

Muonic atom spectroscopy with radioactive targets

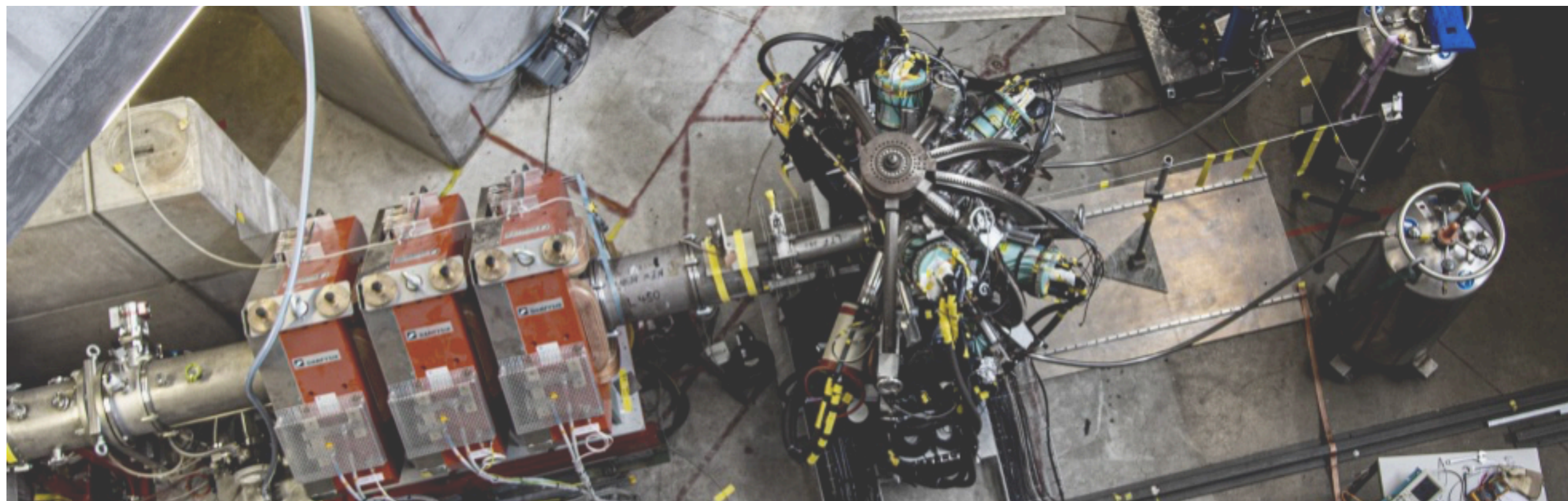
Stella Vogiatzi

On behalf of the muX collaboration*

05.09.2023

Annual Meeting of the Swiss Physical Society

* Institute of Nuclear Physics, Polish Academy of Sciences, Krakow, Poland | Paul Scherrer Institut, Villigen, Switzerland | ETH Zürich, Switzerland | Johannes Gutenberg University Mainz, Germany | KU Leuven, Belgium | GSI Helmholtzzentrum für Schwerionenforschung Darmstadt, Germany | Helmholtz Institute Mainz, Germany | Institut für Kernphysik, Universität zu Köln, Germany | LKB Paris, France | University of Groningen, The Netherlands | University of Pisa and INFN, Pisa, Italy | University of Victoria, Canada | Perimeter Institute, Waterloo, Canada | CSNSM, Université Paris Sud, CNRS/IN2P3, Université Paris Saclay, Orsay Campus, France



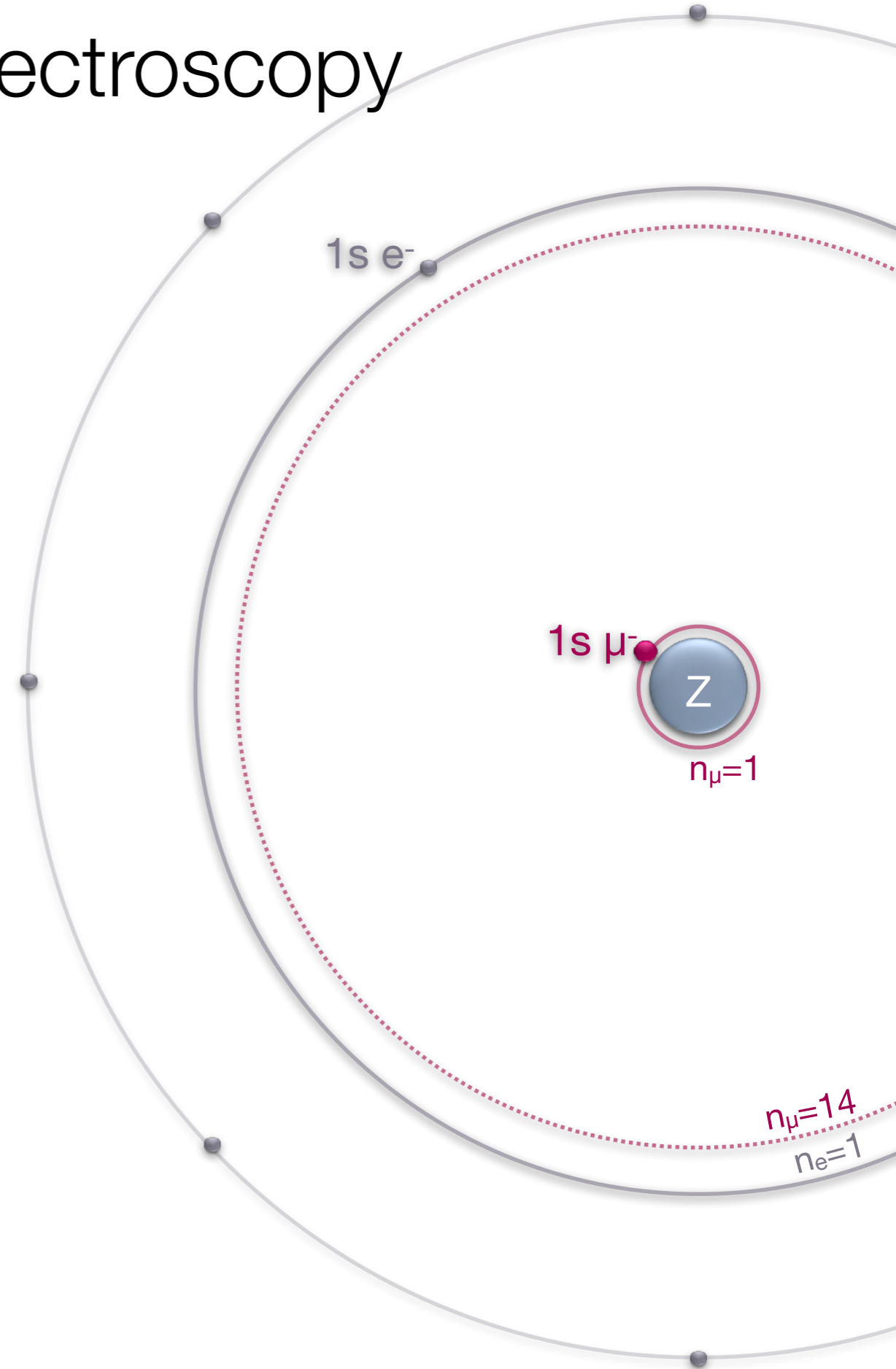
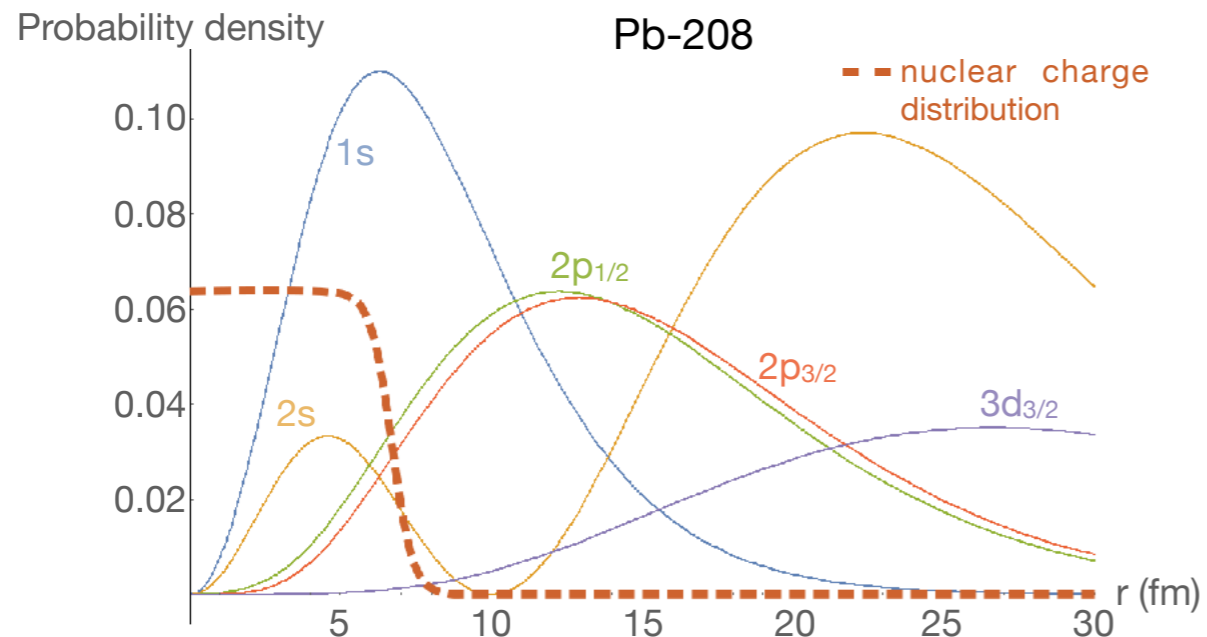
Muonic atom spectroscopy

- Large overlap of the low-lying muonic states with the nuclear charge distribution.
- The measurement of the muonic energy levels allows to extract properties of the nucleus such as the charge radius R_N :

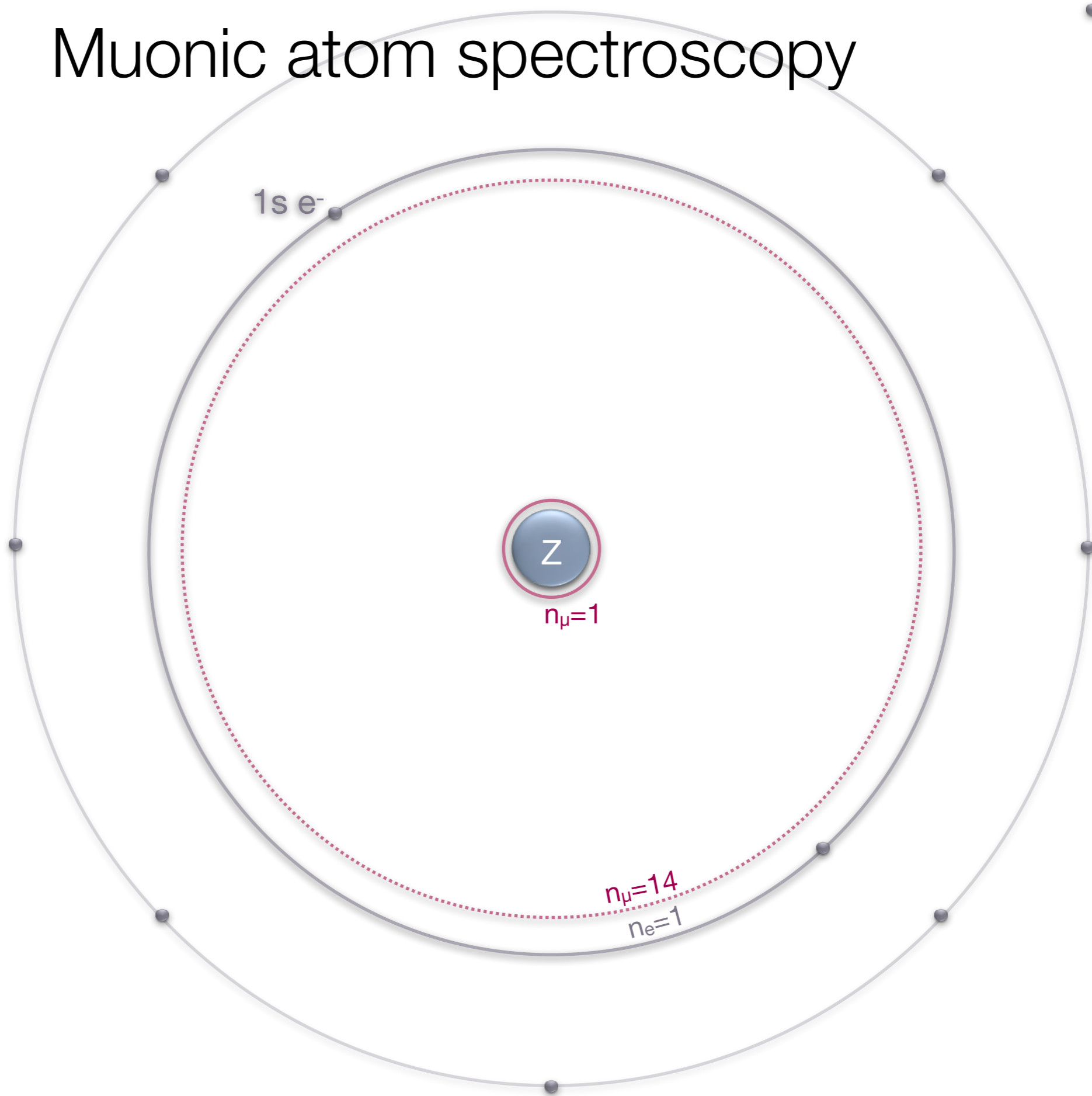
2p-1s transition most sensitive to R_N

$$R_N^2 = \langle r^2 \rangle = \frac{1}{Ze} \int d^3\vec{r} \rho(\vec{r}) r^2$$

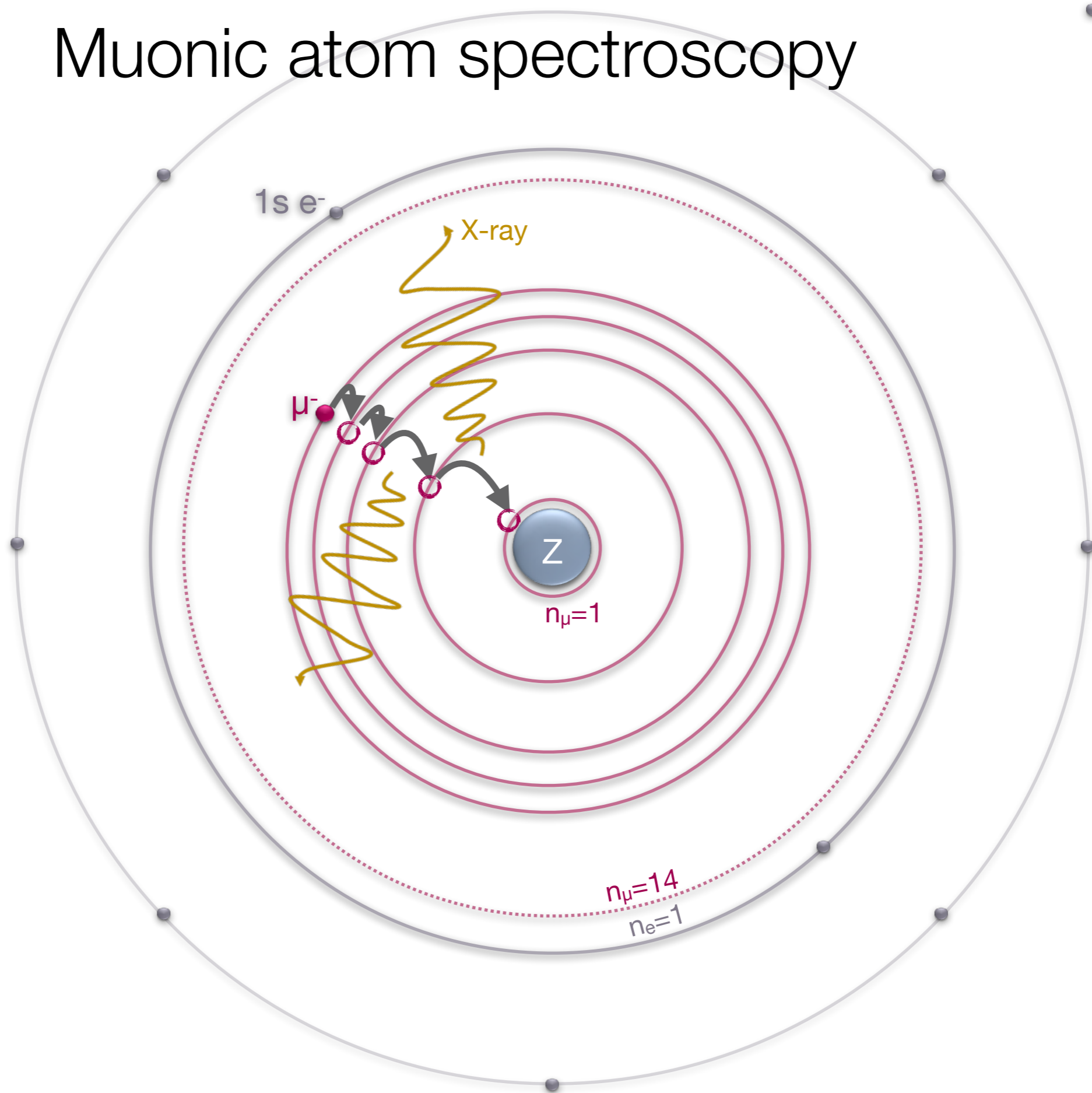
nuclear charge distribution



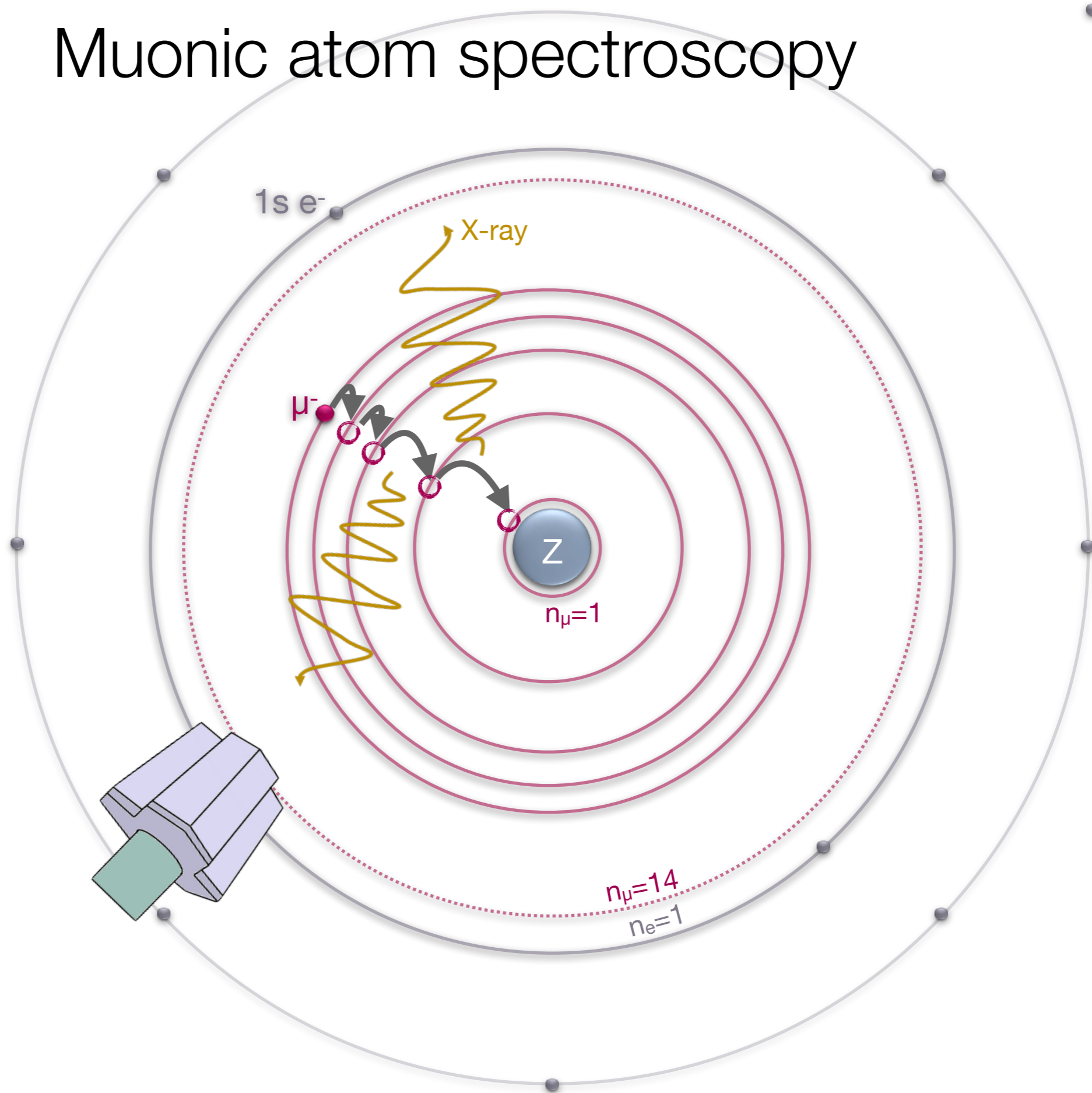
Muonic atom spectroscopy



Muonic atom spectroscopy

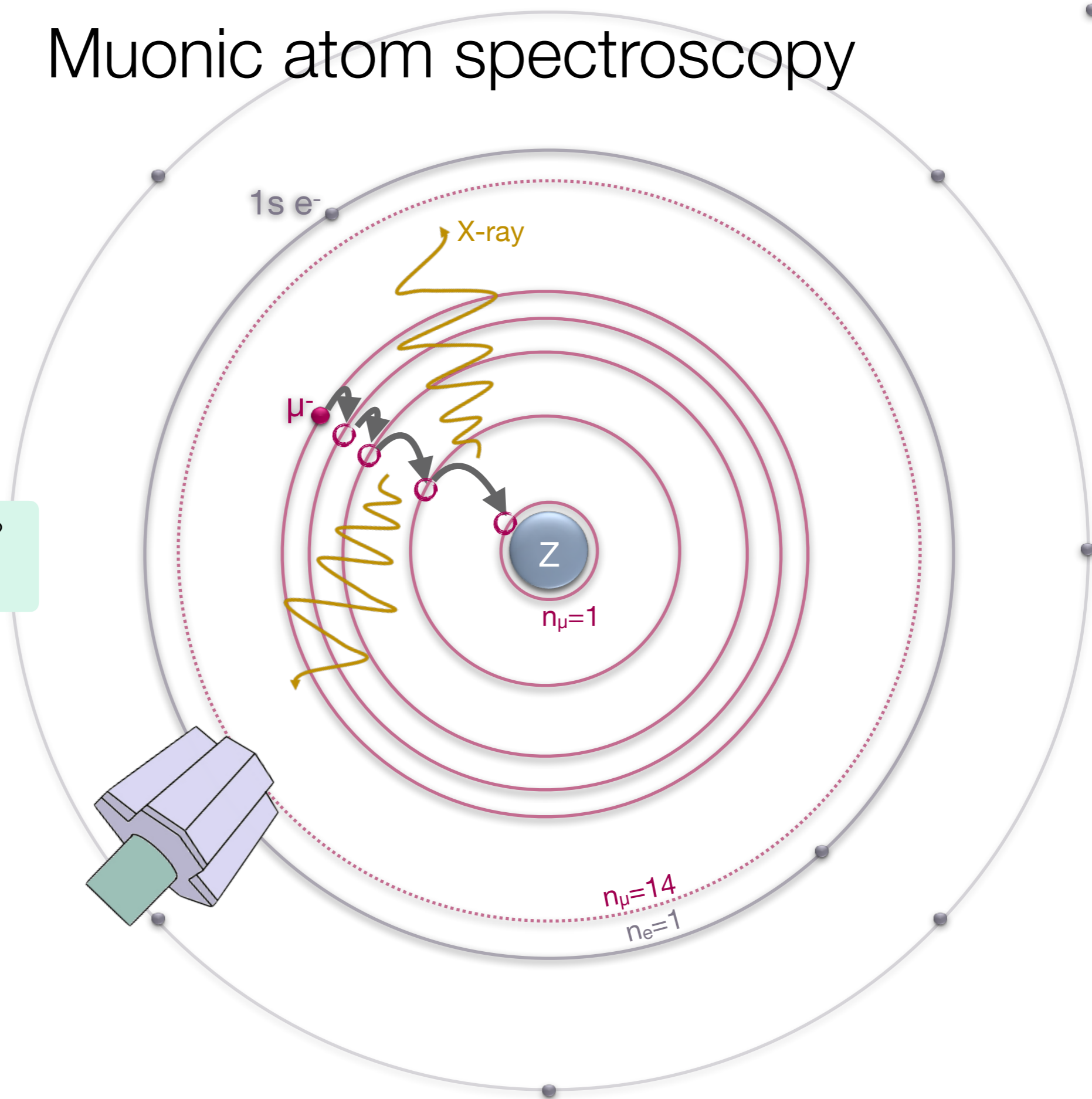


Muonic atom spectroscopy



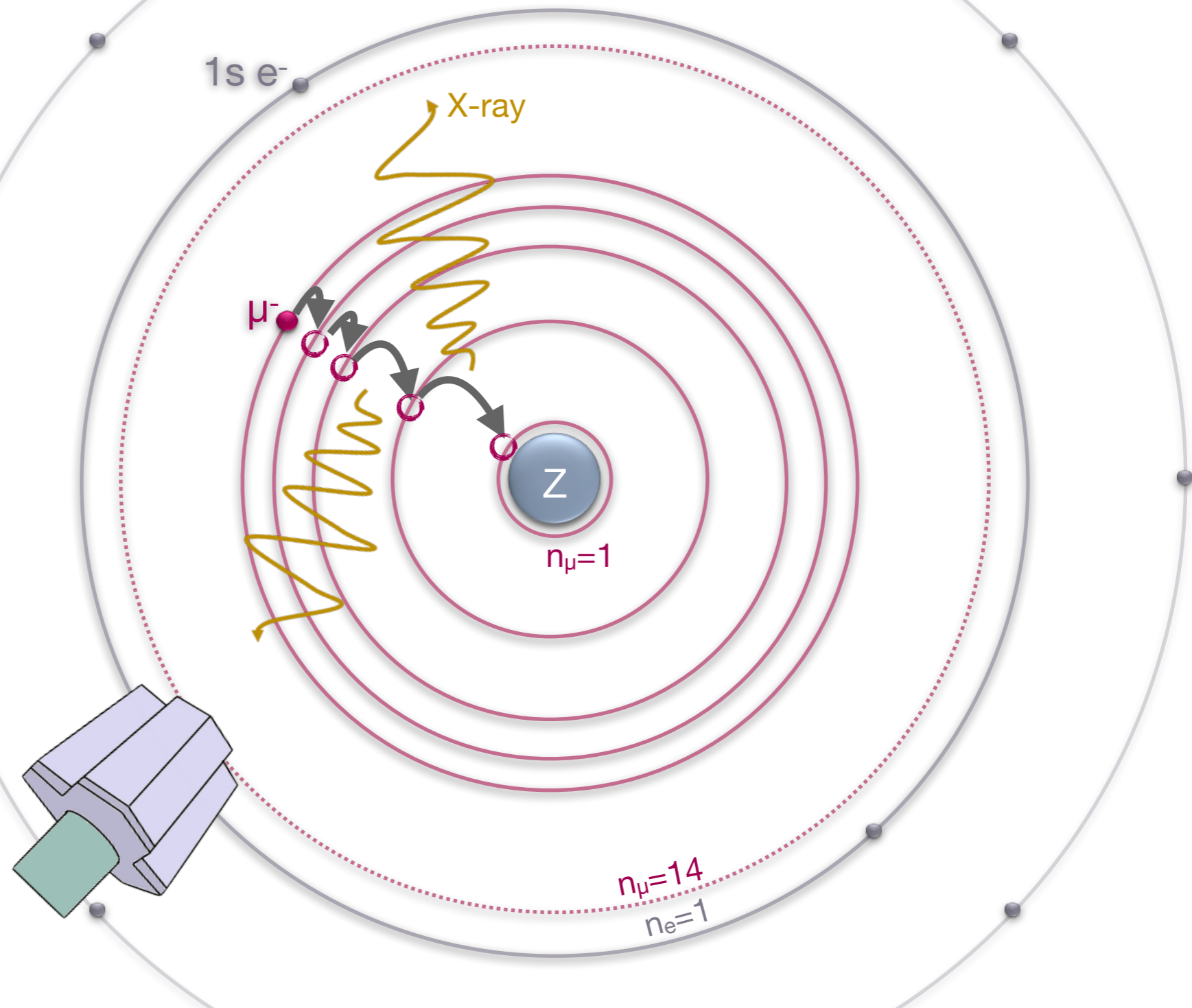
Muonic atom spectroscopy

What to measure?
 ^{248}Cm & ^{226}Ra



Muonic atom spectroscopy

What to measure?
 ^{248}Cm & ^{226}Ra



... only possible with $\gtrsim 100$ mg targets... radioactive targets are typically restricted to $\sim \mu\text{g}$ quantities in the experimental areas

Where does this happen?

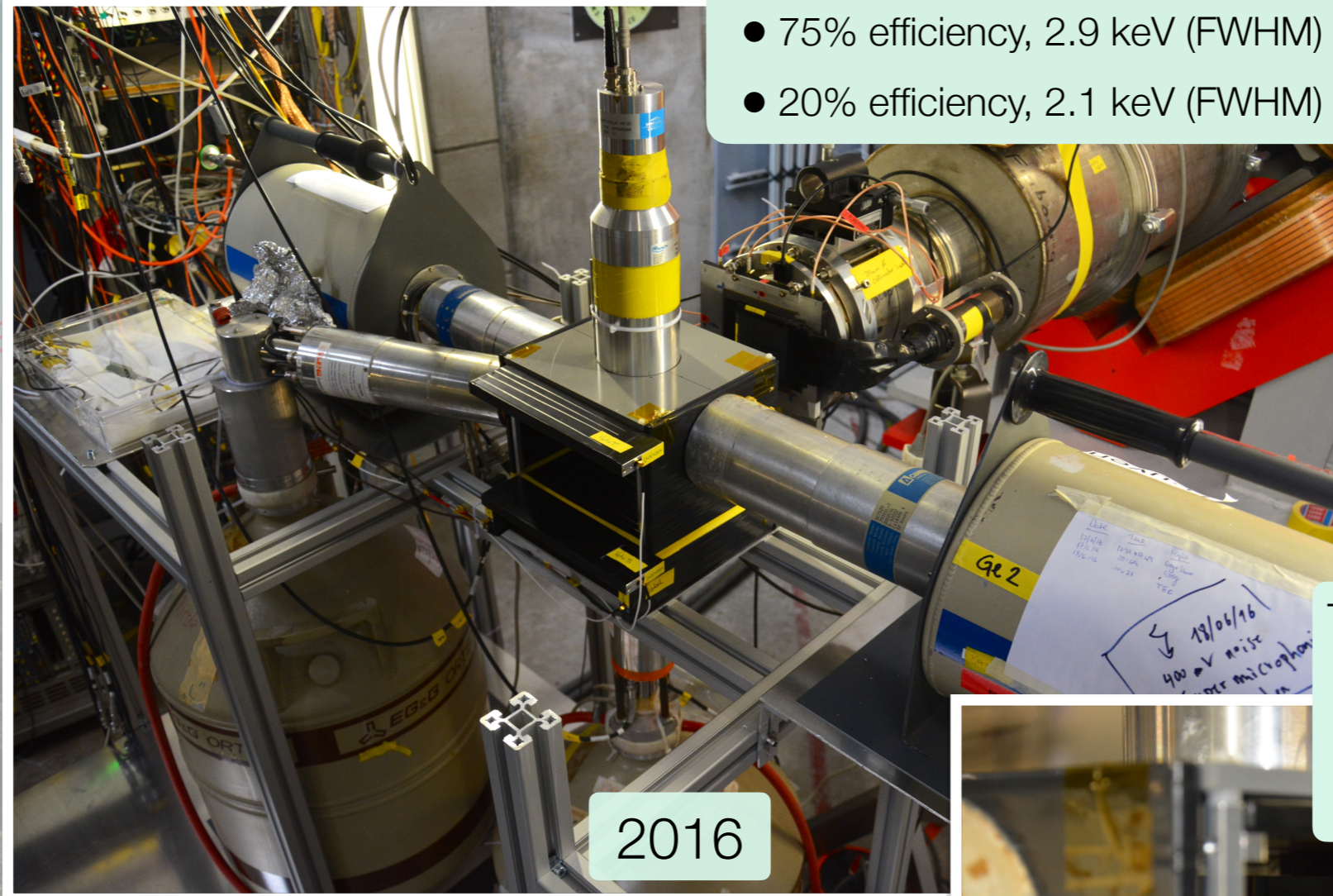


High Intensity Proton
Accelerator (HIPA) facility at Paul
Scherrer Institut

Where does this happen?

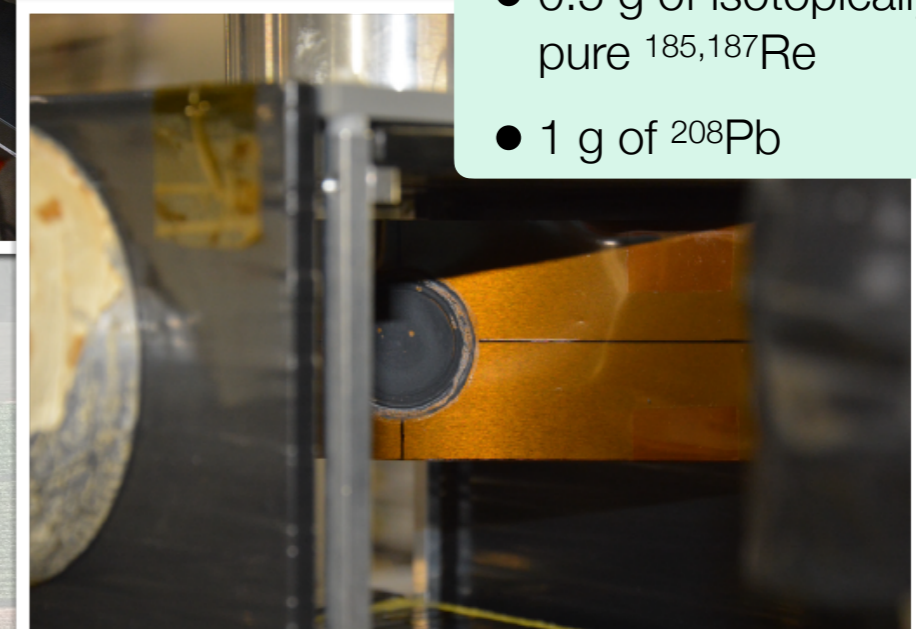
2 x HPGe detectors:

- 75% efficiency, 2.9 keV (FWHM) @1.3 MeV
- 20% efficiency, 2.1 keV (FWHM) @1.3 MeV



Targets:

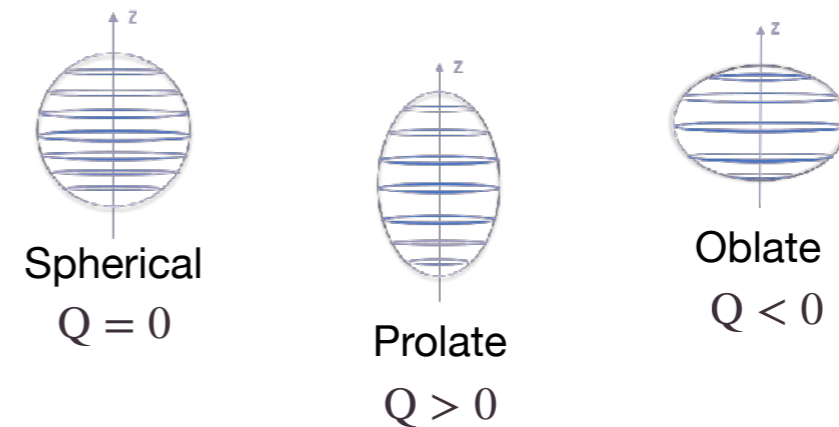
- 0.5 g of isotopically pure $^{185,187}\text{Re}$
- 1 g of ^{208}Pb



2016

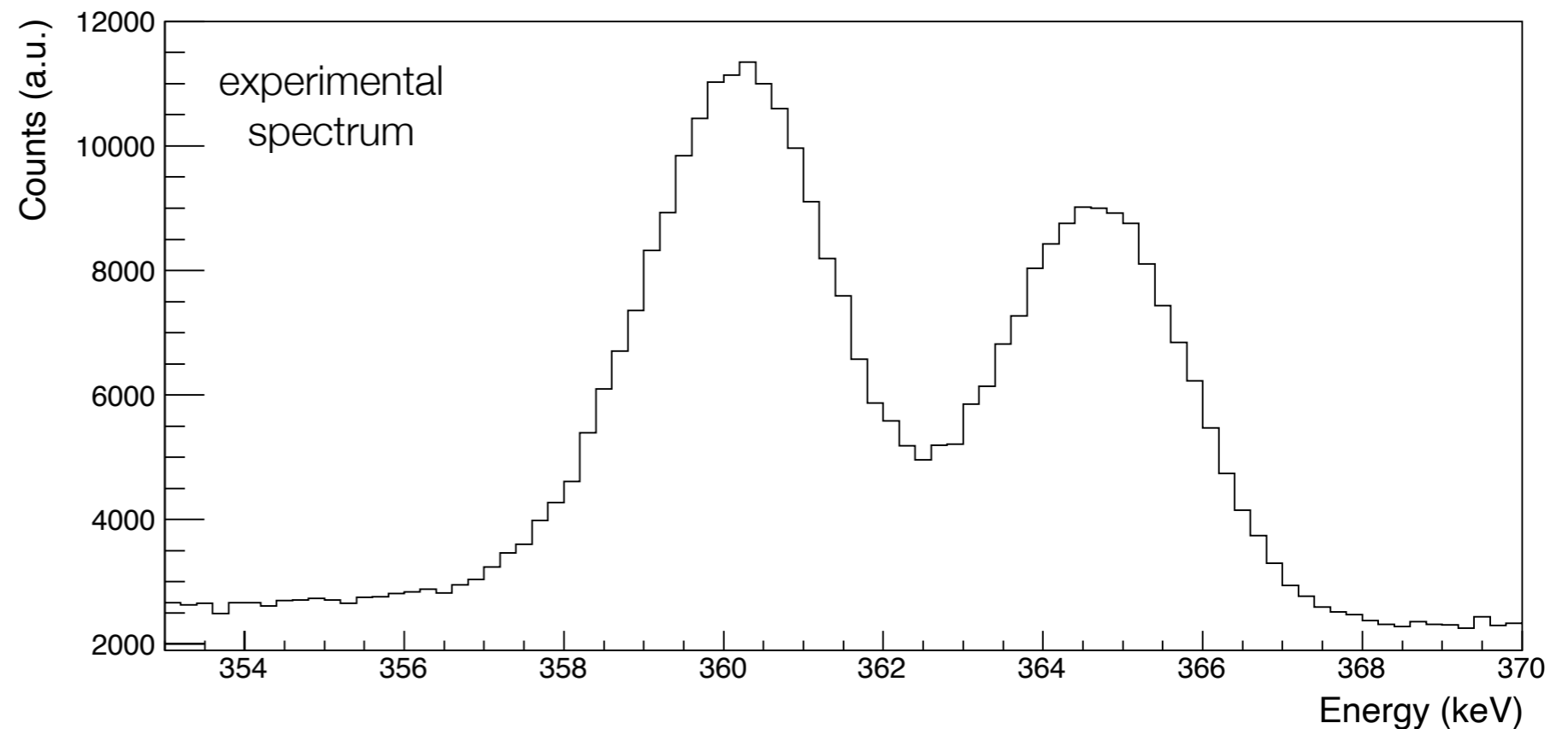
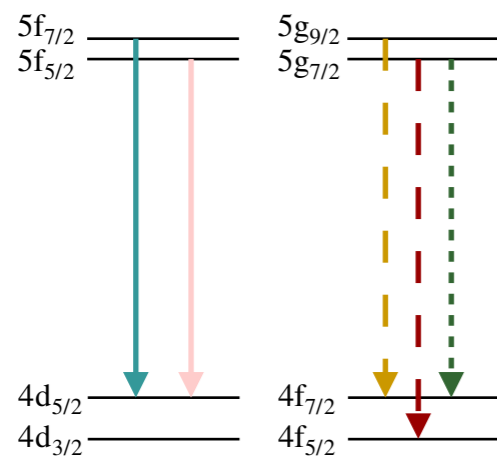
Spectroscopic quadrupole moment in $^{185,187}\text{Re}$

The spectroscopic quadrupole moment Q is extracted from the $5 \rightarrow 4$ transitions being not sensitive to the details of the nuclear charge distribution



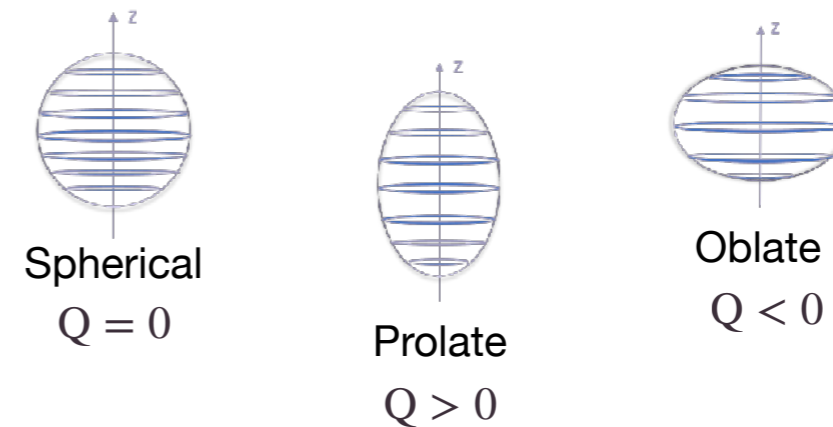
To extract the quadrupole moment:

- Theoretical predictions of hyperfine transitions
- Good knowledge of the instrumental response is required



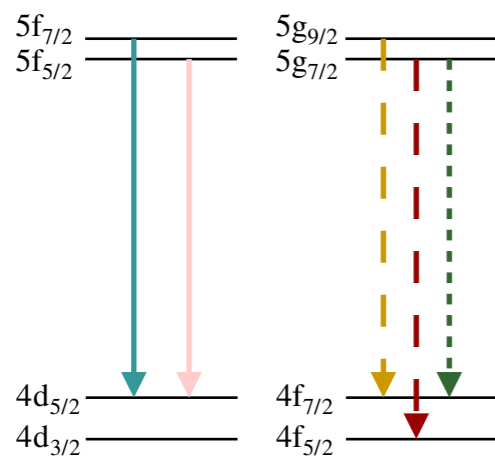
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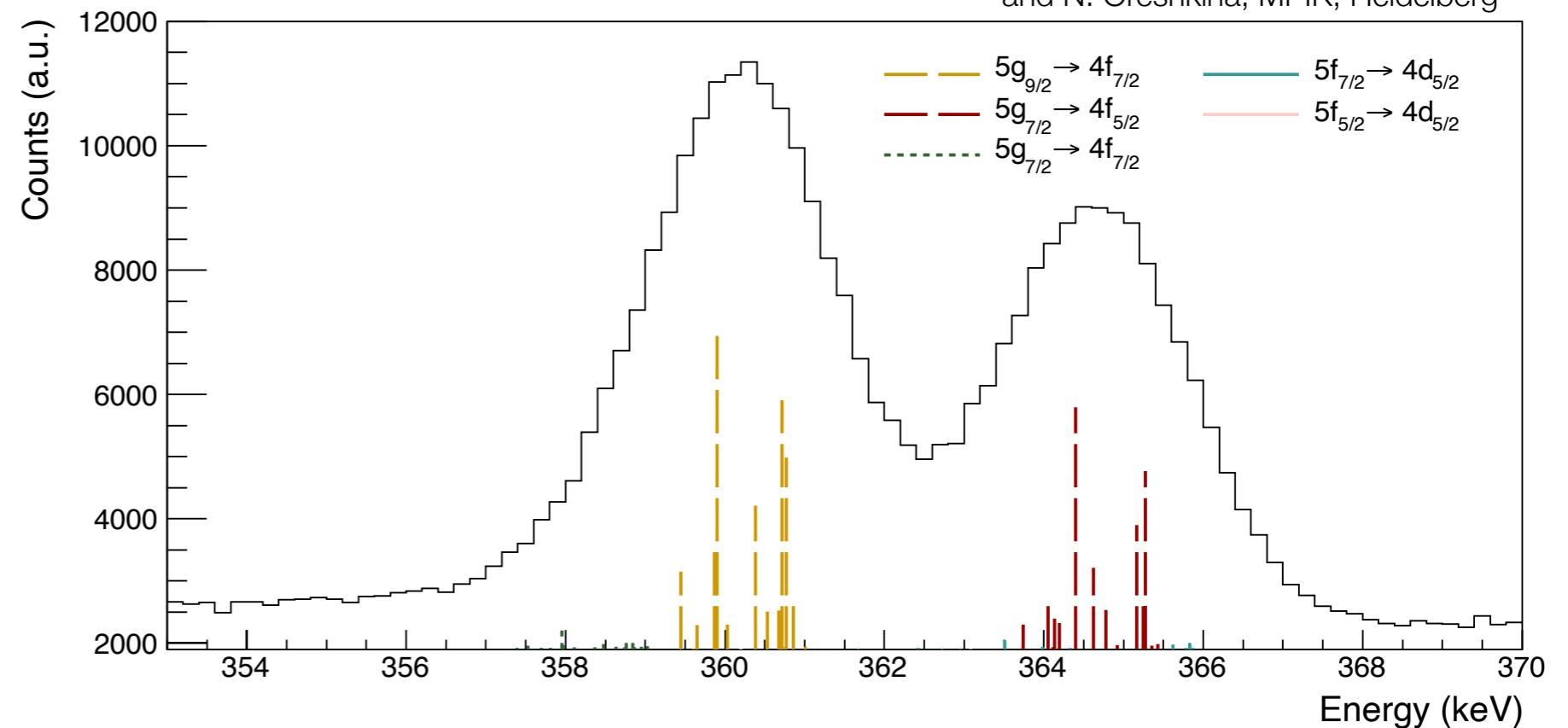


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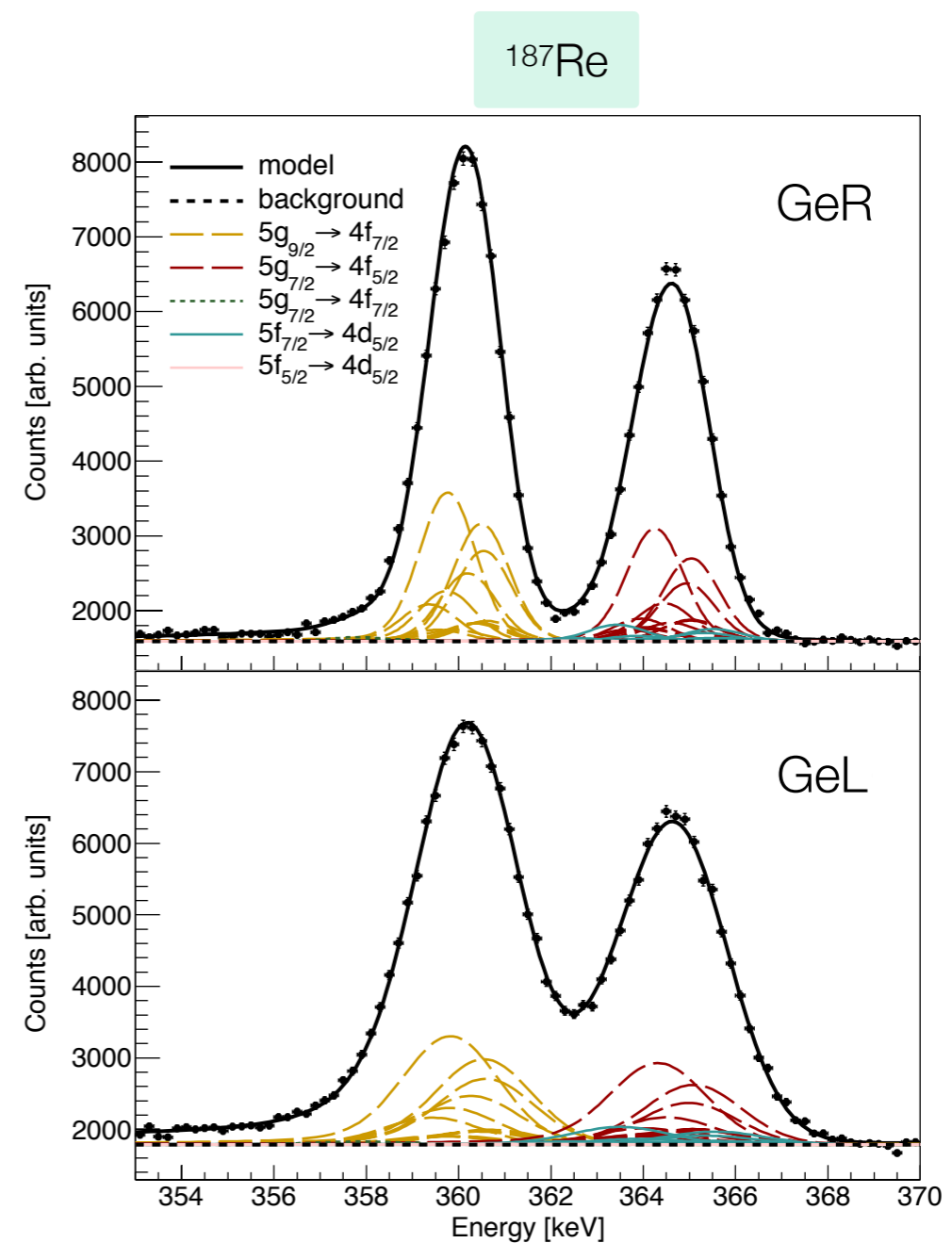
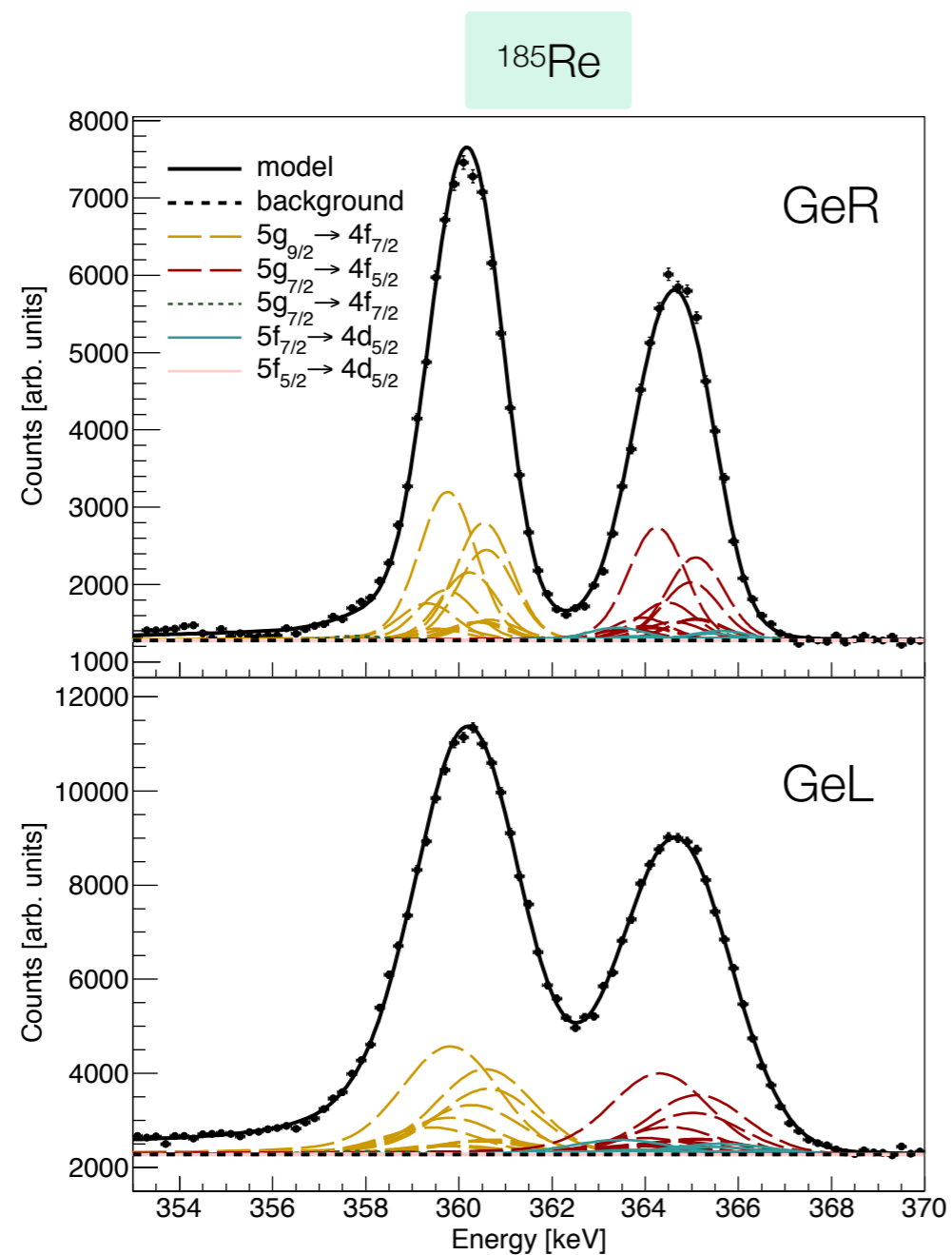


Theoretical calculations by N. Michel and N. Oreshkina, MPIK, Heidelberg



Fit the $5 \rightarrow 4$ transitions in $^{185,187}\text{Re}$

- Fit of the theoretical predictions in the experimental spectrum with the quadrupole moment as a free parameter



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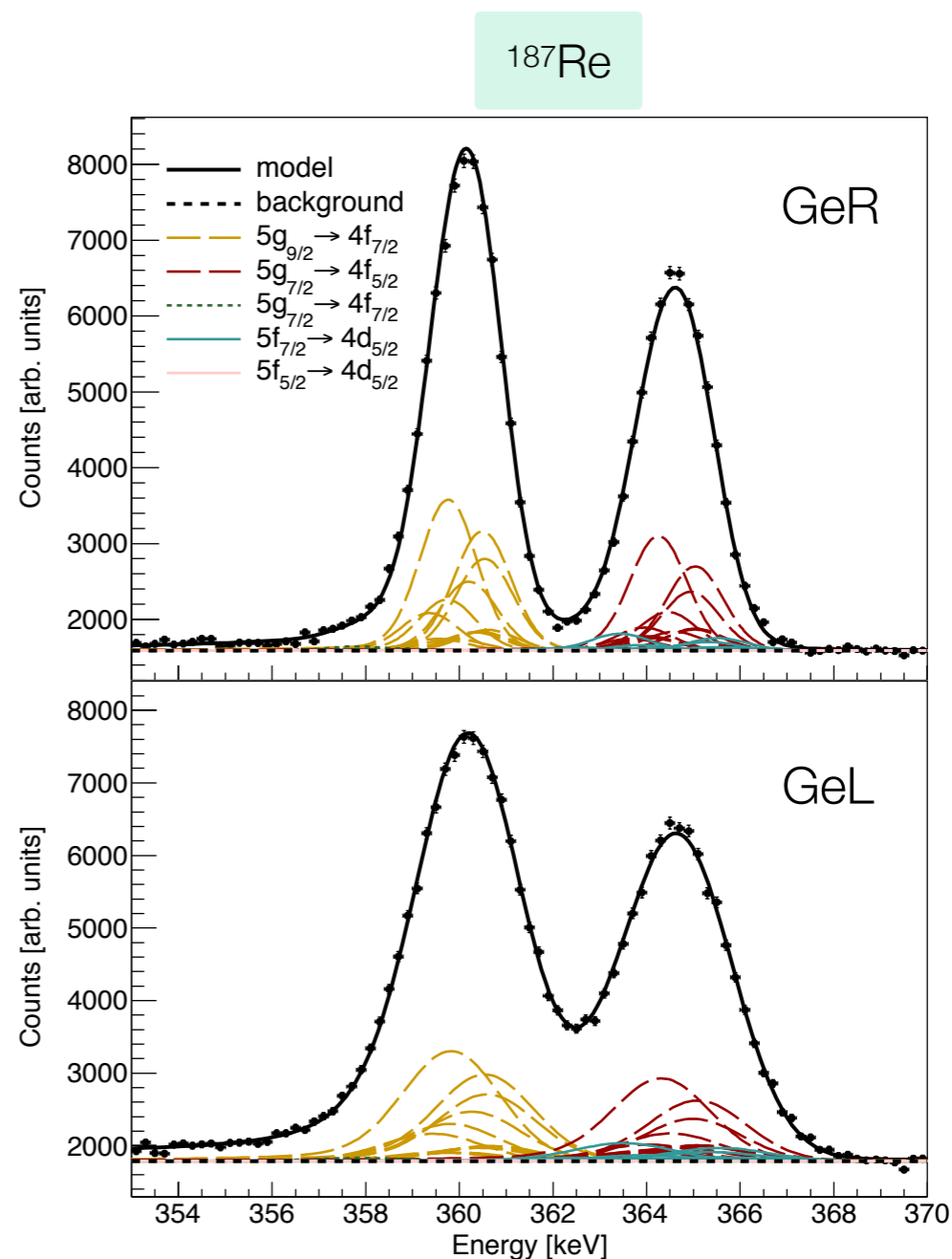
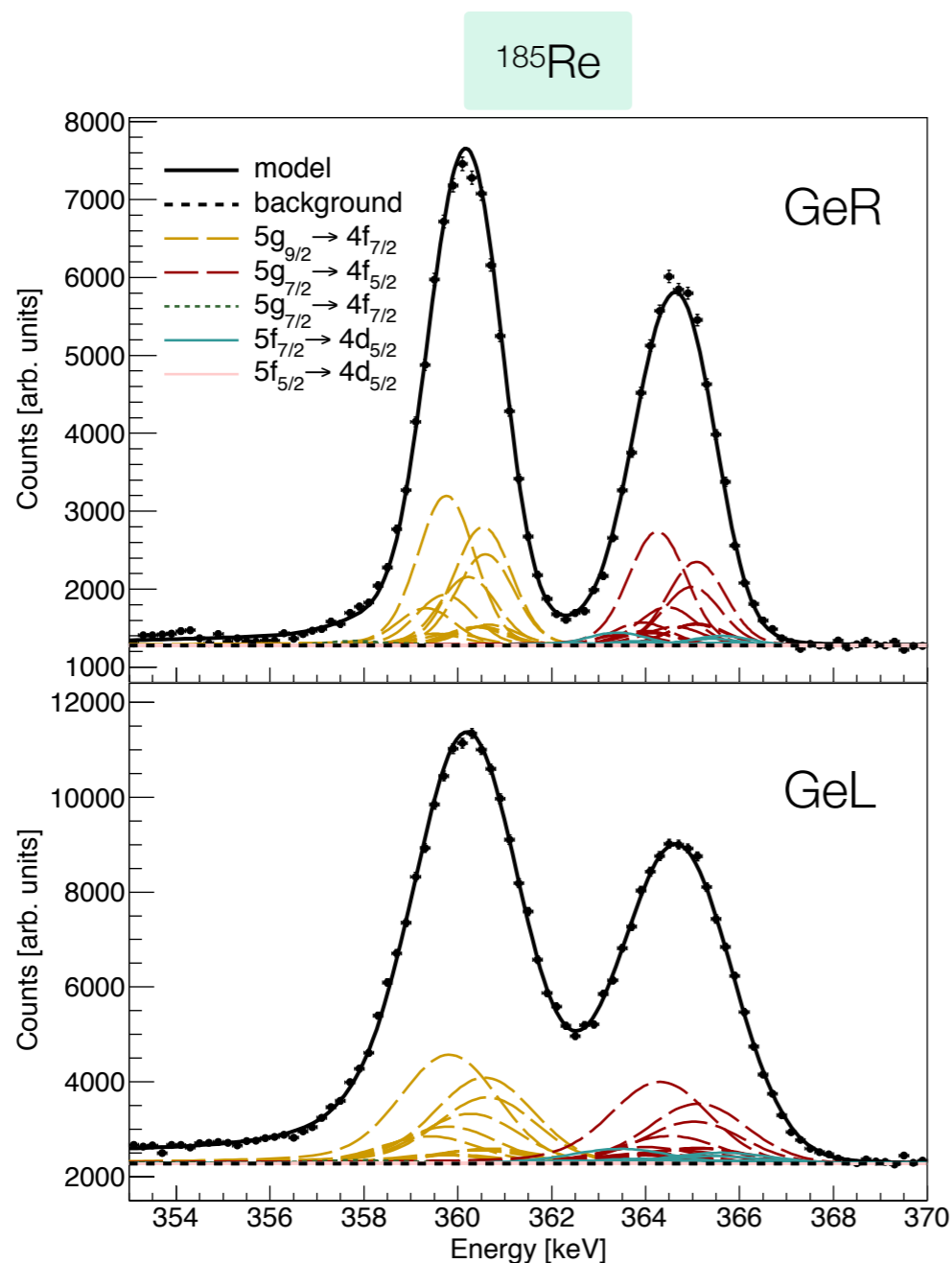
- First physics result:

$$Q(^{185}\text{Re}) = 2.07(5) \text{ b}$$

$$Q(^{187}\text{Re}) = 1.94(5) \text{ b}$$

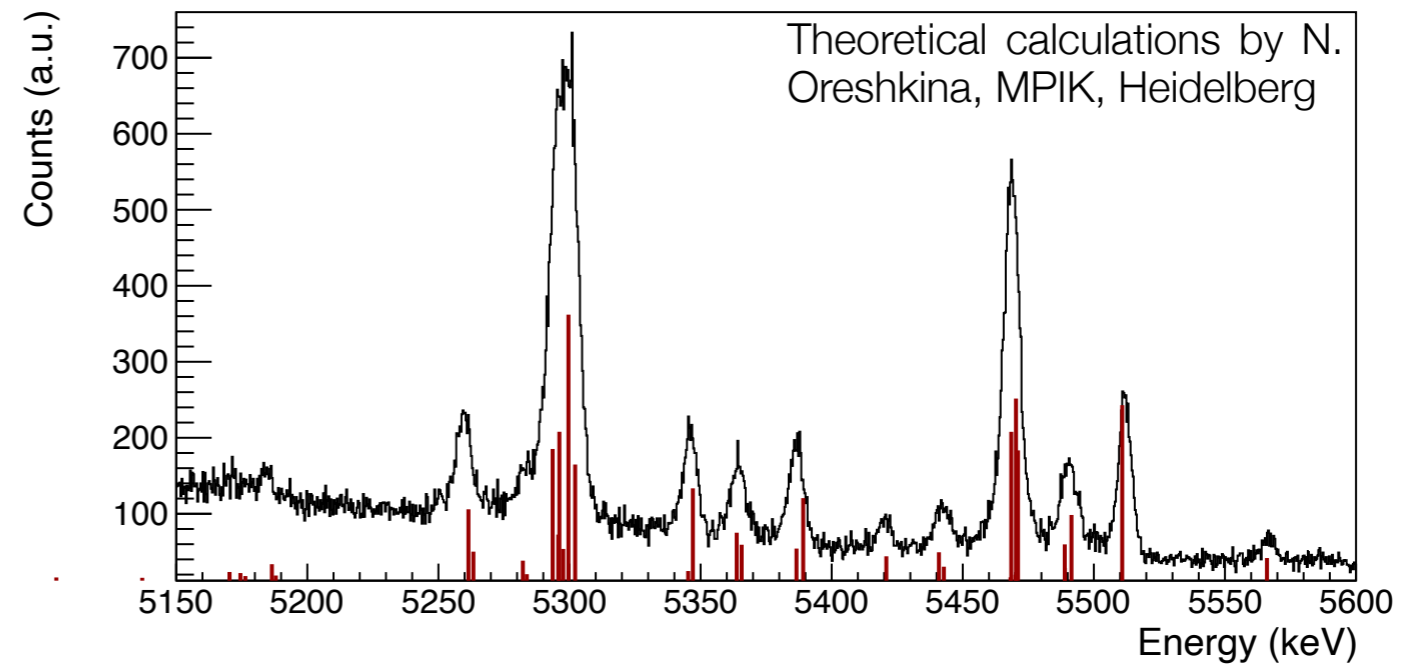
Measurement of the quadrupole moment of ^{185}Re and ^{187}Re from the hyperfine structure of muonic X rays

A. Antognini,^{1,2} N. Berger,³ T. E. Cocolios,⁴ R. Dressler,¹ R. Eichler,¹ A. Eggenberger,² P. Indelicato,⁵ K. Jungmann,⁶ C. H. Keitel,⁷ K. Kirch,^{1,2} A. Knecht,¹ N. Michel,⁷ J. Nuber,^{1,2} N. S. Oreshkina,^{7,*} A. Ouf,⁸ A. Papa,^{1,9} R. Pohl,⁸ M. Pospelov,^{10,11} E. Rapisarda,^{1,1} N. Ritjoho,^{1,2} S. Rocca,^{12,13} N. Severijns,⁴ A. Skawran,^{1,2} S. M. Vogiatzi,^{1,2} F. Wauters,³ and L. Willmann⁶



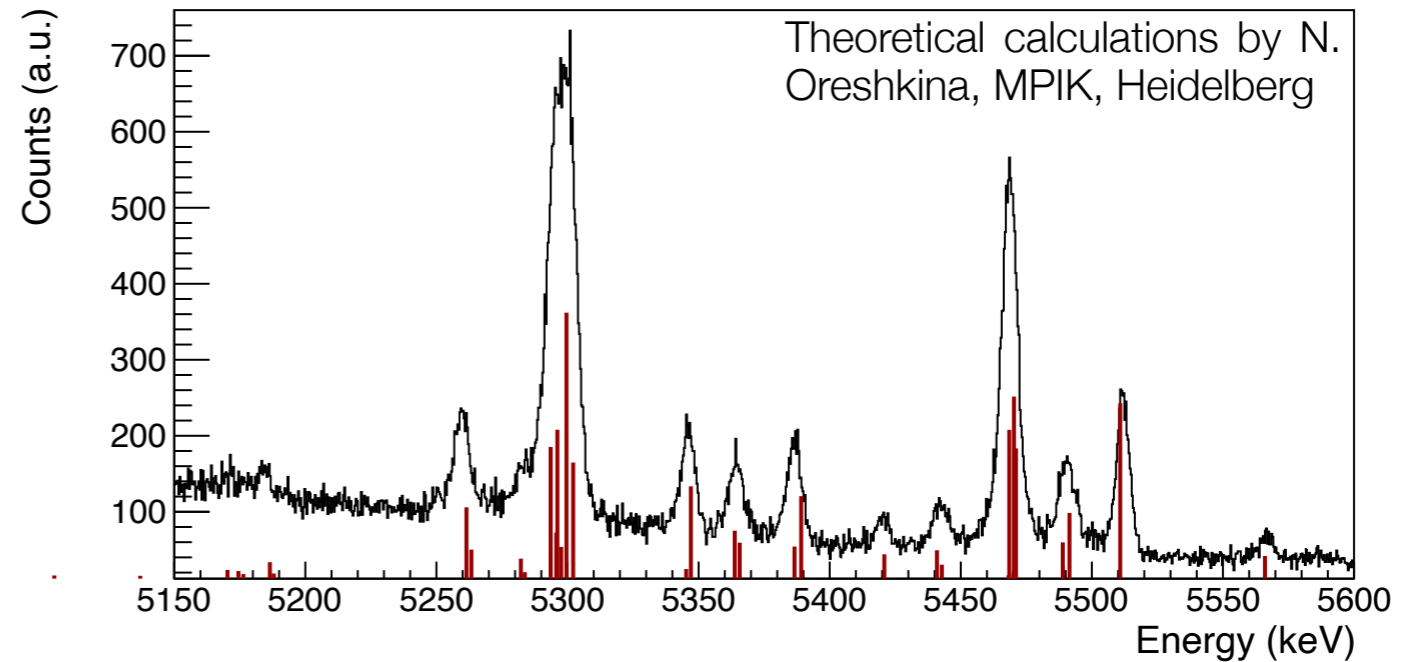
Nuclear charge radius in $^{185,187}\text{Re}$

The extraction of the nuclear charge radius from the analysis of the $2p \rightarrow 1s$ hyperfine transitions



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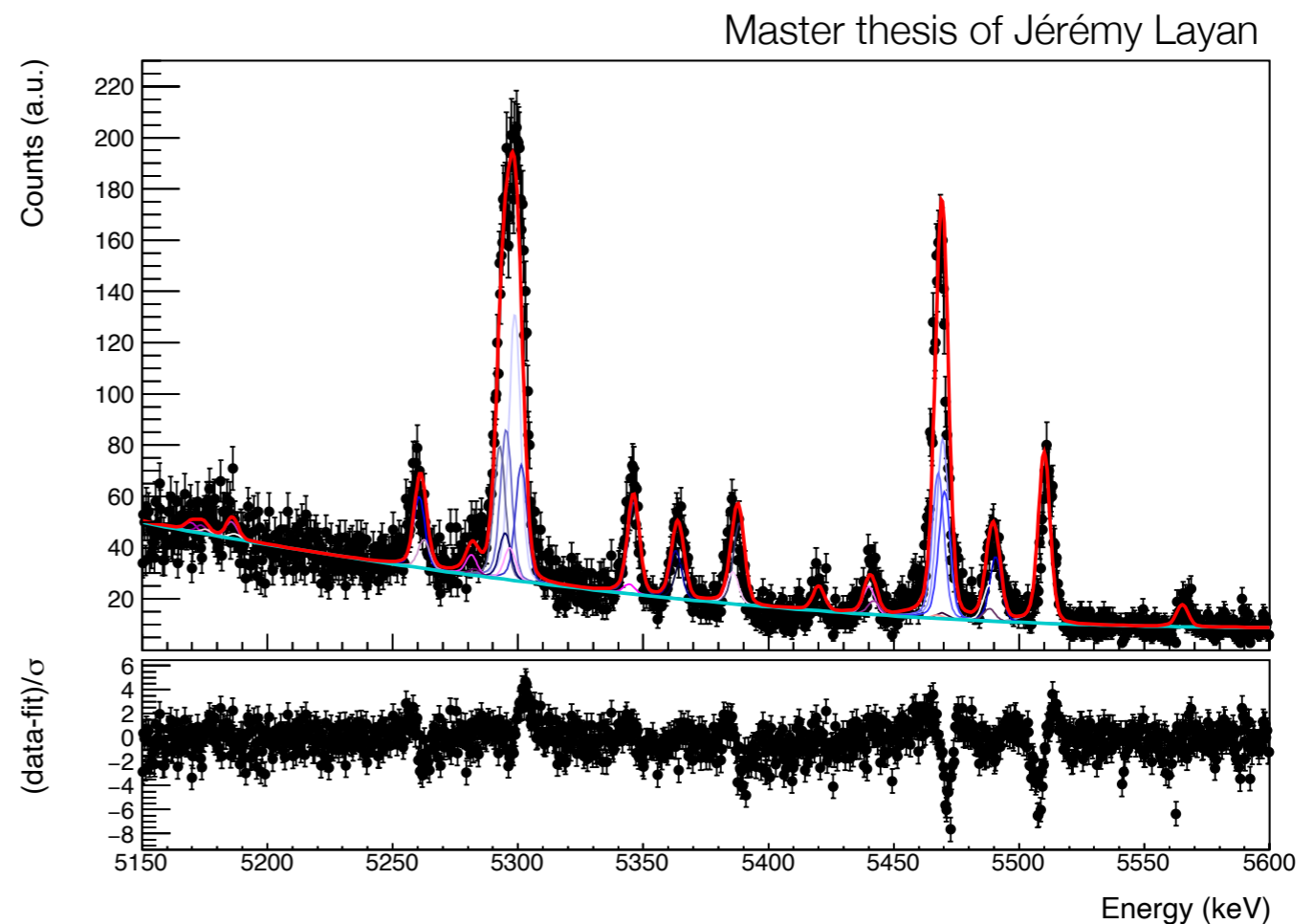


Preliminary results on the nuclear charge radii:

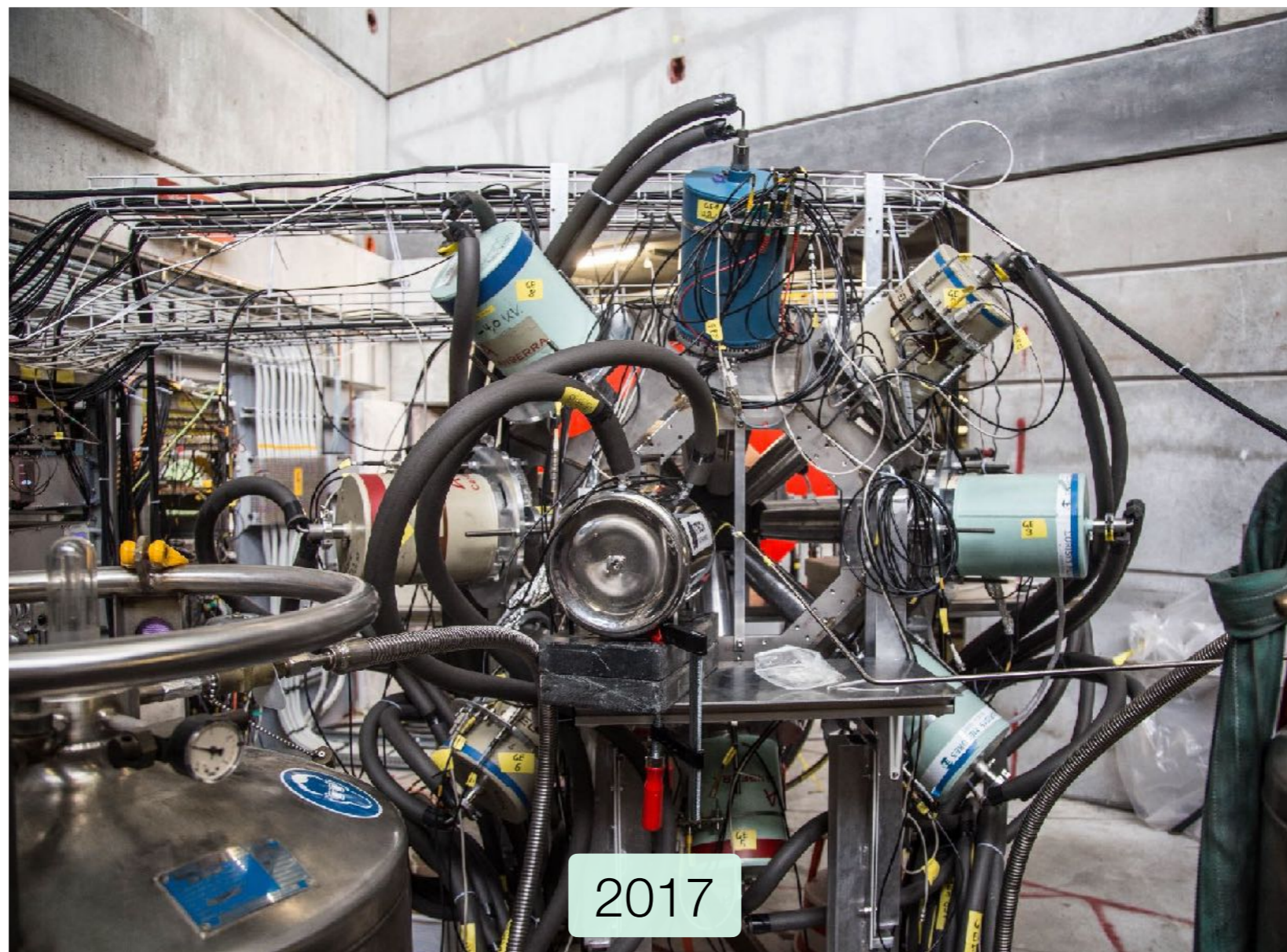
$$R(^{185}\text{Re}) = 5.297(2)_{\text{stat}}(6)_{\text{sys}} \text{ fm}$$

$$R(^{187}\text{Re}) = 5.288(2)_{\text{stat}}(4)_{\text{sys}} \text{ fm}$$

Not all the systematics are taken into account.



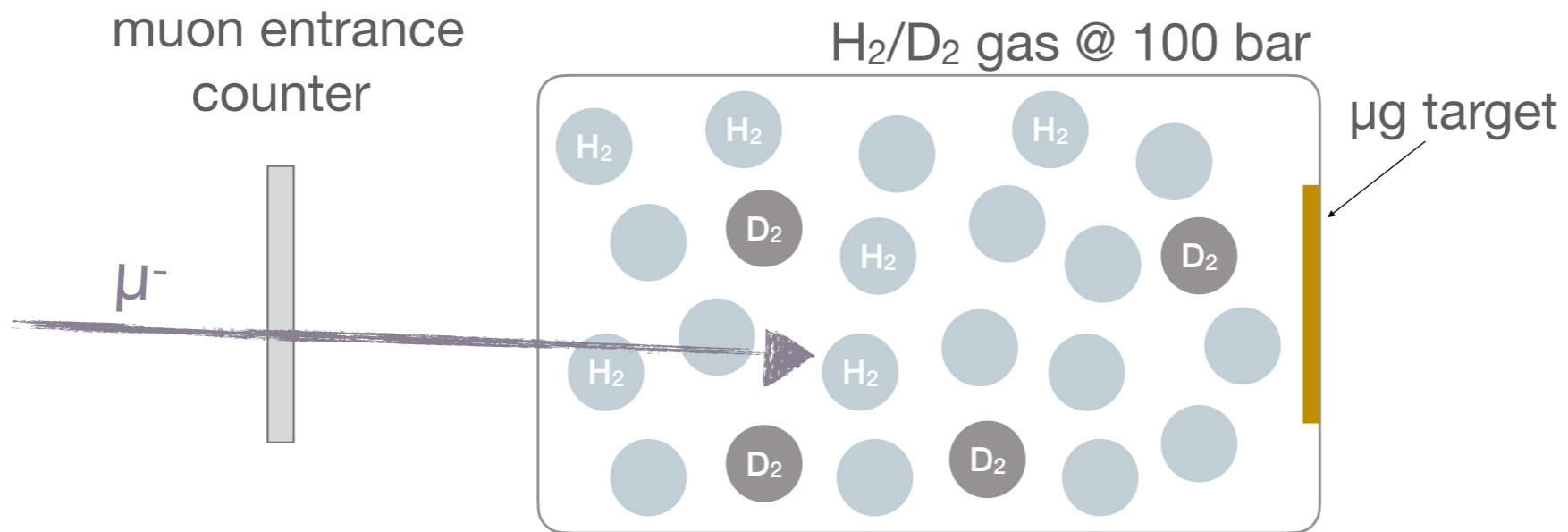
Towards the measurement of μg targets



2017

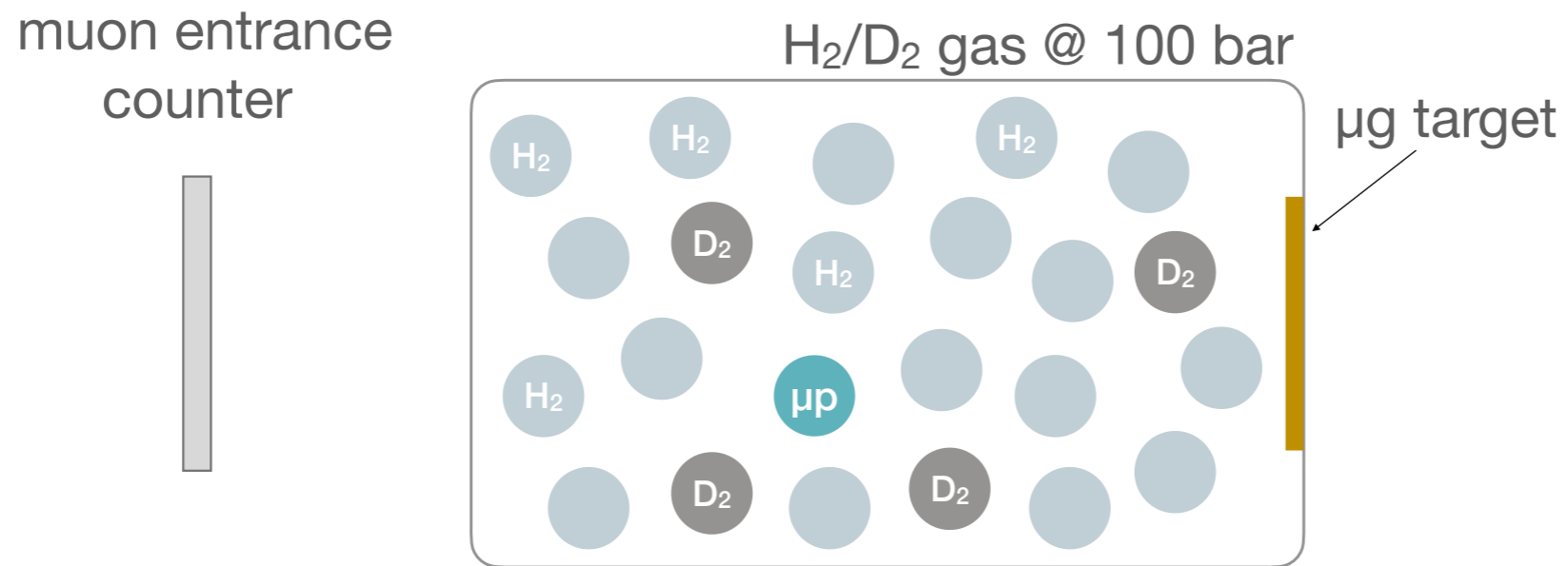
Muon transfer to microgram targets

Inspired by the work of Strasser *et al.* and Kraiman *et al.*



Muon transfer to microgram targets

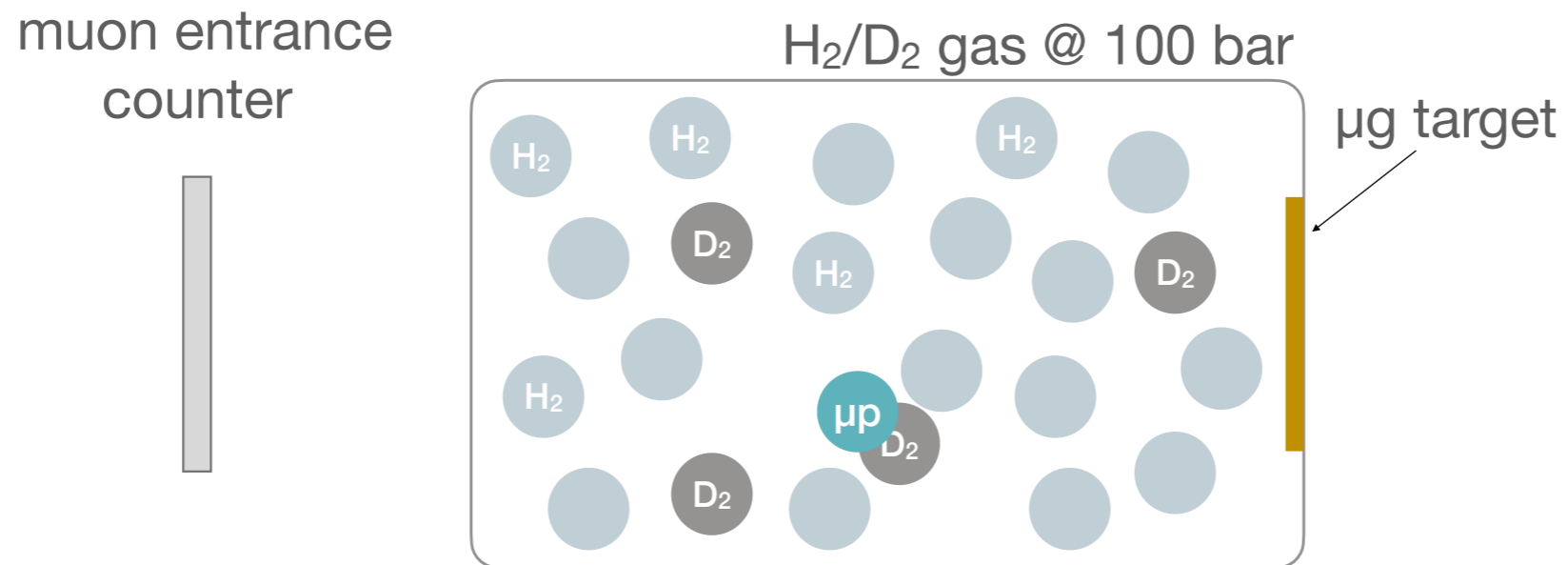
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1. μ^- stops in 100 bar of H₂ + 0.25% D₂ & forms muonic hydrogen μp

Muon transfer to microgram targets

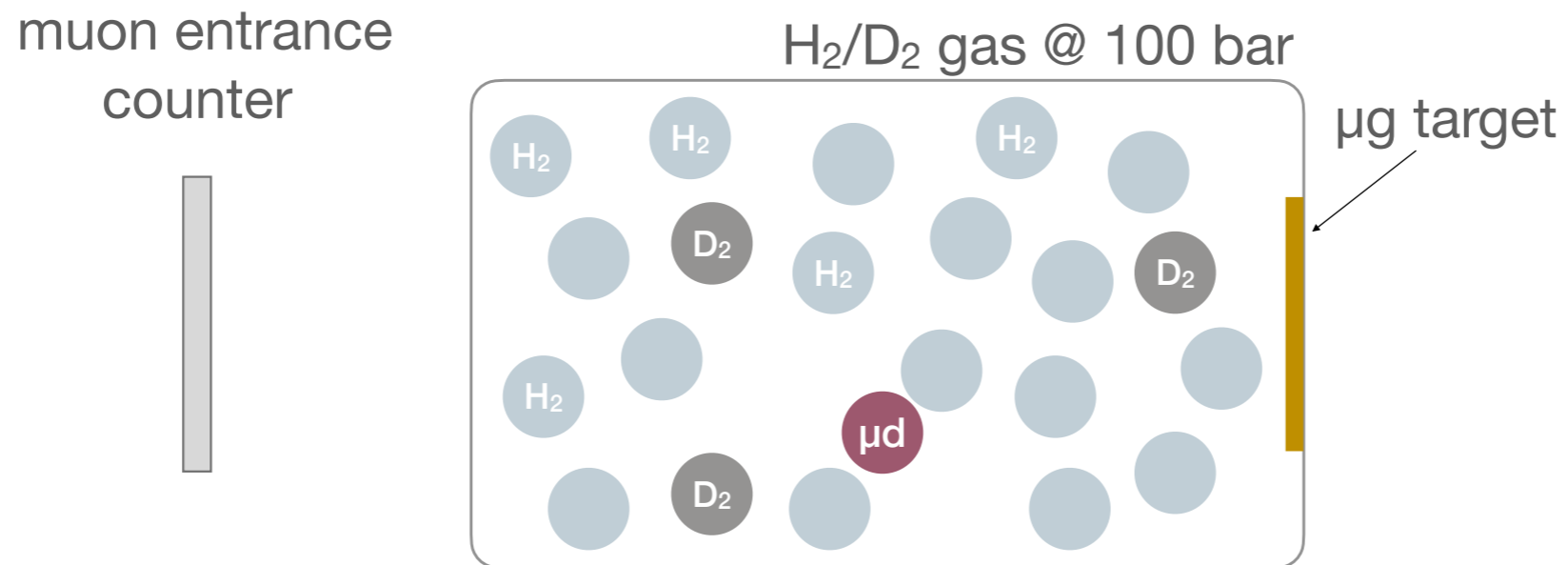
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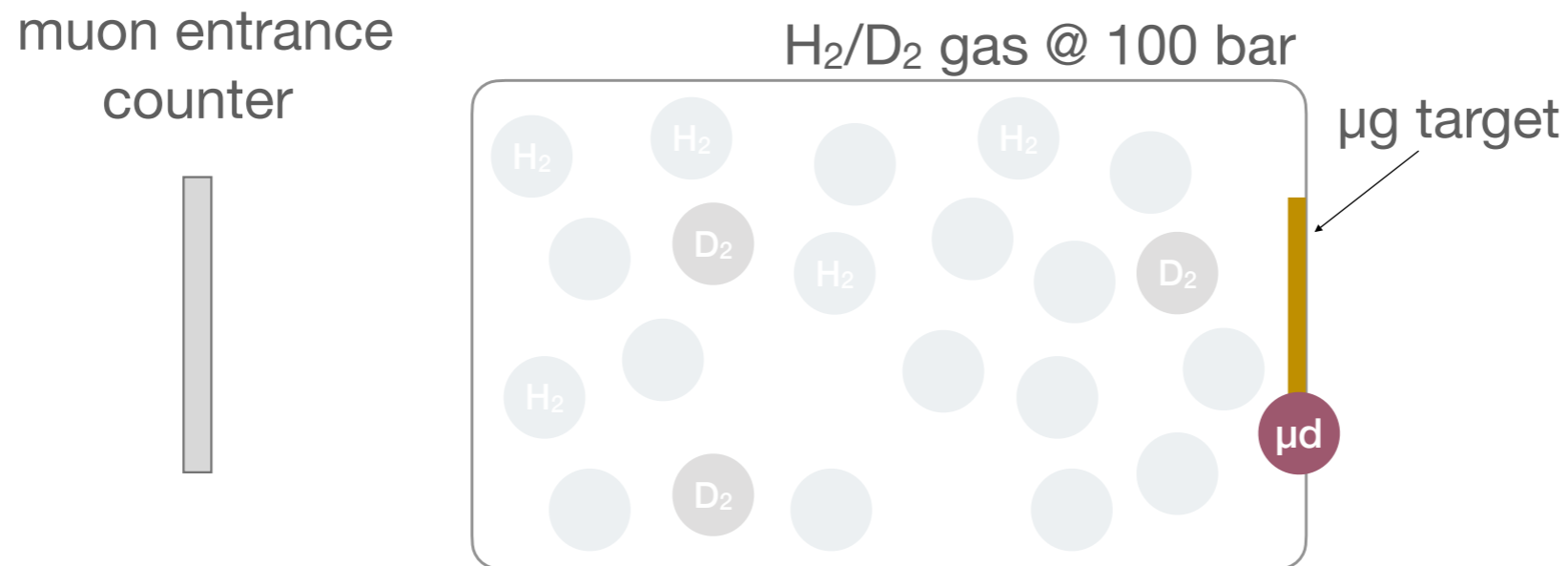
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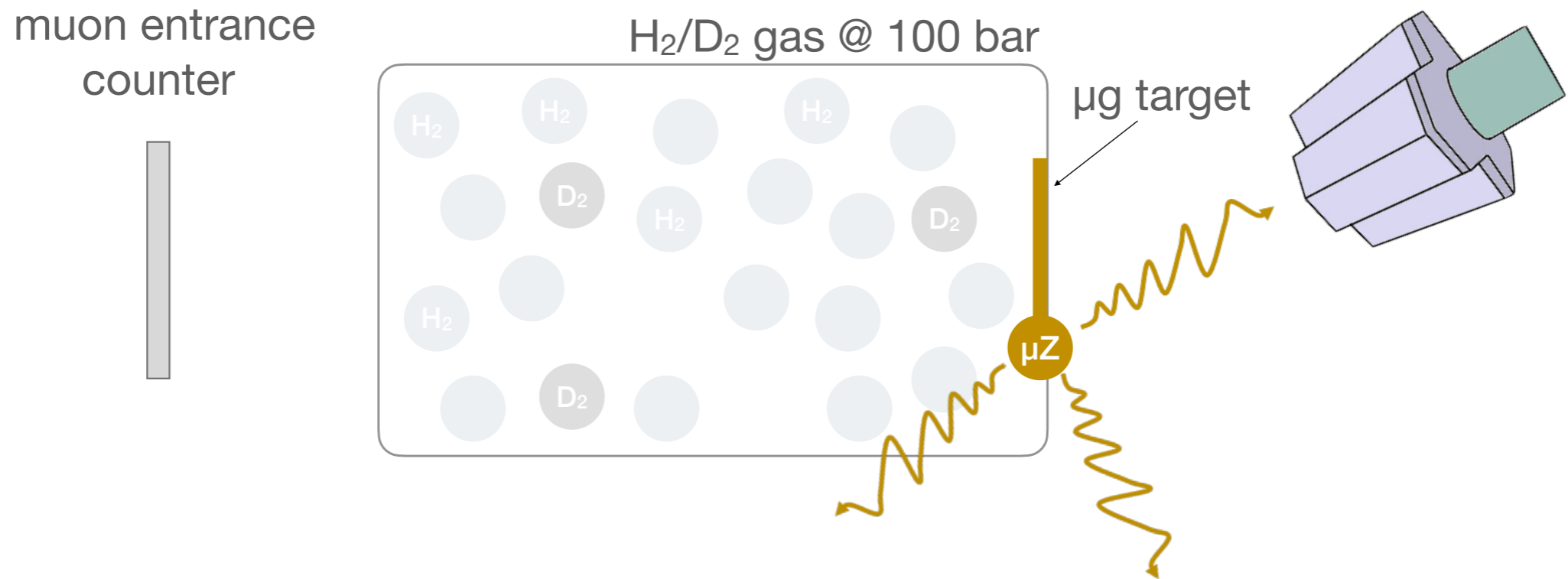
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3. μd moves almost freely in the H₂ gas

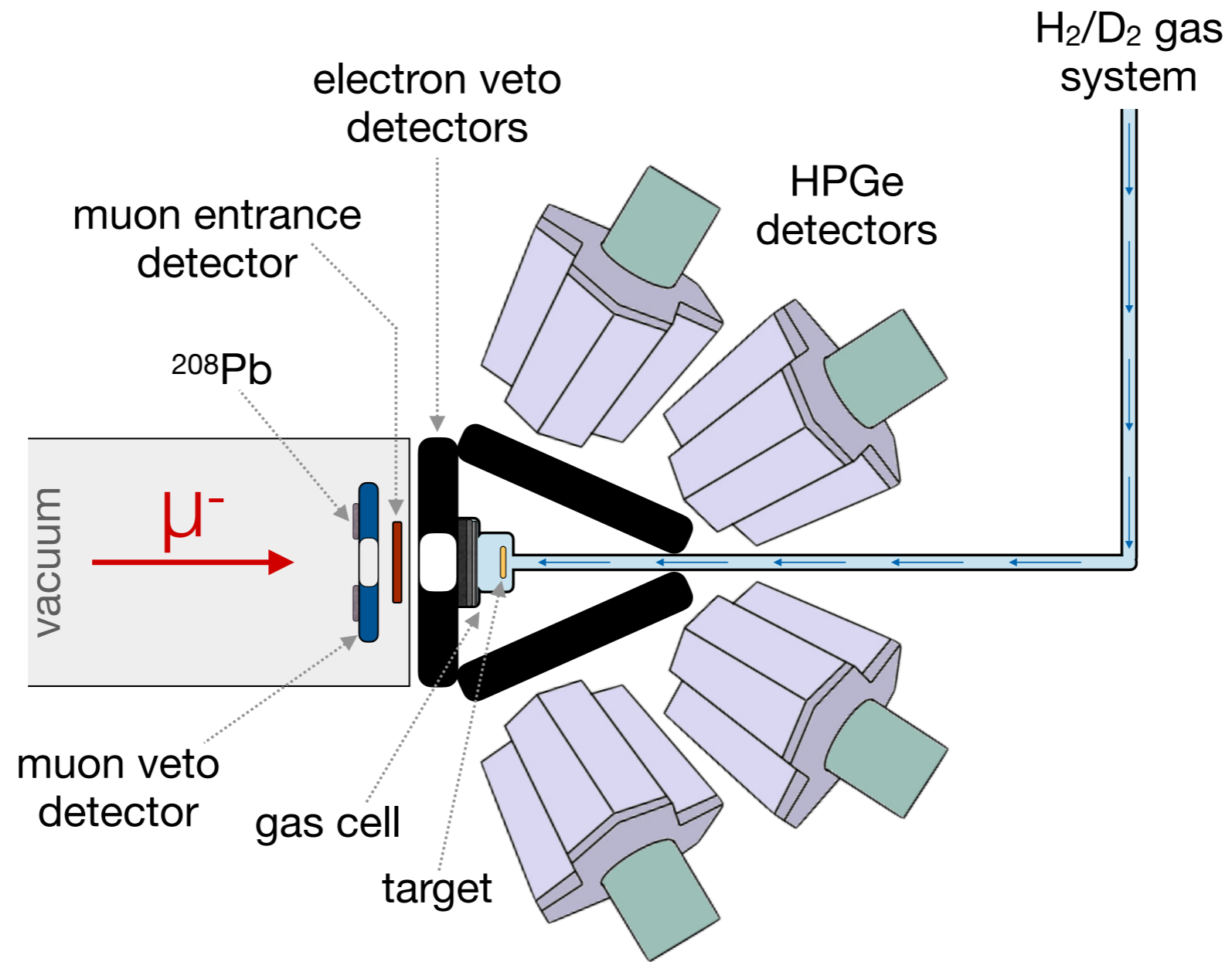
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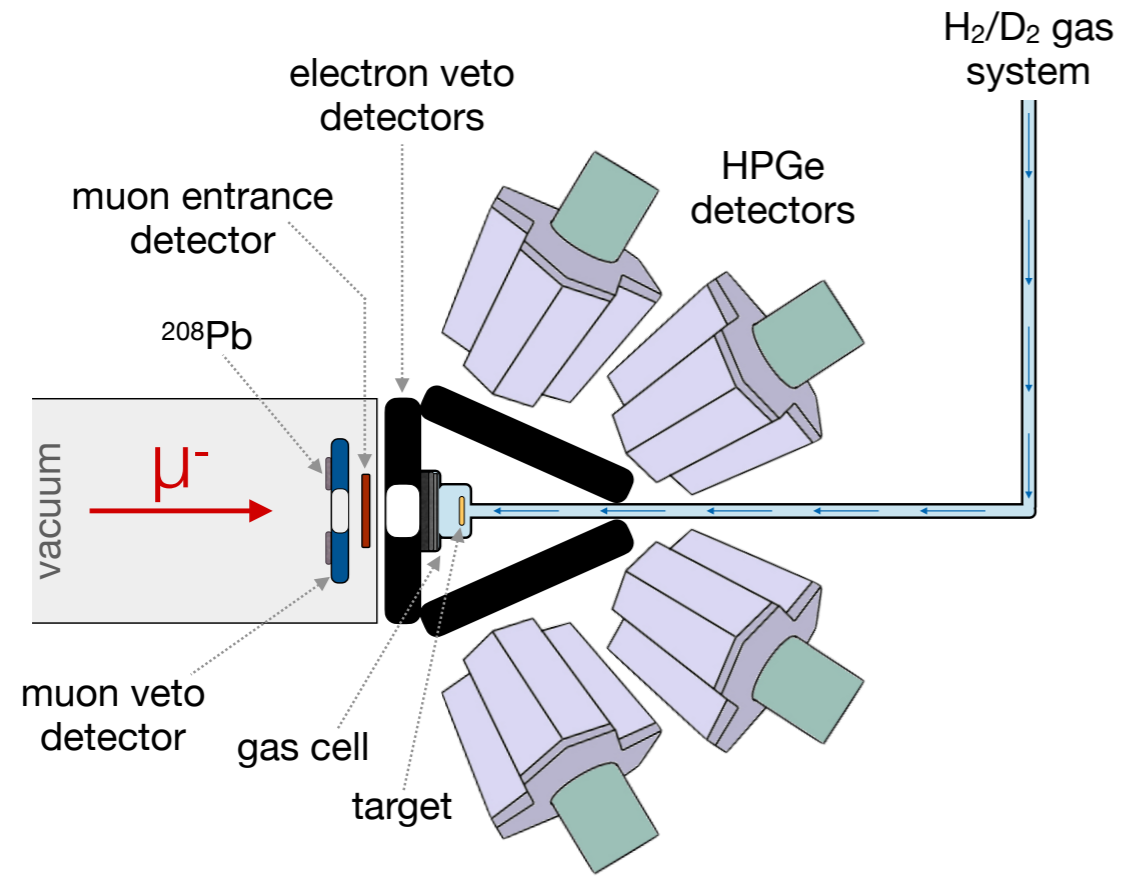


1. μ^- stops in 100 bar of H₂ + 0.25% D₂ & forms muonic hydrogen μp
2. transfer to deuterium $\mu p \rightarrow \mu d$
3. μd moves almost freely in the H₂ gas
4. transfer to high-Z element $\mu d \rightarrow \mu Z$ when hitting target & emission of x rays during the atomic cascade

muX detectors

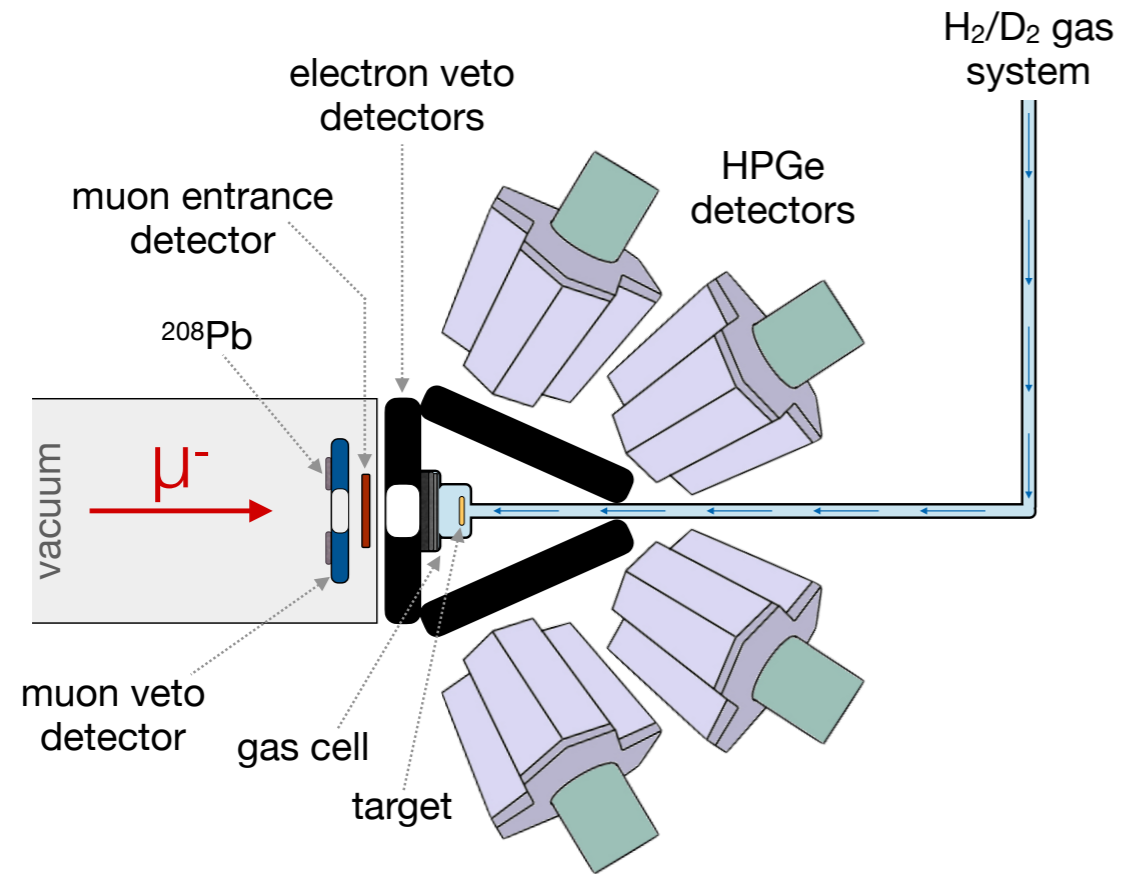


muX detectors

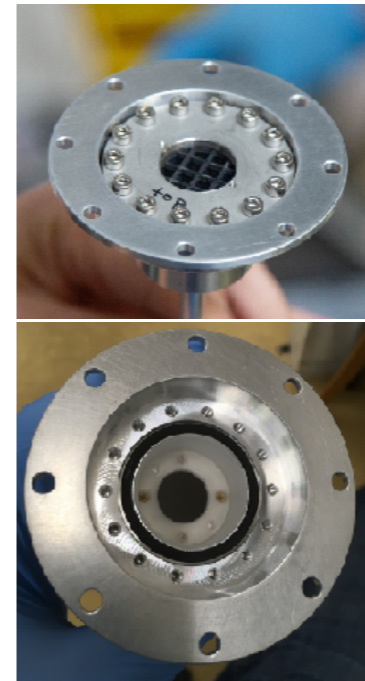


Schematics of detectors setup

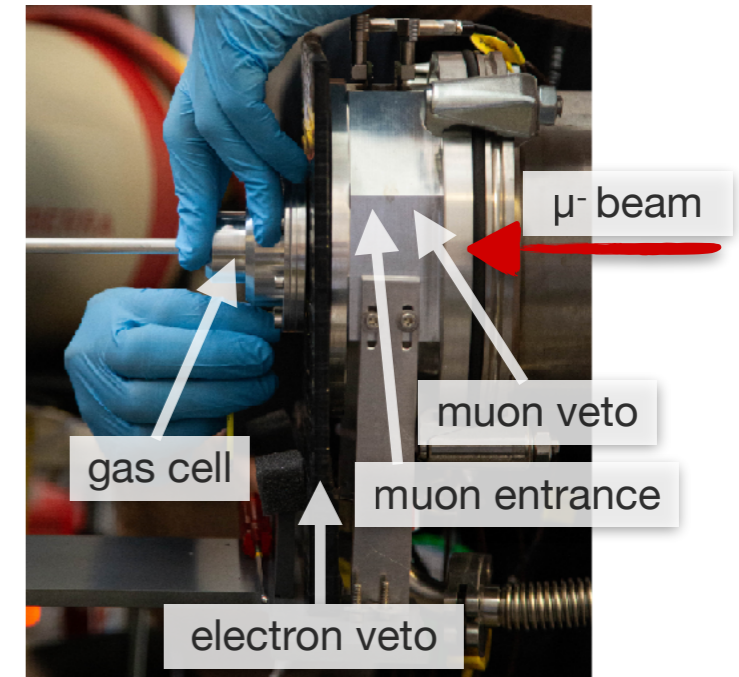
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Schematics of detectors setup

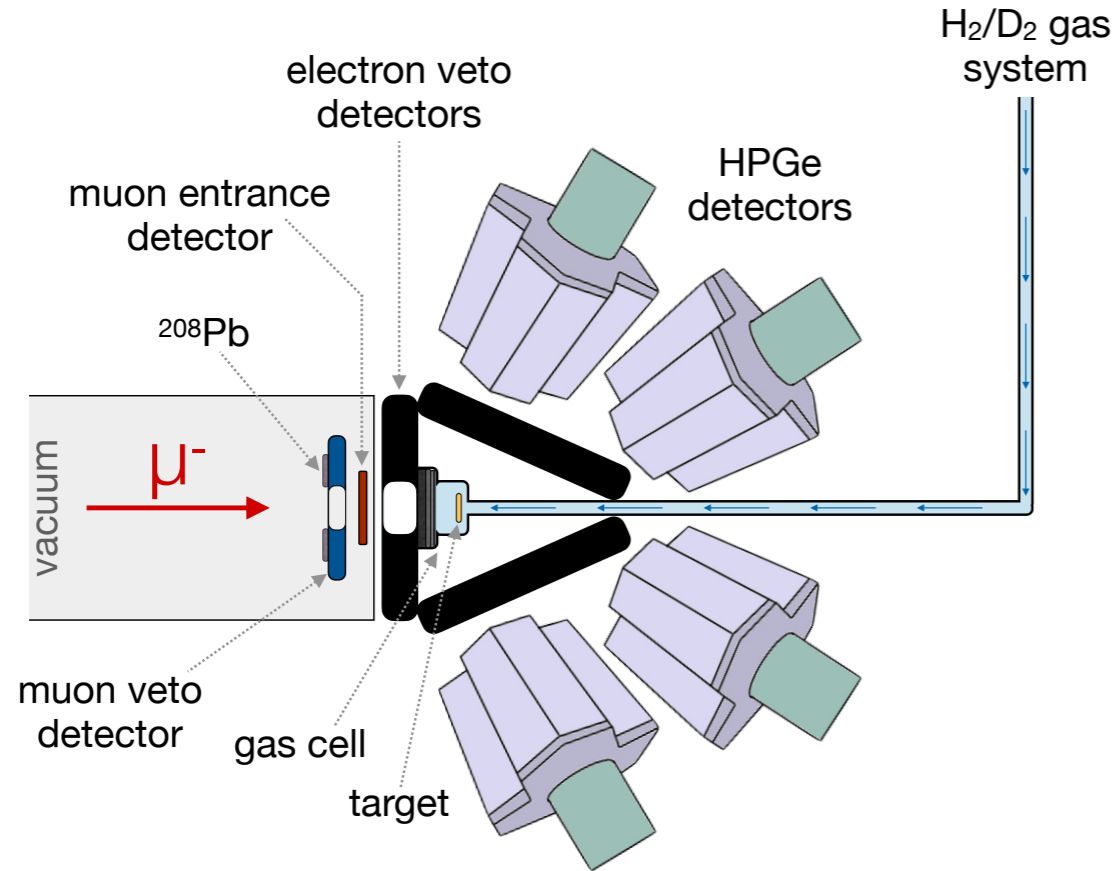


The gas cell

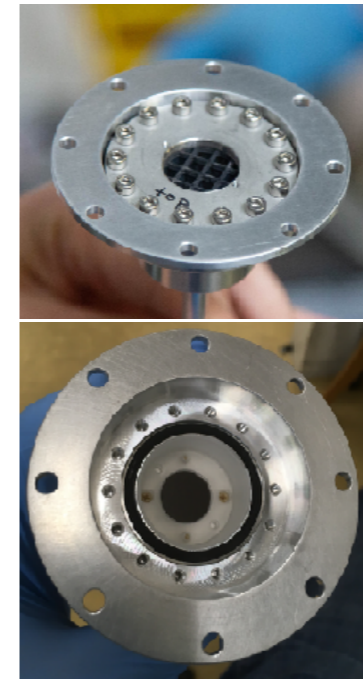


The muon and electron counters

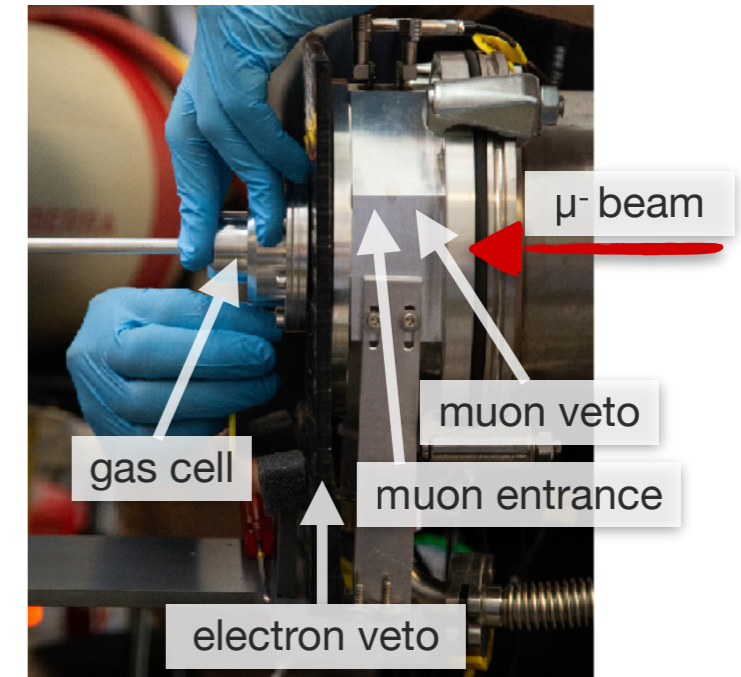
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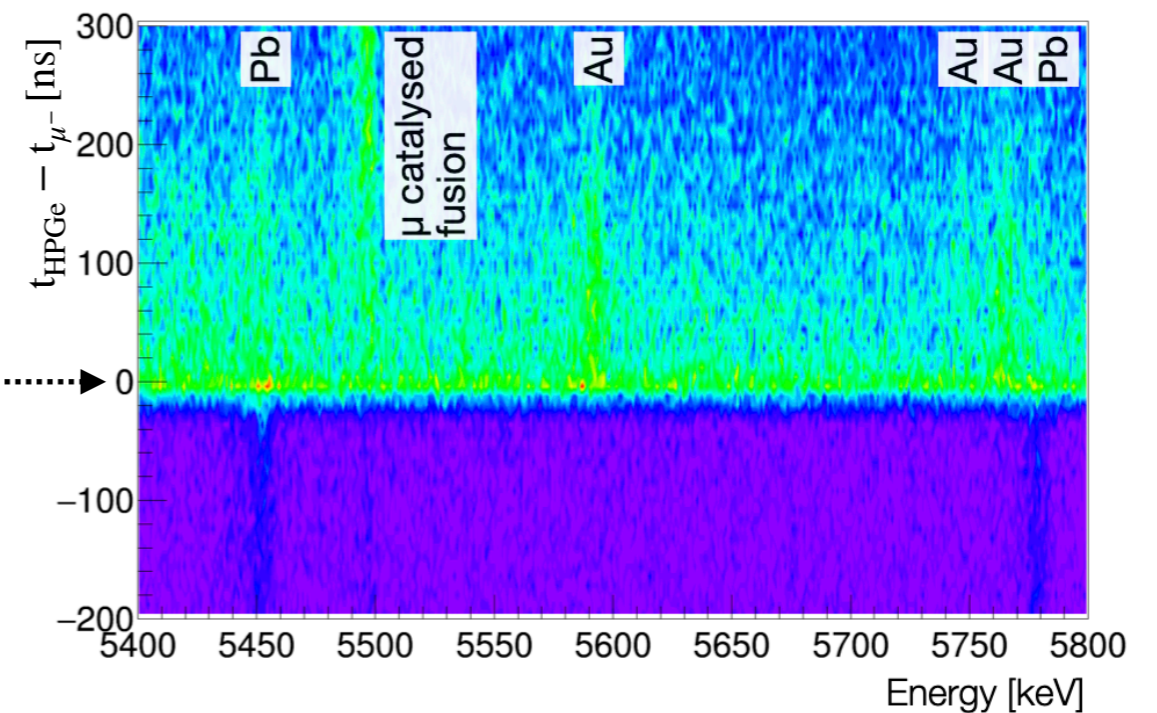


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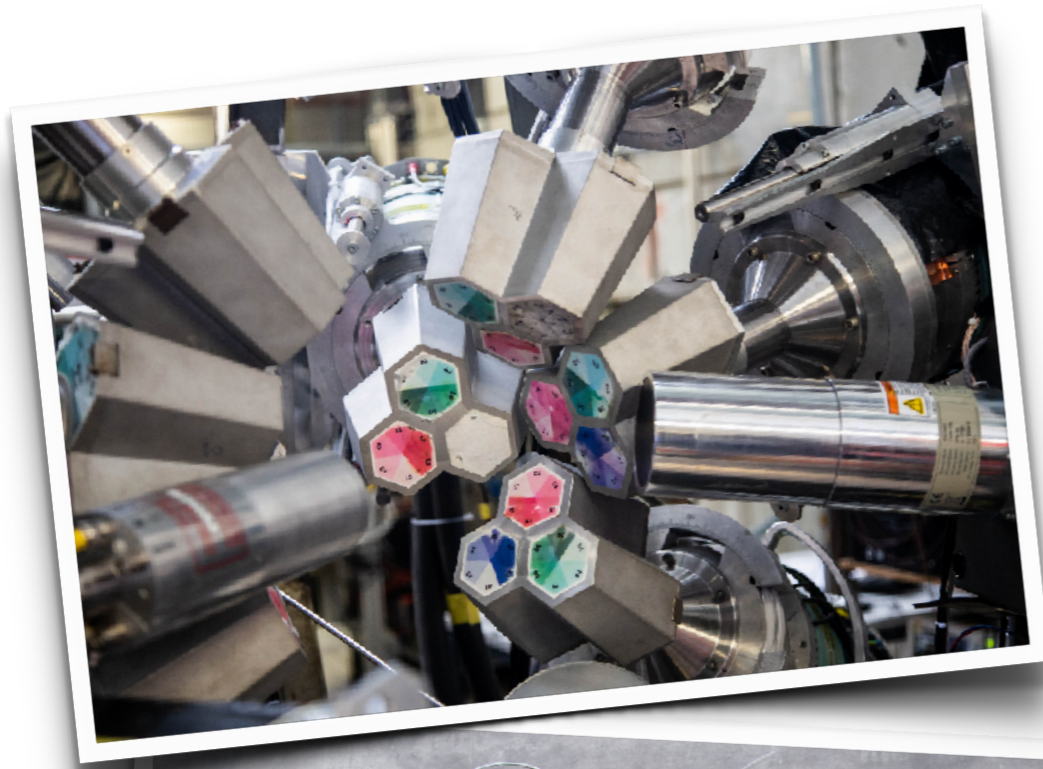
Demonstration of principal in a 5 μg gold target in 2017, 18.5 h of measurement

A. Adamczak *et al.*, *Eur. Phys. J. A* 59, 15 (2023)

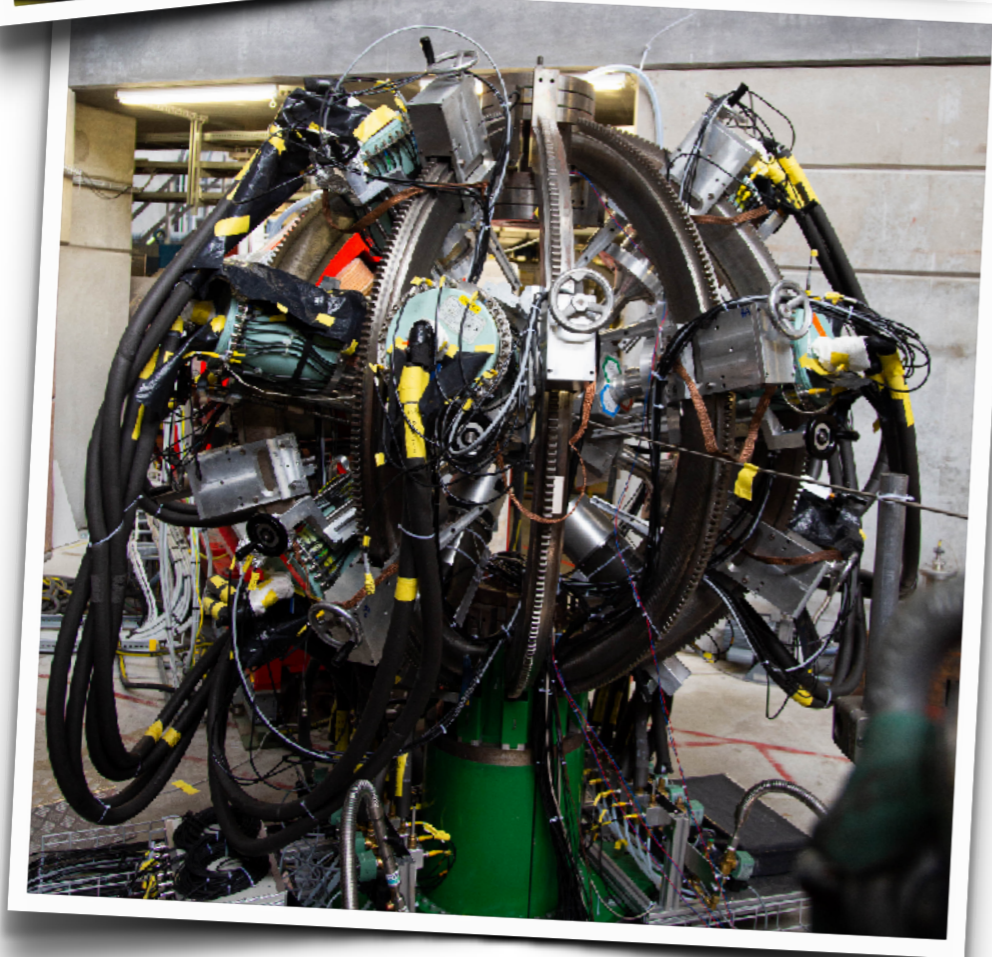
the time the muon enters the gas cell



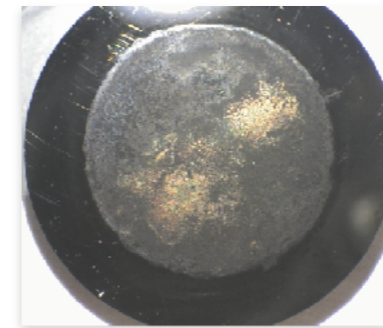
2019 measurement of ^{248}Cm and ^{226}Ra



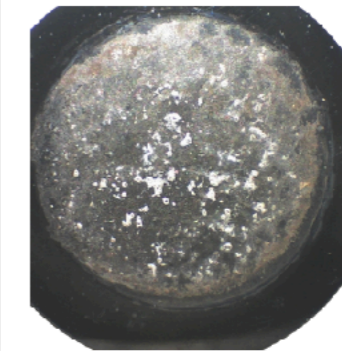
- 8 Miniball^[1] germanium clusters and 2 standalone germanium detectors making a total of 26 HPGe crystals were operating
- Radiation protection restrictions at PSI allow for 32.6 μg of ^{248}Cm and 5.5 μg of ^{226}Ra
- Targets were produced by the radiochemistry group of the University of Mainz



15 μg of ^{248}Cm



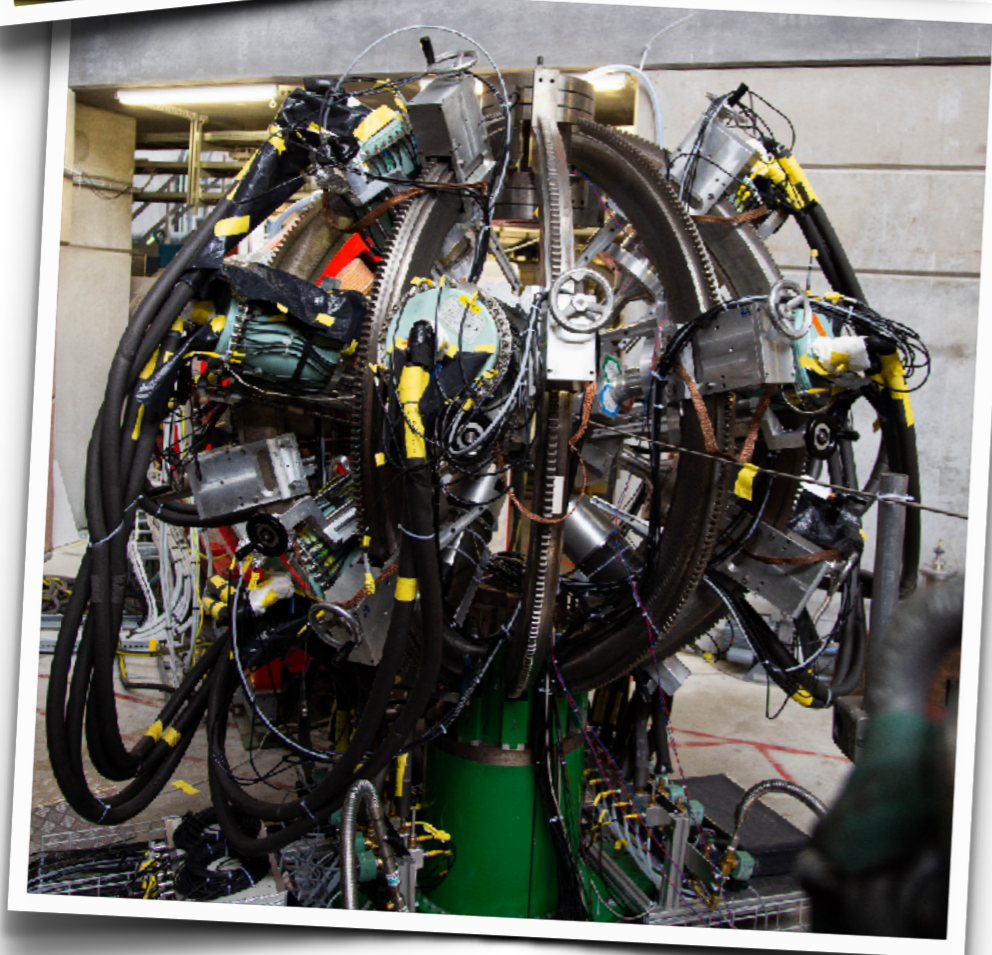
1.4 μg of ^{226}Ra



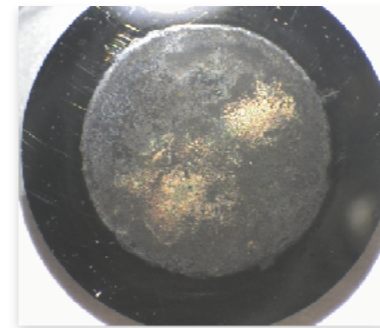
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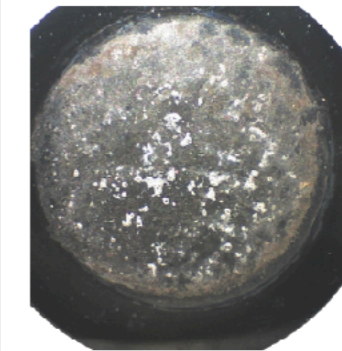
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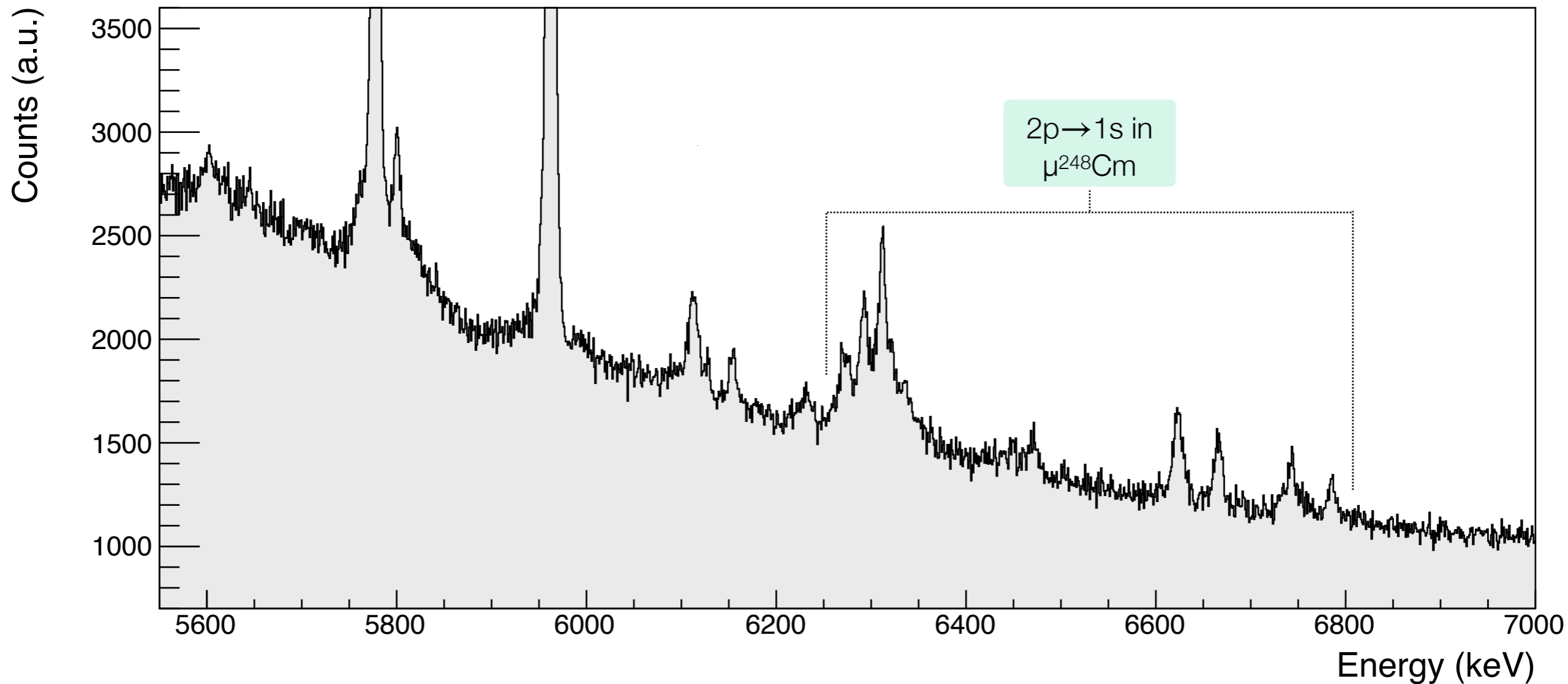


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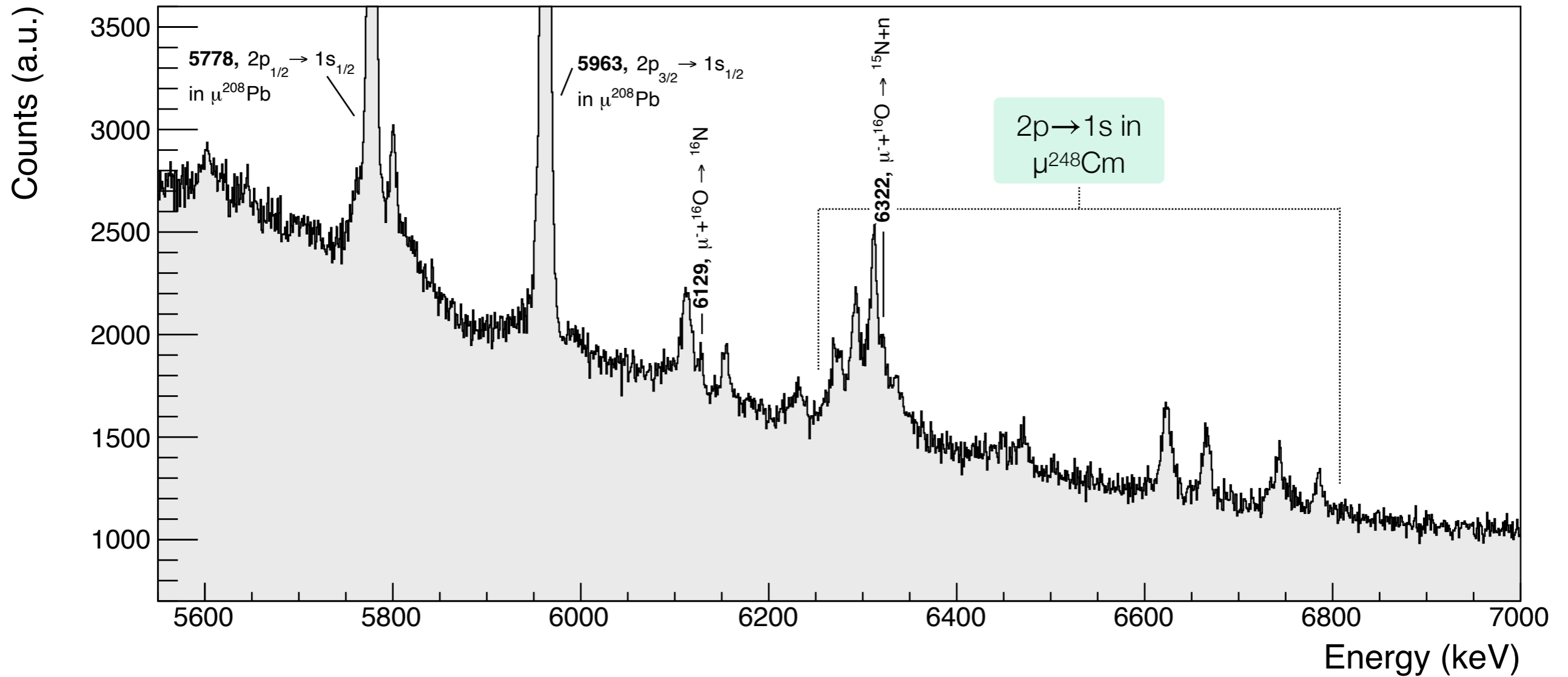


No clear $2p \rightarrow 1s$ hyperfine transitions are observed in $^{226}\text{Ra} \Rightarrow$ Radium needs to be remeasured!

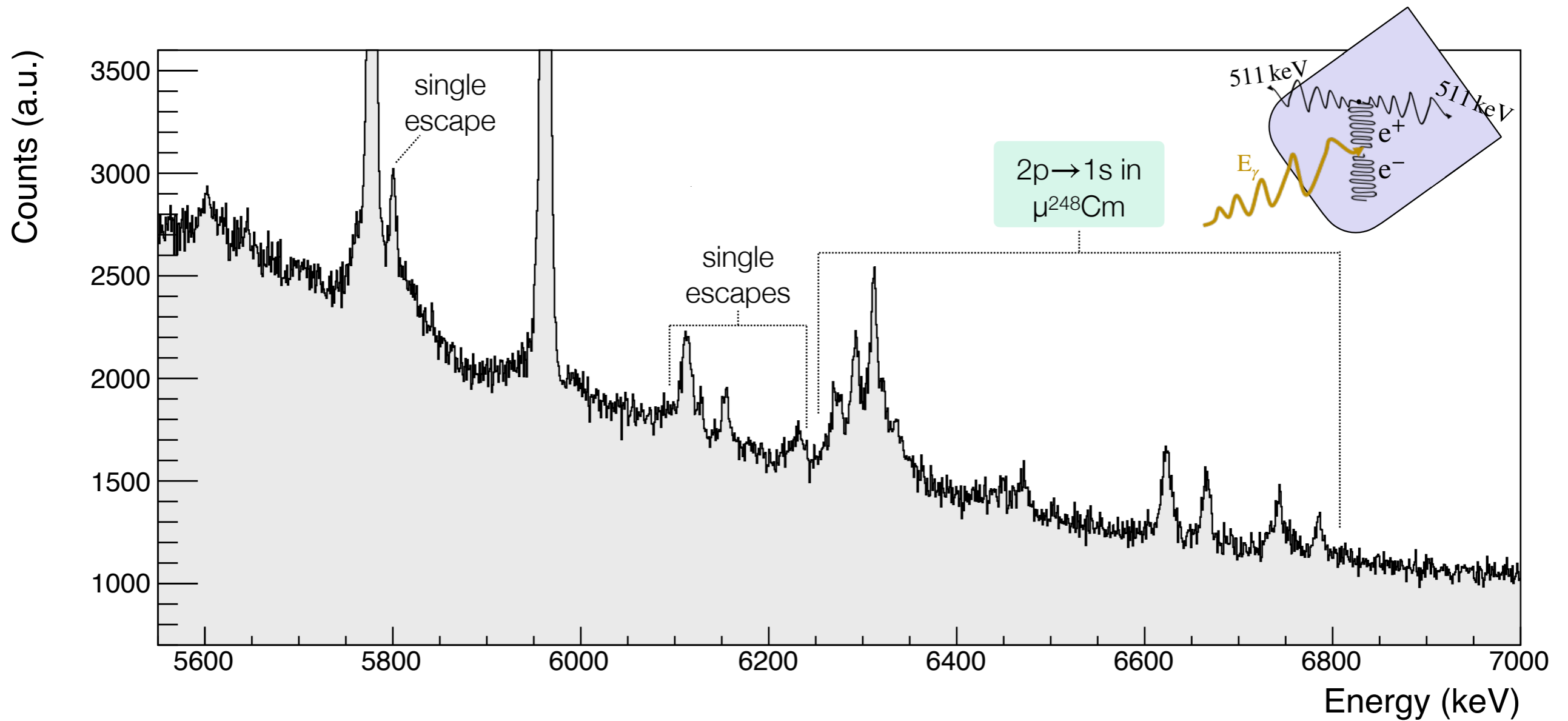
$2p \rightarrow 1s$ transitions in ^{248}Cm



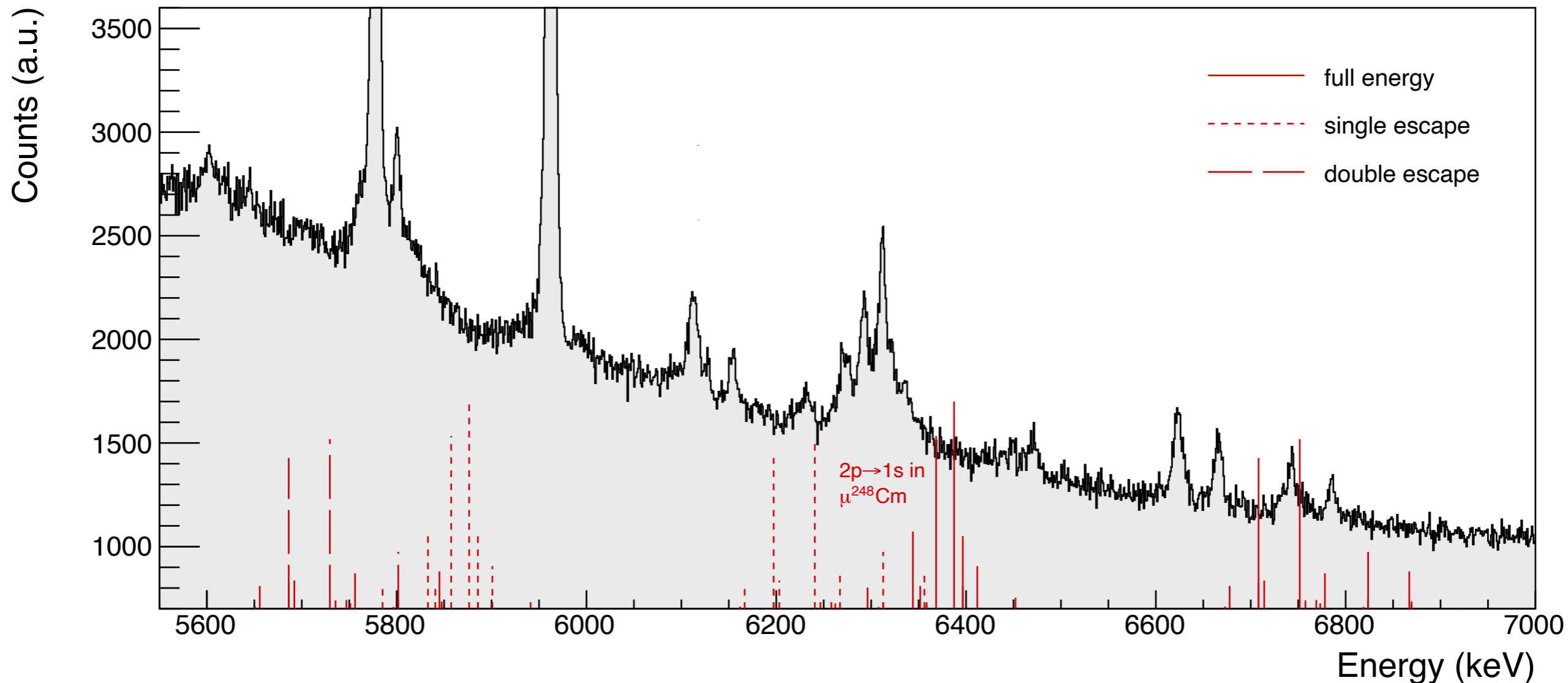
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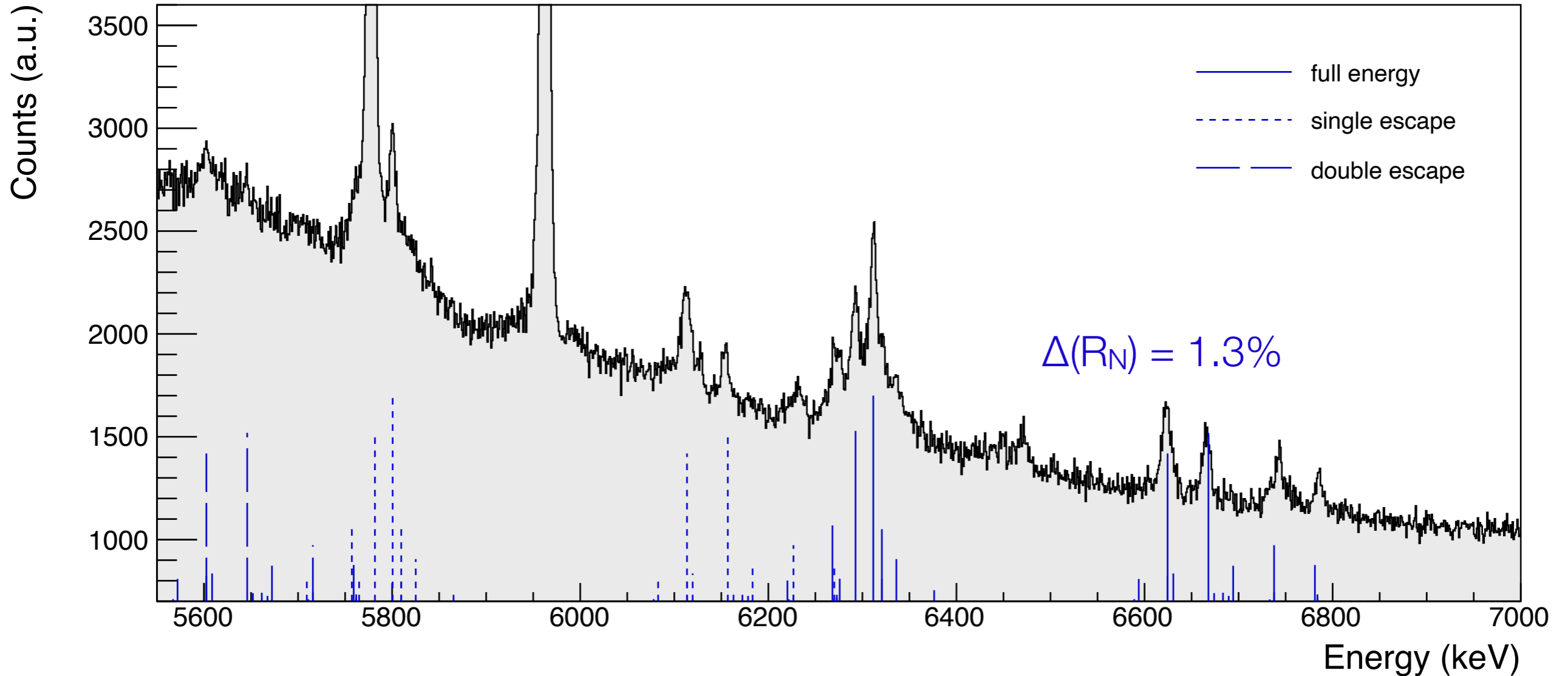
2p → 1s transitions in ^{248}Cm



Predictions shown for a charge radius estimated by interpolating the nuclear charge parameters of neighbouring nuclei^[1]

$$R_N = 5.8687 \text{ fm}$$

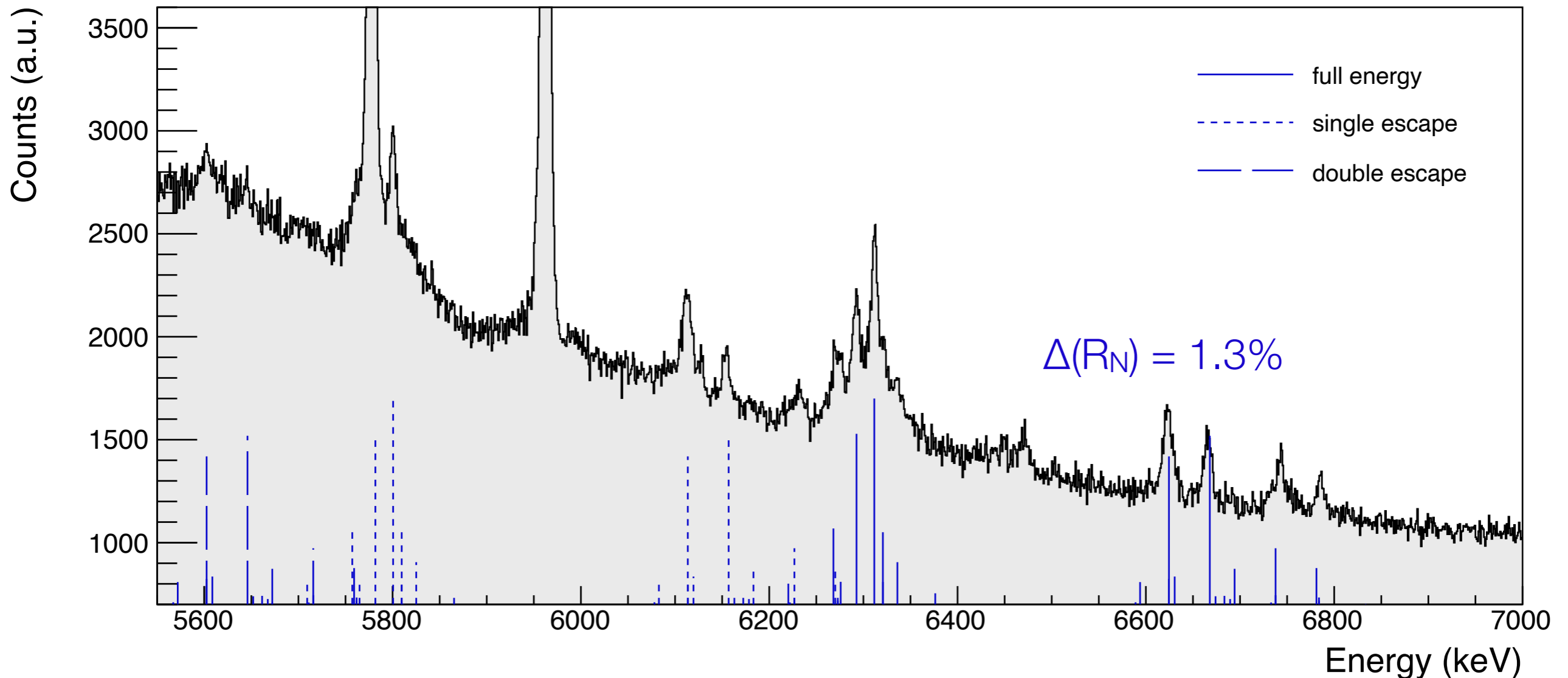
2p → 1s transitions in ^{248}Cm



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Predictions shown for a charge radius estimated by interpolating the nuclear charge parameters of neighbouring nuclei^[1]

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sensitivity on the quadrupole moment Q is also observed

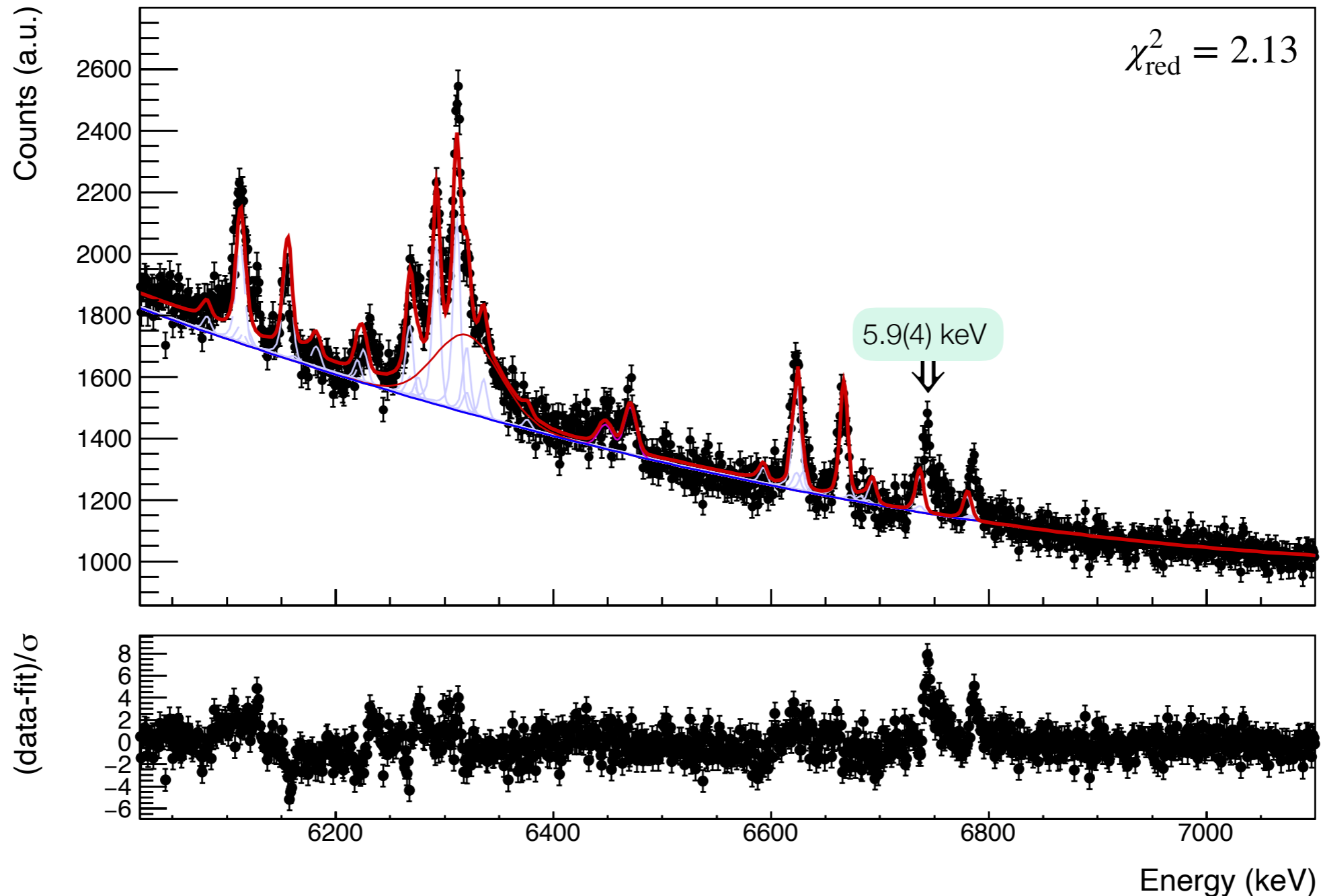
Fit of the $2p \rightarrow 1s$ transitions in ^{248}Cm

Fit results (only statistical):

$$R = 5.94518(9) \text{ fm}$$

$$Q = 12.003(8) \text{ b}$$

+ full study of systematic uncertainties



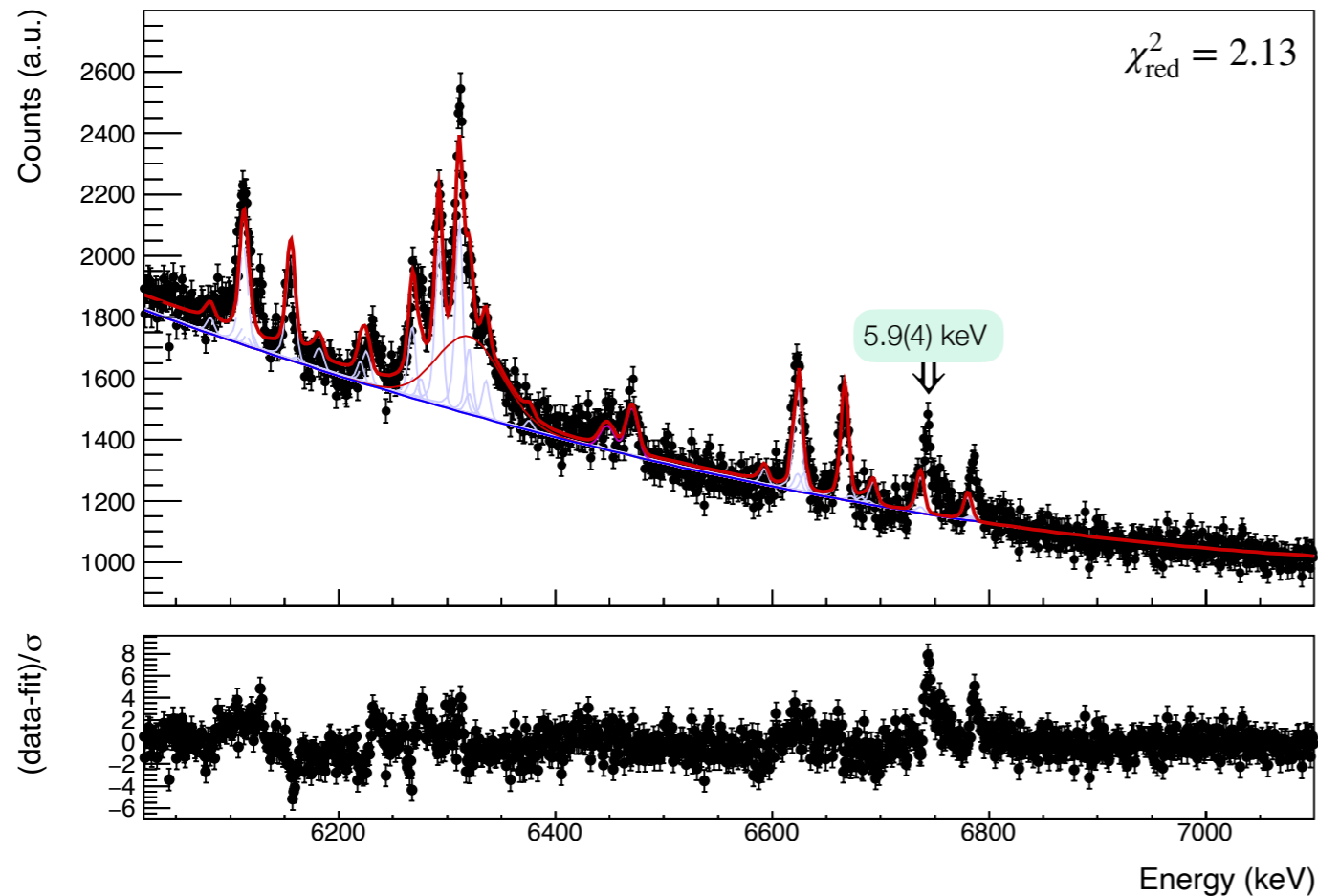
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Results on the nuclear charge radius and quadrupole moment:

$$R = 5.9455(1)_{\text{stat}}(117)_{\text{sys}} \text{ fm}$$

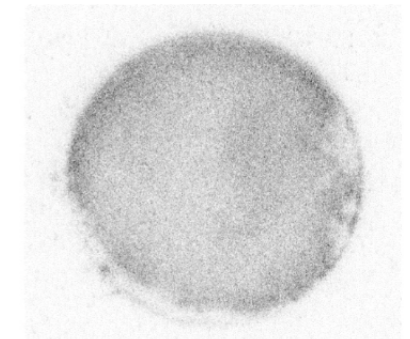
$$Q = 12.003(8)_{\text{stat}}(361)_{\text{sys}} \text{ b}$$

Outlook: Further revision of the theory to resolve the discrepancies between the fitted and measured hyperfine transitions

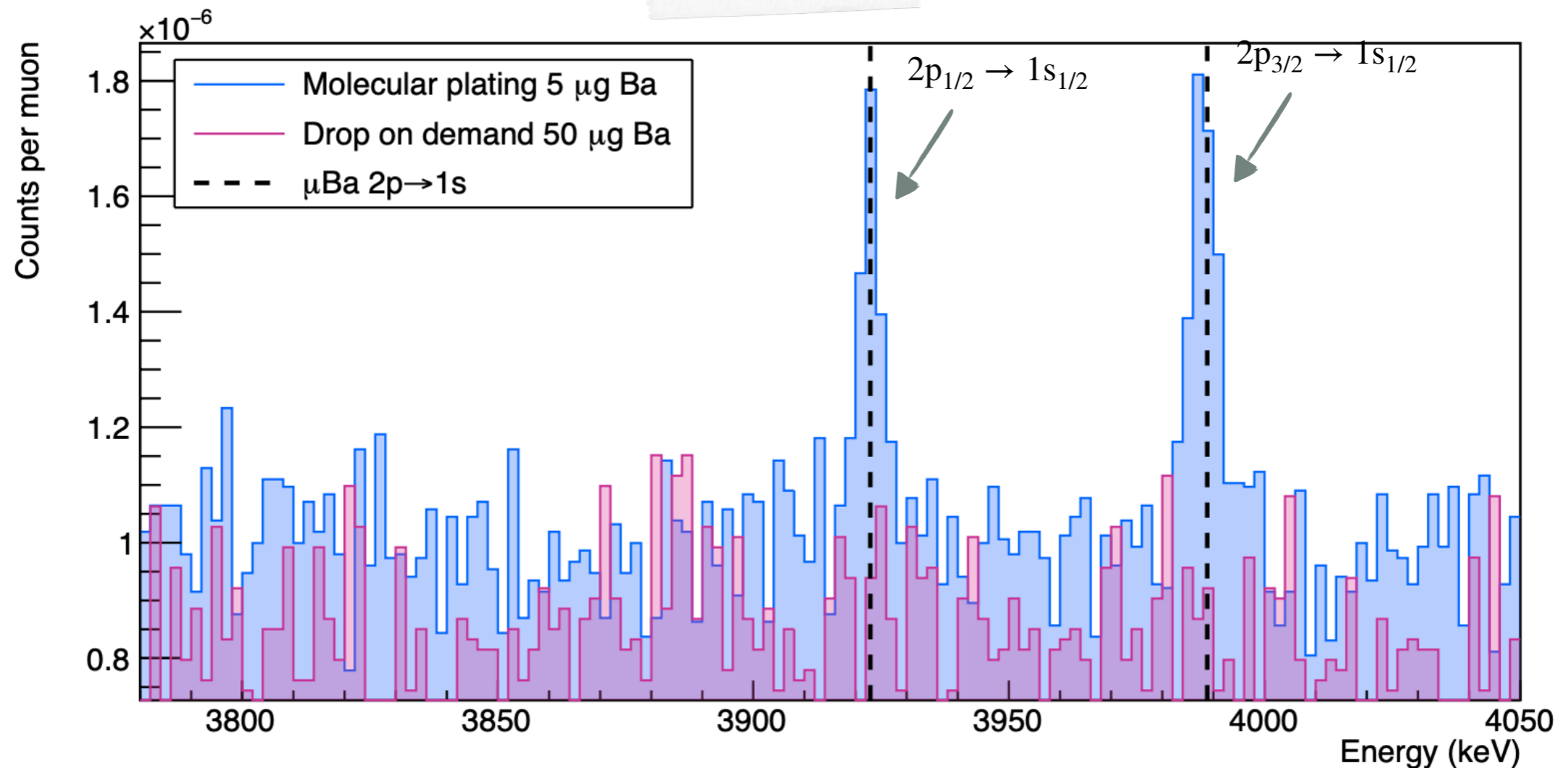
2022 test measurement with ^{nat}Ba

Towards the future measurement of ^{226}Ra

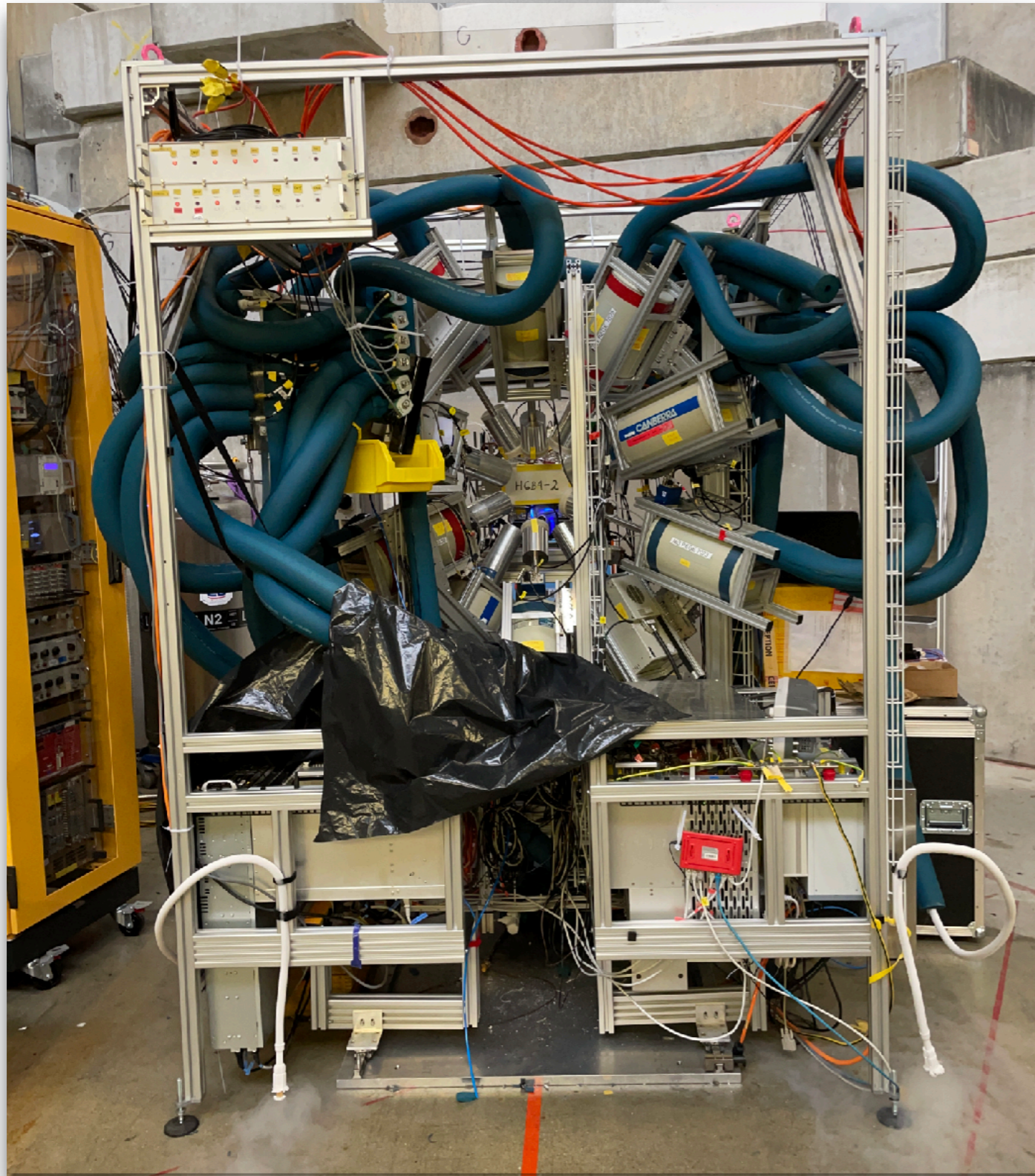
- Measured of barium deposition targets
- Compared two methods: drop-on-demand (DoD) & molecular plating (MP)



^{nat}Ba target with MP



What's next: remeasure ^{226}Ra this September!



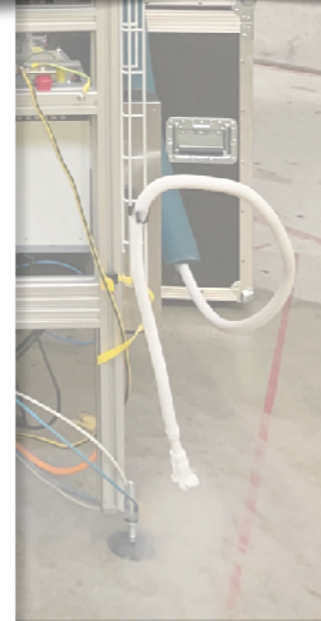
- Measure a ^{226}Ra target produced by the MP & DoD method
- Using 20 HPGe crystals in total!



What's next: remeasure ^{226}Ra this September!



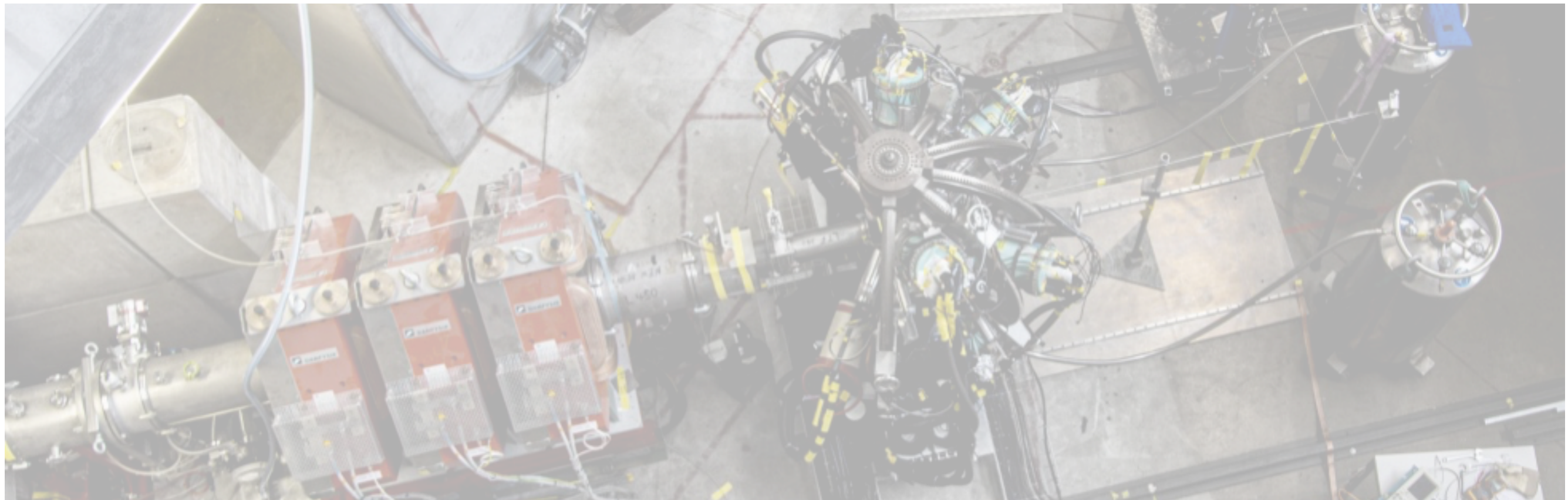
e a 226
& DoD
O HPC



Thank you for your attention!



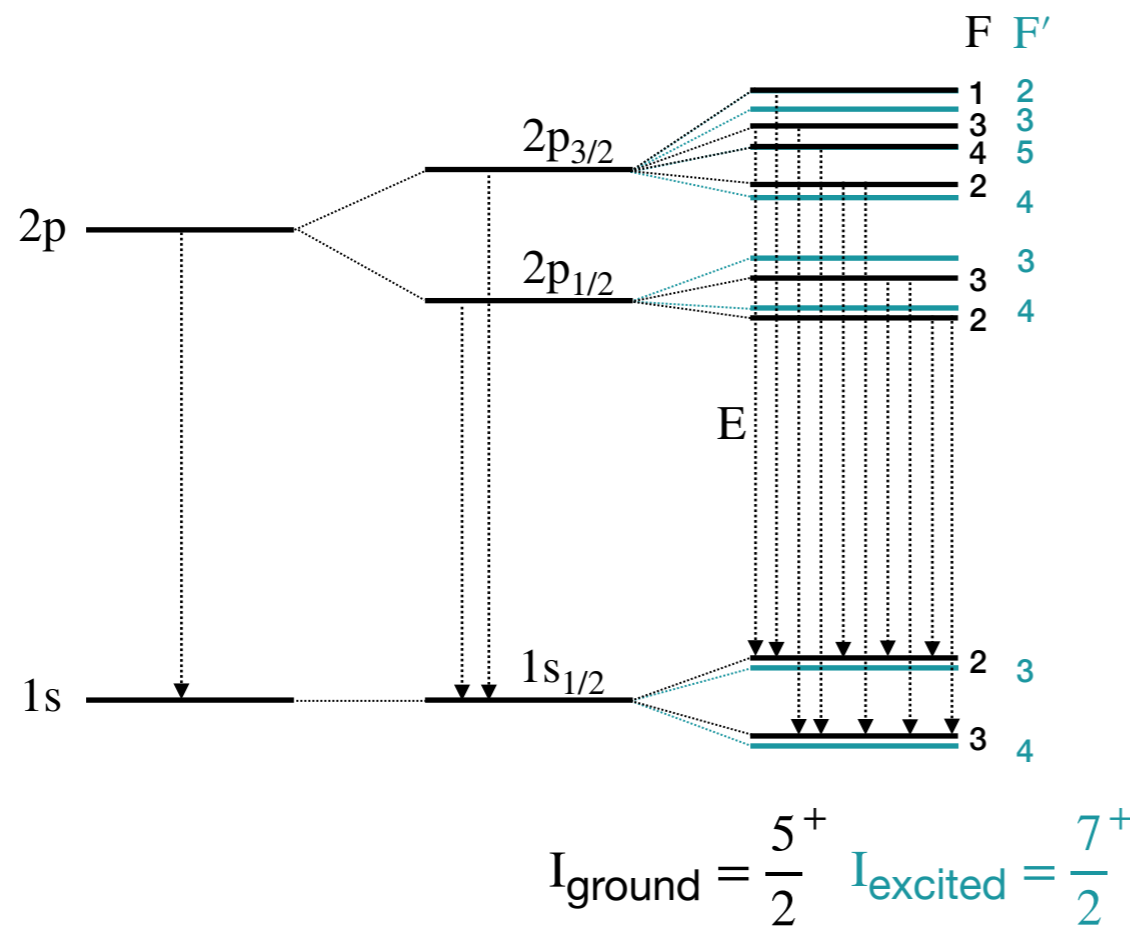
Backup Slides



Fine & Hyperfine splitting in muonic atoms

Fine splitting (FS): $\vec{J} = \vec{L} + \vec{s}$

Static hyperfine splitting (HFS): $\vec{F} = \vec{I} + \vec{J}$

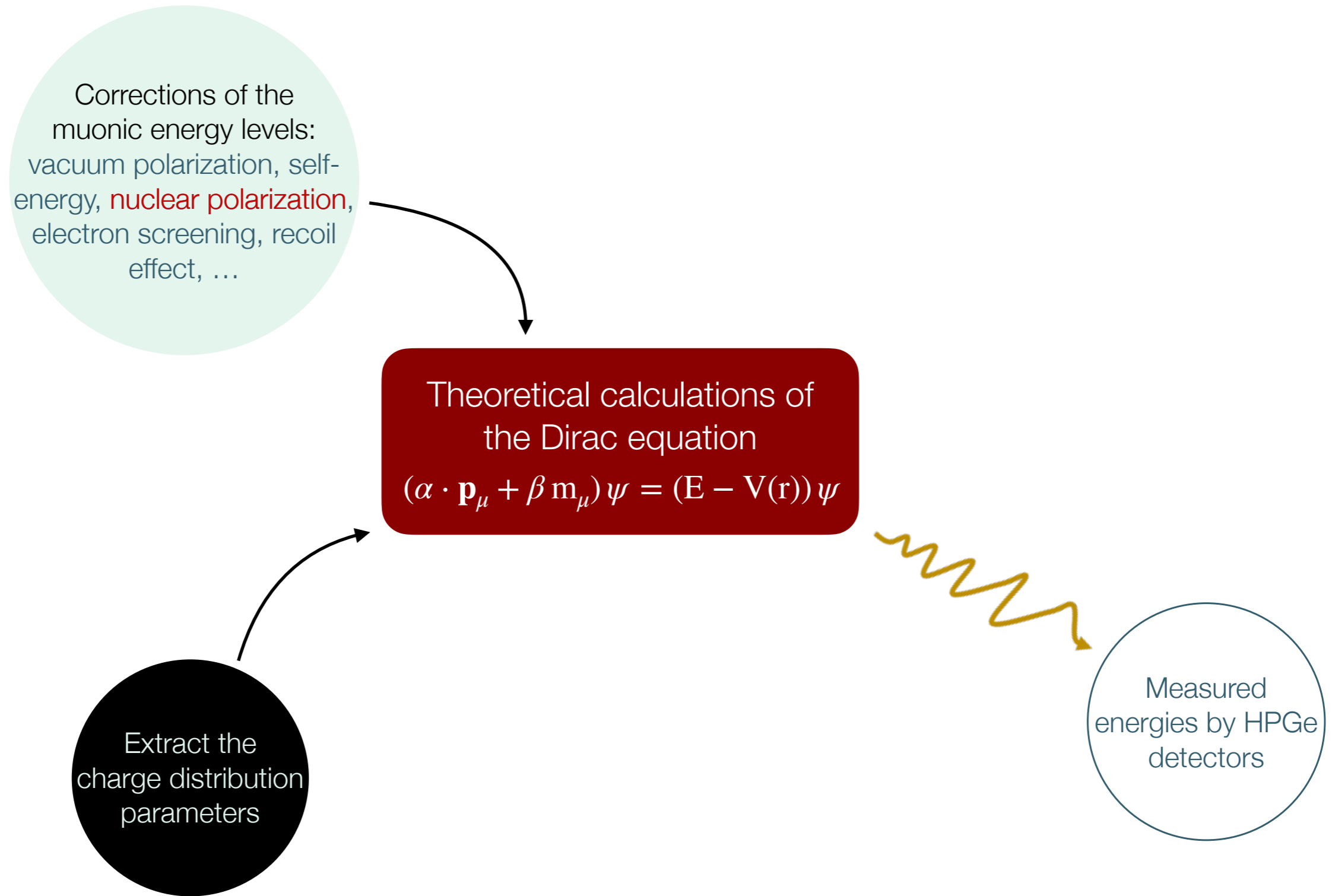


- Energy shift of hyperfine states due to the electric quadrupole (E2) and magnetic dipole (M1) interaction

Dynamic hyperfine splitting

- The hyperfine levels from ground and excited nuclear states are mixed due to the high energy of muonic transitions
- HFS also observed in even-even nuclei with zero spin in the ground state

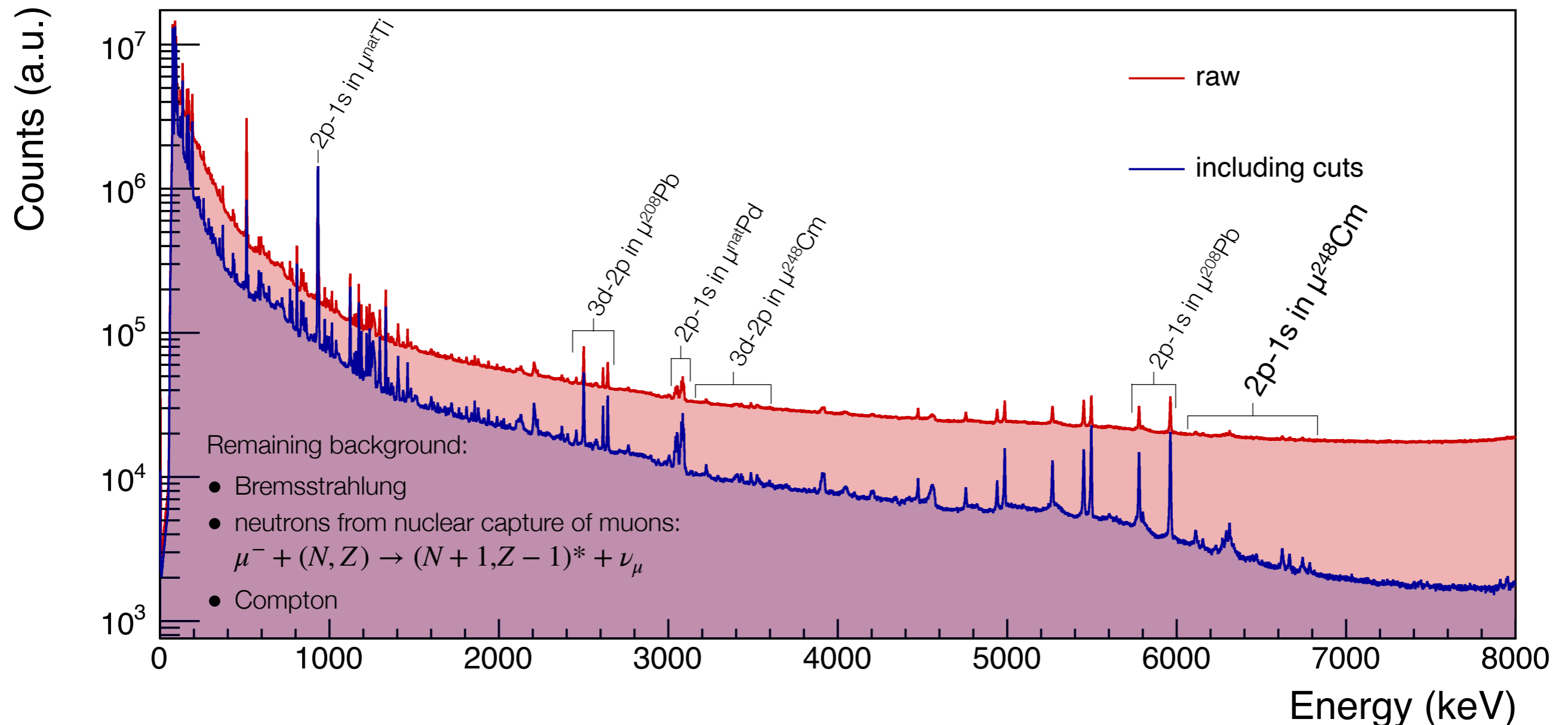
How to extract nuclear charge parameters from measured muonic energies



How to get to good S/B ratio at high energies?

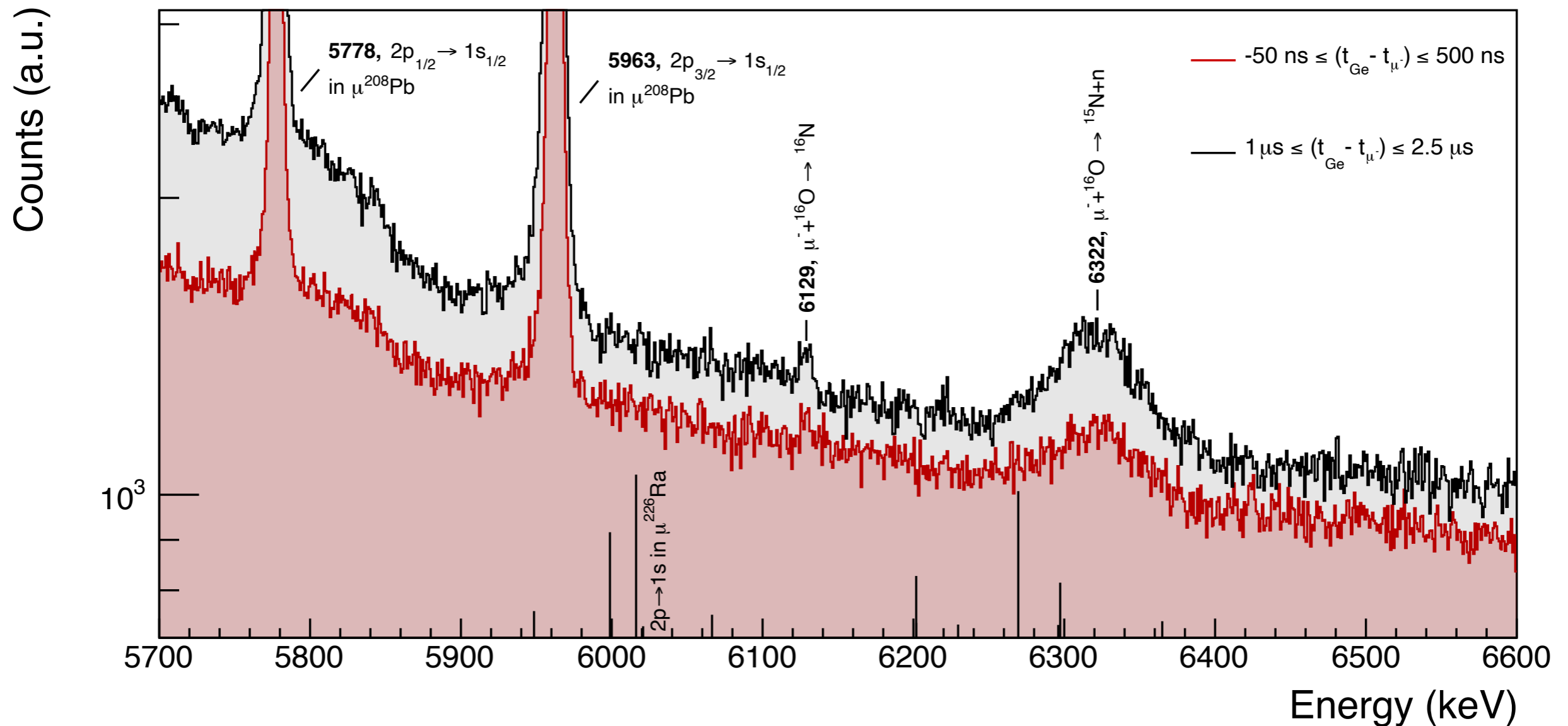
Cuts included:

- Electron veto & HPGe detectors' anti-coincidences
- Clustering of events using the 3-crystal Miniball geometry
- Correction of gain drifts before energy calibration
- Baseline correction to correct for low-energy tails of the peaks



2p → 1s transitions in ^{226}Ra ?

- No clear 2p → 1s hyperfine transitions are observed in ^{226}Ra
- Radium needs to be remeasured

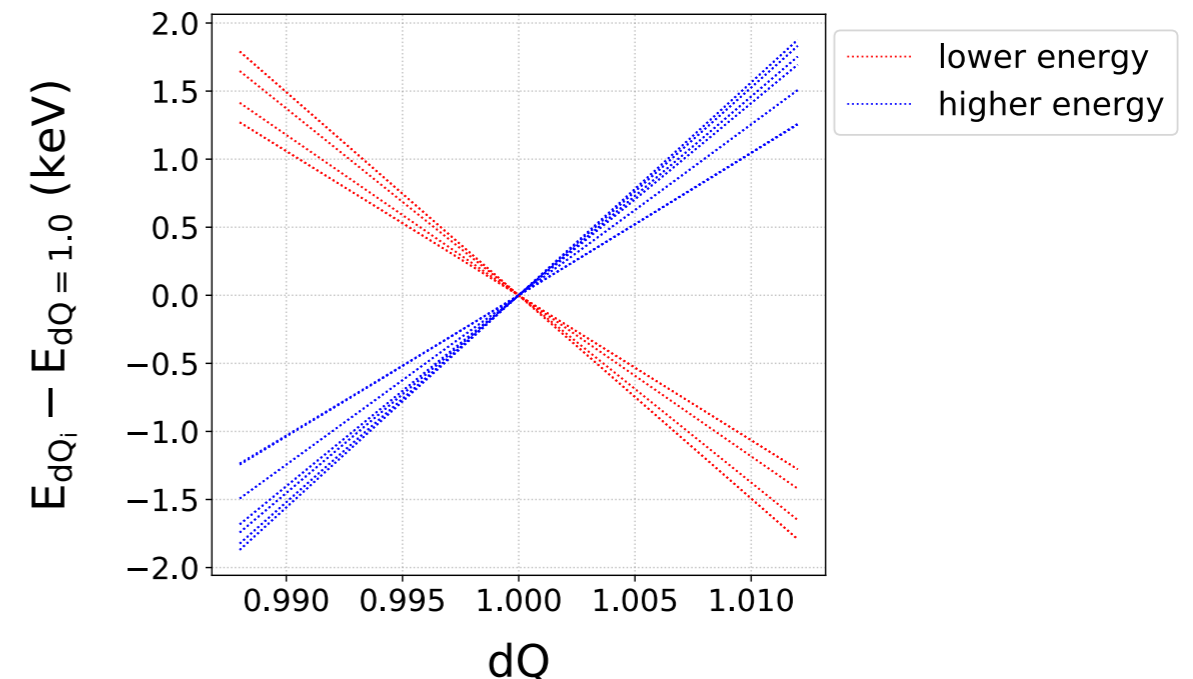
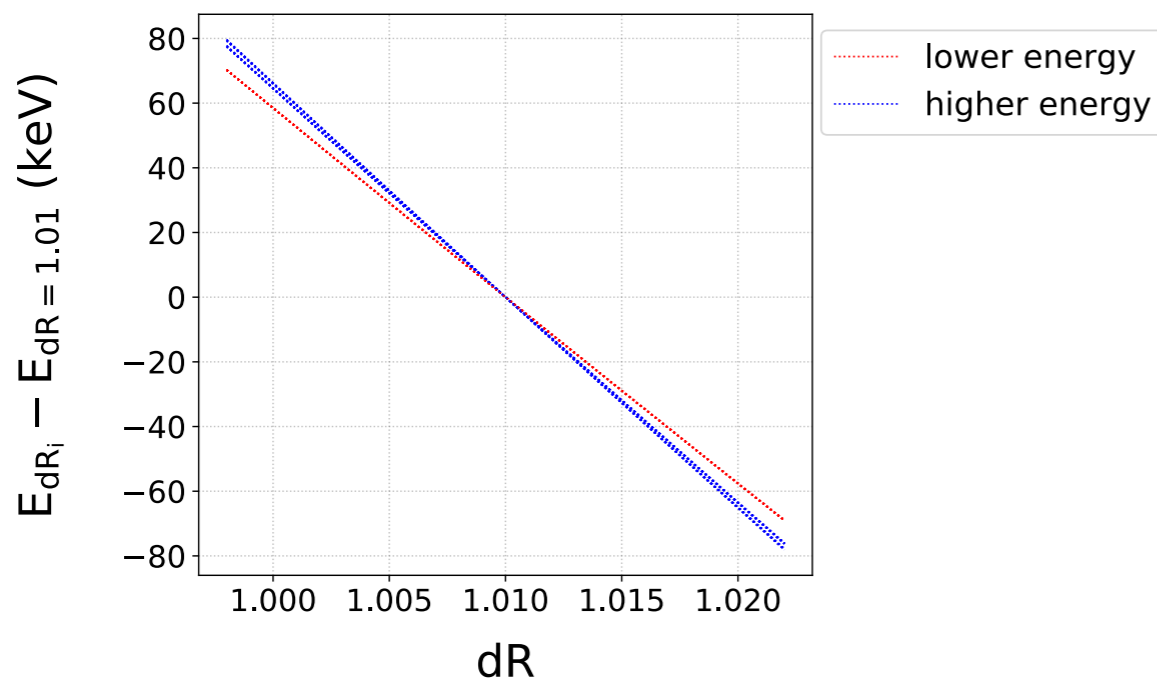
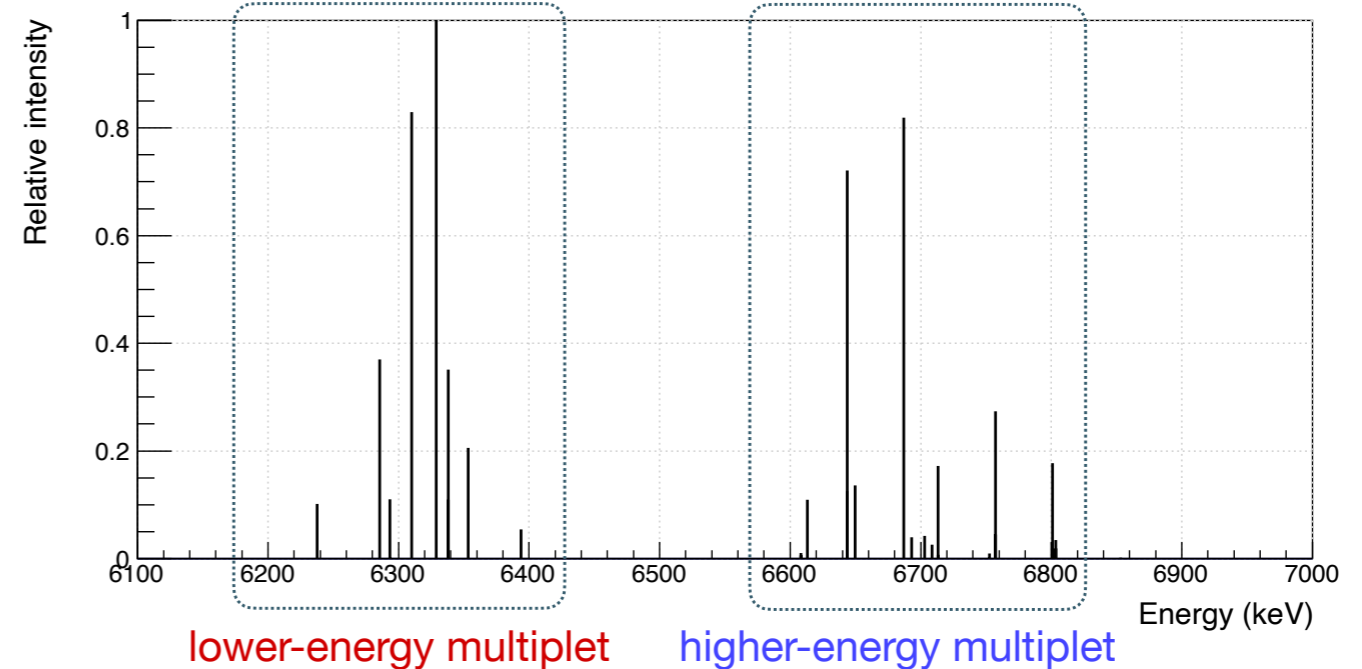


Energy dependence on the charge radius and quadrupole moment

$$dR = \frac{R}{R_N}, \text{ where } R_N = 5.8687 \text{ fm}$$

$$dQ = \frac{Q}{Q_N}, \text{ where } Q_N = 12.04 \text{ b}$$

Theoretical calculations are performed by
N. Oreshkina & I. Valuev, MPIK, Heidelberg



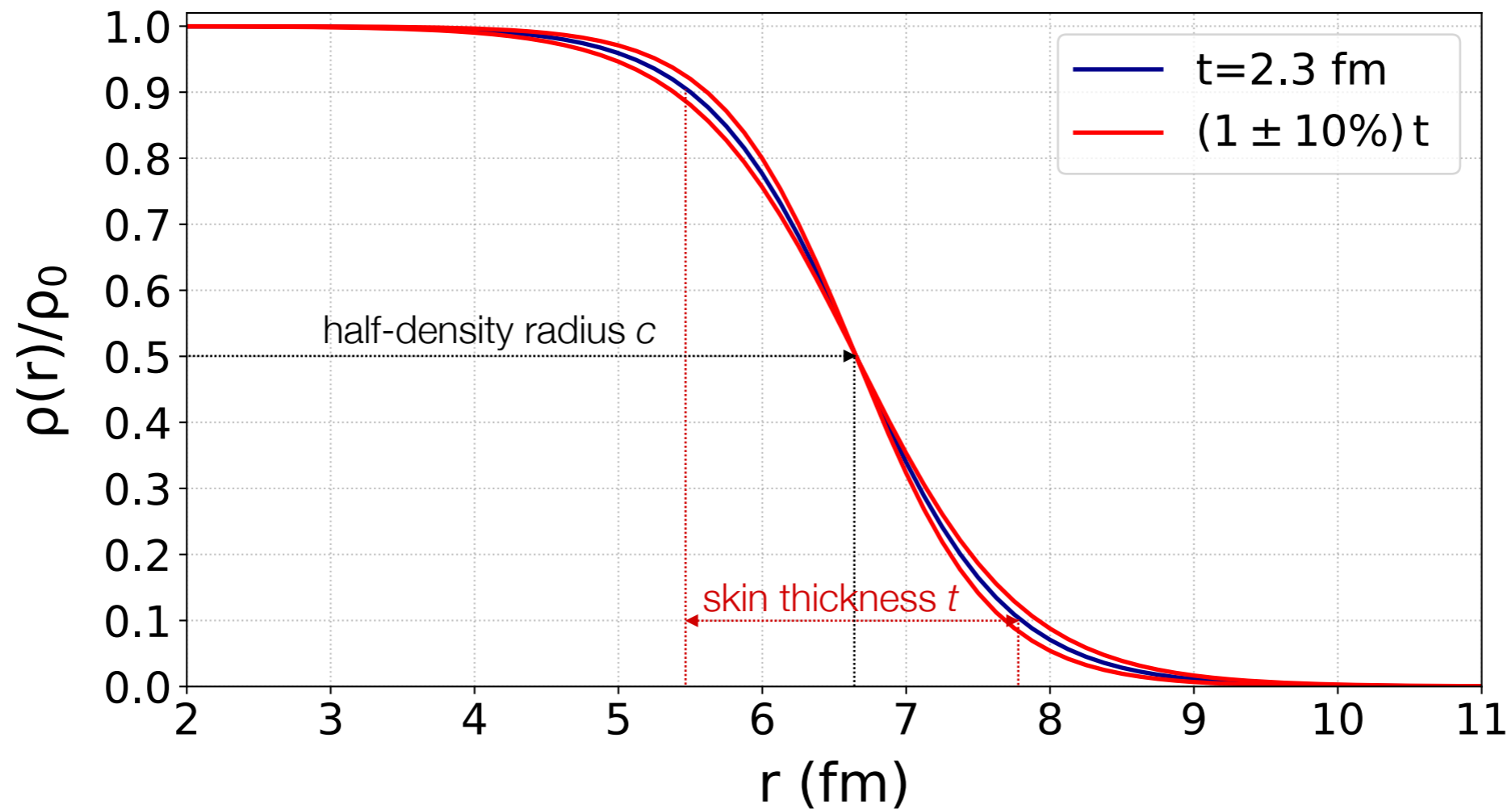
Sources of systematic uncertainties

		Charge radius		Quadrupole moment			
Systematic effect	Description	$\Delta(dR)$	σ_{dR}	$\Delta(dQ)$	σ_{dQ}		
Experimental	Fitting features	Instrumental line-shape	0	0.000 032	0	0.001 4	
		SE/FE ratio	0	0.000 015	0	0.001 3	
		Binning	0	0.000 008	0	0.000 8	
		Fitting energy range	0	0.000 033	0	0.001 3	
		Background model	0	0.000 042	0	0.004 0	
		Free intensities fit	0	0.000 028	0	0.000 4	
		Combined	0	0.000 070	0	0.004 7	
	Energy calibration	Wrong energy of ^{16}N line	0.000 12	0.000 04	-0.000 82	0.000 55	
		Energy calibration scheme	0	0.000 007	0	0.000 67	
		Uncertainty of literature energy	0	0.000 018	0	0.002 25	
		Line-shape for energy calibration	-0.000 068	0.000 038	0	0	
		Combined	0.000 052	0.000 058	-0.000 82	0.002 4	
	Theory	Theory	Uncertainty of nuclear polarization correction	0	0.000 20	0	0.000 11
	Theory input	Charge distribution model	Change of the skin thickness parameter	0	0.002 2	0	0.001 91
Discrepancies of spectrum and fit		Free Gaussian fits	0	0.000 95	0	0.028 2	
Total		0.000 052	0.002 4	-0.000 82	0.028 8		

Charge distribution model systematics: skin thickness

Nuclear charge distribution
for spherical nuclei

$$\rho(r) = \rho_0 \left(1 + e^{4 \ln 3 \frac{r-c}{t}} \right)^{-1}$$



$$\frac{t}{E} \text{ sensitivity} \sim 1.8 \times 10^{-5} \frac{\text{fm}}{\text{eV}} :$$

$$\Delta E_{\pm 10\% t} \sim 13 \text{ keV} \Rightarrow \sigma_{\text{dR}} \approx 0.0022$$

Spectroscopic quadrupole moment results

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Measurement of the quadrupole moment of ^{185}Re and ^{187}Re from the hyperfine structure of muonic X rays

A. Antognini,^{1,2} N. Berger,³ T. E. Cocolios,⁴ R. Dressler,¹ R. Eichler,¹ A. Eggenberger,² P. Indelicato,⁵ K. Jungmann,⁶ C. H. Keitel,⁷ K. Kirch,^{1,2} A. Knecht,¹ N. Michel,⁷ J. Nuber,^{1,2} N. S. Oreshkina,^{7,*} A. Ouf,⁸ A. Papa,^{1,9} R. Pohl,⁸ M. Pospelov,^{10,11} E. Rapisarda,^{1,1} N. Rijubo,^{1,2} S. Roccia,^{12,13} N. Severijns,⁴ A. Skawran,^{1,2} S. M. Vogiatzi,^{1,2} F. Wauters,³ and L. Willmann⁶

- Thorough study of systematic uncertainty studies is performed:
 - background model
 - line-shape model
 - cascade intensities
- Results are compared with previous analysis using muonic atoms by Konijn *et al.*^[1]:
 - natural rhenium
 - weaker multiplets ($5f_{5/2} \rightarrow 4d_{5/2}, 5g_{7/2} \rightarrow 4f_{7/2}$) are not included

Isotopes	Q (b)	
^{185}Re	2.07(5)	our analysis
^{187}Re	1.94(5)	
^{185}Re	2.21(4)	past analysis with muonic atoms by Konijn <i>et al.</i> ^[1]
^{187}Re	2.09(4)	
^{185}Re	2.18(3) _{stat}	our analysis without the weaker multiplets
^{187}Re	2.05(2) _{stat}	

[1] Konijn *et al.*, Nucl. Phys. A 360, 187 (1981)