### **Measurement of D(\*) production cross sections**



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### **Outline :**

### **Introduction & Motivation**

 $D^{*\pm}$ ,  $D^{\pm}$  and  $D_s^{\pm}$  reconstruction Towards x-section measurement  $D^{*\pm}$ ,  $D^{\pm}$  and  $D_s^{\pm}$  production x-sections Summary & Plans

(ATLAS-CONF-2010-034)

(ATLAS-CONF-2011-XXX)

# Introduction & Motivation

*D*-mesons are produced in c and b fragmentation c and b quark production are hard processes ( $m_Q >> \Lambda_{QCD}$ ) Theoretical calculations available till NLO+NLL level Still large theoretical uncertainties (scales, multiple interactions)

Reconstruction of D-mesons already feasible with first ATLAS data due to

- large cross-section values
- clean D-meson signatures
- precise ATLAS tracking and vertexing

Important to measure production of D-mesons

- to evaluate and calibrate tracking performance
- to compare production in pp and heavy ion collisions
- to test theoretical calculations
- to verify  $\rm m_{c}$  value and proton structure functions
- to realistically estimate charm contribution to backgrounds and trigger rates



## *D*<sup>(\*)</sup>-mesons reconstruction at 7 TeV

For *D*-meson reconstruction: 1.1 nb<sup>-1</sup> (March-July, minimum-bias triggers after prescale)

No dE/dx particle identification (not effective for high-p<sub>T</sub> tracks) Selection Strategy:

To select *D*-mesons one can utilise:

hard nature of charm production  $(p_T(D^{(*)}), p_T(K, \pi))$ hard nature of charm fragmentation  $(p_T(D^{(*)})/E_T)$ relatively large *D*-mesons' life-times  $(l_{XY})$ 



use widest kinematic range where signals can be measured

 $p_T(D^{(*)}) > 3.5 \,\text{GeV}, \ |\eta(D^{(*)})| < 2.1$ 

make signals as clean (significant) as possible in the kinematic range use cuts (their values) which do not produce large systematic uncertainties

### To tune actual cut values realistic MC has been used

### Figures from ATLAS-CONF-2010-034



### Unbiased trigger selection

For data starting from period B,

EF slots should be used to take into account all prescales

Largest biased slot in MinBias stream:

EF\_mbSpTrkMh\_MB2 ( $\implies$  L1\_MBTS\_2) high-multiplicity selection

Largest unbiased slot in MinBias stream:

 $EF\_L1ItemStreamer\_L1\_MBTS\_1 \ (\Longrightarrow L1\_MBTS\_1)$ 

Second useful unbiased slot in MinBias stream:

EF\_mbSPTrk ( $\implies$  L1\_RD0\_FILLED), add 10 – 15% to the previos one

Other slots are marginal and not 100% efficient

 $(\Longrightarrow L1\_MBTS\_1\_1, L1\_MBTS\_4\_4)$ 

We can use:

 $EF\_L1ItemStreamer\_L1\_MBTS\_1 ~||~ EF\_mbSpTrk$ 

 $\mathcal{L}_{1||2} = \mathcal{L}_{tot} * (1. - (1. - 1./N_1^{presc}) \cdot (1. - 1./N_2^{rand}/N_2^{presc}))$ 

### Data sample and luminosity

DATA selection for x-section measurement GRL BPhys.tracking\_noBS Period A, runs 152166-153200, May's reprocessing L1\_MBTS\_1  $\mathcal{L} = 417.8 \,\mu b^{-1}$ , OffLumi-7TeV-002 Period B, runs 153565-155160, May's reprocessing EF\_L1ItemStreamer\_L1\_MBTS\_1 || EF\_mbSpTrk  $\mathcal{L} = 514.6 \,\mu b^{-1}$ , OffLumi-7TeV-002 Period C, runs 155228-156682, T0 processing EF\_L1ItemStreamer\_L1\_MBTS\_1 || EF\_mbSpTrk  $\mathcal{L} = 73.3 \,\mu b^{-1}$ , OffLumi-7TeV-002 Period D1, runs 158045-158392, T0 processing EF\_L1ItemStreamer\_L1\_MBTS\_1 || EF\_mbSpTrk  $\mathcal{L} = 31.2 \,\mu \mathrm{b}^{-1}$ , OffLumi-7TeV-002 Period D2, runs 158443-158582, T0 processing EF\_L1ItemStreamer\_L1\_MBTS\_1 || EF\_mbSpTrk  $\mathcal{L} = 35.3 \,\mu b^{-1}$ , OffLumi-7TeV-002  $\mathcal{L}_{ABCD1/2} = 1072.1 \,\mu b^{-1}$ 

Periods D3 (12.3  $\mu b^{-1}),$  D4 (13.1  $\mu b^{-1}),$  D5 (4.7  $\mu b^{-1})$  and D6 (3.4  $\mu b^{-1})$  not used

### $D^{*+} \rightarrow D^0 \pi_s^+ \rightarrow (K^- \pi^+) \pi_s^+ (+\text{c.c.})$ reconstruction

$$\begin{split} N_{hits}^{pixel}(K,\pi,\pi_s) &\geq 1 \\ N_{hits}^{SCT}(K,\pi,\pi_s) &\geq 4 \\ p_T(\pi_s) &> 0.25 \, \text{GeV}, \ p_T(K,\pi) > 1.0 \, \text{GeV}, \ |\eta(K,\pi,\pi_s)| < 2.5 \\ d_0^{PV}(\pi_s) &< 2.0 \, \text{mm}, \ z_0^{PV}(\pi_s) \sin \theta < 2.0 \, \text{mm} \\ M(K\pi) \text{ window: } (1.82 - 1.91) \implies (1.80 - 1.93) \text{ for } p_T(D^*) > 12 \, \text{GeV or } |\eta(D^*)| > 1.3 \end{split}$$

 $D^0$  vertexing:

 $\begin{array}{l} \mathbf{VKalVrt}, \ \chi^2(D^0) < 5, \ l_{XY} > 0 \\ d_0^{PV}(D^0) < 0.2 \ \mathbf{mm}, \ \ z_0^{PV}(D^0) \sin \theta < 0.5 \ \mathbf{mm} \ \ (z_0^{PV}(D^0) == z \ \mathbf{at} \ \mathbf{2D} \ \mathbf{DCA} \ (\mathbf{ATLAS \ default})) \end{array}$ 

kinematic range:

 $p_T(D^{*\pm}) > 3.5 \,\text{GeV}, \ |\eta(D^{*\pm})| < 2.1$ 

"hard fragmentation" cleaning:

$$\frac{p_T(D^{*\pm})}{E_T} > 0.02$$

### $D^+ \rightarrow K^- \pi^+ \pi^+ (+\text{c.c.})$ reconstruction

$$\begin{split} N_{hits}^{pixel}(K,\pi_{1,2}) &\geq 1, \quad N_{hits}^{SCT}(K,\pi_{1,2}) \geq 4 \\ p_T(K) &> 1.0 \, \text{GeV}, \quad p_T(\pi_{1,2}) > 0.8 \, \text{GeV}, \quad max(p_T(\pi_{1,2})) > 1.0 \, \text{GeV}, \quad |\eta(K,\pi_{1,2})| < 2.5 \\ \cos \theta^*(K) > -0.8 \end{split}$$

suppression of  $D^{*\pm}$  and  $D_s^+ \to \phi \pi^+ \to (K^- K^+) \pi^+$  (+c.c.) reflections: remove  $\Delta M_{1,2} < 153 \,\text{MeV}$  and  $|M(K^{\pm}, "K^{\mp}") - M(\phi)_{\text{PDG}}| < 8 \,\text{MeV}$ 

 $D^{\pm}$  vertexing:

VKalVrt,  $\chi^2(D^{\pm}) < 6$ ,  $l_{XY} > 1.2 \text{ mm}$  $d_0^{PV}(D^{\pm}) < 0.15 \text{ mm}$ ,  $z_0^{PV}(D^{\pm}) \sin \theta < 0.3 \text{ mm}$   $(z_0^{PV}(D^{\pm}) == z \text{ at } 2D \text{ DCA (ATLAS default)})$ 

kinematic range:

 $p_T(D^{\pm}) > 3.5 \,\text{GeV}, \ |\eta(D^{\pm})| < 2.1$ 

"hard fragmentation" cleaning:

$$\frac{p_T(D^{\pm})}{E_T} > 0.02$$

$$D_s^+ \rightarrow \phi \pi^+ \rightarrow (K^- K^+) \pi^+$$
 (+c.c.) reconstruction

 $N_{hits}^{pixel}(K_{1,2},\pi) \ge 1, \quad N_{hits}^{SCT}(K_{1,2},\pi) \ge 4$   $p_T(K_{1,2}) > 0.7 \,\text{GeV}, \quad p_T(\pi) > 0.8 \,\text{GeV}, \quad |\eta(K_{1,2},\pi)| < 2.5$   $\cos \theta^*(\pi) < 0.4, \quad |\cos \theta'(K)|^3 > 0.2$   $\frac{dN}{d\cos \theta'(K)} \propto \cos^2 \theta'(K) \Longrightarrow \frac{dN}{d\cos^3 \theta'(K)} \propto const$ 

 $\mathbf{select}~|M(K^+K^-)-M(\phi)_{\rm PDG}|<6\,\mathbf{MeV}$ 

 $D_s^{\pm}$  vertexing:

VKalVrt,  $\chi^2(D_s^{\pm}) < 6$ ,  $l_{XY} > 0.4$  mm  $d_0^{PV}(D_s^{\pm}) < 0.15$  mm,  $z_0^{PV}(D_s^{\pm}) \sin \theta < 0.3$  mm  $(z_0^{PV}(D_s^{\pm}) == z$  at 2D DCA (ATLAS default))

kinematic range:

 $p_T(D_s^{\pm}) > 3.5 \,\text{GeV}, \ |\eta(D_s^{\pm})| < 2.1$ 

"hard fragmentation" cleaning:

$$\frac{p_T(D_s^{\pm})}{E_T} > 0.02$$

### Cross section of $D^{*\pm}$ production



kinematic range:  $pT > 3.5 \,\text{GeV}, |\eta| < 2.1$  $Acc = 28.7 \pm 0.4 \%$  $\sigma = 284.5 \pm 16.3 \,\mu\text{b}$ 

signal: Gauss<sup>mod</sup>

backgr: 
$$A \cdot (x - m_{\pi^+})^B \cdot \exp^{C \cdot (x - m_{\pi^+})}$$

exponential term is used if needed (integrated, 2nd  $p_T$  bin, 3rd and 4th  $\eta$  bins)

### $D_s^{\pm}$ with combined MC sample



signals: Gauss<sup>mod</sup>, backgr: exp  $M(D^+) \equiv M(D^+)^{\text{PDG}}, \quad \sigma(D^{\pm}) \equiv \sigma(D_s^{\pm})$ 

free fit:  $\sigma(D^{\pm}) = 11.6 \pm 1.8 \,\mathrm{MeV}, \ \sigma(D_s^{\pm}) = 10.3 \pm 1.5 \,\mathrm{MeV}$ matched:  $\sigma(D^{\pm}) = 11.3 \pm 0.7 \,\mathrm{MeV}, \ \sigma(D_s^{\pm}) = 10.9 \pm 0.7 \,\mathrm{MeV}$ 

### Cross section of $D_s^{\pm}$ production



kinematic range: pT > 3.5 GeV,  $|\eta| < 2.1$  $Acc = 7.28 \pm 0.75 \%$  $\sigma = 168 \pm 34 \,\mu\text{b}$ 

signals: Gauss<sup>mod</sup>, backgr: exp  $M(D^+) \equiv M(D^+)^{\text{PDG}}, \quad \sigma(D^{\pm}) \equiv \sigma(D_s^{\pm})$ 

### Subtraction of $D_s^{\pm}$ reflection from $D^+ \rightarrow K^- \pi^+ \pi^+$

 $D_s^{\pm}$  reflection to  $D^+ \rightarrow K^- \pi^+ \pi^+$  signal :

 $\mathcal{B}_{D^+_s \to \phi \pi^+, \phi \to K^+ K^-} = 2.32 \pm 0.14\,\%$  : removed by veto on M(KK)

 $\mathcal{B}_{D^+_s \rightarrow K^+ K^- \pi^+} = 5.50 \pm 0.27\,\%$  : > 3% can distort  $D^\pm$  signal

How to remove the reflection ?

MC : by matching to true level

DATA : shape from the matched MC reflection

normalisation =  $\frac{N^{DATA}(D_s^{\pm})}{N^{MC}(D_s^{\pm})} \cdot \frac{r_B^{PDG}}{r_B^{RC}}$ where  $r_B = \frac{\mathcal{B}_{D_s^{\pm} \to K^{\pm}K^{-}\pi^{\pm}} - \mathcal{B}_{D_s^{\pm} \to \phi\pi^{\pm}, \phi \to K^{\pm}K^{-}}}{\mathcal{B}_{D_s^{\pm} \to \phi\pi^{\pm}, \phi \to K^{\pm}K^{-}}}$   $\frac{N^{DATA}(D_s^{\pm})}{N^{MC}(D_s^{\pm})} = \frac{304}{442}$   $\frac{r_B^{PDG}}{r_B^{MC}} = \frac{1.37}{1.92}$   $r_B^{MC}$  is taken from the MinBias MC (105001)

L. Gladilin

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### $D^{\pm}$ with combined MC sample

### before $D_s^{\pm}$ subtraction

after  $D_s^{\pm}$  subtraction



kinematic range:  $pT > 3.5 \,\text{GeV}, |\eta| < 2.1$  $Acc = 6.43 \pm 0.13\%$ 

signal: Gauss<sup>mod</sup>

backgr: exp



kinematic range:  $pT > 3.5 \,\text{GeV}, |\eta| < 2.1$ 

 $\sigma = 238 \pm 13 \,\mu b$ 

signal: Gauss<sup>mod</sup>

backgr: exp

<sup>&</sup>quot;D<sup>(\*)</sup> production cross sections at 7 TeV" L. Gladilin

### Binning for $D^{*\pm}$ and $D^{\pm}$ diff. x-sections

 $p_T$ : 3.5 - 5.0 - 6.5 - 8.0 - 12. - 40.

 $|\eta|$  : 0.0 - 0.2 - 0.5 - 0.8 - 1.3 - 2.1

$$\sigma_{pp \to D^{(*)}X} = \frac{N(D^{(*)})}{\mathscr{A} \cdot \mathscr{L} \cdot \mathscr{B}},$$

where  $N(D^{(*)})$  is the number of reconstructed charmed mesons,  $\mathscr{A}$  is the reconstruction acceptance obtained from the MC sample,  $\mathscr{L}$  is the integrated luminosity and  $\mathscr{B}$  is the branching fraction or the product of the branching fractions [11] for the decay channel used in the reconstruction.

The differential cross sections  $d\sigma/dp_T$  and  $d\sigma/d|\eta|$  were calculated for  $D^{*\pm}$  and  $D^{\pm}$  production in five bins in  $p_T$  (3.5-5; 5-6.5; 6.5-8; 8-12; 12-40 GeV) and five bins in  $|\eta|$  (0-0.2; 0.2-0.5; 0.5-0.8; 0.8-1.3; 1.3-2.1). To obtain the differential cross sections in a given bin, the visible

### **Control plots from CONF-note**



### Systematic uncertainties

- 1) trigger: eff = 100%,  $\delta = 0.0\%$  (only this analysis)
- 2) luminosity:  $\delta = 11\% (\rightarrow 5\% \text{ soon ?})$
- 3) branchings:  $\delta = 1.5\%$ , 4.3%, 6.0% (PDG 2010)
- 4) track reconstruction efficiency
  - 4.1) MC material description:  $\delta = 6.0 9.3\%$  (ATLAS-CONF-2010-046)
  - 4.2) tracks selection:  $\delta = 3\%$  (ATLAS-CONF-2010-046)
- 5) *D*-meson reconstruction basic cuts variations  $(\chi^2, p_T/E_T, L_{XY}, ...)$  (next slide)
- 6) *D*-meson signals extraction (fits) variations of backgr. functions (+1 par.) and ranges (-2 bins from one or another edge) for all signals and bins in addition for  $D_s^{\pm}$ :  $\sigma(D^{\pm})$  to 0.9 and 1.1 of  $\sigma(D_s^{\pm})$ in addition for  $D^{\pm}$ : subtraction of  $D_s^{\pm}$  admixture by  $\pm 1\sigma$  (~ 30%)
- 7) model dependence
  - 7.1) variations of  $p_T$  and  $\eta$  MC spectra
  - (within compatibility with data, see slide)
  - 7.2) variations of beauty contribution (by factor 2)

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"Update D^{(*)\pm} production"
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### NLO predictions to confront with data

Used : MC@NLO, POWHEG-HERWIG, POWHEG-PYTHIA (NLO+PS MC, public codes)

To be involved : FONLL (NLO+NLL), GM-VFNS (variable flavour number) Expected : MC@NLO+PYTHIA, NNLO ?

Is it the same "NLO" in all predictions ?

 $\sqrt{s} = 14 \text{ TeV}, \text{ CTEQ6m}, m_b = 4.75 \text{ GeV}, m_c = 1.5 \text{ GeV}$ 

MC@NLO 3.41:

 $\sigma_{b\bar{b}} = 0.457 \pm 0.001 \,\mathrm{mb} \quad \sigma_{c\bar{c}} = 6.539 \pm 0.020 \,\mathrm{mb} \\ \mu^2 = m_Q^2 + \frac{(p_{T,Q} + p_{T,\bar{Q}})^2}{4}$ POWHEG-hvq 1.01:

 $\sigma_{b\bar{b}} = 0.464 \pm 0.001 \, {\rm mb} \quad \sigma_{c\bar{c}} = 5.312 \pm 0.010 \, {\rm mb}$ 

Comment from Stefano Frixione:

 $\mu^2 = m_Q^2 + (M_{Q\bar{Q}}^2/4 - m_Q^2) \cdot \sin^2(\theta_Q)$ 

MC@NLO returns a total cross section identical to that computed by the underlying NLO computation. This is not the case for Powheg; extra terms, beyond NLO, are included. These are NOT, however, the result of a proper NNLO computation (or beyond). Thus, large differences between MC@NLO and Powheg reflect our ignorance on the predictions beyond NLO.

The above difference is due to differences in the scale choice Total x-sections with fixed scale ( $m_o$ ) agree within 1%

## Hadronisation and theoretical uncertainties

Hadronisation : HERWIG cluster model or Bowler modification of Lund symmetric fragmentation function

Fragmentation fractions set to LEP data :

### Theoretical uncertainties :

	LEP data			
	stat.⊕syst. br.			
$f(c \rightarrow D^{*+})$	$0.235\ \pm 0.007\ \pm 0.003$			
$f(c \rightarrow D^+)$	$0.222 \pm 0.010 \pm 0.009$			
$f(c \rightarrow D_s^+)$	$0.087\ \pm 0.009\ \pm 0.005$			
$f(b \rightarrow D^{*\pm})$	$0.175\ \pm 0.020\ \pm 0.001$			
$f(b \rightarrow D^{\pm})$	$0.227\ \pm 0.016\ \pm 0.010$			
$f(b \rightarrow D_s^{\pm})$	$0.140 \pm 0.016 \pm 0.008$			

- scale uncertainty. The uncertainty was determined by varying μ<sub>r</sub> and μ<sub>f</sub> independently to μ/2 and 2μ, with the additional constraint 1/2 < μ<sub>r</sub>/μ<sub>f</sub> < 2, and selecting the largest positive and negative variations;</li>
- *m<sub>Q</sub>* uncertainty. The uncertainty was determined by varying the charm and bottom quark masses independently by 0.2 GeV and 0.25 GeV, respectively. The total *m<sub>Q</sub>* uncertainty was obtained by adding the positive and negative cross-section variations in quadrature;
- PDF uncertainty. The uncertainty was determined by using the CTEQ6.6 PDF error eigenvectors. The total PDF uncertainty was obtained by adding the positive and negative cross-section variations in quadrature;
- hadronisation uncertainty. This uncertainty was obtained for each D<sup>(\*)</sup> meson as a sum in quadrature of the corresponding fragmentation fraction uncertainty and the fragmentation function uncertainty. The latter uncertainty was determined in frame of the POWHEG-PYTHIA predictions by using the Peterson fragmentation function [26] with extreme choices of the fragmentation parameter: 0.02 and 0.1 for charm fragmentation, and 0.002 and 0.01 for beauty fragmentation. The uncertainties of the fragmentation fractions originating from the uncertainties in the charm meson decay branching ratios were not included into the total hadronisation uncertainties because they affect experimental and theoretical cross-section calculations in the same way and can be ignored in the comparison.

## **Fragmentation Fractions**

The heavy quark hadronisation was performed using the cluster model [16] in case of HERWIG. In PYTHIA, the Lund string model [17] with the Bowler modification [18] of the Lund symmetric fragmentation function [19] for heavy quarks was used. Fragmentation fractions of heavy quarks hadronising as a particular charm meson,  $f(Q \rightarrow D^{(*)})$ , were set to experimental values. The *c*-quark fragmentation fractions, obtained by averaging of the LEP measurements as in [20], and the *b*-quark fragmentation fractions, measured by ALEPH [21] and OPAL [22], were recalculated using updated values [24] of the relevant charm-meson decay branching ratios. They are summarised in Table 1.

	LEP data				
	stat.⊕syst. br.				
$f(c \rightarrow D^{*+})$	$0.235\ \pm 0.007\ \pm 0.003$				
$f(c \rightarrow D^+)$	$0.222\ \pm 0.010\ \pm 0.009$				
$f(c \rightarrow D_s^+)$	$0.087 \pm 0.009 \pm 0.005$				
$f(b \rightarrow D^{*\pm})$	$0.175\ \pm 0.020\ \pm 0.001$				
$f(b \rightarrow D^{\pm})$	$0.227\ \pm 0.016\ \pm 0.010$				
$f(b \rightarrow D_s^{\pm})$	$0.140 \pm 0.016 \pm 0.008$				

Table 1: The fractions of *c* and *b* quarks hadronising as a particular charm meson,  $f(Q \rightarrow D^{(*)})$ , obtained using the LEP measurements (see text).

### [23] L.Gladilin, Charm Hadron Production Fractions, hep-ex/9912064

## **Charm Hadron Production Fractions**

The probabilities that a charm quark fragments into  $D^{*+}$ ,  $D_s^+$  and other charm hadrons are not calculable in perturbative QCD. To obtain them one can use data on charm production in  $e^+e^-$  annihilations. The most comprehensive charm measurements were performed by the CLEO [1, 2] and ARGUS [3, 4, 5] collaborations at centre-of-mass energies of about 10 GeV and by the OPAL [6, 7], ALEPH [8] and DELPHI [9, 10] collaborations in Only LEP measurements used for current averaging because CLEO and ARGUS results are rather old and close to threshold

Particle	OPAL	ALEPH	DELPHI	
	$\frac{\Gamma_{c\bar{c}}}{\Gamma_{b-d}} \cdot f(c \rightarrow D, \Lambda) \cdot B$	$\frac{\Gamma_{c\bar{c}}}{\Gamma_{bad}} \cdot f(c \rightarrow D, \Lambda) \cdot B$	$\frac{\Gamma_{c\bar{c}}}{\Gamma_{b-d}} \cdot f(c \rightarrow D, \Lambda) \cdot B$	
	(%)	(%)	(%)	
$D_s^+$	$0.056\ \pm 0.015\ \pm 0.007$	$0.072\ \pm 0.012\ \pm 0.004$	$0.076\ \pm 0.007\ \pm 0.007$	
$D^0$	$0.389 \pm 0.027 \substack{+0.026 \\ -0.024}$	$0.370\ \pm 0.011\ \pm 0.023$	$0.360\ \pm 0.010\ \pm 0.021$	
$D^+$	$0.358 \pm 0.046  {}^{+0.025}_{-0.031}$	$0.368 \ \pm 0.012 \ \pm 0.020$	$0.349 \ \pm 0.012 \ \pm 0.021$	
$\Lambda_c^+$	$0.041 \pm 0.019 \pm 0.007$	$0.067 \pm 0.007 \pm 0.004$	$0.074 \pm 0.015 \pm 0.009$	

Charm hadron production fractions can be obtained dividing the measured products by the branching ratios and the Standard Model value of  $\frac{\Gamma_{c\bar{c}}}{\Gamma_{had}}$  [13]:

$$\frac{\Gamma_{c\bar{c}}}{\Gamma_{had}} = 0.1719 \pm 0.0017 .$$
(7)

There are four LEP measurements which can be used for the  $f(c \rightarrow D^{*+})$  calculation. The OPAL collaboration has done a double tagged measurement of

Two OPAL's  $f(c \to D^{*+})$  measurements are statistically and systematically correlated. To calculate the correlations one can use the result for  $\frac{\Gamma_{c\bar{c}}}{\Gamma_{had}}$  obtained by the OPAL collaboration [6] from the measurements (8) and (12):

## $f(b \rightarrow D^{\pm})$ Fragmentation Fractions

$f(b \rightarrow D^{*\pm})$	0.175	$\pm 0.020$	$\pm 0.001$
$f(b  ightarrow D^{\pm})$	0.227	$\pm0.016$	$\pm 0.010$
$f(b \rightarrow D_s^{\pm})$	0.140	$\pm 0.016$	$\pm 0.008$

OPAL with new branching ratios ALEPH with new branching ratios ALEPH with new branching ratios

### PDG 2010: $B^{\pm}/B^0/B_s^0/b$ -baryon ADMIXTURE

Г <sub>38</sub>	D*(2010) <sup>+</sup> anything	( 17.3 $\pm$ 2.0 )%	branching ratio not updated
Г <sub>36</sub> Г <sub>37</sub>	D <sup>+</sup> anything D <sup>-</sup> anything	( 22.7 $\pm$ 1.8 )%	old branching ratio uncertainty
Г <sub>46</sub> Г <sub>47</sub>	$D_s^-$ anything $D_s^+$ anything	$(14.7 \pm 2.1)\%$ $(10.1 \pm 3.1)\%$	misinterpretation

## Visible cross sections

The visible cross sections for  $D^{(*)}$  mesons in the kinematic range  $p_T(D^{(*)}) > 3.5$  GeV and  $|\eta(D^{(*)})| < 2.1$  are

$$\begin{aligned} \sigma^{vis}(D^{*\pm}) &= 285 \pm 16(\text{stat.})^{+32}_{-27}(\text{syst.}) \pm 31(\text{lum.}) \pm 4(\text{br.})\,\mu\text{b}, \\ \sigma^{vis}(D^{\pm}) &= 238 \pm 13(\text{stat.})^{+35}_{-23}(\text{syst.}) \pm 26(\text{lum.}) \pm 10(\text{br.})\,\mu\text{b}, \\ \sigma^{vis}(D^{\pm}_{s}) &= 168 \pm 34(\text{stat.})^{+27}_{-24}(\text{syst.}) \pm 18(\text{lum.}) \pm 10(\text{br.})\,\mu\text{b}, \end{aligned}$$

where the last two uncertainties are due to those on the luminosity measurement and the charmed meson decay branching fractions.

The POWHEG-PYTHIA predictions are

$$\begin{aligned} \sigma(D^{*\pm}) &= 153^{+169}_{-80}(\text{scale})^{+13}_{-15}(m_Q)^{+24}_{-21}(\text{PDF})^{+20}_{-16}(\text{hadr.})\,\mu\text{b}, \\ \sigma(D^{\pm}) &= 132^{+137}_{-65}(\text{scale})^{+11}_{-10}(m_Q)^{+20}_{-18}(\text{PDF})^{+21}_{-11}(\text{hadr.})\,\mu\text{b}, \\ \sigma(D^{\pm}_s) &= 59^{+57}_{-28}(\text{scale})^{+4}_{-6}(m_Q)^{+9}_{-8}(\text{PDF})^{+7}_{-8}(\text{hadr.})\,\mu\text{b}. \end{aligned}$$

The corresponding POWHEG-HERWIG predictions are  $\sigma(D^{*\pm}) = 135 \,\mu$ b,  $\sigma(D^{\pm}) = 121 \,\mu$ b and  $\sigma(D_s^{\pm}) = 50 \,\mu$ b, while MC@NLO predicts  $\sigma(D^{*\pm}) = 155 \,\mu$ b,  $\sigma(D^{\pm}) = 138 \,\mu$ b and  $\sigma(D_s^{\pm}) = 57 \,\mu$ b.

## **Differential cross sections**



### Extrapolation, fragm. ratios and and total $c\bar{c}$ x-section

$$\sigma_{tot}(D^{(*)}) = \sigma_{pp \to c\bar{c}X \to D^{(*)}X'} = \sigma_{vis}^{DATA}(D^{(*)}) \ \frac{\sigma_{vis}^{NLO,c\bar{c}}}{\sigma_{vis}^{NLO,c\bar{c}} + \sigma_{vis}^{NLO,b\bar{b}}} \ f_{extr}^{NLO,c\bar{c}}$$

 $f_{extr}^{NLO,c\bar{c}} \sim 14 - 16 \text{ (relatively stable)}$   $\sigma_{tot}(D^{*\pm}) = 3.36 \pm 0.19 \text{ mb}$   $\sigma_{tot}(D^{\pm}) = 3.10 \pm 0.17 \text{ mb}$  $\sigma_{tot}(D_s^{\pm}) = 1.90 \pm 0.38 \text{ mb}$ 

$$\sigma^{dir}(D^{\pm}) = \sigma_{tot}(D^{\pm}) - \sigma_{tot}(D^{*\pm}) \cdot (1 - \mathcal{B}_{D^{*+} \to D^{0}\pi^{+}})$$
$$\gamma_{s} = \frac{\sigma_{tot}(D^{\pm})}{\sigma^{dir}(D^{\pm}) + \sigma_{tot}(D^{*\pm})} = 0.35 \pm 0.07$$
$$P_{v} = \frac{\sigma_{tot}(D^{*\pm})}{\sigma^{dir}(D^{\pm}) + \sigma_{tot}(D^{*\pm})} = 0.63 \pm 0.03$$

 $\sigma_{c\bar{c}} = \sigma_{tot}(D^{(*)}) / f(c \to D^{(*)}) / 2.$ 

weighted mean from  $D^{*\pm}$ ,  $D^{\pm}$  and  $D_s^{\pm}$ :

 $\sigma_{c\bar{c}}=7.13\pm0.28\,\mathrm{mb}$ 

## **Extrapolated cross sections**

The visible  $D^{(*)}$  cross section were extrapolated to the cross sections in the full kinematic phase space after subtraction of the cross sections fractions originating from beauty production. To calculate the fractions and extrapolation factors the POWHEG-PYTHIA calculations were used. The combined uncertainty of the beauty-fraction subtraction and the extrapolation to the full kinematic phase space was determined by adding in quadrature results differences originating from all sources of the theoretical uncertainty (Section 7) and differences due to using POWHEG-HERWIG and MC@NLO calculations.

The total cross sections of the  $D^{(*)}$  production in charm hadronisation are

$$\begin{split} &\sigma_{cc}^{tot}(D^{*\pm}) = 3.36 \pm 0.19(\text{stat.})^{+0.38}_{-0.32}(\text{syst.}) \pm 0.40(\text{lum.}) \pm 0.05(\text{br.})^{+1.76}_{-0.82}(\text{extr.}) \,\text{mb}, \\ &\sigma_{cc}^{tot}(D^{\pm}) = 3.10 \pm 0.17(\text{stat.})^{+0.46}_{-0.30}(\text{syst.}) \pm 0.34(\text{lum.}) \pm 0.13(\text{br.})^{+1.70}_{-0.89}(\text{extr.}) \,\text{mb}, \\ &\sigma_{cc}^{tot}(D^{\pm}_{s}) = 1.90 \pm 0.38(\text{stat.})^{+0.30}_{-0.27}(\text{syst.}) \pm 0.21(\text{lum.}) \pm 0.11(\text{br.})^{+1.23}_{-0.55}(\text{extr.}) \,\text{mb}, \end{split}$$

where the last uncertainties are the combined uncertainties due to the beauty-fraction subtraction and the extrapolation procedure.

To calculate the total cross section of charm production, the total cross sections of a given  $D^{(*)}$  meson production should be divided by the doubled value of the corresponding charm fragmentation fraction from Table 1. The weighted mean of three values calculated from  $D^{*\pm}$ ,  $D^{\pm}$  and  $D_s^{\pm}$  cross sections is

$$\sigma_{cc}^{tot} = 7.13 \pm 0.28 (\text{stat.})^{+0.89}_{-0.66} (\text{syst.}) \pm 0.78 (\text{lum.})^{+3.82}_{-1.90} (\text{extr.}) \text{ mb},$$

where uncertainties of the fragmentation fractions were included into the extrapolation uncertainties. The common for measured cross sections and fragmentation fractions branching-fraction uncertainties do not affect the calculation of the total cross section of charm production.

## **Charm fragmentation ratios**

The strangeness-suppression factor was calculated as the ratio of the  $\sigma_{cc}^{tot}(D_s^{\pm})$  to the sum of  $\sigma_{cc}^{tot}(D^{\pm\pm})$  and that part of  $\sigma_{cc}^{tot}(D^{\pm})$  which originates not from  $D^{\pm}$  decays:

$$\gamma_{s/d} = \frac{\sigma_{cc}^{tot}(D^{\pm})}{\sigma_{cc}^{tot}(D^{*\pm}) + \sigma_{cc}^{tot}(D^{\pm}) - \sigma_{cc}^{tot}(D^{*\pm}) \cdot (1 - \mathscr{B}_{D^{*+} \to D^0 \pi^+})} = \frac{\sigma_{cc}^{tot}(D^{\pm})}{\sigma_{cc}^{tot}(D^{\pm}) + \sigma_{cc}^{tot}(D^{*\pm}) \cdot \mathscr{B}_{D^{*+} \to D^0 \pi^+}},$$

where  $\mathscr{B}_{D^{*+}\to D^0\pi^+} = 0.677 \pm 0.005$  [10] is the branching ratio of the  $D^{*+} \to D^0\pi^+$  decay. Using the extrapolated cross sections, the strangeness-suppression factor is

$$\gamma_{s/d} = 0.35 \pm 0.07 (\text{stat.}) \pm 0.03 (\text{syst.}) \pm 0.03 (\text{br.})^{+0.04}_{-0.03} (\text{extr.}).$$

The fraction of *D* mesons produced in a vector state was calculated as the ratio of  $\sigma_{cc}^{tot}(D^{*\pm})$  to the sum of  $\sigma_{cc}^{tot}(D^{*\pm})$  and that part of  $\sigma_{cc}^{tot}(D^{\pm})$  which originates not from  $D^{*\pm}$  decays:

$$P_{v} = \frac{\sigma_{cc}^{tot}(D^{*\pm})}{\sigma_{cc}^{tot}(D^{*\pm}) + \sigma_{cc}^{tot}(D^{\pm}) - \sigma_{cc}^{tot}(D^{*\pm}) \cdot (1 - \mathscr{B}_{D^{*+} \to D^{0}\pi^{+}})} = \frac{\sigma_{cc}^{tot}(D^{*\pm})}{\sigma_{cc}^{tot}(D^{\pm}) + \sigma_{cc}^{tot}(D^{*\pm}) \cdot \mathscr{B}_{D^{*+} \to D^{0}\pi^{+}}}$$

Using the extrapolated cross sections, the fraction Pv is

$$P_{\rm v} = 0.63 \pm 0.03 (\text{stat.})^{+0.02}_{-0.03} (\text{syst.}) \pm 0.02 (\text{br.})^{+0.04}_{-0.03} (\text{extr.}).$$

The measured charm fragmentation ratios can be compared with those obtained in  $e^+e^-$  annihilations at LEP. The LEP fragmentation ratios were calculated using fragmentation fractions from Table 1:

$$\begin{split} \gamma_{s/d}^{\text{LEP}} &= \frac{f(c \to D_s^+)}{f(c \to D^+) + f(c \to D^{*+}) \cdot \mathscr{B}_{D^{*+} \to D^0 \pi^+}} = 0.23 \pm 0.02 (\text{stat.} \oplus \text{syst.}) \pm 0.02 (\text{br.}), \\ P_v^{\text{LEP}} &= \frac{f(c \to D^{*+})}{f(c \to D^+) + f(c \to D^{*+}) \cdot \mathscr{B}_{D^{*+} \to D^0 \pi^+}} = 0.62 \pm 0.02 (\text{stat.} \oplus \text{syst.}) \pm 0.02 (\text{br.}). \end{split}$$

## Summary & Plans

- Clean  $D^{*\pm}$ ,  $D^{\pm}$  and  $D_s^{\pm}$  signals reconstructed
- Performance ATLAS-CONF note released for Physics@LHC and reported in many conferences (by author in ICHEP 2010, IHEP-LHC 2010, LHC Physics day on HQ)
- New ATLAS-CONF note on D<sup>(\*)</sup> production cross sections is ready and under review by the ATLAS editorial board and Bphysics group

### **Plans:**

- release the new ATLAS-CONF note for winter conferences
- try to extend the measurement to high-p<sub>T</sub> range using lowest L1Calo triggers – another note and, then, combined journal paper
- use *D*<sup>(\*)</sup> reconstruction methodic for *W*+*c* measurement

## **Back-up Slides**

## The ATLAS detector

**Tile Calorimeter** 

**Muon Detectors** 

3-level trigger reducing the rate from 40 MHz to ~200 Hz

Length: ~46 m Radius: ~12 m Weight: ~7 Ktons

> Inner Detector ( $|\eta| < 2.5$ , B=2T): Si Pixels, Si strips, Transition Radiation Tracker (straws). Precise tracking and vertexing, e/n separation.  $p_t$  resolution:  $\sigma/p_t \sim 3.8 \times 10^{-4} p_t$ (GeV) $\pm 0.015$

Muon Spectrometer  $(|\eta| < 2.7)$  : aircore toroids (average 0.5T) with gas-based muon chambers. Muon trigger and measurement with momentum resolution < 10%up to  $E(\mu) \sim 1$  TeV Solenoid Magnet SCT Tracker Pixel Detector TRT Tracker Toroid Magnets

**EM calorimeter**: Pb-LAr Accordion. e/ $\gamma$  trigger, identification and measurement. E-resolution:  $\sigma/E \sim 10\%/\sqrt{E}$ 

**HAD calorimetry** ( $|\eta| < 5$ ): Fe/scintillator Tiles (central), Cu/W-LAr (fwd). Trigger and measurement of jets and missing ET. E-resolution:  $\sigma/E \sim 50\%/\sqrt{E \pm 0.03}$ 

Liquid Argon Calorimeter

# **Tracking with ATLAS Inner Detector**



#### Scatter Plot of Hits on Tracks



Pixel Detector:

3 barrel layers,  $2 \times 3$  end-cap discs

σ<sub>rφ</sub> ~ 10 µm, σ<sub>z</sub> ~ 115 µm<sup>§</sup> Silicon Strip Detector (SCT)

4 barrel layers, 2 x 9 end-cap discs

 $\sigma_{r\phi} \sim 17 \ \mu m, \sigma_z \sim 580 \ \mu m$  $\Re$  Transition Radiation Tracker(TRT)

73 barrel straw layers, 2x160 end-cap radial straw discs

 $\sigma_{r\phi} \sim 130 \ \mu m$ 

#### In this analysis:

tracks with at least 1 hit in Pixel and 4 hits in SCT and pT>250 MeV and  $|\eta|$ <2.5

# Minimum-Bias Trigger



MinBias Trigger Scintillator at  $z=\pm3.56$  m on LAr cryostat; 2 rings with 8 sector in azimuth 2.09< $|\eta|$ <2.82, 2.82< $|\eta|$ <3.84



# At least one hit above threshold in the Minimum-Bias Trigger Scintillators at each end of the detector

Efficiency is ~100% for events with at least 2 tracks passing beam-spot region

MBTS trigger allow us to measure *D*-mesons production cross-sections without uncertainty originating from trigger efficiency

The trigger is heavily prescaled with luminosity increase

## Luminosity and Data Taking



### All ID components operational > 97%

Subdetector	Number of Channels	Approximate Operational Fraction
Pixels	80 M	97.5%
SCT Silicon Strips	6.3 M	99.3%
TRT Transition Radiation Tracker	350 k	98.0%

For *D*-meson reconstruction: 1.4 nb<sup>-1</sup> (March-May, minimum-bias trigger after prescale)

### Resonance reconstruction with secondary and tertiary vertexes



Fitted masses and widths agree with MC and PDG mass values

### Predictions and expected cross sections

Monte Carlo with LO matrix elements and LL parton showering (PYTHIA)

- flavor creation (gg  $\rightarrow Q\bar{Q}$  ,  $q\bar{q} \rightarrow Q\bar{Q}$ )
- flavor excitation (gQ  $\rightarrow$  gQ, qQ  $\rightarrow$  qQ)
- gluon splitting (gg  $\rightarrow Q\bar{Q}$ )

PYTHIA with ATLAS Geant4 simulation (arXiv:1005.4568)

- detector level comparisons
- efficiencies and corrections

NLO+PS codes:

MC@NLO  $3.41 \implies$  HERWIG POWHEG-hvq  $1.01 \implies$  HERWIG, PYTHIA

and for PDFs: LHAPDF 5.8.1

MC@NLO, CTEQ6.6,  $m_b = 4.75 \text{ GeV}, m_c = 1.5 \text{ GeV}, \mu_r = \mu_f = m_T = \sqrt{m_Q^2 + p_T^2}$ 

 $\sqrt{s}$  dependence:

$\sqrt{s}[TeV]$	$\sigma_{b\bar{b}}[mb]$	$\sigma_{c\bar{c}}[mb]$
0.9	0.0225	0.891
2.36	0.757	1.95
7.0	0.243	4.40
10.	0.345	5.68
14.	0.475	7.18

Theor. Uncertainties are large

### $D^{*\pm}$ production x-section at $\sqrt{s} = 900 \,\text{GeV}$



 $\mathcal{L} = 8.6 \pm 1.8 \,\mu \mathrm{b}^{-1}, \quad \mathcal{A} = 13.0 \pm 1.3 \,\%, \quad \mathcal{B}_{D^{*+} \to (K^{-}\pi^{+})\pi^{+}} = 2.57 \pm 0.05 \,\%$ 

In the kinematic range:  $p_T(D^{*\pm}) > 1.8 \text{ GeV}, |\eta(D^{*\pm})| < 2.$ 

$$\sigma_{pp \to D^{*\pm}X} = 1.6 \pm 0.5 \,\mathrm{mb} \,\,(\mathrm{stat.})$$
  
(MC:  $c\bar{c}/b\bar{b} = 95\%/5\%$ )

## $D^{*+} \rightarrow D^0 \pi^+ \rightarrow (K^- \pi^+) \pi^+ (+c.c.)$ reconstruction



Kinematic range:  $p_T(D^{*\pm}) > 3.5 \text{ GeV}, |\eta(D^{*\pm})| < 2.1$   $p_T(D^{*\pm})/\sum E_T > 0.02 \iff \text{hard fragmentation}$   $L_{XY}(D^0) > 0 \iff c\tau(D^0) = 123 \,\mu\text{m}$   $p_T(K,\pi) > 1 \text{ GeV}, p_T(\pi_s) > 0.25 \text{ GeV}$  $\iff 1.83 < M(K\pi) < 1.90 \text{ GeV}$ 

Wrong-charge combinations:  $(K^+\pi^+)\pi_s^-$  (+ c.c.)

 $\Leftarrow 144 < M(K\pi\pi) - M(K\pi) < 147 \,\mathrm{MeV}$ 

Fitted masses and widths consistent with MC and PDG mass values

## $D^+ \rightarrow K^- \pi^+ \pi^+$ (+c.c.) reconstruction



Kinematic range:  $p_T(D^{\pm}) > 3.5 \text{ GeV}, |\eta(D^{\pm})| < 2.1$   $p_T(D^{\pm}) / \sum E_T > 0.02 \iff \text{hard fragmentation}$   $L_{XY}(D^{\pm}) > 1.3 \text{ mm} \iff c\tau(D^+) = 312 \,\mu\text{m}$   $p_T(K) > 1 \text{ GeV}$  $p_T(\pi_{1,2}) > 0.8 \text{ GeV}, p_T(\pi_{1,2}^{\max}) > 1 \text{ GeV}$ 

suppression of  $D^{*\pm}$  and  $D_s^+ \to \phi \pi^+ \to (K^- K^+) \pi^+$  (+c.c.) reflections: remove  $\Delta M_{1,2} < 150 \text{ MeV}$  and  $|M(K^{\pm}, "K^{\mp}") - M(\phi)_{\text{PDG}}| < 8 \text{ MeV}$  $\cos \theta^*(K) > -0.8$  (angle between  $\vec{p}(K)$  in  $D^{\pm}$  rest frame and  $\vec{p}(D^{\pm})$  in the lab)

Fitted mass and width consistent with MC and PDG mass value

## $D_s^+ \rightarrow \phi \pi^+ \rightarrow (K^- K^+) \pi^+$ (+c.c.) reconstruction



Kinematic range:  $p_T(D_s^{\pm}) > 3.5 \text{ GeV}, |\eta(D_s^{\pm})| < 2.1$   $p_T(D_s^{\pm}) / \sum E_T > 0.04 \iff \text{hard fragmentation}$   $L_{XY}(D_s^{\pm}) > 0.4 \text{ mm} \iff c\tau(D_s^{\pm}) = 150 \,\mu\text{m}$   $p_T(K_{1,2}) > 0.7 \text{ GeV}, p_T(\pi) > 0.8 \text{ GeV}$   $\iff |M(KK) - M(\phi)_{\text{PDG}}| < 6 \text{ MeV}$   $\cos \theta^*(\pi) < 0.4$   $(\angle \text{ between } \vec{p}(\pi) \text{ in } D_s^{\pm} \text{ r.f. and } \vec{p}(D_s^{\pm}) \text{ in the lab})$   $|\cos \theta'(K)|^3 > 0.2$  $(\angle \text{ between } \vec{p}(K) \text{ and } \vec{p}(\pi) \text{ in } K^+K^- \text{ r.f.})$ 

 $\Leftarrow= 1.83 < M(KK\pi) < 1.91 \,\mathrm{GeV}$ 

Fitted masses and widths consistent with MC and PDG mass values

### $D^{*\pm}$ acceptances in $p_T$ bins



 $D^{*\pm}$  x-sections in  $p_T$  bins



### $D^{*\pm}$ acceptances in $|\eta|$ bins



 $D^{*\pm}$  x-sections in  $|\eta|$  bins



### $D^{\pm}$ acceptances in $p_T$ bins



### $D^{\pm}$ x-sections in $p_T$ bins



### $D^{\pm}$ acceptances in $|\eta|$ bins



### $D^{\pm}$ x-sections in $|\eta|$ bins



### Differential $D^{\pm}$ and $D^{\pm}$ acceptances



### Systematic uncertainties

- 1) trigger: eff = 100%,  $\delta = 0.0\%$  (only this analysis)
- 2) luminosity:  $\delta = 11\%$  ( $\rightarrow 5\%$  soon ?)
- 3) branchings:  $\delta = 1.5\%$ , 4.3%, 6.0% (PDG 2010)
- 4) track reconstruction efficiency
  - 4.1) MC material description:  $\delta = 6.0 9.3\%$  (ATLAS-CONF-2010-046)
  - 4.2) tracks selection:  $\delta = 3\%$  (ATLAS-CONF-2010-046)
- 5) *D*-meson reconstruction basic cuts variations  $(\chi^2, p_T/E_T, L_{XY}, ...)$  (see slide)
- 6) D-meson signals extraction (fits) variations of backgr. functions (+1 par.) and ranges (-2 bins from one or another edge) for all signals and bins in addition for D<sup>±</sup><sub>s</sub>: σ(D<sup>±</sup>) to 0.9 and 1.1 of σ(D<sup>±</sup><sub>s</sub>) in addition for D<sup>±</sup><sub>t</sub>: subtraction of D<sup>±</sup><sub>s</sub> admixture by ±1σ (~ 30%)
- 7) model dependence
  - 7.1) variations of  $p_T$  and  $\eta$  MC spectra
  - (within compatibility with data, see slide)
  - 7.2) variations of beauty contribution (by factor 2)

### Track reconstruction uncertainties

### from ATLAS-CONF-2010-046 and corresponding internal note:

Table 4: Summary of all tracking efficiency systematics. The tracking efficiency systematics are given in bins of track  $p_{\rm T}$  and  $\eta$ . All numbers are in percent and relative to the corresponding tracking efficiencies. The total uncertainty is the sum in quadrature of the listed components; MC material description, track selection, and (if given) the  $p_{\rm T}$  turn-on systematic.

$p_{\Gamma}$ bin [MeV]	$ \eta  < 1.3$	$1.3 <  \eta  < 1.9$	$1.9 <  \eta  < 2.1$	$2.1 <  \eta  < 2.3$	$2.3 <  \eta  < 2.5$
$100 < p_{\rm T} \le 150$	$8 \oplus 1 \oplus 5$	$8 \oplus 1 \oplus 5$	$10 \oplus 1 \oplus 5$	$10 \oplus 1 \oplus 5$	$15 \oplus 1 \oplus 5$
$150 < p_{\rm T} \leq 200$	$4 \oplus 1$	$6 \oplus 1$	$7 \oplus 1$	$9 \oplus 1$	$13 \oplus 1$
$200 < p_{\rm T} \leq 250$	$3 \oplus 1$	$5 \oplus 1$	$6 \oplus 1$	$7 \oplus 1$	$12 \oplus 1$
$250 < p_{\rm T} \leq 300$	$2 \oplus 1$	$4 \oplus 1$	$6 \oplus 1$	$6 \oplus 1$	$11 \oplus 1$
$300 < p_{\rm T} \leq 350$	$2 \oplus 1$	$4 \oplus 1$	$5 \oplus 1$	6⊕1	$9 \oplus 1$
$350 < p_{\rm T} \leq 400$	$2 \oplus 1$	$4 \oplus 1$	$5 \oplus 1$	$5 \oplus 1$	$8 \oplus 1$
$400 < p_{\rm T} \leq 450$	$2 \oplus 1$	$3 \oplus 1$	$4 \oplus 1$	$5 \oplus 1$	$8 \oplus 1$
$450 < p_{\rm T} \leq 500$	$2 \oplus 1$	$3 \oplus 1$	$4 \oplus 1$	$4 \oplus 1$	$7 \oplus 1$
$500 < p_{\rm T}$	$2 \oplus 1$	$3 \oplus 1$	$4 \oplus 1$	$4 \oplus 1$	$7 \oplus 1$

3.2) tracks selection:  $\delta = 1\%$  for all tracks used: 3% for 3-tracks *D*-mesons

3.1) MC material description:  $\delta = 2 - 11\%$  for tracks in our range

the above table has been ported to the MC and correlated uncert. evaluated: Integrated  $D^{*\pm}$ ,  $D^{\pm}$  and  $D_s^{\pm}$  x-sections :  $^{+7.7}_{-6.6}\%$ ,  $^{+7.7}_{-6.6}\%$  and  $^{+7.6}_{-6.6}\%$   $p_T(D^{*\pm})$  :  $^{+7.9}_{-6.8}\%$ ,  $^{+7.6}_{-6.6}\%$ ,  $^{+7.4}_{-6.5}\%$ ,  $^{+7.4}_{-6.4}\%$   $|\eta(D^{*\pm})|$  :  $^{+6.4}_{-5.7}\%$ ,  $^{+6.4}_{-5.7}\%$ ,  $^{+6.8}_{-6.0}\%$ ,  $^{+10.9}_{-8.9}\%$   $p_T(D^{\pm})$  :  $^{+7.7}_{-6.7}\%$ ,  $^{+7.7}_{-6.6}\%$ ,  $^{+7.6}_{-6.6}\%$ ,  $^{+7.6}_{-6.6}\%$   $|\eta(D^{\pm})|$  :  $^{+6.4}_{-5.7}\%$ ,  $^{+6.4}_{-5.7}\%$ ,  $^{+6.8}_{-6.6}\%$ ,  $^{+10.9}_{-8.9}\%$  $p_T(D^{\pm})$  :  $^{+7.7}_{-6.7}\%$ ,  $^{+7.7}_{-6.6}\%$ ,  $^{+7.6}_{-6.6}\%$ ,  $^{+7.6}_{-6.6}\%$   $|\eta(D^{\pm})|$  :  $^{+6.4}_{-5.7}\%$ ,  $^{+6.4}_{-5.7}\%$ ,  $^{+6.8}_{-6.9}\%$ ,  $^{+10.2}_{-8.5}\%$ 

### Uncertainties of *D*-meson reconstruction

- only MC is varied  $\rightarrow$  check for mismatches between data and MC
- 1) invariant mass resolution
  - 1.1)  $M(D^0)$  window:  $\pm 5 \,\mathrm{MeV}$  on both sides
  - 1.2)  $M(\phi)$  window:  $\pm 0.5 \,\mathrm{MeV}$  on both sides
- 2)  $L_{XY}$  selections:  $\pm 50 \,\mu \mathrm{m}$
- 3)  $\chi^2$  selections:  $\pm 0.5$
- 4)  $d_0(D)$  selections:  $\pm 10\%$
- 5)  $z_0 \cdot \sin \theta(D)$  selections:  $\pm 10\%$
- 6)  $p_T/E_T$  selections: +25%, -100% of it's value

MC shows that the cut  $p_T(D^{(*)})/E_T > 0.02$  removes 0.7%, 0.2% and 0.4% of true (matched)  $D^{*\pm}$ ,  $D^{\pm}$  and  $D_s^{\pm}$ , respectively Let's re-check the cut effect for data (see slides)

### Effect of $p_T(D^{*\pm})/E_T > 0.02$ in data

#### with the cut (nominal)

#### without the cut



no significant signal reduction  $\sim 20\%$  background reduction  $\sim 12\% \sigma(N(D^{*\pm}))$  reduction

#### "D<sup>(\*)</sup> production cross sections at 7 TeV" L. Gladilin

### Effect of $p_T(D^{\pm})/E_T > 0.02$ in data

#### with the cut (nominal)

without the cut



no significant signal reduction  $\sim 5\%$  background reduction  $\sim 3\% \sigma(N(D^{\pm}))$  reduction

#### "D<sup>(\*)</sup> production cross sections at 7 TeV" L. Gladilin

### Effect of $p_T(D_s^{\pm})/E_T > 0.02$ in data

### with the cut (nominal)

without the cut



no significant signal reduction  $\sim 20\%$  background reduction  $\sim 10\% \sigma(N(D_s^{\pm}))$  reduction

#### "D<sup>(\*)</sup> production cross sections at 7 TeV" L. Gladilin

MC reweighting in  $p_T(D^{(*)})$  (systematics)



"D<sup>(\*)</sup> production cross sections at 7 TeV"

MC reweighting in  $|\eta(D^{(*)})|$  (systematics)



central fit :  $p0 = -0.2086 \pm 0.0845$ 

(could be used for central calculations)

$$p0 = -0.2921$$

applied as function of true  $|\eta|$ at both rec. and true levels

" $D^{(*)}$  production cross sections at 7 TeV"

### **Results from independent analyses**



### $M(D^0)$ and $M(D^+)$ in "MinBias || L1Calo" streams



### $\Delta M$ with nDoF=2 (Cascade) and nDoF=3 (3-tracks vertex)



nDoF=2 (actual Cascade)  $D^{*\pm}$  vertex is before  $D^0$  vertex small shift to larger values marginal improvement in resolution

nDoF=3 (3-tracks vertex)  $D^{*\pm}$  vertex is after  $D^0$  vertex sizable shift to smaller values great improvement in resolution

Should we also use 1-vertex approx. for cases when  $D^{*\pm}$  vertex is insignificantly before  $D^0$  vertex ?

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### $\Delta M$ with Cascade and 1-Vertex for low- and high- $p_T$ slow pions



1-vertex approx. works for both low- and high- $p_T$  slow pions

Statistical expectations for  $\chi_3$  (1-Vertex) and  $\chi_2$  (Cascade):  $\chi_3 < \chi_2 + 1$ : 1-Vertex is expected to be better  $\chi_3 > \chi_2 + 1$ : Cascade is expected to be better

<sup>&</sup>quot;D<sup>(\*)</sup> production cross sections at 7 TeV" L. Gladilin

### $\Delta M$ with Cascade and 1-Vertex for various $\chi^2$ changes





### 1-vertex approx. works better in all cases

"D<sup>(\*)</sup> production cross sections at 7 TeV"

L. Gladilin

### $\Delta M$ with Cascade and 1-Vertex for various $L_{XY}(PV - L)$



"D(\*) production cross sections at 7 TeV"

### Momentum scale for $D^{\pm}$ tracks in MC



### Momentum scale for $D^{\pm}$ tracks in MC



MC suggests reconstructed tracks momenta are overestimated by at least 0.1% (for systematics:  $0.1 \pm 0.1\%$ )

"D-mesons production"

e

### Description of $D^{(*)}$ signals' shapes by Gauss



### Signal fitting (matched $D^{\pm}$ in another exp. as example)



### **Description of** $D^{(*)}$ signals' shapes by Modified Gauss



Shapes of  $D^{(*)}$  signals are described much better by the Modified Gaussian function than by usual Gaussian function

"D-mesons production"

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 $\Delta_{\text{shift due to 1-vert.}} = 35 \pm 16 \, \text{keV}$ 

Possible fix:

2-vertex fit with the distance between  $D^{*\pm}$  and  $D^0$  vertices fixed to the expected (from momentum) median distance

"D-mesons production"