

# Three-loop mixing matrix for flavor-singlet operators

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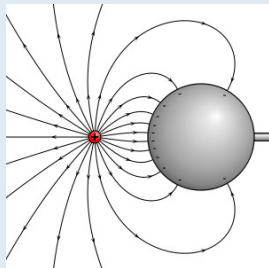
*based on: V. Braun, K. Chetyrkin, A. Manashov, Phys.Lett. B834 (2022) 137409*

QCD Evolution, May 2023



*L.I. Magnus (1831):*  
Inversion transformation

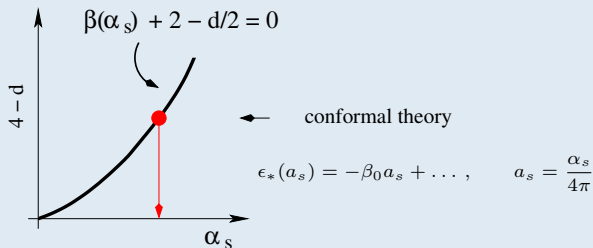
$$x^\mu \rightarrow \frac{x^\mu}{x^2}$$



- Small distances  $\leftrightarrow$  large distances ??
- Asymptotic freedom  $\leftrightarrow$  confinement ??

## QCD?

QCD is not a conformal theory, but



$$\mathcal{A}_{\text{QCD}} = \mathcal{A}_{\text{QCD}}^{\text{conf}} + \frac{\beta(g)}{g} \Delta \mathcal{Q}$$

“Conformal QCD”: QCD in  $d - 2\epsilon$  at Wilson-Fischer critical point  $\beta(\alpha_s) = 0$

V.B., A. Manashov, Eur.Phys.J.C 73 (2013) 2544

- No corrections  $\sim \beta(\alpha_s)$  for the counterterms  
 $\Rightarrow$  exact conformal symmetry of the QCD RG equations in MS-schemes

## Mixing with total derivatives

### Renormalization of off-forward matrix elements

$$\langle p' | \mathcal{O}_{Nk} | p \rangle = \langle p' | \bar{q}(0) \gamma_{\mu_1} \overleftarrow{D}_{\mu_2} \dots \overleftarrow{D}_{\mu_k} \overrightarrow{D}_{\mu_{k+1}} \dots \overrightarrow{D}_{\mu_N} q(0) | p \rangle$$

Mixing matrix with operators containing total derivatives with anomalous dimensions on the diagonal

$$\begin{pmatrix} \gamma_1 & 0 & 0 & 0 \\ * & \gamma_2 & 0 & 0 \\ * & * & \gamma_3 & 0 \\ * & * & * & \gamma_4 \end{pmatrix} \quad \text{need off-diagonal elements}$$

*[evolution kernels in  $x$ -space are functions of two variables]*

- In conformal theories, operators with different scaling dimensions are orthogonal

$$\langle O_1(x_1) O_2(x_2) \rangle = \frac{\text{const}}{|x_1 - x_2|^{2\Delta_1}} \delta_{\Delta_1 \Delta_2}$$

- $n$ -loop ADs from  $n$ -loop two-particle functions

VB, K.Chetyrkin, A.Manashov, PLB 834 (2022) 137409

- Gauge-invariant two-point functions
- Simple algorithmic implementation

simple example:

$$O_2(x) = \partial_+^2 \bar{q}_1(x) C_2^{(3/2)} \left( \begin{array}{c} \overleftarrow{D}_+ - \overrightarrow{D}_+ \\ \overleftarrow{D}_+ + \overrightarrow{D}_+ \end{array} \right) \gamma_+ q_2(x) \quad O_1(x) = \partial_+^2 \bar{q}_1(x) \gamma_+ q_2(x),$$

$$\gamma = \begin{pmatrix} \gamma_{11} & 0 \\ \gamma_{21} & \gamma_{22} \end{pmatrix} \quad \begin{aligned} \gamma_{11} &= 0 \\ \gamma_{22} &= \frac{100}{9} a_s + \left[ \frac{34450}{243} - \frac{830}{81} n_f \right] a_s^2 + \dots \\ \gamma_{21} &= \gamma_{21}^{(2)} a_s^2 + \dots = ? \end{aligned}$$

Let

$$\begin{pmatrix} [\mathbb{O}_1] \\ [\mathbb{O}_2] \end{pmatrix} = \begin{pmatrix} [O_1] \\ [O_2] + \lambda_{21} [O_1] \end{pmatrix}, \quad \lambda_{21} = \frac{\gamma_{21}}{\gamma_{22} - \gamma_{11}}$$

At the critical point

$$(\mu \partial_\mu + \gamma_{11}) \mathbb{O}_1 = 0, \quad (\mu \partial_\mu + \gamma_{22}) \mathbb{O}_2 = 0,$$

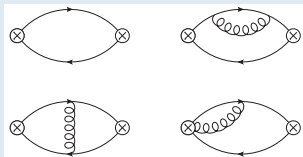
and conformal symmetry requires that to all orders of perturbation theory

$$\langle \mathbb{O}_2(x) \mathbb{O}_1(0) \rangle = \langle [O_2](x) [O_1](0) \rangle + \lambda_{21} \langle [O_1](x) [O_1](0) \rangle = 0$$

— an equation for  $\gamma_{21}$

NB: this equation holds at  $d = 4 - 2\epsilon_*$ , but the result for  $\gamma_{21}$  is valid for  $d = 4$

to leading order  $\mathcal{O}(a_s^2)$



$$\langle \mathbb{T}\{O_1(x)O_1(0)\} \rangle = \mathcal{N} \left[ -105 + \mathcal{O}(a_s, \epsilon) \right],$$

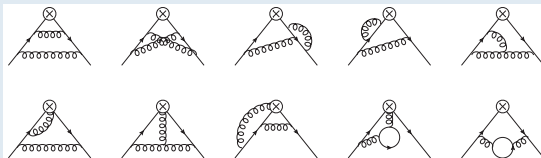
$$\langle \mathbb{T}\{O_2(x)O_1(0)\} \rangle = \mathcal{N} \left[ 63\epsilon + 70a_s + \mathcal{O}(a_s^2, a_s\epsilon, \epsilon^2) \right],$$

$$\mathcal{N} = \frac{(n \cdot x)^6}{(4\pi)^d} \left( \frac{4}{-x^2 + i0} \right)^{d+4}.$$

Using these expressions and the one-loop  $\gamma_{22}$ , expanding to  $\mathcal{O}(a_s)$ , and replacing  $\epsilon \mapsto -\beta_0 a_s$ , one obtains

$$\gamma_{21}^{(2)} = \frac{260}{9} - \frac{40}{9} n_F$$

compare:



- Straightforward to generalize to higher orders, more derivatives and flavor-singlet

## Three-loop flavor singlet anomalous dimensions

$$\mathcal{O}_{nk}^q = i \partial_+^n \sum_{f=1}^{n_f} \bar{q}^f C_k^{(3/2)} \left( \begin{array}{c} \overleftarrow{D}_+ - \overrightarrow{D}_+ \\ \overleftarrow{D}_+ + \overrightarrow{D}_+ \end{array} \right) \gamma_{+q^f},$$

$$\mathcal{O}_{nk}^g = 6 \partial_+^{n-1} F^{\mu,+} C_{k-1}^{(5/2)} \left( \begin{array}{c} \overleftarrow{D}_+ - \overrightarrow{D}_+ \\ \overleftarrow{D}_+ + \overrightarrow{D}_+ \end{array} \right) F_{\mu,+}.$$

$$\gamma_n(a) = \begin{pmatrix} \gamma_{11} & 0 & \cdots & 0 \\ \gamma_{31} & \gamma_{33} & \cdots & 0 \\ \vdots & \vdots & \ddots & \vdots \\ \gamma_{n1} & \gamma_{n2} & \cdots & \gamma_{nn} \end{pmatrix}.$$

$$\gamma_{nk} = \begin{pmatrix} \gamma_{nk}^{qq} & \gamma_{nk}^{qg} \\ \gamma_{nk}^{gq} & \gamma_{nk}^{gg} \end{pmatrix} = a_s \gamma_{nk}^{(1)} + a_s^2 \gamma_{nk}^{(2)} + a_s^3 \gamma_{nk}^{(3)} + \dots$$

- Three-loop off-forward anomalous dimensions

$$\gamma^{(3)} = \gamma^{(3,0)} + n_f \gamma^{(3,1)} + n_f^2 \gamma^{(3,2)}$$

$$\gamma_{31}^{(3,0)} = \begin{pmatrix} \frac{36623912}{54675} & 0 \\ -\frac{2430374}{3645} & \frac{261063}{50} \end{pmatrix}$$

$$\gamma_{51}^{(3,0)} = \begin{pmatrix} \frac{8049304723}{31255875} & 0 \\ -\frac{26632998209}{112521150} & \frac{2829671009}{329280} \end{pmatrix},$$

$$\gamma_{53}^{(3,0)} = \begin{pmatrix} \frac{320657981731}{520931250} & 0 \\ -\frac{29333397389}{20837250} & \frac{14378664569}{6860000} \end{pmatrix}$$

$$\gamma_{71}^{(3,0)} = \begin{pmatrix} \frac{7192640196053}{56710659600} & 0 \\ \frac{52031947546}{506345175} & \frac{49155659027}{3969000} \end{pmatrix},$$

$$\gamma_{73}^{(3,0)} = \begin{pmatrix} \frac{159898280729473}{525098700000} & 0 \\ -\frac{5108698450661}{3750705000} & \frac{832037077}{441000} \end{pmatrix}$$

$$\gamma_{75}^{(3,0)} = \begin{pmatrix} \frac{220023775251709}{396974617200} & 0 \\ -\frac{10780083012803}{7088832450} & \frac{16149051685793}{9724050000} \end{pmatrix}$$

etc.



Summary:

three-loop off-forward vector ADs obtained using three-loop two-particle correlators

- Gauge-invariant two-point functions
  - Simple algorithmic implementation
  - Axial-vector flavor singlets — work in progress
  - Straightforward to extend to four loops (MINCER → FORCER)
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- $2N - 2$  open indices, progressing to high  $N$  computationally expensive
  - Restore evolution kernels from AD matrices?