## Jet SIFT-ing

A Scale-Invariant Jet Clustering Algorithm for the Substructure Era

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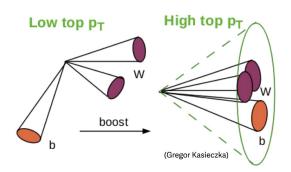


#### with

Andrew Larkoski (UCLA), Denis Rathjens (CMS), and Jason Veatch (ATLAS) arXiv: 2302.08609

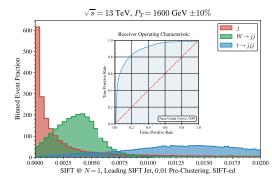
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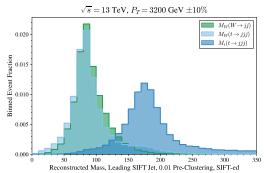
### SIFT: Scale-Invariant Filtered Tree



- Massive resonances decay into hard prongs
- Jet definitions with fixed cones impose a scale
- Boosted objects collimate and structure is lost
- Substructure recovery techniques are complex
- Can we avoid losing resolution in the first place?
- Select proximal objects w/ scale-invariant measure
- Candidate pairs are merged, dropped, or isolated, according to criteria integrated into the SI measure
- SIFT unifies: a) large-radius jet finding, b) filtering of soft wide radiation, and c) substructure axis finding into a single-pass prescription for low/high boosts

$$\delta_{AB} \equiv \frac{\Delta M_{AB}^2}{E_{\mathrm{T}A}^2 + E_{\mathrm{T}B}^2}$$





- *N*-subjet Tree holds superposition of projections onto *N*=1,2,3 prongs
- Hard prongs are preserved to end
- The measure history discriminates
   N=1,2,3 typically above 90% AUC
  - Faithful kinematic reconstruction

## Standard kT Jet Clustering Algorithms

- Debris from showering & hadronization must be reassembled in a manner that preserves correlation with the underlying hard (partonic) event
- 3 related algorithms reference an input angular width  $R_0$  & differ by an index n
- Objects wider than R<sub>0</sub> will never be clustered; Objects inside cone always merge
- n = 0, or "Cambridge/Aachen" favors objects with high angular adjacency
- n = +1, or "kT" additionally favors clustering where one of the pair is soft
- n = -1, or "Anti-kT" prioritizes clustering where one of the pair is hard
- Anti-kT is now the default jet clustering tool at LHC, with  $R_0 \sim 0.5$
- It is robust against "soft" and "collinear" jet perturbations and has regular jet shapes which are favorable for calibration against pileup, etc.

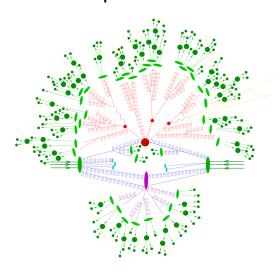


Image: Stefan Höche

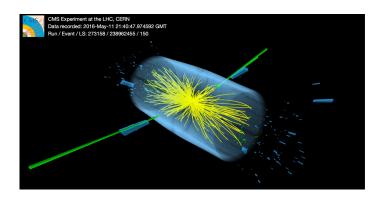
$$\delta_{AB} \equiv \min \left[ P_{\mathrm{T}A}^{2n}, P_{\mathrm{T}B}^{2n} \right] \times \left( \frac{\Delta R}{R_0} \right)^2$$

#### A Scale-Invariant Distance Measure

- It is worth asking whether alternative techniques could provide intrinsic resiliency to boosted event structure; this requires dropping the input scale R<sub>0</sub>
- It would be good to "asymptotically" recover key behaviors of Anti-kT
- Numerator should favor angular collimation; we propose  $\Delta M^2$ , similar to JADE
- Denominator should suppress soft pairings; we propose  $\Sigma E_T^2$ , similar to Geneva
- Result is dimensionless, Lorentz invariant (longitudinally in the denominator), and free from references to external / arbitrary scales

$$\delta_{AB} \equiv \frac{\Delta M_{AB}^2}{E_{TA}^2 + E_{TB}^2}$$

$$\delta_{AB} \equiv rac{\Delta M_{AB}^2}{E_{{
m T}A}^2 + E_{{
m T}B}^2} \qquad \Delta m_{AB}^2 \equiv (p_A^\mu + p_B^\mu)^2 - m_A^2 - m_B^2 = 2p_A^\mu p_\mu^B \\ \simeq 2E^A E^B \times (1 - \cos \Delta \theta_{AB}) \simeq E^A E^B \Delta \, \theta_{AB}^2$$



$$\begin{split} E_{\mathrm{T}} &\equiv \sqrt{M^2 + \vec{P}_{\mathrm{T}} \cdot \vec{P}_{\mathrm{T}}} = \sqrt{E^2 - P_z^2} \\ \lim_{M=0} &\Rightarrow |\vec{P}_{\mathrm{T}}| \end{split}$$

Image: CMS

## Comparison to the Geneva Measure

$$y_{ij} = \frac{8}{9} \frac{E_i E_j (1 - \cos \theta_{ij})}{(E_i + E_j)^2}$$

- Though motivated for new reasons, our measure is similar to "Geneva"
- In addition to normalization, there are three primary differences:
  - Sum of squares rather than square of sum (minor change)
  - Transverse cylindrical coordinates are referenced, as suitable for hadron collider rather than electron collider applications (relevant change)
  - Mass of merger candidates is accounted for (significant change)
- The more novel updates are not to the measure, but relate instead to:
  - Filtering of stray radiation and a related halting criterion
  - The concept of an N-subject Tree (superposition of axis candidates)

## Moving Toward a Geometric Measure

- An efficient algorithm needs something like a "GEOMETRIC" neighbor finding
- We need to refer to the collider coordinates of A & B directly ( $\Delta \eta_{AB}$ ,  $\Delta \phi_{AB}$ , etc.)
- For massive A & B, it will actually be rapidity  $\Delta y_{AB}$  that is relevant
- Boost from the  $P_z=0$  frame into the lab to examine the SIFT numerator:

$$\Delta m_{AB}^2 = 2 \times \left( E^A E^B - p_z^A p_z^B - p_T^A p_T^B \cos \Delta \phi_{AB} \right)$$

$$\begin{pmatrix} E \\ p_z \end{pmatrix} = \begin{pmatrix} \cosh y & \sinh y \\ \sinh y & \cosh y \end{pmatrix} \begin{pmatrix} E_{\rm T} \\ 0 \end{pmatrix} = \begin{pmatrix} E_{\rm T} \cosh y \\ E_{\rm T} \sinh y \end{pmatrix}$$
 • The difference between  $E_T \& P_T$  (i.e. MASS) means that we cannot perfectly

$$E^{A}E^{B} - p_{z}^{A}p_{z}^{B}$$

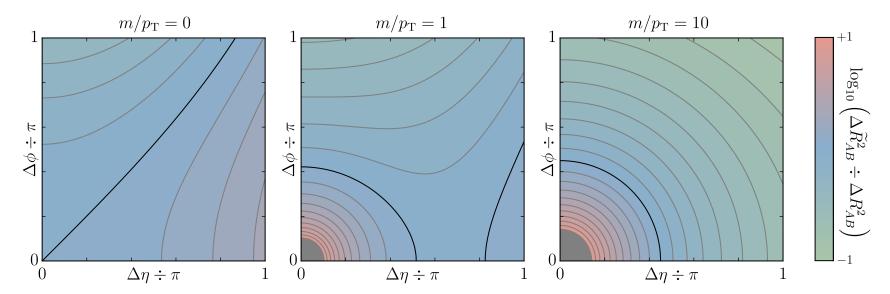
$$= E_{T}^{A}E_{T}^{B} \times \left(\cosh y^{A}\cosh y^{B} - \sinh y^{A}\sinh y^{B}\right)$$

$$= E_{T}^{A}E_{T}^{B} \times \cosh \Delta y_{AB}$$

- factorize kinematics from geometrics
- The role of  $\xi$  is to deemphasize azimuthal differences in the nonrelativistic limit

$$\Delta m_{AB}^{2} = 2 E_{\rm T}^{A} E_{\rm T}^{B} \times \left( \cosh \Delta y_{AB} - \xi^{A} \xi^{B} \cos \Delta \phi_{AB} \right)$$
$$\xi \equiv \frac{p_{\rm T}}{E_{\rm T}} = \left( 1 - \frac{m^{2}}{E_{\rm T}^{2}} \right)^{+1/2} = \left( 1 + \frac{m^{2}}{p_{\rm T}^{2}} \right)^{-1/2}$$

#### Comparative Angular Response



$$\Delta \widetilde{R}_{AB}^{2} \equiv \frac{\Delta m_{AB}^{2}}{E_{T}^{A} E_{T}^{B}}$$

$$= 2 \times \left(\cosh \Delta y_{AB} - \xi^{A} \xi^{B} \cos \Delta \phi_{AB}\right)$$

$$\simeq \Delta \eta_{AB}^{2} + \Delta \phi_{AB}^{2} \equiv \Delta R_{AB}^{2}$$

$$\Delta \widetilde{R}_{AB}^{2} \Rightarrow \Delta R_{AB}^{2} + \left\{ 1 - \frac{\Delta R_{AB}^{2}}{2} \right\} \times \left\{ \left( \frac{m_{A}}{p_{\mathrm{T}}^{A}} \right)^{2} + \left( \frac{m_{B}}{p_{\mathrm{T}}^{B}} \right)^{2} \right\} + \cdots$$

- The  $\Delta R^2$  measure is recovered for zero mass & small angular separations
- Hyperbolic cosine differs from cosine in that all Taylor terms are POSITIVE ... rapidity separations dominate azimuth
- Massive or low-pT objects resist clustering, even at small angles; this is a type of BEAM MEASURE

## Geometrizing the Denominator

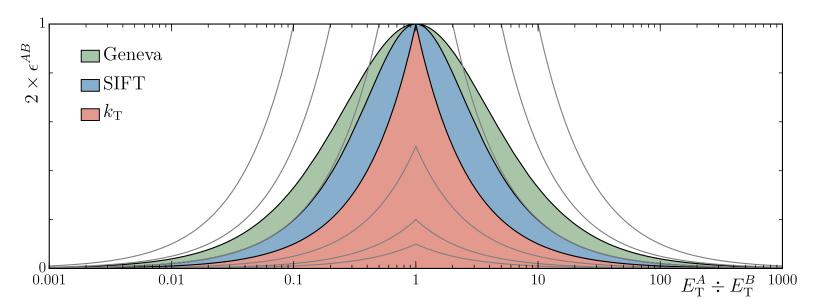
- The remainder of the metric refers only to transverse energy RATIOS
- This factor has a symmetry under  $E_T \rightarrow 1/E_T$
- It asymptotically mimics BOTH kt and anti-kt clustering, preferencing the clustering of pairs with hierarchically DISPARATE transverse scales
- It has the benefit of being ANALYTIC

$$\epsilon^{AB} \equiv \frac{E_{\rm T}^A E_{\rm T}^B}{(E_{\rm T}^A)^2 + (E_{\rm T}^B)^2} = \left\{ \left( \frac{E_{\rm T}^A}{E_{\rm T}^B} \right) + \left( \frac{E_{\rm T}^B}{E_{\rm T}^A} \right) \right\}^{-1}$$

$$u \equiv \ln \left( E_{\rm T} / [{\rm GeV}] \right)$$

$$\epsilon^{AB} = \left( e^{+\Delta u_{AB}} + e^{-\Delta u_{AB}} \right)^{-1} = \left( 2 \cosh \Delta u_{AB} \right)^{-1}$$

#### Comparative Energy-Momentum Response



$$\begin{split} \epsilon^{AB} &= \left\{ \min \left( \frac{E_{\mathrm{T}}^{A}}{E_{\mathrm{T}}^{B}} \right) + \max \left( \frac{E_{\mathrm{T}}^{A}}{E_{\mathrm{T}}^{B}} \right) \right\}^{-1} \\ &\simeq \left\{ \max \left( \frac{E_{\mathrm{T}}^{A}}{E_{\mathrm{T}}^{B}} \right) \right\}^{-1} \\ &= \min \left( \frac{E_{\mathrm{T}}^{A}}{E_{\mathrm{T}}^{B}} \right) \simeq \frac{\min(E_{\mathrm{T}}^{A}, E_{\mathrm{T}}^{B})}{E_{\mathrm{T}}^{A} + E_{\mathrm{T}}^{B}} \end{split}$$

$$\delta_{AB}^{k_{\mathrm{T}}} \, \stackrel{\textstyle \propto}{\sim} \, \left(\frac{E_{\mathrm{T}}^A E_{\mathrm{T}}^B}{E_0^2}\right)^{\!\!n} \! \times \min\left[\frac{E_{\mathrm{T}}^A}{E_{\mathrm{T}}^B}, \frac{E_{\mathrm{T}}^B}{E_{\mathrm{T}}^A}\right]$$

- SIFT & Geneva are scale invariant here
- The kT algorithms SCALE the overall response by a power of the geometric mean of transverse energies
- Grey contours are 0.1, 0.2, 0.5, 2, 5, 10, with reverse ordering for anti-kT

## All Together: the SIFT Measure

$$\delta_{AB} = \epsilon^{AB} \times \Delta \widetilde{R}_{AB}^{2}$$

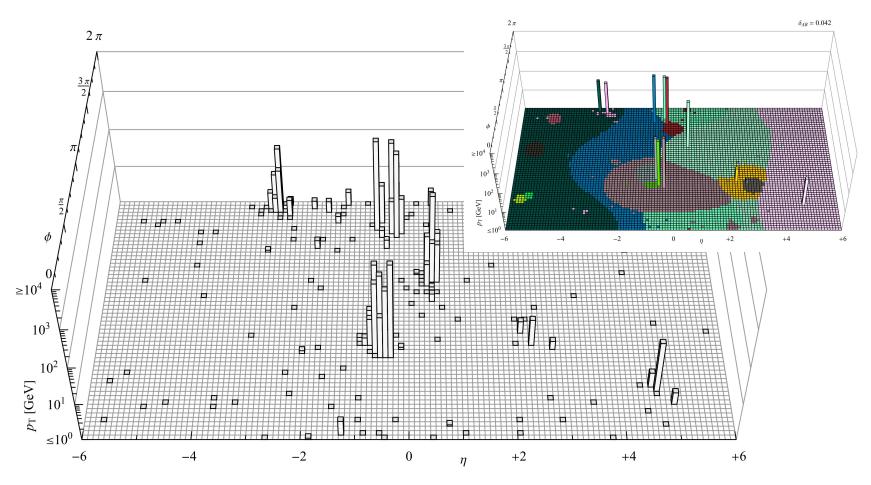
$$= \frac{\cosh \Delta y_{AB} - \xi^{A} \xi^{B} \cos \Delta \phi_{AB}}{\cosh \Delta u_{AB}}$$

- The measure is a simple product of energy and angular-type factors
- Clustering preferences pairs that are (relatively) soft and/or collinear
- Since mutually hard (relative to other available radiation) members will defer clustering, prongy structure is preserved to the end and easily accessed

Several problems remain beyond the measure (read on for the solutions ...)

- Extraneous wide and soft radiation is assimilated very early
- This distorts the kinematic reconstruction (mass especially)
- Moreover, there is no sense of when to \*stop\* clustering

# pp to TTbar (pT ~ 800 GeV) Scale Invariant Clustering with Ghost Radiation



See Video "A" Posted at Indico

## **FILTERING** Stray Radiation

- We know, at least, how to deal with soft, wide-angle radiation
- Take a cue from "Soft Drop" (2014 Larkoski, Marzani, Soyez, Thaler)
- This "Grooming" removes contaminants like ISR, UE, and pileup
- SD iteratively DECLUSTERS C/A, dropping softer object unless & until:

$$\frac{\min(P_{TA}, P_{TB})}{P_{TA} + P_{TB}} > z_{\text{cut}} \left(\frac{\Delta R_{AB}}{R_0}\right)^{\beta}$$

- Typically,  $z_{\rm cut}$  is  $\mathcal{O}(0.1)$ , and  $\beta > 0$  for grooming
- We propose an analog to be applied within the original clustering itself, expressible in the scale invariant language

Cluster: 
$$\frac{\Delta \widetilde{R}_{AB}^2}{2} < \{(2 \epsilon^{AB}) \le 1\}$$

- With factors of 2 in their "natural" places the maximal effective cone size is  $\sqrt{2}$
- This is a DYNAMIC boundary, and the angular size reduces for imbalanced scales

## Dropping vs. Isolating

- This leaves the question of what to do when clustering FAILS ...
- There are two distinct ways to fail the filtering criterion, to be handled differently
- The scale disparity can be too extreme (soft radiation) at O(1) angular separation

$$(\epsilon_{AB} \ll 1)$$
 and  $(\Delta \widetilde{R}_{AB}^2 \simeq 1)$ 

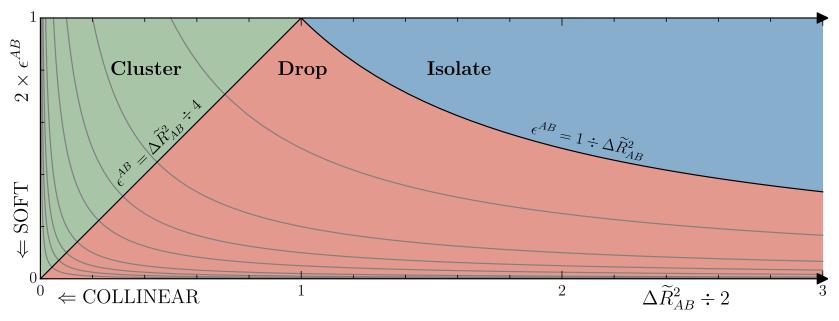
- In this case the metric product is small ... DROP the softer member
- Or, the angular separation can be too large (wide angle) with comparable scales

$$(\Delta \widetilde{R}_{AB}^2 \gg 1)$$
 and  $(\epsilon_{AB} \simeq 1)$ 

In this case the metric product is large ... ISOLATE both objects

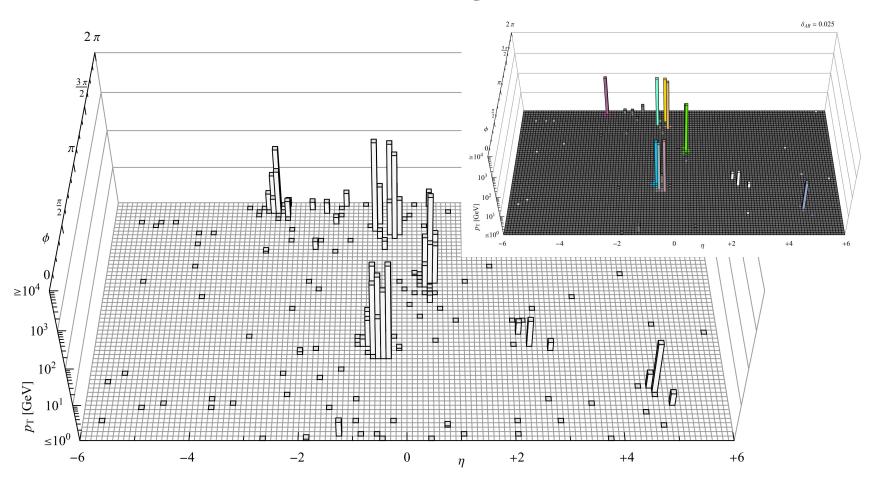
Isolate: 
$$\{1\} \leq \delta_{AB}$$
  
Drop:  $\{(2 \epsilon_{AB})^2 \leq 1\} \leq \delta_{AB} < \{1\}$ 

## Clustering Phase Diagram



- The unification of clustering, filtering, and isolation also provides natural halting
- Grey contours " $y = \delta/x$ " mark constant values of the measure
- Isolation occurs above  $\delta=1$ ; this amounts finding of variable large-radius jets
- The same factors separate clustering from dropping at "y = x"

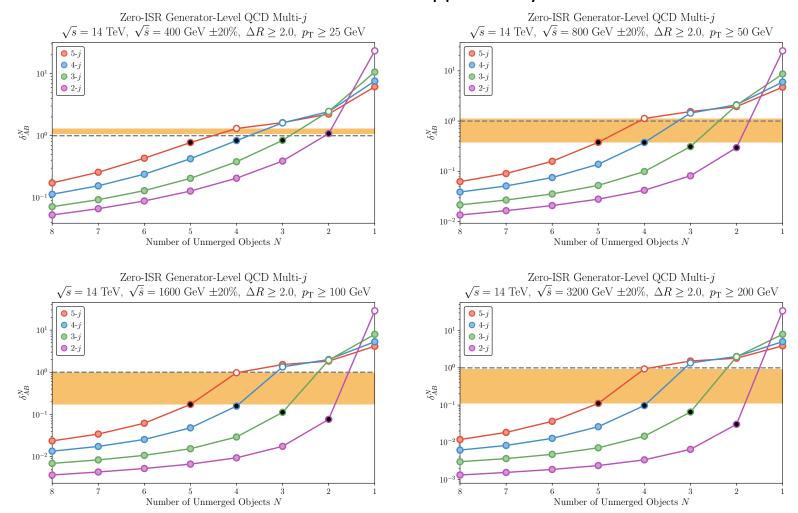
# pp to TTbar (pT ~ 800 GeV) Filtered Scale Invariant Clustering with Ghost Radiation



See Video "E" Posted at Indico

### **Evolution of the Measure**

- The measure "jumps" when it crosses the natural joint count
- The transition to isolation for  $\delta \geq 1$  is supported by simulation



## The N-Subjet TREE

- We observe that:
  - hard structures are preserved
  - wide concentrations of hard objects are isolated
  - soft wide radiation is dropped
- However, hard prongs within a variable radius jet do still cluster
- How do we fix the interior halting criterion to avoid losing structure?
- The most interesting alternative is to not halt at all ...
- We learn more about whether the prongs "want" to merge by merging!
- Hard prongs are the final objects to be merged, and we retain a superposition of projections onto all numbers N of prongs – suitable for computing N-subjettiness
- The record of structure is also directly imprinted on the measure history

## Tagging Jet Substructure

- N-Subjettiness  $\tau_N$  is the leading tool for characterizing how well a given event matches an N-prong hypothesis (axes chosen separately)
- The best discrimination comes from the ratio  $r_N$ , e.g. how much more 3-prong-like is the event than 2-prong like
- However, this procedure is also substantially complicated

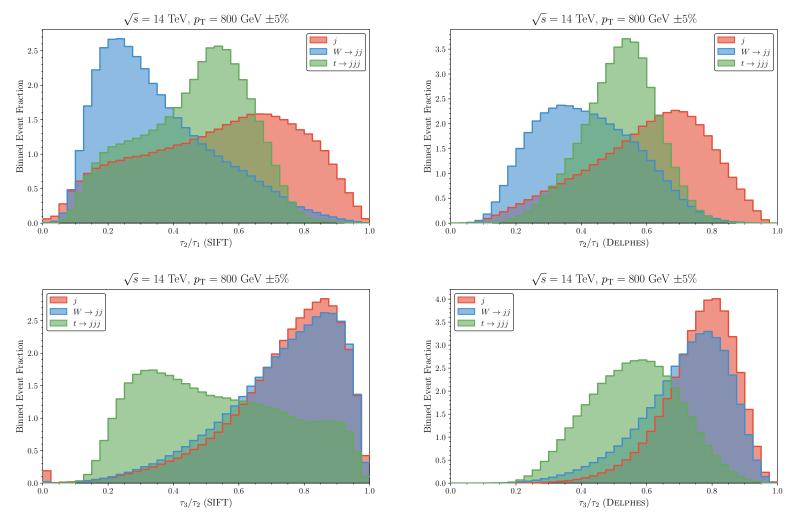
Given 
$$N$$
 axes  $\hat{n}_k$ ,  $\tau_N = \frac{\sum_{i \in J} p_{T,i} \min(\Delta R_{ik})}{\sum_{i \in J} p_{T,i} R_0}$ 

$$r_N = \frac{\tau_N}{\tau_{N-1}}$$

- It is interesting to ask if structure tagging can be incorporated into clustering
- To compare and assess performance, we simulate 1, 2 (W > j j), and 3 (t > j j j)
  jet event samples, at a range of transverse scales

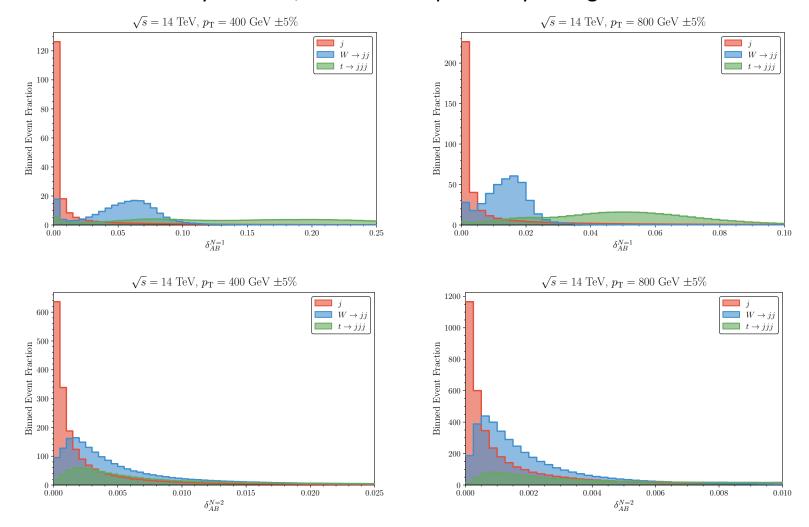
## $\tau_2/\tau_1$ and $\tau_3/\tau_2$ with SIFT Axes

SIFT is very good for N-subjettiness axis finding (Delphes versions on right)



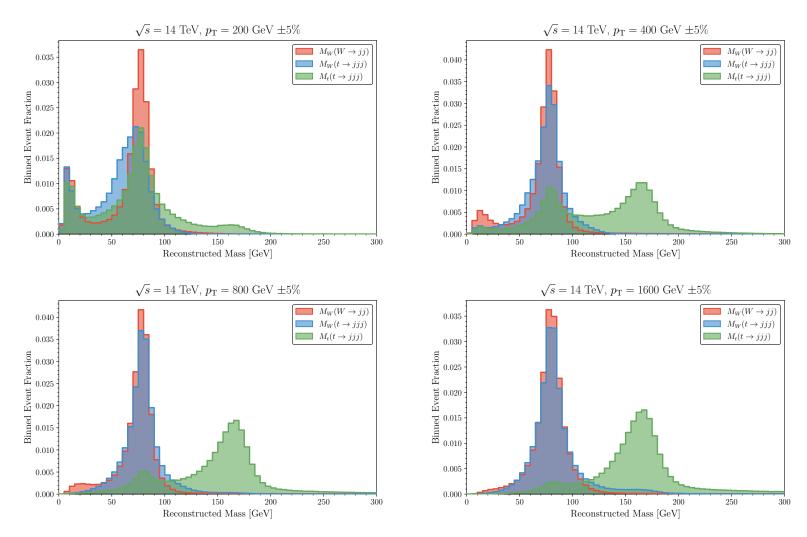
## SIFT Measure at Final Mergers

- We are also interested in whether the SIFT measure tracks jettiness DIRECTLY
- It seems not only to do so, but to excel specifically at large boost



## W & top Mass Reconstruction

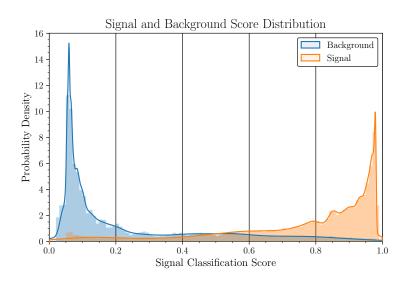
The included filtering also gives sharp accurate mass reconstruction at large boost



## **Assessing Performance**

- A Boosted Decision Tree lets us compare information density in an unbiased way
- The BDT is also completely transparent, since it amounts simply to cascaded binary selection cuts (branchings) with assigned scores
- We feed the BDT Delphes N-subjettiness ratios up to 5/4
- We also provide it with the final values of the SIFT measure
- We compare outcomes in isolation, and with both data sets provided together
- We compare the power of 2/1 and 3/2 discrimination at a range of scales

## 1/2 and 3/2 Discrimination with BDT



$p_{ m T}^{ m GeV\pm 5\%}$	$ au_{ ext{Delphes}}^{N+1/N}$	$ au_{ ext{SIFT}}^{N+1/N}$	$\delta^N_{AB}$	$\delta + \tau$
100	0.62	0.68	0.69	0.70
200	0.91	0.86	0.88	0.89
400	0.89	0.85	0.91	0.92
800	0.82	0.79	0.92	0.93
1600	0.77	0.74	0.91	0.92
3200	0.78	0.76	0.88	0.90

TABLE III. Area under curve ROC scores for discrimination of resonances with hard 1- and 2-prong substructure using a BDT trained on various sets of event observables.

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0.2	1	Are	a Under C	urve: 0.91			
0.0		.4 0 Positive		.8 1.0			

$p_{ m T}^{ m GeV\pm 5\%}$	$ au_{ ext{Delphes}}^{N+1/N}$	$ au_{ ext{SIFT}}^{N+1/N}$	$\delta^N_{AB}$	$\delta + \tau$
100	0.61	0.61	0.63	0.65
200	0.63	0.60	0.71	0.72
400	0.82	0.74	0.90	0.90
800	0.85	0.80	0.94	0.95
1600	0.77	0.77	0.97	0.97
3200	0.77	0.79	0.98	0.99

TABLE IV. Area under curve ROC scores for discrimination of resonances with hard 2- and 3-prong substructure using a BDT trained on various sets of event observables.

## **Summary and Conclusions**

- SIFT is a SCALE INVARIANT clustering algorithm designed to avoid losing substructure
- FILTERING of soft-wide radiation and variable-radius isolation is fully integrated
- The measure history & TREE of N-subjet axis candidates encode structure on the fly

### Thank You!

Mathematica movie-generating notebook is available with arXiV source (2302.08609)

Also, please join us this evening for some homemade gumbo!

## 4502 Lapis Court

(8-mile drive from Mitchell) all are welcome – there will be plenty of food those with cars are encouraged to offer rides

### Software Advertisement



- All data analysis for this project was performed with the indicated set of tools
- The package is available for download & public use from GitHub:
- https://github.com/joelwwalker/AEACuS
- I will help you!

### Automated collider event selection, plotting, & machine learning with AEACuS, RHADAManTHUS, & MInOS

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A trio of automated collider event analysis tools are described and demonstrated, in the form of a quick-start tutorial. AEACuS interfaces with the standard MadGraph/MadEvent, Pythia, and Delphes simulation chain, via the Root file output. An extensive algorithm library facilitates the computation of standard collider event variables and the transformation of object groups (including jet clustering and substructure analysis). Arbitrary user-defined variables and external function calls are also supported. An efficient mechanism is provided for sorting events into channels with distinct features. RHADAManTHUS generates publication-quality one- and two-dimensional histograms from event statistics computed by AEACuS, calling MatPlotLib on the back end. Large batches of simulation (representing either distinct final states and/or oversampling of a common phase space) are merged internally, and per-event weights are handled consistently throughout. Arbitrary bin-wise functional transformations are readily specified, e.g. for visualizing signalto-background significance as a function of cut threshold. MInOS implements machine learning on computed event statistics with XGBoost. Ensemble training against distinct background components may be combined to generate composite classifications with enhanced discrimination. ROC curves, as well as score distribution, feature importance, and significance plots are generated on the fly. Each of these tools is controlled via instructions supplied in a reusable cardfile, employing a simple, compact, and powerful meta-language syntax.