

Accessing entanglement and suppressing background in semileptonic $H \rightarrow WW$

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- **The W pair can effectively be seen as an entangled qutrit system which is rare and interesting compared to the well studied and more usual qubit entanglement.**
- **This is a brand new way to look at the Higgs!**

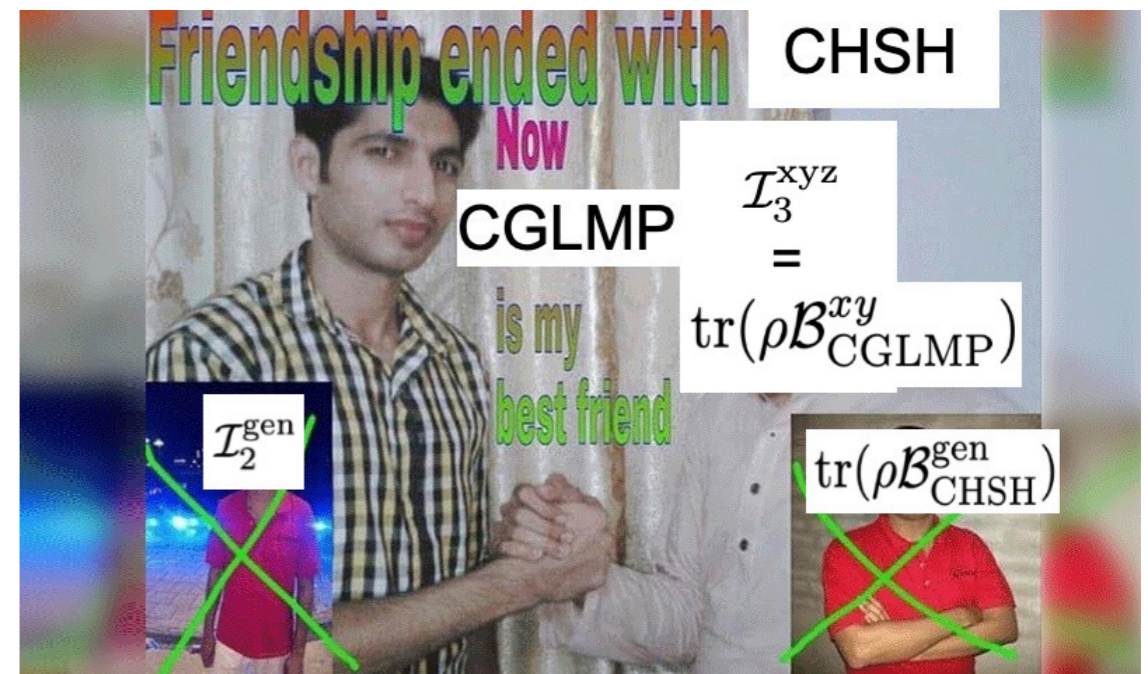
- Everything here is shamelessly stolen from Alan's paper: <https://arxiv.org/pdf/2106.01377.pdf>

$$\begin{aligned} \text{tr}(\rho \mathcal{B}_{\text{CGLMP}}^{xy}) &= \frac{8}{\sqrt{3}} \langle \xi_x^+ \xi_x^- + \xi_y^+ \xi_y^- \rangle_{\text{av}} \\ &+ 25 \langle ((\xi_x^+)^2 - (\xi_y^+)^2) ((\xi_x^-)^2 - (\xi_y^-)^2) \rangle_{\text{av}} \\ &+ 100 \langle \xi_x^+ \xi_y^+ \xi_x^- \xi_y^- \rangle_{\text{av}} \end{aligned}$$

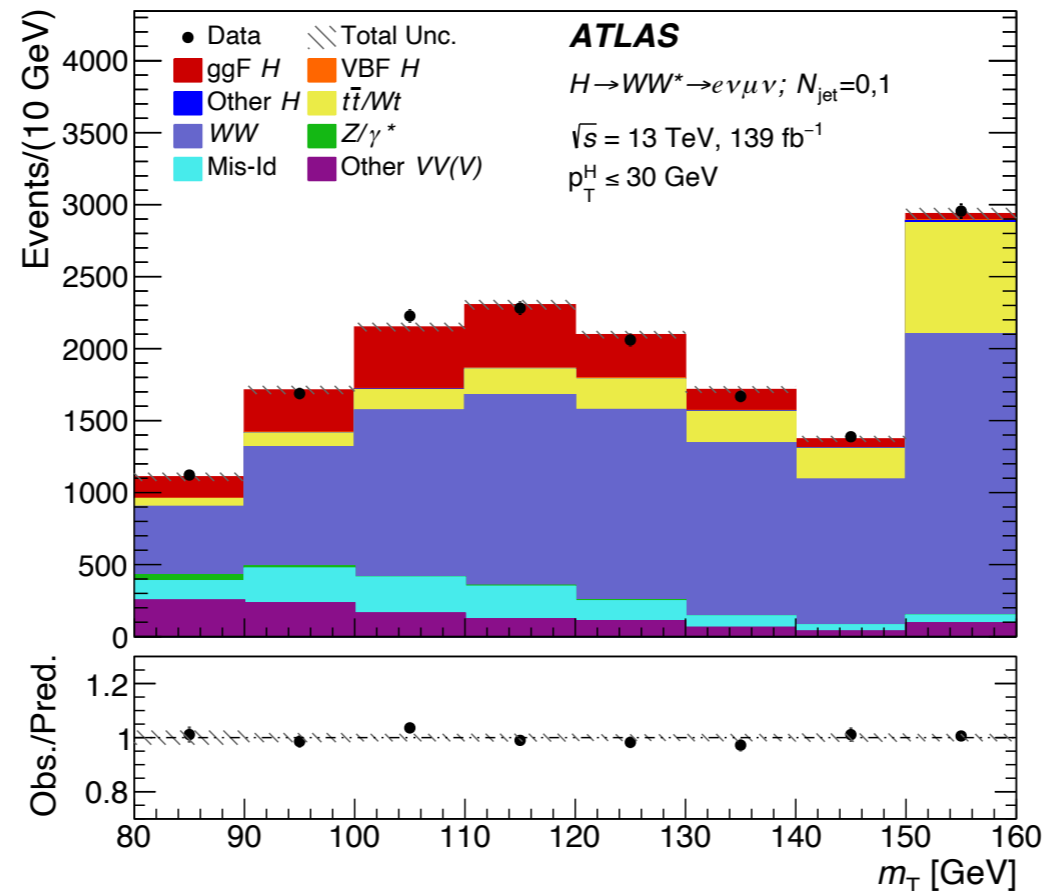
$$\begin{aligned} \xi_i^\pm &= \hat{\mathbf{n}}_i \cdot \mathbf{n}_{l^\pm} \\ i &= x, y, z \end{aligned}$$

$$\mathcal{I}_3^{\text{xyz}} = \max(\langle \mathcal{B}_{\text{CGLMP}}^{xy} \rangle, \langle \mathcal{B}_{\text{CGLMP}}^{yz} \rangle, \langle \mathcal{B}_{\text{CGLMP}}^{zx} \rangle)$$

- We have entanglement if $\mathcal{I}_3 > 2$.

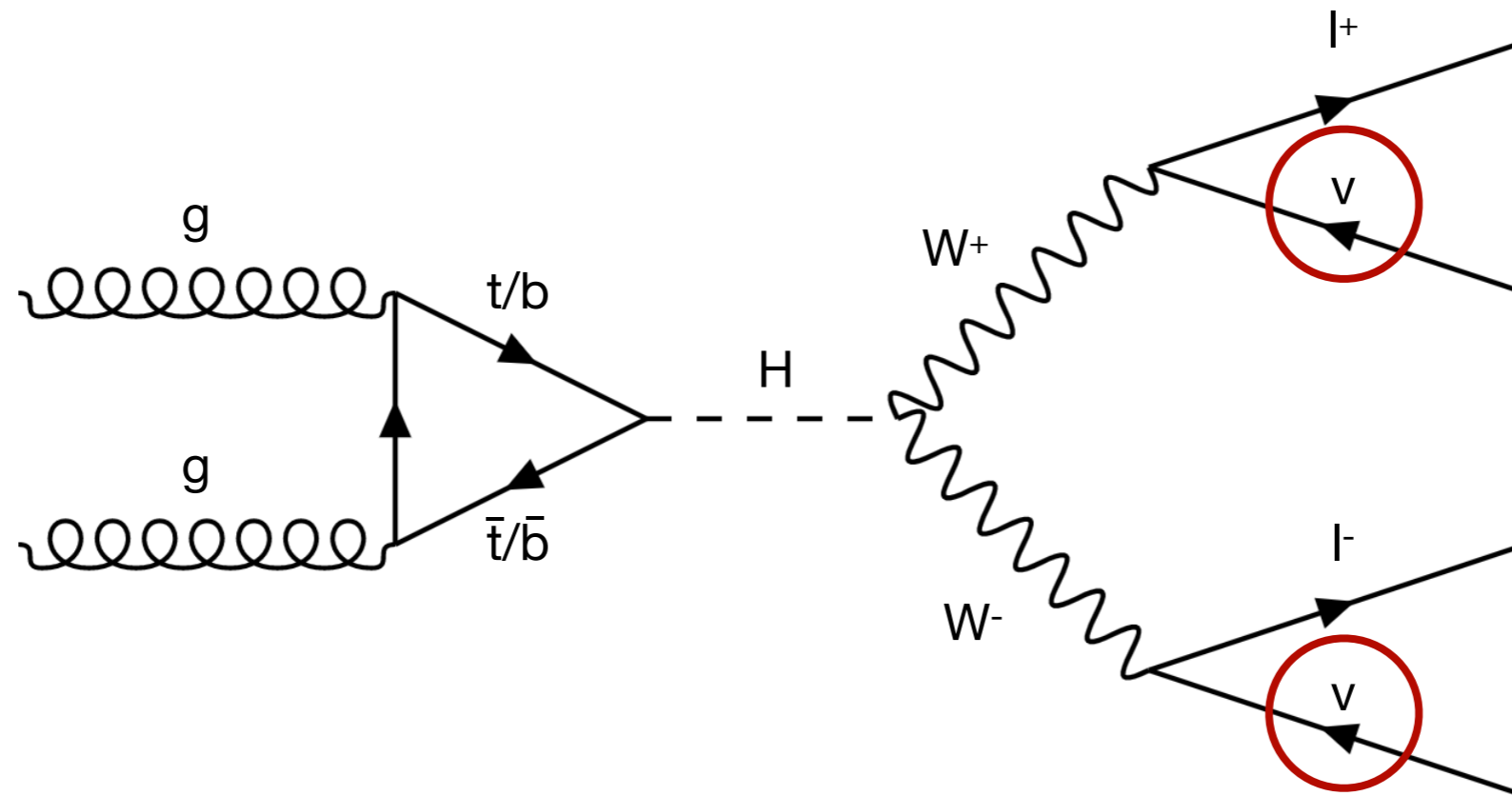


- Generally, $H \rightarrow WW^*$ is only studied in the dilepton channel:



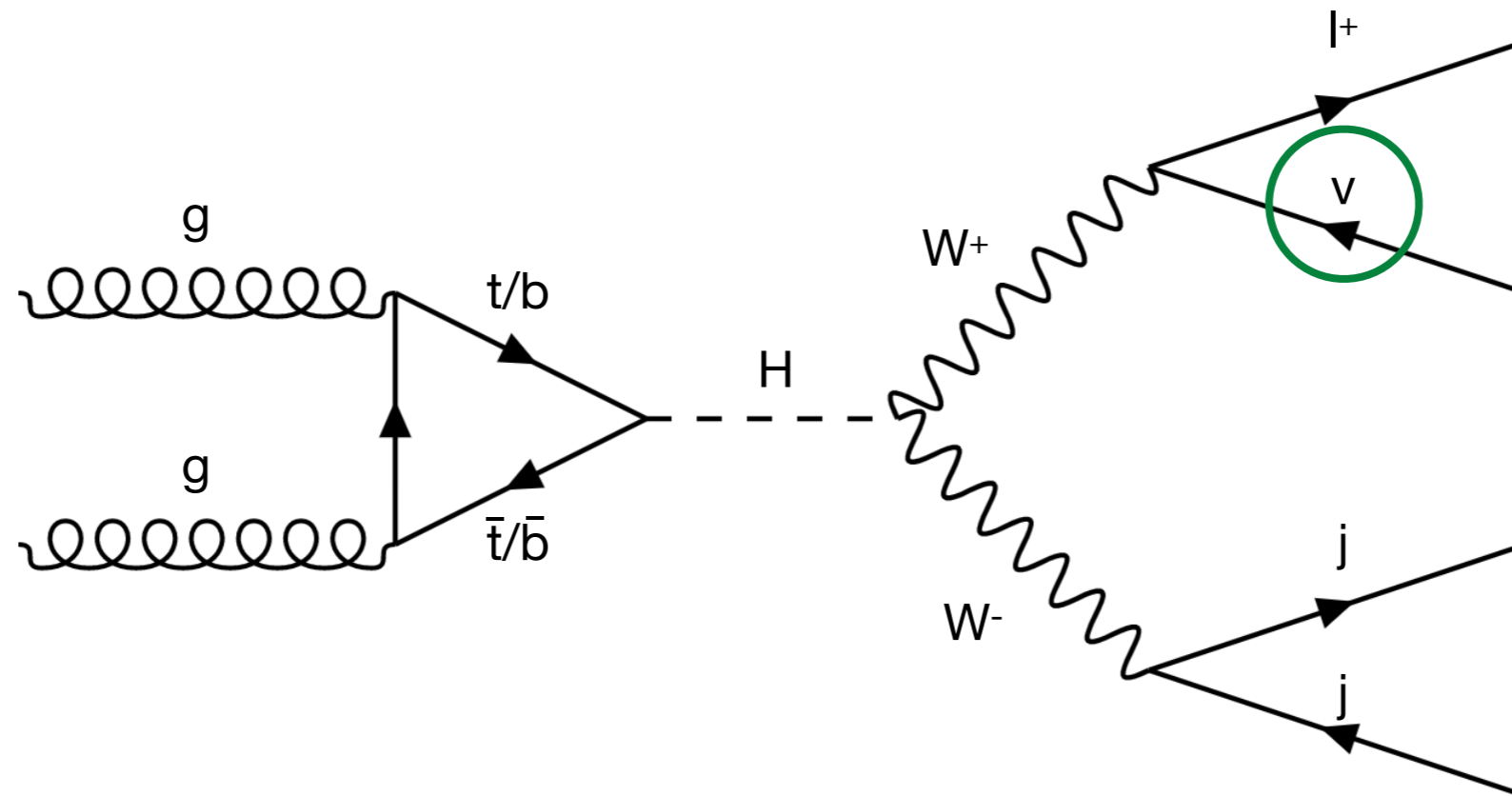
- The semi- and fully- hadronic decay modes are too difficult to isolate from W +jets and QCD backgrounds.

- So, why not use $H \rightarrow WW^* \rightarrow l\nu l\nu$ for Quantum information measurements?



- **Presence of two neutrinos** (and lack of additional mass constraints due to off-shell W) **makes it extremely hard to reconstruct the Higgs.**

- Using the $H \rightarrow WW^* \rightarrow jj\nu$ channel we can reconstruct the Higgs:



- But there are two challenges:
 1. Dealing with overwhelming W^+ -jets background.
 2. Accessing the spin information of the hadronic W .

- **W spin information can be accessed either using charged leptons or down-type quarks:**

“Spin Analysing Power”

degree to which a particle carries parent spin info, 1 = fully, 0 = none, -1 = ‘anti’ fully

	ℓ	u, c	d, s
LO	1.000	-0.310	1.000
NLO	0.998	-0.310	0.930

- **If we can identify the down-type jet in a hadronic W decay, we can access the spin information.**
- **$W \rightarrow cs$ allows us to do this because we can tag the c-flavoured jet and then take the other jet that pairs to make the W mass as our s-jet (spin analyser).**

- **We use the following MC generators to simulate signal and background processes:**
 - ➔ **Higgs (gg fusion): Powheg + Pythia8**
 - ➔ **W+jets: Powheg + Pythia8**
 - ➔ **ttbar: Powheg + Pythia8**
 - ➔ **Diboson: Sherpa**
- **We generate about 900k signal events and ~1 mil background events for each process.**

- Cuts and reconstruction/trigger efficiencies are applied to jets and leptons to simulate detector effects:

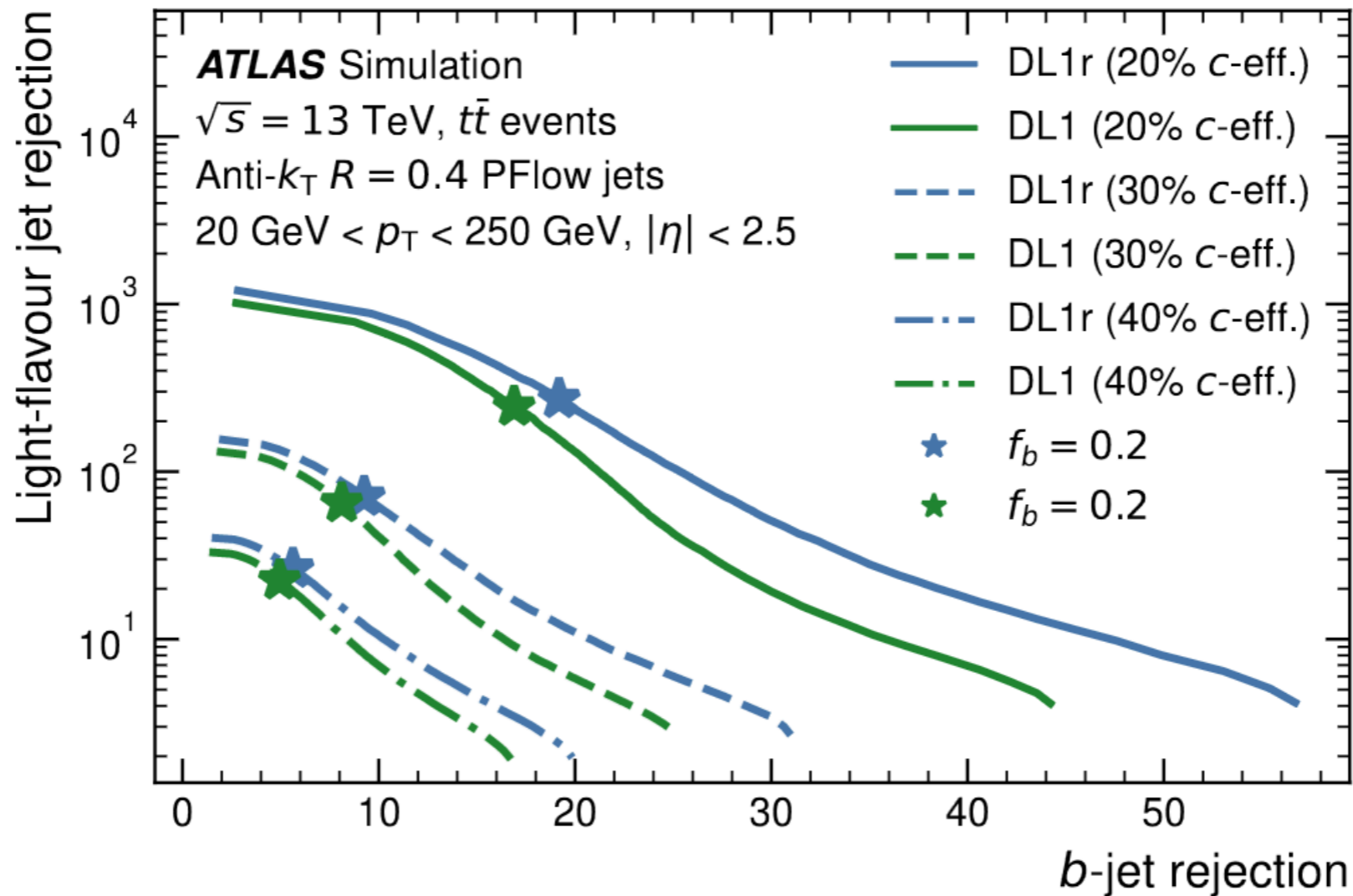
$$p_t > 25 \text{ GeV} \quad |\eta| < 5$$

- Jets are ‘tagged’ to approximate realistic c- and b-tagging efficiencies.

Efficiency type	Efficiencies (%)	
	c-tagger	b-tagger
ϵ_b	14	77
ϵ_c	40	20
ϵ_l	3.3	0.8

- All of these together approximate what we call “reco level” in experiments like ATLAS and CMS.

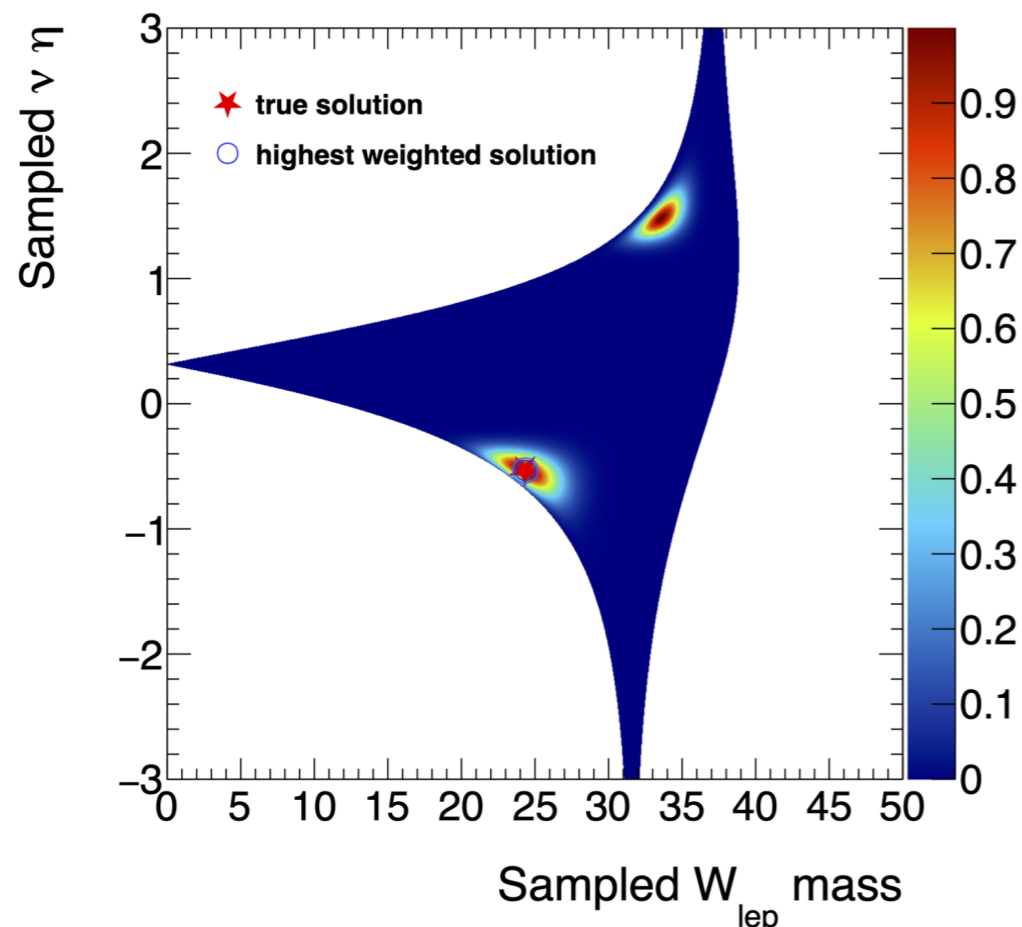
- Key question is how well can we tag charm jets at ATLAS and CMS?



Taken from <https://arxiv.org/pdf/2211.16345.pdf>

- Using $W \rightarrow cs$ decays requires the hadronic W to be on-shell and therefore the leptonic to be off-shell.
- We have two unknowns remaining:
 1. The long. component of the neutrino momentum
 2. The inv. mass of the off-shell W boson ($m_W < 80$ GeV)
- This is similar to a problem in $t\bar{t}$ final states, where we use a tool called ‘Neutrino Weighting’ to reconstruct events by integrating over missing kinematics.
- In this case, we integrate over 1. and 2.

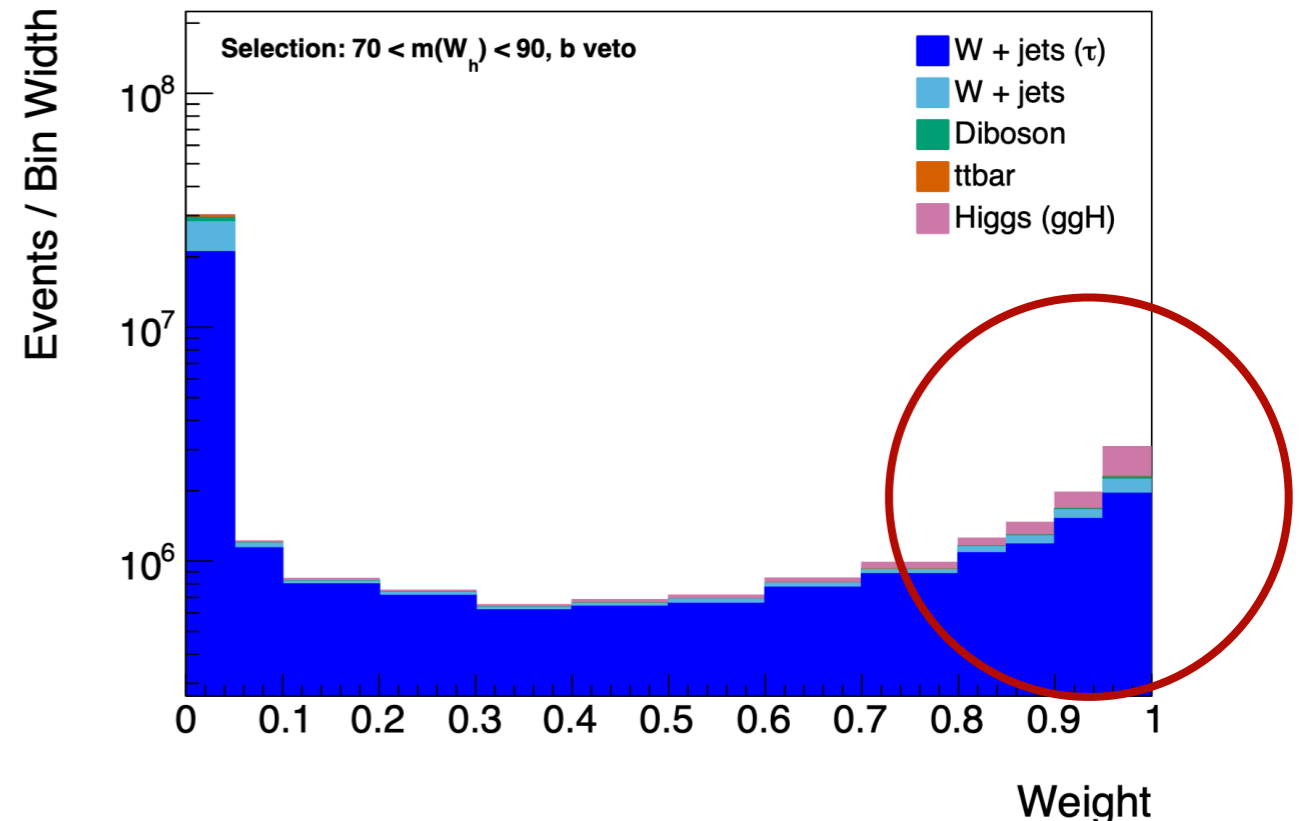
- We reconstruct many Higgs each event under different assumptions of m_{W^*} and η_{ν} .



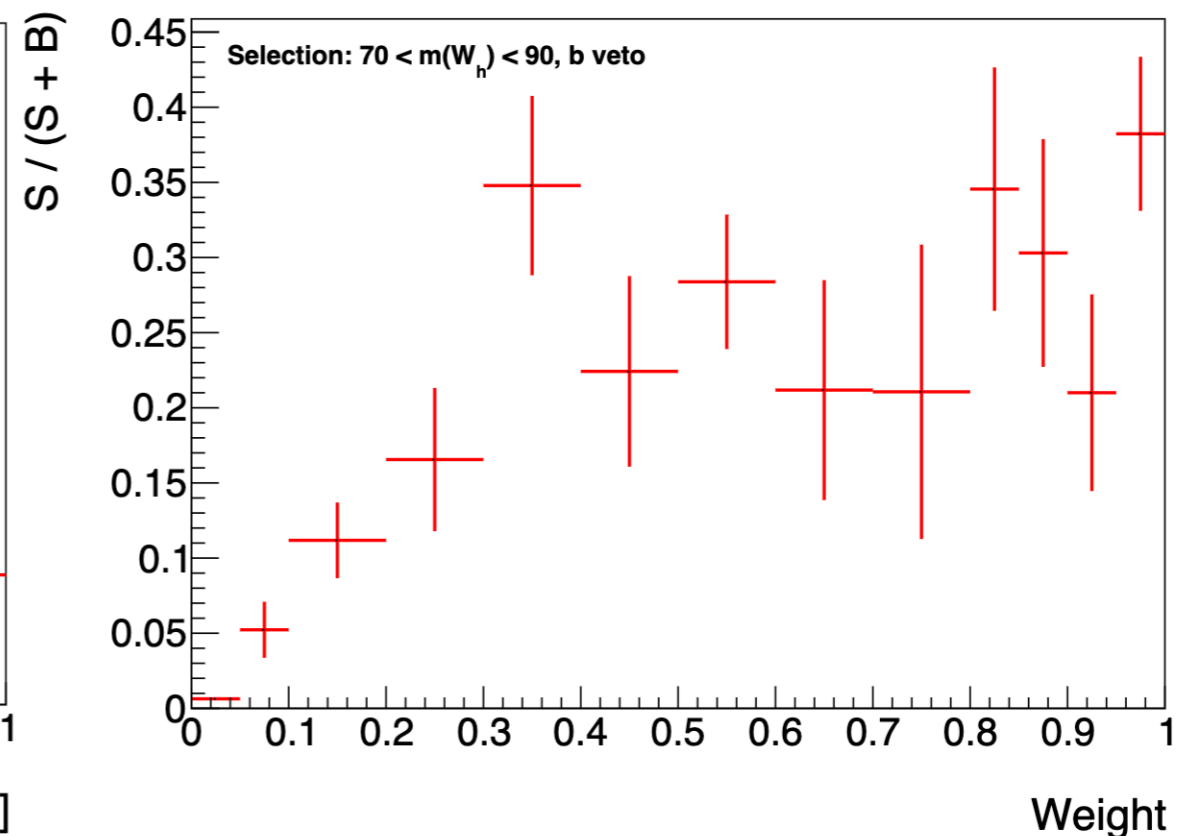
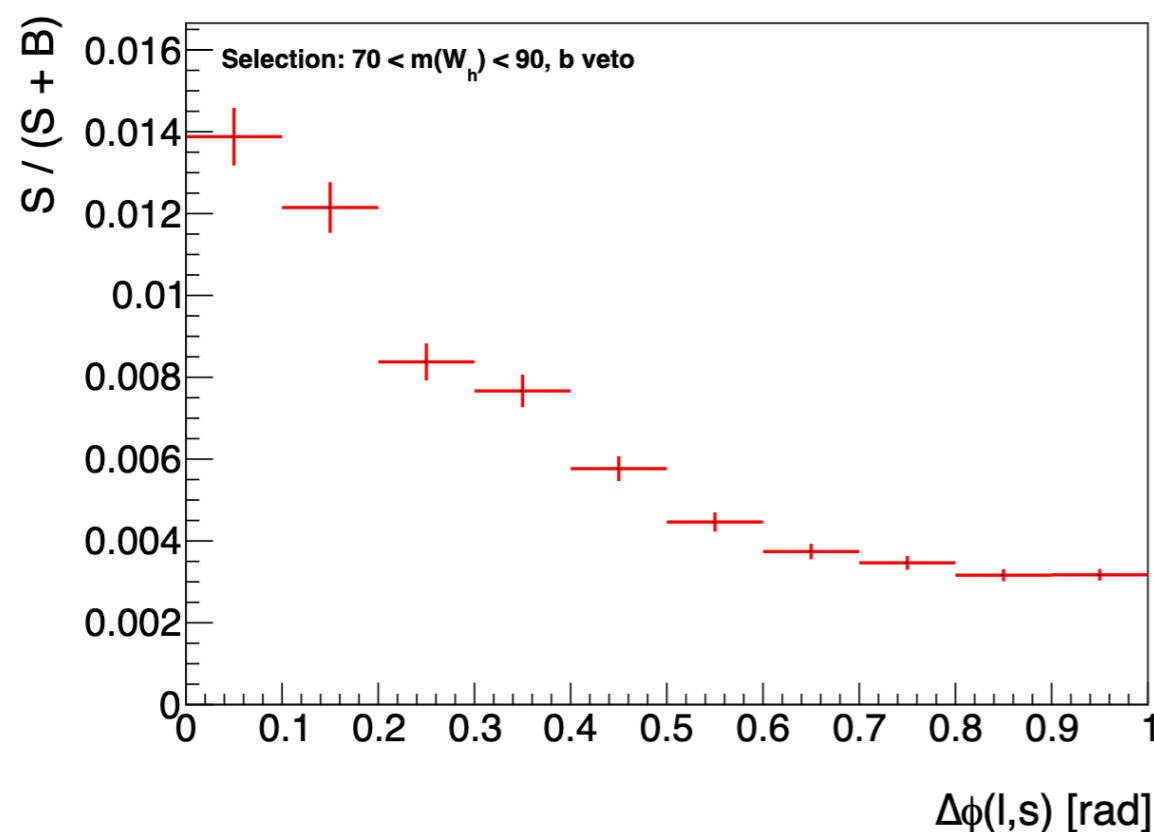
- Each solution yields a weight based on how well the reconstructed ν agrees with the observed MET (missing transverse momentum in the event).

- The solution with the highest weight is taken as the correct one.

- Using $W \rightarrow cs$ decays requires the hadronic W to be on-shell and therefore the leptonic to be off-shell.
- This is a bonus because the primary background ($W \rightarrow lvjj$) has an on-shell leptonic W .
- The weight from NW also acts similarly to an MVA classifier (signal events are more likely to result in higher weights than background events).

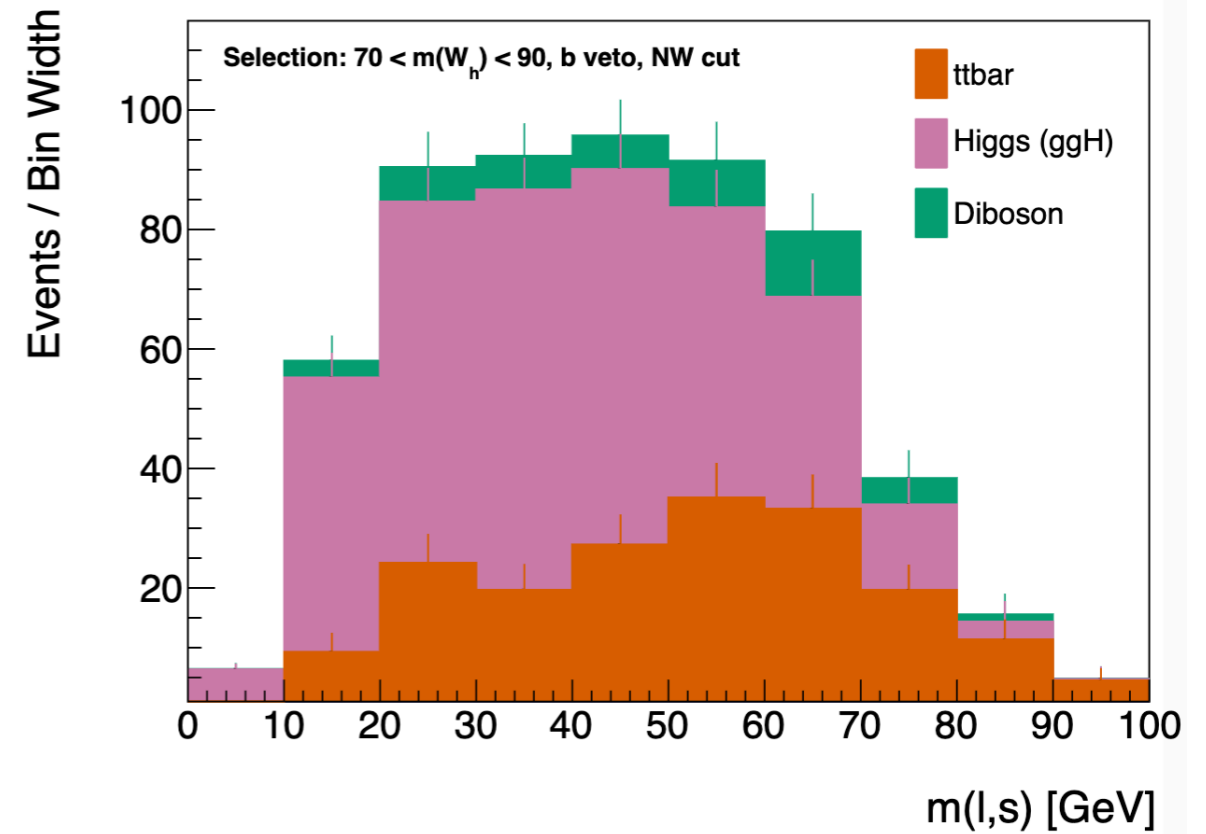
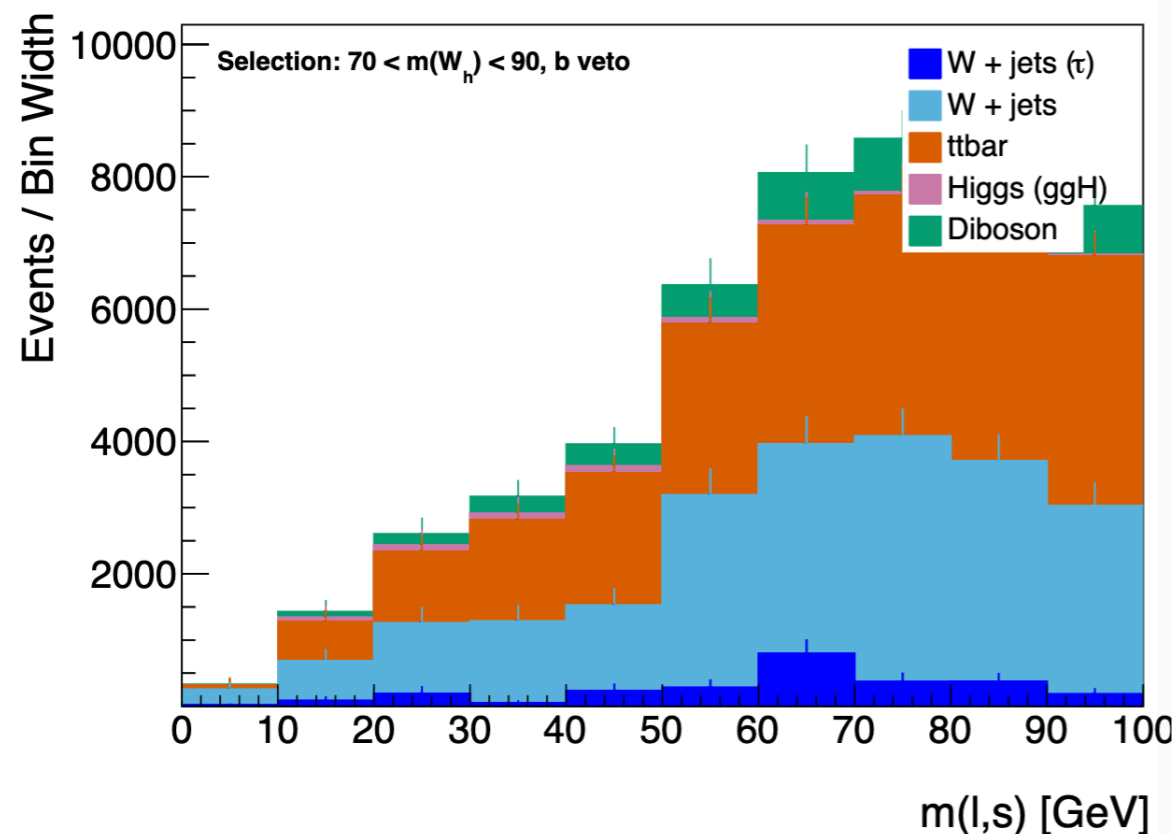


- We identified two parameters to cut on to enhance signal: ($\Delta\Phi_{ls}$) and the weight from the NW.



- Full analyses in ATLAS in CMS can optimise these and find additional sensitive selection cuts.

- The invariant mass of the spin analysing pair (l,s) before and after NW and $\Delta\Phi$ cut. This is at 139 fb^{-1} .



- Can see that the backgrounds are significantly suppressed and the signal becomes more dominant.

- **With these kinematic & NW cuts, we obtain good S/B** (considering this was, previously, an impossible channel!)

Process	idealised		Expected events				$\epsilon_c = 40\%$		
			idealised	lepton eff					
$W + \text{jets}$	42971	± 471	2774	± 112	2365	± 97			
$WW/WZ/ZZ$	5262	± 102	592	± 34	292	± 24			
$t\bar{t}$	1796	± 42	405	± 20	2656	± 52			
Higgs	38236	± 196	6334	± 80	2986	± 55			
S/(S+B)	0.43		0.62				0.36		

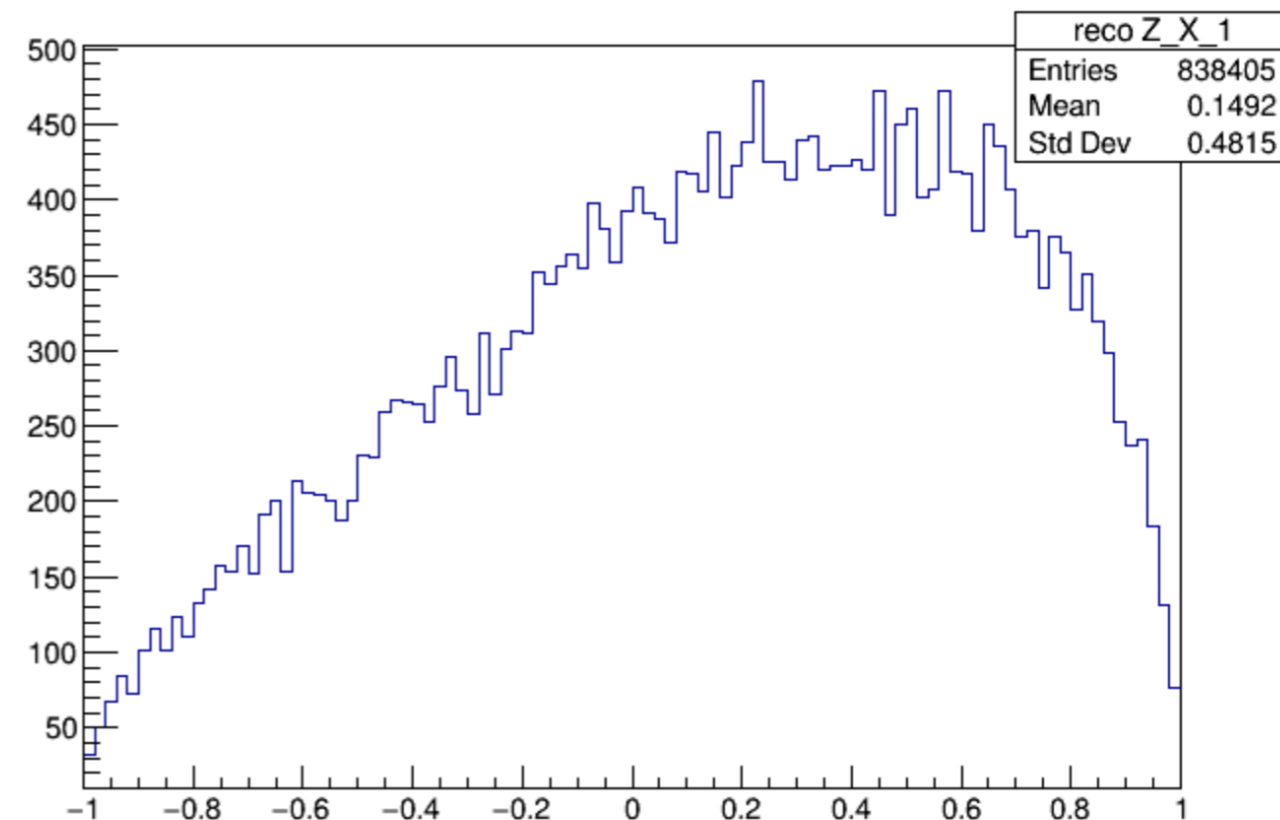
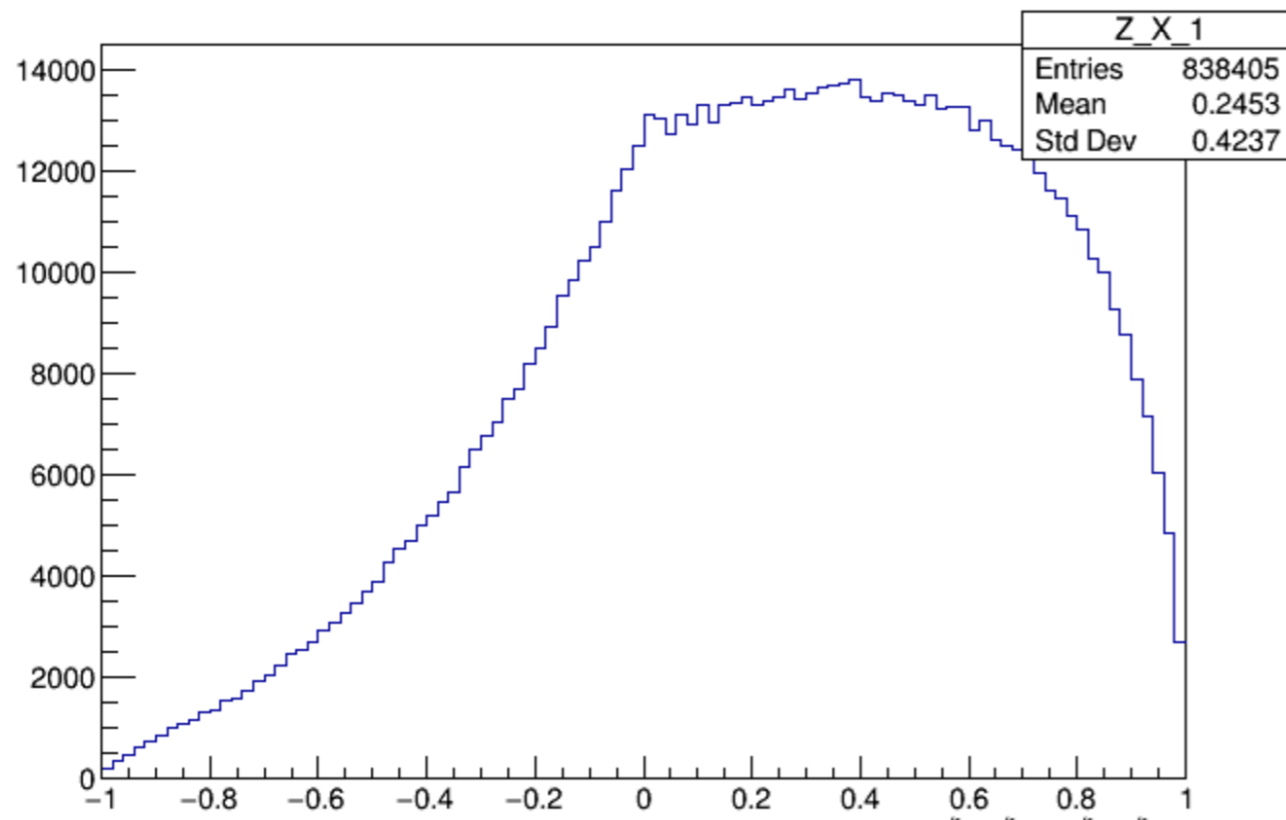
- **Can be used for entanglement measurements, but also a brand new Higgs decay topology for free!**

- Inequality becomes three observables:

$$\begin{aligned} \text{tr}(\rho \mathcal{B}_{\text{CGLMP}}^{xy}) &= \frac{8}{\sqrt{3}} \langle \xi_x^+ \xi_x^- + \xi_y^+ \xi_y^- \rangle_{\text{av}} && \leftarrow \text{1st} \\ &+ 25 \langle ((\xi_x^+)^2 - (\xi_y^+)^2) ((\xi_x^-)^2 - (\xi_y^-)^2) \rangle_{\text{av}} && \leftarrow \text{2nd} \\ &+ 100 \langle \xi_x^+ \xi_y^+ \xi_x^- \xi_y^- \rangle_{\text{av}} && \leftarrow \text{3rd} \end{aligned}$$

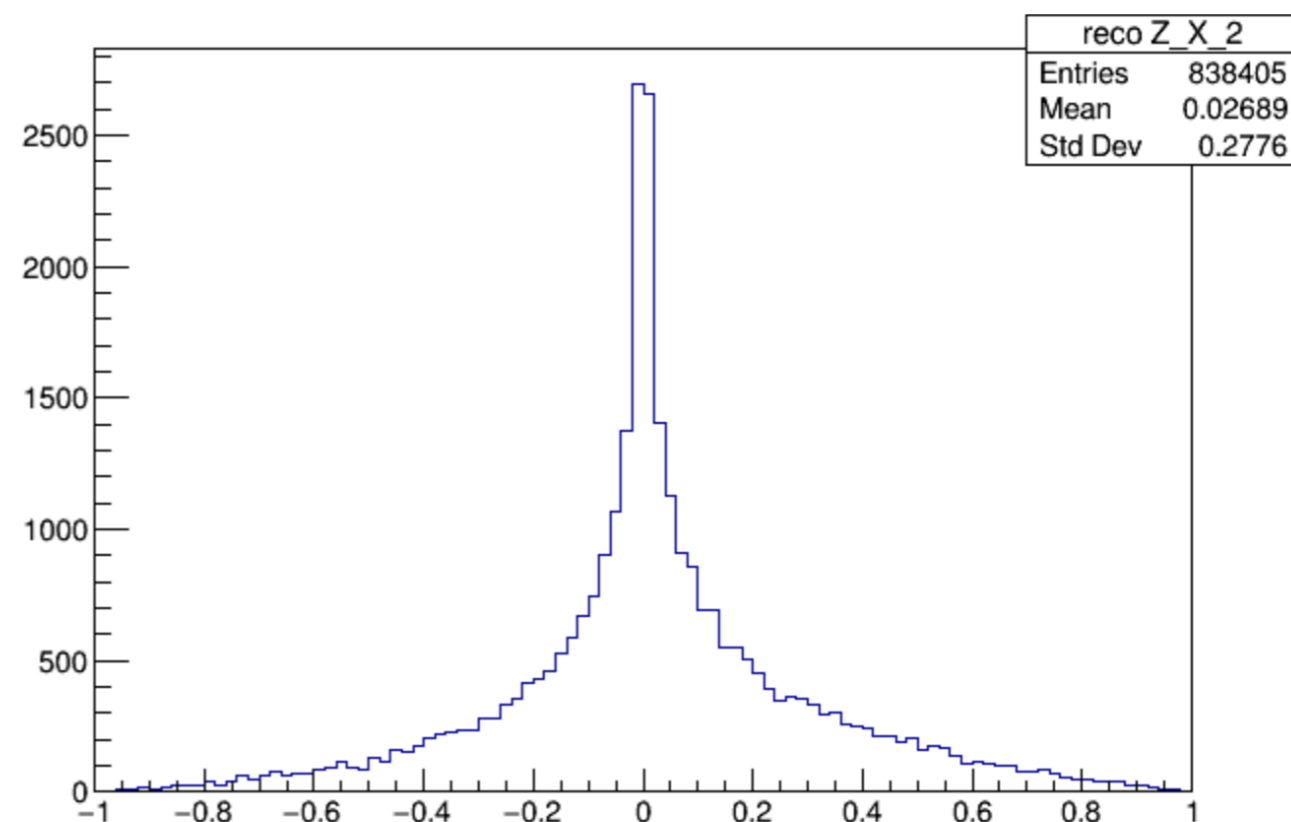
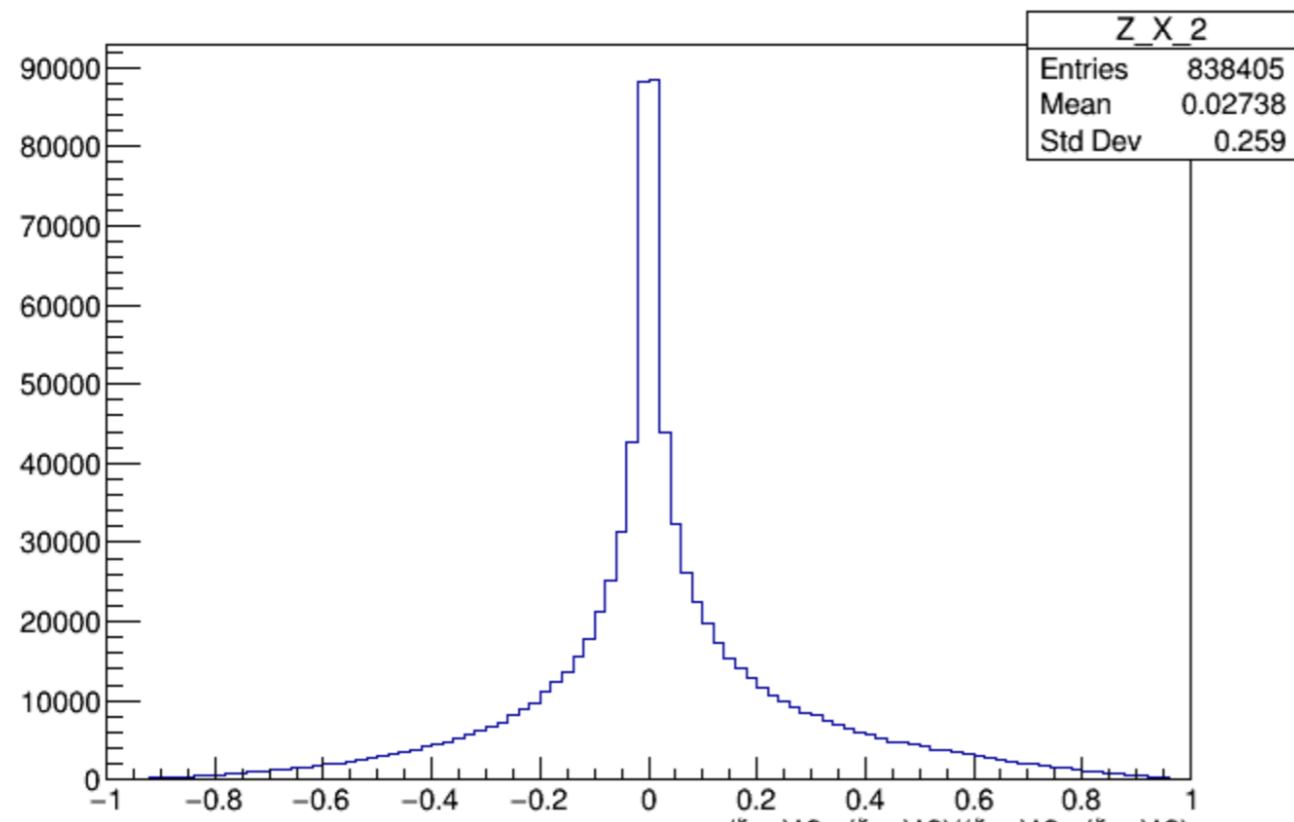
First observable

Left is the parton level and right our reco level



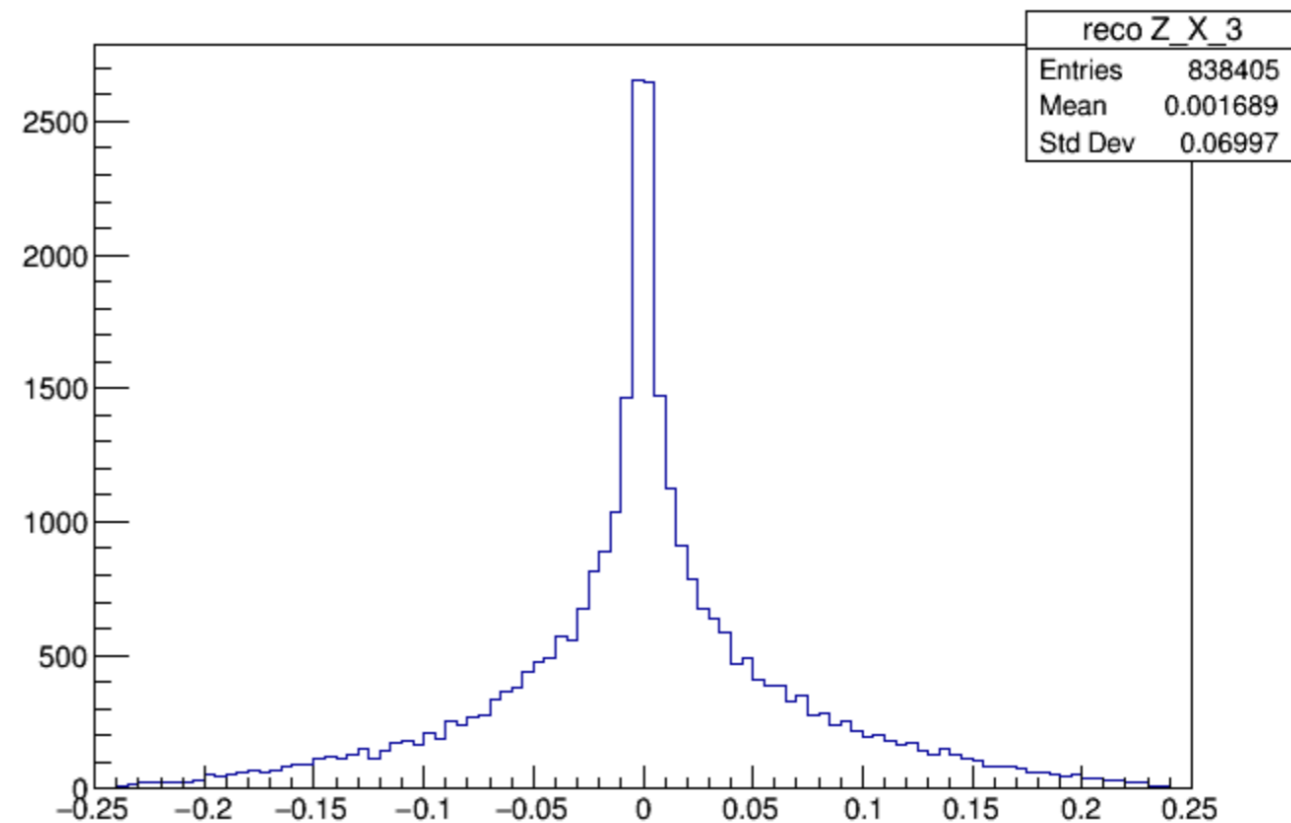
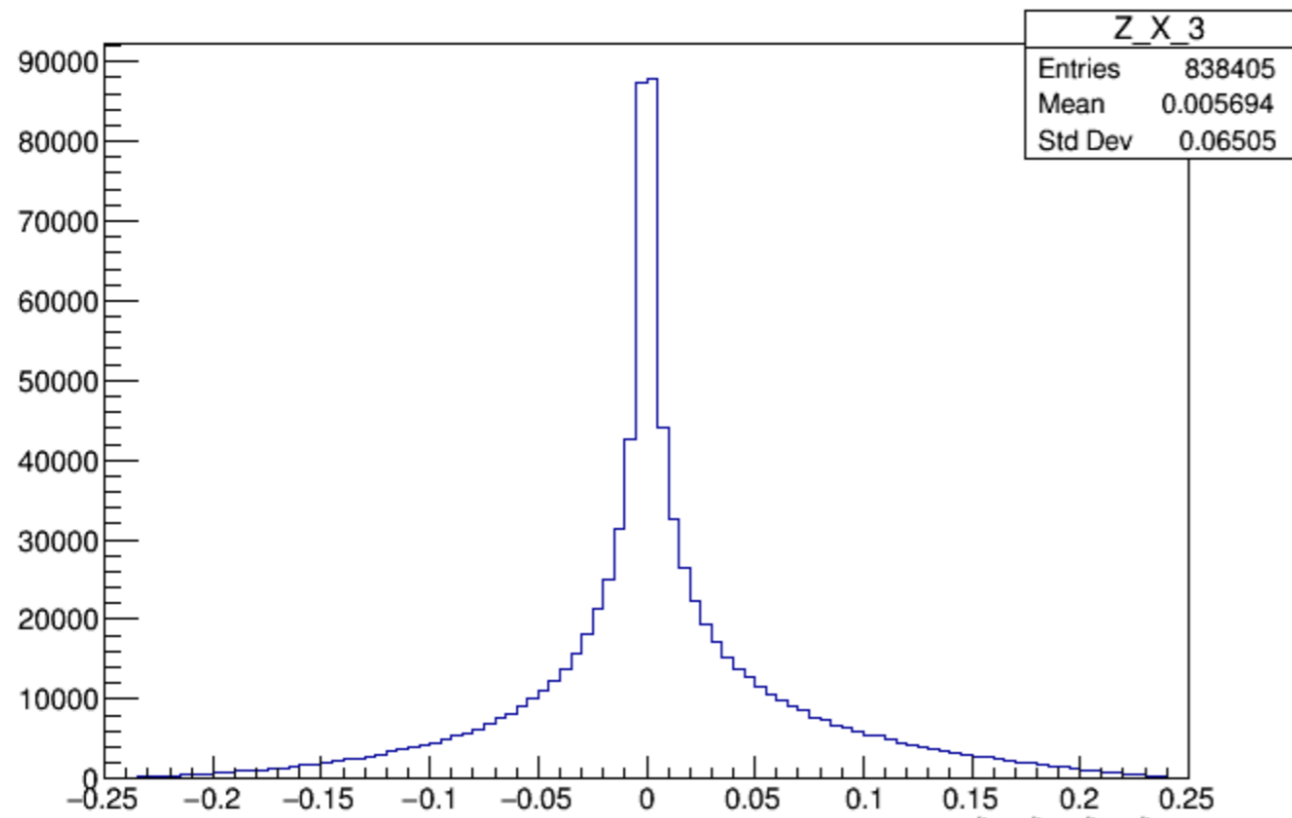
Second observable

Left is the parton level and right our reco level

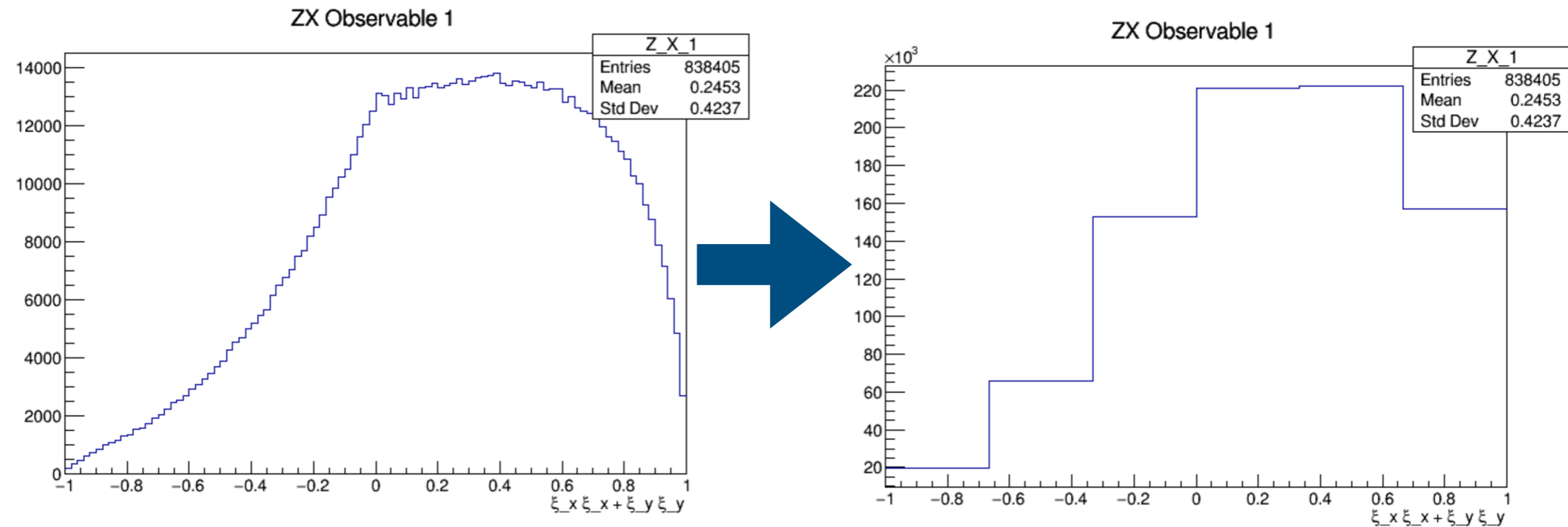


Third observable

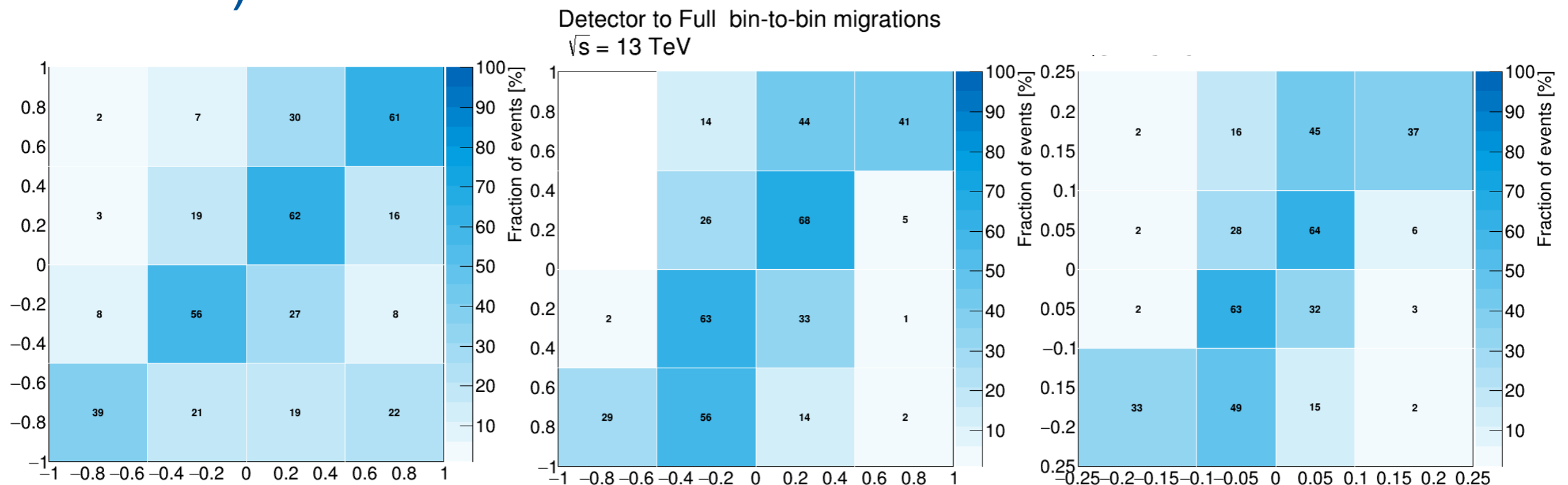
Left is the parton level and right our reco level



- In practice, can't use this many bins in unfolding.
- We reduce them to a more realistic number. E.g.



- We remove ‘detector’ effects using Bayesian Iterative Unfolding (signal deconvolution explained by Ethan).



- Pretty good approximation of what is done in actual experiments.

- Use the unfolded distributions to extract the value for the CGLMP inequality at different integrated luminosities.

- At 139 fb⁻¹:

$$I_3^{xyz} = 2.45 \pm 0.86$$

For N fb⁻¹:

- At 300 fb⁻¹:

$$I_3^{xyz} = 2.45 \pm 0.57$$

$$\sigma \sim \sqrt{\frac{N}{139}} 0.86$$

- At 3000 fb⁻¹:

$$I_3^{xyz} = 2.44 \pm 0.19$$

- **The usage of NW can be improved with Machine Learning. Important to note that NW is somewhat idealised in this study.**
- **Charm tagging can be optimised in an actual analysis. We can expect better S/B ratio and more total signal events.**
- **Custom trigger could improve acceptance of events with low energy leptons.**

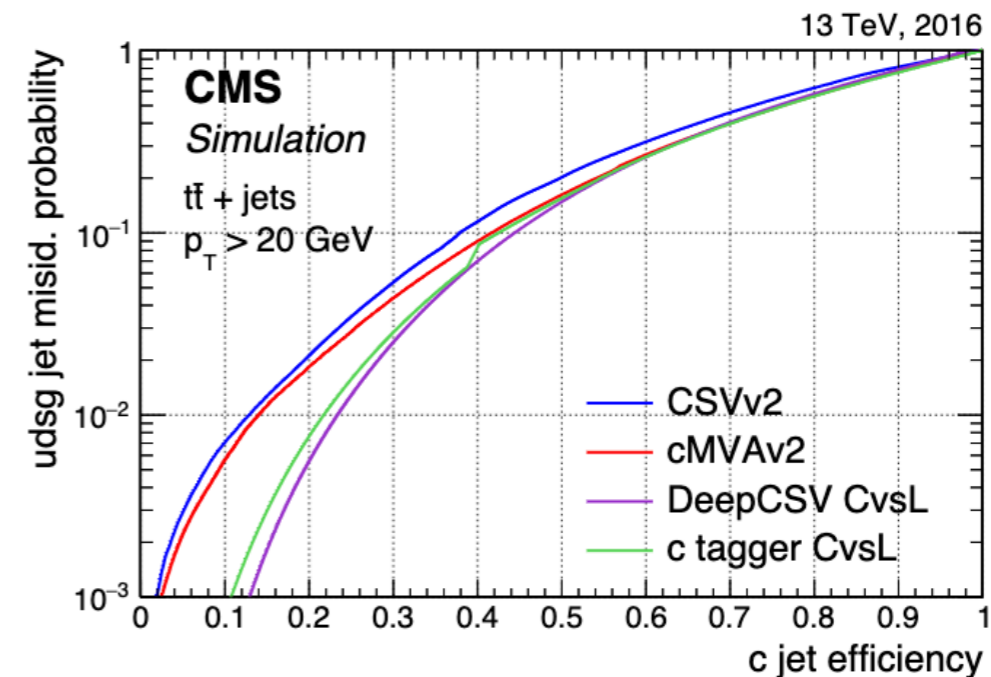
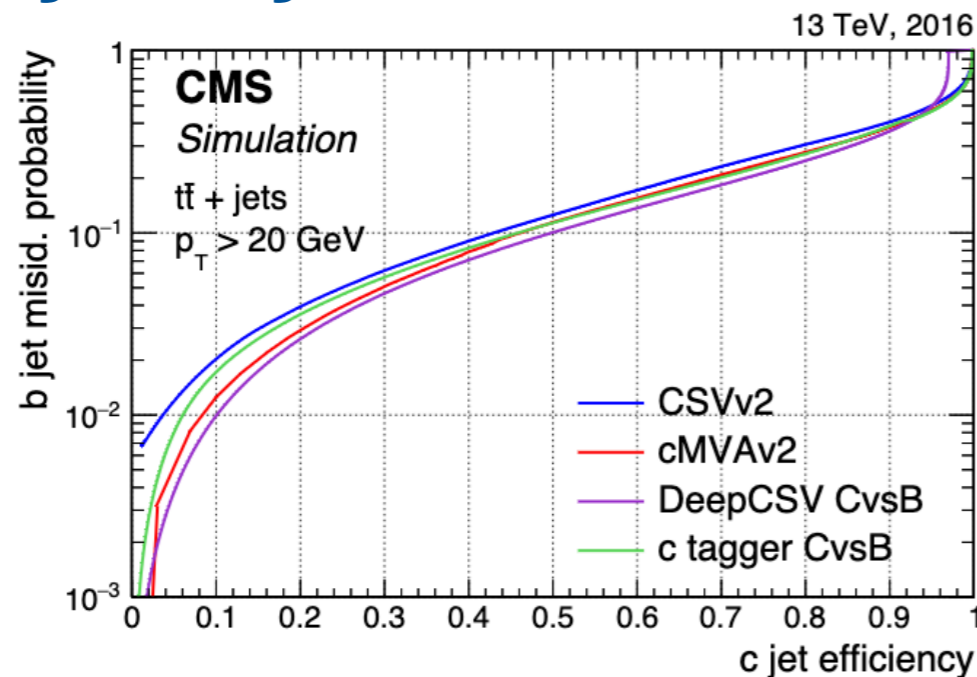
- **Potential for an entirely new Higgs decay mode!**
- **With this preliminary analysis, significance of $I_3 > 2$ is:**
 - ➔ **0.5 s.d. at 140 fb⁻¹**
 - ➔ **0.8 s.d. at 300 fb⁻¹**
 - ➔ **2.3 s.d. at 3000 fb⁻¹**
- **Entanglement in $H \rightarrow WW^* \rightarrow jjlv$ channel is challenging, but most of these challenges are experimental (and we're quite good at beating expectations for these types of problems).**

Backup

- How well can we tag charm jets at ATLAS and CMS?

Working point	ϵ_c (%)	ϵ_b (%)	ϵ_{udsg} (%)
c tagger L	88	36	91
c tagger M	40	17	19
c tagger T	19	20	1.2

Suspiciously high numbers that probably come at the cost of high jet rejection.



Taken from <https://arxiv.org/pdf/1712.07158.pdf>

- NW improvements

