Investigating LHC Electron Cloud Instabilities through Linearized Vlasov Method

Sofia Johannesson, Giovanni Iadarola, Prof. Mike Seidel, Tatiana Pieloni.



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Introduction

E-cloud in a Vlasov formalism

Solutions to the Vlasov equation

First benchmark with measurements



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Outline

Introduction

E-cloud in a Vlasov formalism

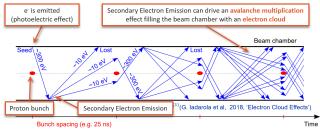
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Electron clouds

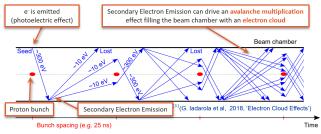


^[1] G. ladarola et al, 2018, Electron Cloud Effects



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Electron clouds



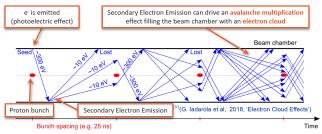
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Depends on:

- Beam Chamber
- Beam Configuration
- Magnetic fields



Electron clouds



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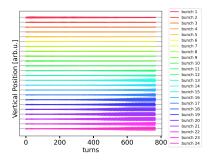
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Unwanted effects:

- Transverse instabilities
- Transverse emittance blow-up
- Particle losses
- Heat Loads
- Vacuum Degradation



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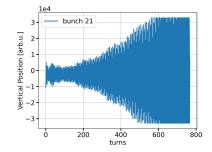


• Electron clouds can drive transverse instabilities



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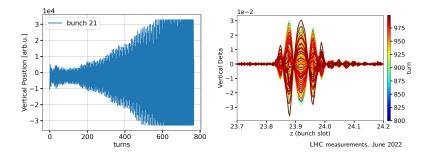


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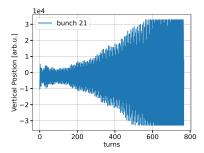




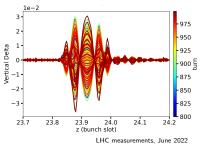
 Electron clouds can drive transverse instabilities, which cannot be mitigated by transverse feedback system due to strong intrabunch motion. [6] F. Zimmermann, 2004, Review of Single bunch instabilities driven by electron cloud



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 Conventional simulations using macroparticle tracking together with the PIC method for e-cloud beam interaction, are very computationally heavy. [7] G. ladarola. et al., 2017, Evolution of Python Tools for the Simulation of Electron Cloud Effects



 Instabilities driven by impedance effects have been modelled using the linearised Vlasov Equation, which identifies the Instability growth rate and betatron frequency shift for each instability mode. [8] N. Mounet, 2017, Vlasov Solvers and Macroparticle Simulations





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 - \Rightarrow Possible to study slow instabilities!



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Previous attempts of using the Vlasov method to model e-cloud driven instabilities have not included these points together. [10] K. Ohmi et al, 2001, Wake-Field and Fast Head-Tail Instability Caused by an Electron Cloud, [11] E. Perevedentsev, 2002, Head-Tail Instability Caused by Electron Cloud





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• ψ_0 is a distribution of particles where each individual particle obeys a Hamiltonian H_0 .

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- Introduce a perturbation, ΔH and $\Delta \psi$, which means that the total Hamiltonian is $H = H_0 + \Delta H$ and the total distribution is $\psi = \psi_0 + \Delta \psi$

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- This leads to the Linearized Vlasov Equation, which truncated to first order and expressed with Poisson brackets is:

$$\frac{\partial \Delta \psi}{\partial t} + [\Delta \psi, H_0] = -[\psi_0, \Delta H]$$
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• The electron cloud forces are contained in ΔH

• The distortion $\Delta \psi$ is the impact of the perturbation and the unknown [3] N. Mounet, 2018, Direct Vlasov Solvers

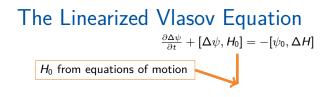


The Linearized Vlasov Equation $\frac{\partial \Delta \psi}{\partial t} + [\Delta \psi, H_0] = -[\psi_0, \Delta H]$



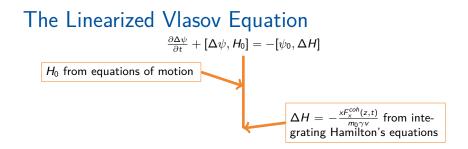
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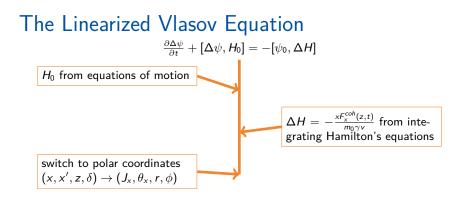






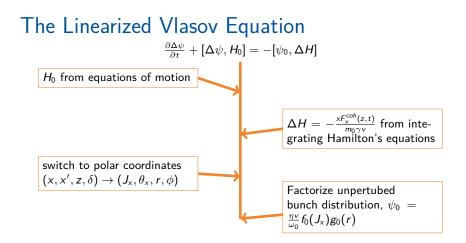






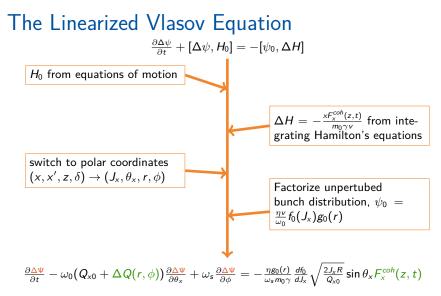












reds are unknowns and greens come from e-cloud forces



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Begin by describing the dipolar e-cloud forces:

[2] G. ladarola, et. al. 2020, Linearized method for the study of transverse instabilities driven by electron clouds



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Begin by describing the dipolar e-cloud forces:

Choose a set of sinusoid beam distortions, $h_n(z)$ where z is the position along the bunch. The sinusoid test functions satisfy the orthogonality condition: $\int h_n(z)h_{n'}(z) = H_n^2 \delta_{n,n'}$

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Each distortion, h_n , corresponds to a response function k_n calculated from the interaction with e-cloud using single-pass PIC simulations.

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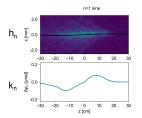
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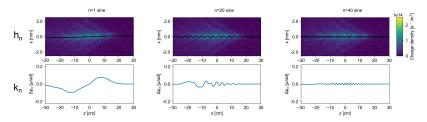
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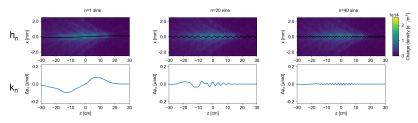
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Each distortion, h_n , corresponds to a response function k_n calculated from the interaction with e-cloud using single-pass PIC simulations.



These calculations use the e-cloud in the superconducting quadrupoles of the LHC for a beam energy of 450GeV.

[2] G. ladarola, et. al. 2020, Linearized method for the study of transverse instabilities driven by electron clouds



Describe the transverse centroid along the bunch, $\bar{x}(z)$, as a linear combination of test functions h_n :

$$\bar{x}(z) = \sum_{n=0}^{\infty} a_n h_n(z); \quad a_n = \frac{1}{H_n^2} \int \bar{x}(z) h_n(z) dz$$
 (2)





E-cloud in the Vlasov Equation - Dipolar Forces

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The kick, $\Delta x'$, of arbitrary distribution $\bar{x}(z)$ is:

$$\Delta x'(z) = \sum_{n=0}^{\infty} a_n k_n(z) \quad (3)$$

 k_n is the resulting electron cloud kick from a bunch distortion h_n .





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The coherent force can be expressed from the transverse kick $\Delta x'$. Assuming the force is distributed uniformly in the accelerator:

$$F_{x}^{coh}(z,t) = \frac{m_{0}\gamma v^{2}}{2\pi R} \Delta x'$$
(4)

 m_0 is the proton mass, γ is the relativistic gamma, v is the velocity of the protons and R is the total radius of the LHC.

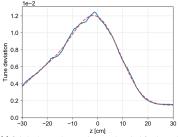


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E-cloud in Vlasov equation - Quadrupolar forces

Model detuning using a polynomial

$$\Delta Q(z) = \sum_{n=0}^{N_{\rho}} A_n z^n \qquad (5)$$



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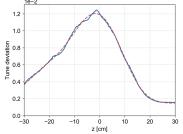
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$$\Delta Q(z) = \sum_{n=0}^{N_{p}} A_{n} z^{n} \qquad (5)$$

Generalize by adding chromaticity

$$\Delta Q(z,\delta) = \sum_{n=0}^{N_p} A_n z^n + B_n \delta^n \quad (6)$$



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Including only linear chromaticity:

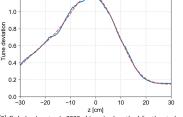
$$\Delta Q(z,\delta) = Q'\delta + \sum_{n=0}^{N_p} A_n z^n \quad (7)$$



1.2

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[2] G. Iadarola, et. al. 2020, Linearized method for the study of transverse instabilities driven by electron clouds

Equation to solve:

 $\frac{\partial \Delta \Psi}{\partial t} - \omega_0 (Q_{x0} + \Delta Q(r, \phi)) \frac{\partial \Delta \Psi}{\partial \theta_x} + \omega_s \frac{\partial \Delta \Psi}{\partial \phi} = -\frac{\eta g_0(r)}{\omega_s m_0 \gamma} \frac{df_0}{dJ_x} \sqrt{\frac{2J_x R}{Q_{x0}}} \sin \theta_x F_x^{coh}(z, t)$

 ω_0 is the angular revolution frequency, Q_0 is the unperturbed tune and Q_s is the synchrotron frequency



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Equation to solve:

$$\frac{\partial \Delta \Psi}{\partial t} - \omega_0 (Q_{x0} + \Delta Q(r, \phi)) \frac{\partial \Delta \Psi}{\partial \theta_x} + \omega_s \frac{\partial \Delta \Psi}{\partial \phi} = -\frac{\eta g_0(r)}{\omega_s m_0 \gamma} \frac{df_0}{dJ_x} \sqrt{\frac{2J_x R}{Q_{x0}}} \sin \theta_x F_x^{coh}(z, t)$$
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detuning from e-cloud and chromaticity dipolar forces from e-cloud

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detuning from e-cloud and chromaticity dipolar forces from e-cloud

First ansatz:

 $\Delta \psi(J_x, \theta_x, r, \phi, t) = e^{j\Omega t} \Delta \psi(J_x, \theta_x, r, \phi).$

 ω_{0} is the angular revolution frequency, Q_{0} is the unperturbed tune and Q_{s} is the synchrotron frequency



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Equation to solve:

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First ansatz:
$$\Delta \psi(J_x, \theta_x, r, \phi, t) = e^{i\Omega t} \Delta \psi(J_x, \theta_x, r, \phi).$$

$$\frac{(Re(\Omega) - Q_0)/Q_s}{\int_{0}^{0} \frac{1}{Q_x} - Q_0} \int_{0}^{0} \frac{1}{Q_x} \int_{0$$

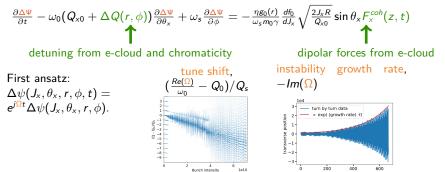
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Equation to solve:



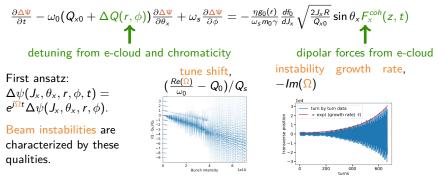
 $\omega_{\rm D}$ is the angular revolution frequency, $Q_{\rm D}$ is the unperturbed tune and $Q_{\rm S}$ is the synchrotron frequency



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Equation to solve:

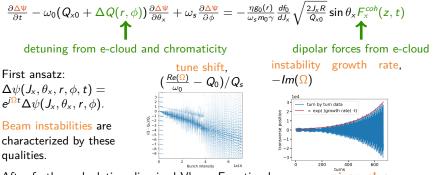


 $\omega_{\rm D}$ is the angular revolution frequency, $Q_{\rm D}$ is the unperturbed tune and $Q_{\rm S}$ is the synchrotron frequency



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Equation to solve:



After further calculations linerized Vlasov Equation becomes an eigenvalue problem:

$$b_{lm}(\Omega - Q_{x0}\omega_0 - l\omega_s) = \sum_{l'm'} (\mathbf{M}_{lm,l'm'} + \tilde{\mathbf{M}}_{lm,l'm'}) b_{l'm'}$$
(8)

Unknowns in red and terms including electron cloud forces in green ω_0 is the angular revolution frequency. Q_0 is the unperturbed tune and Q_c is the synchrotron frequency



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E-cloud in a Vlasov formalism

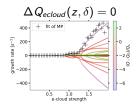
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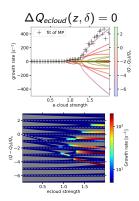


benchmarked against macro-particle simulations using the same formalism of e-cloud forces



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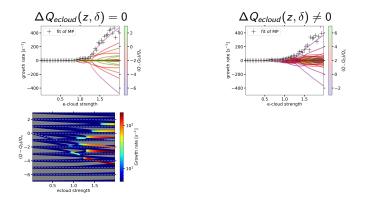


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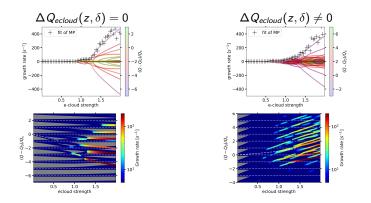




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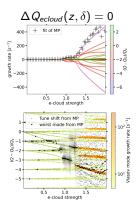


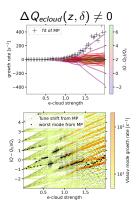


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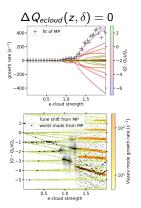


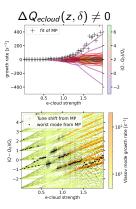




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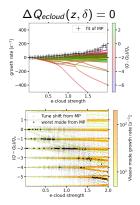




Good agreement with MP simulations!



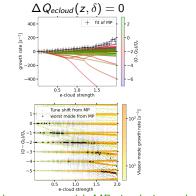
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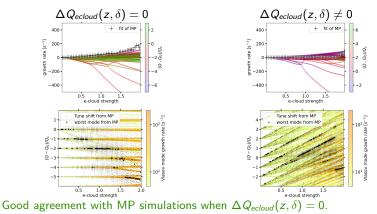


Good agreement with MP simulations when $\Delta Q_{ecloud}(z, \delta) = 0$.



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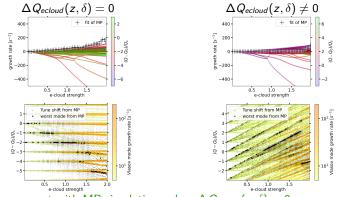






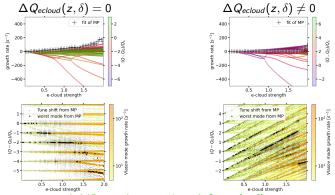
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Good agreement with MP simulations when $\Delta Q_{ecloud}(z, \delta) = 0$. For positive chromaticity, weak Vlasov modes are not visible in the macro-particle simulations.





Good agreement with MP simulations when $\Delta Q_{ecloud}(z, \delta) = 0$. For positive chromaticity, weak Vlasov modes are not visible in the macro-particle simulations.

The tune shift of the Vlasov modes agree well with the macro-particle spectra.





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First benchmark with measurements



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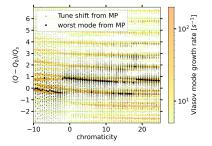
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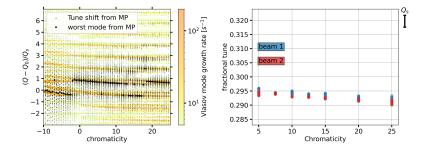


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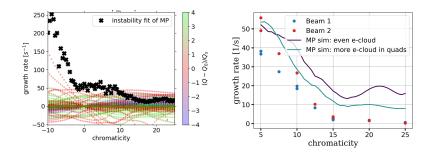


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This is confirmed by measurements.



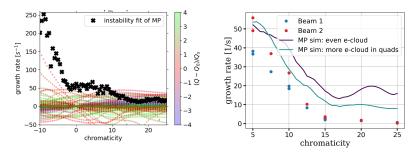
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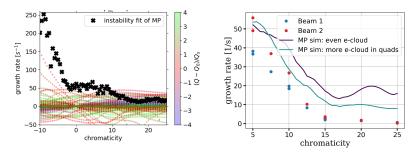


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The macroparticle simulations are damped by incoherent mechanisms not captured by the Vlasov model.





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Thank you for your attention!



References

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