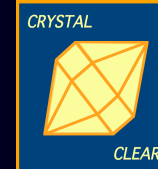


Scintillators - Fundamentals

Paul Lecoq
CERN, Genève



Outlook



- Lecture 1: Scintillator fundamentals

- Organic scintillators
- Inorganic scintillators
 - Scintillation mechanisms
 - Limits to the light yield and decay time
 - Energy resolution and non-proportionality

- Lecture 2: Scintillator applications

- Crystal growth techniques
- High energy physics and dark matter searches
- Medical applications
- Space borne missions
- Geophysical exploration
- Homeland security

Scintillators according to various schemes

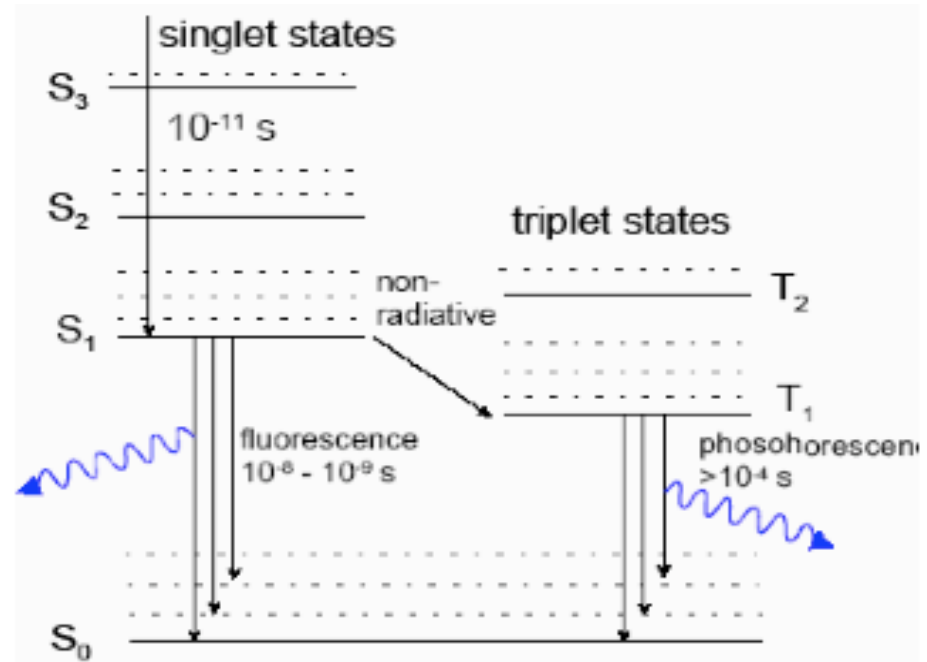
Transform dE/dx of an ionizing particle into light that can be measured by a photodetector

- **Physical state**
 - Solid
 - Liquid
 - Gas
- **Structure**
 - Single crystal
 - Ceramic
 - Glass
- **Composition**
 - Organic
 - Inorganic
- **Scintillation mechanism**
 - Intrinsic
 - Activated
 - Core-valence

Convert PART of the energy of the incident particle

organic scintillators low Z (C,H) \rightarrow
 - low γ -detection efficiency
 - high n-detection efficiency via (np)
 scintillation mechanism:

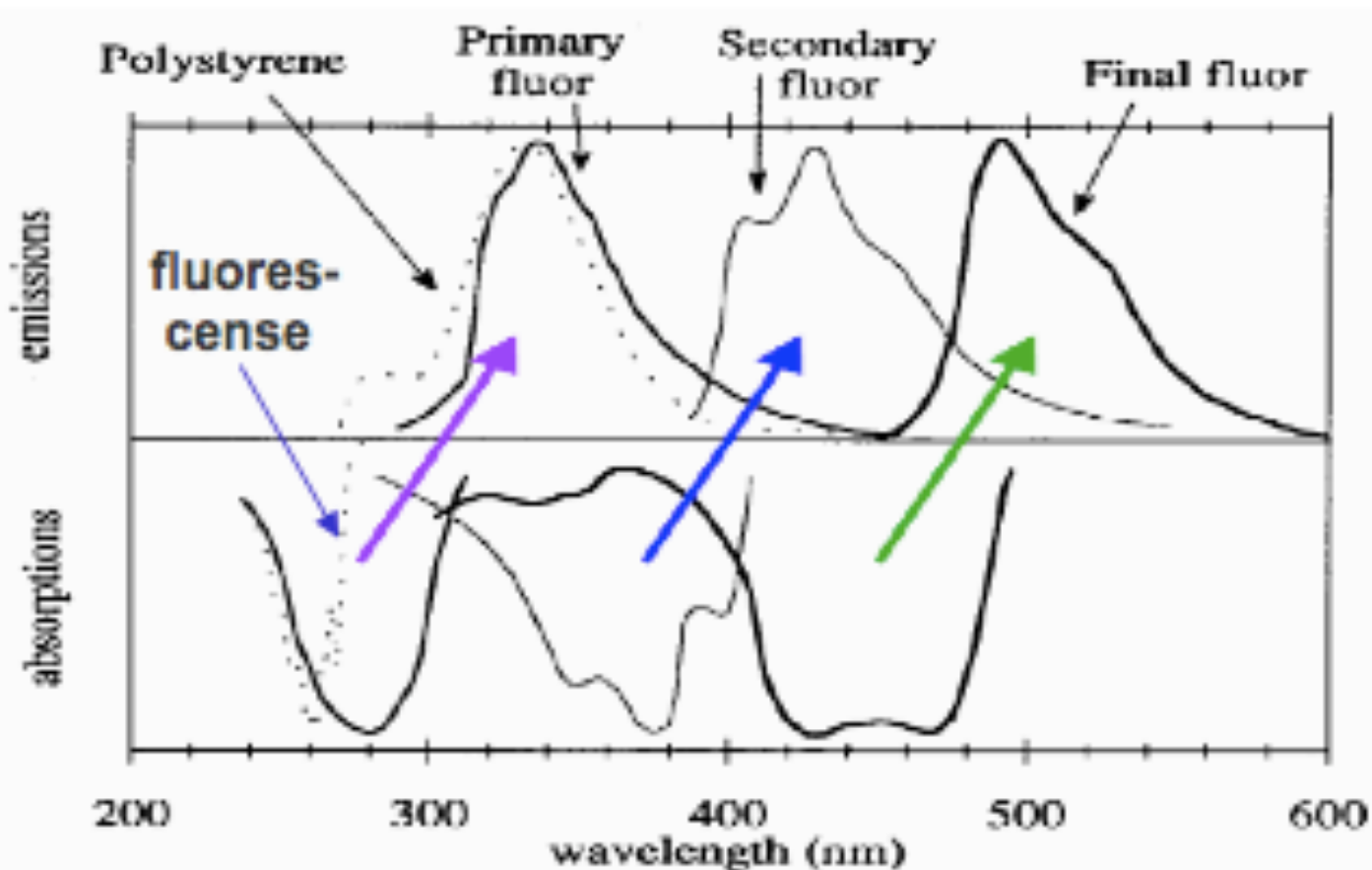
Delocalized π electron states of the Benzene molecule



- Organic crystals**
 Anthracène, Trans-Stilbène, Naphtaline
- Organic liquids**
 Solvent: Xylène, Toluène, benzène
 Solute: p-Terphénil, PBD, PPO, POPOP, 3g/l
- Plastics**
 Solvent: polyvinyletoluène, polyphénilbenzène, polystyrène
 Solute: PBD, pTerphénil, PBO, second soluté POPOP, 10g/l for wavelength shifting

Wavelength shifter

Principle of WLS:

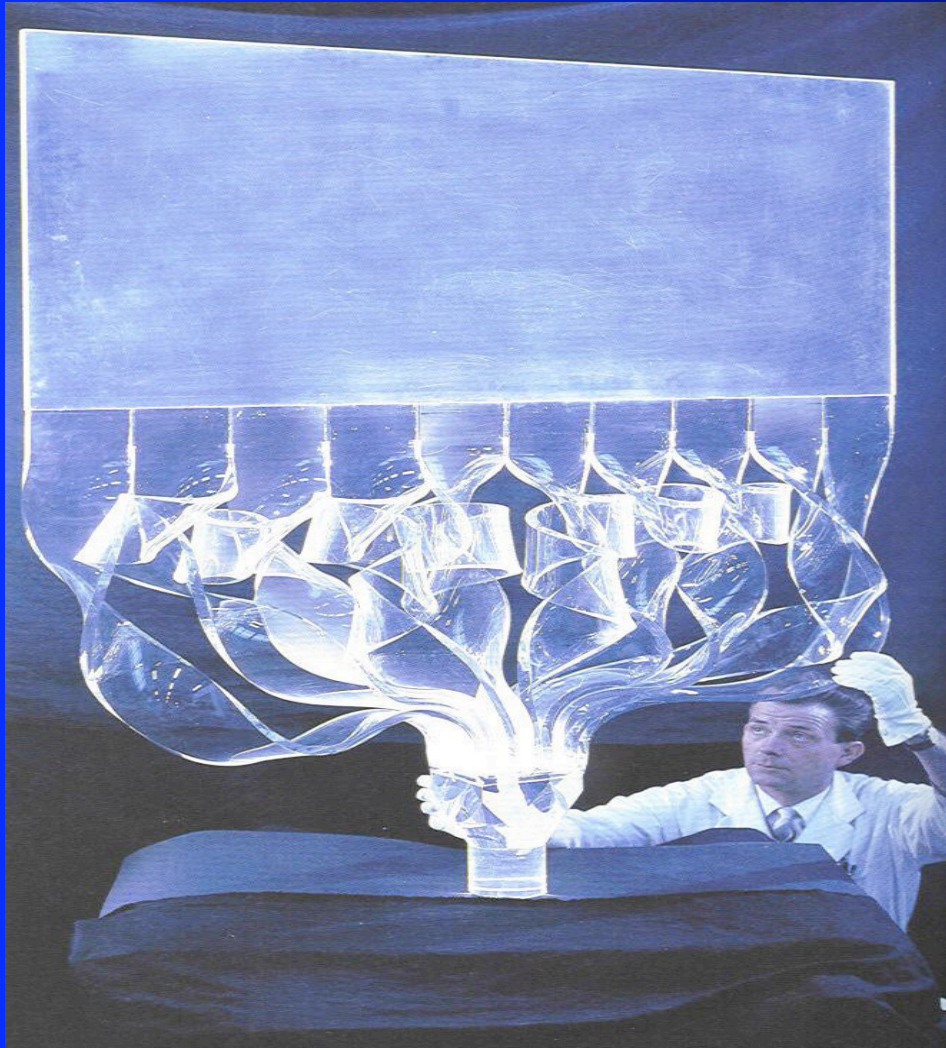


Crystalline organic scintillators

crystal	Chemical formula	density	n	yield	emission wavelength nm
anthracene	$C_{14}H_{10}$	1,25	1,62	100	447
Trans-stilbene	$C_{14}H_{12}$	1,16	1,62	50	410
naphtalene	$C_{10}H_8$	1,162	1,62	30	340

- organic scintillators are usually very fast (a few ns)
- used for fast detection, time tagging, time of flight
- Anthracene has a very good yield: 1 optical photon per 60eV deposited energy

Plastic organic scintillator: plates



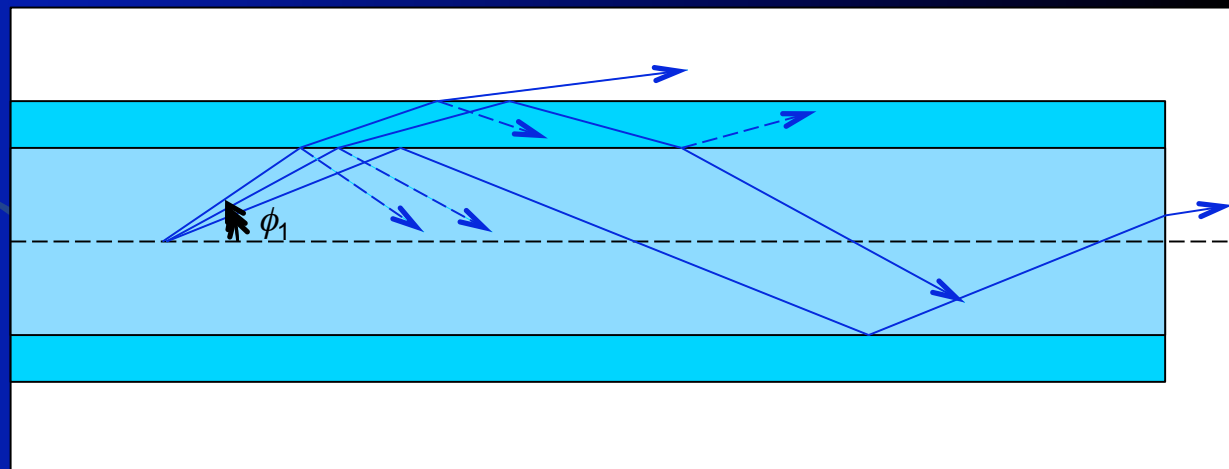
- Easily machined
- Large sizes available
- Good light transport with wavelength shifting using primary and secondary fluors
- Very fast~ns,
- Cheap
- Not very radiation hard

1 optical photon per 100 eV deposited energy

Air: $n_0 = 1.0003$

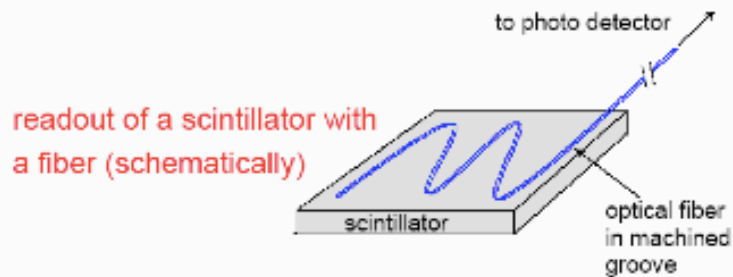
Core, polystyrene: $n_1 = 1.59$

Cladding, acrylic: $n_2 = 1.49$



- Propagation in the core: $\phi_1 < 20.2^\circ$, $f_1 = 1 - n_2/n_1 = 6.2\%$
- Propagation in the cladding: $20.2^\circ < \phi_1 < 51^\circ$, $f_2 = n_2/n_1 - n_0/n_1 = 31\%$
- Lost in air: $\phi_1 > 51^\circ$, $f_0 = n_0/n_1 = 63\%$

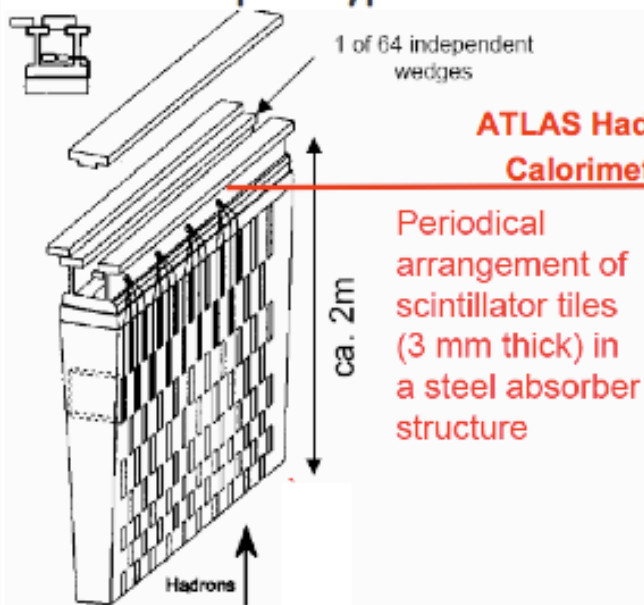
Fibres can be embedded in scintillator:



readout of a scintillator with a fiber (schematically)

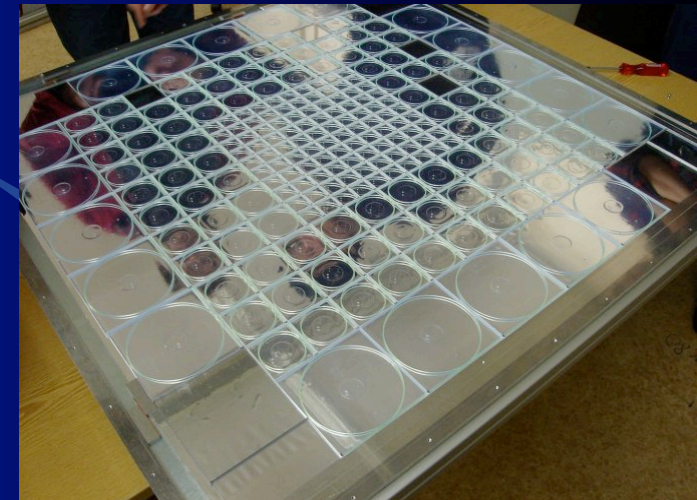
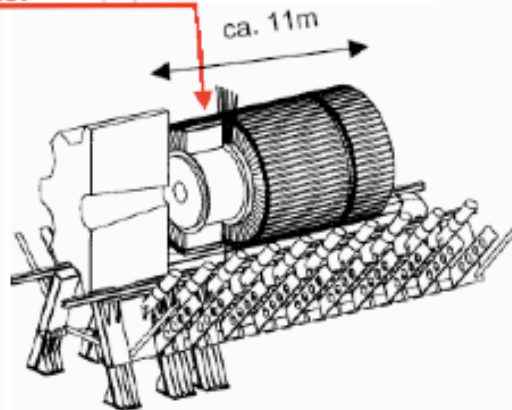
Read out 216 tiles/module
~8000 channels

(with miniaturised Si-PM no transport needed → prototype calorimeter ILC)



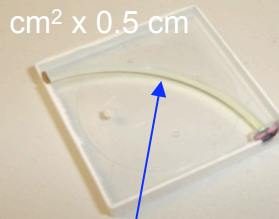
ATLAS Hadron Calorimeter

Periodical arrangement of scintillator tiles (3 mm thick) in a steel absorber structure



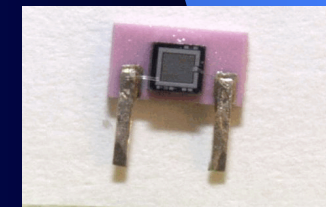
Calorimeter cell

3x3 cm² x 0.5 cm



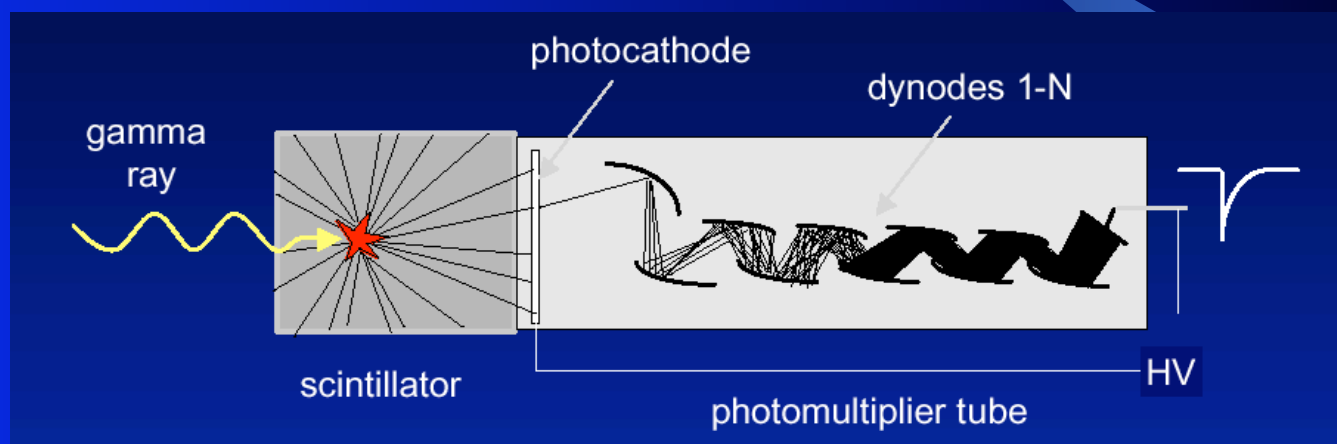
Wavelength shifter fiber

Single tile readout with SiPM



Scintillating crystals for homogeneous calorimeters

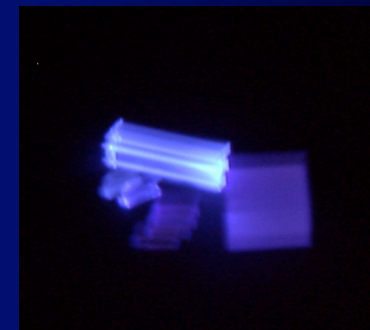
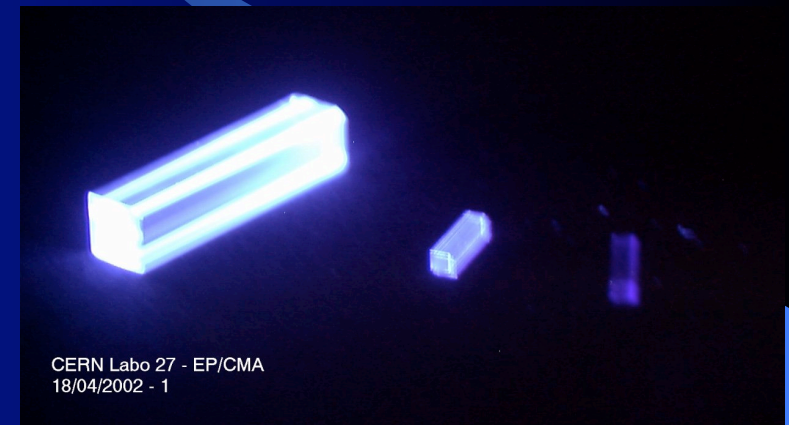
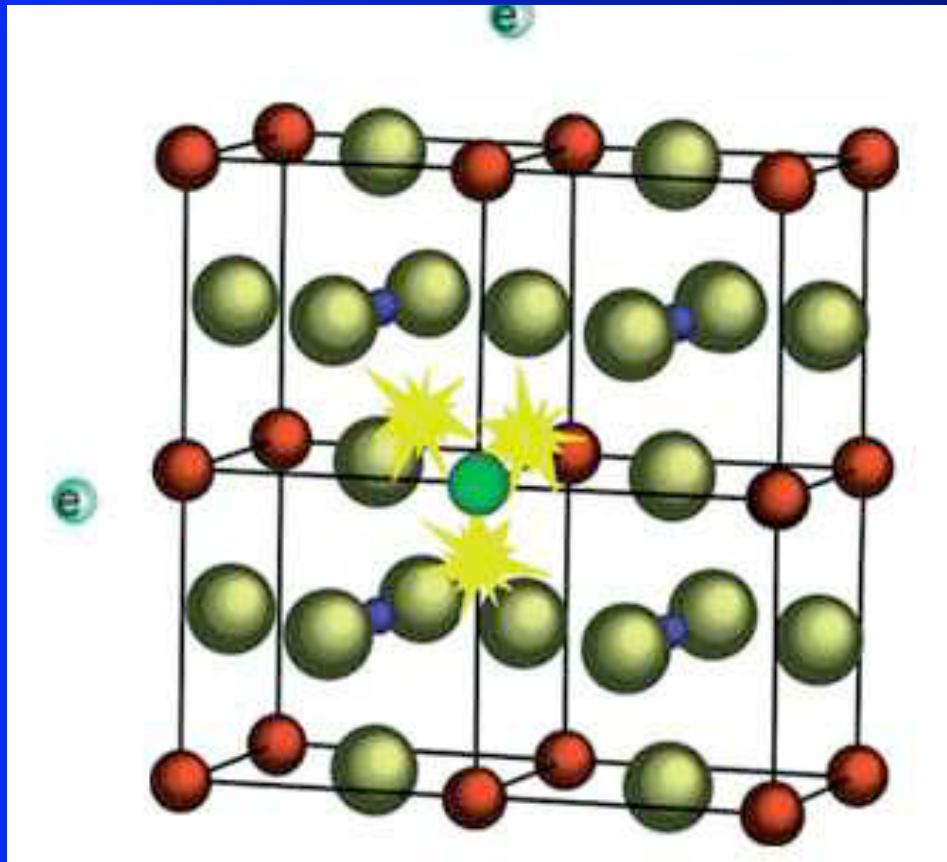
- To convert ALL the energy of the incident particle in to light
- Necessity to use dense materials



- Above certain minimum level most scintillators are linear with respect to the energy deposited
- Light output is directly proportional to energy deposited

Why a crystal?

- Heavy material are rich in electrons, which interact strongly with light
- Only ordered system can confine electrons in well separated energy bands, so that the material is transparent to its scintillation light



Examples of crystals

Quartz



Natural

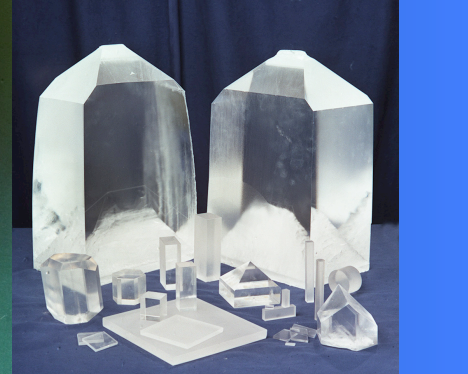
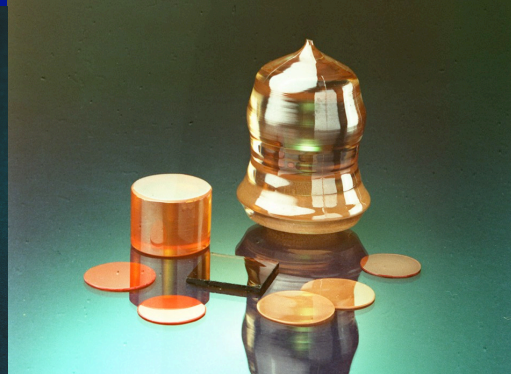
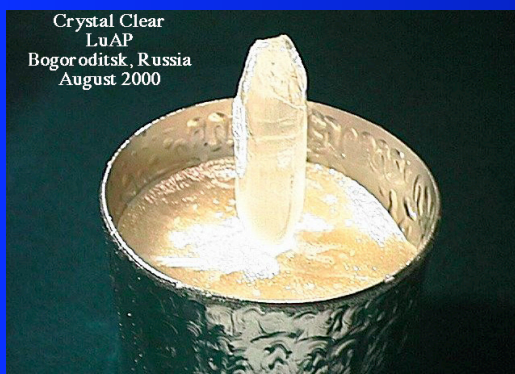
Lead Tungstate



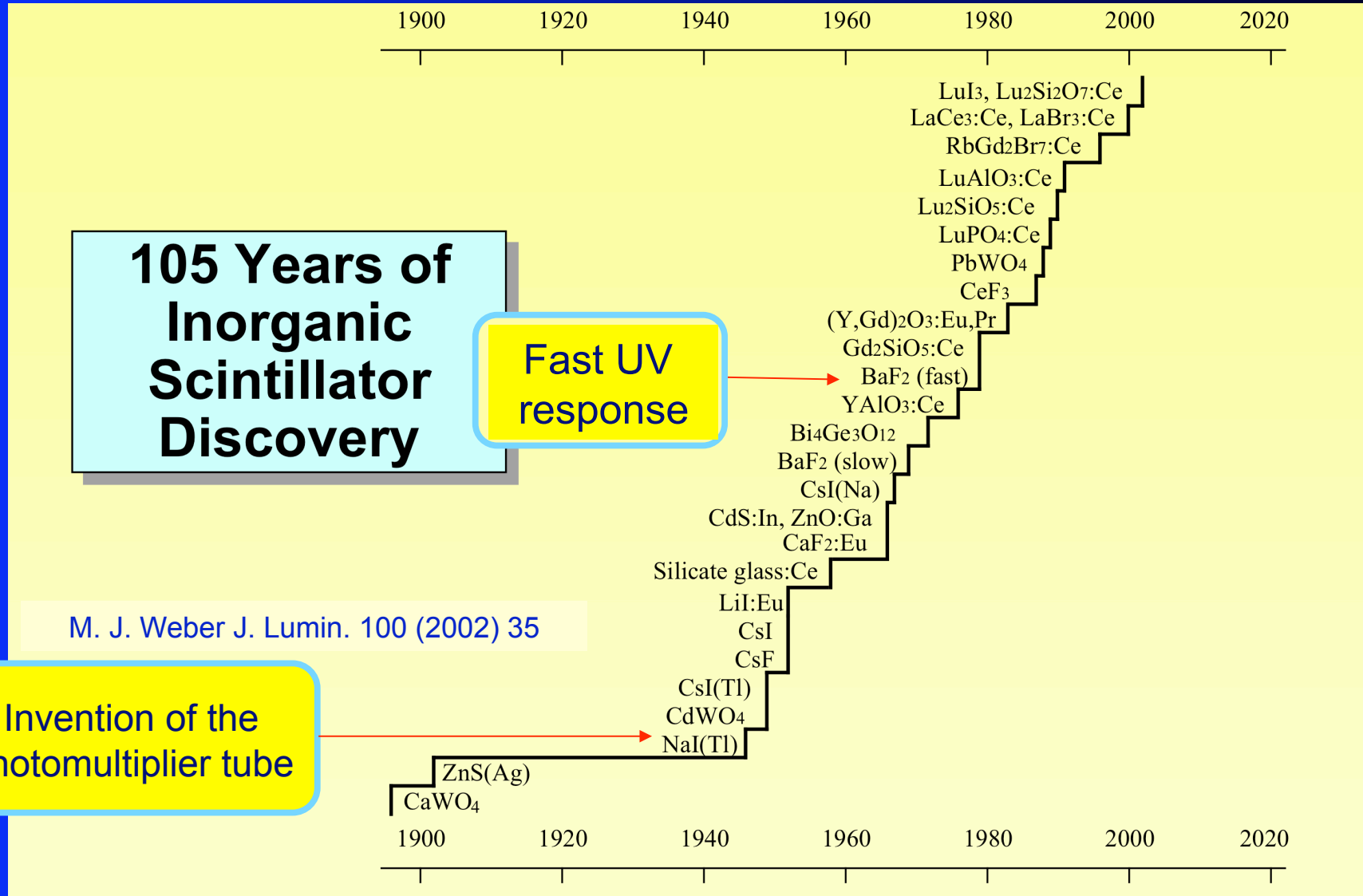
Synthetic



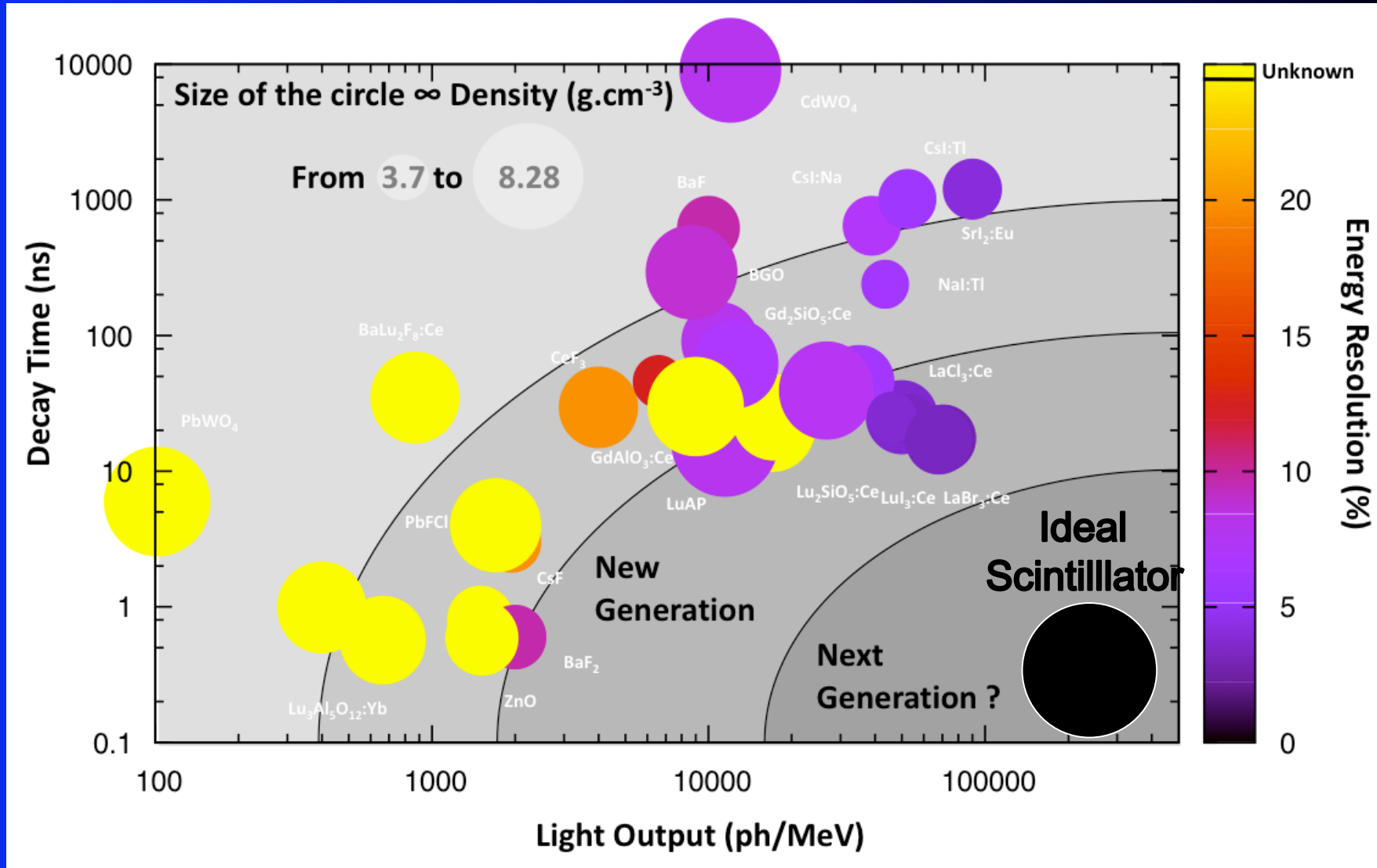
Crystal Clear
LuAP
Bogoroditsk, Russia
August 2000



History of scintillator discovery

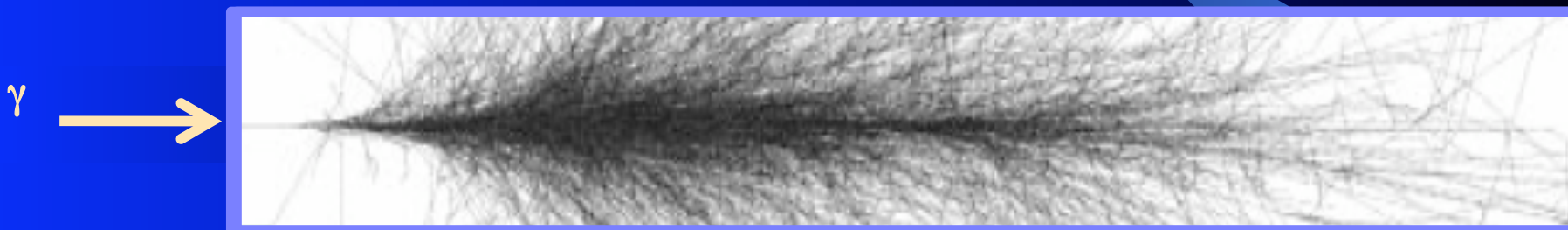


Classification of scintillators



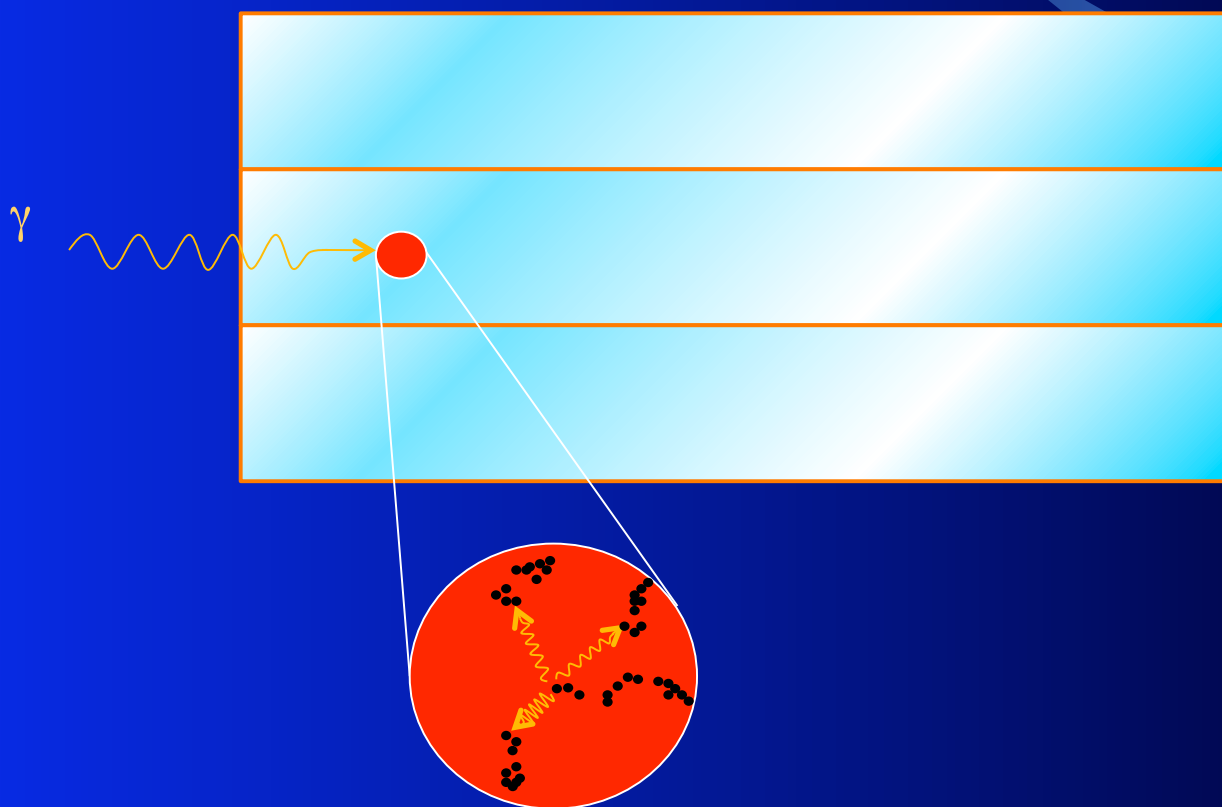
A zoom on the conversion process (HEP)

- The energy conversion from incoming X or γ Rays is a complex process resulting from a cascade of events.

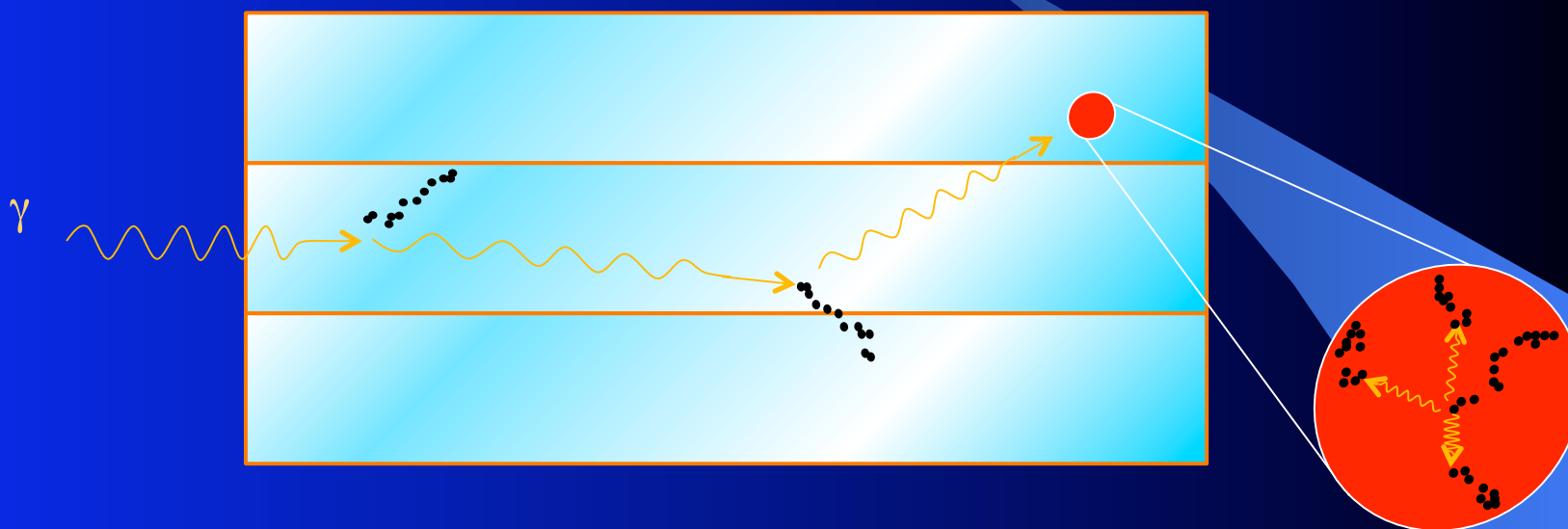


- Hadronic events are even more complex
 - Details of the full cascade for HEP with contributions from different conversion mechanisms: scintillation and Cerenkov, would lead to particle identification within the shower

A zoom on the conversion process (low energy)



A zoom on the conversion process (low energy)



- For charged particles: high ρ materials to increase dE/dx
- For X and γ -rays (but also high energy electrons, which radiate γ -rays by bremsstrahlung) 3 mechanisms:

– Photoelectric:

$$\sigma_{ph} \propto \frac{Z^5}{E_\gamma^{7/2}}$$

– Compton:

$$\sigma_c \propto Z$$

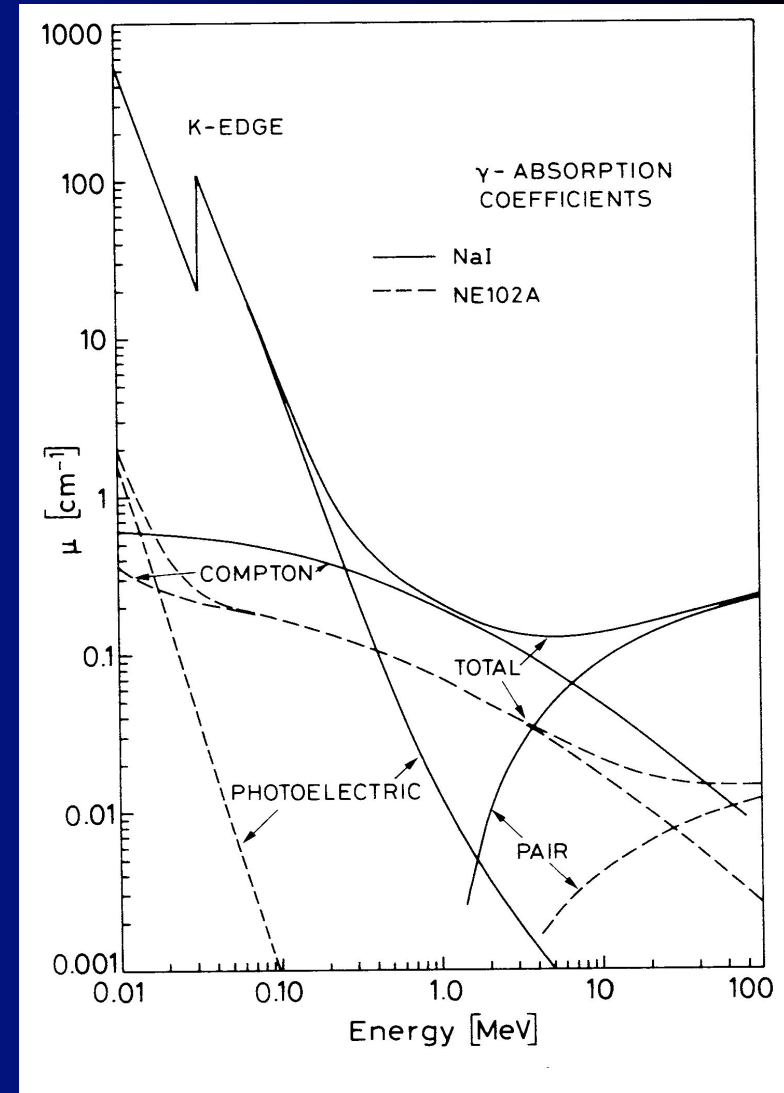
– Pair production:

$$\sigma_{pair} \propto Z^2 \ln(2E_\gamma)$$

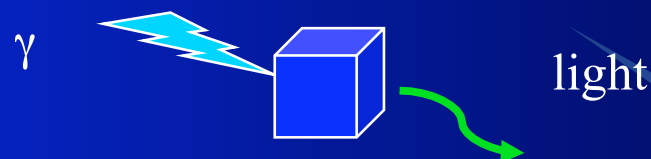
- At low energy high photoelectric cross-section is desired
- At high energy good shower containment requires
 - Small radiation length:
 - Small Moliere radius:

$$X_0 = \frac{A}{\rho} \frac{716.4 \text{ gcm}^{-2}}{Z(Z+1) \ln(287/Z)}$$

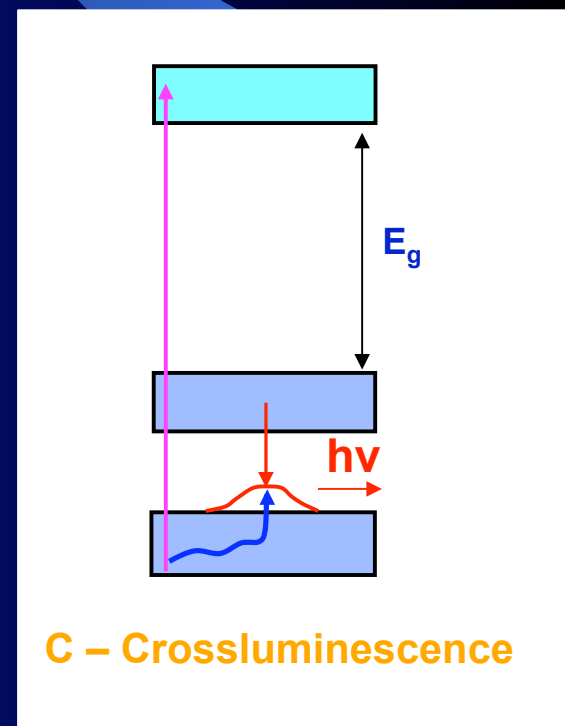
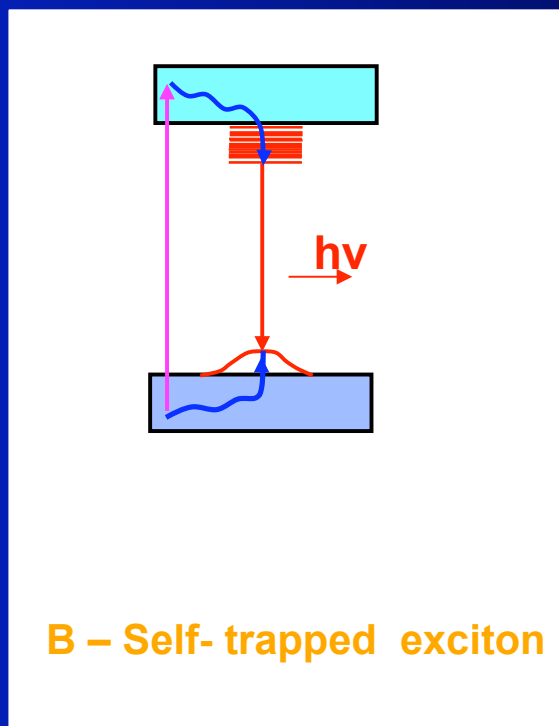
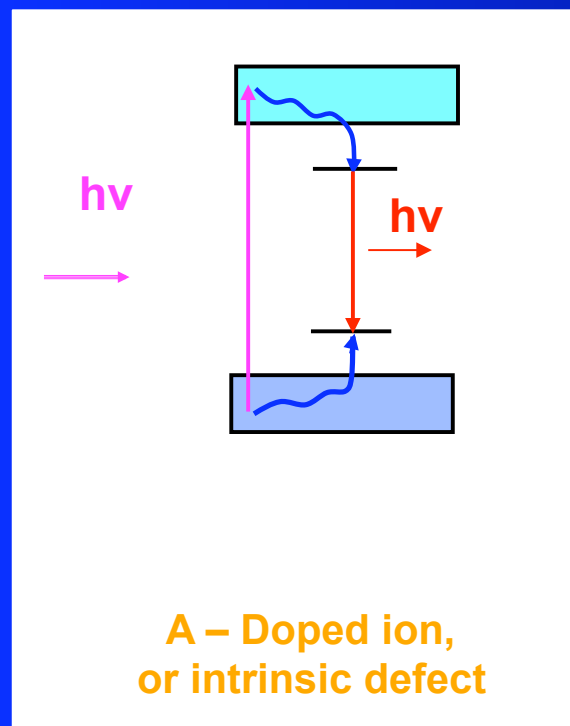
$$R_M \approx X_0 \frac{Z+1.2}{37.74} \propto 1/\rho$$



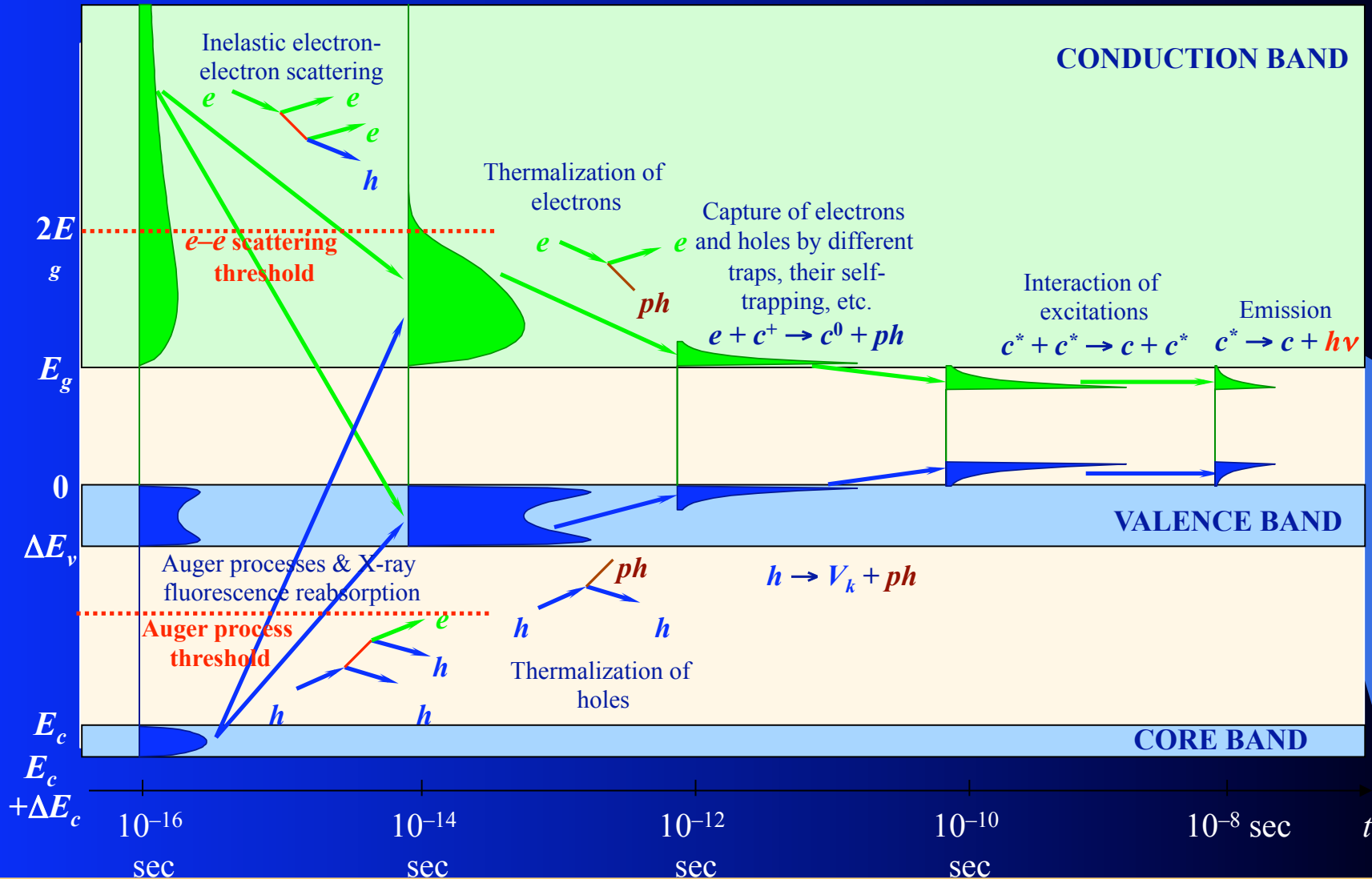
Fundamental aspects of Scintillation



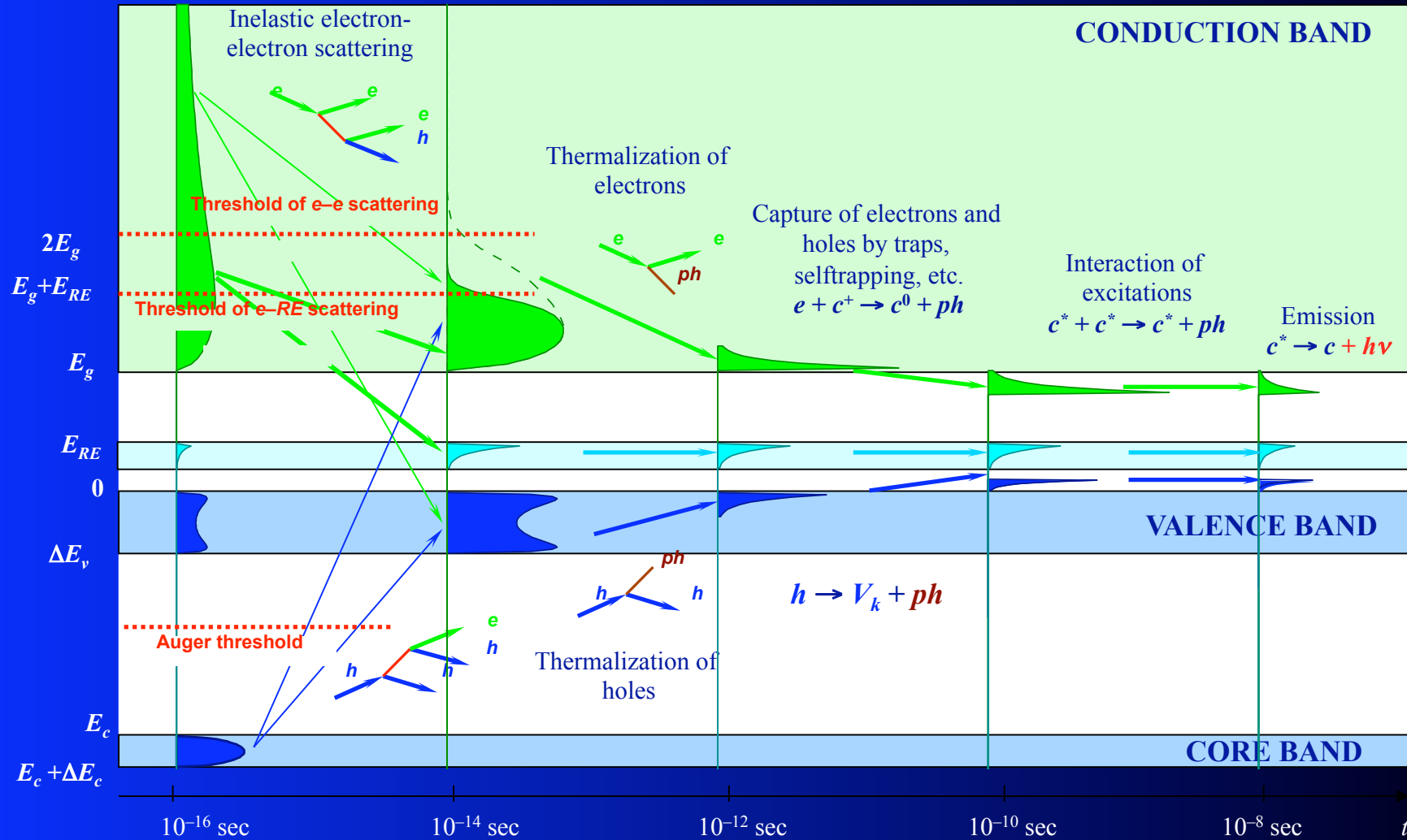
Different scintillation mechanisms



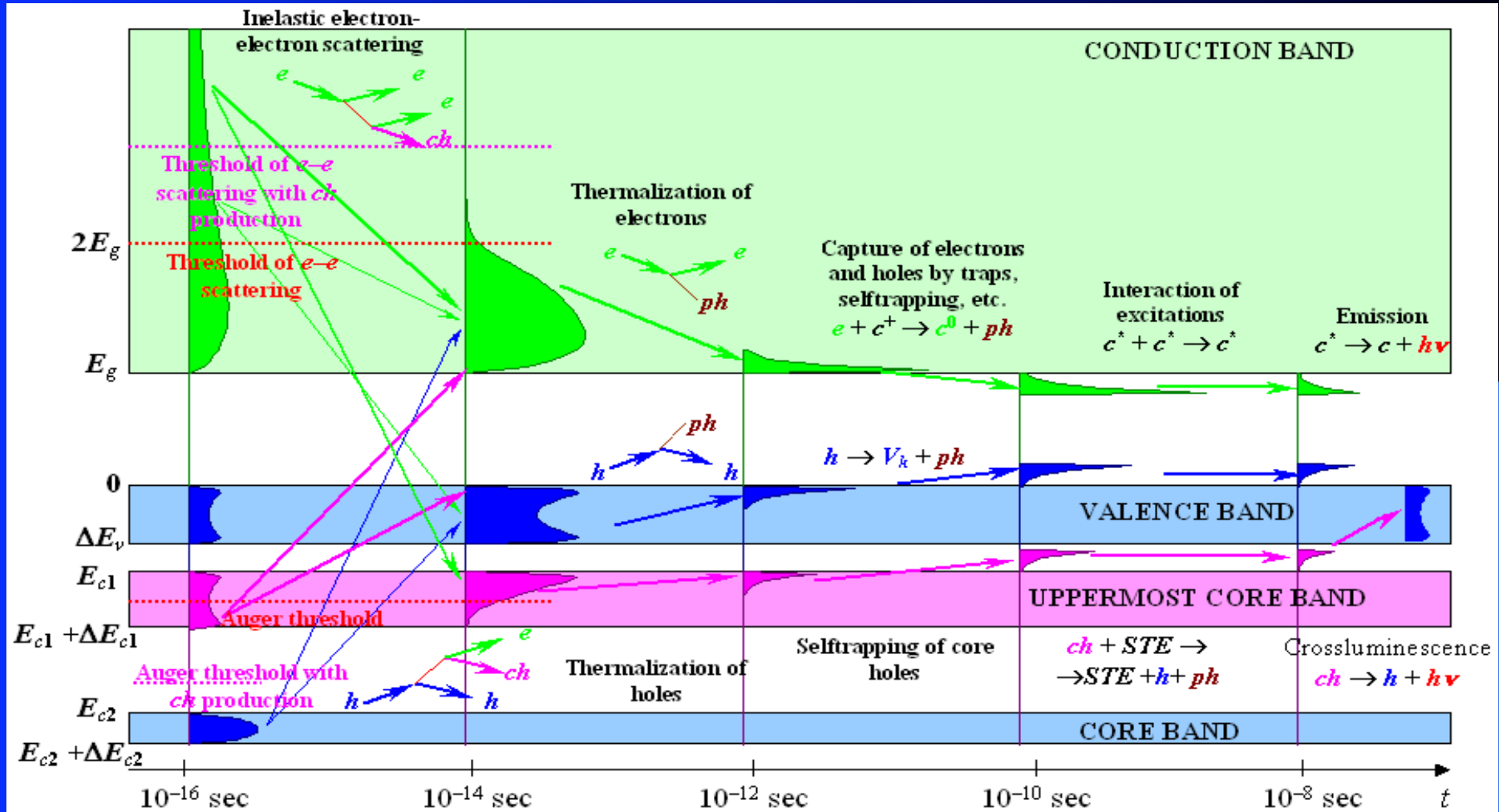
Relaxation of electronic excitations *intrinsic luminescence*



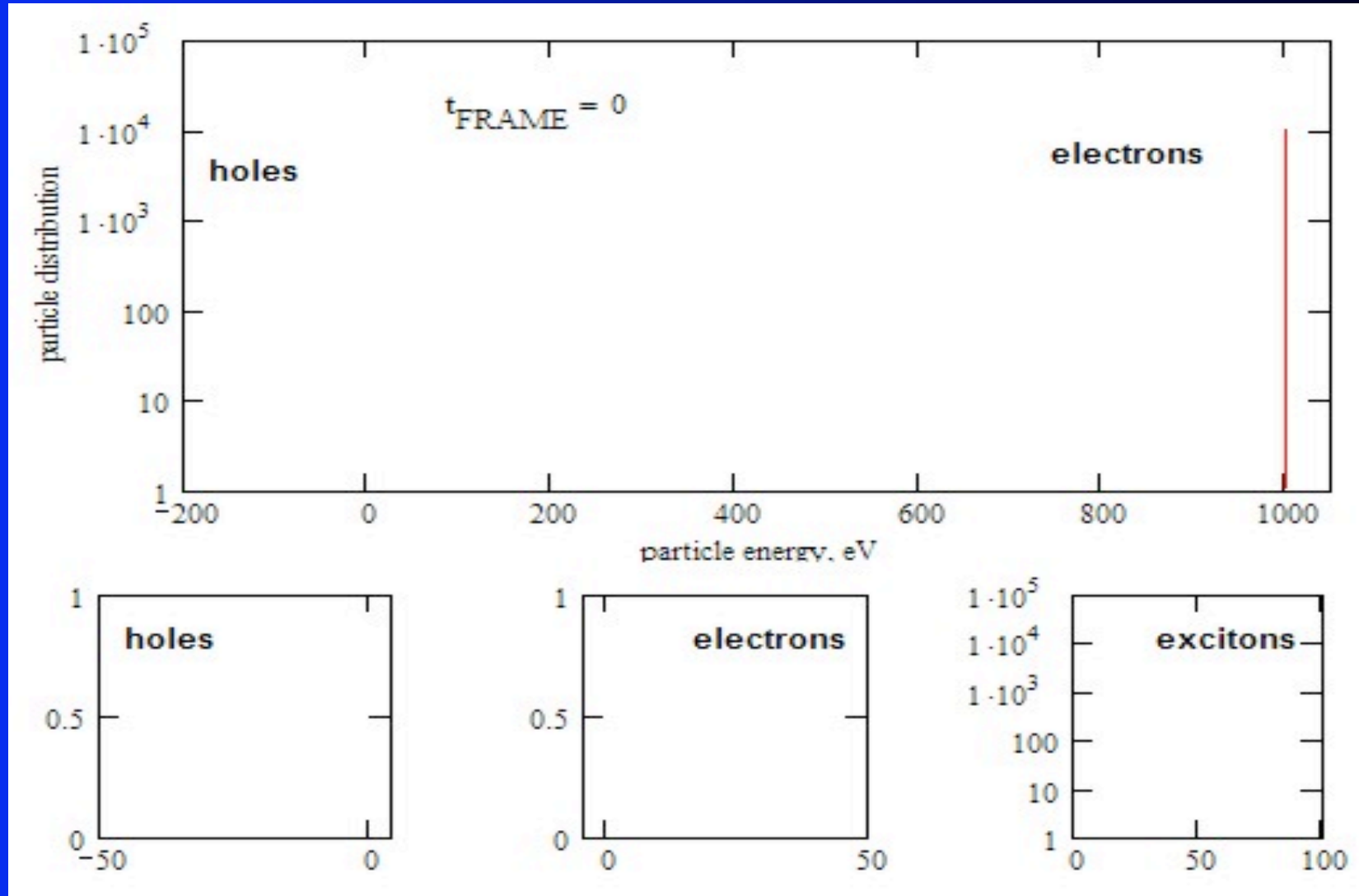
Relaxation of electronic excitations *extrinsic luminescence*



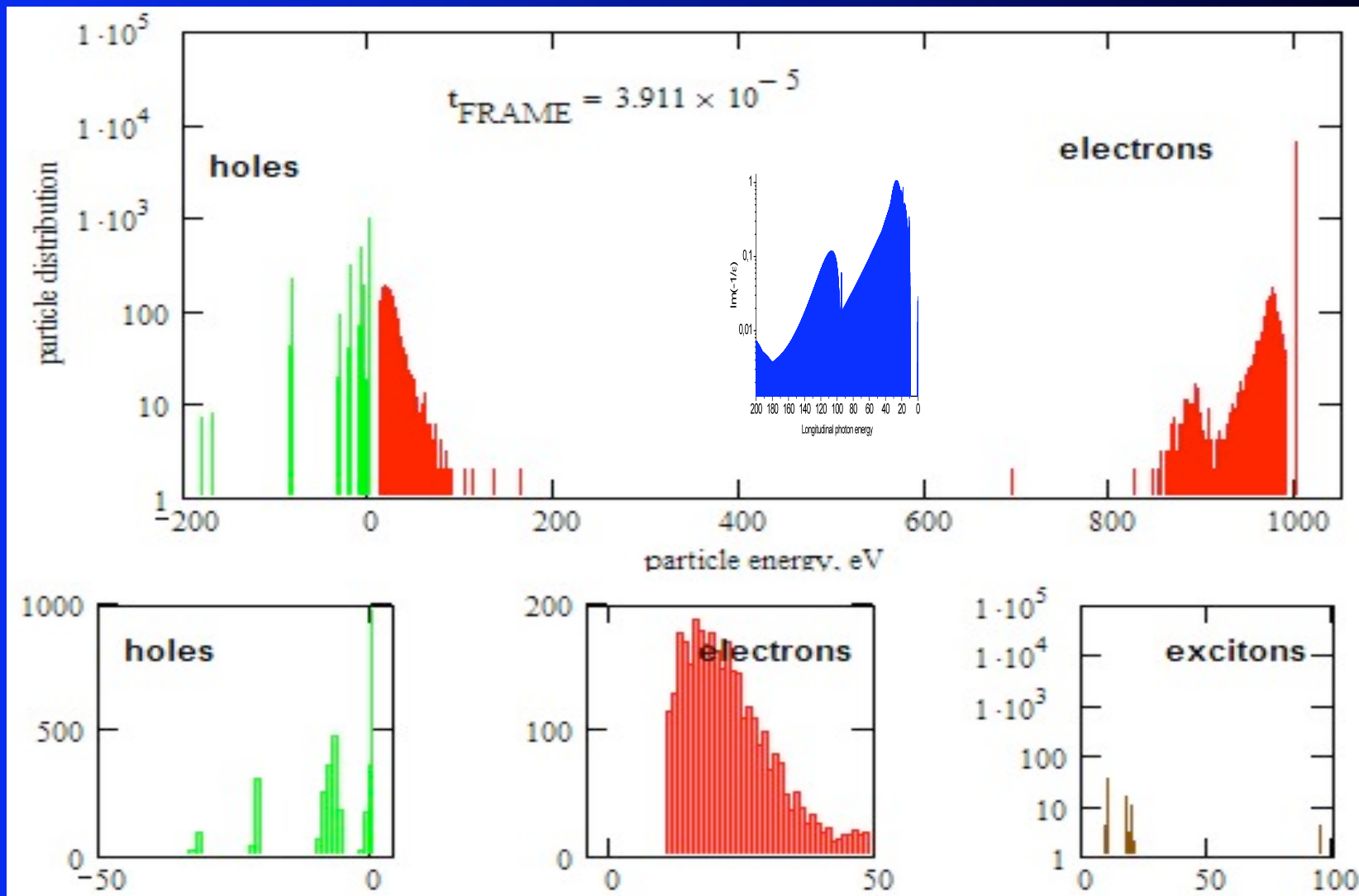
Relaxation of electronic excitations Cross-luminescence



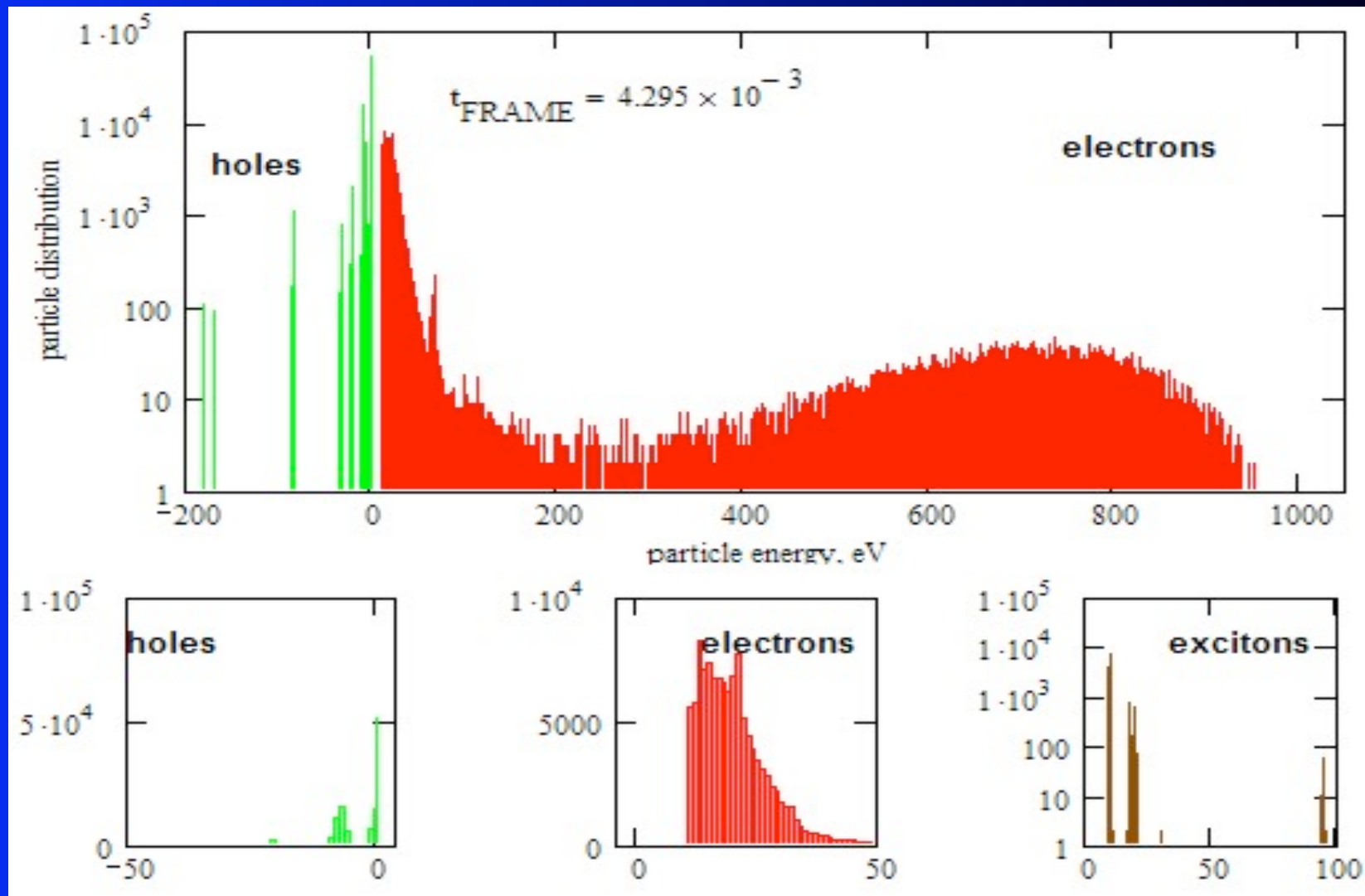
Evolution of energy distribution for 1000 eV electrons



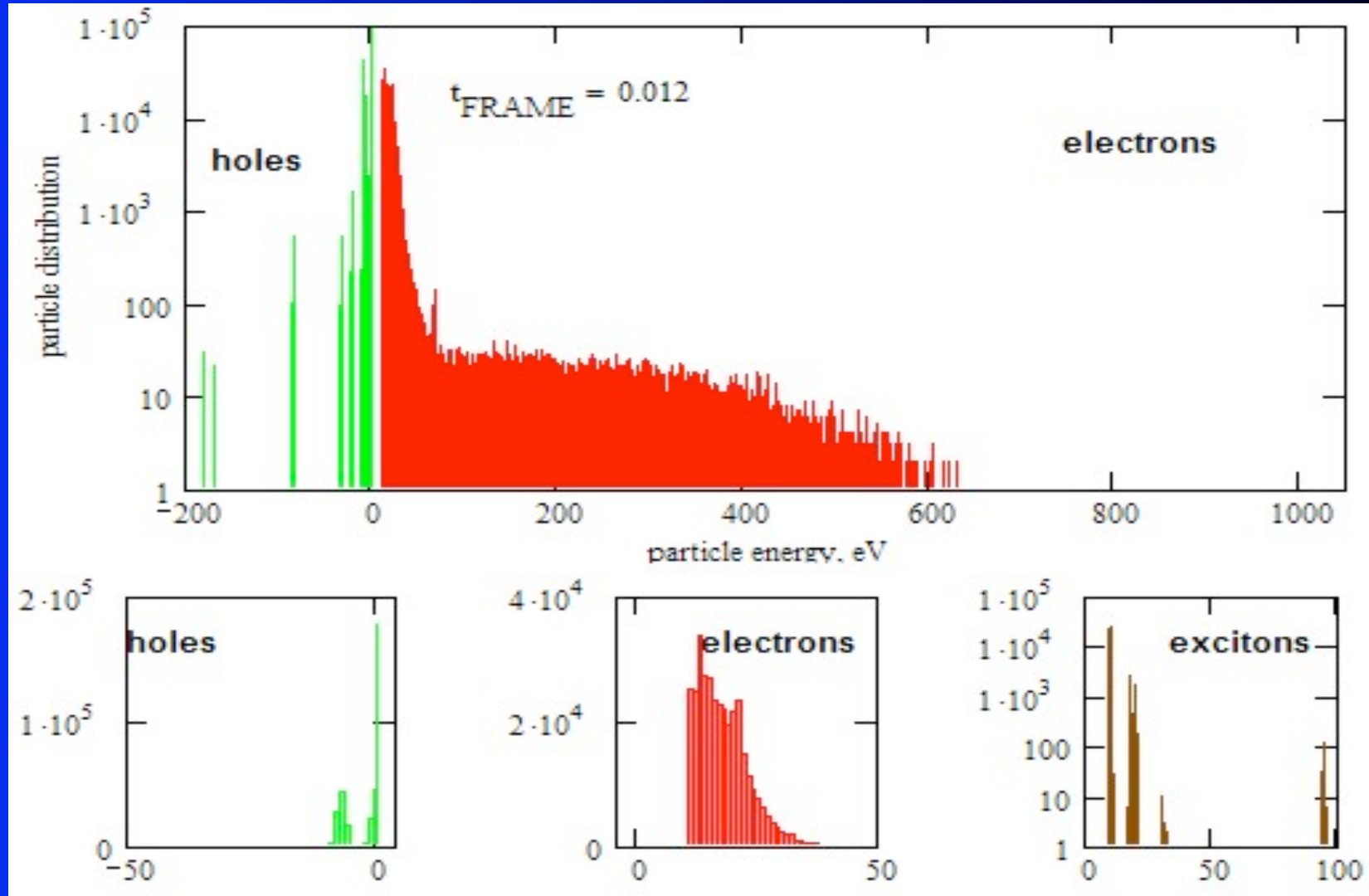
Evolution of energy distribution for 1000 eV electrons



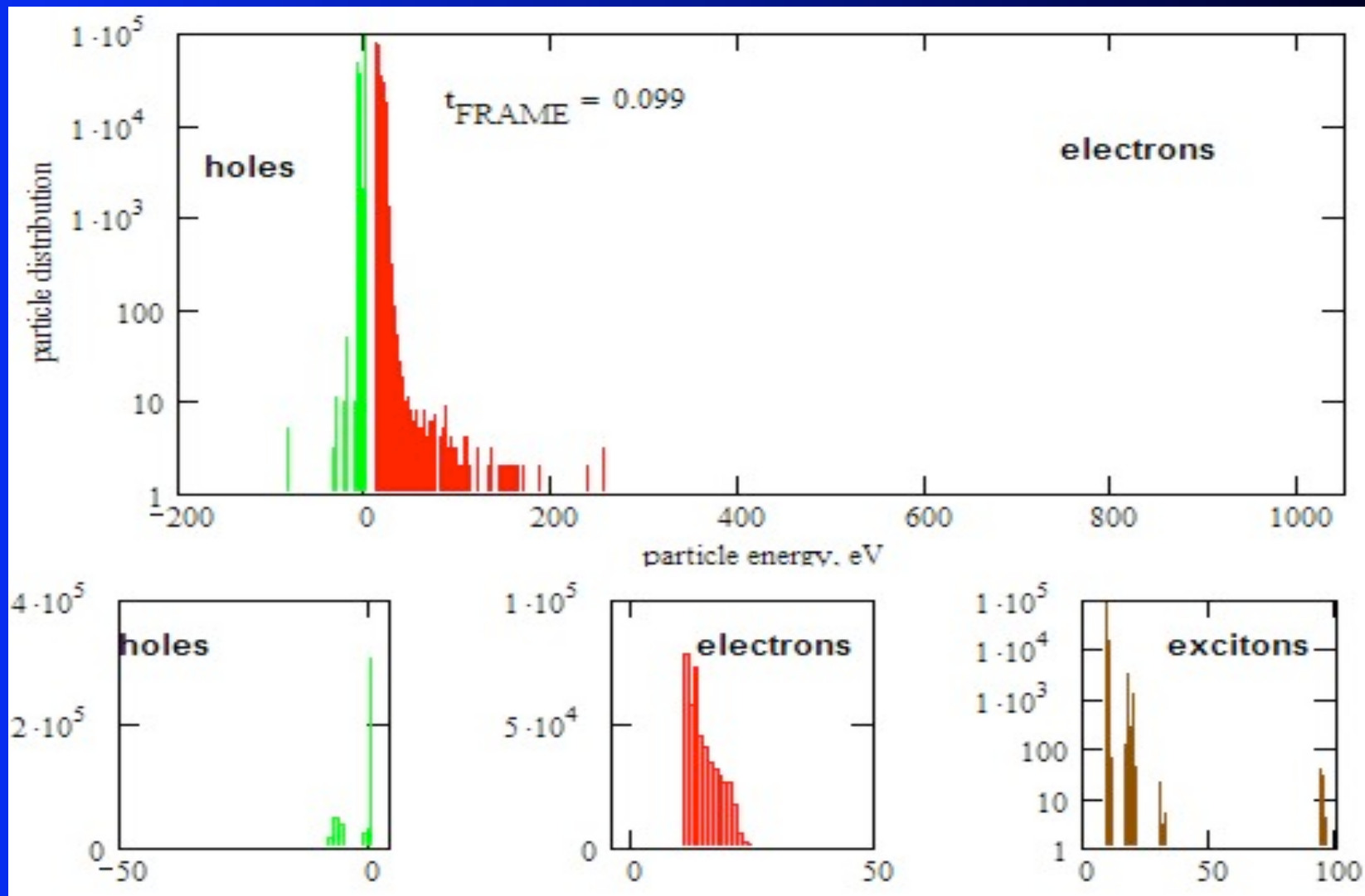
Evolution of energy distribution for 1000 eV electrons

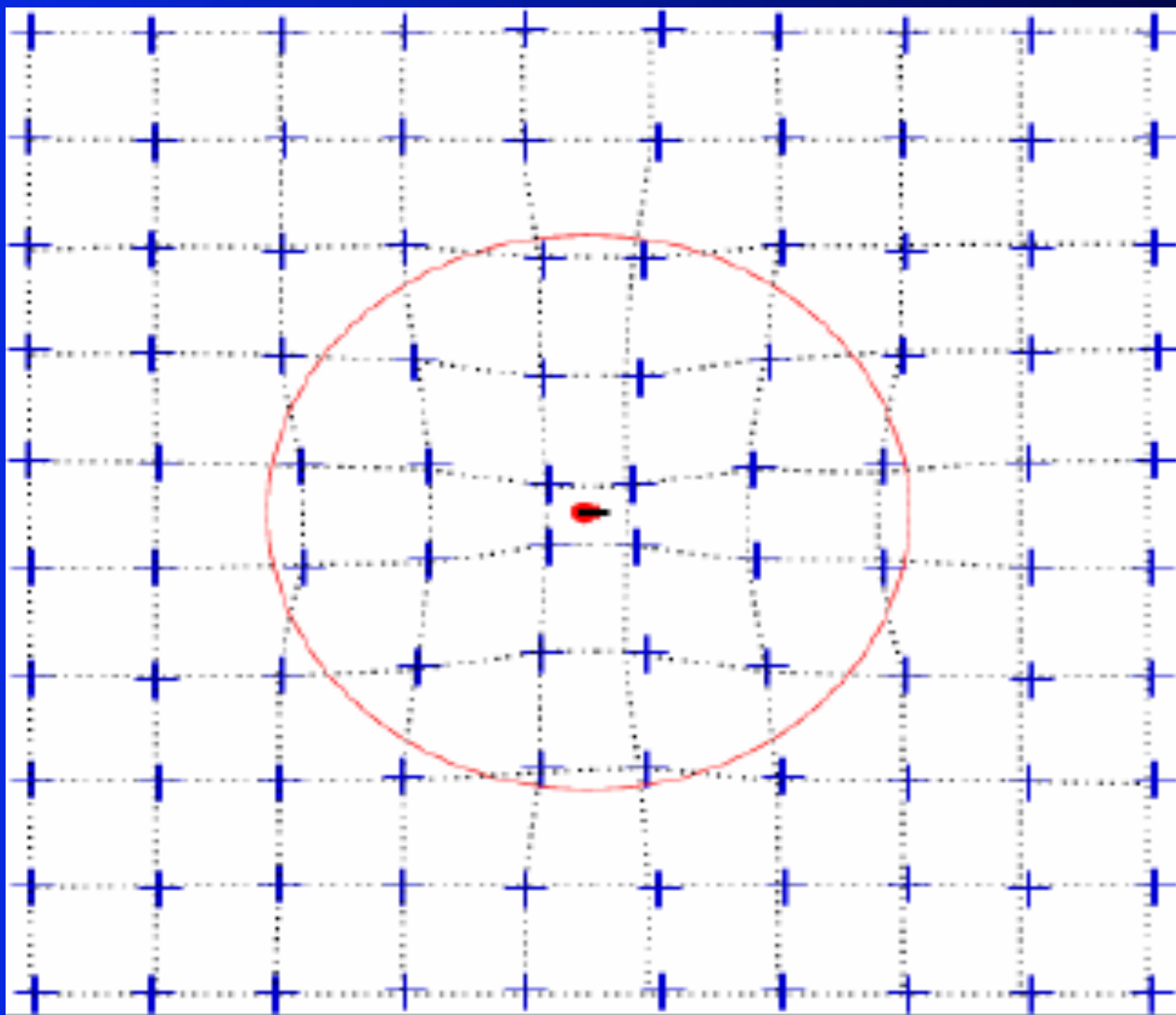


Evolution of energy distribution for 1000 eV electrons



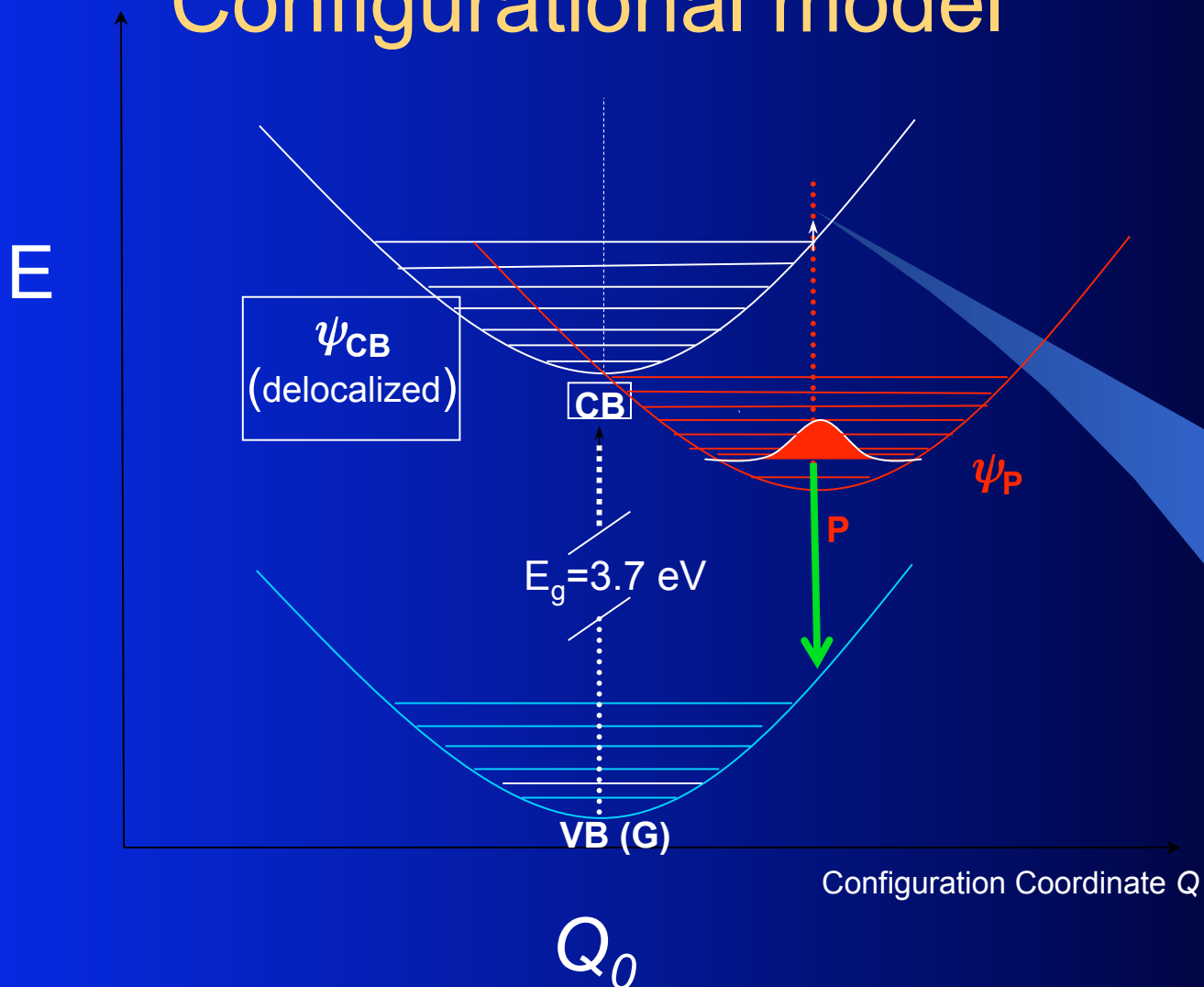
Evolution of energy distribution for 1000 eV electrons





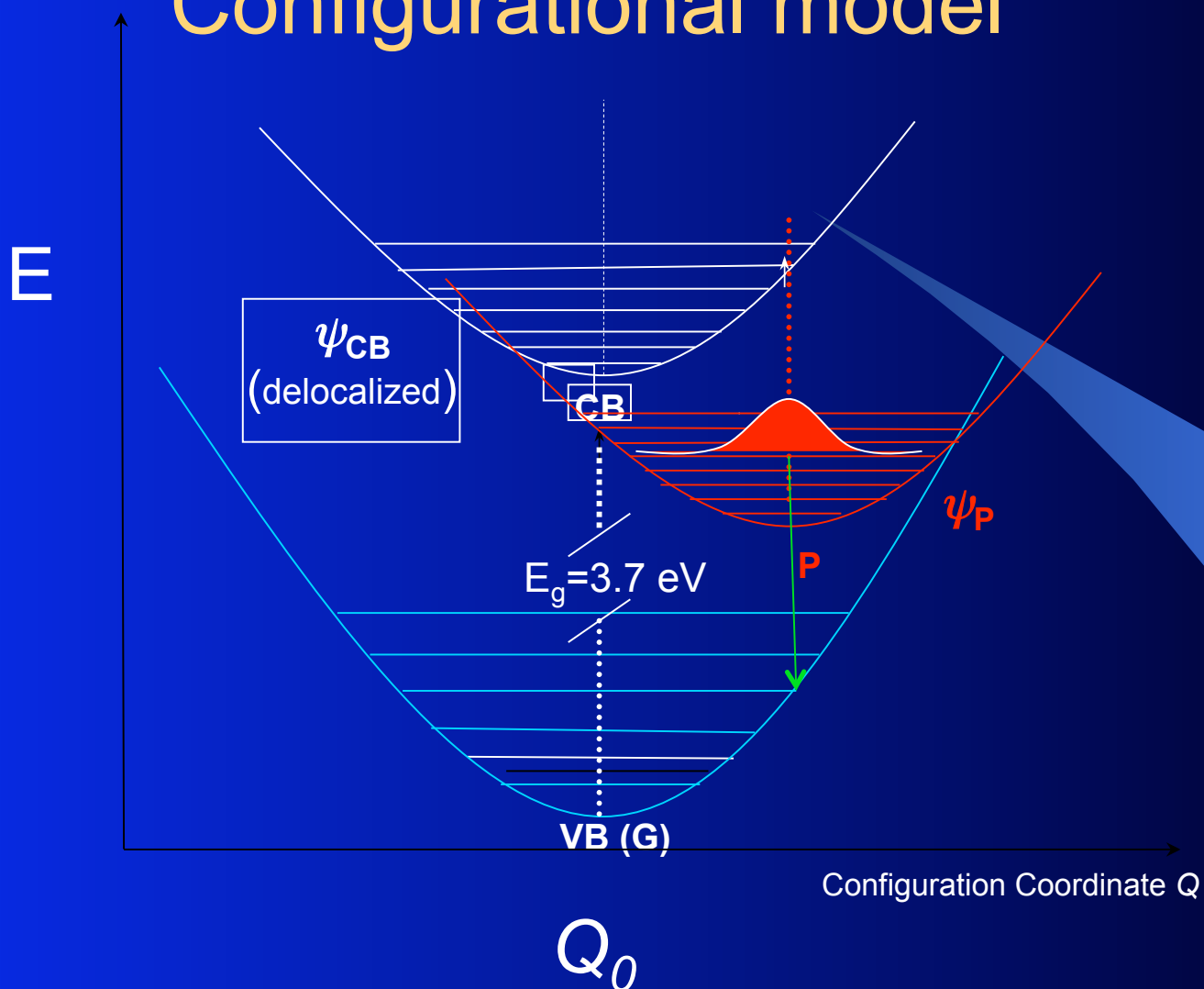
polaron - electron + distorted lattice

Configurational model



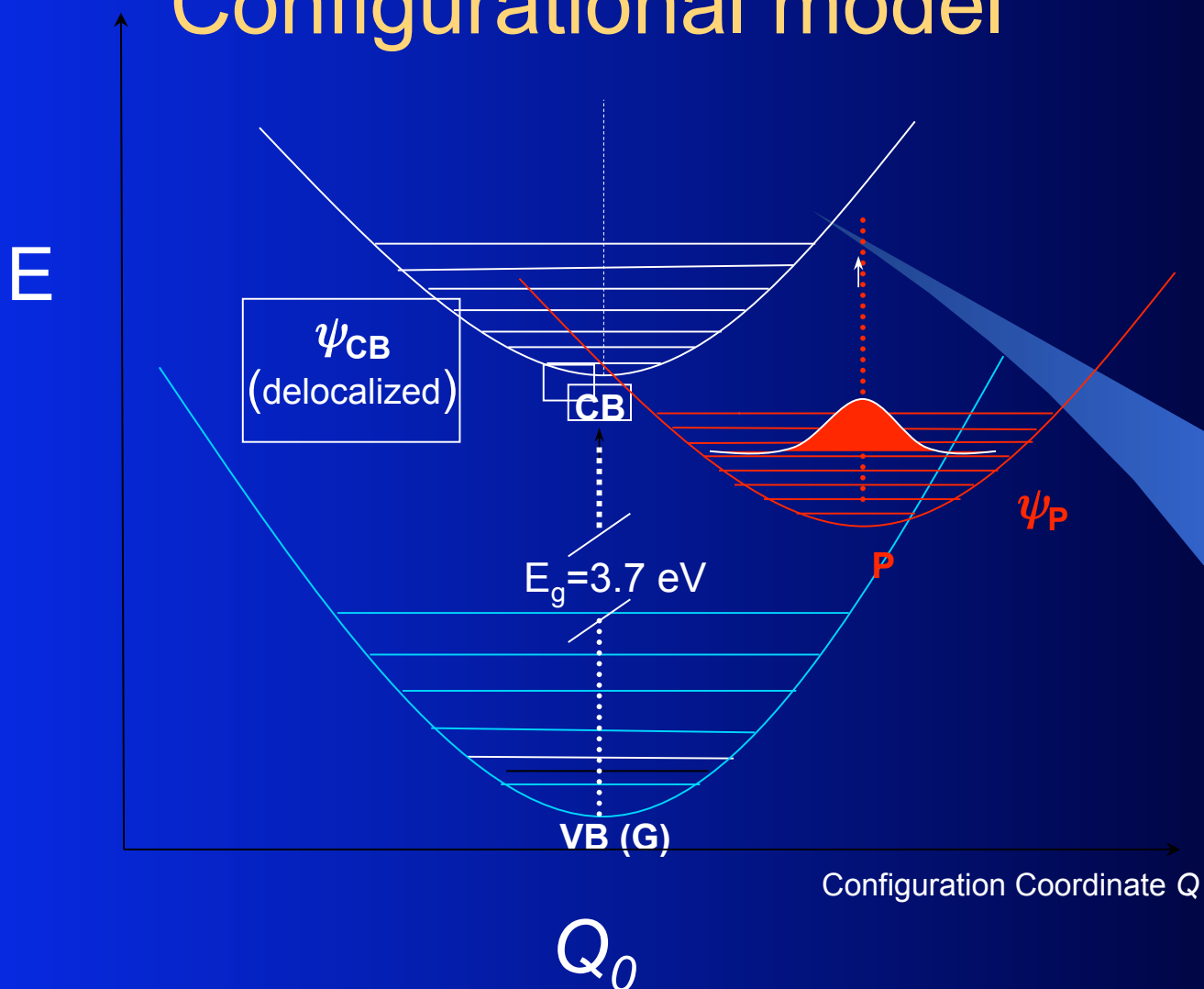
Configuration coordinate model for the local lattice with electron in **valence** and **conduction band** states and in **localized polaron** state.

Configurational model



Configuration coordinate model for the local lattice with electron in **valence** and **conduction band** states and in **localized polaron** state.

Configurational model



Configuration coordinate model for the local lattice with electron in **valence** and **conduction band** states and in **localized polaron** state.

Fundamental aspects of Scintillation

The 3 phases of the scintillation mechanism

1. Absorption : Creation of pair e-h

$$n_{e-h} = \frac{E_\gamma}{\beta E_{gap}}$$

2. Transfer to the luminescence centre

Efficiency of energy transfer :

S

3. Emission

Efficiency of emission :

q

Efficiency of scintillation

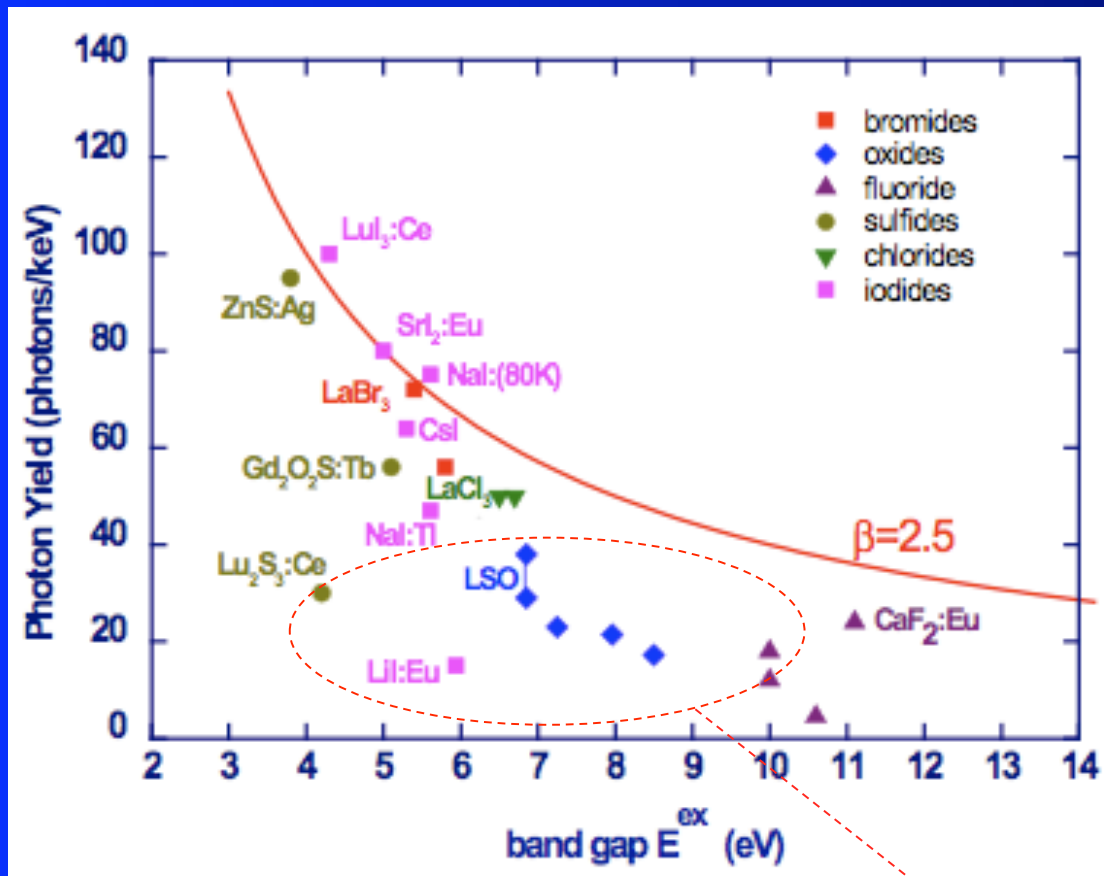
$$n_{photon} = n_{eh} Sq = \frac{E_\gamma}{\beta E_g} Sq$$

Determination of the maximum of light

$$LY_{max} = \frac{n_{photon}}{E_\gamma} = \frac{1}{\beta E_g}$$

Usually $\beta = 2$ to 4

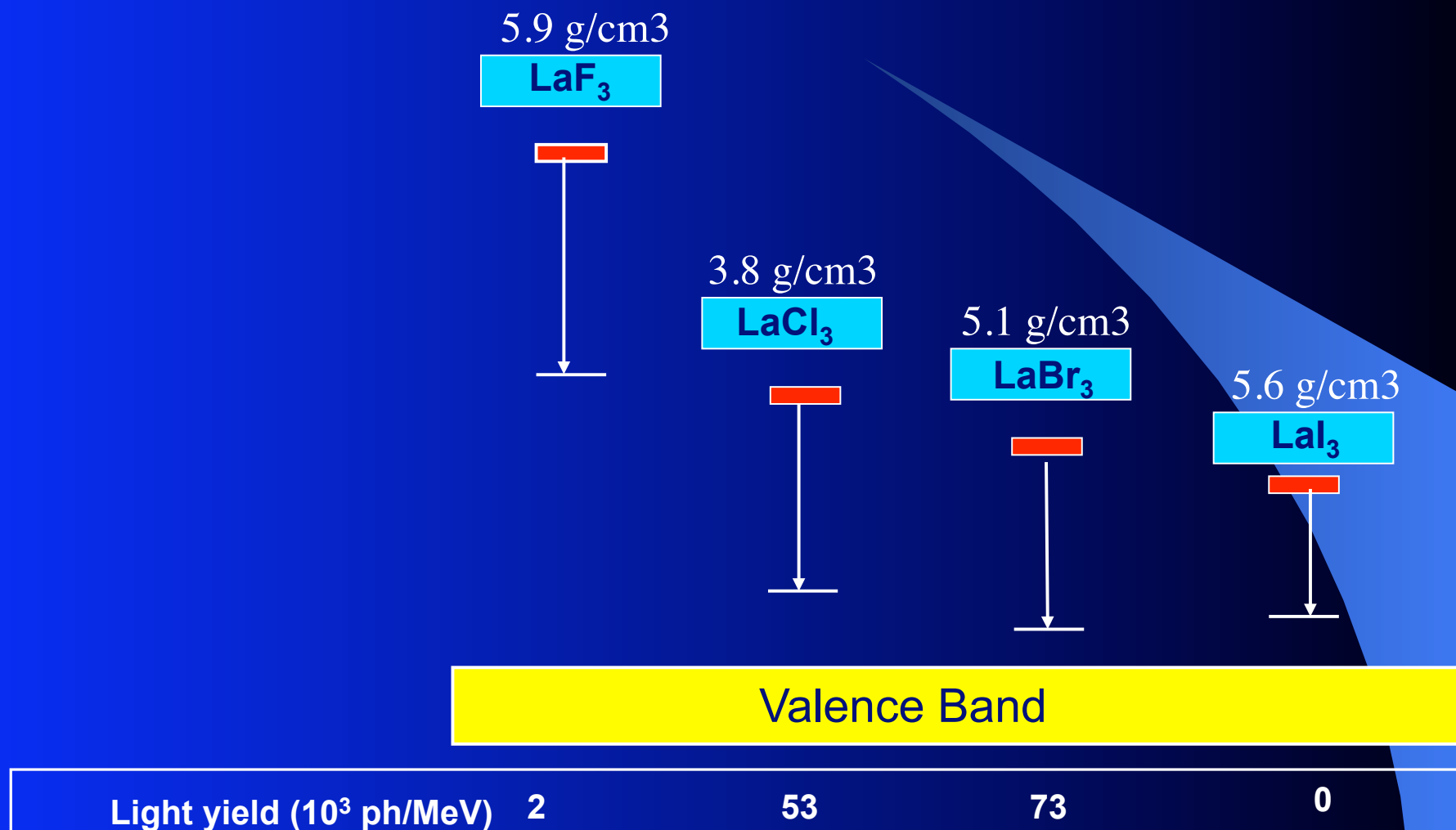
Fundamental limits to the LY



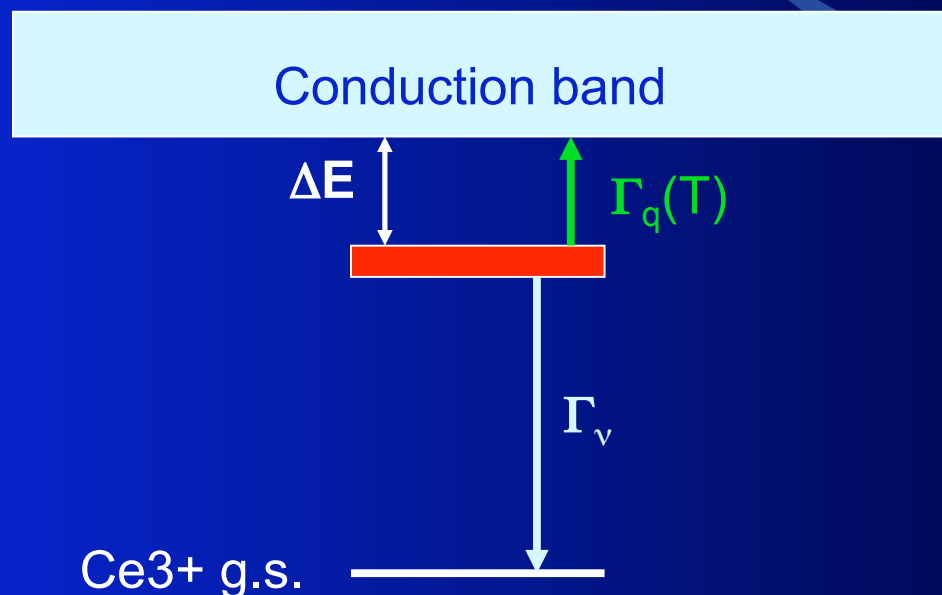
$$N_{ph} \leq N_{eh} = \frac{E_{\gamma}}{\beta E_{gap}}$$

Why?

Towards smaller band gap compounds

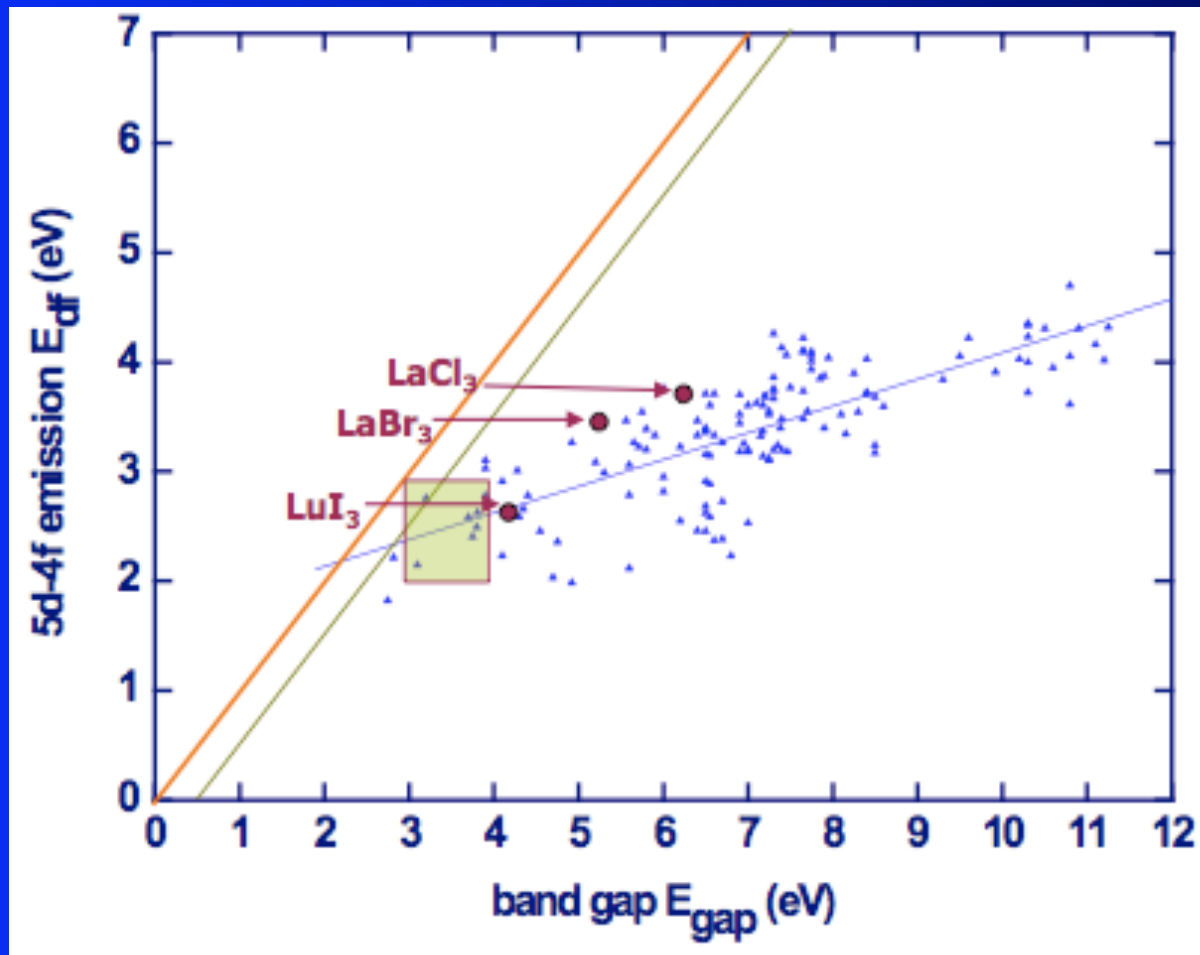


Luminescence quenching in $\text{LaI}_3:\text{Ce}$



Absolute location of doping levels is crucial

The host lattice must be transparent to the Ce^{3+} emission



Ultimate Ce^{3+}
Scintillator

$$E_{gap} > 2.5-3 \text{ eV}$$

$$Y_{ph} \leq 140000/\text{MeV}$$

$$\lambda_{em} \approx 600 \text{ nm}$$

Best result

$\text{LuI}_3: \text{Ce}^{3+}$

$$E_{gap} = 4.2 \text{ eV}$$

$$Y_{ph} = 100000/\text{MeV}$$

$$\lambda_{em} \approx 470 \text{ nm}$$



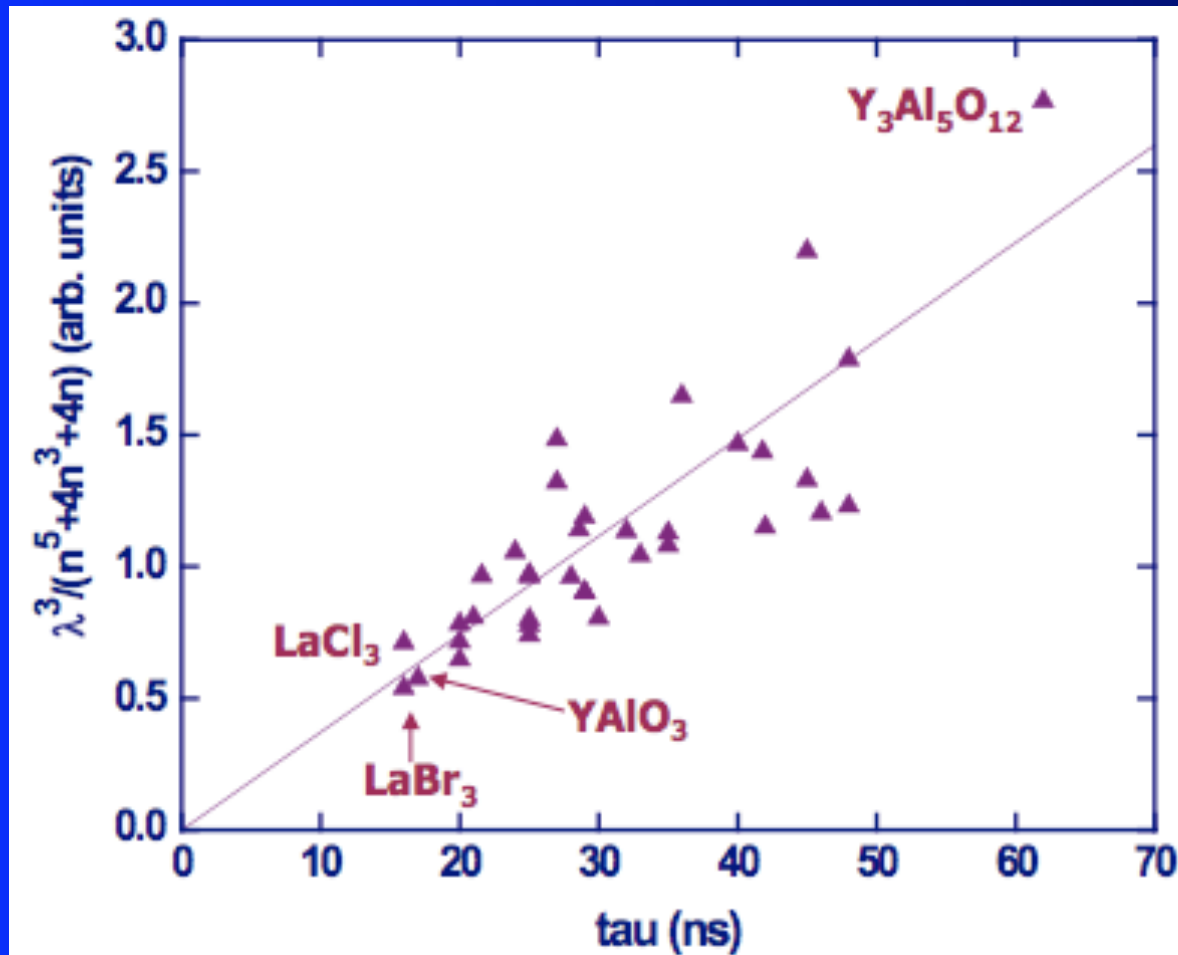
Why is Ce^{3+} so popular? Limits on the scintillation speed



$$\Gamma_v = \frac{1}{\tau_v} \propto \frac{n}{\lambda^3} \left(\frac{n^2 + 2}{3} \right)^2 \sum_f |\langle f | \mu | i \rangle|^2$$

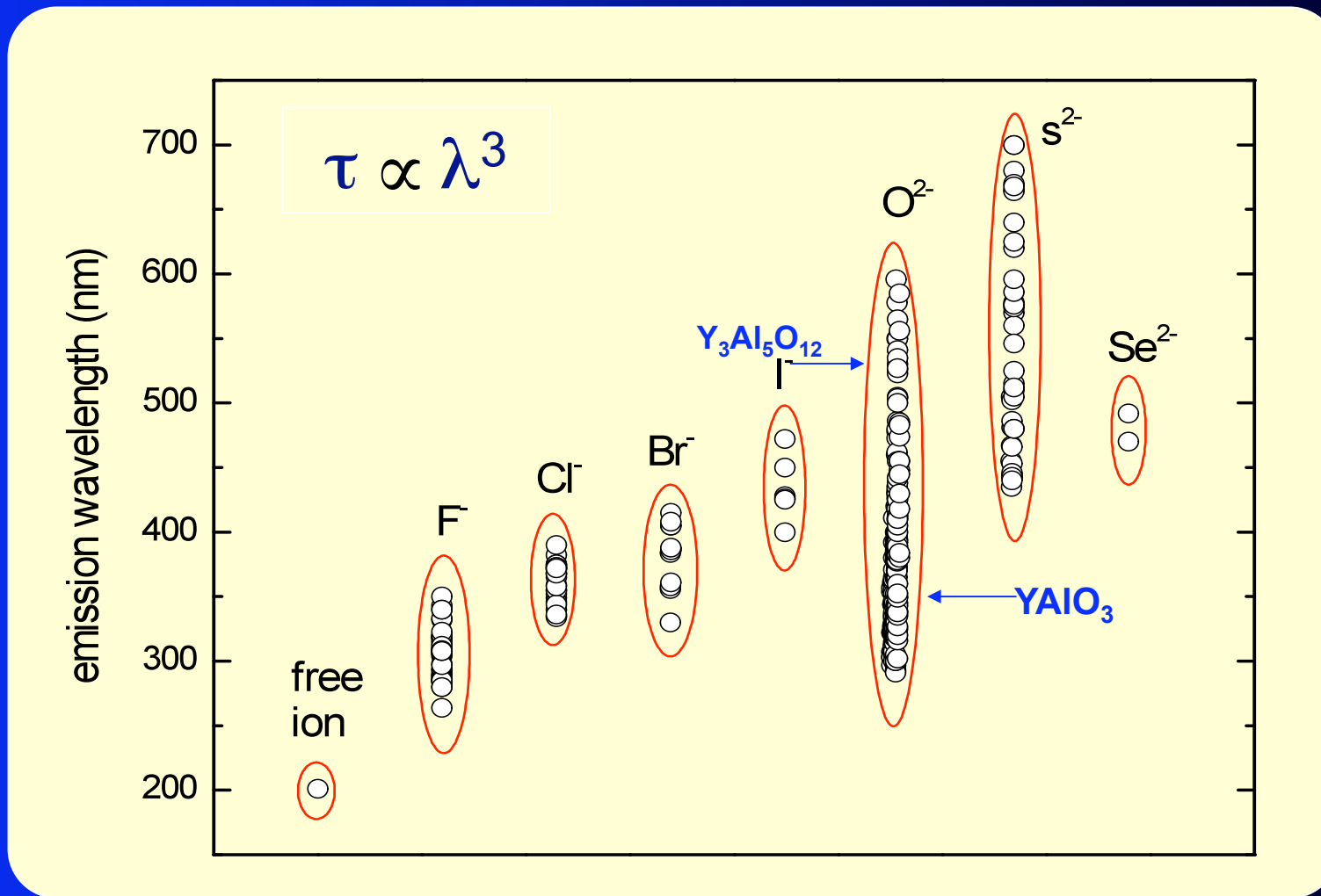
Three important aspects

- Dipole and spin allowed transitions
- Short wavelength of emission
- High refractive index

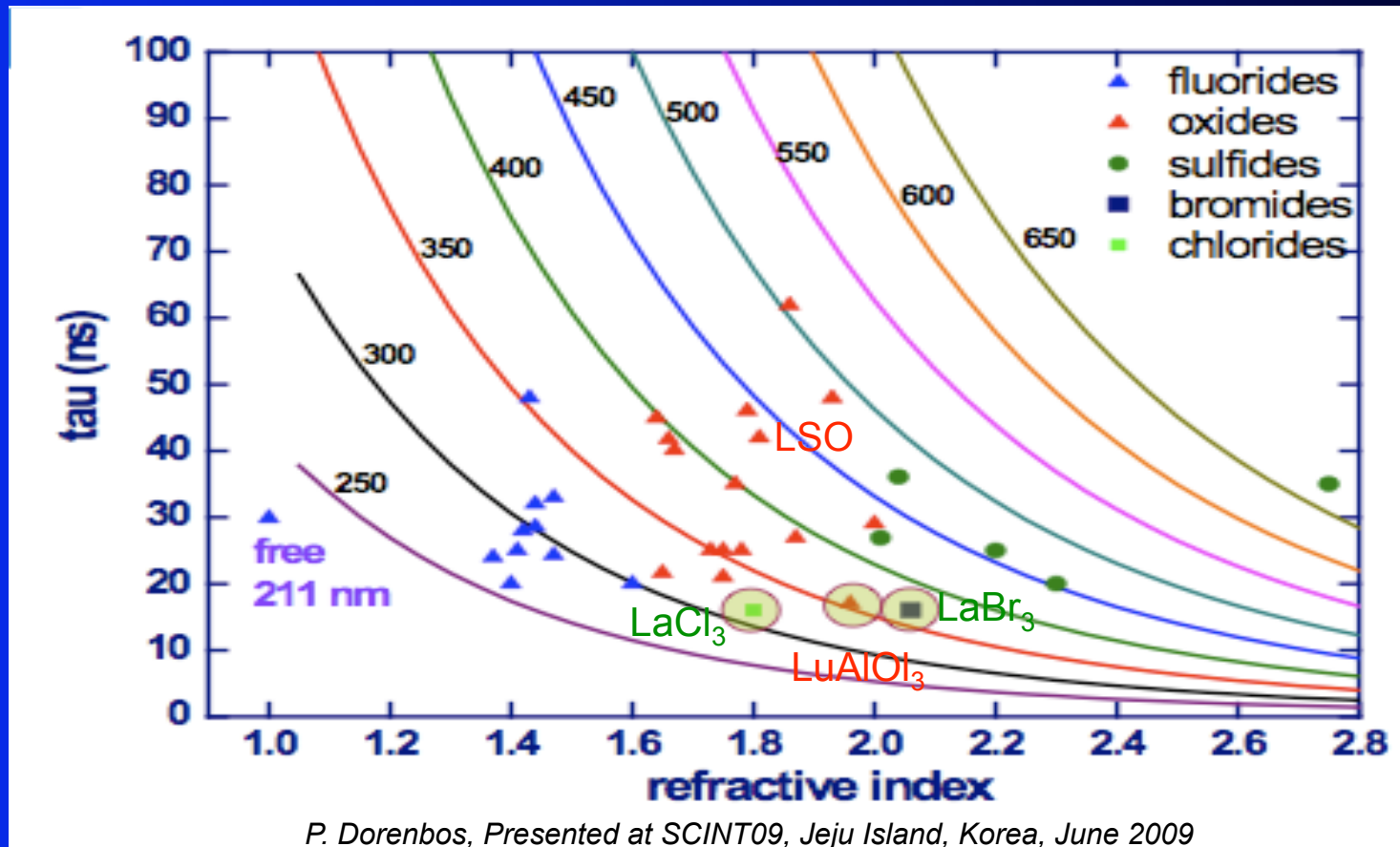


$$\tau \approx A \times \frac{\lambda^3}{n^5 + 4n^3 + 4n}$$

5d-4f emission wavelength of Ce^{3+}

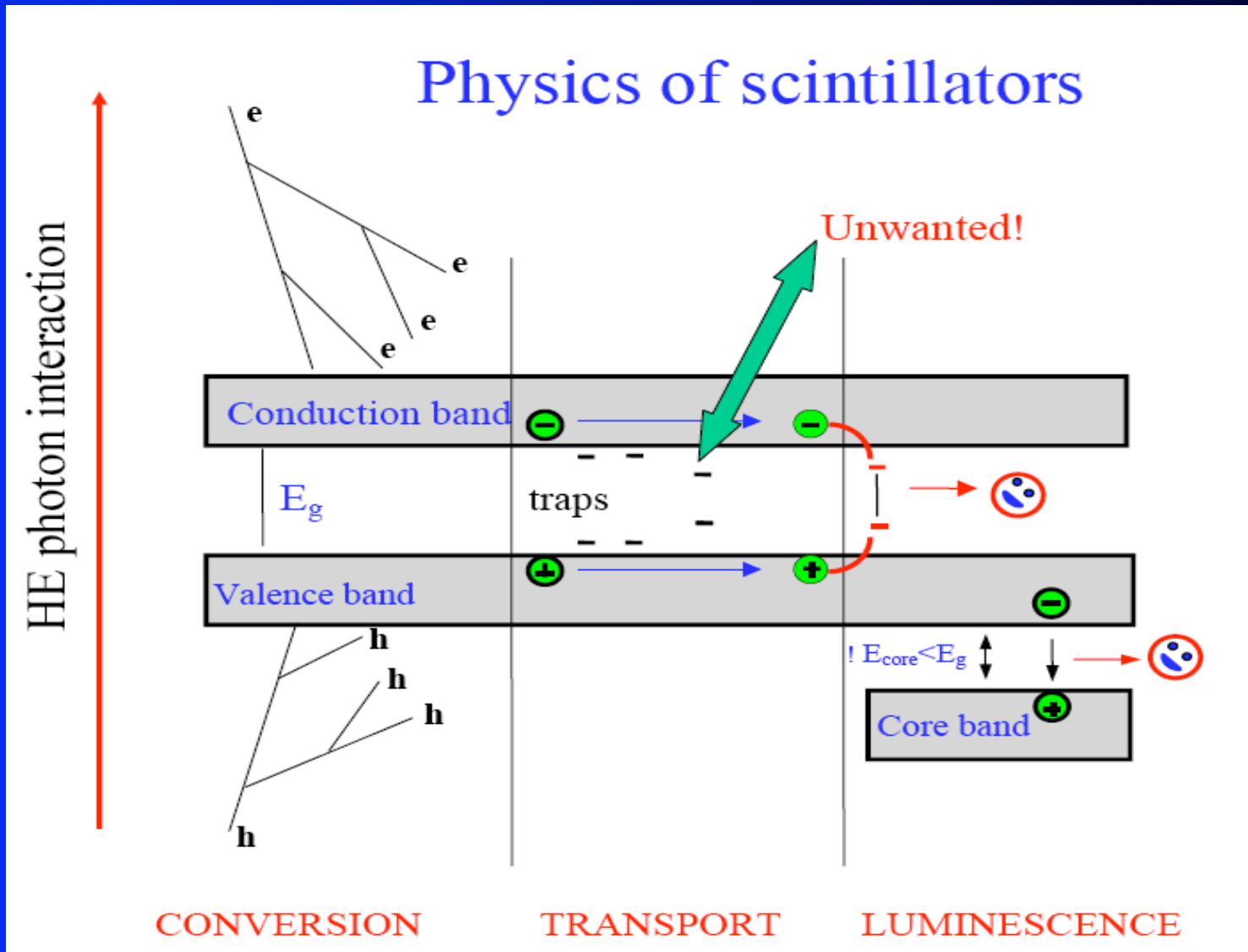


Decay time of Ce^{3+} 5d-4f emission



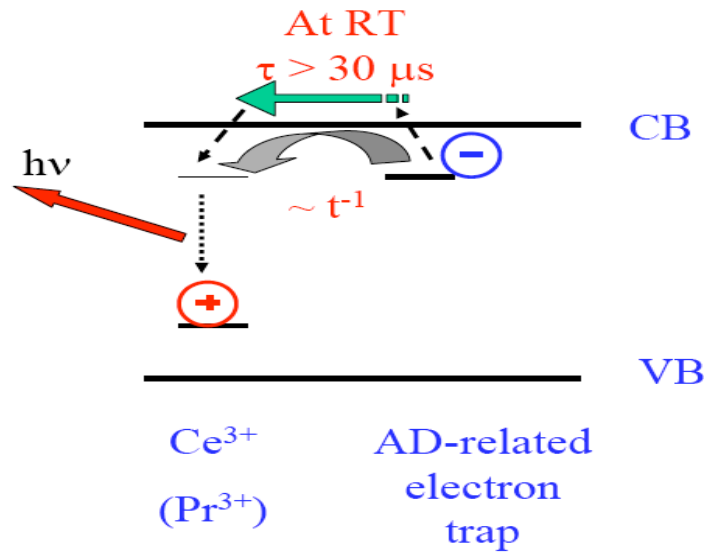
- λ tends to lengthen from F \rightarrow O \rightarrow Cl \rightarrow Br \rightarrow S 300nm-600nm
- n tends to increase from F \rightarrow O \rightarrow Cl \rightarrow Br \rightarrow S 1.4 – 2.4
- shortest τ of 17ns for YAlO_3 , LuAlO_3 , LaCl_3 , LaBr_3

Effect of traps



Effect of traps

Ce³⁺ and Pr³⁺-doped Lu₃Al₅O₁₂



Light yield (1 μs time gate)

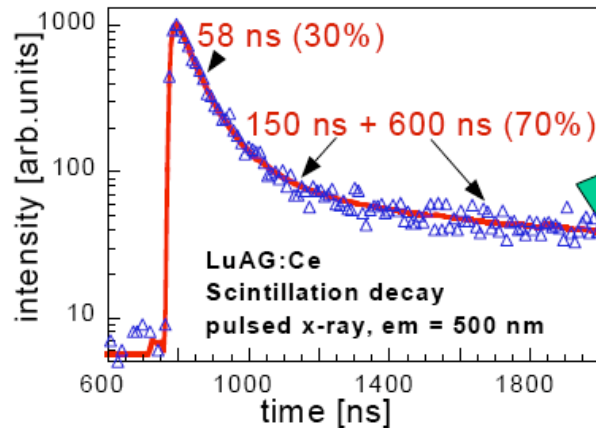
Best YAG:Ce ~ 3x BGO

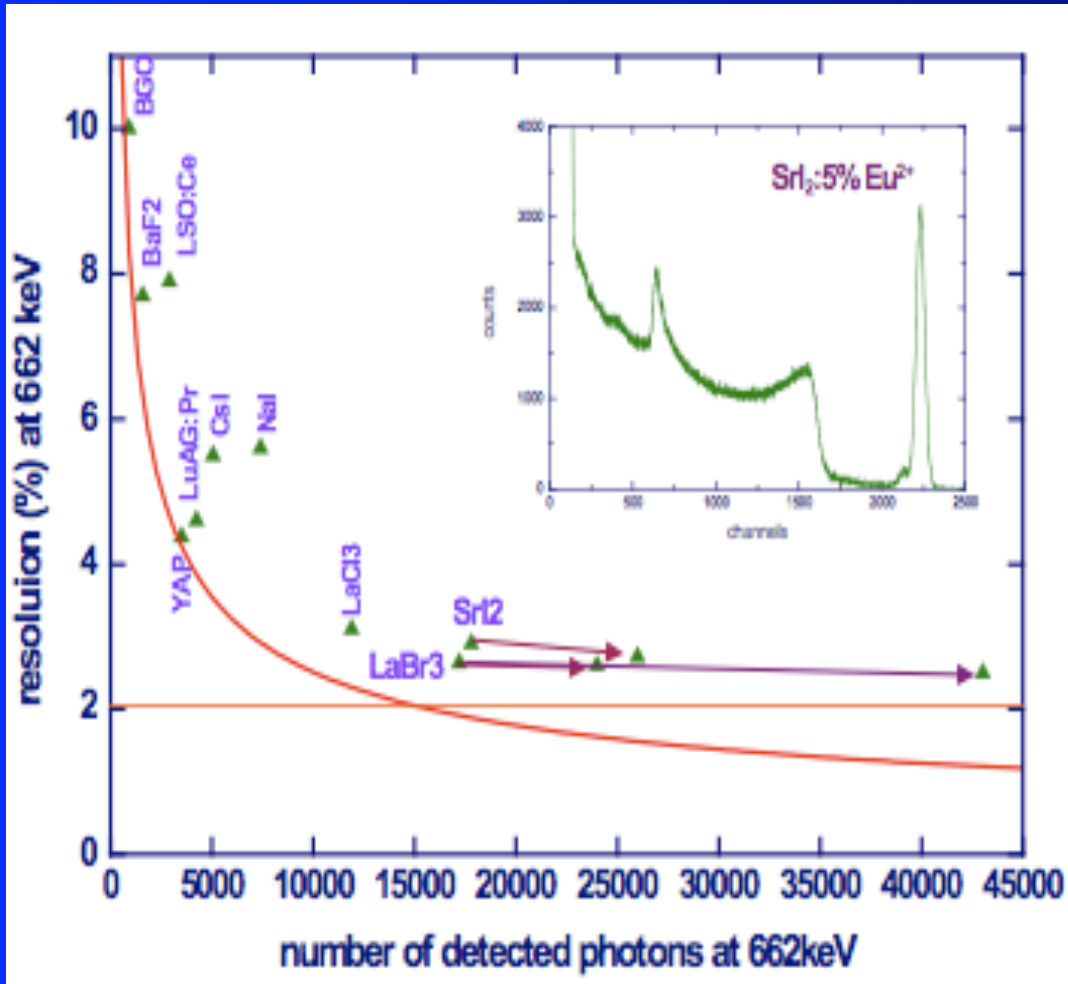
Best LuAG:Ce ~ 60% of YAG:Ce

A lot of “slow light”
in these materials

Retrapping of electrons at
shallow traps before their
radiative recombination at
Ce³⁺ ions

Nikl et al, pss (a) **201**, R41 (2004)

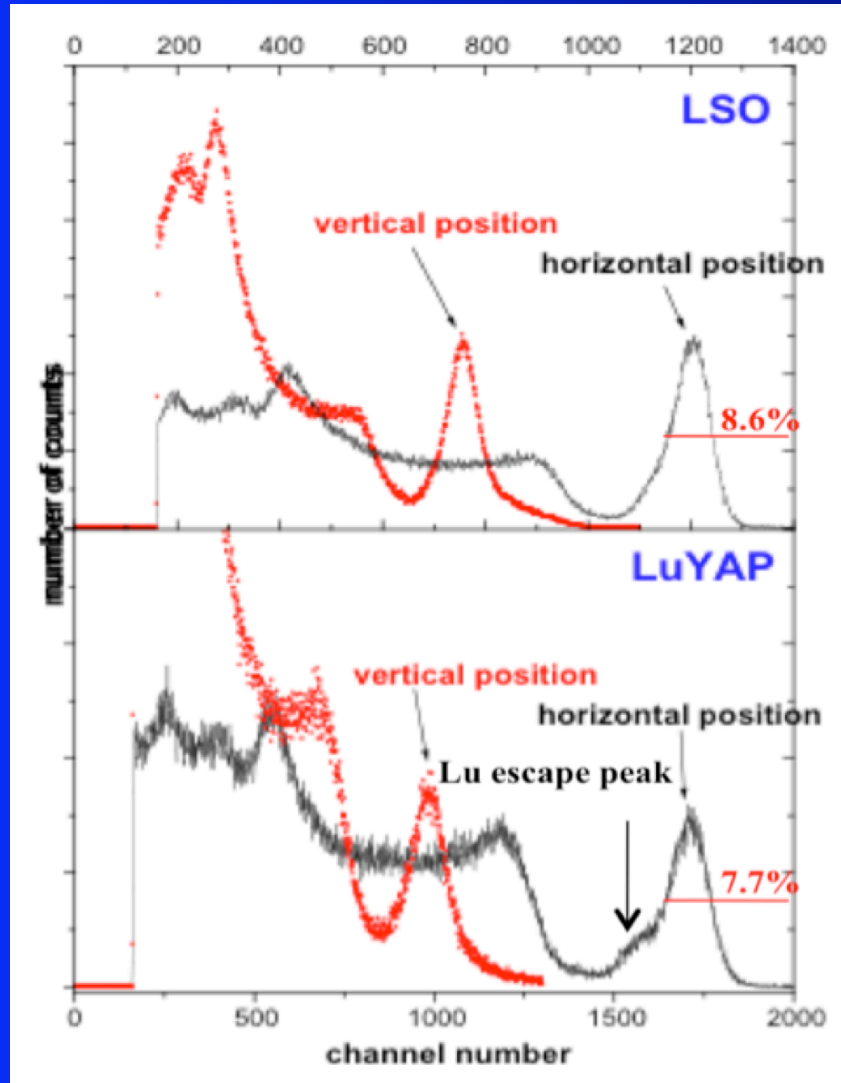




$$R_{\text{stat}} = 2.35 \sqrt{\frac{1+0.15}{N_{\text{dph}}}}$$

- Some crystals are far from the theoretical limit
- LaBr₃ and SrI₂ measured with standart and new high QE PMT
- Significantly higher number of detected photons
- No significant improvement resolution
- Can we pass the 2% barrier?

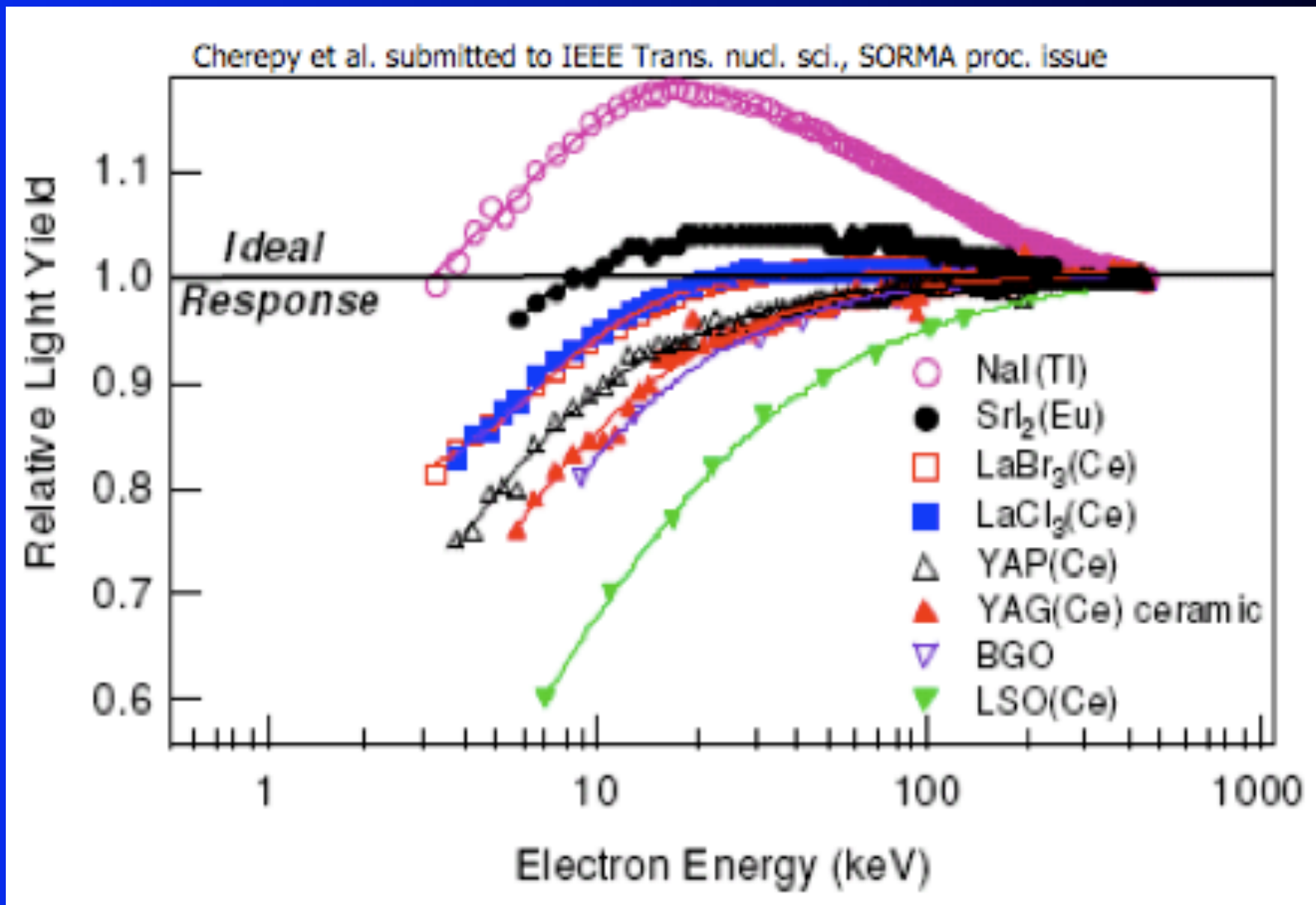
Comparison LSO - LuYAP



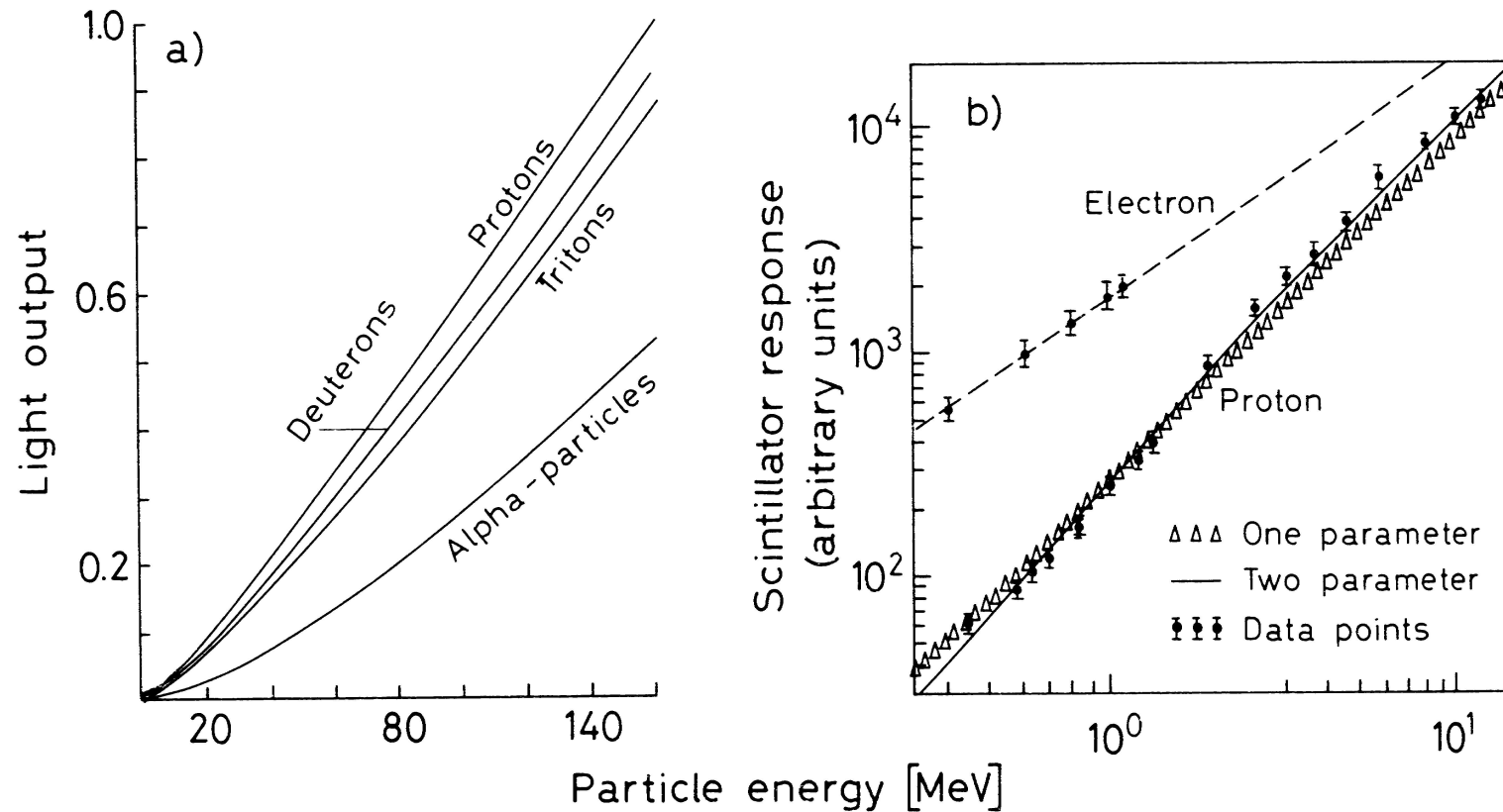
- 2x2x10 mm³ crystals
 - LSO: 6000 pe/MeV
 - LuYAP: 2000 pe/MeV
- Each crystal in 2 positions
 - Vertical
 - Horizontal
- Gain LuYAP=3xgain LSO

Same energy resolution

Low energy non-proportionality

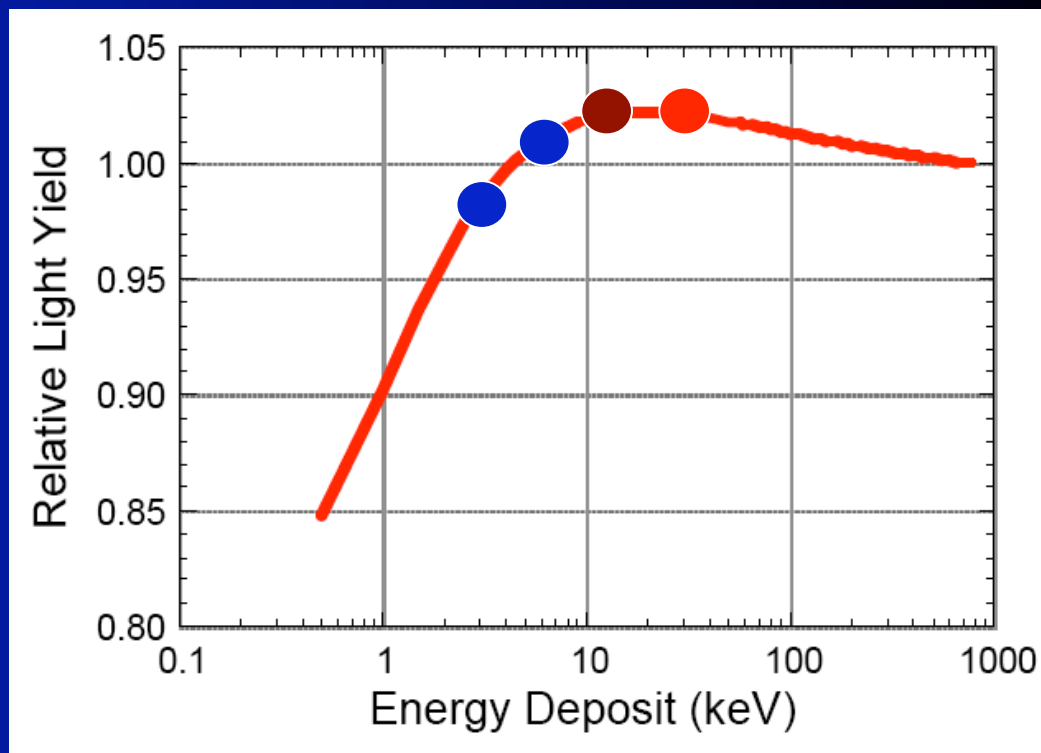
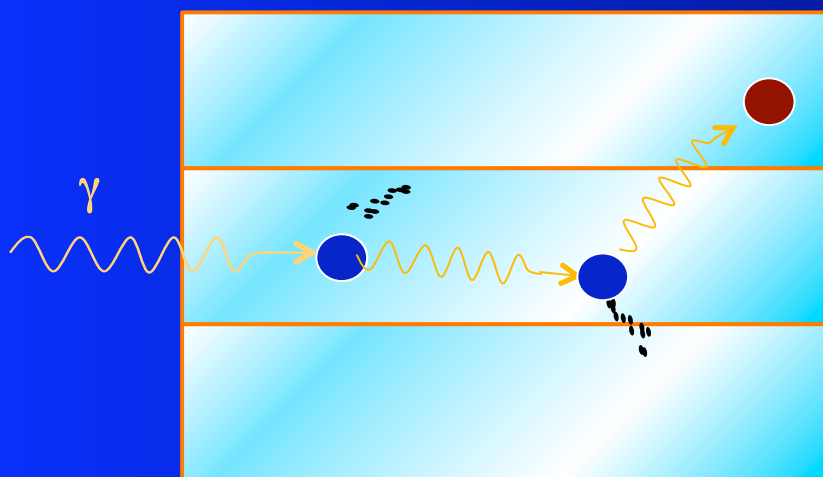


Non-proportionality

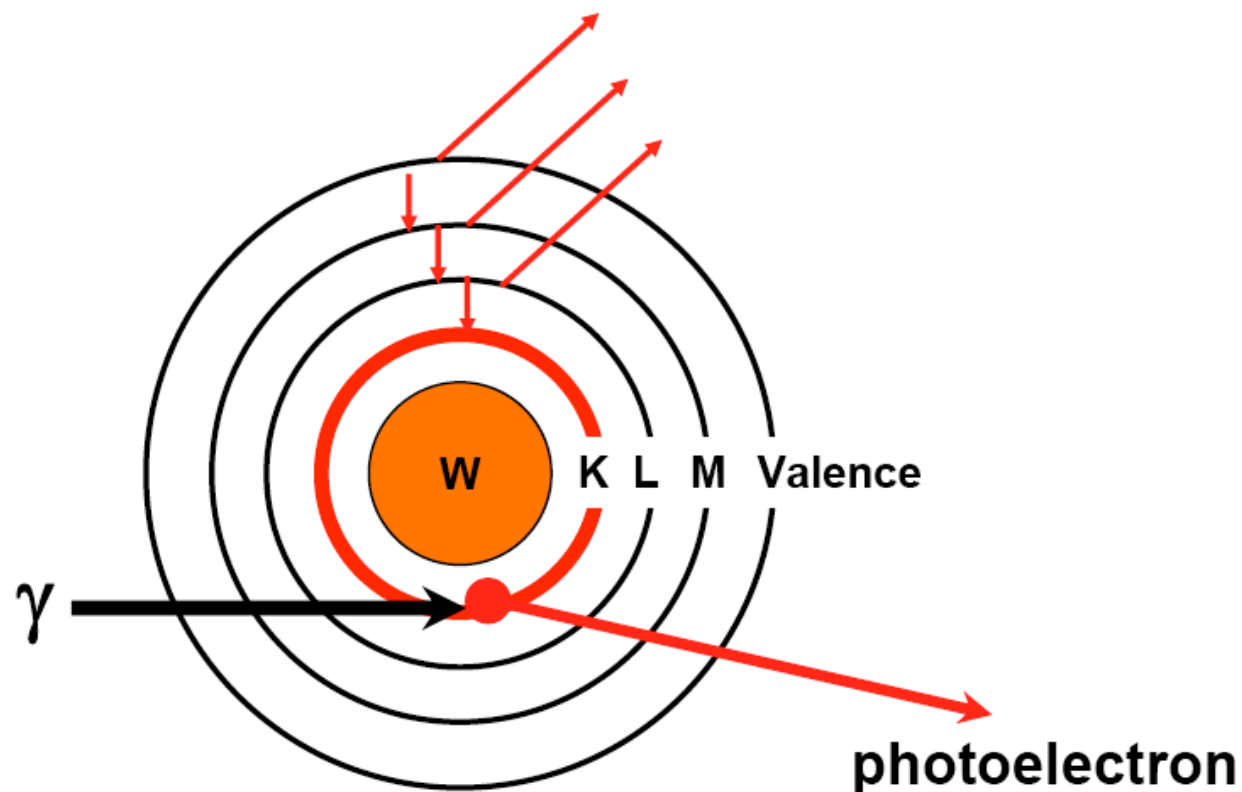


- **Light Yield *Not* Constant**
- **Depends on Particle Energy & Type**

Initial interaction: Compton versus Photoelectric

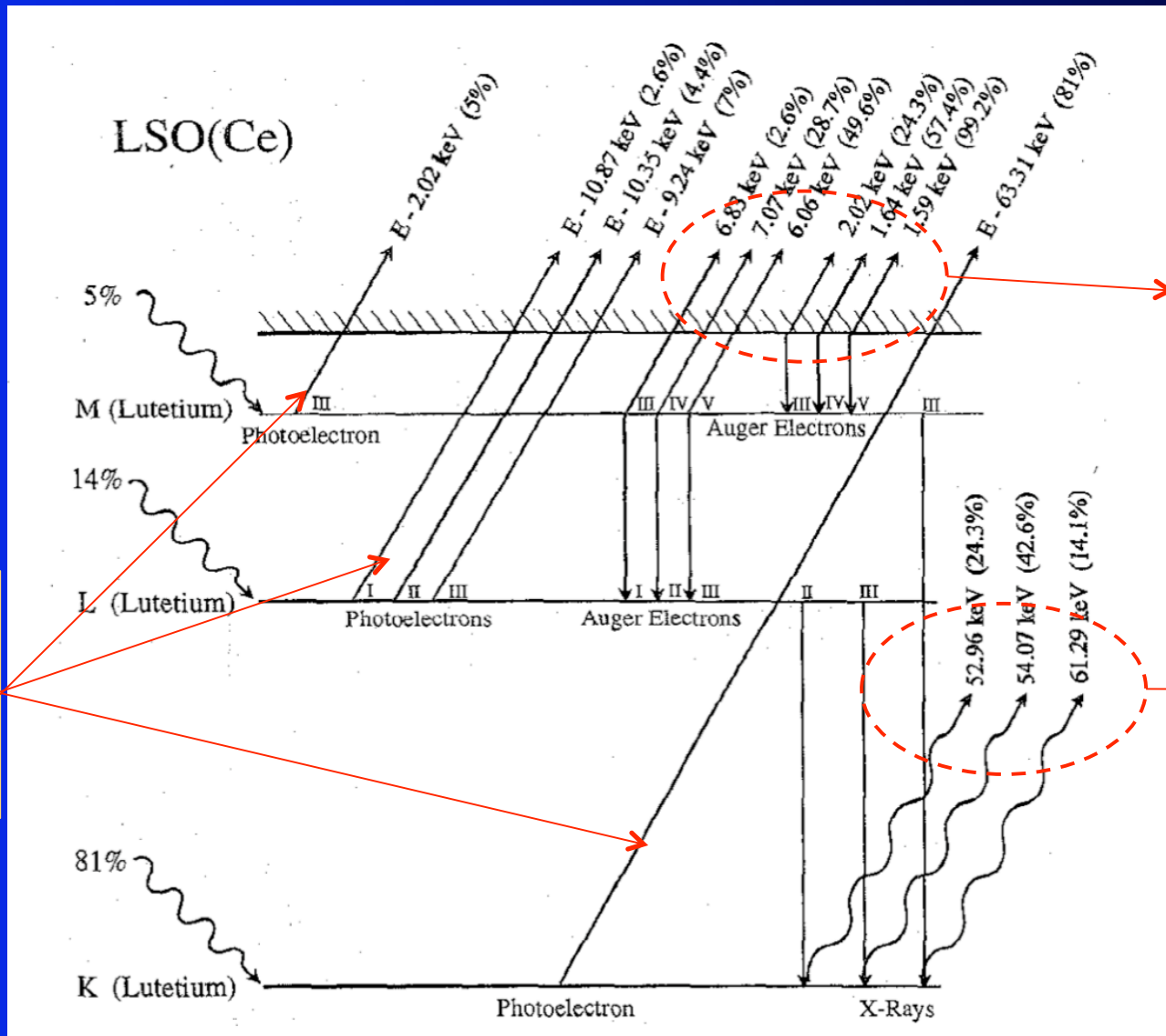


**Non-Proportionality + Multiple Energy Deposit
⇒ Degraded Energy Resolution**



- Usually Occur with *Inner* Shell Electrons
 ⇒ Inner Shell Hole Filled via Cascade

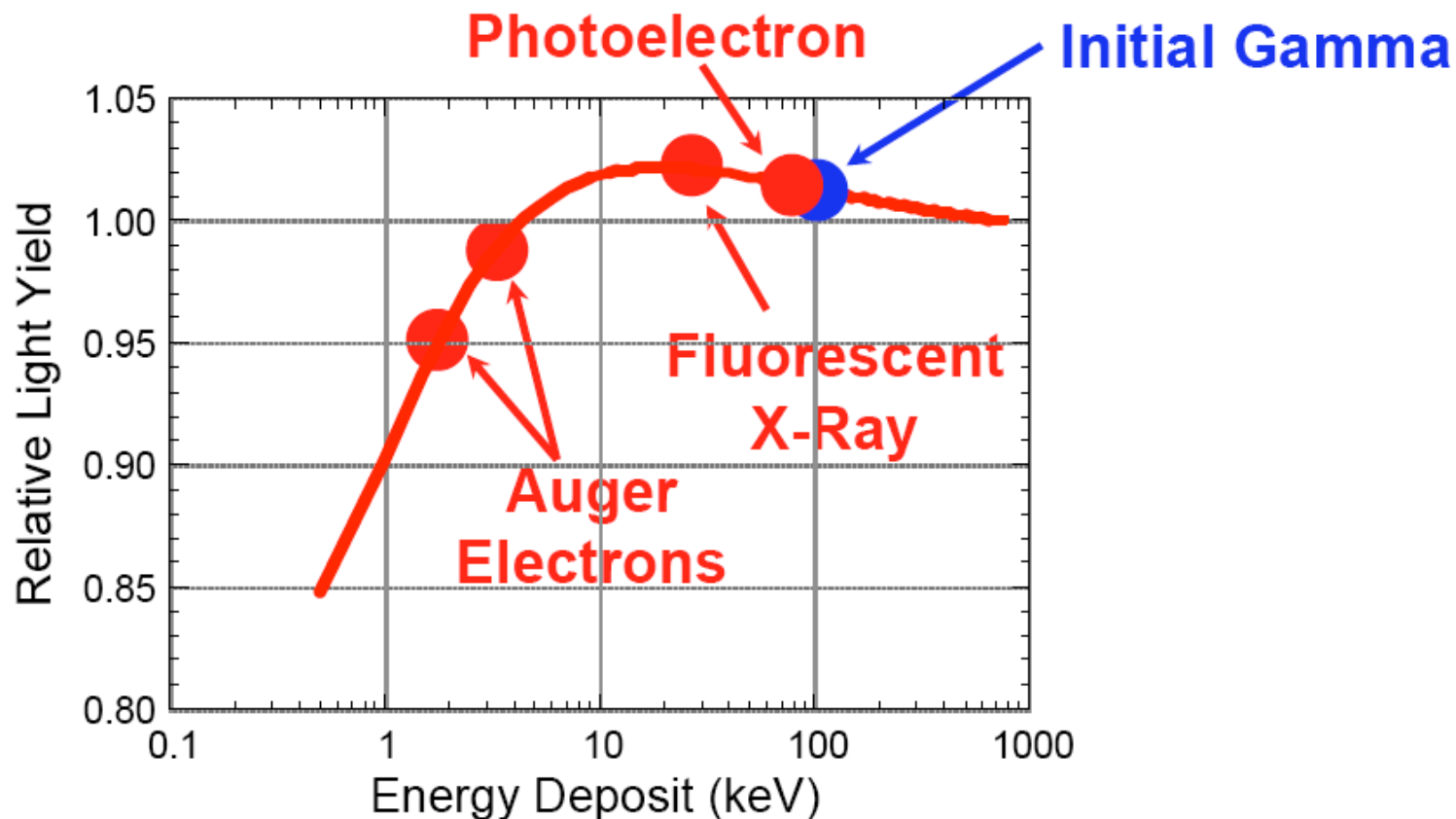
Simplified cascade model for LSO



Different photo-electron energies

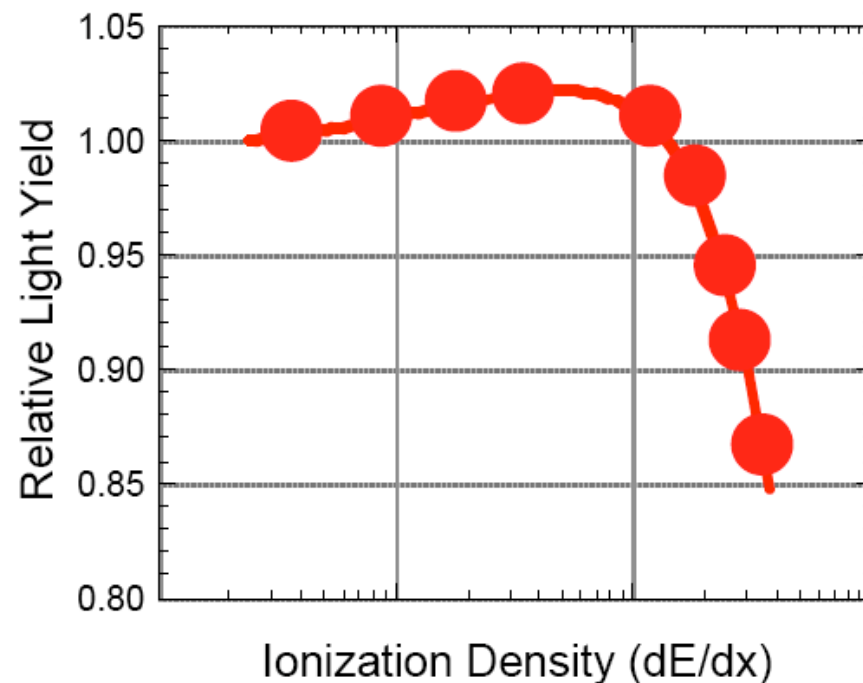
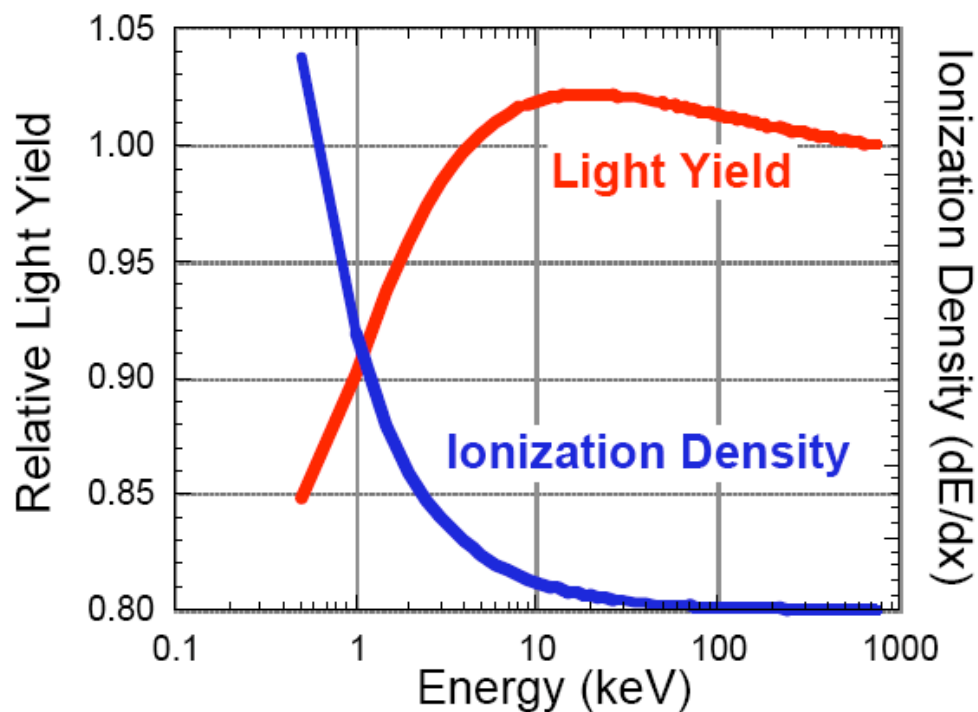
1 to 7 keV Auger electrons

50 to 60 KeV Fluorescent X-rays



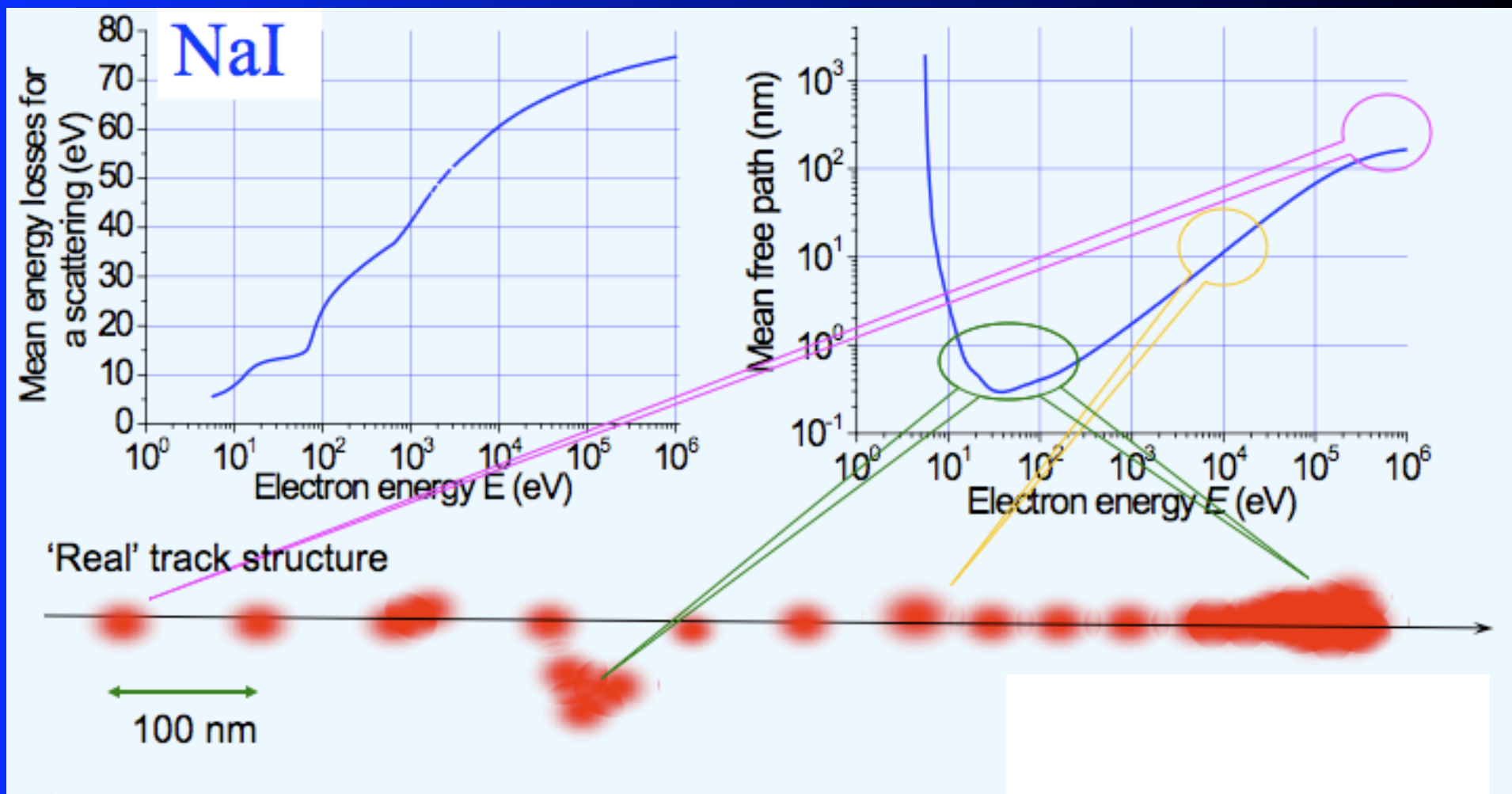
**Non-Proportionality + Multiple Energy Deposit
 ⇒ Degraded Energy Resolution**

Yield depends on electron ionization density



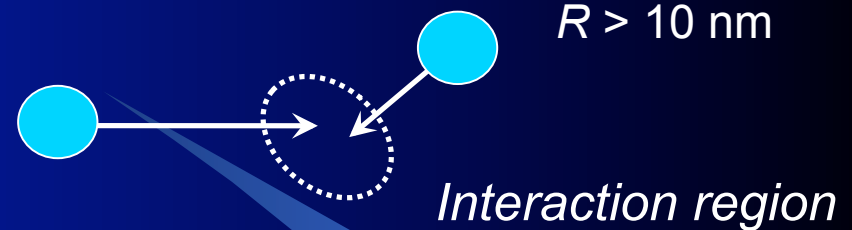
**Non-Proportionality + Non-Uniform Energy Deposit
 ⇒ Degraded Energy Resolution**

Non-uniformity of electron energy deposit

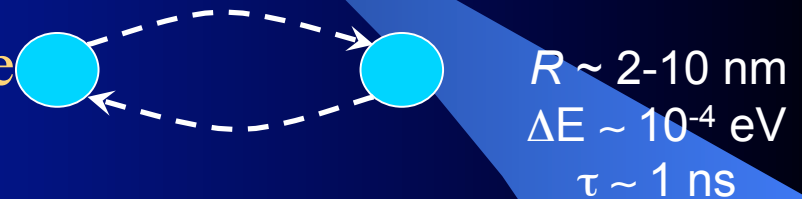


EE interactions in scintillators

- Weak interaction by collision

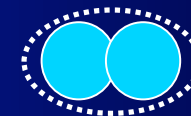


- Interaction due to energy exchange



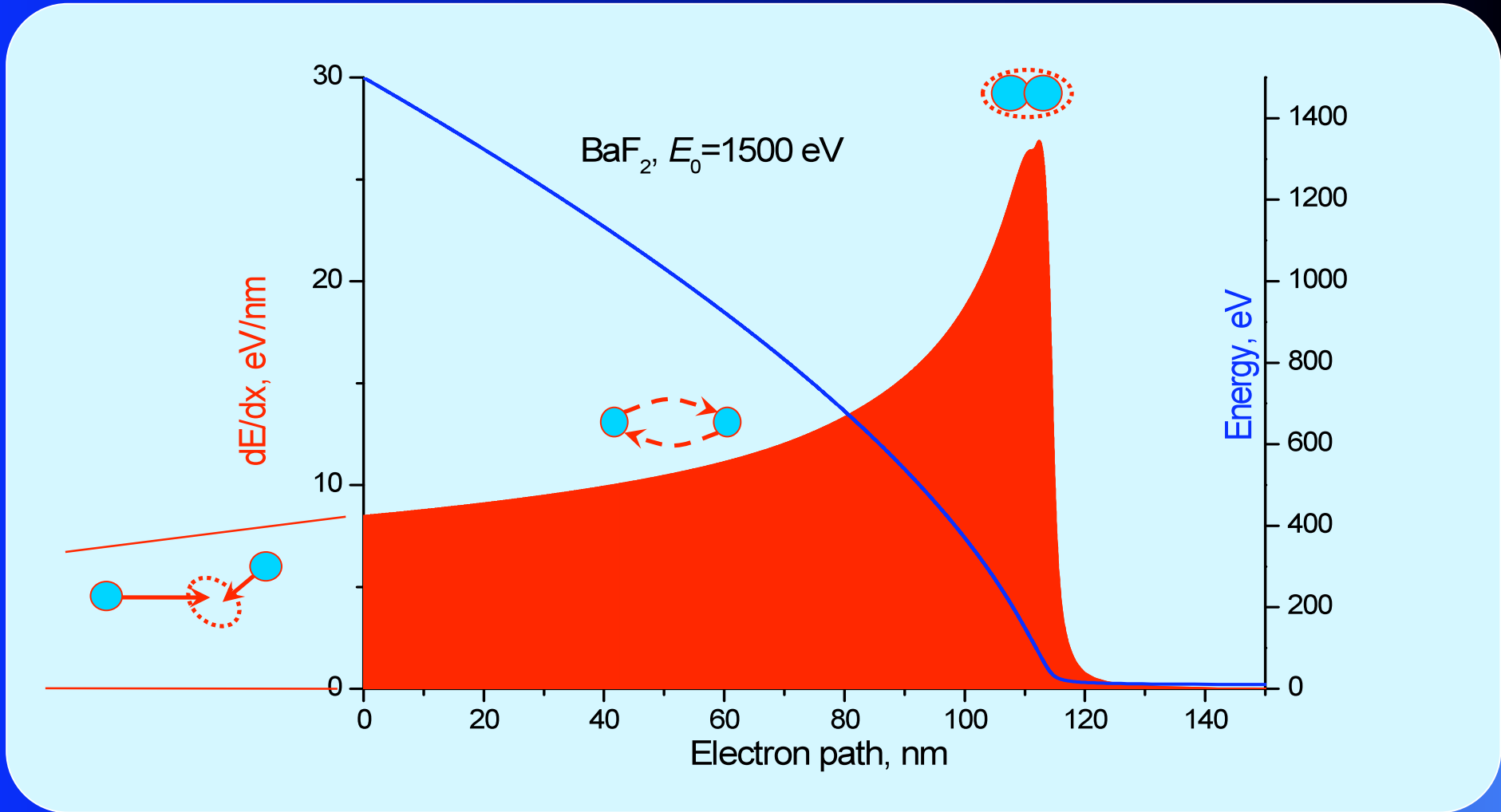
- Strong interaction:

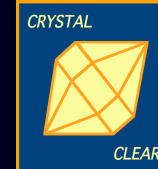
- new type of excitations
- specific paths of localization
- defects creation



$R < 2 \text{ nm}$
 $\Delta E \sim 0.1 \text{ eV}$
 $\tau \sim 10 \text{ fs}$

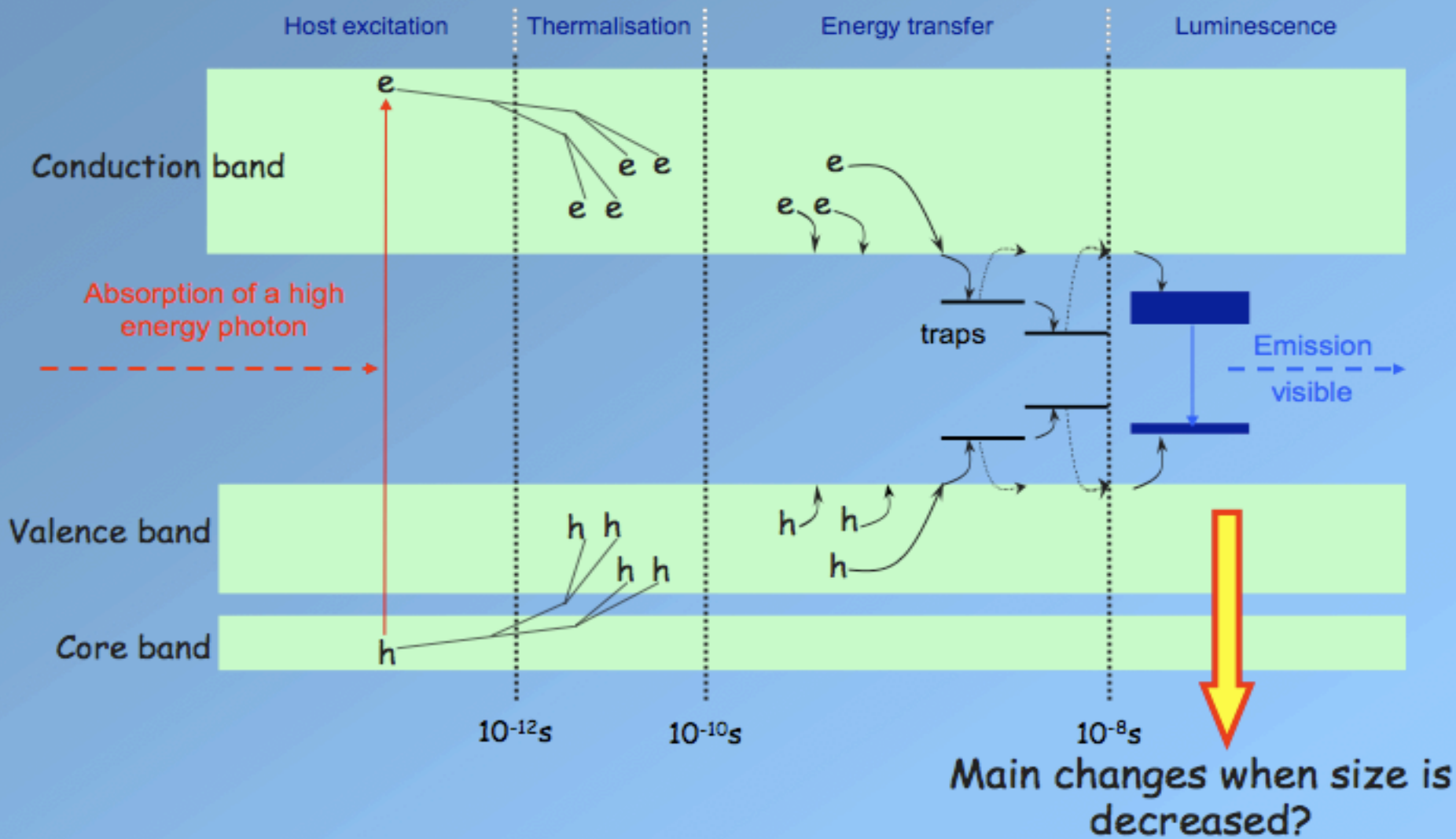
EE density effect in scintillators





End of Part 1

Scintillation process



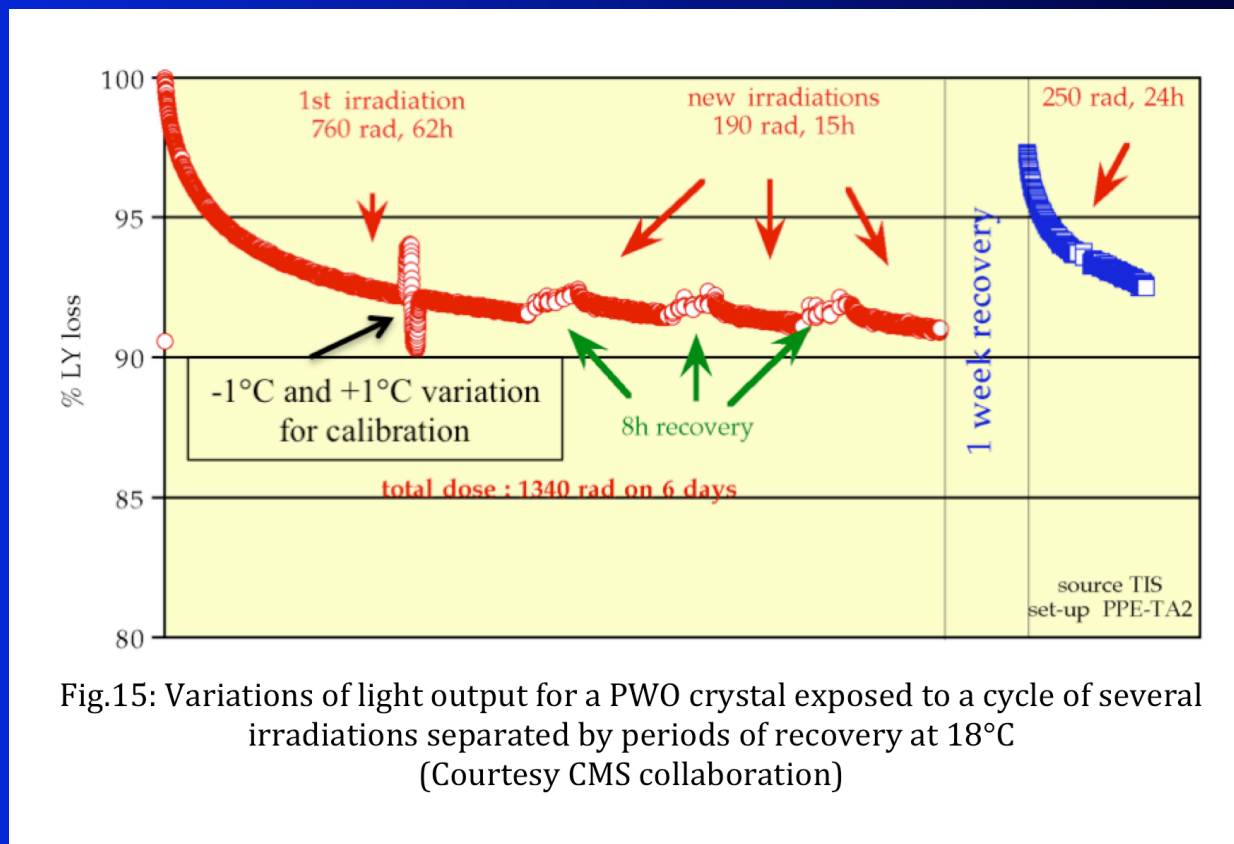
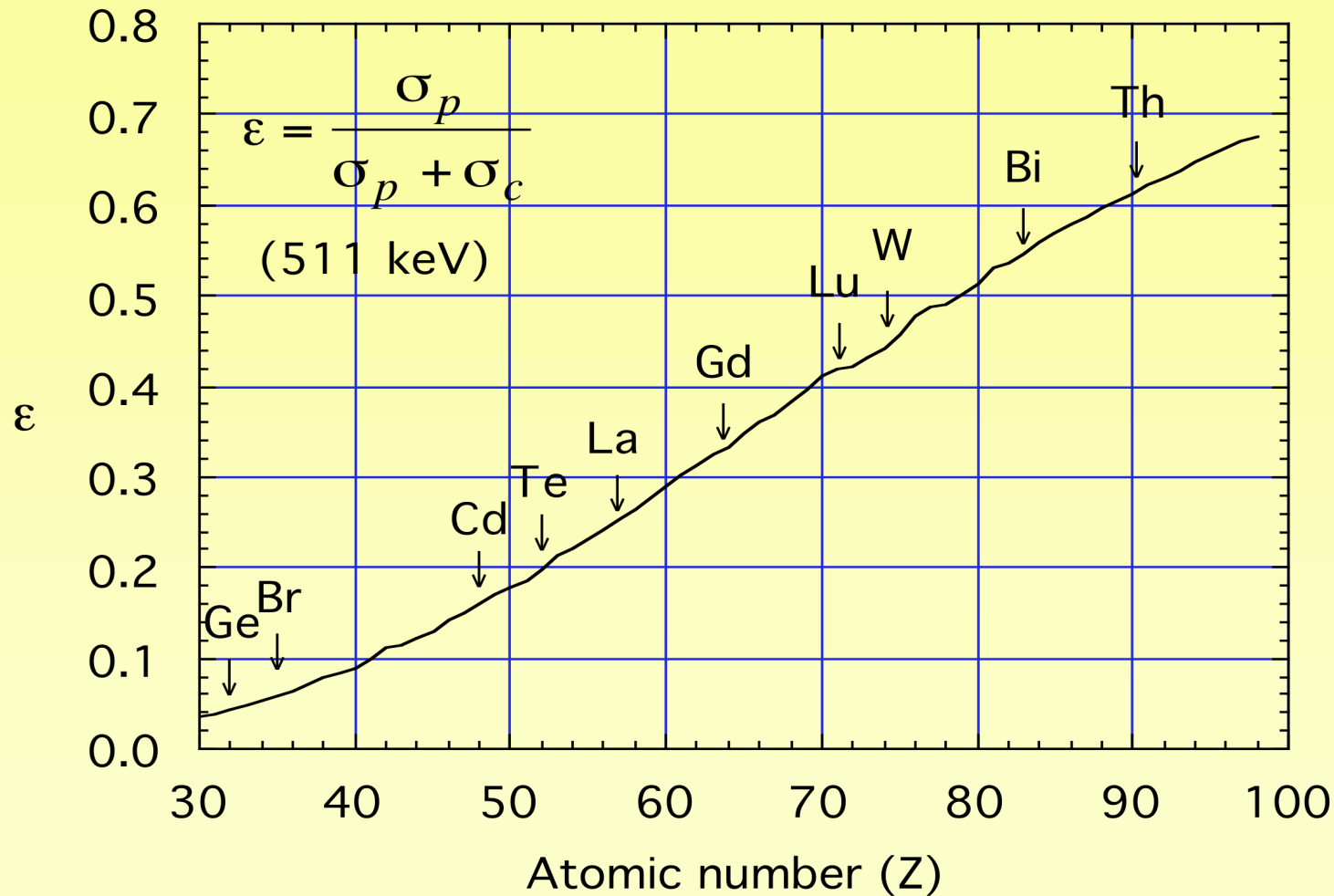


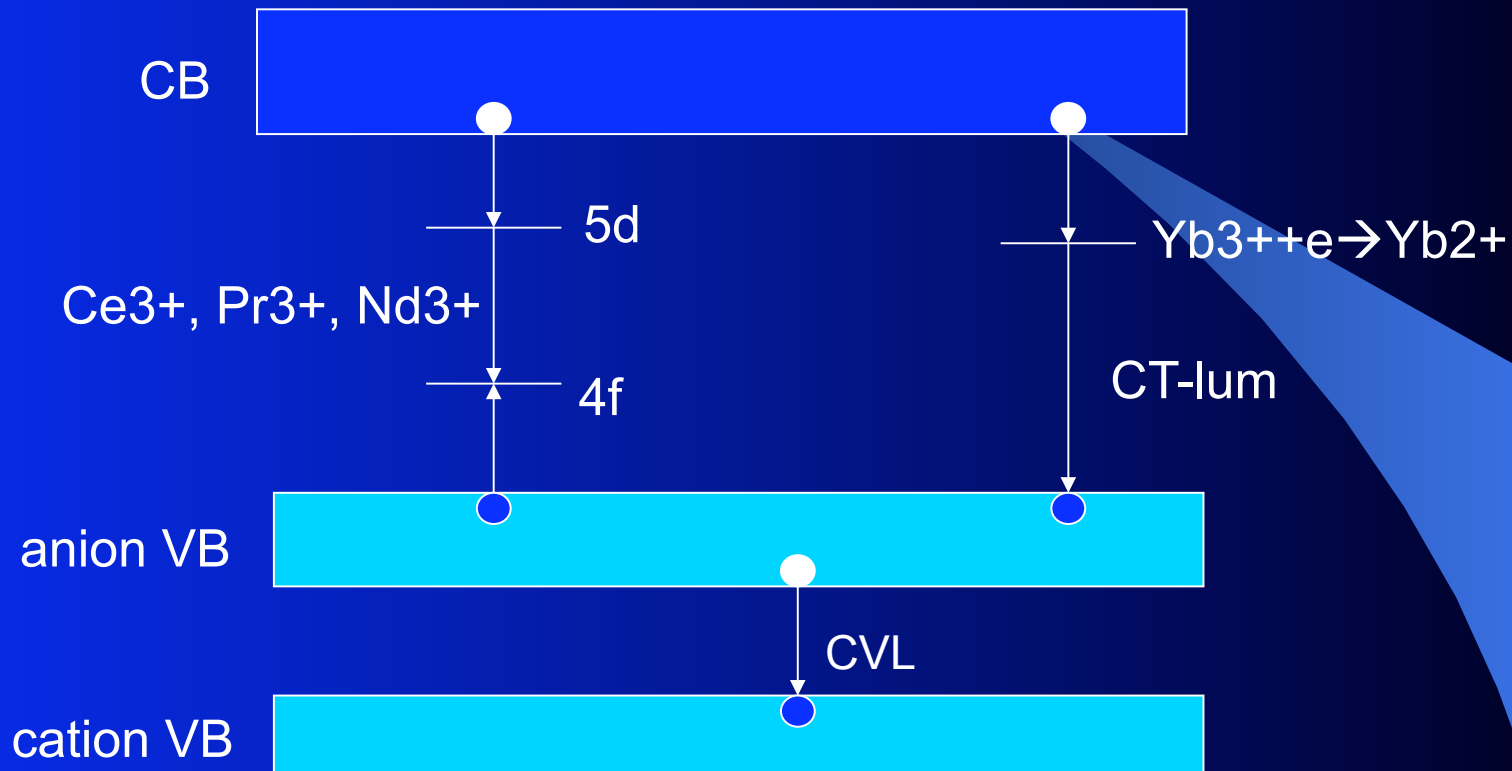
Fig.15: Variations of light output for a PWO crystal exposed to a cycle of several irradiations separated by periods of recovery at 18°C
(Courtesy CMS collaboration)

511 keV Photofraction vs. Z

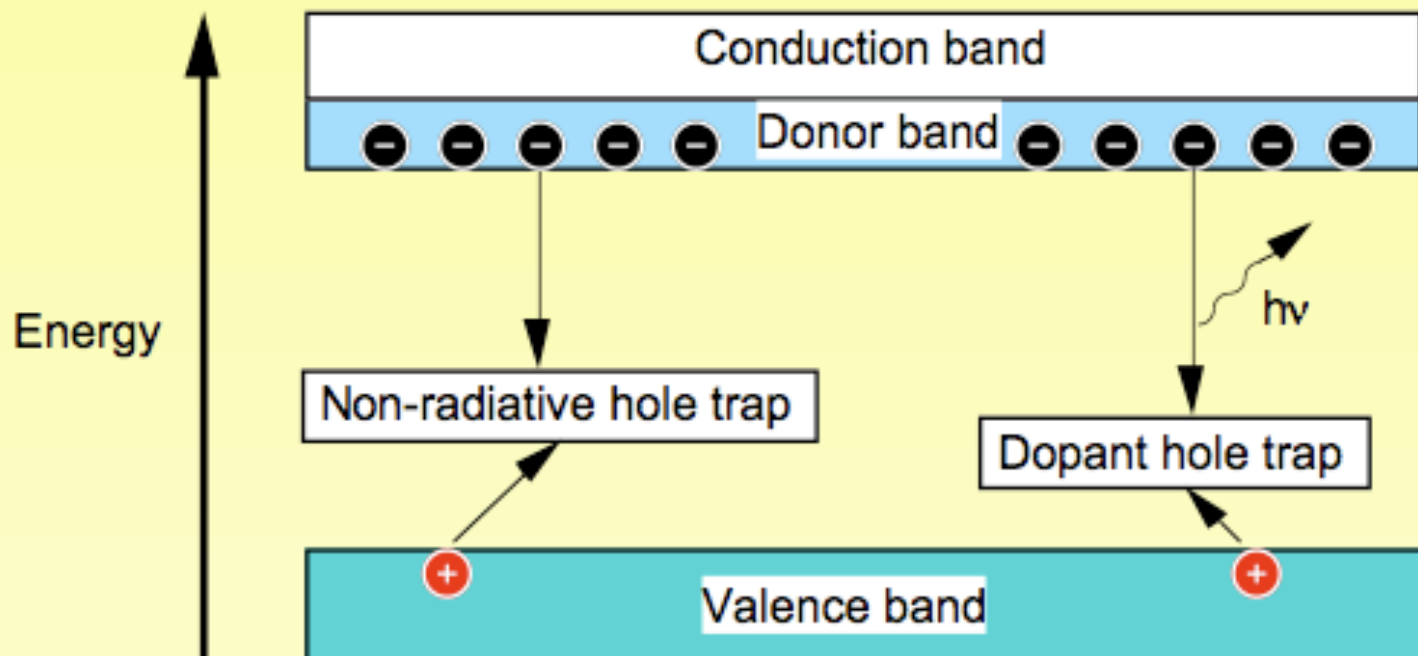


Dipole and spin allowed transitions

Charge transfer or 5d-4f transitions



Co-doped semiconductors Faster than Ce³⁺?



- Direct-gap semiconductor host with $E_g = >2.5$ eV
- Prompt (<50 ps), efficient trapping of hot holes by dopant ions
- Fast (~1 ns) recombination with donor band electrons

