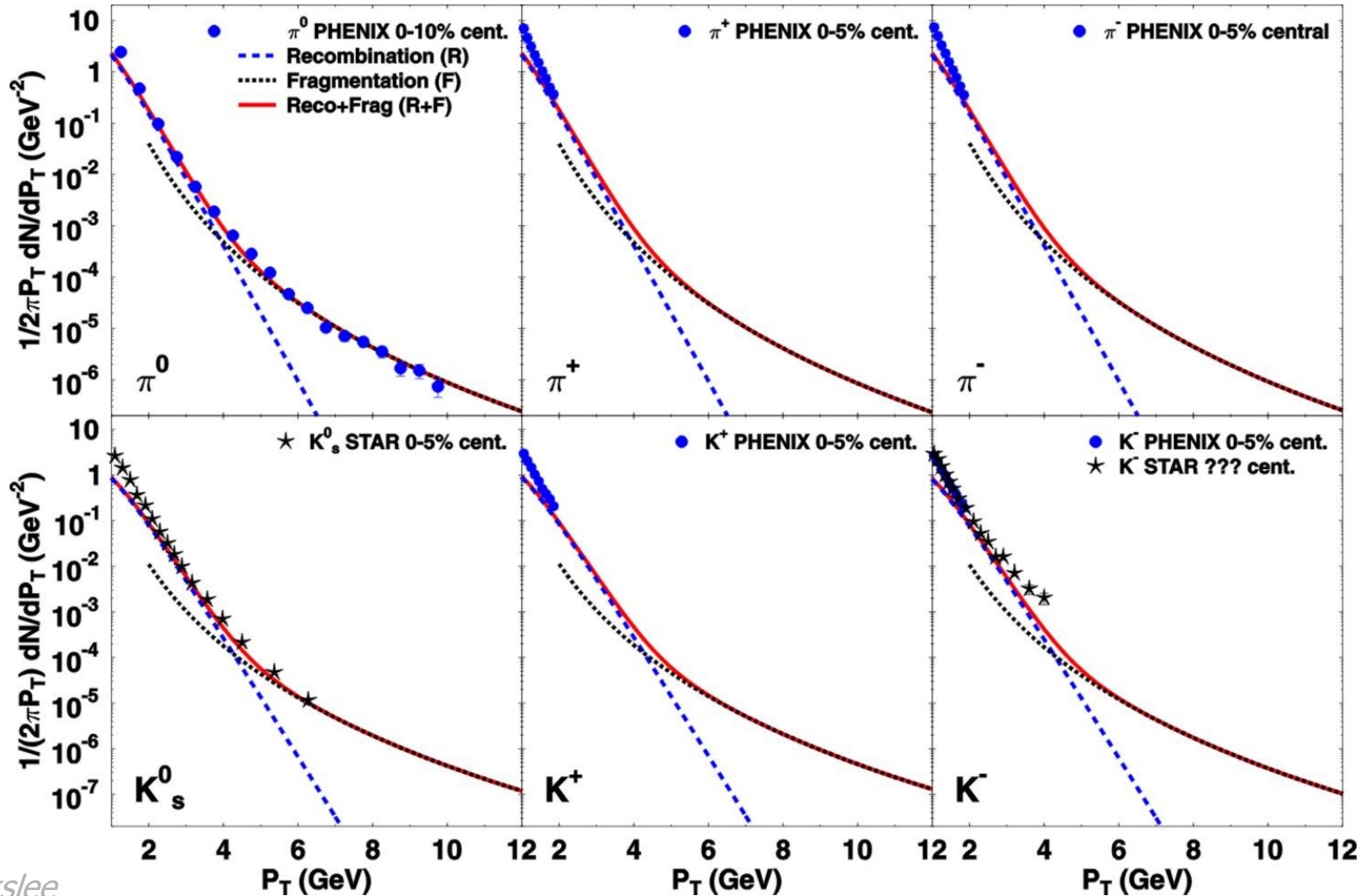
basic assumptions:

- at low  $p_t$ , the quarks and antiquark spectrum is thermal and they recombine into hadrons locally “at an instant”:

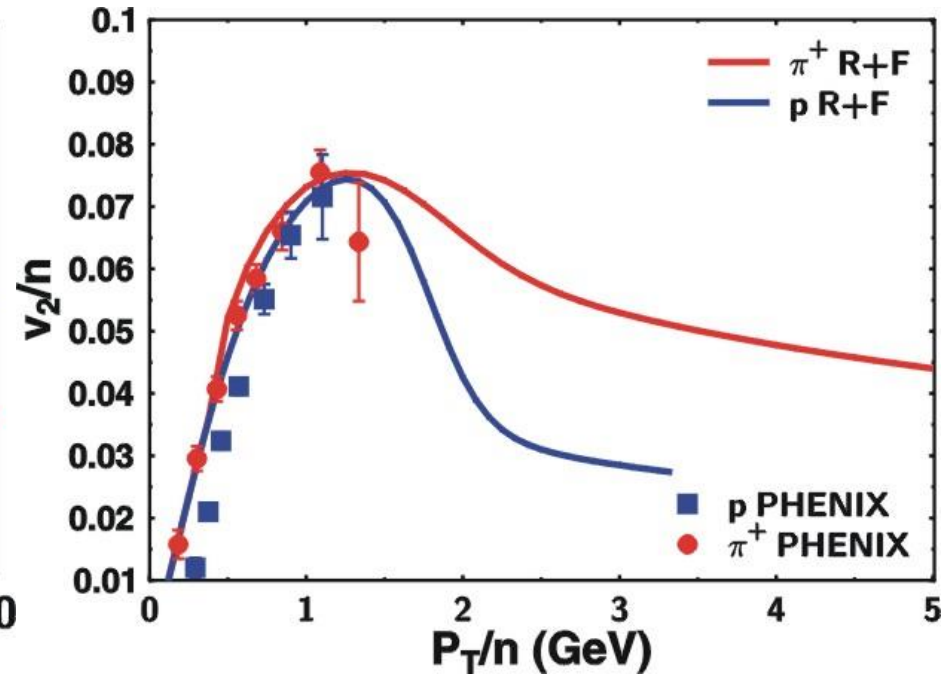
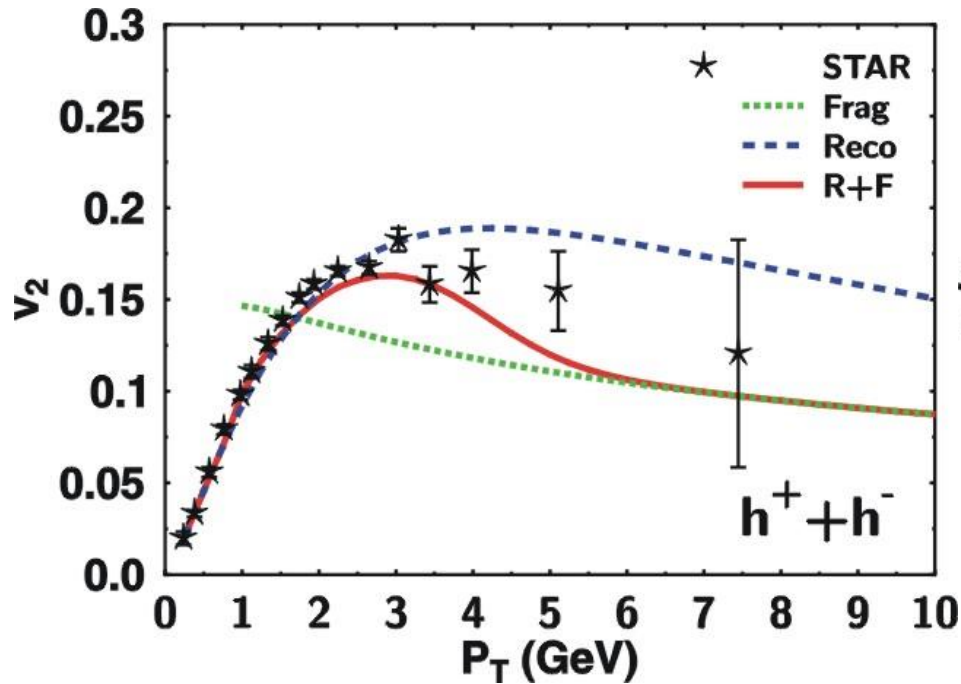


- at high  $p_t$ , the parton spectrum is given by a pQCD power law, partons suffer jet energy loss and hadrons are formed via fragmentation of quarks and gluons

# Hadron Spectra I



# Elliptic Flow: Recombination vs. Fragmentation



- high  $p_t$ :  $v_2$  for all hadrons merge, since  $v_2$  from energy-loss is flavor blind
- charged hadron  $v_2$  for high  $p_t$  shows *universal & limiting* fragmentation  $v_2$
- quark number scaling breaks down in the fragmentation domain

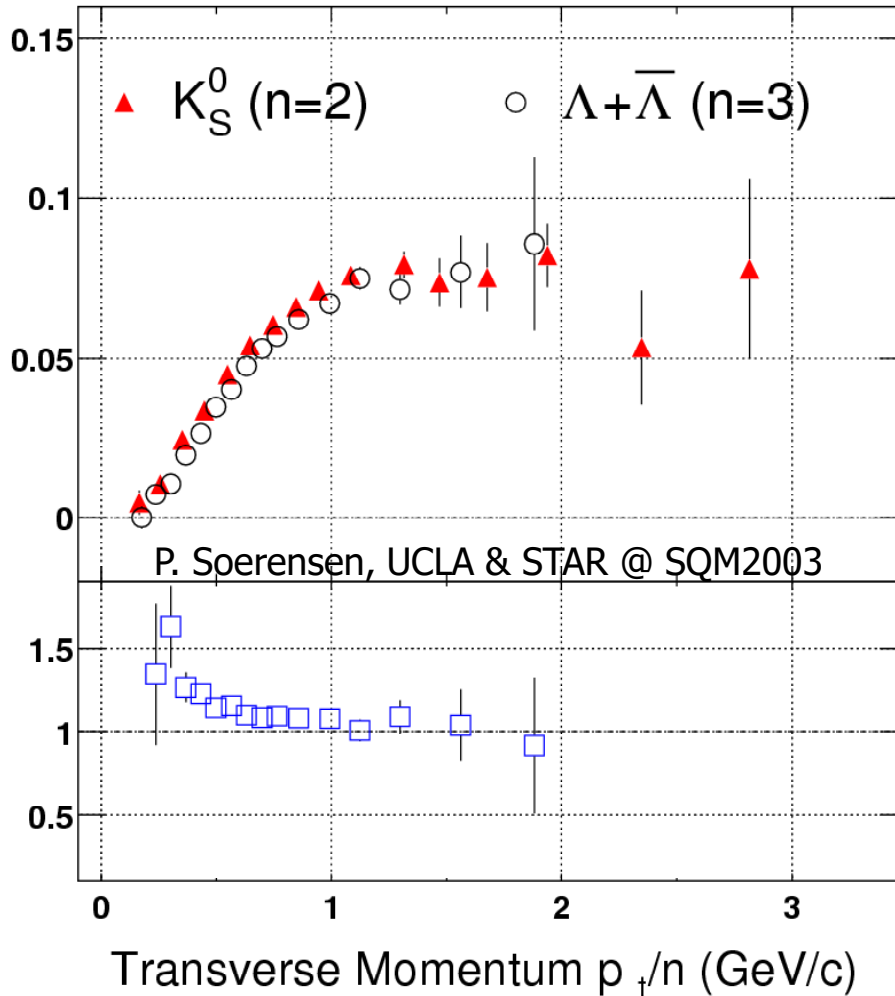
# Parton Number Scaling of $v_2$

- in leading order of  $v_2$ , recombination predicts:

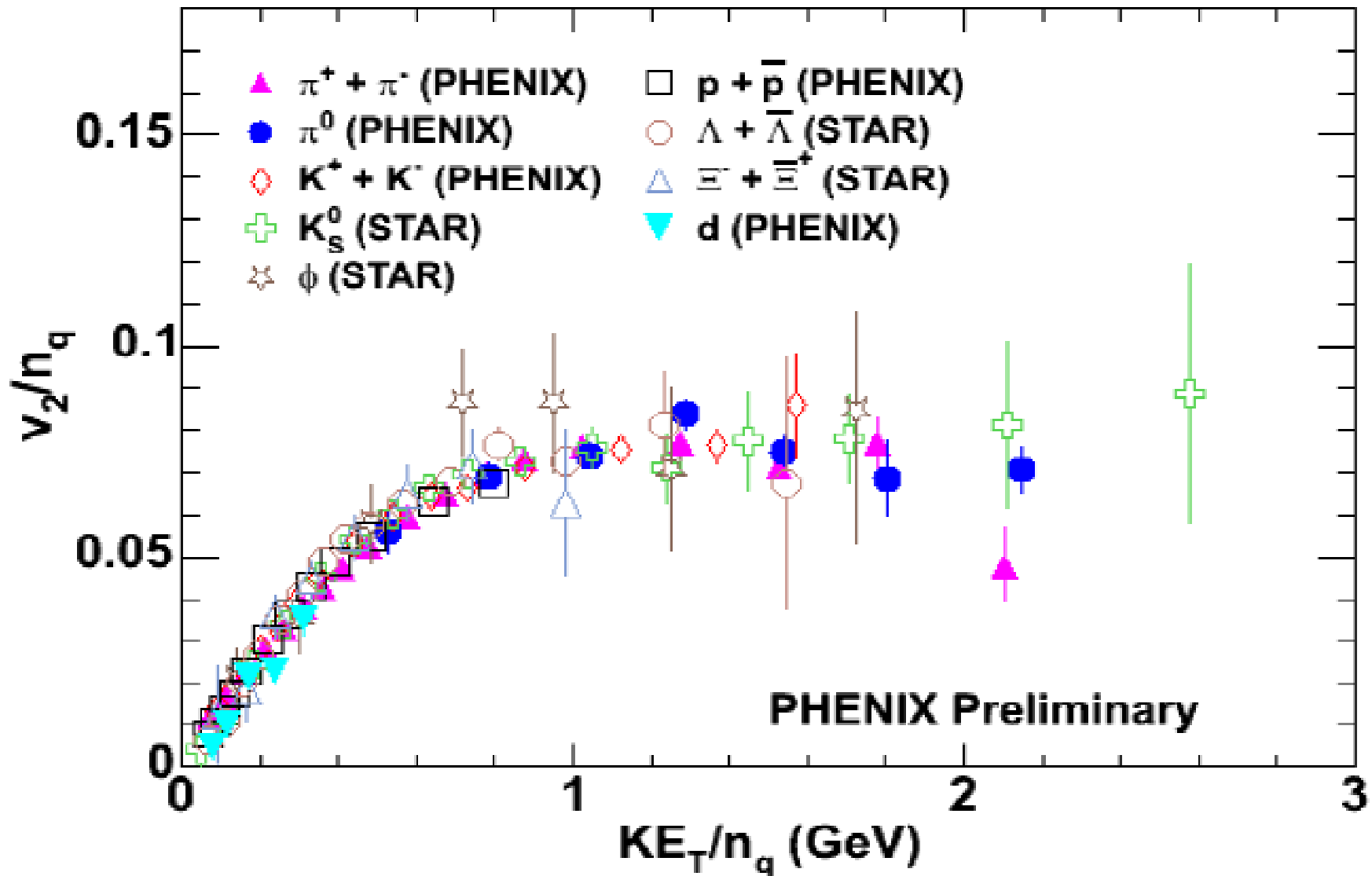
$$v_2^M(p_t) = 2v_2^P \left( \frac{p_t}{2} \right)$$

$$v_2^B(p_t) = 3v_2^P \left( \frac{p_t}{3} \right)$$

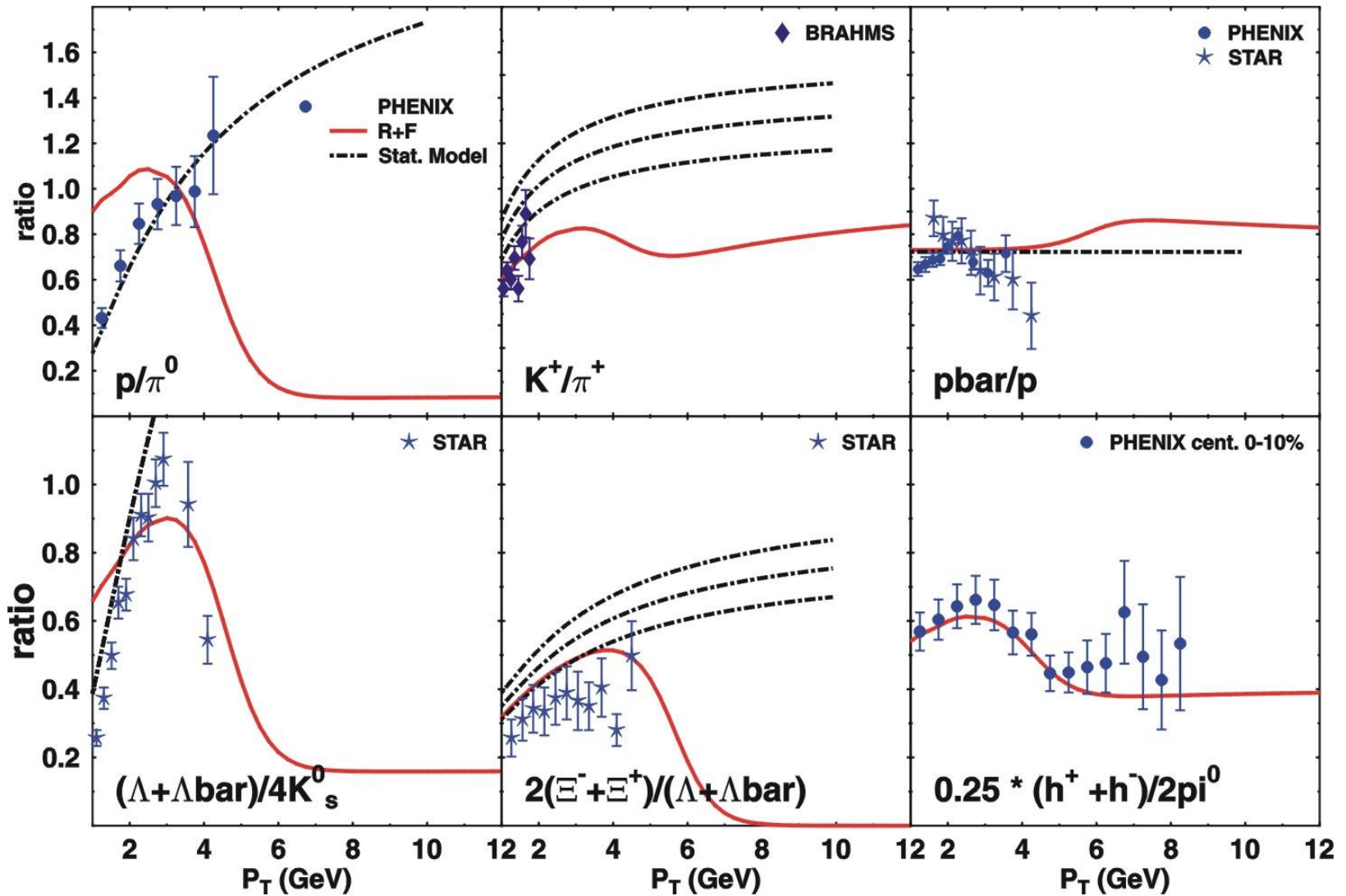
- smoking gun for recombination
- measurement of partonic  $v_2$  !



# Parton Number Scaling of $v_2$



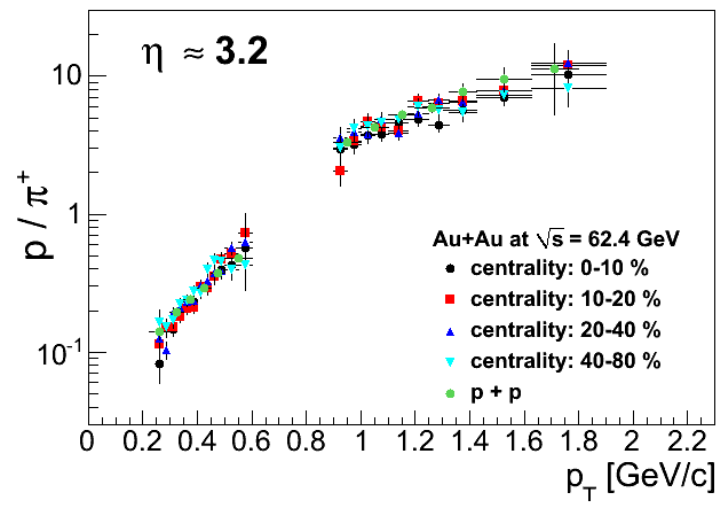
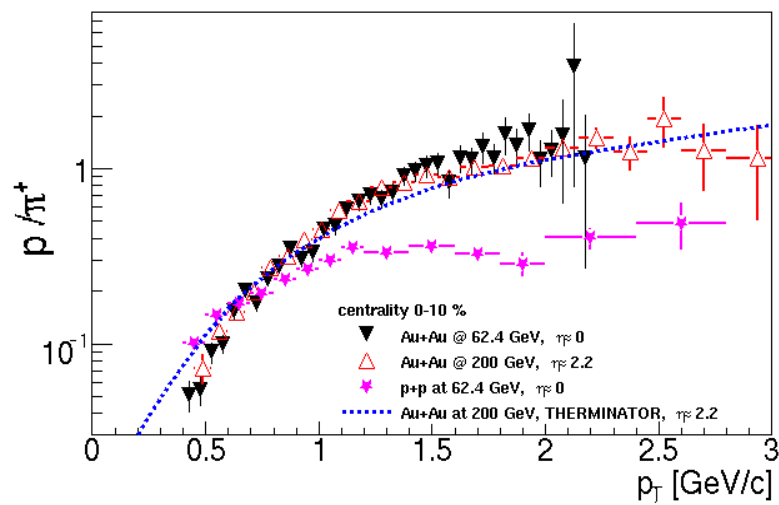
# Hadron Ratios vs. $p_t$



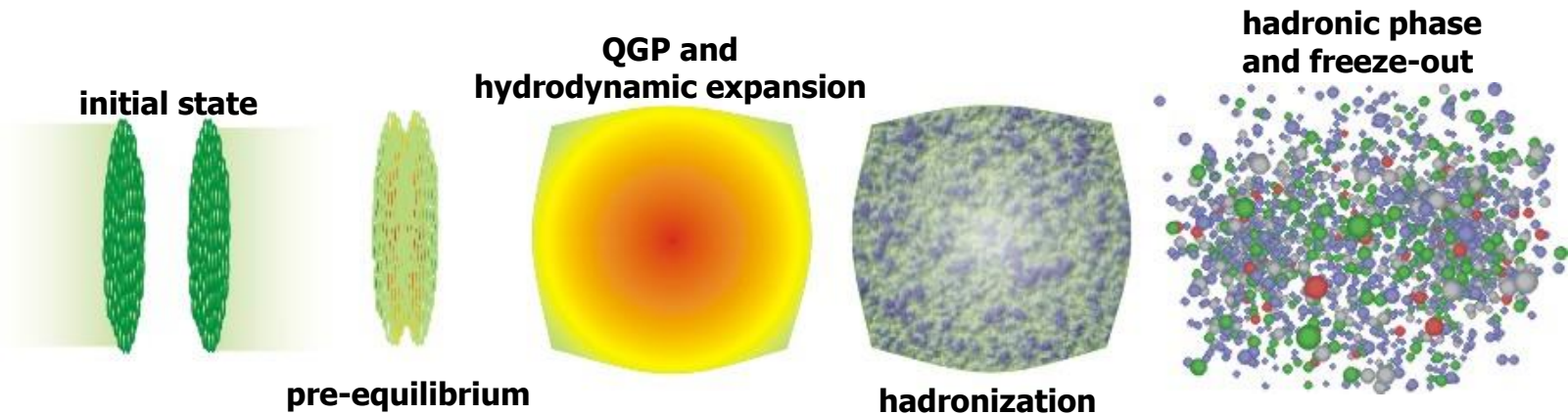


## Energy systematics – 200 and 62

- Mid-rapidity 62 GeV similar to  $\eta \sim 2.2$  at 200 GeV with significant  $P/\pi$  ratios
- At  $y \sim 0$  and  $y \sim 2.2$  significant medium effects
- Close to beam rapidity (forward at 62 GeV)  $P/\pi$  exhibits no centrality dependence and is similar to pp indicating little influence from media (quark coalescence, hadronic re-scattering)



# Exploring QCD Matter at RHIC and LHC



## QGP

- hydrodynamic evolution - C. Nonaka
- reasonable for a perfect fluid

## Hadronization via recombination

- recombination

## Hadronic rescattering

- URQMD - S. Bass

# Recombination of thermal quarks I

$$E \frac{dN^M_i}{d^3 P} = \frac{d_i}{(2\pi)^3} \int_{\Sigma} P^\mu d\sigma_\mu G^M_{ab}(k, P-k)$$

$$E \frac{dN^B_i}{d^3 p} = \frac{d_i}{(2\pi)^3} \int_{\Sigma} P^\mu d\sigma_\mu G^B_{abc}(k_1, k_2, P-k_1-k_2)$$

$$G^M_{ab} = \int d^3 k \bar{w}_a(k) \left| \bar{\phi}_M(k) \right|^2 \bar{w}_b(P-k)$$

$$G^B_{abc} = \int d^3 k_1 d^3 k_2 w_a(k_1) w_b(k_2) w_c(P-k_1-k_2) \left| \bar{\phi}_B(k_1, k_2, P-k_1-k_2) \right|^2$$

- wave functions are best known in the light cone frame.
- for a thermal distribution:  $w(r, k) \propto \exp(-(k \cdot u - \mu_q)/T)$

# Recombination of thermal quarks II

- hadronization hypersurface : mixed phase
- Energy conservation during the recombination

➤ For a cell at hadronization,  $E_i \delta\alpha_\tau$  is the energy available for hadronization where  $\alpha = V_Q / V_{tot}$  is the volume fraction of QGP phase.

$$\sum_i E_i^{cell} \Delta\alpha_\tau^i = \sum_j E_j^{had}$$

- In a hydrodynamic cell,  $i$ , hadronization rate in  $\mu$  direction is proportional to  $\Delta\alpha_\mu^i = \alpha_\mu^i - \alpha_\mu^{i+1}$

$$\int p^\mu d\sigma_\mu \rightarrow \sum_i \sum_\mu p^\mu d\sigma_\mu \Delta\alpha_\mu^i$$

# Recombination of thermal quarks III

- $G^M_{ab}(k_1, k_2)$  depends only on the quark masses and is independent of hadron mass.
- Since  $m_q \neq m_s$ , there are 4 cases for mesons and 8 cases for baryons.

$$q\bar{q} \rightarrow M1, \quad q\bar{s} \rightarrow M2, \quad s\bar{q} \rightarrow M3, \quad s\bar{s} \rightarrow M4$$

$$qqq \rightarrow B1, \quad qq_s \rightarrow B2, \quad q_s s \rightarrow B3, \quad s s s \rightarrow B4$$

$$\bar{q}\bar{q}\bar{q} \rightarrow B5, \quad \bar{q}\bar{q}\bar{s} \rightarrow B6, \quad \bar{q}\bar{s}\bar{s} \rightarrow B7, \quad \bar{s}\bar{s}\bar{s} \rightarrow B8$$

- Inside a group, relative abundances are assumed to be

$$N_i / N_j = e^{(m_i - m_j)/T} C_i / C_j \quad (C_i : \text{degeneracy factors})$$

kslee

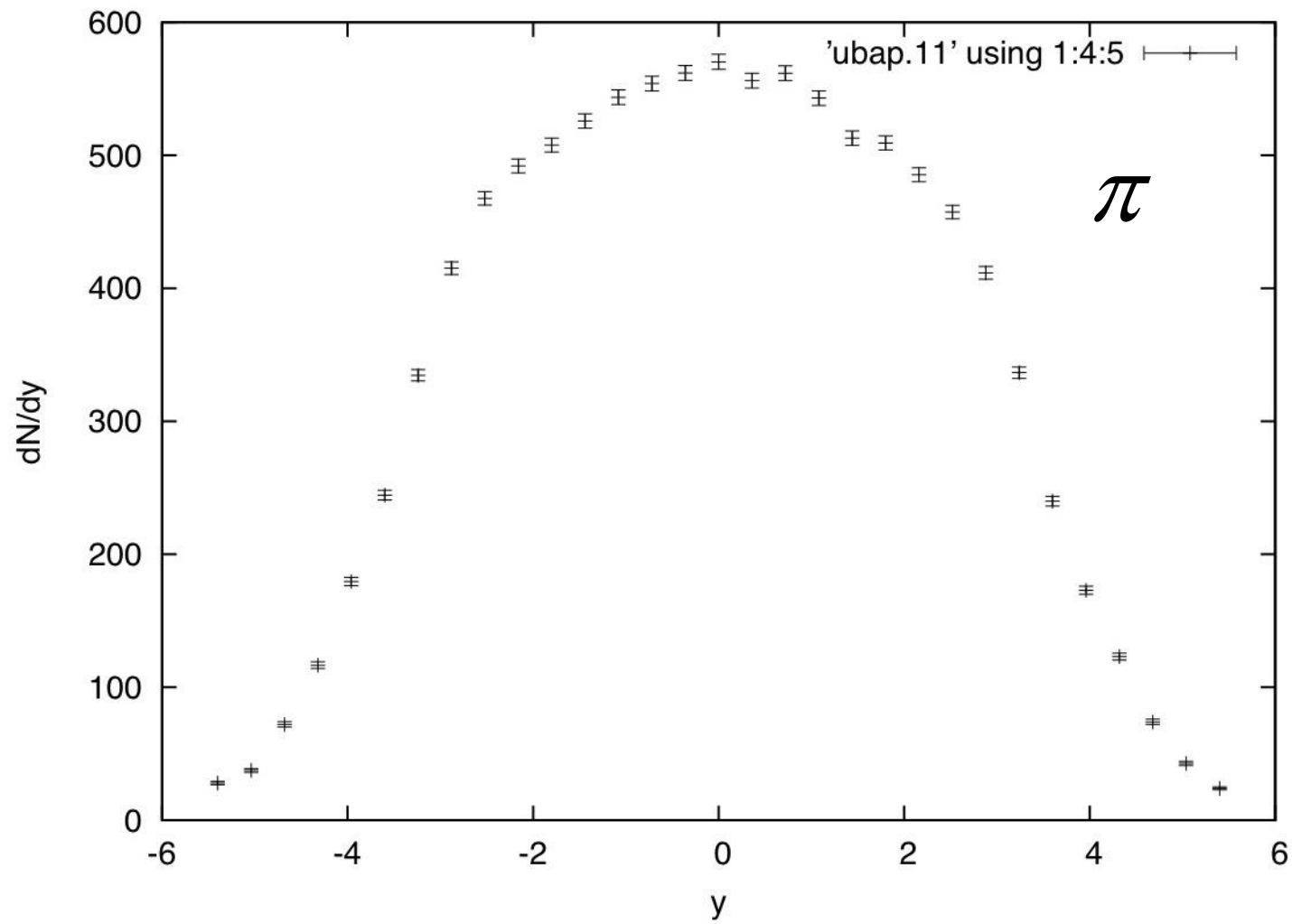
e.g.  $M1 = (\pi, \omega, \rho, \dots)$

# Recombination of thermal quarks IV

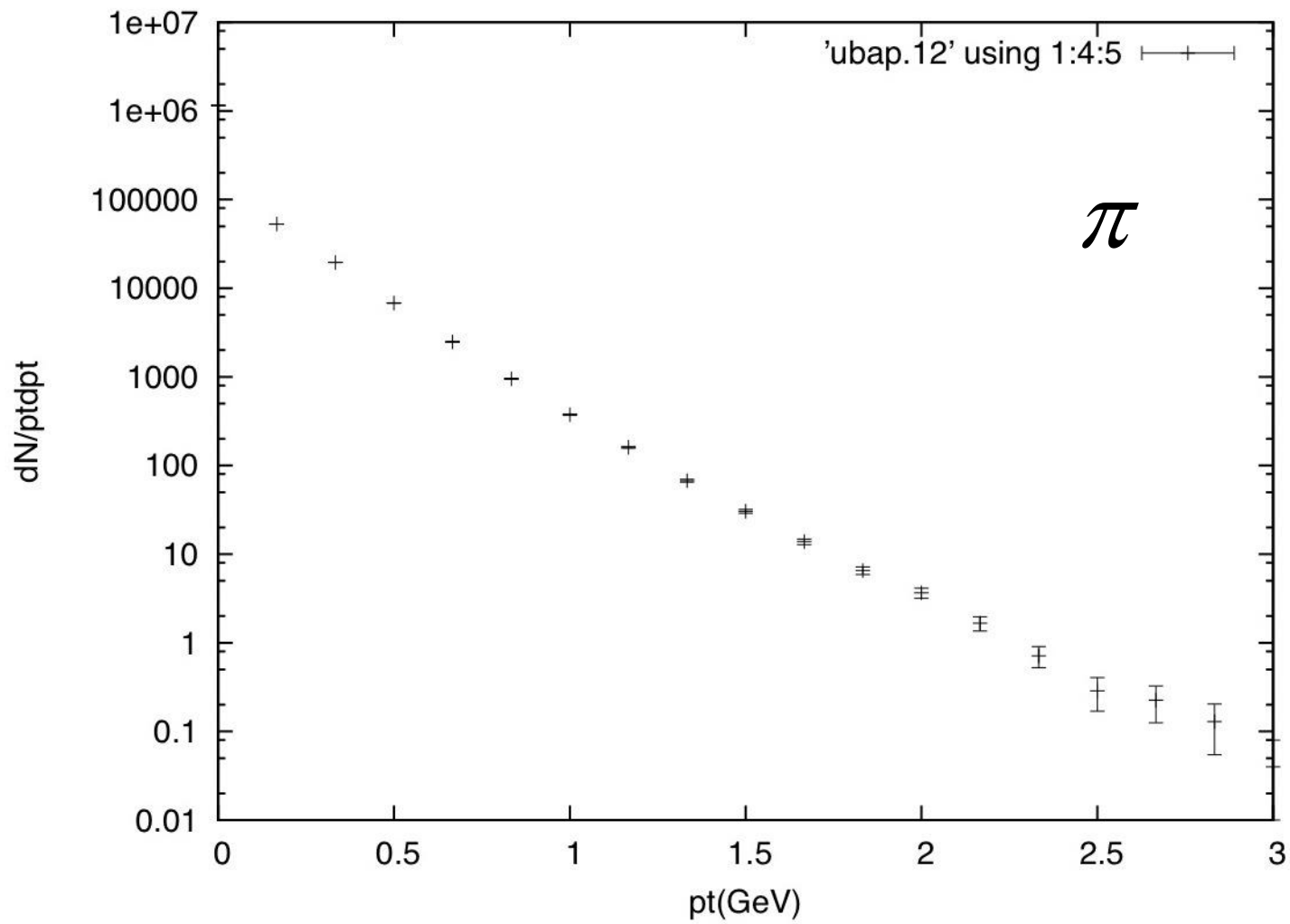
- After hadronization, hadronic rescattering is simulated by URQMD !

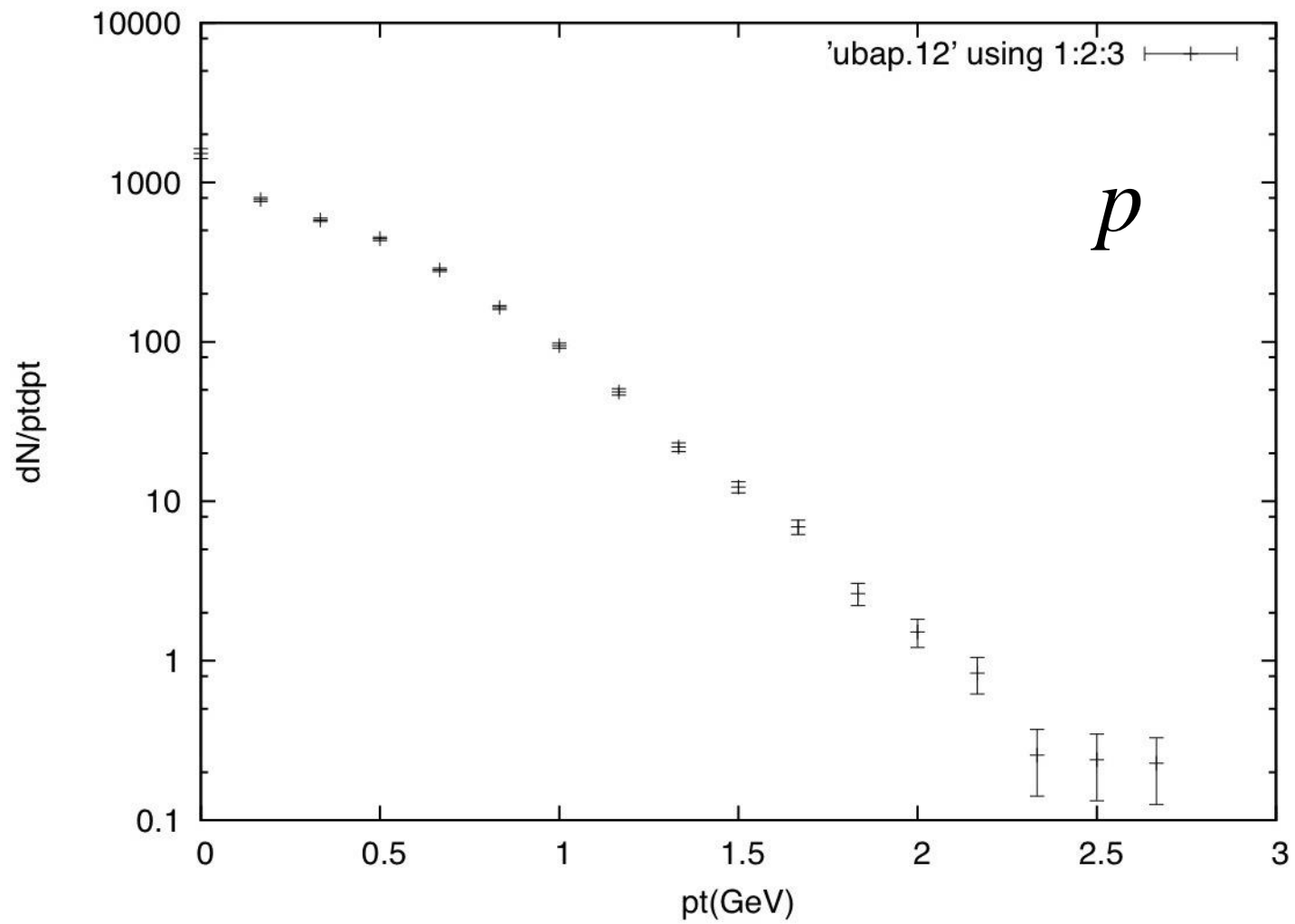
## Results

- hadron spectra
- hadron ratios – coming soon
- $R_{AA}$  – coming soon
- elliptic flow – coming soon
- Comparison with statistical model









# Summary & Outlook

## Hadronization of a QGP via Recombination:

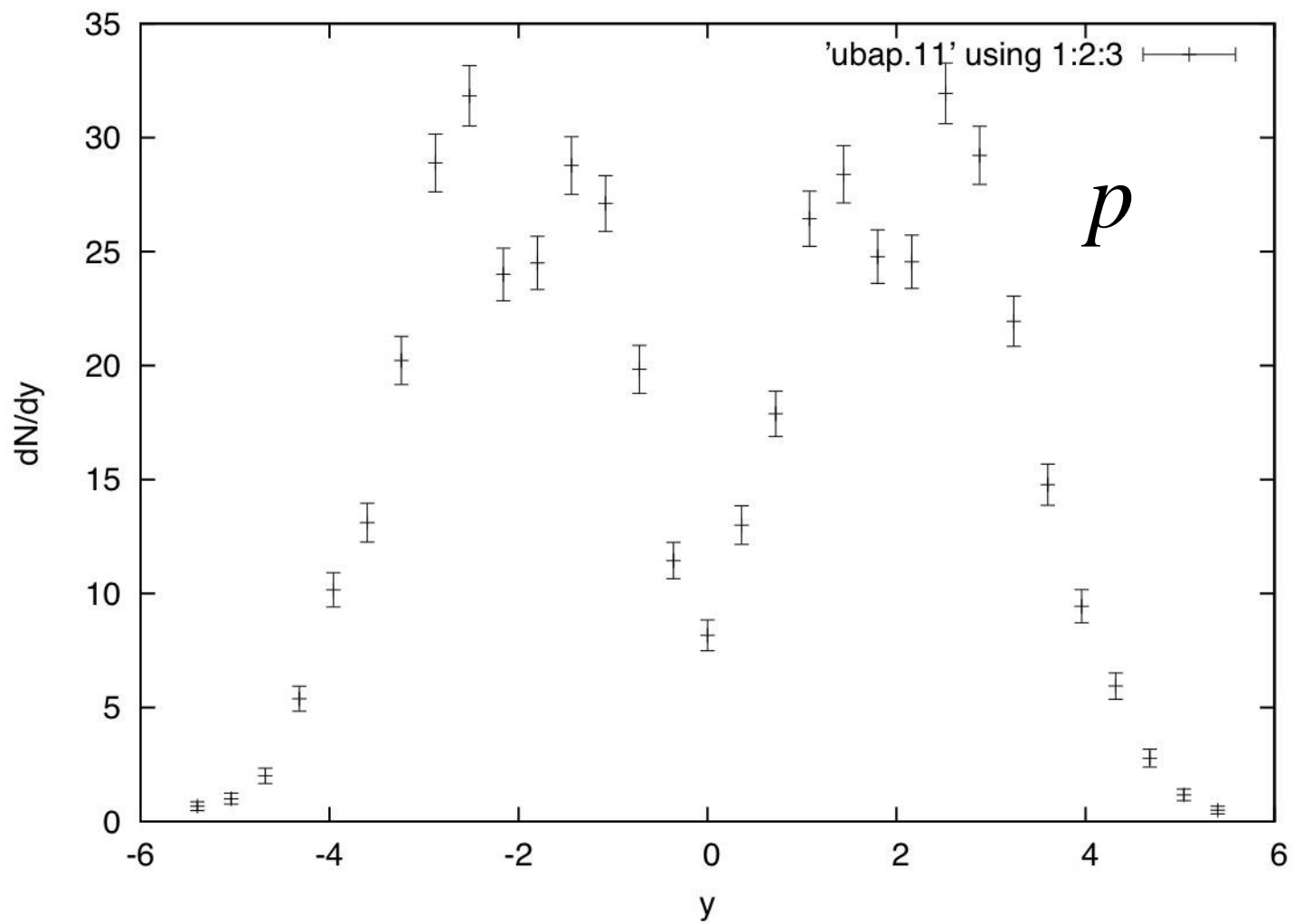
- Hydrodynamic evolution of a QGP until mixed phase
  - reasonable for a perfect fluid
- hadronization via quark recombination
  - quark number scaling of elliptic flow, ...
- hadronic rescattering by URQMD
  - different chemical and thermal freeze-out

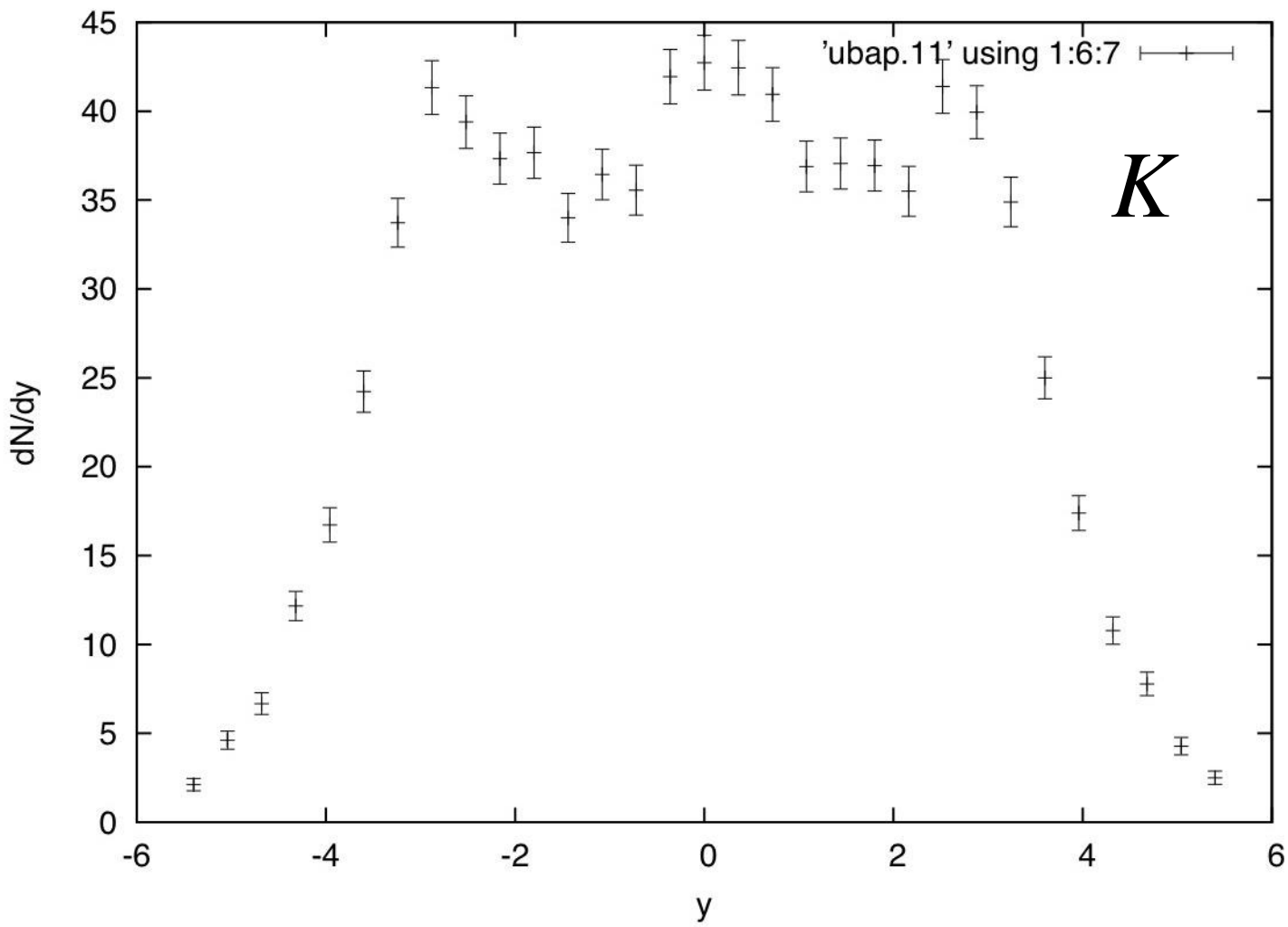
No problem of freeze-out prescription !!!

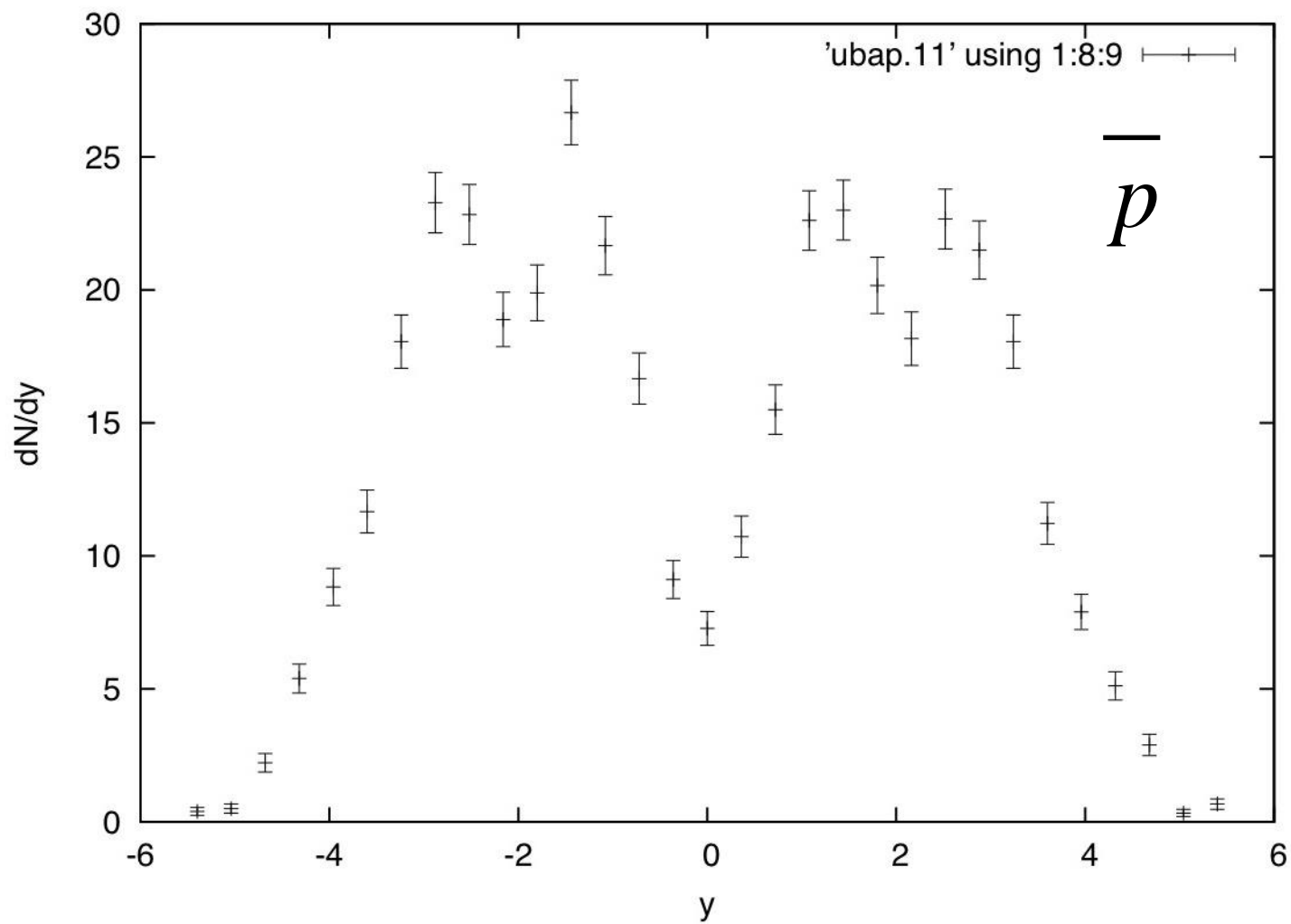
issues to be addressed in the future:

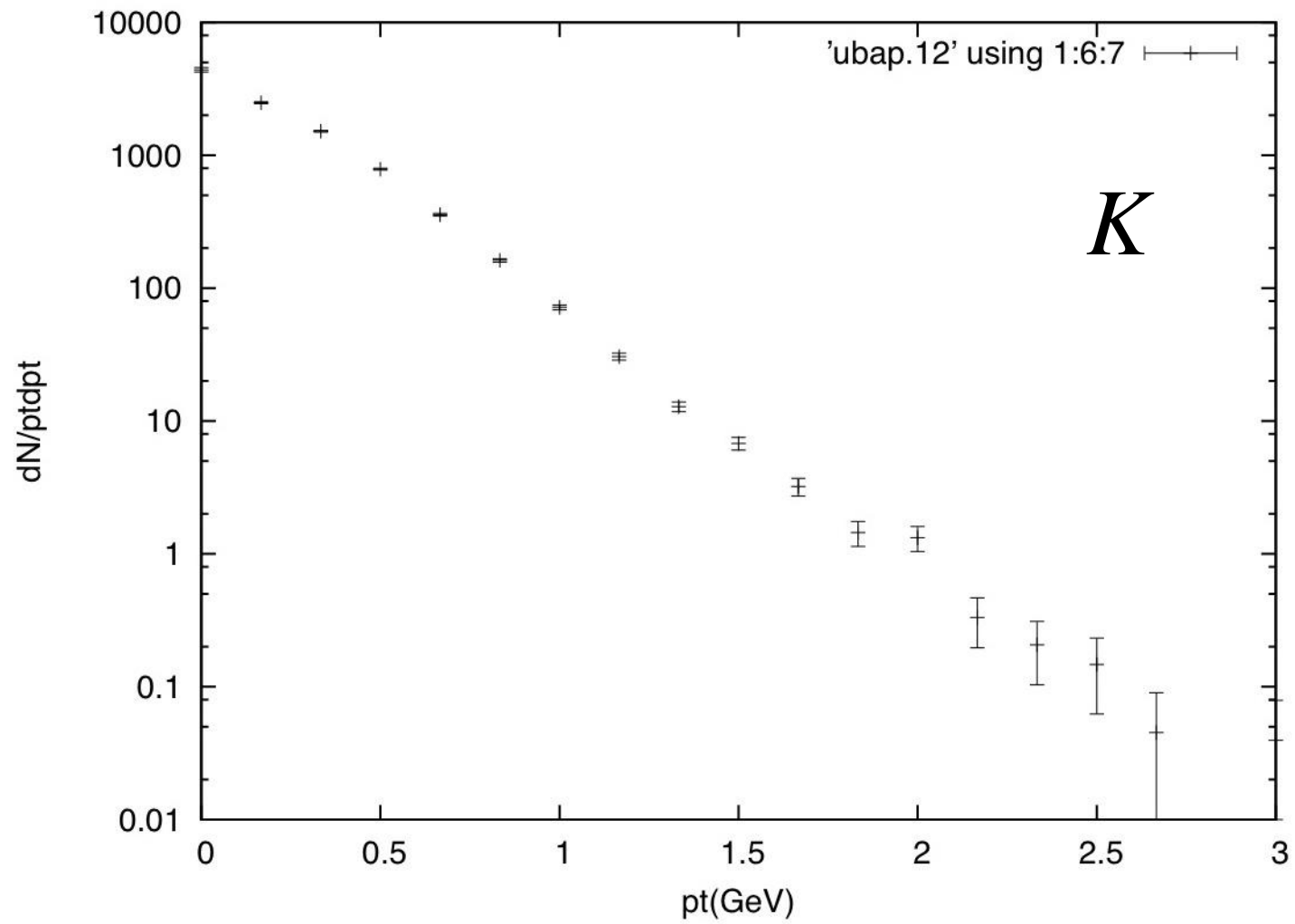
- results are coming soon.
- comparison with statistical hadronization
- combine with viscous hydrodynamics or parton cascade

The End

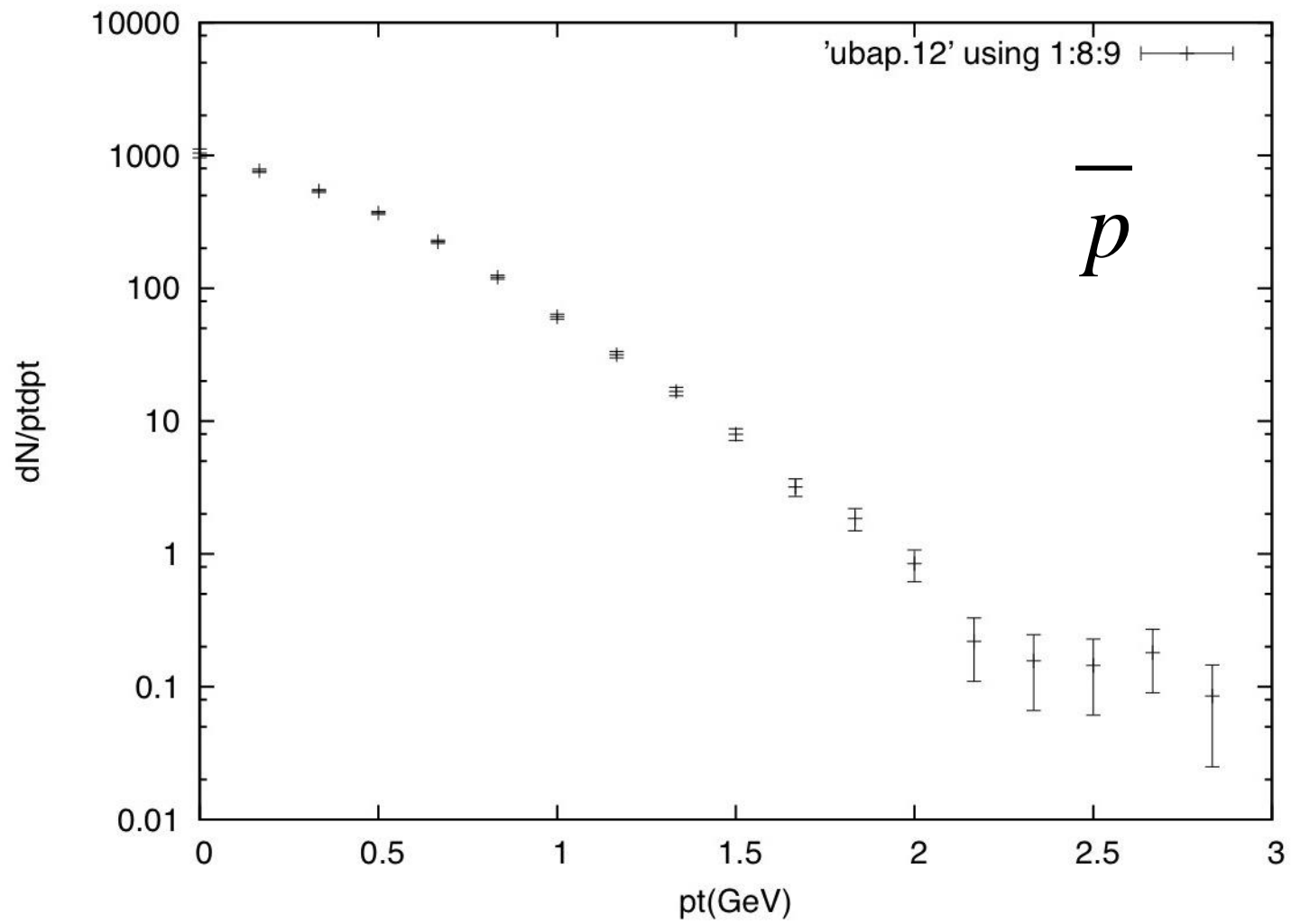












# Statistical Model vs. Recombination

- in the Statistical Model, the hadron distribution at freeze-out is given by:

$$E \frac{d^3 N_i}{d^3 P} = \int_{\sigma} f_i(P \cdot u) P^\lambda d\sigma_\lambda$$

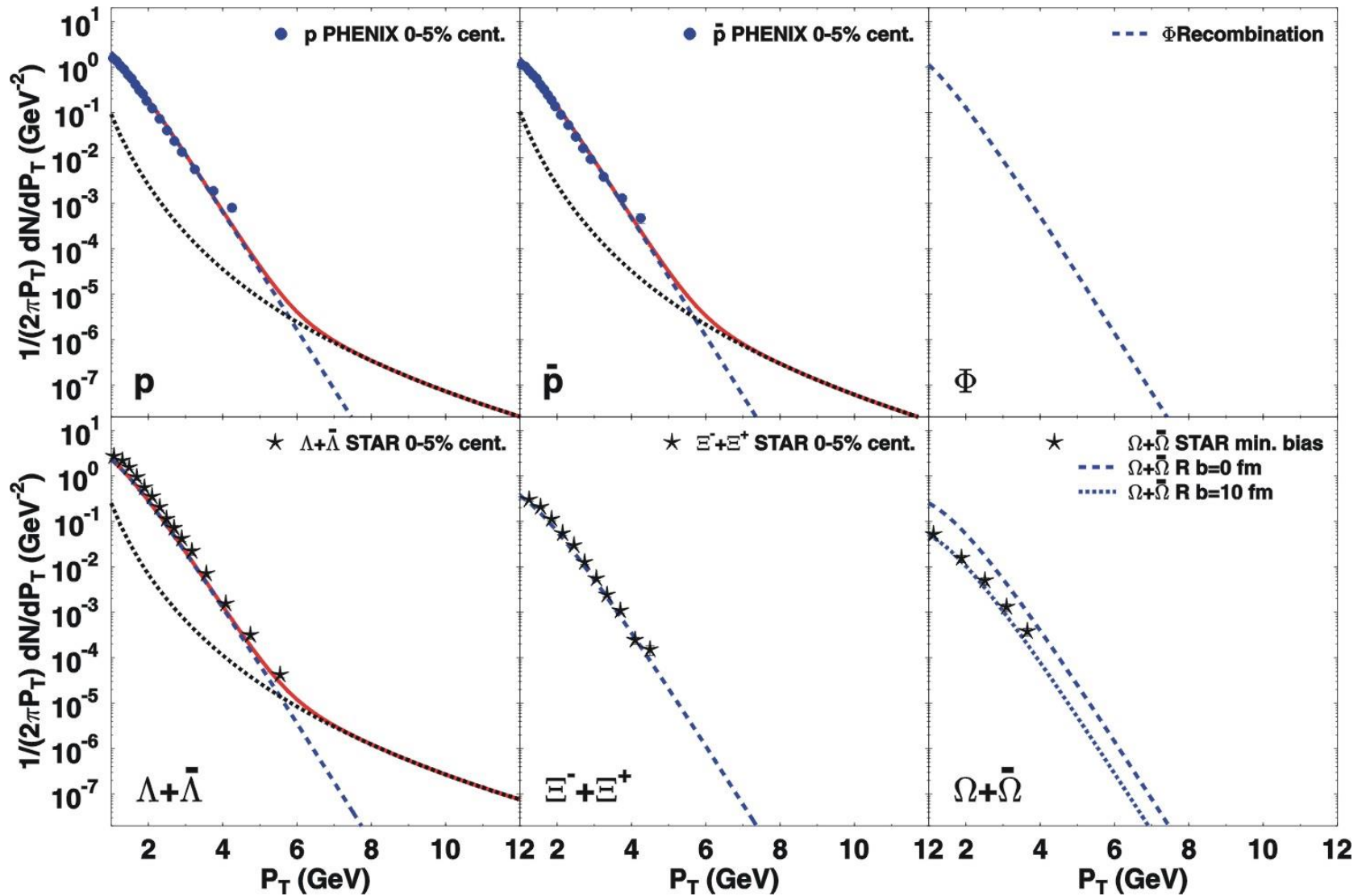
$$\text{with } f_i(P \cdot u) = \frac{g_i}{(2\pi)^3} \frac{1}{\exp\left[\frac{(P \cdot u - \mu_B B_i - \mu_s S_i - \mu_I I_i)}{T} \pm 1\right]}$$

$$G^M_{ab} = \int d^3 k \bar{w}_a(k) \left| \bar{\phi}_M(k) \right|^2 \bar{w}_b(P - k)$$

$$G^B_{abc} = \int d^3 k_1 d^3 k_2 w_a(k_1) w_b(k_2) w_c(P - k_1 - k_2) \left| \bar{\phi}_B(k_1, k_2, P - k_1 - k_2) \right|^2$$

- If  $E_h(P) = E_a(k) + E_b(P - k)$  , recombination is identical to SM!

# Hadron Spectra II



# Elliptic Flow: partons at low $p_t$

- azimuthal anisotropy of parton spectra is determined by elliptic flow:

$$\frac{d^2 N}{p_t dp_t d\phi_p} = \frac{1}{2\pi} \left[ \frac{dN}{p_t dp_t} \right] \left( 1 + 2v_2 \cos(2\phi_p) \right) \quad (\phi_p: \text{azimuthal angle in p-space})$$

- with Blastwave parametrization for parton spectra:

$$v_2(p_t) = \langle \cos(2\phi_p) \rangle = \frac{\int_0^{2\pi} d\phi_s \cos(2\phi_s) I_2 \left( \frac{p_t \sinh(\rho(\phi_s))}{T} \right) K_1 \left( \frac{m_t \cosh(\rho(\phi_s))}{T} \right)}{\int_0^{2\pi} d\phi_s I_0 \left( \frac{p_t \sinh(\rho(\phi_s))}{T} \right) K_1 \left( \frac{m_t \cosh(\rho(\phi_s))}{T} \right)}$$

- azimuthal anisotropy is parameterized in coordinate space and is damped as a function of  $p_t$ :

$$\rho(\phi_s) = \frac{1}{2} \ln \left( \frac{1 + \beta_t}{1 - \beta_t} \right) \left( 1 + \alpha_p(p_t) \cos(2\phi_s) \right) \quad \text{and} \quad \alpha_p(p_t) = -\alpha_0 \frac{1}{1 + (p_t/p_0)^2}$$

# Recombination: relativistic formalism

- choose a hypersurface  $\Sigma$  for hadronization
- use local light cone coordinates (hadron defining the + axis)
- $w_a(\mathbf{r}, \mathbf{p})$ : single particle Wigner function for quarks at hadronization
- $\Phi_M$  &  $\Phi_B$ : light-cone wave-functions for the meson & baryon respectively
- $x, x'$  &  $(1-x)$ : momentum fractions carried by the quarks
- integrating out transverse degrees of freedom yields:

$$E \frac{dN_M}{d^3 P} = \int_{\Sigma} d\sigma \frac{P \cdot u}{(2\pi)^3} \sum_{\alpha, \beta} \int dx w_{\alpha}(R, xP^+) \bar{w}_{\beta}(R, (1-x)P^+) \left| \bar{\phi}_M(x) \right|^2$$

$$E \frac{dN_B}{d^3 p} = \int_{\Sigma} d\sigma \frac{P \cdot u}{(2\pi)^3} \sum_{\alpha, \beta, \gamma} \int dx dx' w_{\alpha}(R, xP^+) w_{\beta}(R, x'P^+) w_{\gamma}(R, (1-x-x')P^+) \left| \bar{\phi}_B(x, x') \right|^2$$