

Towards CP violation in baryons from lattice QCD

Maxwell T. Hansen

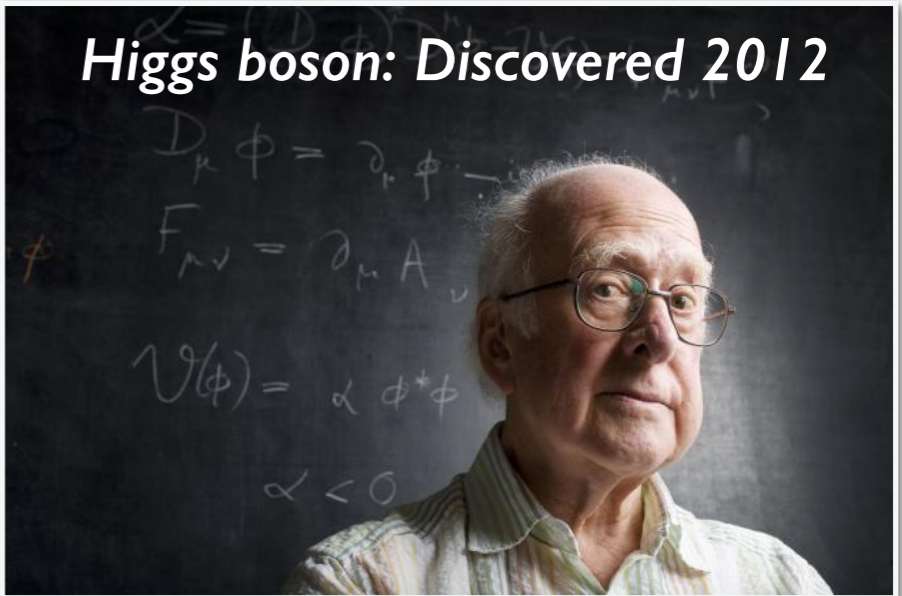
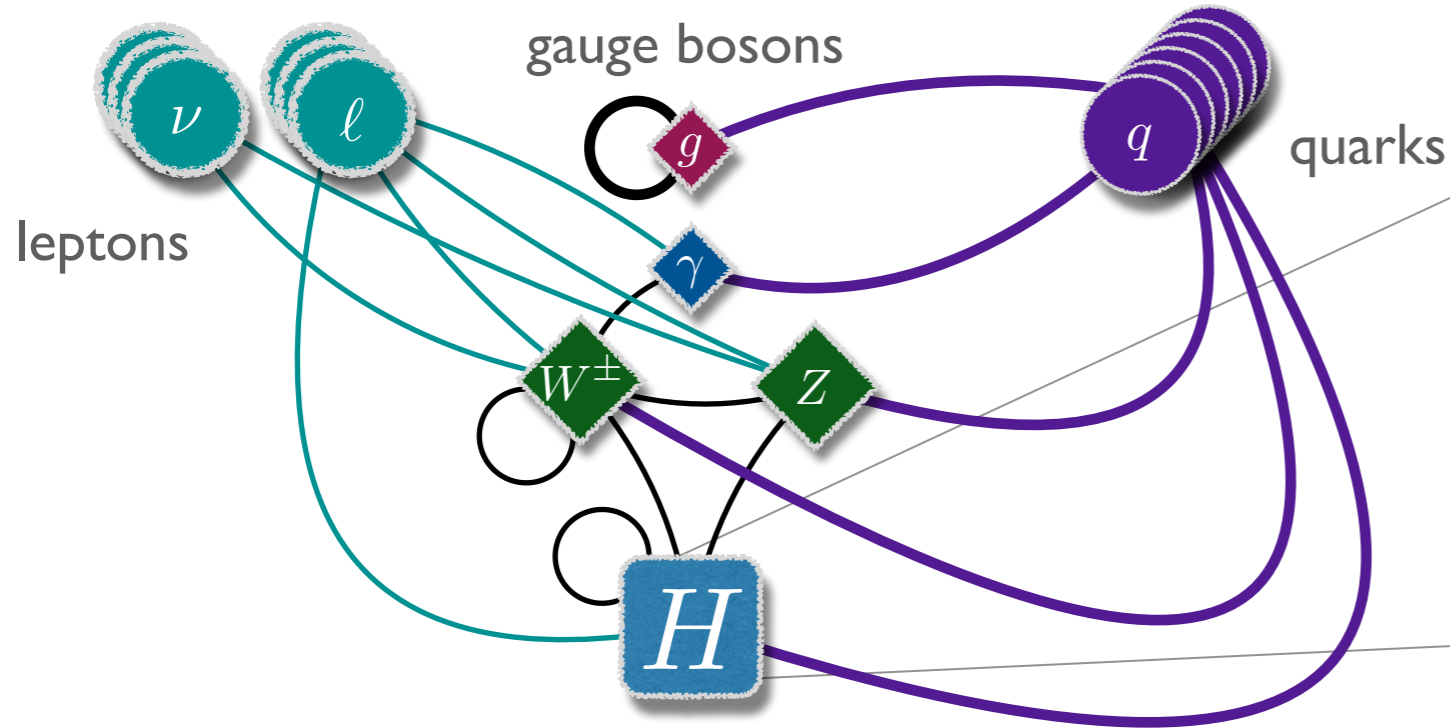
March 6th, 2023



THE UNIVERSITY
of EDINBURGH

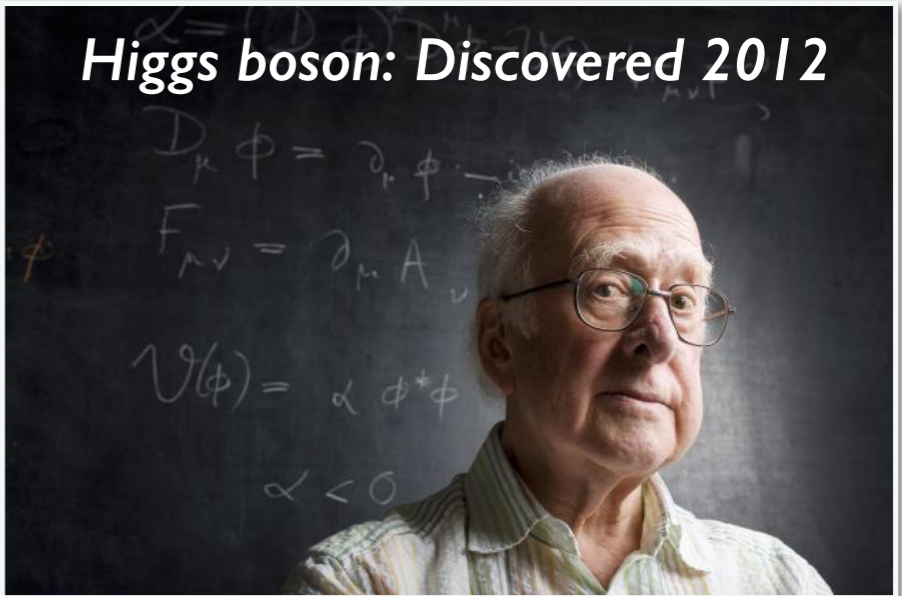
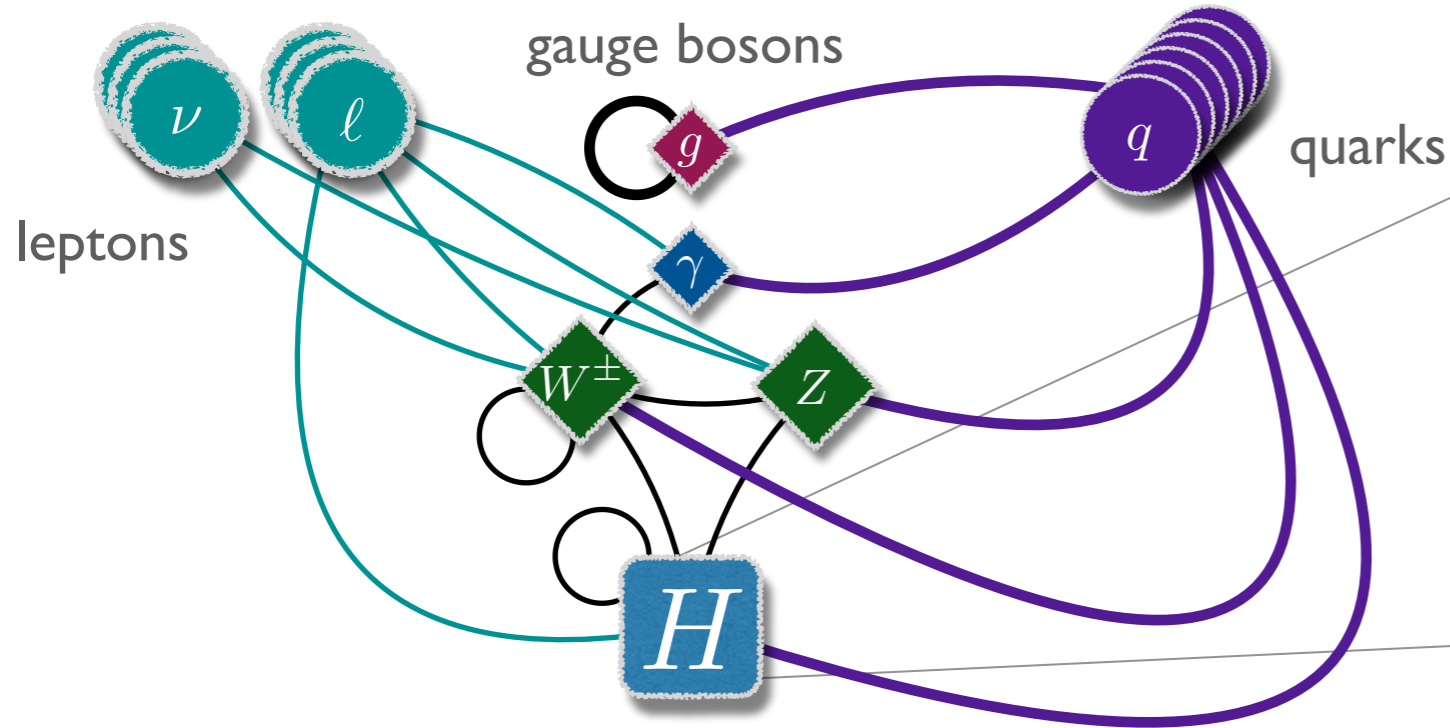
The Standard Model

□ A catalogue of particles and forces, highly constrained by *fundamental symmetries*



The Standard Model

□ A catalogue of particles and forces, highly constrained by *fundamental symmetries*

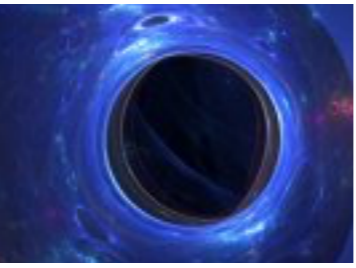


□ SM = a very *effective theory*... but definitely not the whole story

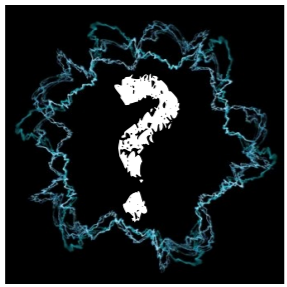
neutrino masses



dark matter



gravity



matter / anti-matter asymmetry
Why is our universe matter dominated?

Sakharov's conditions:

- Baryon number violation
- C and CP symmetry violation
- Interactions out of thermal equilibrium

Taking new physics seriously

□ SM is well known to have CPV

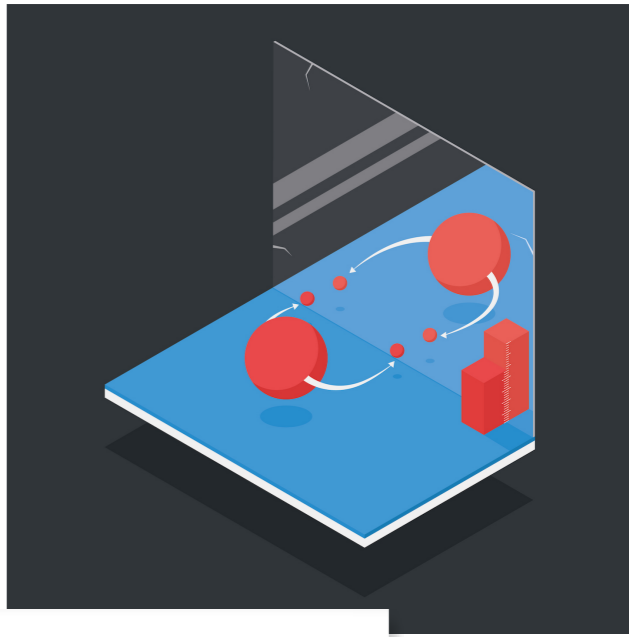
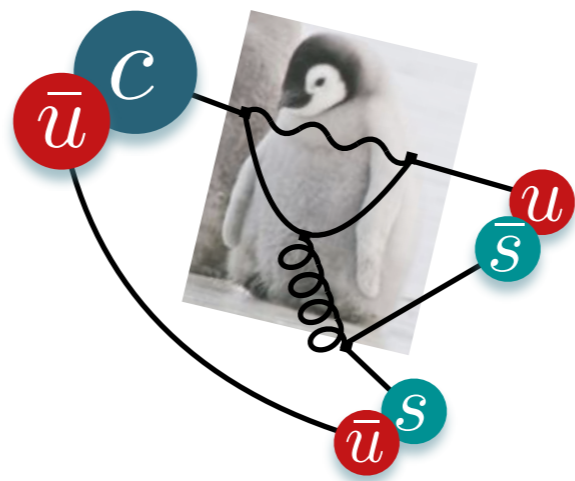
$$\text{Im}[V_{CKM}] \neq 0$$

...but not enough for *baryogenesis*!



□ 2019: LHCb observed CP violation in *charm*

$$D \rightarrow \pi\pi, K\bar{K}$$



$$\Delta A_{CP} = (-15.4 \pm 2.9) \times 10^{-4}$$

• LHCb (PRL, 2019) •

Taking new physics seriously

☐ SM is well known to have CPV

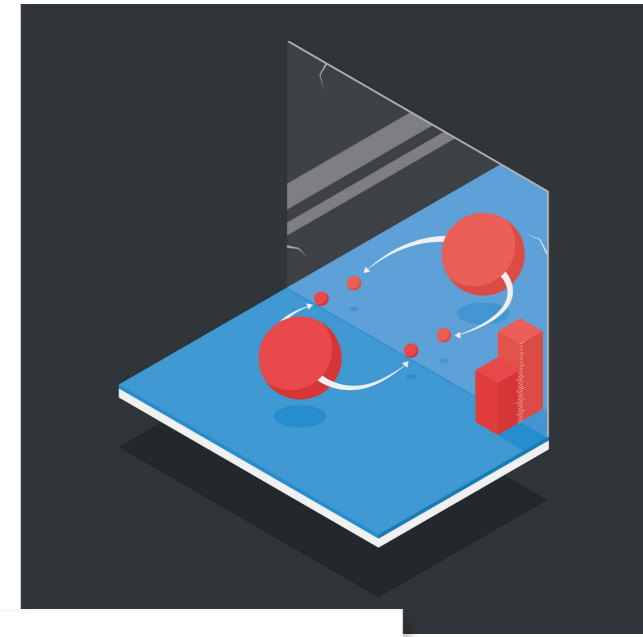
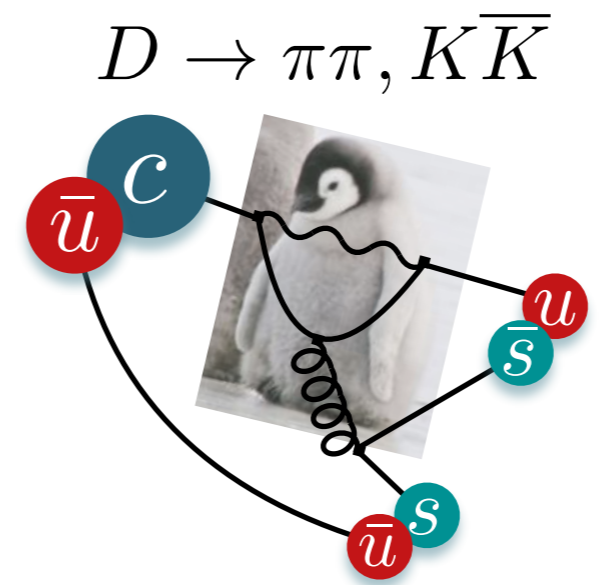
$$\text{Im}[V_{CKM}] \neq 0$$

...but not enough for *baryogenesis!*



Is this new physics?

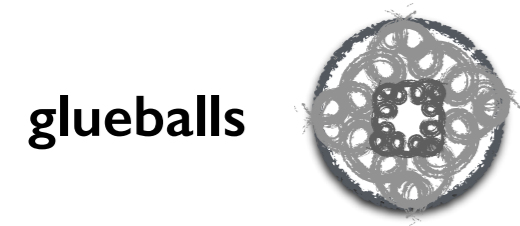
☐ 2019: LHCb observed CP violation in *charm*



$$\Delta A_{CP} = (-15.4 \pm 2.9) \times 10^{-4}$$

• LHCb (PRL, 2019) •

☐ No Standard Model prediction is available, due to complicated QCD dynamics

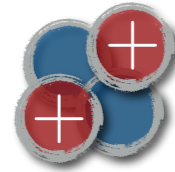


☐ What if CPV is seen in baryons? → same issue = *limited SM predictions*

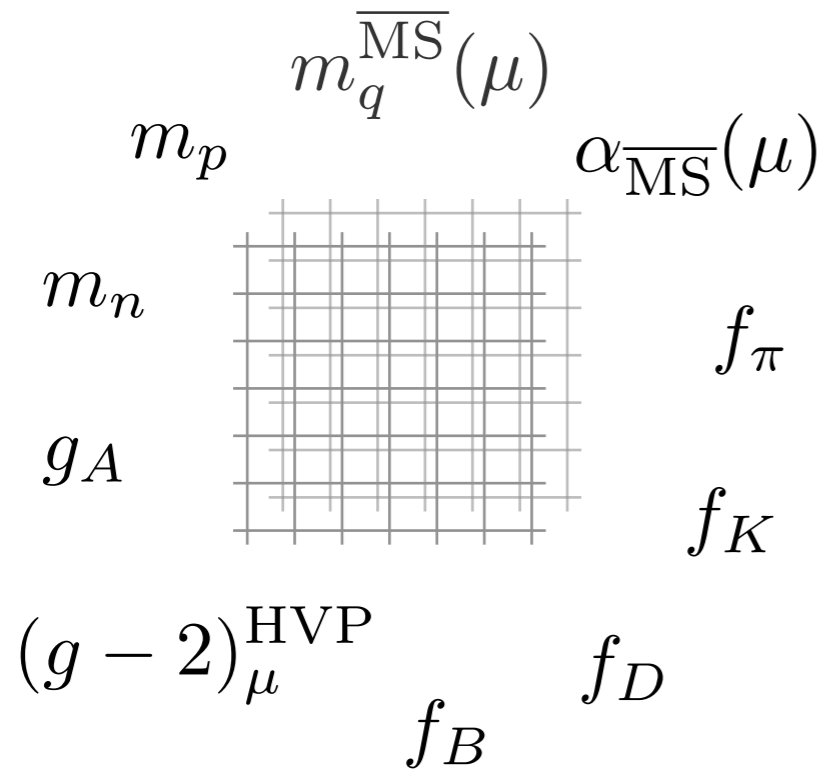
Lattice QCD: *Recipe for strong force predictions*

1. Lagrangian defining QCD
2. Formal / numerical machinery (lattice QCD)
3. A few experimental inputs (e.g. M_π, M_K, M_Ω)

$$\mathcal{L}_{\text{QCD}} = \sum_f \bar{\Psi}_f (i\not{D} - m_f) \Psi_f - \frac{1}{4} G_{\mu\nu}^a G_a^{\mu\nu}$$



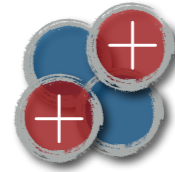
Wide range of precision pre-/post-dictions



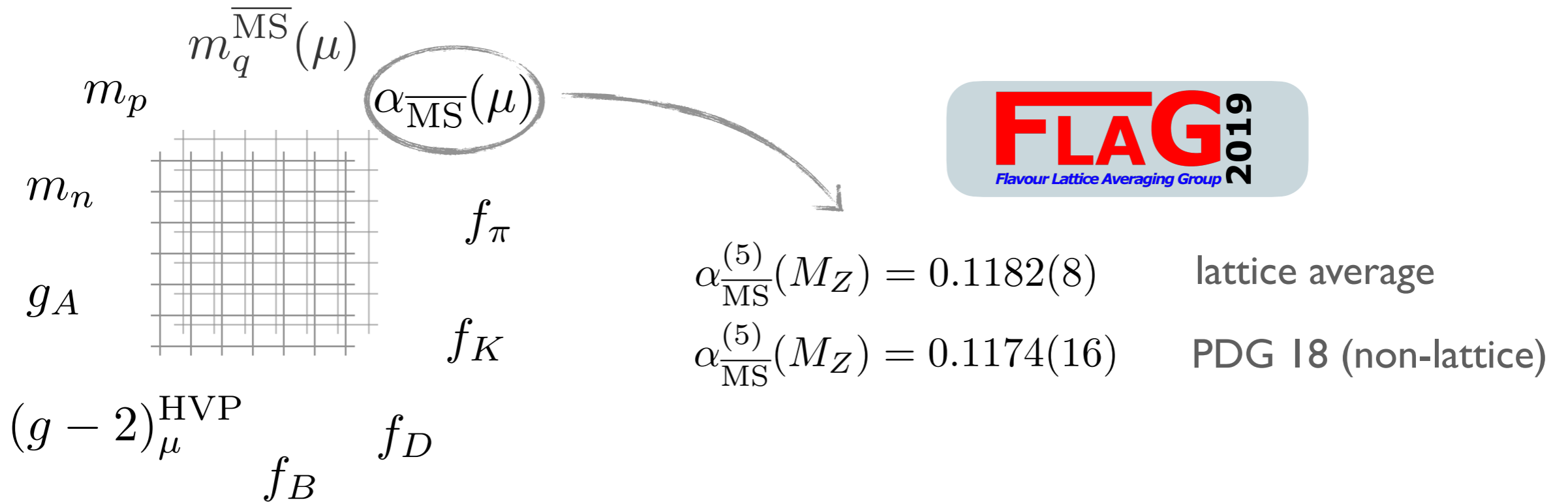
Lattice QCD: *Recipe for strong force predictions*

1. Lagrangian defining QCD
2. Formal / numerical machinery (lattice QCD)
3. A few experimental inputs (e.g. M_π, M_K, M_Ω)

$$\mathcal{L}_{\text{QCD}} = \sum_f \bar{\Psi}_f (i\not{D} - m_f) \Psi_f - \frac{1}{4} G_{\mu\nu}^a G_a^{\mu\nu}$$



Wide range of precision pre-/post-dictions



Overwhelming evidence for QCD ✓

Lattice QCD as a reliable tool ✓

More challenging observables? 😬

Baryonic amplitudes and LQCD

□ Essential to classify by types of states and inserted operators

- 1. Single-hadron form factors

$$n \rightarrow p e^- \bar{\nu}_e$$

$$\langle \text{hadron} | \mathcal{O}(0) | \text{hadron} \rangle$$

- 2. Multi-hadron decays

$$\Lambda_c^+ \rightarrow p K_S^0, \Lambda^0 \pi^+, \Sigma^+ \pi^0, \Sigma^0 \pi^+ \quad \Lambda_c^+ \rightarrow p K_S^0 \pi^+$$

$$\langle \text{multi-hadron state} | \mathcal{O}(0) | \text{hadron} \rangle$$

- 3. Intermediate multi-hadron states

$$\Sigma^+ \rightarrow p \ell^+ \ell^-$$

$$\langle \text{hadron} | \mathcal{J}(x) \mathcal{O}(0) | \text{hadron} \rangle$$

□ No published LQCD baryonic calculations for items 2. and 3.

Can LQCD calculate X for baryonic CPV?

easier

easier

harder

$$\langle \text{hadron} | \mathcal{O}(0) | \text{hadron} \rangle$$



mesons
quark bilinear
light-quarks

baryons
four-quark operators
heavy-quarks

$$\langle \text{multi-hadron state} | \mathcal{O}(0) | \text{hadron} \rangle$$

$$\langle \text{hadron} | \mathcal{J}(x) \mathcal{O}(0) | \text{hadron} \rangle$$

harder

Can LQCD calculate X for baryonic CPV?

easier

easier

harder

$$\langle \text{hadron} | \mathcal{O}(0) | \text{hadron} \rangle$$

mesons
quark bilinear
light-quarks

baryons
four-quark operators
heavy-quarks

easier

harder

$$\langle \text{multi-hadron state} | \mathcal{O}(0) | \text{hadron} \rangle$$

single two-particle channel

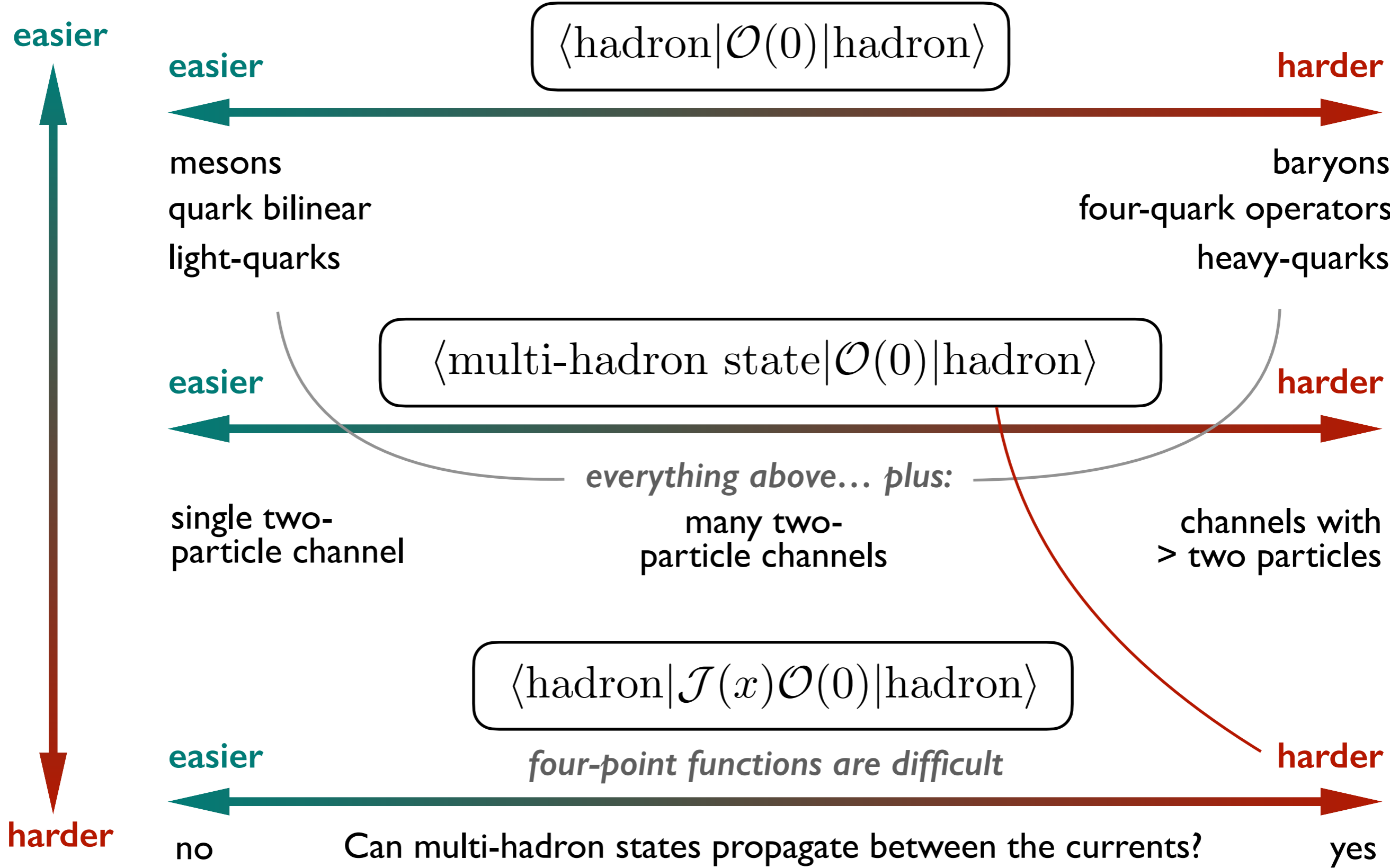
everything above... plus:
many two-particle channels

channels with > two particles

harder

$$\langle \text{hadron} | \mathcal{J}(x) \mathcal{O}(0) | \text{hadron} \rangle$$

Can LQCD calculate X for baryonic CPV?



Resonances

□ If multi-hadron states play a role... *resonances* could be relevant

□ Meson decays

CP violation in strange $K \rightarrow \pi\pi$

What is the role of the $\sigma/f_0(500)$?

CP violation in charm $D \rightarrow \pi\pi, K\bar{K}$

$f_0(1710)$ could enhance ΔA_{CP}
• Soni (2017) •

Resonant B decays $B \rightarrow K^* \ell\ell \rightarrow K\pi \ell\ell$

Resonances

□ If multi-hadron states play a role... *resonances* could be relevant

□ Meson decays

CP violation in strange $K \rightarrow \pi\pi$

What is the role of the $\sigma/f_0(500)$?

CP violation in charm $D \rightarrow \pi\pi, K\bar{K}$

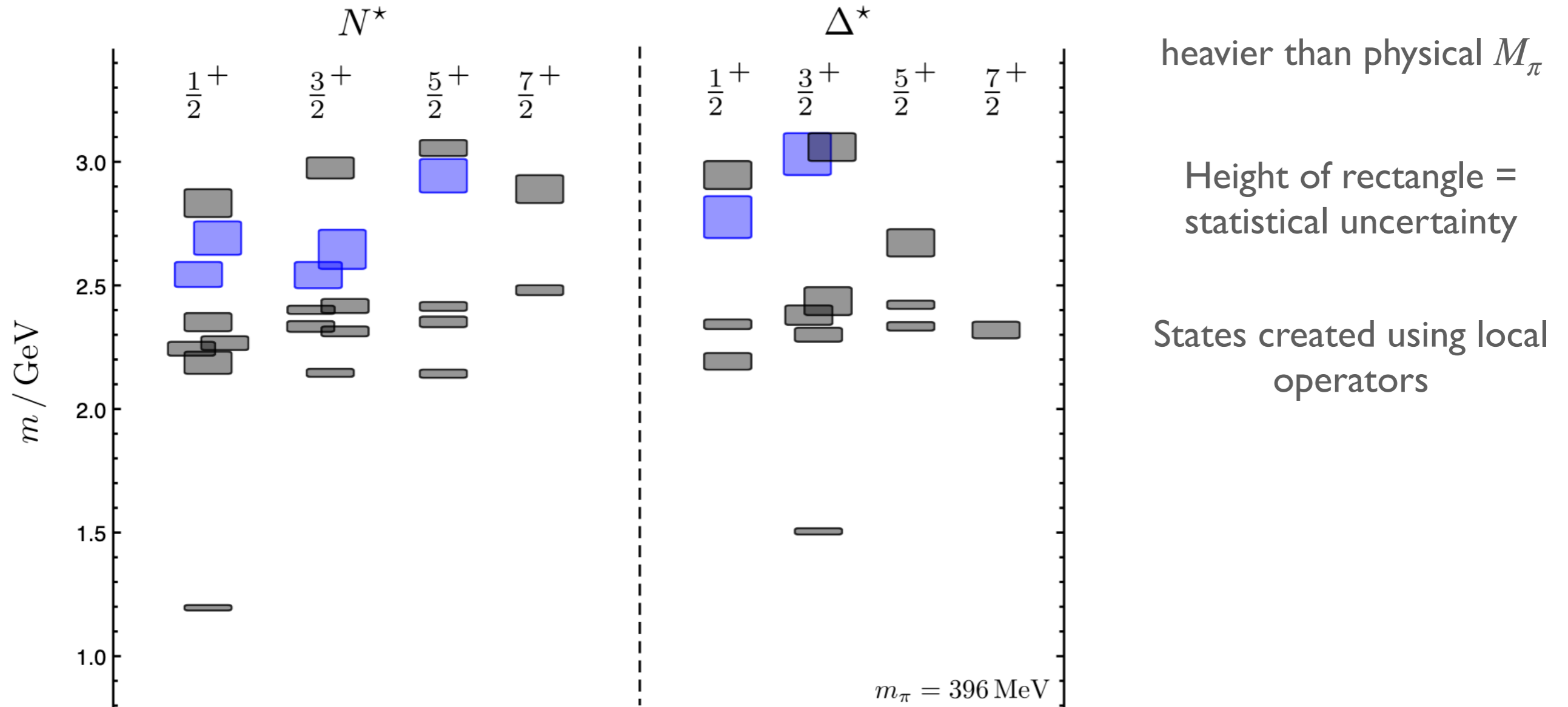
$f_0(1710)$ could enhance ΔA_{CP}
• Soni (2017) •

Resonant B decays $B \rightarrow K^* \ell\ell \rightarrow K\pi \ell\ell$

□ Of course, resonances = also important for baryonic processes

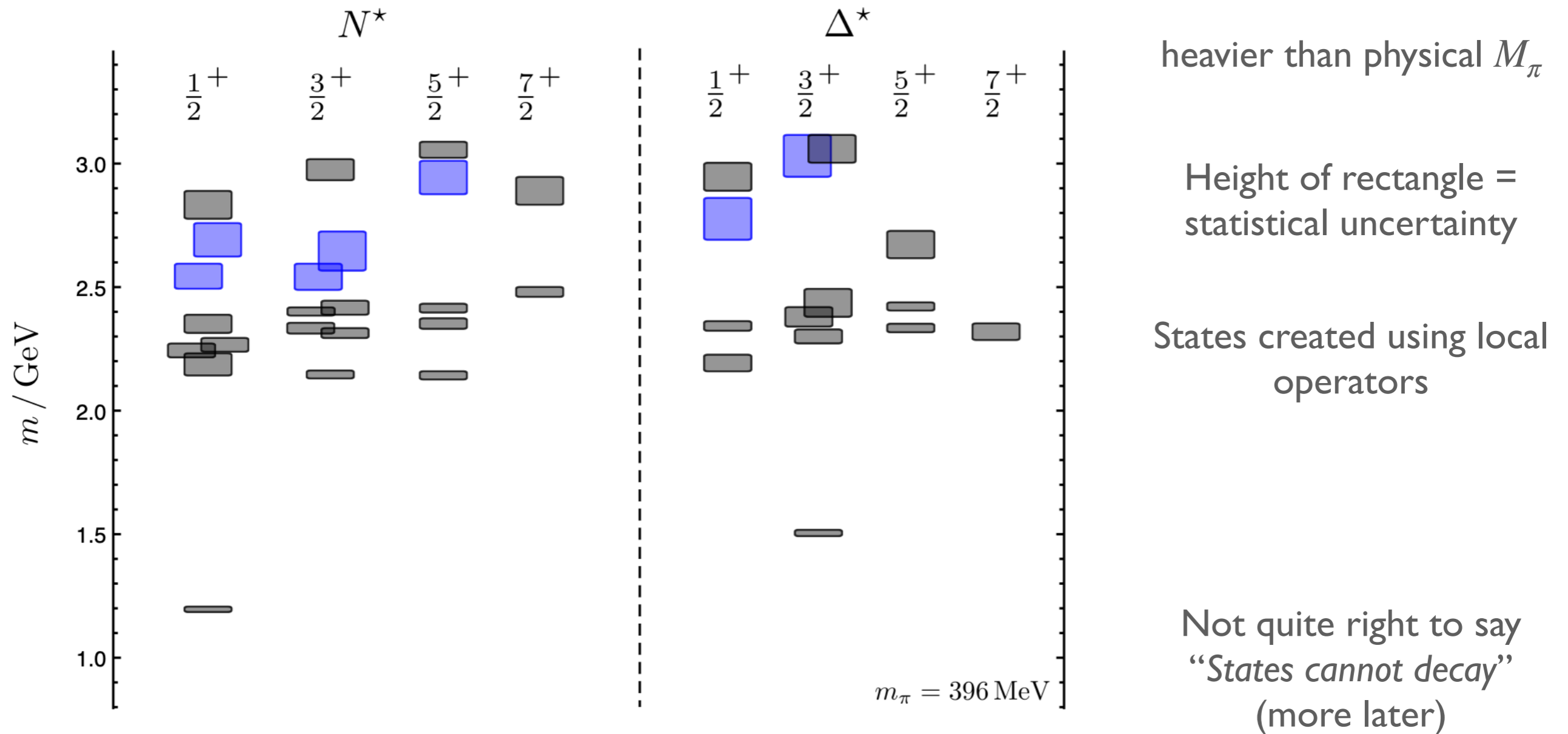
- Any reliable approach should consider whether such effects are relevant
- Can be both a challenge and an opportunity (also for lattice QCD)

Baryonic resonances



- Dudek and Edwards, *Hybrid Baryons in QCD*, PRD, 2012 •

Baryonic resonances

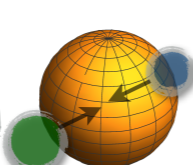
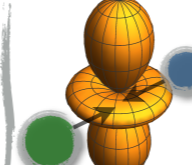
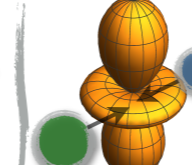
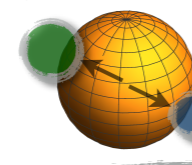
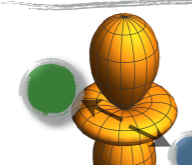
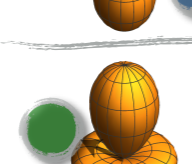


Remarkable progress... but not the complete picture!

- Dudek and Edwards, *Hybrid Baryons in QCD*, PRD, 2012 •

QCD Fock space

- At low-energies QCD = hadronic degrees of freedom $\pi \sim \bar{u}d, K \sim \bar{s}u, p \sim uud$
- Overlaps of multi-hadron *asymptotic states* \rightarrow S matrix

		$ N\pi, \text{in}\rangle$		
				
$S(s) \equiv \langle N\pi, \text{out} $		$e^{2i\delta_{1/2,0}(s)}$	0	0
		0	$e^{2i\delta_{1/2,1}(s)}$	0
		0	0	$e^{2i\delta_{3/2,1}(s)}$

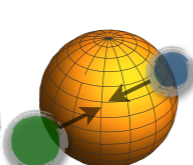
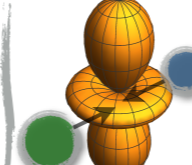
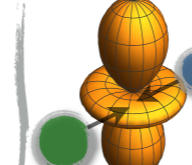
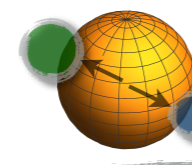
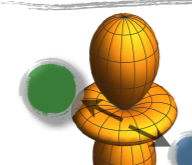
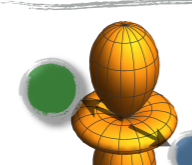
depends on $s = E_{\text{cm}}^2$
and angular variables

diagonal in total angular momentum

$$\mathcal{M}_{J,\ell}(s) \propto e^{2i\delta_{J,\ell}(s)} - 1$$

QCD Fock space

- At low-energies QCD = hadronic degrees of freedom $\pi \sim \bar{u}d, K \sim \bar{s}u, p \sim uud$
- Overlaps of multi-hadron *asymptotic states* \rightarrow S matrix

	$ N\pi, \text{in}\rangle$				
					
	$e^{2i\delta_{1/2,0}(s)}$	0	0	depends on $s = E_{\text{cm}}^2$ and angular variables	
	0	$e^{2i\delta_{1/2,1}(s)}$	0		diagonal in total angular momentum
	0	0	$e^{2i\delta_{3/2,1}(s)}$		$\mathcal{M}_{J,\ell}(s) \propto e^{2i\delta_{J,\ell}(s)} - 1$

- Poles on the second Riemann sheet give resonances

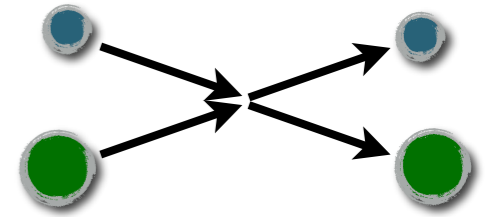
Full QCD demands this description... lattice QCD cannot escape it

Unitarity and analyticity

□ For $s < (M_N + 2M_\pi)^2$, the optical theorem tells us...

$$\rho_{N\pi}(s) |\mathcal{M}_{J,\ell}(s)|^2 = \text{Im } \mathcal{M}_{J,\ell}(s)$$

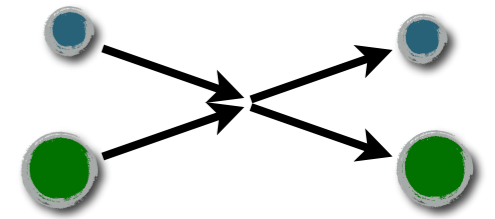
where $\rho_{N\pi}(s) = \sqrt{p^2(s, M_\pi, M_N)}/s$ is the phase space... continuum of $N\pi$ states



Unitarity and analyticity

- For $s < (M_N + 2M_\pi)^2$, the optical theorem tells us...

$$\rho_{N\pi}(s) |\mathcal{M}_{J,\ell}(s)|^2 = \text{Im } \mathcal{M}_{J,\ell}(s)$$

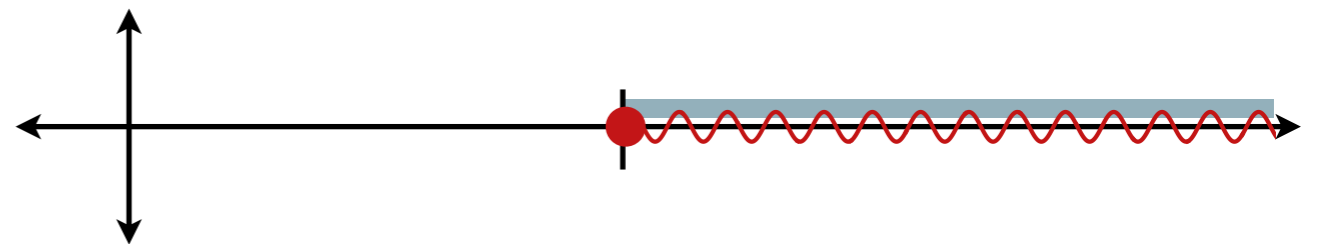


where $\rho_{N\pi}(s) = \sqrt{p^2(s, M_\pi, M_N)}/s$ is the phase space... continuum of $N\pi$ states

- Unique solution is...
$$\mathcal{M}_{J,\ell}(s) = \frac{1}{\mathcal{K}_{J,\ell}(s)^{-1} - i\rho_{N\pi}(s)}$$

K matrix (short distance)

phase-space cut (long distance)



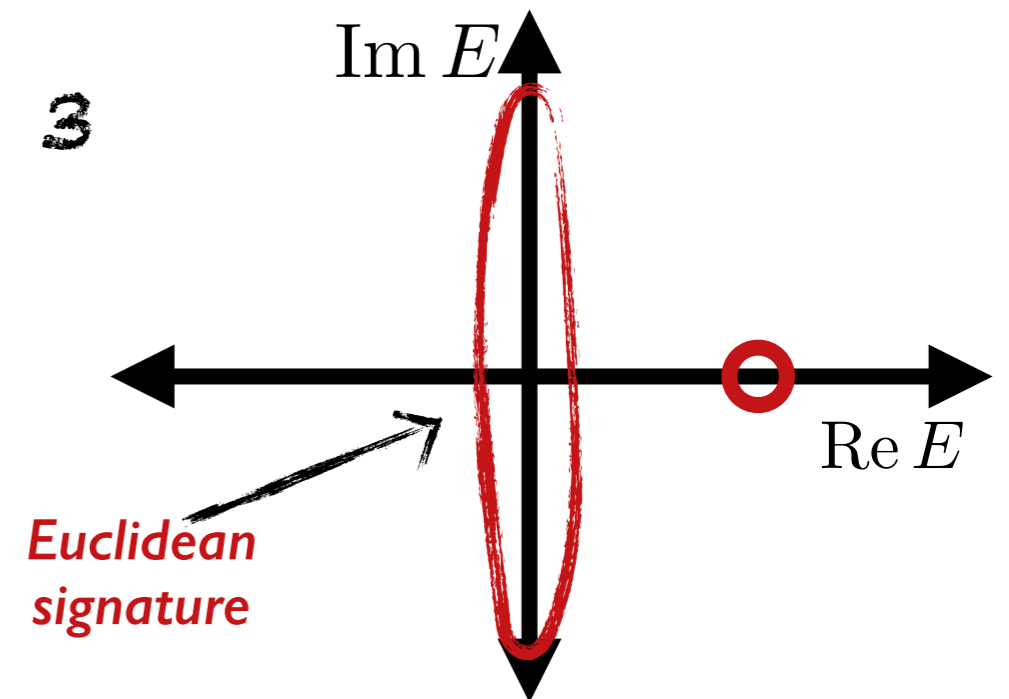
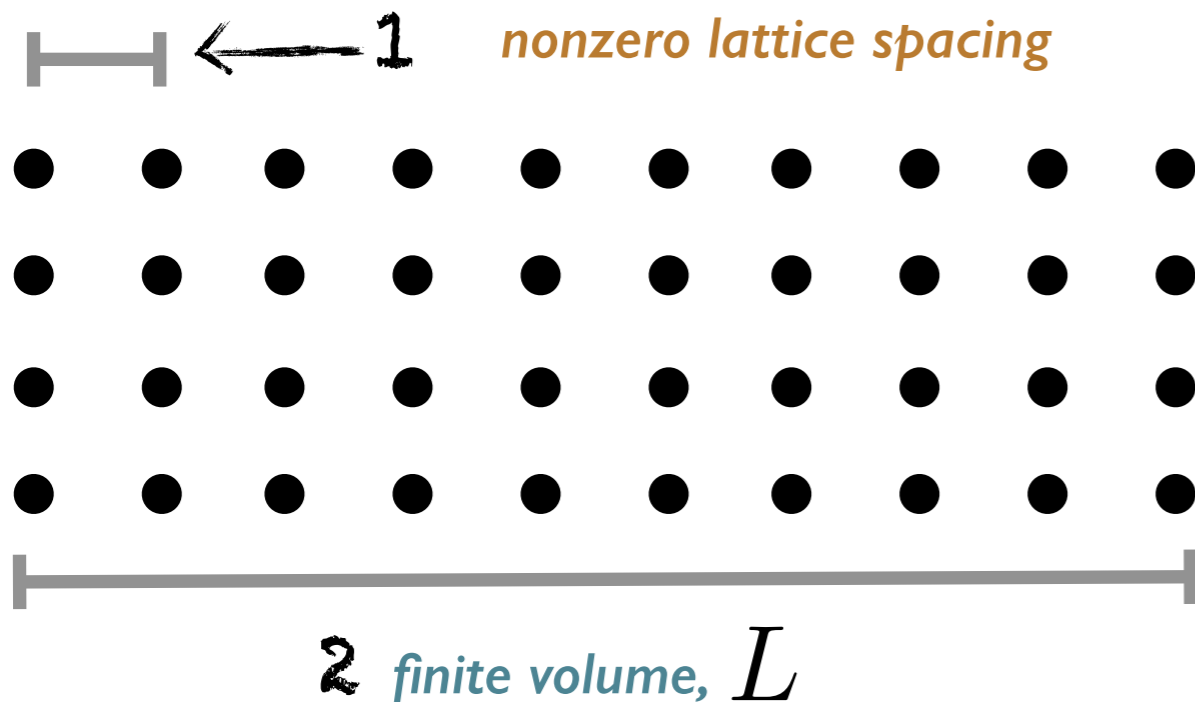
Amplitude has a branch cut ✓

K-matrix is useful for parametrizing ✓

Lattice QCD

$$\text{observable?} = \int d^N \phi e^{-S} \left[\begin{array}{l} \text{interpolator} \\ \text{for observable} \end{array} \right]$$

To proceed we have to make *three modifications*



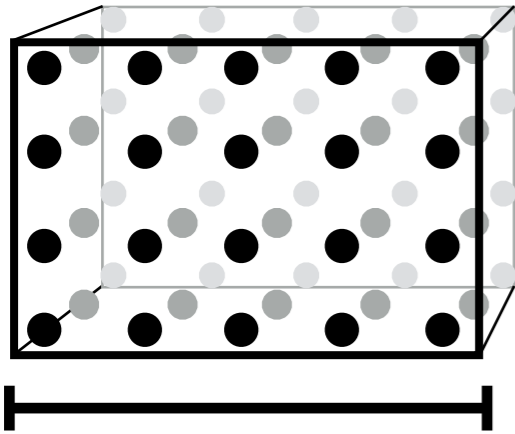
Also... $M_{\pi, \text{lattice}} > M_{\pi, \text{our universe}}$
(but physical masses \rightarrow increasingly common)



Difficulties for multi-hadron observables

□ The *Euclidean signature*...

- Prevents usual on-shell approach (want $p_4^2 = -E(p)^2$, but have only $p_4^2 > 0$)



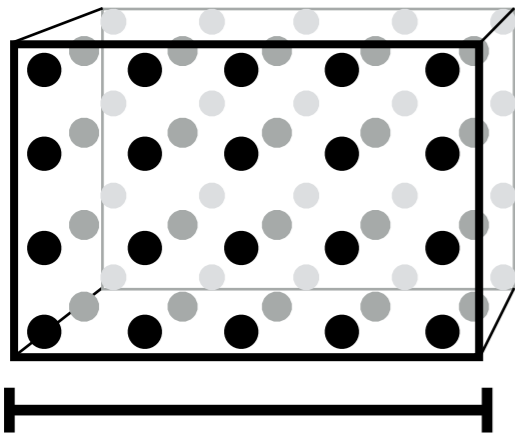
□ The *finite volume*...

- Discretizes the spectrum
- Eliminates the branch cuts and extra sheets
- Hides the resonance poles

Difficulties for multi-hadron observables

□ The *Euclidean signature*...

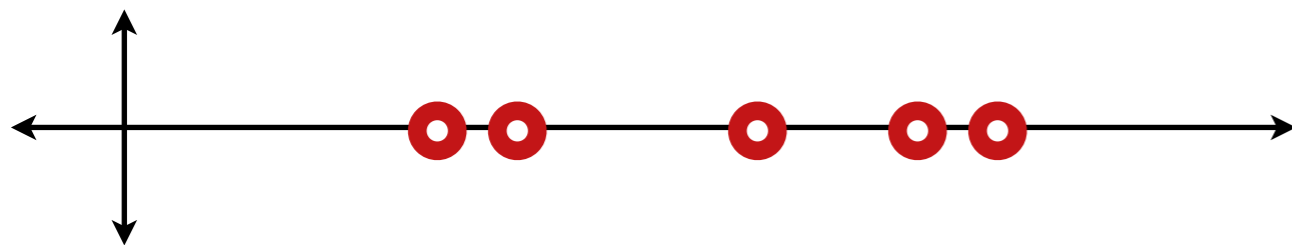
- Prevents usual on-shell approach (want $p_4^2 = -E(p)^2$, but have only $p_4^2 > 0$)



□ The *finite volume*...

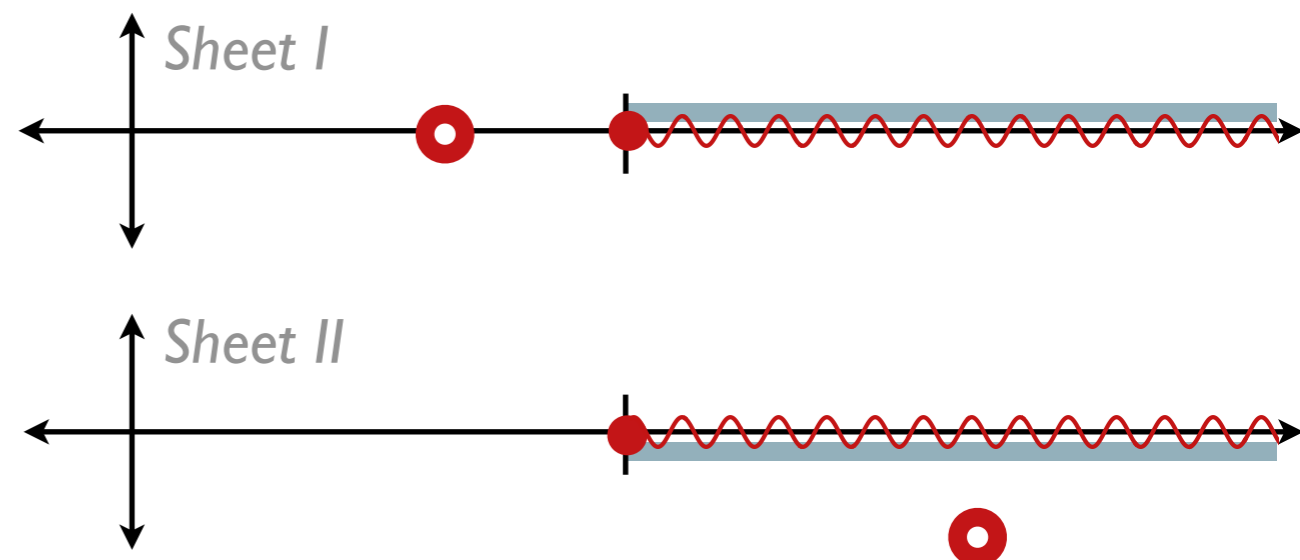
- Discretizes the spectrum
- Eliminates the branch cuts and extra sheets
- Hides the resonance poles

Finite-volume analytic structure



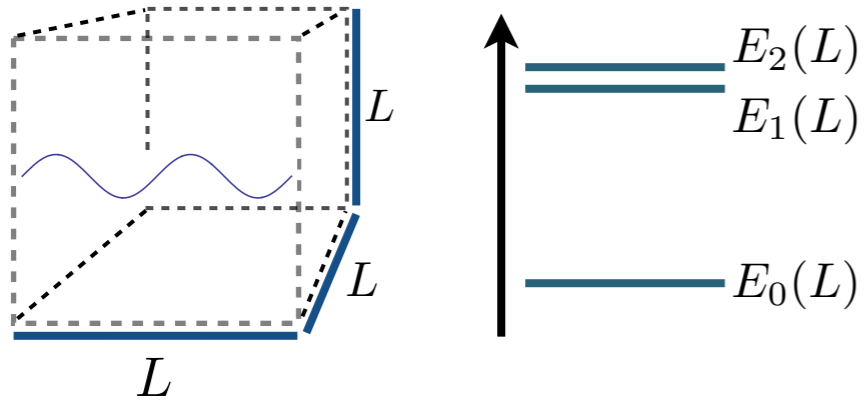
Note: cannot count finite-volume energies to count resonance poles!

Infinite-volume analytic structure



The finite-volume as a tool

□ Finite-volume set-up



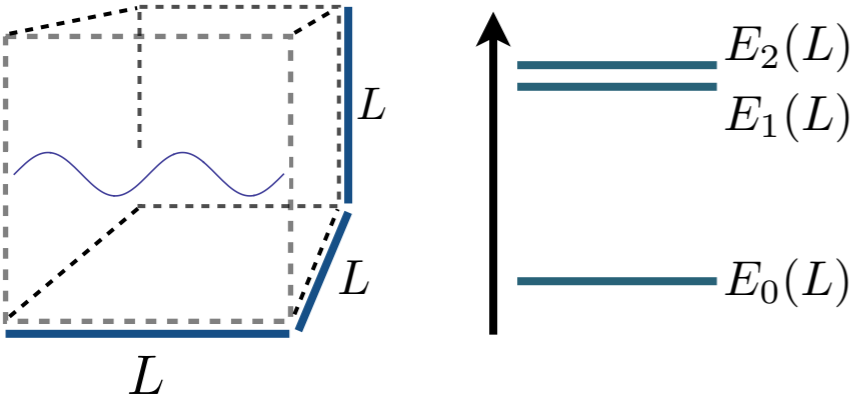
□ **cubic**, spatial volume (extent L)

□ **periodic**

□ L is large enough to neglect $e^{-M_\pi L}$

The finite-volume as a tool

□ Finite-volume set-up

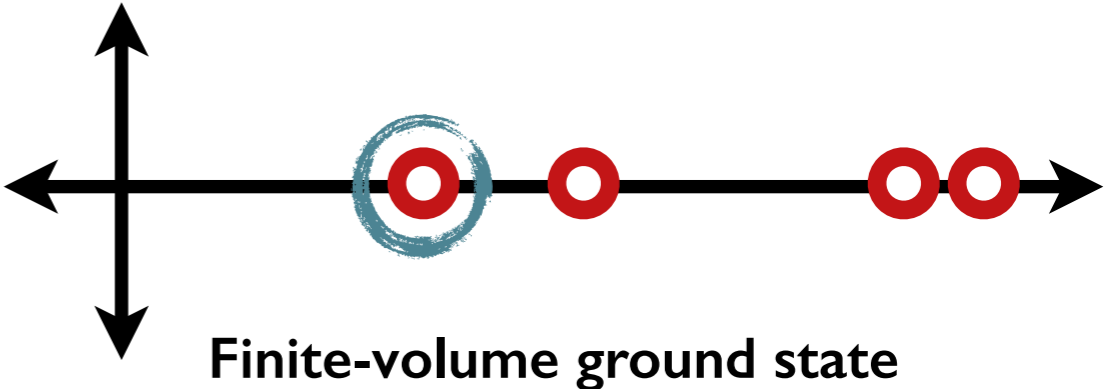
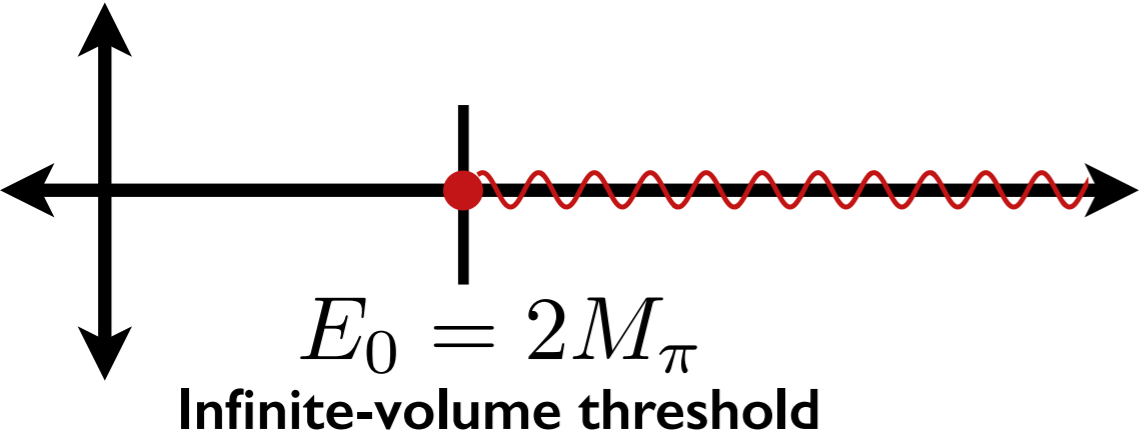


□ **cubic**, spatial volume (extent L)

□ **periodic**

□ L is large enough to neglect $e^{-M_\pi L}$

□ Scattering leaves an *imprint* on finite-volume quantities



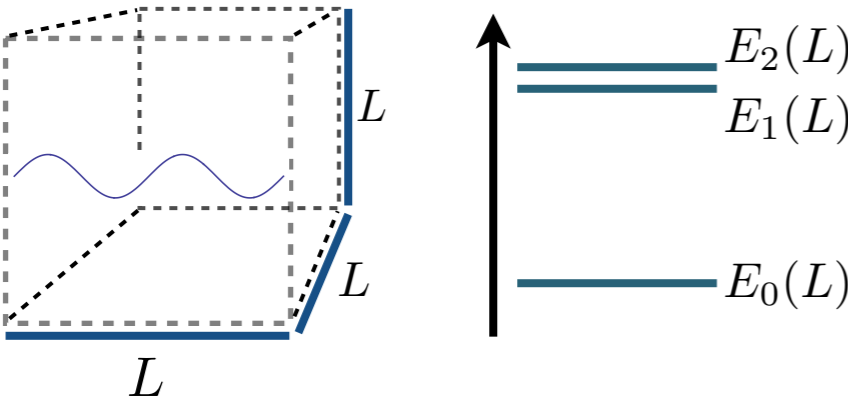
$$\mathcal{M}_{\ell=0}(2M_\pi) = -32\pi M_\pi a$$

$$E_0(L) = 2M_\pi + \frac{4\pi a}{M_\pi L^3} + \mathcal{O}(1/L^4)$$

• Huang, Yang (1958) •

The finite-volume as a tool

□ Finite-volume set-up

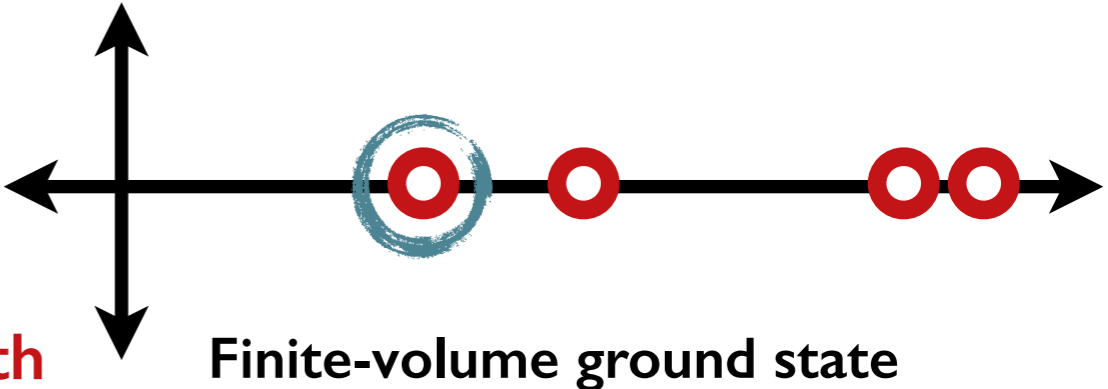
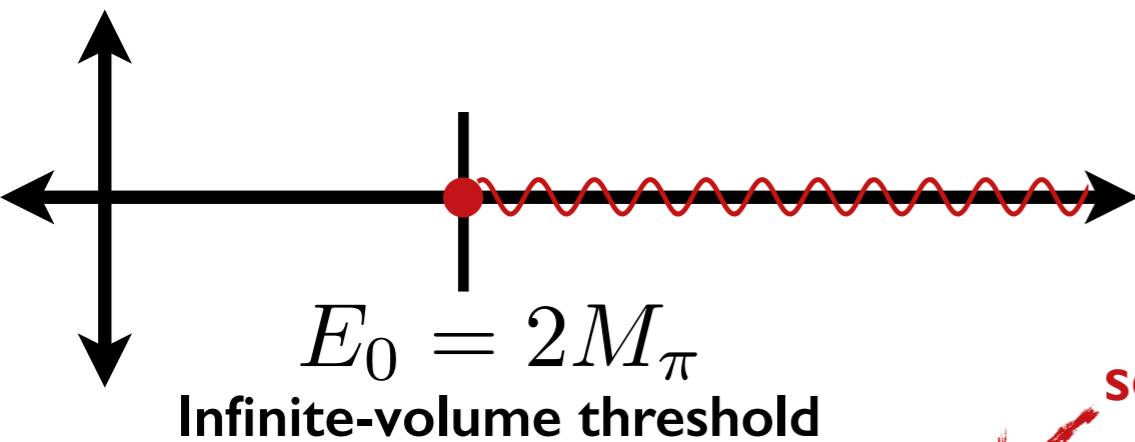


□ **cubic**, spatial volume (extent L)

□ **periodic**

□ L is large enough to neglect $e^{-M_\pi L}$

□ Scattering leaves an *imprint* on finite-volume quantities



$$\mathcal{M}_{\ell=0}(2M_\pi) = -32\pi M_\pi a$$

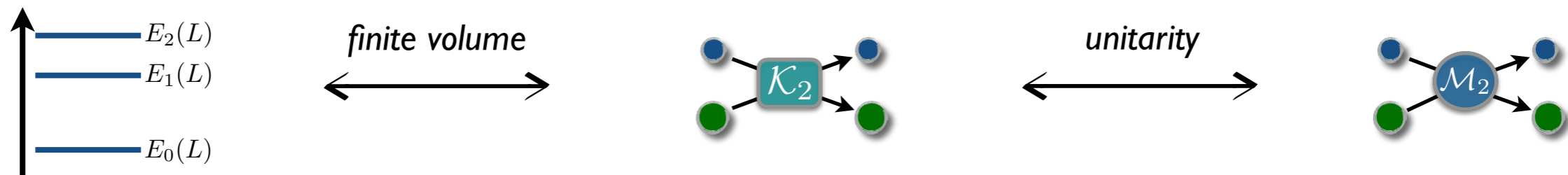
$$E_0(L) = 2M_\pi + \frac{4\pi a}{M_\pi L^3} + \mathcal{O}(1/L^4)$$

• Huang, Yang (1958) •

General method

$$\det[\mathcal{K}^{-1}(s) + F(P, L)] = 0$$

$F(P, L) \equiv$ Matrix of known geometric functions



Holds only for two-particle energies $s < (M_N + 2M_\pi)^2$ Neglects $e^{-M_\pi L}$

Generalized to *non-degenerate masses, multiple channels, spinning particles*

Encodes angular momentum mixing

Huang, Yang (1958) • Lüscher (1986, 1991) • Rummukainen, Gottlieb (1995)

Kim, Sachrajda, Sharpe (2005) • Christ, Kim, Yamazaki (2005) • He, Feng, Liu (2005)

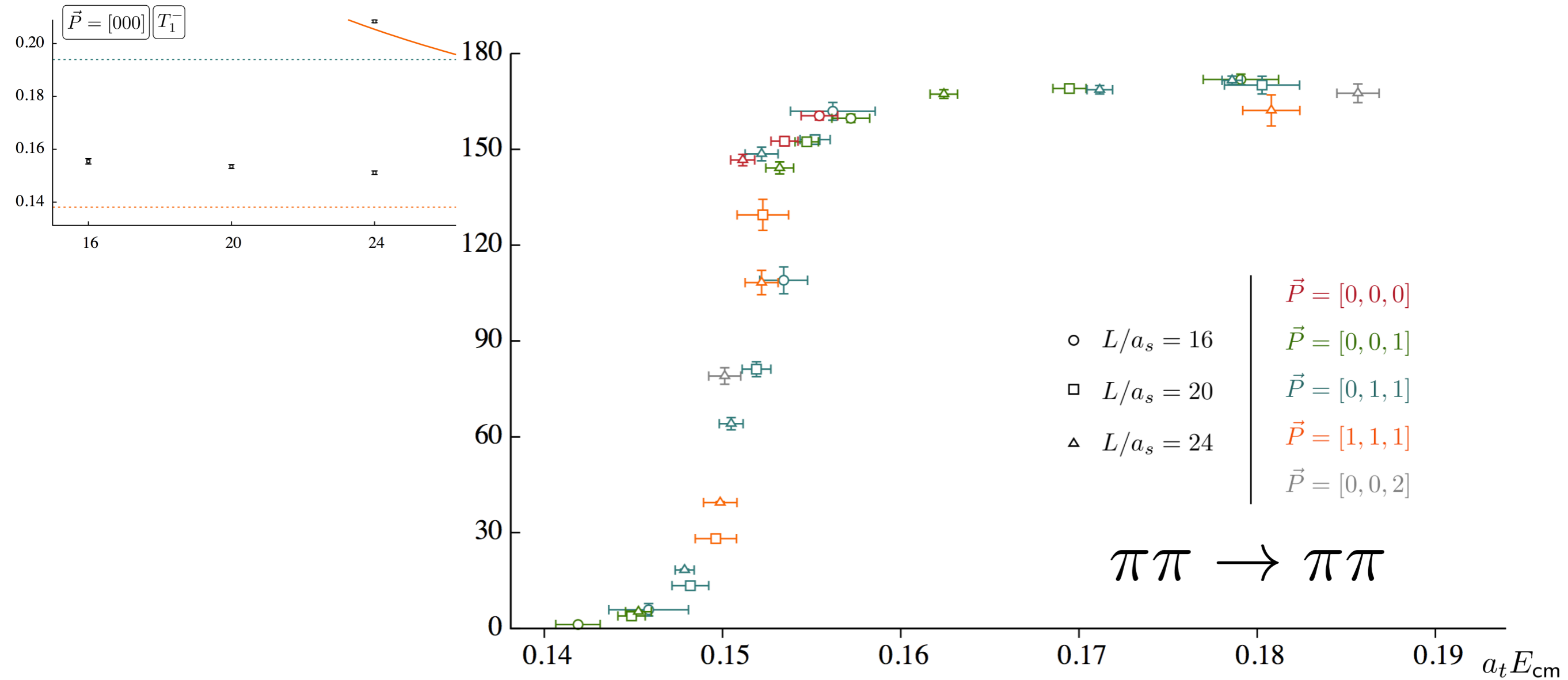
Leskovec, Prelovsek (2012) • Bernard *et. al.* (2012) • MTH, Sharpe (2012) • Briceño, Davoudi (2012)

Li, Liu (2013) • Briceño (2014)

Using the result

□ Single-channel case (*pions in a p-wave*)

$$\mathcal{K}(s_n)^{-1} = \rho \cot \delta(s_n) = -F(E_n, \vec{P}, L)$$

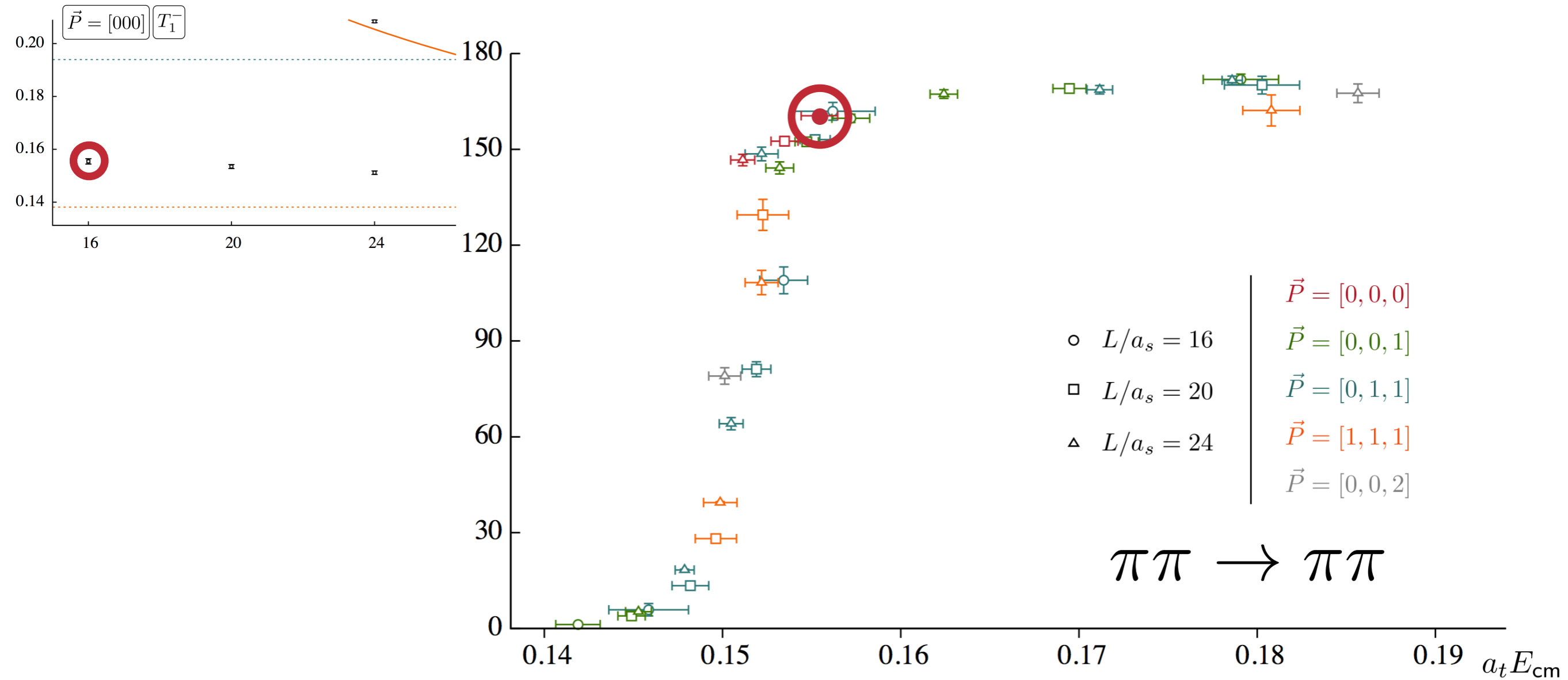


- Dudek, Edwards, Thomas in *Phys.Rev.* D87 (2013) 034505 •

Using the result

□ Single-channel case (*pions in a p-wave*)

$$\mathcal{K}(s_n)^{-1} = \rho \cot \delta(s_n) = -F(E_n, \vec{P}, L)$$

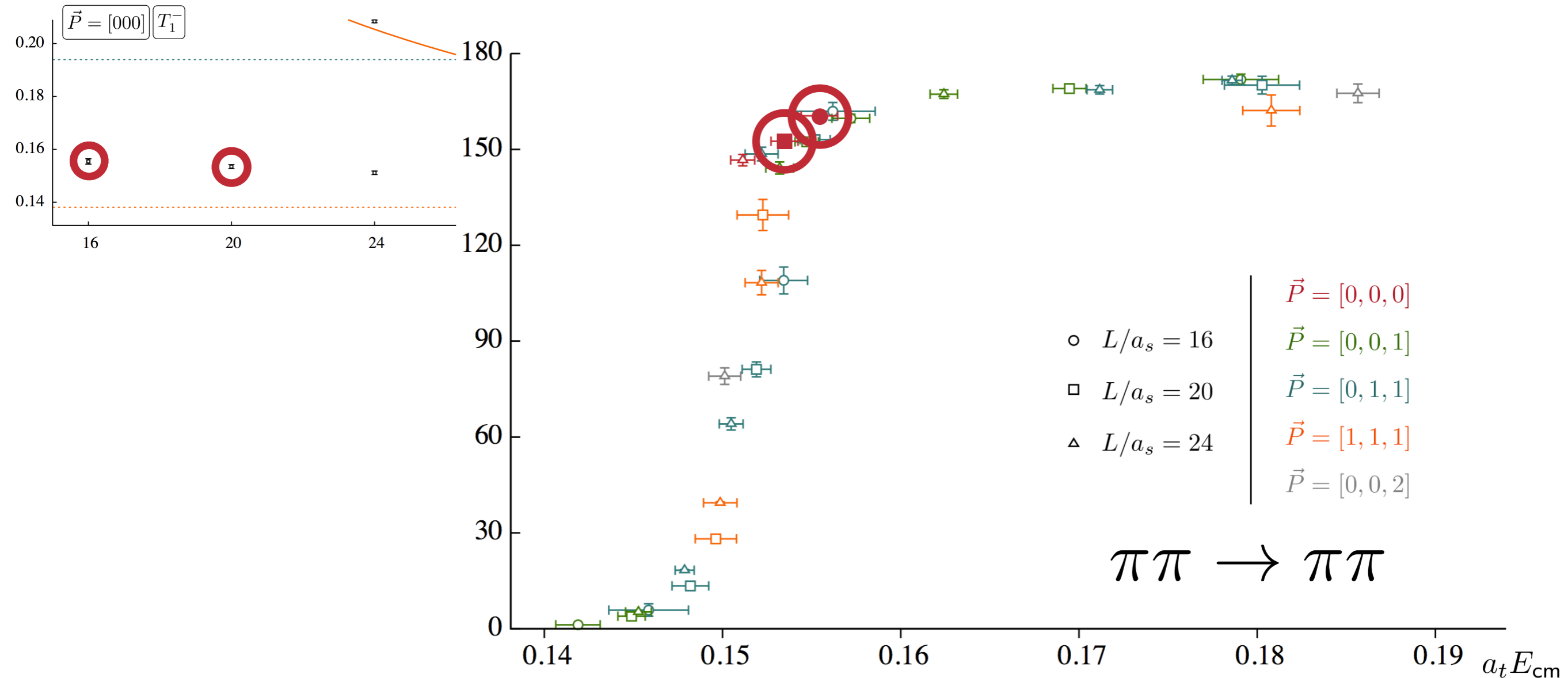


- Dudek, Edwards, Thomas in *Phys.Rev.* D87 (2013) 034505 •

Using the result

□ Single-channel case (*pions in a p-wave*)

$$\mathcal{K}(s_n)^{-1} = \rho \cot \delta(s_n) = -F(E_n, \vec{P}, L)$$

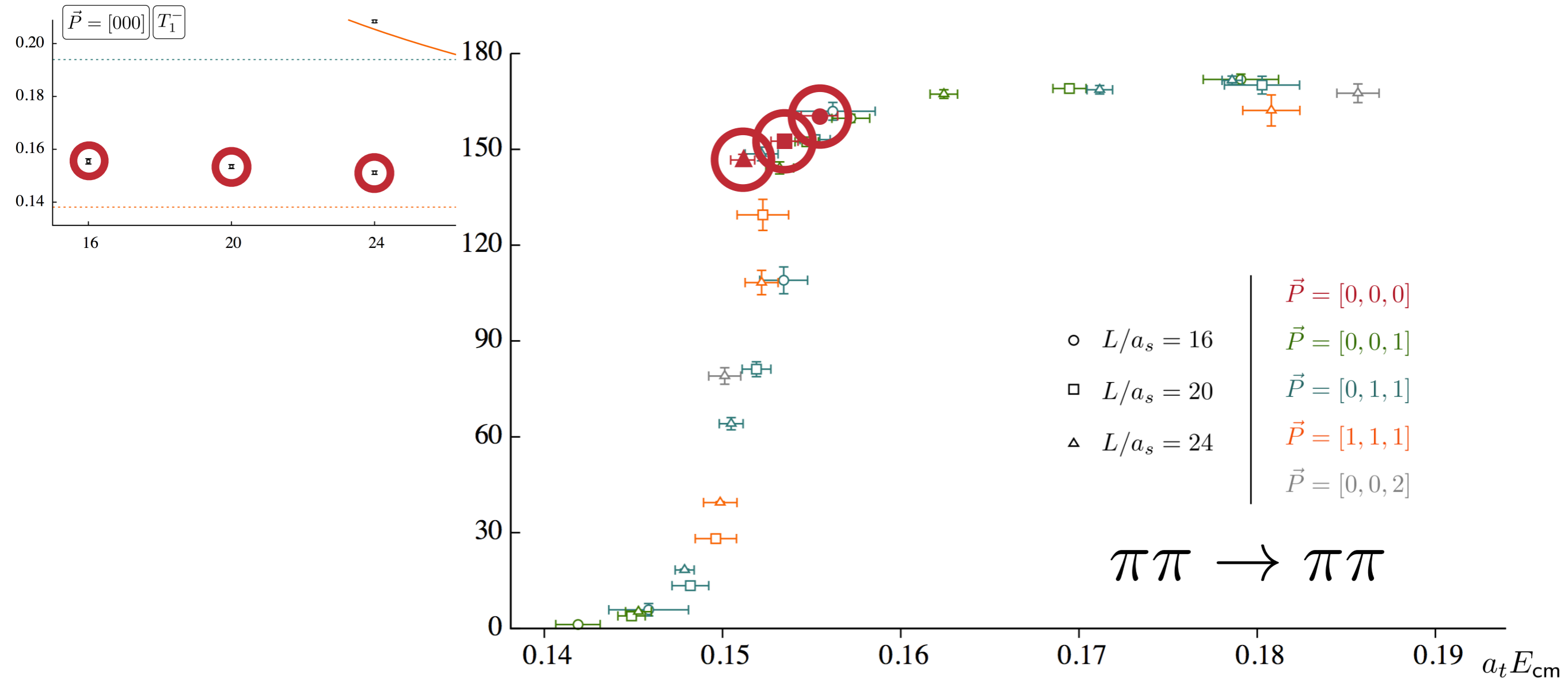


- Dudek, Edwards, Thomas in *Phys.Rev.* D87 (2013) 034505 •

Using the result

□ Single-channel case (*pions in a p-wave*)

$$\mathcal{K}(s_n)^{-1} = \rho \cot \delta(s_n) = -F(E_n, \vec{P}, L)$$

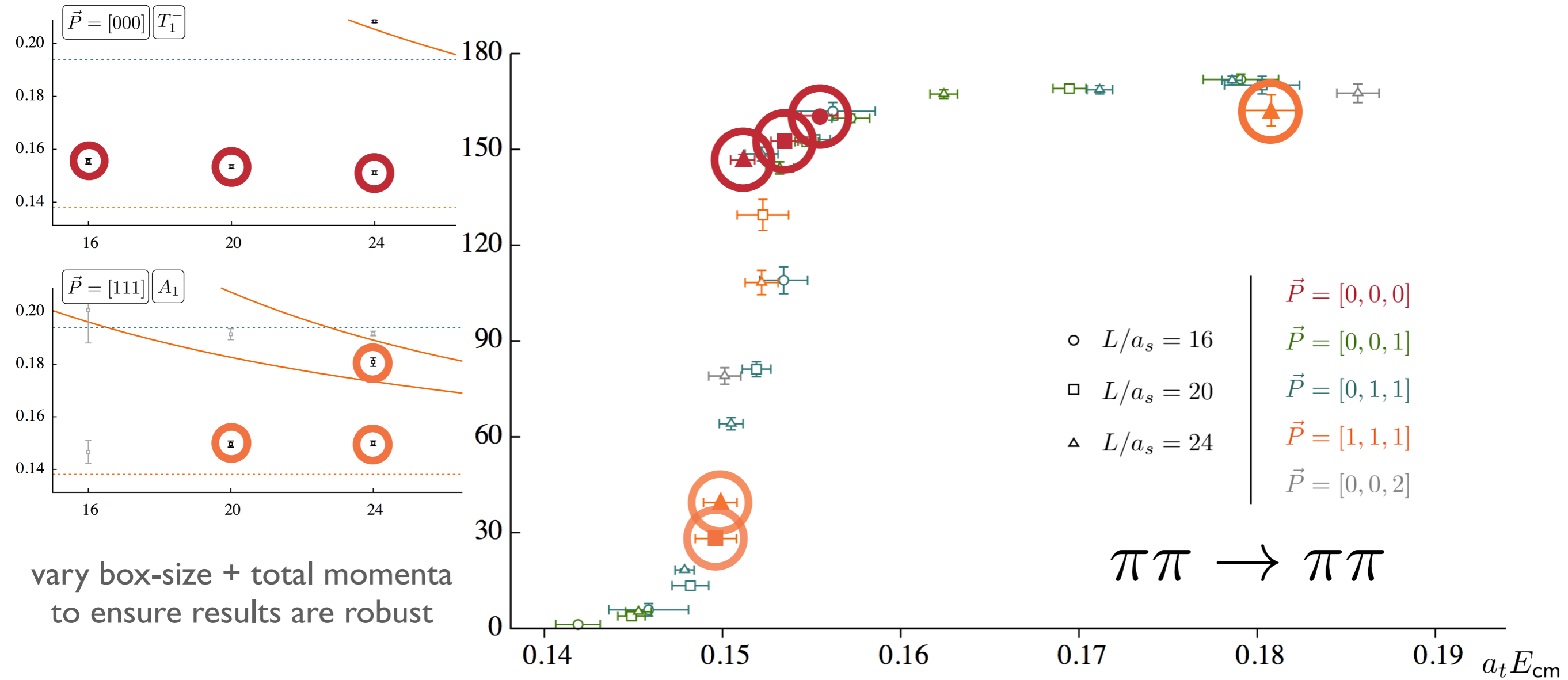


- Dudek, Edwards, Thomas in *Phys.Rev.* D87 (2013) 034505 •

Using the result

□ Single-channel case (*pions in a p-wave*)

$$\mathcal{K}(s_n)^{-1} = \rho \cot \delta(s_n) = -F(E_n, \vec{P}, L)$$

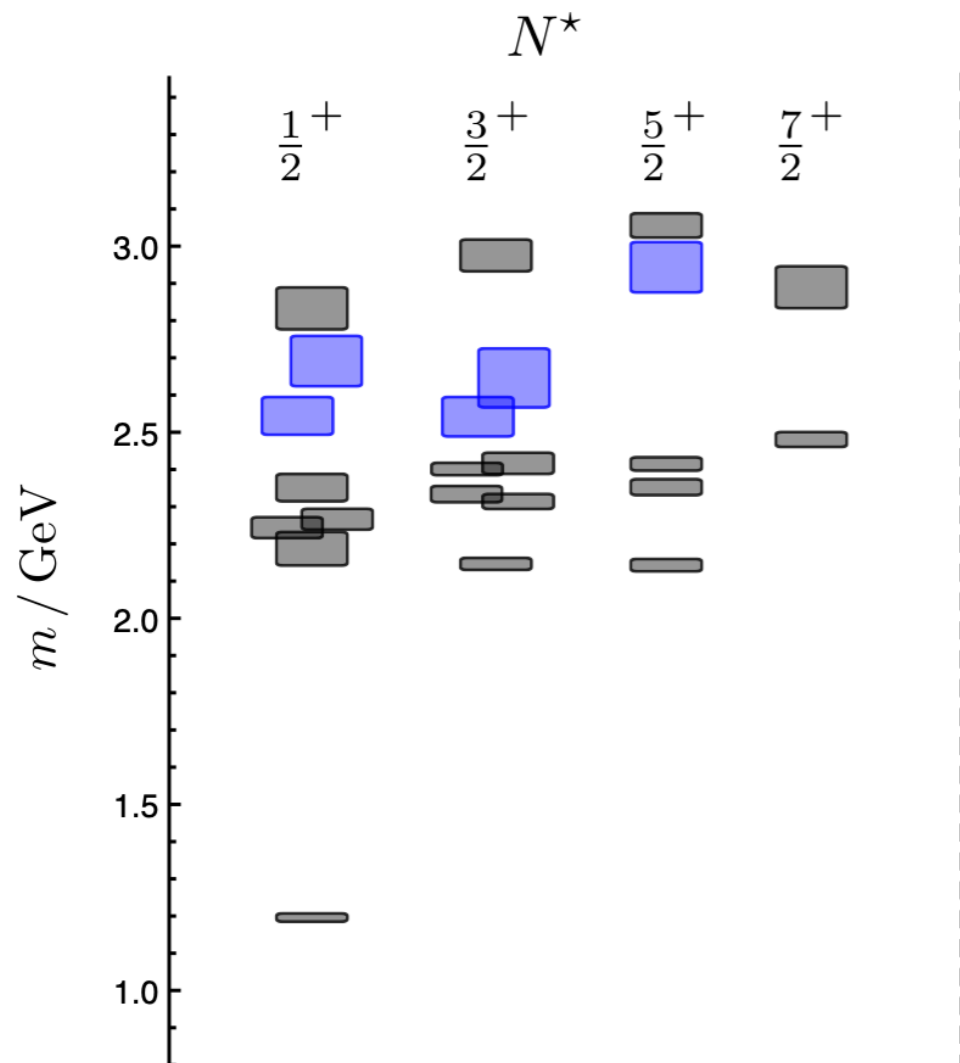


- Dudek, Edwards, Thomas in *Phys.Rev.* D87 (2013) 034505 •

Two types of spectroscopy

Explore the spectrum of compact QCD
excited states

(via quark-model inspired local operators)

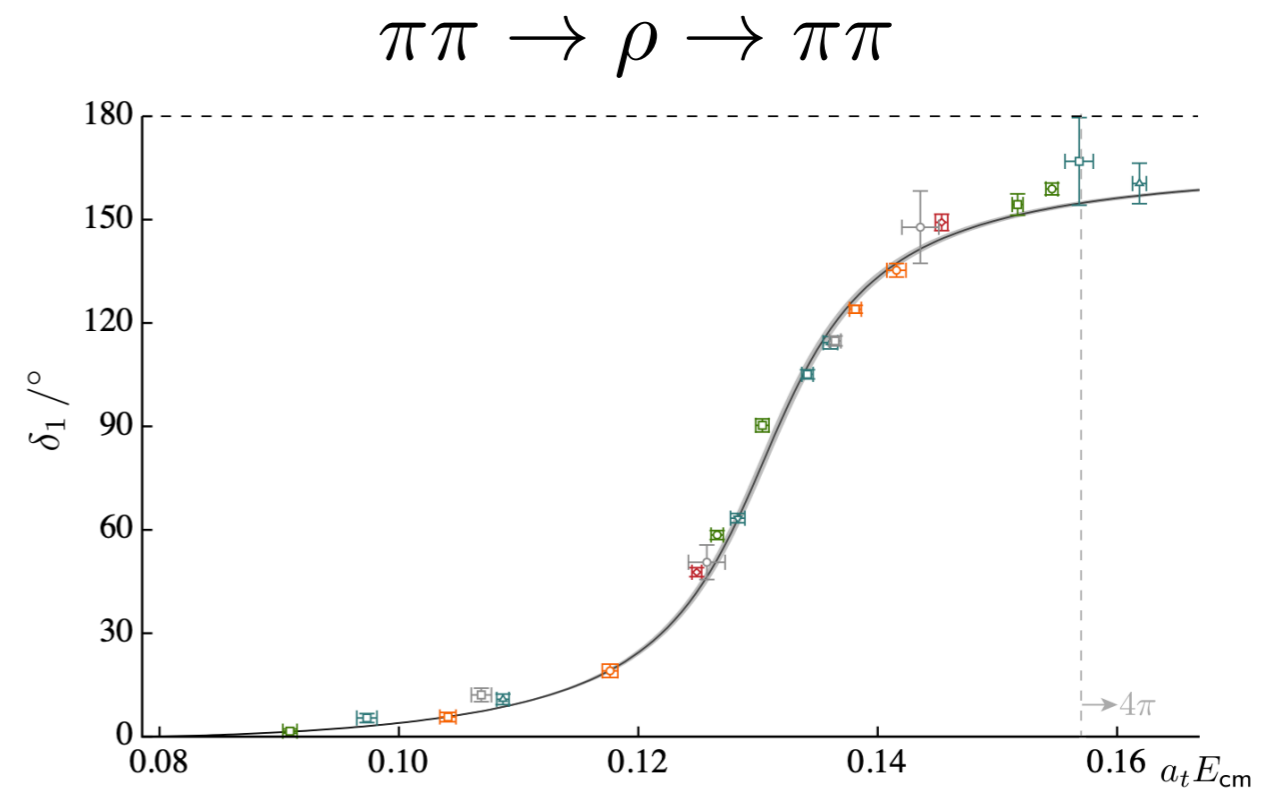


Dudek, Edwards (2012)

Extract the full finite-volume
energy spectrum

local operators

+ many multi-hadron operators

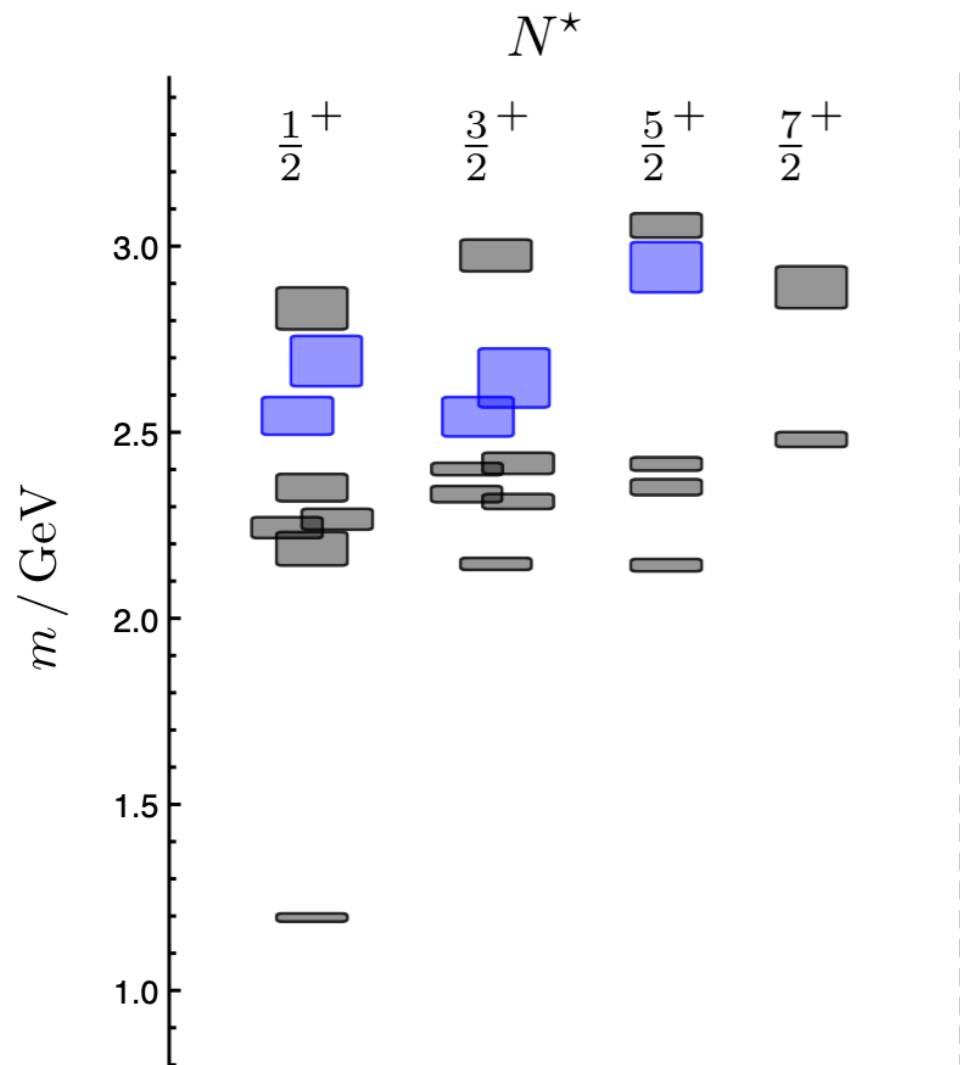


Wilson, Briceño, Dudek, Edwards, Thomas (2015)

Two types of spectroscopy

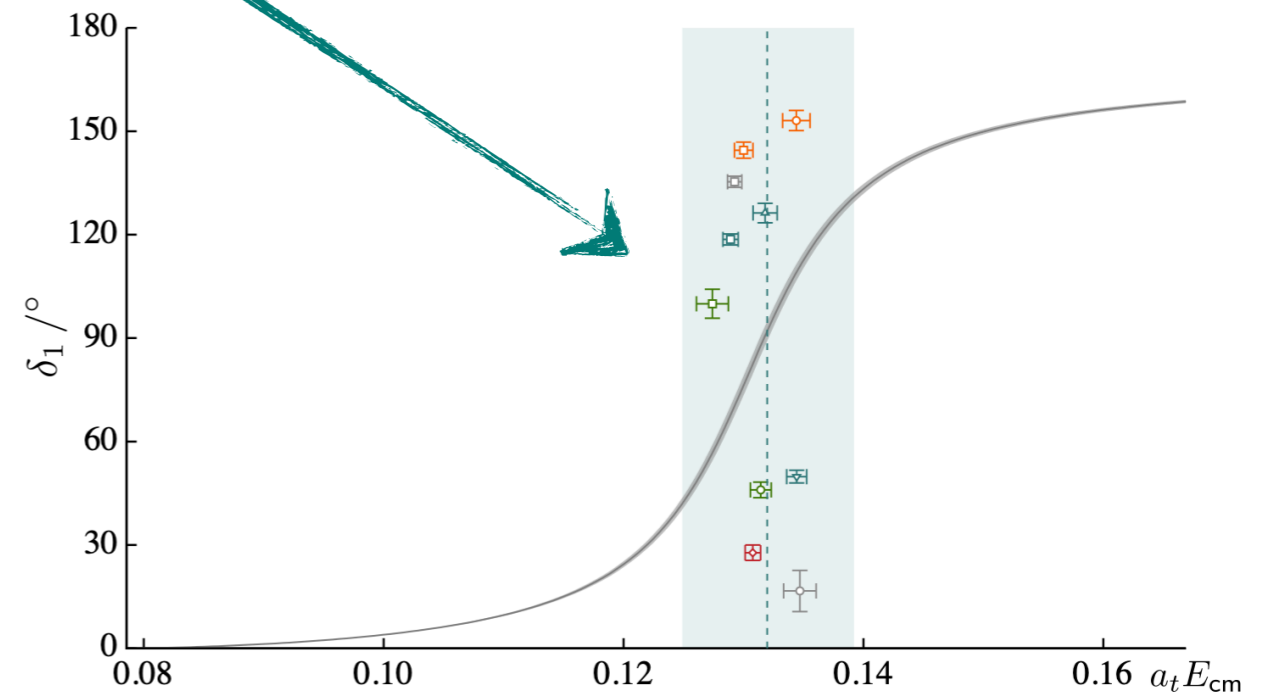
Explore the spectrum of compact QCD
excited states

(via quark-model inspired local operators)



Dudek, Edwards (2012)

local operator spectrum =
not suitable for phase shift extraction



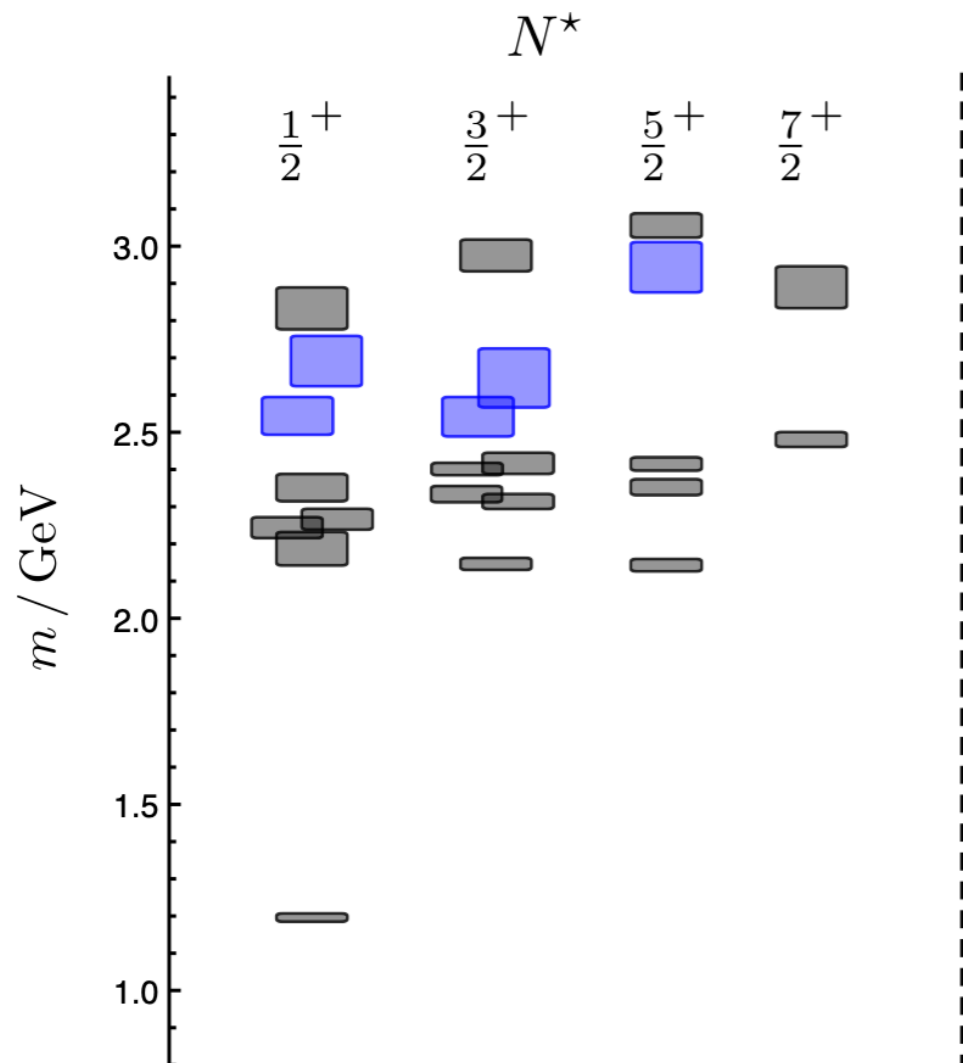
Note: cannot count finite-volume
energies to count resonance poles!

Wilson, Briceño, Dudek, Edwards, Thomas (2015)

Two types of spectroscopy

Explore the spectrum of compact QCD
excited states

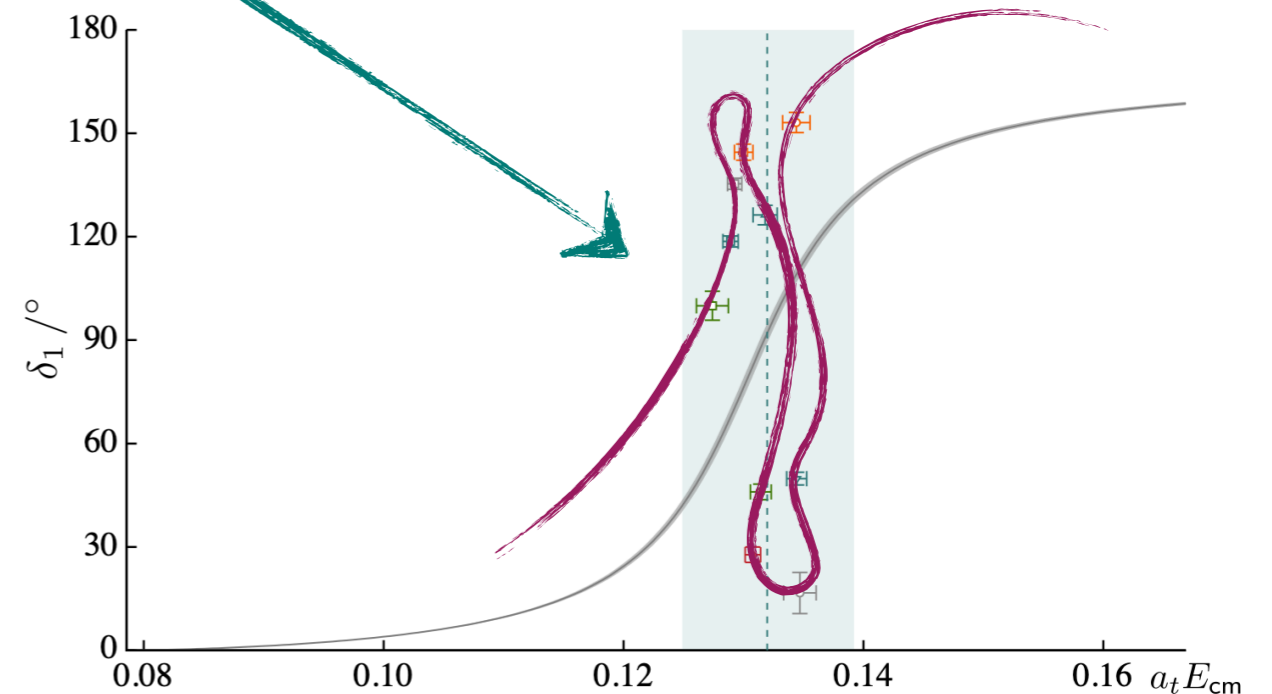
(via quark-model inspired local operators)



Dudek, Edwards (2012)

local operator spectrum =
not suitable for phase shift extraction

$$\pi\pi \rightarrow \rho \rightarrow \pi\pi$$

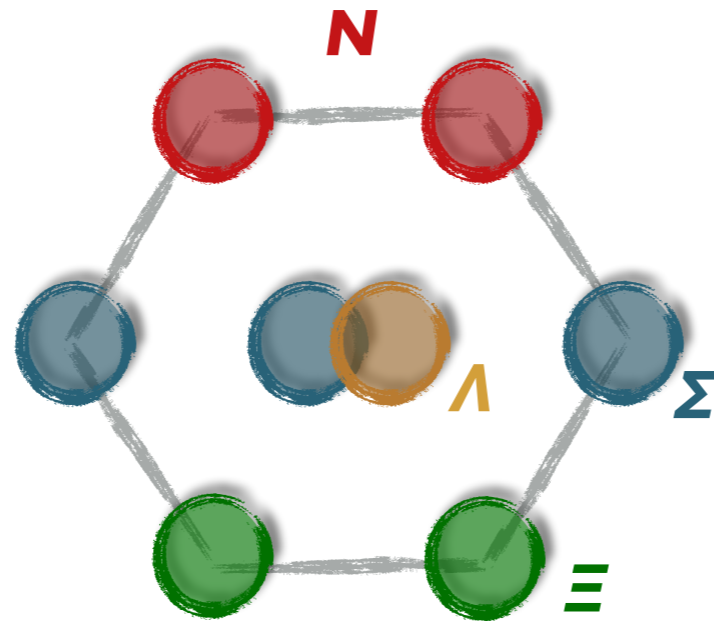


Note: cannot count finite-volume
energies to count resonance poles!

Wilson, Briceño, Dudek, Edwards, Thomas (2015)

$$\Delta \rightarrow N\pi$$

- Andersen et al. 2018
- Andersen et al. 2019
- Silvi et al. 2021
- Pittler et al. 2021
- Bulava et al. 2022

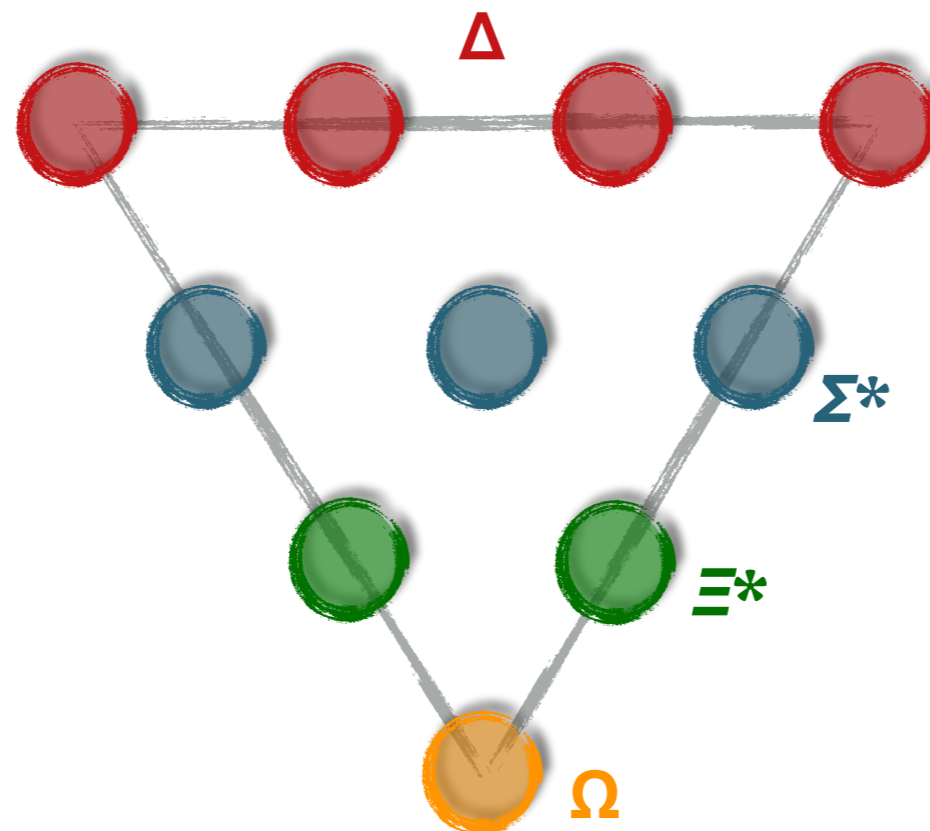


(focusing here on studies with scattering states)

Baryons are difficult!

$$N^* \rightarrow N\pi$$

- Lang et al. 2017
- Wu et al. 2017
- Kiratidis et al. 2017



$$\Lambda \rightarrow \bar{K}N$$

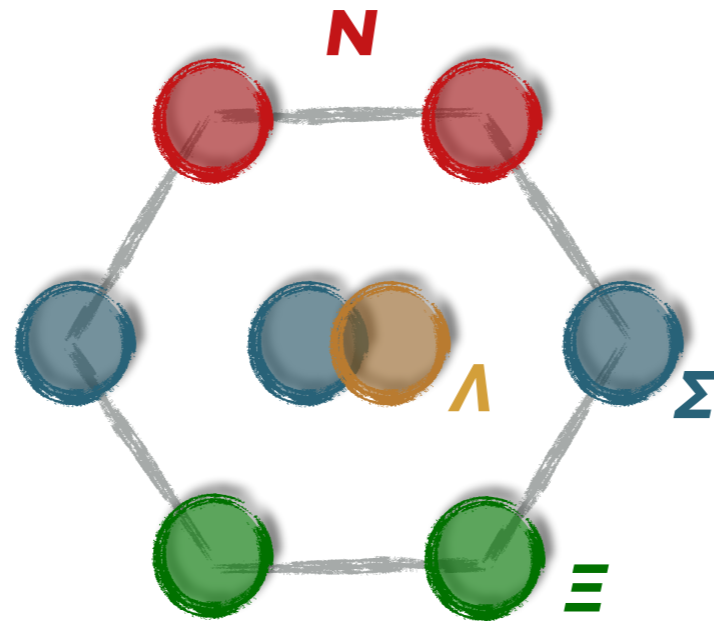
- Hall et al. 2015

See also...

- Detmold and Nicholson 2015
- Wu et al. 2018
- Xing & Liu, LATT2022 (in prep)

$$\Delta \rightarrow N\pi$$

- Andersen et al. 2018
- Andersen et al. 2019
- Silvi et al. 2021
- Pittler et al. 2021
- Bulava et al. 2022

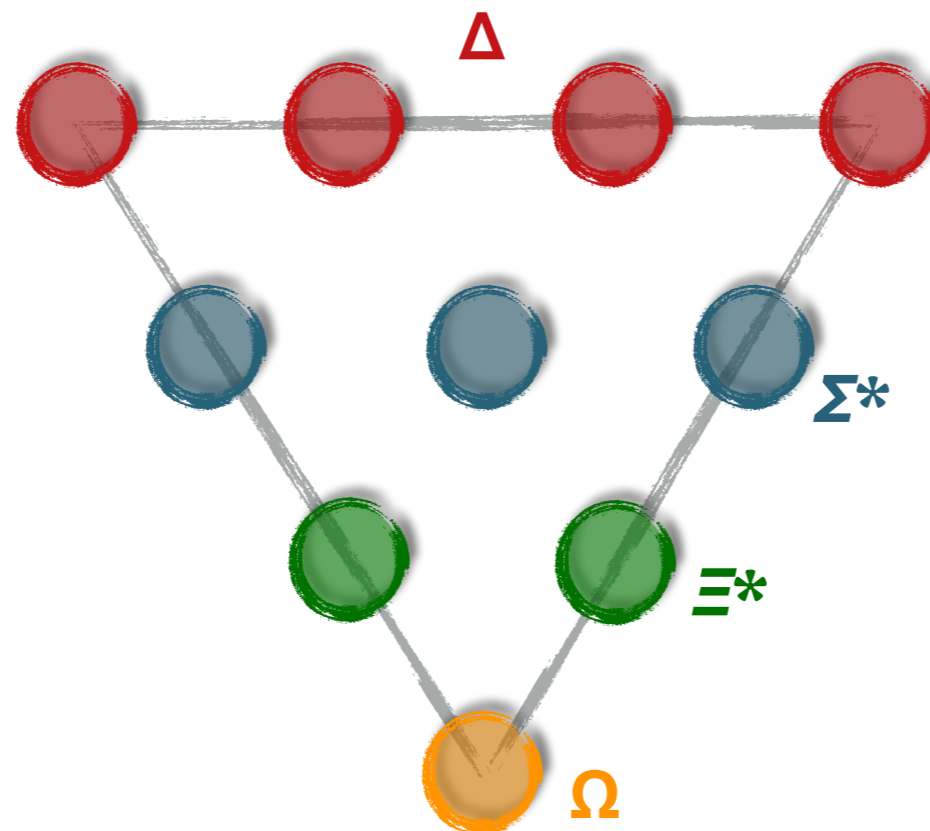


(focusing here on studies with scattering states)

Baryons are difficult!

$$N^* \rightarrow N\pi$$

- Lang et al. 2017
- Wu et al. 2017
- Kiratidis et al. 2017



$$\Lambda \rightarrow \bar{K}N$$

- Hall et al. 2015

See also...

- Detmold and Nicholson 2015
- Wu et al. 2018
- Xing & Liu, LATT2022 (in prep)

$N\pi$ elastic scattering ($M_\pi = 200$ MeV)

arXiv > hep-lat > arXiv:2208.03867

High Energy Physics - Lattice

[Submitted on 8 Aug 2022]

Elastic nucleon-pion scattering at $m_\pi \approx 200$ MeV from lattice QCD

John Bulava, Andrew D. Hanlon, Ben Hörz, Colin Morningstar, Amy Nicholson, Fernando Romero-López, Sarah Skinner, Pavlos Vranas, André Walker-Loud

□ Studied scattering-lengths and the Δ channel

$$I = 1/2 : J^P = 1/2^- (S)$$

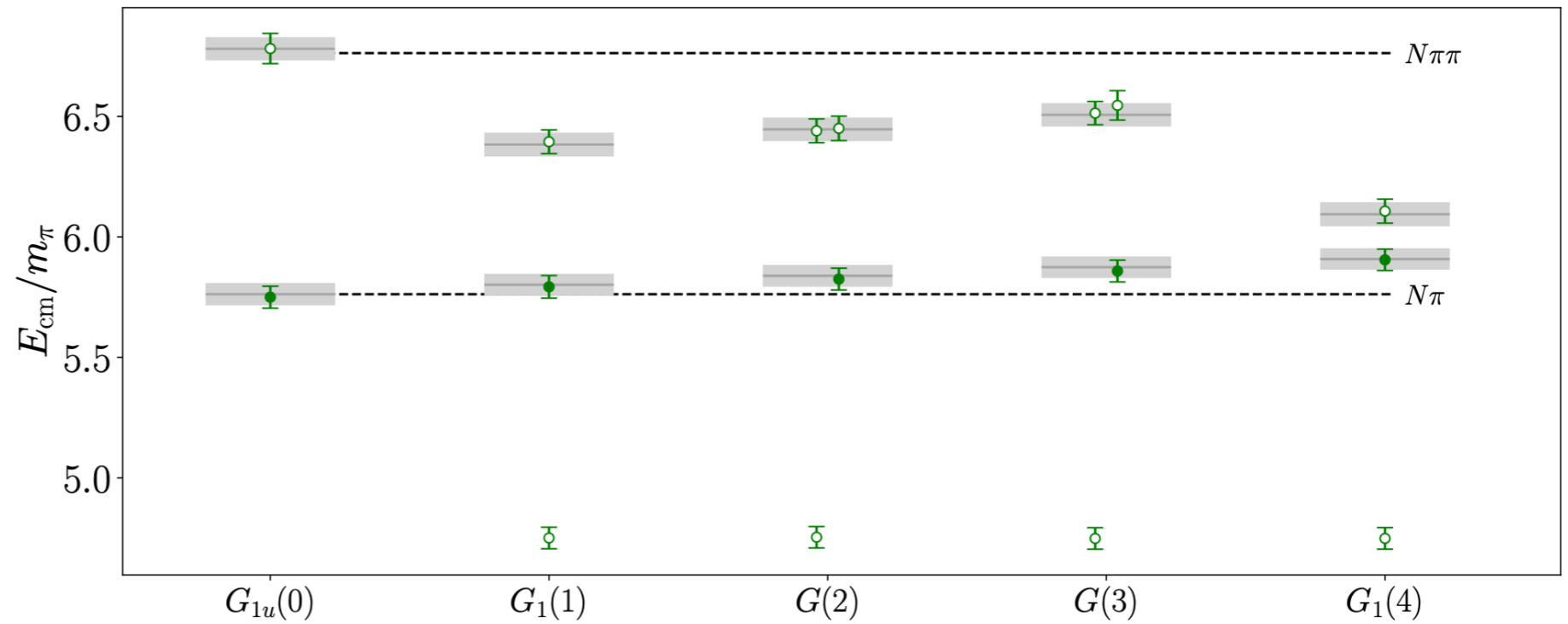
$$I = 3/2 : J^P = 1/2^- (S), 3/2^+ (P) \quad [1/2^+ (P), 3/2^- (D), 5/2^- (D)]$$

□ Advanced operators + fits to extract finite-volume energies

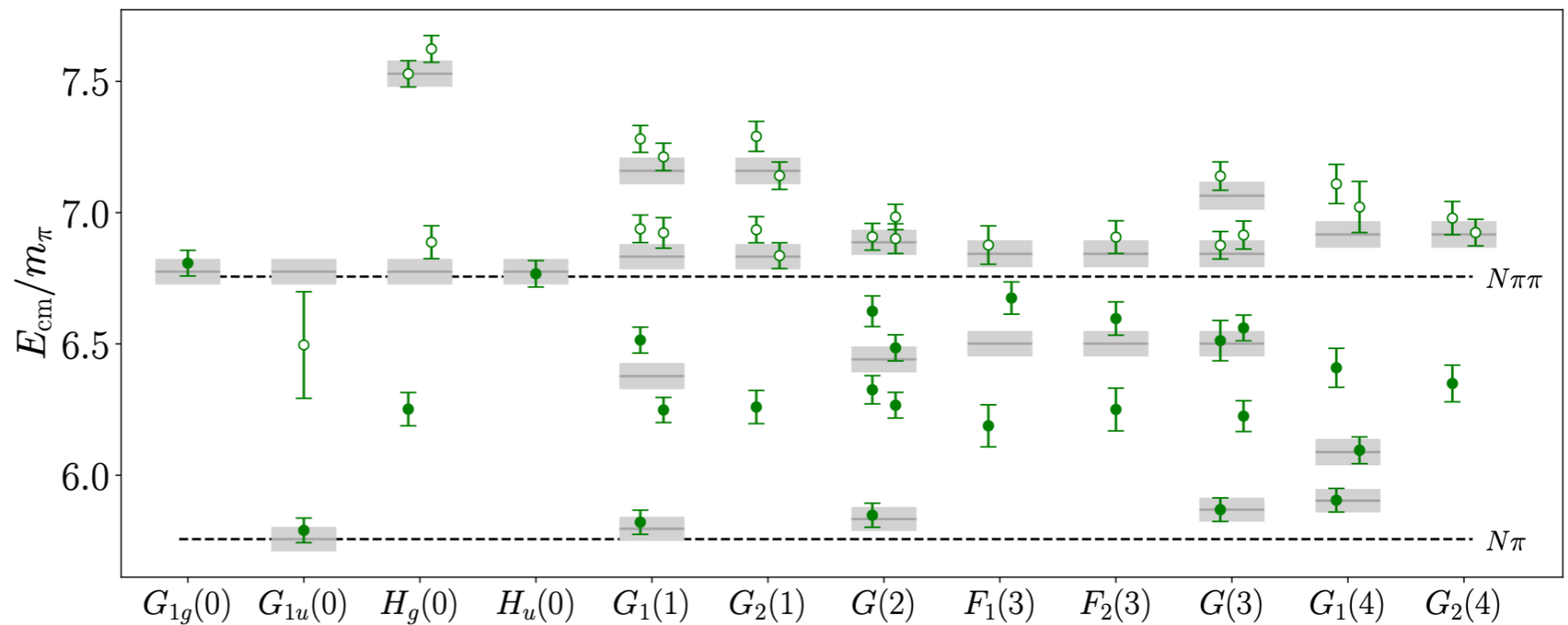
• Bulava et al. (2022) 2208.03867 •

$N\pi$ finite-volume energies ($M_\pi = 200$ MeV)

$$I = 1/2$$



$$I = 3/2$$

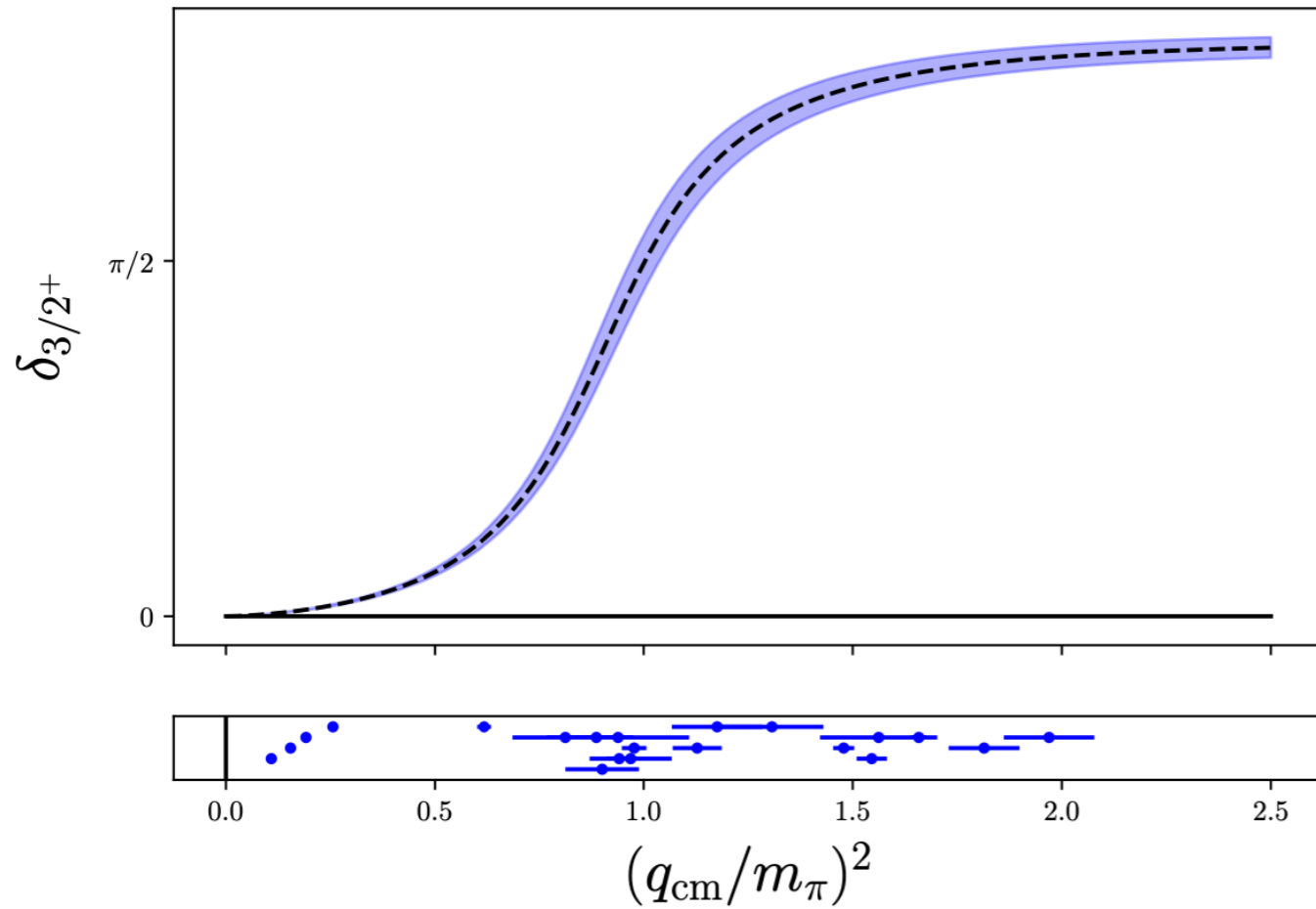


• Bulava et al. (2022) 2208.03867 •

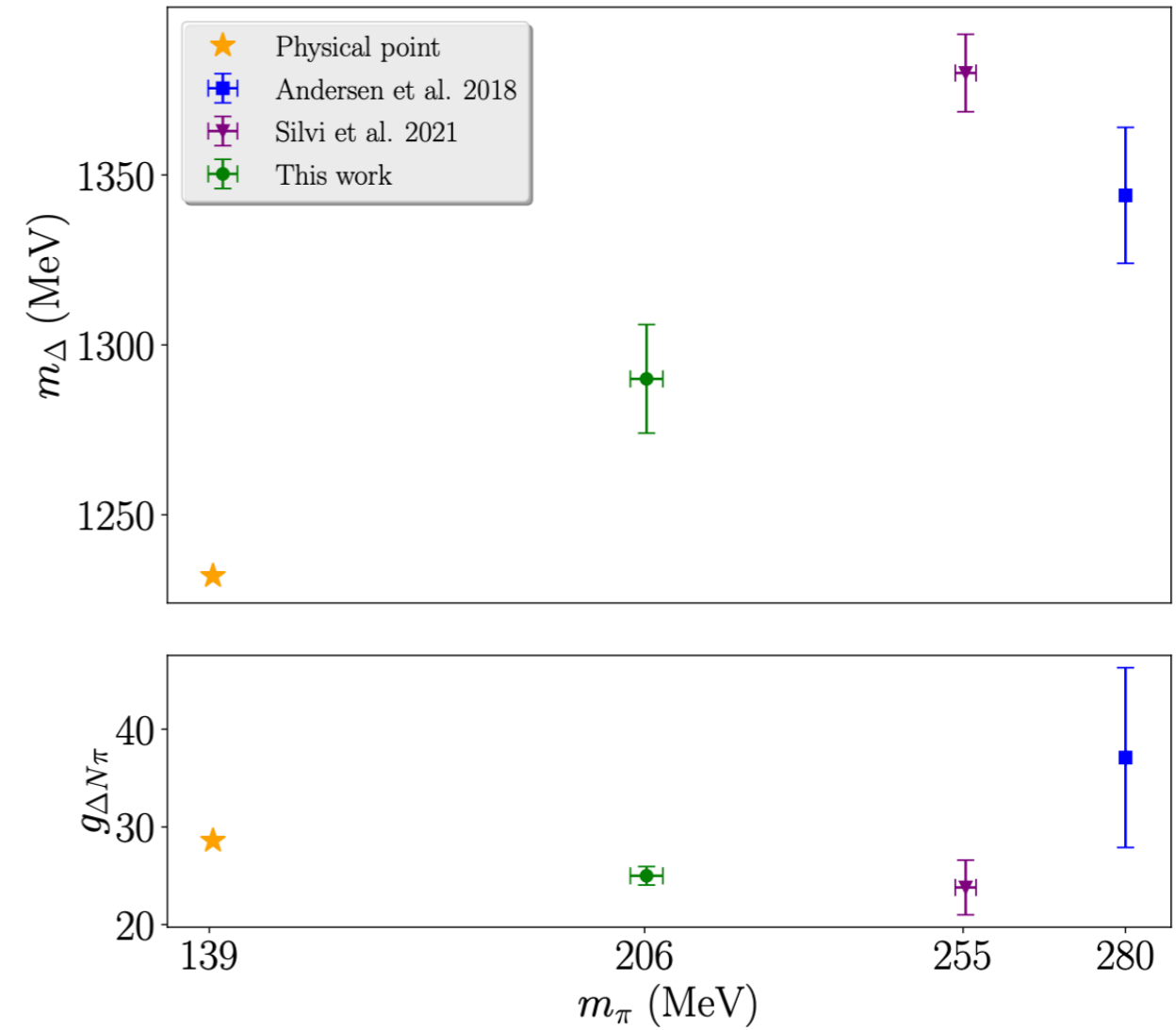


$$N\pi \rightarrow \Delta \rightarrow N\pi$$

Scattering phase shift



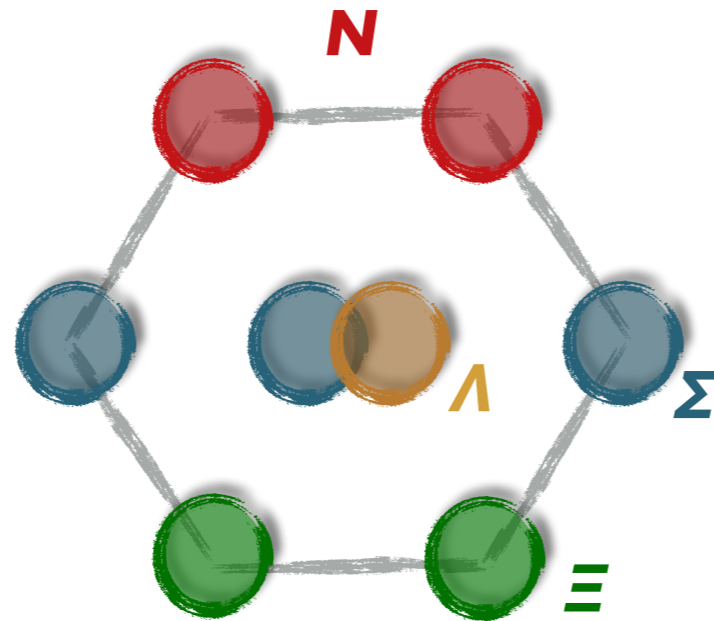
Δ summary plot



• Bulava et al. (2022) 2208.03867 •

$$\Delta \rightarrow N\pi$$

- Andersen et al. 2018
- Andersen et al. 2019
- Silvi et al. 2021
- Pittler et al. 2021
- Bulava et al. 2022

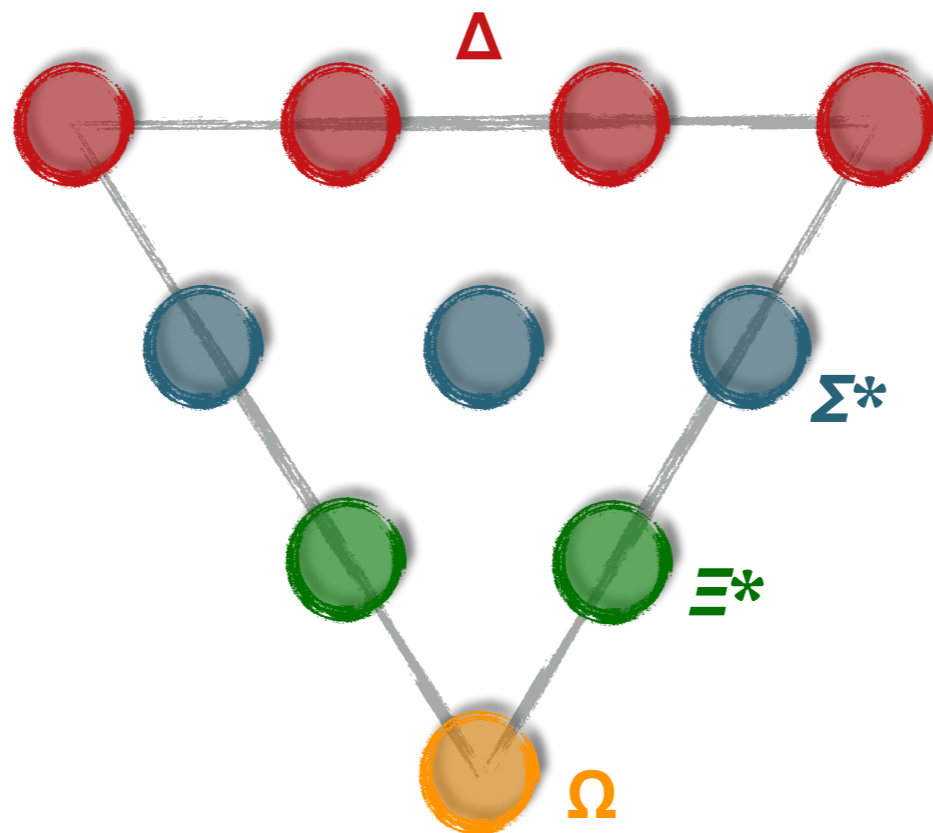


(focusing here on studies with scattering states)

Baryons are difficult!

$$N^* \rightarrow N\pi$$

- Lang et al. 2017
- Wu et al. 2017
- Kiratidis et al. 2017



$$\Lambda \rightarrow \bar{K}N$$

- Hall et al. 2015

See also...

- Detmold and Nicholson 2015
- Wu et al. 2018
- Xing & Liu, LATT2022 (in prep)

Journey of a lattice calculation

for $2 \rightarrow 2$ scattering we are arriving here!

now let's look toward the future...

(iv) Baryonic applications

(iii) Improved control, towards physical masses

(ii) Application in meson sector
(often $M_\pi > M_{\pi, \text{phys}}$)

(i) Formal developments of methods

Can LQCD calculate X for baryonic CPV?

easier

easier

$$\langle \text{hadron} | \mathcal{O}(0) | \text{hadron} \rangle$$

harder

mesons
quark bilinear
light-quarks

baryons

four-quark operators
heavy-quarks

$$\langle \text{multi-hadron} | \text{multi-hadron} \rangle$$



$$\langle \text{multi-hadron state} | \mathcal{O}(0) | \text{hadron} \rangle$$

easier

harder

everything above... plus:

single two-particle channel

many two-particle channels

channels with > two particles

$$\langle \text{hadron} | \mathcal{J}(x) \mathcal{O}(0) | \text{hadron} \rangle$$

four-point functions are difficult

easier

harder

harder

no

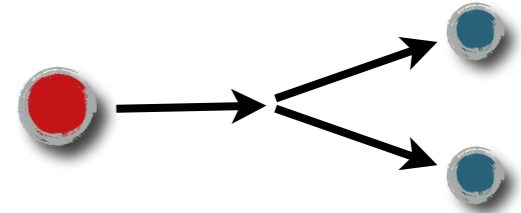
Can multi-hadron states propagate between the currents?

yes

Formal progress: Transition amplitudes

Weak decay

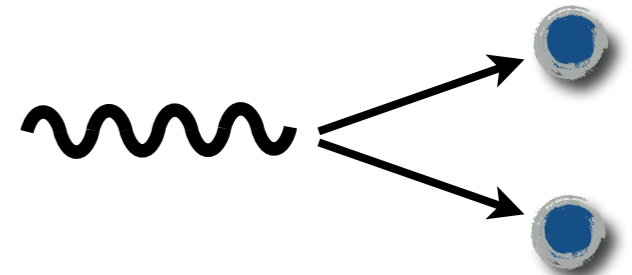
$$\langle \pi\pi, \text{out} | \mathcal{H} | K \rangle \equiv$$



Lellouch, Lüscher (2001) • Kim, Sachrajda, Sharpe (2005) • Christ, Kim, Yamazaki (2005) • MTH, Sharpe (2012)

Time-like form factors

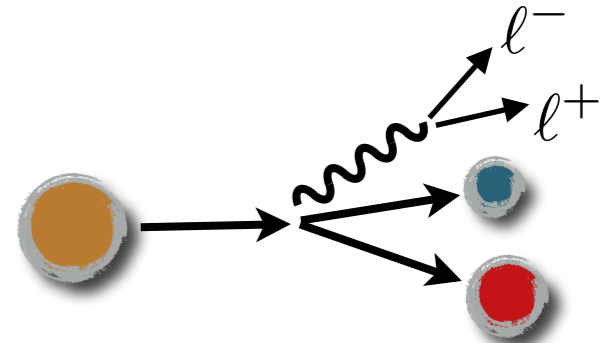
$$\langle \pi\pi, \text{out} | \mathcal{J}_\mu | 0 \rangle \equiv$$



Meyer (2011)

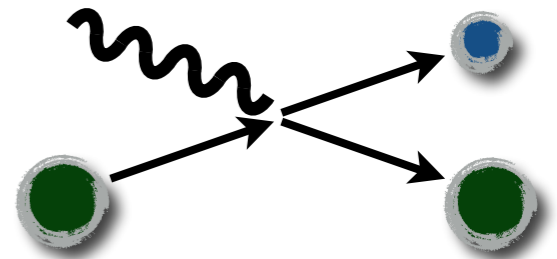
Resonance form factors

$$\langle K\pi, \text{out} | \mathcal{J}_{\alpha\beta} | B \rangle \equiv$$



Particles with spin

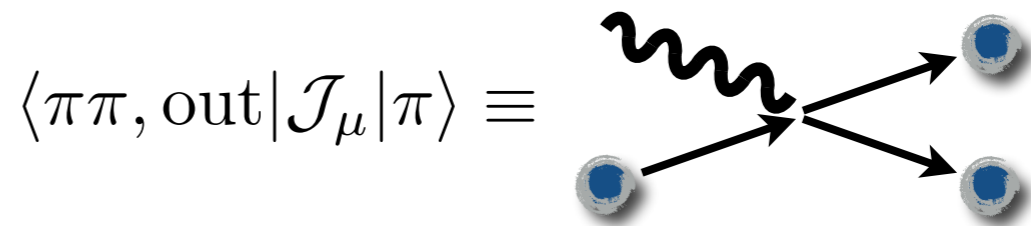
$$\langle N\pi, \text{out} | \mathcal{J}_\mu | N \rangle \equiv$$



Agadjanov *et al.* (2014) • Briceño, MTH, Walker-Loud (2015) • Briceño, MTH (2016)

Pion photo-production

Formal relation



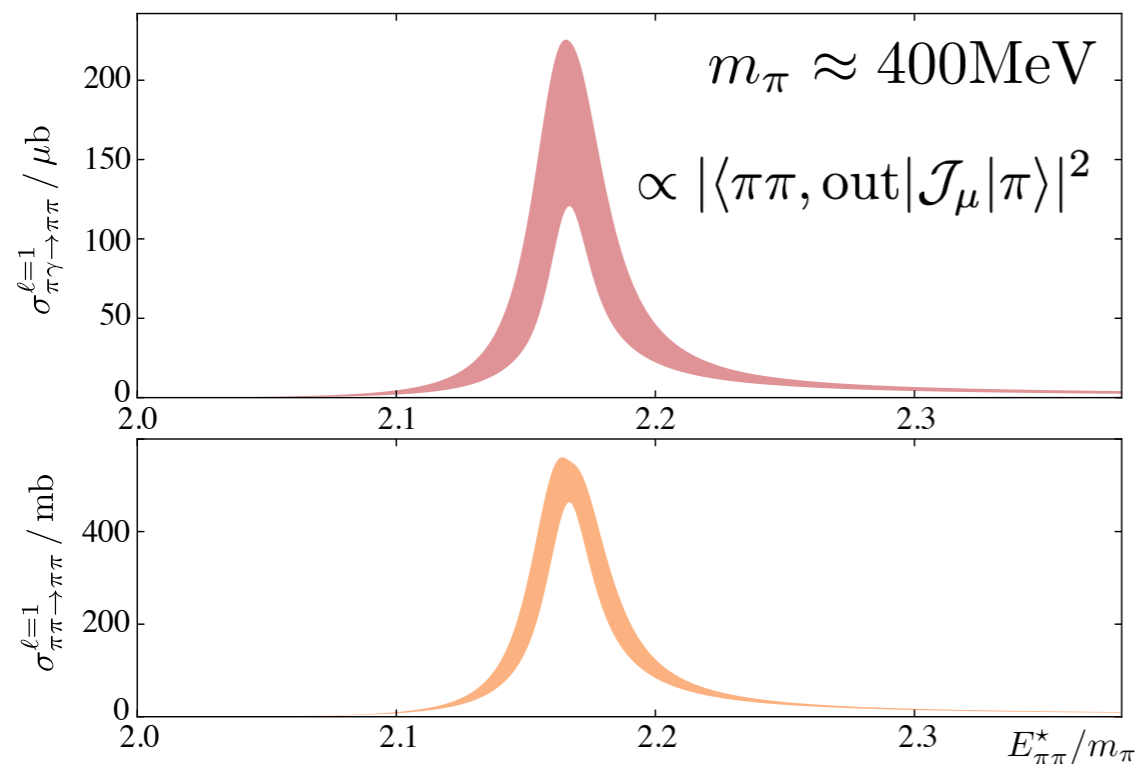
get this from the lattice

experimental observable

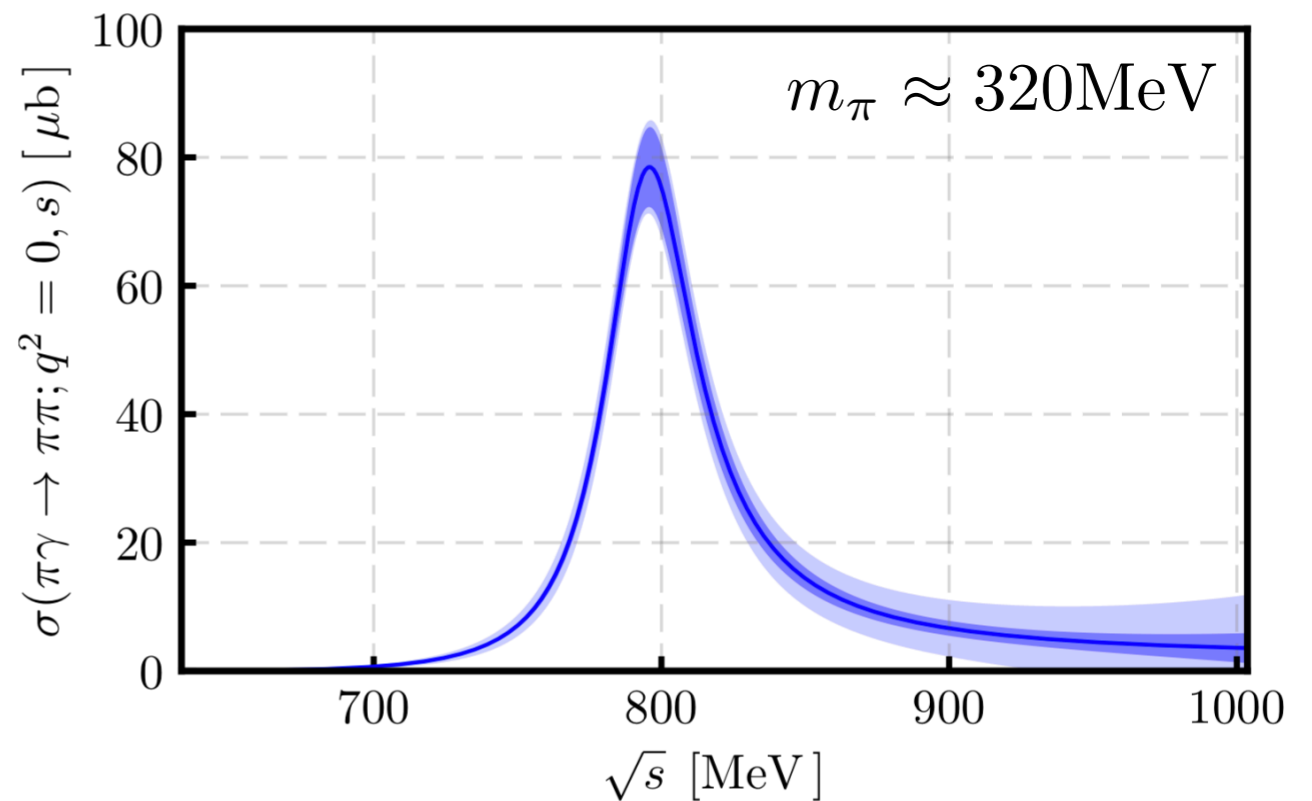
$$|\langle n, L | \mathcal{J}_\mu | \pi \rangle|^2 = \langle \pi | \mathcal{J}_\mu | \pi\pi, \text{in} \rangle \mathcal{R}(E_n, L) \langle \pi\pi, \text{out} | \mathcal{J}_\mu | \pi \rangle$$

Briceño, MTH, Walker-Loud (2015)

Numerical implementation

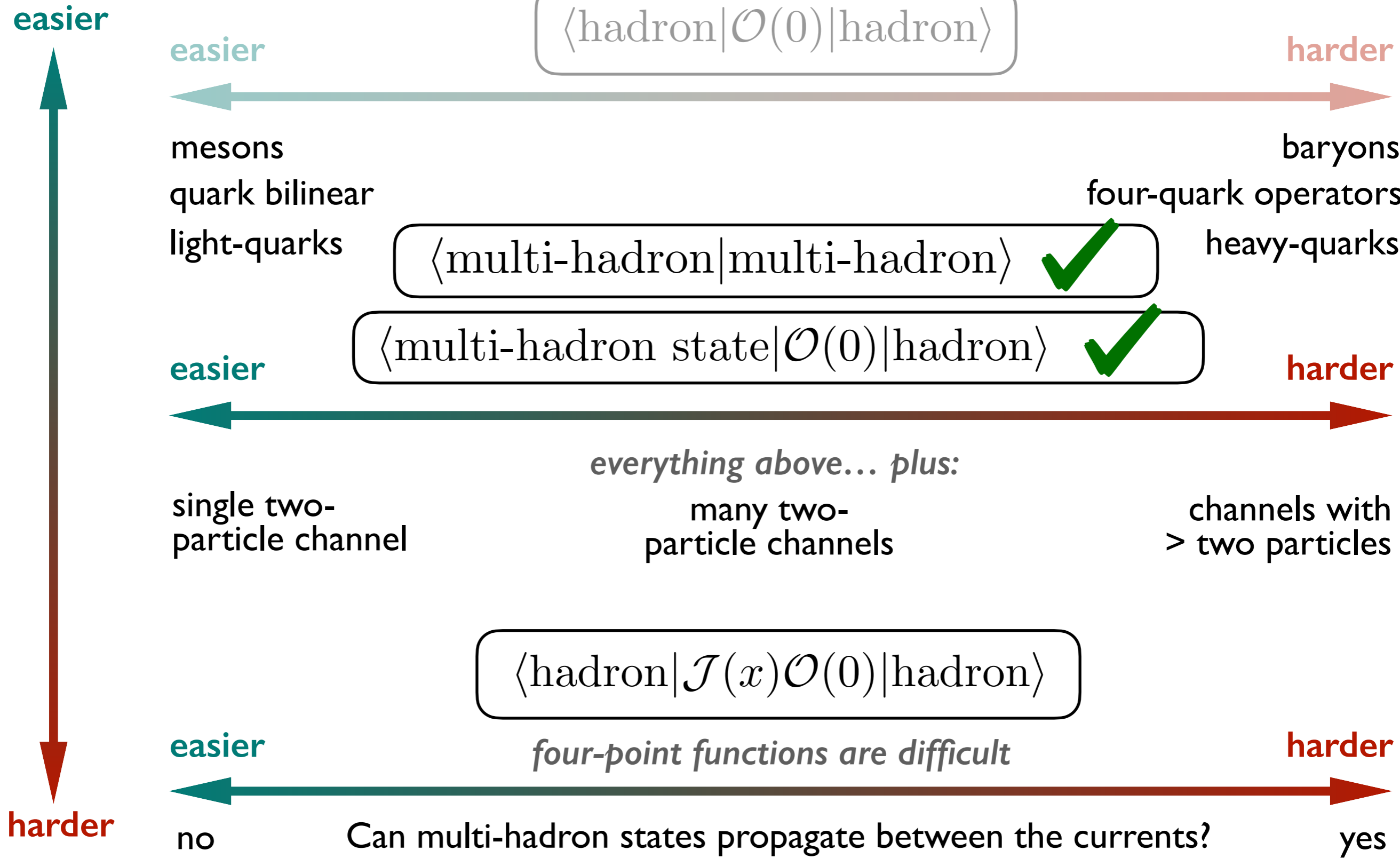


Briceño et. al., Phys. Rev. D93, 114508 (2016)



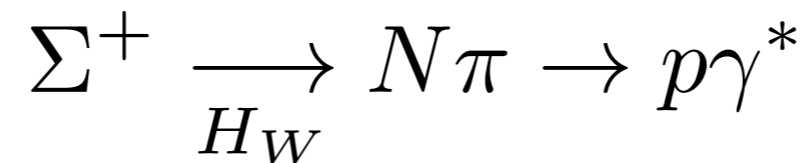
Alexandrou et. al., Phys. Rev. D98, 074502 (2018)

Can LQCD calculate X for baryonic CPV?



Formal & numerical progress: Long-distance matrix elements

- Formal method for extracting these processes is understood
- Key complication is multi-hadron intermediate states



- Christ, Feng, Martinelli, Sachrajda (2015) • Christ *et al.* (2016)
• Briceño, Davoudi, MTH, Schindler, Baroni (2019) • Erben, Gülpers, MTH, Hodgson, Portelli (2022)

Formal & numerical progress: Long-distance matrix elements

- Formal method for extracting these processes is understood
- Key complication is multi-hadron intermediate states

$$\Sigma^+ \xrightarrow{H_W} N\pi \rightarrow p\gamma^*$$

- Issue of growing exponentials (*Christ et al.*)

$$\langle p | j_\mu(0) \mathcal{H}_W(-|\tau|) | \Sigma^+ \rangle_L = \sum_n c_n(L) e^{-(E_n(L) - M_\Sigma)|\tau|} \xrightarrow{\int_{-T}^0 d\tau} \sum_n c_n \frac{1 - e^{-(E_n - M_\Sigma)T}}{M_\Sigma - E_n}$$

Christ, Feng, Martinelli, Sachrajda (2015) • Christ *et al.* (2016)

- Briceño, Davoudi, MTH, Schindler, Baroni (2019) • Erben, Gülpers, MTH, Hodgson, Portelli (2022)

Formal & numerical progress: Long-distance matrix elements

- Formal method for extracting these processes is understood
- Key complication is multi-hadron intermediate states

$$\Sigma^+ \xrightarrow{H_W} N\pi \rightarrow p\gamma^*$$

- Issue of growing exponentials (*Christ et al.*)

$$\langle p | j_\mu(0) \mathcal{H}_W(-|\tau|) | \Sigma^+ \rangle_L = \sum_n c_n(L) e^{-(E_n(L) - M_\Sigma)|\tau|} \xrightarrow{\int_{-T}^0 d\tau} \sum_n c_n \frac{1 - e^{-(E_n - M_\Sigma)T}}{M_\Sigma - E_n}$$

- Issue of power-like finite-volume effects (after discarding exponential)

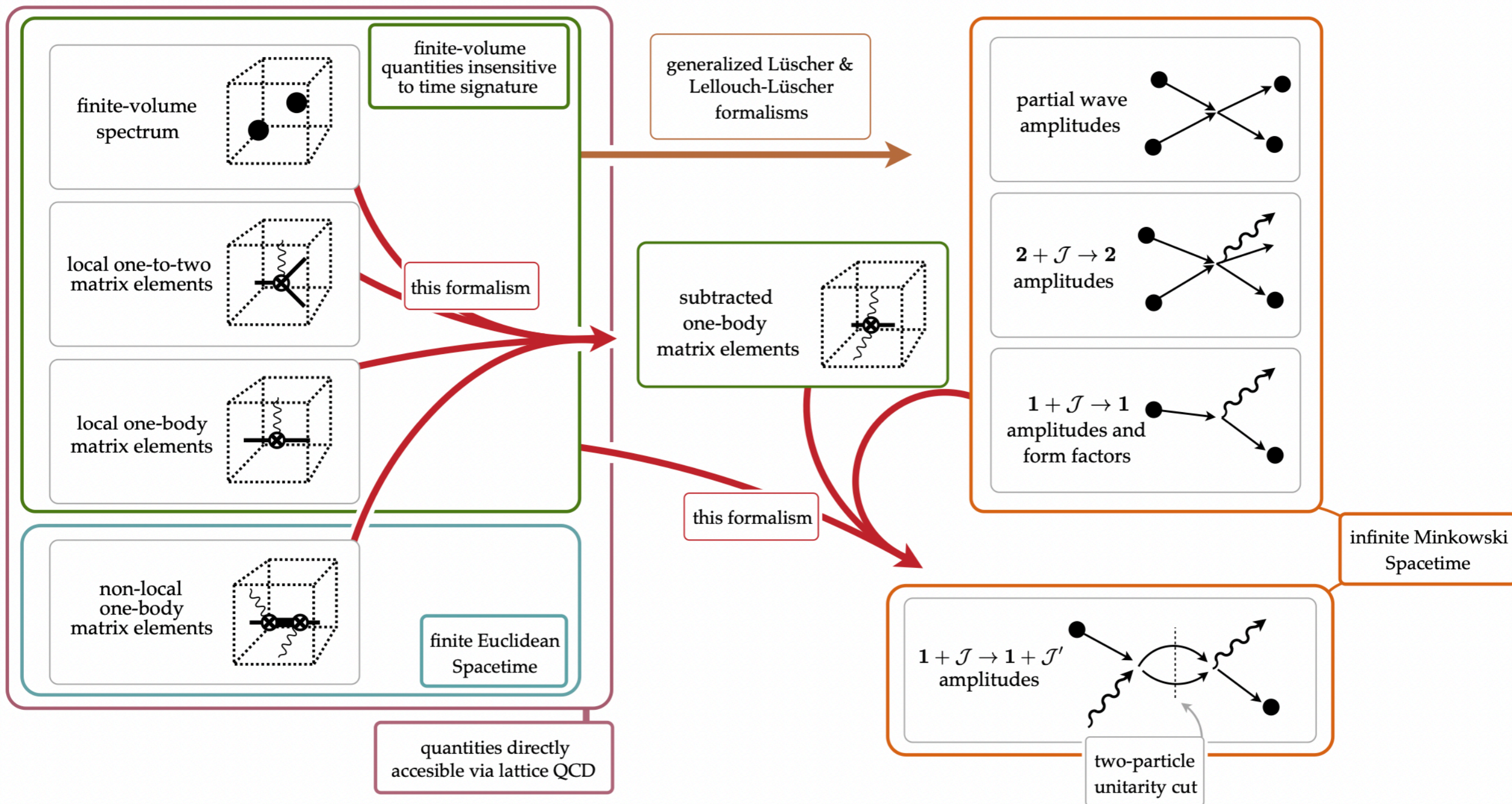
$$F_L = \sum_n \frac{c_n}{M_\Sigma - E_n}$$

- See lattice proceedings (and forthcoming thesis) of Raoul Hodgson

Christ, Feng, Martinelli, Sachrajda (2015) • Christ *et al.* (2016)

• Briceño, Davoudi, MTH, Schindler, Baroni (2019) • Erben, Gülpers, MTH, Hodgson, Portelli (2022)

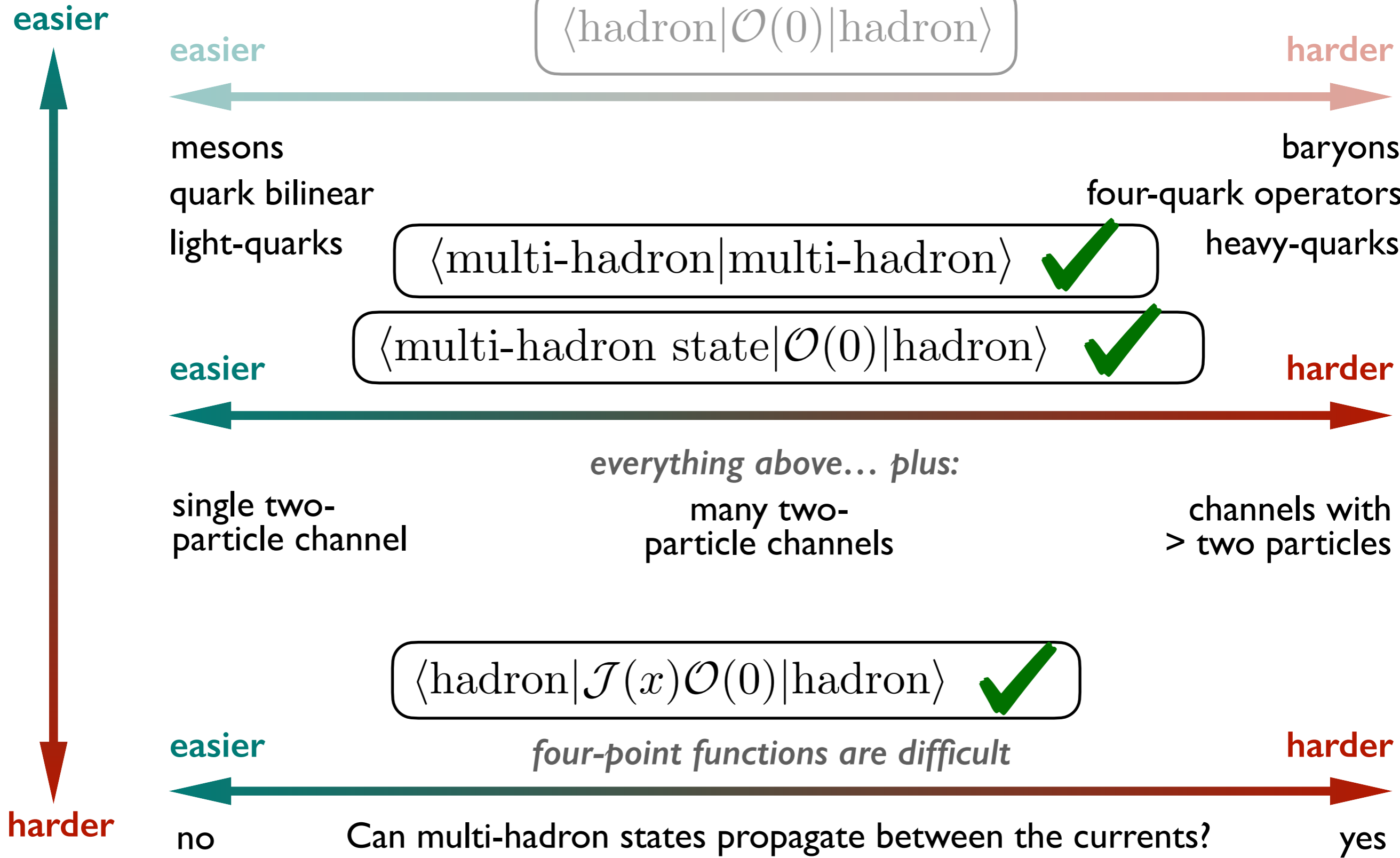
Formal & numerical progress: Long-distance matrix elements



Christ, Feng, Martinelli, Sachrajda (2015) • Christ *et al.* (2016)

• Briceño, Davoudi, MTH, Schindler, Baroni (2019) • Erben, Gülpers, MTH, Hodgson, Portelli (2022)

Can LQCD calculate X for baryonic CPV?



Towards $N\pi\pi$

- Multiple three-particle finite-volume formalisms developed (so far only spin zero)

MTH, Sharpe (2014-2016)

See also Döring, Mai, Hammer, Pang, Rusetsky

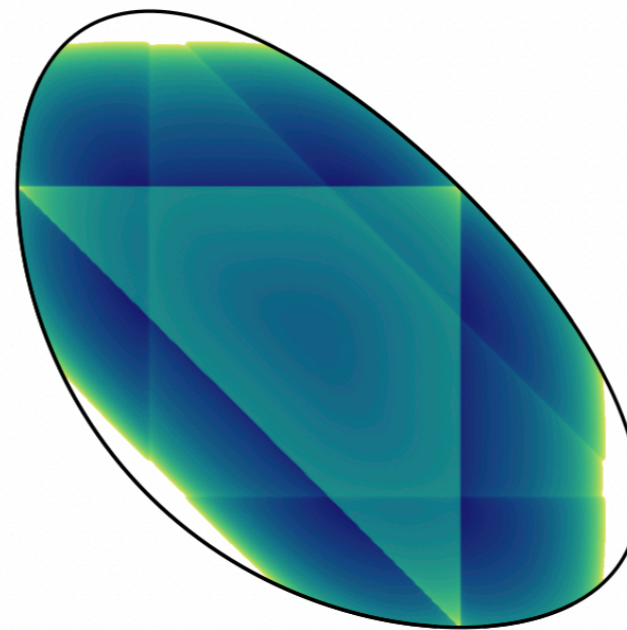
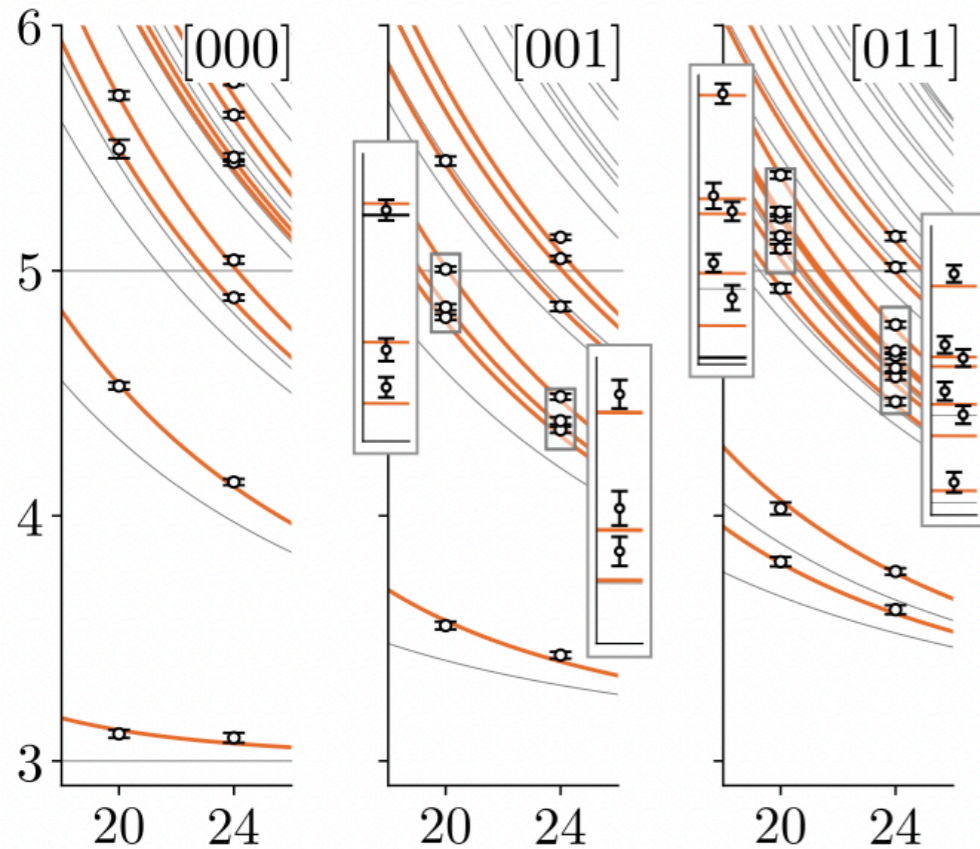
Towards $N\pi\pi$

- Multiple three-particle finite-volume formalisms developed (so far only spin zero)

MTH, Sharpe (2014-2016)

See also Döring, Mai, Hammer, Pang, Rusetsky

- First lattice calculations appearing... e.g. $\pi^+\pi^+\pi^+ \rightarrow \pi^+\pi^+\pi^+$



- Extract reliable spectrum
- Use formalism to fit scheme-dependent K-matrix
- Solve integral equations to reach physical amplitude

MTH, Briceño, Edwards, Thomas, Wilson, *Phys.Rev.Lett.* 126 (2021) 012001

- See Blanton et al. 2022 and 2023 for pion and kaon results
- See Sadasivan et al. 2022 for application to $a_1(1260)$

Not discussed here...

□ Signal-to-noise problem and solutions

- For a single nucleon (signal/noise) $\sim \exp[-(M_N - 3/2M_\pi)\tau_{\text{Euclidean}}]$
- One method to solve this is multi-level = "only update gauge field in regions"

□ HAL-QCD potential method

- Extract effective potential from lattice calculation
- Requires derivative expansion

See... Murakami et al.,

Lattice QCD studies on decuplet baryons as meson-baryon bound states in the HAL QCD method

□ Spectral-function reconstruction methods

- Approach that does not use finite-volume as a tool
- Attempts to solve regulated inverse Laplace transform
- Delivers observables "smeared" over a range of energies

Points for discussion

□ With lattice QCD experts...

- What are the biggest limitations & next steps for **single-baryon observables**?
- Will multi-level algorithms play an important role?
- How far are multi-meson observables from full-error-budget physical point?
 - Will we get there?... When?
- How far are multi-baryon observables from full-error-budget physical point?

□ With other theorists...

- What are the single-baryon observables where lattice QCD can be most impactful?
- What are the simplest multi-hadron observables for lattice QCD?
- Can theory provide scattering data to be used in lattice finite-volume formulae?
 - Also away from physical pion mass?
- How light to pions have to be in lattice QCD calculations to be useful?

□ With experimentalists...

- Which processes are on the horizon?
- Can lattice + experiment identify more lattice-friendly observables?

Points for discussion

□ With lattice QCD experts...

- What are the biggest limitations & next steps for **single-baryon observables**?
- Will multi-level algorithms play an important role?
- How far are multi-meson observables from full-error-budget physical point?
 - Will we get there?... When?
- How far are multi-baryon observables from full-error-budget physical point?

□ With other theorists...

- What are the single-baryon observables where lattice QCD can be most impactful?
- What are the simplest multi-hadron observables for lattice QCD?
- Can theory provide scattering data to be used in lattice finite-volume formulae?
 - Also away from physical pion mass?
- How light to pions have to be in lattice QCD calculations to be useful?

□ With experimentalists...

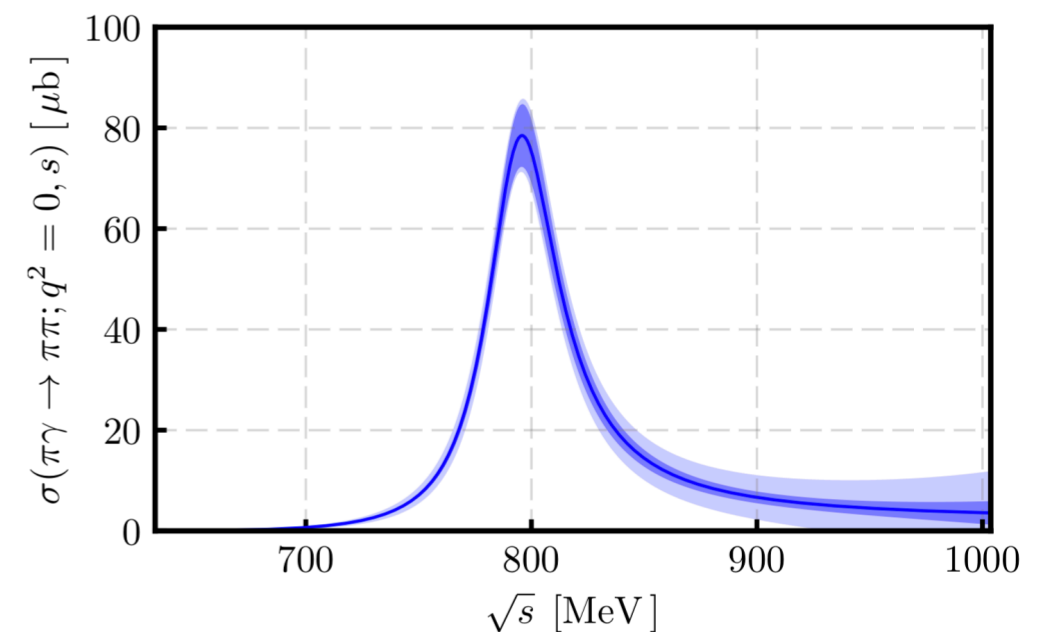
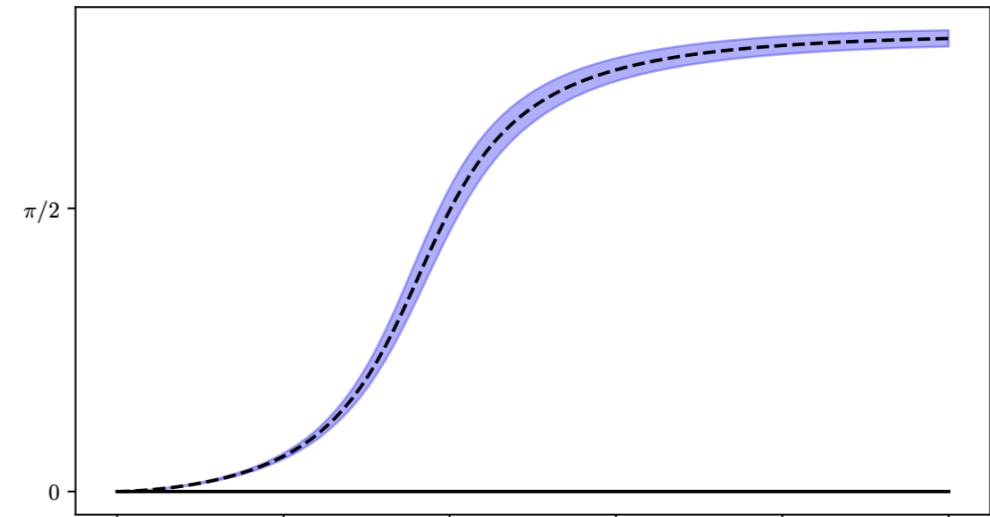
- Which processes are on the horizon?
- Can lattice + experiment identify more lattice-friendly observables?

Summary and outlook

- If amplitude has multi-hadron contributions
→ LQCD requires complicated formalism (*volume + Euclidean*)

- Robust baryonic scattering studies now appearing
complete finite-volume spectrum, field theoretically mapped to amplitudes

- Formal developments ahead of numerical calculations



Supported by UKRI FLF



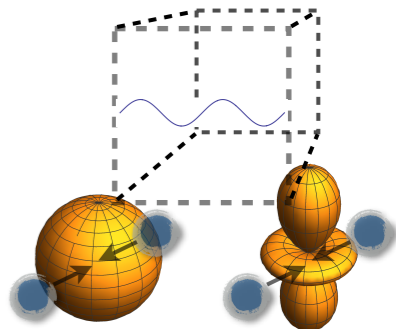
UK Research
and Innovation

Thanks for listening!... questions?

Coupled channels

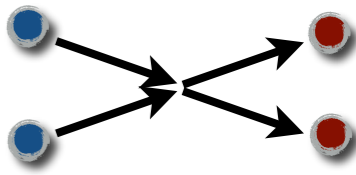
□ The cubic volume mixes different partial waves...

e.g. $K\pi \rightarrow K\pi$
 $\vec{P} \neq 0 \longrightarrow \det \left[\begin{pmatrix} \mathcal{K}_s^{-1} & 0 \\ 0 & \mathcal{K}_p^{-1} \end{pmatrix} + \begin{pmatrix} F_{ss} & F_{sp} \\ F_{ps} & F_{pp} \end{pmatrix} \right] = 0$



...as well as different flavor channels...

e.g. $a = \pi\pi$
 $b = K\bar{K} \longrightarrow \det \left[\begin{pmatrix} \mathcal{K}_{a \rightarrow a} & \mathcal{K}_{a \rightarrow b} \\ \mathcal{K}_{b \rightarrow a} & \mathcal{K}_{b \rightarrow b} \end{pmatrix}^{-1} + \begin{pmatrix} F_a & 0 \\ 0 & F_b \end{pmatrix} \right] = 0$



□ Workflow...

Correlators with a large operator basis
 $\langle \mathcal{O}_a(\tau) \mathcal{O}_b^\dagger(0) \rangle$

Reliably extract finite-volume energies
 $\langle \Omega_m(\tau) \Omega_m^\dagger(0) \rangle \sim e^{-E_m(L)\tau}$

Vary L and P to recover a dense set of energies

$[000], \Delta_1$	○	○	○	○	○
$[001], \Delta_1$		○	○	○	○
$[011], \Delta_1$	○	○	○	○	○

└───────────────────┘ $E_n(L)$



Identify a broad list of K-matrix parametrizations

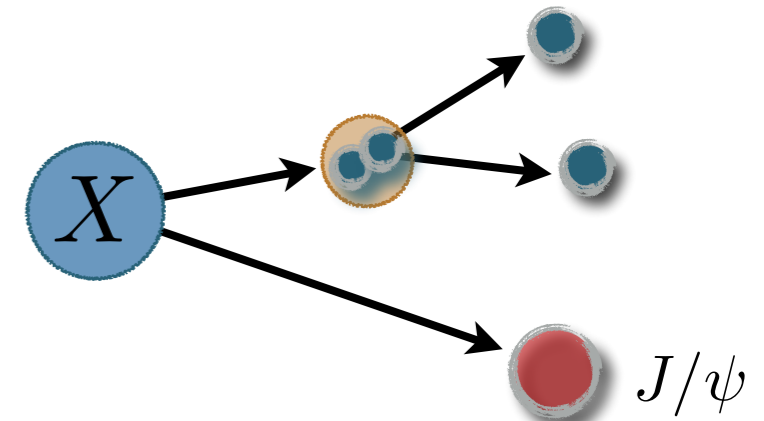
polynomials and poles EFT based dispersion theory based

Perform global fits to the finite-volume spectrum

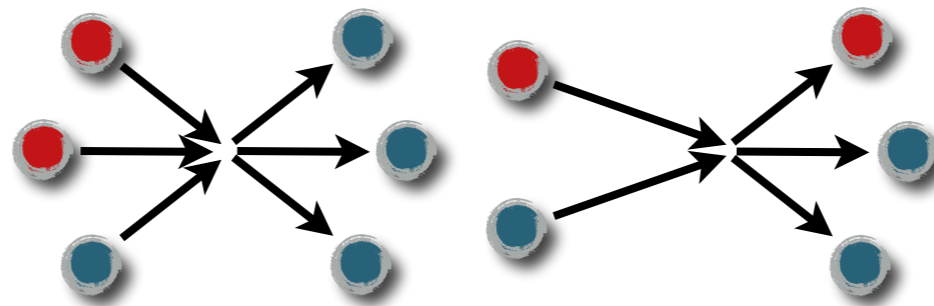
3-particle amplitudes

2-to-2 only samples J^P 0^+ 1^- 2^+ ...

many interesting resonances have significant 3-body decays



Goal: *finite-volume + unitarity formalism* for generic two- and three-particle systems



Applications...

exotic resonance pole positions, couplings, quantum numbers

$$\omega(782), a_1(1420) \rightarrow \pi\pi\pi$$

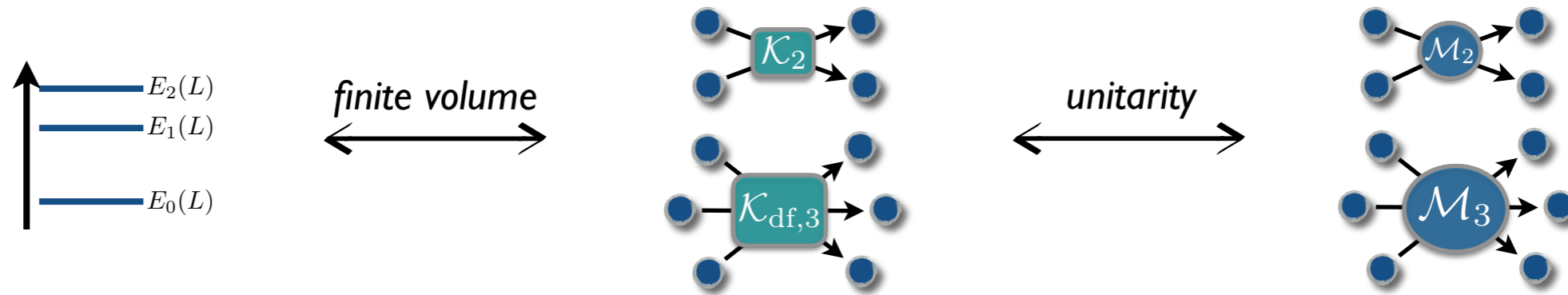
$$X(3872) \rightarrow J/\psi\pi\pi$$

$$X(3915)[Y(3940)] \rightarrow J/\psi\pi\pi$$

form factors and transitions

and much more!... (3-body forces, weak transitions, gluons content)

Status...



Identical spin-zero, no 2-to-3, no K2 poles • MTH, Sharpe (2014, 2015) •

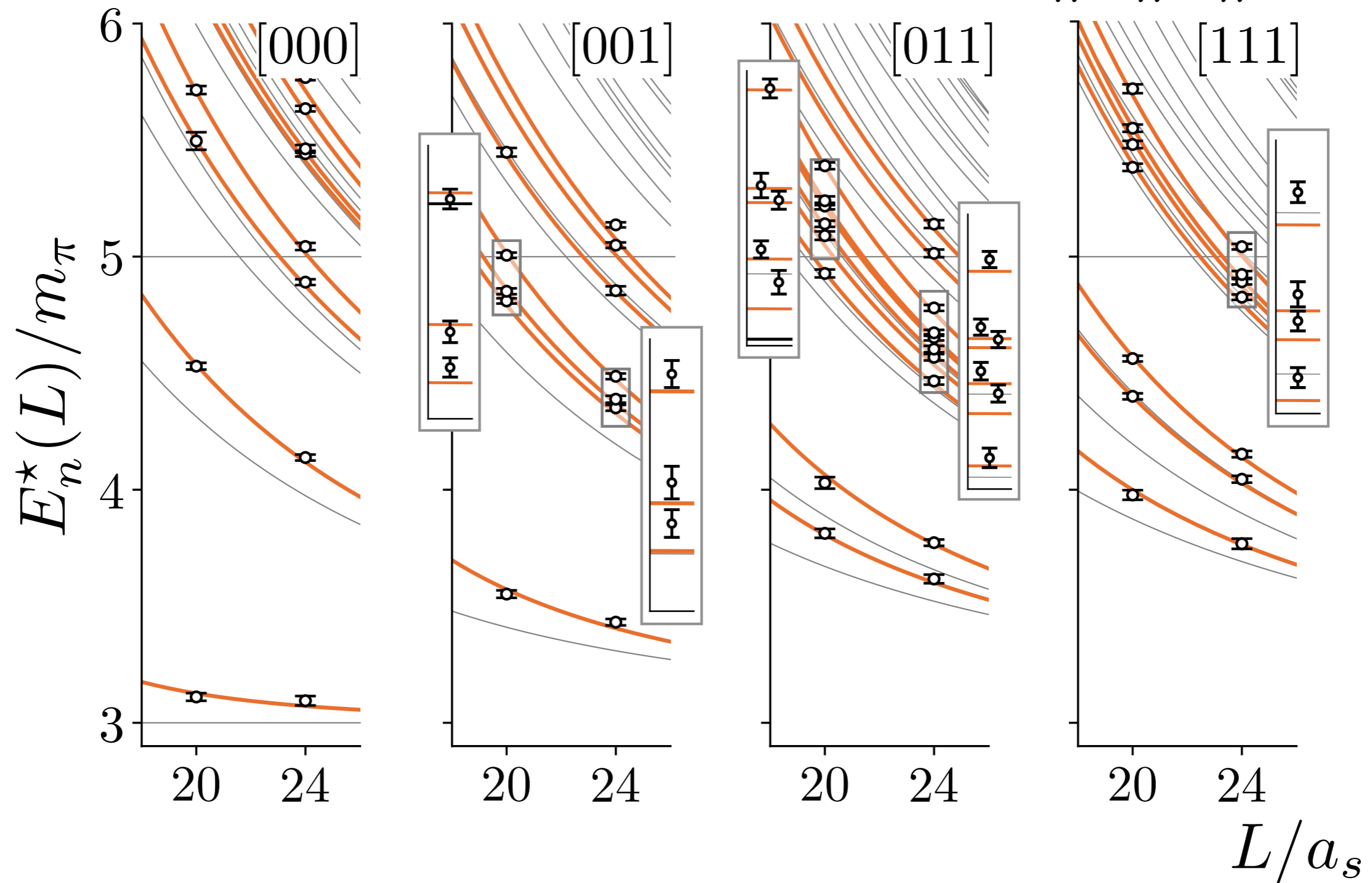
as above... but including 2-to-3 • Briceño, MTH, Sharpe (2017) •

as above... but including K2 poles • Briceño, MTH, Sharpe (2018) •

Non-identical, non-degenerate spin-zero $\pi\pi\pi \rightarrow \rho\pi \rightarrow \omega \rightarrow \rho\pi \rightarrow \pi\pi\pi$
 • MTH, Romero-López, Sharpe (2020) • Blanton, Sharpe (2020, 2021)

Multiple three-particle channels... Spin!

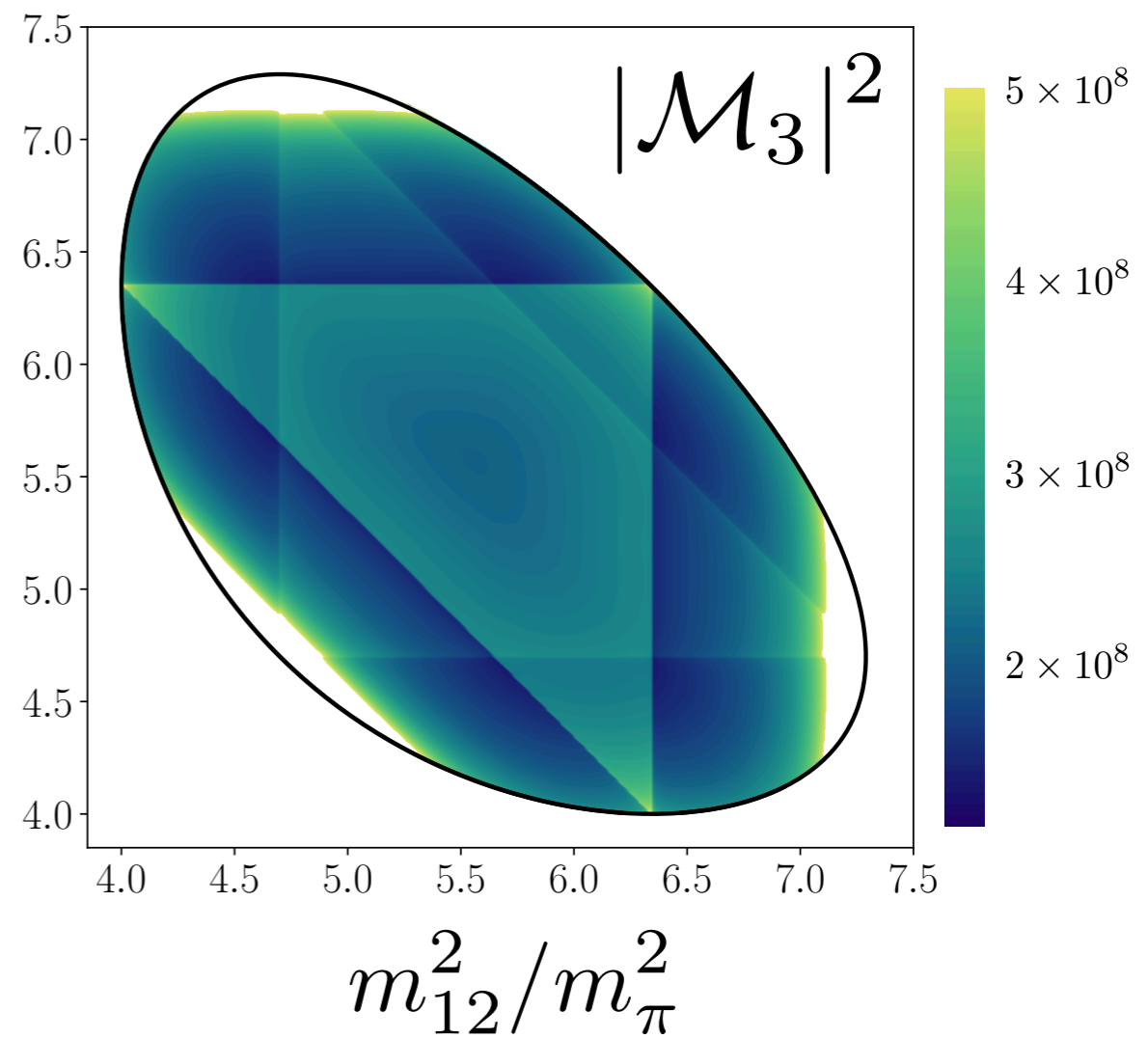
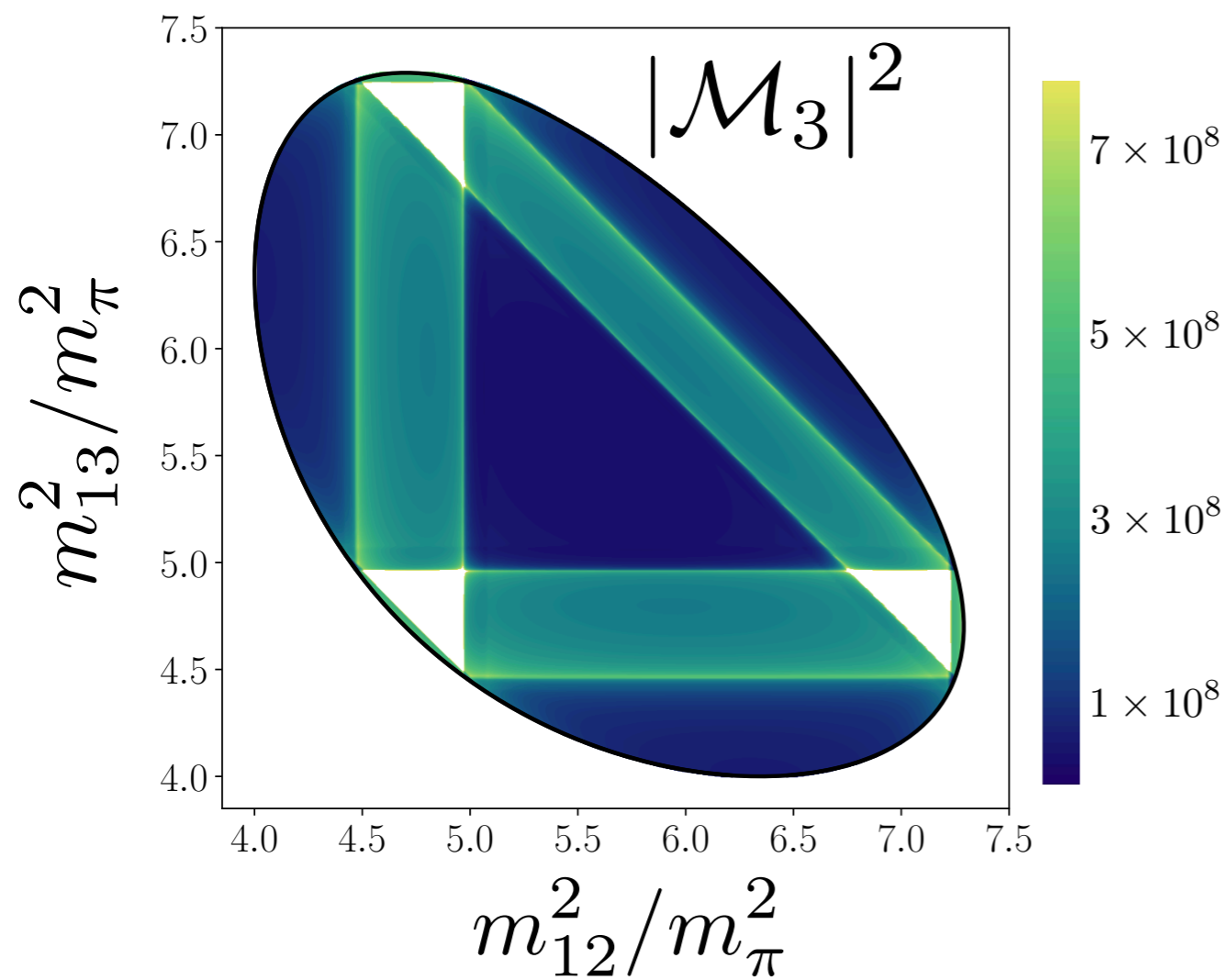
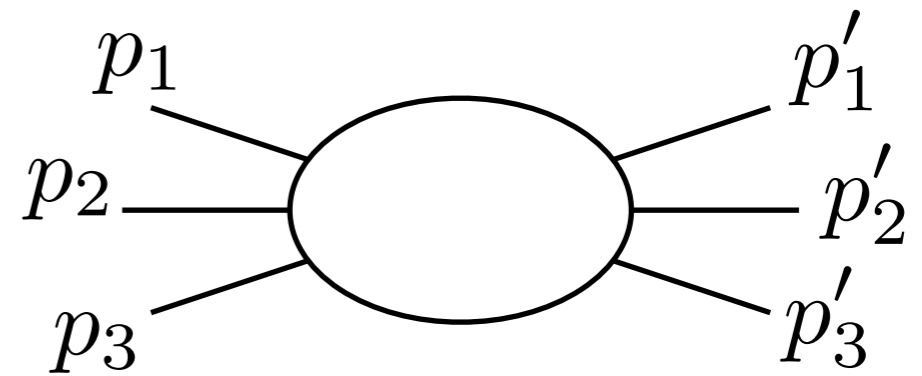
$\pi^+\pi^+\pi^+$ energies



MTH, Briceño, Edwards, Thomas, Wilson, *Phys.Rev.Lett.* 126 (2021) 012001,

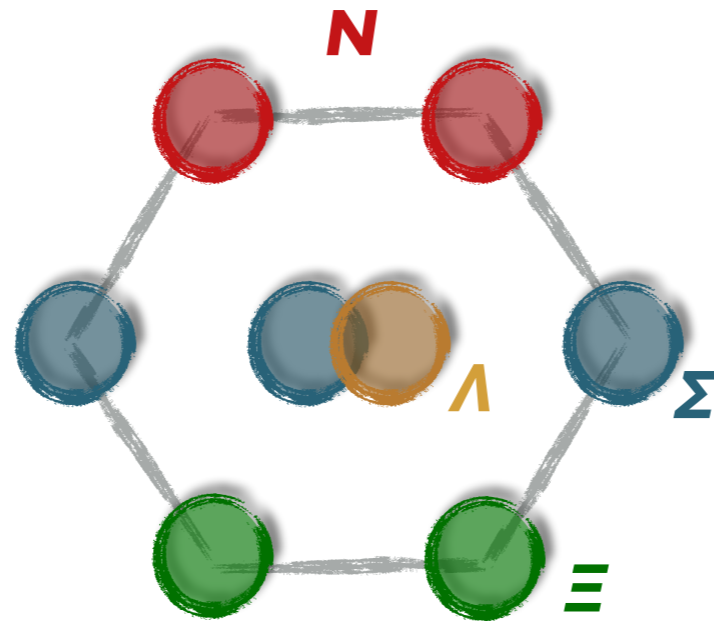
see also work by... Culver, Döring, Hanlon, Hörz, Mai, Morningstar, Romero-Lopez, Sharpe + ETMC

$$\mathcal{M}_3 = \sum_{i,j \in \{1,2,3\}} \mathcal{M}_3^{\text{un}}(p'_i, p_j)$$



$$\Delta \rightarrow N\pi$$

- Andersen et al. 2018
- Andersen et al. 2019
- Silvi et al. 2021
- Pittler et al. 2021
- Bulava et al. 2022

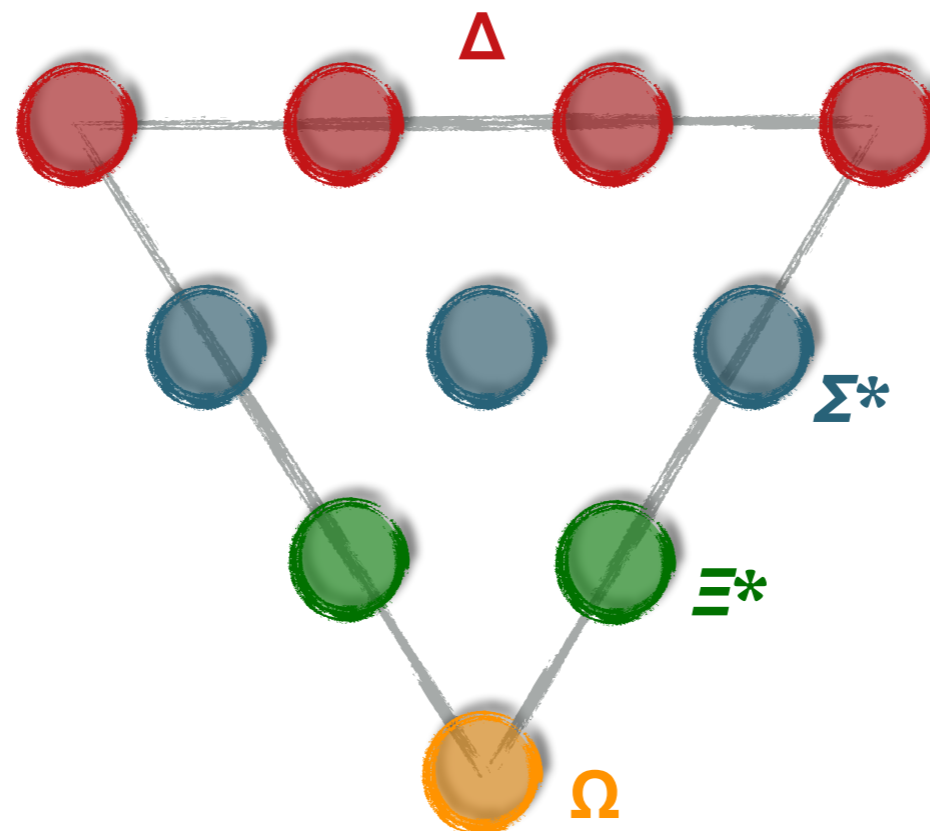


(focusing here on studies with scattering states)

Baryons are difficult!

$$N^* \rightarrow N\pi$$

- Lang et al. 2017
- Wu et al. 2017
- Kiratidis et al. 2017



$$\Lambda \rightarrow \bar{K}N$$

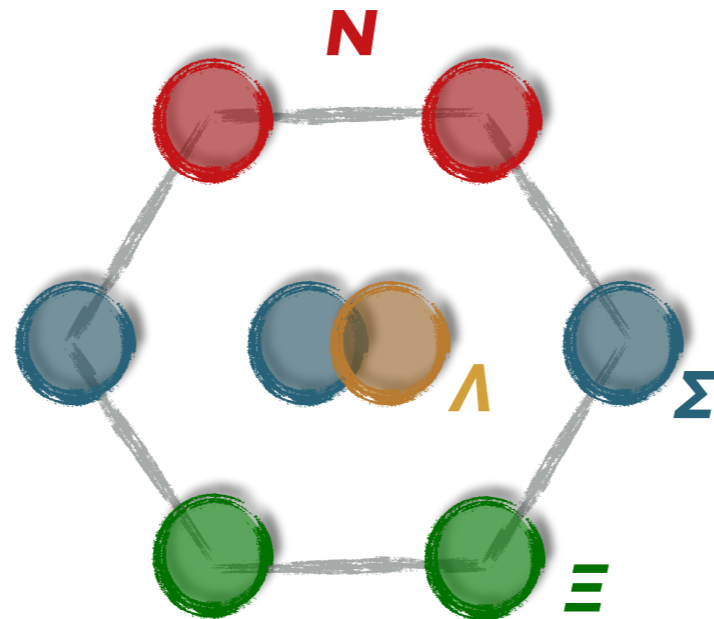
- Hall et al. 2015

See also...

- Detmold and Nicholson 2015
- Wu et al. 2018
- Xing & Liu, LATT2022 (in prep)

$$\Delta \rightarrow N\pi$$

- Andersen et al. 2018
- Andersen et al. 2019
- Silvi et al. 2021
- Pittler et al. 2021
- Bulava et al. 2022

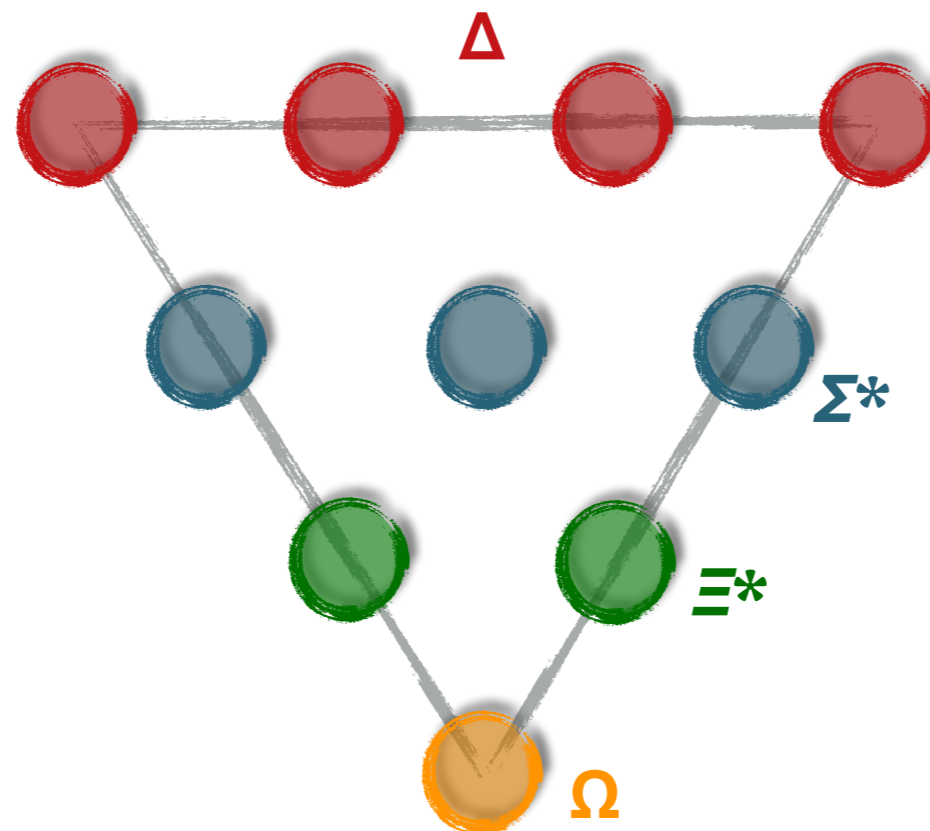


(focusing here on studies with scattering states)

Baryons are difficult!

$$N^* \rightarrow N\pi$$

- Lang et al. 2017
- Wu et al. 2017
- Kiratidis et al. 2017



$$\Lambda \rightarrow \bar{K}N$$

- Hall et al. 2015

See also...

- Detmold and Nicholson 2015
- Wu et al. 2018
- Xing & Liu, LATT2022 (in prep)

$N\pi$ elastic scattering ($M_\pi = 255$ MeV)

PHYSICAL REVIEW D

covering particles, fields, gravitation, and cosmology

P -wave nucleon-pion scattering amplitude in the $\Delta(1232)$ channel from lattice QCD

Giorgio Silvi, Srijit Paul, Constantia Alexandrou, Stefan Krieg, Luka Leskovec, Stefan Meinel, John Negele, Marcus Petschlies, Andrew Pochinsky, Gumaro Rendon, Sergey Syritsyn, and Antonino Todaro

□ $M_\pi = 255$ MeV

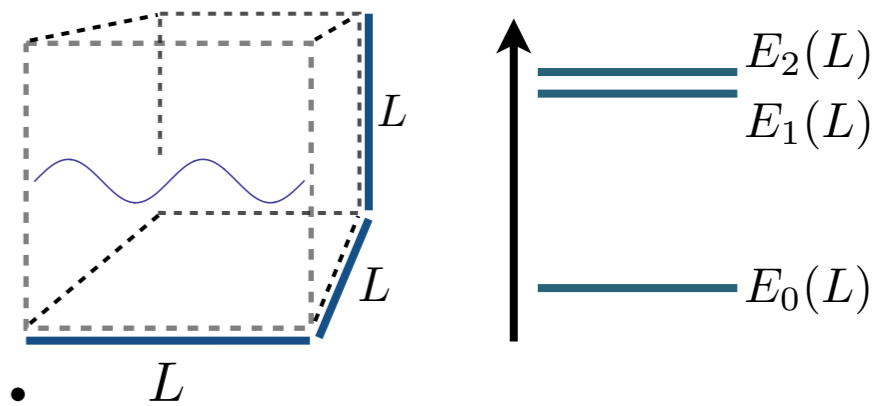
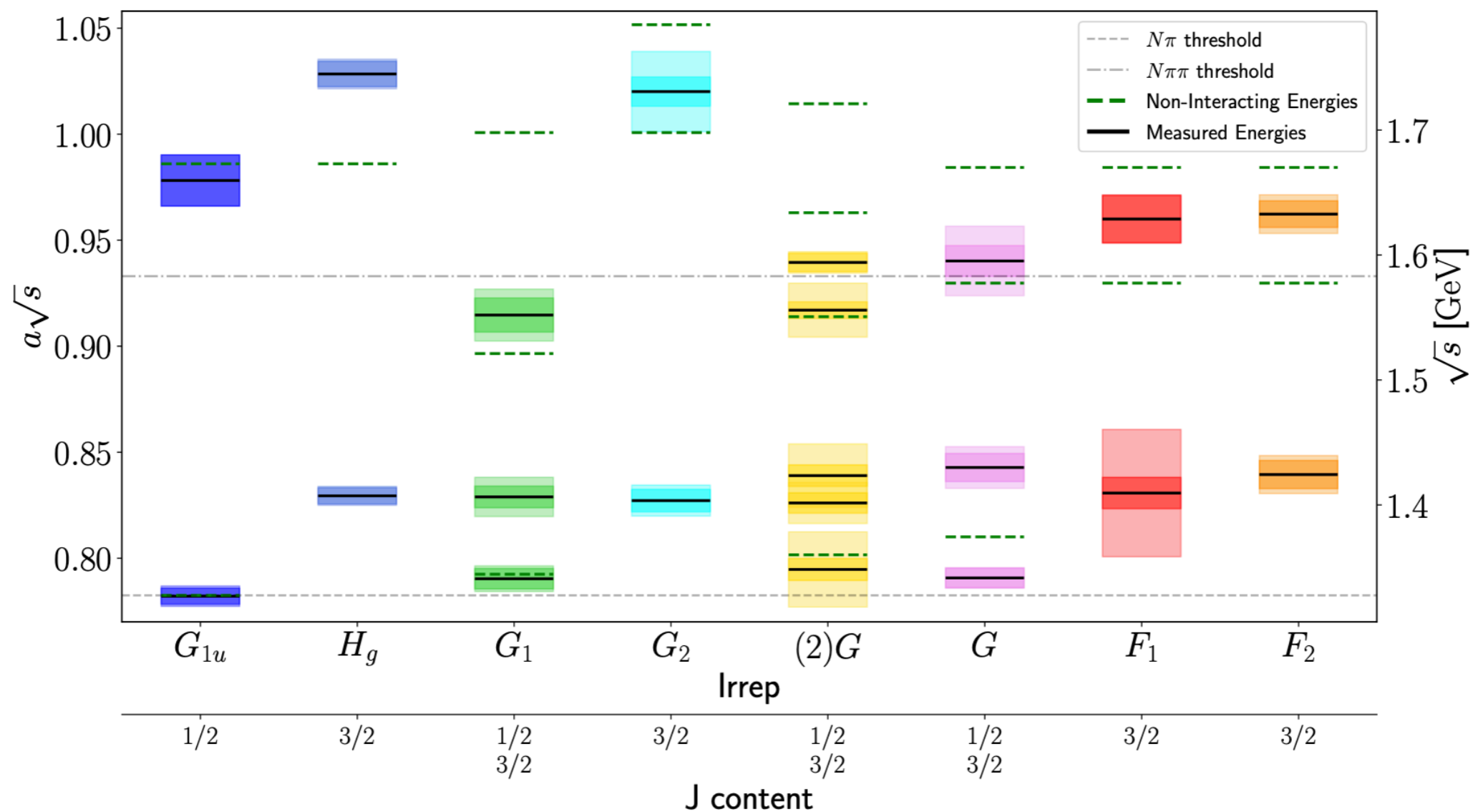
□ Studied scattering-lengths and the Δ channel

$I = 3/2$: $J^P = 1/2^- (S), 3/2^+ (P)$ [$1/2^+ (P), 3/2^- (D), 5/2^- (D)$]

□ Local Δ -like operators + $N(\mathbf{p}_1)\pi(\mathbf{p}_2)$ operators

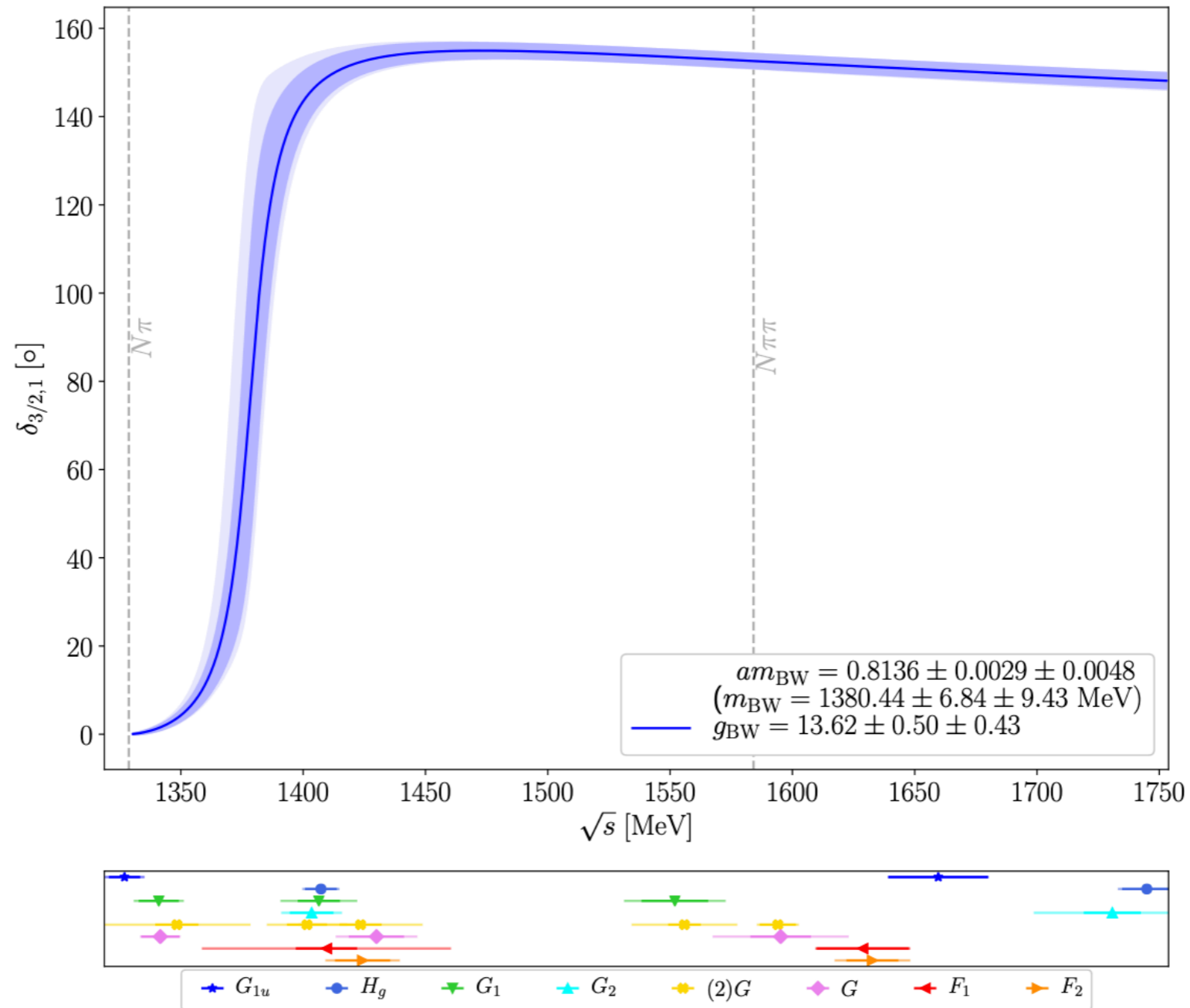
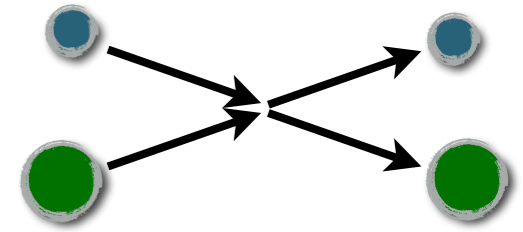
• Silvi et al., PRD 2021, 2101.00689 •

$N\pi$ finite-volume energies ($M_\pi = 255$ MeV)



• Silvi et al., PRD 2021, 2101.00689

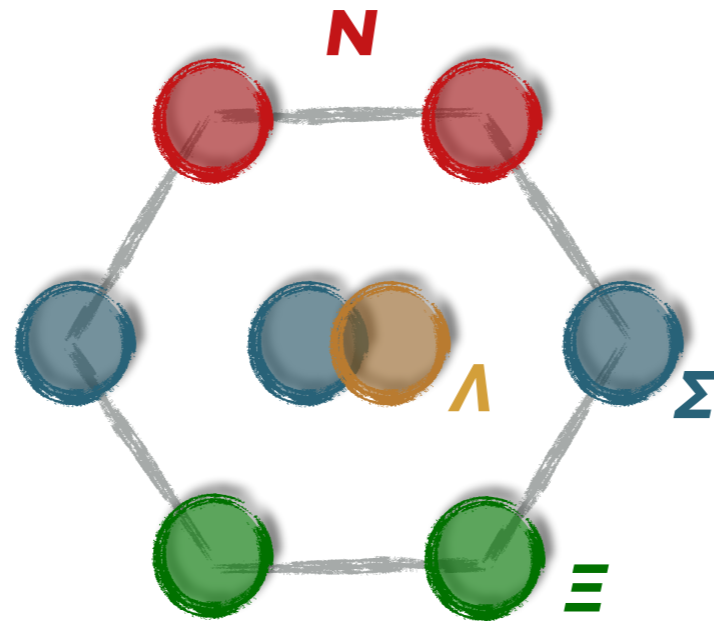
$$N\pi \rightarrow \Delta \rightarrow N\pi \quad (M_\pi = 255 \text{ MeV})$$



- Silvi et al., PRD 2021, 2101.00689 •

$$\Delta \rightarrow N\pi$$

- Andersen et al. 2018
- Andersen et al. 2019
- Silvi et al. 2021
- Pittler et al. 2021
- Bulava et al. 2022

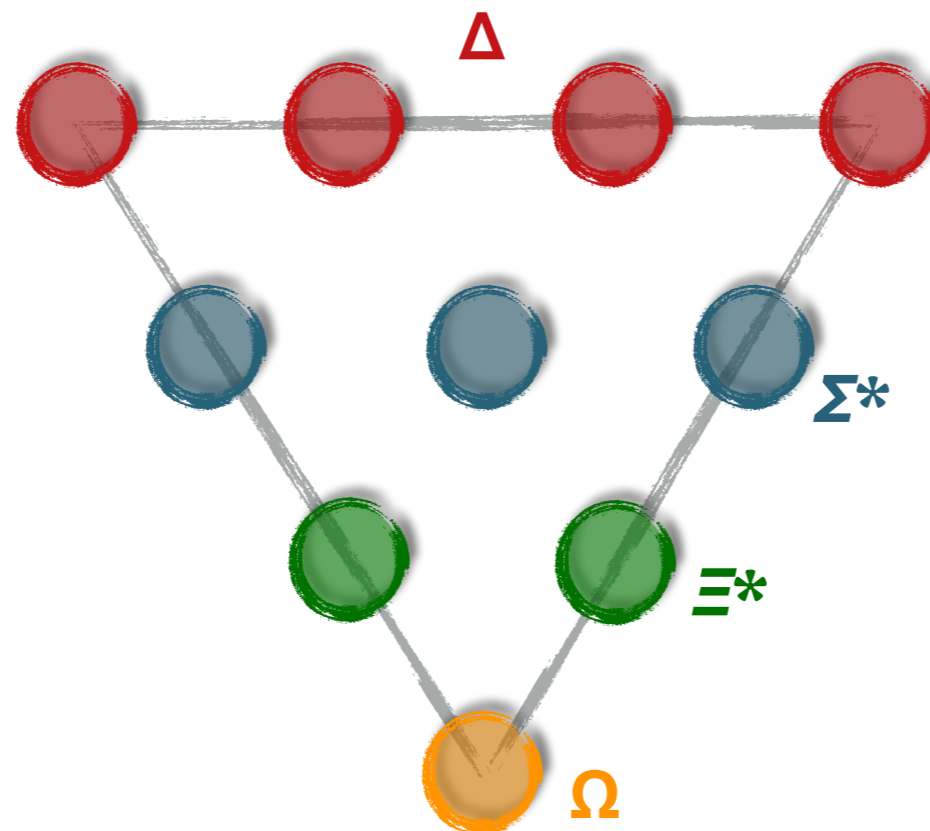


(focusing here on studies with scattering states)

Baryons are difficult!

$$N^* \rightarrow N\pi$$

- Lang et al. 2017
- Wu et al. 2017
- Kiratidis et al. 2017



$$\Lambda \rightarrow \bar{K}N$$

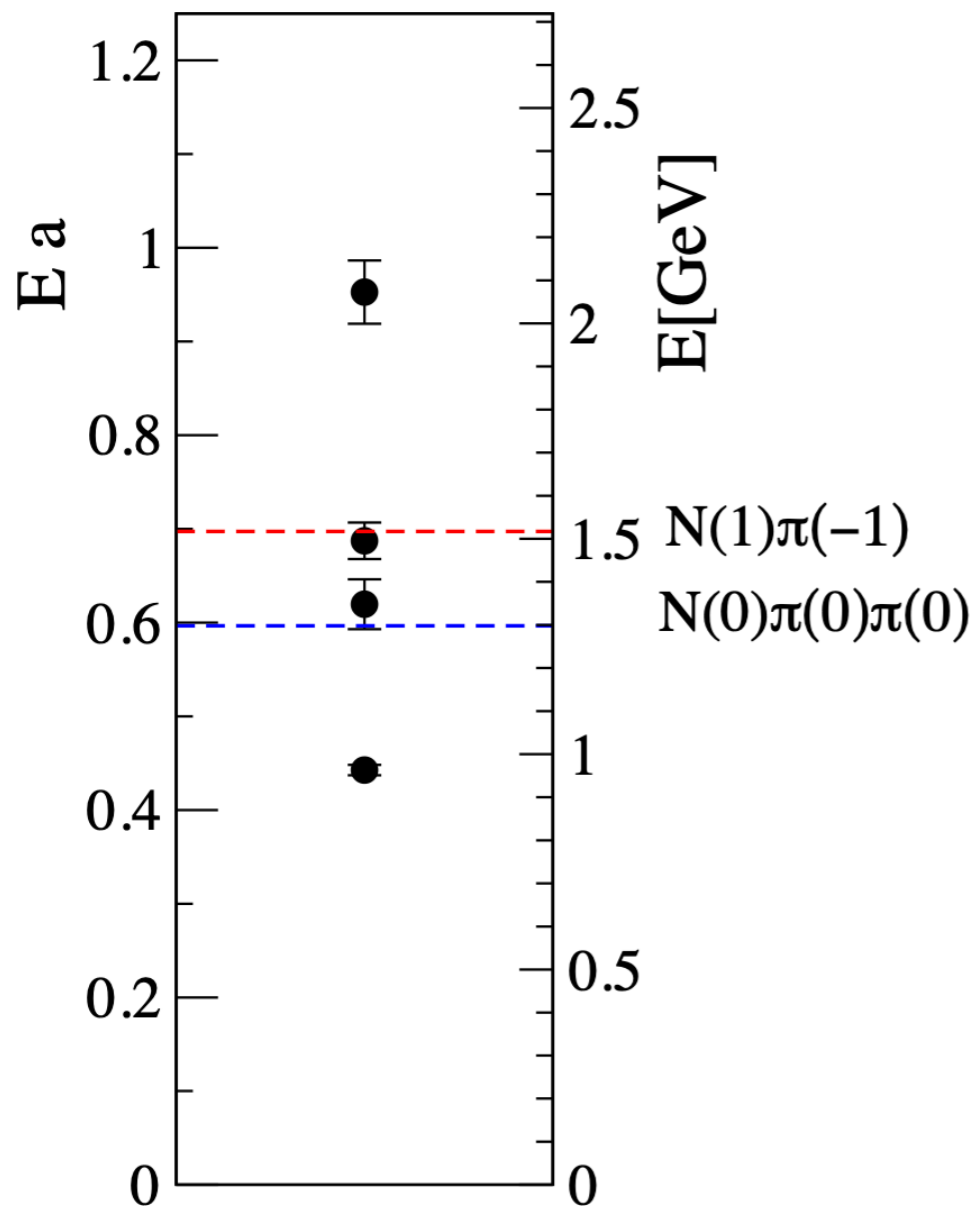
- Hall et al. 2015

See also...

- Detmold and Nicholson 2015
- Wu et al. 2018
- Xing & Liu, LATT2022 (in prep)

Roper resonance

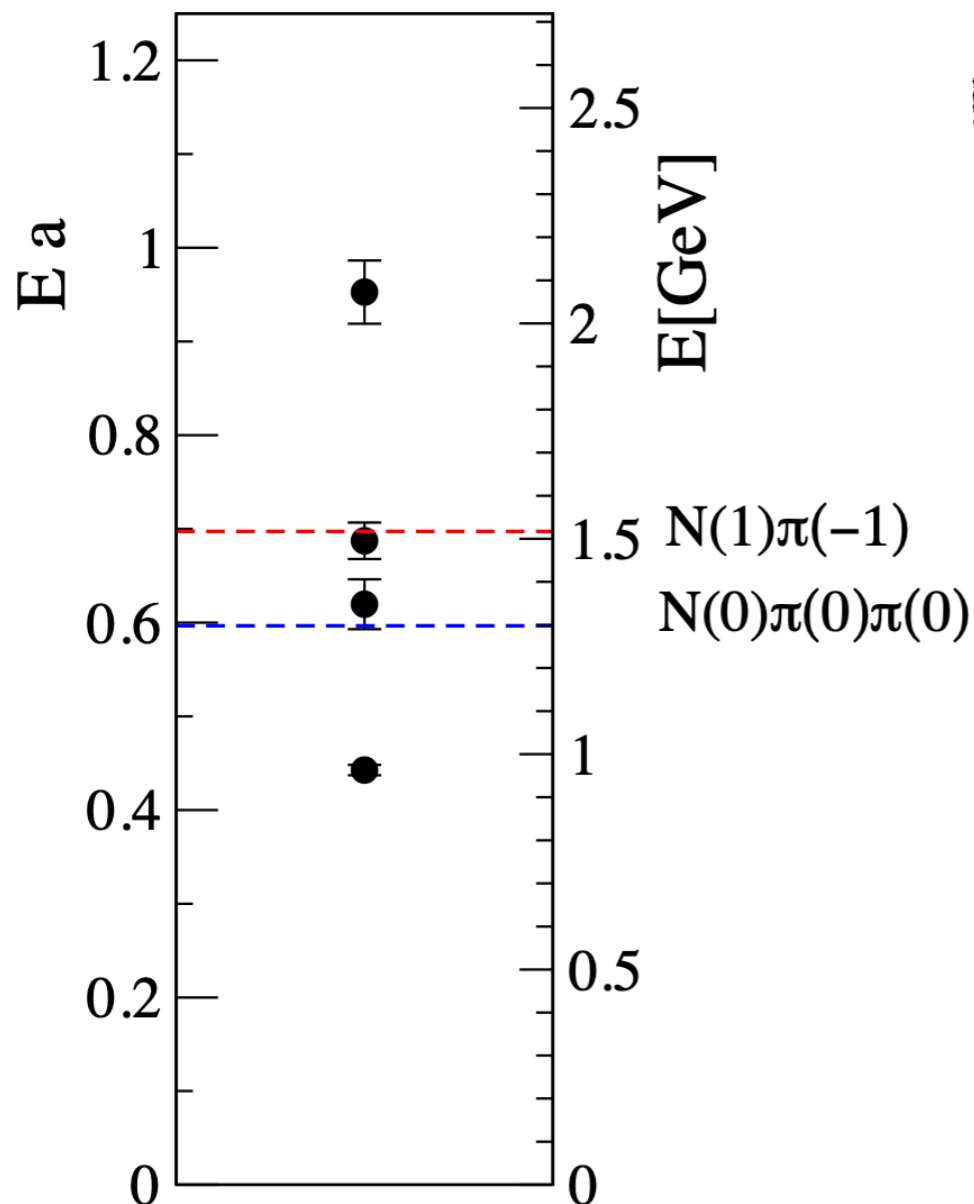
- Results with impressive operators from Lang et al., $M_\pi = 156$ MeV
- Interpretation presented by D. Leinweber this morning



• Lang, Leskovec, Padmanath, Prelovsek (2017) •

Roper resonance

- Results with impressive operators from Lang et al., $M_\pi = 156$ MeV
- Interpretation presented by D. Leinweber this morning



- My personal take, need more data...

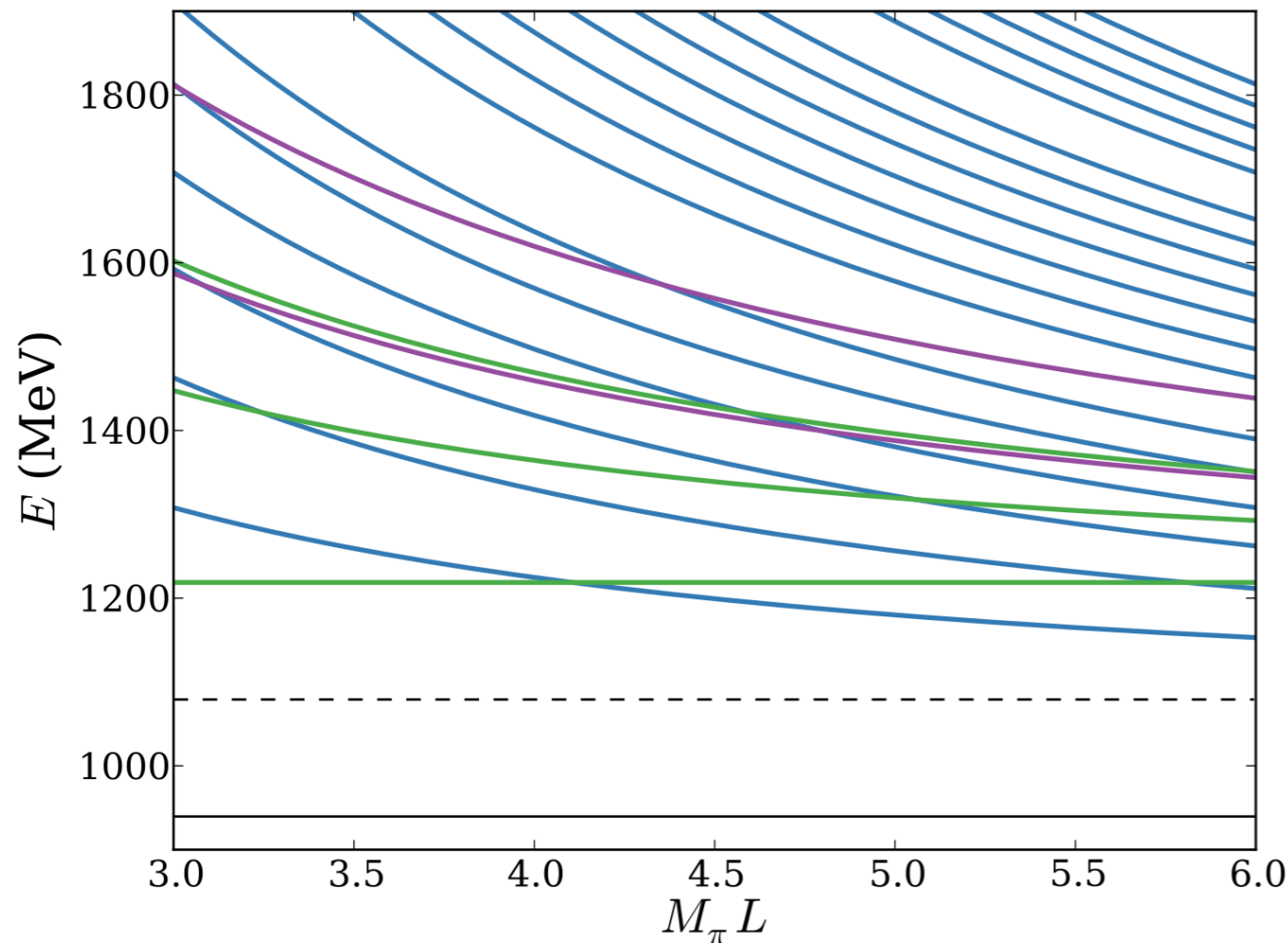
- Lattice “rule-of-thumb” is to take $M_\pi L \sim 4$ or larger
- Here $M_\pi L \sim 2.2$
- First three states are statistically consistent with non-interacting nucleons and pions
- Careful with operator overlaps... discretisation dependent, no guarantee of continuum limit

• Lang, Leskovec, Padmanath, Prelovsek (2017) •

Roper resonance

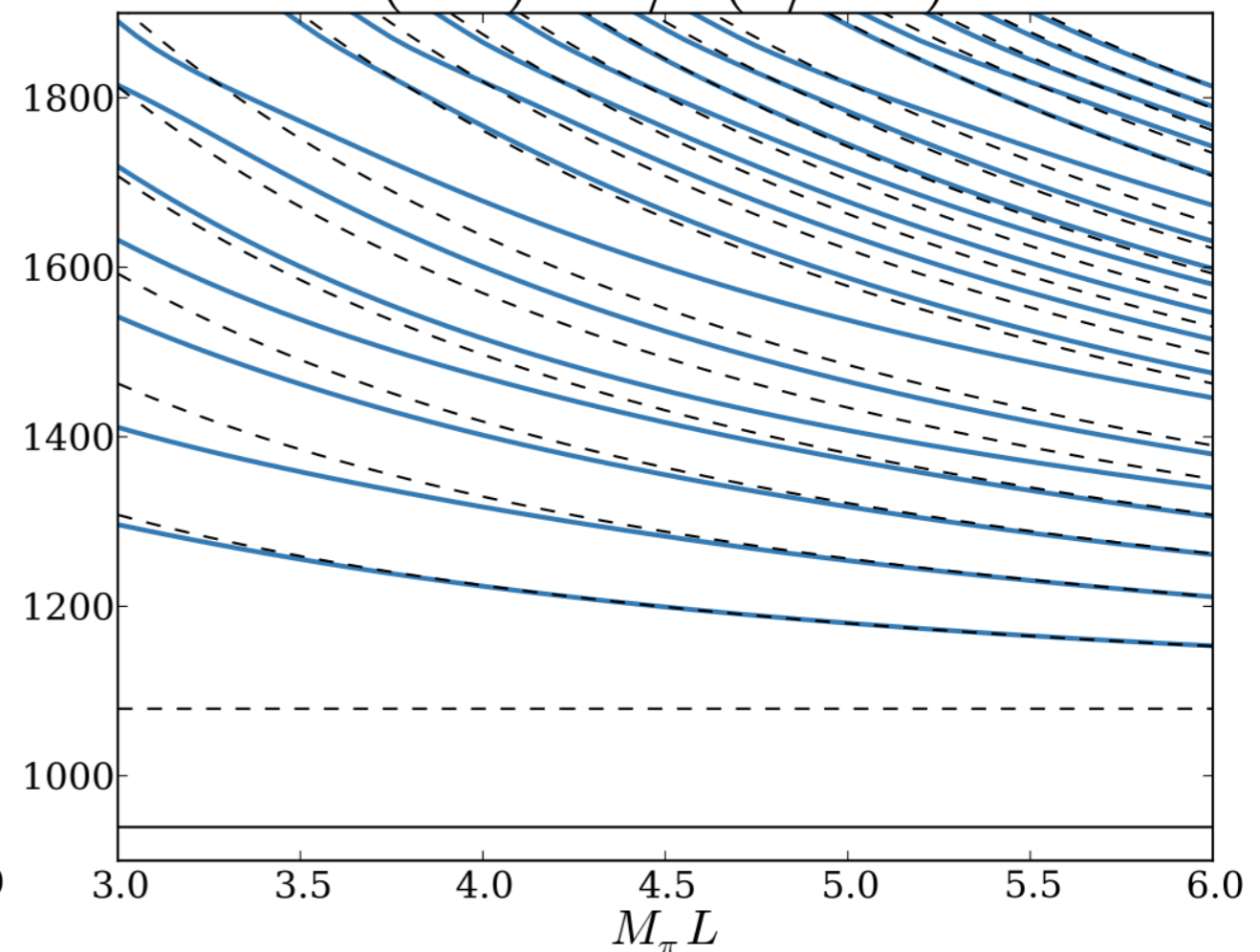
□ Naive spectra from MTH and Meyer 2016

○ Non-interacting



○ GWU WI08 + Lüscher

$$I(J^P) = 1/2(1/2^+)$$

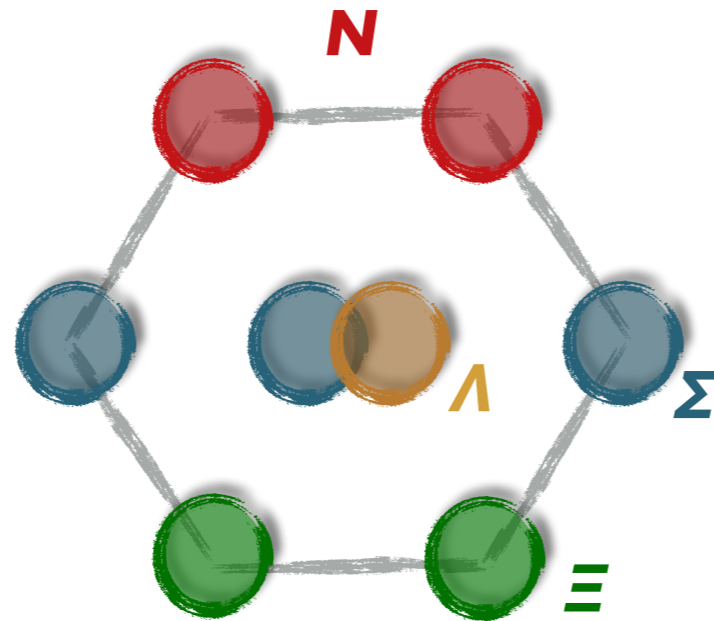


□ But, right panel ignores three-body effects!... I would say work is needed

- MTH and Meyer (2016)
- see also Döring et al. (2013)
- [Daniel Severt LATT2022](#)
-

$$\Delta \rightarrow N\pi$$

- Andersen et al. 2018
- Andersen et al. 2019
- Silvi et al. 2021
- Pittler et al. 2021
- Bulava et al. 2022



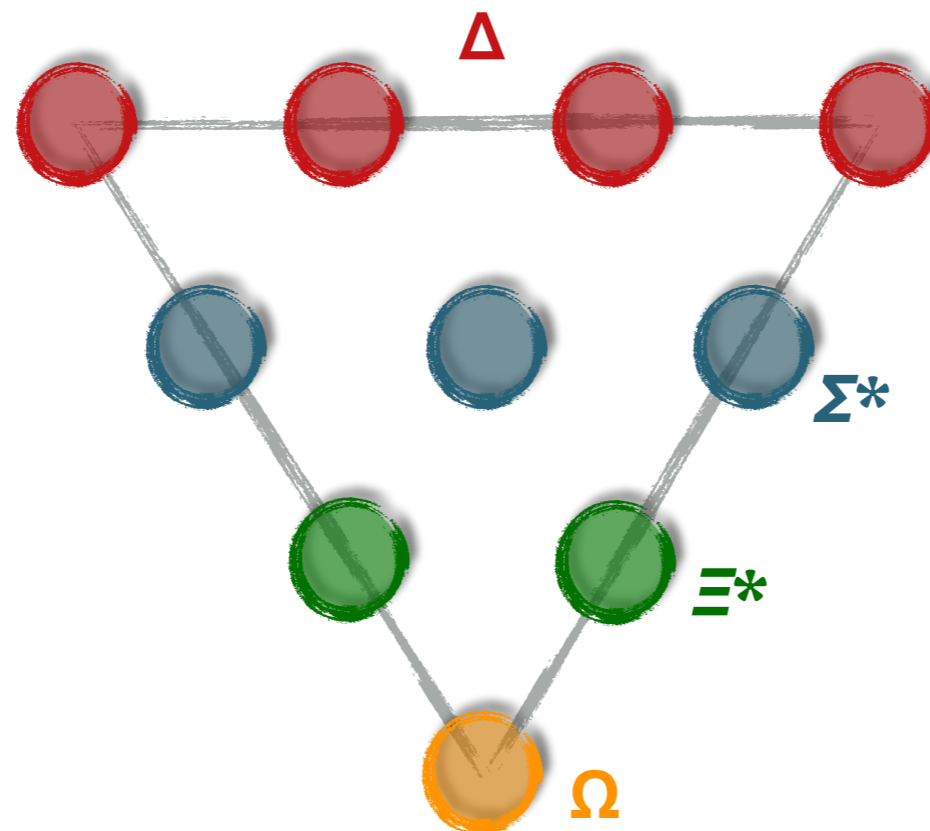
(focusing here on studies with scattering states)

Baryons are difficult...

Stay tuned!

$$N^* \rightarrow N\pi$$

- Lang et al. 2017
- Wu et al. 2017
- Kiratidis et al. 2017



$$\Lambda \rightarrow \bar{K}N$$

- Hall et al. 2015

See also...

- Detmold and Nicholson 2015
- Wu et al. 2018
- Xing & Liu, LATT2022 (in prep)