

Plan for Korea Rare Isotope Accelerator (KoRIA*)

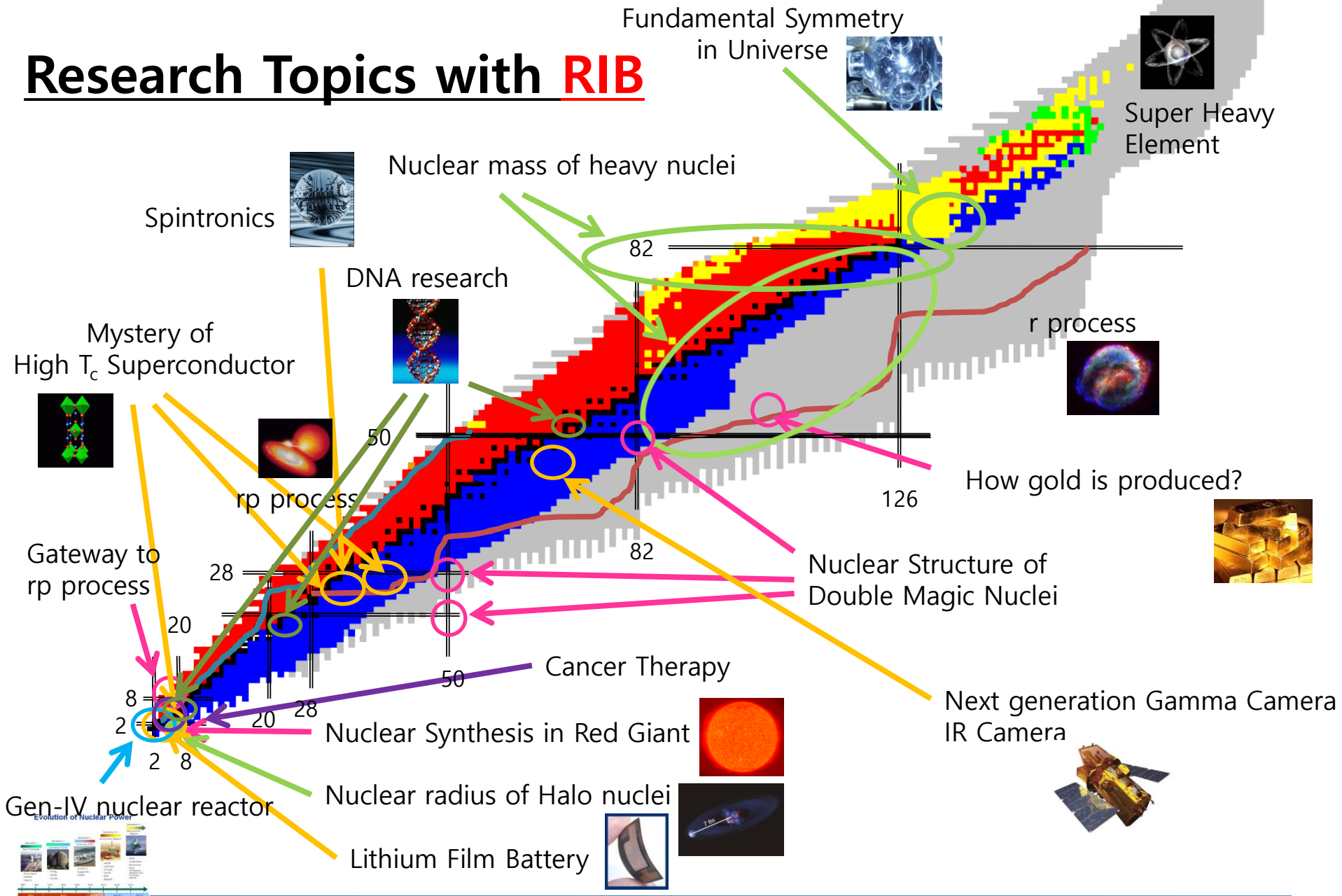
* Tentative

Byungsik Hong (Korea University)

Outline

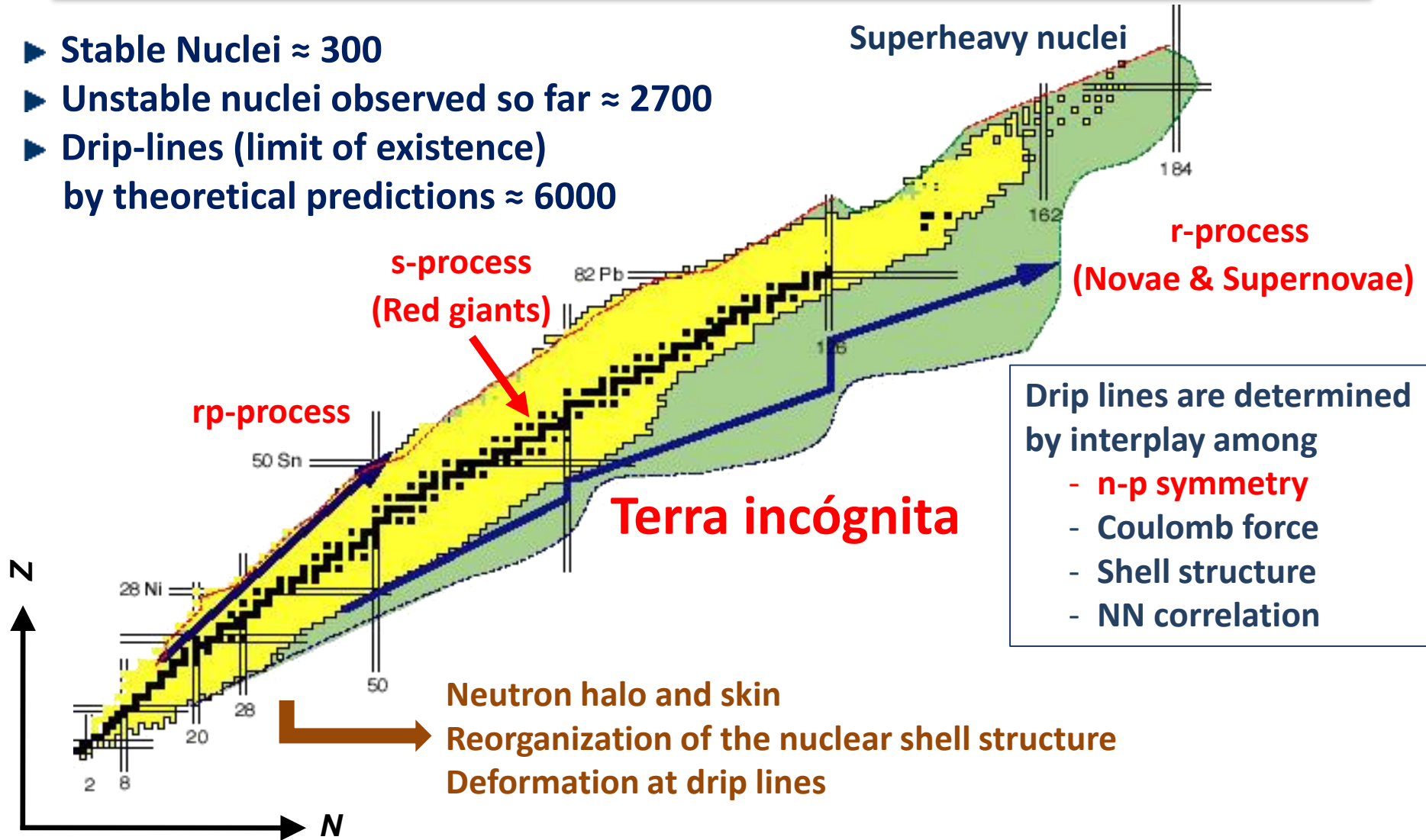
- Introduction
- Physics topics
- Experimental observables
- Summary

Research Topics with RIB



Nuclear Chart for Nuclear Physics

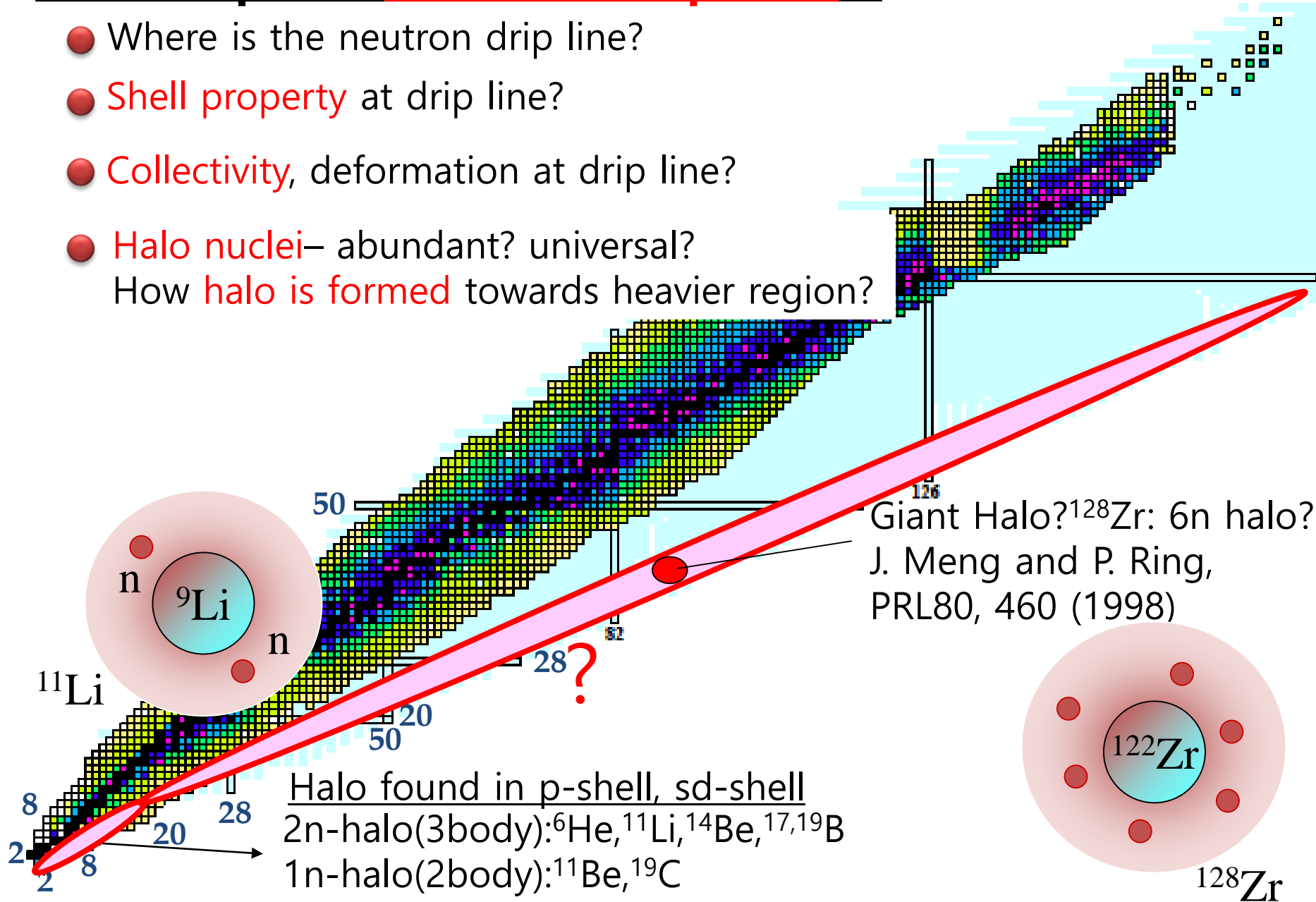
- ▶ Stable Nuclei ≈ 300
- ▶ Unstable nuclei observed so far ≈ 2700
- ▶ Drip-lines (limit of existence) by theoretical predictions ≈ 6000



Landscape of Neutron Drip Line?

From T. Nakamura

- Where is the neutron drip line?
- **Shell property** at drip line?
- **Collectivity**, deformation at drip line?
- **Halo nuclei**– abundant? universal?
How **halo is formed** towards heavier region?

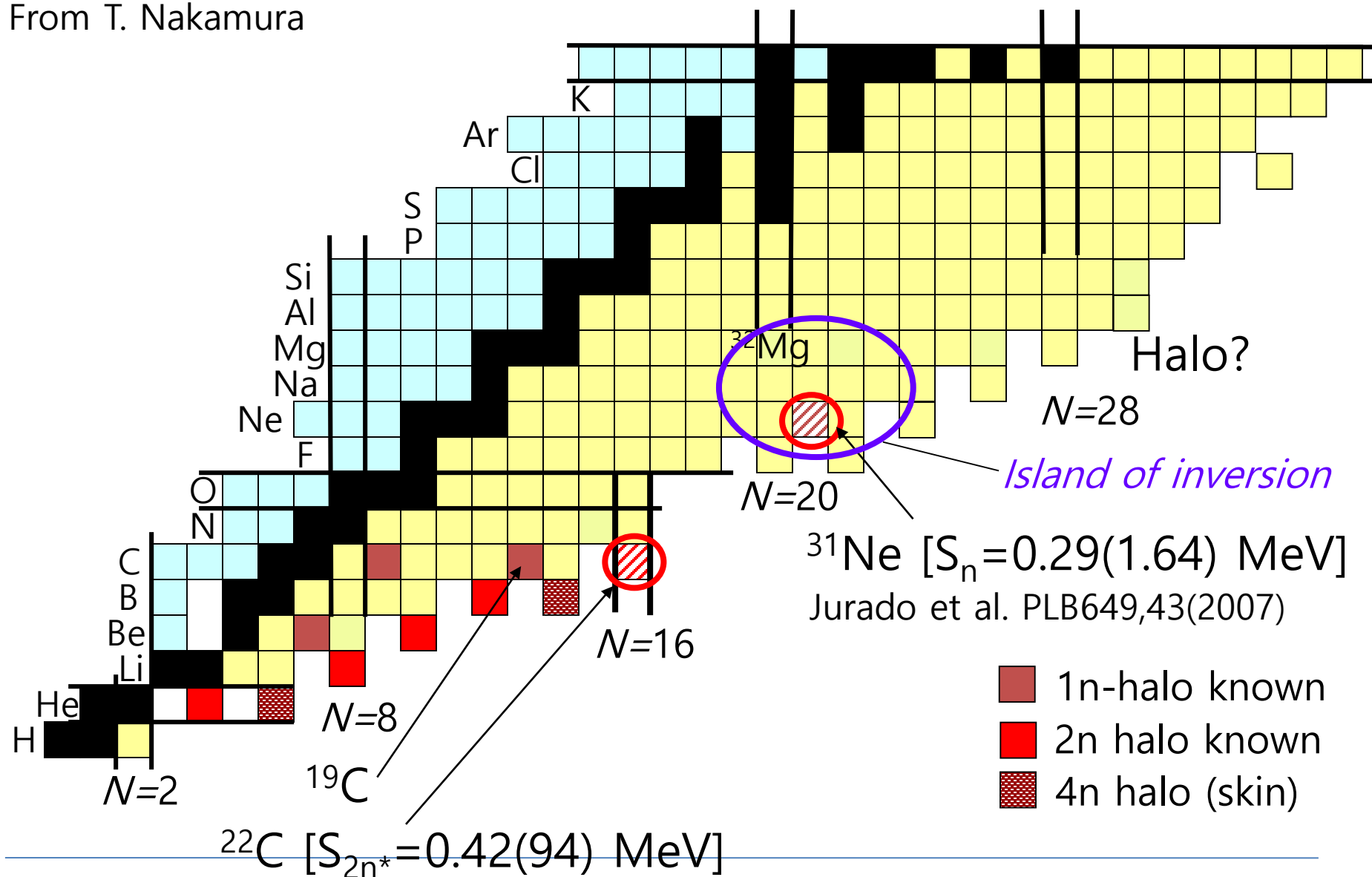


Giant Halo? ¹²⁸Zr: 6n halo?
J. Meng and P. Ring,
PRL80, 460 (1998)

Halo found in p-shell, sd-shell
2n-halo(3body): ⁶He, ¹¹Li, ¹⁴Be, ^{17,19}B
1n-halo(2body): ¹¹Be, ¹⁹C

More neutron halo's along the drip line?

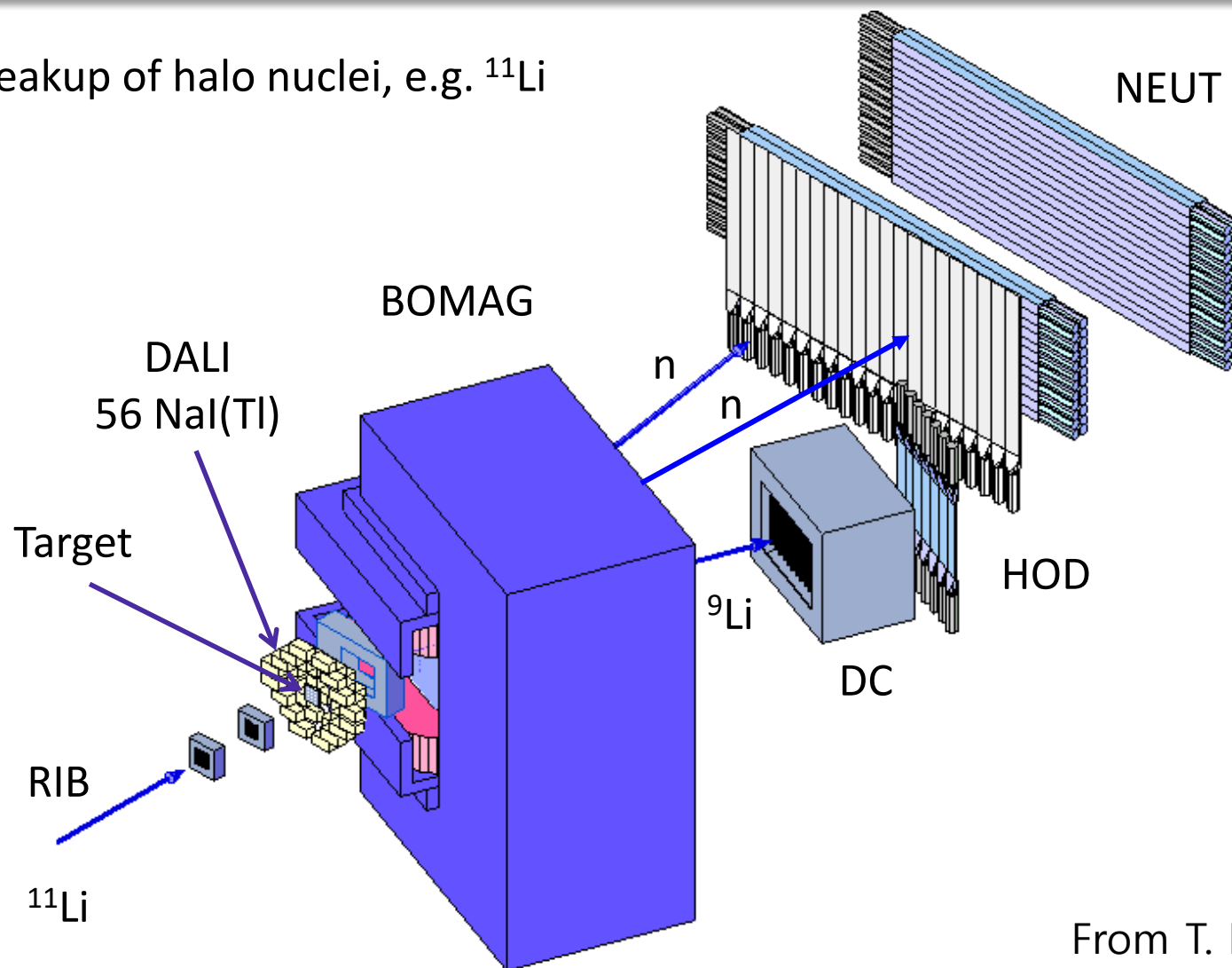
From T. Nakamura



* Estimated value by Audi & Wapstra

RIPS @ RIKEN

2n breakup of halo nuclei, e.g. ^{11}Li

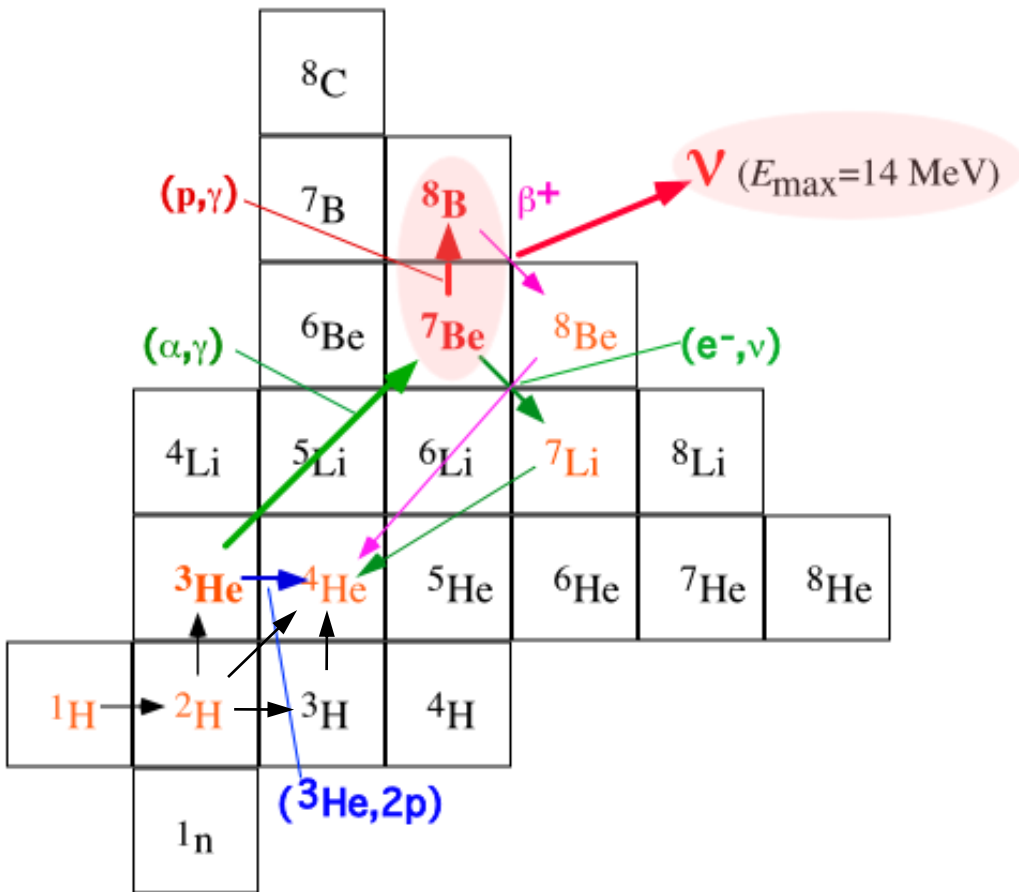


From T. Nakamura

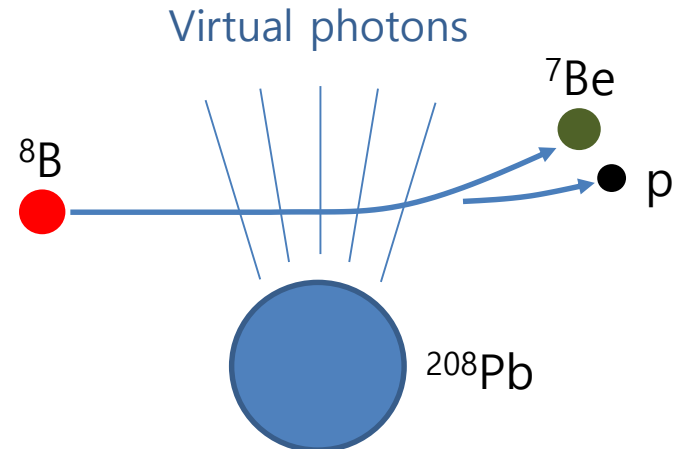
Nuclear Astrophysics

p-p chain in the sun ${}^7\text{Be}(p,\gamma){}^8\text{B}$

From T. Motobayashi



We can measure ${}^8\text{B}$
Coulomb dissociation
with RIB

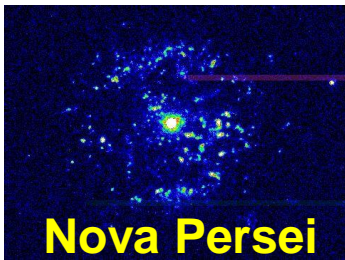
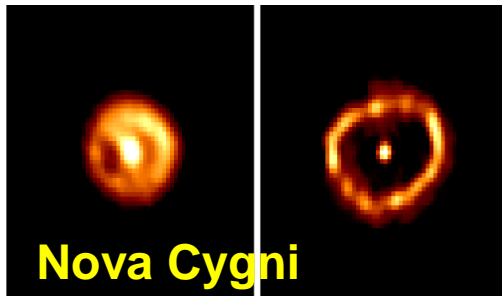


Nova Models

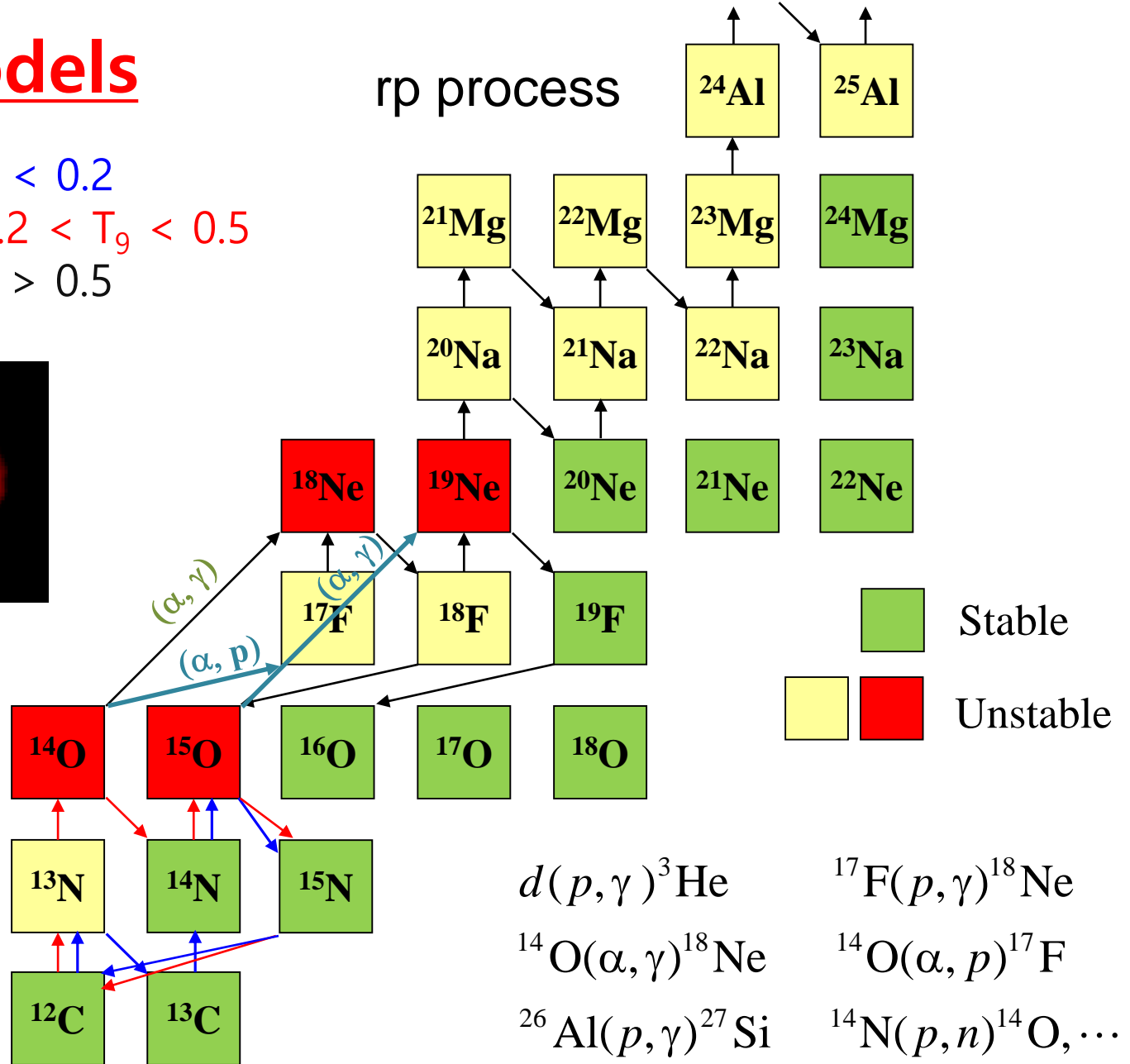
CNO cycle : $T_9 < 0.2$

HCNO cycle: $0.2 < T_9 < 0.5$

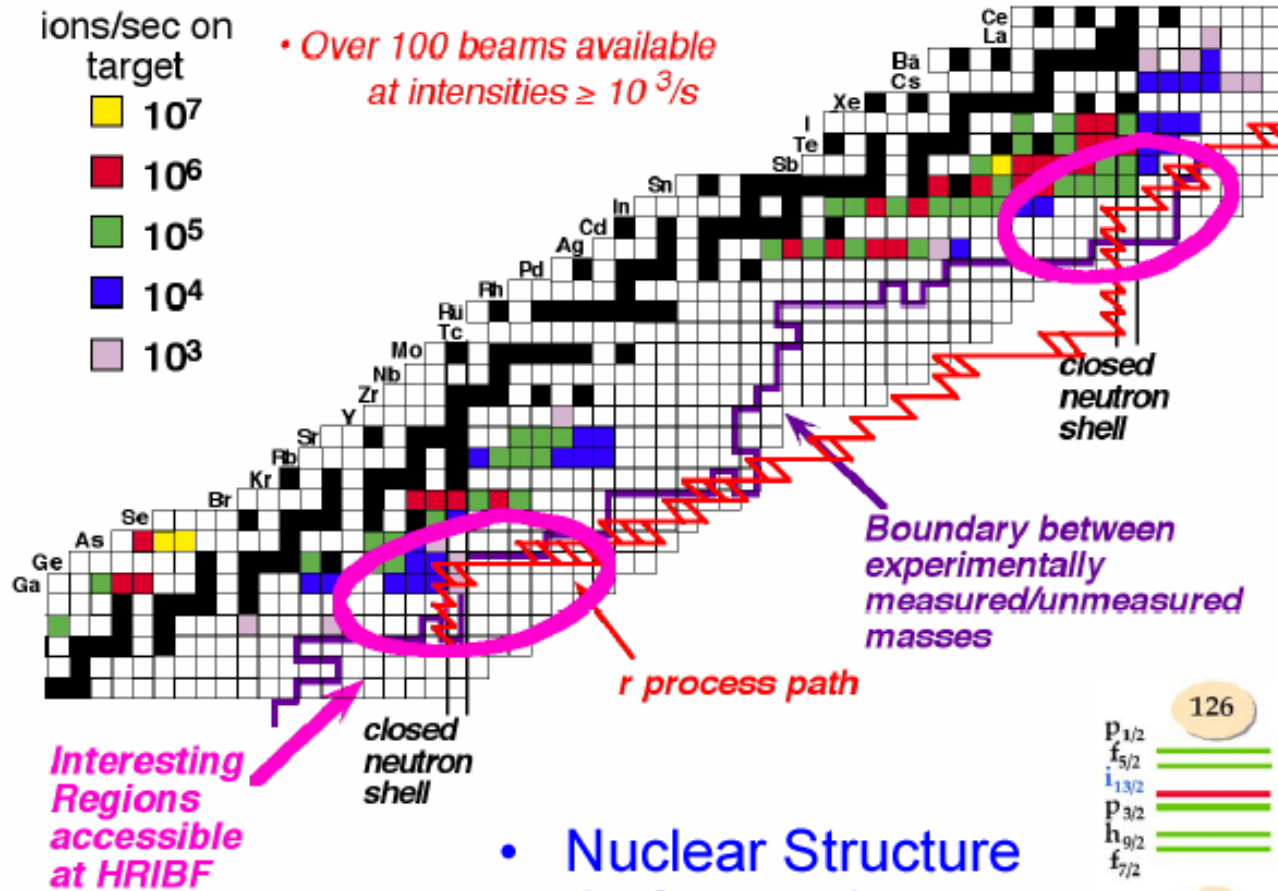
rp process : $T_9 > 0.5$



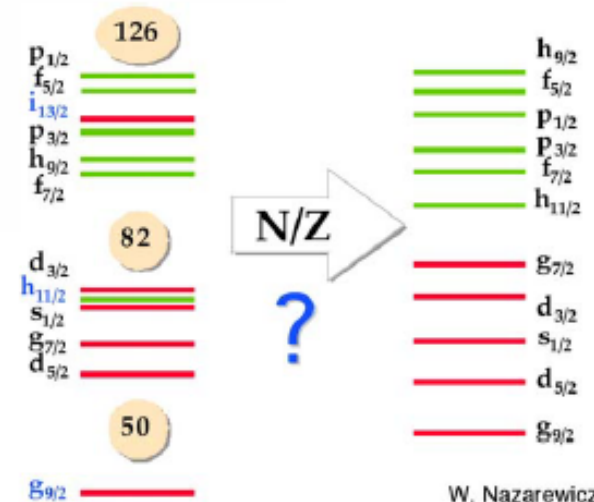
→ CNO cycle
→ HCNO cycle



Nuclear Structure Studies at Large Neutron Excess

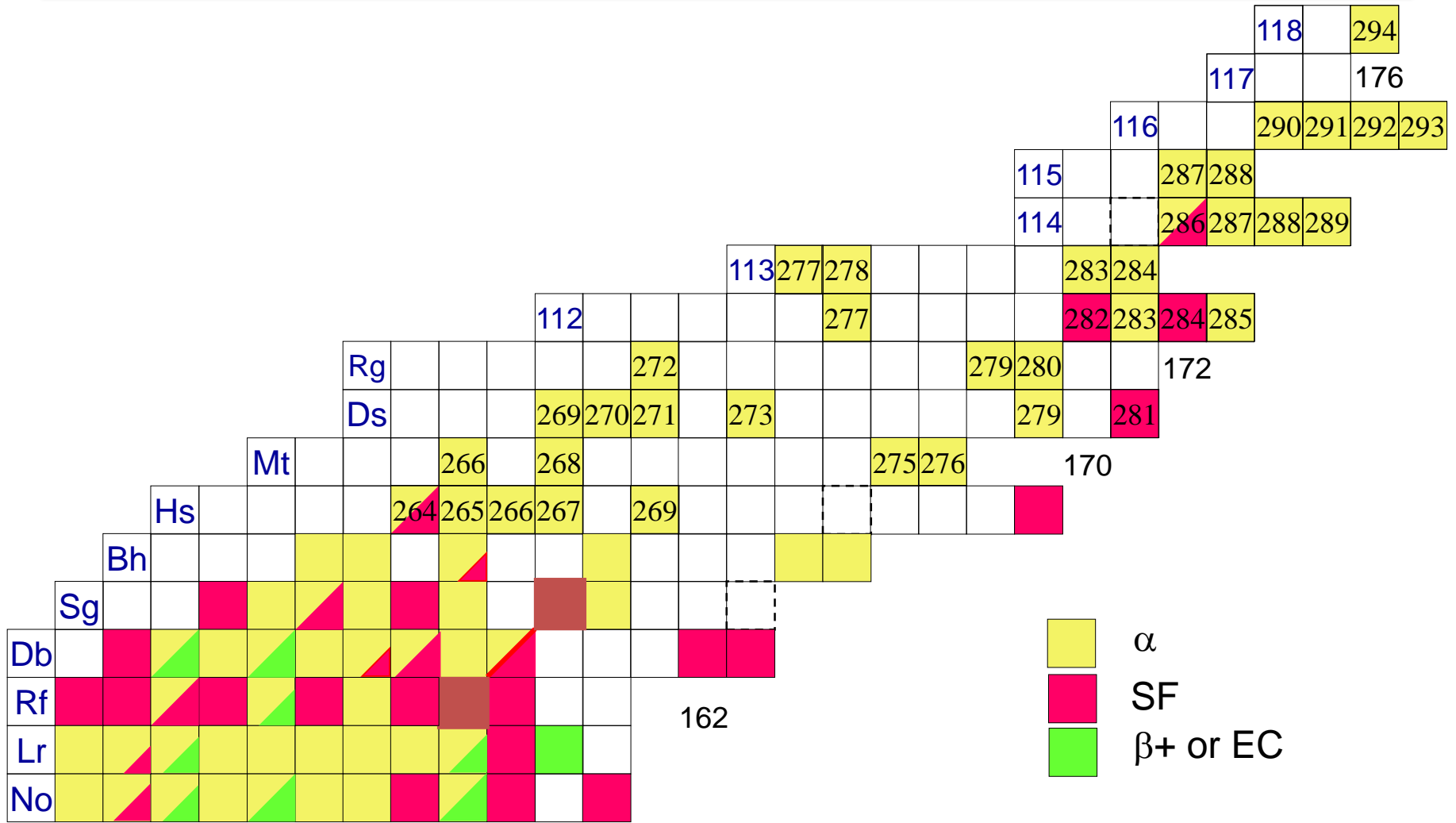


- Nuclear Structure in & near the r-process path



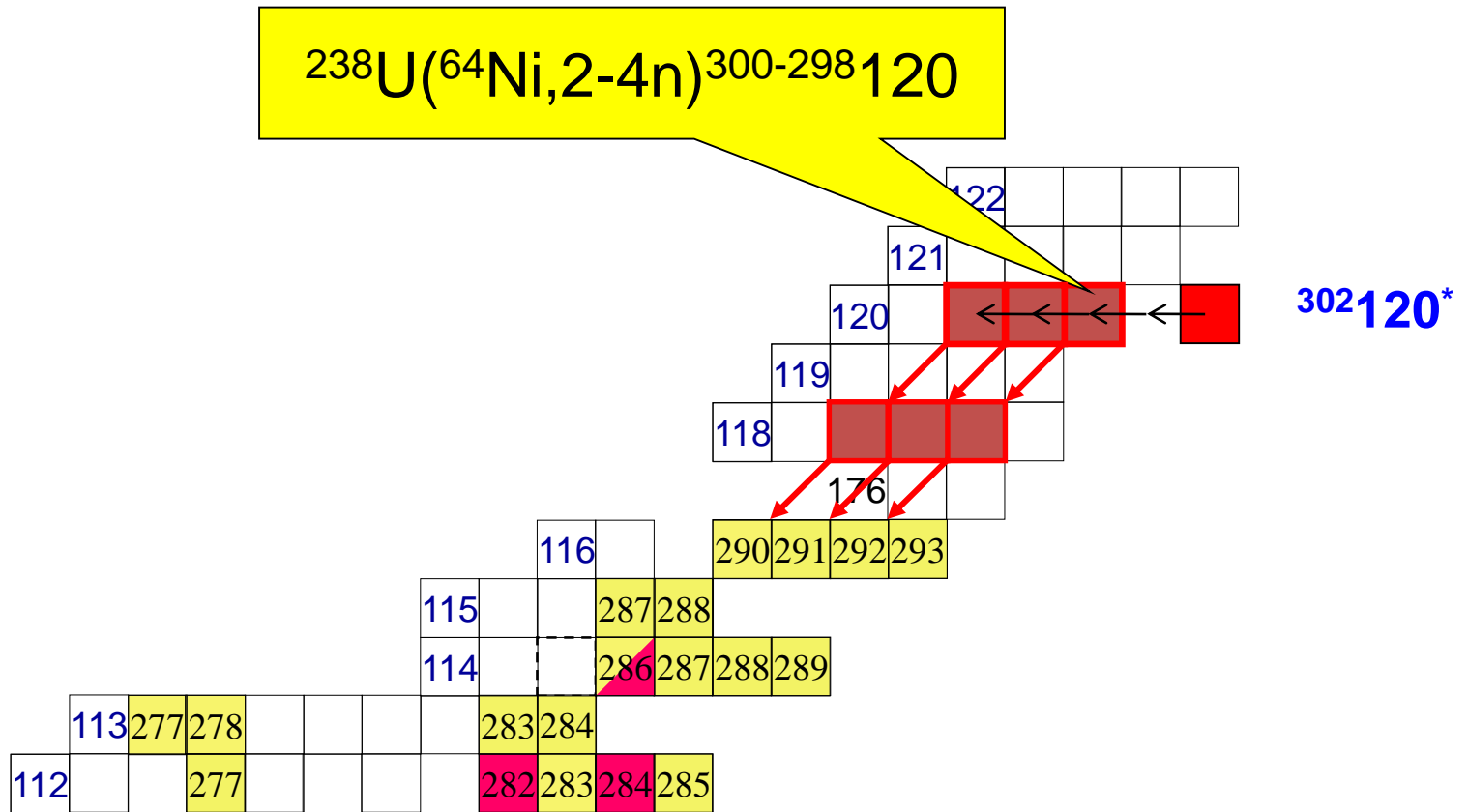
W. Nazarewicz

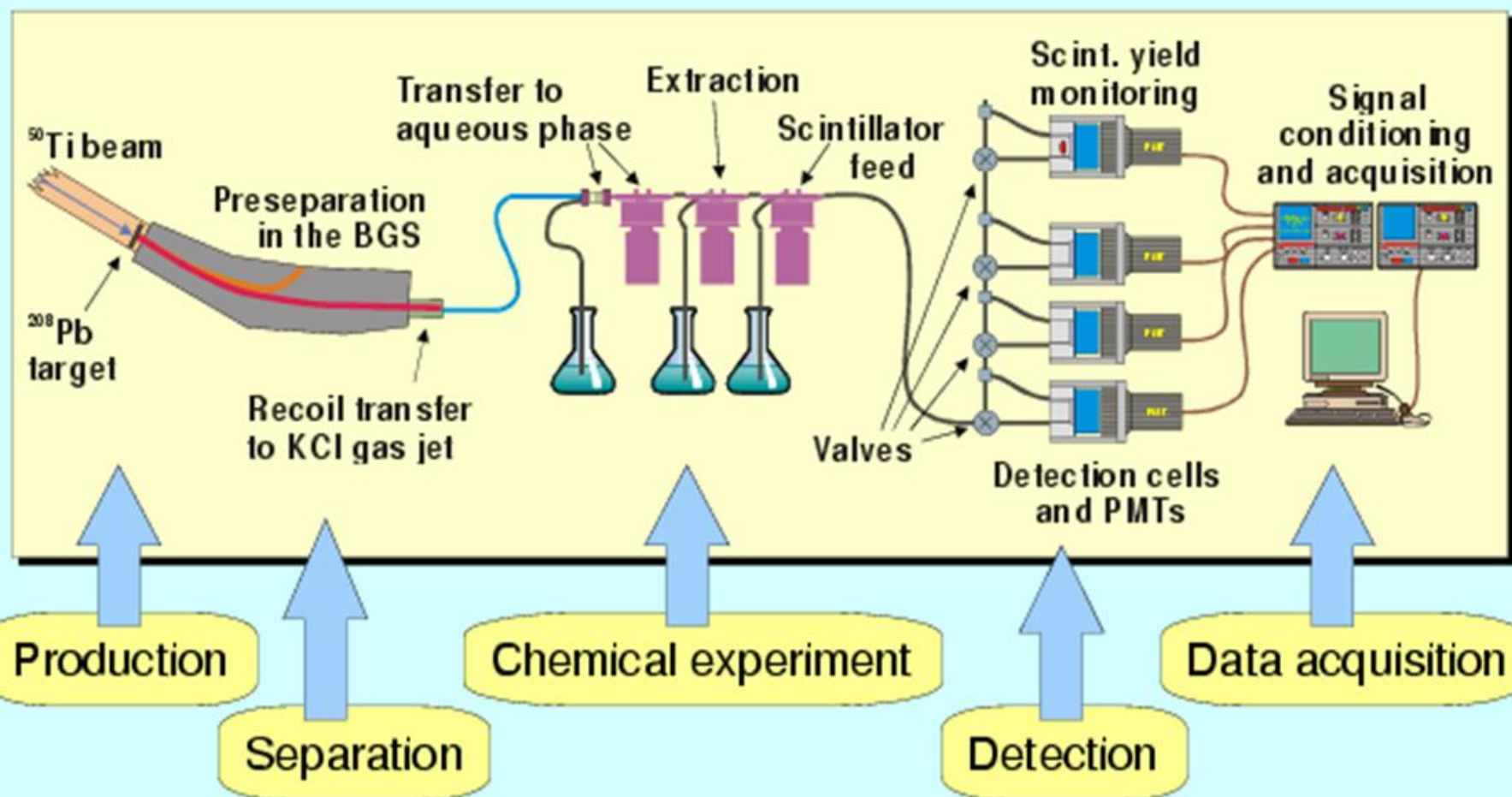
Superheavy Nuclei: Current Status



152

Reactions Tried at GSI in 2007-2008





Spin Structure of Unstable Nuclei

From W.Y. Kim, Kyungbook Nat. Univ.

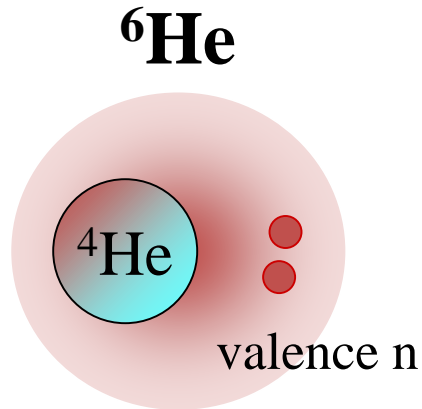
1. Spin-dependent interactions
 - Origin of fundamental properties of nuclei
 - Modification in neutron rich nuclei
2. Spin-orbit couplings and potentials
 - Localized at the nuclear surface

$$V_{LS} \sim \frac{1}{r} \frac{d\rho(r)}{dr}$$

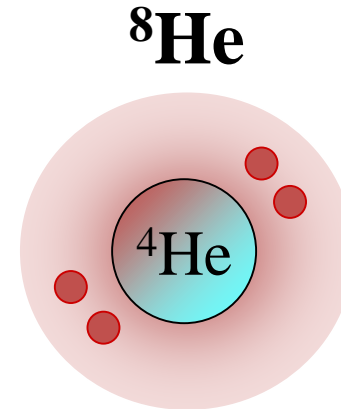
- Will be modified in neutron rich nuclei
 - Should be composed of two parts localized at different positions **if p and n have different $\rho(r)$**
 - Would have extended shape if n has an extended distribution in skin or halo nuclei
3. Need polarized p, d, and ^3He targets

Present Status at RIKEN

S.Sakaguchi Ph.D. Thesis, University of Tokyo (2008)



1. Di-neutron structure
 - Large recoil motion of α -core
 - Large charge radius (2.068 fm)
2. Two valence neutrons
 - Small matter radius (2.45 fm)



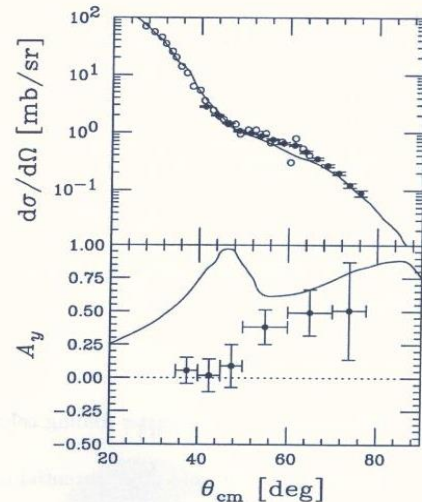
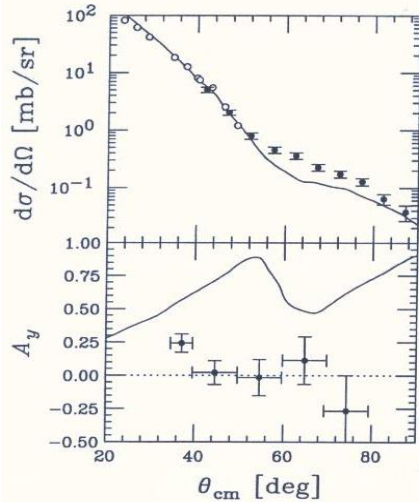
1. Isotropically distributed neutrons
 - Small recoil motion of α -core
 - Small charge radius (1.929 fm)
2. Four valence neutrons
 - Large matter radius (2.53 fm)

Present Status at RIKEN

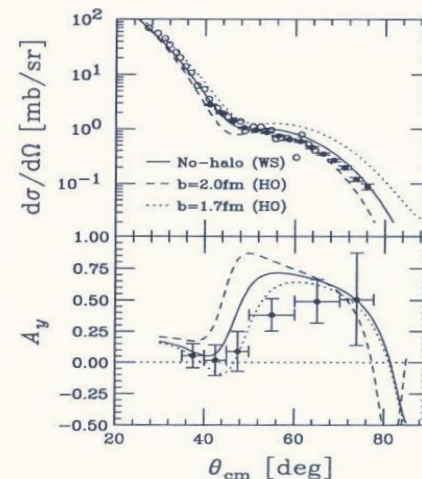
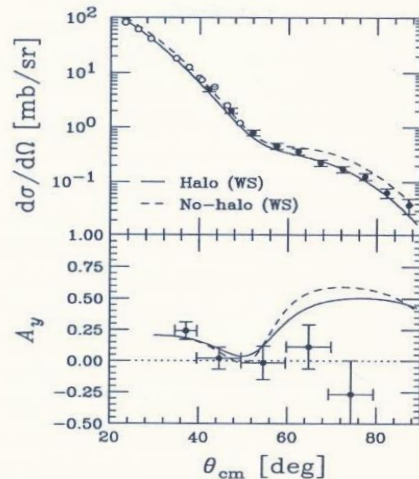
$\vec{p}-{}^6\text{He}$

$\vec{p}-{}^8\text{He}$

S.Sakaguchi (2008)



t-matrix folding calculation



Non-local g-matrix folding calculation

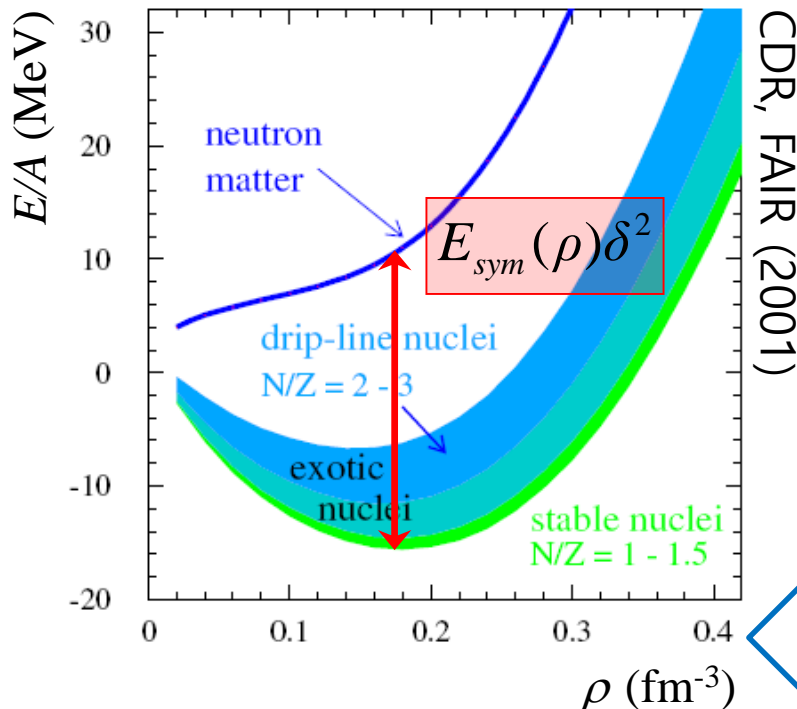
Nuclear Equation of State

$$E(\rho_n, \rho_p) = E(\rho_n = \rho_p) + E_{sym}(\rho)\delta^2 + O(\delta^4)$$

$$E_{sym}(\rho) = \frac{1}{2} \frac{\partial^2 E}{\partial \delta^2} \approx E(\rho)_{\text{pure neutron matter}} - E(\rho)_{\text{symmetric nuclear matter}}$$

with $\rho = \rho_n + \rho_p$, $\delta = (\rho_n - \rho_p) / \rho = (N - Z) / A$

B.-A. Li, L.-W. Chen
& C.M. Ko
Physics Report,
464, 113 (2008)



$$E/A(\rho_n = \rho_p)$$

Symmetric
nuclear matter
($\rho_n = \rho_p$)

Isospin
asymmetry
 δ

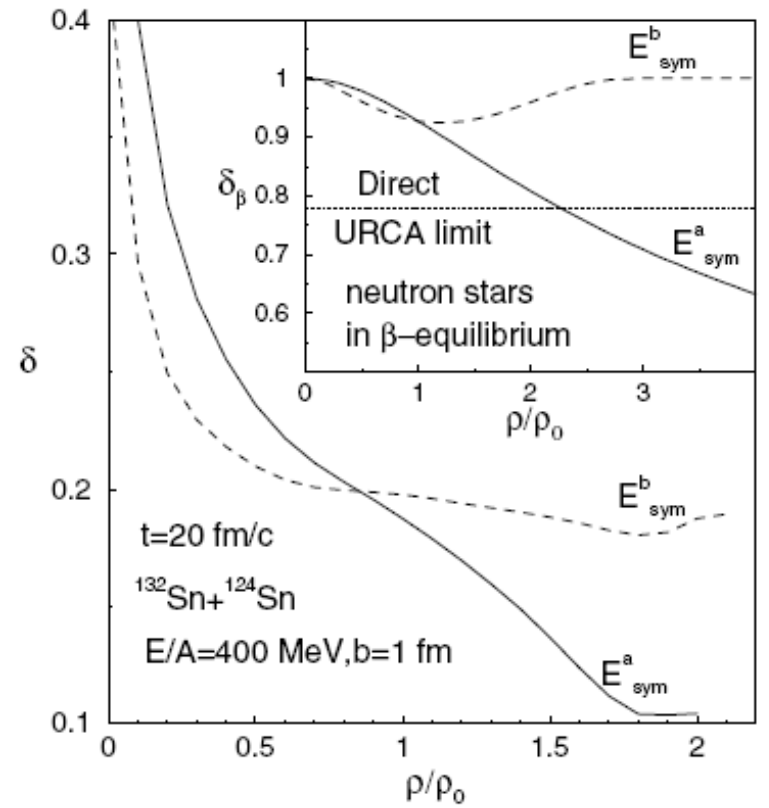
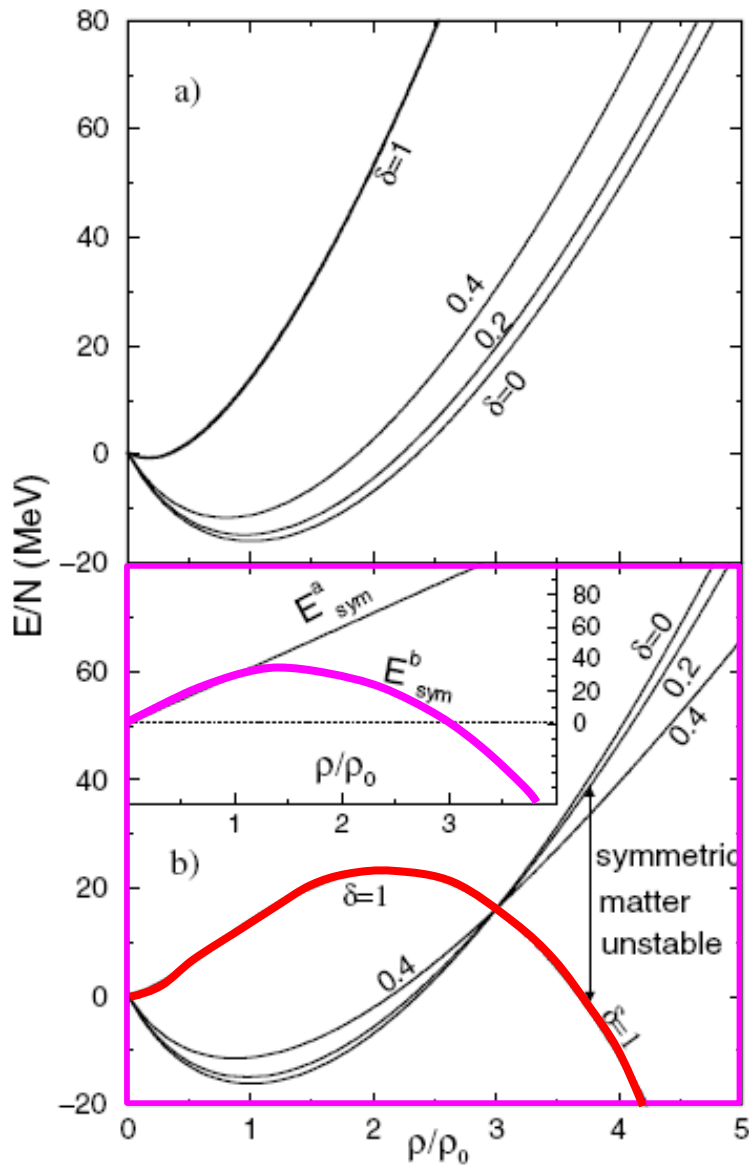
Nucleon
density
 ρ

F. de Jong & H. Lenske, RPC 57, 3099 (1998)

F. Hofman, C.M. Keil & H. Lenske, PRC 64, 034314 (2001)

Nuclear Equation of State

Bao-An Li, PRL 88, 192701 (2002)



High (Low) density matter is more neutron rich with **soft (stiff)** symmetry energy

Importance of Symmetry Energy

RIB can provide crucial input.

Effective field theory, QCD

π/π^+
 K^+/K^0
 n/p +
 ${}^3\text{H}/{}^3\text{He}$
 γ

isodiffusion
 isotransport
 + isocorrelation
 isofractionation
 isoscaling

Isospin Dependence of Strong Interactions

Nuclear Masses
 Neutron Skin Thickness
 Isovector Giant Dipole Resonances
 Fission

Heavy Ion Flows
 Multi-Fragmentation
 Nuclei Far from Stability
 Rare Isotope Beams

Many-Body Theory
 Symmetry Energy
 (Magnitude and Density Dependence)

Supernovae
 Weak Interactions
 Early Rise of $L_{\nu e}$
 Bounce Dynamics
 Binding Energy

Proto-Neutron Stars
 ν Opacities
 ν Emissivities
 SN r-Process
 Metastability

Neutron Stars
 Observational
 Properties

Binary Mergers
 Decompression/Ejection
 of Neutron-Star Matter
 r-Process

QPO's
 Mass
 Radius

NS Cooling
 Temperature
 R_{∞}, z
 Direct Urca
 Superfluid Gaps

X-ray Bursters
 R_{∞}, z

Gravity Waves
 Mass/Radius
 dR/dM

Pulsars
 Masses
 Spin Rates
 Moments of Inertia
 Magnetic Fields
 Glitches - Crust

Maximum Mass, Radius
 Composition:
 Hyperons, Deconfined Quarks
 Kaon/Pion Condensates

■ A.W. Steiner, M. Prakash, J.M. Lattimer and P.J. Ellis, Physics Report 411, 325 (2005)

■ Red boxes: added by B.-A. Li

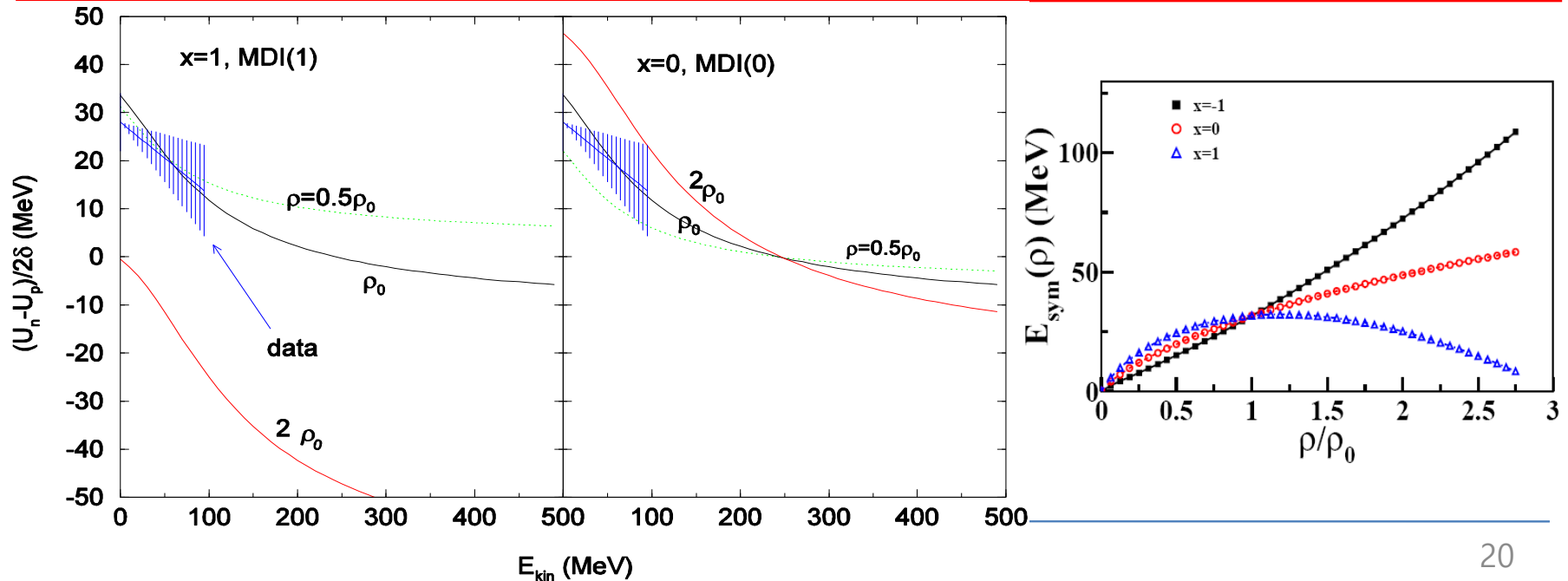
Uncertainty in E_{sym} at high ρ

$$E_{sym} = E_{sym}^{kin} + E_{sym}^{pot1} + E_{sym}^{pot2} = \frac{1}{3}t(k_F) + \frac{1}{6} \frac{\partial U_0(k)}{\partial k} \Big|_{k_F} k_F + \frac{3}{2k_F^3} \int_0^{2^{1/3}k_F} U_{sym}(k) k^2 dk$$

Kinetic **Isoscalar** **Isovector**

$$U_0 = \frac{1}{2}(U_n + U_p) = \frac{1}{4}(3u_{T1} + u_{T0}): \text{Isoscalar particle potential}$$

$$U_{sym} = \frac{1}{2\delta}(U_n - U_p) = \frac{1}{4}(u_{T1} - u_{T0}) \left\{ \begin{array}{l} \text{Isovector potential: Isospin dep. SI in medium} \\ \text{We know very little about this term at high } \rho! \end{array} \right.$$

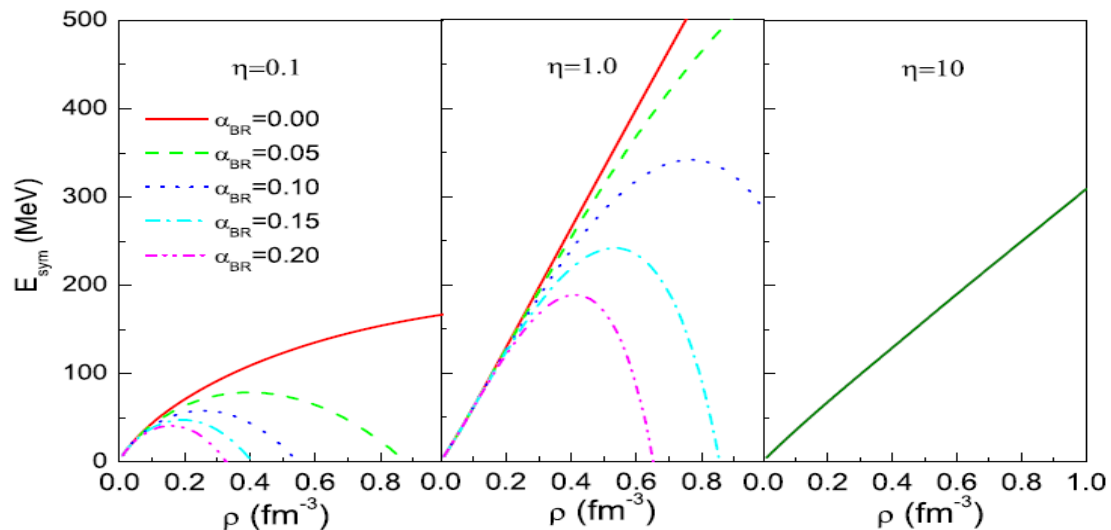


Possible Effects on E_{sym}

Brown-Rho Scaling

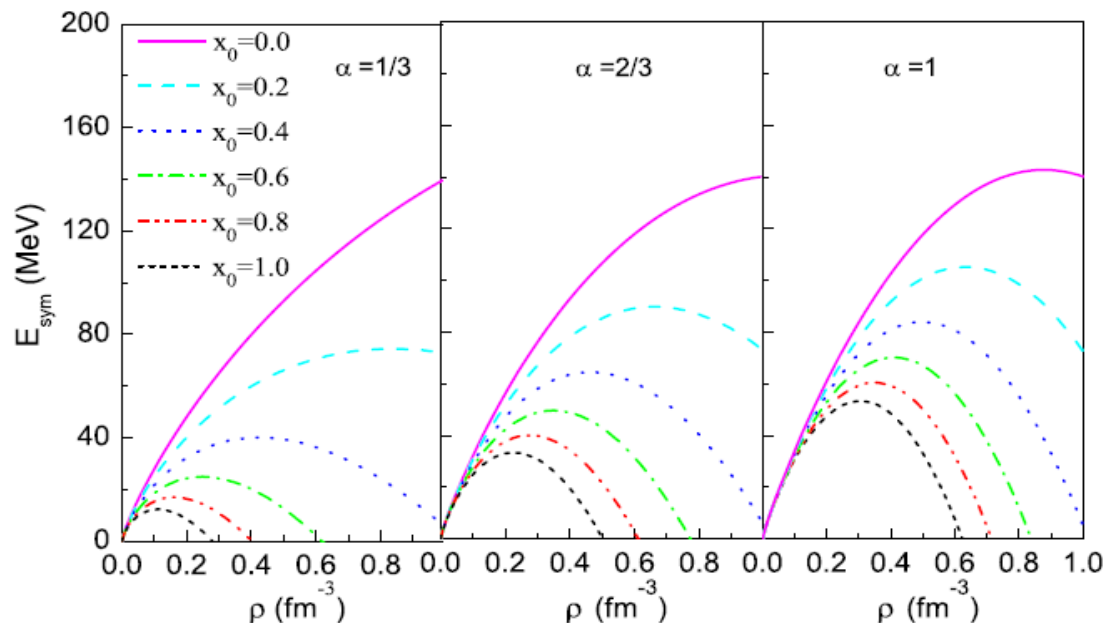
$$\frac{m_{\rho}^*}{m_{\rho}} = 1 - \alpha_{BR} \frac{\rho}{\rho_0}$$

G.E. Brown and M. Rho,
PRL 66, 2720 (1991);
Phys. Rep. 396, 1 (2004)



3-Body Force

$$E_{sym}^{pot2} = -F(\rho) - (1 + 2x_0) \frac{t_0}{8} \rho^{\alpha+1}$$



Is NS Stable with a Super Soft E_{sym} ?

If the symmetry energy is too soft, then a mechanical instability will occur when $dP/d\rho < 0$, neutron stars will, then, collapse while they exist in nature.

Gravity



Nuclear pressure

For npe matter,

$$P(\rho, \delta) = P_0(\rho) + P_{asy}(\rho, \delta) = \rho^2 \left(\frac{\partial E}{\partial \rho} \right)_\delta + \frac{1}{4} \rho_e \mu_e$$

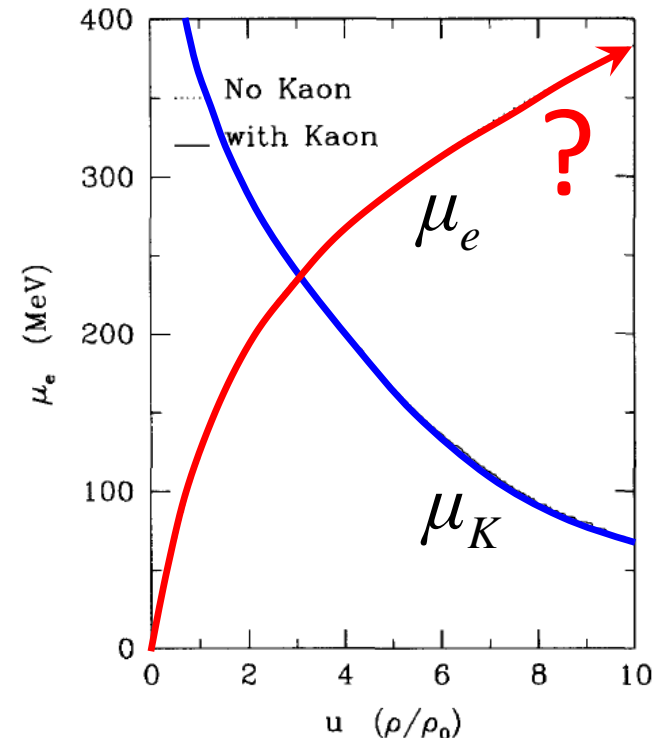
$$= \rho^2 \left[E'(\rho, \delta = 0) + E'_{sym}(\rho) \delta^2 \right] + \frac{1}{2} \delta(1 - \delta) \rho E_{sym}(\rho)$$

$dP/d\rho < 0$, if E'_{sym} is big and negative (super-soft)

$$\frac{dP}{dr} = -(\varepsilon + P) \frac{m_g + 4\pi r^3 P}{r(r - 2m_g)}$$

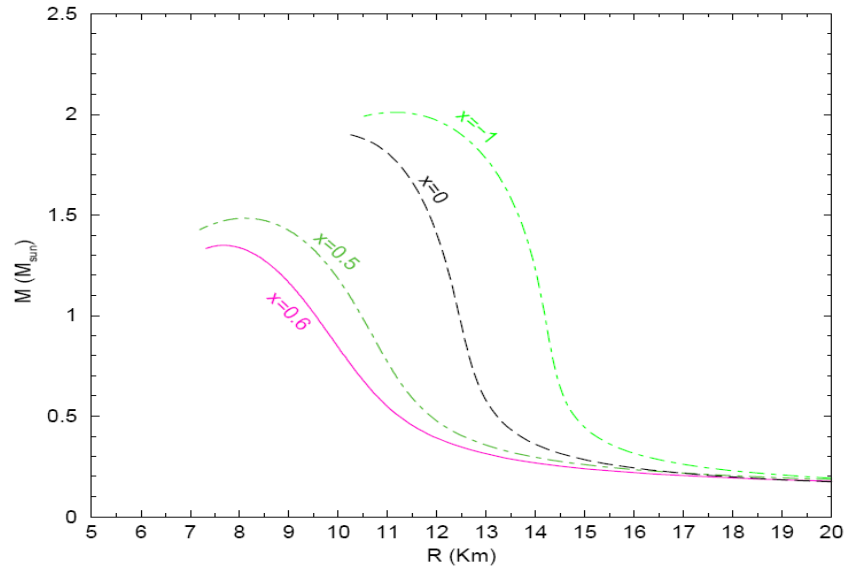
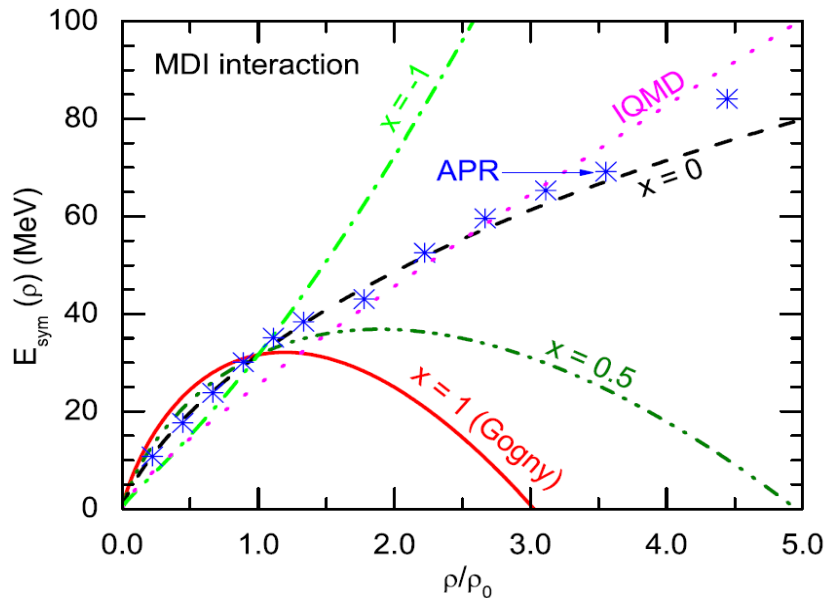
TOV equation: a condition at hydrodynamical equilibrium

$\mu_e(\rho)$ is critical for kaon condensation



G.Q. Li, C.-H. Lee & G.E. Brown
Nucl. Phys. A 625, 372 (1997)

Astrophysical Implication

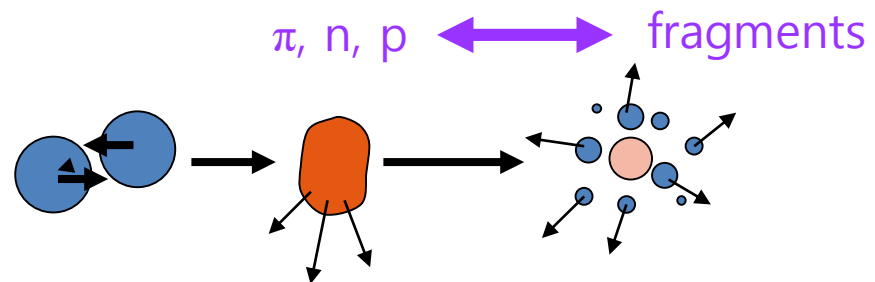
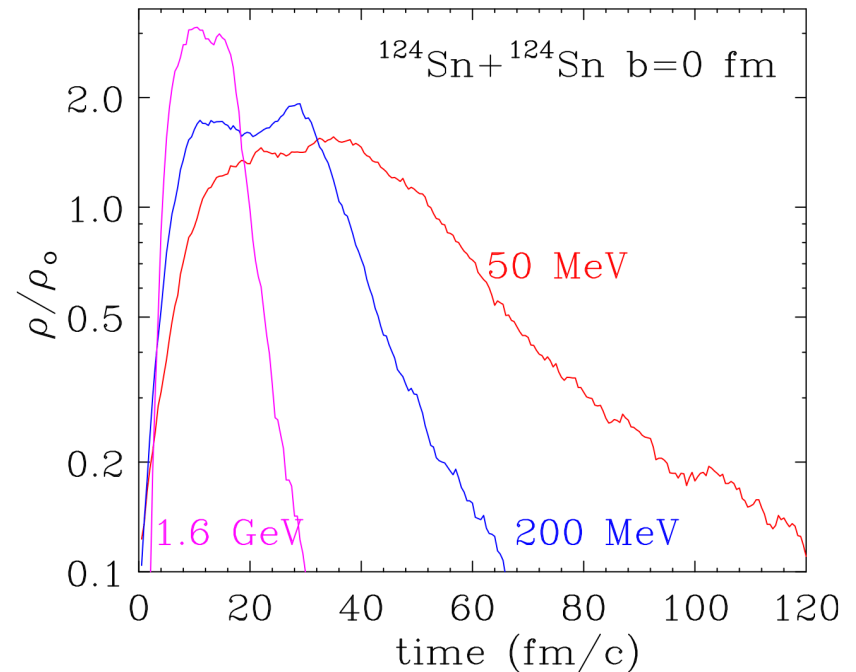


$K_0=211$ MeV is used for this calculation; higher incompressibility for symmetric matter will lead to higher masses, systematically.

The softest symmetry energy that the TOV is still stable is $x = 0.93$, giving $M_{\text{max}} = 0.11 M_{\odot}$ and $R \geq 28$ km.

Experimental Principles

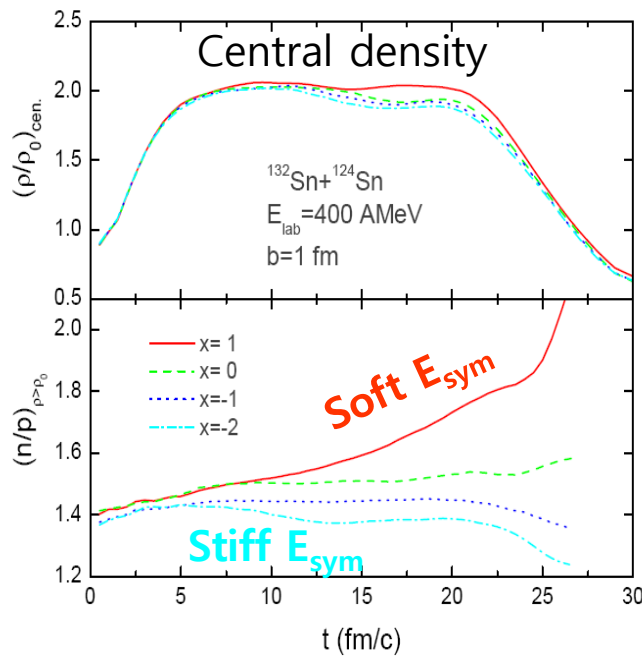
- Range of density in HIC by
 - incident energy
 - impact parameter
- Types of particles formed
 - emission time & density
 - n & p are emitted throughout
 - Fragments (Z=3-20) at sub-saturation densities
- Change N/Z of nuclei
 - larger N/Z ratio is preferable



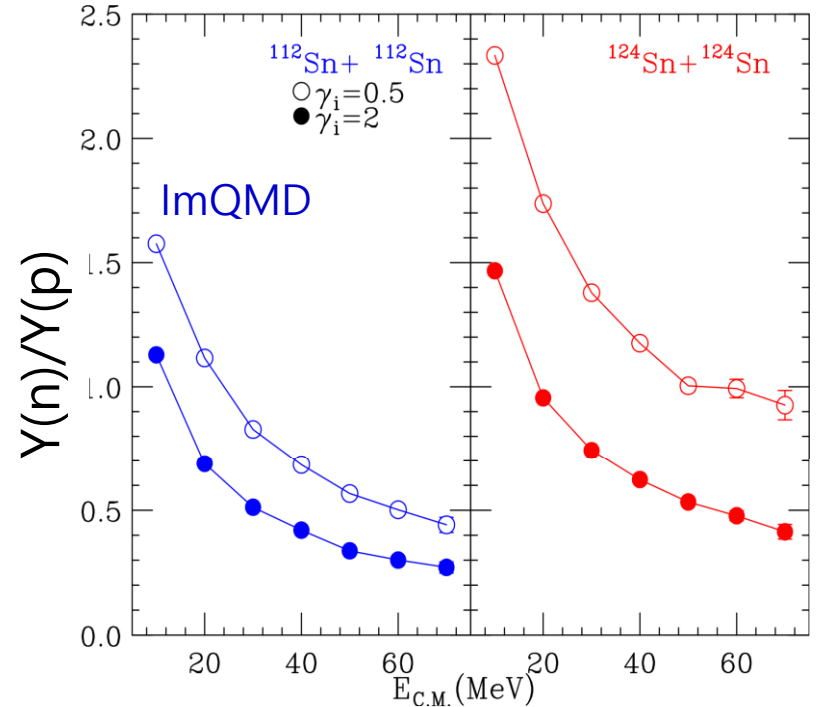
Experimental Observables

- Signals at sub-saturation densities
 - 1) Sizes of n-skins for unstable nuclei
 - 2) n/p ratio of fast, pre-equilibrium nucleons
 - 3) ${}^3\text{H}/{}^3\text{He}$ ratio
 - 4) Isospin fractionation and isoscaling in nuclear multifragmentation
 - 5) Isospin diffusion (transport)
 - 6) Differential collective flows (v_1 & v_2) of n and p
 - 7) Correlation function of n and p
- Signals at supra-saturation densities
 - 1) π^-/π^+ ratio
 - 2) K^+/K^0 ratio
 - 3) Differential collective flows (v_1 & v_2) of n and p
 - 4) Azimuthal angle dependence of n/p ratio with respect to the R.P.
- Correlation of various observables
- Simultaneous measurement of neutrons and charged particles

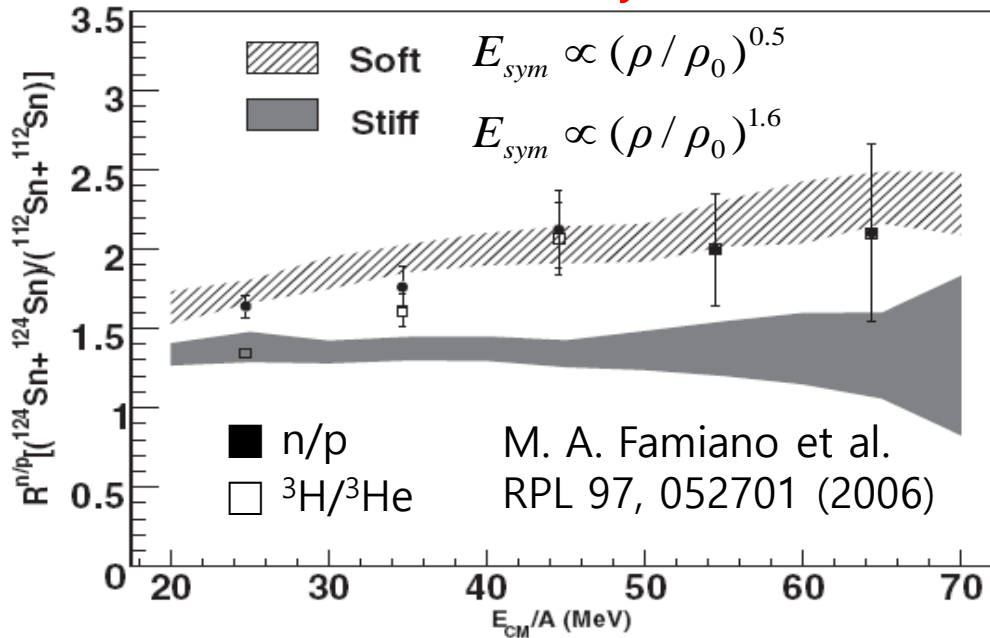
Yield Ratio



$$E_{\text{sym}}(\rho) = 12.7(\rho/\rho_0)^{2/3} + 17.6(\rho/\rho_0)^{\gamma_i}$$



Double ratio: min. systematic error



More neutrons are emitted from the n-rich system and softer symmetry energies.

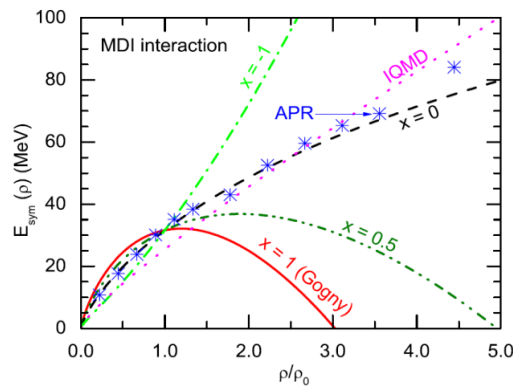
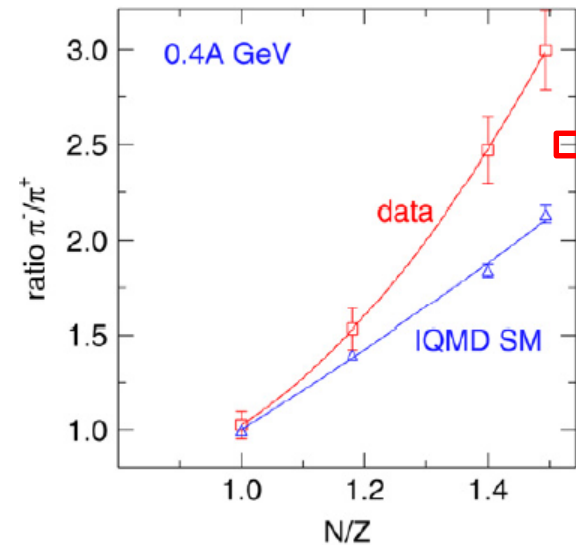
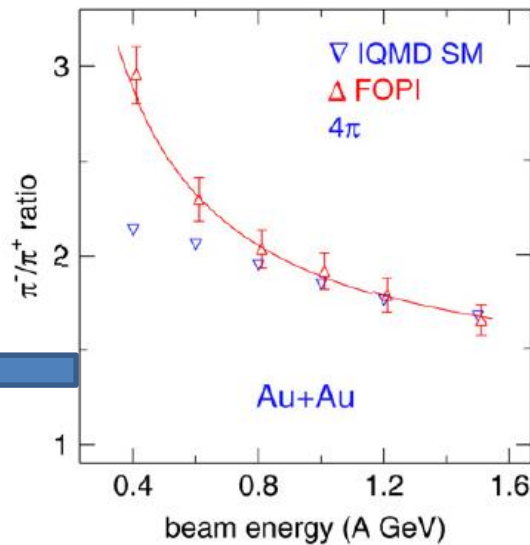
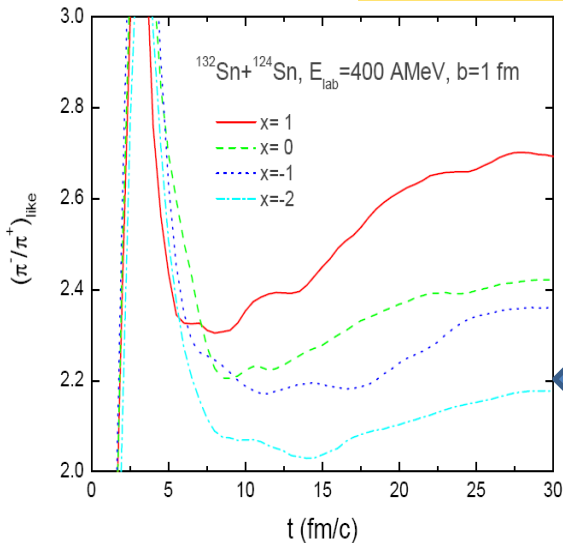
Yield Ratio (π^-/π^+)

Data: FOPI Collaboration, Nucl. Phys. A 781, 459 (2007)

IQMD: Eur. Phys. J. A 1, 151 (1998)

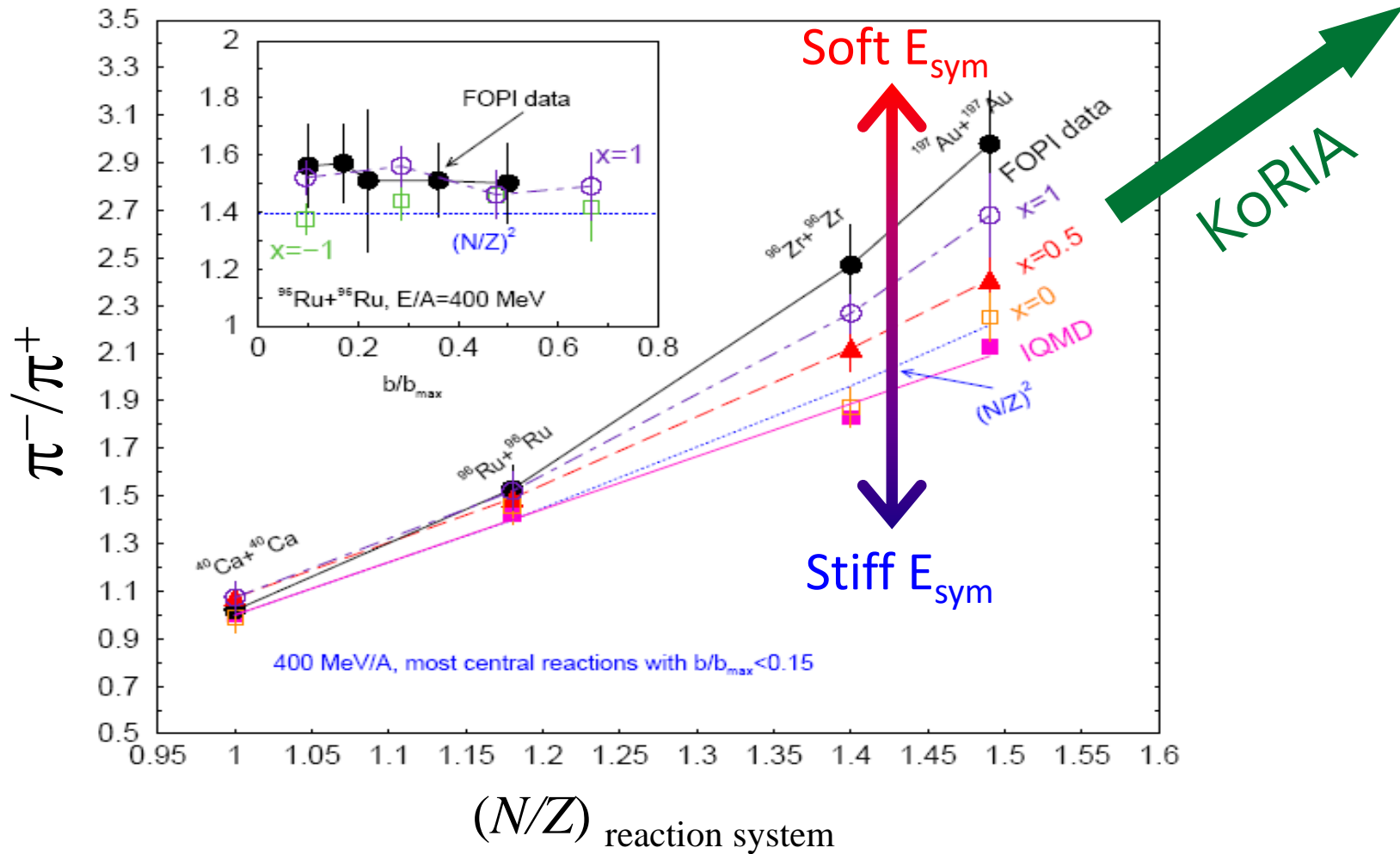


$$\text{corresponding to } E_{\text{sym}}(\rho) = \frac{100}{8} \frac{\rho}{\rho_0} + (2^{2/3} - 1) \frac{3}{5} E_F^0 \left(\frac{\rho}{\rho_0}\right)^{2/3}$$



Need a symmetry energy softer than the above to make the pion production region more neutron-rich!

π^-/π^+ Ratio



Isospin Diffusion Parameter

Isospin diffusion occurs only in asymmetric systems $A+B$

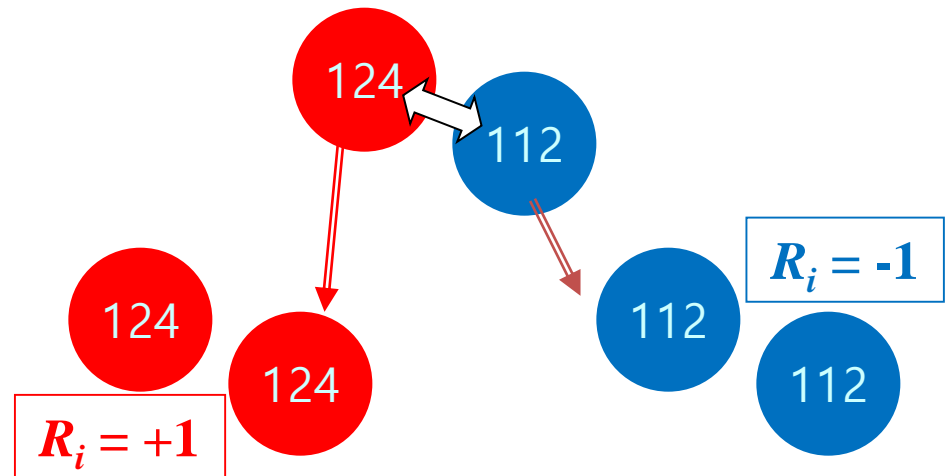
No isospin diffusion between symmetric systems

Non-isospin diffusion effects are the same for A in $A+B$ & $A+A$ and also for B in $B+A$ & $B+B$

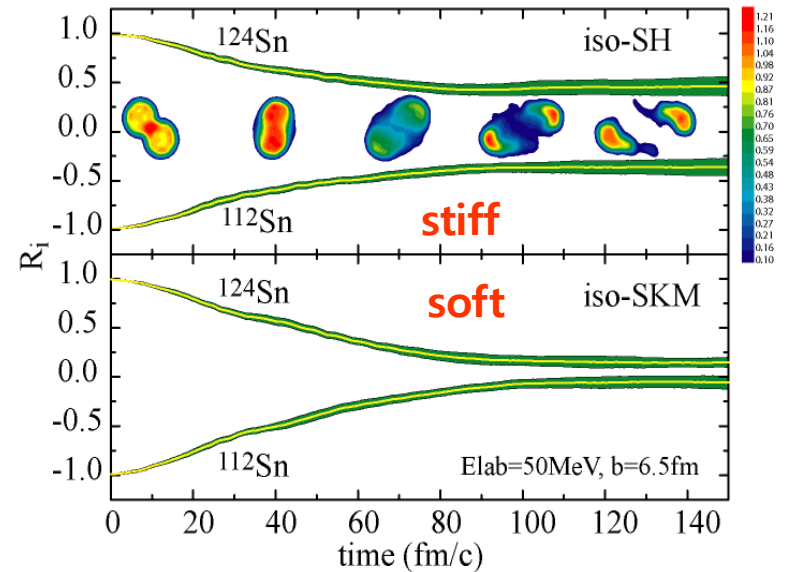
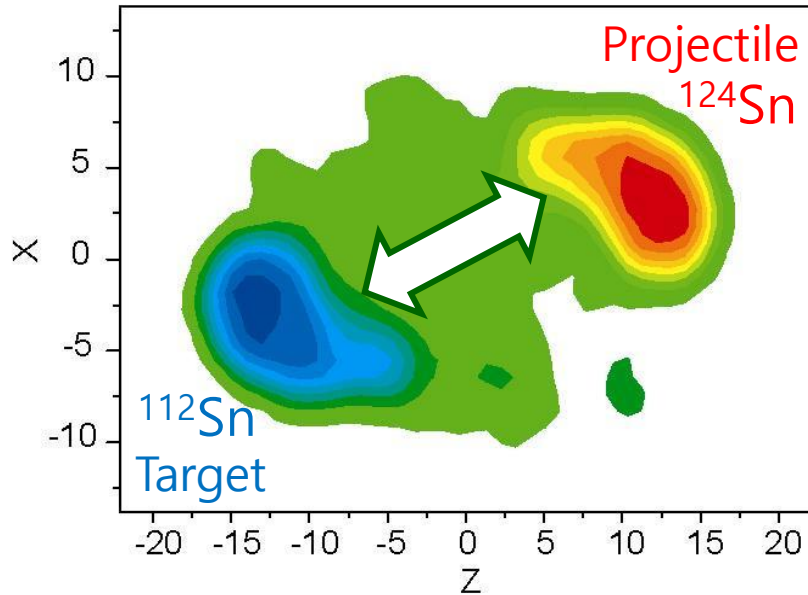
$$R_i = 2 \frac{x_{AB} - (x_{AA} + x_{BB})/2}{x_{AA} - x_{BB}}$$

F. Rami et al., FOPI, PRL 84, 1120 (2000)
B. Hong et al., FOPI, PRC 66, 034901 (2002)
Y.-J. Kim & B. Hong, in preparation

$R_i = 0$ for complete mixing



Isospin Diffusion Parameter



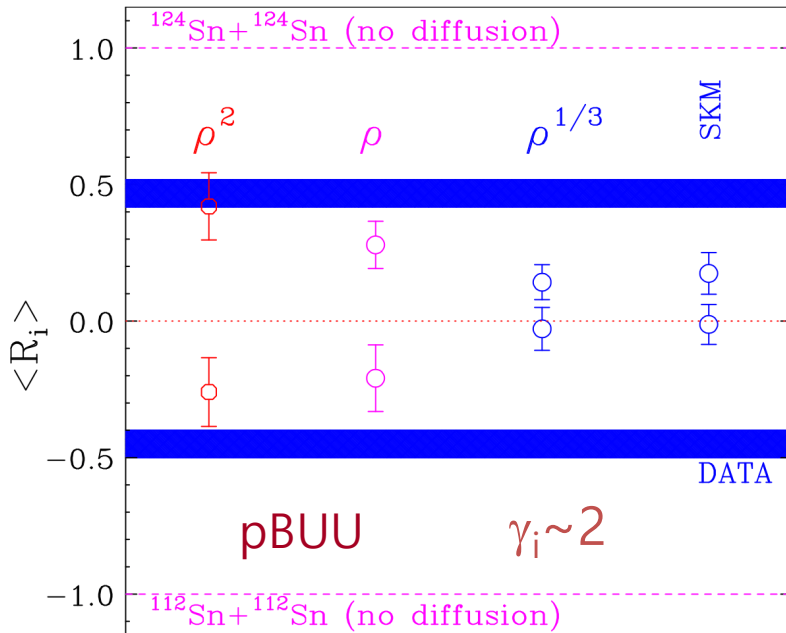
- Symmetry energy drives system towards equilibrium
 - stiff EOS : small diffusion ($|R_i| \gg 0$)
 - soft EOS : large diffusion & fast equilibrium ($R_i \rightarrow 0$)

M.B. Tsang et al., PRL 92, 062701 (2004)

Isospin Diffusion Parameter

M.B. Tsang et al., PRL 92, 062701 (2004)

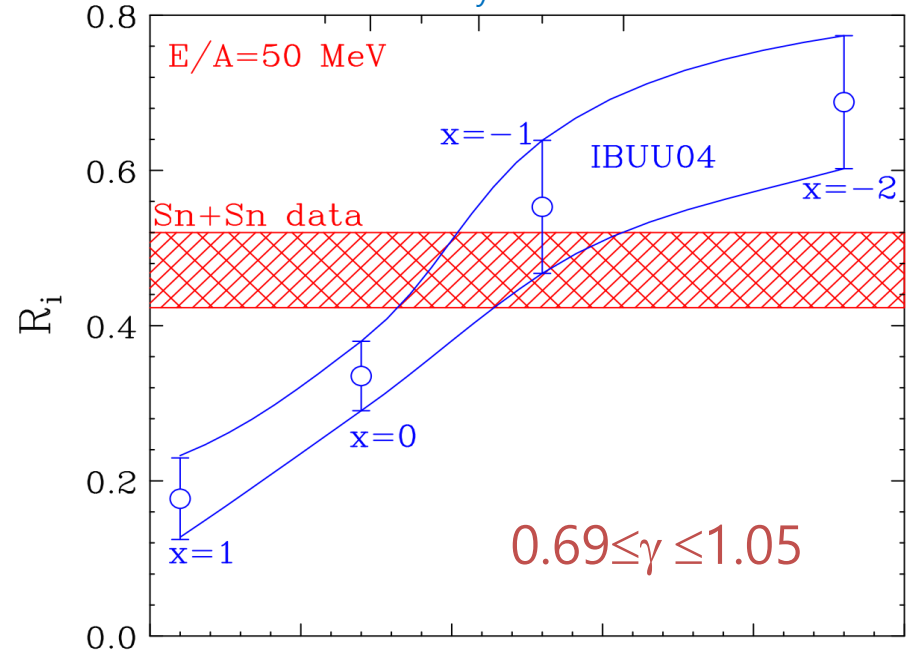
$$E_{\text{sym}}(\rho) = 12.7(\rho/\rho_0)^{2/3} + 12.5(\rho/\rho_0)^{\gamma_i}$$



stiff \longleftrightarrow soft

L.W. Chen et al., PRL 94, 032701 (2005)

$$\text{IBUU04: } E_{\text{sym}}(\rho) \sim 31.6(\rho/\rho_0)^\gamma$$



stiff \longleftrightarrow soft

Observable in HIC is sensitive to the ρ dependence of E_{sym} and should provide constraints to the symmetry energy.

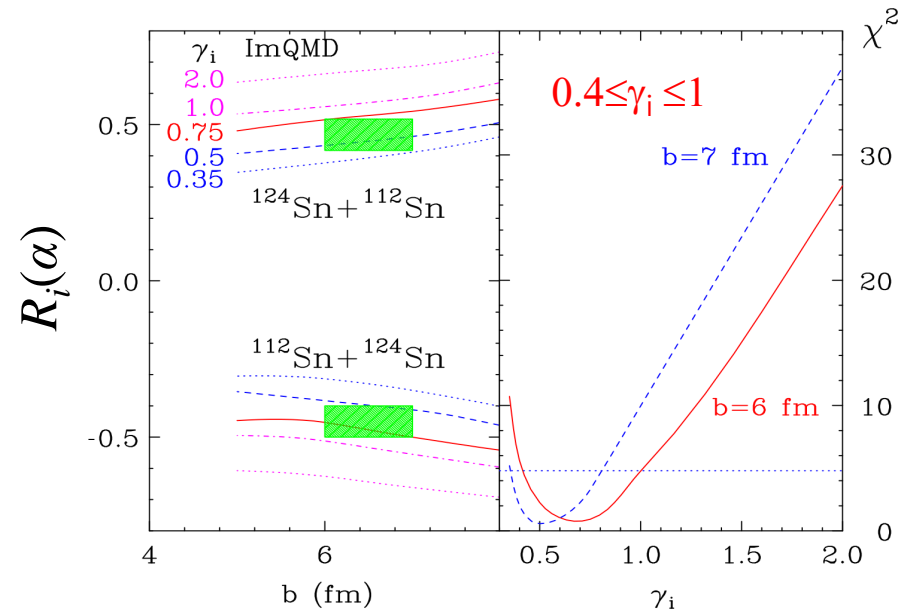
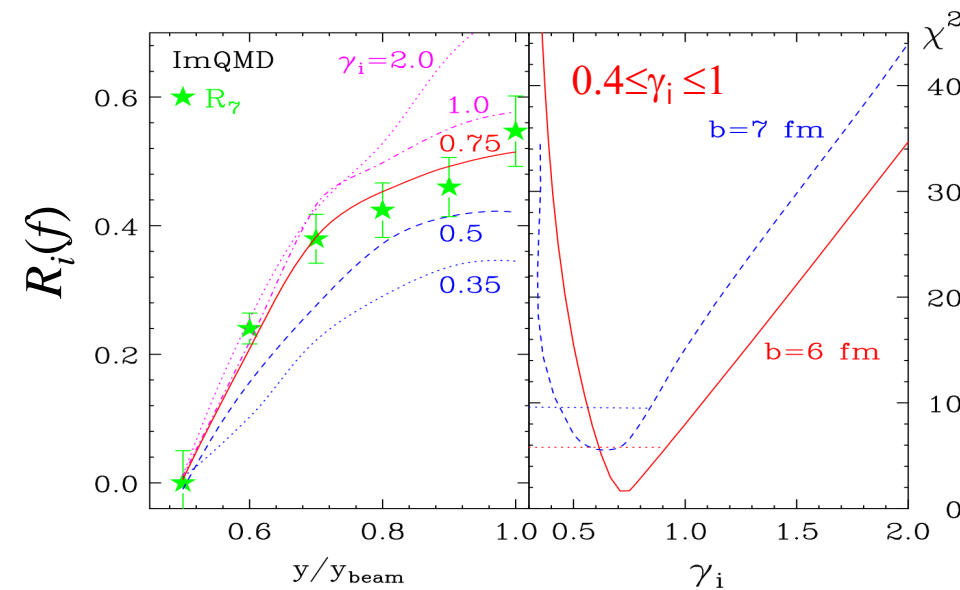
Isospin Diffusion Parameter

NSCL/MSU Data at Low Energy

$$E_{\text{sym}}(\rho) = 12.5(\rho/\rho_0)^{2/3} + 17.6(\rho/\rho_0)^{\gamma_i}$$

$$f = \frac{Y_{124}(^7\text{Li}) / Y_{124}(^7\text{Be})}{Y_{112}(^7\text{Li}) / Y_{112}(^7\text{Be})}$$

$$f = \frac{Y_{124+124}(Z = 3 \sim 8)}{Y_{112+112}(Z = 3 \sim 8)} \propto \exp(\alpha N)$$



Collective Flow

B.-A. Li,
PRL 85, 4221
(2000)

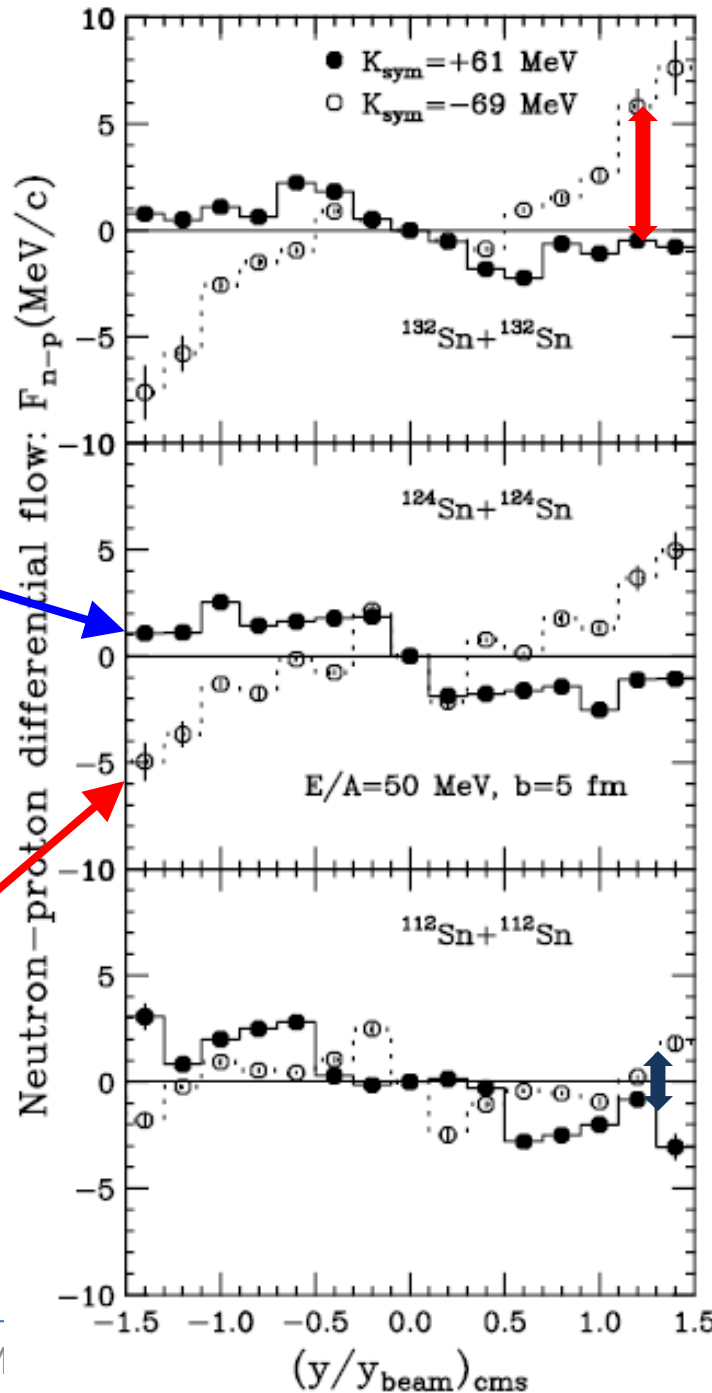
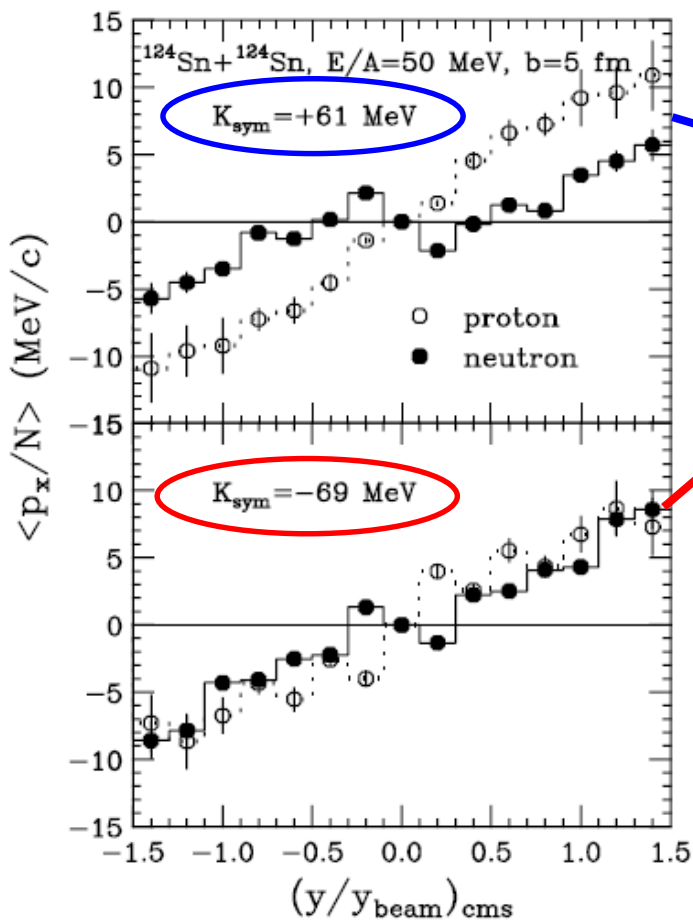
$$K_{sym} \equiv 9\rho_0^2 \left. \frac{\partial^2 E_{sym}(\rho)}{\partial \rho^2} \right|_{\rho=\rho_0}$$

Stiff



Also known as ν_1

Super
Soft

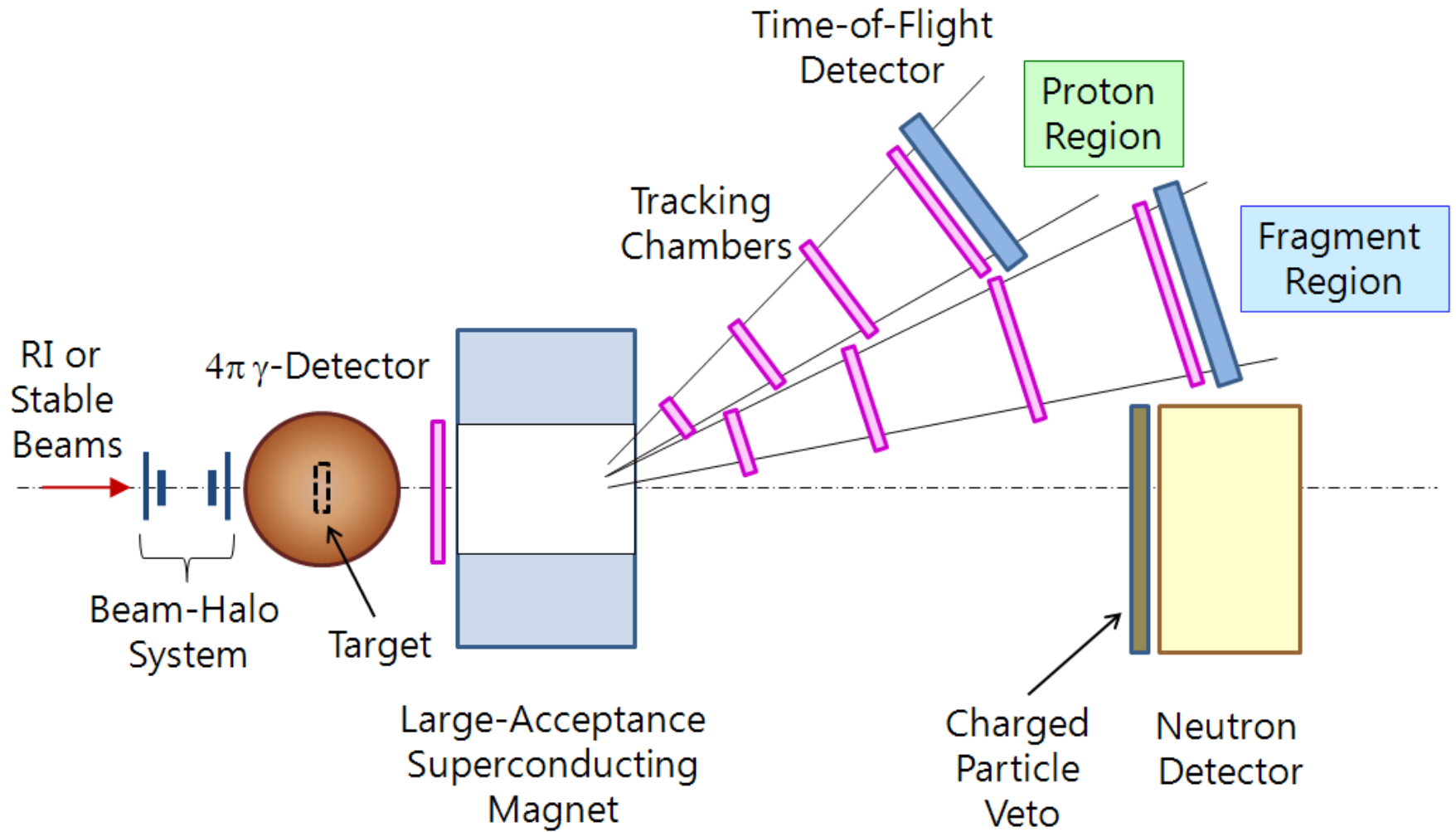


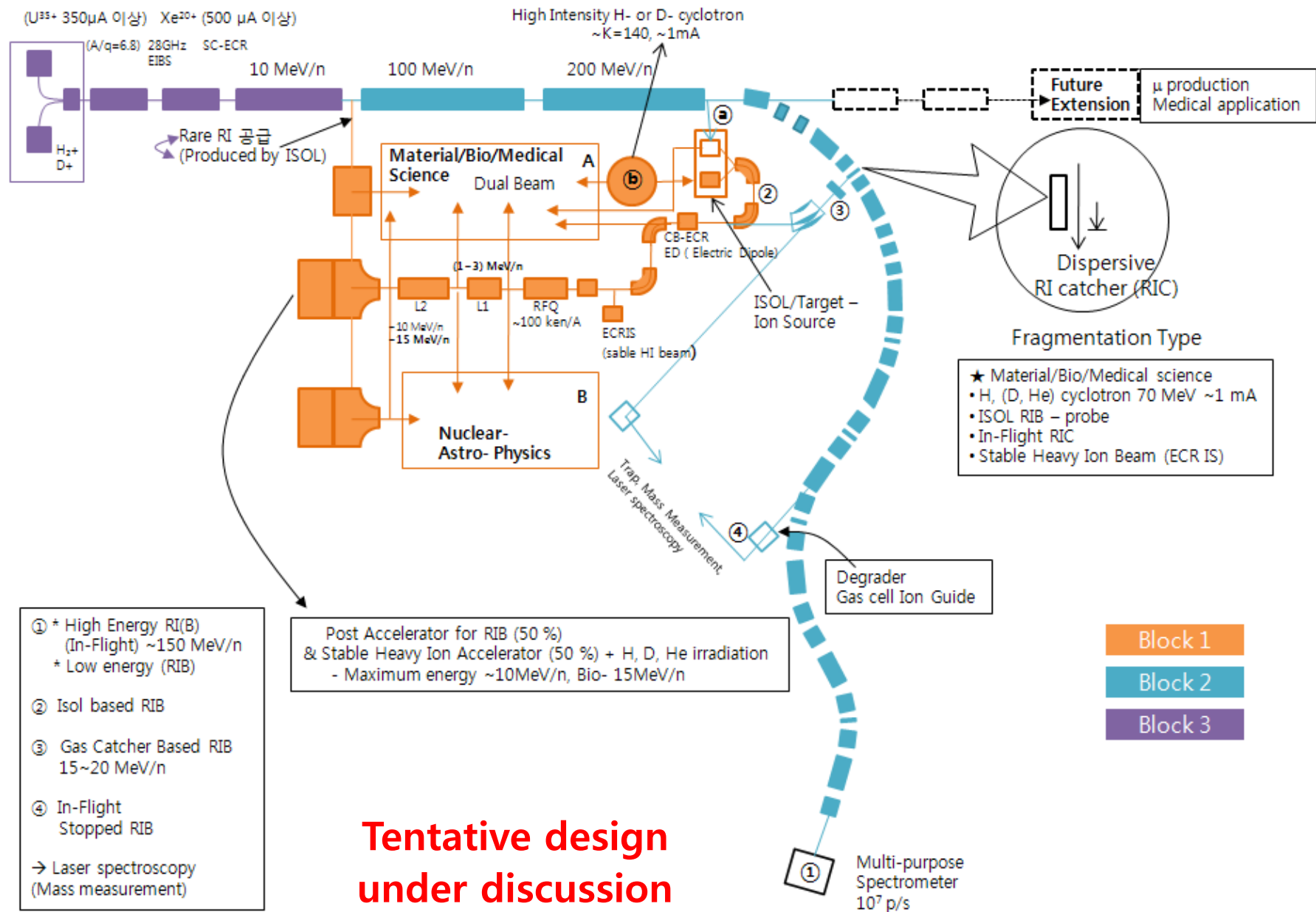
Large
N/Z



Small
N/Z

Multi-Purpose Detector





Summary

1. Rich physics with RI beams

- Neutron & proton drip lines
- Neutron halo & skin structures
- Nuclear Astrophysics
- Super-heavy elements
- Fundamental symmetries
- Nuclear symmetry energy
 - Long-standing problem in nuclear physics
 - Crucial to understand the neutron matter
 - Crucial to understand the astrophysical objects

2. KoRIA

- First large scale accelerator for basic science in Korea
- We need your help & participation!