



Production and detection of nuclear fragments at the Electron-Ion Collider (EIC)

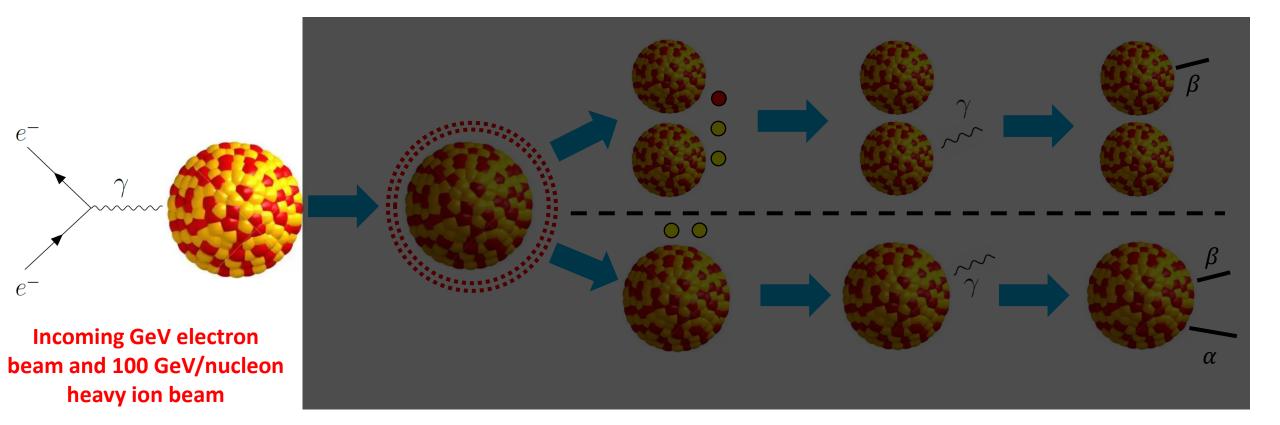
Barak Schmookler (for the EIC Rare Isotopes team)

Production of nuclear fragments

Motivating questions

- Can we use high-energy electron-heavy nucleus scattering at the future EIC to produce nuclear fragments, including exotic nuclei (i.e. undiscovered rare isotopes)?
- Can we go on to detect and correctly identify the produced nuclei? Can we also study the level structure of the nuclei by detecting gamma rays? What requirements does this place on the far-forward detection area?
- □If we can produce, detect, and identify nuclear fragments at the EIC, how can these results complement the work being done at dedicated rare isotope facilities?

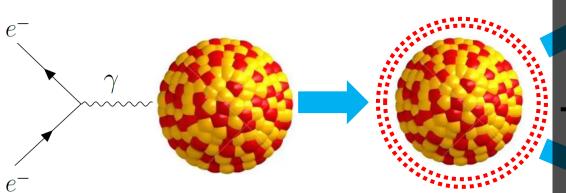
Nuclear fragment production at the EIC



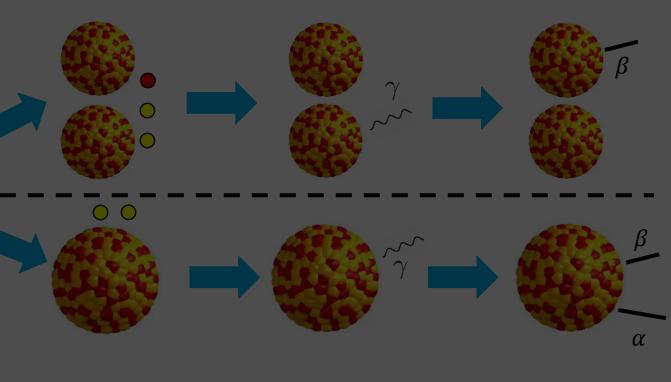
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Nuclear fragment production at the EIC

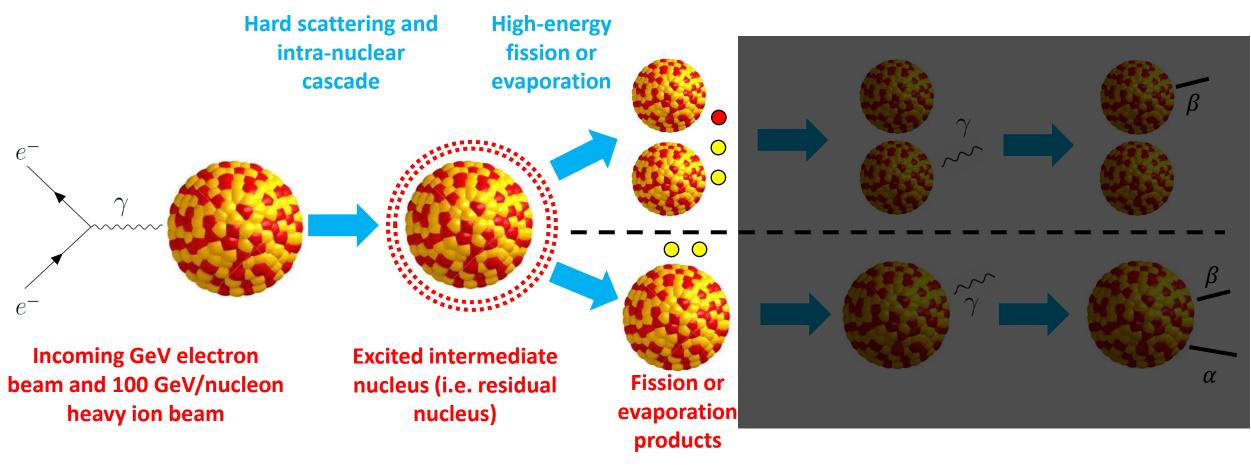
Hard scattering and intra-nuclear cascade



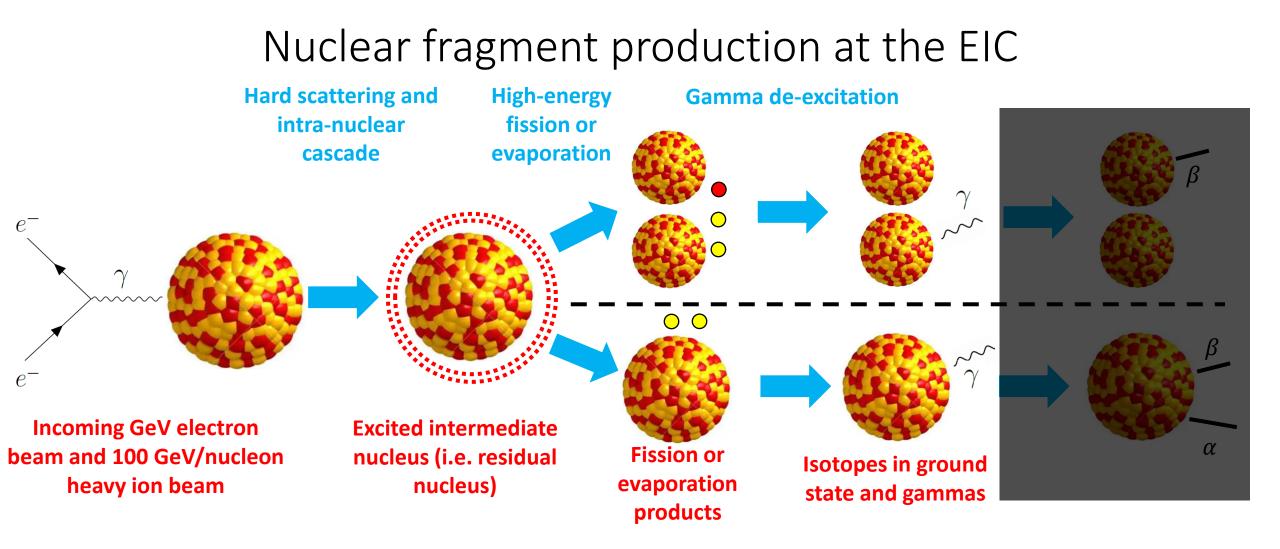
Incoming GeV electron beam and 100 GeV/nucleon heavy ion beam Excited intermediate nucleus (i.e. residual nucleus)



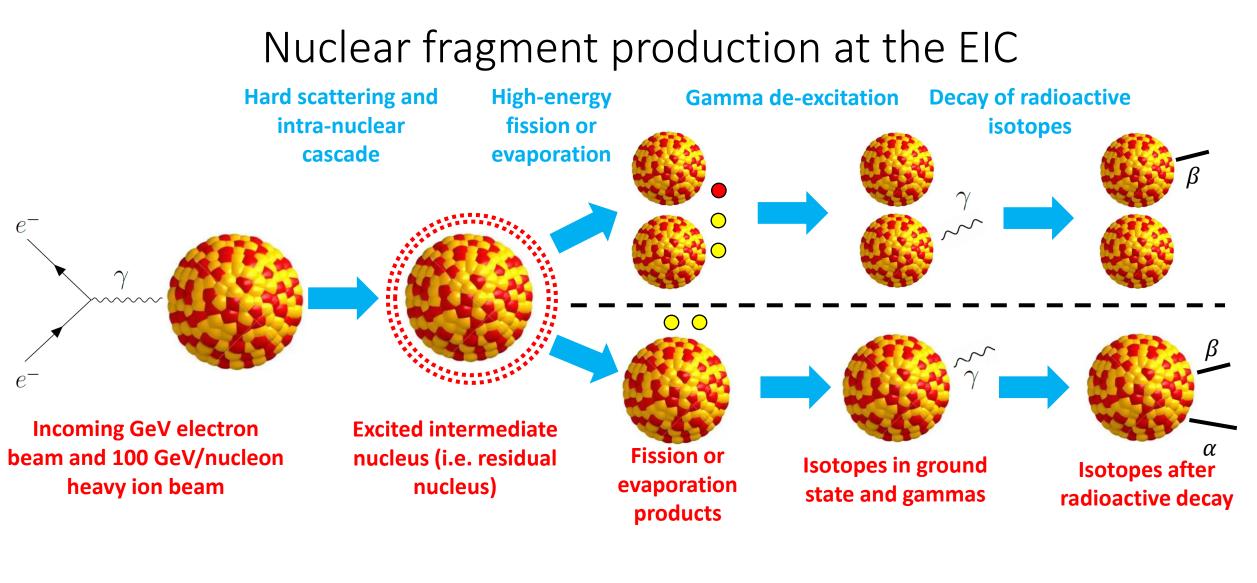
Nuclear fragment production at the EIC



t = 0



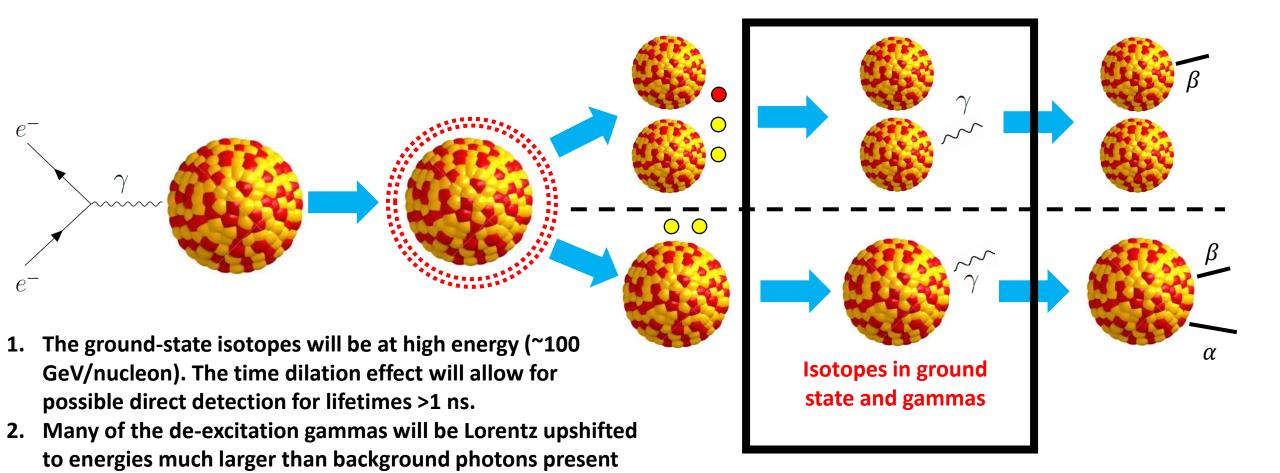
t = 0
$$t = 10^{-22} s$$
 t = $10^{-20} - 10^{-17} s$ t = $10^{-14} s$



t = 0

t = ? – never (stable)

Where the EIC can potentially contribute



in the detector area. This will allow for clean

used to study the level-structure of the isotopes.

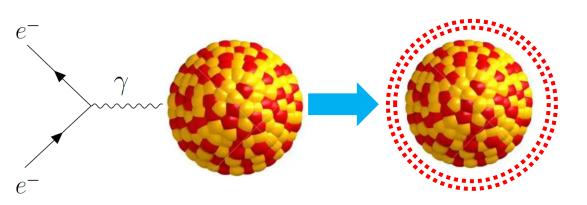
detection/identification of these gamma rays, which can be

Where the EIC can potentially contribute – specifics

Subject	Details
Reaction mechanism	 Excitation energy distribution. Improvement of the fast Abrasion-Fission model and a better understanding of the reaction mechanism. Simultaneous detection of two fission fragments and no target contribution to fragment kinematics. Improvement of production models.
Production of new isotopes	Production of new neutron-deficient isotopes in the Z=89-94 range. Advantages over RIB facilities due to short flight time and possibly higher production cross section.
Nuclear structure	Coincidence measurement of isotopes and de- excitation gammas.
Hadron formation time	Sensitivity of residual nucleus excitation energy distribution to formation time parameters.

How can we study this?

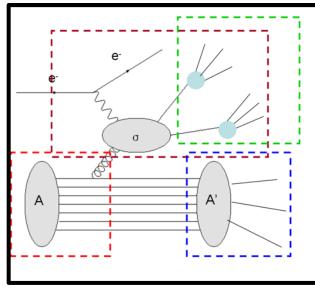
Hard scattering and intra-nuclear cascade



Incoming GeV electron beam and 100 GeV/nucleon heavy ion beam Excited intermediate nucleus (i.e. residual nucleus)

Step 1

The hard scattering (primary interaction) and the intranuclear cascade which follows are modelled using the *Benchmark eA Generator for Leptoproduction – BeAGLE* (Phys. Rev. D 106, 012007). This leaves us with the residual nucleus in an excited state.



A hybrid model consisting of DPMJet and PYTHIA with nPDF EPS09.

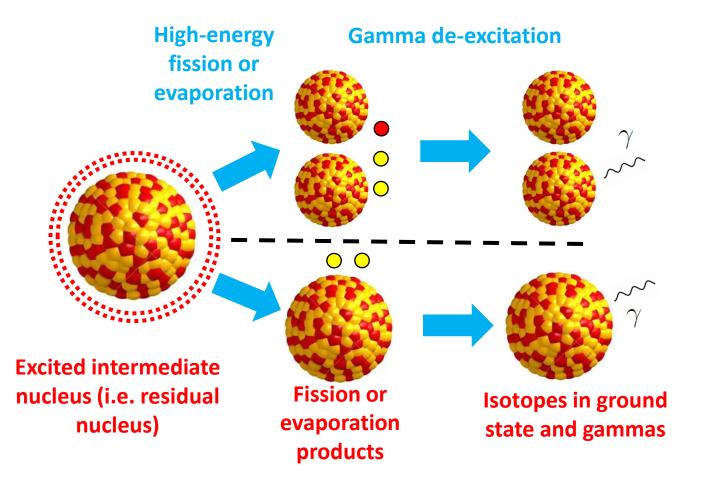
Nuclear geometry by DPMJet and nPDF provided by EPS09.

Parton level interaction and jet fragmentation completed in PYTHIA.

Nuclear evaporation (gamma dexcitation/nuclear fission/fermi break up) treated by DPMJet

Energy loss effect from routine by Salgado&Wiedemann to simulate the nuclear fragmentation effect in cold nuclear matter

How can we study this?

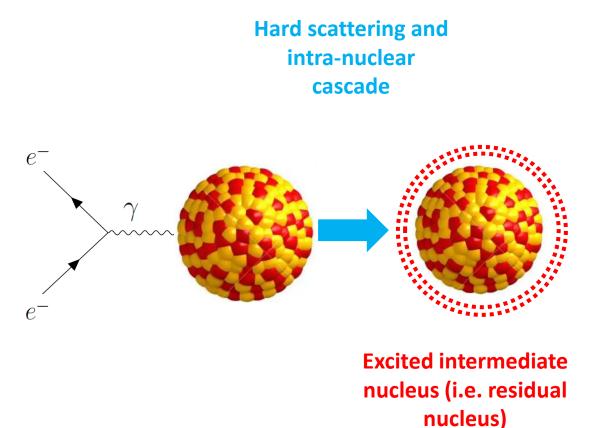


Step 2

For each event, the residual nucleus with a given A, Z, and excitation energy is then handed over to either *FLUKA* (Annals of Nuclear Energy 82, 10-18 (2015)) or ABLA07 for decay (evaporation or fission) followed by gamma de-excitation. We are left with the decay products of the residual nucleus.

FLUKA is used extensively in high-energy physics but has not been used for the study of rare isotope production.
 ABLA07 is used extensively in the rare isotope community – and is the second part of the abrasion-ablation code ABRABLA07. We run the BeAGLE events though both these codes and study the results.

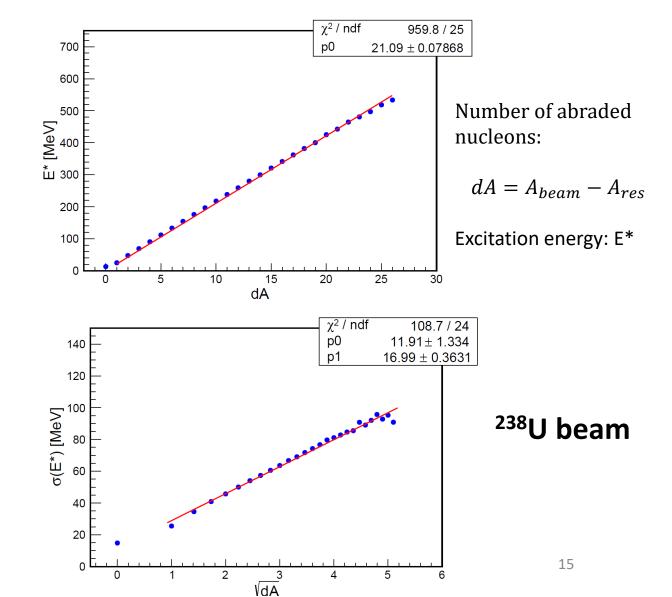
- □Using *BeAGLE*, we simulate an 18 GeV electron beam colliding with a 110 GeV/nucleon ²³⁸U or ²⁰⁸Pb beam.
- □We then study the excited residual nucleus that is created following the hard scattering and intra-nuclear cascade.
- The only relevant quantities are the A and Z and excitation energy of the residual nucleus. (The residual nucleus is assumed to have zero angular momentum.)



■We find that the production of the residual nucleus in *BeAGLE* manifests as a very simple abrasion model:

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 - The excitation energy shows a linear dependence on the number of abraded nucleons.

We plot the statistical mean and standard deviation here.



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The excitation energy shows a linear dependence on the number of abraded nucleons.

The E* distribution at fixed dA may be described with a Log-normal distribution, with some dependence the relative number of protons and neutrons abraded.

σ = 0.32 Amp = 42300.5 $\chi^2 \text{ LogN} = 1.866$ meanT = 21.634 Amp EED = 38923.346 250 $\chi^2 EED = 7.172$. FED Lognor Data Energy Bins (MeV) Log-Scale $\mu = 5.082$ σ = 0.32 Amp = 42300.5 χ² LogN= 1.866 meanT = 21.634104 Amp EED = 38923.34 $\chi^2 EED = 7.172$ • EED Lognorr Data Energy Bins (MeV) LogNorm Median / dA Sum of Median Column Labels

Excitation Energy Distribution for Z=88, A=230 with dN =4 and dZ = 4 Bin Number: 1 Sum: 42301

μ = 5.082

Number of abraded nucleons:

$$dA = A_{beam} - A_{res}$$

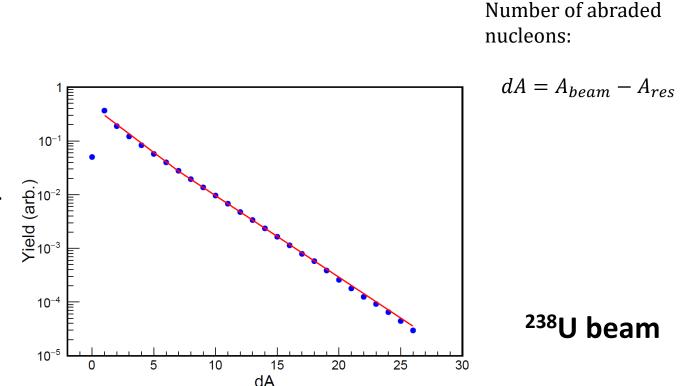
Excitation energy: E*

238

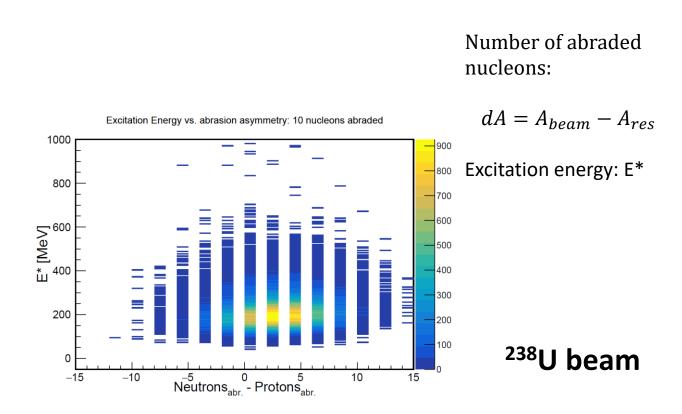


beam

- ■We find that the production of the residual nucleus in *BeAGLE* manifests as a very simple abrasion model:
 - The excitation energy shows a linear dependence on the number of abraded nucleons.
 - The cross section for abrading a given number of nucleons (for dA>1) shows a (piecewise) exponential dependence.



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Note: The observed simple abrasion model comes out of *BeAGLE* 'naturally'. The simulation uses an intra-nuclear cascade model and a nuclear potential model to determine the A, Z and excitation energy of the residual nucleus. The ground state mass model comes from *FLUKA*.

Intra-nuclear cascade hadron formation time:

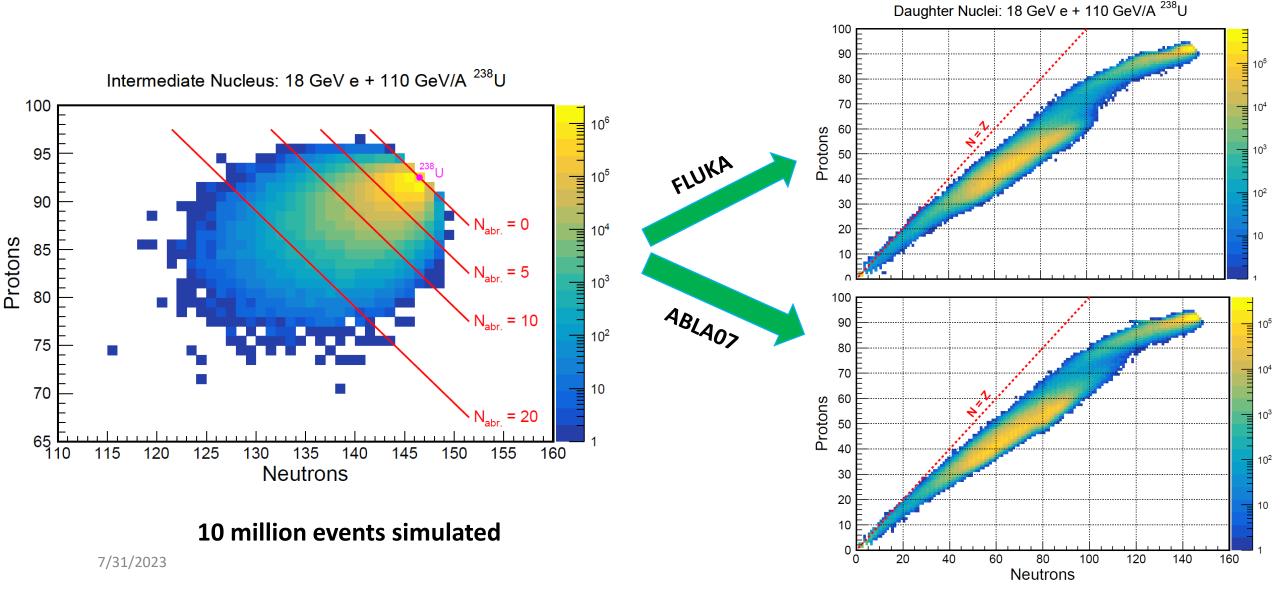
$$\tau_{Lab} = \tau_0 \frac{E_s}{m_s} \frac{m_s^2}{m_s^2 + p_{s\perp}^2}$$

Mass (excitation energy) of the residual nucleus:

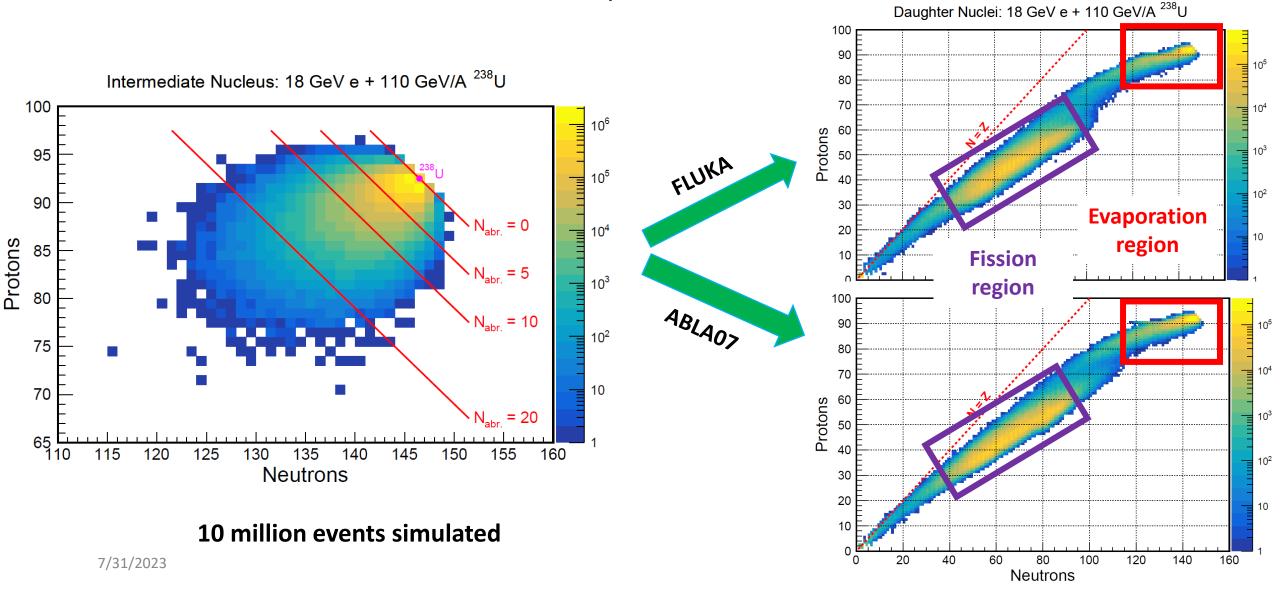
$$(E_{res}, \mathbf{p_{res}}) = (M_A, \mathbf{0}) - \sum_{i=1}^{N_w} \left(E_i^F, \mathbf{p_i^F} \right) + (E_{rec}, \mathbf{p_{rec}})$$

Z.Phys. C70 (1996) 413-426 Z. Phys. C 71, 75-86 (1996)

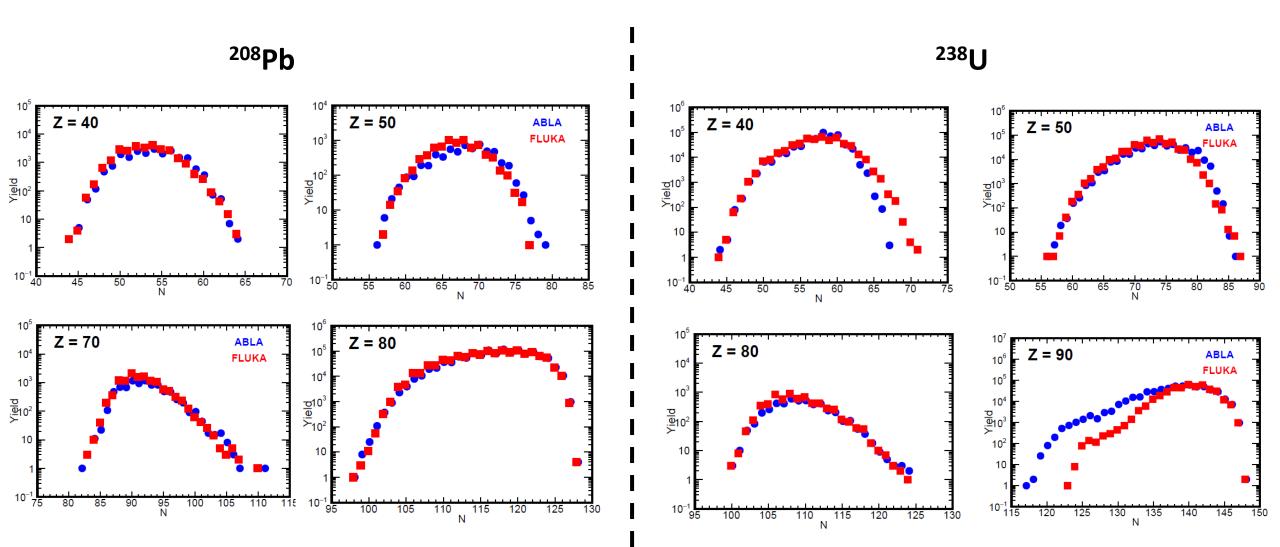
We can then decay the residual nucleus



We can then decay the residual nucleus



FLUKA and ABLA07 are largely in agreement about EIC production rates

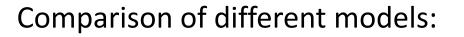


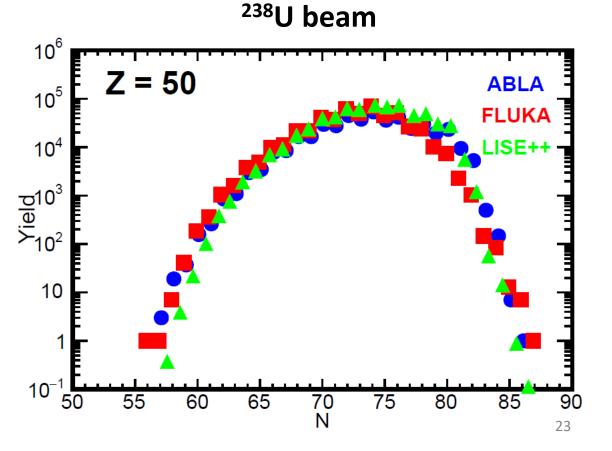
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Fission fragment production can also be studied with LISE++

Based on *BeAGLE* findings above, an Exponential Abrasion Model has been implemented in *LISE++*:

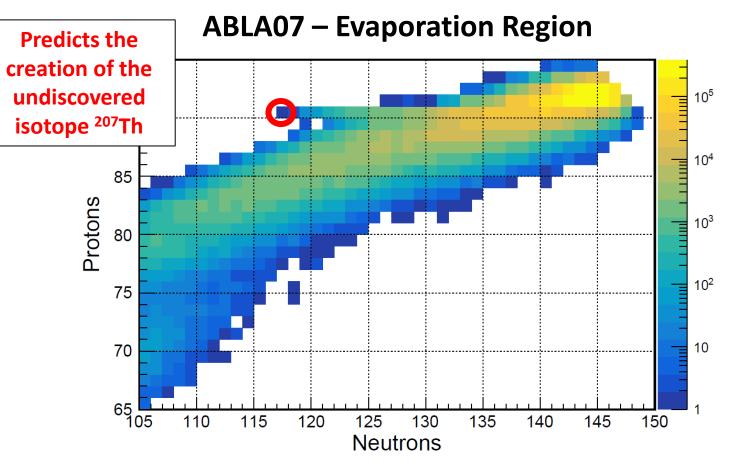
Excitation Energy of prefragment		
A Flement 7	Global Abrasion Cross-Section Factor = 1 (default 1)	
132 Sn 50 🗁 Z 🖨	A. J.W.Wilson, L.W.Towsend, F.F.Badavi, NIM B18 (1986) 225-231 geometrical model	
B [°] decay	Excitation Energy = 0 MeV π^* (μ , c , h , c) , π	
	Exclasion Energy = $E^* = (\gamma \cdot f \cdot \Delta S)_{geom} + E_{friction}$ Standard deviation = 98.84 MeV	
Reaction 236U (110.0 GeV/u) + H	gamma = 0.95 MeV/tm ² Excitation Energy Transfer (friction)	
Excitation Energy in the code = 1562.84 MeV	sigma = 9.6 *d. abr ¹² [MeV] $E_{\text{tricion}} = \text{coef}_1 * C_p + \text{coef}_2 * C_p * C_1$	
Abrasion model	Correction factor of Surface distortion avoidation	
O Geometrical : J.Gosset et al., PRC 16 (1977) 629	Correction factor of Surface distortion excitation surface interface, C ₁ is the chord of intersection coef. = 6.5 C _n = 14.6 fm	
Exponentional (Y~ exp(-k*d_abr) k= 0.363	f = 1 + c ₁ * d_abr / Ap + c ₂ * (d_abr / Ap) ²	
C A. J.W.Wilson et al., NIM B18 (1987) 225-231	c1= 1.5 c2= 2.5 f= 2.16 Use Friction = MeV	
B. JJ.Gaimard and KH.Schmidt, NPA531 (1991) 709 B. JJ.Gaimard an		
C. Parametrized Gaussian distribution	(MeV)	
O D. Exponential excitation-energy distribution	TO sigma = 16.5 *d_abr ^{1/2} [MeV] Standard deviation = 168.91 MeV	
Apply the limiting temperature threshold: T=min(T.Tlim)	C. Parametrizied Gaussian distribution – simplified combination from NPA710 (2002) 157	
"Isospin-thermometer model", corresponds to Fig.9 K-H.Schmidt et al., NPA710 (2002) 157	<pre><pre></pre> sigma Mean Excitation Energy = 1562.84 MeV</pre>	
	Signa Otenderd devicing a dan da	
Use LISE++ corrections for Geometric A-A model	-0.0737 *d_abr ²⁺ -1.1644 *d_abr + Standard deviation = 133.44 MeV	
Apply thermalization for Excitat energy according to JJ.Gaimard KH.Schmidt. NPA531 (1991) 709:	22.556 *d_abr+ 24.949 *d_abr ¹² +	
see Equation 3.4	0 [MeV] 0 [MeV] 4-bit is the projectile mass, d-abr is the number of abraded nucleons	
D. Exponential excitation-energy distribution - LAudriac et al., PRC88, 041602(R) (2013)		
Mean Temperature <= 13 *d abr (MeV) Mean Excitation Energy = 1378 MeV		
Plot as f (Z_pf) Make default		
V OK X Cancel ? Help		





Using our small simulation sample, we see hints of interesting physics

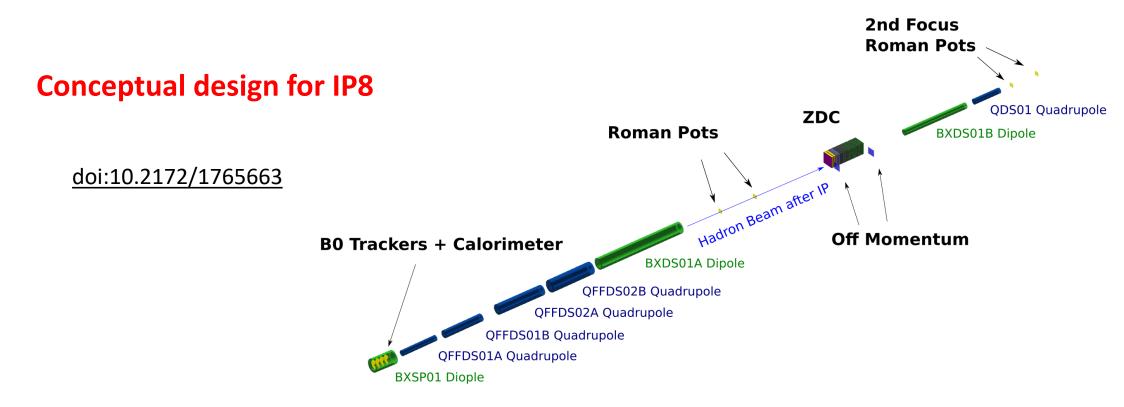
Production of new neutron-deficient isotopes in the Z=89-94 range.Advantages over RIB facilities due to short flight time and possibly higher production cross section.



We need to simulate many more events to model the production rates at the EIC; or use LISE++ for an analytical approach

Detection of nuclear fragments

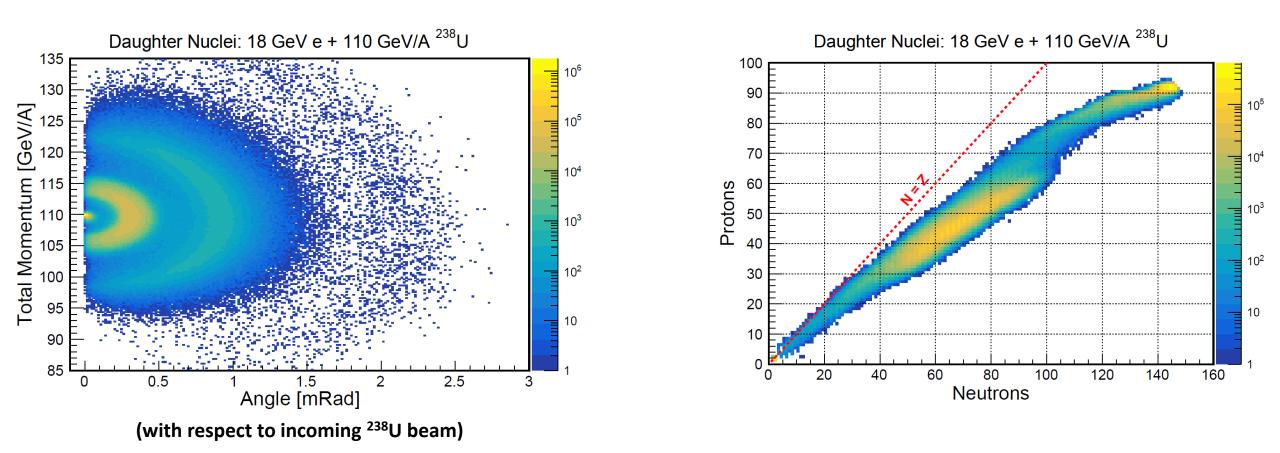
EIC Detectors – far-forward region

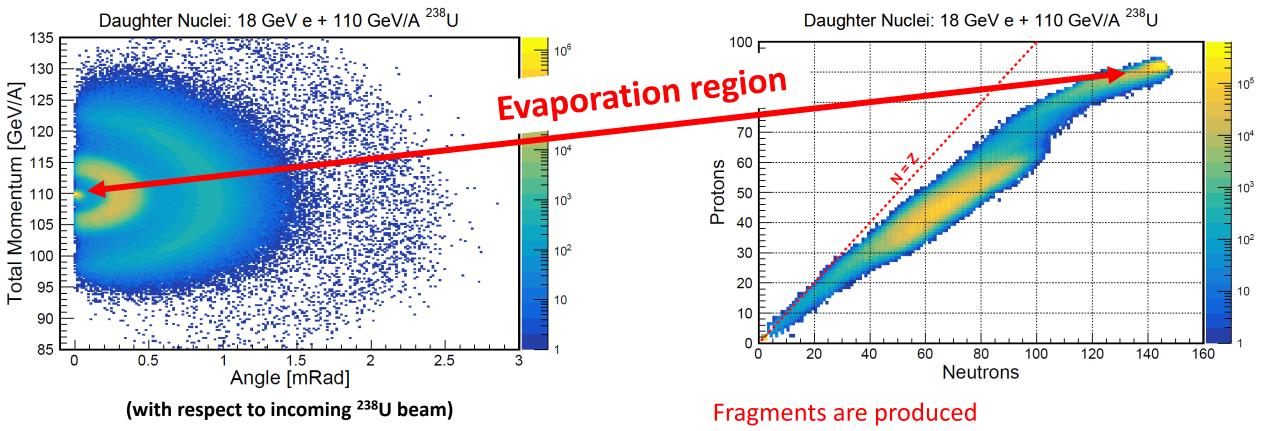


The nuclear fragments can be measured using detectors in the (second set of) Roman Pots (RP) – two tracking planes to measure local positions and angles.

Gamma rays can be detected using the Zero-Degree Calorimeter (ZDC).

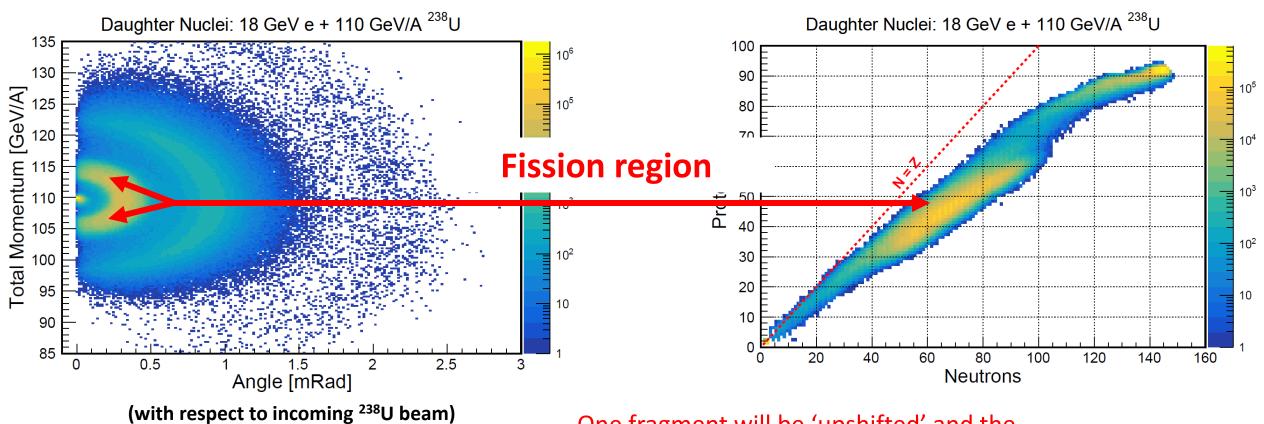
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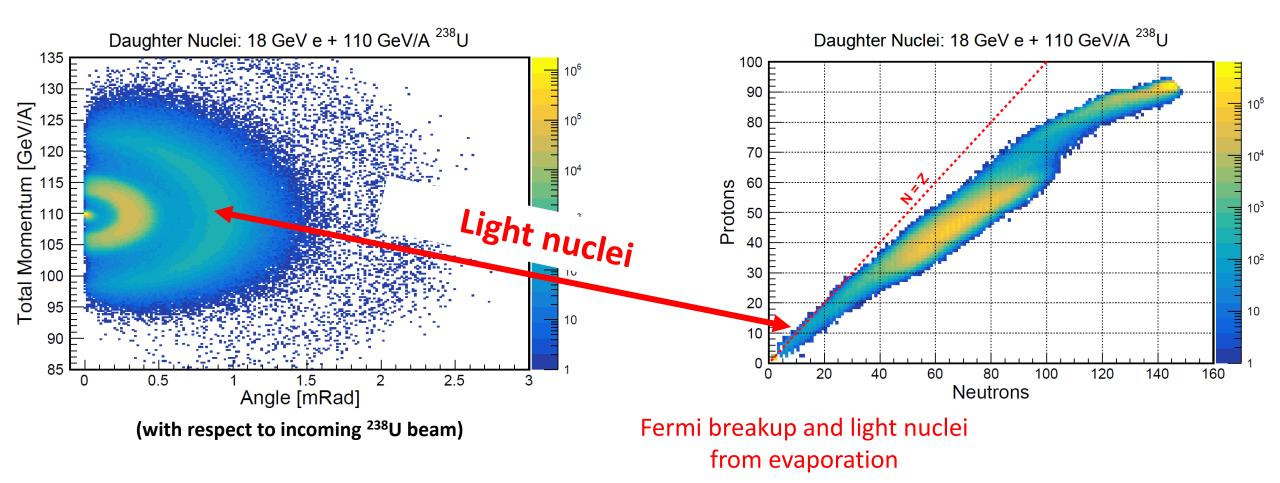
parallel to beam with same momentum per nucleon.

BeAGLE + FLUKA



One fragment will be 'upshifted' and the other 'downshifted'. Both fission fragments can be registered in coincidence.

BeAGLE + FLUKA



Principle of detection – rigidity measurement

At first approximation the momentumper-nucleon of the outgoing fragment (p_N) is the same as the momentum-pernucleon of the incoming beam $(p_{N,beam})$.

$$x_{L} = \frac{R}{R_{beam}} = \left[\frac{\left(\frac{Ap_{N}}{Z}\right)}{\left(\frac{A_{beam}p_{N,beam}}{Z_{beam}}\right)} \right]$$
$$= \left[\frac{\left(\frac{A}{Z}\right)}{\left(\frac{A_{beam}}{Z_{beam}}\right)} \right]$$
Measurement of rigidity (x_L) determines the fragment A/Z ratio

Some definitions

$$\overrightarrow{Fragment Rigidity (R)} = \frac{p}{Z}$$

$$\overrightarrow{R} = \frac{R}{R_{beam}}$$

$$\overrightarrow{R} = \frac{R}{R_{beam}}$$

$$\overrightarrow{R} = \frac{R - R_{beam}}{R_{beam}} = x_L - 1$$

Principle of detection – rigidity measurement

The hit position at the Roman Pot (RP) detectors in the dispersive direction:

 $x_{RP} = D_x(-R_{Rel}) = D_x(1-x_L)$

Additional definitions



At Roman Pots:

Dispersion (D_x) Beta Function (β_x)

Minimum allowed hit position at the RPs to exclude beam envelope:

$$x_{RP}^{min} = 10\sigma_x = 10\sqrt{\beta_x\varepsilon_x + D_x^2\sigma_p^2}$$

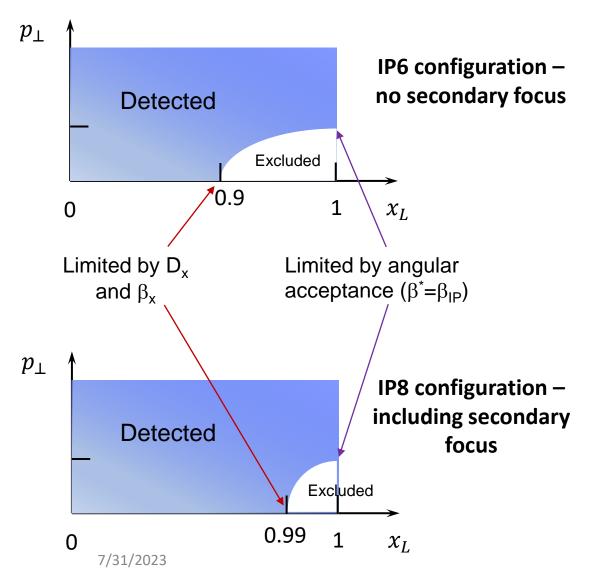


Accelerator parameters (EIC CDR Table 3.5):

Beam Emittance (ε_x) = 43.2 *nm*

Momentum spread (σ_p) = 6.2 × 10⁻⁴

Acceptance for fragments in IP6 and IP8



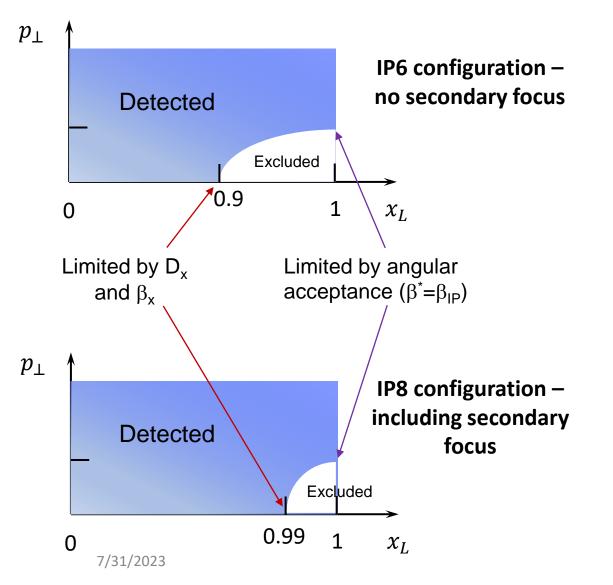
IP6 acceptance at first RP (using the highdivergence 10x100 GeV shifted lattice):

 $\beta_x = 865 m$ $D_x = 16.7 cm$ $\rightarrow x_{RP1}^{min} = 6.11 cm$

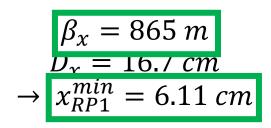
IP8 acceptance at first RP:

$$\beta_x = 2.28 m$$
$$D_x = 38.2 cm$$
$$\rightarrow x_{RP1}^{min} = 0.39 cm$$

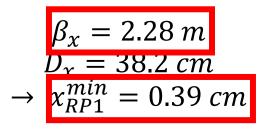
Acceptance for fragments in IP6 and IP8



IP6 acceptance at first RP (using the highdivergence 10x100 GeV shifted lattice):



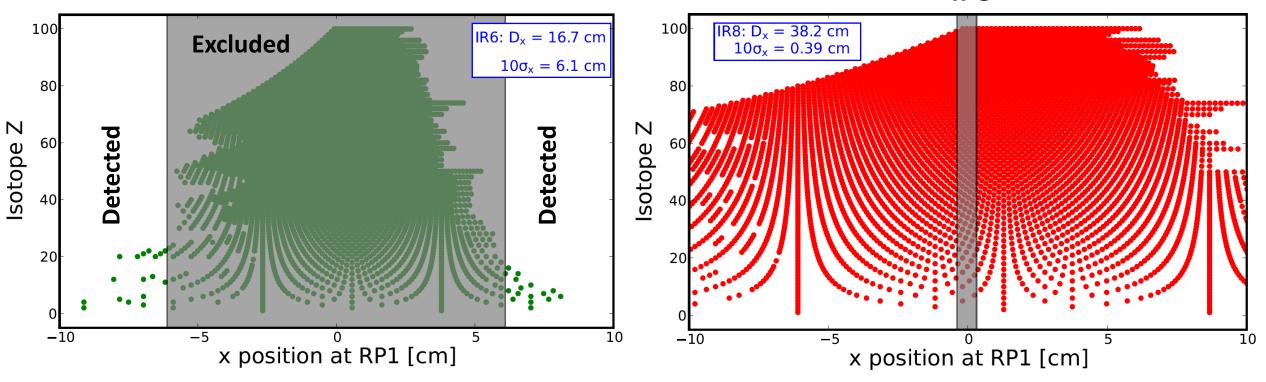
IP8 acceptance at first RP:



Acceptance for fragments in IP6 and IP8

IP6

IP8



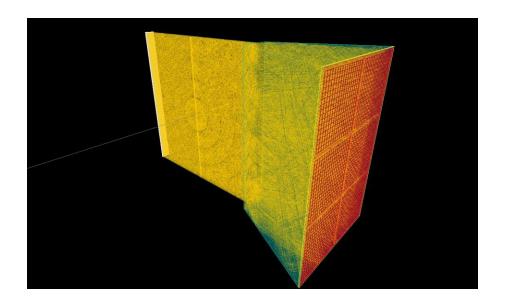
Each point is an individual isotope. All known and potential isotopes which come from a combined *NNDC* and *LISE++* database are included.

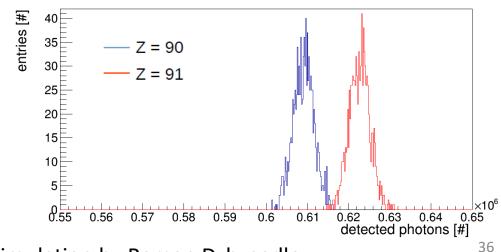
Assuming a RP position resolution of 10-100 microns, isotopes with the same Z are well separated.

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Full reconstruction of the fragments

- 1. The charge of the isotope (Z) must be determined. This can potentially be done using a thin (few mm thick) quartz bar placed inside the RP (behind the tracker) at the second focus. The quartz bar would be perpendicular to the beam, extended along the dispersive (x) direction. The number of Cherenkov photons produced will be quite large (proportional to Z²).
- 2. In the fission region, the outgoing isotopes do not have the same momentum-per-nucleon as the ion beam. This can be corrected for by measuring the angles at the RP detectors and registering both fission fragments in coincidence.

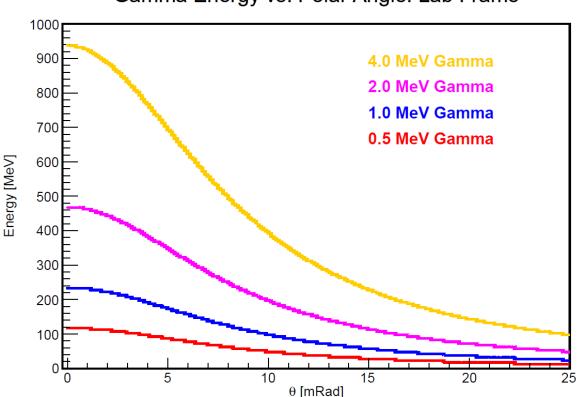




Simulation by Roman Dzhygadlo

Detection of gamma rays

Single gamma simulation – 110 GeV/A beam



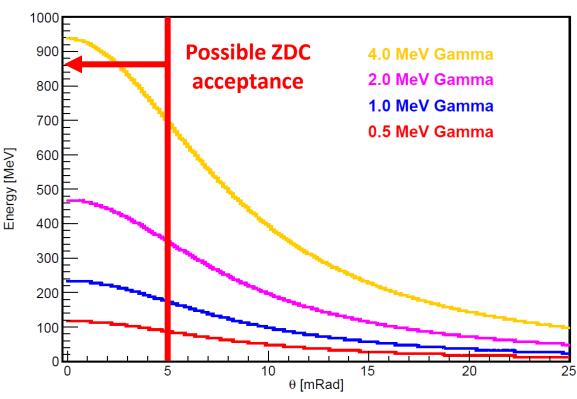
Gamma Energy vs. Polar Angle: Lab Frame

- Gamma rays from nuclear de-excitations can be detected in the Zero-Degree Calorimeter (ZDC). The ZDC acceptance range will be approximately 0-5 mRad.
- The energy resolution of the ZDC for photon detection may be as good as $2\%/\sqrt{E (GeV)}$ if a material such as LYSO crystals are used.
- □We will therefore be able to measure gamma rays which are Lorentz upshifted and moving very close to the ion beam direction.
- ■A 1 MeV gamma will have an energy of ~240 MeV at zero degrees in the lab frame. For the ZDC resolution above, this gamma will have its energy measured to 4% in the lab frame. At first approximation, the energy resolution in the nucleus' rest frame is equivalent – that is, a 40 keV resolution for a 1 MeV gamma.

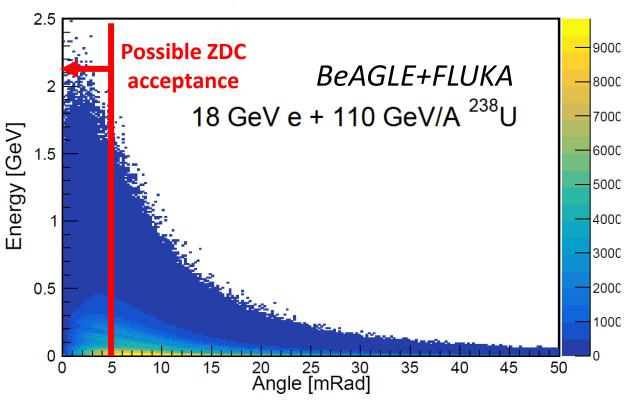
Detection of gamma rays

Single gamma simulation – 110 GeV/A beam

Gamma Energy vs. Polar Angle: Lab Frame



De-excitation gammas: full simulation results



Z.Phys. C70 (1996) 413-426 Z. Phys. C 71, 75-86 (1996)

Summary and ongoing work

Our simulation studies suggest the EIC has the potential to produce nuclear fragments using various heavy-ion beams. We believe that measuring these fragments can complement current and future work being done at dedicated rare isotope facilities.

□We are working to implement the current official IR8 lattice (see <u>https://wiki.bnl.gov/eic-detector-2/index.php?title=Project Information</u>) into *Geant4* to conduct some detailed simulation studies. We are conducting additional simulations with BeAGLE and working to compare the residual nucleus results to published data.

□With the right combination of detectors, these nuclei can be reconstructed using the proposed optics design of the 2nd interaction point using detectors located at a secondary focus.

- □Our studies also suggest that de-excitation gamma rays can be measured in coincidence with the nuclear fragments to quite high resolution.
- □Given the time scales for the EIC project and the 2nd interaction region in particular there is sufficient time to conduct further studies on the potential of the EIC to contribute to this physics, as well as place requirements on the far-forward spectrometer optics and detector design.

The EIC Rare Isotopes Team



- I. Mark Harvey
- II. Mark Ddamulira

Facility for Rare Isotope Beams at Michigan State University

I. Oleg TarasovII. Isaiah Richardson



- I. Abhay Deshpande
- II. Pawel Nadel-Turonski
- III. Brynna Moran
- IV. Charles Joseph Naim
- V. Niseem Magdy
- VI. Wenliang Li



I. Barak Schmookler



I. Ketevi Assamagan

Thanks!

Acknowledgements

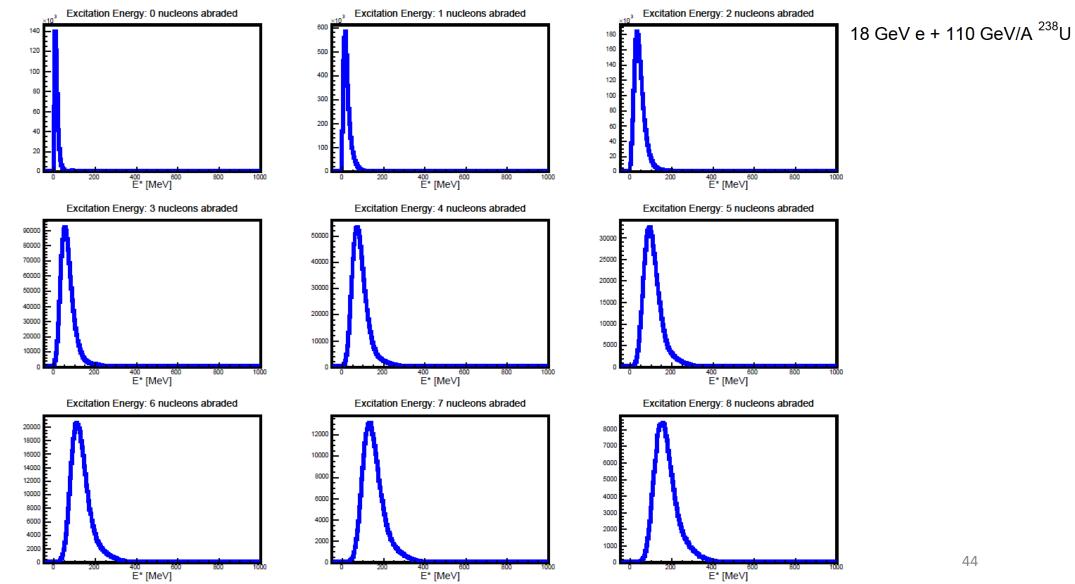
- Thanks to Mark Baker and Kong Tu for help with the BeAGLE event generator!
- Thanks to Aleksandra Kelic-Heil for providing access to the *ABRABLA07* code, as well as instructions on running the ablation portion!



Backup Slides

Residual nucleus excitation energy distributions

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Residual nucleus excitation energy distributions Excitation Energy: 9 nucleons abraded Excitation Energy: 10 nucleons abraded Excitation Energy: 11 nucleons abraded 18 GeV e + 110 GeV/A ²³⁸U ⁴⁰⁰ E* [MeV][∞] ⁴⁰⁰ E* [MeV] ⁴⁰⁰ E* [MeV] Excitation Energy: 12 nucleons abraded Excitation Energy: 13 nucleons abraded Excitation Energy: 14 nucleons abraded E* [MeV] E* [MeV] E* [MeV] Excitation Energy: 15 nucleons abraded Excitation Energy: 16 nucleons abraded Excitation Energy: 17 nucleons abraded 25(

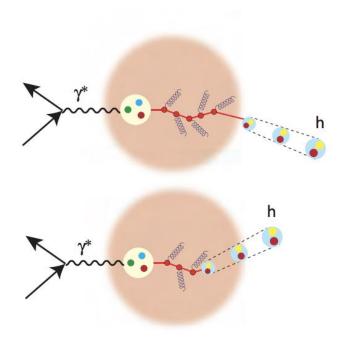
E* [MeV]

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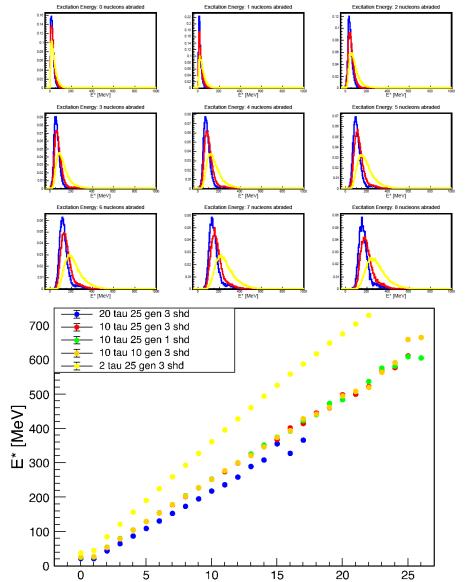
E* [MeV]

E* [MeV]

Residual nucleus sensitivity to formation time parameter



$$\tau_{Lab} = \tau_0 \frac{E_s}{m_s} \frac{m_s^2}{m_s^2 + p_{s\perp}^2}$$



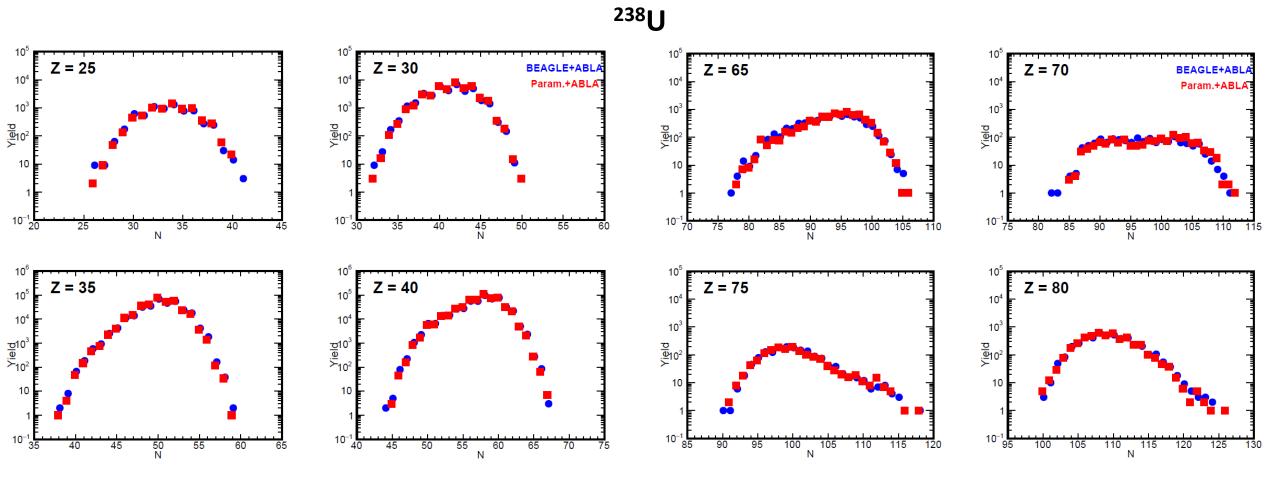
dA

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Expected EIC event counts

- Event rates at the EIC will be on the order of 10,000 events per second. Most of these events are at very low Q² (the photoproduction region of the e-p/A total cross section), but nuclear fragments can still be produced and detected in for these kinematics.
- The 10 million event sample which we generated may correspond to less than an hour of EIC running. Generating a larger number of events with *BeAGLE* becomes computationally expensive.
- □Since all we care about here is the production of the residual nucleus, we can create a simple empirical parameterization of the abrasion model observed in *BeAGLE*.

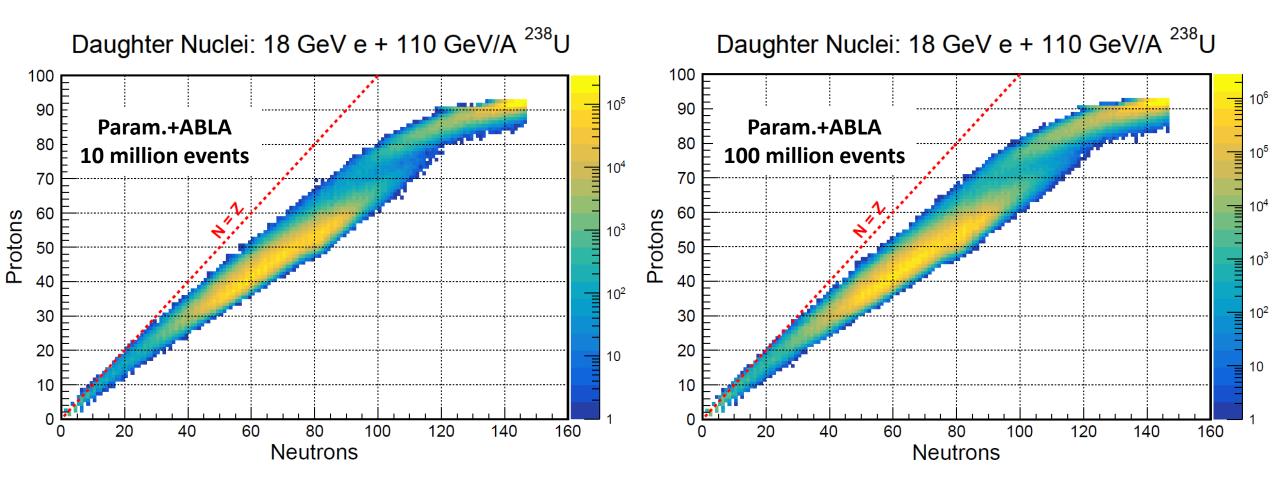
Comparison of *BeAGLE* results and parameterized distribution



Using our parameterized model for the residual nucleus, we can simulate 10 million events in a few minutes.

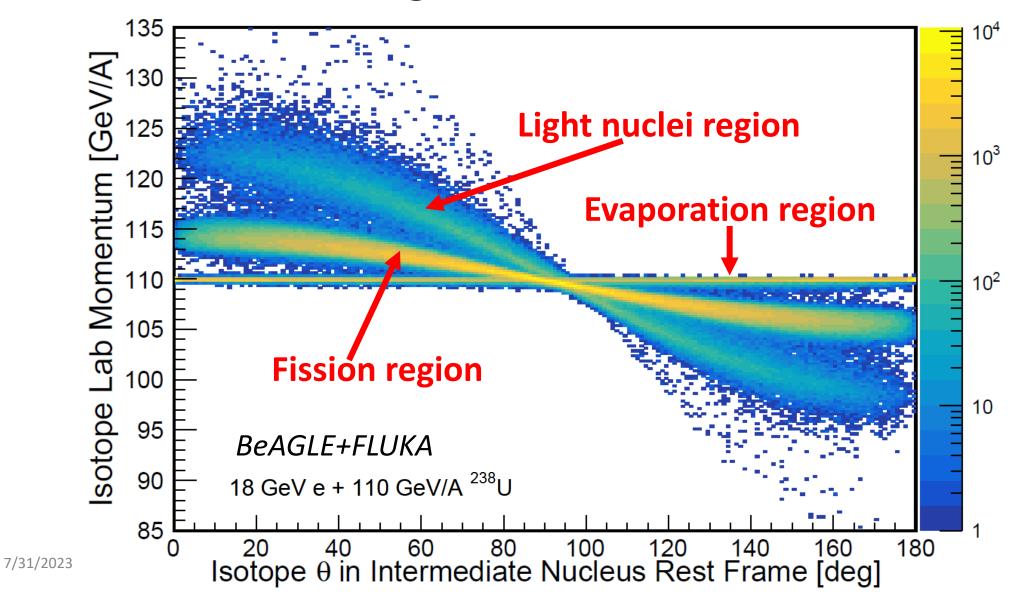
The results are very consistent with using the full *BeAGLE* simulation.

Towards higher statistics simulations



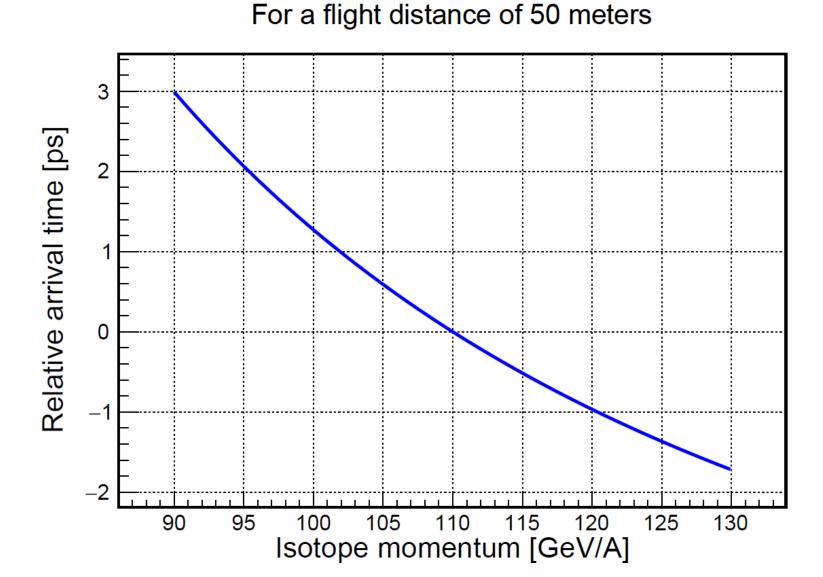
Note how borders expand towards more unstable isotopes as additional events are generated.

Fragment kinematics

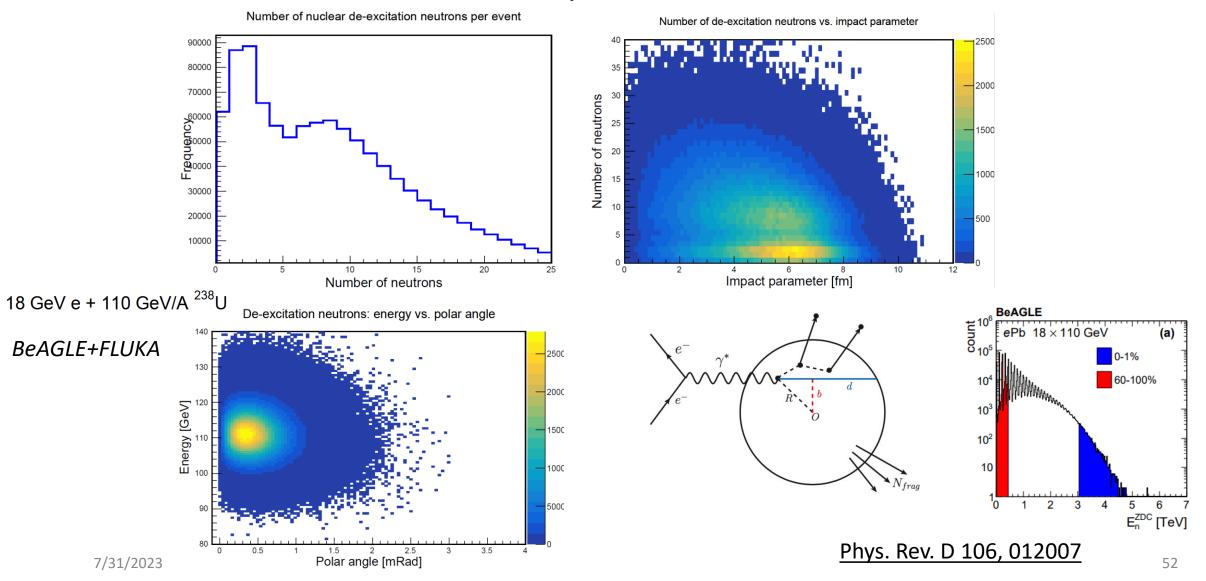


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Time-of-flight measurements would require picosecond resolution



Centrality determination



Centrality determination – model sensitivity

