A new universality class describes Vicsek's flocking phase in physical dimensions

P. Jentsch, C.F. Lee, Phys. Rev. Lett. 133, 128301 (2024)

Patrick Jentsch Chiu Fan Lee



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Flocking



https://www.youtube.com/watch?v=q6iXT4-Oc2Q



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https://www.quantamagazine.org/cells-blaze-their-own-trails-tonavigate-through-the-body-20220328/

The Vicsek Model



The Vicsek Model



Update rule:

$$\mathbf{x}_{i}(t + \Delta t) = \mathbf{x}_{i}(t) + \begin{pmatrix} \cos(\theta_{i}(t)) \\ \sin(\theta_{i}(t)) \end{pmatrix} v_{m} \Delta t$$

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$$\theta_{i}(t + \Delta t) = \langle \theta_{i}(t) \rangle_{R} + \xi_{i}(t)$$





Disorder

T. Vicsek, et al., PRL (1995).

Banding

H. Chaté, et al., PRE (2008).

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Two Key Results

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1. Nonequilibrium Continuous Order-Disorder T. Vicsek, et al., PRL (1995) Transition



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TAISIUU



2. Spontaneous Sym. Breaking in 2D

Two Key Results

G. Grégoire and H. Chaté, PRL (2004)

1. Nonequilibrium Continuous Order-Disorder T. Vicsek, et al., PRL (1995) Transition

2. Spontaneous Sym. Breaking in 2D

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3. Nonequilibrium Ordered Phase w/ nontrivial Toner, Tu PRL (1995) Exponents





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Target: Scaling Exponents of Flocking Phase

$$\langle \delta \mathbf{g}(t, x, \mathbf{r}_{\perp}) \delta \mathbf{g}(0, 0, \mathbf{0}) \rangle = r_{\perp}^{2 \chi} S_g \left(\frac{t}{r_{\perp}^{\mathbf{Z}}}, \frac{x - v_0 t}{r_{\perp}^{\boldsymbol{\zeta}}} \right)$$

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 $\eta > 0 \rightarrow$ Stable Ordered Phase in d = 2

Toner, Tu, Phys. Rev. Lett (1995)

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Under assumptions of:

- Mass conservation
 - Translation sym.
 - Rotation sym.
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Continuity equation for density

$$\partial_t \rho = -\boldsymbol{\nabla} \cdot \boldsymbol{g}$$

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$$\begin{split} \gamma \partial_t \boldsymbol{g} &+ \lambda_1 \boldsymbol{g} \cdot \boldsymbol{\nabla} \boldsymbol{g} + \lambda_2 \boldsymbol{g} \boldsymbol{\nabla} \cdot \boldsymbol{g} + \lambda_3 \boldsymbol{\nabla} (|\boldsymbol{g}|^2) \\ &= -U(\rho, g^2) \boldsymbol{g} - P_1(\rho, g^2) \boldsymbol{\nabla} \rho + \mu_1 \nabla^2 \boldsymbol{g} + \mu_2 \boldsymbol{\nabla} (\boldsymbol{\nabla} \cdot \boldsymbol{g}) \\ &+ \mu_3 (\boldsymbol{g} \cdot \boldsymbol{\nabla})^2 \boldsymbol{g} + P_2 \boldsymbol{g} (\boldsymbol{g} \cdot \boldsymbol{\nabla} \rho) + \dots + \boldsymbol{f} \end{split}$$

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No fluctuations in flocking direction

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$$(\boldsymbol{g} = g_0 \hat{\boldsymbol{x}} + \boldsymbol{g}_{\perp}, \boldsymbol{g}_{\perp} = \delta \boldsymbol{g}_L + \delta \boldsymbol{g}_T)$$

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No nonlinear density terms

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$$\begin{split} \gamma \partial_t \boldsymbol{g}_{\perp} + \lambda_1 g_0 \partial_x \boldsymbol{g}_{\perp} + \lambda_1 \boldsymbol{g}_{\perp} \cdot \boldsymbol{\nabla}_{\perp} \boldsymbol{g}_{\perp} + \lambda_2 \boldsymbol{g}_{\perp} \boldsymbol{\nabla}_{\perp} \cdot \boldsymbol{g}_{\perp} + \lambda_3 \boldsymbol{\nabla}_{\perp} (|\boldsymbol{g}_{\perp}|^2) \\ &= -\beta |\boldsymbol{g}_{\perp}|^2 \boldsymbol{g}_{\perp} - \kappa_1 \boldsymbol{\nabla} \rho + \mu_1 (\boldsymbol{\nabla}_{\perp}^2 + \partial_x^2) \boldsymbol{g}_{\perp} + \mu_2 \boldsymbol{\nabla}_{\perp} (\boldsymbol{\nabla}_{\perp} \cdot \boldsymbol{g}_{\perp}) \\ &+ \mu_3 g_0^2 \partial_x^2 \boldsymbol{g}_{\perp}^2 \boldsymbol{g}_{\perp} \boldsymbol{f} \end{split}$$

1

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 Martin-Siggia-Rose-de Dominicis-Janssen → Ansatz for effective action

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$$\Gamma_{k}\left[\overline{\boldsymbol{g}}_{\perp},\boldsymbol{g}_{\perp},\overline{\rho},\rho\right] = \int \{\overline{\rho}[\partial_{t}\rho + \boldsymbol{\nabla}_{\perp}\cdot\boldsymbol{g}_{\perp}] - D|\overline{\boldsymbol{g}}_{\perp}|^{2} + \overline{\boldsymbol{g}}_{\perp}\cdot[\gamma\partial_{t}\boldsymbol{g}_{\perp} + \lambda_{1}g_{0}\partial_{x}\boldsymbol{g}_{\perp}] \}$$

$$+\lambda_1 \boldsymbol{g}_{\perp} \cdot \boldsymbol{\nabla}_{\perp} \boldsymbol{g}_{\perp} + \lambda_2 \boldsymbol{g}_{\perp} \boldsymbol{\nabla}_{\perp} \cdot \boldsymbol{g}_{\perp} + \lambda_3 \boldsymbol{\nabla}_{\perp} (|\boldsymbol{g}_{\perp}|^2) + \beta |\boldsymbol{g}_{\perp}|^2 \boldsymbol{g}_{\perp} + \kappa_1 \boldsymbol{\nabla} \rho$$
$$-\mu_1 (\boldsymbol{\nabla}_{\perp}^2 + \partial_x^2) \boldsymbol{g}_{\perp} - \mu_2 \boldsymbol{\nabla}_{\perp} (\boldsymbol{\nabla}_{\perp} \cdot \boldsymbol{g}_{\perp}) - \mu_3 g_0^2 \partial_x^2 \boldsymbol{g}_{\parallel} \}$$

Regulator

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• Sharp regulator for q_{\perp} only

$$R_k(\boldsymbol{q}_{\perp}, \boldsymbol{q}_x, \omega) = \Gamma_k^{(2)}(\boldsymbol{q}_{\perp}, \boldsymbol{q}_x, \omega) \left(\frac{1}{\Theta_{\epsilon}(|\boldsymbol{q}_{\perp}| - k)} - 1\right)$$

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- q_x and ω remain unregulated
- $\rightarrow \omega$ and q_{\perp} Integrals can be performed analytically





Two possible fixed points



Two possible fixed points • $\frac{\overline{\kappa}_1^7}{\overline{\lambda}_g^{13}} = 0, \, \eta_{\chi} = 0 \rightarrow \text{TT95}$



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Two possible fixed points

$$\frac{\overline{\kappa}_{1}^{7}}{\overline{\lambda}_{g}^{13}} = 0, \ \eta_{x} = 0 \rightarrow \text{TT95}$$

$$\partial_{k} \frac{\overline{\kappa}_{1}^{7}}{\overline{\lambda}_{g}^{13}} = 0 \quad \eta_{x} > 0 \rightarrow \text{New Fixed point}$$

$$\text{Nontrivial Scaling Relation}$$

$$7(2z - 2) = 13(z - \zeta)$$



 \rightarrow 3 scaling relations fix exponents



2 other vanishing loop corrections

 \rightarrow 3 scaling relations fix exponents

$$\chi = \frac{13(1-d)}{40}, \qquad z = \frac{27+13d}{40}, \qquad \zeta = \frac{41-d}{40}$$

Scaling Exponents



Scaling Exponents





Fixed points*



*Qualitative Flow: Flow lines are fictitious

Gauss.



• We used nonperturbative, functional RG to study a simplified TT model



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- TT UC applies for

$$\frac{11}{3} (\approx 3.67) < d < 4$$

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- We used nonperturbative, functional RG to study a simplified TT model
- TT UC applies for

 $\frac{11}{3} (\approx 3.67) < d < 4$

 Below d = 11/3, a new UC emerges, whose scaling exponents agree remarkably well with simulation in 2D & 3D



Thank you!



IMPERIAL



Andrew Killeen Sulaimaan Lim Sam Whitby Alastar Phelan John-Antonio Argyriadis Adam Kline

Chiu Fan Lee Léonie Canet Gunnar Pruessner

P. Jentsch, C.F. Lee, Phys. Rev. Lett. 133, 128301 (2024)