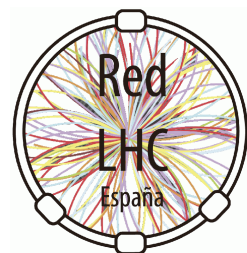


Measuring quantum entanglement in $t\bar{t}$ events at the LHC



7th RedLHC workshop - May 11, 2023

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The **Standard Model** is a **Quantum Field Theory**: Special Relativity + Quantum Mechanics

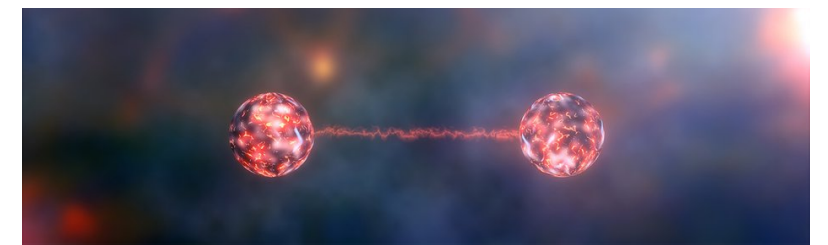
**Standard Model
measurements**



**Test fundamental properties
of Quantum Mechanics**

Entanglement is perhaps the most genuine and essential feature of Quantum Mechanics

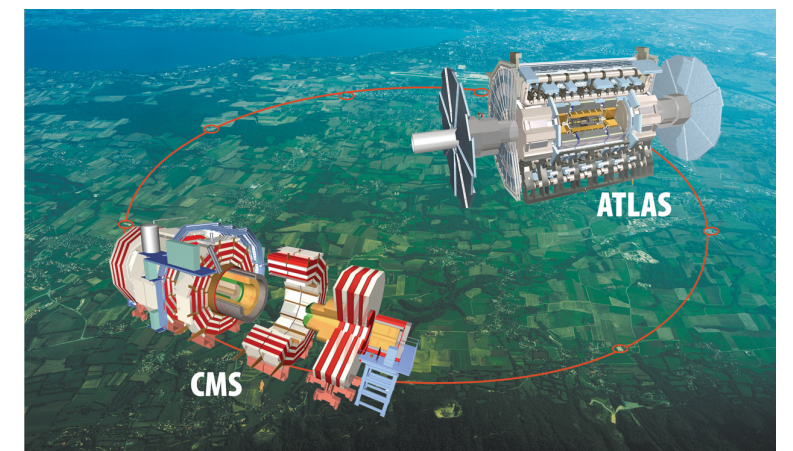
- If two (or more) particles become entangled, they remain connected even when separated by vast distances
- In other words, the quantum state of one particle cannot be described independently of the quantum state(s) of the other(s)



$$|\psi\rangle = |a_1\rangle_A \otimes |b_1\rangle_B + |a_2\rangle_A \otimes |b_2\rangle_B$$

The **LHC** has the potential to explore fundamental properties of Quantum Mechanics such as **Entanglement**!

- It can be measured with data already recorded at the LHC
→ Run 2 dataset
- Measuring experimentally this fundamental property requires a **very precise understanding of our detectors**



The **LHC** is a **top-quark factory**... and the top quark is the ideal candidate for measuring spin correlations:

- **Lifetime ($\sim 10^{-25}$ s) \ll hadronization ($\sim 10^{-23}$ s) \ll depolarization ($\sim 10^{-21}$ s)**
 - Decays before forming bound states
 - **Spin information preserved in the angular distribution of its decay products**

In SM, $t\bar{t}$ production:

- General form:

$$\rho = \frac{I_4 + \sum_i (B_i^+ \sigma^i + B_i^- \bar{\sigma}^i) + \sum_{i,j} C_{ij} \sigma^i \bar{\sigma}^j}{4}$$

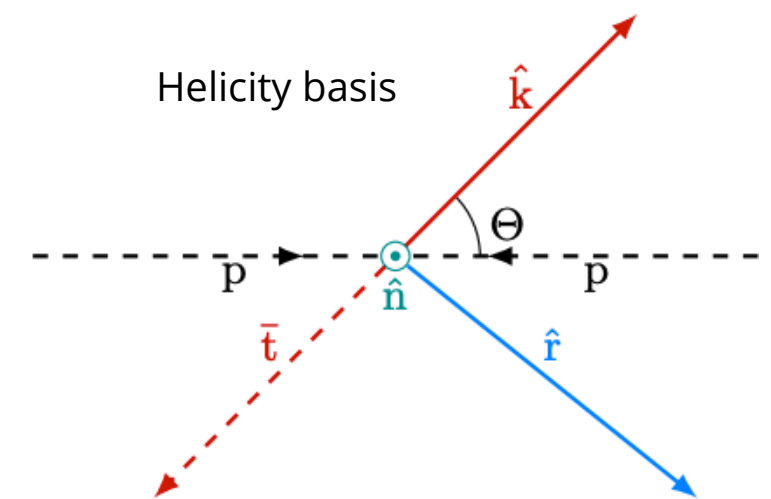
- $\sigma^i/2, \bar{\sigma}^i/2$ - spin operators of the top, antitop.
- B_i^+, B_i^- characterize the spin polarizations, $B_i^+ = \langle \sigma^i \rangle$, $B_i^- = \langle \bar{\sigma}^i \rangle$.
- At LO $B_i^\pm = 0$.
- C_{ij} the $t\bar{t}$ spin correlations, $C_{ij} = \langle \sigma^i \bar{\sigma}^j \rangle$.

- **Top-quark not polarised (at LO) in $t\bar{t}$ production in SM (parity invariant)**
- **But spins of t and \bar{t} strongly correlated (rich structure of spin correlations)**

Spin correlations

- Top-quark spins cannot be measured directly!
- This is done by measuring the **angle between spin axis and lepton** in parent top-quark rest frame

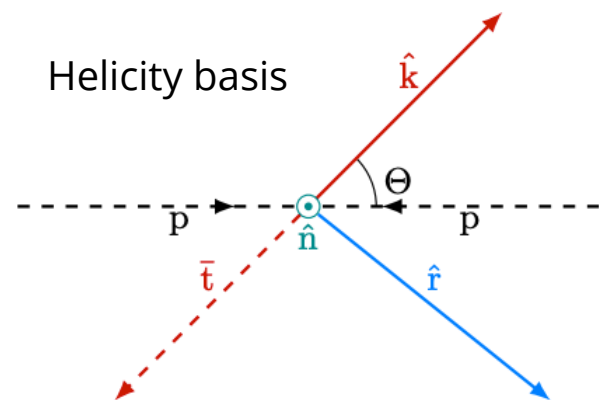
- Helicity basis: $\{\hat{k}, \hat{r}, \hat{n}\}$:
 - \hat{k} - direction of the top in the $t\bar{t}$ CM frame.
 - \hat{p} - direction of the beam.
 - $\cos \Theta = \hat{k} \cdot \hat{p}$.
 - $\hat{r} = (\hat{p} - \cos \Theta \hat{k}) / \sin \Theta$.
 - $\hat{n} = \hat{r} \times \hat{k}$.
 - Describe each individual process with a fixed direction.
- Beam basis: $\{\hat{x}, \hat{y}, \hat{z}\}$:
 - \hat{z} along the beam axis.
 - \hat{x}, \hat{y} transverse directions to the beam.
 - After averaging: $C_x = C_y = C_\perp$.
 - Studying the total quantum state.



Spin correlations

- Coefficients measured by CMS and ATLAS from diff. cross-section:

$$\frac{1}{\sigma} \frac{d\sigma}{d\cos\theta_+ d\cos\theta_-} = \frac{1}{4} \left(1 + B_1 \cos\theta_+ + B_2 \cos\theta_- - C \cos\theta_+ \cos\theta_- \right)$$



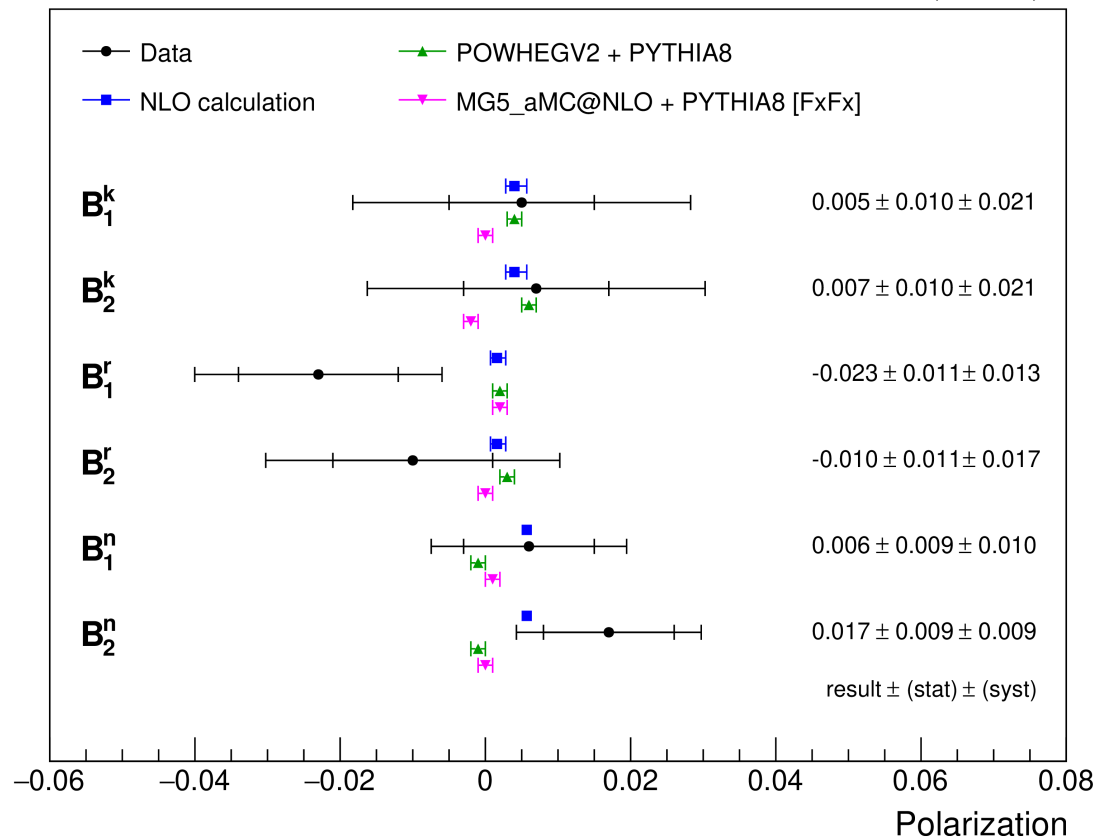
Observable	Measured coefficient	Coefficient function	Symmetries
$\cos\theta_1^k$	B_1^k	b_k^+	P-odd, CP-even
$\cos\theta_2^k$	B_2^k	b_k^-	P-odd, CP-even
$\cos\theta_1^r$	B_1^r	b_r^+	P-odd, CP-even
$\cos\theta_2^r$	B_2^r	b_r^-	P-odd, CP-even
$\cos\theta_1^n$	B_1^n	b_n^+	P-even, CP-even
$\cos\theta_2^n$	B_2^n	b_n^-	P-even, CP-even
$\cos\theta_1^{k*}$	B_1^{k*}	b_k^+	P-odd, CP-even
$\cos\theta_2^{k*}$	B_2^{k*}	b_k^-	P-odd, CP-even
$\cos\theta_1^{r*}$	B_1^{r*}	b_r^+	P-odd, CP-even
$\cos\theta_2^{r*}$	B_2^{r*}	b_r^-	P-odd, CP-even
$\cos\theta_1^k \cos\theta_2^k$	C_{kk}	c_{kk}	P-even, CP-even
$\cos\theta_1^r \cos\theta_2^r$	C_{rr}	c_{rr}	P-even, CP-even
$\cos\theta_1^n \cos\theta_2^n$	C_{nn}	c_{nn}	P-even, CP-even
$\cos\theta_1^r \cos\theta_2^k + \cos\theta_1^k \cos\theta_2^r$	$C_{rk} + C_{kr}$	c_{rk}	P-even, CP-even
$\cos\theta_1^r \cos\theta_2^k - \cos\theta_1^k \cos\theta_2^r$	$C_{rk} - C_{kr}$	c_n	P-even, CP-odd
$\cos\theta_1^n \cos\theta_2^r + \cos\theta_1^r \cos\theta_2^n$	$C_{nr} + C_{rn}$	c_{nr}	P-odd, CP-even
$\cos\theta_1^n \cos\theta_2^r - \cos\theta_1^r \cos\theta_2^n$	$C_{nr} - C_{rn}$	c_k	P-odd, CP-odd
$\cos\theta_1^n \cos\theta_2^k + \cos\theta_1^k \cos\theta_2^n$	$C_{nk} + C_{kn}$	c_{kn}	P-odd, CP-even
$\cos\theta_1^n \cos\theta_2^k - \cos\theta_1^k \cos\theta_2^n$	$C_{nk} - C_{kn}$	$-c_r$	P-odd, CP-odd
$\cos\varphi$	D	$-(c_{kk} + c_{rr} + c_{nn})/3$	P-even, CP-even
$\cos\varphi_{\text{lab}}$	$A_{\cos\varphi}^{\text{lab}}$	—	—
$ \Delta\phi_{\ell\ell} $	$A_{ \Delta\phi_{\ell\ell} }$	—	—



Spin correlations

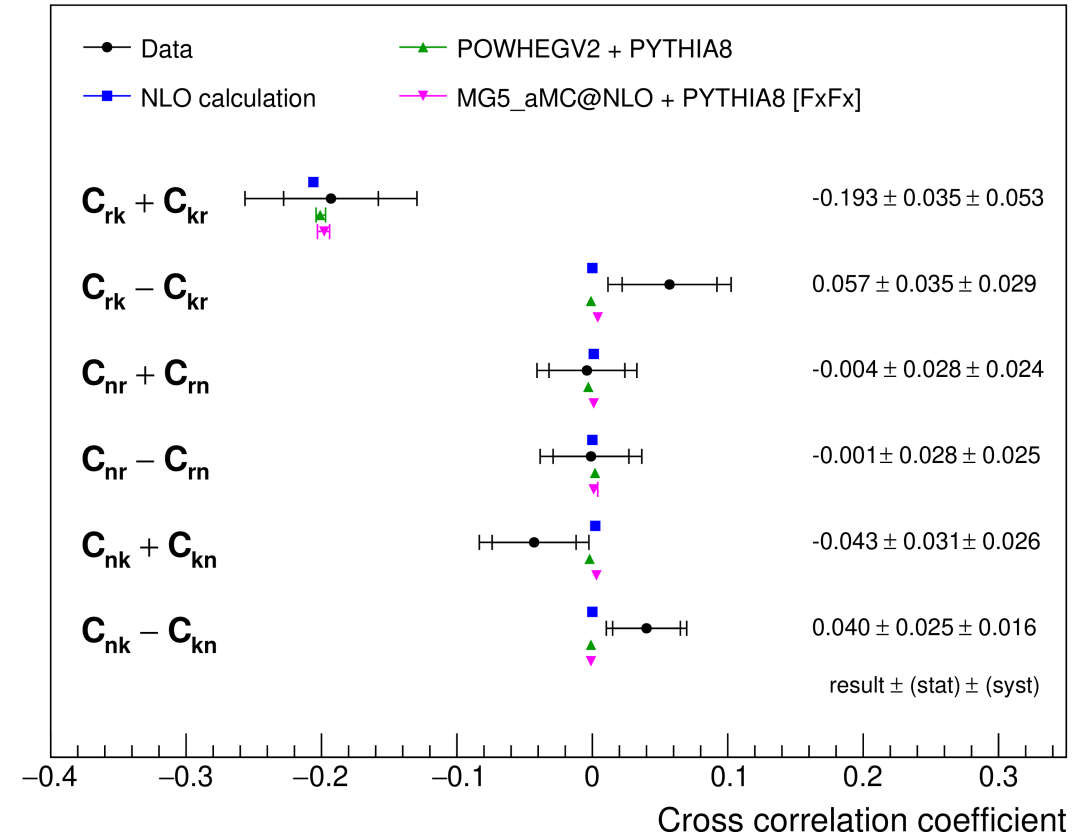
CMS

35.9 fb⁻¹ (13 TeV)



CMS

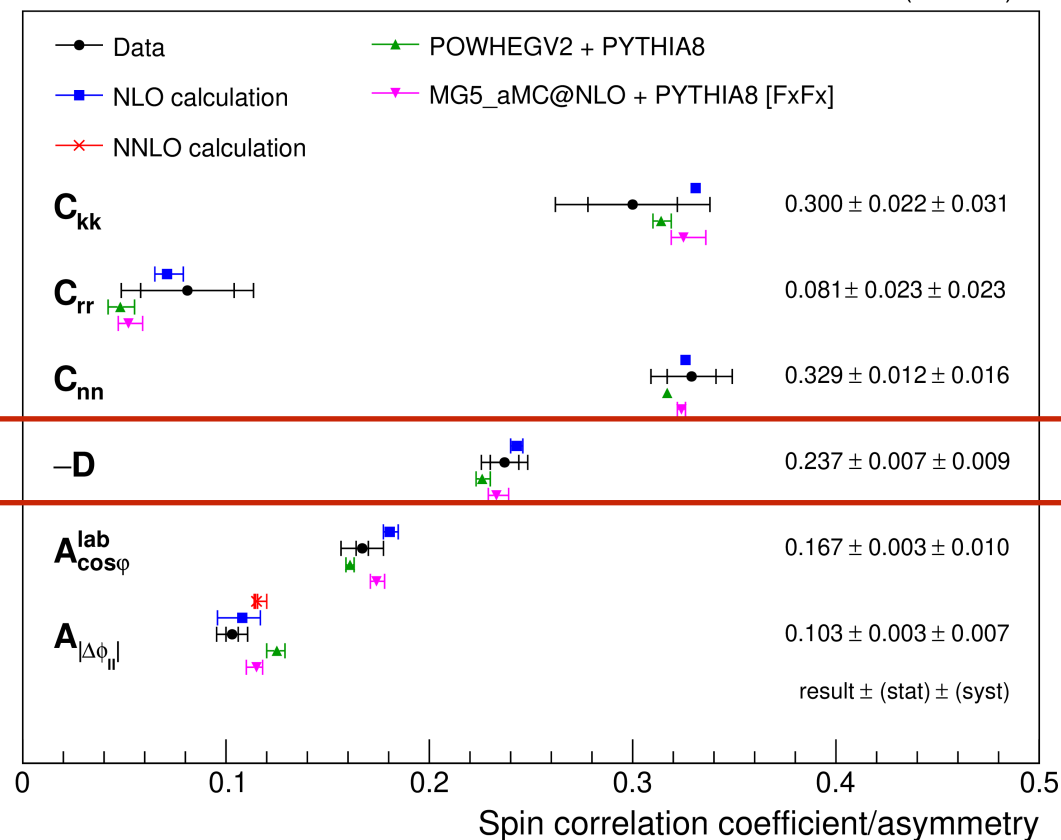
35.9 fb⁻¹ (13 TeV)



[Phys. Rev. D 100, 072002](https://arxiv.org/abs/1905.07202)

CMS

35.9 fb⁻¹ (13 TeV)



- Spin Correlations can be a classical property
 - **Spin Correlations \neq Quantum Entanglement!**
 - However, **Quantum Entanglement \subset Spin-Correlations**
- Indeed, the link between spin correlations of top quarks and Quantum Information is recent ([Eur. Phys. J. Plus 136 \(2021\) 9, 907...](#))
 - $t\bar{t}$ process represents a simple entangled system composed by two qubits
 - $t\bar{t}$ events \rightarrow good candidates to test entanglement and Bell inequalities at high energy!

Entanglement criterion:
$$D = \frac{\text{tr}[C]}{3}$$

where
$$\text{tr}[C] = 2C_{\perp} + C_z = C_{rr} + C_{nn} + C_{kk}$$

No requirement to measure the spin density matrix!

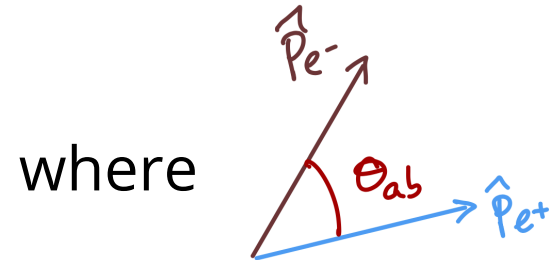
Entanglement condition

Upper bound on trace of spin density matrix:

$$D = \frac{\text{tr}[C]}{3} < -\frac{1}{3}$$

- **Theory: simple observable** from single differential normalized cross-section:

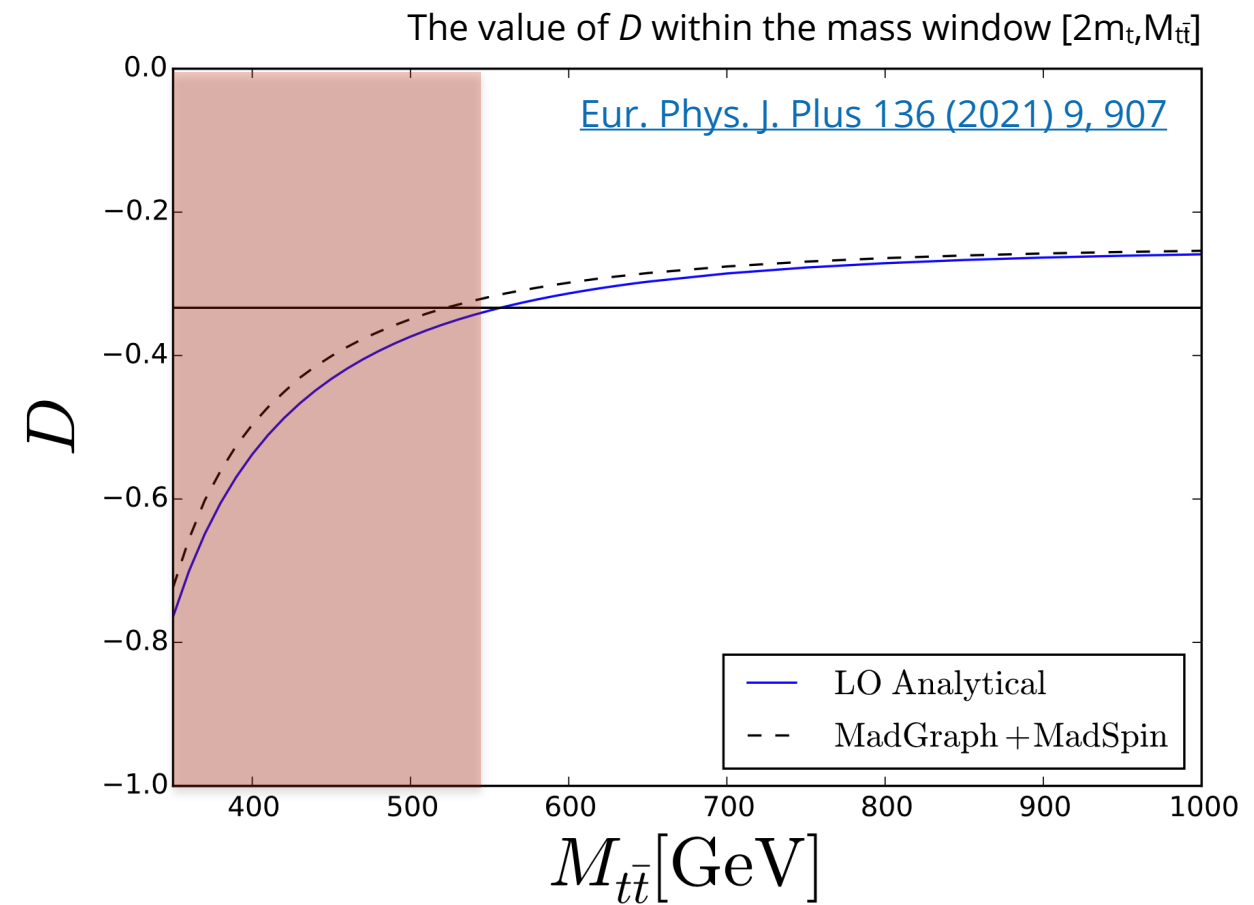
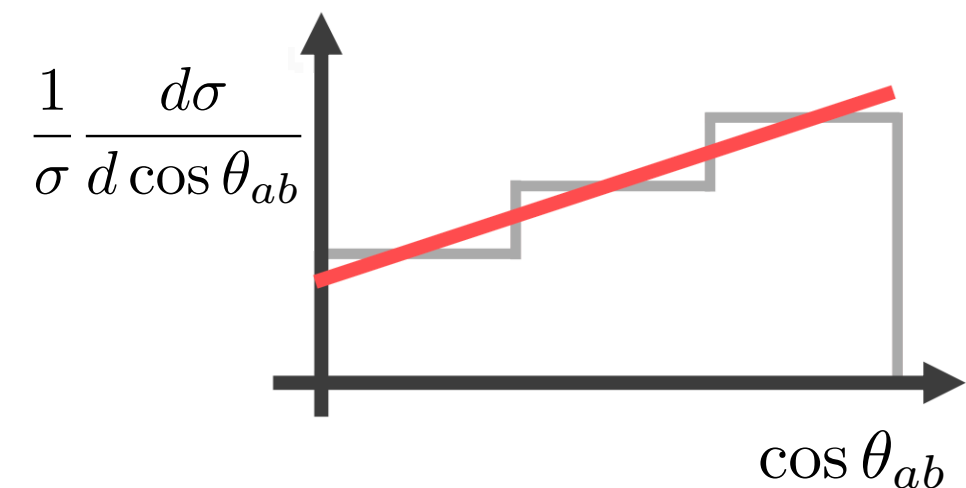
$$\frac{1}{\sigma} \frac{d\sigma}{d \cos \theta_{ab}} = \frac{1}{2} (1 - \alpha_a \alpha_b D \cos \theta_{ab})$$



where θ_{ab} is the angle between the two leptons in their respective top-quark parent rest-frame

- **Theory: “just” measure $\cos(\theta_{ab})$ in a low $M(t\bar{t})$ region**

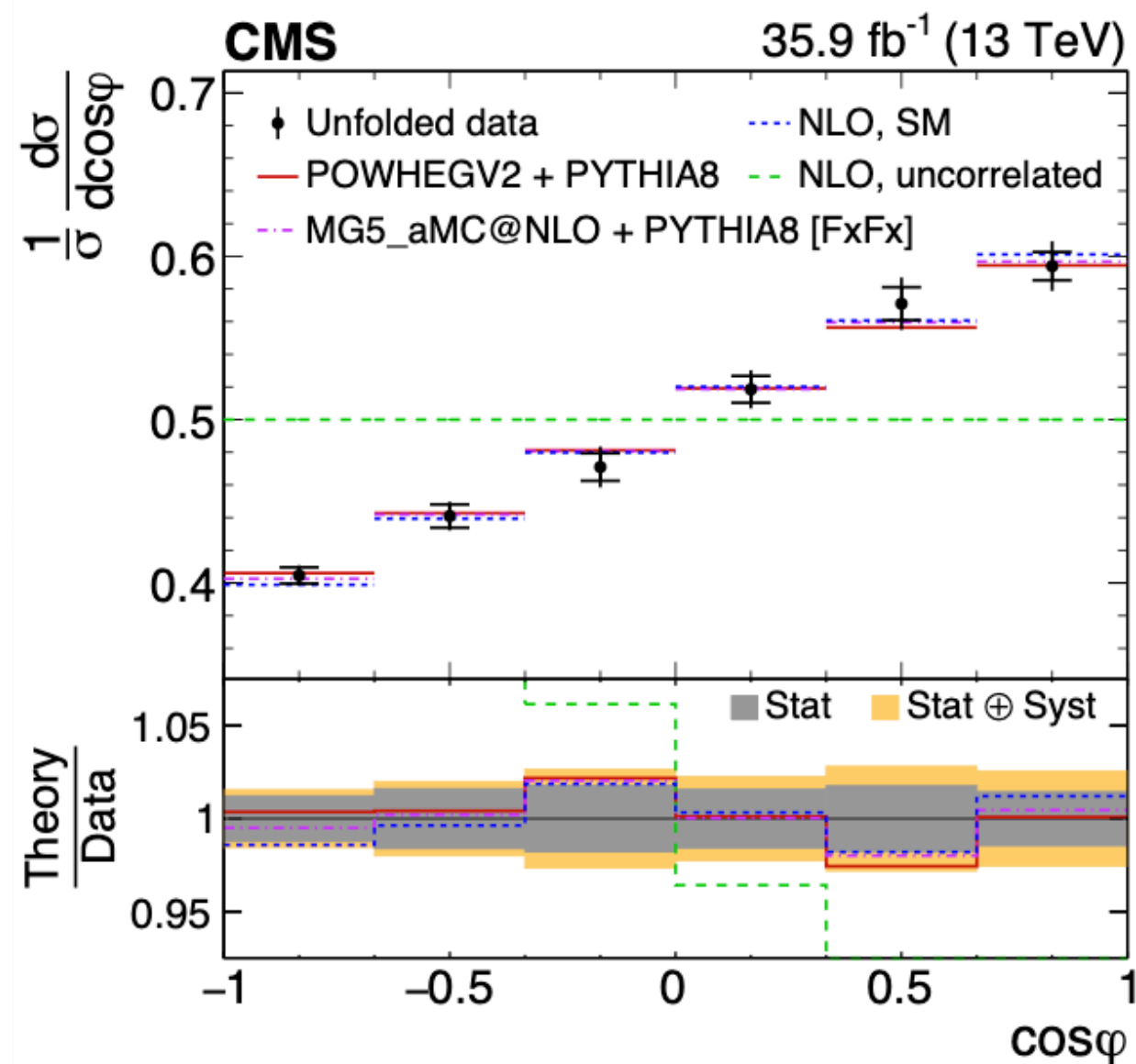
$$\left. \frac{1}{\sigma} \frac{d\sigma}{d \cos \theta_{ab}} \right|_{m_{t\bar{t}} < M} = \frac{1}{2} (1 - D \cos \theta_{ab})$$



D already measured though inclusively

Recently, D was measured inclusively on $M(t\bar{t})$ by CMS:

- $D = -0.237 \pm 0.011 > -1/3$
- $\Delta D/D = 4.6\%$



[Phys. Rev. D 100, 072002](#)

- Spin Correlations \neq Quantum Entanglement!



No public results yet neither from ATLAS nor CMS

- I cannot show you any result (though they exist) 🙄
- Good news: many people currently working on this topic!

ATLAS and CMS are working on two $t\bar{t}$ final states:

Dileptonic

2 opposite-sign leptons
 ≥ 2 jets (b-tagged)
 ≥ 1 b-tagged jets
Exclude Z-mass window
 $M_{t\bar{t}} < \sim 400$ GeV



Challenge: reconstruct the neutrinos



Several techniques are available:

- Roots of quartic polynomial
- NeutrinoWeighter
- Sonnenschein method
- Ellipse method
- Use ML (new): Transformers (SPANet)

Lepton+jets

- 1 lepton
- ≥ 4 jets
- 2 jets must be b-tagged
- $MET > 20$ GeV
- $M_{t\bar{t}} < \sim 400$ GeV



Challenge: tagging the correct jet(s)

- High-multiplicity jet final state



Use ML for multi-jet final state:

- Transformers (SPANet)

How we are trying to measure entanglement in $t\bar{t}$:

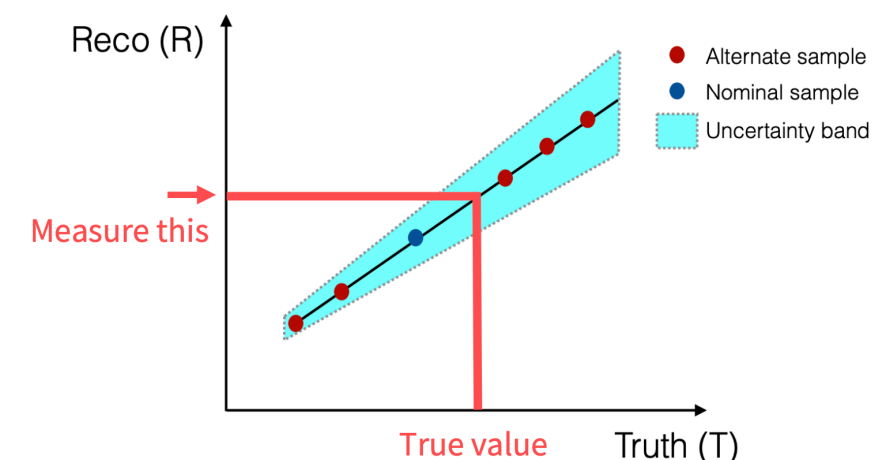
Unfold $\cos(\theta_{ab})$, then extract D

- Correction of the detector effects is needed since efficiency decreases when the leptons approach collinearity
- Many unfolding techniques exist:
 - Iterative Bayesian unfolding
 - Profile likelihood unfolding
 - SVD unfolding
- Stress tests (SM bias) needed → **Issues appear!**



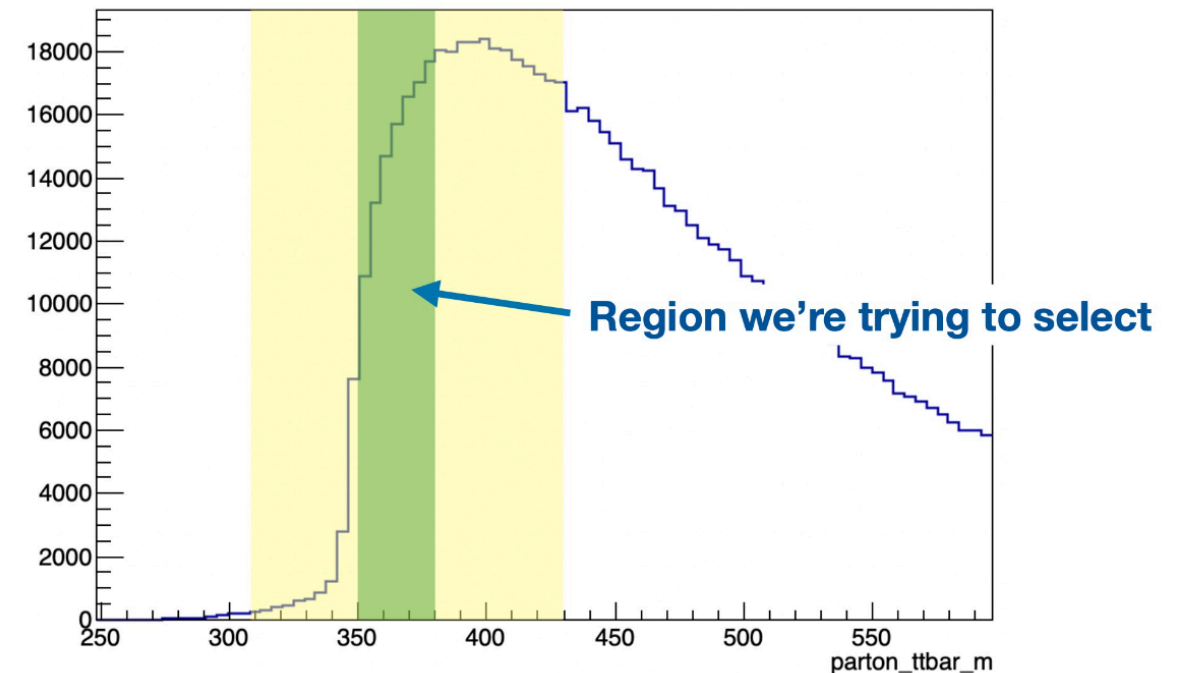
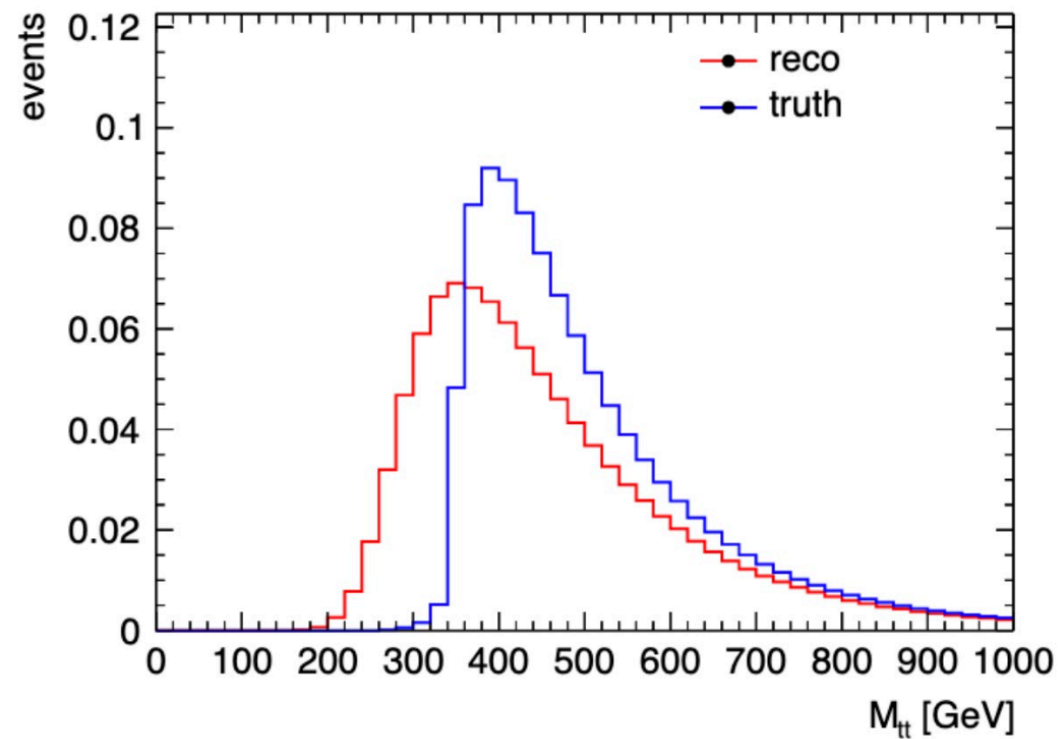
Extract D from reconstructed events, then extract D

- Calibration Curve
- Problem how to produce alternate samples
 - This is not a free parameter in the SM
 - In MC event generators, not an input parameter which we can alter
- Reweighting? → **Issues appear!**
 - Alter slope of $\cos(\theta_{ab})$ artificially
 - Did not preserve linearity
 - Preserved inclusive value of D



Measuring entanglement in $t\bar{t}$ events at the LHC

The problem is the resolution of $M_{t\bar{t}}$



- Measuring the entanglement in $t\bar{t}$ events requires better (great) top-quark reconstruction when considering such narrow phase-space
- Modelling effects affect systematics estimation

Theory: “just” measure $\cos(\theta_{ab})$
in a low $M(t\bar{t})$ region



Experiment: wait,
this is not so easy



Conclusions

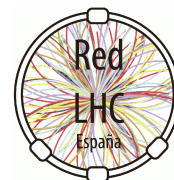
- First studies of the measurement of entanglement between quarks
- Can be detected at the LHC with current recorded data (Run2)
- Simple observable from single differential cross-section:

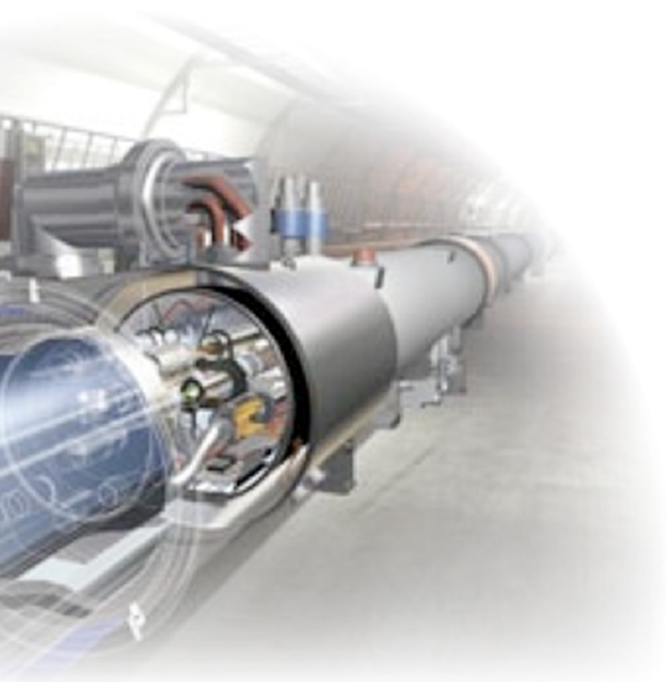
$$D = \frac{\text{tr}[C]}{3} < -\frac{1}{3} \longrightarrow \frac{1}{\sigma} \frac{d\sigma}{d\cos\theta_{ab}} \Big|_{m_{t\bar{t}} < M} = \frac{1}{2} (1 - D \cos\theta_{ab})$$

- Several $t\bar{t}$ final states being studied
- Several techniques being explored to extract D
- Real requirement for superior top-quark reconstruction
- Narrow phase-space is problematic

“Quantum entanglement is theoretically clean, but experimentally quite nasty” — Alan Barr —

- **ATLAS and CMS: work in progress, please stay tuned!**



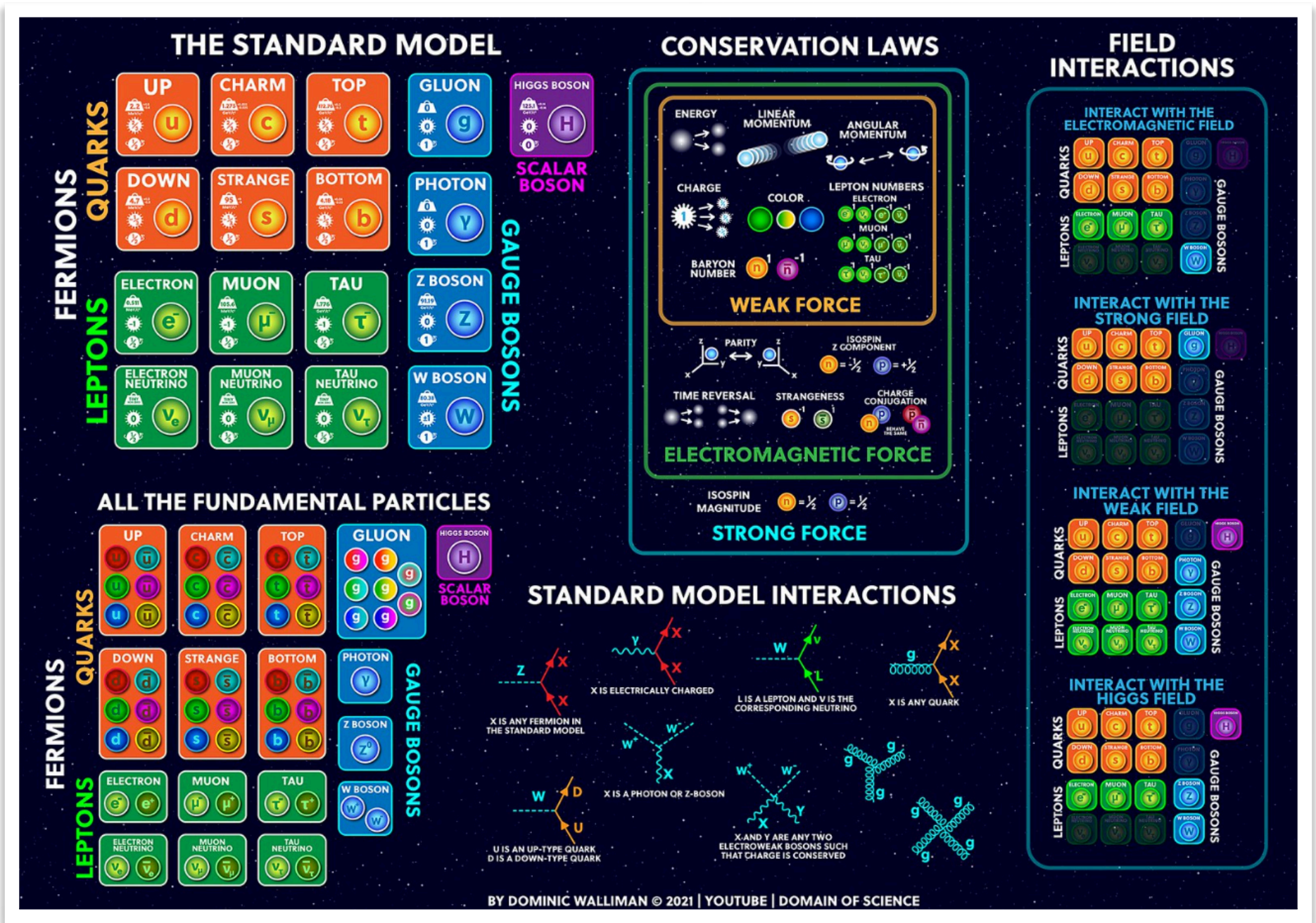


BACKUP



The **Standard Model** is a **Quantum Field Theory**:

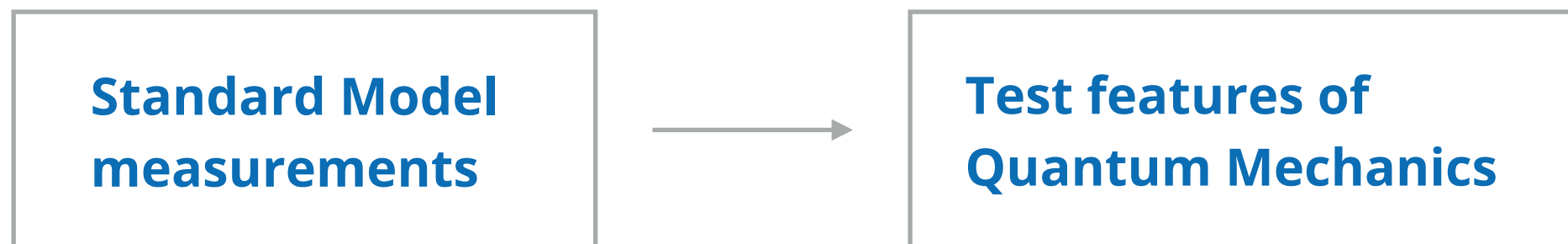
- Special Relativity
- Quantum Mechanics



The **Standard Model** is a **Quantum Field Theory**:

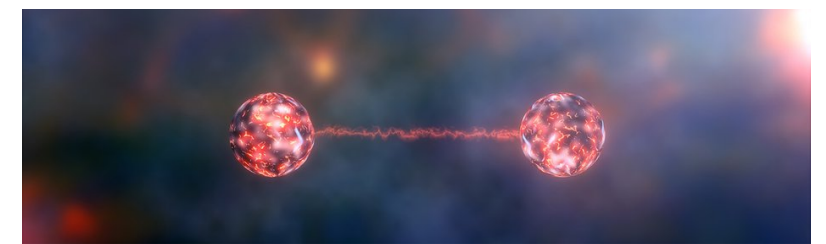
- Special Relativity
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Fundamental properties of Quantum Mechanics can be tested via the Standard Model



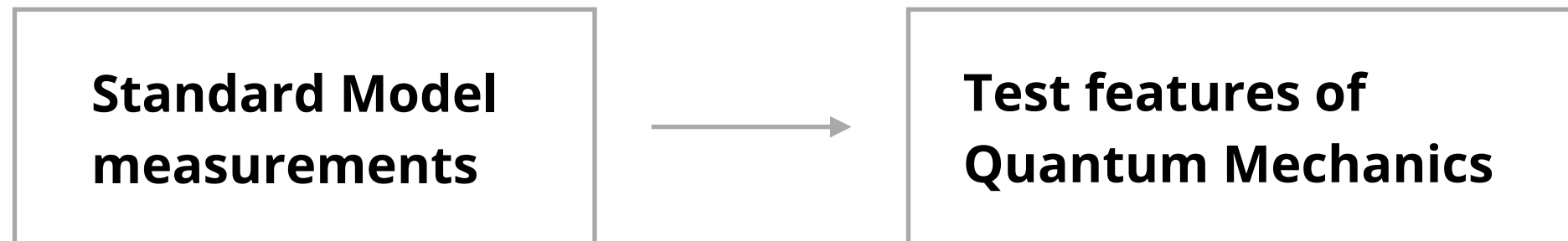
Entanglement is perhaps the most genuine and essential feature of Quantum Mechanics

- If two (or more) particles become entangled, they remain connected even when separated by vast distances
- In other words, the quantum state of one particle cannot be described independently of the quantum state(s) of the other(s)



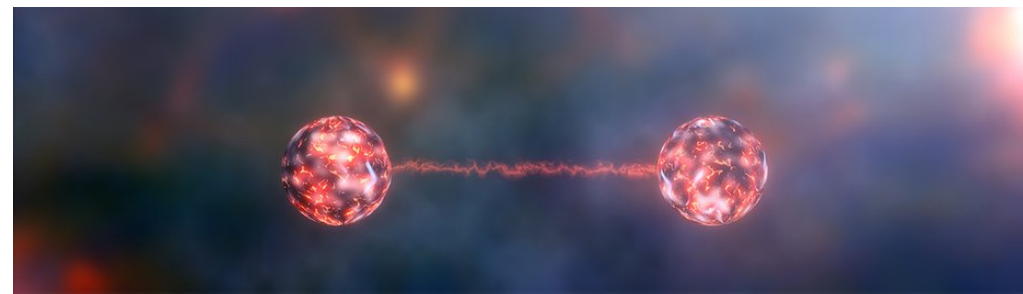
$$|\psi\rangle = |a_1\rangle_A \otimes |b_1\rangle_B + |a_2\rangle_A \otimes |b_2\rangle_B$$

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Entanglement is perhaps the most genuine and essential feature of Quantum Mechanics

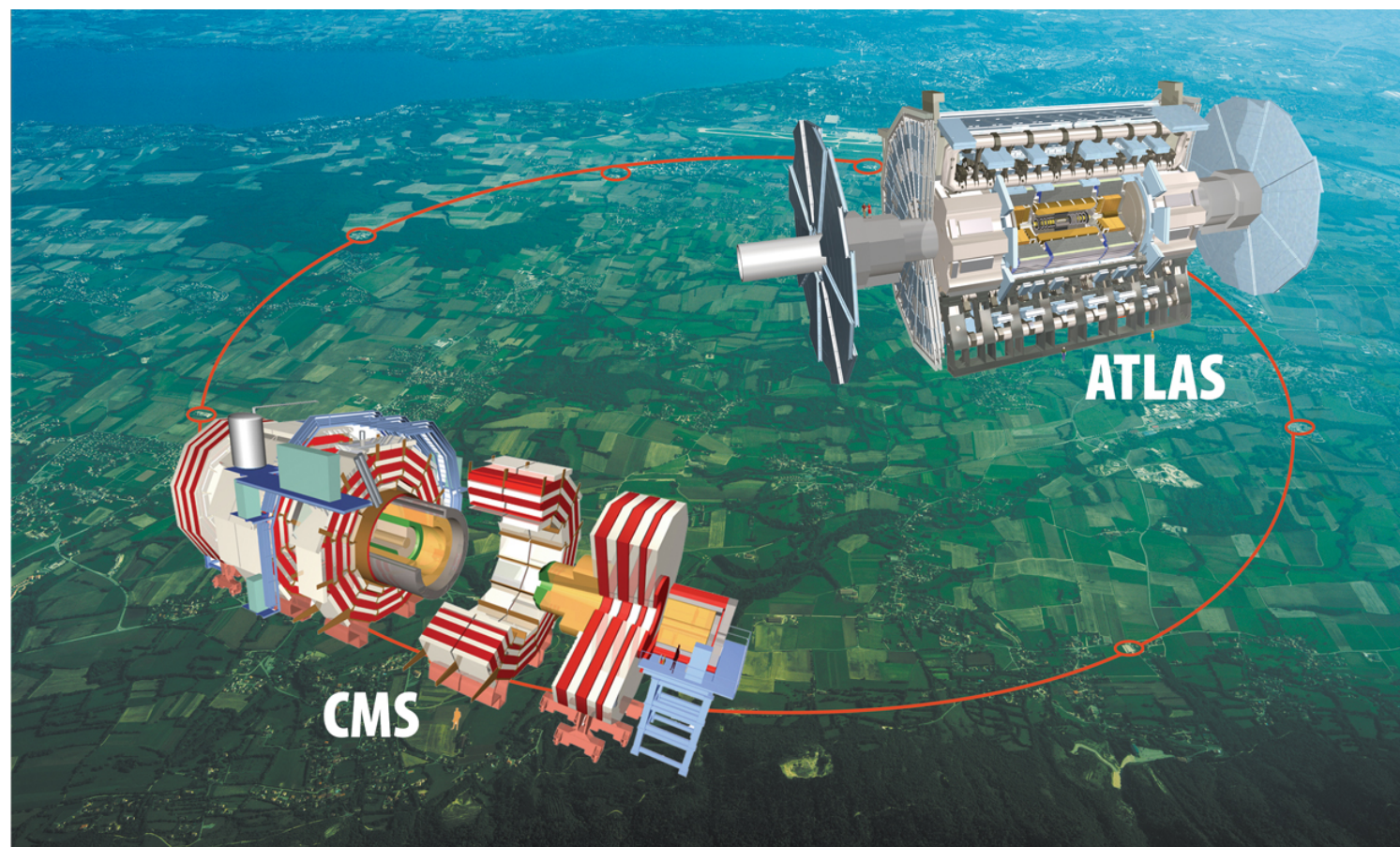
- In other words, the quantum state of one particle cannot be described independently of the quantum state(s) of the other(s)
- EPR paradox → Information travel faster than light?
 - Contradicts the theory of relativity
 - Conclusion: the theory of Quantum Mechanics is incomplete



$$|\psi\rangle = |a_1\rangle_A \otimes |b_1\rangle_B + |a_2\rangle_A \otimes |b_2\rangle_B$$

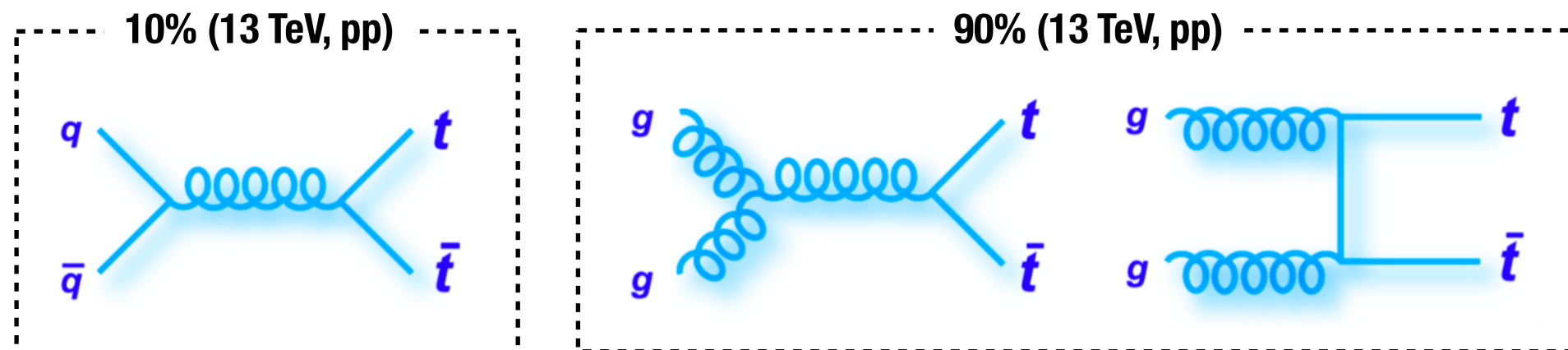
The **LHC** has the potential to explore fundamental properties of Quantum Mechanics such as **Entanglement**!

- It can be measured with data already recorded at the LHC → Run 2 dataset
- Measuring experimentally this fundamental property requires a **very precise understanding of our detectors**



The **LHC** is a **top-quark factory**:

- Top quarks are abundantly produced at the LHC
 - At low luminosity (i.e. $10^{32} \text{ cm}^{-2} \text{ s}^{-1}$ @ 7 TeV): $\sim 60 \text{ } t\bar{t}$ every hour
 - At design luminosity (i.e. $10^{34} \text{ cm}^{-2} \text{ s}^{-1}$ @ 14 TeV): $\sim 8 \text{ } t\bar{t}$ every second



Unique properties of the top quark:

- Heaviest known elementary particle
- Strongest Yukawa coupling (almost unity)
- Smallest cross-section of all of the SM particles
- **Lifetime ($\sim 10^{-25} \text{ s}$) \ll hadronization ($\sim 10^{-23} \text{ s}$) \ll depolarization ($\sim 10^{-21} \text{ s}$)**
 - **The decay products preserve the spin information of the top quark**
 - Only quark whose most of its properties can be directly measured!
- The top quark decays almost exclusively ($>99\%$, i.e. $|V_{tb}| \approx 0.999$) to $t \rightarrow Wb$ (@LO)
- New physics: a key player in most solutions to the problems of the SM

Hidden variables

By EPR, each particle "carries" variables that knows the state before the measurement

⇒ There are some hidden variables that are missing in order to have a full theory

The Copenhagen Interpretation: superposition of states until a measurement is done

Bell's Inequality

- If local hidden variables holds, they should satisfy some inequality
- $C(x, y)$ are the correlations between different measurements at different detectors
- The parameters a, b, c are different directions for the measurement
- Original form: $1 + C(b, c) \geq |C(a, b) - C(a, c)|$

Top quark is the ideal candidate for measuring spin correlations:

- extremely short lifetime \rightarrow decays before forming bound states
- spin information preserved in the angular distribution of its decay products
- top-quark spin observables (expected to be) well predicted by perturbative QCD

In SM, $t\bar{t}$ production:

- General form:

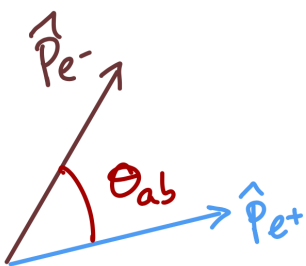
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- **Theory:** **simple observable** from single differential normalized cross-section:

$$\frac{1}{\sigma} \frac{d\sigma}{d \cos \theta_{ab}} = \frac{1}{2} (1 - \alpha_a \alpha_b D \cos \theta_{ab})$$

where  is the angle between the two leptons in their respective top-quark parent rest-frame

	<i>b</i> -quark	<i>W</i> ⁺	<i>l</i> ⁺	<i>d</i> [−] -quark or <i>s</i> [−] -quark	<i>u</i> -quark or <i>c</i> -quark
α_i (LO)	-0.41	0.41	1	1	-0.31
α_i (NLO)	-0.39	0.39	0.998	0.97	-0.32

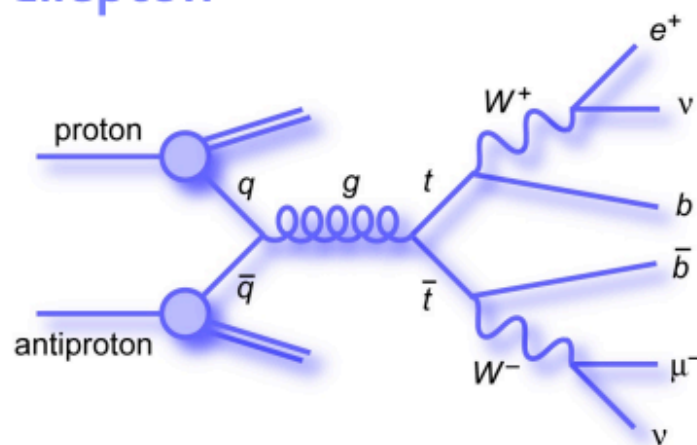


No public results yet neither from ATLAS nor CMS

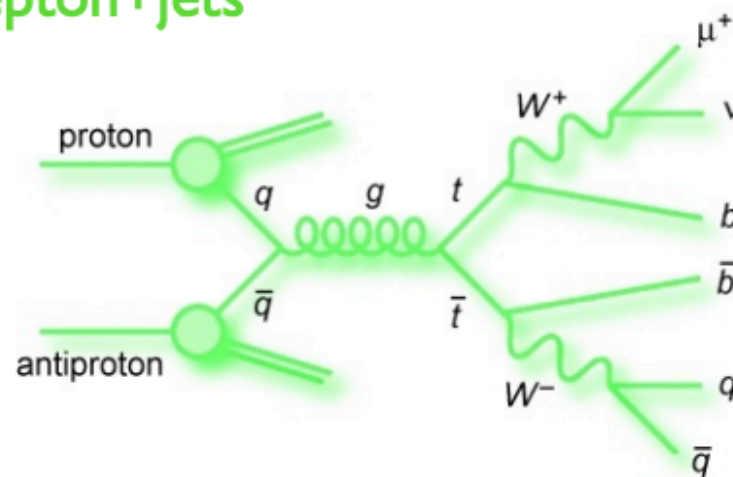
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From the experimental point of view, we have three options:

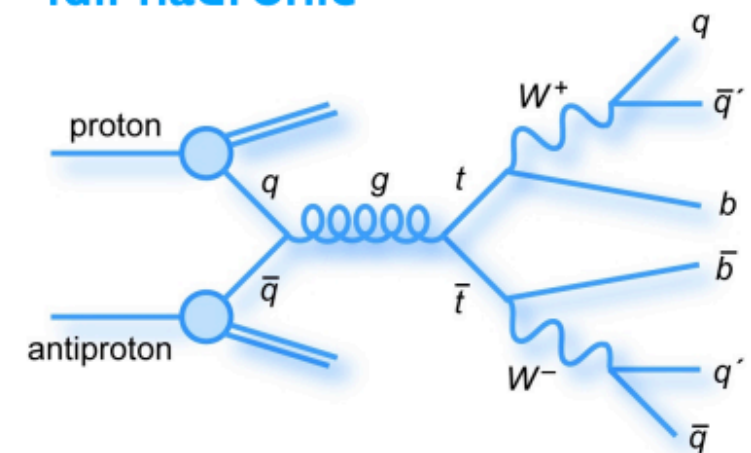
dilepton



lepton+jets



full hadronic



BR increases.... but background too!

Challenge: reconstruct the neutrinos

Challenge: tagging the correct jet(s)

- High-multiplicity jet final state