

Equations of state for cold dense matter in D4/D6 model

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based on the work with Y. Kim, Y. Seo and S. J. Sin

JHEP 1106 (2011) 011 [arXiv:1011.0868 [hep-ph]] & arXiv:1108.2751 [hep-ph]

and with Y. Kim, C. -H. Lee and M. -B. Wan

JHEP 1110 (2011) 111 [arXiv:1108.6139 [hep-ph]]

Holographic QCD

- Quantum ChromoDynamics
: Theory of Strong Interaction
⇒ Usual tools such as perturbation are not valid
- Holography
 $(d+1)$ -dim. (bulk) gravity $\sim d$ -dim. (boundary) field theory
ex) AdS/CFT J. M. Maldacena, Adv. Theor. Math. Phys. 2, 231 (1998)
 $\mathcal{N}=4$ SYM on 4 dim. \longleftrightarrow Type IIB SUGRA on $AdS_5 \times S^5$
- AdS/CFT dictionary

Gauge theory (boundary)	Gravity (bulk)
operator \mathcal{O}	field ϕ
source J	non-normalizable mode ϕ_0
expectation value $\langle \mathcal{O} \rangle$	normalizable mode
...	...

Holographic QCD

- Toward to QCD ??
 - * $\exists \Lambda_{\text{QCD}} \Rightarrow$ Not conformal
 - * No SUSY
 - * Fermions in *fundamental* representation

⇒ To find an appropriate *dual* geometry to real QCD !!
- Holographic QCD
 - * Bottom-up approach :
 - Looking at QCD, gathering field contents
 - constructing (DBI) action
 - :: Hard-wall model, Soft-wall model, ...
 - * Top-down approach :
 - From string theory, setting brane configuration with DBI action
 - reproducing QCD-like theory
 - :: D3/D7, **D4/D6**, Sakai-Sugimoto (D4/D8/ $\overline{\text{D}8}$), ...

From String to D-brane

- String

- * one-dimensional object

$$X^\mu(\tau, \sigma) = X^\mu(\sigma^0, \sigma^1)$$

- * Nambu-Goto action

$$S_{\text{NG}} = -T \int dA = -T \int d\tau d\sigma \sqrt{(\dot{X} \cdot X')^2 - \dot{X}^2 X'^2}$$

using an induced metric γ_{ab} on the world-sheet

$$\gamma_{ab} = \eta_{\mu\nu} \frac{\partial X^\mu}{\partial \sigma^a} \frac{\partial X^\nu}{\partial \sigma^b}$$

in the reparametrization invariant form

$$S_{\text{NG}} = -T \int d^2\sigma \sqrt{-\gamma} \quad \text{where } \gamma = \det(\gamma_{ab})$$

From String to D-brane

- D p -brane

- * multi-dimensional object

$$X^\mu(\sigma^0, \sigma^1) \rightarrow X^\mu(\sigma^0, \sigma^1, \dots, \sigma^p)$$

- * DBI action

$$S_{D_p} = -T_p \int d^{p+1}\sigma \sqrt{-\det P[g]}$$

using a pull-back metric $P[g]_{ab}$ on the world-volume

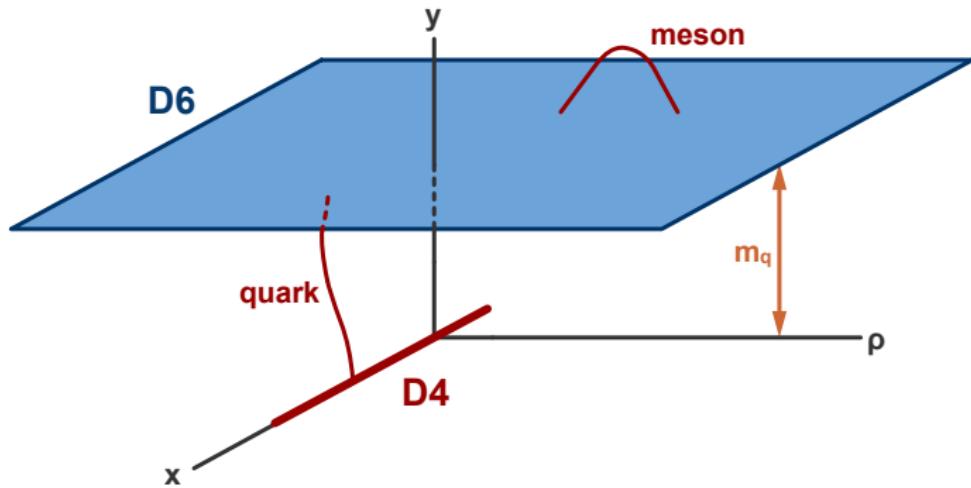
$$P[g]_{ab} = g_{\mu\nu} \frac{\partial X^\mu}{\partial \sigma^a} \frac{\partial X^\nu}{\partial \sigma^b}$$

including a gauge field A_a on the world-volume

$$S_{DBI} = -T_p \int d^{p+1}\sigma \sqrt{-\det (P[g] + 2\pi\alpha' F)}$$

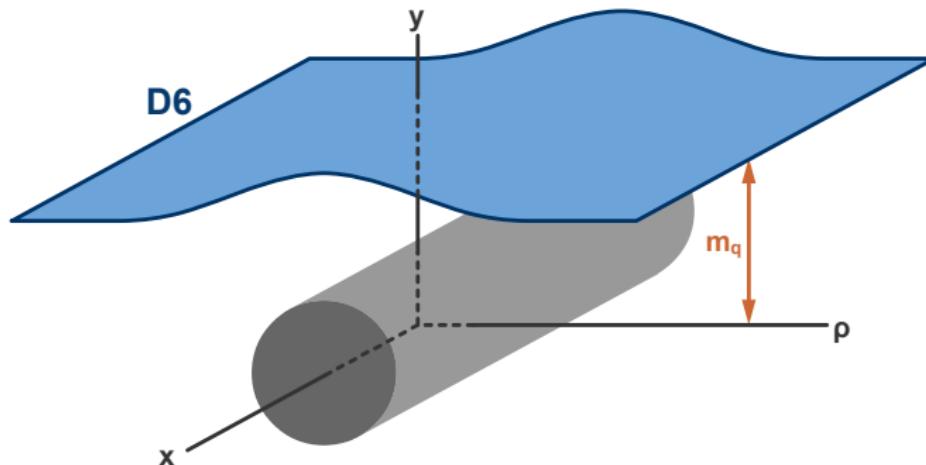
Configuration

- background D4 + probe D6



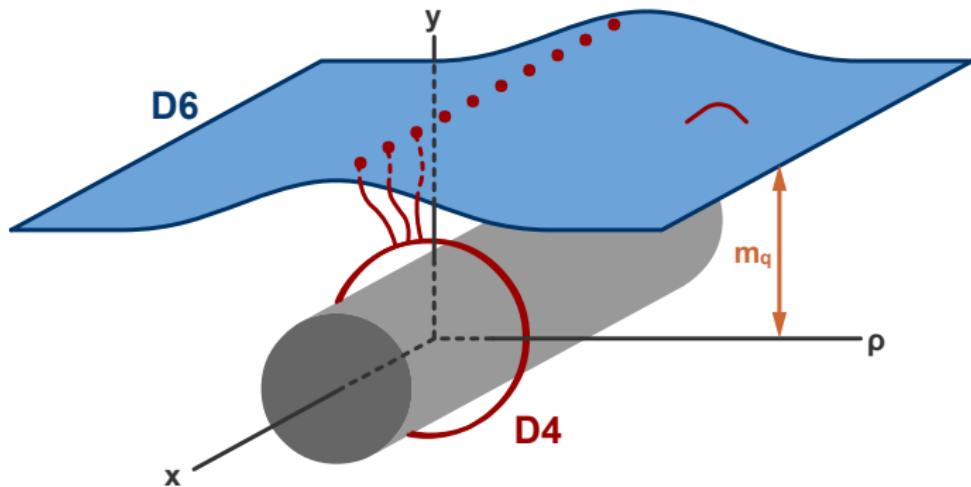
Configuration

- confining D4 background



Configuration

- add compact D4 branes \Rightarrow Baryon



D4/D6

- background D4 + probe D6

	t	1	2	3	(τ)	ρ	ψ_1	ψ_2	y	ϕ
D4	•	•	•	•	•					
D6	•	•	•	•		•	•	•		

The brane configurations : the background D4 and the probe D6

$$ds^2 = \left(\frac{U}{L}\right)^{3/2} (dx_{1,3}^2 + f(U)d\tau^2) + \left(\frac{L}{U}\right)^{3/2} \left(\frac{dU^2}{f(U)} + U^2 d\Omega_4^2 \right)$$

$$e^\phi = g_s \left(\frac{U}{L}\right)^{3/4}, \quad F_4 = dC_3 = \frac{2\pi N_c \epsilon_4}{\Omega_4}$$

$$\text{where } f(U) = 1 - (U/U_{\text{KK}})^3$$

D4/D6

- background D4 + probe D6

- * DBI action for D6 brane

$$\begin{aligned} S_{\text{D6}} &= -\mu_6 \int d^7\sigma e^{-\phi} \sqrt{-\det(P[g] + 2\pi\alpha' F)} \\ &= -\tau_6 \int dt d\rho \rho^2 (1 + 1/\xi^3)^{4/3} \sqrt{(1 + 1/\xi^3)^{4/3} (1 + y'^2) - \tilde{F}^2} \end{aligned}$$

- * equation of motion with respect to \tilde{A}_t

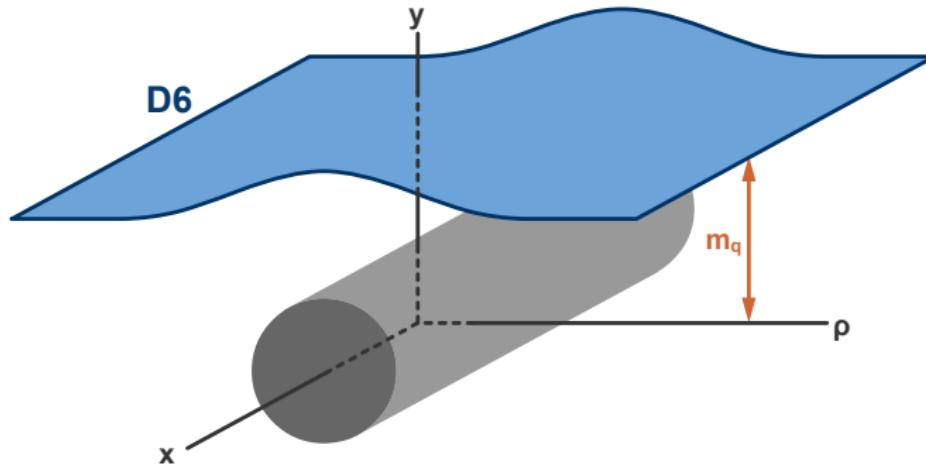
$$\frac{\partial \mathcal{L}_{\text{D6}}}{\partial \tilde{F}} = \frac{\rho^2 (1 + 1/\xi^3)^{4/3} \tilde{F}}{\sqrt{(1 + 1/\xi^3)^{4/3} (1 + y'^2) - \tilde{F}^2}} \equiv \tilde{Q}$$

- * through Legendre transformation

$$H_{\text{D6}} = \tau_6 \int d\rho \sqrt{(1 + 1/\xi^3)^{4/3} [\rho^4 (1 + 1/\xi^3)^{8/3} + \tilde{Q}^2] (1 + y'^2)}$$

D4/D6

- confining D4 background



$$\lim_{\rho \rightarrow \infty} y(\rho) \sim m_q + \frac{c}{\rho} + \dots$$

Finite density

- background D4 + probe D6 + compact D4

	t	1	2	3	(τ)	ρ	ψ_1	ψ_2	y	ϕ
D4	•	•	•	•	•					
D6	•	•	•	•		•	•	•		
	t	1	2	3	(τ)	U	Ω_4			
cD4	•						•	•	•	•

The brane configurations : the compact D4

- * baryon \sim compact D4 transverse to N_c D4 branes
on which N_c fundamental strings are attached

Finite density

- compact D4
 - * DBI action for compact D4 brane

$$\begin{aligned} S_{\text{D4}} &= -\mu_4 \int d^5\sigma e^{-\phi} \sqrt{-\det(P[g] + 2\pi\alpha' F)} \\ &= -\tau_4 \int dt d\rho \rho^2 (1 + 1/\xi^3)^{4/3} \sqrt{(1 + 1/\xi^3)^{4/3} (1 + y'^2) - \tilde{F}^2} \end{aligned}$$

- * equation of motion with respect to \tilde{A}_t

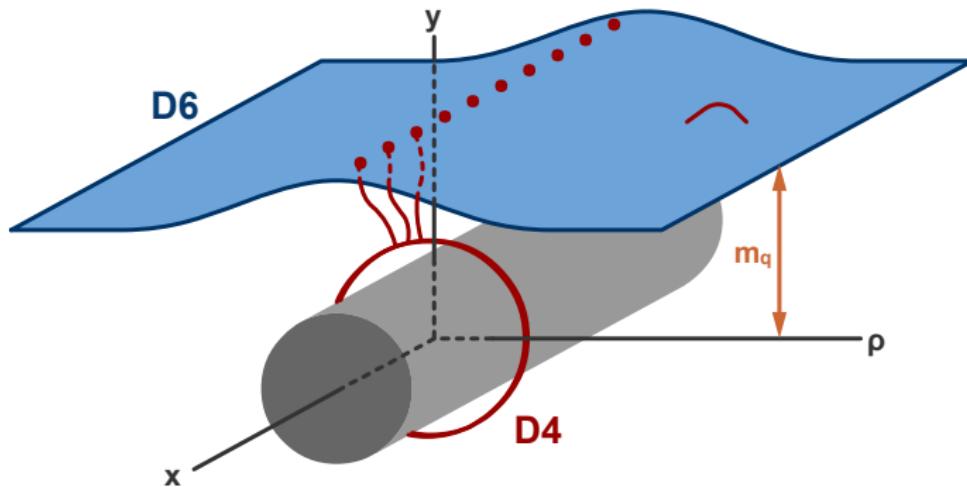
$$\frac{\partial \mathcal{L}_{\text{D4}}}{\partial \tilde{F}} = \frac{\sin^3 \theta \tilde{F}}{\sqrt{(1 + 1/\xi^3)^{4/3} (\xi^2 + \dot{\xi}^2) - \tilde{F}^2}} \equiv \tilde{D}$$

- * through Legendre transformation

$$H_{\text{D4}} = \tau_4 \int d\theta \sqrt{(1 + 1/\xi^3)^{4/3} (\xi^2 + \dot{\xi}^2)} \sqrt{\tilde{D}^2 + \sin^6 \theta}$$

Finite density

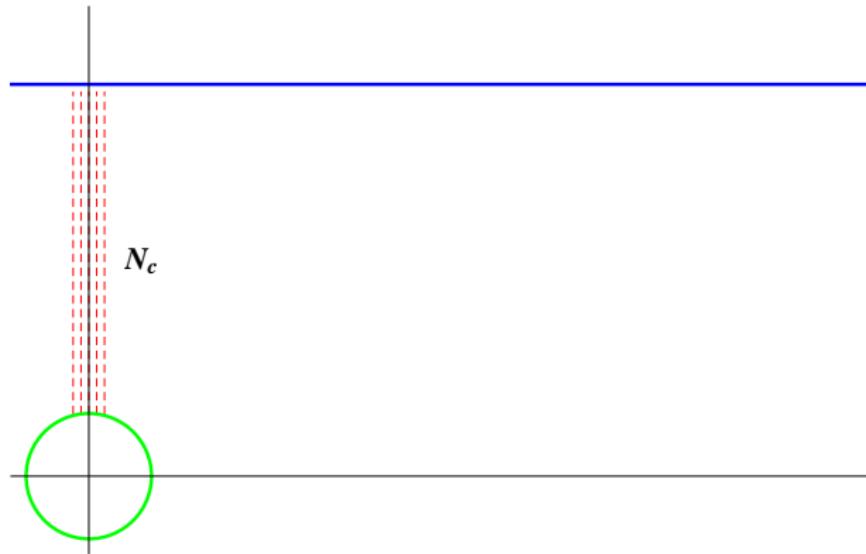
- add compact D4 branes \Rightarrow Baryon



$$\lim_{\rho \rightarrow \infty} A_t(\rho) \sim \mu - \frac{Q}{\rho} + \dots$$

Force balancing condition

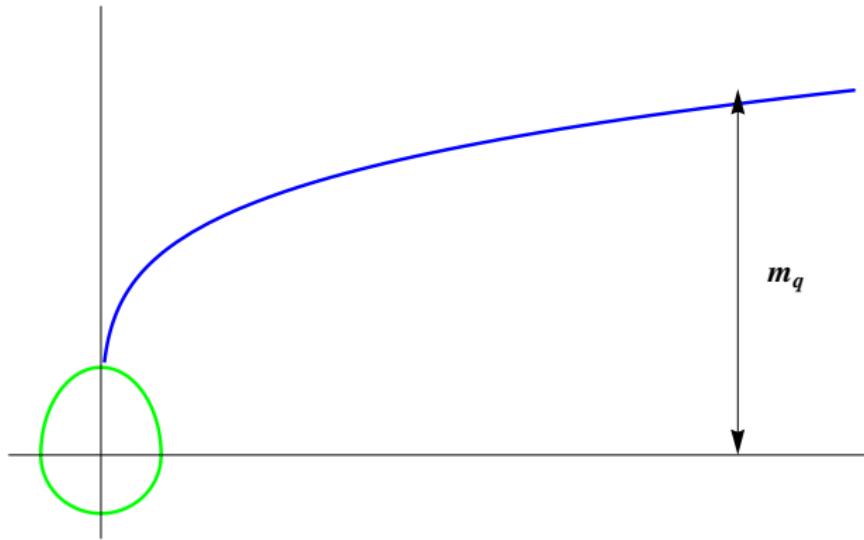
- confining D4 background



$$f \equiv \left. \frac{\partial H}{\partial \xi_c} \right|_{\text{fixed other values}}$$

Force balancing condition

- confining D4 background



$$\frac{Q}{N_c} f_{D4} = f_{D6}$$

Energy density ϵ and pressure p

- Gibbs relation:

$$\Omega = U - TS - \mu N = -pV$$

- The grand potential is identified with the on-shell DBI action:

$$\Omega = S_{\text{DBI}}|_{\text{on-shell}}$$

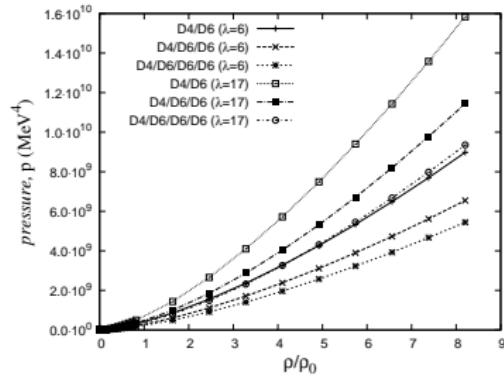
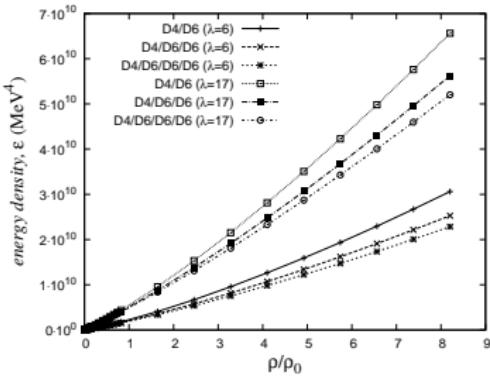
- Finally we can get the thermodynamic quantities:

$$\epsilon = \frac{1}{V_3} H|_{\text{on-shell}}$$

$$p = -\frac{1}{V_3} S_{\text{DBI}}|_{\text{on-shell}}$$

Holographic EoS

- Equation of States, $\epsilon(\rho)$ and $p(\rho)$, are given from D4/D6 model.

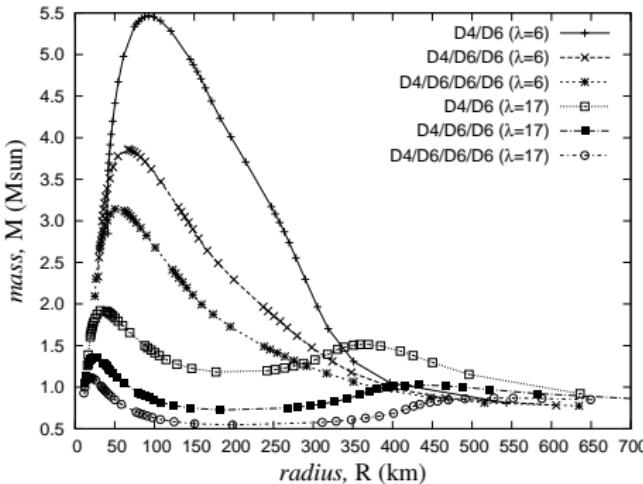


energy density $\epsilon(\rho)$ and pressure $p(\rho)$

Compact star with Holographic EoS

- TOV equation

$$\frac{dP(r)}{dr} = -\frac{G_N (\epsilon(r) + P(r)) (m(r) + 4\pi r^3 P(r))}{r (r - 2G_N m(r))}$$



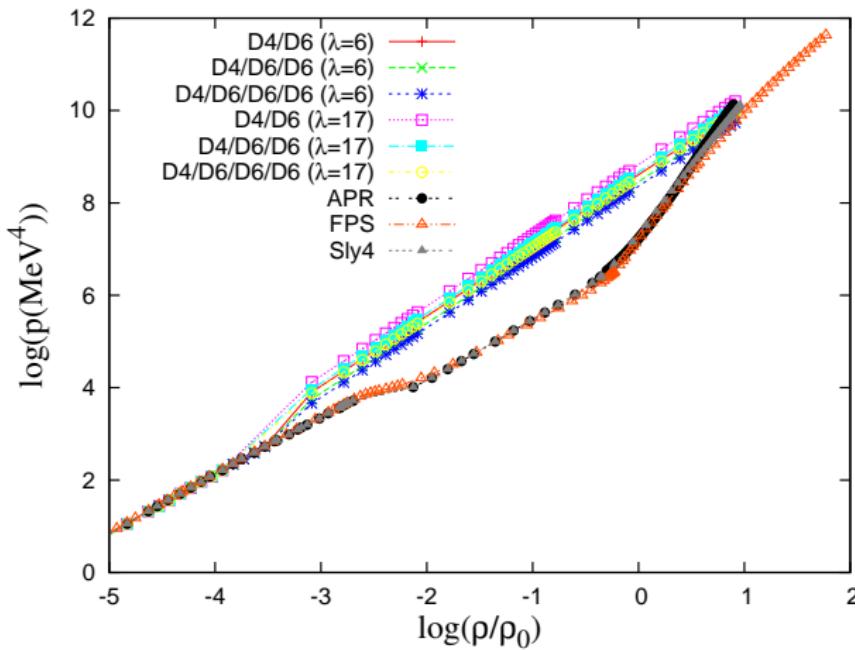
Mass-Radius relation using holographic EoS

Summary

- There is a restriction to deal with strongly coupled system perturbatively.
- AdS/CFT correspondence is a powerful theoretical tool to study strongly-coupled gauge theories.
- Here equation of states are obtained from D4/D6 model as an example.

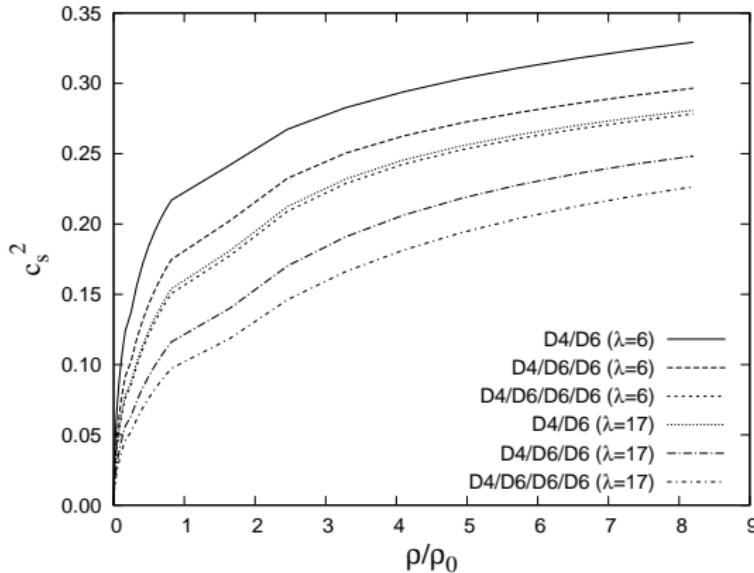
If needed

• Holographic EoS



If needed

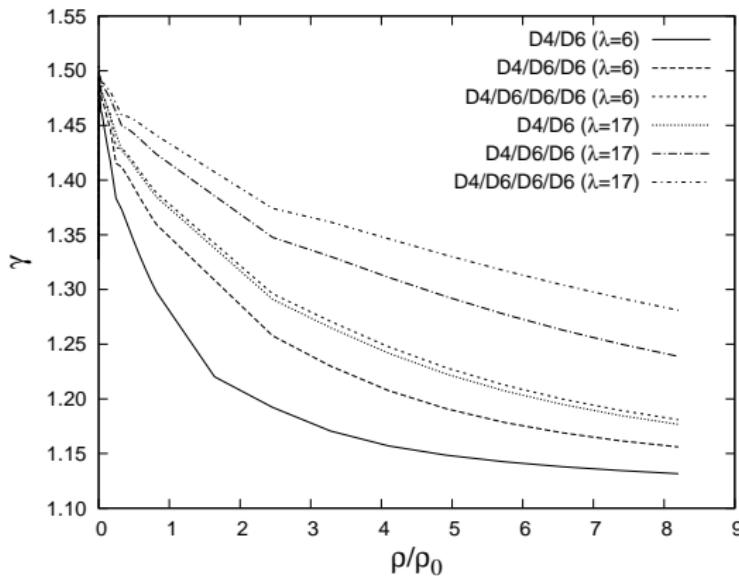
- Holographic EoS



$$c_s^2 \equiv dp/d\epsilon$$

If needed

- Holographic EoS



$$\gamma \equiv c_s^2(\epsilon/p)$$