

# Forward photon measurement and generic detector R&D

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# Contents

- CGC-related  $\gamma$ -physics at forward
  - Related to RHIC & LHC
  - Calculable from the 1<sup>st</sup> principle.
  - To be addressed 5-10 years down the road.
- Relevant generic detector R&D
  - Si/W sandwich calorimeter
  - SiPM (Silicon photomultiplier)
  - 무엇보다도 한국이 우수한 잠재력을 갖춘 영역에서

# CGC-related $\gamma$ -physics at forward

- Application domain
  - Perturbative QCD :  $\Lambda_{\text{QCD}} \ll p_T, p_T \sim Q_s$
  - High parton density at small  $x$  : bigger effect at forward/backward
- Parton kinematics
  - Projectile : large  $x$  ( $\sim$ valence quarks)
  - Target : Small  $x$  ( $\sim$ gluons  $\rightarrow$  CGC)
- Study based on Jamal Jalilian-Marian's work (PRD66(2002)014021, PRD66(2002)094014, PRD67(2003)074019, NPA753(2005)307)

# Baseline prediction

- Forward hadron suppression
  - Jet quenching?
  - Parton coalescence model?

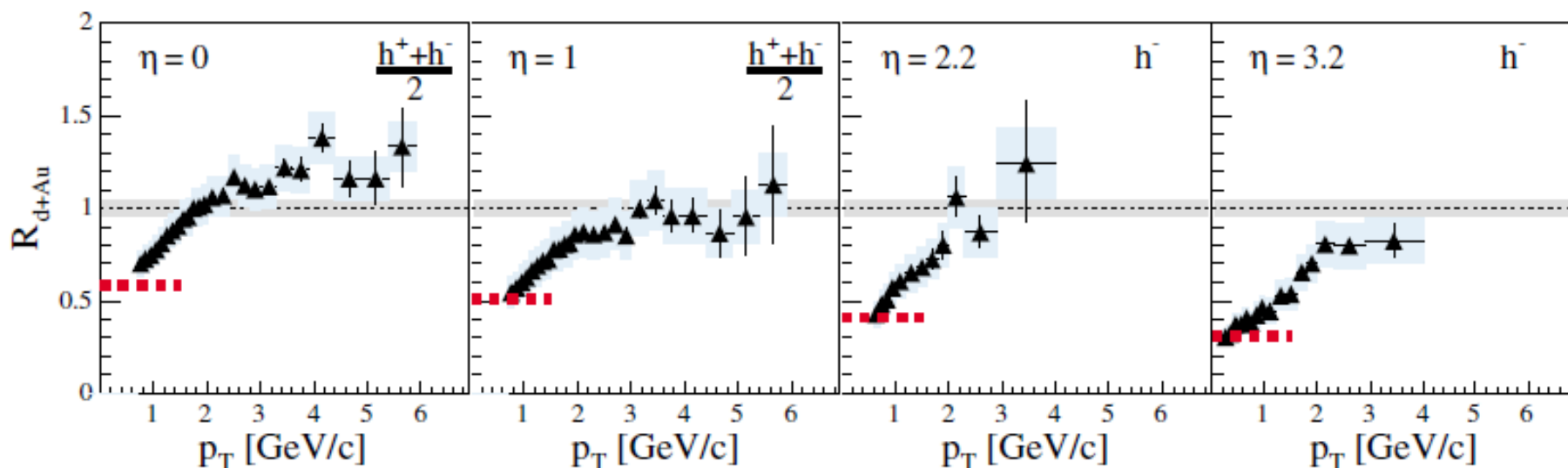
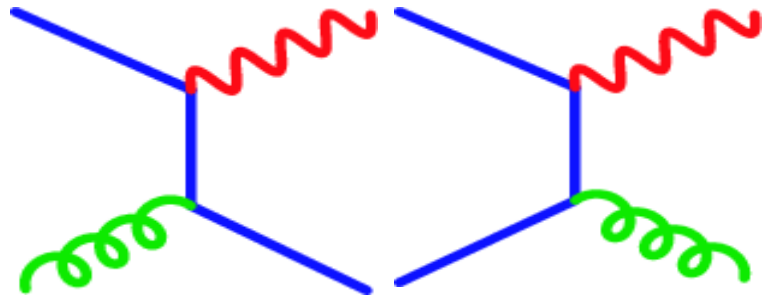


FIG. 2 (color online). Nuclear modification factor for charged hadrons at pseudorapidities  $\eta = 0, 1.0, 2.2, 3.2$ . One standard deviation statistical errors are shown with error bars. Systematic errors are shown with shaded boxes with widths set by the bin sizes. The shaded band around unity indicates the estimated error on the normalization to  $\langle N_{coll} \rangle$ . Dashed lines at  $p_T < 1.5$  GeV/c show the normalized charged-particle density ratio  $\frac{1}{\langle N_{coll} \rangle} \frac{dN/d\eta(\text{Au})}{dN/d\eta(\text{pp})}$ .

# Motivation : Direct $\gamma$ production

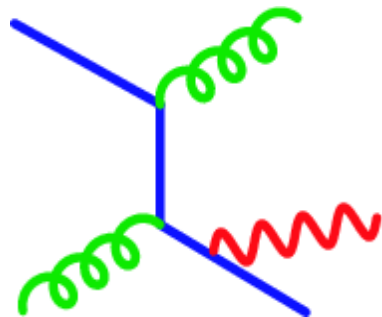
- Direct  $\gamma$  production in p+p

→ One of the best known QCD process...



Really?

→ Leading order diagram  
in perturbation theory



Hard photon : Higher order pQCD  
Soft photon : Initial/final radiation,  
Fragmentation function

# Baseline approach

$$q(p) + A \rightarrow q(q) + \gamma(k) + X, A$$

A : Background field, Color Glass Condensate at small x

$$\langle \{q(\vec{q})\gamma(\vec{k})\}_{out} | q(\vec{p})_{in} \rangle = \langle 0_{out} | a_{out}(\vec{k}) b_{out}(\vec{q}) b_{in}^+(\vec{p}) | 0_{in} \rangle$$

Using reduction formalism,

$$\begin{aligned} & \langle 0_{out} | a_{out}(\vec{k}) b_{out}(\vec{q}) b_{in}^+(\vec{p}) | 0_{in} \rangle \\ &= \frac{e}{Z_2 \sqrt{Z_3}} \int d^4x d^4y d^4z \cdot e^{i(k \cdot x + q \cdot z - p \cdot y)} \times \bar{u}(\vec{q})(i\vec{\partial}_z - m) \\ & \quad \times \langle 0_{out} | T \psi(z) \varepsilon \cdot J(x) \bar{\psi}(y) | 0_{in} \rangle \times (i\vec{\partial}_y + m) u(p), J_\mu(x) \equiv \bar{\psi}(x) \gamma_\mu \psi(x) \\ & \langle 0_{out} | T \psi(z) \bar{\psi}(x) \not{x} \psi(x) \bar{\psi}(y) | 0_{in} \rangle = -G_F(z, x) \not{x} G_F(x, y) \end{aligned}$$

$$\begin{aligned}
G_F(x, y) &\equiv \langle 0_{out} | T \bar{\psi}(y) \psi(x) | 0_{in} \rangle \\
&= G_F^0(x-y) + \int d^4 z \delta(z^-) [\theta(x^-) \theta(-y^-) (U^+(z_T) - 1) \\
&\quad - \theta(-x^-) \theta(y^-) (U(z_T) - 1)] G_F^0(x-z) \gamma^- G_F^0(z-y) \\
U(x_T) &\equiv T e^{-ig^2 \int_{-\infty}^{\infty} dz^- \frac{1}{\nabla_T^2} \rho_a(z^-, z_T) t^a}
\end{aligned}$$

$U(z_T)$  : a unitary matrix containing the interactions between the quark and CGC.

in coordinate space and in the "singular" gauge.

L. McLerran and R. Venugopalan, Phys. Rev. D59, 094002(1999)

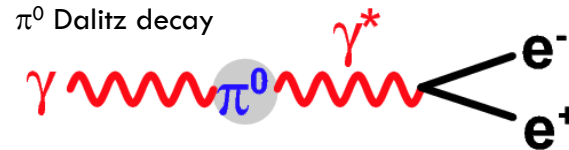
$$\begin{aligned}
\langle \{q(\vec{q})\gamma(\vec{k})\}_{out} | q(\vec{p})_{in} \rangle &= -ie\bar{u}(q) \left[ \frac{\gamma^-(p-k+m)\not{\epsilon}}{(p-k)^2 - m^2} + \frac{\not{\epsilon}(q+k+m)\gamma^-}{(q+k)^2 - m^2} \right] u(p) \\
&\times 2\pi\delta(q^- + k^- - p^-) \int d^2x_T e^{i(\vec{q}_T + \vec{k}_T - \vec{p}_T)\vec{x}_T} (U(\vec{x}_T) - 1) \\
&= 2\pi\delta(q^- + k^- - p^-) M(\vec{p} | \vec{q}\vec{k})
\end{aligned}$$

$$d\sigma = \frac{d^3\vec{k}}{(2\pi)^3 2k_0} \frac{d^3\vec{q}}{(2\pi)^3 2q_0} \frac{1}{2p^-} \times M(\vec{p} | \vec{q}\vec{k}) M^*(\vec{p} | \vec{q}\vec{k}) \cdot 2\pi\delta(p^- - q^- - k^-)$$



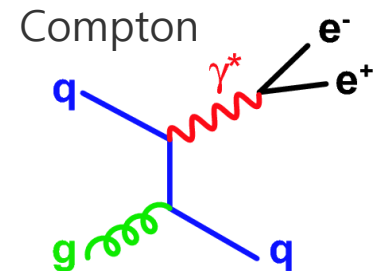
# $\gamma^*(e^+e^- \text{ pair})$ measurement

- Motivation:  
Low  $p_T$  : large  $\pi^0$ ,  $\eta$  background
- Internal conversion
  - Any source of real photons also emits virtual photons
  - Well known example:



- ◆ Rate and  $m_{ee}$  distribution calculable in QED (Kroll-Wada formula)
- Hadron decays:  $m_{ee} < M_{\text{hadron}}$
- Essentially not such limit for point-like source.

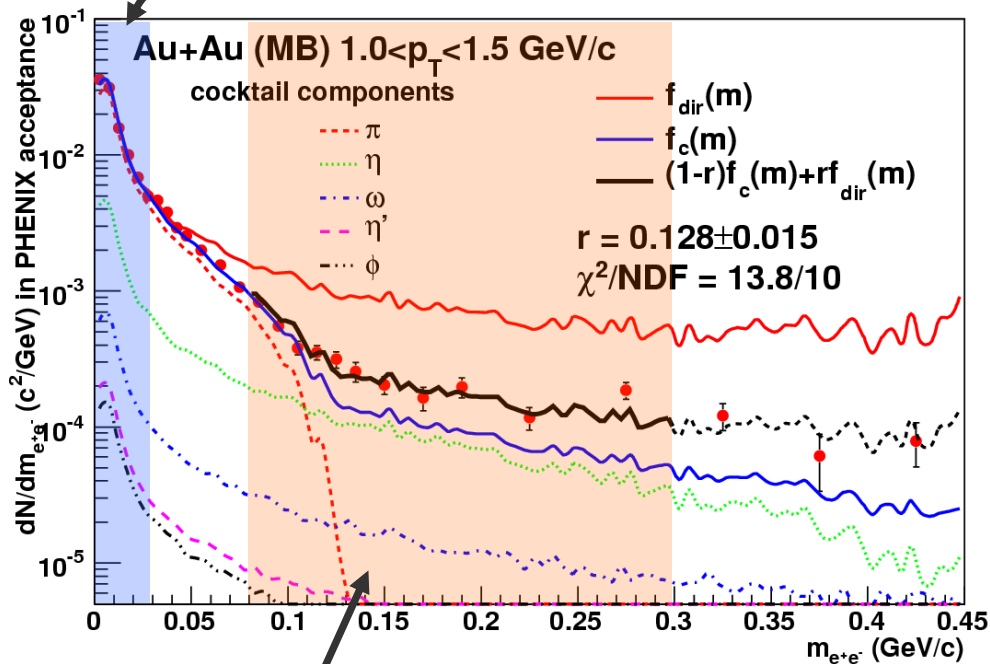
Improve signal-to-background ratio by measuring  $e^+e^-$  pairs with  $m_{ee} > \sim M_{\text{pion}}$



# Extraction of $\gamma$ signal

$$f(m_{ee}) = (1 - r) \cdot f_{\text{cocktail}}(m_{ee}) + r \cdot f_{\text{direct}}(m_{ee})$$

Separately normalized  
to data at  $m_{ee} < 30$  MeV



Fit range:  $80 < m_{ee} < 300$  MeV

- Interpret deviation from hadronic cocktail ( $\pi, \eta, \omega, \eta', \phi$ ) as signal from virtual direct photons

- Extract fraction  $r$  with two-component fit

$$r = \frac{\gamma_{\text{direct}}^*}{\gamma_{\text{inclusive}}^*} \Big|_{m_{ee} < 30 \text{ MeV}}$$

- Fit yields good  $\chi^2/\text{NDF}$  (13.8 / 10)

# $\gamma$ Production vs $\gamma^*$ production

$$d\sigma = \frac{d^3\vec{k}}{(2\pi)^3 2k_0} \frac{d^3\vec{q}}{(2\pi)^3 2q_0} \frac{1}{2p^-} \times M(\vec{p} | \vec{q}\vec{k}) M^*(\vec{p} | \vec{q}\vec{k}) \cdot 2\pi\delta(p^- - q^- - k^-)$$

$$M(\vec{p} | \vec{q}\vec{k}) = -ie\bar{u}(q) \left[ \frac{\gamma^-(p-k+m)\not{\epsilon}}{(p-k)^2 - m^2} + \frac{\not{\epsilon}(q+k+m)\gamma^-}{(q+k)^2 - m^2} \right] u(p)$$

$$\times \int d^2x_T e^{i(\vec{q}_T + \vec{k}_T - \vec{p}_T)\vec{x}_T} (U(\vec{x}_T) - 1)$$

$$d\sigma = \frac{d^4\vec{k}}{(2\pi)^4} \frac{d^3\vec{q}}{(2\pi)^3 2q_0} \frac{1}{2p^-} \frac{2\alpha_{em}}{3k^2} \times M^\mu(\vec{p} | \vec{q}\vec{k}) M_\mu^*(\vec{p} | \vec{q}\vec{k}) \cdot 2\pi\delta(p^- - q^- - k^-)$$

$$M^\mu(\vec{p} | \vec{q}\vec{k}) = -ie\bar{u}(q) \left[ \frac{\gamma^-(p-k+m)\gamma^\mu}{(p-k)^2 - m^2} + \frac{\gamma^\mu(q+k+m)\gamma^-}{(q+k)^2 - m^2} \right] u(p)$$

$$\times \int d^2x_T e^{i(\vec{q}_T + \vec{k}_T - \vec{p}_T)\vec{x}_T} (U(\vec{x}_T) - 1)$$

# Universality (DIS)

$$\sigma_{DIS} = \int_0^1 dz \int d^2 r_T \left| \Psi(k^\pm, k_T | z, r_T) \right|^2 \sigma_{dipole}(r_T)$$

$$= \int_0^1 dz \int d^2 r_T F_{known} \cdot \sigma_{dipole}(r_T)$$

$$\sigma_{dipole}(r_T) \equiv \frac{2}{N_c} \int d^2 X_T \text{Tr} \left\langle 1 - U \left( X_T + \frac{r_T}{2} \right) U^\dagger \left( X_T - \frac{r_T}{2} \right) \right\rangle_\rho$$

$$d\sigma_{\gamma^*} = \frac{d^4 \vec{k}}{(2\pi)^4} \frac{d^3 \vec{q}}{(2\pi)^3 2q_0} \frac{1}{2p^-} \frac{2\alpha_{em}}{3k^2} \times M^\mu(\vec{p} | \vec{q} \vec{k}) M_\mu^*(\vec{p} | \vec{q} \vec{k}) \cdot 2\pi \delta(p^- - q^- - k^-)$$

$$= \frac{2}{N_c} \int d^2 X_T \text{Tr} \left\langle 1 - U \left( X_T + \frac{r_T}{2} \right) U^\dagger \left( X_T - \frac{r_T}{2} \right) \right\rangle_\rho \cdot G_{known}$$

# Issues...

Phenomenological models for the dipole cross sections...  
(baseline JIMWLK equations)

Dipole cross section

$$N(x, r_T, b_T) = \frac{1}{N_c} \text{Tr} \langle 1 - V^+(x_T) V(y_T) \rangle$$

$$x_g = 1.6 \times 10^{-4}$$

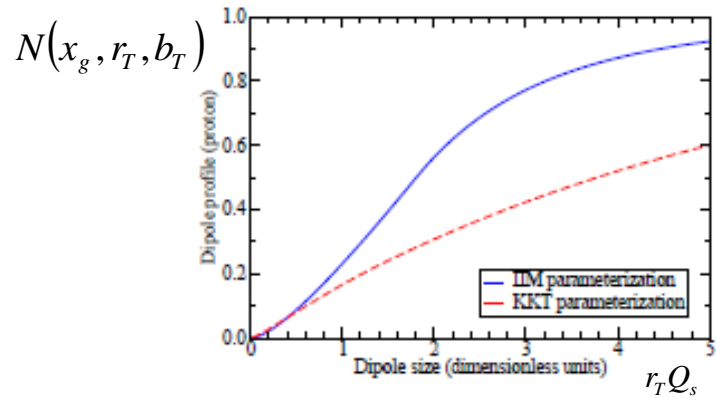


Figure 1: Quark anti-quark dipole profile for a proton target.

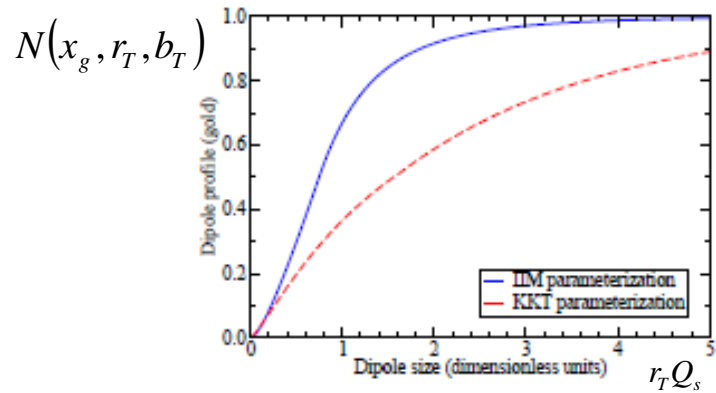
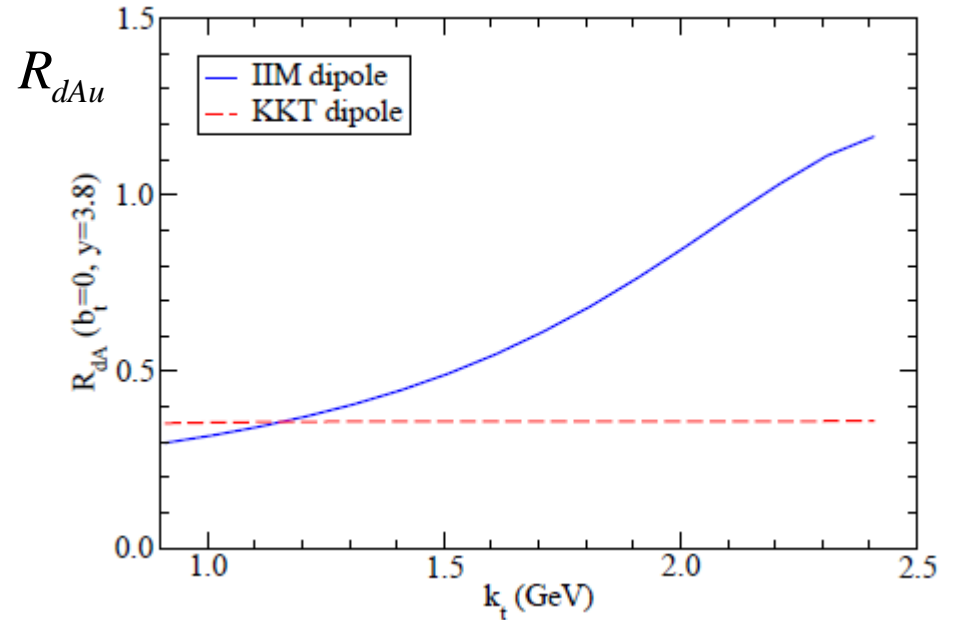


Figure 2: Quark anti-quark dipole profile for a nuclear target.

$$R_{dAu} \equiv \frac{d\sigma^{dA \rightarrow \gamma X}}{dy d^2 k_T d^2 b_T} \frac{2A}{d\sigma^{pp \rightarrow \gamma X}}$$

RHIC



J. Jalilian-Marian, NPA753 (2005) p307-315

# Summary

- Forward  $\gamma$  production at RHIC/LHC
  - Extension of pQCD possible
  - Extreme higher twist (large  $Q_s$ ) and CGC.
  - Universality in dipole cross section (DIS & Forward  $\gamma$  production)
  - Dipole cross section is not yet calculated from baseline JIMWLK equation, but models for it exist.

# Generic detector R&D



# 한국의 잠재역량

R&D environment



6 inch fabrication line



8 inch fabrication line

## MEMORANDUM

- (1) Youngil Kwon, Mann-Ho C
- (2) Edward Kistenev, Andrey S
- (3) John Lajoie, Physics and As
- (4) Yongsun Yoon, BT division,
- (5) Kwun-bum Chung, Electrop
- (6) Zheng Li, SDDPL, Instrum
- (7) Jinsoo Kim, National Nano

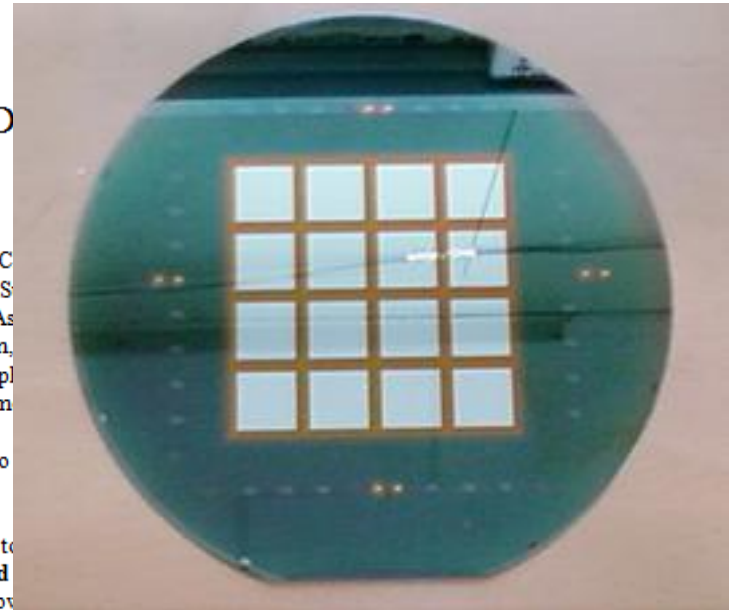
### I. Purpose & Scope<sup>1)</sup>

The purpose of this MOU is to... the 'Radiation damage and... planning to contribute their ov... studies. This MOU clarifies the areas of pe... collaborate by sharing their expertise and se... parties for the stated academic goal. <sup>1)</sup>

300 cm<sup>2</sup> ~ \$ 500 ... rties will ticipating

### II. Responsibilities Under this MOU<sup>1)</sup>

- A. Dr. E. Kistenev, and Prof. J. Lajoie, and Prof. Y. Kwon will propose silicon semiconductor detectors/devices to achieve academic goals in their field of interest in the experimental nuclear and high energy physics.<sup>1)</sup>
- B. Dr. Z. Li will design the proposed silicon semiconductor detectors/devices using standards approved by industry for large area radiation hard Si devices and will advice on the radiation induced defects in Si devices. <sup>1)</sup>
- C. Dr. A. Sukhanov will advice on the electronic design and implementation of the readout electronics for silicon semiconductor device testing.<sup>1)</sup>
- D. Dr. Yoon will inspect designs of the proposed detectors/devices and advise on matching design ideas to fabrication technologies. He will also perform his own radiation hardness testing of the devices he develops.<sup>1)</sup>
- E. Prof. M.-H. Cho, Prof. G. T. Park, and Prof. K. B. Chung will advise on possible defects in silicon sensors/devices and will study radiation defects in the produced sensors/devices exposed to different kinds of radiation.<sup>1)</sup>
- F. Mr. Kim, leader of nano|patterning process team in National Nanofab Center, will assist in fabrication of the silicon sensors/devices with university discount program and consult on details of silicon detector/device fabrication process.<sup>1)</sup>

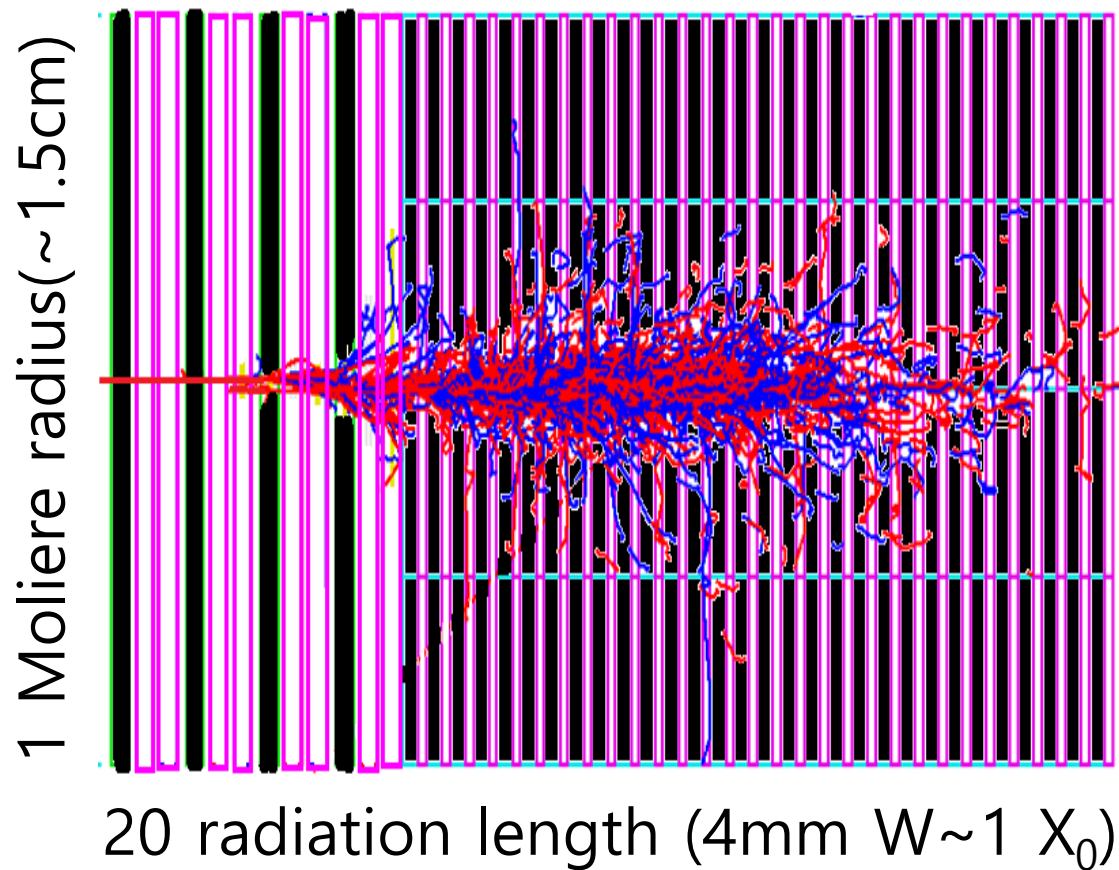


# FOCAL

- A compact sampling electromagnetic calorimeter
  - deals with large particle flux at forward
  - observes part of the energy deposited by particle for the optimized shower
  - measures key particles,  $e$ ,  $\gamma$  and  $\pi^0$ , within small space
  - measures energy.

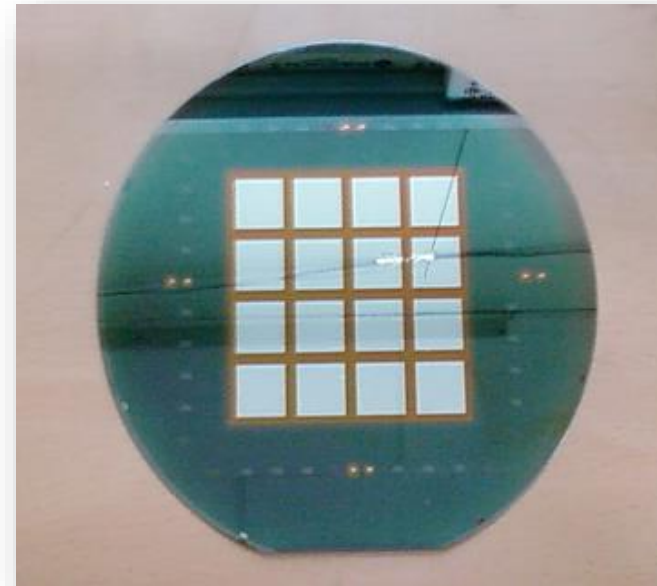
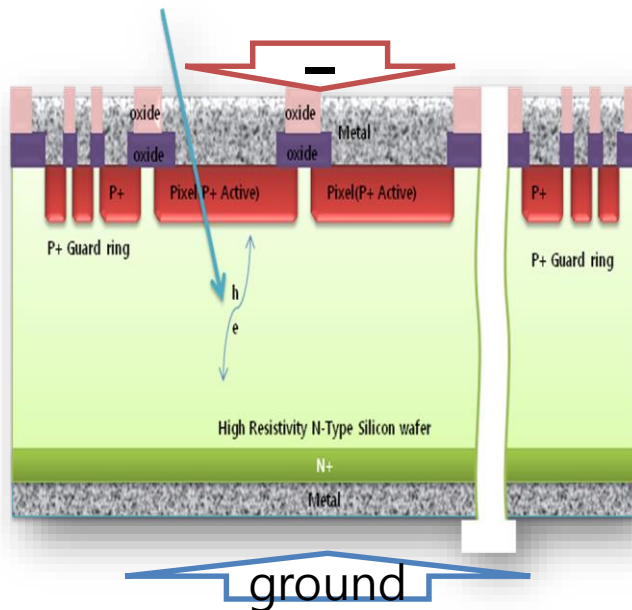
# A compact EMCal, Si/W sandwich calorimeter

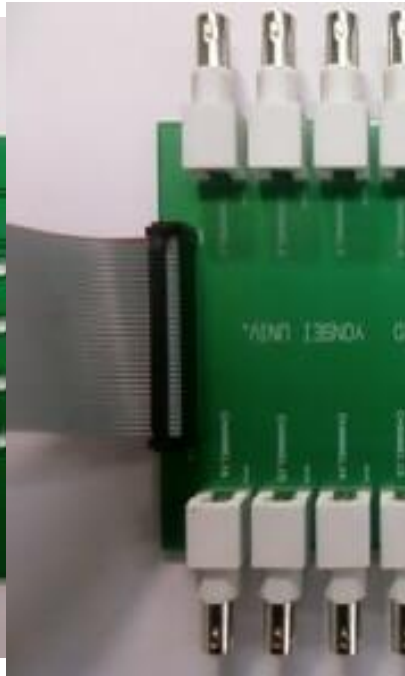
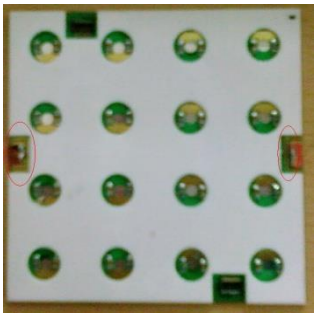
50 (GeV) electron



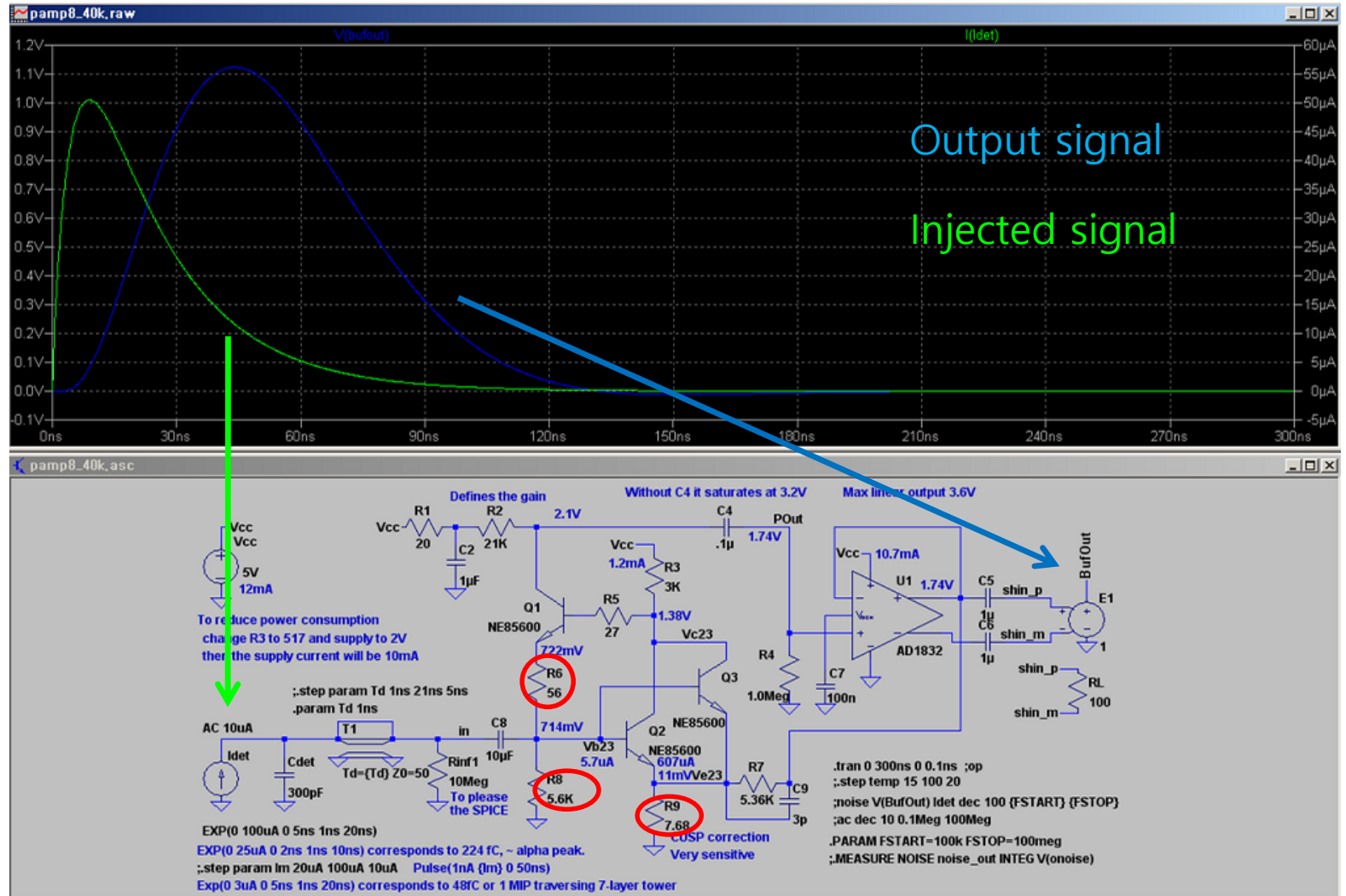
# Silicon pad sensor

- **Basically** PN junction diode in reverse bias mode.
  - N-type substrate and p-type pattern for high energy application => **electrons are carriers**
- **16 square(1.5cm×1.5cm) pads in one micro-module**



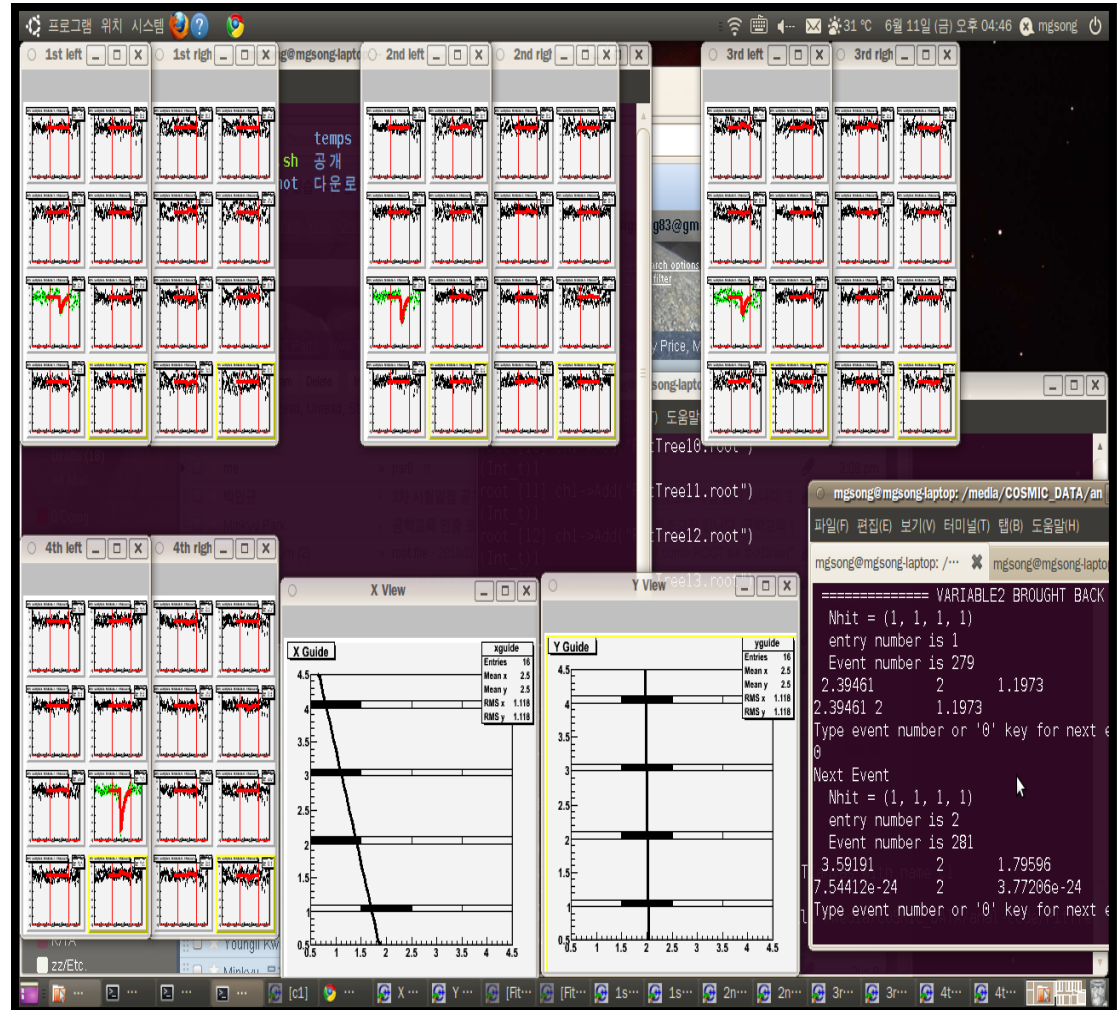
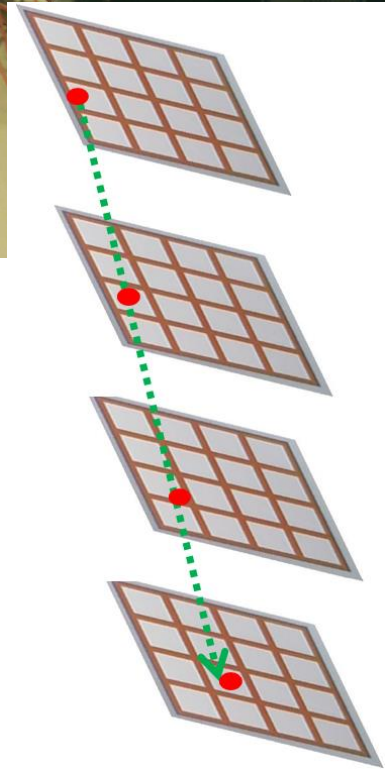
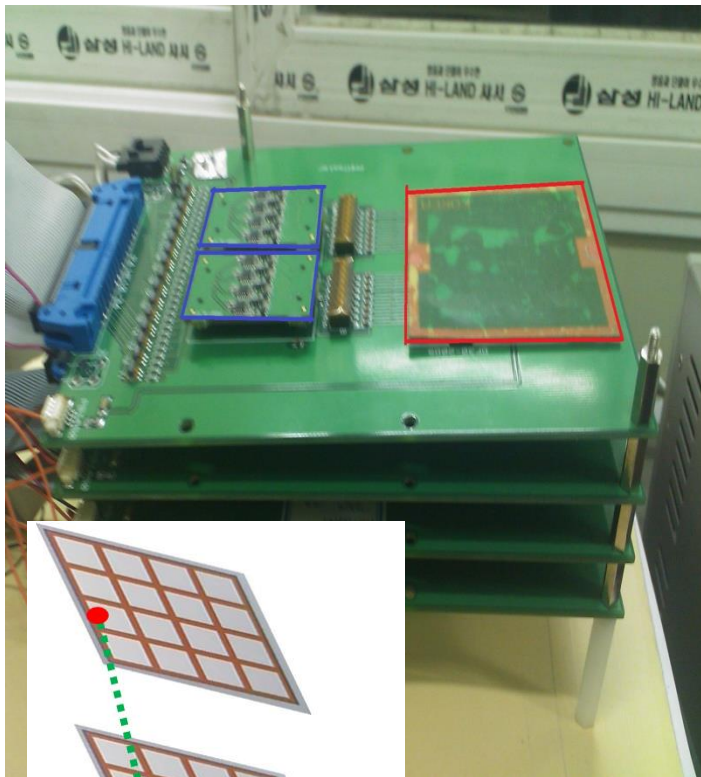


# Preamplifier

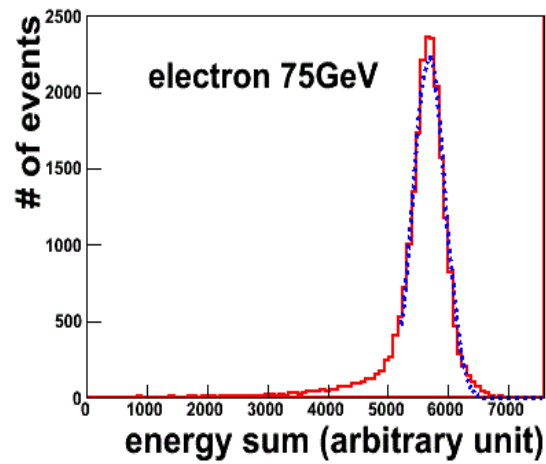
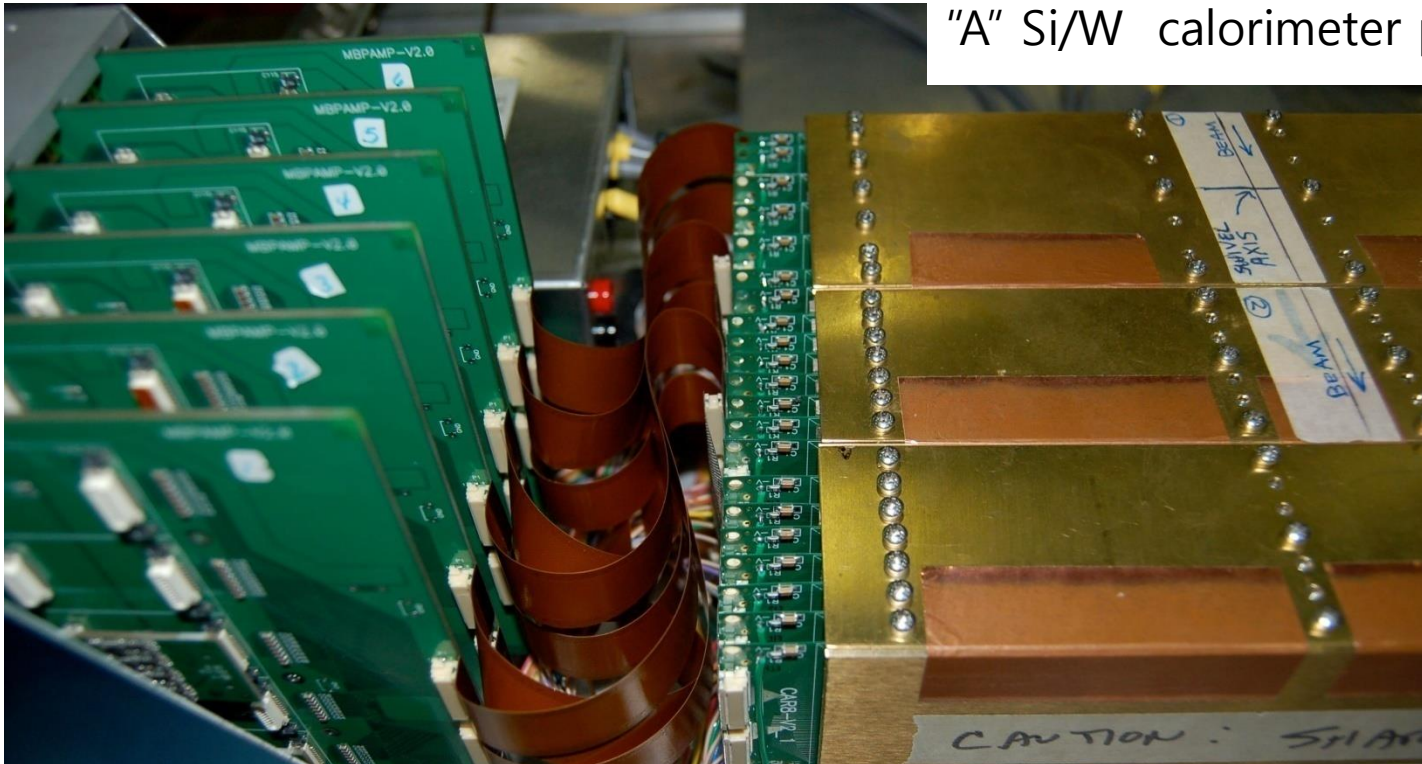




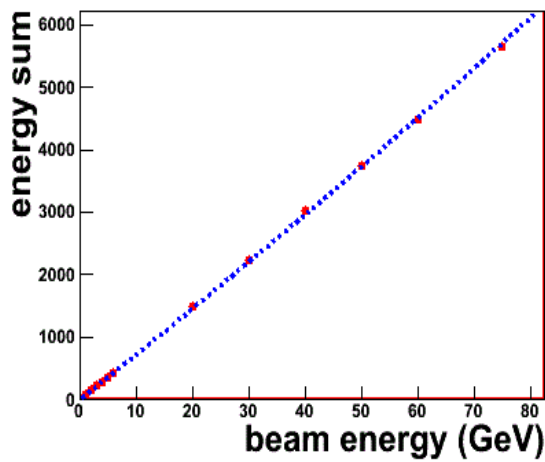
# Cosmic muon test



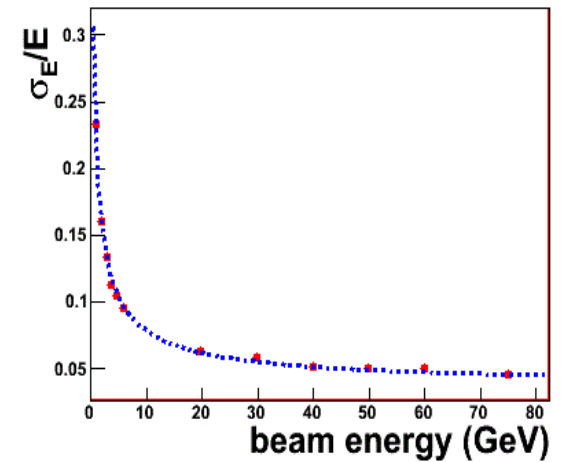
# "A" Si/W calorimeter prototype



(a)



(b)



(c)

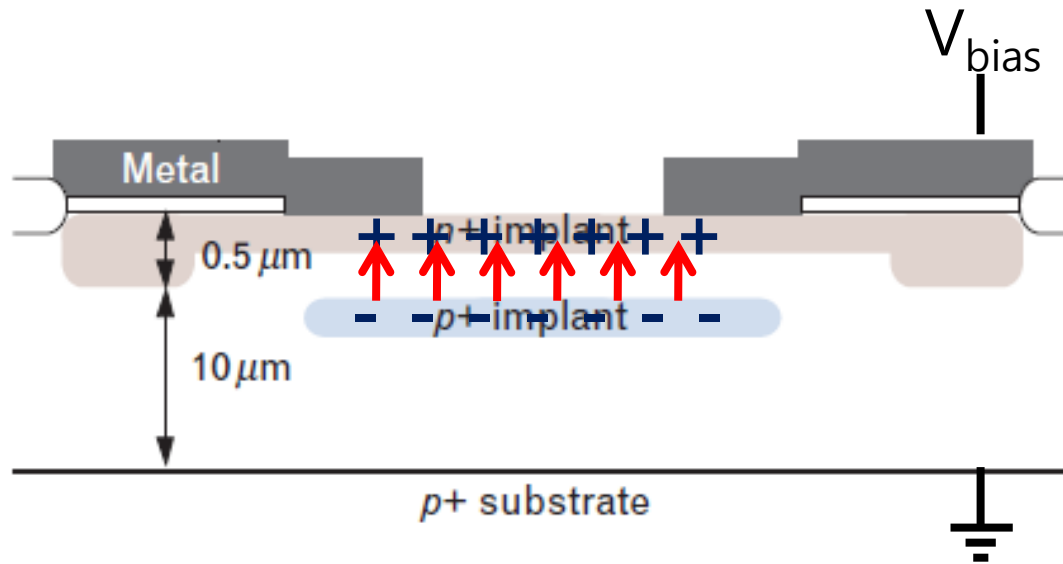


# High sensitivity photon sensor (SiPM)



$\sim 10^5$ - $10^6$  electrons for a photon in visible range

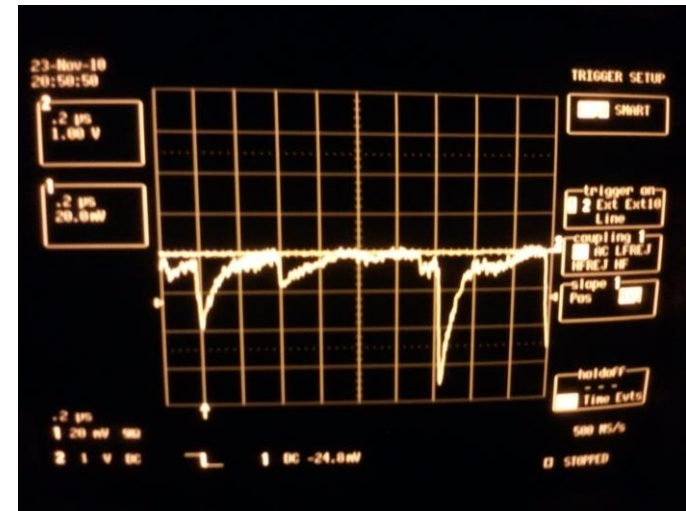
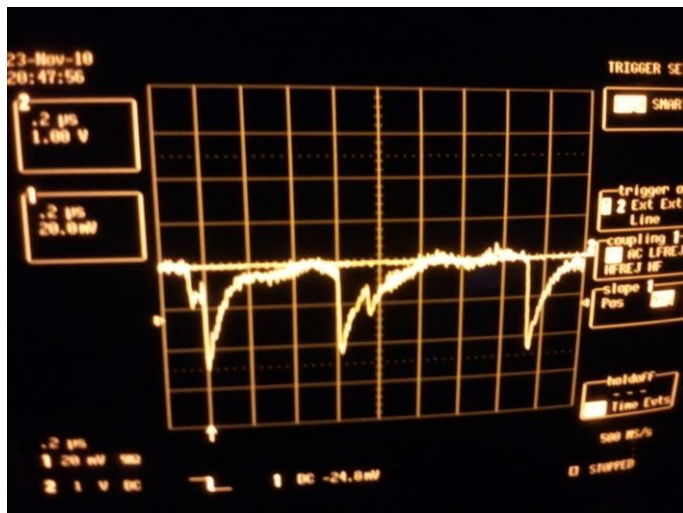
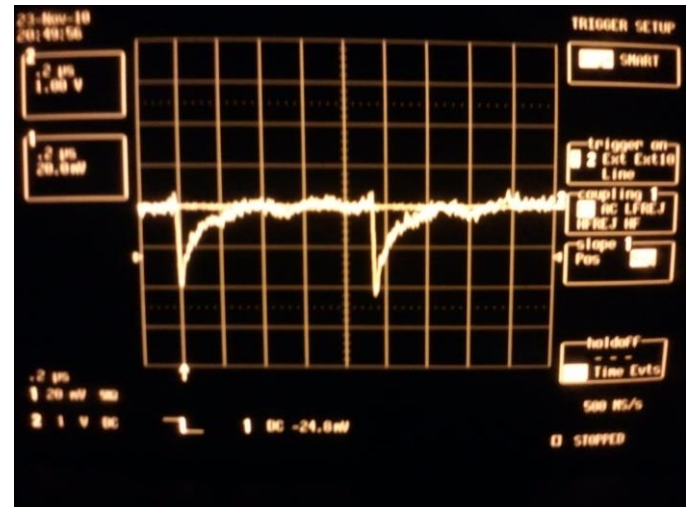
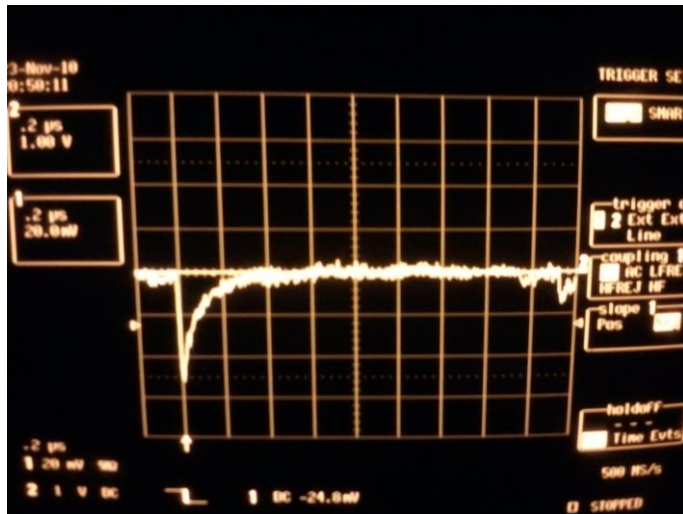
# 구조와 동작



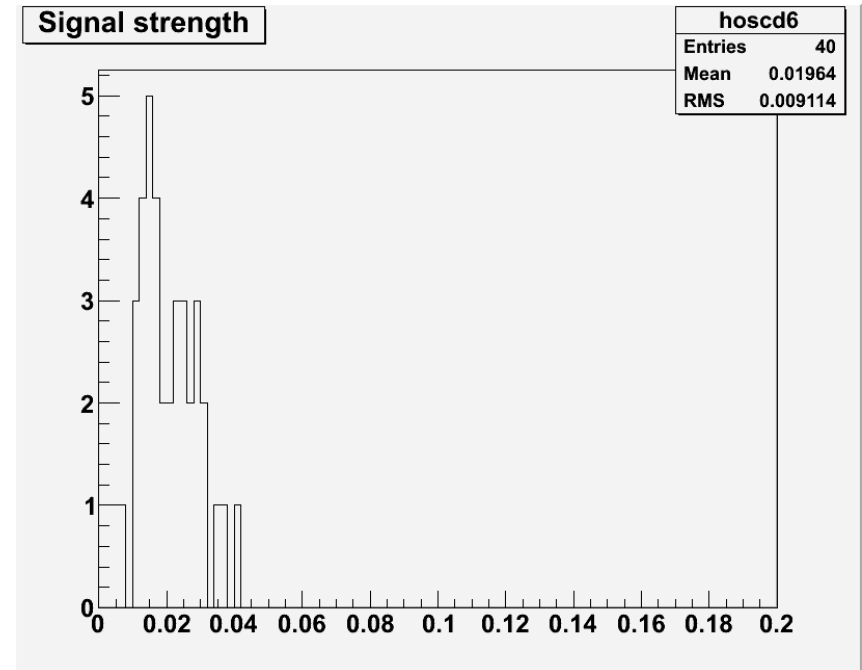
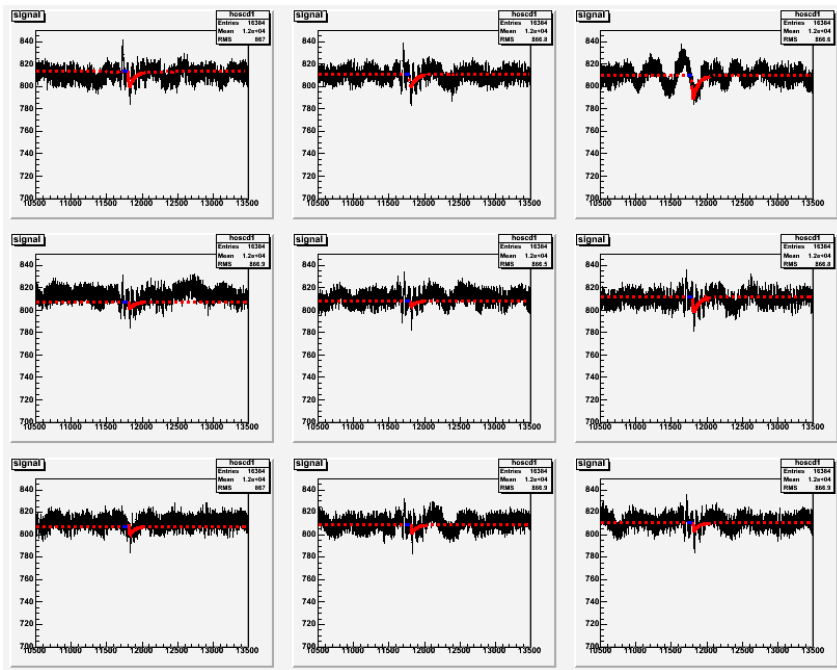
**FIGURE 7.** Cross-section of the basic APD device structure. The APD is fabricated in a lightly  $p$ -doped epitaxial layer grown on a heavily  $n$ -doped ( $n^+$ ) substrate and the substrate is grounded. The  $p^+$  implant creates a sheath electric field that divides the bias voltage into regions of low electric field (below the implant), where photon absorption takes place, and high electric field (above the implant), where impact ionization occurs.

$$\Delta E = \frac{1}{2} C V_{bias}^2 - \frac{1}{2} C V_{B.D.}^2$$

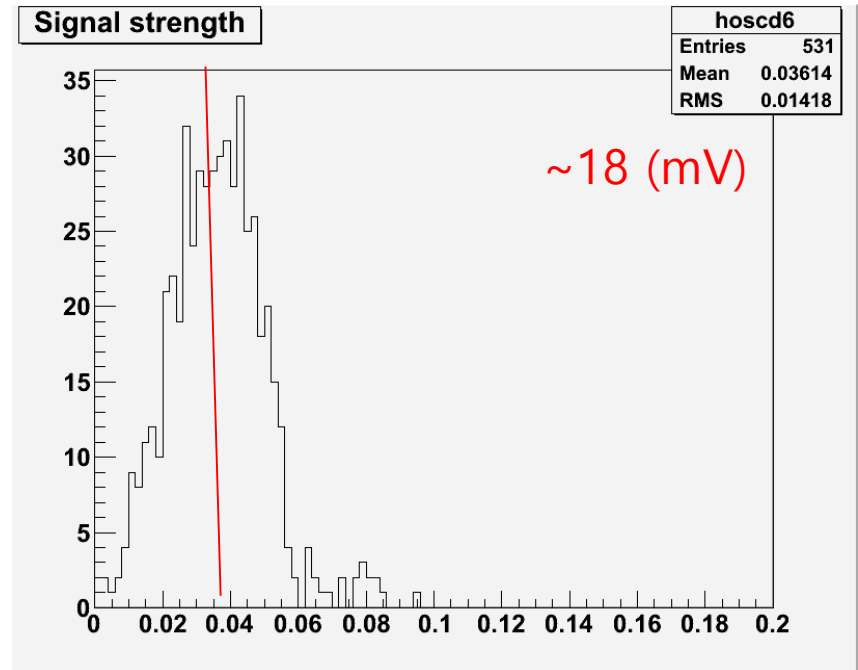
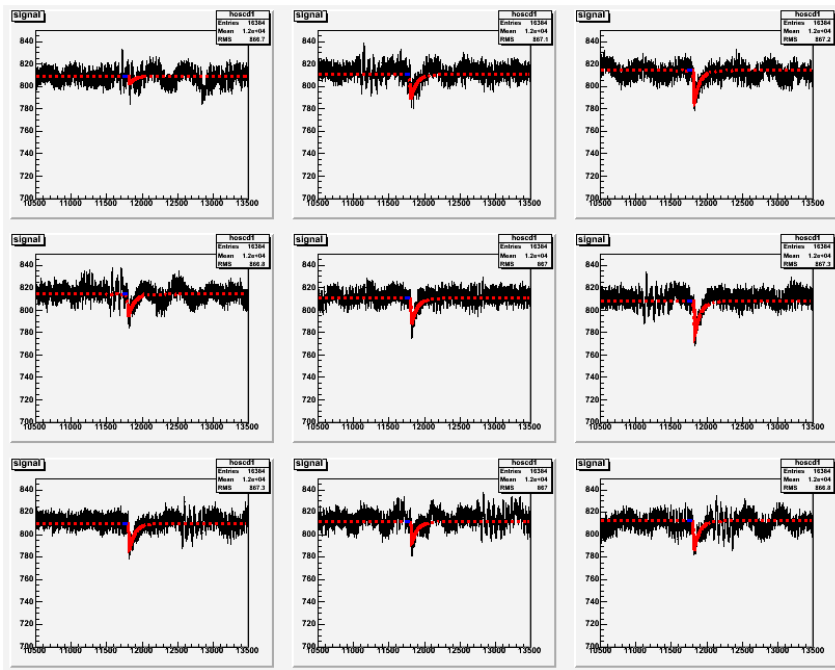
# Noise signal



# 24.5 (V) : Under breakdown



# 25.5 (V) : Breakdown



# 27.0 (V) : Over breakdown

