$\label{eq:information} Information\ geometry\ and\ the\ renormalization\ group$

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12th International Conference on the Exact Renormalization Group, September 2024 Les Diablerets, Suisse



The ideas of information geometry

[Ronald A. Fisher, Calyampudi R. Rao, Shun'ich Amari, Nikolai N. Chentsov, ...]

- studies spaces of probability distributions $p(x,\xi)$ with parameters ξ^{α}
- Fisher information metric (symmetric, positive semi-definite)

$$G_{\alpha\beta}(\xi) = \int dx \, p(x,\xi) \left(\frac{\partial}{\partial \xi^{\alpha}} \ln p(x,\xi) \right) \left(\frac{\partial}{\partial \xi^{\beta}} \ln p(x,\xi) \right)$$

- unique Riemannian metric that is invariant under sufficient statistics [Chentsov 1972]
- higher geometric structure: pair of dual connections, non-metricity etc. [Amari, Chentsov, ...]
- extension to quantum states $ho(\xi)$
- geometric structure follows from a *divergence* or *relative entropy*

$$D(p||q) = \int dx \ p(x) \ln(p(x)/q(x))$$

Sufficient statistics and Chentsov's theorem

- start from random variable x with probability distribution $p(x,\xi)$ where ξ^{α} are parameters
- \bullet consider map to new random variable $x \to y = f(x)$ with probability distribution $q(y,\xi)$
- information about ξ^{lpha} could get lost in the map
- new random variable y is called sufficient statistic for ξ when no information about ξ is lost:

 $p(x,\xi) = p(x|y,\xi)q(y,\xi) = r(x)q(y,\xi)$ factorizes

or

$$p(x|y,\xi) = rac{p(x,\xi)}{q(y,\xi)} = r(x)$$
 independent of ξ^{lpha}

Chentsov's theorem: unique invariant metric for sufficient statistic is

$$\begin{aligned} G_{\alpha\beta}(\xi) &= \int dx \, p(x,\xi) \left(\frac{\partial}{\partial \xi^{\alpha}} \ln p(x,\xi) \right) \left(\frac{\partial}{\partial \xi^{\beta}} \ln p(x,\xi) \right) \\ &= \int dy \, q(y,\xi) \left(\frac{\partial}{\partial \xi^{\alpha}} \ln q(y,\xi) \right) \left(\frac{\partial}{\partial \xi^{\beta}} \ln q(y,\xi) \right) \end{aligned}$$

$Relative \ entropy$

• classical relative entropy or Kullback-Leibler divergence

$$D(p||q) = \sum_{j} p_j \ln(p_j/q_j)$$

• not symmetric distance measure, but a divergence

 $D(p||q) \ge 0$ and $D(p||q) = 0 \iff p = q$

• quantum relative entropy of two density matrices (also a divergence)

$$D(\rho \| \sigma) = \mathsf{Tr} \left\{ \rho \left(\ln \rho - \ln \sigma \right) \right\}$$

- ullet signals how well state ρ can be distinguished from a model σ
- Gibbs inequality: $D(\rho \| \sigma) \ge 0$
- $D(\rho \| \sigma) = 0$ if and only if $\rho = \sigma$

Significance of Kullback-Leibler divergence

Uncertainty deficit

- true distribution p_j and model distribution q_j
- uncertainty deficit is expected surprise $\langle -\ln q_j \rangle = -\sum_j p_j \ln q_j$ minus real information content $-\sum_j p_j \ln p_j$

$$D(p||q) = -\sum_{j} p_{j} \ln q_{j} - \left(-\sum_{j} p_{j} \ln p_{j}\right)$$

Asymptotic frequencies

- true distribution q_j and frequency after N drawings $p_j = \frac{N(x_j)}{N}$
- probability to find frequencies p_j for large N (similar: Sanov theorem)

 $\sim e^{-ND(p\|q)}$

• probability for fluctuation around expectation value $\langle p_j \rangle = q_j$ tends to zero for large N and when divergence D(p||q) is large

Advantages of relative entropy: continuum limit

• consider transition from discrete to continuous distributions

$$p_j \to f(x) dx$$
 $q_j \to g(x) dx$

not well defined for entropy

$$S = -\sum p_j \ln p_j \xrightarrow{i} - \int dx f(x) \left[\ln f(x) + \ln dx \right]$$

• relative entropy remains well defined

$$D(p||q) \rightarrow D(f||g) = \int dx f(x) \ln(f(x)/g(x))$$

Information geometry for Euclidean quantum fields

[S. Floerchinger, 2303.04081 and Phys. Lett B 846, 138244 (2023)]

- classical statistical field theories
- bosonic quantum fields with real action in Euclidean space
- extension of information geometry
- new flow equation for divergence functional

Probabilities for Euclidean fields: exponential family

• probability density for Euclidean field theory with respect to measure $D\chi$

$$p[\chi, J] = \exp\left(-I[\chi] + J^{\alpha}\phi_{\alpha}[\chi] - W[J]\right)$$

• uses abstract index notation

$$J^{\alpha}\phi_{\alpha} = \int_{x} \sum_{n} J_{n}(x)\phi_{n}(x)$$

partition function

$$e^{W[J]} = \int D\chi \exp\left(-I[\chi] + J^{\alpha}\phi_{\alpha}[\chi]\right)$$

- \bullet sources J^{α} could also compromise coupling constants
- will be considered as coordinates on space of probability distributions
- known as exponential family in information geometry

Affine geometry for sources

• exponential family is closed with respect to affine transformations

 $J^{\alpha} \to J^{\prime \alpha} = M^{\alpha}_{\ \beta} J^{\beta} + c^{\alpha}$

- affine transformations respect convexity of W[J]
- so-called *e*-geodesics

$$J^{\alpha}(t) = (1-t)J^{\prime \alpha} + tJ^{\prime \prime \alpha}$$

characterized by differential equation

$$\frac{d^2}{dt^2}J^{\alpha}(t) + \left(\Gamma_{\mathsf{E}}\right)_{\beta}{}^{\alpha}{}_{\gamma}[J] \,\left(\frac{d}{dt}J^{\beta}(t)\right)\left(\frac{d}{dt}J^{\gamma}(t)\right) = 0$$

where the connection vanishes in terms of source coordinates

 $(\Gamma_{\mathsf{E}})^{\ \alpha}_{\beta\ \gamma}[J] = 0$

Fisher information metric

• Fisher information metric

$$\begin{aligned} G_{\alpha\beta}[J] &= \int D\chi \, p[\chi,J] \, \frac{\delta}{\delta J^{\alpha}} \ln p[\chi,J] \, \frac{\delta}{\delta J^{\beta}} \ln p[\chi,J] \\ &= -\int D\chi \, p[\chi,J] \, \frac{\delta^2}{\delta J^{\alpha} \delta J^{\beta}} \ln p[\chi,J] \end{aligned}$$

• Fisher-Rao distance between nearby probability distributions

$$ds^2 = G_{\alpha\beta}[J]dJ^{\alpha}dJ^{\beta}$$

• for the exponential family

$$G_{\alpha\beta}[J] = \frac{\delta^2}{\delta J^\alpha \delta J^\beta} W[J] = \langle \phi_\alpha[\chi] \phi_\beta[\chi] \rangle - \langle \phi_\alpha[\chi] \rangle \langle \phi_\beta[\chi] \rangle$$

- equal to connected two-point correlation function !
- generalization of Zamolodchikov metric for conformal field theories

Expectation value coordinates

• can also use field expectation values as coordinates for $p[\chi,\Phi]$

$$\Phi_{\alpha} = \langle \phi_{\alpha}[\chi] \rangle = \frac{\delta}{\delta J^{\alpha}} W[J] = \int D\chi \, p[\chi, J] \, \phi_{\alpha}[\chi]$$

• best described in terms of quantum effective action

$$\Gamma[\Phi] = \sup_{J} \left(J^{\alpha} \Phi_{\alpha} - W[J] \right) = -\inf_{J} \left(-\int D\chi \, p[\chi, J] \ln p[\chi, J] \right)$$

Fisher-Rao distance

$$ds^{2} = G_{\alpha\beta}[J] \,\delta J^{\alpha} \delta J^{\beta} = G^{\alpha\beta}[\Phi] \,\delta \Phi_{\alpha} \delta \Phi_{\beta} = \delta J^{\alpha} \delta \Phi_{\beta}$$

• Fisher metric in expectation value coordinates

$$G^{\alpha\beta}[\Phi] = -\int D\chi \, p[\chi,\Phi] \, \frac{\delta^2}{\delta\Phi_\alpha\delta\Phi_\beta} \ln p[\chi,\Phi] = \frac{\delta^2\Gamma[\Phi]}{\delta\Phi_\alpha\delta\Phi_\beta}$$

• another affine structure, dual to the one for sources

$$\Phi_{\alpha} \to \Phi_{\alpha}' = N_{\alpha}^{\ \beta} \Phi_{\beta} + d_{\alpha}$$

• defines so-called *m*-connection

Divergence functional in source coordinates

• functional generalization of Kullback-Leibler divergence

$$D[J||J'] = \int D\chi \, p[\chi, J] \ln \left(p[\chi, J] / p[\chi, J'] \right)$$

- compares two probability distributions in asymmetric way
- non-negative

$D[J\|J'] \ge 0$

• equals Fisher-Rao distance for close-by distributions

$$D[J||J'] = \frac{1}{2} G_{\alpha\beta}[J] \delta J^{\alpha} \delta J^{\beta} + \dots$$

- characterizes probabilities for large deviations (Sanovs theorem)
- can be written as Bregman divergence

$$D[J||J'] = (J^{\alpha} - J'^{\alpha})\frac{\delta W[J]}{\delta J^{\alpha}} - W[J] + W[J']$$

• functional derivatives w.r.t. second argument yield connected correlation functions !

Divergence functional in expectation value coordinates

• Divergence functional in terms of expectation values

$$D[\Phi \| \Phi'] = \int D\chi \, p[\chi, \Phi] \ln \left(p[\chi, \Phi] / p[\chi, \Phi'] \right)$$
$$= \Gamma[\Phi] - \Gamma[\Phi'] - \frac{\delta \Gamma[\Phi']}{\delta \Phi'_{\lambda}} (\Phi_{\lambda} - \Phi'_{\lambda})$$

• functional derivatives w.r.t. first argument yield one-particle irreducible correlation functions (for $n \ge 2$)

 $D^{(n,0)}[\Phi \| \Phi'] = \Gamma^{(n)}[\Phi],$

• mixed representation generates connected and 1-P.I. correlation functions

$$D[\Phi||J'] = \Gamma[\Phi] + W[J'] - J'^{\alpha}\Phi_{\alpha}$$

Functional integral representations

• divergence functional in source coordinates

$$e^{-D[J||J']} = \frac{e^{W[J] - J^{\alpha}\Phi_{\alpha}}}{e^{W[J'] - J'^{\alpha}\Phi_{\alpha}}} = \frac{\int D\chi \exp\left(-I[\chi] + J^{\alpha}(\phi_{\alpha}[\chi] - \Phi_{\alpha})\right)}{\int D\tilde{\chi} \exp\left(-I[\tilde{\chi}] + J'^{\alpha}(\phi_{\alpha}[\tilde{\chi}] - \Phi_{\alpha})\right)}$$

- well defined as ratio of functional integrals
- similar in expectation value coordinates

Geometry from divergence

• Fisher metric from functional derivative of divergence

$$G_{\alpha\beta}[J] = -\frac{\delta^2}{\delta J^{\alpha} \delta J^{\prime\beta}} D[J||J']|_{J=J'}$$

- transforms automatically as a metric under coordinate changes $J \rightarrow K[J]$
- *m*-connection symbols

$$(\Gamma_{\mathsf{M}})_{\alpha\beta\gamma}[J] = -\frac{\delta^2}{\delta J^{\alpha}\delta J^{\gamma}} \frac{\delta}{\delta J'^{\beta}} D[J||J']\Big|_{J=J'}.$$

e-connection symbols

$$(\Gamma_{\mathsf{E}})_{\alpha\beta\gamma}[J] = -\frac{\delta}{\delta J^{\beta}} \frac{\delta^2}{\delta J'^{\alpha} \delta J'^{\beta}} D[J||J']|_{J=J'}$$

- automatically transform like connections under $J \rightarrow K[J]$
- information geometry nicely encoded in divergence functional !
- expectation values are another useful coordinate choice

Regularized probability distribution

• introduce now quadratic regulator in probability density

$$p_k[\phi, J] = \exp\left(-S[\phi] - \frac{1}{2}R_k^{\alpha\beta}\phi_\alpha\phi_\beta + J^\alpha\phi_\alpha - W_k[J]\right),$$

• with modified partition function

$$e^{W_k[J]} = \int D\phi \exp\left(-S[\phi] - \frac{1}{2}R_k^{\alpha\beta}\phi_{\alpha}\phi_{\beta} + J^{\alpha}\phi_{\alpha}\right).$$

• regulator can be chosen to suppress fluctuations, e. g.

$$R_k^{\alpha\beta} = k^2 \delta^{\alpha\beta}$$

Divergence functionals with regulator

• divergence functional with regulator

$$\begin{split} \tilde{D}_k[J||J'] &= \int D\phi \, p_k[\phi, J] \ln(p_k[\phi, J]/p_k[\phi, J']) \\ &= (J^\alpha - J'^\alpha) \frac{\delta \, W_k[J]}{\delta J^\alpha} - W_k[J] + \, W_k[J'], \end{split}$$

• flowing divergence in expectation value coordinates with regulator terms subtracted

$$D_{k}[\Phi \| \Phi'] = \tilde{D}_{k}[\Phi \| \Phi'] - \frac{1}{2} R_{k}^{\alpha\beta} (\Phi_{\alpha} - \Phi'_{\alpha}) (\Phi_{\beta} - \Phi'_{\beta})$$
$$= \Gamma_{k}[\Phi] - \Gamma_{k}[\Phi'] - \frac{\delta \Gamma_{k}[\Phi']}{\delta \Phi'_{\lambda}} (\Phi_{\lambda} - \Phi'_{\lambda}).$$

Limit of large and small regulator

 \bullet for large k saddle point approximation becomes valid

$$\lim_{k \to \infty} D_k[\Phi \| \Phi'] = S[\Phi] - S[\Phi'] - \frac{\delta}{\delta \Phi'_{\alpha}} S[\Phi'](\Phi_{\alpha} - \Phi'_{\alpha})$$

• for small k the full Kullback-Leibler divergence functional is recovered

 $\lim_{k \to 0} D_k[\Phi \| \Phi'] = D[\Phi \| \Phi']$

Flow equation for the divergence functional

[S. Floerchinger, Phys. Lett B 846, 138244 (2023)]

exact flow equation

$$\frac{\partial}{\partial k} D_k[\Phi \| \Phi'] = \frac{1}{2} \left(\frac{\partial}{\partial k} R_k^{\alpha\beta} \right) \left[(G_k[\Phi])_{\alpha\beta} - (G_k[\Phi'])_{\alpha\lambda} (\tilde{D}_k^{(0,2)}[\Phi \| \Phi'])^{\lambda\kappa} (G_k[\Phi'])_{\kappa\beta} \right]$$

- close relative of Polchinskis and Wetterichs equations
- starting point for approximate solutions
- can be used to flow from large to small regulators
- flow vanishes when $\Phi = \Phi'$
- general coordinates changes possible

Conclusions

- information geometry concepts can be applied to quantum and statistical field theories
- divergence functional encodes the information about geometry: metric, e-connection, m-connection etc.
- divergence functional is generating functional for connected and one-particle irreducible correlation functions
- new exact flow equation for divergence functional

Backup

Advantages / disadvantages of divergence functional

- information theoretic meaning
- positivity $D[\Phi \| \Phi'] \ge 0$ instead of convexity for $\Gamma[\Phi]$
- geometric realization
 - connected correlation functions: e-connection
 - \bullet one-particle irreducible: *m*-connection
- general coordinate changes $\Phi o \Psi[\Phi]$

 $D[\Psi \| \Psi'] = D[\Phi[\Psi] \| \Phi'[\Psi']]$

preserve geometric structure

 \bullet equilibrium expectation value $\Phi_{\rm eq}$ corresponding to J=0 must be known in addition

Square roots of probabilities

• Fisher information metric

$$\begin{aligned} G_{\alpha\beta}(\xi) &= \int dx \, p(x,\xi) \left(\frac{\partial}{\partial \xi^{\alpha}} \ln p(x,\xi) \right) \left(\frac{\partial}{\partial \xi^{\beta}} \ln p(x,\xi) \right) \\ &= 4 \int dx \left(\frac{\partial}{\partial \xi^{\alpha}} \sqrt{p(x,\xi)} \right) \left(\frac{\partial}{\partial \xi^{\beta}} \sqrt{p(x,\xi)} \right) \end{aligned}$$

• for discrete random variable, take coordinates

$$p_j = \xi_j^2, \qquad j = 1, \dots, N.$$

normalization implies

$$\xi_1^2 + \ldots + \xi_N^2 = 1$$

• Fisher information metric is just induced Euclidean metric on the sphere!

