

# Symmetry energy of dense matter in holographic QCD

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CQUeST

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Based on

arXiv:1011.0868 : Youngman Kim, YS, Ik Jae Shin, Sang-Jin Sin

JHEP 0804:010, 2008: YS, Sang-Jin Sin

JHEP 1003:074, 2010: Youngman Kim, YS, Sang-Jin Sin

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- One of main research project in KoRIA(Korea Rare Isotope Accelerator).
- QCD phenomena at finite density is hard to understand.
- AdS/CFT may give clues to understand non-perturbative phenomema.

- String theory

Open Strings

||

Closed Strings

- String theory

Open Strings

massless excitation

Gauge Field  $A_\mu$



Closed Strings

massless excitation

Graviton  $G_{\mu\nu}$



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D-branes

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Curved spacetime

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low energy limit  
 $\mathcal{N} = 4, D = 4$  SYM

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10d Supergravity

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Near horizon limit  
 $AdS_5 \times S^5$

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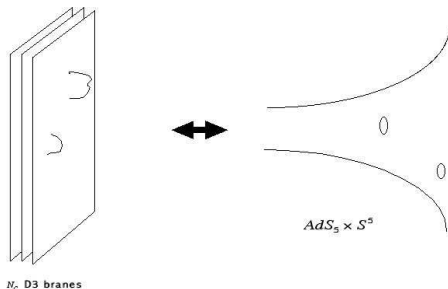
Curved spacetime

low energy limit  
10d Supergravity

Near horizon limit  
 $AdS_5 \times S^5$

- There is Open-Closed string duality

# AdS/CFT correspondence



- Weak coupling limit ( $\lambda \ll 1$ ):  $\mathcal{N} = 4$ ,  $D = 4$ ,  $SU(N_C)$  SYM
- Strong coupling limit ( $\lambda \gg 1$ ): Classical gravity in  $AdS_5 \times S^5$
- From calculating classical gravity, we can obtain some quantities in gauge theory with strong coupling.

# AdS/CFT correspondence

- AdS/CFT dictionary

Gauge Theory(boundary)      ||      Gravity(bulk)

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Global symmetry

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Gauge Theory(boundary)		Gravity(bulk)
Operator $\mathcal{O}$		Field $\phi$
Energy momentum tensor $T_{\mu\nu}$		Graviton $g_{\mu\nu}$
Conformal dimension $\Delta_\phi$		mass of field $m_\phi$
Global symmetry		Gauge symmetry
...		...

- But it is too far from realistic QCD !!
  - SUSY, zero temperature, no fundamental field ...

- Toward to QCD
  - Break SUSY

- D4 brane background with circle compactification

$$ds^2 = \left(\frac{U}{R}\right)^{3/2} (-dt^2 + d\vec{x}^2 + dx_4^2) + \left(\frac{R}{U}\right)^{3/2} (dU^2 + r^2 d\Omega_4^2)$$

$$e^\phi = g_s \left(\frac{U}{R}\right)^{3/4}, \quad F_4 = \frac{2\pi N_c}{\Omega_4} \epsilon_4, \quad R^3 = \pi g_s N_c l_s^3$$

By imposing anti-periodic condition to fermion along  $x_4$  direction, we can break supersymmetry

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- Finite temperature

- Schwarzschild type solution

$$ds^2 = \left(\frac{U}{R}\right)^{3/2} (-f(U)dt^2 + d\vec{x}^2 + dx_4^2) + \left(\frac{R}{U}\right)^{3/2} \left(\frac{dU^2}{f(U)} + U^2 d\Omega_4^2\right)$$

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- Confining phase

- Double Wick rotation

$$t \leftrightarrow i x_4, \quad x_4 \leftrightarrow i \tau$$

$$ds^2 = \left(\frac{U}{R}\right)^{3/2} (-d\tau^2 + d\vec{x}^2 + f(U)dx_4^2) + \left(\frac{R}{U}\right)^{3/2} \left(\frac{dU^2}{f(U)} + U^2 d\Omega_4^2\right)$$

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    - Flavor degrees of freedom
      - Add probe branes
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$$A_t(\rho) \sim \mu + \frac{Q}{\rho^\#} + \dots$$
  - Source of gauge field
    - endpoint of fundamental strings
    - baryon vertex (spherical D-brane with  $N_c$  fundamental strings)

# (A)symmetric dense matter

- Basic Setup
  - Confining phase

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  - Two flavor  $\leftrightarrow$  Two probe brane
  - Finite density

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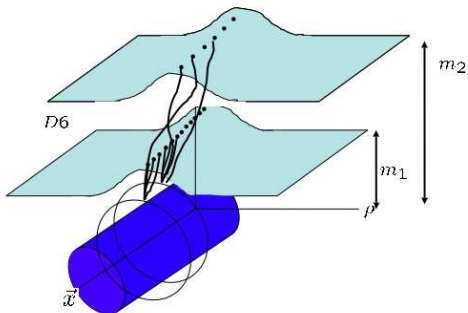
- Basic Setup
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# (A)symmetric dense matter

- DBI action
  - D4 brane(baryon vertex)

$$S_{D4} = -\mu_4 \int e^{-\phi} \sqrt{\det(g + 2\pi\alpha' F)} + \mu_4 \int A_{(1)} \wedge G_{(4)}$$

- D6 brane(flavor brane)

$$S_{D6} = \int dt \mathcal{L}_{D6} = -\mu_6 \int e^{-\phi} \sqrt{\det(g + 2\pi\alpha' F)}$$

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- Legendre transformed 'Hamiltonian'

- D4 brane(baryon vertex)

$$\mathcal{H}_{D4} = \tau_4 \int d\theta \sqrt{\omega_+^{4/3} (\xi^2 + \xi'^2)} \sqrt{D(\theta)^2 + \sin^6 \theta}$$

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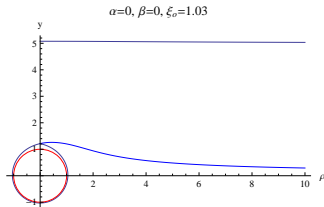
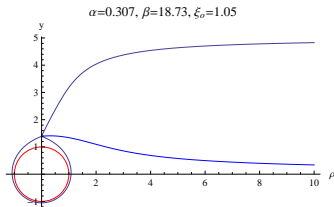
$$\mathcal{H}_{D6} = \tau_6 \int d\rho \sqrt{\omega_+^{4/3} (\tilde{Q}^2 + \rho^4 \omega_+^{8/3})} \sqrt{1 + \dot{y}^2}$$

# (A)symmetric dense matter

- Solve equation of motion for DBI action with force balance condition

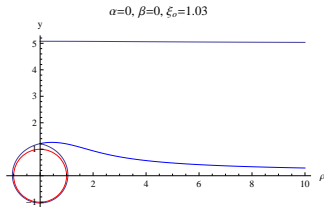
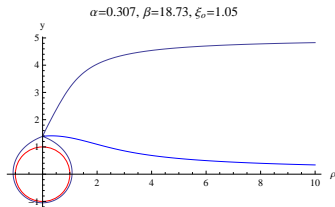
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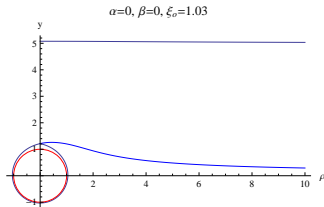
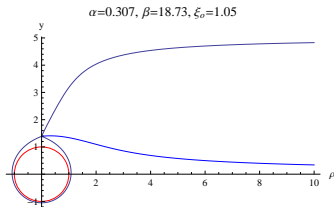


- Total free energy

$$E_{tot} = \frac{Q}{N_C} \mathcal{H}_{D4} + \mathcal{H}_{D6}((1-\alpha)Q) + \mathcal{H}_{D6}(\alpha Q)$$

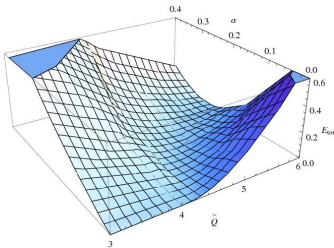
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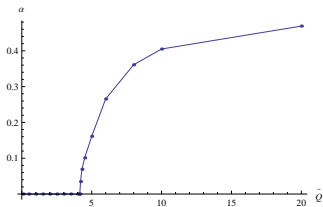
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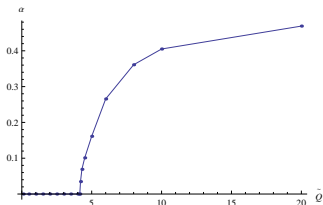
- Transition from nuclear matter to strange matter



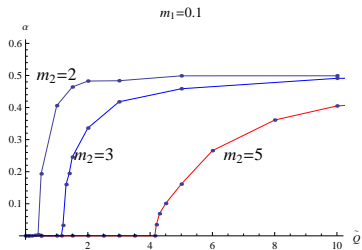


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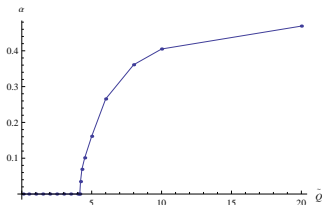


- Different mass ratio

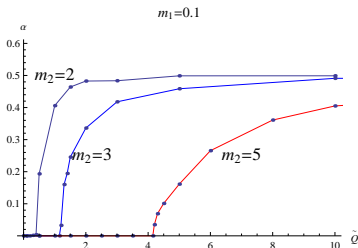


# (A)symmetric dense matter

- Transition from nuclear matter to strange matter



- Different mass ratio



- $m_2/m_1 = 1$ , physical system is  $\alpha = 0.5$  for any density  $\rightarrow$  symmetric matter

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$$\tilde{\alpha} \equiv (N - Z)/A = 1 - 2\alpha$$

$$S_2 \sim \frac{\partial^2 E_{tot}}{\partial \alpha^2}$$

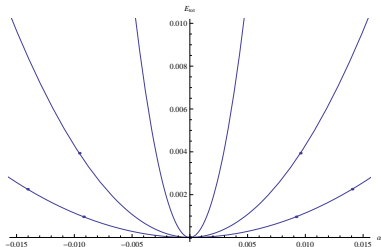
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- Total energy

$$\begin{aligned} E(\rho, \alpha) &= E(\rho, 0) + \left. \frac{\partial E}{\partial \tilde{\alpha}} \right|_{\tilde{\alpha}=0} \tilde{\alpha} + \frac{1}{2} \left. \frac{\partial^2 E}{\partial \tilde{\alpha}^2} \right|_{\tilde{\alpha}=0} \tilde{\alpha}^2 \dots \\ &= E_0(\rho, 0) + E_1 \tilde{\alpha} + E_2 \tilde{\alpha}^2 \dots \end{aligned}$$

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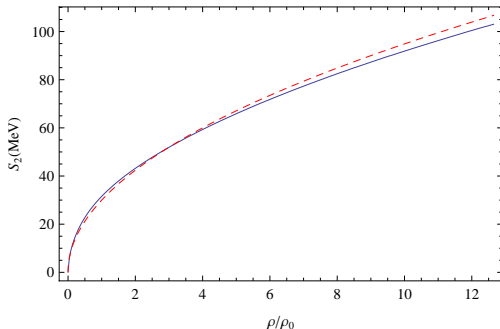
$$E_1 = 0$$

# Symmetry energy

$$E_2 = \frac{2\tau_6}{N_B} \int d\rho \frac{\sqrt{1 + \dot{y}^2} \tilde{Q}^2 \omega_+^{10/3} \rho^4}{(\tilde{Q}^2 + 4\omega_+^{8/3} \rho^4)^{3/2}}$$

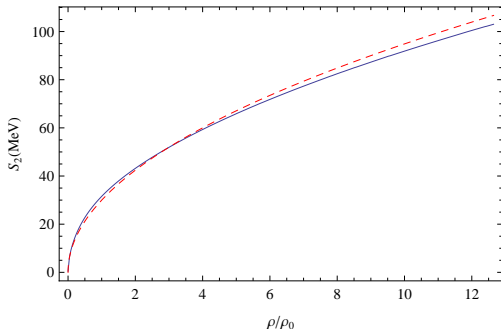
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- Analytic calculation ( $m_q \gg 1$ ,  $\tilde{Q} \ll 1$ )

$$E_2 = \left( \Gamma \left( \frac{5}{4} \right) \right)^2 \sqrt{\frac{\lambda \rho_0}{2M_{KK}}} \sqrt{\frac{\rho}{\rho_0}}$$

- Physical quantities

$$\begin{aligned}M_q &= \frac{m_\infty \lambda M_{KK}}{9\pi} \\ \rho &\equiv \frac{N_b}{V_3} = \frac{Q}{N_c V_3} \\ &= \frac{2 \cdot 2^{2/3}}{81(2\pi)^3} \lambda M_{KK}^3 \tilde{Q}\end{aligned}$$

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- From  $m_{\eta'} \sim 390\text{MeV}$  and  $M_q = 7\text{MeV}$ ,

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# Conclusion and Discussion

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- We are calculating fluctuation spectrum for goldstone mode for  $m_2/m_1 = 50$ 
  - Turn on matrix valued fluctuation
  - Off-diagonal fluctuation spectrum corresponds to mass of 'Kaon'

# Conclusion and Discussion

- We calculate symmetry energy in dense matter in D4/D6/D6 model
  - Symmetry energy is monotonically increase as density increase
  - In low density

$$E_2 \sim \rho^{1/2}$$

- In the case of  $m_2/m_1 = 2$ , result of almost same
- We are calculating (a)symmetry energy for large mass ratio.
- We are calculating fluctuation spectrum for goldstone mode for  $m_2/m_1 = 50$ 
  - Turn on matrix valued fluctuation
  - Off-diagonal fluctuation spectrum corresponds to mass of 'Kaon'
  - We can study density dependence of mass of Kaon



Thank you !!!