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Index

- $-1.$ AdS/QCD
- 2. Computing Correlators in AdS/CFT
- **3. Hydrodynamics**
	- Perturbation in RN AdS5 background
- 4. Spectral function
	- Photoemission rate
- **4. Remarks**

AdS/QCD

AdS/CFT : AdS_{5} X S⁵ ~ N=4 SYM in 4 dim. **AdS/QCD:** $aAdS_d \sim low$ energy Yang-Mills in d-1 dim.

Phase transition in AdS/QCD

Assume that confinement and chiral phase transition occurs at same scale, Tc.

SU(2)L X SU(2)R flavor sym. -> SU(2)V

Two point function

Intermediate energy regime

 $\delta^2 Z_{gauge}$ $\delta \varphi_i \delta \varphi_j$ $=G_{ij}$

Two point function

 $\delta \varphi_i \delta \varphi_j$

Two point function

Strongly coupled gauge theory weakly coupled gravity theory

How to compute Green function

Find:
$$
g_{ij}^{(0)}
$$
, $A_t^{(0)}$ / Perturb: $g_{ij}^{(0)} + \delta g_{ij}$, $A_t^{(0)} + \delta A_t$

Find and solve the equations of motion for δq_{ij} , δA_t

plug it into the original action : S eff

Differentiate Seff with respect to the source field

Correlators

th/0205051

Holography : Bulk(d+1 dim) projected onto the Boundary(d dim)

$$
e^{-I_{SUGRA}} \simeq Z_{\text{string}} = Z_{\text{gauge}} = e^{-W}
$$

AdS/CFT says, on-shell partition function of Gravity(Bulk) = CFT partition function(Boundary)

solving Bulk e.o.m. and integrating extra dim of bulk action : Boundary action (On-shell action).

$$
S_{bdry}[\phi_0] = \int \frac{d^4k}{(2\pi)^4} \phi_0(-k) \mathcal{G}(k, u) \phi_0(k) \Big|_{u=0}^{u=1}
$$

$$
G_{\mu\nu}^R(k) = \frac{\delta^2 S}{\delta A_{\mu}^0(-k) \delta A_{\nu}^0(k)}
$$

Boundary conditions

- **2**nd order linear differential e.o.m.,
	- two integration constants(boundary conditions)

At the horizon $\phi = (1-\nu)^{-i w/4\pi T}$ At the boundary : $\dot{\Phi} = \dot{\Phi}_0$

Black hole eat everything!

Bottom up approach

Five dimensional action : Einstein + Maxwell + cosmological constant

$$
S = \frac{1}{2G_5^2} \int d^5x \sqrt{-g} (\mathcal{R} - 2\Lambda) + \frac{1}{4g_5^2} \int d^5x \sqrt{-g} F^2
$$

But $T \neq 0$, There is no hadrons in hQCD but quark-gluon plasma.

Recipe for Green function

EOM for perturbed fields : δg_{ij} , δA_t

$$
E''_{\alpha}(u) + P(\mathfrak{w}, \mathfrak{q}, u)E'_{\alpha}(u) + Q(\mathfrak{w}, \mathfrak{q}, u)E_{\alpha}(u) = 0
$$

$$
S_{bdry}[\phi_0] = \int \frac{d^4k}{(2\pi)^4} \phi_0(-k) \mathcal{G}(k, u) \phi_0(k) \Big|_{u=0}^{u=1}
$$

On-shell action and Green function

$$
G_{ij}^{R}(K) = -2\mathcal{G}_{ij}(K, u = 0) \quad i = j
$$

$$
= -\mathcal{G}_{ij}(K, u = 0) \quad i \neq j
$$

Hydrodynamics

- **Hydrodynamics describes the system** at large distance and time scale : (small w, k limit).
- **Basic equation :** conservation of energy momentum

$$
\partial_{\mu}T^{\mu\nu}=0
$$

$$
T^{\mu\nu} = (\epsilon + P)u^{\mu}u^{\nu} + Pg^{\mu\nu} - \sigma^{\mu\nu}
$$

$$
\sigma^{\mu\nu} = P^{\mu\alpha}P^{\nu\beta} \left[\eta \left(\partial_{\alpha}u_{\beta} + \partial_{\beta}u_{\alpha} - \frac{2}{3}g_{\alpha\beta}\partial_{\lambda}u^{\lambda} \right) + \zeta g_{\alpha\beta}\partial_{\lambda}u^{\lambda} \right]
$$

 $P^{\mu\nu}=q^{\mu\nu}+u^{\mu}u^{\nu}$

Hydrodynamics

Transport coefficient : viscosity, thermalization time, vorticity etc.

$$
\omega = c_s q - i\Gamma q^2 + \frac{\Gamma}{c_s} \left(c_s^2 \tau^{eff} - \frac{\Gamma}{2} \right) q^3 + \mathcal{O}(q^4) \quad \text{where} \quad \Gamma = \frac{\eta}{sT} \left(\frac{2}{3} + \frac{\zeta}{2\eta} \right)
$$

Hydrodynamics

Transport coefficient : viscosity, thermalization time, vorticity etc.

$$
E''_{\alpha}(u) + P(\mathfrak{w}, \mathfrak{q}, u)E'_{\alpha}(u) + Q(\mathfrak{w}, \mathfrak{q}, u)E_{\alpha}(u) = 0
$$

$$
E_{\alpha}(u) = (1 - u)^{-i\frac{w}{4\pi T}} F(u, w, q)
$$

$$
F(u, w, q) = F^{(0)}(u) + w F^{(1,0)}(u) + q^2 F^{(0,2)}(u) + O(w, q)
$$

Taylor expansion of the equation of motion in w, q near 0. The solution of each order in w, q is the $\,F^{(0)}(u)$, $F^{(i,j)}(u).$ Hydrodynamics: fundamental d.o.f. = densities of conserved charges

Need to add constitutive relations!

Example: charge diffusion

Conservation law

Constitutive relation

[Fick's law (1855)]

Diffusion equation

 $\partial_t i^0 + \partial_i i^i = 0$ $j_i = -D \partial_i j^0 + O[(\nabla)^2]$ $\partial_{t}j^{0} = D\nabla^{2}j^{0}$

Dispersion relation

 $\omega = -i D q^2 + \cdots$

Expansion parameters:

 $\omega \ll T$, $q \ll T$

Transport coefficients in N=4 SYM in the limit $N_c \rightarrow \infty$, $g_{YM}^2 N_c \rightarrow \infty$

- $\eta = \frac{\pi}{8} N_c^2 T^3 \left[1 + O\left(\frac{1}{(a^2 N_c)^{3/2}}, \frac{1}{N_c^2} \right) \right]$ • Shear viscosity
- Bulk viscosity $\zeta = 0$
- Charge diffusion constant
- Thermal conductivity
- Electrical conductivity

$$
D_R = \frac{1}{2\pi T} + \cdots
$$

$$
\frac{\kappa_T \mu^2}{\eta T} = 8\pi^2 + \cdots
$$

$$
\sigma = e^2 \frac{N_c^2 T}{16 \pi} + \cdots
$$

Viscosity/entropy ratio in QCD: current status

Theories with gravity duals in the regime where the dual gravity description is valid

Kovtun, Son & A.S; Buchel; Buchel & Liu, A.S

QCD: RHIC elliptic flow analysis suggests

QCD: (Indirect) LQCD simulations

H.Meyer, 0805.4567 [hep-th]

Trapped strongly correlated cold alkali atoms 6 Li, ${}^{40}K$ T.Schafer, 0808.0734 [nucl-th]

Liquid Helium-3

$$
\frac{\eta}{s} = \frac{1}{4\pi} \approx 0.08
$$

(universal limit)

 $0.08 < \frac{\eta}{s} < 0.16$ $1.2 T_c < T < 1.7 T_c$ $\left(\frac{\eta}{s}\right)_{\text{min}} \approx 0.5$ $\left(\frac{\eta}{s}\right)_{\text{min}} \approx 0.7$

Hydrodynamics, ex1)

Black D3 metric and T_{xy} mode ($\phi = \delta g_{xy}$)

EOM

Solution

Greens function $G_{xy,xy}^{R}(\omega,k) = -\frac{\pi^2 N^2 T^4}{4} i w$

 $\phi''_p - \frac{1+u^2}{u f} \phi'_p + \frac{w^2-q^2 f}{u f^2} \phi_p = 0$

 $f_n(z) = (1 - u^2)^{-iw/2} + O(w^2, q^2)$

Kubo formular

$$
\eta = -\lim_{\omega \to 0} \frac{1}{\omega} \operatorname{Im} G_{xy,xy}^{R}(\omega, 0) \qquad \eta = \frac{\pi}{8} N^2 T^3
$$

Shear viscosity

Eta/s

Hydrodynamics, ex2)

Black D₃ metric and R charge current $δ$ Ao mode

R-charge Diffusion constant

Hydrodynamics, ex3)

Black D3 metric and T_{01} mode, $a_0 = \delta g_{01}$

Spectral function

Imaginary part of Green function.

$$
G = \frac{1}{p - m - i\,\Gamma} = \frac{p - m + i\,\Gamma}{(p - m)^2 + \Gamma^2}
$$

$$
\chi = Im\,G = \frac{\Gamma}{(p - m)^2 + \Gamma^2}
$$

Maximum value $1/\Gamma$ at p=m

Correlator (Finite temperature)

N-point function = differentiating N times partition function with respect to the source

Spectral function and quasiparticles

 $\chi_{\mu\nu,\alpha\beta}(k)=\int d^4x e^{-ikx}\langle \left[T_{\mu\nu}(x)T_{\alpha\beta}(0)\right]\rangle=-2\,\mathrm{Im}\,G^R_{\mu\nu,\alpha\beta}(\omega,q)$

A: scalar channel

B: scalar channel - thermal part

C: sound channel

Spectral function (Im G)

In medium effects of QCD

Particle production In RHIC

In medium effects of QCD

The w, k is scale by temperature

$$
w = \frac{\omega}{2\pi T}, \qquad q = \frac{k}{2\pi T}
$$

By changing μ, we will see the density effect on hydrodynamic quantities and spectral function

RN AdS5

 ${\sf N}$

$$
\text{etric}\n\begin{cases}\nds^2 = \frac{r^2}{l^2} \left(-f(r)dt^2 + d\vec{x}^2 \right) + \frac{l^2}{r^2 f(r)} dr^2 \\
f(r) = 1 - \frac{ml^2}{r^4} + \frac{q^2 l^2}{r^6}, \quad A_t = -\frac{Q}{r^2} + \mu\n\end{cases}
$$

$$
\fbox{\underline{\hspace{1.5em}Temperature}}
$$

Chemical pot.

$$
T = \frac{r_+^2 f'(r_+)}{4\pi l^2} = \frac{r_+}{\pi l^2} \left(1 - \frac{q^2 l^2}{2r_+^6} \right)
$$

$$
\mu = \frac{4Qb^2}{l^4} = \frac{1}{2b} \frac{g_5 l}{G_5} \sqrt{\frac{3a}{2}}
$$

SO(p-1) classification

$$
AdSp+2 \rightarrow p+1 CFT
$$

SO(p) rotational sym. ---(wave propagating)--> SO(p-1) rotational sym.

 h_{ai} : bi-vector $SO(1,1)$ X $SO(p-1)$

$$
x^M = (x^{\mu}, r) = (x^{\alpha}, x^{\iota}, r)
$$

$$
a = t, x \qquad i = y, z
$$

Metric and gauge field perturbation : vector mode 1005.0200

$$
0 = h_t^{x''} - \frac{1}{u}h_t^{x'} - \frac{(1 - \frac{a}{2})^2}{uf}(\omega_0 h_z^x + \omega_0^2 h_t^x) - 3auB'
$$

\n
$$
0 = \mathfrak{q}_f h_z^{x'} + \omega_0 h_t^{x'} - 3awuB
$$

\n
$$
0 = h_z^{x''} + \frac{(f/u)'}{f/u}h_z^{x'} + \frac{(1 - \frac{a}{2})^2}{uf^2}(\omega^2 h_z^x + \omega_0 h_t^x)
$$

\n
$$
0 = B'' + \frac{f'}{f}B' + \frac{(1 - \frac{a}{2})^2}{uf^2}(\omega^2 - \mathfrak{q}^2 f)B - \frac{h_t^{x'}}{f}
$$

\n
$$
\text{With finite density } \mathbf{h} \text{ and } \mathbf{a} \text{ is coupled}
$$

\n
$$
\Psi_{\pm} = -\frac{f\tilde{\mathfrak{q}}}{\tilde{\mathfrak{w}}^2 - \tilde{\mathfrak{q}}^2 f}Z'_1 + \left(-3au\frac{f\tilde{\mathfrak{q}}^2}{\tilde{\mathfrak{w}}^2 - \tilde{\mathfrak{q}}^2 f} + C_{\pm}\right)B
$$

\n
$$
\Psi''_{\pm} + \frac{f'}{f}\Psi'_{\pm} + \left(\frac{\tilde{\mathfrak{w}}^2 - \tilde{\mathfrak{q}}^2 f}{uf^2} - \frac{f'}{uf} - \frac{C_{\pm}}{f}\right)\Psi_{\pm} = 0
$$

\n
$$
\text{master variable}
$$

Near boundary series solution of master variable

$$
\Psi_{\pm} = \hat{\Psi}_{\pm} (1 + \dots + \hat{\Pi}_{\pm} u + \dots)
$$

From the definition of master variables, get the transformation matrix R to convert original boundary value and their conjugates

$$
\begin{pmatrix}\n\Psi_{+} \\
\Psi_{-}\n\end{pmatrix} = R \begin{pmatrix}\n(Z_{1})'(u=0) \\
\hat{B}\n\end{pmatrix} = R \begin{pmatrix}\n(\tilde{\mathbf{w}}^{2} - \tilde{\mathbf{q}}^{2})\hat{Z}_{1} \\
\hat{B}\n\end{pmatrix} \qquad R = \begin{pmatrix}\n-\frac{\tilde{\mathbf{q}}}{\tilde{\mathbf{w}}^{2} - \tilde{\mathbf{q}}^{2}} C_{+} \\
-\frac{\tilde{\mathbf{q}}}{\tilde{\mathbf{w}}^{2} - \tilde{\mathbf{q}}^{2}} C_{-}\n\end{pmatrix}
$$
\n
$$
\begin{pmatrix}\n\Psi_{+} \Pi_{+} \\
\Psi_{-} \Pi_{-}\n\end{pmatrix} = R \begin{pmatrix}\n\pi_{Z} \\
\pi_{B}\n\end{pmatrix}
$$

$$
\begin{pmatrix} \pi_Z \\ \pi_B \end{pmatrix} = R^{-1} \text{Diag}(\Psi_+, \Psi_-) \begin{pmatrix} \Pi_+ \\ \Pi_- \end{pmatrix}
$$

$$
\boxed{\mathcal{G}_{x,x}=\frac{l}{4e^2b^2}\frac{C_+\hat{\Pi}_+-C_-\hat{\Pi}_-}{C_+-C_-}}
$$

Photoemission rate

th/0607237

$$
d\Gamma_{\gamma} = \frac{d^3k}{(2\pi)^3} \frac{e^2}{2|\mathbf{k}|} \eta^{\mu\nu} C_{\mu\nu}^{\langle\langle}(K)\rangle_{k^0 = |\mathbf{k}|}
$$

of emitted photon / Vol4 Leading in e[^]2 expandsion (e is the EM coupling)

Wightman function of EM currents (not time-ordered) $C_{\mu\nu}^{\le}(K) = \int d^4X \; e^{-iK \cdot X} \; \langle J_{\mu}^{\text{EM}}(0) J_{\nu}^{\text{EM}}(X) \rangle$

Other way to get the Wightman function
 $C_{\mu\nu}^<(K) = n_{\rm b}(k^0) \chi_{\mu\nu}(K)$ $\chi_{\mu\nu}(K) = -2\,{\rm Im}\, C_{\mu\nu}^{\rm ret}(K)$
 $n_{\rm b}(k^0) = 1/(e^{\beta k^0} - 1)$ $C_{\mu\nu}^{\le}(K) = n_{\rm b}(k^0) \chi_{\mu\nu}(K)$

 $\frac{d\Gamma_{\gamma}}{dk} = \frac{\alpha_{EM}}{\pi} k \eta^{\mu\nu} C^{<}_{\mu\nu}(K)|_{\omega=k}$

Photo emission rate

Photoemission rate

th/0607237 ph/0111107

3 flavor massless QCD (dashed curve) / SYM (Solid Blue)

SYM computation gives quite good hint of hot QGP

Photoemission rate

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N=4 Super Yang-Mills computation with small t' Hooft coupling

Photo emission rate

Photo emission rate with unit $\alpha_{\text{EM}}(N_c^2-1)T^3$ Maximum at $\left(\frac{d\Gamma_{\gamma}}{dk}\right)_{\text{max}} \approx 0.01567 \, \alpha_{\text{EM}} N_{\text{c}}^2 T^3$

Remarks

- **Holography: two point function of QGP**
- Small w, k limit : Hydrodynamics
	- transport coefficients
- **Full w, k regime : Spectral function**
	- quasi particle peaks, photoemission
- **Finite T and** μ **, RN AdS is holographic dual.**
- **Thermal photon production is affected by** density effects.