

Identifying an event topology with a quantum annealer

Dong Woo Kang (KIAS)

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Leveraging **Quantum Annealer** to identify an Event-topology at High Energy Colliders

Minho Kim,^{1,*} Pyungwon Ko,^{2,†} Jae-hyeon Park,^{2,‡} and Myeonghun Park^{2,3,§}

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(Dated: November 15, 2021)

With increasing energy and luminosity available at the Large Hadron collider (LHC), we get a chance to take a pure bottom-up approach solely based on data. This will extend the scope of our understanding about Nature without relying on theoretical prejudices. The required computing resource, however, will increase exponentially with data size and complexities of events if one uses algorithms based on a classical computer. In this letter we propose a simple and well motivated method with a quantum annealer to identify an event-topology, a diagram to describe the history of particles produced at the LHC. We show that a computing complexity can be reduced significantly to the order of polynomials which enables us to decode the “Big” data in a very clear and efficient way. Our method achieves significant improvements in finding a true event-topology, more than by a factor of two compared to a conventional method.

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Quantum Annealer

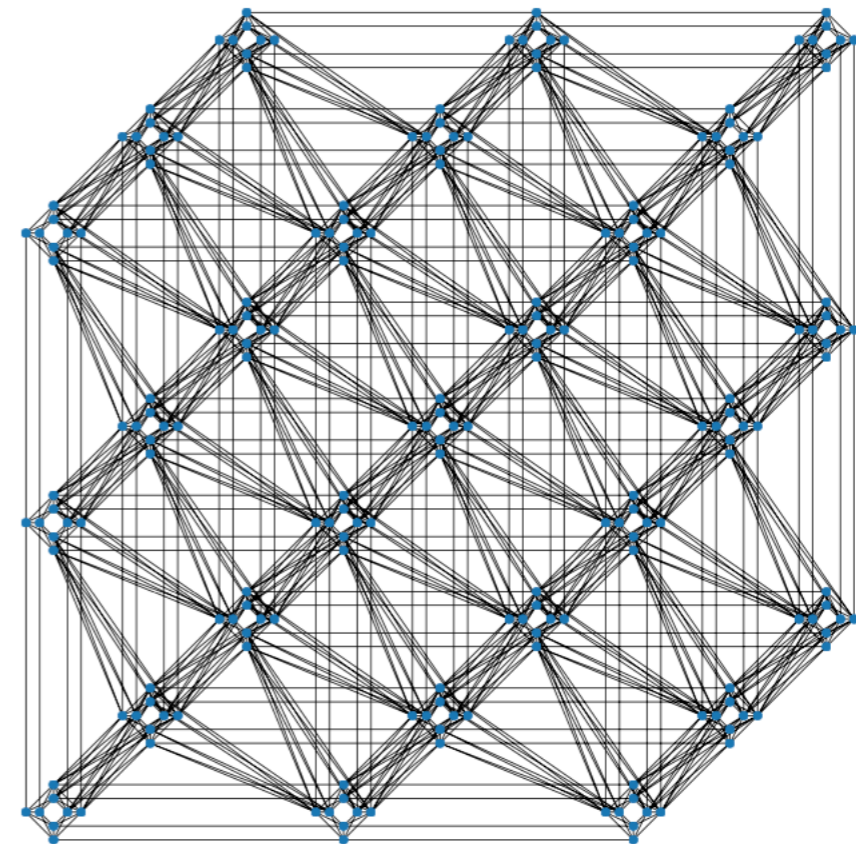
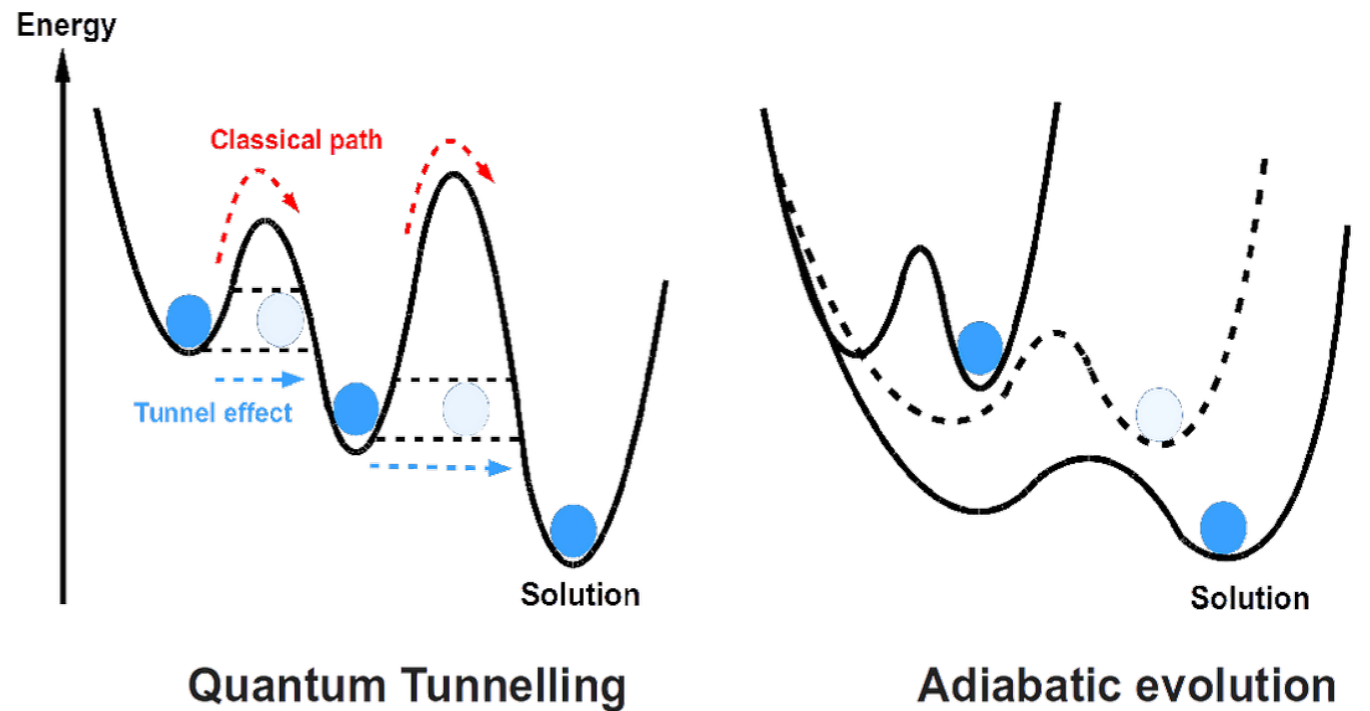
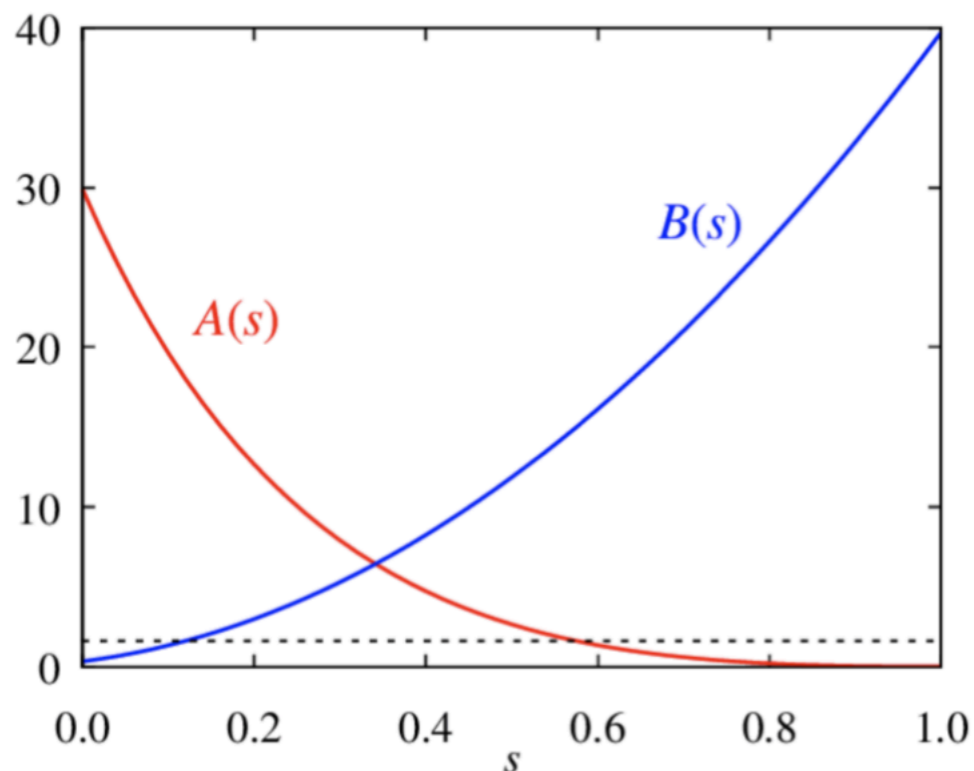


FIG. 7. A Pegasus graph with 27 unit cells.

Quantum Annealing

With adiabatic theorem, we can find the ground state of a complicated H_{QUBO} starting from simple Hamiltonian H_0

$$H_{\text{QA}}(t) = A(t) H_0 + B(t) H_{\text{QUBO}}$$



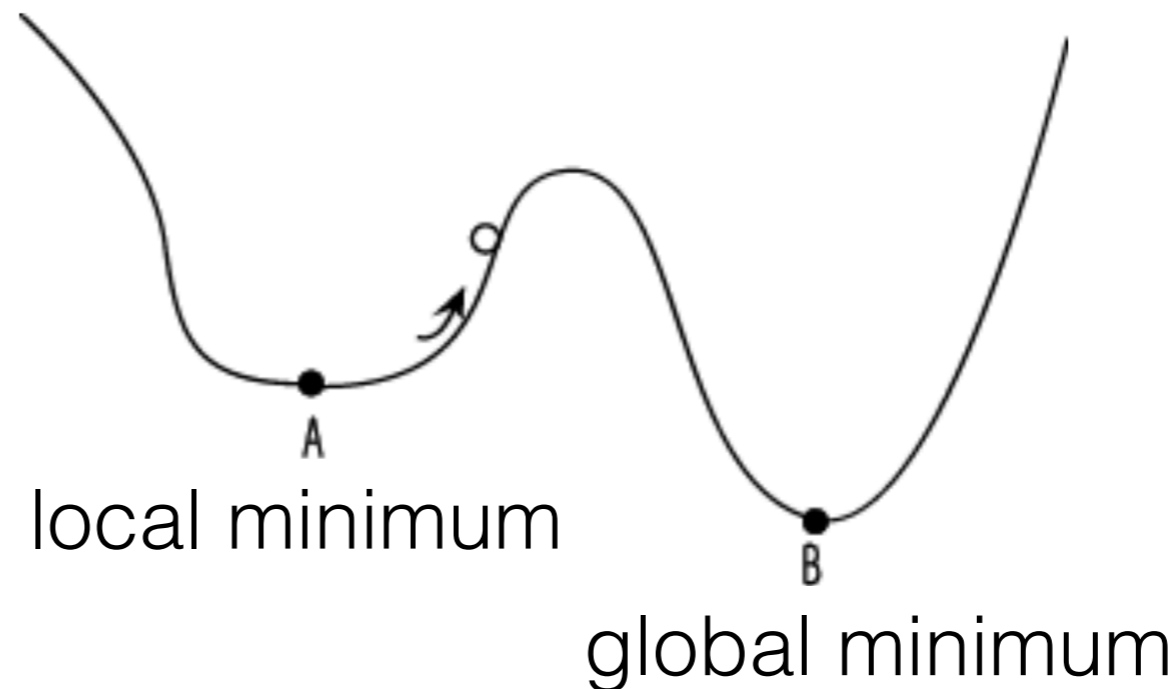
QUBO?

Claims to utilize “quantum tunneling” to find the minimum of the Hamiltonian

$$H_{\text{QUBO}} = \sum J_{ij} \sigma_i \sigma_j + \sum h_i \sigma_i$$

for **Quadratic Unconstrained Binary Optimization** problems

"Classic" minimization method (for Ising hamiltonian)

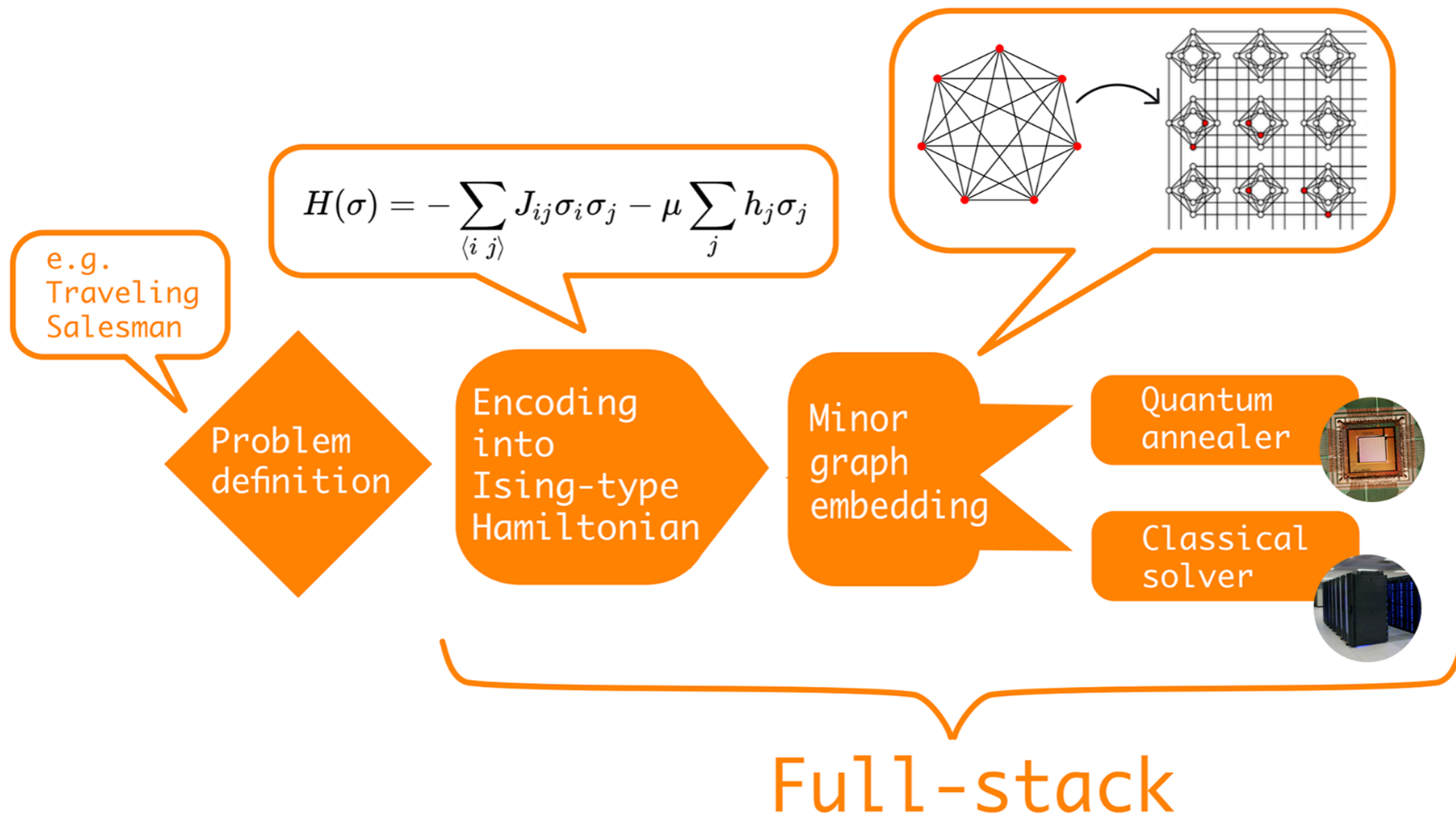


Simulated annealing

- Go to the next spin state $s_n \rightarrow s_{n+1}$
 - 1) If $E_n > E_{n+1}$: go to the lower energy
 - 2) If $E_n < E_{n+1}$, go with a probability of $e^{-\frac{E_{n+1} - E_n}{k_B T}}$ to **jump out**

(A "temperate $T \rightarrow 0$. With large T, SA can jump out local minimum)

QA workflow of Quantum Annealer



Problem in High Energy Physics

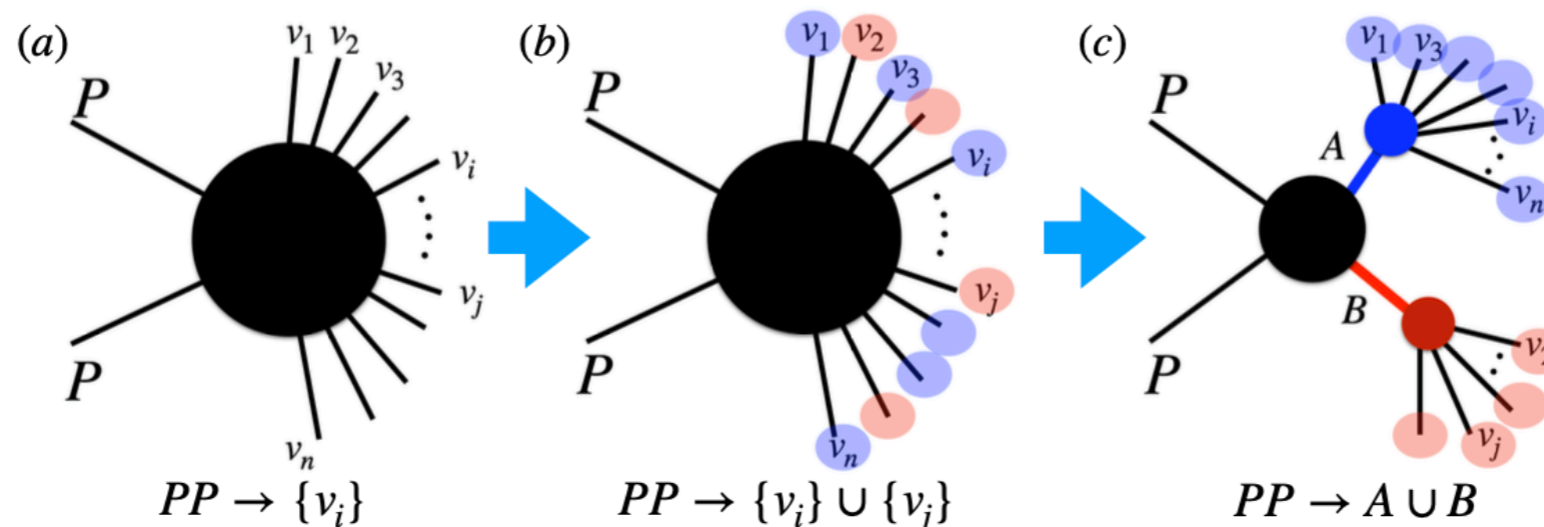


FIG. 1. (a) n -observed particles (b) Dividing n particles into two groups for $2 \rightarrow 2$ process (c) Identified event-topology with A and B .

Q1. What is the relevant event-topology behind anomalous events?

Q2. Mass spectrum?

Q3. Spin configuration?

Problem in High Energy Physics

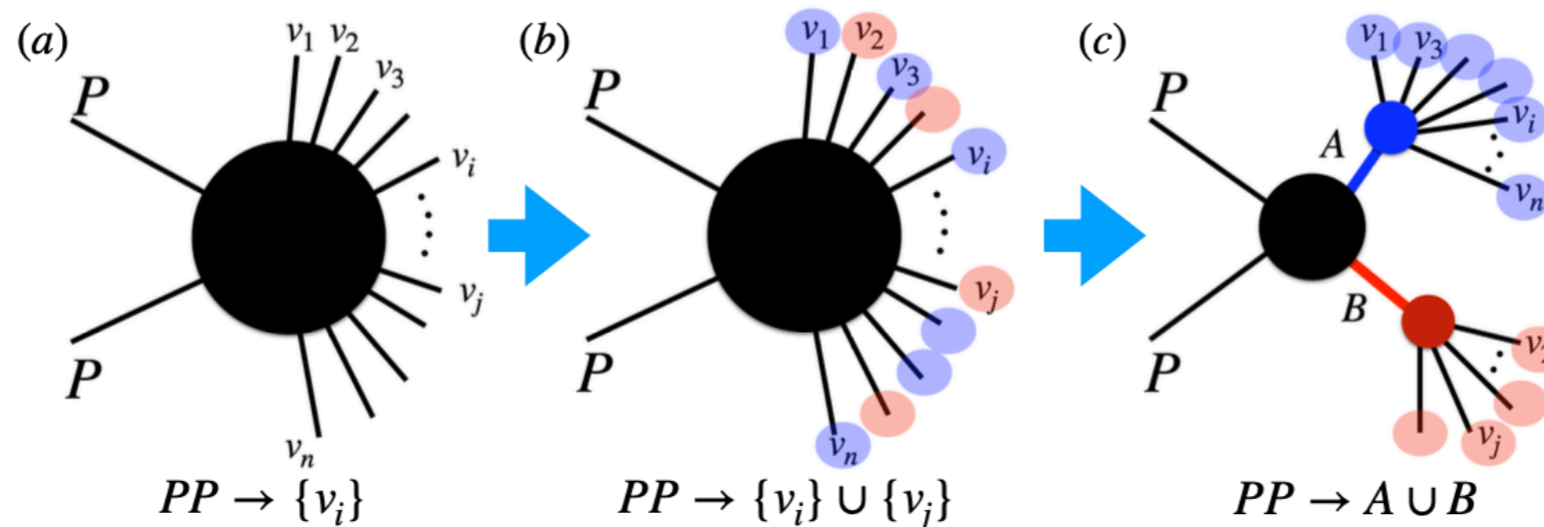
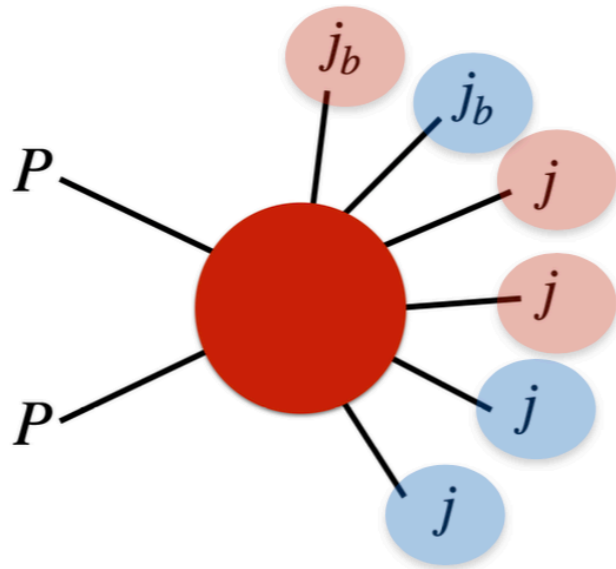


FIG. 1. (a) n -observed particles (b) Dividing n particles into two groups for $2 \rightarrow 2$ process (c) Identified event-topology with A and B .

1. Under the simple assumption $PP \rightarrow AB \rightarrow \{v_i\}\{v_j\}$
2. Find a right combination to reconstruct A and B particles
3. Read out the masses and spins from event reconstruction.

Example event topologies



- Standard example of six jets

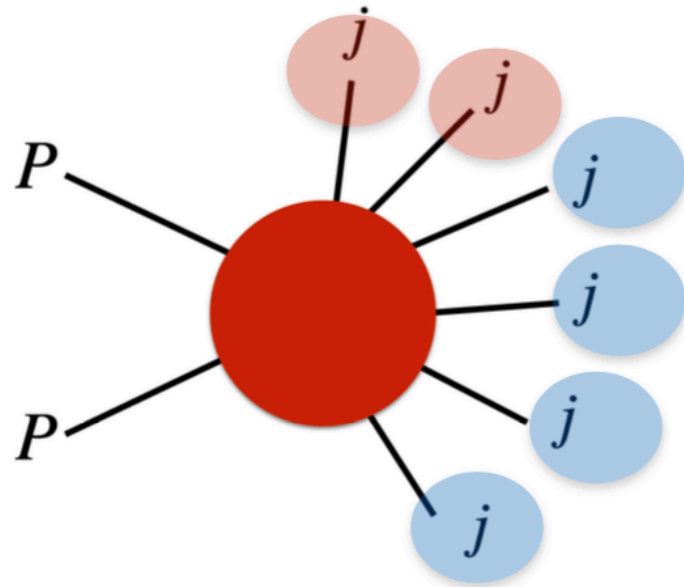
$$pp \rightarrow t\bar{t} \rightarrow \{j_b, (W \rightarrow jj)\} \cup \{j_b, (W \rightarrow jj)\}$$

(when A and B have same mass)

- Right answer is $(n_A, n_B) = (3, 3)$

$2^6 = 64$ cases,
no special assignment
for b-jet

Example event topologies



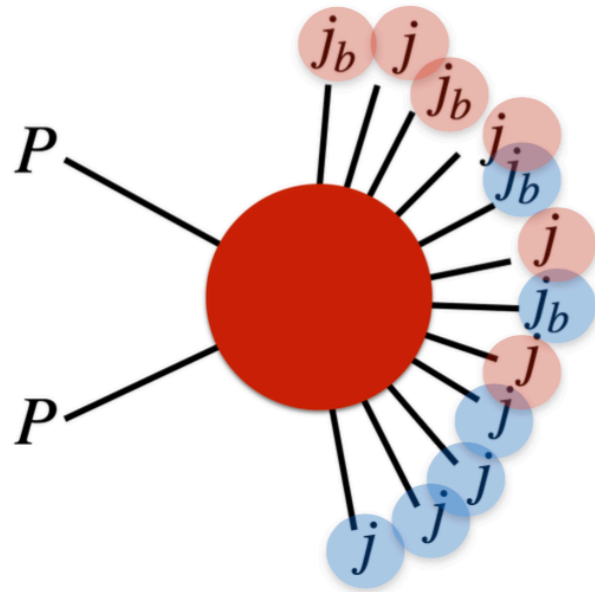
- Different mother particles

$$pp \rightarrow ZH \rightarrow \{j, j\} \cup \{(W \rightarrow jj), (W^* \rightarrow jj)\}$$

- Right answer is $(n_A, n_B) = (2, 4)$

$2^6 = 64$ cases,
no special assignment
for b-jet

Example event topologies



- Complicate situation (12 jets)

$$pp \rightarrow o\tilde{o} \rightarrow \{t, \bar{t}\} \cup \{t, \bar{t}\}$$

$$o \rightarrow t\bar{t} \rightarrow \{j_b, (W \rightarrow jj)\} \cup \{j_b, (W \rightarrow jj)\}$$

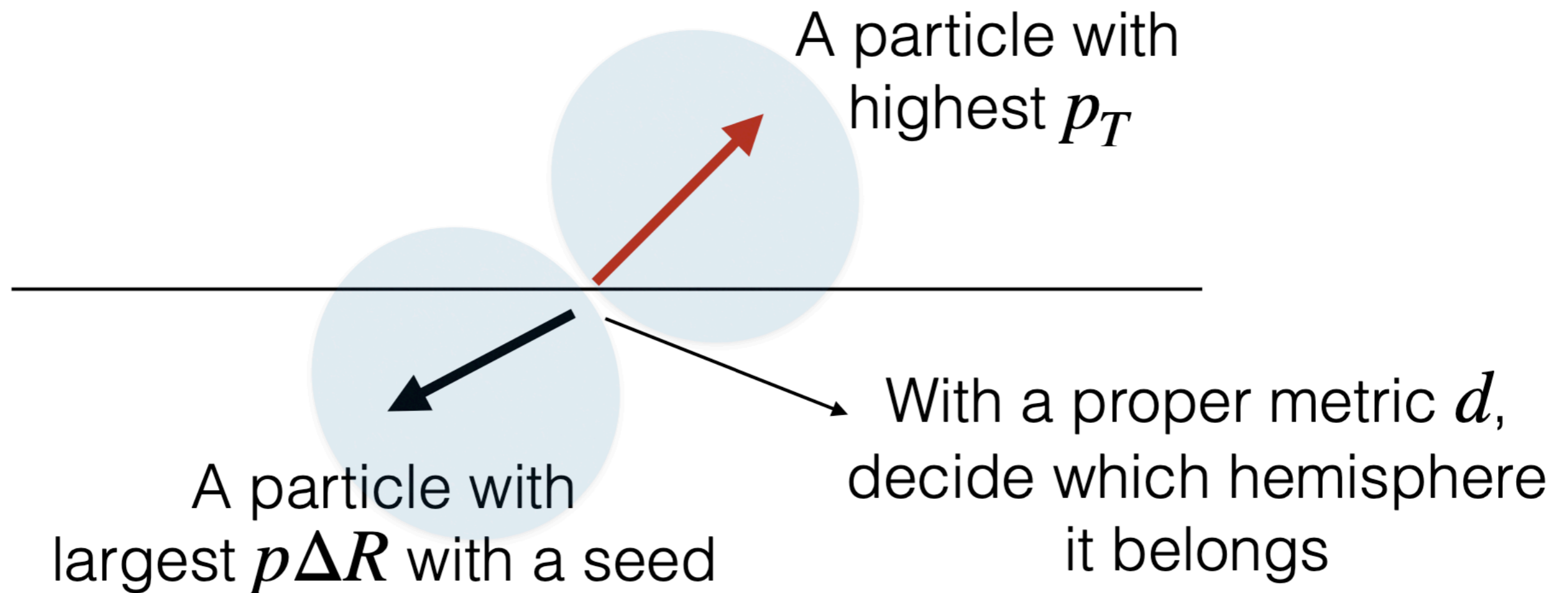
$$\tilde{o} \rightarrow t\bar{t} \rightarrow \{j_b, (W \rightarrow jj)\} \cup \{j_b, (W \rightarrow jj)\}$$

$2^{12} = 4096$ cases,
no prior knowledge on
a decay-structure

Combinatorial explosion exponentially

Classical algorithm

Hemisphere method: a **seed**-based method (iterative and converge)



CMS hemisphere TDR,
Shigeki Matsumoto, Mihoko M Nojiri, and Daisuke Nomura (2006)

Encoding in to H_{QUBO}

$$P_1 = \sum_i p_i x_i, \quad P_2 = \sum_i p_i (1 - x_i),$$

$$H = (P_1^2 - P_2^2)^2 = (M_A^2 - M_B^2)^2$$

Try to minimize the mass difference

How can we deal with the gas of $M_A \neq M_B$?

Add regularized term $\lambda(P_1^2 + P_2^2)$

Encoding in to H_{QUBO}

$$P_1 = \sum_i p_i x_i, \quad P_2 = \sum_i p_i (1 - x_i),$$

$$H = (P_1^2 - P_2^2)^2 + \lambda(P_1^2 + P_2^2)$$

$$H_{\text{QUBO}} = \sum J_{ij} \sigma_i \sigma_j + \sum h_i \sigma_i$$

$$J_{ij} = \sum_{kl} P_{ik} P_{jl},$$

$$h_i = 2 \sum_j [\sum_{kl} (P_{ik} P_{jl} - P_{kl} P_{ij})],$$

$$H_{\text{QUBO}} \rightarrow H_{\text{QUBO}} + \lambda(P_1^2 + P_2^2)$$

$$= H_{\text{QUBO}} + \lambda \sum_{ij} P_{ij} [s_i s_j + (1 - s_i)(1 - s_j)]$$

$$= \sum_{ij} J'_{ij} s_i s_j + \sum_i h'_i s_i, \quad (6)$$

with $J'_{ij} = J_{ij} + 2\lambda P_{ij}$ and $h'_i = h_i - 2\lambda \sum_j P_{ij}$. Here

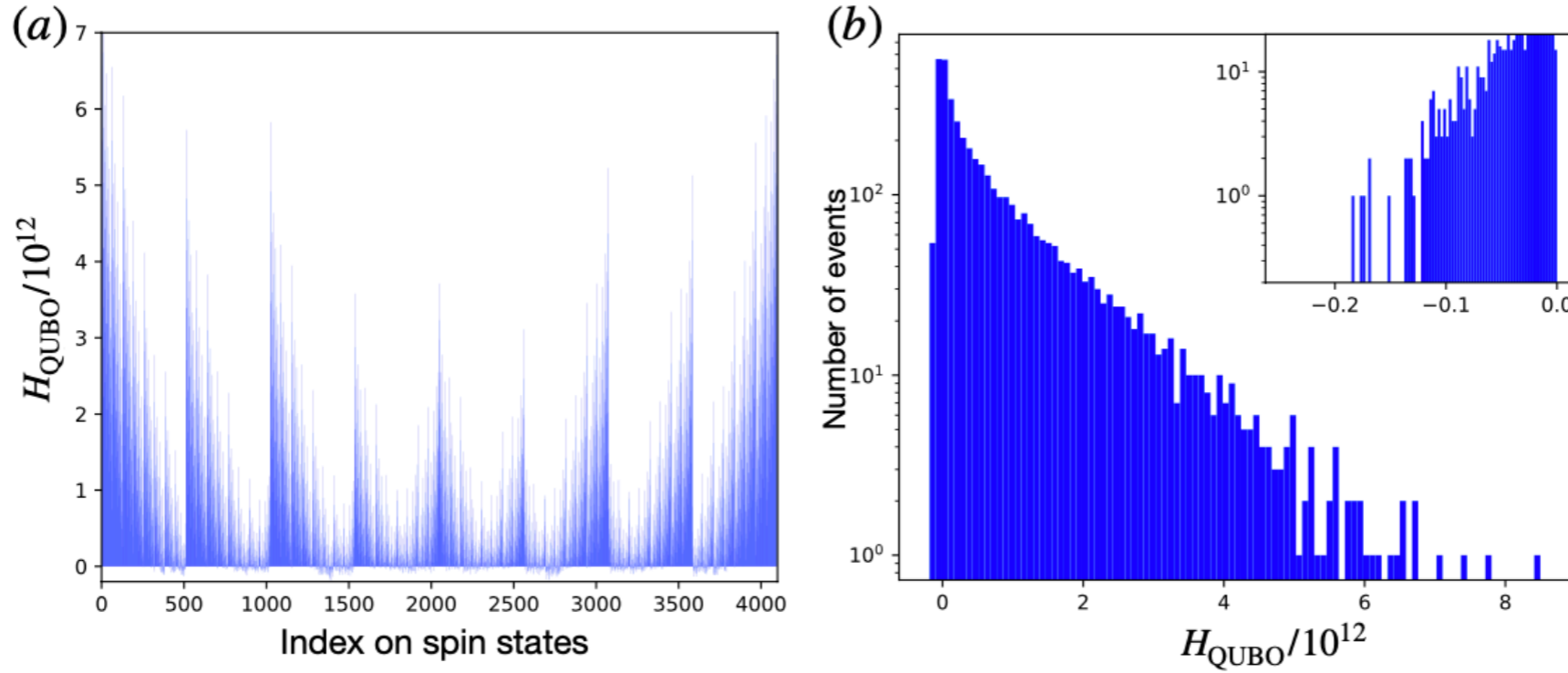


FIG. 2. We choose an event from detector level MC samples of EEq.q. (7c) to calculate (a) energy spectrum of H_{QUBO} with increasing indices of spin states, (b) histogram of energy spectrum of H_{QUBO} for all possible 2^{12} (= 4096) spin states.

Process	$pp \rightarrow t\bar{t}$ (2 → 6)	$pp \rightarrow HZ$ (2 → 6)	$pp \rightarrow \tilde{o}\tilde{o}^*$ (2 → 12)
Quantum annealing	100%	100%	74.3%
Simulated annealing	36.7%	45.7%	1%

Results

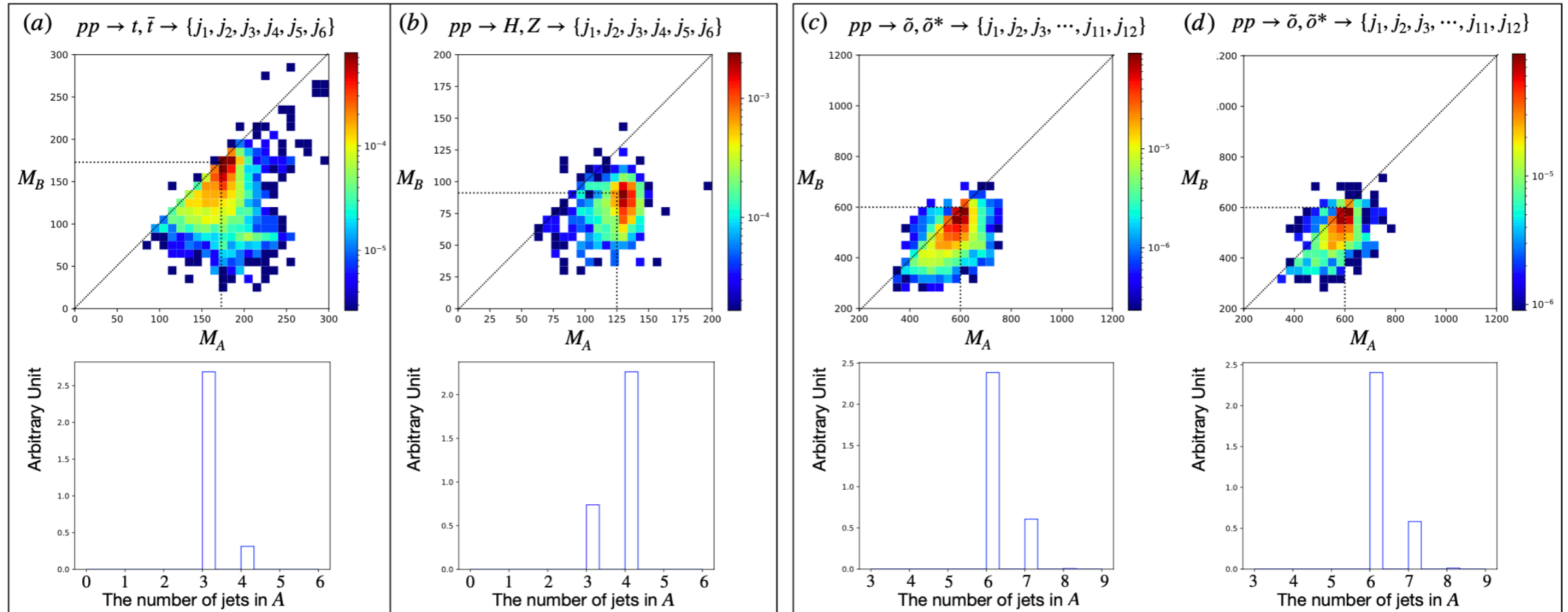
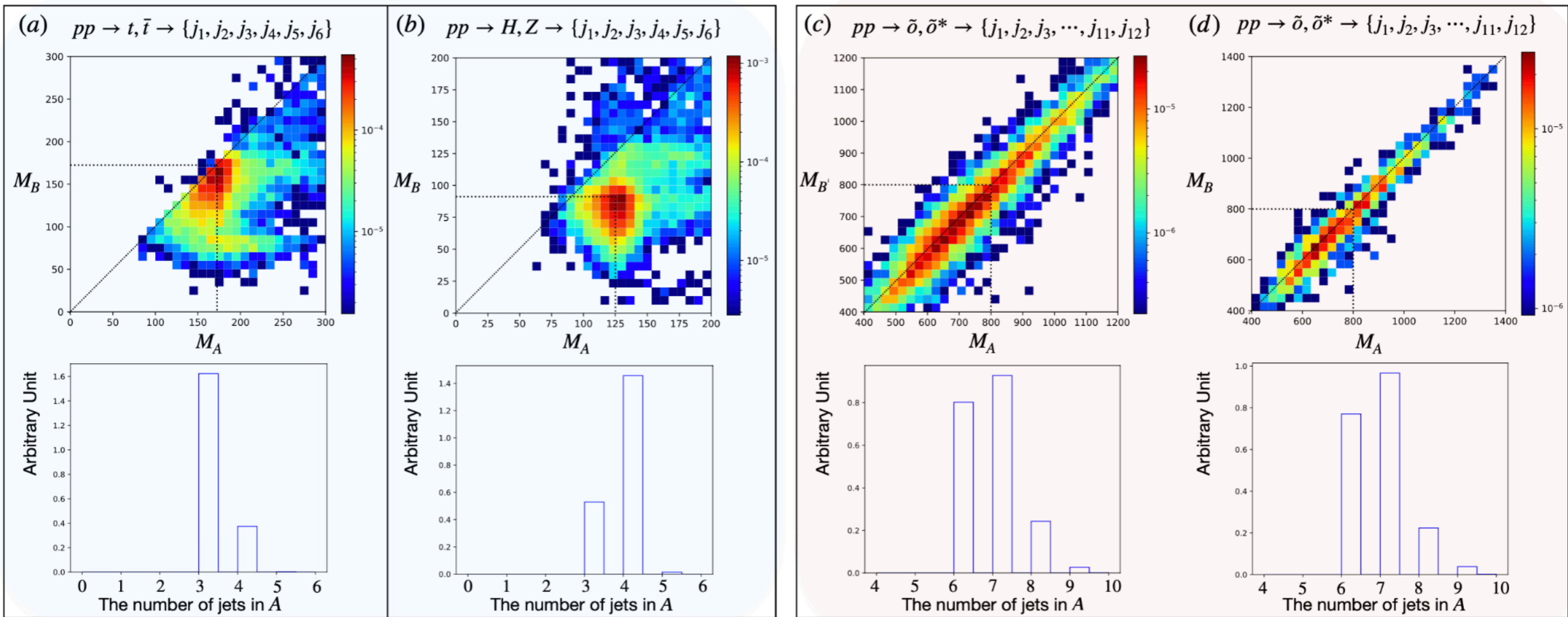


FIG. 3. **Common:** H_{QUBO} is calculated with Monte Carlo samples. To find a global minimum of H_{QUBO} , we use a brute force scanning in (a)-(c), and quantum annealer in (d). For (d), we randomly choose 1000 events from MC samples we used in (c). **Top:** Normalized density histogram for reconstructed mass M_A and M_B **Bottom:** The number of jets clustered into A .

Madgraph \rightarrow Pythia (ISR/FSR/MPI turned off) \rightarrow Delphes

Results



Madgraph \rightarrow Pythia (**ISR/FSR/MPI turned ON**) \rightarrow Delphes

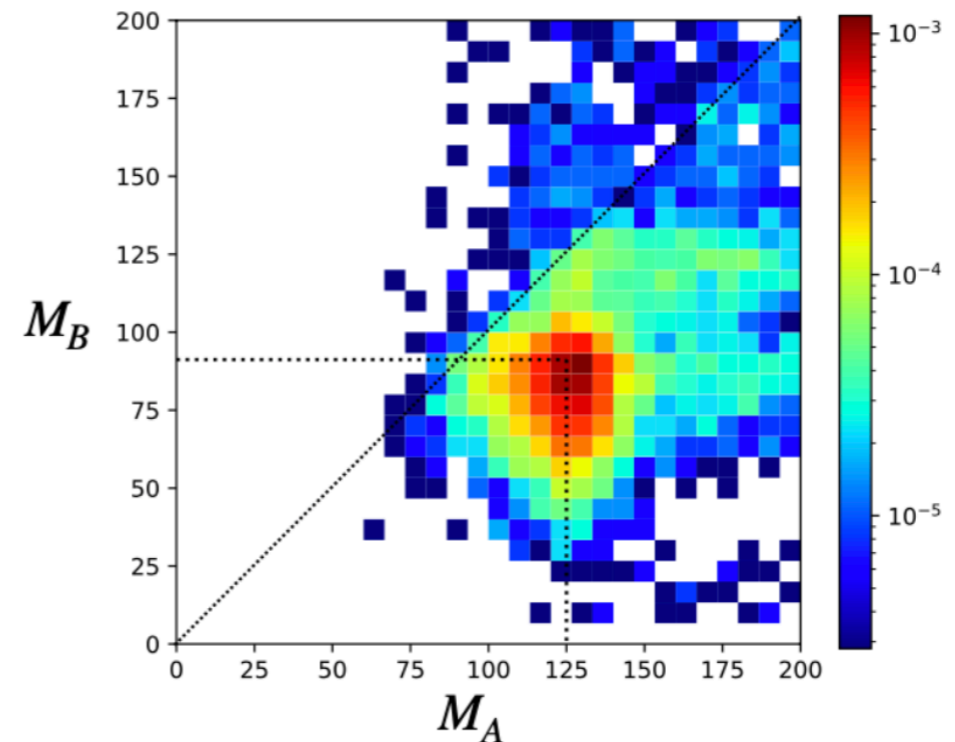
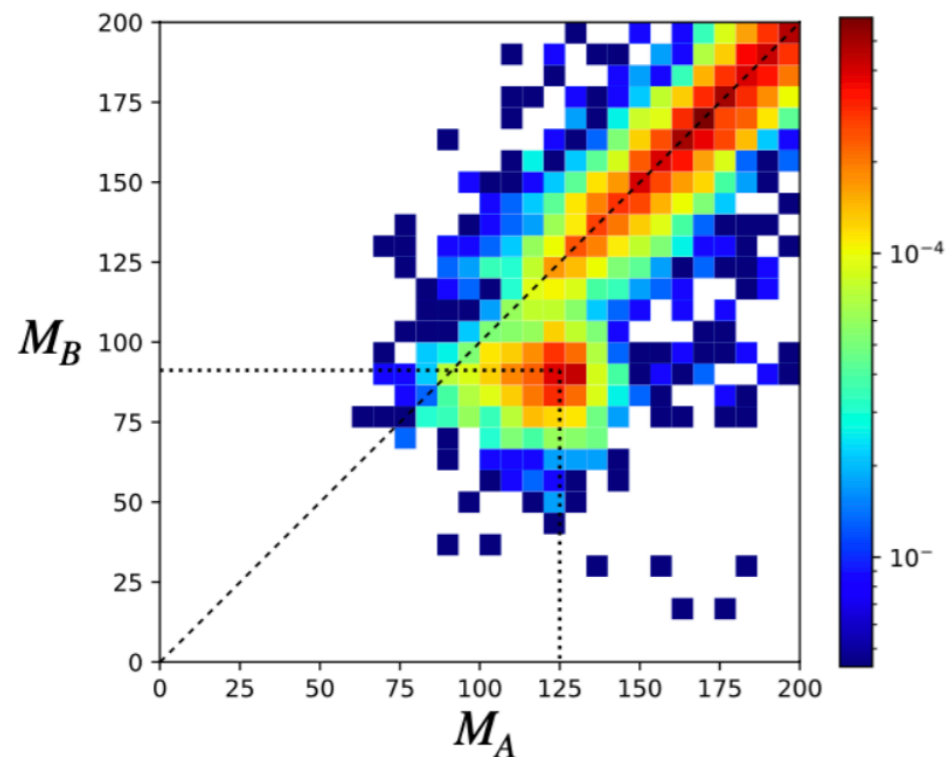
Role of the λ term

$$H = (P_1^2 - P_2^2)^2 \rightarrow H + \lambda (P_1^2 + P_2^2)$$

For different mother particle cases: $pp \rightarrow HZ$

$$H = (P_1^2 - P_2^2)^2$$

$$H \rightarrow H + \lambda (P_1^2 + P_2^2)$$



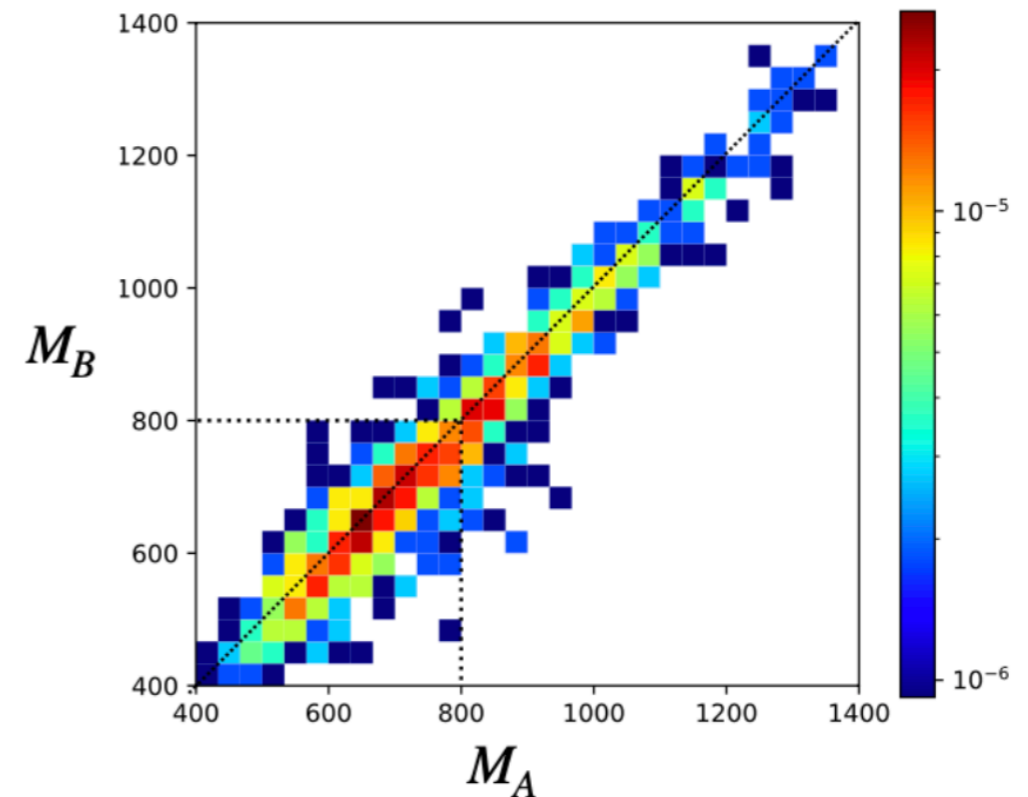
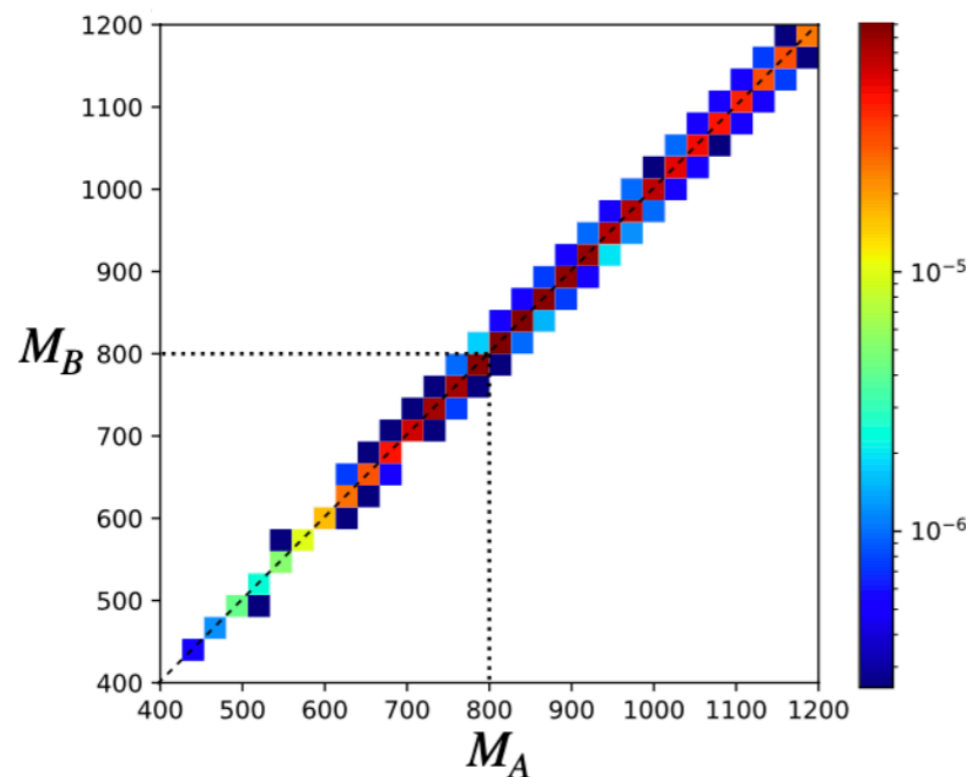
Role of the λ term

$$H = (P_1^2 - P_2^2)^2 \rightarrow H + \lambda (P_1^2 + P_2^2)$$

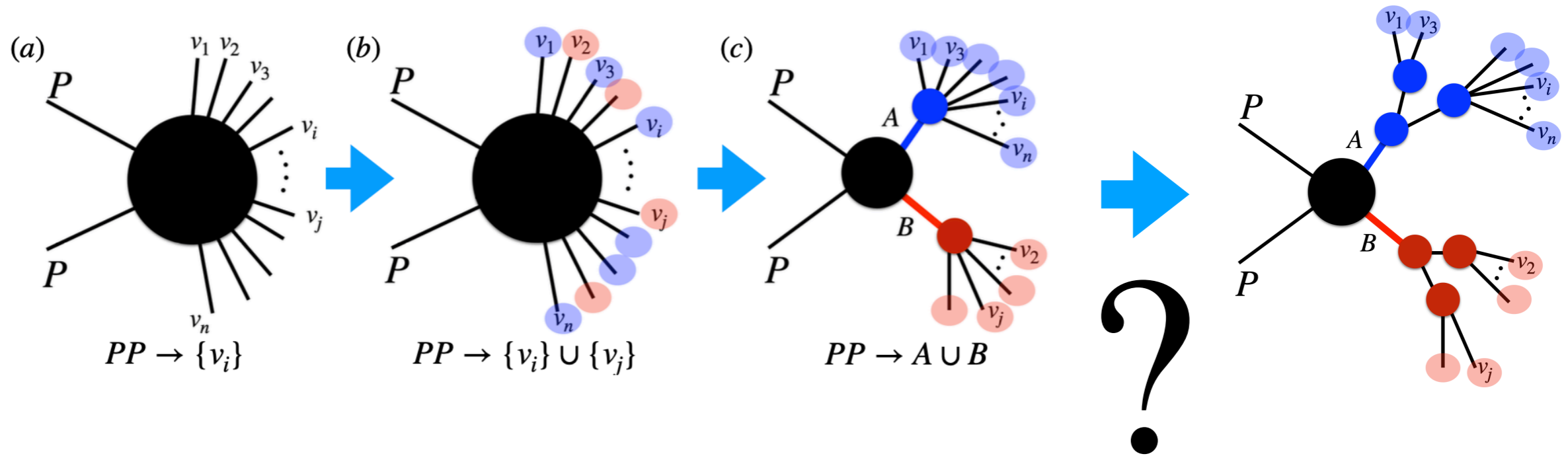
For smearing effects : $pp \rightarrow \tilde{o}\tilde{o} \rightarrow t\bar{t}t\bar{t}$

$$H = (P_1^2 - P_2^2)^2$$

$$H \rightarrow H + \lambda (P_1^2 + P_2^2)$$



Sequential algorithm



$$H_{\text{QUBO}}^{(A)} = \sum_{ij=1}^{\ell} J_{ij}^{\prime\alpha} s_i^{\alpha} s_j^{\alpha} + \sum_{i=1}^{\ell} h_i^{\prime\alpha} s_i^{\alpha},$$

$$H_{\text{QUBO}}^{(B)} = \sum_{ij=1}^m J_{ij}^{\prime\beta} s_i^{\beta} s_j^{\beta} + \sum_{i=1}^m h_i^{\prime\beta} s_i^{\beta},$$

- For 12 hard-jets production, it would be worthy if we can check whether this is four-tops events or not !

Results

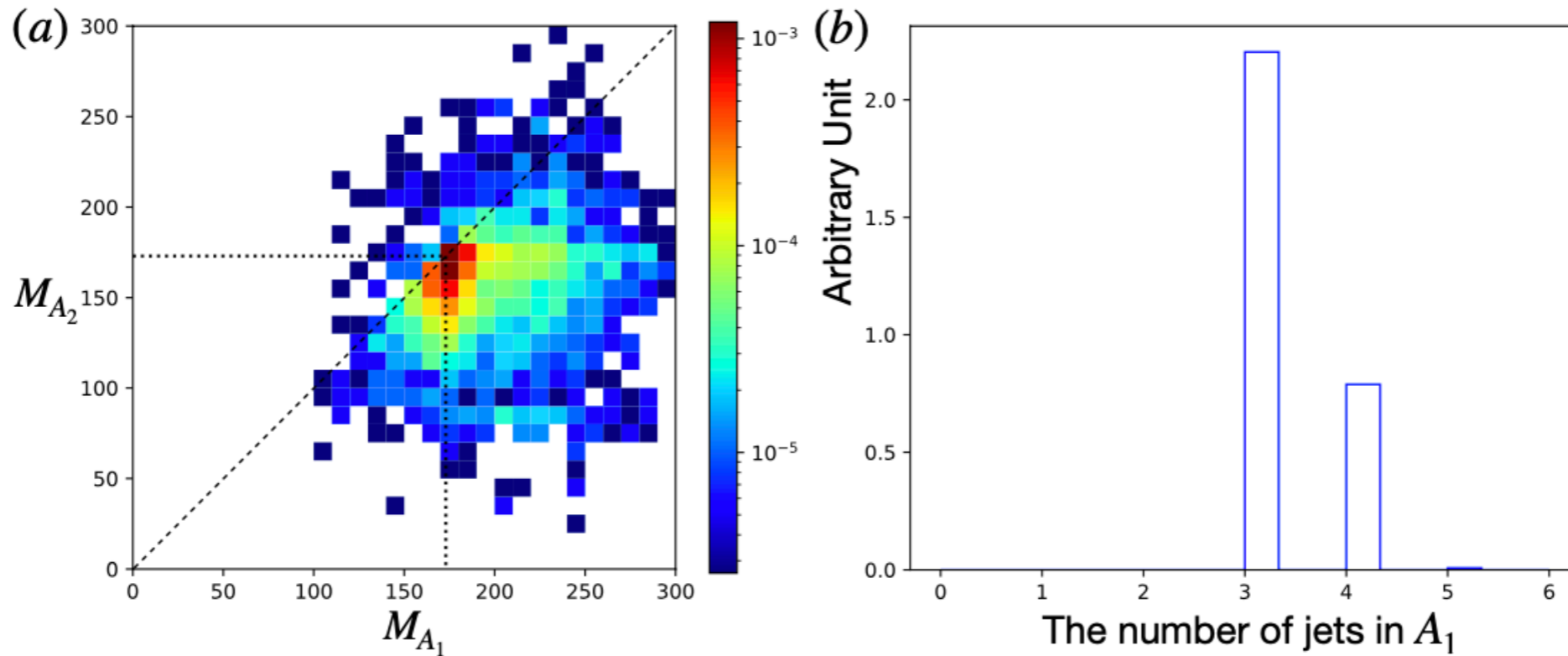


FIG. 4. Sequential application of H_{QUBO} to zoom in on the structure of A in Eq. (7c). After 12 jets are divided into two groups A and B , H_{QUBO} further investigates the structure of A . It identifies that $A \rightarrow t\bar{t}$ by measuring (a) masses and (b) the number of decaying particles of A_1, A_2 for $A \rightarrow A_1, A_2$.

Compare with Hemisphere method

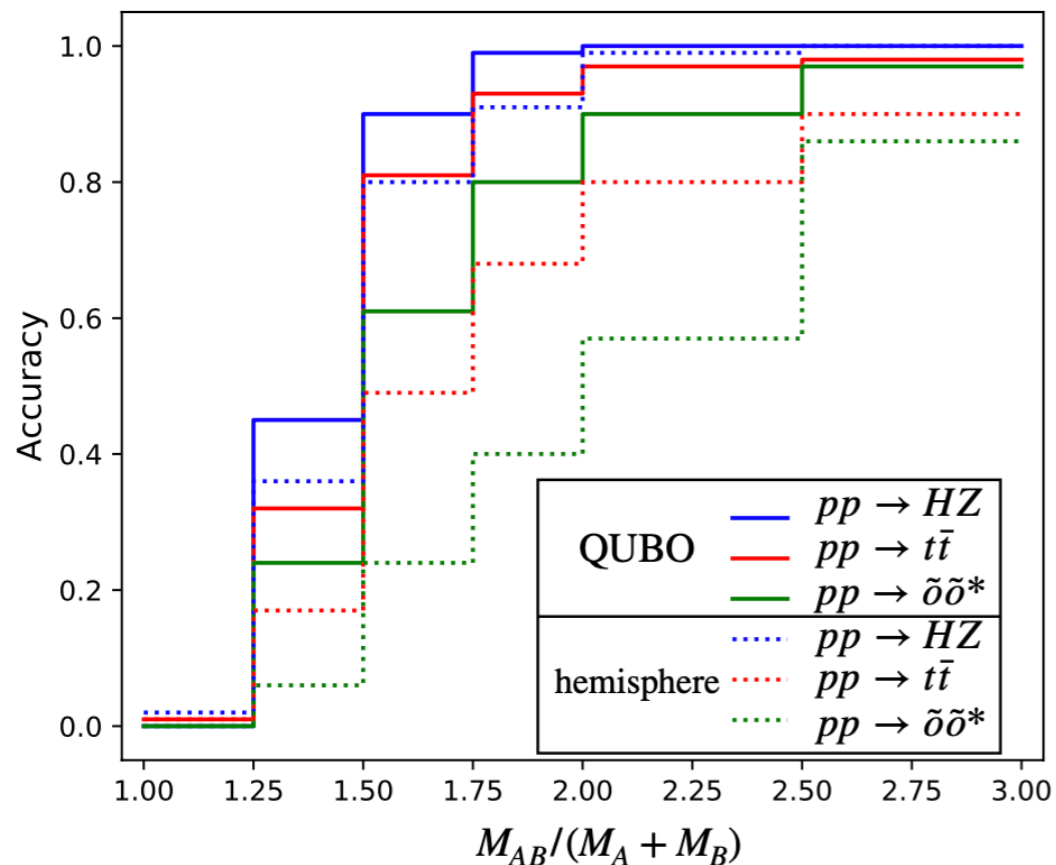


FIG. 6. Matching accuracy same as in Tab. II in terms of their boost factor $M_{AB}/(M_A + M_B)$. Here M_{AB} is an invariant mass of A and B .

Process		$pp \rightarrow t\bar{t}$ Eq. (7a)	$pp \rightarrow HZ$ Eq. (7b)	$pp \rightarrow \tilde{o}\tilde{o}^*$ Eq. (7c)
Algorithm	QUBO	47.3%	89.5%	17.4%
	Hemisphere	33.6%	86.2%	7.2%

TABLE II. Matching accuracy for the reconstructed momenta of particles A and B using a clustering algorithm to an actual momenta of A and B (parton level analysis but with same basic cuts as in previous Monte Carlo samples).

Performance of an algorithm based on Hemisphere method becomes weak when particles are **not boosted enough to develop structure**

Conclusion

With increasing energy and luminosity available at the Large Hadron collider (LHC), we get a chance to take a pure bottom-up approach solely based on data. This will extend the scope of our understanding about Nature without relying on theoretical prejudices. The required computing resource, however, will increase exponentially with data size and complexities of events if one uses algorithms based on a classical computer. In this letter we propose a simple and well motivated method with a quantum annealer to identify an event-topology, a diagram to describe the history of particles produced at the LHC. We show that a computing complexity can be reduced significantly to the order of polynomials which enables us to decode the “Big” data in a very clear and efficient way. Our method achieves significant improvements in finding a true event-topology, more than by a factor of two compared to a conventional method.

Questions for future direction

1. How about $2 \rightarrow n$ ($n > 2$)
2. Off-shell case?
3. Missing daughter particle
4. Contamination from the ISR
5. Systematic error from Quantum Annealer for large qubit
6. Another quantum computing method
7. And more