

HIGGS COUPLINGS TO FERMIONS

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The good, the bad and the ugly

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the flavor hierarchy, the theory

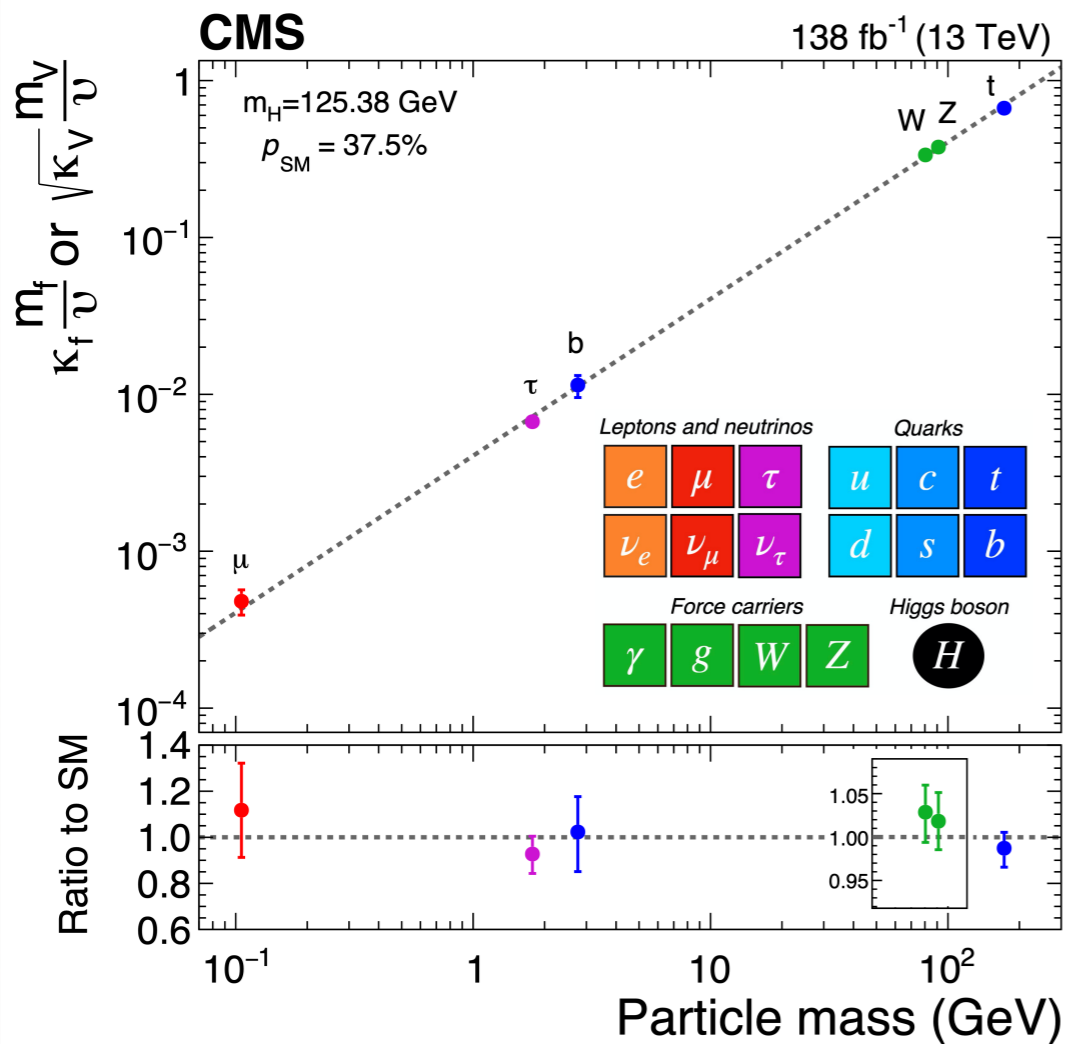
(C). Higgs couplings at colliders:

the practice: LHC, ILC & a μ C

y_t , y_c , y_μ and even y_e

(A). THE HIGGS MAGIC

$m_H = 125.38 \pm 0.14 \text{ GeV}$ *PLB 805 (2020) 135425*



The Standard Model – first time ever!

- Quantum mechanical
- Relativistic
- Renormalizable
- Perturbatively unitary to exponentially high scales, perhaps to the Planck scale!

All known physics

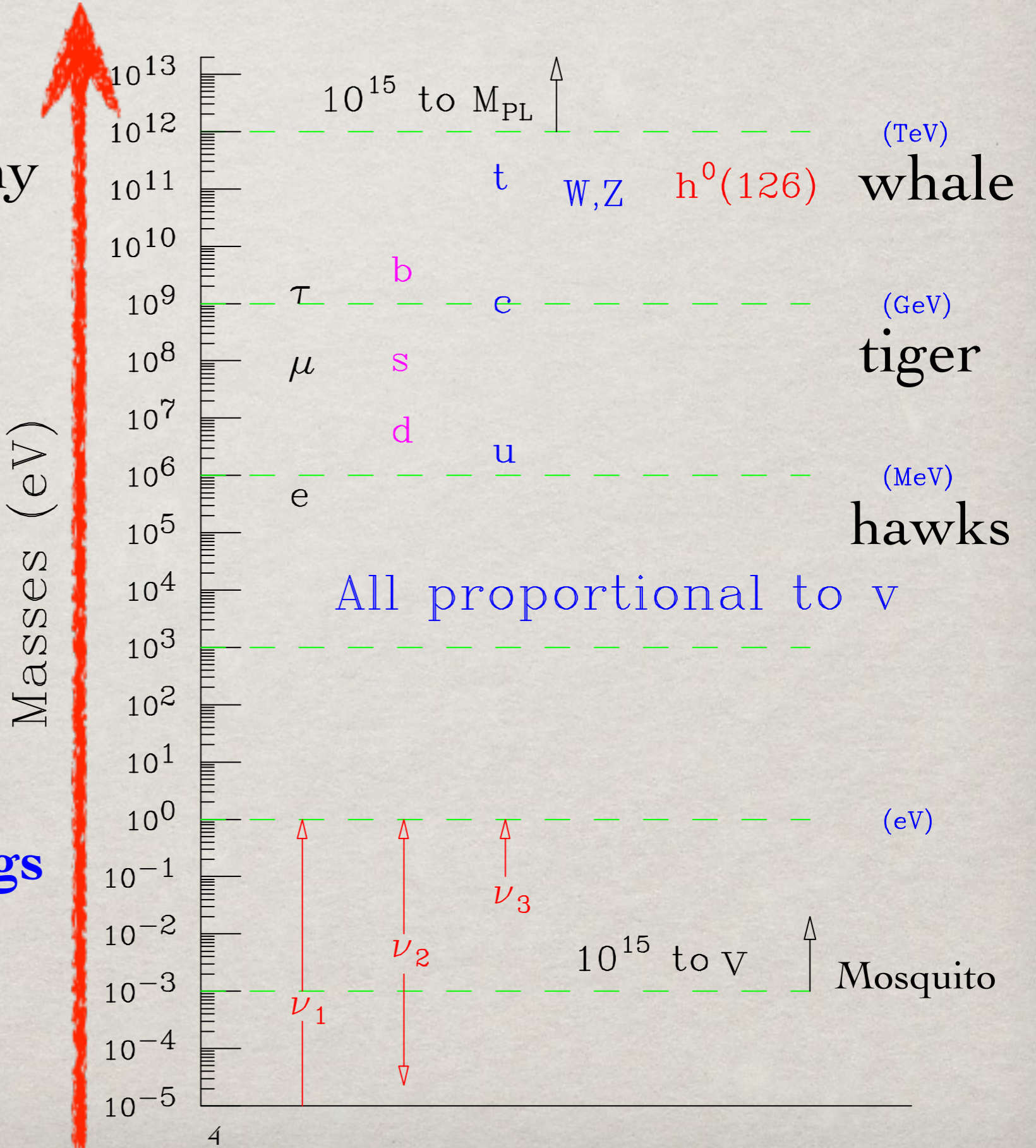
$$W = \int_{k < \Lambda} [\mathcal{D}g \dots] \exp \left\{ \frac{i}{\hbar} \int d^4x \sqrt{-g} \left[\frac{1}{16\pi G} R - \frac{1}{4} F^2 - \bar{\psi} i \not{D} \psi - \lambda \phi \bar{\psi} \psi + |D\phi|^2 - V(\phi) \right] \right\}$$



THE HIGGS BLEMISHES

- Particle mass hierarchy
- Patterns of quark, neutrino mixings
- New CP-violation sources?
- Tiny neutrino masses!

Higgs Yukawa couplings as the pivot for all !



THREE TYPES OF MASSES IN SM

$M_{W,Z}$ versus m_H versus m_f :

The good, the bad, and the ugly

(1). $M_{W,Z}$: the good! In the SM: $\phi = \frac{1}{\sqrt{2}} e^{i \sum \xi^i L^i} \begin{pmatrix} 0 \\ \nu + H \end{pmatrix}$

$$V(|\Phi|) = -\mu^2 \Phi^\dagger \Phi + \lambda (\Phi^\dagger \Phi)^2 \quad \langle |\Phi| \rangle = v = (\sqrt{2} G_F)^{-1/2} \approx 246 \text{ GeV}$$

$$M_W^2 W^{\mu+} W_\mu^- \left(1 + \frac{H}{v}\right)^2 + \frac{1}{2} M_Z^2 Z^\mu Z_\mu \left(1 + \frac{H}{v}\right)^2$$

$\rightarrow M_W, M_Z = g v/2$ predicted, and:

$$\delta m_w^2 \sim m_w^2 \ln(\Lambda/m_w)$$

BSM: **easy to break** $SU(2)_L$ gauge sector:

- Fundamental scalars (SUSY)
- Dynamical breaking (TC, composite ...) $\langle \bar{q}_L q_R + \bar{q}_R q_L \rangle \sim v^3$
- Non-linear realization (or even "Higgsless")

$$\Phi = \frac{1}{\sqrt{2}} (v + H) U, \quad U = \exp[i\pi^a \tau^a / v]$$

- All calculable and predictable \rightarrow e.g. CDF M_W

$M_{W,Z}$ versus m_H versus m_f :

(2). m_H : the bad! $m_H = \sqrt{2} \mu = (2\lambda)^{1/2} v = 125 \text{ GeV}$

In the SM, all fixed: $\rightarrow \lambda \approx 0.13$; $\frac{1}{2}m_H^2 H^2 + \frac{m_H^2}{2v} H^3 + \frac{m_H^2}{8v^2} H^4$

The value itself doesn't matter much \sim EW scale

\rightarrow Not too light (vacuum stable), not too heavy (unitarity).

But quantum corrections: $\delta\mu^2 = -\frac{3y_t^2}{8\pi^2}\Lambda^2$ ($M_{\text{PL}}^2 \dots$)

Quadratically sensitive to the new physics cutoff scale.

\rightarrow "Naturalness" or "Large hierarchy problem" ?

BSM: easy to construct a scalar model / potential,

but model-parameters quadratically sensitive to a

new physics scale: $\delta m_H^2 \propto -\frac{k^2}{4\pi^2}\Lambda^2$ ($M_{\text{SUSY}}^2, M_{\text{comp}}^2, \dots$)

Not seeing other states nearby, we do not understand m_H

\rightarrow "Little hierarchy problem" !



Note: the quadratic mass corrections are NOT experimentally observable!

$M_{W,Z}$ versus m_H versus m_f :

(3). m_f : not governed by the gauge couplings

$$\mathcal{L}_Y \sim - \sum_{i,j} (Y_{ij}^d \bar{Q}_{iL} \Phi d_{jR} + Y_{ij}^u \bar{Q}_{iL} \tilde{\Phi} u_{jR} + Y_{ij}^e \bar{L}_{iL} \Phi e_{jR} + Y_{ij}^\nu \bar{L}_{iL} \tilde{\Phi} \nu_{jR})$$



- Couplings are fixed by the masses:

$$\mathcal{L}_Y \sim \sum_f m_f \bar{f} f (1 + H/v) \quad Y_f = \frac{\sqrt{2} m_f}{v}$$

- Technically “natural” against quantum corrections:

$$\delta m_f \sim m_f \ln(\Lambda/m_w) \quad (\text{chiral symm})$$



Higgs is responsible for our existence!

- Atoms/chemistry/biology governed by $Y_e \sim m_e$:



- Y_t / m_t : not too large for vacuum stability!

$M_{W,Z}$ versus m_H versus m_f :

(3). m_f : the ugly!



- Vastly different hierarchical masses
- ad hoc flavor mixings and the CPv phase(s)
- Neutrino masses: Dirac vs. Majorana?

BSM: **much harder** to accommodate!

- Generate multiple mass scales
- Avoid FCNC
- Avoid Excessive CP violation
- Why the flavor mixing aligned with the SM Yukawa form?
→ Minimal Flavor Violation (MFV)

(B). YUKAWA COUPLINGS:

(1). Generate flavor hierarchies

- **Horizontal flavor symmetry:** Froggatt-Nielsen mechanism
SM fermions charged $[q_i, u_i, d_i]$ under $U(1)_{FN}$ symmetry
broken by $\langle\phi\rangle/M \sim 0.2$ Froggatt & Nielsen (1979)

$$(Y_u)_{ij} \sim \left(\frac{\langle\phi\rangle}{M}\right)^{[q_i]-[u_j]}, \quad (Y_d)_{ij} \sim \left(\frac{\langle\phi\rangle}{M}\right)^{[q_i]-[d_j]}$$

- **Warped extra-dimension:**
Yukawa couplings determined by
the overlapping with the Higgs brane.
→ dual to (partial) composite model.

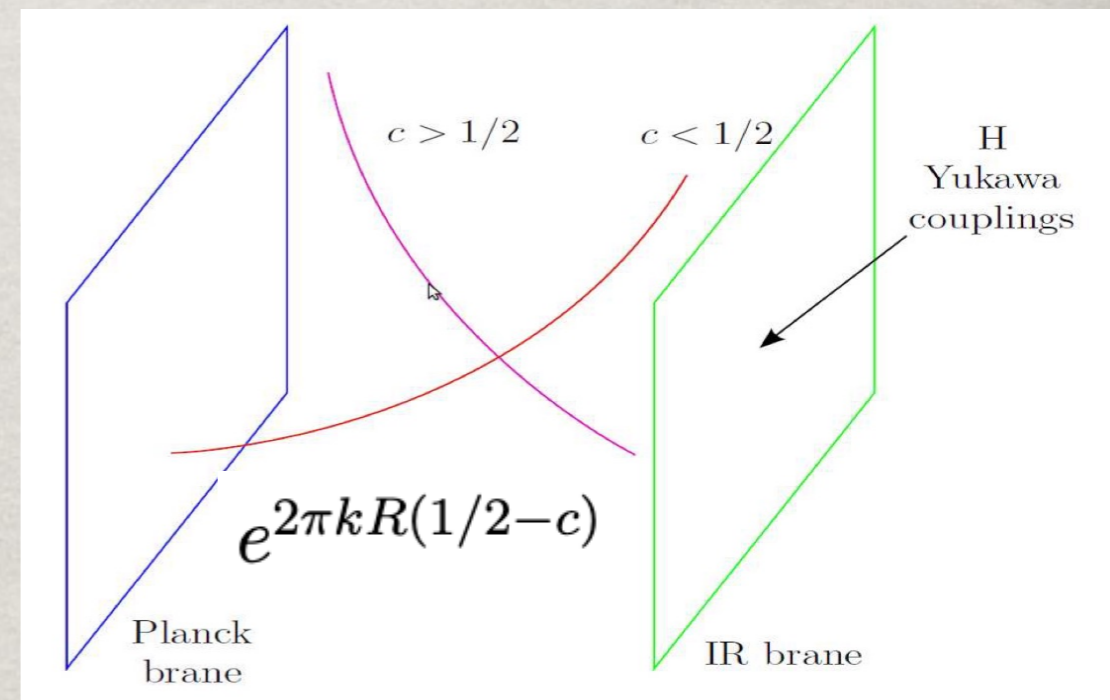
Randall & Sundrum (1999); Huber & Shafi (2001);
Agashe et al. (2005)

- **Radiative generation** of m_f :

The 3rd generation @ tree-level

Light generations by new particle loops $\sim 1/16\pi^2 \sim 10^{-2}$.

S. Weinberg (1972)



(2). The Higgs sector extension

- 2HDM (MSSM): well-motivated
($\tan\beta = v_2/v_1$; α the neutral Higgs mixing)

Talks in HPNP2023:
Pilaftsis; Haber; Takeuchi;
Yagyu; Ferreira; De Curtis;
Song; Dey; Kartayama;
Heinemeyer; Muta; Ivanov;
Sakurai

	Tree-level Normalized Higgs couplings			
	κ_h^u	κ_h^d	κ_h^e	κ_h^V
Type-I	$\frac{\cos\alpha}{\sin\beta}$	$\frac{\cos\alpha}{\sin\beta}$	$\frac{\cos\alpha}{\sin\beta}$	$\sin(\beta - \alpha)$
Type-II	$\frac{\cos\alpha}{\sin\beta}$	$-\frac{\sin\alpha}{\cos\beta}$	$-\frac{\sin\alpha}{\cos\beta}$	$\sin(\beta - \alpha)$
Type-L	$\frac{\cos\alpha}{\sin\beta}$	$\frac{\cos\alpha}{\sin\beta}$	$-\frac{\sin\alpha}{\cos\beta}$	$\sin(\beta - \alpha)$
Type-F	$\frac{\cos\alpha}{\sin\beta}$	$-\frac{\sin\alpha}{\cos\beta}$	$\frac{\cos\alpha}{\sin\beta}$	$\sin(\beta - \alpha)$

Decoupling/
Alignment limit:

$$\kappa's \rightarrow 1$$

H. Haber & Y. Nir (1990)

- Plus a singlet S (NMSSM):
more mixing & flavor physics, connect to dark sector

- Add a triplet Φ (Type-II seesaw):
 $\phi^{\pm\pm}, \phi^{\pm}, \phi^0$; connect to neutrino Majorana mass

For a review, see, i.e., G.C. Branco, M. Sher et al., arXiv:1106.0034 ...

(3). SMEFT

SM Effective Field Theory: a linear representation

$$\text{SM-like Higgs } \Phi = \frac{1}{\sqrt{2}} \begin{pmatrix} \sqrt{2}\phi^+ \\ v + h + i\phi^0 \end{pmatrix}$$

$$\mathcal{L}_Y \sim \sum_{n=0} \frac{Y_{ij}^n}{\Lambda^{2n}} (\Phi^\dagger \Phi)^n \bar{L}_{iL} \Phi e_{jR} \rightarrow m_f = \frac{v}{\sqrt{2}} \sum_{n=0} Y_n^f \frac{v^{2n}}{\Lambda^{2n}}$$

Yukawa coupling deviates from the mass relation!

At the dim-6 leading order:

$$\rightarrow \delta\kappa_f \sim Y_1 \frac{v^2}{\Lambda^2} \sim O(\text{a few}\%) \text{ for } \Lambda \sim 2 \text{ TeV!}$$

This is the immediate target @ LHC!

(4). HEFT

Higgs Effective Field Theory: a non-linear representation

$$U = e^{i\phi^a \tau_a / v} \quad \text{with} \quad \phi^a \tau_a = \sqrt{2} \begin{pmatrix} \frac{\phi^0}{\sqrt{2}} & \phi^+ \\ \phi^- & -\frac{\phi^0}{\sqrt{2}} \end{pmatrix}$$
$$L_Y \sim -\frac{v}{2\sqrt{2}} \left[\sum_{n \geq 0} y_n \left(\frac{H}{v} \right)^n (\bar{\nu}_L, \bar{\mu}_L) U (1 - \tau_3) \begin{pmatrix} \nu_R \\ \mu_R \end{pmatrix} + \text{h.c.} \right]$$

$$Y_f(H) = \frac{\sqrt{2}m_f}{v} + \sum_{n=1} y_{fn} \left(\frac{H}{v} \right)^n$$

- The scale for new dynamics is at $\Lambda \sim 4\pi v$
 \rightarrow close by! The deviation can be sizable:

$$\rightarrow \delta\kappa_f \sim Y_1 \frac{H}{v} \sim O(1)$$

- Multiple Higgs couplings may be sizeable!

(C). HIGGS COUPLINGS @ COLLIDERS

theories in practice

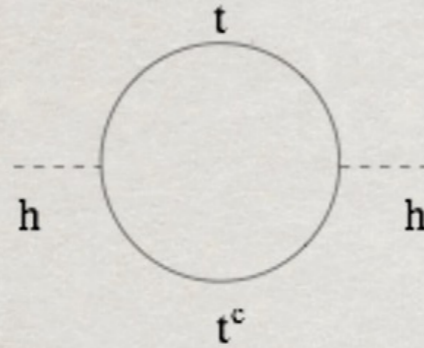
- Seek for more Higgs bosons: H^0, A^0, H^\pm ,
 - Continue to search for Z', W^\pm, T' in light of FN, L-R, composite H, etc.
 - Flavor Changing Neutral Currents:
 $H \rightarrow tc, tq; \& \mu\tau, e\mu, e\tau!$
- & possible new CPv phases in Yukawa ...

Cheng & Sher, PRD (1987); G.W.S. Hou, PLB (1992);
TH, D. Marfatia, PRL 86, 1442 (2001); Harnik, Kopp, Zupan, arXiv:1209.1397.

In the absence of the new physics discovery:

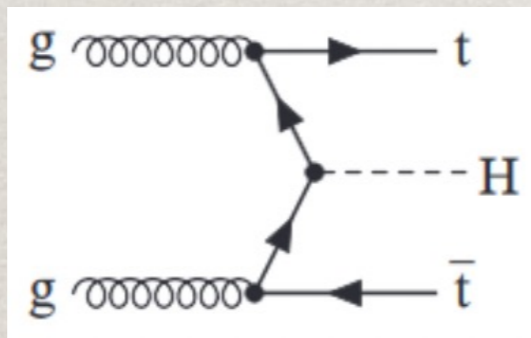
- Exploring flavor physics is complementary & rewarding.
- Measuring Higgs Yukawa couplings is indispensable:
the smaller the coupling is, the more sensitive to deviations!

(1). The most-wanted: y_t



$$m_H^2 = m_{H0}^2 - \frac{3}{8\pi^2} y_t^2 \Lambda^2$$

The current LHC sensitivity:



$$\delta\kappa_t = 0.35^{+0.36}_{-0.34} \quad (\text{ATLAS})$$

(ICHEP 2022)

Also, see Sanmay Ganguly & Xiaohu Sun

Future lepton collider sensitivity:

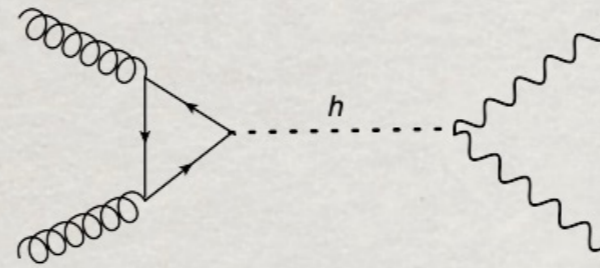
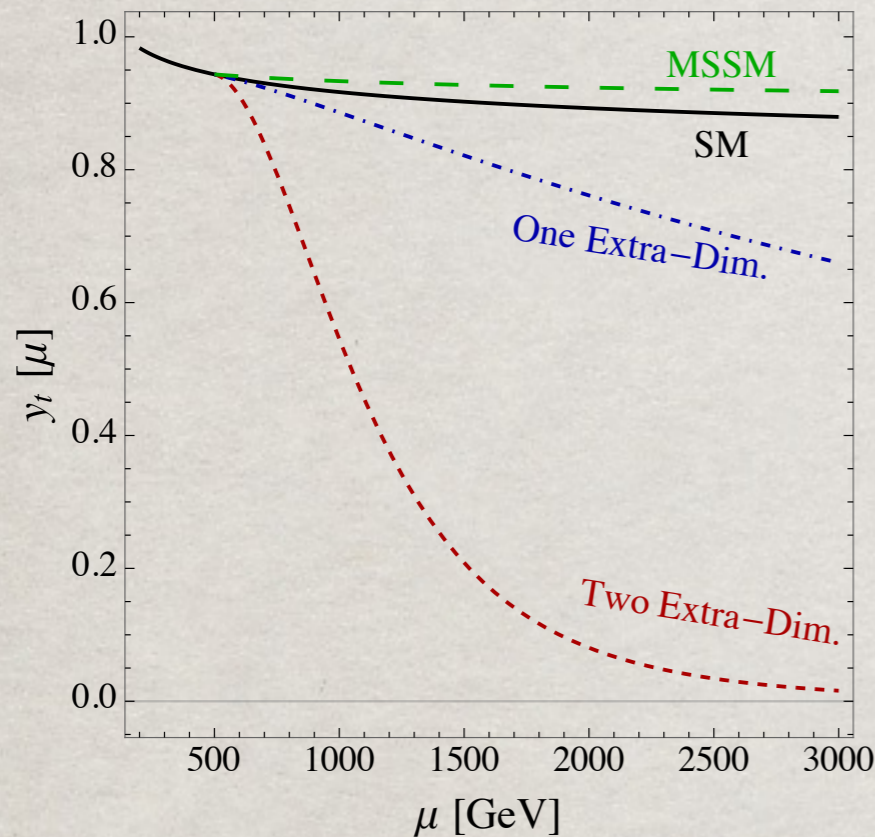
Values in % units		LHC	HL-LHC	ILC500	ILC550	ILC1000	CLIC
δy_t	Global fit	6.12	2.53	2.08	1.30	0.739	1.48
	Indiv. fit	5.08	1.85	1.80	1.17	0.705	1.26

Table 8: Uncertainties for the top-quark Yukawa coupling at 68% probability for different scenarios, in percentage. The ILC500, ILC550 and CLIC scenarios also include the HL-LHC. The ILC1000 scenario includes also ILC500 and HL-LHC.

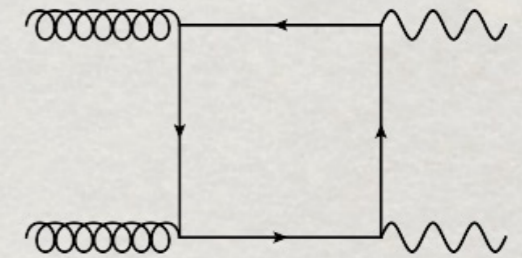
EF04 report: <https://arxiv.org/pdf/2209.08078.pdf>

ttH coupling @ high scales:

1. Yukawa $y_t(Q)$ RGE running: 2. Off-shell probe of EFT operators:



$$\sigma_{\text{on}} \propto \frac{g_i^2(m_h^2)g_f^2(m_h^2)}{m_h\Gamma_h} \quad \text{and} \quad \sigma_{\text{off}} \propto \frac{g_i^2(Q^2)g_f^2(Q^2)}{Q^2}$$



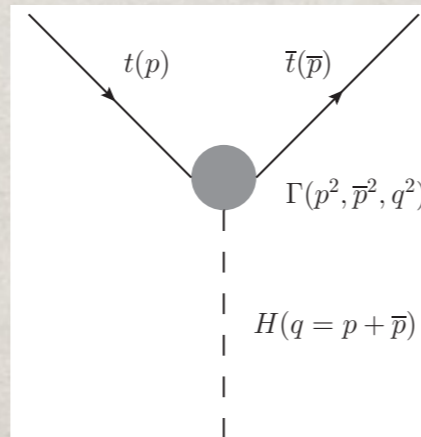
CMS/ATLAS: [arXiv:2202.06923](https://arxiv.org/abs/2202.06923); [2304.01532](https://arxiv.org/abs/2304.01532)

- $3.6\sigma/3.3\sigma$ observation for off-shell Higgs signal
- SM width bound: $3.2^{+2.4}_{-1.7} / 4.5^{+3.3}_{-2.5} \text{ MeV}$

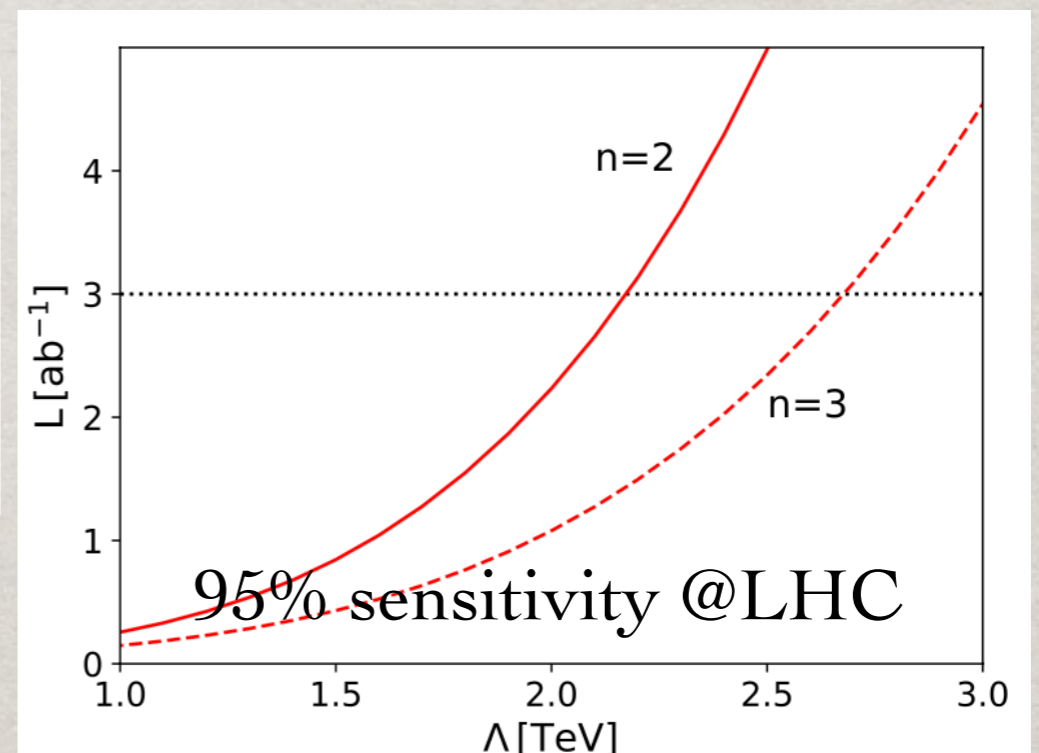
(See Sanmay Ganguly & Xiaohu Sun)

3. Composite form factors:

$$\Gamma(q^2/\Lambda^2) = \frac{1}{(1 + q^2/\Lambda^2)^n}$$



D. Goncalves, TH, S. Mukhopadhyay, [arXiv:1710.02149](https://arxiv.org/abs/1710.02149) (PRL, 2017); [arXiv:1803.09751](https://arxiv.org/abs/1803.09751);
 D. Goncalves, TH, I. Leung, H. Qin, [arXiv:2012.05272](https://arxiv.org/abs/2012.05272);
 R. Abraham, D. Goncalves, TH, S.C.I. Leung, H. Qin, [arXiv:2012.05272](https://arxiv.org/abs/2012.05272).



(2). The real challenge: y_c

Test the 2nd generation couplings:

The current LHC sensitivity: $BR_{H \rightarrow c\bar{c}}^{\text{SM}} = (2.88^{+0.16}_{-0.06})\%$

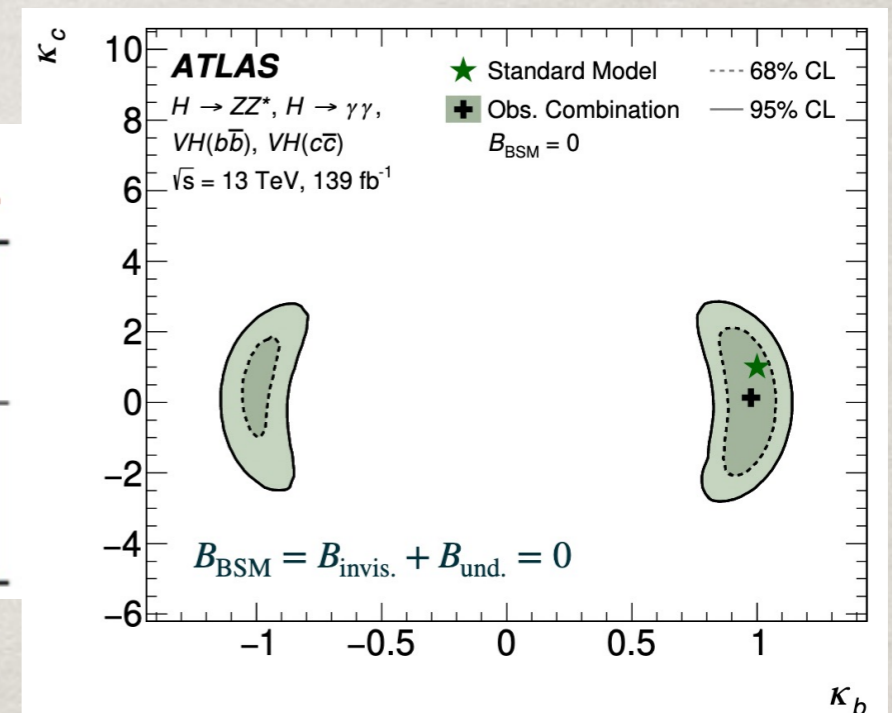
LHC Run 2: ATLAS $\kappa_c \leq 8.5$ [2201.11428], CMS $1.1 < |\kappa_c| < 5.5$ [2205.05550]

(See Sanmay Ganguly & Xiaohu Sun)

Submitted to JHEP [2207.08615]

Upper limit on κ_c of $4.8 \times \text{SM}$ at 95% CL

Scenario	Observed 68% confidence interval	Observed 95% confidence interval
$B_{\text{BSM}} = 0$	[-1.61, 1.70]	[-2.47, 2.53]
No assumption	[-2.63, 3.01]	[-4.46, 4.81]



HL-LHC sensitivity projection: a factor of few from SM

Future HL-LHC: $\kappa_c \leq 3$. [2201.11428]

EF01/02 report: <https://arxiv.org/pdf/2209.07510.pdf>

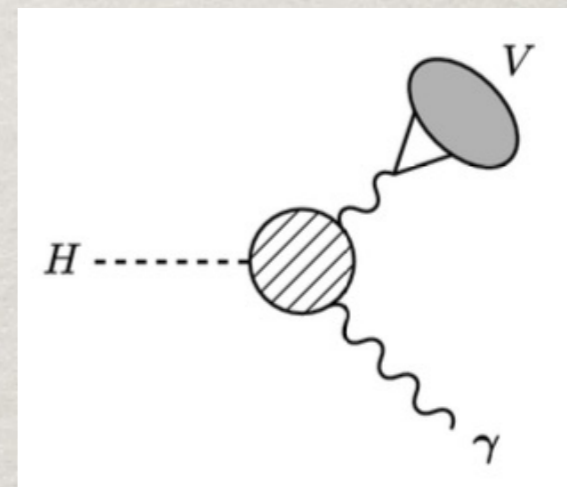
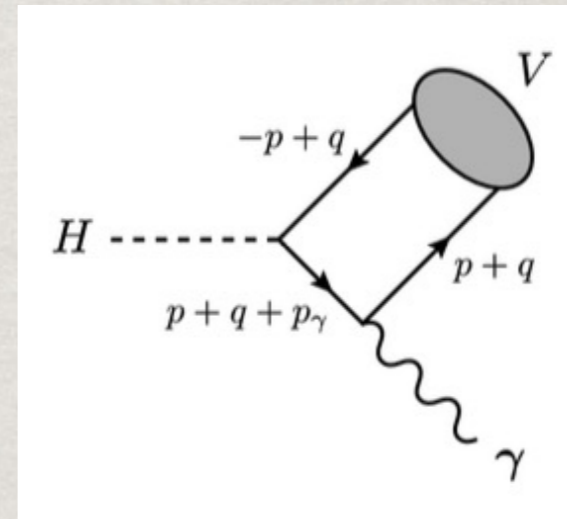
Higgs production rate is high: $\#H@LHC \sim 50 \text{ M /ab} !$
 Need new ideas!

- $H \rightarrow J/\psi + \gamma$
 $\rightarrow \mu^+ \mu^- + \gamma$

for Higgs coupling to charm

Note: $BR(H \rightarrow J/\psi + \gamma) = 2.8 \times 10^{-6}$

➤ Dominated by VMD $\gamma^* \rightarrow J/\psi$,
 not $H cc$ coupling.



➔ This is no use to probe $y_c !$

Bodwin, Petriello et al. (2013, 2014, 2017); Konig, Neubert (2015)

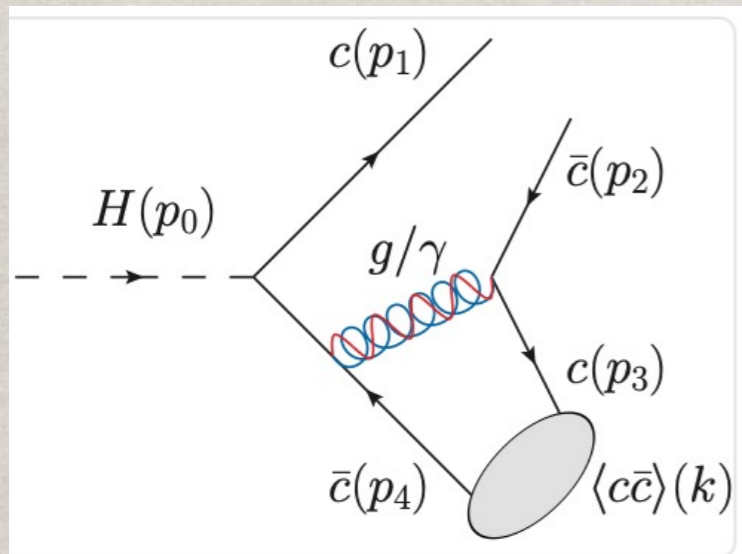
Higgs production rate is high: #H@LHC ~ 50 M /ab ! A new idea!

- $H \rightarrow J/\psi$ via charm-quark fragmentation:

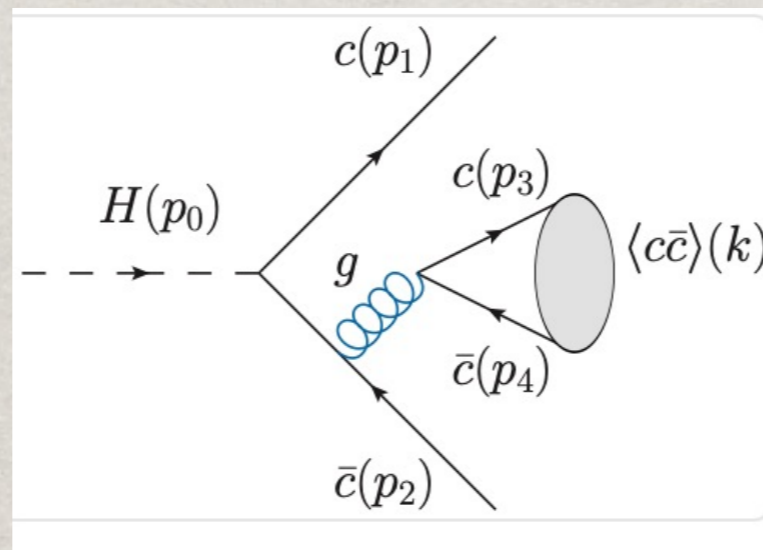
$$H \rightarrow c + \bar{c} + J/\psi \text{ (or } \eta_c)$$

- Enhanced from the fragmentation
- Direct coupling to charm!

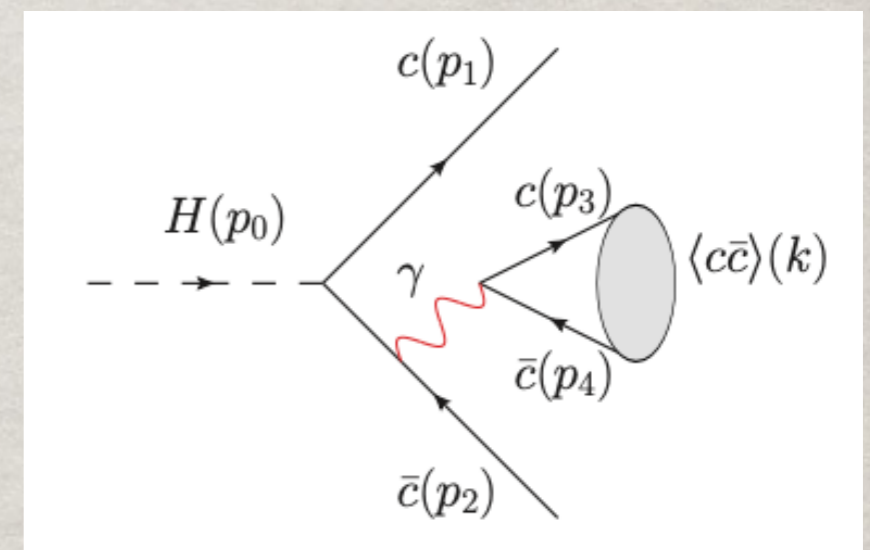
Color-singlet (CS)
(leading)



Color-octet (CO)
(sub-leading, 1/2 of CS)



QED
(sub-leading, 1/4 of CS)



TH, A. Leibovich, Y. Ma, X.Z. Tan: aXive:2202.08273

- $H \rightarrow J/\psi$ via charm-quark fragmentation:

$$H \rightarrow c + \bar{c} + J/\psi \text{ (or } \eta_c)$$

TH, A. Leibovich, Y. Ma, X.Z. Tan:
aXive:2202.08273

$$\Gamma = \sum_{\mathbf{N}} \hat{\Gamma}_{\mathbf{N}}(H \rightarrow (Q\bar{Q})[\mathbf{N}] + X) \times \langle \mathcal{O}^h[\mathbf{N}] \rangle,$$

$$d\hat{\Gamma}_{\mathbf{N}} = \frac{1}{2m_H} \frac{|\mathcal{M}|^2}{\langle \mathcal{O}^{Q\bar{Q}} \rangle} d\Phi_3$$

Calculating the short-distance decay rates,
fit the long-distance hadronic matrix elements,
we obtain:

$$\text{BR}(H \rightarrow c\bar{c} + J/\psi) = (2.0 \pm 0.5) \times 10^{-5}$$

$$\text{BR}(H \rightarrow c\bar{c} + \eta_c) = (6.0 \pm 1.0) \times 10^{-5}$$

- ▶ Sensitivity $S \simeq N_{\text{signal}} / \sqrt{N_{\text{Background}}}$
 \Rightarrow It is possible to reach 2σ for $\kappa_c \approx 2.4$.

\rightarrow At the end, should be better than $J/\psi + \gamma$: $\kappa_c \sim 50$

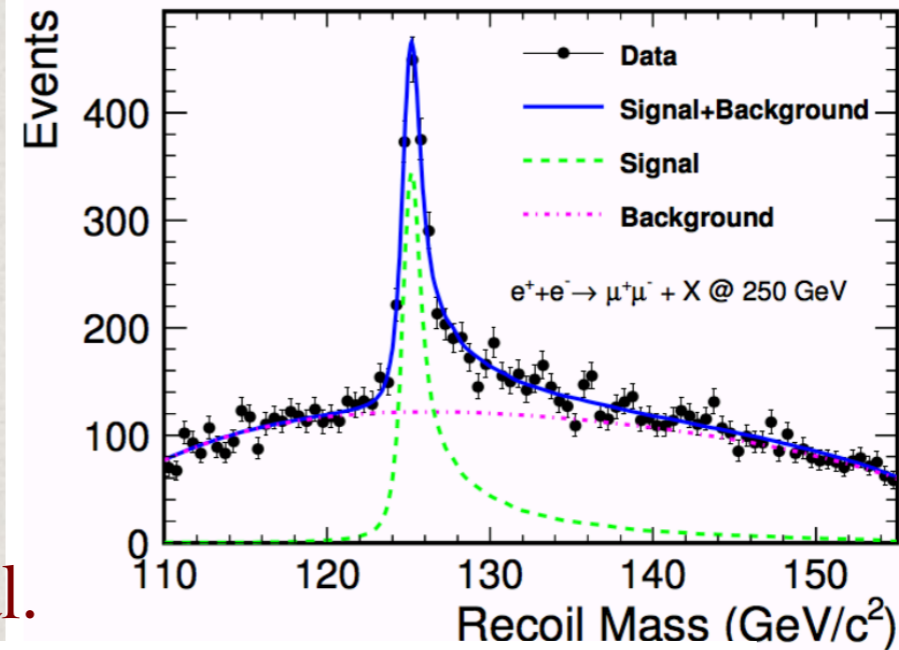
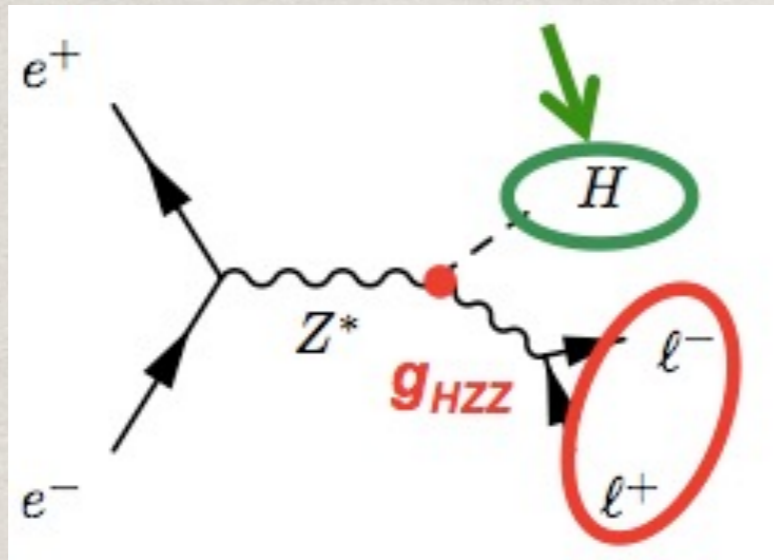
\rightarrow May not beat $W/Z+H \rightarrow W/Z+cc$: $\kappa_c \sim 3$

Active study/simulation on-going!

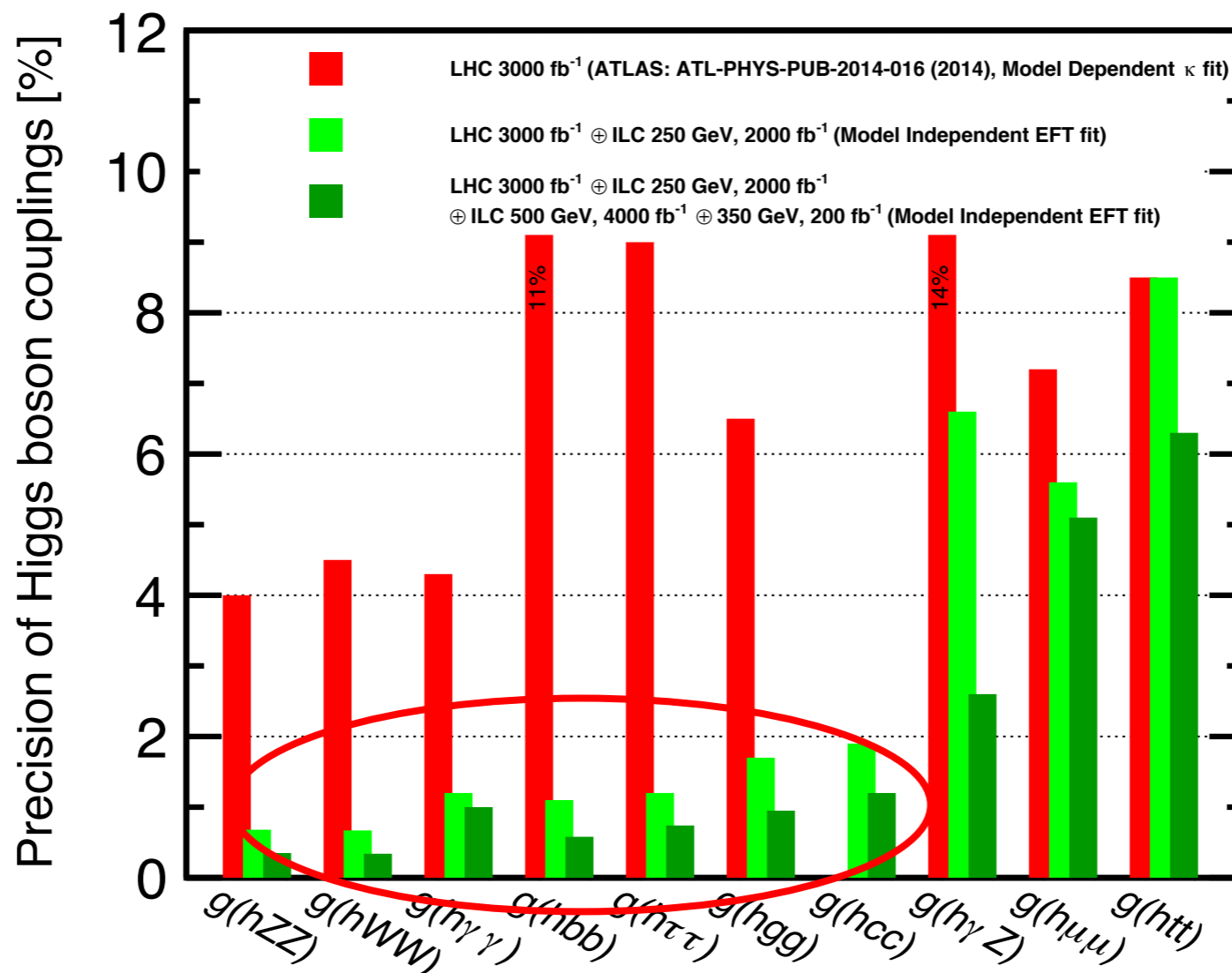
ILC as the Higgs factory

See Junping Tian's talk

$$m_h^2 = (p_{e^+}^2 + p_{e^-}^2 - p_{\ell^+}^2 - p_{\ell^-}^2)^2$$



arXiv:1710.07621, LCC, Fujii et al.



(3). A promising channel: y_μ

The 2nd generation Υ_μ : The next hope!

The current LHC sensitivity: $BR_{H \rightarrow \mu^+ \mu^-}^{\text{SM}} = (2.17 \pm 0.04) \times 10^{-4}$

Current search result:

ATLAS: 2.0σ ; CMS 3.0σ

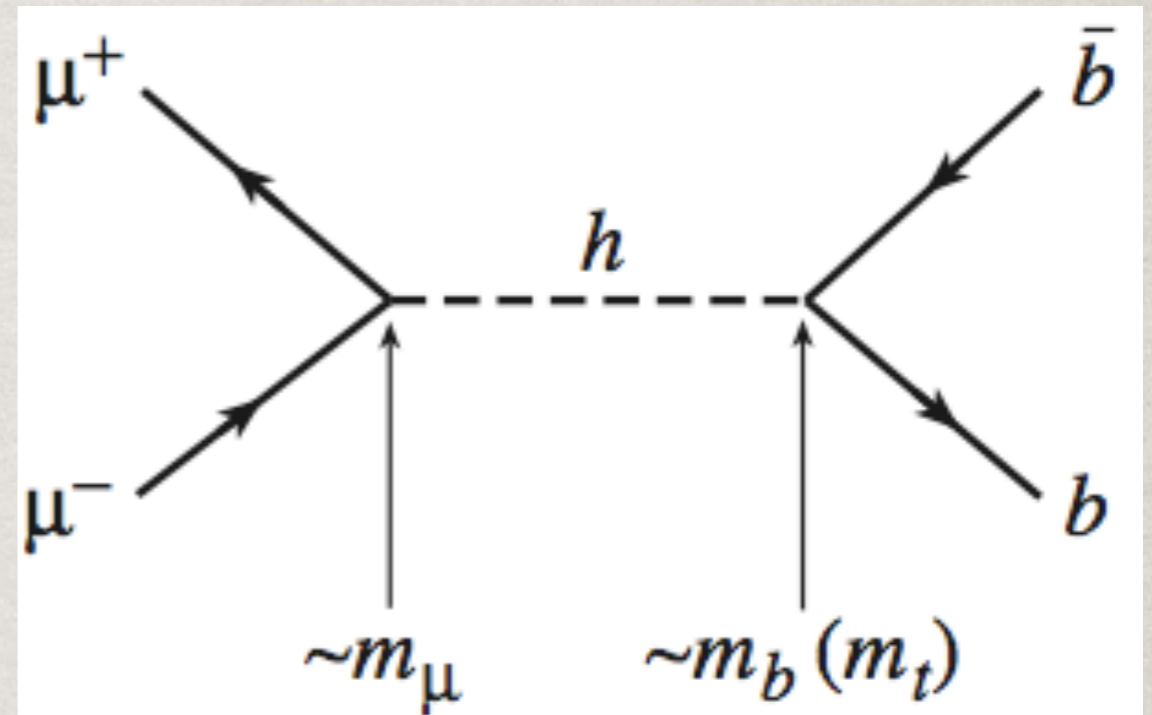
HL-LHC sensitivity projection: $BR(H \rightarrow \mu\mu) < 10\%$

- Assuming the SM width, but won't know it better than a factor of 2-ish
- ILC may not improve this much (low rate)

*Plehn & Rainwater, arXiv:hep-ph/0107180;
TH & McElrath, arXiv:hep-ph/0201023*

(4). Model-independent measurement on y_μ

A muon collider Higgs factory: Resonant Production:



$$\sigma(\mu^+ \mu^- \rightarrow h \rightarrow X) = \frac{4\pi \Gamma_h^2 \text{Br}(h \rightarrow \mu^+ \mu^-) \text{Br}(h \rightarrow X)}{(\hat{s} - m_h^2)^2 + \Gamma_h^2 m_h^2}.$$

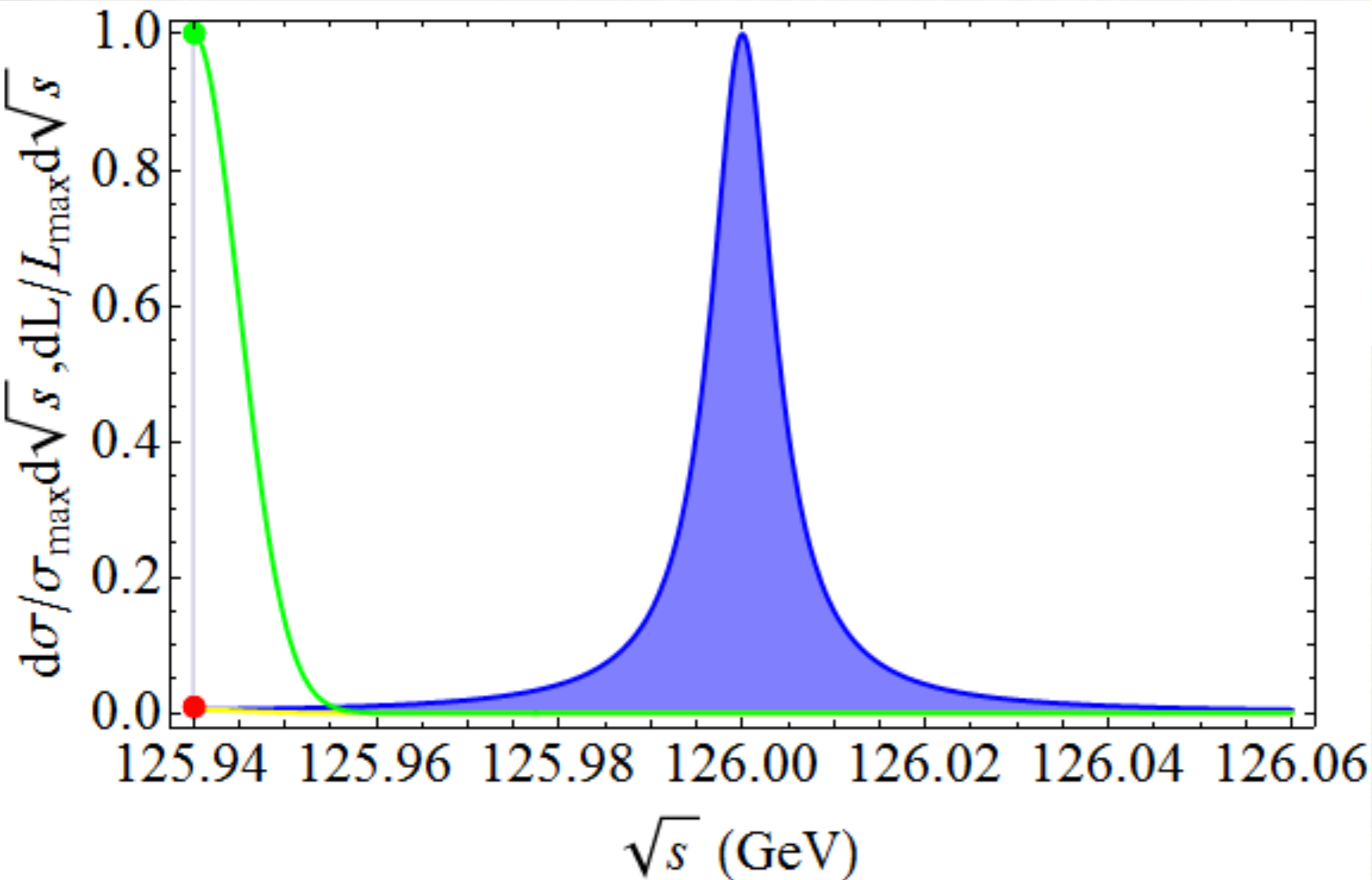
$$\begin{aligned} \sigma_{peak}(\mu^+ \mu^- \rightarrow h) &= \frac{4\pi}{m_h^2} \text{BR}(h \rightarrow \mu^+ \mu^-) \\ &\approx 71 \text{ pb at } m_h = 125 \text{ GeV.} \end{aligned}$$

About $\mathbf{O(70k)}$ events produced per $\mathbf{fb^{-1}}$

Requirement: $E_{\text{cm}} = m_H$, $\Delta E_{\text{cm}} \sim 5 \text{ MeV}$, $L \sim 1 \text{ fb}^{-1}/\text{yr}$.

Ideal, conceivable case:

$$(\Delta = 5 \text{ MeV}, \quad \Gamma_h \approx 4.2 \text{ MeV})$$



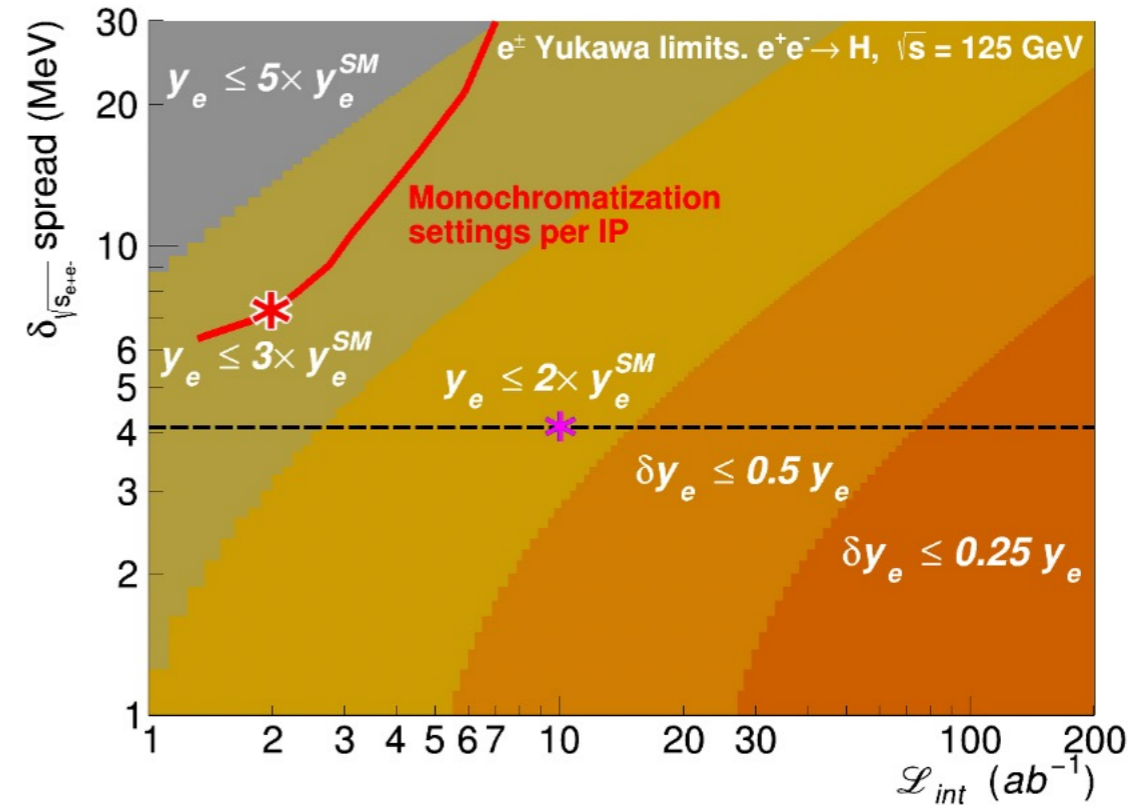
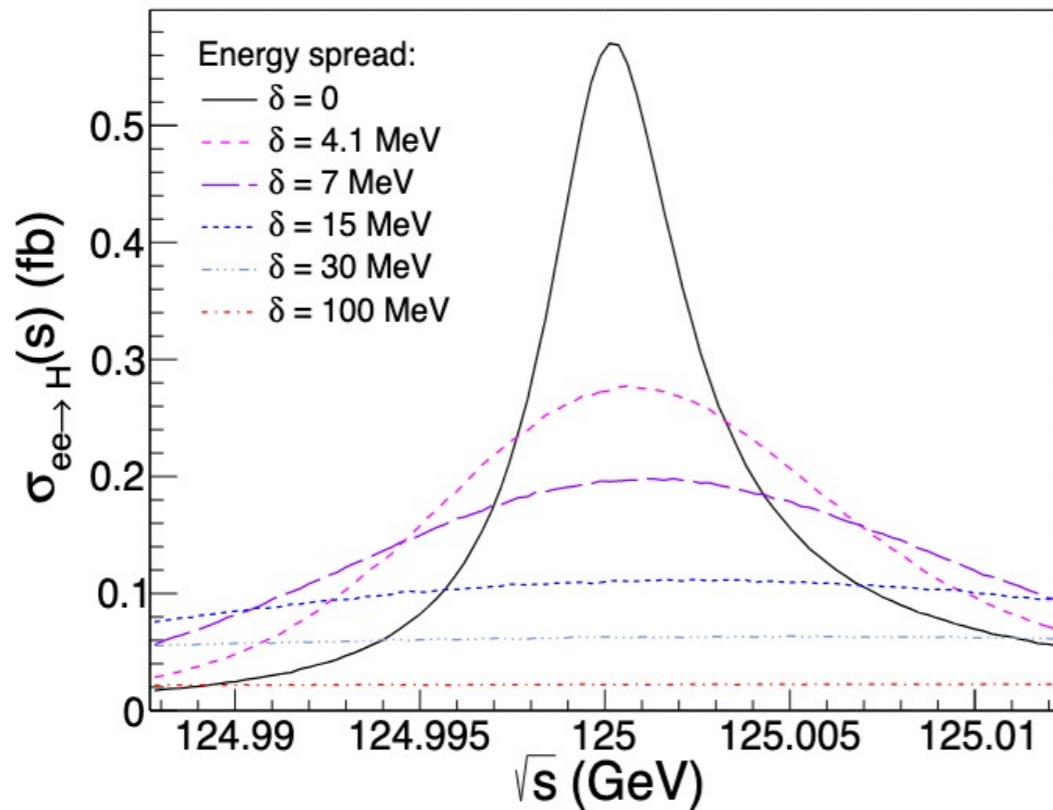
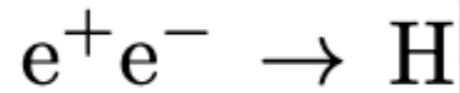
An optimal fitting could reach $\delta\Gamma_h \sim 0.15 \text{ MeV}$, or 3.5%

TH, Liu: 1210.7803; Greco, TH, Liu: 1607.03210

(5). The ultimate test : y_e

The 1st generation y_e : There is a chance!

$$y_e = \sqrt{2}m_e/v = 2.9 \times 10^{-6} \quad \text{SM BR}(H \rightarrow ee) \sim 5 \cdot 10^{-9}$$

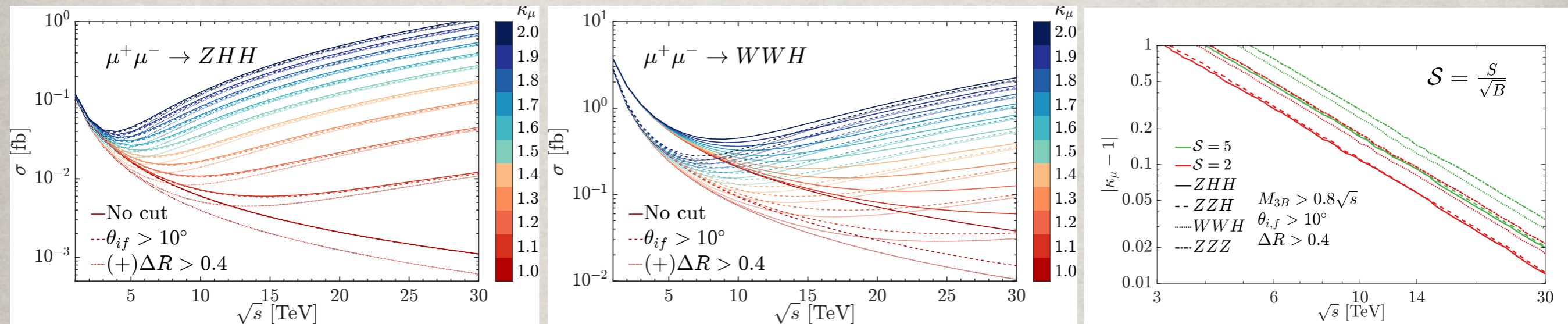
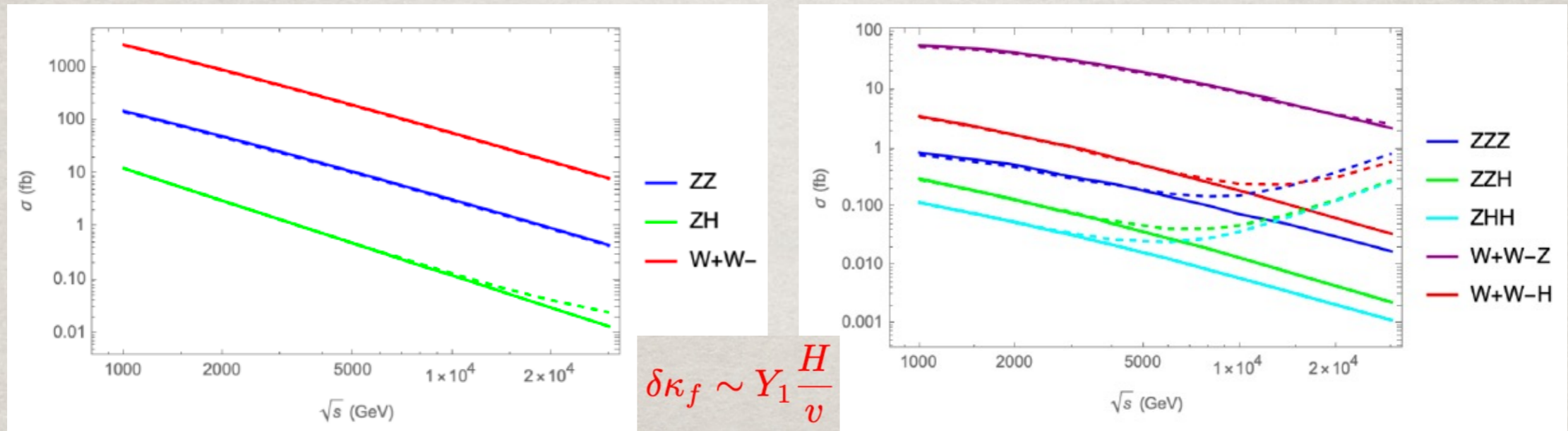


M. Greco, TH, Z. Liu:
<https://arxiv.org/abs/1607.03210>

Accel. Frontier report:
<https://arxiv.org/pdf/2203.06520.pdf>

(6). High-energy option on y_μ

To enhance the Yukawa coupling effects, multiple Higgs/Goldstone boson production more beneficial.



At 30 TeV: $\delta\kappa_\mu \sim 1\% - 4\%$, corresponding to $\Lambda \sim 30 \text{ TeV} - 100 \text{ TeV}$.

TH, W. Kilian, N. Kreher, Y. Ma, J. Reuter, T. Striegl, K. Xie: <https://arxiv.org/abs/2108.05362>;

E. Celada, TH, W. Kilian, N. Kreher, Y. Ma, F. Maltoni, D. Pagani, J. Reuter, T. Striegl, K. Xie; to appear.

Sensitivities to Yukawa couplings at future colliders

Achieving percentage/sub-percentage level!

Symbols of sensitivities:

		y_u	y_d	y_s	y_c	y_b	y_t	y_e	y_μ	y_τ
EF benchmarks										
Higgs + HL-LHC	LHC/HL-LHC	□	□	□	◇	◇	◇	□	◇	◇
	ILC/C ³ 250	□	□	□*	◇	◇	◇	□	◇	◇
	CLIC 380	□	□	?	◇	◇	◇	□	◇	◇
	FCC-ee 240	□	□	?	◇	◇	◇	□	◇	◇
	CEPC 240	□	□	?	◇	◇	◇	□	◇	◇
	μ -Collider	□	□	?	◇	★	◇	□	◇	◇
	FCC-hh/SPPC	?	?	?	?	◇	◇	?	◇	◇

EF01/02 report: <https://arxiv.org/pdf/2209.07510.pdf>

Conclusions:

- Higgs couplings to fermions most mysterious: least understood in theory, but rich phenomenology!
- Continue to search for more Higgses, Z' , T' ...
- Must look for rare Higgs decays: flavor changing, invisible channels ...
- Measuring Higgs Yukawa couplings: indispensable
- SMEFT sets a target: $\delta\kappa_f \sim Y_1 \frac{v^2}{\Lambda^2} \sim O(\text{a few}\%)$
- HEFT could be close by: $\delta\kappa_f \sim Y_1 \frac{H}{v} \sim O(1)$
- Immediate targets on Yukawa couplings:
 ttH @high scale; 2nd generations $H\mu\mu$ & Hcc !

Push for the next discovery for NP from HP!