

# $m_W$ in the SMEFT

*[E. Bagnaschi, J. Ellis, M. Madigan, KM, V. Sanz, T. You; JHEP 08 (2022) 308]*

*J. Ellis, KM, F. Zampedri; 2304.06663]*

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MWDays23 workshop, CERN

19<sup>th</sup> of April 2023

# Outline

Predicting  $m_W$  & interpreting measurements in the SMEFT

Global fits with EWPO, Higgs & Diboson data

Mapping to single field extensions of the SM

Interplay with CKM unitarity

Beyond dimension-6: scalar triplet model

# SMEFT is...

$$\mathcal{L}_{\text{eff}} = \sum_i \frac{c_i \mathcal{O}_i^D}{\Lambda^{D-4}}$$

## Model independent

- Underlying assumptions

*Heavy new physics:  $M > E_{\text{exp}}$   
SM field content & gauge symmetries  
Linear EWSB: Higgs = doublet*

## Systematically improvable

- Double expansion *higher dim.*  $\frac{E^2}{\Lambda^2}$  &  $\{g_s, g, g'\}$  *more loops*

## Global

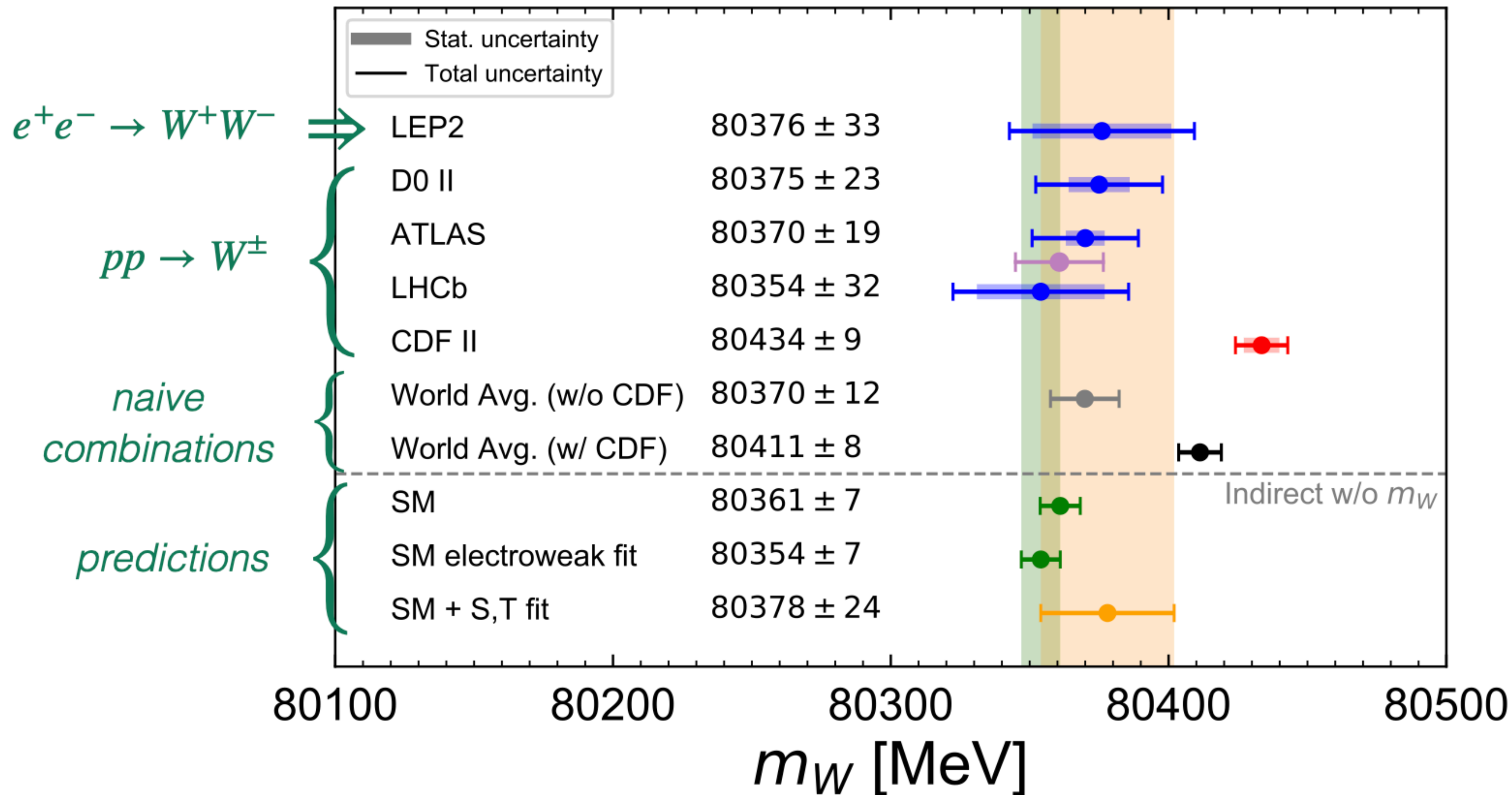
- Model independence: we don't know what operators NP will generate
- *Patterns & correlations* among observables are key
- Ultimate goal: complete SMEFT likelihood confronted with HEP data

EWPO, *Higgs*, *multiboson*, *top*, DY, *flavor*,...

$\mathcal{L}(c_i) \Rightarrow$  **indirectly constrain many UV models**

# $m_W$ measurements

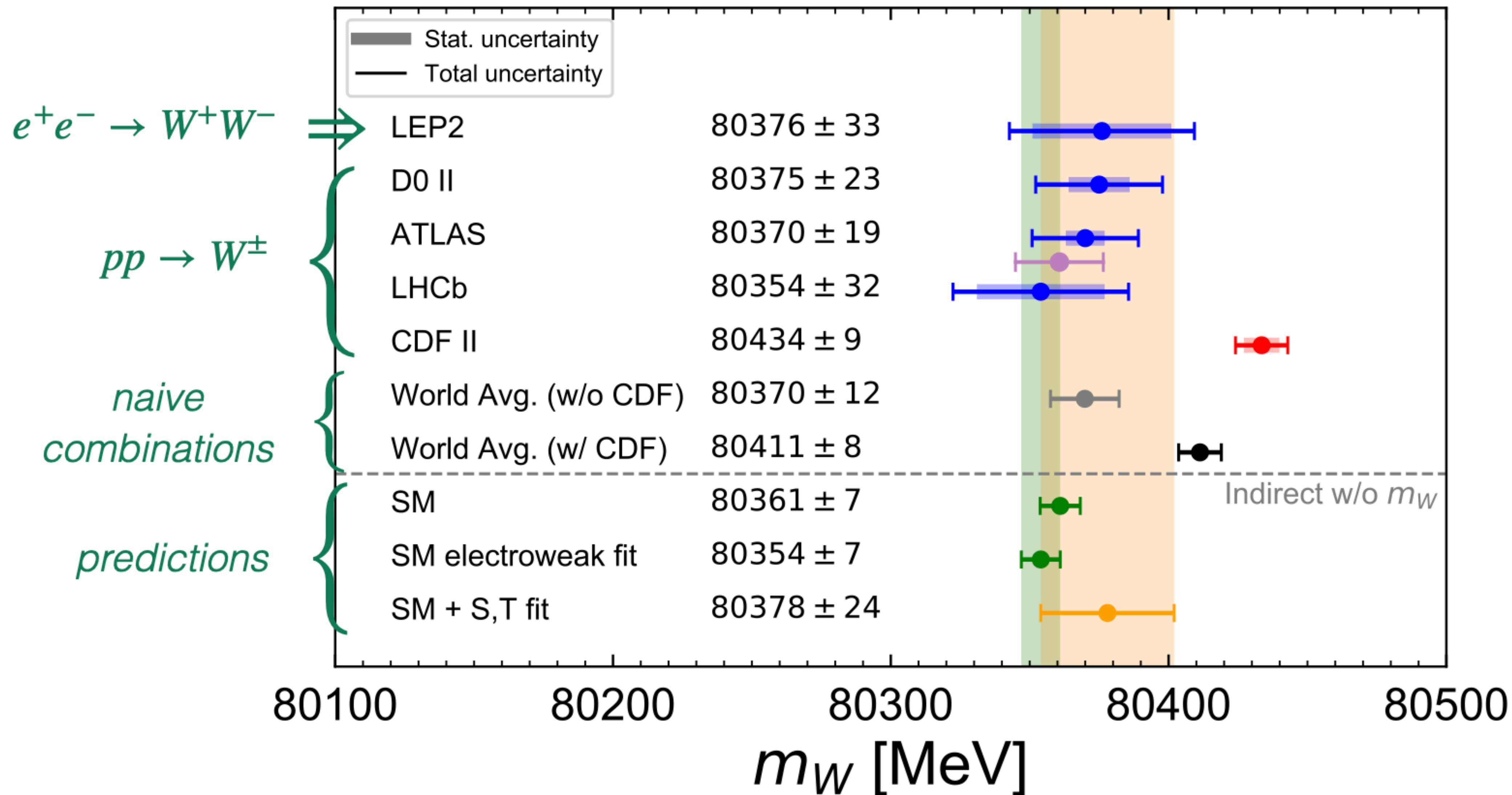
[Bagnaschi, Ellis, Madigan, KM, Sanz & You; JHEP 08 (2022) 308]



ATLAS update, Run 1 re-analysis:  $m_W = 80360 \pm 16$  MeV

# $m_W$ measurements

[Bagnaschi, Ellis, Madigan, KM, Sanz & You; JHEP 08 (2022) 308]



SM + fit to S & T parameters:  $m_W = 80378 \pm 24$  MeV


# $m_W$ prediction in SMEFT


$$\frac{\delta m_W^2}{m_W^2} = -\frac{\sin 2\hat{\theta}_w}{\cos 2\hat{\theta}_w} \frac{\hat{v}^2}{4\Lambda^2} \left( \frac{\cos \hat{\theta}_w}{\sin \hat{\theta}_w} C_{HD} + \frac{\sin \hat{\theta}_w}{\cos \hat{\theta}_w} (4C_{HI}^{(3)} - 2C_{ll}) + 4C_{HWB} \right)$$

Dimension-6: 4 operators modify  $m_W(G_F, \alpha_{EW}, m_Z)$

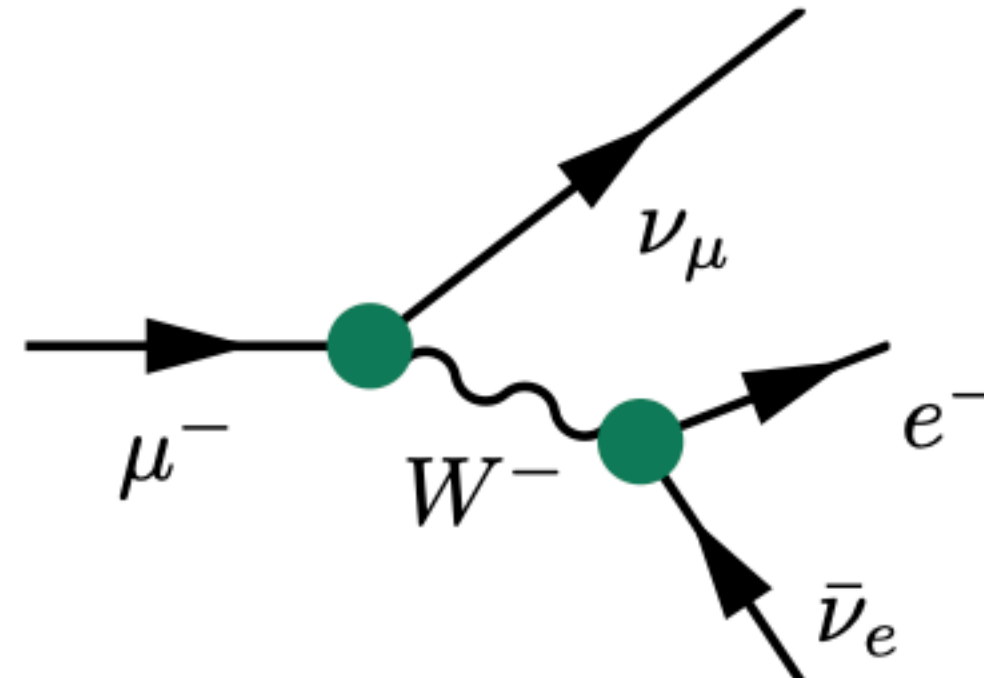
- Shift input parameter relations w.r.t SM  $\Rightarrow m_W(G_F, \alpha_{EW}, m_Z, C_i)$

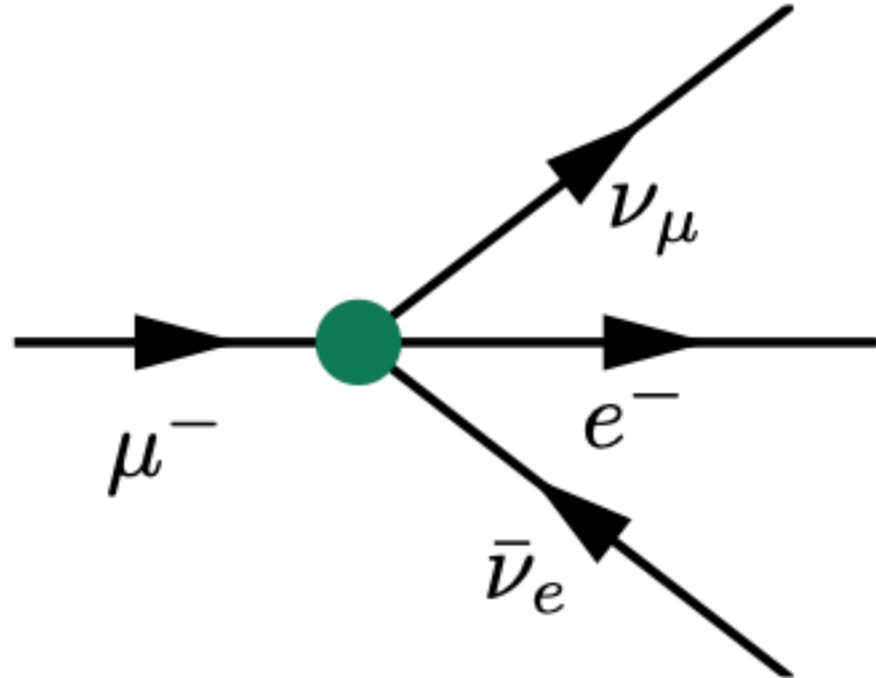
$\Delta m_Z^2$

$\mathcal{O}_{HWB} \equiv H^\dagger \tau^I H W_{\mu\nu}^I B^{\mu\nu} \propto S$ 


$\mathcal{O}_{HD} \equiv (H^\dagger D^\mu H)^\star (H^\dagger D_\mu H) \propto T$ 


$\Delta G_F$

$\mathcal{O}_{H\ell}^{(3)} \equiv (H^\dagger i \overleftrightarrow{D}_\mu^I H) (\bar{\ell}_{1,2} \tau^I \gamma^\mu \ell_{1,2})$ 


$\mathcal{O}_{\ell\ell} \equiv (\bar{\ell}_1 \gamma_\mu \ell_2) (\bar{\ell}_2 \gamma^\mu \ell_1)$ 


# $m_W$ interpretation in SMEFT

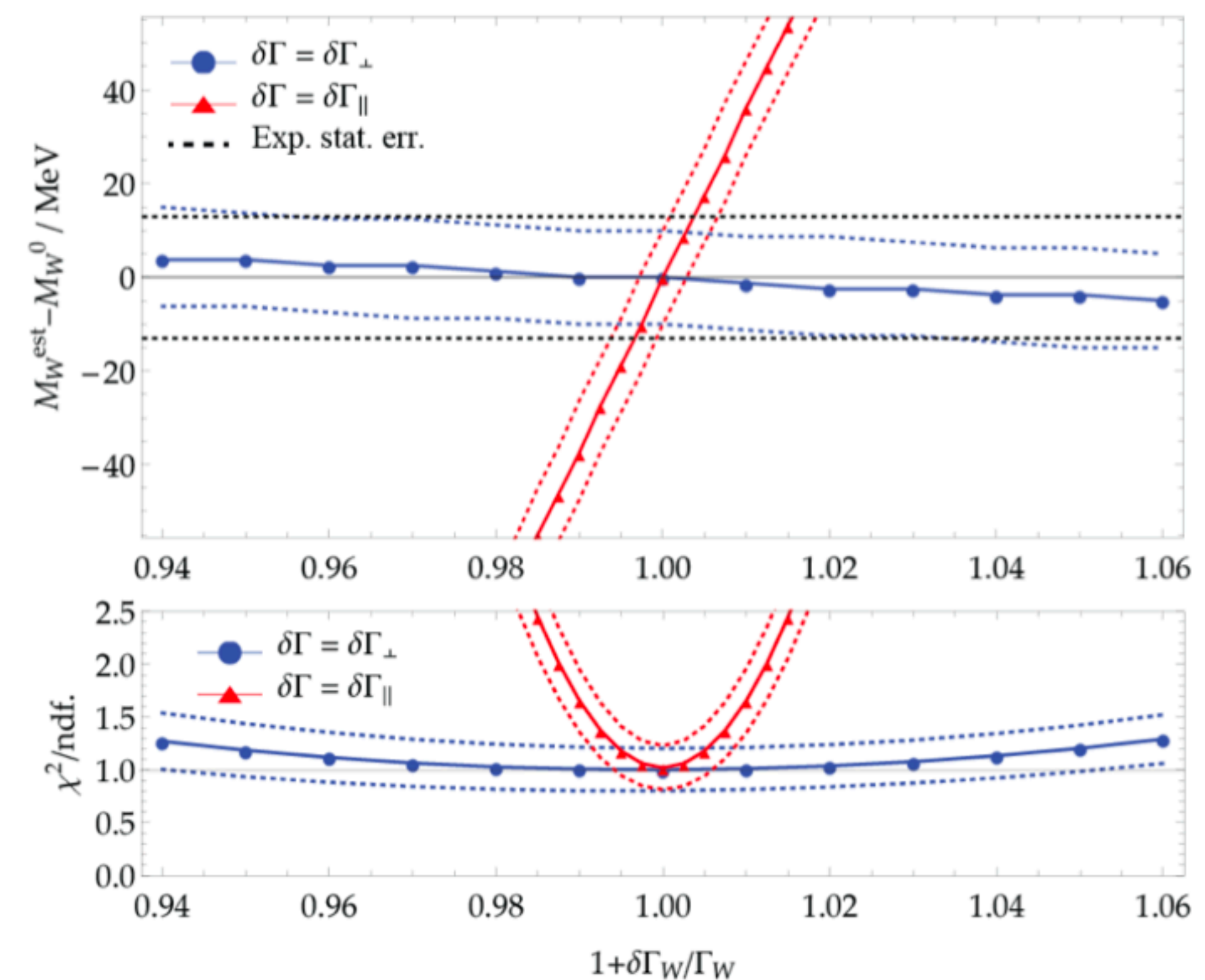
New physics can yield correlated effects in many places

- SMEFT motivates global approach  $\Rightarrow \delta m_W^2$  does not occur in isolation
- *e.g.* modified widths/branching fractions, production & decay kinematics, ...
- Biases in interpreting data that hasn't accounted for such effects?

$W$  boson width shift: 
$$\frac{\delta\Gamma_W}{\Gamma_W} = \frac{\hat{v}^2}{\Lambda^2} \left( \frac{8}{3} C_{Hl}^{(3)} + \frac{4}{3} C_{Hq}^{(3)} - C_{ll} \right) - \frac{3}{2} \frac{\delta m_W^2}{m_W^2}$$

[Bjorn & Trott; PLB 762 (2016) 426-431]

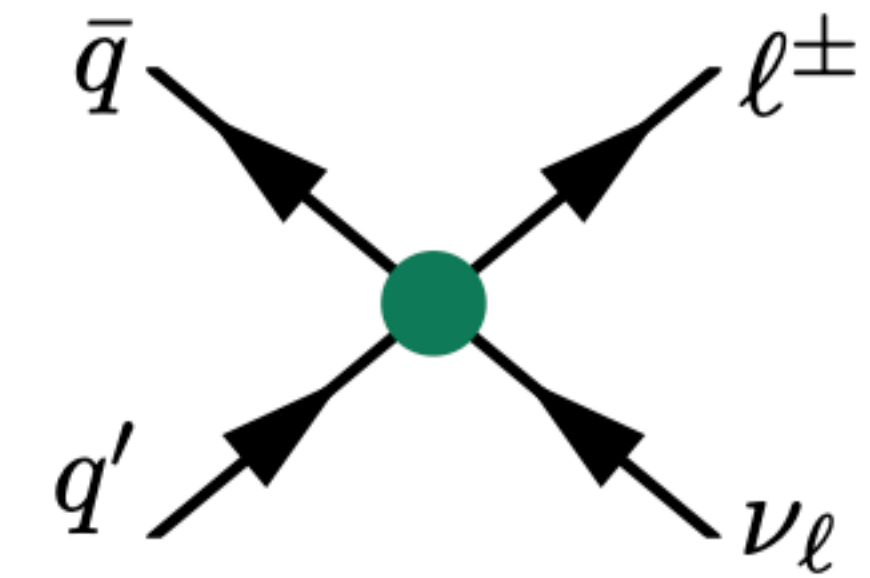
- Affects overall normalisation & shape of template  $m_T$  &  $p_T^l$  distributions
- Estimated bias dominated by normalisation, which is left to float in experimental analyses
- No significant bias expected from  $\Gamma_W$



# $m_W$ interpretation in SMEFT

## Other contact interactions?

- Interference effects with charged-current 4F operators
- May be overwhelmed by narrow  $W$  peak...
- Constrained by charged current Drell Yan. Not studied so far...



## LEP II extraction: $\sigma(e^+e^- \rightarrow W^+W^-)$

*As discussed in [Bjorn & Trott; PLB 762 (2016) 426-431]*

- Various methods to extract  $m_W$ , assuming the SM for  $\sigma(W^+W^-)$
- Other possible SMEFT effects not accounted for
- Triple gauge coupling modifications ( $\mathcal{O}_W$ ),  $W$ - coupling shifts, dipoles?

Safest to do a simultaneous fit to  $m_W$  &  $C_i$



# Input scheme

$m_W$  is a prediction when using, e.g.,  $\{\alpha_{EW}, G_F, m_Z\}$  as inputs

- We use:  $\alpha_{EW}^{-1} = 127.95$  ,  $G_F = 1.16638 \times 10^{-5} \text{ GeV}^{-2}$  ,  $m_Z = 91.1876 \text{ GeV}$

If  $m_W$  is part of the input set...

- Other data must be used to compensate
- e.g.,  $\{m_W, G_F, m_Z\}$  input scheme: use  $a_{EW}(m_Z)$  instead

$$\left. \frac{\delta\alpha_{EW}}{\alpha_{EW}} \right|_{\{m_W, G_F, m_Z\}} \propto \left. \frac{\delta m_W^2}{m_W^2} \right|_{\{\alpha_{EW}, G_F, m_Z\}}$$

$\alpha_{EW}$  prediction deviates from SM if CDF  $m_W$  is input

- $\alpha_{EW}$  uncertainty dominated by parametric  $m_W$  error
- Similar constraints expected in this case

# Global context

## Global new physics searches via high precision/energy

- **Z & W-pole data:** handle on the EW gauge sector [Han & Skiba; PRD 71 (2005) 075009]  
[Falkowski & Riva; JHEP 02 (2015) 039]
- **LHC:** thriving Higgs & top programmes
- Probing gauge interactions at high energy (**VV, VBS, VVV, ...**)

## Cross-talk between EWPO, Diboson & Higgs data

- Significant correlations among operators  $\Rightarrow$  data
- Allows for a **closed fit** to flavor-universal SMEFT

[Corbett et al.; PRD 87 (2013) 015022]

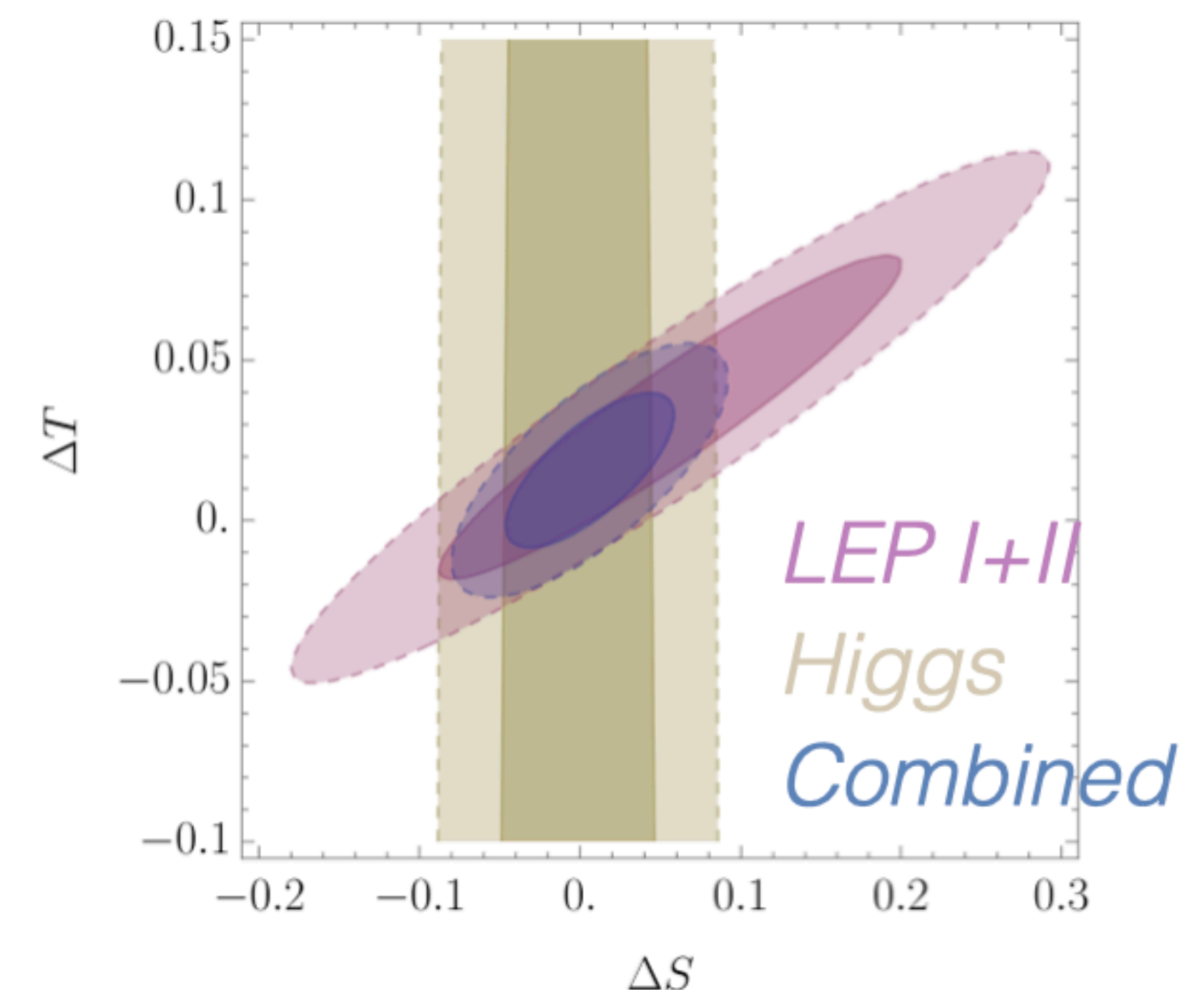
[Pomarol & Riva; JHEP 01 (2014) 151]

[Ellis, Sanz & You; JHEP 03 (2015) 157]

[Biekötter, Corbett & Plehn; SciPost Phys 6 (2019) 6, 064]

[Ellise, Madigan, KM, Sanz, You; JHEP 04 (2021) 279]...

**Can the SMEFT accommodate the  $m_W$  measurements alongside other data?**



[Ellis et al.; JHEP 06 (2018) 146]

# The fit

*fitmaker* <https://gitlab.com/kenmimasu/fitrepo>  
public-friendly version w/ example notebooks in progress

## Top, Higgs, Diboson and Electroweak Fit to the Standard Model Effective Field Theory

John Ellis,<sup>a,b,c</sup> Maeve Madigan,<sup>d</sup> Ken Mimasu,<sup>a</sup> Veronica Sanz<sup>e,f</sup> and Tevong You<sup>b,d,g</sup> [*JHEP* 04 (2021) 279]

## Global SMEFT interpretation of 4 categories of data

*Based on*

- 14 • Electroweak Precision Observables (EWPO): Z-pole & W-mass [*Ellis et al.; JHEP* 06 (2018) 146]
- 118 • LEP2 & LHC diboson production: differential WW, WZ, Zjj
- 72 • Higgs measurements: signal strengths & STXS
- ~~137 • Top data: single-top, ttbar & asymmetries, ttV, tZ, tW~~

## 204 measurements across categories

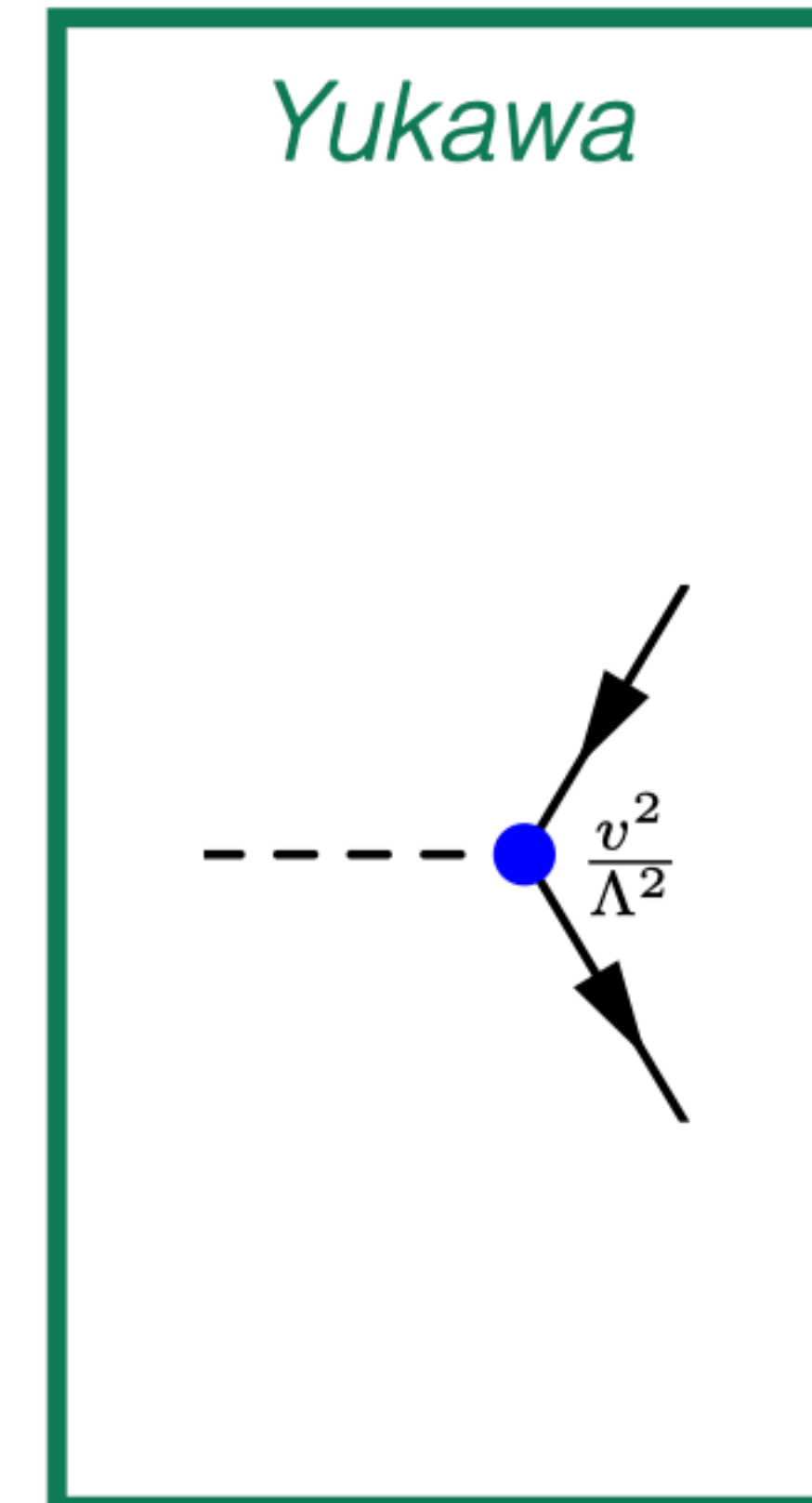
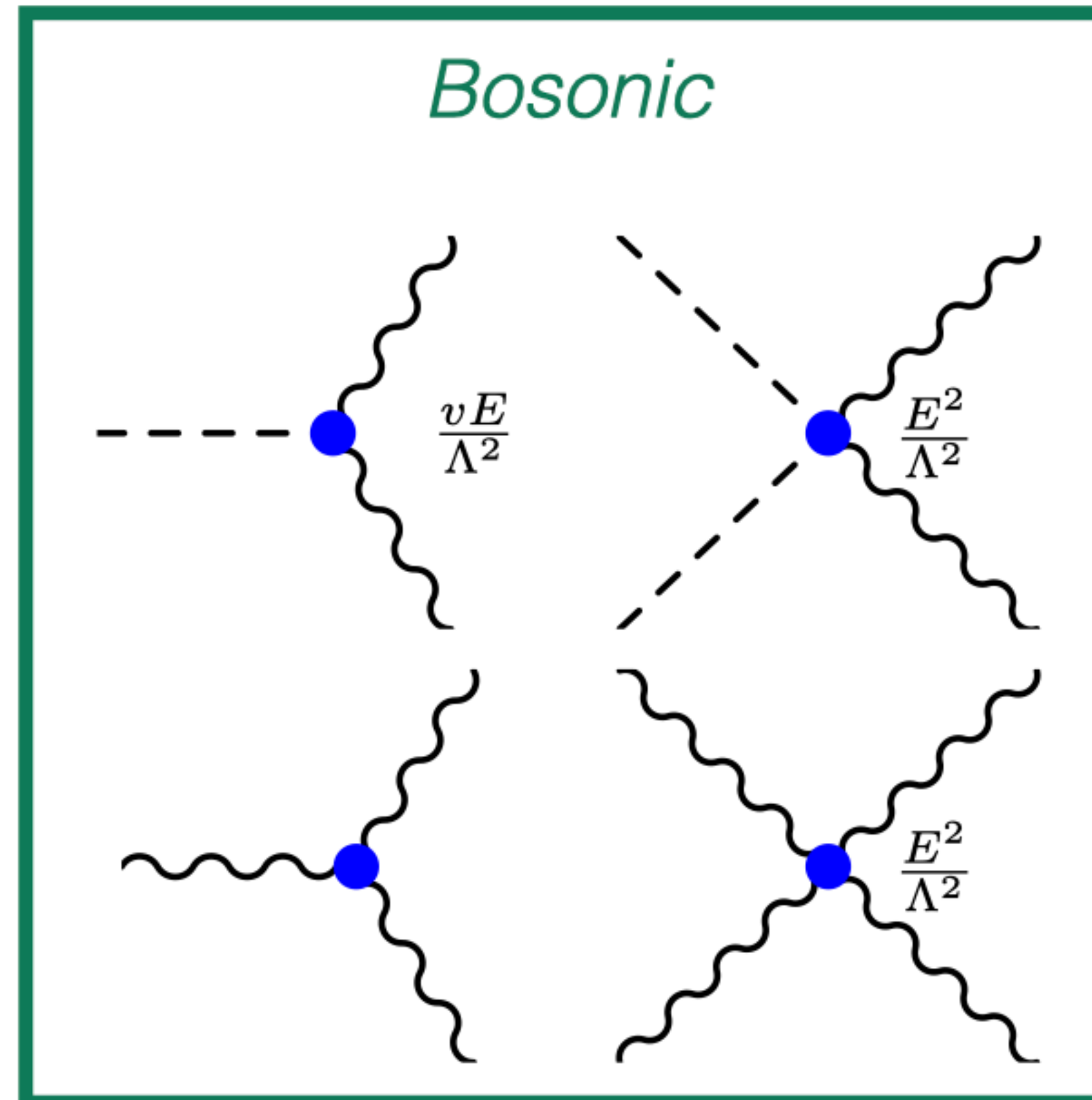
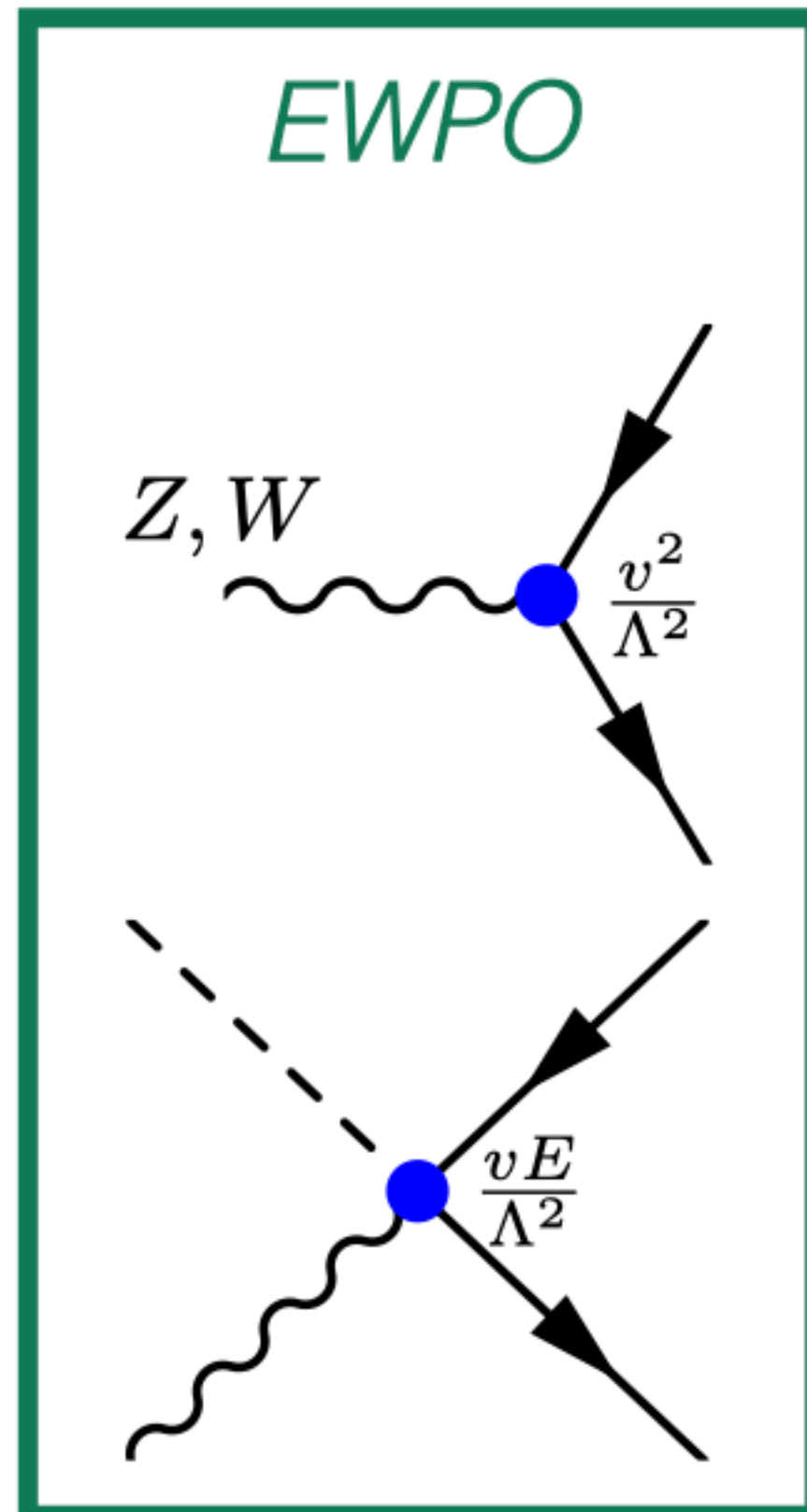
- Chosen to be statistically independent & maximise reach
- Correlations included when publicly available (mostly are)

Linear EFT approximation: 
$$\mu_X \equiv \frac{X}{X_{SM}} = 1 + \sum_i a_i^X \frac{C_i}{\Lambda^2} + \mathcal{O}\left(\frac{1}{\Lambda^4}\right)$$

# Degrees of freedom

EWPO:	$\mathcal{O}_{HWB}, \mathcal{O}_{HD}, \mathcal{O}_U, \mathcal{O}_{Hl}^{(3)}, \mathcal{O}_{Hl}^{(1)}, \mathcal{O}_{He}, \mathcal{O}_{Hq}^{(3)}, \mathcal{O}_{Hq}^{(1)}, \mathcal{O}_{Hd}, \mathcal{O}_{Hu}$
Bosonic:	$\mathcal{O}_{H\Box}, \mathcal{O}_{HG}, \mathcal{O}_{HW}, \mathcal{O}_{HB}, \mathcal{O}_W, \mathcal{O}_G$
Yukawa:	$\mathcal{O}_{\tau H}, \mathcal{O}_{\mu H}, \mathcal{O}_{bH}, \mathcal{O}_{tH}$
	20

Universal  $U(3)^5$  flavor scenario + Yukawa



# S & T parameter fit

$$\frac{v^2}{\Lambda^2} C_{HWB} = \frac{gg'}{16\pi} S$$

$$\frac{v^2}{\Lambda^2} C_{HD} = -\frac{g^2 g'^2}{2\pi(g^2 + g'^2)} T$$

No  $m_W$  data

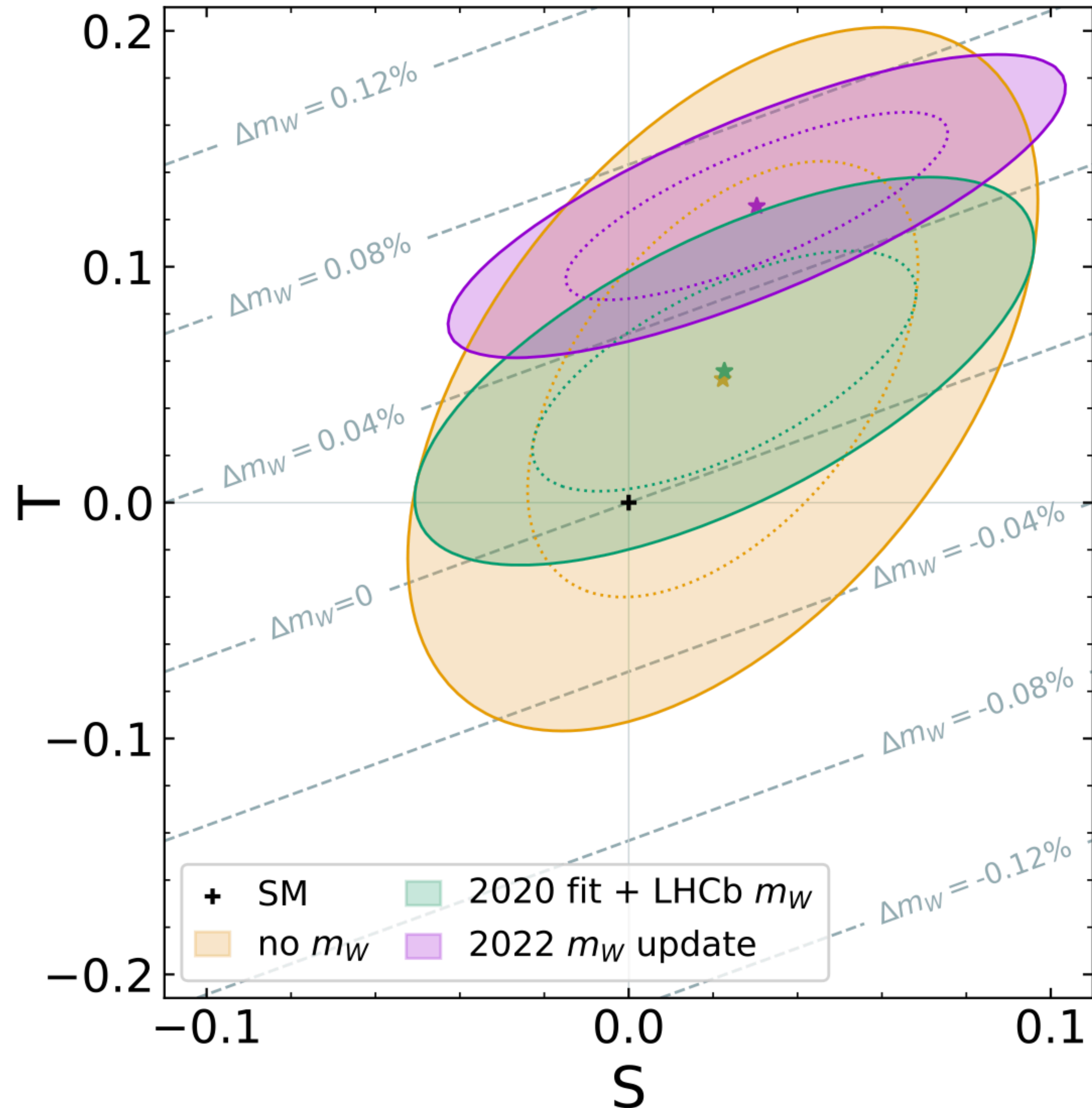
SM compatible  $< 1\sigma$

Without CDF  $m_W$

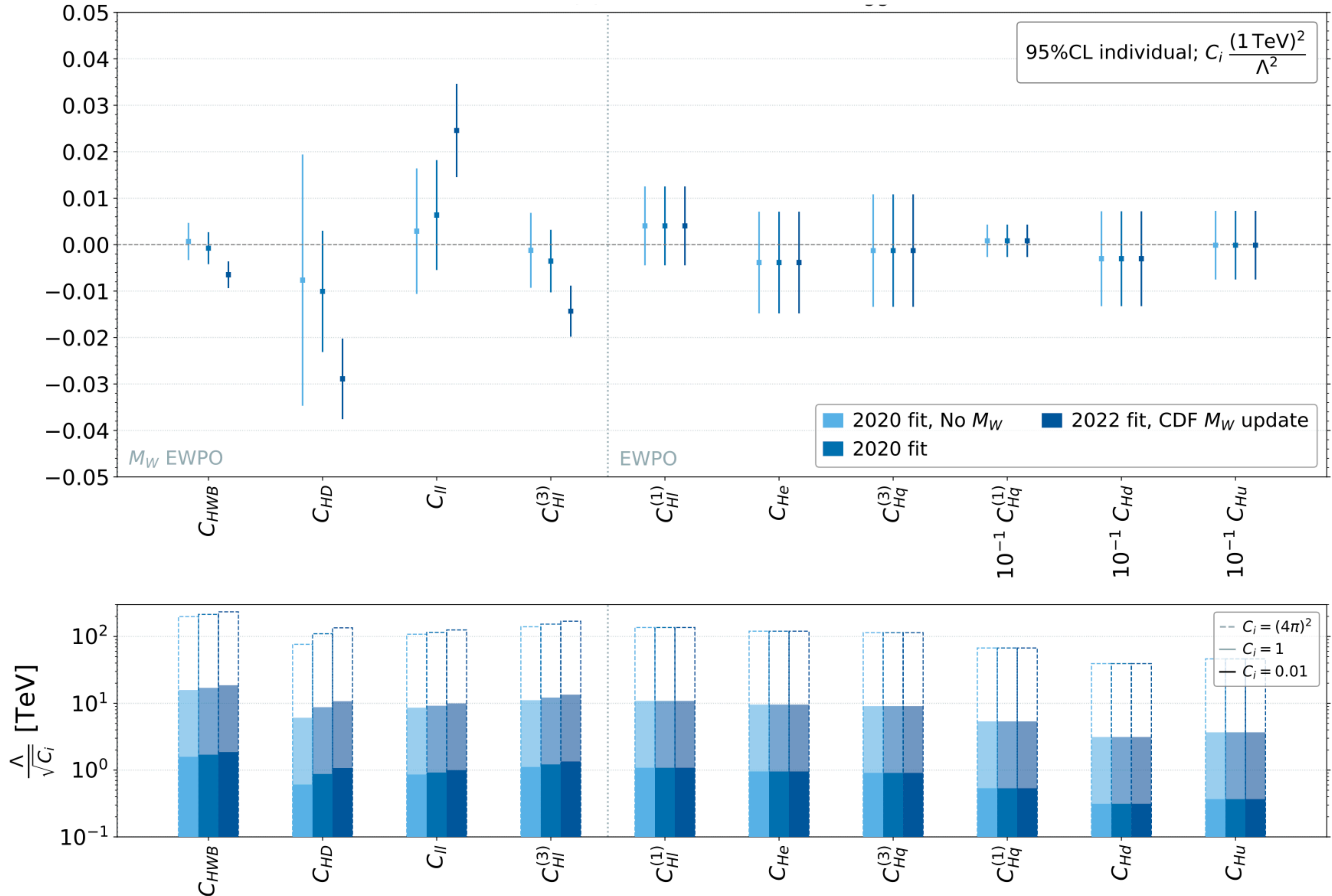
SM compatible  $< 2\sigma$

With CDF  $m_W$

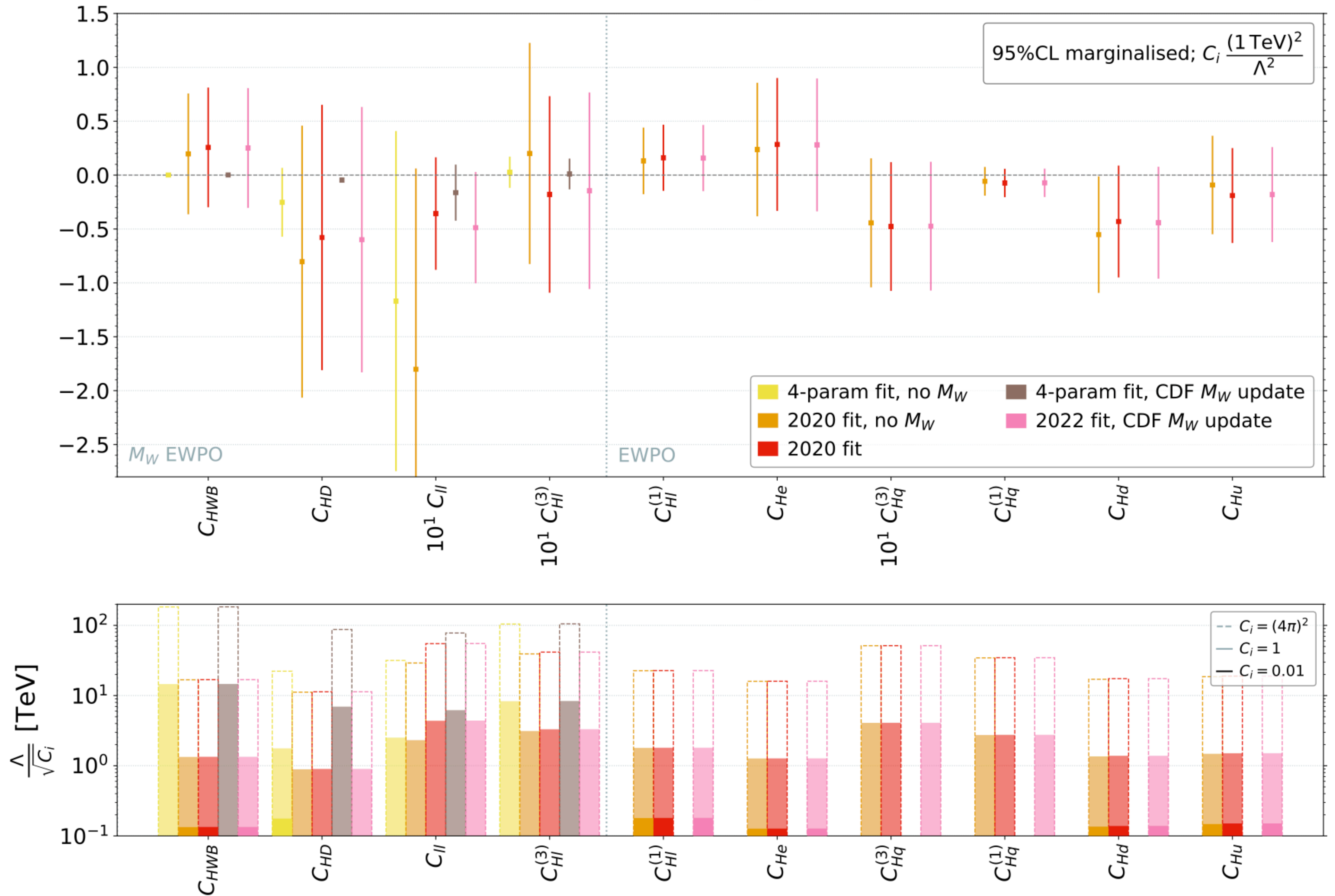
SM incompatible  $\gg 3\sigma$



# SMEFT fit (individual)

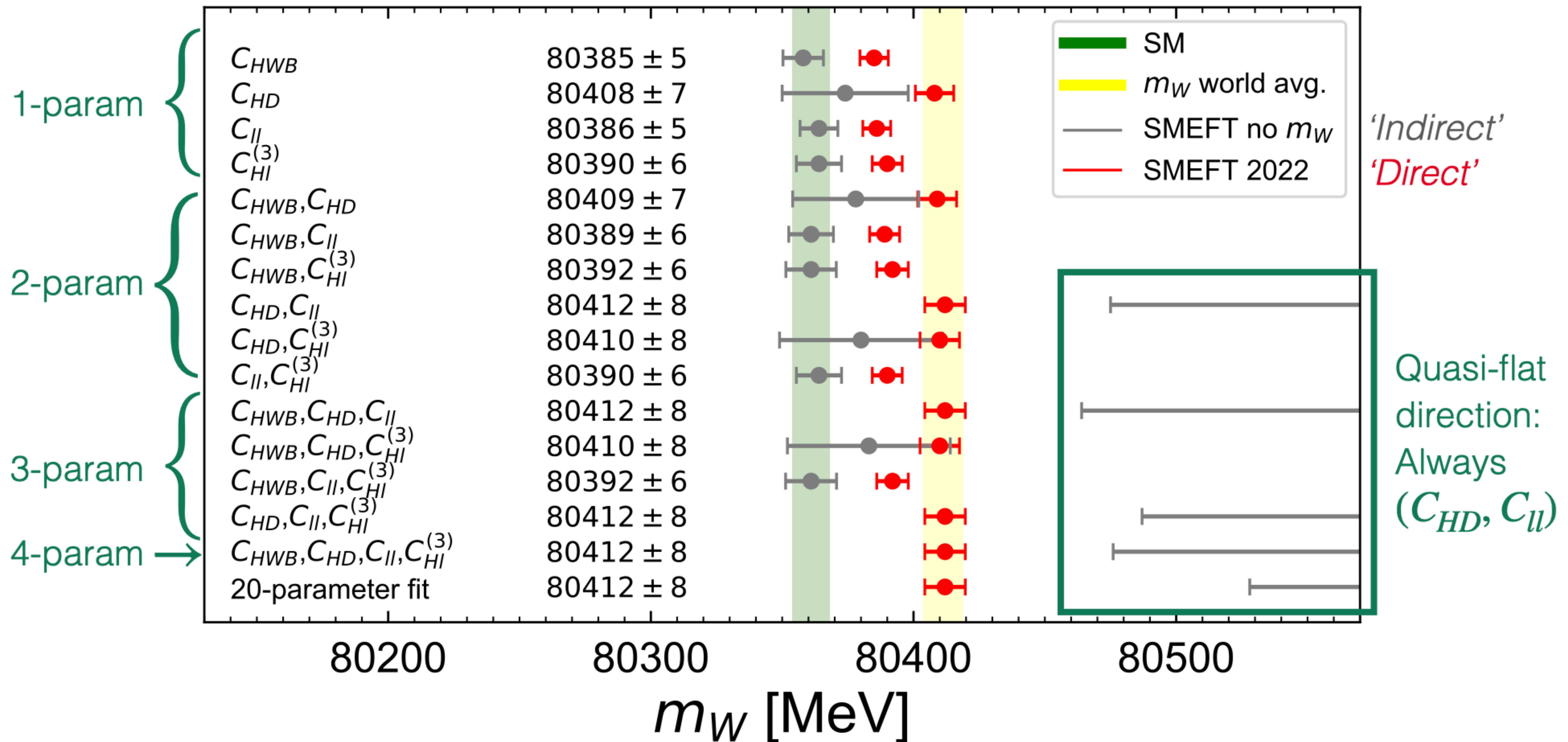


# SMEFT fit (marginalised)

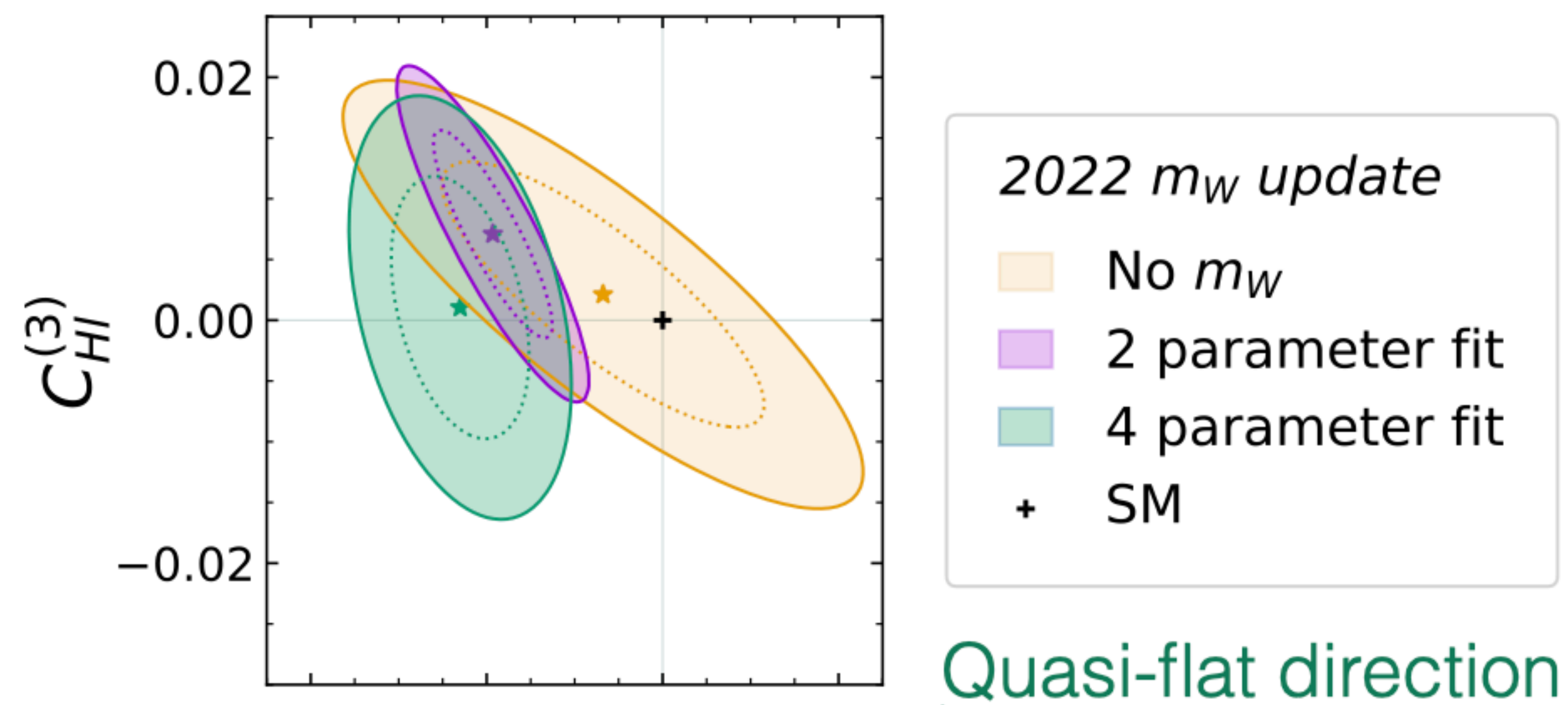


# Direct vs indirect $m_W$

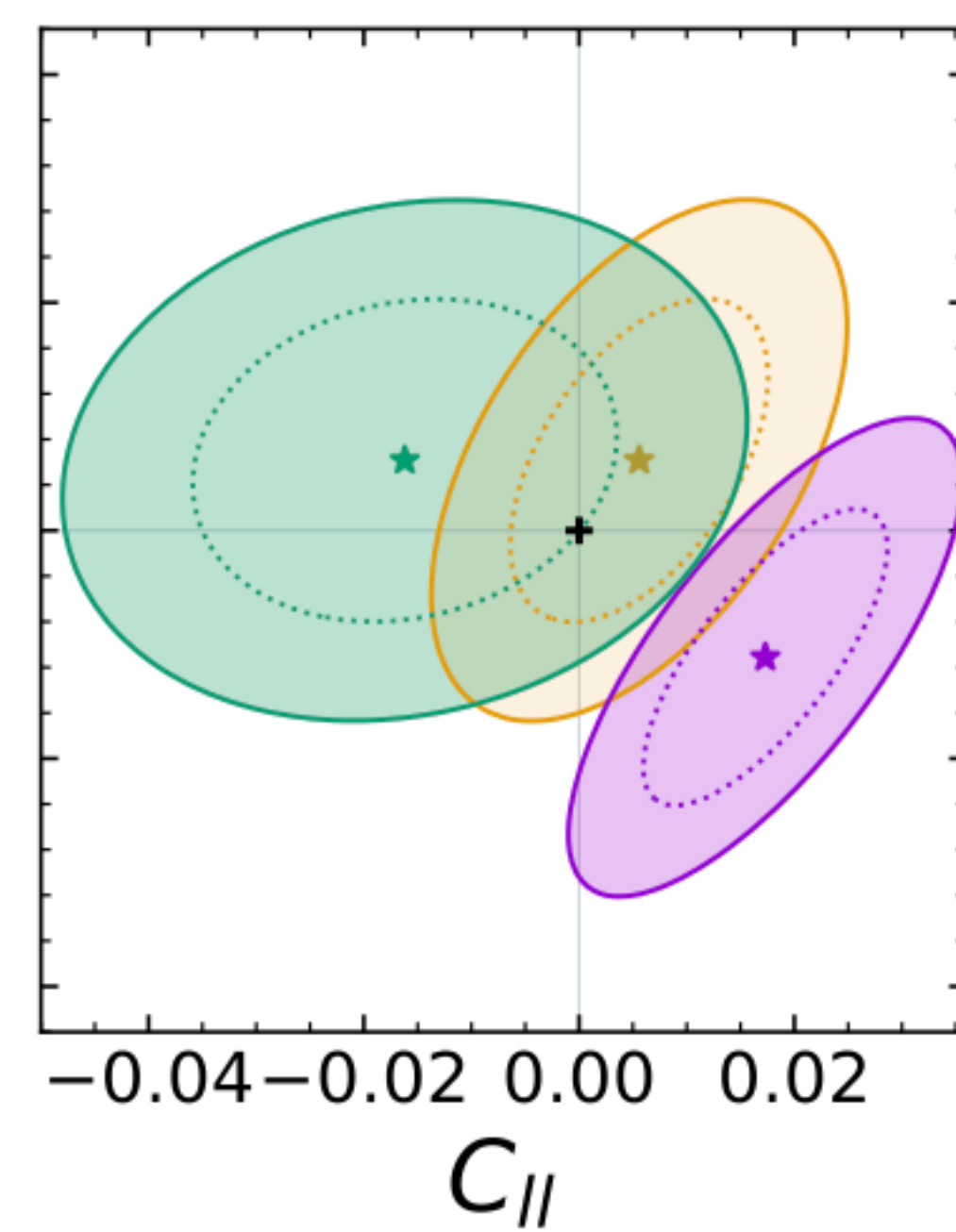
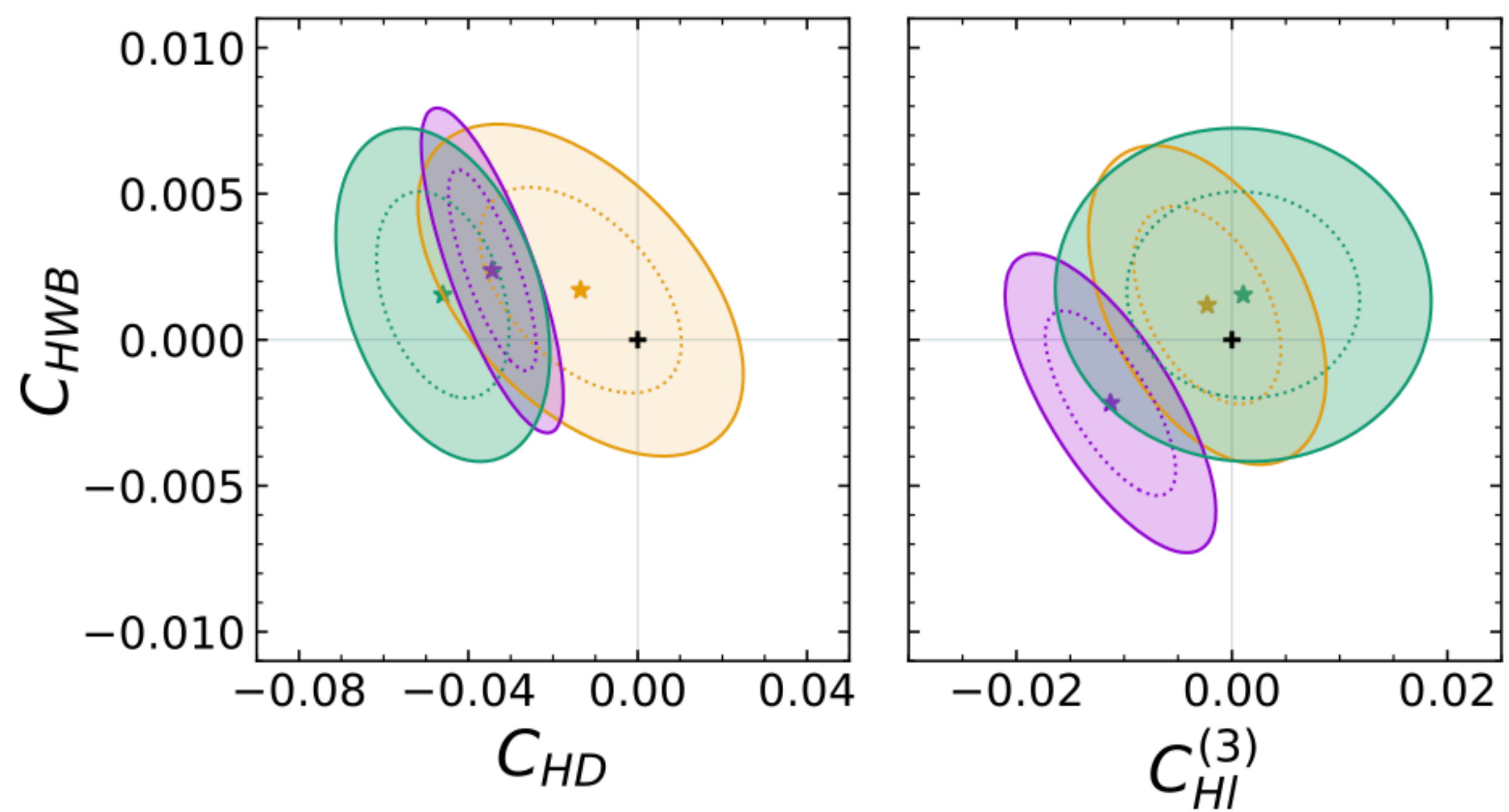
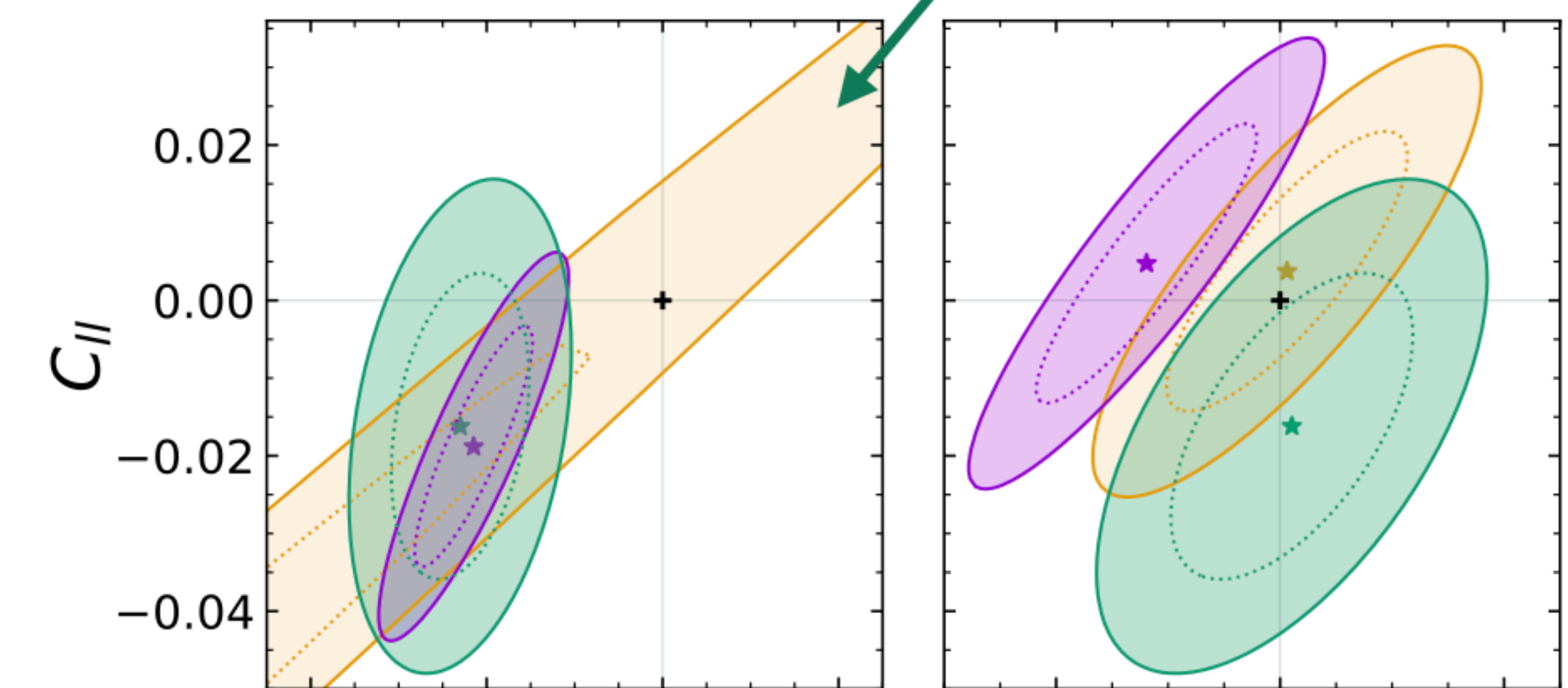
Focus on 4-parameter subspace: all possible combinations







$$\left( C_{HWB}, C_{HD}, C_{HI}^{(3)}, C_{II} \right)$$



# Single-field SM extensions

## What SM extensions can account for the anomaly at tree-level?

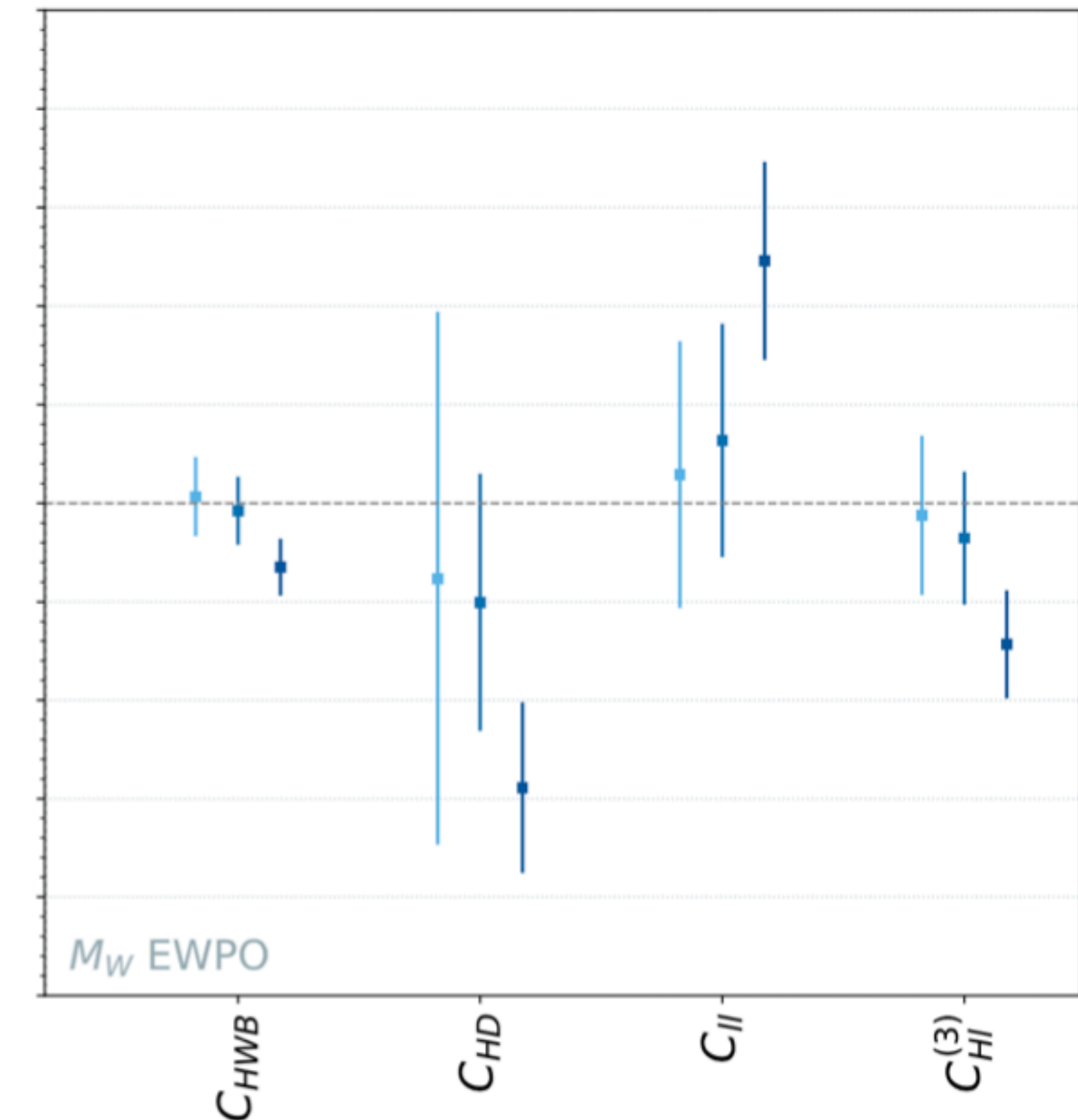
Fit preference:

- Negative  $C_{HWB}$ ,  $C_{HD}$ ,  $C_{HI}^{(3)}$  & positive  $C_{II}$

## Considered single field extensions

[de Blas et al.; JHEP 03 (2018) 109]

- Complete tree-level matching dictionary is known
- Interpret in terms of simplified 1 parameter versions of the models



Model	Spin	SU(3)	SU(2)	U(1)	Parameters
$S_1$	0	1	1	1	$(M_S, \kappa_S)$
$\Sigma$	$\frac{1}{2}$	1	3	0	$(M_\Sigma, \lambda_\Sigma)$
$\Sigma_1$	$\frac{1}{2}$	1	3	-1	$(M_{\Sigma_1}, \lambda_{\Sigma_1})$
$N$	$\frac{1}{2}$	1	1	0	$(M_N, \lambda_N)$
$E$	$\frac{1}{2}$	1	1	-1	$(M_E, \lambda_E)$

Model	Spin	SU(3)	SU(2)	U(1)	Parameters
$B$	1	1	1	0	$(M_B, \hat{g}_H^B)$
$B_1$	1	1	1	1	$(M_{B_1}, \lambda_{B_1})$
$\Xi$	0	1	3	0	$(M_\Xi, \kappa_\Xi)$
$W_1$	1	1	3	1	$(M_{W_1}, \hat{g}_{W_1}^\varphi)$
$W$	1	1	3	0	$(M_W, \hat{g}_W^H)$

# Dimension-6 matching

$$\delta m_W$$

Model	$C_{HD}$	$C_{ll}$	$C_{Hl}^{(3)}$	$C_{Hl}^{(1)}$	$C_{He}$	$C_{H\Box}$	$C_{\tau H}$	$C_{tH}$	$C_{bH}$
$S_1$		-1							
$\Sigma$			$\frac{1}{16}$	$\frac{3}{16}$			$\frac{y_\tau}{4}$		
$\Sigma_1$			$\frac{1}{16}$	$-\frac{3}{16}$			$\frac{y_\tau}{8}$		
$N$			$-\frac{1}{4}$	$\frac{1}{4}$					
$E$			$-\frac{1}{4}$	$-\frac{1}{4}$			$\frac{y_\tau}{2}$		
$B_1$	1					$-\frac{1}{2}$	$-\frac{y_\tau}{2}$	$-\frac{y_t}{2}$	$-\frac{y_b}{2}$
$B$	-2						$-y_\tau$	$-y_t$	$-y_b$
$\Xi$	$-2 \left(\frac{1}{M_\Xi}\right)^2$					$\frac{1}{2} \left(\frac{1}{M_\Xi}\right)^2$	$y_\tau \left(\frac{1}{M_\Xi}\right)^2$	$y_t \left(\frac{1}{M_\Xi}\right)^2$	$y_b \left(\frac{1}{M_\Xi}\right)^2$
$W_1$	$-\frac{1}{4}$					$-\frac{1}{8}$	$-\frac{y_\tau}{8}$	$-\frac{y_t}{8}$	$-\frac{y_b}{8}$
$W$	$\frac{1}{2}$					$-\frac{1}{2}$	$-y_\tau$	$-y_t$	$-y_b$

VLL

Z'

Spin-0,1  
Triplets

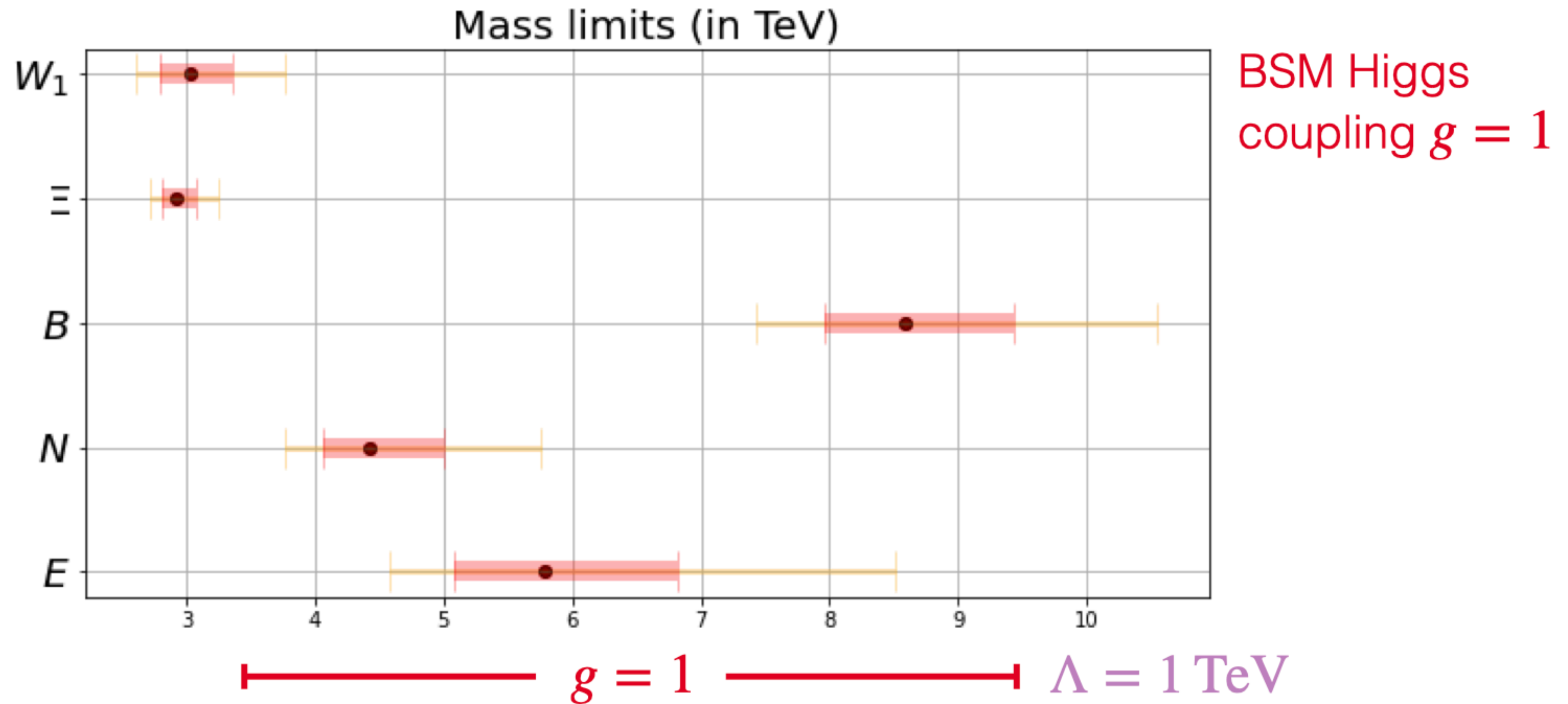
$C_{HWB}$  is loop-generated by weakly coupled models 

[Arzt, Einhorn & Wudka; Nucl.Phys.B 433 (1995) 41-66]

All models can generate one of the required coefficients

\*Only 5 predict them with the right sign

# Masses and couplings



Model	Pull	Best-fit mass (TeV)	1- $\sigma$ mass range (TeV)	2- $\sigma$ mass range (TeV)	1- $\sigma$ coupling <sup>2</sup> range
$W_1$	6.4	3.0	[2.8, 3.6]	[2.6, 3.8]	[0.09, 0.13]
$B$	6.4	8.6	[8.0, 9.4]	[7.4, 10.6]	[0.011, 0.016]
$\Xi$	6.4	2.9	[2.8, 3.1]	[2.7, 3.2]	[0.011, 0.016]
$N$	5.1	4.4	[4.1, 5.0]	[3.8, 5.8]	[0.040, 0.060]
$E$	3.5	5.8	[5.1, 6.8]	[4.6, 8.5]	[0.022, 0.039]

# CKM unitarity

$\beta$ -decay + CKM unitarity imposes significant constraint on one combination of coefficients in  $U(3)^5$

- Semi-leptonic analogue of muon decay:

$$\Delta_{CKM} \equiv |V_{ud}|^2 + |V_{us}|^2 - 1 = \frac{2\hat{v}^2}{\Lambda^2} \left( C_{Hq}^{(3)} - C_{Hl}^{(3)} + C_{ll} - C_{lq}^{(3)} \right)$$

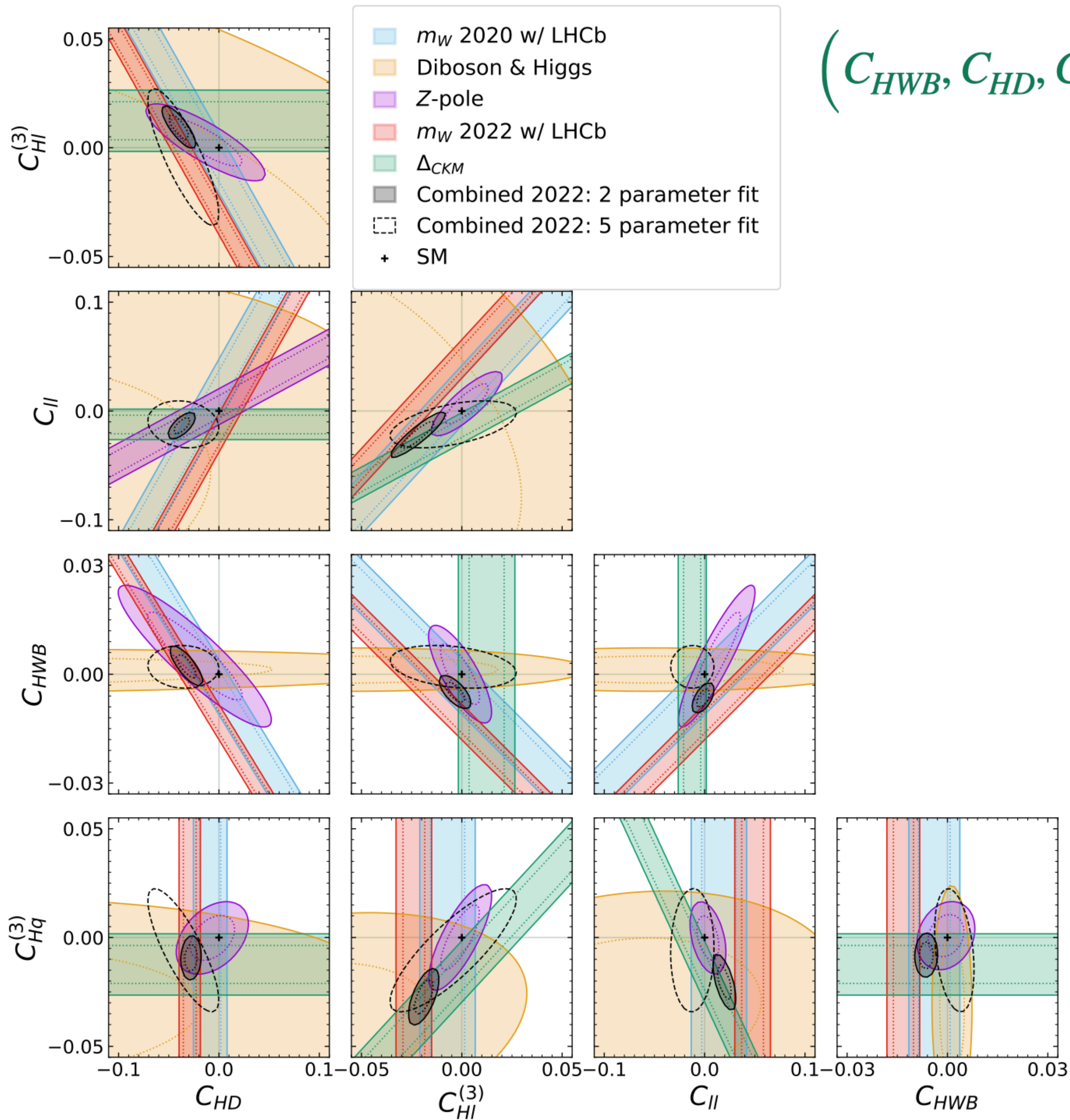
extra  $q\bar{q}\ell\bar{\ell}$   
4F operator

- Measurements of nuclear transitions and kaon decays indicate:

$$\Delta_{CKM} = -0.0015 \pm 0.0007$$

$\Delta_{CKM}$  probes a direction that is correlated with  $\delta m_W^2$

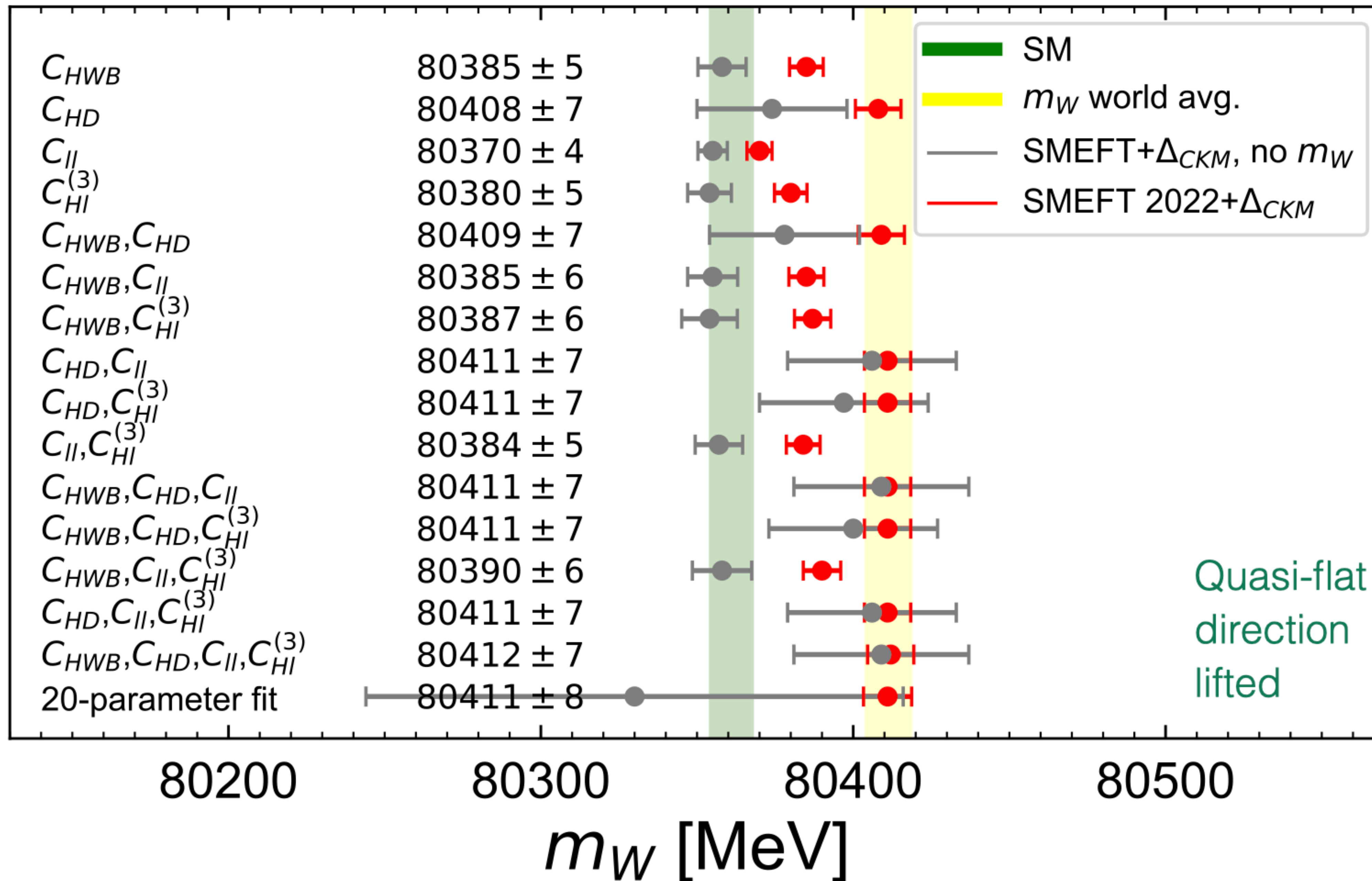
- Per-mille level constraint should compete with  $m_W \Rightarrow$  new information
- Also sensitive to irrelevant parameters for  $m_W$ , bring correlations with, e.g., Drell Yan & other EWPO, Diboson rates...



$$\left( C_{HWB}, C_{HD}, C_{HI}^{(3)}, C_{II}, C_{Hq}^{(3)} \right)$$

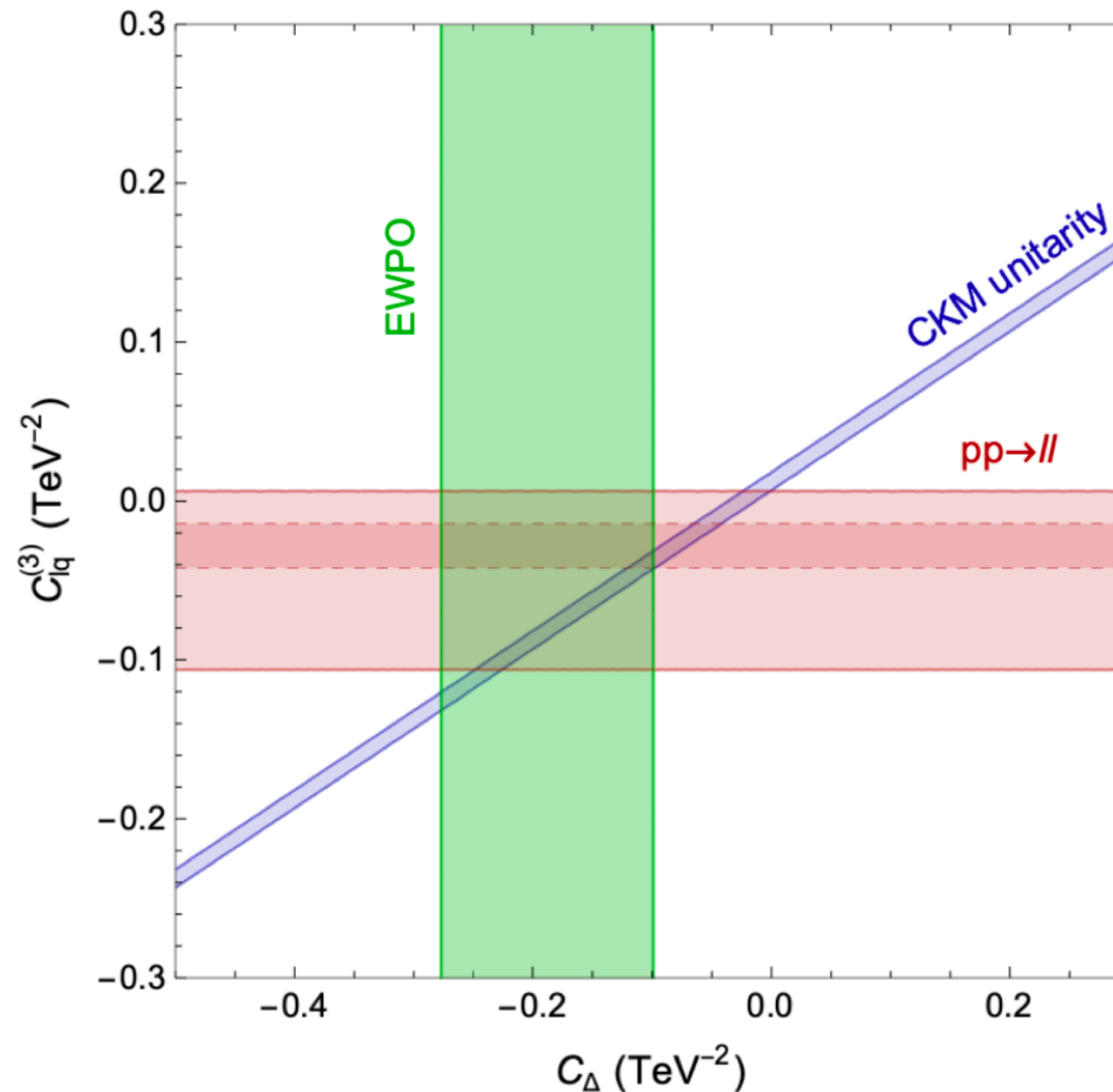
$$C_{ql}^{(3)} = 0$$

# $m_W$ after CKM unitarity



# CKM unitarity & DY

[Cirigliano, Dekens, de Vries, Mereghetti & Tong; 2204.08440]





# Beyond Dimension-6

Are dimension-8 effects important?

Precision?

EFT validity?

- Explicit study of  $Y=0$  triplet scalar,  $\Xi$

$$\mathcal{L}_\Xi = \frac{1}{2}(D_\mu \Xi^a)(D^\mu \Xi^a) - \frac{1}{2}M_\Xi^2(\Xi^a \Xi^a) - \kappa_\Xi H^\dagger \Xi^a \sigma^a H - \lambda_\Xi (\Xi^a \Xi^a)(H^\dagger H) - \frac{1}{4}\eta_\Xi (\Xi^a \Xi^a)^2$$

- $m_W$  dominates the constraints on this model

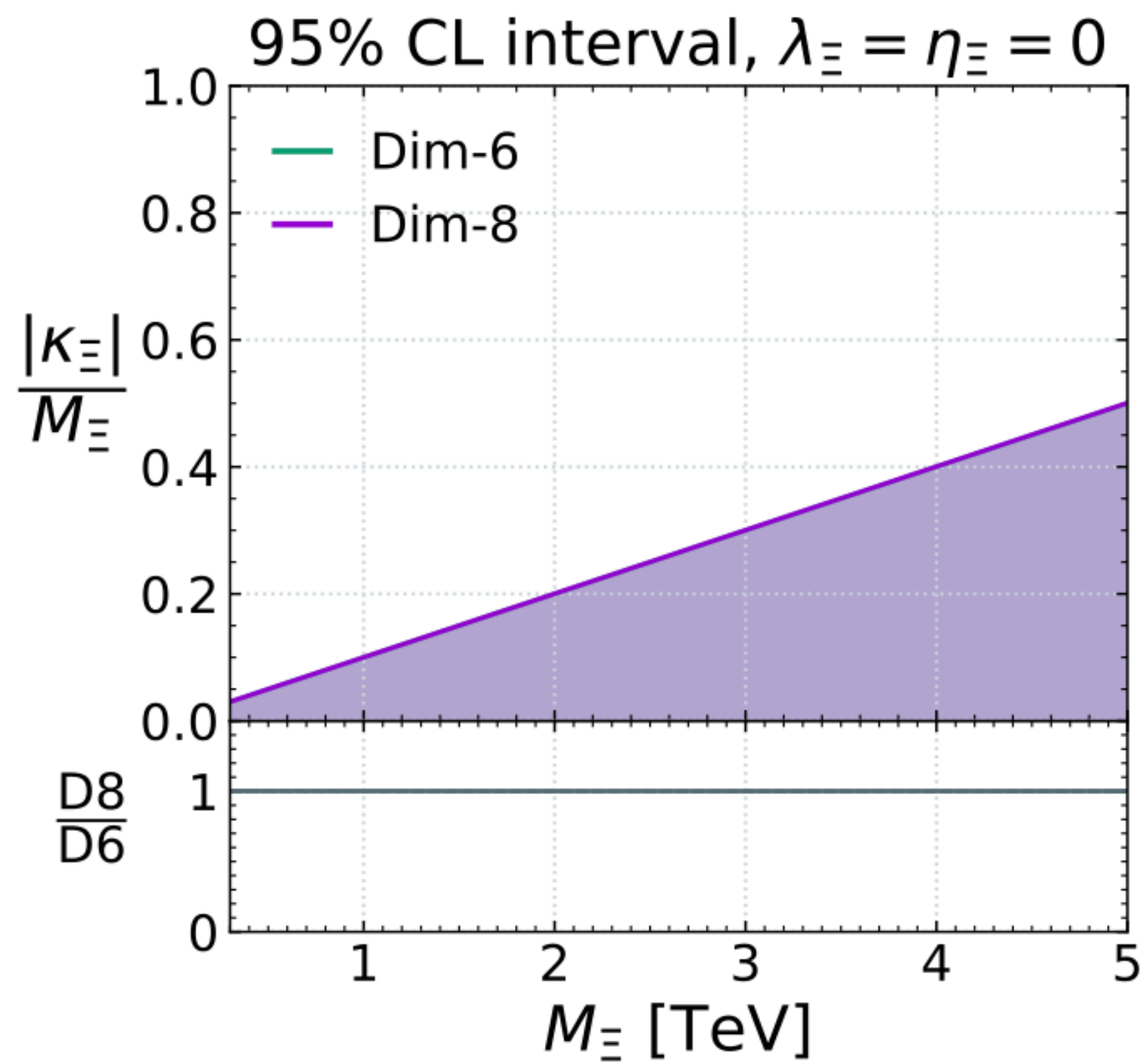
**Dim-6:**  $\mathcal{O}_{HD} \equiv (H^\dagger D^\mu H)^\star (H^\dagger D_\mu H)$       **Dim-8:**  $\mathcal{O}_{H^6}^{(2)} \equiv (H^\dagger H)(H^\dagger \sigma^I H)(D_\mu H^\dagger \sigma^I D^\mu H)$

$$\left. \frac{\delta m_W^2}{m_W^2} \right|_{D=6} = -\frac{\hat{c}_W}{\hat{c}_{2W}} \frac{C_{HD}}{2} \qquad \left. \frac{\delta m_W^2}{m_W^2} \right|_{D=8} = \frac{\hat{c}_W}{2\hat{c}_{2W}} \left( \frac{\hat{c}_W^2(\hat{c}_{2W} - \hat{s}_W^2)}{\hat{c}_{2W}^2} \frac{C_{HD}^2}{2} - C_{H^6}^{(2)} \right)$$

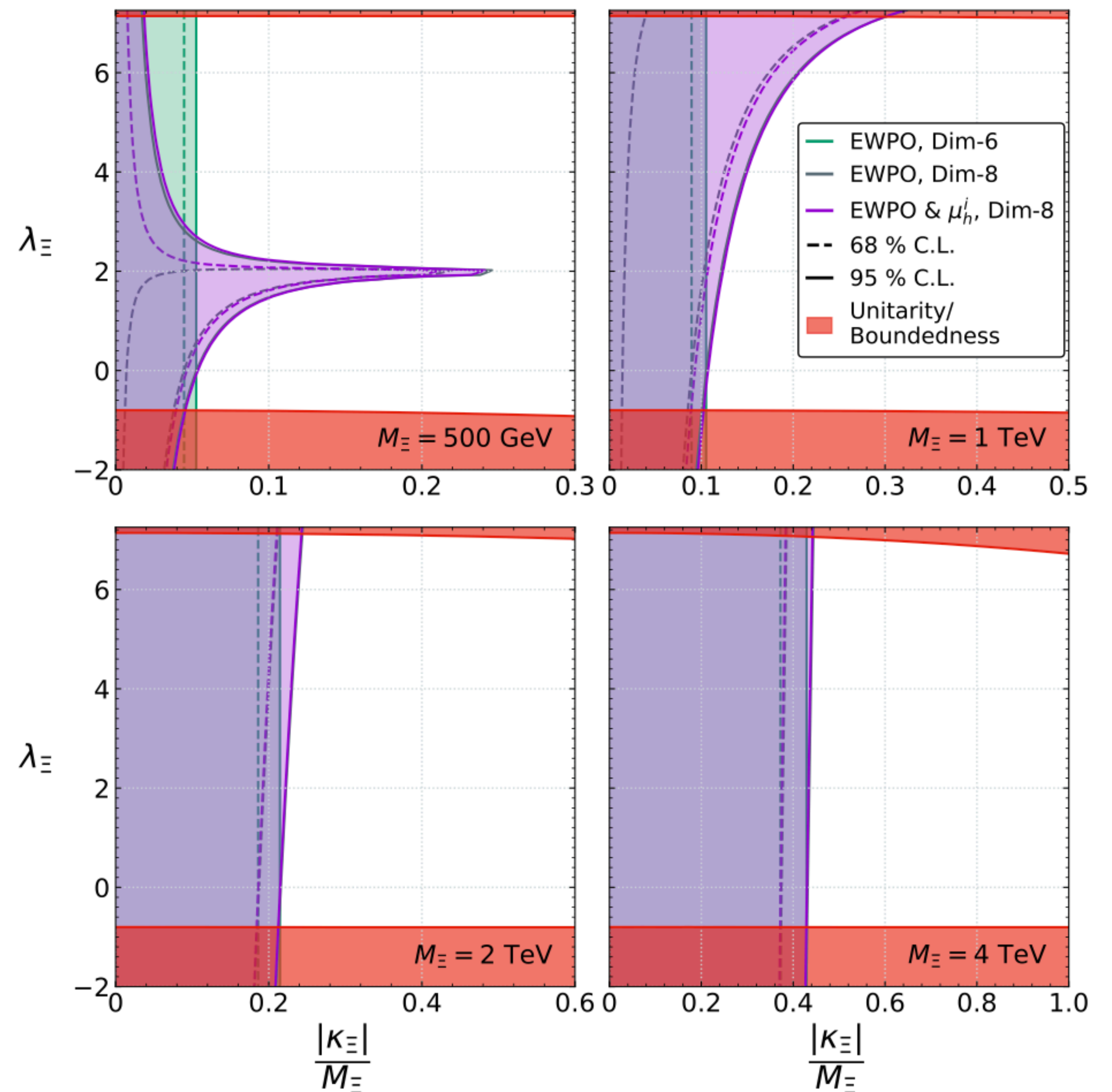
**Dim-6:**  $C_{HD} = -2 \frac{\kappa_\Xi^2}{M_\Xi^2}$       **Dim-8:**  $C_{HD} = -2 \frac{\kappa_\Xi^2}{M_\Xi^2} \left( 1 - \frac{4\mu^2}{M_\Xi^2} \right)$        $C_{H^6}^{(2)} = 4 \frac{\kappa_\Xi^2}{M_\Xi^2} \left( \lambda_\Xi - 2\lambda + \frac{\kappa_\Xi^2}{M_\Xi^2} \right)$

# Pre-CDF

Single parameter:  $\kappa_{\Xi}$

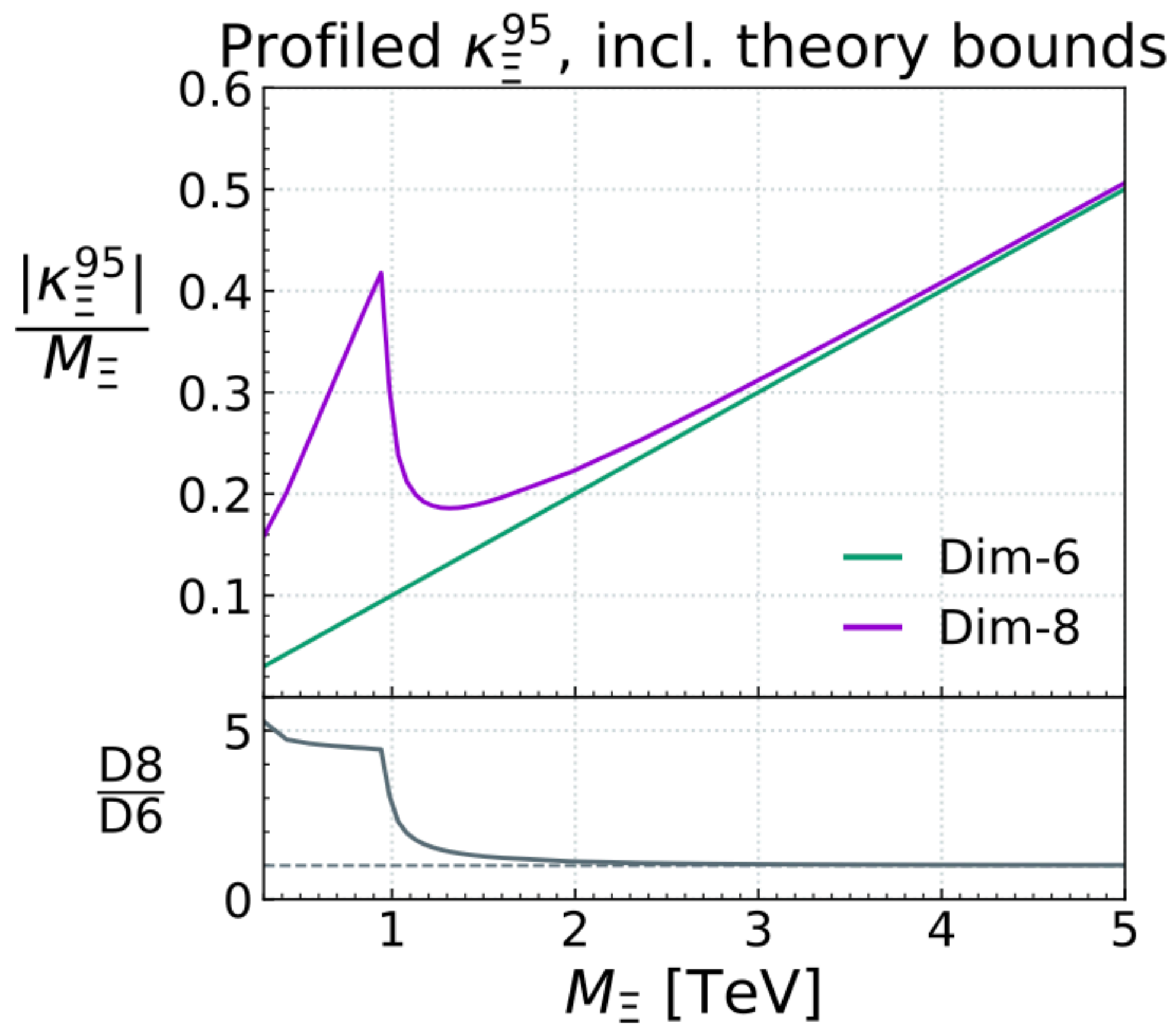


Two parameter:  $\kappa_{\Xi}, \lambda_{\Xi}$

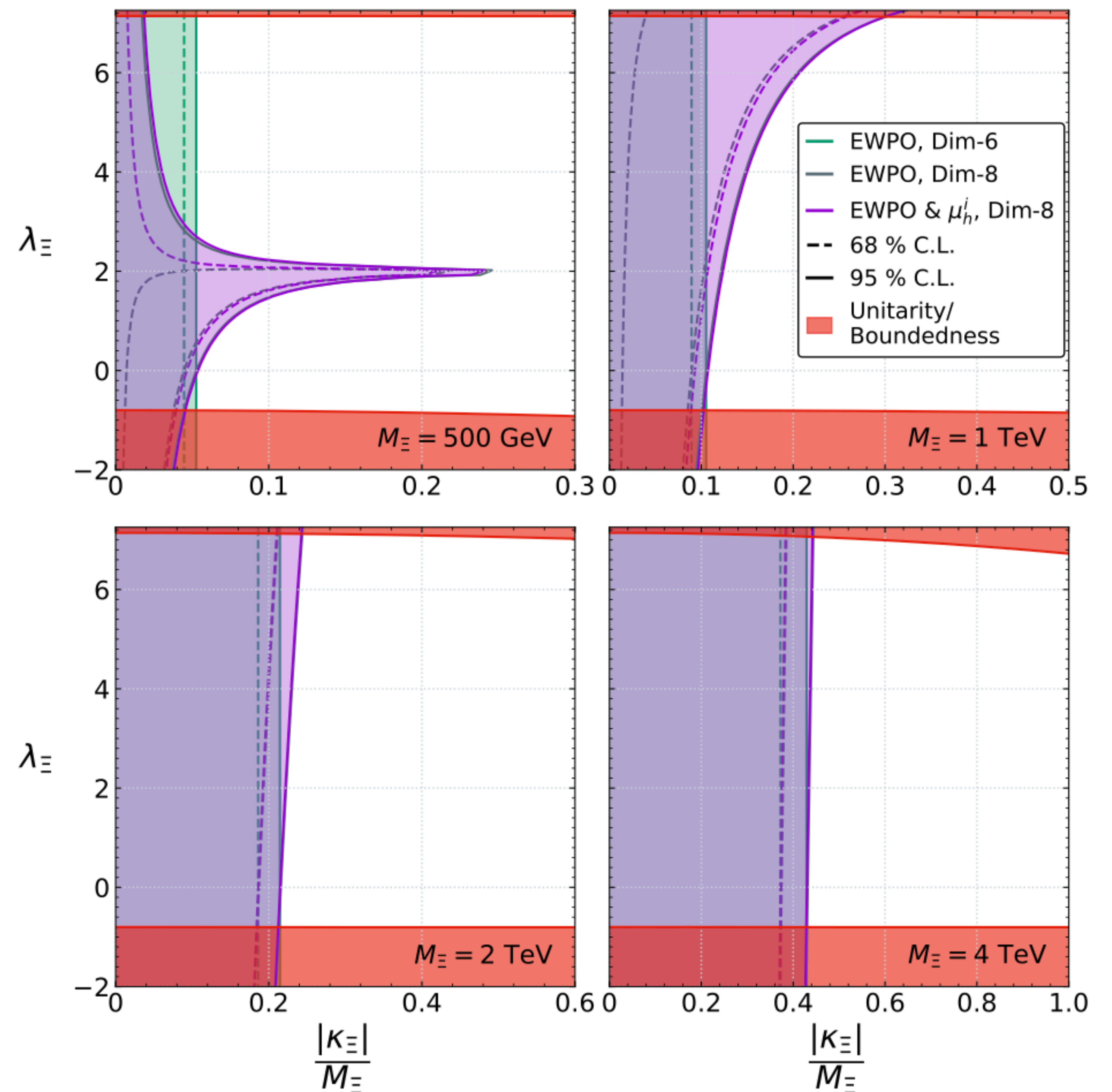


# Profiled

Single parameter:  $\kappa_{\Xi}$

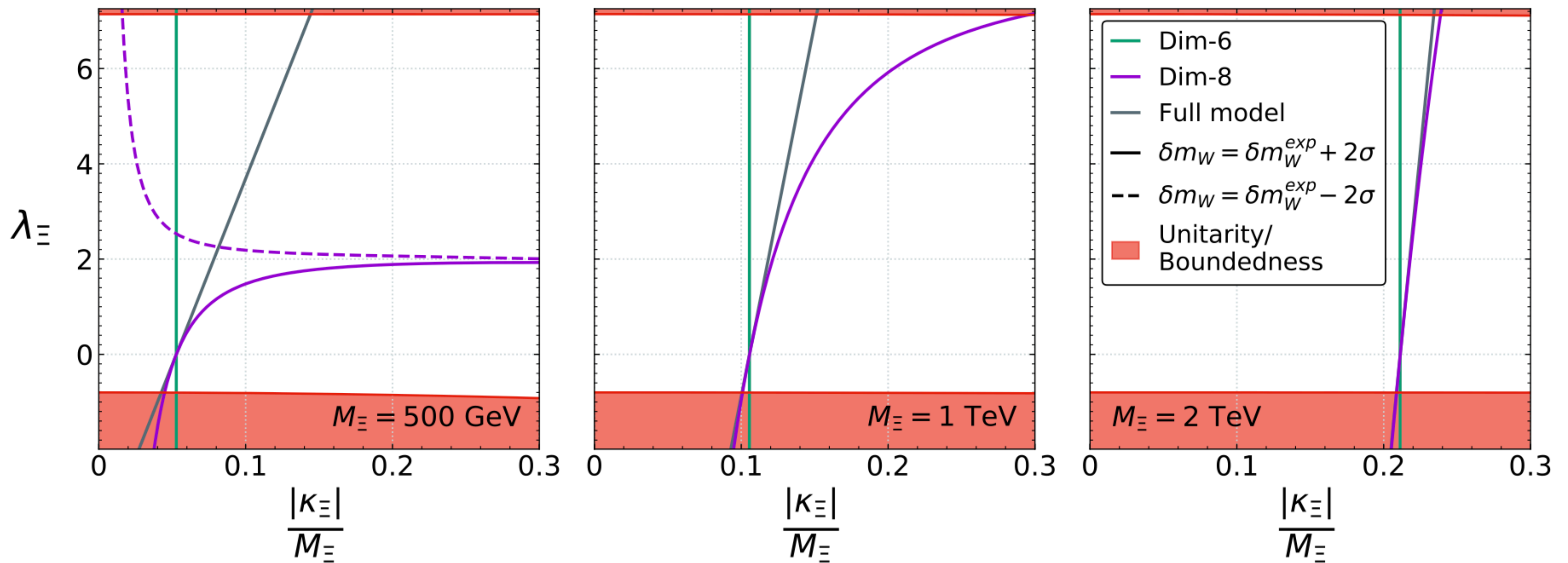


Two parameter:  $\kappa_{\Xi}, \lambda_{\Xi}$

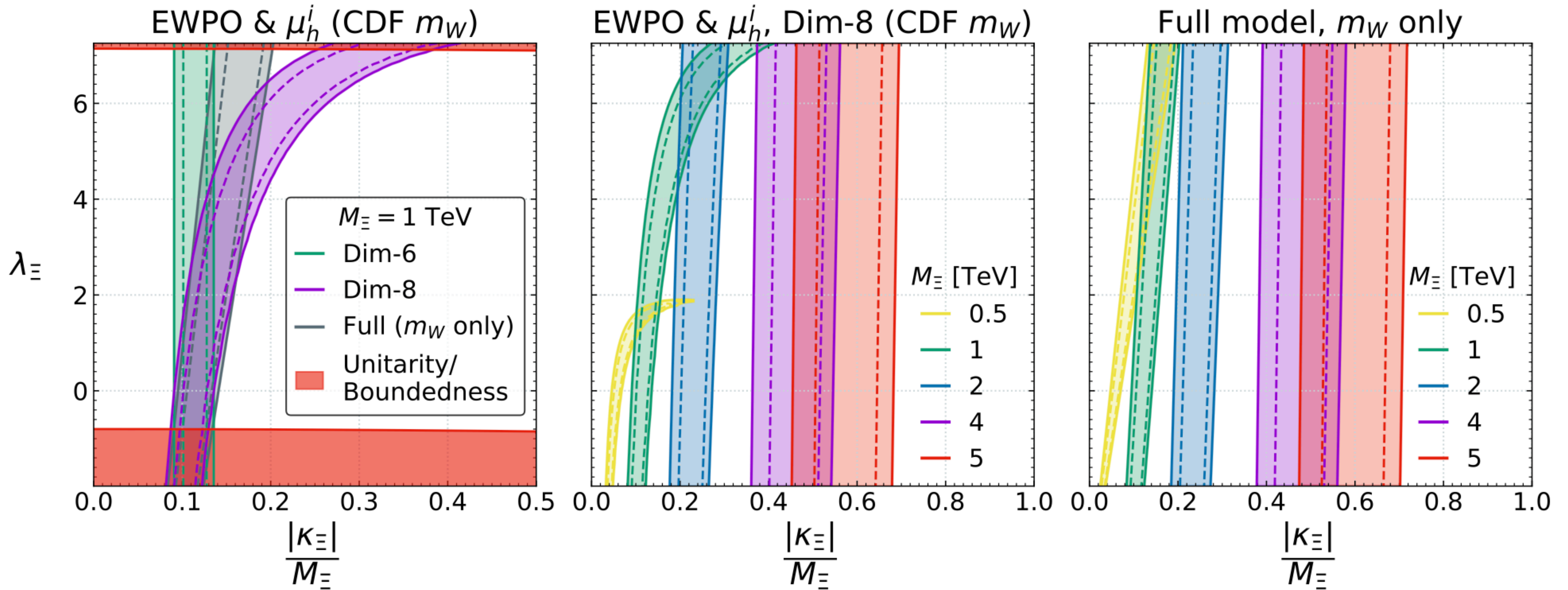


# EFT Validity

Relative  $W$ -mass shift,  $\delta m_W$ ;  $\eta_{\Xi} = 0$



# Post-CDF



# Conclusions

$m_W$  is a precise probe of a direction in SMEFT space

- Crucial input to global fits, lifts a quasi-flat direction

SMEFT can globally accommodate  $m_W^{\text{CDF}}$

- Higgs & Diboson data relatively less precise  $\Rightarrow$  Full 20-parameter fit remains consistent with SM
- SMEFT can help to pinpoint specific tree-level UV completions

- Beyond tree-level? *EWPO @ NLO* [Dawson & Giardino; PRD 101 (2020) 1, 013001]

*RGE effects in interpreting CDF  $m_W$*   
[Gupta; 2204.13690]

Interplay with CKM unitarity

Study of Dimension-8 effects in the scalar triplet model

- Dim-8 better reflects the constraints w.r.t full model
- Until the SMEFT expansion breaks down in strong coupling regions

# Backup

+



# SMEFT: SM v2.0

$$\mathcal{L} = \mathcal{L}_{\text{SM}} + \sum_{i,D} \frac{c_i^{(D)} \mathcal{O}_i^{(D)}}{\Lambda^{D-4}}$$

*BSM particle masses  $M$*



*Generic new physics scale  $\Lambda$*

*Taylor expansion of  $\mathcal{A}_{\text{BSM}}$*



*Tower of operators  $\mathcal{O}_i^{(D)}$*

$\mathcal{O}_i^{(D)} \supset$



*Low energy (SM) fields & symmetries*

*Model parameters  $\{g_{\text{BSM}}^i, M_k\}$*



*Wilson coefficients  $\frac{c_j^{(D)}}{\Lambda^{D-4}} (g_{\text{BSM}}^i, M_k)$*

*measure  $g_i$  : new physics model parameters*

“Matching”

*measure  $c_i$  : coupling strengths of new BSM interactions*



# SMEFT: SM v2.0

$$\mathcal{L}_{\text{eff}} = \sum_i \frac{c_i \mathcal{O}_i^D}{\Lambda^{D-4}}$$

## SM = low energy effective description

- New physics = tower of irrelevant ( $D > 4$ ) operators
- Respecting low energy field content & symmetries

$$\text{SU}(3)_c \times \text{SU}(2)_L \times \text{U}(1)_Y$$

$$\varphi = \begin{pmatrix} G^+ \\ v + h + iG^0 \end{pmatrix} : \mathbf{2}_{\frac{1}{2}}$$

aTGC

$$X^3 : \epsilon_{IJK} W_{\mu\nu}^I W^{J,\nu\rho} W_{\rho}^{K,\mu}$$

$$X^2 H^2 : (\varphi^\dagger \varphi)^2 G_{\mu\nu}^a G_a^{\mu\nu}$$

ggh(h)

$\lambda_h$

$$H^6 : (\varphi^\dagger \varphi)^3$$

$$H^4 D^2 : (\varphi^\dagger D^\mu \varphi)^* (\varphi^\dagger D^\mu \varphi)$$

$\delta M_Z$

$y_f$

$$\psi^2 H^3 : (\varphi^\dagger \varphi)^2 (\bar{q}_i u_j \tilde{\varphi})$$

$$\psi^2 XH : (\bar{q}_i \sigma^{\mu\nu} u_j \tilde{\varphi}) B_{\mu\nu}$$

'dipole'

ffV

$$\psi^2 H^2 D : (\varphi^\dagger \overleftrightarrow{D}_\mu \varphi) (\bar{q}_i \gamma^\mu q_j)$$

$$\psi^4 : (\bar{q}_i \gamma^\mu q_j) (\bar{q}_k \gamma_\mu q_l)$$

4F

## More than 'just' a parametrisation of ignorance

- Unlike anomalous couplings
- Renormalisable QFT (order-by-order)
- Finite energy range ( $\sim \Lambda$ )
- Well defined matching procedure

# SMEFT interpretation

Improving sensitivity means improving...

$$\Delta o_n = o_n^{\text{EXP}} - o_n^{\text{SM}} = \sum_i \frac{a_{n,i}^{(6)}(\mu) c_i^{(6)}(\mu)}{\Lambda^2} + \mathcal{O}\left(\frac{1}{\Lambda^3}\right)$$

## Global nature

As many observables as possible

Identify patterns & correlations in fits

Exploit energy-growth

## Sensitivity

*Experiment:*

Best measurements & understanding of uncertainties and correlations

*Theory:*

Best available predictions for observables (NLO, NNLO, N3LO,...)

## Interpretation

Relies on accurate knowledge of the size & correlation among  $a_i$

Determining  $c_i^{(6)}$  requires most precise available SMEFT predictions

# Theory

[Grzadkowski et al.; JHEP 10 (2010) 085]

$X^3$		$H^6$ and $H^4 D^2$		$\psi^2 H^3$		$(\bar{L}L)(\bar{L}L)$		$(\bar{R}R)(\bar{R}R)$		$(\bar{L}L)(\bar{R}R)$	
$\mathcal{O}_G$	$f^{ABC} G_{\mu\nu}^A G_{\nu\rho}^B G_{\rho\mu}^C$	$\mathcal{O}_H$	$(H^\dagger H)^3$	$\mathcal{O}_{eH}$	$(H^\dagger H)(\bar{l}_p e_r H)$	$\mathcal{O}_{ll}$	$(\bar{l}_p \gamma_\mu l_r)(\bar{l}_s \gamma^\mu l_t)$	$\mathcal{O}_{ee}$	$(\bar{e}_p \gamma_\mu e_r)(\bar{e}_s \gamma^\mu e_t)$	$\mathcal{O}_{le}$	$(\bar{l}_p \gamma_\mu l_r)(\bar{e}_s \gamma^\mu e_t)$
$\mathcal{O}_{\bar{G}}$	$f^{ABC} \tilde{G}_{\mu\nu}^A G_{\nu\rho}^B G_{\rho\mu}^C$	$\mathcal{O}_{H\Box}$	$(H^\dagger H)\Box(H^\dagger H)$	$\mathcal{O}_{uH}$	$(H^\dagger H)(\bar{q}_p u_r \tilde{H})$	$\mathcal{O}_{qq}^{(1)}$	$(\bar{q}_p \gamma_\mu q_r)(\bar{q}_s \gamma^\mu q_t)$	$\mathcal{O}_{uu}$	$(\bar{u}_p \gamma_\mu u_r)(\bar{u}_s \gamma^\mu u_t)$	$\mathcal{O}_{lu}$	$(\bar{l}_p \gamma_\mu l_r)(\bar{u}_s \gamma^\mu u_t)$
$\mathcal{O}_W$	$\epsilon^{IJK} W_{\mu\nu}^I W_{\nu\rho}^J W_{\rho\mu}^K$	$\mathcal{O}_{HD}$	$(H^\dagger D^\mu H)^*(H^\dagger D_\mu H)$	$\mathcal{O}_{dH}$	$(H^\dagger H)(\bar{q}_p d_r H)$	$\mathcal{O}_{qq}^{(3)}$	$(\bar{q}_p \gamma_\mu \tau^I q_r)(\bar{q}_s \gamma^\mu \tau^I q_t)$	$\mathcal{O}_{dd}$	$(\bar{d}_p \gamma_\mu d_r)(\bar{d}_s \gamma^\mu d_t)$	$\mathcal{O}_{ld}$	$(\bar{l}_p \gamma_\mu l_r)(\bar{d}_s \gamma^\mu d_t)$
$\mathcal{O}_{\bar{W}}$	$\epsilon^{IJK} \tilde{W}_{\mu\nu}^I W_{\nu\rho}^J W_{\rho\mu}^K$					$\mathcal{O}_{lq}^{(1)}$	$(\bar{l}_p \gamma_\mu l_r)(\bar{q}_s \gamma^\mu q_t)$	$\mathcal{O}_{eu}$	$(\bar{e}_p \gamma_\mu e_r)(\bar{u}_s \gamma^\mu u_t)$	$\mathcal{O}_{qe}$	$(\bar{q}_p \gamma_\mu q_r)(\bar{e}_s \gamma^\mu e_t)$
						$\mathcal{O}_{lq}^{(3)}$	$(\bar{l}_p \gamma_\mu \tau^I l_r)(\bar{q}_s \gamma^\mu \tau^I q_t)$	$\mathcal{O}_{ed}$	$(\bar{e}_p \gamma_\mu e_r)(\bar{d}_s \gamma^\mu d_t)$	$\mathcal{O}_{qu}^{(1)}$	$(\bar{q}_p \gamma_\mu q_r)(\bar{u}_s \gamma^\mu u_t)$
$X^2 H^2$		$\psi^2 XH$		$\psi^2 H^2 D$		$(\bar{L}R)(\bar{R}L)$ and $(\bar{L}R)(\bar{L}R)$		$B$ -violating			
$\mathcal{O}_{HG}$	$H^\dagger H G_{\mu\nu}^A G^{A\mu\nu}$	$\mathcal{O}_{eW}$	$(\bar{l}_p \sigma^{\mu\nu} e_r) \tau^I H W_{\mu\nu}^I$	$\mathcal{O}_{Hi}^{(1)}$	$(H^\dagger i \overleftrightarrow{D}_\mu H)(\bar{l}_p \gamma^\mu l_r)$	$\mathcal{O}_{ledq}$	$(\bar{l}_p^j e_r)(\bar{d}_s^k q_t^j)$	$\mathcal{O}_{duq}$	$\epsilon^{\alpha\beta\gamma} \epsilon_{jk} [(d_p^\alpha)^T C u_r^\beta] [(q_s^{\gamma j})^T C l_t^k]$		
$\mathcal{O}_{H\bar{G}}$	$H^\dagger H \tilde{G}_{\mu\nu}^A G^{A\mu\nu}$	$\mathcal{O}_{eB}$	$(\bar{l}_p \sigma^{\mu\nu} e_r) H B_{\mu\nu}$	$\mathcal{O}_{Hi}^{(3)}$	$(H^\dagger i \overleftrightarrow{D}_\mu^I H)(\bar{l}_p \tau^I \gamma^\mu l_r)$	$\mathcal{O}_{quqd}^{(1)}$	$(\bar{q}_p^j u_r) \epsilon_{jk} (\bar{q}_s^k d_t)$	$\mathcal{O}_{qqu}$	$\epsilon^{\alpha\beta\gamma} \epsilon_{jk} [(q_p^{\alpha j})^T C q_r^{\beta k}] [(u_s^\gamma)^T C e_t]$		
$\mathcal{O}_{HW}$	$H^\dagger H W_{\mu\nu}^I W^{I\mu\nu}$	$\mathcal{O}_{uG}$	$(\bar{q}_p \sigma^{\mu\nu} T^A u_r) \tilde{H} G_{\mu\nu}^A$	$\mathcal{O}_{He}$	$(H^\dagger i \overleftrightarrow{D}_\mu H)(\bar{e}_p \gamma^\mu e_r)$	$\mathcal{O}_{quqd}^{(8)}$	$(\bar{q}_p^j T^A u_r) \epsilon_{jk} (\bar{q}_s^k T^A d_t)$	$\mathcal{O}_{qqq}$	$\epsilon^{\alpha\beta\gamma} \epsilon_{jkn} \epsilon_{km} [(q_p^{\alpha j})^T C q_r^{\beta k}] [(q_s^{\gamma m})^T C l_t^n]$		
$\mathcal{O}_{H\bar{W}}$	$H^\dagger H \tilde{W}_{\mu\nu}^I W^{I\mu\nu}$	$\mathcal{O}_{uW}$	$(\bar{q}_p \sigma^{\mu\nu} u_r) \tau^I \tilde{H} W_{\mu\nu}^I$	$\mathcal{O}_{Hq}^{(1)}$	$(H^\dagger i \overleftrightarrow{D}_\mu H)(\bar{q}_p \gamma^\mu q_r)$	$\mathcal{O}_{lequ}^{(1)}$	$(\bar{l}_p^j e_r) \epsilon_{jk} (\bar{q}_s^k u_t)$	$\mathcal{O}_{duu}$	$\epsilon^{\alpha\beta\gamma} [(d_p^\alpha)^T C u_r^\beta] [(u_s^\gamma)^T C e_t]$		
$\mathcal{O}_{HB}$	$H^\dagger H B_{\mu\nu} B^{\mu\nu}$	$\mathcal{O}_{uB}$	$(\bar{q}_p \sigma^{\mu\nu} u_r) \tilde{H} B_{\mu\nu}$	$\mathcal{O}_{Hq}^{(3)}$	$(H^\dagger i \overleftrightarrow{D}_\mu^I H)(\bar{q}_p \tau^I \gamma^\mu q_r)$	$\mathcal{O}_{Hu}$	$(H^\dagger i \overleftrightarrow{D}_\mu H)(\bar{u}_p \gamma^\mu u_r)$				
$\mathcal{O}_{H\bar{B}}$	$H^\dagger H \tilde{B}_{\mu\nu} B^{\mu\nu}$	$\mathcal{O}_{dG}$	$(\bar{q}_p \sigma^{\mu\nu} T^A d_r) H G_{\mu\nu}^A$	$\mathcal{O}_{Hu}$	$(H^\dagger i \overleftrightarrow{D}_\mu H)(\bar{u}_p \gamma^\mu u_r)$	$\mathcal{O}_{Hd}$	$(H^\dagger i \overleftrightarrow{D}_\mu H)(\bar{d}_p \gamma^\mu d_r)$				
$\mathcal{O}_{HWB}$	$H^\dagger \tau^I H W_{\mu\nu}^I B^{\mu\nu}$	$\mathcal{O}_{dW}$	$(\bar{q}_p \sigma^{\mu\nu} d_r) \tau^I H W_{\mu\nu}^I$	$\mathcal{O}_{Hd}$	$(H^\dagger i \overleftrightarrow{D}_\mu H)(\bar{d}_p \gamma^\mu d_r)$						
$\mathcal{O}_{H\bar{W}B}$	$H^\dagger \tau^I H \tilde{W}_{\mu\nu}^I B^{\mu\nu}$	$\mathcal{O}_{dB}$	$(\bar{q}_p \sigma^{\mu\nu} d_r) H B_{\mu\nu}$	$\mathcal{O}_{Hud}$	$i(\tilde{H}^\dagger D_\mu H)(\bar{u}_p \gamma^\mu d_r)$						

## Warsaw basis with CP & B conservation

- Full ‘bosonic’ sector: Higgs, triple-gauge & gauge-Higgs
- Scenario 1: Flavor-**universal** degrees of freedom

$$U(3)_L \times U(3)_e \times U(3)_Q \times U(3)_u \times U(3)_d \quad + \text{Yukawas: } \mathcal{O}_{tH}, \mathcal{O}_{bH}, \mathcal{O}_{\tau H}, \mathcal{O}_{\mu H}$$

- Scenario 2: **top**-centric flavor symmetry

$$U(3)_L \times U(3)_e \times U(2)_Q \times U(2)_u \times U(3)_d \quad \begin{array}{l} \text{cf. Minimal flavor violation} \\ \text{[Buras et al.; PLB 500 (2001) 161]} \\ \text{[D'Ambrosio et al.; NPB 645 (2002) 155]} \\ \text{[Aguilar-Saavedra et al.; arXiv:1802.07237]} \end{array}$$

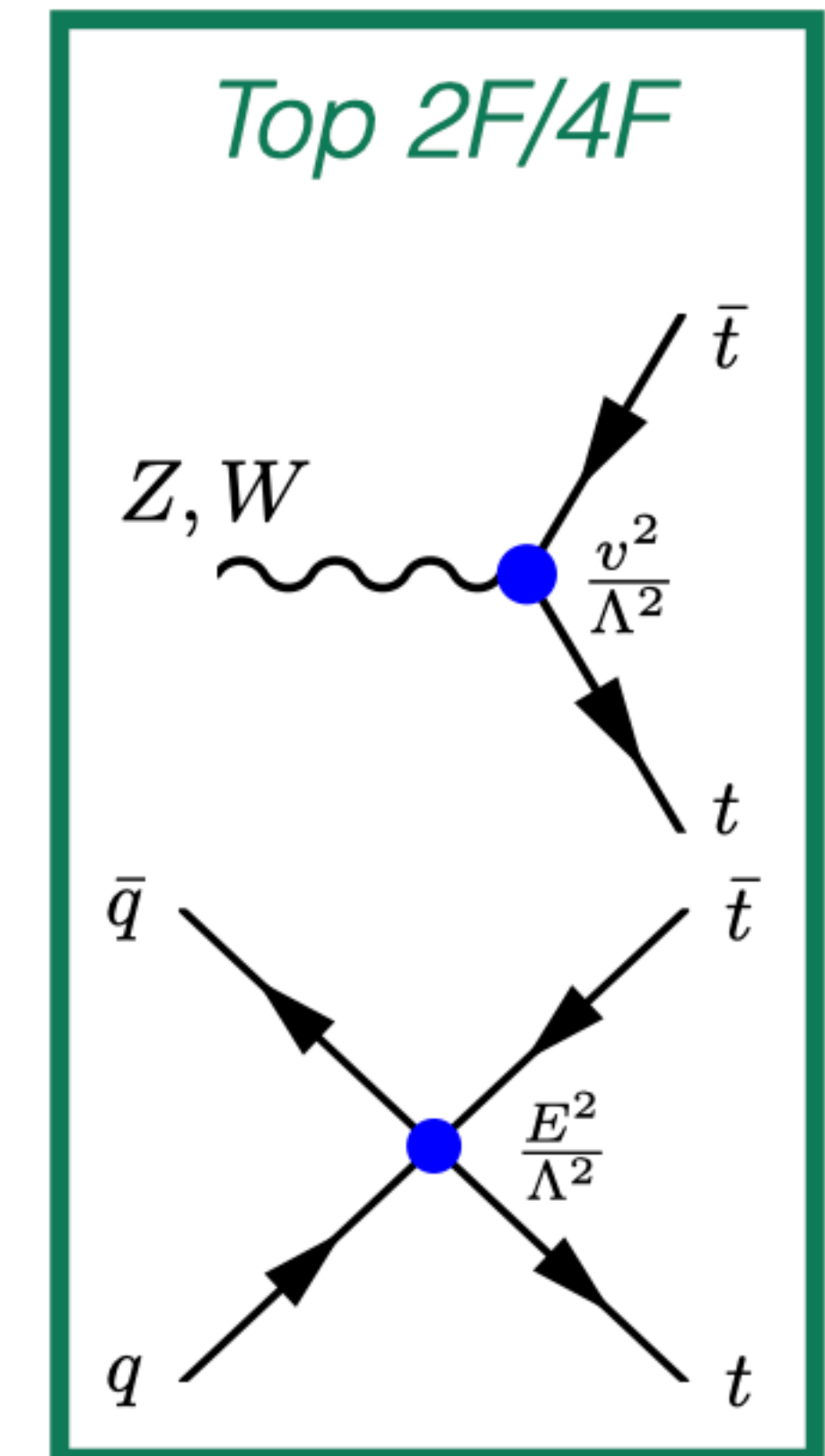
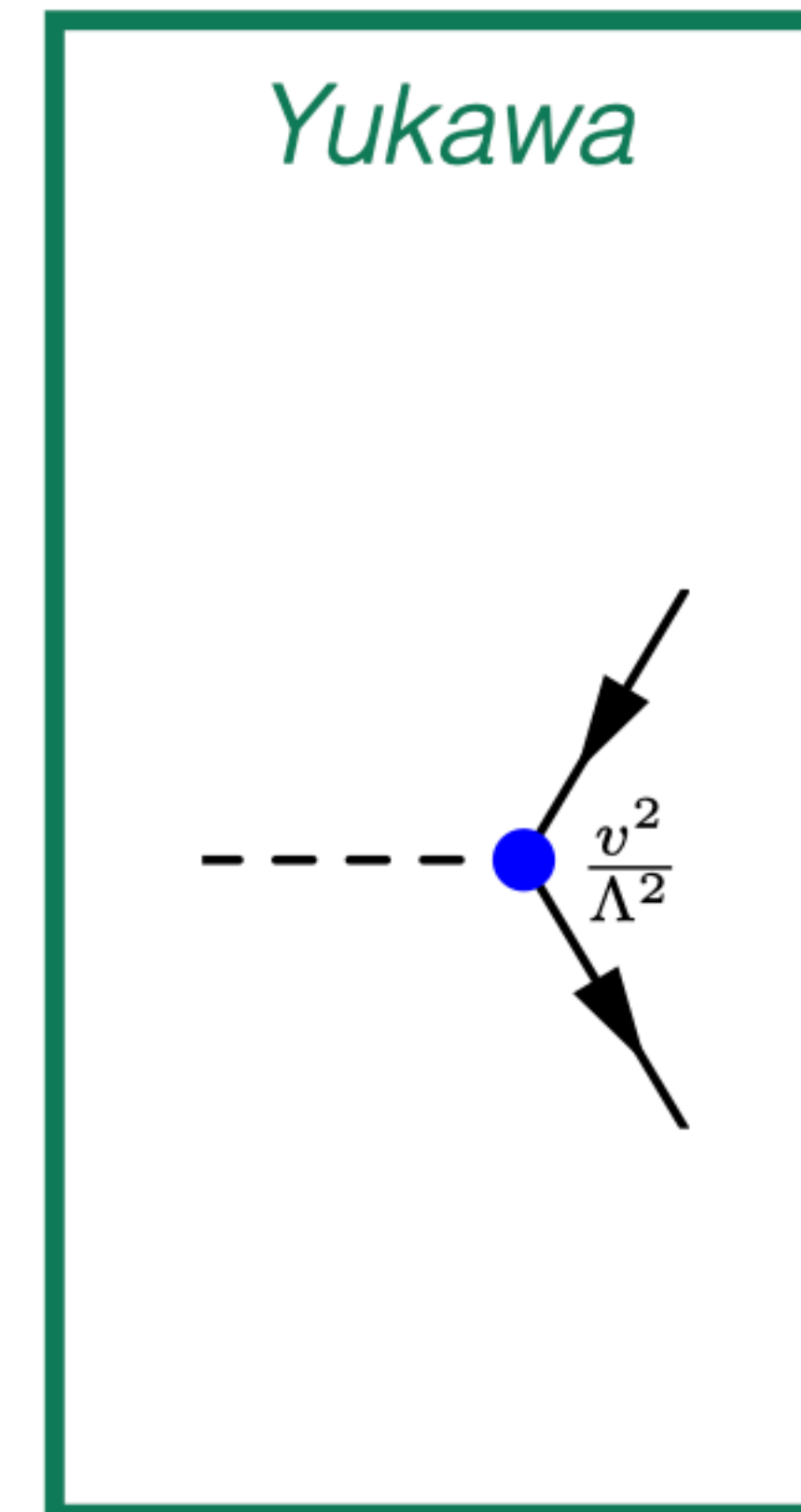
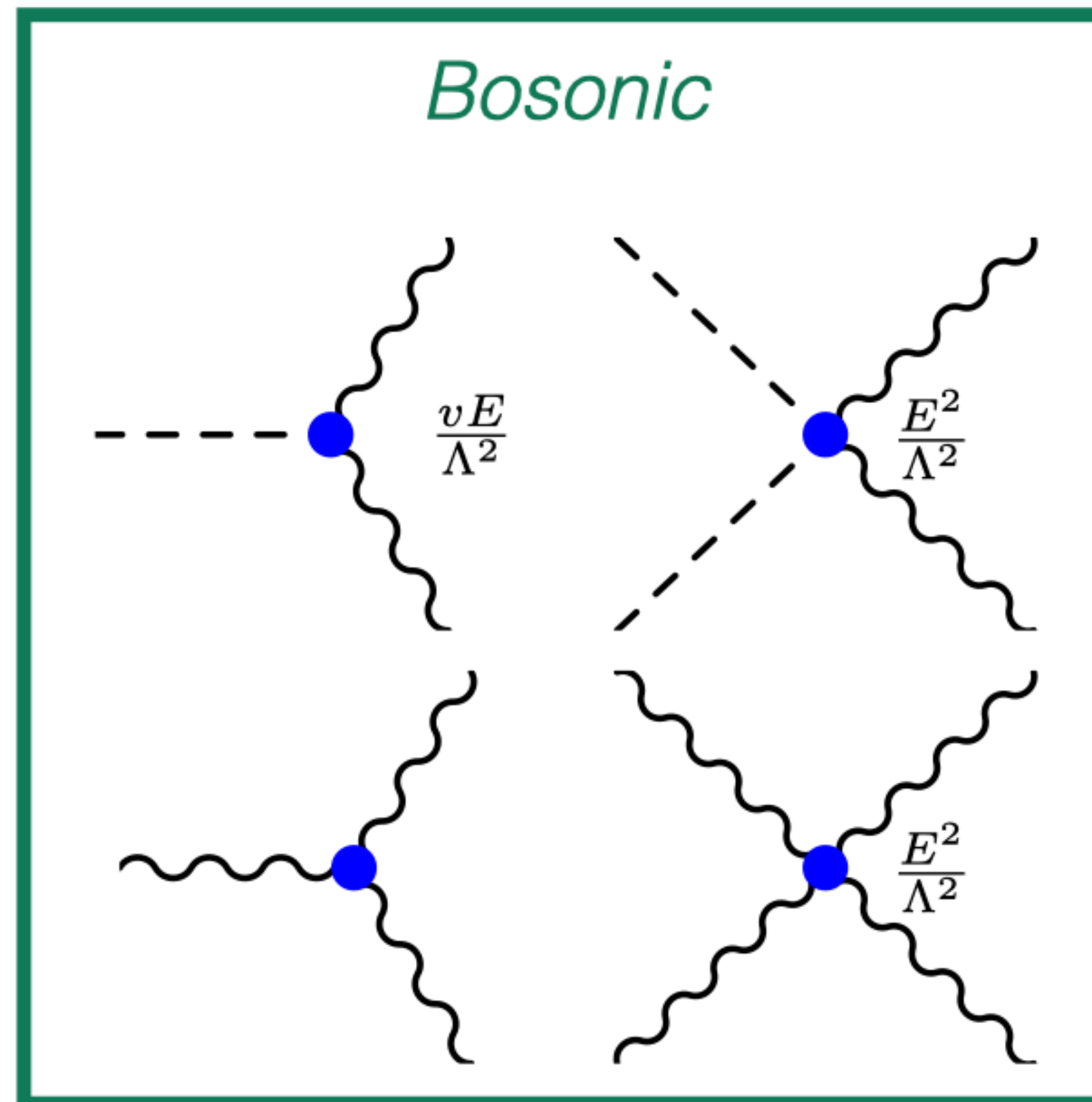
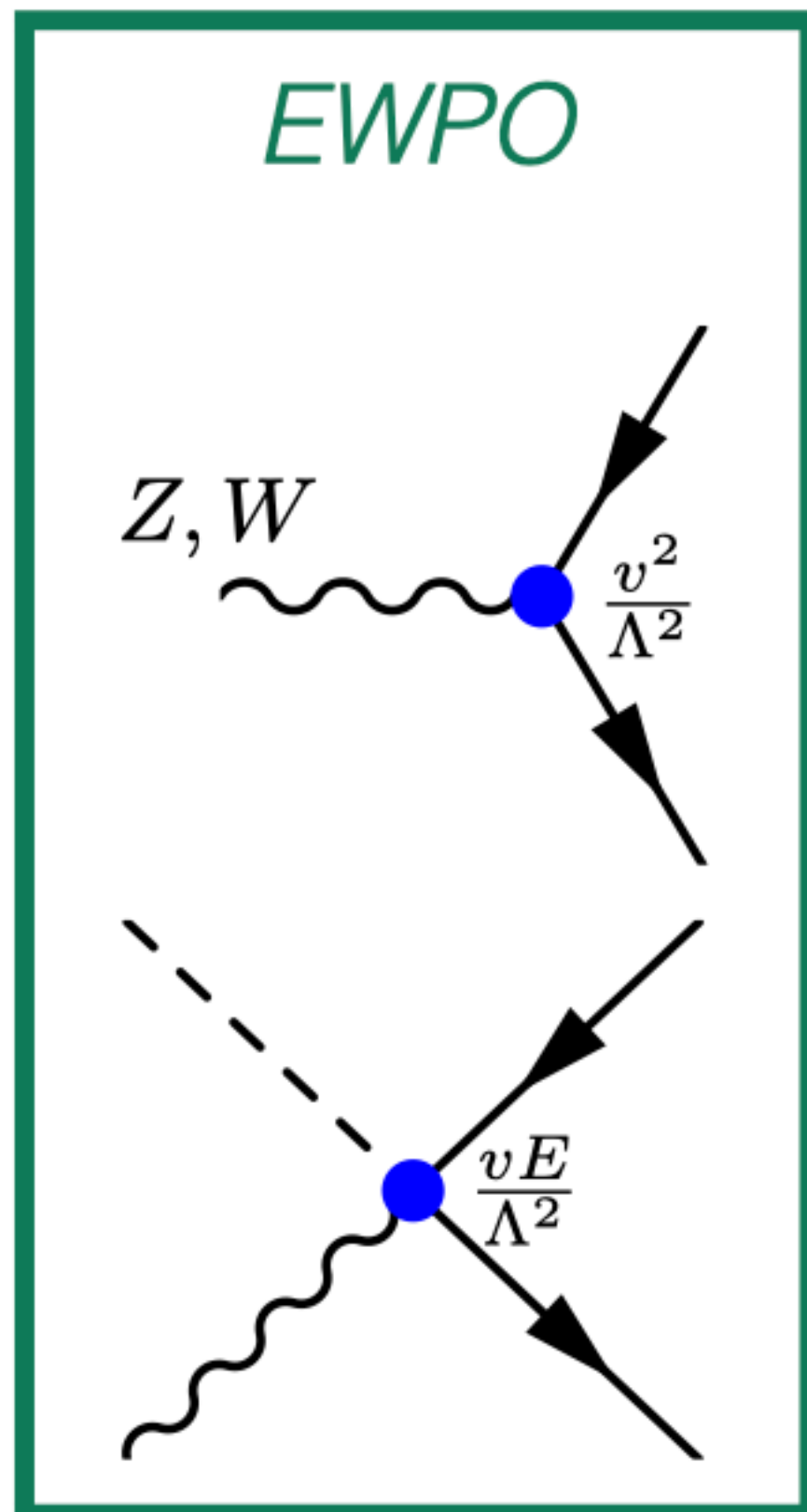
# Degrees of freedom

**Flavor scenario**

*Universal*

*'Top specific'*

EWPO:	$\mathcal{O}_{HWB}, \mathcal{O}_{HD}, \mathcal{O}_U, \mathcal{O}_{Hl}^{(3)}, \mathcal{O}_{Hl}^{(1)}, \mathcal{O}_{He}, \mathcal{O}_{Hq}^{(3)}, \mathcal{O}_{Hq}^{(1)}, \mathcal{O}_{Hd}, \mathcal{O}_{Hu},$	
Bosonic:	$\mathcal{O}_{H\Box}, \mathcal{O}_{HG}, \mathcal{O}_{HW}, \mathcal{O}_{HB}, \mathcal{O}_W, \mathcal{O}_G,$	
Yukawa:	$\mathcal{O}_{\tau H}, \mathcal{O}_{\mu H}, \mathcal{O}_{bH}, \mathcal{O}_{tH},$	20
Top 2F:	$\mathcal{O}_{HQ}^{(3)}, \mathcal{O}_{HQ}^{(1)}, \mathcal{O}_{Ht}, \mathcal{O}_{tG}, \mathcal{O}_{tW}, \mathcal{O}_{tB},$	
Top 4F:	$\mathcal{O}_{Qq}^{3,1}, \mathcal{O}_{Qq}^{3,8}, \mathcal{O}_{Qq}^{1,8}, \mathcal{O}_{Qu}^8, \mathcal{O}_{Qd}^8, \mathcal{O}_{tQ}^8, \mathcal{O}_{tu}^8, \mathcal{O}_{td}^8.$	+14



# SMEFT fit

SMEFT Coeff.	Individual			Marginalised		
	Best fit [ $\Lambda = 1$ TeV]	95% CL range	Scale $\frac{\Lambda}{\sqrt{C}}$ [TeV]	Best fit [ $\Lambda = 1$ TeV]	95% CL range	Scale $\frac{\Lambda}{\sqrt{C}}$ [TeV]
$C_{HWB}$	-0.01	[ -0.009, -0.0034 ]	19.0	0.25	[ -0.3, +0.81 ]	1.3
$C_{HD}$	-0.03	[ -0.035, -0.019 ]	11.0	-0.6	[ -1.8, +0.63 ]	0.9
$C_{ll}$	0.02	[ +0.014, +0.034 ]	10.0	-0.05	[ -0.099, +0.0043 ]	4.4
$C_{Hl}^{(3)}$	-0.01	[ -0.019, -0.0083 ]	14.0	-0.01	[ -0.11, +0.076 ]	3.3
$C_{Hl}^{(1)}$	0.00	[ -0.0045, +0.013 ]	11.0	0.16	[ -0.15, +0.47 ]	1.8
$C_{He}$	0.00	[ -0.015, +0.0071 ]	9.6	0.28	[ -0.34, +0.9 ]	1.3
$C_{Hq}^{(3)}$	0.00	[ -0.013, +0.011 ]	9.1	-0.05	[ -0.11, +0.012 ]	4.1
$C_{Hq}^{(1)}$	0.01	[ -0.027, +0.043 ]	5.4	-0.07	[ -0.2, +0.06 ]	2.8
$C_{Hd}$	-0.03	[ -0.13, +0.072 ]	3.1	-0.44	[ -0.96, +0.079 ]	1.4
$C_{Hu}$	0.00	[ -0.075, +0.073 ]	3.7	-0.18	[ -0.62, +0.26 ]	1.5
$C_{HBox}$	-0.27	[ -1, +0.47 ]	1.2	-1.1	[ -3.2, +1 ]	0.69
$C_{HG}$	0.00	[ -0.0034, +0.0032 ]	17.0	-0.01	[ -0.026, +0.013 ]	7.2
$C_{HW}$	0.00	[ -0.012, +0.006 ]	11.0	0.18	[ -0.33, +0.7 ]	1.4
$C_{HB}$	0.00	[ -0.0034, +0.002 ]	19.0	0.09	[ -0.074, +0.24 ]	2.5
$C_W$	0.18	[ -0.072, +0.42 ]	2.0	0.15	[ -0.1, +0.4 ]	2.0
$C_G$	-0.75	[ -4, +2.5 ]	0.56	1.3	[ -6.1, +8.7 ]	0.37
$C_{\tau H}$	0.01	[ -0.015, +0.025 ]	7.1	0.00	[ -0.017, +0.027 ]	6.7
$C_{\mu H}$	0.00	[ -0.0057, +0.005 ]	14.0	0.00	[ -0.0056, +0.0052 ]	14.0
$C_{bH}$	0.00	[ -0.016, +0.024 ]	7.1	0.02	[ -0.027, +0.058 ]	4.8
$C_{tH}$	-0.09	[ -1, +0.84 ]	1.0	-2.7	[ -8.8, +3.3 ]	0.41

# Operators

Dim - 4	$\mathcal{O}_{H4}$	$(H^\dagger H)^2$
Dim - 6	$H^6$ and $H^4 D^2$	
	$\mathcal{O}_H$	$(H^\dagger H)^3$
	$\mathcal{O}_{HD}$	$(H^\dagger D^\mu H)^*(H^\dagger D^\mu H)$
	$\mathcal{O}_{H\Box}$	$(H^\dagger H)\Box(H^\dagger H)$
	$\psi^2 H^3$	
	$\mathcal{O}_{eH}$	$(H^\dagger H)(\bar{l}_p e_r H)$
$\mathcal{O}_{uH}$	$(H^\dagger H)(\bar{q}_p u_r \tilde{H})$	
$\mathcal{O}_{dH}$	$(H^\dagger H)(\bar{q}_p d_r H)$	

Dim - 8	$H^8, H^6 D^2$ and $H^4 D^4$		$(\bar{L}R)(\bar{L}R)H^2 + \text{h.c.}$	
	$\mathcal{O}_{H^8}$	$(H^\dagger H)^4$	$\mathcal{O}_{lequH^2}^{(1)}$	$(\bar{l}_p^j e_r)\epsilon_{jk}(\bar{q}_s^k u_t)(H^\dagger H)$
	$\mathcal{O}_{H^6}^{(1)}$	$(H^\dagger H)^2(D_\mu H^\dagger D^\mu H)$	$\mathcal{O}_{lequH^2}^{(2)}$	$(\bar{l}_p^j e_r)(\sigma^I \epsilon)_{jk}(\bar{q}_s^k u_t)(H^\dagger \sigma^I H)$
	$\mathcal{O}_{H^6}^{(2)}$	$(H^\dagger H)(H^\dagger \sigma^I H)(D_\mu H^\dagger \sigma^I D^\mu H)$	$\mathcal{O}_{q^2 udH^2}^{(1)}$	$(\bar{q}_p^j u_r)\epsilon_{jk}(\bar{q}_s^k d_t)(H^\dagger H)$
	$\mathcal{O}_{H^4}^{(1)}$	$(D_\mu H^\dagger D_\nu H)(D^\nu H^\dagger D^\mu H)$	$\mathcal{O}_{q^2 udH^2}^{(2)}$	$(\bar{q}_p^j u_r)(\sigma^I \epsilon)_{jk}(\bar{q}_s^k d_t)(H^\dagger \sigma^I H)$
	$\mathcal{O}_{H^4}^{(3)}$	$(D_\mu H^\dagger D^\mu H)(D_\nu H^\dagger D^\nu H)$	$\mathcal{O}_{leqdH^2}^{(3)}$	$(\bar{l}_p e_r H)(\bar{q}_s d_t H)$
	$\psi^2 H^5$		$\mathcal{O}_{l^2 e^2 H^2}^{(3)}$	$(\bar{l}_p e_r H)(\bar{l}_s e_t H)$
	$\mathcal{O}_{leH^5}$	$(H^\dagger H)^2(\bar{l}_p e_r H)$	$\mathcal{O}_{q^2 u^2 H^2}^{(5)}$	$(\bar{q}_p u_r \tilde{H})(\bar{q}_s u_t \tilde{H})$
	$\mathcal{O}_{quH^5}$	$(H^\dagger H)^2(\bar{q}_p u_r \tilde{H})$	$\mathcal{O}_{q^2 d^2 H^2}^{(5)}$	$(\bar{q}_p d_r H)(\bar{q}_s u_t H)$
	$\mathcal{O}_{qdH^5}$	$(H^\dagger H)^2(\bar{q}_p d_r H)$	$(\bar{L}L)(\bar{R}R)H^2$	
	$\psi^2 H^3 D^2 + \text{h.c.}$		$\mathcal{O}_{l^2 e^2 H^2}^{(1)}$	$(\bar{l}_p \gamma^\mu l_r)(\bar{e}_s \gamma_\mu e_t)(H^\dagger H)$
	$\mathcal{O}_{leH^3 D^2}^{(1)}$	$(D_\mu H^\dagger D^\mu H)(\bar{l}_p e_r H)$	$\mathcal{O}_{l^2 e^2 H^2}^{(2)}$	$(\bar{l}_p \gamma^\mu \sigma^I l_r)(\bar{e}_s \gamma_\mu e_t)(H^\dagger \sigma^I H)$
	$\mathcal{O}_{leH^3 D^2}^{(2)}$	$(D_\mu H^\dagger \sigma^I D^\mu H)(\bar{l}_p e_r \sigma^I H)$	$\mathcal{O}_{q^2 u^2 H^2}^{(1)}$	$(\bar{q}_p \gamma^\mu q_r)(\bar{u}_s \gamma_\mu u_t)(H^\dagger H)$
	$\mathcal{O}_{quH^3 D^2}^{(1)}$	$(D_\mu H^\dagger D^\mu H)(\bar{q}_p u_r H)$	$\mathcal{O}_{q^2 u^2 H^2}^{(2)}$	$(\bar{q}_p \gamma^\mu \sigma^I q_r)(\bar{u}_s \gamma_\mu u_t)(H^\dagger \sigma^I H)$
$\mathcal{O}_{quH^3 D^2}^{(2)}$	$(D_\mu H^\dagger \sigma^I D^\mu H)(\bar{q}_p u_r \sigma^I H)$	$\mathcal{O}_{q^2 d^2 H^2}^{(1)}$	$(\bar{q}_p \gamma^\mu q_r)(\bar{d}_s \gamma_\mu d_t)(H^\dagger H)$	
$\mathcal{O}_{qdH^3 D^2}^{(1)}$	$(D_\mu H^\dagger D^\mu H)(\bar{q}_p d_r H)$	$\mathcal{O}_{q^2 d^2 H^2}^{(2)}$	$(\bar{q}_p \gamma^\mu \sigma^I q_r)(\bar{d}_s \gamma_\mu d_t)(H^\dagger \sigma^I H)$	
$\mathcal{O}_{qdH^3 D^2}^{(2)}$	$(D_\mu H^\dagger \sigma^I D^\mu H)(\bar{q}_p d_r \sigma^I H)$	$(\bar{L}R)(\bar{R}L)H^2 + \text{h.c.}$		
$\mathcal{O}_{leqdH^2}^{(1)}$	$(\bar{l}_p^j e_r)(\bar{d}_s q_{tj})(H^\dagger H)$	$\mathcal{O}_{lequH^2}^{(5)}$	$(\bar{l}_p e_r H)(\tilde{H}^\dagger \bar{u}_s q_t)$	
$\mathcal{O}_{leqdH^2}^{(2)}$	$(\bar{l}_p e_r)\sigma^I(\bar{d}_s q_t)(H^\dagger \sigma^I H)$	$\mathcal{O}_{q^2 u^2 H^2}^{(5)}$	$(\bar{q}_p d_r H)(\tilde{H}^\dagger \bar{u}_s q_t)$	

# Triplet matching

Dim - 6	$C_H$	$\frac{\kappa_{\Xi}^2}{M_{\Xi}^2} \left( (4\lambda - \lambda_{\Xi}) \left( 1 - \frac{4\mu^2}{M_{\Xi}^2} \right) - \frac{5\mu^2 \kappa_{\Xi}^2}{M_{\Xi}^6} \right)$
	$C_{HD}$	$-\frac{2\kappa_{\Xi}^2}{M_{\Xi}^2} \left( 1 - \frac{4\mu^2}{M_{\Xi}^2} \right)$
	$C_{H\Box}$	$\frac{\kappa_{\Xi}^2}{2M_{\Xi}^2} \left( 1 - \frac{4\mu^2}{M_{\Xi}^2} \right)$
	$[C_{\psi H}]_{wx}$	$[y_{\psi}]_{wx} \frac{\kappa_{\Xi}^2}{M_{\Xi}^2} \left( 1 - \frac{4\mu^2}{M_{\Xi}^2} \right); \psi = u, d, e$

Dim - 8	$C_{H^8}$	$\frac{2\kappa_{\Xi}^2}{M_{\Xi}^2} \left( (2\lambda - \lambda_{\Xi})^2 + \frac{\kappa_{\Xi}^2}{M_{\Xi}^2} (3\lambda_{\Xi} - 5\lambda - \frac{\eta_{\Xi}}{8}) \right)$
	$C_{H^6}^{(1)}$	$-\frac{\kappa_{\Xi}^4}{M_{\Xi}^4}$
	$C_{H^6}^{(2)}$	$\frac{4\kappa_{\Xi}^2}{M_{\Xi}^2} \left( \lambda_{\Xi} - 2\lambda + \frac{\kappa_{\Xi}^2}{M_{\Xi}^2} \right)$
	$C_{H^4}^{(1)}$	$\frac{4\kappa_{\Xi}^2}{M_{\Xi}^2}$
	$C_{H^4}^{(3)}$	$-\frac{2\kappa_{\Xi}^2}{M_{\Xi}^2}$
	$[C_{l\psi H^5}/C_{q\psi H^5}]_{wx}$	$-[y_{\psi}]_{wx} \frac{2\kappa_{\Xi}^2}{M_{\Xi}^2} \left( \lambda_{\Xi} - 2\lambda + \frac{\kappa_{\Xi}^2}{2M_{\Xi}^2} \right); \psi = u, d, e$
	$[C_{l^2\psi^2 H^2}/C_{q^2\psi^2 H^2}]_{wxyz}$	$-[y_{\psi}]_{wz} [y_{\psi}^{\dagger}]_{yx} \frac{3\kappa_{\Xi}^2}{4M_{\Xi}^2}; \psi = u, d, e$
	$[C_{l^2 e^2 H^2}/C_{q^2 d^2 H^2}]_{wxyz}$	$[y_{\psi}]_{wz} [y_{\psi}^{\dagger}]_{yx} \frac{\kappa_{\Xi}^2}{4M_{\Xi}^2}; \psi = d, e$
	$[C_{q^2 u^2 H^2}]_{wxyz}$	$-[y_u]_{wz} [y_u^{\dagger}]_{yx} \frac{\kappa_{\Xi}^2}{4M_{\Xi}^2}$
	$[C_{l^2\psi^2 H^2}/C_{q^2\psi^2 H^2}]_{wxyz}^{(3)}$	$[y_{\psi}]_{wx} [y_{\psi}]_{yz} \frac{\kappa_{\Xi}^2}{2M_{\Xi}^2}; \psi = u, d, e$
	$[C_{lequH^2}^{(1)}]_{wxyz}$	$[y_e]_{wx} [y_u]_{yz} \frac{5\kappa_{\Xi}^2}{2M_{\Xi}^2}$
	$[C_{leqdH^2}^{(1)}]_{wxyz}$	$[y_e]_{wx} [y_d^{\dagger}]_{yz} \frac{5\kappa_{\Xi}^2}{2M_{\Xi}^2}$
	$[C_{q^2 udH^2}^{(1)}]_{wxyz}$	$-[y_u]_{wx} [y_d]_{yz} \frac{5\kappa_{\Xi}^2}{2M_{\Xi}^2}$
	$[C_{lequH^2}^{(2)}]_{wxyz}$	$[y_e]_{wx} [y_u]_{yz} \frac{\kappa_{\Xi}^2}{2M_{\Xi}^2}$
	$[C_{leqdH^2}^{(2)}]_{wxyz}$	$-[y_e]_{wx} [y_d^{\dagger}]_{yz} \frac{\kappa_{\Xi}^2}{2M_{\Xi}^2}$
	$[C_{q^2 udH^2}^{(2)}]_{wxyz}$	$[y_u]_{wx} [y_d]_{yz} \frac{\kappa_{\Xi}^2}{2M_{\Xi}^2}$
	$[C_{lequH^2}^{(5)}]_{wxyz}$	$[y_e]_{wx} [y_u^{\dagger}]_{yz} \frac{\kappa_{\Xi}^2}{M_{\Xi}^2}$
	$[C_{leqdH^2}^{(3)}]_{wxyz}$	$[y_e]_{wx} [y_d]_{yz} \frac{\kappa_{\Xi}^2}{M_{\Xi}^2}$
	$[C_{q^2 udH^2}^{(5)}]_{wxyz}$	$[y_d]_{wx} [y_u^{\dagger}]_{yz} \frac{\kappa_{\Xi}^2}{M_{\Xi}^2}$
	$[C_{l\psi H^3 D^2}/C_{q\psi H^3 D^2}]_{wx}$	$-[y_{\psi}]_{wx} \frac{4\kappa_{\Xi}^2}{M_{\Xi}^2}; \psi = u, d, e$