

Sudakov Logs at the Muon Collider

Alfredo Glioti

Institut de Physique Théorique

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Based on:

Learning from radiation at a very high energy lepton collider

Chen, Glioti, Rattazzi, Ricci, Wulzer

2202.10509



Introduction

Chen, AG, Rattazzi, Ricci, Wulzer (2022)

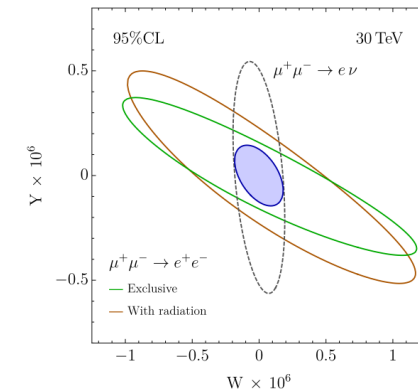
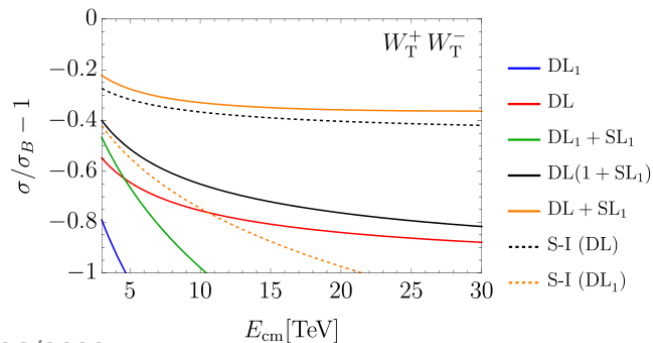
Double (Sudakov) Logs appear in cross-sections involving **soft** and **collinear EW bosons**

$$\frac{\alpha_w}{4\pi} \log^2 E^2 / m_W^2$$

This talk has **two** objectives

Show why understanding Double (and Single) Electroweak Logs is fundamental for a future Muon Collider

Show how these effects can be exploited in order to probe new physics much more effectively



EW Double Logs

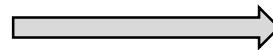
Chen, AG, Rattazzi, Ricci, Wulzer (2022)

Why are Electroweak Double Logs **special**?

- The EW group is **non-abelian** and **spontaneously broken**:
 - Double Logs do **not cancel out** even for inclusive observables
 - Initial/final states are **not EW singlets**
- EW theory is **weakly coupled**
- IR cutoff (m_W) is **physical**
- Double Logs are **extremely large** at a Muon Collider

BN Theorem violation
Ciafaloni, Ciafaloni, Comelli (2000)

$$E \gg m_w$$



$$\frac{\alpha_w}{4\pi} \log^2 E^2/m_W^2 \sim O(1)$$

$$10 \text{ TeV} \implies \sim 0.25$$

+ group theory factors...

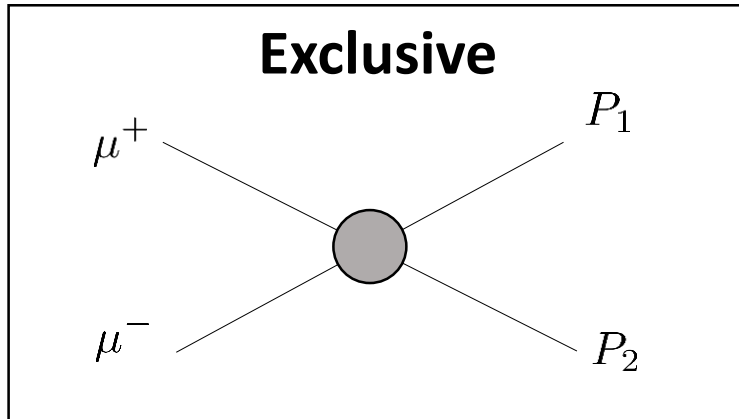
These effects must be **resummed** in order to reach % level precision!

Part 1 – Resumming Double Logs

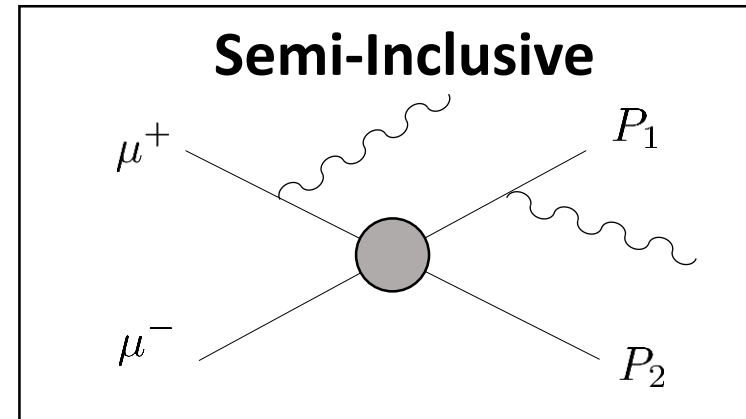
Which observables?

Chen, AG, Rattazzi, Ricci, Wulzer (2022)

We consider **two representative observables** for which we know how to perform the DL resummation



- **Two hard** final particles with **definite EW color**
- **Veto** on soft/collinear EW radiation
- Inclusive on soft photons/gluons



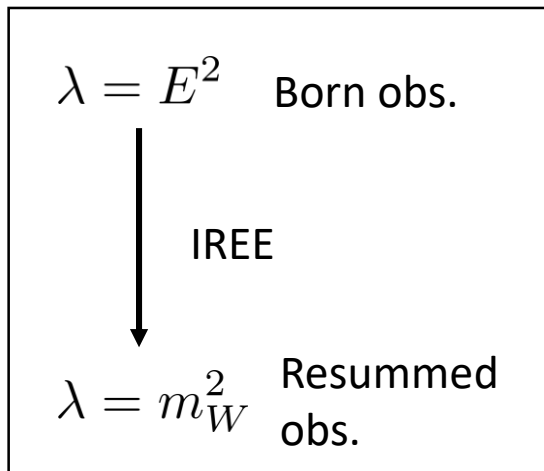
- **Two hard** final particles with **definite EW color**
- All radiation is **allowed** (up to some hardness threshold)
- “**Semi**” since we don’t **sum** over external legs colors

IREE strategy

Chen, AG, Rattazzi, Ricci, Wulzer (2022)

Our resummation strategy is based on an **InfraRed Evolution Equation (IREE)**

Fadin, Lipatov, Martin, Melles, 1999



- Introduce an unphysical **IR cutoff** λ

$$\lambda < \min \left| \frac{(k_i \cdot q)(k_j \cdot q)}{(k_i \cdot k_j)} \right|$$

- Compute derivative of the observable wrt λ through **diagrammatic** techniques

$$\frac{d}{d\lambda} \mathcal{O}^\lambda = \mathcal{K} \cdot \mathcal{O}^\lambda$$

- Solve the IREE with the **boundary condition**

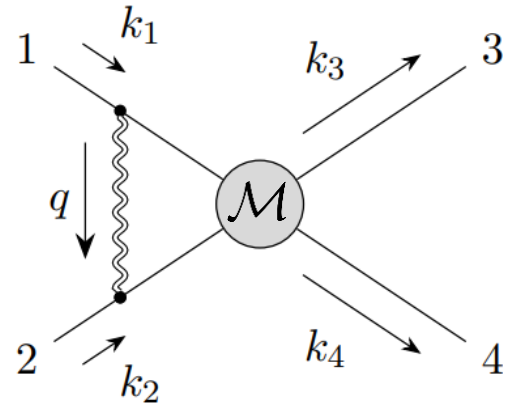
$$\mathcal{O}^{\lambda=E^2} \equiv \mathcal{O}^{\text{Born}}$$

Exclusive observables

Chen, AG, Rattazzi, Ricci, Wulzer (2022)

Exclusive on EW radiation, but inclusive on soft photons

$$\mu^+ \mu^- \rightarrow P_1 P_2 + n\gamma$$



IREE defined at the level of **Amplitudes**

Kernel is **color-diagonal**

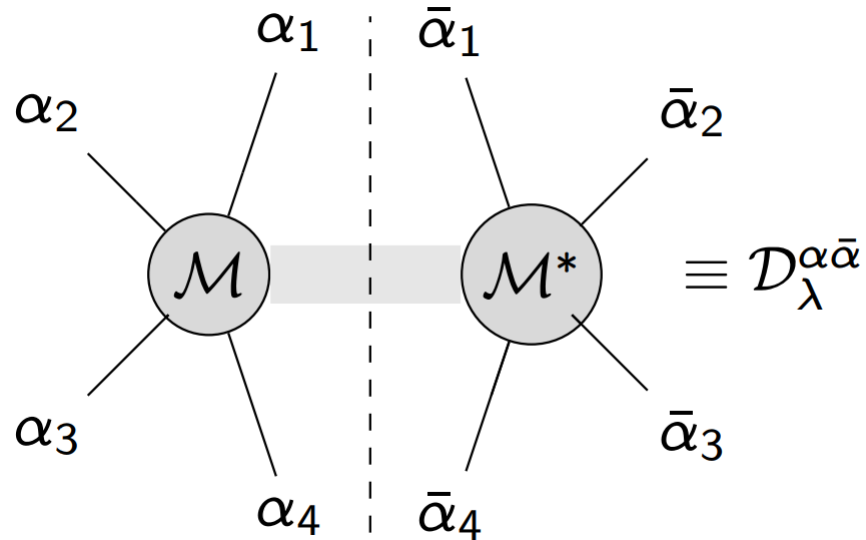
$$\frac{d\mathcal{M}_\lambda^\alpha}{d \log^2(E^2/\lambda)} = -\frac{1}{2} \mathcal{K} \mathcal{M}_\lambda^\alpha, \quad \mathcal{K} = \frac{1}{16\pi^2} \sum_i [g^2 c_i + g'^2 y_i^2]$$

Semi-inclusive observables

Chen, AG, Rattazzi, Ricci, Wulzer (2022)

$$\mu^+ \mu^- \rightarrow P_1 P_2 + n(W, Z, \gamma)$$

P1 and P2 carry ~90%
of COM energy



IREE defined at the level
of **Density Matrix**

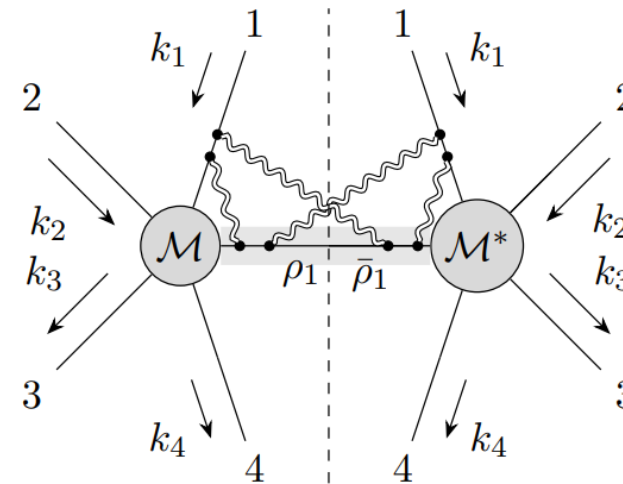
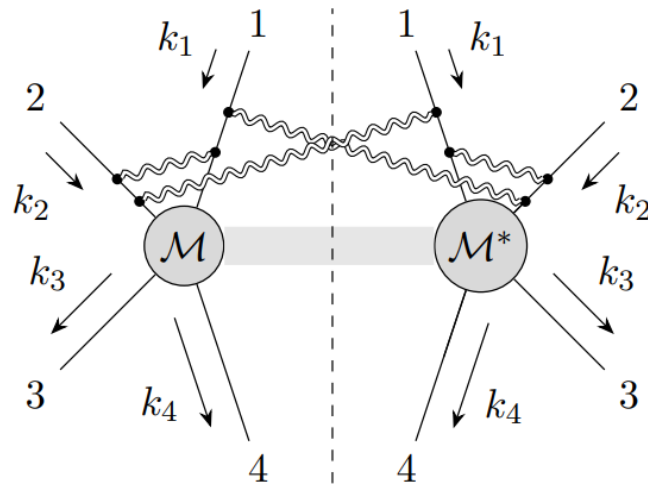
$$\mathcal{D}_\lambda^{\alpha \bar{\alpha}} \equiv \mathcal{M}_\lambda^\alpha (\mathcal{M}_\lambda^{\bar{\alpha}})^* + \sum_{N=1}^{\infty} \int d\text{Ph}_{N,\lambda}^{\mathcal{H}} \sum_{\rho_1 \dots \rho_N} \mathcal{M}_\lambda^{\alpha; \rho} (\mathcal{M}_\lambda^{\bar{\alpha}; \rho})^*$$

Semi-inclusive observables

Chen, AG, Rattazzi, Ricci, Wulzer (2022)

$$\mu^+ \mu^- \rightarrow P_1 P_2 + n(W, Z, \gamma)$$

P1 and P2 carry ~90% of COM energy



These do **not** contribute at DL level

$$\frac{d\mathcal{D}_\lambda^{\alpha\bar{\alpha}}}{d \log^2(E^2/\lambda)} = -\mathcal{K}_{\beta\bar{\beta}}^{\alpha\bar{\alpha}} \mathcal{D}_\lambda^{\beta\bar{\beta}}$$

Kernel is **off-diagonal**
All **Born hard processes** **mix** with each other

Summary of Results

Chen, AG, Rattazzi, Ricci, Wulzer (2022)

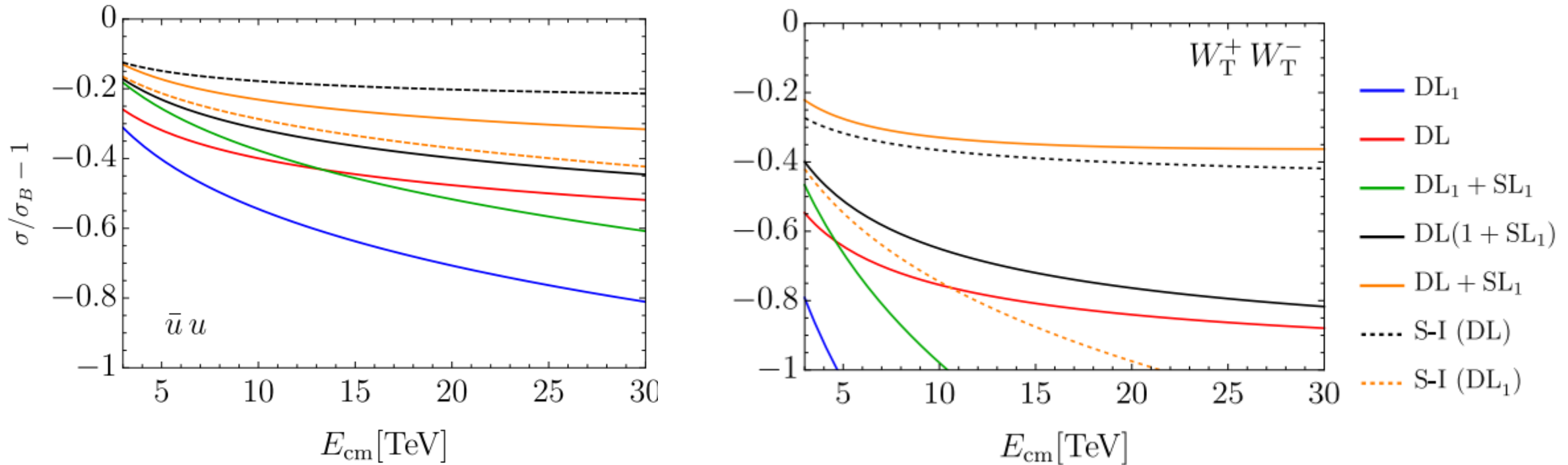
- In both cases the Kernel and double logs simply **exponentiate**
- **Exclusive** Kernel is **color-diagonal**. Gives an exponential **suppression** factor for each external leg **proportional** to its EW **Casimir**
- **Semi-inclusive** Kernel is **off-diagonal**. All allowed Born amplitude are **mixed** together
- SI-Kernel has a **leg-by-leg structure** that vanishes when summed over leg color (similar to KLN theorem, but stronger result)

$$\mathcal{K}_{\beta\bar{\beta}}^{\alpha\bar{\alpha}} = \frac{g^2}{16\pi^2} \sum_i [c_i \delta_{\beta_i}^{\alpha_i} \delta_{\bar{\beta}_i}^{\bar{\alpha}_i} + \sum_{A=1,2,3} (T_i^A)_{\beta_i}^{\alpha_i} (T_i^A)_{\bar{\beta}_i}^{\bar{\alpha}_i}] \left[\prod_{j \neq i} \delta_{\beta_j}^{\alpha_j} \delta_{\bar{\beta}_j}^{\bar{\alpha}_j} \right]$$

- Abelian **hypercharge** term **cancels out** in semi-inclusive observables as expected from BN theorem

Effect of Double and Single Logs

Chen, AG, Rattazzi, Ricci, Wulzer (2022)



Single-Logs (virtual only) added at fixed-order from Denner, Pozzorini (2000)

Effect of Double and Single Logs

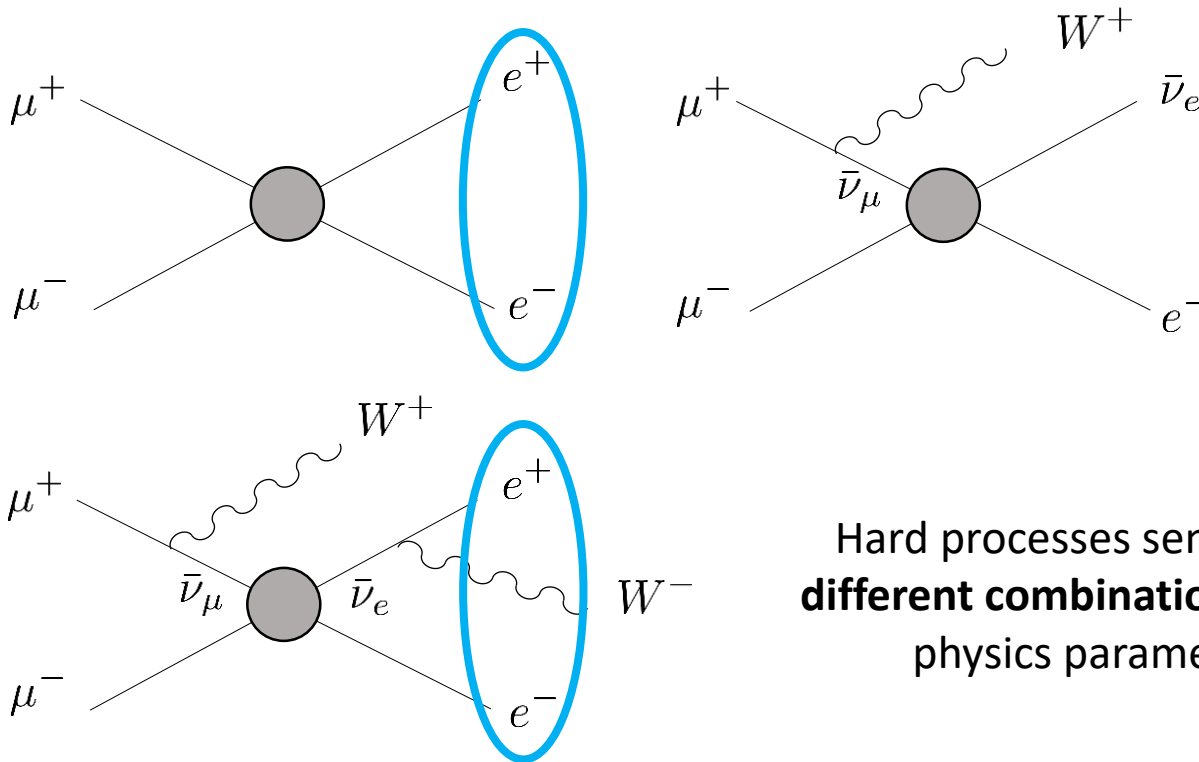
	10 TeV		
	DL	$e^{\text{DL}} - 1$	$\text{SL}(\frac{\pi}{2})$
$\ell_L \rightarrow \ell'_L$	-0.82	-0.56	0.33
$\ell_L \rightarrow q_L$	-0.78	-0.54	0.34
$\ell_L \rightarrow e_R$	-0.56	-0.43	0.17
$\ell_L \rightarrow u_R$	-0.48	-0.38	0.15
$\ell_L \rightarrow d_R$	-0.43	-0.35	0.13
$\ell_R \rightarrow \ell'_L$	-0.56	-0.43	0.17
$\ell_R \rightarrow q_L$	-0.53	-0.41	0.16
$\ell_R \rightarrow \ell'_R$	-0.30	-0.26	0.09
$\ell_R \rightarrow u_R$	-0.22	-0.20	0.07
$\ell_R \rightarrow d_R$	-0.17	-0.16	0.05

Part 2 – Probing New Physics

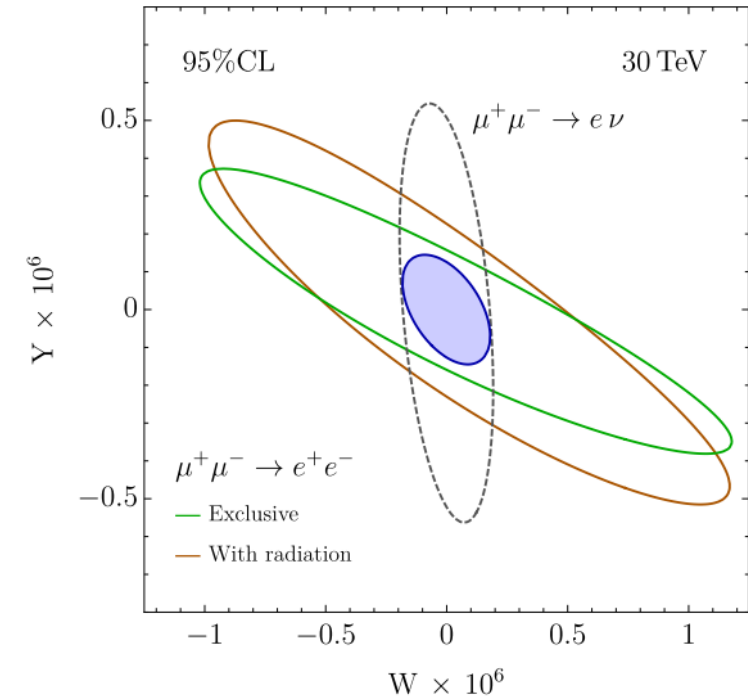
Radiation for BSM

Chen, AG, Rattazzi, Ricci, Wulzer (2022)

Thanks to the large double **logs** processes with **soft emission** become as **likely** as processes allowed at **tree-level**



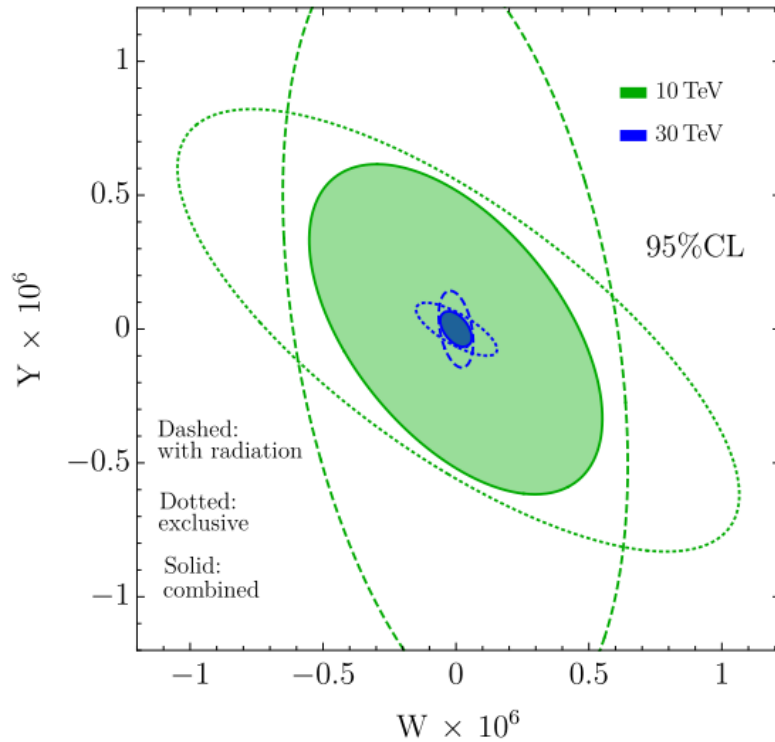
Hard processes sensitive to **different combinations** of new physics parameters



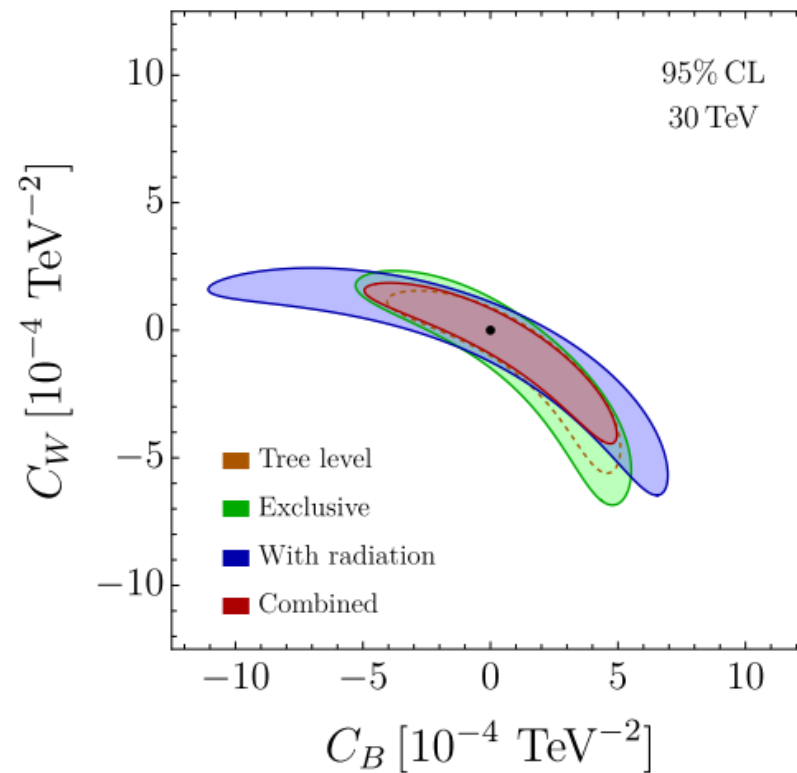
Results on effective operators

Chen, AG, Rattazzi, Ricci, Wulzer (2022)

Difermion production



Diboson production

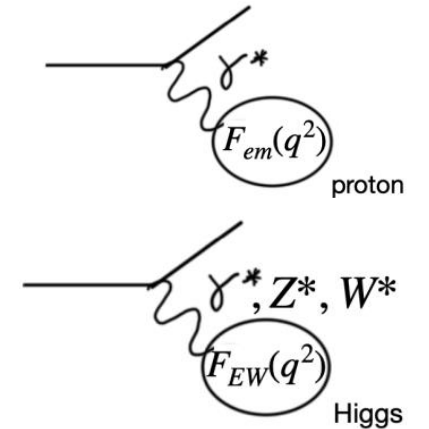


$$O_{2W} = (D_\mu W^{\mu\nu,a})^2 \sim W$$

$$O_{2B} = (\partial_\mu B^{\mu\nu})^2 \sim Y$$

$$O_W = \frac{ig}{2} (H^\dagger \sigma^a \overleftrightarrow{D}_\mu H) D^\nu W_{\mu\nu}^a$$

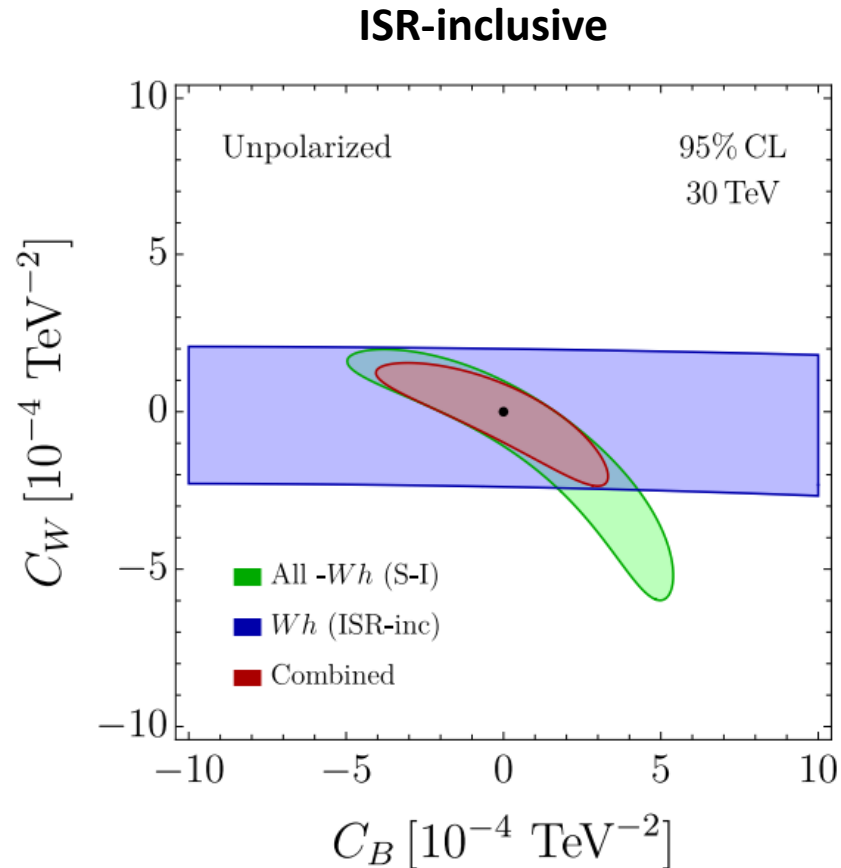
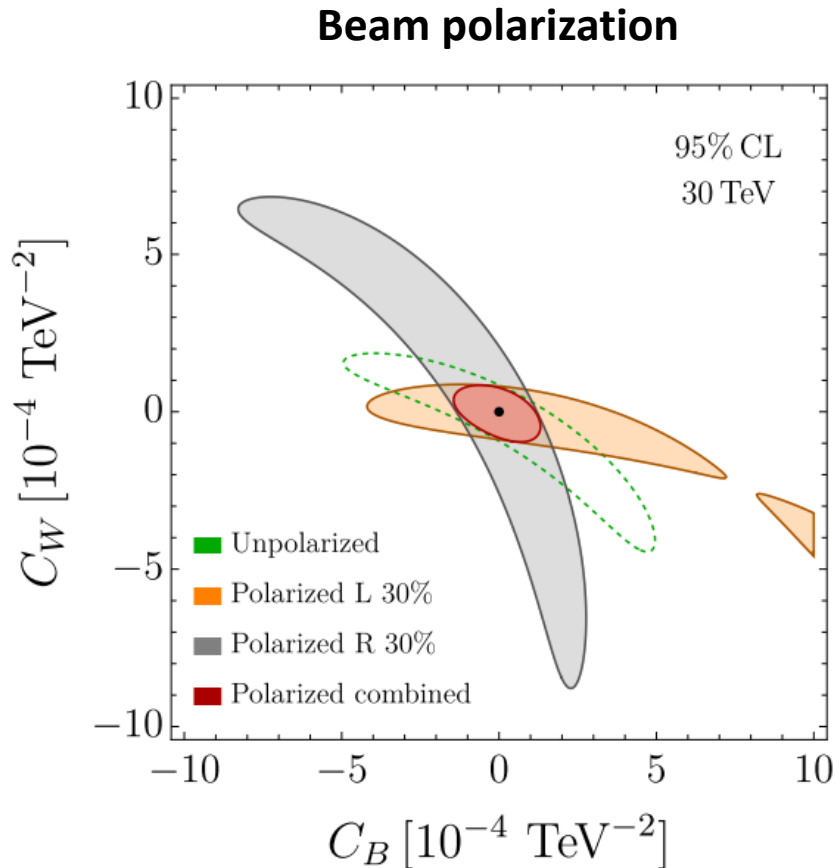
$$O_B = \frac{ig'}{2} (H^\dagger \overleftrightarrow{D}_\mu H) \partial^\nu B_{\mu\nu}$$



Avoiding flat direction

Chen, AG, Rattazzi, Ricci, Wulzer (2022)

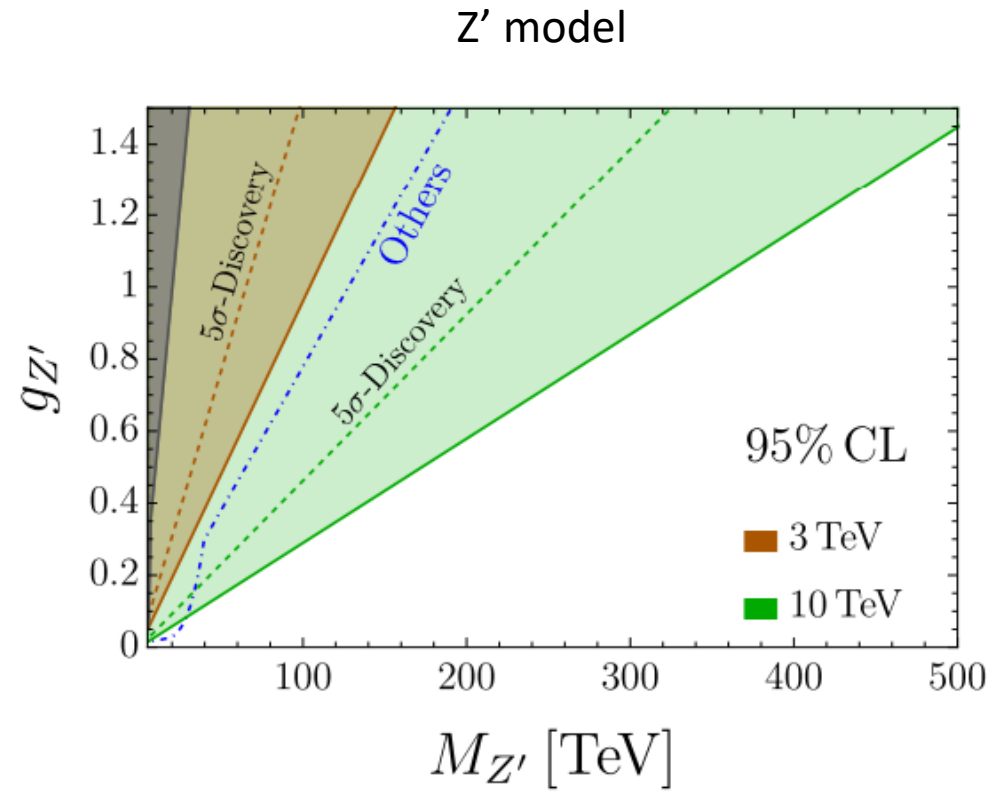
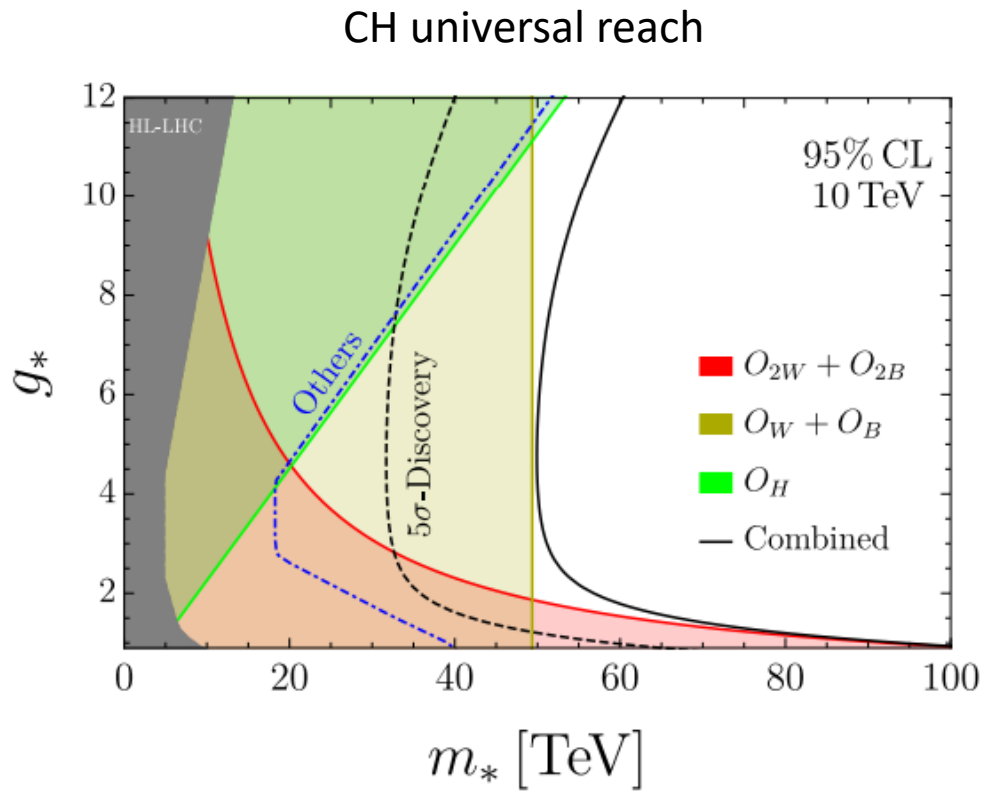
Beam polarization reduces the effect of the flat direction



Wh exclusive on the final state but ISR inclusive effectively left-polarizes the beam

Reach on BSM models

Chen, AG, Rattazzi, Ricci, Wulzer (2022)

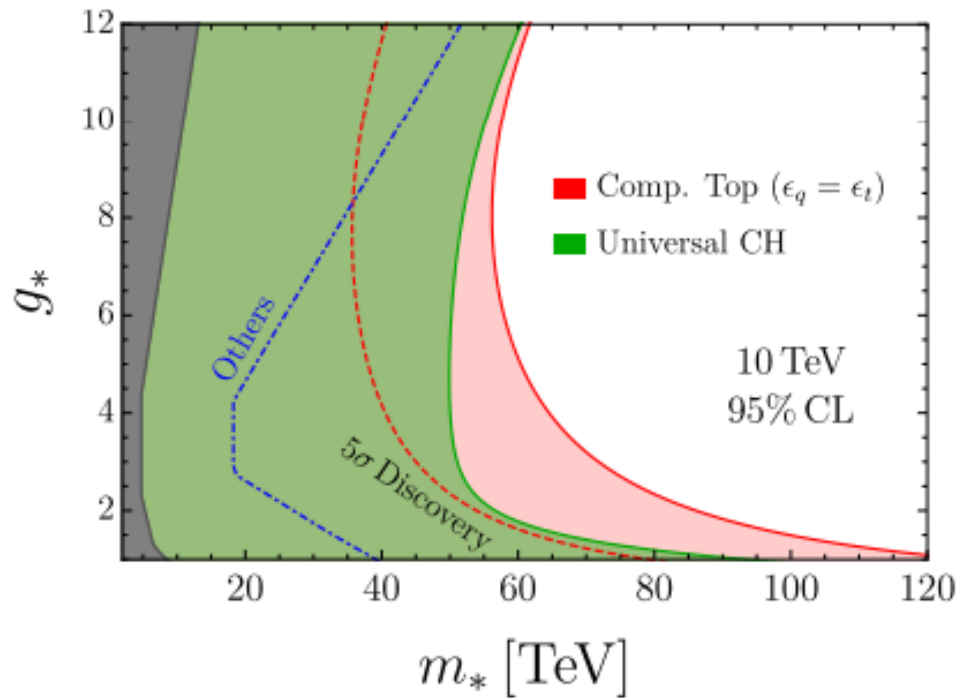


Tree-level prediction for OH from Buttazzo, Franceschini, Wulzer, 2020

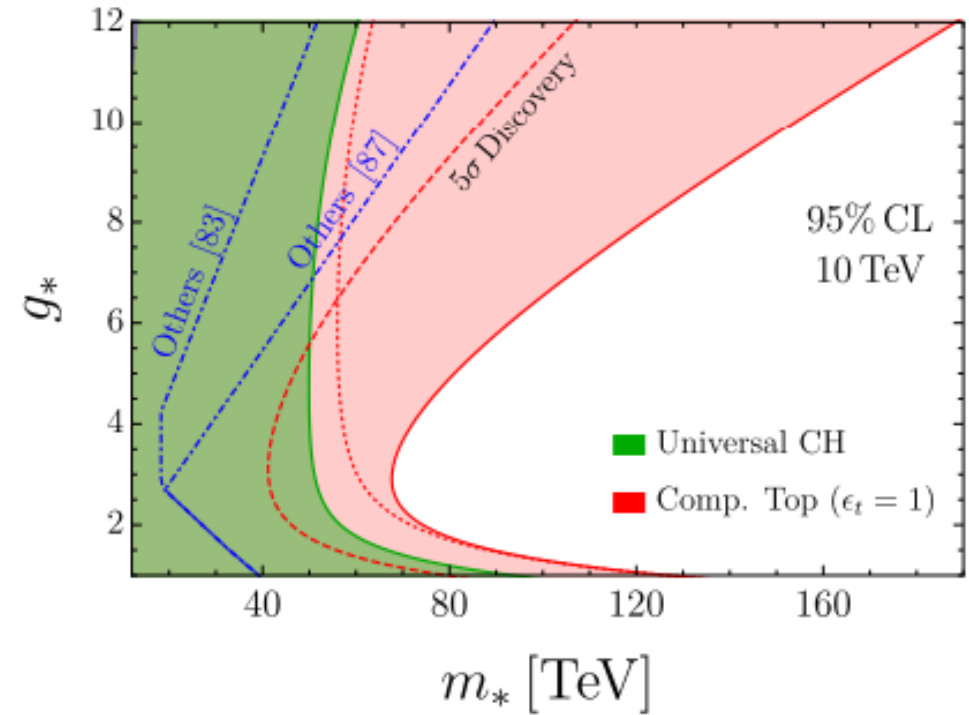
Composite top

Chen, AG, Rattazzi, Ricci, Wulzer (2022)

Equally composite top



Right composite top



Conclusions

Chen, AG, Rattazzi, Ricci, Wulzer (2022)

- The effects of **EW radiation** are extremely **large** and must be **resumed** to reach % level of precision at a Muon Collider
- A **better understanding of EW radiation is needed**: how to compute more general observables? (eg. ISR inclusive) how to include single-logs?
- Large radiation effect allows to **probe New Physics** in a **richer way** compared to tree-level prediction
- A **10 TeV Muon Collider** can indirectly probe New Physics up to **100+ TeV**
- **WIP**: a user-friendly implementation of these computations for these processes + production of heavy EW multiplets

BACKUP

	3 TeV			10 TeV		
	DL	$e^{\text{DL}}-1$	$\text{SL}(\frac{\pi}{2})$	DL	$e^{\text{DL}}-1$	$\text{SL}(\frac{\pi}{2})$
$\ell_L \rightarrow \ell'_L$	-0.46	-0.37	0.25	-0.82	-0.56	0.33
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$\ell_L \rightarrow e_R$	-0.32	-0.27	0.13	-0.56	-0.43	0.17
$\ell_L \rightarrow u_R$	-0.27	-0.24	0.11	-0.48	-0.38	0.15
$\ell_L \rightarrow d_R$	-0.24	-0.21	0.10	-0.43	-0.35	0.13
$\ell_R \rightarrow \ell'_L$	-0.32	-0.27	0.13	-0.56	-0.43	0.17
$\ell_R \rightarrow q_L$	-0.30	-0.26	0.12	-0.53	-0.41	0.16
$\ell_R \rightarrow \ell'_R$	-0.17	-0.16	0.07	-0.30	-0.26	0.09
$\ell_R \rightarrow u_R$	-0.12	-0.12	0.05	-0.22	-0.20	0.07
$\ell_R \rightarrow d_R$	-0.09	-0.09	0.04	-0.17	-0.16	0.05

