

Isotope Shifts to Search for New Light Bosons

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Phys. Rev. D 96, 093001 (2017) w/ Ozeri, Perez, Soreq
Phys. Rev. Lett. **120**, 091801 (2018) w/ Berengut, Budker, Flambaum,
Frugiuele, Fuchs, Grojean, Harnik, Ozeri, Perez, Soreq
Phys.Rev.Res. 2, 043444 (2020) w/ Berengut, Geddes, Soreq

Light New Physics

- The SM **does not** completely describe Nature ($m_\nu, \eta_B, \Omega_{\text{DM}}$)
- The Higgs sector points to NP scales \sim TeV or (much) heavier, but **no experimental hint of it so far** (LHC, direct detection..)
- The motivation for light NP (below \sim GeV) is plenty:
 - Alternative solutions to the hierarchy problem? (like relaxion)
 - Axions
 - Light mediators for dark matter
 - ...
- If such NP couples significantly to **electrons/nucleons**, for $m \sim$ [eV,MeV] the **atom** is the natural place to search for it

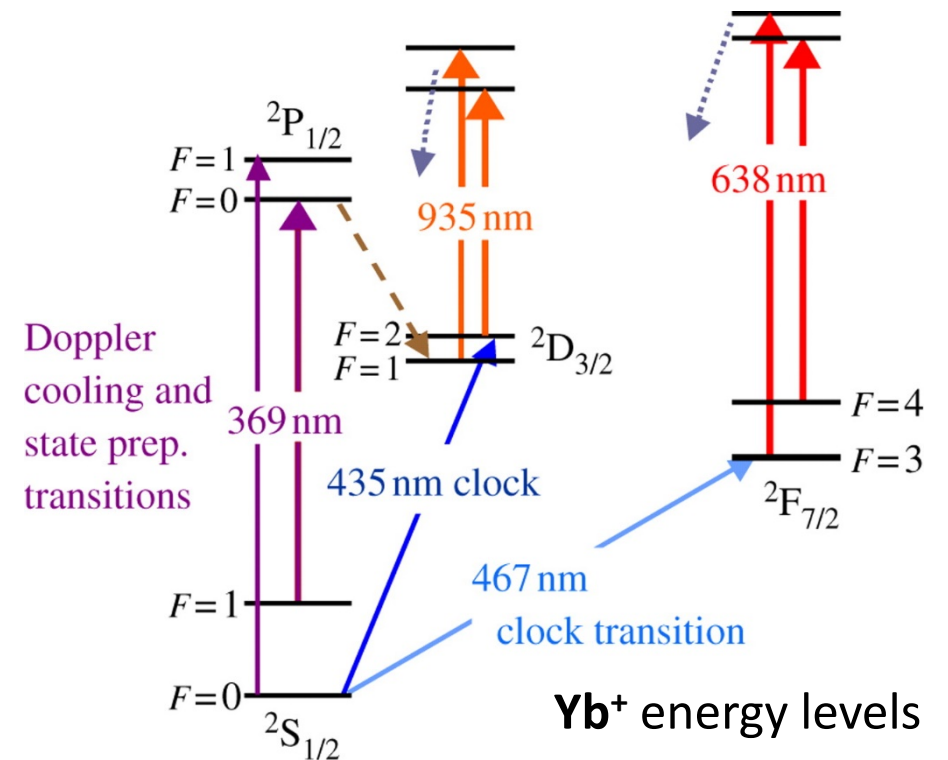
The Precision Atomic Frontier

- AMO techniques allow measurements of atomic frequencies with ultra high precision

- **Optical clock transitions**

$$\nu_{467\text{nm}} = 642\,121\,496\,772\,645.36(25) \text{ Hz}$$

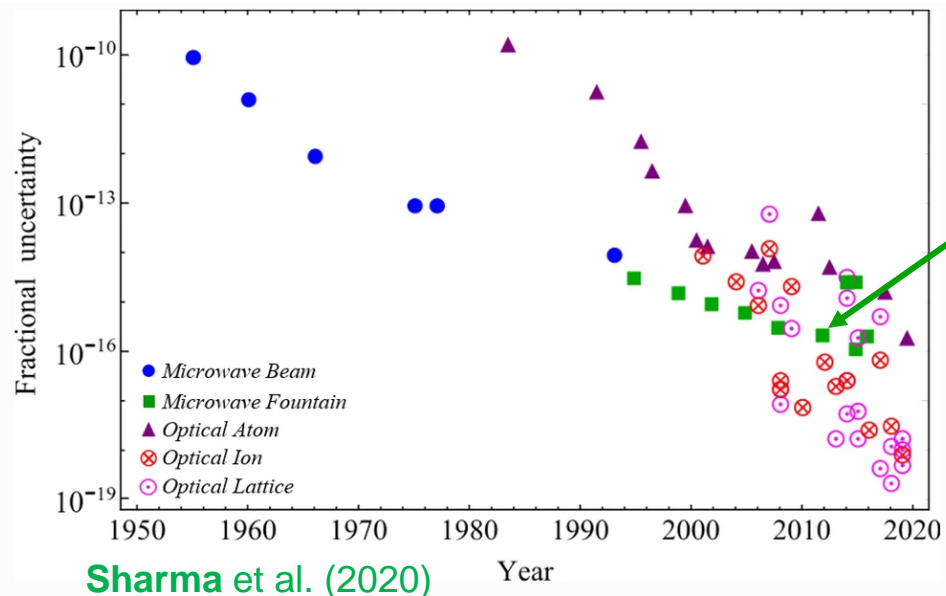
$$u_\nu = 3.9 \times 10^{-16} \quad \text{Huntermann et al. (2014)}$$



The Precision Atomic Frontier

- Precision on frequencies is not limited by technology, but by our definition of the second from the Cesium clock:

$1 \text{ Hz}^{-1} = 1 \text{ second} = 9\,192\,631\,770$ oscillations between the 2 hyperfine levels of the Cs 133 ground state.



Fractional uncertainty in Yb atom demonstrated at the 10^{-18} level

Huntermann et al. (2016)

Present record at 7.6×10^{-21} with lattice clock of 100 000 Sr atoms

Bothwell et al. (2022)

Why isotope shifts?

- Atomic lines are well measured, but theory is not accurate (~1%)
- Frequencies are set by the charge of the nucleus Z
- Hence dominant contributions from EM cancel out in frequency differences between isotopes: $(\nu - \nu')/\nu \sim 10^{-6}$
- Spin-independent NP couples to the entire nucleus A and thus is only mildly suppressed in isotope differences: $(A - A')/A \sim 0.1$
- Yet isotope shifts are also challenging to calculate...

Isotope Shift Theory

- Isotopes (same Z , different A) have the same atomic lines up to small nuclear effects:

$$\nu_i^{AA'} \equiv \nu_i^A - \nu_i^{A'} = K_i \mu_{AA'} + F_i \delta \langle r^2 \rangle_{AA'}$$

- **Mass shift** (MS) due to change in the nuclear mass $\mu_{AA'} \equiv m_A^{-1} - m_{A'}^{-1}$ modifying the global atomic center-of-mass (**normal MS**) and the electron-electron repulsion terms (**specific MS**)
- **Field shift** (FS) due to change in the nuclear charge distribution \sim size, which typically dominates for heavy elements

King Linearity

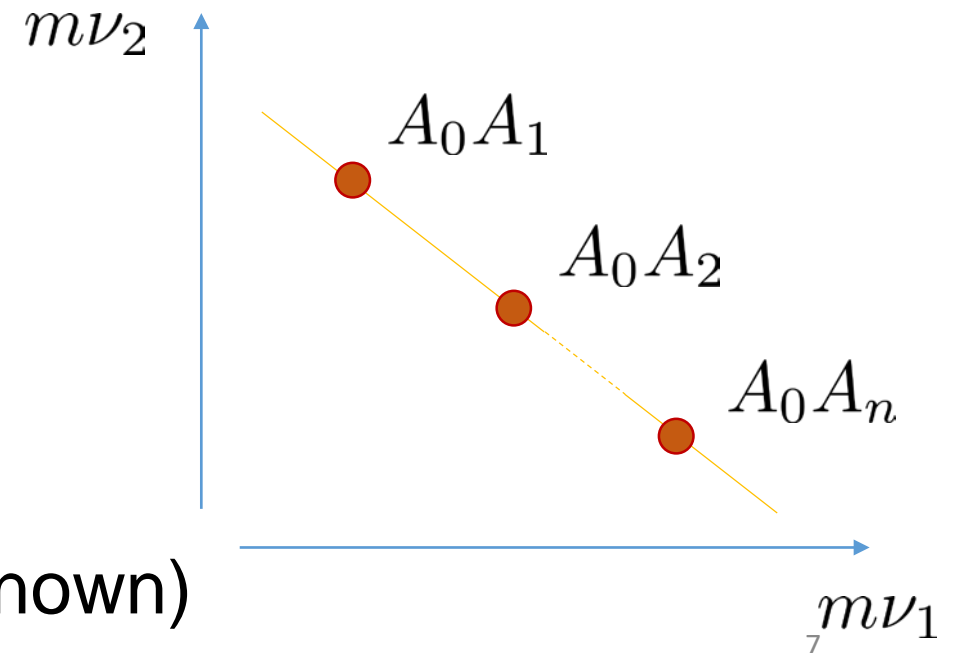
King (1963)

- Combining 2 transitions to eliminate the poorly known $\delta\langle r^2\rangle$ yields a **linear relation among modified IS**, $m\nu_i^{AA'} \equiv \nu_i^{AA'} / \mu_{AA'}$:

$$m\nu_2^{AA'} = F_{21} m\nu_1^{AA'} + K_{21}$$

$\equiv F_2/F_1$ $\equiv K_2 - F_{21}K_1$

- Linearity follows from having only 2 independent nuclear parameters
- $K, F, \delta\langle r^2\rangle$ are not needed to test it experimentally (only masses must be known)



New Physics breaks linearity

- **New Physics** comes with its own independent nuclear/electronic parameters:

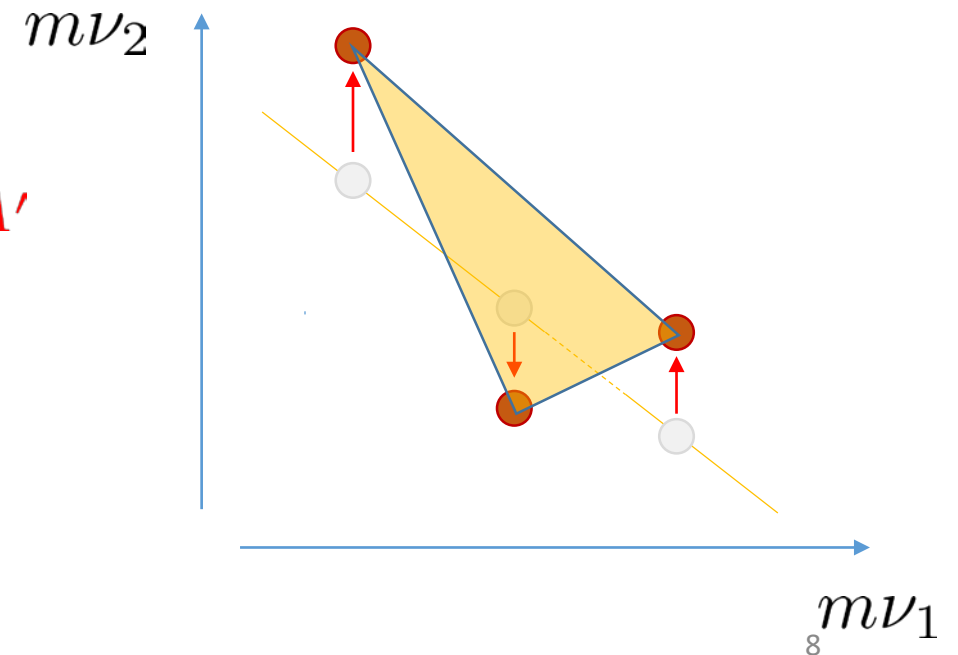
$$\nu_i^{AA'} = K_i \mu_{AA'} + F_i \delta \langle r^2 \rangle_{AA'} + \alpha_{\text{NP}} X_i \gamma_{AA'}$$

thus inducing **nonlinearities** (NL):

$$m\nu_2^{AA'} = F_{21} m\nu_1^{AA'} + K_{21} + \alpha_{\text{NP}} X_{21} h_{AA'}$$

with invariant measure : $\vec{m}\mu = (1, 1, 1)$

$$\text{NL} = \det(\vec{m}\nu_1, \vec{m}\nu_2, \vec{m}\mu)$$



Bounding NP

- Solving King equation for the NP coupling gives:

NL in data

$$\alpha_{\text{NP}} = \frac{\det(\vec{m}\nu_1, \vec{m}\nu_2, \vec{m}\mu)}{\det[X_1 \vec{m}\nu_2 - X_2 \vec{m}\nu_1, \vec{h}, \vec{m}\mu]}$$

NL predicted
by theory

$$= \epsilon_{ij} F_i X_j \times \det(m\delta\langle r^2 \rangle, \vec{h}, \vec{m}\mu)$$

Electronic alignment

→ strong suppression
for large $m_{\text{NP}} : X_i \propto F_i$

Nuclear alignment

→ suppression of $\delta m_A^{\text{max}} / m_A \sim \mathcal{O}(10)$
for NP coupling $\propto A$

New Yukawa-like force

- KL is sensitive to new **spin-independent** interaction between **electron and neutron** (e-p and e-e cancel out in the IS)
- Consider a new boson ϕ with y_e and y_n couplings
- ϕ -exchange in the atom gives rise to a new force described by the non-relativistic Yukawa potential:

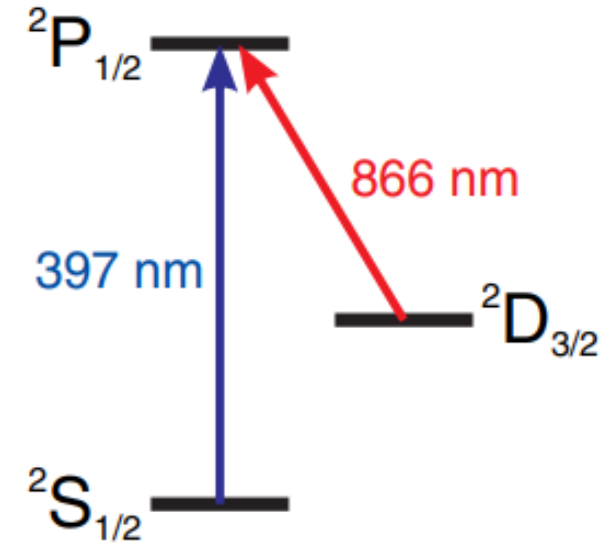
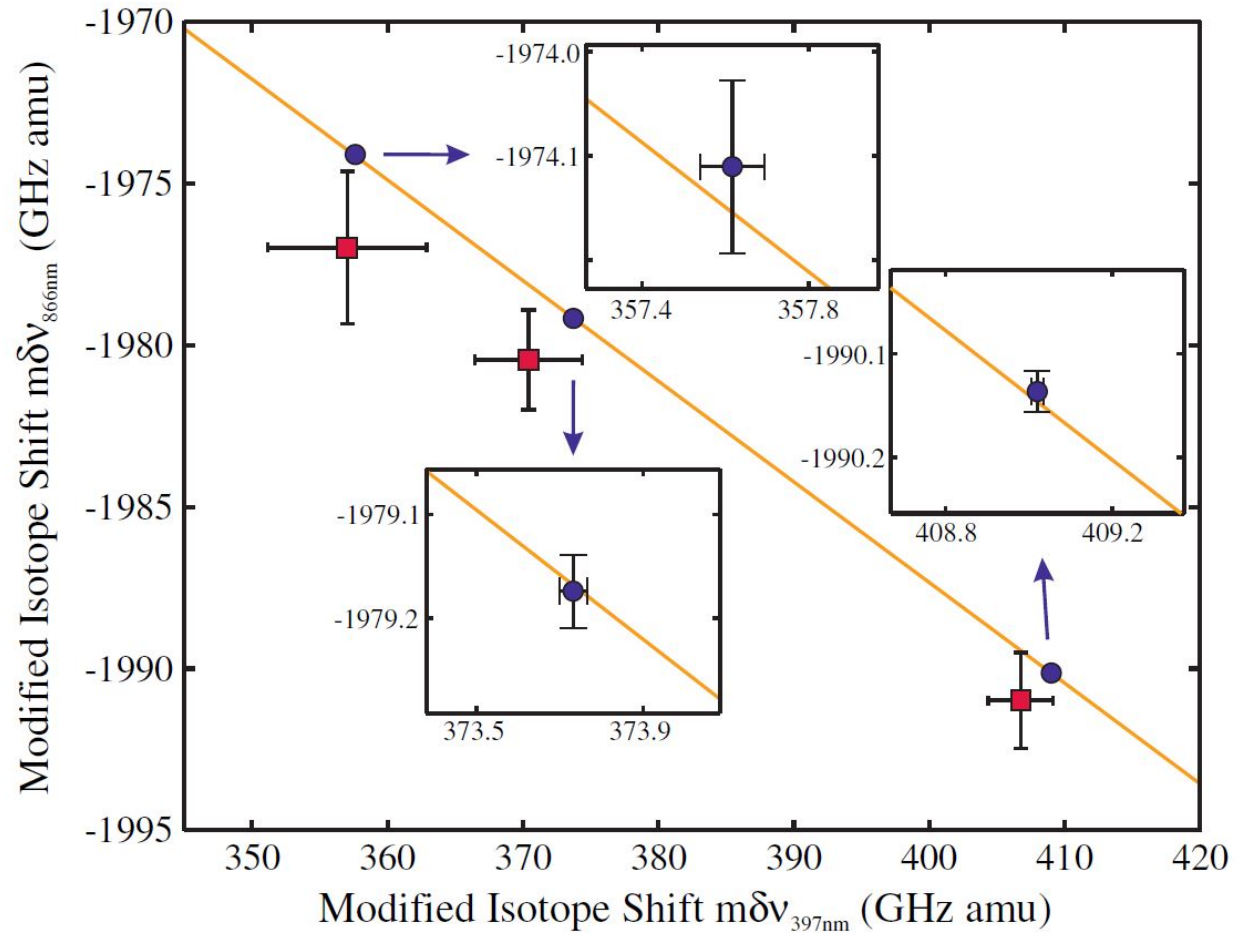
$$V_\phi(r) = \frac{(-1)^{s+1}}{4\pi} y_e y_n (A - Z) \frac{\exp(-m_\phi r)}{r}$$

$$\alpha_{\text{NP}}$$

$$\gamma_{AA'} = A - A'$$

Ca+ Example

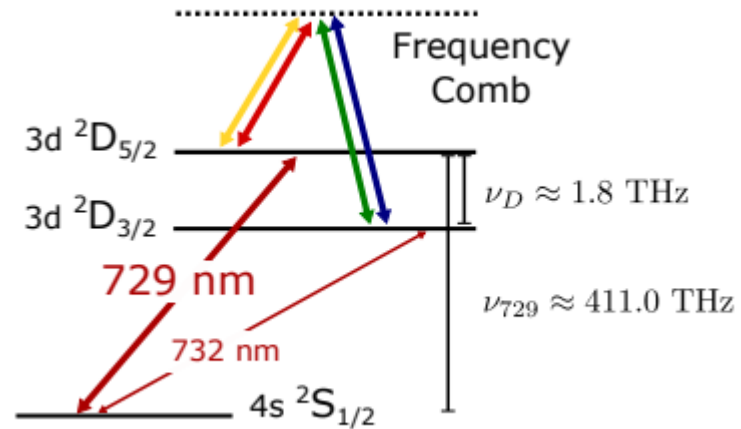
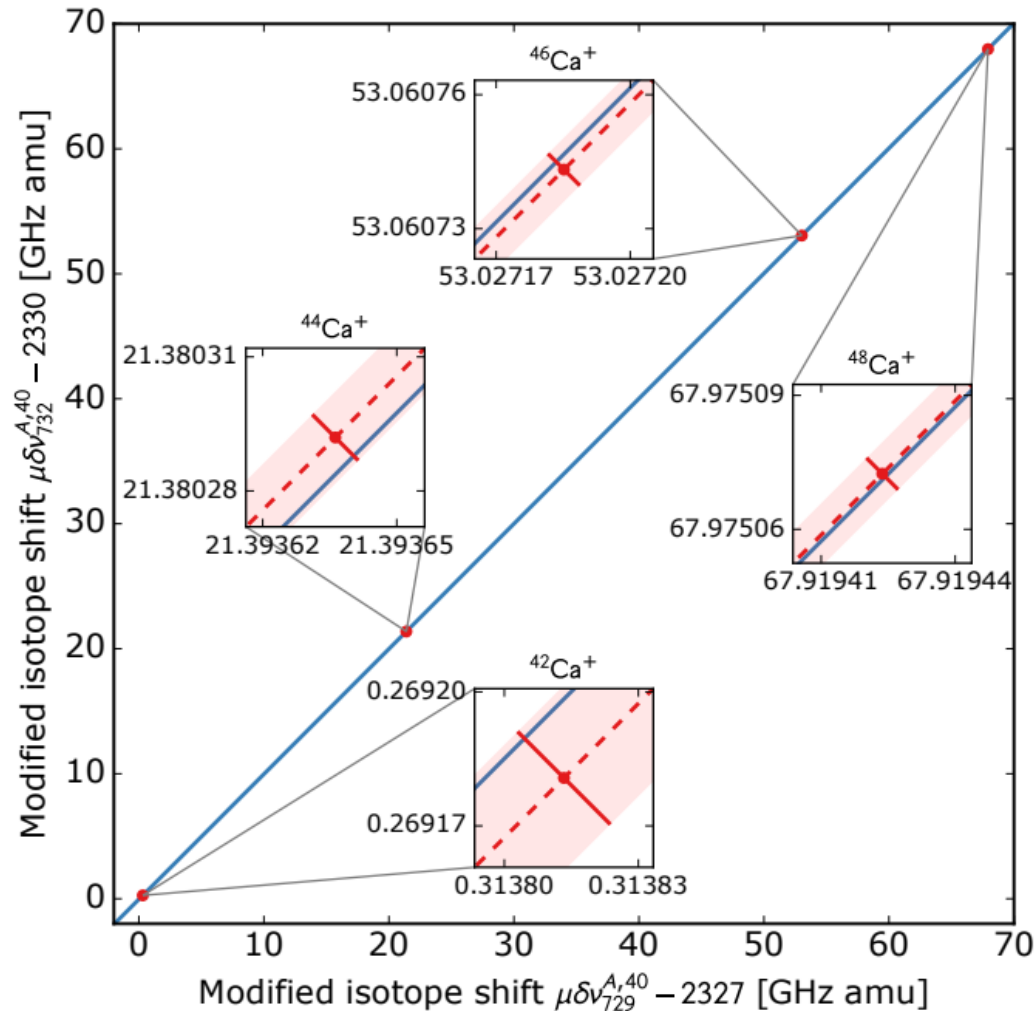
Gebert-Wan-Wolf-Angstmann-Berengut-Schmidt | Phys. Rev. Lett. 115, 053003 (2015)



based on broad transitions
precision $\sim 100\text{kHz} \sim IS/10^4$

Recent Ca+ results

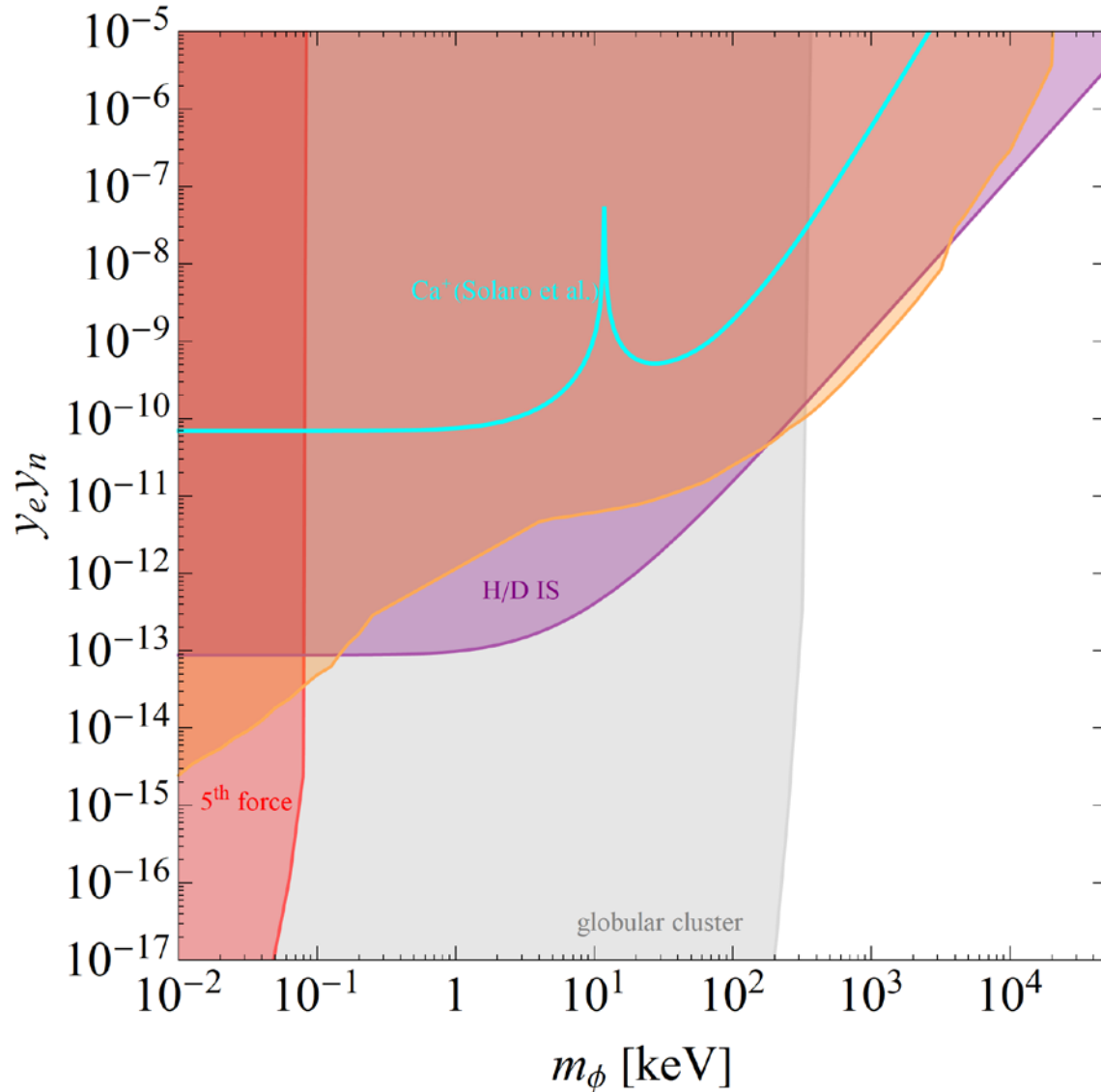
Solaro-Meyer-Fisher-Berengut-Fuchs-Drewsen | Phys. Rev. Lett. **115**, 123003 (2020)



narrower transitions
~ 20Hz accuracy (from DtoD)

consistent with linearity
within uncertainties

NP bound from Ca+



Resulting bound is weaker than

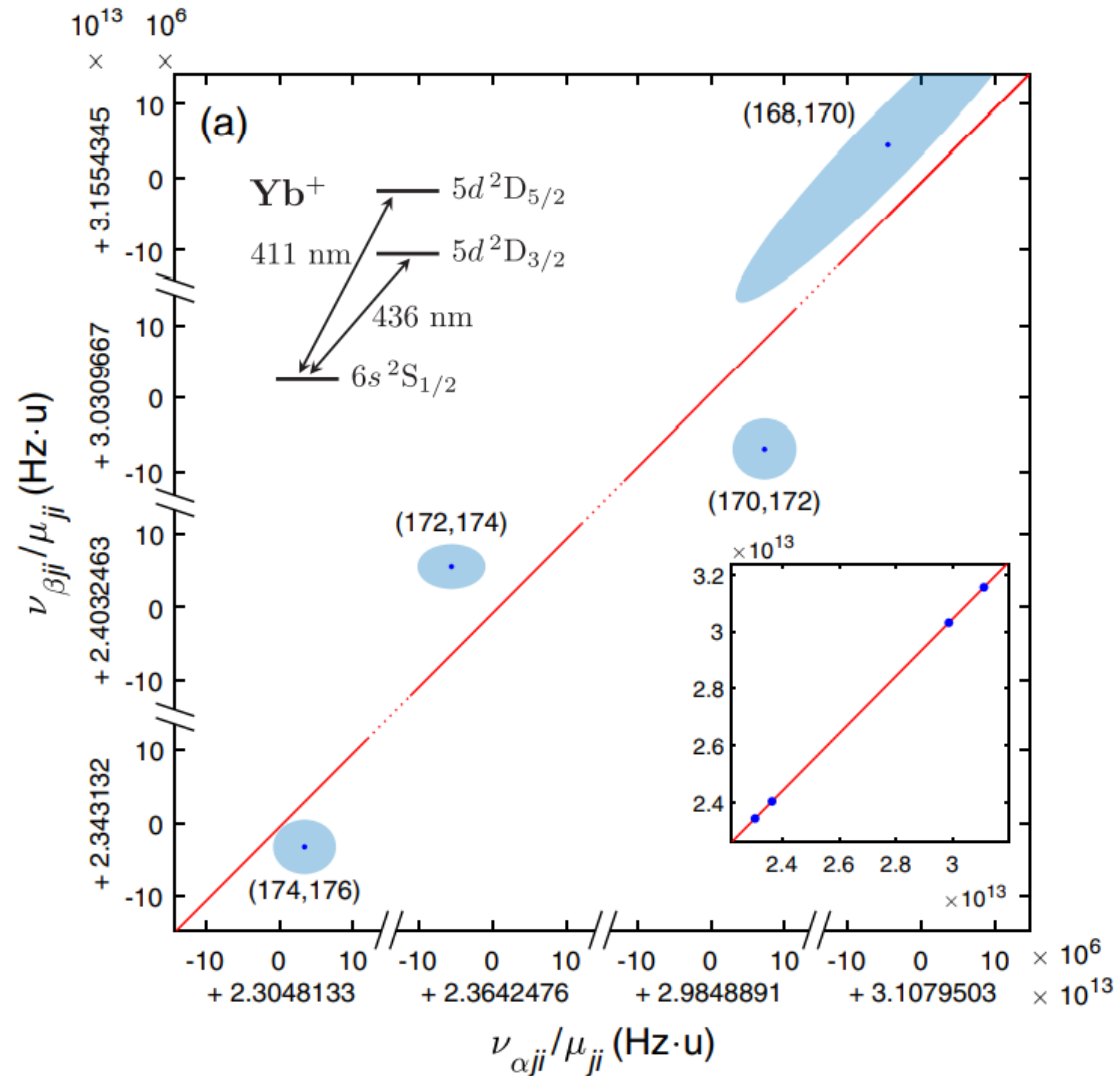
1) **Hydrogen-Deuterium IS** from a direct comparison between theory and measurements (**purple**)

CD-Frugiuele-Fuchs-Soreq (2017)

2) combination of bounds from **electron g-2** and **nuclear scattering** (**orange**)

Nonlinear Yb+ results

Counts-Hur-Aude Craik-Jeon-Leung-Berengut-Geddes-Kawasaki-Jhe-Vuletic | Phys. Rev. Lett. 125, 123002 (2020)



narrow quadrupole transitions
precision $\sim 300\text{Hz}$

nonlinearities observed at 3σ