



Update about the Accelerator and IR

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Electron-Ion Collider

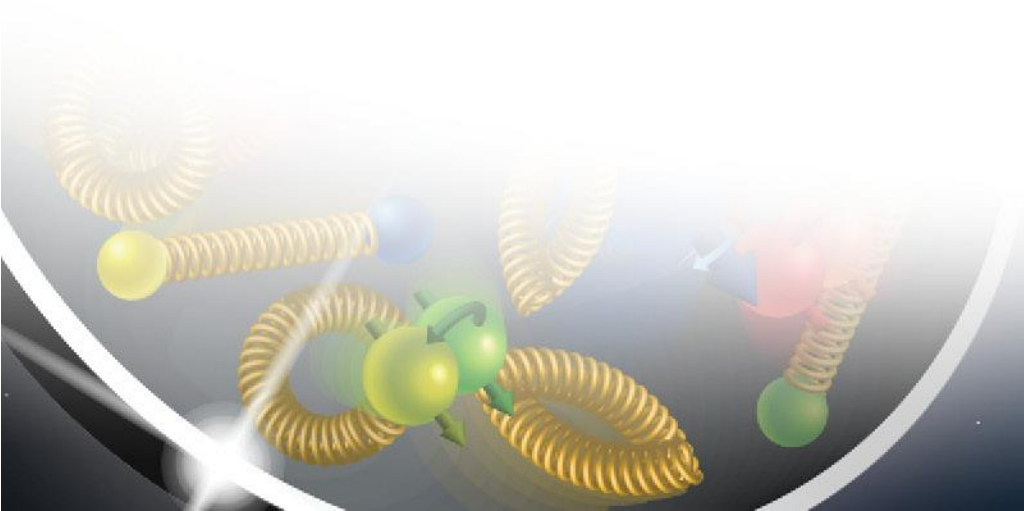
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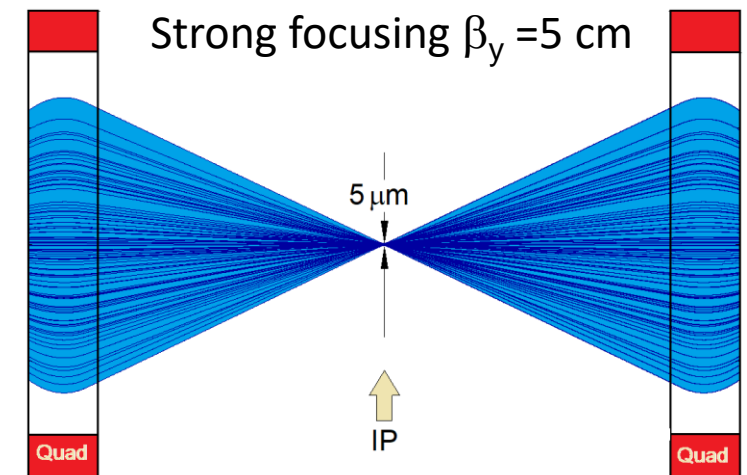
EIC Interaction Region Requirements

- High luminosity
- High p_t acceptance
- Detection of neutral particles (neutrons, photons,...)
- Longitudinal polarization
- Safely pass synchrotron radiation through the detector



Luminosity and Focusing

- Luminosity $\sim 1/(\text{spot size})$
- A **smaller spot size** at the IP means **more luminosity**
- At the IP, **(beam size)X(beam divergence)= const.** in each plane (**emittance ϵ**)
- Beam-beam force scales as $1/\epsilon$ – beam dynamics prefers **large emittance**
- For a given beam (= fixed emittance), a **smaller IP beam size means larger divergence**
- **Two configurations: High luminosity vs. high p_t acceptance**
- A larger beam divergence leads to a larger beam size at the nearest focusing magnets – **(size at magnet)=(divergence)X(distance)**
- **Magnets need to have larger aperture** while gradient (= focusing strength) remains the same – peak field at magnet poles is **technically limited**



Focusing elements for both beams need to be as close as possible to the IP

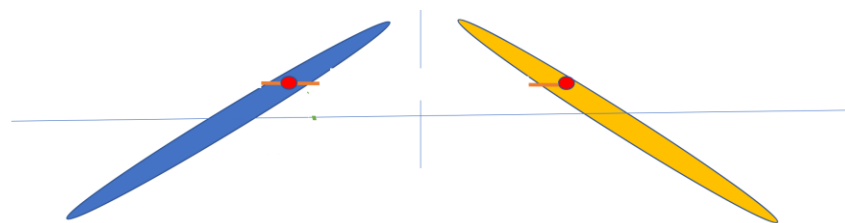
Crossing Angle Collisions

- **Beam energies** of electrons and hadrons are **vastly different** in EIC
- Focusing elements for electrons would have only little effect on hadrons, while hadron magnets would overfocus electrons
- **Beams need to be separated** into their respective focusing systems as close **as possible to the IP**
- A **separator dipole** would have to deflect the (“weaker”) electrons and would therefore generate a **wide synchrotron radiation fan** that would need to pass through the detector – requires **large beam pipe diameter** (HERA-II)
- Best solution: **Crossing angle collisions!**

Crossing Angle and Luminosity



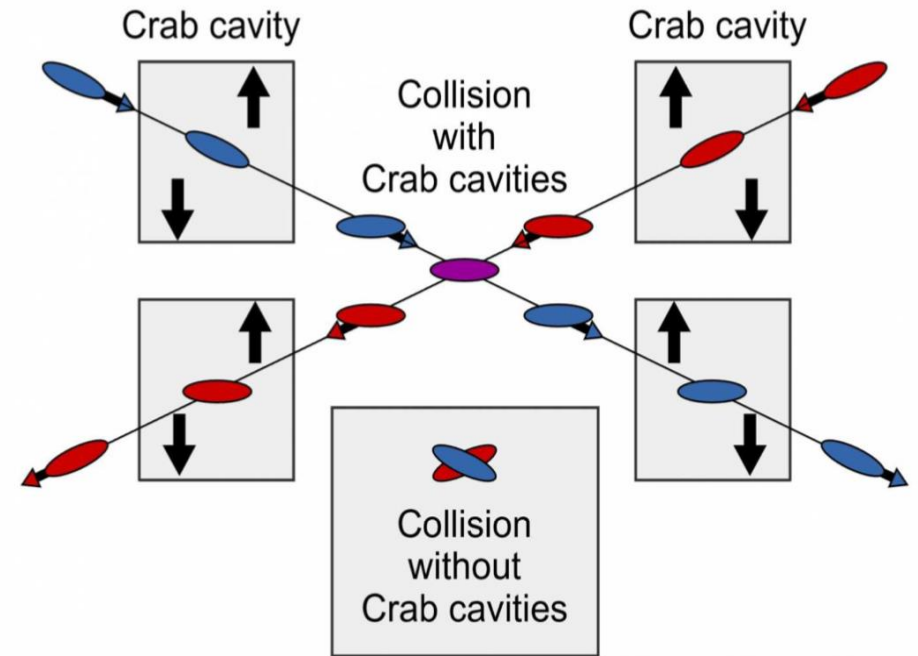
- In **head-on** collisions, **every beam particle** in one beam can potentially interact with **every particle in the other beam**



- **Long** ($\sim \pm 6$ cm), **skinny** (100 μm) **bunches** colliding at an angle have **very little overlap**
- With **25 mrad crossing angle**, each particle can only interact with a **± 4 mm thick slice** of the ± 6 cm long oncoming bunch

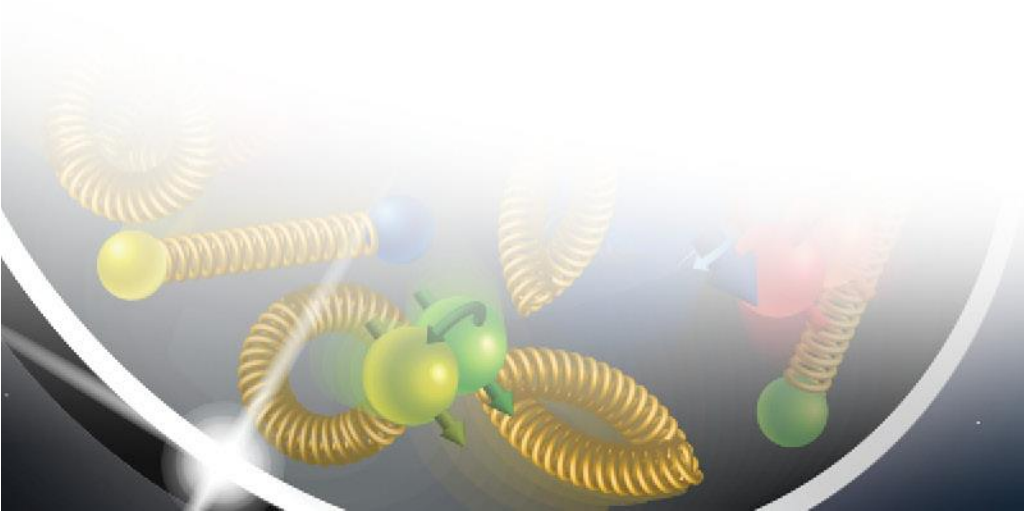
Crab Crossing to the Rescue

- Head-on collision geometry is **restored** by rotating the bunches before colliding (“**crab crossing**”)
- Bunch rotation (“crabbing”) is accomplished by **transversely deflecting RF resonators** (“**crab cavities**”)
- Actual collision point moves laterally during bunch interaction

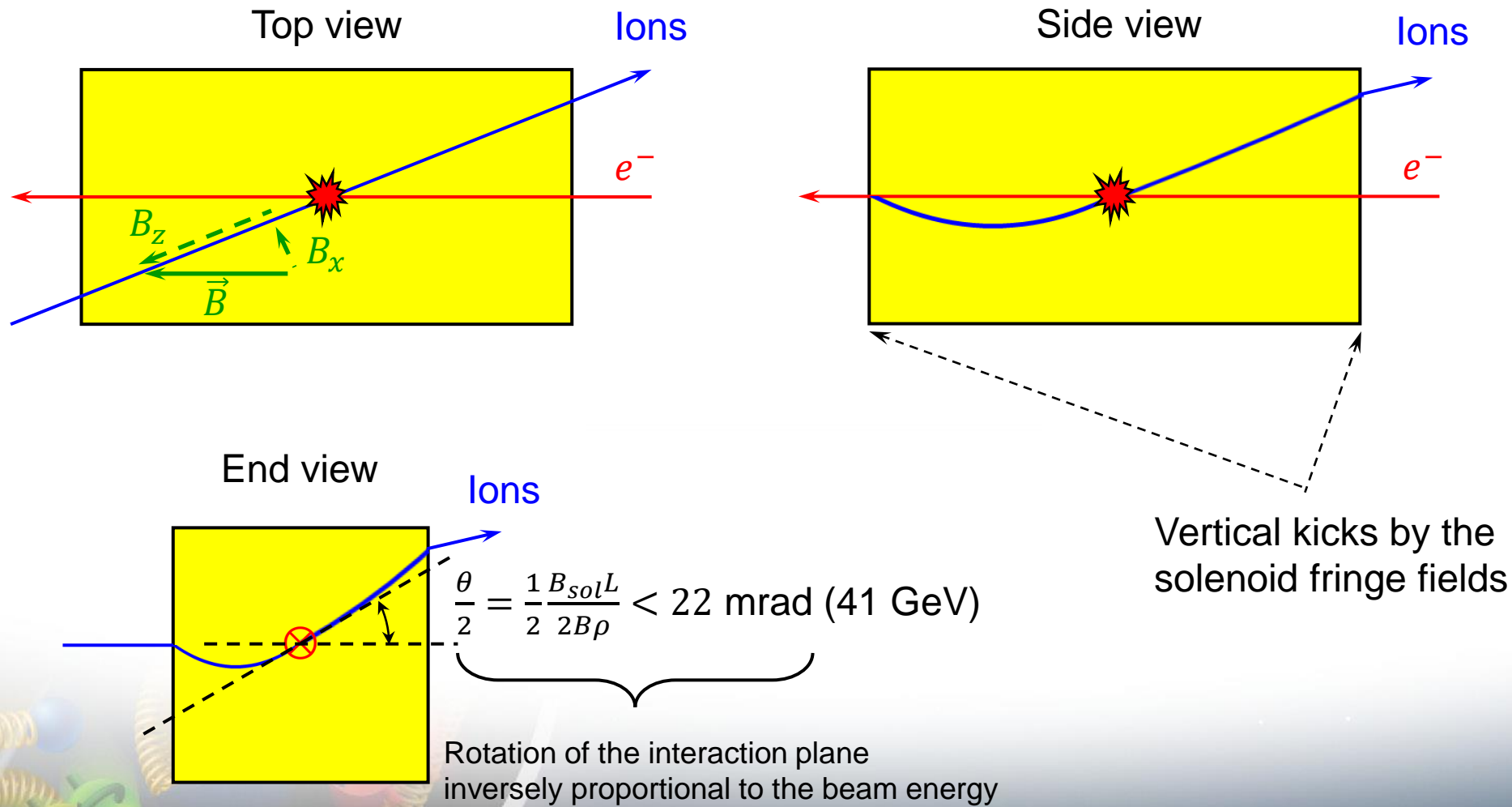


Detector Solenoid Effects

- Coherent orbit distortion
- Transverse coupling
- Rotation of the crabbing plane
- Polarization tilt

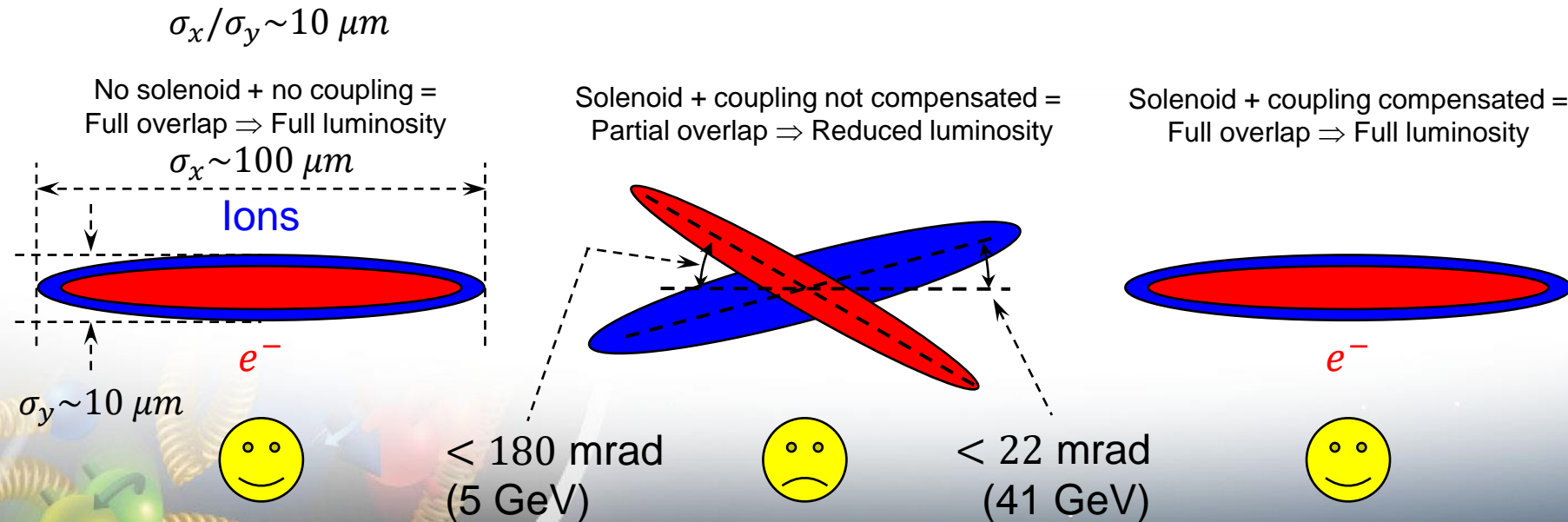


Coherent Distortion of Ion Orbit



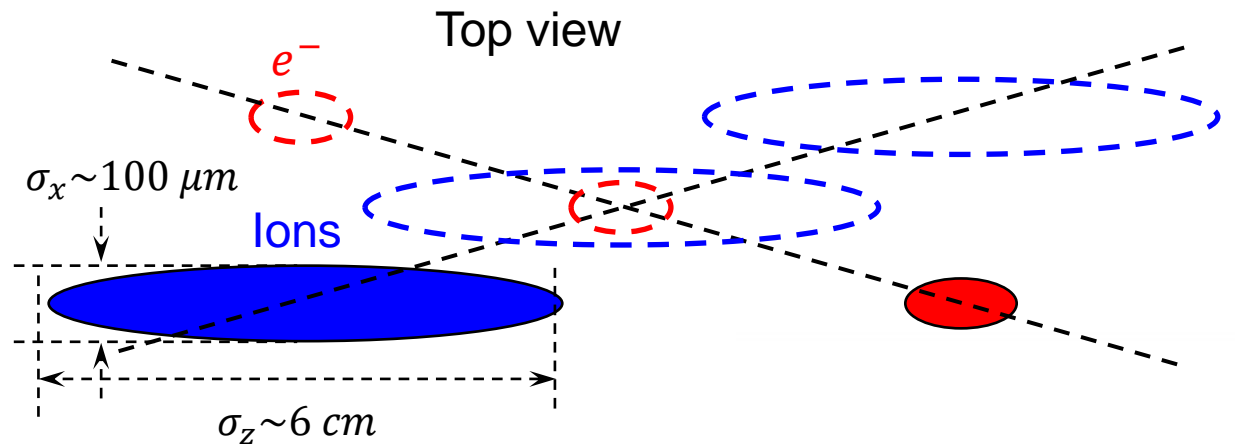
Transverse Coupling

- Coupling is in general a global effect and involves consideration of the entire ring or the entire coupled section
- Potential for negative dynamic effects
- Can lead to redistribution of the horizontal and vertical emittances \Rightarrow smaller beam “flatness” \Rightarrow luminosity reduction
- If not locally compensated at the IP can lead to change in the transverse beam shape \Rightarrow potential for luminosity reduction (beams enter solenoid uncoupled):



Rotation of Crabbing Plane

- Another aspect of coupling
- Potential for negative dynamic effects
- Impact on luminosity



No solenoid + no coupling =
Full overlap \Rightarrow Full luminosity

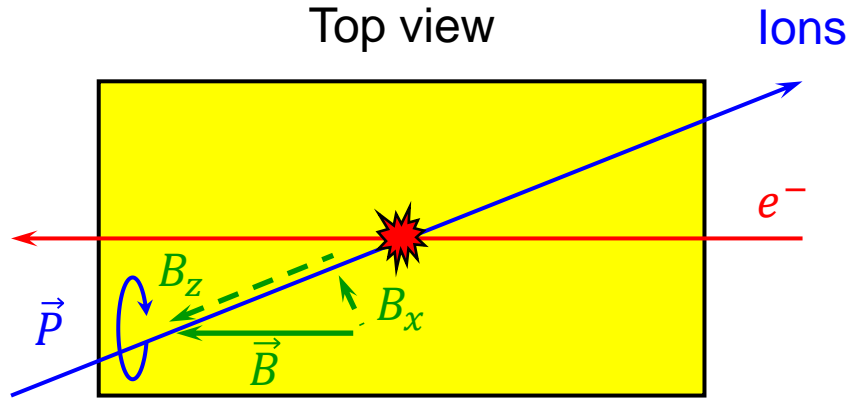


Solenoid + rotation not compensated =
Partial overlap = Reduced luminosity



Polarization Tilt

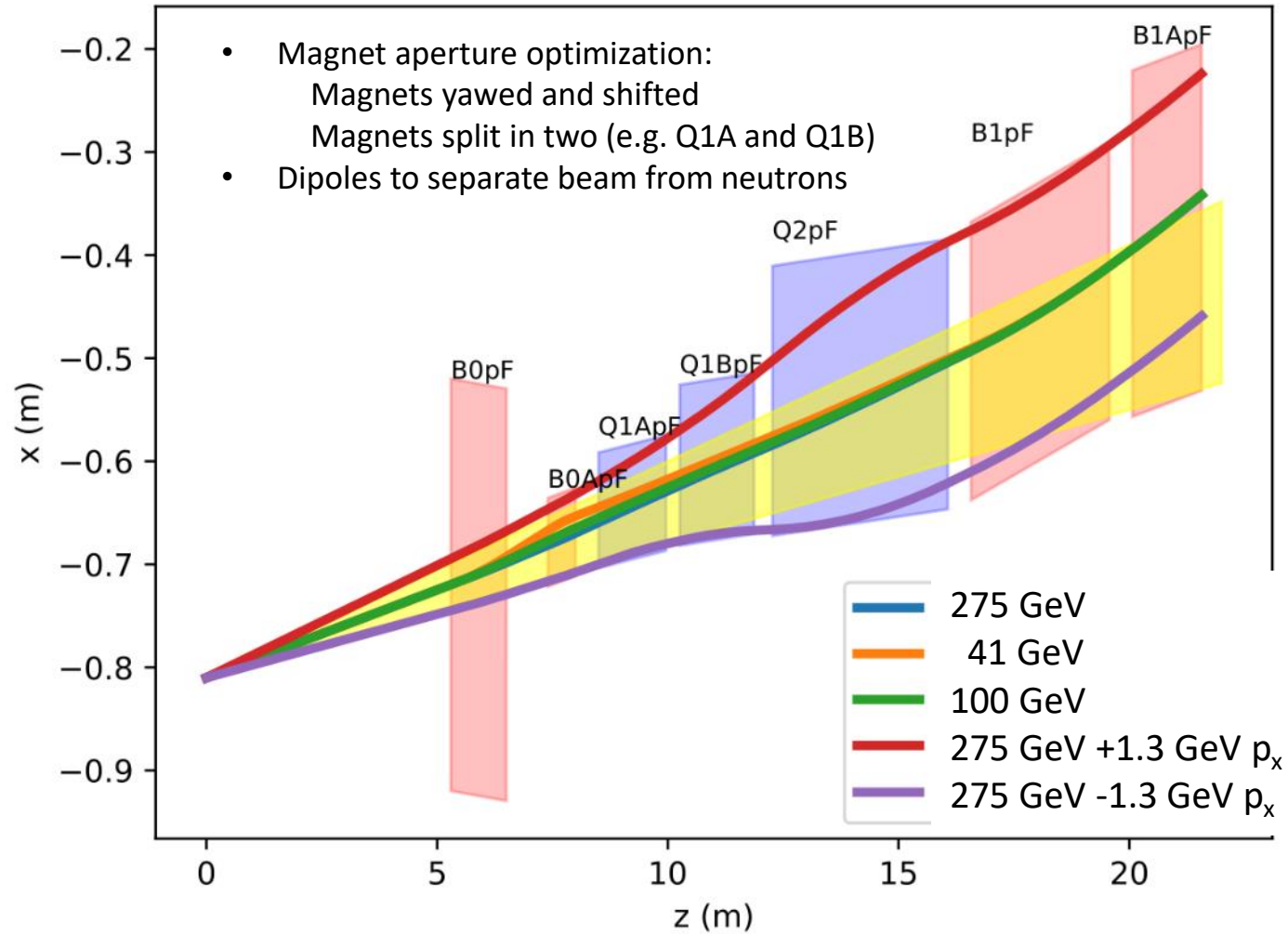
- The baseline scheme has no strong correctors on the rear side, so the polarization tilt is due to the rear side of the solenoid



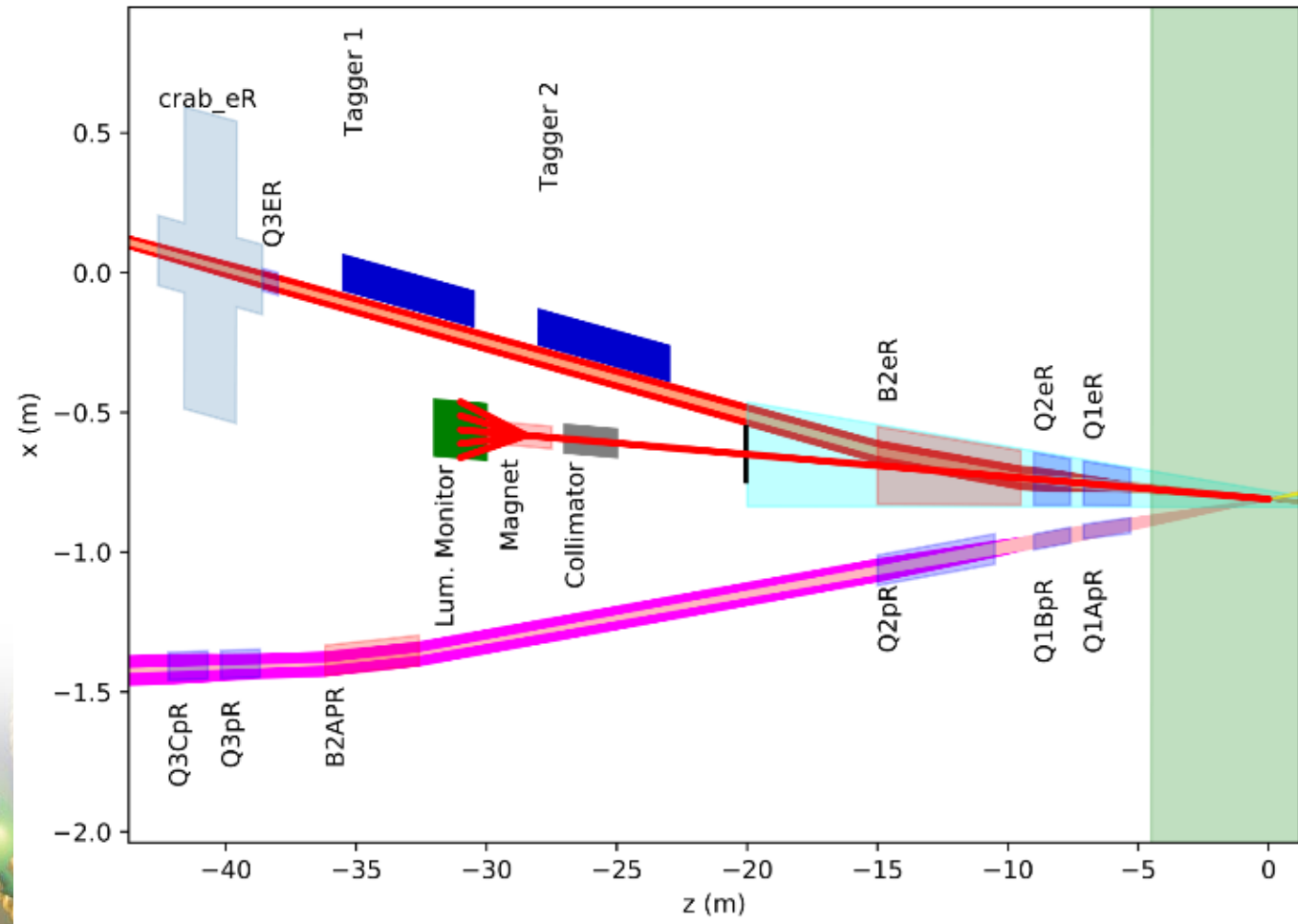
Energy [GeV]	Long. n_x	Rad. n_y	Vert. n_z	φ_s [2π]
41	0.94	-0.34	0.02	0.01
100	0.76	-0.65	0.01	0.01
200	0.50	-0.87	0.01	0.008
275	0.39	-0.92	0.007	0.007

- Polarization orientation at the IP in this scenario is compensated by the rear ion spin rotator
- The net spin effect of the entire IR involves account of the contribution of the forward side
- The forward spin rotator compensation the remaining net effect

Hadron Forward – Large Apertures for low p_t Acceptance

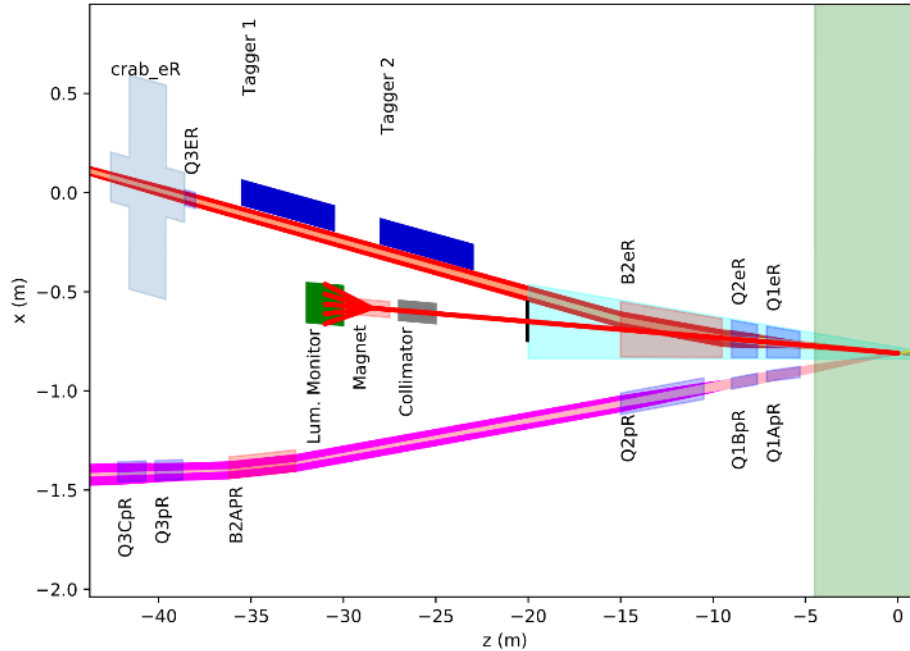


Electron Rear: Large Apertures to Pass Synchrotron Radiation Fan

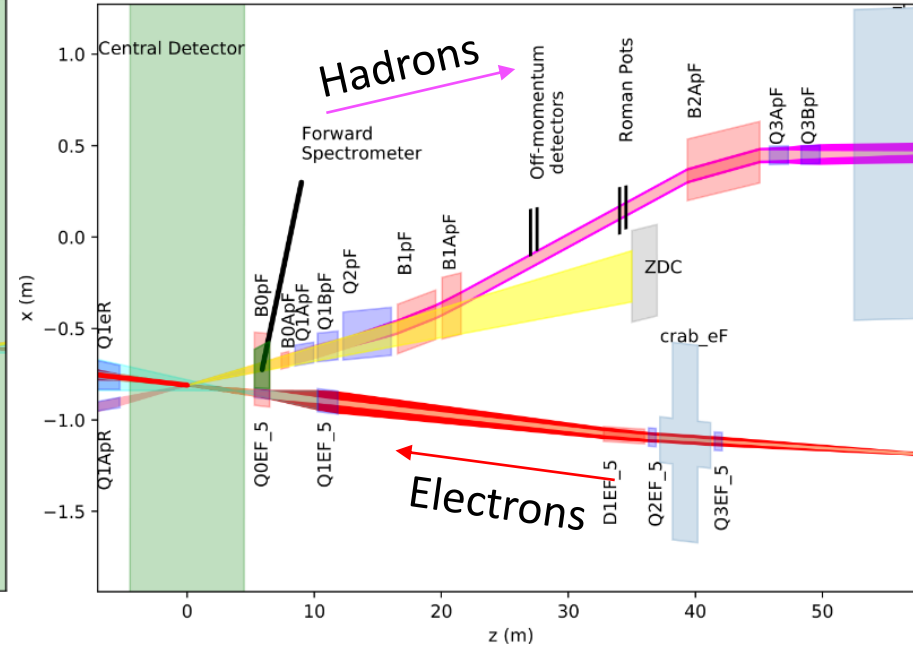


IR Magnets

Rear: RHIC sec 6



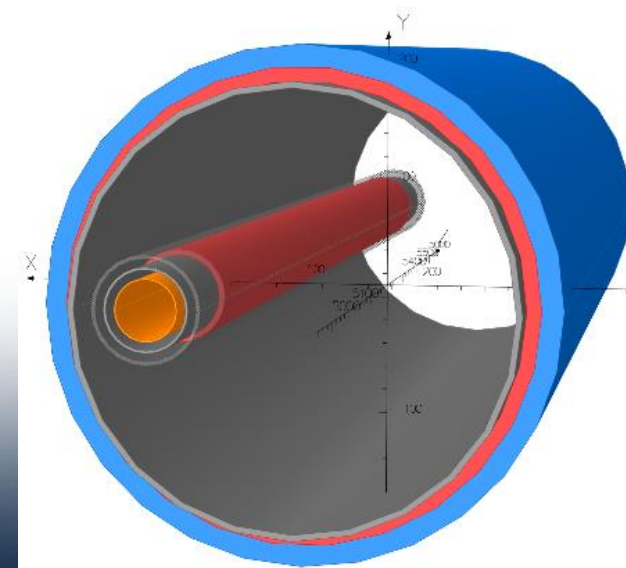
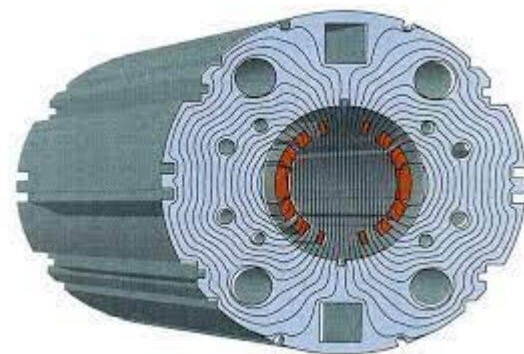
Forward: RHIC sec 5



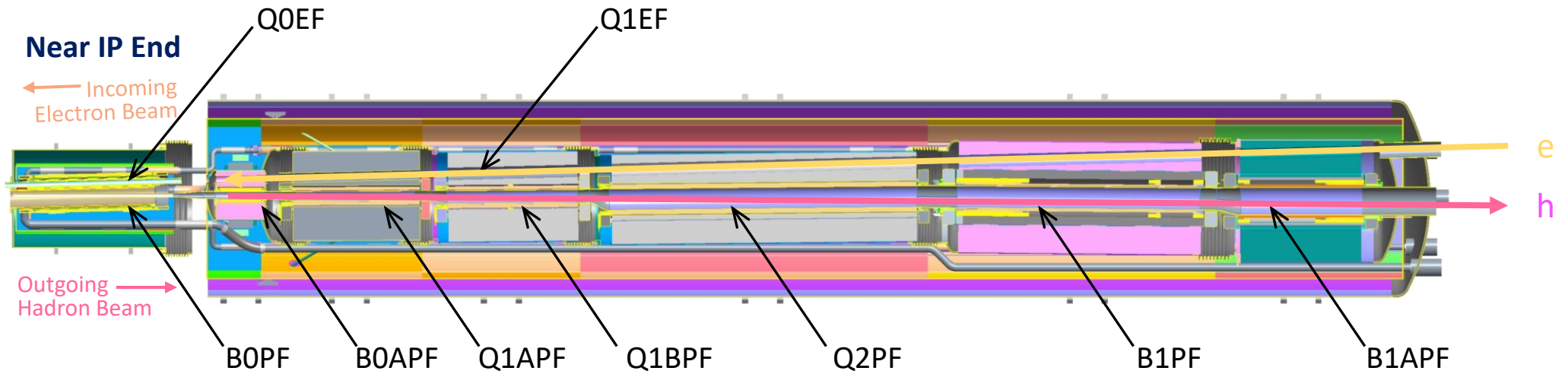
- 16 SC magnets to be built (4.5 K and 1.9 K)
- 10 direct wind magnets:
 - Q1eR, Q2eR, B2eR, Q1ApR, Q1BpR, Q2pR, Q0eF, B0pF, B0ApF, Q1eF
- 6 collared magnets (incl. B2pF)
 - Q1ApF, Q1BpF, Q2pF, B1pF, B1ApF, B2pF

Superconducting IR magnets

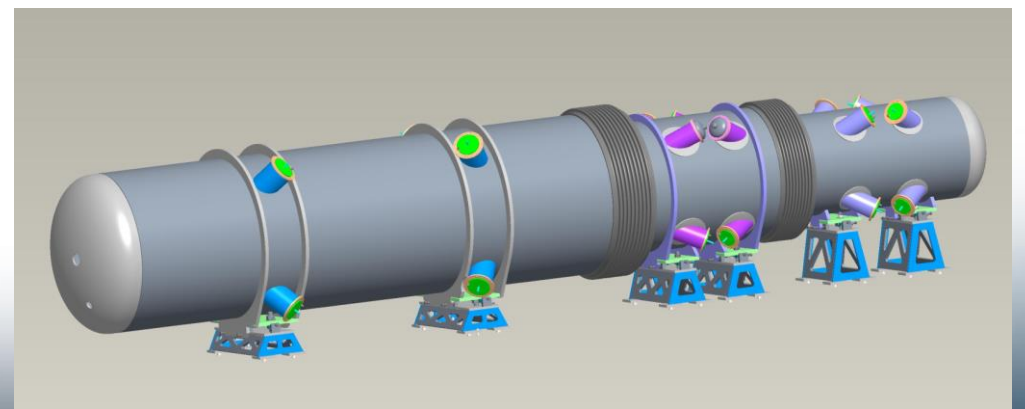
- 10 direct-wind magnets
- 6 collared magnets
- B0 – an electron quadrupole inside a hadron spectrometer dipole



IR Magnet cryostats



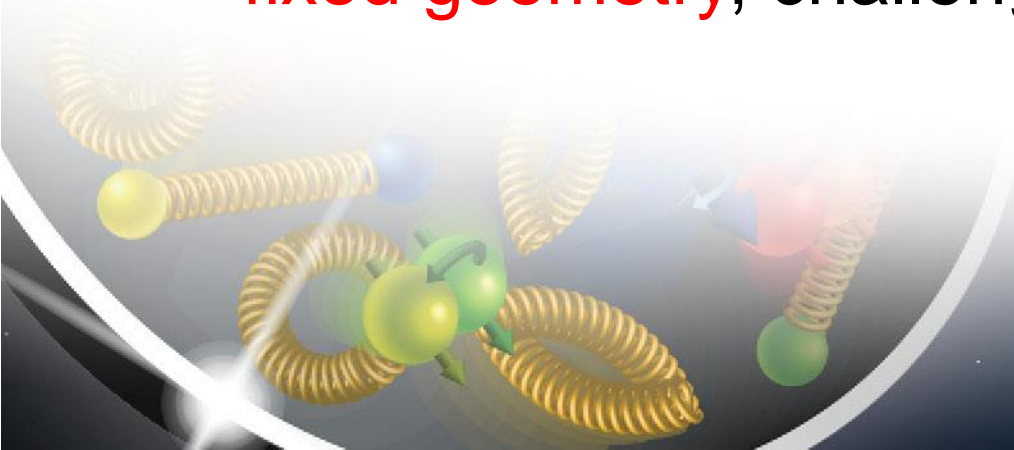
Electron and hadron magnets densely packed, side-by-side in a common cryostat on either side of the detector



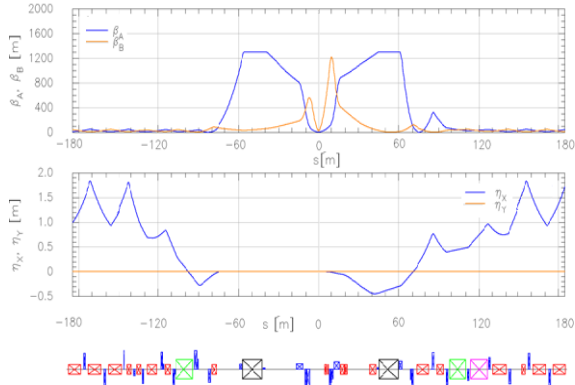
Electron-Ion Collider

Spin Rotators

- Both electrons and protons will have **longitudinal polarization** at the IP
- Hadron spin rotators will be taken from present RHIC (helical dipoles)
- Electron spin rotators are based on solenoid magnets with subsequent dipole – **large ($> \pm 100$ m) chunk of beamline with fixed geometry**, challenging to fit into existing tunnel

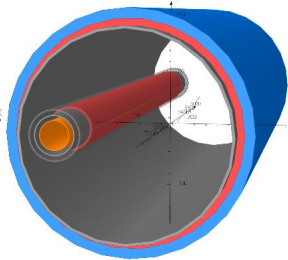
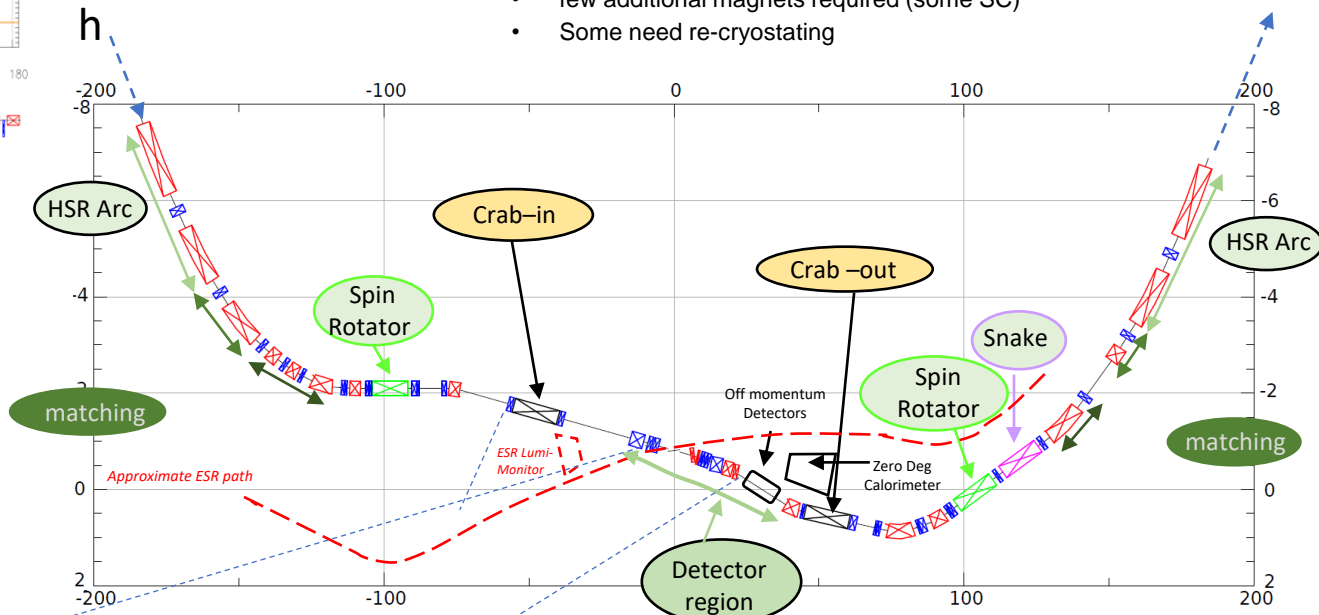


HSR layout in IR6



- Forward and rear hadron lattice matched into RHIC

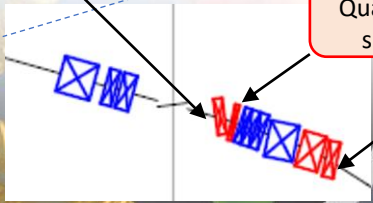
- Snake at correct angle
- Beta = 1300m at crab cavities
 - Hor. phase advance 90°
- Matching Magnets
 - Mostly repurposed RHIC magnets
 - few additional magnets required (some SC)
 - Some need re-cryostating



B0pF spectrometer

Nov 2021 layout

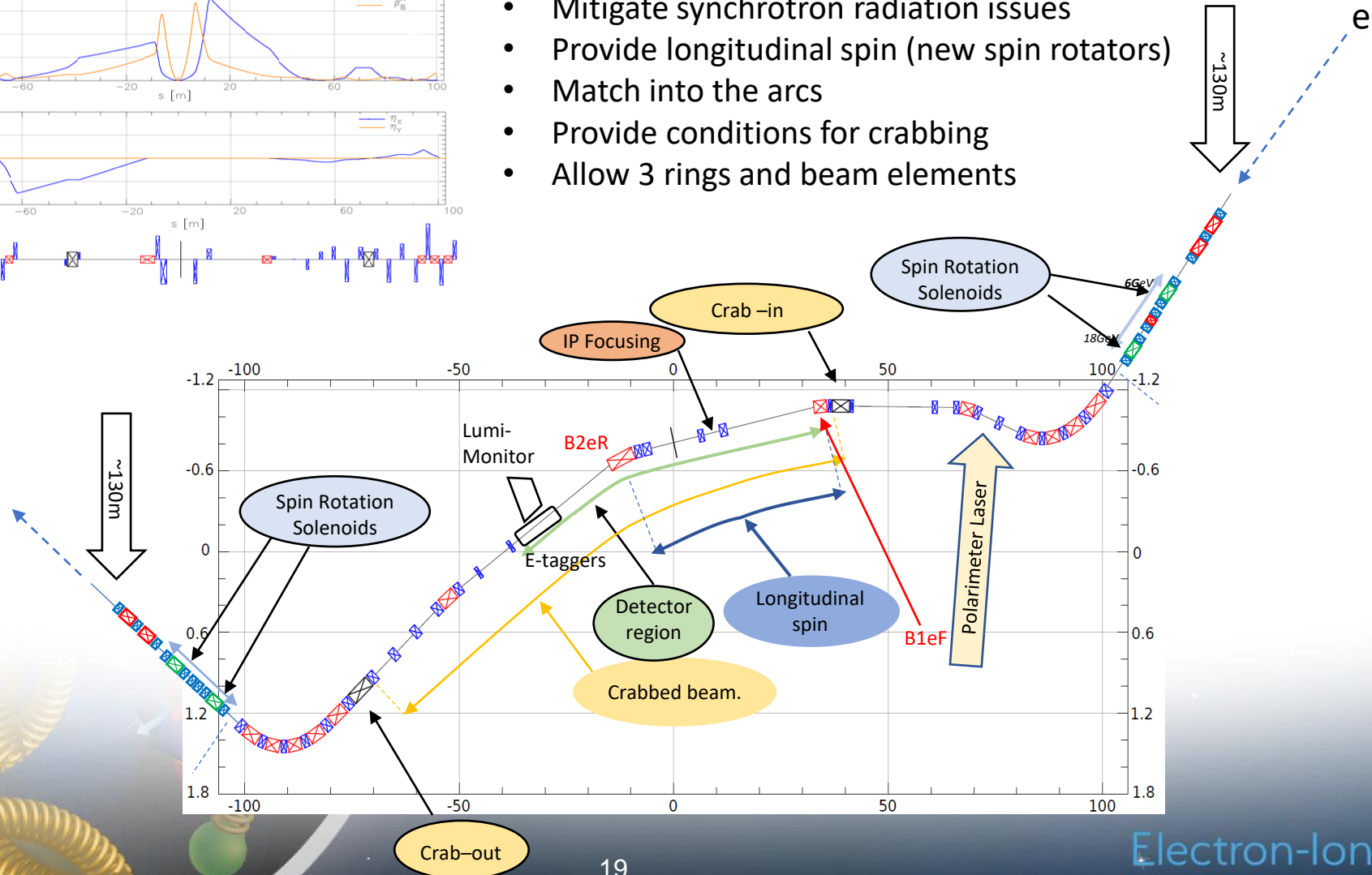
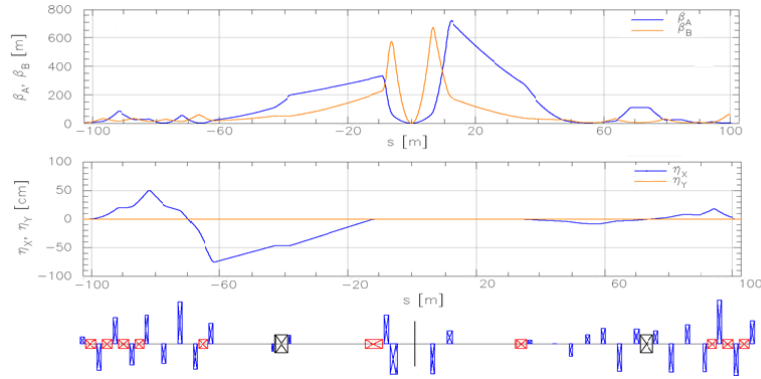
3 Dipoles and 3 Quadrupoles in one shared cryostat



ESR layout in IR6

Design to:

- Provide room for detector components
- Mitigate synchrotron radiation issues
- Provide longitudinal spin (new spin rotators)
- Match into the arcs
- Provide conditions for crabbing
- Allow 3 rings and beam elements

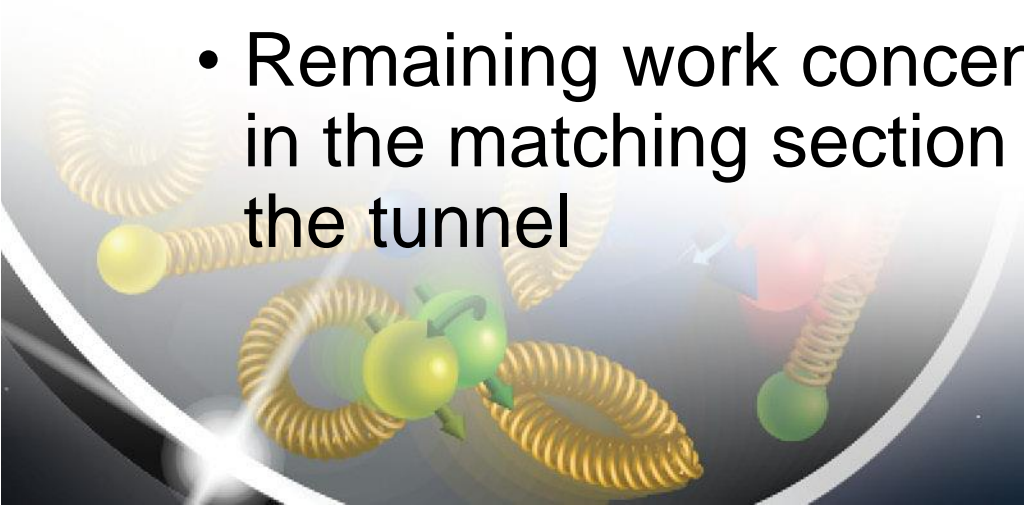


Luminosity Sharing with two IRs

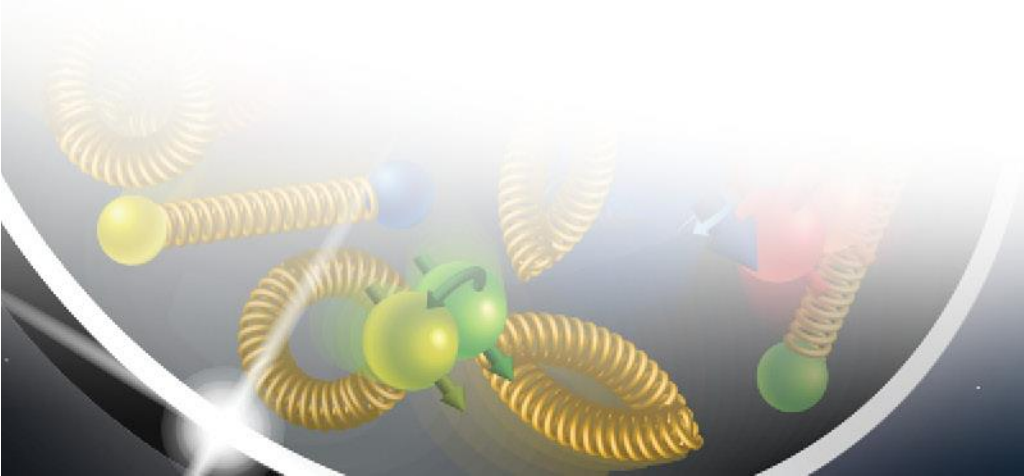
- Both electrons and hadrons are at the **beam-beam limit** with one collision point – they would not “survive” a second IR
- To enable **two collision points**, both electron and hadron bunch **intensity would have to be reduced by a factor two** – resulting luminosity at each IR would be **factor 4 smaller**
- Instead, we modify the fill pattern such that half the bunches collide in IR6, while the other half collides in IR8
- As a result, total luminosity is preserved, and **each detector gets half of the total** – a maximum $5e33$ each instead of $1e34$ with a single IR

Summary

- EIC interaction region is highly optimized, but any IR design is always a compromise between luminosity, acceptance, cost, risk, ...
- Design is driven by physics and detector needs, in close collaboration with experimenters
- Inner IR is practically frozen
- Remaining work concentrates on small geometric adjustments in the matching section towards the arcs, to fit everything into the tunnel

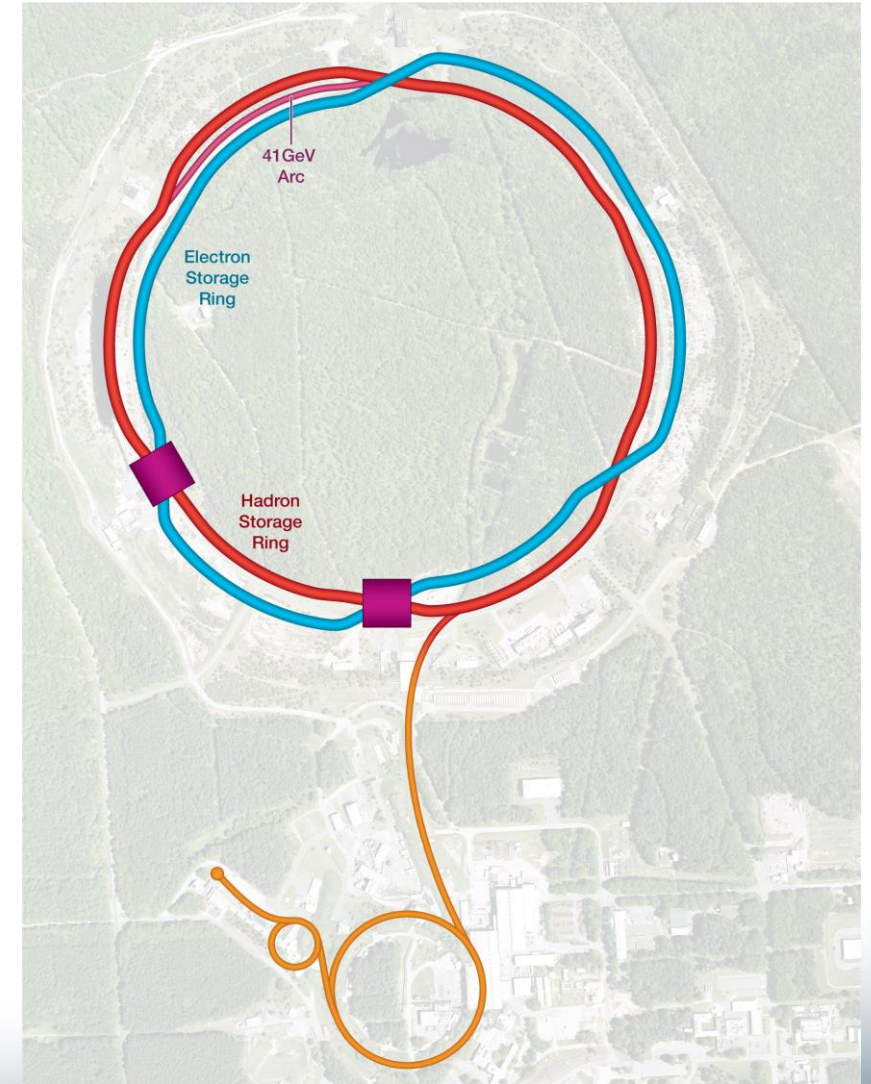


Supplemental Slides



Collision Synchronization

- HSR needs to operate over a **wide energy range**
- Changing the beam energy in the HSR causes a **significant velocity change**
- To **keep the two beams in collision**, they have to be synchronized so bunches arrive at the detector(s) at the same time
- Synchronization accomplished by **path length change**
- Between **100 and 275 GeV (protons)**, this can be done by a **small radial shift** – there is enough room in the beampipe
- For lower energies, use an inner instead of an outer arc as a **shortcut**. 90 cm path length difference corresponds to **41 GeV** proton beam energy



Electron-Ion Collider

Emittance Control in the ESR

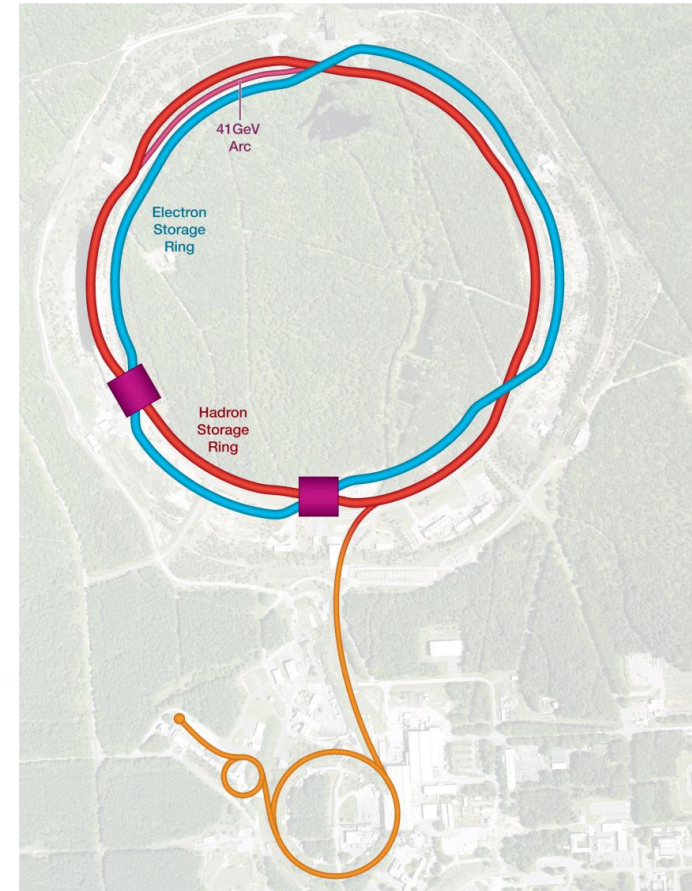
- EIC needs **24 nm emittance from 5 to 18 GeV for optimum luminosity**, but equilibrium emittance in an electron storage ring depends on beam energy:

$$\epsilon_x = C_q \frac{\gamma^2}{J_x} \frac{I_5}{I_2}, \quad \text{with} \quad C_q = \frac{55}{32\sqrt{3}} \frac{\hbar c}{mc^2}$$

- Betatron phase advance μ per FODO cell is the “knob” to adjust the emittance
- **60 degrees at 10 GeV and 90 degrees at 18 GeV both yield ~24 nm**
- **“super-bends” for emittance generation below 10 GeV**

Beam Energies

- γ range for hadrons:
 - $\gamma = 43.7$ through “41 GeV arc”
 - $107 < \gamma < 293$ with radial shift
- Maximum hadron energy:
 - $E [\text{eV}] < 916 * c [\text{m/sec}] * Z/A$
- Electron energies:
 - 5 to 10 GeV, with 60 degree lattice and super-bends
 - 18 GeV, with 90 degree lattice
 - Energies between 10 and 18 GeV are feasible, but at somewhat reduced luminosity due to non-optimum emittance, scaling as γ^2



High Average Electron Polarization

- **Frequent injection** of bunches with high initial polarization of 85%
- Initial **polarization decays** towards P_∞
- At 18 GeV, every **bunch is replaced** (on average) after 2.2 min with RCS cycling rate of 2Hz

B P



Refilled every 1.2 minutes

B P

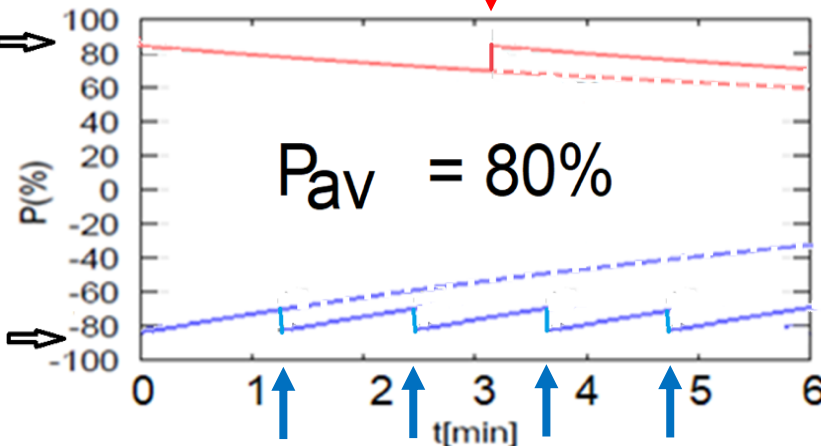


Refilled every 3.2 minutes

$P(0) = 85\%$



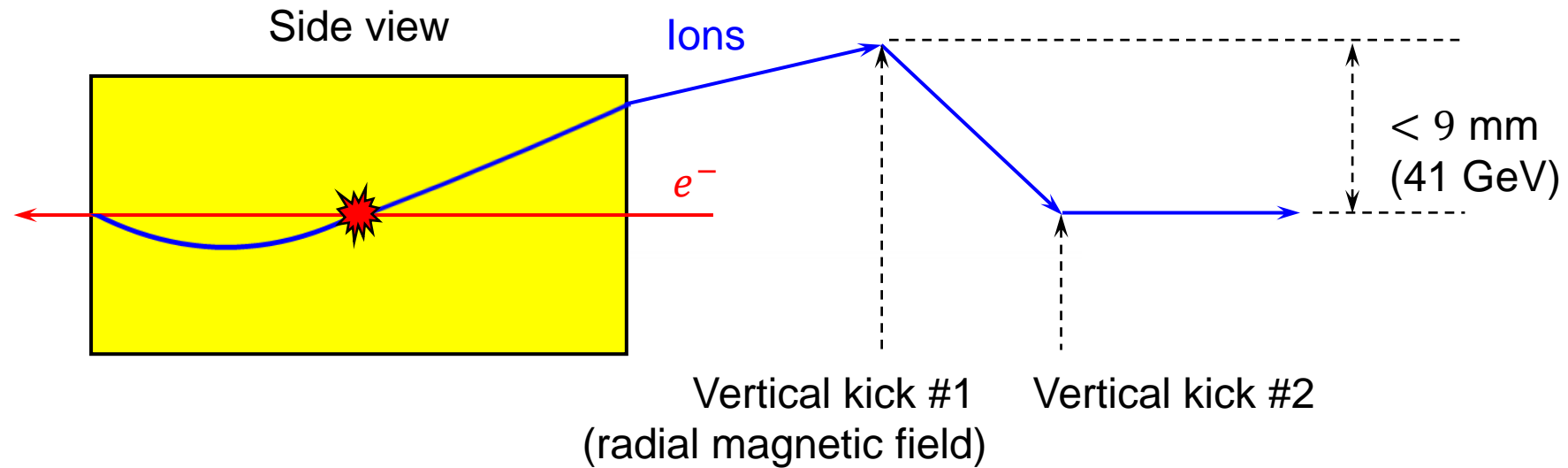
$P(0) = -85\%$



Re-injections

Correction of Ion Orbit

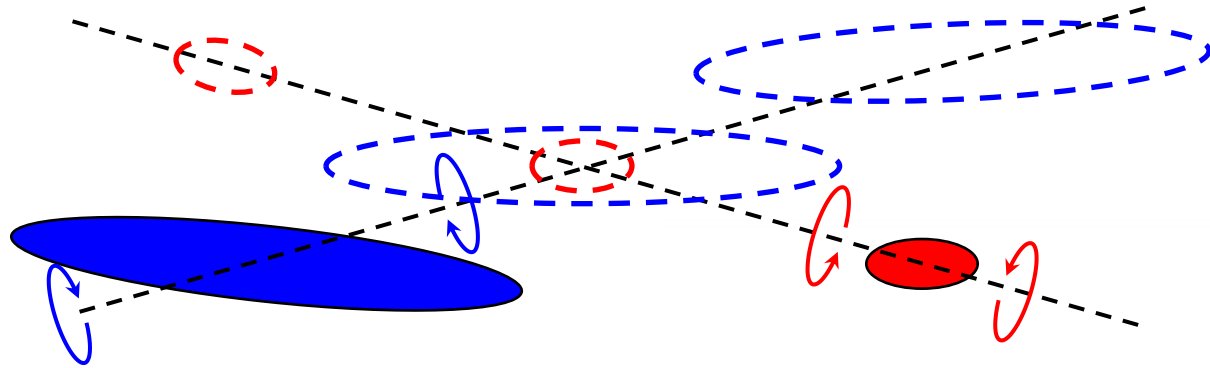
- The closer the kicks to the IP, the smaller the orbit excursion
- Orbit excursion inversely proportional to the beam momentum
- Concern for field non-linearity at large offsets



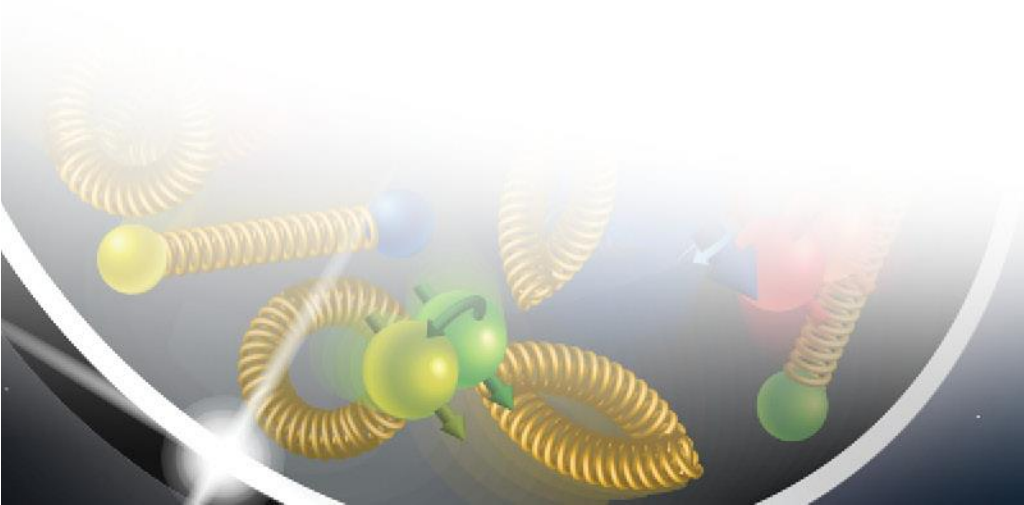
Rotation of Crabbing Plane Solution

- Pre-rotated the crabbing plane using vertical crabbing kicks: ~140 kV for ions and ~430 kV for electrons

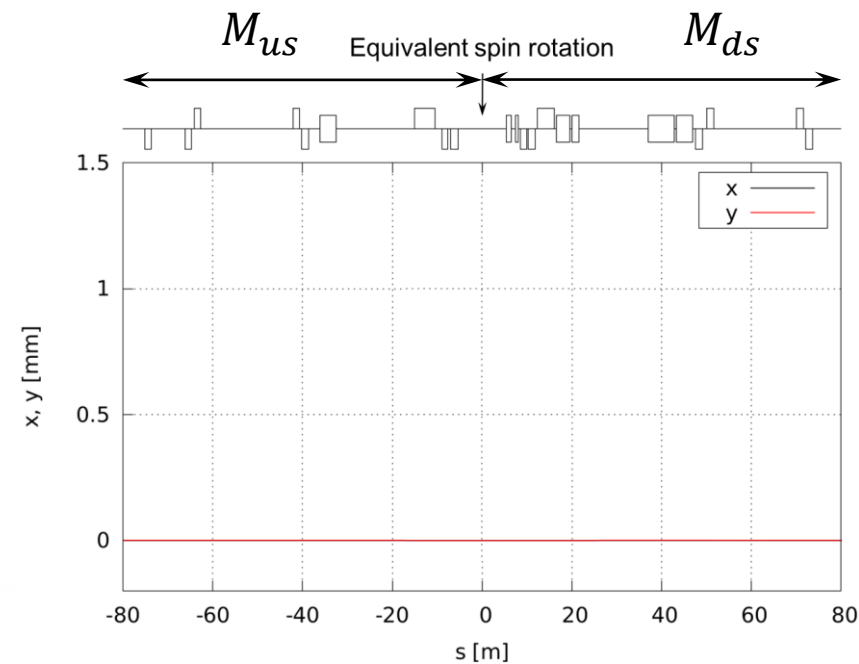
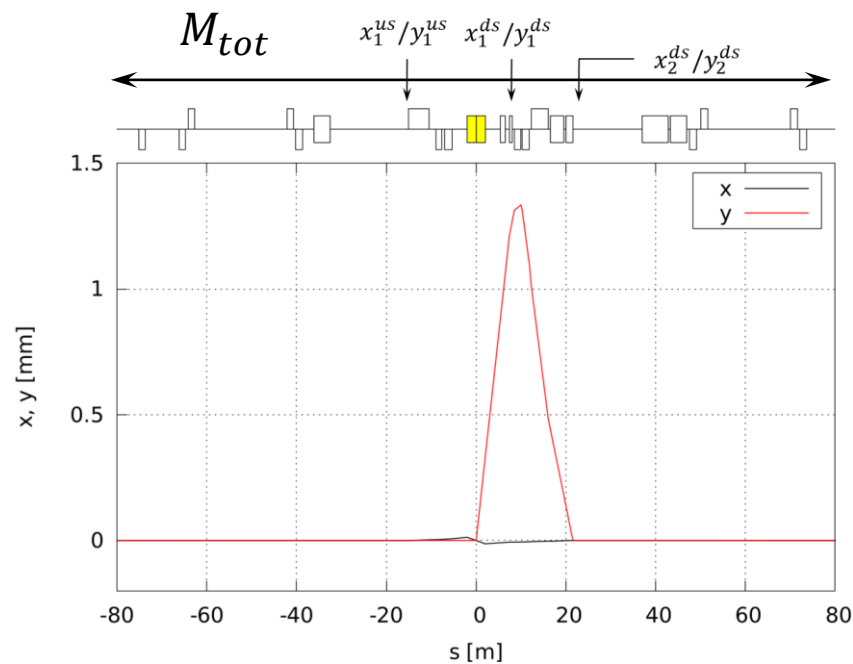
Top view



Solenoid + rotation compensated =
Full overlap \Rightarrow Full luminosity



Net Effect of IR on Ion Spin



Energy [GeV]	Long. n_x	Rad. n_y	Vert. n_z	φ_s [2π]
41	0.98	0.18	-0.076	0.038
100	0.97	0.23	-0.047	-0.035
200	0.93	-0.37	-0.029	0.021
275	0.88	-0.48	-0.025	-0.035