

# General review of Higgs properties

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(On behalf of ATLAS experiment)

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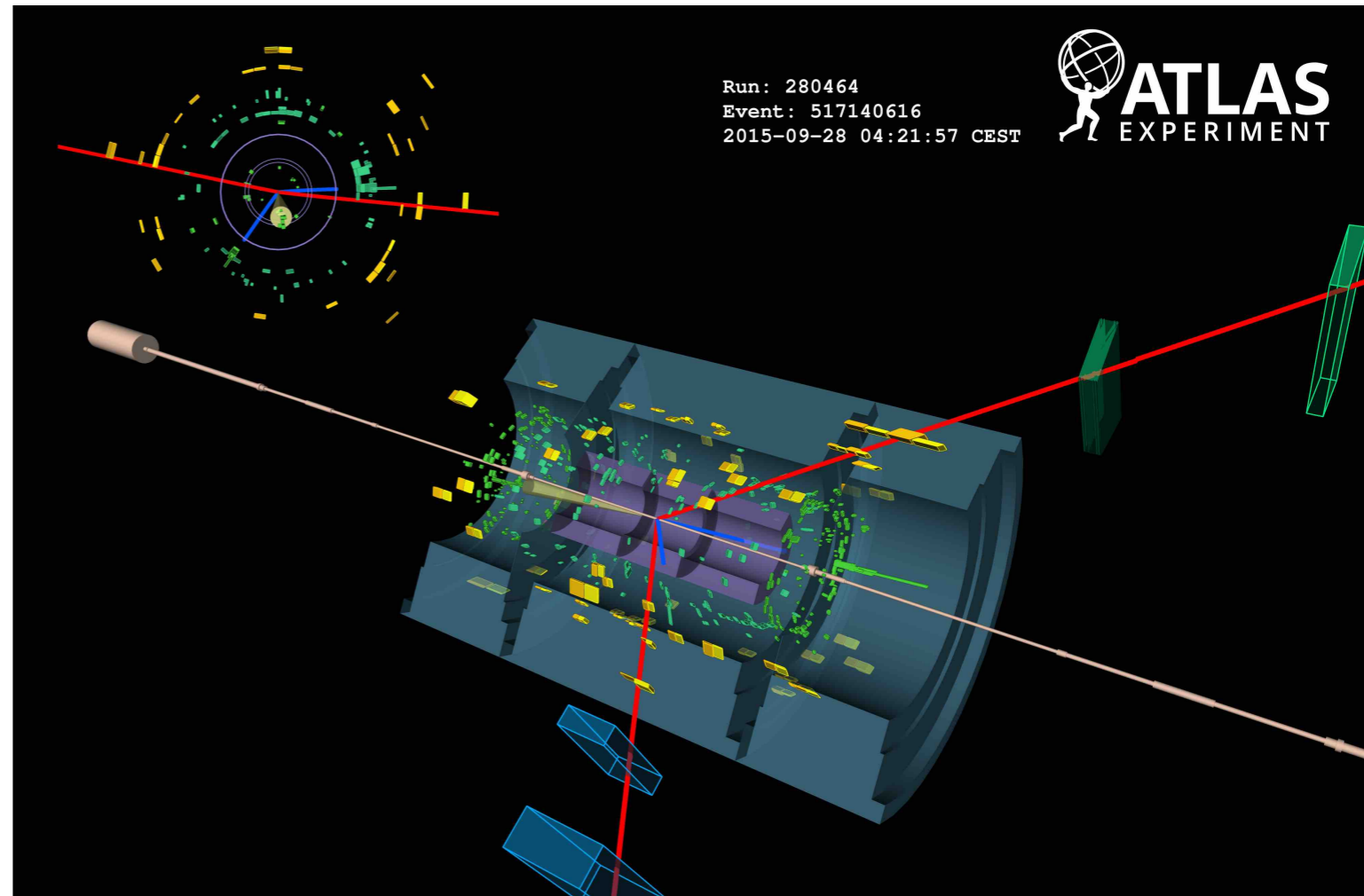
# How well do we know the 10 years old

$$\mathcal{L} = -g_{Hf\bar{f}}\bar{f}fH + \frac{g_{HHH}}{6}H^3 + \frac{g_{HHHH}}{24}H^4 + \delta_V V_\mu V^\mu \left( g_{HVV}H + \frac{g_{HHVV}}{2}H^2 \right)$$
$$g_{Hf\bar{f}} \equiv y_f = \frac{m_f}{v}, \quad g_{HVV} = \frac{2m_V^2}{v}, \quad g_{HHVV} = \frac{2m_V^2}{v^2}, \quad g_{HHH} = \frac{3m_H^2}{v}, \quad g_{HHHH} = \frac{3m_H^2}{v^2}.$$

## Outline

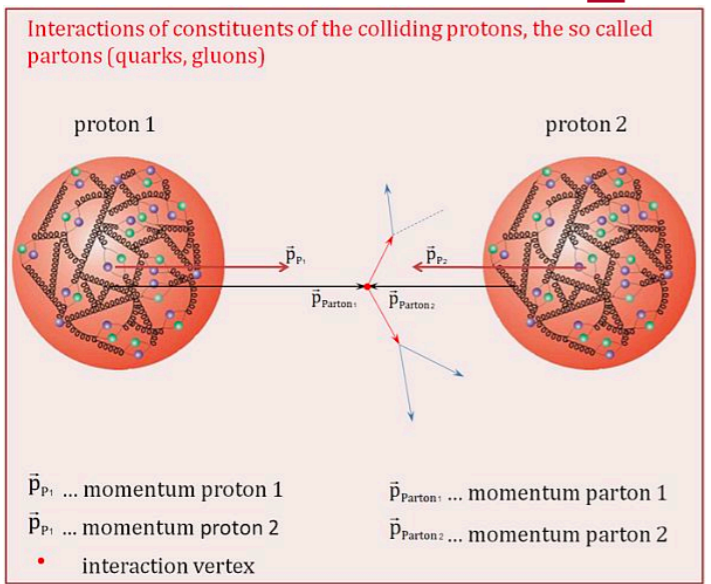
### ► Introduction

- ❑  $H \rightarrow \gamma\gamma$  STXS measurement
- ❑  $H$  width from off-shell  $ZZ^*$
- ❑ Higgs CP from  $H \rightarrow ZZ^* \rightarrow 4l$
- ❑ Search for HH in 4b final state
- ❑ Charm Yukawa in VH mode
- ❑ Summary



<https://twiki.cern.ch/twiki/bin/view/AtlasPublic/EventDisplayRun2Physics>

# We need to produce some H to detect



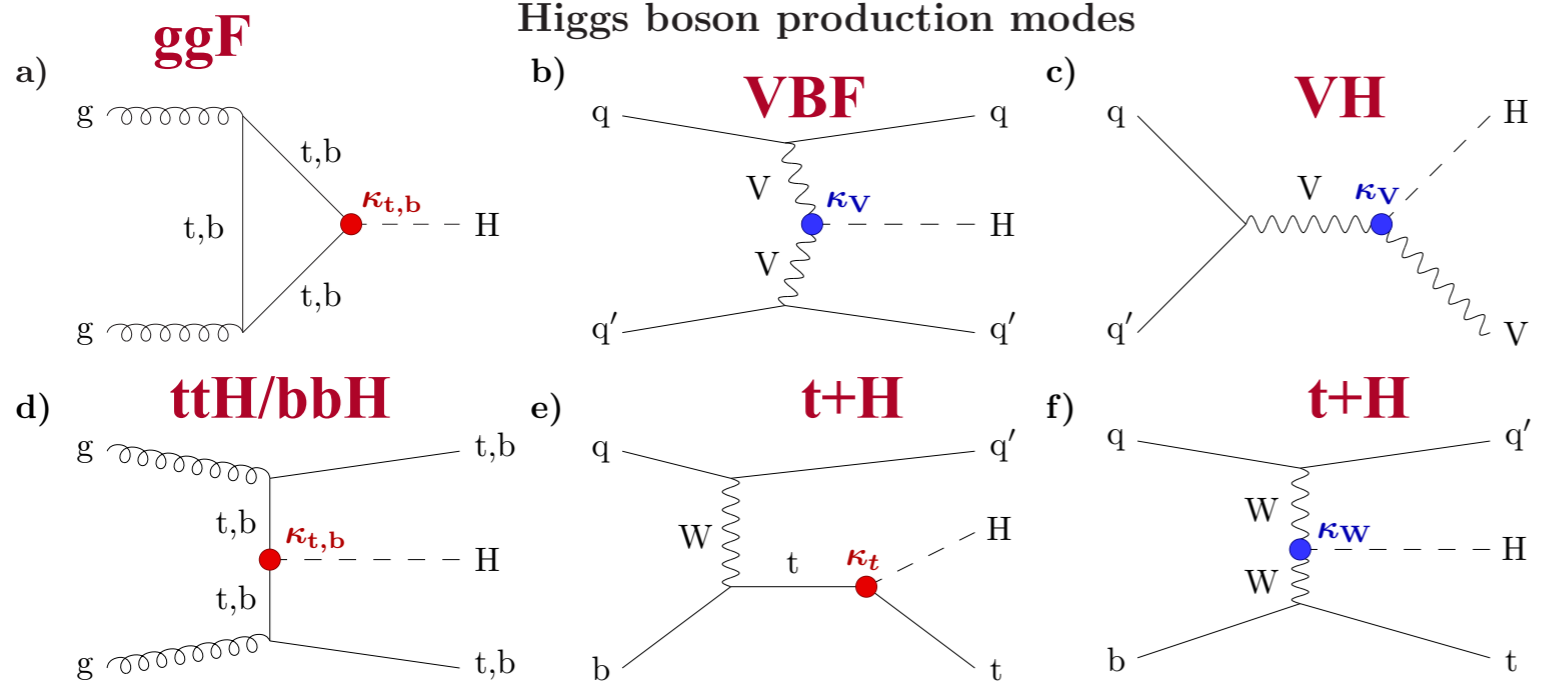
$$C = \kappa_C C_{SM}$$

$$\bar{N}(pp \rightarrow H \rightarrow ab) \sim$$

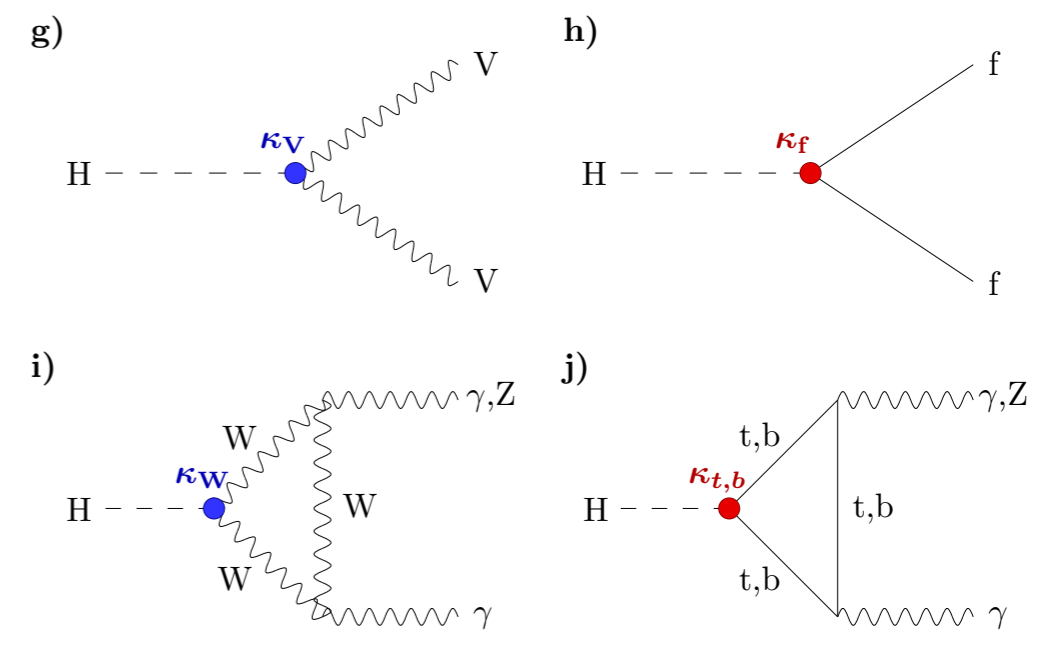
$$\int dx_1 dx_2 \sum_{ij} f_i(x_1) f_j(x_2) \hat{\sigma}(i + j \rightarrow H) \times \mathcal{BR}(H \rightarrow a + b)$$

Nature volume 607, pages 60–68 (2022)

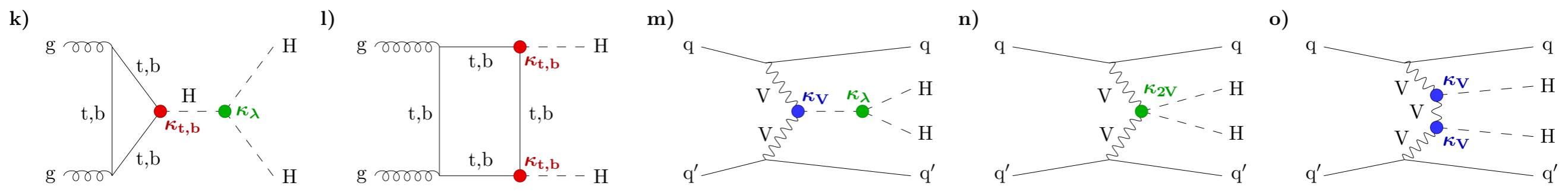
Higgs boson production modes



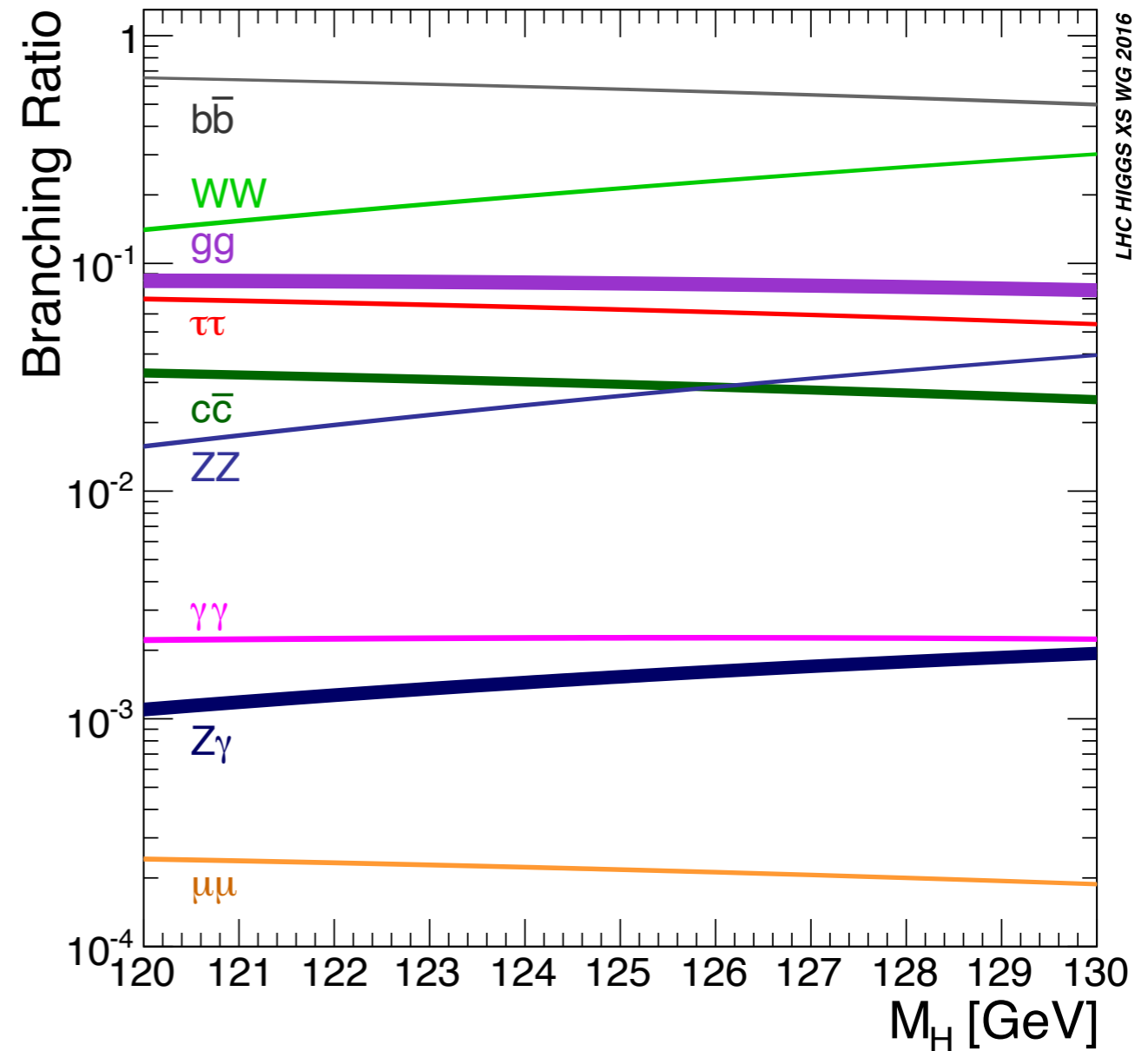
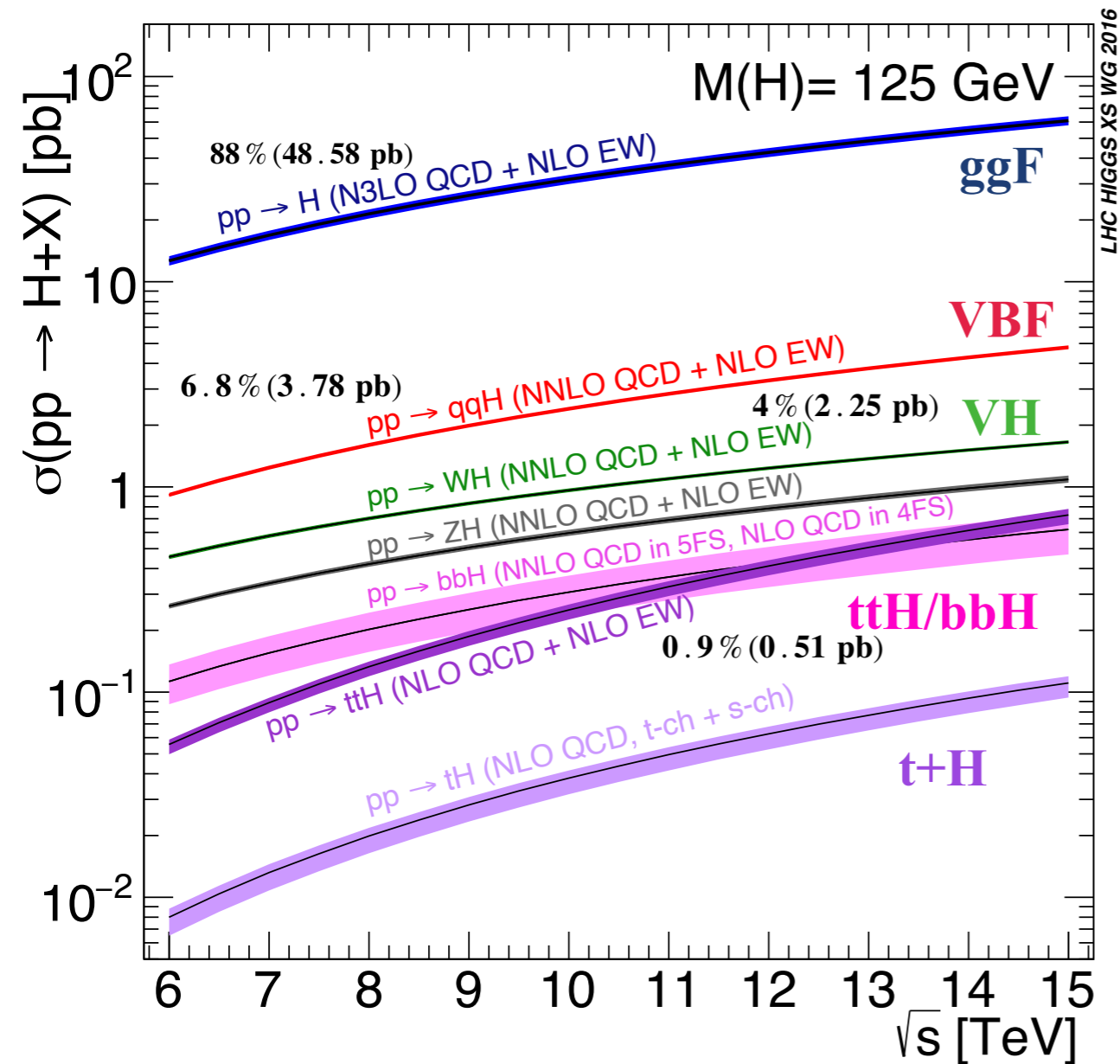
Higgs boson decay channels



Higgs boson pair production

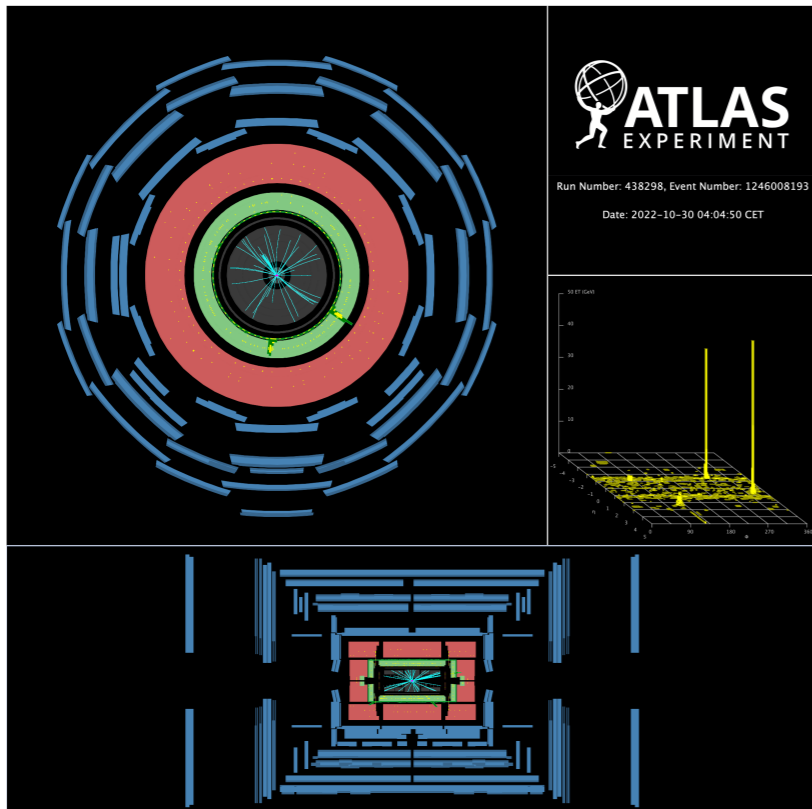
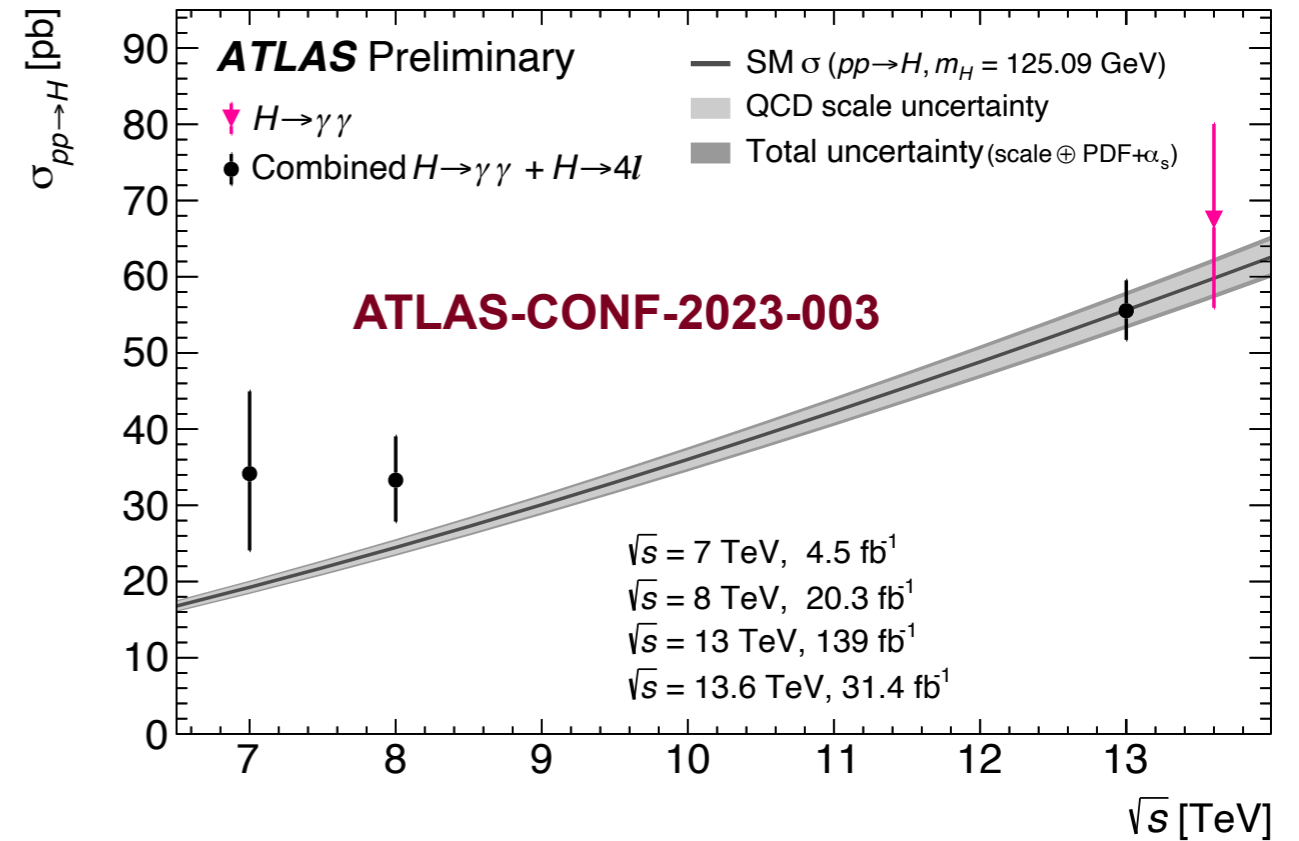
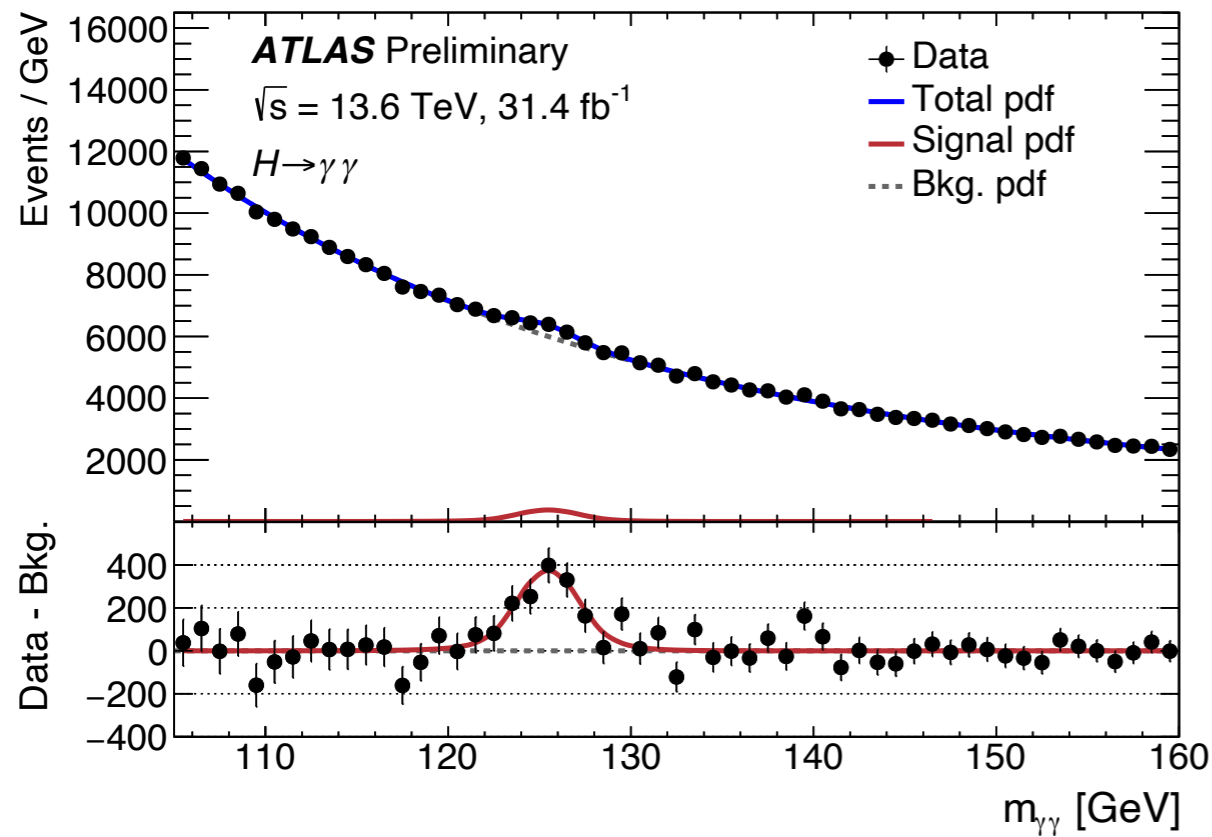


# We measure a combination



The major production and decay modes of Higgs boson  
<https://twiki.cern.ch/twiki/bin/view/LHCPhysics/LHCHWG>

# First Higgs measurement at 13.6 TeV



Di-photon trigger with  $E_T > 25, 35$  GeV with medium selection criteria.  
 Trigger efficiency  $> 99.4\%$ .  
 The NN based PV reconstruction efficiency is  $71.4\%$

Di-photon fiducial region :  $|\eta| < 2.37$ , modulo  $1.37 < |\eta| < 1.52$   
 $105 \text{ GeV} < m_{\gamma\gamma} < 160 \text{ GeV}$  for isolated photons.  
 This fiducial region is  $50\%$  of the total phase space.

$$\sigma_{\text{fid}}(pp \rightarrow H \rightarrow \gamma\gamma) = 76_{-13}^{+14} \text{ fb} = 76 \pm 11(\text{stat})_{-7}^{+9}(\text{syst}) \text{ fb}$$

$67.5 \pm 3.4 \text{ fb.}$  SM prediction

# Outline

Introduction

▶  $H \rightarrow \gamma\gamma$  STXS measurement

H width from off-shell  $ZZ^*$

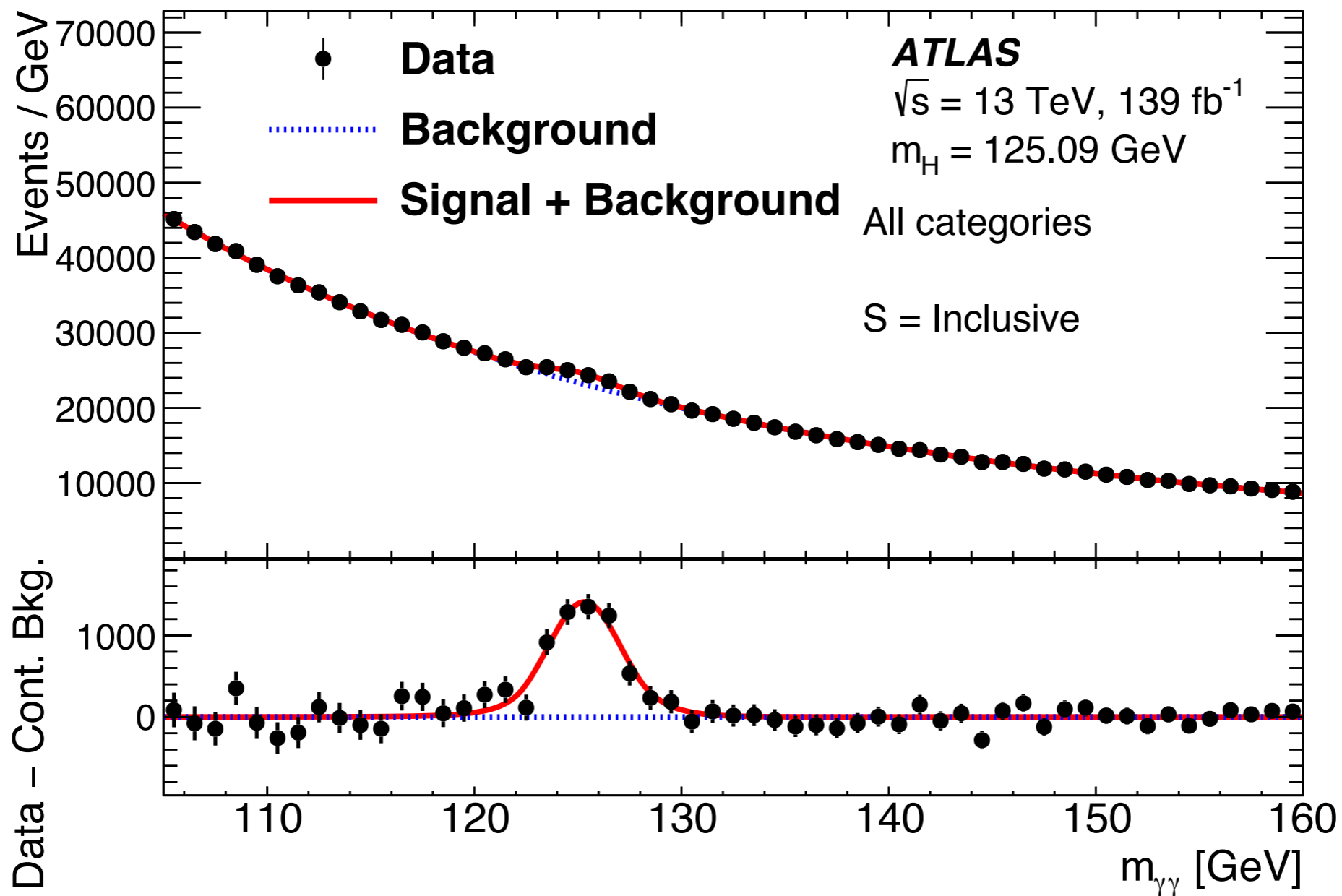
Higgs CP from  $H \rightarrow ZZ^* \rightarrow 4l$

Search for HH in  $4b$  final state

Charm Yukawa in VH mode

Summary

# H properties from $H \rightarrow \gamma\gamma$ : STXS

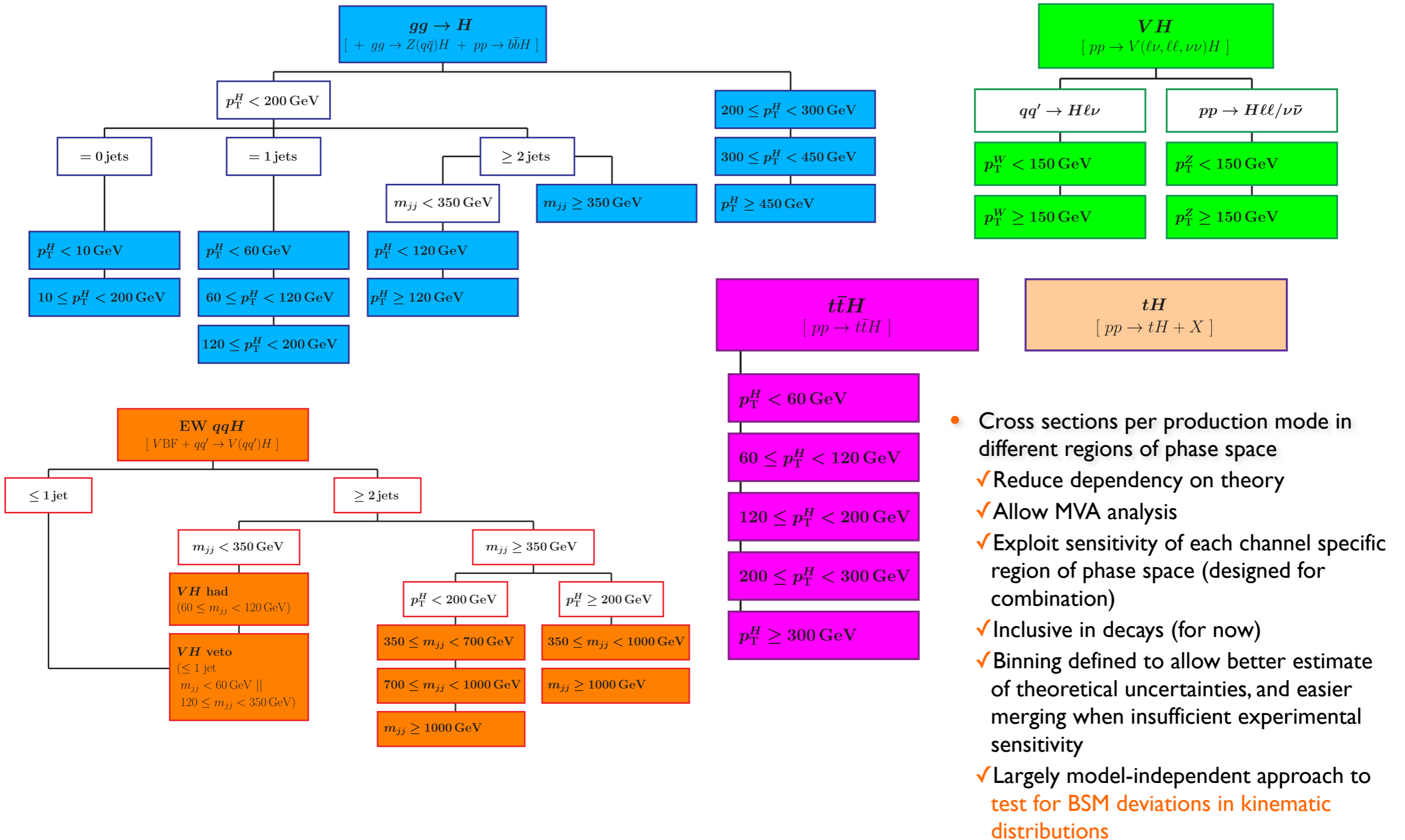


The strategy of this analysis is to divide the Higgs productions measurements into different categories, emulating different physics process.

# H properties from $H \rightarrow \gamma\gamma$ : STXS

arXiv 2207.00348

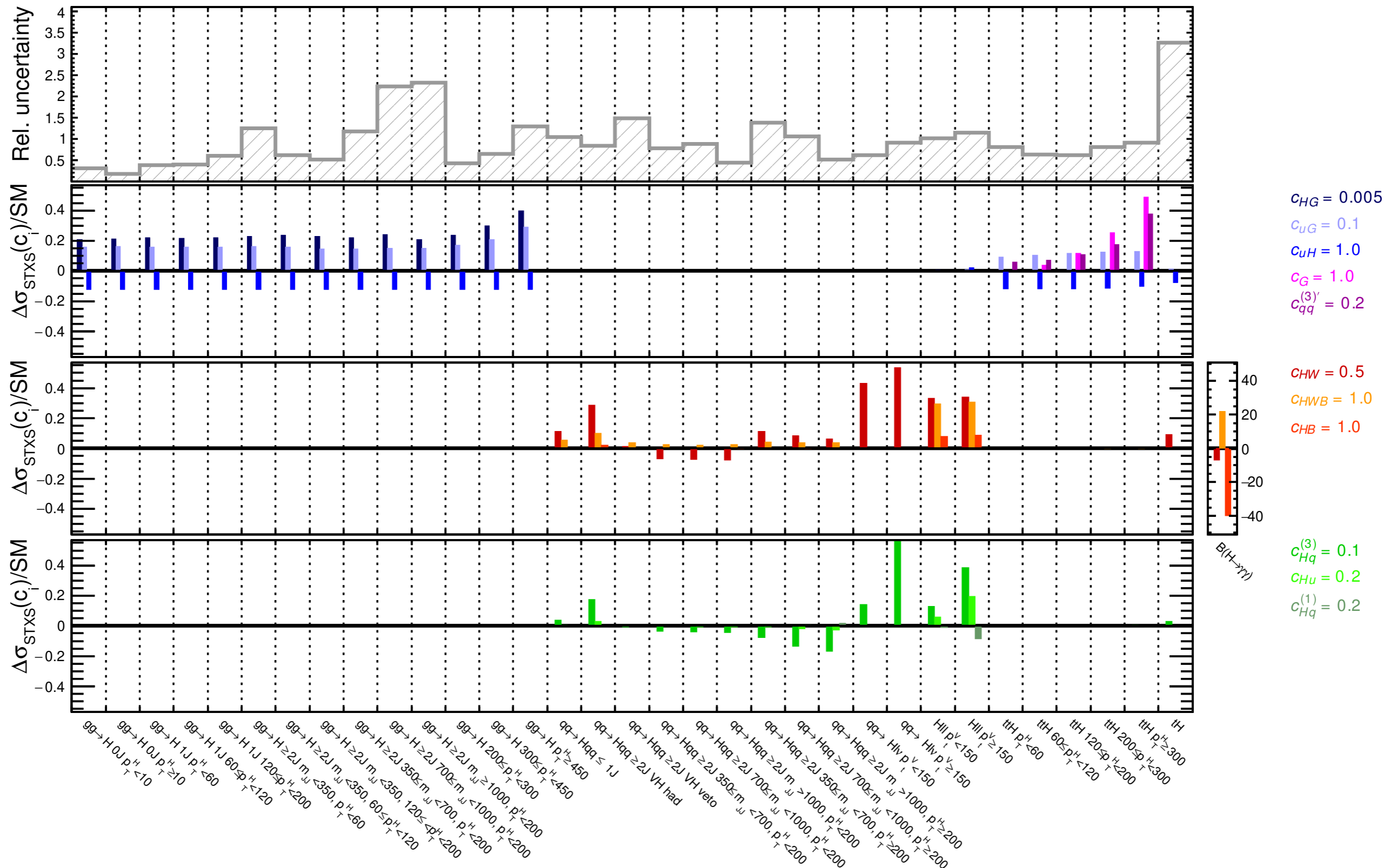
$H \rightarrow \gamma\gamma$  production x-section is measured in 48 different STXS regions.





# Impact on SMEFT parameters

**ATLAS** Simulation  $\sqrt{s}=13$  TeV  $139\text{fb}^{-1}$   $H \rightarrow \gamma\gamma$ ,  $m_H = 125.09$  GeV,  $\Lambda = 1$  TeV



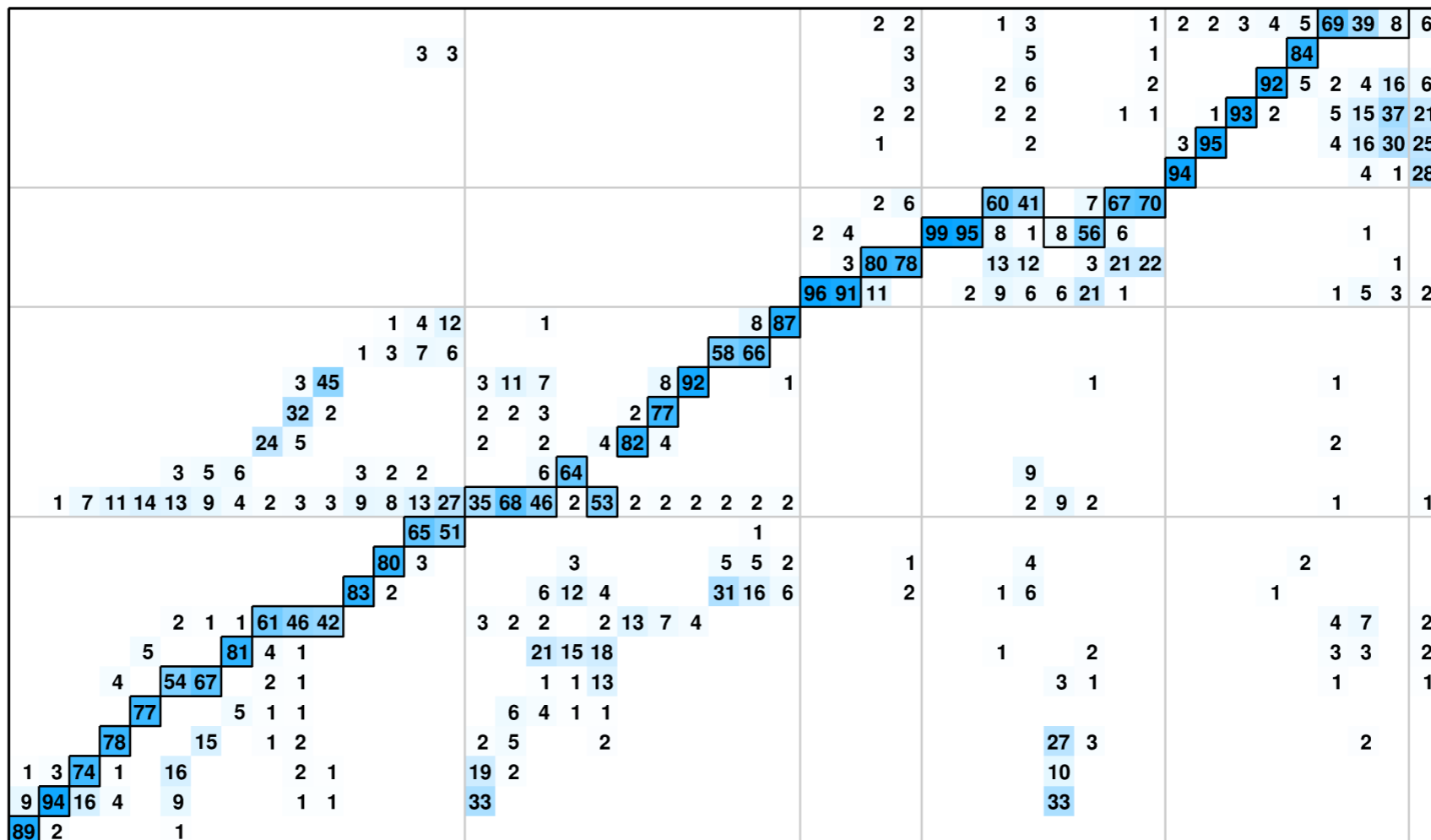
# Contribution of STXS regions to categories

ATLAS Simulation 139 fb<sup>-1</sup>

H → γγ, √s = 13 TeV

STXS Region

- tH
- tH, p<sub>T</sub><sup>H</sup> ≥ 300 GeV
- tH, 200 ≤ p<sub>T</sub><sup>H</sup> < 300 GeV
- tH, 120 ≤ p<sub>T</sub><sup>H</sup> < 200 GeV
- tH, 60 ≤ p<sub>T</sub><sup>H</sup> < 120 GeV
- tH, p<sub>T</sub><sup>H</sup> < 60 GeV
- Hll, p<sub>T</sub><sup>V</sup> ≥ 150 GeV
- Hll, p<sub>T</sub><sup>V</sup> < 150 GeV
- qq → Hlv, p<sub>T</sub><sup>V</sup> ≥ 150 GeV
- qq → Hlv, p<sub>T</sub><sup>V</sup> < 150 GeV
- qq → Hqq, ≥ 2-jets, m<sub>jj</sub> ≥ 1000 GeV, p<sub>T</sub><sup>H</sup> ≥ 200 GeV
- qq → Hqq, ≥ 2-jets, 350 ≤ m<sub>jj</sub> < 1000 GeV, p<sub>T</sub><sup>H</sup> ≥ 200 GeV
- qq → Hqq, ≥ 2-jets, m<sub>jj</sub> ≥ 1000, p<sub>T</sub><sup>H</sup> < 200 GeV
- qq → Hqq, ≥ 2-jets, 700 ≤ m<sub>jj</sub> < 1000 GeV, p<sub>T</sub><sup>H</sup> < 200 GeV
- qq → Hqq, ≥ 2-jets, 350 ≤ m<sub>jj</sub> < 700 GeV, p<sub>T</sub><sup>H</sup> < 200 GeV
- qq → Hqq, VH hadronic
- qq → Hqq, ≤ 1-jet, VH veto
- gg → H, p<sub>T</sub><sup>H</sup> ≥ 450 GeV
- gg → H, 300 ≤ p<sub>T</sub><sup>H</sup> < 450 GeV
- gg → H, 200 ≤ p<sub>T</sub><sup>H</sup> < 300 GeV
- gg → H, ≥ 2-jets, m<sub>jj</sub> ≥ 350 GeV, p<sub>T</sub><sup>H</sup> < 200 GeV
- gg → H, ≥ 2-jets, m<sub>jj</sub> < 350 GeV, 120 ≤ p<sub>T</sub><sup>H</sup> < 200 GeV
- gg → H, ≥ 2-jets, m<sub>jj</sub> < 350 GeV, p<sub>T</sub><sup>H</sup> < 120 GeV
- gg → H, 1-jet, 120 ≤ p<sub>T</sub><sup>H</sup> < 200 GeV
- gg → H, 1-jet, 60 ≤ p<sub>T</sub><sup>H</sup> < 120 GeV
- gg → H, 1-jet, p<sub>T</sub><sup>H</sup> < 60 GeV
- gg → H, 0-jet, p<sub>T</sub><sup>H</sup> ≥ 10 GeV
- gg → H, 0-jet, p<sub>T</sub><sup>H</sup> < 10 GeV

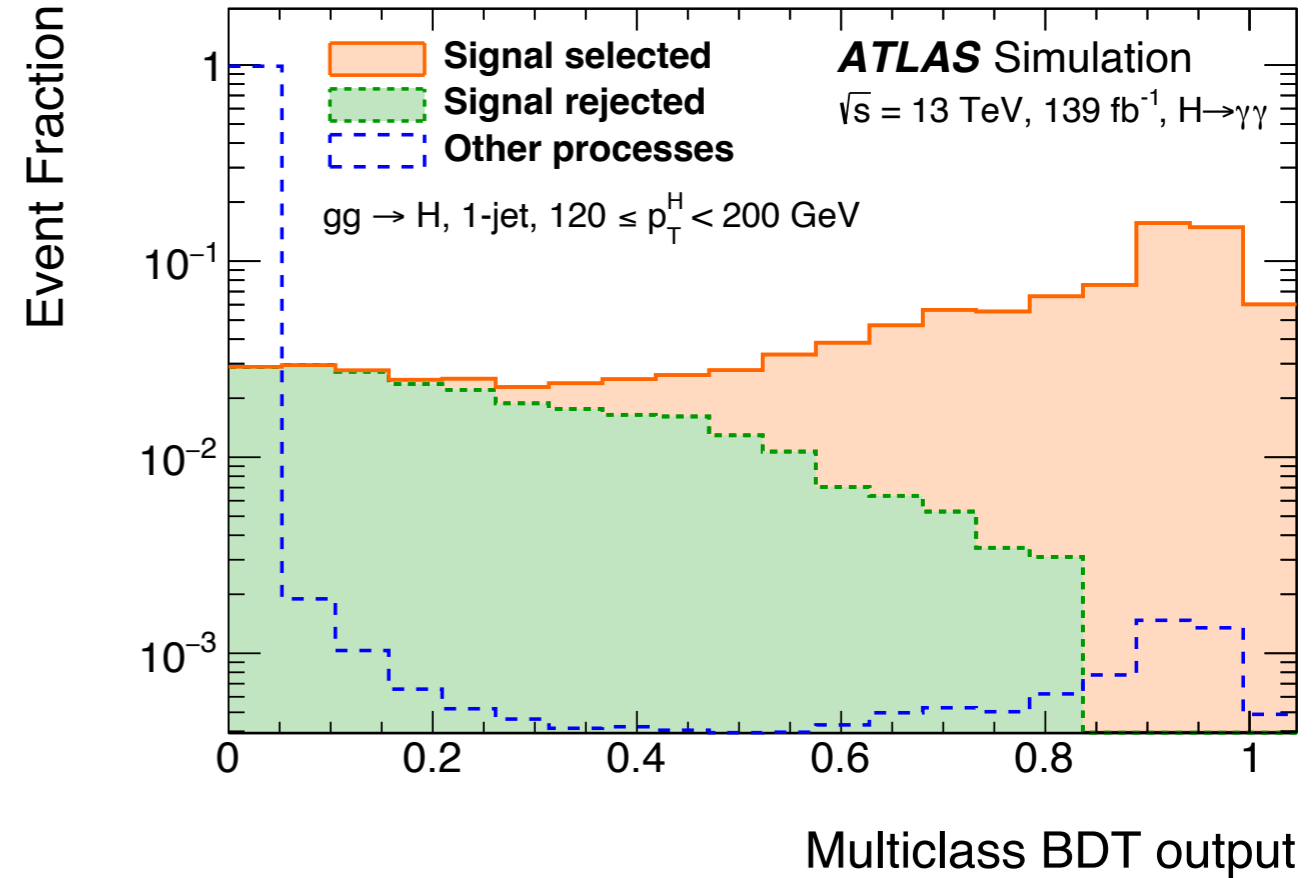
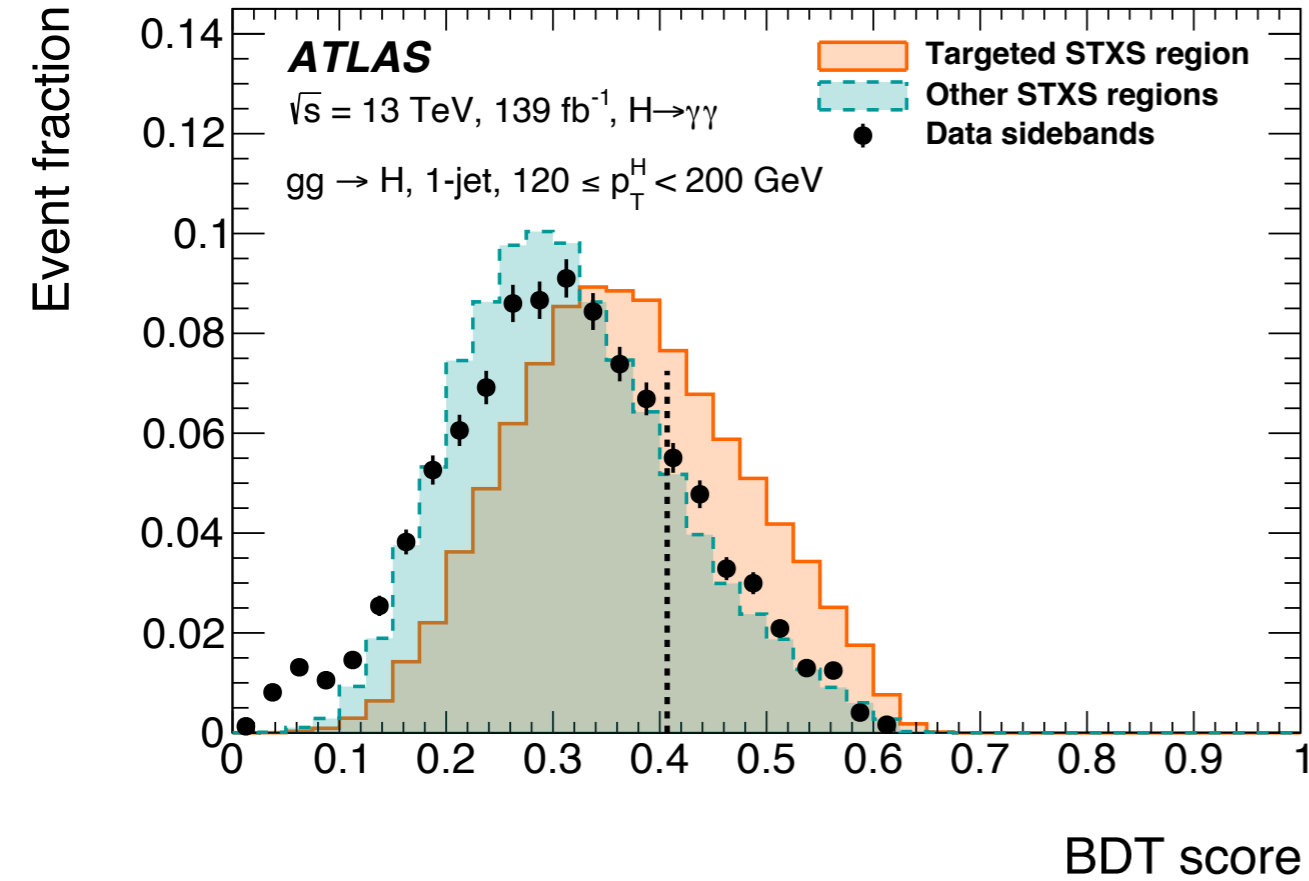


- gg → H, 0-jet, p<sub>T</sub><sup>H</sup> < 10 GeV
- gg → H, 0-jet, p<sub>T</sub><sup>H</sup> ≥ 10 GeV
- gg → H, 1-jet, p<sub>T</sub><sup>H</sup> < 60 GeV
- gg → H, 1-jet, 60 ≤ p<sub>T</sub><sup>H</sup> < 120 GeV
- gg → H, 1-jet, 120 ≤ p<sub>T</sub><sup>H</sup> < 200 GeV
- gg → H, ≥ 2-jets, m<sub>jj</sub> < 350 GeV, p<sub>T</sub><sup>H</sup> < 60 GeV
- gg → H, ≥ 2-jets, m<sub>jj</sub> < 350 GeV, 60 ≤ p<sub>T</sub><sup>H</sup> < 120 GeV
- gg → H, ≥ 2-jets, m<sub>jj</sub> < 350 GeV, 120 ≤ p<sub>T</sub><sup>H</sup> < 200 GeV
- gg → H, ≥ 2-jets, m<sub>jj</sub> < 350 GeV, p<sub>T</sub><sup>H</sup> < 200 GeV
- gg → H, ≥ 2-jets, 350 ≤ m<sub>jj</sub> < 700 GeV, p<sub>T</sub><sup>H</sup> < 200 GeV
- gg → H, ≥ 2-jets, 700 ≤ m<sub>jj</sub> < 1000 GeV, p<sub>T</sub><sup>H</sup> < 200 GeV
- gg → H, ≥ 2-jets, m<sub>jj</sub> ≥ 1000 GeV, p<sub>T</sub><sup>H</sup> < 200 GeV
- gg → H, 200 ≤ p<sub>T</sub><sup>H</sup> < 300 GeV
- gg → H, 300 ≤ p<sub>T</sub><sup>H</sup> < 450 GeV
- gg → H, 450 ≤ p<sub>T</sub><sup>H</sup> < 650 GeV
- gg → H, p<sub>T</sub><sup>H</sup> ≥ 650 GeV
- qq → Hqq, 0-jet
- qq → Hqq, 1-jet
- qq → Hqq, ≥ 2-jets, m<sub>jj</sub> < 60 GeV
- qq → Hqq, ≥ 2-jets, 60 ≤ m<sub>jj</sub> < 120 GeV
- qq → Hqq, ≥ 2-jets, 120 ≤ m<sub>jj</sub> < 350 GeV
- qq → Hqq, ≥ 2-jets, 350 ≤ m<sub>jj</sub> < 700 GeV, p<sub>T</sub><sup>H</sup> < 200 GeV
- qq → Hqq, ≥ 2-jets, 700 ≤ m<sub>jj</sub> < 1000 GeV, p<sub>T</sub><sup>H</sup> < 200 GeV
- qq → Hqq, ≥ 2-jets, m<sub>jj</sub> ≥ 1000 GeV, p<sub>T</sub><sup>H</sup> < 200 GeV
- qq → Hqq, ≥ 2-jets, 350 ≤ m<sub>jj</sub> < 700 GeV, p<sub>T</sub><sup>H</sup> ≥ 200 GeV
- qq → Hqq, ≥ 2-jets, 700 ≤ m<sub>jj</sub> < 1000 GeV, p<sub>T</sub><sup>H</sup> ≥ 200 GeV
- qq → Hqq, ≥ 2-jets, m<sub>jj</sub> ≥ 1000 GeV, p<sub>T</sub><sup>H</sup> ≥ 200 GeV
- qq → Hlv, p<sub>T</sub><sup>V</sup> < 75 GeV
- qq → Hlv, 75 ≤ p<sub>T</sub><sup>V</sup> < 150 GeV
- qq → Hlv, 150 ≤ p<sub>T</sub><sup>V</sup> < 250 GeV
- qq → Hlv, p<sub>T</sub><sup>V</sup> ≥ 250 GeV
- Hll, p<sub>T</sub><sup>V</sup> < 75 GeV
- Hll, 75 ≤ p<sub>T</sub><sup>V</sup> < 150 GeV
- Hll, 150 ≤ p<sub>T</sub><sup>V</sup> < 250 GeV
- Hll, p<sub>T</sub><sup>V</sup> ≥ 250 GeV
- Hvv, p<sub>T</sub><sup>V</sup> < 75 GeV
- Hvv, 75 ≤ p<sub>T</sub><sup>V</sup> < 150 GeV
- Hvv, 150 ≤ p<sub>T</sub><sup>V</sup> < 250 GeV
- Hvv, p<sub>T</sub><sup>V</sup> ≥ 250 GeV
- tH, p<sub>T</sub><sup>H</sup> < 60 GeV
- tH, 60 ≤ p<sub>T</sub><sup>H</sup> < 120 GeV
- tH, 120 ≤ p<sub>T</sub><sup>H</sup> < 200 GeV
- tH, 200 ≤ p<sub>T</sub><sup>H</sup> < 300 GeV
- tH, p<sub>T</sub><sup>H</sup> ≥ 300 GeV
- tHjb, SM-like
- tHjb, BSM-like
- tHW
- low-purity top

Category based on reco-objects while STXS regions refers to particle level selections.

Analysis Category

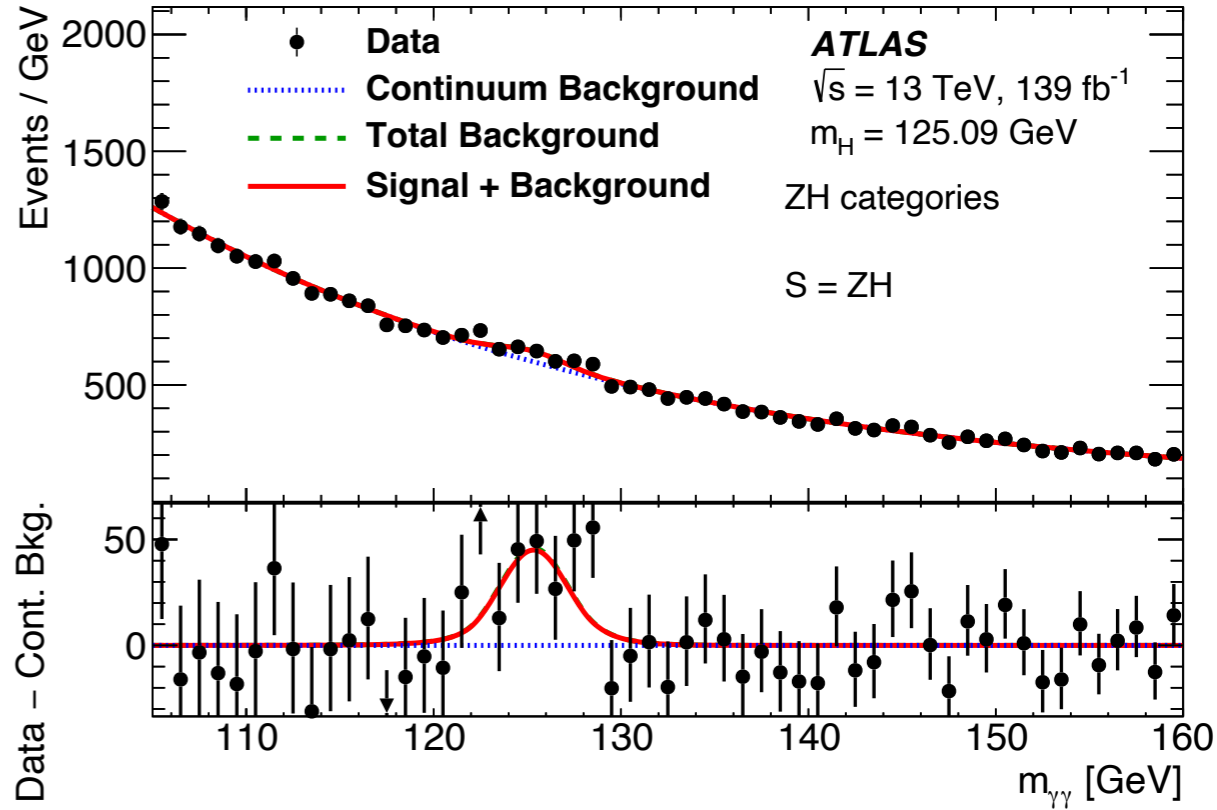
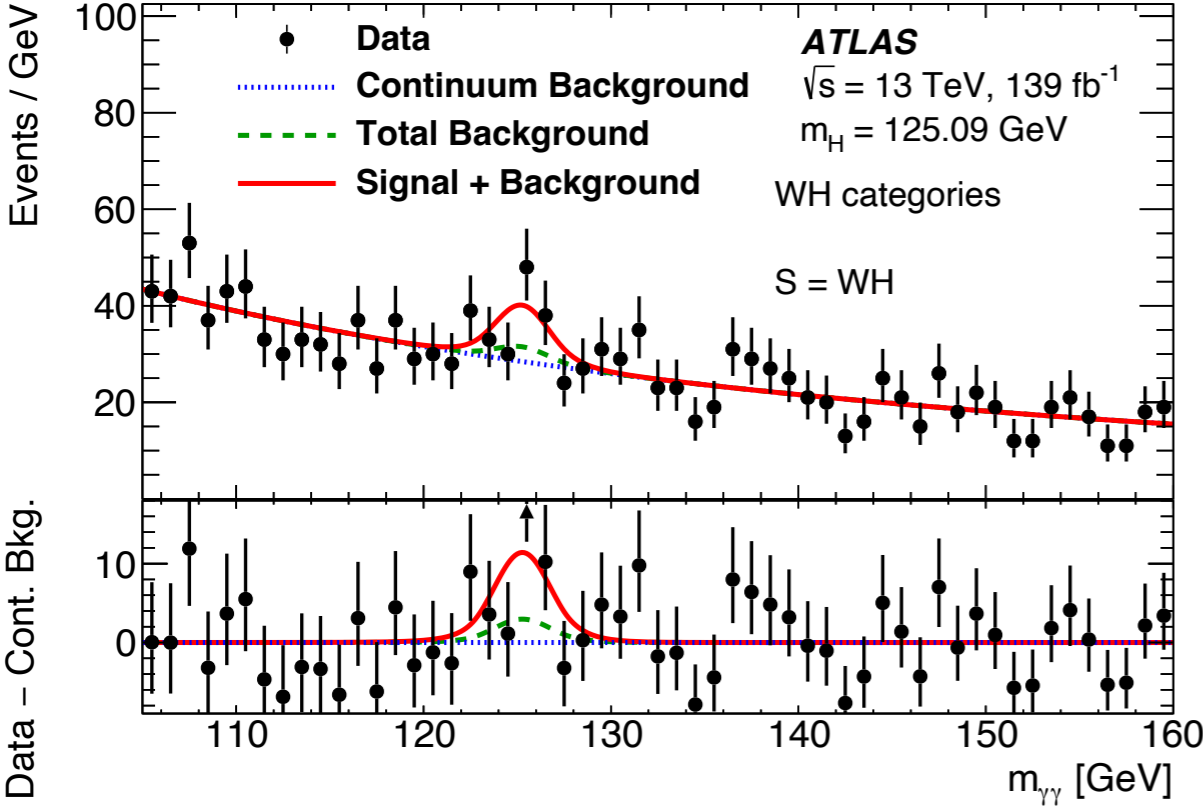
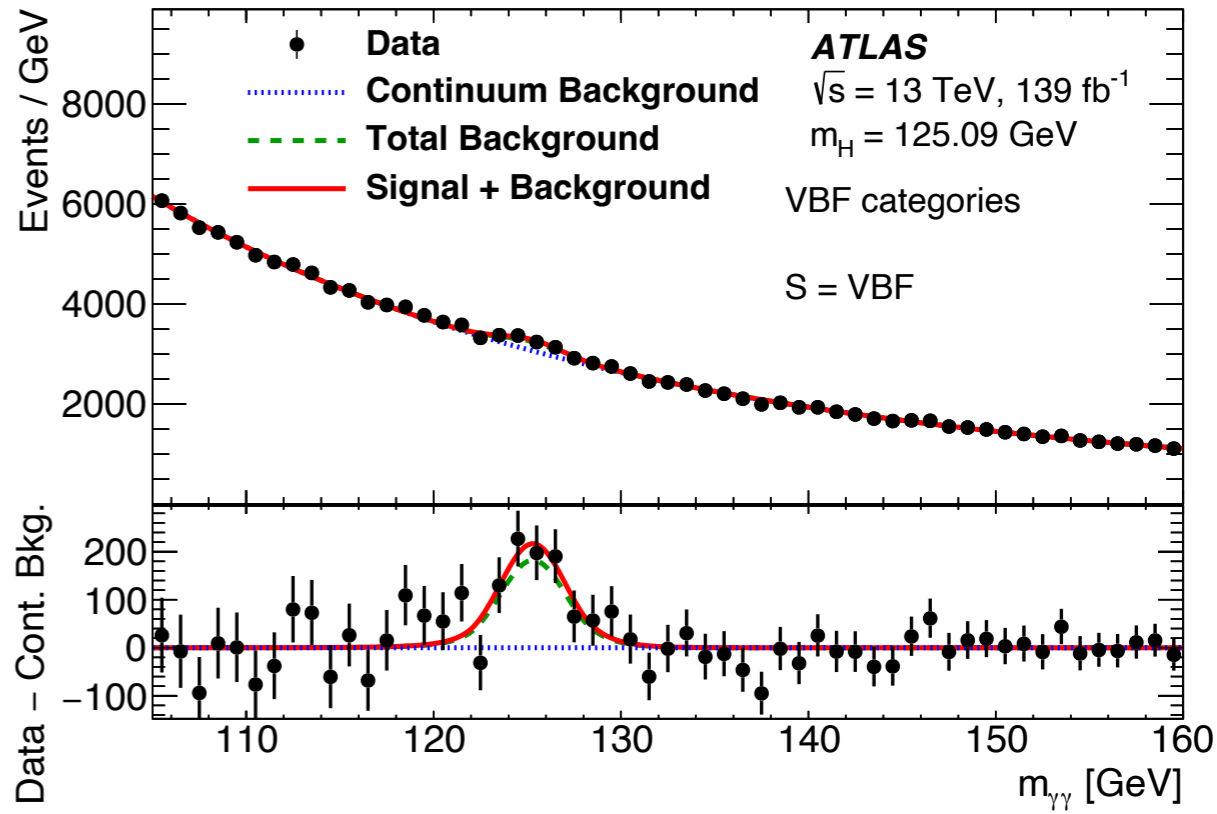
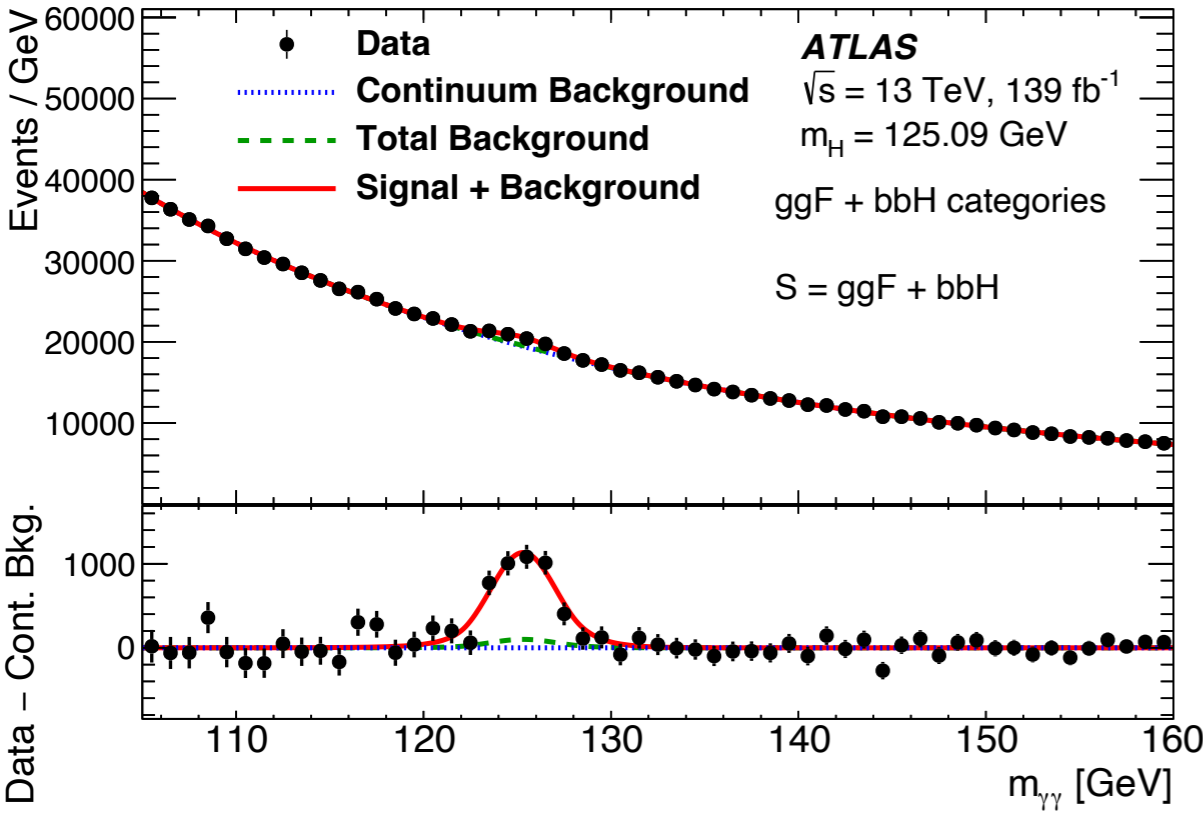
# BDT based region assignment



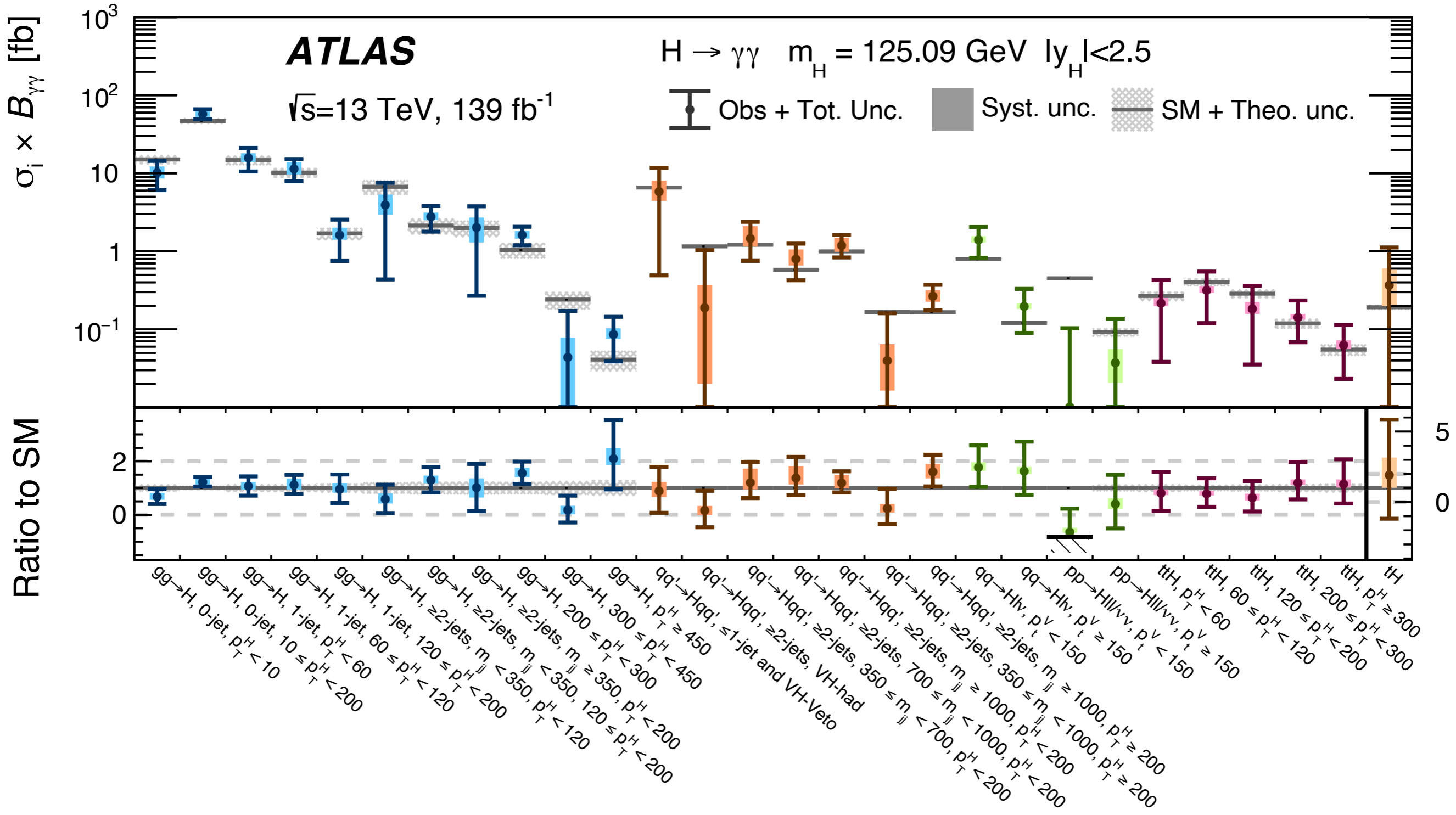
STXS classes	Variables
Individual STXS classes from $gg \rightarrow H$ $qq' \rightarrow Hqq'$ $qq \rightarrow H\ell\nu$ $pp \rightarrow H\ell\ell$ $pp \rightarrow H\nu\bar{\nu}$	All multiclass BDT variables, $p_T^{\gamma\gamma}$ projected to the thrust axis of the $\gamma\gamma$ system ( $p_T^{\gamma\gamma}$ ), $\Delta\eta_{\gamma\gamma}, \eta^{Z_{\text{epp}}} = \frac{\eta_{\gamma\gamma} - \eta_{jj}}{2}$ , $\phi_{\gamma\gamma}^* = \tan\left(\frac{\pi -  \Delta\phi_{\gamma\gamma} }{2}\right) \sqrt{1 - \tanh^2\left(\frac{\Delta\eta_{\gamma\gamma}}{2}\right)}$ , $\cos\theta_{\gamma\gamma}^* = \left  \frac{(E^{\gamma_1} + p_z^{\gamma_1}) \cdot (E^{\gamma_2} - p_z^{\gamma_2}) - (E^{\gamma_1} - p_z^{\gamma_1}) \cdot (E^{\gamma_2} + p_z^{\gamma_2})}{m_{\gamma\gamma} + \sqrt{m_{\gamma\gamma}^2 + (p_T^{\gamma\gamma})^2}} \right $ Number of electrons and muons.
all $t\bar{t}H$ and $tHW$ STXS classes combined	$p_T, \eta, \phi$ of $\gamma_1$ and $\gamma_2$ , $p_T, \eta, \phi$ and $b$ -tagging scores of the six highest- $p_T$ jets, $E_T^{\text{miss}}, E_T^{\text{miss}}$ significance, $E_T^{\text{miss}}$ azimuthal angle, Top reconstruction BDT scores of the top-quark candidates, $p_T, \eta, \phi$ of the two highest- $p_T$ leptons.
$tHqb$	$p_T^{\gamma\gamma} / m_{\gamma\gamma}, \eta_{\gamma\gamma}$ , $p_T$ , invariant mass, BDT score and $\Delta R(W, b)$ of $t_1$ , $p_T, \eta$ of $t_2$ , $p_T, \eta$ of $j_F$ , Angular variables: $\Delta\eta_{\gamma\gamma t_1}, \Delta\theta_{\gamma\gamma t_2}, \Delta\theta_{t_1 j_F}, \Delta\theta_{t_2 j_F}, \Delta\theta_{\gamma\gamma j_F}$ Invariant mass variables: $m_{\gamma\gamma j_F}, m_{t_1 j_F}, m_{t_2 j_F}, m_{\gamma\gamma t_1}$ Number of jets with $p_T > 25$ GeV, Number of $b$ -jets with $p_T > 25$ GeV*, Number of leptons*, $E_T^{\text{miss}}$ significance*

$\eta_{\gamma_1}, \eta_{\gamma_2}, p_T^{\gamma\gamma}, y_{\gamma\gamma}$ , $p_{T,jj}^\dagger, m_{jj}, \text{ and } \Delta y, \Delta\phi, \Delta\eta \text{ between } j_1 \text{ and } j_2$ , $p_{T,\gamma\gamma j_1}, m_{\gamma\gamma j_1}, p_{T,\gamma\gamma j_1}^\dagger, m_{\gamma\gamma j_1}^\dagger$ $\Delta y, \Delta\phi$ between the $\gamma\gamma$ and $jj$ systems, minimum $\Delta R$ between jets and photons, invariant mass of the system comprising all jets in the event, dilepton $p_T$ , di- $e$ or di- $\mu$ invariant mass (leptons are required to be oppositely charged), $E_T^{\text{miss}}, p_T$ and transverse mass of the lepton + $E_T^{\text{miss}}$ system, $p_T, \eta, \phi$ of top-quark candidates, $m_{t_1 t_2}$ Number of jets $^\dagger$ , of central jets ( $ \eta  < 2.5$ ) $^\dagger$ , of $b$ -jets $^\dagger$ and of leptons, $p_T$ of the highest- $p_T$ jet, scalar sum of the $p_T$ of all jets, scalar sum of the transverse energies of all particles ( $\sum E_T$ ), $E_T^{\text{miss}}$ significance, $\left  E_T^{\text{miss}} - E_T^{\text{miss}}(\text{primary vertex with the highest } \sum p_{T,\text{track}}^2) \right  > 30 \text{ GeV}$ Top reconstruction BDT of the top-quark candidates, $\Delta R(W, b)$ of $t_2$ , $\eta_{j_F}, m_{\gamma\gamma j_F}$ Average number of interactions per bunch crossing.
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# Obtained results : diphoton mass

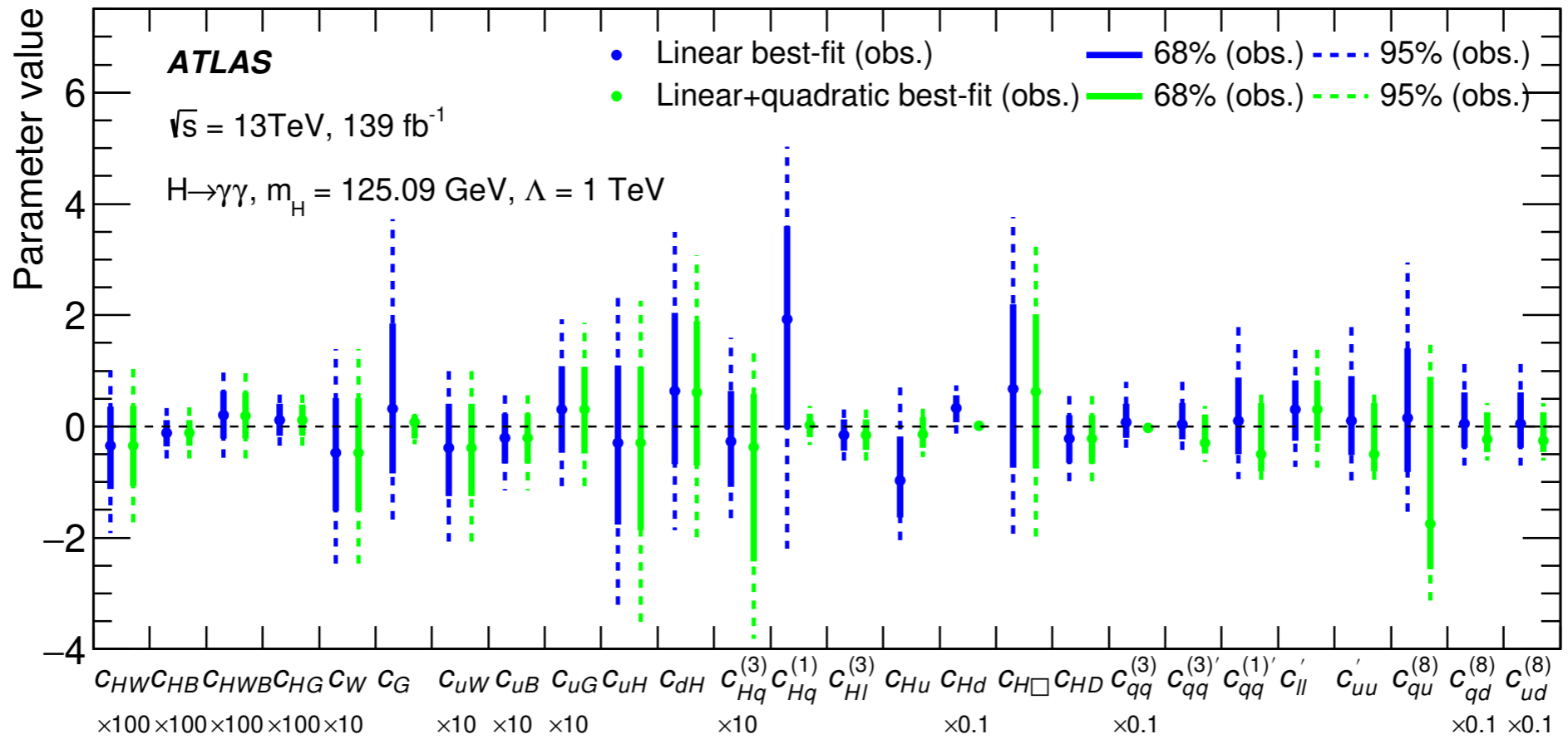
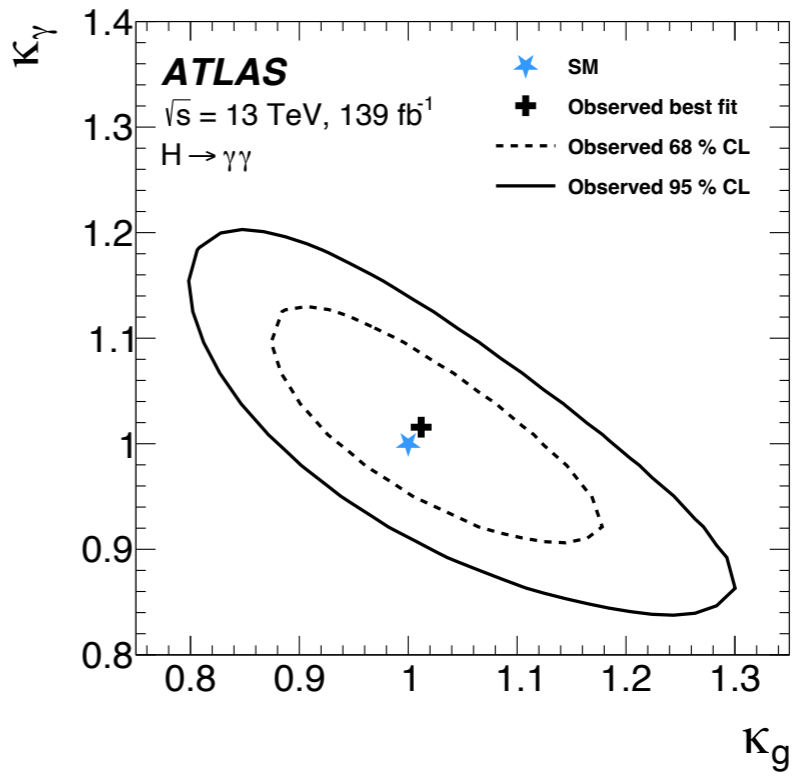
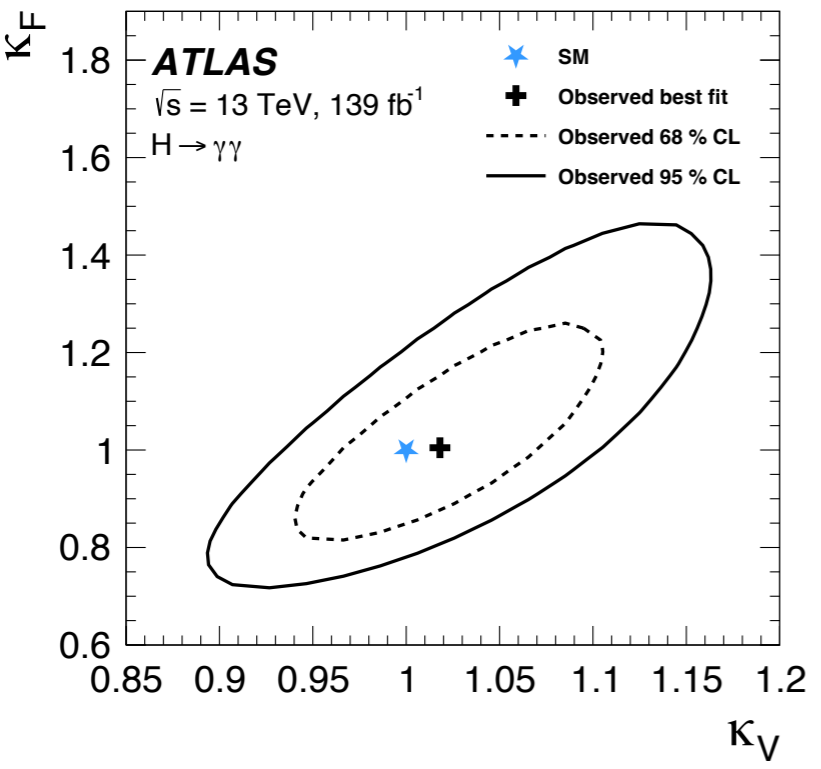


# Obtained results : signal strengths



$$\mu = 1.04^{+0.10}_{-0.09} = 1.04 \pm 0.06 \text{ (stat.)}^{+0.06}_{-0.05} \text{ (theory syst.)}^{+0.05}_{-0.04} \text{ (exp. syst.)}$$

# Obtained results : SMEFT & $\kappa$ parameters



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- ▶  $H$  width from off-shell  $ZZ^*$
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# H width from off shell HZZ

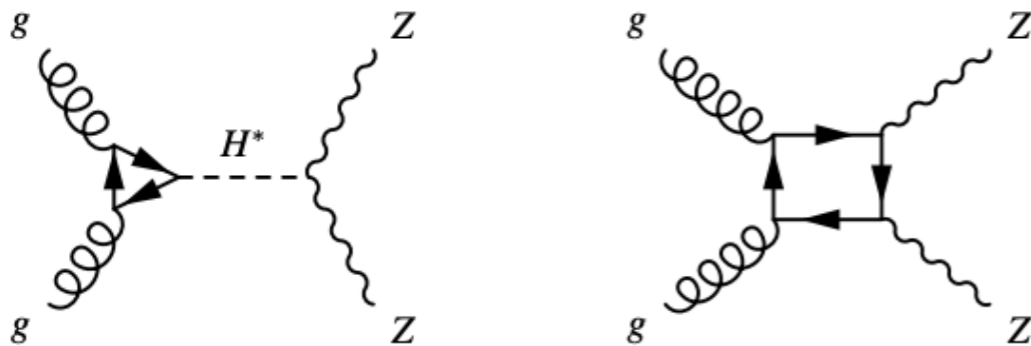
arXiv 2304.01532

Higgs width can be measured by taking the ratio of off-shell to on-shell production x-section

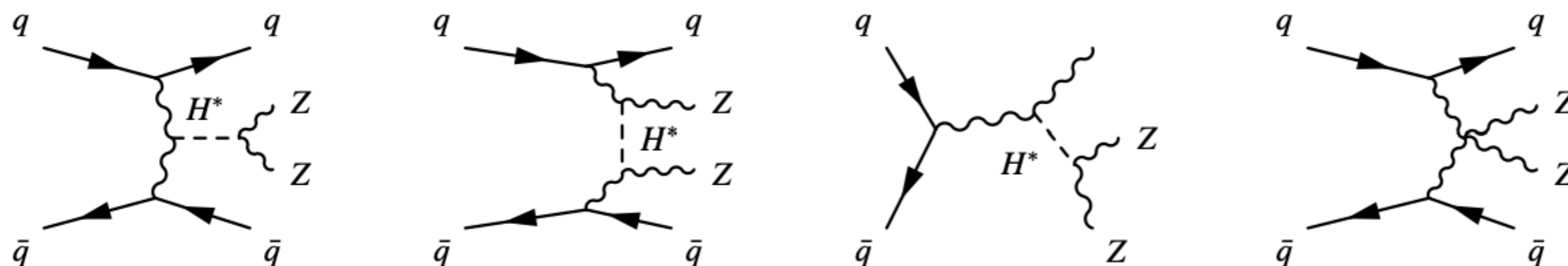
$$\sigma_{gg \rightarrow H \rightarrow ZZ}^{\text{on-shell}} \sim \frac{g_{ggH}^2 g_{HZZ}^2}{m_H \Gamma_H}$$

$$\sigma_{gg \rightarrow H \rightarrow ZZ}^{\text{off-shell}} \sim \frac{g_{ggH}^2 g_{HZZ}^2}{m_{ZZ}^2}$$

Off shell regime  $\sqrt{s} \gg M_Z, m_t$



ggF signal & bkg.

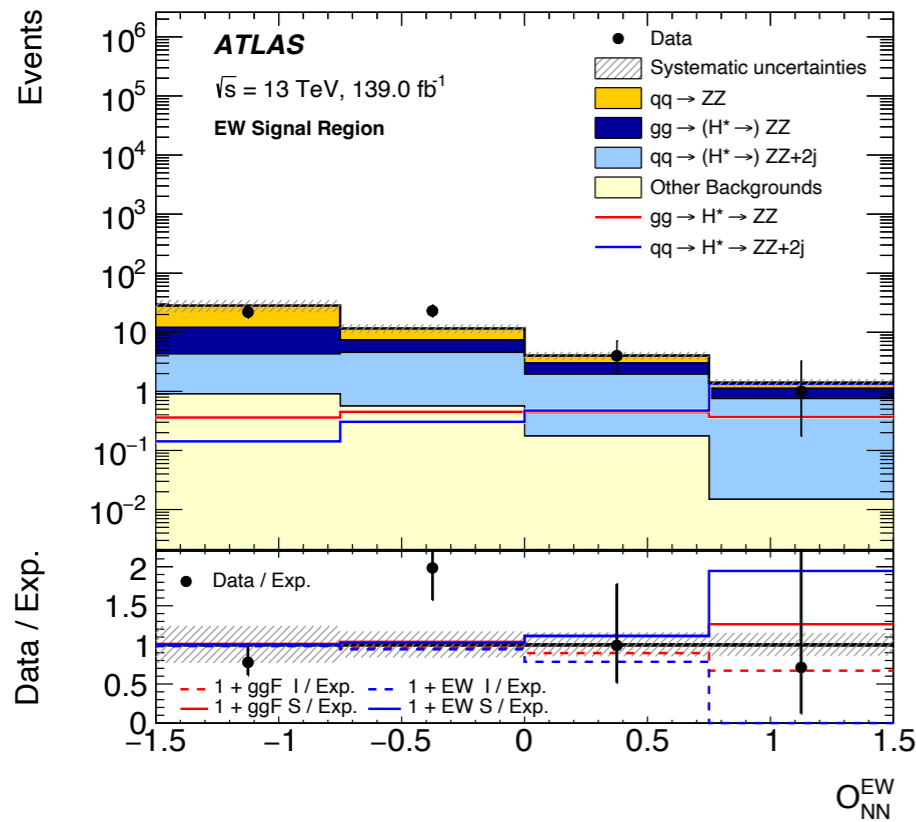
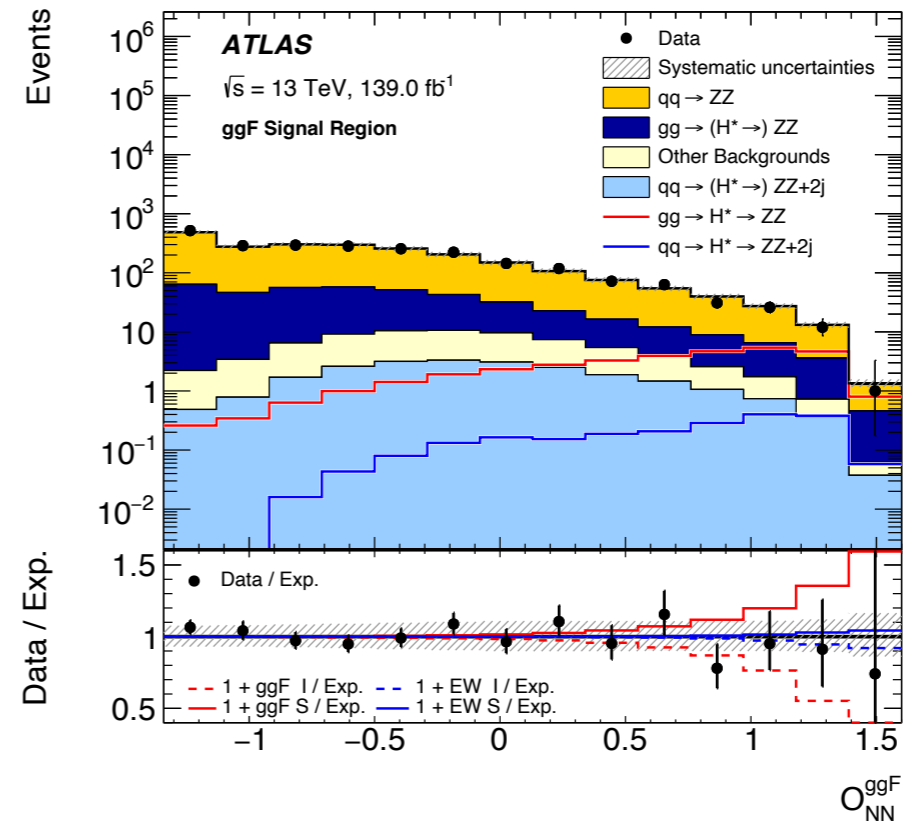
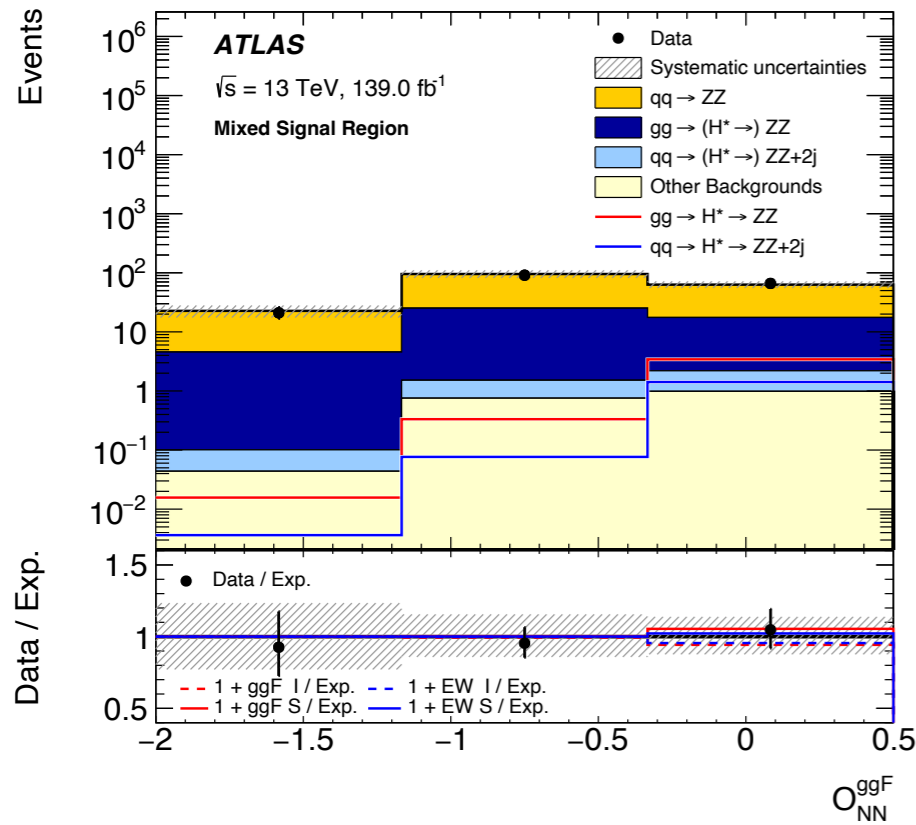


VBF signal & bkg.



# ZZ to 4l analysis

arXiv 2304.01532



$$m_{4l} > 220 \text{ GeV}$$

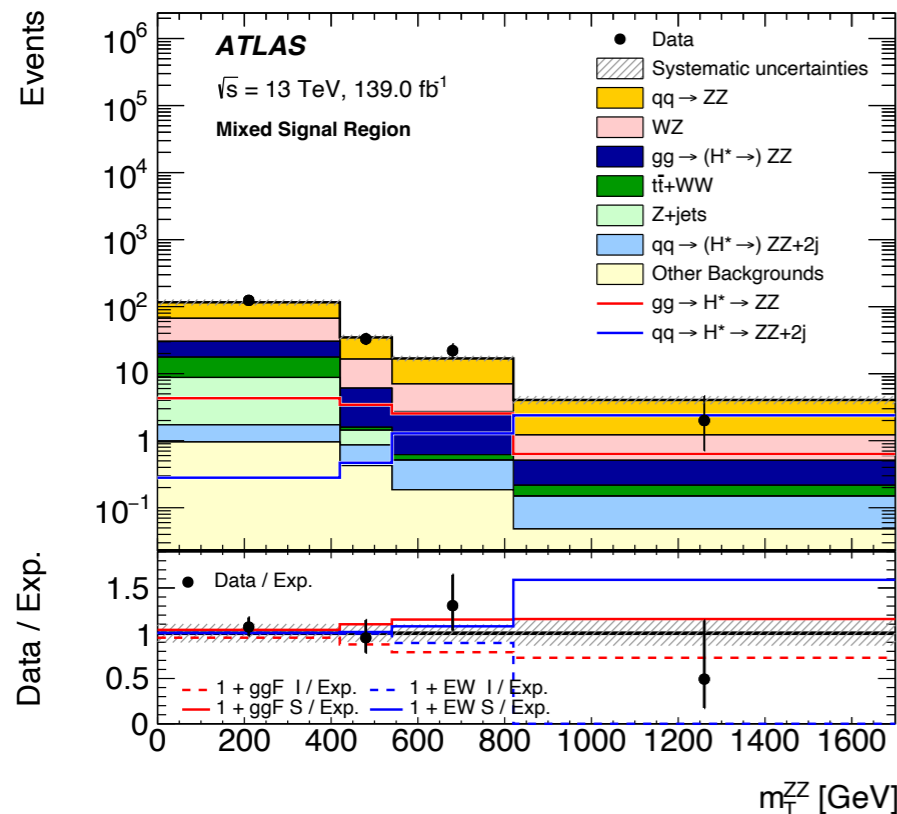
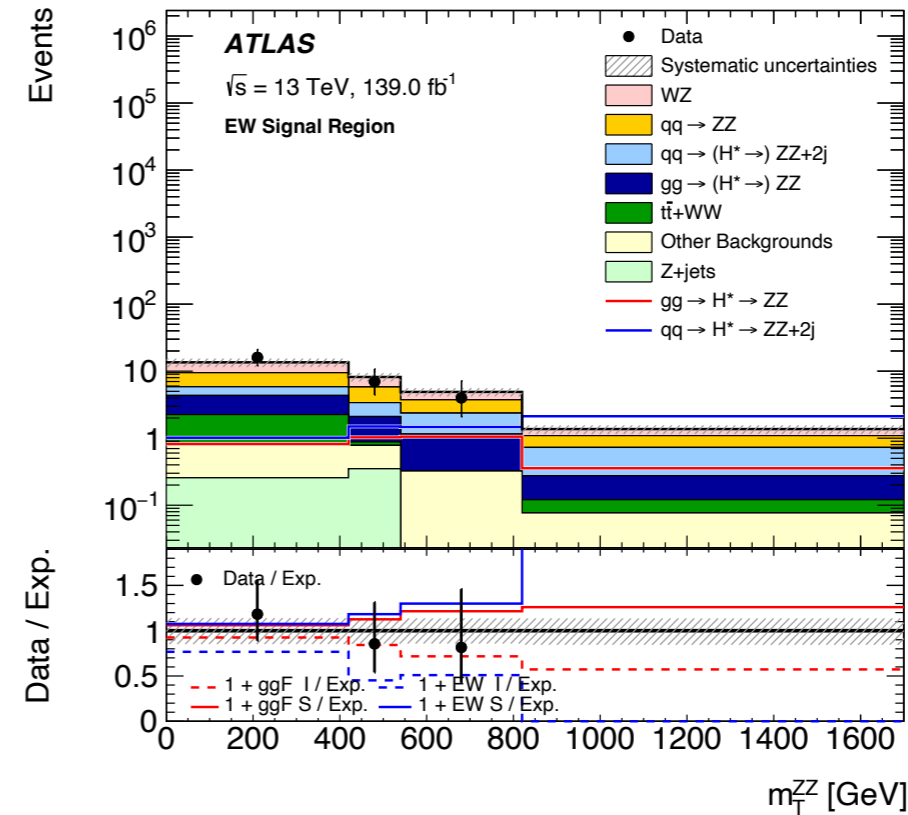
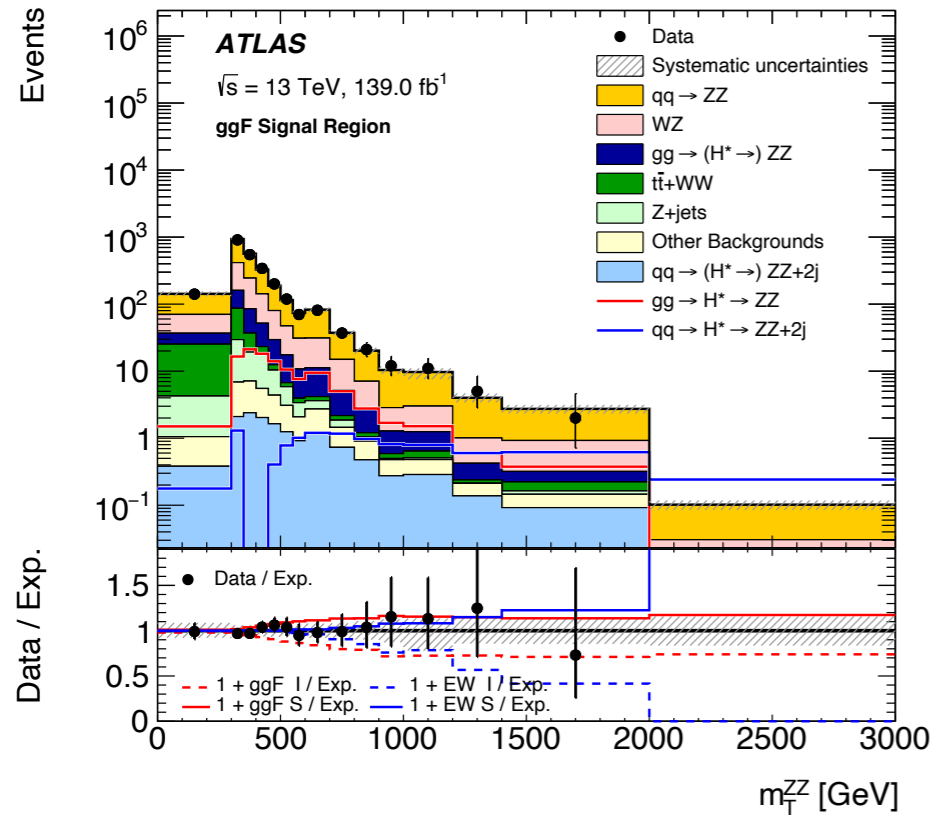
Main bkg :  $q\bar{q} \rightarrow ZZ$ ;  $180 < m_{4l} < 220 \text{ GeV}$  0, 1, > 2J

$$O_{NN} = \log_{10} \left( \frac{P_S}{P_B + P_{NI}} \right)$$

Two different NN for ggF & EW SR

# H to ZZ to 2l + 2 nu

arXiv 2304.01532



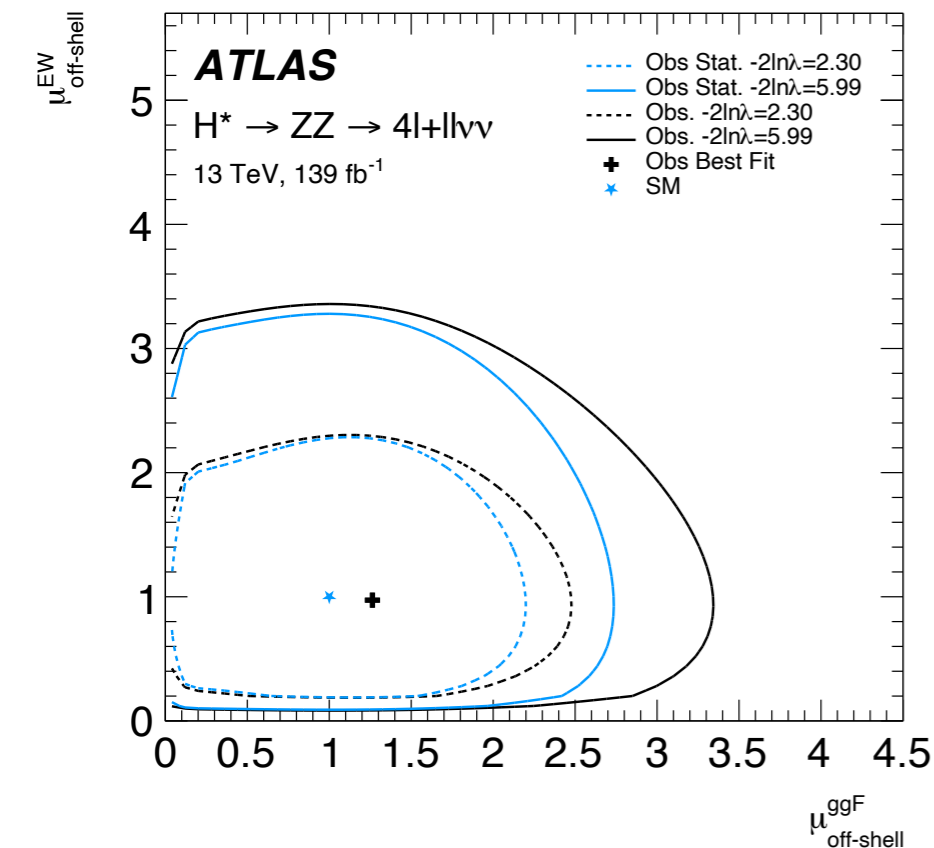
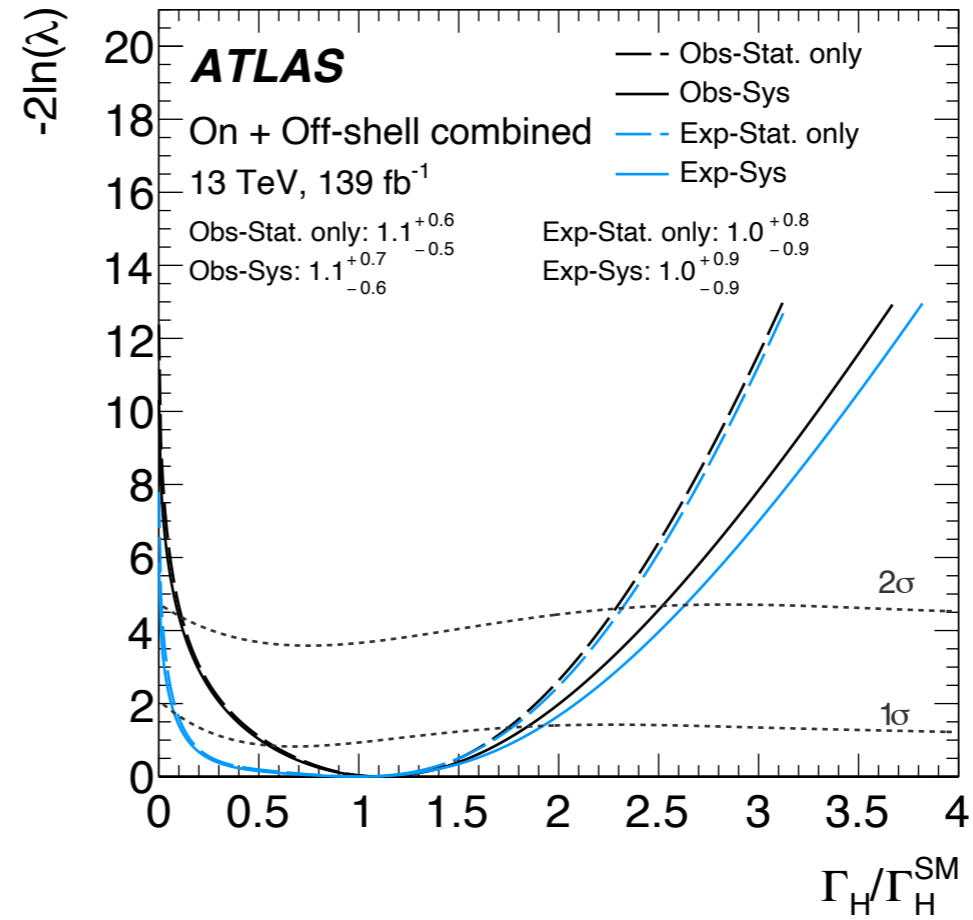
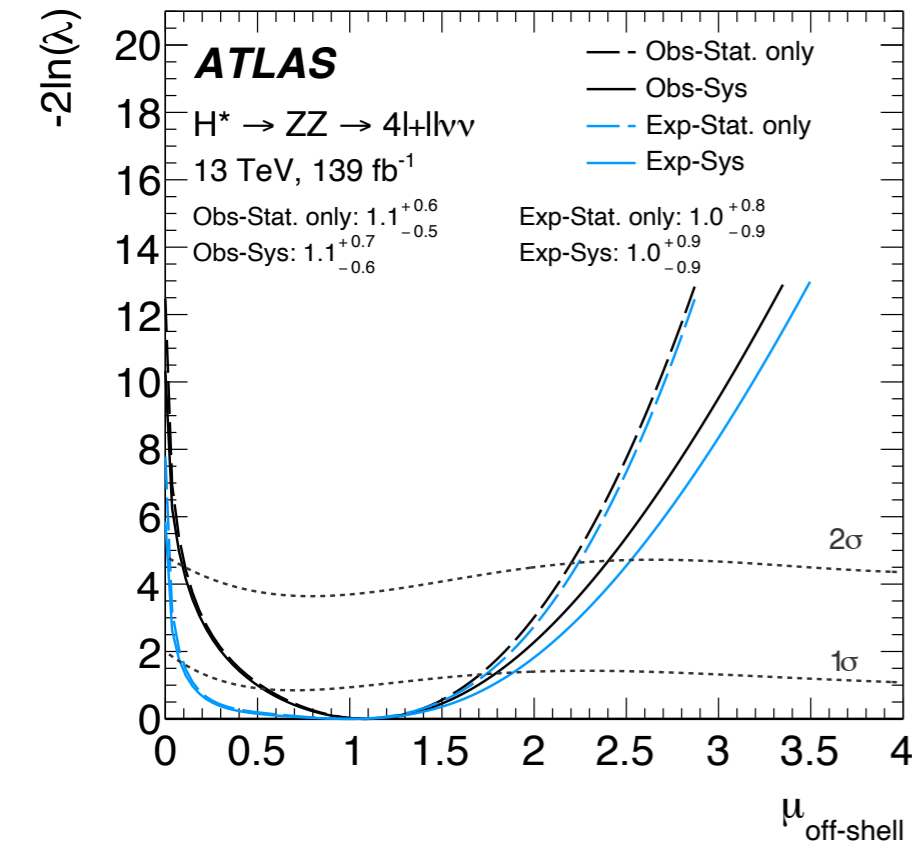
$$m_T^{ZZ} \equiv \sqrt{\left[ \sqrt{m_Z^2 + (p_T^{\ell\ell})^2} + \sqrt{m_Z^2 + (E_T^{\text{miss}})^2} \right]^2 - \left| \vec{p}_T^{\ell\ell} + \vec{E}_T^{\text{miss}} \right|^2}$$

**Reducible bkg's (ttbar, s-top, qq to WW) are killed by  $76 < m_{ll} < 106 \text{ GeV}$ .**

**Scale uncertainty associated with  $q\bar{q} \rightarrow ZZ$  is one of the largest sources of uncertainty, can be upto 40%.**

# H width from off shell HZZ

arXiv 2304.01532



$$\mu_{\text{off-shell}}^{\text{ggF}} = \kappa_{g,\text{off-shell}}^2 \kappa_{V,\text{off-shell}}^2$$

$$\mu_{\text{off-shell}}^{\text{EW}} = \kappa_{V,\text{off-shell}}^4$$

Combination of 4l and 2l+2mu channels are presented.

Width found to be  $4.5^{+3.3}_{-2.5}$  MeV.

# Outline

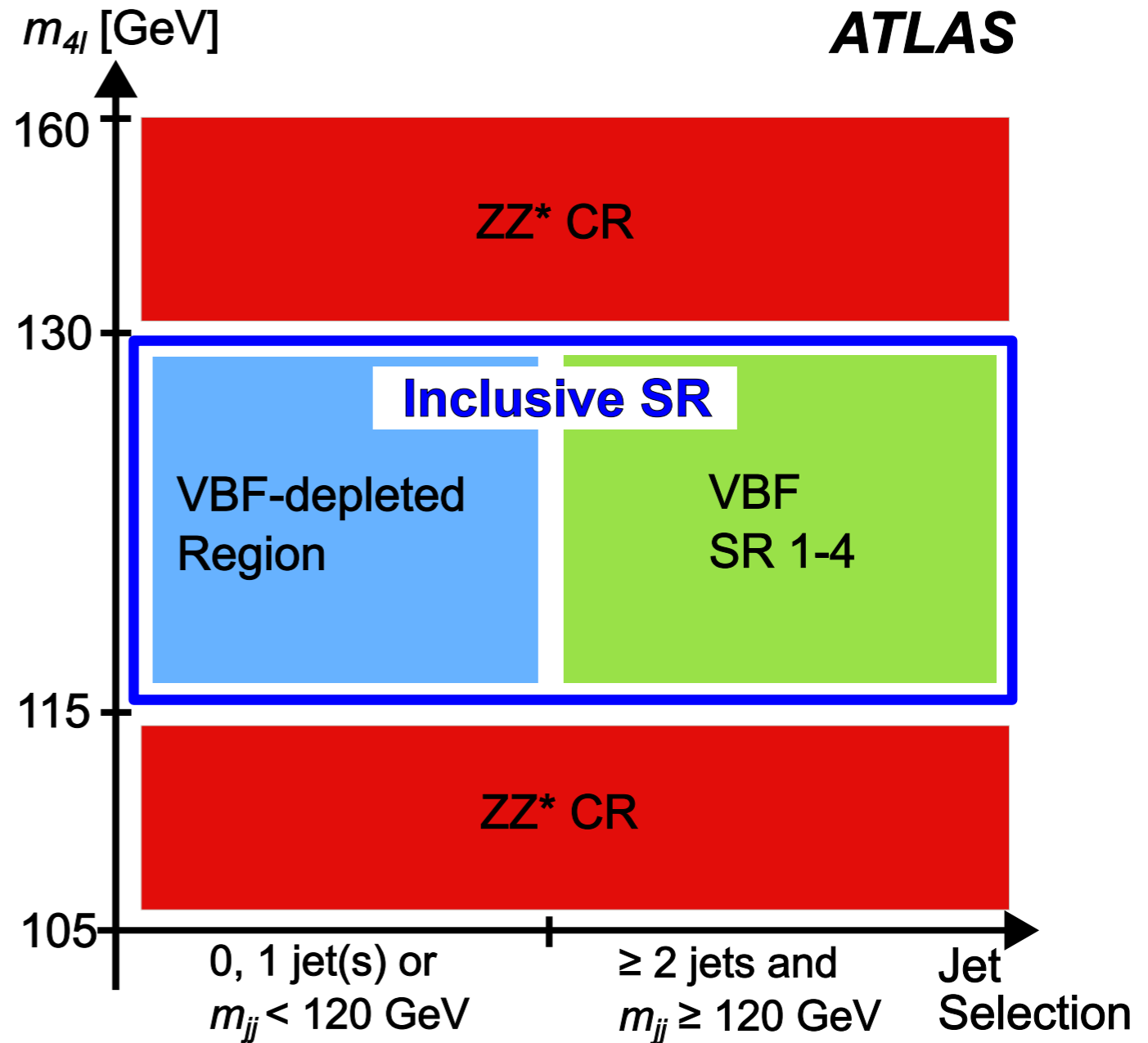
- Introduction
- $H \rightarrow \gamma\gamma$  STXS measurement
- $H$  width from off-shell  $ZZ^*$
- ▶ Higgs CP from  $H \rightarrow ZZ^* \rightarrow 4l$
- Search for HH in  $4b$  final state
- Charm Yukawa in VH mode
- Summary

# Test of CP invariance from $H \rightarrow ZZ^* \rightarrow 4 \text{ lep}$

$$\mathcal{L}_{\text{SMEFT}} = \mathcal{L}_{\text{SM}} + \sum_i \frac{c_i}{\Lambda^2} O_i^{(6)}$$

Operator	Structure	Coupling
Warsaw Basis		
$O_{\Phi\tilde{W}}$	$\Phi^\dagger \Phi \tilde{W}_{\mu\nu}^I W^{\mu\nu I}$	$c_{H\tilde{W}}$
$O_{\Phi\tilde{W}B}$	$\Phi^\dagger \tau^I \Phi \tilde{W}_{\mu\nu}^I B^{\mu\nu}$	$c_{H\tilde{W}B}$
$O_{\Phi\tilde{B}}$	$\Phi^\dagger \Phi \tilde{B}_{\mu\nu} B^{\mu\nu}$	$c_{H\tilde{B}}$
Higgs Basis		
$O_{hZ\tilde{Z}}$	$h Z_{\mu\nu} \tilde{Z}^{\mu\nu}$	$\tilde{c}_{ZZ}$
$O_{hZ\tilde{A}}$	$h Z_{\mu\nu} \tilde{A}^{\mu\nu}$	$\tilde{c}_{Z\gamma}$
$O_{hA\tilde{A}}$	$h A_{\mu\nu} \tilde{A}^{\mu\nu}$	$\tilde{c}_{\gamma\gamma}$

SMEFT CP-odd dim-6 operators relevant for  $H \rightarrow ZZ \rightarrow 4l$  channel



Regions used for direct BSM coupling measurement

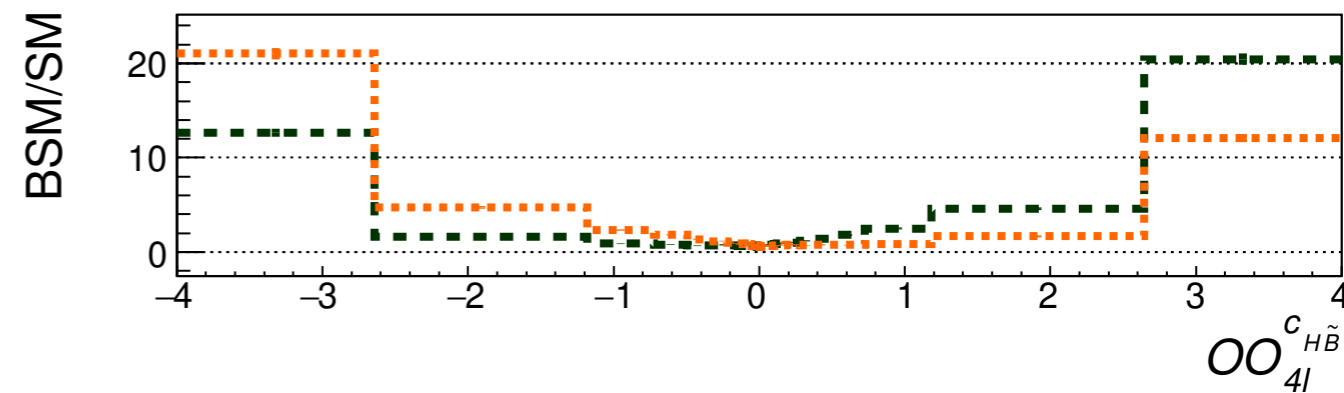
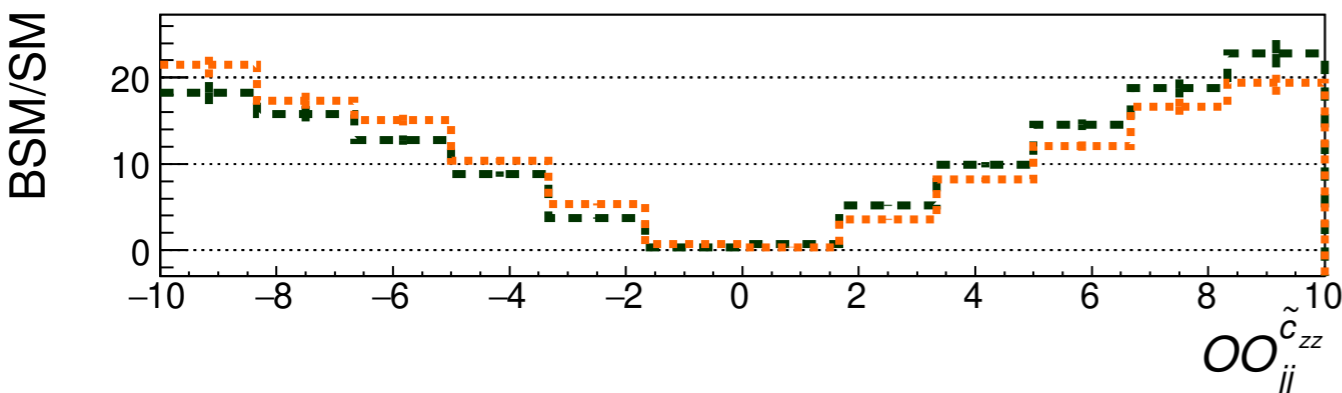
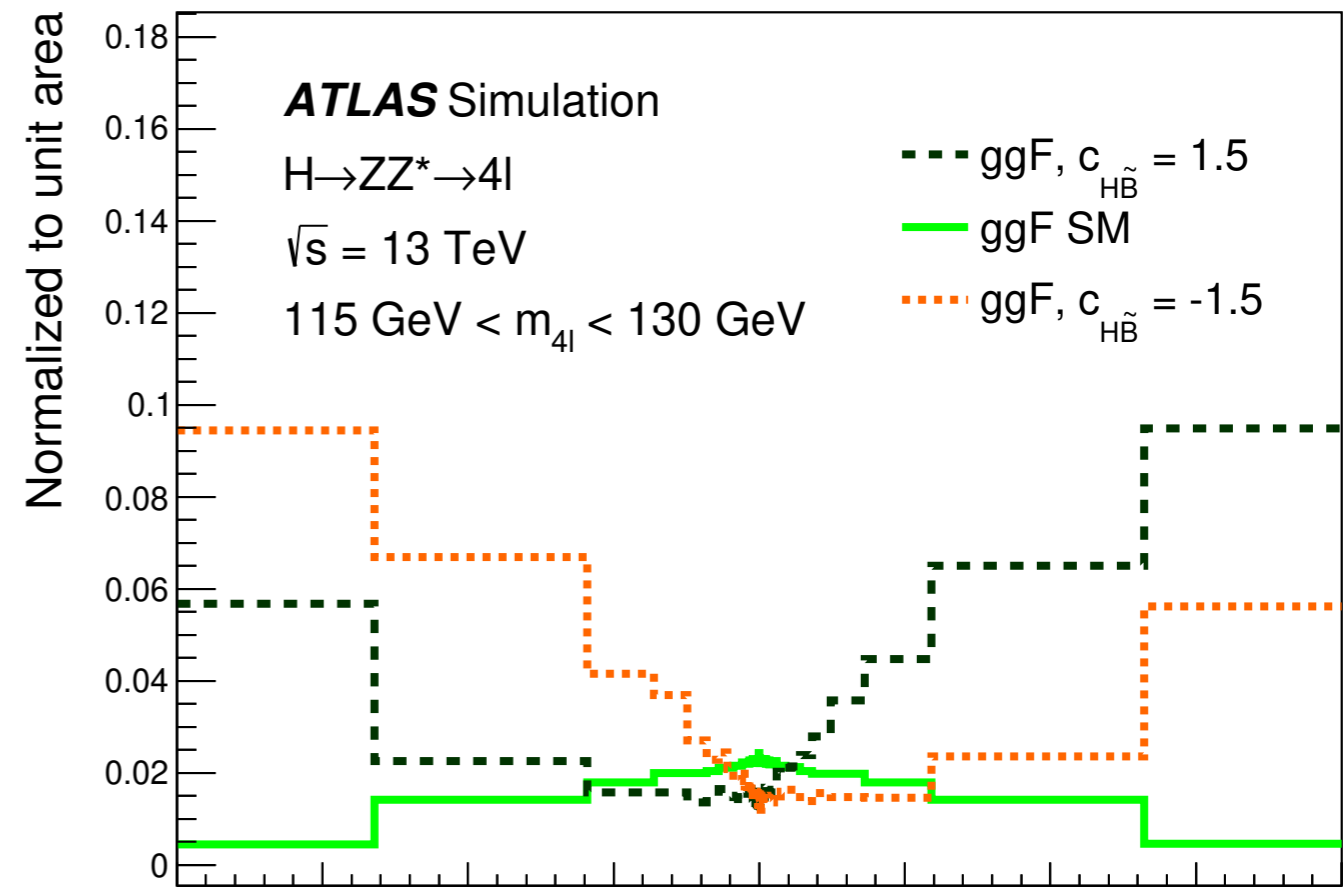
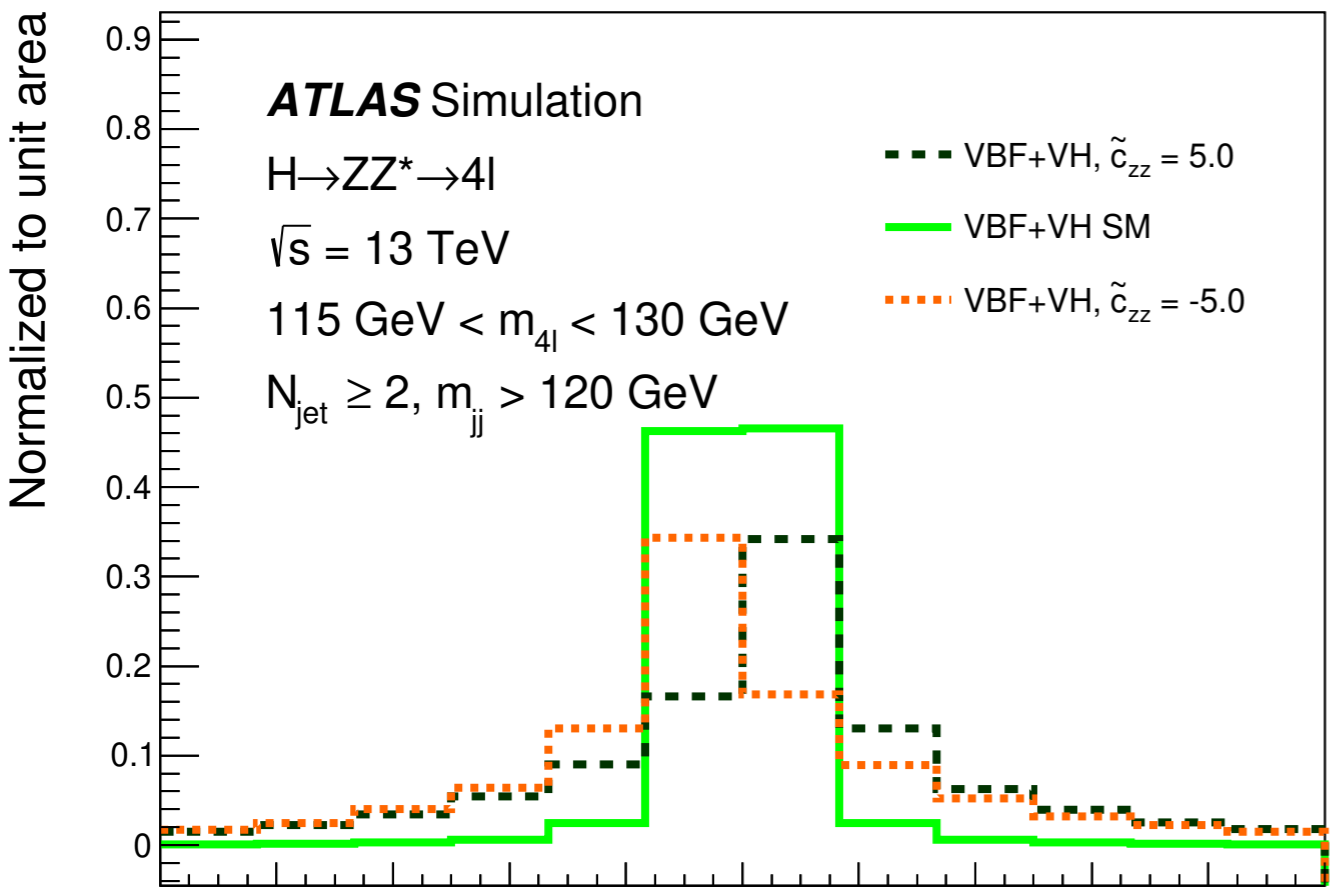
arXiv 2304.09612

Matrix element based optimal observables are used to constrain CP-odd couplings in SMEFT

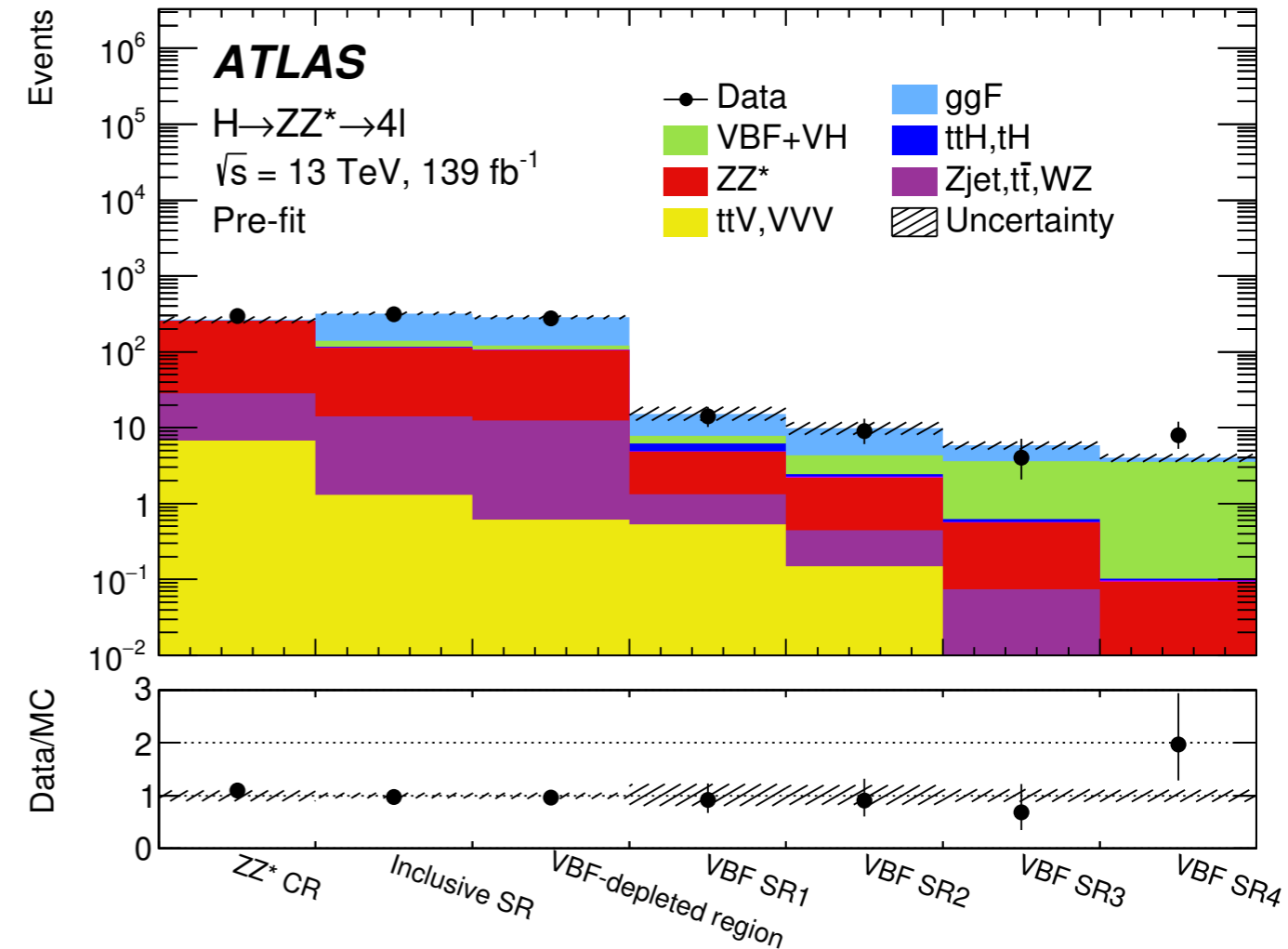
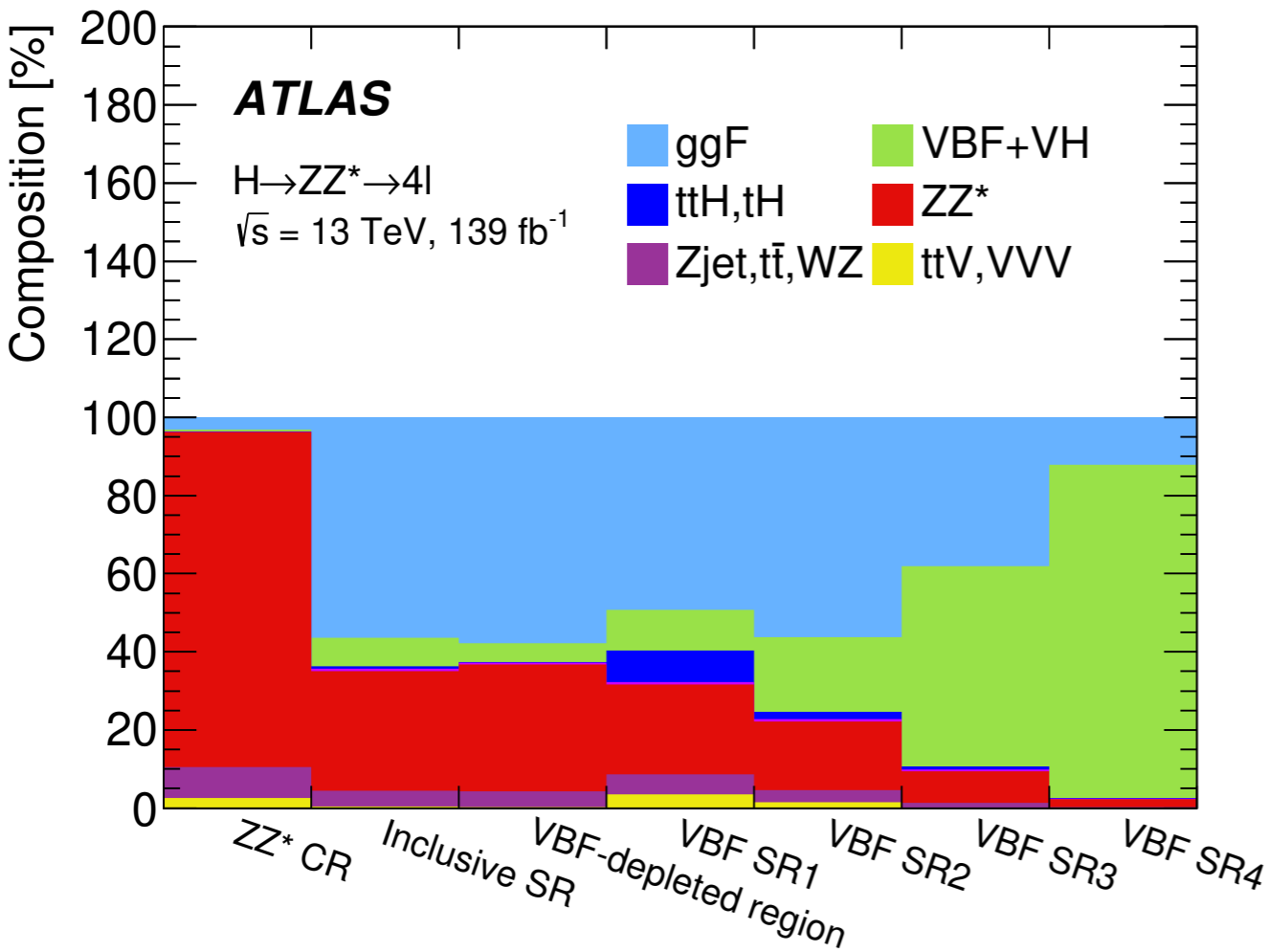
# Test of CP invariance from $H \rightarrow ZZ^* \rightarrow 4 \text{ lep}$

$$\begin{aligned}
 |\mathcal{M}|^2 &= \left| \mathcal{M}_{\text{SM}} + \sum_i \frac{c_i}{\Lambda^2} \mathcal{M}_{\text{BSM},i} \right|^2 \\
 &= |\mathcal{M}_{\text{SM}}|^2 + 2 \sum_i \frac{c_i}{\Lambda^2} \Re(\mathcal{M}_{\text{SM}}^* \mathcal{M}_{\text{BSM},i}) + \sum_i \sum_j \frac{c_i c_j}{\Lambda^4} \Re(\mathcal{M}_{\text{BSM},i}^* \mathcal{M}_{\text{BSM},j})
 \end{aligned}$$

$$\mathcal{O}\mathcal{O} = \frac{2\Re(\mathcal{M}_{\text{SM}}^* \mathcal{M}_{\text{BSM}})}{|\mathcal{M}_{\text{SM}}|^2}$$



# The expected & observed event yields

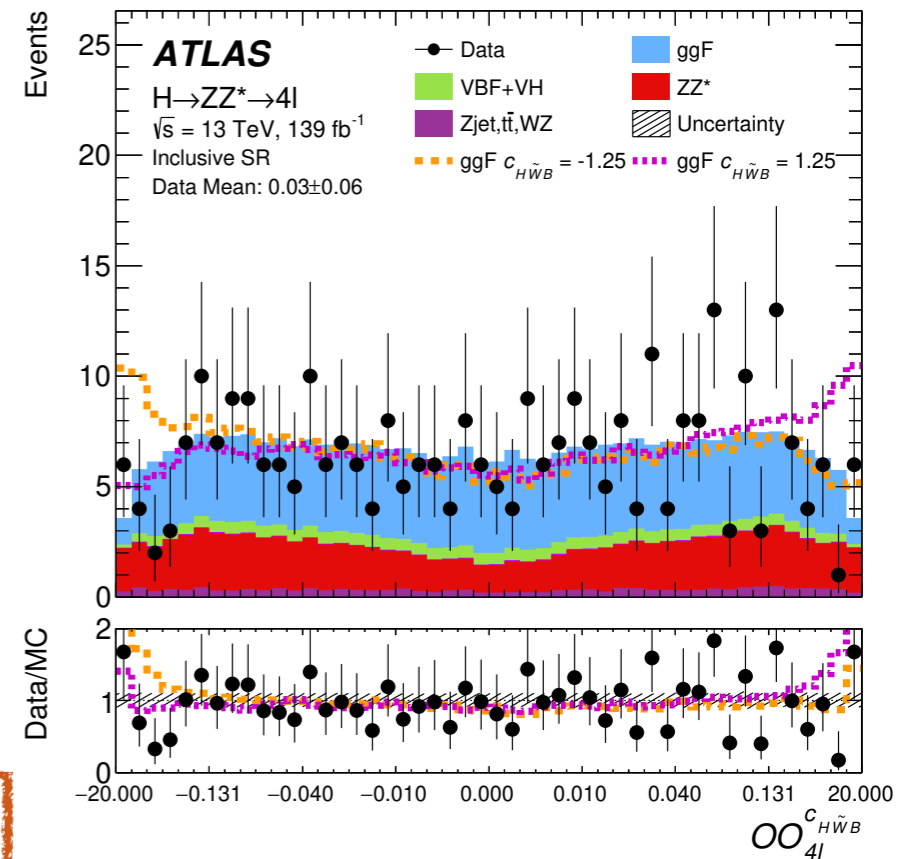
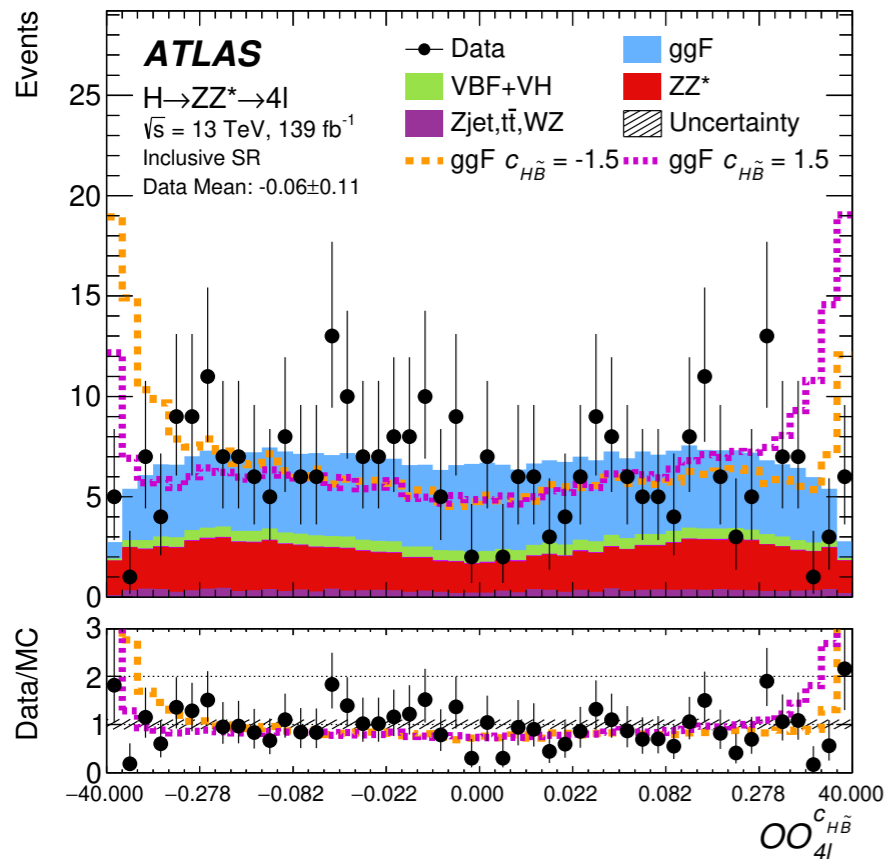


Rel. composition of the predicted event yields

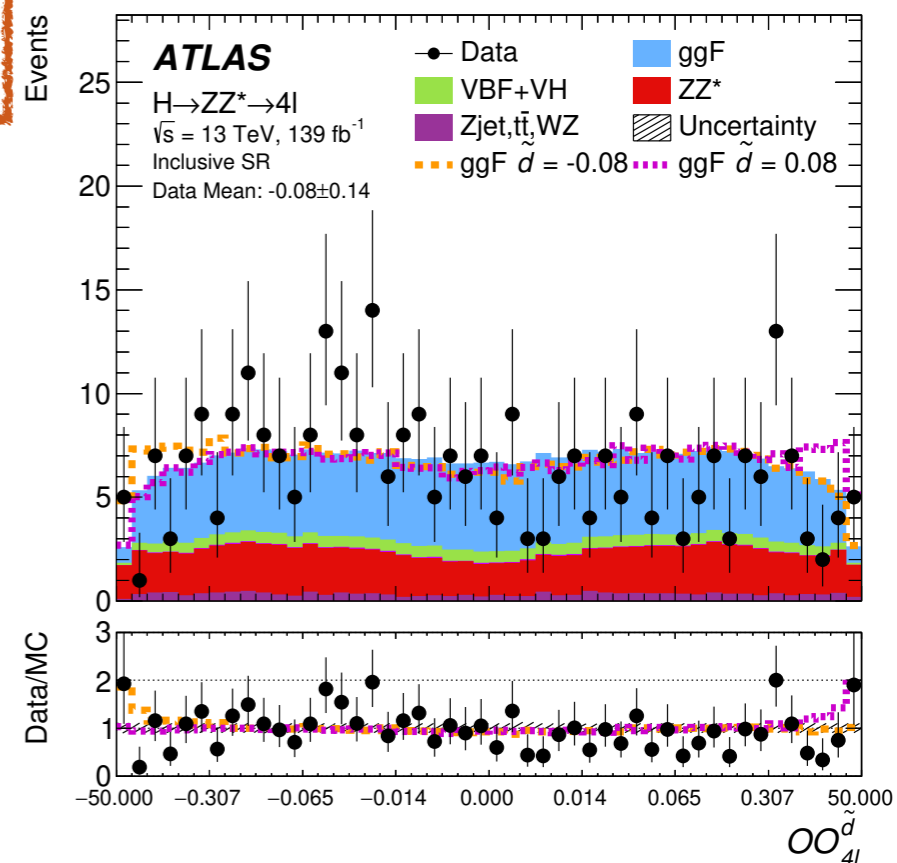
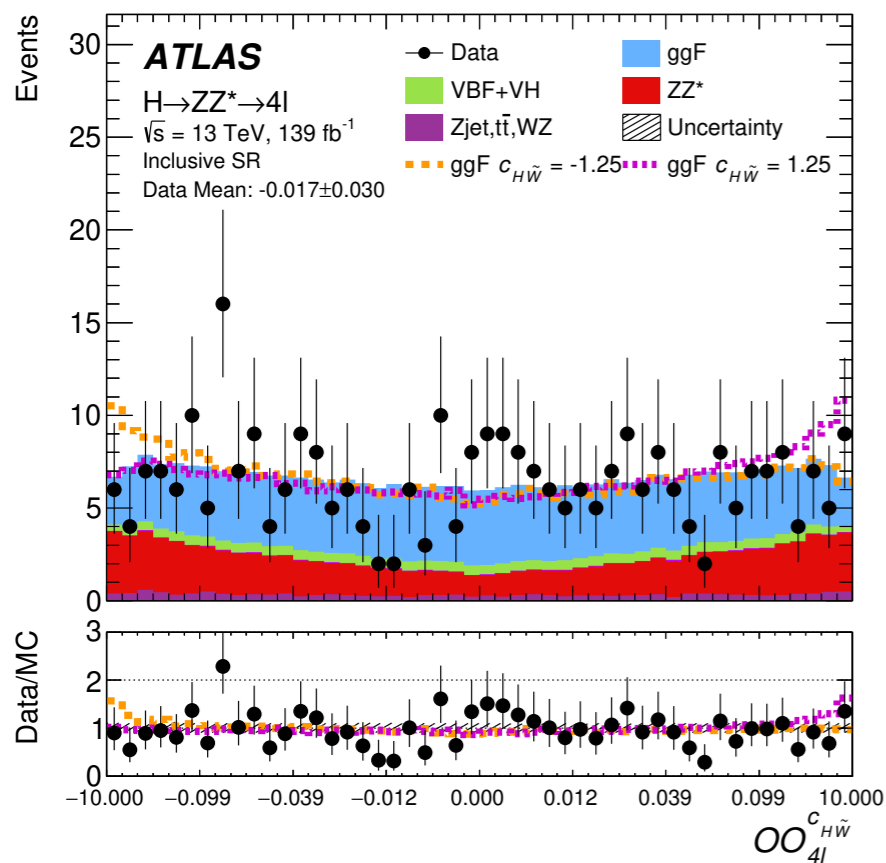
The number of observed events and expected contributions in each event category.

- ✓ The dominant background is non-resonant  $ZZ^*$  production  $\sim 30\%$  :  
 MC based estimation + data-driven normalization.
- ✓ Z+jets, t-tbar, WZ are reducible backgrounds : estimated in data-driven way.
- ✓ The tri-boson big (WWZ, WZZ and ZZZ) and ttX are estimated from simulation.

# The expected & observed event yields

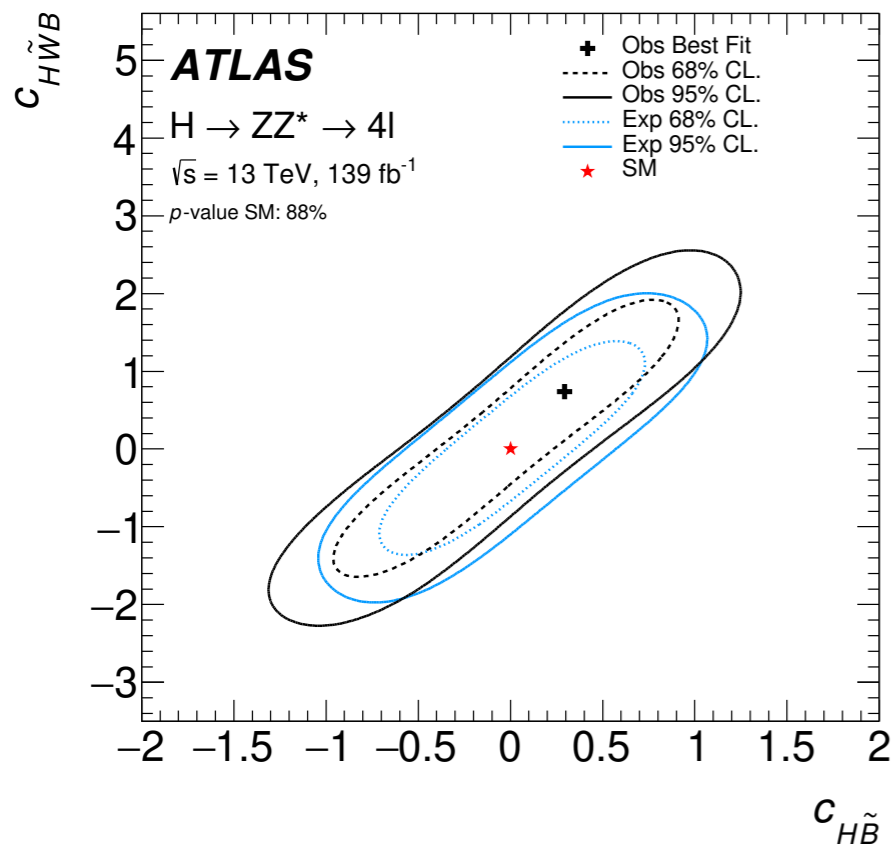


No statistically significant deviation from 0 is observed in any of the distributions

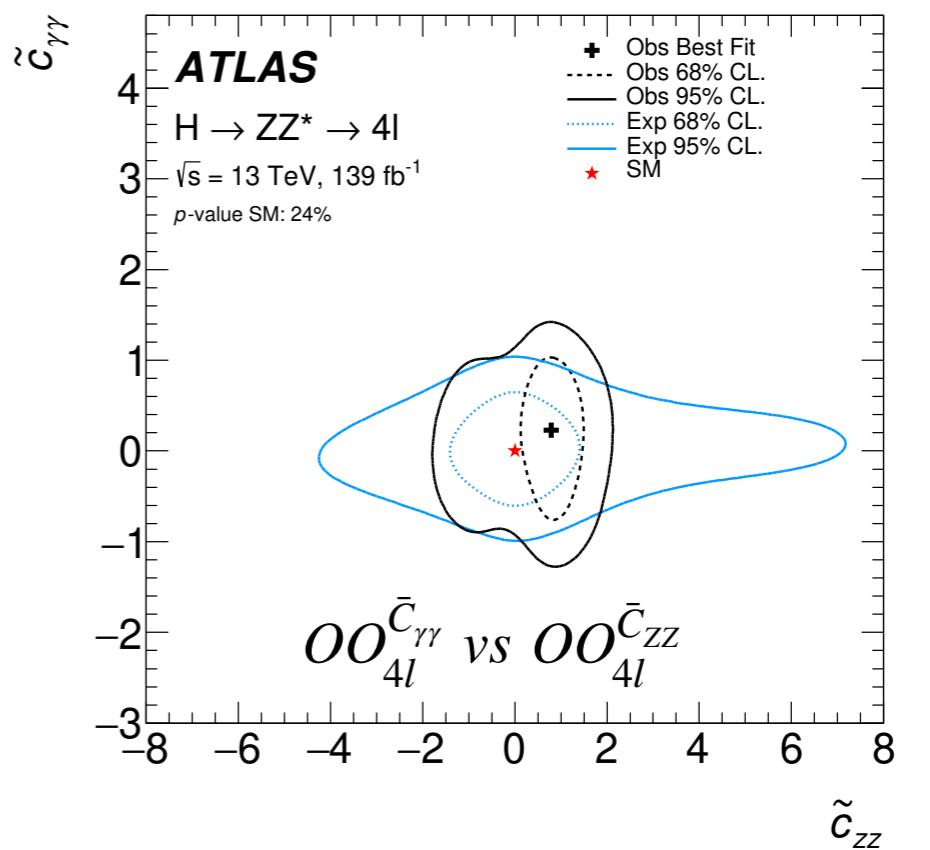
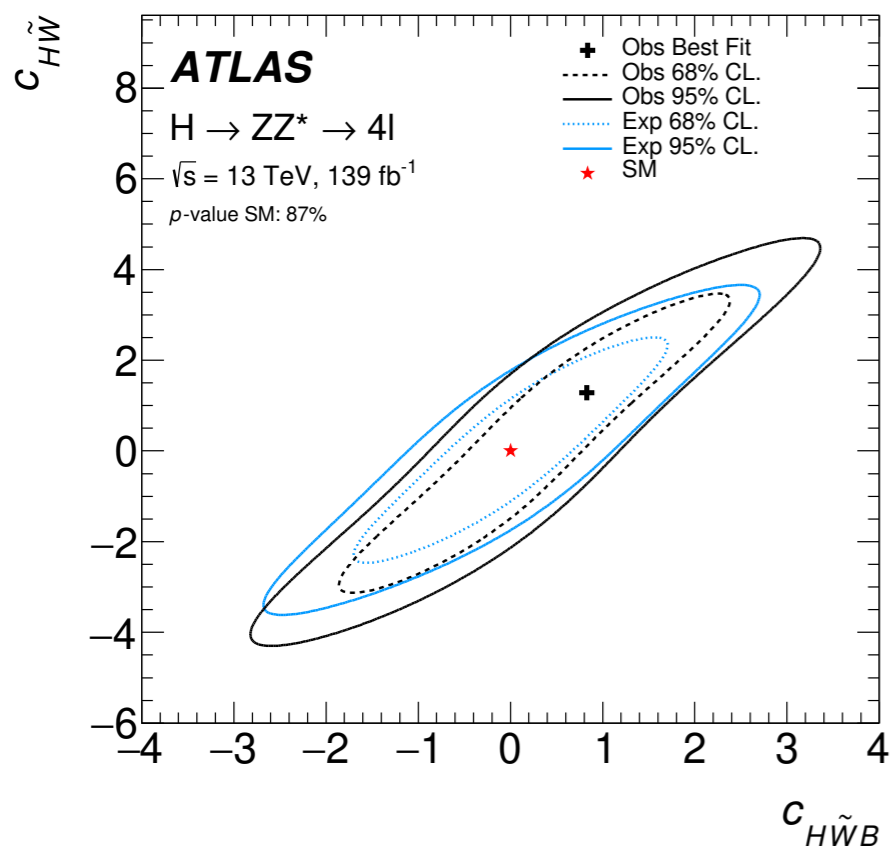
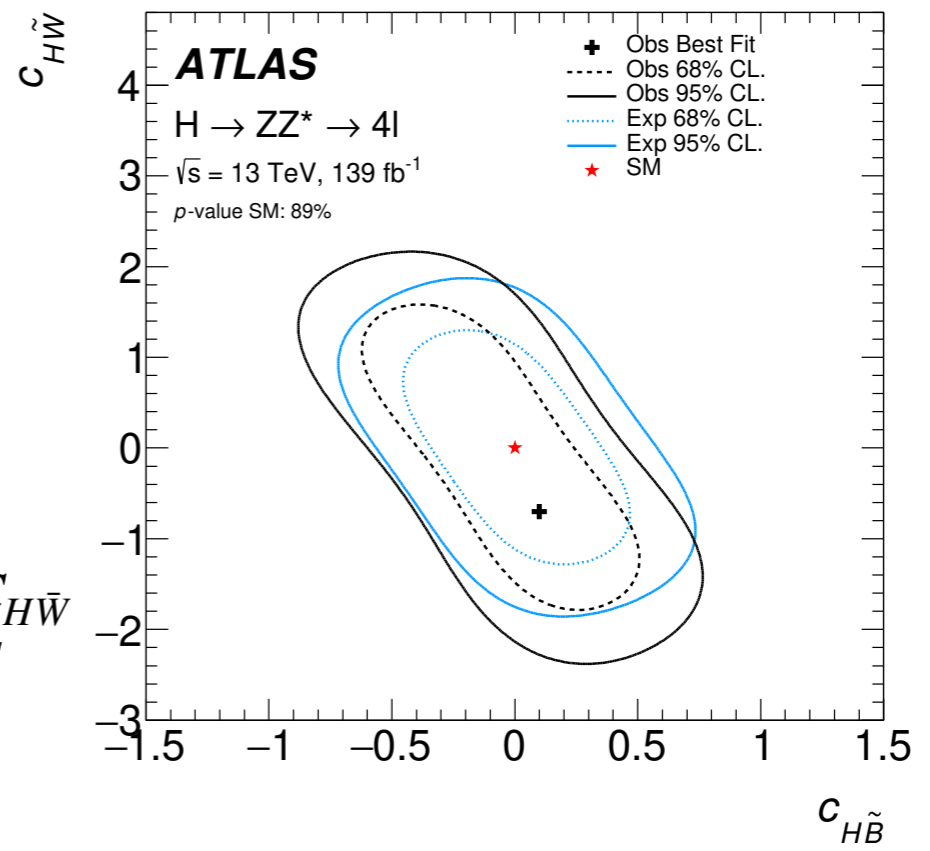




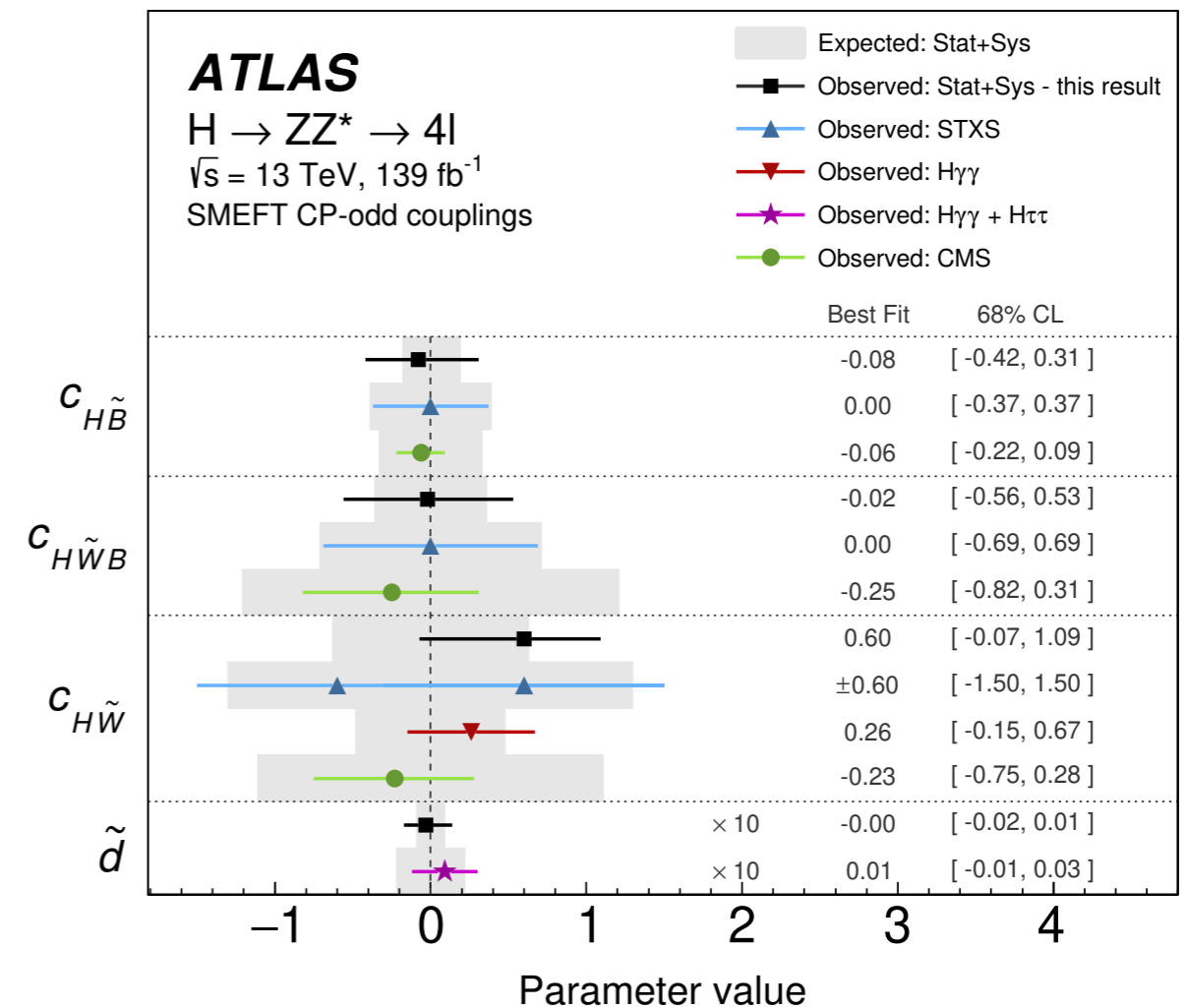
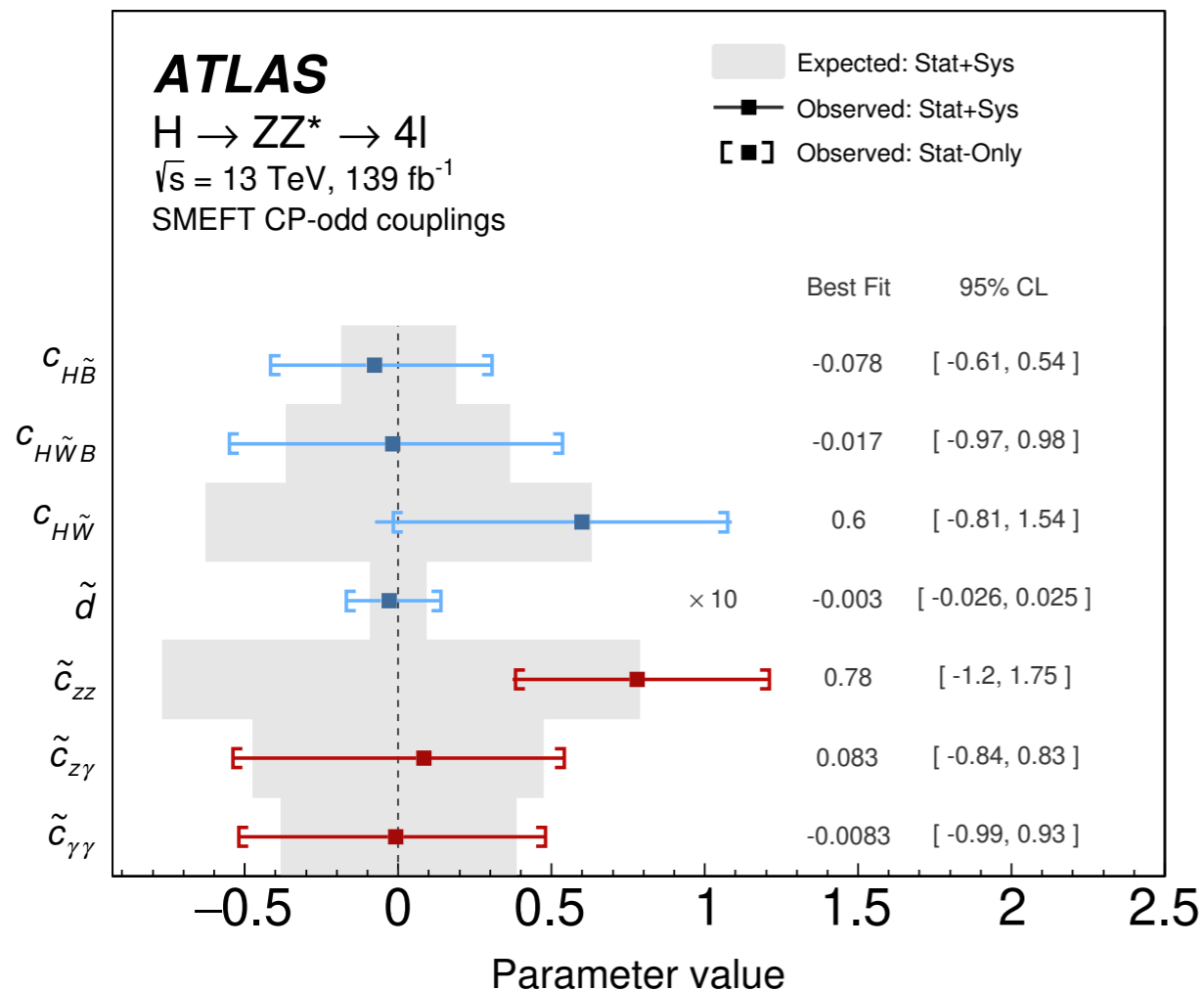
# CL contours



$OO_{4l}^{C_{H\tilde{B}}} \text{ vs } OO_{4l}^{C_{H\tilde{W}}}$



# Extracted effective couplings



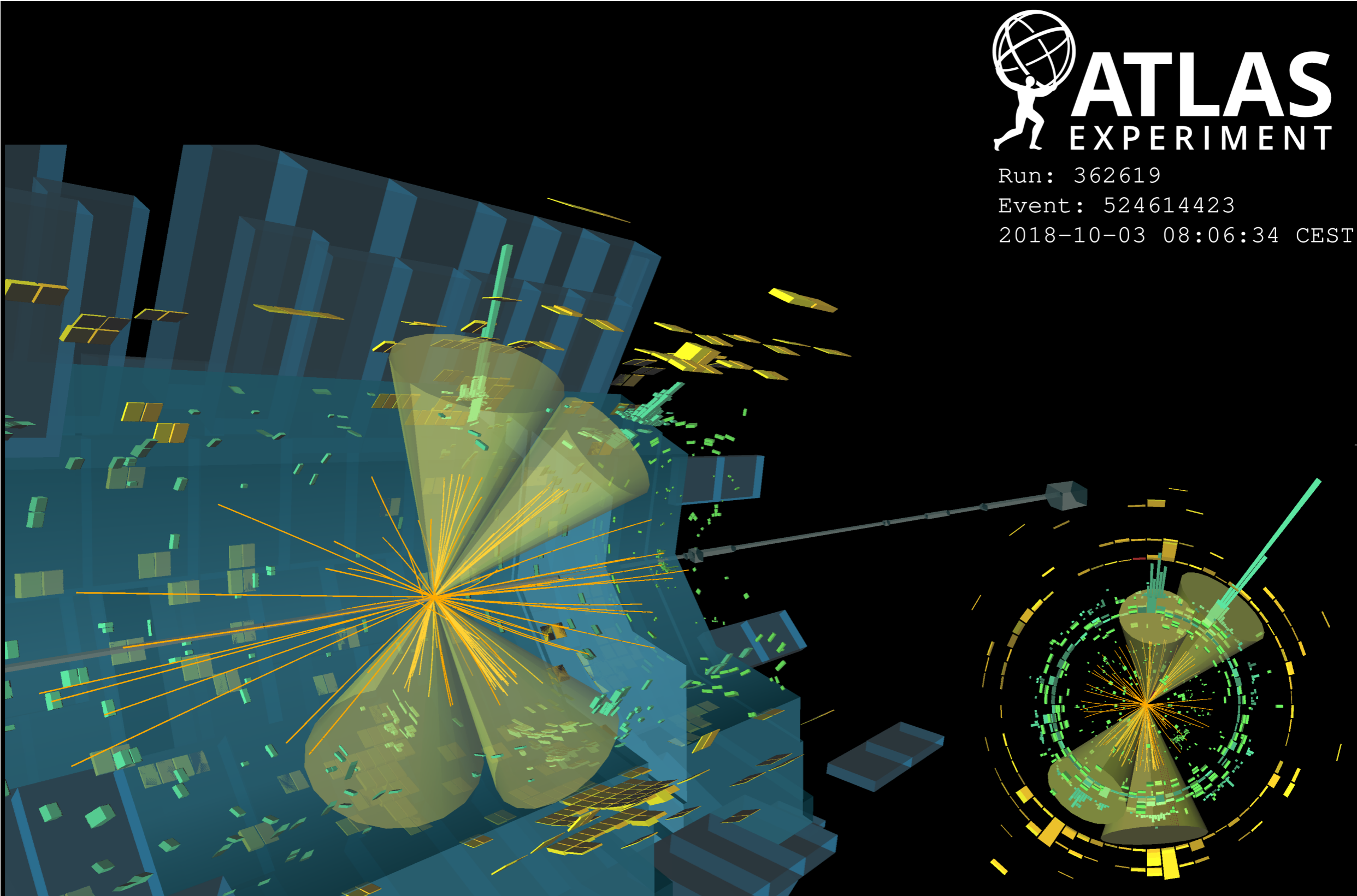
Warsaw basis couplings and  $\tilde{d}$  are in blue.  
 Higgs basis couplings are in red.

$C_{H\tilde{W}}$  is from prod + decay fit.  
 $\tilde{C}_{ZZ}$  is from production only fit.

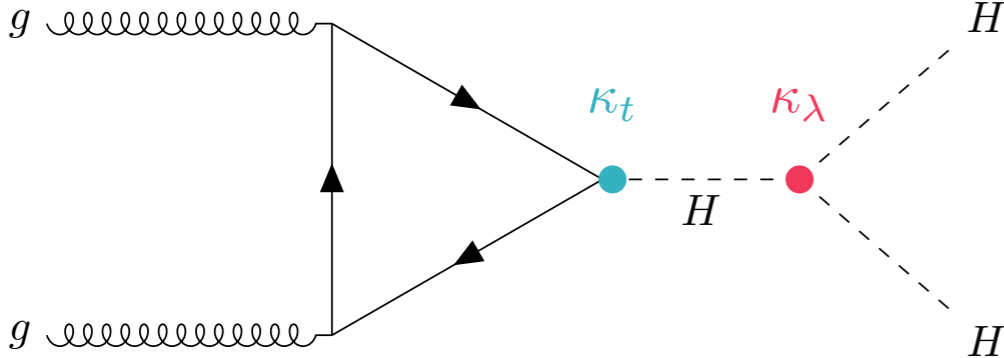
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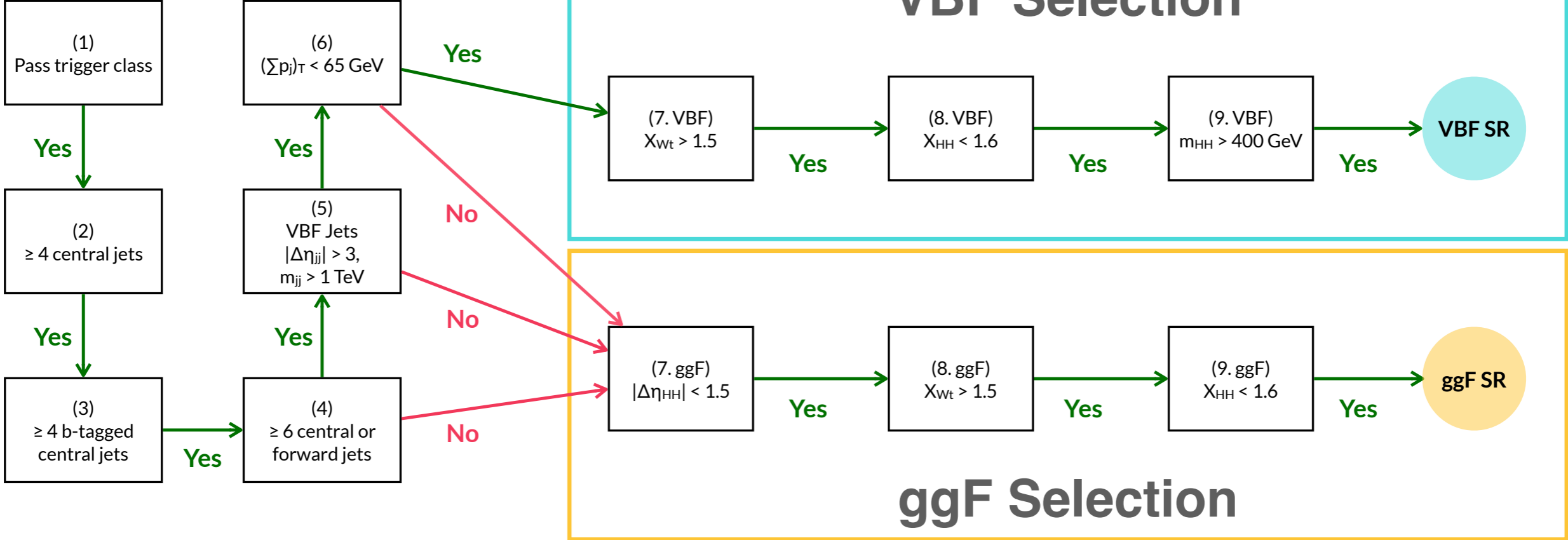
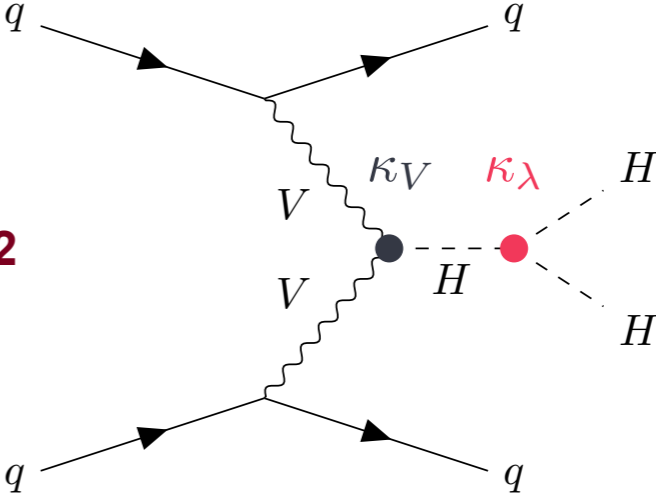
# Search for HH production in 4b final state



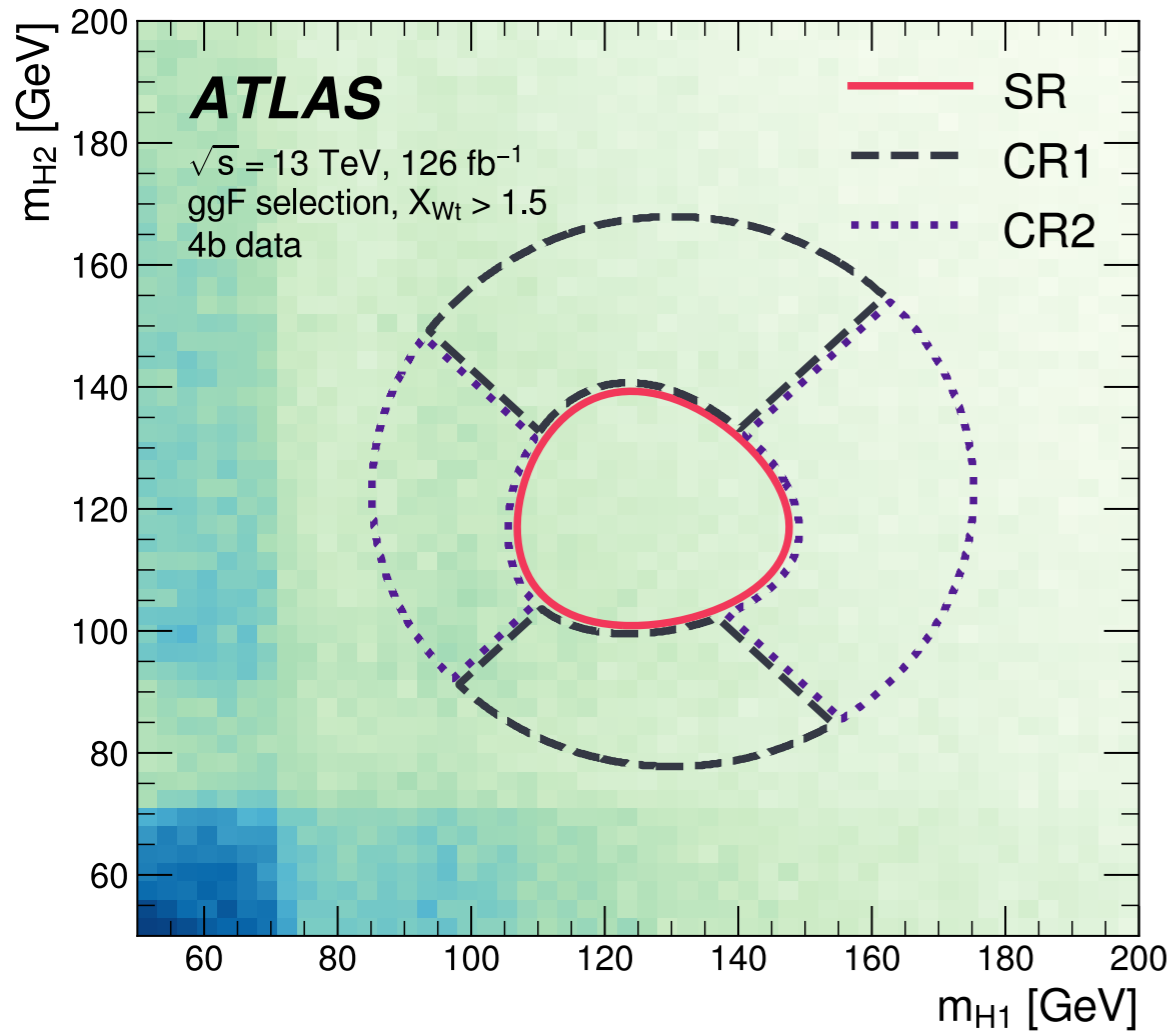
# Search for HH production in 4b final state



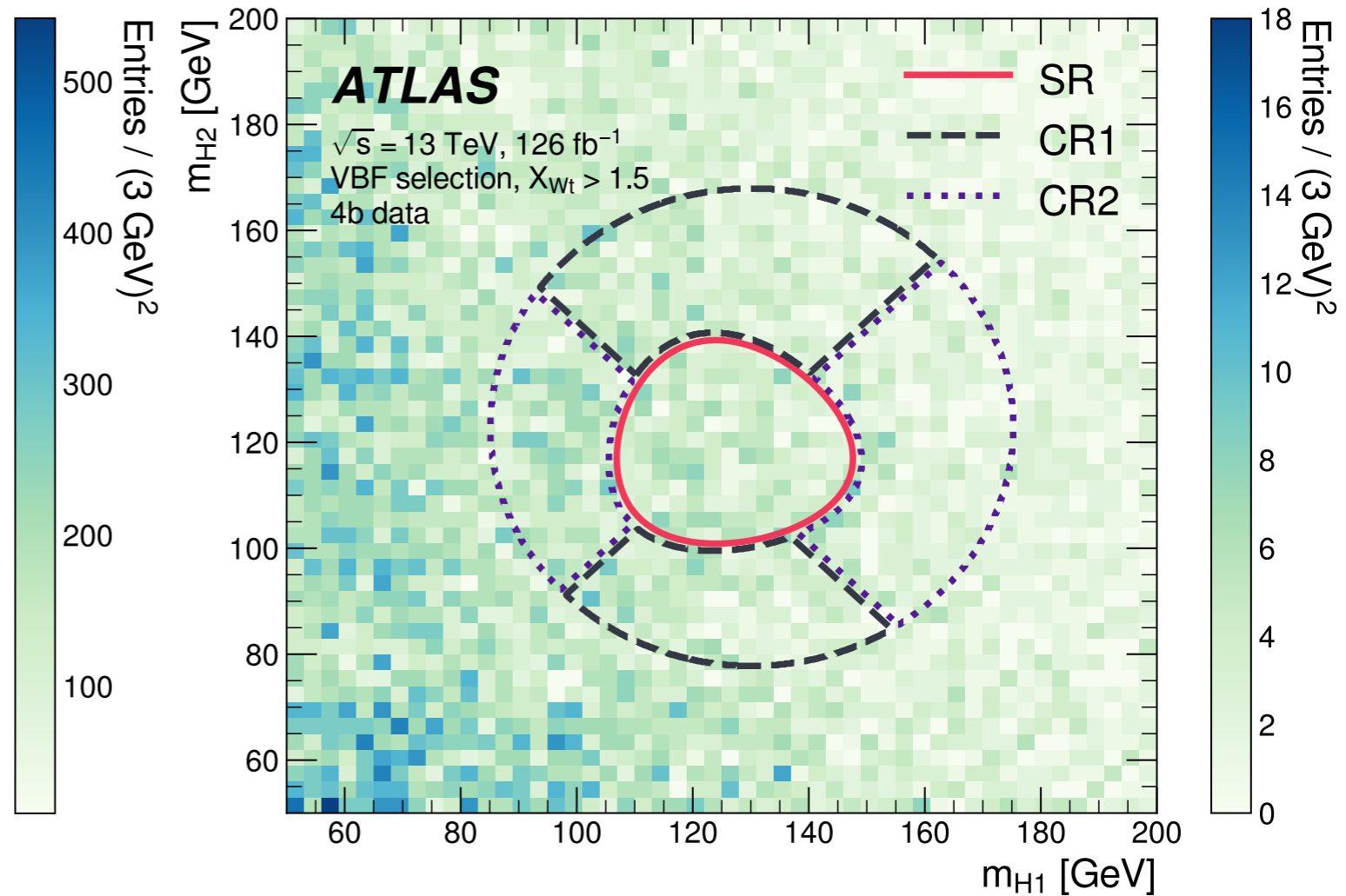
arXiv 2301.03212



# Definition of the analysis region



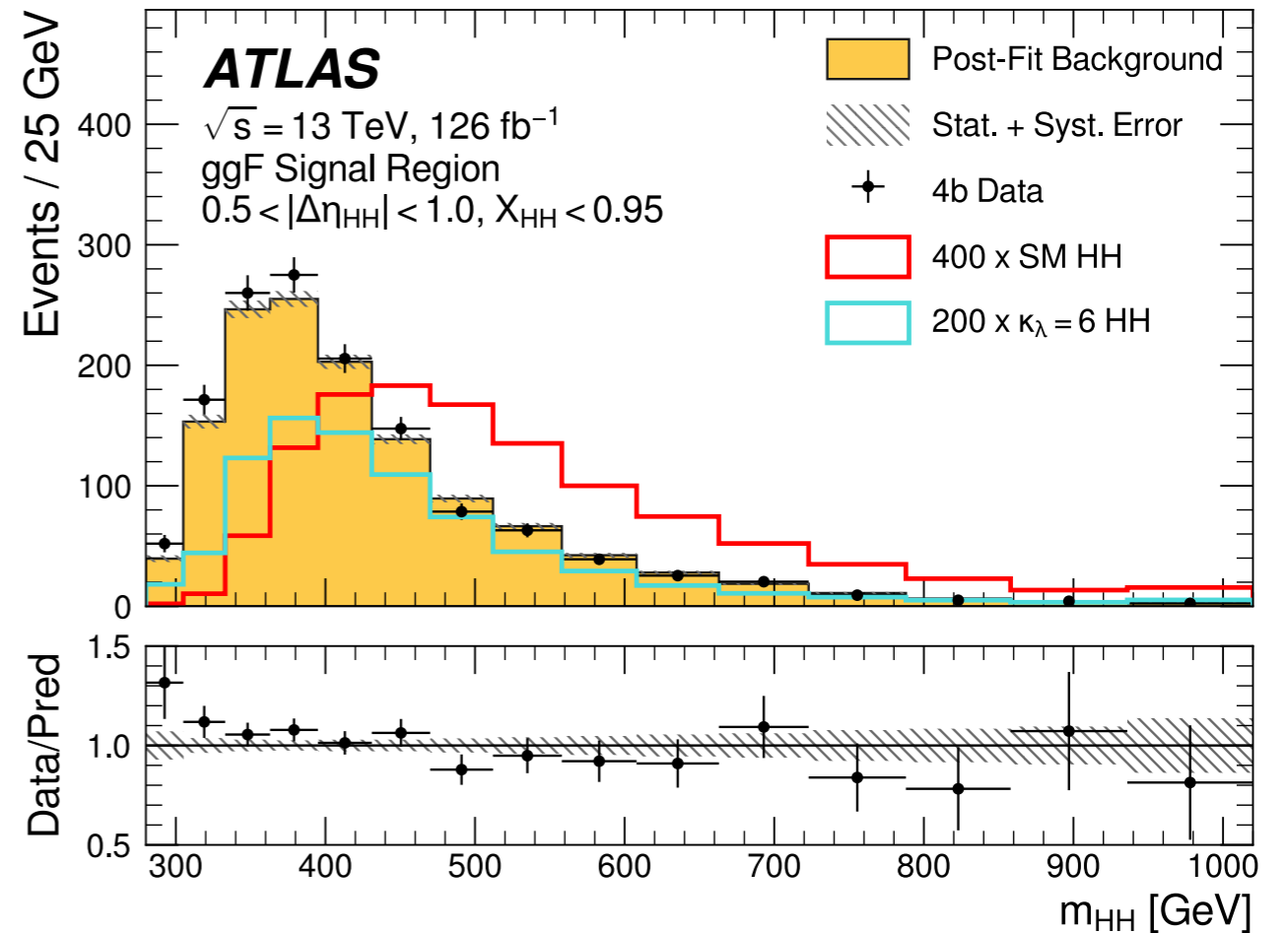
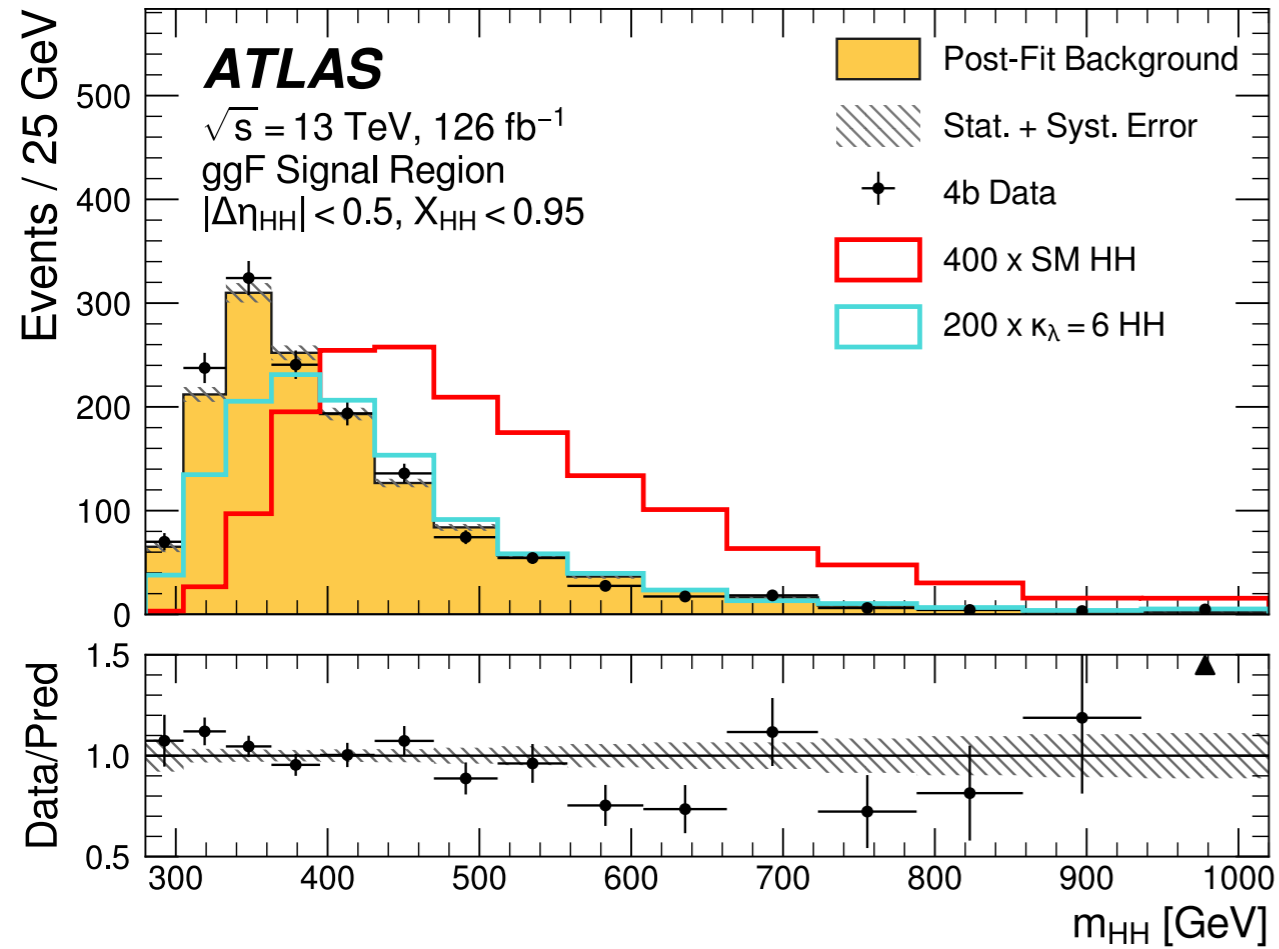
ggF signal region



VBF signal region

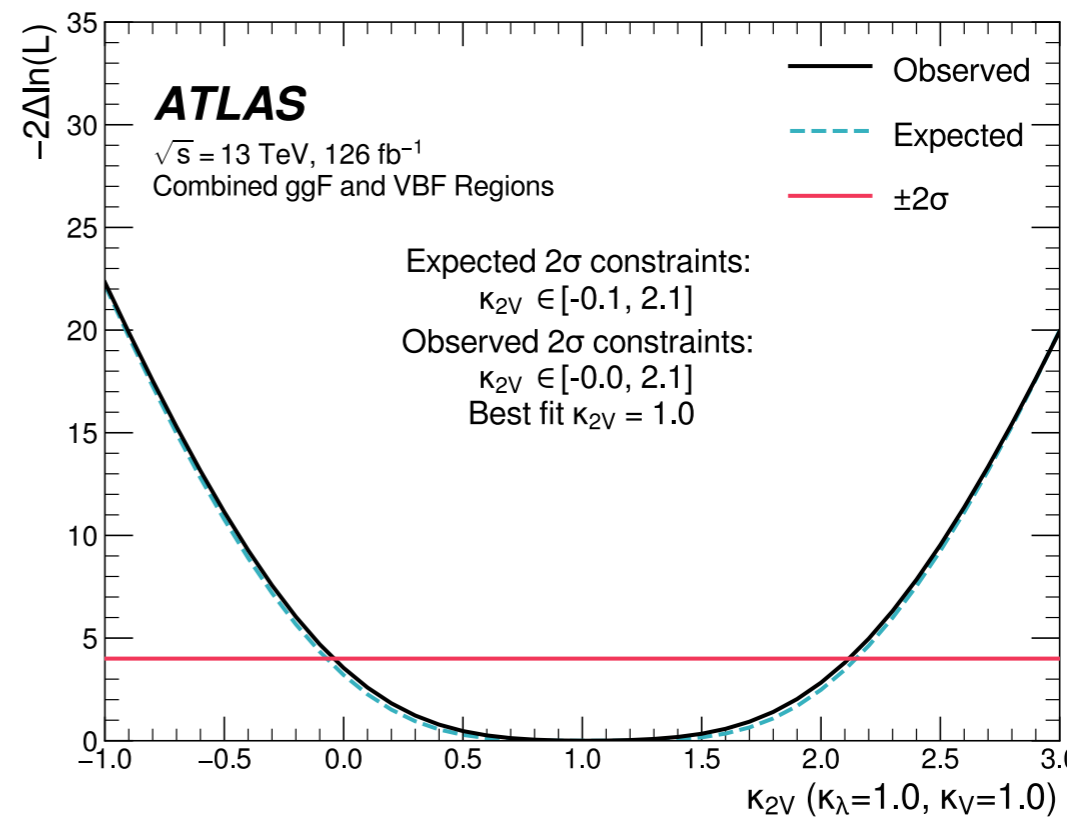
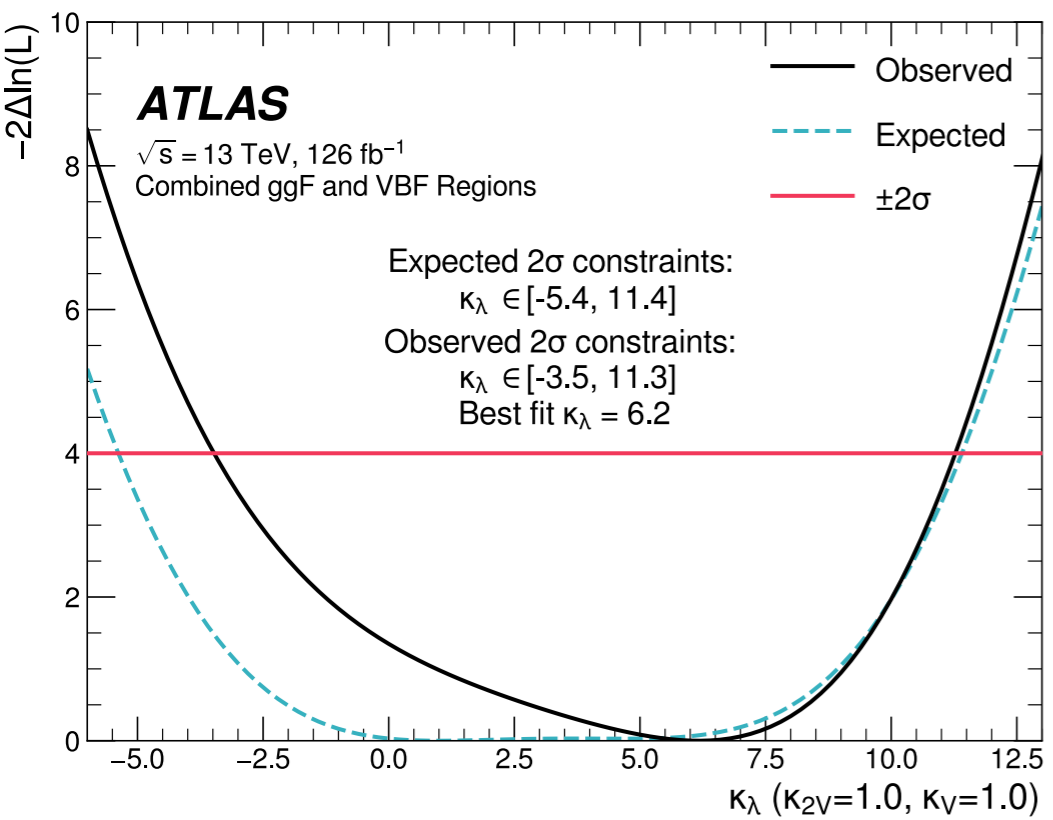
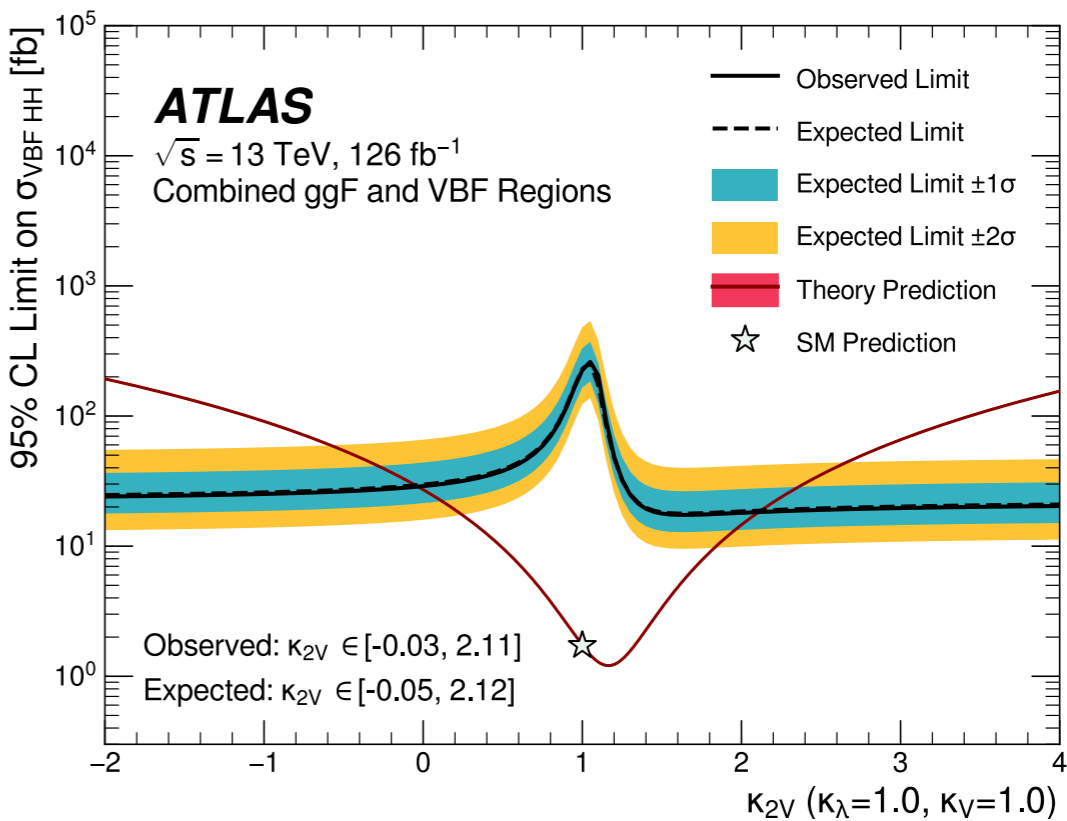
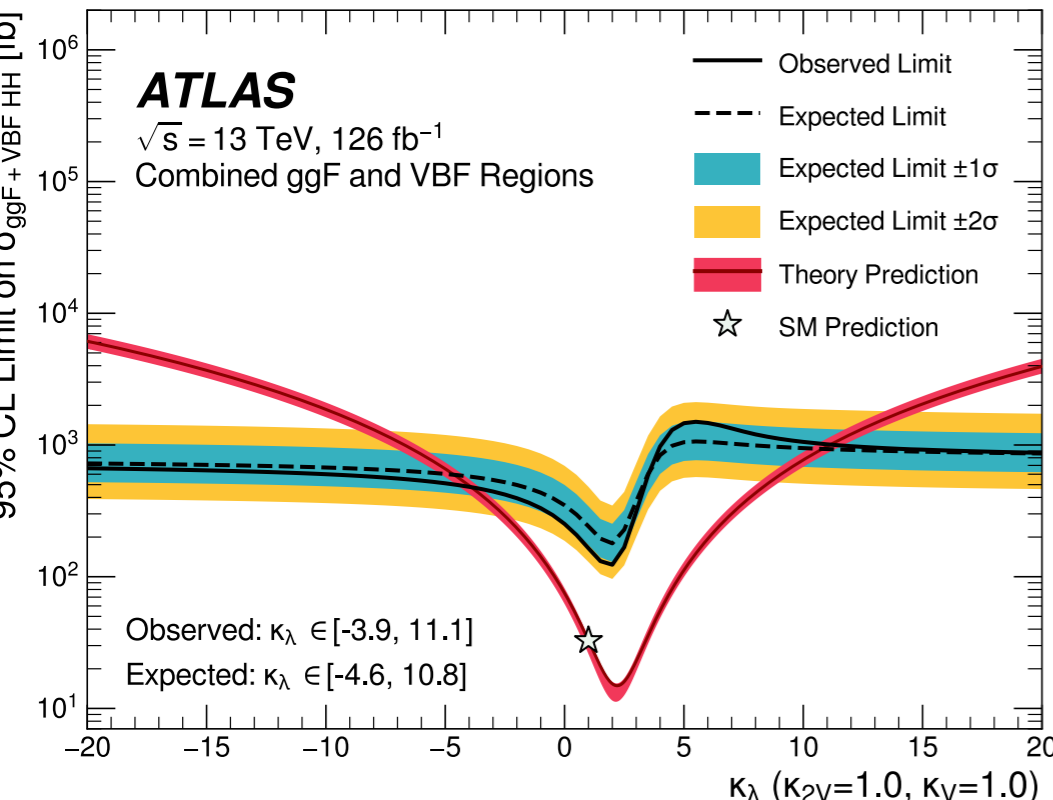
The mass planes of the reconstructed Higgs boson candidates.

# Search for HH production in 4b final state



$$X_{HH} = \sqrt{\left(\frac{m_{H1} - 124 \text{ GeV}}{0.1 m_{H1}}\right)^2 + \left(\frac{m_{H2} - 117 \text{ GeV}}{0.1 m_{H2}}\right)^2}$$

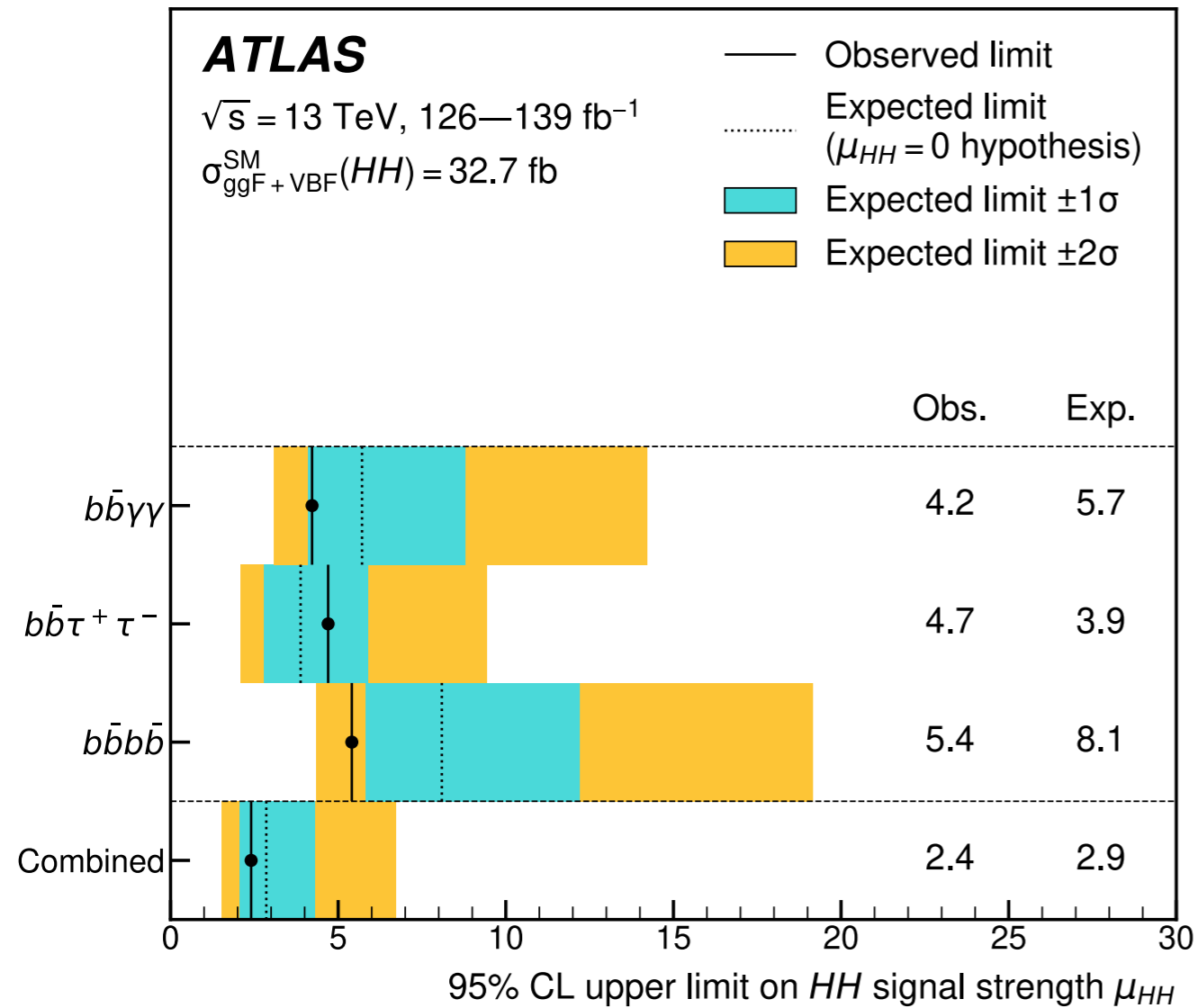
# Search for HH production in 4b final state



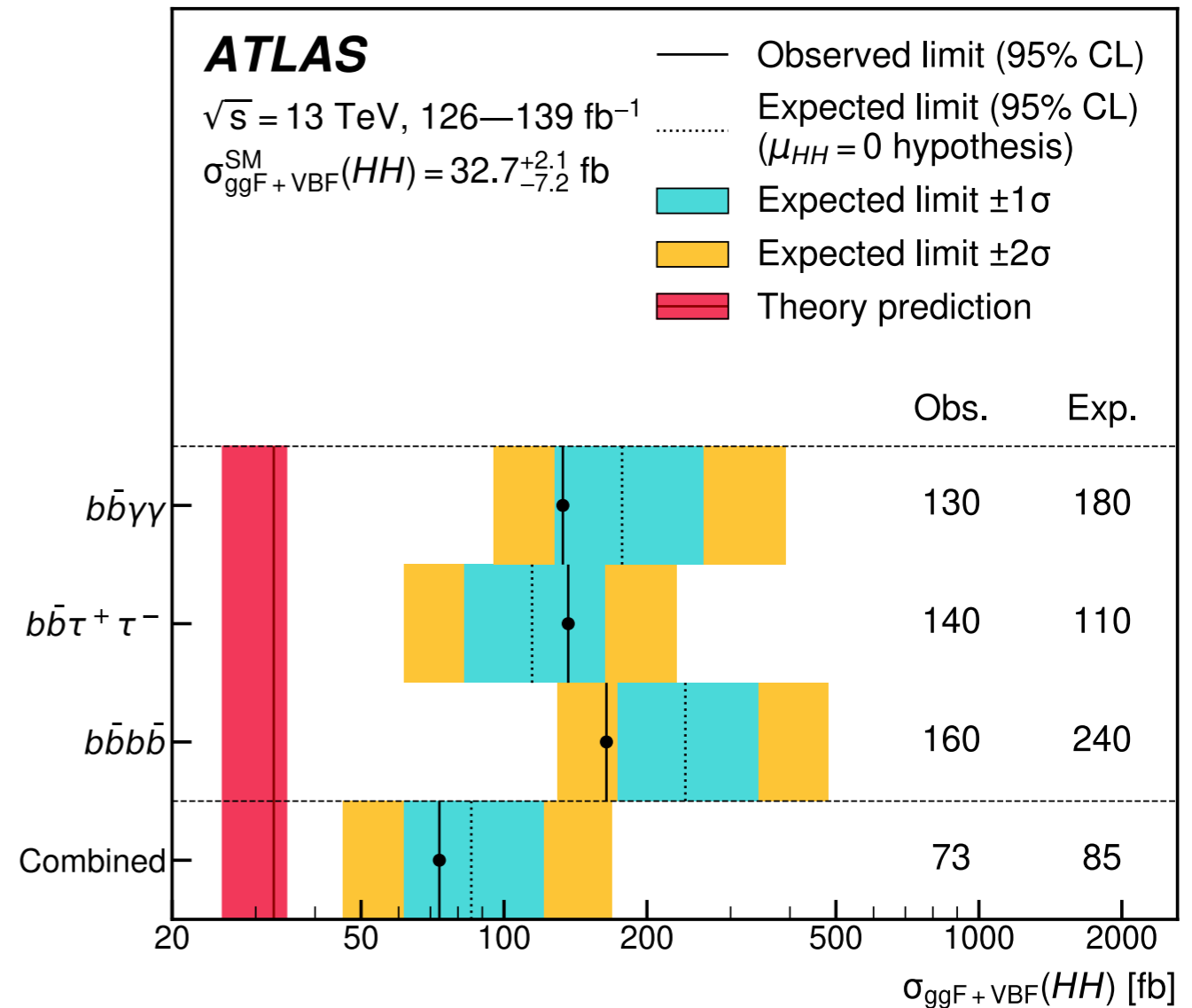


# Double & Single Higgs combination

arXiv 2211.01216

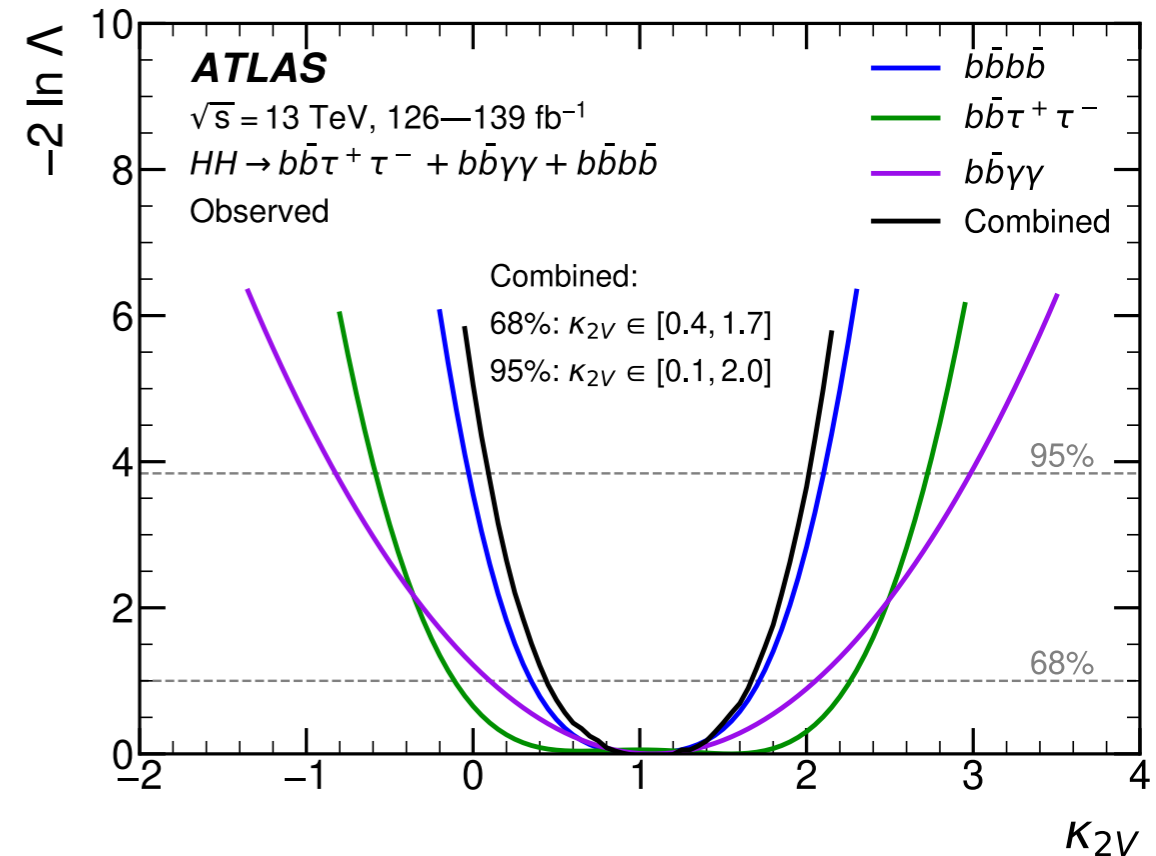
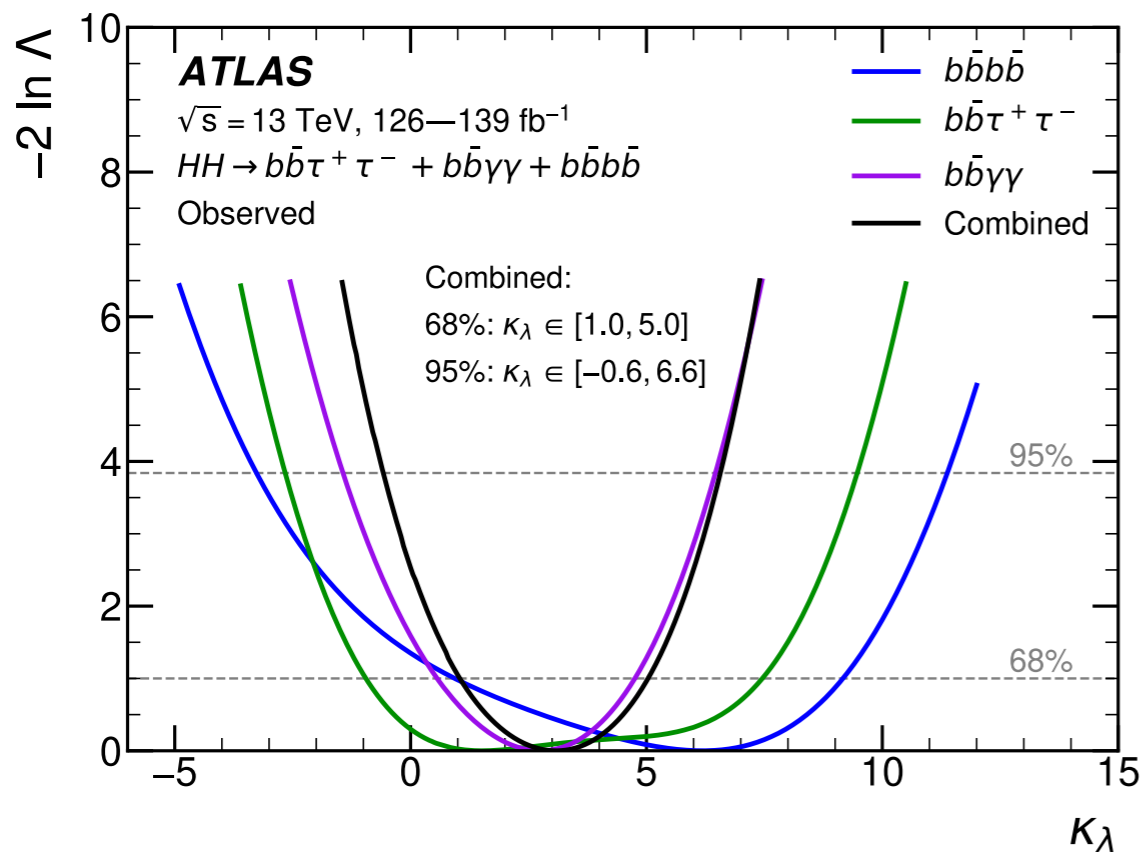
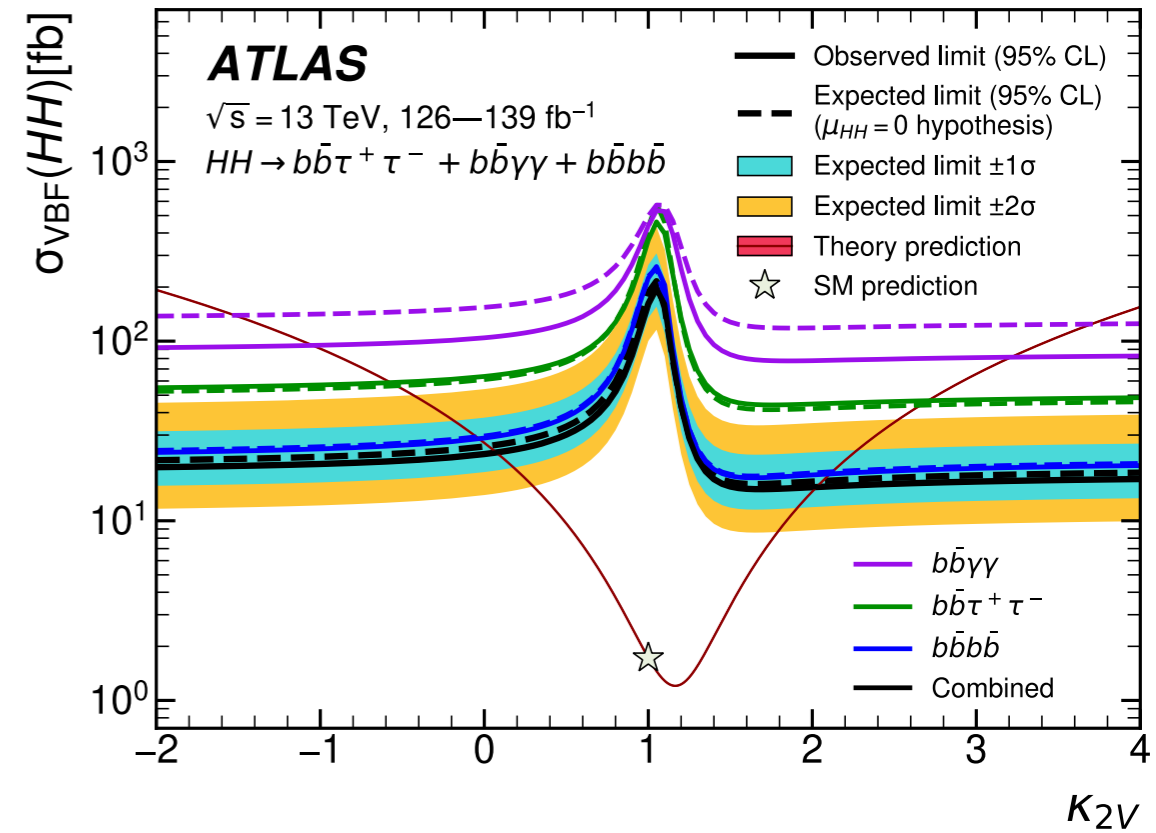
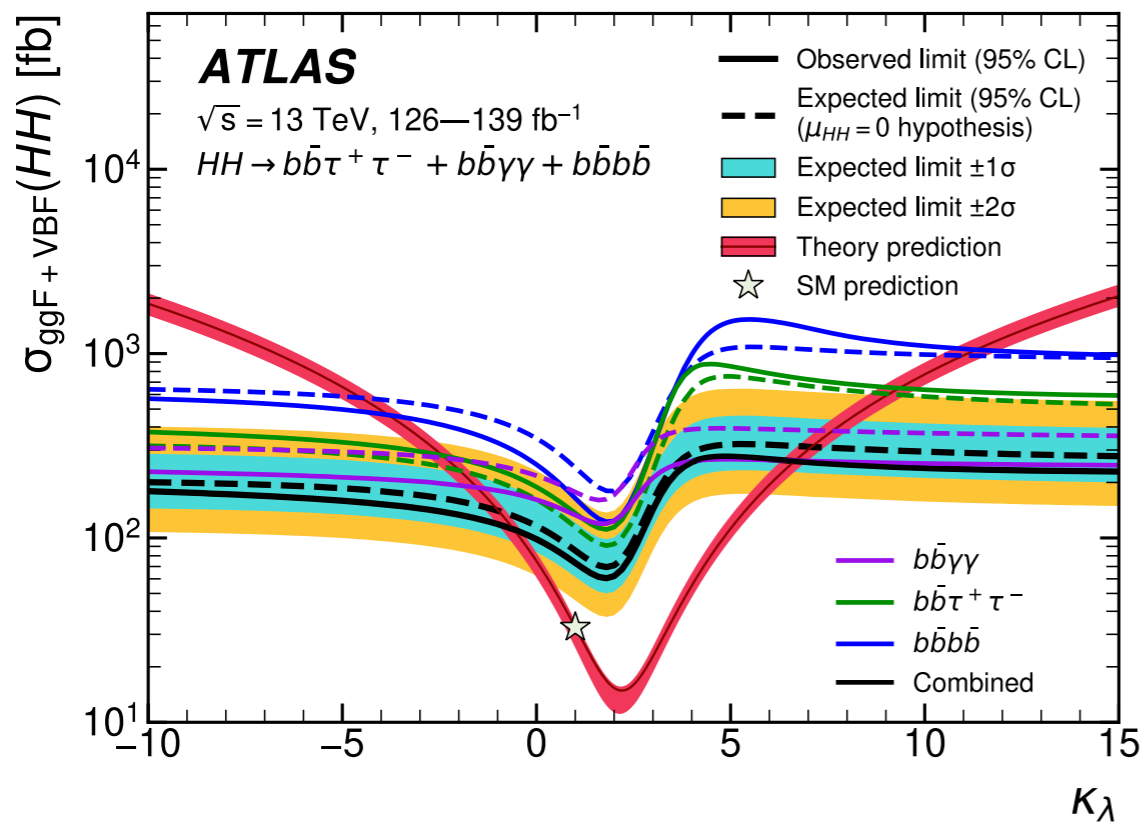


Observed and expected 95% CL upper limits on the signal strength. The SM prediction is with Higgs mass  $m_H = 125.09 \text{ GeV}$ .



Observed and expected 95% CL upper limits on HH x-section. The SM prediction is with Higgs mass  $m_H = 125.09 \text{ GeV}$ .

# Double & Single Higgs combination

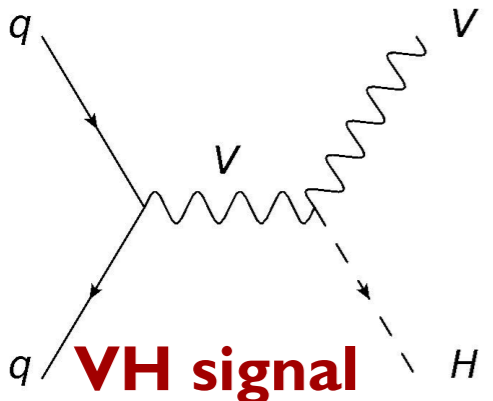


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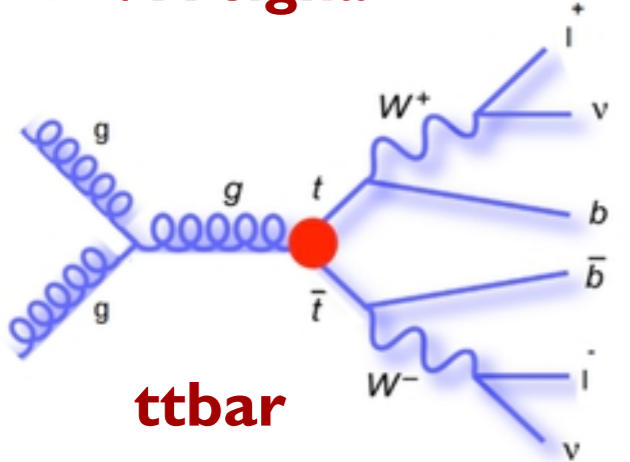
# V + H to cc : Charm Yukawa probe

Signal :  $pp \rightarrow V + H ; H \rightarrow c\bar{c}$

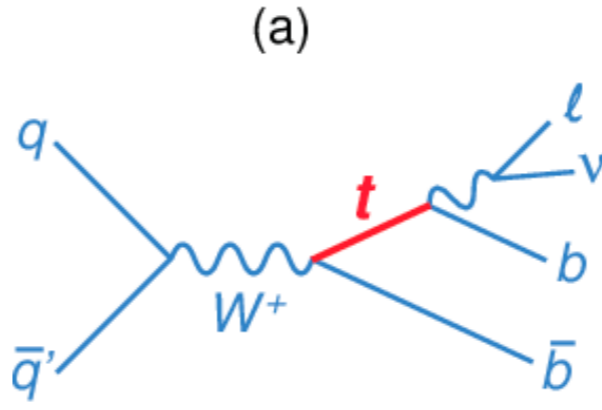


**VH signal**

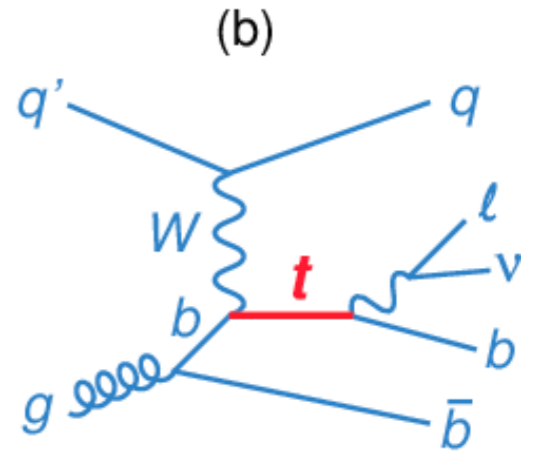
**single-top**



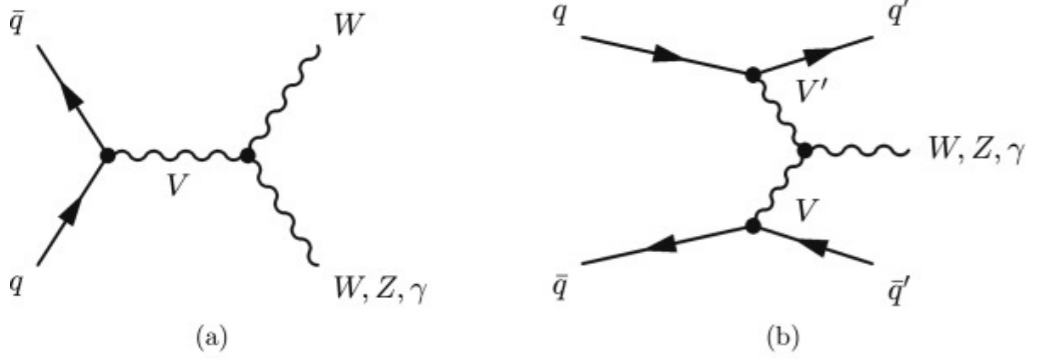
**ttbar**



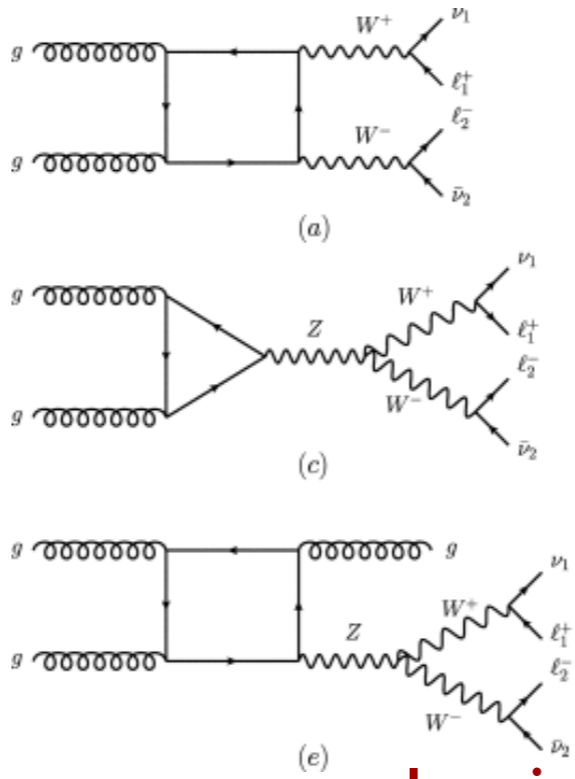
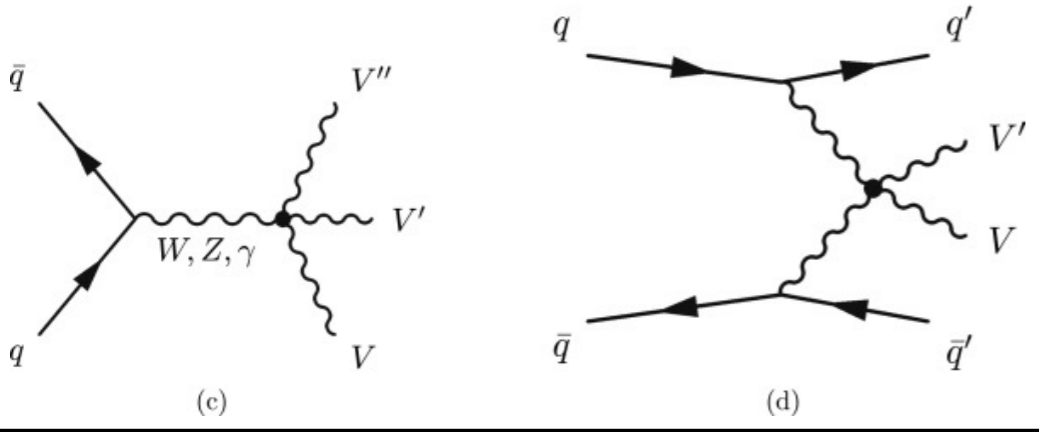
(a)



(b)

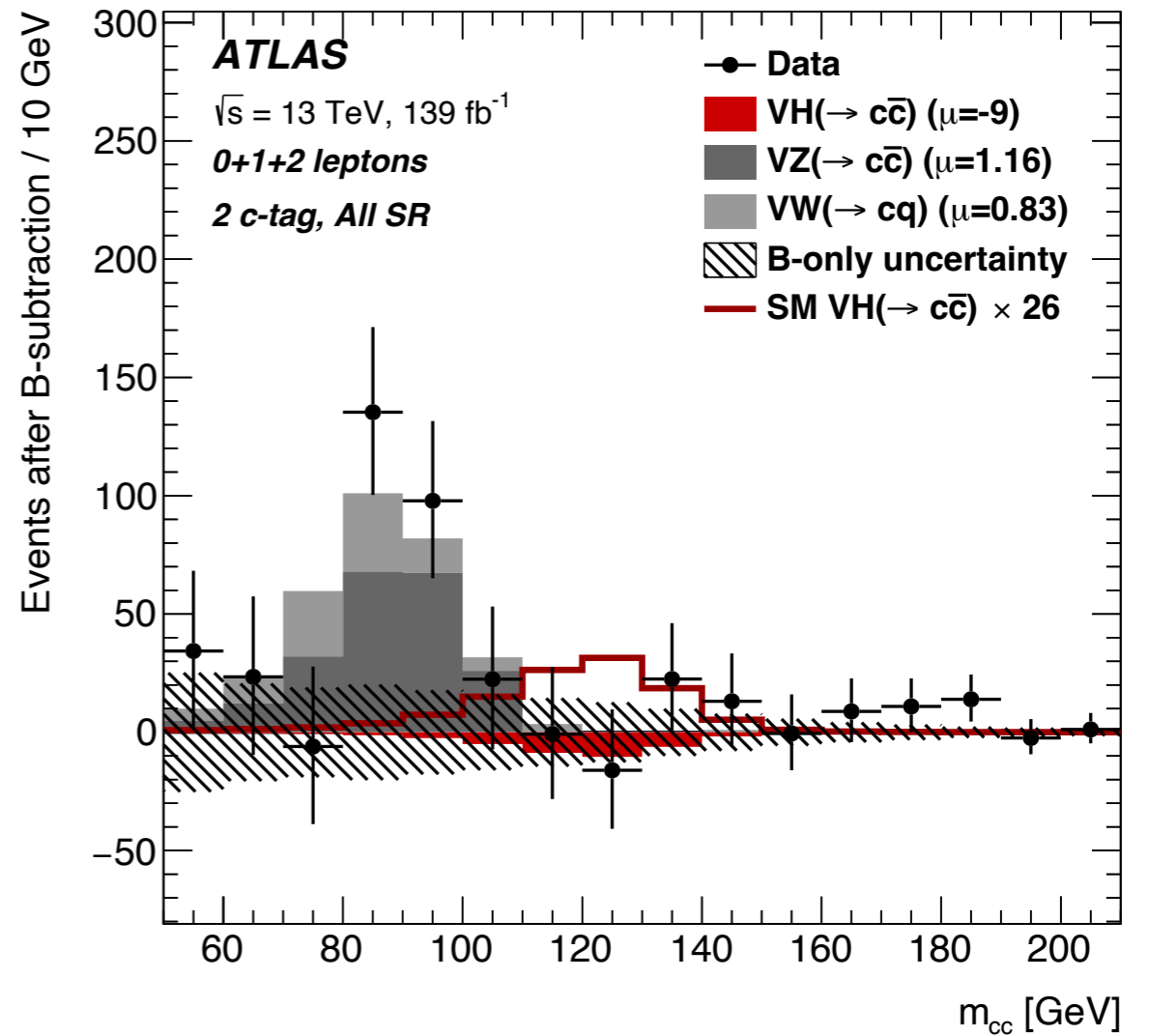
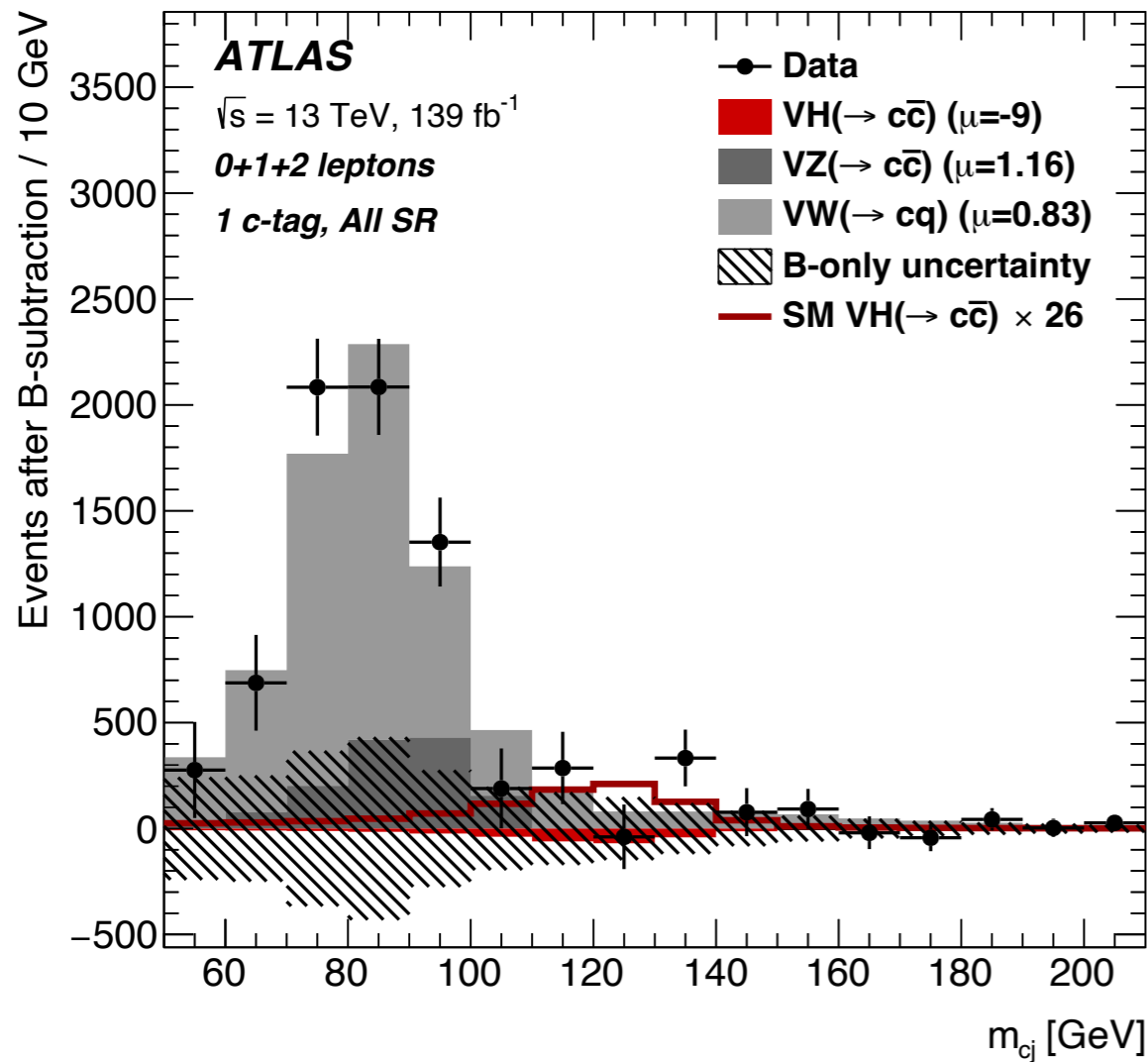


**quark initiated multi-boson/V+jets**



**gluon initiated multi-boson/V + jets**

# Extraction of 3 POI

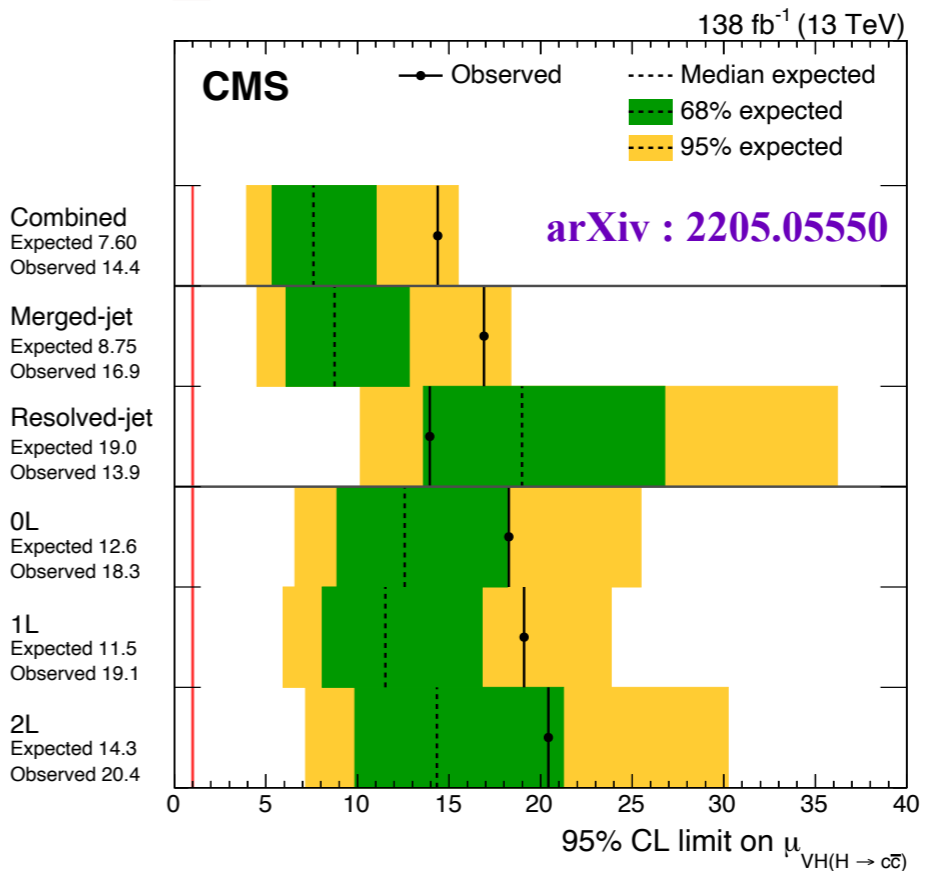
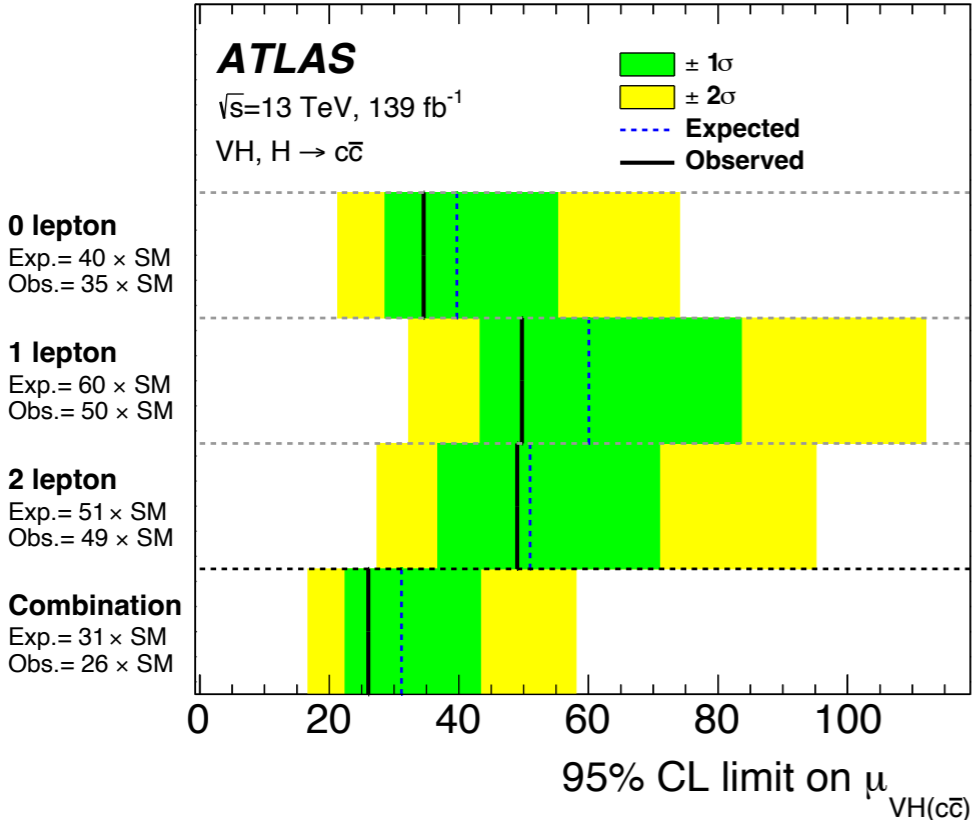


The 3 POI fit after background subtraction combining the three channels

Establishing the consistency of nuisance parameter pull across different regions and understanding the correlation across regions is a key part of this search strategy in order to reliably extract the POI from likelihood fit.

# Where do we stand compared to CMS

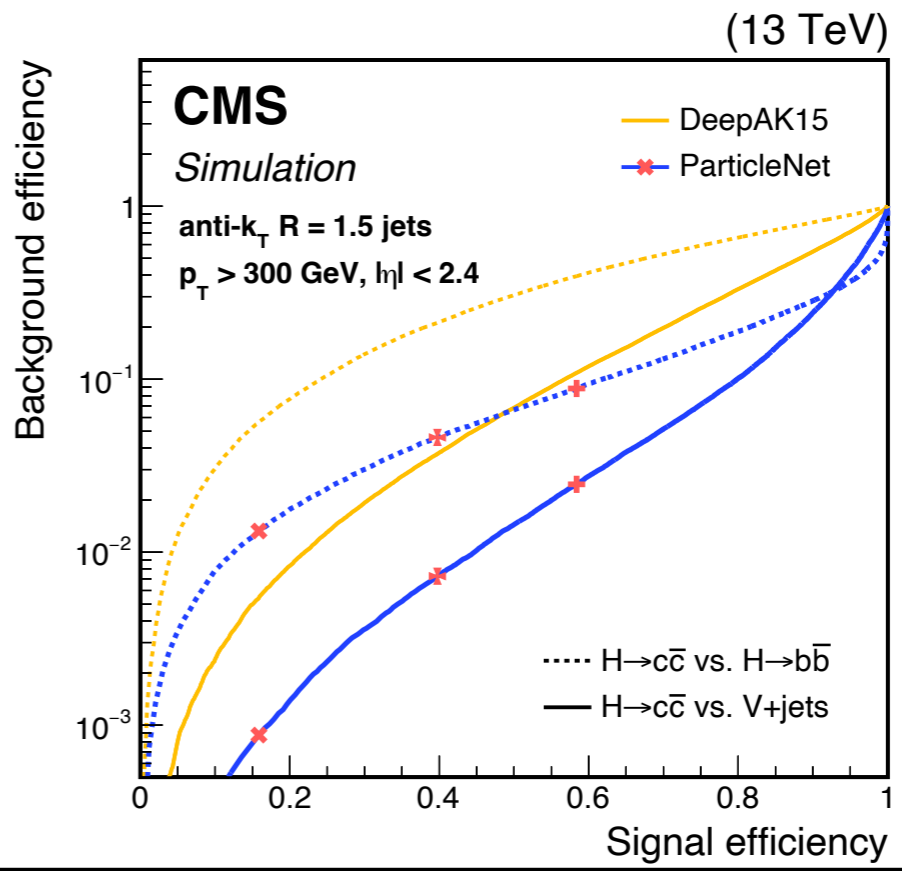
Eur. Phys. J. C (2022) 82:717



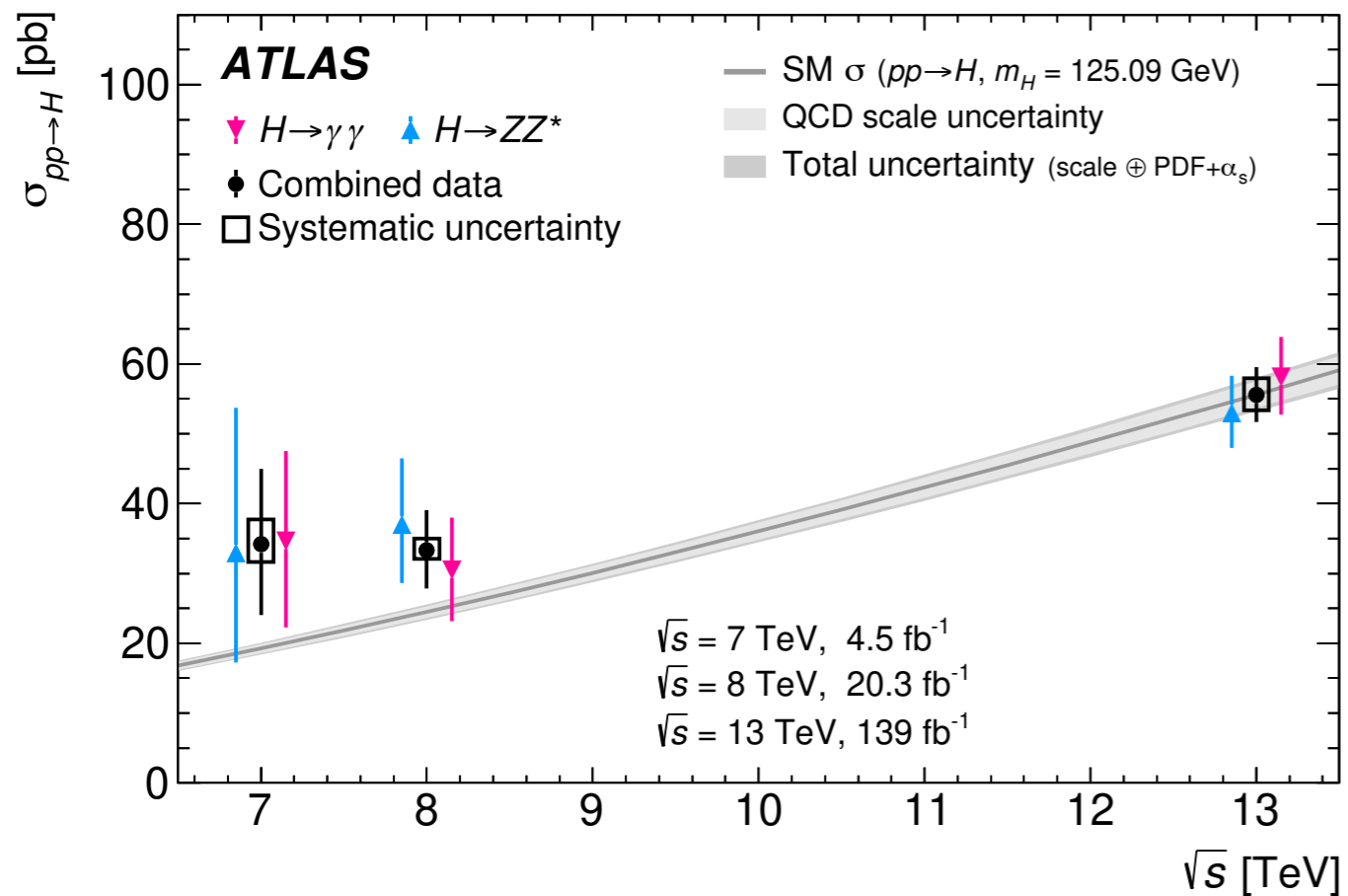
**ATLAS bound :  $|\kappa_c| < 8.5$**

**CMS bound :  $1.1 < |\kappa_c| < 5.5$**

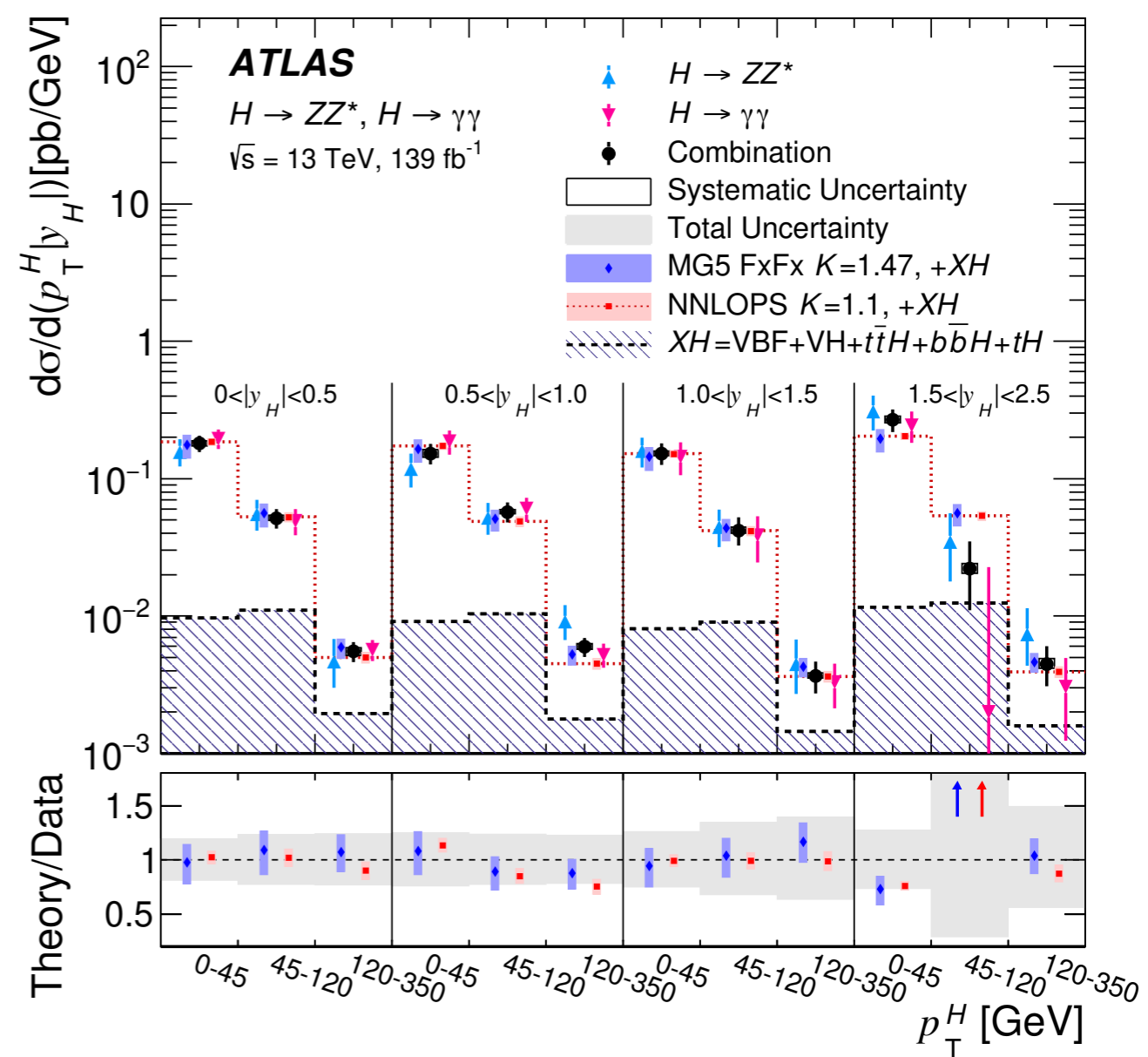
- Probable reasons for CMS got better limit :**
1. They merged the boosted channel
  2. They don't anti-tag the b-jet : higher statistical significance.
  3. Have a dedicated GNN based Xcc tagger



# Extraction of $\kappa_b$ & $\kappa_c$ from $p_T^H$ ( $\gamma\gamma + ZZ$ )



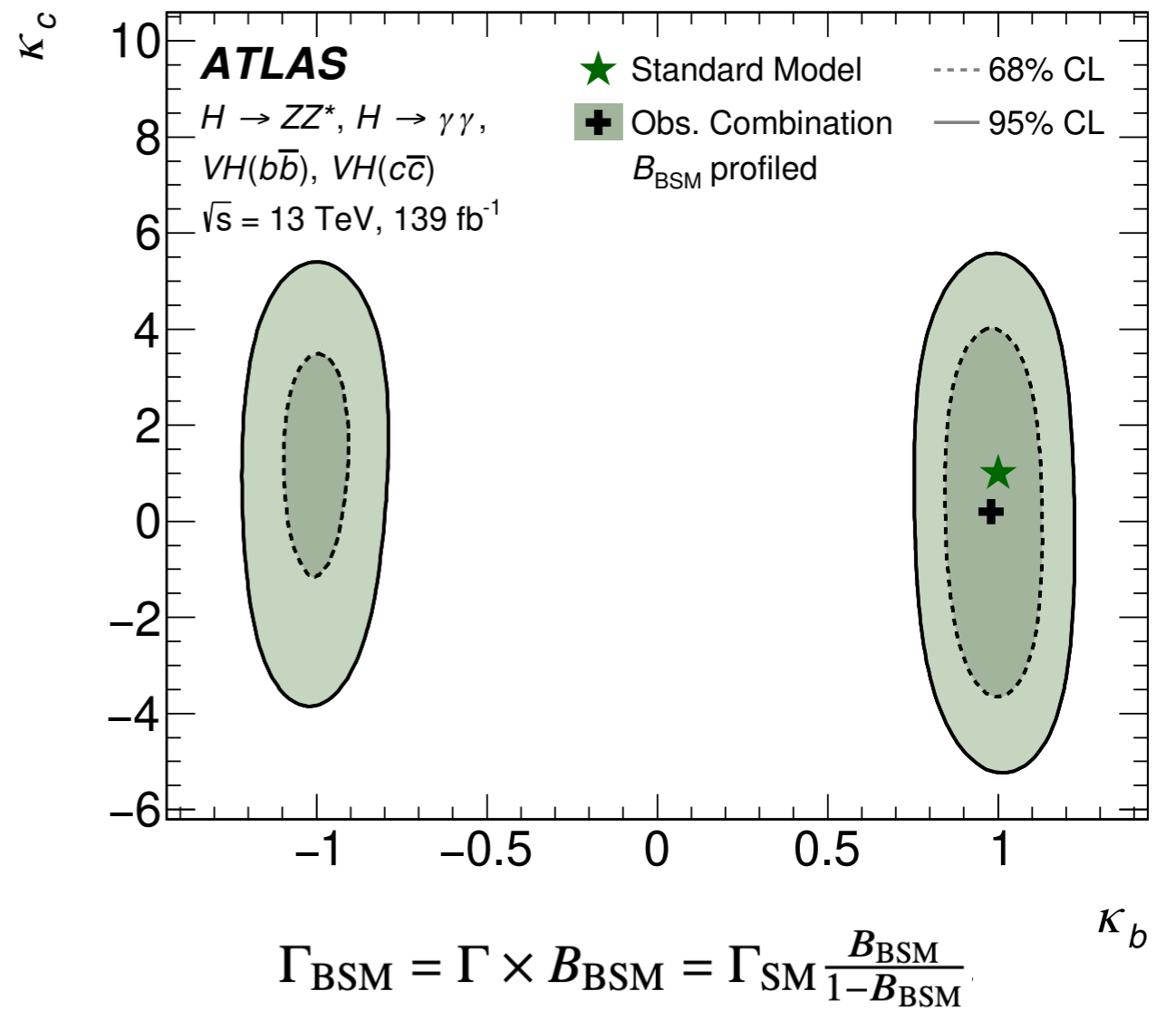
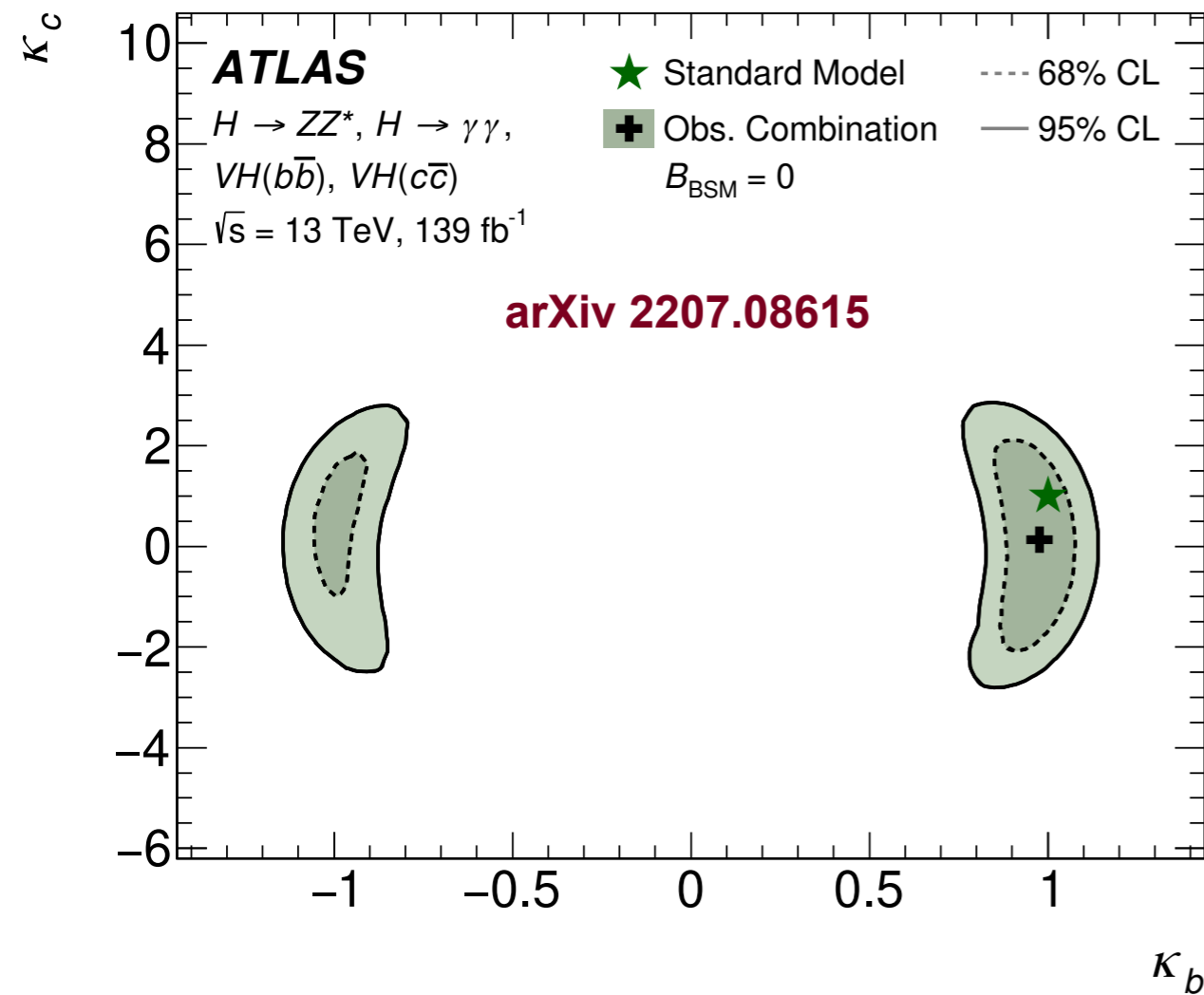
arXiv 2207.08615



The combined measurement of H  $p_T$  from  $H \rightarrow ZZ^* \rightarrow 4l$  and  $H \rightarrow \gamma\gamma$  is made.

Measured inclusive x-section :  $55.5_{-3.8}^{+4.0}$  pb, predicted  $55.6 \pm 2.8$  pb.

# Extraction of $\kappa_b$ & $\kappa_c$ from $p_T^H$



Channel	Parameter	Observed 95% confidence interval	Expected 95% confidence interval
$H \rightarrow ZZ^* \rightarrow 4\ell$	$\kappa_b$	$[-1.14, -0.88] \cup [0.80, 1.17]$	$[-1.23, -0.87] \cup [0.82, 1.20]$
	$\kappa_c$	$[-2.94, 2.99]$	$[-3.33, 3.14]$
$H \rightarrow \gamma\gamma$	$\kappa_b$	$[-1.12, -0.78] \cup [0.78, 1.07]$	$[-1.18, -0.87] \cup [0.83, 1.19]$
	$\kappa_c$	$[-2.46, 2.32]$	$[-3.03, 3.09]$
Combined	$\kappa_b$	$[-1.09, -0.86] \cup [0.81, 1.09]$	$[-1.14, -0.92] \cup [0.86, 1.15]$
	$\kappa_c$	$[-2.27, 2.27]$	$[-2.77, 2.75]$

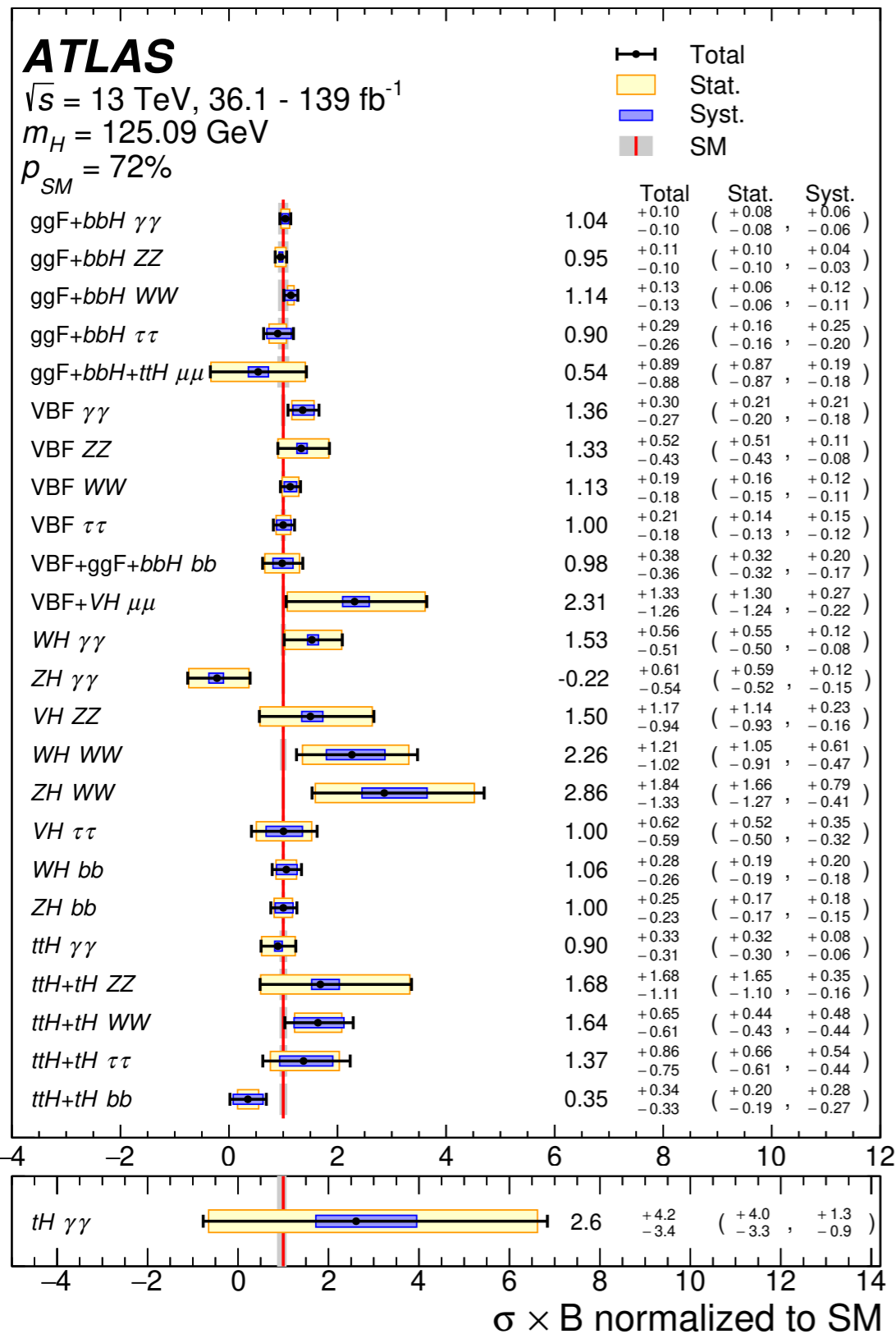


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# Take away

Nature volume 607, pages 52–59 (2022)



✓ Higgs precision measurement is must in order to understand dynamics of SM and probe beyond standard model physics.

✓ The Run-3 and HL-LHC will give us golden opportunity to do Higgs physics and it's worth exploiting it.

✓ There are still plenty of room to innovate new techniques : a few % improvement of H width measurement has large implications on physics understanding.

✓ The up-to-date Higgs Mass, CP measurements & width are presented.

✓ The SMEFT & kappa framework fits are yet to show any hints of BSM models.

**THANK YOU!!**

# Take away

**Backup**

# EFT couplings

Wilson coefficient	Operator definition	Example diagram		Wilson coefficient	Operator definition	Example diagram
$c_{HG}$	$\Phi^\dagger \Phi G_{\mu\nu}^a G^{a\mu\nu}$			$c_{Hl3}$	$(i\Phi^\dagger \overleftrightarrow{D}_\mu^I \Phi)(\bar{\ell}\sigma^I \gamma^\mu \ell)$	
$c_{HB}$	$\Phi^\dagger \Phi B_{\mu\nu} B^{\mu\nu}$			$c_{Hu}$	$(i\Phi^\dagger \overleftrightarrow{D}_\mu^I \Phi)(\bar{u}\gamma^\mu u)$	
$c_{HW}$	$\Phi^\dagger \Phi W_{\mu\nu}^I W^{I\mu\nu}$			$c_{Hd}$	$(i\Phi^\dagger \overleftrightarrow{D}_\mu^I \Phi)(\bar{d}\gamma^\mu d)$	
$c_{HWB}$	$\Phi^\dagger \Phi W_{\mu\nu}^I B^{I\mu\nu}$			$c_{He}$	$(i\Phi^\dagger \overleftrightarrow{D}_\mu^I \Phi)(\bar{e}\gamma^\mu e)$	
$c_{Hq1}$	$(i\Phi^\dagger \overleftrightarrow{D}_\mu^I \Phi)(\bar{q}\gamma^\mu q)$			$ c_{uG} $	$(\bar{q}\sigma^{\mu\nu} T^a \Phi u) G_{\mu\nu}^a$	
$c_{Hl1}$	$(i\Phi^\dagger \overleftrightarrow{D}_\mu^I \Phi)(\bar{\ell}\gamma^\mu \ell)$			$ c_{eH} $	$(\Phi^\dagger \Phi)(\bar{\ell}e\Phi)$	
$c_{Hq3}$	$(i\Phi^\dagger \overleftrightarrow{D}_\mu^I \Phi)(\bar{q}\sigma^I \gamma^\mu q)$			$ c_{dH} $	$(\Phi^\dagger \Phi)(\bar{q}d\Phi)$	