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MADGRAPH SPRING 2011 @ FERMILAB

## CONTENTS

- Motivations
- Basics of NLO computations
- Automation within MadGraph v4
- \* Results
- Work within MadGraph v5 and future plans

# MOTIVATIONS

#### NLO is important because

- \* NLO corrections are large in QCD
- NLO corrections significantly affect the shape of distributions
- \* It reduces the scale dependence inherent to tree-level cross-sections

Les Houches '06 Wish list

Wish list for one-loop high multiplicity processes

Process	background to
1. pp → VV jet	ttH, new physics
2. pp → H+ 2 jet	H production by boson fusion (VBF)
3. $pp \rightarrow t\bar{t}b\bar{b}$	tīH
4. $pp \rightarrow t\bar{t} + 2 \text{ jets}$	tīH
5. pp → VVbb	$VBF \rightarrow H \rightarrow VV$ , $t\bar{t}H$ , new physics
6. pp → VV+ 2 jets	$VBF \rightarrow H \rightarrow VV$
7. $pp \rightarrow V + 3$ jets	various new physics signatures
8. pp → VVV	SUSY trilepton

New production channels open at NLO

### Automation would help!

Accurate theoretical prediction are necessary for the search of signals events in large background samples.

## **AUTOMATION ADVANTAGES**

### Save time

Trade time spent on computing a process with time on studying the physics behind it.

## \* Avoid bugs

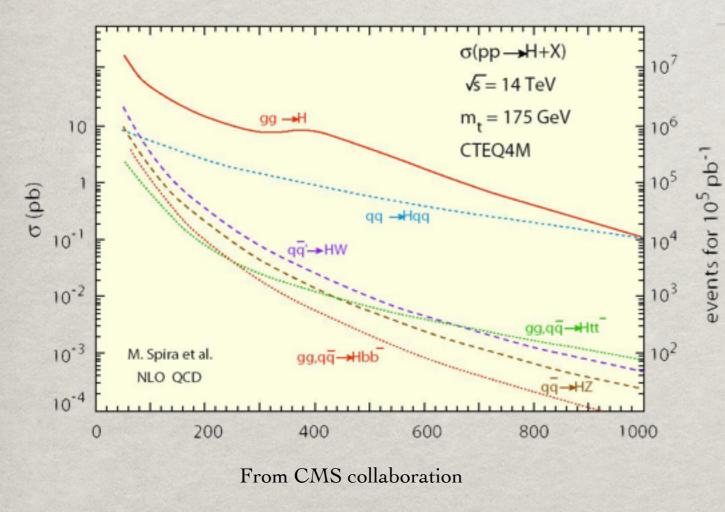
Having a trusted program extensively checked once and for all, eliminates obvious bugs when running different processes.

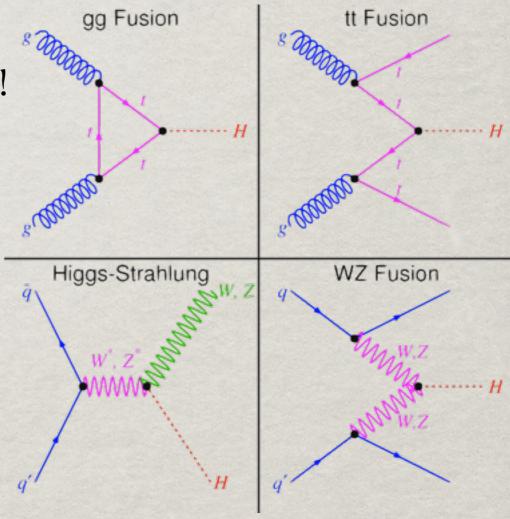
\* Use of the same framework for all processes

It only requires to know how to efficiently use one single program to do all NLO phenomenology.

## HIGGS PRODUCTION

- \* Gluon fusion exclusively loop induced!
- Still very relevant compared to other production channels.

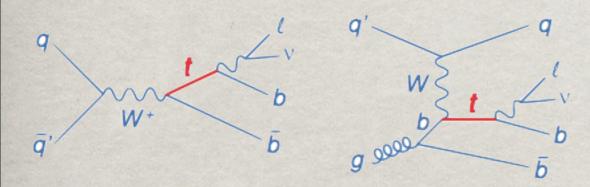


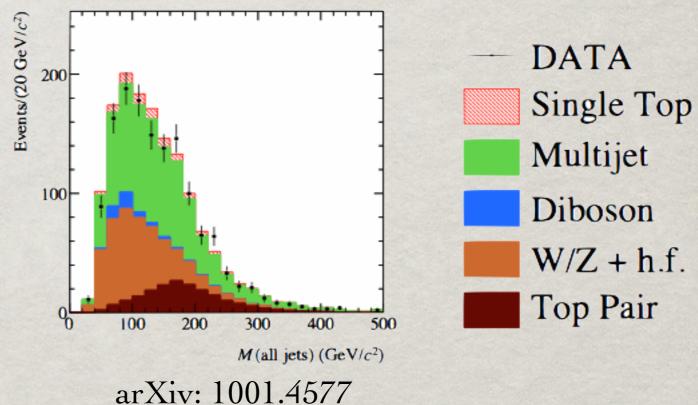


- But consider  $H \to b\bar{b}$  and compare to  $\sigma_{LHC}(b\bar{b}) = \mathcal{O}(10^8 \ [pb])$  !
- So accurate estimation of background process is crucial.

## SINGLE-TOP SEARCHES

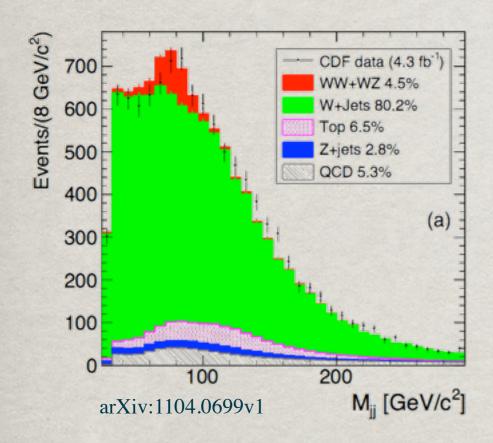
This process closely resemble to many background events which have to be carefully evaluated





Single top events have been observed at Tevatron, giving the first and only direct estimation of  $V_{tb}$ !

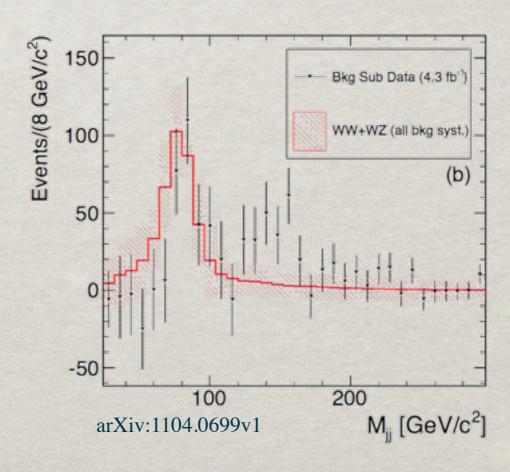
## Wjj at Tevatron



Having NLO computations by default lead to more conclusive observations.

CDF observes 3-0 deviation to the SM signal.

- New Physics, stat. fluctuations?
- Unreliable prediction?
  - → W+jets treated at LO!
  - → Mistreatment of background?



Valentin Hirschi, 4 May 2011

# NLO BASICS

## NLO contributions have two parts

$$\sigma^{\text{NLO}} = \int_{m} d^{(d)} \sigma^{V} + \int_{m+1} d^{(d)} \sigma^{R} + \int_{m} d^{(4)} \sigma^{B}$$

#### Virtual part

- \* Current bottleneck of NLO computations
- \* Algorithms for automation known in principle but not yet efficiently implemented
- \* This work brings automation using MadGraph and CutTools interfaced through MadLoop.

#### Real emission part

- \* Automated for different methods
- \* Challenge is the systematic extraction of singularities
- \* MadFKS using the FKS subtraction method successfully implemented on MGv4

## SUBTRACTION TERMS

IR divergences are dealt with using subtraction terms

$$\sigma^{\text{NLO}} = \int_{m} d^{(d)} \sigma^{V} + \int_{m+1} d^{(d)} \sigma^{R} + \int_{m} d^{(4)} \sigma^{B}$$



$$\sigma^{\text{NLO}} = \int_{m} \left[ d^{(4)} \sigma^{B} + \int_{l} d^{(d)} \sigma^{V} + \int_{1} d^{(d)} \sigma^{A} \right] + \int_{m+1} \left[ d^{(4)} \sigma^{R} - d^{(4)} \sigma^{A} \right]$$

- \* Each of the two integrals are separately finite.
- The only missing input required from MadFKS is the *finite* part of the virtual amplitude.
- \* This is the part MadLoop provides!

# TOWARDS AUTOMATION

- \* MadGraph is a tree-level matrix element generator
- \* MadFKS is built on MG and computes everything but the finite part of the amplitude.
- \* CutTools is a code which computes very general loop integrals using the OPP method
  - → Idea: To create a third-party program, MadLoop which interfaces these three tools to automatically compute some loop x-sections.

## MADFKS

#### PHASE-SPACE: DIVIDE AND CONQUER

- Real emission part:  $d\sigma^R = |M^{n+1}|^2 d\phi_{n+1}$
- $|M^{n+1}|^2$  diverges as  $\frac{1}{\chi_i^2} \frac{1}{1 y_{ij}}$  with  $\begin{cases} \chi_i = \frac{D_i}{\sqrt{\hat{s}}} \\ y_{ij} = \cos \theta_{ij} \end{cases}$
- Divide phase-space so that each partition has at most one soft and one collinear singularity

$$d\sigma^{R} = \sum_{ij} S_{ij} |M^{n+1}|^{2} d\phi_{n+1} \qquad \sum_{ij} S_{ij} = 1$$

 $d\sigma^{R} = \sum_{ij} S_{ij} |M^{n+1}|^{2} d\phi_{n+1} \qquad \sum_{ij} S_{ij} = 1$   $\text{Use plus distribution to regulate the singularities} \int d\chi \left(\frac{1}{\chi}\right)_{+} f(\chi) = \int d\chi \frac{f(\chi) - f(0)}{\chi}$ 

$$d\tilde{\sigma}^{R} = \sum_{ij} \left(\frac{1}{\chi_{i}}\right)_{+} \left(\frac{1}{1 - y_{ij}}\right)_{+} \chi_{i}^{2} (1 - y_{ij}) S_{ij} |M^{n+1}|^{2} d\phi_{n+1}$$

# CUTTOOLS

OR HOW TO COMPUTE LOOPS WITHOUT DOING SO

\* CutTools uses the OPP method for loop reduction at the integrand level

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$$\bar{q}^2 = q^2 + \tilde{q}^2 \qquad (q \cdot \tilde{q}) = 0 \qquad N$$

$$\bar{D}_i = (\bar{q} + p_i)^2 - m_i^2, \quad p_0 \neq 0.$$

$$\int d^{(d)} \sigma^V = \int d^{(4+\epsilon)} \left( A(\bar{q}) + \tilde{A}(\bar{q}) \right)$$

$$A(\bar{q}) = \frac{N(q)}{\bar{D}_0 \bar{D}_1 \cdots \bar{D}_{m-1}} \quad (\tilde{A}(\bar{q}) \to \mathbb{R}2)$$

- R2 can be obtained with a tree-level-like computation with special Feynman-Rules.
- Evaluation of N(q) for different specific q's allows to algebraically obtain the coefficients a, b, c and d
- Reconstruction of the  $\tilde{q}$  dependance of the numerator gives the cut-constructible part R1 of the finite part of the virtual amplitude

$$\bar{q}^2 = q^2 + \tilde{q}^2 \qquad (q \cdot \tilde{q}) = 0 \qquad N(q) \qquad = \sum_{i_0 < i_1 < i_2 < i_3}^{m-1} \left[ d(i_0 i_1 i_2 i_3) + \tilde{d}(q; i_0 i_1 i_2 i_3) \right] \prod_{i \neq i_0, i_1, i_2, i_3}^{m-1} \\ \bar{D}_i = (\bar{q} + p_i)^2 - m_i^2, \quad p_0 \neq 0. \qquad + \sum_{i_0 < i_1 < i_2}^{m-1} \left[ c(i_0 i_1 i_2) + \tilde{c}(q; i_0 i_1 i_2) \right] \prod_{i \neq i_0, i_1, i_2}^{m-1} D_i \\ \bar{d}^{(d)} \sigma^V = \int d^{(4+\epsilon)} \left( A(\bar{q}) + \tilde{A}(\bar{q}) \right) \\ \bar{q}^{(i)} = \frac{N(q)}{\bar{D}_0 \bar{D}_1 \cdots \bar{D}_{m-1}} \left( \tilde{A}(\bar{q}) \rightarrow \mathbb{R}^2 \right) \\ = \sum_{i_0 < i_1 < i_2 < i_3}^{m-1} \left[ b(i_0 i_1) + \tilde{b}(q; i_0 i_1) \right] \prod_{i \neq i_0, i_1}^{m-1} D_i \\ = \sum_{i_0 < i_1 < i_2 < i_3}^{m-1} \left[ b(i_0 i_1) + \tilde{b}(q; i_0 i_1) \right] \prod_{i \neq i_0, i_1}^{m-1} D_i \\ = \sum_{i_0 < i_1 < i_2 < i_3}^{m-1} \left[ b(i_0 i_1) + \tilde{b}(q; i_0 i_1) \right] \prod_{i \neq i_0, i_1}^{m-1} D_i \\ = \sum_{i_0 < i_1 < i_2 < i_3}^{m-1} \left[ b(i_0 i_1) + \tilde{b}(q; i_0 i_1) \right] \prod_{i \neq i_0, i_1, i_2}^{m-1} D_i \\ = \sum_{i_0 < i_1 < i_2 < i_3}^{m-1} \left[ b(i_0 i_1) + \tilde{b}(q; i_0 i_1) \right] \prod_{i \neq i_0, i_1, i_2}^{m-1} D_i \\ = \sum_{i_0 < i_1 < i_2 < i_3}^{m-1} \left[ b(i_0 i_1) + \tilde{b}(q; i_0 i_1) \right] \prod_{i \neq i_0, i_1, i_2}^{m-1} D_i \\ = \sum_{i_0 < i_1 < i_2 < i_3}^{m-1} \left[ b(i_0 i_1) + \tilde{b}(q; i_0 i_1) \right] \prod_{i \neq i_0, i_1, i_2}^{m-1} D_i \\ = \sum_{i_0 < i_1 < i_2 < i_3}^{m-1} \left[ b(i_0 i_1) + \tilde{b}(q; i_0 i_1) \right] \prod_{i \neq i_0, i_1, i_2}^{m-1} D_i \\ = \sum_{i_0 < i_1 < i_2 < i_3}^{m-1} \left[ b(i_0 i_1) + \tilde{b}(q; i_0 i_1) \right] \prod_{i \neq i_0, i_1, i_2}^{m-1} D_i \\ = \sum_{i_0 < i_1 < i_2 < i_3}^{m-1} \left[ b(i_0 i_1) + \tilde{b}(q; i_0 i_1) \right] \prod_{i \neq i_0, i_1, i_2}^{m-1} D_i \\ = \sum_{i_0 < i_1 < i_2, i_2, i_3}^{m-1} \left[ b(i_0 i_1) + \tilde{b}(q; i_0 i_1) \right] \prod_{i \neq i_0, i_1, i_2}^{m-1} D_i$$

Finite part = R1 + R2

#### HELPING MADGRAPH AND CUTTOOLS TO BE FRIENDS

\* Tree-like operations: Generate and obtain Born and R2 amplitudes

Loop diagrams: Generate cut-loop diagrams

Select a non-redundant basis

Compute the color factors

Provide numerator of the loop integrant

Handle exceptional PS points

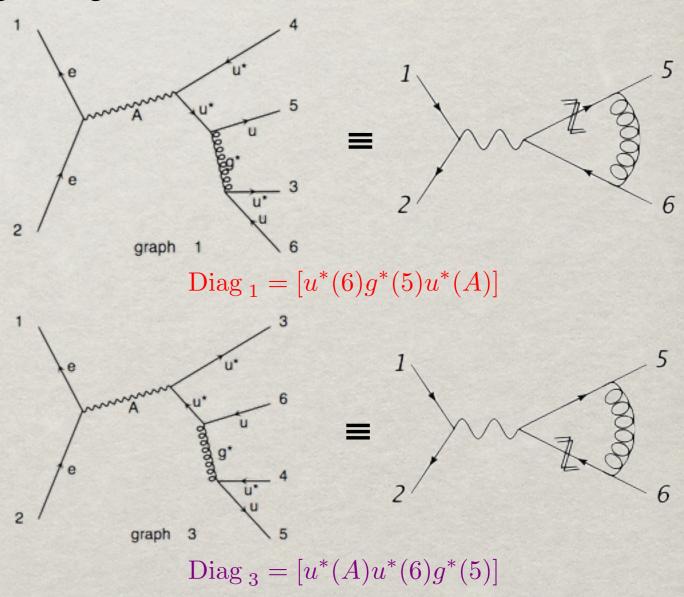
- Compute R2 contribution
- Square born against virtual amplitude
- Carry UV renormalization.
- Perform Sanity checks (Double pole, Ward Identity, ...)

# CUT-LOOP DIAGRAMS

#### WITH A SPECIFIC EXAMPLE

## Consider $e^+e^- \to \gamma \to u\bar{u}$ :

- Loop particles are denoted with a star. When MG is asked for  $e^+e^- \to u^*\bar{u}^*u\bar{u}$  it gives back eight diagrams. Two of them are:
- Selection is performed to keep only one cut-diagram per loop contributing in the process
- \* Tags are associated to each cut-diagram. Those whose tags are mirror and/or cyclic permutations of tags of diagram already in the loop-basis are taken out.
- Additional custom filter to eliminate tadpoles and bubbles attached to external legs.



#### NEW FEATURES IMPLEMENTED

- Recognition of the loop topologies in order to filter L-cut diagrams.
- Structure to deal with two MG process simultaneously (the L-cut and the born-like).
- Treat color to obtain the squaring of the loop color structure against the born one.
- \* Right form of integrant to CutTools: no denominators, complex momenta and reconstruction of missing propagator for the sewed particles.

$$\sum_{i=-1,1} \epsilon_i^{\mu}(p) \epsilon_i^{\nu*}(p) \to -g^{\mu\nu}$$

- \* Implementation of QCD ghosts.
- \* Implementation of the special R2 vertices and automatic UV renormalization.

#### FIGHTING EXCEPTIONAL PHASE SPACE POINTS

- \* CutTools can asses the numerical stability of the computation of a loop by
  - $\Rightarrow$  By sending  $m_i^2 \rightarrow m_i^2 + M^2$ , CT has an independent reconstruction of the numerator and can check if both match.
  - → CT ask MadLoop to evaluate the integrand at a given loop momentum and check if the result is close enough to the one from the reconstructed integrand.
- When an EPS occurs, MadLoop tries to cure it:
  - → Check if Ward Identities hold at a satisfactory level
  - ⇒ Shift the PS point by rescaling momenta:  $k_i^3 = (1 + \lambda_{\pm})k_i^3$
  - → Provide an estimate of the virtual for the original PS point with uncertainty:

$$v_{\lambda_{\pm}}^{FIN} = \frac{V_{\lambda_{\pm}}^{FIN}}{|\mathcal{A}_{\lambda=0}^{born}|^2} \quad c = \frac{1}{2} \left( v_{\lambda_{+}}^{FIN} + v_{\lambda_{-}}^{FIN} \right) \quad \Delta = \left| v_{\lambda_{+}}^{FIN} - v_{\lambda_{-}}^{FIN} \right| \quad V_{\lambda=0}^{FIN} = \left| \mathcal{A}_{\lambda=0}^{born} \right|^2 (c \pm \Delta)$$

→ If nothing works, then use the median of the results of the last 100 stable points

#### WHAT IT CANNOT DO YET

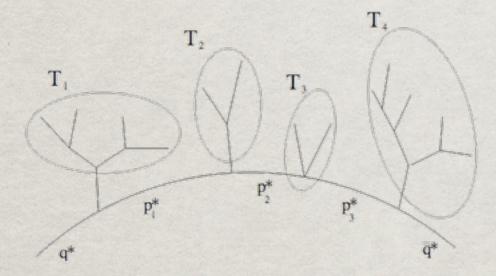
No four-gluon vertex at born level:

$$\begin{split} \mu_{1}, a_{1} & \bigoplus_{\mu_{2}, a_{2}} \\ \mu_{3}, a_{3} & = -\frac{ig^{4}N_{col}}{96\pi^{2}} \sum_{P(234)} \left\{ \left[ \frac{\delta_{a_{1}a_{2}}\delta_{a_{3}a_{4}} + \delta_{a_{1}a_{3}}\delta_{a_{4}a_{2}} + \delta_{a_{1}a_{4}}\delta_{a_{2}a_{3}}}{N_{col}} \right. \\ & + 4Tr(t^{a_{1}}t^{a_{3}}t^{a_{2}}t^{a_{4}} + t^{a_{1}}t^{a_{4}}t^{a_{2}}t^{a_{3}}) \left( 3 + \lambda_{HV} \right) \\ & - Tr(\{t^{a_{1}}t^{a_{2}}\}\{t^{a_{3}}t^{a_{4}}\}) \left( 5 + 2\lambda_{HV} \right) \right] g_{\mu_{1}\mu_{2}}g_{\mu_{3}\mu_{4}} \\ & + 12\frac{N_{f}}{N_{col}}Tr(t^{a_{1}}t^{a_{2}}t^{a_{3}}t^{a_{4}}) \left( \frac{5}{3}g_{\mu_{1}\mu_{3}}g_{\mu_{2}\mu_{4}} - g_{\mu_{1}\mu_{2}}g_{\mu_{3}\mu_{4}} - g_{\mu_{2}\mu_{3}}g_{\mu_{1}\mu_{4}} \right) \right\} \end{split}$$

- \* If EW bosons in the loop, it might be that CutTools cannot handle certain loops.
- \* All born contribution must factorize the same power of all coupling orders.
- No finite-width effects of unstable massive particles also appearing in the loop.

# MADLOOP HOW TO MAKE IT FASTER

Recycling of the tree-structures attached to the loop.



- \* Identify identical contributions (i.e. massless fermion loops of diff. flavors)
- \* Call CT not per diagram, but per set of diagrams with the same loop kinematics.
- \* Use of recursion relations: Big gain in reals, not so much in the virtual.

# LOCAL CHECKS

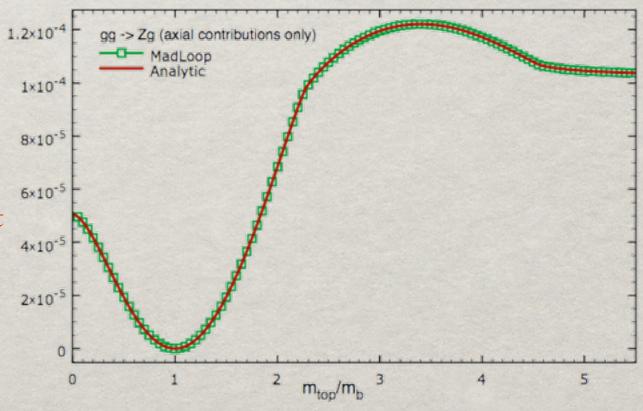
#### YOU DON'T WANT THE EXHAUSTIVE LIST ...

$u\bar{u} \to W^+W^-b\bar{b}$	MADLOOP	Ref. [33]		
$a_0$	2.338047209268890E-008	2.338047130649064E-008		
C=2	-2.493920703542680E-007	-2.493916939359002E-007		
$c_{-1}$	-4.885901939046758E-007	-4.885901774740355E-007		
c <sub>0</sub>	-2.775800623041098E-007	-2.775787767591390E-007		
$gg \to W^+W^-b\bar{b}$				
$a_0$	1.549795815702494E-008	1.549794572435312E-008		
C=2	-2.686312747217639E-007	-2.686310592221201E-007		
$c_{-1}$	-6.078687041491385E-007	-6.078682316434646E-007		
CO	-5.519004042667462E-007	-5.519004727276688E-007		

Ref. [33]: A. van Hameren et al.

We believe the code is very robust - e.g.,
MadLoop helped spot mistakes in published loop computations (Zjj, W+W+jj)

- The numerics are pin-point on analytical data, even with several mass scales.
- Analytic computations from an independent implementation of the helicity amplitudes by J.J van der Bij *et al*.



MadGraph Spring 2011 @ FermiLab

## INTEGRATED RESULTS

- Running time: Two weeks on a 150+ node cluster
- \* Proof of efficient EPS handling with  $Zt\bar{t}$
- Successful cross-check against known results
- Large K-factors sometimes
- No cuts on b, robust numerics with small  $P_T$

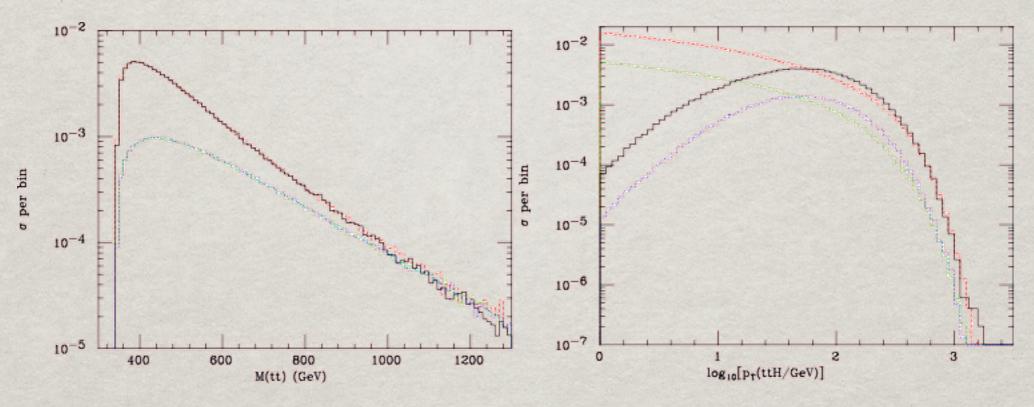
	Process	μ	$n_{if}$	Cross section (pb)	
				LO	NLO
a.1	$pp \rightarrow t\bar{t}$	$m_{top}$	5	$123.76 \pm 0.05$	$162.08 \pm 0.12$
a.2	$pp \rightarrow tj$	$m_{lop}$	5	$34.78\pm0.03$	$41.03\pm0.07$
a.3	$pp \rightarrow tjj$	$m_{top}$	5	$11.851 \pm 0.006$	$13.71\pm0.02$
a.4	$pp \rightarrow t \bar{b} j$	$m_{top}/4$	4	$25.62 \pm 0.01$	$30.96 \pm 0.06$
a.5	$pp \rightarrow tbjj$	$m_{top}/4$	4	$8.195 \pm 0.002$	$8.91 \pm 0.01$
b.1	$pp \rightarrow (W^+ \rightarrow) e^+ \nu_e$	$m_W$	5	$5072.5 \pm 2.9$	$6146.2 \pm 9.8$
b.2	$pp \rightarrow (W^+ \rightarrow) e^+ \nu_e j$	$m_W$	5	$828.4\pm0.8$	$1065.3 \pm 1.8$
b.3	$pp \rightarrow (W^+ \rightarrow) e^+ \nu_e jj$	$m_W$	5	$298.8 \pm 0.4$	$300.3 \pm 0.6$
b.4	$pp \rightarrow (\gamma^*/Z \rightarrow) e^+ e^-$	$m_Z$	5	$1007.0 \pm 0.1$	$1170.0 \pm 2.4$
b.5	$pp \rightarrow (\gamma^{*}/Z \rightarrow)e^{+}e^{-}j$	$m_Z$	5	$156.11 \pm 0.03$	$203.0 \pm 0.2$
b.6	$pp \to (\gamma^*/Z \to) e^+e^- jj$	$m_Z$	5	$54.24\pm0.02$	$56.69\pm0.07$
c.1	$pp \rightarrow (W^+ \rightarrow) e^+ \nu_e b \bar{b}$	$m_W+2m_b$	4	$11.557 \pm 0.005$	$22.95 \pm 0.07$
c.2	$pp \rightarrow (W^+ \rightarrow) e^+ \nu_e t \bar{t}$	$m_W + 2m_{top}$	5	$0.009415 \pm 0.000003$	$0.01159 \pm 0.00001$
c.3	$pp \rightarrow (\gamma^*/Z \rightarrow) e^+ e^- b\bar{b}$	$m_Z + 2m_b$	4	$9.459 \pm 0.004$	$15.31\pm0.03$
c.4	$pp \! \to \! (\gamma^*/Z \to) e^+e^-tt$	$m_Z + 2m_{top}$	5	$0.0035131 \pm 0.0000004$	$0.004876 \pm 0.000000$
c.5	$pp \rightarrow \gamma t \bar{t}$	$2m_{top}$	5	$0.2906 \pm 0.0001$	$0.4169 \pm 0.0003$
d.1	$pp \to W^+W^-$	$2m_W$	4	$29.976 \pm 0.004$	$43.92 \pm 0.03$
d.2	$pp \to W^+W^-j$	$2m_W$	4	$11.613 \pm 0.002$	$15.174 \pm 0.008$
d.3	$pp \to W^+W^+jj$	$2m_{W}$	4	$0.07048 \pm 0.00004$	$0.1377 \pm 0.0005$
e.1	$pp \rightarrow HW^+$	$m_W + m_H$	5	$0.3428 \pm 0.0003$	$0.4455 \pm 0.0003$
e.2	$pp \rightarrow HW^+j$	$m_W + m_H$	5	$0.1223 \pm 0.0001$	$0.1501 \pm 0.0002$
e.3	$pp \rightarrow HZ$	$m_Z + m_H$	5	$0.2781 \pm 0.0001$	$0.3659 \pm 0.0002$
e.4	$pp \rightarrow HZj$	$m_Z + m_H$	5	$0.0988 \pm 0.0001$	$0.1237 \pm 0.0001$
e.5	$pp \rightarrow Ht\bar{t}$	$m_{top} + m_H$	5	$0.08896 \pm 0.00001$	$0.09869 \pm 0.00003$
e.6	$pp \rightarrow Hb\overline{b}$	$m_b + m_H$	4	$0.16510 \pm 0.00009$	$0.2099 \pm 0.0006$
e.7	$pp \rightarrow Hjj$	$m_H$	5	$1.104 \pm 0.002$	$1.036 \pm 0.002$

## DISTRIBUTIONS

#### FULL MACHINERY AT WORK

angle Case study of  $[H/A]t\bar{t}$  with starring actors:

MGv4, CT, MadFKS, MadLoop and aMC@NLO interfaced to Herwig6!



Solid: aMC@NLO scalar. Dashed: aMC@NLO pseudoscalar

Dotted: NLO scalar. Dotdashed: NLO pseudoscalar

Left:  $t \bar{t}$  invariant mass. Right:  $t \bar{t} H p_{\scriptscriptstyle T}$ 

## **FUTURE PLANS**

#### MGV5 TO THE RESCUE

## \* MG5 much more advantageous:

Extreme programming and python leads to modular and flexible code

Test routines makes maintenance easy

Designed from the beginning for being fast and user-friendly.

Output in many formats and automatic writing of HELAS subroutines.

## Short term plans:

MadLoop @ MGv5 with the recycling optimizations

Get rid of ALL existing constraints.

## \* Further term plans:

Automate generation UV and R2 terms in FeynRules.

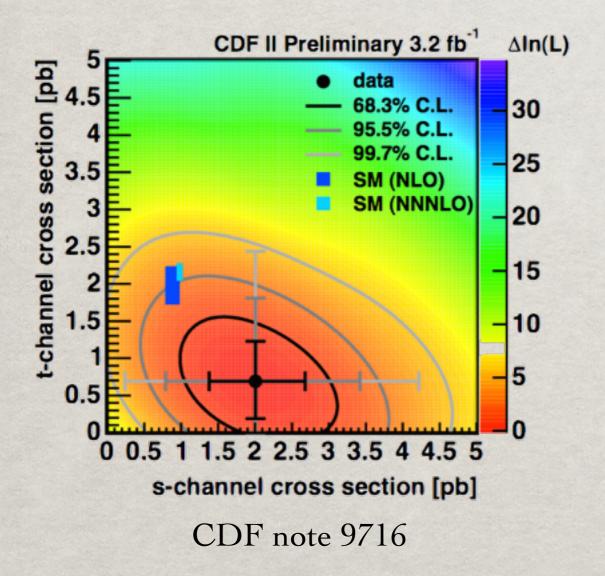
Recursion relations and computation of color ordered amplitudes.

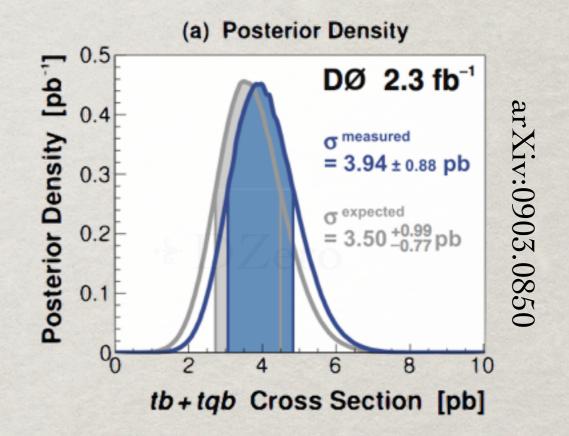
# THANKS!

# ADDITIONAL SLIDES

## SINGLE-TOP ANALYSIS

DZero and CDF both observe significant signal for single-top





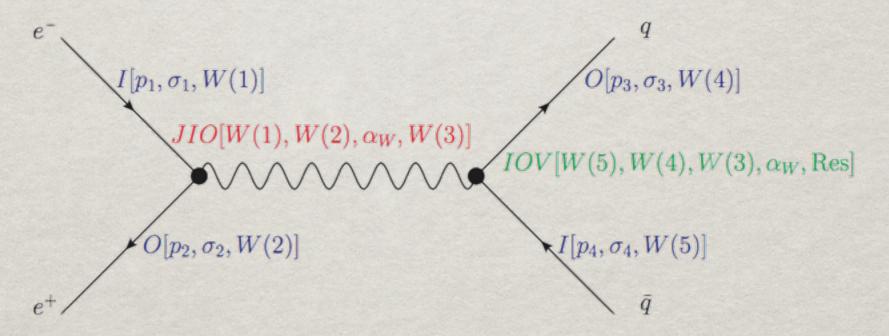
2σ-discrepancy pointed out by CDF on the {s,t}-channel plane

- New Physics, stat. fluctuations?
- \* Mistake in prediction?
  - → Mistreatment of t-channel in 4F/5F scheme?

## MADGRAPH

#### THE EVOLUTIVE WAY OF COMPUTING TREE-DIAGRAMS

- First generates all tree-level Feynman Diagrams
- \* Compute the amplitude of each diagram using a chain of calls to HELAS subroutines



Finally square all the related amplitude with their right color factors to construct the full LO amplitude