



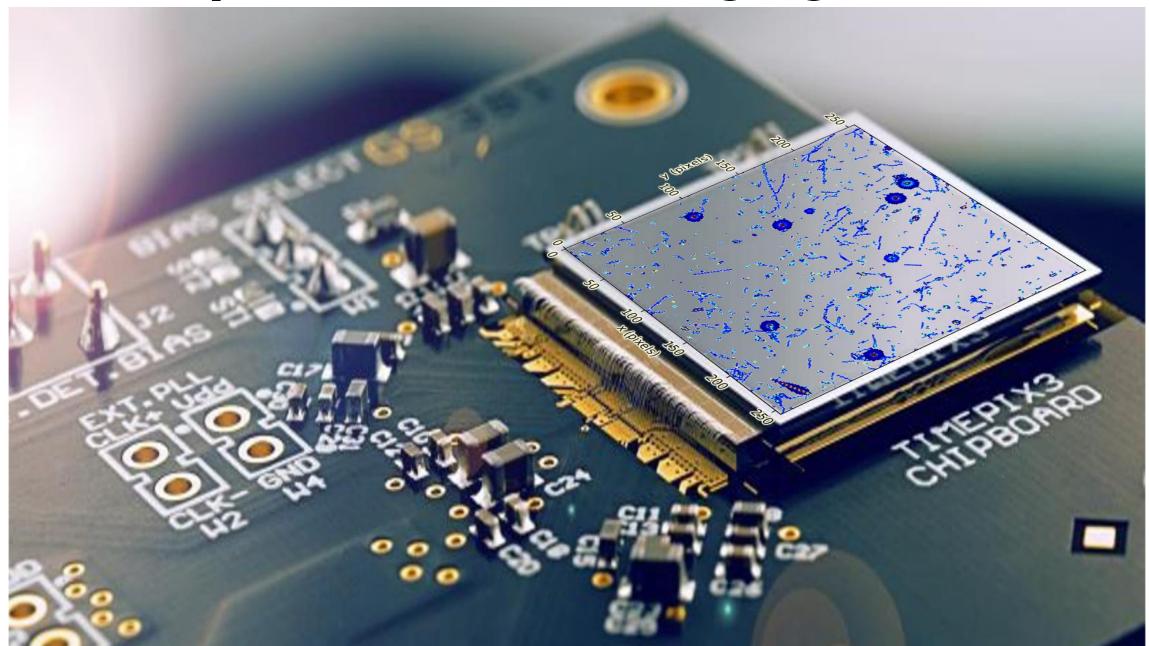
Application of single-layer particle tracking for radiation field decomposition and interaction point reconstruction at MoEDAL

Declan Garvey, Benedikt Bergmann, Petr Mánek, Stanislav Pospíšil, Petr Smolyanskiy,

On behalf of MoEDAL

<u>Declan.Garvey@utef.cvut.cz</u>

Timepix3: Radiation Imaging Detector



Timepix and Timepix3

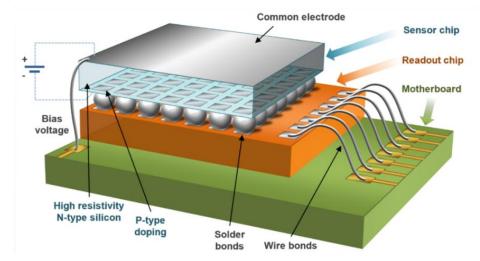
- 256 x 256 pixels with 55 μm pitch (1.98 cm²)
- Sensor layer (silicon, GaAs, CdTe, ...) flip-chip bump bonded to the ASIC

Timepix

- Frame-based readout (92 fps) dead time > 11 ms
 - Information is read out on a frame-by-frame basis
- Measurement of energy or time ($\Delta t = 20.8 \text{ ns}$)
- Minimal detectable energy: ~ 3.5 keV

Timepix3

- Data-driven readout (max. count rate 40 Mpix/s)
 - ⇒ Pixels are continuously read out throughout the entire measurement
- Per pixel dead time: 475 ns
- Measurement of energy and time ($\Delta t = 1.56 \text{ ns}$)
- Minimal detectable energy per pixel: 3 keV





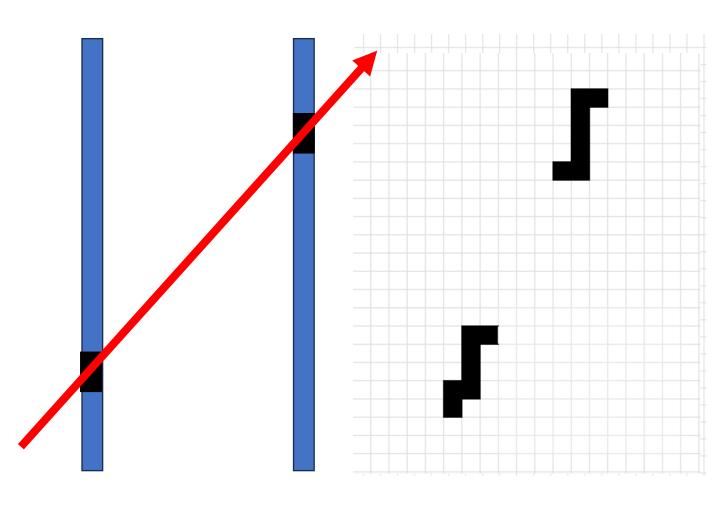
Timepix3 with chipboard

Particle Tracking by Connecting the Dots

- Traditionally, particle tracking is performed using multiple layers of thin detectors
- Tracking information is then calculated by recording coincidence measurements in each detector and "connecting the dots"

So, the question is:

Can we achieve similar results using only a single-layer detector?

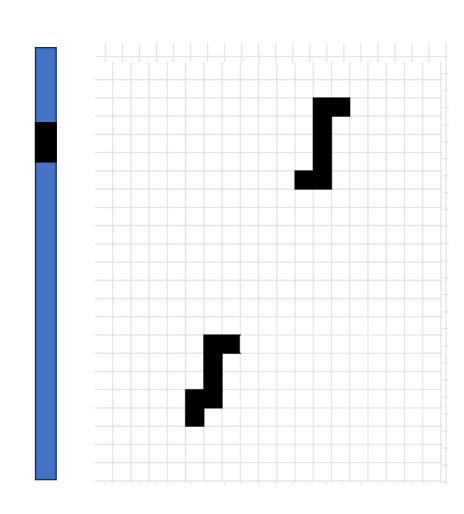


Particle Tracking by Connecting the Dots

- Traditionally, particle tracking is performed using multiple layers of thin detectors
- Tracking information is then calculated by recording coincidence measurements in each detector and "connecting the dots"

So, the question is:

Can we achieve similar results using only a single-layer detector?

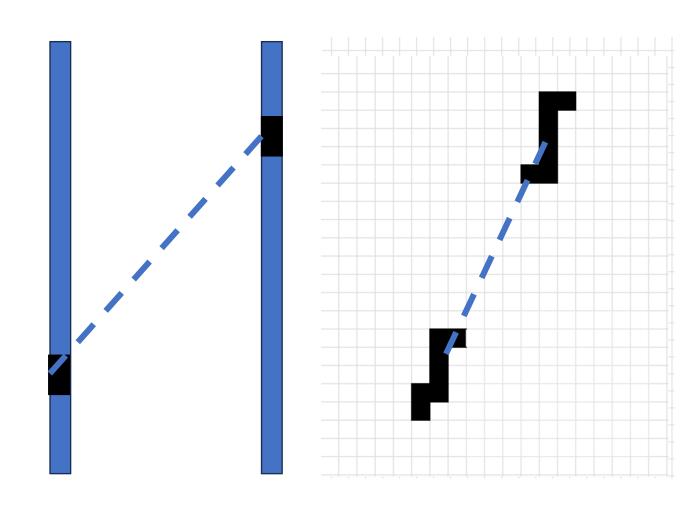


Particle Tracking by Connecting the Dots

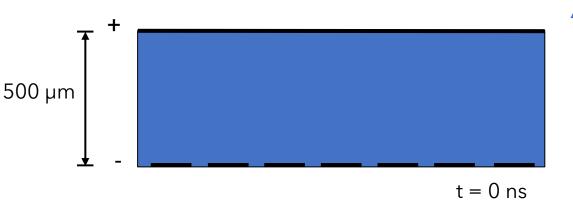
- Traditionally, particle tracking is performed using multiple layers of thin detectors
- Tracking information is then calculated by recording coincidence measurements in each detector and "connecting the dots"

So, the question is:

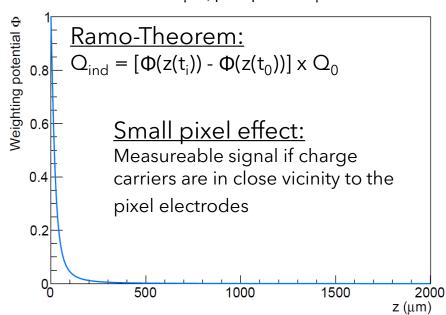
Can we achieve similar results using only a single-layer detector?



3D Reconstruction of Particle Tracks



2000 μm, pixel pitch: 55 μm



↑ Z

Charge carrier drift motion:

e⁻ and h⁺ drift described by

$$v_e = -\mu_e \times E(z)$$

 $v_h = \mu_h \times E(z)$

 $\mu_{e/h}$: Mobility of e⁻/h⁺

Electric field parametrisation:

Si:
$$\vec{E}(z) = \frac{U_B}{d} \vec{e_z} + \frac{2U_{dep}}{d^2} \left(\frac{d}{2} - z\right) \vec{e_z}$$
;

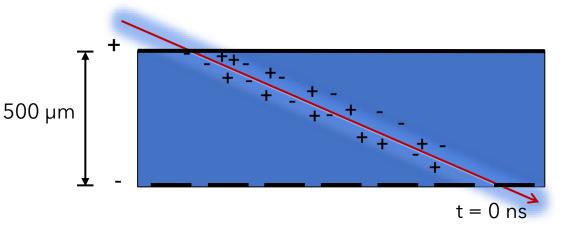
CdTe:
$$\vec{E}(z) = \frac{U_B}{d} \vec{e_z}$$

 U_B : Bias voltage; U_{dep} : Depletion voltage; d: Sensor thickness

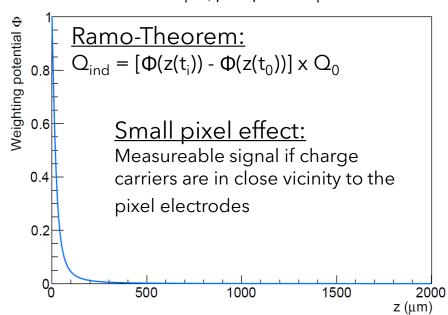
 \rightarrow Look up table: $z(t_{meas.}, E_{meas.})$

Bergmann et al. Eur. Phys. J. C (2017) 77: 421. https://doi.org/10.1140/epjc/s10052-017-4993-4

3D Reconstruction of Particle Tracks



2000 μm, pixel pitch: 55 μm



Z

Charge carrier drift motion:

e⁻ and h⁺ drift described by

$$v_e = -\mu_e \times E(z)$$

 $v_h = \mu_h \times E(z)$

 $\mu_{e/h}$: Mobility of e⁻/h⁺

Electric field parametrisation:

Si:
$$\vec{E}(z) = \frac{U_B}{d} \vec{e_z} + \frac{2U_{dep}}{d^2} \left(\frac{d}{2} - z\right) \vec{e_z}$$
;

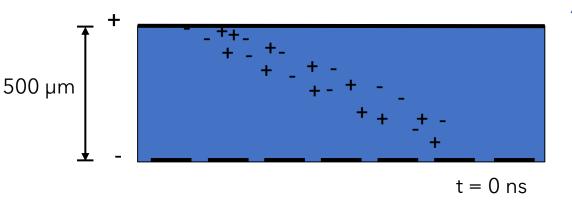
CdTe:
$$\vec{E}(z) = \frac{U_B}{d} \vec{e_z}$$

 U_B : Bias voltage; U_{dep} : Depletion voltage; d: Sensor thickness

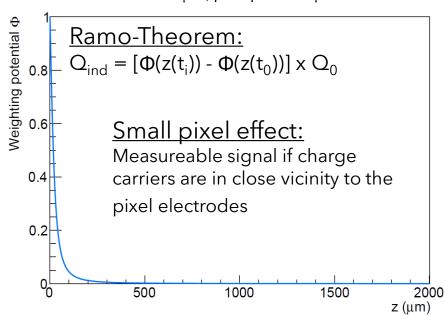
 \rightarrow Look up table: $z(t_{meas.}, E_{meas.})$

Bergmann et al. Eur. Phys. J. C (2017) 77: 421. https://doi.org/10.1140/epjc/s10052-017-4993-4

3D Reconstruction of Particle Tracks



2000 μm, pixel pitch: 55 μm



Z

Charge carrier drift motion:

e⁻ and h⁺ drift described by

$$v_e = -\mu_e \times E(z)$$

 $v_h = \mu_h \times E(z)$

 $\mu_{e/h}$: Mobility of e⁻/h⁺

Electric field parametrisation:

Si:
$$\vec{E}(z) = \frac{U_B}{d} \vec{e_z} + \frac{2U_{dep}}{d^2} \left(\frac{d}{2} - z\right) \vec{e_z}$$
;

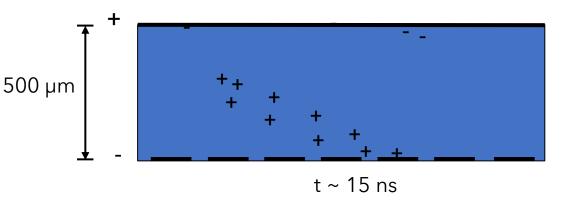
CdTe:
$$\vec{E}(z) = \frac{U_B}{d} \vec{e_z}$$

 U_B : Bias voltage; U_{dep} : Depletion voltage; d: Sensor thickness

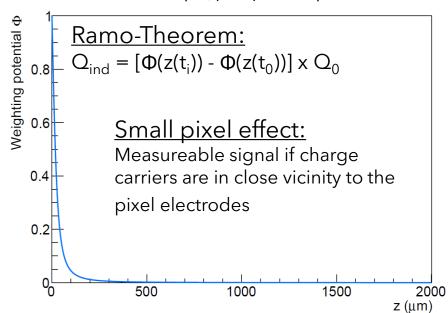
 \rightarrow Look up table: $z(t_{meas.}, E_{meas.})$

Bergmann et al. Eur. Phys. J. C (2017) 77: 421. https://doi.org/10.1140/epjc/s10052-017-4993-4

3D Reconstruction of Particle Tracks



2000 μm, pixel pitch: 55 μm



Z

Charge carrier drift motion:

e⁻ and h⁺ drift described by

$$v_e = -\mu_e \times E(z)$$

 $v_h = \mu_h \times E(z)$

 $\mu_{e/h}$: Mobility of e⁻/h⁺

Electric field parametrisation:

Si:
$$\vec{E}(z) = \frac{U_B}{d} \vec{e_z} + \frac{2U_{dep}}{d^2} \left(\frac{d}{2} - z\right) \vec{e_z}$$
;

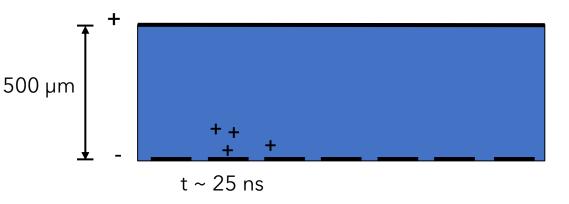
CdTe:
$$\vec{E}(z) = \frac{U_B}{d} \vec{e_z}$$

 U_B : Bias voltage; U_{dep} : Depletion voltage; d: Sensor thickness

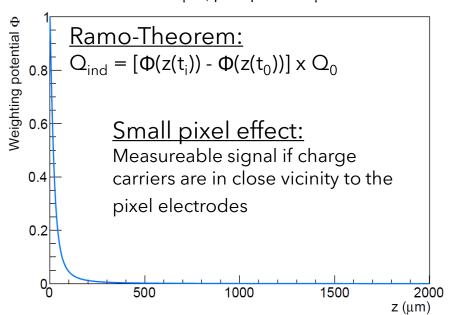
 \rightarrow Look up table: $z(t_{meas.}, E_{meas.})$

Bergmann et al. Eur. Phys. J. C (2017) 77: 421. https://doi.org/10.1140/epjc/s10052-017-4993-4

3D Reconstruction of Particle Tracks



2000 μm, pixel pitch: 55 μm



↑ Z

Charge carrier drift motion:

e⁻ and h⁺ drift described by

$$v_e = -\mu_e \times E(z)$$

 $v_h = \mu_h \times E(z)$

 $\mu_{e/h}$: Mobility of e⁻/h⁺

Electric field parametrisation:

Si:
$$\vec{E}(z) = \frac{U_B}{d} \vec{e_z} + \frac{2U_{dep}}{d^2} \left(\frac{d}{2} - z\right) \vec{e_z}$$
;

CdTe:
$$\vec{E}(z) = \frac{U_B}{d} \vec{e_z}$$

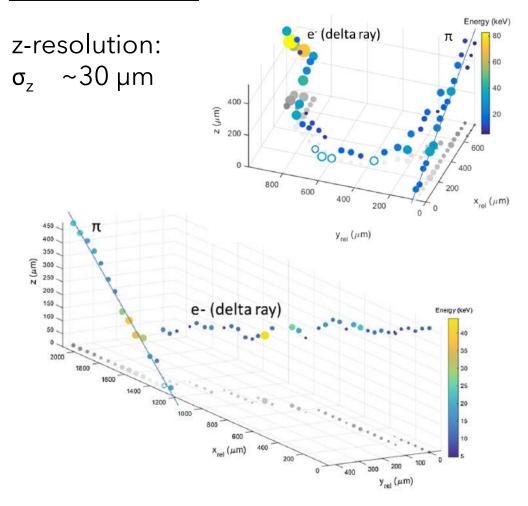
 U_B : Bias voltage; U_{dep} : Depletion voltage; d: Sensor thickness

 \rightarrow Look up table: $z(t_{meas.}, E_{meas.})$

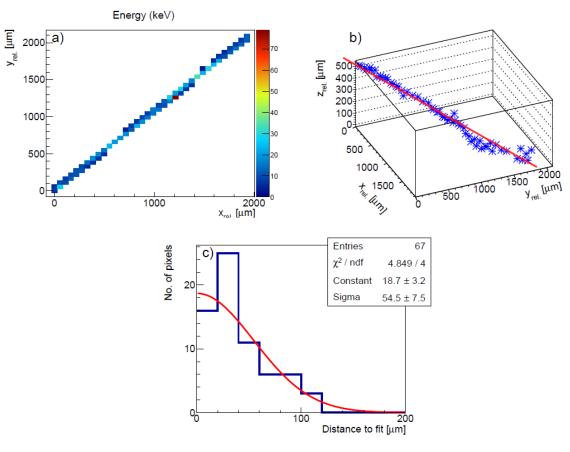
Bergmann et al. Eur. Phys. J. C (2017) 77: 421. https://doi.org/10.1140/epjc/s10052-017-4993-4

3D Tracking in 500 µm Si Timepix3

120 GeV/c pion tracks accompanied by δ-rays measured at SPS:

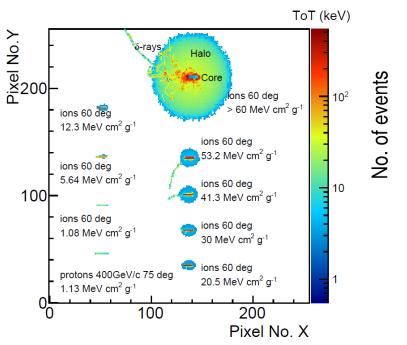


A cosmic µ measured in the Prague laboratory:



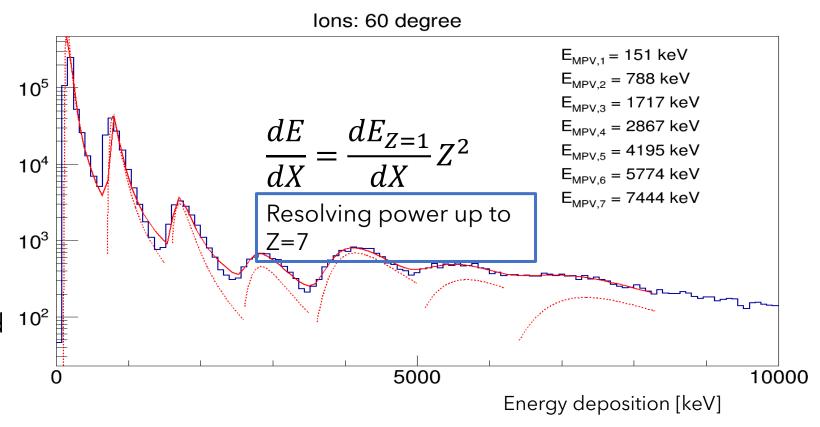
Bergmann et al. Eur. Phys. J. C (2017) 77: 421. https://doi.org/10.1140/epjc/s10052-017-4993-4

Energy Deposition Spectra of Charged Energetic Particles



Different track shapes observed in a mixed relativistic ion beam 330 GeV/c Pb on target

Energy loss of energetic charged particles described by Landau-Vavilov distribution with Gaussian smearing

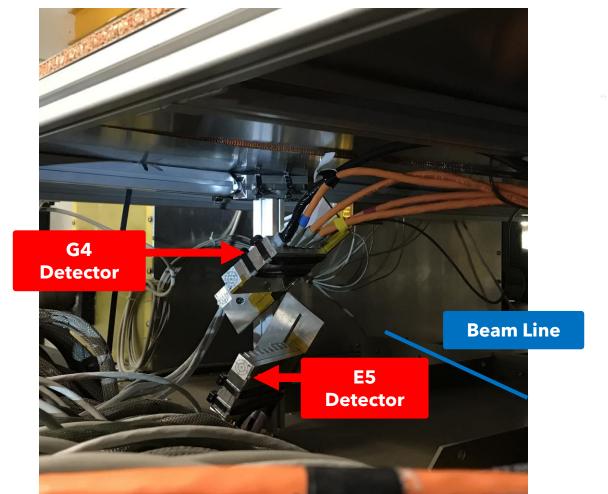


P. Smolyanskiy *et al. JINST* **16** P01022, 2021. https://dx.doi.org/10.1088/1748-0221/16/01/P01022

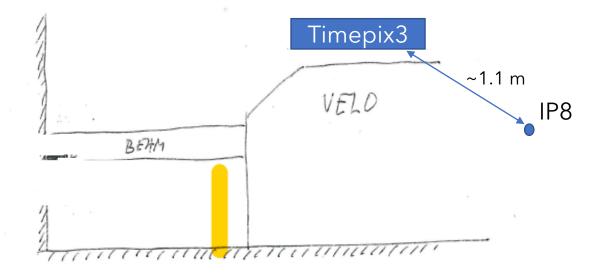
Single-Layer Particle Tracking in MoEDAL

Timepix3 in MoEDAL at Run 2 - Feasibility Study

Installation of two Timepix3 detectors in MoEDAL in **September 2018**. Timepix3 are placed at 1.1 m distance to IP8 with a relatively unobstructed view



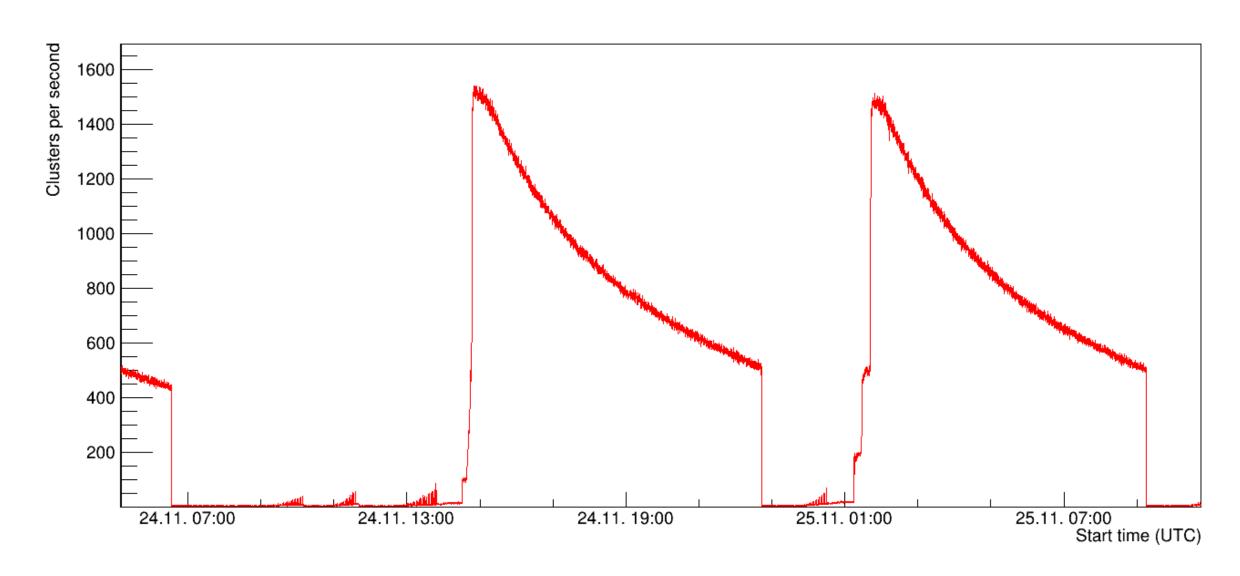
MOEDAL TPX



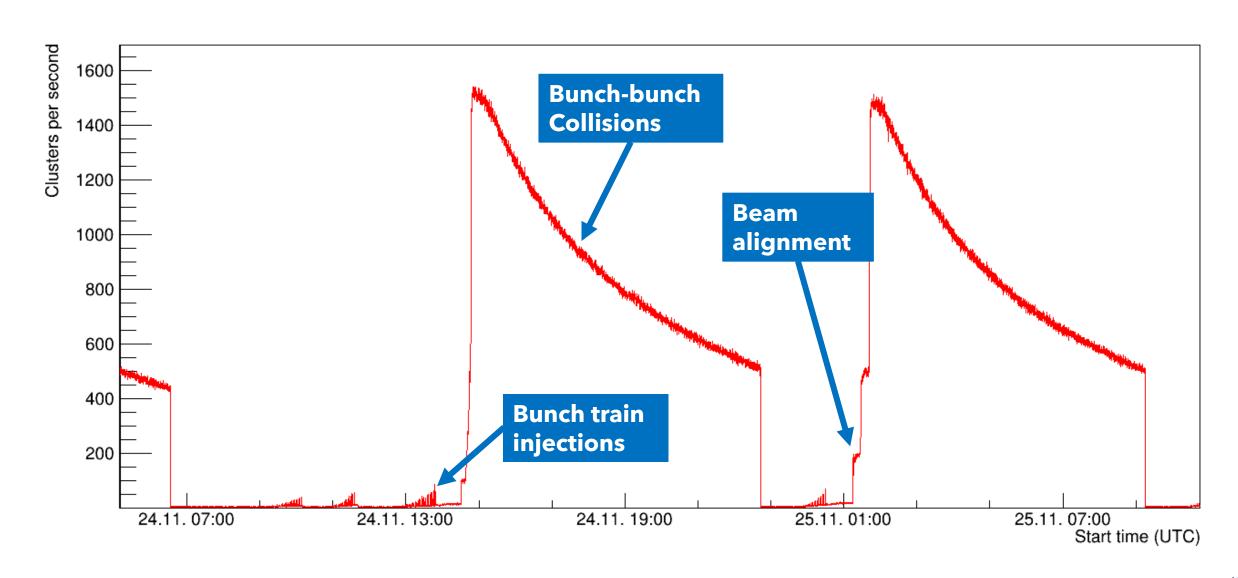
Continuous quasi dead-time free measurement (in real time) keeping a permanent record of **all particle traces**

- Tracking and identification of all particles
- Online outlier detection to search for exotics (highly ionising events)
- Bunch-by-bunch luminosity measurement
- Search for exotic signatures, e.g., "soup" events requiring timing information

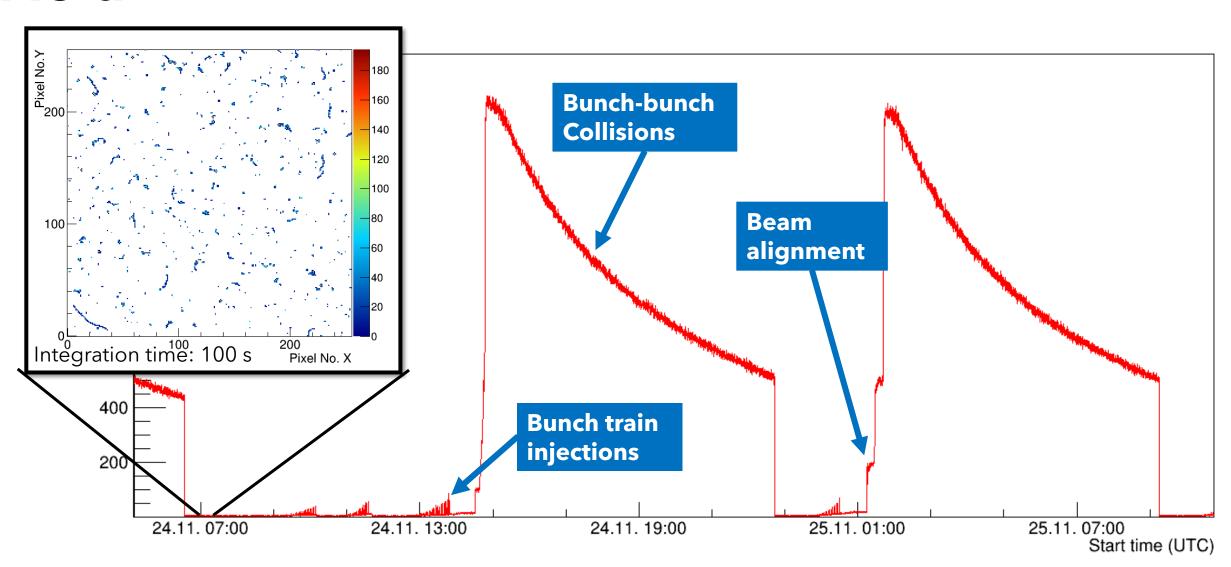
Radiation Levels and Radiation Field

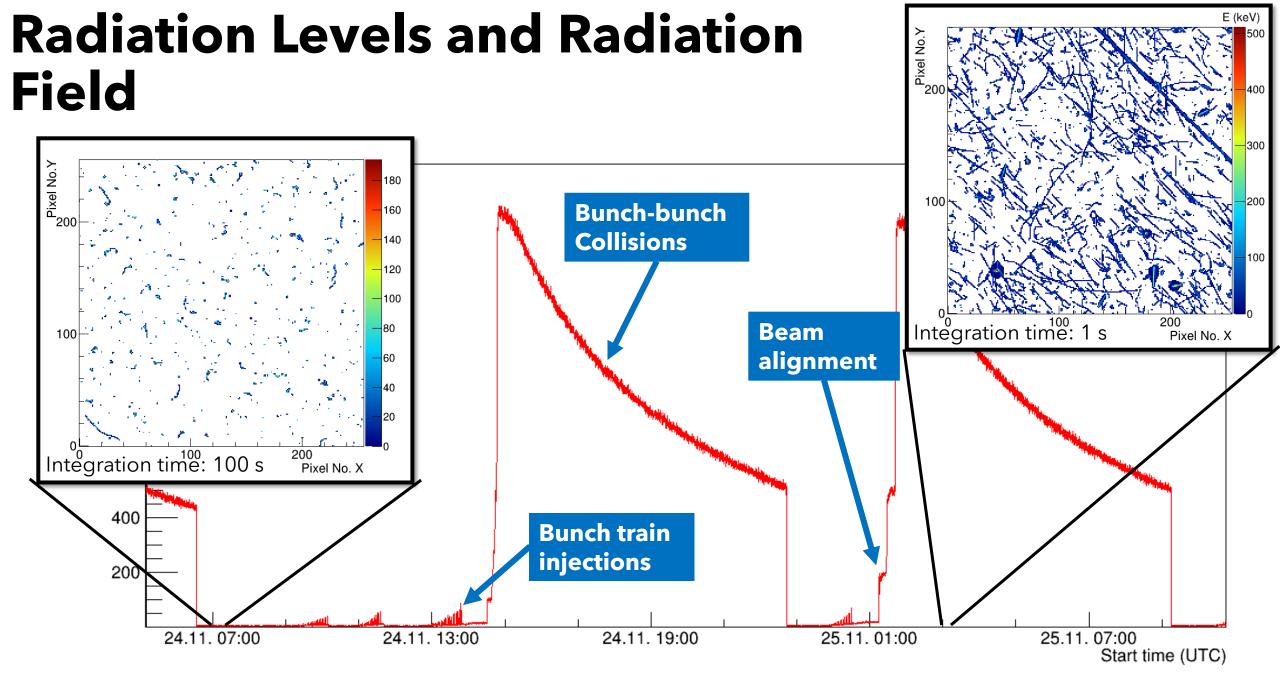


Radiation Levels and Radiation Field

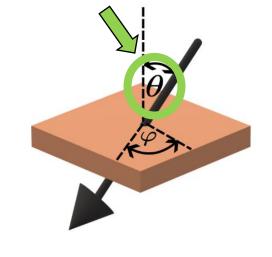


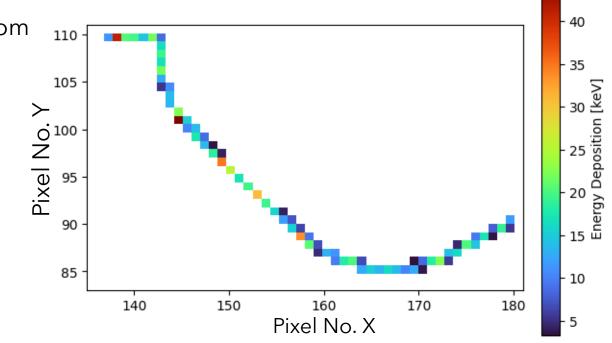
Radiation Levels and Radiation Field





- To reduce the complexity of the problem, only the following two incidence angles are predicted: Azimuthal and perpendicular. Reducing it to regression problem
- Datasets used for model development:
 - ⇒ 0.01-10 MeV electrons (curly tracks)
 - ⇒ 10-500 MeV protons (thick, straight tracks)
 - 40 GeV pions (thin, straight tracks)
- It was found through an extensive study that a Random Forrest Regressor with a selection of input features produced optimal results

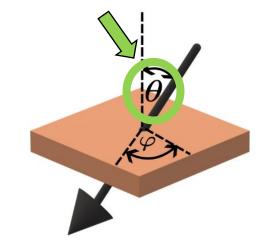


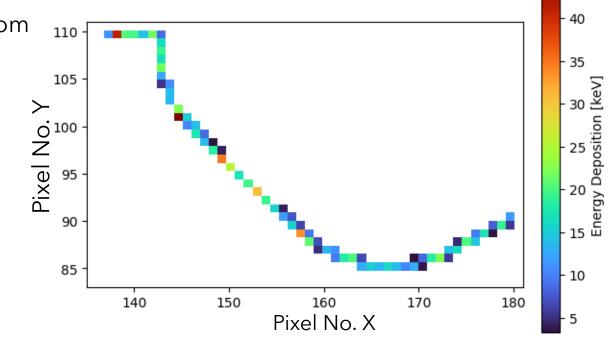


- To reduce the complexity of the problem, only the following two incidence angles are predicted: Azimuthal and perpendicular. Reducing it to regression problem
- Datasets used for model development:
 - ⇒ 0.01-10 MeV electrons (curly tracks)
 - ⇒ 10-500 MeV protons (thick, straight tracks)
 - 40 GeV pions (thin, straight tracks)
- It was found through an extensive study that a Random Forrest Regressor with a selection of input features produced optimal results

Input Features:

1. Size [no. of pixels]



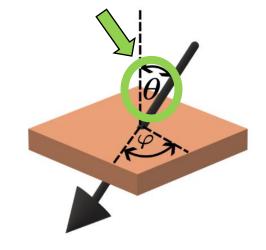


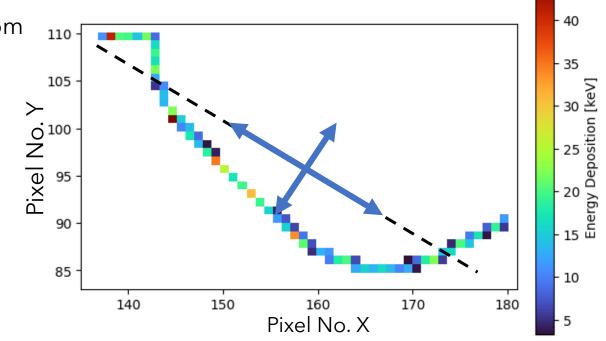
 To reduce the complexity of the problem, only the following two incidence angles are predicted: Azimuthal and perpendicular. Reducing it to regression problem



- ⇒ 0.01-10 MeV electrons (curly tracks)
- ⇒ 10-500 MeV protons (thick, straight tracks)
- 40 GeV pions (thin, straight tracks)
- It was found through an extensive study that a Random Forrest Regressor with a selection of input features produced optimal results

- 1. Size [no. of pixels]
- 2. Line standard deviation



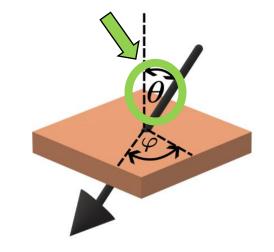


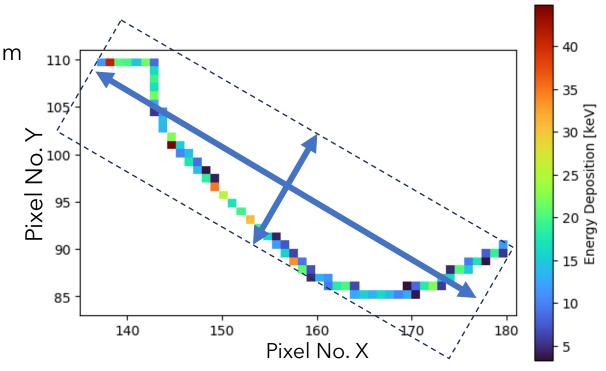
 To reduce the complexity of the problem, only the following two incidence angles are predicted: Azimuthal and perpendicular. Reducing it to regression problem



- ⇒ 0.01-10 MeV electrons (curly tracks)
- ⇒ 10-500 MeV protons (thick, straight tracks)
- ⇒ 40 GeV pions (thin, straight tracks)
- It was found through an extensive study that a Random Forrest Regressor with a selection of input features produced optimal results

- 1. Size [no. of pixels]
- 2. Line standard deviation
- 3. Box dimensions



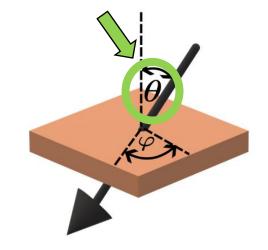


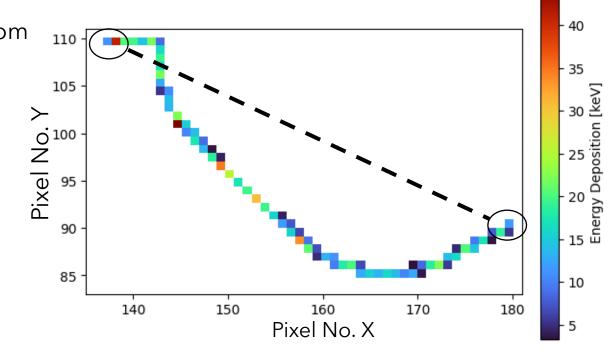
 To reduce the complexity of the problem, only the following two incidence angles are predicted: Azimuthal and perpendicular. Reducing it to regression problem



- ⇒ 0.01-10 MeV electrons (curly tracks)
- ⇒ 10-500 MeV protons (thick, straight tracks)
- 40 GeV pions (thin, straight tracks)
- It was found through an extensive study that a Random Forrest Regressor with a selection of input features produced optimal results

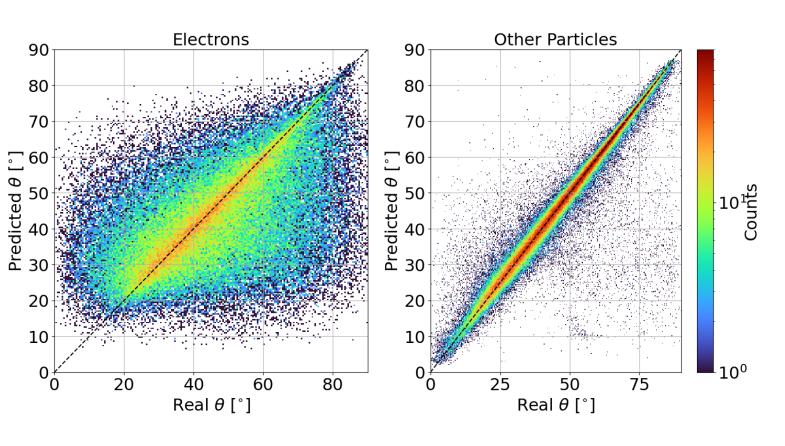
- 1. Size [no. of pixels]
- 2. Line standard deviation
- 3. Box dimensions
- 4. End point distance

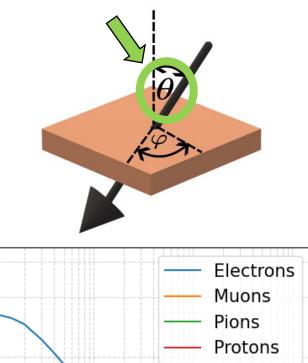


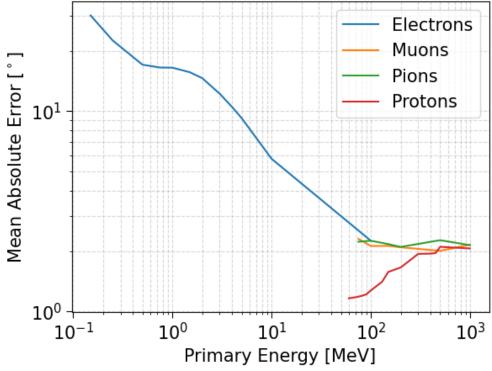


Theta Prediction Results

For testing separate simulation datasets were created

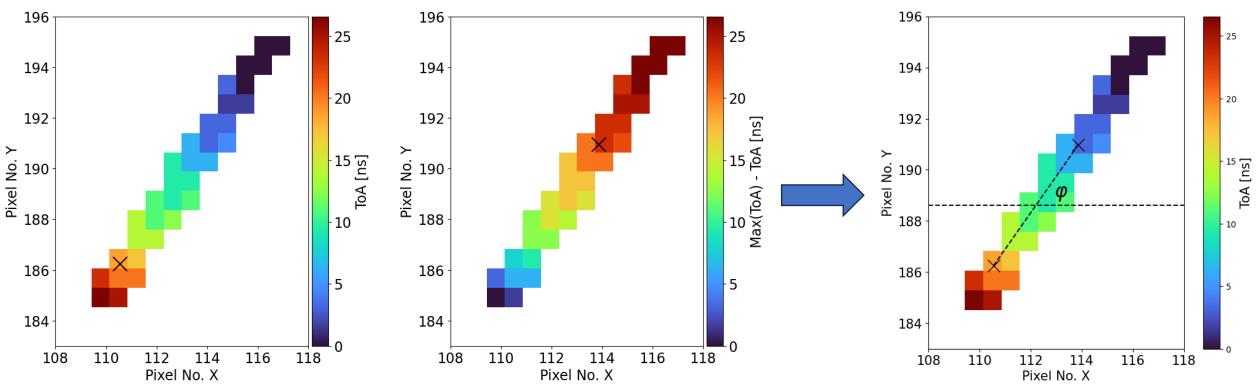


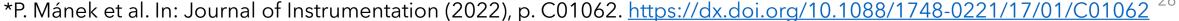




Phi Determination

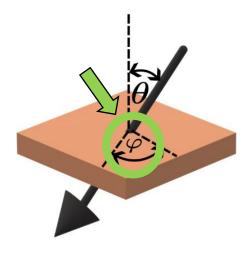
- An analytic approach was used as no improvement using machine learning algorithms was achieved
- Following an extensive study, Time of Arrival [ToA] weighted averaging was found to produce the most accurate results*
- In this method, we utilise the effect that drift time within the sensor has on ToA to approximate the "mean" entry and exit points

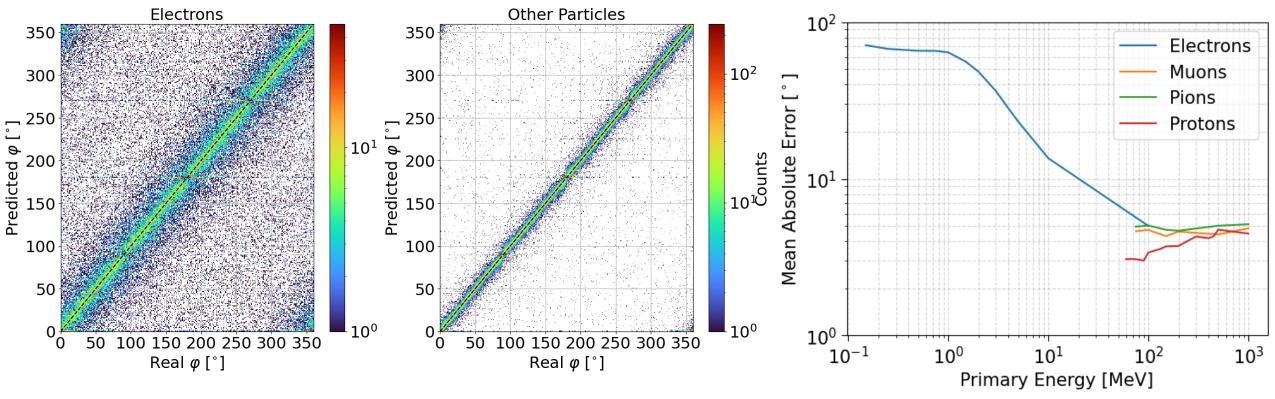




Phi Model Results

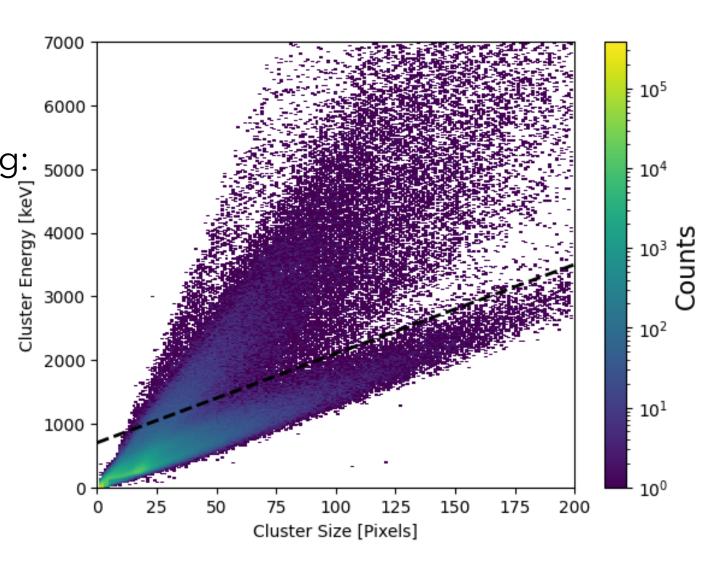
For testing separate simulation datasets were created





Rudimentary Classification

Particle tracks can be divided into two groups, with variable linearity and scattering:

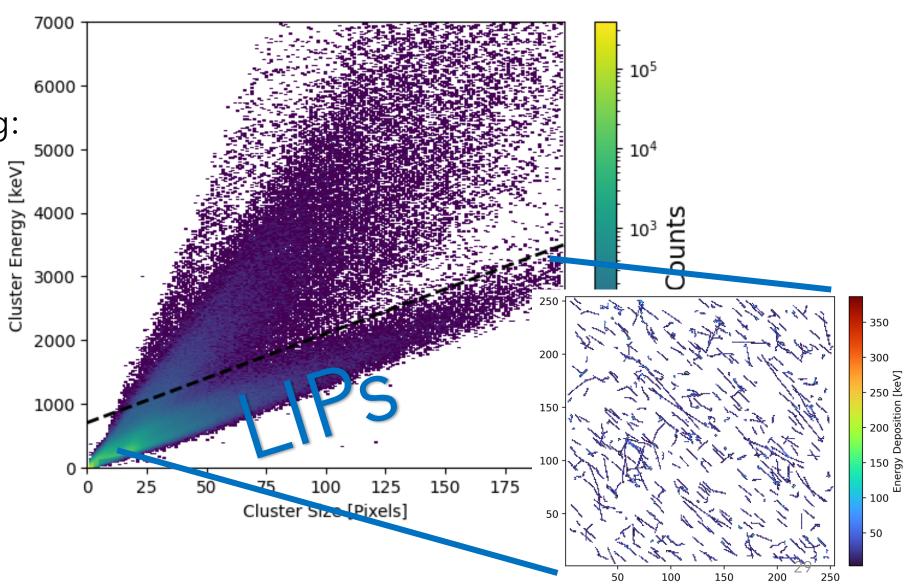


Rudimentary Classification

Particle tracks can be divided into two groups, with variable linearity and scattering:

1. Lowly Ionising Particles (LIPs):

- ⇒Electrons
- ⇒Photons
- Relativistic singly charged particles



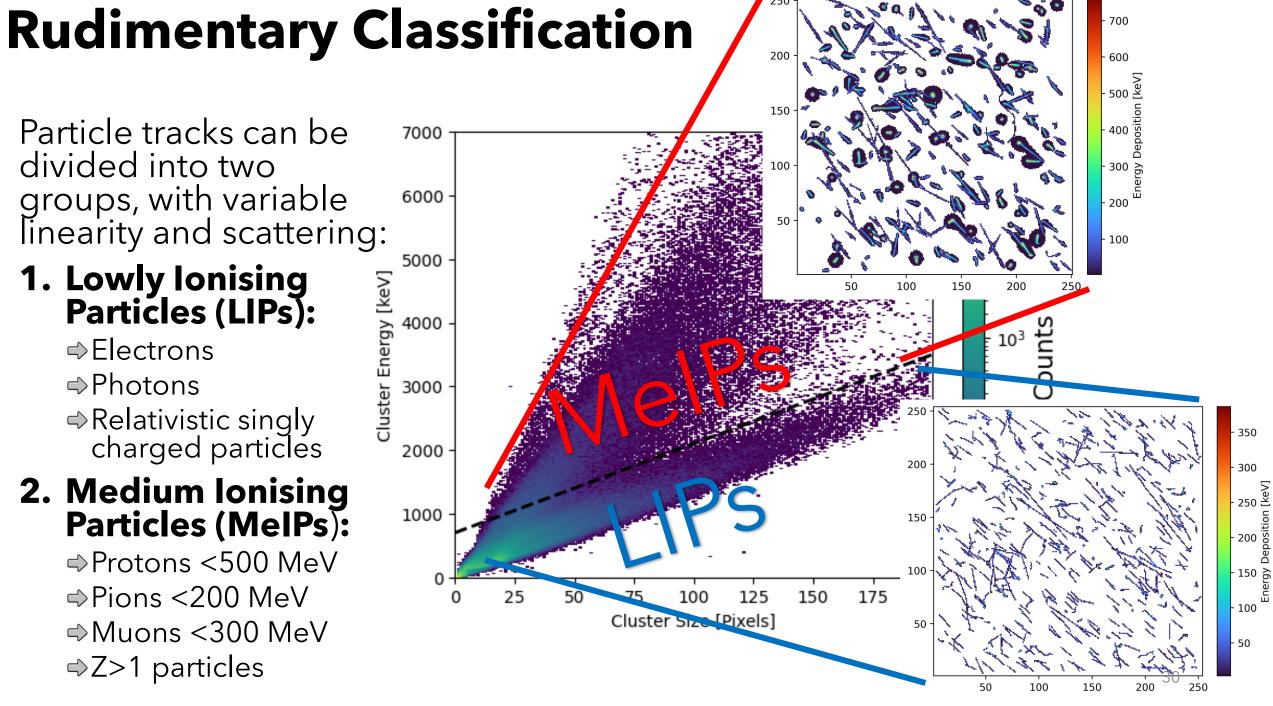
Particle tracks can be divided into two groups, with variable linearity and scattering:

1. Lowly lonising Particles (LIPs):

- ⇒Electrons
- ⇒ Photons
- Relativistic singly charged particles

2. Medium Ionising Particles (MelPs):

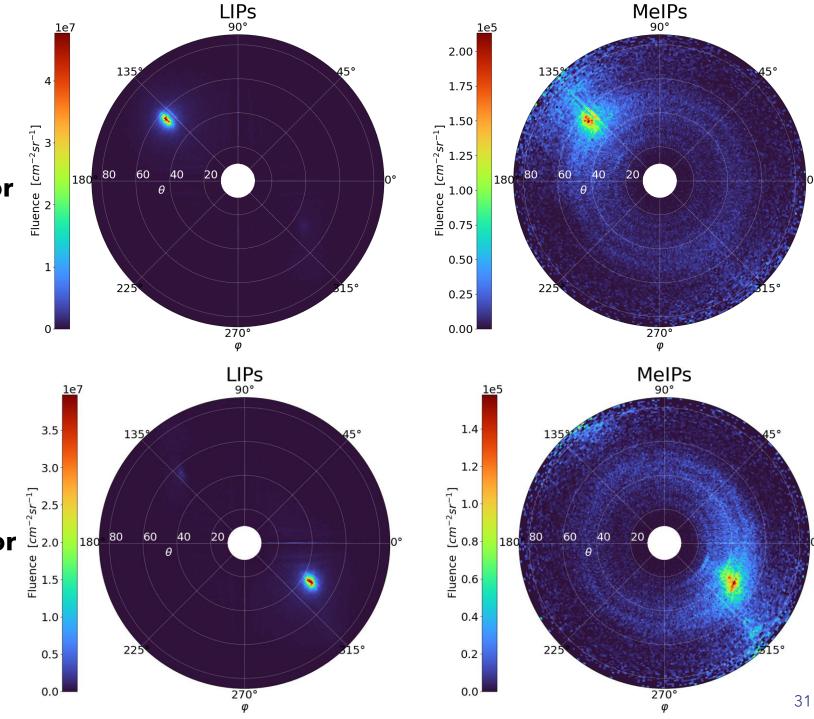
- ⇒Protons <500 MeV
- ⇒Pions <200 MeV
- ⇒Muons <300 MeV
- ⇒Z>1 particles



E5 Detector

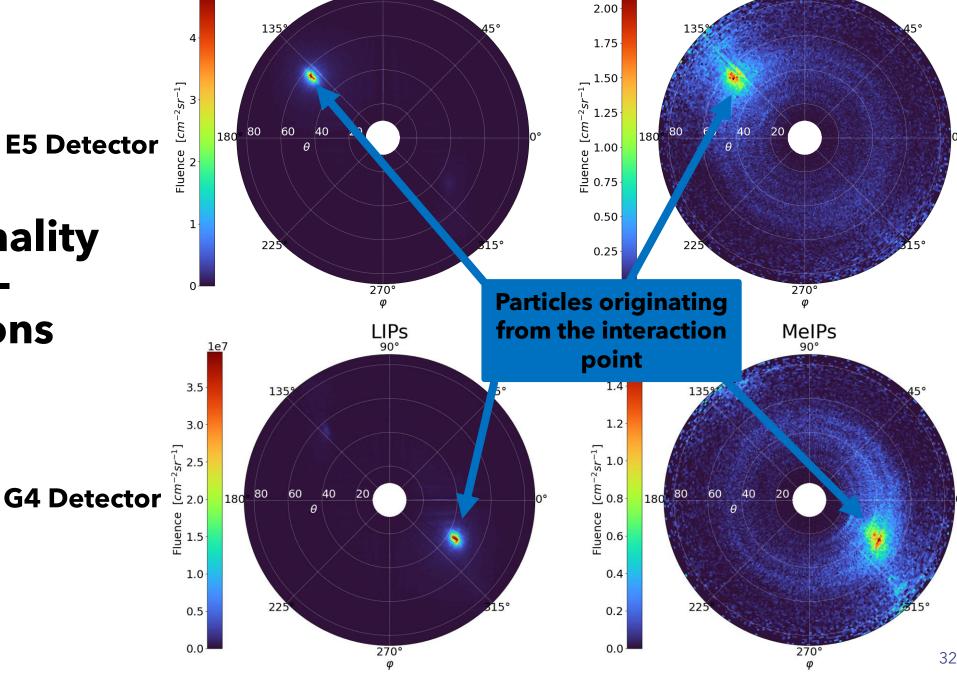
Field Directionality During Proton-Proton Collisions





E5 Detector

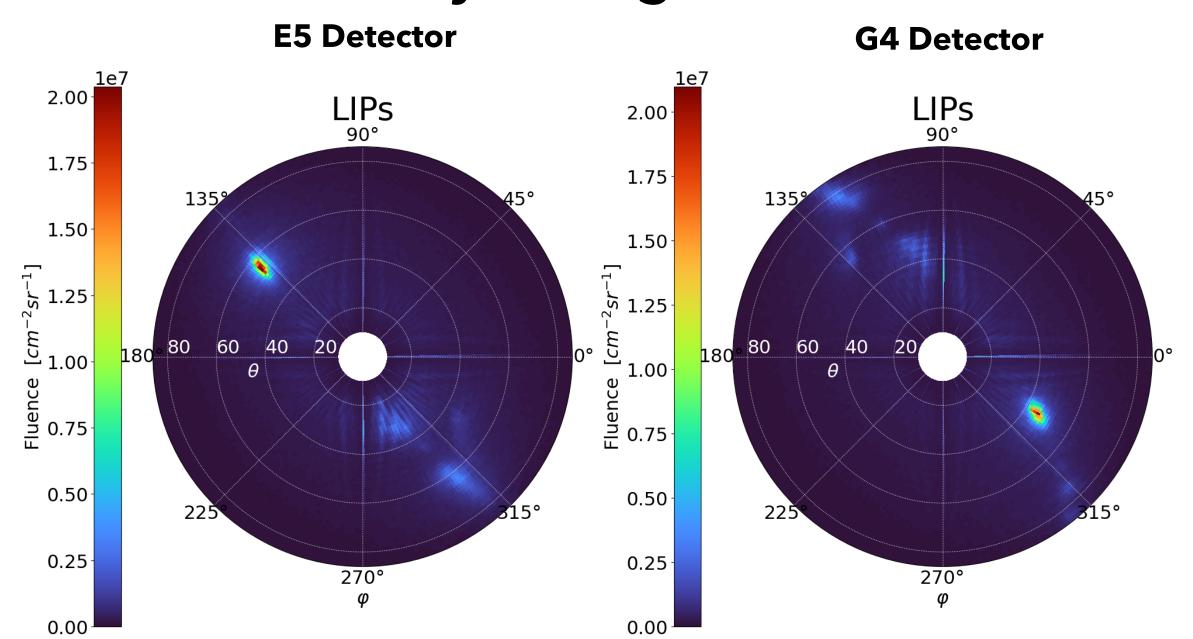
Field Directionality During Proton- Proton Collisions



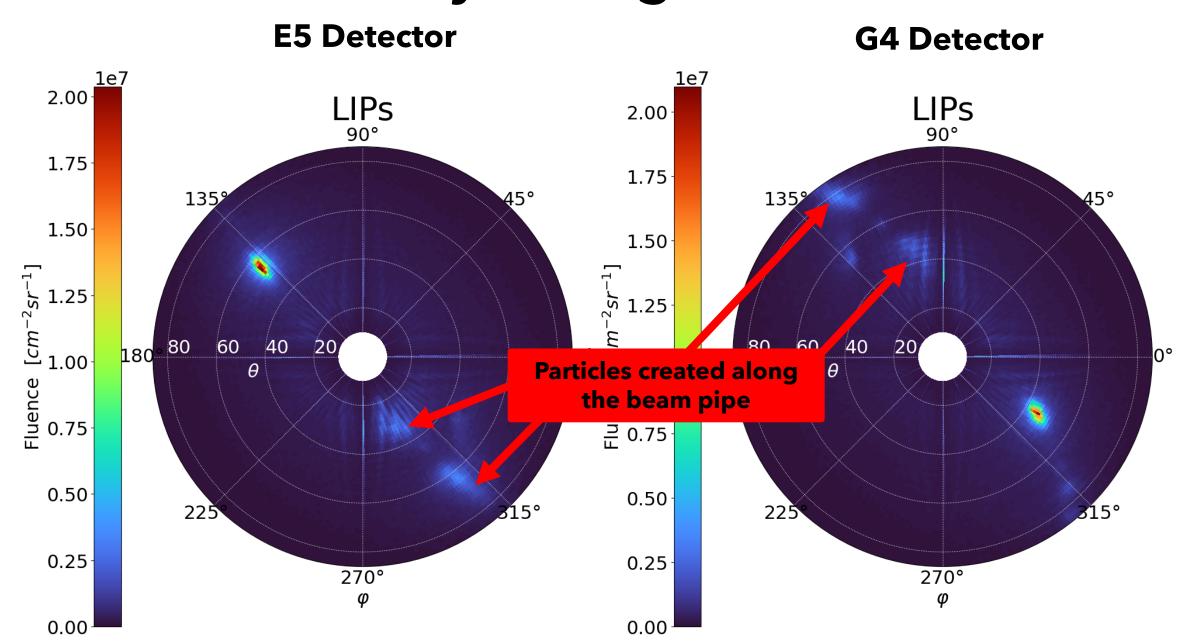
MeIPs

LIPs 90°

Field Directionality During Pb-Pb Collisions

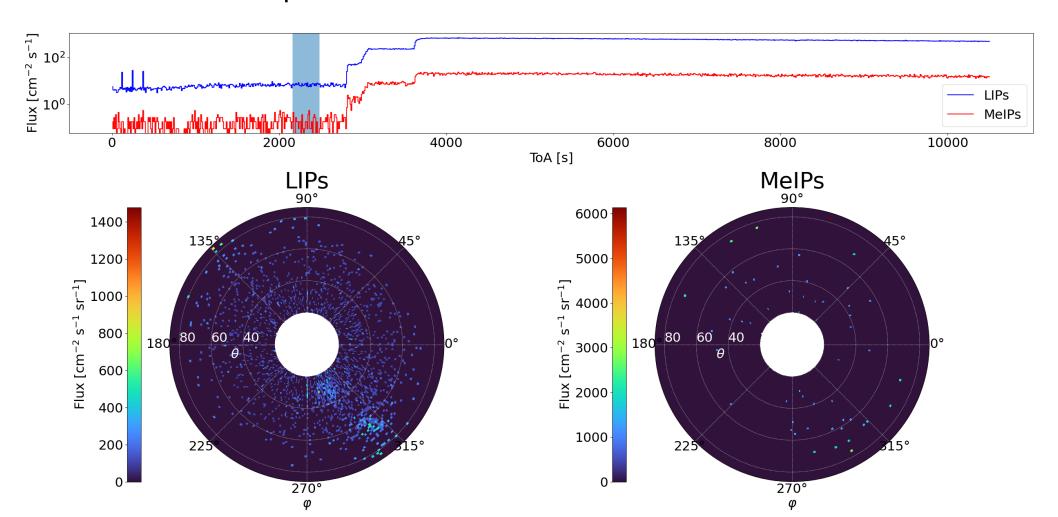


Field Directionality During Pb-Pb Collisions



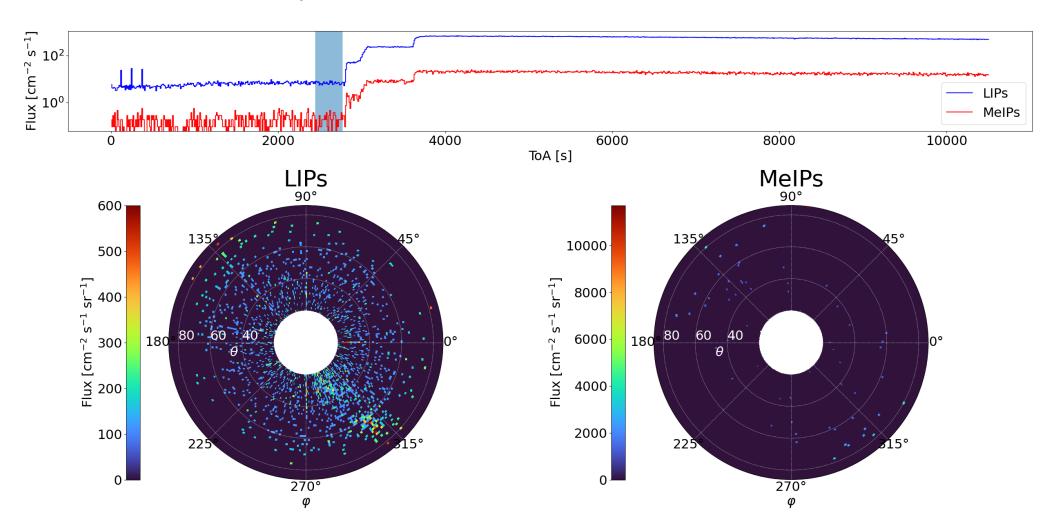
Time-Resolved Measurement of the Directionality Map Pb-Pb Collision Period

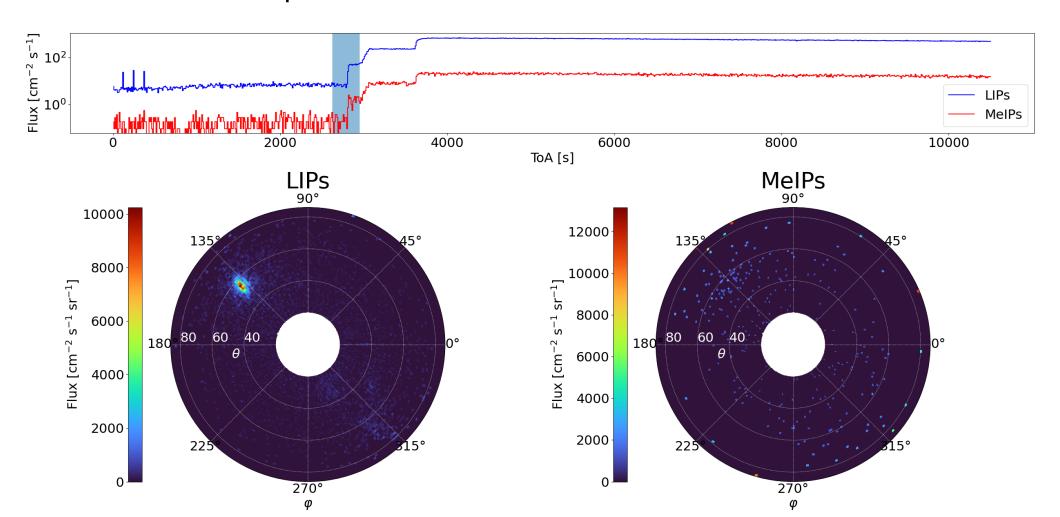
For increased accuracy only 5 minutes of data acquisition is need to produce reasonable statistics

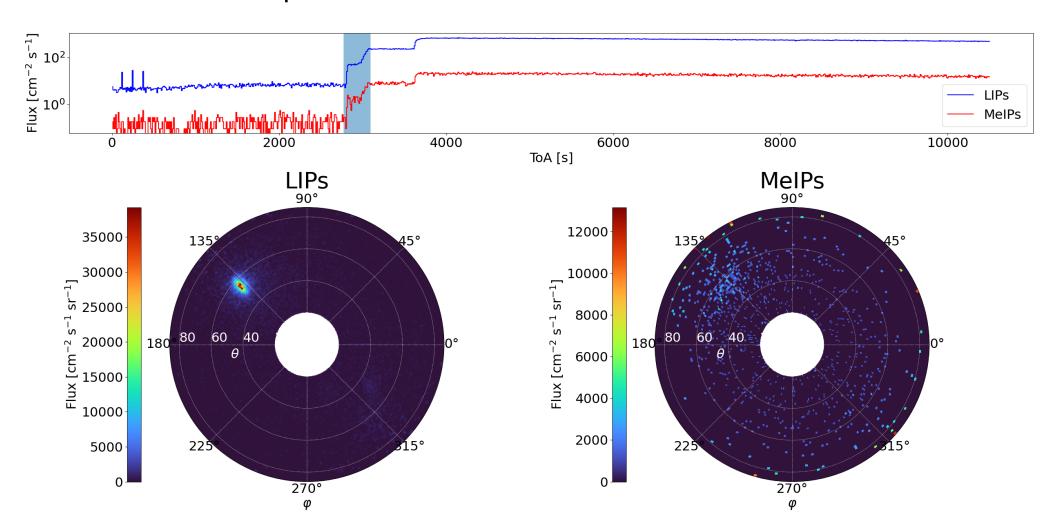


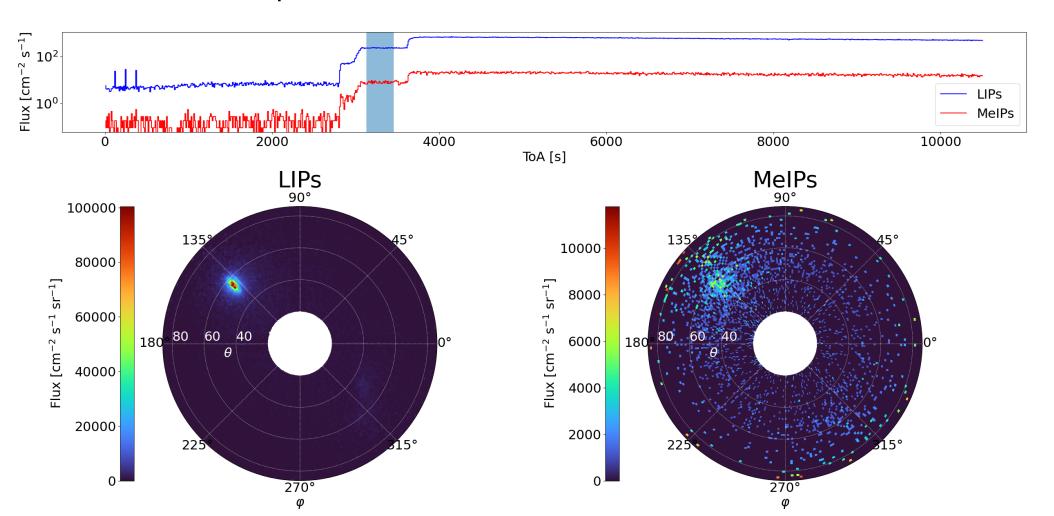
Time-Resolved Measurement of the Directionality Map Pb-Pb Collision Period

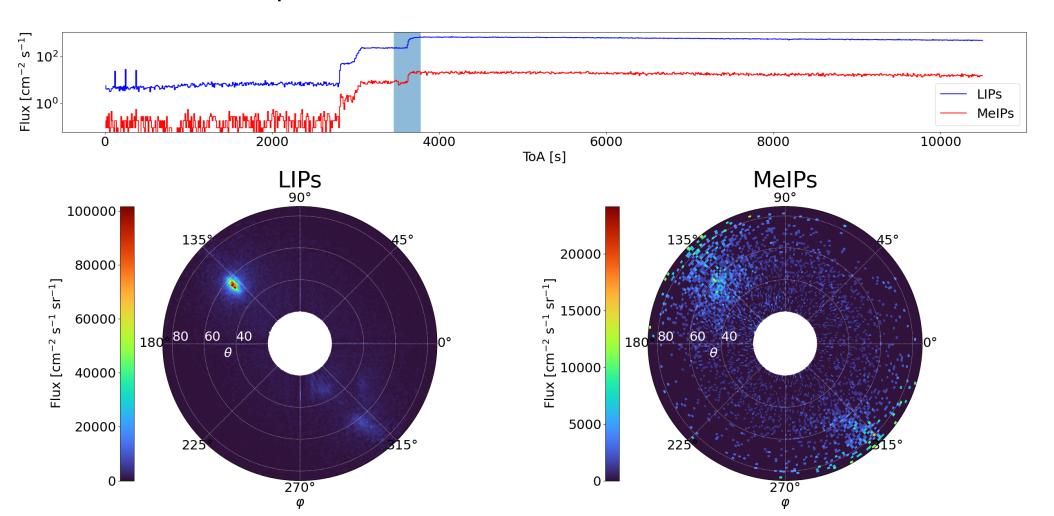
For increased accuracy only 5 minutes of data acquisition is need to produce reasonable statistics

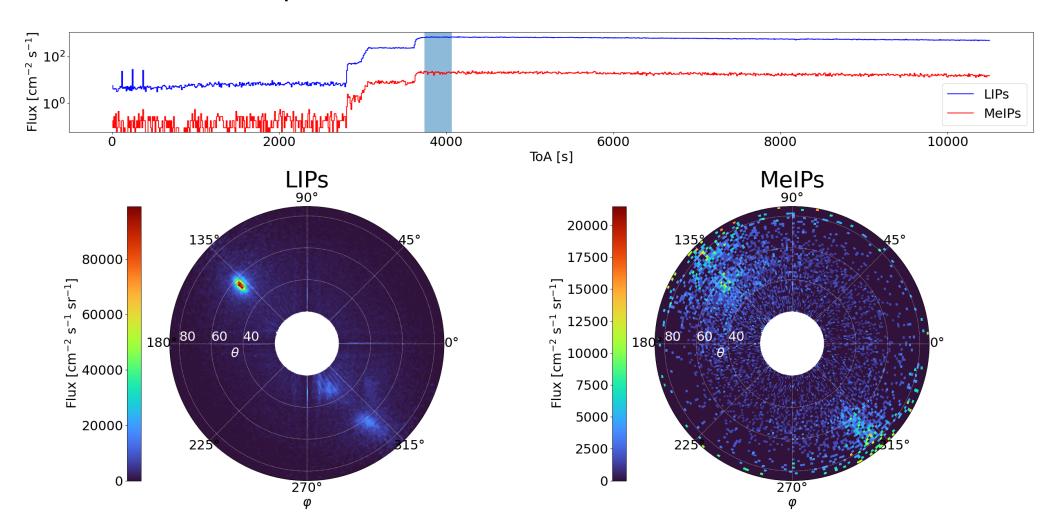












Radiation Field Decomposition According to Stopping Power

SATRAM And MOEDAL

A Brief Excursion: Radiation Field Decomposition in Space

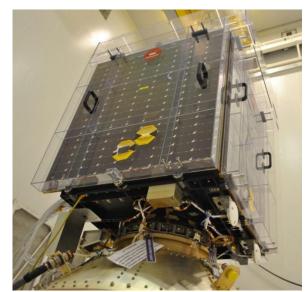
- Why the detour?
 - The mixed radiation field present in space has much lower complexity allowing us to more effectively develop and test algorithms for classification
- The data that will be used will come from the Space Application
 Timepix Radiation Monitor (SATRAM)
- The acquired knowledge can then be applied to the MoEDAL experiment

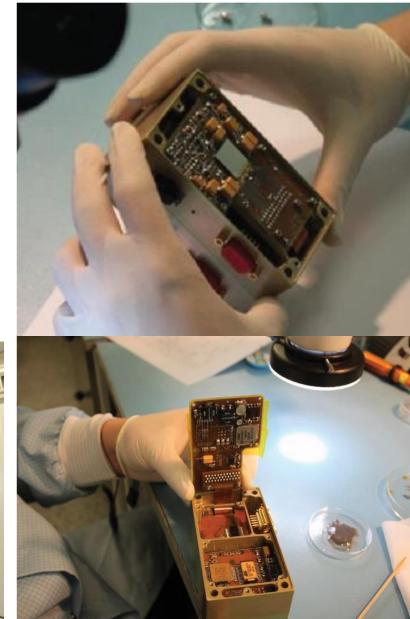
Space Application of Timepix Radiation Monitor (SATRAM)

- First Timepix in open space
- Power consumption of 2.5 W
- Total mass **380 g** (107 x 70 x 55 mm)

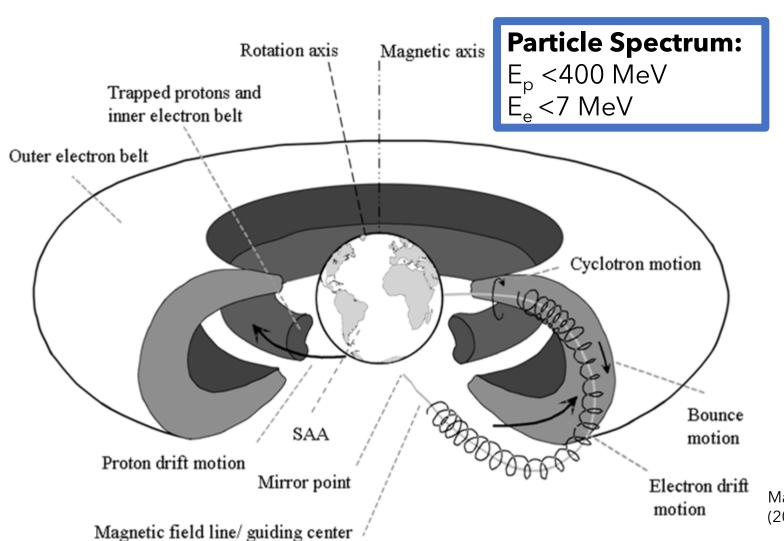
Proba-V

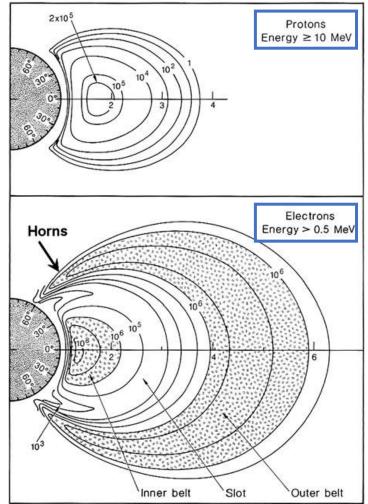
- Minisatellite (158 kg)
- Altitude ~820 km (LEO)
- 101.21 minutes orbit duration
- Inclination 98.6°
- Sun-synchronous
- Launched 7th March 2013





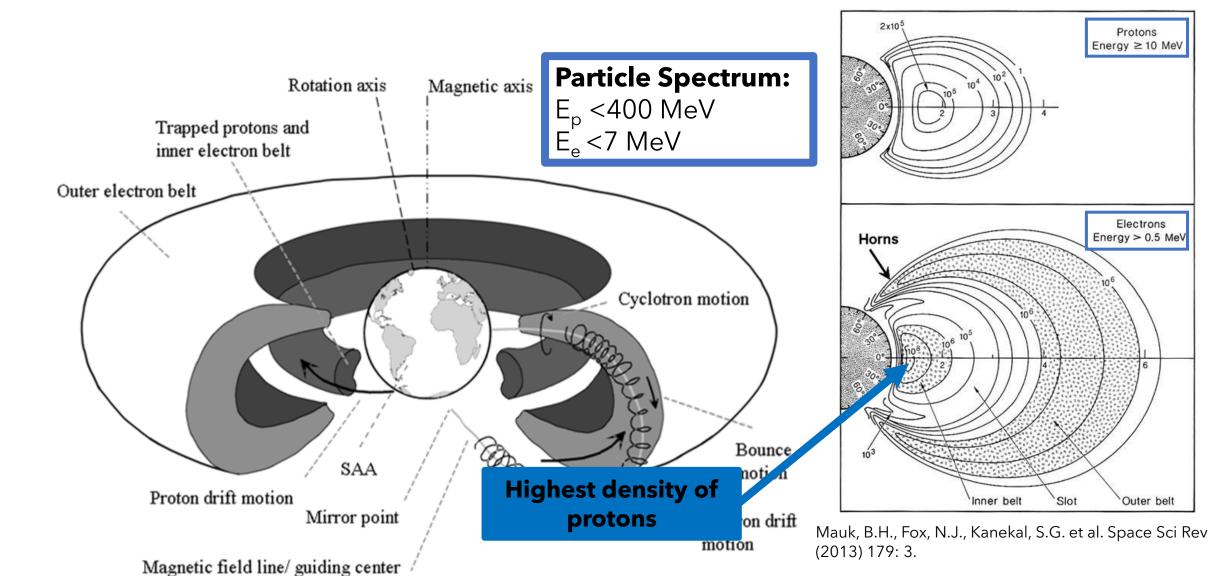
The Radiation Environment

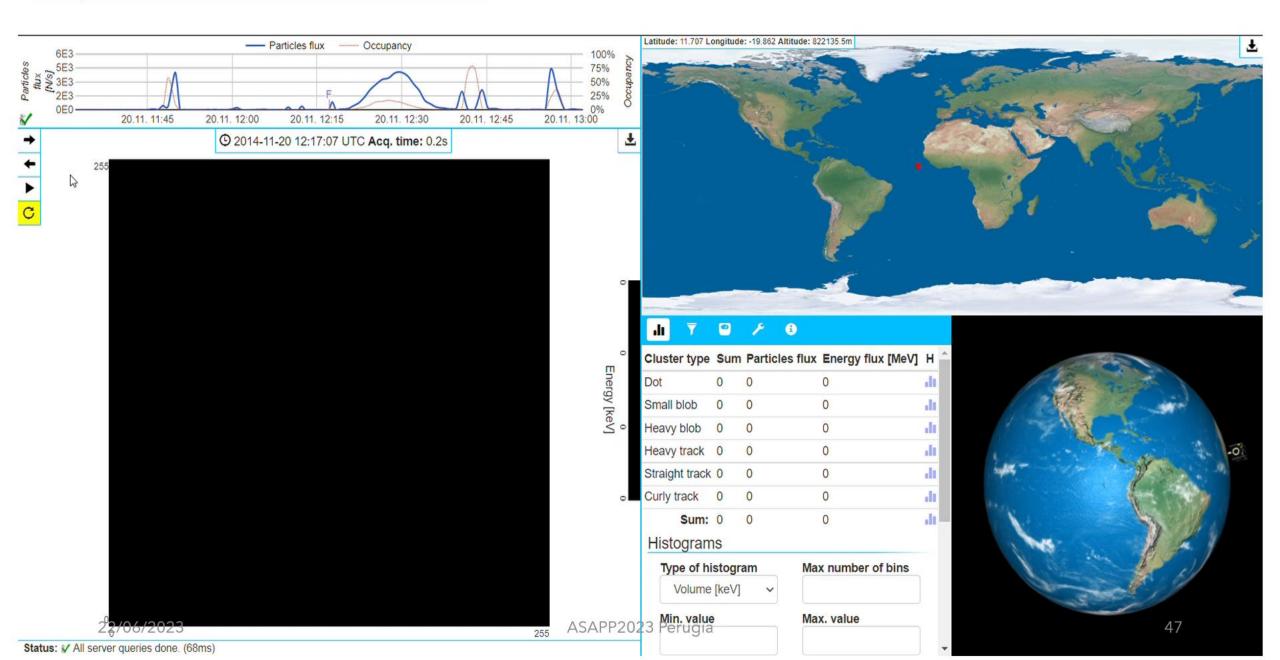


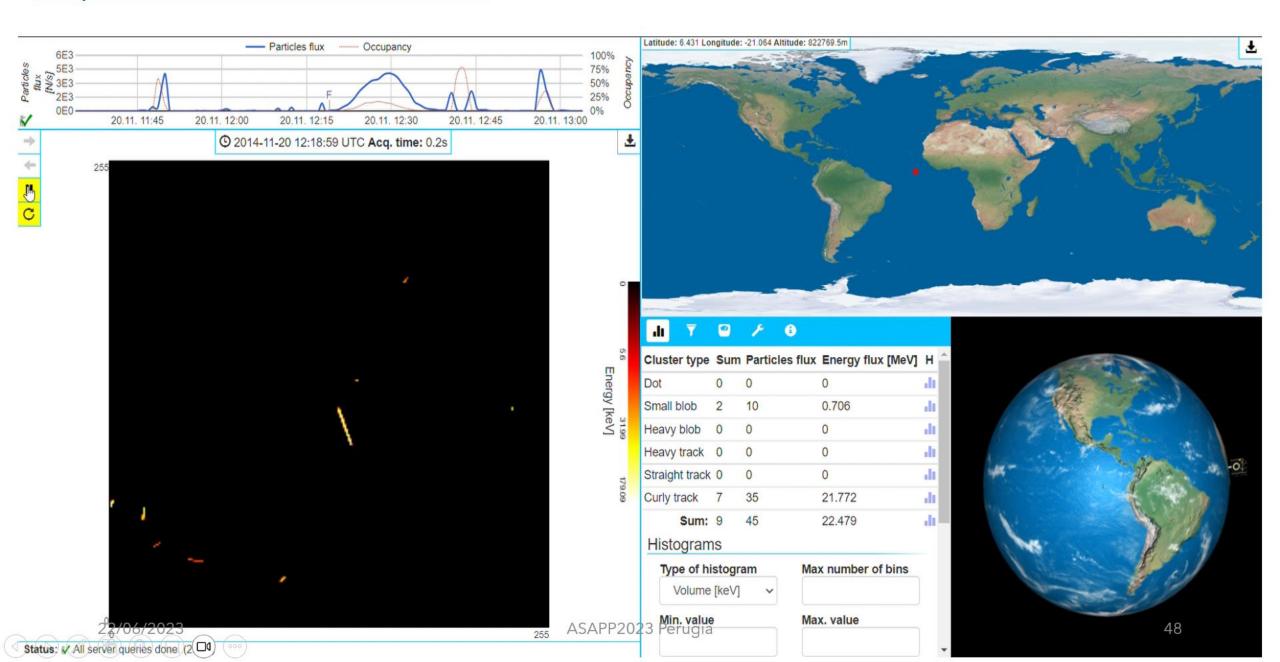


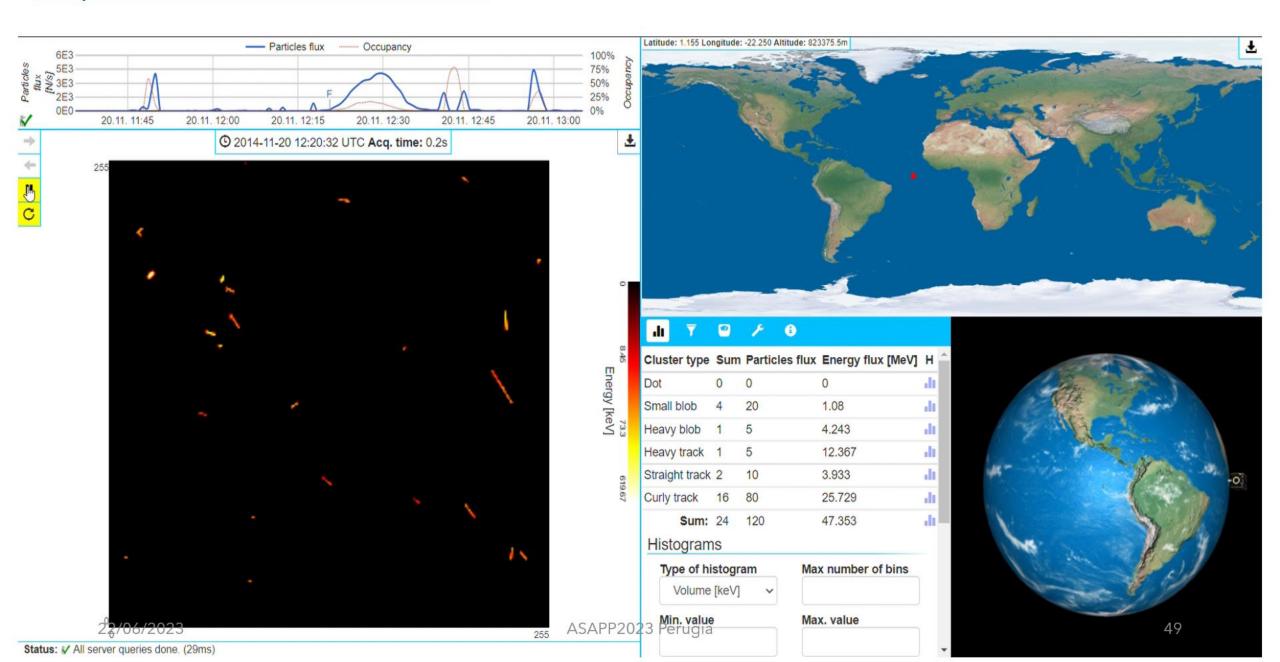
Mauk, B.H., Fox, N.J., Kanekal, S.G. et al. Space Sci Rev (2013) 179: 3.

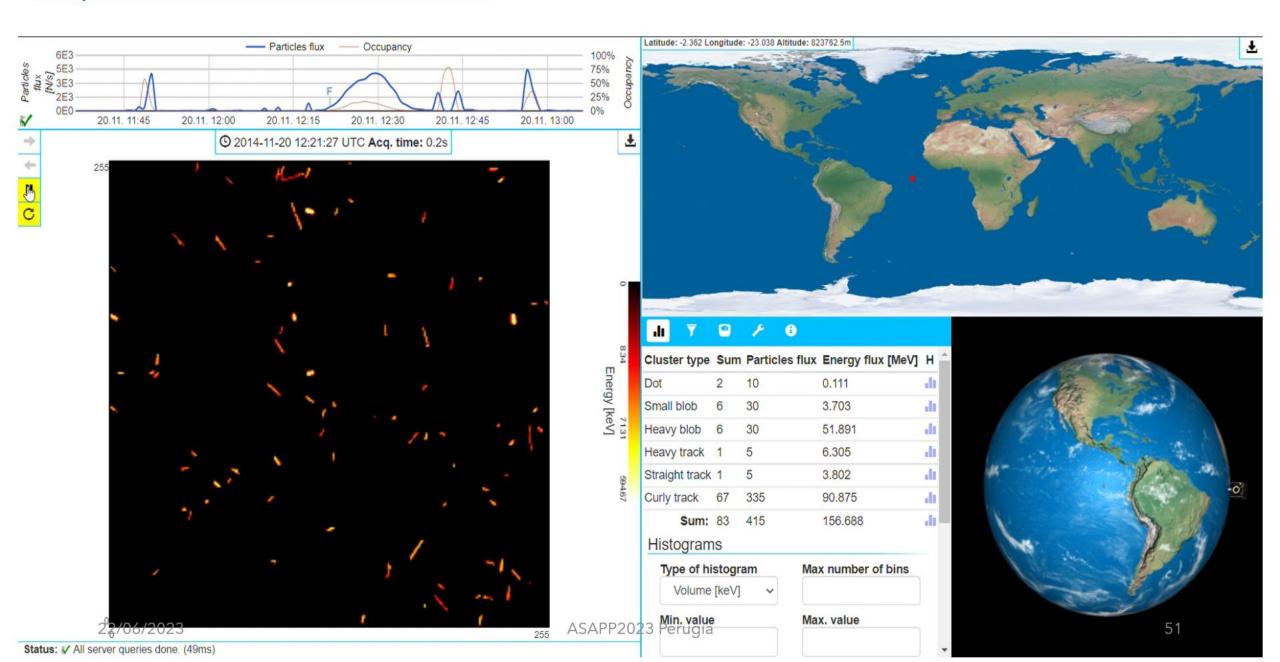
The Radiation Environment

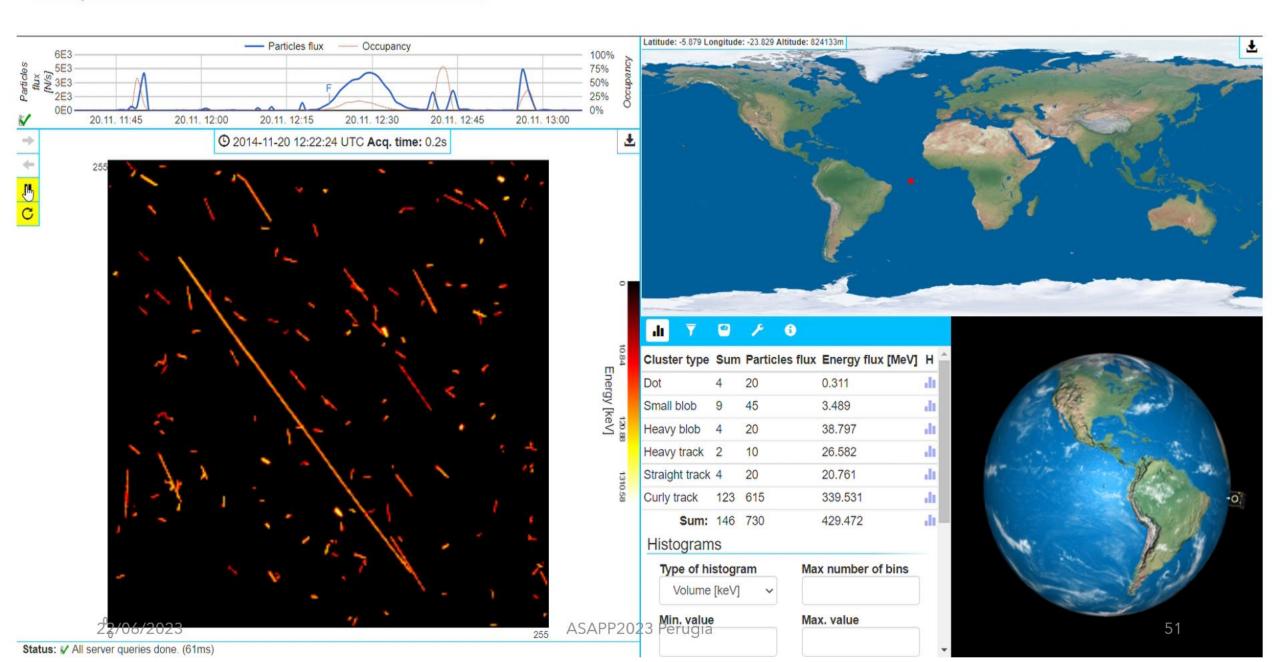


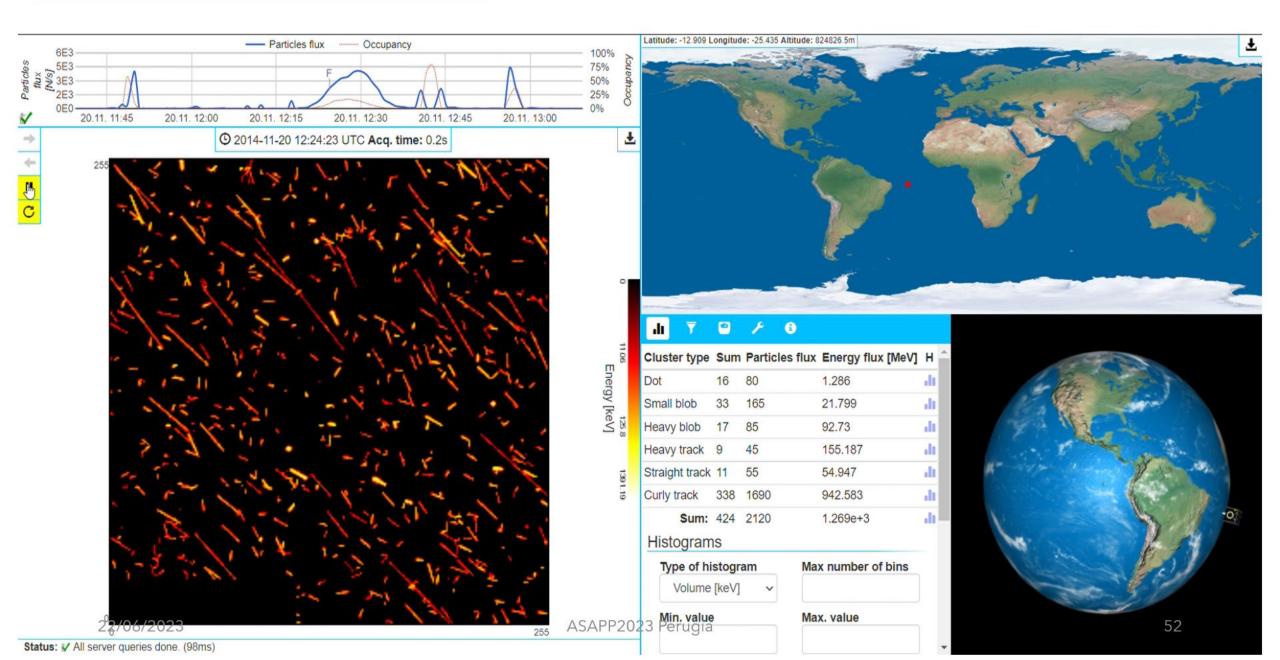


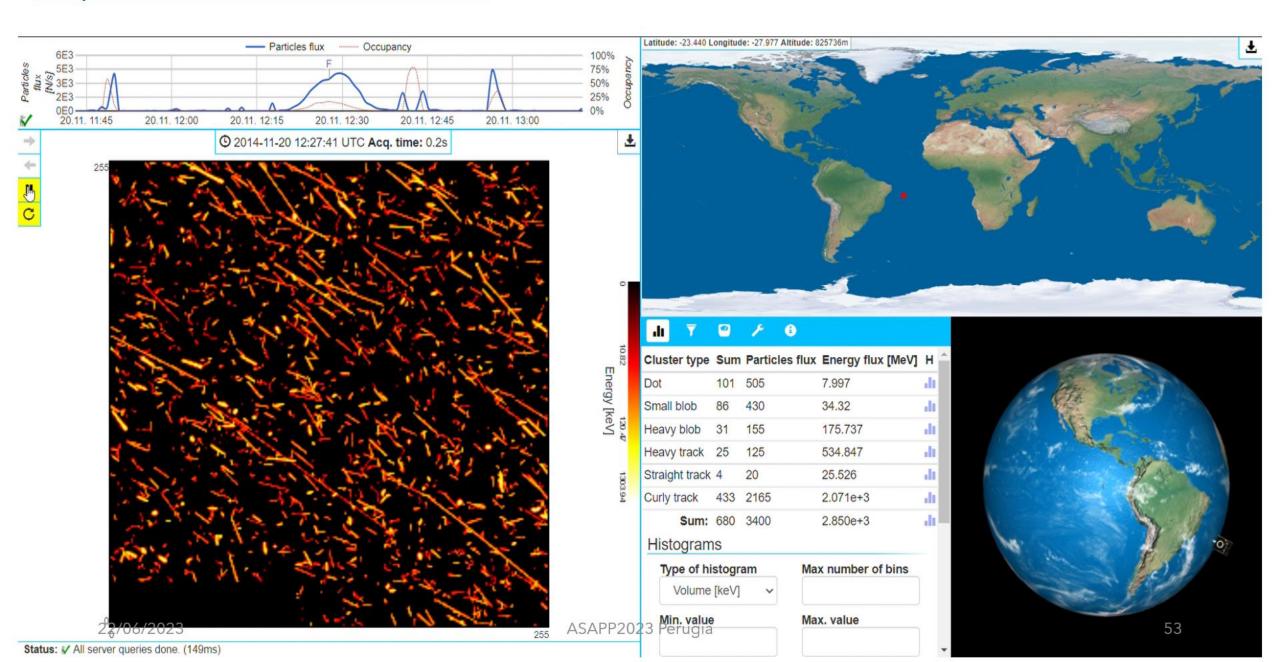


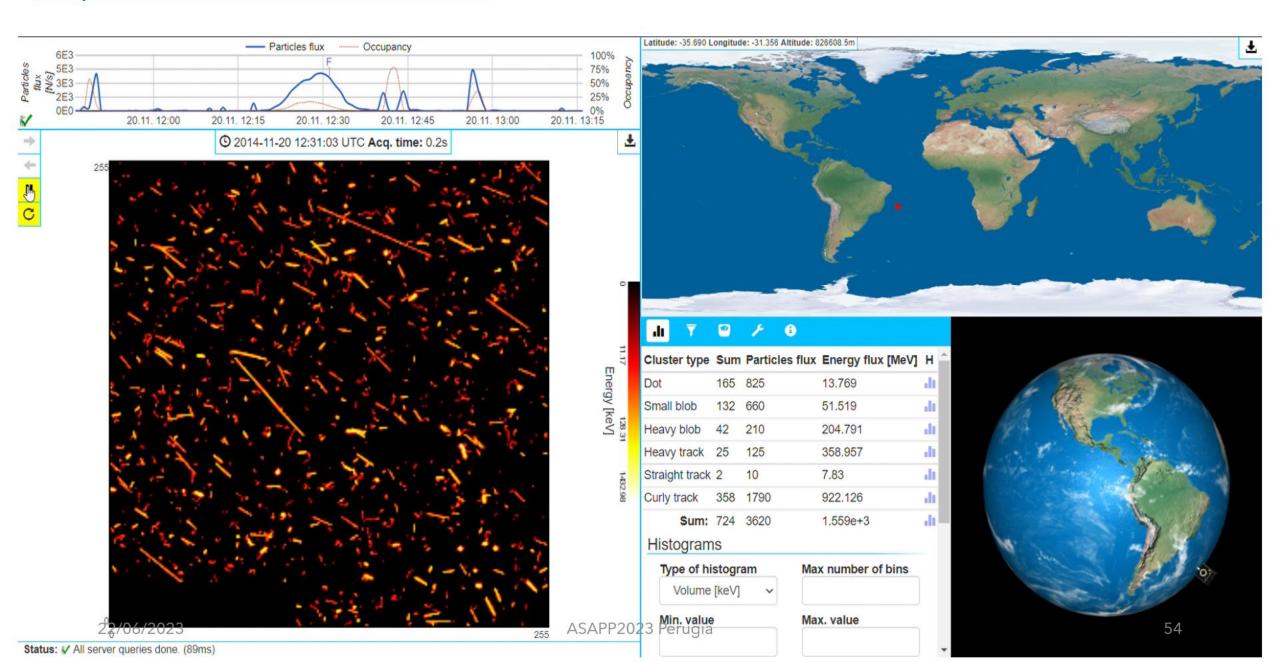


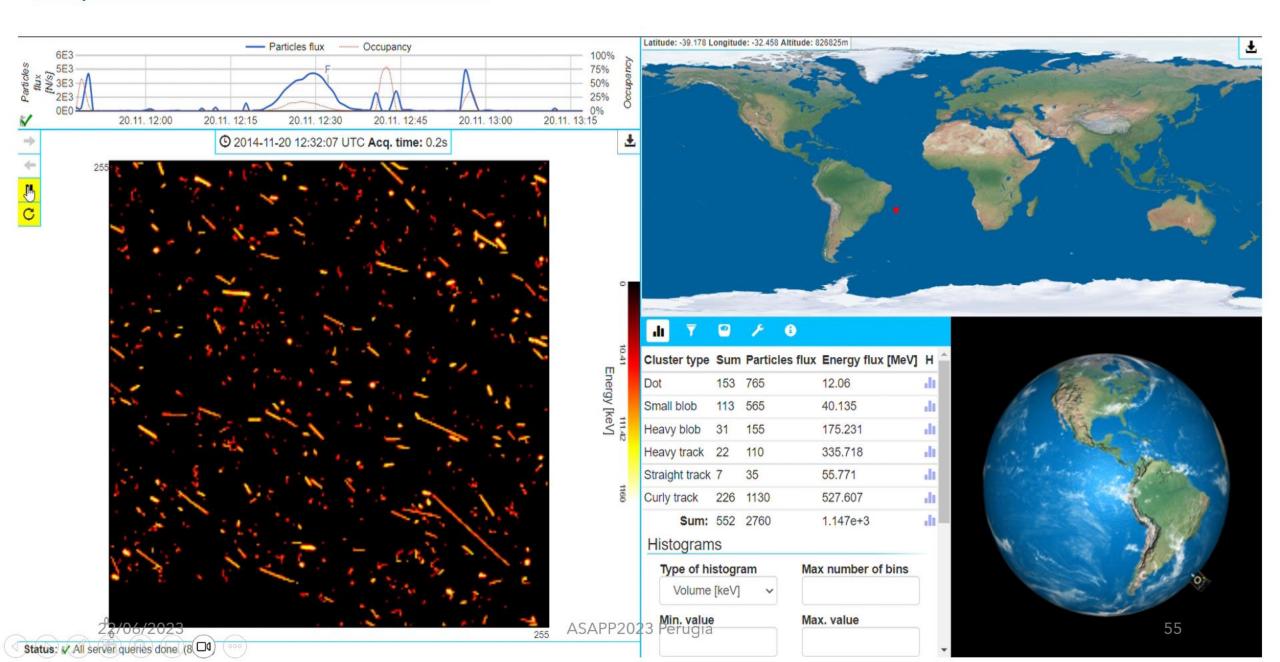


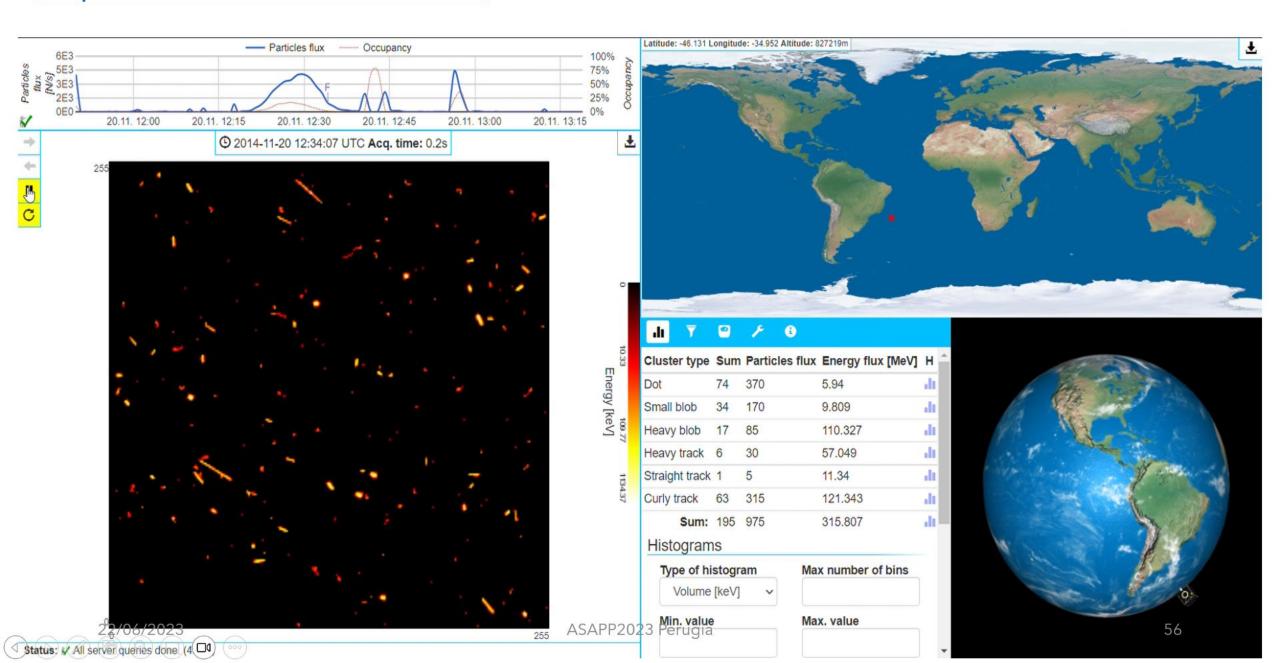


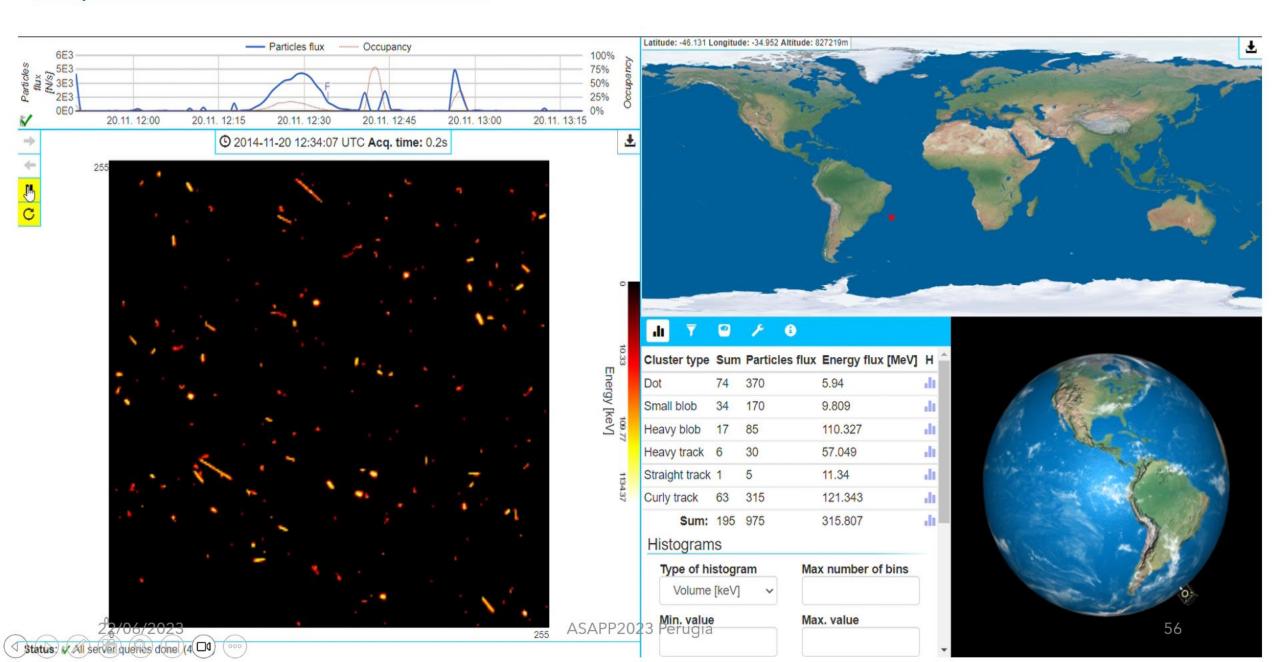


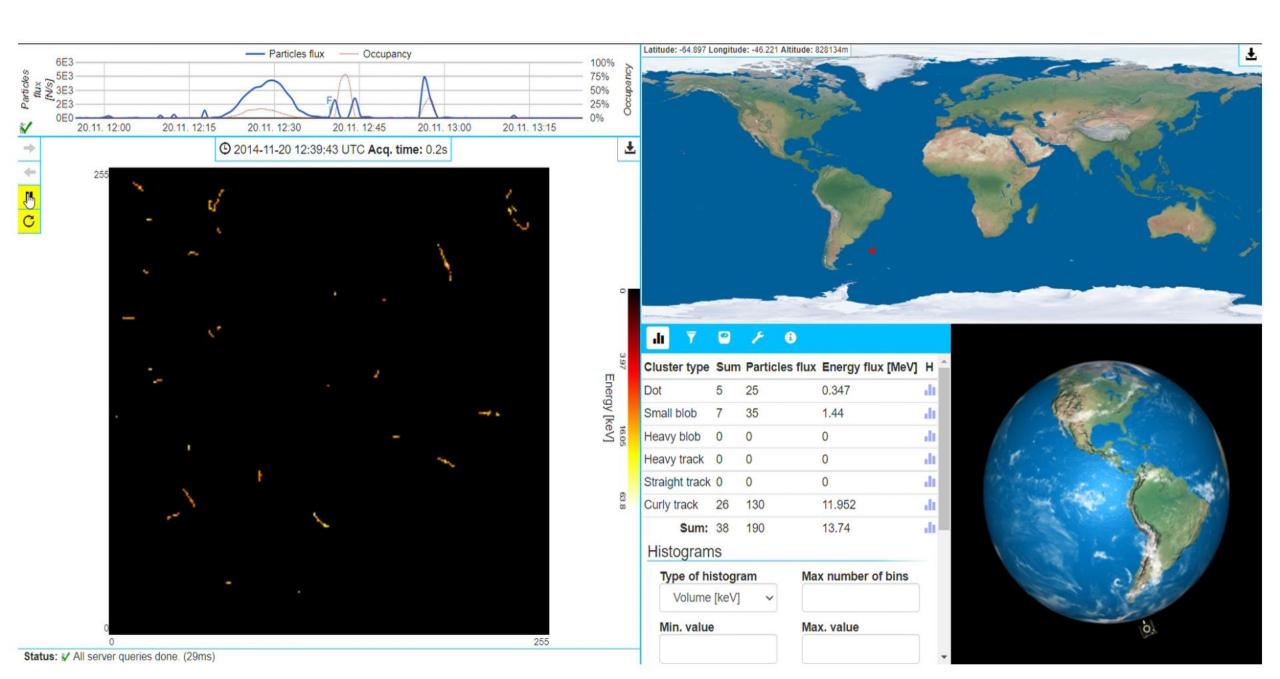






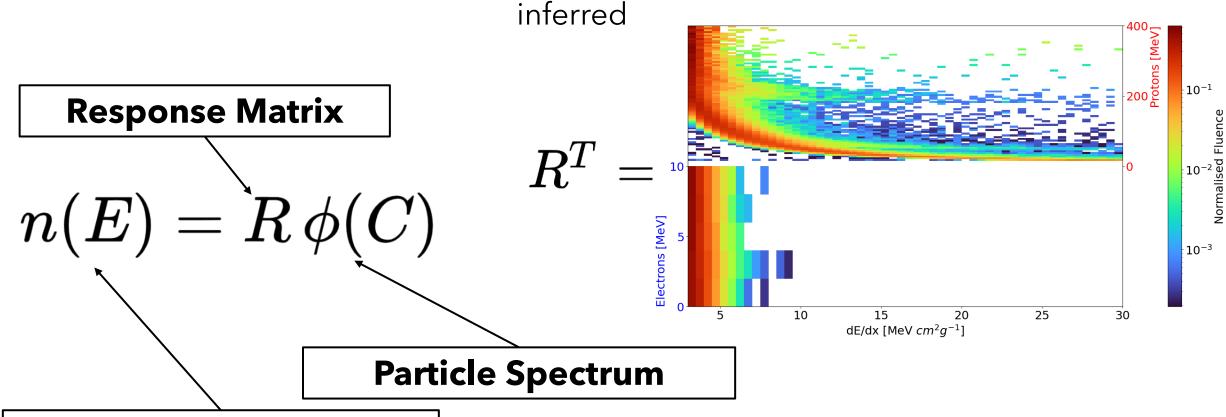






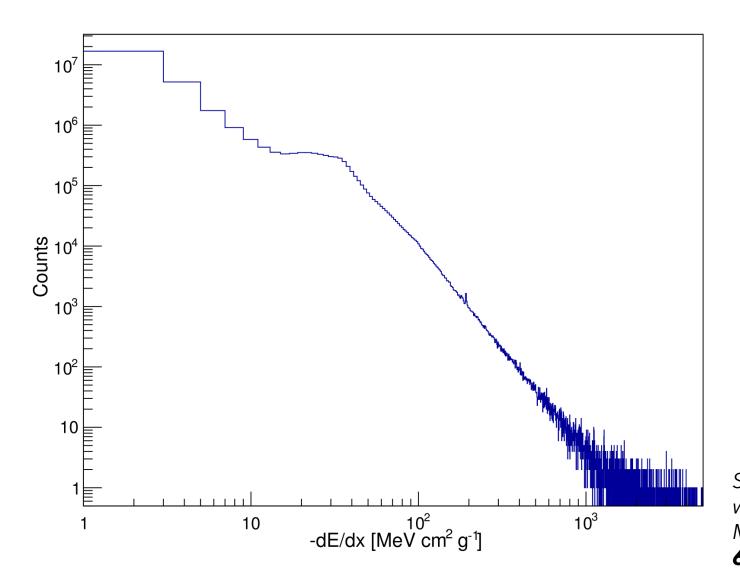
Particle Classification - Bayesian Deconvolution

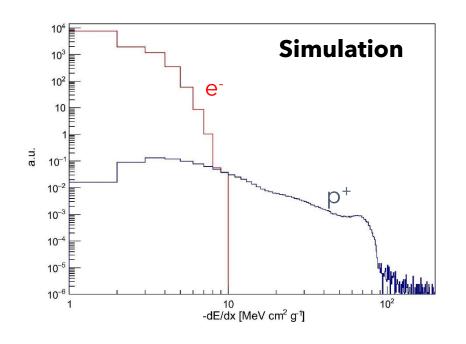
This method works by decomposing the stopping power "signal" of the field into its contributing particle signals, from which the particle's distributions can be



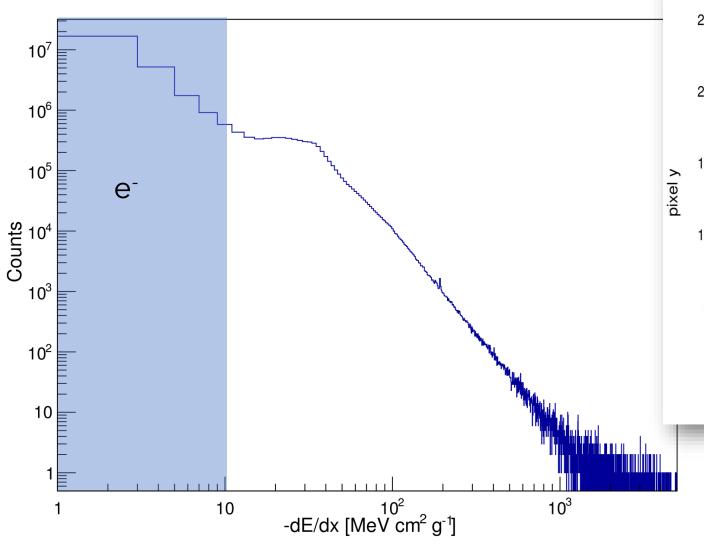
Measured dE/dx Spectrum

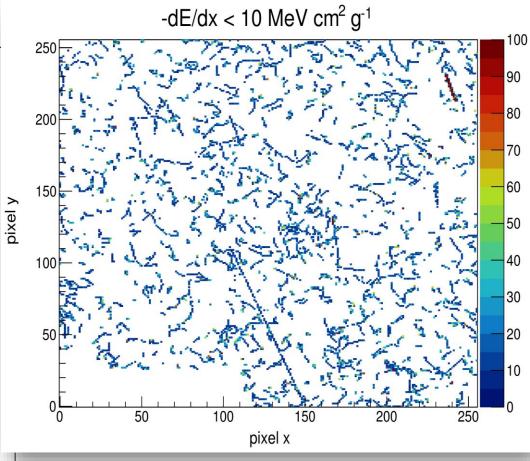
G. D'Agostini. Improved iterative Bayesian unfolding. 2010. https://doi.org/10.48550/arXiv.1010.0632



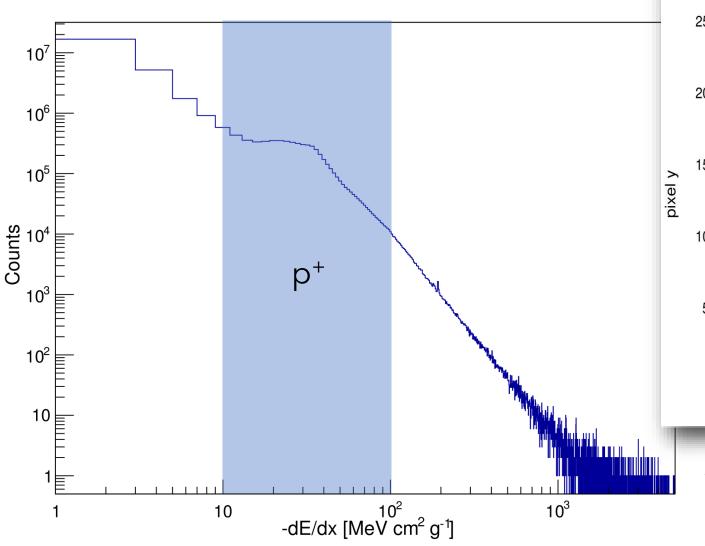


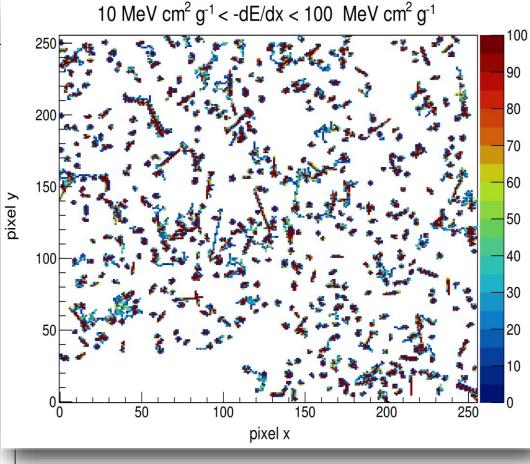
St. Gohl et al., "Study of the radiation fields in LEO with the Space Application of Timepix Radiation Monitor (SATRAM)", Advances in Space Research **63**, Issue 5, pp. 1646-1660, (2019).



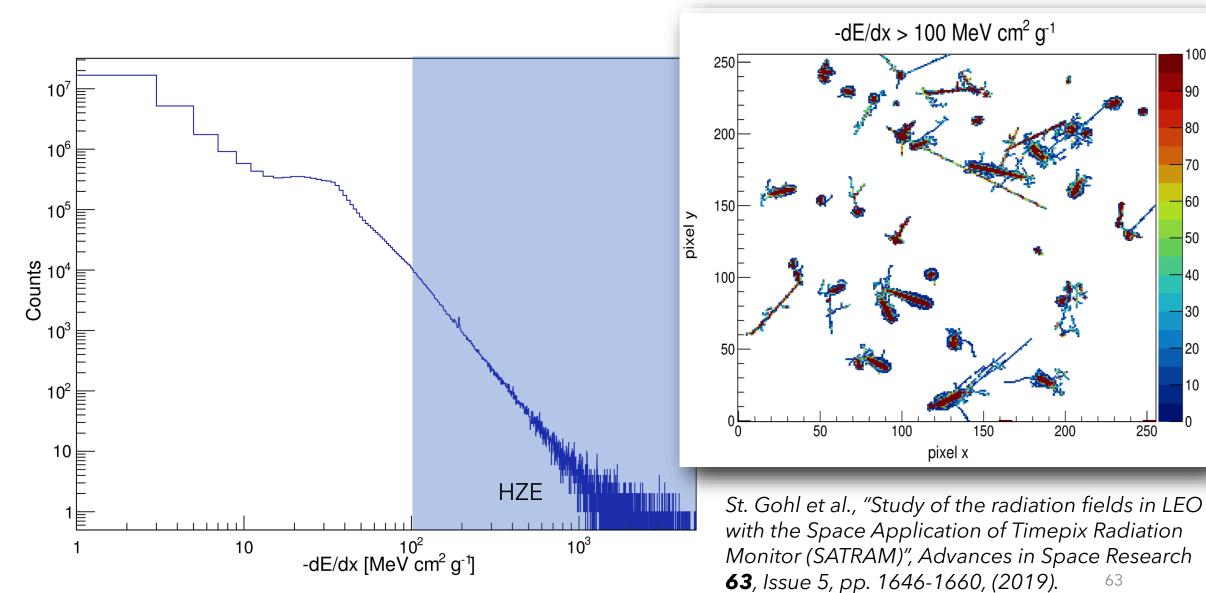


St. Gohl et al., "Study of the radiation fields in LEO with the Space Application of Timepix Radiation Monitor (SATRAM)", Advances in Space Research **63**, Issue 5, pp. 1646-1660, (2019).





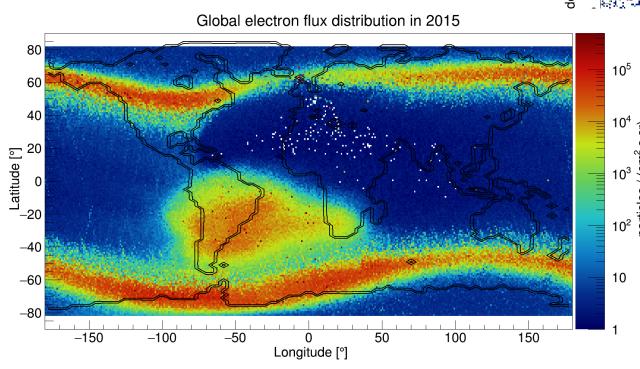
St. Gohl et al., "Study of the radiation fields in LEO with the Space Application of Timepix Radiation Monitor (SATRAM)", Advances in Space Research **63**, Issue 5, pp. 1646-1660, (2019).

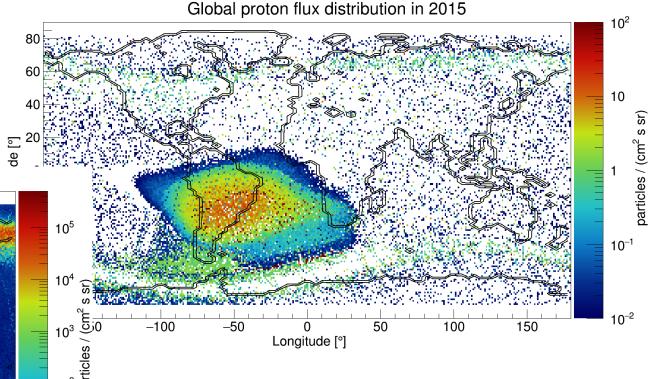


Electron and Proton Flux Maps

e⁻ fluxes 3 orders of magnitude larger than p⁺ fluxes

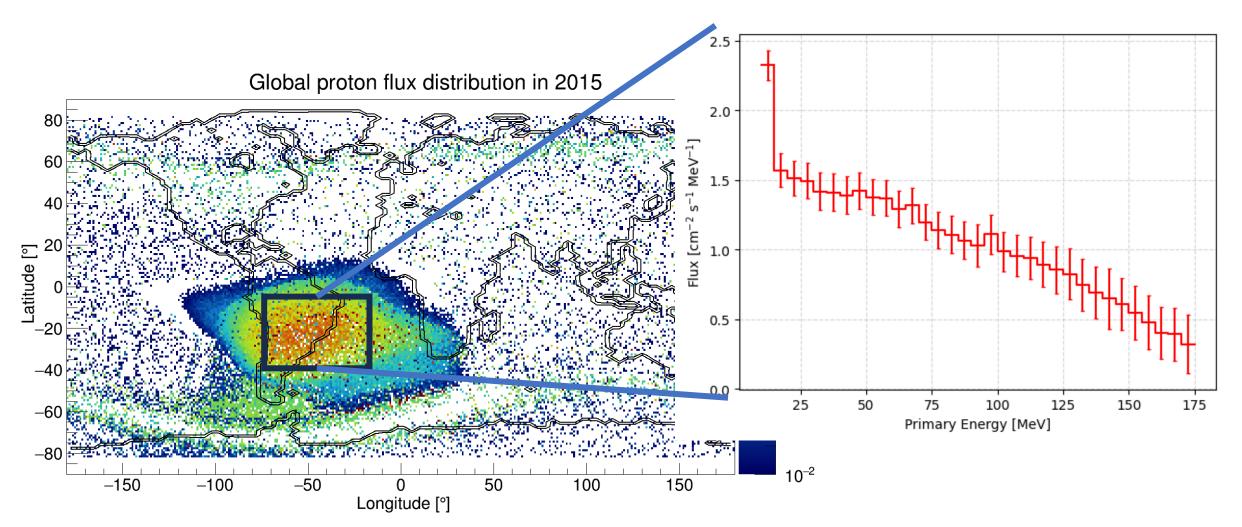
➡ Even small e⁻ misclassification distorts p⁺ flux measurement





St. Gohl et al., "Study of the radiation fields in LEO with the Space Application of Timepix Radiation Monitor (SATRAM)", Advances in Space Research 63, Issue 5, pp. 1646-1660 (2019).

First Measurement of the Trapped Proton Energy Spectrum with a Single-Layer Detector



Stopping Power Classification at MoEDAL

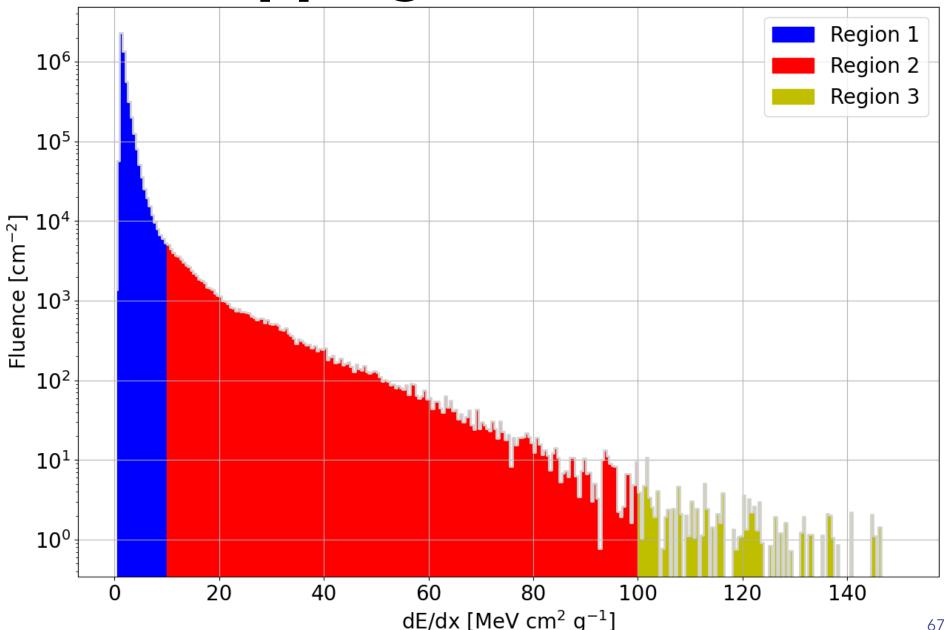
Region 1:

- ⇒Electrons
- ⇒Gammas
- Singly charged relativistic particles

Region 2:

- ⇒Protons <250 MeV
- ⇒Pions <15 MeV
- ⇒Muons <10 MeV</p>

- ⇒Fragmentations
- ⇒ Z>1 particles



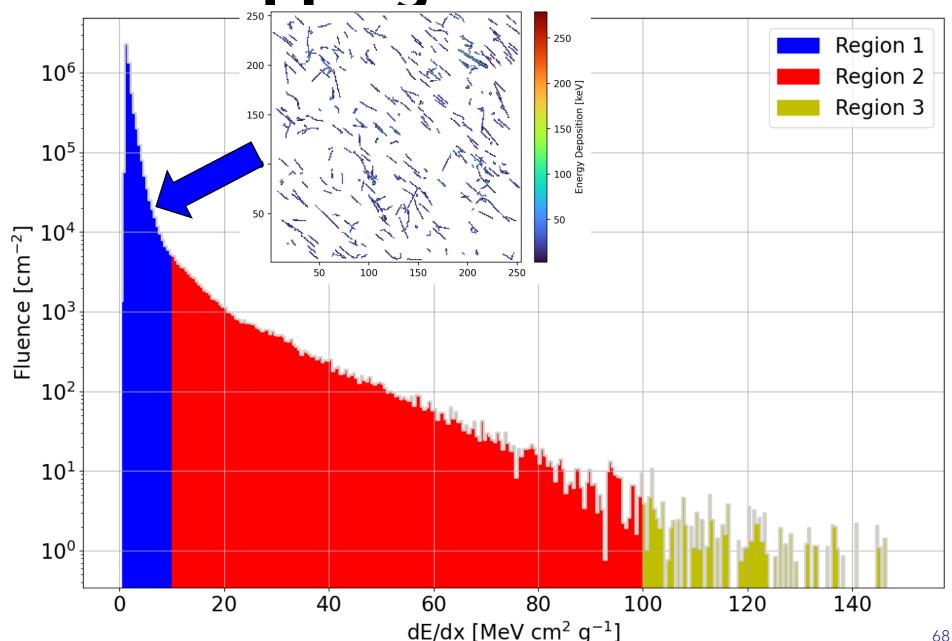
Region 1:

- ⇒Electrons
- ⇒Gammas
- Singly charged relativistic particles

Region 2:

- ⇒Protons <250 MeV
- ⇒Pions <15 MeV
- ⇒Muons <10 MeV</p>

- ⇒Fragmentations
- ⇒ Z>1 particles



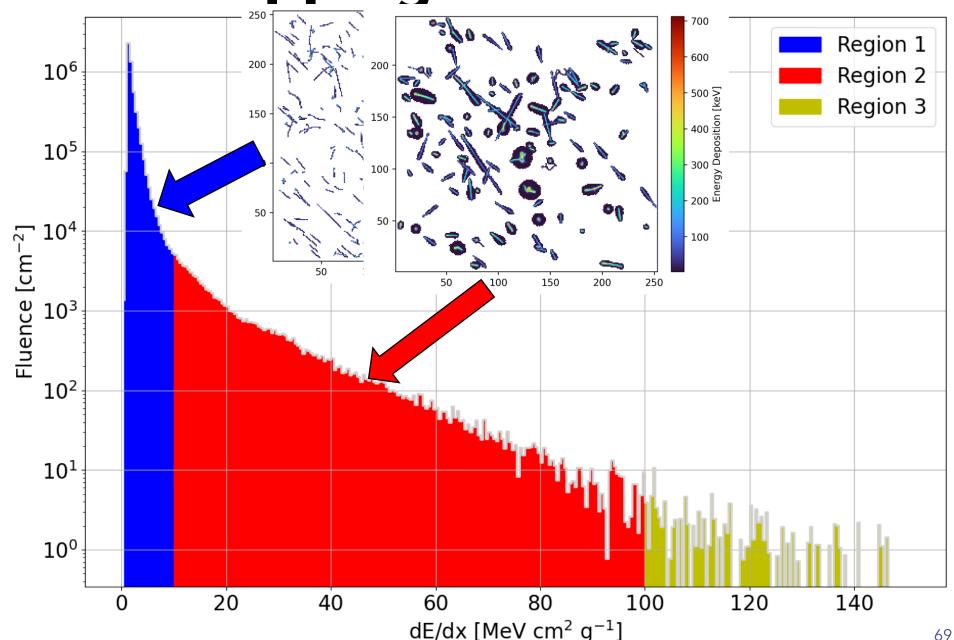
Region 1:

- ⇒Electrons
- ⇒Gammas
- Singly charged relativistic particles

Region 2:

- ⇒Protons <250 MeV
- ⇒Pions <15 MeV
- ⇒Muons <10 MeV</p>

- ⇒Fragmentations
- ⇒ Z>1 particles



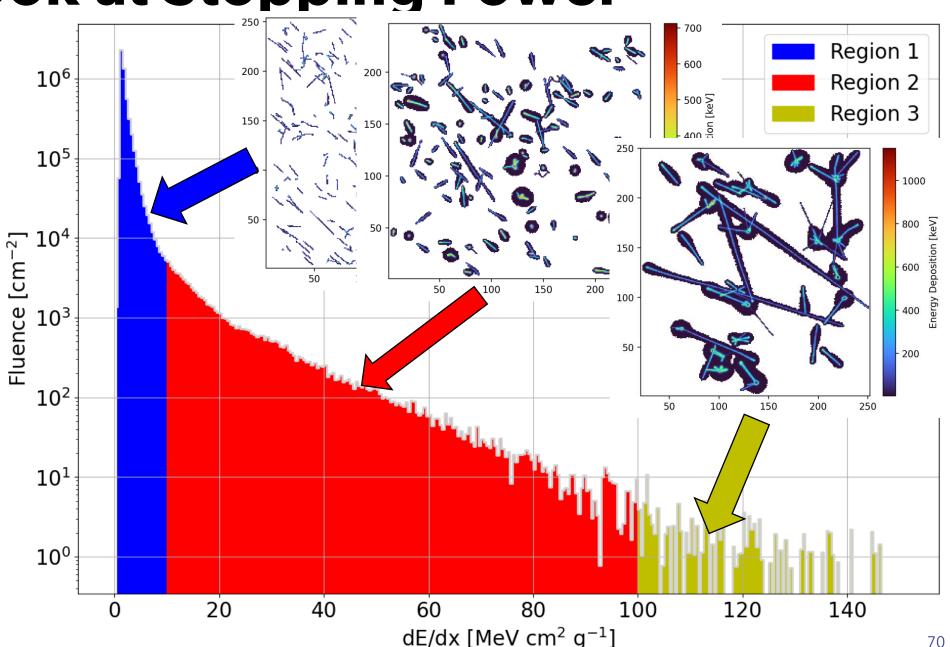
Region 1:

- ⇒Electrons
- ⇒Gammas
- Singly charged relativistic particles

Region 2:

- ⇒Protons <250 MeV
- ⇒Pions <15 MeV
- ⇒Muons <10 MeV</p>

- ⇒Fragmentations
- ⇒ Z>1 particles



Conclusions

- 3D tracking
 - The capabilities of reconstructing particle trajectories from IP8 in real time and with single layer Timepix3 detectors have been demonstrated
- Stopping power and radiation field decomposition
 - The current are state-of-the-art algorithms in particle classification were shown and validated in the environment found in low earth orbit
 - Future work will be to improve particle classification by using advanced machine learning algorithms

Can we use Timepix3 in Search for the Avatars of New Physics?

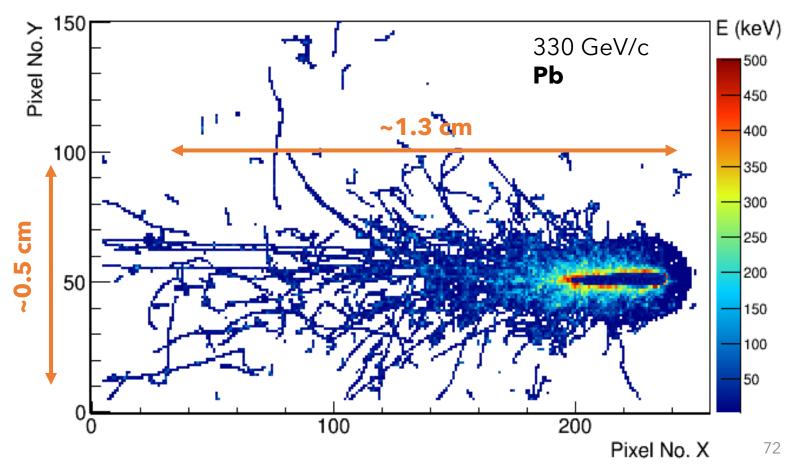


Highly charged energetic particles are known to produce distinct cluster pattern within Timepix3

Associated with:

- High Stopping Power
- Large number of delta rays

$$\left\langle -\frac{dE}{dx} \right\rangle = Kz^2 \frac{Z}{A} \frac{1}{\beta^2} \left[\frac{1}{2} \ln \frac{2m_{\rm e}c^2 \beta^2 \gamma^2 W_{\rm max}}{I^2} - \beta^2 - \frac{\delta(\beta \gamma)}{2} \right]$$



Thank you for your attention

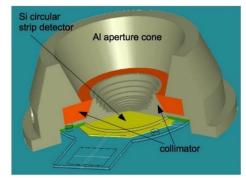
This research was funded by the Czech Science Foundation grant number GM23-04869M.

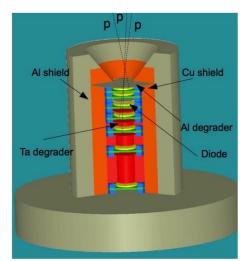
Instruments for measurements in LEO



Next Generation Radiation Monitor (NGRM)

- Mass ~ 1 kg
- Consumption ~1-2 W





EPT (Energetic Particle Telescope)

• Mass: **4.6 kg**

Consumption: 5.6 W



ICARE-NG:

• Mass: **2.4 kg**

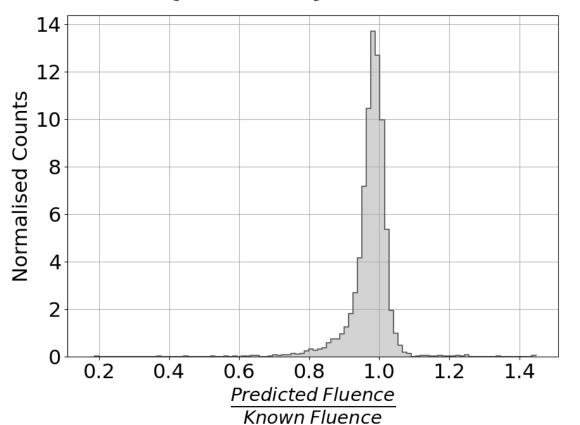
• Consumption: 3

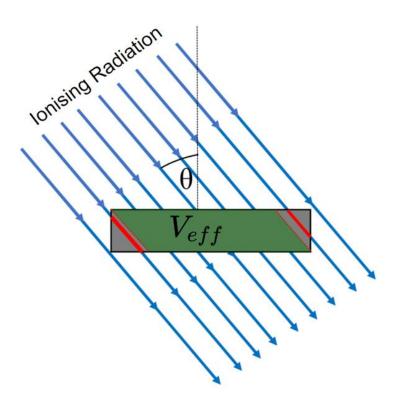


Particle Fluence

The formula was derived analytically to account for the removal of edge clusters

$$F = \sum_{i \in \{particles\}} rac{d/cos(heta_i)}{l \cdot d \cdot (l - d \cdot tan(heta_i))}$$





The new algorithm was tested in a Geant4 omnidirectional field simulation.

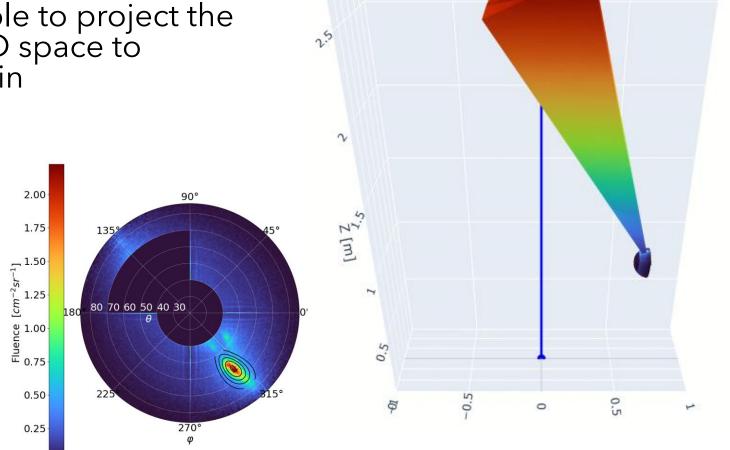
Testing Accuracy: 97%

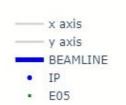
Second Peak Localisation

Follow correct alignment of the interaction of point

• It becomes possible to project the second peak in 3D space to determine its origin

 Projection shows the origin along the beam line further the theory of scattering beam due to wider beam





Bayesian Deconvolution Mathematical Background

Response Matrix

$$n(E) = R \phi(C)$$
—

Particle Spectrum

Measured dE/dx Spectrum

$$Mpprox R^{-1}\Rightarrow \phi(C)=M\,n(E)$$

Bayesian Formula \Rightarrow $P(C_{\mu}|E_{j}) = rac{P(E_{j}|C_{\mu})\,P(C_{\mu})}{\sum_{
u}^{n_{C}}\,P(E_{j}|C_{
u})\,P(C_{
u})}$

Probability of a cause (particle class present) given an effect (measurement)

$$\Rightarrow \phi(C_{\mu}) = \sum\limits_{i=1}^{n_E} P(C_{\mu}|E_i) n(E_i)$$