



STRUCTURES
CLUSTER OF
EXCELLENCE



HGS-HIRe for FAIR
Helmholtz Graduate School for Hadron and Ion Research

Towards the phase diagram of QCD and its critical endpoint

Based on
[arXiv:2408.08413]

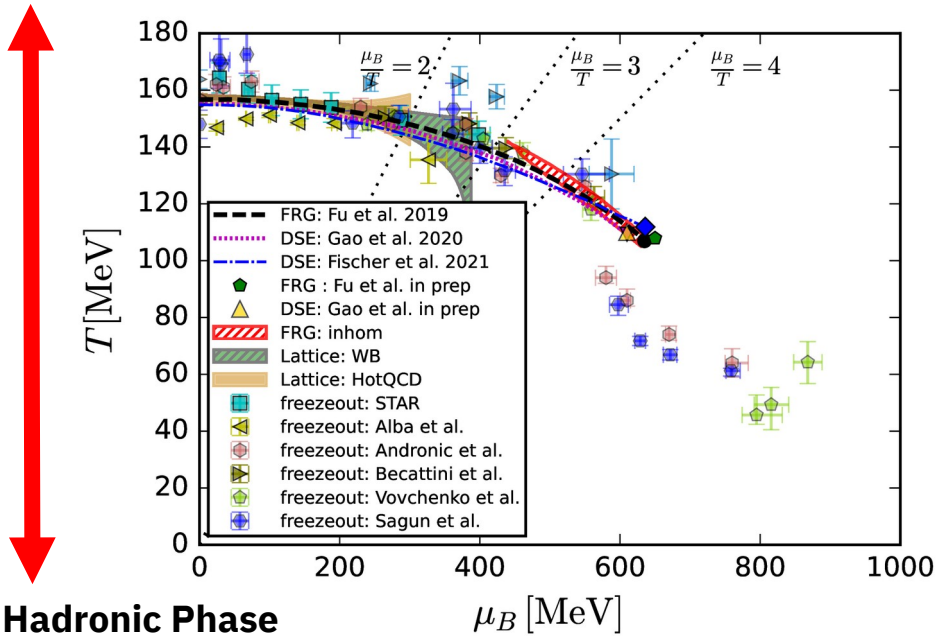
Franz R. Sattler
ITP Heidelberg

In Collaboration with
Friederike Ihssen, Jan M. Pawłowski, Nicolas Wink

ERG2024,
Les Diablerets
September 23, 2024

The QCD phase diagram

Quark-Gluon Plasma



Phase diagram shows
Chiral symmetry breaking
 i.e. condensation of $\langle \bar{q}q \rangle$

Fu, Pawłowski, Rennecke [Phys. Rev. D 101 (2020), 054032]

Gao, Pawłowski [Phys.Lett.B820(2021) 136584]

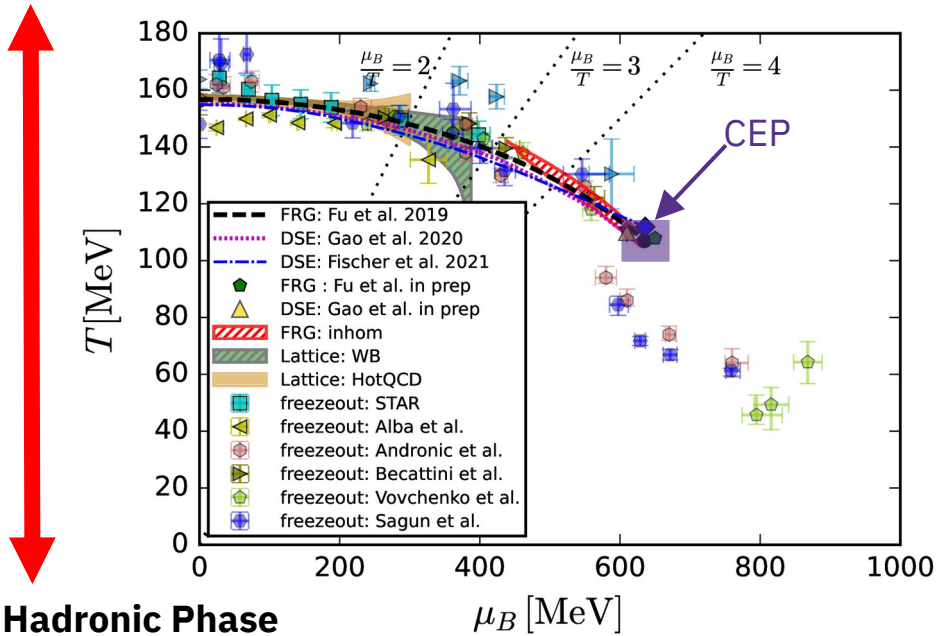
Gunkel, Fischer [Phys.Rev.D 104 (2021) 5, 054022]

Bellwied et al. (WB) [Phys.Lett.B 751 (2015) 559-564]

Bazavov et al. (HotQCD) [Phys.Lett.B 795 (2019) 15-21]

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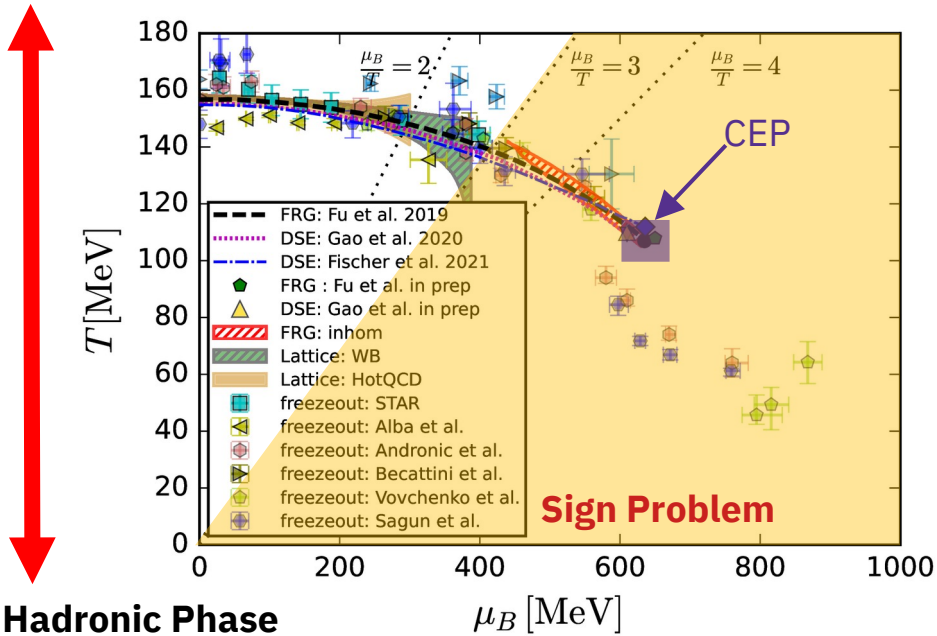
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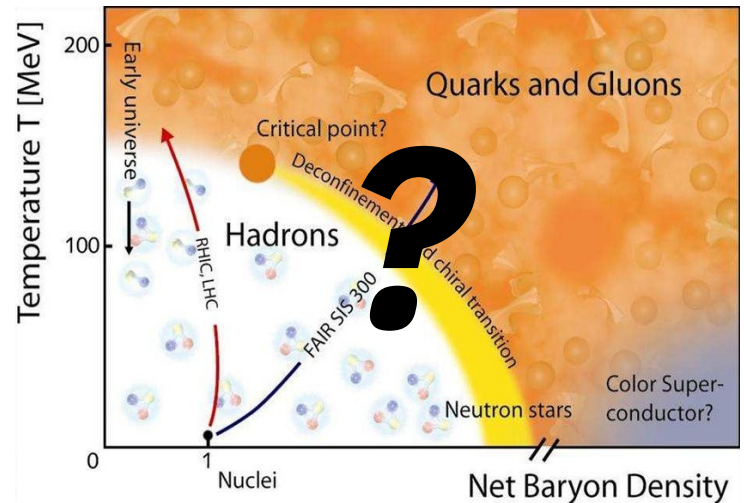
The QCD phase diagram

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Phase Diagram for intermediate μ not know

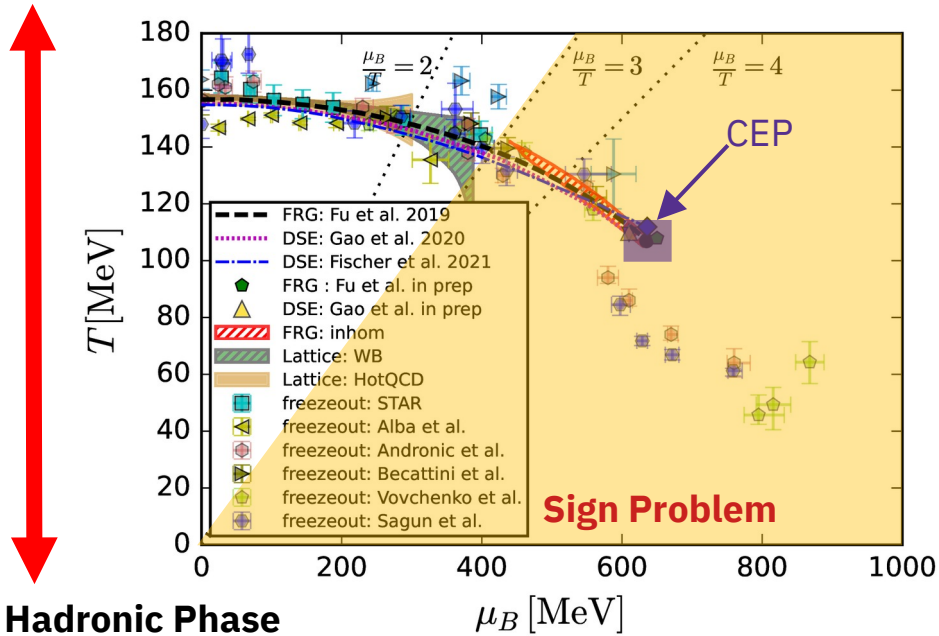


Sketch by GSI Darmstadt

Fu, Pawłowski, Rennecke [Phys. Rev. D 101 (2020), 054032]
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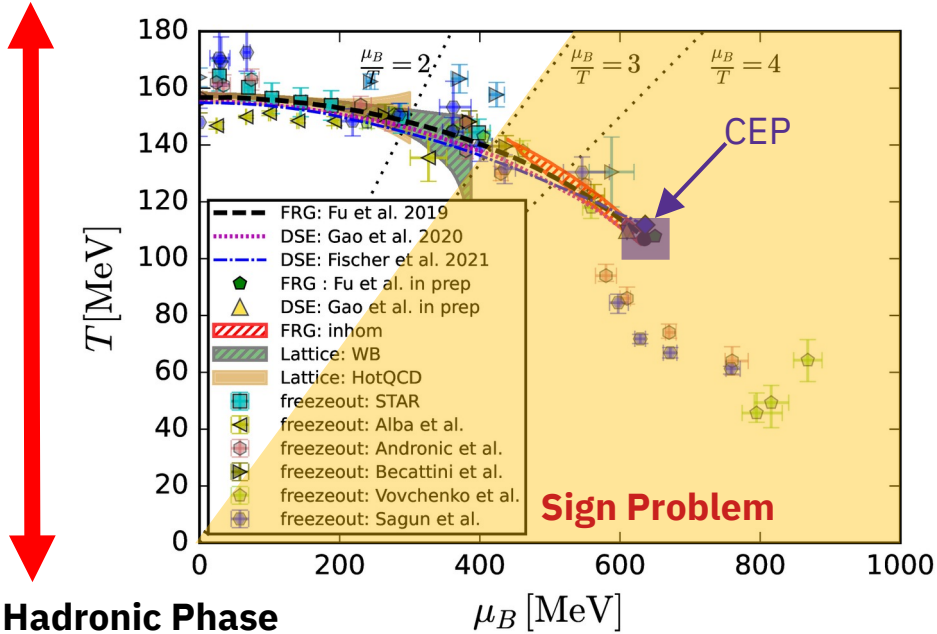
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The QCD phase diagram

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Direct access to phase structure using the

functional Renormalization Group

Non-perturbative, with **direct access to finite μ** .

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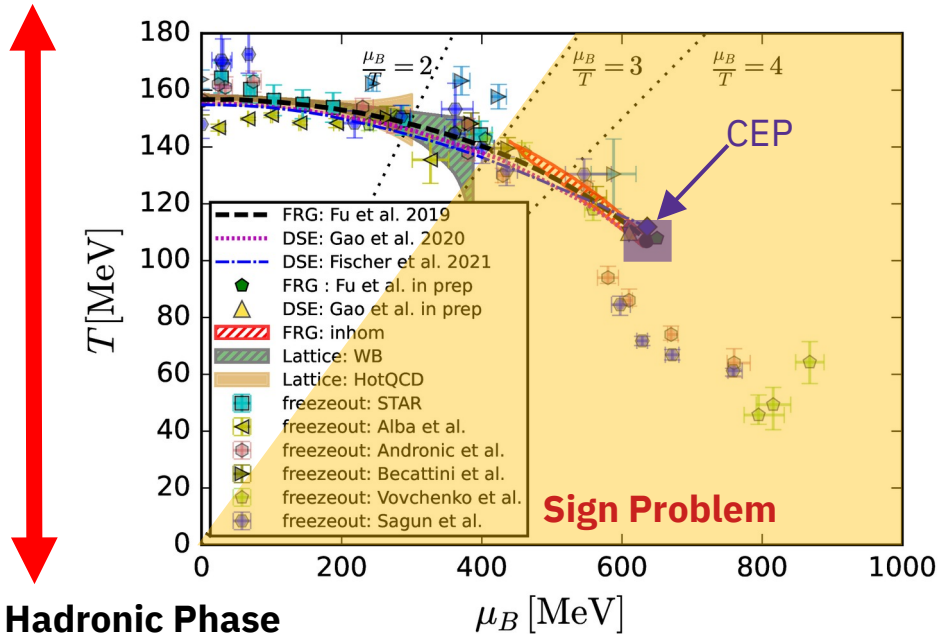
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The QCD phase diagram

Quark-Gluon Plasma



Direct access to phase structure using the

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Non-perturbative, with **direct access to finite μ** .

Here, first step:

- Setup
- Systematics

Vacuum

Ihssen, Pawłowski, Sattler, Wink
[arXiv:2408.08413]

Fu, Pawłowski, Rennecke [Phys. Rev. D 101 (2020), 054032]

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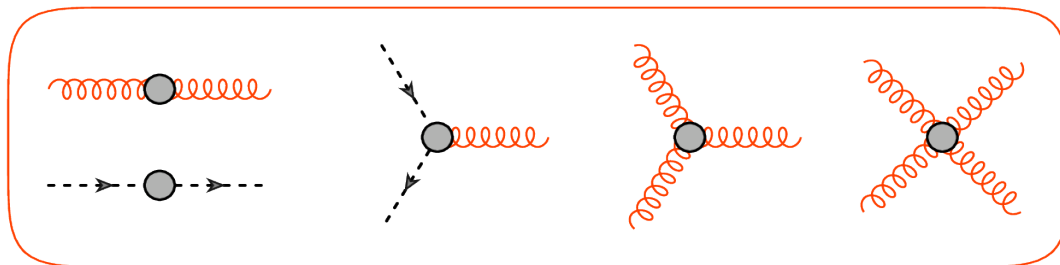
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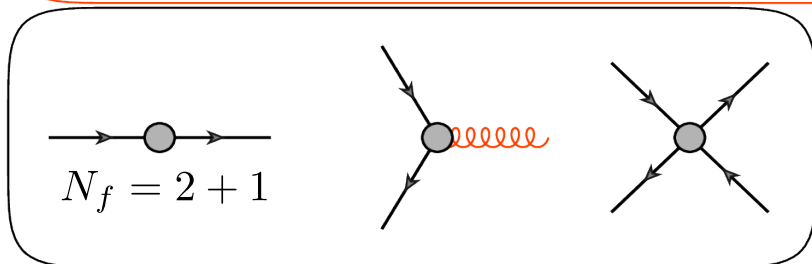
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Current vertex expansion

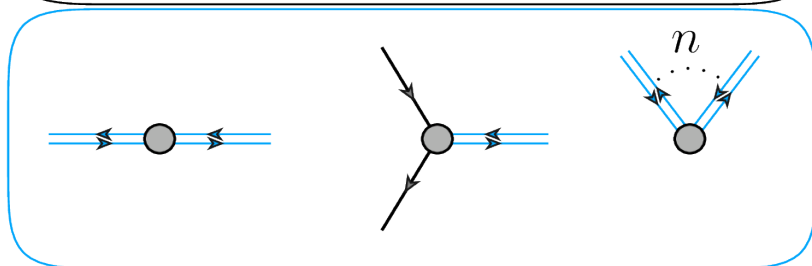
Glue sector



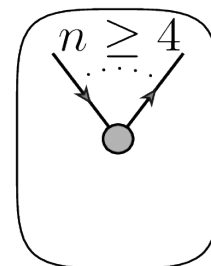
Quark-gluon sector



Quark-meson sector



— EoM →



Matter sector

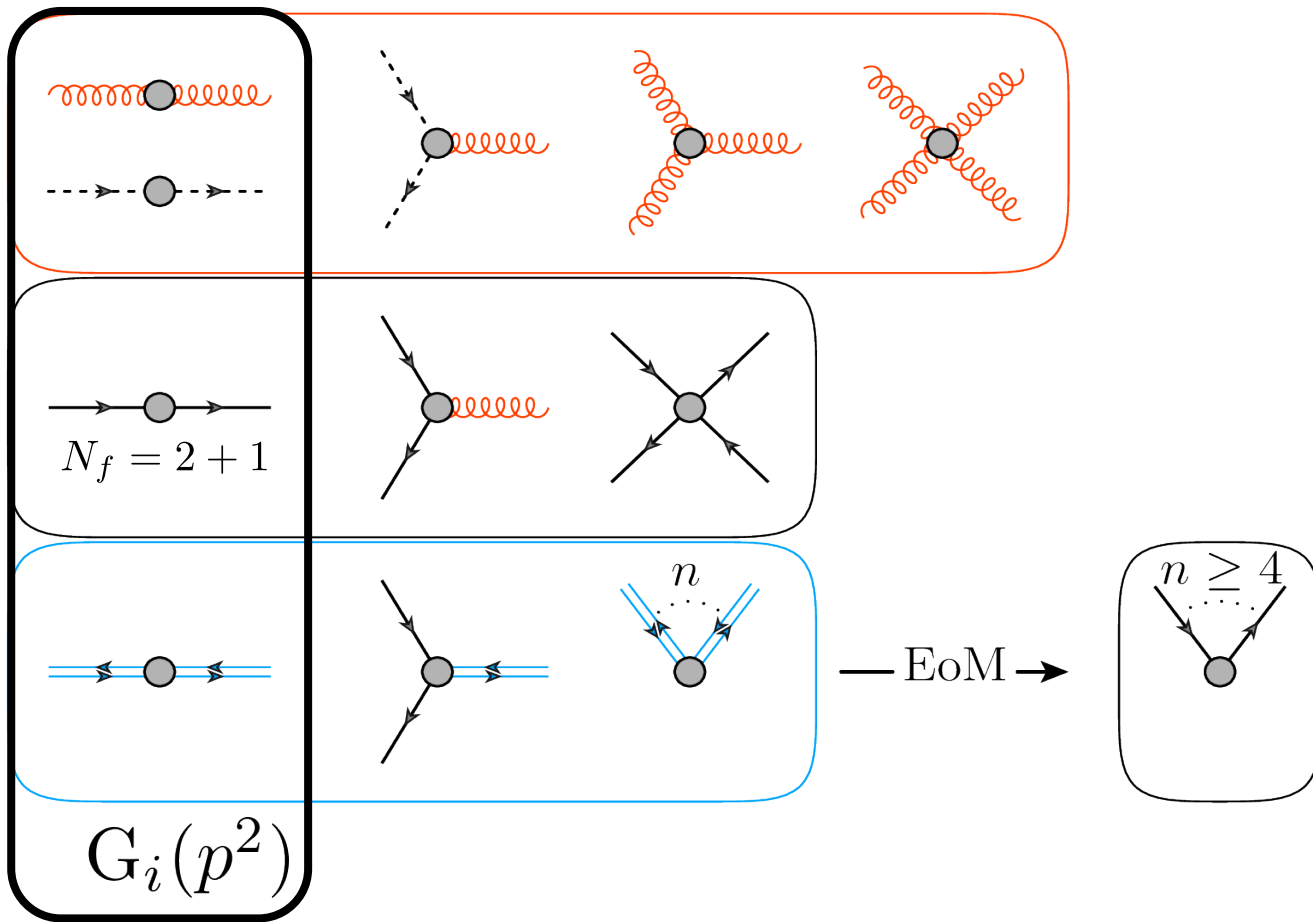
Current vertex expansion

Matter sector

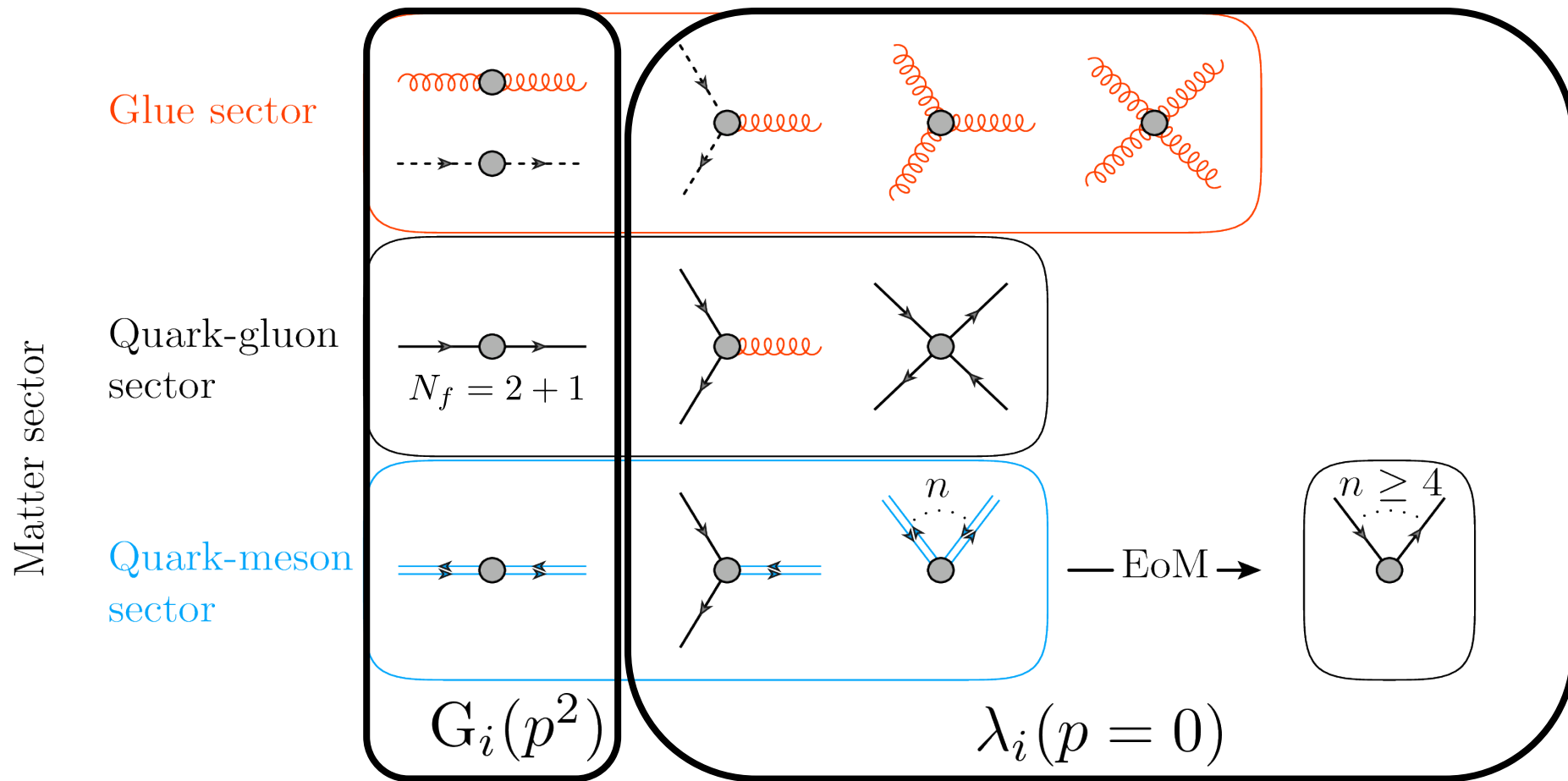
Glue sector

Quark-gluon sector

Quark-meson sector



Current vertex expansion



Matter sector

Glue sector

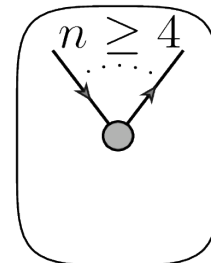
Quark-gluon sector

Quark-meson sector

$$G_i(p^2)$$

$$\lambda_i(p=0)$$

— EoM →



TensorBases Mathematica package

*With J. Braun,
J. Pawlowski, A. Geißel,
N. Wink*

- Automatically derived projectors

- Library of tensor bases, extendable by everyone

```
Needs["TensorBases`"]
```

```
In[2]:= TBGetProjector["transAqbq", 1, {p1, mu, a}, {p2, d2, A2, F2}, {p3, d3, A3, F3}]
```

```
TBGetInnerProduct["transAqbq"][TBGetProjector, 1, TBGetBasisElement, 1] // FormTrace // Simplify
```

```
TBGetInnerProduct["transAqbq"][TBGetProjector, 1, TBGetBasisElement, 8] // FormTrace // Simplify
```

```
Out[2]= 
$$\frac{1}{6 N_f - 6 N_c^2 N_f} i \text{deltaFundFlav}[F2, F3] \times \text{gamma}[\text{nu}\$20834, d2, d3] \times \text{TCol}[a, A2, A3] \times \text{transProj}[p1, mu, \text{nu}\$20834]$$

```

```
Out[3]= 1
```

```
Out[4]= 0
```


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- Automatic derivation and code generation for large fRG systems
- Hydrodynamic methods for full field dependences
- GPU accelerated

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```

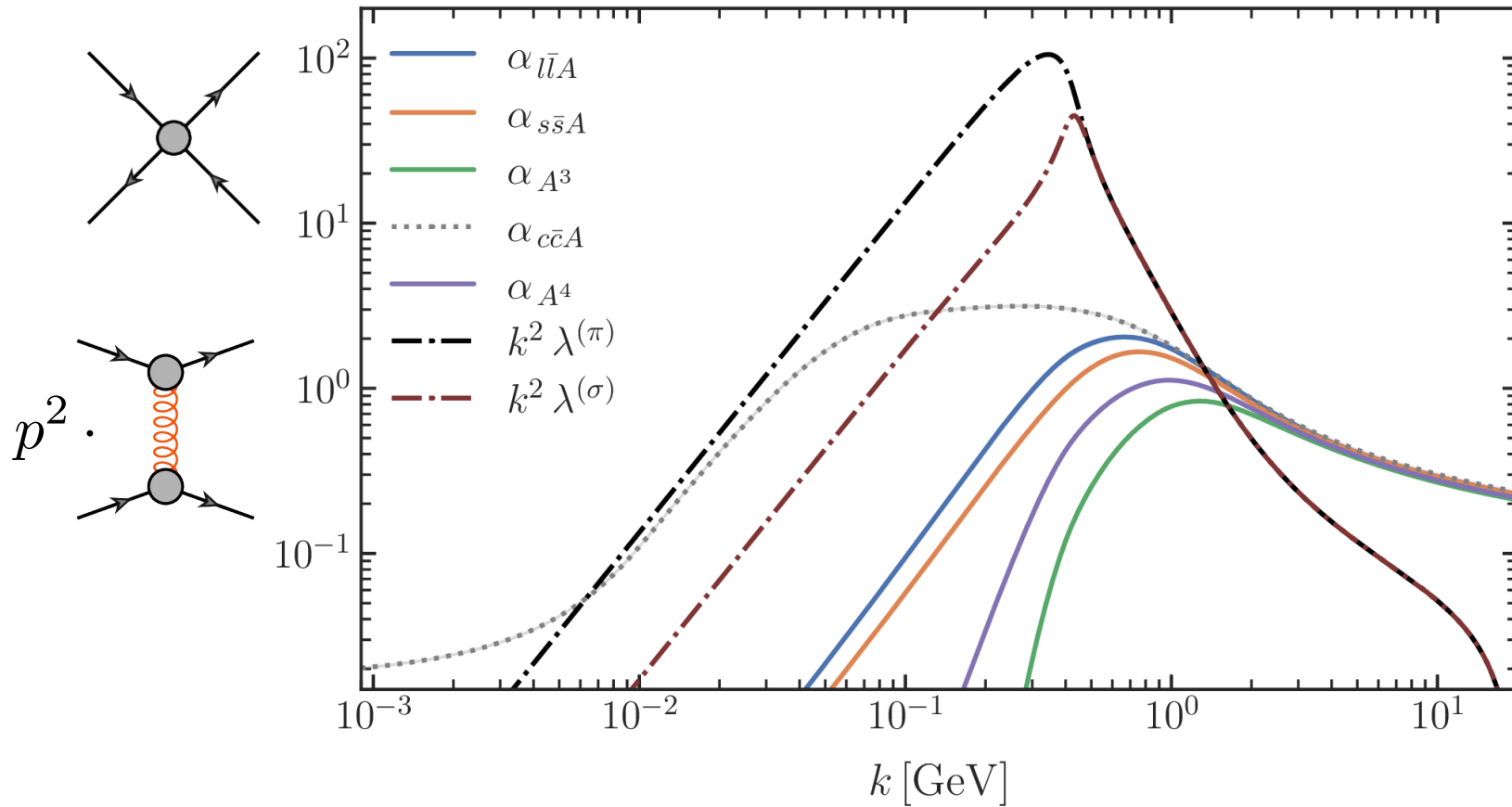
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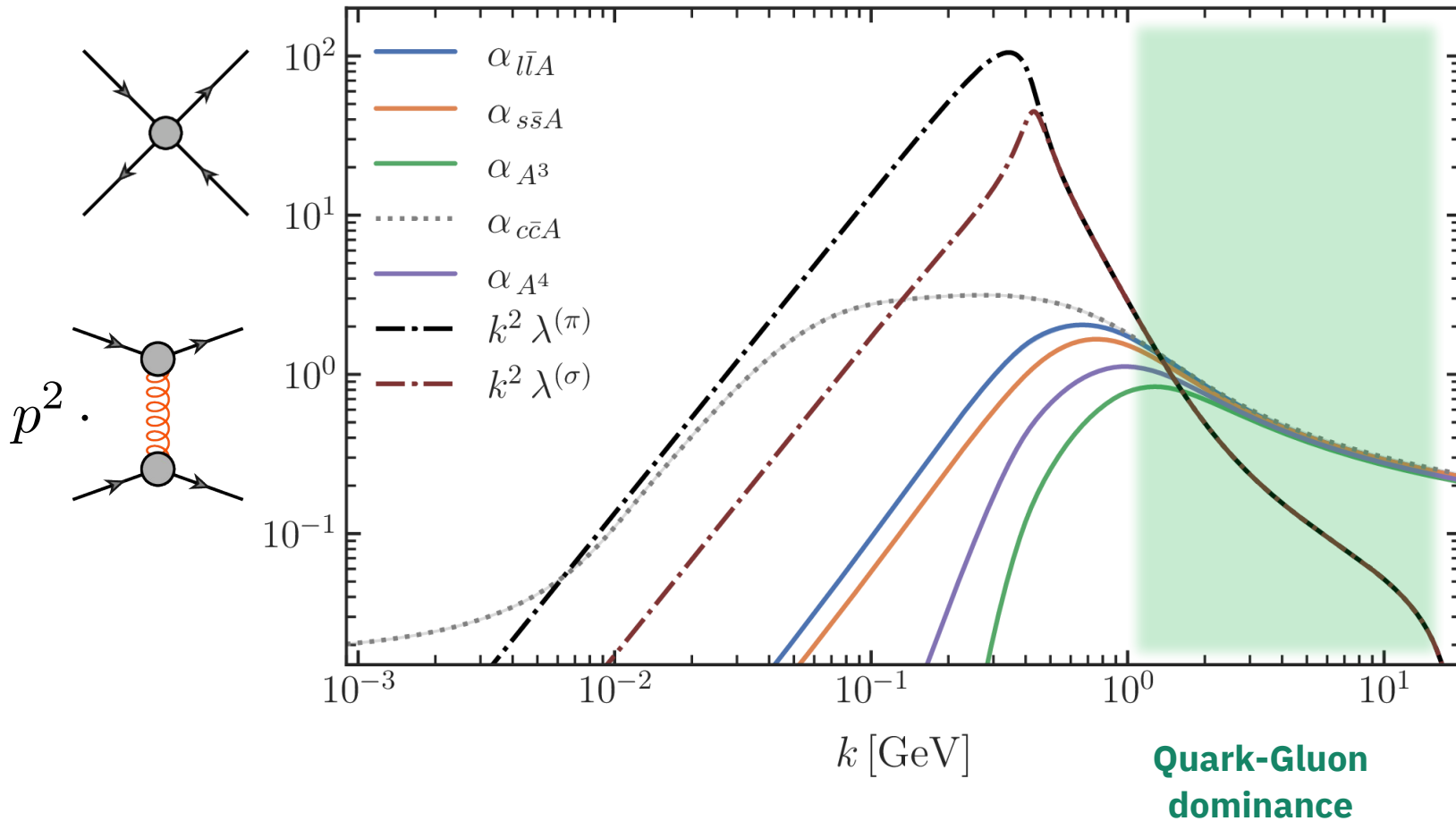
Open Source available around November

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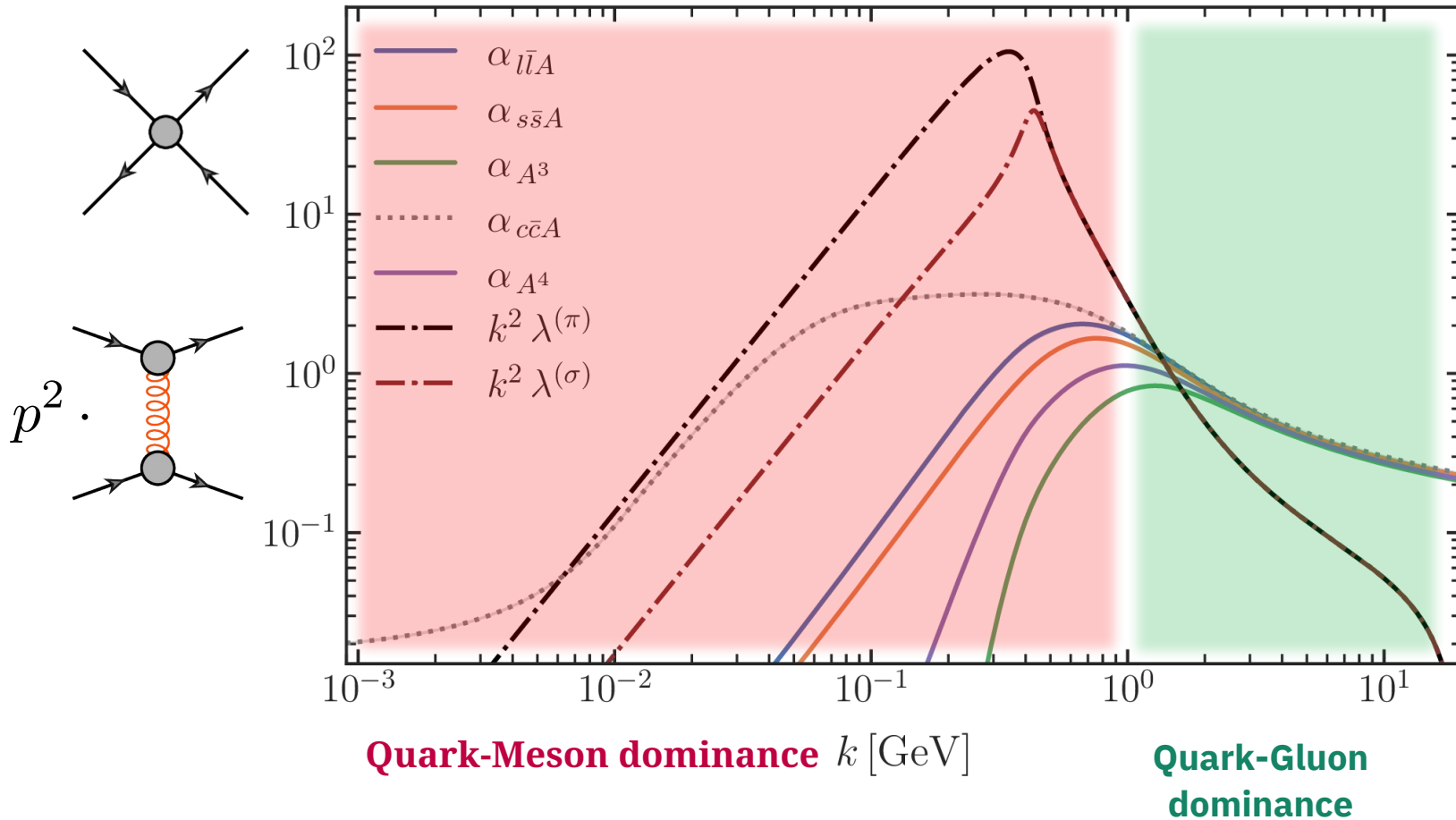
Dynamical Hadronisation in fRG



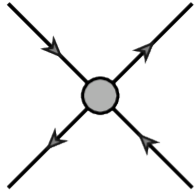
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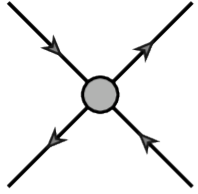
Dynamical Hadronisation in fRG



$\sigma - \pi$ - four-quark flow



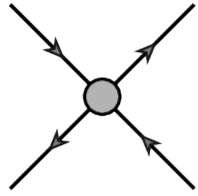
$\sigma - \pi$ - four-quark flow



(Resonant)

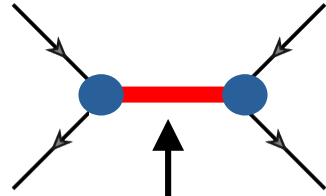
Fu, Huang, Pawłowski, Tan
[SciPost Phys. 14 (2023) 4, 069]
[arxiv:2401.07638 (2024)]

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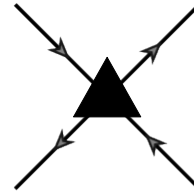
(Resonant)

=



Mesons

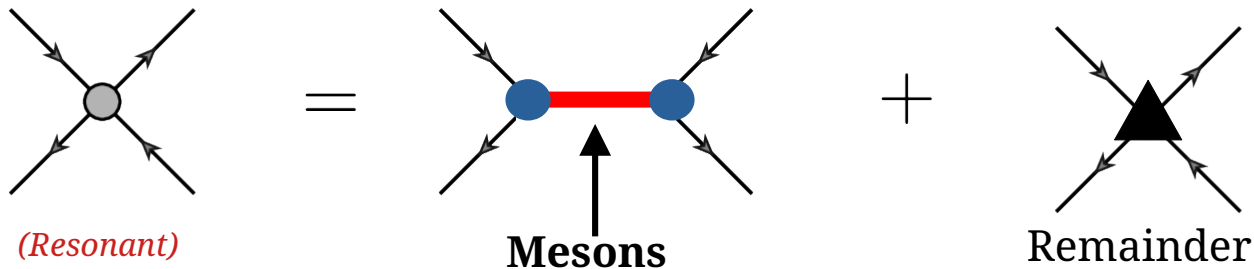
+



Remainder

Fu, Huang, Pawłowski, Tan
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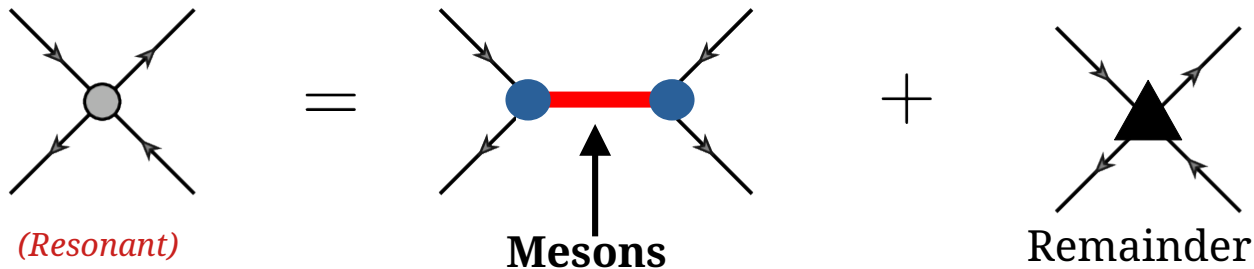
Mesons

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$$\begin{array}{c}
 \text{BS WF} \\
 \swarrow \quad \searrow \\
 h_\phi(p, q) \cdot \frac{1}{Z_\phi(t^2)(t^2 + m_\phi^2)} \cdot h_\phi(p, q) \\
 \text{Meson propagator}
 \end{array}$$

$\sigma - \pi$ - four-quark flow



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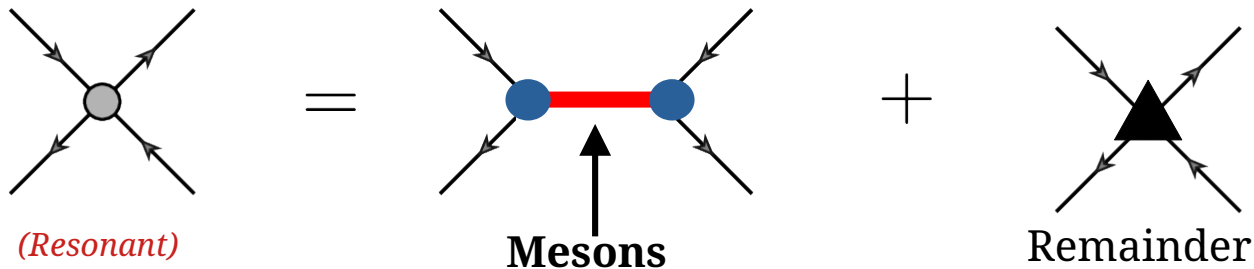
Meson propagator

Full scattering potential

$$V \left(\frac{\sigma^2 + \pi^2}{2} \right)$$

All orders of
n-meson
scatterings

$\sigma - \pi$ - four-quark flow



(Resonant)

Mesons

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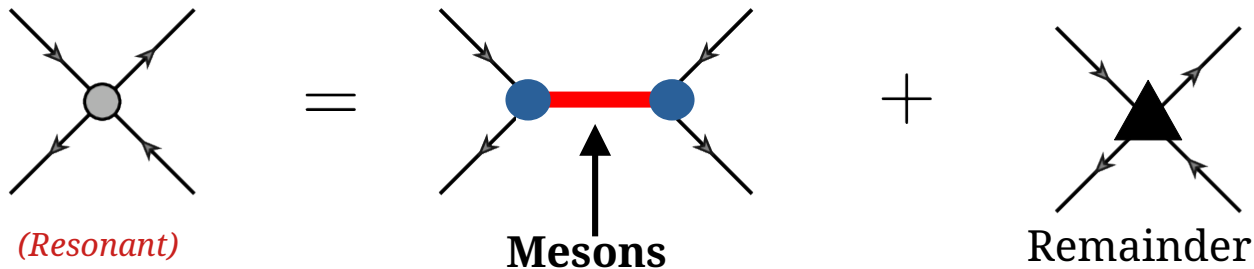
$$V \left(\frac{\sigma^2 + \pi^2}{2} \right)$$

**All orders of
n-meson
scatterings**

**Physical DoFs
emergent from
full QCD**

**Dynamical hadronization
/
Emergent Composites**

$\sigma - \pi$ - four-quark flow



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Physical DoFs
 emergent from
 full QCD

Dynamical hadronization
 /
 Emergent Composites

Full mesonic potential of QCD

Field space:

Finite element method

+ sensible RG-scale integration

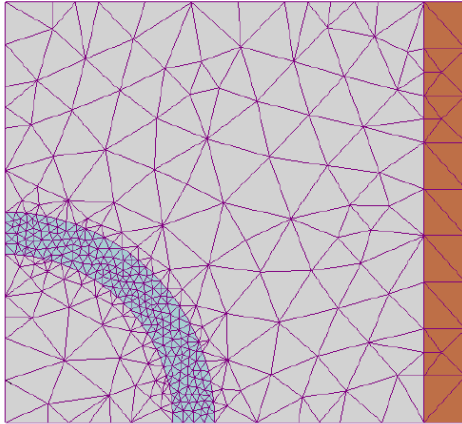
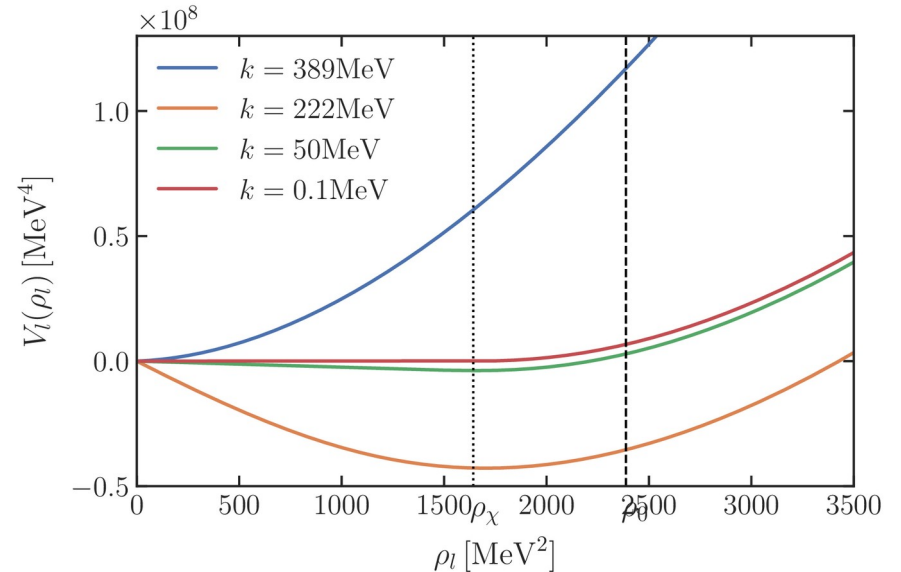


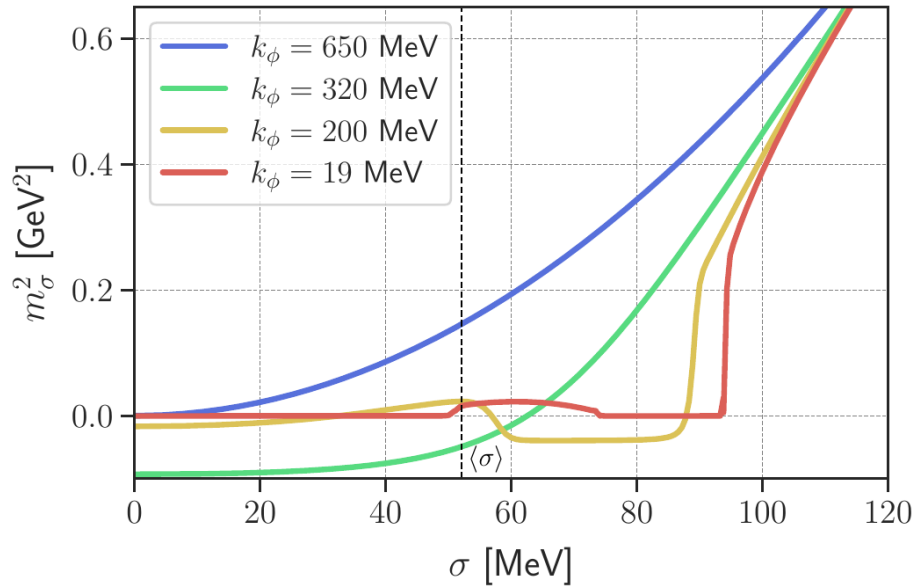
Image source:
wikipedia.org

Grossi, Wink [SciPost Phys.Core 6 (2023) 071]
Grossi, Ihssen, Pawłowski, Wink [Phys.Rev.D
104 (2021) 1, 016028]
Ihssen, Pawłowski, Sattler, Wink
[arXiv:2309.07335], [Comput.Phys.Commun.
300 (2024) 109182]

Ihssen, Sattler, Wink
(Phys.Rev.D 107 (2023) 11, 114009)



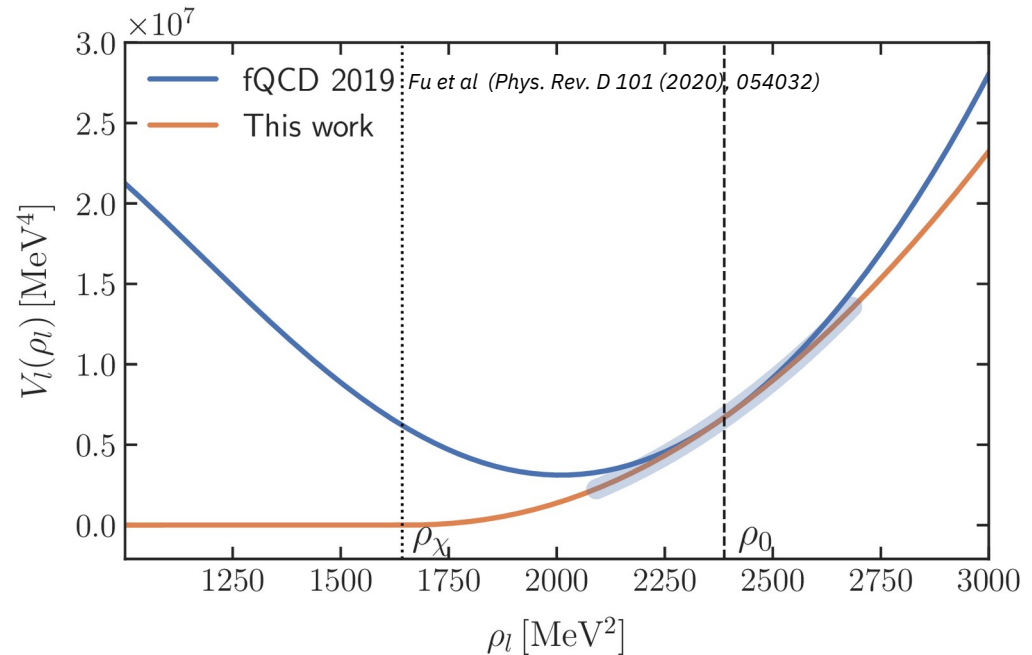
Full mesonic potential of QCD



Ihssen, Pawłowski, Sattler, Wink [arXiv:2309.07335]

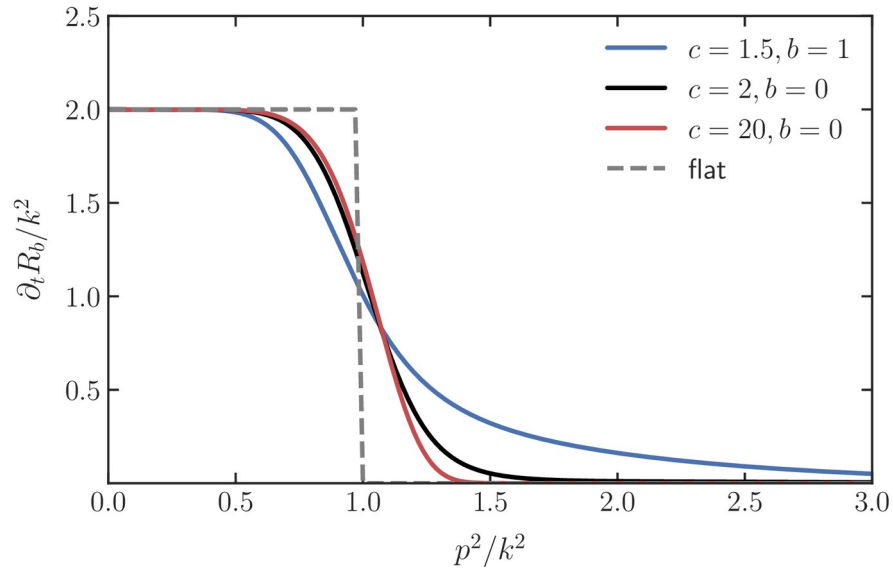
Hydro methods allow to access the full Potential.
Important for phase transitions at high μ

Quantitative access to chiral limit!



Systematic errors I: Regulator dependence

Easy regulator variation thanks to
numerical framework:

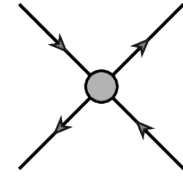


Observable	Value
$(f_K/f_\pi)_\chi$	$1.2168^{+0.0006}_{-0.0007}$
$f_{\pi,\chi}$ [MeV]	$93.2^{+3.5}_{-3.1}$
$m_{l,\chi}$ [MeV]	$311.6^{+0.3}_{-0.1}$
$m_{s,\chi}$ [MeV]	$446.7^{+0.3}_{-0.2}$
$m_{\sigma,\chi}$ [MeV]	$214.7^{+5.4}_{-9.3}$
$\sigma_{l,0,\chi}$ [MeV]	$67.1^{+1.2}_{-0.0}$

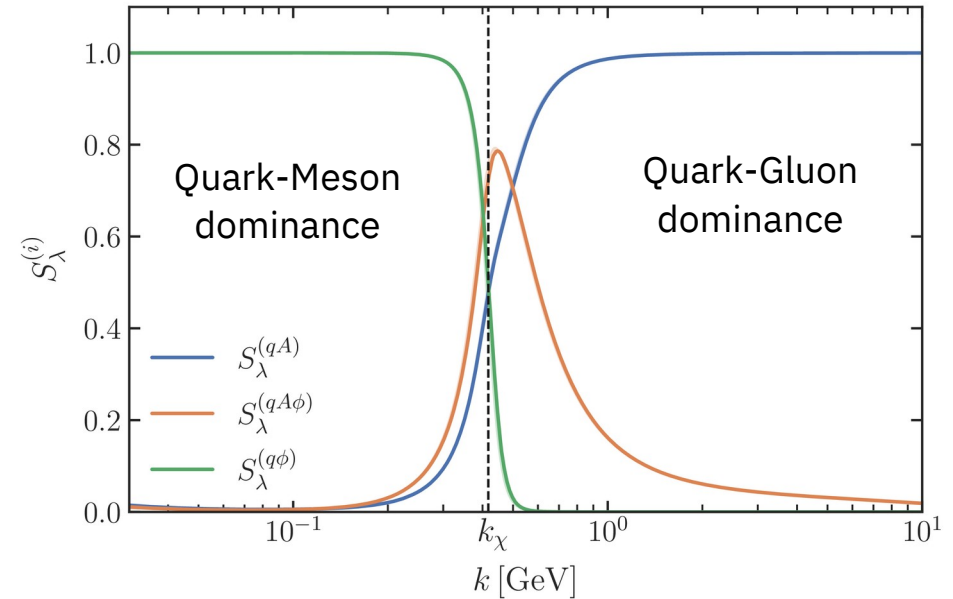
Chiral limit observables

Sattler et al. (in preparation)

Systematic errors II: The LEGO[®] principle

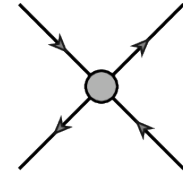


$(\sigma - \pi)$ -four-quark flow



$$S_{\lambda}^{(i)} = \frac{|\text{Flow}_{\lambda}^{(i)}|}{\sqrt{\sum_j (\text{Flow}_{\lambda}^{(j)})^2}}, \quad i, j = qA, q\phi, qA\phi$$

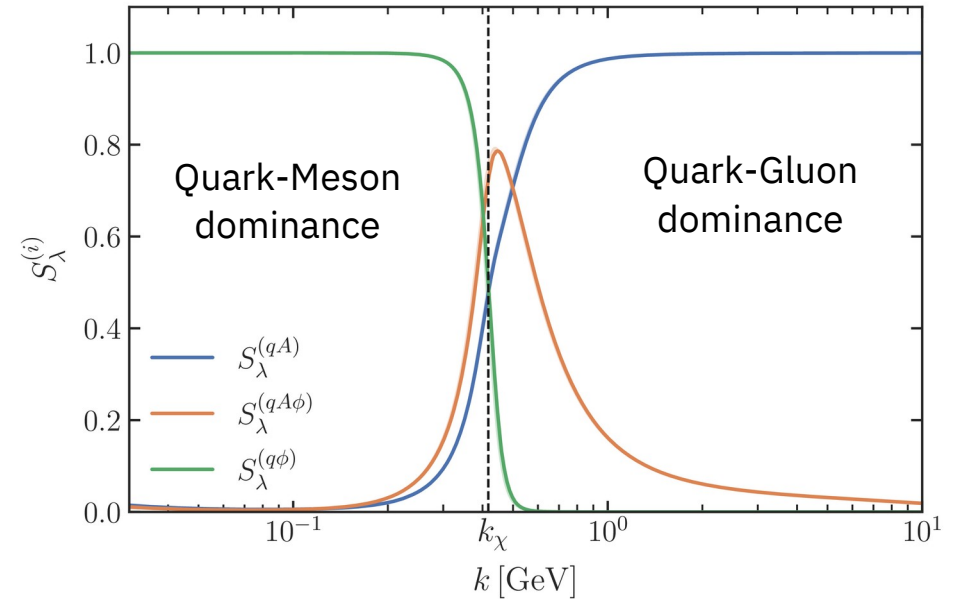
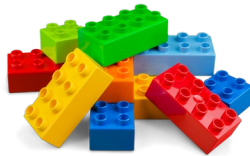
Systematic errors II: The LEGO[®] principle



($\sigma - \pi$)-four-quark flow

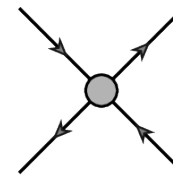
Separate LEGO[®] blocks:

- Glue subsystem $\{\lambda_{\text{glue}}\} = \{\alpha_{A^3}, \alpha_{A^4}, \alpha_{c\bar{c}A}\}$
- Matter subsystem $\{\lambda_{\text{mat}}\} = \{h_\phi(\rho_0), \lambda_{\phi,n}(\rho_0)\}$
- Interface blocks $\{\lambda_{\text{inter}}\} = \{\alpha_{l\bar{l}A}\}$



$$S_\lambda^{(i)} = \frac{|\text{Flow}_\lambda^{(i)}|}{\sqrt{\sum_j (\text{Flow}_\lambda^{(j)})^2}}, \quad i, j = qA, q\phi, qA\phi$$

Systematic errors II: The LEGO® principle

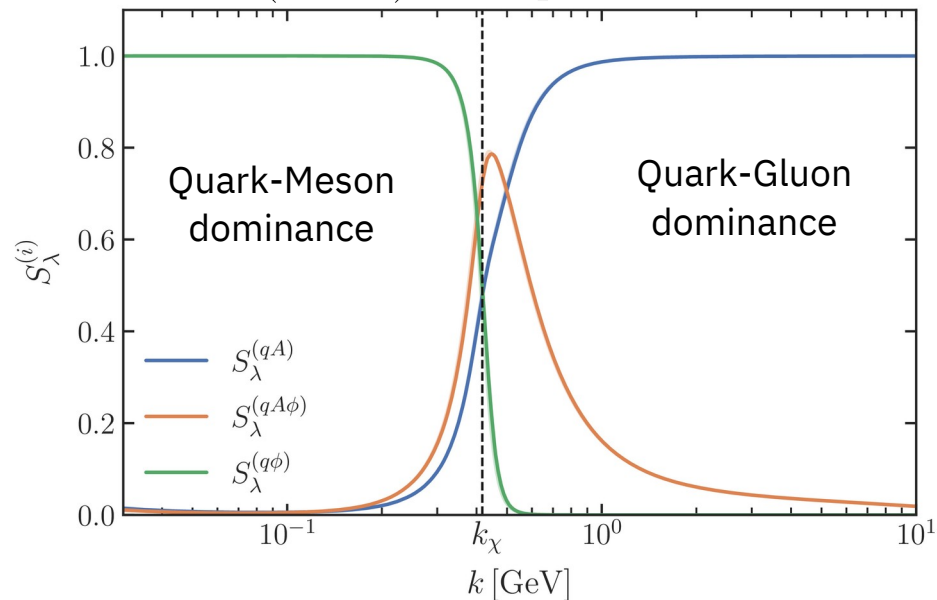


($\sigma - \pi$)-four-quark flow

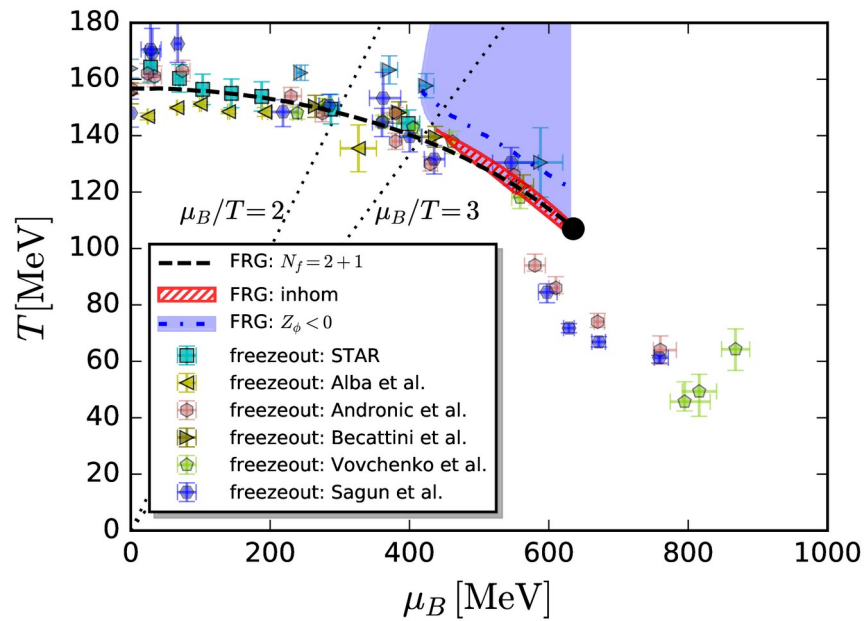
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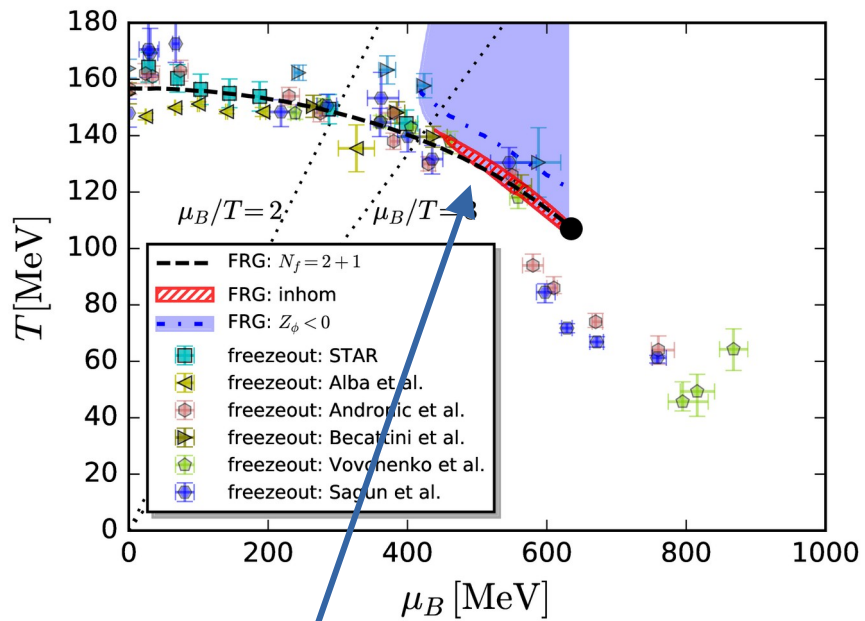
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- ➔ Systematic error estimates from subsystems; preliminary estimate 10%.
- ➔ Low-energy effective theories.



$$S_\lambda^{(i)} = \frac{|\text{Flow}_\lambda^{(i)}|}{\sqrt{\sum_j (\text{Flow}_\lambda^{(j)})^2}}, \quad i, j = qA, q\phi, qA\phi$$





Moat regime at high μ
 (possible inhomogeneous phase)

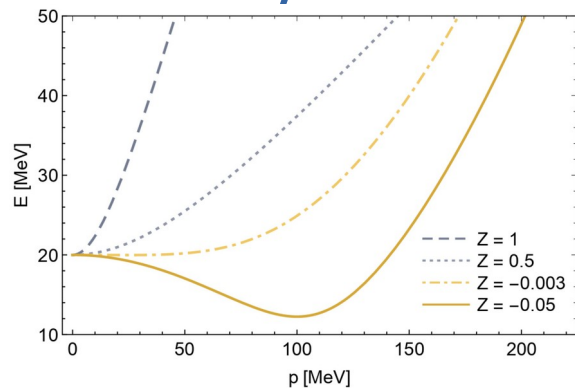
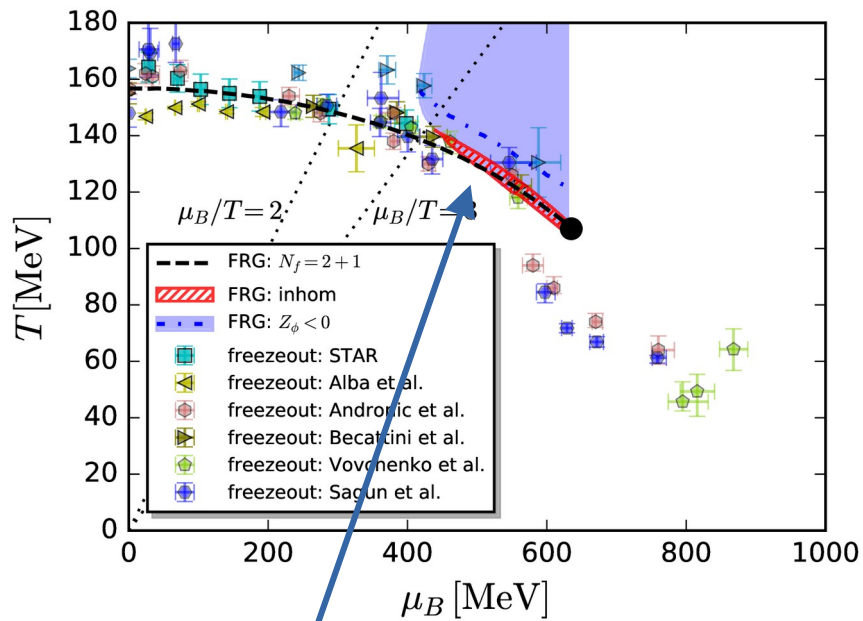


Figure by Rennecke, Pisarski [PoS CPOD2021 (2022) 016]

Moat regime at high μ
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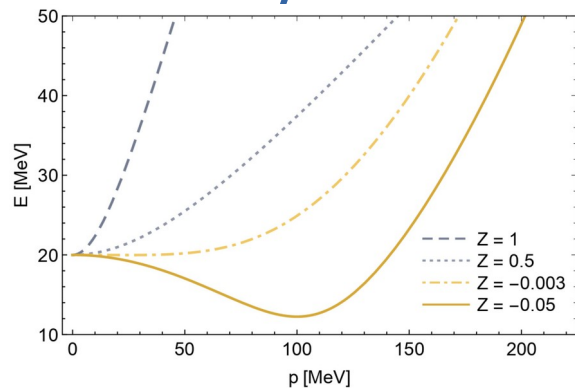
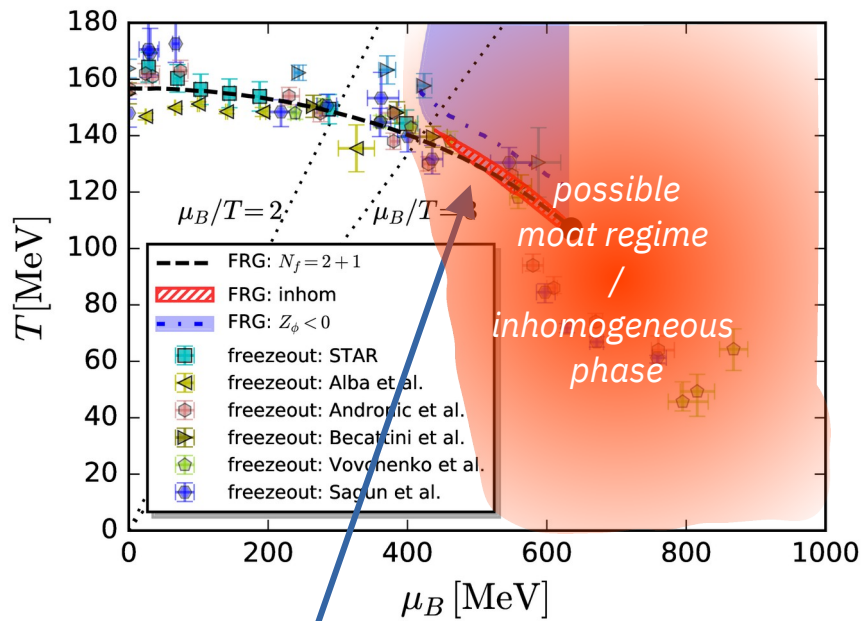
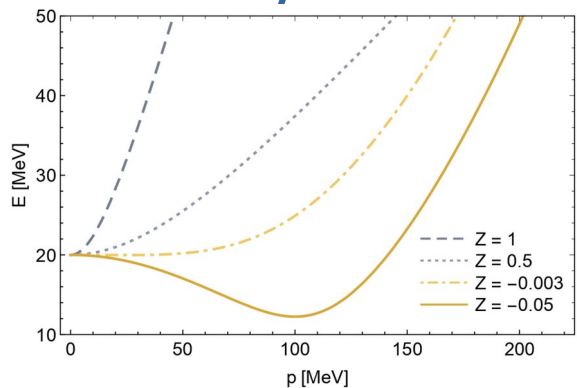
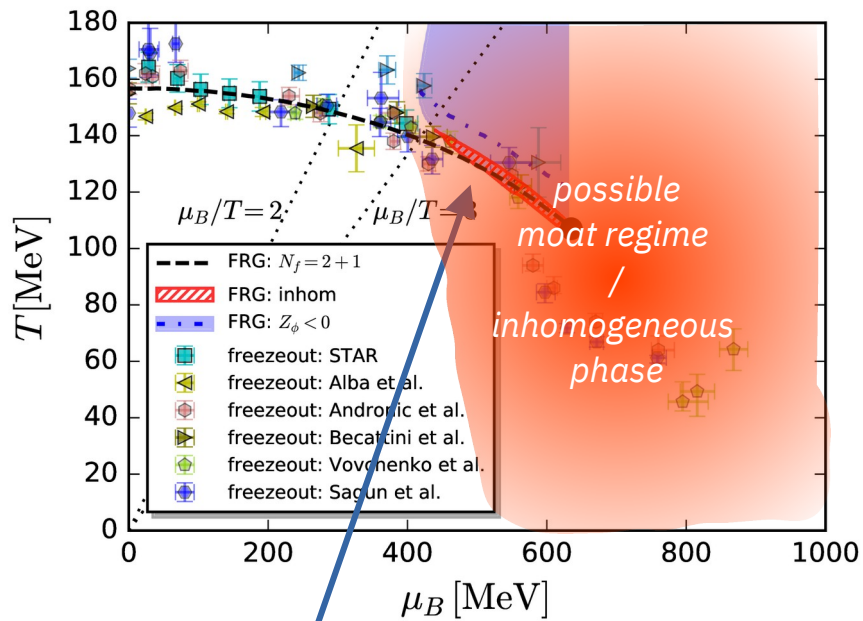


Figure by Rennecke, Pisarski [PoS CPOD2021 (2022) 016]

Moat regime at high μ
(possible inhomogeneous phase)

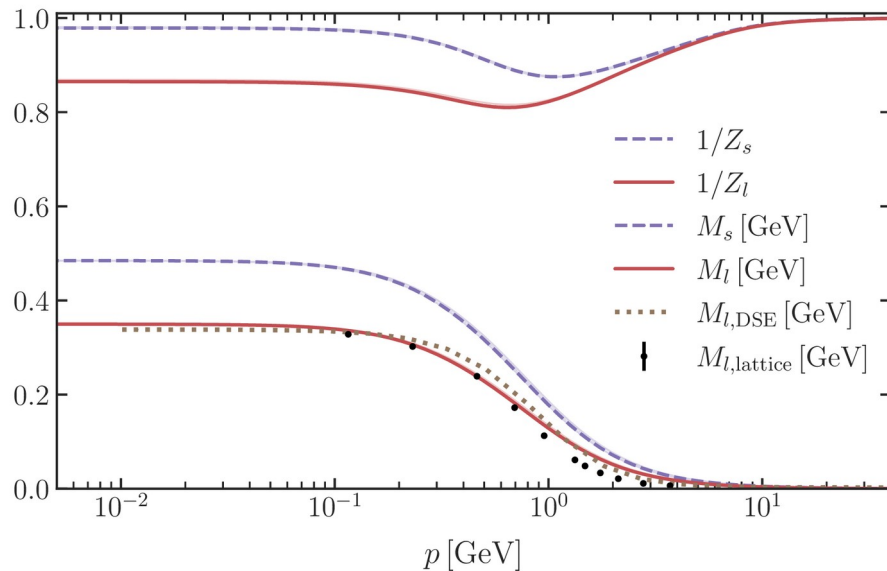


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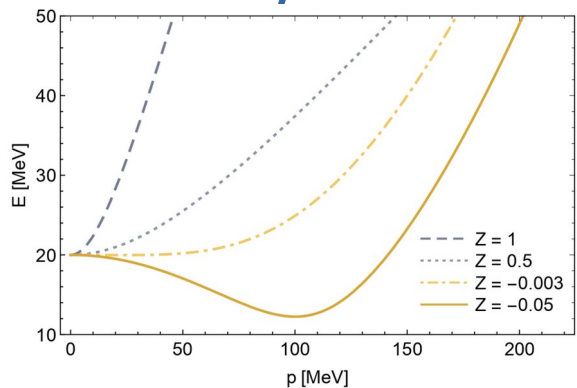
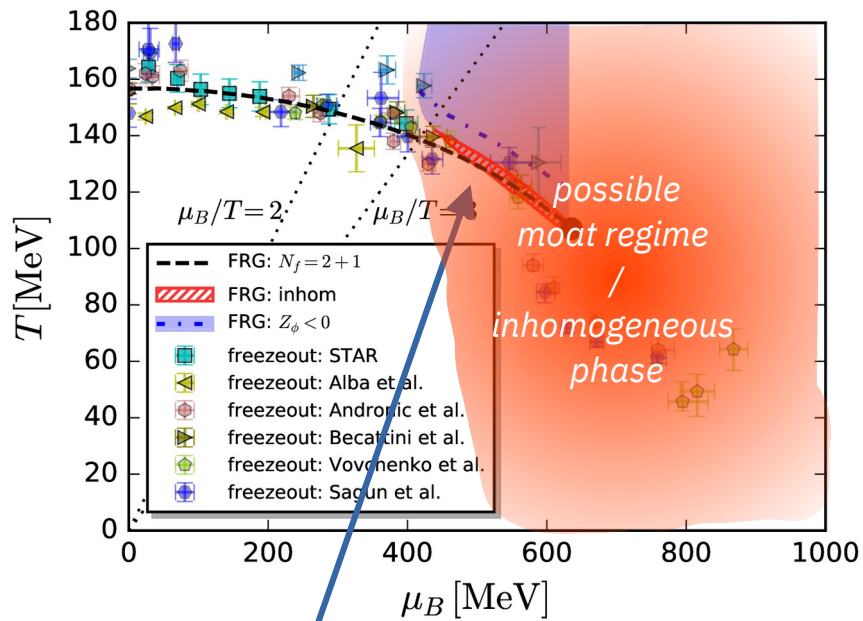
Figure by Rennecke, Pisarski
[PoS CPOD2021 (2022) 016]



Full momentum resolution of propagators



Lattice data: Cheng et al
[Phys. Rev. D 104 (2021), 094509]
DSE data: Gao, Papavassiliou, Pawłowski
[Phys. Rev. D 103 (2021), 094013]

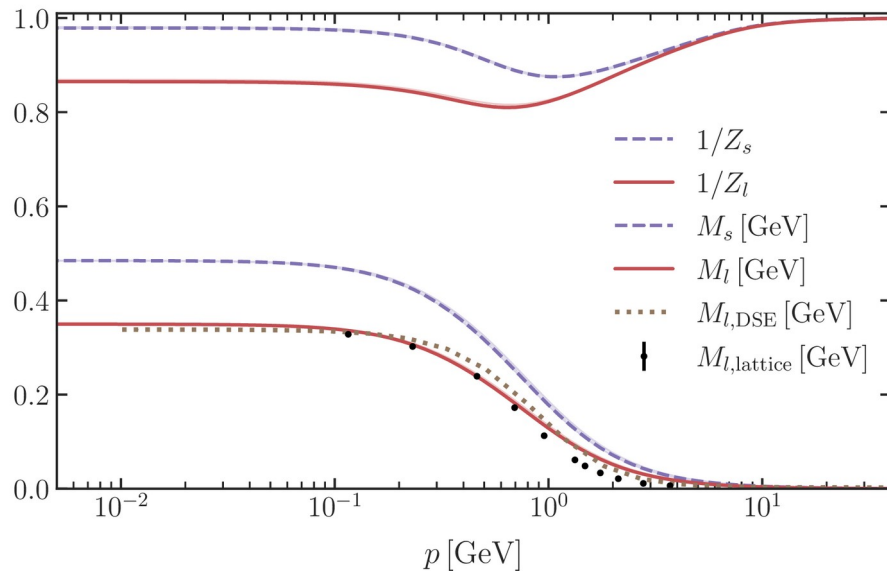


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Access to pole masses:

$$m_{\pi, \text{vacuum}}^{(N_f=2)} = 139(12) \text{ MeV}$$

$$m_{\pi, \text{vacuum}}^{(N_f=2+1)} = 138(9) \text{ MeV}$$

Conclusions

- **Motivation:**
Direct access to phase structure of QCD through fRG
- **Quantitative Vacuum** results in agreement with Lattice & other functional approaches
- **Systematic error estimates**
- **Easily extendable setup**

Outlook

- Results at **finite (T, μ)** (*in progress*)
- More **momentum dependences** (*done in vacuum*)
- Rebosonisation of **further channels** (*in progress*)
- Increase **number of tensor structures** (*done in vacuum*)

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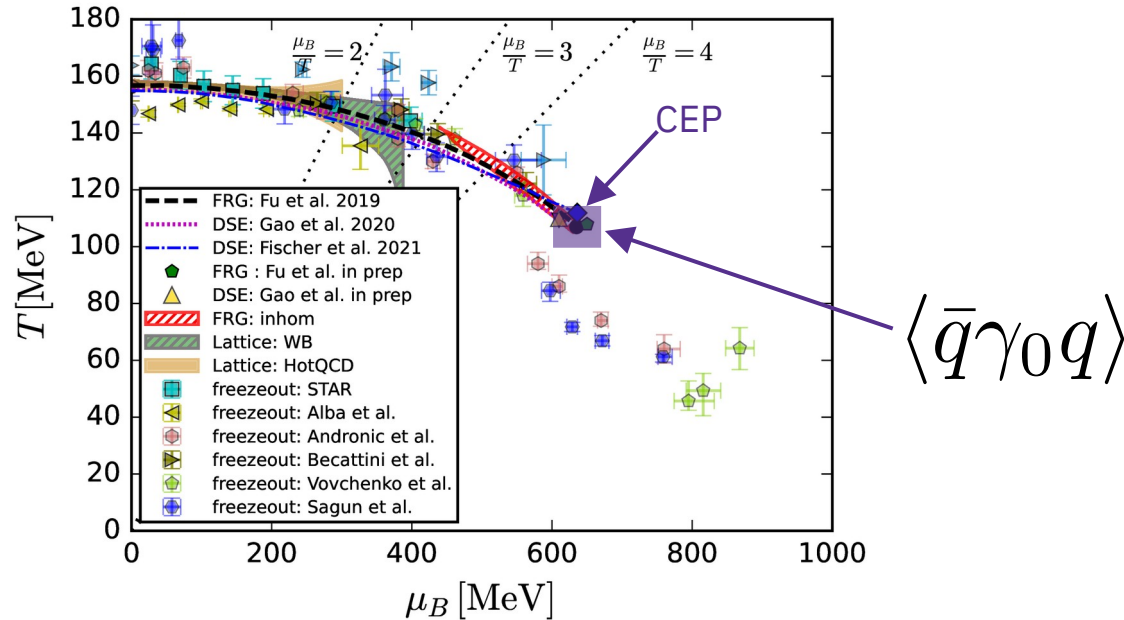
fQCD Collaboration:

Braun
Chen
Fu
Gao
Geissel
Huang
Ihssen
Lu
Pawlowski
Rennecke
Sattler
Schallmo
Stoll
Tan
Töpfel
Turnwald
Wen
Wessely
Wink
Yin
Zorbach

**Thank you
for your attention!**

Backup slides

Mapping out the phase diagram



Fu, Pawłowski, Rennecke (Phys. Rev. D 101 (2020), 054032)

Gao, Pawłowski (Phys.Lett.B820(2021) 136584)

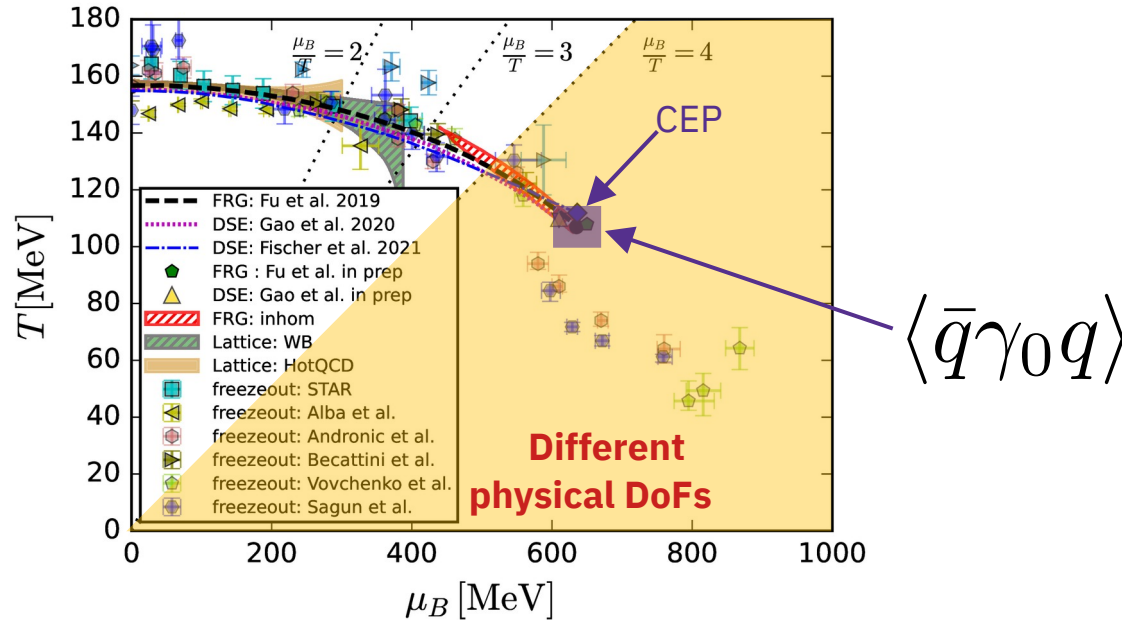
Gunkel, Fischer (Phys.Rev.D 104 (2021) 5, 054022)

Bellwied et al. (WB) (Phys.Lett.B 751 (2015) 559-564)

Bazavov et al. (HotQCD) (Phys.Lett.B 795 (2019) 15-21)

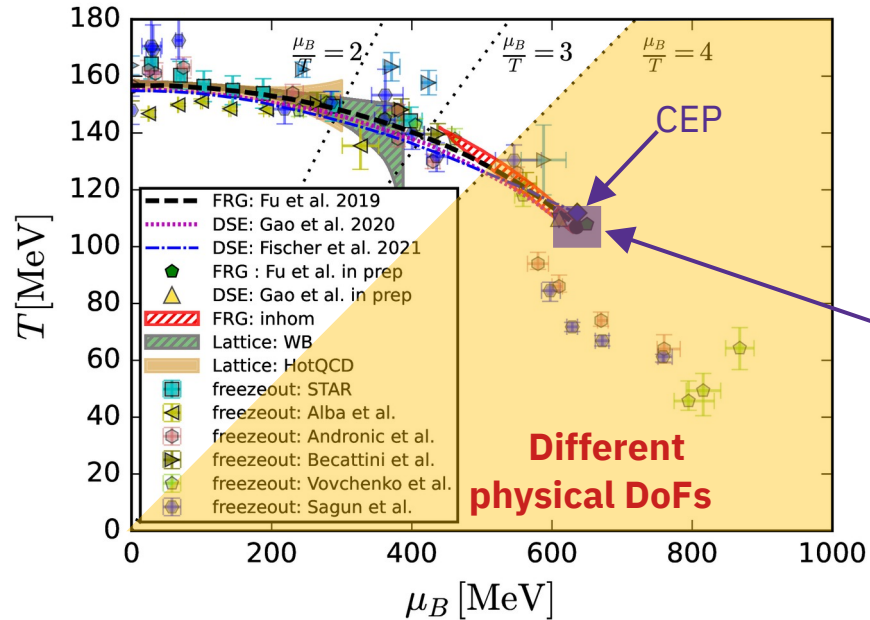
Sattler et al. (in preperation)

Mapping out the phase diagram

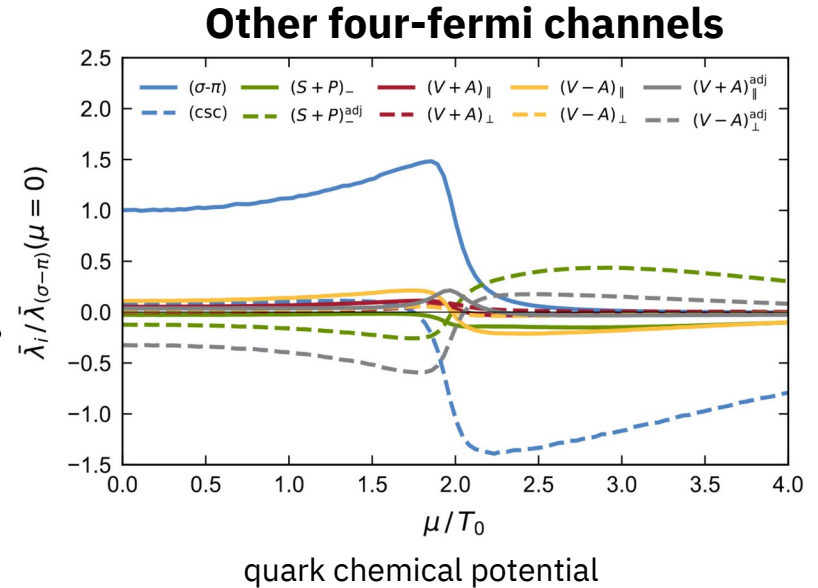


Fu, Pawłowski, Rennecke (Phys. Rev. D 101 (2020), 054032)
 Gao, Pawłowski (Phys.Lett.B820(2021) 136584)
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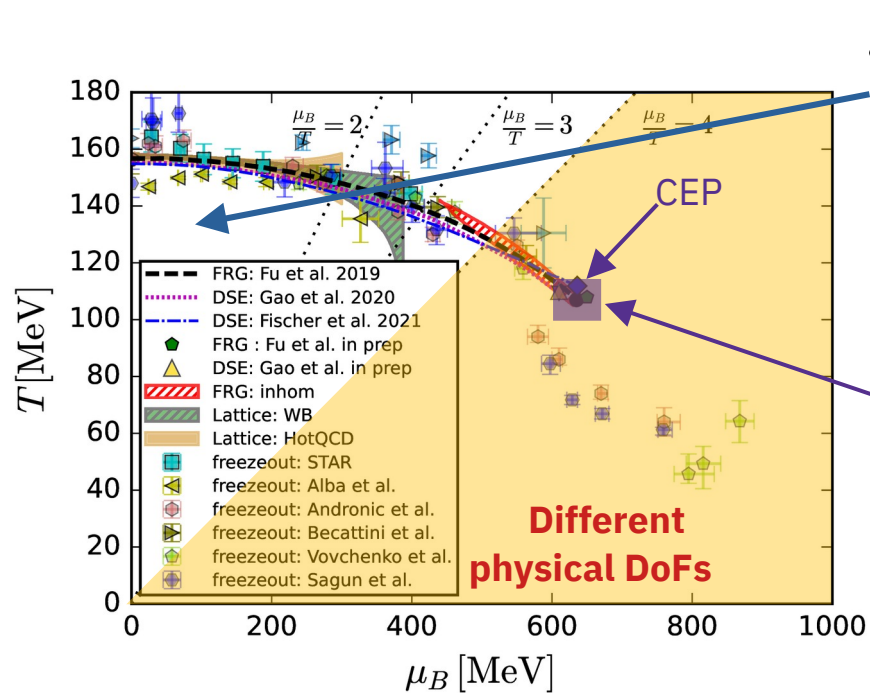
$\langle \bar{q} \gamma_0 q \rangle$



Fu, Pawłowski, Rennecke (Phys. Rev. D 101 (2020), 054032)
 Gao, Pawłowski (Phys.Lett.B820(2021) 136584)
 Gunkel, Fischer (Phys.Rev.D 104 (2021) 5, 054022)
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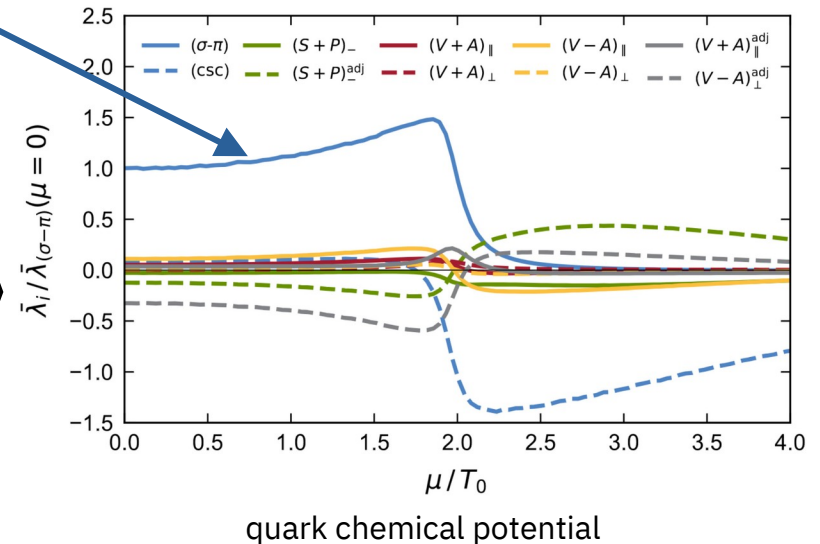
Braun, Leonhardt, Pospiech (Phys.Rev.D 101 (2020) 3, 036004)

Mapping out the phase diagram



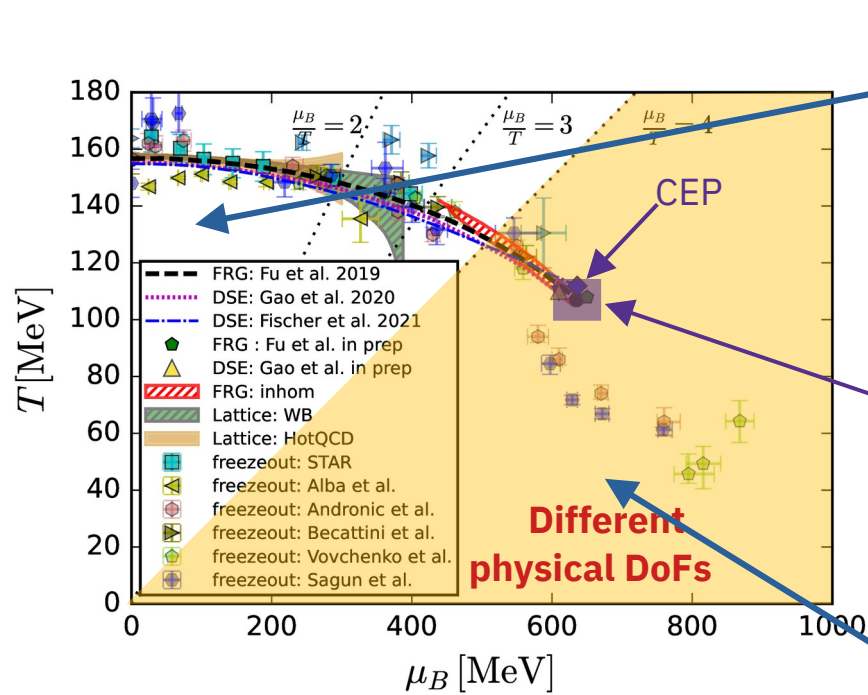
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 Bazavov et al. (HotQCD) (Phys.Lett.B 795 (2019) 15-21)
 Sattler et al. (in preperation)

Other four-fermi channels



Braun, Leonhardt, Pospiech (Phys.Rev.D 101 (2020) 3, 036004)

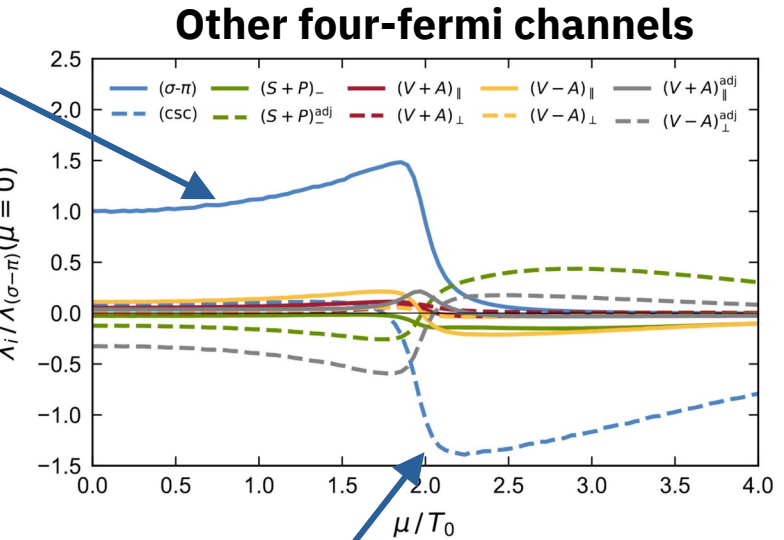
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Fu, Pawłowski, Rennecke (Phys. Rev. D 101 (2020), 054032)
Gao, Pawłowski (Phys.Lett.B820(2021) 136584)
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Sattler et al. (in preperation)

$\langle \bar{q}q \rangle$

$\langle \bar{q}\gamma_0 q \rangle$

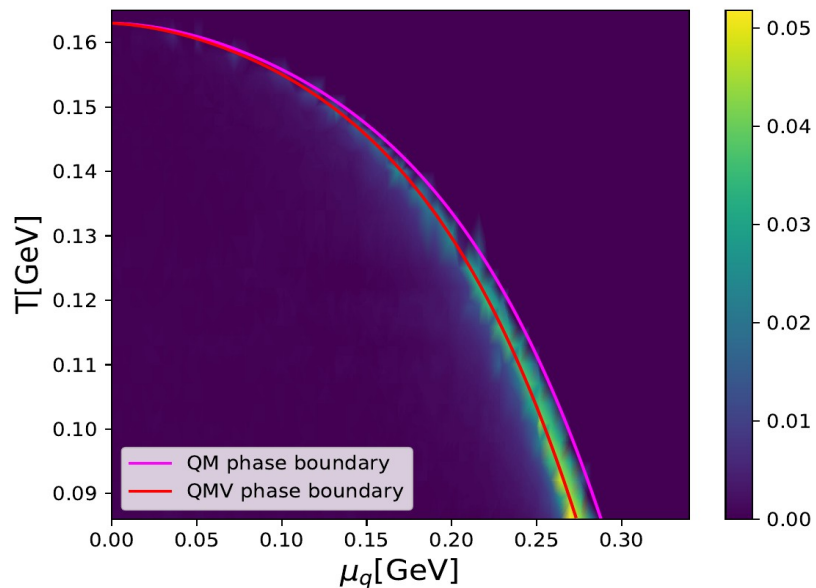


Braun, Leonhardt, Pospiech (Phys.Rev.D 101 (2020) 3, 036004)

$\langle d^* d \rangle$

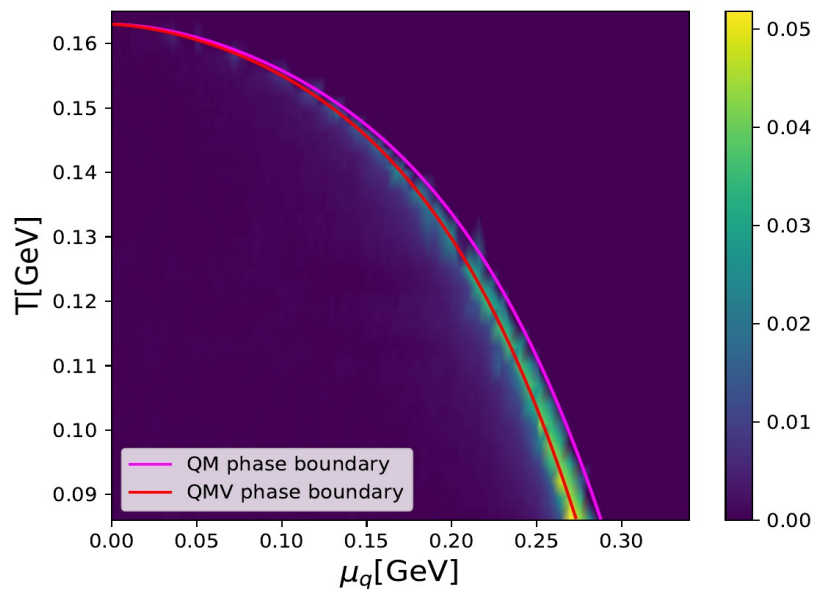
Inclusion of Density mode $\langle \bar{q}\gamma_0 q \rangle$ and Diquarks

*Ihssen, Hendricks, Pawłowski, Sattler
(in preparation)*



Inclusion of Density mode $\langle \bar{q}\gamma_0 q \rangle$ and Diquarks

Ihssen, Hendricks, Pawłowski, Sattler
(in preparation)



Full $N_f = 2+1$

Pawłowski, Sattler, Steck
(in preparation)

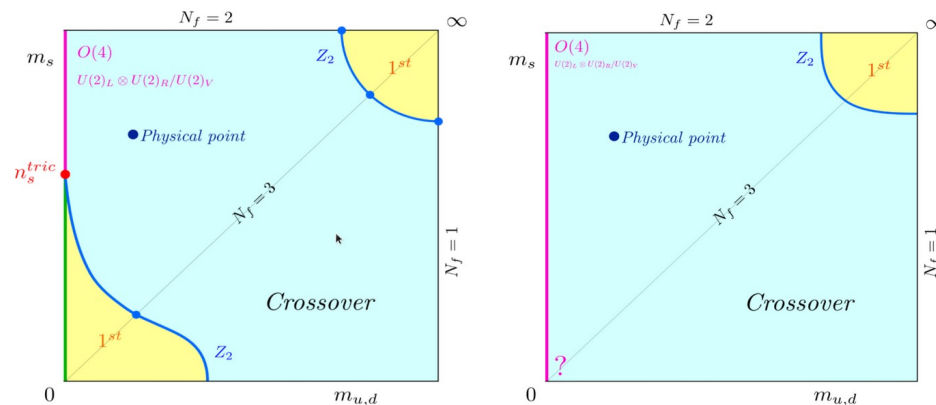
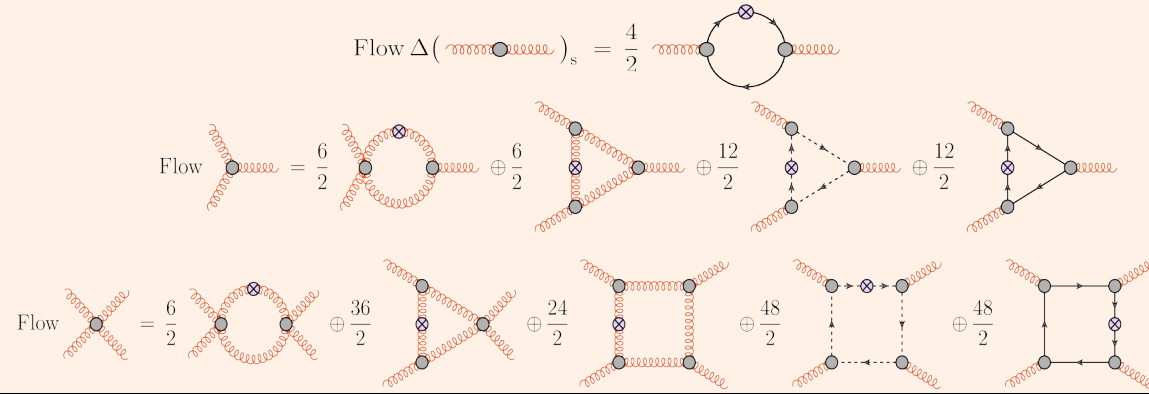


Figure from
Owe Philipsen (arXiv:2111.03590)

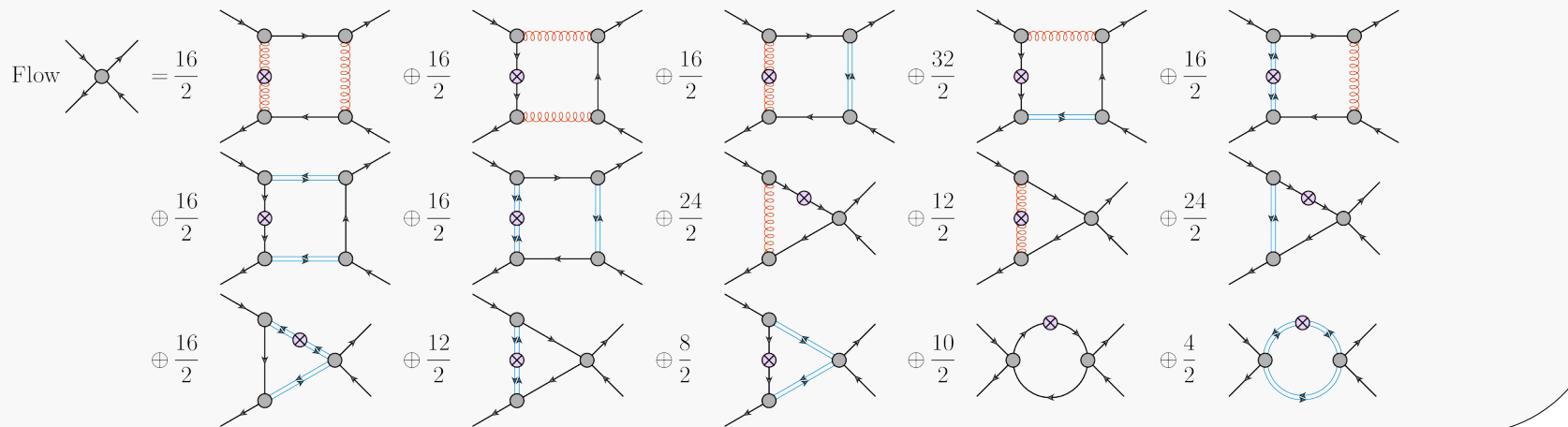
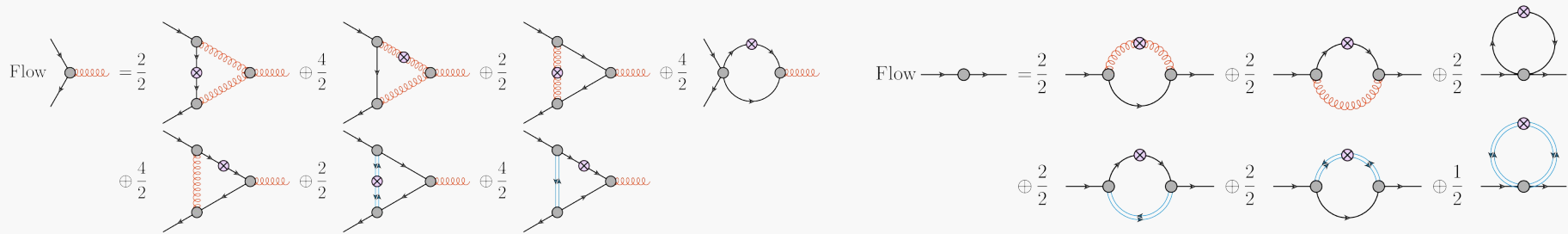
Gluons



Glueons

$$\text{Flow } \Delta(\text{wavy line})_s = \frac{4}{2} \text{ (loop diagram)}$$

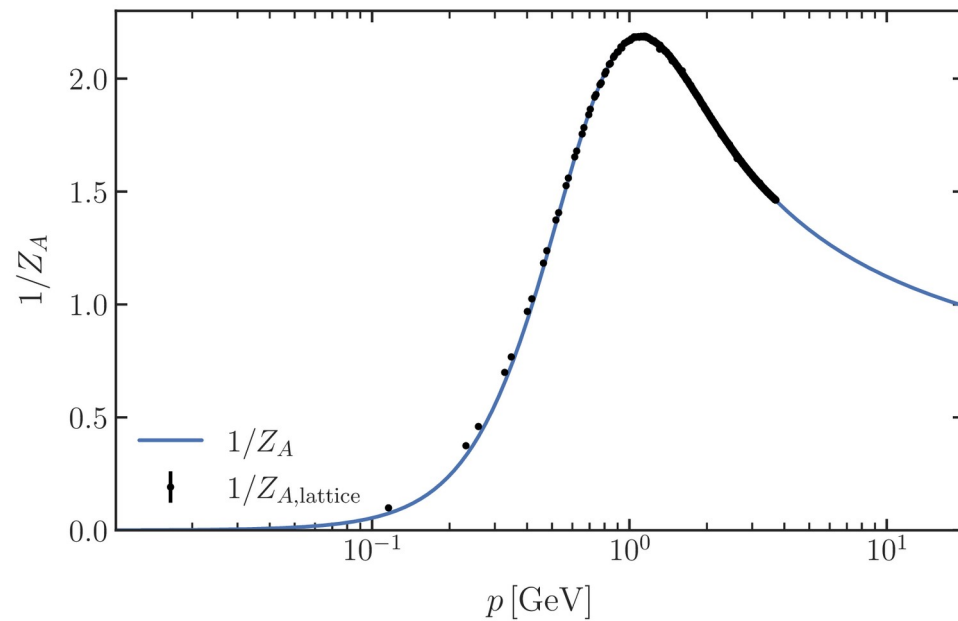
$$\text{Flow } \text{wavy line} = \frac{6}{2} \text{ (loop diagram)} \oplus \frac{6}{2} \text{ (triangle diagram)} \oplus \frac{12}{2} \text{ (triangle diagram)} \oplus \frac{12}{2} \text{ (triangle diagram)}$$



Quarks

Observables	Value	Parameter in $\Gamma_{\Lambda_{UV}}$
$m_{\pi, \text{pol}}$ [MeV]	138(9)	$c_{\sigma_l} = 4.67 \text{ GeV}^3$
f_K/f_π	1.1914	$\Delta m_{sl} = 134.2 \text{ MeV}$
$\alpha_{l\bar{l}A, \Lambda_{UV}}$		$\alpha_{l\bar{l}A, \Lambda_{UV}} = 0.227$
m_l [MeV]	350	$a = 0.0251 \quad b = 2 \text{ GeV}$
f_π [MeV]	$97.2^{+4.0}_{-2.2}$	_____
m_s [MeV]	$485.0^{+0.0}_{-0.3}$	_____
$m_{\pi, \text{cur}}$ [MeV]	138	_____
m_σ [MeV]	$388.1^{+0.0}_{-1.1}$	_____
$\sigma_{0,l}$ [MeV]	$69.^{+1.2}_{-0.2}$	_____

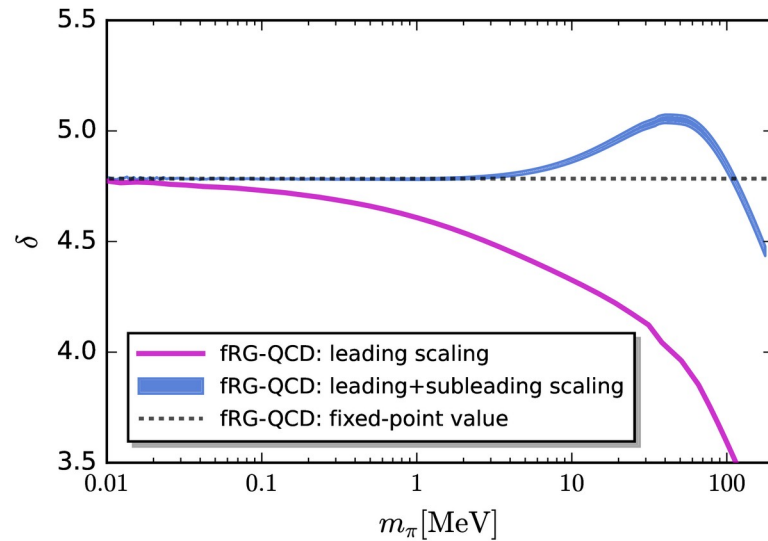
Results on physical point of QCD



Lattice data from Boucaud et al.
[Phys.Rev.D 98 (2018) 11, 114515]

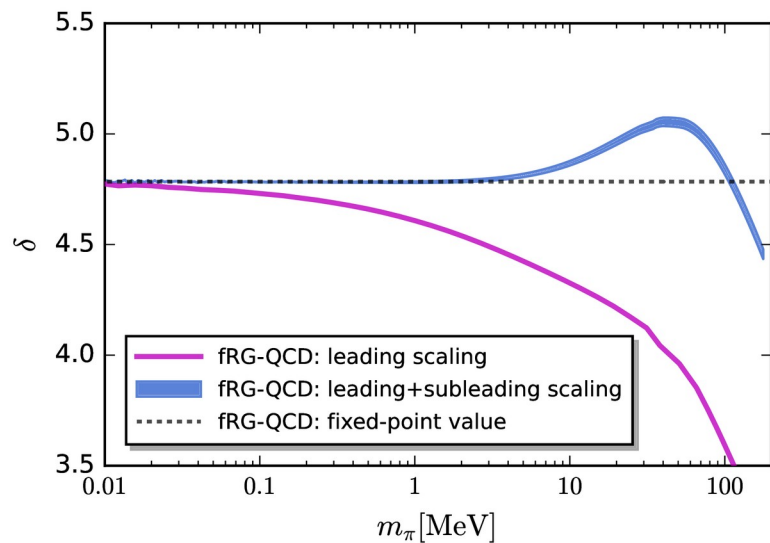
Soft modes in hot QCD matter

Braun, Chen, Fu, Gao, Huang, Ihssen,
Pawlowski, Rennecke, **Sattler**, Tan, Wen, Yin
(arXiv:2310.19853)



Soft modes in hot QCD matter

Braun, Chen, Fu, Gao, Huang, Ihssen,
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Columbia Plot

Pawlowski, **Sattler**, Steck
(in preparation)

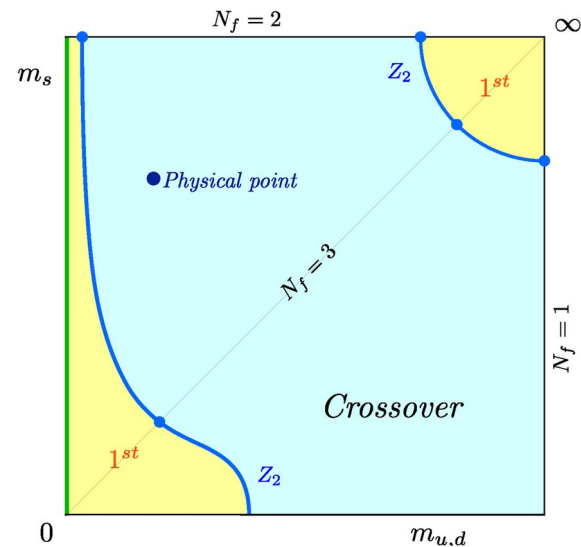
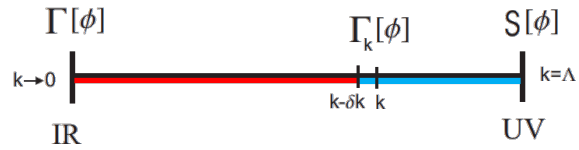


Figure from
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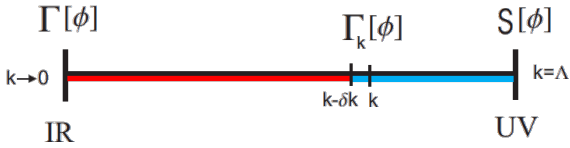
Wilsonian approach:

Integrate out momentum shells



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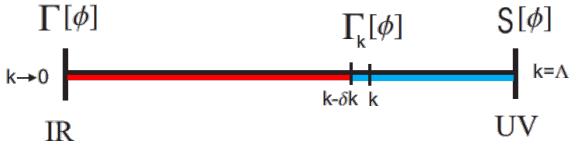
Integrate out momentum shells



$$Z_k[J] = \int [D\varphi]_k e^{-S[\varphi] + \int d^d x J^a(x) \varphi_a(x)}$$

Wilsonian approach:

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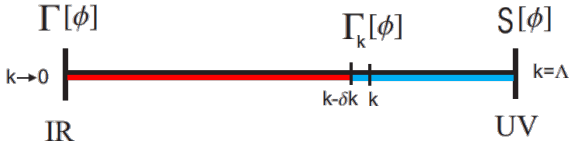
Introduce mass-like
“Regulator”

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$$[D\varphi]_k = [D\varphi]_{\text{ren}} e^{-\frac{1}{2} \int d^d x \varphi_a(x) R_k^{ab}(x) \varphi_b(x)}$$

Wilsonian approach:

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Introduce mass-like
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Obtain Flow equation

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$$\partial_t \Gamma_k[\phi] = \frac{1}{2} \sum_{a,b} \int \frac{d^d p}{(2\pi)^d} G_{ab}^{(2)}[\phi](p) \partial_t R_k^{ab}(p)$$

Infinite Tower of Functional equations

$$\partial_t \Gamma[\bar{\phi}] = \frac{1}{2} \text{Tr} G_k \partial_t R_k,$$

$$\partial_t \Gamma^{(1)}[\bar{\phi}] = -\frac{1}{2} \text{Tr} \Gamma_k^{(3)} (G_k \partial_t R_k G_k),$$

$$\partial_t \Gamma^{(2)}[\bar{\phi}] = -\frac{1}{2} \text{Tr} [\Gamma_k^{(4)} - 2 \Gamma_k^{(3)} G_k \Gamma_k^{(3)}] (G_k \partial_t R_k G_k),$$

$$\partial_t \Gamma^{(3)}[\bar{\phi}] = -\frac{1}{2} \text{Tr} [\Gamma_k^{(5)} - 6 \Gamma_k^{(4)} G_k \Gamma_k^{(3)} + 6 \Gamma_k^{(3)} G_k \Gamma_k^{(3)} G_k \Gamma_k^{(3)}] (G_k \partial_t R_k G_k),$$

$$\begin{aligned} \partial_t \Gamma^{(4)}[\bar{\phi}] = & -\frac{1}{2} \text{Tr} [\Gamma_k^{(6)} - 8 \Gamma_k^{(5)} G_k \Gamma_k^{(3)} - 6 \Gamma_k^{(4)} G_k \Gamma_k^{(4)} + 18 \Gamma_k^{(4)} G_k \Gamma_k^{(3)} G_k \Gamma_k^{(3)} \\ & + 12 \Gamma_k^{(3)} G_k \Gamma_k^{(4)} G_k \Gamma_k^{(3)} - 24 G_k \Gamma_k^{(3)} G_k \Gamma_k^{(3)} G_k \Gamma_k^{(3)} \cdot G_k \Gamma_k^{(3)}] (G_k \partial_t R_k G_k), \end{aligned}$$

⋮ ⋮

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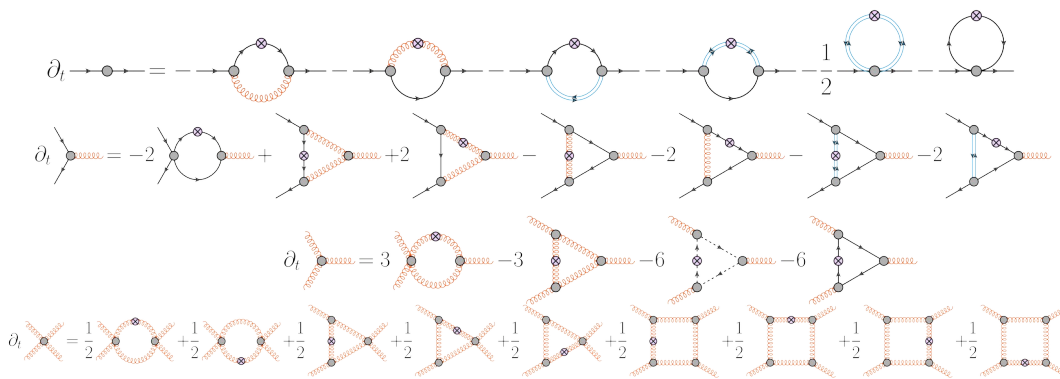
⋮

⋮

=

=

Infinite Tower of Diagrams



⋮

⋮

⋮

⋮