

# Comparison of Harmonic Spin Matching Schemes using Orbit Bumps in the FCC-ee

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Swiss Accelerator  
Research and  
Technology

- FCC-ee will operate on 4 centre-of-mass energies  
 $Z^0$  bosons (91 GeV), WW pairs (160 GeV), Higgs bosons (240 GeV) and top quark pairs (350-365 GeV)
- High-precision centre-of-mass energy calibration
  - basis for precise measurements of the standard model particle properties
  - make it possible for the new rare process detection
  - precise measurements in FCC-ee will contribute to the measurements in FCC-hh

The current precision targets for the energy calibration:  
4 keV at Z mass and 100 keV at W mass

the most promising way to achieve this target: **resonant depolarization**

## **Ensure a sufficient spin polarization level (at least 5 – 10%)**

1. Estimate the achievable polarization under various lattice conditions  
e.g. misalignments+field errors
2. Use special structure to improve polarization  
e.g. closed orbit bumps

- $\hat{n}_0(s)$ : one-turn periodic solution of the T-BMT equation on closed orbit
- Spins on the closed orbit precess around  $\hat{n}_0$  for  $\nu_0$  turns in every revolution  
 $\Rightarrow \nu_0$ : closed orbit spin tune
- $\nu_0 = a\gamma$  in the perfectly aligned flat ring without solenoids  
 $\nu_0 \approx a\gamma$  in misaligned lattice  
 $|\nu_0 - a\gamma|$  impacts the measurement precision
- ST effect + radiative depolarization  $\Rightarrow$  equilibrium polarization

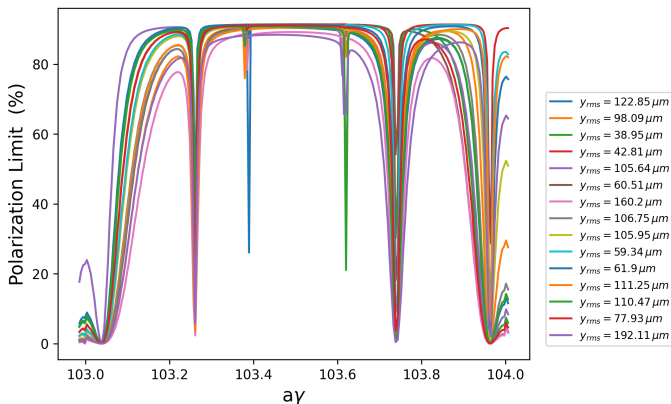
- Use an effective model to simulate **residual orbits after lattice correction**
- Random small errors generated from truncated Gaussian distributions (truncated at  $2.5\sigma$ )

Type	$\sigma_{\Delta X}$ (nm)	$\sigma_{\Delta Y}$ (nm)	$\sigma_{\Delta Z}$ (nm)	$\sigma_{\Delta \text{PSI}}$ ( $\mu\text{rad}$ )	$\sigma_{\Delta \text{THETA}}$ ( $\mu\text{rad}$ )	$\sigma_{\Delta \text{PHI}}$ ( $\mu\text{rad}$ )
Arc quadrupole	120	120	120	2	2	2
Arc sextupole	120	120	120	2	2	2
Dipoles	120	120	120	2	0	0
IR quadrupole	120	120	120	2	2	2
IR sextupole	120	120	120	2	2	2

## Setting 1

$\sigma = 120 \text{ nm}$  for x,y,z misalignments

$\sigma = 2 \mu\text{rad}$  for angular deviations

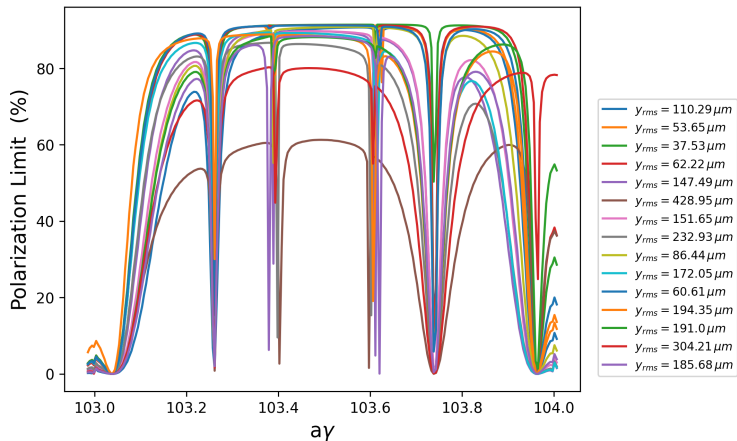


First order energy scan showing equilibrium polarization levels near Z energy

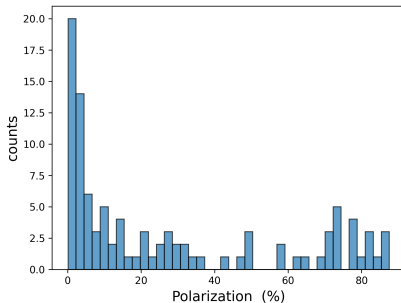
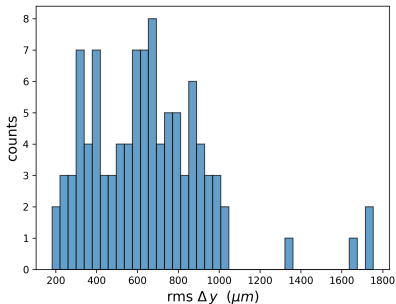
## Setting 2

$\sigma = 200$  nm for x,y,z misalignments

$\sigma = 2$   $\mu$ rad for angular deviations

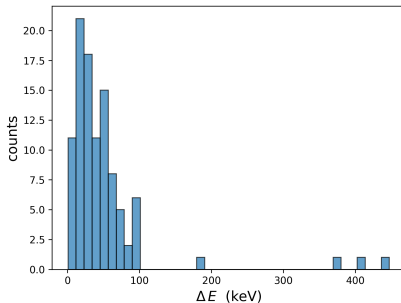
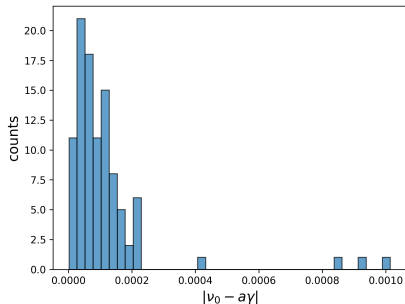


quadrupoles+dipoles:  
 $\sigma = 1 \mu\text{m}$  x,y,z misalignments  
 $\sigma = 1 \mu\text{rad}$  angular deviations





quadrupoles+dipoles:  
 $\sigma = 1 \mu\text{m}$  x,y,z misalignments  
 $\sigma = 1 \mu\text{rad}$  angular deviations



Larger error seeds ( $100\ \mu\text{m}$ ) + orbit correction (correctors and BPMs)

How much polarization?

**Ongoing!**

## Conventional lattice correction + harmonic spin matching

- Misaligned ring  $\rightarrow \hat{n}_0(s)$  not vertical  $\rightarrow$  stronger spin diffusion
- Random vertical quadrupole misalignments are difficult to control
- HSM: use multiple vertical orbit correctors to create an additional controllable  $\hat{n}_0$  tilt and reduce  $(\delta \hat{n}_0)_{\text{rms}} \rightarrow$  elevate polarization
- Use closed vertical bumps to avoid disturbing the orbits outside bumps

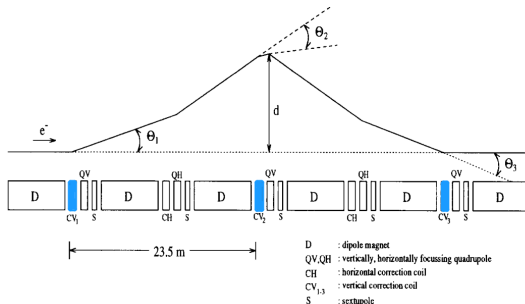
# Harmonic bump composition

1st corrector: give a vertical kick

2nd corrector: kick  $y$  back to initial value at the 3rd corrector

3rd corrector: kick  $y'$  back to initial value

Kicks of 2nd and 3rd correctors are adjusted to make the bump **CLOSED**  
**one independent variable for each bump**



a schematic of a bump in HERA\*

\*D. P. Barber, et al. "High spin polarization at the HERA electron storage ring." Nuclear Instruments and Methods in Physics Research Section A: Accelerators, Spectrometers, Detectors and Associated Equipment 338.2-3 (1994): 166-184.

## 1. HERA formalism (used in HERA)

D. P. Barber, et al. A general harmonic spin matching formalism for the suppression of depolarisation caused by closed orbit distortion in electron storage rings. No. DESY-85-044. DESY, 1985.

## 2. Rossmanith-Schmidt scheme (used in PETRA)

R. Rossmanith and R. Schmidt, Compensation of depolarizing effects in electron-positron storage rings. Nuclear Instruments and Methods in Physics Research Section A: Accelerators, Spectrometers, Detectors and Associated Equipment 236.2 (1985): 231-248.

## 3. LEP method (Deterministic) (used in LEP)

R. W. Assmann, Optimierung der transversalen Spin-Polarisation im LEP-Speicherring und Anwendung für Präzisionsmessungen am Z-Boson. Diss. Munich U., 1994.

$$\delta \hat{n}_0 = \alpha \hat{m} + \beta \hat{l}$$

$\hat{l}(s)$ ,  $\hat{n}_0(s)$ ,  $\hat{m}(s)$  are periodic spin axes that form a right-hand coordinate system. Expand  $\alpha$  and  $\beta$  to Fourier series

$$(\alpha - i\beta)(s) = -i \frac{C}{2\pi} \sum_k \frac{f_k}{k - \tilde{\nu}} e^{i2\pi ks/C}$$

$f_k$ : Fourier coefficients, related to the closed orbit and perturbing fields

Make additional orbit corrections using orbit bumps to reduce the rms tilt by minimizing the Fourier coefficients

- Extract  $n_0$  direction at the end of all elements
- Expand  $n_{0x}(s) + in_{0z}(s)$  into Fourier series
- Minimize target coefficients using four bumps

$$n_{0x} + in_{0z} \approx \sum_{k=-N}^N c_k \cdot e^{i2\pi ks/C}$$

$$c_k = \frac{1}{C} \int (n_{0x} + in_{0z}) \cdot e^{-i2\pi ks/C} ds \approx \frac{\sum_{j=1}^M [n_{0x}(s_j) + in_{0z}(s_j)] \cdot e^{-i2\pi ks_j/C}}{M}$$

Each closed bump can be represented using a single variable, and each bump has an independent and linear contribution to the Fourier coefficients

$$\mathbf{MK} = \mathbf{C}$$

**K**: amplitudes (the first kick value) of the bumps

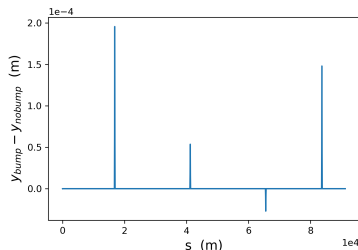
**C**: real and imaginary parts of the required harmonics coefficients

$$[C_{0\text{real}}, C_{0\text{imag}}, C_{1\text{real}}, C_{1\text{imag}}]$$

If the harmonics 0 and 1 of a misaligned lattice is **A**, the bumps should generate  $-\mathbf{A}$ , and bump amplitudes can be estimated via inverting matrix

$$\mathbf{K} = \mathbf{M}^{-1}(-\mathbf{A})$$





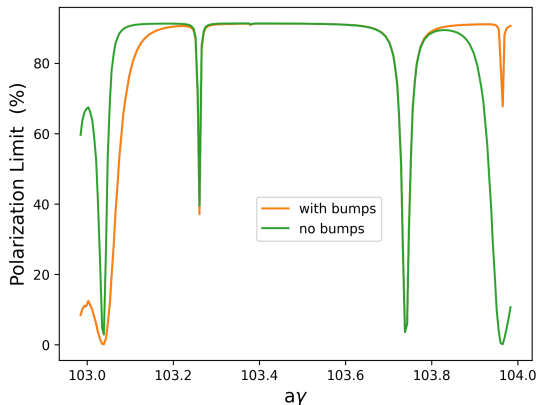
- How does orbit change?
  - Orbits outside of the bumps are unaffected.
  - max 0.2 mm  $\Delta y$  within bumps is at the level of rms orbit of the effective lattice
- How does vertical dispersion change?
  - $(\eta_y)_{\text{rms}}$  7.602 mm  $\rightarrow$  7.675 mm, 0.96% increase
  - $[\Delta\eta_y(s)]_{\text{rms}} = 0.4$  mm
- How does vertical emittance change?
  - $\varepsilon_y$  0.703 pm  $\rightarrow$  0.719 pm, 2.16% increase
- How does  $\nu_0$  shift?
  - $a\gamma = 103.983116$ ,  $\nu_0 = 103.983100158 \rightarrow 103.983100156$ , insignificant

At 45.82 GeV ( $a_\gamma = 103.983$ )

	$(\delta n_0)_{\text{rms}}$ (mrad)	Polarization (%)
no correction	2.28	10.68
a set of four random bumps	0.912	90.58
optimized four bumps	0.9	90.96
a set of eight random bumps	0.903	90.79

Not necessary to use optimized locations, but better to have a symmetric layout

Using 4 bumps which are optimized at 45.82 GeV ( $a\gamma = 103.983$ )



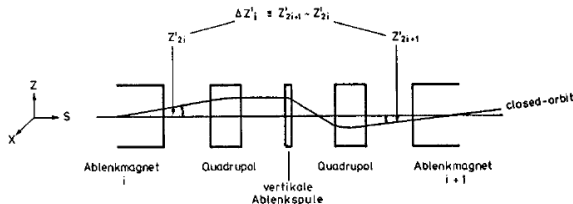
Assume that spin precessions around vertical direction only happen in bending magnets, and the radial perturbing fields on the closed orbit only exist between bending magnets

$$|\delta \vec{n}_0(s)| = \frac{1/c^2}{2(1-\cos 2\pi\nu)} \left[ \left( \int_s^{s+L} \delta \Omega_x \cos \phi ds \right)^2 + \left( \int_s^{s+L} \delta \Omega_x \sin \phi ds \right)^2 \right]$$

$$\delta \Omega_x = \frac{e}{m_0 c \gamma} (1 + a\gamma) B_x \text{ and } \phi = \gamma a \alpha$$

$$\text{also } \frac{e}{m_0 c \gamma} \int_{s_{2i}}^{s_{2i+1}} B_x(s) ds = -\Delta y'_i$$

$$|\delta \vec{n}_0| = \frac{1/c^2}{2(1-\cos 2\pi\nu)} (1 + \gamma a) \left[ \left( \sum_{i=1}^N \sin(\gamma a \alpha_i) \Delta y'_i \right)^2 + \left( \sum_{i=1}^N \cos(\gamma a \alpha_i) \Delta y'_i \right)^2 \right]$$



Spindrehung um die Z-Achse \* Spindrehung um die X-Achse \* Spindrehung um die Z-Achse

Expand  $\Delta y'(\alpha)$  into Fourier series

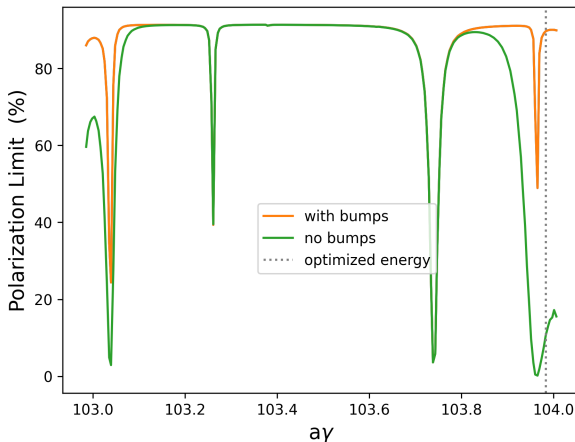
$$\Delta y'(\alpha) = \sum_{k=1}^{\infty} (a_k \cos k\alpha + b_k \sin k\alpha)$$

$$\begin{matrix} a_k \\ b_k \end{matrix} = \frac{1}{N} \sum \Delta y'_i(\alpha_i) \begin{matrix} \cos k\alpha_i \\ \sin k\alpha_i \end{matrix}$$

The harmonics which are adjacent to  $a_\gamma$  contribute most to the sum. FCC-ee (Z) operates between  $a_\gamma$  103 and 104, so that  $a/b_{103}$  and  $a/b_{104}$  are to be suppressed using four closed bumps.

45.82 GeV ( $a\gamma = 103.983$ )

$\delta n_0 : 2.28 \text{ mrad} \Rightarrow 0.90 \text{ mrad}$  ,  $P_{DK} : 10.68\% \Rightarrow 89.65\%$



$$\int B_x(s)ds \propto -\Delta y' \approx \sum_{i=1}^{N_{\text{quad}}} k_{1i} y_i L_i + \sum_{j=1}^{N_{\text{V kicker}}} \text{kick}_j$$

- Use  $y_{\text{eff}}$  to avoid the errors from thin lens approximation  
 $y_{\text{eff}} \approx \frac{1 - \cos \sqrt{k}L}{\sqrt{k}L \sin \sqrt{k}L} (y_1 + y_2)$  ( $k < 0$ ) or  $y_{\text{eff}} \approx \frac{\cosh \sqrt{k}L - 1}{\sqrt{k}L \sinh \sqrt{k}L} (y_1 + y_2)$  ( $k > 0$ )
- If two BPMs are installed at both ends of each quadrupole  
Polarization 10.68%  $\rightarrow$  84.71%
- 3056 dipoles, 1856 quadrupoles  $\rightarrow$  fewer BPMs required

Assume radial fields in quadrupoles and proportional to the beam position.  
Analyze unweighted vertical BPM readings and minimize critical harmonics

$$a_k = \frac{1}{\pi} \sum_{i=1}^{N_{BPM}} y_i \cdot \Delta\theta_i \cdot \cos(k \cdot \theta_i)$$
$$b_k = \frac{1}{\pi} \sum_{i=1}^{N_{BPM}} y_i \cdot \Delta\theta_i \cdot \sin(k \cdot \theta_i)$$

If there is a BPM next to each quadrupole and functions properly

$$45.82 \text{ GeV } (a_\gamma = 103.983)$$

$$\delta n_0 : 2.28 \text{ mrad} \Rightarrow 2.03 \text{ mrad} , P_{DK} : 10.68\% \Rightarrow 13.72\%$$



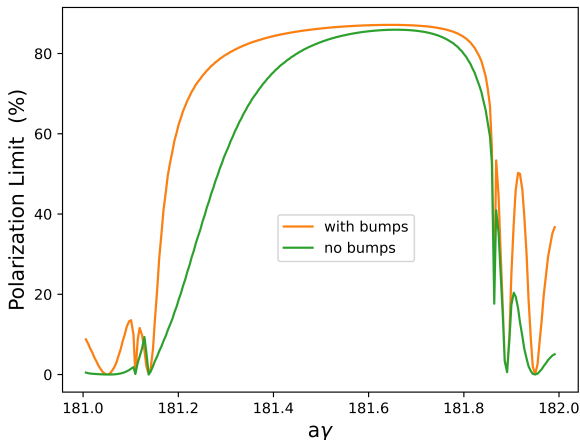
At 45.82 GeV ( $a_\gamma = 103.983$ )

Method	$(\delta n_0)_{rms}$ (mrad)	Polarization (%)
no correction	2.28	10.68
HERA formalism	0.90	90.96
Rossmannith-Schmidt scheme	0.90	89.65
Modified R-S scheme	1.01	84.71
LEP method	2.03	13.72

Many questions remained regarding all three schemes

80.192 GeV ( $a\gamma = 181.986$ ),  $(\Delta y)_{\text{rms}} = 117.44 \mu\text{m}$

$\delta n_0 : 5.87 \text{ mrad} \Rightarrow 2.88 \text{ mrad}$  ,  $P_{DK} : 4.69\% \Rightarrow 34.65\%$



## 1. HERA formalism

- Pro: systematic and rigorous mathematical derivation
- Con: empirically setting the bumps will be inevitable

## 2. Rossmanith-Schmidt scheme

- Pro: based on the acquisition of a more measurable quantity
- Con: BPMs at both ends of each dipole/quadrupole → extra cost
- Con: restricted by BPM misalignments and calibration errors

## 3. LEP method

- Pro: based on the real observables
- Con: restricted by BPM misalignments and calibration errors

- Model the lattice using multiple **larger** error seeds, and estimate the maximum acceptable orbits that guarantee a sufficient polarization
- Complete the harmonic spin matching schemes
  - solve the remaining questions regarding the three schemes
  - find an effective method that relies on the analysis of real observables
  - test its effectiveness under different lattice conditions
- Possible polarization at Higgs and  $t\bar{t}$  energies in FCC-ee  
( 1% at H in CEPC simulation)

# Thank you!

## Remained problems of HSM

- What's the harmonics that should be corrected in the simplified HERA formalism
- Whether there is a way to extract  $\Delta y'$  information from vertical BPM readings in quadrupoles
- How to make LEP method work
- If it's possible to correct vertical resonance (not HSM)
- What will happen if the errors are much larger
- What will happen near other integers besides 103 and 104

	$\tau_{ST}$	$\tau_{BKS}$	$\tau_{dep}$	
FCC-ee $(\Delta y)_{\text{rms}} = 72 \mu\text{m}$	11 779 min	11 773 min	$4.26 \times 10^6 \text{ min}$	90 min for 10% with wigglers (CDR)
HERA (26.7 GeV)	$\sim 43 \text{ min}$	$\sim 40 \text{ min}$	$\sim 10 \text{ min}$	$\tau_{dk} \sim 8 \text{ min}$
LEP	$\sim 310 \text{ min}$ 46 GeV	—	$\sim 24 \text{ min}$ 46.5 GeV	30 min for 10% no wigglers

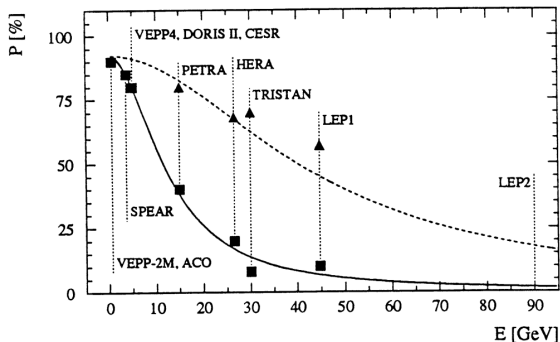
in LEP  $(\Delta y)_{\text{rms}} = 530 \mu\text{m}$ ,  $(\eta_y)_{\text{rms}} = 13 \text{ cm}$

$$P(t) = P_{dk} [1 - e^{-t/\tau_{dk}}] + P_0 e^{-t/\tau_{dk}} \simeq P_0 e^{-t/\tau_{dep}}$$



## Systematic errors of the average beam energy determination

- Energy dependent momentum compaction
- Vertical orbit distortions (radial fields)
- Longitudinal fields



Maximum measured polarization in different storage rings with HSM (triangles) and without HSM (squares)

## Possible problems with LEP method

- radial fields not only exist in quads
- the radial field seen by the particle is not fully proportional to the y position
- even if it's proportional, each quad has different strength ( $k_y$ )
- how much spin rotates is an integration of radial field within the element ( $k_y L$ )

- Search for all possible locations in a perfectly aligned lattice
- Add one bump in one possible location
- Adjust 2nd and 3rd kicks to make  $y$  and  $y_p$  be all 0 outside the bump
- Analyze its contribution to the target harmonics
- Change the bump position and redo the analysis
- Select four bumps that build a matrix with the largest determinant
- Add four bumps into a misaligned machine with estimated bump amplitudes, and match the orbits to make bumps closed