

Luka Leskovec

A lattice QCD study of $B \rightarrow \rho \ell \bar{\nu}$

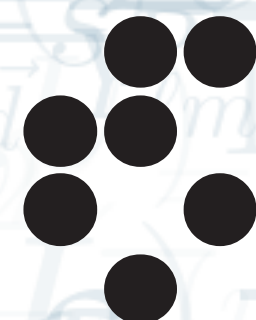
ACHT 2023

Non-perturbative aspects of
Nuclear, Particle and Astroparticle Physics

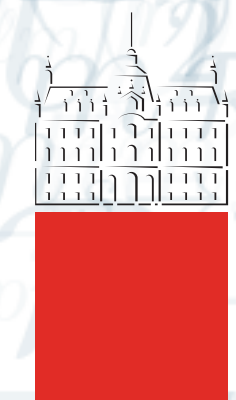
Schloss Retzhof, Austria

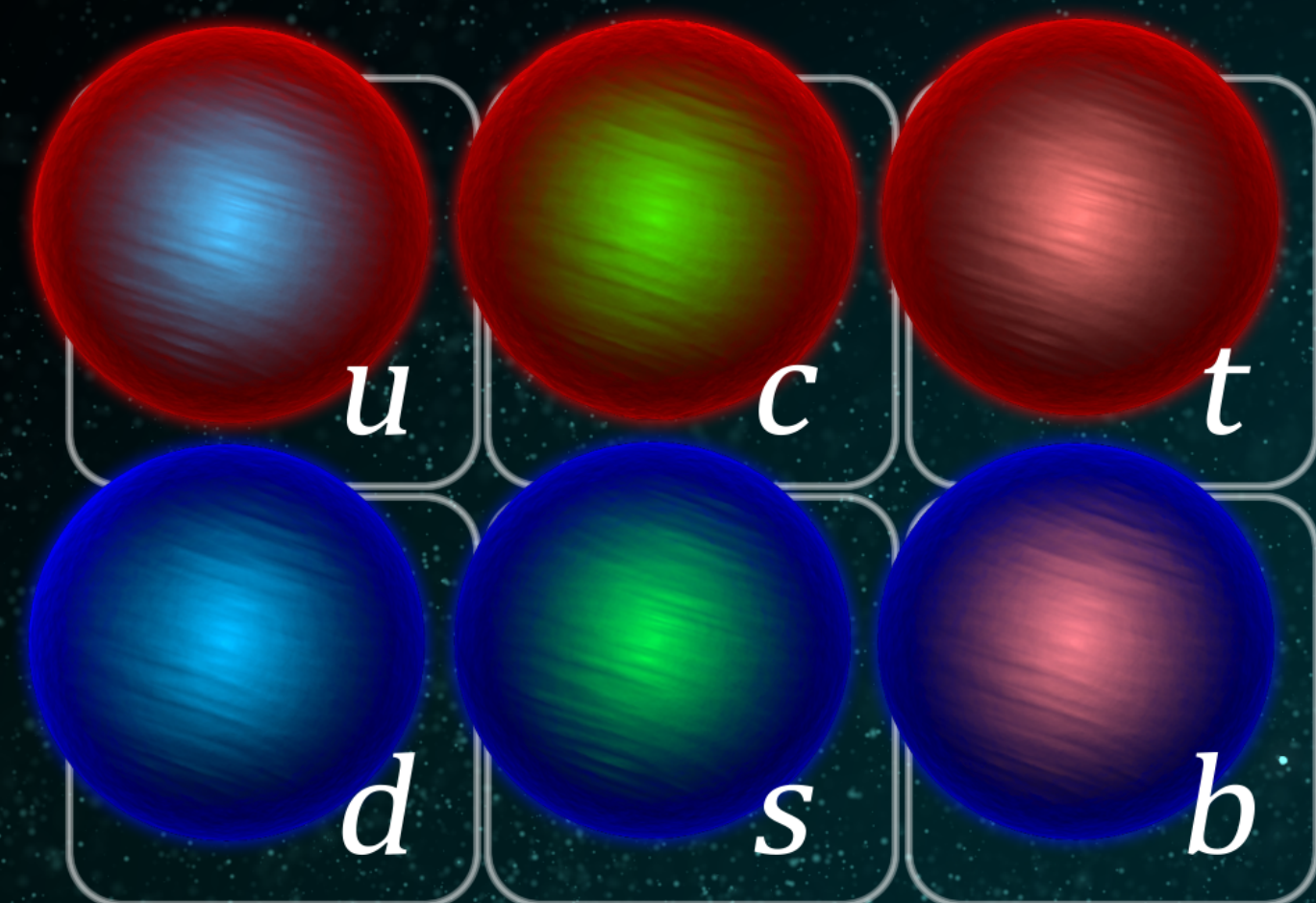
September 27-29, 2023

University of Ljubljana
Faculty of *Mathematics and Physics*

 **Institut**
"Jožef Stefan"
Ljubljana, Slovenija

in collaboration with:
S. Meinel, M. Petschlies, C. Alexandrou,
J.W. Negele, S. Paul, A. Pochinsky

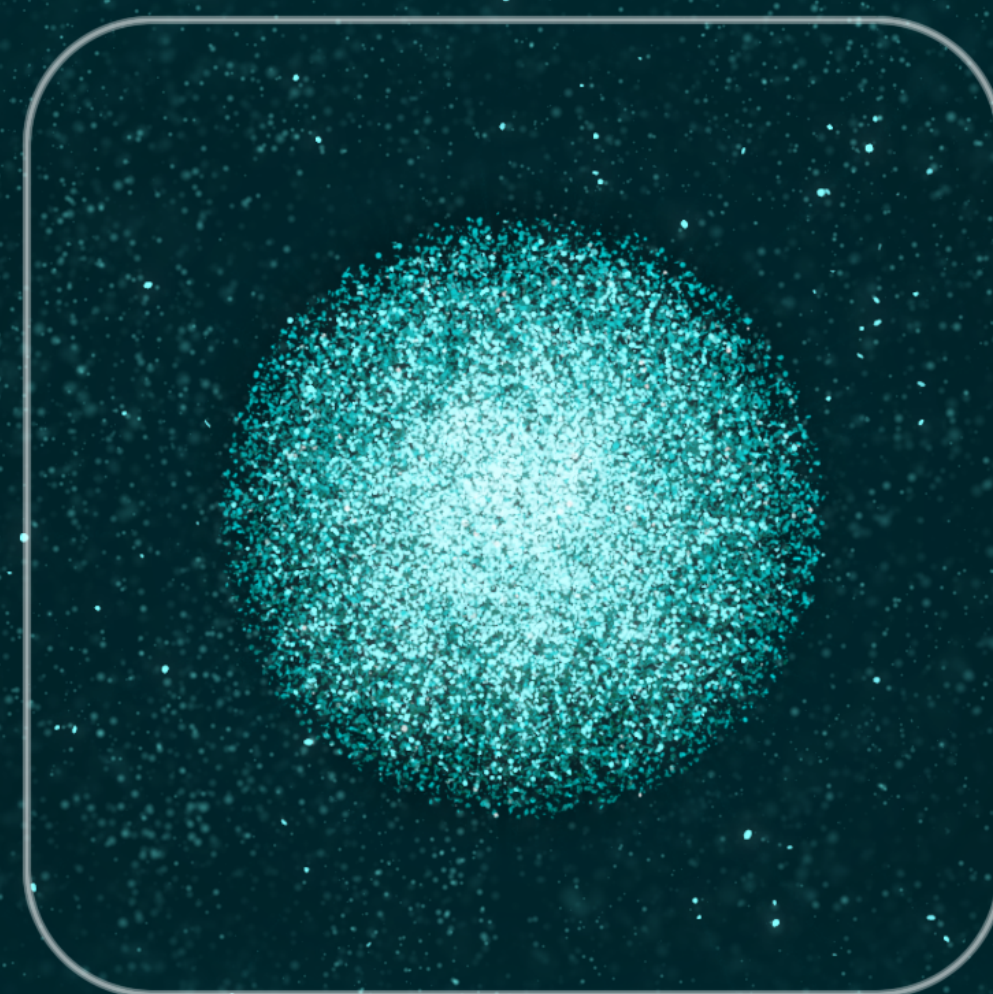




Quarks



Leptons



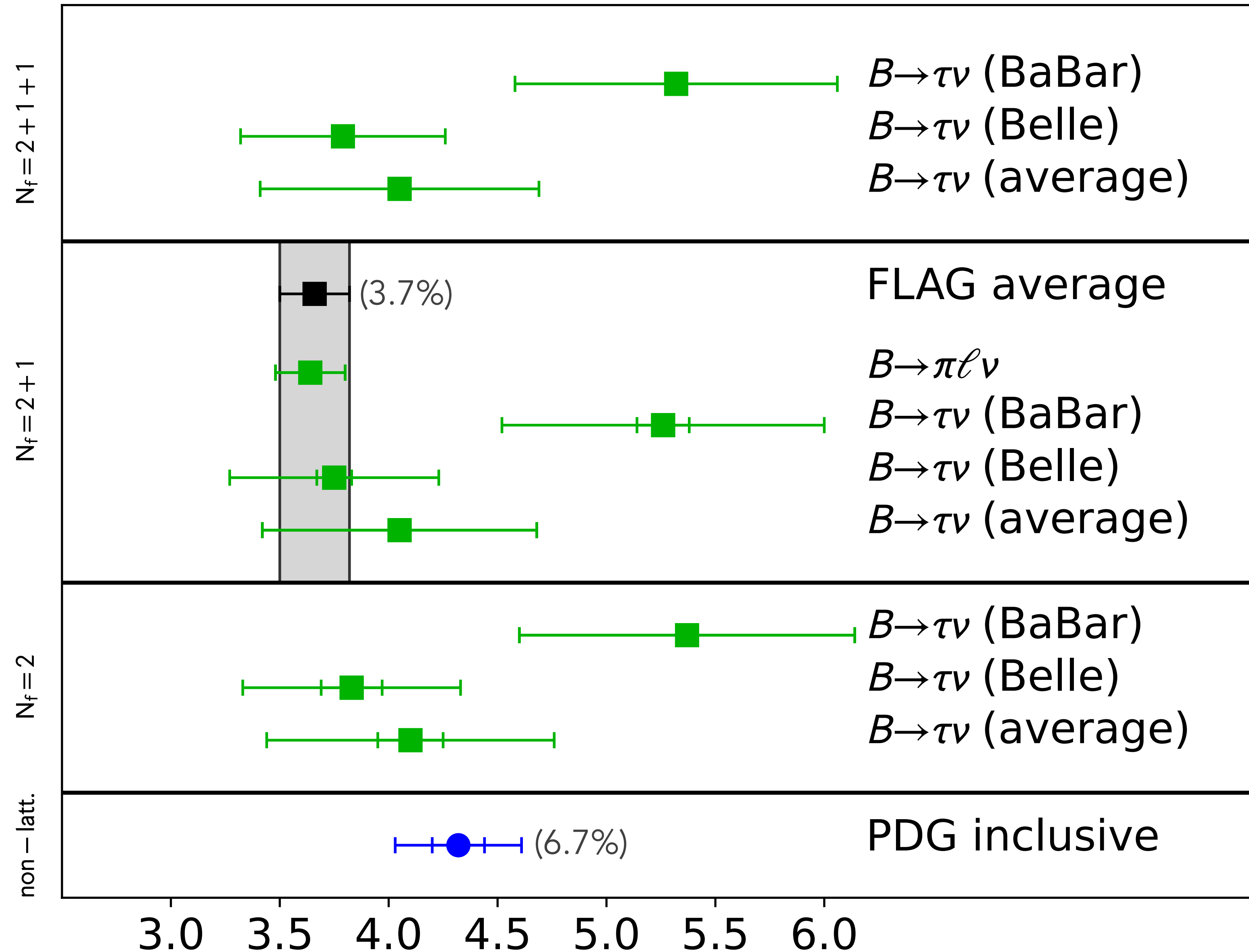
Higgs boson



Forces

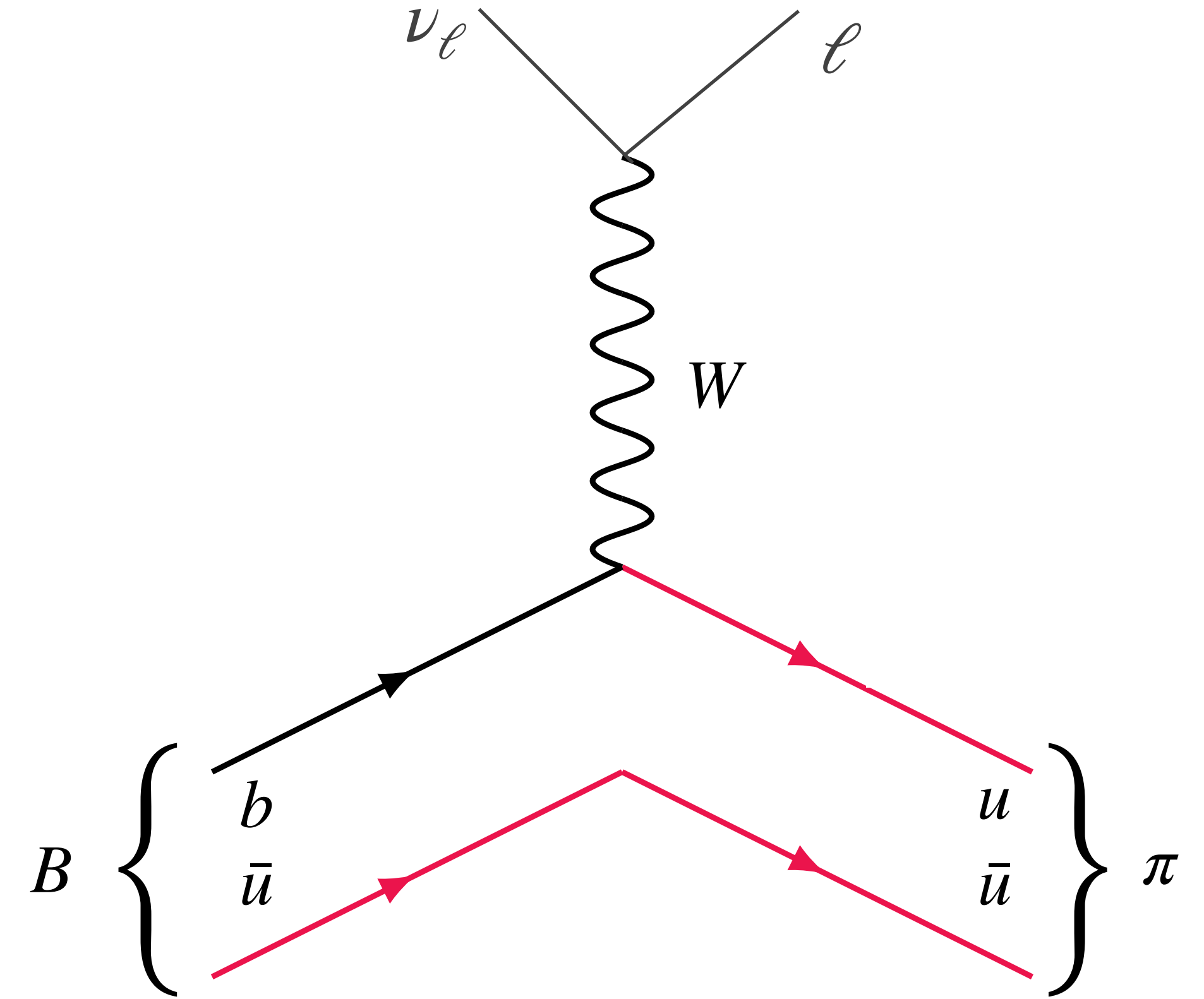
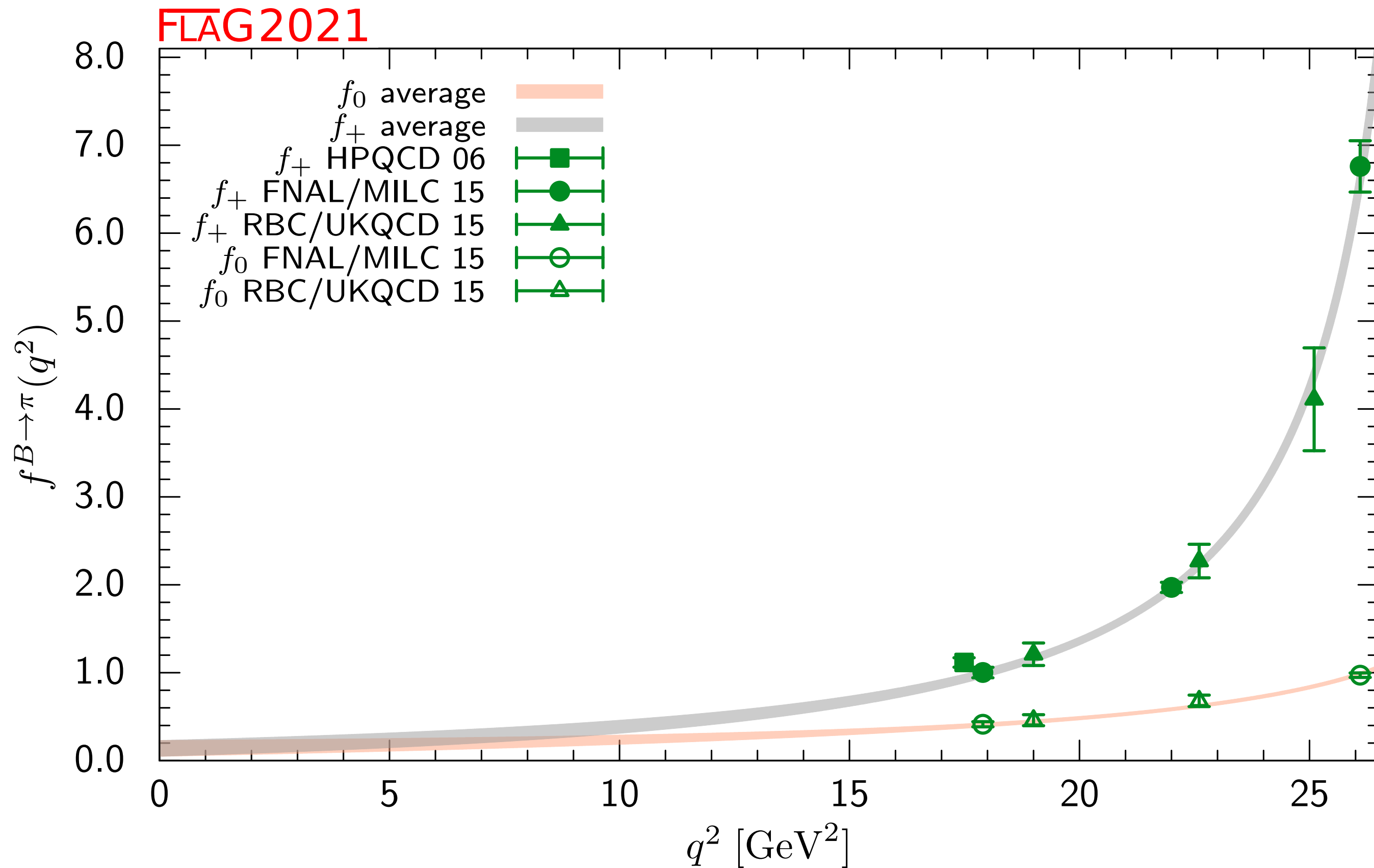
FLAG 2023

$|V_{ub}| \times 10^3$



example: $B \rightarrow \pi \ell \nu$

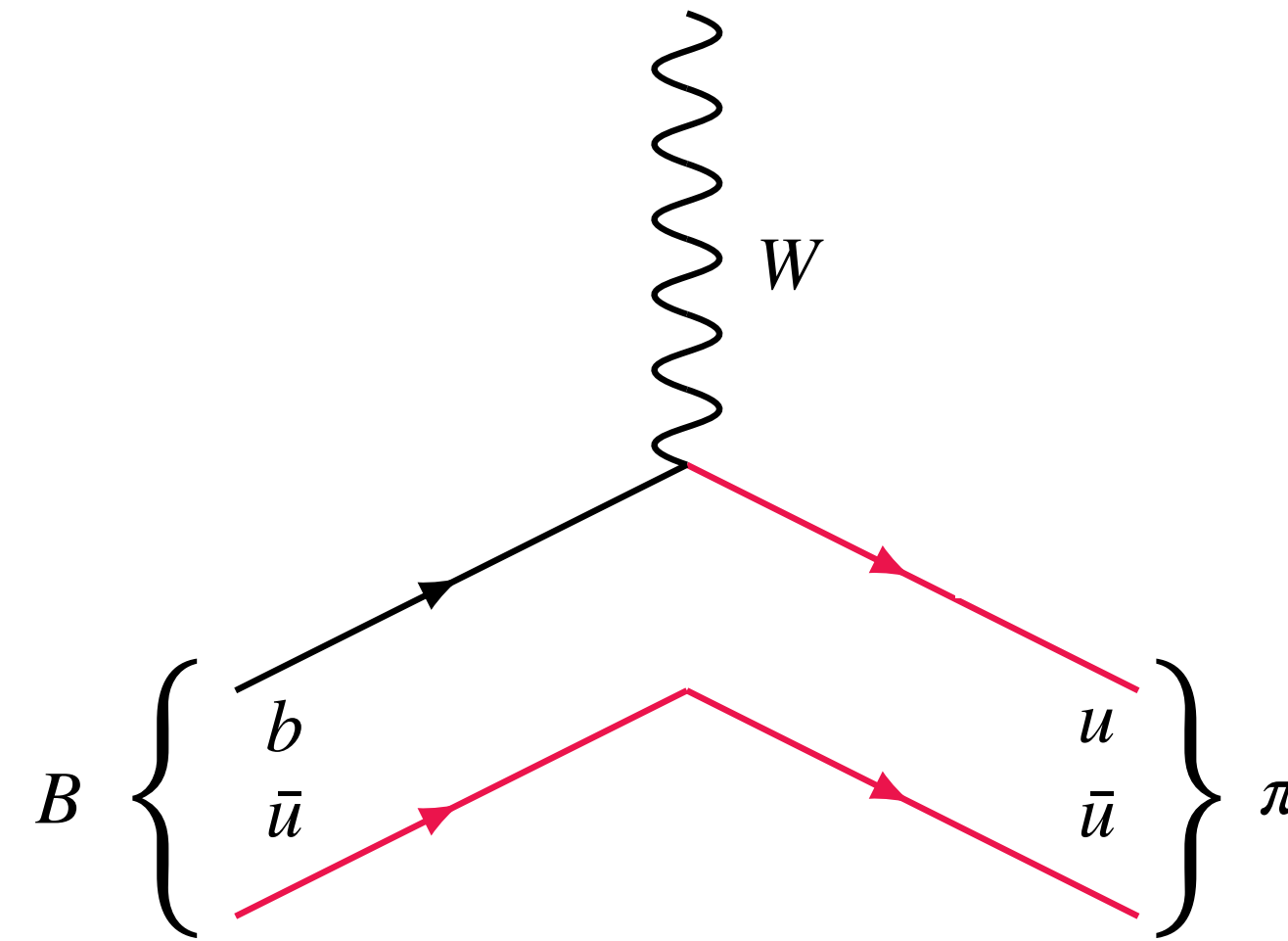
$$\langle \pi | J_{V-A}^\mu | B \rangle = f_+(q^2) \left[p_B^\mu + p_\pi^\mu + \frac{M_B^2 - m_\pi^2}{q^2} q^\mu \right] + f_0(q^2) \frac{M_B^2 - m_\pi^2}{q^2} q^\mu$$



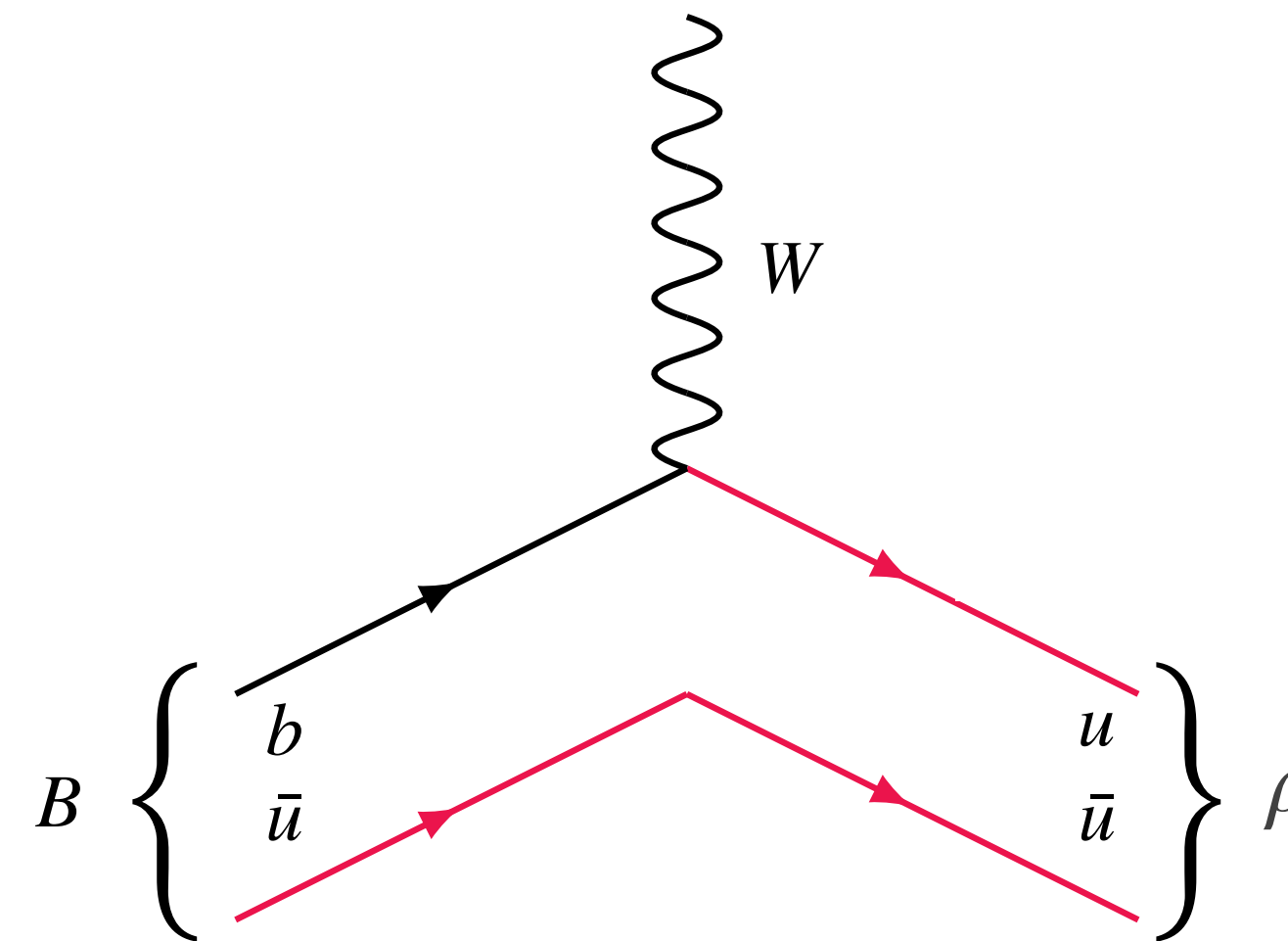
$$\frac{d\mathbf{B}}{dq^2} = \frac{G_F^2 V_{ub}^2}{24\pi^3 q^2} \vec{p}_\pi^3 f_+(q^2)^2$$

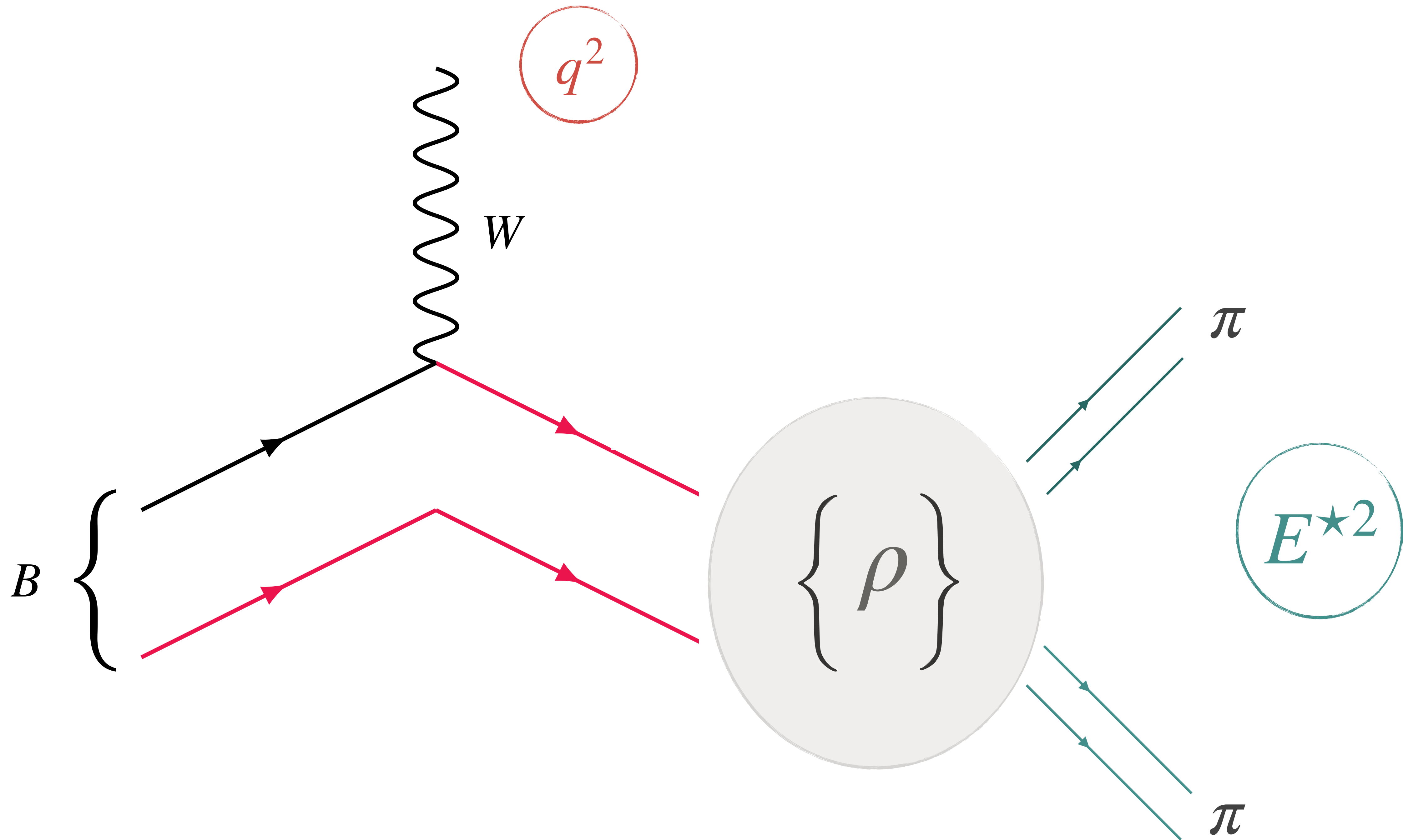
a novel process: $B \rightarrow \rho \ell \bar{\nu}$?

- $B \rightarrow \pi \ell \nu$
 - f_+, f_0
 - established
 - very precise

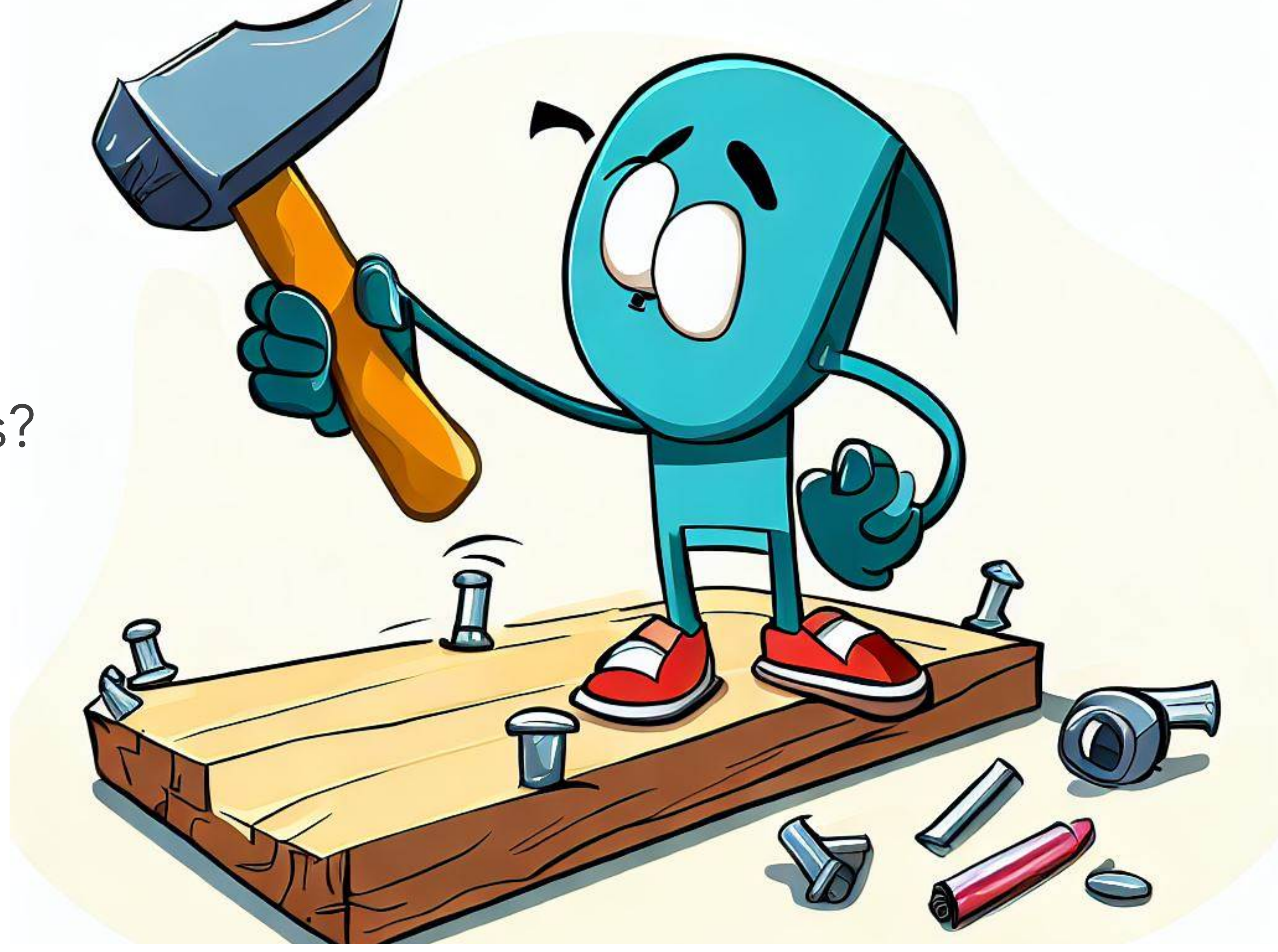


- $B \rightarrow \rho \ell \nu$
 - V, A_0, A_1, A_2
 - new



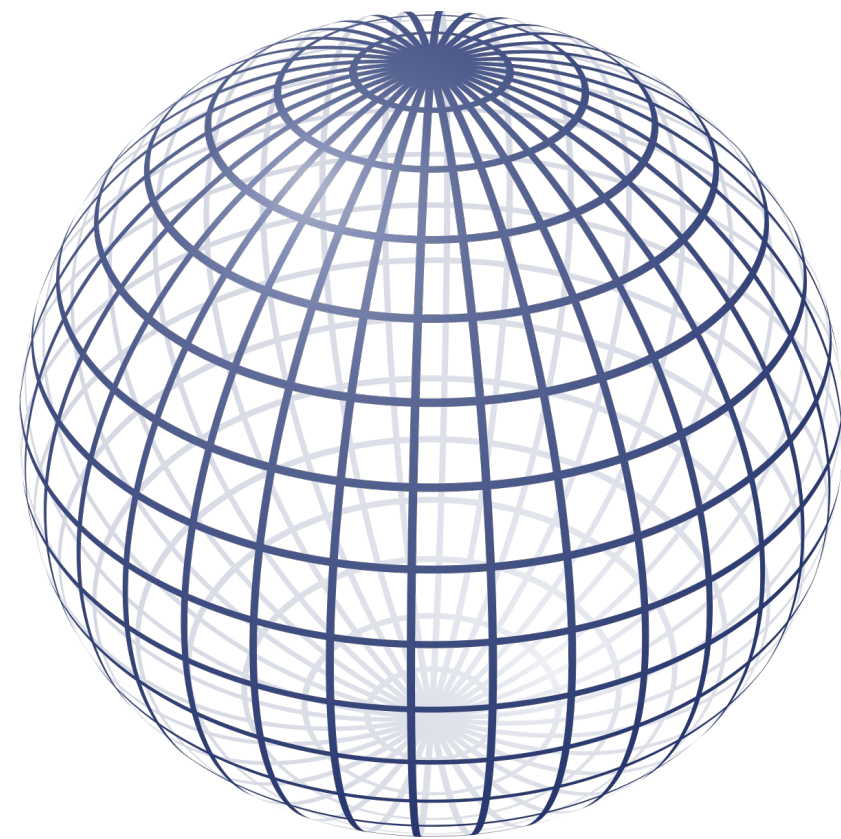


how to
achieve this?

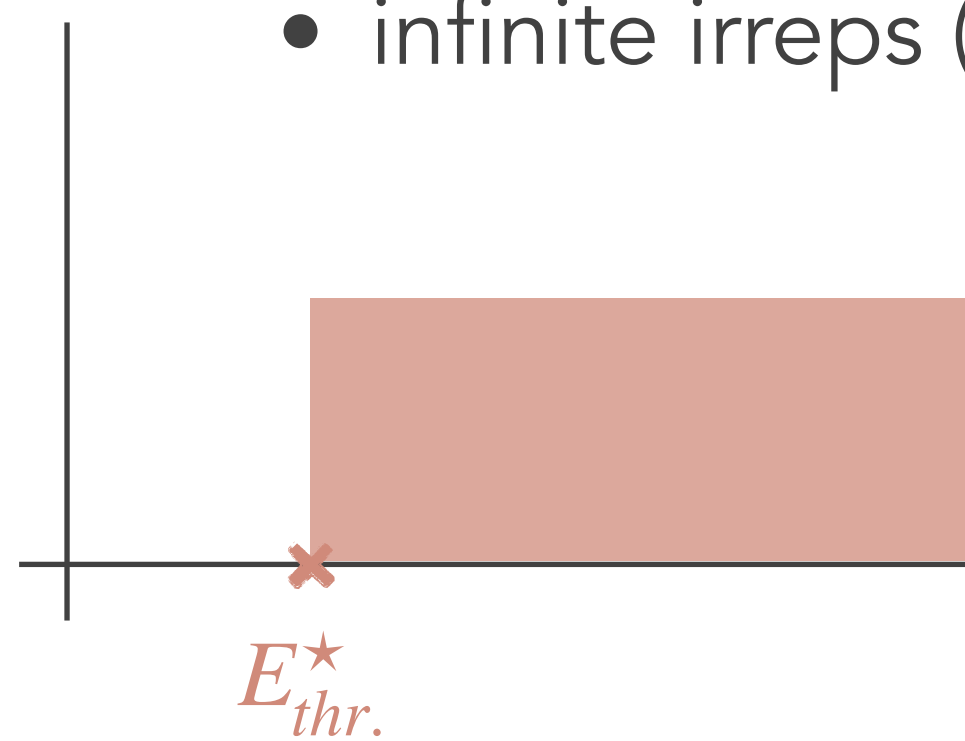


scattering on the lattice

infinite volume:

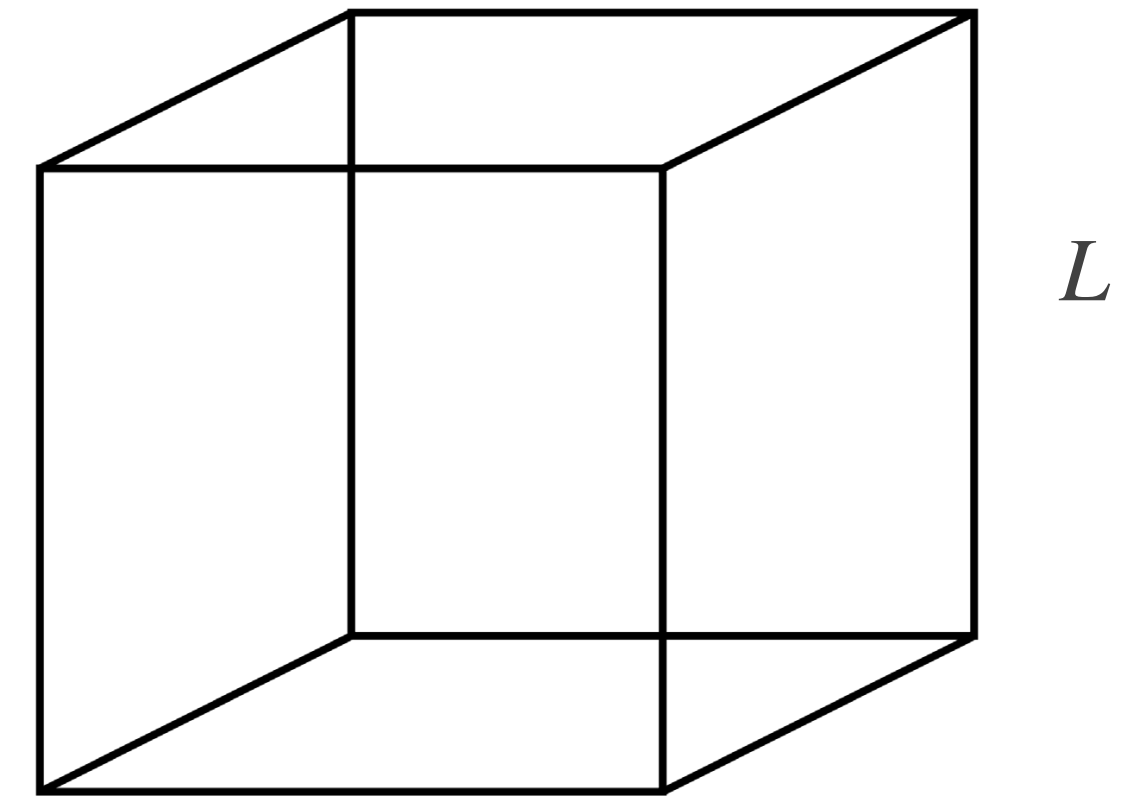


- $O(3)$ symmetry
- infinite irreps (J^P)

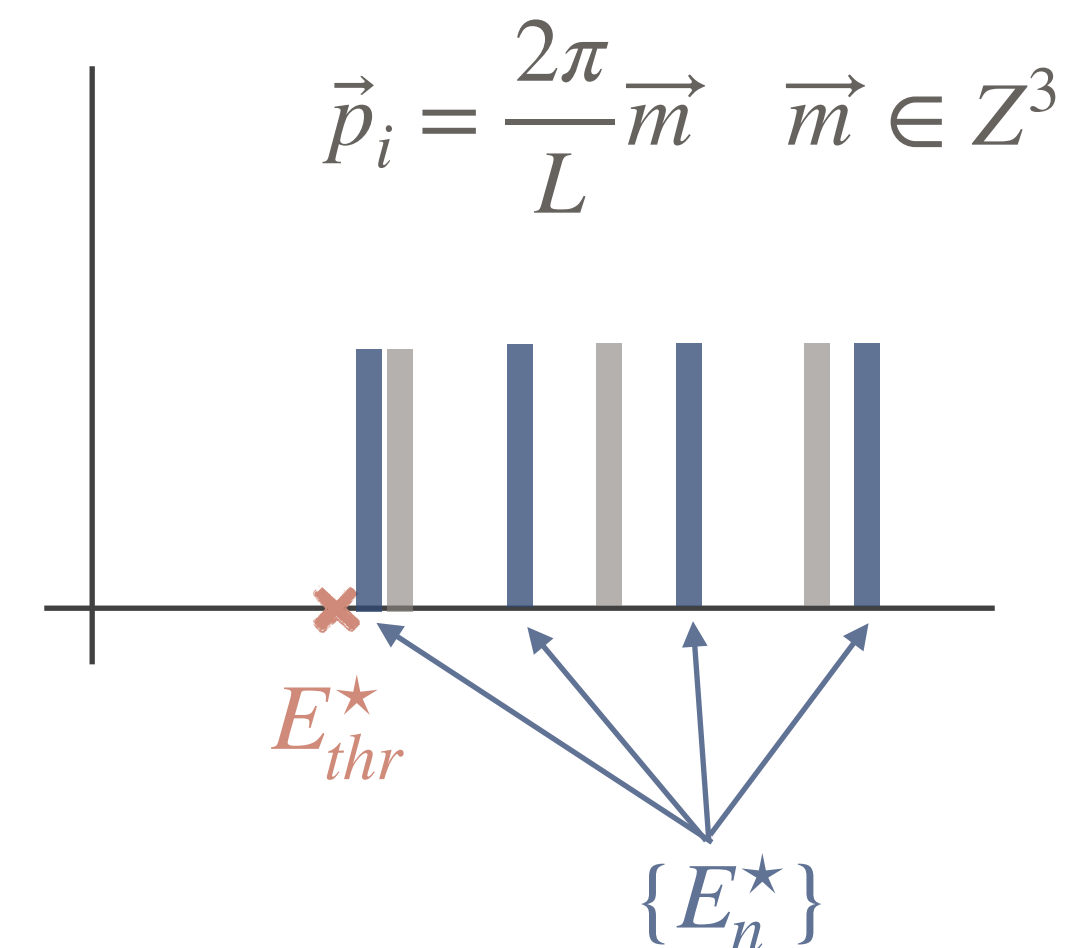


$$\langle p_1 p_2 \rangle = 2E \delta(p_1 - p_2)$$

finite volume:



- discrete symmetries, Λ



$$\langle p_1 p_2 \rangle \neq 2E \delta(p_1 - p_2)$$

many-to-one

lattice states: energies

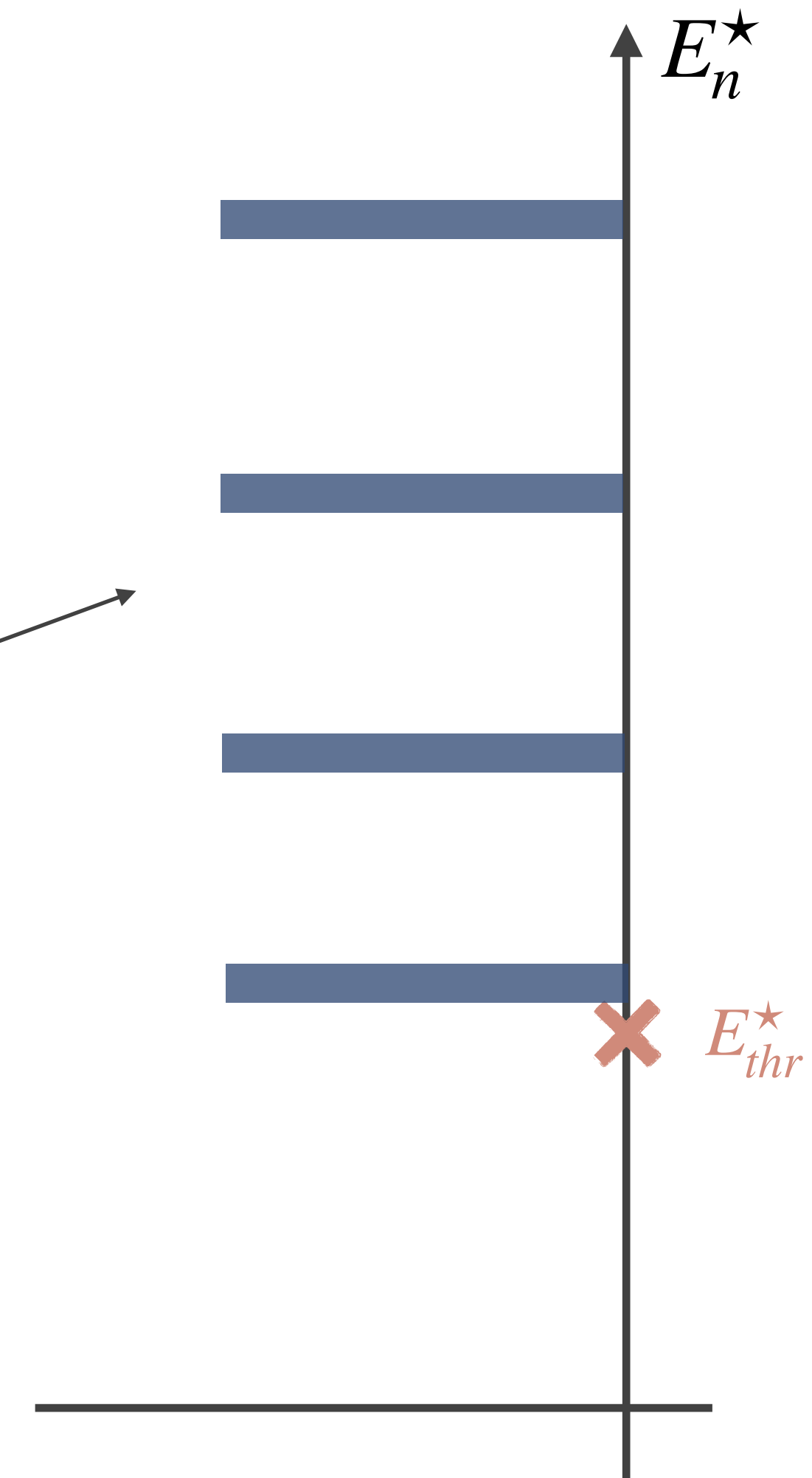
$$C_L^{(2)} = \text{O} \text{---} \text{O} + \text{O} \text{---} \text{●} \text{---} \text{O} + \dots$$

$$C_L^{(2)} = C_\infty^{(2)} - A' \frac{1}{F^{-1}(E^\star) + T(E^\star)} A$$

poles!

discrete spectrum where:

$$\det [F^{-1}(E^\star) + T(E^\star)] = 0$$



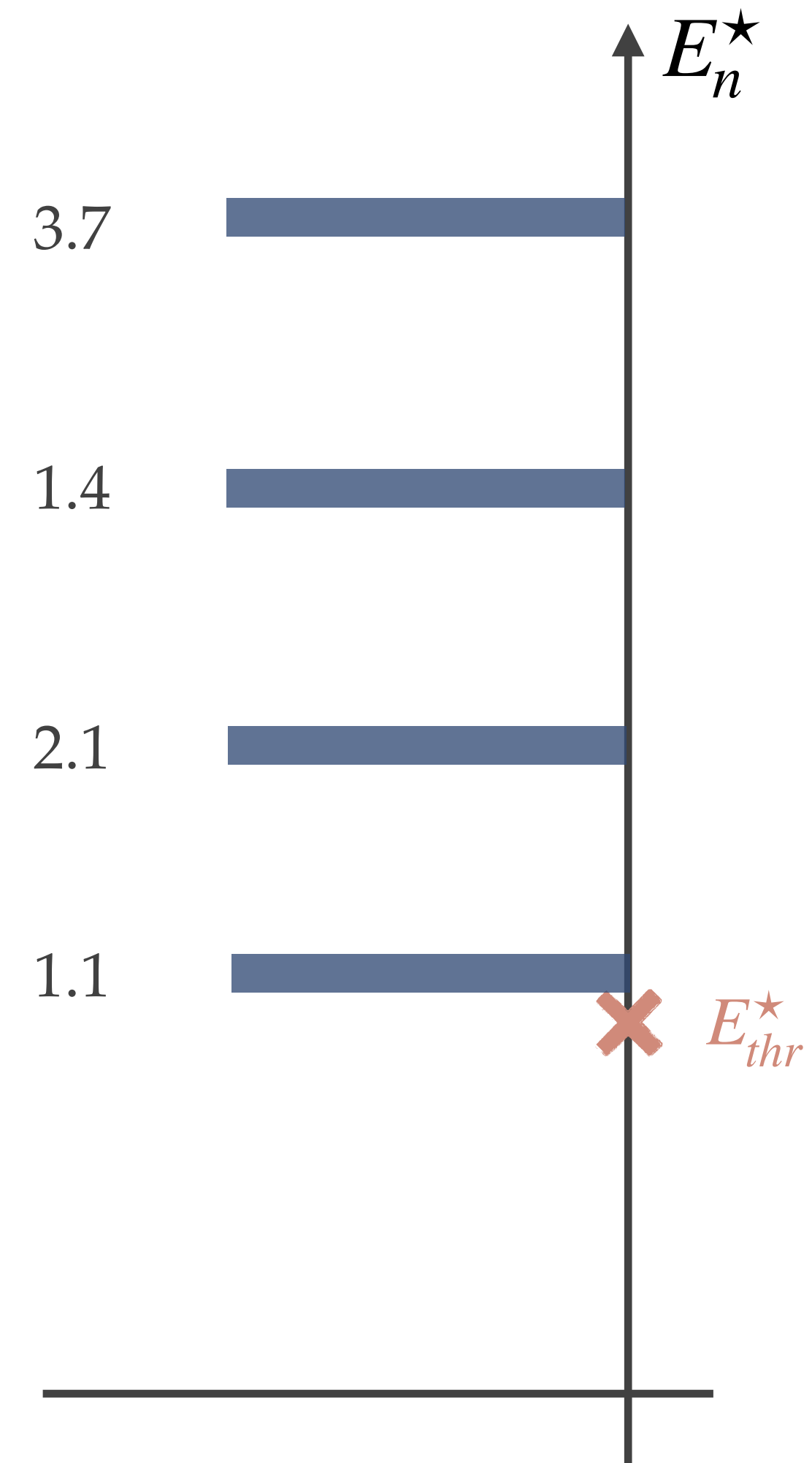
lattice states: normalization

$$C_L^{(3)} = \text{diagram 1} + \text{diagram 2} + \dots$$

$$C_L^{(3)} = C_\infty^{(3)} - A'RA$$

$$R = \lim_{E \rightarrow E_n} \frac{E - E_n}{F^{-1} + T}$$

residue of pole!

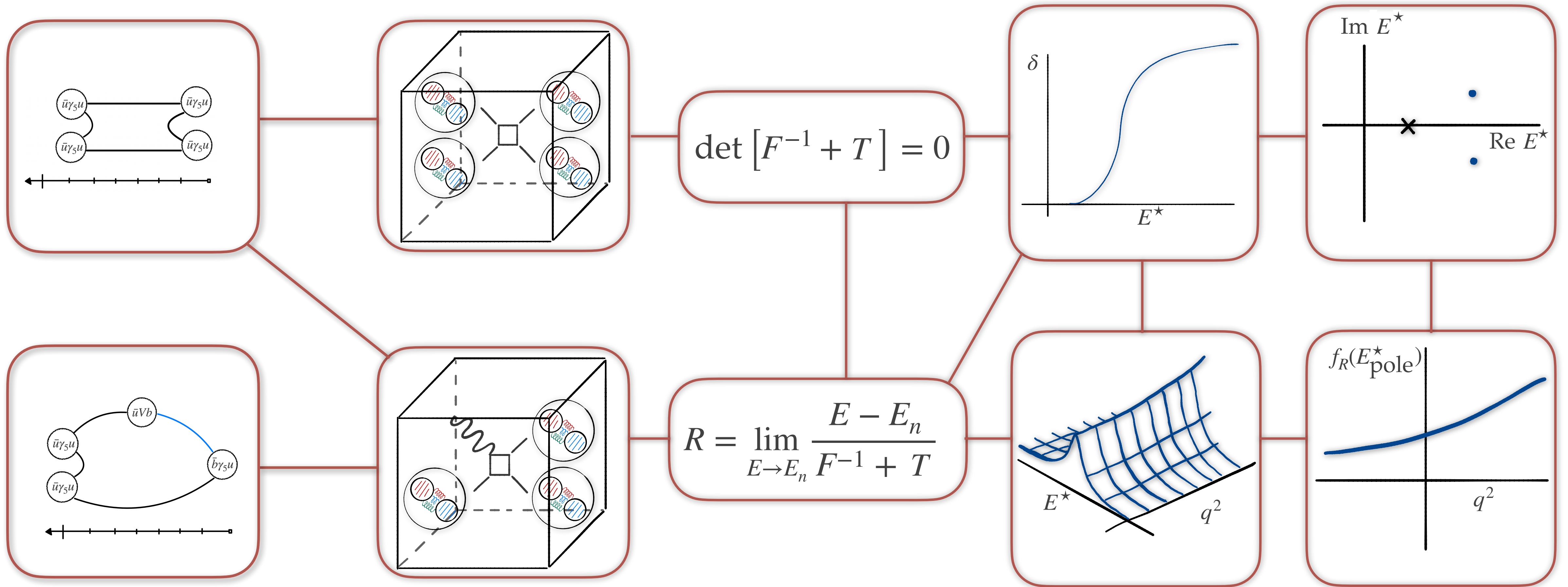


normalization of finite-volume states

$$\langle E_n^{\Lambda^*} \rangle_L \sim \sqrt{R} p_1 p_2 \langle E^* = E_n^{\Lambda^*} \rangle_\infty$$

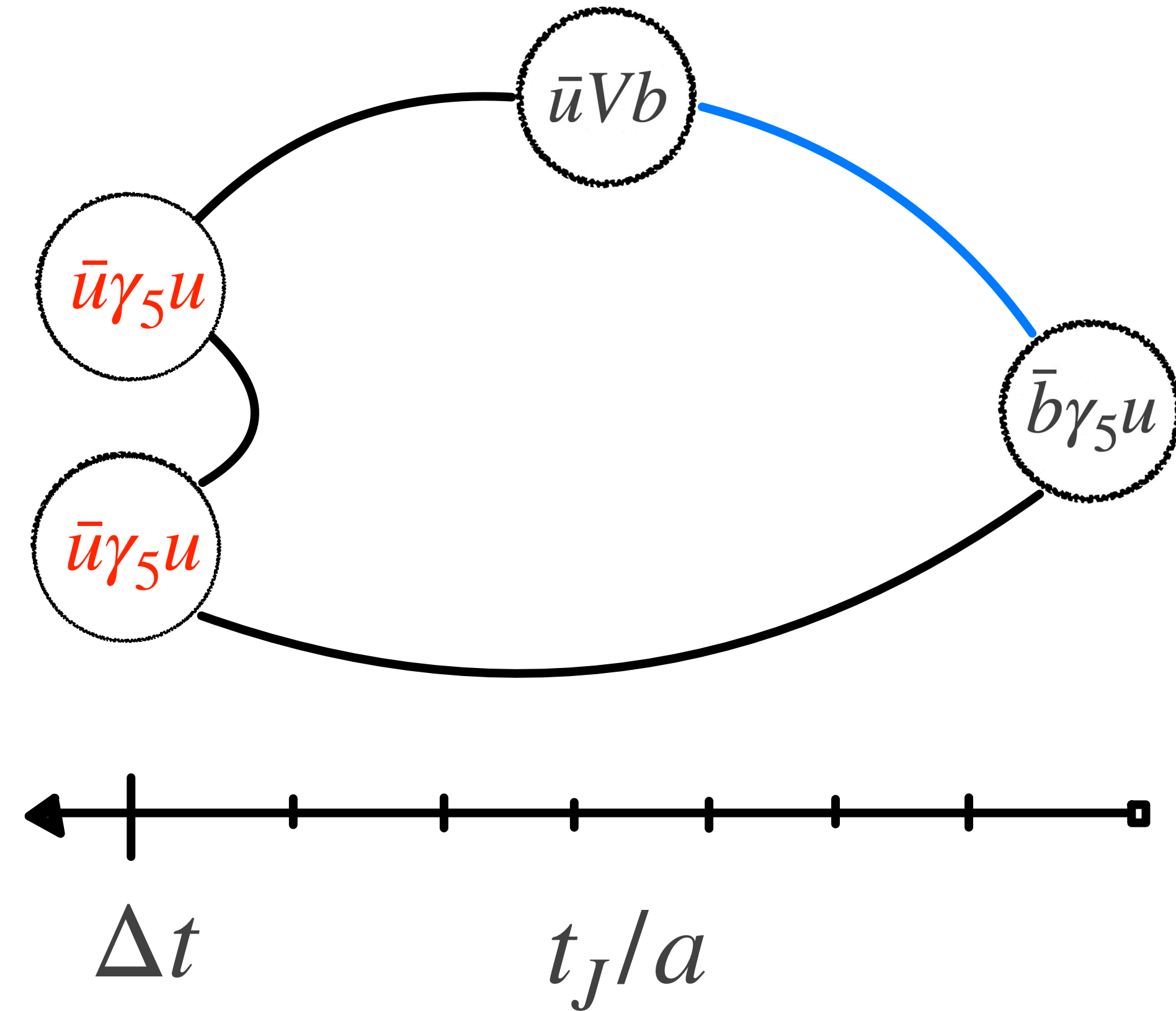
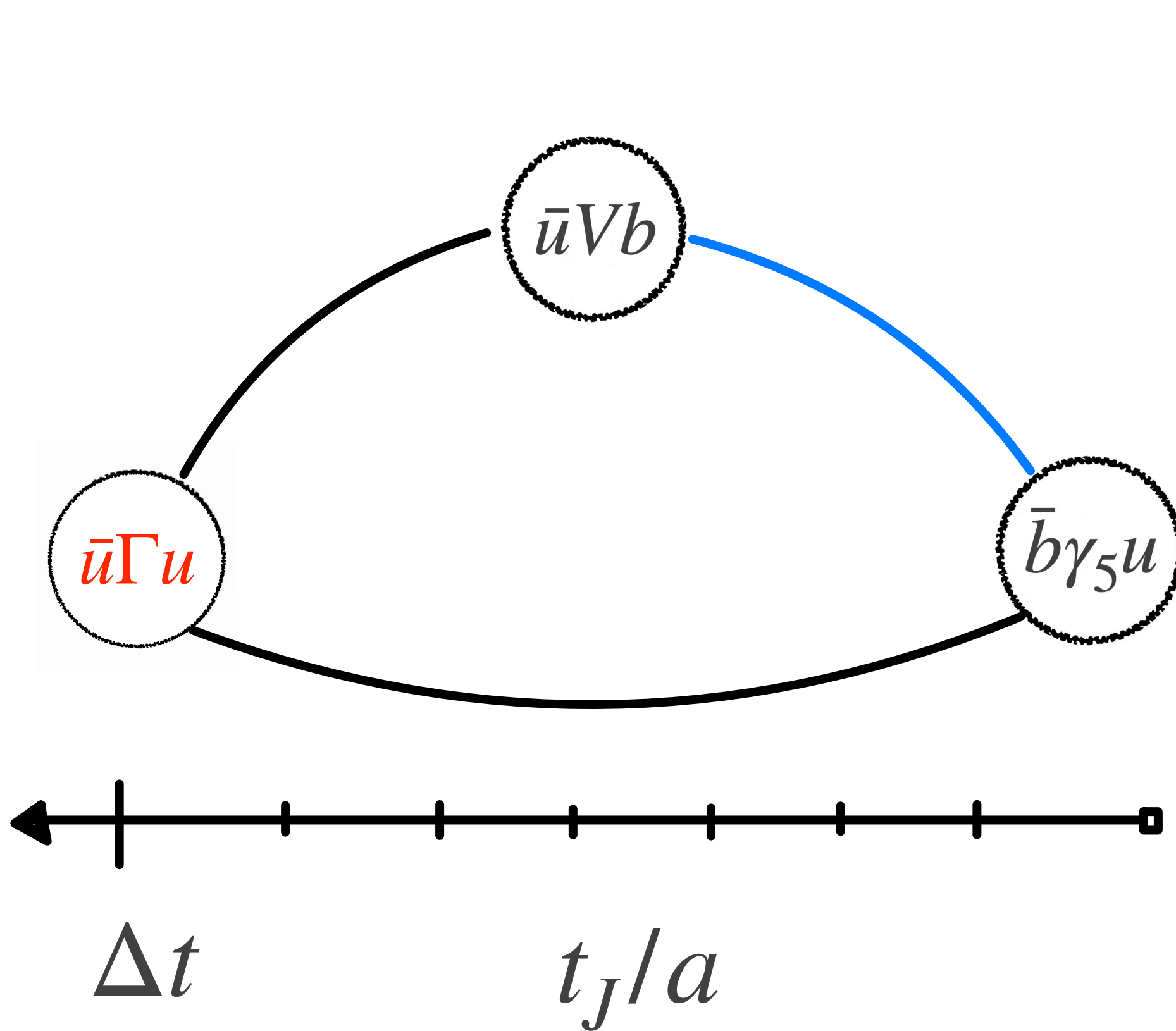
the "Lellouch-Lüscher" factor

a cartoon scheme of $B \rightarrow \pi\pi\ell\bar{\nu}$



3-point functions

$$C_{3,i} = \langle O_i(\vec{p}, \Lambda) V^\mu O_B(\vec{p}_B) \rangle$$



3-point function lattice data

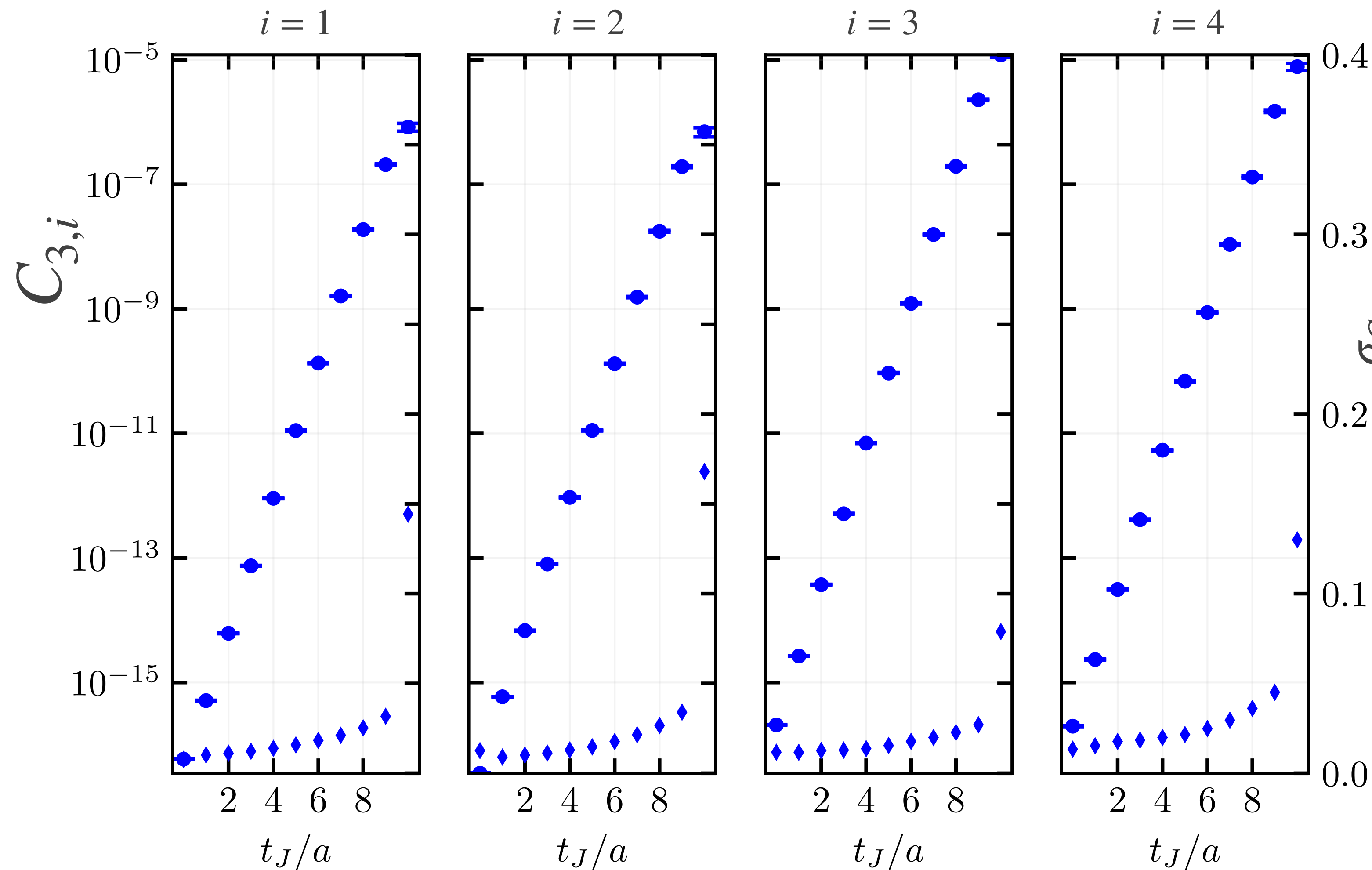
$$\vec{p} = \frac{2\pi}{L}[1,0,1]$$

$$\Lambda = B_2$$

$$\mu = z$$

$$\vec{p}_B = \frac{2\pi}{L}[0,0,1]$$

$$C_{3,i} = \langle O_i(\vec{p}, \Lambda) V^\mu O_B(\vec{p}_B) \rangle$$



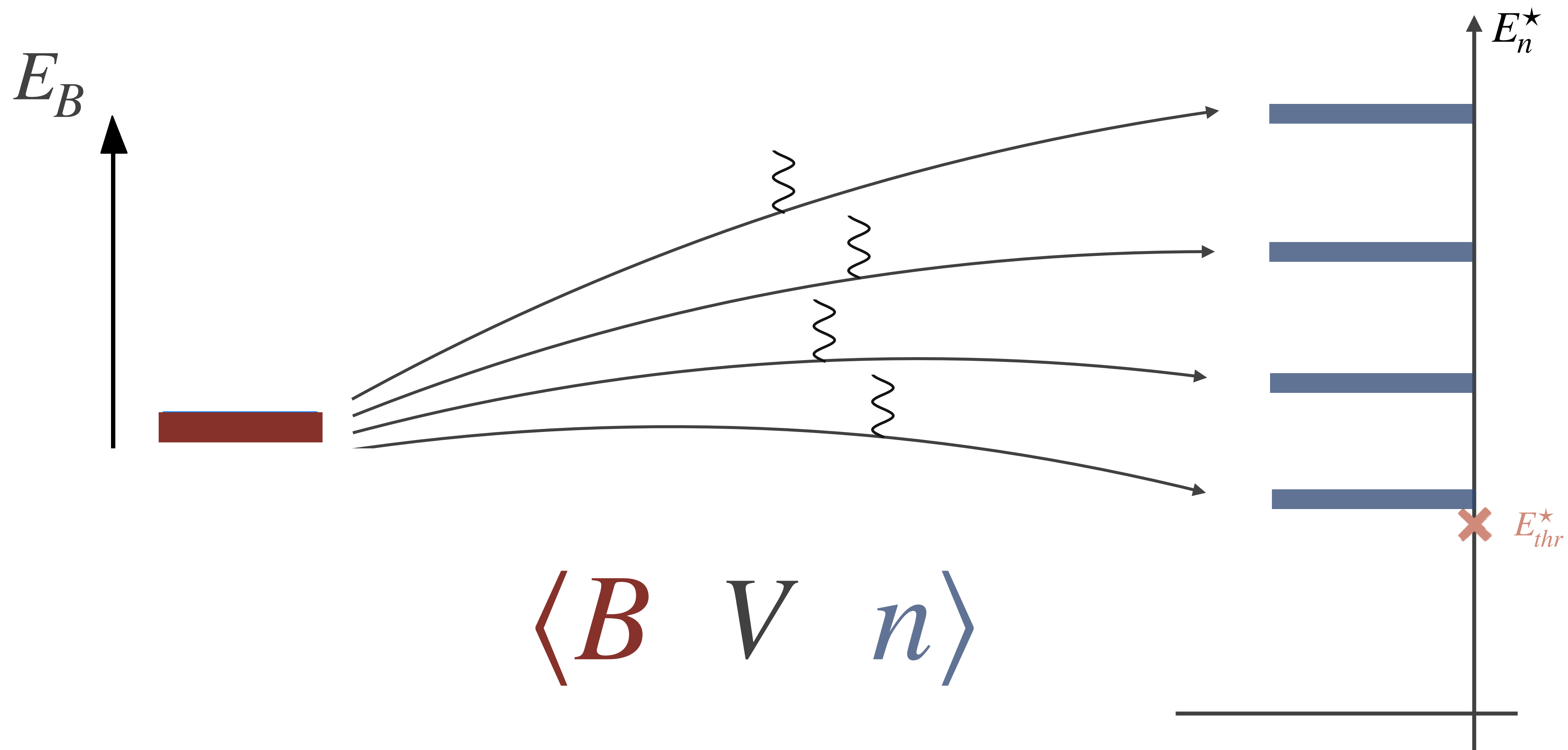
$$O_1 = \bar{u}\Gamma_{B_1}u$$

$$O_2 = \bar{u}\gamma_t\Gamma_{B_1}u$$

$$O_3 = \pi(\vec{p}_1)\pi(\vec{p}_2)$$

$$O_4 = \pi(\vec{p}_1)\pi(\vec{p}_2)$$

matrix elements



state projection

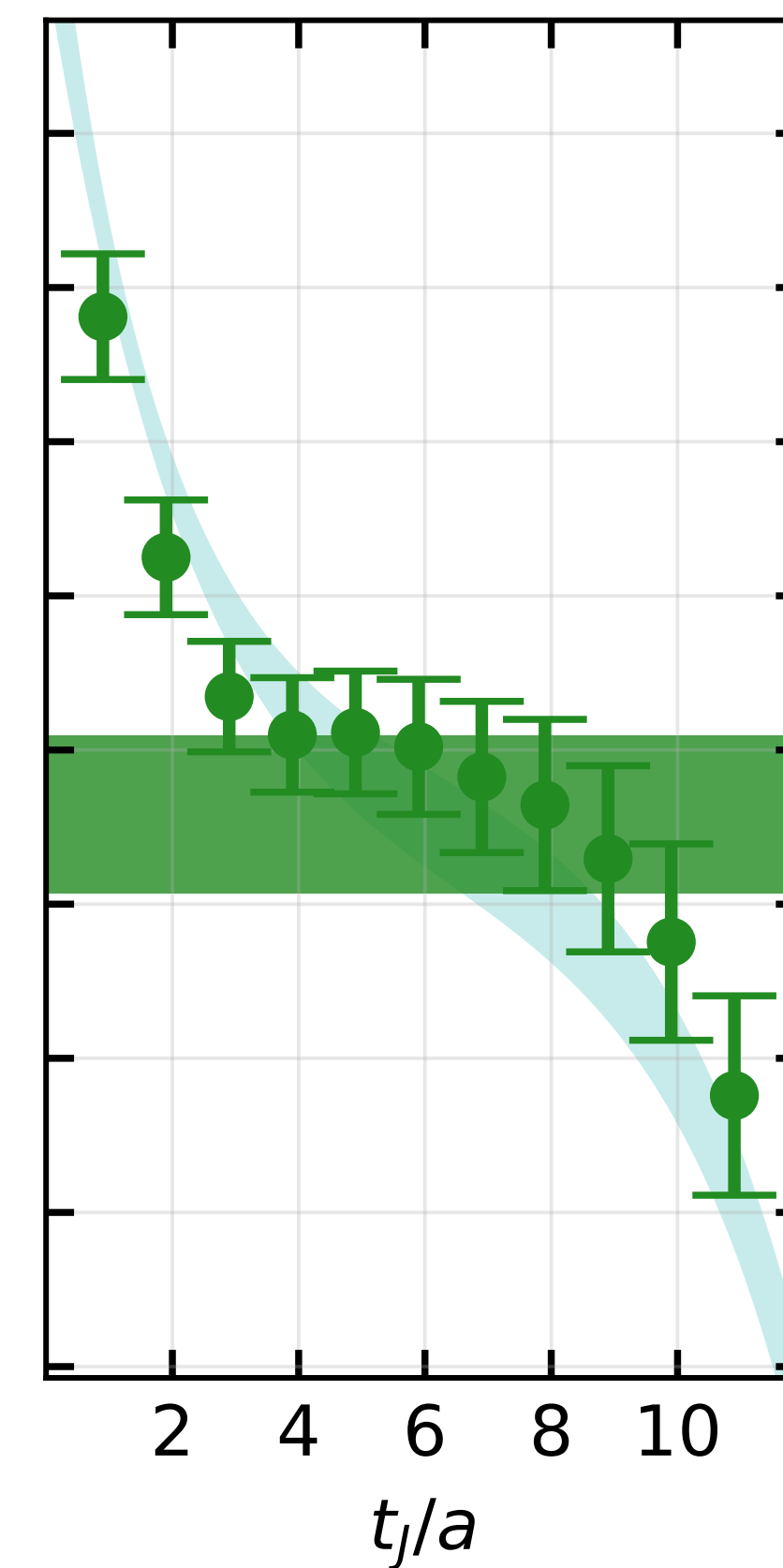
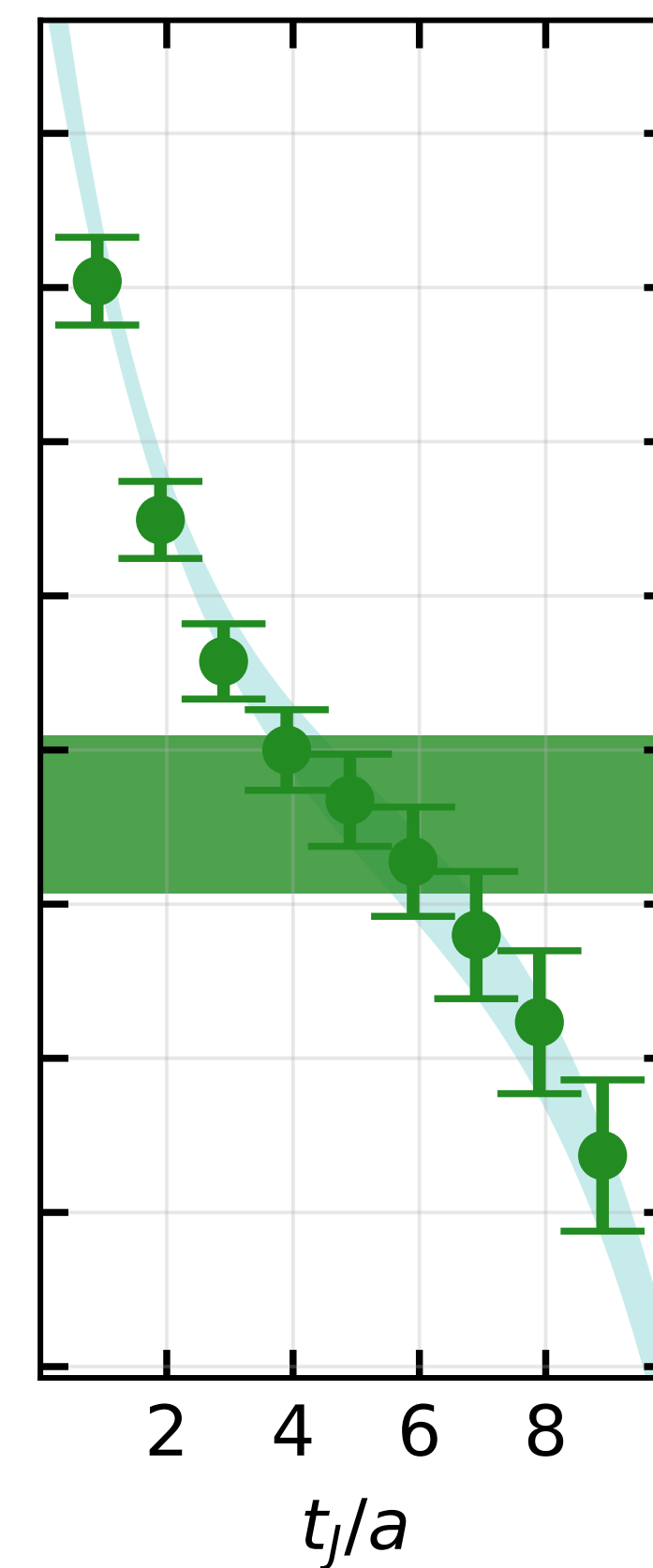
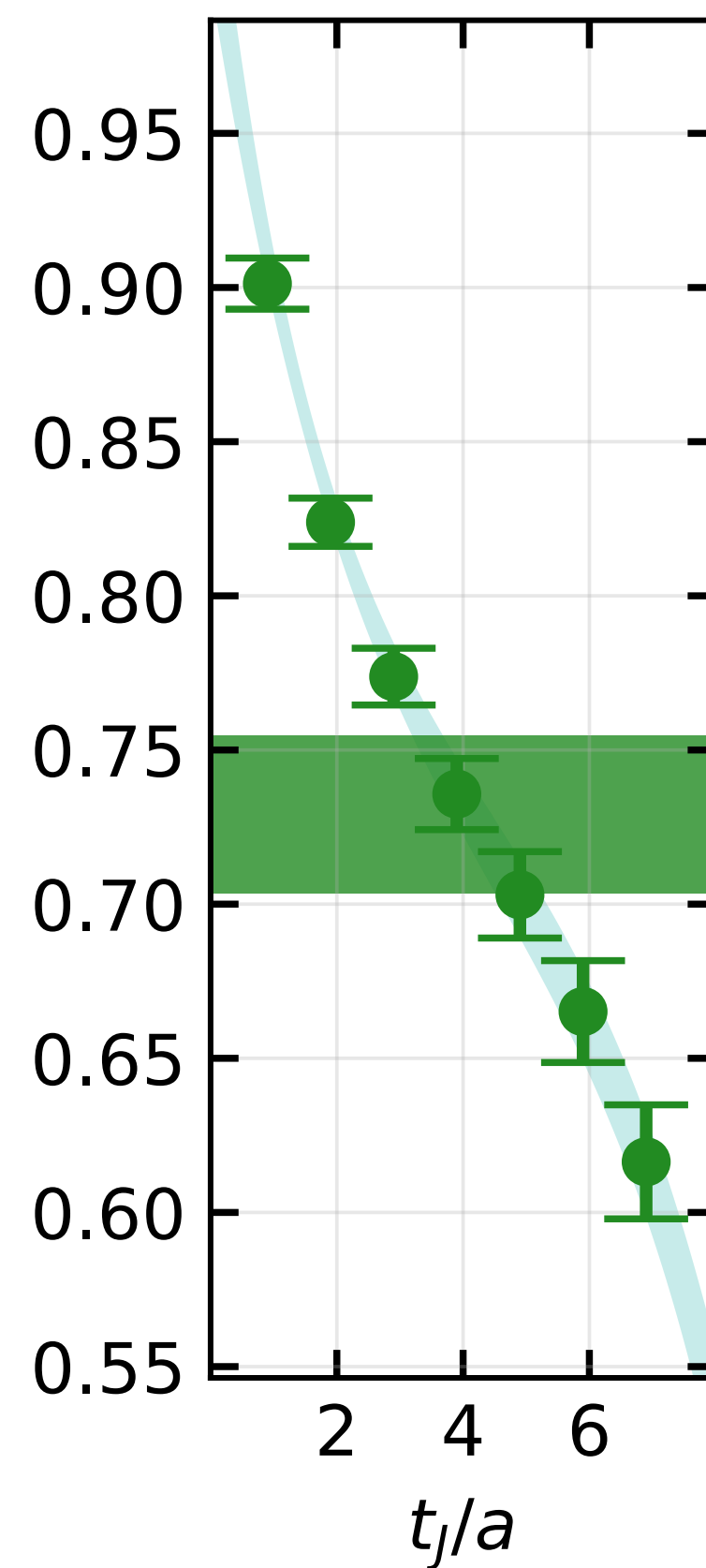
$$C_{3,i} = \sum_{n \in [\pi\pi]} Z_i^n \langle B | V | n \rangle Z_B \frac{e^{-E_n(\Delta t - t)} e^{-E_B t}}{2E_n 2E_B}$$

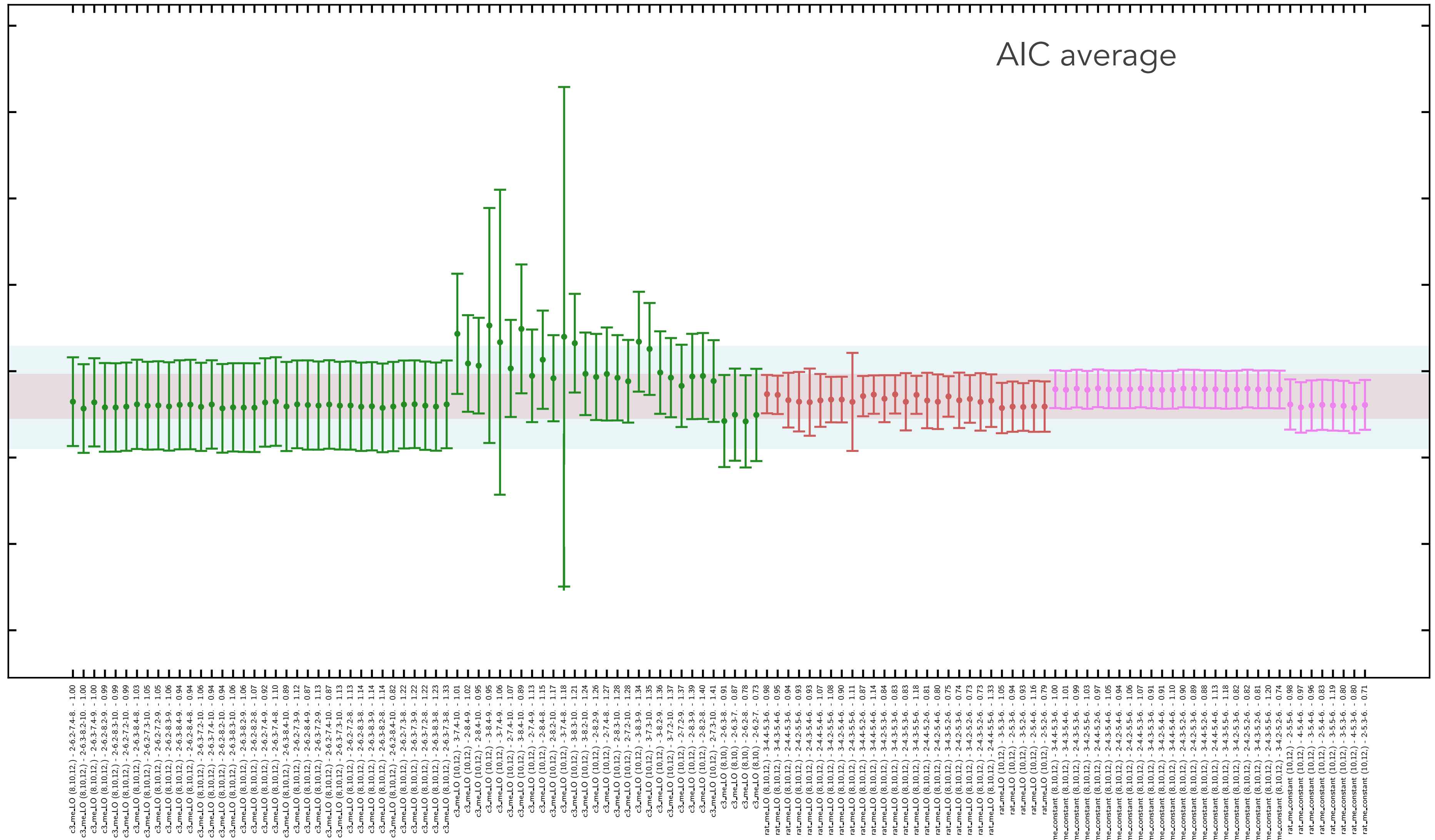
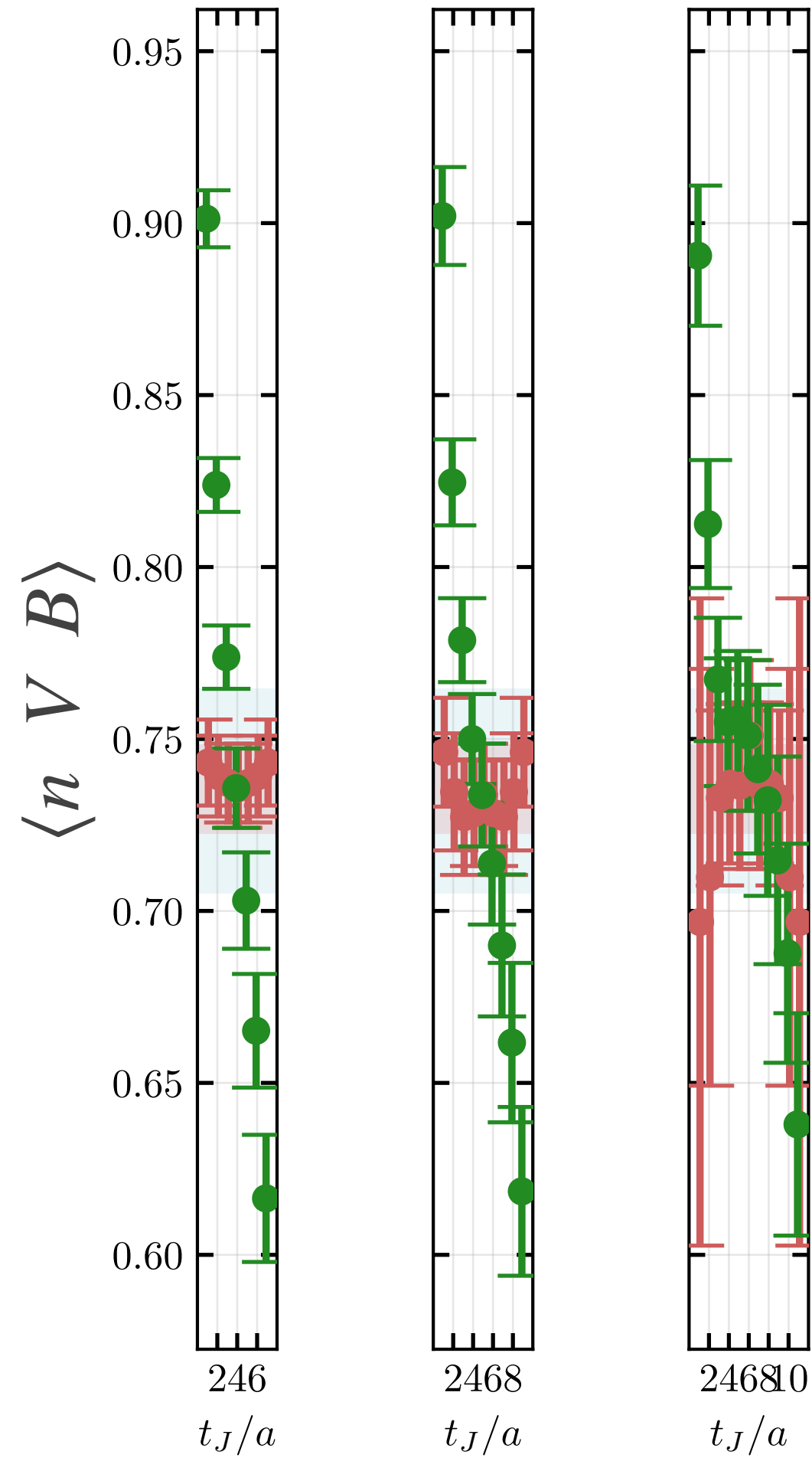
$$C_3^n = u_i^n C_{3,i}$$

because in GEVP

$$u_i^{n\star} Z_i^m = 2E_n e^{E_n t_0} \delta^{nm}$$

$\langle n | V | B \rangle$





transition amplitude

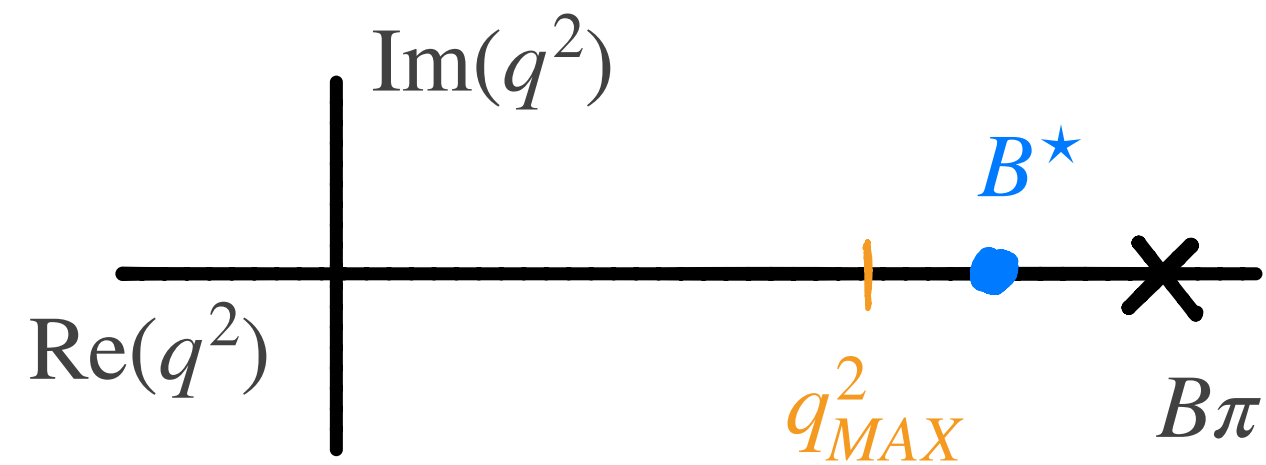
Boyd, Grinstein, Lebed [hep-ph/9412324](https://arxiv.org/abs/hep-ph/9412324)
 Bourrely, Caprini, Lellouch [0807.2722](https://arxiv.org/abs/hep-ph/0807.2722)
 Alexandrou, LL, Meinel et al. [1807.08357](https://arxiv.org/abs/1807.08357)

$$\langle \pi\pi, E^* \mid V \mid B, p_B \rangle_\infty = \frac{2iV(E^*, q^2)}{m_B + 2m_\pi} \epsilon_{\mu\nu\alpha\beta} \epsilon^{\nu*} p^\alpha p_B^\beta$$

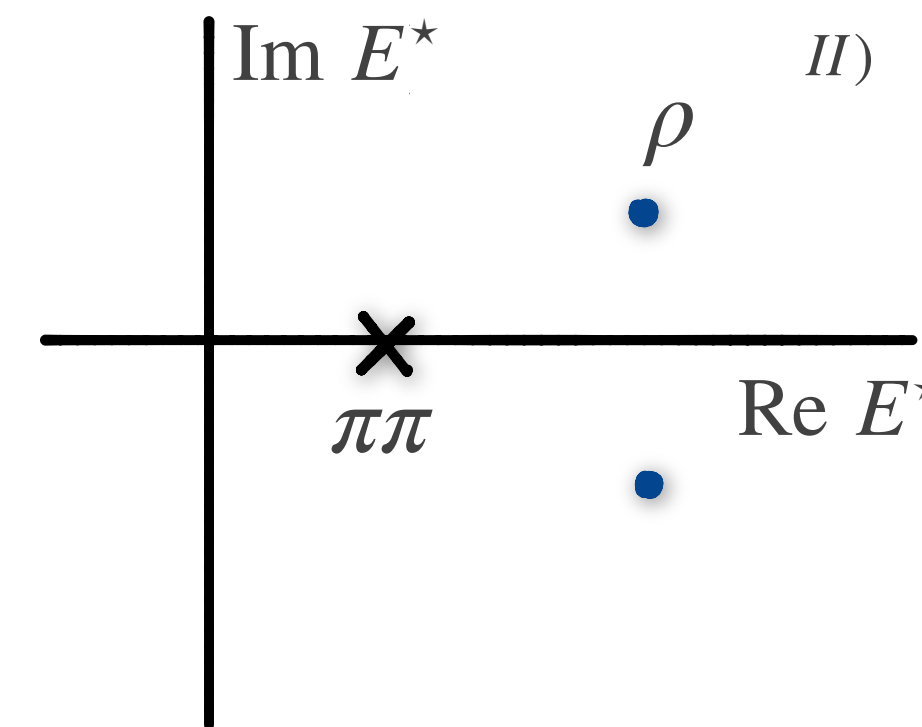
$$q = p_f - p_i$$

$$E^* = 2\sqrt{m_\pi^2 + k^2}$$

$$V(E^*, q^2) = F(E^*, q^2) \frac{T(E^*)}{k}$$



$$F(E^*, q^2) = \frac{1}{1 - \frac{q^2}{m_p^2}} \sum_{n,m} A_{n,m} z^n(q^2) (E^{*2} - E_{thr}^2)^m$$



$$T = \frac{E^* \Gamma_i}{m_R^2 - E^{*2} - iE^* \Gamma_i}$$

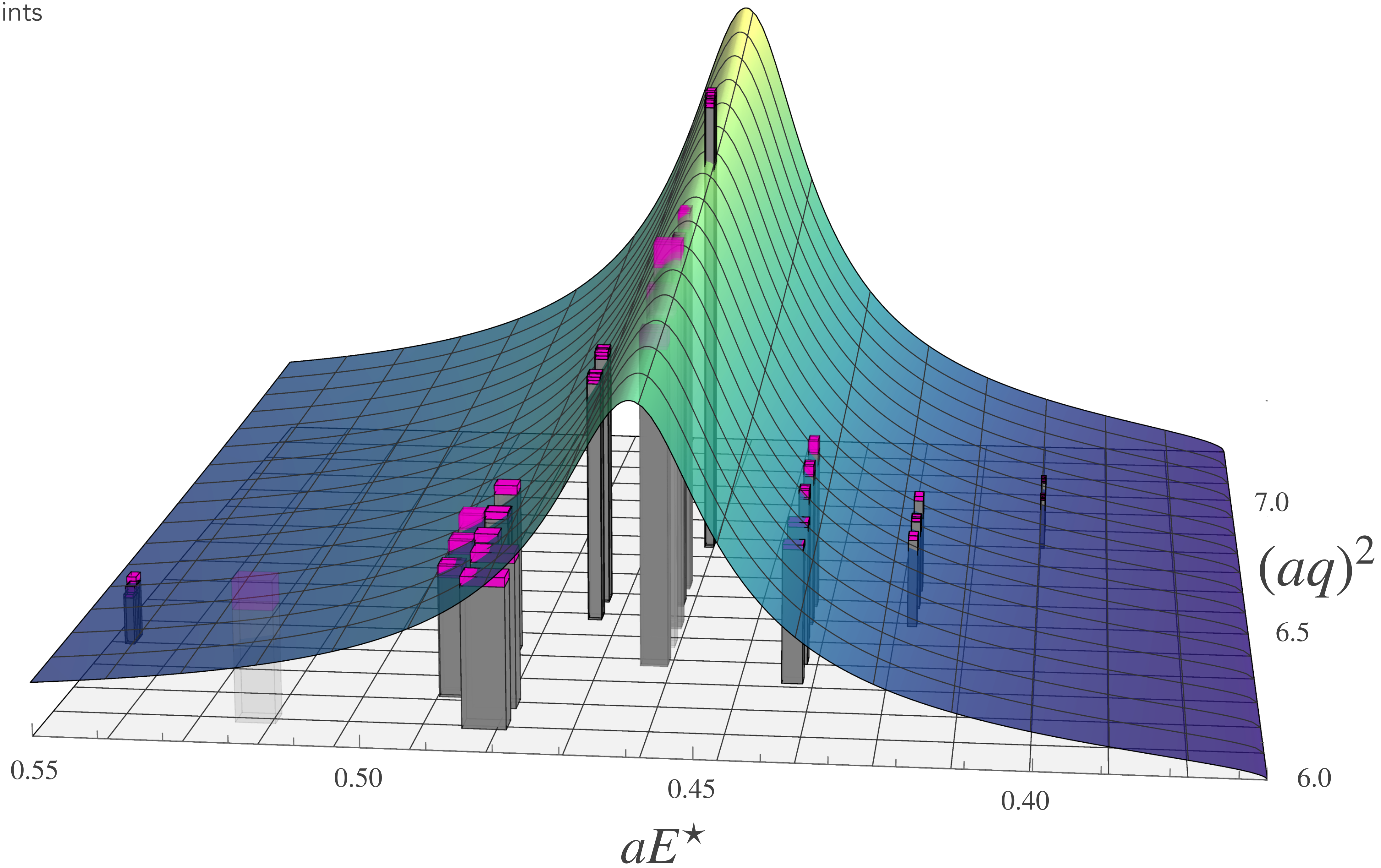
$$\Gamma_I = \frac{g_{\rho\pi\pi}^2 p^3}{6\pi E^{*2}}$$

$$\Gamma_{II} = \frac{g_{\rho\pi\pi}^2 p^3}{6\pi E^{*2}} \frac{1 + (k_R r_0)^2}{1 + (k r_0)^2}$$

$$\langle B \mid V \mid n \rangle_L = \sqrt{R_n} \langle B, p_B \mid V \mid \pi\pi, E^* \rangle_\infty$$

"Lellouch-Lüscher" factor

64 data points

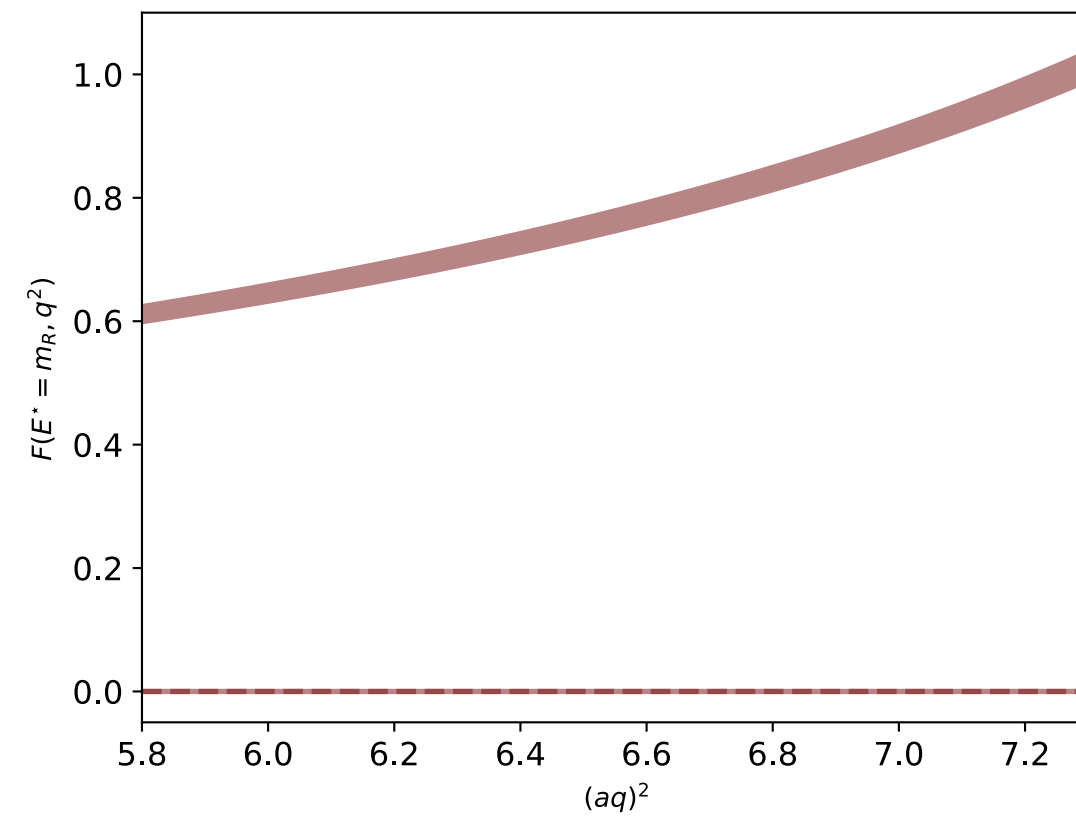


details I

❖ FF3N0M0_TBWI

❖ $\chi^2/\text{dof} = 74.3/63 = 1.179$

❖ $A_{0,0} = 0.2413(66)$

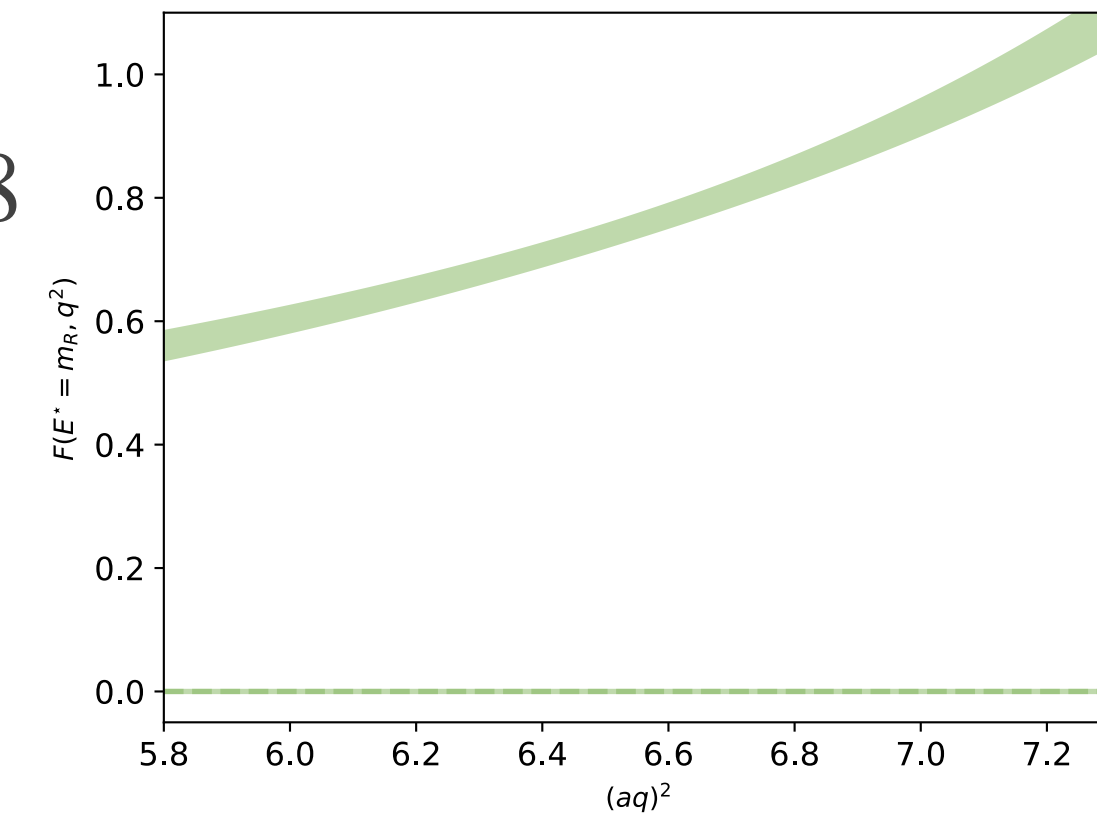


❖ FF3N1M0_TBWI

❖ $\chi^2/\text{dof} = 66.8/62 = 1.078$

❖ $A_{0,0} = 0.2256(87)$

❖ $A_{1,0} = -0.41(18)$

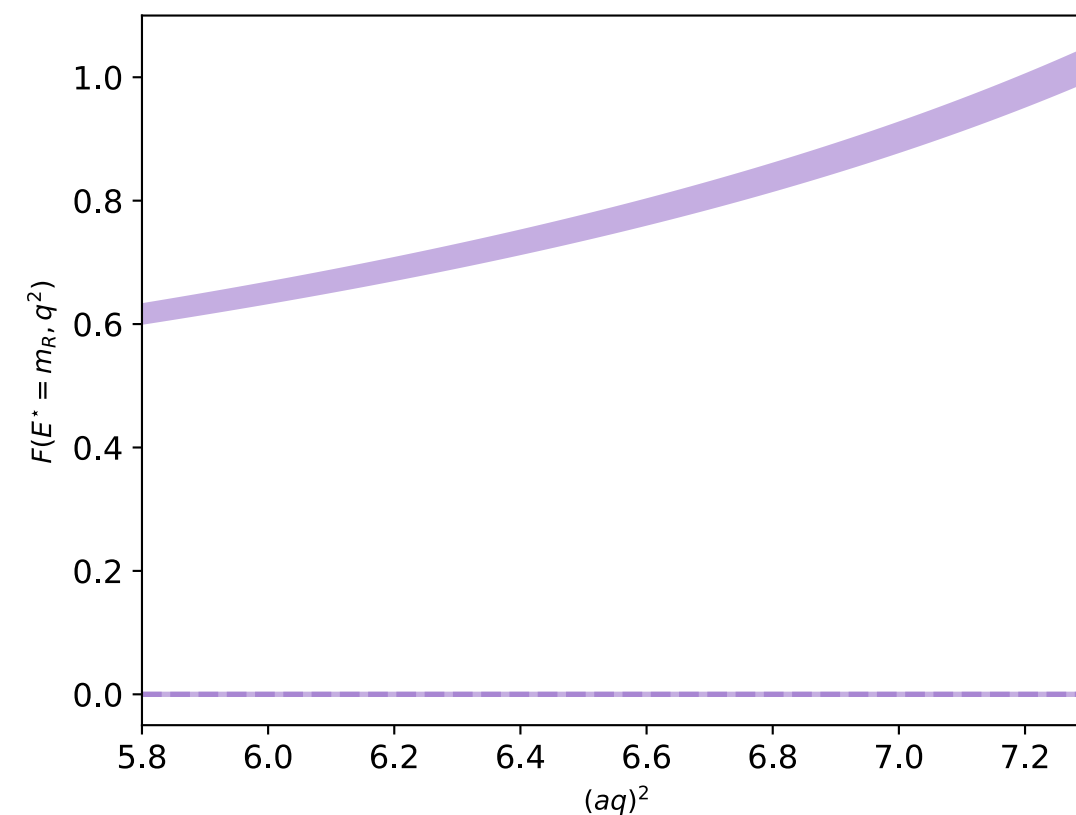


❖ FF3N0M1_TBWI

❖ $\chi^2/\text{dof} = 39.1/62 = 0.630$

❖ $A_{0,0} = 0.279(12)$

❖ $A_{0,1} = -0.061(13)$



❖ FF3N1M1_TBWI

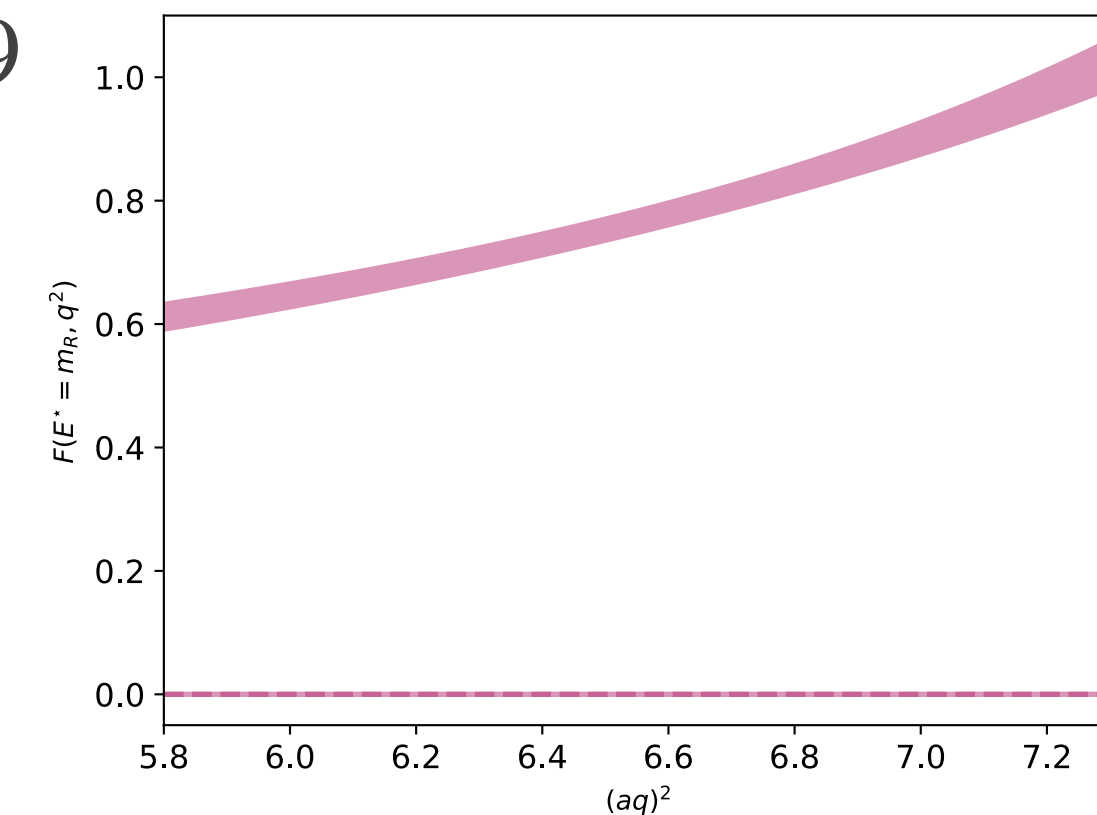
❖ $\chi^2/\text{dof} = 29.9/60 = 0.499$

❖ $A_{0,0} = 0.260(16)$

❖ $A_{1,0} = -0.88(50)$

❖ $A_{0,1} = -0.031(21)$

❖ $A_{1,1} = 1.46(82)$

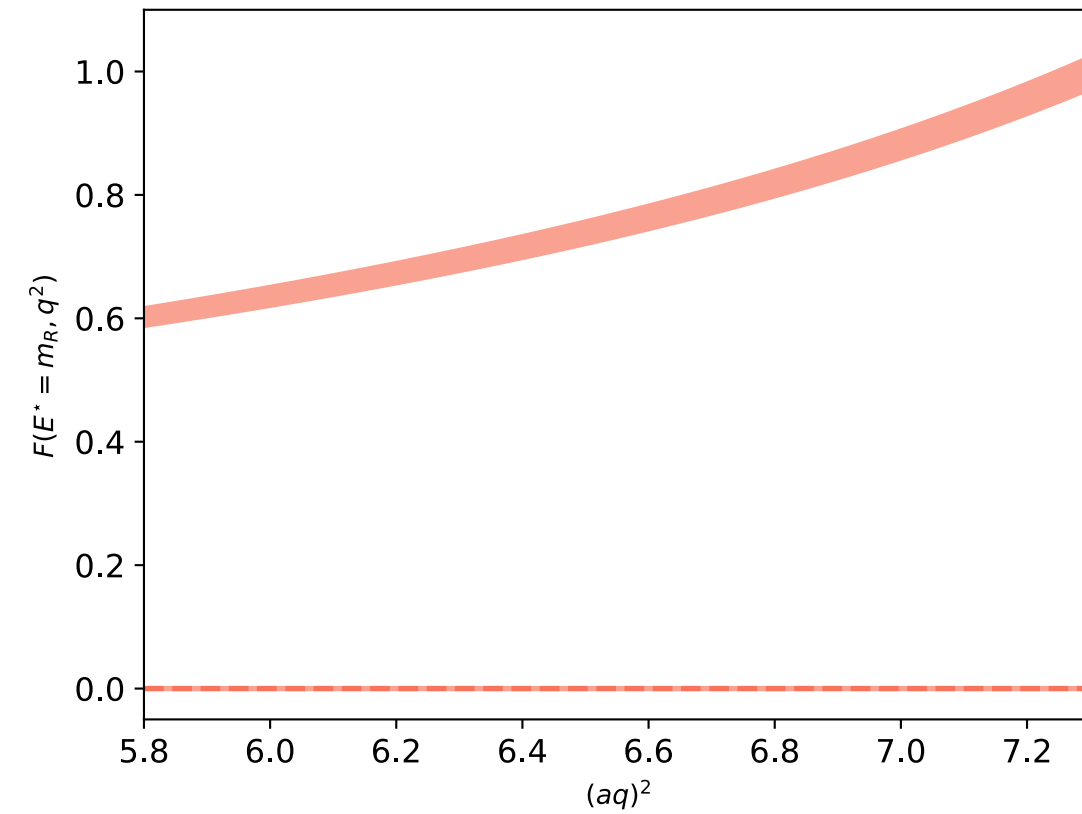


details II

❖ FF3N0M0_TBWII

❖ $\chi^2/\text{dof} = 34.3 / 63 = 0.545$

❖ $A_{0,0} = 0.2376(71)$

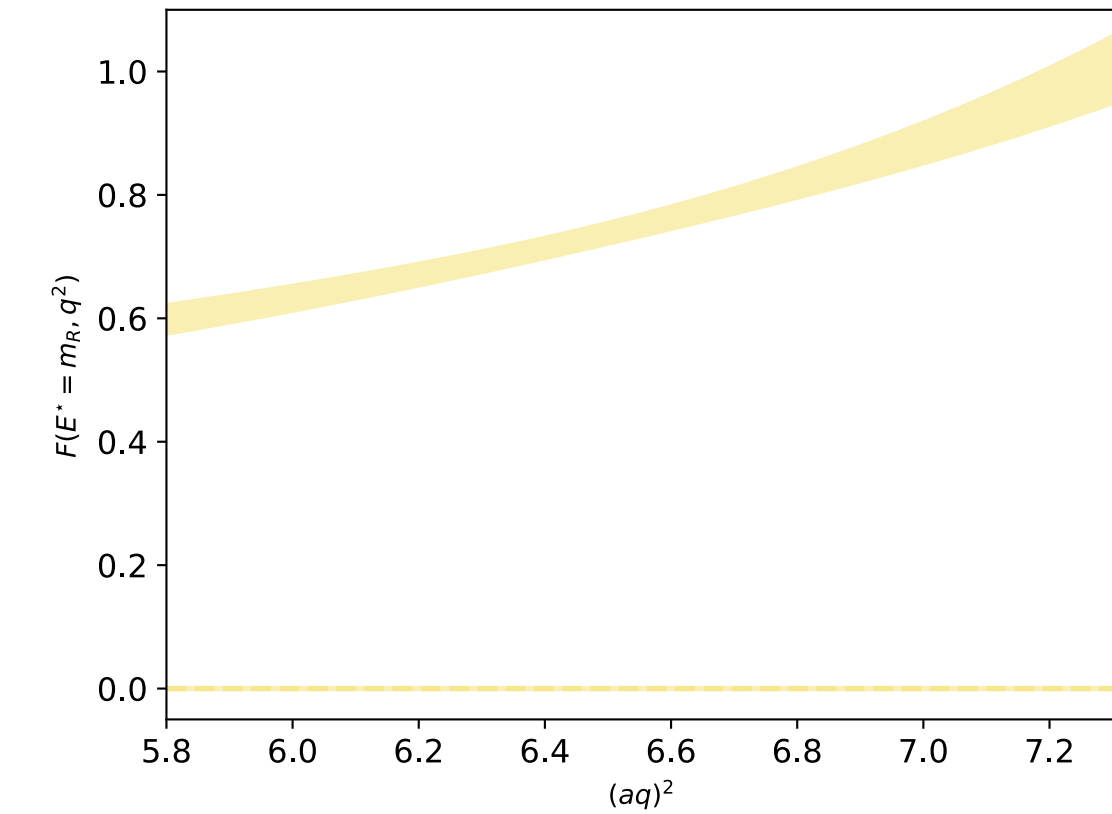


❖ FF3N1M0_TBWII

❖ $\chi^2/\text{dof} = 34.0 / 62 = 0.549$

❖ $A_{0,0} = 0.2365(89)$

❖ $A_{1,0} = -0.03(22)$

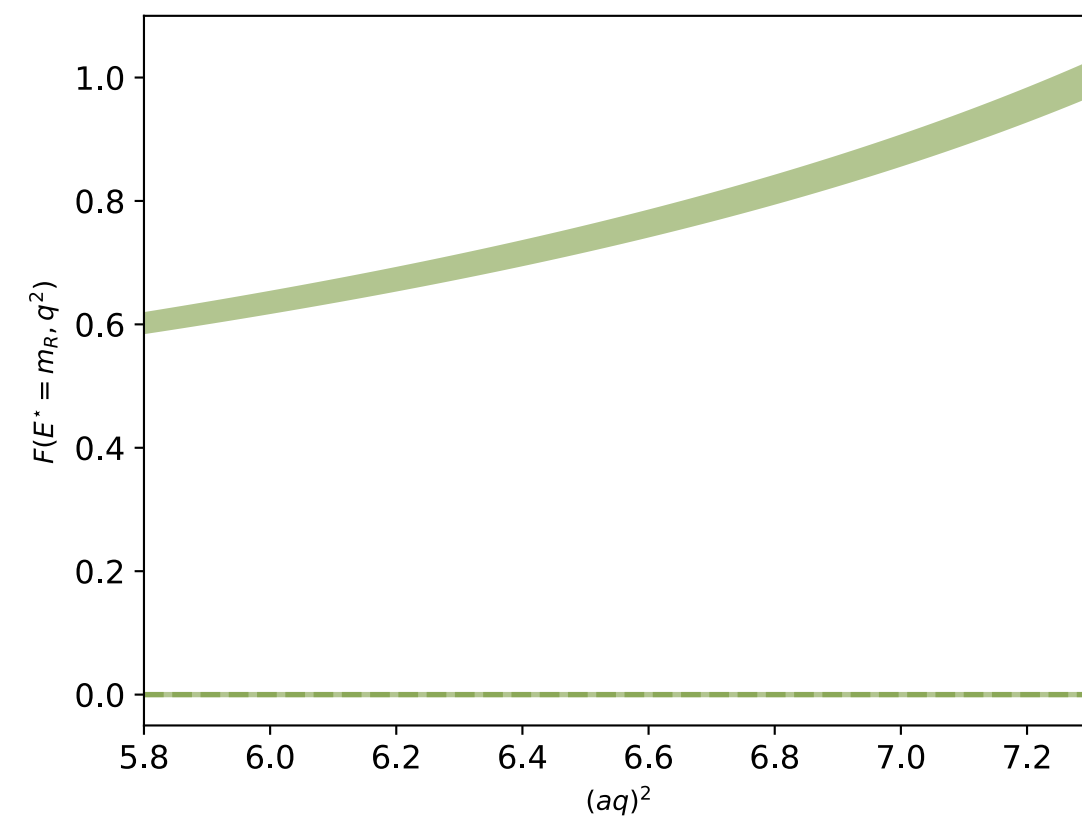


❖ FF3N0M1_TBWII

❖ $\chi^2/\text{dof} = 32.2 / 62 = 0.520$

❖ $A_{0,0} = 0.240(23)$

❖ $A_{0,0} = -0.003(33)$



❖ FF3N1M1_TBWII

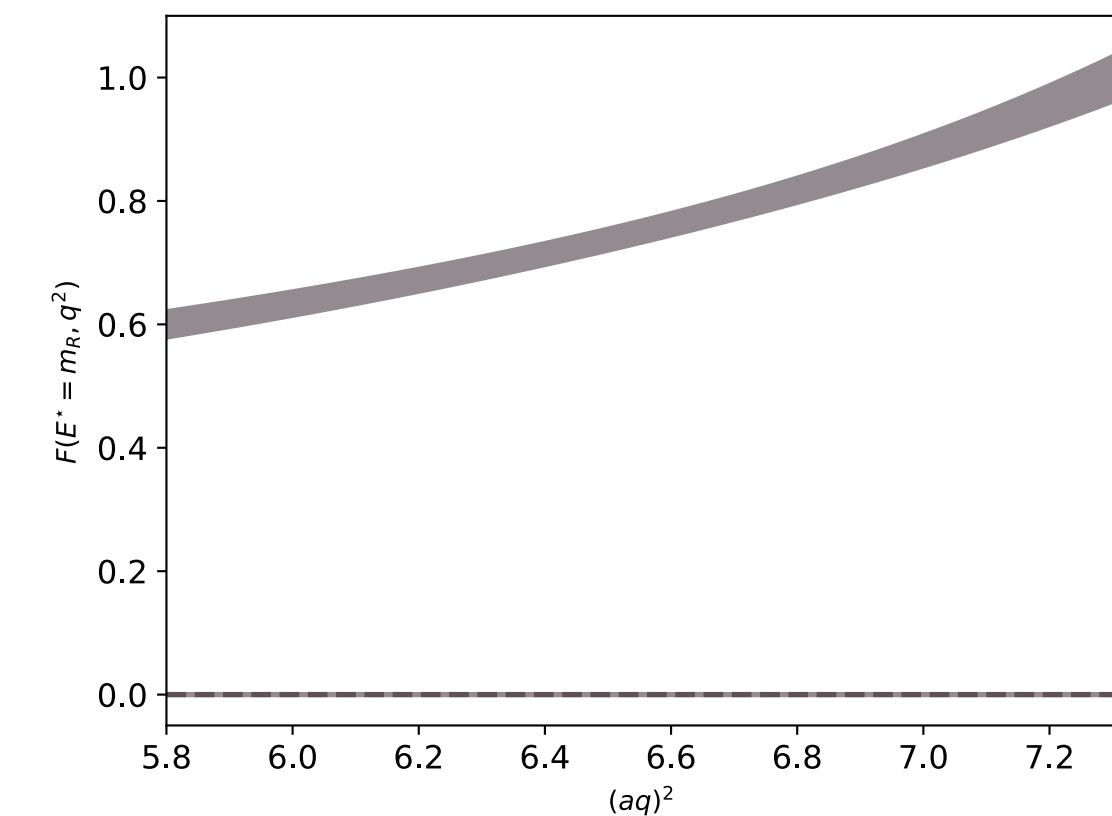
❖ $\chi^2/\text{dof} = 30.4 / 60 = 0.507$

❖ $A_{0,0} = 0.232(20)$

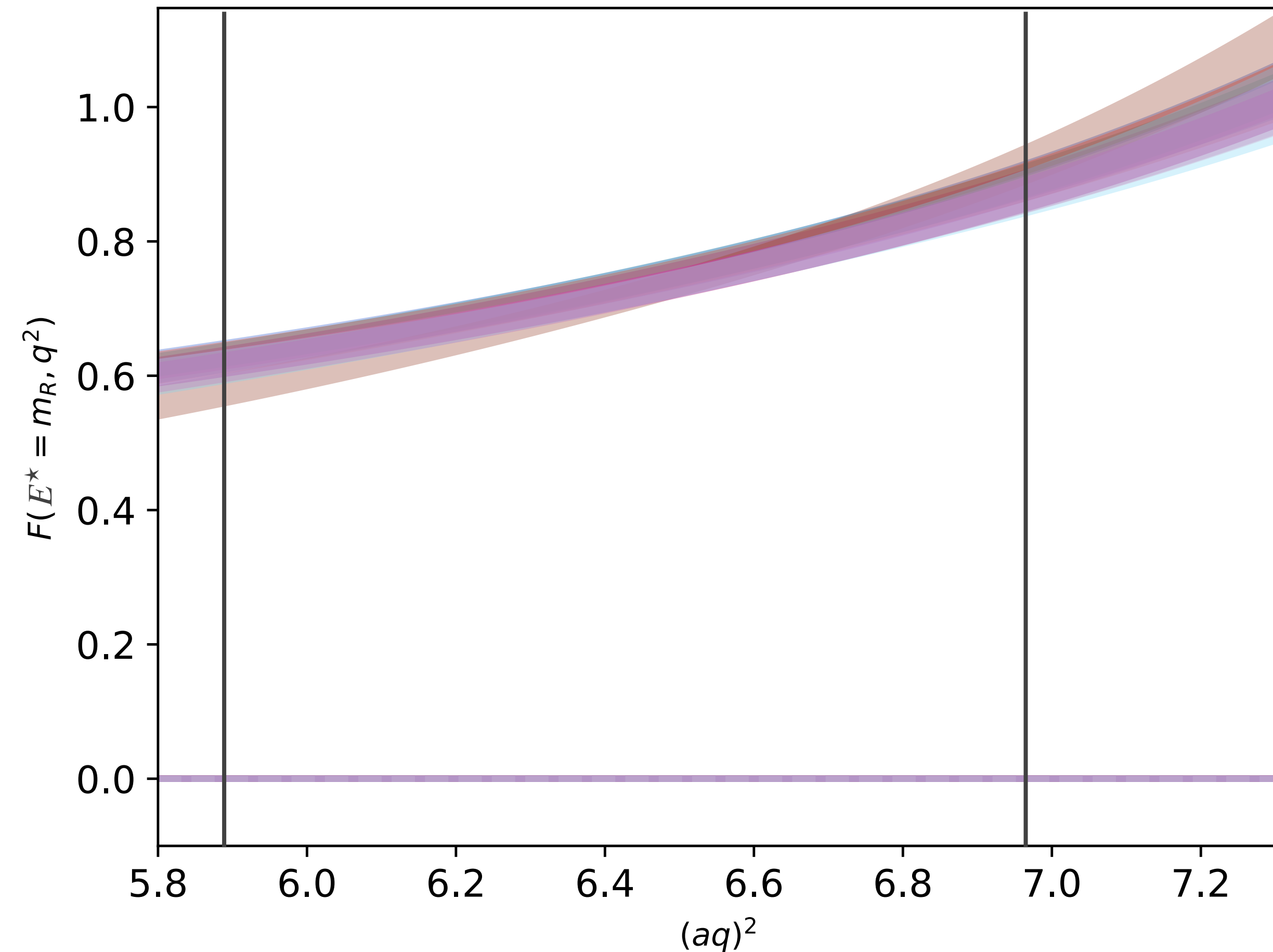
❖ $A_{1,0} = -0.27(60)$

❖ $A_{0,1} = 0.008(28)$

❖ $A_{1,1} = 0.45(1.02)$



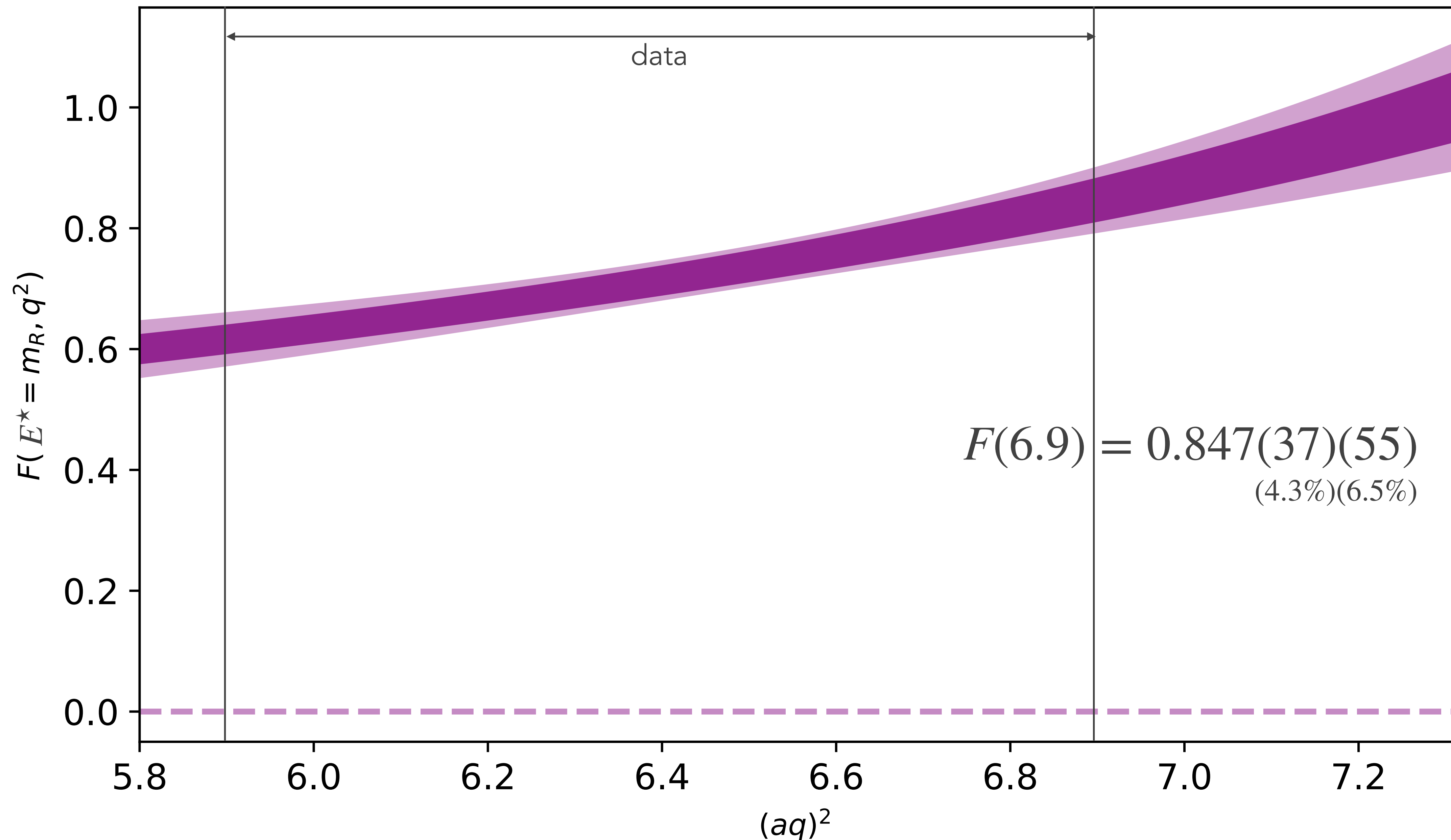
parameterization dependence

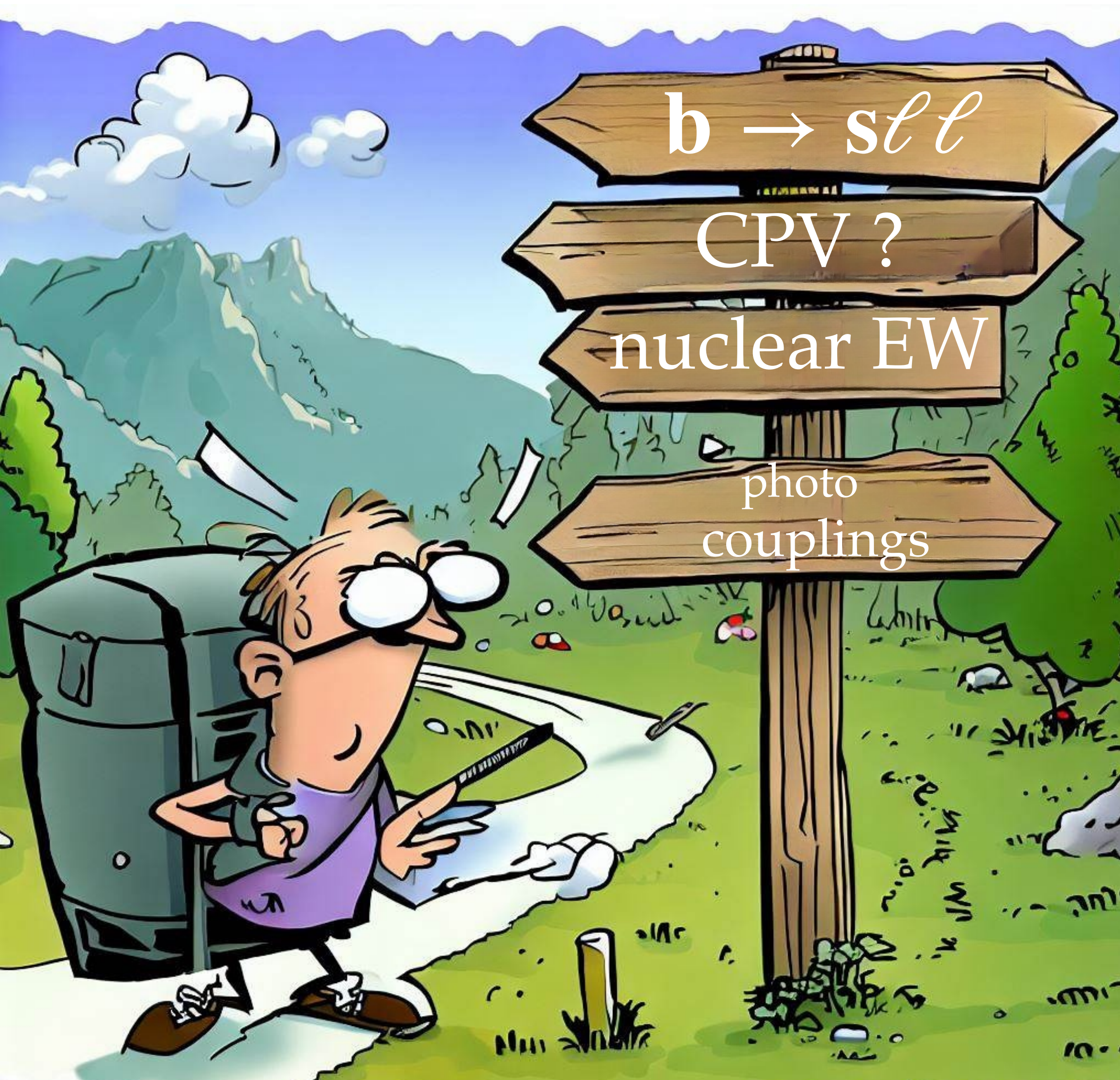


FF1K1_TBWI 0.641	FF1K1_TBWII 0.529
FF3N0M0_TBWI 1.179	FF3N0M0_TBWII 0.545
FF3N0M1_TBWI 0.630	FF3N0M1_TBWII 0.520
FF3N1M0_TBWI 1.078	FF3N1M0_TBWII 0.549
FF3N1M1_TBWI 0.499	FF3N1M1_TBWII 0.507

- ❖ E^* dependence
- ❖ different E^* dependence in BWI and BWII
- ❖ **FF3N1M1_TBWII - central**

combining the parameterizations





- ❖ Vector transition amplitude example
- ❖ similar uncertainties as FLAG!
- ❖ a great start to a (more) complete understanding of SM
- ❖ todo:
 - ❖ m_π continuation to physics!
 - ❖ $a \rightarrow 0$ limit!